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OFF-CENTER-LINE SHOCK-INTERFERENCE HEATING PATTERNS ON BASIC SHAPES IN HYPERSONIC FLOWS

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SUMMARY

Experimental off-center-line shock-interference heating results on basic bodies are presented for a range of hypersonic Mach numbers, specific heat ratios, Reynolds numbers, and impinging shock strengths. Three interference flow patterns explicitly defined in previous studies were generated by the intersection between a plane impinging shock and the bow shock created by a hemisphere or a cylindrical fin.

Heat-transfer rates higher than stagnation values were measured over much of the off-center-line model surface. For example, heating levels equal to or much higher than stagnation existed over most of the lower half of the hemisphere. Heating amplifications up to 17 times that for stagnation were measured on the hemisphere because of jet impingement. The results of this study imply that shock interference on a high-speed flight vehicle can cause severe heating over large areas near the interaction region.

INTRODUCTION

Shock-interference heating can be a serious problem in the design of high-speed vehicles such as the space shuttle and hypersonic cruise aircraft. Heat-transfer values several times the stagnation-point values can occur on vehicle surfaces because of the interfering flow fields. (See refs. 1 to 3.) These interfering flow fields are influenced by free-stream flow conditions, real-gas effects, and angle-of-attack changes as the vehicle moves along its trajectory.

Edney (ref. 1) made a detailed study of the entire spectrum of interference flows and defined six types of shock-interference patterns. The local pressure and heating peaks were a result of either shock—boundary-layer interaction, free shear-layer attachment, or supersonic jet impingement. The results of an extensive analytical and experimental study of shock-interference heating at hypersonic speeds are reported in references 2 to 6. Reference 2 presented theoretical flow models and semiempirical methods of predicting the peak pressure and heat-transfer levels for the six types of shockinterference patterns; also included were center-line experimental pressure and heattransfer distributions on hemispheres, a 30° wedge, and a cylindrical leading-edge fin for different interference patterns.

This paper presents the effects of free-stream Mach number and Reynolds number, specific heat ratio (shock density ratio), and impinging shock strength on the interference heating amplification away from the plane of symmetry. Heat-transfer data were obtained over a range of nominal Mach numbers (6 to 20), specific heat ratios (1.40 and 1.67), and impinging shock strengths (5° to 25°). This investigation supplements the study reported in reference 2.

SYMBOLS

The measurements and values are given in SI units; calculations were made in U.S. Customary Units.

A	cylinder length (see fig. 2(a))
с	specific heat of model material
h	heat-transfer coefficient
k	thermal conductivity of model material
L	fin length
M	Mach number
^N Re	Reynolds number
р	pressure
R _b	radius of hemisphere
Т	temperature
X = x/L	
x	axial coordinate
x., y.	coordinates of shock intersection in plane of symmetry (see fig. 2)

 x_i, y_i

γ	ratio of specific heats
θ	body angle for hemisphere (see figs. $2(a)$ and $5(a)$)
θ _i	shock generator angle
^θ SL	shear-layer angle relative to surface (see fig. $3(c)$)
λ .	sweep angle of fin
ρ	density of model material
^p 2 ^{/p} 1	normal-shock density ratio
φ	angle in a plane perpendicular to axis of hemisphere or cylindrical leading- edge fin (see figs. 2 and 5)
Subscripts	:
aw	adiabatic wall
cyl	stagnation-line value
pk	peak value
ref	reference value

stagnation value at outer edge of boundary layer

total value wall value w

stag

t

free-stream value œ

Abbreviations (see fig. 3):

bow shock BS

IP	impingement point
IS	impinging shock
\mathbf{SL}	shear layer
TS ·	transmitted shock

APPARATUS AND TEST CONDITIONS

Test Facilities

The wind tunnels used in this study are summarized in table I:

Facility	Test gas	γ	M _∞	Facility described in reference –
Langley 20-inch Mach 6 tunnel	Air	1.40	6	- 7
Langley 22-inch helium tunnel	Helium	1.67	20	8
,				•

TABLE I. - TEST FACILITIES

Models

The experimental setup shown in figure 1 was designed to use interchangeable models. These interchangeable models (shown in fig. 2) consisted of a 2.54-cm- and a 5.08-cm-diameter hemisphere and a 2.54-cm-diameter cylindrical leading-edge fin. A 15.24-cm-wide by 25.40-cm-long sharp leading-edge plate which can be set at angles up to 30° was used as a shock generator. The main series of models were made of a silica-base epoxy material. Additional models of solid stainless steel (type 347) were used to measure peak heating in some tests.

The model thermophysical properties (density, specific heat, and thermal conductivity) were measured on samples of material for each model. The value of the product ρ ck used was approximately 1.798 × 10³ W $\sqrt{\text{sec}/\text{m}^2}$ -K for the epoxy models and 7.260 × 10³ W $\sqrt{\text{sec}/\text{m}^2}$ -K for the stainless-steel models. Further details are available in reference 2.

Test Parameters

A summary of the models, nominal flow conditions, and types of interference patterns investigated is presented in table II:

Test	Model	${ m M}_\infty$	γ	^{\$2/\$2} 1	$p_{t, \infty},$ N/cm ²	T _{t,∞} , K	N _{Re,∞} /m	Type of interference (ref. 1)
6352	5.08-cm-diameter hemisphere	6	1.40	5.5	83 290	478	7.9×10^{6} 25.6	III and IV
6363	5.08-cm-diameter hemisphere	6	1. 40	5.5	290	478	25.6	Ш and IV
6363	Fin	6	1.40	5.5	2 90	478	25.6	IVa and V
373	2.54-cm-diameter hemisphere	20.2	1.67	4.0	690	433	9.8	III and IV

TABLE II. - SUMMARY OF TEST PARAMETERS

The types of interference are discussed in a later section.

EXPERIMENTAL TECHNIQUE AND DATA ACCURACY

The methods of testing were similar for both facilities. The phase-change coating technique described in detail in references 9 and 10 was used to measure the heating. Because of the highly nonuniform heat-transfer distribution on the models, some models were tested with two or more temperature coatings. In some cases, peak heating was so intense that the coating temperature that would be needed was near T_{aw} . (This condition would indicate short melt times and the possibility of large errors.) To measure the peak heating in these cases, the stainless-steel model with its larger value of ρck was used; this model permitted the use of a coating with a lower melt temperature.

Phase-change heat-transfer data are subject to numerous errors. Errors in measuring the thermophysical properties of the model material, the melt temperature and time, the initial model wall temperature, and the initial time of exposure to the free stream can affect the accuracy of the heat-transfer data. Another important source of error is the determining of the adiabatic wall temperature distribution. In the present tests, T_{aw} was assumed to be equal to the free-stream total temperature. An analysis of these combined errors indicated that the heating amplification measured at Mach 6 in air may have a maximum error of about 29 percent (25 percent at Mach 20 in helium).

A semi-infinite slab solution of the general heat conduction equation was used to reduce the heating data. (See ref. 9.) In order for this assumption to be valid, the radius of curvature must be much greater than the depth of heat penetration. As stated in reference 10, the heat penetration depth for the semi-infinite slab is approximately independent of the aerodynamic heat-transfer coefficient and depends only on the thermal diffusivity of the wall material and the thermal diffusion time. The minimum distance

from peak heating for which the semi-infinite slab solution is valid can also be determined by using the melt time. (See ref. 10.) Although the phase-change method does not account for lateral conduction errors, these errors can be important. (See ref. 10 for additional details.)

Details of the interference flow regions were observed by using the schlieren technique. Surface flow patterns were obtained by using a mixture of titanium dioxide and silicone oil in the Mach 6 facility. Further details of the flow-visualization technique are given in reference 2.

TYPES OF SHOCK-INTERFERENCE PATTERNS

In estimating the effects of shock-interference heating, it is necessary to determine the type of interference pattern that will exist when two shocks of different strength intersect. The pattern that will occur depends on body geometry, Mach number, specific heat ratio, and strength of the impinging shock and its position relative to the body. (See ref. 1.) The six possible shock-interference-flow patterns as defined by Edney are shown in figure 3.

These interference patterns can be further grouped according to the flow mechanism that causes the heat-transfer change at the surface. Types I, II, and V are associated with a shock—boundary-layer interaction and type III is characterized by an attaching free shear layer. Type IV is characterized by an impinging or grazing supersonic jet (type IVa) and type VI by an expansion-fan—boundary-layer interaction. In the present study, only types III, IV, IVa, and V were investigated. Types III and IV occurred on the hemispheres and type IVa and V on the fin.

For a weak two-dimensional shock impinging on the bow shock of a hemisphere, the interference pattern in the plane of symmetry depends on the relative shock locations as shown in figure 4. The type of interference can be different away from the plane of symmetry where the bow-shock strength is weaker. A detailed description of these six types of shock-interference patterns and the transition from one type to another is given in references 1 and 2.

PRESENTATION OF RESULTS

The heat-transfer data are presented as nondimensional heat-transfer contours (h/h_{ref}) along with off-center-line peak heating distributions and the center-line heating distributions from reference 2. Selected schlieren photographs and oil-flow patterns are also shown. The data are presented according to the type of interference occurring on the center line, the free-stream Mach number and Reynolds number, the strength and

location of the impinging shock, and the model geometry. Table III provides an index to the figures presented in this paper.

TABLE III.- DATA FIGURES

Figure	Туре	$\mathrm{M}_{\mathrm{\infty}}$	^N Re,∞/m	γ	^{\$2/\$1}	^ℓ i nominal, deg	x _i /R _b	y _i /R _b	Model diameter, cm	Test-run numbers (ref. 2)
a ₆	ш	5.94	$7.9 imes 10^6$	1.40	5.5	15	0.214	-0.621	5.08	
7	Ш		7.7			10	.228	417		6352-16,17,30
8	ш		7.9			15	.214	621		6352-12,13
9	ш		7.6			20	.262	680		6352-21,22
10	ш		7.8			· 25	.262	786		6352-25,26
11	ш	6.00	25.9			15	.233	548		6363-36,37
12	Î	6.00	25.8	1	↓ `	15	.209	650	V	6363-39,40
13	ш	20.2	9,9	1.67	4.0	5			2.54	373-19,20
14	ш		9.7		1	5	-			373-22,23,24
15	ш		9.8			10	.310	370		373-11,12,26
16	ш		9.7			10	.260	480		373-8,9
17	ш		9.8			15				373-5,6
18	ш	↓	9.8	↓	↓	15			↓	373-2,3
^a 19	IV	6.00	25.7	1.40	5.5	15	.252	427	5.08	
20	IV	5.94	7.6			5	.200	141		6363-48
21	IV		8.0			10	.218	340		6352-14,15
22	IV		7.5			15	.257	340		6352-27,28
23	·IV		7.7			15	.248	513		6352-8,9,29
24	IV		7.6			20	.288	557		6352-18,19,20
25	IV	↓	7.8			25	.291	640		6352-23,24
26	IV	6.00	25.6			5	.204	199		6363-46
27	IV		25.0			10	.223	252		6352-31,32,33
28	IV		25.7			15	.252	427		6352-1,2,3
29	IV		25.6			20	.282	417		6363-41,42
30	IV	l ₩	25.4	V	₩	25	.301	534	l V	6363-43,44

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Figure	Туре	M∞	N _{Re,∞} /m	γ	^p 2/ ^p 1	^ℓ i nominal, deg	x _i /L	y _i /L	λ nominal, deg	Test-run numbers (ref. 2)
^a 31	IVa	6.00	25.8×10^6	1.40	5.5	20	0.305	0.105	0	
.32	IVa	5:94	8.1			10	.355	.080	0	6363-24
33	IVa		7.7			15	.435	.095	0	6363-22
34	IVa		7.9			20	.310	.110	0	6363-18
35	IVa	6.00	26.2			10	.340	.075	0	6363-23
36	IVa	1	25.4		[[15	.425	.090	0	6363-20
37	IVa		25.8			20	.305	.105	0	6363-19
^a 38	v	↓	25.9			20	.300	.07 <u>0</u>	25	
39	v	5.94	7.2			10	.330	.060	25	6363-6
40	v		7.3			15	.360	.065	25	6363-3
41	v		7.6			20	.310	.065	25	6363-14
42	v	6.00	25.7			10	.300	.050	25	6363-9
43	v		26.2			15	.340	.060	25	6363-10
44	v	↓	25.9	l v ∣	₩	20	.300	.070	25	. 6363-13

(b) Cylindrical fin

^aSchlieren and/or oil-flow photographs.

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The off-center-line peak heating data on the hemisphere are plotted by using the coordinate system defined in figure 5(a); values of θ and ϕ were measured by using the intersections of the individual heating contours h/h_{stag} with the attachment line or jet impingement line. (See fig. 7(a), for example.) Since the body is axisymmetric, this system collapses the off-center-line data on to the center line in the vertical plane of symmetry. The values of θ and ϕ are listed on the heating-contour figures. Off-center-line peak heating data on the fin were plotted by using the coordinate ϕ as shown in figure 5(b). In this case ϕ is located in a plane perpendicular to the fin center line. The undisturbed heating levels on the hemisphere and cylindrical fin are included for comparison purposes.

Reference values h_{stag} for the hemispheres are obtained from an expression for laminar stagnation-point heating in reference 11. Laminar stagnation-line heating h_{cyl} calculated from a method in reference 12 is used for the cylindrical leading-edge fin.

DISCUSSION OF RESULTS

Type III Interference

An extensive investigation of type III interference (shear-layer attachment) on hemispheres was conducted at Mach 6 in air ($\gamma = 1.40$) and Mach 20 in helium ($\gamma = 1.67$). Typical schlieren and oil-flow photographs for a type III interference pattern on a hemisphere are shown in figure 6. The reshaped bow shock and the attaching free shear layer which emanates from the shock intersection can be seen in figure 6(a). The attaching shear layer with its accompanying high shear flow appears as a "feather" pattern on the surface with the flow moving away from the attachment line (fig. 6(b)).

Hemisphere (5. 08-cm diameter). - Heat-transfer results for the 5. 08-cm-diameter hemisphere at Mach 6, $\gamma = 1.40$ ($\rho_2/\rho_1 = 5.5$) are presented in figures 7 to 12 for a range of shock generator angles and two free-stream Reynolds numbers. The attaching shear layer has a strong influence on the heating along the lower surface of the hemisphere as shown in part (a) of figures 7 to 12. For example, the attachment line (dashed line on contour plots) extends from the vertical plane of symmetry around to the tangent point of the hemisphere and cylindrical afterbody. The extent of the high heating region near the nose decreases as the local unit Reynolds number increases for a given θ_i and θ_{pk} . (See, for example, figs. 8(a) and 12(a).) The size of this region is also influenced by the location of the impinging-shock—bow-shock intersection (x_i/R_b , y_i/R_b) relative to the sonic point ($\theta \approx 47^{\circ}$) for a given θ_i and Reynolds number. (See figs. 11(a) and 12(a).)

Comparisons of the undisturbed heating on the model with the vertical center-line heating and the off-center-line peak heating (along the attachment line) are presented in parts (b) and (c), respectively, of figures 7 to 12. Peak heating values from 7.5 to 14 times the stagnation values were measured on the vertical center line. It was shown in reference 2 that this peak heating is a function of shock generator angle θ_i , the shearlayer angle relative to the surface θ_{SL} , and the state of the shear layer at attachment (laminar, transitional, or turbulent). The shear layers at Mach 6 are either transitional or turbulent. (See ref. 2.) The entire heating on the center line is considerably higher than the local undisturbed level (for example, 10 to 15 times at $\theta = -80^{\circ}$). Off-centerline attachment heating distributions indicate that heating levels of comparable severity to center-line heating are also observed over large parts of the hemisphere surface. These results show that even though the maximum peak is located near the center line, severe heating can extend over much of the model surface.

<u>Hemisphere (2.54-cm diameter)</u>. - Heating contours, center-line, and off-centerline peak heating distributions are presented in figures 13 to 18 at $\theta_1 = 5^{\circ}$, 10° , and 15° for various shear-layer attachment locations on the 2.54-cm-diameter hemisphere at Mach 20, $\gamma = 1.67$ ($\rho_2/\rho_1 = 4.0$). In general, the heating data in helium followed the same trends as the Mach 6 air data. However, the peak heating ratios were lower in helium as a result of the lower normal-shock—density ratio and the laminar shear-layer attachment (refs. 5 and 6) as opposed to transitional or turbulent attachment for the Mach 6 data.

Type IV Interference

Of the six types of interference, type IV (supersonic jet impinging in a subsonic region) results in the most severe heating (refs. 1 and 2); therefore, considerable interference heating data were obtained. A typical schlieren photograph of the shock and jet patterns is shown in figure 19(a) at Mach 6 in air. The effect of jet impingement on the surface oil flow on a hemisphere is shown in figure 19(b). At $\theta_i = 15^{\circ}$, the surface flow pattern appears to be similar to the pattern for shear-layer attachment (fig. 6(b)). The model stagnation region is severely displaced since the jet stagnation pressure can be as high as 6 to 8 times the model stagnation value. (See ref. 2.)

<u>Hemisphere (5.08-cm diameter)</u>. - Heating data at two Reynolds numbers for shock generator angles from 5° to 25° are presented in figures 20 to 30 for the 5.08-cm-diameter hemisphere at Mach 6, $\gamma = 1.40 \ (\frac{\rho}{2}/\frac{\rho}{1} = 5.5)$. The contours and center-line distribution at $\theta_i = 5^\circ$ in figures 20 and 26 are very similar to the shear-layer attachment patterns of type III; however, as θ_i increases, the heating patterns become very complex as shown in figures 21 to 25 and 27 to 30. The secondary areas of high heating appearing in these patterns are a result of flow distortions caused by shocks and expansions in the jet flow

moving along the surface. (See ref. 2.) Peak heating values 17 times the stagnation heating values were measured on the vertical center line. The variation of the maximum off-center-line peak heating along the jet impingement line (dashed line in part (c) of figs. 20 to 30) indicates that high heating extends far back on the model similar to type III off-center-line heating. The interference pattern may also change from one type to another with increasing distance from the vertical center line.

<u>Cylindrical fin (2.54-cm-diameter leading edge)</u>. - A typical schlieren photograph of a type IVa grazing jet pattern is presented in figure 31 for a 2.54-cm-diameter leadingedge cylindrical fin at Mach 6, $\gamma = 1.40 \ (\rho_2/\rho_1 = 5.5)$. In this case the jet is turned upward and flows parallel to the model surface since the pressure in the lower subsonic region is higher than that in the upper region. The dark and light areas near the surface above the shock interaction are compressions and expansions in the jet flow. The flow in this region is unsteady as indicated by the wavy bow shock.

Heating contours, center-line, and off-center-line heating distributions for type IVa interference are shown in figures 32 to 37. The peaks and dips observed in the heating distributions are caused by the compressions and expansions in the grazing jet flow interacting with the fin boundary layer. The sharp increase in heating and the spike appearing near the lower end of the fin (figs. 33 and 36) are a result of the expansion waves from the trailing edge of the shock generator interacting with the lower segment of the bow shock. Heating levels as high as the stagnation-line values were measured on the cylindrical leading edge as far back as $\phi \approx 60^{\circ}$. (See part (c) of figs. 32 to 37.)

Type V Interference

A schlieren photograph showing the complex flow pattern on the 2.54-cm-diameter leading-edge cylindrical fin at Mach 6, $\gamma = 1.40 \ (\rho_2/\rho_1 = 5.5)$ is shown in figure 38(a) at $\lambda = 25^{\circ}$ and $\theta_i = 20^{\circ}$. This photograph also shows that the jet flow is unsteady as for type IVa. The oil-flow pattern for the fin leading edge is shown in figure 38(b).

Heating contours, stagnation-line heating distributions, and off-stagnation-line peak values are presented in figures 39 to 44 for various θ_i values and two free-stream Reynolds numbers. In general, the heating contours and stagnation-line distributions exhibit similar shapes and trends, respectively, regardless of the impinging shock strength. The secondary peaks observed in the stagnation-line distribution are the result of the grazing jet and shear layer interacting with the boundary layer. (For example, compare the interactions shown in fig. 38(a) with the heating of fig. 39(c).) The peak-off-stagnation-line distributions show that heating as high as stagnation-line values can exist as far back as $\phi = 85^{\circ}$. (See fig. 30(c).) Small areas of high heating which appear to be caused by vortices from the lower corner were observed in some runs. (See figs. 43(a) and 44(a).)

CONCLUDING REMARKS

The results of an extensive experimental study of off-center-line shock interference heating on simple shapes are presented for three interference flow patterns. These patterns were formed by the intersection of a plane impinging shock and the bow shock created by a hemisphere or cylindrical fin. The study covered a range of hypersonic Mach numbers, specific heat ratios, Reynolds numbers, and impinging shock strengths.

Heating levels higher than stagnation values were found over a large percentage of the off-center-line body surface. For example, in most cases the lower half of the hemisphere was exposed to levels equal to or much greater than stagnation heating. Maximum amplifications up to 17 times the stagnation heating were measured on the hemisphere near the center line for an impinging jet. These results indicate that shock interference on high-speed flight vehicles can cause severe heating over significant areas in the vicinity of the interaction region.

Langley Research Center,

National Aeronautics and Space Administration, Hampton, Va., September 19, 1973.

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Figure 1. - Sketch of model test assembly.



(a) Hemispheres.



(b) Cylindrical fin.

Figure 2. - Sketch of models and center-line coordinate system.. Dimensions are in cm.



Figure 3. - Six types of shock-interference patterns.







(a) Hemisphere.



(b) Fin.

Figure 5. - Off-center-line coordinate system.



(a) Schlieren photograph.

Figure 6.- Shock and oil-flow patterns on a 5.08-cm-diameter hemisphere at Mach 5.94 in air for a type III interference. $\theta_i = 14.8^{\circ}$; $N_{\text{Re},\infty}/\text{m} = 7.9 \times 10^{6}$.



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(b) Oil-flow pattern.Figure 6.- Concluded.



Figure 7.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air.

 $\theta_{i} = 9.9^{\circ}; N_{Re,\infty}/m = 7.7 \times 10^{6}.$







(c) Off-center-line peak heating.





(a) Heating contours. $h_{stag} = 306.42 \text{ W/m}^2 \text{-K}$.

Figure 8.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 14.8^{\circ}$; $N_{Re,\infty}/m = 7.9 \times 10^{6}$.



(b) Center-line heat-transfer distribution.







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Figure 9.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 19.7^{\circ}$; $N_{Re,\infty}/m = 7.6 \times 10^{6}$.













Figure 10.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 24.6^{\circ}$; $N_{Re,\infty}/m = 7.8 \times 10^{6}$.



(b) Center-line heat-transfer distribution.





Figure 10. - Concluded.



(a) Heating contours. $h_{stag} = 592.41 \text{ W/m}^2 \text{-K}$.

Figure 11.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_i = 14.8^{\circ}$; $N_{Re,\infty}/m = 25.9 \times 10^{6}$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line peak heating.

Figure 11. - Concluded.



(a) Heating contours. $h_{stag} = 571.98 \text{ W/m}^2 \text{-K}$.

Figure 12.- Type III interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_i = 14.8^\circ$; $N_{Re,\infty}/m = 25.8 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line peak heating.









(a) Heating contours. $h_{stag} = 1205.25 \text{ W/m}^2 \text{-K}$.

Figure 13.- Type III interference on a 2.54-cm-diameter hemisphere at Mach 20.2 in helium. $\theta_{i} = 5.1^{\circ}; N_{Re, \infty}/m = 9.9 \times 10^{6}.$



(b) Center-line heat-transfer distribution.







Off-center-line neak locations

ns			
carlo	ϕ , deg	164 152 158 158 145 145 142	
ar 10	θ , deg	46 55 58 68 69 72 69 72	
he	h/h _{stag}	2.53 2.05 0.97 0.82 0.86 0.45 0.38	
	Isotherm	H0040068	
	Model	Steel Epoxy	
	Run	22	
	Test	373	





Figure 14.- Type III interference on a 2.54-cm-diameter hemisphere at Mach 20.2 in helium. $\theta_1 = 5.1^{\circ}$; $N_{Re,\infty}/m = 9.7 \times 10^{6}$.


(b) Center-line heat-transfer distribution.







Off-center-line





(a) Heating contours. $h_{stag} = 1205.25 \text{ W/m}^2\text{-K}$.

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Figure 15.- Type III interference on a 2.54-cm-diameter hemisphere at Mach 20.2 in helium. $\theta_1 = 10^{\circ}$; $N_{Re,\infty}/m = 9.8 \times 10^{6}$.











 ϕ , deg

49 40 34 30

(a) Heating contours. $h_{stag} = 1205.25 \text{ W/m}^2\text{-K}$.

Figure 16.- Type III interference on a 2.54-cm-diameter hemisphere at Mach 20.2 in helium. $\theta_{\rm i} = 10^{\rm o}; \ \rm N_{\rm Re, \infty}/m = 9.7 \times 10^{\rm 6}.$



(b) Center-line heat-transfer distribution.



(c) Off-center-line peak heating.





TINO	ϕ , deg	150 136 127 127 121 117
	$^{\theta}$, deg	34 34 51 56 63
4	h/h _{stag}	6.84 4.59 2.58 0.88 0.59
	Isotherm	
	Model	Steel
	Run	ທ≁ຍ>
	Test	373





Figure 17.- Type III interference on a 2.54-cm-diameter hemisphere at Mach 20.2 in helium. $\theta_i = 15^0$; $N_{Re,\infty}/m = 9.8 \times 10^6$.













(b) Center-line heat-transfer distribution.





Figure 18.- Concluded.



L-73-6873

(a) Schlieren photograph.







(a) Heating contours. $h_{stag} = 306.42 \text{ W/m}^2 \text{-K.}$

Figure 20.- Type IV interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 5^0$; $N_{Re, \infty}/m = 7.6 \times 10^6$.









Figure 20.- Concluded.



(a) Heating contours. $h_{stag} = 306.42 \text{ W/m}^2\text{-K}$.

Figure 21. - Type IV interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 9.8^{\circ}$; $N_{Re,\infty}/m = 8.0 \times 10^{6}$.



(b) Center-line heat-transfer distribution.







Figure 22. - Type IV interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_i = 14.8^\circ$; $N_{Re,\infty}/m = 7.5 \times 10^6$.













(a) Heating contours. $h_{stag} = 306.42 \text{ W/m}^2 \text{-K}.$

Figure 23. - Type IV interference on a 5.08-cm-diameter hemisphere at Mach 5.94 in air. $\theta_{i} = 14.8^{\circ}; N_{Re,\infty}/m = 7.7 \times 10^{6}.$













 $\theta_{i} = 19.7^{\circ}; N_{Re,\infty}/m = 7.6 \times 10^{6}.$



(b) Center-line heat-transfer distribution.











(b) Center-line heat-transfer distribution.



Figure 25.- Concluded.





Figure 26.- Type IV interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_1 = 5^0$; $N_{Re,\infty}/m = 25.6 \times 10^6$.



(b) Center-line heat-transfer distribution.









 $\theta_{i} = 9.8^{\circ}; N_{Re,\infty}/m = 25.0 \times 10^{6}.$











Figure 28.- Type IV interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_{i} = 14.8^{\circ}; N_{Re,\infty}/m = 25.7 \times 10^{7}.$











Figure 29.- Type IV interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_{\rm i} = 19.8^{\rm o}; \ {\rm N}_{{\rm Re},\infty}/{\rm m} = 25.6 \times 10^{\rm 6}.$











Figure 30.- Type IV interference on a 5.08-cm-diameter hemisphere at Mach 6 in air. $\theta_{i} = 24.8^{\circ}; N_{Re,\infty}/m = 25.4 \times 10^{6}.$



(b) Center-line heat-transfer distribution.









Figure 31.- Shock pattern for a type IVa interference on the cylindrical fin at Mach 6 in air. $\theta_i = 19.9^\circ$; $\lambda = 0^\circ$; $N_{\text{Re},\infty}/\text{m} = 25.8 \times 10^6$.



(a) Heating contours. $h_{cyl} = 367.70 \text{ W/m}^2\text{-K}$. Figure 32.- Type IVa interference on a fin at Mach 5.94 in air. $\theta_i = 9.9^{\circ}; \ \lambda = 0^{\circ}; \ N_{\text{Re},\infty}/\text{m} = 8.1 \times 10^{\circ}$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.57.

Figure 32.- Concluded.


(a) Heating contours. $h_{cyl} = 347.28 \text{ W/m}^2\text{-K}$. Figure 33.- Type IVa interference on a fin at Mach 5.94 in air. $\theta_i = 14.9^{\circ}; \ \lambda = 0^{\circ}; \ N_{\text{Re},\infty}/\text{m} = 7.7 \times 10^{6}$.











(a) Heating contours. $h_{cyl} = 367.70 \text{ W/m}^2\text{-K}$. Figure 34.- Type IVa interference on a fin at Mach 5.94 in air. $\theta_i = 19.9^{\circ}; \ \lambda = 0^{\circ}; \ N_{\text{Re},\infty}/\text{m} = 7.9 \times 10^{\circ}$.









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(a) Heating contours. $h_{cyl} = 653.70 \text{ W/m}^2 \text{-K.}$ Figure 35.- Type IVa interference on a fin at Mach 6 in air. $\theta_i = 9.9^\circ$; $\lambda = 0^\circ$; $N_{\text{Re},\infty}/\text{m} = 26.2 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.57.

Figure 35.- Concluded.



(a) Heating contours. $h_{cyl} = 633.27 \text{ W/m}^2\text{-K}$. Figure 36.- Type IVa interference on a fin at Mach 6 in air. $\theta_i = 14.9^\circ$; $\lambda = 0^\circ$; $N_{\text{Re},\infty}/\text{m} = 25.4 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.59. Figure 36.- Concluded.



(a) Heating contours. $h_{cyl} = 653.70 \text{ W/m}^2\text{-K}$. Figure 37.- Type IVa interference on a fin at Mach 6 in air. $\theta_i = 19.9^\circ$; $\lambda = 0^\circ$; $N_{\text{Re},\infty}/\text{m} = 25.8 \times 10^6$.











(a) Schlieren photograph.

Figure 38.- Shock and oil-flow patterns for a type V interference on a fin at Mach 6 in air. $\theta_i = 19.9^{\circ}$; $\lambda = 24.8^{\circ}$; $N_{\text{Re},\infty}/\text{m} = 25.9 \times 10^{6}$.







(a) Heating contours. $h_{cyl} = 306.42 \text{ W/m}^2\text{-K}$. Figure 39.- Type V interference on a fin at Mach 5.94 in air. $\theta_i = 9.9^\circ$; $\lambda = 25^\circ$; $N_{\text{Re},\infty}/\text{m} = 7.2 \times 10^6$.



(b) Center-line heat-transfer distribution.





Figure 39.- Concluded.



(a) Heating contours. $h_{cyl} = 306.42 \text{ W/m}^2\text{-K}$. Figure 40.- Type V interference on a fin at Mach 5.94 in air. $\theta_i = 15^{\circ}; \ \lambda = 25^{\circ}; \ N_{\text{Re},\infty}/\text{m} = 7.3 \times 10^{6}$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.39.

Figure 40.- Concluded.



(a) Heating contours. $h_{cyl} = 306.42 \text{ W/m}^2\text{-K}$. Figure 41.- Type V interference on a fin at Mach 5.94 in air. $\theta_i = 19.8^\circ$; $\lambda = 24.9^\circ$; $N_{\text{Re},\infty}/\text{m} = 7.6 \times 10^6$.









Figure 41.- Concluded.



(a) Heating contours. $h_{cyl} = 592.41 \text{ W/m}^2\text{-K}$. Figure 42.- Type V interference on a fin at Mach 6 in air. $\theta_i = 9.9^\circ$; $\lambda = 25^\circ$; $N_{\text{Re},\infty}/\text{m} = 25.7 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.34.

Figure 42.- Concluded.



(a) Heating contours. $h_{cyl} = 692.41 \text{ W/m}^2\text{-K}$. Figure 43.- Type V interference on a fin at Mach 6 in air. $\theta_i = 14.9^\circ$; $\lambda = 24.9^\circ$; $N_{\text{Re},\infty}/\text{m} = 26.2 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.37.

Figure 43.- Concluded,



(a) Heating contours. $h_{cyl} = 592.41 \text{ W/m}^2\text{-K}$. Figure 44.- Type V interference on a fin at Mach 6 in air. $\theta_i = 19.8^\circ$; $\lambda = 24.9^\circ$; $N_{\text{Re},\infty}/\text{m} = 25.9 \times 10^6$.



(b) Center-line heat-transfer distribution.



(c) Off-center-line heating. X = 0.33.

Figure 44.- Concluded.

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