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# THE EFFECT OF WIND-TUNNEL WALL INTERFERENCE ON THE PERFORMANCE OF A FAN-IN-WING VTOL MODEL

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#### 16. Abstract

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The present investigation shows that the rules for choosing model sizes to produce negligible wall effects, as given by Cook and Hickey in NASA SP-116, are considerably in error and permit the use of excessively large models. Even simple momentum theory appears to yield more nearly correct performance estimates in transition flight than uncorrected wind-tunnel data when the model span approaches one-half of the tunnel width. The "fan-induced" lift indicated by a number of previous studies in which the model was of similar relative span appears to be largely the result of wall interference that was not accounted for in reducing the data.

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## THE EFFECT OF WIND-TUNNEL WALL INTERFERENCE ON THE PERFORMANCE OF A FAN-IN-WING VTOL MODEL

By Harry H. Heyson Langley Research Center

#### SUMMARY

A fan-in-wing model with a 1.07-m (42-in.) span was tested in seven different test sections with cross-sectional areas ranging from  $2.2 \text{ m}^2$  to  $265 \text{ m}^2$  (24 ft<sup>2</sup> to  $2857 \text{ ft}^2$ ). The data from the different test sections are compared both with and without correction for wall interference. The results demonstrate that extreme care must be used in interpreting uncorrected VTOL data since the wall interference may be so large as to invalidate even trends in the data. The wall interference is particularly large at the tail, a result which is in agreement with recently published comparisons of flight and large-scale wind-tunnel data (NASA CR-2135) for a propeller-driven deflected-slipstream configuration. The data of the present investigation verify the wall-interference theory of NASA TR R-124 even under conditions of extreme interference. A method given by Tyler and Williamson in AGARD CP-91-71 yields reasonable estimates for the onset of Rae's minimum-speed limit.

The present investigation shows that the rules for choosing model sizes to produce negligible wall effects, as given by Cook and Hickey in NASA SP-116, are considerably in error and permit the use of excessively large models. Even simple momentum theory appears to yield more nearly correct performance estimates in transition flight than uncorrected wind-tunnel data when the model span approaches one-half of the tunnel width. The ''fan-induced'' lift indicated by a number of previous studies in which the model was of similar relative span appears to be largely the result of wall interference that was not accounted for in reducing the data.

#### INTRODUCTION

Despite considerable theoretical study (e.g., refs. 1 to 6) of wind-tunnel interference for VTOL and STOL aircraft, it is not a general practice to correct all such data for wall effects. This failure to correct is due in part to conflicting reports of the efficacy of such corrections (e.g., refs. 7 and 8); it is due in part to some confusion between the effects of corrections and of the minimum-speed limits proposed by Rae (ref. 9); and finally, it is due in part to the rather considerable effort required to program corrections for data reduction when the programing may be significantly different for different types of model.

The magnitude of wall interference and the extent to which the data may be corrected for such interference become of paramount importance in the design of a new wind tunnel because the required test-section dimensions must be selected so that the data from the tunnel will be representative of the model operating in free air. Consequently, in connection with the design of a new full-scale subsonic wind tunnel (refs. 10 to 12), a major experimental study of wall effects was undertaken. This program was a joint effort of the Langley and Ames Research Centers of NASA. The model chosen was a simplified fan-in-wing aircraft differing from the model of reference 13 only in the addition of a large tail and a slight increase in wing-section thickness ratio. This model was chosen because it is considerably more complex, from a wall-effects viewpoint, than the models which have heretofore been used in V/STOL wall-effects investigations (e.g., refs. 8 and 14 to 18), and because the general type of configuration was representative of the configurations of reference 19 and therefore could provide an evaluation of the conclusions of reference 19. The model was tested, with and without smaller test-section inserts, in a 2.13- by 3.05-m (7- by 10-ft) wind tunnel at Ames as well as briefly in the 12.2- by 24.4-m (40- by 80-ft) wind tunnel at Ames. It was also tested with and without a testsection insert in the 9.14- by 18.3-m (30- by 60-ft) Langley full-scale wind tunnel.

The immediate objectives of the test program were twofold: First, since the tests conducted by Rae (ref. 9), which defined the problem of minimum-speed limits, were all conducted using relatively large single rotors, it was desired to examine the differences in these limits which might be caused by distributing much of the lift into two discrete highly loaded fans. Second, it was desired to obtain some experimental indication of the magnitude and correctability of the wall interference engendered by a model of this type. The approach used was to correct all the data to the maximum extent possible, and then to examine the differences in the data (both corrected and uncorrected) from tests under controlled conditions in the various test sections.

Examination of the data indicates that the first of the aforementioned objectives was only partially achieved. Some insight was obtained into the relative magnitude of the minimum-speed limits in different size test sections; however, the results are not adequate to distinguish any order of relative merit between the different cross-sectional shapes of the test sections. The second objective was met in a more satisfying manner. The data presented herein demonstrate the extent to which V/STOL data from different wind tunnels can be correlated, even in the face of extraordinarily large wall interference.

This wall-interference study is of particular interest since it demonstrates that, for V/STOL flight conditions, the interference may be of such magnitude that even the trends of the data may be incorrect. An example of correlation of wind-tunnel and flight-test data

for an entirely different aircraft (ref. 20) is presented to demonstrate that this observation is also true for configurations totally different from that of the present investigation. Comparisons are made between the present work and previously published theoretically and experimentally chosen limits for V/STOL wind-tunnel testing (refs. 6 and 7).

In correcting wind-tunnel data of the nature of those presented herein, the biggest problem is the lack of uniformity over the model of the wall-induced interference (ref. 21). Some compensation must be made for the varying effective angles of attack and dynamic pressures over the different components of the model. Thus, in order to correct the data in a complete manner, it is necessary to have at least a rudimentary theoretical treatment of the performance of each component as affected by changes in velocity and angle of attack. In the present case, a simple momentum theory for the lifting fan in cross flow was used. This theoretical treatment, based largely on reference 22, is presented in its entirety in a separate paper (ref. 23). Throughout the present paper, the theoretical predictions of reference 23 are compared with the measured model performance as obtained both with and without wall interference.

Reference 13 has noted that a vortex-density correction is needed in applying the theory of references 2 and 3 to the correction of data obtained for fan-supported models. A justification of this vortex-density correction is presented in appendix A. A sample of the FORTRAN programs used in correcting the data obtained in the present investigation is presented as appendix B.

#### SYMBOLS

Because of the limited font of characters available in the automatic figure-plotting equipment, certain symbols may vary between the text and the figures. Where this variation occurs, the symbol used in the figures is shown parenthetically at the beginning of the definition.

A aspect ratio,  $b^2/S_W$ 

A<sub>I.</sub> momentum area of VTOL elements

 $A_{M}$  momentum area of wing,  $\frac{\pi}{4} b^2$ 

A<sub>T</sub> cross-sectional area of wind-tunnel test section

b span of wing

 $C_D$  drag coefficient,  $D/qS_W$  or  $D/q_cS_W$ 

CL	lift coefficient, $L/qS_W$ or $L/q_cS_W$
$\boldsymbol{c}_{\mathbf{L},j}$	lift coefficient based on fan-area and fan-exit dynamic pressure, $\frac{L}{\frac{1}{2}\rho V_i^2 S_F}$
$c_{L_{\alpha}}$	lift-curve slope of wing, $\partial C_L / \partial \alpha$ , per rad
$c_{N,T}$	tail normal-force coefficient, $N_T/qS_T$ or $N_T/q_cS_T$
с	chord
D	drag due to lift, total drag less drag at $\alpha = 0^{\circ}$
$\mathbf{D}_{\mathbf{E}}$	sum of D and D <sub>se</sub>
D <sub>se</sub>	drag equivalent to shaft power, $P_s/V$
d <sub>e</sub>	equivalent fan diameter, $\sqrt{4S_F/\pi}$
h	height of fan exit above test-section floor
L	lift
NT	tail normal force
Ps	shaft power
q	(Q) dynamic pressure of test-section flow, $~{ m  ho\over 2}~{ m V}^2$
q <sub>c</sub>	$\left( Q_{\mathbf{C}}  ight)$ corrected dynamic pressure at wing
${}^{q}\mathbf{F}$	$\left( \mathbf{Q}_{\mathbf{F}} \right)$ corrected dynamic pressure at fans
$^{q}T$	$\left( \mathbf{Q}_{\mathbf{T}} ight)$ corrected dynamic pressure at tail
R	body radius
$s_{F}$	fan area
$s_{T}$	tail area
$s_W$	wing area
4	

т <sub>s</sub>	static thrust
v	test-section, or forward, velocity
$\mathbf{v_j}$	fan-exit velocity in static thrust, $\sqrt{T_S/\rho S_F}$
w <sub>0</sub>	vertical induced velocity in forward flight, positive upward
α	angle of attack, angle between relative wind axis and longitudinal axis of model, positive nose up, deg
ΔD	increment in fan external drag resulting from changes in $\alpha$ and V (see eq. (24))
∆i <sub>F</sub>	difference in $\Delta \alpha$ at wing and fans, $(\Delta \alpha)_{\rm F} - (\Delta \alpha)_{\rm W}$ , deg except rad in equations (24) and (25)
$\Delta i_{\mathrm{T}}$	difference in $\Delta \alpha$ at wing and tail, $(\Delta \alpha)_{\mathrm{T}}^{-} - (\Delta \alpha)_{\mathrm{W}}^{-}$ , deg
$\Delta L$	increment in fan lift resulting from changes in $lpha$ and V (see eq. (25))
$\Delta L_i$	so-called "fan-induced" lift, total lift less the independent lifts of the fans and the model with the fans covered
Δu	longitudinal component of wall-induced interference velocity, positive rearward
Δw	vertical component of wall-induced interference, positive upward
Δα	change in angle of attack caused by wall interference, referred to wing angle of attack unless otherwise subscripted, positive nose up, deg except rad in equations (13), (15), and (16)
$\delta_{e}$	elevator deflection angle, positive trailing-edge down, deg
<sup>δ</sup> u,D	interference factor for longitudinal interference due to drag
$^{\delta}$ u,L	interference factor for longitudinal interference due to lift

<sup>δ</sup> w,D	interference factor for vertical interference due to drag
$^{\delta}\!$	interference factor for vertical interference due to lift
e	downwash angle at tail, positive downward, deg
ρ	mass density of air
x	wake skew angle, angle measured from vertical axis of test section to center of wake, positive rearward, deg
x <sub>e</sub>	effective wake skew angle, deg
Subscripts:	
С	corrected
F	fans
Т	tail
u	uncorrected
W	wing

#### APPARATUS AND TESTS

#### Model

The model used in this investigation is shown in figure 1 and pertinent dimensions are further detailed in tables I and II. The model consisted of a symmetrical streamline body 2.13 m (84 in.) long with a maximum diameter of 0.2 m (8 in.). A symmetrical tapered wing with a 1.07-m (42-in.) span was mounted at the midpoint of the body. The airfoil section at the wing tip was NACA 16-015 and the section increased in thickness to NACA 16-017 at the centerline of the body; straight-line fairings were used between these two stations.

Two commercially available 0.2-m (8-in.) tip-turbine-driven fans were mounted on centers spaced 0.56 m (22 in.) apart at the midchord position of the wing. The inlets to these fans were of the simple bellmouth type obtained by providing a reasonable radius at the intersection of the fan duct and the upper wing surface.

A slab tail with a 0.76-m (30-in.) span and a 0.32-m (12.5-in.) chord was mounted symmetrically so that its trailing edge was coincident with the rearmost end of the fuse-lage. This tail was installed during all tests for which the data are presented herein.

The model was mounted on a pivot at the midpoint of the fuselage and 6.67 cm  $\left(2\frac{5}{8} \text{ in.}\right)$  below the centerline of the model. A linear actuator, installed between the mounting strut and a point farther rearward on the model, provided remote control of angle of attack.

#### Model Instrumentation

The model was designed to be operated on the normal external mechanical balances of the wind tunnels; thus, it was not necessary to provide a sting balance for measurement of the overall forces and moments. The mechanical balances involved are all of the simple platform type and have relatively poor resolution of moments for model forces of the magnitude encountered during the tests. The expected balance accuracy, together with some anticipated difficulty in setting precisely the same powered-lift flight conditions, precluded the possibility of obtaining accurate measurements of the effect of the tail on the moments by comparison of tail-on and tail-off tests. Consequently, the tail and tailcone were mounted on the body by means of a commercial 1.9-cm-diameter (3/4-in.) six-component strain-gage balance. The primary measurement desired was the tail normal force, and the balance had a maximum load capability of 445 N (100 lb) for this component of force.

Numerous pressure and temperature transducers were provided in the independent pneumatic systems powering the two fans. The only measurement pertinent to the final results was the rotational speeds of the fans, for which magnetic pickups were provided in the fan casings. Considerable difficulty was experienced with this system because of 60-Hz pickup during the tests. The initial series of tests was actually conducted by setting the fan rotational speeds with a stroboscopic tachometer. For subsequent series of tests, a discriminator circuit was constructed to minimize the pickup problem, and magnetic tachometers with higher output were used. This aspect of the testing is discussed more completely in a later section of this paper.

Angle of attack was measured by an accelerometer-type transducer mounted within the model, except in the 12.2- by 24.4-m (40- by 80-ft) tunnel, where a selsyn indicator mounted at the actuator strut was used. Differences in the data-acquisition systems of the other two tunnels required the use of different transducers in each tunnel. In the smaller two tunnels, the accuracy of the overall system was approximately the same. In the largest tunnel the overall accuracy was somewhat less.

#### Wind Tunnels

2.13- by 3.05-m (7- by 10-ft) tunnel.- The smallest of the three tunnels used in this investigation was the Ames 7-foot by 10-foot Subsonic Wind Tunnel No. 2. This tunnel is described on pages 1-32 and 1-33 of reference 24. The model was mounted in the tunnel on a single unfaired strut (fig. 2). The pivot point at the top of the strut was on the centerline of the tunnel; thus, the model aerodynamic center was slightly above the tunnel centerline.

Air was supplied to the fans by means of two 5.1-cm-diameter (2-in.) hoses which were dressed closely to the front and back of the strut by means of guide rings. Below the floor of the tunnel, and above the balance frame, some slack was provided in the air lines to provide for the motion which occurred as a result of changes in angle of attack. An elaborate trapeze connection was provided between the balance frame and the main air supply. Tests of this system under pressure, with the model hoses blocked, indicated no measurable effect on the loads as seen by the balance.

Instrument leads were taped tightly to the sides of the strut and were connected to the data-acquisition system by means of a large hanging loop of wiring below the balance frame. The gap between the strut and the floor was closed to a minimum by specially trimmed sheet metal screwed to the floor of the tunnel.

Tunnel airspeed was measured by means of a pitot-static tube mounted from the ceiling of the tunnel. Corrections for position error are discussed in a later section of this paper. The tube was mounted 0.254 m (10 in.) below the tunnel ceiling, and the static-pressure holes of the tube were 1.33 m (52.5 in.) ahead of the model pivot point. (At zero angle of attack this location is 0.267 m (10.5 in.) ahead of the nose of the model.) The dynamic pressure measured by this tube was passed through a pressure transducer and then to both the data-acquisition system and a digital indicator, which was used as a speed reference during the tests.

Since this tunnel has continuous speed control, it had been hoped to maintain a close control over tunnel speed during each set of data points; however, this was not possible in practice. At high speed, blockage of the tunnel caused by the powerful variation of fan momentum drag with speed and angle of attack resulted in excessive time losses in attempting to set the tunnel speed precisely. At the lowest speeds, recirculation effects became so severe that the tunnel speed was found to lope; the pulsations in the tunnel flow were obvious even to the ear. Consequently, the tunnel speed was taken as the average of three readings, each of which in turn was averaged over a time of 1.25 seconds.

Tuft boards were placed on the floor of the tunnel for visual observations of the flow when recirculation began (refs. 9 and 13).

Inserts in 2.13- by 3.05-m (7- by 10-ft) tunnel.- In order to simulate still smaller test-section sizes, the insert technique of references 9, 13, 18, and 25 was used. Two rectangular test sections were simulated by means of two vertical walls (one of plywood and the other of transparent plastic to permit observation of the tufts) between which two horizontal surfaces were suspended to simulate the floor and ceiling of the small test sections. The entire assembly in each case was generously braced to insure stability and dimensional constancy.

The first of these simulated test sections had a width of 1.83 m (6 ft) and a height of 1.22 m (4 ft) providing a width-height ratio of 1.5. The second test section had a width of 2.24 m (88 in.) and a height of 1.12 m (44 in.) providing a width-height ratio of 2.0. A third test section was obtained by fitting the 2:1 insert internally with sheet metal ends which were rolled to a semicircular cross section; thus a flat-oval test section having a width-height ratio of 2.0 was provided. All these inserts were 3.66 m (12 ft) long. The model was centered longitudinally within each insert. Photographs of the model mounted in these test sections are given in figure 3.

The cross-sectional areas of the 1.5:1 rectangular insert and the 2:1 flat-oval insert were essentially identical with each other at 2.23 m<sup>2</sup> (24 ft<sup>2</sup>). The 2:1 rectangular test section, with a cross-sectional area of 2.50 m<sup>2</sup> (26.89 ft<sup>2</sup>), was approximately 12 percent larger in cross-sectional area. The choice of these sizes was not accidental. The dimensions of the 2:1 flat-oval test section were specifically chosen to represent the wing-span to tunnel-width ratios used in several Ames full-scale tunnel tests of fan and fan-in-wing models (e.g., refs. 26 to 32).

Speed measurement in the inserts was by means of the same pitot-static tube used in the basic wind tunnel. The longitudinal location of this pitot-static tube was constant, and in each insert the vertical location was adjusted so that the tube was 25.4 cm (10 in.) from the insert ceiling.

Each insert was generously tufted for visual flow observations; however, the curved sheet metal walls of the flat-oval section severely limited the field of view.

Compressed air to drive the fans was supplied from a large high-pressure storage tank. This air supply was adequate to drive the fans at nominal rotational speed of 10 000 and 12 000 rpm.

No evidence of any flow inclination was found in the data. Thus, the wind-tunnel stream angle was zero irrespective of the presence or absence of the inserts. Under these conditions, it was possible to set the model angle of attack directly to the desired values throughout the tests.

In any wind-tunnel wall-effects investigation the relative sizes of the test sections are of vital importance. A sketch illustrating the relative sizes is presented in figure 4. <u>12.2- by 24.4-m (40- by 80-ft) tunnel.</u>- In order to obtain conditions essentially free of wall constraints, the model was tested briefly in the 12.2- by 24.4-m (40- by 80-ft) Ames full-scale tunnel. This wind tunnel is described in reference 24. The external balance of this tunnel was not designed to measure loads as small as those which were produced by the present model. In order to gain some increase in precision, the model was mounted with its span vertical (fig. 5) so that the lift could be measured by the side-force scales, which have a greater sensitivity than the lift scales.

The model was mounted on the same strut that was used in the tests conducted in the previously described tunnel; however, the mechanical arrangements did not allow the air hoses to be dressed closely to the strut. Instead, angled fittings were provided at the model and at the base of the strut. The required motion of the hoses with angle of attack was obtained by bending the supply hoses. As may be seen in figure 5, the resulting installation was substantially less clean than the installation in the smaller Ames tunnel.

A different pitch actuator was installed, and angle of attack was measured as a function of the actuator extension. The least division of the angle-of-attack indicator was  $0.25^{\circ}$ , and this reading was manually inserted into the data-acquisition system. A straingage balance was inserted into the actuator linkage in order to measure pitching moments; however, these measurements were invalidated by the omission of a static tare accounting for the moments imposed on the model by bending the air-supply hoses.

The discriminator circuit and large magnetic pickups were used in measuring the fan rotational speeds. This arrangement substantially reduced the amount of 60-Hz noise accepted by the counters; however, the static thrust measurements indicated that some spurious counts were still obtained. The counters were not connected directly to the data-acquisition system; the readings were manually inserted into the system.

The air supply in this tunnel was not adequate for continuous operation of the fans at 12 000 rpm. Therefore, the tests were conducted at 10 000 rpm and at the maximum available rotational speed, which tended to be on the order of 11 500 rpm.

Three different systems of tunnel flow-velocity measurement were employed. There were substantial disagreements in the data measured by the three systems. The staff of the tunnel provided their best estimates of the true velocities, and these values were punched into the data cards at a later date.

The tail-balance readings were recorded on a second data-acquisition system. This second system proved troublesome, with obviously mispunched cards being obtained even while recording zeros. It is believed that reasonably accurate readings were obtained during initial tests with the fans covered; however, the data obtained with the fans operating were so different from the data obtained in all the other test sections that they were rejected. Approximately half of the powered phase of the testing was complete when this system failed completely and no further tail data were obtained. Initial tests with the fans covered indicated a very large stream-angle correction. Inasmuch as this tunnel is not equipped with flow-survey apparatus, it was not possible to obtain direct measurements of the flow inclination. Subsequent tests, with the model removed and the air hoses taped tangentially to the top of the strut, indicated that a large lift tare was also present in the data. No complete sequence of tare tests were conducted to obtain the precise magnitude of the tare. In analyzing the data, the stream angle and the lift and drag tares were obtained by finding those values that yielded the same performance as in all the other test sections when the fans were covered. These values were assumed to be unaltered by fan operation.

The stream angle obtained in the foregoing manner was significantly different from that presumed to exist during the conduct of the tests. Consequently, the maximum true angle of attack obtained in this tunnel was several degrees less than that obtained in the other tunnels.

The tests in this tunnel were conducted under the direction of Kenneth W. Mort, of the NASA Ames Research Center.

<u>9.14- by 18.3-m (30- by 60-ft) tunnel.</u> The deficiencies inherent in the tests conducted in the Ames 12.2- by 24.4-m (40- by 80-ft) tunnel were such that the resulting data were too ambiguous to be accepted as defining the free-air characteristics of the model. Consequently, more complete tests were conducted in the 9.14- by 18.3-m (30by 60-ft) Langley full-scale tunnel. This tunnel is described in reference 33. Some later information on the wind tunnel is presented in references 24 and 34.

The ground board normally used in the Langley full-scale tunnel was in place during these tests. The upper surface of this ground board is approximately 0.61 m (2 ft) above the lower edge of the jet boundary and thus reduces the cross-sectional area of the test section to 141.8 m<sup>2</sup> (1527 ft<sup>2</sup>). By comparison, the model is very small; its wing area is less than one-half of 1 percent of the test-section cross-sectional area.

Because of the size of this tunnel, it was necessary to prepare a new mounting strut for the model. The new strut was designed so that the model was mounted vertically on the centerline of the active region of the tunnel (4.26 m (14 ft) above the ground board). As nearly as possible, the uppermost 1.07 m (3.5 ft) of the strut was identical with the strut used in the smaller tunnel. The end fitting on this strut, the hoses and their arrangement, and the angle-of-attack actuator were the same as those which were used in the smallest tunnel. A close-fitting fairing was installed around the strut starting 1.07 m (3.5 ft) below the model and continuing downward to meet the ground board. All hoses and electrical leads were dressed to the strut in, as closely as possible, the identical manner in which they were installed in the smaller tunnel. Photographs of the model installed in the tunnel are presented in figure 6.

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The air-pressure lines were brought across the balance in a trapeze arrangement. Tests conducted under pressure with the hoses blocked at the model indicated no effect on the balance readings. The instrument leads were carried across the balance by means of a large hanging loop.

Prior to mounting the model on the strut, the region occupied by the model was surveyed with a pitot-static-pitch-yaw head. The dynamic pressure measured by this survey instrument was used to calibrate the velocity at the model as a function of static depression in the tunnel settling chamber; this static depression in turn was used to determine the tunnel velocity during the tests. The survey also disclosed the presence of a significant stream angle (approximately  $0.7^{\circ}$ ) at the model location. The presence of this stream angle was confirmed later by the raw data from the symmetrical model when it was tested with the fans covered. The effects of this stream angle have been removed from all the data presented herein.

The Langley full-scale tunnel does not have continuous speed control throughout the velocity range covered in these tests; instead, it has some 24 discrete power settings, or "points." A number of these points appropriate to the prior tests in the smaller tunnel were selected. The actual velocity presented herein was determined from the average of no fewer than 10 samplings, spaced 1 second apart, of the static pressure.

In order to accommodate the different data-acquisition systems in this tunnel, it was necessary to use a different type of angle-of-attack transducer within the model. Again, the values presented result from the average of no fewer than 10 samplings of the transducer output.

Insert in 9.14- by 18.3-m (30- by 60-ft) tunnel.- It was desired to insure continuity of the test results between the tests conducted in the two wind-tunnel facilities. Consequently, a 2.13- by 3.05-m (7- by 10-ft) insert, 6.4 m (21 ft) long, was built up around the model in the Langley tunnel without disturbing the mounted model on the strut. The insert was fitted with a simple 15.2-cm-diameter (6-in.) semicircular sheet metal bell-mouth inlet to discourage separation of the flow at the inlet.

The insert was constructed of 1.9-cm-thick (3/4-in.) plywood and was rigidly braced by angle iron to insure dimensional stability during the tests. It was supported by pipe columns and cable bracing so that the model pivot point was on the centerline of the insert, that is, in the same location as in the tests at the Ames Research Center, and so that it was centered longitudinally on the model. Photographs of this installation are presented in figure 7.

Within the insert, the gap at the floor of the tunnel was reduced to minimal size by means of closely trimmed sheet metal plates screwed to the floor. The fairing around the lower portion of the strut was sealed to the exterior of the insert.

The flow velocity within the insert was measured from the average of four sets of total- and static-pressure measurements. The probes for these measurements were located 45.7 cm (18 in.) behind the leading edge of the insert and 30.5 cm (12 in.) inward from the walls of the insert. Since no divergence was built into the insert, a small correction (approximately 4 percent) was made to the velocity in order to account for the difference in boundary-layer displacement thickness between the probe and the model locations.

It was not possible to survey the flow within the insert walls with the existing equipment in the Langley full-scale tunnel. Stream angle was determined by finding the angle which was required to reduce the lift of the symmetrical model with the fans covered to zero at an angle of attack of zero. In this regard, a root-mean-square average of such angles for all tunnel speeds was used. The resulting stream angle was approximately  $-0.2^{\circ}$  and is accounted for in all data presented herein.

Air to power the fans was provided by a permanent compressor in the tunnel. It would have been desirable to maintain the same rotational speeds as were used in the earlier Ames tests. Unfortunately, the compressor proved inadequate in capacity for continuous operation at 12 000 rpm. Consequently, the tests at Langley were conducted at lower rotational speeds, 8000 and 10 000 rpm, which overlapped those in the other test sections.

All the tests in the Langley full-scale tunnel, as well as all the tests in the Ames 7-foot by 10-foot Subsonic Wind-Tunnel No. 2, were conducted under the personal supervision and direction of Frank A. Lazzeroni, of the U.S. Army Air Mobility R&D Laboratory, Ames Directorate.

#### Procedure

The same test procedure was used in all the tunnels and test-section inserts. First, the fans were started and brought to the required rotational speed. Generally, static thrust was measured, usually throughout the same angle-of-attack range as in the subsequent forward-flight tests. Then, the tunnel was started and brought to the desired velocity. Data were recorded in the following angle-of-attack sequence:  $0^{\circ}$ ,  $-10^{\circ}$ ,  $-5^{\circ}$ ,  $0^{\circ}$ ,  $5^{\circ}$ ,  $10^{\circ}$ , and  $16^{\circ}$ . The tunnel speed was then altered to the next desired speed. Although the angle-of-attack sequence was constant, the progression of tunnel speeds was not constant. The sequence of speeds was often reversed so that the test commenced with the highest speed and ended with static thrust. Even more erratic velocity sequences were used in the Langley tunnel, where, because of a pole change in the motor-control system at approximately 48 knots, it was often more convenient to descend in velocity to that speed, drop to the smallest velocity, and then increase tunnel speed to obtain velocities up to 48 knots.

Data recording procedures differed in the three tunnel facilities because of differences in the data-acquisition systems. In the smallest tunnel, the data were obtained as three sets of time-averaged data (over a 1.25-second period) and were punched on cards for off-line reduction. In this tunnel, angle of attack was set manually and "dialed" into the data system manually. A similar set of two independent systems was used in the largest tunnel. In the Langley full-scale tunnel the data were obtained as at least 10 (and often 25) sets of samplings (with essentially no time averaging on each of the sets); the data were stored on magnetic tape for off-line processing. In this latter system, angle of attack was included as one of the directly recorded variables.

In each case, essentially no data were available during the actual testing. While this "blindness" may be a disadvantage during tests of a specific configuration, it is an advantage during tests of the present type because it eliminates any tendency to tinker with the model in order to obtain a preconceived result.

#### **Precision of Measurement**

Detailed examination of the data, together with the known capabilities of the external balances, indicates that the forces should be accurate to within the values shown in the following table:

	Force					
Tunnel facility	Lift		Drag		Tail normal force	
	N	lb	N	lb	N	lb
2.13 by 3.05 m (7 by 10 ft)	±8.9	±2.0	±2.2	±0.5	±4.4	±1.0
9.14 by 18.3 m (30 by 60 ft)	±13.3	±3.0	±4.4	±1.0	±4.4	±1.0
12.2 by 24.4 m (40 by 80 ft)	±22.2	±5.0	±22.2	±5.0	±4.4	±1.0

The values given for the 12.2- by 24.4-m (40- by 80-ft) tunnel include an allowance for the ambiguous nature of the stream angle and the tares.

It will be observed that these accuracies vary percentagewise according to the overall level of forces observed, and further that they will be reflected in the zeros for the data as well as in the data points themselves. As a proportionate point of reference for those figures in which the data have been nondimensionalized with respect to static thrust, it should be noted that the static thrust for the complete model ranges from about 196 N (44 lb) at a nominal speed of 8000 rpm to about 480 N (108 lb) at a nominal speed of 12 000 rpm.

In the Ames full-scale tunnel the dynamic pressure is believed accurate within 5 percent; angle of attack, to within  $2.0^{\circ}$ . In all the other test sections, dynamic pressure is believed accurate to within 1 percent and angle of attack to within  $0.1^{\circ}$ .

All data presented herein have been corrected for the effects of stream angle on the forces and the angle of attack, where such correction is appropriate. Where adequate rotational-speed data were recorded, the quantities  $V_j$  and  $T_S$  used in nondimensionalizing much of the data have been corrected for the actual rotational speed. Forces, where presented directly rather than as coefficients, have been corrected to standard density from the density at which the data were obtained.

Corrections for wall effects are discussed separately at appropriate points in the discussion of the results.

### RESULTS AND DISCUSSION OF DATA FROM MODEL WITH FANS COVERED

#### **Uncorrected Data**

The uncorrected data for the model operating with the fans covered are presented in terms of lift, drag, and tail normal-force coefficients as a function of angle of attack in figure 8. Because of the small loads and the coarse sensitivity of the external balances employed, only the data from the highest dynamic pressure run in each test section are presented.

The strut used to mount the model during these tests was not faired, and a different length of this strut was exposed to the full dynamic pressure of the tunnel in each test section. No series of tare runs was made to determine the tare loads in the data; however, as noted earlier, the mounting arrangements near the model were as identical as possible during most of the tests. Consequently, an amount of drag equal to the entire drag of the model with fans covered at zero angle of attack has been removed from the data in this figure and in all subsequent figures. The resulting values of drag and drag coefficients may be considered to be approximately those due to lift.

The data for each coefficient, as obtained in the 9.14- by 18.3-m (30- by 60-ft) tunnel (where boundary interference is negligible due to the extremely small size of the model compared with the test section), were subjected to least-squares analysis. The resulting expressions for a quartic fit to the data are displayed as a curve on each figure.

It will be observed that even though the model was symmetrical, the data do not quite possess the expected symmetries and antisymmetries with angle of attack. This result is rational, for the rearmost portions of the model at positive angle of attack were immersed in a region of lowered dynamic pressure behind the mounting strut and were free of this region when at negative angle of attack. It is clear that under such conditions the emphasis placed on maintaining the mounting conditions as identical as possible in most of the test sections was entirely justified and necessary.

#### Considerations in Correcting Data

The data of figure 8 contain several types of boundary-induced interference. First, there is solid blockage. This interference is easily evaluated to a sufficient degree of accuracy from the compilation of studies presented in section 6:10 of reference 35. Next, there are the boundary-induced effects due to the presence of the lifting model within the tunnel. In the present case this last-named effect was obtained using the method of references 2 and 3 as implemented by the FORTRAN programs given in reference 36.

The use of references 2 and 3 presents two problems when the theory is applied to a model for which the lift may be zero. First, the momentum theory (refs. 2 and 22) used to obtain the wake skew angle appears to fail when the lift is negative. This difficulty is resolved by calculating the skew angle using the absolute value of the lift in the equations and subsequently choosing the proper quadrant for the wake according to whether the lift is positive or negative. The second problem is that the computer programs of reference 36 are arranged in such a manner that they yield the correct interference factors only when the wake skew angle is greater than  $-90^{\circ}$  and less than or equal to  $90^{\circ}$  (that is, the wake cannot pass upward as it passes rearward). Some rules for treating the calculation by symmetries are presented in reference 5; however, in the present case, where the wake skew angles are only slightly greater than  $90^{\circ}$  (slightly upward), it is more convenient merely to extrapolate from the values calculated for the first quadrant.

The model was somewhat unusual in that the wing was closely coupled to an extraordinarily large tail. Furthermore, the tail had a greater aspect ratio than the wing (2.4 compared with 1.6), and thus would be expected to have a higher lift-curve slope than the wing. Under such conditions, it would be expected that the model would behave more nearly as a tandem-wing system than as a simple wing-tail combination; this expectation is confirmed by the nonlinear character of the lift-curve slope (fig. 8(a)). It is important to consider this tandem-wing-like character of the model in the corrections; that is, the effect of the interference at the tail must be considered not only with respect to tail normal force, but also with respect to the overall lift and drag of the model.

The appropriate interference factors for the wing due to its own presence may be obtained from the FORTRAN program given as appendix B of reference 36. It was assumed that the wing had an elliptic load distribution. Since the quarter chord of the model is displaced from the pivot point, both vertically and longitudinally, these interference factors will be a function of angle of attack by virtue of the different vertical location of the wing within the tunnel at each angle of attack. (See eq. (58) of ref. 3.) The effect of the presence of the tail and its loads on the tail itself is also significant and can be obtained from the same program; it is imperative that the location of the tail as a function of angle of attack be considered. Observe that this effect would have been difficult to consider if it had not been for the use of a tail balance to measure the tail loads. The interference factors at the tail due to the presence of the wing may be obtained from appendix D of reference 36. The interference at the wing due to the presence of the tail could be obtained from the same program (by considering the wing to be a canard tail); however, the rapid decrease of interference with distance upstream from the causative lifting element precludes any significant effect from this source and it may be ignored safely.

In the tests of the model installed in the inserts in the 2.13- by 3.05-m (7- by 10-ft) wind tunnel at the Ames Research Center, one other feature must be considered. This feature is the tunnel velocity measurement by means of a pitot-static tube near the nose of the model. At this location, the pitot-static tube is affected by the direct field of the model (both due to the body shape and to the lifting system) as well as by wall effects caused by the presence of the model. These effects must be evaluated in order to obtain the proper tunnel velocity to use in the interference calculations and in forming the corrected force coefficients.

The solid blockage at the pitot locations is caused primarily by the body because the body is the portion of the model closest to the pitot tube. The blockage is not the same as a classical blockage correction (ref. 35), since the classical blockage calculation is for the model location. In the present analysis, this blockage effect was approximated by setting up a calculation based on the use of a source and a sink to represent a Rankine ovoid (ref. 37, p. 208), and using the technique of reference 38 to obtain the strengths and spacing of these elements to produce an ovoid which, in free-air, has the same length and diameter as the fuselage. The ovoid was then reflected both horizontally and vertically to produce a pattern which represents the boundary conditions at the walls. The interference velocities at the pitot-static tube location can be obtained from this field of elemental sources and sinks.

It is the usual practice in such wall-effects calculations to omit the central image which represents the model itself on the basis that this is the portion to be measured and corrected. In the present calculations, the central image is retained since it is desired to include the direct field of the model as well as the blockage interference. The level of this correction is approximately 1 percent of the free-stream velocity.

The foregoing treatment was used in the present analysis; however, it does contain certain inadequacies. First, the existence of the images representing the boundary conditions at the wall results in an overall velocity at the model which is somewhat greater than the free-air condition for which the source and sink were chosen. Thus, the body for which the interference is obtained will be somewhat more slender than the desired body shape, and, furthermore, it will be slightly different in shape than a Rankine ovoid. These effects probably result in an underestimate of the actual interference. Second, the calculation method will produce a streamlined symmetrical body only at an angle of attack of zero. Thus, it is not possible to examine the effect of the angle of attack of the body on the calculated interference. Such effects would be expected to be large at positive angles of attack, where the nose approached more closely the pitot-static tube location; however, one would expect only smaller changes at negative angles of attack, where the model nose moved farther away from the tube. On an overall average basis, the actual effects of solid blockage at the pitot location probably are underestimated from the omission of angle-of-attack effects.

The wall interference at the pitot location can be obtained from appendix D of reference 36 by considering the pitot to be a canard tail of zero span. Considerable care must be exercised in choosing the tail length and height as a function of model angle of attack in order to retain the correct pitot location. The direct field of the lifting model is obtained by retaining the central image. This is accomplished most simply by altering the subroutine DLTAS given in appendix Q of reference 36. (Delete lines (Q13) and (Q67) through (Q105).)

#### Procedure in Correcting Data

The first step in correcting the data is to divide the loads between the wing and the tail. This is possible only because the present model was fitted with a tail balance. Then, the loads assigned to the wing are used to solve the momentum quartic (ref. 2 or 22) for  $V/w_0$  and the wake skew angle  $\chi$ . In those test sections where the tunnel velocity was measured near the model, the measured tunnel velocity is then corrected, and the corrected value is used to recompute  $V/w_0$  and  $\chi$ .

The value of  $\chi$  obtained in this manner is the momentum-theory value and, as pointed out by reference 8, is not the value that should be used in wall-interference calculations. Because of wake roll-up, the wake vorticity will be concentrated at some higher location in the tunnel given by an effective average value of the skew angle  $\chi_e$ . As discussed in references 6 and 39, the most appropriate choice for a simple wing is that given by

$$\tan \chi_{\rm e} = \frac{\pi^2}{4} \tan \chi \tag{1}$$

The values of the interference factors (previously obtained from ref. 36) are then interpolated to obtain the values corresponding to this value of  $\chi_{e}$ . In this range of skew

angles, the effect of  $\chi_e$  on the interference factors for the effect on any element due to its own presence is small; thus,  $\chi$  may be assumed to be 90<sup>o</sup> when considering the effect of the tail on itself.

At this point the lift and drag of the wing and tail may be used to compute the individual vertical and horizontal increments of interference velocity separately at the wing and at the tail (eqs. (40) to (43) of ref. 2). At the wing, the total values of  $\Delta w$  and  $\Delta u$ are simply the sums of the respective components occasioned by the lift and drag of the wing; however, at the tail, the contributions of both the wing (which are different at the tail than at the wing) and the tail itself must be summed to obtain the total components of interference.

The total values of  $\Delta w$  and  $\Delta u$  are then used to obtain separately, at the wing and at the tail, the values of  $\Delta \alpha$  and  $q_c/q$  by use of equations (48b) and (49b) of reference 2, which are

$$\Delta \alpha = \tan^{-1} \frac{\Delta w/V}{1 + \frac{\Delta u}{V}}$$
(2)

$$\frac{q_{c}}{q} = \left(1 + \frac{\Delta u}{V}\right)^{2} + \left(\frac{\Delta w}{V}\right)^{2}$$
(3)

The values of  $\Delta \alpha$  and  $q_c/q$  at the wing are used as a first correction to the data; however, it is also necessary to account for the differences in  $\Delta \alpha$  and  $q_c/q$  at the wing and the tail. These differences are conveniently expressed as

 $\Delta i_{T} = (\Delta \alpha)_{T} - (\Delta \alpha)_{W}$ (4)

$$\frac{\mathbf{q}_{\mathbf{T}}}{\mathbf{q}_{\mathbf{C}}} = \frac{\left(\frac{\mathbf{q}_{\mathbf{C}}}{\mathbf{q}}\right)_{\mathbf{T}}}{\left(\frac{\mathbf{q}_{\mathbf{C}}}{\mathbf{q}}\right)_{\mathbf{W}}}$$
(5)

It will be noted that the difference in the two values of  $\Delta \alpha$  is effectively a change in tail incidence (ref. 40), and  $q_T/q_c$  is an alteration in the effective dynamic-pressure ratio at the tail. If these effects were not removed from the data, the model would not be aerodynamically equivalent to the model under test. The procedure used herein was to resolve the forces of the <u>entire</u> model around a new effective stream direction given by  $\Delta \alpha$ , where

$$(\alpha)_{c} = (\alpha)_{u} + \Delta \alpha \tag{6}$$

$$\left( \mathbf{C}_{\mathbf{L}} \right)_{\mathbf{C}} = \frac{ \left( \mathbf{C}_{\mathbf{L}} \right)_{\mathbf{u}} \cos \Delta \alpha - \left( \mathbf{C}_{\mathbf{D}} \right)_{\mathbf{u}} \sin \Delta \alpha}{q_{\mathbf{c}}/q}$$
 (7)

$$\left( C_{\mathbf{D}} \right)_{\mathbf{C}} = \frac{ \left( C_{\mathbf{D}} \right)_{\mathbf{u}} \cos \Delta \alpha + \left( C_{\mathbf{L}} \right)_{\mathbf{u}} \sin \Delta \alpha}{q_{\mathbf{c}}/q}$$
 (8)

Next the tail forces as measured were resolved, and then the tail lift and drag were adjusted for  $\Delta i_t$  and  $q_T/q_c$ . This adjustment requires a knowledge of the lift-curve slope of the tail and the free-air dynamic-pressure ratio at the tail. It would be desirable to have test results for the tail in the presence of the body, but without the wing, as a guide in estimating these values. Unfortunately, the construction of the model did not allow for such tests; thus, the lift-curve slope was taken as 0.03 per degree (approximately the value given in fig. 5-5 of ref. 41) and the dynamic-pressure ratio at the tail was rather arbitrarily selected to be 0.9. A small correction to the induced drag of the tail was made to account for the difference in the measured and adjusted lift. A correction to the profile drag would be appropriate; however, insufficient data were available to make such an adjustment. In any event, such a profile-drag adjustment probably would be significant only if the tail were to stall during the test.

The foregoing adjustments were sufficient to provide the corrected values of  $C_{N,T}$ . As a final step, the differences in the lift and drag of the tail were applied as adjustment to the overall lift and drag of the model in order to obtain the final corrected values of  $C_L$  and  $C_D$ .

#### Corrections

The corrections obtained for the model with the fans covered are shown in figure 9. The dynamic pressure ratios differ from unity only by 2 or 3 percent and thus have only a comparatively small effect on the data. However,  $\Delta \alpha$  and  $\Delta i_T$  assume significant proportions in the smallest test sections at large angles of attack.

One common rule of thumb in wind-tunnel testing (e.g., ref. 42) is that  $\Delta \alpha$  should not exceed 2<sup>o</sup>. It is evident from figure 9(a) that  $\Delta \alpha$  has assumed almost this value in the smaller inserts even with the fans covered, and that  $\Delta i_T$  (fig. 9(c)) is well in excess of 2<sup>o</sup>, yielding a total correction angle at the tail on the order of 4<sup>o</sup> (eq. (4)).

#### **Corrected Data**

After the application of corrections, the data for the model with the fans covered appear as shown in figure 10. The solid line shown in figure 10 is again a least-squares quartic faired through the data obtained in the 9.14- by 18.3-m (30- by 60-ft) tunnel.

On the basis of the force accuracies previously given and the dynamic pressures of the tests, the anticipated agreement should be on the order of 0.03 for  $C_L$ , 0.008 for  $C_D$ , and 0.03 for  $C_{N,T}$ . Examination of figure 10 indicates that the correlation between  $C_L$  and  $C_D$  is generally within these limits but that  $C_{N,T}$  appears to be overcorrected to a somewhat greater extent than would be anticipated by a simple examination of the measurement accuracy at the highest values of lift.

One possible cause of the poorer correlation in the case of the tail normal-force coefficient could lie in the required estimates of the tail lift-curve slope and tail dynamic-pressure ratio. These estimates are far more critical in correcting tail normal force than in correcting the overall lift and drag of the model. Another possible cause could be the effect of the wall-induced velocities in relocating the wake to a higher position in the small test sections than in the large tunnel. References 43 and 44 have examined this latter effect theoretically. The maximum ratios of  $C_L/A$ , references 43 and 44 indicate that the correction to the tail should increase when the tail moves with the model. Such an effect would further degrade the present correlation. Finally, the assessment of test accuracy may be excessively optimistic since the value quoted represents only 1 percent of the full normal-force capability of the tail balance and considerable vibration and buffeting of the tail was obvious during the tests.

#### RESULTS AND DISCUSSION OF DATA FROM

#### MODEL WITH FANS OPERATING

#### **Uncorrected Data**

<u>Presentation of data</u>.- In view of the difficulties experienced with the measurement of fan rotational speed during the tests, it is not possible to present the data directly for constant rotational speeds. Instead, the forward velocity has been nondimensionalized with respect to  $V_j$ , which is the fan efflux velocity in static thrust, defined from simple incompressible momentum theory as

$$\mathbf{V_j} = \sqrt{\frac{\mathbf{T_S}}{\rho \mathbf{S_F}}}$$

Similarly, forces are presented only in nondimensional quantities, generally referenced either to the static thrust or to each other. The static thrust used in these nondimensionalizations is always the value obtained in the largest test section used in each series of tests. It is also measured at zero angle of attack. These two conditions insure that the value of static thrust used is the best available from the viewpoint of minimum flow recirculation in the tunnel during the measurement. Indeed, comparisons of the static-thrust data, with and without the insert, in the Langley full-scale tunnel, indicate that any errors caused by recirculation in the 2.13- by 3.05-m (7- by 10-ft) tunnel are within the accuracy of the data.

The uncorrected data in the form of lift, drag, and normal-force coefficients are presented in figures 11 to 13. They are presented in the form of the ratios of  $L/T_S$ ,  $D/T_S$ , and D/L in figures 14 to 16. Finally, the ratio  $L/D_E$  (where  $D_E$  is the sum of the external drag and a drag equivalent to the power supplied to the fans (ref. 23)) is presented in figure 17.

For lift coefficients and the ratios of lift to static thrust, a line on the figures indicates the values which would be obtained if the lift were simply the direct sum of the vertical component of the fan static thrust and the lift of the wing with the fans covered. In a similar manner, the momentum-theory values of all the other parameters (with the exception of  $C_{N,T}$ , for which momentum theory is inappropriate) have been computed by means of the equations of reference 23, and these calculated values are compared with the corrected data.

When examining figure 17 it should be noted that no measurements adequate for the calculation of the power supplied directly to the fans were actually made. Indeed, in view of the small size and fairly low efficiency of the model fan turbines, such measurements of power would be meaningless in relation to flight hardware. Instead, the momentum-theory value of shaft power, as computed from reference 23, has been converted into an effective drag by means of the relationship

$$D_{se} = \frac{P_s}{V}$$
(10)

The values of  $D_{se}$  obtained from equation (10) have been added to both the experimental data and the theoretical curve. Although figure 17 presents no measured data that were not available in the preceding figures, it does serve the purpose of illustrating the effects of wall interference on the efficiency of this type of aircraft in transition. Effect of wall interference on lift.- Figure 11 indicates clearly that at any constant angle of attack, the measured lift coefficient increases as the cross-sectional area of the test section decreases. The magnitude of this effect is disguised somewhat by the logarithmic scales and by the effect of the variation in dynamic pressure in computing the coefficient when a large part of the lift (from the fans) is essentially independent of forward speed. Figure 14 presents a truer picture of the influence of the walls by presenting the lift in the form of a ratio to the static thrust. Reference 23 has shown that the ratio  $L/T_S$  is proportional to a lift coefficient based on fan area and fan-exit velocity; that is

$$\frac{\mathrm{L}}{\mathrm{T}_{\mathrm{S}}} = \frac{1}{2} \, \mathrm{C}_{\mathrm{L},\mathrm{j}} \tag{11}$$

In figure 14 the differences in the data, as measured in the various test sections, are demonstrated to represent very significant differences in lift. For example, at an angle of attack of zero (fig. 14(b)) and a speed of  $V/V_j = 0.4$  (which would represent a speed near the high-speed end of transition), the data from the small inserts indicate that a lift of about 25 percent more than the static thrust would be obtained; the data from the moderately larger 2.13- by 3.05-m (7- by 10-ft) test section would indicate that the gain in lift would be only 10 percent; and the data from the largest test section indicate that a small loss in lift would be encountered. Indeed, the data from the largest test section indicate that this model would have a loss of lift (from that expected from a simple addition of lift components) for all angles of attack and for all forward speeds less than  $V/V_j = 0.5$ . This speed range encompasses the entire feasible transition range of lifting fans with modern pressure ratios.

Reference 7 presents a set of charts which define relative proportions between model and test section which were believed to yield negligible wall effects at a speed of 30 knots. Figure 18 shows the degree to which the present tests meet these size limits. Only the highest and the lowest disk loadings encountered are shown. The test conditions include points between these two disk loadings. In particular, the present tests in the 1.12- by 2.24-m (44- by 88-in.) insert meet these limits at least as well as many of the tests reported in references 26 to 32. In the 2.13- by 3.05-m (7- by 10-ft) tunnel, the size of the present model falls well within the size limitations of reference 7. In contrast, the data in figure 14 clearly indicate unacceptably large overestimates of "faninduced" lift in these test sections. Therefore, it must be concluded that the size limits proposed by reference 7 are not valid for configurations such as that of the present investigation; indeed, since there is nothing very unusual about this configuration except the relative size of the tail, it would be presumed that these size limits are equally inapplicable to other configurations as well. Reference 7 attempts to limit its conclusions to conditions for which the overall drag of the model is trimmed. This limitation to zero net drag is based upon the the-oretical results of reference 2, which the authors of reference 7 claim to be incorrect and even in the wrong direction. In fact, reference 2 was misinterpreted in arriving at the limitation to zero net drag. As will become clear in the subsequent discussion of correcting the present fan-in-wing data, each element of the aircraft must be considered individually. Thus, the fans, except for a few isolated conditions, always have a drag; also, the wing always has an induced drag. The addition of a centered jet exhausting directly rearward ( $\chi = 90^{\circ}$ ) could balance the drag of the model under any condition; however, the thrusting jet would contribute nothing to the interference at the model (from ref. 2,  $\delta_{u,L} = \delta_{w,D} = \delta_{u,D} = 0$ , and  $\delta_{w,L}$  has no effect since the lift of the thrusting jet would be zero). Consequently, the limitation to trimmed drag in reference 7 is meaningless.

The comparisons between flight and wind-tunnel data given in reference 7 have already been discussed in reference 39, which shows that the conclusions of reference 7 were based upon faulty comparisons between flight and wind tunnel. Such an error should have been anticipated, since one of the conclusions was that both Glauert's corrections and those of reference 2 were in the wrong direction. Since both Glauert and reference 2 predict upwash interference in a closed tunnel, this result of reference 7 could only be obtained if a downwash interference was produced by the walls. Such a result is physically impossible for an overall correction in a closed tunnel. Indeed, references 4 and 39 have already demonstrated that the calculated flow of reference 2 is in the correct direction.

Since the only guide in choosing model sizes for the fan-in-wing tests of references 26 to 32 has been the set of limits given in reference 7, the data shown in figures 14 and 18 should lead to serious concern regarding the highly favorable "fan-induced" lift reported as one of the main advantages to the fan-in-wing configuration in those studies which have produced and correlated uncorrected wind-tunnel data (e.g., refs. 19 and 26 to 32). One such correlation (from ref. 19) is presented in figure 19, where the "faninduced" lift is correlated as a function of the ratio of fan area to wing area.

The present model has a ratio of fan area to wing area of 0.094 and, as may be seen in figure 14, has a "fan-induced" lift in the 1.12- by 2.24-m (44- by 88-in.) flat-oval insert which lies very near the lower boundary of the correlation region. This value does not really correspond to the other data in figure 19 because all those data were obtained with models having either no tail or a small tail; whereas the present model has an extremely large tail which carries a significant download (fig. 13) under almost all conditions. It is easily shown from the definition of  $C_{N,T}$  and  $V_j$  that

$$\frac{N_{T}}{T_{S}} = \frac{1}{2} C_{N,T} \frac{S_{T}}{S_{F}} \left( \frac{V}{V_{j}} \right)^{2}$$
(12)

The value of  $C_{N,T}$  at  $\alpha = 0^{\circ}$  and  $V/V_j = 0.4$  is obtained from figure 13(b) as -0.53;  $S_T/S_F$  is 3.73; thus, from equation (12) the ratio  $N_T/T_S$  is equal to -0.16 for the conditions of figure 19. Removing the tail load from the data of figure 14 increases  $\Delta L_i$  (= L - T<sub>S</sub>) from 0.25 to 0.41, which is near the upper edge of the correlation band of reference 19. (See fig. 19.) On the other hand, in the largest test section,  $\Delta L_i/T_S$  is negative for the complete model at this value of  $V/V_j$ .

The slope of the correlation band of reference 19 which is reproduced in figure 19 deserves some comment. It is obvious in examining the data points of figure 19 that a band of the same width, drawn parallel to the abscissa, and thus indicating total independence from the ratio  $S_F/S_W$ , would have encompassed a larger number of data points than the band which was drawn. In either event, the major exceptions to the correlation band are those configurations in which the fans are displaced far from the center of pressure of the wing. Such aircraft would be unflyable as VTOL configurations without the provision of additional fans to provide moment balance.

Since reference 7 presents several different criteria upon which to scale windtunnel tests for wall interference, it is advisable to perform a first-order analysis in order to determine which parameters really are significant. For this first-order analysis, examine the zero-angle-of-attack case, for which the wing of the present model would have no lift in free air. Then assume that the horizontal components of wallinduced interference have only a second-order effect and that  $\Delta \alpha$  is sufficiently small to let  $\Delta \alpha \approx \tan \Delta \alpha$ . Under these assumptions, following references 1 to 3,

$$\Delta \alpha = \frac{\Delta w}{V} = \left( \delta_{w,L} + \frac{D}{L} \delta_{w,D} \right) \frac{S_F}{A_T} \frac{w_0}{V}$$
(13)

where  $\Delta \alpha$  is in radians and where  $\delta_{w,L}$  and  $\delta_{w,D}$  are calculated for the fans.

From reference 23, momentum theory shows that for the fans at  $\alpha = 0$ ,

$$\begin{array}{c} \mathbf{w}_{0} = \mathbf{w}_{h} = -\mathbf{V}_{j} \\ \\ \frac{\mathbf{D}}{\mathbf{L}} = \frac{\mathbf{V}}{\mathbf{V}_{j}} \end{array} \right\}$$
(14)

Substitution of equations (14) into equation (13) yields

$$\Delta \alpha = -\left(\delta_{w,L} + \frac{V}{V_j} \delta_{w,D}\right) \frac{S_F}{A_T} \frac{1}{V/V_j}$$
(15)

Observe that both  $\delta_{w,L}$  and  $\delta_{w,D}$  are negative in a closed tunnel; thus,  $\Delta \alpha$  will be positive (upwash). Reference 23 shows that the lift of the fan is virtually insensitive to angle of attack for angles near zero; therefore, the increase in lift will be essentially all on the wing. This increase in lift may be written as

$$\Delta \mathbf{L} = \Delta \alpha \mathbf{C}_{\mathbf{L}_{\alpha}} \mathbf{q} \mathbf{S}_{\mathbf{W}} \tag{16}$$

Substitute equation (16) into equation (15) to obtain

$$\Delta \mathbf{L} = -\left(\delta_{\mathbf{w},\mathbf{L}} + \frac{\mathbf{V}}{\mathbf{V}_{j}} \delta_{\mathbf{w},\mathbf{D}}\right) \frac{\rho}{2} C_{\mathbf{L}_{\alpha}} \mathbf{V}^{2} S_{\mathbf{W}} \frac{S_{\mathbf{F}}}{A_{\mathbf{T}}} \frac{1}{\mathbf{V}/\mathbf{V}_{j}}$$
(17)

Divide both sides of equation (17) by  $T_S = \rho S_F V_j^2$  to yield

$$\frac{\Delta \mathbf{L}}{\mathbf{T}_{S}} = -\left(\delta_{W,L} + \frac{\mathbf{V}}{\mathbf{V}_{j}} \delta_{W,D}\right) \frac{\mathbf{C}_{L_{\alpha}}}{2} \frac{\mathbf{S}_{W}}{\mathbf{A}_{T}} \frac{\mathbf{V}}{\mathbf{V}_{j}}$$
(18)

Consider the product  $C_{L_{\alpha}}S_{W}$  in equation (18). Since  $C_{L_{\alpha}} = \frac{2\pi A}{A+2}$ , this product may be rewritten as

$$C_{L_{\alpha}}S_{W} = \frac{2\pi A}{A+2}S_{W} = \frac{2\pi \frac{b^{2}}{S_{W}}S_{W}}{A+2} = 2\pi \frac{b^{2}}{A+2}$$
(19)

Finally, substitute equation (19) into equation (18) to yield

$$\frac{\Delta \mathbf{L}}{\mathbf{T}_{S}} = -\left(\delta_{w, \mathbf{L}} + \frac{\mathbf{V}}{\mathbf{V}_{j}} \delta_{w, \mathbf{D}}\right) \frac{\pi}{\mathbf{A} + 2} \frac{\mathbf{b}^{2}}{\mathbf{A}_{T}} \frac{\mathbf{V}}{\mathbf{V}_{j}}$$
(20)

which is the wall-induced lift.

Observe that the only term of equation (20) which explicitly involves the model dimensions is  $b^2/A_T$ . For test sections having approximately the same width-height ratios (in the present case, from 1.4 to 2.0),  $b^2/A_T$  will be approximately proportional to the square of the ratio of the wing span to the test-section width. The ratio of fan area to test-section cross-sectional area is completely immaterial.

One more significant point is evident in the preceding analysis. It is generally believed that wall effects are greatest at low speeds. For a constant model configuration and for data presented in terms of coefficients based on free-stream dynamic pressure, wall effects are greatest at low speed. (Note that  $\Delta C_L = \Delta \alpha C_{L\alpha}$  and, from eq. (15), that  $\Delta \alpha$  has a 1/V component.) On the other hand, again for a constant model configuration, when the data are presented in terms of forces or force ratios, equation (20) clearly shows that the greatest effect of wall interference will be at high speed. This conclusion is confirmed by the data presented in figure 14.

It is obvious that the present results from the small insert lead to a gross overestimate of "fan-induced" lift. The correlation with the data presented in reference 19 indicates that the data presented therein also include substantial overestimates of "faninduced" lift which would not be obtained in flight. (Observe that, with two exceptions, the models of ref. 19 have essentially the same span-to-width ratio as the present model in the smallest inserts. Of the two exceptions, one is anomalous because of its thin delta wing; the second was notable for producing the smallest "fan-induced" lift of any of the models of ref. 19.) Further, the model of reference 19 in which the fans are behind the trailing edge of the wing would be expected to show a far smaller "fan-induced" lift than indicated in figure 19, and, similarly, the model with the fans well forward of the wing would be expected to show far greater "fan-induced" losses than indicated. In either of these two cases, the results would be affected substantially by provision of the additional fans required for moment control in the VTOL mode. This latter effect is evident in figure 19 when these two configurations are combined into one. (See the appendix of ref. 23 for a further discussion of the effect of fan location on mutual interference.) Irrespective, however, of whether or not VTOL moment control is feasible for the configurations of reference 19, it is obvious that all the data of that paper contain a large increment of wallinduced, rather than 'fan-induced,' interference. The 'good' configurations will be far less "good" in free air; the "poor" configurations will be even worse in free air.

It will be observed that there are differences in notation between the present paper and reference 19. In the present paper,  $\Delta L_i$  is defined (at  $\alpha = 0^{\circ}$ ) as L - T<sub>S</sub> since T<sub>S</sub> is equal to the thrust in forward flight according to ideal momentum theory (ref. 23). Similarly, V<sub>j</sub> is defined (see Symbols) as the fan efflux velocity in static thrust. In reference 19,  $\Delta L_i$  is defined as total lift, less any wing lift (which is zero at  $\alpha = 0^{\circ}$  in the present tests), less the thrust in forward flight as measured by rakes in the fan exit; and  $V_i$  is defined in relation to this measured thrust.

Because of the square root involved in determining  $V_j$  from the thrust, as well as the relatively flat character of  $L/T_S$  near  $V/V_j = 0.4$  and  $\alpha = 0^{\circ}$  (the conditions chosen by ref. 19; see fig. 14(b)), there will be little effect of the difference in definition of  $V_j$ . The difference in definition of  $\Delta L_i$  has a more serious effect. In practice, because of inlet efficiency, the actual value of thrust in forward flight will be somewhat less than the theoretical value of  $T_S$ . Figure 6 of reference 32 indicates that at  $V/V_j$ of 0.4, a loss of 10 to 15 percent of  $T_S$  may be expected for a typical lift-fan model. For complete comparability, this loss should be added to the present results; that is, in figure 19 the values of  $\Delta L_i/T_S$  should be about 0.1 to 0.15 greater than indicated therein for the present model. Thus, the effect of wall interference on the data of reference 19 may be even greater than indicated by the data shown in figure 19.

Effect of wall interference on drag.- The drag of the model also increases as the tunnel size decreases (figs. 12 and 15); however, the increases in drag are not commensurate with the increases in lift. Indeed, at the lower angles of attack, the increases in drag are minimal. The disparity between the increases in lift and drag may be seen more clearly in figure 16, which presents the external drag-lift ratio for the model. At all angles of attack, D/L is greater in the larger test sections, thus indicating poorer efficiency.

The apparent gain in efficiency in the smallest test sections is retained even when the data are presented in terms of  $L/D_E$  (fig. 17). For example, at  $V/V_j = 0.4$  and at  $\alpha = 0^{\circ}$  (fig. 17(b)),  $L/D_E$  as measured in the smallest test sections is approximately 25 percent greater than the same values measured in the largest test sections. Thus, wall effects are sufficient to indicate a 25-percent decrease in the power required to fly in the transition speed range.

The values of  $L/D_E$  shown in figure 17 appear at first glance to be remarkably small. They are confirmed however by the momentum theory presented in reference 23. This confirmation is demonstrated by the theoretical curves (from ref. 23) given in figure 17. They are further confirmed by calculations made using the measured shaft powers given in reference 45, as well as by the extraordinary fuel consumption in low-speed flight found in design studies of fan-supported aircraft (ref. 46).

Effect of wall interference on tail normal force.- The uncorrected measurements of tail normal force, as a function of  $V/V_j$ , are shown in figure 13. At low speed, the trends shown for the various test sections are observed to scatter. This effect is probably due to Rae's limit (refs. 4, 9, 13, and 18), and it will be discussed in the next section.

At the higher speeds and for angles of attack less than  $10^{\circ}$  (figs. 13(a) to 13(c)), the observed tail normal-force coefficient is essentially independent of the test-section size or shape. As the angle of attack becomes greater, the data from the various test sections show greater differences (fig. 13(e)), with the tail normal force becoming more positive as the test-section size decreases.

Even constancy in tail normal-force coefficient would indicate a serious degree of wall interference since, in free air, the increased lift (shown in fig. 14) in the small sections would increase the downwash at the tail, reduce the tail angle of attack, and result in a more negative tail normal-force coefficient. However, the data of figure 13 indicate that the wall-induced interference at the tail is of sufficient magnitude to negate, or even to reverse, the trend that would be expected with increased lift.

Wall-induced effects at the tail, of course, are not confined to this configuration. Large wall effects have also been noted in comparing large-scale wind-tunnel and flight data; for example, consider the comparison of flight-test data and uncorrected wind-tunnel test data (Ames full-scale tunnel test 388) presented in reference 20 for a YOV-10 aircraft fitted with a rotating-cylinder flap. Serious differences were found in maximum lift and the angle of attack at which it was obtained; however, by far the greatest disagreement between wind tunnel and flight was with regard to the effects at the tail.

Figure 20 shows these differences (as presented in ref. 20) in terms of the elevator angle required to trim the aircraft as a function of forward speed. The uncorrected wind-tunnel data indicate positive speed stability with the stick moving rearward (the elevator moving trailing edge upward) as the speed decreases; the elevator is  $20^{\circ}$  trailing edge up when 55 knots is reached. In contrast, the flight data indicate a speed instability; the elevator angle is always in the opposite sense (trailing edge down); and at 55 knots the elevator angle is  $13^{\circ}$  trailing edge down. The total disagreement between tunnel and flight at 55 knots is  $33^{\circ}$ , and this disagreement is in the same direction as that indicated in figure 13.

The trends shown in figure 20 are given further import by the flight-test data when extended to slightly lower speeds (fig. 21). Here the speed instability became more dramatic, and the minimum speed in many cases was determined by the speed at which the elevator contacted the limit of travel in the trailing-edge-down direction and not by maximum lift. Needless to say, under such circumstances, the pilot finds himself in somewhat compromised circumstances because he has no control left for any unanticipated maneuvering requirement. The point here, of course, is that not even full-scale windtunnel tests of the actual aircraft gave any indication that the pilot would find himself in these circumstances because wall effects were not properly accounted for in the data reduction. It is noted that the YOV-10 with the rotating-cylinder flap, as tested in the tunnel, also fell well within the boundaries of reference 7, which, according to that paper, would indicate negligible wall effects (fig. 22). This evidence confirms the previous conclusion that the testing boundaries of reference 7 are erroneous.

The opinion is sometimes voiced that wind-tunnel interference does not affect the trends shown by the data, or, as expressed in reference 47 (p. 7-1): 'Informative results, even when the model lifting system spans 2/3 to 3/4 of the wind tunnel test section width, will still be obtained.'' Neither the model of the present investigation, nor the aircraft of reference 20, approached so great a size relative to the test section; and, in each case, the wall interference was so great that even the trends shown by the data were in the opposite direction from ''free air'' results at low speed. Such results clearly show that extreme caution must be used when interpreting uncorrected wind-tunnel data.

Effect of test-section size on Rae's limit.- The separately instrumented tail was installed on the model in the hope that measurements of tail normal force would provide a sensitive indication of the onset of the recirculation which results in Rae's minimum-speed limit (ref. 9). This procedure was chosen because of the dramatic alterations in tail lift which were observed behind a rotor in reference 18.

Although the tail normal-force-coefficient data presented in figure 13 do show marked effects as a function of tunnel configuration at low forward speed, effects as definitive as those of reference 18 were not always observed. One reason may be the magnitude of the wall interference in the present tests. This aspect of the problem will be discussed in subsequent sections of the present paper. A second reason is that the tail on this model, as may be seen by comparisons between the individual parts of figure 13, has very nearly zero tail effectiveness (that is,  $\frac{d\epsilon}{d\alpha} \approx 1$ , so that  $1 - \frac{d\epsilon}{d\alpha} \approx 0$ ) until the highest angle of attack (16<sup>o</sup>) is reached. At  $\alpha = 16^{\circ}$ , the tail normal-force coefficient suddenly turns upward as the speed is reduced in the small test sections. In the small inserts,  $C_{N,T}$  departs from the trends shown in the data from the 9.14- by 18.3-m (30- by 60-ft) tunnel at a value of  $V/V_1$  below 0.38; a similar departure may be observed in the data from the 2.13- by 3.05-m (7- by 10-ft) tunnel at a value of  $V/V_1$  below about 0.2. The values correspond approximately to conditions at which visual tuft observations indicated substantial flow reversal on the floor; furthermore, the values are roughly in proportion to the height of the various test sections, as might be expected from the correlation rules presented in references 6, 9, and 21. Unfortunately, those rules are expressed in terms of the momentum wake angle. Since the momentum wake angle for the fan is always along its axis and is actually negative for the data of figure 13(e), those rules cannot apply to the present case.

Tyler and Williamson (refs. 48 and 49) have conducted a systematic program to determine minimum-speed test limits for jet lifting systems. Their results indicate

incipient stagnation (near  $\alpha = 0^{\circ}$ ) on the floor of the test section when

$$\frac{\mathbf{V}}{\mathbf{V}_{j}} = 1.59 \, \frac{\mathbf{d}_{e}}{\mathbf{h}} \tag{21}$$

for single and tandem-paired jets, and when

$$\frac{\mathbf{V}}{\mathbf{V}_{j}} = 1.31 \, \frac{\mathrm{d}_{\mathbf{e}}}{\mathrm{h}} \tag{22}$$

for a laterally paired system of two jets spaced 4.3 nozzle diameters apart. The spacing of the two fan nozzles in the present model is considerably closer (2.75 diameters); nevertheless, using equation (22) yields  $V/V_j = 0.67$  for the two 1.12- by 2.24-m (44- by 88-in.) test sections;  $V/V_j = 0.62$  for the 1.22- by 1.83-m (4- by 6-ft) test section; and  $V/V_j = 0.39$  for the 2.13- by 3.05-m (7- by 10-ft) test section. The corresponding values for the largest test section are below the smallest velocities at which tests were run.

To define the point of incipient stagnation and to define the point at which the data will be affected are two different things, as is noted in reference 49. Tyler and Williamson suggest that test speeds as small as 55 percent of the speed for incipient stagnation may be acceptable for single jets and 65 percent of this speed may be acceptable for widely spaced lateral pairs of jets. If the values obtained in the preceding paragraph are reduced by a multiplying factor of 0.6 (an average of 0.55 and 0.65), they will be observed to agree closely with the previously noted points of figure 13(e). Therefore, it would appear that the Tyler and Williamson relations (eqs. (21) and (22)) provide a reasonable means of estimating the minimum speed for wind-tunnel testing of jet- and fansupported models.

The value observed for the degradation of data due to recirculation in the 1.12- by 2.24-in. (44- by 88-in.) flat-oval test section was about  $V/V_j = 0.38$ , and that obtained from equation (22) reduced by 40 percent was 0.40. Not only the correlation between these values is of interest; their magnitude is significant in itself. Observe that the correlation of "fan-induced" lift in reference 19 was obtained at  $V/V_j = 0.4$ . It is entirely possible that some of the data upon which reference 19 is based are suspect because of recirculation effects, since the model to tunnel-size ratios in those data are comparable with those obtained in the present small flat-oval insert. Furthermore, at 30 knots, it is clear that much of the data used to prepare the testing limits defined in reference 7 were obtained for flow conditions which were unrepresentative of flight in free air because of flow breakdown induced by the model in the wind tunnel.

Two primary requirements exist in planning wind-tunnel tests. One is simulation of the aircraft, and, given a drawing of the aircraft, it is simple to produce a reasonable model of it. Equally important, however, is that the basic free-air flow must also be simulated. At speeds less than Rae's limit, a powerful cylindrical sheet of vorticity is formed ahead of the intersection of the wake on the floor (refs. 4, 13, 18, and 39). This sheet ultimately extends across the floor and up the sides of the test section. Except in ground effect, no equivalent vortex formation exists in actual flight. Under such conditions, the flow in the test section does not simulate free air and almost any result may be obtained.

It is particularly important to realize that the existence of the basic alteration of the flow does not depend upon the presence, or the absence, of a tail on the model. The flow alteration is caused by the presence within the walls of the main lifting system. Indeed, the models used by Rae (ref. 9) when he discovered this effect had no tails; neither did the models used by Tyler and Williamson (refs. 48 and 49).

<u>Comparison of simple momentum theory with experimental results</u>.- Reference 23 develops a simple incompressible-flow momentum theory for the fan-in-wing configuration based upon the assumption that there is no mutual interference between the fans and the wing. Momentum theory, by itself, is incapable of calculating the actual lift of the model, since the lift depends intimately on the local angles of attack of the wing and of the fan blades. However, once the lift is given, momentum theory is capable of estimating the remaining performance items. Momentum theory, obviously, also is incapable of predicting the tail normal-force coefficient because this coefficient depends upon a detailed calculation of the flow field in the vicinity of the tail.

In the present case, it is assumed (following ref. 23) that the thrust of the fan is unaltered by forward speed or angle of attack. This assumption is verified by figure 6 of reference 32, which shows that the actual thrust for a typical lift-fan model (at  $V/V_j$  as great as 0.6) is only 10 or 15 percent less than the static thrust. When the normal component of static thrust is added to the lift of the wing with covered, inoperative fans, the results previously presented in figures 11 and 14 are obtained. Evidently, at high speed, significant fan-wing interaction effects are present; however, throughout the usable transition speed range ( $0 \leq V/V_j < 0.5$ ), the assumption of zero interaction yields values close to the observed total lift. The differences in notation between reference 19 and the present paper have no effect herein, since both theory and experiment are presented in the identical manner.

All the remaining curves in figures 12 and 15 to 17 follow directly from the equations of reference 23 once the lifts are assumed. It will be observed that, for transition speeds, the observed performance is predicted more closely by even this simple momentum theory than by the data from the small inserts. Note that the model in the present investigation spanned only a little less than half the width of the smallest inserts. This relative size is essentially the same as that used in references 26 to 32, and a similar result may be implied to be true for those tests as well.

#### Correcting the Data

Considerations in correcting the data.- Correcting the data with the fans operating follows the same general procedure described earlier for the data with the fans covered. Obviously, the procedure is complicated to a degree in accommodating the presence of the fans. In this case, the interference factors are obtained from appendices O and P of reference 36.\* The previously discussed modifications to subroutine DLTAS (appendix Q of ref. 36) were used to obtain the interference factors at the pitot-static tube location. The solid blockage factors are identical with those used when the fans were covered. Since there was no independent measurement of the thrust of each of the fans, there is no alternative but to deal with them simultaneously. Therefore, the appropriate interference factors for the pair of fans are the average of those for the fan due to its own presence and those for the fan due to the presence of the other fan. In all cases, slight changes to the programs of reference 36 allowed data cards containing the interference factors to be punched automatically as they were calculated. These cards were used as input data to the data correction program to be discussed shortly. This procedure eliminates the possibility of errors in transcription when preparing the input to the correction program.

Reference 36 offers choices of wing load distribution and rotor-disk load distribution. In the absence of definitive measurements to the contrary, an elliptic load distribution was chosen for the wing. In order to ascertain the degree to which this choice might affect the corrections, the data were also reduced using the interference factors for a uniform load distribution; no significant effect was found for this model, perhaps because of the magnitude of the corrections. Because of the large central boss in the fans, the disk load distribution over the faces of the fans is not uniform. Consequently, the triangular disk load distribution. In any event, the fans were so small compared with the test-section dimensions that little effect of this choice should be evident.

<sup>\*</sup>Three known errors exist in the programs given in reference 36. Two of these affect the work contained herein. The following lines should be corrected to read as follows:

805	XDELTA(L1)=XDELTA(L1)+DELTA(L1)*XLOAD(N1)	(E 79)
	SUML=0.063052	(P 113)
	SUML=0.252208	(P 136)
The first problem is to divide the measured loads between the elements which produce them. The tail presents no problem since it was mounted on its own strain-gage balance; however, there was no balance to separate the independent forces of the wing and fans when they were operating in unison. In the absence of specific information, the wing was assumed to produce the same lift and drag as it did when the fans were covered. Thus, the fan lift and drag are assumed to be the main balance readings, less the measured tail loads, and less the aforementioned assumed lift and drag of the wing. Then the fan lift and drag thus obtained were used to solve the momentum quartic and to calculate the fan wake skew angle  $\chi$  and the fan velocity ratio  $V/w_0$ . The resulting values of  $\chi$ were within 2<sup>o</sup> or 3<sup>o</sup> of being equal to  $-\alpha$ , as they should be (ref. 23). For the fans, it is more appropriate to use the effective skew angle

$$\chi_{\rm e} = \frac{90^{\rm o} + \chi}{2} \tag{23}$$

as given in references 8 and 39. For the wing and for the tail, the effective skew angle was assumed to be 90°. In the face of the powerful downwash field generated by the fans, the use of the momentum quartic given in references 2 and 22 is not strictly applicable because it would be necessary to include the local effective downwash angle in the vector diagram defining  $\chi$ . In any event, the use of  $\chi = 90^{\circ}$  for the wing and the tail is a great simplification in the calculation.

At this point, everything is in hand to compute the average interference velocity components over the wing, over the tail, and at the pitot tube. Observe that the fans now contribute substantially to the interference velocities at each location.

Reference 13 has noted that it is necessary to apply a vortex-density correction to the theoretical interference factors when applying them to ducted fans. This correction was used in the present calculations. A justification of the vortex-density correction is presented in appendix A.

The first step in applying corrections to the data from the smaller Ames tunnel is to correct the pitot-static tube reading of tunnel velocity. The forces that were charged to the wing are directly dependent on dynamic pressure; thus, it is necessary at this point to return to the original division of loads and redo that division with the corrected dynamic pressure, and then repeat all the steps to this point.

The calculation now proceeds as before, correcting the overall performance to the corrected flight condition at the wing and then adjusting the tail loads to account for the substantially different wall-induced interference at the tail location. Despite the fact that the fans are mounted within the wing planform, there is a significant difference in the average wall-induced interference over the full span of the wing and the similar average

over the faces of the fans. It is necessary to remove this difference by adjustments to the fan lift and drag. The adjustment is accomplished by the use of the following equations from reference 23:

$$\Delta D = T_{S} \left[ \left( \frac{\Delta V}{V_{j}} \right)_{F} + \Delta i_{F} \cos \alpha \right]$$
(24)

 $\Delta L = -T_S \Delta i_F \sin \alpha \tag{25}$ 

where  $\Delta i_F$  is in radians. It will be observed from figure 15 that an  $\alpha$ -dependent multiplying factor, generally greater than 1.0, could have been applied legitimately to equation (24) (that is, the actual drag of the fans is generally greater than the momentum-theory value). This was not done; equation (24) was used directly as given above.

Corrections at zero tunnel velocity are particularly suspect. As noted in reference 8, a hovering condition in the tunnel leads to an interference which is a pure upwash. Proceeding in a formal manner, this upwash is equivalent to change in angle of attack of  $90^{\circ}$ . While true in a sense, it is more rational to consider the model to be at the same angle of attack, but with a rate of sink equal to the upwash velocity. The corrections to the data then would depend upon the effect of a rate of sink. Unfortunately, sufficient data are not present to make such a correction for this model. Furthermore, at zero speed, the test conditions always violate Rae's minimum-speed limit (ref. 9).

In view of the foregoing observations, no attempt has been made herein to correct data obtained at zero velocity. When such points are shown in the corrected data, they are identical with the uncorrected data, and they are presented only to preserve the continuity of the data set.

<u>Computer program</u>.- Appendix B presents one of the computer programs used in correcting the data obtained in the present investigation. Because of differences in the measurement of fan rotational speed, it was necessary to write slightly different programs for the two separate sets of tests in the Ames tunnel. The absence of a stream-angle correction in that tunnel provides one simplification in that it was possible to compute the interference factors specifically for the angles of attack which were used; thus, only a single interpolation against  $\chi_e$  is required to obtain the proper factors for each data point. The presence of a small stream angle in the Langley 2.13- by 3.05-m (7- by 10-ft) insert was accommodated by means of a double interpolation against both  $\chi$  and  $\alpha$  in a third modified version of the computer program.

The program of appendix B is substantially more complex than would normally be required because it simultaneously treats four different test sections. This feature requires considerably more storage and additional steps (which increase running time) than would a program written for a single test section. Nevertheless, in the Langley Computer Complex, the program requires only  $54000_8$  ( $\approx 22500_{10}$ ) storage locations for compilation,  $46000_8$  ( $\approx 19500_{10}$ ) storage locations to run, and completely corrects more than 360 data points in about 30 seconds (including compilation time) at a cost of only \$7. The storage lengths, the time, and the cost obviously would be different in almost any other computer; however, it also is obvious that only minimal costs and computer capabilities are required to fully correct data for wall effects even for a fairly complex model.

The program of appendix B produces several files of output data and a sequence of punched-card sets for subsequent plotting of the data. In sequence, the written files present the interference factors used in the correction routines, the uncorrected data together with a preliminary breakdown of the division of loads (in the Langley insert, the presence of a stream angle requires interpolation in order to obtain this file), a point-bypoint listing of the corrected data, a listing of the corrected data at fixed angles of attack obtained by interpolation of the previous listing, an interpolated listing of the corrections themselves at a series of fixed corrected angles of attack. Punched-card decks of the last four listings are provided for subsequent automatic plotting of the data.

The data herein are presented as a function of forward speed. If it is desired to obtain polar plots of the performance at fixed speeds, an interpolation against  $V/V_j$  would be required. The addition of one more interpolation should not present any significant difficulty.

## Interference in Uncorrected Data

The corrections in the uncorrected data of figures 13 to 17, as calculated from the foregoing considerations, are presented in figures 23 to 28. The corrections are distinguished by their enormous magnitude, which far exceeds the more reasonable values suggested as the maximum practical limits in references 6 and 42. Depending upon  $\alpha$  and  $V/V_j$ , the average interference angle  $\Delta \alpha$  at the wing varies from about  $2\frac{10}{2}$  to over  $14^{\circ}$  in the smaller inserts (fig. 23). Similarly, the effective dynamic pressure at the wing is reduced by 5 to 22 percent (fig. 24). The effective tail incidence in the smaller inserts is increased by from 5° to 12° (fig. 25) and the dynamic pressure at the tail varies from 1.15 to almost 3 times that at the wing (fig. 26). In the small inserts, even the fans are operating at effective angles of attack as much as 7° more than the wing in which they were mounted (fig. 27). Only the ratios of the dynamic pressures at the fans to those at the wing remain relatively small (fig. 28). The wall interference in the 2.13- by 3.05-m

(7- by 10-ft) test section is generally of a lesser magnitude, although, even in that test section,  $\Delta \alpha$  and  $\Delta i_T$  (figs. 23 and 25) tend to be larger than would be desirable.

It is noticeable that the values for Rae's limit, which were obtained earlier, were for speeds so low that a prudent investigator would have discontinued testing long before this limit was a serious concern. This result is in accordance with the results of reference 6, which show that Rae's limit is of primary concern only for those models which are very small with respect to the test-section size. In all other cases, the magnitude of the wall-induced distortions of flow over the model are the controlling factors in choosing the maximum permissible model size.

#### **Corrected Data**

It is obvious from the magnitude of the indicated corrections (figs. 23 to 28) that perfect correlation of the data from all the test sections cannot be expected. Nevertheless, the corrected data show remarkably improved agreement (figs. 29 to 35). This agreement is poorest at those speeds previously determined to be less than Rae's limit  $(V/V_j \approx 0.38)$  in the smallest test sections). It is also somewhat poorer in the smaller inserts than in the more moderately sized 2.13- by 3.05-m (7- by 10-ft) test section. (See particularly fig. 33(a).) On the other hand, considering the magnitude of the required corrections, the data presented in figures 29 to 35 are an impressive verification of wall-interference theory (refs. 1 to 6).

#### Interference in Corrected Data

It is obvious that the angular range of model settings in figures 29 to 35 differs from that in figures 11 to 17 because the geometric angle of attack in figures 11 to 17 has been decreased by  $\Delta \alpha$  to obtain the corrected values of  $\alpha$  in figures 29 to 35. Since the wall-induced interference increases as the lift and drag increase, and since these forces increase with  $\alpha$ , the corrections are somewhat less in figures 29 to 35 than in figures 11 to 17. For completeness, figures 36 to 41 have been prepared to indicate the magnitude of the corrections that are actually present in the corrected data. Although slightly less than the values presented earlier, the corrections in the final data are still extremely large.

#### Maximum Model Size

It is clear from the data presented in this paper that the model was so large in relation to the small inserts that the corrections were excessive. At low speeds, the model was excessively large in the 2.14- by 3.05-m (7- by 10-ft) test section as well. It would appear that prudent model sizing would have led to a model having a span of about a quarter of the test-section width. A similar conclusion as to an appropriate model size would have been reached by examining the charts of reference 6. The overall corrections  $(\Delta \alpha \text{ and } q_c/q)$  and the wall-induced tail incidence  $(\Delta i_T)$  are about the same in those charts as were found herein; however, the charts of reference 6 fail totally to indicate the magnitude of the wall-induced dynamic-pressure ratio  $(q_T/q_c)$  at the tail. Such a discrepancy is understandable. The charts of reference 6 are based on the assumption of a winglike model having a uniformly loaded span and a single, blended wake; these assumptions are grossly violated by the present fan-in-wing model. Thus, in using reference 6 to size models of unusual VTOL configurations which do not approximate the assumptions of that paper, it is best to err on the safe side by selecting a model size even smaller than indicated therein.

It will be observed that nonuniformities in wall-induced interference decrease more rapidly (ref. 6) with decreases in span (in a given tunnel) than do the overall corrections, which are roughly proportional to the square of the span. Thus, a small model does not require the same rigor in applying corrections as is required for a large-span model. Furthermore, considerably more confidence in the final results is justified when the model is small enough to require only minimal correction to the data.

While the present model should have been about half its present size, in the smallest insert, this conclusion should not be extended to indicate that all VTOL models should span about one-quarter of the test-section width. The allowable size of a VTOL or STOL model will depend upon the configuration, upon the minimum speed for which useful data are required, and upon the degree of correction applied to the data. Reference 6 should be some help in this regard; however, the only real safety will be found in correcting, as fully as possible, all wind-tunnel data as a standard practice. As noted earlier, the additional computing cost is minor in comparison with the total cost of a wind-tunnel investigation.

#### **CONCLUSIONS**

The results of this investigation of wind-tunnel wall interference on the performance of a fan-in-wing model are as follows:

1. Extreme caution must be used in interpreting uncorrected wind-tunnel data obtained at low speeds. Unless the model is extremely small in relation to the testsection size, the wall interference can be so large that even the trends in the data may be opposite to those which would be obtained in flight.

2. Wall-induced interference is particularly large at the model tail. This result confirms recently published (NASA CR-2135) conclusions based on the correlation of flight and wind-tunnel data for a YOV-10 aircraft.

3. The theory of wall interference for VTOL and STOL models, presented in NASA TR R-124 and subsequent papers, has been verified under conditions of extreme wall interference.

4. The rules for choosing model sizes to produce negligible wall effects, as given by Cook and Hickey in NASA SP-116 (also in AGARD Rep. 520), appear to be considerably in error and to permit the use of models which are significantly too large for the tunnel.

5. The method presented by Tyler and Williamson in AGARD CP-91-71 yields reasonable estimates of the onset of Rae's minimum-speed limit for jet- and fan-supported models; however, for reasonably large models, wall interference becomes so great that testing should be discontinued at a speed significantly greater than Rae's limit.

6. The "fan-induced" lift indicated by a number of previous investigations appears to be largely the result of wind-tunnel wall interference which was not accounted for in reducing the data. The uncorrected results obtained herein, when the model spanned almost half of the tunnel width, fall directly on a previously published correlation of "faninduced" lift (Hickey and Cook, AGARD CP-22, paper No. 15); however, the increase in lift for the model under conditions which approach testing in free air was small or negative for the actual transition speed range.

7. The simple incompressible-flow momentum theory presented in NASA TN D-7498 appears to yield reasonable estimates of fan-in-wing performance in transition flight; indeed, the theoretical predictions are more accurate than uncorrected wind-tunnel test data in which the model span is approximately half of the tunnel width.

Langley Research Center,

National Aeronautics and Space Administration, Hampton, Va., December 12, 1973.

#### APPENDIX A

# JUSTIFICATION OF VORTEX-DENSITY CORRECTION FOR FANS AND JETS

The theoretical treatment of wall interference for VTOL-STOL aircraft (refs. 1, 2, and 5) sets up the inclined wake of the aircraft in free air as a doublet string extending from the aircraft to infinity. If the model is small, this doublet string might represent a rotor or a lifting fan or a jet. Indeed, the relationship between the doublet strength and the induced velocity was obtained directly from an earlier analysis of a rotor wake in the wind tunnel (ref. 50). The basis of the relationship was the doublet strength required to match the vorticity along the edge of the vortex cylinder comprising the wake.

Now, if one considers a vortex cylinder cutting through an otherwise unbounded flow, and then takes the line integral of  $\overline{\mathbf{v}} \cdot d\overline{\mathbf{s}}$  (where  $\overline{\mathbf{v}}$  is the total vectorial velocity and  $\overline{\mathbf{s}}$  the path length), the eventual result is that the vorticity along the edge of the cylinder is precisely equal to the velocity jump across the cylinder. Unfortunately, integration, by means of the Biot-Savart law, of all the vorticity in the wake, leads to a velocity only one-half this great at the origin of the cylinder. In order to obtain the correct velocity at the end of the cylinder, it is necessary to double the vorticity. Since the corrections for wall interference depend upon the velocity  $w_0$ , which is the mean vertical induced velocity, it appears appropriate to take the corrections due to the fan as being twice as great as those of references 1 to 3. The changes need not be made in the interference factors, but can be made most simply by doubling the wall-induced interference components caused by the presence of the fans.

This effect was first noted in reference 13 and the vortex-density correction was used both in that paper and in the present analysis. The results of both papers appear to justify its use.

It will be observed that no similar correction is required for rotors or propellers. In those configurations, the induced velocity at the origin of the wake <u>should</u> be equal to one-half the vortex density.

#### FORTRAN PROGRAM FOR CORRECTING DATA FROM A FAN-IN-WING MODEL

#### TESTED IN FOUR DIFFERENT TEST SECTIONS

THIS PROGRAM WAS WRITTEN IN CDC FORTRAN, VERSION 2.1, TO RUN ON THE CDC 6000 SERIES COMPUTERS IN THE LANGLEY RESEARCH CENTER COMPUTER COMPLEX. MINOR MODIFICATIONS MAY BE NECESSARY PRIOR TO USE IN OTHER COMPUTERS. A DESCRIPTION OF THIS PROGRAM IS GIVEN IN THE TEXT OF THIS PAPER. A COMPLETE LISTING OF THE INTERFERENCE FACTORS USED IS INCLUDED. EACH LINE IS CODED AT THE END BY THE ANGLE OF ATTACK FOR WHICH THE FACTORS WERE COMPUTED, AN INTEGER CODE SPECIFYING THE TEST SECTION, AND THE CODE WORD DESCRIBED WITHIN THE PROGRAM. THESE INTER-FERENCE FACTORS WERE OBTAINED USING THE COMPUTER PROGRAMS OF NASA TM X-1740 (REF 36). CERTAIN ERRORS IN THAT REFERENCE ARE DISCUSSED IN THE TEXT.

THE SUBROUTINE DISCOT IS INCLUDED FOR COMPLETENESS. THIS IS A RELATIVELY STANDARD SINGLE OR DOUBLE INTERPOLATION ROUTINE. IT, OR ITS EQUIVALENT, WILL BE FOUND IN MOST COMPUTER SYSTEM LIBRARIES.

PROGRAM AARL6A (INPUT,OUTPUT,TAPE5=INPUT,TAPE6=OUTPUT,TAPE1,TAPE3, (B 1) 1 TAPE4,PUNCH) (B 2)

TUNNEL CODE ITUN≖1 IS 44X88 INCH WITH ROUND ENDS ITUN=2 IS 44X88 INCH WITH RECTANGULAR ENDS ITUN=3 IS 48X72 INCH WITH RECTANGULAR ENDS ITUN=4 IS 7X10 FOOT TUNNEL

FORTRAN WORDS REPRESENTING INTERFERENCE FACTORS ARE ALL CODED BY THE LAST FOUR CHARACTERS OF THE WORD. STARTING FROM THE RIGHT-HAND SIDE OF THE WORD, THE FIRST CHARACTER REPRESENTS THE ELEMENT ACTED UPON AND THE SECOND CHARACTER THE ELEMENT WHICH CAUSES THE WALL INTERFERENCE, WHERE: W=WING; F=FANS; T=TAIL; AND P=PITOT. THE NEXT TWO CHARACTERS ARE THE SUBSCRIPTS OF THE INTERFERENCE FACTORS AS DEFINED IN NASA TR R-124. VARIOUS PREFIXES ARE APPENDED TO THESE CODE LETTERS TO DISTINGUISH SPECFIC CHOSEN VALUES.

REAL LOTS	(B 3)
DIMENSION WLFF(4,6,8),ULFF(4,6,8),WDFF(4,6,8),UDFF(4,6,8),	(B 4)
1 WLWF(4,6),ULWF(4,6),WLWW(4,6),ULWW(4,6),WLFW(4,6,8),ULFW(4,6,8),	(B 5)
2 WDFW(4,6,8),UDFW(4,6,8),WLFT(4,6,8),ULFT(4,6,8),WDFT(4,6,8),	(B 6)
3 UDFT(4,6,8),WLWT(4,6),ULWT(4,6),WLTT(4,6),ULTT(4,6),ANAMF(28)	(B 7)
DIMENSION CWLFF(8),CULFF(8),CWDFF(8),CUDFF(8),CWLFW(8),CULFW(8),	(B 8)
1 CWDFW(8),CUDFW(8),CWLFT(8),CULFT(8),CWDFT(8),CUDFT(8)	(B 9)
DIMENSION TDALPF(6), TDALPT(6), TQOQF(6), TQOQT(6), TDVOVJ(6), TALP(6),	(8 10)
1 TQ(6), TDALPW(6), TQ0QW(6), TQC(6), TALPC(6), TCNC(6), TLIFT(6), SLT(6),	(8 11)
2 TVOVJ(6),TLOTS(6),TDOL(6),A(5),TQOQJ(6),EPSILON(4),BDT(6),BLT(6),	(B 12)
3 CHI(8),AT(4),SLAT(4,6),SDAT(4,6),JTUN(4),JTYPE(3),SLAT1(6),	(B 13)
4 SLAT2(6),SLAT3(6),SLAT4(6),SDAT1(6),SDAT2(6),SDAT3(6),SDAT4(6),	(B 14)
5 UOVPIT(4),WOVPIT(4)	(8 15)

С С С С С С £ С С С С С C C С С С

DIMENSION WLFP(4,6,8),ULFP(4,6,8),WDFP(4,6,8),UDFP(4,6,8), 1 WLWP(4,6),ULWP(4,6),WDWP(4,6),UDWP(4,6),CWLFP(8),CULFP(8), 2 CWDFP(8),CUDFP(8)	(B 16) (B 17) (B 18)
NOTE THAT WLFP;ULFP;WDFP;UDFP;WLWP;ULWP;WDWP; AND UDWP ARE ALL COMPUTED RETAINING THE N=M=O TERMS IN THE WALL-EFFECTS CALCULA- TION. THUS; THEY CONTAIN THE DIRECT EFFECT OF THE MODEL FLOW FIELD AS WELL AS THE WALL EFFECTS AT THE PITOT LOCATION.	
DATA (A(I),I=1,5)/-5.,0.,5.,10.,16./ DATA (AT(I),I=1,4)/24.004,26.889,24.,70./ DATA (CHI(I),I=1,8)/20.,30.,40.,50.,60.,70.,80.,90./ DATA (JTUN(I),I=1,4)/10H2-1 ROUND ,10H2-1 RECT. ,10H 1.5-1 , 1 10H 7X10 / DATA (JTYPE(I),I=1,3)/10H WT TARE ,10H AERO TARE,10H DATA / DATA (EPSILON(I),I=1,4)/.02705,.02445,.02565,.00716/	(B 19) (B 20) (B 21) (B 22) (B 23) (B 24) (B 25)
UDVPIT AND WOVPIT ARE THE SOLID BLOCKAGE EFFECTS (INCLUDING THE DIRECT MODEL FIELD) OF THE BODY AT THE PITOT LOCATION. THESE ARE COMPUTED FROM A SOURCE AND SINK SPACED SO AS TO PRODUCE IN FREE AIR A RANKINE OVOID OF THE SAME LENGTH AND DIAMETER AS THE MODEL FUSELAGE. IN THE TUNNEL THE BODY WILL BE SOMEWHAT SLENDERER AND THE CALCULATED BLOCKAGE WILL BE SOMEWHAT UNDERESTIMATED.	
DATA (UOVPIT(I),I=1,4)/0)672,00707,00443.0.00035/ DATA (WOVPIT(I),I=1,4)/2*0.00685,0.00622,0.CC191/	(B 26) (B 27)
STATIC WEIGHT TARE DATA	
DATA (BDT(I),I=1,6)/-1.7383,-0.8929,-C.0238,0.8546,1.7515,2.7357/ DATA (BLT(I),I=1,6)/0.1566,0.0113,0.0081,0.0197,0.0835,0.3732/ DATA (SLT(I),I=1,6)/0.1999,-0.0750,0.2250,0.1000,0.0333,0.0374/ AERODYNAMIC TARE DATA	(B 28) (B 29) (B 30)
DATA (SLAT1(1), I=1,6)/477,200,.026,.258,.518,.897/ DATA (SLAT2(1), I=1,6)/482,207,.021,.251,.523,.902/ DATA (SLAT3(I), I=1,6)/433,179.036,.256,.508,.882/ DATA (SLAT4(I), I=1,6)/.198,.152,.137,.144,.178,.273/ DATA (SDAT1(I), I=1,6)/.198,.152,.137,.144,.178,.273/ DATA (SDAT2(I), I=1,6)/.182,.140,.128,.135,.169,.258/ DATA (SDAT4(I), I=1,6)/.182,.140,.128,.135,.169,.258/ DATA (SDAT4(I), I=1,6)/.208,.164,.151,.155,.182,.246/ DO 6 I=1,6 SLAT(1,I)=SLAT1(I) SLAT(2,I)=SLAT2(I) SLAT(2,I)=SLAT2(I) SLAT(2,I)=SLAT3(I) SDAT(1,I)=SDAT1(I) SDAT(1,I)=SDAT1(I) SDAT(2,I)=SDAT2(I) SDAT(2,I)=SDAT2(I) SDAT(3,I)=SDAT3(I) 6 SDAT(4,I)=SDAT4(I) PI=3.14159265358979 RAD=PI/180. SW=7.41125 ST=2.6042 SF=2.*PI/9. REWIND 1 REWIND 3 REWIND 4	<pre>(8 31) (8 32) (8 33) (8 34) (8 35) (8 36) (8 37) (8 38) (8 39) (8 40) (8 40) (8 42) (8 43) (8 44) (8 45) (8 44) (8 45) (8 46) (8 47) (8 48) (8 49) (8 50) (8 51) (8 52) (8 53) (8 55)</pre>

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	READ IN INTE	RFERENCE FACTORS	
	WRITE (6,120)		(B 56)
	READ (5,2024)	WINGLDG	(B 57)
	WRITE (6,2025)	WINGLDG	(8 58)
	DO 20 ITUN=1,	4	(8 59)
	WRITE (6,127)	JTUN(ITUN)	(8 60)
	DO 20 IALPHA=	1,6	(B 61)
	READ (5,121) (	WLFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(1)	(B 62)
	IF (EOF,5) 99	9,22	(8 63)
22	READ (5,121) (	ULFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(2)	(B 64)
	READ (5,121) (	WDFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(3)	(B 65)
	READ (5,121) (	UDFF(ITUN, IALPHA, ICHI), ICHI=1.8), ANAME(4)	(8 66)
	READ (5.122)	WLWF(ITUN, IALPHA), ANAME(5)	(8 67)
	READ (5,122)	ULWF(ITUN, IALPHA), ANAME(6)	(8 68)
	READ (5,2022)		(8 69)
	READ (5,2022)		(8 70)
	READ (5,122)	WLWW(ITUN, IALPHA), ANAME(7)	(B 71)
	READ (5,122)	ULWW(ITUN, IALPHA), ANAME(8)	(8 72)
	READ (5,2022)		(B 73)
	READ (5.2022)		(8 74)
	READ (5.121)	(WLFW(ITUN.TALPHA.ICHI).ICHI=1.8).ANAME(9)	(B 75)
	READ (5.121)	(ULEWLITUN.IALPHA.ICHI).ICHI=1.8).ANAME(10)	(8 76)
	READ (5.121)	(WDFW(ITUN.IALPHA.ICHI).ICHI=1.8).ANAME(11)	(8 77)
	READ (5.121)	(UDEW(ITUN • TAL PHA. ICHI). ICHI=1.8). ANAME(12)	(8 78)
	READ (5.121)	(WEFT(ITUN.IALPHA.ICHI).ICHI=1.8).ANAME(13)	(B 79)
	READ (5.121)	(ULFT(ITUN.IALPHA.ICHI).ICHI=1.8).ANAME(14)	(8.80)
	READ (5,121)	(WDFT(ITUN.IALPHA.ICHI).ICHI=1.8).ANAME(15)	(B 81)
	READ (5.121)	(UDET(ITUN.IALPHA.ICHI).ICHI=1.8).ANAME(16)	(B 82)
	READ (5,122)	(WLWT(ITUN,IALPHA),ANAME(17))	(8 83)
	READ (5.122)	(ULWT(ITUN,IALPHA),ANAME(18))	(B 84)
	READ (5,2022)		(8 85)
	READ (5.2022)		(8 86)
	READ (5,122)	(WLTT(ITUN, IALPHA), ANAME(19))	(8 87)
	READ (5,122)	(ULTT(ITUN,IALPHA),ANAME(20))	(B 88)
	READ (5,2022)		(8 89)
	READ (5,2022)		(B 90)
	READ (5,121)	(WLFP(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(21)	(B 91)
	READ (5,121)	(ULFP(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(22)	(B 92)
	READ (5,121)	(WDFP(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(23)	(8 93)
	READ (5,121)	(UDFP(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(24)	(B 94)
	READ (5,122)	(WLWP(ITUN, IALPHA), ANAME(25))	(8 95)
	READ (5,122)	(ULWP(ITUN, IALPHA), ANAME(26))	(8 96)
	READ (5,122)	(WDWP(ITUN+IALPHA)+ANAME(27))	(8 97)
	READ (5,122)	(UDWP(ITUN, IALPHA), ANAME(28))	(8 98)
	WRITE OUT IN	TERFERENCE FACTORS	
	WRITE (6,123)	(CHI(I),I=1,8)	(8 99)
	WRITE (6,123)	(WLFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(1)	(8 100)
	WRITE (6,123)	(ULFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(2)	(8 101)
	WRITE (6,123)	(WDFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(3)	(8 102)
	WRITE (6,123)	(UDFF(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(4)	(8 103)
	WRITE (6,124)	(WLWF(ITUN, IALPHA), ANAME(5))	(8 104)
	WRITE (6,124)	(ULWF(ITUN, IALPHA), ANAME(6))	(8 105)
	WRITE (6,124)	(WLWW(ITUN, IALPHA), ANAME(7))	(B 106)
	WRITE (6,124)	(ULWW(ITUN, IALPHA), ANAME(8))	(8 107)
	WRITE (6,123)	(WLFW(ITUN, IALPHA, ICHI), ICHI=1,8), ANAME(9)	(8 108)
	WRITE (6,123)	(ULFW(ITUN.IALPHA.ICHI).ICHI=1.8).ANAMF(10)	(B 109)

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	<pre>WRITE (6,123) (WDFW(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(11) WRITE (6,123) (UDFW(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(12) WRITE (6,123) (WLFT(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(13) WRITE (6,123) (ULFT(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(14) WRITE (6,123) (UDFT(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(15) WRITE (6,124) WLWT(ITUN,IALPHA),ANAME(17) WRITE (6,124) ULWT(ITUN,IALPHA),ANAME(18) WRITE (6,124) ULWT(ITUN,IALPHA),ANAME(19) WRITE (6,124) ULTT(ITUN,IALPHA),ANAME(20) WRITE (6,123) (ULFP(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(21) WRITE (6,123) (ULFP(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(22) WRITE (6,123) (ULFP(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(23) WRITE (6,123) (UDFP(ITUN,IALPHA,ICHI),ICHI=1,8),ANAME(24) WRITE (6,123) (UDFP(ITUN,IALPHA,ICHI),ICHI] WRITE (6,123) (UDFP(ITUN,IALPHA</pre>	( B ( B ( B ( B ( B ( B ( B ( B ( B ( B	110) 111) 112) 113) 114) 115) 116) 117) 118) 119) 120) 121) 122) 123)
с	WRITE (6,124) (ULWP(ITUN,IALPHA),ANAME(26)) WRITE (6,124) (WDWP(ITUN,IALPHA),ANAME(27)) WRITE (6,124) (UDWP(ITUN,IALPHA),ANAME(28)) WRITE (6,2022) 20 CONTINUE WRITE (6,110) LINE=0	(B (B (B (B (B (B (B	125) 126) 127) 128) 129) 130) 131)
	READ IN TEST DATA 1 READ (5,100) IRUN, IRPM, ITUN, ITYPE, IFRAME, Q, BALDRAG, BALLIFT, SCLDRAG 1 , SCLLIFT, ALPHA, RPM, DENS ITY IF (EDF, 5) 998, 2 NOTE THAT PITOT MEASUREMENT OF TUNNEL VELOCITY IS CORRECTED ON THE FIRST PASS THROUGH THE INTERFERENCE CALCULATIONS. ON THE SECOND PASS, THE LOADS ARE REDIVIDED ACCORDING TO THE CORRECTED PITOT READING AND ALL THE CORRECTIONS ARE CALCULATED.	( B ( B ( B	132) 133) 134)
	2 ITRIP=1 AQ=Q I=1 IF (ALPHA.GT6.) I=2 IF (ALPHA.GT1.) I=3 IF (ALPHA.GT.4.) I=4 IF (ALPHA.GT.9.) I=5 IF (ALPHA.GT.15.) I=6 LINE=LINE+3 IF (LINE.LE.40) GO TO 75 LINE=O WRITE (6.110) SUBTRACT OUT WEIGHT TARES	( B ( B ( B ( B ( B ( B ( B ( B ( B ( B	135) 136) 137) 138) 139) 140) 141) 142) 144) 144) 145) 146)
C C	75 BALDRAG=BALDRAG-BDT(I) BALLIFT=BALLIFT-BLT(I) SCLLIFT=SCLLIFT-SLT(I) DIVISION OF FORCES BETWEEN WING, FANS, AND TAIL	( B ( B ( B	147) 148) 149)
L	IF (ITYPE-2) 1.1.5 5 TL=SCLLIFT-BALLIFT*COS(ALPHA*RAD)+ BALDRAG*SIN(ALPHA*RAD) TD=SCLDRAG-BALLIFT*SIN(ALPHA*RAD)- BALDRAG*COS(ALPHA*RAD) FL=SCLLIFT-SLAT(ITUN,I)*Q*SW FD=SCLDRAG-SDAT(ITUN,I)*Q*SW	( B ( B ( B ( B ( B	150) 151) 152) 153) 154)

.

		CN=BALLIFT/(Q*ST)	( B	155)
		CA=BALDRAG/(Q*ST)	( B	156)
		IF (ITRIP.EQ.2) GO TO 42	(8	157)
С				
С		WRITE OUT UNCORRECTED DATA		
С				
		WRITE (6,111) JTUN(ITUN), IKUN, IFRAME, IRPM, ALPHA, Q, SCLLIFT, SCLORAG,	(8	158)
	1		(8	1597
			18	1001
		$IF IIRPM \cdot EQ \cdot 0 I ISIAIIU = 40 \cdot 2$		1011
		$IF (IRPM \cdot EQ \cdot IU) = ISIA(IU = 00 \cdot 3)$	10	1021
		$\frac{1}{1} \left( \frac{1}{1} + \frac{1}{1} + \frac{1}{1} \right) = \frac{1}{1} = $		1661
				104/
			10	1441
			10	1471
				1401
		V9V30N-V0V3	18	1601
			10	1701
			18	171)
			10	1721
			18	1731
			18	174)
			(B)	1751
r			10	1121
č		PUNCH CARDS FOR SUBSEQUENT PLOTTING OF UNCORRECTED DATA.		
č				
-		PUNCH 2023, IRUN, IFRAME, RPM, ALPHA, Q, SCLLIFT, CL, CN, QQQJ, VOVJ, DOL,	<b>(</b> B	176)
	1	1 TOTS.ITUN.ISIZE	( B	177)
1(	000	CONTINUE	( B	178)
c				
Ċ		ELIMINATE STATIC THRUST POINTS FROM THE DATA TO BE CORRECTED.		
С				
		IF (Q.LE.0.8) GO TO 1	( B	179)
С				
C		CORRECT FOR SOLID BLOCKAGE		
С				
	42	QFAC=(1.0+EPSILUN(ITUN))**2	(8	1801
			(8	1817
		CLW=SLAT(ITUN, I)/QFAC	(8	1821
		CDW=SDAT([IUN, I])/QFAC	18	1831
		CDO=SDAT(ITUN,3)/QFAC	18	184)
			18	1851
r			(8	1801
C		SOLVE HOMENTUM QUADTLE FOD EVEN ANCLE AND VELOCITY DATLD		
c c		SULVE MUMENIUM QUARIIC FUR SKEW ANGLE AND VELOCIIT RATIO.		
C		XI-0	18	1971
			18	1881
				1901
		IF (CL) 1.44.45	IR	1901
	45		(8	1911
	. ,	VWH=-SQRT(2,*Q*SE/EL)	(B	1921
	46	X=XT+DELTX	(B	1931
	. •	XT=X	( B	1941
				1951
		$IF (XI \cdot GI \cdot I \cdot O2) = GO   IO   SI   IF (XI \cdot LT - O \cdot O1) = GO   IO   S2   IF (XT - LT - O \cdot O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - LT - O - O1) = GO   IO   S2   IF (XI - O - O1) = GO   IO   S2   IF (XI - O - O1) = GO   IO   S2   IF (XI - O - O1) = GO   IO   S2   IF (XI - O1) = GO   IO   S1   IF (XI - O1) = GO   IO   S1   IF (XI - O1) = GO   IO   S2   IF (XI - O1) = GO   IO   S2   IF (XI - O1) = GO   IO   S2   IF (XI - O1) = GO   IO   S1   IF (XI - O1) = GO   IO   S1   IF (XI - O1) = GO   IO   S1   IF (XI - O1) = GO   IO   IF (XI - O1) = GO   IF (XI - O1) = GO   IF (XI - O1) = GO   IO   IF (XI - O1) = GO   IF (XI - O1) = GO   IO   IF (XI - O1) = GO   IF (X$	(B	1951
		IF (XI+GI+I+O2) GO IO 51 IF (XT+LT+-0+O1) GO TO 52 X=(1+0+D0L*D0L)*X*X*X*X+2+0*D0L*VWH*X*X*X+VhH*VWH*X*X-1+0	(B (B	195) 196) 197)
		IF (XI-GI-I-02) GO TO 51 IF (XT-LT0.01) GO TO 52 X=(1.0+DOL*DOL)*X*X*X*X+2.0*DOL*VWH*X*X*X+VhH*VWH*X*X-1.0 IF (ABS(X).LT.0.000001) GO TO 47	(B (B (B	195) 196) 197) 198)

	48 49 47 44 51 52	DELTX=-0.5*DELTX AX=X GO TO 46 WWH=XT VWO=VWH/WWH TANCHI=-VWO-DOL CHIM=ATAN(TANCHI)/RAD CHIEFF=45.0+CHIM/2. GO TO 53 CHIM=CHIEFF=90. VWO=10E10 GO TO 53 WRITE (6,151) IRUN, IFRAME GO TO 1 WRITE (6,152) IRUN, IFRAME GO TO 1	(B 200) (B 201) (B 202) (B 203) (B 204) (B 205) (B 205) (B 207) (B 207) (B 208) (B 209) (B 210) (B 212) (B 213) (B 214) (B 214) (B 215)
с с		SELECT TABLES OF INTERFERENCE FACTORS FOR APPROPRIATE TUNNEL.	
с с	53	DD 60 ICHI=1,8 CWLFP(ICHI)=WLFP(ITUN,I,ICHI) CUUFP(ICHI)=UUFP(ITUN,I,ICHI) CWDFP(ICHI)=UDFP(ITUN,I,ICHI) IF (ITRIP.EQ.1) GD TO 60 CWLFF(ICHI)=WLFF(ITUN,I,ICHI) CULFF(ICHI)=WLFF(ITUN,I,ICHI) CWDFF(ICHI)=UDFF(ITUN,I,ICHI) CWDFF(ICHI)=WLFW(ITUN,I,ICHI) CULFW(ICHI)=WLFW(ITUN,I,ICHI) CULFW(ICHI)=ULFW(ITUN,I,ICHI) CULFW(ICHI)=UDFW(ITUN,I,ICHI) CULFW(ICHI)=UDFW(ITUN,I,ICHI) CUDFW(ICHI)=UDFW(ITUN,I,ICHI) CULFT(ICHI)=WLFT(ITUN,I,ICHI) CULFT(ICHI)=WDFT(ITUN,I,ICHI) CUDFT(ICHI)=UDFT(ITUN,I,ICHI) CUDFT(ICHI)=UDFT(ITUN,I,ICHI) CUDFT(ICHI)=UDFT(ITUN,I,ICHI)	(B 216) (B 217) (B 218) (B 219) (B 220) (B 221) (B 222) (B 223) (B 224) (B 225) (B 226) (B 227) (B 227) (B 227) (B 228) (B 223) (B 231) (B 233)
č		CALL DISCOT (CHIEFF, CHIEFF, CHI, CWLFP, CWLFP, -030, 8, 0, DWLFP)	(B 234)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CULFP, CULFP, -030, 8, 0, DULFP)	(B 235)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CUDFP, CUDFP, -030, 8, 0, DUDFP)	(8 237)
		IF (ITRIP.EQ.1) GO TO 39	(B 238)
		CALL DISCUT (CHIEFF, CHIEFF, CHI, CWLFF, CWLFF, -030, 8, 0, DWLFF)	(B 239)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CWDFF, CWDFF, -030, 8, 0, DWDFF)	(B 241)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CUDFF, CUDFF, -030, 8, 0, DUDFF)	(B 242)
		CALL DISCUT (CHIEFF,CHIEFF,CHI,CULFW,CWLFW,-030,8,0,DULFW) CALL DISCOT (CHIEFF,CHIEFF,CHI,CULFW,CULFW,-030,8,0,DULFW)	(B 244)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CWDFW, CWDFW, -03C, 8, 0, DWDFW)	(B 245)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CUDFW, CUDFW, -030, 8, 0, DUDFW)	(8 246)
		LALL UISCUT (CHIEFF;CHIEFF;CHI;CWLFT;CWLFT;CWLFT;D30;8;0;DWLFT) CALL DISCOT (CHIEFF;CHIEFF;CHI;CHIEF;CHIEF;CHIEF;CHIEF;CHIEFT;	(8 247) (8 248)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CWDFT, CWDFT, -030, 8, 0, DWDFT)	(B 249)
		CALL DISCOT (CHIEFF, CHIEFF, CHI, CUDFT, CUDFT, -030, 8, 0, DUDFT)	(B 250)

CONDECT TUNNEL NATOR MEASUREMENT FOR UNLY SECTOR		
CURRECT TUNNEL PITUT MEASUREMENT FUR WALL EFFECTS		
	10	2511
$\begin{array}{c} 3 \\ 1 \\ + \omega P(TIN,TIX,TUN) \\ \end{array}$		2011
D(P=2 + x)(1 + P + SF / (AT(1T)) + 2 + x)(1) + 2 + x)(1) + 2 + x)(1) + 2 + x + y + y + y + y + y + y + y + y + y	(8	2521
1 + 11  WP(1110.13)(1+(1+0.13))(1+(1+		2551
	(8	2551
	18	2551
0 = 10 + 10 + 10 + 10 + 10 + 10 + 10 + 1	(0	2501
	(0)	2591
	18	2501
IE (ITRIP.EQ.2) GO TO 40	(B)	2601
	(8	2611
	18	2621
	18	2621
	(8	2641
	(0)	2071
	(8	2651
	10	2001
FIND INTERFERENCE VELOCITY RATIOS		
DWF=2.*DWLFF*SF/(AT(ITUN)*VWO)+2.*DWDFF*SF*DOL/(AT(ITUN)*VWO)	( B	267)
1 +WLWF(ITUN,I)*CLW*SW*(-0.25)/AT(ITUN)	(B	268)
DUF=2.*DULFF*SF/(AT(ITUN)*VWO)+2.*DUDFF*SF*DOL/(AT(ITUN)*VWO)	(B	269)
1 + $ULWF(ITUN,I)$ + $CLW$ + $SW$ + (-0,25) / AT(ITUN)	(B	270)
DWW=2.*DWLFW*SF/(AT(ITUN)*VWG)+2.*DWDFW*SF*DDL/(AT(ITUN)*VWD)	(B	271)
1 + $WLWW(ITUN,I)$ +CLW+SW+(-0.25)/AT(ITUN)	(8	272)
DUW=2.*DULFW*SF/(AT(ITUN)*VWO)+2.*DUDFW*SF*DOL/(AT(ITUN)*VWO)	(8	273)
1 +ULWW(ITUN,I)*CLW*SW*(-0.25)/AT(ITUN)	(B	274)
DWT=2.*DWLFT*SF/(AT(ITUN)*VWO)+2.*DWDFT*SF*DDL/(AT(ITUN)*VWO)	(B	275)
1 +WLWT(ITUN,I)*CLW*SW*(-0.25)/AT(ITUN)	(B	276)
2 +WLTT(ITUN,I)*CN*ST*(-0.25)/AT(ITUN)	(8	277)
DUT=2.*DULFT*SF/(AT(ITUN)*VWO)+2.*DUCFT*SF*DOL/(AT(ITUN)*VWO)	( B	278)
1 +ULWT(ITUN,I)*CLW*SW*(-0.25)/AT(ITUN)	(8	279)
2 +ULTT(ITUN,I)*CN*ST*(-0.25)/AT(ITUN)	(8	280)
CALCULATE CORRECTIONS		
TANAL $PE=DWE/(1 + DUE)$	(B	281)
TANAL $PW = DWW/(1 + DUW)$	18	2821
TANAL $PT = DWT / (1 + DUT)$	(B	283)
	(B	284)
	18	2851
	18	2861
	18	2871
	18	2881
	18	2891
AL PCF=AI PHA+DAI PF	18	2901
ALPCW=ALPHA+DALPW	18	2911
	19	2021
	(B	2021
0C₩=0 <b>+</b> 0 <b>∩</b> 0₩	18	2041
	18	2951
	1 B	2961
	18	2971
CNTC=CN/DODT	Í B	2981

c c c

C C C

C C C

	ADJUST FOR DIFFERENCES IN CORRECTIONS AT WING, FANS, AND TAIL Assume a tail efficiency factor (qt/q) of 0.9	
L	DDALPT=DALPT-DALPW CNTCA=(CNTC-0.030*DDALPT*0.9) TDTL=SLIFTC -CN*ST*Q *COS(ALPHA*RAD) 1 +CNTCA*QCW*ST*COS(ALPHA*RAD) CDIUN=CN*CN/(PI*2.40) CDIC=CNTCA*CNTCA/(PI*2.40) SDRAGC=SDRAGC-CDIUN*Q*ST+CDIC*QCW*ST DDAFW=DALPF-DALPW TSTATIC=40.2 IF (IRPM.EQ.10) TSTATIC=66.3 IF (IRPM.EQ.12) TSTATIC=95.0 QJ=TSTATIC/(2.*SF) QCQJ=QCW/QJ QCQJF=QCF/QJ VOVJ=SQRT(QQQJF) DVOVJ=VCVJF-VOVJ SA=SIN(ALPCW*RAD) CA=COS(ALPCW*RAD) TOTL=TOTL+TSTATIC*DDAFW*RAD*SA SDRAGC=SDRAGC-TSTATIC*(DVOVJ+DDAFW*RAD*CA) TOTDDL=SDRAGC/TOTL	<pre>(B 299) (B 300) (B 301) (B 302) (B 303) (B 304) (B 305) (B 306) (B 307) (B 307) (B 309) (B 310) (B 311) (B 312) (B 312) (B 313) (B 314) (B 315) (B 316) (B 317) (B 318) (B 319) (B 320)</pre>
c c	STORE CORRECTED VALUES ON TAPE 1 AND CORRECTIONS ON TAPE 4	
c c	<pre>wRITE (1) ITUN ,IRUN,IFRAME,IRPM,ALPHA,Q,CHIM,DALPF,QOQF, 1 DALPW,QOQW,DALPT,QOQT,CCW,ALPCW,CNTCA,TOTL,TOTDOL,DENSITY wRITE (4) ITUN,IRUN,ALPHA,ALPCW,DALPW,DALPF,DALPT,QOQW,QOQF, 1 QOQT,VOVJ,DVOVJ,VOVJUN GO TO 1 998 ENDFILE 1 RFWIND 1 ENDFILE 4 REWIND 4 wRITE (6,125)</pre>	(B 321) (B 322) (B 323) (B 324) (B 325) (B 326) (B 327) (B 328) (B 329) (B 330)
с с	WRITE OUT CORRECTED VALUES FROM TAPE 1	
С	LINE=0 DD 61 K=1,1000 READ (1) ITUN ,IRUN,IFRAME,IRPM,ALPHA,Q,CHIM,DALPF,QOQF, 1 DALPW,QOQW,DALPT,QOQT,QCW,ALPCW,CNTCA,TOTL,TOTDOL,DENSITY IF (EOF,1) 999,62 62 WRITE (6,126) JTUN(ITUN),IRUN,IFRAME,IRPM,ALPHA,Q,CHIM,DALPF, 1 QOQF,DALPW,QOQW,DALPT,QOQT,QCW,ALPCW,CNTCA,TOTL	(8 331) (8 332) (8 333) (8 334) (8 335) (8 336) (8 337)
Č C	NONDIMENSIONALIZE CORRECTED VALUES.	
_	IF (IRPM.EQ.0) GO TO 1002 IF (IRPM.EQ.8) TSTATIC=40.2 IF (IRPM.EQ.10) TSTATIC=66.3 IF (IRPM.EQ.12) TSTATIC=95.0 QJ=TSTATIC/(2.*SF) QOQJ=QCW/QJ VOVJ=SQRT(QOQJ) TOTS=(TOTL/TSTATIC)*(0.002378/DENSITY)	(B 338) (B 339) (B 340) (B 341) (B 342) (B 342) (B 343) (B 344) (B 345)

i

C       STORE CORRECTED NONDIMENSIONAL VALUES ON TAPE 3         C       WRITE (3) ITUN, IRUN, IRPM, ALPHA, Q, DALPM, QOQW, QCW, ALPCW, CNTCA, TOTL, (6 3         1 000J, VOVJ, TOTS, TOTDOL       (8 3         LINE-LINE+3       (8 3         1002 CONTINUE       (8 3         10102 CONTINUE       (8 3         1002 CONTINUE       (8 3         10102 CONTINUE       (8 3         10102 CONTINUE       (8 3         10102 CONTINUE       (8 3         909 ENDFILE 3       (8 3         REWIND 3       (8 3         C       INTERPOLATE CATA ON TAPE 3 TO OBTAIN CORRECTED VALUES AT FIXED         C       WRITE (6,2020)       (8 3         C       REWIND 3       (8 3         C       INTERPOLATE CATA ON TAPE 3 TO OBTAIN CORRECTED VALUES AT FIXED         C       WRITE (6,2020)       (8 3         IF (16,2020)       (8 3         IF (16,2020)       (8 3         IF (16,01,1AND.TALPC(1),TLIFT(1),TQ00J(1),TV0VJ(1),TLDTS(1),TD0L(1)       (8 3         DO 202 10:1;       ITALPC(1),TCNC(1),TLIFT(1),TQ00J(1),TV0VJ(1),TLDTS(1),TD0L(1)       (8 3         IF (1.E0,1,AND.TALPC(1),CT4.5) GO TO 2002       (8 3         IF (1.E0,1,AND.TALPC(1),CT4.5) LINE+LINE+1       (8 3         CALL DISCOT (AA,A	c	WRITE (6,1001) QOQJ,VOVJ,TOTS,TOTDOL	(8 346)
<pre>write (3) frum, rRPM, ALPHA, Q, DALPW, QOQH, QCW, ALPCW, CNTCA, TOTL, [6 3 1 000 J, VOVJ, TOTS, TOTDOL</pre>	C	STORE CORRECTED NONDIMENSIONAL VALUES ON TAPE 3	
C INTERPOLATE CATA ON TAPE 3 TO OBTAIN CORRECTED VALUES AT FIXED C CORRECTED ANGLES OF ATTACK. C () WRITE (6,2020) (8 3 LINE=0 (8 3 2003 D0 2001 I=1,6 (8 3 READ (3) ITUN,IRUN,IRPM,TALP(I),TQ(I),TDALPW(I),TQOQW(I),TQC(I), (8 3 I TALPC(I),TCNC(I),TLIFT(I),TQOQJ(I),TVOVJ(I),TLOTS(I),TDDL(I) (8 3 2001 CONTINUE (8 3 DO 2002 I=1,5 (8 3 IF (I.EQ.1.AND.TALPC(I).GT4.5) GO TO 2002 (8 3 IF (I.EQ.1.AND.TALPC(I).GT4.5) LINE=LINE+1 (8 3 AA=A(I) (8 3 CALL DISCOT (AA,AA,TALPC,TALP,TALP,-010,6,0,0UN) (8 3 CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,0UN) (8 3 CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,0UN) (8 3 CALL DISCOT (AA,AA,TALPC,TQOQH,TCOQH,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TQC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TQC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TCC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TCC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TCC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TCC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TCNC,TCNC,TON,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TQO,TQC,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TCNC,TCNC,TON,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOU,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,TO,0,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 CALL DISCOT (AA,AA,TALPC,TOO,TCN,TOO,-010,6,0,0CA) (8 3 DOTS=DL+LOTS (8 3 DOTS=DL+LOTS (8 3 DOTS=DL+LOTS (8 3 I CNC,DL,DOTS,AUN,QUN,DALH,QOO (8 3 I SIZE=1 (8 3 I PUNCH 2023, IRUN,IRUN,RPM,AA,QC,TLFT,CL,CNC,	100 6 99	<pre>wRITE (3) ITUN, IRUN, IRPM, ALPHA, Q, DALPW, QOQW, QCW, ALPCW, CNTCA, TOTL, 1 QOQJ, VOVJ, TOTS, TOTDOL LINE=LINE+3 2 CONTINUE IF (LINE.LT.40) GO TO 61 LINE=0 wRITE (6,125) 1 CONTINUE 9 ENDFILE 3 REWIND 3</pre>	(B 347) (B 348) (B 349) (B 350) (B 351) (B 352) (B 353) (B 354) (B 355) (B 356)
C WRITE (6,2020) (8 3 LINE=0 (8 3 2003 DD 2001 I=1,6 (8 3 READ (3) ITUN,IRUN,IRPM,TALP(I),TO(I),TDALPW(I),TQOQW(I),TQC(I), (8 3 I TALPC(I),TCNC(I),TLIFT(I),TQOQJ(I),TVOVJ(I),TLOTS(I),TDOL(I) (8 3 IF (EDF,3) 9999,2001 (8 3 DD 2002 I=1,5 (8 3 IF (I.EQ.I.AND.TALPC(I).GT4.5) GO TO 2002 (8 3 IF (I.EQ.I.AND.TALPC(I).GT4.5) LINE=LINE+1 (8 3 AA=A(I) (8 3 CALL DISCOT (AA,AA,TALPC,TALP,TALP,-010,6,0,AUN) (8 3 CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,QUN) (8 3 CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,QUN) (8 3 CALL DISCOT (AA,AA,TALPC,TQOQW,TCOV,0,0,DALW) (8 3 CALL DISCOT (AA,AA,TALPC,TQC,TQC,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQC,TQC,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQC,TQC,-010,6,0,CC) (8 3 CALL DISCOT (AA,AA,TALPC,TQC,TQC,-010,6,0,CC) (8 3 CALL DISCOT (AA,AA,TALPC,TQU,TTVOVJ,-010,6,0,CC) (8 3 CALL DISCOT (AA,AA,TALPC,TQOQJ,TQOQJ,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOQJ,TQOQJ,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOQJ,TQOQJ,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOV,TUVTJ,-010,6,0,CC) (8 3 CALL DISCOT (AA,AA,TALPC,TQOV,TUTS,-010,6,0,CVJ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOJ,TQOQJ,-010,6,0,QOQ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOJ,TQOJ,-010,6,0,CUJ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOJ,TQOJ,-010,6,0,CUJ) (8 3 CALL DISCOT (AA,AA,TALPC,TQOJ,TDOL,-010,6,0,CL) (8 3 DOTS=DL*LOTS (8 3 DOTS=DL*LOTS (8 3 C WRITE (0,2021) JTUN(ITUN),IRUN,IRPM,AA,QC,QOOJ,VOVJ,TLFT,LOTS, (8 3 I SIZE=1 (8 3 PUNCH 2023, IRUN,IRUN,RPM,AA,QC,TLFT,CL,CNC,QQQJ,VOVJ,DL,LOTS, (8 3 I ITUN,ISIZE (8 3 LINE=LINE+1 (1) I TUN,ISIZE (8 3 LINE=LINE+1 (1) I TUN,ISIZE (8 3 LINE=LINE+1 (1) I CLNE+LT.35) GO TO 2002 (8 3 C C		INTERPOLATE CATA ON TAPE 3 TO OBTAIN CORRECTED VALUES AT FIXED CORRECTED ANGLES OF ATTACK.	
C WRITE OUT, AND PUNCH FOR SUBSEQUENT PLOTTING, THE CORRECTED NONDIMENSIONAL VALUES. WRITE (6,2021) JTUN(ITUN),IRUN,IRPM,AA,QC,QOQJ,VOVJ,TLFT,LOTS, (B 3 I CNC,DL,DOTS,AUN,QUN,DALW,QOQ (B 3 RPM=1030.*FLOAT(IRPM) (B 3 CL=TLFT/(QC*SW) (B 3 I SIZE=1 (B 3 PUNCH 2023, IRUN,IRUN,RPM,AA,QC,TLFT,CL,CNC,QOQJ,VOVJ,DL,LOTS, (B 3 I ITUN,ISIZE (B 3 LINE=LINE+1 (B 3 IF (LINE.LT.35) GD TO 2002 (B 3	200 200	<pre>wRITE (6,2020) LINE=0 3 D0 2001 I=1,6 READ (3) ITUN,IRUN,IRPM,TALP(I),TQ(I),TDALPW(I),TQ0QW(I),TQC(I), 1 TALPC(I),TCNC(I),TLIFT(I),TQ0QJ(I),TV0VJ(I),TLOTS(I),TDOL(I) IF (E0F,3) 9999,2001 1 CONTINUE D0 2002 I=1,5 IF (I.EQ.1.AND.TALPC(1).GT4.5) G0 T0 2002 IF (I.EQ.1.AND.TALPC(I).GT4.5) LINE=LINE+1 AA=A(I) CALL DISCOT (AA,AA,TALPC,TALP,TALP,-010,6,0,AUN) CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,QUN) CALL DISCOT (AA,AA,TALPC,TQ,TQ,-010,6,0,QUN) CALL DISCOT (AA,AA,TALPC,TDALPW,TDALPW,-010,6,0,DALW) CALL DISCOT (AA,AA,TALPC,TQC,TQC,-010,6,0,QC) CALL DISCOT (AA,AA,TALPC,TCNC,TCNC,-010,6,0,CNC) CALL DISCOT (AA,AA,TALPC,TLIFT,TLIFT,-010,6,0,TLFT) CALL DISCOT (AA,AA,TALPC,TQOQJ,TQOQJ,-010,6,0,QQJ) CALL DISCOT (AA,AA,TALPC,TUOT,TUOT,-010,6,0,UC) CALL DISCOT (AA,AA,TALPC,TUOT,TUOT,-010,6,0,UC) CALL DISCOT (AA,AA,TALPC,TLOT,TUOT,-010,6,0,UC) CALL DISCOT (AA,AA,TALPC,TLOT,TUOT,-010,6,0,UC) CALL DISCOT (AA,AA,TALPC,TUOT,TUOT,-010,6,0,UC) CALL DISCOT (AA,AA,TALPC,TUOT,TUOT,-010,6,0,UC) DOTS=DL*LOTS</pre>	<pre>(B 357) (B 358) (B 359) (B 360) (B 361) (B 362) (B 363) (B 364) (B 365) (B 366) (B 367) (B 367) (B 370) (B 371) (B 374) (B 375) (B 377) (B 377) (B 377) (B 377)</pre>
LINE=0 (B 3 WRITE (6,2020) (B 3 2002 CONTINUE (B 3 WRITE (6,2022) (B 3 GO TO 2003 (B 3	c c c 200	<pre>wRITE OUT, AND PUNCH FOR SUBSEQUENT PLOTTING, THE CORRECTED NONDIMENSIONAL VALUES. wRITE (6,2021) JTUN(ITUN), IRUN, IRPM, AA, QC, QDQJ, VOVJ, TLFT, LOTS, CNC, DL, DOTS, AUN, QUN, DALW, QOQ RPM=1030.*FLOAT(IRPM) CL=TLFT/(QC*SW) ISIZE=1 PUNCH 2023, IRUN, IRUN, RPM, AA, QC, TLFT, CL, CNC, QOQJ, VOVJ, DL, LOTS, ITUN, ISIZE LINE=LINE+1 IF (LINE.LT.35) GD TO 2002 LINE=0 wRITE (6,2020) CONTINUE wRITE (6,2022) GO TO 2003</pre>	(B 380) (B 381) (B 382) (B 383) (B 384) (B 385) (B 386) (B 387) (B 388) (B 389) (B 390) (B 391) (B 393)

c	WRITE (6,3008) ICOR=1 'LINE=0	(B (B (B	395) 396) 397)
	INTERPOLATE DATA ON TAPE 4 TO OBTAIN VALUES OF CORRECTION ANGLES AND VELOCITY RATIOS AT FIXED CORRECTED ANGLES OF ATTACK.		
3003 1 3001 C C	DO 3001 I=1.6 READ (4) ITUN, IRUN, TALP(I), TALPC(I), TDALPW(I), TDALPF(I), TDALPT(I), 1 TCOQW(I), TQOQF(I), TQOQT(I), TVOVJ(I), TDVOVJ(I), VOVJUN IF (EOF,4) 3002, 3001 CONTINUE DO 3004 I=1,5 AA=A(I) CALL DISCOT (AA, AA, TALPC, TALP, TALP, -010,6,0, AUN) CALL DISCOT (AA, AA, TALPC, TDALP, TDALPW, -010,6,0, DALW) CALL DISCOT (AA, AA, TALPC, TDALPF, TDALPF, -010,6,0, DALF) CALL DISCOT (AA, AA, TALPC, TDALPF, TDALPF, -010,6,0, DALF) CALL DISCOT (AA, AA, TALPC, TQQW, TQQW, -010,6,0, TQW) CALL DISCOT (AA, AA, TALPC, TQQQF, TQQQF, -010,6,0, TQF) CALL DISCOT (AA, AA, TALPC, TQQQT, TQQQT, -010,6,0, TQF) CALL DISCOT (AA, AA, TALPC, TQOQT, TQQQT, -010,6,0, TQT) CALL DISCOT (AA, AA, TALPC, TDVOVJ, TVOVJ, -010,6,0, DVOVJ) DIT=DALT-DALW DIF=DALF-DALW QFQQC=TQF/TQW LINE=LINE+1 WRITE OUT, ANC PUNCH FOR SUBSEQUENT PLCTTING, THE CORRECTION	(8) (8) (8) (8) (8) (8) (8) (8) (8) (8)	398) 399) 400) 401) 402) 403) 404) 405) 406) 405) 406) 407) 408) 407) 408) 410) 411) 412) 413) 414) 415) 416) 416) 417) 418)
с С	ANGLES AND VELOCITY RATIOS.	4.0	(10)
1 3004	WRITE (8,3005) JTUN(ITUN),IRUN,AA,AUN,VUVJ,DALW,DIT,DIF,TUW, QTOOC,QFOQC,DVOVJ PUNCH 3006, IRUN,AA,VOVJ,DALW,DIT,DIF,TQW,QTOQC,QFOQC,DVOVJ, I ITUN,ICOR CONTINUE WRITE (6,3007) LINE=LINE+1 IF (LINE.LT.36) GO TO 3003 GO TO 9999	(B (B (B (B (B (B (B (B	419) 420) 421) 422) 423) 424) 425) 425) 426) 427)
	WRITE OUT, AND PUNCH FOR SUBSEQUENT PLCTTING, THE CORRECTION ANGLES AND VELOCITY RATIOS IN THE UNCORRECTED DATA.		
3002 3011 3014 1 3013	REWIND 4 ICOR=3 WRITE (6,3009) WRITE (6,3008) LINE=0 DO 3010 I=1,6 PEAD (4) ITUN, IRUN, ALPHA, ALPCW, DALPW, DALPF, DALPT, QDQW, QDQF, I QOQT, VOVJ, DVOVJ, VOVJUN IF (EDF,4) 3012,3013 DIT=DALPT-DALPW QTOQC=QOQT/QDQW QFOQC=QOQF/QOQW	(B (B (B (B (B (B (B (B (B (B (B (B (B))))))))	428) 429) 430) 431) 432) 433) 434) 435) 436) 435) 436) 437) 438) 439) 440)

	(8	441)
WRITE (6.3005) JTUN(ITUN), IRUN, ALPCW, ALPHA, VOVJ, DALPW, DIT, DIF,	(8	442)
	(8	443)
	(8	444)
	1 B	445)
	(8	446)
	(B)	447)
	10	4411
LINE=LINE+1		4401
IF (LINE.LI.3) GU IU 3014		4471
GO TO 3011	18	4501
3012 STOP	(B	451)
	19	4521
LUU FURMAI (212,211,144,73,24,770,24,777,11,70,0)	10	4521
110 FURMAL (IHI/ 3/X*FAN-IN-WING LESIS IN AMES / ATO TONNEL (AARL LESI	(D)	4221
1 NO. 6)*//4X*1UNNEL KUN FRAME RPM*5X*ALPHA*4X*U*6X*L*8X*D*	(8	4541
2 5X*CN(T) CA(T)*4X*L(WB)*4X*D(WB)*5X*L(F)*5X*D(F)*2X7)	(B	4551
111 FORMAT (2X,A10,2X,I2,I6,I4*,000*F7.1,F7.2,2F9.3,2F8.4,4F9.3,F8.2)	(B	456)
120 FORMAT (1H1//30X*INTERFERENCE FACTORS*)	( B	457)
121 FORMAT (8F7.4,16X,A8)	( B	458)
122 FORMAT (49X,F7.4,16X,A8)	( B	4591
123 FORMAT (8F10.4,5X,A10)	(8	460)
124 FORMAT (70X,F10.4,5X,A10)	( B	461)
125 FORMAT (1H1//50X*CORRECTED DATA*//4X*TUNNEL RUN FRAME RPM*5X	( B	462)
1 *ALPHA*4X*Q*3X*CHI DALPF QOQF DALPW QOQW DALPT QOQT*	( B	4631
$2 + 4 \times 0$ CW ALPC(W) + 4 X + CNTC TOT L + /)	( B	464)
126 FORMAT (2X.A10.I4.I6.I4*.000*F7.1.2F7.2.3(F7.2.F7.4)	(8	465)
1 • F7-2-F8-2-F9-4-F8-2)	( B	466)
127 FORMAT (//35% A10//)	(8	467)
151 FORMAT (10/#APUN #12#, FRAME #14# FAILS, WWH TOO GREAT#)	(B	468)
152 FORMAT (10X+ROW) $\pm 124$ FRAME $\pm 144$ FATLS, WHI TOD SMALL $\pm$ )	(B	4691
152 FORMAT (2004) THEY TRADE STATE ALLS AND THE FALLS AND	18	470)
100 FORMAT (204+001 - +F0.27) CHI = +F0.27) $4/80$ -+10-27 $0/2$ 10-7)	18	4707
$1001 \ \text{FURMAL} (100+0/0) = +F0.444 \ \text{VV} = +F0.444 \ \text{V} = -F0.444 $	18	4721
	10	4721
202J FURMAT (IHI//49X*INTERPOLATED CURRECTED VALUES+//12X+TUNNEL+JATRUN	10	4761
	10	4751
23X#0/L#3X#0/IS#2X#ALPHA#4X#0#4X#0 ALP#2X#0/0#/1	10	4721
2021 FURMAT (10X,A10,16,13*000*2F7.2,2F6.4,F7.2,2F7.3,F6.3,F7.3,F6.2,	10	4101
1 2F7.2,F8.4)	18	4(1)
2022 FORMAT ()	(8	4/8)
2023 FORMAT (13,14,F5.0,3F6.2,2F7.4,4F6.4,211)	(8	479)
2024 FORMAT (A8)	(8	480)
2025 FORMAT (/30X,A8* WING LOADING*)	( B	481)
3000 FORMAT (1H1//43X*CORRECTIONS ACCORDING TO CORRECTED ALPHAS*/)	( B	482)
3005 FORMAT (4X,A10,I7,2F10.2,F10.4,3F10.2,4F10.4)	(8	483)
3006 FORMAT (13,9F8.4,211)	( B	484)
3007 FORMAT ()	( B	485)
3008 FORMAT (6X*TUNNEL*6X*RUN*4X*ALPHAC*4X*ALPHAU*6X*V/VJ*4X*D ALPW*	( B	486)
1 6X*D IT*6X*D	( B	487)
3039 FORMAT (1H1//42X*CORRECTIONS ACCORDING TO UNCORRECTED ALPHAS*/)	( B	488)
END	( B	4891

	SUBROUTINE DISCOT (XA,ZA,TABX,TABY,TABZ,NC,NY,NZ,ANS)	(B 490	)
	DIMENSION TABY (2), TABY (2), TAB7 (2)	18 491	1
		(0 1)1	Ś.
	DIMENSION NEXTON NETTON	10 492	
	CALL UNS (NC,1A,1DX,1DZ,1MS)	(B 493	
	IF(NZ-1) 5,5,10	(B 494	)
5	CALL DISSER (XA.TABX(1).1.NY.IDX.NN)	(B 495	1
-		18 496	, i
	CALL LAGRAN (XA+TABXINNI+TABTINNI+NNI+ANS)	18 497	1
	GOTO 70	(B 498	
10	ZARG=ZA	(B 499	))
	[P]X=[DX+]	(B 5)0	1
		(B 501	ŝ
	IF (IA) 15,25,15	18 502	
15	IF (ZARG-TABZ(NZ)) 25,25,20	(8 503	)
20	ZARG=TABZ(NZ)	(B 504	)
25	CALL DISSER (ZARG.TARZ(1)).1.NZ.IDZ.NPZ)	(8.505	1
22		18 506	, i
	NPZL=NPZ+IUZ	(8 507	1
	I = 1	(B 508	()
	IF (IMS) 30-30-40	(B 509	) }
3.)	CALL DISSER (XA-TABY(1)-1-NY-IDY-NPY(1))	(B 510	1
50		(0 51)	1
	00 35 JJ = NPZ + NPZL		
	NPY(I) = (JJ-I) * NX + NPX(I)	(8 512	; <b>}</b>
	NPX(I) = NPX(1)	(B 513	5)
35	I = I + 1	(B 514	)
• •	6010.50	(B 515	÷.
		(D 514	
<b>4</b> 0			
	$1 S = (JJ - I) \times NX + I$	(8 517	1
	CALL DISSER (XA,TABX(1),IS,NX,IDX,NPX(I))	(B 518	))
	NPY(I) = NPX(I)	(B 519	))
45	1=1+1	(B 520	11
50		(0.52)	
20		(0.521	
		(8 522	1
	NLOCY=NPY(LL)	(B 523	3)
55	CALL LAGRAN(XA,TABX(NLOC),TABY(NLOCY), [P1X, YY(LL))	LB 524	• )
	CALL LAGRAN (ZARG. TARZ (NPZ). YY(1). TP17. ANS)	(8 525	51
70		18 524	
10	RE LOKN	18 920	
	END	(8 527	)
	SUBROUTINE UNS (IC, IA, IDX, IDZ, IMS)	(B 528	3)
	IF (IC) 5.5.10	(B 529	))
5		10 520	•••
,		10 530	
	NC = -1C	(8 231	. )
	GOTO 15	(8 532	2)
10	IMS=0	(B 533	3)
		(R 534	1
15	LE (NC-100) 20.25.25	10 525	5
20		10 55	/ <b> </b>
20		(8 536	21
	GUTO 30	(B 537	1)
25	I A = 1	LB 538	31
	NC=NC-100	(B 539	<b>;</b> ;
30		19 547	11
50			
		(8 54)	11
	RETURN	(8 542	2)
	END	(B 543	3)

```
SUBROUTINE DISSER (XA, TAB, I, NX, ID, NPX)
                                                                           (B 544)
   DIMENSION TAB(2)
                                                                            (B 545)
   NPT=ID+1
                                                                            (B 546)
   NPB=NPT/2
                                                                            (B 547)
   NPU=NPT-NPB
                                                                            (B 548)
   IF (NX-NPT)
                10,5,10
                                                                            (8 549)
 5 NPX=I
                                                                            (B 550)
   RETURN
                                                                            (8 551)
10 NLCW=I+NPB
                                                                            (8 552)
   NUPP = I + NX - (NPU + 1)
                                                                            (B 553)
   DO 15 II=NLOW, NUPP
                                                                            (B 554)
   NLOC=II
                                                                            (8 555)
   IF (TAB(II)-XA)
                      15,20,20
                                                                            (B 556)
15 CONTINUE
                                                                            (8 557)
   NPX=NUPP-NPB+1
                                                                            (8 558)
   RETURN
                                                                            (B 559)
20 NL=NLOC-NPB
                                                                            (B 560)
   NU=NL+ID
                                                                            (B 561)
   DO 25 JJ=NL,NU
                                                                            (B 562)
   NDIS=JJ
                                                                            (8 563)
   IF (TAB(JJ)-TAB(JJ+1)) 25,30,25
                                                                            (B 564)
25 CONTINUE
                                                                            (B 565)
   NPX=NL
                                                                            (B 566)
   RETURN
                                                                            (B 567)
30 IF (TAB(NDIS)-XA) 40,35,35
                                                                            (B 568)
35 NPX=NDIS-ID
                                                                            (B 569)
   RETURN
                                                                            (B 570)
40 NPX=NDIS+1
                                                                            (B 571)
   RETURN
                                                                            (B 572)
   END
                                                                            (B 573)
   SUBROUTINE LAGRAN (XA, X, Y, N, ANS)
                                                                           (B 574)
   DIMENSION X(2),Y(2)
                                                                            (B 575)
   SUM=0.0
                                                                            (8 576)
   DO 3 I=1.N
                                                                            (B 577)
   PROD=Y(I)
                                                                            (B 578)
   DO 2 J=1,N
                                                                            (B 579)
   A = X(I) - X(J)
                                                                            (B 580)
   IF (A)
          1,2,1
                                                                            (B 581)
1 B = (XA - X(J))/A
                                                                            (B 582)
   PROD=PROD*B
                                                                            (B 583)
 2 CONTINUE
                                                                            (B 584)
 3 SUM=SUM+PROD
                                                                            (8 585)
   ANS=SUM
                                                                            (B 586)
   RETURN
                                                                            (B 587)
   END
                                                                            (8 588)
```

#### INTERFERENCE FACTORS

#### ELLIPTIC WING LOADING

#### 2-1 ROUNE

#### ALPHA = -10

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	SC.C000	Сні
-1-0118	8363	6547	5091	4268	4055	4103	4C63	-1C1WLFF
- 3881	. 4843	.4934	.4245	.304E	.1654	•0271	0776	-101ULFF
- 7516	5460	3678	2319	1354	0619	.0069	.0776	-1C1WDFF
11556	.1213	.1216	.0855	.0416	.0091	CC46	C.CO00	-101UDFF
•0550	••••						6552	-101wLwF
							0592	-101ULwF
							- 5024	-101wLww
							(383	-1016L h m
7761	6036	- 1117	- 3449	- 2651	- 2843	2 + 7 3	- 2821	-101mLFW
//51				2483	.1350	-02+0	0508	-1010LFW
.4994	.5029	•4490	• 3 3 7 0	- 1011	- (51)	- 0027	6469	-101WDEW
5651	3865	2555	1044	1011			1563	-10100Ew
.3149	. 3108	•2145	•2303	. 1914	- 9017		- 9208	-101WLET
4//2	6266	1998	9492	9402	0037	0552	- 0266	-10101 61
2087	1603	0490	.1622	• 421 7	. 4299	-1926	0204	-1010007
4997	53/1	5528	5183	3967	2341	1020	• (43)	-1010061
8854	8167	7219	0086	5218	5313	5674	0020	
							-1.1440	-IUINENI
							.035	-ICIULWI
							5791	-ICIWLII
							4450	-ICIULII
033b	0327	0322	0324	0329	35ذ0. –	0335	0339	-101WLFP
C+28	0408	0397	0394	0395	0345	0370	0083	-101ULFP
0910	0012	0729	0657	0591	0520	046t	C405	-101wDFP
.8766	.8740	.8715	.8694	.8677	. 8068	.8666	.8674	-101UDFP
							6684	-101 WL NP
							0408	-101UL WP
							0493	-1ClwDwP
							1.1376	-1CluDwP
26.0000	<b>1</b> 0 - 0000	40,0000	50,0000	£0.000C	70.0000	4 <b>0.0</b> 000	90.0000	CH1
20.0000	30.0000	40.0000						
-1.0023	8370	6550	5092	4270	4059	4107	4064	-51WLFF
.3865	• 4835	•4929	.4238	.3037	.1642	•0258	0788	-51ULFF
7506	5452	3667	2307	1342	0608	.0081	•C788	-51wDFF
.0537	.1212	.1208	.0848	.C41C	.0687	0048	<b>c.</b> ccco	-51UDFF
							6602	-51wLwF
							C758	-51ULWF
							5048	-51WLWW
							0688	-51ULWW
7023	5879	4451	3475	2978	2060	2876	2811	-51wLfw
.4752	.4811	.4311	.3435	.2363	.1237	.C146	0620	-51ULFw
	3676	2399	1504	C875	0376	.0106	.C600	-51 wDFw
2954	2959	.2647	.2240	.1875	.1641	.1548	.1588	-51UDFw
- 5010	- 6778	- 8915	-1.0795	-1.0384	8222	8215	8819	-51WLFT
2785	- 2400	1262	.1221	. 4447	4445	.2019	C140	-51ULFT
- 5302	- 5918	- 62.58	- 5592	- 4545	2621	1215	.0225	-51m0FT
- Cinf	8614	7557	- 6181	- 5066	5245	5773	5932	-51UDWT
• , , , 0 0 0	.0014	••••		.,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,	• 52 7 5		-1.1188	-51wLnT
							.084	-514L nT
							- 4037	-51 mi TT
							-, 2463	=5111 TT
1.7.15	- 0333	- 0316	_ 0210	- 0334	- (22)	- (330	- 6334	-51WLEP
1335	0323	- 0303	- 0340	- 0325	1331	(350	- 0076	-5101 60
0403	0403	0:52		- 0590	0389	0303	- 0404	= JIUCFF _ 61W1FP
0403	0805	0724	0003	0588	0525	0464	0404	
• 8104	• 8/3/	•8/11	.8089	.00/2	.8061	•8025	• 8000	- 5100FF
							06/3	- 514644
							0394	- 210L WP
							0492	- 21 WUWP
							1.1361	- SILUWP

#### ALPHA = 0

20.000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
9887	8382	6555	5093	427C	4061	4108	4065	01WLFF
. 3858	.4837	.4933	.4238	.3035	.1638	.0254	(792	CIULEE
7513	5455	3666	2304	1338	0604	.0084	.0792	01 mDFF
.0529	.1201	.12C8	.0847	.0388	.0080	0648	C.C000	01UDFF
							6642	CINLWF
							0917	OTULWE
							- 5079	
							0990	0101.00
7520	5851	4468	3509	3011	2884	2887	- 2808	Cluben
.4528	.4611	.4150	.3307	.2255	.1133	.0040	0727	0101 FW
5171	3507	2259	1376	0749	- 0249	.0233	. (727	GINDEN
.2763	.2812	.2545	.2173	.1838	-162U	.1544	.1603	OLUDEW
5261	7369	-1.0049	-1.2490	-1.1605	- 8497	8191	- 8525	01 #1 FT
3580	3335	2187	.0809	.4898	• 4662	.2108	0015	
~.5773	6556	7143	7000	- 522 9	- 2936	1424	.0015	OIWDET
5922	9170	7994	6298	4860	5115	- 5686	5861	0100FT
							-1.0952	01.01.01
							. 0132	C1LL LT
							- 4489	
							0934	
0330	0319	0314	0315	0320	+ . (326	- 0325	- 0328	C141ED
0417	0397	0387	0383	0384	(383	(357	070	
(895	0799	0719	- 0648	0584	0523	6462	0404	010000
.8761	.8733	.8707	.8683	- 866.5	. 8654	8651	8656	
			• • • • • •		• 0024	*UCJI	- 0662	010077
							0382	
								0160 #
							1.1345	010000
								01004
			ALPHA =	5				
20.0000	30.0000	40.0000	50.0000	60.CCOC	70.000	80.0000	90.0000	СНІ
-1.0162	8400	6562	5093	4570	- 4060	4108	- 4664	51 M EE
.3860	.4849	.4947	.4248	- 204C	.1642	0258	0788	514166
7536	5471	3675	- 2308	- 1343	- 0600	-0040	.0788	5160255
. 0532	.1210	.1217	.0852	. 0411	.0087	0048	C. C000	511066
							6671	510010
							- 1069	5111 wF
							-,5117	51 Wi hh
							1287	5144.66
7442	5840	4496	3552	3051	2914	2905	- 2815	5161.56
•4320	.4430	.4005	. 3194	.2159	.1040	0058	0828	51UL FW
4972	3359	2135	1259	0633	0131	.0352	. ( 84 7	51 #DF#
.2574	.2666	.2441	.2102	.1790	.159u	.1530	1606	51UDEw
5518	8048	-1.1464	-1.4726	-1.3066	- 8848	- 8254	8320	51 mi FT
4506	4467	3341	.0381	-5656	4948	.2220	.0113	510LET
6235	7301	8227	- 8280	6030	3286	- 1652	0198	51WDET
-1.0606	9872	- 8567	- 6434	4570	- 4974	5610	5807	511.0FT
				• • • • • •			-1.0729	51 81 87
							.0181	5111
							- 4369	51 51 51
							. 0409	5110 TT
0326	0314	- 03.09	0310	0315	0321	0319	(323	516150
0411	0391	-,0381	0377	0378	0376	-,0351	(^^3	516650
- 0888	- 0793	0714	0645	0582	0521	0461	(404	516050
.8758	.8720	.8701	.8677	. 8658	- 2561	10701 18447	• UTUT . 8646	511060
		-0101		•0000	- 30+0	-0172	(652	516660
							(369	5110 60
							- (493	51 61 60
							1,1328	511042
							******	210085

#### ALPHA = 10

20.000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.000	СНІ
-1.0196	8422	6569	5092	4266	4056	4105	4063	101WLFF
- 3871	.4872	.4971	.4266	.3053	.1654	.0270	C776	1Clulff
- 6830	5499	3692	2321	1355	0620	-CC68	.0776	1C1wDFF
.0546	.1229	.1233	•0863	.0417	.0091	0046	C.CO00	101UDFF
							6687	1C1WLWF
							1213	101ULWF
							5160	1ClwLww
							1579	101ULWW
7385	5844	4535	3602	3098	2950	2931	2829	101wLFw
. 4129	.42.65	.3877	.3096	.2076	.0957	0147	(922	1 CI UL FW
- 4800	3230	2026	1155	(527	0023	.0464	.0962	101wUFw
- 2369	2522	.2336	.2027	.1736	.1551	.1506	.1598	101UDFW
5771	8625	-1.3255	-1.7749	-1.4724	9266	8399	8202	101WLFT
- 5604	5871	- 4841	- 0044	.6840	.5301	.2378	.0247	101ULFT
	8173	9576	9946	6948	3680	1904	C420	101wDFT
-1 1445	-1.0768	- 9373	6581	4158	4837	5545	5769	101UDFT
101772	1.0700	••••••					-1.0523	1ClwLwT
							•C234	101ULWT
							4547	101 mL TT
							.1767	101ULTT
	0309	0305	0306	0311	0316	0314	0318	101 WL FP
0404	0385	0374	0371	0372	C370	0344	C057	1C1ULFP
- 0881	0787	- 0709	- 0641	0579	0520	0461	0404	1C1WDFP
- • 0001 8753	.8723	.8695	-8670	.8651	.8638	.8633	.8636	101UDFP
•0100	.0125	•••••					0642	101wLwP
							C358	101UL wP
							C495	101 wD wP
							1.1311	1C1UDhP

ALPHA = 16

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
-1-0249	8454	6581	5090	426C	405U	4100	4060	161wLF+
.3896	.4912	.5011	•4299	.3080	.1678	.0296	C751	161ULFF
- 7639	5547	3726	2348	1379	0644	•0044	.0751	161wDFF
.0578	.1265	.1263	.0883	.0427	.0100	0042	C.COOO	161UDFF
							6689	161wLwF
							1374	161ULwF
							5219	161WLWW
							1919	161UL WW
- 7344	5869	4596	3672	3161	3002	2971	2858	161hLFh
.3921	40.91	.3745	.2998	.1994	. 6873	0241	1026	161ULFW
- 4629	31.04	1918	1049	0416	.0094	•0586	.1089	161wDFw
.2170	2350	.2208	.1931	.1662	.1492	.1463	.1572	161UDFw
- 6040	- 9897	-1.6116	-2.3032	-1.6787	9850	8686	8179	161wLFT
7229	8075	7369	0439	.9005	.5804	.2609	.0421	161ULFT
7436	9418	-1.1686	-1.2737	8174	4212	2246	0707	161wDFT
-1.2716	-1.2202	-1.0644	6723	3438	4657	5481	5747	161LDFT
							-1.0308	161WLwT
							•C306	161UL#T
							5208	161wLTT
							.3715	161ULTT
0315	03 04	0299	0300	C305	0310	0368	0312	161wLFP
0396	0377	0367	0363	0364	0362	0337	051	161ULFP
0873	0781	0705	0638	0577	0519	0461	- <u>.</u> C405	161wDFP
. 8746	.8716	.8686	.8661	.8640	. 8627	.8621	.8622	161UDFP
							0630	161WLWP
							0345	161UL mP
							0498	161WDWP
							1.1288	161UDwP

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## 2-1 RECT.

#### ALPHA = -10

20.000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	96.000	CHI
-1.1143	9164	7110	5449	4485	4196	4207	4160	-102WLFF
•4351	• 54 30	•5535	.4768	•3431	.1870	.0330	0862	-102ULFF
8233	5931	3939	2426	1365	0577	.0138	.0862	-102wDEE
.1121	.1757	.1666	.1173	.0598	.0162	0038	0000.0	-1.02UUEE
							- 6718	-1020011
							- 6450	-1020000
								-IUZUL MF
							5064	-1C2WLWW
							0430	-1 C2UL ww
8488	6433	4761	3620	3030	2872	2877	2820	-102wLFW
.5600	.5641	•5039	•4020	•2795	•1528	•C306	C574	-102ULFW
6131	4137	2679	1678	(993	0409	.0020	.0515	-1C2hDFw
.3806	•3654	.3157	.2571	.2060	.1709	.1545	<b>.</b> 1561	-102UDFw
5264	6925	8845	-1.0485	-1.0325	8689	8874	5808	-102 WL FT
2366	1826	0583	.1779	-468t	-4794	.2213	0276	-102111 FT
5531	5940	61.04	5705	- 4328	2498	- 1047	0504	-1020017
9040	- 8245	- 7202	- 5956	- 5015	- 5234	- 6963	- 6063	-1020001
.,,,,	.0215	••••	• > > > 0	• 201 2			-1 3102	-10200FT
							-1-2103	-ICZWLWI
							.0055	-ICZULWI
							6270	-102wLTT
							4956	-102LLTT
0362	0349	0344	0345	0350	0357	0356	C363	-1C2WLFP
0465	0443	0431	0426	0427	0425	0399	0093	-1C2ULFP
0973	0866	0777	0699	0629	0561	0495	C431	-1C2mDFP
. 9045	.9003	.8967	.8937	.8915	. 8903	-89u4	.8916	-102UDEP
							0733	-102wLwP
							6448	-102111 MP
							0524	-102604
							1 1750	-1021040
							1.1750	-IUZCDWP
			ALPHA =	-5				
20.0000	30.0000	40.0000	50.0000	60.0000	70.000	80.0000	90.000	Сні
-1.1147	- 91 71	- 7112	- 5450	- 4497	- 4394	- 4310	- 4141	- 5241 66
. 4333	5451	5520		2416	1043	0217	- 0075	
- 2222	- 50.22	• 3023	• 1 2 3	• 3 • 1 7	• 100 5	-0517	0075	-520675
0225		3926	2413	1352	0554	-0151	• 6875	-52NUFF
•1100	•1745	• 16 58	.1165	• ( 591	.0150	0041	<b>C.CCCO</b>	-52UDFF
							6775	-52WL NF
							C837	-526L NF
							5096	-52 WL WW
				-			0764	-52ULww
d346	6383	4768	3651	3062	2893	2883	2811	-52wLFw
.5332	•5399	•4842	•3863	.2665	.1407	.C185	0693	-52ULFW
5847	3926	2507	1522	(842	0319	.0169	. 6663	-52hDFw
.3585	. 3486	.3042	-2500	-2020	. 1693	1551	.1588	-52UDEw
5522	7488	9859	-1.1927	-1.1399	8856	8685	- \$314	-52 =1 FT
3140	- 2710	- 1/27	1263	4955	4965	2277	- 0150	-52111.51
- 5627	- 6540	1457	• 1 3 4 1		• 4 5 0 5	• 2 2 1 1	0150	- 520LFT
- 3551	- 0049	0718	0004	4900	2002	1209	.0203	- DZWUFI FOLDET
- • > > > > > > >	8/3/	1512	-+ 6053	4835	~.5085	5/38	5940	-5200FT
							-1.1739	-52WLWT
							.0097	-52UL HT
							5258	-52wLTT
							2741	-52UL TT
0358	0345	0339	0340	0346	0352	0351	C357	-52WLFP
0459	•0437	0425	0421	0421	0419	0392	085	-52ULFP
0965	0859	0771	0694	0625	(558	0493	0430	-52WDFP
.9041	.8999	.8961	.893.0	8908	8896	8896	8906	- 52UDEP
		-0701	-0,50	-0,00			0721	- 5291 - 0
							- ()477	-5211
							- 0522	_ 62WDWD
							1 1 7 2 2	
							1.1122	->2UUMP

#### ALPHA = 0

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0CCL	SC.C000	CHI
-1.1164	9185	7118	5450	4487	4200	4211	4162	02 WLFF
. 4325	.5424	• 5534	.4760	•3417	·1059	.0312	0880	C2ULFF
	5926	3925	2409	1347	0560	.0156	.0883.	C2wDFF
1 391	.1744	-1658	.1164	.0589	.6156	0041	C.0000	C2UDFF
•1071	••••		••••				6823	C2WLWF
							- 1016	C2III bE
								020640
							5157	OZNENN
						_	1094	UZULAW
0233	6354	4789	3692	3102	2922	- • 298	2811	UZWLEW
.5082	.5177	•4663	.3723	.2548	.1296	.0073	C8C6	CZULFN
- 5595	3738	2352	1379	0701	0178	.0312	.0806	02wDFW
3471	. 3321	-2928	.2426	.1976	.1670	.1547	.1603	CZUDEW
	- 9142	-1 1120	-1 3813	-1.2749	- 9136	8621	8952	02WLFT
	0172	2663	0003	5473	5216	2270	- 0024	02ULET
4021	3/48	2462	• 0092	• 5 7 5 7	• 121 0	1401	0024	C2=0FT
6396	1262	/906	1130	5121	3148	1491	-CU24	
-1.0178	9352	8052	6174	4596	4530	3642	5801	OZUDFI
							-1.1432	CZWLWI
							.C136	OZULNI
							4729	02 mL T T
							1045	C2ULTT
- 0353	- 0340	0335	0336	0341		- 0345	0351	C2WLFP
- 0453	- 0431	- 0619	- 0414	- 0414	- 0412	- 0385	0078	02ULEP
	~•0451	0419	0414	- 0414	0666	- 0461	- 6429	026057
0956	0852	0765	0690	0621	0556	0471		020010
• 5036	.8993	• 89 55	•8923	.8900	. 888 /	•888c	.8890	020069
							0709	UZWLWP
							0419	02UL MP
							0523	C2WDWP
							1.1714	C2UDwP
			ALPHA =	5				
20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
-1.1191	9204	7125	5449	448ć	4200	4211	4161	52WLFF
.4327	. 5438	.555C	.4771	. 3423	.1863	.0316	0875	52ULFF
8256	- 5943	- 3935	2414	1352	0564	.0151	.0875	52WDFF
1055	. 1754	-1668	.1170	- 0592	.0158	0041	C.C000	52UDFF
•1075	• • • • •	•1000	•••••				6860	52 hE hE
							- 1188	5211 44
							£107	5262 81
							2107	JZWL NN
							1420	52ULWW
ºl47	6343	4823	3742	3150	2959	2920	2820	SZWLFW
.4051	•4975	•4503	• 3599	.2444	.1196	0031	0913	52ULFW
5373	3573	2213	1248	0571	0046	.0446	•C942	52wDFw
.3161	.3160	.2814	.2349	.1925	.1638	.1533	.1608	52UDFW
6081	8899	-1-2700	-1.6312	-1.4376	9516	8677	8714	52WLFT
- 5047	5004	3743	.0426	. 6336	. 5545	-2495	. (105	52ULFT
- 6914	- 9064	- 91.20	- 9163	6622	- 3540	- 1747	0218	52WDET
0914	0030		(310)	- 4345		- 5541	- 5903	520061
-1.0936	-1.0129	-+0043	0310	4200	4103		1 1101	520011
							-1.1101	DZHLWI COMUNT
							.0176	52 UL NI
							4583	52WLTT
							. C439	52ULTT
48 د ل	0335	0330	0331	C336	0341	0340	0346	52WLFP
0440	0424	0412	0408	C408	0405	0378	0072	52UL FP
0946	0846	0760	0686	C61 8	- 0554	0490	0429	52WUFP
. 2021	HQA7	8948	.8916	8492	. 8878	_RA74	.8884	52LDFP
							0A9A	5211 HP
							0405	52111 40
							- 0534	E2MDHD
								524044
							1.1094	SZUDWP

#### ALPHA = 10

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
-1.1230	9229	7134	5449	4482	4196	4208	4160	102wlff
.4340	•5463	.5576	.4791	.3438	.1876	.0330	C862	1024L FF
8298	5974	3954	2429	1365	0577	.0138	.0862	102mDFF
.1111	.1775	.1686	.1182	• C6O C	.0162	0038	c.ccoo	1020066
							6884	1020011
							- 1351	10206-66
							- 5244	10206 #F
							- 1740	102016.000
8086	6351	4869	3801	- 3205	- 2002	- 2051	- 2020	1020688
- 4638	-4763	-4361	3491	- 2355		- 0125	- 1012	1028658
5182	- 3430	- 2092	- 1132	- (452	• 1100		1012	IUZULFW
2956	3002	2700	1155	1046	.0010	.0572	•1671	LUZNDEN
	- 0770	-1 4704	-1 0409	+1000	•1233	.1509	.1600	I CZUDEW
264	- 4544	-1.4700	-1+9090	-1+0234	9986		2001	LOZWEFT
0200		5412	0037	. 1010	. 5949	.2657	.0242	1C2ULFT
	9074	-1.0632	-1.1031	7653	3983	2633	0469	102WDFT
-1.1000	-1.1125	90:0	04/0	38UC	4625	5495	5766	1020011
							-1.0986	102wLWT
							• C221	1 C 2 U L N T
							4795	102WLTT
							.1941	1 G2 UL TT
0343	0330	0325	0326	C33C	0330	0334	C340	102wLFP
0439	0417	0405	0401	0401	0398	0371	0065	102ULFP
0941	0840	0755	0682	0615	0552	0490	C429	1C2WDFP
.9024	.8980	.8940	.8907	.8882	.8868	.8865	.8872	102UDFP
							0687	102mL bP
							(393	1 C2 UL 10 P
							0526	1026062
							1,1673	1 C2UDWP
				•.				
			ALPHA =	16				
20.0000	30.0000	40.0000	50.0000	60.000C	70.0000	80.0000	90.0000	СНІ
-1.1289	9266	- 7146	- 5447	4475	4189	- 4203	- 4156	1624166
-4368	.5508	- 5621	. 4828	- 3467	. 1904	-0365	- 0834	1620155
8372	6028	- 3992	- 2458	- 1392	- 6604	0111	6834	142000
1146	.1815	1719	1205	1372	0004	- 0(22	0000	16280FF
•••••	• • • • • • • • • • • • • • • • • • • •		•••	• (01)	• • • • • •	0052		10200FF
							0070	
							- 5220	
							520	LOZWLWW
- 9442	- 4303	1010	2 102	2270	70.44	2000	2114	I CZUL WW
	0382	4540	3082	3219	3064	2999	2873	162WLFW
- 4003	.4000	• 4215	• 3 3 8 4	• 2 20 1	.1020	0224	1122	LEZULFW
4992	3290	1971	1013	0328	.0208	.0710	.1215	162wDFw
.2/10	.2816	•2563	.2169	•1792	•1540	.1466	.1576	162UDFw
6671	-1.0976	-1.7918	-2.5625	-1.8559	-1.0666	9217	630	162wLFT
8076	9022	8229	0463	1.0117	•6520	.2905	.C420	162ULFT
4260	-1.0475	-1.3002	-1.4164	9035	4589	2426	0792	162WDFT
-1.3284	-1.2722	-1.0997	663C	2993	4432	5436	5751	162UDFT
							-1.0835	162WLWT
							.0284	162UL mT
							5578	162wLTT
							.4104	162LL TT
0336	0324	0319	0319	0324	0329	0328	- 0334	16261 50
0429	04C8	0397	- 0392	- 0342	- 0390	0363	( . 58	1621150
0432	0833	0750	- 0679	0613	- (55)	0490	6430	] 62 001 0
.9015	.8970	20100	- 100 - A	- 20013	. 9985	-0770 2061	- • 0730 6057	1621058
		.0730	.0070	•00TU	.0000	•0031	• C C J I - 0475	14200FP
							0013	ICZWEWP
							03/9	1 CZULWP
							0530	162WUWP
							1.1647	162UD wP

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#### 1.5-1

ALPHA = -10

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
9307	8193	6189	4954	4301	4205	4325	4308	-103wLFF
1435	.4276	.4341	.3721	.2661	.1442	.0248	0635	-103ULFF
- 7085	5263	3673	2445	1548	0826	0115	.C635	-103wDFF
	06.05	0724	. 0516	-0231	-0028	0045	0.0000	-103UDFF
0104	.0007	.0124					6806	-103 mi hF
							- 0507	-103ULWE
							0.001	103664
							5450	TUSHLWW
							0304	-1030LWW
7394	5771	4461	3594	3194	3161	3243	3208	-103WLFW
.4381	.4449	.3989	.3183	.2212	.1212	.0258	C4C3	-1C3ULFW
5004	3969	2744	1870	1233	0699	0161	.0401	-103wDFw
2260	.2429	.2244	.1956	.1695	.1533	.1465	.1483	-103UDFW
	- 6363	- 7894	- 9028	8638	7591	8089	8868	-103WLFT
4920	- 1294	- 0134	1906	. 396.8	. 36 4 1	.1518	0357	-103ULFT
1854	- 1204	0114	- 4956	- 3638	- 2197	0977	- 0438	-103WDFT
	5201	5295	4000	3030	- 5205	- 5747	- 5924	-103UDET
8742	7985	1024	5945	• 2231		3 14 1	1.1.251	-103-41 -
							-1.1251	-103WLWT
							0106	-1030LWI
							5329	-103NL11
							3457	-103ULTT
0 145	0384	0381	0386	0395	0406	0408	C410	-1C3WLFP
	0542	0531	0530	0536	0542	0524	0264	-103ULFP
- 0500	- 0858	- 0762		0600	0525	0450	0376	-103WDFP
	04.05	0102	9507	8618	8533	.8555	.8586	-1 03UDEP
.8465	.8482	.0471	.0501	.010		• • • • • • •	- 0767	-103-1-0
							- 0704	-10314 60
								-1030LWP
							0470	-ICSNDWP
							1.1224	-10300 MP
			ALPHA =	-5				
20.0000	30.0000	40.0000	50.0000	60.0000	76.000	80.0000	90.0000	CHI
					1.2.10		(300	53-41 55
9310	7771	6192	4956	4304	4209	4328	4309	TOWLFF
•3422	.4268	.4336	.3714	.2651	.1432	•C238	645	- 53ULFF
7077	- 5255	3664	2436	1538	0817	0105	.0645	-53WDFF
0119	.05 55	.0717	.0510	.0226	.0025	0047	C.CCCO	-53UUFF
							6844	-53 WL WF
							C638	-53UL NF
							5471	-53WLWW
							0562	- 53UL WW
	6701		2404	- 2210	- 2170	- 3242	3168	-53al Fb
1212	5/21	4455	3000			0154	- 0506	-53ULEW
.4182	.4265	.3830	. 3060	.2100	.1110	-0134	0505	-536064
5387	3809	2615	1755	1124	(592	0056	.0505	
•2140	•2306	.2155	.1902	.1668	.1520	.1470	.1505	
5203	- • 6888	8804	-1.0197	9418	7792	8688	8652	-53WEF1
2483	1992	0736	.1686	•4220	.3761	.1611	0204	-53ULFT
- 5236	5758	6008	5596	<b>414</b> C	2469	1175	•C243	-53w0FT
171	83 81	7295	- 5982	5090	5291	5662	- 5757	-53UDFT
• / • • •			•••••				-1.1175	-53 WL WT
							0017	-53ULWT
							- 4767	-53WL TT
							_ 1054	-5311L TT
			<u> </u>	<b>~~</b>		o / 00	1424	- 500011
91دن	0360	0377	0381	0390	0401	0403	0404	
0560	0536	0525	0524	0530	0535	0510	0256	-53ULFP
0463	0850	0756	0672	0596	0521	0447	0375	-53wDFP
.8467	.8485	.8496	.8504	.8514	.8529	.8545	.6579	-53LUFP
							0785	-53WLWP
							0688	-53ULwP
							0468	-53 NDWP
							1.1211	-53UDWP

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			ALPHA =	0				
20.0000	30.0000	40.00CC	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
9322	7781	6196	4957	4305	4210	4329	4310	<b>U3WLFF</b>
. 241.5	4209	4339	.3714	2649	-1428	.0234	(648	03LLEE
- 7082	- 5258	- 3663	2434	- 1536	0814	0102	.0648	C3WDEE
- 0127	0593	0716	- 0509	.0224	.0023	0047	6.000	031055
	• • • • • • • •	•0710	•••••	I ULL I	•0025		6871	0341 65
							- 0763	Calline
							- 6480	COULWE
								CONENN CONENN
	5		3/36	2221	21/1	3340	0815	USULWW
/1/3	5686	4460	3028	3251	3184	3248	3194	CONLEN
.3951	.4055	.3697	.2949	.2011	.1017	.0057	0604	CBULEW
5193	3666	2499	1650	1022	0491	.0044	.C6C4	C3hDFh
•1981	.21 82	•2070	•1843	·163C	.1499	•1464	.1517	C3UDFw
5507	7530	9938	-1.1693	-1.0344	8047	8129	8470	03wLFT
3228	2827	1472	.1507	•4643	• 393a	<b>.</b> 1723	0051	03ULFT
5674	6409	6871	6500	4711	2761	1377	.0051	03wDFT
9690	8873	7639	6013	4903	5182	5584	5699	C3UDFT
							-1.1049	03nLhT
							.0070	03LL HT
							4481	COWLTT
							0737	C3ULTT
- 0387	- 0376	(1373	0377	0386	- 11295	0397	(399	03wLEP
- 0553	- 0526	- 0519	- 0517	- 0523	- 0527	- 0509	- (248	A 3LL EP
- 0955	- 0844	- 0750	- 0658	- (592)	- (514	- 0446	- (374	C3WDEP
		0150	0000				8571	030060
.0400	• 0 4 0 4	• 0 4 7 4	•050L	•0)10	.0525	•0.742	- 0774	(3001 F
							0013	03ULWP
								USHUNP
							1.1198	U 3U D WP
			ALPHA =	5				
20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.000	CHI
9341	7743	6201	4957	4305	4200	4329	4309	53WLFF
. 3417	.4279	.4350	. 3721	.2653	.1431	.0237	(645	53ULFF
7100	- 5271	- 3670	2438	- 1539	0818	0105	.0645	53WDEE
0124	.06.00	.0722	-0512	.0226	.0025	0047	C.CCC0	53UDEE
							6886	53WL NE
							0881	5311 66
							5510	5341 88
							1064	5301.66
- 7093	- 5666	- 4474	- 3656	- 3258	- 3204	- 3259	3198	FBULFW
.3416	3940	3572	2850	.1926	. 0933	0042	(697	531.1 FW
- 5022	- 3540	- 2266	- 1555	- (920	- (395	0138	.0698	53-0254
-1022			-•1770	1504	1470	1450	1610	52UUEW
•1025	+2021	+1900	•1119	• 1000	- 4363	- 4201	• 1010	5300FW
5851	0209	-1+1505	-1.3023	-1+1392	0342	0201	0314	5 DUNE T
4100	3838	2313	.1400	.5291	• 4104	+1007	-0104	530051
6170	/1//	1929	1622	5351	3074	1588	0145	53 WUFT
-1.0324	9492	8081	֥6024		5066	5509	5649	5300FT
							-1.0867	53WL NT
							•0158	53ULWT
							4406	53WLTT
							.0360	53ULTT
0382	0371	0368	0372	C38C	0390	0392	0393	53WLFP
0546	0522	0511	0510	0515	0520	0501	0241	53ULFP
0947	0837	0745	0663	0589	0516	0444	0374	53wDFP
. 0467	•8482	.8491	.8497	.8505	.8517	.8534	.8562	53UDFP
							0763	53WLWP
							C658	53UL NP
							0468	53wDWP
							1.1182	53UD mP

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#### ALPHA = 10

20.0000	30.0000	40.000C	50.0000	60.0000	70.0000	0000.00	90.0000	СНІ
9367	7811	6707	4956	4302	4208	4328	4301	103wLFF
.3426	.4297	•4368	• 3735	•2664	.1441	•0247	0635	103ULFF
- 7130	- 5253	3685	2449	1549	C828	C115	.0635	1C3wDFF
0113	-0615	.0735	.0521	•C231	.0026	0045	0.000	103UDFF
.0115		••••					6889	1 C3WL WF
							- (992	10301 WE
							- 5525	1030281
							1207	10311 64
							1307	LOBULWW
7032	5658	4497	3690	3290	3229	3216	3208	LUSALEW
.3005	• 3800	• 3462	•2765	•1852	<u>.</u> C858	0114	0783	IC3ULFW
4074	3432	2306	1471	0845	0313	•C224	.0786	103WDFW
-1066	.1932	.1886	.1711	1534	.1432	•1425	.1508	103UDFw
6192	9189	-1.3197	-1.6148	-1.2510	8664	8298	8183	103wLFT
- 5143	5095	3516	.1441	.6224	. 4439	.2022	.C264	103ULFT
- 4721		- 4252	- 9037	6053	3404	1810	0345	103wDFT
	1 03 01	- 9455	- 5091	- 4310	- 4640	- 5438	5608	103LDET
-1.1105	-1.0201	- • • • • • • • •		4510		- • • • • • • •	-1 0622	103651
							-1.0035	1020167
							+ 6 2 4 7	
							4522	
							.1467	1030L11
0377	0366	0363	0367	0375	0385	0386	C387	1C3WLFP
0538	0514	0504	0502	C507	0512	0493	0233	1C3ULFP
0539	0831	0740	0660	0586	0515	0444	C374	103wDFP
- 4466	. 8480	.8487	.8492	.8499	. 8509	.8526	.8552	1C3UDFP
							0752	103WLWP
							0644	103UL WP
							0469	103wDwP
							1,1166	1 C 3 L D 5 P
			ALPHA =	16				
								C LL Z
20.0000	30.0000	40.0000	50.0000	60.000	70.0000	80.0000	50.000	CHI
9408	7836	6215	4954	4297	4202	4323	4306	163wLFF
. 3446	. 4329	-440C	. 3762	.2686	.1463	.0269	0614	163ULFF
7183	- 5332	3698	2471	1570	0848	0136	.0614	163wDFF
	. 0644	.0755	-0538	. 0242	.0035	0042	C.0000	1630DFF
		••••					6875	163mL mF
							1114	163UL WF
							5569	16301 66
							- 1588	16311 66
			2720	222/	33. 6	3.0.	- 2221	1420160
0481	50 00	4001			3205			14211164
.3478	• 3651	.3348	.2679	• 1779	.0781	0201	00/9	1COULEN
4727	3325	2216	1385	0757	0223	.0317	.0883	IC3WUFW
.1477	.1781	.1771	.1622	•1463	.1374	.1381	<b>.</b> 1480	16300FW
6040	-1.0508	-1.6188	-2.0318	-1.3796	9080	8448	8069	163wLFT
6702	7075	5375	.1924	.7780	.4835	.2260	•C467	1636LFT
7503	9426	-1.1345	-1.1289	6948	3846	2100	C6C4	163WDFT
-1-2297	-1-1544	9602	5760	3758	4774	5356	5560	163UDFT
					• • • • •		-1.0307	163 ML WT
							(367	16300 HT
							- 4954	16341 77
								16201 TT
4 A			<b>63</b> / 6	0.24 0	a <b>a a</b>	0.70	• 2777	1430111
0371	0360	0357	0360	0369	0378	0319	0381	103WLFP
0527	0505	0494	0493	C498	0502	0483	C225	163ULFP
0431	0824	0735	0656	0583	0513	C443	0375	163WDFP
.8464	• 84 76	.8481	.8485	•8491	.8500	.8515	.8540	163UDFP
							0740	163wLhP
							0628	163UL mP
							0472	163WDWP
							1.1145	163UDmP

#### 7X10

## ALPHA = -10

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
-1.1129	9093	7039	5440	4580	4396	4519	4403	-1C4wLFF
.4294	.5293	.5312	.4527	.3291	.1947	.0673	0298	-1 C4ULEE
4990	6676	- 4672	- 3150	- 2071	1252	- (49)	.0298	-104-026
- 086.0	1539	1489	.1058	. 056 7	0209	0030	C CCOO	-1041055
	• • • • • • • • • • • • • • • • • • • •	•1+07	•1050	• • • • • • •	.0203	.0030		10400FF
							0345	-104WLWF
							0208	-104ULWF
							5698	-1C4WLWW
							0069	-104ULww
9765	7615	5768	4484	3844	3738	3851	3871	-104wLFw
.5157	• 55 22	•5104	.4163	.2984	.1784	.0661	0173	-1 C4ULFW
7923	5698	3963	2718	1841	1158	- <u>.</u> 0508	. C165	-104WDFW
2454	-2647	.2335	.1836	.1372	.1050	-0882	- 0840	-104UDEM
9457	-1.1313	-1.1673	-1-0441	- 8348	7314	- 7523	7852	-10451 FT
- 1746	01.05	2513	4525	4697	2104	1111	- 0534	-1041167
- 0639	- 9390	- 7393	- 5544	- 2517	- 1074	- 0720		-10402FF
	0307	-+1363			19/0	0129	• 6969	-1048061
-•1428	5701	4082	3109	3118	3602	-•3211	1950	-10400FT
							-1.0301	-104WLWT
							0422	-104ULWT
							4987	-104mLTT
							1713	-1 C4UL TT
0688	0709	0761	0828	051 C	1001	1081	1143	-1C4mLFP
2041	2795	2808	2876	2991	3132	3243	3058	-1C4ULEP
2217	1966	1748	- 1551	1366	1182	( 992	0794	-104#0EP
.6975	. 74 54	. 7873	. 8261	8645	9050	0495	6601	
•0 71 3	• • • • • •	•1013	.0201	•0045	• 9000	.7772	1011	-10400FP
							1011	-IU4WLWP
							5131	-IC4UE MP
							1314	-1C4WDwP
							1.2980	-1C4UUwP
<b>2</b> p. douo	20.0000	60.0000	50.0000	<b>60</b> 0000	70 0000	SU 0000	<b>60.00</b> 0	
20.0000	30.0000	40.0000	50.0000	60.0000	10-0000	80.0000	90.0000	LFI
-1.1126	9317	7038	5445	4581	4397	4495	4552	-54 WLFF
.4286	• 52 87	.5307	.4521	.3284	.1941	.0566	0305	-54LLFF
8981	6668	4664	3143	2065	- 1240	- C484	.0305	-54WDFF
-0850	.1532	.1483	-1053	- 0563	-0207	0029	0.0000	- 54110FF
	••••	.1405	•1000	• • • • • • •	.0207	••••	- 6564	- 5461 65
							- 0204	- 54111 45
							0293	- 54UL NF
							5768	- 34 ML NH
							0228	-54ULWW
9263	7542	5742	4482	3849	3742	3851	3867	-54wLFw
.5008	•5380	•4987	.4073	•2913	.1720	·C597	(236	-54ULFW
7731	5564	3864	2636	1767	1087	0440	•C232	-54WDFW
.2333	.2548	•2265	.1793	.1347	.1039	<b>.</b> C881	.0851	-54LDFW
-1.0703	-1.2331	-1.2740	-1.1166	8616	7377	7503	7749	-54wLFT
- 2289	0202	-2610	. 4936	.5052	. 3308	.1296	( 329	-54ULET
9365	9331	8257	- 6163	3892	2244	0560	- 0354	-54wDFT
7780	- 5854	- 3004	- 2995	- 2056	- 1521	- 3868	- 3940	- 54UDET
			2099	27,77	-• 3321	5808	-1 (235	- 546011
							-1.0250	- 34WLN1
							0253	-54UL NI
							4769	-54wLTT
	_						0997	-540LTT
0668	0708	0760	0825	0906	0996	1075	1136	-54WLFP
2825	2778	2791	2835	2971	3110	3218	3031	-54ULFP
2202	1952	1735	1539	1355	1172	0583	0786	-54mDFP
.6996	.7471	.7887	.8272	•8653	.9054	.9496	.5987	- 54UUFP
	2						1799	-54mLWP
							- 5084	- 5411 MP
								<i></i>
							- 1299	-54-0-0
							1299	-54 NDNP

#### ALPHA = 0

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	90.0000	CHI
-1.1130	9045	7039	5445	4581	4396	4521	4552	04wLFF
-1-11-50 	-5287	.5307	.4520	.3282	.1939	.0664	0307	C4ULFF
- 4202 - 4981	6668	- 4663	3141	- 2062	1244	0483	.0307	04wDFF
	1531	1482	-1052	. (562	0206	-0028	c.c000	04UDFF
.041	• 1 7 71	• I +U C	•1072				6577	04WLWF
							0374	04LL hF
							5721	64 HI HH
							0383	CALL WW
	<b>-</b>	c 7 0 7	4407	2050	- 3750	- 3664	- 3968	0452 84
9501	1484	5/2/	4487		3750		- (296	
.4865	• 5246	.4880	• 399 3	• 204 9	- 1002	- 0374		040014
7559	5445	3/15	2562	1695	1022	0570	+LZ90	
.2211	• 2449	.2195	•1747	.1319	.1023	.0870	• 6 6 5 5	
-1.1010	-1.3568	-1.4012	-1.1980	8502	/461	/500	1000	
2906	0524	.2790	•5463	• 545 2	• 3522	.1484	0127	
-1.0327	-1.0427	9255	6835	4280	2516	1190	.0127	
8195	6011	3853	2621	2773	3427	3814	3919	0400FT
							-1.0158	C4WENT
							0087	04ULWI
							4648	04WLTT
							0330	C4ULTT
0667	0707	0757	0822	0902	0991	1069	1128	04WLFP
2807	2760	2772	2837	2949	3086	3191	3003	C4 UL FP
2187	1938	1722	1527	1344	1162	0474	0779	04WDFP
- 701.6	.7487	.7899	.8281	.8659	.9056	.9493	.9981	04UDFP
							1787	C4 WL NP
							5036	C4ULWP
							1285	C4wDWP
							1.2963	C4UDwP
			ALPHA ≠	5				
20.000	30.0000	40.0000	50.0000	60.0000	70.0000	8 <b>0.0</b> 000	96.0000	CHI
		30/0		(501	(34)	- 4521	- 4552	544166
-1.1141	9103	7042	5445	4581	4598	4 321	- (305	5410 55
.4204	• 52 92	•5313	• 4524	.3285	.1941	•Ucoc	(3(5	540677
	6675	4667	3144	2065	1247	0485		540066
.6449	<b>.</b> 1535	.1486	.1054	• C 5 6 3	.0207	.0029		5400FF
							6584	54 WL WF
							0449	54 UL MF
							5736	54WLWW
							0534	54ULWW
9397	7439	5720	4498	<b></b> 3875	3701	3664	3875	54WLFW
.4725	.5122	.4782	.3921	.2793	.1610	.0483	C352	54ULFW
7408	5341	3698	2497	1638	0962	0316	.C356	54kDFk
2090	.2351	.2125	.1700	.1288	.1002	.0364	.0854	54UDF#
			1 2 5 4 5	0105	36.3	7613	- 7594	5441 67

-.9195

.5900

-.4680

-.2556

-.7563

-.2794

-.3317

.3740

.9056

-.7513 .1678

-.1422

-.3749

-.1061 -.3164

-.0967

.9489

-.7584

.CC74

-.0099

-.3889

-1.0068

.078 -.4615

.0316

-.1120

-.2975

.9971

-.1773 -.4987

-.1272

1.2947

-.0773

54WLFT

54ULFT

54wDFT 54UDFT

54WLNT 54UL HT

54 mL TT

54UL TT

54WLFP

54ULFP

54wDFP

54UDFP

54WLWP 54UL NP

54wDwP

54LD%P

-.0819 -.2815 -.1516 -. 0005 -. 0705 -. 0754 -. 2786 -. 2739 -. 2750 -. 2172 -. 1925 -. 1710 . 7034 .7501 .7909 -.0897 -.0985 -.3060 -.1153 -.2925 -.1334 .8288 .8662

-1.2714 -1.5072 -1.5515 -1.2865 -.3011 -.0853 .3092 .6133 -1.1457 -1.1716 -1.039E -.7557 -.8664 -.6164 -.3639 -.2273

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#### ALPHA = 10

20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	80.0000	50.0000	СНІ
-1 1158	- 9114	7047	- 5446	458C	4397	4520	4630	104wLFF
-1.1150	= <b>5</b> 204	5275	4533	2242	. 1947	-0673	(298	1 G4ULFF
.4291	.5504	• 5525	- 2151	- 2072	- 1253	0501	.0298	1C4wDFF
9012	6690	4677	3191	- • 2012		0030	C C000	104UDEE
.0857	.1545	•1494	.1050	. 6567	. 0209	.0050	6.000	1040071
							0 204	I CAMENE
							0520	104ULWP
							5754	1C4wLww
							0679	104ULWW
. (31.2	- 74.09	- 5724	4515	3892	3777	2876	3885	104 NL FW
			2450	2746	1565	-0435	0404	104ULFW
.4600	.5008	.4092	• • • • • • • • • • • • • • • • • • • •	1 6 9 /	- 0909	- 0262	. (412	104h0En
7277	5252	3633	2441	1050	0303	0667	(845	10400Em
•1926	•2254	•2055	.1651	.1255	.0977	.0.47	2624	1040010
-1.4067	-1.6910	-1.7274	-1.3794	54/5	1680	1540		
4419	1170	.3572	.6971	.6392	.3989	.1880	.0277	104CLFT
-1-2797	-1.3244	-1.1704	8322	5089	3081	1659	0327	1 C4WDF (
- 9203	- 6300	3321	1839	2308	3192	3673	3849	104UDFT
9203	.0300						9967	1C4WLnT
							.0246	104UL HT
							- 4668	104mLTT
							-4000	10404 TT
				_				1040211
0664	0702	0751	0814	0892	0978	1054	1111	
2764	2716	2727	2791	2895	3033	3136	2940	1046LFP
- 2158	1912	1698	1505	1324	1145	0960	0767	104mDFP
7051	7514	.7918	.8292	.8603	.9053	.9481	. 5959	1 C4UDFP
.7051	•1314	•••					1759	1C4WLWP
							- 4939	10446
							1261	1 C4mDWP
							1 2027	104UDWP
							1+2721	TOTODAL
			ALPHA =	16				
20.0000	30.0000	40.0000	50.0000	60.0000	70.0000	90.000	90.0000	CHI
	0177	7055	- 5447	- 4578	4394	4518	- 4550	164wLFF
-1.1186	9132	7055			1661	- 6686	0285	164ULEE
<b>.</b> 4306	•5325	.5345	•4221	. 3507	• 1 901	0506	0205	1644085
9048	6717	4697	3167	2085	1200	0504	.0200	1440066
.0875	.1564	.1511	.1072	• 0575	.0214	.0032	0.0000	1C4UUFF
							65/4	104 WL WF
							0596	164UL WF
							5777	164 ML MM
							0845	164ULWW
	7700	5720	- 4543	- 3019	3801	3899	3905	164wLfm
9251		- • 51 3 5		2600	.1520	-0384	0460	164ULFw
•4455	•4884	.4604	.3191	• 2077	+1720	- 0204	0472	164mDEW
7145	5166	3569	2385	1529	0095		C977	1641056
.1826	.2139	.1971	.1591	-1208	.0941	.0819	. 621	144018
-1.6116	-1.9668	-1.9741	-1.4897	9786	7832	/590	1412	LC4WLF1
5551	1486	.4496	.8228	.7031	.4300	.2138	• C527	LC4ULF1
-1.4766	-1-5474	-1.3506	9274	5585	344C	1954	0609	164WDFT
9957	- 64.09	- 2730	+ 1187	1972	3021	3567	3789	164uDFT
	0408	- • 2 / 3 /		• • • • •			9835	164WLWT
							.0435	164UL hT
							- 4848	16461 TT
							1770	1641.1 77
							.1110	1440611
0661	0698	0746	0808	0885	0970	1044	1101	IC4WLFP
- 2734	2687	-,2697	2760	2867	3000	3102	2912	164ULFP
	- 1807	-,1685	- 1494	1314	1136	0952	0762	164wDFP
7040	76.24	7025	8795	.8661	.9046	.9469	.9942	164UDFP
. 1009	•1520	.1723					1742	164mLwP
							- 4881	164ULWP
							- 1251	1646060
							1 2000	1641:0-0
							1.2070	TOTODEL

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NACA 16-015

Modified NACA 16-017

Tip sta.: c = 45.72 cm (18.00 in.)			¢ Sta.:	c = 83.414	14 cm (32.84 in.)		
X		±Ζ		x		$\pm \mathbf{Z}$	
cm	in.	cm	in.	cm in.		cm	in.
-22.860	-9.000	0	0	-41.707	-16.42	0	0
-22.288	-8.775	.739	.291	-40.640	-16.00	1.521	.599
-21.717	-8.550	1.031	.406	-39.624	-15.60	2.126	.837
-20.574	-8.100	1.435	.565	-37.567	-14.79	2.957	1.164
-19.431	-7.650	1.732	.682	-35.458	-13.96	3.571	1.406
-18.288	-7.200	1,976	.778	-33.350	-13.13	4.074	1.604
-16.002	-6.300	2.362	.930	-29.210	-11.50	4.862	1.914
-13.716	-5.400	2.664	1.049	-25.019	-9.85	5.497	2.164
-9.144	-3.600	3.096	1.219	-16.688	-6.57	6 <b>.38</b> 6	2.514
-4.572	-1.800	3.345	1.317	-8.331	-3.28	6 <b>.89</b> 6	2.715
0	0	3.429	1.350	0	0	7.069	2.783
4.572	1.800	3.335	1.313	8.331	3.28	6 <b>.8</b> 76	2.707
9.144	3.600	3.012	1.186	16.688	6.57	6.208	2.444
13.716	5.400	2.400	.945	25.019	9.85	4.945	1.947
18.288	7.200	1.438	.566	33.350	13.13	2.967	1.168
20.574	8.100	.808	.318	37.567	14.79	1.651	.650
22.860	9.000	.069	.027	41.707	16.42	.142	.056
L.E. radius: 0.503 cm (0.198 in.)			L.E. radius: 0.917 cm (0.361 in.)				

## TABLE II.- FUSE LAGE ORDINATES



X		R			
cm	in.	cm	in.		
0	0	0	0		
1.70	.67	2.18	.86		
3.38	1.33	3.05	1.20		
6.78	2.67	4.24	1.67		
10.16	4.00	5.13	2.02		
13.54	5.33	5.87	2.31		
20.32	8.00	7.01	2.76		
27.10	10.67	7.90	3.11		
40.64	16.00	9.17	3.61		
54.18	21.33	9.40	3.70		
67.74	26.67	10.16	4.00		
145.62	57.33	10.16	4.00		
159.18	62.67	9.88	3.89		
172.72	68.00	8.92	3.51		
<b>186.2</b> 6	73.33	7.11	2.80		
199.82	78.67	4.27	1.68		
<b>206.58</b>	81.33	2.39	.94		
213.36	84.00	.20	.08		
Nose radius: 1	l.50 cm (0.59 in.)				



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Figure 3.- Fan-in-wing model mounted in inserts in Ames 2.13- by 3.05-m (7- by 10-ft) wind tunnel.



Figure 3.- Continued.



Figure 3.- Concluded.

(c) 1.22- by 1.83-m (4- by 6-ft) rectangular test section.



has been omitted because of its vastly greater size. (See fig. 5 for relative size of Ames full-scale tunnel and fig. 6 for relative size of Langley full-scale tunnel.)



A-72-1215

(a) Overall view.

Figure 5.- Fan-in-wing model mounted in Ames 12.2- by 24.4-m (40- by 80- ft) tunnel.



A-72-1216

(b) View from ''below'' model.Figure 5.- Continued.



(c) View from "above" model.

Figure 5.- Concluded.



L-72-2844

(a) Near view.

Figure 6.- Fan-in-wing model mounted in 9.14- by 18.3-m (30- by 60-ft) Langley full-scale tunnel.



## Figure 6.- Concluded.



Figure 7.- Fan-in-wing model mounted in 2.13- by 3.05-m (7- by 10-ft) insert in Langley full-scale tunnel.

## Figure 7.- Continued.

## (b) Side view.





Figure 7.- Concluded.

model at zero angle of attack has been removed from the data. The solid curve is a least-squares quartic faired Figure 8.- Uncorrected data for the model with the fans covered in several different test sections. The drag of the through the data from the 9.14- by 18.3-m (30- by 60-ft) tunnel. (a) C<sub>L</sub>. <u></u>С. ес



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C<sub>n,T</sub>







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Figure 9.- Continued.





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vertical component of the static infusitiation use the vine wing there is no interference between the wing and the fans.







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TUNNELS 1.12X2.24-M FLAT OVAL (444X88-IN) FLAT OVAL 1.12X2.24-M RECTANGLE (44X6-FT) RECTANGLE (4X6-FT) RECTANGLE 2.13X3.05-M RECTANGLE 2.13X3.05-M RECTANGLE (30X60-FT) FLAT OVAL 12.2X24.4-M FLAT OVAL





Figure 11.- Continued.





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Figure 11.- Continued.





Figure 11.- Continued.

	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL
TUNNELS	(NI-88Xft) (H4X88-24-N)	1.12X2.24-M (44X88-IN)	1.22X1.83-M (4X6-FT)	2.13X3.05-M (7X10-FT)	9.14X18.3-M (30X60-FT)
	0		$\diamond$	٥	⊿



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Figure 11.- Concluded.



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(a)  $\alpha = -5^{\circ}$ .

Figure 12.- Uncorrected values of drag coefficient as a function of  $V/V_j$ . The solid curve represents the sum of the momentum-theory value of fan drag and the drag of the wing expressed in coefficient form. It is assumed that there is no interference between the wing and the fans.



Figure 12.- Continued.

(b)  $\alpha = 0^{0}$ .

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TUNNELS	1.12X2.24-M FLAT OVAL (44X88-IN)	1.12X2.24-M RECTANGLE (44X88-IN)	1.22X1.83-M RECTANGLE (4X6-FT)	2.13X3.05-M RECTANGLE (7X10-FT)	9.14X18.3-M FLAT OVAL (30X60-FT)	12.2X24.4-M FLAT OVAL (40X80-FT)
	0		$\diamond$	٥	⊿	۵



Figure 12.- Continued.

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Figure 13.- Uncorrected values of tail normal-force coefficient as a function of  $V/V_j$ .

(a)  $\alpha = -5^{\circ}$ .









Figure 13.- Continued.









Figure 13.- Concluded.



	FLAT OVAL	RECTANGLE	RECTANGLE	RECTRNGLE	FLAT OVAL	FLAT OVRI.
TUNNELS	1.12X2.24-M	1.12X2.24-M (44X88-IN)	1.22X1.83-M (4X6-FT)	2.13X3.05-M (7X10-FT)	9.14X18.3-M (30X60-FT)	12.2X24.4-M (40X80-FT)

0



Figure 14.- Continued.





Figure 14.- Continued.

Figure 14.- Continued.

(d)  $\alpha = 10^{\circ}$ .





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It is assumed that there is Figure 15.- Uncorrected values of ratios of drag to static thrust as a function of  $V/V_j$ . The curve is calculated from the momentum-theory values of the fan forces and the lift and drag of the wing. no interference between the fans and the wing.







Figure 15.- Continued.









Figure 15.- Concluded.

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Figure 16.- Uncorrected values of external drag-lift ratios as a function of  $V/V_j$ . The curve is calculated from the momentum-theory values of the fan forces and the lift and drag of the wing. It is assumed that there is no interference between the fans and the wing.

	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL	FLAT OVAL
TUNNELS	1.12X2.24-M	1.12X2.24-M	1.22X1.83-N	2.13X3.05-M	9.14X18.3-M	12.2X24.4-N
	(44X88-IN)	[44X88-IN]	(4X6-FT)	(7X10-FT)	[30X60-FT]	(40X80-FT)

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Figure 16.- Continued.







Figure 16.- Concluded.

the momentum-theory values of the fan lift, drag, and shaft power together with the lift and drag of the wing. It is assumed that there is no interference between the fans and the wing.







TUNNELS	1.12X2.24-M FLAT OVAL (444X88-IN)	1.12X2.24-M RECTANGLE (44X88-IN)	1.22X1.83-N RECTANGLE (4X6-FT)	2.13X3.05-M RECTANGLE (7X10-FT)	9.14X18.3-M FLAT OVAL (30X60-FT)	12.2X24.4-M FLAT OVAL (40X80-FT)
	0		$\diamond$	⊲	⊿	۵







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Figure 17.- Continued.





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show the values measured during the present tests in the 1.12- by 2.24-m (44- by 88-in.) flat-oval test section. Figure 19.- Influence of the ratio of fan area to wing area on "fan-induced" lift (from ref. 19). The solid symbols  $\alpha = 0^{0}; \quad V/V_{j} = 0.4.$ 



Figure 20.- Comparison between flight test and uncorrected wind-tunnel test (Ames test 388) measurements of the elevator deflection required to trim the YOV-10 air-craft when fitted with a rotating-cylinder flap (ref. 20). Flap set at  $60^{\circ}/30^{\circ}$ ; c.g. at 0.219c.









function of  $V/V_j$  in several different wind tunnels.



1.12X2.24-M RECTANGLE

1.12X2.24-M FLAT OVAL (444X88-1N)

TUNNELS

1.22X1.83-M RECTANGLE

2.13X3.05-M RECTANGLE (7X10-FT)

Figure 23.- Continued.

TUNNELS 1.12X2-24-M FLAT OVAL (44X80-IN) FLAT OVAL 1.12X2-24-M FECTANGLE (44X80-IN) FECTANGLE 1.22X1-83-M FECTANGLE (4X6-FT) FECTANGLE 2.13X3-05-M FECTANGLE

$$\circ \Box \diamond \triangleleft$$











0 0 0 0



Figure 23.- Concluded.



Figure 24.- Wall-induced values of  $q_{\rm c}/q$  in the uncorrected data from the fan-in-wing model as a function of  $V/\dot{V}_j$  in several different wind tunnels.
















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Figure 24.- Continued.

FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE
( NI -88Xff )	( N] -88244 )	1.22X1.83-N	2.13X3.05-M
N-fi2, 5X5 [ , 1	M-45. 5251 . 1	(4X6-FT)	(7X10-FT)







.85

.80

.75

Figure 24.- Continued.

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1.00



0 0 0 0



Figure 24.- Concluded.

(44X88-IN) RECTANGLE 2.13X3.05-M RECTANGLE (7X10-FT) I.22X1.83-M RECTANGLE 1.12X2.24-M FLAT OVAL TUNNELS 0  $\Diamond$ 4 קי ġ . ₽0∆ Φ 5 4 (a)  $\alpha = -5^{\circ}$ . ڢ ഗ് **H** 4 ņ **600** 0  $\Box$ 'n <del>0</del>0  $\mathbf{A}$  $\triangleleft^{\triangleleft}$ , Ο <u>1</u>7 10 θ g Ţ Ś 0 ∆ ï <sub>⊺</sub>

Figure 25.- Wall-induced tail incidence  $\Delta i_T$  in the uncorrected data from the fan-in-wing model as a function of  $V/V_j$  in several different wind tunnels.



 $\Delta \ddot{l}_{T}$ 













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TUNNELS





















 $Q_T/Q_c$ 

Figure 26.- Continued.



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 $0_7/0_c$ 

Figure 26.- Continued.





 $Q_T/Q_C$ 





 $Q_T/Q_C$ 









Figure 27.- Continued.







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1.12X2.24-M RECTANGLE (444X88-IN) 1.22X1.83-M RECTANGLE (4X6-FT) 2.13X3.05-M RECTANGLE (7X10-FT) 1.12X2.24-M FLAT OVAL [444X88-IN] TUNNELS



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Figure 27.- Concluded.

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Figure 28.- Continued.







Figure 29.- Corrected values of lift coefficient as a function of  $V/V_j$ . The solid curve represents the sum of the vertical component of the static thrust and the lift of the wing expressed in coefficient form. It is assumed that there is no interference between the wing and the fans.





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Figure 29.- Continued.

TUNNELS	1.12X2.24-M FLAT OVAL (444X88-IN)	I.12X2.24-M RECTANGLE	1.22X1.83-M RECTANGLE (4X6-FT)	2.13X3.05-M RECTANGLE (7X10-FT)	9.14X18.3-M FLAT OVAL (30X60-FT)	12.2X24.4-M FLAT OVAL (40X80-FT)
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Figure 29.- Continued.

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Figure 29.- Continued.





Figure 29.- Concluded.





	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL	FLAT OVAL
TUNNELS	1.12X2.24-M	1.12X2.24-M	1.22X1.83-M	2.13X3.05-M	9.14X18.3-M	12.2X24.4-M
	(44X88-IN)	(44X88-IN)	(4X6-FT)	(7X10-FT)	[30X60-FT]	[40X80-FT]

0 · 0 ◊ ◊ △ △



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	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL	FLAT OVAL
TUNNELS	(N]-88X44) (N]-88X44)	(N]-88X44) (N]-88X44)	1 .22X1 .83-M (4X6-FT)	2.13X3.05-M (7X10-FT)	9.14X18.3-M (30X60-FT)	12.2X24.4-M (40X80-FT)
	0		$\diamond$	4	⊿	۵



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Figure 30.- Continued.

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	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL
TUNNELS	1 , 12X2 , 24-M M-H2, 2X2 , 1	1 . 12X2 . 24-M [44X88- [N]	1 .22X1 .83-M (4X6-FT)	2.13X3.05-M (7X10-FT)	9.14X18.3-M (30X60-FT)
	0		$\diamond$	4	⊿



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Figure 31.- Corrected values of tail normal-force coefficient as a function of  $V/V_j$ .



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Figure 31.- Concluded.

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vertical component of static thrust and the lift of the wing. It is assumed that there is no interference between Figure 32.- Corrected values of ratios of lift to static thrust as a function of  $\sqrt{V_{i}}$ . The curve is the sum of the the fans and the wing.

	FLAT OVAL	RECTANGLE	RECTANGLE	RECTANGLE	FLAT OVAL	FLAT OVAL
TUNNELS	1.12X2.24-M	1.12X2.24-M	1.22X1.83-M	2.13X3.05-M	9.14X18.3-M	12.2X24.4-M
	[44X88-[N]	(44X88-[N)	(4X6-FT)	(7X10-FT)	(30X60-FT)	(40X80-FT)

0



Figure 32.- Continued.













Figure 33.- Corrected values of ratios of drag to static thrust as a function of  $\sqrt{V_j}$ . The curve is calculated from the momentum-theory values of the fan forces and the lift and drag of the wing. It is assumed that there is no interference between the fans and the wing.



Figure 33.- Continued.



Figure 33.- Continued.



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Figure 33.- Continued.





Figure 33.- Concluded.



interference between the fans and the wing.

	FLAT OVAL	RECTANGLE	RECTRNGLE	RECTANGLE	FLAT OVAL	FLAT OVAL
TUNNELS	1.12X2.24-M	1.12X2.24-M	1.22X1.83-M	2.13X3.05-M	9.14X18.3-M	12.2X24.4-M
	(44X88-IN)	(44X88-IN)	(4X6-FT)	(7X10-FT)	(30X60-FT)	(40X80-FT)



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Figure 34.- Continued.





Figure 34.- Continued.





Figure 35.- Corrected values of equivalent lift-drag ratios as a function of  $V/V_i$ . The curve is calculated from the momentum-theory values of the fan lift, drag, and shaft power together with the lift and drag of the wing. It is assumed that there is no interference between the fans and the wing.

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Figure 35.- Continued.



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Figure 35.- Continued.

(c)  $\alpha = 5^{\circ}$ .

9.14X18.3-M FLAT OVAL (30X60-FT) FLAT OVAL 12.2X24.4-M FLAT OVAL (40X80-FT) FLAT OVAL 1.12X2.24-M FLAT DVAL (44X88-IN) 1.12X2.24-M RECTANGLE (444X88-IN) 1.22X1.83-M RECTANGLE (4X6-FT) 2.13X3.05-M RECTANGLE (7X10-FT) TUNNELS

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L/DE





L/D<sub>E</sub>

Figure 35.- Concluded.



Figure 36.- Wall-induced angle of attack  $\Delta \alpha$  in the corrected data from the fan-in-wing model as a function of  $\sqrt{V_j}$  in several different wind tunnels.



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Z 2.13X3.05-M RECTANGLE

1.22X1.83-M RECTANGLE (4X6-FT)

(HHX88-IN) RECTANGLE

1.12X2.24-M FLAT OVAL

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Figure 36.- Continued.

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a function of  $\sqrt{V_j}$  in several different wind tunnels.







TUNNELS 1.12X2.24-M (441X88-IN) 1.12X2.24-M 1.12X2.24-M rectanGLE 1.22X1.83-M rectanGLE 1.22X1.83-M rectanGLE 2.13X3.05-M rectanGLE

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Figure 37.- Continued.





	FLAT OVAL	RECTRNGLE	RECTANGLE	RECTANGLE
IUNNELS	1.12X2.24-M	1.12X2.24-M (44X88-IN)	1.22X1.83-M (4X6-FT)	2.13X3.05-M (7X10-FT)







Figure 37.- Concluded.





a function of  $V/V_j$  in several different wind tunnels.



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Figure 38.- Continued.





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Figure 39.- Wall-induced values of  $q_T/q_c$  in the corrected data from the fan-in-wing model as  $\sqrt{V_j}$  in several different wind tunnels. a function of





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Figure 39.- Continued.

 $Q_{T}/Q_{C}$ 

(44X88-IN) RECTANGLE 1.22X1.83-M RECTANGLE (4X6-FT) 2.13X3.05-M RECTANGLE 1.12X2.24-M FLAT OVAL [444X88-1N] TUNNELS





Figure 39.- Concluded.



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Figure 40.- Wall-induced fan incidence  $\Delta i_F$  in the corrected data from the fan-in-wing model as a function of  $V/V_j$  in several different wind tunnels.



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Figure 40.- Continued.

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2.13X3.05-M RECTANGLE (7X10-FT)

1.22X1.83-M RECTANGLE (4X6-FT)

1.12X2.24-M RECTANGLE (44X88-IN)

1.12X2.24-M FLAT DVAL (444X88-1N)

TUNNELS

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Figure 40.- Continued.



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Figure 40.- Concluded.



a function of  $V/V_j$  in several different wind tunnels.









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Figure 41.- Continued.



