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# ANALYSIS OF THE VERTICAL TEMPERATURE PROFILE RADIOMETER (VTPR) RADIOMETRIC PROBLEM 

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\section*{ANALYSIS OF THE VERTICAL}

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\begin{abstract}
The VTPR instruments that are presently in orbit exhibit a significant variation in bias radiation on the detector as a function of scan mirror position. Tests in the laboratory show that this bias radiation disappears when optical baffles are added. A detailed analysis explains quantitatively the observed bias variation with is due to (1) the extraneous field-of-view of the detector and (2) the variation in magnitude of the far field-of-view solid angle as a function of mirror position.
\end{abstract}

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ANALYSIS OF THE VERTICAL TEMPERATURE PROFILE RADIOMETER (VTPR) RADIOMETRIC PROBLEM
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\section*{INTRODUCTION AND SUMMARY}

The VTPR telescopes that are presently in orbit (as of July 1974) do not contain optical baffles. Without baffles there is a significant variation in bias radiation on the detector as a function of scan mirror position. This is true even when the entire external scene is cold space. When scanning space the largest signal output occurs at or near scan step 10 and the smallest signal occurs near step 23. At first, it was believed that the varying bias was due to temperature gradients within the VTPR and parts of the spacecraft toward the front of the VTPR. The explanation of the varying bias is as follows:
a. The detector has an extraneous solid angle field-of-view (FOV) much larger than the nominal \(2.1^{\circ} \times 2.1^{\circ} \mathrm{FOV}\) ( 0.00134 steradians). Recent tests indicate that the total extraneous FOV is about three orders of magnitude higher than the nominal solid angle FOV. Not all of the total extraneous FOV plays a significant role in the bias variation. Only that portion that represents the scene beyond the spacecraft causes the error. Calculations show that this portion is larger than the nominal solid angle FOV by a factor of 30 to 82 , depending on scan mirror position. Although the signals from points within the extraneous field is greatly attenuated the large field size produces an appreciable radiant input.
b. When no baffles are employed, the rotation and displacement of the scan mirror changes the detector's extraneous view of the outside world. This happens because the detector can view parts of the scan mirror directly, even though a large portion is obscured by the secondary mirror. In addition, a portion of the outside scene, as viewed by direct reflection from the scan mirror, is obstructed by the VTPR housing and "race track" shaped port in front of the instrument. We shall refer to this unobstructed outside scene as the "far field" scene. As the scan mirror steps from position 1 to 23 the far field solid angle varies considerably. It is primarily this variation that causes the bias to change as a function of scan mirror position! Quantitatively, this is confirmed by the fact that the far field solid angle is greatest at step 10 and least at step 23 (see Figures 2 through 7). Thus when scanning space the largest far field solid angle will produce the least radiant input to the detector (the extraneous radiant input comes solely from the near field when the instrument scans space). Plots made by RCA AED of the far and near fields of the VTPR detector for 6 scan step positions show that there is excellent qualitative correlation between the far field solid angle and observed signal amplitude
when the radiometer scans space. This is true for the instruments in orbit and those tested in the laboratory.

A quantitative analysis has been made using test data taken at Barnes Engineering with VTPR S/N 007. When baffles are installed the bias variation essentially disappears. Without baffles the bias variation was calculated using:
1. FOV measurements over a \(60^{\circ} \times 42^{\circ}\) field.
2. Calculated geometric FOV taken from RCA plots for 6 scan mirror positions of the VTPR.
3. Correction factors due to the out-of-focus condition in the test configuration.

Calculations were made for 6 scan mirror positions. Scan step 23 was used as a reference to obtain differences in signals. The calculated values were compared with the \(S / N 007\) measurements for the \(833 \mathrm{~cm}^{-1}\) channel, without baffle, and the results are shown in Figure 1. The bias radiation tests were made at Barnes Engineering with a "cold space" target at a temperature of about 100 K . The actual value of the maximum signal for the two tests shown was 837 digital counts (average). Note that only differences ( \(\triangle \mathbf{C}\) ) are plotted, scan step 23 being used as the reference. Considering the sensitivity of the calculations to small errors and the fact that the \(1000^{\circ} \mathrm{C}\) source used for the FOV tests has a temperature drift of \(\pm 30^{\circ} \mathrm{C}\) the correlation between measured and calculated values is excellent.

\section*{DETAILED ANALYSIS}

\section*{Calculations Summary}

The calculated bias voltage variations can be converted to equivalent digital count variations \(\Delta \mathrm{C}\) by
\[
\frac{\Delta C}{C_{p}}=\frac{L}{V_{p}}\left(A_{n} V_{n}-A_{23} V_{23}\right)
\]
where
\(C_{p}=\) peak value in digital counts, measured with VTPR S/N 007 in the \(833 \mathrm{~cm}^{-1}\) channel in the bias radiation tests.


Figure 1. VTPR S/N 007-Bias Variation
\(\mathrm{L}=\) correction factor due to out-of-focus condition during FOV tests
\(\mathrm{V}_{\mathrm{p}}=\) on-axis peak signal voltage
\(A_{n}=\) number of \(3^{\circ} \times 3^{\circ}\) elements in the far field scene for scan position \(n\)
\(A_{23}=\) number of \(3^{\circ} \times 3^{\circ}\) elements in the far field scene for scan position 23
\(\mathrm{V}_{\mathrm{n}}=\) average voltage within \(\mathrm{A}_{\mathrm{n}}\)
\(\mathrm{V}_{23}=\) average voltage within \(\mathrm{A}_{23}\)
As explained below, \(\mathrm{C}_{\mathrm{p}}=837\) counts, \(\mathrm{L}=0.603\) and \(\mathrm{V}_{\mathrm{p}}=26.0\) volts. Therefore,
\[
\Delta C=0.0194\left(A_{n} V_{n}-A_{23} V_{23}\right)
\]
where \(V_{n}\) and \(V_{23}\) are in millivolts. The values of \(A_{n}, V_{n}\) and \(\Delta C\) are shown in Table 1. The measured values of \(\Delta C\) shown in the table are the average readings from the two FOV tests with VTPR S/N 007 for the \(833 \mathrm{~cm}^{-1}\) channel.

The field-of-view tests and the method by which \(V_{n}\) was calculated is described in the next section.

\section*{FOV Tests}

The FOV tests were made using a 2-inch diameter, \(1000^{\circ} \mathrm{C}\) blackbody source with a radiation chopper in front of the source. The internal chopper of the VTPR was not used. The signals were taken from a test point before the A/D

Table 1
Calculated Bias Variations ( \(833 \mathrm{~cm}^{-1}\) Channel)
\begin{tabular}{|c|c|c|c|c|}
\hline \multirow[b]{2}{*}{Scan Step} & \multirow{2}{*}{\(\mathrm{A}_{\mathrm{n}}\)} & \multirow[b]{2}{*}{\[
\begin{gathered}
V_{n} \\
(\mathrm{mV})
\end{gathered}
\]} & \multicolumn{2}{|c|}{\(\Delta \mathrm{C}\) (Counts)} \\
\hline & & & Calc & Meas \\
\hline 1 & 22.2 & 26.5 & 3.0 & 3.6 \\
\hline 6 & 33.5 & 24.4 & 7.4 & 7.9 \\
\hline 10 & 39.6 & 24.3 & 10.2 & 9.3 \\
\hline 12 & 35.5 & 25.4 & 9.1 & 8.1 \\
\hline 18 & 20.0 & 28.0 & 2.4 & 3.9 \\
\hline 23 & 14.4 & 30.2 & 0 & 0 \\
\hline
\end{tabular}
converter and therefore the measured values are in volts rather than digital counts. For the extraneous FOV measurements an amplifier with a gain of 10 was employed. This external amplifier was not used for the on-axis peak signal measurement. All FOV data described in this report take amplifier gain into account.

The FOV was measured over a \(60^{\circ} \times 42^{\circ}\) field at \(3^{\circ}\) intervals. The peak radiation signal, collected by the primary mirror at the center of the FOV was 26 volts. The average value of the extraneous signals within the far field scene varied from 24.3 mV for scan step 10 to 30.2 mV for scan step 23 . These values are "above noise" voltages, i. e., the noise voltage was subtracted from the actual voltage reading for each point measured. This is important to do since we are interested in differences only. For example, if this were not done and if there were no extraneous signals, i.e., only noise, the calculations would indicate an erroneous bias variation because the far field areas are different for different scan mirror positions. This is because the calculated bias variation is proportional to
\[
A_{n} V_{n}-A_{23} V_{23}
\]

So that if
\[
\mathrm{V}_{\mathrm{n}}=\mathrm{V}_{23}=\text { noise voltage }=\mathrm{V}_{\mathrm{N}}
\]
the calculated variation would be proportional to
\[
V_{N}\left(A_{n}-A_{23}\right)
\]

This is obviously wrong. The explanation is that no account has been taken of the total field, in which all points have a value \(V_{N}\). Since the total field (near and far fields) does not vary (it is always \(2 \pi\) steradians) the bias would not change with mirror position. Therefore, in order to eliminate this source of error in the calculations, the noise voltage has been subtracted out.

Earlier FOV tests contained erroneous measurements of the peak voltage because the amplifier used had saturated. The proper value of \(V_{p}\) was measured during tests run on June 12, 1974 and is 26.0 volts.

The far field views for scan steps 1, 6, 10, 12, 18 and 23 are shown in Figures 2 through 7. They are simplified versions of RCA plots of the calculated geometric FOV of the VTPR. The signal voltage readings within the usable portions of the far field views for each of the above scan mirror positions are shown in Tables 2 through 7. The first entry in each \(3^{\circ} \times 3^{\circ}\) box element is the relative


Figure 2. Scan Step 1


Figure 3. Scan Step 6


Figure 4. Scan Step 10


Figure 5. Scạn Step 12


Figure 6. Scan Step 18


Figure 7. Scan Step 23

Table 2
FOV VTPR S/N \(007833 \mathrm{~cm}^{-1}\) Channel Scan Step 1
\(A_{1}=22.15 \quad A_{1 A}=16.71=\) NO. OF ELEMENTS USED FOR OBTAINING \(V_{1}\)

\(V_{1}=\frac{442.37}{16.71}=26.47 \mathrm{MV}=\) AVE. VOLT. FOR STEP 1

NOTE: ELEMENTS MARKED WITH AN \(X\) NOT USED FOR OBTAINING \(V_{1}\)

Table 3
FOV VTPR S/N 007 Scan Step 6


Table 4
FOV VTPR S/N 007 Scan Step 10
\(\mathrm{A}_{10}=39.59\)
\[
\mathrm{A}_{10 \mathrm{~A}}=31.53
\]
\(V_{10}=\frac{764.86}{31.53}=\mathbf{2 4 . 2 6} \mathrm{MV}\)

Table 5
FOV VTPR S/N 007 Scan Step 12
\(\mathrm{A}_{12}=35.54\)
\(A_{12 A}=29.26\)

SUBTOTAL
\begin{tabular}{|c|c|c|c|c|}
\hline \(-18^{\circ}\) & \(-15^{\circ}\) & \(-12^{\circ}\) & -9 \({ }^{\circ}\) & \(-6^{\circ}\) \\
\hline \multicolumn{2}{|l|}{\multirow[t]{3}{*}{}} & \(0.32 \times\) & \(0.47 \times\) & 0.03 \\
\hline & & 23= & 20.5= & \\
\hline & & 7.36 & 9.64 & \\
\hline & \(0.15 \times\) & \(0.99 \times\) & 0.36 & \\
\hline & 41- & 25= & & \\
\hline & 6.15 & 24.75 & & \\
\hline & 0.47 X & \(0.9 \times\) & \multicolumn{2}{|l|}{\(\rightarrow\)} \\
\hline & 44= & 22= & & \\
\hline , & 20.68 & 19.8 & & \\
\hline - & \(0.6 \times\) & 0.75 X & \multicolumn{2}{|l|}{\(\square\)} \\
\hline & 44= & 20.5= & & \\
\hline & 26.4 & 15.38 & \multicolumn{2}{|l|}{} \\
\hline - & 0.47 X & \(0.9 \times\) & \multicolumn{2}{|l|}{\(\rightarrow>\)} \\
\hline & 39= & 19.5= & , & \\
\hline - & 18.33 & 17.55 & \multicolumn{2}{|l|}{} \\
\hline & 0.15 X & \(0.99 \times\) & 0.36 & \multirow[t]{2}{*}{} \\
\hline & 28= & 25= & & \\
\hline \(\square\) & 4.2 & 24.75 & \multicolumn{2}{|l|}{\(\rightarrow\)} \\
\hline & - & \(0.32 \times\) & 0.47 X & 0.03 \\
\hline & & 36= & 19.5= & \\
\hline \[
L
\] &  & 11.52 & 9.17 & \\
\hline & 75.76 & 121.11 & 18.81 & \\
\hline
\end{tabular}

TOTAL: \(\mathbf{7 4 3 . 4 3}\)

\(V_{12}=\frac{743.43}{29.26}=25.41 \mathrm{MV}\)

Table 6
FOV VTPR S/N 007 Scan Step 18
\begin{tabular}{|c|c|c|c|c|}
\hline \(-18^{\circ}\) & \(-15^{\circ}\) & \(-12^{\circ}\) & \(-9^{\circ}\) & \(-6^{\circ}\) \\
\hline , & \(0.12 \times\) & \(0.52 \times\) & 0.47 X & 0.02 \\
\hline & 25.5= & 23= & 20.5= & \\
\hline & 3.06 & 11.96 & 9.64 & \\
\hline , & \(0.7 \times\) & \(1 \times\) & 0.36 & \\
\hline & 41= & 25= & & \\
\hline , & 28.7 & 25 & & \\
\hline  & \(0.98 \times\) & \(0.9 \times\) & & \\
\hline & 44= & 22= & & \\
\hline  & 43.12 & 19.8 & & \\
\hline 0.08 X & 1 X & \(0.75 \times\) & & \\
\hline 33= & 44= & 20.5= & & \\
\hline 2.64 & 44 & 15.38 & & \\
\hline  & 0.98 X & \(0.9 \times\) & & \\
\hline & 39= & 19.5= & , & \\
\hline , & 38.22 & 17.55 & & \\
\hline - & \(0.7 \times\) & 1 X & 0.36 & \\
\hline X & 28= & 25= & & \\
\hline \[
75
\] & 19.6 & 25 & , & \\
\hline & \(0.12 \times\) & 0.52 X & \(0.47 \times\) & 0.02 \\
\hline  & 32= & 36= & 19.5 \(=\) & \\
\hline  & 3.84 & 18.72 & 9.17 & \\
\hline 2.64 & 180.54 & 133.41 & 18.81 & \\
\hline
\end{tabular}

TOTAL: 507.0
\(A_{18 \mathrm{~A}}=18.10\)
\(+9^{\circ}\)
\(+6^{\circ}\)
\(+3^{\circ}\)
\(0^{\circ}\)
\(-3^{\circ}\)
\(-6^{\circ}\)
\(-9^{\circ}\)

\(V_{18}=\frac{507.0}{18.1}=28.01 \mathrm{MV}\)

Table 7
FOV VTPR S/N 007 Scan Step 23

\(V_{23}=\frac{411.73}{13.64}=30.19 \mathrm{MV}\)
amount of the elemental area that lies within the far field of view. The second entry is the measured signal in millivolts (less noise voltage) and the third entry is the product of the first two entries. Elements that lie within the far field•view but were not used for obtaining the average signal voltage contain one entry only (the portion of the elemental area within the far field view). Only those measurements that were made with a completely unobscured source were used. Two types of FOV measurements were made. The left side measurements were made using a large auxiliary flat mirror inserted between the scan and secondary mirrors. This gave good readings for the left side, but the secondary mirror obstructed most of the far field view on the right side. The right side readings that were used in the calculations were made with the scan mirror in the scan step 10 position (without the auxiliary mirror). The sum of the first entries is the value of \(A_{n}\).

\section*{Bias Variation Tests}

Two independent bias variation tests were made with \(\mathrm{S} / \mathrm{N} 007\) without baffles. In the first test the average peak signal in the \(833 \mathrm{~cm}^{-1}\) channel was 833.35 counts at scan step 9 (for scan step 10 the average signal was 833.1 counts). In the second test the average peak signal was 840.55 counts at scan step 10 (at step 9 the average signal was 840.45 counts). The average peak signal for both tests is 837 counts. The average readings for both tests are shown in Table 8.

Table 8
Bias Variation Tests Without Baffle ( \(833 \mathrm{~cm}^{-1}\) Channel). Space
Target \(\approx 100 \mathrm{~K}\).
\begin{tabular}{|c|c|c|}
\hline \multirow{2}{*}{\begin{tabular}{c} 
Scan \\
Step
\end{tabular}} & Test 1 & Counts \\
\cline { 2 - 3 } & & 827.3 \\
Test 2 \\
\hline 1 & 828.5 & 835.05 \\
2 & 831.0 & 836.05 \\
\hline \(\mathbf{5}\) & 833.35 & 838.95 \\
10 & 833.1 & 840.45 \\
12 & 831.5 & 840.55 \\
15 & 830.75 & 839.83 \\
19 & 826.8 & 837.8 \\
22 & 823.2 & 834.2 \\
23 & 823.4 & 831.45 \\
\hline
\end{tabular}

\section*{Out-of-Focus Corrections}

The need for extremely high signal-to-noise ratios necessitated the use of an uncollimated 2 -inch diameter source at a distance of only 76 inches. Because of this, signal energy is lost due to (1) increase in obscuration of the secondary mirror, (2) decrease in the angular size of the telescope entrance pupil, and (3) out-of-focus blurring of the image.

\section*{1. Secondary Mirror Obscuration}

The shadow of the secondary mirror on the primary is increased when the source is moved from infinity to 76 inches. For a source at infinity the obscuration diameter is 1.125 inches. Since the back rim of the secondary is located 2.377 inches from the vertex of the primary, which has a 7.268 -inch radius of curvature, the obscuration diameter is 1.161 inches (see Appendix 1).
2. Entrance Pupil Size

For an on-axis point source at infinity the entrance pupil diameter is 2.970 inches. For a source at 76 inches the pupil diameter is 2.720 inches (see Appendix 2).

The correction factor due to the change in angular size of both the primary and secondary mirrors is
\[
\frac{(2.720)^{2}-(1.161)^{2}}{(2.970)^{2}-(1.125)^{2}}=0.801
\]
3. Out-of-Focus Blur

The image of every source point is blurred by an angular diameter \(\theta_{B}\) given by (see Appondix 2 for more precise calculation)
\[
\theta_{\mathrm{B}} \approx \frac{\mathrm{D}}{\mathrm{~S}}
\]
where \(\quad D=\) effective entrance pupil diameter \(S=\) distance of source from primary mirror
\[
\begin{aligned}
\therefore \theta_{\mathrm{B}} & \approx \frac{2.720}{76}=0.0358 \mathrm{radian} \\
& \approx 2.05^{\circ}
\end{aligned}
\]

This blur size should be compared to the in-focus image size of \(1.5^{\circ}\) and the nominal \(2.1^{\circ} \mathrm{FOV}\) of the VTPR. Because of this large blur the detector collects only \(75.3 \%\) of the energy from the source, compared to what it would receive from an in-focus source. See Appendix 3 for detailed analysis of this energy loss.

The total correction factor is therefore
\[
L=(0.801)(0.753)=0.603
\]

The optical system of the VTPR is summarized in Appendix 4.

\section*{APPENDIX 1}

SECONDARY MIRROR OBSCURATION

APPENDIX 1
SECONDARY MIRROR OBSCURATION


NOTE: NOT DRAWN TO SCALE
\(\mathrm{y}=\) radius of circular shadow on primary
Primary radius of curvature \(=7.268\) inches
Secondary obscuration radius \(=0.5625\) inch
\(\tan \varphi=\frac{0.5625}{76-2.377}=0.0076403\)
\(\varphi=0.0076401\) radian \(=0.4378^{\circ}\)
\(=\) subtended angle at source
\[
\begin{gather*}
y=(76-a) \tan \varphi  \tag{1}\\
y^{2}=(7.268)^{2}-(7.268-a)^{2} \tag{2}
\end{gather*}
\]

Solving Equations 1 and 2 we get
\(y=0.5807\) inch
\(\mathrm{a}=0.0232\) inch
The secondary obscuration diameter is therefore 1.161 inches.

APPENDIX 2

\section*{ENTRANCE PUPIL SIZE}

APPENDIX 2
ENTRANCE PUPIL SIZE


NOTE: NOT DRAWN TO SCALE
\(y_{S}=\) clear aperture radius of secondary \(=0.540\) inch
\(y_{P}=\) entrance pupil radius for point source at a distance \(S\) (on-axis)
\(R=\) radius of curvature of primary \(=7.268\) inches
\(b=\) distance of front of secondary mirror from vertex of primary \(=2.317\) inches
\(S=\) distance of source to primary vertex \(=76\) inches
\[
\begin{gather*}
a+\sqrt{R^{2}-y_{P}^{2}}=R  \tag{1}\\
\tan \alpha=\frac{y_{P}-y_{S}}{b-a}  \tag{2}\\
\sin \beta=\frac{y_{P}}{R}  \tag{3}\\
\gamma=\alpha-\beta \tag{4}
\end{gather*}
\]
\[
\begin{gather*}
\varphi=\alpha-2 \gamma=2 \beta-\alpha  \tag{5}\\
\frac{\mathrm{S}-\mathrm{R}}{\sin \gamma}=\frac{\mathrm{R}}{\sin \varphi} \tag{6}
\end{gather*}
\]

For \(y_{S}=0.540\) inch, \(R=7.268\) inches, \(b=2.317\) inches and \(S=76\) inches, the solution of the above six equations is
\[
\begin{aligned}
\alpha & =0.35848 \mathrm{rad} \\
\beta & =0.18823 \mathrm{rad} \\
\gamma & =0.170245 \mathrm{rad} \\
\varphi & =0.01798 \mathrm{rad} \\
\mathrm{a} & =0.1284 \mathrm{inch} \\
\mathrm{y}_{\mathbf{P}} & =1.360 \text { inches }
\end{aligned}
\]

The entrance pupil diameter for a source at 76 inches is therefore 2.720 inches. The angular radius of the entrance pupil is given by \(\varphi\left(0.01798\right.\) radian \(\left.=1.0302^{\circ}\right)\),

\section*{APPENDIX 3}

\section*{ENERGY LOSS DUE TO OUT-OF-FOCUS BLUR OF IMAGE}

\section*{APPENDIX 3 \\ ENERGY LOSS DUE TO OUT-OF-FOCUS BLUR OF IMAGE}

The energy falling on the detector was found by convolving the image blur annulus with the square field-of-view of the detector. The annular shape is due to the image of the circular primary mirror obscured by the circular secondary mirror. The convolution was averaged along three paths,
1. horizontally through the center of the detector's nominal FOV (along' the x -axis)
2. diagonally across the square FOV (at an angle of \(45^{\circ}\) with the x -axis)
3. at an angle of \(22.5^{\circ}\) with the \(x\)-axis

The results of the calculations are shown below.
\(\delta\) (degrees)
0
22.5
45
p
0.7773
0.7615
0.7213
\(\delta\) is the convolution path angle with respect to the \(x\)-axis and \(p\) is the relative energy falling on the detector. The average value for the three paths is 0.753 .

The calculation details are described below.

\section*{1. Horizontal Convolution Path}

Along this path the convolution is divided into three parts where
a. portions of both primary and secondary images lie outside the nominal FOV
b. only the primary mirror image extends beyond the nominal FOV
c. both primary and secondary images lie inside the nominal FOV.

The following symbols and values will be used for all convolution paths
\[
\begin{aligned}
& \mathrm{r}=\text { angular radius of secondary mirror }=0.4378^{\circ}(\text { see Appendix } 1) \\
& \mathrm{R}=\text { angular radius of primary mirror }=1.0302^{\circ}(\text { see Appendix 2) } \\
& \rho=\text { angular radius of in-focus image }=\tan ^{-1} 1 / 76=0.7538^{\circ}
\end{aligned}
\]
\(2 \alpha=\) angular, in-focus length of one side of the square field-of-view
\[
=2.1^{\circ}\left(\alpha=1.05^{\circ}\right)
\]
a. \(\alpha-\rho \leqslant \mathrm{x} \leqslant \mathrm{r}\)


The point shown at an angular distance x from the side of the nominal FOV lies within the in-focus image of radius \(\rho\).

If \(p_{x}\) is the relative image blur area that lies inside the square FOV for an in-focus image point at \(x\) then
\[
\mathrm{p}_{\mathrm{x} 1}=\frac{2\left(\mathrm{~A}_{1}-\mathrm{A}_{0}\right)+2\left(\mathrm{~A}_{\phi_{1}}-\mathrm{A}_{\phi_{0}}\right)}{\mathrm{A}_{\mathrm{R}}-\mathrm{A}_{\mathrm{T}}}
\]
where \(A_{0}\) and \(A_{1}\) are the triangular areas shown in the figure; \(A_{\phi_{0}}\) and \(A_{\phi_{1}}\) are the areas of the sectors shown; \(A_{R}\) is the area of the large circle (radius \(R\) ) and \(A_{r}\) is the area of the small circle (radius r).
\[
\begin{aligned}
& 2 \mathrm{~A}_{0}=\mathrm{x} \sqrt{\mathrm{r}^{2}-\mathrm{x}^{2}} \\
& 2 \mathrm{~A}_{1}=\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}} \\
& 2 \mathrm{~A}_{\phi_{0}}=\phi_{0} \mathrm{r}^{2} \\
& 2 \mathrm{~A}_{\phi_{1}}=\phi_{1} \mathrm{R}^{2} \\
& \theta_{0}=\cos ^{-1} \frac{\mathrm{x}}{\mathrm{r}} \\
& \theta_{1}=\cos ^{-1} \frac{\mathrm{x}}{\mathrm{R}} \\
& \phi_{0}=\pi-\theta_{0} \\
& \phi_{1}=\pi-\theta_{1} \\
& \mathrm{p}_{\mathrm{x} 1}=\frac{\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}}-\mathrm{x} \sqrt{\mathrm{r}^{2}-\mathrm{x}^{2}}+\left(\pi-\theta_{1}\right) \mathrm{R}^{2}-\left(\pi-\theta_{0}\right) \mathrm{r}^{2}}{\pi\left(\mathrm{R}^{2}-\mathrm{r}^{2}\right)}
\end{aligned}
\]
b. \(r \leqslant x \leqslant R\)
\[
\mathrm{p}_{\mathrm{x} 2}=\frac{\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}}+\left(\pi-\dot{\theta}_{.1}\right) \mathrm{R}^{2}-\pi \mathrm{r}^{2}}{\pi\left(\mathrm{R}^{2}-\mathrm{r}^{2}\right)}
\]
c. \(\mathrm{R} \leqslant \mathrm{x} \leqslant \alpha\)
\[
\mathrm{p}_{\mathrm{x} 3}=1
\]

Since there is more energy from a circular ring toward the edge of the source than toward the center the total energy falling on the detector must be weighted accordingly. If \(v\) is the (angular) distance of the in-focus image point from the center of the image of the source the total integrated relative energy \(p\) received by the detector is
\[
\mathrm{p}=\int_{0}^{\rho} \frac{2 \pi v}{\pi \rho^{2}} \mathrm{p}_{\mathrm{x}} \mathrm{dv}=\frac{2}{\rho^{2}} \int_{0}^{\rho} \mathrm{vp}_{\mathrm{x}} \mathrm{dv}
\]

Since
\[
\begin{gathered}
\mathrm{v}=\alpha-\mathrm{x} \\
\mathrm{dv}=-\mathrm{dx} \\
\mathrm{p}= \\
=-\frac{2}{\rho^{2}} \int_{\alpha}^{\alpha-\rho}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x}} \mathrm{dx}=\frac{2}{\rho^{2}} \int_{\alpha-\rho}^{\alpha}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x}} \mathrm{dx} \\
=\frac{2}{\rho^{2}} \int_{\alpha-\rho}^{\mathrm{r}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 1} \mathrm{dx}+\frac{2}{\rho^{2}} \int_{\mathrm{T}}^{\mathrm{R}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
\\
+\frac{2}{\rho^{2}} \int_{\mathrm{R}}^{\alpha}(\alpha-\mathrm{x}) \mathrm{dx}
\end{gathered}
\]

Let
\[
\begin{aligned}
I_{1} & =\frac{2}{\rho^{2}} \int_{\alpha-\rho}^{\mathrm{r}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 1} \mathrm{dx} \\
& =\frac{2}{\rho^{2}} \int_{0.2962}^{0.4378}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 1} \mathrm{dx} \\
\mathrm{I}_{2} & =\frac{2}{\rho^{2}} \int_{\mathrm{r}}^{\mathrm{R}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
& =\frac{2}{\rho^{2}} \int_{0.4378}^{1.0302}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
\mathrm{I}_{3} & =\frac{2}{\rho^{2}} \int_{\mathrm{R}}^{\alpha}(\alpha-\mathrm{x}) \mathrm{dx}
\end{aligned}
\]
\[
\begin{aligned}
& =\frac{2}{\rho^{2}} \int_{1.0302}^{1.05}(\alpha-\mathrm{x}) \mathrm{dx} \\
& =\frac{2}{(0.7538)^{2}}\left[\alpha \mathrm{x}-\frac{\mathrm{x}^{2}}{2}\right]_{1.0302}^{1.05} \\
& =0.0007
\end{aligned}
\]
\(I_{1}\) and \(I_{2}\) were integrated numerically and the results are
\[
\begin{aligned}
& I_{1}=0.2275 \\
& I_{2}=0.5491 \\
& p=I_{1}+I_{2}+I_{3}=0.7773
\end{aligned}
\]
2. Diagonal Path

Along the diagonal the image of the secondary remains entirely within the square FOV of the detector. However, it is still necessary to consider the following three cases:
a. The image of the primary extends beyond the corner of the square FOV.
b. The image of the primary extends beyond the sides of the FOV but not the corner.
c. The entire image of the primary lies within the detector's FOV.
a. \(\quad \alpha-\frac{\rho}{\sqrt{2}} \leqslant \mathrm{x} \leqslant \frac{\mathrm{R}}{\sqrt{2}}\)
\[
\begin{aligned}
\mathrm{p}_{\mathrm{x} 1} & =\frac{2\left(\mathrm{~A}_{1}+\mathrm{A}_{0}+\mathrm{A}_{\phi}\right)-\pi \mathrm{r}^{2}}{\pi\left(\mathrm{R}^{2}-\mathrm{r}^{2}\right)} \\
2 \mathrm{~A}_{1} & =\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}} \\
2 \mathrm{~A}_{0} & =\mathrm{x}^{2} \\
2 \mathrm{~A}_{\phi} & =\phi \mathrm{R}^{2}
\end{aligned}
\]

b. \(\frac{R}{\sqrt{2}} \leqslant x \leqslant R\)
\[
\mathrm{p}_{\mathrm{x} 2}=\frac{2 \mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}}+(\pi-2 \theta) \mathrm{R}^{2}-\pi \mathrm{r}^{2}}{\pi\left(\mathrm{R}^{2}-\mathrm{r}^{2}\right)}
\]
c. \(\mathrm{R} \leqslant \mathrm{x} \leqslant \alpha\)
\[
\mathrm{p}_{\mathrm{x} 3}=1
\]

The integrated energy along the diagonal is
\[
\mathrm{p}=\frac{2}{\rho^{2}} \int_{0}^{\rho} \mathrm{p}_{\mathrm{x}} \mathrm{vdv}
\]

For this case,
\[
\begin{aligned}
& v=\sqrt{2}(\alpha-x) \\
& p=-\frac{4}{\rho^{2}} \int_{\alpha}^{\alpha-\frac{\rho}{\sqrt{2}}} p_{x}(\alpha-x) \mathrm{dx} \\
&=\frac{4}{\rho^{2}} \int_{\alpha-\frac{\rho}{\sqrt{2}}}^{\alpha}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x}} \mathrm{dx} \\
&=\frac{4}{\rho^{2}} \int_{\alpha-\frac{\rho}{\sqrt{2}}}^{\frac{\mathrm{R}}{\sqrt{2}}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 1} \mathrm{dx}+\frac{4}{\rho^{2}} \int_{\frac{\mathrm{R}}{\sqrt{2}}}^{\mathrm{R}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
&+\frac{4}{\rho^{2}} \int_{\mathrm{R}}^{\alpha}(\alpha-\mathrm{x}) \mathrm{dx} \\
&= \mathrm{I}_{1}+\mathrm{I}_{2}+\mathrm{I}_{3} \\
& \mathrm{I}_{1}=\frac{\mathrm{R}}{\rho^{2}} \int_{\alpha}^{\frac{\mathrm{R}}{\sqrt{2}}}\left(\alpha-\frac{\rho}{\sqrt{2}}\right.
\end{aligned}
\]
\[
\begin{aligned}
& \mathrm{I}_{2}=\frac{4}{\rho^{2}} \int_{\frac{\mathrm{R}}{\sqrt{2}}}^{\mathrm{R}}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx}=\frac{4}{\rho^{2}} \int_{0.7285}^{1.0302}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
& \mathrm{I}_{3}=\frac{4}{\rho^{2}} \int_{\mathrm{R}}^{\alpha}(\alpha-\mathrm{x}) \mathrm{dx}=\frac{4}{\rho^{2}} \int_{1.0302}^{1.05}(\alpha-\mathrm{x}) \mathrm{dx} \\
& \mathrm{I}_{3}=\frac{4}{(0.7538)^{2}}\left|\alpha x-\frac{x^{2}}{2}\right|_{1.0302}^{1.05}=0.0014
\end{aligned}
\]

Numerical integration of \(I_{1}\) and \(I_{2}\) yields
\[
\begin{aligned}
\mathrm{I}_{1} & =0.4179 \\
\mathrm{I}_{2} & =0.3020 \\
\therefore \mathrm{p} & =0.7213
\end{aligned}
\]
3. Convolution Path Angle of \(22.5^{\circ}\)

Along this path five cases are considered.
a. The primary mirror image extends beyond the corner of the nominal FOV and the secondary mirror image extends beyond one side of the FOV.
b. Same as above except that the secondary mirror image lies entirely within the square FOV.
c. The primary mirror image extends beyond two sides of the square FOV but not beyond the corner.
d. The primary mirror image extends beyond only one side of the FOV.
e. The primary mirror image lies entirely within the FOV.
a. \(\alpha-\rho \cos \delta \leqslant \mathrm{x} \leqslant \mathrm{r}\)

\[
p_{x 1}=\frac{y / 2 \sqrt{R^{2}-y^{2}}+y x+x / 2 \sqrt{R^{2}-x^{2}}+\phi / 2 R^{2}-\left(\pi-\theta_{0}\right) r^{2}}{2.7321}
\]
where
\[
\begin{aligned}
& y=\alpha-(\alpha-x) \tan \delta \\
& \phi=\frac{3 \pi}{2}-(\theta+\beta) \\
& \theta=\cos ^{-1} \frac{y}{R}
\end{aligned}
\]
\[
\begin{aligned}
& \beta=\cos ^{-1} \frac{x}{R} \\
& \theta_{0}=\cos ^{-1} \frac{x}{r}
\end{aligned}
\]

Note that \(\pi\left(\mathrm{R}^{2}-\mathrm{r}^{2}\right)=2.7321\)
b. \(r \leqslant x \leqslant x_{0}\)
\(\mathrm{x}_{0}\) is the value of x for which w (see preceding figure) is equal to R. \(x_{0}\) can be found from the following equation:
\[
\mathrm{w}^{2}=\mathrm{x}^{2}+\mathrm{y}^{2}=\mathrm{R}^{2}=\mathrm{x}^{2}+[\alpha-(\alpha-\mathrm{x}) \tan \delta]^{2}
\]

For \(\delta=22.5^{\circ}, \alpha=1.05, \mathrm{R}=1.0302\), the solution is
\[
\begin{gathered}
x_{0}=0.5764 \\
p_{x 2}=\frac{y / 2 \sqrt{R^{2}-y^{2}}+y x+x / 2 \sqrt{R^{2}-x^{2}}+\phi / 2 R^{2}-\pi r^{2}}{2.7321}
\end{gathered}
\]
c. \(\quad 0.5764 \leqslant \mathrm{x} \leqslant \mathrm{x}_{1}\)
\(x_{1}\) is the value of \(x\) for which \(y=R\). It can be found using the equation
\[
y=R=\alpha-(\alpha-x) \tan \delta
\]

For \(\delta=22.5^{\circ}, \alpha=1.05, \mathrm{R}=1.0302\)
\[
\begin{gathered}
x_{1}=1.0022 \\
p_{x 3}=\frac{x \sqrt{R^{2}-x^{2}}+y \sqrt{R^{2}-y^{2}}+\phi / 2 R^{2}+\gamma / 2 R^{2}-\pi r^{2}}{2.7321}
\end{gathered}
\]
where
\[
\begin{aligned}
& \phi=\frac{3 \pi}{2}-(\theta+\beta) \\
& \gamma=\frac{\pi}{2}-(\theta+\beta)
\end{aligned}
\]

\[
\begin{gathered}
\therefore \frac{\phi}{2} \mathrm{R}^{2}+\frac{\gamma}{2} \mathrm{R}^{2}=(\pi-\theta-\beta) \mathrm{R}^{2} \\
\mathrm{p}_{\mathrm{x} 3}=\frac{\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}}+\mathrm{y} \sqrt{\mathrm{R}^{2}-\mathrm{y}^{2}}+(\pi-\theta-\beta) \mathrm{R}^{2}-\pi \mathrm{r}^{2}}{2.7321}
\end{gathered}
\]
d. \(\quad 1.0022 \leqslant x \leqslant R\)
\[
\mathrm{p}_{\mathrm{x} 4}=\frac{\mathrm{x} \sqrt{\mathrm{R}^{2}-\mathrm{x}^{2}}+(\pi-\beta) \mathrm{R}^{2}-\pi \mathrm{r}^{2}}{2.7321}
\]
e. \(R \leqslant x \leqslant \alpha\)
\[
\mathrm{p}_{\mathrm{x} 5}=1
\]

The integrated energy along \(\delta\) is
\[
\mathrm{p}=\frac{2}{\rho^{2}} \int_{0}^{\rho} \mathrm{p}_{\mathrm{x}} \mathrm{dx}
\]
where now
\[
\begin{aligned}
& v=\frac{\alpha-x}{\cos \delta} \\
& \mathrm{~d}= \frac{2}{\rho^{2} \cos ^{2} \delta} \int_{\alpha-\rho \cos \delta}^{\alpha}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x}} \mathrm{dx} \\
&= \frac{2}{\rho^{2} \cos ^{2} \delta} \int_{0.3535}^{0.4378}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 1} \mathrm{dx}+\frac{\mathrm{dx}}{\rho^{2} \cos ^{2} \delta} \int_{0.4378}^{0.5764}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 2} \mathrm{dx} \\
&+\frac{2}{\rho^{2} \cos ^{2} \delta}-\int_{0.5764}^{1.0022}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 3} \mathrm{dx}+\frac{2}{\rho^{2} \cos ^{2} \delta} \int_{1.0022}^{1.0302}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 4} \mathrm{dx} \\
&+\frac{2}{\rho^{2} \cos ^{2} \delta} \int_{1.0302}^{1.05}(\alpha-\mathrm{x}) \mathrm{dx} \\
&= \mathrm{I}_{1}+\mathrm{I}_{2}+\mathrm{I}_{3}+\mathrm{I}_{4}+\mathrm{I}_{5}
\end{aligned}
\]
\[
\text { For } \delta=22.5^{\circ} \text {, }
\]
\[
\begin{aligned}
& I_{1}=4.1237 \int_{0.3535}^{0.4378}(\alpha-x) p_{x 1} d x=0.1490 \\
& I_{2}=4.1237 \int_{0.4378}^{0.5764}(\alpha-x) p_{x 2} d x=0.2148
\end{aligned}
\]
\[
\begin{gathered}
I_{3}=4.1237 \int_{0.5764}^{1.0022}(\alpha-x) \mathrm{p}_{\mathrm{x} 3} \mathrm{dx}=0.3930 \\
\mathrm{I}_{4}=4.1237 \int_{1.0022}^{1.0302}(\alpha-\mathrm{x}) \mathrm{p}_{\mathrm{x} 4} \mathrm{dx}=0.0039 \\
I_{5}=4.1237 \int_{1.0302}^{1.05}(\alpha-\mathrm{x}) \mathrm{dx}=0.0008 \\
\therefore \mathrm{p}=0.7615
\end{gathered}
\]

\section*{APPENDIX 4}

\section*{SUMMARY OF VTPR OPTICAL SYSTEM}

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\section*{1. Primary Mirror}

Clear aperture diameter: 2.970 inches
Radius of curvature: 7.268 inches
Vertex thickness: 0.060 inches
Hole diameter: 0.60 inches
Distance of rim ray focus from vertex: 3.555 inches
Effective focal length: 3.736 inches
2. Secondary Mirror

Clear aperture diameter: 1.080 inches
Obscuration diameter: 1.125 inches
Radius of curvature: 4.886 inches
Rim thickness: 0.060 inches
Distance between primary and secondary: 2.317 inches (variedforbestfocus)
3. Telescope

Effective focal length: 7.955 inches
Distance between focus and back of primary: 0.465 inches
Distance between focus and front of primary: 0.525 inches
4. Field stop size: \(0.700 \times 0.700\) inches

The optical layout is shown in the following figure.


Optical Layout - VTPR```

