# SOLPRO: A COMPUTER CODE TO CALCULATE PROBABILISTIC ENERGETIC SOLAR PROTON FLUENCES 

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# SOLPRO: A Computer Code to Calculate Probabilistic Energetic Solar Proton Fluences 

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#### Abstract

A code has been developed for the calculation of interplanetary solar proton iluences at 1 A . U. for the active years 1977-1983. The fluences . are presented as functions of mission duration $t$, energy threshold $E$, and confidence level $Q$. For a given combination of $\tau$ and $Q$, the routine determines whether ordinary or anomalously large events are to be considered, and in the latter case, the number of anomalously large events that are predicted by probabilistic theory for the specified mission duration. The code is described in detail. A listing and sample calculations are attached.


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The code was originally developed as a supplement to UNIFLUX, a unified orb'tal flux integration system (Stassinopoulos and Gregory, 1975), from an Initial solar proton model (Stassinopoulos and King, 1974) designed exclusively for the evaluation of flux levels to be encountered during the active phase of the next solar cycle (1977-1983) by earth-orviting spacecraft on missions involving partial magnetospheric shielding.

This first empirical model was characteristically simple and easy to use. It was derived entirely from experimental satellite measurements made during the active years of solar cycle 20 ; it included only crude confidence level estimates, that is, the probability of actual cycle 21 fluxes exceeding the predicted intensities.

A more probabilistic analysis of solar cycle 20 fluxes was later given by King (1974). This formal statistical treatment was subsequently used by King and Stassinopoulos (1975) in their computation of ratios of solar to tripped proton fluences for circular-orbit geocentric space missions to be flown during the active years of the next solar cycle (1977-1983). In this more sophisticated model the probability of exceeding mission fluence levels is given as a function of mission duration and energy.

The new SOLPRO code incorporates all the important features of the King (1974) analysis, which includes the distinction betwén "ordinary" (OR) and "anomalously large" (AL) events, and the probability of occurrence of the latter. Be it noted, that only one anomalously large event
was observed during solar cycle 20: the event of August, 1972; according to King (1974), the proton fluences of that one event alone constituted about $69 \%(E>10 \mathrm{MeV})$ and $84 \%(E>60 \mathrm{MeV})$ of the entire cycle fluence. Consequently, whenever at least one AL event is predicted, the or-eventcontribution is negligible.

The referenced solar proton models are based on the most continuous set of satellite data available to date, covering the period of enhanced solar activity during the 20 th solar cycle (1964-1975) through the 1972 AL event, after which no significant events occurred. In both models, all fluxes were taken to be isotropic. As noted in Stassinopoulos and King (1974), on an event-integrated basis, departures from isotropy are typically only a few percent or less. This applies to the interplanetary medium and to much of the magnetosphere, but may not apply at low altitudes where anisotropy should result from atmospheric loss mechanisms.

When using the SOLPRO routine, two important items should be kept in mind. First, there is no assurance that the overall flux levels observed during the 20 th solar cycle will also occur during the 21 st cycle. However, as indicated in King (1974), a comparison of the annual mean sunspot numbers of past solar cycles suggests that cycle 21 will most likely be similar to the very ordinary cycle 20 than to the very extraordinary cycle 19. Second, there is no reliable way of predicting the distribution of individual solar events in time, in flux level, and in spectra through the $2 l s t$ cycle. That pertains to ordinary as well as to anomalously large events. Therefore, the results obtained from the routine, although the best predictions available of solar proton fluences expected to occur during 1977-1983, should be considered guideline spproximations only.

In the code, the solar proton fluences are expressed as functions of the parameters "confidence level", $Q^{*}$ (in percent), "mission duration", $T$ (in months), and energy, $E$ (in MeV). Within the limits established by the routine the user is free to select his own values of $Q$ and $\boldsymbol{f}$ for process?ng, but $E$ is not an input parameter and the calculations are performed at ten fixed energy thresholds.

For missions whose trajectories will involve significant time away from a heliocentric distance of 1 A.U. (astronomical unit), a helforadial dependence of event fluences must be allowed for, as discussed by King (1974).

Similarly, for missions whose trajectories will involve partial magnetospheric shielding, the fractional exposure factor has to be' determined, as discussed by Stassinopoulos and King (1974). An independcat code performing these shielding calculations has been developed (Stas inopoulos, 1975) and is available from the National Space Science Data Center, Grcenbelt, Maryland.

A detailed description of the solar proten routine is given in the next section on "Method and Procedure", with an analysis of the method employed and a review of the organization and structure of the code. In section "Application and Use" the routine is briefly discussed, its operation is described, and the arguments of its transfer vector are presented. Coje listing and sample outputs are given in the attachments.

[^0]SOLPRO is a short, compact routine which evaluates omnidirectional, unattenuated, interplanetiry energetic solar proton fluences in terms of three variables: $Q, \tau$, and $E$. The first two are input parameters, the last is internally programed. As defined elsewhere in this report, $Q$ denotes the degree of confidence one wishes to assign to the results, namely that for the specified mission duration the calculated fluences at ten preset energy values ranging from 10 to 100 MeV in increments or 10 MeV , are the smallest values which will not be exceeded by actually encountered intensities.

In order to hold the size and complexity of the code to a miniram, the desirable elements of the King (1974) analysis were empirically formulated by convenient approximation techniques, which are able to reproduce the statistically treated King results with relatively high accuracy. Specifically, tie maximum error introduced into the statistics by this method is less than 50\%. This error is indeed insignificant in view of the large intrinsic statistical uncertainties of the models themselves and the much greater uncertainties involved in predicting cycle 21 fluxes from cycle 20 data.

The most important feature of the code is the distinction between OR and AL events. This capability was developed re?ying entirely on King's probabilistic treatment and using the data given therein (King, 1974; Figure 5). For simplicity let:

$$
\begin{equation*}
n=f(Q, T) \tag{1}
\end{equation*}
$$

where $n$ is the number of $A L$ events expected to occur with a confidence level $Q$ and for a mission duration $t$. It is then convenient to effect an apparent separation of variables by evaluating $n$ independently for every value of $Q$ within the regime of interest, and by expressing $n$ in terms of $t$ alone. Thus:

$$
\begin{equation*}
n_{i}=c+\sum_{j=0}^{6} a_{i j}(Q)_{\tau} j \tag{2}
\end{equation*}
$$

where the index $i$ relates $n$ to $Q$ through:

$$
\begin{equation*}
1=100-Q \tag{3}
\end{equation*}
$$

and where the polynomial expansion coefficients $a_{i j}$ are determined by curve fitting. The constant, $c=1.001$, is a dimensionless range adjustment term.

Table $l$ contains the numerical values of the $a_{i j}$ coefficients.

The code calculates $n$ at the very beginning from equation 2 (see Table 2 for obtained values of $n$ ) and, depending on whether AL events are expected to occur or not, it branches to its appropriate section, where it computes the corresponding $A L$ or $O R$ solar proton fluences for the specified $Q$ and $T$.

For AL-event conditions, that is $n>0$, the fluences are obtained from the expression:

$$
\begin{equation*}
F_{A L}(>E ; Q, T)=S_{A L}(>E) * \cap(Q, \tau) \tag{4}
\end{equation*}
$$

where the function:

$$
\begin{equation*}
S_{A L}(>E)=7.9 \times 10^{9} \exp \{(30 .-E) / 26.5\} \tag{5}
\end{equation*}
$$

is King's (1974) exponential-in-energy representation of the integral flux and spectrum of the Augusts. 1972, anomalously large event, for energy thresholds between 10 and 200 MeV .

In case $O R$ events only are predicted ( $n=0$ ), the respective fluences are obtained from:

$$
\begin{equation*}
F_{O R}(>E ; Q, \tau)=S_{O R}(>E) * J(Q, \tau) \tag{6}
\end{equation*}
$$

where now King's (1974) expression:

$$
\begin{equation*}
S_{\because \cdot R}(>E)=\exp \{.0158(30 .-E)\} \tag{7}
\end{equation*}
$$

for the integral energy spectrum of ordinary events is used, reliable over the range $10 \leq E(\mathrm{MeV}) \leq 100$.

The function $J(Q, \tau)$ gives the ordinary fluence of protons above 30 MeV in terms of the two input variables (see Table 3 for obtained values of $J$ ). It was developed along similar lines as equation 1 , with data from the sarie source (King, 1974; Figure 8).

Again, through the independent evaluation of $J$, (but this time for every $t$ within the relevant interval of $1-9$ months for these events) the same apparent separation of variables is effected, resulting in:

$$
\begin{equation*}
J_{k}=d \sum_{m=0}^{4} b_{k m}(\tau) Q^{m} \tag{8}
\end{equation*}
$$

where the index $m$ relates $J$ to $T$ through:

$$
\begin{equation*}
m=\text { integer }(T) \tag{9}
\end{equation*}
$$

and where the coefficients, contained in Table 4, are also determined
by curve fitting. The constant, $d=10^{7}$, is a scaling factor.

The code is structurally complete and self-contained. All expansion coefficients for the statistical treatment are stored in data statements at the beginning of the routine, while the expressions for the AL and OR models are included in the respective sections of the code.

To guard against accidental misuse and so as to insure that calculations are only performed within the valid ranges of the independent variables $Q$ and $\tau$, SOLPRO tests the values of these input parameters in its first executable statement. If the parameter(s) exceeds the program limits, it prints a warning message before retirning control to the calling routine.

Finally, the code bypasses ordinary event fluence calculations, although it branches to the appropriate section, whenever $\tau=1$ and $Q$ has values ranging from 80 to 83. In tha: domain the flux gradients are very high, the intensities dfcrease very rapidly towards $Q=80$ and the galactic cosmic ray component becomes predominant (see discussion by King, 1974). For all practical purposes the OR-event solar protons are negligible for $Q<80$ and $\tau>1$.

## Results and Discussion


#### Abstract

Sample calculations performed with the code indicate good agreement with probabilistic predictions. That is, the results represent a reasonable approximation to these data, considering the large intrinsic limitations and uncertainties inherent in this type of estimation process.


Thus, as shown in Figure 1 , the ordinary event fluence $J$ of protons with energies $\mathrm{E}>30 \mathrm{MeV}$, computed by the routine for the applicable $Q$ and T ranges, compares well with the original data in King (197.t; Figure 8); the solid curves indicate the former, the doted curves the latter. Apparently, the curves overlap for most $Q$ and $\tau$ values. In the few places where they do not agree, the differences are truly insignificant; for eximple, the average error for all $\tau>1$ is about $3 \%$, while the maximum error, which occurs at $\tau=1$ for $Q=85$ is less than $50 \%$.

Similarly, the number of anomalously large events $n$ given by the code agrees favorably with the model (King, 1974; Figure 5). This 15 shown in Figure 2, where the dots indicate model data and the triangles SOLPRO predictions. Since the trarsition from one $n$ value to another is probabilisticaliy (and not deterministically) defined, the differences between the King model and the SOLPRO approximation are inconsequential.

Finally, the $A L$ event exponential-in-energy function (Equation 5) and the ordinary-event integral-energy spectrum (Equation 7 ) are identical to the formulae given by King (1974). Their respective jalues are listed in Table 5.

Subroutine SOLPRO is written in standard FORTRAN-IV computer language and card decks are available in either the 029 model IBM keypunch format (EBCDIC) or the 026 keypunch format (BCD), for use with FORTRAN-compatible compilers.

The variables in the calling sequence are:
SOLPRO(TAU,IO,F)
The first two arguments are input data:
TAU : mission duration in months, with a valid range
from 1. to 72. (REAL*4)
IQ : percent confidence level that calculated fluence will not be exceeded, with a valid range from 80 to 99 (INTEGER*4)

The last argument is output data:
$F(N)$ : solar proton spectrum for energies $E>10 * N M e V$, where index $N$ goes from 1 to 10.

All ainments except $I Q$ are single precision floating point variables. Besides their range limits, there are no other restrictions or limitations on the values of the input arguments.

There are no READ or other input statements in SOLPRO. A WRLTE statement is executed each time an input variable exceeds the stipulated limits. In addition, a WRITE statement is executed whenever for a given combination of $T$ and $Q$ no significant solar proton fluences are to be expected.

Figure 3 shows a flow diagram of the program. A complete listing is given in the attachment. The cards of the deck are appropriately labeled in columns $73-80$ as SOLPR xxx where the three last columns (xxox) contain the sequential numbering, which is incremented by 10.

## Acknowledgements

The author wishes to thank Das. J. I. Vette and J. H. King for many helpful suggestions and comments.

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Table 1

## NUMBER OF ANOMALOUSLY LARGE EVENTS

## Polynomial Expansion Coefficients

| Q (\%) | ${ }^{2} 0$ | $\mathrm{a}_{1}$ | ${ }^{2}$ | ${ }^{3}$ | ${ }_{4}$ | ${ }^{5}$ | ${ }^{6}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 99 | -. 1571 | . 2707 | -. $1269^{-1}$ | $.4428^{-3}$ | -.8185 ${ }^{-5}$ | $.7754^{-7}$ | -. $2939-9$ |
| 98 | -. 1870 | . 1951 | -.6559-2 | $.1990^{-3}$ | -. $3618^{-5}$ | $.3740^{-7}$ | -. $1599^{-9}$ |
| 97 | -. 2007 | . 1497 | -. $3179-2$ | $.5730^{-4}$ | -. $4664^{-6}$ | . $1764^{-8}$ | 0 |
| 96 | -. 1882 | . 1228 | -. $1936{ }^{-2}$ | $.2660^{-4}$ | -. $1022^{-6}$ | 0 | 0 |
| 95 | -. 2214 | . 1149 | -. 1871-2 | . $2695^{-4}$ | -. $11116^{-6}$ | 0 | 0 |
| 94 | -. 2470 | . 1062 | -. $16588^{-2}$ | $.2367^{-4}$ | -. $9465^{-7}$ | 0 | 0 |
| 93 | -. 2509 | $.8710^{-1}$ | $-.8300^{-3}$ | $.8438^{-5}$ | 0 | 0 | 0 |
| 92 | -. 2923 | . $8933{ }^{-1}$ | -. $1023^{-2}$ | $.1029^{-4}$ | 0 | 0 | 0 |
| 91 | -. 3222 | . $8648^{-1}$ | -. $99922^{-3}$ | $.9935^{-5}$ | 0 | 0 | 0 |
| 90 | -. 3518 | $.8417^{-1}$ | -. $1000{ }^{-2}$ | . $9956^{-5}$ | 0 | 0 | 0 |
| 89 | -. 3698 | $.7951^{-1}$ | -. $89833^{-3}$ | $.8940^{-5}$ | 0 | 0 | 0 |
| 88 | -. 2771 | $.5473^{-1}$ | -. 1543-4 | 0 | 0 | 0 | 0 |
| 87 | -. 2818 | $.5072^{-1}$ | $.2511^{-4}$ | 0 | 0 | 0 | 0 |
| 86 | -. 2845 | $.4717^{-1}$ | . $5664^{-4}$ | 0 | 0 | 0 | 0 |
| 85 | -. 2947 | $.4405^{-1}$ | $.8507^{-4}$ | 0 | 0 | 0 | 0 |
| 84 | -. 2923 | $.4111^{-1}$ | $.1106^{-3}$ | 0 | 0 | 0 | 0 |
| 83 | -. 2981 | $.3853^{-1}$ | $.1312^{-3}$ | 0 | 0 | 0 | 0 |
| 82 | -. 3002 | . $3585{ }^{-1}$ | $.1529{ }^{-3}$ | 0 | 0 | 0 | 0 |
| 81 | -. 3001 | $.3312^{-1}$ | $.1781^{-3}$ | 0 | 0 | 0 | 0 |
| 80 | -. 3141 | $.3248^{-1}$ | $.1654^{-3}$ | 0 | 0 | 0 | 0 |

## Table 2

Number of Anomalously Large Events Predicted by SOLPRO

*NO CALCULATIONS PERFORMED

Table 3

OR-Event Fluences ( $\mathrm{E}>30 \mathrm{MeV}$ ) given by SOLPRO Routine

| 0 | $\tau=1$ | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 99 | ----- | ----- | ----- | -----* | ----- | ----- | ----- | ----- | -----* |
| 98 | ----- | ---*- | -- | ----- | ----- | -- | - | -- | -* |
| 97 | $5.07 \times 10^{8}$ | --..-- | ------ | -.---- | ---.-* | - | - | - | -- |
| 96 | $3.42 \times 10^{8}$ | - | - | - | -- | -* | -- | - | --.-.- |
| 95 | $2.34 \times 10^{8}$ | --- | ---- | -- | -- | - | - | -- | ----- |
| 94 | $1.67 \times 10^{8}$ | $4.83 \times 10^{8}$ | ----- | ----- | ---- | - | -- | - | -.-.... |
| 93 | $1.30 \times 10^{8}$ | $3.91 \times 10^{8}$ | ----- | - | -- | -- | -...-.. | - | ----- |
| 92 | $1.09 \times 10^{8}$ | $3.23 \times 10^{8}$ | $6.46 \times 10^{8}$ | ----- | - | ---- | - | - | -...-- |
| 91 | $9.78 \times 10^{7}$ | $2.74 \times 10^{8}$ | $5.09 \times 10^{8}$ |  | - | - | - | -.-.-- | ---.-- |
| 90 | $8.85 \times 10^{7}$ | $2.39 \times 10^{8}$ | $4.15 \times 10^{8}$ | $7.82 \times 10^{8}$ | ------ | -----. | -...-- | .....- | -...-. |
| 89 | $7.72 \times 10^{7}$ | $2.12 \times 10^{8}$ | $3.53 \times 10^{8}$ | $6.02 \times 10^{8}$ | --.--- | ----* | --...- | ------ | --..-- |
| 88 | $6.20 \times 10^{7}$ | $1.91 \times 10^{8}$ | $3.14 \times 10^{8}$ | $4.84 \times 10^{8}$ | $9.16 \times 10^{8}$ | --..-- | -..... | --.... | ---.-- |
| 87 | $4.34 \times 10^{7}$ | $1.73 \times 10^{8}$ | $2.88 \times 10^{8}$ | $4.12 \times 10^{8}$ | $6.76 \times 10^{8}$ | ---.-- | ------ | ------ | --..-- |
| 86 | $2.40 \times 10^{7}$ | $1.56 \times 10^{8}$ | $2.70 \times 10^{8}$ | $3.68 \times 10^{8}$ | $5.44 \times 10^{8}$ |  | ------ | --..-. | --...- |
| 85 | $8.86 \times 10^{6}$ | $1.38 \times 10^{8}$ | $2.53 \times 10^{8}$ | $3.42 \times 10^{8}$ | $4.74 \times 10^{8}$ | $7.52 \times 10^{8}$ | - | .....- | -..... |
| 84 | $5.23 \times 10^{6}$ | $1.19 \times 10^{8}$ | $2.36 \times 10^{8}$ | $3.23 \times 10^{8}$ | $4.32 \times 10^{8}$ | $6.25 \times 10^{8}$ |  | ------ | -...-. |
| 83 | * | $1.00 \times 10^{8}$ | $2.17 \times 10^{8}$ | $3.04 \times 10^{8}$ | $3.98 \times 10^{8}$ | $5.39 \times 10^{8}$ | $8.06 \times 10^{8}$ | ----- | -----. |
| 82 | * | $8.16 \times 10^{7}$ | $1.95 \times 10^{8}$ | $2.83 \times 10^{8}$ | $3.60 \times 1 C^{8}$ | $4.80 \times 10^{8}$ | $6.94 \times 10^{8}$ | $1.01 \times 10^{9}$ | ------ |
| 81 | * | $6.56 \times 10^{7}$ | $1.74 \times 10^{8}$ | $2.58 \times 10^{8}$ | $3.22 \times 10^{8}$. | $4.38 \times 10^{8}$ | $6.08 \times 10^{8}$ | $8.38 \times 10^{8}$ | ----- |
| 80 | * | $5.48 \times 10^{7}$ | $1.56 \times 10^{8}$ | $2.32 \times 10^{8}$ | $2.98 \times 10^{8}$ | $4.04 \times 10^{8}$ | $5.41 \times 10^{8}$ | $7.34 \times 10^{8}$ | $1.01 \times 10^{9}$ |

- No Calculations performed: see couments page 7.

Table 4

## ORDINARY EVENT FLUENCE

Polynomial Expansion Coefficients

| $\boldsymbol{t}$ (mths) | $b_{0}$ | $b_{1}$ | $b_{2}$ | $b_{3}$ | $b_{4}$ |
| :--- | :--- | :--- | :--- | :--- | :--- |


| 1 | $.154047 \times 10^{3}$ | $-.522258 \times 10^{4}$ | $.714275 \times 10^{5}$ | $-.432747 \times 10^{6}$ | $.955315 \times 10^{6}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | $.198004 \times 10^{3}$ | $-.448788 \times 10^{4}$ | $.438148 \times 10^{5}$ | $-.196046 \times 10^{6}$ | $.325520 \times 10^{6}$ |
| 3 | $.529120 \times 10^{3}$ | $-.122227 \times 10^{5}$ | $.112869 \times 10^{6}$ | $-.465084 \times 10^{6}$ | $.710572 \times 10^{6}$ |
| 4 | $.121541 \times 10^{4}$ | $-.266412 \times 10^{5}$ | $.226778 \times 10^{6}$ | $-.857230 \times 10^{6}$ | $.120444 \times 10^{7}$ |
| 5 | $.452062 \times 10^{4}$ | $-.103248 \times 10^{6}$ | $.896085 \times 10^{6}$ | $-.346028 \times 10^{7}$ | $.499852 \times 10^{7}$ |
| 6 | $.272028 \times 10^{4}$ | $-.499088 \times 10^{5}$ | $.353050 \times 10^{6}$ | $-.111929 \times 10^{7}$ | $.133386 \times 10^{7}$ |
| 7 | $.275595 \times 10^{4}$ | $-.469718 \times 10^{5}$ | $.314729 \times 10^{6}$ | $-.960383 \times 10^{6}$ | $.111650 \times 10^{7}$ |
| 8 | $.570997 \times 10^{4}$ | $-.799689 \times 10^{5}$ | $.381074 \times 10^{6}$ | $-.610714 \times 10^{6}$ | 0 |
| 9 | $.101000 \times 10^{3}$ | 0 | 0 | 0 | 0 |

## Table 5

| Integral Flux \& Spectrum |  | Distribution of |  |
| :---: | :---: | :---: | :---: |
| of the August 1972 AL Event $\quad$ OR-Event Integral Spectrum |  |  |  |
| (Equation 5) |  | (Equation 7) |  |
| $\frac{\text { Energy }}{(\mathrm{MeV})}$ | $\frac{\text { Solar Protons ( }>E \text { ) }}{\left(\# / \mathrm{cm}^{2} / \text { AL-event }\right)}$ | $\frac{\text { Energy }}{(\mathrm{MeV})}$ | $\frac{\text { Spectral }}{}$ Function |
| 10 | $1.680 \times 10^{10}$ | 10 | 1.37163 |
| 20 | $1.152 \times 10^{10}$ | 20 | 1.17117 |
| 30 | $7.900 \times 10^{9}$ | 30 | 1.00000 |
| 40 | $5.417 \times 10^{9}$ | 40 | 0.85385 |
| 50 | $3.714 \times 10^{9}$ | 50 | 0.72906 |
| 60 | $2.547 \times 10^{9}$ | 60 | 0.62251 |
| 70 | $1.746 \times 10^{9}$ | 70 | 0.53153 |
| 80 | $1.197 \times 10^{9}$ | 80 | 0.45384 |
| 90 | $8.210 \times 10^{8}$ | 90 | 0.38752 |
| 100 | $5.629 \times 10^{8}$ | 100 | 0.33088 |



Figure 1. OR-Event Fhuences ( $\mathbf{E} \mathbf{>} \mathbf{3 0} \mathrm{MoV}$ ): Comperison of Predictions


Figure 2. Number of Anomalously Lerge Events: Comparison of Predictions


Figure 3. Flow Diagram for SOLPRO

## Sample Program Listing



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## Sample Calculations: Sola* Proton Fluences $F(>E ; Q, \tau)$

## ORIGRAL rAGE LS OF POOR QUALITY





[^0]:    ${ }^{*} Q$ is the degree of confidence a mission planner wishes to have that the actiduley encountered solar proton fluence will not exceed the predicicd fluence, as given by the statistical solar proton model used.

