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## EXPERIMENTAL STUDIES OF FLOW SEPARATION AND STALLING ON A TWO-DIMENSIONAL AIRFOIL AT LOW SPEEDS

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## SUMMARY

Detailed measurements of flow fields associated with low speed turbulent boundary layers have been made for the 17\% thick GA(W)-1 airfoil section at nominal angles of attack of $10^{\circ}, 14^{\circ}$, and $18^{\circ}$, Reynolds number $2.2 \times 10^{6}$, and Mach number 0.13. The data includes extensive pressure and velocity surveys of the pre- and post-separated regions on the airfoil and the associated wake. The boundary layer characteristics including regions of separation on the airfoil are also presented.

The results indicate steep gradients of displacement thickness, momentum thickness, shape factor and the separation streamline from the point of separation to the trailing edge of the airfoil. The present tests reveal that the region of flow reversal terminates within a surprisingly short distance of less than $20 \%$ chord downstream from the trailing edge for the test range of angle of attack.

## INTRODUCTION

The complex phenomena of the interaction between potential and viscous flow fields associated with airfoils has been mathematically modeled by various research groups (Refs. 1, 2, 3). The results of these computer programs show that with the exception of drag estimation the modeling is highly accurate for prediction of the performance of airfoils up to the onset of flow separation. References 1 and 2, however, cannot predict the post-separated boundary layer behavior.

Several methods have been proposed to handle the postseparated boundary layer behavior (Refs. 3, 4, 5). References 3 and 4 base their analysis on the observed constant pressure distribution in the separated regions, whereas Reference 5 approaches the problem using an extension of boundary layer
theory to include regions of separation. These models work reasonably well when the depths of separated regions are small, but fail with large regions of separated flow.

The present experimental investigation is carried out in order to furnish a set of detailed boundary layer characteristics under the influence of extensive separated flow conditions on a low speed airfoil. This data is intended to help the development of better mathematical models for separated flows.

## SYMBOLS

To the maximum extent possible, physical measurements are presented in non-dimensional form. Where dimensional quantities are required, they are given in both International (SI) Units and U. S. Customary Units. Measurements were made in U. S. Customary Units. Conversion factors between SI Units and U.S. Customary Units are given in Reference 6. The symbols used in the present report are defined as follows:

| c | Wing chord |
| :---: | :---: |
| ${ }^{c^{\prime}}{ }_{s}$ | Static pressure coefficient, $\frac{\mathrm{P}_{\mathrm{s}}-\mathrm{P}_{\infty}}{\mathrm{q}_{\infty}}$ |
| ${ }^{c} p_{T}$ | Total pressure coefficient, $\frac{\mathrm{P}_{T}-\mathrm{P}_{\infty}}{\mathrm{q}_{\infty}}$ |
| H | Shape factor ( ${ }^{*} / \delta_{\text {** }}$ ) |
| $\mathrm{p}_{\mathrm{s}}$ | Local static pressure |
| $\mathrm{p}_{\text {T }}$ | Local total pressure |
| $\mathrm{P}_{\infty}$ | Free stream static pressure |
| $\mathrm{q}_{\infty}$ | Free stream dynamic pressure |
| RN | Reynolds number based upon wing chord |
| U | Velocity at the edge of the boundary layer, non-dimensionalized with respect to free stream velocity. |

Local velocity non-dimensionalized with respect to free stream velocity, $\sqrt{\frac{\mathrm{P}_{T}-\mathrm{P}_{S}}{\mathrm{q}_{\varphi}}}$

Component of local velocity in the free stream direction, non-dimensionalized with respect to free stream velocity.
Streamwise coordinate
Spanwise coordinate
Vertical coordinate
Angle of attack, degrees
Boundary layer thickness
Boundary layer displacement thickness, $\delta_{0}^{\delta}\left(1-\frac{u}{\bar{U}}\right) d z$
Boundary layer momentum thickness, $\int_{0}^{\frac{u}{U}}\left(1-\frac{u}{U}\right) d z$

APPARATUS AND PROCEDURE
TEST

The experimental investigations were carried out in the WSU $213 \mathrm{~cm} \times 305 \mathrm{~cm}\left(7^{\prime} \mathrm{x} 10^{\prime}\right)$ low speed wind tunnel fitted with a $213 \mathrm{~cm} x 91.4 \mathrm{~cm}\left(7^{\prime} \mathrm{x} 3^{\prime}\right)$ two-dimensional insert employing a $17 \%$ thick $G A(W)-1$ airfoil section having 61.0 cm (24") chord and a 91.4 ( $36^{\prime \prime}$ ) span. Details of the construction of the model, supporting disks and the surface pressure taps are given in Reference 7. All tests were conducted with flap nested and at angles of attack of $10.3^{\circ}, 14.4^{\circ}$, and $18.4^{\circ}$, which represent pre-stall, near-stall, and post-stall conditions respectively. Reynolds number of the test was $2.2 \times 10^{6}$ based on the airfoil chord and Mach number was 0.13. In order to ensure transition, $2.5 \mathrm{~mm}(0.1 ")$ wide strips of $\# 80$ carborundum grit were employed at $5 \%$ chord on both upper and lower surfaces (Ref. 8). In this test series details of the flow field were investigated only on the upper surface of the model. At each angle of attack pressure measurements were made at ten chordwise stations in order to study the behavior of the boundary layer upstream, downstream and near the separation point on
the airfoil. Pressure measurements for determining velocities were also made at four stations downstream of the trailing edge of the airfoil in order to obtain a complete picture of the development of the wake with flow reversal.

## INSTRUMENTATION

Velocities were obtained by employing three types of probes (Fig. 1). The five tube pressure sensing pitch-yaw probe of 3.2 mm (.125") diameter was used to obtain velocity profiles for heights more than 2.5 mm (.10") above the local surface of the airfoil. For heights of 5.4 mm (.213") and less, velocities were obtained by employing either a flat tube probe or a cylindrical tube probe for total pressure measurements and model surface orifices for local static pressure measurements. At certain downstream chordwise stations with shallow regions of separation, difficulty was experienced while scanning the reversed flows with the flat tube probe due to the interference of the probe stem on the upstream flow. Hence a cylindrical tube of a 0.46 mm diameter, sealed at one end, with a .3 mm dia. hole drilled at a distance of .38 mm from•the tip was used to scan the stations at and downstream of the separation point. The probes were fixed to a mechanism which has provisions for remote traversing in the vertical axis and remote rotation of $\pm 180^{\circ}$ about the probe axis. The traversing gear was mounted to an aluminum channel fixed to the tunnel ceiling. Thus the probes could be locked in position at any chordwise station between the leading edge and 200\% chord. The maximum vertical range available for probe travel varied from about 11.4 cm (4.5") to 39.4 cm (12") depending on the chordwise probe location. Unbonded strain gage pressure transducers with a range of $\pm 17.2$ Kilo-newtons $/ \mathrm{m}^{2}$ ( $\pm 2.5 \mathrm{psi}$ ) were utilized for all pressure measurements.

## METHODS

Surface pressures were obtained through a system of pressure switches and transducers with digital data recorded on punch cards.

Tuft surveys and oil flow methods were employed for observation of the surface flow patterns. In order to obtain qualitative information regarding the depth and the longitudinal growth of the reversed flow regions, tufted spokes of about 1.5 mm (.0625") diameter and 15.2 cm (6") long were fixed to the surface at every $10 \%$ chordwise station on the airfoil. These spokes were offset by about $3.8 \mathrm{~cm}\left(1.5^{\prime \prime}\right)$ in the spanwise direction in order to avoid interference of the tufts. The flow pattern downstream of the trailing edge and the wake was observed by employing tuft spokes at $100 \%$ chord, 110\% chord, and $120 \%$ chord stations. These spokes were located at the midspan and were supported by horizontal rods fixed to the endplate disks to which the model was fitted.

Flow field velocity data was acquired by initially tilting the five-tube probe to align with the local airfoil slope. The probe was then yawed into the plane of the local flow and all five pressures were recorded on punched cards along with probe position. Velocities close to the wall were obtained by aligning either the flat tube probe or the cylindrical tube probe with the direction of the local surface. While scanning the wake regions, the five tube probe was initially aligned horizontally (the free stream direction) and yawed as required into the local flow.

## DATA REDUCTION

Surface pressures, local velocities, and flow inclinations were calculated by the computer routine developed for the IBM 1130 computer at WSU. The local velocity is expressed in a non-dimensional form as the ratio of local to free stream velocity. Experimental velocity profiles were plotted by employing a computer plotting routine written for the IBM 1130.

The five tube probe was calibrated through a pitch angle range of $\pm 45^{\circ}$, under dynamic pressures of 479 and 2874 newtons per meter ${ }^{2}$ (10 and 60 psf) with probe extensions of 25.4 cm and $40.6 \mathrm{~cm}\left(10^{\prime \prime}\right.$ and $16^{\prime \prime}$ ) in order to explore the effects of dynamic
pressure and the support interference. No noticeable deviations were observed. Calibration curves were extrapolated up to a maximum pitch angle of $\pm 50^{\circ}$ to evaluate velocities. Based on the calibration data, polynomial curve fits were employed for subsequent use in the data reduction program. The inputs for the data reduction program are the five pressures, which are fed into the calibration equations to compute the pitch angle, static pressure, total pressure and velocity. For probe measurements near the airfoil surface (within 5 probes diameters), the image effect of the probe due to the proximity of the wall influences probe readings. Corrections for these effects are included in the computing routine. A more detailed account of the calibration procedure and the results is given in Reference 9.

The flow in the turbulent separated boundary layer and the wake is basically unsteady. The instrumentation employed in the present tests is heavily damped and, therefore, records time-averaged quantities. However, the readings from the five tube probe were repeatable within $\pm 5 \%$ even under extreme poststall conditions.

## RESULTS AND DISCUSSIONS

Separation point: Exploration of the point of separation was carried out at angles of attack of $10.3^{\circ}$, $14.4^{\circ}$ and $18.4^{\circ}$. Surface flow visualization methods and the tuft spoke survey were found to agree very well with the separation point obtained from the surface pressure distribution. However, the tuft spoke survey did not show realistic depths of separated regions. It is believed that mutual interference of the tufts resulted in magnified depth of the reversed flow regions.

Velocity plots: Computer plots of the velocity profiles at the mid-span section are shown in Figures 2, 3 and 4. Velocity
profiles obtained from flat tube and cylindrical tube are shown in Figures 5, 6 and 7. While measuring the total pressure values with flat tube and cylindrical tube probes, a problem was encountered for positions very close to the wall (distances less than 1.0 mm (.04")). This was attributed to the differential deflections of the wing and the probe under aerodynamic loads, similar to those reported in Reference 10. Appropriate corrections have been applied to the present data for these effects, based upon the method of Reference 10 and special static load deflection calibrations. The details of these corrections are given in the Appendix.

The velocity profiles at the point of separation did not in general exhibit the vertical slope characteristic of separation profiles. This probably reflects on the capability and . the limitation of the probe. The region near the wall in which the velocity gradient approaches to zero is very small (about $0.05 \%$ c) as indicated by the velocity profiles at one or two chordwise stations downstream from the separation point. (For example $90 \%$ and 1008 c at $10.3^{\circ}, 75 \% \mathrm{c}$ at $14.4^{\circ}$, and $60 \% \mathrm{c}$ at $18.4^{\circ}$.) Neither of the two probes employed could detect the reversed flow, thus very shallow regions of separation are indicated. Surface tufts, on the other hand, very clearly indicate flow reversal at these stations. At the chordwise station of $100 \% \mathrm{c}$ for $10.3^{\circ}$ and $18.4^{\circ}$ angles of attack, satisfactory total pressure measurements could not be made either with the flat tube probe or with the cylindrical tube. This is attributed to very low velocities (about 5 or $10 \%$ of free stream velocity) with violent fluctuations around the zero velocity point in the region less than about $.05 \%$ chord above the surface. The probe interference might be contributing to the violent fluctuation. In these regions the indicated dynamic pressure was negative for probe yaw directions of both $0^{\circ}$ and $180^{\circ}$.

Static pressure profiles: Static pressure profiles obtained from the five tube probe at various chordwise stations on the airfoil are shown in Figures 8, 9 and 10, together with the surface static pressures. A close examination of these profiles
reveals distinctly the region of viscous layer characterized by constant static pressure. The inviscid outer flow region is indicated by static pressure gradually decreasing toward the free stream value $\left(c_{p_{s}}=0\right)$. Thus the general character of the boundary layer and potential flow fields are clearly seen, with the exception of some scatter within a height of about . 05\% chord at 5\% chord station. At this station the presence of the transition strip caused a rise in the local static pressure, and influenced the local profile.

Overall static pressure field: Static pressure contours derived from the pressure distributions obtained at ten chordwise stations on the airfoil and four chordwise stations beyond the trailing edge ( $\mathrm{x}=105.6,110,120 \& 150 \%$ c) are shown in Figures 11, 12 and 13. A close look at the wake regions indicates a distinct high pressure ridge a short distance behind the airfoil. A relatively steep vertical pressure gradient from lower surface region to upper surface is also observed indicating a vertical flow tendency. The velocity plots (Figure 2, 3 and 4) substantiate this tendency for upward turning of the lower surface flow.

Boundary layer characteristics: Figures 14 and 15 show the growth of displacement and momentum thickness on the airfoil, indicating substantial increases between pre-stall and poststall conditions. The shape factor variations at these conditions are shown in Figure 16. All three curves indicate the characteristic behavior of the turbulent separated boundary layer, i.e., rapid growth prior to separation.

For the case of $18.4^{\circ}$ angle of attack, the value of shape factor at the separation location (45\% chord) is l.67, which is a little below the theoretical range of 1.8 to 2.2 ordinarily associated with separation (Reference ll). This discrepancy is believed to be caused by the fluctuations associated with the post-stall turbulent flow fields. This results in some fluctuation of the point of separation. However, the value of the
shape factor of 1.8 is seen to be attained at the $47.5 \%$ chord station. Shape factor values for separation points observed for $10.3^{\circ}$ and $14.4^{\circ}$ angle of attack ( $80 \%$ and $65 \%$ chord) agree very well with the expected values (Table 1).

TABLE 1 - Shape Factors at Separation
$\left.\begin{array}{cccc} & \begin{array}{c}\text { Separation Point } \\ \text { From Tuft } \&\end{array} & \begin{array}{c}\text { Shape }\end{array} & \begin{array}{c}\text { Normally } \\ \text { Expected } \\ \text { Fressure Data }\end{array} \\ \text { Falue of } \mathrm{H}\end{array}\right]$

The boundary layer displacement thickness superimposed on the airfoil is shown in Figure l7. The augmented surface follows the airfoil surface very closely up to the point of separation and then diverges away with varying degree of slope depending on the depth of the separated layer. A similar trend is shown by the separation streamline, which is obtained by equating the mass flow in the reversed flow region to that of the flow in the forward direction. This is plotted for angles of attack of $14.4^{\circ}$ and $18.4^{\circ}$ in Figure 18.

Velocity and pressure distribution in the wake: Velocity profiles, static, and total pressure profiles are shown in Figures 19 through 21 for a $z$ range of about $-8 \%$ to $20 \%$ at each chordwise station. The progressive growth of the width of the wake with increasing longitudinal station is seen as expected. Progressive reduction of static pressure gradients in the wake can also be seen.

Total pressure profiles at the higher angles of attack show irregular shapes near the trailing edge, but become progressively smoother downstream. Contour plots of the total pressure (Figure 22) appear to be basically different between pre- and post-stall conditions. At $18.4^{\circ}$ (post-stall) angle of attack,
two low pressure regions can be seen, whereas the $10.3^{\circ}$ and $14.4^{\circ}$ cases have only one low pressure region.

Reattachment point location in the wake: An examination of the wake velocities, (Figures 19, 20 and 21) indicates a very interesting feature. The regions of reversal in the wake terminate within a relatively short distance downstream from the trailing edge. This occurs in the region of high pressure ridge discussed earlier. A detailed interpolation was carried out in order to establish the point at which the reversed flow terminates. This point, which is characterized by a single zero velocity point in the velocity profile, is referred to as the "reattachment point". Before the determination of the reattachment point, wake vertical centers at various chordwise stations were obtained by picking out the minimum velocity points from the velocity profiles (Figure 23).

Since at the reattachment point the static pressure and the total pressure have the same value, total and static pressure distributions were drawn along the wake center line (Figure 24). Unfortunately these comparisons did not result in accurate reattachment point estimation for the case of $18.4^{\circ}$ angle of attack.

A second scheme for obtaining location of the $18.4^{\circ}$ reattachment point is to plot $u_{x}$ versus $x$ for constant $z$ values. These curves, as shown in Figure 25, permit easy identification of the most downstream point at which reversed flow (negative $u_{x}$ ) is present. A summary of feattachment point locations is given in Table 2.

| TABLE 2 | Reattachment Point Location |
| :---: | :---: |
| Angle of Attack | Reattachment Point |
| $10.3^{\circ}$ | $103 \% \mathrm{c}$ |
| $14.4^{\circ}$ | $106 \% \mathrm{C}$ |
| $18.4^{\circ}$ | $117 \% \mathrm{C}$ |

Integrated boundary layer characteristics of the wake: The boundary layer characteristics of the wake depend on the free shear layer generated at both top and bottom separations from the airfoil surface. The wall shear stress, of course, vanishes in the wake. For flow situations where the wake is symmetric, numerical integration can be carried out for the top or bottom half of the wake alone. In such cases the entrainment of mass from the upper and lower jet sheets would be equal.

When the wake is asymmetric as in the present case, the entrainments of mass by the upper and lower jet sheets are no longer equal. Since mass is transferred between the upper and lower streams, it is not possible to distinguish upper side displacement or momentum thicknesses aft of the trailing edge. Consequently, estimation of the dividing streamline between the wake jet sheets generated by the airfoil upper and lower surfaces is not possible. Also, the recirculation region, which terminates at the reattachment point, is characterized by relatively strong vertical and horizontal pressure gradients (Figs. 23 through 25). Vertical pressure gradients violate the underlying assumptions of Von Karman's momentum integral approach (Ref. ll).

For the reasons just outlined, the integrated boundary layer characteristics of the wake are difficult to assess. In any case, they should be expressed in terms of total width of the wake.

Unfortunately, in the present research, measurement were not obtained to the lower edge of the wake viscous region because of the limitations of the stroke of the probe traversing mechanism. Thus, even though the outer edge of the viscous region of the upper surface wake was surveyed, calculations of longitudinal variations of the displacement thickness, momentum thickness and the separation streamline were not carried out for positions downstream from the trailing edge.

Skin friction computations: Several attempts were made to obtain the skin friction distribution up to the point of separation, employing the Von Karman momentum integral approach (Reference ll). The results were not consistent, expecially near $5 \%$ and $10 \%$ chord stations, where the velocity gradients are steep and momentum thicknesses are small. These discrepancies are attributed to inaccurate estimation of velocity and momentum thickness gradients with relatively few data points. Experimental methods like the Preston tube and razor blade techniques (References 12 \& 13) should be exploited in measuring the local skin friction, especially in the region of rather large pressure gradients. Also, the possibility should be explored of using these techniques under separated flow conditions, where the velocity gradients are low.

## CONCLUSIONS

1. Experimental velocity profiles, flow inclinations and static pressure distributions have been obtained for the GA(W)-l airfoil upper surface and wake at typical pre- and post-stall angle of attack conditions. The data include reversed flow regions of separated boundary layers.
2. Integrations of velocity profiles to obtain displacement thickness, momentum thickness and the separation streamline show that all of these parameters exhibit a rather steep diverging trend from the separation point to the airfoil trailing edge.
3. The present tests reveal that the region of reversed flow associated with separation terminates at a reattachment point which is located a short distance downstream from the airfoil trailing edge.
4. Attempts to determine skin friction distributions by indirect methods did not result in meaningful values.

## RECOMMENDATIONS

1. Boundary layer characteristics on the airfoil lower surface and outer edge of the associated viscous region of the wake should be obtained. These will aid in establishing meaningful mathematical models for separated flows.
2. More direct methods for measurement of local skin friction employing the Preston tube or razor blade should be employed as possible means for obtaining skin friction distributions.
3. A rather detailed scanning of the wake at closer chordwise intervals should be obtained in order to refine the interpolated reattachment point location obtained from the current series of experiments.
4. The experimental techniques utilized in the present research should be applied to an airfoil-flap combination to determine the important features of separated flows with multi-element configurations.

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Figure 1 - Probes

GA(W)-1 AIRFOIL
RN: $2.2 \times 10^{6}$
Mach No.: 0.13


Figure 2 - Experimental Velocity Plot, $\alpha=10.3^{\circ}$

GA(W)-1 AIRFOIL


Figure 2 - Concluded


Figure 3 - Experimental Velocity Plot, $\alpha=14.4^{\circ}$

GA(W)-1 AIRFOIL


Figure 3 - Concluded


Figure 4 - Experimental Velocity Plot, $\alpha=18.4^{\circ}$

GA(W)-1 AIRFOIL


Figure 4 - Concluded


Figure 5 - Near-Wall Velocity Profiles, $\alpha=10.3^{\circ}$

| Symbol | $\underbrace{x /{ }^{(8)}}$ |
| :---: | :---: |
| $\bigcirc$ | 5 |
| $\square$ | 10 |
| $\stackrel{\rightharpoonup}{ }$ | 20 |
| $\triangle$ | 30 |
| $\triangle$ | 40 |
| $\bigcirc$ | 50 |
| $\bigcirc$ | 65 |
| $\odot$ | 75 |
| $\bigcirc$ | 90 |
| $\bigcirc$ | 100 |

$$
\begin{aligned}
& =\text { Cylindrical Probe } \\
& \text { Five Tube Probe }
\end{aligned}
$$

Unprimed: Flat Tube Probe


Figure 6 - Near-Wall Velocity Profiles, $\alpha=14.4^{\circ}$


Figure 7 - Near-Wall Velocity Profiles, $\alpha=18.4^{\circ}$

| Symbo1 | $\underline{x / c}{ }^{(\%}$ |
| :---: | :---: |
| $\odot$ | 5 |
| $\square$ | 10 |
| $\odot$ | 20 |
| $\triangle$ | 30 |
| $\Delta$ | 40 |
| $\bigcirc$ | 50 |
| $\bigcirc$ | 70 |
| $\odot$ | 80 |
| $\bigcirc$ | 90 |
| $\bigcirc$ | 100 |



Figure 8 - Static Pressure Profiles, $\alpha=10.3^{\circ}$

| Symbol | $x_{C}{ }^{(\%)}$ |
| :---: | :---: |
| $\bigcirc$ | 5 |
| $\square$ | 10 |
| $\bigcirc$ | 20 |
| $\triangle$ | 30 |
| $\Delta$ | 40 |
| $\square$ | 50 |
| $\bigcirc$ | 65 |
| $\odot$ | 75 |
| 8 | 90 |
| 0 | 100 |



Figure 9 - Static Pressure Profiles, $\alpha=14.4^{\circ}$


Figure 10 - Static Pressure Profiles, $\alpha=18.4^{\circ}$


Figure 11 - Static Pressure Field Contours, $\alpha=10.3^{\circ}$


Figure 12 - Static Pressure Field Contours, $\alpha=14.4^{\circ}$


Figure 13 - Static Pressure Field Contours, $\alpha=18.4^{\circ}$


Figure 14 - Boundary Layer Displacement Thickness


Figure 15 - Boundary Layer Momentum Thickness


Figure 16 - Boundary Layer Shape Factor
dISPLACEMENT THICKNESS DISTRIBUTION
$\alpha=10.3^{\circ}$


Figure 17 - Augmented Airfoil Profile

# dISPLACEMENT THICKNESS DISTRIBUTION 

$\alpha=14.4^{\circ}$


Figure 17 - Continued

# displacement thickness distribution 

$$
\alpha=18.4^{\circ}
$$



Figure 17 - Concluded
$\alpha=14.4^{\circ}$


Figure 18 - Separation Streamline

$$
\alpha=18.4^{\circ}
$$



Figure 18 - Concluded.

$$
x=105.6 \% \mathrm{c}
$$



Figure 19 - Wake Velocity and Pressure Profiles, $\alpha=10.3^{\circ}$

$$
x=110 \% c
$$



Figure 19 - Continued


Figure 19 - Continued

$$
x=150 \% c
$$



Figure 19 - Concluded

$$
x=105.6 \% \mathrm{c}
$$



Figure 20 - Wake Velocity and Pressure Profiles, $\alpha=14.4^{\circ}$

$$
x=110 \% \mathrm{c}
$$



Figure 20 - Continued.
$x=120 \% c$


Figure 20 - Continued.
$x=150 \% \mathrm{c}$


Figure 20 - Concluded.

$$
x=105.6 \% \mathrm{c}
$$




Figure 21 - Wake Velocity and Pressure Profiles, $\alpha=18.4^{\circ}$

$$
x=110 \% c
$$



Figure 21 - Continued.

$$
x=120 \% c
$$



Figure 21 - Continued.

$$
x=150 \% c
$$



Figure 21 - Concluded.


Figure 22 - Wake Total Pressure Contours, $\alpha=10.3^{\circ}$


Figure 22 - Continued, $\alpha=14.4^{\circ}$


| $\frac{\text { Symbol }}{}$ |  | $\underline{\alpha}$ |
| :---: | :---: | :---: |
|  |  | $10.3^{\circ}$ |
| $\square$ |  | $14.4^{\circ}$ |
| $\Delta$ |  | $18.4^{\circ}$ |



Figure 23 - Wake Center Locations


Figure 24 - Total and Static Pressures Along Wake Centerlines


Figure 25 - Determination of Reattachment Point Locations

## APPENDIX

## Corrections to Velocity Probe Positions <br> for Model Aeroelastic Effects

Pitot probe positions were set with the probe initially touching the surface of the airfoil, before the start of the tunnel run. When the required speed is established in the tunnel, aerodynamic loads cause the wing and balance system to deflect relative to the probe. For upward deflections of the wing, the flexible probe tip deflects such that contact with the wing surface continues for two or three "nominal" probe settings. This effect appears as a discontinuity in slope of the indicated velocity profiles. Examples of this are shown in Figure Al, which show constant or nearly constant velocities very near the surface. This effect and the method of correcting for it have been described in Reference 10 , and are illustrated by the figure. Corrections obtained in this manner have been applied to velocity profiles forward of the 50\% c location.

A second method for determining model deflections was to carry out a static loading calibration, simulating test forces and moments. This loading results in vertical translation, horizontal translation, and balance rotation about the 50\% c point. Since the airfoil was restrained in torsion at only one end, it is also necessary to account for elastic deflections due to wing spar torsion. An additional correction due to mechanical backlash in the pitch mechanism has also been included.

This calibration results in the deflections tabulated in Tables A1, A2 and A3. Maximum upward deflection is about $0.15 \%$ c near the leading edge, and maximum downward deflection is $0.125 \%$ at the trailing edge. For positions forward of mid-chord, the deflections obtained from static loading agree within $.025 \%$ c with those obtained from the method of Reference lo. For positions aft of $50 \% \mathrm{c}$, the airfoil deflects
downward, opening the probe-airfoil gap. Corrections to these profiles have been made based upon the static loading data.

TABLE Al: Airfoil Deflections Under Load, $\alpha=10.3^{\circ}$
Lift $=890$ Newtons (200 lbs)
Pitching Moment: 107 Newton-Meters (950 lb-ins)
Drag: 40 Newtons (9 lbs)


|  |  |  |
| :---: | :---: | :---: |
| $\Delta \mathrm{z}$ | Backlash | $\Delta z$ |
| Torsion | ( $\alpha=.1^{\circ} \mathrm{Max}$ ) | Total |
| $\% \mathrm{c}$ | $\%$ c | \% C |
| 0.0092 | 0.0792 | 0.1113 |
| 0.0079 | 0.0708 | 0.1275 |
| 0.0063 | 0.0525 | 0.1167 |
| 0.0042 | 0.0350 | 0.0917 |
| 0.0021 | 0.0173 | 0.0608 |
| 0. | 0. | 0.0317 |
| -0.0042 | -0.0350 | -0.0163 |
| -0.0063 | -0.0525 | -0.0475 |
| -0.0079 | -0.0708 | -0.0829 |
| -0.0104 | -0.0875 | -0.1192 |

```
TABLE A2: Airfoil Deflections Under Load, \alpha = 14.4*
Lift = l068 Newtons (240 1bs)
Pitching Moment: }129\mathrm{ Newton-Meters (1140 lb-ins)
Drag: 58 Newtons (13 lbs)
```

|  | $\begin{aligned} & x_{1}^{\prime} \\ & 0_{0} \end{aligned}$ | $\begin{gathered} \Delta z \text { Drag } \\ \% \quad c \\ \hline \end{gathered}$ | $\begin{gathered} \Delta z \text { Lift } \\ \% \quad \mathrm{C} \\ \hline \end{gathered}$ | ```\| Pitching Moment % c``` | $\begin{gathered} \Delta z \\ \text { Torsion } \\ \text { \% c } \end{gathered}$ | $\begin{gathered} \Delta z \\ \text { Mechanical } \\ \text { Backlash } \\ \left(\alpha=.1^{\circ} \text { Max }\right) \\ \% \mathrm{c} \\ \hline \end{gathered}$ | - <br> Total <br> \% c |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 5 | -0.1125 | 0.0292 | 0.1067 | 0.0113 | 0.0792 | 0.1137 |
|  | 10 | -0.0667 | 0.0292 | 0.0946 | 0.0104 | 0.0708 | 0.1383 |
| $\stackrel{\sim}{*}$ | 20 | -0.0325 | 0.0292 | 0.0708 | 0.0079 | 0.0525 | 0.1279 |
|  | 30 | -0.0175 | 0.0292 | 0.0475 | 0.0050 | 0.0350 | 0.0992 |
|  | 40 | -0.0041 | 0.0292 | 0.0238 | 0.0025 | 0.0173 | 0.0688 |
|  | 50 | 0.0083 | 0.0292 | 0. | - | - | 0.0375 |
|  | 65 | 0.0367 | 0.0292 | -0.0354 | -0.0038 | -0.0267 | 0.0004 |
|  | 75 | 0.0521 | 0.0292 | -0.0592 | -0.0063 | -0.0433 | -0.0283 |
|  | 90 | 0.0608 | 0.0292 | -0.0946 | -0.0113 | -0.0708 | -0.0833 |
|  | 100 | 0.0608 | 0.0292 | -0.1183 | -0.0125 | -0.0875 | -0.1292 |

```
TABLE A 3: Airfoil Deflections Under Load, \alpha = 18.4*
Lift = l068 Newtons (240 lbs)
Pitching Moment: }129\mathrm{ Newton-Meters (1140 lb-ins)
Drag: lll Newtons (25 lbs)
```

|  | $\begin{aligned} & x_{1} \\ & \%_{c} \end{aligned}$ | $\begin{gathered} \Delta z \text { Drag } \\ \% \quad \mathrm{c} \\ \hline \end{gathered}$ | $\begin{gathered} \Delta z \operatorname{Lift} \\ \% \end{gathered}$ | ```\Deltaz Pitching Moment % C``` | $\begin{gathered} \Delta z \\ \text { Torsion } \\ \% \mathrm{c} \end{gathered}$ | $\begin{gathered} \text { Mechanical } \\ \text { Backlash } \\ \left(\alpha=.1^{\circ} \text { Max }\right) \\ 8 \mathrm{c} \end{gathered}$ | Total \% c |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 5 | -0.1375 | 0.0292 | 0.1067 | 0.0113 | 0.0792 | 0.1000 |
|  | 15 | -0.0550 | 0.0292 | 0.0917 | 0.0092 | 0.0617 | 0.1363 |
| 䍐 | 25 | -0.0396 | 0.0292 | 0.0654 | 0.0063 | 0.0433 | 0.1054 |
|  | 35 | -0.0133 | 0.0292 | 0.0392 | 0.0038 | 0.0267 | 0.0833 |
|  | 45 | 0.0025 | 0.0292 | 0.0129 | 0.0013 | 0.0092 | 0.0546 |
|  | 60 | 0.0325 | 0.0292 | -0.0263 | -0.0025 | -0.0175 | 0.0154 |
|  | 75 | 0.0638 | 0.0292 | -0.0592 | -0.0063 | -0.0433 | -0.0229 |
|  | 82.5 | 0.0696 | 0.0292 | -0.0850 | -0.0079 | -0.0567 | -0.0500 |
|  | 90 | 0.0742 | 0.0292 | -0.0946 | -0.0113 | -0.0708 | -0.0829 |
|  | 100 | 0.0746 | 0.0292 | -0.1183 | -0.0125 | -0.0875 | -0.1258 |



Figure Al - Determination of airfoil deflection from velocity profile data.

