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#### NÁSA TM X-62,458 (NASA-TM-X-62458)CHEMISORPTION KINETICS OF N75-30297 HYDROGEN ON EVAPORATED IRON FILMS (NASA) 13 p HC \$3.25 CSCL 11F Unclas 33054

## CHEMISORPTION KINETICS OF HYDROGEN ON

# EVAPORATED IRON FILMS

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## **CHEMISORPTION KINETICS OF HYDROGEN ON EVAPORATED IRON FILMS\***

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### Abstract

Measurements have been made of the isothermal adsorption-desorption kinetics for H<sub>2</sub> chemisorbed onto Fe films. The chemisorption process is observed to proceed via a precursor state of adsorbed molecular hydrogen similar to the H<sub>2</sub>-Ni system. Reported here are the first measurements of the activation energy for desorption ( $\epsilon_d = 5.3 \times 10^{-20}$  J/molecule) and estimates of the values of the fast kinetic rates between the precursor and chemisorbed states. Adsorption into the precursor state does not appear to be activated, but the process connecting the precursor state with the chemisorbed state will, under certain circumstances, be a rate limiting step for adsorption. The effects of contamination of the surface are evidenced in the measurements.

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From measurements conducted with adsorbate temperatures below  $100^{\circ}$ K, Porter and Tompkins [1] interpreted the chemisorption of H<sub>2</sub> onto Fe surfaces as occurring via an activated process. In studies of the angular distribution of H<sub>2</sub> desorbed from Fe surfaces at  $1000^{\circ}$ K, Bradley and Stickney [2] concluded that activated adsorption is associated with contamination of the surface and that a clean surface shows little indication of activated adsorption. Reported here are the results of measurements of the isothermal adsorption-desorption kinetics for H<sub>2</sub> chemisorbed onto thin (20 to 50 Å) polycrystalline Fe films at temperatures near  $300^{\circ}$ K. This work was undertaken to expand our understanding of the general kinetic processes involved and to clarify the previous work, since the chemisorption kinetics of the H<sub>2</sub>-Fe system are postulated to be relevant to such phenomena as the gaseous embrittlement of steels [3].

The results of this investigation indicate that chemisorption in the H<sub>2</sub> Fe system occurs via a precursor state of molecularly adsorbed H<sub>2</sub> similar to the process previously found for the H<sub>2</sub>-Ni system [4]. Reported here are the first measurements of the activation energy for desorption from the precursor state and estimates of the values of the absolute rate constants connecting the precursor and chemisorbed state at  $\sim 300^{\circ}$ K. Adsorption into the precursor state does not appear to be an activated process. However, the intermediate step, going from the precursor state to the chemisorbed state, was observed to be rate limiting under certain circumstances. This observation suggests that the results of Porter and Tompkins [1] and of Bradley and Stickney [2] are not necessarily in disagreement, since their measurements were conducted under different conditions. The influence of contamination of the surface from unknown impurities in the gas phase is evidenced in this work as affecting (1) the number of available adsorption sites, and (2) modifying the prefactor for the absolute desorption rate constant for the precursor state.

The general experimental technique and procedures have been reported elsewhere [4]; however, for continuity pertinent points will be mentioned briefly while emphasizing specific procedures relevant to this particular work. The kinetic measurements were made by observing the change in the resistance of thin Fe films produced by the chemisorption process. At low coverages, the change in the film's resistance,  $\delta R$ , is directly proportional to the concentration of chemisorbed adatoms [5]. Isotherms, obtained from the resistance changes measured in this study, are consistent with the dissociation of molecular hydrogen into an atomic species in the chemisorbed state. Assuming that the resistance change responds instantaneously with the chemisorption process, than the resistance change reflects directly the temporal behavior of the average surface concentration of chemisorbed adatoms during the various kinetic processes under study [6].

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Iron films were produced by sublimation onto glass substrates from high purity (99.998%) Fe wire; typical film dimensions are  $25 \times 3 \text{ mm}^2$ . Some films were evaporated onto substrates heated to  $200^{\circ}$ C and subsequently annealed at  $250^{\circ}$ C for one hour to stabilize the resistance for measurements at various temperatures below the annealing temperature. Other films were evaporated onto substrates at room temperature ( $21^{\circ}$ C) with no anneal. The resistance of these latter films was sufficiently stable at room temperature that they could be utilized shortly after evaporation for room temperature adsorption studies on fresh surfaces. With regard to the kinetic measurements, no significant difference was noted between films prepared by the two different procedures. Values of film thickness, d, presented here are estimates based on the value of the film's room temperature resistance with corrections for finite size effects, assuming entirely diffuse scattering of the electrons at the film surface,

The base pressure in the vacuum system after bakeout was typically in the range between  $6.7 \times 10^{-8}$  Nm<sup>-2</sup> and  $1.3 \times 10^{-7}$  Nm<sup>-2</sup>. During the period of measurement, the total pressure of residual gases other than hydrogen is estimated to have been below  $1.3 \times 10^{-8}$  Nm<sup>-2</sup>, based on previous experience with the apparatus. The total pressure rose to a typical maximum of  $\sim 5 \times 10^{-7}$  Nm<sup>-2</sup> during the initial part of the evaporation process, dropping significantly during the evaporation as the Fe film deposited on the chamber walls gettered residual gases.

The main experimental data consist of the following: (1) measurements of the isothermal kinetics of adsorption from the gas phase to the chemisorbed state, and (2) measurements of the isothermal kinetics of desorption from the chemisorbed state to the gas phase. The adsorption kinetics were obtained by observing  $\delta R$  as equilibrium was approached following an increase in PH<sub>2</sub>, the hydrogen gas phase pressure, from a low value where adsorption was negligible (~ 7 × 10<sup>-7</sup> Nm<sup>-2</sup>) to a higher value where adsorption was enhanced. A characteristic adsorption curve is presented in fig. 1. Analysis of this type of adsorption data for a variety of hydrogen pressures shows an exponential approach to equilibrium with a characteristic rate constant,  $\tau_a^{-1}$ .  $\tau_a^{-1}$  is observed to be linearly proportional to PH<sub>2</sub> at low pressures and extrapolates to nonzero values at PH<sub>2</sub> = 0 [7]. The value of the zero pressure intercept was found to be in reasonable numerical agreement with the value obtained for the isothermal rate of desorption, when measured under comparable circumstances. The relationship between the rate,  $\tau_a^{-1}$ , and PH<sub>2</sub> deviates from linearity at higher values of PH<sub>2</sub> and shows signs of saturating; an example of this is presented later.

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The initial sticking coefficient,  $s_0$ , for adsorption was estimated for low pressures (linear region) by dividing the measured value of  $\tau_a^{-1}$  by the flux of molecules per adsorption site. In estimating the flux of molecules per adsorption site, the value of  $1.2 \times 10^{15}$  sites/cm<sup>2</sup> was employed [8]. The variation in  $s_0$  measured as a function of film age at 294°K is shown in fig. 2. Attempts to estimate  $s_0$  at higher temperatures for freshly evaporated films were not successful as significant measurement problems arise from the instability of the film's resistance due to thermally activated annealing. A few measurements on older surfaces (films greater than 10 hours old) over the temperature range from 20°C to 60°C did not reveal any significant variation in  $s_0$  with temperature when compared with anticipated experimental fluctuations.

Data on the desorption kinetics were obtained by observing the decay of  $\delta R$  efter rapidly evacuating the gas phase H<sub>2</sub> to a pressure where readsorption was negligible. Analysis of the data is consistent with an exponential decay for  $\delta R$  with a temperature dependent decay rate denoted by  $\tau_d^{-1}$ . Within experimental fluctuations,  $\tau_d^{-1}$ , as well as the functional form of the decay, is independent of both film thickness and the initial surface concentration of chemisorbed adatoms,  $\Theta$ , for  $\Theta \leq 0.4$  ( $\Theta$  was estimated from the resistance isotherm [4]). Variations of  $\tau_d^{-1}$  with temperature are presented in fig. 3 for relatively "new" films (approximately 2 to 6 hours old) with a few measurements on "older" films to indicate the effect of aging.

The general features of the data, namely, 1st-order adsorption-desorption kinetics (exponential behavior) for a dissociative chemisorption process and a small activation energy for desorption compared to the heat of adsorption [9], suggest that chemisorption is occurring via a precursor state as demonstrated for the H<sub>2</sub>-Ni system [4]. In this model, the chemicorption of hydrogen is schematically represented by the following reaction:

$$H_2 + \Sigma^* \stackrel{k_a P_{H_2}}{\rightleftharpoons} (H_2 - \Sigma^*) + 2\Sigma \stackrel{k_1}{\rightleftharpoons} 2(H - \Sigma) + \Sigma$$

Here,  $k_a P_{H_2}$  and  $k_d$  are, respectively, the absolute rate of adsorption and desorption between the gas phase  $H_2$  and the precursor state of molecularly adsorbed hydrogen  $(H_2 - \Sigma^*)$ ;  $k_1$  and  $k_2$ are absolute rates connecting the precursor state with the chemisorbed state of dissociated adsorbed hydrogen  $(H - \Sigma)$ ; and  $\Sigma^*$  and  $\Sigma$  are physically distinct adsorption sites for the precursor and chemisorbed state. The solutions to the coupled differential equations which describe the kinetics of this model have been shown to be [10] exponential in nature (1st order) if the fractional coverage in the chemisorbed state,  $\Theta$ , and the physisorbed state,  $\Theta^*$ , is small compared with unity so that

.

it is possible to linearize the equations by neglecting quadratic and crossterms. Assuming further that there is a rate limiting step between the precursor state and the gas phase, i.e.,  $k_a P_{H_2}$ ,  $k_d \ll k_1$ ,  $k_2$ , then the slow kinetic rate is given by

$$\tau^{-1} = k_a P_{H_2} + k_d \tag{1}$$

where a small systematic correction has been neglected [4]. Since  $P_{H_2} \approx O$  when measuring the desorption kinetics, the measured net desorption rate  $\tau_d^{-1}$  is interpreted as being equal to  $k_d$ , the rate for desorption from the precursor state to the gas phase. The measured adsorption rate  $\tau_a^{-1}$  is interpreted as being equal to  $\tau^{-1}$ . The measured form of  $\tau_a^{-1}$  at low pressures is in agreement with this relationship.

The initial sticking coefficient,  $s_0$ , presented in fig. 2 represents the sticking coefficient associated with the precursor state. The nearly unity value of  $s_0$  for new films may be fortuitous and could be in error by at most a factor of two; the correct value depends critically upon the actual average number of adsorption sites per unit area. The fact that  $s_0$  is so large for new films strongly suggests that adsorption into the precursor state is not activated, since an activation energy of any significance would greatly reduce  $s_0$  (at 300° K an activation energy of only 7 × 10<sup>-21</sup> J/molecule would reduce  $s_0$  from unity to a value of ~ 0.2).

Assuming the variation of  $s_0$  with film age to be due to contamination of the surface by adsorption to a equilibrium coverage of residual impurities, the variation of  $s_0$  with time was fit to an exponential decay as shown by the solid line in fig. 2. From the decay constant the estimated pressure of the impurities (based on  $O_2$ ) is approximately  $2.7 \times 10^{-8}$  Nm<sup>-2</sup>. This value is somewhat higher than anticipated, but is not unreasonable considering the typical residual base pressures observed during the course of these measurements. Since there was no obvious temperature dependence associated with  $s_0$  for older films, we conclude that the reduction of  $s_0$  represents a loss of available chemisorption sites rather than the development of an activation barrier.

Deviations from linearity in  $\tau_a^{-1}$  as a function of  $P_{H_2}$  would be expected as  $k_a P_{H_2}$  is increased to the point where it is large compared with  $k_1$ ;  $k_1$  would then become the rate limiting step for adsorption. This behavior can be observed in fig. 4 where, for an aged film with  $s_0$ essentially constant, the difference ( $\tau_a^{-1} - k_d$ ) is plotted against  $P_{H_2}$ . At high incident fluxes (large  $k_a P_{H_2}$ ), the precursor state is rapidly occupied, while the chemisorbed state is, comparatively, more slowly occupied. Assuming that the backflow from the chemisorbed state to the precursor can be neglected (the chemisorbed state is initially unoccupied) and that steady state conditions exist with regard to the precursor state, then

$$\tau_a^{-1} \simeq k_1 \left[ \frac{k_a P_{H_2}}{k_1 + k_d + k_a P_{H_2}} \right] + k_d$$

This approximation for the net adsorption rate for the coupled system reduces to the limit expressed by eq (1) when the rate limiting step is between the gas phase and the precursor state. The solid line in fig. 4 is a least means square fit of the above expression to the data where  $k_a$  is required to be consistent with the low pressure data for  $\tau_a^{-1}$ . The value for  $k_1$  which gives the best fit is  $k_1 \approx 0.15 \text{ sec}^{-1}$  at 294°K. From the resistance isotherm, the quantity  $k_1/k_2$  can be evaluated [4] and, in this instance, was found to be approximately 0.034 at 294°K. Based on the value for  $k_1$  and the ratio  $k_1/k_2$ ,  $k_2$  is computed to be  $\sim 4.4 \text{ sec}^{-1}$  at 294°K. Although there are no other estimates for  $k_2$ , the value obtained here is not unreasonable when compared with similar values obtained for the H<sub>2</sub>-Ni system where fluctuation spectroscopy was employed to observe the fast process [11]. The values for  $k_1$  and  $k_2$  should only be considered as representative since the presence of nearby chemisorbed impurities may alter these rates.

The activation energy for desorption from the precursor state,  $\epsilon_d$ , was obtained from the data for  $\tau_d^{-1}$  presented in fig. 3. The low value of  $\epsilon_d = 5.3 \times 10^{-20}$  J/molecule suggests that the precursor state corresponds to a physisorbed state of molecular hydrogen. For comparison, the value of  $\epsilon_d$  for the H<sub>2</sub>-Fe system is only slightly larger that for the H<sub>2</sub>-Ni system ( $\epsilon_d = 4 \times 10^{-20}$  J/molecule), but the prefactor for  $k_d (\approx \tau_d^{-1})$  is larger for the H<sub>2</sub>-Fe system by more than two orders of magnitude [4].

As will be noted in fig. 3, the principal effect of the contaminated surface on  $k_d$  is to decrease the magnitude of the prefactor without significantly modifying the activation energy. The chemisorbed impurities may modify the shape of the potential for the precursor state without significantly modifying the binding energy. This may not be unreasonable, assuming the precursor state to be that of an essentially physisorbed molecule, where the dominant force giving rise to the binding energy is an electrostatic interaction between the adsorbed molecule and the adsorbate.

In conclusion, we find that the chemisorption process for the  $H_2$ -Fe system proceeds via a precursor state of molecularly adsorbed hydrogen; the low activation energy for desorption from this state suggests that it corresponds to a physisorbed state. Adsorption into the precursor state is not activated, whereas the step from the precursor state to the chemisorbed state may be activated and can, under certain circumstances, be the rate limiting adsorption step. Additional work will

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reveal more about the kinetics between the precursor and chemisorbed states; in particular, the determination of the activation energies for  $k_1$  and  $k_2$  should be useful in establishing potential energy curves for the states involved.

## **References and Footnotes**

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- [7] This behavior is similar to that shown in fig. 2 of ref. 4.
- [8] This value is a simple average of the number of sites per unit area for the three principal directions, (100), (110), and (111).
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[11] M. R. Shanabarger, Bull. Am. Phys. Soc. 18 (1973) 1589.

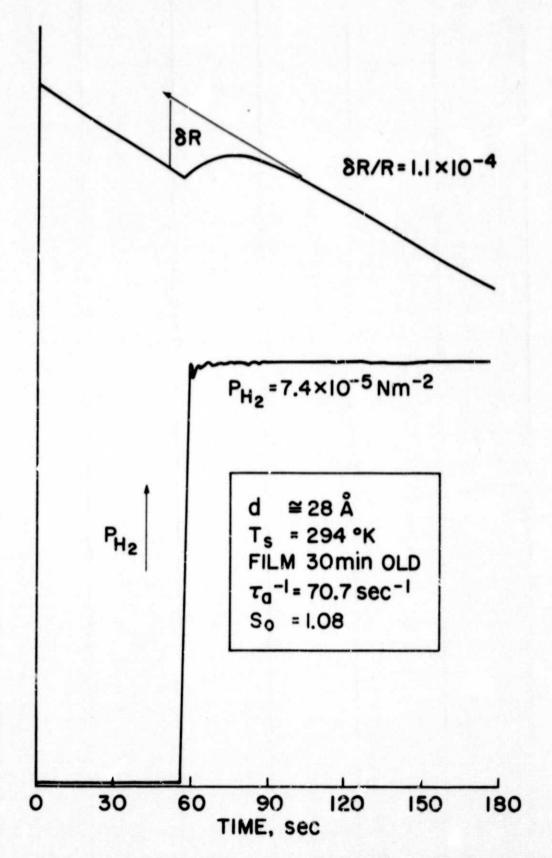


Fig. 1. Tracing of a typical recording of the change in resistance,  $\delta R$ , of a thin Fe film due to the dissociative chemisorption of H<sub>2</sub>. The data were taken on a two-channel recorder with the hydrogen pressure, P<sub>H<sub>2</sub></sub>, presented simultaneously with the resistance change. The baseline drift is due to self-annealing of the film. The film substrate temperature is indicated as T<sub>s</sub> and d is the estimated film thickness.

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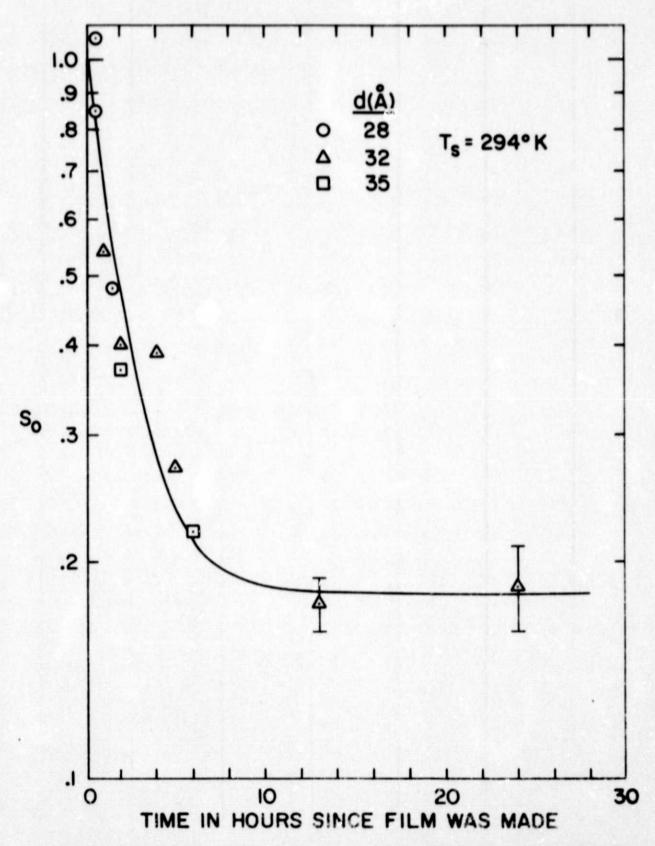


Fig. 2. The variation with film age of the initial sticking coefficient,  $s_0$ , for a series of films measured at room temperature ( $T_s = 294^{\circ}$ K). The values for  $s_0$  were calculated from the measured values of the adsorption rate  $\tau_a^{-1}$  and the gas phase hydrogen pressure, PH<sub>2</sub>. The solid curve represents an empirical fit to the data of the form:  $s_0 = 0.18 + 0.82 e^{-\alpha t}$ , where  $\alpha = 0.533$  hr<sup>-1</sup>.

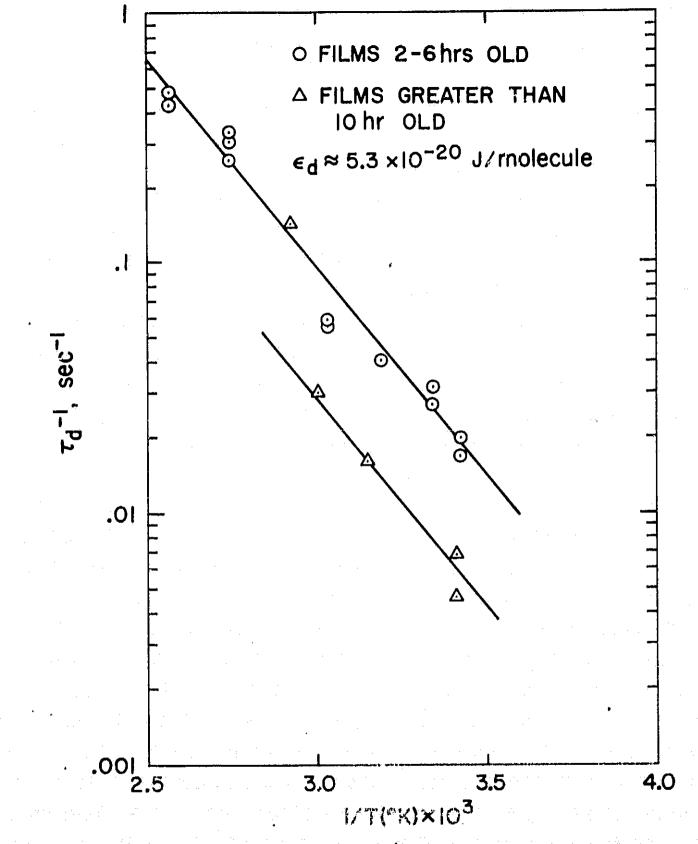


Fig. 3. Arrhenius plot of the desorption rate,  $\tau d^{-1}$ , as a function of reciprocal temperature for relatively new films (uncontaminated surface) and older film (contaminated surface). The activation energy for desorption,  $\epsilon_d = 5.3 \times 10^{-20}$  J/molecule, was obtained from the slope of the indicated straight line. Within small systematic errors,  $\tau d^{-1}$  equals kd, the absolute rate of desorption from the precursor state. A reasonable fit to the data is described by

 $k_d \approx (9.42 \times 10^3) e^{-\epsilon_d/k_BT} (sec^{-1}).$ 

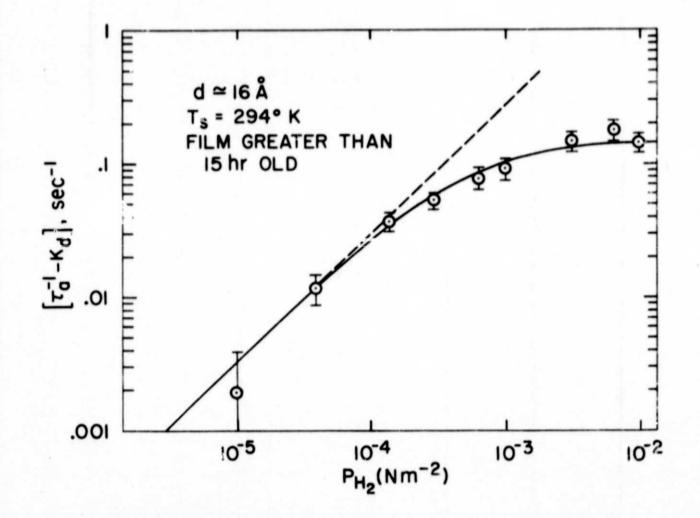


Fig. 4. Plot of the quantity  $[\tau_a^{-1} - k_d]$  as a function of hydrogen gas phase pressure,  $P_{H_2}$ , demonstrating the crossover in the adsorption rate limiting step from between the gas phase and the precursor state to between the precursor and chemisorbed state. The dashed line is based on the predicted low  $P_{H_2}$  behavior, whereas the solid line is the fit of an approximate theoretical result spanning the entire pressure range. For this fit, the only adjustable parameter is the value of  $k_1$ , the rate to go from the physisorbed state to the chemisorbed state; a least means square fit. gives  $k_1 \approx 0.15 \text{ sec}^{-1}$  at  $294^{\circ}$ K.