

## FORMATION FLYING BENEFITS BASED ON VORTEX LATTICE CALCULATIONS

by Brian Maskew

Prepared by
analytical methods, inc. Bellevue, Washington 98004

for Ames Research Center

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## FOREWORD

The research reported here was performed under Purchase Order A-26151-B(DG) between Analytical Methods, Inc., Belleview, Washington, and the National Aeronautics and Space Administration. The NASA project monitor was Remus N. Bretoi, who also prepared Appendix A. The AMI principal investigator was Brian Maskew.

A quadrilateral vortex-lattlce method was applied to a formation of three wangs to calculate force and moment data for use in estimating potential benefits of flying aircraft in formation on extended range missions, and of anticipating the control problems which may exlst. The investigation led to two types of formation having virtually the same overall benefits for the formation as a whole, i.e., a "V" or echelon formation and a double row formation (with two staggered rows of aircraft). These formations have unequal savings on aircraft within the formation, but this allows large longitudinal spacings between aircraft which is preferable to the small spacing required in formations having equal benefits for all aırcraft. A reasonable trade-off between a practical formation size and range benefit seems to lie at about three to five arrcraft with corresponding maximum potential range increases of about $46 \%$ to 67\%. At this time it is not known what fraction of thas potential range increase is achievable in practice.

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# FORMATION FLYING BENEFITS <br> BASED ON VORTEX-LATHICE CALCULATIONS 

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## SUMMARY

A quadrılateral vortex-lattice method was applied to a formation of three wings to calculate the force and moment data for use in estimating potential benefits of flying alrcraft in formation on extended range missions, and of anticipating the control problems which may exist.

The anvestigation led to two types of formation having virtually the same overall benefits for the formation as a whole. Calculations indicate that "V" or echelon formations have an induced drag reduction of $80 \%$ on all but the leading alrcraft. Double-row formations (with two staggered rows of alrcraft) have virtually no induced drag savings in the leading row, but aircraft in the second row have an induced thrust of $48 \%$ of their free-air induced drag. Both formations have large longitudinal distances between aurcraft (about three wing spans), and are, therefore, preferable to formations having equal benefits for all aircraft. The latter formations require vary small spacings as observed in migrating bird formations. The double-row formation has an advantage over the "V" formation in that rolling moment required for trim is zero. Both forma tions are insensitıve to longitudinal spacing, but benefits (and trimming problems) decrease rapıdly with lateral and vertical movements from the optimum locatıons.

A reasonable trade-off between a practical formation size and range beneflt seems to lie at about three to five aurcraft with corresponding maximum potential range increases of about $46 \%$ to $67 \%$. At this time it is not known what fraction of this potential range increase 15 achıevable in practice.

The optimum locations require the trailing vortices from the leading wangs to pass close to the thps of the followers, and this poses a problem for exlsting calculation techniques. The problem affects not only the prediction of the maximum benefits, but also the stability and control aspects of holding formation. Further work, therefore, is recommended to investigate force and moment calculations with close vortex/wing interactions and to refine the calculation method. The results should then be used to analyse the control, stability and range performance of specific aircraft in formation.

[^0]
## INTRODUCTION

Recently there has been renewed interest in the concept of flying aircraft in formation to increase thelr range or to reduce the amount of fuel required for a given range. Modern avionics and control augmentation systems (including alternative autopilot modes) would facilıtate holding accurate formation positions over long periods and would thereby help to achieve the maximum benefits. The benefits arise from an induced drag saving that is caused by favorable aerodynamic interference between the aircraft. So far, formation flying calculations, such as references 1 through 4, have been based mainly on a simple horseshoe vortex representation of each wang; Shollenberger's calculations (ref. 4), however, indicate important differences between the results from that simple approach and from one based on elliptic loading when the spanwise spaclng between wing tips is less than about a semispan. In fact, significant benefits are obtained only with spanwise spacings considerably smaller than a wing semispan, and so a detailed wing representation is clearly required for accurate calculations. This requirement is especially mmportant in practical calculations in which each aircraft is trumed in roll. Formations considered in previous calculations were mannly based on achieving the same induced drag saving for each aircraft in the formation -- a situation that evidently exists in mıgrating bird formations. For aircraft, however, such formations require dangerously small spacings.

The present objective is to apply a vortex-lattice method (ref. 5) to calculate formation flyıng benefits and to consıder a range of formations to find safer alternatives to the "equal benefits" formation. Use of the vortex-lattice method allows a more reallstıc calculation than is possible with the sumple horseshoe-vortex model; the matual interference effects can be represented in greater detall, and each wing can be trimmed in roll. Themaln aims of this study are to find optimum formation configurations to assess the loss in benefits caused by small movements from the optima, and to assess the degree of sensing and control required to maintain range benefits within given limits. The wing planform used in the calculation $1 s$ based on a possible military S.T.O.L. transport aircraft configuration.

| $\mathrm{X}, \mathrm{Y}, \mathrm{z}$ | Cartesıan coordinate system (see fig. 1) |
| :---: | :---: |
| $C_{L}$ | lift coefficient $=1$ ift $/\left(q_{\infty} s\right)$ |
| $C_{D_{i}}$ | induced drag coefficlent $=$ induced drag/( $q_{\infty} s$ ) |
| $c_{\ell}$ | rolling moment coefficuent $=$ rolling moment/( $\left.\mathrm{q}_{\infty} \mathrm{sb}\right)$ |
| $q_{\infty}$ | free-stream dynamic pressure |
| $\mathrm{V}_{\infty}$ | true air speed of formation, m/s |
| S | wing planform area, $\mathrm{m}^{2}$ |
| $\alpha$ | wing incidence relative to the flight vector, deg. |
| $\xi$ | flaperon deflection, posıtıve downwards, deg. |
| b | wing span, m |

## Subscrapts:

$\left.\begin{array}{ll}\text { A } & \text { inner wing } \\ \text { B } & \text { outer wing }\end{array}\right\} \quad$ See figure I

## Basic Method

The method (ref. 5) used here $3 s$ based on vortex-lattice theory (e.g., ref. 6), but the wing lattice is formed into quadrilateral vortzces (fig. 1) instead of horseshoe vortices. Each quadrilateral has a control polnt at which the boundary condition of zero normal velocity is specified; the normal veloci-
 the trailing vortices and quadrilateral vortices in the lattices (i.e., on all the wings in the formation in this case). This gives a set of simultaneous equations in the unknown quadrilateral-vortex strengths; for example,

$$
\begin{array}{ll}
\frac{1}{4 \pi} \Gamma_{k} A_{j k}+B_{J}=0 ; & \begin{array}{l}
J \\
k
\end{array}=1, \ldots, N  \tag{1}\\
& =1, \ldots, N
\end{array}
$$

where $\quad \Gamma_{k}$ is the $k^{\text {th }}$ quadrilateral-vortex strength,
$A_{j k}$ is the influence coefficient for the normal component of $j k$ is the lnfluence coefficlent for the normal component of point by the $k$ th quadrılateral vortex;
$B_{j}$ is the normal component of free-stream velocity at the $]^{\text {th }}$ control pount; and
$N$ is the total number of quadrilateral vortzces in the system.

Inıtially, the trailing vortices are assumed semı-infinite in the streamwise direction, but the method incorporates an iterative procedure for calculatıng the trajectories of the vortices for a force-free wake. In this procedure, the first part of each tralling vortex is divided into a number of short, stralght segments, and each segment is aligned with the local mean velocity vector. As the wake geometry changes, the trailing vortex contributions to the coefficlents, $A$, in Equation (1) change, and so a new vortex strength solution is calculated at each iteration.

For most of the present calculations, the procedure was terminated at the end of the first pass, i.e., with the straight undeflected wake. One case was attempted with the wake iteration to assess the effect of wake roll-up.

When the quadrilateral vortex strengths are known, the forces and moments are calculated by applying the Kutta Joukowski law to each bound vortex segment in the lattices, viz.:

$$
\begin{equation*}
\mathrm{C}_{\mathrm{E}}=\frac{2 \mathrm{~V} \operatorname{s} \Gamma}{s v_{\infty}^{2}} \tag{2}
\end{equation*}
$$



Figure 1. Vortex-lattice Arrangement and Formation Definition.
> where
> $\mathrm{C}_{\mathrm{f}}$ is the force vector coefficient which is resolved into the lift and drag direction;

> V is the calculated local velocity vector at the bound vortex segment mid-point; and
> s is the vortex segment vector.

The force is assumed to act at the segment mid-point for the purpose of evaluating moments.

Formation Representation

Symmetrical formations are assumed; this halves the total number of unknown quadrilateral vortex strengths for a given problem. The formations consist of three wings and each wing has a lattice of 4 (chordwise) by 18 (spanwise) quadrilaterals, figure 1. The total number of unknowns, therefore, is 108.

Ideally, more quadrilateral elements should be used, e.g., $6 \times 30$ on each wing, but the computer program is limited to 150 unknowns and 30 trailing vortices at thas time. Also, the lattice should be set in from each wing tip by a quarter of a quadrilateral "span" for best results (ref. 5); in the program coding prevented this from being applied directly on the inner tips of the outer wangs, and so tip inset was not used in these calculations. Because the main interest in this study is in ratios of induced drag, etc., (i.e., formation values/free-air values) rather than absolute values, nexther of the above factors will be a serious drawback to the present calculations.

## Trimming Procedure

Each wing is trimmed in pitch so that all the wings in the formation have the required lift coefficient, $C_{L}$. The outer wings are also trimmed in roll using a downward-deflecting "flaperon" (fig. 1). It is assumed that the flap mechanısm could be arranged to allow the flaps on one side only to be deflected to provide the basic asymmetry required in, some of the formations; the alleron would be treated similarly, and would be used for periodic trimming purposes to maıntain formation position.

In the calculations, the wings are set initially at incidence $\alpha$ (i.e., $\alpha_{A}=\alpha_{B}=\alpha$ ) giving the required lift coefficient in free-alr conditions, and the outer wings have zero flaperon deflection, i.e., $\xi_{B}=0$. The initial solution gives luft coefficients $C_{L_{A}}$ and $C_{L_{B}}$ for the inner and outer wings, respectavely, and an out-of-balance rolling moment, $C_{\ell B}$, for the outer wings. Using superscripts $o$ and $n$ for the old and new conditions, respectively, the
trimmed incldence for wing $A$ is approximately:

$$
\begin{equation*}
\alpha_{A}^{n}=\alpha_{A}^{o}+\left(C_{L}-C_{L_{A}}^{0}\right) / \frac{d C_{L}}{d \alpha} \tag{3}
\end{equation*}
$$

where the lift curve slope, $\frac{\mathrm{dC}_{\mathrm{L}}}{\mathrm{d} \mathrm{\alpha}}$, is calculated for a single wing with the same vortex lattice.

For the outer wang, the flaperon $1 s$ first deflected through angle $\xi_{B}$ to tram out the rolling moment:

$$
\begin{equation*}
\xi_{B}^{\mathrm{n}}=-\mathrm{C}_{\ell}^{0} / \frac{\mathrm{dC}_{\ell}}{\mathrm{d} \xi} \tag{4}
\end{equation*}
$$

Because of asymmetry, the deflected flaperon gives a lift increment. The trimmed incıdence for wing $B$, therefore, is approximately:

$$
\begin{equation*}
\alpha_{B}^{n}=\alpha_{B}^{o}+\left(C_{L}-C_{I_{B}}^{o}-\frac{d C_{L}}{d \xi} \cdot \xi_{B}^{n}\right) / \frac{d C_{L}}{d \alpha} \tag{5}
\end{equation*}
$$

The trimming steps are repeated until the $C_{I}$ value and zero rolling moment are achieved.

## BASIC DATA

## Wang Geometry

Figure 1 shows the wing planform ancluding the flaperon. The basic detazls are given in table 1. For one calculation the flaperon span was extended over the full semıspan.

Formations

The coordinate system for defining the formations is shown in figure 1 , and figure 2 shows a summary of the formations considered. Wing $A$, the inner wang, as fixed at the origin, while Wing $B$ (and its image) moves along five scan lines, one streamwise, two spanwise and two vertical. The strearmise scan is in the plane $z=0$, and goes from three spans upstream of the origin to three spans downstream. Except for one case where all three wings are tip to tip, the spanwise position of Wing $B$ in the streamwise scan is at $y / b=0.89$; this places the tip of the following wing near to the centrold of shed vorticrty from the leading wang semıspan.

There are two spanwise scans in the plane $z=0$, one at three spans upstream of the origan and one at three spans downstream. Spanwise positions range from $y / b=0.83$ to 2 in each scan.

There are two vertical scans, one at three spans upstream and the other at three spans downstream of the origan. These scans anclude the common poants in the streamwise and spanwise scans, viz., $y / b=0.89$. The vertical movement from $z=0$ is 0.5 b .

In formations where the trailing vortex-lattice from the leading wing overlaps the lattice on the following wing, (i.e., formations with $y / b<$ I) a numerical problem occurs in the calculations if the lattices do not match (ref. 7); only spanwise positions whlch line up the lattices are therefore considered (see fig. l).

SIngle Wing Aerodynamıc Characteristics

Figure 3 shows the basic lift and induced drag characteristics calculated for a single wing represented by the $4 \times 18$ lattice. Because of the limitatıons mentioned under "Formation Representation", these characteristics should be regarded as approximate in the absolute sense; they are used for trimming the calculations here and for the base in the ratio between formation and freeair conditıons. (However, a calculatıon was made at $\alpha=7^{\circ}$ with a tip inset of a quarter quadrılateral span and a $4 \times 40$ lattice, i.e., a $4 \times 20$ lattice on


# Table 1. Basic Wing Geometry <br> (Based on a Semispan of 1.0) 


the half span. The results (fig. 3) show only a small displacement from the $4 \times 18$ lattice characteristics.)

The lift curve slope to be used in equations (3) and (5) is, therefore:

$$
\begin{equation*}
\frac{d C_{I}}{d \alpha}=0.0851 \text { per degree } \tag{6}
\end{equation*}
$$

The flaperon effectuveness (for equations (4) and (5)) calculated at $\alpha=0$ with the $4 \times 18$ lattice is:

$$
\frac{d C_{\ell}}{d \xi}=0.002982 \text { per degree }
$$

$$
\frac{d C_{L}}{d \xi}=0.01060 \text { per degree }
$$

(The corresponding quantitıes calculated with a $4 \times 20$ lattice were 0.003014 and 0.01065, respectavely. Because of asymmetry, this calculation considered the full wang and the lattice was, therefore, limited by the number of trailing vortaces allowed in the computer program, viz., 30 at this time.)

## Data Points

Calculations performed for each formation configuration (after trimming) evaluate the induced drag for each wing and for the formation, together with the incidence (for the required lift coefficient) and the flaperon angle. The induced drag and the incldence are then divided by the corresponding free-air values (at the same $C_{L}$ ). The single-wing data are:

$$
\text { for } C_{L}=0.6, ~ \begin{aligned}
C_{D_{1}} & =0.01296 \\
\alpha & =7.051^{\circ}
\end{aligned}
$$




Figure 3. Lift and Induced Drag Characteristics for Wing Alone.

$$
\begin{aligned}
& \text { and for } C_{L}=1.2, \\
& \qquad \begin{aligned}
C_{D_{i}} & =0.05189 \\
\alpha & =13.92^{\circ}
\end{aligned}
\end{aligned}
$$

(Both sets of data are for the $4 \times 18$ lattice.)

A few data cases considered at the higher $C_{I}$ values gave essentially the same incidence and induced drag ratios as for the $C_{I}=0.6$ case because of the approxmate linear relationships of $C_{L} \sim \alpha$ and $C_{D_{i}} \sim C_{I}{ }^{2}$ (fig; 3).

The incidence and flaperon angles for each formation are related back to approximate lift and rolling moment increments, respectively (using the derivatives from equations (6) and (7)). These quantities vary approximately linearIy with $C_{L}$ and are presented divided by the $C_{L}$ value of the calculation (i.e., 0.6) .

## RESULTS AND DISCUSSION

Streamwlse Scan

Figure $4(a)$ shows the induced drag ratıo varıations as wing $B$ (and its image) moves from three spans upstream to three spans downstream of wing A. The corresponding varlations in incidence ratio, lift increment, flaperon angle and rolling moment are given in parts (b), (c) (d) and (e) of figure 4, respectuvely.

The position for equal induced drag benefit for all three wings occurs at a very small $x / b$ value, $1 . e ., 0.15$, glving a streamwise distance of approximately 2.5 tip chords between the tip trailing-edge of wing $A$ and the tip leading-edge of wang $B$.

The induced drag ratio for the formation of three wings remanns virtually constant at 0.468 over the complete scan (and therefore complies with Munk's stagger theorem). The individual wing induced drag values change markediy in the region where wings $A$ and $B$ change place as leader, but, by three spans upstream and downstream of the origin, the values have almost reached steady conditions. Wing $B$ achleves an induced drag ratio of 0.21 at $x / b=3.0$; the limiting value $1 s 0.20$ for the larger downstream distances assumng the total for the three wings remains constant and that wing $A$ goes to free-air conditions.

At $x / b=-3$, wing $B$ has an induced drag ratio of 0.94 , but wing $A$ ns now in an upwash field from both wings $B$ and receives a negative induced drag ratio (1.e., a thrust) of -0.48 .

The angle of incidence for a given $C_{J}$ is reduced appreciably in formation, figure $4(b)$. The incidence ratio for wing $B$ is $0.83 a t x / b=3$ while that for wing $A$ with $x / b=-3$ goes down to 0.767 . The corresponding lift increment ratios (before trimming) are 0.127 and 0.235 respectively (fig. 4 (c)).

The flaperon angle required to trim the rolling moment on wing $B$ is $2.7^{\circ}$ when $x / b=3$ (flg. $4(d)$ ), and represents an induced rolling moment coefficlent of $-0.0133 C_{\text {( }}$ (fig. 4(e)). Formations with wings $B$ ahead of wing $A$ have almost negligable rolling moment.

One case was considered at $x / b=3$ with a full semispan flaperon to see if this would reduce the roll trim drag. The larger flaperon, rather than reducing trim drag, gives a small increase (fig. 4(a)). (But the smaller deflection required by the larger flaperon should reduce any profile drag penalty.) Evidently, deflecting the shorter flaperon must have improved the basıc loading distribution towards the ideal elliptical shape, giving a negative trim drag for this wing. Deflecting the larger flaperon essentially restored the basic loading distrıbution and induced drag.

(a) Induced Drag Ratio.

Figure 4. Effect of Streamwise Position: $z / b=0.0$;

$$
\mathrm{y} / \mathrm{b}=0.89 \text { (except } \mathrm{y} / \mathrm{b}=1.0 \text { at } \mathrm{x} / \mathrm{b}=0.0 \text { ). }
$$


(b) Incidence Ratio.


Figure 4. Continued.

(e) Rolling Moment Increment in Formation.

Figure 4. Concluded.

## Spanwise Scans

The induced drag ratio variation with spanwise movement at $x / b=3$ is shown in figure 5(a) for wangs $A$ and $B$ and for the formation of three wings. The minımum value of 0.21 for wing $B$ at the spanwise position of 0.89 (fig. 4 (a)) increases to 0.28 at $y / b=1.0$, then increases rapidly and is 0.49 at $y / b=1.05$. The variations in the incıdence, induced lift, flaperon angle and rolling moment ratios, figure $5(\mathrm{~b})$, (c), (d) and (e), respectively, become very steep as $y / b$ decreases towards 1 ; $i . e .$, as the inner tip of wing $B$ approaches the edge of wing A's trailing lattice. As $\mathrm{y} / \mathrm{b}$ increases from 1.0 to 1.05 , the incadence ratio increases by 0.07 (fig. $5(b)$ ), the induced lift decreases by $0.037 \mathrm{C}_{\mathrm{L}}$ (fig. $5(\mathrm{c})$ ), and the induced roliıng moment decreases by $0.01 \mathrm{C}_{\mathrm{L}}$ (fig. $5(\mathrm{e})$ ). The rapid turnover that occurs near $\mathrm{y} / \mathrm{b}=1$ is a limitation of the lattice approach at this time, and is discussed later under "Effect of Wake RollUp".

This formation $1 s$ roll stable, $1 . e .$, with flaperon fixed in a trimmed positıon, a spanwise movement of wing B would cause changes in induced rolling moment that would tend to return the wing to the trimmed conditon. The results for the spanwase scan with wang $B$ at $x / b=-3$ are shown in figure 6. The induced drag ratio for wing A (fig. 6 (a)) shows a rapıd variation from the thrust value seen in figure $4(a)$; as $y / b$ increases from 1.0 to 1.05 the ratio increases by 0.46 . The kink at $y / b=1.0$ seen in figure $4(a)$ for wing $B$ is present here also for wing $A$.

The incidence and induced lift ratios for wing A (figs. 6 (b) and (c)) also have rapid variation near $\mathrm{y} / \mathrm{b}=1.0$; as wing B moves from $\mathrm{y} / \mathrm{b}=1.0$ to 1.05 , the incldence ratio for wing A increases by 0.07 and the corresponding induced lift decreases by $.068 \mathrm{C}_{\text {I }}$

The rolling moment on wing $B$ is not plotted for this case; the maximum flaperon deflection calculated at $\mathrm{x} / \mathrm{b}=-3$ is about $0.1^{\circ}$.

Vertical Scans

The results for a vertical scan at $\mathrm{x} / \mathrm{b}=3.0, \mathrm{y} / \mathrm{b}=0.89$, are given in figure 7. These show a strong sensltivity to vertical movement. The 0.21 induced drag ratio for wing $B$ (fig. 4(a)) increases to 0.35 for a vertical movement of 0.05 b , see figure 7 (a). The incidence ratio, meanwhile, increases by 0.04 (fig. $7(\mathrm{~b})$ ) and the induced lift ratio, $\Delta \mathrm{C}_{\mathrm{I}} / \mathrm{C}_{\mathrm{I}}$, decreases by 0.025 (fig. 7 (c)).

The formations considered here keep the following wings (i.e., wing B) above the wake of the leader. This is stable as far as lift is concerned, ie.e, a vertical displacement upwards would reduce the induced lift, (fig. $7(c))$, causing the wing to come back down, and vice-versa. With wing B below the wake, a downward displacement would reduce the induced lift and so the wing would continue going down.


(a) Induced Drag Ratio.

Figure 5. Effect of Spanwise Position: $x / b=3$; $z / b=0.0$.

(b) Incidence Ratıo.

(c) Lift Increment in Formation.

Figure 5. Continued.

(e) Rolling Moment Incremont in Formation.

Figure 5. Concluded.

(a) Induced Drag Ratio.

Figure 6. Effect of Spanwisc Position:

$$
x / b=-3 ; z / b=0.0 .
$$



Figure 6. Concluded.


Figure 7. Effect of Vertical Posıtion: $x / b=3 ; y / b=.89$.

(e) Rolling Moment Increment in Formation

Figure 7. Concluded.

The effect of vertical position on flaperon deflection is shown in figure 7 (d). A vertical movement of $0.05 b$ from $z=0$ aauses a reduction in rolling moment from $-.0134 C_{L}$ to $-.0085 \mathrm{C}_{\mathrm{I}}$, (fig. $7(\mathrm{e})$ ). In fact, the induced rolling moment changes sign (goes positive) above a $\mathrm{z} / \mathrm{b}$ value of about 0.2 before returning towards zero. This is because the position of maximum vertical induced velocity on wing $B$ moves outboard as wing $B$ moves vertically; whereas the inner tip of wing $B$ recelves the maximum upwash on the wing when $z / b=0$, the upwash there rapıdly goes toward zero as the tip moves above wing A's trailing vortex sheet.

The results for the vertical scan at $\mathrm{x} / \mathrm{b}=-3$ are shown in figure 8. In this case, the results are shown for negatuve $z / b$ values, i.e., wing $A$ is in the stable luft posıtion above the wakes from wings $B$. Agaln, rapid changes with vertical displacement are indicated; a $0.05 b$ vertical movement from $z=0$ causes an increase of 0.38 in the induced drag ratio on wing A (fig. 8(a)) while the incidence ratio increases by 0.05 (fig. $8(b)$ ) and the induced lift increment decreases by $0.45 \mathrm{C}_{\text {L }}$, (fig. 8(c)). The rolling moment for wing $B$ is not presented for this case since it is constant and very small ( $\xi_{B} \simeq 0.1^{\circ}$ ).

## Effect of Wake Roll-up

In the calculations discussed so far in this Section, the trailing vortex sheet shed by the leading wang (or wangs in the case of negative $\mathrm{x} / \mathrm{b}$ values) is represented by a parallel lattice of semi-infinıte vortices. In real flow, the sheet would be almost completely rolled up at three spans downstream, the vortex cores being at the centrold of shed vorticity for each semispan. In the calculations where the lattices overlap (i.e., $y / b<1$, see fig. 1), the trailing vortices that pass over the wing induce a downwash over the tip region, whereas they would induce an upwash there if they were at the rolled-up vortex position. Neglecting roll-up, therefore, gives reductions in calculated induced lift and induced thrust; these effects are indicated in the calculated spanwise distributions of lift and induced drag in figure 9 (a) and (b), respectuvely. The calculations discussed so far are, therefore, pessimistic for the maximum induced drag saving, but larger rolling moments would be expected to occur with wake roll-up accounted for.

The loss in induced upwash as the lattices overlap causes the kinks observed in the spanwise scan results at $y / b=1.0$ (figs. 5 and 6). Also the turnover in the curves of incldence ratio, unduced lift, flaperon angle and induced rolling moment (figs. 5(b), (c), (d), (e) and 6 (b) and (c)) would occur at a smaller $y / b$ value as the wing tip passed through the vortex core position.

Two attempts were made to try to allow for the effects of roll-up in the formation with $x / b=3.0$ and $y / b=0.89$. The first attempt used the iterative vortex roll-up capability of the present calculation method. Several computer runs were carried out, but the calculations were not successful because of the close interaction between the opposing tip vortices and the wing tip. After two iterations, the-calculated vortex trajectories passed too close to the followang wing control points for rellable force calculations to be made.

(a) Induced Drag Ratio.
$\begin{array}{ll}\text { Figure 8. Effect of Vertical posation: } & x / b=-3 ; \\ & y / b=.89 .\end{array}$


Figure 8. Concluded.

(a) Spanwise Distribution of Lift.

(b) Spanwise Distribution of Induced Drag.

Figure 9. Spanwisc Distrabutions of Lift and Induced Drag: $x / b=3 ; y / b=.89 ; z / b=0$.

The problem is essentially a close vortex/wing interaction, and this poses numerical difficulties for a vortex-lattice method. Recent developments in near-field modeling (refs. 7 and 8 ), would remove these difficulties, but the developments are not yet incorporated in the method.

In a second attempt to allow for wake roll-up, the leading wing and its trailing vortex lattıce were replaced by a single vortex having the total curculation of all the trazling vortices on the leading wing semispan. The vortex was placed $0.1 b$ away from the following wing thp. This was done to try to avoid numerical difficultzes that occur when calculating the induced thrust on a wing in the presence of a vortex using equation (2) and the vortex-lattice method. Even so, an unduly large induced thrust ratio, -0.62 , was calculated.

Further work is clearly required to calculate the effect of roll-up of these formation calculations. Such work should consider near-field modeling techniques such as developed in xeferences 7 and 8.

The present calculations indicate that formation flyıng with all aurcraft having the same induced drag saving is not practical because of the small distances separating the wings. However, comparable savings for the formation as a whole are obtained by having streamwise separation of three or more spans, (fig. 4(a)). This would allow very open formation and should lead to safer operation. With large streamwise spacing, the leading aircraft is essentially in free air, so each aircraft in the formation would take a turn in the lead position to equalize fuel used.

The reason for the improvement in benefit for the following wing as the spacing increases is that the upwash induced by the trailing vortices from the leading wing increases with downstream distance and quickly achleves double the inztial value. (In effect, it is the transition from a semi-infinite vortex to an infinite one, and is demonstrated for a single vortex in figure 10.) Wake roll-up would enhance this effect.

Two basic types of open formation appear feasible. The first type is similar to the "V" vormation used by birds, except larger streamwise spacings would be used. The formation could start in echelon (fig. Il(a)), then the leading aircraft could drop back after a time, and so on, until the echelon was reversed. Each aircraft would then have been in the lead position. All the aircraft would require trimming in roll in this formation.

The second type of formation 15 based on the negatave $\mathrm{x} / \mathrm{b}$ positions, and would have two rows of aircraft. The rows could be swept (fig. li(b)) to increase the separation between the aircraft in each row. Aircraft in the leading row would have virtually no drag saving, but the second row would achreve double the saving. This type of formation has no basic roll-trim problem, and could be used for ferrying small aırcraft. Large, long range aircraft could form the first rows, and each "slot" in between could support several smaller aircraft (e.g., R.P.V.'s) in line astern, (fig. ll(c)). Positions in each "slot" are roll-stable; they are also altitude stable if the following wings are above the level of the trailing vortices.

For the echelon or " V " type of formation, the induced drag ratio for N aircraft would be

$$
\begin{equation*}
\frac{C_{D_{I_{F}}}}{C_{D_{1}}}=0.2+\frac{0.8}{N} \tag{8}
\end{equation*}
$$



Figure 10. Upwash Induced by a Smı-infinite Vortex.


Figure ll. Possible Formations.

Cone leader with no saving and $\mathrm{N}-1$ followers with an induced drag ratio of '0.2.)

For the double row type, the induced drag ratio for $N$ (odd) aircraft would be

$$
\begin{equation*}
\frac{C_{D_{F}}}{C_{D_{i_{S}}}}=0.26+\frac{0.74}{N} \tag{9}
\end{equation*}
$$

( $\frac{(N+1)}{2}$ leaders with no saving and $\frac{(N-1)}{2}$ followers with a thrust ratio of 0.48.)

The induced drag ratio for the double row formation could be reduced by adding an aircraft to each end of the second row. These alrcraft would have the 0.2 induced drag-ratio, and would require the trimming in roll.

## RANGE EXTENSION

Any general conclusions relative to increase in crulsing range which may be expected from reduced induced drag must be approached with caution because factors such as buffet boundaries and drag rise with Mach number are aircraftdependent. Further, because of lack of simulation and flight experience to assess the feasibllity of accurate stationkeeping in formation, it is not known what fraction of the potentıal reduction in induced drag is achıevable under operatıonal conditions. However, automatic statıonkeeping using the aircraft automatlc flight control system and laser or microwave position determination may be necessary, especially for flights of long duration.

In view of this, it should be recognnzed that the estimated increases in crulsing range using the relationshıps derived in Appendix A are potential increases, and must be scaled down by some factor depending on the characteristics or constraints of the aircraft considered.

In estimating potentıal increases in range by flying in formation the following assumptions were made:

* Crulse Mach number is the same in formation flight as for the case of a single aircraft not flying in formation.
* Induced drag can be approximated by the expression
$C_{D_{i}}=\left(C_{D_{i}} / C_{I}{ }^{2}\right) C_{I}{ }^{2}$, where $\left(C_{D_{I}} / C_{I}{ }^{2}\right)$ is a constant.
* Specrfic fuel consumption is constant.
* Altitude $1 s$ allowed to change in oxder to achieve the desired values of
$C_{L} / C_{D}$ or $C_{L}^{\frac{1}{2}} / C_{D}$.
* Mach number effects and buffet boundaries are neglected.
* Zero-lıft drag, $C_{D_{0}}$, is constant and the same for each aircraft.
* Aircraft in formation change position frequently enough that an average induced drag coefficient, $C_{D_{i}}$, can be used for the formation, and each aircraft will consume the same amount of fuel during the mussion.

With these assumptions, four formation flight strategies were considered, and estumates were made for range of the formation relative to a range of a
single alrcraft not flying formation. These four strategies, and the corresponding single-aurcraft strategies used for determining relative ranges, are as Follows:

A Formation at its $\max C_{I} / C_{D}$; single aircraft at its max $C_{I} / C_{D}$.

B Formation at its max $C_{L_{1}}^{\frac{1}{2} / C_{D}}$; single aircraft at its max $C_{L^{\prime}}^{\frac{3}{2}} / C_{D}$.

C Formation flyang flight profile similar to that of a single aircraft flyıng at its max $C_{L} / C_{D}$.

D Formation flying flight profile simalar to that of a single aircraft flying at its $\max C_{L}{ }^{\frac{1}{2} / C_{D}}$.

Equations derived in Appendix A for each of the above cases were used to estimate the relative ranges for each of the above cases, and plotted in the figure below for an echelon formation; double-row formations yıeld very nearly the same results.

In Figure 12, it can be seen that greatest range increases occur when the formation flies flight profules optimized for the formation (cases $A$ and B) rather than when the formation flies flight profiles similar to those optimized for an alrcraft not flying in formation (cases $C$ and D).

The figure shows slgnıficant befit can be derıved even for formations of only two alrcraft (as much as $30 \%$ increase for cases A and B). A point of duminishing return is reached beyond formations of ten aircraft. A reasonable trade-off between a practical formation size and range benefit seems to lie at about three to five aircraft, with corresponding maximum potentıal range increases of about $46 \%$ to $67 \%$ (cases $A$ and $B$ ). Corresponding increases for case D for three to five aircraft formations are about 15\% to 19\%. Even these increases could be significant for strategic augmentation missions. Very large formations, say 15 aircraft, show maximum potential increases of about $100 \%$.

Again, a word of caution: it is not known what fraction of this potential range increase is achievable in practice.


Figure 12. RELATIVE RANGE VS. NUMBER OF AIRCRAFT IN FORMATION Echelon Formation

## Conclusions

1) Formations giving equal induced drag saving on all aircraft require dangerously small spacings, but comparable savings for the formation as a whole are obtained with large streamwise spacing. Leading aircraft in the open formations have vartually no induced drag saving.
2) Two types of open formations are possible. The first type is a "v" or echelon type having one leader. Each of the following aircraft would require tramming in roll. The second type of formation has two rows of aircraft. The leadıng row has virtually no induced drag benefit, but the second row has double the benefit of the echelon type. The second type of formation has no basic out-of-balance rolling moment.
3) Significant induced drag savings are acheived in both formations. Alrcraft in the echelon formation have an induced drag of the order of 20\% of the free-alr value at the same $C_{L}$. The second row of the double row formation has an induced thrust of $48 \%$ of the free-air induced drag at the same $C_{L}$.
4) Both formatıons are insensitıve to longitudinal movements, but benefits (and tramming problems) decrease rapıdiy with lateral or vertacal movements from the optimum position.
5) A reasonable tradeoff between a practical formation size and range benefit seems to lie at about three to five aurcraft with corresponding maxumum potential range increases of about $46 \%$ to $67 \%$ wath a flight profile optimized for the formation (rather than one optimized for aircraft flying alone). At this time it is not known what fraction of this potential range increase is achievable in practice.

## Recommendations

Maximum benefits are obtained when the vortex core from the leader passes close to the tip of the follower. This is a close vortex/wang interaction problem which is beyond the scope of present calculation methods. It is recommended that the numerical problems of calculating close vortex/wing interaction. be investigated, based on recent developments in neax-field techniques (refs. 7 and 8). Such an investigation should lead to a more accurate assessment of the maximum benefits, and of the sensing and control aspects near the optimum locations.

Addıtıonal calculatıons should be done for different lattıce densities. The computer program size lumited the number of vortex quadrilaterals for the present calculations. Extensions of the program to allow more quadrilaterals should be considered to anvestigate the convergence of the results. The method should then be used to assess the range benefits and stabliity and control aspects of a specific aircraft in formation.

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## APPENDIX A

DERIVATION OF CRUISING RANGE EQUATION FOR A JET AIRCRAFT

NOTATION:
$S=$ wing area
$R=$ aircraft cruising range, miles
$V=$ true airspeed, MPH
$C=$ specific fuel consumption, lbs. per hour per pound of thrust
$L=1 i f t$
$D=\mathrm{drag}$
$W=$ aircraft weight
$W_{B}=$ aircraft weight at beginning of cruise
$W_{E}=$ aircraft weight at end of cruise
$\sigma=$ air density at altitude divided by sea level density
$\rho=$ air density
()$_{F}=$ Denotes formation of aircraft
( $)_{\text {FOL }}=$ Denotes following aircraft
()$_{L}=$ Denotes Lead aircraft
()$_{S}=$ Denotes a single aircraft (flying alone)

The purpose of this appendix is to develop some measure of potential range increase which may result from reduction in induced drag due to flying in formation.

It is recognized that the benefits of flying in formation are very much aircraft dependent, and that constraints such as buffet boundaries and drag rise as a function of Mach number together with control problems will probably prevent achieving full potential benefit. However, some generalized computations would still be useful to provide some insight for establishing range performance goals.

## A. 1 FORMATION CRUISE AT ITS MAXIMUM LIFT-DRAG RATIO

Consider first the case of a single aircraft not flying in formation. Assume that it normally cruises, for maximum range, at a Mach number below the drag rise, perhaps Mach 0.6 to 0.7 , and at or near its maximum lift-drag ratio. With these assumptions the cruise altitude will be a function of wing loading (or gross weight, since wing area is fixed).

The classical range equation for jet aircraft is given by the following equation: $R=-\int_{W_{B}}^{W_{E}} \frac{V}{C} \frac{C^{\prime}}{C_{D}} \frac{d W}{W}$

Assume $V, C$, and $C_{L} / C_{D}$ are constant, then
$R=\frac{V}{C} \frac{C_{L}}{C_{D}} \log \left(W_{B} / W_{E}\right)$

Now, define relative range as the maximum range of a formation of aircraft divided by the maximum range of a single aircraft. Assuming changes
in cruise Mach number, specific fuel comsumption and zero-lift drag are negligible, the range of a formation of aircraft at its optimum cruise altitude relative to that of a single aircraft at its optimum cruise altitude becomes:

$$
\begin{equation*}
\frac{\mathrm{R}_{\mathrm{F}}}{\mathrm{R}_{\mathrm{S}}}=\left(\mathrm{C}_{\mathrm{L}} / \mathrm{C}_{\mathrm{D}}\right)_{\mathrm{F}} /\left(\mathrm{C}_{\mathrm{L}} / \mathrm{C}_{\mathrm{D}}\right)_{\mathrm{S}} \tag{A-3}
\end{equation*}
$$

The task now is to determine $C_{L} / C_{D}$ for the single aircraft and for the formation. Assuming that the equation for drag coefficient takes the form

$$
\begin{align*}
C_{D} & =C_{D_{O}}+C_{D_{i}} \\
& =C_{D_{0}}+\left(C_{D_{i}} / C_{L}^{2}\right) C_{L}^{2} \tag{A-4}
\end{align*}
$$

it can be shown that $C_{L} / C_{D}$ is maximum when $C_{D_{i}}=C_{D_{O}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}$

The drag coefficient for this condition is then written $C_{D}=2 C_{D_{O}}=2\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}$

And the corresponding lift coefficient is
$C_{L}=\frac{C_{D_{0}}^{1 / 2}}{\left(C_{D_{i}} / C_{L}{ }^{2}\right)^{\frac{1}{2}}}$

## Subsituting

$$
\begin{equation*}
\left(C_{L} / C_{D}\right)_{\text {max }}=\frac{1}{2}\left(C_{D_{0}}\left(C_{D_{i}} / C_{L}^{2}\right)\right)^{-\frac{1}{2}} \tag{A-5}
\end{equation*}
$$

Using the subscripts $S$ and $F$ to distinguish the lift-drag ratio for a single aircraft not flying in formation, and for the average of a number of alrcraft in formation, respectively, the relative range equation becomes, after substitution for $\left(C_{L} / C_{D}\right)_{\max }$,

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{\left(C_{D_{i}} / C_{L}^{2}\right)_{S}^{1 / 2}}{\left(C_{D_{i}} / C_{L}^{2}\right)_{F}^{3 / 2}} \tag{A-6}
\end{equation*}
$$

Assuming that the lead aircraft in the formation derive no change in Induced drag from the follower aircraft, i.e., their aerodynamic characteristics are basically the same as if they were flying as single aircraft, and assuming the formation may contain $N_{L}$ lead aircraft, $N_{\text {FOL }}$ follower aircraft, and $N$ total aircraft (lead plus follower aircraft), the average induced drag coefficient for the formation is given by,

$$
\begin{equation*}
\mathrm{C}_{\mathrm{D}_{1_{F}}}=\frac{\mathrm{N}_{\mathrm{L}} \mathrm{C}_{\mathrm{D}_{i_{L}}}+\mathrm{N}_{\mathrm{FOL}} \mathrm{C}_{\mathrm{D}_{\mathrm{i}_{\mathrm{FOL}}}}}{\mathrm{~N}} \tag{A-7}
\end{equation*}
$$

Since the entire formation will fly at the same lift coefficient, and since the drag coefficient of the lead aircraft will be the same as for a single aircraft not flying in formation,

$$
\begin{equation*}
\left(C_{D_{i}} / C_{L}^{2}\right)_{F}=\frac{N_{L}\left(C_{D_{i}} / C_{L}^{2}\right)_{S}+N_{F O L}\left(C_{D_{i}} / C_{L}^{2}\right)_{F O L}}{N} \tag{A-8}
\end{equation*}
$$

Substituting into Equation $A-6$,

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\left[\frac{N_{L}+N_{F O L}\left(C_{D_{i}} / c_{L}^{2}\right)_{F O L} /\left(C_{D_{i}} / C_{L}^{2}\right)_{S}}{N}\right]^{-\frac{3 / 2}{2}} \tag{A-9}
\end{equation*}
$$

A. 2 FORMATION CRUISE AT ITS MAXIMUM ${ }^{C_{L}}{ }^{\frac{1}{2} / C_{D}}$

Recall the range equation (Equation $A-1$ )
$R=-\int \frac{V}{C} \frac{C_{L}}{C_{D}} \frac{d w}{\mathrm{~V}}$

In equilibrium flight
$V\left[\frac{2 W}{\rho_{0} \sigma S C_{L}}\right]^{\frac{1}{2}}$

Assume specific fuel consumption, $C$, is constant. Substituting for $V$ in the range equation

$$
\begin{equation*}
R=-\frac{1}{C} \int\left(\frac{2 W}{\rho_{O} S \sigma}\right)^{1 / 2} \frac{C_{L} L^{\frac{1}{2}}}{C_{D}} \frac{d W}{W} \tag{A-11}
\end{equation*}
$$

If drag can be approximated by the expression $C_{D}=C_{D_{O}}+\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}$, the ratio $\left(C_{L}^{\frac{1}{2}} / C_{D}\right)$ is maximum when

$$
\begin{equation*}
C_{L}=\left[\frac{C_{D_{0}}}{3\left(C_{D_{i}}\right) / C_{L}^{2}}\right]^{\frac{1}{2}} \tag{A-12}
\end{equation*}
$$

The corresponding drag coefficient can then be written
$C_{D}=(4 / 3) C_{D_{0}}$
and, combining equations $A-12$ and $A-13$,

$$
\begin{equation*}
\left(C_{L}^{\frac{1}{2}} / C_{D}\right)_{\max }=(1 / 4)\left(3 / C_{D_{0}}\right)^{3 / 4}-\left(C_{D_{i}} / C_{L}^{2}\right)^{-1 / 4} \tag{A-14}
\end{equation*}
$$

Assume that as fuel burns the aircraft is allowed to climb at constant angle of attack, i.e., $C_{L}$ and $\left(C_{L}{ }^{\frac{1}{2}}\right) / C_{D}$ remain constant, and that cruise Mach number and velocity remain constant. This being the case $W / \sigma$ is constant and

$$
\begin{equation*}
R=\frac{1}{C}\left(\frac{S}{\rho_{0} S} \frac{W_{B}}{\sigma_{B}}\right)\left(C_{L}^{1 / 2} / C_{D}\right) \log \left(W_{B} / W_{E}\right) \tag{A-15}
\end{equation*}
$$

Range of a formation of aircraft rélative to a single aircraft then becomes

$$
\begin{align*}
& \mathrm{R}_{\mathrm{F}} / R_{\mathrm{S}}=\left(\sigma_{B_{S}} / \sigma_{B_{F}}\right)^{\frac{1}{2}}\left(\left(C_{L}^{\frac{1}{2}} / C_{D}\right)_{F} /\left(C_{L}^{1 / 2} / C_{D}\right)_{S}\right)  \tag{A-16}\\
& \sigma_{B_{S}} / \sigma_{B_{F}} \text { is determined from the relationship } \\
& W_{B}=\frac{1}{2} \rho_{O} V^{2} \sigma_{B} C_{L}
\end{align*}
$$

If $V$ is essentially the same in going from single aircraft flight to formation flight

$$
\left(\sigma_{B} C_{L}\right)_{F}=\left(\sigma_{B} C_{L}\right)_{S}
$$

and

$$
\begin{equation*}
\sigma_{B_{S}} / \sigma_{B_{F}}=C_{L_{F}} / C_{L_{S}} \tag{A-18}
\end{equation*}
$$

Recall that for maximum $C_{L} \frac{1 / 2}{2} / C_{D}$
$C_{L}=\left[\frac{C_{D_{0}}}{3\left(C_{D_{i}} / C_{L}{ }^{2}\right)}\right]$

Then

$$
\begin{equation*}
\frac{\sigma_{B_{S}}}{\sigma_{B_{F}}}=\frac{\left(C_{D_{i}} / C_{L}^{2}\right)_{S}^{\frac{1}{2}}}{\left(C_{D_{i}}^{\prime} / C_{L}^{2}\right)_{F}^{\frac{1}{2}}} \tag{A-19}
\end{equation*}
$$

Substituting expressions for $C_{L} \frac{1 / 2}{2} / C_{D}$ and $\sigma_{B} / \sigma_{B_{F}}$ into the relative range (Equation (A-16),

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{\left(C_{D_{i}} / C_{L}^{2}\right)_{S}^{1 / 2}}{\left(C_{D_{i}} / C_{L}^{2}\right)_{F}^{1 / 2}} \tag{A-20}
\end{equation*}
$$

Interestingly, the relative range obtained by maximizing $C_{L}^{\frac{1 / 2}{/ 2}} / C_{D}$ is the same as that obtained by flying maximum $C_{L} / C_{D}$, and the relative range for the formation can again be written $\frac{R_{F}}{R_{S}}=\left[\frac{N_{L}+N_{F O L}\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F O L} /\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{S}}{N}\right]^{-\frac{1 / 2}{2}}$ (A-21)

## A. 3 FORMATION CRUISE PROFILE SIMILAR TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS MAXIMUM LIFT DRAG RATIO

Assume that a single aircraft not flying in formation flies at a maximum, constant lift-drag ratio, and that the formation flies a flight profile similar to that of a single aircraft not flying in formation. This means
that the range of the formation relative to that of an aircraft flying single at its maximum lift-drag ratio can be expressed as follows:

$$
\begin{align*}
\frac{R_{F}}{R_{S}} & =C_{D_{S}} / C_{D_{F}} \\
& =\left(C_{D_{0}}+C_{D_{i_{S}}}\right) /\left(C_{D_{0}}+C_{D_{i_{F}}}\right)  \tag{A-22}\\
\frac{R_{F}}{R_{S}}= & \frac{1+C_{D_{i_{S}}} / C_{D_{0}}}{1+C_{D_{i_{F}}} / C_{D_{0}}}
\end{align*}
$$

For maximum lift-drag ratio (single aircraft), $C_{D_{i S}}=C_{D_{0}}$ Hence,

$$
\begin{equation*}
\mathrm{R}_{\mathrm{F}} / \mathrm{R}_{\mathrm{S}}=\frac{2}{1+\left(\mathrm{C}_{\mathrm{D}_{\mathrm{i}_{F}}} / \mathrm{C}_{\mathrm{D}_{\mathrm{i}_{\mathrm{S}}}}\right)} \tag{A-23}
\end{equation*}
$$

Assuming that $C_{D_{i}}$ can be approximated by $C_{D_{i}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}$ and that, because the flight profile of the formation is assumed to be similar to that of a single aircraft, the lift coefficients are the same, then

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{2}{1+\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F} /\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{S}} \tag{A-24}
\end{equation*}
$$

A. 4 FORMATION CRUISE PROFILE SIMILAR TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS IAXIMUM $\mathrm{C}_{[ }^{\frac{1}{2}} / \mathrm{C}_{\mathrm{D}}$
By analogue to the case of a formation cruise profile similar to that of a single aircraft programmed to fly at a constant, maximum $C_{L} / C_{D}$, relative range för flight at maximum $C_{L}^{\frac{1}{2}} / C_{D}$ can also be written

$$
\begin{equation*}
\frac{\mathrm{R}_{F}}{\mathrm{R}_{\mathrm{S}}}=\frac{1+{ }^{\mathrm{C}_{\mathrm{D}_{\mathrm{i}_{S}}} / \mathrm{C}_{\mathrm{D}_{0}}}}{1+\mathrm{C}_{\mathrm{D}_{\mathrm{i}_{F}}} / \mathrm{C}_{\mathrm{D}_{0}}} \tag{A-25}
\end{equation*}
$$

From Equation $A-13$, for maximum $C_{L}^{\frac{1}{2}} / C_{D}$,

$$
c_{D}=4 / 3 C_{D_{0}}
$$

$$
\begin{align*}
\text { Writing } C_{D_{i_{S}}}=1 / 3 C_{D_{O}} \\
C_{D_{0}}=3 C_{D_{i_{s}}} \\
C_{D_{i_{F}}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F} c_{L}^{2}  \tag{b}\\
C_{D_{1_{s}}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{S} C_{L}{ }^{2} \tag{c}
\end{align*}
$$

and substituting into Equation A-24
$\frac{R_{F}}{R_{S}}=\frac{4 / 3}{1+(1 / 3)\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F} /\left(C_{D_{i}} / C_{L}{ }^{2)} S\right.}$

The equations derived in this appendix are sumnarized below for convenience and comparison. Assumptions made in deriving them are:

Specific fuel consumption is constant.

Velocity and Mach number are constant and the same for fornation and single aircraft flight.

Altitude is allowed to change to achieve the desired value of $C_{L}^{3 / 2} / C_{D}$ or $C_{L}{ }^{\frac{1}{2}} / C_{D}$.

Change in sound velocity with altitude is neglected.

Mach number effects and buffet boundaries are neglected.

Induced drag can be approximated by the expression $C_{D_{i}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}$, where $\left(C_{D_{i}} / C_{L}{ }^{2}\right)$ is a constant.

Profile drag, $C_{D_{0}}$, is constant and the same for each aircraft.

Induced drag of the lead aircraft of the formation is not affected by other aircraft in the formation; follower aircraft do experience changes in induced drag from aircraft immediately preceding.

$$
\begin{align*}
& C_{D_{i_{F}}}=\frac{N_{L} C_{D_{i_{I}}}+N_{F O L} C_{D_{i_{F O L}}}}{N}, \\
& \left(C_{D_{i}} / C_{L}^{2}\right)_{F}=\frac{N_{L}\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{L}+N_{F O L}\left(C_{D_{i}} / C_{L}^{2}\right)_{F O L}}{N}  \tag{A-28}\\
& \text { Assuming } C_{D_{i}}=\left(C_{D_{i}} / C_{L}{ }^{2}\right) C_{L}{ }^{2}, \text { and } C_{D_{i_{L}}}=C_{D_{i_{S}}} \text {, then } \\
& \frac{\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F}}{\left(C_{D_{i}} / C_{L}\right)_{S}}=\frac{N_{L}+N_{F O L}\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F O L} /\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{S}}{N} \tag{A-29}
\end{align*}
$$

RAFGE OF A FORMATION CRUISING AT ITS YAXIMUII FORIAATION LIFT-DRAG RATIO RELATIVE TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS IRAXIMUM LIFT-DRAG RATIO

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{\left(C_{D_{i}} / C_{L}^{2}\right)_{S}^{1 / 2}}{\left(C_{D_{i}} / C_{L}^{2}\right)_{F}^{1 / 2}} \tag{A-30}
\end{equation*}
$$

RANGE OF A FORMATION CRUISING AT ITS MAXIMUM $C_{T}{ }^{\frac{1}{2}} / C_{D}$ RELATIVE TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS MAXIMMM $\mathrm{C}_{T} \frac{1}{2} / \mathrm{C}_{D}$

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{\left(C_{D_{i}} / c_{L}{ }^{2}\right)_{S}{ }^{\frac{1 / 2}{2}}}{\left(C_{D_{i}} / C_{L}^{2}\right)_{F}^{1 / 2}} \tag{A-31}
\end{equation*}
$$

RANGE OF A FORMATION FLYING A CRUISE PROFILE SIMULAR TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS MAY $C_{I} / C_{D}$

$$
\begin{equation*}
\frac{R_{F}}{R_{S}}=\frac{2}{1+\left(C_{D_{i}} / C_{L}{ }^{2}\right)_{F} /\left(C_{D_{i}} / C_{L}\right)_{S}} \tag{A-32}
\end{equation*}
$$

RAINGE OF A FORMATION FLYING A FLIGHT PROFILE SIMILAR TO THAT OF A SINGLE AIRCRAFT FLYING AT ITS MAX $C_{T} \frac{1}{2} / C_{D}$

$$
\begin{equation*}
\frac{\mathrm{R}_{\mathrm{F}}}{\mathrm{R}_{\mathrm{S}}}=\frac{4 / 3}{1+(1 / 3)\left(\mathrm{CD}_{\mathrm{i}} / \mathrm{C}_{\mathrm{L}} 2\right)_{\mathrm{F}} /\left(\mathrm{CD}_{\mathrm{i}} / \mathrm{C}_{\mathrm{L}} 2\right)_{\mathrm{S}}} \tag{A-33}
\end{equation*}
$$

Representative values obtained from Section 5 for use in estimating relative range potential are as follows:

|  | FORMATION |  |
| :---: | :---: | :---: |
|  | ECHELON | DOUBLE ROW |
|  | 0.2 | -0.48 |

Substitution of these values into Equations A-27 through A-32 yields the results summarized in Table $A-1$. Some comments on these results are as follows:

- Echelon formations provide slightly greater benefits than double row formations.
- Relative ranges are significantly greater for formations cruising at flight conditions optimized for maximizing range of the formation (cases $A$ and $B$ in Table $A-1$ ) rather then flying flight profiles optimized for maximizing range of a single aircraft.
- Potential cruising range for a formation of five aircraft in estimated to be 67 per cent greater than that of a single aircraft. Constraints, such as buffet boundaries, and possible difficulties in stationkeeping control would probably reduce this potential benefit substantially.

| $N$ | $\mathrm{R}_{\mathrm{F}} / \mathrm{R}_{S}$ |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ECHELON FORMATION |  |  |  | DOUBLE ROW FORMATION |  |  |  |
|  | $\frac{\left(\frac{C_{D_{i}}}{\left.C_{L}{ }^{2}\right)_{F}}\right.}{\left(\frac{C_{D_{j}}}{C_{L}{ }^{2}}\right)}$ | $A \& B$ | C | D | $\frac{\left(\frac{C_{D_{i}}}{C_{L}{ }^{2}}\right)_{F}}{\left(\frac{C_{i}}{C_{L}{ }^{2}}\right)}$ | $A \& B$ | C | D |
| 1 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 | 1.0 |
| 2 | . 6 | 1.29 | 1.25 | 1.11 | . 63 | 1.26 | 1.23 | 1.10 |
| 3 | . 467 | 1.46 | 1.36 | 1.15 | . 507 | 1.4 | 1.33 | 1.14 |
| 5 | . 36 | 1.67 | 1.47 | 1.19 | . 408 | 1.57 | 1.42 | 1.17 |
| 7 | . 314 | 1.78 | 1.52 | 1.21 | . 366 | 1.65 | 1.46 | 1.19 |
| 9 | . 298 | 1.86 | 1.55 | 1.22 | . 342 | 1.71 | 1.49 | 1.20 |
| 15 | . 253 | 1.99 | 1.60 | 1.23 | . 309 | 1.80 | 1.53 | 1.21 |
| 39 | . 221 | 2.13 | 1.64 | 1.25 | . 279 | 1.89 | 1.56 | 1.22 |
| 99 | . 208 | 2.19 | 1.66 | 1.25 | . 267 | 1.94 | 1.58 | 1.22 |
| INFINITE | . 20 | 2.24 | 1.67 | 1.25 | . 26 | 1.96 | 1.59 | 1.23 |

TABLE A-1. RELATIVE RANGE AS A FUNCTION OF NUMBER OF AIRCRAFT IN FORMATION
$A=$ Formation cruising at max formation $C_{L} / C_{D}$
Single airćraft cruising at its max $C_{L} / C_{D}$.
$B=$ Formation cruising at max formation $C_{L}^{\frac{1}{2}} / C_{D}$.
Single aircraft cruising at its $\max C_{L}^{\frac{1}{2} / C_{D}}$.
$C=$ Single aircraft cruising at its $\max C_{L} / C_{D}$.
Formation cruising at single aircraft's max $C_{L} / C_{D}$
$D=$ Single aircraft cruising at its max $C_{L}^{\frac{1}{2}} / C_{D}$.
Formation cruising at single aircraft's $C_{L}^{\frac{1}{2}} / C_{D}$.


[^0]:    Senior Research Scientist, Analytical Methods, Inc., Bellevue, Washington

