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BOEING AIRCRAFT COMPANY

SEATTLE, WASHINGTON

MODEL General

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TITLE

POSSIBILITIZS FOR DRAG REDUCTION

BY BOUNDARY LAYER CONTROL

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TABLE OF CONTENTS

References	1
Summery	3
Introduction	4
Transition and Boundary Layer Control	5
Symbols	8
Analysis of Possibilities	10
Analysis of Experimental Work	15
Research Recommendations	19
Figure Index	
1. Types of Intake Slot	21
2. Drag Possibilities with Boundary Layer Control, Duct Loss = 0	22
3. Drag Possibilities with Boundary Layer Control, Duct Loss = .59	23
4. Drag Possibilities with Boundary Layer Control, Duct Loss = lq	24
Annandiz	25
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Figure 5. Co vs. R./R.	28

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SUMMARY

The existing literature on the mechanism of transition is reviewed. Oscillations in the laminar boundary layer are found to be amplified until the disturbance is so great as to cause a breakdown of the flow. This breakdown is known as transition. An analysis is made of the drag possibilities by means of boundary layer control assuming certain conditions of transition Reynolds Number, inlet loss, number of slots, blower efficiency and duct losses. The results appear to be highly favorable. However, experimental investigations on this matter give conflicting results, showing only small gains, and sometimes losses. An analysis of this data indicates that there is a lower limit to the quantity of air which must be removed at a slot in order to stabilize the laminar flow. The removal of insufficient air permits transition to occur. The removal of excessive amounts of air results in high power costs, with a net drag increase. With the estimated value of flow coefficient and duct losses equal to half the dynamic pressure, drag reductions of 50% may be obtained; with twice this flow coefficient, the drag saving is reduced to 25%. It is recommended, therefore, that further research be instigated in order to determine what amount of air must be removed at a slot in order to (1) stabilize the laminar flow, and (2) recstablish laminar flow; also further studies on slot design should be made. The recent work on boundary layer effects on compressibility shock should also be continued.

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<u>INTRODUCTION</u>

While a considerable amount of work has been done on boundary layer removal, almost all of it has been concerned with improvement of maximum lift. Very little work has been done on reduction of drag. Reference 1 gives some information on drag reduction but this appears to be principally due to the jet effect of the discharged air. References 2 and 3, while primarily concerned with lift effects, also include some drag data. Few specific applications for drag reduction have been found. Reference 4 gives results of tests on a glove fitted to one part of the wing of a B-18 dirplane. The drag and the nature of the flow was investigated for various numbers of slots and volumes of air flow. Leminar flow was mainteined up to about 45% chord, the peak pressure position. The power cost was found to be less than the reduction in drag so that a net gain was obtained.

Reference 5 gives the results of some German tests on drag reduction which also show a not gain. Reference 6 gives some further tests but a not gain was not realized.

The body of this paper is divided into three sections: (1) a review of existing literature on the mechanism of transition and the seams of boundary k yer control for reduction of drag; (2) a theoretical analysis to determine the possibilities of drag reduction by boundary layer control; (3) analysis of existing experimental work on this matter.



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TRAMBITION AND FOUNDARY LAYSH CONTROL

The problem of reduction of dreg has been advrowed to one of meintainint leminer flow as far as possible on the body. This means the delay of transition as long as reasible. The understanding of the mechanism of transition has been advanced by research on the stability of the leminar boundary layer.

Stability of laminer boundery layer.--Early theoretical work by Tollmien and Schlichting (described in reference 7) on the stability of the laminar boundary layer indicated that amplification of oscillations was responsible for transition. Recent experimental studies by Schubauer and Skramstad (reference 7) and Liepmann (references 8 and 9) have confirmed this explanation. The phenomenon may be described briefly as follows.

Uscillations of some nort are always present in the sirstream due to sound, surface vibrations, etc. In the laminar layer these vibrations are initially damped but after the boundery layer has grown to a certain thickness, depending upon the frequency of the disturbance, emplification occurs. The theory indicates that after further thickening of the boundary layer, a stable region is reached in which demping once more occurs. However, unless the disturbance is initially very weak, the emplification causes the oscillation to reach a magnitude so great as to cause a complete breakdown of the flow before the stable region is regained. This breakdown area is known as the transition region.

Important elements in this picture are the effects of pressure gradient and suction on transition. It has long been known experimentally that negative pressure gradient (accelerated flow) is favorable to maintenance of laminar flow, whereas

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positive pressure gradient (decelerated flow) is unfavorable. Recent German theoretical work (reference 10) confirms the effect of pressure gradient and also shows that suction has a marked stabilizing effect upon the laminar layer. Pressure discharge has great destabilizing effect causing transition to occur immediately. Ras (reference 11) has confirmed experimentally the stabilizing effect of suction.

<u>Boundary layer control</u>. -- The actual mechanical features of boundary layer control equipment may be considered as composed of inlets, manifolding, ducting, and discharge slots. The design of the ducting and discharge slots need not be further considered here as their design will be based only secondarily on aerodynamic requirements. Besides, these are well-'mown and nothing new is involved.

Very little work has been done on the inlet slot design for drag control. In reference 12 slots are described which were developed by smoke flow observations. Inasmuch as these tests were necessarily performed at very low velocities the validity of the results is questionable. Heference 13 gives quantitative results obtained at higher air speeds. Both references indicate that the direction of entry is not critical, (though backward slots tend to be inferior), that the edges should be rounded, and that the inner part of the slot should expand slowly as in a diffuser (see Figure 1). Under these circumstances the pressure in the inner chamber will be very nearly the same as the surface static pressure. Poorer designs may incur a considerable pressure loss. It appears, therefore, that for a good inlet slot, no pressure recovery is to be expected.

It is usually necessary to use more than one suction slot along the surface and in order to economize on ducting, several slots are manifolded. On wing surfaces

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difficulty usually arises, because the surface static pressure is not the same at each slot. The chamber pressure must be sufficiently low to assure that there is no outward flow at a low pressure slot with resulting breakdown of laminar flow. Under these circumstances a slot in a region of relatively high pressure may have too great a rate of flow unless throttling is resorted to. In any event, the power cost is increased.

<u>Compressibility effects</u>. — Research has been instituted recently to determine the effect of boundary layer upon the shock wave and the resultant drag effects (references 14 - 16). The preliminary results indicate that the shock wave characteristics are different in the presence of a laminar boundary layer than with a turbulent layer, though which is preferable is not yet known. It has also been found that at high velocities in the region of rapid drag increase boundary layer removal will allow a higher Mach Number to be reached before the drag coefficient reaches the same value. Further work is undoubtedly in progress on both these results.

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	<i>e</i>		
x	position of suction slot		
7	distance above surface	and the second sec	
1	chord length		
s	wing area		
δ	boundary layer thickness		·
R	critical Reynolds number, V	'x/ν	
R1	airfoil Reynolds number, Vl	/v	
¥	flight velocity		
۷.	discharge velocity of induc	ted air	
q, q'	dynamic pressures		
ν	kinematic viscosity	ORIGINAL PAGE OF POOR QUALIT	IS TY
c _D	drag coefficient of flat pl	ate at R	l e
с _{ъм}	wing drag coefficient with	boundary layer control assuming	₩ = ₩
с _{DЭ}	equivalent drag coefficient	for power expenditure of blowe	Γ,
°DT	equivalent total drag, C _{DW}	+ c _{DB}	
ନ୍	volume of air inducted per	unit time	
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volume coefficient, 2/SV

 C_{wk} volume coefficient, $\sqrt{V\delta}$

internal pressure of wing, referred to atmospheric static

 Δp pressure loss in duct

Op pressure coefficient, P/q

K duct loss coefficient, AD/q

η_B blower efficiency

Subscripts

11

C.

F

1, 1, 1,n number of suction slot



AMALYSIS OF POSSIBILITITS

The enalysis of drag reduction possible with boundary layer removal is nacesserily brased upon several assumptions. The essimptions used here are as follows: 1. The law of surface friction is that for a flat plate.

2. The laminar velocity profile is assumed to be parabolic.

Both of these assumptions introduce some error but the results of the analysis should be fairly close to the correct values.

Given certain initial conditions such as turbulance of the stream , vibrations, etc., transition from laminer to tarbulent flow will tend to occur at a certain Reynolds Number. If a suction slot is placed at this point, so that air is removed just as transition is about to occur, then the boundary layer thickmass will be reduced and the flow will continue laminer. When the thickmass has again increased to that corresponding to transition, another suction slot again reduces the boundary layer thickness.







layer in the wake is of width 2δ , the same as at the first slot. The velocity profile is indicated at the right of the figure. The inducted air is discharged elsewhere to the rear at stream velocity. The drag coefficient for laminar flow on a flat plate at Reynolds Number R_x may be designated as C_{D_x} .

The wing drag is the same as that accumulated up to the first slot and the drag coefficient is then given by

$$C_{D_{W}} = C_{D_{X}} - \frac{X}{1} = C_{D_{X}} - \frac{R_{X}}{R_{1}}$$
 (1)

¥, 9

To the above drag figure must be added the drag equivalent of the power expended by the blower. This power cost may be considered in two parts: the power te restore the stream total pressure to the inducted air, and the power to compensate for ducting losses.



In the literature the internal wing pressure is referred to the stream static in terms of dynamic pressure.

$$C_{\rm P} = \frac{P}{q}$$

As stated earlier, experimentally it is found that very little energy is recovered

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upon induction of the air, and therefore, Cp is zero or negative. The inlet loss is given by q-P, or non-dimensionally

Inlet loss =
$$1 - C_P$$

The additional pressure loss in the ducting is a function of duct design, length, etc. and can be treated only generally here. This loss may be given as $\Delta p = kq$. The pressure added by the blower is Δp . The discharged air then has a dynamic pressure q' given by

$$q' = q - (q - P) - \Delta \overline{p} + \Delta p$$
$$= \Delta p + P - \Delta \overline{p}$$

or

If the air is discharged at stream velocity (as is usually assumed) then, with $q^{\dagger} = q$,

$$\Delta p = q - P + \Delta \overline{p}$$

$$\Delta p = 1 - C_{p} + \varepsilon$$

$$\Delta P = 1 - C_{p} + \varepsilon$$

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The blower power required is given as $\frac{Q}{N_B}$ where Q is the volume of air inducted per second and η_B is the blower efficiency. The equivalent drag coefficient is given by

$$C_{D_{B}} = \frac{1}{\eta_{B}} \frac{Q_{AD}}{Q_{SV}} = \frac{1}{\eta_{B}} C_{q} \left(1 - C_{p} + \kappa \right)$$
(3)

The total drag as the result of application of boundary layer control is, therefore, given by

 $C_{p_{T}} = C_{p_{W}} + C_{p_{R}} \tag{4}$

Figure 2 gives the drag possibilities with boundary layer control, assuming no duct losses, $\kappa = 0$. Transition Reynolds Number is assumed to be $4 \ge 10^6$, inlet loss is equal to q, i.e. $C_p = 0$, exhaust velocity = V, $\gamma_B = .75$. The figure gives the results for 1, 2, 3, 4, 5, 10, and co number of slots. Contours of

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constant C_{i} are indicated. Figure 3 gives the same results for k = .5 and figure 4 for k = 1. The method of computation is given in the Appendix. Examination of the figures indicate that for a given Reynolds Number there is a minimum number of slots required for drag reduction. The use of more slots with closer spacing reduces the total drag by reducing the volume requirements. The limiting condition is given by the curve representing an infinite number of slots. It would appear desirable, therefore, to use as many slots as possible so as to reduce the volume requirements and thereby reduce the losses and power requirements.

Note on optimum discharge velocity. -- In the usual 'reatments on boundary layer control it is assumed that the air is discharged at stream volocity. This is done for convenience as it eliminates momentum questions. However, with a blower efficiency of less than unity, this is not the optimum condition.

The momentum condition gives

$$\Delta D_{w} = \rho Q (V - V')$$

where V' is the discharge velocity, or

$$\Delta C_{D_w} = 2C_Q \left(1 - \frac{V}{V} \right).$$

The corrected energy equation gives

$$C_{P_{B}} = \frac{1}{\eta_{B}} \frac{Q_{AP}}{qSV} = \frac{1}{\eta_{B}} C_{Q} \left(\frac{q}{q} - C_{P} + \kappa \right)$$

and the drag increment is







The total drag increment is therefore

()

$$\Delta C_{D_{T}} = \Delta C_{D_{W}} + \Delta C_{D_{B}}$$
$$= C_{Q} \left\{ 2 \left(1 - \frac{V}{V} \right) + \frac{1}{\eta_{B}} \left[\left(\frac{V}{V} \right)^{2} - 1 \right] \right\}$$
(5)

When $V^{\bullet} = V$, equation (5) reduces to zero. The optimum is obtained by setting the derivative equal to zero

$$\frac{d\Delta C_{D_{T}}}{d\frac{V}{V}} = 2C_{Q}\left(\frac{1}{\eta_{B}}\frac{V}{V}-1\right) = 0$$
or
$$\frac{V'}{V} = \eta_{B}$$

The optimum discharge velocity is therefore obtained when the ratio of discharge to stream velocity is equal to the blower efficiency.

With high blower efficiency this effect is small, but with low blower efficiencies it becomes important ospecially with large volume coefficients.

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AMALYSIS OF EXPERIMENTAL WORK

Existing experimental results give conflicting results on the reduction of drag by boundary layer control. Some writers profess to have been successful, whereas others have not been. The previous analysis, as presented in figures 1 to 3, indicates that successful results should be obtained if the value of Cq is kept small. Some further enalysis is therefore indicated.

The published reports seldom give the individual slot pressure requirements and give no deta on the boundary layer thickness. The volume coefficient for the individual slots is, however, reported. This information can be used to apgraise the boundary layer condition at the first slot.

We may define
$$C_{\alpha_{\delta}} \equiv \frac{Q_{1}}{V\delta} = \frac{Q_{1}}{V} = \frac{C_{\alpha_{1}}}{\delta/l}$$

Now if x_1 is the position of the first suction slot; and if it is assumed that the boundary layer has the same characteristics as for a flat plate; and again assuming a parabolic velocity distribution:

Then

$$\frac{\delta}{\lambda_{i}} = \frac{5.50}{\sqrt{R_{\lambda_{i}}}}$$

$$C_{\varphi_{\delta}} = \frac{C_{\varphi_{i}}}{\delta/\lambda_{i}} \frac{\lambda_{i}}{\lambda_{i}}$$

$$= \frac{C_{\varphi_{i}}}{5.50} \cdot \frac{\sqrt{R_{\lambda_{i}}}}{\lambda_{i}/\lambda}$$
(6)

In this manner the volume intake at the first slot of the various experiments may be placed upon a cormon basis, the boundary layer thickness δ .

In reference 4 are described flight tests upon a glove fitted on the B-18 sirrlane. The maximum test Heynolds Number, $P_1 = 30 \times 10^6$; $x_1/1 = .20$; C_{Q_1} max. = .00004. This gives $C_{Q_5} = .089$. Nine suction slots from x = .201 to .60 1

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were fitted but laminar flow could not be maintained beyond the fifth slot, corresponding to the peak pressure point. Vibration difficulties were encountered which produced cracks in the airfoil surfaces. An attempt to utilize more slots by halving the spacing was unsuccessful, no laminar flow being obtained. However, with nine slots, at the optimum condition a total drag reduction, after allowance for blower power, was obtained equal to $\Delta C_{D_T} = -.00023$. The drag results were not obtained by wake measurements, but computed from surface measurements of pressure distribution, transition point, and velocity pattern. The magnitudes of the results are, therefore, questionable.

In reference 5 are described wind tunnel tests on an airfoil. The tests were made on single slots and combinations at several angles of attack. $R_1 = 1.6 \times 10^6$; $x_4/1 = .42$; $C_{Q_4} = .0004$ (slot 4 only) and $C_{Q_4} = .0002$ (slots 4 and 6). These values were at minimum total drag. This gives $C_{Q_5} = .142$ and .071, respectively. For slot 4 only at $C_{Q_4} = .0002$, $\Delta C_{D_T} = -.0006$ and at $C_Q = .0004$, $\Delta C_{D_T} = -.0009$; for slots 4 and 6, $C_{Q_4} = .0002$, $\Delta C_{D_T} = -.0009$. In both cases Cp was approximately -0.4, surface pressure coefficient -.32. Somewhat larger drag savings were accomplished with more slots and combinations.

In reference 6 are presented further wind tunnel tests on a model with 15 slots. The report presents results using only six slots; one presumes the selected six gave the optimum distribution. However, no reduction in total drag was obtained. The author ascribes this to the high Reynolds number and states that the slots were too far apart. $R_1 = 3.2 \times 10^6$; $x_1/1 = .075$; $C_{Q1 opt} = .00055$. This gives $C_{-55} = .655$, a figure much larger than in the two preceding references. Because the slots were not close enough, very large quantities of air were exhausted before the wing drag was reduced. The power cost was so high, because of the large

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quantities of air handled, that the total drag was increased by ΔC_{DT} = .0007.

The tests described in reference 3 were made primarily for lift increase, but measurements were also made at low lift coefficients. Only two slots were used, the first one being tested in two positions $x_1/1 = 0.30$ and 0.45, and the second one at $x_2/1 = 0.75$. At the forward position at low volume coefficients apparently all the air was inducted at the first slot. $R_1 = 1.9 \times 10^6$; $x_1/1 = 0.30$; $C_1 = .0020$; therefore, $C_{15} = .915$. No measurements were made at lift coefficients less than 0.4 at this volume coefficient but the total drag was increased by .0035. For the other position tests were made at two Reynolds Numbers. At $R_1 = 1.9 \times 10^6$; $x_1/1 = 0.45$; $C_{21} = .0004$; $C_{15} = .150$; $\Delta C_{27} = .001$. At $R_1 = 6 \times 10^6$; $x_1/1 = 0.45$; $C_{11} = .00067$; $C_{15} = .443$; $\Delta C_{27} = .003$. The value of C_{25} at $R_1 = 1.9 \times 10^6$ is questionable because at higher volume coefficients, $C_{11} = .00067$, the same as at higher Reynolds Number. If this correction is assumed C_{45} is found to be .250 instead of .150.

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C s	°D _T	n	C _L	Extent of Laminar Flow
.071 .071 .089 .142 .150 (?) .250 (?) .443 .655 .915	0006 0009 00023 0009 +.001 .003 .0007 .0035	1 2 9 1 2 2 6 2	0 0 0 0 0 0	? To peak pressure point, .45 1 ? At least .9 1 ?

The above results may be cummarized as follows:

It is evident that for small values of $C_{\mathcal{Q}_{\mathcal{S}}}$ drag reduction is possible. However,

at too small a value it is not possible to maintain laminar flow against the

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pressure gradient. This is reasonable because adverse pressure gradient is destabilizing and suction is stabilizing. There is some point at which balance is achieved. With the meager amount of data on hand, it may be assumed that this value is, say, $C_{ij} = .150$. Assuming a parabolic velocity distribution for the laminar boundary layer, this value of C_{Q_j} corresponds to $y/\delta = .42$. The value of $C_{ij} = 2/3$ represents complete removal of the boundary layer.

Theoretical and experimental investigations of the distribution across the boundary layer of the amplitude of the oscillations (reference 7) indicate that the peak value is reached at approximately y/s = 0.2. No disturbance exists at y/s = 0.7, at which point a phase shift occurs, the oscillations beyond this point being 180° out of phase with the oscillations close to the surface. At y/s = .42 the oscillation emplitude is approximately half the maximum. If the volume of the inducted air is too small, the large amplitude disturbances are not removed and transition is not delayed appreciably. The removal of this highly disturbed air is therefore necessary to preserve the laminar flow.

If it is assumed that a value of $C_{Q_{\delta}} = .150$ is sufficient for the purpose of stabilizing the flow, then a lower feasible limit is set. This value can be translated into C_Q and C_D values as before and is indicated in figures 2 to 4 by the dashed curve.



RESEARCH RECOMMENDATIONS

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From the preceding discussion, it is now evident that some further research is necessary before drag reduction by boundary layer control can be attempted with any real hope of success. Specific work must be undertaken to provide the required information in the following categories: (1) the method of stabilizing the boundary layer so as to maintain laminar flow, (2) the method of reestablishing laminar flow when turbulence has arisen, and (3) further investigations on the design of intake slots.

Stabilizing the laminar flow. - The problem of stabilizing the laminar flow may be considered in two parts. The first part is the consideration of the amount of air that must be withdrawn in order to maintain stability. The experimental results indicate that there is some minimum quantity, apparently an amount sufficiently large so as to include the region in which the oscillations have leached their maximum amplitude. While this quantity has been tentatively set at $C_{\chi\delta} = .150$ (or about 22% of the boundary layer), the exact value is not known and further specific research is necessary to determine this quantity exactly. The second part of this problem is the spacing of the slots. Can the first slot and each succeeding slot be placed at the position where transition is just about to becur? Or is it necessary to place the slots somewhat ahead of this point, and, if so, how much?

<u>Reestablishment of laminar flow</u>. - Another problem closely allied to the above is the reestablishment of laminar flow after the boundary layer has become turbulent. How much of the turbulent layer must be removed for this purpose? It would seem reasonable to assume that the entire removal of the turbulent

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boundary layer is necessary to establish laminar flow. No definite information exists on this matter, however, and experiment may indicate a smaller or larger amount to be required. Specific research should be instituted to secure the answer to this question.

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<u>Design of intake slots</u>. - Existing information on design of intake slots for drag reduction is very meager. The ideal slot will not produce a disturbance in the external flow and will reduce the inlet loss to a minimum value. There is some question as to the flow conditions in the slot under these circumstances. It has been suggested that it is necessary for transition to occur just within the entry so that the internal flow is turbulent. This again is a matter requiring complete investigation.

The above research should preferably be performed in a wind tunnel which is quite free of turbulence, such as the Bureau of Standards, the NACA, or the experimental tunnel at GALCIT. While some of this information could be obtained here, the question of spacing requires turbulent free air, perhaps flight tests.

<u>Compressibility</u>. - The research now in progress on the influence of the boundary layer on the nature of the shock wave should be continued. The delay of the critical region by means of boundary layer control also requires further investigation. This information is invaluable in order to appraise the possible gains at high Mach Numbers.

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APPENDIX

.C.PUTATIONAL METHOD

If the perebolic velocity distribution is assumed for the boundary layer, the ratio of local to stream velocity is given by

$$\frac{v}{V} = 1 - \left(1 - \frac{v}{\delta}\right)^2 \tag{7}$$

The quantity of sir flowing for unit time in the strip between the surface and 9 distince y above the surface is given by

$$Q = \int_{a}^{y} v \, dy$$

or in coefficient form

$$C_{\omega_{\delta}} \equiv \frac{Q}{\sqrt{\delta}} = \int_{0}^{\frac{1}{4}} \frac{\sqrt{\delta}}{\sqrt{\delta}} \frac{dv}{\delta}$$
$$= \left(\frac{y}{\delta}\right)^{2} \left(1 - \frac{1}{3}\frac{y}{\delta}\right)$$
(B)

The momentum loss per unit time in this space is given by

$$= \rho \sqrt{\delta} \left[\frac{1}{\delta} - \frac{1}{\delta} \left(1 - \frac{1}{\delta} \right)^{\delta} + \frac{1}{\delta} \left(1 - \frac{1}{\delta} \right)^{\delta} \right]$$

$$= \rho \sqrt{\delta} \left[\frac{1}{\delta} - \frac{1}{\delta} \left(1 - \frac{1}{\delta} \right)^{\delta} + \frac{1}{\delta} \left(1 - \frac{1}{\delta} \right)^{\delta} \right]$$

$$(9)$$

The drag is given by

$$D = M \Big|_{0}^{\delta} = \frac{2}{15} \rho V^{2} \delta$$

end the momentum thickness is given by

$$\theta = \frac{M}{\rho V_{L}} = \frac{2}{15} \delta$$

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With boundary layer control a certain amount of sir, C_{Q_5} , is removed at a slot corresponding to a layer of width y. The momentum loss in the remaining boundary layer is given by

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$$(10)$$

$$\int_{y}^{\delta} (V - v) v \, dy$$

$$= \int_{y}^{\delta} \left[\frac{1}{2} (v - v) \frac{1}{2} - \frac{1}{2} \left(\frac{1}{2} - v \right)^{2} \right] = \int_{0}^{\delta} V^{2} dy$$

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Further downstream there exists a new boundary layer with a thickness greater than that normally created in the distance from the slot, because of the presence of the unremoved boundary layer. It is <u>essured</u> that immediately after the slot the unremoved min is rearranged into a parabolic velocity distribution having a new thickness δ' but with the same <u>momentum</u> as that represented by equation (10). That is,

$$M|_{y}^{\delta} = M|_{z}^{\delta} = \frac{2}{13}PV^{2}\delta'$$

then

$$\frac{\delta}{\delta} = \frac{5}{2} \left(1 - \frac{\gamma}{\delta} \right)^3 - \frac{3}{2} \left(1 - \frac{\gamma}{\delta} \right)^6$$
(11)

The boundary layer develops along a flat plate, with the parabolic law as follows:

$$\delta = 5.50 \sqrt{\frac{\nu x}{V}}$$

and the thickness δ' is equivalent to a reduction of local surface Reynolds Number to a value F_{x}' , then

$$\frac{\mathsf{R}_{x}}{\mathsf{R}_{x}} = \frac{\mathsf{x}'}{\mathsf{x}} = \left(\frac{\mathsf{s}'}{\mathsf{s}}\right)^{\mathsf{z}} \tag{12}$$

Using y δ as parameter, $C_{\chi_{\delta}}$ can be computed as a function of R_{χ} / R_{χ} using equations (c), (11), and (12). This function is presented in figure \leq .

Now if r wing is to be operated at a design value F_{l} , the minimum number of slots required is the integer in the ratio P_{l}/R_{x} . If it is decided to use n slots the required value of R'_{x} is given by

 $R'_{x} = R_{x} - \frac{R_{1} - R_{x}}{n}$ ORIGINAL PAGE IS $\frac{R'_{x}}{R_{x}} = 1 - \frac{R_{1}/R_{x} - 1}{n}$ OF POOR QUALITY
(13)

or

and from figure 4 the required value of $C_{Q_{S}}$ may be determined.

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The total number of slots is 2 n, cllowing for both surfaces, and the total volume coefficient is therefore

$$C_{q} = 2n C_{q_{\delta}} \frac{\delta}{1}$$

$$= 2n C_{q_{\delta}} \frac{\delta}{2} \frac{\pi}{1}$$

$$= 11n \frac{C_{q_{\delta}}}{\sqrt{R_{x}}} \frac{R_{x}}{R_{1}}$$
(14)

The equivalent drag coefficient for the power expended by the blower is by equation (3)

$$C_{D_{B}} = \frac{1}{\eta_{B}} C_{q} \left(1 - C_{p} + \kappa \right)$$

which assumes the discharge velocity to be equal to flight velocity. Figures 2. to 4 were constructed by assuming the inlet loss to be equal to q, i.e. $C_P = 0$, and using duct loss coefficients of $K_{\pm} 0$, .5, and 1., with $\eta_{\rm E} = 0.75$.







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