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EXPERIMENTAL AND ANALYTICAL TOOLS FOR EVALUATION OF STIRLING ENGINE ROD SEAL BEHAVIOR

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February 1979

Prepared for
NATIONAL AERONAUTICS AND SPACE ADMINISTRATION
Lewis Research Center
Cleveland, OH 44135
Under Contract DEN3-22

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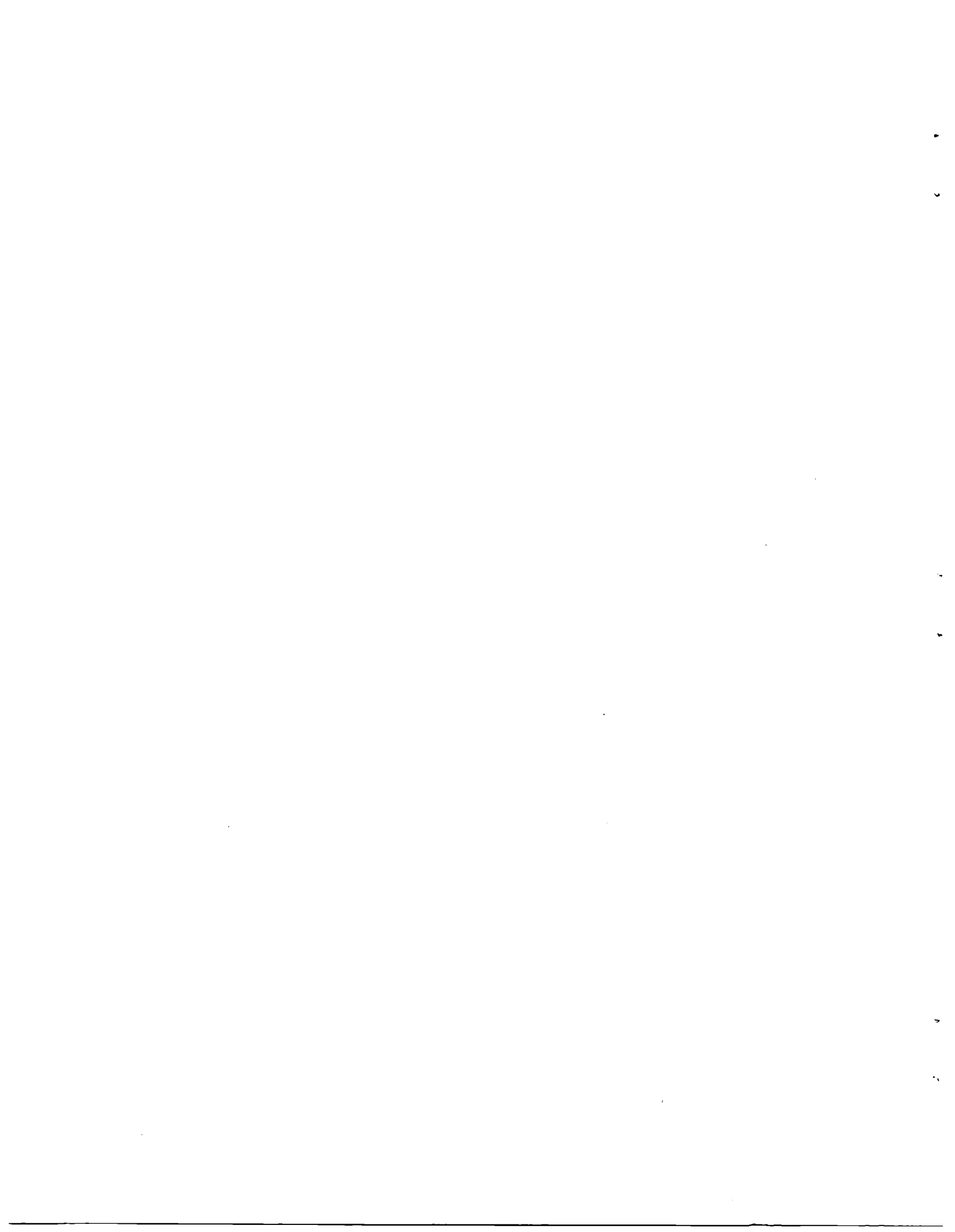
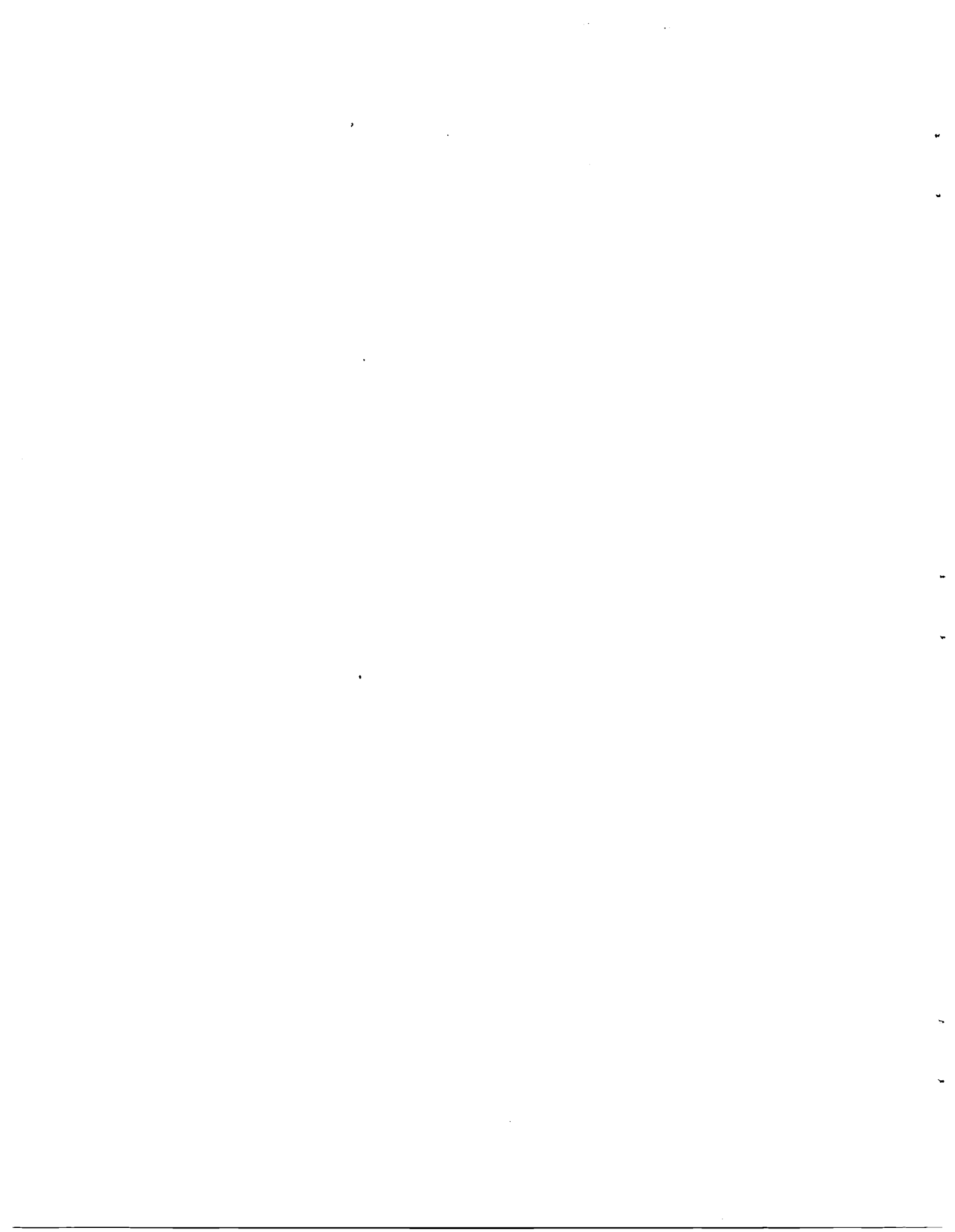


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INTRODUCTION

The Stirling engine is currently the subject of intensive study by the Department of Energy. This engine has the potential to provide clean and efficient propulsive power for the automobile and for other highway vehicles. In order to do so, however, several technical problems associated with such use of the Stirling engine must be solved.

One of the technical problems that must be addressed is the rod seal. This seal has the function of separating the high pressure gas from the low pressure crank case oil. The seal can be either a sliding seal or a roll sock seal. The sliding seal is simpler, less expensive, and possesses a more gradual failure mode than the roll sock seal. However, the sliding seal, unlike the roll sock seal, is not hermetic. Additionally, the sliding seal is subject to wear and can produce undesirable levels of friction force.

The performance of the sliding seal with respect to leakage, wear, and friction can be improved from its present level. To do so requires that the details of the behavior of the seal in the Stirling engine be understood. Additionally, it is necessary to be able to treat the behavior of the seal in a quantitative manner. It is the objective of the work under the current NASA Contract DEN3-22 to progress substantially toward these ends. Specifically, the contract is directed at applying hydrodynamic and elastohydrodynamic theory to the rod seal. The contract is also concerned with the experimental determination of film thickness, fluid leakage, and power loss. The contract entails both producing correlation of experimental and theoretical results and developing tools appropriate for evaluation of rod seal behavior.

The present document is an interim report of work completed in the first year of the two year effort. During the year tools for evaluating rod seal behavior have been developed. Analytically, a computer model of the elastohydrodynamic behavior of the rod seal for reciprocating motion is available. The model permits the large initial radial squeeze, which is typically used for elastomeric seals, to exist. The model determines both the pressure distribution and the oil film thickness distribution in the seal/cylinder contact zone. In this, the calculations are made with time as the independent parameter; consequently, the computer model provides film thickness and pressure results as a function of cyclic position.

Experimentally, a rugged apparatus capable of providing film thickness, friction force, and seal leakage data has been completed. The apparatus contains a moving transparent cylinder (guided by precision hydrostatic bearings) and a stationary elastomeric test seal. A pressure gradient of 690 kPa (100 psi) can be applied across the seal. Frequencies from 10 Hz to 50 Hz with a 0.254 m (1.000 inch) total stroke can be employed. Film thickness is measured with interferometry,

fluid leakage by level and pressure changes, and power (friction) loss by force cells.

This report describes both the experimental and the analytical work completed to date. The experimental effort is discussed in the next section. Included in the section are descriptions of the experimental apparatus, of its capabilities, and of preliminary results from its use. Also included in the section is a discussion of the technique, developed under the contract work, for producing interferometry fringes at the seal/cylinder interface. The special attention given in this report to interferometry is due to the importance of film thickness measurement in the contract work.

The analytical effort is covered both in the section following that for the experimental work and in the Appendices. The section presents the rod seal analysis which is implemented in the computer program given in the Appendixes. (A description of how the program is used is also given in the Appendixes.) In addition, the section presents results obtained from use of the computer program.

The Conclusions and Recommendations section concerns the application during the ensuing year of the analytical and experimental tools developed during the past year and described in this report.

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EXPERIMENTAL REPORT

This section discusses the experimental apparatus that was designed and constructed during the past year. The capabilities of the device and the results obtained from its preliminary use are described. In addition, measurement of film thickness, accomplished by interferometric techniques, is considered. The separate treatment is given to interferometry because of the importance of film thickness measurement to the contract work.

The experimental device has not as yet been used to evaluate the behavior of the rod seal or to produce experimental/analytical correlations. Work in these areas is planned for the ensuing year and is described in the Conclusions and Recommendations section.

The Experimental Apparatus

The experimental apparatus is designed to facilitate the measurement of film thickness, seal friction force, and fluid leakage. The greatest experimental difficulties are associated with determining film thickness. Consequently, the design of the apparatus is based on the technique selected for this measurement.

A review of the available techniques for film thickness measurements was made in order to select that method most appropriate to the elastomeric seal/cylinder wall interface. Some details of that review are presented in the subsection below entitled "Interferometry." The result of the review was the selection of optical interferometry. Accordingly, the experimental apparatus contains a transparent cylinder. In addition, in order to minimize the difficulties of using the interferometric technique, it is this cylinder, rather than the rod seal, that reciprocates. The use of a reciprocating transparent cylinder and a stationary elastomeric rod seal was the most important feature that affected the design of the device.

Other features also affected the design. These features can be grouped in three categories: Measurement capabilities, performance capabilities and characteristics, and ease of use. Each of these categories is considered below.

In the measurement category, it is desirable to be able to determine the instantaneous seal friction force and the instantaneous pressure on the gas side of the seal. It is also desirable that some indication be obtained of the temperature near the seal/cylinder interface. Finally, it is desirable for the apparatus to produce direct macroscopic indications of oil and gas leakage.

In the category of performance capabilities and characteristics, a rather lengthy list of desirable features can be developed. Some of the more significant ones are:

- . capability for a reasonable magnitude of pressure drop across the seal
- . compression (rather than tension) in the transparent cylinder (to utilize its strength characteristics)
- . minimum variation of the gas pressure during each cycle of the cylinder
- . precise guidance of the cylinder and minimum cylinder runout
- . infinitely variable speed in the range from 10 to 50 Hz
- . provision for water cooling
- . accurate alignment of the seal with respect to the transparent cylinder
- . long life span (low stress levels)
- . rigid mounting structure
- . few static seals
- . oil above the test seal (to insure fully flooded conditions)
- . vertical rod axis (to insure axisymmetric seal behavior)
- . balanced dynamic forces, and
- . synchronization of optics with cylinder motion.

In the category of ease of use, significant features are:

- . ease of testing various types of seals
- . ease of replacing the seal and maintaining alignment
- . ease of removing the test seal from the transparent cylinder, and
- . ease of parts replacement.

The experimental apparatus constructed during the past year was designed to incorporate the features in the three categories above. The design is centered around the transparent cylinder which reciprocates with a 25.4 mm (1.000 inch) peak to peak displacement. The transparent cylinder has a 76 mm (3 inch) I. D. This I. D. was chosen large enough to prevent excessive difficulties in producing optical interferometry measurements, and small enough to keep the size of the apparatus (and the associated internal forces) reasonable. In addition, the cylinder bore was chosen such that the seal size is standard and relatively common.

The overall apparatus is shown in the photograph of Figure 1. In the photograph can be seen the stationary plunger which holds the test seal and the oscillating assembly which contains the transparent cylinder. The photograph also shows, at the lower left, a pump and pressure tank. This pump and pressure tank provide high pressure oil to precision hydrostatic bearings which guide the reciprocating parts. The large motor in the center of the photograph is the prime mover which drives the apparatus. Its output speed can be varied continuously from less than 10 Hz to more than 50 Hz. The lattice cage which encloses the rotating and reciprocating parts provides a measure of operator safety. The framework for the apparatus is constructed of thick steel plate and I beams with relatively large cross sections.

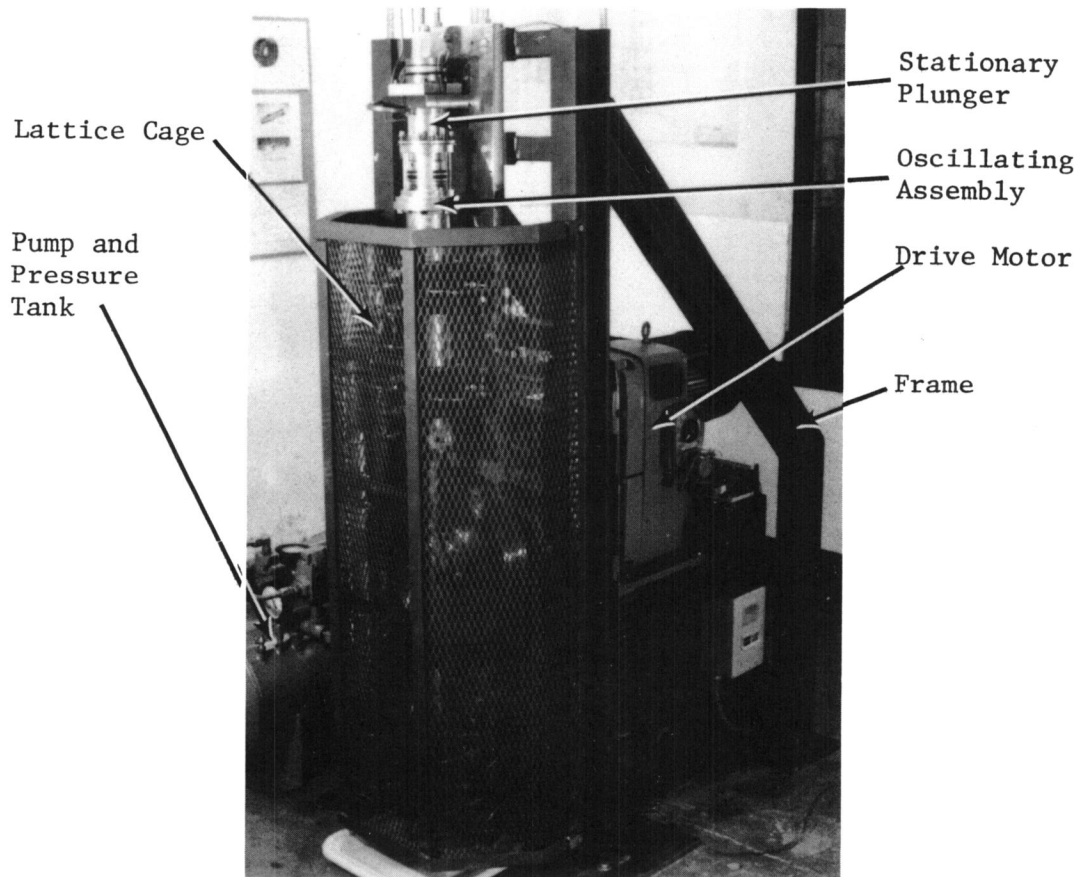


Figure 1 Overall View of
Experimental Apparatus

The frame stands about 1.8 m (6 feet) tall. As pictured in Figure 1, the device has a mass of approximately 500 kg (1,100 pounds mass).

Removal of the guard cage allows components of the drive system to be seen. This has been done in Figure 2, which shows the flywheel, two connecting rods, and 4 hydrostatic bearings with their associated plumbing. These components can also be identified by referring to Figure 3, which is a schematic drawing of the apparatus.

The lower connecting rod is attached to a shaft which moves downward when the upper assembly (which contains the transparent cylinder) moves upward. The sole function of this lower shaft is to balance the vertical forces in the rig. Each shaft is guided by two precision hydrostatic oil bearings which are of a four pad design and contain rubber seals at top and bottom. The shafts are constructed of aluminum. They are attached to their respective connecting rods by needle bearings.

Figure 4 shows some details of the upper end of the apparatus. A schematic drawing which corresponds generally to this photograph is presented in Figure 5. The schematic shows both the exterior outline of this end of the apparatus (on the left in the figure) and some internal details (on the right in the figure).

Referring to both Figures 4 and 5, the hydrostatic bearings and the plunger which holds the test seal are mounted to a backing plate. This backing plate and the one for the lower shaft were assembled to the framework (against the machined tabs, Figure 4) and the bores of all four hydrostatic bearings were machined together. This process insured that the axes of all four hydrostatic bearings became coincident upon assembly of the apparatus. During this machining process, the stationary plate was milled. This procedure resulted in the top surface of the stationary plate being perpendicular to the axis of the four hydrostatic bearings. Doweling of the stationary plate to the backing plate and of the hydrostatic bearings to the backing plate allowed for disassembly of the apparatus subsequent to the machining operation for insertion of the moving shaft.

Several additional components in Figure 4 are worthy of discussion. The vertical pipe at the lower left of the photograph is the manifold for the hydrostatic oil supply. The three small pipes just to the left of center at the top of the photograph are the supply lines to the oil side of the seal, to the gas side of the seal, and for the water cooling. The return water leaves the apparatus through the threaded pipe at the top center of the photograph.

The plunger is connected to the stationary plate with four bolts (three of which can be seen in Figure 4). This arrangement can be readily described by referring to Figures 6 and 7. Figure 6 shows the

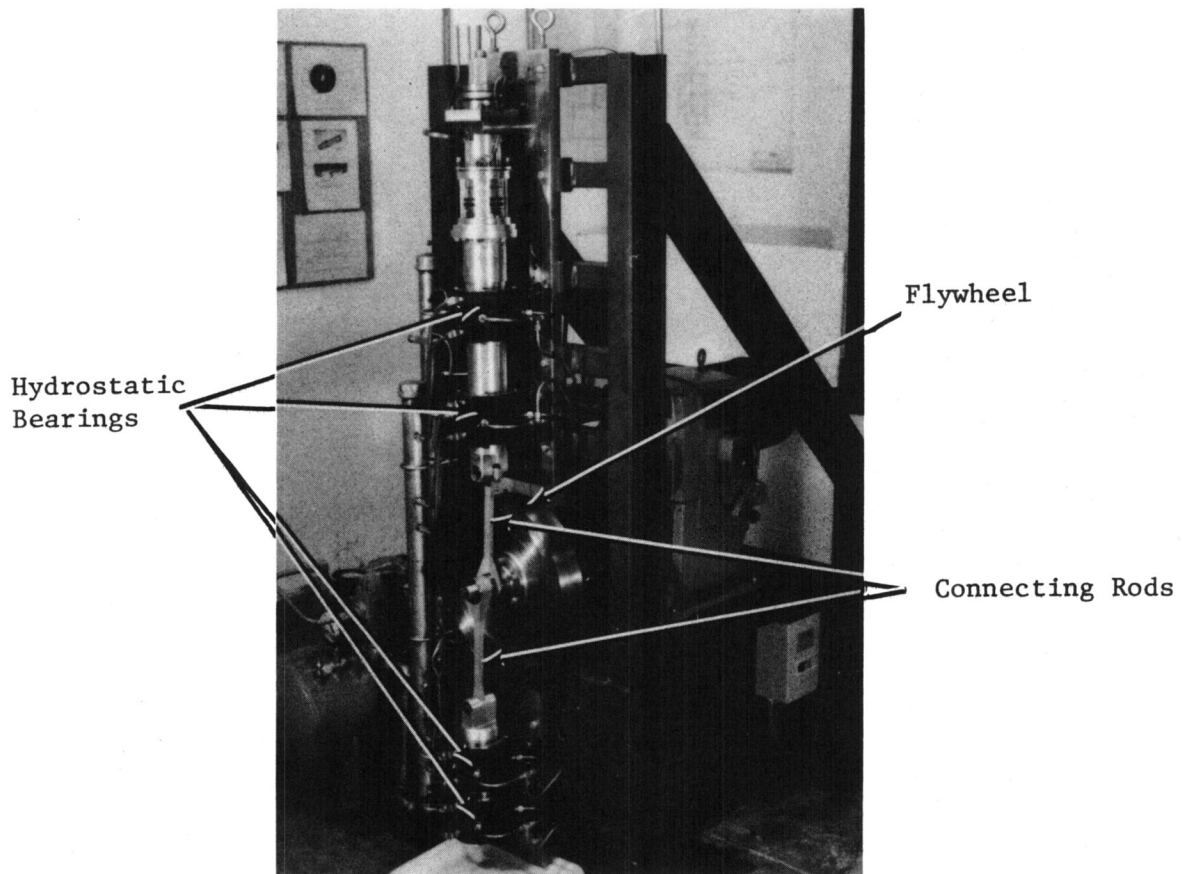


Figure 2 Overall View of Experimental Apparatus
With Guard Cage Removed

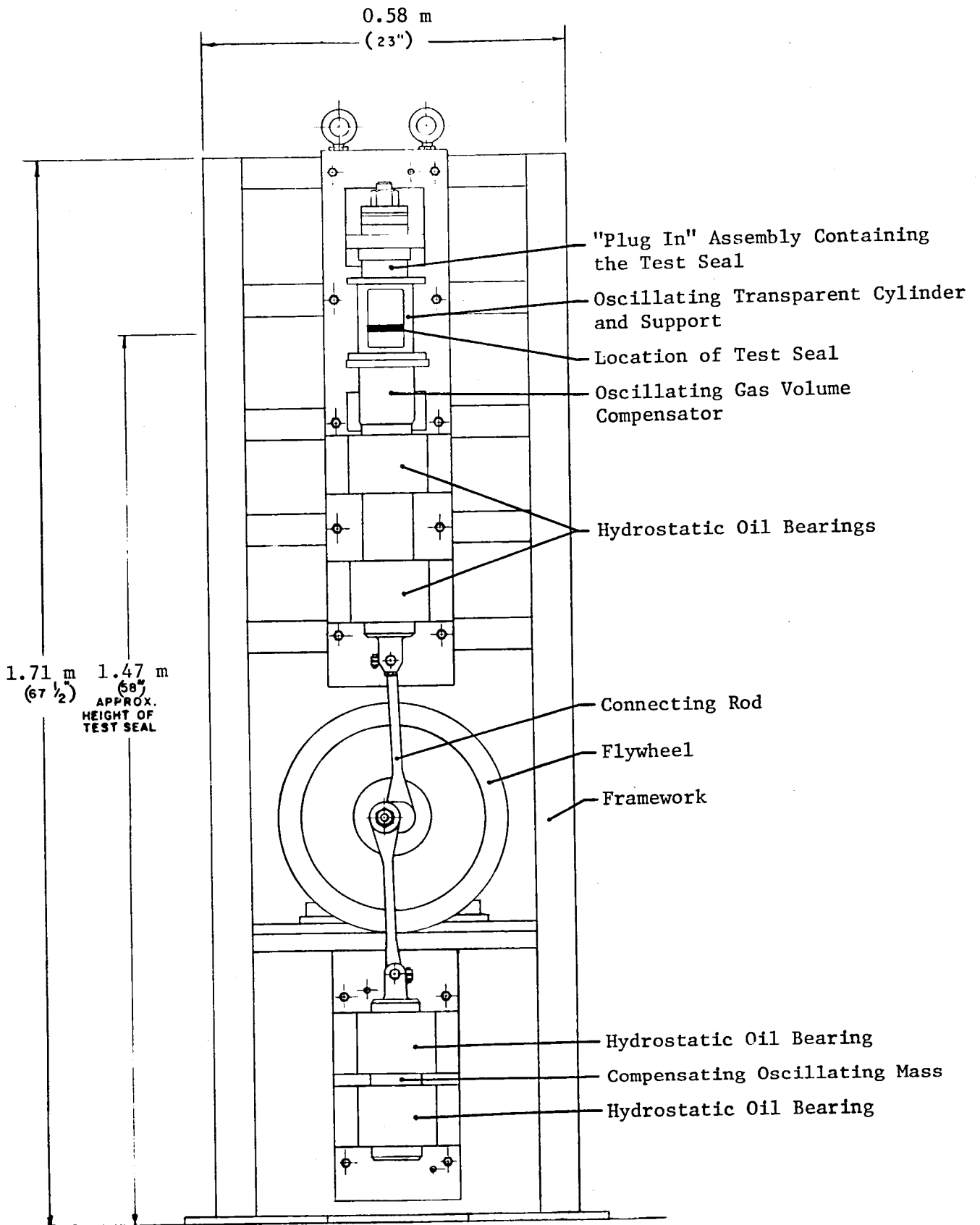


Figure 3
Layout of the Experimental Device

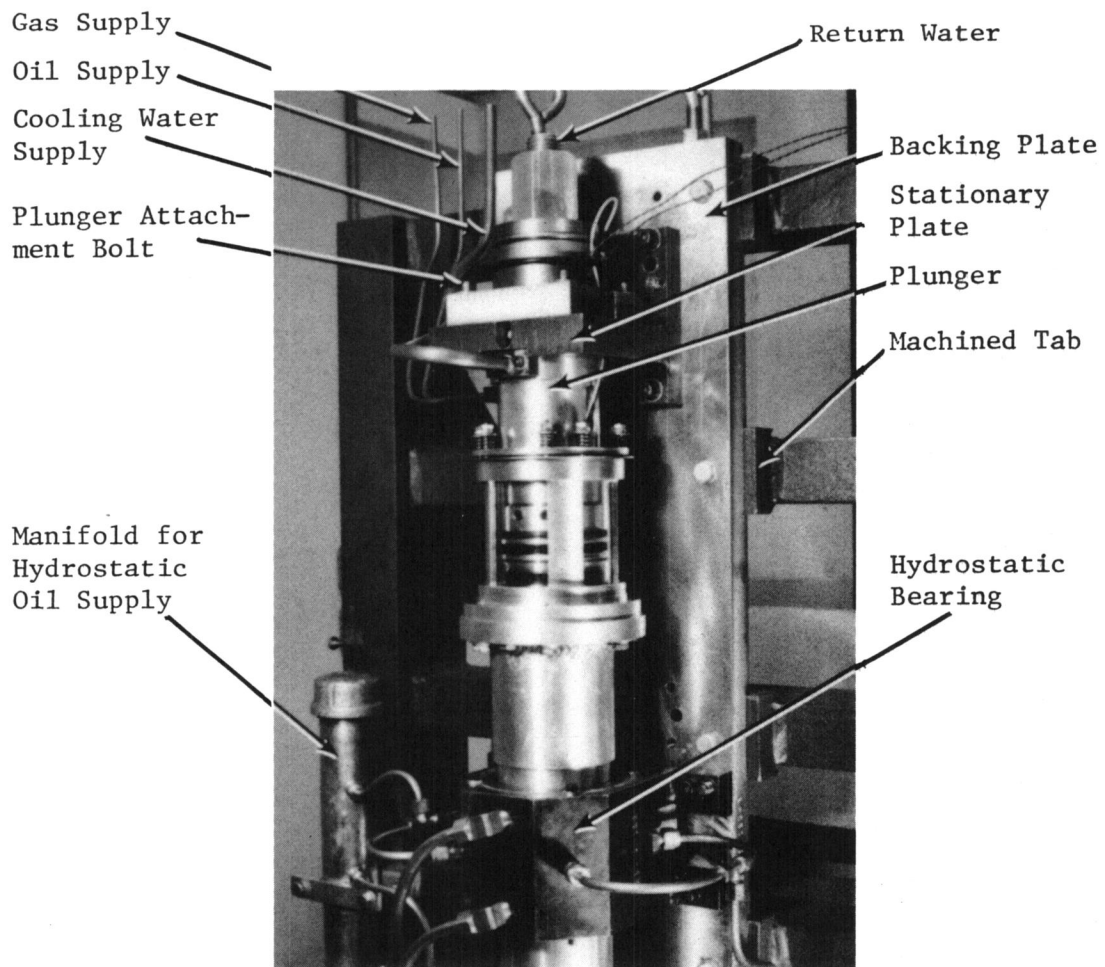


Figure 4 Close-up of Upper End of Experimental Apparatus

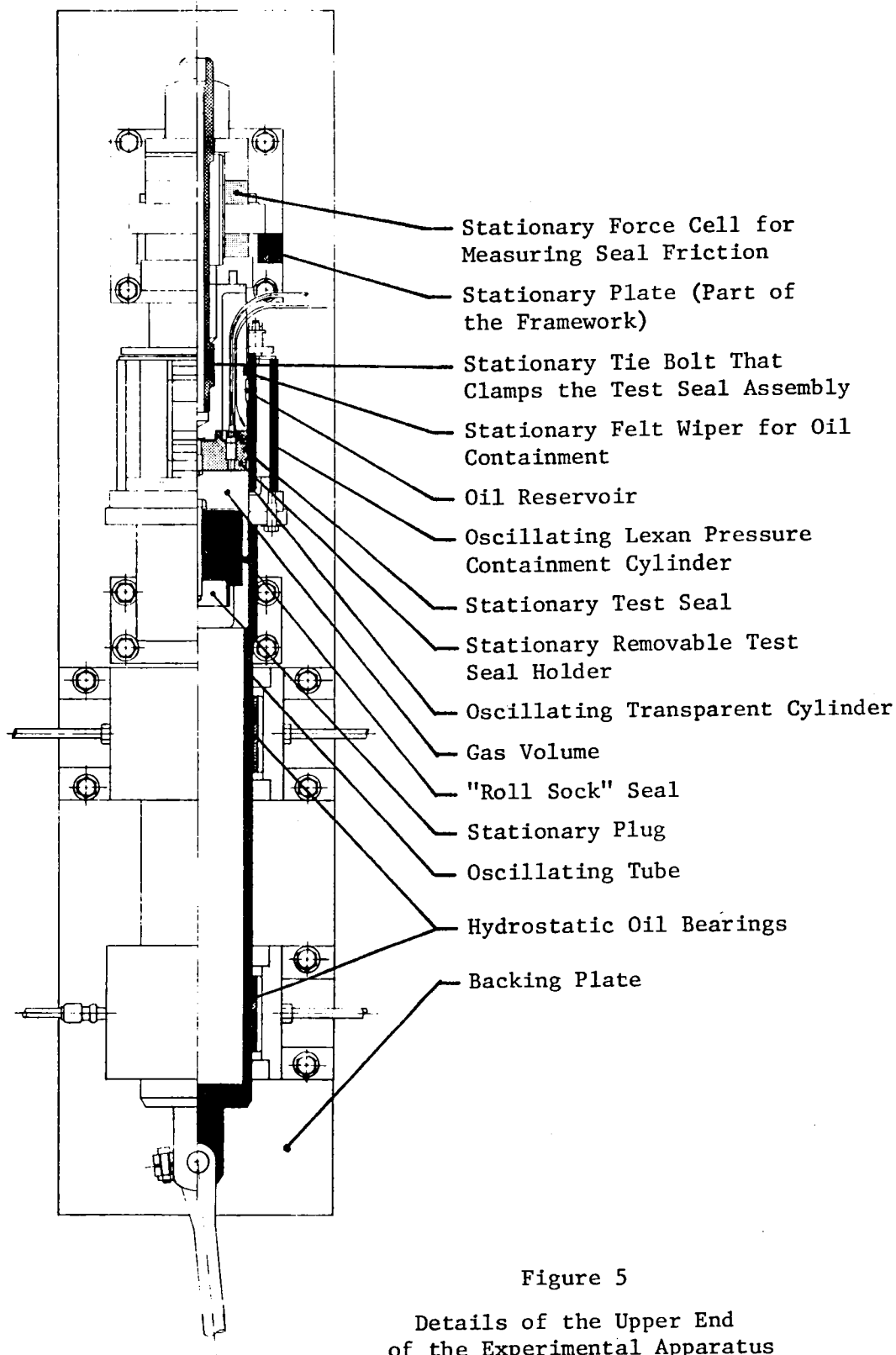


Figure 5

Details of the Upper End of the Experimental Apparatus

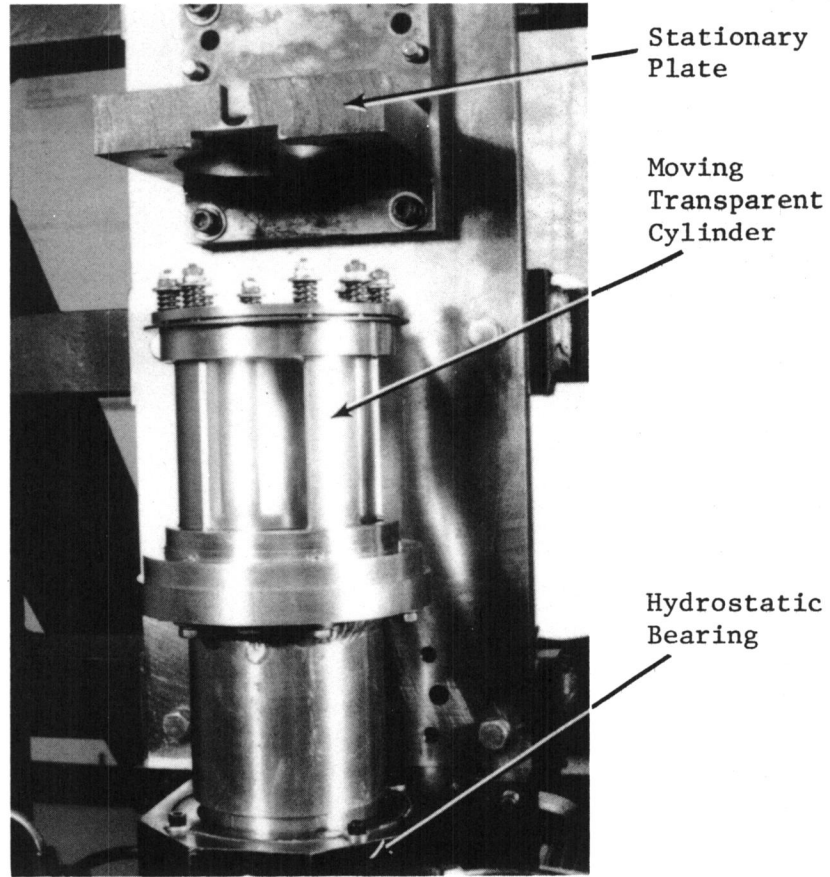


Figure 6 Transparent Cylinder Region of Apparatus With Plunger Removed

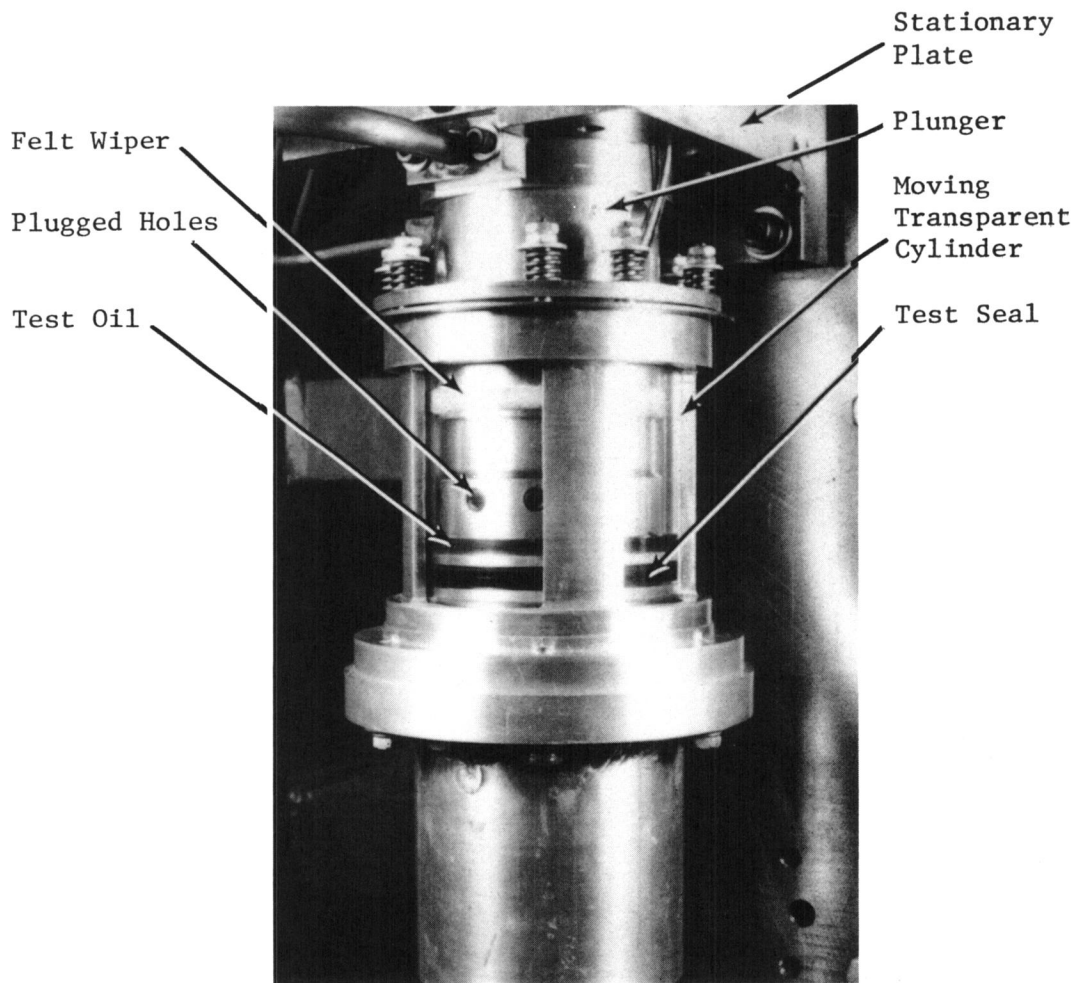


Figure 7 Transparent Cylinder Region of Apparatus With Plunger Installed

apparatus with the plunger removed. The stationary plate is at the top of the picture, the moving transparent cylinder (and its housing) are at the center of the picture, and the uppermost hydrostatic bearing is at the bottom of the picture. Figure 7 is a similar view with the plunger installed. In Figure 7, the bottom of the stationary plate can be seen at the very top of this photograph.

Figures 6 and 7 (as well as the schematic of Figure 5) indicate how the transparent cylinder is supported in its housing. The cylinder is inside the webs of this aluminum housing and is compressed axially by the 8 springs. Alignment is produced by boring the cylinder at the same time as the bottom face is cut. Not shown in the photographs (but included in Figure 5) is an outer transparent cylinder. This outer cylinder is also loaded axially by the 8 springs and can sustain 690 kPa (100 psi) of gas pressure. Use of the outer cylinder results in compressive, rather than tensile, gas loading on the inner cylinder. The apparatus can be run without the outer cylinder (for zero pressure drop across the seal) or with the outer cylinder (for pressure drops up to 690 kPa (100 psi)) across the seal. Also not shown in the photographs (but included in Figure 5) is the containment for the gas volume. This containment is effected by attaching a Bellofram rollsock seal to the inside of the moving upper shaft. The other end of this rollsock seal is attached to a stationary plug which is in turn, attached to the backing plate. The design of the rollsock seal is such that the gas volume changes little during the reciprocating shaft motion. When the gas volume is connected to an external gas reservoir, the peak pressure variations are estimated to be less than 1.70 kPa (0.25 psi) at 690 kPa (100 psi) gas pressure.

Figure 7 shows several significant parts of the stationary plunger. Through the transparent Lexan cylinder, one can see, proceeding from top to bottom, the felt wiper which contains the seal oil, plugged holes which are associated with the water cooling system, a volume of test oil (upper dark band), and the test seal (lower dark band). The parts can be seen more easily in Figure 8, which shows only the plunger.

In Figure 8, the test seal can be seen near the lower right end of the plunger. This end of the plunger (the test seal holder) can be disassembled from the body of the plunger by removing the 4 bolts at the lower right of the picture. This is accomplished by first straightening the locking tabs. Having the test seal holder replaceable allows various types of seals to be tested. For each type of seal, a holder is required in which a groove appropriate to the seal has been cut. For the holder in Figure 8, the groove has been cut for the T seal shown in the photograph.

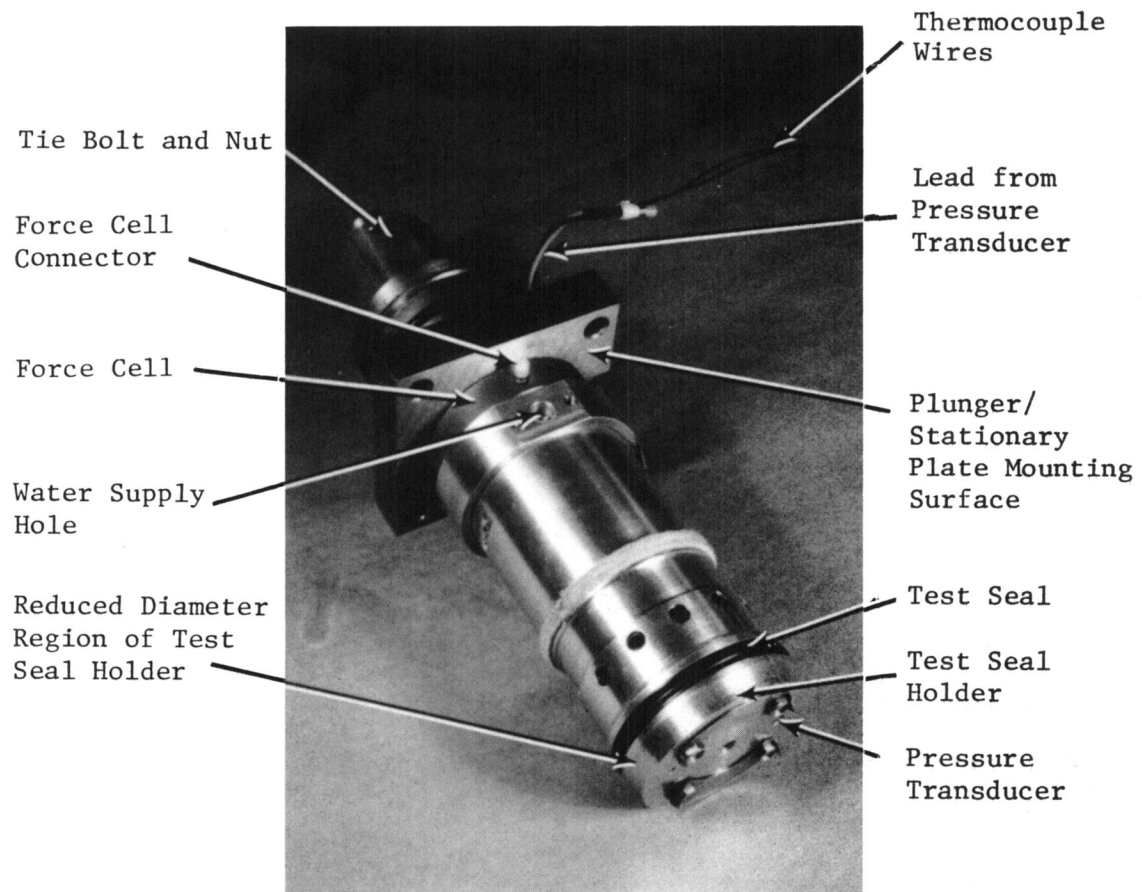
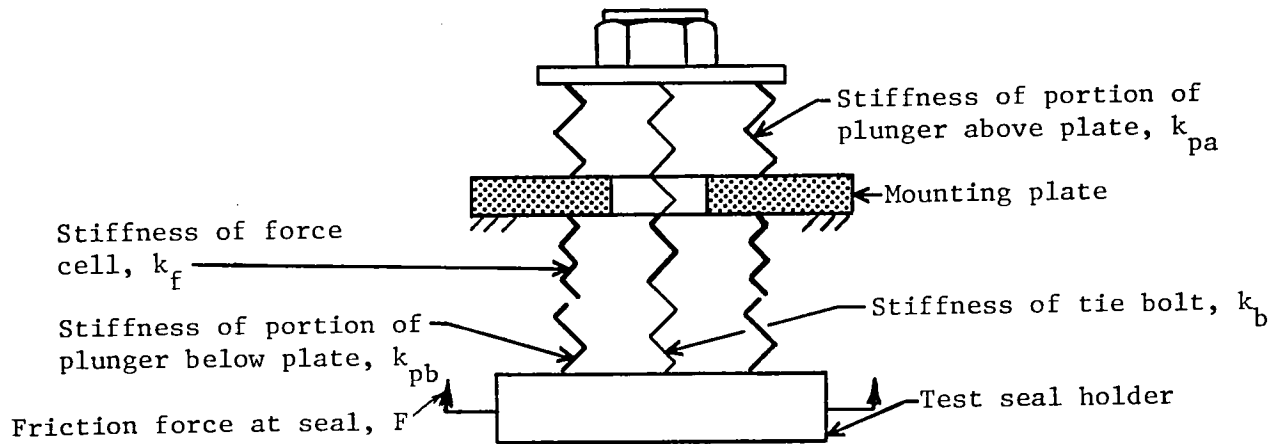


Figure 8 Close-up of the Plunger

The plunger mounts to the stationary plate at the lower surface of the plate having the holes (Figure 8). Just below this plate is a force cell (Kiag type 9071) for measuring the friction force produced by the seal. The connector for the lead is the white protrusion just above the cooling water supply hole. The force cell is preloaded between the plate and the remainder of the plunger body by a hollow tie bolt (the nut that controls the magnitude of the preload is at the upper left in Figure 8). The assembly of the plunger to produce the preload is a one-time operation--the tie bolt portion of the plunger need never be disassembled at any time during all planned use of the apparatus.

Use of the tie bolt to preload the force cell is desirable for several reasons. The first reason is that the hollow tie bolt is a convenient way to remove the cooling water from the plunger. The second reason is that with a properly designed tie bolt only one force cell need be used to measure seal friction. To see this, the following diagram, which shows the axial elastic behavior of the plunger, can be used:



The diagram indicates four stiffnesses, those of the tie bolt, of the plunger below the plate, of the plunger above the plate, and of the force cell. A friction force, F , at the seal/cylinder interface must produce an axial displacement, d , of the test seal holder such that

$$F = \left(\frac{k_b k_{pa}}{k_{pa} + k_b} + \frac{k_f k_{pb}}{k_{pb} + k_f} \right) \cdot d$$

Also, the force, F_c , measured by the force cell is

$$F_c = \frac{k_f k_{pb}}{k_{pb} + k_f} \cdot d$$

so that

$$\frac{F_c}{F} = \frac{k_f k_{pb} / (k_{pb} + k_f)}{k_b k_{pa} / (k_{pa} + k_b) + k_f k_{pb} / (k_{pb} + k_f)}$$

or

$$\frac{F_c}{F} = \frac{1}{k_b k_{pa} (k_{pb} + k_f) / (k_f k_{pb} (k_{pa} + k_b)) + 1}$$

Now, if the design of the plunger is such that k_b is small compared to k_{pa} , the above expression reduces to

$$\frac{F_c}{F} = \frac{1}{k_b/k_f + k_b/k_{pb} + 1}$$

which further reduces to $F_c/F = 1$ if k_b/k_f and k_b/k_{pb} are small compared to 1. The plunger was designed to meet these conditions so that the applied and the measured forces are nearly identical. Evidence that this design objective has been met is presented in the preliminary results section below.

Figure 8 can be used to point out other design features of the plunger. One such feature is the reduced diameter region below the test seal. This has the function of easing the installation of the test seal on the test seal holder. Another feature is the pressure transducer for measuring the small cyclic variations of gas pressure. This measurement can be made on a real time basis and can be used to subtract the gas pressure force variations from the force measured by the load cell. The lead from the pressure transducer can be seen (white wire) emerging from the plunger at the upper center of the photograph. The other two wires are thermocouple leads for measuring temperatures in the oil above the seal and in the test seal holder near the seal groove.

The present test seal holder (that shown in Figures 1, 2, 4, 7 and 8) was designed to accept a T seal. This T seal consists of an elastomeric portion, one part of which contacts the cylinder, and two fiber back-up rings. The mounted seal is shown schematically in Figure 9.

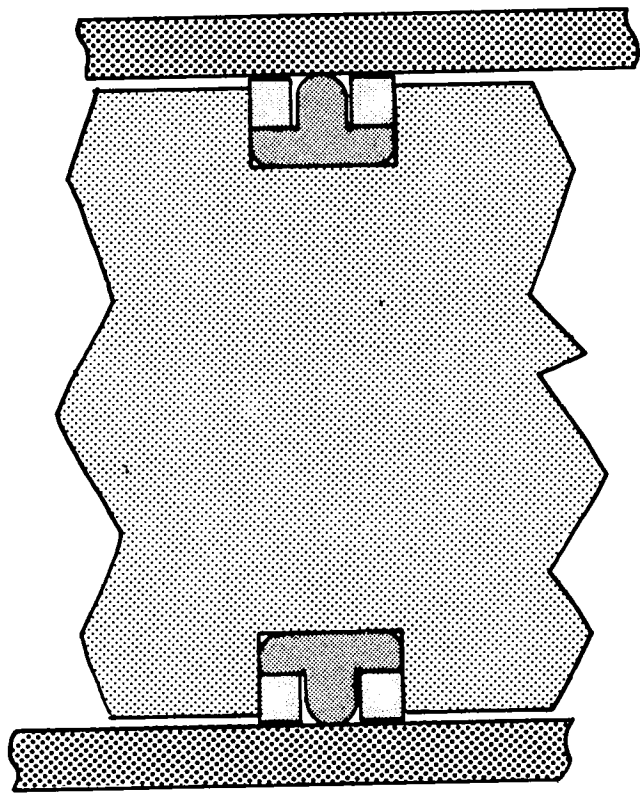


Figure 9 Schematic of "T" Seal as Installed

Figure 10 is a photograph of the seal prior to installation on the test seal holder. The elastomeric material of this seal is Viton, a fluoroelastomer.

Preliminary Results

In this section, results obtained from preliminary use of the experimental apparatus are presented. These results concern the operation of the apparatus rather than the behavior of the seal. The results show that the device performs well mechanically and that its force measurement capability is as intended.

Upon completing the assembly of the rig, the first measurement made was of the stiffness of the hydrostatic bearings. A high degree of stiffness is desirable in order that dynamic forces not result in appreciable horizontal motion of the transparent cylinder with respect to the seal. The measurement was made by simply applying a steady horizontal force between the two hydrostatic bearings of the upper shaft. This force was measured with a spring scale. The displacement was measured with a dial indicator. The results showed the both upper hydrostatic bearings together have a horizontal stiffness of 1.2×10^5 N/mm (700,000 pounds per inch). The stiffness of each bearing is, therefore, at least 0.6×10^5 N/mm (350,000 pounds per inch).

To show what the effects of this stiffness are on the seal/cylinder interface, two rough calculations can be made. For these calculations, the following conservative assumptions are used:

- The horizontal stiffness of each upper hydrostatic bearing is 0.6×10^5 N/mm (350,000 pounds per inch).
- The horizontal dynamic force in the apparatus is 290 N (65 pounds). This force is calculated using the geometry of the connecting rod (the ratio of crank length to connecting rod length is 0.05), the maximum frequency (50 Hz), and the mass of the upper reciprocating shaft. This mass is less than 4.5 kg (10 pounds mass).
- The horizontal dynamic force is applied 0.13 m (5.0 inches) below the lower hydrostatic bearing of the upper shaft and these hydrostatic bearings are 0.21 m (8.1 inches) apart.
- The seal is 0.22 m (8.5 inches) above the upper hydrostatic bearing of the upper shaft.

The displacement of the transparent cylinder results from two effects-- translation of the upper shaft due to the 290 N (65 pound) force, and rocking in the vertical plane due to the same 290 N (65 pound) force. The translation, t , can be computed simply from

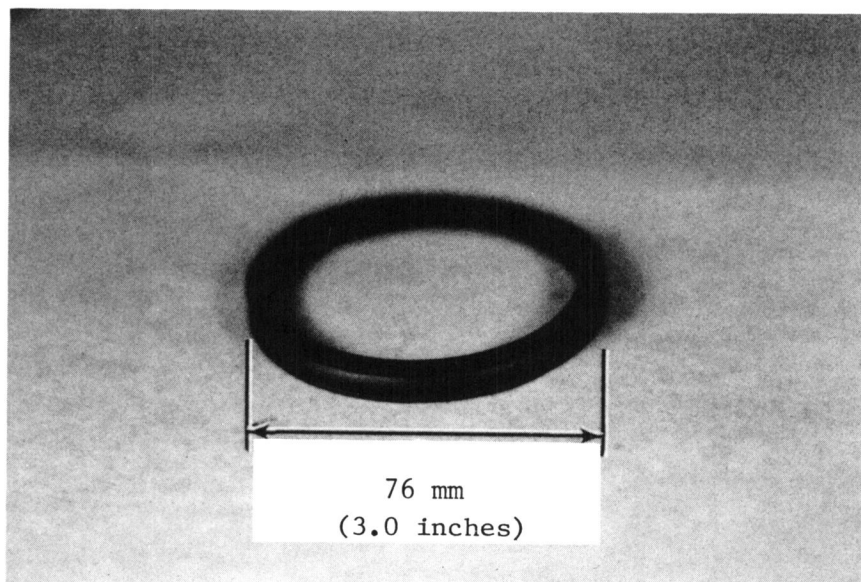


Figure 10 Test "T" Seal

$$t = \frac{290}{2(0.6 \times 10^8)} = 2.4 \times 10^{-3} \text{ mm (0.1 mils)}$$

This motion is clearly negligible when compared to a typical seal preload. The displacement, r , due to rocking can be computed from:

$$r = \frac{290}{0.6 \times 10^8} \left[\frac{0.13 + 0.105}{0.21} \right] \left[\frac{0.22 + 0.105}{0.105} \right]$$

$$= 0.17 \times 10^{-2} \text{ mm (0.67 mils)}$$

This motion is also negligible when compared to a typical seal preload.

The next measurements to be made after assembly of the experimental apparatus were concerned with the smoothness of the device during operation. To obtain an evaluation of this, an accelerometer was placed on the backing plate at a height approximately corresponding to that of the seal. The sensitive axis of the accelerometer was oriented first in the vertical direction and then in the horizontal direction (such that the axis was in the plane of the backing plate). The drive motor was operated at about 900 rpm. The plunger was removed.

The results obtained are shown in Figures 11 and 12, where Figure 11 is for the horizontal acceleration and Figure 12 is for the vertical acceleration. Each figure gives an average spectrum made from 8 spectrum analyses of the measured acceleration. The analyses were made in the 0-200 Hz range and therefore cover the flywheel frequency of 15 Hz as well as the drive motor frequency of about 30 Hz.

The plots show that the vibration levels in the apparatus are relatively low. Reference to typical charts* used in vibration analysis indicates that both the horizontal and vertical vibration levels fall in the "smooth" range for machine operation. This shows that a satisfactory dynamic environment exists in the experimental apparatus for conducting the optical interferometry studies.

The plots show that the horizontal acceleration is greater than the vertical acceleration, particularly in the regions near 60 Hz and 30 Hz. This may have been anticipated, since the vertical, but not the horizontal, forces are balanced by the lower reciprocating shaft. However, these two frequencies are not directly associated with the rotation speed of the flywheel, which is about 15 Hz.

*One such chart is the "Vibration Computer," a paper slide rule, published by Endevco Corp.

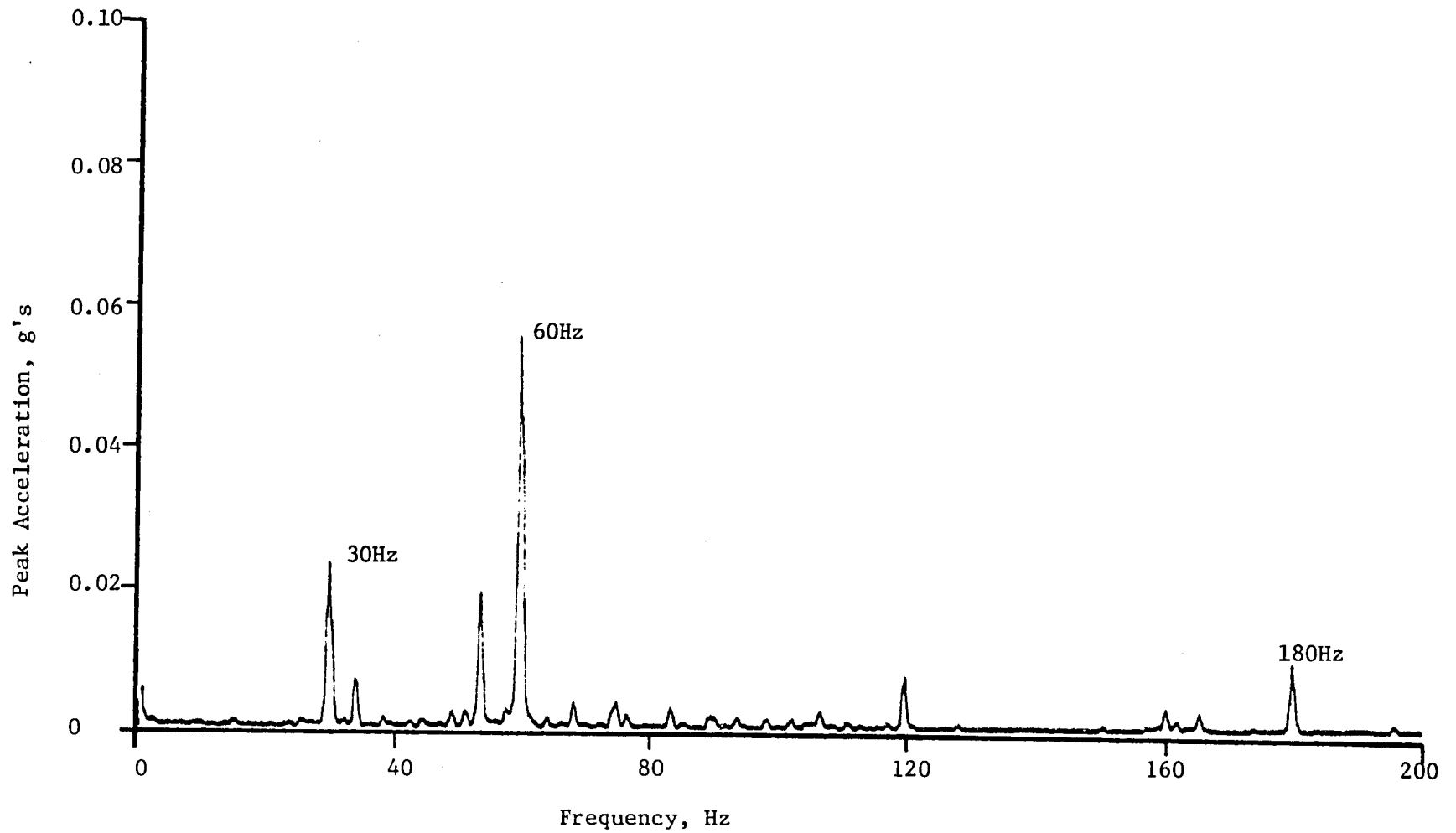


Figure 11 Average of 8 Spectra for Horizontal Acceleration

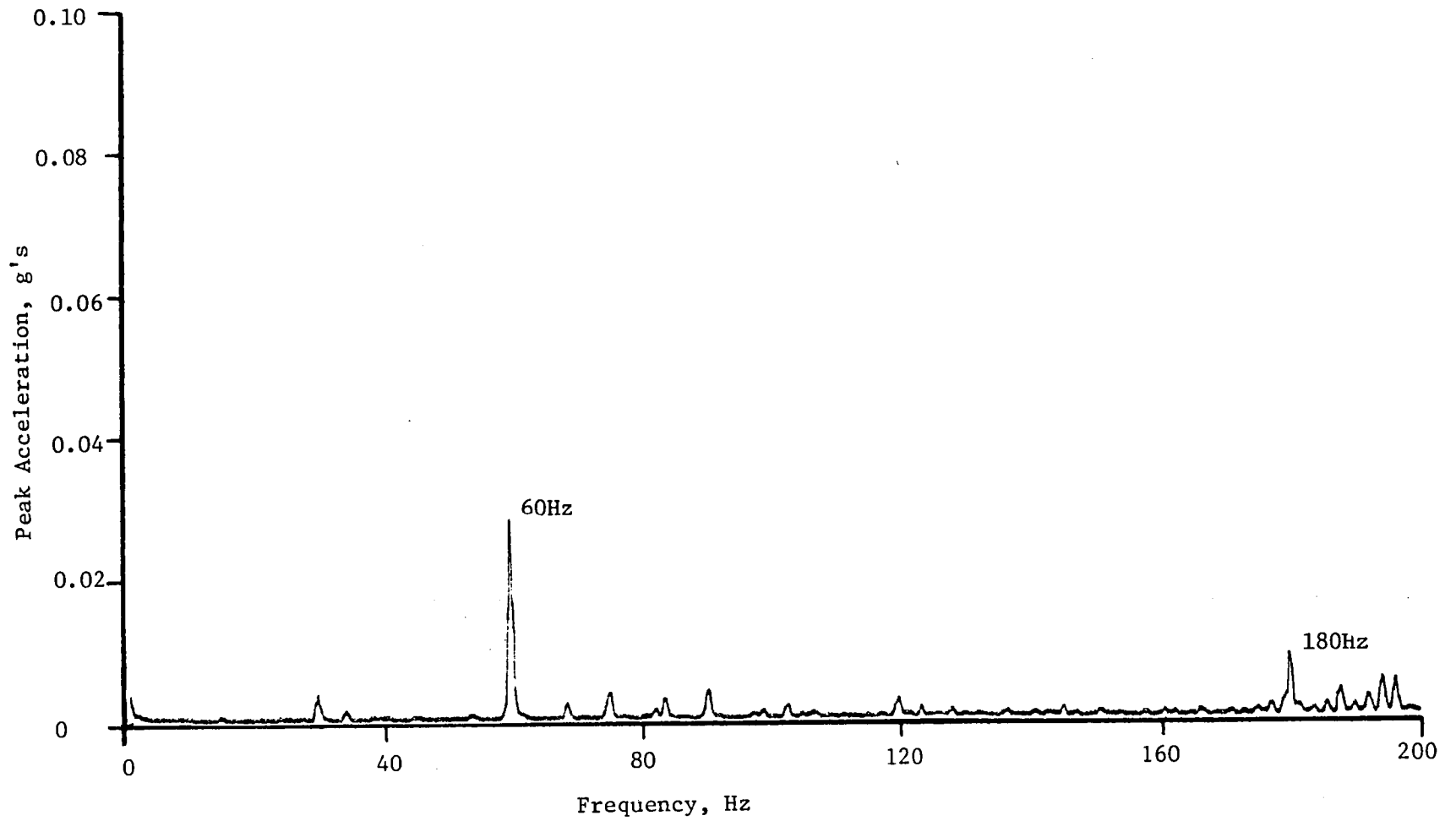


Figure 12 Average of 8 Spectra for Vertical Acceleration

The origin of the acceleration signals at these frequencies is of interest, so additional spectra, Figures 13 and 14, were produced. For these spectra, all conditions were identical to those of Figures 11 and 12; however, the coupling between the drive motor and the flywheel was disconnected. Consequently, the spectra of Figures 13 and 14 were obtained under the condition of a stationary flywheel (and stationary connecting rods) but a rotating drive motor whose output speed was about 900 rpm (15 Hz).

Comparison of Figures 11 and 13 shows that most of the horizontal acceleration at 30 Hz, 60 Hz, and 180 Hz is not due to the flywheel and reciprocating masses; in fact, the only significant flywheel/reciprocating mass signal is the one at about 53 Hz. Comparison of Figures 12 and 14 gives a similar result. For vertical vibrations, it can be observed that the flywheel/reciprocating mass portion of the apparatus contributes essentially nothing to the vibration spectrum.

One further plot was made. The intent of this plot was to determine whether the 30, 60, and 180 Hz signals were due to the drive motor or, perhaps, due to some electrical source. The plot, Figure 15, is an instantaneous spectrum made from the accelerometer with the drive motor switched off. The plot shows that significant portions of the 60 Hz and 180 Hz signals are still present. These residual portions are, therefore, not associated with mechanical vibrations produced either by the drive motor or by the flywheel/reciprocating mass combination. This indicates that the drive motor contributes much less vibration to the experimental apparatus than Figures 13 and 14 suggest. Consequently, of the vibration levels of Figure 11, half of the 60 Hz and almost all of the 180 Hz signals are not mechanically produced. Also, most of the remainder of the 60 Hz signal and most of the 30 Hz signal are produced by the drive motor.

One other mechanical aspect of rig operation is of concern. That aspect is whether a vertical force (i.e., the seal friction force) applied to the plunger at the seal is accurately measured by the force transducer. (The stiffness conditions for this to occur are considered in the section above.) To evaluate the performance of the force transducer, a reference force transducer was attached to the bottom of the plunger at the hole on the vertical centerline. To this reference transducer was fastened a spring mass combination. Schematically, the arrangement is as shown in Figure 16. The spring mass combination was arranged so that the natural frequency was about 12 Hz.

Displacement and subsequent release of the mass resulted in an oscillation which applied a sinusoidal vertical force to the reference cell. The force was also transmitted to the bottom of the plunger. Both the force measured by the reference force cell and that measured

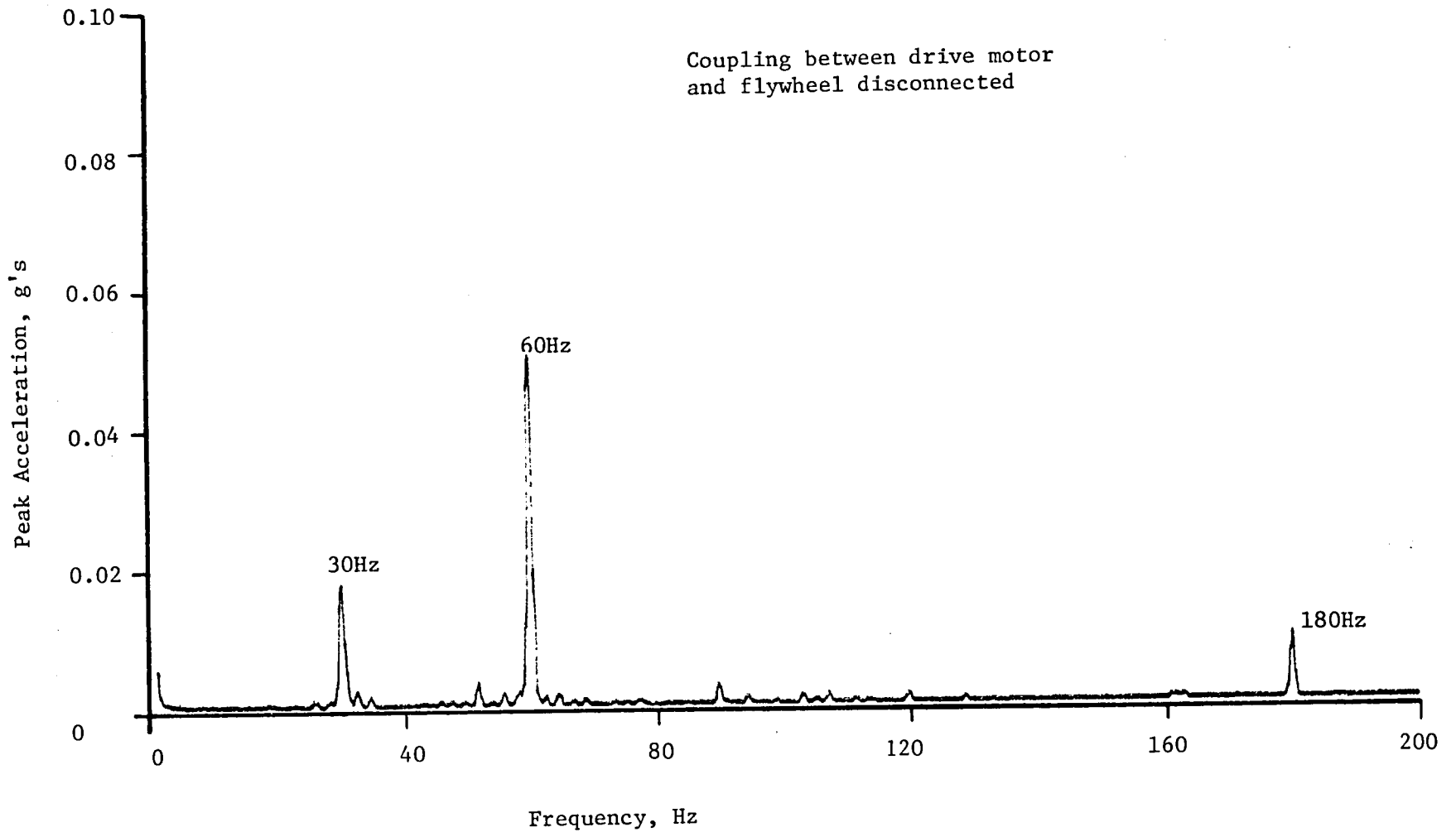


Figure 13 Average of 8 Spectra for Horizontal Acceleration

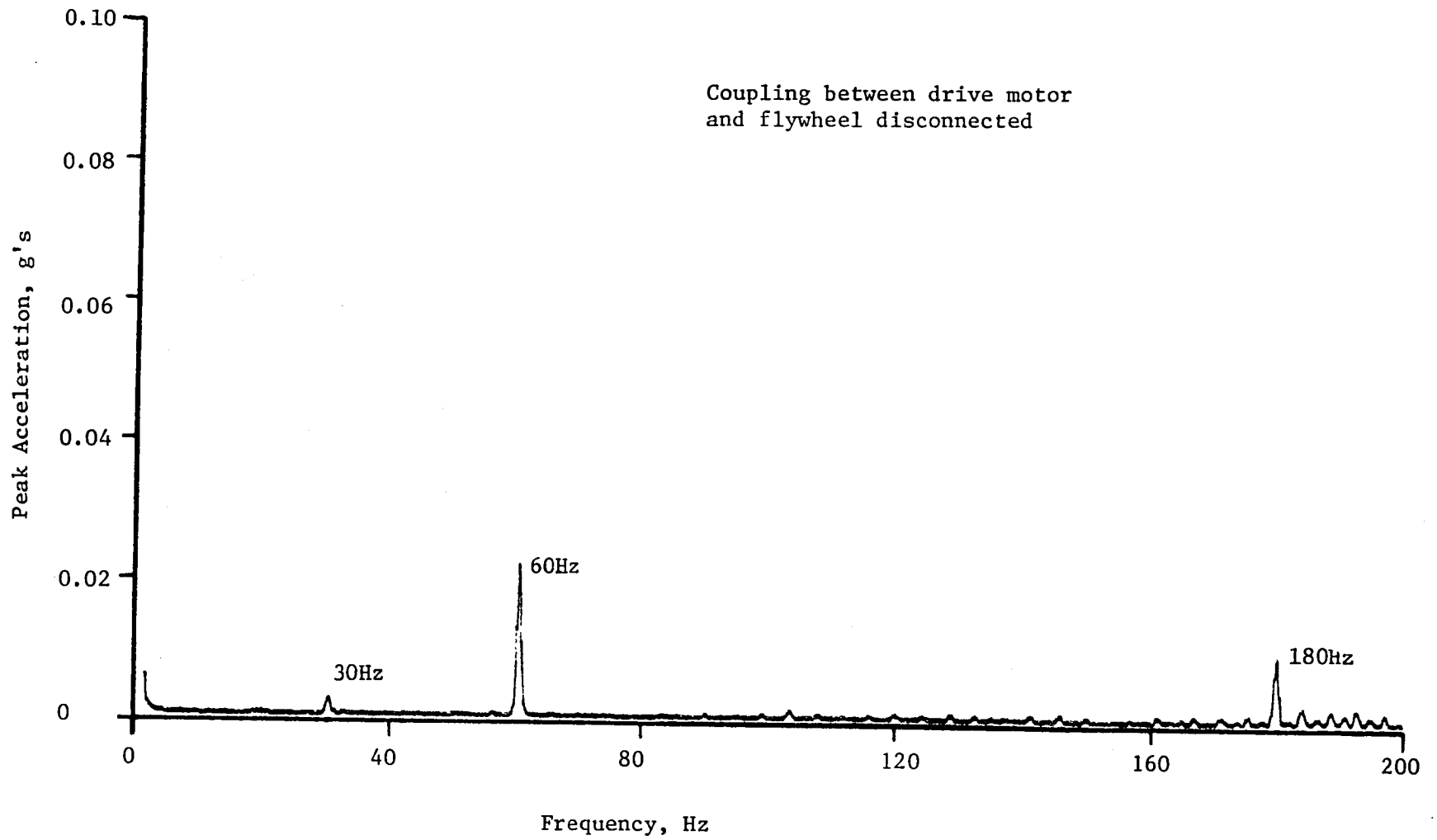


Figure 14 Average of 8 Spectra for Vertical Acceleration

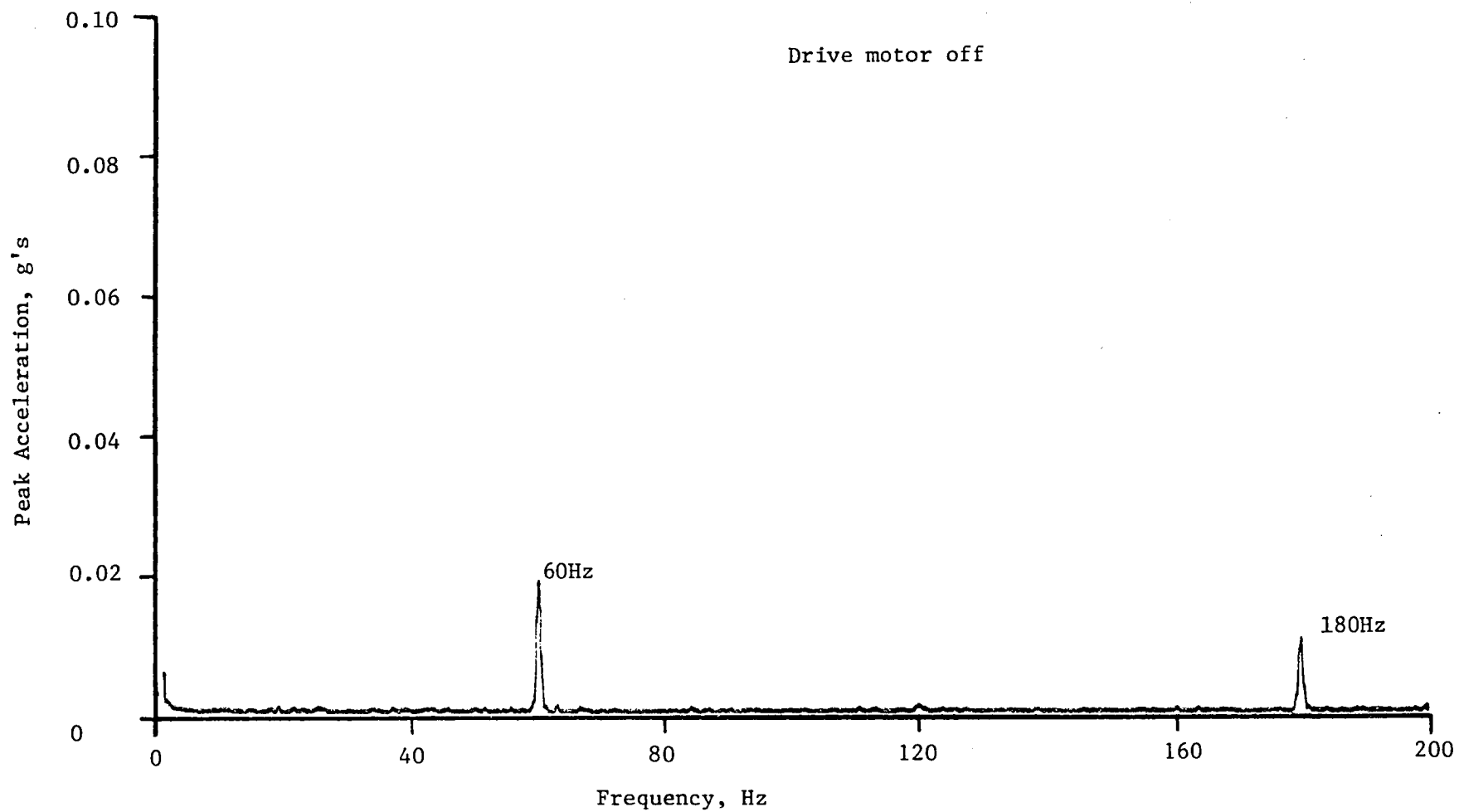


Figure 15 Instantaneous Spectrum for Vertical Acceleration

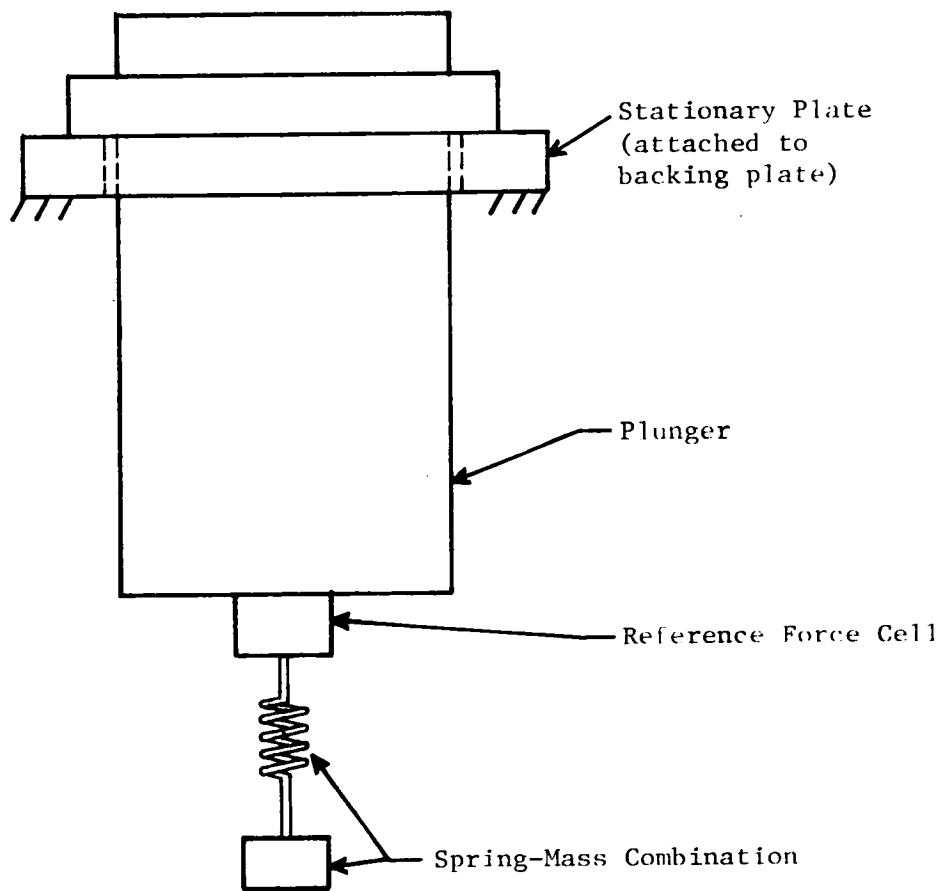


Figure 16 Location of Reference Force Cell and Spring-Mass with Respect to the Plunger

by the force cell in the plunger were recorded by photographing the voltage outputs from their electronic drivers. These drivers were set so that each load cell had the same sensitivity 2.25 mV/N (10 mV/pound). The results are given in the photographs of Figure 17.

The three sets of time traces in Figure 17 are for segments from three separate releases of the mass. The first set corresponds to a vibration amplitude such that no spring bottoming occurred. The second set involved some bottoming near the start of the trace shown. The third set was made under a condition of relatively hard spring bottoming at the start of the oscillation. In each set of traces, two curves are given--the top one is for the reference force transducer and the bottom one is for the force cell in the plunger. The vertical and horizontal scales are as shown in the figure.

All three photographs indicate very good correspondence between the forces measured by the two cells. The time traces are in phase, indicating that mass effects are not significant, even for the higher frequencies that occurred. The magnitudes correspond well, suggesting that the required stiffness relationships for the tie bolt, plunger assembly and plunger force cell have been met. It therefore, appears that the force cell in the plunger does satisfactorily measure vertical forces (i.e., seal friction forces) that are applied to the test seal holder.

Interferometry

The accurate measurement of the film thickness at the seal/cylinder interface is central to the success of the current program. Film thickness measurements are the most sensitive means for performing analytical and experimental correlations. In addition, knowledge of film thickness can be used in studying seal friction behavior and can provide insight into how leakage can be controlled.

Historically, many techniques of measurement have been used to obtain EHD film thickness. Successful methods of film thickness detection have been developed with the aid of optical equipment, capacitive probes, inductive probes (eddy current devices), and resistance measurement sensors. The optical approach, however, is the only direct scheme which can provide a full point-by-point topographical film measurement within a reciprocating seal interface. In addition, the optical approach allows gradients within the contact region to be observed and film thicknesses as small as 2 microinches to be measured reliably. For these reasons, an optical approach to film thickness measurement has been taken and the moving transparent cylinder concept has been used.

Application of optical procedures to measuring thicknesses such as those in elastohydrodynamic contacts invariably involves the interferometry technique. This technique, which relies on the interference

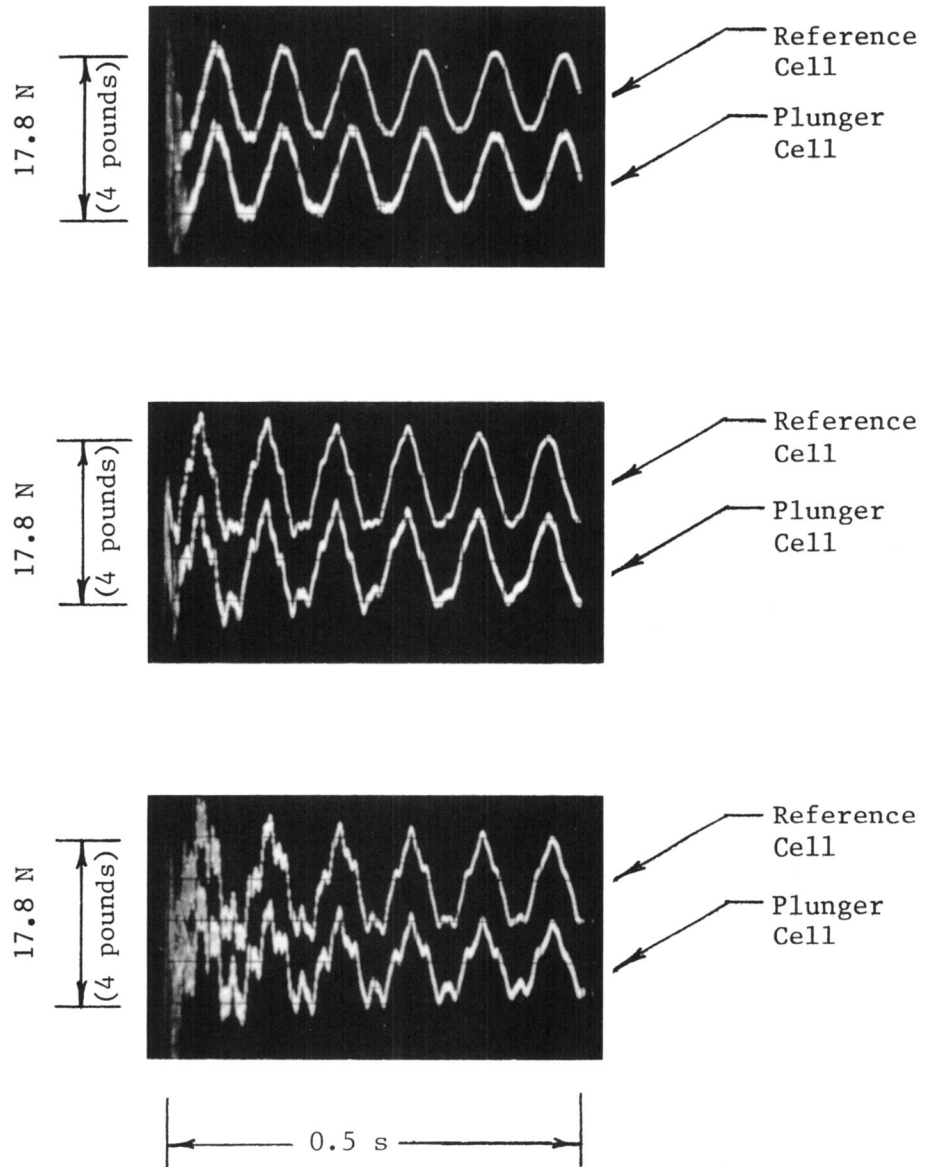


Figure 17 Comparison of Outputs from Reference Force Cell and from Plunger Force Cell

between two coherent beams of light, is described briefly below.

Figure 18 shows a greatly expanded drawing of the seal/cylinder contact zone. The orientation of the seal and cylinder in the drawing is the same as that in the experimental apparatus--the section is made by passing a vertical plane through the vertical axis of the transparent cylinder. The drawing includes a light source which sends a light beam, A, toward the outer surface of the cylinder. This beam, which also lies in this vertical plane, has an intensity I_a .

At the outer surface of the cylinder, this beam splits into two parts; one part, B, having intensity I_b , is reflected. The other part

enters the cylinder wall and proceeds to the inner surface of the cylinder. At this surface, the beam splits such that one part is reflected internally in the glass and the other part passes into the oil, to the surface of the elastomer, and then back to the inner wall of the glass. The viewer, looking in the direction indicated sees beams B, C, and D, where C is from the glass/oil interface and D is from the seal/oil interface. If certain physical conditions are met, interference between beams C and D will occur. One of these conditions is that the ratio of the thickness h to the wavelength of the incident light is $\frac{1}{2}$, 1, $1\frac{1}{2}$, etc. If the film thickness changes with axial position, the viewer will see alternating light and dark bands. These bands are termed interference fringes. It is evident that the axial spacing of these fringes gives directly a map of the changes in film thickness throughout the contact region.

Several physical conditions must be met, in addition to the proper relationship between wavelength and film thickness, for observable fringes to exist. One of these conditions is that the beam intensities I_c and I_d be approximately equal. Another condition is that the intensities I_c and I_d be sufficiently high to allow them to be seen and photographed. A third condition is that I_b be sufficiently small (or be somehow eliminated) so that it does not overwhelm the interference between C and D. Finally, the interface between the glass and the oil and the interface between the oil and the seal must be sufficiently specular that beams C and D are well defined.

Satisfying the above conditions for an elastomeric seal is not straightforward. The most difficult problem to overcome is the creation of a specular (optically smooth) surface at the oil/seal interface. Additional, but more readily overcome, problems are to minimize the effects of beam B, and to produce sufficiently intense beams C and D. Because of these problems, a little optical interferometry is done using elastomers such as rubber, and few reports of such work can be found in the literature. In those reports that exist (see, for example, [1] - [3]*), the problems are dealt with as follows:

*Numbers in brackets denote references.

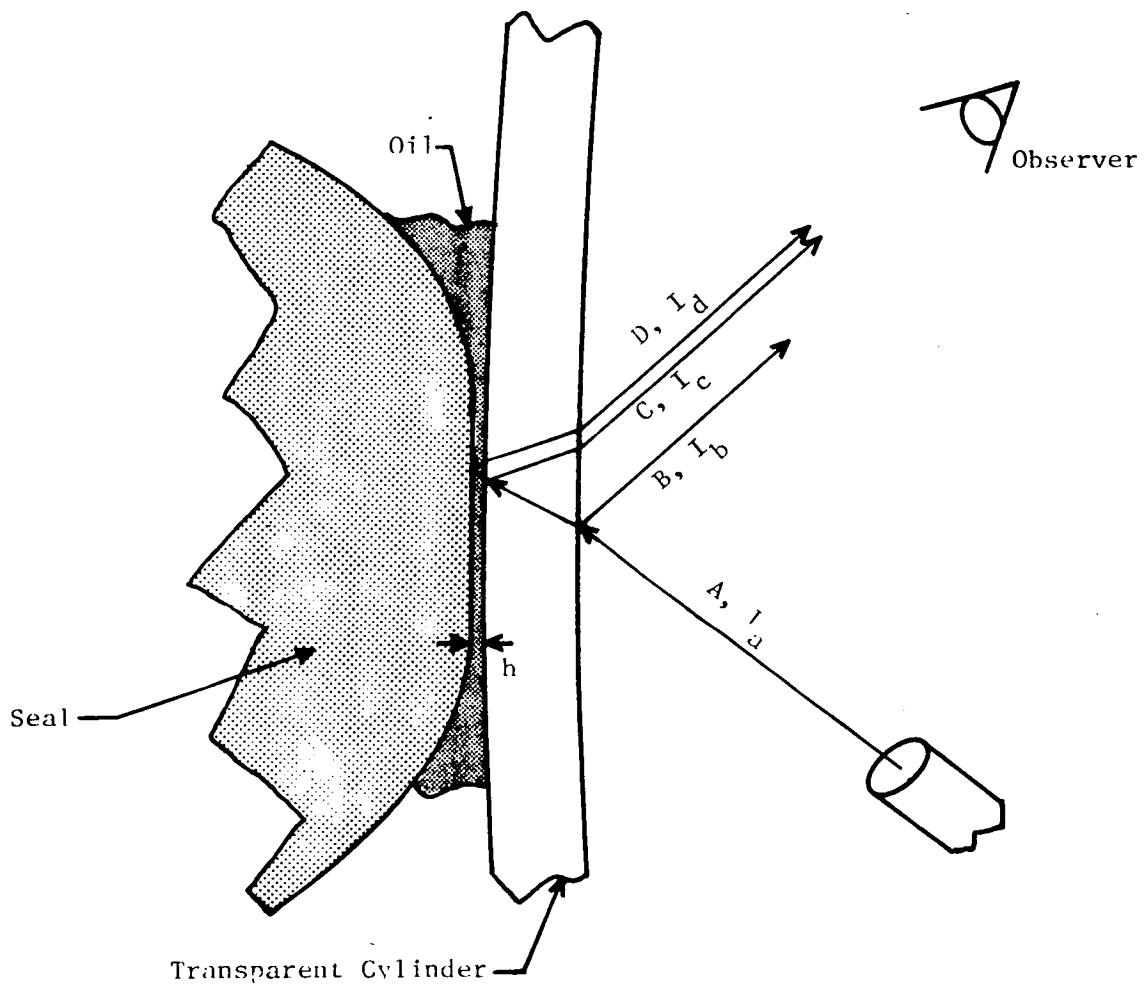


Figure 18 Schematic Drawing of Contact Zone

- . Creation of a specular surface is accomplished by molding the elastomer in an optically smooth mold, thereby creating an optically smooth surface on the elastomer.
- . Creation of a specular surface by coating the elastomer with a thin specular material such as a metal.
- . Reducing the effects of beam B by coating the outer cylinder surface with a thin film which causes destructive interference of B in the coating.
- . Increasing the intensity I_c by coating the interior wall of the cylinder and increasing the intensity I_d with a coating on the elastomer.

For the present work, none of the above techniques is entirely satisfactory. It is desirable, for example, not to have to coat either surface of the cylinder. Such coatings are expensive and the one on the inner surface is prone to scratching and wear. It is also desirable not to coat the seal surface with a material which affects its deformation characteristics. Finally, it is very desirable that the experimental apparatus employ production seals rather than seals especially produced with optically smooth surfaces. The cost for the specially produced seals can be very high, their behavior may not be typical of production seals, and the optically smooth surfaces can be prone to scratching during testing.

Experimental work conducted during the course of the present effort resulted in a technique which gives good interference fringes and which is appropriate to the experimental apparatus. Interference fringes were observed and photographed with a laboratory bench setup by using:

- . A Parker Viton T Seal which was coated with shiny black lacquer. The lacquer was employed because it is flexible (to conform to the deformed rubber shape), strongly adherent to the rubber, and insoluble in the production oil used. Additionally, the lacquer coating does not crack easily or wear rapidly. It can be applied such that the thickness is small and relatively uniform.
- . An optical flat having a refractive index in the range of 1.5 to 1.6. The optical flat could be moved manually to simulate the motion of the glass cylinder.
- . An automatic transmission oil which is a candidate for use in actual Stirling engines. This oil (Mobil XRL 1032 AR-1) has a refractive index of 1.4605.
- . A fiber optics white light source having a high intensity. The end of the fiber optics tube was mounted close to the optical flat and oriented so that the angle between the incident light A and the normal to the surface of the optical flat was about 30° . This mounting method eliminates the effects of beam B by taking advantage of the (axial) distance between

beams C and B; i.e., when sighting along C and D, the beam B does not enter the field of view.

A stereo microscope with a maximum magnification of about 50X.

The fringes obtained were of sufficient optical quality that 35 mm color slides could be taken and a 16 mm color movie could be made. A black and white photograph made from a 35mm slide is shown in Figure 19. The dark zones are interference fringes which appear as multicolor bands in the original slide. The bands are very apparent in the color movies. The bands show the presence of a lubricant film between the glass and rubber at relatively low sliding speeds (less than 0.25 m/s (10 inches/second)) and at relatively large seal preloads (contact width of the seal approximately equal to half the seal width).

Once good quality interference fringes had been produced using the laboratory bench setup, additional optical work was postponed until completion of the experimental apparatus. The optical instrumentation was then installed on the device as shown in Figure 20. In the figure, the fiber optics tube can be seen at the lower center of the photograph. The microscope is at the left center.

A low magnification photograph, taken through the microscope is given as Figure 21. The photograph was taken with the cylinder stationary. The figure shows the aluminum bands on the test seal holder, the two fiber backing rings for the T seal, and the T seal itself. At this low magnification, the contact zone cannot be seen but is approximately equidistant from the two fiber backing rings. A photograph for which a higher magnification was used is given as Figure 22. For this figure, a scale was taped to the outside of the transparent cylinder. Each of the smallest divisions is one millimeter. In Figure 22, the contact zone can be seen bounded by horizontal light lines in the dark region of the T seal.

Interferometry was attempted with both a stationary transparent cylinder and for a slowly moving transparent cylinder (the flywheel was turned by hand). Fringes were obtained, but were not sufficient quality to be photographed. The lack of optical quality was caused by the transparent Lexan cylinder, whose optical characteristics were poor. Specifically, both the inner and outer surfaces of the cylinder had circumferential scratches and grooves which appeared to have been produced by turning during manufacture. These scratches and grooves tended to dissipate both the incident and emerging light. In addition, the contact region, during motion of the cylinder, was rather irregular on a microscopic level. This caused the fringes to be numerous, close together, and non-uniform.

An attempt was made to polish the Lexan cylinder. This did improve the surface quality somewhat, but not sufficiently that a smooth contact zone resulted and that sufficient optical quality was attained.

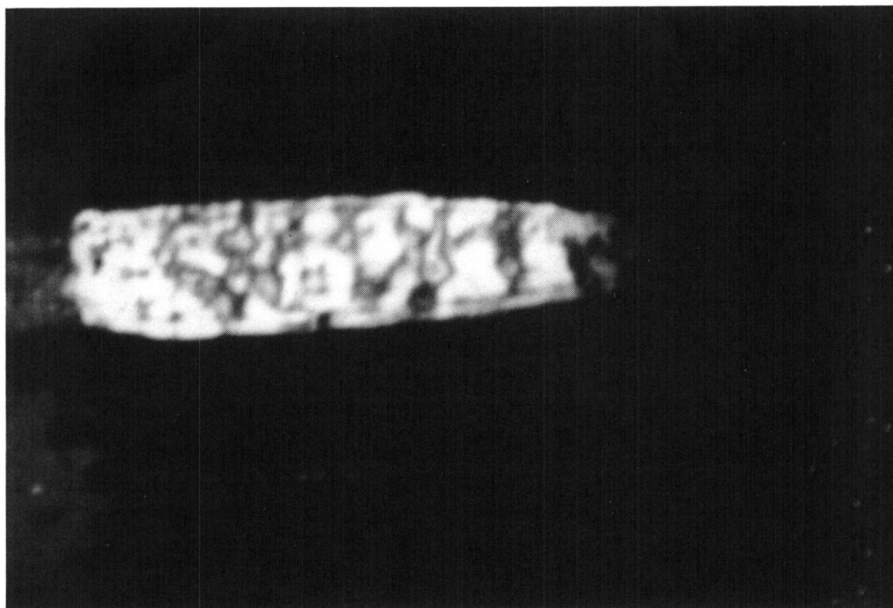


Figure 19 Interference Fringes Produced with Laboratory
Bench Setup (Original Color Photograph Made
Using White Light)

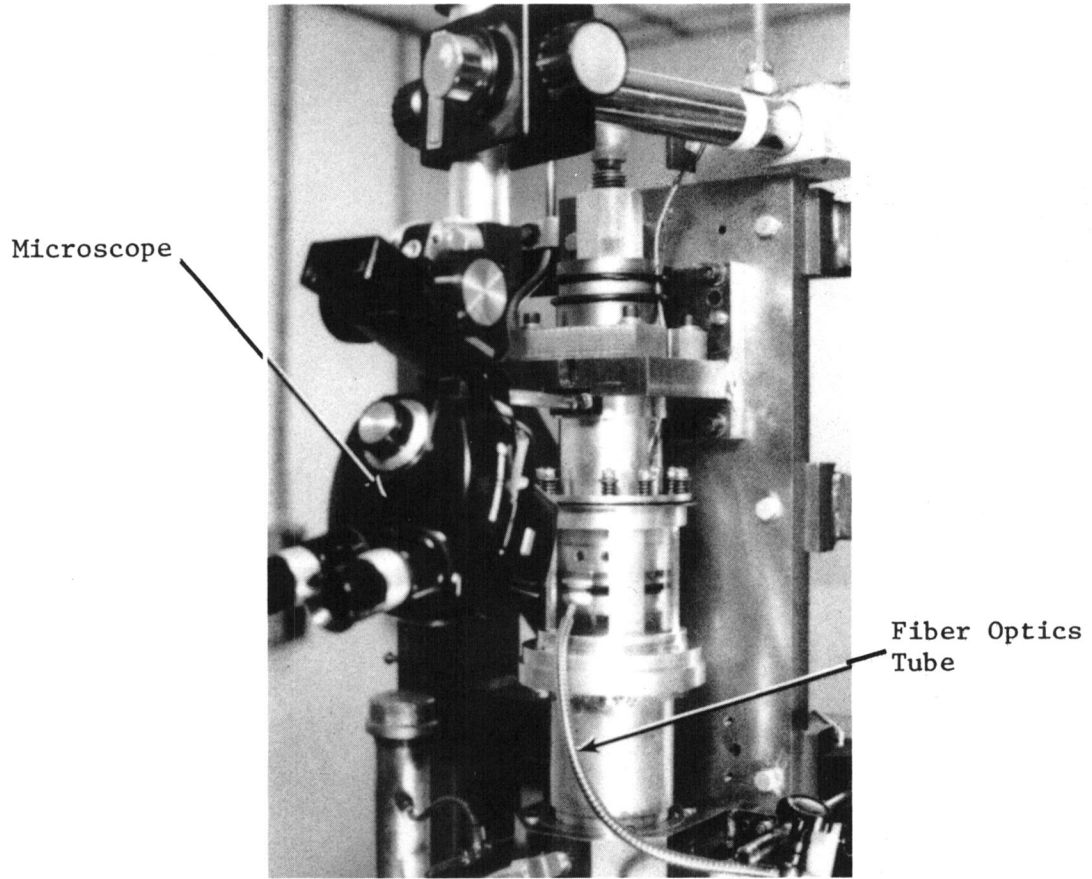


Figure 20 Experimental Apparatus With
Optical Instrumentation Installed

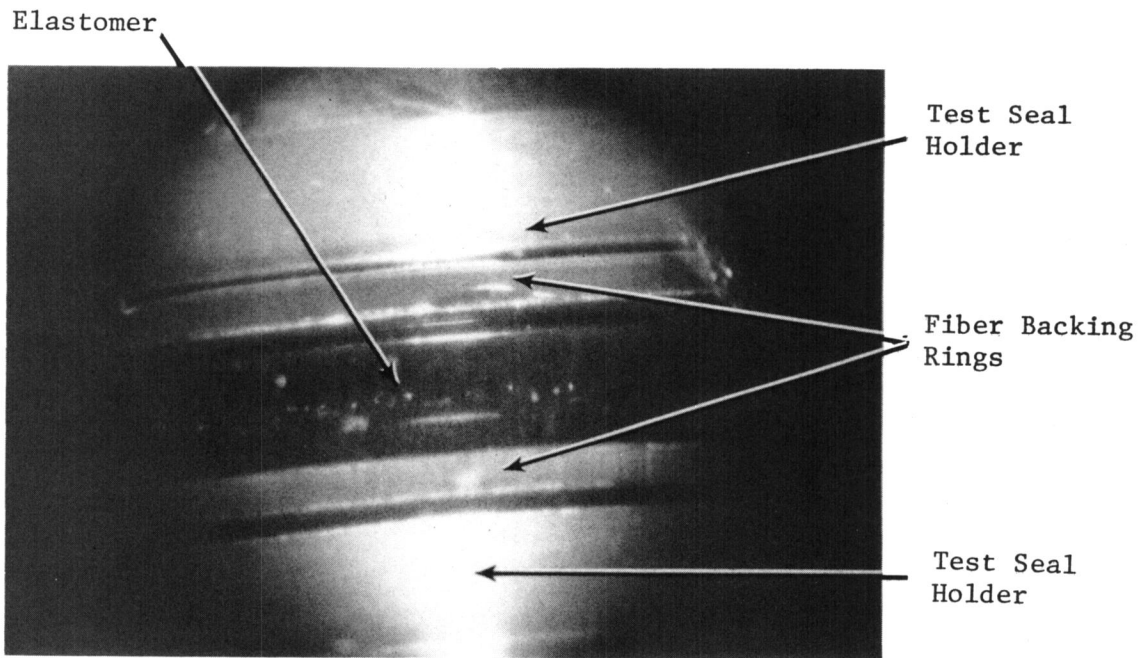


Figure 21 "T" Seal in Contact
With Transparent Cylinder

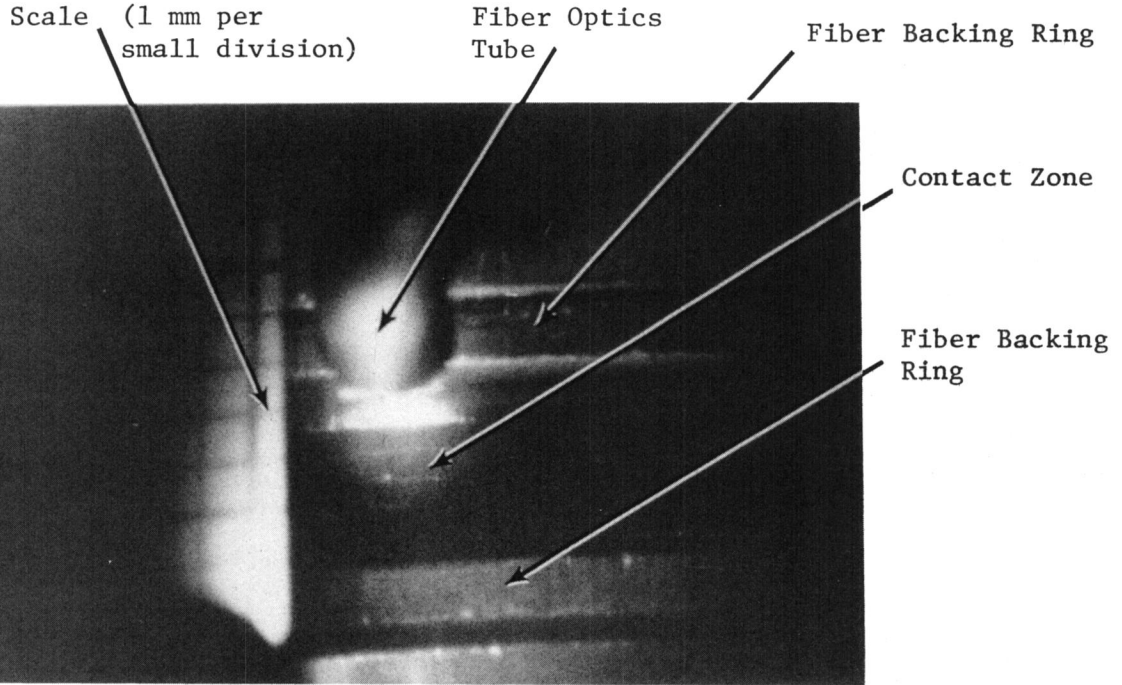


Figure 22 Close-up of "T" Seal in Contact With Transparent Cylinder

The attempt at polishing did indicate that Lexan does not polish well because of its inherent softness. In addition, the material does not maintain its degree of polish well -- again due to its inherent softness.

The above work led to the conclusion that Lexan is unsuitable for use in the transparent cylinder. Two alternative materials are glass and acrylic. At the time of this writing, a glass cylinder is on order and due for installation in the experimental device before the end of March. The manufacturing of the glass, however, is risky and it is not certain that a glass cylinder having the proper dimensions and tolerances can be secured. Nevertheless, the glass cylinder is the most desirable, since the inner and outer surfaces will be optically smooth and will be very scratch resistant.

Acrylic is the second choice--it is desirable because it is relatively easy to polish and relatively scratch resistant. In the event that substantial problems arise in manufacturing the glass cylinder, this material will be used in the apparatus. However, since acrylic is susceptible to scratching, designs are being considered currently in which the acrylic member is expendable and easily replaced.

ANALYTICAL REPORT

In this section, the analytical work completed during the past year is discussed. The analytical work was directed at determining the film thickness distribution at the seal/cylinder interface. This film thickness distribution is important because it controls the net leakage during each cycle of the reciprocating motion. The film thickness distribution also affects the frictional forces produced by the sliding seal.

Determination of the film thickness distribution for a reciprocating elastomeric seal requires that two elements of seal behavior be considered. Firstly, due to the low modulus of the seal, the film thickness and the pressure distributions are coupled. These must be determined together by theories developed in elastohydrodynamic (EHD) lubrication. Secondly, the film thickness at the seal/cylinder interface fluctuates during each reciprocating cycle. This is unlike conventional EHD problems for bearings and cams where steady-state EHD theories provide accurate film thickness predictions. Because of the reciprocating motion, the film thickness is influenced by both squeeze film and sliding effects - an accurate determination of film thickness behavior therefore requires a solution of the time-dependent Reynolds equation coupled with the elasticity equation. However, a reasonable estimate of the average film thickness level in rod seals can be had by using the existing steady-state EHD theories based on some effective velocity over each cycle of the reciprocating motion. Contributions in this category include the earlier analysis by Dowson and Swales [4, 5], and later by Hooke, et. al. [6, 7].

The analytical work given in this section includes two parts. First, an approximate method is presented to estimate the mid-stroke and the end stroke film thickness. In the second part, a time-dependent EHD analysis is developed to determine the temporal behavior of the film thickness distribution during the entire cycle of the reciprocating motion. The nomenclature used for both parts of the analytical work is given in Table 1.

Estimate of the Mid-Stroke and End-Stroke Film Thickness

The variation of film profile in a rod seal during a reciprocating cycle was observed first by Blok and Koens [8] using the interferometric technique. Figure 23 shows the observed film profiles at twelve different stroke positions. At mid-stroke positions, i.e., $\omega t = \frac{\pi}{2}$, $\omega t = \frac{3\pi}{2}$, the sliding velocity is at its maximum and the film profile is similar to those predicted by the steady-state EHD analyses. At the end-stroke positions, i.e., $\omega t = 0$ and $\omega t = \pi$, the sliding velocity is zero, and surfaces are separated only by squeeze-film action. As the shaft travels from mid-stroke to end-stroke position, the minimum film shifts forward from the exit side of the contact to the entrance region. There also seems to be a considerable variation of the average

Table 1
Nomenclature for Analytical Report

a = semi-major axis of the contact ellipse
 b = semi-minor axis of the contact ellipse
 c = an integration constant used in Equation (32)
 c' = an integration constant used in Equation (36)
 c_I' , c_O' = integration constant c' in inlet and exit half of the contact region respectively

$$C_1 = 16 \left(\frac{P_0}{E'} \right)^2 / \left(\frac{h_0}{R} \right)$$

$$C_7 = 16 \left(\frac{P_0}{E'} \right)^2 / \sqrt{48U}$$

d_{1,k}, d_{2,k}, d_{3,k} = derivative coefficients for uneven grids

E₁, E₂ = elastic moduli for bodies 1 and 2

$$\frac{1}{E'} = \frac{1}{2} \left(\frac{1-\nu_1^2}{E_1} + \frac{1-\nu_2^2}{E_2} \right)$$

f = frequency of reciprocating cycle

h = film thickness

h_o = center film thickness

h_{min} = minimum film thickness

$$\bar{h}_o = \left(\frac{h_0}{R} \right) / \sqrt{48U}$$

$\bar{h}_{o,L}$ = \bar{h}_o for the last time interval or last stroke position

$$H = \frac{h}{h_0}$$

$$H^* = \frac{h_{\min}}{R} \left(\frac{E'R}{2\mu_0 u} \right)^{1/2}, \text{ used by Herrebrugh [9]}$$

H_k = H at kth grid

H_{k,L} = H_k for the last stroke position

j = grid designation

k = grid designation, also used as the ellipticity ratio by Hamrock and Dowson [10]

k_o = grid number at the contact center

k_{ff} = grid number at the film termination

$$M = \left(\frac{W}{E'R} \right) \left(\frac{E'R}{2\mu_o u} \right)^{1/2}, \text{ used by Herrebrugh [9]}$$

p = pressure

p_o = pressure at the contact center or maximum Hertzian pressure

$$P = \frac{p}{p_o}$$

P_k = P at kth grid

$Q_{k,j}$ = a kernel function, see Appendix B

R = radius of "0" ring

R_x, R_y = effective radius in the x and y direction respectively

S = stroke of the shaft

t = time

$$T = \frac{u_o t}{b}$$

u = entrainment velocity

u_o = entrainment velocity at mid-stroke = $\frac{\omega S}{4}$

$$U = \frac{\mu_o u_o}{E'R}$$

V = normal approach velocity

W = load per unit circumferential length of "0" ring

W_T = total load

$$\bar{W} = \frac{W_T}{E'R_x^2}$$

x = coordinate in the direction of motion

x_i = inlet coordinate

x_{ff} = coordinate at the film termination

$\bar{x} = \frac{x}{b}$

μ_o = viscosity

ν_1, ν_2 = Poisson's ratio of bodies 1 and 2

$\omega = 2\pi f$

ψ = equations to be solved for P_k , etc., see Equations (37)

ξ = dummy coordinate for x

$\bar{\xi} = \frac{\xi}{b}$

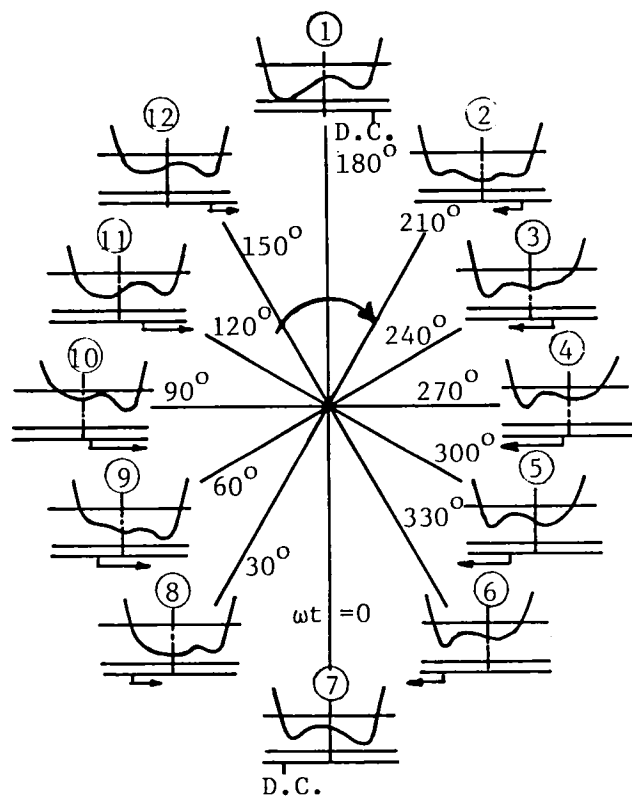


Figure 23 Twelve Successive Phases During the Cycle of a Breathing Lubricant Film

film thickness during the reciprocating cycle. It will be seen later in the detailed time-dependent EHD analysis that these observed film characteristics are completely in accord with the predicted film profiles. It should be noted that if the kinematics are the same during the forward and backward stroke, the film profiles should also have the same characteristics. Figure 24 depicts the qualitative features of the film profiles relative to the four stroke positions in Stirling cycle rod seals.

Approximate Mid-Stroke Film Predictions. Even though results of steady-state EHD analyses are not exactly applicable to a reciprocating rod seal, they are nevertheless suitable for making approximate predictions of the mid-stroke film thickness and the level of film thickness for each cycle. For this purpose, results obtained by Herrebrugh [9] for line contacts as well as those recently developed by Hamrock and Dowson [10] for low modulus point contacts have been adopted.

The numerical solution by Herrebrugh, Figure 25, gives the variation of a dimensionless film parameter H^* with an EHD parameter M . Since the load parameter, W/ER , in rod seals is usually very large, the portion of the curve in Figure 25 applicable to rod seals is towards right in the heavily loaded region. In this region, the curve can be best fitted by

$$H^* = 2.45 M^{-1/4} \quad (1)$$

where

$$H^* = \frac{h_{\min}}{R} \left(\frac{E'R}{2\mu_o u} \right)^{1/2} \quad (2)$$

$$M = \left(\frac{W}{E'R} \right) \left(\frac{E'R}{2\mu_o u} \right)^{1/2} \quad (3)$$

Equation (1) can be rearranged to give

$$\frac{h_{\min}}{R} = 2.45 \left(\frac{2\mu_o u}{E'R} \right)^{5/8} \left(\frac{W}{E'R} \right)^{-1/4} \quad (4)$$

Since in line contacts

$$\frac{W}{E'R} = 2\pi \left(\frac{P_o}{E'} \right)^2 \quad (5)$$

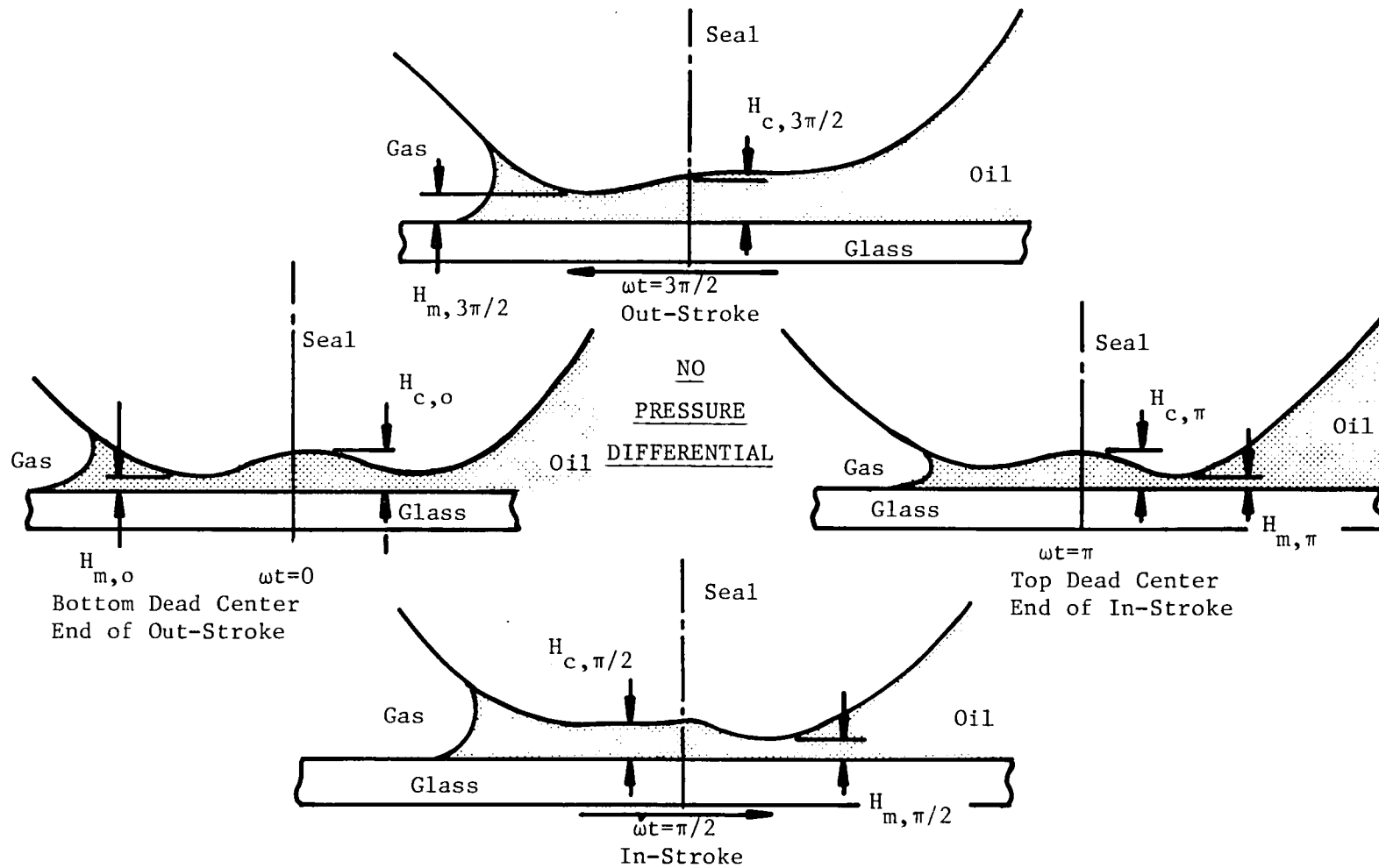


Figure 24 Film Thickness Profiles at Four Piston Positions

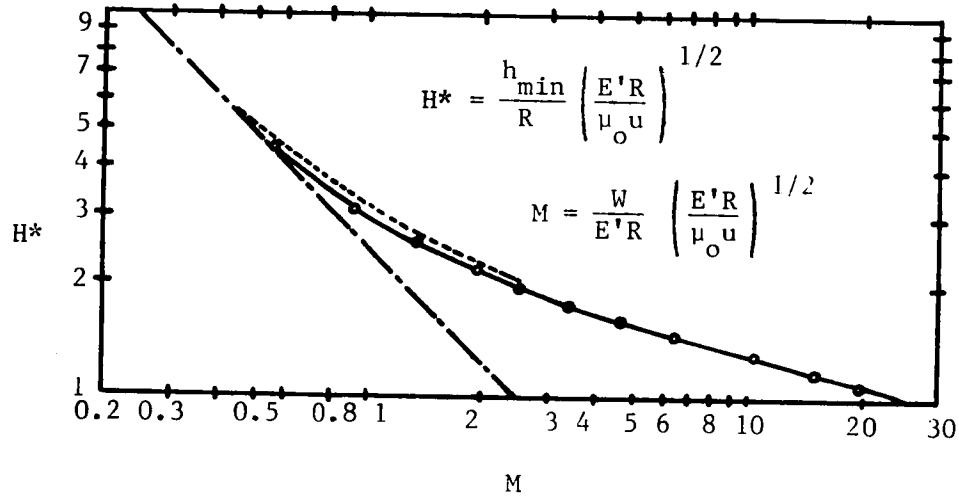


Figure 25 The Functional Relationship: $H^* = F(M)$.
Dash and dot — Gumbel-Martin;
Dashed line — Weber and Saalfeld;
Circled points — author's results.

Equation (4) in terms of a pressure parameter $\frac{p_o}{E'}$, becomes

$$\frac{h_{\min}}{R} = 2.45 \left(\frac{2^{5/8}}{2\pi^{1/4}} \right) \left(\frac{\mu_o u}{E'R} \right)^{5/8} \left(\frac{p_o}{E'} \right)^{-1/2}$$

or

$$\frac{h_{\min}}{R} = 2.385 \left(\frac{\mu_o u}{E'R} \right)^{5/8} \left(\frac{p_o}{E'} \right)^{-1/2} \quad (6)$$

An alternate method can also be used to estimate the minimum film at the mid-stroke by modifying the minimum film thickness formula given recently by Hamrock and Dowson [10] for low-modulus point contacts. Reference [10] gives

$$\frac{h_{\min}}{R_x} = 7.43 (1 - 0.85 e^{-0.31k}) U^{0.65} \bar{W}^{-0.21} \quad (7)$$

where

h_{\min} = minimum film thickness

$$k = \frac{a}{b} = 1.03 \left(\frac{R_y}{R_x} \right)^{0.64}$$

a = semi-major axis of contact ellipse

b = semi-minor axis of contact ellipse

R_x = effective radius in the direction of rolling

R_y = effective radius perpendicular to rolling

$$U = \frac{\mu_o u}{E'R_x}$$

μ_o = viscosity

u = entrainment velocity

$$\frac{1}{E'} = \frac{1}{2} \left(\frac{1-\nu_1^2}{E_1} + \frac{1-\nu_2^2}{E_2} \right)$$

$$\bar{W} = \frac{W_T}{E'R_x^2}$$

W_T = total load of the contact

From Hertz theory for an elliptical contact, it can be shown that

$$\frac{P_o}{E'} = \frac{3}{2\pi} \frac{1}{\left(\frac{a}{b}\right)\left(\frac{b}{R_x}\right)^2} \left(\frac{W_T}{E'R_x^2}\right) \quad (8)$$

or

$$\frac{W_T}{E'R_x^2} = \frac{2\pi}{3} \frac{a}{b} \left(\frac{b}{R_x}\right)^2 \frac{P_o}{E'} \quad (9)$$

Substitution of [9] into [7] gives

$$\frac{h_{\min}}{R_x} = 7.43 (1 - 0.85 e^{-0.31k}) U^{0.65} (k)^{-0.21} \left[\frac{2\pi}{3} \frac{P_o}{E'} \left(\frac{b}{R_x}\right)^2 \right]^{-0.21} \quad (10)$$

It was suggested in Reference [10] that for low-modulus contact Equation (10) can be converted to line contact if k is taken to be 11. Setting $k = 11$ and using the following relation for line contacts,

$$\frac{b}{R_x} = 4 \frac{P_o}{E'} \quad (11)$$

Equation (10), for line contacts, becomes

$$\frac{h_{\min}}{R_x} = 7.43 (1 - 0.85 e^{-0.31 \times 11}) (11)^{-0.21} \left(16 \cdot \frac{2\pi}{3}\right)^{-0.21} U^{0.65} \left(\frac{P_o}{E'}\right)^{-0.63} \quad (12)$$

or

$$\frac{h_{\min}}{R_x} = 2.09 U^{0.65} \left(\frac{P_o}{E'}\right)^{-0.63} \quad (13)$$

Equation (13) is the minimum film for low-modulus line contacts based

on Hamrock and Dowson's point contact theory. Film thickness curves based on Equation (6) and Equation (13) for conditions used in the Stirling engine rod seals are plotted in Figure 26. The agreement between these two theories is seen to be reasonably good at low loads. At higher loads, the larger exponent in Equation (13), deduced from the point contact theory, yields a much lower estimate compared to Herrebrugh's theory. Nevertheless, for pressures encountered in Stirling engine seals $3.5 \times 10^3 - 14 \times 10^3$ kPa (500-2000 psi) the discrepancy is not serious; therefore, both theories appear to be satisfactory.

Herrebrugh's theory was recently correlated experimentally by Swales et. al. with rubber pads [5]. Comparison in Reference [5] showed that the load dependence in the experimental data is considerably smaller than in Herrebrugh's formula. Since the values of p_0/E' in Swales' experiment were considerably smaller than would be encountered in Stirling engine seals, it is believed that film predictions based on Equation (6) or Equation (13) for rod seals are likely to be on the pessimistic side.

End-Stroke Film Thickness. Based on the steady state analysis, the film thickness at both ends of the stroke would be zero because the entrainment velocity is zero. However, due to squeeze-film effects, the end-stroke film thickness will not vanish, but will attain a finite value which depends on the squeeze-film velocity.

To estimate the values of end-stroke film thickness due to squeeze-film effects, the results developed by Herrebrugh [11] for two normally approaching cylinders with an isoviscous lubricant are used. Based on his findings, the center film thickness during a normal approach of two cylinders in the heavily loaded region is

$$\frac{h_o}{R} = 4.85 \left(\frac{W}{E'R} \right)^{1/6} \left(\frac{\mu_o V}{E'R} \right)^{1/3} \quad (14)$$

where

V = normal approaching velocity

It is assumed that the film thickness at the end-stroke does not deviate greatly from the mid-stroke film thickness; i.e.,

$$h_o = h_{o,-/2} \quad (15)$$

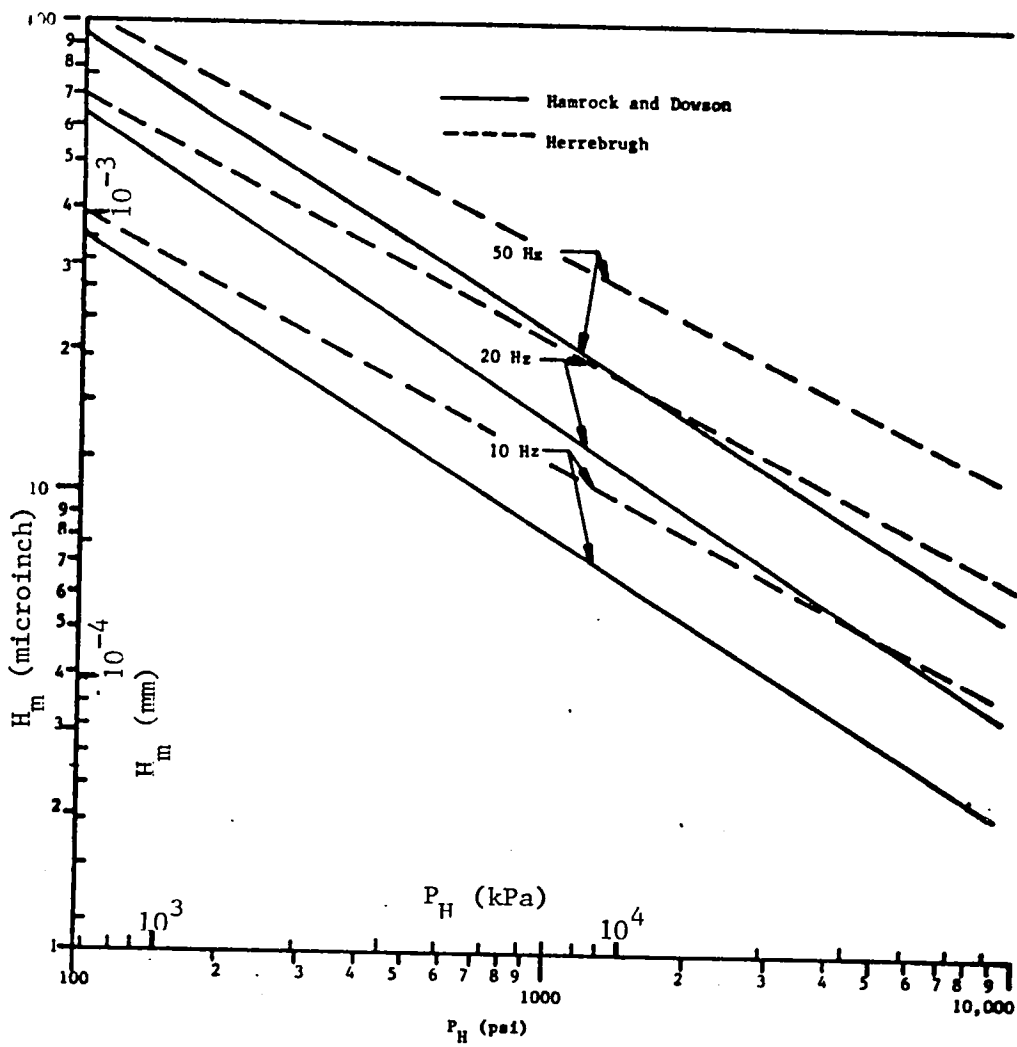


Figure 26 Comparison of Approximate Analytical Film Thickness Results

The normal approaching velocity can then be approximated by

$$V = \frac{\Delta h_o}{\Delta t} = \frac{h_{o,\pi} - h_{o,\pi}}{2 \Delta t} \quad (16)$$

where the time increment can be taken as one quarter of the period; i.e.,

$$t = \frac{1}{4} \left(\frac{1}{f} \right) \quad (17)$$

In Equation (17), f is the frequency of the strokes.

Rearranging of Equation (14) gives

$$\frac{\mu_o V}{E'R} = \left(\frac{h_o}{R} \right)^3 \frac{1}{(4.85)^3} \frac{1}{\left(\frac{W}{E'R} \right)^{1/2}} \quad (18)$$

By substituting $\frac{\Delta h_o}{\Delta t}$ for V one obtains

$$\frac{\Delta h_o}{h_o} = \frac{E'R}{\mu_o} \left(\frac{\Delta t}{h_o} \right) \left(\frac{h_o}{R} \right)^3 \frac{1}{(4.85)^3} \frac{1}{\sqrt{2\pi} \left(\frac{P_o}{E'} \right)}$$

or

$$\frac{\Delta h_o}{h_o} = \left(\frac{E'R}{\mu_o u_o} \right) \left(\frac{u_o \Delta t}{R} \right) \left(\frac{h_o}{R} \right)^2 \frac{1}{285.8 \left(\frac{P_o}{E'} \right)} \quad (19)$$

Here, u_o is the entrainment velocity at the mid-stroke, so that

$$u_o = \left(\omega \frac{S}{2} \right) \frac{1}{2} = \pi f S / 2$$

from which

$$u_o \Delta t = \pi f \frac{S}{2} \cdot \frac{1}{4f} = \frac{\pi S}{8} \quad (20)$$

Equation (6) can now be used to express h_c/R in terms of the speed and load parameter. Substituting Equations (20) and (6) into (19) yields

$$\frac{\Delta h_o}{h_o} = \left(\frac{\pi S}{8R}\right) 2.385 \times U^{0.25} \frac{1}{285.8} \left(\frac{P_o}{E'}\right)^{-2.0} \quad (21)$$

$$= 0.00327 \left(\frac{S}{R}\right) \times U^{0.25} \left(\frac{P_o}{E'}\right)^{-2.0} \quad (22)$$

When the following ranges for $\frac{S}{R}$, U , and $\frac{P_o}{E'}$ are used,

$$\frac{S}{R} = 4 - 12$$

$$U = 2 \times 10^{-6} - 5 \times 10^{-7}$$

$$\frac{P_o}{E'} = 0.05 - 0.2$$

the ratio $\frac{\Delta h_o}{h_o}$ is found to be within the range from 0.59 - 0.0087

This shows that the squeeze-film effect can be quite substantial under conditions in rod seal lubrication. The variation of the center film h_c during the entire cycle of the reciprocating motion is expected to be small for high values of p_o/E' .

Transient Film Thickness Analysis

It was shown in the last section that the mid-stroke film thickness can be estimated by equations derived from existing low-modulus EHD theories. The reduction of film thickness at the end-stroke can be estimated by a relation deduced from Herrebrugh's normal approach EHD analysis. In this section, an analysis is developed to determine the temporal behavior of the film thickness by solving simultaneously the time-dependent Reynolds' equation and the deformation equation. Since it was anticipated that this transient EHD solution for severely loaded rod seal contacts might encounter unforeseen difficulties, it was decided that the first attempt at producing the transient analysis be based on simple Hertzian contact. The analysis below, therefore, does not include the deformation due to contact pressures exerted by groove surfaces on the seal ring. The analytical effort was focused on developing a

numerical method to handle the transient EHD solution under severe pressure loading. Other deformation profiles, including those produced by groove contact pressures, can be incorporated into the perfected method at a later date.

Governing Equations. The time-dependent Reynolds' equation for a pure Hertzian contact is

$$\frac{\partial}{\partial x} \left(h^3 \frac{\partial p}{\partial x} \right) = 12\mu_0 \left(u(t) \frac{\partial h}{\partial x} + \frac{\partial h}{\partial t} \right) \quad (23)$$

The film thickness is governed by

$$h = h_0 + \frac{x^2}{2R} - \frac{4}{\pi E'} \int_{x_i}^{x_{ff}} p(\xi) \ln \left| \frac{\xi - x}{|\xi|} \right| d\xi \quad (24)$$

where

- x = coordinate along the sliding surface
- h = film thickness
- p = pressure
- μ_0 = viscosity
- u = time varying velocity
- h_0 = center film at $x = 0$
- R = seal radius
- x_i = inlet boundary of the lubricant film
- x_{ff} = exit boundary of the lubricant film
- ξ = dummy variable for x

Introducing the following non-dimensional quantities:

- \bar{x} = $\frac{x}{b}$
- b = half contact width
- H = $\frac{h}{h_0}$
- P = p/p_0
- p_0 = maximum Hertzian pressure

$$\bar{h}_o = \frac{(h_o/R)}{\sqrt{48 \frac{\mu_o u_o}{E'R}}}$$

u_o = mid-stroke entrainment velocity, half of the mid-stroke sliding velocity ($u_o = \omega S/4$)

$$C_1 = 16(p_o/E')^2/(h_o/R)$$

S = stroke of the piston

$$T = \frac{u_o t}{b}$$

ω = angular velocity = $2\pi f$

$$\bar{\xi} = \xi/b$$

Equation (23) and Equation (24) in non-dimensional form become

$$\frac{\partial}{\partial \bar{x}} \left(H^3 \frac{\partial P}{\partial \bar{x}} \right) = \frac{1}{\bar{h}_o^2} \left[\sin \omega t \frac{\partial H}{\partial \bar{x}} + \frac{1}{\bar{h}_o} \frac{\partial (\bar{h}_o H)}{\partial T} \right] \quad (25)$$

$$H = 1 + C_1 \left[\frac{-2}{2} - \frac{1}{\pi} \int_{\bar{x}_i}^{\bar{x}_{ff}} P(\bar{\xi}) \ln \frac{|\bar{\xi} - \bar{x}|}{|\bar{\xi}|} d\bar{\xi} \right] \quad (26)$$

In the above, the shaft velocity is assumed to be sinusoidal. The boundary conditions for Equation (25) are:

$$P = 0 \text{ at } \bar{x} = \bar{x}_i \quad (27)$$

$$P = 0 \text{ and } \frac{dP}{d\bar{x}} = 0 \text{ at } \bar{x} = \bar{x}_{ff} \quad (28)$$

$$P = 1 \text{ at } \bar{x} = 0 \quad (29)$$

$$\left(\frac{h_o}{R} \right)_{\omega t = 0} = \left(\frac{h_o}{R} \right)_{\omega t = 2\pi} \quad (30)$$

For a given set of parameters $\frac{p_o}{E'}$, $\frac{\mu_o u_o}{E'R}$, and $\frac{S}{R}$, one can determine the center film thickness $\frac{h_o}{R}$ and the film thickness and pressure profiles at each time interval using numerical schemes. The two numerical

schemes considered in solving Equation (25) and (26) and the experience gathered with each are documented below.

Numerical Treatment--Direct Discretization (Scheme "A"). This scheme is the conventional way of handling the Reynolds' equation. In Equation (25), the second order derivative of P is directly discretized at a grid point k. The discretized equation appears as

$$\begin{aligned} \psi_k = & \left[-H_{k-1/2}^3 \frac{P_k - P_{k-1}}{\Delta \bar{x}_{k-1}} + H_{k+1/2}^3 \frac{P_{k+1} - P_k}{\Delta \bar{x}_k} \right] \bar{h}_o^2 \\ & - \sin \omega t \left[d_{3,k} H_{k+1} + d_{2,k} H_k - d_{1,k} H_{k-1} \right] \\ & - \frac{1}{\Delta T} \left[H_k - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{k,L} \right] = 0 \end{aligned} \quad (31)$$

where

$$H_k = 1 + C_1 \left[\frac{\bar{x}_k^2}{2} - \frac{1}{\pi} \sum_{j=1}^{k_{ff}} P_j Q_{kj} \right]$$

$$C_1 = 16(P_o/E')^2 / \left(\frac{h_o}{R} \right)$$

\bar{x}_h = \bar{x} coordinate at grid k

P_j = P at grid j

Q_{kj} = A Kernel function developed in Reference [10], and reproduced in Appendix B

$$H_{k-1/2} = (H_k + H_{k-1}) / 2$$

$$H_{k+1/2} = (H_k + H_{k+1}) / 2$$

$$\Delta \bar{x}_{k-1} = \bar{x}_k - \bar{x}_{k-1}$$

$$\bar{h}_o = (h_o/R) / \left(48 \frac{\mu_o u_o}{E'R} \right)^{1/2}$$

$$d_{1,k} = \frac{\Delta \bar{x}_k}{\Delta \bar{x}_{k-1} (\Delta \bar{x}_k + \Delta \bar{x}_{k-1})}$$

$$d_{2,k} = \frac{\Delta \bar{x}_k - \Delta \bar{x}_{k-1}}{\Delta \bar{x}_k \Delta \bar{x}_{k-1}}$$

$$d_{3,k} = \frac{\Delta \bar{x}_{k-1}}{\Delta \bar{x}_k (\Delta \bar{x}_k + \Delta \bar{x}_{k-1})}$$

$$\bar{h}_{o,L} = \bar{h}_o \text{ for the previous time interval}$$

$$H_{k,L} = H_k \text{ for the previous time interval}$$

$$\Delta T = \frac{u_o \Delta t}{b}$$

At each time step, Equation (31) is a set of equations written at grids 2 to $k_{ff}-1$ for a total of $k_{ff}-2$ equations. These $k_{ff}-2$ equations can be solved by the Newton-Raphson procedure to obtain values for $k_{ff}-2$ unknowns consisting of P_2 to P_{k_o-1} , P_{k_o+1} to $P_{k_{ff}-1}$, and \bar{h}_o . In this solution,

$$P_1 = 0$$

$$P_{k_o} = 1$$

$$P_{k_{ff}} = 0$$

which are the boundary conditions for Equation (31).

The location of $\bar{x}_{k_{ff}}$ is not fixed, and must be such to make the slope $\frac{dP}{d\bar{x}}$ vanish at $\bar{x}_{k_{ff}}$. This is accomplished by moving $\bar{x}_{k_{ff}}$ towards the contact center one grid point at each iteration whenever a negative pressure appears near k_{ff} .

Scheme "A" was tried and failed to yield satisfactory results particularly at high loads. It begins to experience difficulties at

$\frac{P_0}{E'} = 0.01$. It is not known whether the difficulties arose because of the inherent discretization error associated with the interpolations factors $d_{1,k}$, $d_{2,k}$, $d_{3,k}$ for uneven grids or because of a possible error in the complicated algebraic expression of $\frac{\partial \psi_k}{\partial \bar{h}_0}$ which has to be

evaluated in the Newton-Raphson procedure. Scheme "A" was finally abandoned in favor of an alternative approach, Scheme "B", which is described next.

Numerical Treatment--Discretization of Integrated Reynolds Equation (Scheme "B"). This scheme follows closely that developed previously by Lee and Cheng [12] for studying the effect of a single asperity entering the inlet zone of an EHD contact. In this formulation, the Reynolds equation is first integrated once with respect to \bar{x} , and then discretized.

Integrating Equation (25) once, one obtains

$$\bar{h}_0^2 H^3 \frac{\partial P}{\partial \bar{x}} = \sin \omega t \cdot H + \frac{1}{\bar{h}_0} \int_{\bar{x}_i}^{\bar{x}} \frac{\partial (\bar{h}_0 H)}{\partial T} d\bar{\xi} + c \quad (32)$$

where c is an integration constant.

If $\frac{dP}{d\bar{x}}$ is zero at $\bar{x} = 0$, then

$$c = -\sin \omega t - \frac{1}{\bar{h}_0} \int_{\bar{x}_i}^0 \frac{\partial (\bar{h}_0 H)}{\partial T} d\bar{\xi} \quad (33)$$

For general cases, $\frac{dP}{d\bar{x}}$ at $\bar{x} = 0$ is small but not exactly zero; therefore, a small residual constant c' must be provided to permit any small non-zero pressure gradient to exist at $\bar{x} = 0$. Thus, for general cases,

$$c = c' - \sin \omega t - \frac{1}{\bar{h}_0} \int_{\bar{x}_i}^0 \frac{\partial (\bar{h}_0 H)}{\partial T} d\bar{\xi} \quad (34)$$

Substituting (34) in (32), and observing that

$$\int_{\bar{x}_i}^0 \frac{\partial(\bar{h}_o H)}{\partial T} d\bar{\xi} = \int_{\bar{x}_i}^{\bar{x}} \frac{\partial(\bar{h}_o H)}{\partial T} d\bar{\xi} + \int_{\bar{x}}^0 \frac{\partial(\bar{h}_o H)}{\partial T} d\bar{\xi} \quad (35)$$

Equation (32) becomes

$$\bar{h}_o^2 H \frac{\partial P}{\partial \bar{x}} = \sin \omega t (H - 1) - \frac{1}{\bar{h}_o} \int_{\bar{x}}^0 \frac{\partial(\bar{h}_o H)}{\partial T} d\bar{\xi} + c' \quad (36)$$

Equation (36) is discretized at the mid-point, $k + 1/2$, between two pressure nodes, P_k and P_{k+1} . The discretized equation is

$$\psi_{k+1/2} = \bar{h}_o^2 \left(H_{k+1/2}^3 \frac{P_{k+1} - P_k}{\Delta \bar{x}_k} \right) - \sin \omega t (H_{k+1/2}^{-1}) + \frac{1}{\Delta T} \text{SUMS}(K, TH, 0) + c' = 0 \quad (37)$$

where

$$\begin{aligned} \text{SUMS}(K, TH, 0) &= \frac{\Delta \bar{x}_k}{4} \left[(H_{k+1/2} - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{k+1/2,L}) + (H_{k+1} - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{k+1,L}) \right] \\ &+ \sum_{\ell = k+1}^{k_o - 1} \frac{\Delta \bar{x}_\ell}{2} \left[(H_\ell - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{\ell,L}) + (H_{\ell+1} - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{\ell+1,L}) \right] \end{aligned} \quad (38)$$

For $k < k_o - 1$

$$\text{SUMS}(K, TH, 0) = \frac{\Delta \bar{x}_{k\phi-1}}{4} \left[(H_{k\phi-1/2} - \frac{\bar{h}_{o,L}}{\bar{h}_o} H_{k\phi-1/2,L}) + (1 - \frac{\bar{h}_{o,L}}{\bar{h}_o}) \right] \quad (39)$$

For $k = k_o - 1$

$$\text{SUMS (K, TH, 0)} = \frac{\Delta \bar{x}_{k\phi}}{4} \left[\left(1 - \frac{\bar{h}_{0,L}}{\bar{h}_0}\right) + \left(H_{k\phi+1/2} - \frac{\bar{h}_{0,L}}{\bar{h}_0} H_{k\phi+1/2,L}\right) \right]$$

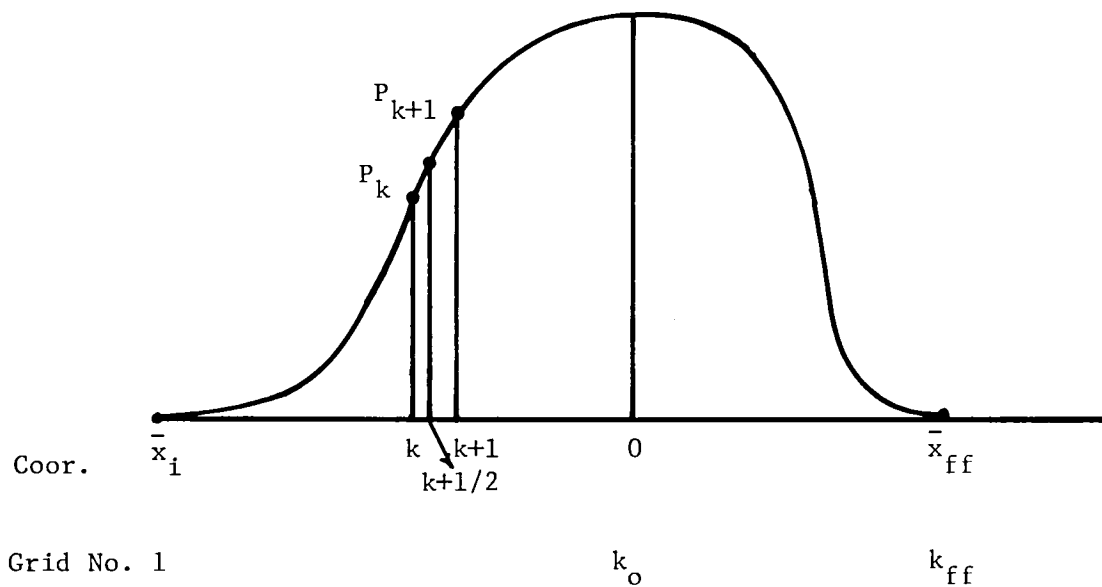
For $k = k_0$ (40)

$$\text{SUMS (K, TH, 0)} = \sum_{\ell=k_0}^k \frac{\Delta \bar{x}_\ell}{2} \left[\left(H_\ell - \frac{\bar{h}_{0,L}}{\bar{h}_0} H_{\ell,L}\right) + \left(H_{\ell+1} - \frac{\bar{h}_{0,L}}{\bar{h}_0} H_{\ell+1,L}\right) \right]$$

$$+ \frac{\Delta \bar{x}_k}{4} \left[\left(H_k - \frac{\bar{h}_{0,L}}{\bar{h}_0} H_{k,L}\right) + \left(H_{k+1/2} - \frac{\bar{h}_{0,L}}{\bar{h}_0} H_{k+1/2,L}\right) \right]$$

For $k \geq k_0 + 1$ (41)

By referring to the following diagram,



Equations (37) can be written at $k = 1$ to $k = k_{ff} - 1$ for a total number of $k_{ff} - 1$ equations. In solving Equations (37), the following

pressures are known,

$$P_1 = 0$$

$$P_{k_o} = 1$$

$$P_{k_{ff}} = 0$$

Thus, the total number of unknown pressure nodes is $k_{ff} - 3$, and they are P_2 to P_{k_o-1} and P_{k_o+1} to $P_{k_{ff}-1}$. These $k_{ff}-3$ unknown pressure nodes together with \bar{h}_o and c' form the right number of unknowns for the $k_{ff}-1$ equations.

As mentioned in Scheme "A", \bar{x}_{ff} must satisfy the free boundary condition $\frac{dP}{d\bar{x}} = 0$. This is accomplished at each iteration by searching for the point near the exit region where the following term vanishes

$$\sin\omega t (H - 1) - \frac{1}{\bar{h}_o} \int_{\bar{x}_{ff}}^0 \frac{\partial(\bar{h}_o H)}{\partial T} d\bar{\xi} + c' = 0 \quad (42)$$

It was found the inclusion of \bar{h}_o as an unknown in Equations (37) in the Newton-Raphson procedure gave rise to some difficulties in convergence. This was remedied by allowing the integration constant c' in the inlet half to be different from the value of c' in the exit half. In so doing, c' is now replaced with two unknown constants, c'_I for the inlet region, and c'_O for the exit region; however, \bar{h}_o is now considered as known. Thus, the total number of unknowns still matches the total number of equations.

At each time step, one begins the first iteration with an initial guess of \bar{h}_o , which is usually the value of \bar{h}_o for the previous time step. With the first guess of \bar{h}_o , Equations (37) can be solved for P_k and c'_I and c'_O . If c'_I is not equal to c'_O , \bar{h}_o is readjusted until c'_I matches c'_O . This procedure permits \bar{h}_o values to be calculated during half of the reciprocating cycle. However, \bar{h}_o thus determined in the first half cycle may not satisfy the periodic conditions, i.e.,

$$(\bar{h}_o)_{\omega t = \pi/2} = (\bar{h}_o)_{\omega t = 3\pi/2}$$

The cycle must be repeated until the periodic boundary condition in time is satisfied.

The procedure seems to work well for a wide range of load and speed parameters, even for conditions beyond those occurring in Stirling engine rod seals. Results obtained from use of the procedure are given in the next section. In the Appendices, details and use of the procedure are discussed. Specifically, the expressions for the derivatives of $\psi_{k+1/2}$ with respect to P_j for solving Equations (37) can be found in Appendix C. The computer program for Scheme "B" is entitled "RODSLJ" and is documented in Appendix A.

Results

Three input parameters are required for each run. These include:

$$\frac{P_o}{E'} : \text{load parameter}$$

$$U = \frac{\mu_o u_o}{E'R} : \text{speed parameter}$$

$$\frac{S}{R} : \text{stroke to seal radius ratio}$$

For the conditions to be used in the experimental rig, the range of the above dimensionless parameters was found to be:

$$\frac{P_o}{E'} = 0.05 \text{ to } 0.25$$

$$U = 10^{-7} \text{ to } 2.0 \times 10^{-6}$$

$$\frac{S}{R} = 4 \text{ to } 16$$

Table 2 lists the input data for seven runs. Runs 50, 51, and 52 are for finding the effect of speed; Runs 51, 53, and 54 for the effect of load; and Runs 51, 55, and 56 for the effect of stroke to radius ratio. Figures 27 to 34 show dimensionless film profiles at six different stroke positions during the forward stroke ($\omega t = 0$ to $\omega t = \pi$) of the reciprocating cycle. Since the sealing pressure differential is not considered in the present analysis, the film profiles during the backward stroke ($\omega t = \pi$ to $\omega t = 2\pi$) should be a mirror image of the forward stroke profiles, which are the only ones presented here. All film

Table 2 List of Input Data

Run No.	$\frac{P_o}{E'}$	U	$\frac{S}{R}$	$\frac{P_o}{E'} \sqrt{U}$
50	0.1	10^{-7}	8	31.6
51	0.1	5×10^{-7}	8	14.14
52	0.1	2×10^{-6}	8	7.07
53	0.25	5×10^{-7}	8	88.3
54	0.05	5×10^{-7}	8	3.54
55	0.1	5×10^{-7}	16	14.14
56	0.1	5×10^{-7}	4	14.14

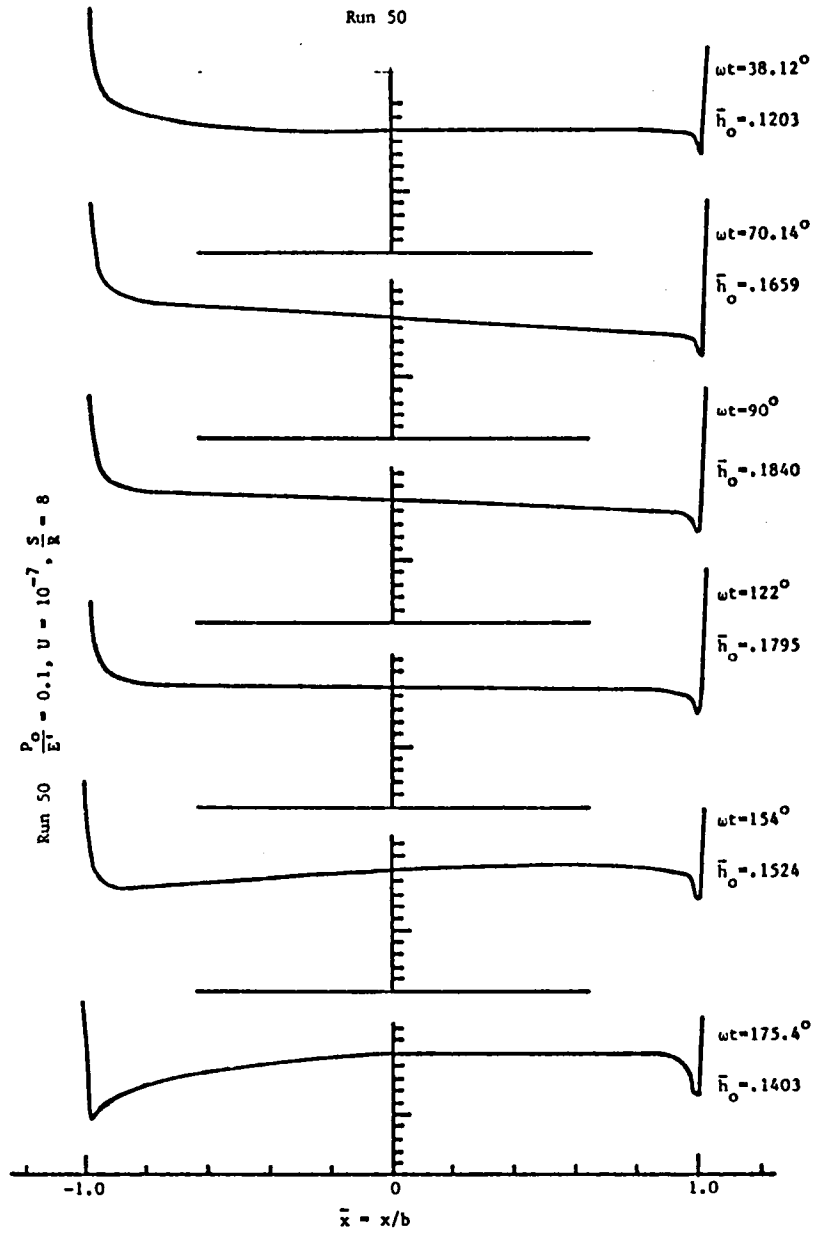


Figure 27 Film Profiles for Run 50, $\frac{P_0}{E^*} = .1, U = 10^{-7}, \frac{S}{R} = 8$

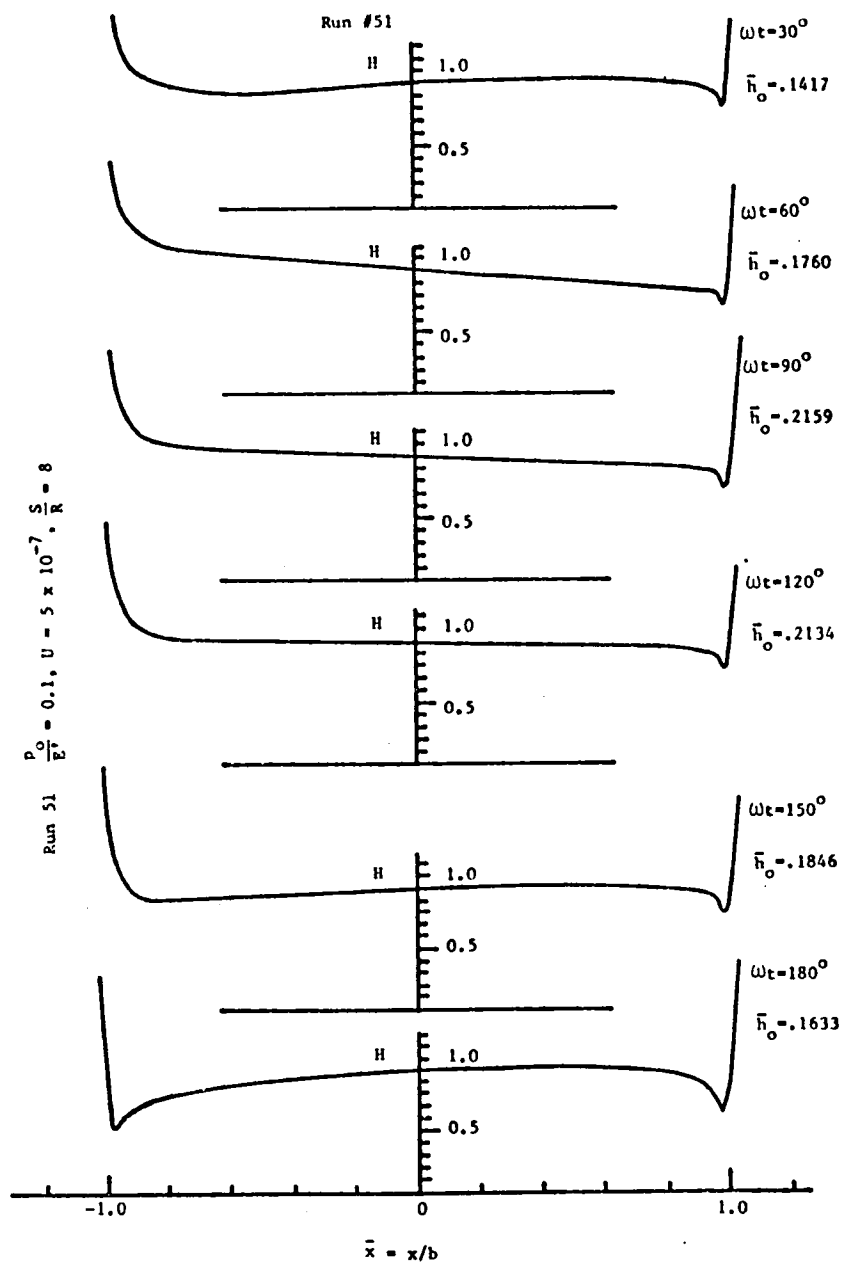


Figure 28 Film Profiles for Run 51, $\frac{P_0}{E^*} = 0.1, U = 5 \times 10^{-7}, \frac{S}{R} = 8$

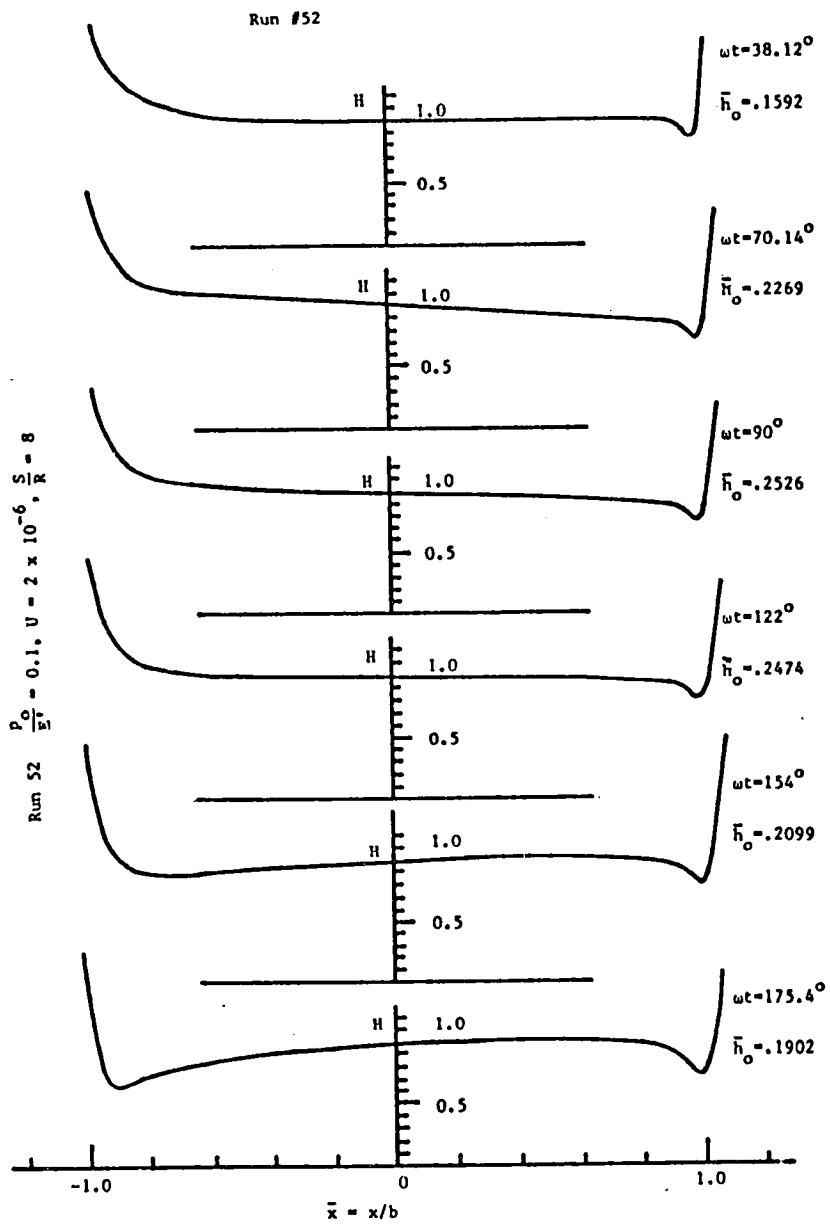


Figure 29 Film Profiles for Run 52, $\frac{P_0}{E\Gamma} = 0.1, U = 2 \times 10^{-6}, \frac{S}{R} = 8$

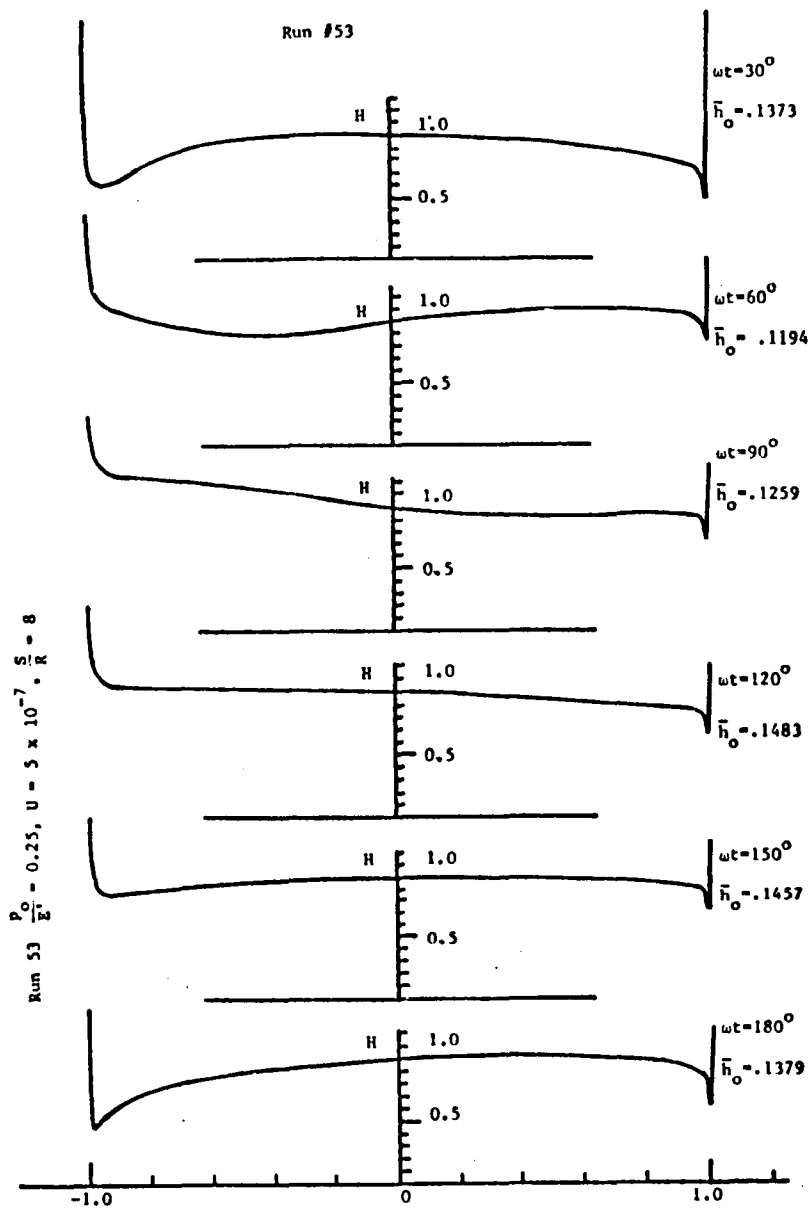


Figure 30 Film Profiles for Run 53, $\frac{P_0}{E^2} = 0.25$, $U = 5 \times 10^{-7}$, $\frac{S}{R} = 8$

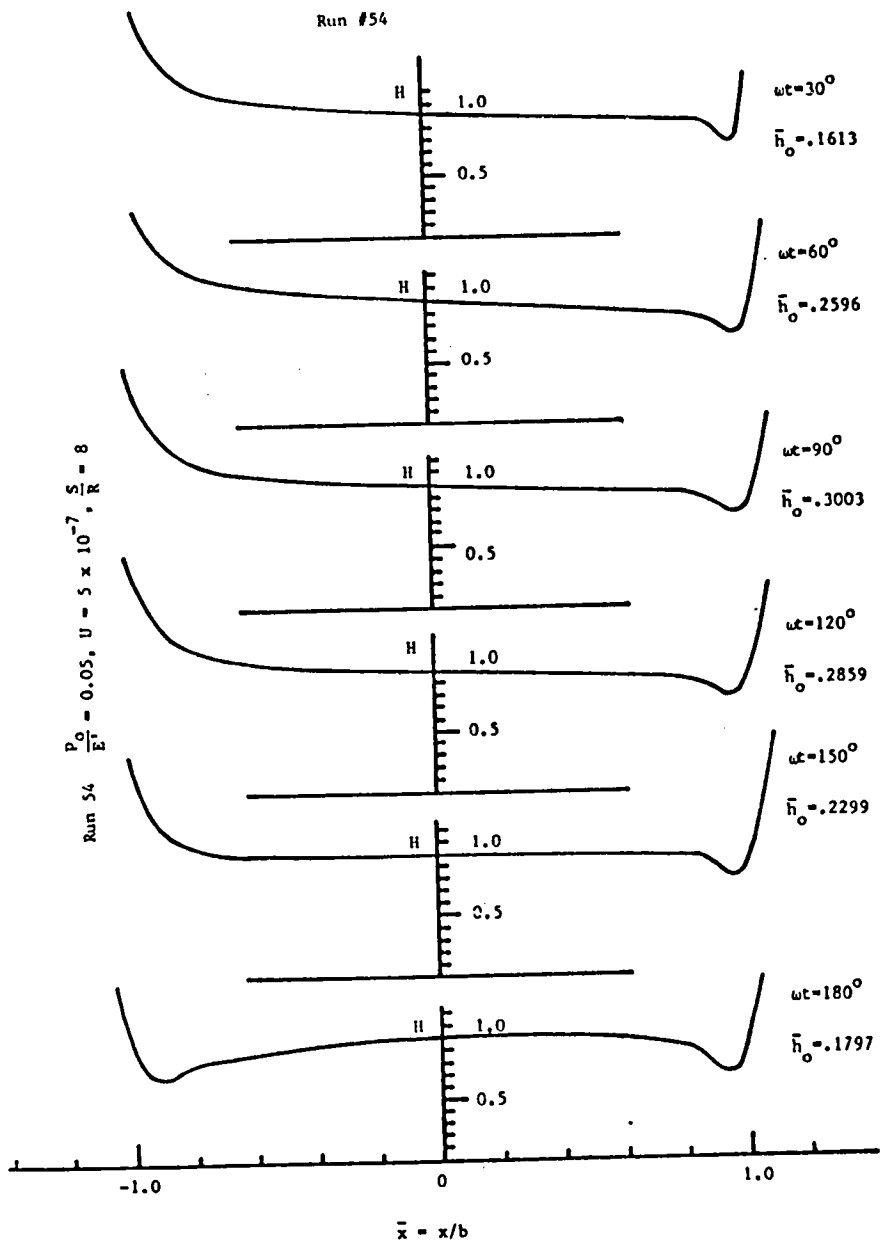


Figure 31 Film Profiles for Run 54, $\frac{P_0}{E^*} = 0.05, U = 5 \times 10^{-7}, \frac{S}{R} = 8$

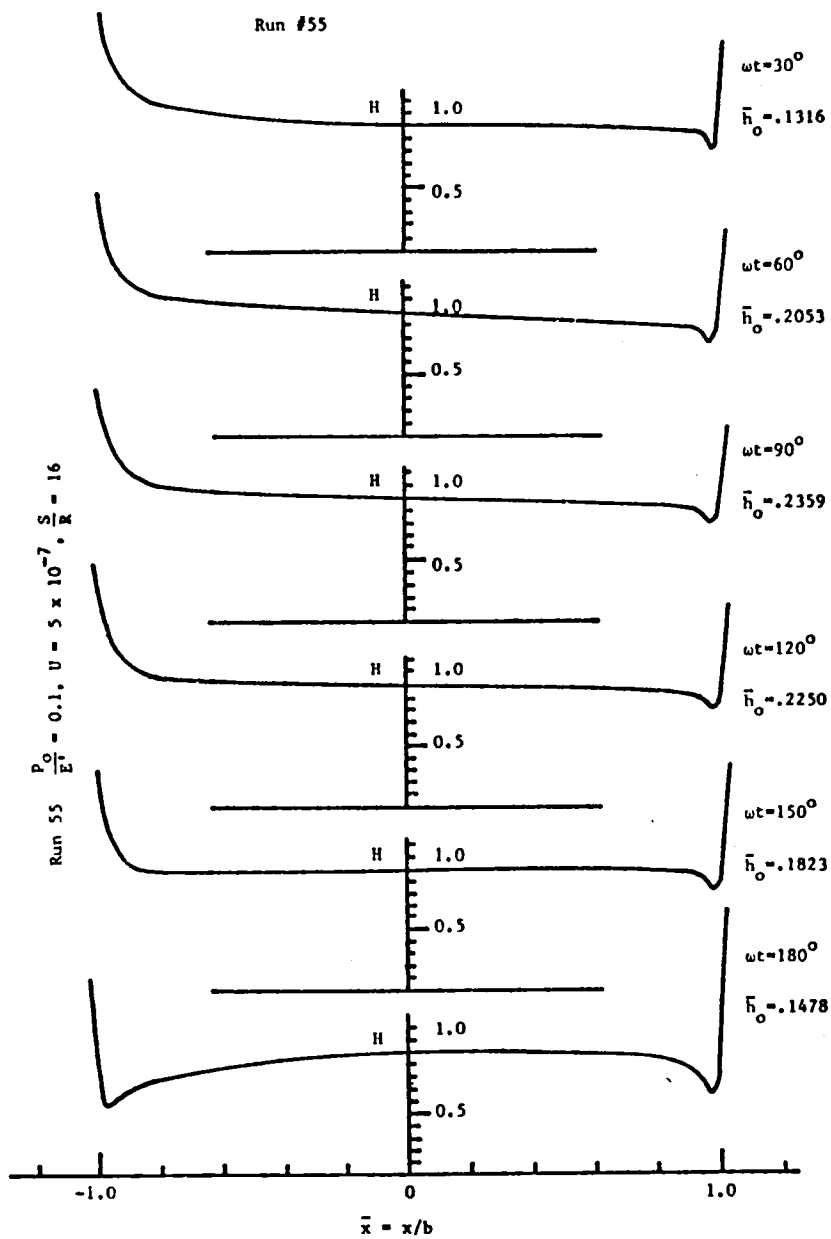


Figure 32 Film Profiles for Run 55, $\frac{P_0}{E^*} = 0.1, U = 5 \times 10^{-7}, \frac{S}{R} = 16$

Run #56A Fine Grid RODA30 RODSLJ

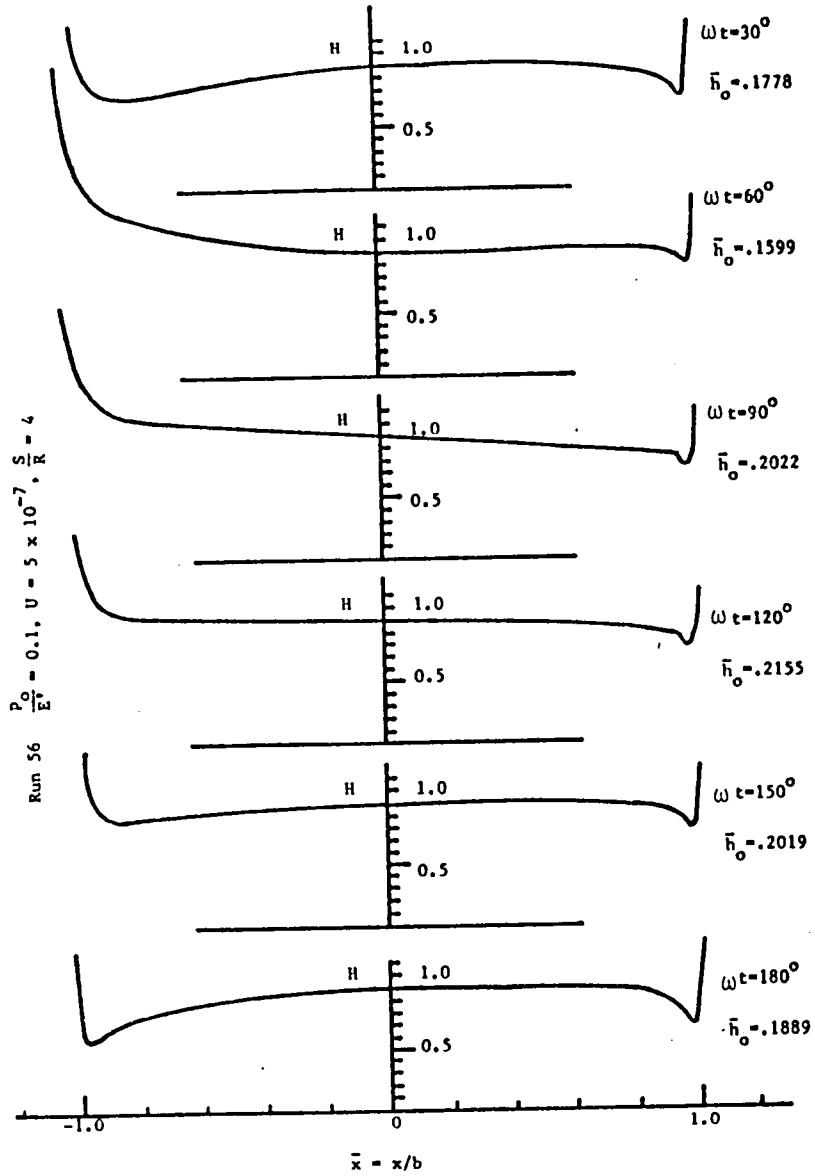


Figure 33 Film Profiles for Run 56, $\frac{P_0}{E_1} = .1, U = 5 \times 10^{-7}, \frac{S}{R} = 4$

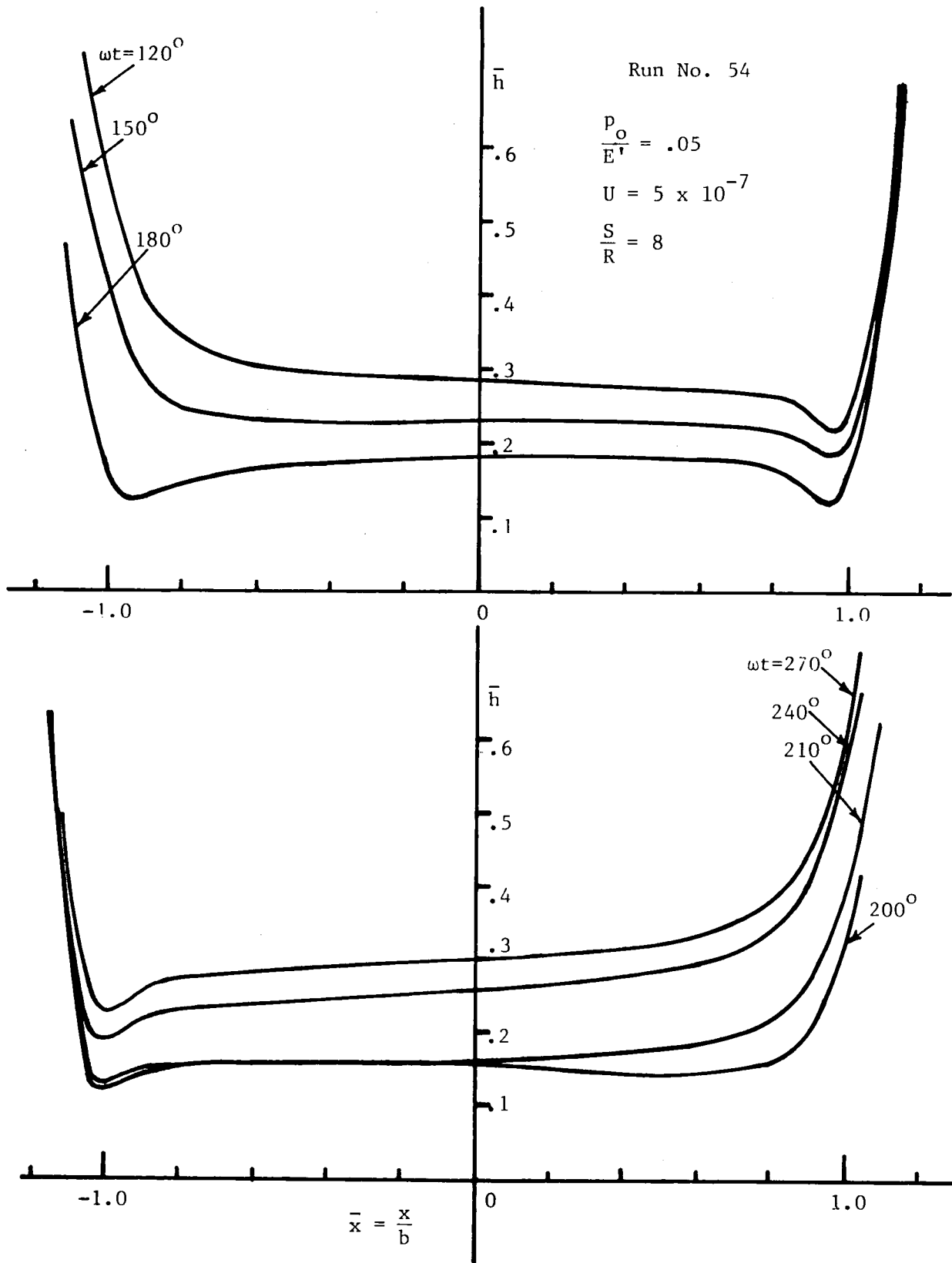


Figure 34 Film Profiles for Run No. 34

profiles are normalized with respect to the film thickness at the center. The value of \bar{h}_0 , the non-dimensional film thickness at the center, at each stroke position is labelled for each profile.

For Run 54, the most lightly loaded case, the film profile between the last half of the forward stroke ($\omega t = \frac{\pi}{2}$ to $\omega t = \pi$) and the first half of the backward stroke ($\omega t = \pi$ to $\omega t = \frac{3\pi}{2}$) are replotted on the same scale to show their relative positions. These curves show a trend consistent with those observed by Blok and Koens [8]. It would be of interest to ascertain the exact conditions used in their experimental runs so that more quantitative comparisons can be made.

Aside from the fact that the edge protrusions are much more pronounced for heavily loaded or low speed cases, the general features of film profiles are very similar for all runs. It is interesting to note that at the end-stroke position ($\omega t = \pi$), where the sliding velocity vanishes, the film droops down more in the inlet half region, whereas in the exit half the film remains reasonably flat or even rises slightly above the center film in some cases. This feature seems to be due to the sliding action producing a very effective suction once a negative wedge is formed at the inlet. As can be seen in Figure 23, this same feature also exists in the observed film profiles at the end-strokes.

The variation of the center film thickness during half of the reciprocating cycle ($\omega t = \frac{\pi}{2}$ to $\omega t = \frac{3\pi}{2}$) is shown in Figure 35. At the mid-stroke ($\omega t = \frac{\pi}{2}$), the shaft velocity reaches a maximum and is decreasing. However, the center film thickness does not reach a maximum until about 10-15 degrees after the mid-stroke. This shift is due to the fact that part of the sliding action at the mid-stroke is nullified by the squeeze-film action $\frac{\partial H}{\partial T}$. The stronger the squeeze-film term, the larger the shift of the maximum film from the mid-stroke position. As the shaft reaches the end-stroke position ($\omega t = \pi$), the sliding action ceases, and the pressure within the film is supported entirely by the squeeze-film action. Perhaps the most interesting feature in this curve is the sudden reduction of \bar{h}_0 soon after the end-stroke position, where the motion reversal takes place. This reduction is attributable to the penetration of the exit protrusion into the center of the contact after the shaft reverses its direction. This penetration causes a sharp reduction \bar{h}_0 for low-speed or high-pressure cases. It becomes more pronounced as the squeeze-film effect becomes more influential.

Figure 36 shows influence of speed on \bar{h}_0 . It is seen that the squeeze-

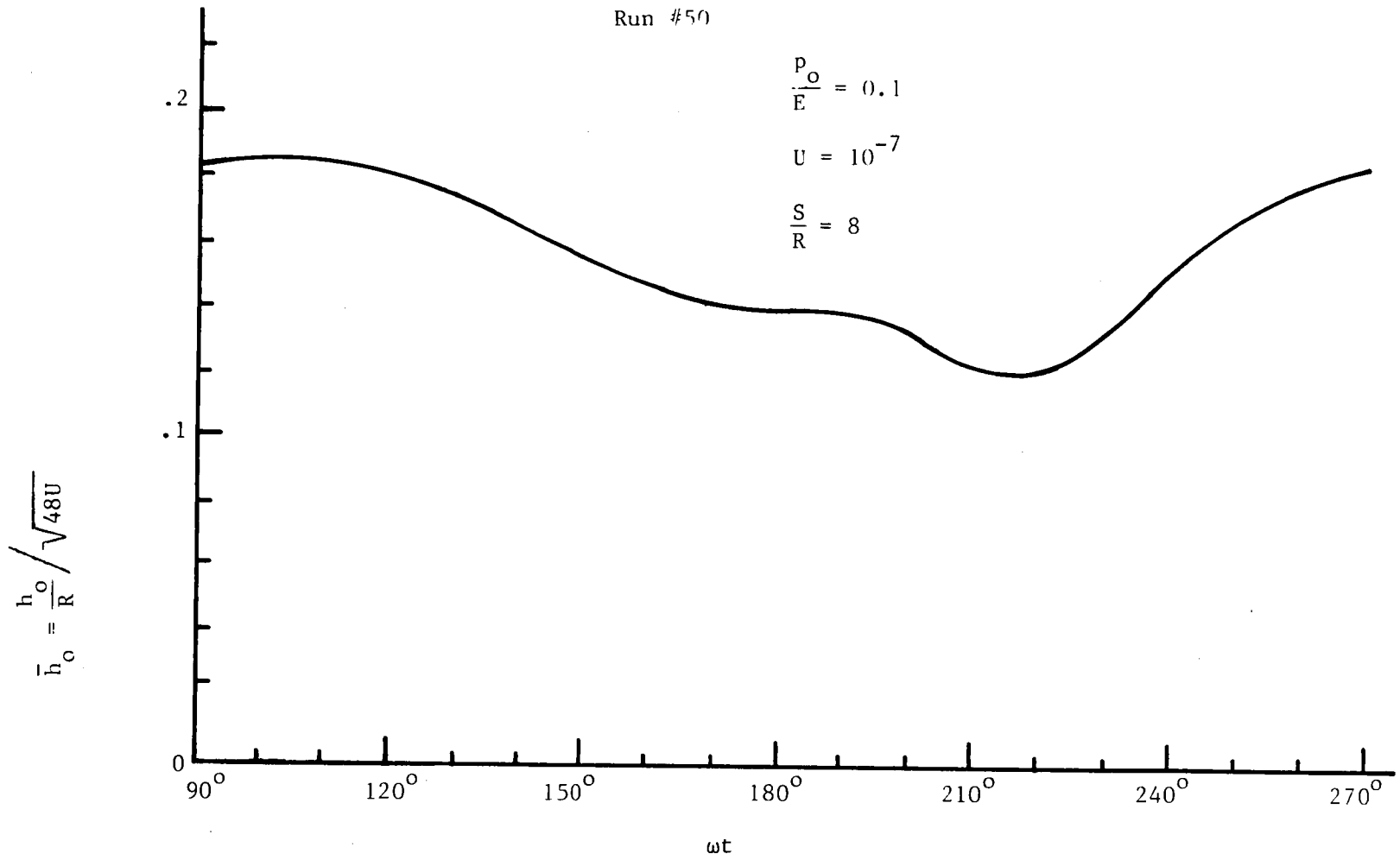


Figure 35 Variation of \bar{h}_0 with the Stroke Position

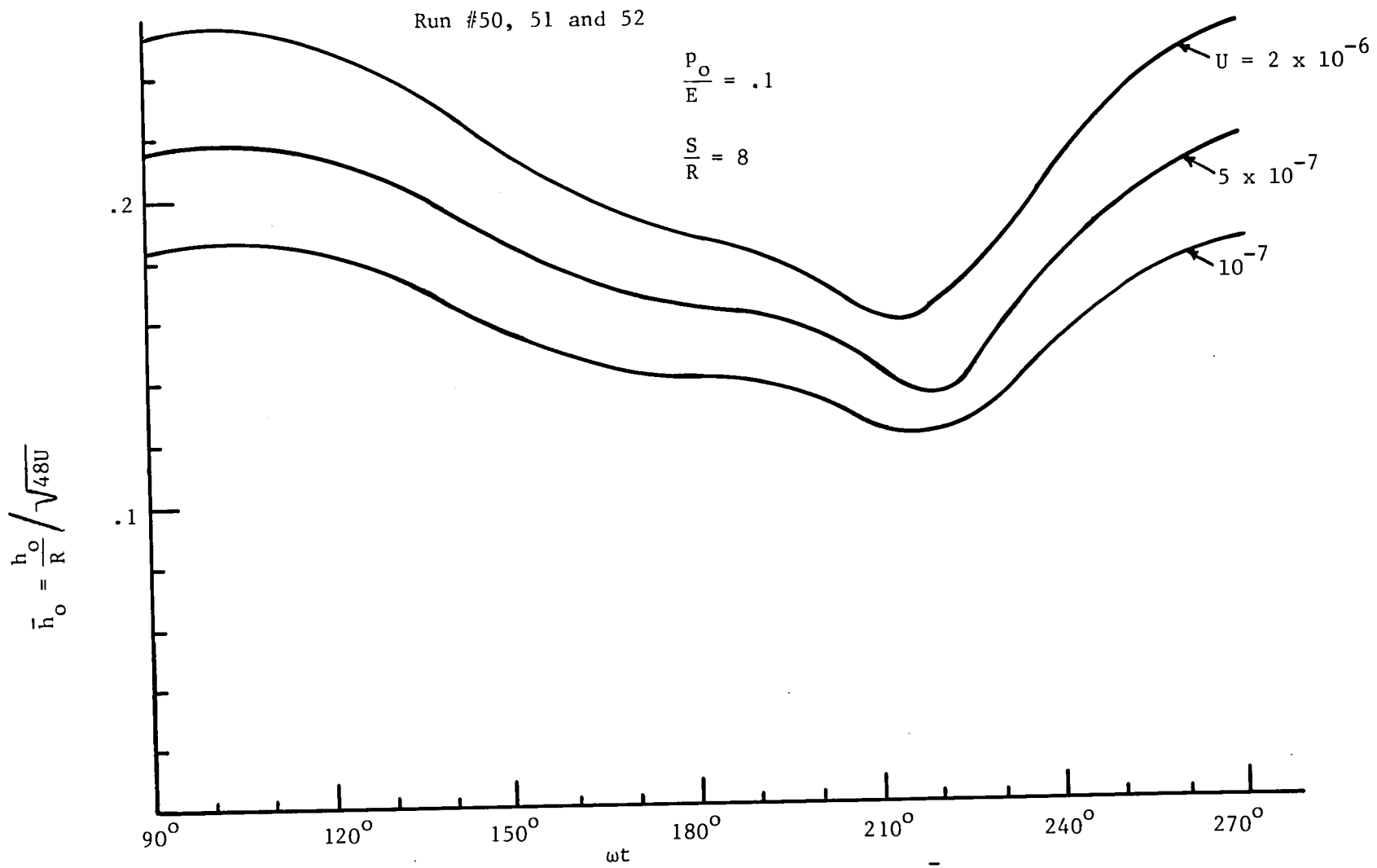


Figure 36 Effect of Speed on the Variation of \bar{h}_0

film action is not altered significantly even for a five-fold increase in speed. This trend agrees well with the approximate analysis of squeeze-film effect represented by Equation (20).

Figure 37 shows the effect of load parameter on \bar{h}_0 . It is seen that the fluctuation of \bar{h}_0 is most strongly influenced by the change in pressure. A five-fold increase in Hertzian pressure from 0.05 to 0.25 yields more than ten times reduction of \bar{h}_0 during the last half of the forward cycle ($\omega t = \frac{\pi}{2}$ to $\omega t = \pi$). This result again confirms qualitatively the approximate formula Equation (20) derived from Herrebrugh's squeeze-film analysis. It should be noted that for the highest Hertzian pressure $\left(\frac{P_0}{E'} = .25\right)$ the reduction of \bar{h}_0 right after the end-stroke position ($\omega t = \pi$) is quite sudden. It should also be noted that the lowest film thickness at this position is accompanied by a substantial sliding velocity. If the seal surface is to fail by sliding wear, it is likely to begin first at this position.

Figure 38 shows the effect of stroke length on the variations of \bar{h}_0 . As expected, a longer stroke would lessen the squeeze-film effect and allow \bar{h}_0 to approach more to the steady-state value based solely on the sliding velocity. With a shorter stroke, the squeeze-film action dominates, and the fluctuation due to the sinusoidal sliding velocity is considerably suppressed.

Concluding Remarks on Analyses

A time-dependent numerical EHD analysis for low-modulus line contacts was developed to investigate the temporal behavior of the film profiles of sliding rod seals during a reciprocating cycle. Resulting film profiles reveal all essential features of the temporal film behavior observed by Blok and Koens [8] between a rubber cylinder and a transparent reciprocating surface.

The results indicate that the squeeze-film effect dampens the film thickness fluctuations caused by the oscillatory sliding velocity. It also produces a shift of the maximum film from the mid-stroke position, where the maximum velocity occurs. Likewise, the minimum center film thickness occurs not at the end-stroke, but a few degrees thereafter. For cases of pronounced squeeze-film effect, usually associated with high-load cases, the reduction of the center film after the end-stroke is rather abrupt, particularly at a position where there is substantial sliding velocity.

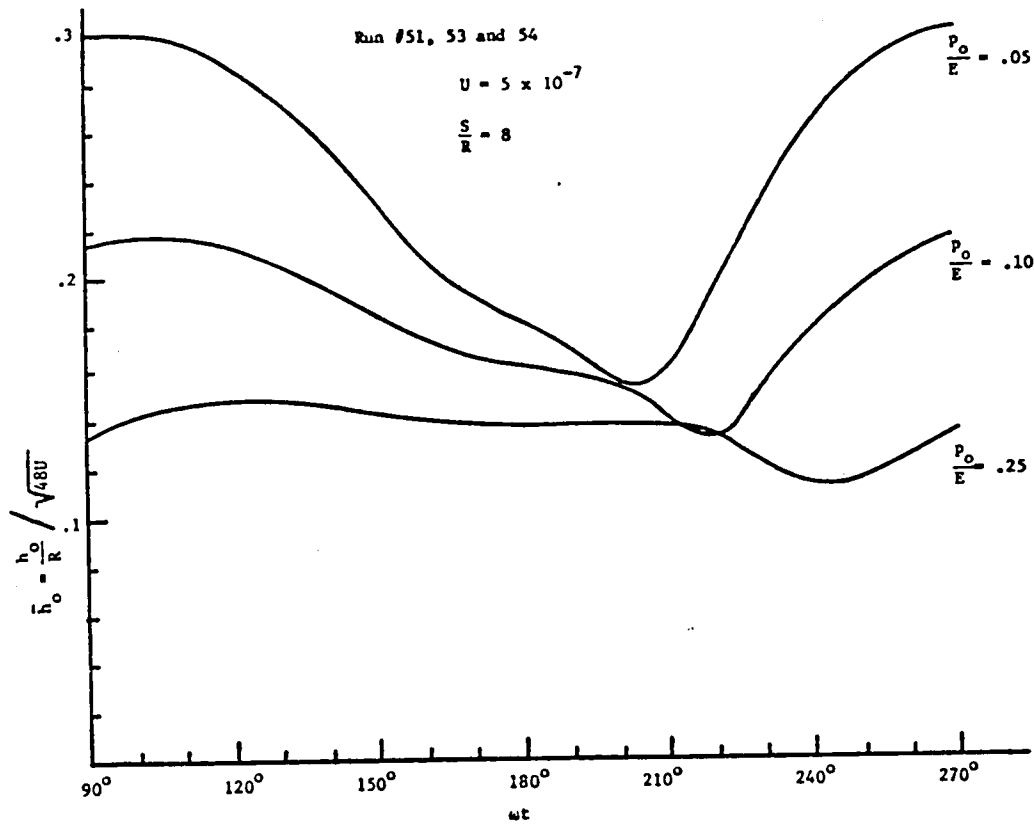


Figure 37 Effect of Load on the Variation of \bar{h}_o

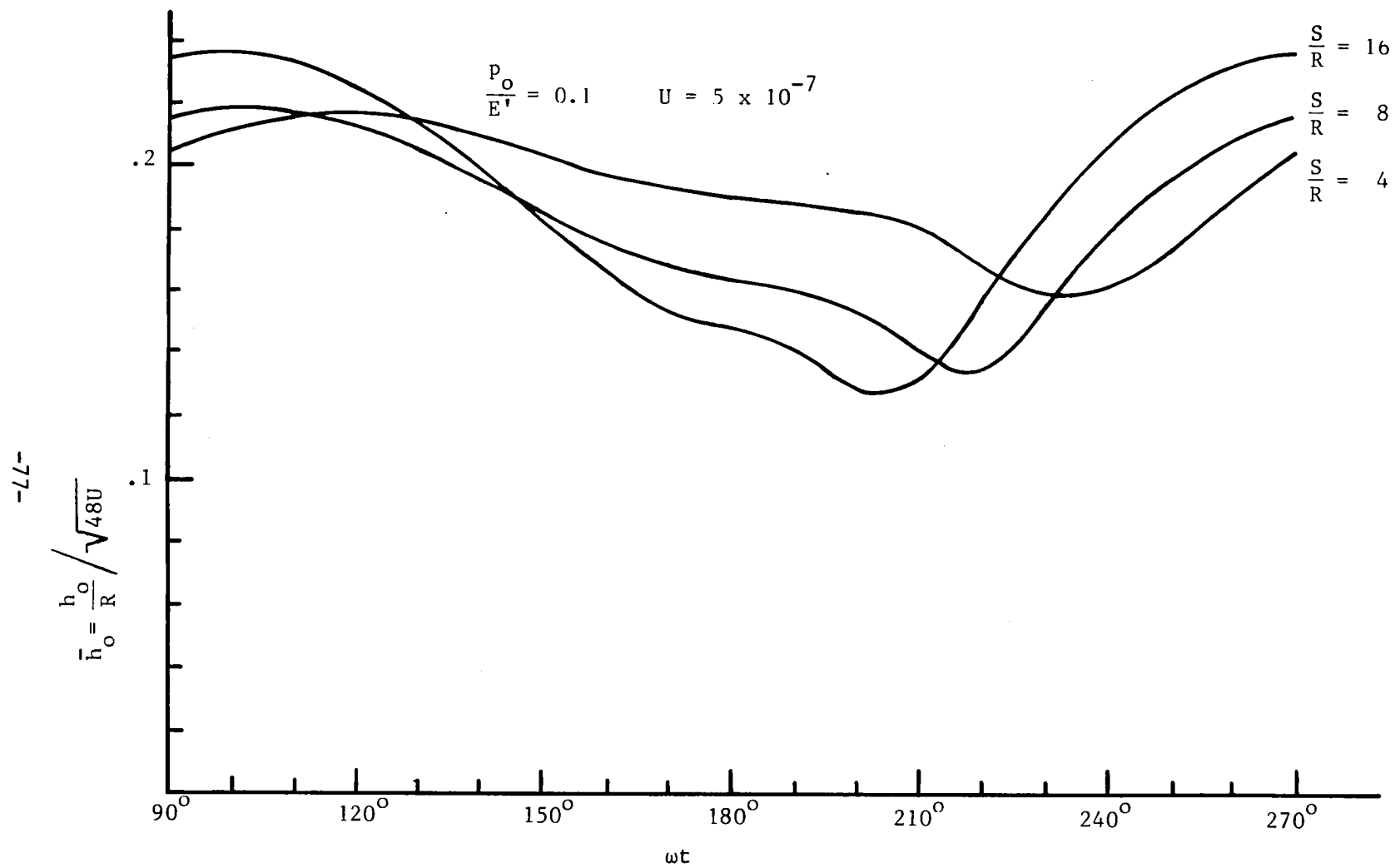


Figure 38 \bar{h}_0 vs ωt , Effect of Stroke to Radius Ratio

The squeeze-film effect increases strongly with an increase in contact pressure, but is almost uninfluenced by a significant change in speed. These findings correlate well with an approximate formula for the change of film thickness due to squeeze-film action based on an existing normal approach analysis [11].

The results developed in the present analysis are based on a pure Hertzian contact uninfluenced by any contact pressure between the seal ring and the groove walls or by the effects of a large sealing pressure differential. The analysis should be valid for cases where the sealing pressure differential is small compared to the contact pressure. The transient analysis developed is an essential building block for further study of the temporal film behavior of rod seals under large fluid sealing pressures. The analysis can be readily modified to include an inlet deformation profile obtained from a static finite element deformation computation. When such a profile is combined with the present analysis, the resulting analysis will be capable of treating the wall pressure effects and the pressure drop across the contact region.

CONCLUSIONS AND RECOMMENDATIONS

On the basis of work done to date, the following conclusions have been reached regarding the experimental apparatus and the analytical effort:

Experimental Apparatus

- . The experimental apparatus, in design and performance, is suitable for evaluating Stirling engine rod seal behavior. The apparatus produces a well-defined reciprocating motion of the transparent cylinder. It exhibits little vibration. The built-in force cell measures properly axial forces applied to the test seal holder. In addition, the device is rugged and easy to use.
- . The interferometric measurements, necessary in the evaluation of seal behavior, can be made with the device. However, the Lexan presently used for the transparent cylinder must be replaced. The new material should be optically smooth both on the inner and outer diameters and should be scratch resistant. The first choice for this material is glass and the second choice is acrylic.

Analysis

- . A time dependent numerical method was developed which determines the fluid film thickness and pressure distribution at the contact zone between an elastomeric seal and a reciprocating cylinder. The method is appropriate for the high values of static contact pressure to seal modulus ratio found in Stirling engine rod seal applications. At its present stage of development, the method applies for a Hertzian static contact pressure profile and for zero pressure gradient across the seal.
- . Application of the numerical method produces results which correlate well with experimental results in the literature and with appropriate approximate analyses. The results obtained show that
 - a) the squeeze-film effect dampens the film thickness fluctuations caused by the oscillatory sliding velocity
 - b) the maximum film does not occur at the instant of maximum velocity
 - c) the minimum film does not occur at the instant of zero velocity
 - d) the squeeze-film effect is affected strongly by the contact pressure but is almost independent of the sliding velocity.

Continuation of the contract work should be directed at maximizing the usefulness of the analytical and experimental tools developed to date. The following recommendations are based on this objective.

It is recommended that the experimental apparatus be used in an in-depth study of Stirling engine rod seal behavior. To do this requires that some refinement to the apparatus be done, that a test program be developed and conducted, that the resulting data be reduced appropriately, and that some runs of the analytical seal model be performed for cases corresponding to the experimental situations.

The refinement of the experimental apparatus entails the addition of hardware to synchronize a stroboscopic photographic flash with the cylinder position.

The test program should include use of the apparatus with the transparent cylinder, with an aluminum cylinder, and with a coated and uncoated seal. Using the transparent cylinder and the coated seal can provide, primarily, film thickness data for comparison with theoretical results. Parameters which can be considered include:

1. Seal type and seal material
2. Seal preload and seal groove geometry
3. Gas pressure to 690 kPa (100 psi), (zero pressure for at least one set of tests)
4. Frequency (10 to 50 Hz)
5. Temperature and lubricant type
6. Seal surface finish (variations to be obtained by changing the optical coating on the seal)

Operation of the apparatus with the glass cylinder and an uncoated seal can reveal the role of the optical coating in the results with the coated seal. To accomplish this, measurements of leakage and friction-force can be made and compared to these obtained for the coated seal.

Operation of the apparatus with the aluminum cylinder and an uncoated seal can reveal the role of cylinder surface finish on friction and leakage. In this, measurements of friction and leakage can be taken and compared to those obtained for the uncoated seal/glass cylinder combination. Also, aluminum cylinders having various surface finishes can be employed in the comparisons.

Regarding data reduction, the most important aspect is quantifying the interferometry results to determine the film thickness and the film thickness profile. Techniques for the film thickness measurement include white light with color tables, two discrete wavelengths of light, or monochromatic light with counting of fringes born. It is desirable that the most suitable technique or techniques be determined and used in the above experimental work.

Additional runs of the elastohydrodynamic seal model will be needed for those situations to be treated experimentally. Such situations, not yet considered analytically, include various combinations of seal properties, seal preload, cyclic frequency, and lubricant viscosity. Comparing similar experimental/analytical cases (to the extent that the model can represent the experimental situation) will provide insight into both the important aspects of the elastohydrodynamic seal behavior and additional capabilities desirable for eventual addition to the computer model. Capabilities already identified for eventual addition to the computer model are a general (non-Hertzian) static pressure distribution and a pressure drop across the seal.

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APPENDIX A

COMPUTER CODE "RODSLJ"



APPENDIX A
Computer Code "RODSLJ"

Flow Chart.

The computer code "RODSLJ" was written to calculate the time-dependent pressure and film thickness profiles, P_k and H_k , at various stroke positions during half of the reciprocating cycle ($\omega t = \pi/2$ to $\omega t = 3\pi/2$). The flow diagram shown in Figure A-1 is based on the numerical treatment described as Scheme "B" in the Analytical Report. For a set of given input data, the calculation of P_k and H_k for the entire half cycle is nested in four loops. In the innermost loop (IT=1, ITP), P_k and H_k are calculated for a prescribed \bar{h}_o to satisfy Equations (37). The second inner loop (ITM=1, ITH) is an iterative loop for solving the center film \bar{h}_o to satisfy the condition, $c_I' - c_o' = 0$. The third inner loop (MT=MII, MTMAX) performs the calculation of \bar{h}_o , P_k , and H_k in equal time intervals from the one mid-stroke position ($\omega t = \pi/2$) to the other mid-stroke position ($\omega t = 3\pi/2$). The outermost loop (IE=1, ITE) repeats the MT loop until $\bar{h}_{o, \pi/2} = \bar{h}_{o, 3\pi/2}$.

The nomenclature used in the computer program is listed in Table A-1.

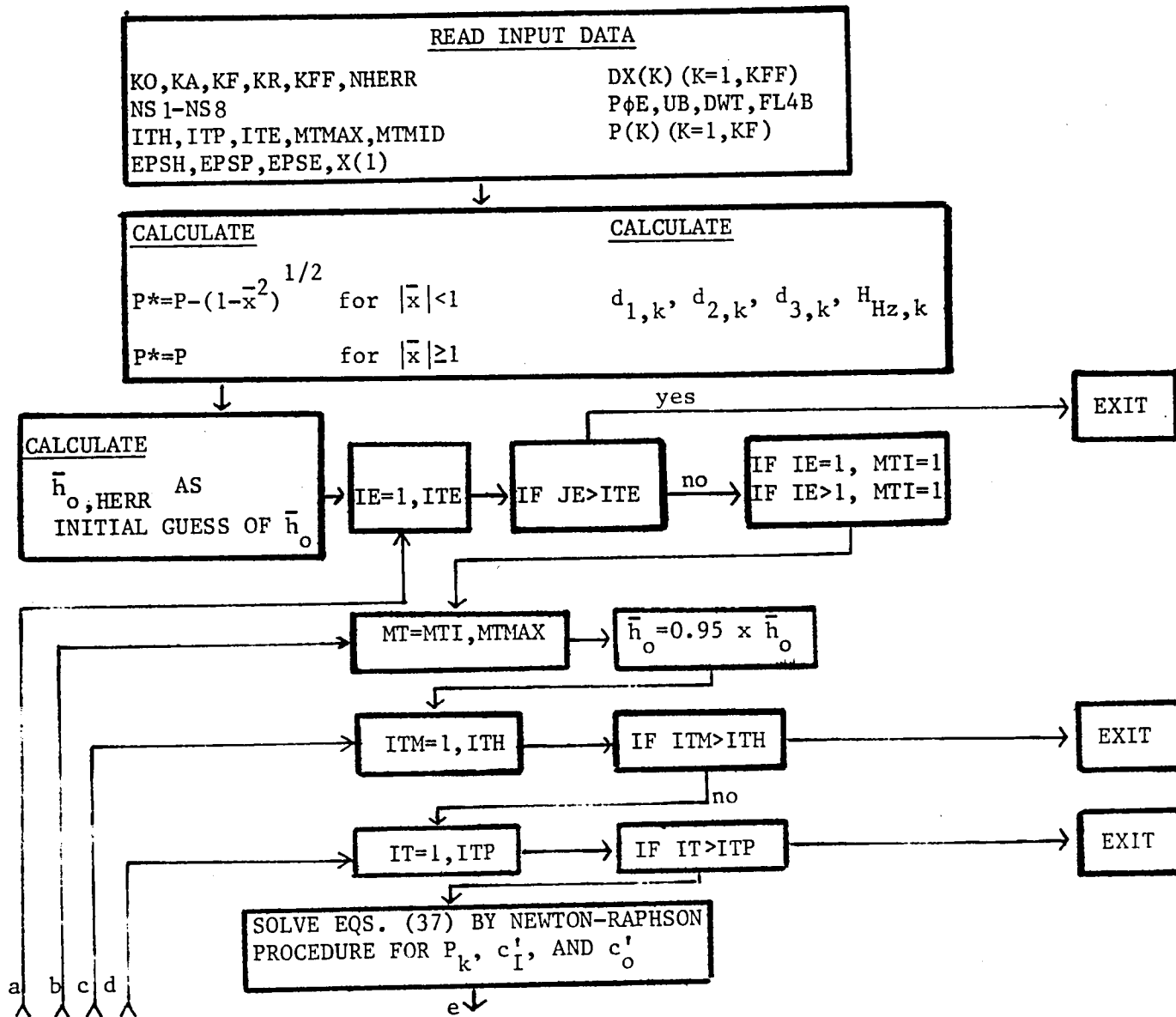


Figure A-1 Flow Diagram for Code "RODSLJ" (Continued on next page)

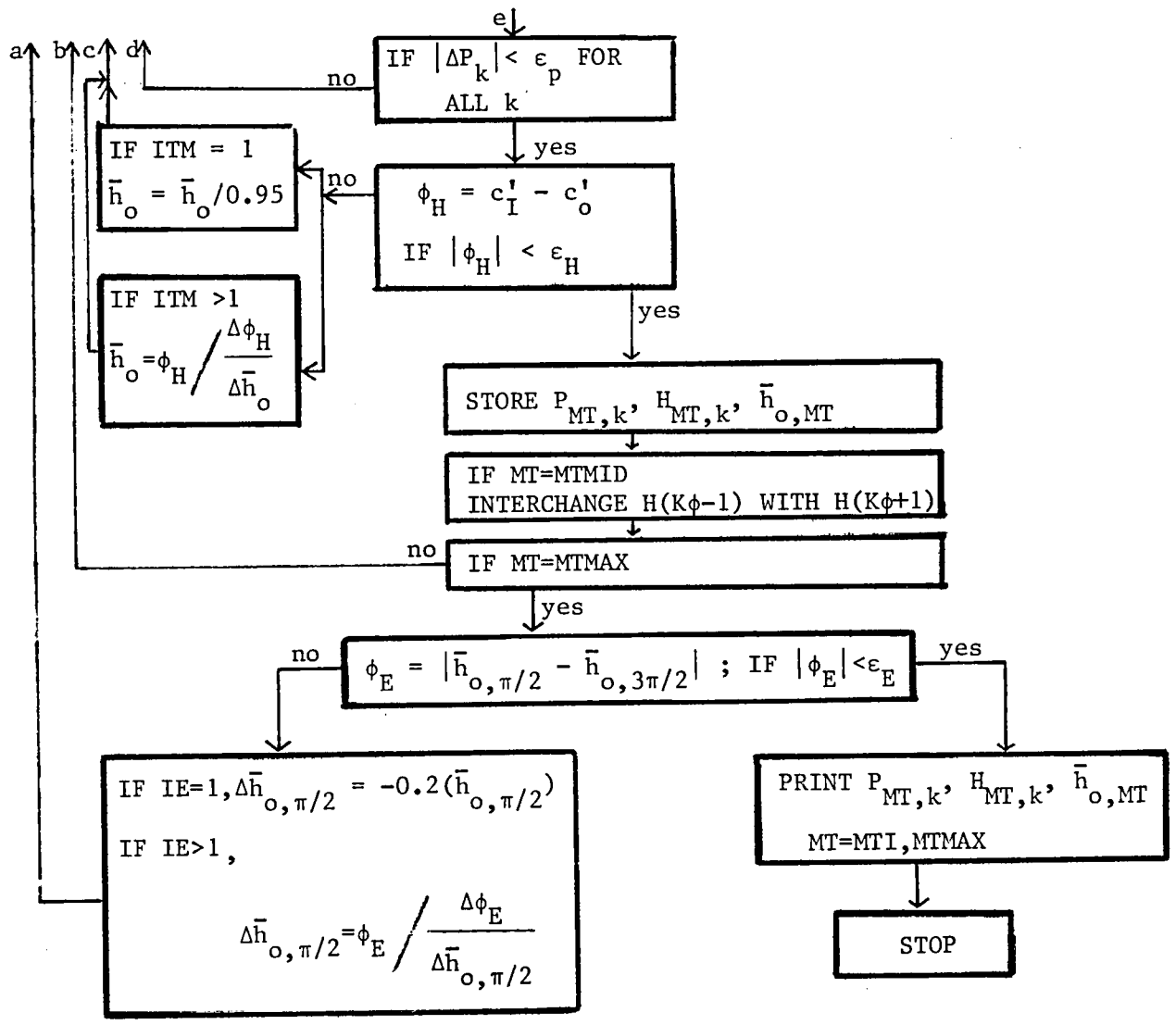


Figure A-1 Flow Diagram for Code "RODSLJ" (Continued from previous page)

Table A-1
List of Major Computer Symbols

Computer Symbol	Symbol in the Analysis
KA	k_A
K ϕ	k'_0
KF	k_f
KFF	k_{ff}
DX(K)	$\Delta \bar{x}_k$
X(K)	\bar{x}_k
P ϕ E	p_0/E'
UB	U
WT	ωt
DWT	$\Delta \omega t$
FL4B	S/(4b)
P(K)	P_k
C7	C_7
DT	ΔT
Q(I,J)	$Q_{i,j}$
H ϕ B	\bar{h}_0
H(K)	H_k
CKI	c'_I
CK ϕ	c'_0
HL(K)	$H_{k,L}$
H ϕ BL	$\bar{h}_{0,L}$
SUMS(K,TH,0)	See Eqs. (38) to (41)
MT	Time interval designation
DP(K)	$(P_{k+1} - P_k) / \Delta \bar{x}_k$

Fortran Listings of Code "RODSLJ"

```
10 PROGRAM RODSLJ(INPUT,CONSOL,OUTPUT,TAPE5=INPUT,TAPE6=CONSOL,
11 1 TAPE7=OUTPUT)
20 DIMENSION A(60,60),C(60),WORK(60),HINT(80),SUMTK(80)
30 DIMENSION HM(80),HP(80),HM2(80),HP2(80),HM3(80),HP3(80)
31 DIMENSION PTEM(80),HTEM(80)
32 DIMENSION PSTOR(25,80),HSTOR(25,80),HOBST(25),SUMKFF(25)
40 COMMON HOB,HOBA(25),WT,DT,KA,KO,KF,KR,KFF,UB,POE
50 COMMON X(80),DX(80),Q(80,80),QJA(80)
60 COMMON H(80),HD(80),HS(80),HL(80),HSL(80)
70 COMMON P(80),PD(80),PS(80),DP(80),DPS(80),PSAVE(80)
80 COMMON FR(80),FS(80),D1(80),D2(80),D3(80)
90 COMMON TA(80),TB(80),TC(80),TD(80),TE(80),TF(80),TG(80),TH(80)
100 COMMON QR(80),HZ(80)
103 C
104 C LAST STATEMENT NO. USED, 600
105 C LAST KFN USED, KF5
106 C
110 C
120 C READ BASIC INPUT DATA
130 C
140 C RODSLC WITHOUT ITERATION OF HOB
141 ITMP=50
142 MTP=50
150 NR=5
160 NW=7
161 NWOUT=7
170 READ(NR,1)
180 WRITE(NWOUT,1)
181 WRITE(NWOUT,1)
190 READ(NR,2)NRUN
200 DO 2000 NNR=1,NRUN
210 READ(NR,2) KA,KO,KF,KR,KFF,NHERR
211 WRITE(NWOUT,2)KA,KO,KF,KR,KFF
220 READ(NR,2)NS1,NS2,NS3,NS4,NS5,NS6,NS7,NS8
221 WRITE(NWOUT,2)NS1,NS2,NS3,NS4,NS5,NS6,NS7,NS8
230 READ(NR,2)ITH,ITP,ITE,MTMAX,MTMID
231 WRITE(NWOUT,2)ITH,ITP,ITE,MTMAX,MTMID
240 READ(NR,3)EPSH,EPSP,EPSE,X(1)
241 WRITE(NWOUT,11)EPSH,EPSP,EPSE
250 WRITE(NWOUT,11)EPSH,X(1)
260 KKF=KF-1
270 READ(NR,3) (DX(K),K=1,KKF)
280 DO 100 K=2,KF
290 100 X(K)=X(K-1)+DX(K-1)
291 WRITE(NWOUT,5)(K,X(K),K=1,KF)
300 PI=3.141593
310 READ(NR,3)POE,UB,DWT,FL4B
320 READ(NR,3)(P(K),K=1,KF)
330 WRITE(NWOUT,11)POE,UB,DWT,FL4B
340 WRITE(NWOUT,5)(K,P(K),K=1,KF)
350 C7=16.0*POE**2/SQRT(48.0*UB)
355 DT=DWT*FL4B
```

```

360 C
370 C          CALCULATE PD(K)
380 C
390      DO 105 K=1,KF
400      PD(K)=0
410      AXK=ABS(X(K))
420      IF(AXK .LT. 1.0) PD(K)=SQRT(1.0-X(K)**2)
430 105 PS(K)=P(K)-PD(K)
440      WRITE(NWOUT,5)(K,PS(K),K=1,KF)
450 C
460 C          CALCULATE Q(I,J)
470 C
480      CALL KERCAL
490      IF (NS5 .EQ. 0)GO TO 104
500      WRITE(NWOUT,7)
510      DO 106 I=KO,KI
520 106 WRITE(NWOUT,5)(J,Q(I,J),J=KO,KF)
530 104 DO 115 K=2,KKF
540      DXS=DX(K)+DX(K-1)
550      D1(K)=DX(K)/(DX(K-1)*DXS)
560      D2(K)=(DX(K)-DX(K-1))/(DX(K)*DX(K-1))
570      D3(K)=DX(K-1)/(DX(K)*DXS)
580 115 CONTINUE
581      DO 510 K=1,KF
583 510 FS(K)=0.0
585      DO 610 K=1,KF
586      AXK=ABS(X(K))
587      SXK=SQRT(AXK**2-1.0)
588      HZ(K)=0.0
589      IF(AXK .LE. 1.0) GO TO 610
590      HZ(K)=ABS(AXK*SXK)-ALOG(AXK+SXK)
591 610 CONTINUE
630 C
640 C          CALCULATE INITIAL HOB
650 C
660      HERR=2.374*UB**0.125/SQRT(48.*POE)/0.8
661      DO 1990 IE=1,ITE
662      MTI=2
663      IF(IE .EQ. 1) MTI=1
664      IF(IE .EQ. 1)HOB=HERR
665      DO 665 MT=MTI,MTMAX
666      FMT=MT
667      WT=(FMT-1.0)*DWT+1.5707963
670 C
680 C          START TIME LOOP
690 C
700      SINWT=SIN(WT)
705      SINWT=ABS(SINWT)
710      ITM=1
711      HOB2=HOB
712      HOB1=HOB*0.95
713      KF1=1
714      HOB=HOB1

```

```

720 C
730 C          CALCULATE F(1) TO P(KA)
740 C
741 122      KFM=KF-1
750          DO 123 K=1,KFM
760 123      PSAVE(K)=F(K)
761          CALL HCAL(1)
766          CALL HCAL(2)
771          IF(NS1 .EQ. 0)GO TO 110
776          WRITE(NWOUT,6)HOB
781          WRITE(NWOUT,5)(K,H(K),K=1,KF)
786 110      CONTINUE
791          DO 120 K=1,KF
796 120      HS(K)=H(K)-HD(K)
800          KKO=KO-1
805          DO 581 K=1,KA
810          HH=(H(K)+H(K+1))*0.5
815          TG(K)=1.0/HH**3
820          FR(K)=(HH-1.)*TG(K)
830 581      CONTINUE
835          IF(MT .EQ. 1) GO TO 585
840          HDH=HOB/L/HOB
845          DO 582 K=1,KO
850 582      TH(K)=(H(K)-HDH*HL(K))
870          FS(KO)=0.0
875          DO 583 K=1,KKO
880          KK=KKO-K+1
885 583      FS(KK)=0.5*(TH(KK)+TH(KK+1))*DX(KK)+FS(KK+1)
890 585      CONTINUE
891          F(1)=0.0
892          DO 586 L=2,KA
893          P(L)=P(L-1)+(SINWT*FR(L-1)-(1.0/DT)*(FS(L-1)+FS(L)))*0.5
894          1*TG(L-1))*DX(L-1)/HOB**2
895 586      CONTINUE
896 124      IT=1
897          IF(ITM .GT. 1) GO TO 590
898          CKI=0.0
899          CKO=0.0
900 590      CONTINUE
905          IF(MT .EQ. 1) GO TO 372
906          HDH=HOB/L/HOB
907          DO 371 K=KO,KF
908 371      TH(K)=(H(K)-HDH*HL(K))
909 372      CONTINUE
910          NNN=KF-KO
911          DO 340 K=1,NNN
912          KK=KF-K
913          TEMP1=0.0
914          IF(MT .EQ. 1) GO TO 373
920          TEMP1=SUMS(KK,TH,1)
922 373      TEMP=(H(KK)-1.0)*SINWT-TEMP1/DT
924          IF(TEMP .GT. 0.0) GO TO 340
926          WRITE(NWOUT,18) KK,H(KK),TEMP1,TEMP
928          KFF=KK+1
930          GO TO 345
932 340      CONTINUE
940 345      CONTINUE
941          DO 417 K=KFF,KF
942 417      F(K)=0.0

```

```

1030 C
1040 C                                CALCULATE FR(K) AND FS(K)
1050 C
1055 125 CONTINUE
1060 126 KKO=KO-1
1070 DO 130 K=1,KO
1075 HH=(H(K)+H(K+1))*0.5
1080 TG(K)=1.0/HH**3
1090 FR(K)=(HH-1.)*TG(K)
1097 130 CONTINUE
1100 IF(MT .EQ. 1) GO TO 140
1110 HDH=HOB/L/HOB
1120 DO 135 K=1,KFF
1121 135 TH(K)=(H(K)-HDH*HL(K))
1131 IF(ITM .LE. ITMP)GOTO 560
1132 IF(MT .NE. MTP)GOTO 560
1133 WRITE(NWOUT,16)
1137 WRITE(NWOUT,5) (K,TH(K),K=1,KFF)
1140 560 FS(KO)=0.0
1150 DO 136 K=1,KKO
1160 KK=KKO-K+1
1170 136 FS(KK)=0.5*(TH(KK)+TH(KK+1))*DX(KK)+FS(KK+1)
1180 140 CONTINUE
1190 IF(NS4 .EQ. 0 ) GO TO 139
1191 IF(MT .EQ. 1) GO TO 139
1200 WRITE(NWOUT,5)(K,FR(K),K=1,KO)
1210 WRITE(NWOUT,5)(K,FS(K),K=1,KO)
1220 139 CONTINUE
1340 IF(NS1 .EQ. 0 ) GO TO 143
1350 WRITE(NWOUT,9)IT, CKI, CKO
1351 IF(NS6. EQ. 0) GO TO 143
1360 WRITE(NWOUT,5)(K,FS(K),K=1,KFF)
1370 WRITE(NWOUT,5)(K,H(K),K=1,KFF)
1380 143 CONTINUE
1390 NN=KFF-KA
1400 C
1410 C                                A(1,1),A(1,2)
1420 C
1490 C
1500 C                                C(1)
1510 C
1520 C
1530 C                                C(N),N=1,NN
1540 C
1545 KFM =KFF-1
1560 DO 170 K=1,KFM
1570 HP(K)=(H(K)+H(K+1))*0.5
1580 HP2(K)=HP(K)**2
1590 HP3(K)=HP2(K)*HP(K)
1595 DP(K)=(P(K+1)-P(K))/DX(K)
1605 170 CONTINUE
1610 DO 175 N=1,NN
1620 K=N+KA-1
1621 C(N)=-HOB**2*HP3(K)*DP(K)+SINWT*(HP(K)-1.0)

```



```

1622      IF(K .LT. KO) C(N)=C(N)-CKI
1623      IF(K .GE. KO) C(N)=C(N)-CKO
1637  561  CONTINUE
1640      IF(MT .EQ. 1) GO TO 175
1641      SUMTK(K)=SUMS(K,TH,0)
1650      C(N)=C(N)-1.0/DT*SUMTK(K)
1670  175  CONTINUE
1680  C
1690  C      A(N,M),N=1,NN,M=1,NN
1700  C
1710      T1=1.5*C7*HOB/PI
1720      T2=HOB**2
1730      T3=C7/PI/HOB
1740      DO 200 N=1,NN
1741      K=N+KA-1
1742      A(N,1)=0.0
1743      A(N,NN)=0.0
1744      IF(K .LT. KO) A(N,1)=1.0
1745      IF(K .GE. KO) A(N,NN)=1.0
1759      NNM=NN-1
1760      DO 200 M=2,NNM
1770      J=M+KA-1
1772      IF(J .GE. KO) J=J+1
1790      DO 505 KQ=1,KFF
1800  505  QQ(KQ)=Q(KQ,J)
1860      A(N,M)=-HP2(K)*DP(K)*(QQ(K+1)+QQ(K))*T1
1870      IF(J .EQ. K+1)A(N,M)=A(N,M)+T2*HP3(K)/DX(K)
1880      IF(J .EQ. K)A(N,M)=A(N,M)-T2*HP3(K)/DX(K)
1890      A(N,M)=A(N,M)+T3*SINWT*0.5*(QQ(K)+QQ(K+1))
1891      IF(NS7 .EQ. 0) GO TO 562
1895      WRITE(NWOUT,17) N,M,A(N,M)
1900  562  IF(MT .EQ. 1)GO TO 200
1910      A(N,M)=A(N,M)-T3*SUMS(K,QQ,0)/DT
1911      IF(NS7 .EQ. 0) GO TO 200
1920      WRITE(NWOUT,17) N,M,A(N,M)
1930  200  CONTINUE
1940  201  IF(NS2 .EQ. 0)GO TO 210
1950      DO 205 KK=1,NN
1960  205  WRITE(NWOUT,5)(JJ,A(KK,JJ),JJ=1,NN)
1970      WRITE(NWOUT,14)
1980      WRITE(NWOUT,5)(KK,C(KK),KK=1,NN)
1990  210  CONTINUE
2000      IDGT=0
2010      CALL LEQT1F(A,1,NN,60,C,IDGT,WORK,IER)
2020      IF(NS1 .EQ. 0)GO TO 220
2021      NNM=NN-1
2030      WRITE(NWOUT,12)
2031      WRITE(NWOUT,11) C(1), C(NN)
2040      WRITE(NWOUT,5)(KK,C(KK),KK=2,NNM)
2050  220  CONTINUE
2060      CV=1.0
2070      CKI=CKI+C(1)
2071      CKO=CKO+C(NN)
2072      NNM=NN-1

```

```

2090      DO225 N=2,NNM
2091      K=N+KA-1
2092      IF(K .GE. KO) K=K+1
2100
2120      DPS(K)=C(N)
2121      IF(DPS(K) .LT. -3.0)GO TO 499
2130      IF(ABS(DPS(K))-EPSP) 222,222,221
2140  221  CV=0.0
2150  222  PS(K)=PS(K)+DPS(K)
2160      P(K)=PD(K)+PS(K)
2180  225  CONTINUE
2220  C
2230  C                      CHECK FOR N-R CONVERGENCE
2240  C
2250      IF(CV .EQ. 1.0) GO TO 240
2260      IF(IT .GT. ITP)GO TO 500
2270  C
2280  C                      N-R NEXT ITERATION
2290  C
2300      IT=IT+1
2310      CALL HCAL(1)
2311      NS2=0
2312      NS4=0
2315      GO TO 125
2320  240  CONTINUE
2340  C
2350  C                      N-R CONVERGED, NEXT MAIN INLET ITERATION
2360  C
2390      PHI=CKI-CKO
2400      GO TO(594,595),KF1
2410  594  PF1=PHI
2411      WRITE(NWOUT,21) ITM, HOB, PHI
2420      HOB=HOB2
2430      KF1=2
2450      ITM=ITM+1
2460      GO TO 122
2470  595  PF2=PHI
2490      DPHI=(PF2-PF1)/(HOB2-HOB1)
2500      DHOB=-PHI/DPHI
2510      WRITE(NWOUT,10)ITM, HOB, PHI, DPHI, DHOB, IT
2520      HOB=HOB+DHOB
2530      IF(ABS(PHI) .LT. EPSH) GO TO 460
2540      IF(ITM .GT. ITH) GO TO 495
2541      ITM=ITM+1
2542      HOB1=HOB2
2543      PF1=PF2
2544      HOB2=HOB
2545      GO TO 122
3350  460  CONTINUE
3351      WRITE(NW,8)MT,WT,HOB
3950      WRITE(NWOUT,8)MT,WT,HOB
3951      NWQ=NWOUT
3952      IF(MT .EQ. MTMID) NWQ=NWOUT
3953      IF(MT .EQ. MTMAX) NWQ=NWOUT

```

```

3960     WRITE(NWQ,5)(K,P(K),K=1,KF)
3961     WRITE(NWQ,5)(K,H(K),K=1,KF)
3962     SUMKFF(MT)=SUMTK(KFM)
3965     HOBST(MT)=HOB
3966     DO 465 K=1,KF
3967     PSTOR(MT,K)=P(K)
3968     465 HSTOR(MT,K)=H(K)
3971     IF(NS8 .EQ. 0)GOTO 640
3972     WRITE(NWOUT,5)(K,TH(K),K=1,KFF)
3973     WRITE(NWOUT,5)(K,SUMTK(K),K=KA,KFM)
3974     640 CONTINUE
3990     CALL HCAL(2)
4000     DO 270 K=1,KF
4010     HS(K)=H(K)-HD(K)
4020     HSL(K)=HS(K)
4030     HL(K)=H(K)
4031     270 CONTINUE
4032     HDBL=HOB
4033     GO TO 121
4042     495 IF(KF5 .EQ. 1) GO TO 500
4043     WRITE(NW,26)
4044     NWOUT=6
4051     NS1=1 $NS2=0 $ NS3=1 $NS4=1 $NS6=1
4052     KF5=1
4053     GO TO 122
4054     499 IF(KF5 .EQ. 1) GO TO 500
4055     WRITE(NW,27)
4056     NWOUT=6
4060     NS1=1
4061     NS4=1
4062     KF5=1
4063     GO TO 125
4064     121 IF(MT .GT. 1) GO TO 655
4065     HIN=HOB
4066     DO 650 K=1,KF
4067     650 HINT(K)=H(K)
4068     655 IF(MT .NE. MTMAX) GO TO 660
4069     HLAST=HOB
4070     PHIA=HIN-HLAST
4071     660 IF(MT .NE. MTMID) GO TO 665
4072     DO 666 K=1,KF
4073     PTEM(K)=P(KF+1-K)
4074     HTEM(K)=H(KF+1-K)
4075     666 CONTINUE
4076     DO 667 K=1,KF
4077     F(K)=PTEM(K)
4078     667 HL(K)=HTEM(K)
4079     WRITE(NWOUT,5)(K,HL(K),K=1,KF)
4080     665 CONTINUE
4104     DO 680 K=1,KF
4105     680 HL(K)=H(K)
4106     IF(IE .EQ. 1) GO TO 670
4107     QF2=PHIA
4108     DPHIA=(QF2-QF1)/(HIN2-HIN1)
4109     DHIN=-PHIA/DPHIA
4110     WRITE(NW,22)
4111     WRITE(NW,23) IE,HIN1,HIN2,PHIA,DPHIA,DHIN
4112     QF1=QF2

```

```

4113      HIN1=HIN2
4114      HIN=HIN+DHIN
4115      HIN2=HIN
4116      IF (ABS(DHIN) .LT. EPSE) GO TO 2000
4117      HOBL=HIN
4118      HOB=HIN
4119      GO TO 1990
4120 670   HIN1=HIN
4121      HIN2=HIN/1.25
4122      QF1=PHIA
4123      WRITE(NW,24) IE,HIN,PHIA
4124      HIN=HIN2
4125      HOBL=HIN
4126      HOB=HIN
4127 1990 CONTINUE
4163 500  CONTINUE
4164 2000 CONTINUE
4165      SUMKFF(1)=0.0
4166      DO 2010 MT=1,MTMAX
4167      FMT=MT
4168      WT=(FMT-1.)*DWT+1.5707963
4169      WRITE(NW,25) MT, WT, HOBST(MT), SUMKFF(MT)
4170      WRITE(NW,5)(K, PSTOR(MT,K),K=1,KF)
4171      WRITE(NW,5)(K, HSTOR(MT,K),K=1,KF)
4172 2010 CONTINUE
4175      STOP
4176 1     FORMAT(72H
4177      1
4178 2     FORMAT(16I5)
4179 3     FORMAT(7E10.3)
4180 4     FORMAT(/4X,4HPOE=,E12.5,5X,3HUB=,E12.5,5X,4HDWT=,E12.5/)
4181 5     FORMAT(8(1X,I2,E12.5))
4182 6     FORMAT(/4X,4HH(K),5X,5HHD(K),4HHOB=,E12.5/)
4200 7     FORMAT(/4X,6HQ(I,J)/)
4210 8     FORMAT(/4X,3HMT=,I5,5X,3HWT=,E12.5,5X,4HHOB=,E12.5/)
4220 9     FORMAT(/4X,3HIT=,I5,4X,4HCKI=,E12.5,4X,4HCKO=,E12.5/)
4230 10    FORMAT(/4X,4HITM=,I5,5X,4HHOB=,E12.5,5X,4HPHI=,E12.5,5X,
4240      1 5HDPHI=,E12.5,5X,5HDHOB=,E12.5,3HIT=,I2)
4250 11    FORMAT(8(1X,E12.5))
4260 12    FORMAT(4X,16HDCKI,DCKO,DPS(K)
4265 14    FORMAT(2X,5HC(KK))
4270 15    FORMAT(/5X,4HKFF=,I5/)
4273 16    FORMAT(2X,5HTH(K))
4277 17    FORMAT(7(1X,2I2,E12.5))
4278 18    FORMAT(I5,3E12.5)
4279 19    FORMAT(2X,1HN,4X,1HM,4X,1HK,4X,1HJ,4X,6HQ(K,J),
4280      17X,8HQ(K+1,J),5X,6HA(N,M))
4281 20    FORMAT(4I5,6(1X,E12.5))
4282 21    FORMAT(/4X,4HITM=,I5,5X,4HHOB=,E12.5,5X,4HPHI=,E12.5
4283      22 FORMAT(/4X,2HIE,3X,4HHIN1,9X,4HHIN2,9X,4HPHIA,9X,5HDPHIA,
4284      1 8X,4HDHIN)
4285 23    FORMAT(4X,I5,5(1X,E12.5)/)
4286 24    FORMAT(/4X,3HIE=,I5,3HHIN,E12.5,4HPHIA,E12.5/)
4287 25    FORMAT(/4X,3HMT=,I5,5X,3HWT=,E12.5,5X,4HHOB=,
4288      1 E12.5,7HSUMKFF=,E12.5/)
4289 26    FORMAT(13HITMEXCEEDED /)
4290 27    FORMAT(13HDPS NEGATIVE /)
4291      END

```

```

4292     SUBROUTINE KERCAL
4293     COMMON HOB,HOBA(25),WT,DT,KA,KO,KF,KR,KFF,UB,POE
4310     COMMON X(80),DX(80),Q(80,80),QJA(80)
4320     COMMON H(80),HD(80),HS(80),HL(80),HSL(80)
4330     COMMON F(80),FD(80),FS(80),DF(80),DPS(80),PSAVE(80)
4340     COMMON FR(80),FS(80),D1(80),D2(80),D3(80)
4350     COMMON TA(80),TB(80),TC(80),TD(80),TE(80),TF(80),TG(80)
4360     DO 1 I=1, KF
4370     DO 1 J=1, KF
4380     1   Q(I,J)=0.0
4390     KKF=KF-2
4400     DO 8 K=1, KF
4410     Q(K,1)=0.0
4420     F5=X(K)
4430     DO 8 J=1, KKF, 2
4440     U=X(J)-F5
4450     U2=X(J+2)-F5
4460     AU=ABS(U)
4470     AU2=ABS(U2)
4480     IF(AU) 50, 51, 50
4490     50 AU=ALOG(AU)
4500     51 IF(AU2) 52, 6, 52
4510     52 AU2=ALOG(AU2)
4520     6  DJ=X(J+1)-X(J)
4530     F2=3.0*DJ
4540     UQ=U*U
4550     U2Q=U2*U2
4560     FK=UQ*(AU-1.5)*0.5
4570     FK2=U2Q*(AU2-1.5)*0.5
4580     FKB=U*(FK-UQ/6.0)-U2*(FK2-U2Q/6.0)
4590     Q(K,J)=((-3.0*FK-FK2)/2.0-FKB/F2)/DJ-U*(AU-1.0)+ Q(K,J)
4600     Q(K,J+1)=(2.0*(FK+FK2)+2.0*FKB/F2)/DJ
4610     Q(K,J+2)=((-FK-3.0*FK2)/2.0-FKB/F2)/DJ+U2*(AU2-1.0)
4620     8  CONTINUE
4630     KFF=KF+1
4640     DO 100 J=1,KF
4650     100 Q(KFF,J)=Q(KR,J)
4660     DO 300 K=1, KF
4670     DO 300 J=1, KF
4680     300 Q(K,J)=Q(K,J)-Q(KFF,J)
4690     RETURN
4700     END

```

```

4710      SUBROUTINE HCAL (KFORK)
4720      COMMON HOB,HOBA(25),WT,DT,KA,KO,KF,KR,KFF,UB,POE
4730      COMMON X(80),DX(80),Q(80,80),QJA(80)
4740      COMMON H(80),HD(80),HS(80),HL(80),HSL(80)
4750      COMMON P(80),PD(80),PS(80),DP(80),DPS(80),PSAVE(80)
4760      COMMON FR(80),FS(80),D1(80),D2(80),D3(80)
4770      COMMON TA(80),TB(80),TC(80),TD(80),TE(80),TF(80),TG(80),TH(80)
4771      COMMON QQ(80),HZ(80)
4780      PI=3.141593
4790      C7=16.0*POE**2/SQRT(48.*UB)
4800      C8=C7/HOB
4810      IF(KFORK .EQ. 2) GOTO 20
4820      DO 10 K=1,KF
4821      IF(ABS(X(K)) .LE. 1.00001) GO TO 5
4830      H(K)=0.
4840      DO 1 J=1,KF
4850      1 H(K)=H(K)+(P(J)-PD(J))*Q(K,J)
4860      H(K)=1. +C8*(0.5*HZ(K)-H(K)/PI)
4861      GO TO 10
4862      5 H(K)=0.0
4863      DO 6 J=1,KF
4864      6 H(K)=H(K)+(P(J)-PD(J))*Q(K,J)
4865      H(K)=1.0-C8*H(K)/PI
4866      10 CONTINUE
4870      GO TO 30
4880      20 DO 25 K=1,KF
4890      HD(K)=0.
4900      DO 21 J=1,KF
4910      21 HD(K)=HD(K)+PD(J)*Q(K,J)
4920      25 HD(K)=1.0+C8*(0.5*X(K)**2-HD(K)/PI)
4930      30 RETURN
4940      END

```

```

4950     FUNCTION SUMN(K1,K2,A)
4960     DIMENSION A(80)
4970     COMMON HOB,HOBA(25),WT,DT,KA,KO,KF,KR,KFF,UB,POE
4980     COMMON X(80),DX(80),Q(80,80),QJA(80)
4990     KK2=K2-1
5000     SUM=0.0
5010     DO 10 K=K1,KK2
5020 10    SUM=SUM+DX(K)*A(K)
5030     SUMN=SUM
5040     RETURN
5050     END
5060

```

```

5070     FUNCTION SUMS(K1,A,NEXIT)
5080     DIMENSION A(80)
5090     COMMON HOB,HOBA(25),WT,DT,KA,KO,KF,KR,KFF,UB,POE
5100     COMMON X(80),DX(80),Q(80,80),QJA(80)
5110     KKO=KO-1
5120     KK1=K1+1
5130     KK2=K1-1
5140     IF(K1 .LT. KKO) GO TO 10
5150     IF(K1 .EQ. KKO) GO TO 20
5160     IF(K1 .EQ. KO) GO TO 30
5170     IF(K1 .GT. KO) GO TO 40
5180     SUMS=0.0
5190     RETURN
5200 10    SUM=DX(K1)/4.0*(0.5*A(K1)+1.5*A(K1+1))
5210     DO 100 K=KK1,KKO
5220 100    SUM=SUM+DX(K)/2.0*(A(K)+A(K+1))
5230     SUMS=SUM
5240     RETURN
5250 20    SUMS=DX(K1)/4.0*(0.5*A(K1)+1.5*A(K1+1))
5260     RETURN
5270 30    SUMS=-DX(KO)/4.0*(1.5*A(KO)+0.5*A(KO+1))
5280     RETURN
5290 40    SUM=0.0
5300     DO 200 K=KO,KK2
5310 200    SUM=SUM-DX(K)/2.0*(A(K)+A(K+1))
5315     IF(NEXIT .EQ. 1) GO TO 210
5320     SUM=SUM-DX(K)/4.0*(1.5*A(K1)+0.5*A(K1+1))
5321 210    SUMS=SUM
5330     RETURN
5340     END

```

Input Instructions

Card 1 Format(72H . . .)
 Identification card for the run

Card 2 Format(16I5)
 NRUN - Number of runs to be made

Card 3 Format(16I5)
 KA - k_A , the grid which divides the contact into two regions.
 For $k < k_A$, the pressure is determined by integrating the
 Reynold's equation directly [13]. For $k > k_A$, the Newton-
 Raphson method is used.

 K ϕ - k_o , the grid number at the contact center.

 KF - k_f , the last grid.

 KR - k_R , the grid at which $H = 1$.

 KFF - k_{ff} , the grid at which the film terminates.

 NHERR - Not used in RODSLJ, set NHERR = 1.

Card 4 Format(16I5)
 NS1 to NS8 - diagnostic controls, set NS1 to NS8 = 0.

Card 5 Format(16I5)
 ITH - Maximum iterations allowed for ITM. Recommended ITH = 15.
 ITP - Maximum iterations allowed for IT. Recommended ITP = 25.
 ITE - Maximum iterations allowed for IE. Recommended ITE = 4.
 MTMAX - (MTMAX-1) represents the total number of time intervals
 within half of the reciprocating cycle.
 MTMID - The timestep MT corresponding to $\omega t = \pi$, or the end-
 stroke position.

Card 6 Format(7E10.3)
 EPSH - Maximum error allowed for ϕ_H , ($\phi_H = |c'_I - c'_O|$). Recommended
 EPSH = 0.0005.
 EPSP - Maximum error allowed for ΔP_k . Recommended EPSP = 0.0005.
 EPSE - Maximum error allowed for ϕ_E , ($\phi_E = |\bar{h}_{o,\pi} - \bar{h}_{o,3\pi/2}|$).
 Recommended EPSE = 0.005.

X(1) — The \bar{x} coordinate for the first grid.

Card 7 Format(7E10.3)

DX(K) (K=1, (KF-1)) — A set of cards for $\Delta\bar{x}_k = \bar{x}_{k+1} - \bar{x}_k$.

Card 8 Format(7E10.3)

PφE — p_o/E'

UB — $\frac{\mu_o u_o}{E'R}$

DWT — $\Delta\omega t$ (radians)

FL4B — S/4b

Card 9 Format(7E10.3)

P(K) (K=1, KF) — A set of cards for the initial guess of pressure profile P_k .

Sample Output Data

A sample of output data based on the input data for Run 51 is included. The first part lists the values of \bar{h}_0 for different stroke positions for each iteration IE. For this sample run, the IE loop converges after three iterations. Thus, results of \bar{h}_0 for IE = 3 are taken to be the final solution for Run 51. In the second part of the output, there are two arrays of data listed under each stroke position ωt . The first array is the pressure profile, and it is followed by the film thickness profile.

The value listed as SUMKFF represents SUMS(KFF,TH,0) which can be used in the future for calculating the leakage rate at each stroke position.

WT=	.15708E+01	HOB=	.20225E+00
WT=	.17453E+01	HOB=	.22136E+00
WT=	.19199E+01	HOB=	.241845E+00
WT=	.20944E+01	HOB=	.26360E+00
WT=	.22689E+01	HOB=	.28682E+00
WT=	.24435E+01	HOB=	.311579E+00
WT=	.26180E+01	HOB=	.33789E+00
WT=	.27925E+01	HOB=	.36479E+00
WT=	.29671E+01	HOB=	.39224E+00
WT=	.31416E+01	HOB=	.42023E+00
WT=	.33161E+01	HOB=	.44876E+00
WT=	.34907E+01	HOB=	.47783E+00
WT=	.36652E+01	HOB=	.50744E+00
WT=	.38397E+01	HOB=	.53759E+00
WT=	.40143E+01	HOB=	.56828E+00
WT=	.41888E+01	HOB=	.59951E+00
WT=	.43633E+01	HOB=	.63128E+00
WT=	.45379E+01	HOB=	.66359E+00
WT=	.47124E+01	HOB=	.69644E+00

WT=	.17453E+01	HOB=	.21902E+00
WT=	.19199E+01	HOB=	.21828E+00
WT=	.20944E+01	HOB=	.21348E+00
WT=	.22689E+01	HOB=	.20572E+00
WT=	.24435E+01	HOB=	.19572E+00
WT=	.26180E+01	HOB=	.18465E+00
WT=	.27925E+01	HOB=	.17436E+00
WT=	.29671E+01	HOB=	.16702E+00
WT=	.31416E+01	HOB=	.16329E+00
WT=	.33161E+01	HOB=	.16116E+00
WT=	.34907E+01	HOB=	.15440E+00
WT=	.36652E+01	HOB=	.14171E+00
WT=	.38397E+01	HOB=	.13966E+00
WT=	.40143E+01	HOB=	.15446E+00
WT=	.41888E+01	HOB=	.17601E+00
WT=	.43633E+01	HOB=	.19485E+00
WT=	.45379E+01	HOB=	.20817E+00
WT=	.47124E+01	HOB=	.21591E+00

2	WT=	.21591E+00	HOB=	.35671E-07
---	-----	------------	------	------------

00	0	WT=	.17453E+01	HOB=	.19258E+00
01	1	WT=	.19199E+01	HOB=	.20439E+00
02	2	WT=	.20944E+01	HOB=	.20760E+00
03	3	WT=	.22689E+01	HOB=	.20356E+00
04	4	WT=	.24435E+01	HOB=	.19477E+00
05	5	WT=	.26180E+01	HOB=	.18412E+00
06	6	WT=	.27925E+01	HOB=	.17401E+00
07	7	WT=	.29671E+01	HOB=	.16675E+00
08	8	WT=	.31416E+01	HOB=	.16303E+00
09	9	WT=	.33161E+01	HOB=	.16087E+00
10	10	WT=	.34907E+01	HOB=	.15404E+00
11	11	WT=	.36652E+01	HOB=	.14138E+00
12	12	WT=	.38397E+01	HOB=	.13947E+00
13	13	WT=	.40143E+01	HOB=	.15438E+00
14	14	WT=	.41888E+01	HOB=	.17599E+00
15	15	WT=	.43633E+01	HOB=	.19485E+00
16	16	WT=	.45379E+01	HOB=	.20817E+00
17	17	WT=	.47124E+01	HOB=	.21591E+00

00	0000	0000	0000	0000	0000
3	00000000	.17780E+00	-.39109E-01	.99999E+00	.38108E-01

MT= 7 WT= .26180E+01 HOB= .18465E+00SUMKFF= .48946E-01

1	2	.12441E-02	3	.92645E-02	4	.20495E-01	5	.43369E-01	6	.56304E-01	7	.72789E-01	8	.93398E-01
9	10	.14732E+00	11	.17970E+00	12	.21383E+00	13	.24805E+00	14	.31269E+00	15	.36856E+00	16	.49977E+00
17	18	.79817E+00	19	.91554E+00	20	.97977E+00	21	.10000E+01	22	.98038E+00	23	.91701E+00	24	.90113E+00
25	26	.51197E+00	27	.39597E+00	28	.34798E+00	29	.28399E+00	30	.23890E+00	31	.17591E+00	32	.94201E-01
33	34	-.67663E-02	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	42	0.	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.
1	2	.29913E+02	3	.69273E+01	4	.47858E+01	5	.31133E+01	6	.26713E+01	7	.22910E+01	8	.19732E+01
9	10	.15062E+01	11	.13390E+01	12	.12217E+01	13	.11361E+01	14	.10309E+01	15	.97637E+00	16	.93287E+00
17	18	.94312E+00	19	.96478E+00	20	.98713E+00	21	.10000E+01	22	.10145E+01	23	.10177E+01	24	.10209E+01
25	26	.98685E+00	27	.96774E+00	28	.95176E+00	29	.90966E+00	30	.87212E+00	31	.82324E+00	32	.80081E+00
33	34	.10324E+01	35	.13506E+01	36	.17582E+01	37	.22363E+01	38	.40265E+01	39	.62642E+01	40	.29532E+02
41	42	.64139E+02	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.

MT= 8 WT= .27925E+01 HOB= .17436E+00SUMKFF= .51364E-01

1	2	.87684E-03	3	.65641E-02	4	.15280E-01	5	.34306E-01	6	.45920E-01	7	.61383E-01	8	.81608E-01
9	10	.13771E+00	11	.17261E+00	12	.20949E+00	13	.24607E+00	14	.31358E+00	15	.37042E+00	16	.50034E+00
17	18	.79834E+00	19	.91565E+00	20	.97983E+00	21	.10000E+01	22	.98037E+00	23	.91691E+00	24	.90108E+00
25	26	.51181E+00	27	.39589E+00	28	.34715E+00	29	.28135E+00	30	.23495E+00	31	.17112E+00	32	.90937E-01
33	34	-.30973E-02	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	42	0.	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.
1	2	.31508E+02	3	.70720E+01	4	.47803E+01	5	.29861E+01	6	.25132E+01	7	.21089E+01	8	.17750E+01
9	10	.13006E+01	11	.11400E+01	12	.10342E+01	13	.96184E+00	14	.88176E+00	15	.84607E+00	16	.83299E+00
17	18	.91753E+00	19	.94150E+00	20	.97791E+00	21	.10000E+01	22	.10217E+01	23	.10288E+01	24	.10355E+01
25	26	.99752E+00	27	.97403E+00	28	.95281E+00	29	.90233E+00	30	.86068E+00	31	.81002E+00	32	.79269E+00
33	34	.10551E+01	35	.13913E+01	36	.18231E+01	37	.23302E+01	38	.42288E+01	39	.66010E+01	40	.31253E+02
41	42	.47999E+02	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.

MT= 9 WT= .29671E+01 HOB= .16702E+00SUMKFF= .48080E-01

1	2	.45541E-03	3	.35265E-02	4	.88070E-02	5	.21649E-01	6	.30598E-01	7	.43558E-01	8	.62169E-01
9	10	.12107E+00	11	.16077E+00	12	.20291E+00	13	.24369E+00	14	.31540E+00	15	.37382E+00	16	.50175E+00
17	18	.79879E+00	19	.91592E+00	20	.97983E+00	21	.10000E+01	22	.98026E+00	23	.91672E+00	24	.90090E+00
25	26	.51163E+00	27	.39552E+00	28	.34521E+00	29	.27656E+00	30	.22842E+00	31	.16408E+00	32	.97043E-01
33	34	-.13270E-02	35	-.44577E-02	36	0.	37	0.	38	0.	39	0.	40	0.
41	42	0.	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.
1	2	.32695E+02	3	.70684E+01	4	.46389E+01	5	.27245E+01	6	.22199E+01	7	.17912E+01	8	.14430E+01
9	10	.97742E+00	11	.93838E+00	12	.76028E+00	13	.71619E+00	14	.68451E+00	15	.68363E+00	16	.72390E+00
17	18	.86769E+00	19	.92606E+00	20	.97127E+00	21	.10000E+01	22	.10246E+01	23	.10323E+01	24	.10385E+01
25	26	.98952E+00	27	.95475E+00	28	.92350E+00	29	.86050E+00	30	.81434E+00	31	.76389E+00	32	.75527E+00
33	34	.10477E+01	35	.13980E+01	36	.18546E+01	37	.23880E+01	38	.43771E+01	39	.68585E+01	40	.32604E+02
41	42	.70843E+02	43	0.	44	0.	45	0.	46	0.	47	0.	48	0.

MT= 10 WT= .31416E+01 HOB= .16329E+00SUMKFF= .42561E-01

1	1	2	.58496E-05	3	.10985E-03	4	.61054E-03	5	.26038E-02	6	.50123E-02	7	.99954E-02	8	.20944E-01
2	13303E-01	10	.88138E-01	11	.14990E+00	12	.21019E+00	13	.25936E+00	14	.33192E+00	15	.38624E+00	16	.50713E+00
17	19513E+00	18	.79937E+00	19	.91614E+00	20	.97980E+00	21	.10000E+01	22	.98007E+00	23	.91660E+00	24	.80056E+00
25	6003E+00	26	.51126E+00	27	.39255E+00	28	.33693E+00	29	.26015E+00	30	.20804E+00	31	.14546E+00	32	.83261E-01
33	49265E-01	34	.17989E-01	35	.90616E-02	36	.37814E-02	37	.18259E-02	38	.37005E-03	39	.51213E-04	40	.21407E-05
41	1	2	.33272E+02	3	.69053E+01	4	.43652E+01	5	.23353E+01	6	.17957E+01	7	.13393E+01	8	.98041E+00
9	74209E+00	10	.59241E+00	11	.54258E+00	12	.55458E+00	13	.58024E+00	14	.63027E+00	15	.66558E+00	16	.73006E+00
17	27635E+00	18	.87155E+00	19	.92913E+00	20	.97231E+00	21	.10000E+01	22	.10199E+01	23	.10236E+01	24	.10190E+01
25	97255E+00	26	.94459E+00	27	.86909E+00	28	.81336E+00	29	.73116E+00	30	.68804E+00	31	.66139E+00	32	.70182E+00
33	9498E+00	34	.10839E+01	35	.14426E+01	36	.18992E+01	37	.24387E+01	38	.44673E+01	39	.70050E+01	40	.33353E+02
41	1	2	.45178E-03	3	.34576E-02	4	.83580E-02	5	.20175E-01	6	.28388E-01	7	.40340E-01	8	.57688E-01
9	82035E-01	10	.11439E+00	11	.15378E+00	12	.19609E+00	13	.23729E+00	14	.31081E+00	15	.37059E+00	16	.50127E+00
17	5927E+00	18	.79999E+00	19	.91657E+00	20	.97997E+00	21	.10000E+01	22	.97982E+00	23	.91610E+00	24	.79941E+00
25	19369E+00	26	.50838E+00	27	.38917E+00	28	.33963E+00	29	.27716E+00	30	.23445E+00	31	.17122E+00	32	.80196E-01
33	55176E-02	34	-.10609E-01	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	1	2	.33891E+02	3	.72946E+01	4	.47637E+01	5	.27611E+01	6	.22305E+01	7	.17778E+01	8	.14087E+01
9	11312E+01	10	.91364E+00	11	.76730E+00	12	.68709E+00	13	.64473E+00	14	.62455E+00	15	.63928E+00	16	.73290E+00
17	94383E+00	18	.95557E+00	19	.99416E+00	20	.10065E+01	21	.10000E+01	22	.98282E+00	23	.94915E+00	24	.90664E+00
25	32862E+00	26	.79506E+00	27	.75049E+00	28	.73014E+00	29	.69148E+00	30	.65550E+00	31	.60115E+00	32	.57074E+00
33	64384E+00	34	.87478E+00	35	.12637E+01	36	.17471E+01	37	.23079E+01	38	.43884E+01	39	.69740E+01	40	.33710E+02
41	1	2	.87567E-03	3	.45505E-02	4	.15126E-01	5	.33943E-01	6	.45508E-01	7	.60982E-01	8	.81330E-01
9	10708E+00	10	.13818E+00	11	.17350E+00	12	.21056E+00	13	.24700E+00	14	.31377E+00	15	.36773E+00	16	.49843E+00
17	19503E+00	18	.79746E+00	19	.91550E+00	20	.97977E+00	21	.10000E+01	22	.98011E+00	23	.91637E+00	24	.80009E+00
25	6003E+00	26	.50942E+00	27	.39048E+00	28	.34176E+00	29	.28143E+00	30	.24003E+00	31	.17694E+00	32	.84041E-01
33	49265E-02	34	0.	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	1	2	.35494E+02	3	.78826E+01	4	.52891E+01	5	.32567E+01	6	.27207E+01	7	.22622E+01	8	.18838E+01
9	17932E+01	10	.13456E+01	11	.11627E+01	12	.10404E+01	13	.95450E+00	14	.85160E+00	15	.79437E+00	16	.75457E+00
17	9159E+00	18	.84365E+00	19	.92469E+00	20	.97657E+00	21	.10000E+01	22	.10085E+01	23	.99336E+00	24	.97177E+00
25	21721E+00	26	.89284E+00	27	.86039E+00	28	.84739E+00	29	.81767E+00	30	.78406E+00	31	.72762E+00	32	.68297E+00
33	1098E+00	34	.10106E+01	35	.13983E+01	36	.18902E+01	37	.24473E+01	38	.46224E+01	39	.73098E+01	40	.33710E+02
41	1	2	.15440E+00SUMKFF=	3	-.18799E-01	4		5		6		7		8	

MT= 11 WT= .33161E+01 HOB= .16116E+00SUMKFF= -.19066E-01

1	1	2	.87567E-03	3	.45505E-02	4	.15126E-01	5	.33943E-01	6	.45508E-01	7	.60982E-01	8	.81330E-01
9	10708E+00	10	.13818E+00	11	.17350E+00	12	.21056E+00	13	.24700E+00	14	.31377E+00	15	.36773E+00	16	.49843E+00
17	19503E+00	18	.79746E+00	19	.91550E+00	20	.97977E+00	21	.10000E+01	22	.98011E+00	23	.91637E+00	24	.80009E+00
25	6003E+00	26	.50942E+00	27	.39048E+00	28	.34176E+00	29	.28143E+00	30	.24003E+00	31	.17694E+00	32	.84041E-01
33	49265E-02	34	0.	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	1	2	.35494E+02	3	.78826E+01	4	.52891E+01	5	.32567E+01	6	.27207E+01	7	.22622E+01	8	.18838E+01
9	17932E+01	10	.13456E+01	11	.11627E+01	12	.10404E+01	13	.95450E+00	14	.85160E+00	15	.79437E+00	16	.75457E+00
17	9159E+00	18	.84365E+00	19	.92469E+00	20	.97657E+00	21	.10000E+01	22	.10085E+01	23	.99336E+00	24	.97177E+00
25	21721E+00	26	.89284E+00	27	.86039E+00	28	.84739E+00	29	.81767E+00	30	.78406E+00	31	.72762E+00	32	.68297E+00
33	1098E+00	34	.10106E+01	35	.13983E+01	36	.18902E+01	37	.24473E+01	38	.46224E+01	39	.73098E+01	40	.33710E+02
41	1	2	.15440E+00SUMKFF=	3	-.18799E-01	4		5		6		7		8	

MT= 12 WT= .34907E+01 HOB= .15440E+00SUMKFF= -.18799E-01

1	1	2	.87567E-03	3	.45505E-02	4	.15126E-01	5	.33943E-01	6	.45508E-01	7	.60982E-01	8	.81330E-01
9	10708E+00	10	.13818E+00	11	.17350E+00	12	.21056E+00	13	.24700E+00	14	.31377E+00	15	.36773E+00	16	.49843E+00
17	19503E+00	18	.79746E+00	19	.91550E+00	20	.97977E+00	21	.10000E+01	22	.98011E+00	23	.91637E+00	24	.80009E+00
25	6003E+00	26	.50942E+00	27	.39048E+00	28	.34176E+00	29	.28143E+00	30	.24003E+00	31	.17694E+00	32	.84041E-01
33	49265E-02	34	0.	35	0.	36	0.	37	0.	38	0.	39	0.	40	0.
41	1	2	.35494E+02	3	.78826E+01	4	.52891E+01	5	.32567E+01	6	.27207E+01	7	.22622E+01	8	.18838E+01
9	17932E+01	10	.13456E+01	11	.11627E+01	12	.10404E+01	13	.95450E+00	14	.85160E+00	15	.79437E+00	16	.75457E+00
17	9159E+00	18	.84365E+00	19	.92469E+00	20	.97657E+00	21	.10000E+01	22	.10085E+01	23	.99336E+00	24	.97177E+00
25	21721E+00	26	.89284E+00	27	.86039E+00	28	.84739E+00	29	.81767E+00	30	.78406E+00	31	.72762E+00	32	.68297E+00
33	1098E+00	34	.10106E+01	35	.13983E+01	36	.18902E+01	37	.24473E+01	38	.46224E+01	39	.73098E+01	40	.33710E+02
41	1	2	.15440E+00SUMKFF=	3	-.18799E-01	4		5		6		7		8	

HT= 19 WT= .47124E+01 HOB= .21591E+00SUMKFF= -.46652E-01

1 0.	2 .25296E-02	3 .17753E-01	4 .34376E-01	5 .63940E-01	6 .78593E-01	7 .96036E-01	8 .11644E+00
9 .13972E+00	10 .16567E+00	11 .19383E+00	12 .22335E+00	13 .25331E+00	14 .31192E+00	15 .36536E+00	16 .49769E+00
17 .50146E+00	18 .79802E+00	19 .91548E+00	20 .97967E+00	21 .10000E+01	22 .98015E+00	23 .91686E+00	24 .80067E+00
25 .60112E+00	26 .51147E+00	27 .39420E+00	28 .34688E+00	29 .28739E+00	30 .24632E+00	31 .18631E+00	32 .10079E+00
33 .13666E-01	34 -.19919E-01	35 0.	36 0.	37 0.	38 0.	39 0.	40 0.
41 0.							
1 .55359E+02	2 .25982E+02	3 .65160E+01	4 .47211E+01	5 .33180E+01	6 .29425E+01	7 .26143E+01	8 .23334E+01
9 .21043E+01	10 .18957E+01	11 .17254E+01	12 .15951E+01	13 .14917E+01	14 .13472E+01	15 .12596E+01	16 .11539E+01
17 .11047E+01	18 .10603E+01	19 .10360E+01	20 .10183E+01	21 .10000E+01	22 .98393E+00	23 .96389E+00	24 .94419E+00
25 .91458E+00	26 .90017E+00	27 .88287E+00	28 .87507E+00	29 .85205E+00	30 .82621E+00	31 .78512E+00	32 .75463E+00
33 .69306E+00	34 .91882E+00	35 .11995E+01	36 .15528E+01	37 .19634E+01	38 .34957E+01	39 .54090E+01	40 .85226E+02
41 .98875E+02							

APPENDIX B

DERIVATION OF $Q_{k,j}$

APPENDIX B

Derivation of $Q_{k,j}$

The kernel functions $Q_{k,j}$ were developed originally in Reference [14] and are reproduced as follows:

$$Q_{k,j} = K_1(k_o, j) - K_1(k_o, j) \quad \text{for } j = 1$$

$$Q_{k,j} = K_2(k, j-1) - K_2(k_o, j-1) \quad \text{for } j = 2, 4, 6, \dots, k_f-1$$

$$Q_{k,j} = \left[K_1(k, j) + K_3(k, j-2) \right] - \left[K_1(k_o, j) + K_3(k_o, j-2) \right] \\ \text{for } j = 3, 5, 7, \dots, k_f-2$$

$$Q_{k,j} = K_3(k, j-2) - K_3(k_o, j-2) \quad \text{for } j = k_f$$

where

$$K_1(k, j) = \frac{1}{2\delta_j} (-3v_j - v_{j+2}) - \frac{\bar{v}_j}{3\delta_j^2} - u_j (\ln|u_j| - 1)$$

$$K_2(k, j) = \frac{2}{\delta_j} (v_j + v_{j+2}) + \frac{2\bar{v}_j}{3\delta_j^2}$$

$$K_3(k, j) = \frac{1}{2\delta_j} (-v_j - 3v_{j+2}) - \frac{\bar{v}_j}{3\delta_j^2} + u_{j+2} (\ln|u_{j+2}| - 1)$$

and

$$\delta_j = \bar{\xi}_{j+1} - \bar{\xi}_j$$

$$u_j = \bar{\xi}_j - \bar{x}_k$$

$$v_j = \frac{u_j^2}{2} \left(\ln|u_j| - \frac{3}{2} \right)$$

$$\bar{v}_j = u_j \left[v_j - \frac{u_j^2}{6} \right] - u_{j+2} \left[v_{j+2} - \frac{u_{j+2}^2}{6} \right]$$

APPENDIX C

DERIVATION OF $\frac{\partial \psi_{k+1/2}}{\partial P_j}$ FOR EQUATIONS (37)

APPENDIX C

Derivation of $\frac{\partial \psi_{k+1/2}}{\partial P_j}$ for Equations (37)

Equations (37) are linearized to form the following system of equations which are solved for P_{k_A} to P_{k_o-1} , P_{k_o+1} to $P_{k_{ff}-1}$, c'_I , and c'_o .

$$\frac{\partial \psi_{k+1/2}}{\partial c'_I} \Delta c'_I + \sum_{j=k_A}^{k_o-1} \left(\frac{\partial \psi_{k+1/2}}{\partial P_j} \right)^n \Delta P_j^{n+1} + \sum_{j=k_o+1}^{k_{ff}-1} \left(\frac{\partial \psi_{k+1/2}}{\partial P_j} \right)^n \Delta P_j^{n+1} + \frac{\partial \psi_{k+1/2}}{\partial c'_o} \Delta c'_o =$$

$$= \left(-\psi_{k+1/2} \right)^n$$

$$k = k_A, \dots, k_{ff}-1 \quad (C-1)$$

where

$$\left. \begin{aligned} \frac{\partial \psi_{k+1/2}}{\partial c'_I} &= 1 \\ \frac{\partial \psi_{k+1/2}}{\partial c'_o} &= 0 \end{aligned} \right\} \text{for } k = k_A \text{ to } k = k_o-1 \quad (C-2)$$

$$\left. \begin{aligned} \frac{\partial \psi_{k+1/2}}{\partial c'_I} &= 0 \\ \frac{\partial \psi_{k+1/2}}{\partial c'_o} &= 1 \end{aligned} \right\} \text{for } k = k_o \text{ to } k = k_{ff}-1 \quad (C-3)$$

Also,

$$\frac{\partial \psi_k}{\partial P_j} = 3\bar{h}_o^{-2} H_{k+1/2}^2 \left(\frac{P_{k+1} - P_k}{\Delta \bar{x}_k} \right) \left(- \frac{(QQ)_{k+1/2} C_7}{\pi \bar{h}_o} \right)$$

$$\begin{aligned}
& + \delta \frac{\bar{h}_o^2 H_{k+1/2}^3}{\Delta \bar{x}_k} + \sin \omega t (QQ)_{k+1/2} \frac{C_7}{\pi \bar{h}_o} \\
& - \frac{1}{\Delta T} \text{SUMS}(K, QQ, 0)
\end{aligned} \tag{C-4}$$

where $(QQ)_k = Q_{k,j}$ for $j \neq k_A$

$$(QQ)_k = \sum_{\ell=1}^{k_A} \frac{P_\ell}{P_{k_A}} Q_{k,\ell} \text{ for } j \neq k_A$$

$$\left. \begin{aligned}
\delta &= 1 \quad \text{for } j = k+1 \\
\delta &= -1 \quad \text{for } j = k \\
\delta &= 0 \quad \text{for } j \neq k \text{ or } j \neq k+1
\end{aligned} \right\} \tag{C-5}$$

$$C_7 = 16 \left(\frac{P_o}{E^*} \right)^2 / \sqrt{48U}$$

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16. Abstract This report concerns the first year of a two year experimental and analytical program. The program is directed at the elastohydrodynamic behavior of sliding elastomeric rod seals for the Stirling engine. During the year, experimental and analytical tools were developed for evaluating seal leakage, seal friction, and the fluid film thickness at the seal/cylinder interface. Experimentally, a rugged apparatus containing a moving transparent cylinder (guided by precision hydrostatic bearings) and a stationary elastomeric test seal has been constructed. A pressure gradient of 690 kPa (100 psi) can be applied across the seal. Frequencies from 10 Hz to 50 Hz with a 25 mm (1 in.) total stroke can be employed. Film thickness is measured with optical interferometry, fluid leakage by level and pressure changes, and power (friction) loss by a force cell. Analytically, a computer model of the elastohydrodynamic behavior of the rod seal for reciprocating motion was produced. The model permits the large initial radial squeeze, which is typically used for elastomeric seals, to exist. The model, based on a time domain analysis, determines both the pressure distribution and the oil film thickness distribution in the seal/cylinder contact zone. Conclusions and recommendations concern use during the ensuing year of these experimental and analytical tools.					
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