## NOTICE

THIS DOCUMENT HAS BEEN REPRODUCED FROM MICROFICHE. ALTHOUGH IT IS RECOGNIZED THAT CERTAIN PORTIONS ARE ILLEGIBLE, IT IS BEING RELEASED IN THE INTEREST OF MAKING AVAILABLE AS MUCH INFORMATION AS POSSIBLE

# ASNAPIOTIC EIGENSTRUCTURES* 

## by

Peter M, Thompson Gunter Stein

## ABSTRACT

The behavior of the closed loop aigenstructure of a linear system with output feedback is analyzed as a single parameter multiplying the feedback gain is varied. An algorithm is presented that computes the asymptotically infinite eigenstructure, and it is shown how a system with high gain feedback decouples into single input single output systems. Then a synthesis algorithm is presented which uses full state feedback to achieve a desired asymptotic eigenstructure.


[^0]

## I. Introduction

The closed looop eigenstructure of a insear syatem has long been recognized as an important conaideration in the design of control systems. Most attention in the control literature has been given to the behavior of closed loop eigenvalues as parameters of the control system are varied, and the root locus has emerged as a tool to study this behavior. The behavior of closed loop eigenvectors has also received a lot of attention because for multi-input multi-output (MIMO) systems the placement of eigenvectors can be used to shape transient responses, decouple inputs and outputs, and reject certain types of disturbances.

The design procedure we have in mind is to specify an asymptotic eigenstructure that can be achieved with high gain state feedback. Then the design can be simplified to choosing a single gain so that bandwidth constraints are satisfied. A way to do this using the linear quadratic regulator (LQR) is given in $[1,2]$. The design procedure suggested there is to specify an asymptotic eigenstructure, choose quadratic weights based on these specifications, and then vary a single parameter multiplying the control weights. In this paper we propose the following alternative procedure: specify an asymptotic eigenstructure, choose a state feedback matrix based on these specifications, and then vary a single parameter multiplying this matrix. The alternative procedure does not have the advantages of guarenteed stability, phase and gain margins of the LQR; but is important none-the-less. Both procedures suffer the serious drawback of requiring full state feedback. A precursor to both of these procedures is in $[3,4]$ where state feedback is used to achieve
a non-asymptotic eigenstructure.
This paper starts with a review of the finite and infinite zezo structure of the open loop systam, and treate the infinite zeros as 11 defined quantities. The behavior of the closed loop eigenetructure is reviewed, and then an algorithm is presented to compute the asymptotically infinite eigenstructure. Next we show how a MiMO system can be decomposed into separate single-input single-output (SISO) systems that have the same asymptotically infinite eigenstructures. Then we present the previously mentioned algorithm for achieving a desired asymptotic elgenstructure with full state feedback, and then we finish with several examples.

Previous analysis techniques for multivariable root loci have appeared in $[5,6]$. In the former an algorithm is presented to compute the asymptotically infinite patterns of the closed loop eigenvalues. The first of our algorithms is similar to this, but differs importantiy in that we compute the asymptotic behavior of the eigenvectors and we introduce a subspace decomposition of $\mathbf{R}^{\mathrm{m}}$ which can be used decompose the syatem into SISO parts. In the latter reference the root iocus is interpreted as living on a Riemann surface. We have not used this approach because of computational reasons and because we are not yet convinced that an engineer trying to design a control system need concern himself with Riemann surfaces.

The linear systems we consider are restricted to having the same number of inputs and outputs, being controllable and observable, being nondegenerate, and having distinct finite zeros. Further restrictions are placed on the allowable asymptotically infinite behavior of the root locus.

In later work we hope to ranove some or all of these restrictions and to study in more detail the synthesis of output feedback.

## Notation

Matrices are denoted by capital letters, scalars and vectors by lower case letters. $A^{T}$ and $Y^{H}$ are the transpose of $A$ and the Hermitian transpose of $y$. Subspaces are denoted by script letters, with the exception of $R^{n}$. "Im A" and "ker $A$ " are the image and kernel of the linear map $A$. The dimension of $U$ is dim $U$, subspace inclusion is $C$, subspace intersection is $\cap$, and linear combination of subspaces is $U+V$.
II. Pole zero Configuration of an Open Loop System

We are concerned with the following linear time invariant aystem:

$$
\begin{align*}
& \dot{x}=A x+B u  \tag{1}\\
& y=C x \tag{2}
\end{align*}
$$

where

$$
\begin{aligned}
& x \in \mathrm{R}^{\mathrm{n}} \\
& \text { ue } \mathrm{m}^{\mathrm{m}} \\
& y e \mathrm{R}^{\mathrm{m}} .
\end{aligned}
$$

The number of inputs and outputs are equal. We assume that the realization is minimal, which is equivalent to assuming that ( $A, B$ ) is controllable and ( $A, C$ ) is observable. The open loop system has $n$ poles and $n$ zeros associated with it, and in this section we review their dynamic interpretations.

The open loop poles are the $n$ eigenvalues of A. The open loop poles are the complex frequencies that can appear in the output without appearing in the input.

There are $n$ zeros associated with the open loop system, of which $p$ are finite and $n-p$ are infinite. The zeros are sometimes called "transmission zeros." The finite zeros are defined to be those finite values of $s$ which reduce the rank of

$$
\left[\begin{array}{ll}
A-s I & B \\
-C & 0
\end{array}\right]
$$

We assume that the finite zeros are distinct and that the system is not degenerate in the sense that not all values of $s$ in the complex plane reduce the rank of the matrix. The finite zeros are the finite solutions $s_{i}^{0}$ of the generalized eigenvalue problem [7].

$$
\left[\begin{array}{cc}
A-s_{i}^{0} I & B  \tag{3}\\
-C & 0
\end{array}\right]\left[\begin{array}{l}
x_{1}^{0} \\
v_{i}^{0}
\end{array}\right]=0 \quad 1=1 \ldots \ldots p
$$

Under our assumptions the number of finite zeros is $0 \leq p \leq n-m$. Associated with each finite zero are right zero directions $x_{i}^{0}$ and $\nu_{i}^{0}$. It is also possible to define the generalized eigenvalue problem

$$
\left[y_{i}^{O H} n_{i}^{O H}\right]\left[\begin{array}{ll}
A-8_{i}^{O} I & B  \tag{4}\\
-C & 0
\end{array}\right]=0 \quad i=1, \ldots, p
$$

and associated with each finite zero are also the left zero directions $y_{i}^{\mathrm{OH}}$ and $\eta_{i}^{\mathrm{CH}}$. The finite zeros are the complex frequencies "absorbed" by the system in the following sense. If at time t=0 the system is at state $x(0)=x_{i}^{0}$, then an input of the form $u(t)=v_{i}^{0} e^{8_{i}^{t}}$ for $t \geq 0$ will result in $Y(t)=0$ for $t \geq 0$ [8].

The n-p zeroes at infinity are well defined and can be given the following interpretation [9]. If the input is of the form

$$
\begin{equation*}
u(t)=\sum_{i=0}^{\ell} u_{i} \delta^{(i)}(t) \tag{5}
\end{equation*}
$$

where $\delta^{0}(t)$ is an impulse, $\delta^{l}(t)$ is a doublet, and so on; and if the initial state of time $t=0^{-}$is

$$
\begin{equation*}
x\left(0^{-}\right)=-\sum_{i=0}^{\ell} x^{i} \mathrm{Bu}_{i} \tag{6}
\end{equation*}
$$

where $u_{0}$ is arbitrany and the remaining $u_{i}$ vectore satisfy

then the output $y(t)=0$ for $t \geq 0$. The large matrix in (7) is called a Toeplitz matrix. The integer $\ell$ is characteristic of the open loop aystem and will be interpreted in the next section.

## III. The Asymptotic Eigenstructure of the Closed Loop Syetem

We use output feedback of the form

$$
\begin{equation*}
u=-\frac{1}{k} x y \tag{8}
\end{equation*}
$$

where $K$ is an mom invertiable matrix and $k$ is a real number in the range $0 \leq k \leq \infty$. The closed loop system matrix if

$$
\begin{equation*}
A_{C l}-A-\frac{1}{k} B K C . \tag{9}
\end{equation*}
$$

and its eigenvalues, right, and left eigenvectors are defined in the usual way by

$$
\begin{array}{ll}
\left(A_{c l}-s_{i} I\right) x_{i}=0 & i=1, \ldots, n \\
y_{i}^{H}\left(A_{c l}-n_{i} I\right)=0 & i=1, \ldots, n \tag{11}
\end{array}
$$

As $k$ is varied from infinity down to zero the $n$ closed loop eigenvalues trace out a root locus on the complex plane and the right and left eigenvectors "spin" in $\mathbf{R}^{n}$. We are particularly interested in asymptotic behavior as $k+0$, which we now review.

As $k \rightarrow 0$ p of the $n$ branches of the root locus approach finite zeros. The right and left eigenvectors associated with the eigenvalues on the finite branches approach the right and left zero directions. The finite zeros and zero directions can be computed using the generalized eigenvalue problems $(3,4)$, and for notational convience we group the solutions into

$$
s^{0}=\operatorname{diag}\left(s_{1}^{0}, \ldots, s_{p}^{0}\right)
$$

$$
\begin{aligned}
& x^{0}=\left[x_{2}^{0} \ldots \ldots, x_{p}^{0}\right\} \\
& \mathbf{x}^{0}=\left[y_{1}^{0} \ldots, y_{p}^{0}\right\} .
\end{aligned}
$$

The ramaining $n-p$ branches approach the infinite zeros. The $n-p$ right eigenvectors associated with the eigenvalues on the infinite branches will anymptotically span the same subspace of $\mathrm{In}_{\mathrm{n}}$ spanned by all possible $x\left(0^{-}\right)$of (6), and a similar property can be derived for the 1eft eiganvectors.

The information we have given so far about the finite and infinite eigenstructure can be computed knowing only the open loop aystom, and does not depend on the output faedback gain matrix $K$. Unfortunately this information is not mough for a good analysis of a control mystam, and thia is especially true for an analyaie of the asymptotically infinite aigenatructure. In the rest of this section we list properties of the asymptoticaliy infinite digenstructure, most of which dapend on $K$, and in the next zaction we show how these prope:tiles can be computed.

The infinite branches of the roct locus break into $m$ patterns. The order of the $i{ }^{\text {th }}$ pattern is $n_{i}$, where "order" is defined to be the number of closed loop eigenvalues in the pattern. The following identity must be true:

$$
\sum_{i=1}^{m} n_{i}=n-p
$$

and the highest order pattern is $r$. It turns out that $r=\ell+1$, where $i$ is given in $(5,6,7)$. There are $n_{i}$ asymptotes in the $i^{\text {th }}$ pattern. The closed loop eigenvalues approach these asymptotes and have a magnitude approximately equal to

$$
\left(\frac{\left|e_{1}^{\infty}\right|}{k}\right)^{\frac{1}{n_{1}}}
$$

The asymptotes are spaced $\left(360 / n_{i}\right)^{\prime}$ apart, arks angles with the positive real axis of

$$
\left(\frac{\arg \left(-s_{i}^{\infty}\right)+\ell 360}{n_{i}}\right)^{\bullet} \quad \ell=0,1, \ldots, n_{i}-1 .
$$

and have a center of gravity $\psi_{i}$. The $n_{i}$ right and left eigenvectors associated with the $i^{\text {th }}$ pattern asymptotically span the same subspaces of $x^{n}$ as

$$
\begin{aligned}
& B v_{i}^{\infty}, A B v_{i}^{\infty}, \ldots, A^{n_{i}-1} B V_{i}^{\infty} \\
& C^{T} n_{1}^{\infty},(A C)^{T} n_{i}^{\infty}, \ldots,\left(A^{n_{i}-1} C\right)^{T} n_{i}^{\infty}
\end{aligned}
$$

For notational convenience we group the asymptotic properties into

$$
\begin{aligned}
s^{\infty}= & \operatorname{diag}\left(s_{1}^{\infty}, \ldots, s_{m}^{\infty}\right) \\
N^{\infty}= & {\left[\nu_{1}^{\infty}, \ldots, \nu_{m}^{\infty}\right] } \\
\mathcal{N}^{\infty}= & {\left[\eta_{1}^{\infty} \ldots, \eta_{m}^{\infty}\right] } \\
\psi= & \left(\psi_{1}, \ldots, \psi_{m}\right) \\
i= & \left(01,11,21, \ldots,\left[n_{1}-1\right] 1,\right. \\
& \left.\left.02,12,22, \ldots,\left[n_{2}-1\right] 2, \ldots\left(n_{m}-1\right] m\right) .\right]
\end{aligned}
$$

Each part if of the malti-index $\gamma$ describes the $\lambda^{i} v_{j}^{\infty}$ vector. The complex numbers $s_{i}^{\infty}$ and $j_{i}^{\infty}$ and the complex m-vectore $\nu_{i}^{\infty}$ and $\eta_{i}^{\infty}$ must occur in complex conjugate paire.
IV. An Alcorithm to Compute the Anymptotically Infinite Eipenatructure

The problem of finding the asymptotically infinite eigenstructure
 $Y$. We firat 90 through a anries of definitions and then solve the problem. The algorithem can be programmed on a computer uaing the stable numarical algorithes of EIspAcx [10] (apecifically ainguiar velue decomposition (gVD) and the aigenvalue problem), but we make no cladian that our algorithm is the "fastant" or "best". We apologize for the notational nighemare that followa, but we do not yet know of any way to avoid it.
. Define the matrices

$$
c_{i}=c A^{i-1_{B}} \quad i=1, \ldots, r
$$

where $r$ is the highest order pattern, and define the subspaces of $\mathbb{R}^{n}$

$$
\begin{aligned}
& U_{0}=\mathbb{R}^{m} \\
& L_{1}=\operatorname{ker} G_{1} \\
& \vdots \\
& U_{x}=\operatorname{ker} G_{1} \cap \ldots \cap \operatorname{ker} G_{r} \\
& V_{0}=\mathbb{R}^{m} \\
& V_{1}=\operatorname{ker} G_{1}^{T} \\
& \vdots \\
& V_{r}=\operatorname{ker} G_{1}^{T} \cap \ldots \cap \operatorname{ker} G_{r}^{T} .
\end{aligned}
$$

$$
\begin{aligned}
& 0=u_{2} \subseteq \cdots \subseteq u_{0}=2^{m} \\
& 0=v_{5} \subseteq \cdots \subseteq v_{0}=\underbrace{m} .
\end{aligned}
$$

and their dimensions are

$$
l_{1}=\operatorname{dim} U_{1}=\operatorname{din} V_{i} \quad i=0, \ldots, x
$$

Define the indices

$$
m_{i}=l_{i-1}-l_{i} \quad 1=1, \ldots, x
$$

and it follows that

$$
\sum_{i=1}^{n} m_{i}=m
$$

We restrict our attention to system for which $m_{i}$ is the number of $i$ th order patterns and $l_{i}$ is the number of patterns greater than $i^{\text {th }}$ order. Sext we define the matrices

$$
v_{1}, v_{1} \quad 1=0,1, \ldots, 5-1
$$

where the columns of $U_{i}$ form basis for $U_{i}$ and the column of $v_{i}$ form a basis for $V_{i}$. The dimensions of both $U_{i}$ and $V_{i}$ are ma ${ }_{i}$. we decompose $8^{\infty}, N^{\infty}$, and $M^{\infty}$ into

$$
\begin{aligned}
& s^{\infty}=\operatorname{diag}\left(s_{1}^{\infty}, \ldots, s_{r}^{\infty}\right) \\
& N^{\infty}=\left[N_{2}^{\infty}, \ldots, N_{r}^{\infty}\right] \\
& M^{\infty}=\left[M_{1}^{\infty} \ldots \ldots M_{r}^{\infty}\right] \quad
\end{aligned}
$$

Is general each $S_{i}^{\infty}$ is a $m_{i} x_{i}$ block diagonal matrix, but we restrict our attention to systems for which each $S_{i}^{\infty}$ is diagonal. The diagonal elements
are the $s_{j}^{\infty} s$ associated with $i^{\text {th }}$ order patterns. Each of the $\mathbb{N}_{i}^{\infty}$ and $M_{i}^{\infty}$ matrices have dimensions $\mathrm{mxm}_{i}$ and have as columns the $v_{j}^{\infty}$ 's and $\eta_{j}^{\infty}$ 's associated with $i^{\text {th }}$ order patterns. As we will see later,

$$
\left.\begin{array}{l}
u_{i}=\operatorname{Im} N_{i+1}^{\infty}+\ldots+\operatorname{Im} N_{r}^{\infty} \\
U_{i}=\operatorname{Im} M_{i+1}^{\infty}+\ldots+\operatorname{Im} M_{r}^{\infty}
\end{array}\right\}\left\{\begin{array}{l}
1=0, \ldots, r-1
\end{array}\right.
$$

and in general

$$
\left.\begin{array}{l}
\operatorname{Im} \mathbb{N}_{i}^{\infty} \cap \operatorname{Im} \mathbb{N}_{j}^{\infty} \neq 0 \\
\operatorname{Im} M_{i} \cap \operatorname{Im} M_{j}^{\infty} \neq 0
\end{array}\right\} \quad i \neq j
$$

In words the top relationships tell us that one of the possible sets of basis vectors of $U_{i}$ is the $\nu_{j}^{\infty}$ 's associated with patterns of greater than $i^{\text {th }}$ order, and that in general the $\nu_{j}^{\infty}$ 's are not orthogonal. A similar statement can be made about the $\eta_{j}^{\infty}$ 's. The last definitions are

$$
T_{i}=\left(v_{i-1}^{H} K^{-1} U_{i-1}\right)^{-1}\left(v_{i-1}^{H} G_{i} U_{i-1}\right) \quad i=1, \ldots, r
$$

and the Jordan form decompositions

$$
T_{i}=\left[\begin{array}{ll}
W_{i 1} & w_{i 2}
\end{array}\right]\left[\begin{array}{ll}
\Lambda_{i} & 0 \\
0 & 0
\end{array}\right]\left[\begin{array}{c}
W_{i 3}^{H} \\
W_{i 4}^{H}
\end{array}\right] \quad i=1, \ldots . r .
$$

The matrices $W_{i 1}$ and $W_{i 3}$ have as many columns as there are $i^{\text {th }}$ order patterns, and $W_{i 2}$ and $W_{i 4}$ have as many columns as there are greater than $i^{\text {th }}$ order patterns. The dimensions of $T_{i}$ are $\ell_{i-1} \times \ell_{i-1}$, $W_{i 1}$ and $W_{i 3}$ are $\ell_{i-1} \times m_{i}$, and $W_{i 2}$ and $W_{i 4}$ are $\ell_{i-1} \times \ell_{i}$. Finally
we conclude this long paragraph with the fact that if there are no ith order patterns then $m_{1}=0, l_{1-1}=l_{i}, u_{1-1}=u_{i}, V_{1-1}=V_{1}, T_{i}=0$, and the $s_{i}^{\infty}, N_{i}^{\infty}, N_{i}^{\infty}, N_{11}$, and $W_{13}$ matrices vanish.

The algorithm to find $s^{\infty}, N^{\infty}, N^{\infty}, Y$, and $\psi$ is:

1) Use SWD to compute $v_{i}, v_{i}$, and $m_{i}$ for $i=1, \ldots, t$.
2) Use $m_{i}$ for $1-1, \ldots, r$ to compute $\gamma$.
3) Coupute $x_{i}$ for $i=1, \ldots, x$.
4) Compute the Jordan form decompositions of $\mathrm{T}_{\mathrm{i}}$ for $\mathrm{iml}, \ldots, \mathrm{r}$.
5) Compute $s_{i}^{\infty}, N_{i}^{\infty}$, and $M_{i}^{\infty}$ for $1=1, \ldots, r$ by

$$
\begin{aligned}
& s_{i}^{\infty}=\Lambda_{i} \\
& N_{i}^{\infty}=U_{i-1} W_{11} \\
& M_{i}^{\infty}=v_{i-1} N_{13}
\end{aligned}
$$

6) Compute $\psi_{i}$ for $1=1, \ldots, m$ by

$$
\begin{aligned}
\psi_{i} & =\frac{\beta_{i}}{n_{i}}, \\
\text { where } B_{i} & =\frac{\eta_{i}^{\infty^{H}} C A^{n_{1}}{ }_{B v_{i}^{\infty}}}{n_{i}^{\infty^{H} C A^{n}} n_{i}^{-1} B v_{i}^{\infty}}
\end{aligned}
$$

The following explanation of the first step may be helpful. The columns of $U_{i}$ form a basis for the karnel of

$$
\left[\begin{array}{c}
C B \\
\vdots \\
C A^{1} B
\end{array}\right] \quad i=1, \ldots, r-1
$$

$s 0$ to compute $U_{i}$ we can use SVD to compute an orthonormal basis of this tall skinny matrix. Likewise $V_{i}$ can be found by computing the kernel of

$$
\left[\begin{array}{c}
(C B)^{T} \\
\vdots \\
\left(C A^{i} B\right)^{T}
\end{array}\right]
$$

$$
i=1, \ldots, r-1 .
$$

It is not necessary to fill $U_{i}$ and $V_{i}$ with orthonormal vectors, but this is a convenient by-product of SVD. There are many other ways to compute the kernel of a matrix, and $S V D$ is almost certainly not the method to use when computing the $U_{i}$ 's and $v_{i}$ 's by hand. However, there may be difficulty determining the dim $U_{i}$, and SVD is the most reliable way to do this. If there is difficulty determining the dim $U_{i}$ then the root locus may have "strange" behavior such as asymptotically infinite patterns that shift cicers at large radii.

We note that we have used the last column of a Toeplitz matrix to find the order of each of the asymptotically infinite root locus patterns. We have restricted our attention to systems for which this can be done. In general the entire Toeplitz matrix must be used.

The generic case is when Rank (CB) $=m$. This should be viewed as a mathematical property, see [11] for a precise definition of "generic." If the system to be analyzed is generic then it has $n-m$ transmission zeros
-17-
and $m$ firat order infinite patterns. The $s_{i}^{\infty} s, v_{i}^{\infty} \prime s$, and $n_{i}^{\infty}$ make up the eigenstructure of KCB. For design purposes the generic case is too restrictive to be of interest. This is easient to see for SISO systems, for which root loci with second and higher order infinite patterns are commonplace.

## V. Proof of Analyais Algorithm

In this section we prove the fifth step of the algorithm to find the asymptotically infinite eigenstructure. The proof of the sixth step we defer until the next section. The first four steps do not need to be proved.

The proof of the first stop is by induction and uses the fact that all of the closed loop eiganvalues $s_{j}$ and the associated $v_{j}$ vectors must satisfy

$$
\begin{equation*}
\left[k x^{-1}+\Phi\left(s_{j}\right)\right] v_{j}=0 \quad j=1, \ldots, n \tag{12}
\end{equation*}
$$

where

$$
\begin{aligned}
\Phi(s) & =C(s I-A)^{-1} B \\
& =\sum_{i=1}^{\infty} \frac{1}{s^{i}} G_{i}
\end{aligned}
$$

To show that (12) is true we note from ( $1,2,8,9,10$, and 11) that

$$
\left[\begin{array}{ll}
\lambda_{j}-s_{j}^{I} & B \\
-C & -k k^{-1}
\end{array}\right] \quad\left[\begin{array}{l}
x_{j} \\
v_{j}
\end{array}\right]=0 \quad j=1, \ldots . n
$$

and therefore for $j=1, \ldots, n$ we have that

$$
\begin{aligned}
& \left(A-s_{j} I\right) x_{j}+B V_{j}=0 \\
& C x_{j}=-k K^{-1} \nu_{j}=C\left(s_{j}^{I-A}\right)^{-1} B V_{j} \\
& {\left[k x^{-1}+C\left(s_{j} I-A\right)^{-1} B\right] V_{j}=0 .}
\end{aligned}
$$

The first step of the induction proof is to show that the fifth step is valid for $S_{1}^{\infty}, N_{1}^{\infty}$, and $N_{1}^{\infty}$. We assume without loss of generality that first order patterns exist. Equation (12) can be rewritten

$$
\left[k x^{-1}+\frac{1}{s_{j}} G_{1}+o\left(\frac{1}{s_{j}}\right)\right] v_{j}=0 \quad j=1, \ldots, n
$$

As $k \rightarrow 0$ this becomes

$$
\left[8_{j}^{\infty} I-T_{1}\right] \nu_{j}^{\infty}=0 \quad j=1, \ldots, m
$$

where

$$
s_{j}^{\infty}=-k s_{j}
$$

The nonzero eigenvalues of $T_{1}$ correspond to first order patterns because the closed loop eigenvalues $s_{j}$ are solutions of

$$
s_{j}=\frac{-s_{j}^{\infty}}{k}
$$

The single solution $s_{j}$ traces out a first order pattern. The right and left eigenvectors of $T_{1}$ are the $\nu_{j}^{\infty}$ 's and $\eta_{j}^{\infty}$ 's associated with first order patterns, and therefore the fifth step is true for $i=1$. The $V_{j}^{\infty}$ 's associated with second and higher order pattern: lie in the kernel of $T_{1}$, which is $U_{1}$. Heuristically speaking, there $v_{j}^{\infty}$ 's are not "trapped" by the $s_{j}^{-1}$ term.

The next step in the induction proof is to assume that $s_{i-1}^{\infty}, N_{i-1}^{\infty}$, and $M_{i-1}^{\infty}$ are valid and then show that $s_{i}^{\infty}, N_{i}^{\infty}$, and $M_{i}^{\infty}$ are valid. If $N_{i-1}^{\infty}$ is valid then the $v_{j}^{\infty}$ vectors associated with $\geq i^{\text {th }}$ order patterns form a basis for $u_{i-1}$. Therefore for each of these $\nu_{j}^{\infty}$ 's there exists a $\omega_{j}$ such that $\nu_{j}^{\infty}=U_{i-1} \omega_{j}$. Substituting into (12) we get

$$
\left[k x^{-1}+\Phi\left(\varepsilon_{j}\right)\right] u_{i-1} \omega_{j}=0 \quad j=1, \ldots, \ell_{i-1}
$$

Multiply on the left to get

$$
v_{i-1}^{H}\left[k x^{-1}+\Phi\left(s_{j}\right)\right] u_{i-1} \omega_{j}=0 \quad j=1, \ldots, l_{i-1},
$$

which reduces to

$$
\begin{aligned}
& {\left[x v_{i-1}^{H} X^{-1} D_{i-1}+\frac{1}{E_{j}^{1}} v_{i-1}^{H} G_{i} V_{i-1}+o\left(\frac{1}{v_{j}^{1}}\right)\right] \omega_{j}=0} \\
& j=1, \ldots, l_{i-1} .
\end{aligned}
$$

As $k=0$ this becomes

$$
\left(s_{j}^{\infty} I-T_{j}\right) \omega_{j}=0 \quad j=1, \ldots, l_{i-1}
$$

where

$$
s_{j}^{\infty}=-k s_{j}^{i}
$$

The nonzero eigenvalues of $T_{i}$ correspond to $i^{\text {th }}$ ordex patterns because the closed loop eigenvalues $s_{j}$ are the $i$ solutions of

$$
s_{j}^{i}=\frac{-s_{j}^{\infty}}{k}
$$

The $\omega_{j}$ 's can be used to compute $v_{j}^{\infty}=v_{i-1} \omega_{j}$ and therefore $N_{i 1}^{\infty}=U_{i-1} W_{i 1}$. Using a similar argument $M_{i}^{\infty}=V_{i-1} W_{i 3}$. The $V_{j}^{\infty}$ 's associated with greater than $i^{\text {th }}$ order patterns are not "trapped" by the $s_{j}^{-i}$ term of (12) and therefore lie in $\mathrm{U}_{1}$. This completes the proof.

## 

There are m anymptotically infinite patterns of the root loous, and as we will now show we can decompose the MINO systam into m SISO aystems, each of which has one of the infinite patterns. After doing this we will prove the sixth step of the analyais algorithm (finding the center of gravity of the infinite patterns).

Use the $\nu_{i}^{\infty}$ and $\eta_{i}^{\infty}$ vectors to define the following sIsO systems:

$$
\begin{aligned}
& \left.\begin{array}{l}
\dot{x}=A x+b_{i} u \\
y=c_{i} x \\
u=-\frac{1}{k_{i}} y
\end{array}\right\} \\
& 1=1, \ldots, m
\end{aligned}
$$

where

$$
\begin{aligned}
& b_{i}=B v_{i}^{\infty} \\
& c_{i}=\eta_{i}^{\infty E_{i}} c \\
& k_{i}=k n_{i}^{\infty} x^{-1} v_{i}^{\infty} .
\end{aligned}
$$

The return difference equation of each of the SISO systems is $1+g_{i}(s) / k_{i}$ where

$$
\begin{aligned}
g_{i}(s) & =c_{i}(s I-A)^{-1} b_{i} \\
& =\sum_{j=n_{i}}^{\infty} \frac{1}{j} n_{i}^{\infty} C_{A}^{j-1} B V_{i}^{\infty} \quad i=1, \ldots, m .
\end{aligned}
$$

Since the closed loop poles are the zeroes of the return difference equation we have that

$$
1+g_{i}(a) / k_{1}=0 \quad 1=1, \ldots, m
$$

which can be rowritten as

$$
\frac{k_{i}}{g_{i}(s)}+1=0 \quad i=1, \ldots, m
$$

Carry out the long division and the result is

$$
\begin{equation*}
n^{n_{i}}-\beta_{i} n^{n_{i}-1}+\ldots+\frac{1}{x} s_{i}^{\infty}+\ldots=0 \quad i=1, \ldots, m \tag{13}
\end{equation*}
$$

whare

$$
\begin{aligned}
& B_{i}=\frac{n_{1}^{\infty} A^{n_{1}} B v_{i}^{\infty}}{n_{i}^{\infty} C A^{n_{1}-1} B v_{i}^{\infty}} \\
& \varepsilon_{i}^{\infty}=\frac{n_{1}^{\infty} C^{n_{i}-1} B v_{i}^{\infty}}{n_{1}^{\infty} x^{-1} v_{i}^{\infty}}
\end{aligned}
$$

To verify that the amymptotically infinite pattern of each of the SISO systems is the same as one of the patterns of the larger system we need only note that as $k+0$ (13) can be approximated by

$$
s_{i}^{n_{i}}=\frac{-1}{k} e_{i}^{\infty} \quad i=1, \ldots, m
$$

To verify the sixth step of the analysis algorithm of sestion IV we approximate (13) by

$$
s^{n_{i}}-\beta_{i} s^{n_{i}-1}+\frac{1}{k} s_{i}^{\infty}=0 \quad i=1, \ldots, m \quad .
$$

The sum of the $n_{i}$ infinite eiganvaiuas in this pattern munt equal $B_{i}$ and therefore the center of gravity is

$$
\psi_{1}=\frac{\beta_{i}}{n_{i}} \quad i=1, \ldots, m
$$

For $n_{i}=1$ the term "center of gravity" could be more appropriately replaced by "starting point."

Some of the SISO systems may have complex coefficients. Whan this happens a second zystem exists which has the complex conjugate coefficients. To work with only real coefficients a $2 \times 2$ system must be used. The closed loop eigenvalues of the combined asymptotic pattern are solutions of

$$
\left(s^{n_{i}}+\frac{1}{k} \varepsilon_{i}^{\infty}\right)\left(s^{n_{i}}+\frac{1}{i}-_{i}^{\infty}\right)=0 .
$$

which is

$$
s^{2 n_{i}}+\frac{2}{k} \operatorname{Re}\left(s_{i}\right) s^{n_{i}}+\frac{1}{k^{2}}\left|s_{i}^{\infty}\right|^{2}=0 .
$$

The center of gravity of the combined asymptotic pattern can be found to be Re $\left(\psi_{1}\right)$.

## VII. An Alcorith to synthegin an Agyototic Bigenatructure with Tull

 state FecelbackWe considar a systean with full state foedback:

$$
\begin{aligned}
& \dot{x}=x x+2 x \\
& x-\frac{1}{x} I x
\end{aligned}
$$

whare

$$
x \in \quad \frac{p}{2}
$$

$40 \boldsymbol{m}^{m}$ 。

We assume that $(\lambda, B)$ is controllable. The problen is to choose a feedback matrix $F$ auch that as $k+0$ the closed loop syatam has achieved $a$ described eigenstructure. Concisely stated the problem is: Given $s^{0}, x^{0}, s^{\infty}, R^{\infty}$, and $Y$ sind $F$.

The colution is

$$
\begin{equation*}
F=\Delta^{\infty} s^{\infty}[0 \quad I]\left[x^{0}, 0^{\gamma}\right]^{-1} \tag{14}
\end{equation*}
$$

where

$$
\begin{aligned}
U^{X}= & {\left[B v_{1}^{\infty}, \lambda B v_{1}^{\infty}, \ldots, A^{n_{1}-1} B v_{1}^{\infty}\right.} \\
& \left.B v_{2}^{\infty}, \lambda B v_{2}^{\infty}, \ldots, A^{n_{2}-1} B v_{2}^{\infty} \ldots \ldots A^{n_{m}-1} B v_{m}^{n}\right]
\end{aligned}
$$

$P$ = Parmutation matrix that rearranges the colums of $0^{\gamma}$ in an arbitary way except that the last m columas of $U^{\gamma}$ are

$$
A^{n_{1}^{-1}} B v_{1}^{\infty}, A^{n_{2}-1} B v_{2}^{\infty}, \ldots, A^{n_{m}^{-1}} B v_{m}^{\infty}
$$

In general the $X^{0}, s^{\infty}, \forall^{\infty}$, and $U^{\gamma}$ contain complax numbers. To work juat with real numbers then raplece the complex conjugate colume of $x^{0}$. $\|^{\omega \prime}$, and $0^{\gamma}$ with thair real and imaginary parte, and raplace the $2 \times 2$ blocks

$$
\left[\begin{array}{ll}
s_{i} & \\
& \\
& s_{i}
\end{array}\right]
$$

in $8^{\infty} \mathrm{by}$

$$
\left[\begin{array}{cccc}
\operatorname{sen} & s_{i}^{\infty} & \min & s_{i}^{\infty} \\
-m & s_{i} & \operatorname{se} & x_{i}^{\infty}
\end{array}\right]
$$

One way to interpret this result is that state feedback is used to place the finite and infinice zaros and zero directions, and an the feedback gain is increased the cloced loop eigenvalues and eigenvectors approach these zeros and zero directions.

The deaired asymptotic eigenstructure is not completely arbitrary. We restrict our attention to cases where the finite zaros are distinct and the $s^{\infty}$ matrix is diagonal. More fundamental are the restrictions that
 conjugate pairs;
(2) Eor each zero direction $x_{i}^{0}$ there must exist a $v_{i}^{0}$ such that $\left(\lambda-B_{i}^{0} I\right) x_{i}^{0}+B v_{i}^{0}=0_{i}$
(3) the number of finite zeros must be $0 \leq p \leq n-m$,
(4) the columss of $X^{0}$ and $\mathrm{U}^{\gamma}$ must be linearly independent; and
(5) the multi-index $\gamma$ must apecify m acymptotically infinite patterns whose ordars totel to $\mathrm{n}-\mathrm{p}$.

Io prove that (14) is true we comprite the anympotic elgenatzucture using the previowe methode deseribed in this paper. Once I is sixed we can treat the atate feecback problen as an output facoback probion

$$
\begin{aligned}
& \dot{x}=\lambda x+80 \\
& y=I x \\
& u=-\frac{1}{x} I y .
\end{aligned}
$$

The asymptoticaliy Einite oigenstructure mat satisiy

$$
\left[\begin{array}{cc}
A-s_{1}^{0} & 1 \\
-F & 0
\end{array}\right]\left[\begin{array}{l}
x_{1}^{0} \\
v_{1}^{0}
\end{array}\right]=0 \quad 1=1, \ldots . p .
$$

This is tzue because wo have ascumed that a $v_{1}^{0}$ exists such that $\left(A-8_{i}^{0} I\right) x_{i}^{0}+8 \nu_{i}^{0}=0$, and becauce $I$ is constructed in such a way that $I x_{i}^{0}=0$. naxt we show (for $i=1, \ldots$, m) that the $i^{\text {th }}$ anymptotically infinite pattern is $n_{i}^{\text {th }}$ order and has $s_{i}^{\infty}$ and $v_{i}^{\infty}$ ascociated with it. By the way $F$ is constructed $F 0^{\gamma_{P}}=\left[0 \quad x^{\infty} 8^{\infty}\right]$ and in particular $E \lambda^{n_{1}^{-1}} 8 v_{1}^{\omega}=v_{i}^{\infty} \varepsilon_{1}^{\infty}$. so if we start with

$$
\begin{equation*}
\left[k I+\emptyset\left(s_{1}\right)\right] v_{i}=0 \quad i=1, \ldots, n \tag{15}
\end{equation*}
$$

where

$$
\Phi(s)=F(s I-\lambda)^{-1} s \text {, and if we let } v_{1}=v_{1}^{\infty} \text { for } i=1, \ldots, m,
$$ then as $x \rightarrow 0$ (15) can be rewritten

$$
\begin{equation*}
\left[k I+\frac{s_{i}^{\infty}}{n_{i}} I\right] v_{i}=0 \quad i=1, \ldots \ldots m \tag{16}
\end{equation*}
$$

Equation (16) is terve $12 \varepsilon_{i}^{n_{1}}=-\xi_{i}^{\infty} / k$, which is the equation for an $i^{\text {th }}$ order pattern. This complatea the proof of the synthasis alyorithm.

We note that once wh have ohosen $\mathrm{s}^{0}, \mathrm{x}^{0}, \mathrm{~B}^{\infty}, \mathrm{m}^{0}$, and $\gamma$ thare does not exiat any extra frection to place the center of gravitios $\psi_{i}$. This is aasiest to see for the siso onse because it in a well known clagsical zoot locus reault that the open 100 p polen and Elnite seron dotezaine the center of gravity of the single ammpotionily infinite pattern.

Equation (14) can be simplitied in the generic case (Rank (CB) m). Then we have $0^{\gamma}=33^{\circ \prime}$ and (24) beccene

$$
y=x^{\infty} s^{\infty}\left(x^{\infty}\right)^{-1}\left[\begin{array}{ll}
0 & I]\left[x^{0} 3\right]^{-1} . \tag{17}
\end{array}\right.
$$

Equation (14) can be simplified even more in the 8180 case when the state apace realization is in controllable canonical form. For 8 sso system there is only one aymptotically infinite pattern, it suffices to set $V_{1}=1$, and there does not exist any Eroedion to place the sero directions $x_{i}^{0}$, $1=1, \ldots, p$. The procedure is to choose the finite zeros $s_{i}^{0}$ for $1=1, \ldots, p$ where $0 \leq p \leq n-1$; then form the polynominal

$$
\prod_{i=1}^{p}\left(s-s_{i}^{0}\right)=s^{p}+\beta_{p-1} s^{p-1}+\ldots+B_{0},
$$

and then the state Eoadback matrix is

$$
\begin{equation*}
F=\left[\beta_{0}, \beta_{1}, \ldots, \beta_{p-1}, 0, \ldots, 0\right] . \tag{18}
\end{equation*}
$$

VIII. Examples

In the first example we analyze the asymptotically infinite eigenstructure of a linear system. Then we use full state feedback to achieve a specified asymptotic aigenstructure.

## Example 1

We use the same system as example 2 of [5]. Given the following $A, B, C$, and $K$ matrices we use the algorithm of section IV to find the asymptotically infinite eigenstructure.

$$
\begin{aligned}
& A=\left[\begin{array}{ccccccr}
12 & -60 & 28 & -4 & 32 & 36 & -54 \\
28 & -75 & 12 & 30 & 0 & 45 & -92 \\
-8 & -134 & 33 & 48 & 12 & 71 & 0 \\
-8 & -14 & -6 & 21 & -18 & 5 & 24 \\
-4 & 62 & -21 & -12 & -18 & -35 & 24 \\
64 & -90 & 18 & 24 & 12 & 60 & -200 \\
9 & -22 & 10 & -2 & 12 & 14 & -33
\end{array}\right] \\
& B=\left[\begin{array}{lll}
6 & 0 & 8 \\
2 & 1 & 3 \\
6 & 4 & 0 \\
0 & 0 & 1 \\
-3 & -2 & 1 \\
2 & 1 & 5 \\
2 & 0 & 3
\end{array}\right] C^{T}=\left[\begin{array}{rrrr}
-3 & 2 & 3 \\
11 & 7 & -1 \\
1 & 4 & 1 \\
-9 & -11 & -1 \\
5 & 10 & 2 \\
-6 & -3 & 1 \\
9 & -7 & -9
\end{array}\right]
\end{aligned}
$$

$$
x=\left[\begin{array}{lll}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right]
$$

The open loop eigenvalues are $-3,-2,-1,0,1,2,3$. To find the $v_{i}$ and $v_{i}$ matrices we need

$$
\begin{array}{ll}
C B=\left[\begin{array}{ccc}
1 & -1 & 2 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] & C A B=\left[\begin{array}{ccc}
-2 & 3 & -5 \\
0 & 1 & -1 \\
0 & 0 & 0
\end{array}\right] \\
C A^{2} B=\left[\begin{array}{ccc}
4 & -5 & 17 \\
0 & -1 & -7 \\
0 & 0 & -8
\end{array}\right] & C A^{3} B=\left[\begin{array}{ccc}
-8 & -3 & -23 \\
0 & -11 & 35 \\
0 & 0 & 24
\end{array}\right]
\end{array}
$$

It can easily be verified that

$$
\begin{array}{ll}
u_{1}=\left[\begin{array}{ll}
1 & 0 \\
1 & 2 \\
0 & 1
\end{array}\right] & u_{2}=\left[\begin{array}{c}
-1 \\
1 \\
1
\end{array}\right] \\
v_{1}=\left[\begin{array}{ll}
0 & 0 \\
1 & 0 \\
-2 & 1
\end{array}\right] & v_{2}=\left[\begin{array}{l}
0 \\
0 \\
1
\end{array}\right] \\
m_{1}=1 & m_{2}=1
\end{array}
$$

At this point we know there will be a first, second, and third order infinite pattern. The multi-index $\gamma$ is therefore

$$
\gamma=(01,02,12,03,13,23)
$$

where

$$
n_{1}=1, n_{2}=2, \text { and } n_{3}=3
$$

The $T_{i}$ 's and their Jordan form decompositions are

$$
\begin{aligned}
& T_{1}=\operatorname{XCB}=\left[\begin{array}{ccc}
1 & -1 & 2 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right] \\
& =\left[\begin{array}{ccc}
1 & 1 & -1 \\
0 & 1 & 1 \\
0 & 0 & 1
\end{array}\right]\left[\begin{array}{lll}
1 & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right]\left[\begin{array}{ccc}
1 & -1 & 2 \\
0 & 1 & -1 \\
0 & 0 & 1
\end{array}\right] \\
& T_{2}=\left(v_{1}^{H} X^{-1} U_{1}\right)^{-1} v_{1}^{H} C A B U_{1}=\left[\begin{array}{ll}
1 & 0 \\
1 & 0
\end{array}\right] \\
& =\left[\begin{array}{cc}
1 & 1 \\
0 & -1
\end{array}\right]\left[\begin{array}{ll}
1 & 0 \\
0 & 0
\end{array}\right]\left[\begin{array}{ll}
1 & 1 \\
0 & -1
\end{array}\right] \\
& T_{3}=\left(V_{2}^{H} X^{-1} U_{2}\right)^{-1} V_{2}^{H} C^{2} B_{2}=-8 .
\end{aligned}
$$

-31-

Therefore we have that

$$
\begin{array}{ll}
s_{1}^{\infty}=1 & v_{1}^{\infty}=\left[\begin{array}{l}
1 \\
0 \\
0
\end{array}\right] \quad n_{1}^{\infty}=\left[\begin{array}{r}
1 \\
-1 \\
2
\end{array}\right] \\
s_{2}^{\infty}=1 & v_{2}^{\infty}=v_{1} W_{21}=\left[\begin{array}{l}
1 \\
1 \\
0
\end{array}\right] \quad n_{2}^{\infty}=v_{1} w_{23}=\left[\begin{array}{c}
0 \\
1 \\
-1
\end{array}\right] \\
s_{3}^{\infty}=-8 & v_{3}^{\infty}=\left[\begin{array}{r}
-1 \\
1 \\
1
\end{array}\right] \quad
\end{array}
$$

The center of gravities are

$$
\begin{aligned}
& \psi_{1}=-2 \\
& \psi_{2}=-1 / 2 \\
& \psi_{3}=-1
\end{aligned}
$$

In sumary, there is one first order pattern along the negative real axis with a radius of $k^{-1}$; there is one second order pattern with angles $\pm 90^{\circ}$, radius $k^{-1 / 2}$, and center of gravity $-1 / 2$; and there is one third order pattern with angles $0^{\circ}, \pm 120^{\circ}$, radius $2 k^{-1 / 2}$, and center of gravity -1.

## Example 2

consider a system with

$$
A=\left[\begin{array}{cccc}
0 & 1 & 0 & 1 \\
-5 & -4 & 0.1 & 1 \\
0.1 & 0 & -1 & 1 \\
0 & 0 & 0 & -5
\end{array}\right] \quad B=\left[\begin{array}{ll}
0 & 0 \\
0 & 0 \\
1 & 0 \\
0 & 1
\end{array}\right]
$$

For three different cases we specify the asymptotically infinite eigenstructure and then compute the state feedback matrix $F$. Since this is a generic example we use equation (17) to compute $F$. The asymptotically finite eigenstructure is the same in each case.

$$
s_{1,2}^{0}=-3 \pm 2 i \quad x_{1,2}^{0}=\left[\begin{array}{c}
0 \\
2 \\
0 \\
-4
\end{array}\right] \pm\left[\begin{array}{c}
-1 \\
3 \\
0 \\
2
\end{array}\right] i
$$

The infinite eigenstructure differs for each case. The specifications and the resulting $F$ matrices are shown in Table 1. The root loci are shown in Figure 1.

In the first case we have $s_{1}^{\infty}=1$ and $s_{2}^{\infty}=2$. The two first order patterns stay on the negative real axis. In the second case $s_{1,2}^{\infty}=1 \pm \sqrt{3} 1$ The two first order patterns make angles of $\pm 120^{\circ}$ with the positive real axis. In the third case $s_{1,2}^{\infty}=-\sqrt{3} \pm i$. Again there are two first order patterns, this time making angles of $\pm 30^{\circ}$, and the system is unstable for high gain.
-33-
Table 1
Matrices Used in Example 2

Case $\quad \mathrm{N}^{\infty}{ }^{\infty}\left(\mathrm{N}^{\infty}\right)^{-1}$
$F$
$1\left[\begin{array}{ll}1 & 0 \\ 0 & 2\end{array}\right]$
$\left[\begin{array}{cccc}0 & 0 & 1 & 0 \\ 16 & 4 & 0 & 2\end{array}\right]$
$2\left[\begin{array}{cc}0 & 1 \\ -4 & 2\end{array}\right]$
$\left[\begin{array}{cccc}8 & 2 & 0 & 1 \\ 16 & 4 & -4 & 2\end{array}\right]$
$3\left[\begin{array}{cc}0 & 1 \\ -4 & -2 \sqrt{3}\end{array}\right] \quad\left[\begin{array}{cccc}8 & 2 & 0 & 1 \\ -27.713 & -6.928 & -4 & -3.464\end{array}\right]$


## IX. Concluaions

For a restricted class of linear syatems we have shown how to analyze the aymptotic aiganstructure as control gains get very large. More importantly, we have shown how to decompose MIMO systems into SISO system that have the same asymptotically infinite eigenstructures, and we have shown how to use state feedback to achieve a desired asymptotic eigenstructure.

Acknowledgement:
We would like to thank Alan J. Laub, Bernard Levy, and George Verghese for their help in the development of this paper.

## RESERGUCES

1. Hazvey, C.A., and G. Stein, "Quadratic Weights for Asymptotic Regulator Properties," Isse trans. Auto. Control, Vol. 23, No. 3, June 1978.

2 Stein, G., "Generalized Quadratic Weights for Aaymptotic Regulator Properties," IEse Trans. Auto. Control, Vol. 24, No. 4, Aug. 1979.
3. Moore, B.C., "On the Flexibility Offered by state Feedback in Multivariable Systams Beyond Closed Loop Eigenvalue Assignment," IERE Trans. Auto. Control, Vol. 21, No. 5, Oct. 1976.
4. Klein, G., and B.C. Moore, "Eigenvalue-Generalized Eigenvector Assignment with State Feedback," IsEE Trans. Auto. Control, Vol. 22, No. 1, Feb. 1977.
5. Kouvaritakas, B., and U. Shaked, "Asymptotic Behavior of Root-Loci of Linear Multivariable Systems," Int. J. Control, Vol. 23, No. 3, 1976.
6. MacFariane, A.G.J., and I. Postlethwaite, "The Generalized Nyquist Stability Criterion and Multivariable Root Loci," Int. J. Control, Vol. 25, No. 1, 1977.
7. Laub, A.J., and B.C. Moore, "Calculation of Transmission zeros Using QZ Techniques," Automatica, Vol. 14, 1978.
8. MacFarlane, A.G.J., and N. Karcanias,"Poles and Zeros of Linear Multivariable systems: A Survey of the Algebraic, Geometric, and Complex Variable Theory, " Int. J. Control, Vol. 24, No. 1, 1976.
9. Verghese, G., "Infinite Frequency Behavior in Geneneralized Dynamical Systems," Ph.D. Thesis, Stanford University, Dec. 1978.
10. Garbow, B.S., et al., Matrix Eigensyatem Routines - EISPACK Guide Extension, Lecture Notes in Computer Science, vol. 51, SpringerVerlag, Berlin, 1977.
11. Wonham, W.M., Linear Multivariable Control: A Geometric Approach, Springer-Verlag, Berlín, 1974.


[^0]:    The research was conducted at the MIT Laboratory for Information and Decision Systems, with support provided by the NASA Langley Research Center under grant NASA/NAG1-2.

