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# A Rapid Perturbation Procedure for Determining Nonlinear Flow Solutions: Application to Transonic Turbomachinery Flows

Stephen S. Stahara, James P. Elliott,  
and John R. Spreiter

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A Rapid Perturbation Procedure  
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Solutions: Application to  
Transonic Turbomachinery Flows

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## TABLE OF CONTENTS

<u>Section</u>	Page
SUMMARY	1
1. INTRODUCTION	3
2. ANALYSIS	5
2.1 Perturbation Concept and Methods	5
2.2 Previous Applications	7
2.3 Coordinate Straining	8
2.4 Theoretical Formulation for Single Parameter Perturbations	9
2.5 Current Applications: Surface Pressures	12
3. RESULTS	15
3.1 Perturbation Results for Supercritical Single-Shock Flows and Subcritical Flows	15
3.1.1 Supercritical applications	15
3.1.2 Subcritical applications	18
3.2 Comparison of Continuous and Piecewise-Continuous Straining Function Perturbation Results	19
3.3 Perturbation Applications to Complex Supercritical Flows	21
3.3.1 Multi-shock supercritical flows	21
3.3.2 Supercritical compressor cascade flows	22
4. CONCLUSIONS AND RECOMMENDATIONS	23
APPENDIXES	
A - USER'S MANUAL FOR COMPUTER PROGRAM PERTURB	25
A.1 INTRODUCTION	25
A.2 PROGRAM DESCRIPTION	26
A.3 PROGRAM FLOW CHART	29
A.4 DICTIONARY OF INPUT VARIABLES	30
A.5 PREPARATION OF INPUT DATA	33
A.5.1 Description of Input	33
A.5.2 Format of Input Data	34

A.6	DESCRIPTION OF OUTPUT	37
A.7	ERROR MESSAGES	37
A.8	SAMPLE CASE	39
	FIGURES A.1 TO A.3	41
B	- LISTING OF COMPUTER PROGRAM PERTURB	64
C	- LIST OF SYMBOLS	73
	REFERENCES	75
	FIGURES	76

A RAPID PERTURBATION PROCEDURE FOR DETERMINING  
NONLINEAR FLOW SOLUTIONS: APPLICATION TO  
TRANSONIC TURBOMACHINERY FLOWS

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SUMMARY

An investigation was conducted to develop perturbation procedures and associated computational codes for determining nonlinear flow solutions, with the objective of establishing a method for minimizing computational requirements associated with parametric studies of transonic flows in turbomachines. The theoretical analysis involved the development of a rapid method for calculating first-order changes in nonlinear flow solutions due to variations of an arbitrary geometrical or flow parameter.

The procedure developed and evaluated, referred to as the direct correction method, was found to be capable of determining highly accurate approximations to families of strongly nonlinear solutions which are either continuous or discontinuous, and which represent variations in some arbitrary parameter. The method consists of defining a unit perturbation by employing two nonlinear solutions which differ from one another by a nominal change in some geometric or flow parameter, and then using that unit perturbation to predict a family of related nonlinear solutions over a range of parameter variation. Coordinate straining is used in determining the unit perturbation to account for the movement of discontinuities and maxima of high-gradient regions due to the perturbation. While simultaneous multiple-parameter perturbations can be treated by the method, the theoretical development and results presented in this initial study are for the single-parameter perturbation problem.

Although the procedure is generally applicable, the results reported here have been directed toward nonlinear aerodynamic applications. Attention is focused in particular on transonic

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flows which are strongly supercritical and exhibit large surface shock movement over the parametric range studied; and on subsonic flows which display large pressure variations in the stagnation and peak suction pressure regions. Flows past both isolated airfoils and compressor cascades involving a wide variety of flow and geometry parameter changes are reported. Comparisons with the corresponding 'exact' nonlinear solutions indicate a remarkable accuracy and range of validity of such a procedure. Computational time of the method, beyond the determination of the base solutions, is trivial.

## 1. INTRODUCTION

Given the remarkable growth in capability of advanced computational methods for the determination of a spectrum of nonlinear phenomena in such diverse disciplines as fluid dynamics, structures, and nuclear physics to name just a few - a capability which has already made many difficult calculations routine and which is certain to improve in the future - it is apparent that a need exists for complementary methods capable of alleviating, at least in part, the usage limitations imposed on these methods by their run times. The need becomes particularly compelling when large numbers of related cases are required as in parametric or design studies. Techniques such as direct acceleration procedures provide an important means of reducing computer time by improving computational efficiency of the solution algorithm, but these and similar methods, which enhance the solution algorithm itself, represent only a partial answer. What is most desirable is a means to minimize the actual number of separate calculations required in a particular application by extending, over some parametric range, the usefulness of each individual solution determined by these computationally expensive procedures.

Consequently, the basic motivation underlying this study is to extend the usefulness of such numerical solutions computed for specific turbomachinery configurations and flow conditions with a view toward reducing the computational requirements now necessary. The nature of the present investigation is both exploratory and developmental in the sense that aspects of the procedure such as validity, range of application, and economy will be investigated, and a computational code embodying all the results of the study will be developed.

Two fundamental methods for accomplishing such a perturbation procedure are available: a classical approach involving posing and solving linear perturbation equations; and a direct correction method employing two or more nonlinear base solutions. In this report, both of these methods are discussed; and an evaluation of the latter method, based on a large number of different applications, is made.

A crucial aspect of such perturbation methods is their ability to accurately treat regions where either discontinuities or high gradients exist. For the results presented here coordinate straining is introduced as a means of accounting properly for the displacement of discontinuities due to an arbitrary change in some solution parameter. This is shown to result in highly accurate perturbation predictions in the vicinity of the discontinuity. That idea has also been extended to improve predictions in the vicinity of other high-gradient regions.



Although the procedures developed are generally applicable, the specific results reported here are for aerodynamic applications. Since one of the primary objectives of this study was to provide a definitive proof-of-concept of such a perturbation method, a large variety of perturbation results based on transonic small-disturbance and full potential solutions were studied and are presented for nonlinear subsonic and transonic flows past both isolated airfoils and compressor cascades. In order to enable a critical evaluation of the range of validity and accuracy of the straining procedure, emphasis was placed on transonic flows which are strongly supercritical and exhibit large surface shock movement over the parametric range studied; and on subsonic flows which display large pressure variations in the stagnation and peak suction pressure regions.

## 2. ANALYSIS

### 2.1 Perturbation Concept and Methods

The basic hypothesis underlying the present procedure is that a range of solutions in the vicinity of a previously determined or base solution can be calculated to first-order accuracy in the incremental change of the varied parameter by determining a linearized unit perturbation solution  $Q_p$  defined according to the relation

$$\underbrace{Q}_{\substack{\text{Approximate solution for} \\ \text{conditions differing from} \\ \text{those of the base solution} \\ \text{by an amount characterized} \\ \text{by } \epsilon}} = \underbrace{Q_0}_{\substack{\text{Base solution} \\ \text{for some flow} \\ \text{quantity } Q}} + \underbrace{\epsilon \cdot \{Q_p\}}_{\substack{\text{Linearized perturbation} \\ \text{solution for a unit} \\ \text{change of } \epsilon}} \quad (1)$$

The effectiveness of such a method, of course, depends upon the ability of the relationship defined by equation (1) to remain accurate over a range  $\epsilon$  of practical significance, and the fact that the unit perturbation  $Q_p$  need be determined only once. The significance of the unit perturbation  $Q_p$  is obvious. It represents the local rate of change of the base<sup>p</sup> flow solution  $Q_0$  with respect to the particular quantity, say  $q$ , perturbed; that is  $Q_p = (\partial Q / \partial p)_0$ .

Two generic methods exist for determining  $Q_p$ , each differing in philosophy and having its own particular strengths and weaknesses. We refer to these methods simply as the linear perturbation equation method and the direct correction method.

The linear perturbation equation method represents the classical approach for performing a perturbation analysis and proceeds by establishing and solving a linear differential equation for the perturbation. Although in the present application, we confine our interest solely to the first-order term, the complete procedure represents a rational approximation scheme capable of continuation to any order. The method proceeds by expanding the dependent variables in an ascending power series in the incremental change  $\epsilon$  of the varied parameter, inserting that representation into the full governing equations and then assembling the result into a corresponding series of linear equations in ascending orders in  $\epsilon$ . Higher-order solutions

in general depend on both base flow plus lower-order solutions. Determination of the appropriate boundary conditions is done in a similar fashion.

The power of the linear perturbation equation method is that it requires the calculation of only one nonlinear base solution. With that information, any number of individual perturbations can then be calculated, subject to the particular governing linear partial differential equations and boundary conditions which apply. The disadvantages are that each perturbation problem must be posed individually, including differential equations and boundary conditions. Furthermore, it may be necessary to simplify the governing equations and boundary conditions to a point where they can be solved rapidly relative to rerunning the base flow procedure. Moreover, the perturbation solutions themselves may be quite sensitive to the base flow solutions which usually enter into the perturbation problem through the differential equation and sometimes through the boundary conditions as well.

The fundamental limitation of the method is the restriction of the range over which the perturbation procedure remains valid to a linear one. Since this characteristic depends upon the local behavior of the base flow with respect to the varied parameter, no general statement regarding range of validity is possible. Typical behavior for a given class of flows must be ascertained by checks with the base flow procedure. Initially unknown at the outset of an application with this technique, then, are the accuracy requirements imposed on the base solution by the perturbation procedure and the range of parameter variation over which the linear assumption is valid.

For the alternative method, the perturbation solution per unit change of the varied parameter,  $Q_p$ , is determined simply by differencing two nonlinear base flow solutions removed from one another by some nominal change of a particular flow or geometrical quantity. A unit perturbation solution is then obtained by dividing that result by the change in the perturbed quantity. Related solutions are determined by multiplying the unit perturbation by the desired parameter change and adding that result to the base flow solution. This simple procedure, however, only works directly for continuous flows for which the perturbation change does not alter the solution domain. For those perturbations which change the flow domain, coordinate stretching (usually obvious) is necessary to insure proper definition of the unit perturbation solution. Similarly, for discontinuous flows, coordinate straining is necessary to account for movement of discontinuities due to the perturbation solution.

The attractiveness of the correction method is that it is not restricted to a linear variation range but rather replaces the nonlinear variation between two base solutions with a linear

fit. This de-emphasizes the dependence and sensitivity inherent in the linear perturbation equation method on the local rate of change of the base flow solution with respect to the varied quantity. For many applications, particularly at transonic speeds, the flow is highly sensitive, and the linear range of parameter variation can be sufficiently small to be of no practical use. Furthermore, other than the approximation of a linear fit between two nonlinear base solutions, the direct correction method is not restricted by further approximations with respect to the governing differential equations and boundary conditions. Rather, it retains the full character of the original methods used to calculate the base flow solutions. Most importantly, no perturbation differential equations have to be posed and solved, only algebraic ones. In fact, it isn't even necessary to know the exact form of the perturbation equation, only that it can be obtained by some systematic procedure and that the perturbations thus defined will behave in some 'generally appropriate' fashion so as to permit a logical perturbation analysis. For situations involving perturbations of physical parameters, such as reported here, the governing perturbation equations are usually transparent, or at least readily derivable. Finally, in applying this method it isn't necessary to work with primitive variables; rather the procedure can be applied directly to the final quantity desired.

The primary disadvantage of this method is that two base solutions are required for each parameter perturbation considered. Furthermore, both flows must be topologically similar, i.e., discontinuities or other characteristic features must be present in both base solutions used to establish the unit perturbation.

## 2.2 Previous Applications

Detailed studies of the linear perturbation equation method to sensitive transonic flows, with a view toward testing the method as an effective tool for reducing computational requirements, have not been done. The primary reason is that such studies quickly become overwhelming. Each perturbation problem must be posed individually, subject to its own particular governing equations and boundary conditions; and then a separate computational code for the perturbation established. Generally, the governing equations and boundary conditions of the perturbation, even though they are linear, are more involved than those for the base solution. Additionally, the computational and convergence characteristics can pose similar or additional problems from those of the base flow procedure.

In an attempt to examine some of these problems for transonic applications in at least a preliminary fashion, an application of the linear perturbation equation method to

transonic turbomachinery flows was made in reference 1. The conclusions obtained from that study were that reasonable results could be anticipated from the method for blade geometry changes, such as blade thickness and angle of attack. Less satisfactory results were obtained for perturbation changes in overall quantities, such as blade spacing and free-stream Mach number, a result that could be anticipated a priori since such perturbations alter the basic character of the flow more rapidly. The most significant conclusion of that study was the demonstration of the primary limitation of the linear perturbation equation method. That is, for sensitive flows such as occur in transonic situations, the basic linear variation assumption fundamental to the technique is sufficiently restrictive that the permissible range of parameter variation becomes so small as to be of limited practical use. Some preliminary applications of the direct correction method, however, displayed a significantly wider range of perturbation solution validity, in particular for strongly supercritical flows when coordinate straining was employed to account for shock movement.

### 2.3 Coordinate Straining

The concept of employing coordinate straining to remove nonuniformities from perturbation solutions of nonlinear problems is well established and originally suggested by Lighthill (ref. 2) three decades ago. The basic idea of the technique is that a straightforward perturbation solution may possess the appropriate form, but not quite at the appropriate location. The procedure is to strain slightly the coordinates by expanding them as well as the dependent variables in an asymptotic series. It is often unnecessary to actually solve for the straining. It can generally be established by inspection. The final uniformly valid solution is then found in implicit form, with the strained coordinate appearing as a parameter.

In the original applications of the method (ref. 3), it was applied in the 'classical' sense; that is, series expansions of the dependent and independent variables in ascending powers in some small parameter were inserted into the full governing equation and boundary conditions, and the individual terms of the series determined. An ingenious variation in the application of the method was made by Pritulo (ref. 4) who demonstrated that if a perturbation solution in unstrained coordinates has been determined and found to be nonuniform, the coordinate straining required to render that solution uniformly valid can be found by employing straining directly in the known non-uniform solution, and then solving algebraic rather than differential equations. The idea of introducing strained coordinates a posteriori has since been applied to a variety of different problems (see ref. 3), and forms the basis of the current applications.

The fundamental idea underlying coordinate straining as it relates to the application of perturbation methods to supercritical transonic flows is illustrated geometrically in figure 1. In the upper plot on the left, two typical transonic pressure distributions are shown for a highly supercritical flow about a nonlifting symmetric profile. The distributions can be regarded as related nonlinear flow solutions separated by a nominal change in some geometric or flow parameter. The shaded area between the solutions represents the perturbation result that would be obtained by directly differencing the two solutions. We observe that the perturbation so obtained is small everywhere except in the region between the two shock waves, where it is fully as large as the base solutions themselves. This clearly invalidates the perturbation technique in that region and most probably somewhat ahead and behind it as well. The key idea of a procedure for correcting this, pointed out by Nixon (refs. 5,6), is first to strain the coordinates of one of the two solutions in such a fashion that the shock waves align, as shown in the upper plot on the right of figure 1, and then determine the unit perturbation. Equivalently, this can be considered as maintaining the shock wave location invariant during the perturbation process, and assures that the unit perturbation remains small both at and in the vicinity of the shock wave. Obviously, shock points are only one of a number of characteristic high-gradient locations such as stagnation points, maximum suction pressure points, etc., in which the accuracy of the perturbation solution can degrade rapidly. The plots in the lower left part of the figure 1 indicate such a situation and display typical transonic pressure distributions which contain multiple shocks and high-gradient regions. Simultaneously straining at all these locations, as indicated in the lower right plot, serves to minimize the unit perturbation over the entire domain considered, and provides the key to maximizing the range of validity of the perturbation method.

#### 2.4 Theoretical Formulation for Single-Parameter Perturbations

In order to provide the theoretical essentials of the correction method, consider the formulation of the procedure at the level of the full potential equation, as most of the results presented here are based on that level. We denote the operator  $L$  acting on the velocity potential  $\phi$  as that which results in the two-dimensional full potential equation for  $\phi$ , i.e.

$$L[\phi] = 0 \quad (2)$$

If we now expand the potential in terms of zero- and higher-order components in order to account for the variation of some arbitrary geometrical or flow parameter  $q$

$$\begin{aligned}\phi &= \phi_0 + \epsilon \phi_1 + \dots \\ q &= q_0 + \Delta q\end{aligned}\tag{3}$$

and then insert this into the governing equation (2), expand the result, order the equations into zero- and first-order components, and make the obvious choice of expansion parameter  $\epsilon = \Delta q$ , we obtain the following governing equations for the zero- and first-order components

$$\begin{aligned}L[\phi_0] &= 0 \\ L_1[\phi_1] + \frac{\partial}{\partial q} L[\phi_0] &= 0\end{aligned}\tag{4}$$

Here  $L_1$  is a linear operator whose coefficients depend on zero-order quantities and  $\partial L[\phi_0]/\partial q$  represents a 'forcing' term due to the perturbation. Actual forms of  $L_1$  and the 'forcing' term are provided in reference 1 for a variety of flow and geometry parameter perturbations of a two-dimensional turbomachine, and in reference 7 for profile shape perturbations of an isolated airfoil. An important point regarding equation (4) for the first-order perturbation  $\phi_1$  is that the equation represents a unit perturbation independent of the actual value of the perturbation quantity  $\epsilon$ .

Appropriate account of the movement of discontinuities and maxima of high-gradient regions due to the perturbation is now accomplished by the introduction of strained coordinates  $(s,t)$  in the form

$$\begin{aligned}x &= s + \epsilon x_1(s,t) \\ y &= t + \epsilon y_1(s,t)\end{aligned}\tag{5}$$

where

$$\begin{aligned}x_1(s,t) &= \sum_{i=1}^N \delta x_i x_{1i}(s,t) \\ y_1(s,t) &= \sum_{i=1}^N \delta y_i y_{1i}(s,t)\end{aligned}\tag{6}$$

and  $\epsilon \delta x_i$ ,  $\epsilon \delta y_i$  represents individual displacements of the  $N$  strained points, and  $x_{1i}(s,t)$ ,  $y_{1i}(s,t)$  are straining functions associated with each of the  $N$  strained points. Introducing the strained coordinate equations (5) and (6) into the expansion formulation leaves the zero-order result in equation (4) unchanged, but results in a change of the following form for the perturbation

$$L_1[\phi_1] + L_2[\phi_0] + \frac{\partial}{\partial q} L[\phi_0] = 0 \quad (7)$$

Here the operators are understood to be expressed in terms of the strained (s,t) coordinates, and the additional operator  $L_2$  arises specifically from displacement of the strained points. In references 6 and 7, specific expressions for  $L_2$  are provided for selected perturbations involving transonic small-disturbance and full potential equation formulations. The primary point, however, with regard to perturbation equation (7) expressed in strained coordinates is that it remains valid as before for a unit perturbation and independent of  $\epsilon$ .

In employing the correction method, equation (7) for the unit perturbation is solved by taking the difference between two solutions obtained by the full nonlinear procedure after appropriately straining the coordinates. If we designate the two solutions for some arbitrary flow quantity  $Q$  as base  $Q_0$  and calibration  $Q_c$ , respectively, of the varied parameter, we have for the predicted flow at some new parameter value  $q$  (ref. 8)

$$Q(x,y) = Q_0(s,t) + \frac{\epsilon}{\epsilon_0} [Q_c(\bar{x},\bar{y}) - Q_0(s,t)] \quad (8)$$

where

$$\begin{aligned} \bar{x} &= s + \epsilon_0 x_1(s,t) \\ \bar{y} &= t + \epsilon_0 y_1(s,t) \\ x &= s + \frac{\epsilon}{\epsilon_0} [\bar{x} - s] \\ y &= t + \frac{\epsilon}{\epsilon_0} [\bar{y} - t] \\ \epsilon_0 &= q_c - q_0 \\ \epsilon &= q - q_0 \end{aligned} \quad (9)$$

In the following section, applications of the correction procedure are made to predict surface properties. Also provided are the particular forms of the straining functions equation (6) for those applications.



## 2.5 Current Applications: Surface Pressures

For the current applications, we have employed coordinate straining with the correction method to predict surface pressure distributions for a wide variety of single-parameter geometrical flow perturbations of isolated airfoils and cascades. In that instance where flow properties are required along some contour, the solutions can be represented by

$$\begin{aligned} Q(x;\epsilon) &\sim Q_0(s) + \epsilon Q_1(s) + \dots \\ x &\sim s + \epsilon x_1(s) + \dots \end{aligned} \tag{10}$$

where  $x$  is the independent variable measuring distance along the contour or a convenient projection of that distance,  $s$  is the strained coordinate, and  $\epsilon$  a small parameter representing the change in some flow or geometrical variable which we wish to vary.

In order to determine the first-order corrections  $Q_1(s)$ , we require a base and calibration solution in which the calibration solution is determined by varying an arbitrary parameter  $q$  by some nominal amount from the base flow value.

In this way, the first-order correction  $Q_1(s)$  can be determined as

$$Q_1(s) = \frac{Q_c(\bar{x}) - Q_0(s)}{q_c - q_0} \tag{11}$$

where  $Q_c$  is the calibration solution corresponding to changing the parameter  $q$  to a new value  $q_c$ ,  $\bar{x}$  is the strained coordinate pertaining to the  $Q_c$  calibration solution, and  $q_c - q_0$  represents the change in the  $q$  parameter from its base flow value. If we now desire to keep invariant during the perturbation process a total of  $N$  points corresponding to discontinuities or high-gradient maxima, we can represent the solution by:

$$Q(x;\epsilon) = Q_0(s) + \epsilon Q_1(s) \tag{12}$$

where

$$Q_1(s) = \frac{Q_c(\bar{x}) - Q_0(s)}{\epsilon_c} \tag{13}$$

$$\bar{x} = s + \sum_{i=1}^N \epsilon_c (\delta x_i^C) \cdot x_{1_i}(s) \quad (14)$$

$$x = s + \sum_{i=1}^N \epsilon (\delta x_i^C) \cdot x_{1_i}(s) \quad (15)$$

$$\epsilon_c = q_c - q_o \quad (16)$$

$$\epsilon = q - q_o \quad (17)$$

$$\epsilon_c (\delta x_i^C) = (x_i^C - x_i^O) \quad (18)$$

$$\epsilon (\delta x_i^O) = \frac{\epsilon}{\epsilon_c} (x_{1_i} - x_i^O) \quad (19)$$

Here  $\epsilon_c (\delta x_i^C)$  given in equation (18) represents the displacement of the  $i$ th invariant point in the calibration solution from its base flow location due to the selected change  $\epsilon_c$  in the  $q$  parameter given by equation (16),  $\epsilon (\delta x_i^O)$  given in equation (19) represents the predicted displacement of the  $i$ th invariant point from its base flow location due to the desired change  $\epsilon$  in the  $q$  parameter given by equation (17), and  $x_{1_i}(s)$  is a unit-order straining function having the property that

$$x_{1_i}(x_k^O) = \begin{cases} 1 & k = i \\ 0 & k \neq i \end{cases} \quad (20)$$

which assures alignment of the  $i$ th invariant point between the base and calibration solutions.

In addition to the single condition equation (20) on the straining function, it may be convenient or necessary to impose additional conditions at other locations along the contour. For example, it is usually necessary to hold invariant the end points along the contour, as well as to require that the straining vanish in a particular fashion in those locations. All of these conditions, however, do not serve to determine the straining uniquely. The nonuniqueness of the straining, nevertheless, can often be turned to advantage, either by selecting particularly simple classes of straining functions or by requiring the straining to satisfy further constraints convenient for a particular application. An example of the effect of employing two different straining functions for a strongly-supercritical flow was

provided in reference 6. Here we provide additional results demonstrating some of the limitations of various polynomial straining functions and provide some comparisons with piecewise-continuous functions. The particular classes of straining functions employed were continuous polynomial and linear piecewise-continuous. For these two classes, the functional forms of the straining can be compactly written. For example, equation (14) becomes, for continuous polynomial straining

$$\bar{x} = s + \sum_{i=2}^{N-1} L_i(s) \cdot (x_i^C - x_i^O) \quad (21)$$

where  $L_i$  are Lagrangian coefficients given by

$$L_i(s) = \prod_{\substack{k=1 \\ k \neq i}}^N \frac{(s - x_k^O)}{(x_i^O - x_k^O)} \quad (22)$$

whereas for linear piecewise-continuous straining,  $\bar{x}$  is given by

$$\bar{x} = s + \sum_{i=2}^{N-1} \left\{ \frac{x_{i+1}^O - s}{x_{i+1}^O - x_i^O} \cdot (x_i^C - x_i^O) + \frac{s - x_i^O}{x_{i+1}^O - x_i^O} \cdot (x_{i+1}^C - x_{i+1}^O) \right\} H(x_{i+1}^O - s) \cdot H(s - x_i^O) \quad (23)$$

where  $H$  denotes the Heaviside step function. As discussed above, it is usually necessary to hold invariant both of the end points along the contour in addition to the points corresponding to discontinuities or high-gradient maxima. Consequently, for the results reported here, the array of invariant points in the base and calibration solutions have been taken as

$$\begin{aligned} x_i^O &= \{0, x_1^O, x_2^O, \dots, x_n^O, 1\} \\ x_i^C &= \{0, x_1^C, x_2^C, \dots, x_n^C, 1\} \end{aligned} \quad (24)$$

where the contour length has been normalized to unity. Figure 2 provides a summary of the various combinations of flows and straining functions employed.

### 3. RESULTS

One of the primary objectives of the present investigation is to explore the accuracy and range of validity of such perturbation procedures to determine to what extent they are capable of providing results useful in an engineering analysis. To this end, we have tested the correction method with coordinate straining over a wide variety of different geometrical and flow condition perturbations, including applications to both isolated airfoils and compressor cascades. In particular, since the ability of the method to account accurately for the movement of discontinuities and maxima of high-gradient but continuous regions is essential if such procedures are to be of general use, emphasis was placed on transonic flows which are strongly supercritical and exhibit large surface shock movement over the parametric range studied. Base flow theoretical solutions were determined from small-disturbance transonic potential (ref. 9) and full potential solutions (refs. 10, 11, 12). In the results to follow, which were selected as typical from systematic calculations of a much larger number of cases, the choice of base and calibration solutions was often made at the limits of validity of the procedure to observe how well the method works under such conditions.

#### 3.1 Perturbation Results for Supercritical Single-Shock Flows and Subcritical Flows

3.1.1 Supercritical applications. - In figure 3, we present results for a thickness-ratio perturbation of strongly supercritical flows past a nonlifting cascade of biconvex profiles at  $M_\infty = 0.80$  having a spacing-to-chord ratio of  $H/C = 1.0$ . The dotted and dashed results on the figure represent the base and calibration surface pressure distribution for  $\tau = (0.075, 0.065)$ , respectively, and were obtained by solving the transonic small-disturbance potential equation using the code TSFOIL (ref. 9). An x-grid having 48 points on the blade profile was used. These solutions were then used to determine the unit perturbation. The open circles represent the perturbation solution for  $\tau = 0.073$  in the plot on the left and for  $\tau = 0.070$  in the plot on the right. Those perturbation results are meant to be compared with the solid lines in the plots which are the corresponding nonlinear solutions obtained by rerunning TSFOIL at the new thickness ratios. Quadratic straining was used with shock point and leading and trailing edges held invariant. The base and calibration flow shock-point locations for this example, as well as for all of the supercritical cases presented here, were determined as the point where the pressure coefficient passed through critical with compressive gradient.

With regard to the results, several points are noteworthy. Selection of a cascade rather than an isolated airfoil provides a more sensitive transonic flow situation. Additionally, the choice of a highly supercritical base and almost subcritical calibration solution provides both an example of extreme separation between the two nonlinear solutions used to define the unit perturbation, as well as a situation where one solution is near the limits of validity of the perturbation analysis. Recall that both solutions must be topographically similar, i.e., must contain the same number of discontinuities (shocks) and other characteristic features.

We note that comparisons of the perturbation results with the nonlinear calculations are very satisfactory for both thickness ratios, with the only discrepancy being a slight disagreement at the lower thickness ratio ( $\tau = 0.070$ ) at several points in the post-shock region. Additional calculations not presented here in which a more reasonable choice of calibration solution is made, say at  $\tau = 0.070$ , removes that discrepancy as well. The main point provided by the results of figure 3 is that for certain classes of supercritical flows even widely separated base solutions can be used to provide reasonable perturbation predictions.

In figure 4, we provide similar strongly supercritical results again for interpolation-only perturbation solutions, but in this instance on a somewhat finer grid. These results employed full potential base solutions (ref. 10), and represent thickness ratio perturbations of nonlifting symmetric free-air flows past NACA four-digit thickness-only airfoils at  $M_\infty = 0.820$ . The body-fitted mesh employed had 75 points on both upper and lower surfaces, which is half again as many as in the preceding example. For the base and calibration flows, the thickness ratios were  $\tau = 0.120$  and  $0.080$ , respectively. Comparisons between the perturbation predictions and the full nonlinear calculation are exhibited in figure 4 for  $\tau = 0.110$ ,  $0.105$ ,  $0.100$ , and  $0.095$ . We note that the comparisons are remarkably good, in particular, in the region of the shock. The first-order perturbation accurately predicts both shock location and the post-shock expansion behavior. Reference to the coarser grid results given in figure 3 indicates that the finer grid resolution clearly enhances the perturbation result, indicating that better accuracy and a larger range of validity of the perturbation solutions can be anticipated when fine-grid base solutions are used to define the unit perturbation.

In the two preceding examples, perturbation results were provided for interpolation-only between widely spaced base and calibration solutions. In figure 5, we provide similar strongly supercritical thickness-ratio perturbation results for extreme solution extrapolation using very closely spaced base and calibration solutions (ref. 10). The upper plots display results

for extrapolation downward from base and calibration flows past nonlifting NACA 00XX profiles with  $\tau = 0.115$  and  $0.120$  at  $M_\infty = 0.820$ . Perturbation predictions are shown for  $\tau = 0.105$  and  $0.100$ , which represent  $\Delta\tau$  excursions from the base flow ( $\tau = 0.115$ ) that are two and three times the parameter change between the base and calibration solutions ( $\Delta\tau = 0.005$ ) used to define the unit perturbation. For these results, comparisons with the full nonlinear calculations are very good. The lower plots display similar results for extreme extrapolation upward from base and calibration solutions have  $\tau = 0.095$  and  $0.090$ . Perturbation predictions are shown for  $\tau = 0.105$  and  $0.110$ , which again represent excursions from the base flow that are two and three times the parameter change between the base and calibration solutions. In this instance, while comparisons of the perturbation results and the full nonlinear solutions for both cases are good, the results at  $\tau = 0.110$  are beginning to display some not surprising discrepancies near the shock wave, indicating that the perturbation result is nearing the limit of its range of validity for this particular choice of base and calibration flows.

The results indicated in figure 5, however, clearly demonstrate that not only is accurate solution extrapolation possible, but that for some situations even closely spaced nonlinear solutions can be used to cover a wide range of related solutions. Additionally, the range of parameter variation in this example over which the perturbation results remain accurate - i.e., parameter changes three times the difference between the two nonlinear solutions used to define the unit perturbation - is remarkable, and far beyond what one would anticipate for a first-order correction.

Perturbation results using a more reasonable choice of base and calibration solutions are provided in figure 6. Those results involve Mach number perturbations of highly supercritical full potential (ref. 10) flows past a NACA 0012 airfoil at  $\alpha = 0^\circ$ . The base and calibration results are for  $M_\infty = 0.800$  and  $0.820$ , and the comparisons indicated are for perturbation results interpolated to  $M_\infty = 0.810$  and extrapolated downward to  $M_\infty = 0.790$ . As in the case of the geometric perturbations given in figures 4 and 5, these perturbation results are also in very good agreement with the nonlinear calculations at the new Mach numbers. For this perturbation, as well as for a number of other Mach number perturbations, we have separately determined the perturbation result in two different ways. First, we have taken cognizance of the fact that a Mach number perturbation alters the governing differential equation for the first-order perturbation from that of other geometric or flow parameter changes; and have used the suggestion of reference 6 to consider such perturbations via a transonic small-disturbance approximation, whereby the same perturbation equation can be preserved by employing a modified expansion parameter  $\epsilon$ . An alternative procedure is to treat a

Mach perturbation directly and interpret  $\epsilon$  as the difference in Mach number. We have done these calculations and compared the perturbation results for a number of cases using both full potential solutions, as for the results shown in figure 6, and transonic small-disturbance solutions, and have observed no essential difference between the two sets of results. The perturbation results presented in figure 6 correspond to those for  $\epsilon$  equal to the difference in Mach number.

All of the supercritical perturbation results presented in figures 3 to 6 have been for symmetric flows and have employed a quadratic straining function. In figure 7, we present results for an angle of attack perturbation of lifting flows past a NACA 0012 profile at  $M_\infty = 0.70$ . The full potential (ref. 10) base and calibration solutions are at  $\alpha = 3.0^\circ$  and  $4.0^\circ$ , with comparisons of the perturbation and full nonlinear results shown for  $\alpha = 3.5^\circ$  and  $2.5^\circ$ . Cubic straining has been used with the invariant points corresponding to the lower trailing edge, stagnation point, shock point, and the upper trailing edge (see fig. 2). We note that  $\alpha = 3.5^\circ$ , the perturbation results are very good everywhere, in particular, in the vicinity of the shock and stagnation regions. At  $\alpha = 2.5^\circ$ , the perturbation results are still very good in the shock and stagnation regions and on most of the upper and lower surface, but near the trailing edge a discrepancy has occurred. The cause of this discrepancy lies solely with the cubic straining function used. It is due to the fact that although the straining vanishes identically at the trailing edge, for the particular choice of base and calibration solutions in this example, the straining in the near vicinity of the trailing edge becomes sufficiently large to introduce a misalignment in the unit perturbation in that high-gradient region. The correction to this is discussed in the section describing piecewise-continuous straining functions.

3.1.2 Subcritical applications.- Although supercritical flows are clearly of central concern in any transonic analysis for which the perturbation methods presented here would be used, applications to subcritical nonlinear flows are also of significance. To this end, we have applied these same techniques to a variety of subcritical flows to examine their accuracy and range of validity for such applications.

In figure 8, we present some summary results for four different subcritical perturbation applications to an isolated airfoil. All of these results are based on full potential solutions (ref. 10) with quadratic straining holding invariant the stagnation point and the trailing edge points. The plot on the upper left displays comparisons for a camber line perturbation of a lifting flow with  $M_\infty = 0.50$  and  $\alpha = 2^\circ$  past an airfoil having a NACA 0012 thickness distribution and a parabolic-arc camber line having the maximum camber located at midchord. Base

and calibration flows with camber ratio  $h/c = 0.02$  and  $0.01$  were used to extrapolate perturbation results to  $h/c = 0.05$ . Comparisons with the full result is essentially exact. The plot on the upper right provides similar results for a thickness-ratio perturbation of a lifting flow with  $M_\infty = 0.50$  and  $\alpha = 2.0^\circ$  past NACA 00XX thickness-only airfoils. Base and calibration flows with  $\tau = 0.12$  and  $0.04$  were used to provide interpolation results at  $\tau = 0.08$ . Again, the agreement is essentially exact even in the peak suction pressure region. The plot on the lower left provides angle-of-attack perturbation results for  $M_\infty = 0.50$  flow past a NACA 0012 airfoil, using base/calibration results for  $\alpha = 4.0^\circ$ ,  $2.0^\circ$  to predict results at  $\alpha = 3.0^\circ$ , with the agreement again being quite good. The final comparisons given in the plot on the lower left are for a Mach number perturbation of a lifting flow at  $\alpha = 2^\circ$  past an airfoil having a NACA 0012 thickness distribution and a parabolic-arc camber line with camber ratio  $h/c = 0.03$  at midchord. Base/calibration results at  $M_\infty = 0.40$ ,  $0.60$  were used to predict results at  $M_\infty = 0.55$ , with good agreement with the full nonlinear calculation.

In figure 9, we present similar summary results for subcritical perturbation applications to a compressor cascade having a 4% biconvex thickness distribution and a 1% parabolic-arc camber line blade, a pitch of  $t/c = 0.37$ , and oncoming Mach number  $M_\infty = 0.770$ . These results are based on the full potential solution procedure of reference 11 and have also used quadratic straining to hold the trailing edge points and stagnation point invariant. The plots in the upper part of the figure represent an inflow angle perturbation, with base/calibration inflow angles  $\beta_i = 47.8^\circ$ ,  $49.8^\circ$  used to predict extrapolation results in the plot on the left for  $\beta_i = 48.8^\circ$  and interpolation results in the plot on the right for  $\beta_i = 48.8^\circ$ . In the lower left plot, interpolation results are displayed for an outflow angle perturbation with base/calibration outflow angles  $\beta_o = 31.5^\circ$ ,  $39.5^\circ$  used to predict the flow at  $\beta_o = 35.5^\circ$ . The lower right plot provides interpolation results for a rotational speed perturbation with base/calibration rotational speeds  $\omega = 967,667$  rad/sec used to predict the flow at  $\omega = 827$  rad/sec. In all of these results, the perturbation results are good, including the regions near the leading and trailing edge where a peaky behavior due to local grid resolution is observed.

### 3.2 Comparison of Continuous and Piecewise-Continuous Straining Function Perturbation Results

The results presented in figures 10 to 13 illustrate the effect of using different straining functions to determine the perturbation results. Comparisons are provided for several strongly supercritical flows, demonstrating the differences in perturbation solutions between using quadratic and cubic straining



functions and corresponding piecewise-continuous straining functions.

Figure 10 displays a comparison for a symmetric supercritical thickness-ratio perturbation at  $\tau = 0.110$  for which results based on quadratic straining were given in figure 4. In that figure the open circles denote the previously obtained perturbation results using quadratic straining, while the asterisks denote the corresponding result when using linear piecewise-continuous straining. The points held invariant are the leading and trailing edges and the shock point. For this case there is virtually exact agreement everywhere between the two perturbation results as well as with the nonlinear result. An analogous comparison with a cubic straining result is provided in figure 11 where the invariant points are the lower trailing edge, stagnation point, shock point, and upper trailing edge. Displayed in that figure as open circles are the cubic-straining supercritical angle-of-attack perturbation results at  $\alpha = 2.5^\circ$  which were previously given in figure 7. Asterisks denote the corresponding linear piecewise-continuous straining perturbation result. We note that the discrepancy near the trailing edge caused by the cubic straining has been effectively removed in the piecewise-continuous result. Moreover, the good agreement with the full nonlinear result which the cubic result displayed near the shock and stagnation regions, as well as over the remainder of the airfoil surface, is also obtained with the piecewise-continuous result.

Finally, we have found that when employing quadratic, cubic, and higher-order polynomials as straining functions, for certain combinations of base flow shock location and shock movement between base and calibration solutions, particularly when large shock movements are involved, the polynomial straining functions will strain some points off the airfoil surface. This, of course, invalidates the determination of the unit perturbation, and requires that a different straining function be employed. Piecewise-continuous straining functions provide a simple means of avoiding such difficulties.

In figures 12 and 13, we have provided examples illustrating this effect for both quadratic and cubic straining functions. Figure 12 provides a comparison of perturbation results obtained using quadratic (open circles) and linear piecewise-continuous (asterisks) straining applied to a supercritical Mach number perturbation for symmetric nonlifting flow past a NACA 0012 airfoil. Widely separated base/calibration flows (ref. 10) at  $M_\infty = 0.820$  and  $0.750$  were used to predict the flow at  $M_\infty = 0.810$ . The spurious behavior near the leading edge displayed by the open circles is due to the quadratic function moving points in the strained calibration solution off the airfoil surface. The piecewise-continuous results indicated by the asterisks display a smooth variation in that region, and provide good agreement

everywhere with the full nonlinear result. Figure 13 provides a corresponding comparison for cubic straining. Angle-of-attack perturbation results at  $M_\infty = 0.70$  for flow past a NACA 0012 profile using base/calibration results (ref. 10) at  $\alpha = 2.25^\circ$  and  $4.00^\circ$  are used to predict the flow at  $\alpha = 3.25^\circ$ . The unusual results displayed by the open symbols near the trailing edge indicate that the cubic function has strained points off the airfoil surface in that region. However, the linear piecewise-continuous result corrects that problem and displays good agreement with the nonlinear calculation in that region as well as at the shock and stagnation point.

### 3.3 Perturbation Applications to Complex Supercritical Flows

In order to provide a severe test of the perturbation procedure, we have applied the method to a number of transonic flows that are characterized by surface pressure distributions having multiple shock and/or high-gradient locations, such as those typified schematically in the lower plots of figure 1. Demonstration of the ability of the perturbation method to predict accurately such classes of flows, which are typical of those encountered in certain transonic turbomachinery applications, is crucial to the present study. In order to accomplish such a demonstration, we have investigated two separate classes of sensitive supercritical transonic flows, i.e. those with multiple-shock waves, and those having a single shock together with multiple high-gradient regions. Examples of perturbation results for such flows are provided below.

3.3.1 Multi-Shock Supercritical Flows.- In figure 14, we present results for an angle-of-attack perturbation of supercritical lifting flows past a NACA 0012 profile at  $M_\infty = 0.80$ . These highly sensitive flows exhibit two shocks, one on each the upper and lower surface. The full potential (ref. 10) base and calibration flows employed are at  $\alpha = 0.50^\circ$  and  $0.20^\circ$ , with comparisons of the perturbation and full nonlinear results shown for  $\alpha = 0.0^\circ, 0.1^\circ, 0.4^\circ,$  and  $0.6^\circ$ . Piecewise-continuous linear straining has been used with the invariant points corresponding to the lower trailing edge, lower surface shock point, stagnation point, upper surface shock point, and upper trailing edge (see fig. 2). We note that the symmetrical extrapolation result at  $\alpha = 0.0^\circ$  is separately predicted from both the upper surface and lower surface pressure distributions, and, as can be seen, the results are quite good. The remaining results at  $\alpha = 0.1^\circ, 0.4^\circ,$  and  $0.6^\circ$ , which represent both extrapolation and interpolation from the base and calibration flows, are in excellent agreement with the full nonlinear result. As an indication of the sensitivity of these flows, we have found that the lower surface shock

disappears at an angle of attack of approximately  $0.8^\circ$ ; yet the lower surface pressure distribution is well predicted by the perturbation result over the parametric range studied.

3.3.2 Supercritical Compressor Cascade Flows.- As an example of the ability of the method to predict a complex supercritical flow, in figure 15 we provide results for oncoming Mach number perturbation of supercritical flows past a cascade composed of Jose Sanz (ref. 12) profiles. For these results, the oncoming and exit flow angles are  $30.81^\circ$  and  $0.09^\circ$ , respectively, the blade twist is  $9.33^\circ$ , while the gap to chord ratio is 1.028. The full potential (ref. 12) base and calibration flow oncoming Mach numbers are  $M_\infty = 0.77$  and  $0.81$ , with comparisons of perturbation and full nonlinear results shown at  $M_\infty = 0.75, 0.79, 0.89,$  and  $0.83$ . Piecewise-continuous linear straining was employed with invariant points at the lower trailing-edge, stagnation point, shock point and upper trailing edge. As with the multiple-shock example shown in figure 14, we note that the perturbation predictions are in excellent agreement with the nonlinear results. In particular, we note that the perturbation procedure captures the variation of the plateau-like pressure distribution on the upper surface near the leading edge, the location and strength of the shock, the post-shock expansion region, the rapid expansion near the trailing edge, and the expansion on the lower surface near the stagnation point, indicating a capability for treating very general flow situations.

#### 4. CONCLUSIONS AND RECOMMENDATIONS

An evaluation has been made of a perturbation procedure for determining highly accurate approximations to families of nonlinear solutions which are either continuous or discontinuous, and which represent variations in some arbitrary parameter. The procedure employs a unit perturbation, determined from two nonlinear solutions which differ from one another by a nominal change in some geometric or flow parameter, to predict a family or related nonlinear solutions. Coordinate straining is used in determining the unit perturbation in order to account properly for the motion of discontinuities and maxima of high-gradient regions. Extensive perturbation calculations based on full potential nonlinear solutions have been carried out. These calculations cover a variety of flow and geometric parameter perturbations involving isolated airfoils and compressor cascades at both subsonic and transonic flow conditions. Particular emphasis was placed on supercritical transonic flows which exhibit large surface shock movements over the parameter range studied; and on subsonic flows which display large pressure variations in the stagnation and peak suction pressure regions. Perturbation results for single-parameter perturbations, characterized by both extreme solution interpolation using widely separated base flow solutions and extreme solution extrapolation using closely spaced based flow solutions, were obtained in order to determine the accuracy and range of validity of the method. Additionally, calculation of perturbation results were made to investigate the effectiveness of employing piecewise-continuous straining functions rather than polynomial (quadratic, cubic, quartic) functions. Multi-shock and other complex flow situations were studied in order to examine the capability for treating general transonic flows.

Comparisons of the perturbation results with the corresponding 'exact' nonlinear solutions indicate a remarkable accuracy and range of validity of the perturbation method across the spectrum of examples reported. Geometry and flow parameter perturbations are treatable with equal success. Solution interpolation and extrapolation are both feasible. Results evaluating the polynomial and piecewise-continuous straining functions indicate that the piecewise-continuous functions are superior. The latter class of straining functions eliminate both the problem of unwanted straining in the domain of interest, as well as the problem of spurious straining out of the domain. Finally, it was demonstrated that this procedure can successfully treat flows containing multiple shocks and high-gradient regions by simultaneously straining all of these characteristic points. Computational time of the method, beyond the determination of the base solutions, is trivial. A code encompassing these developments has been written for the single-parameter perturbation problem

and is included as part of this report. Based on these results, we conclude that such a perturbation procedure can provide a means for substantially reducing computational requirements in design studies or other applications where large numbers of related nonlinear solutions are needed. Further development is needed, however, to provide a computational tool of wide utility. Because of the practical need in design or parametric studies to consider variations in several parameters simultaneously, we suggest the development of the capability for multiple-parameter perturbations, making full use of the current developments of the single parameter procedure. That procedure should incorporate a limiting-parameter calculation whereby the parameter bounds with respect to each varied parameter are determined. Finally, in order to demonstrate their ultimate power and utility, these procedures should now be tested by actual application to a practical problem which involves the high-frequency use of expensive computational codes in order to determine a large number of related flow solutions. We suggest transonic turbomachinery blade design optimization studies as both feasible and of high current importance.

## APPENDIX A - USER'S MANUAL FOR COMPUTER PROGRAM PERTURB

### A.1 INTRODUCTION

The purpose of this appendix is to describe the operation of the computer code which was developed in conjunction with the theoretical work presented in this report, and to provide sufficient detail to permit convenient use and change of the program. The program computes and plots an arbitrary flow variable on a contour surface by employing the strained-coordinate perturbation method previously discussed. The plot package included in this version refers to system routines at the Stanford University Center for Information Processing facility. In general, the plotting software must be supplied by the user according to the requirements of his operating system. This can be accomplished directly by replacing or modifying the subroutines PLOT, LIMITS, and ROUND.

A description of the general operating procedure of the program is given, together with complete description of both input and output. The program is written in FORTRAN IV and has been developed on an IBM 3033 computer. Typical run times are 1 to 3 seconds. The storage requirements are 50K<sub>10</sub>.

## A.2 PROGRAM DESCRIPTION

The program calculates both continuous and discontinuous nonlinear perturbation solutions which represent a single-parameter change in either geometry or flow conditions by employing a strained-coordinate procedure. The method utilizes a unit perturbation, determined from two previously calculated solutions ('base' and 'calibration' solutions) obtained from an 'expensive' computational procedure and displaced from one another by some reasonable change in geometry or flow variable, to predict new nonlinear solutions over a range of parameter variation.

This version of the procedure is configured to predict and plot an arbitrary flow variable (e.g., pressure coefficient) on the surface of a blade or airfoil, and can account for the motion of:

1. one or more critical points (shock points),
2. a stagnation point,
3. a maximum-suction-pressure point,

or simultaneously for any combination of these.

The program is also configured to compare the perturbation-predicted solutions with the corresponding 'exact' solutions obtained by employing the same 'expensive' computational procedure used to determine the base and calibration solutions.

The coordinate straining employed is piecewise linear with the end points and up to six interior points held invariant. At the option of the user, these additional interior points may be arbitrarily preselected, or chosen from among the minimum, maximum, and critical points automatically located by the program itself.

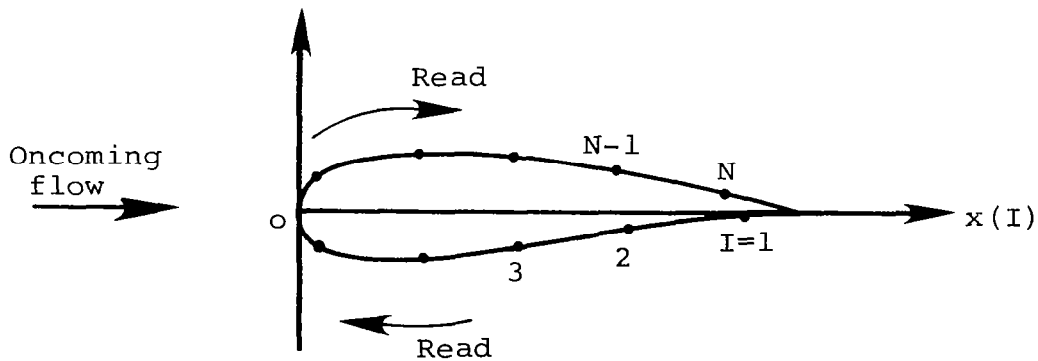
Critical or shock points are located on the basis of a user-supplied statement function defining the critical value of the dependent variable as a function of some single flow variable. The program default is with dependent variable  $y$  defined as pressure coefficient, with the independent variable being Mach number. In this case, the critical value is defined as

$$y_{\text{crit}} = C_p^* = \frac{2}{\gamma M_\infty^2} \left[ \left( \frac{2 + (\gamma - 1)M_\infty^2}{\gamma + 1} \right)^{\frac{\gamma}{\gamma - 1}} - 1 \right] \quad (\text{A-1})$$

where  $\gamma$  is the ratio of specific heats. If instead of surface pressure coefficient, the surface velocity distribution were used, then the value of  $y_{\text{crit}}$  would be given by

$$y_{\text{crit}} = \frac{V^*}{V_\infty} = \left( \frac{\gamma + 1}{2 + (\gamma - 1)M_\infty^2} \right)^{\frac{1}{\gamma - 1}} \quad (\text{A-2})$$

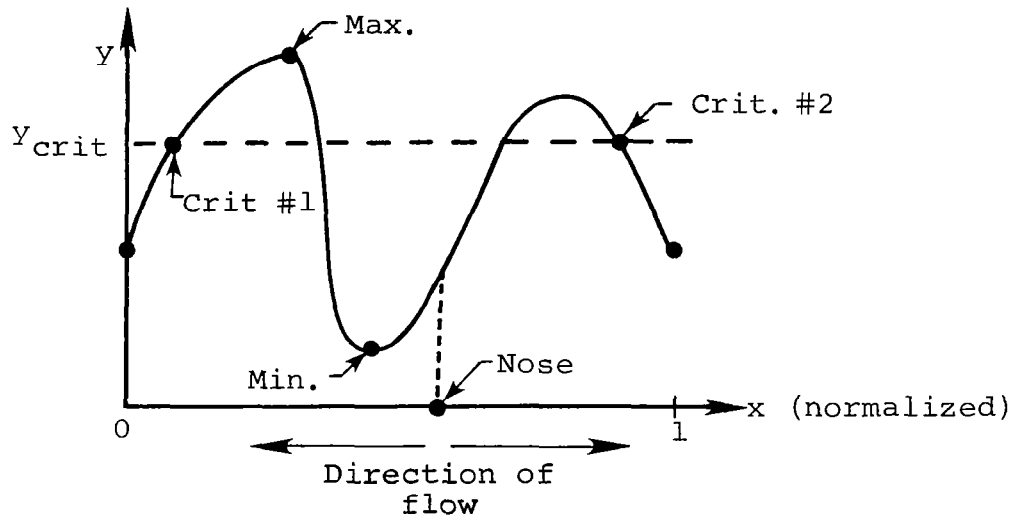
Data for base, calibration, and comparison solutions (if available) are input as an array  $x(I)$  of coordinates and a corresponding array  $y(I)$  giving the dependent variable at each coordinate location, where  $1 \leq I \leq N$  and  $N \leq 200$ .



The leading edge is at  $x = 0$ ; the data are read in beginning on the lower surface at the point farthest from the leading edge and proceeding clockwise around the surface as shown in the sketch. Data for the different solutions need not correspond to identical locations on the surface, except for the initial and final points, i.e.,  $x(1)$  and  $x(N)$  must be the same for all cases. The program normalizes the  $x$  coordinates ( $0 \leq x \leq 1$ ) such that  $x = 0$  corresponds to  $I = 1$  and  $x = 1$  to  $I = N$ .



The base and calibration solutions are searched for minimum, maximum, and critical points, e.g.,

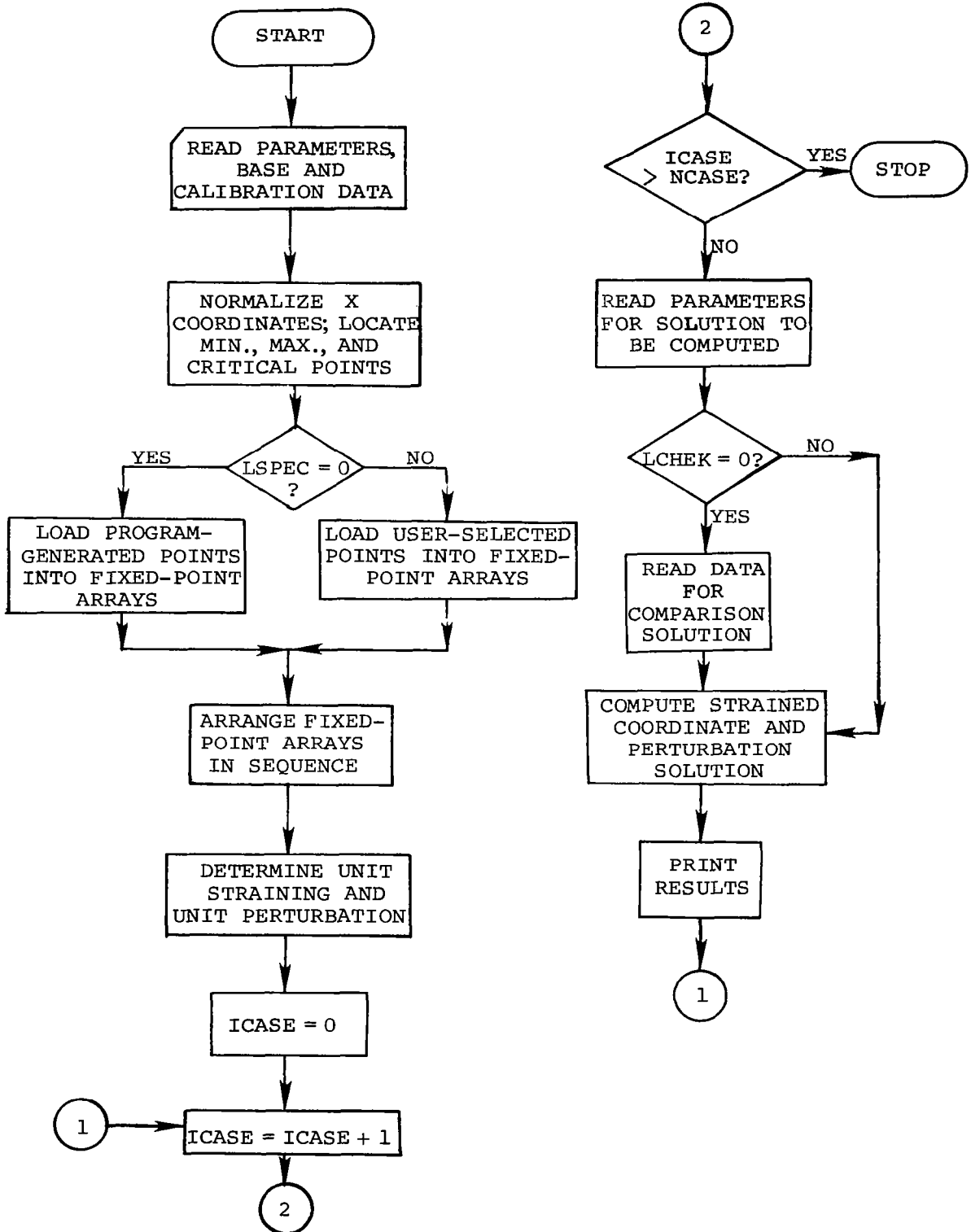


Note that the sign of  $dy/dx$  in physical coordinates is used in determining the critical points. For example, both critical points indicated on the above figure correspond to  $dy/dx > 0$  in physical coordinates, since at point #1 the physical coordinate increases in the direction from right to left, whereas at point #2 it increases from left to right.

The points to be held invariant in straining are either selected from among those located by the program or individually specified by the user, after which the unit coordinate straining and unit perturbation are computed.

Data for the test cases is then read in and nonlinear perturbation solutions constructed from the unit perturbation.

A.3 PROGRAM FLOW CHART



#### A.4 DICTIONARY OF INPUT VARIABLES

- A           Scaling parameter in straining procedure.  $A = -x(1)$ , where  $x(1)$  is location of first data point on lower surface (see PROGRAM DESCRIPTION).
- B           Scaling parameter in straining procedure.  $B = x(N)$ , where  $x(N)$  is location of last data point on upper surface (see PROGRAM DESCRIPTION).
- LCHEK       Specifies whether or not perturbation solution is to be compared with an exact solution.
- LCHEK = 0 ... no comparison  
            LCHEK = 1 ... comparison
- LECHO       Controls whether or not input deck is printed.
- LECHO = 0 ... no print  
            LECHO = 1 ... print
- LOC0(I)     Array of length 6 containing subscripts of user-specified invariant points in base solution; operational only when LSPEC = 1.
- LOC1(I)     Array of length 6 containing subscripts of user-specified invariant points in calibration solution; operational only when LSPEC = 1.
- LPERT       Specifies type of perturbation; operational only when LCHEK = 1 and only affects output from plot subroutine.
- LPERT = 1 ... thickness-ratio perturbation  
            LPERT = 2 ... angle-of-attack perturbation  
            LPERT = 3 ... Mach-number perturbation
- LSELCT(I)   Array of length 6 of which NSELCT elements are read in; operational only when LSPEC = 0, and specifies nature of points to be held invariant according to the code:
- 1 ... minimum point held invariant  
            2 ... maximum point held invariant  
            3 ... 1st critical point held invariant  
            4 ... 2nd critical point held invariant  
            5 ... 3rd critical point held invariant  
            6 ... 4th critical point held invariant

Note that critical point ordering is determined from order of occurrence starting at the lower surface at the point furthest from the leading edge and proceeding clockwise around the surface (see PROGRAM DESCRIPTION).

Note that the code numbers can be assigned in any order, e.g.,

LSELCT(1) = 1		LSELCT(1) = 4
LSELCT(2) = 3	and	LSELCT(2) = 1
LSELCT(3) = 4		LSELCT(3) = 3

are equivalent, both corresponding to NSELCT = 3, with the minimum, and first and second critical points held invariant.

LSPEC

Controls how invariant points in straining are specified.

LSPEC = 0 ... invariant points selected from among those located by the program, using the array LSELCT(I)

LSPEC = 1 ... invariant points preselected by user, using the arrays LOC0(I), LOC1(I)

LUNIT

Controls whether or not unit coordinate straining and unit perturbation are printed.

LUNIT = 0 ... no print

LUNIT = 1 ... print

M0,M1,M2

Oncoming Mach numbers in base, calibration, and perturbation solutions.

N

Number of locations for which data are input for base, calibration, and comparison solutions.

NAME

Character string of length 2 which symbolizes dependent variable, e.g., "CP" for pressure coefficient.

NCASE

Number of cases for which perturbation solutions are to be computed.

NSELCT

Number of points (in addition to end points) to be held invariant in straining; note:  $1 \leq \text{NSELCT} \leq 6$ .

Q0,Q1,Q2

Values of perturbation parameter in base, calibration, and perturbation solutions.

TITLE

Character string of length 80; identifies job and is printed as headline on first page of output.

XBASE(I),XCALB(I),XCHEK(I)...

Arrays of surface coordinates in base, calibration,  
and comparison solutions.

YBASE(I),YCALB(I),YCHEK(I)...

Arrays of dependent variables in base, calibration,  
and comparison solutions.

## A.5 PREPARATION OF INPUT DATA

### A.5.1 Description of Input

- Item 1 One card, containing the parameters N, NCASE, LSPEC, LECHO, LUNIT, LCHEK, LPERT.
- Item 2 One card, containing either
- (a) NSELCT, (LSELCT(I), I=1,NSELCT)
  - (b) NSELCT, (LOC0(I), I=1,NSELCT),  
(LOC1(I), I=1,NSELCT)
- where (a) and (b) correspond to LSPEC = 0 and LSPEC = 1, respectively.
- Item 3 One card, containing the character string TITLE.
- Item 4 One card, containing the character string NAME.
- Item 5 One card, containing the scaling parameters A and B.
- Item 6 One card, containing M0(real) and Q0.
- Item 7 One set of K cards, where  $K = 1 + \text{INT}(N/8)$ , containing data for x coordinate in base solution.
- Item 8 One set of K cards, K as above, containing data for dependent variable in base solution.
- Item 9 One card, containing M1(real) and Q1.
- Item 10 One set of K cards, K as above, containing data for x coordinate in calibration solution.
- Item 11 One set of K cards, K as above, containing data for dependent variable in calibration solution.

- Item 12      One card, containing M2(real) and Q2.
- Item 13      One set of K cards, K as above, containing data for  
x coordinate in comparison solution. This item is  
required only when LCHEK = 1.
- Item 14      One set of K cards, K as above, containing data for  
dependent variable in comparison solution. This item  
is required only when LCHEK = 1.

Note: Items 12-14 are required, in sequence, as many times as  
specified by NCASE.

### A.5.2 Format of Input Data

34

Item no. 1: 1 card

Variable	N	NCASE	LSPEC	LECHO	LUNIT	LCHEK	LPERT
Card column	5	10	15	20	25	30	35
Format type	I	I	I	I	I	I	I

Item no. 2a (LSPEC = 0): 1 card

Variable	NSELCT	LSELCT(1)	LSELCT(2)	----	LSELCT(NSELCT)		
Card column	5	10	15	20	25	30	35
Format type	I	I	I	I	I		

Item no. 2b (LSPEC = 1): 1 card

Variable	NSELCT	LOCO(1)	----	LOCO(NSELCT)	LOCI(1)	----	LOCI(NSELCT)
Card column	5	10	15	20	25	30	35
Format type	I	I	I	I	I	I	I

Item no. 3: 1 card

Variable	TITLE							
Card column	10	20	30	40	50	60	70	80
Format type	A							

Item no. 4: 1 card

Variable	NAME
Card column	2
Format type	A

Item no. 5: 1 card

Variable	A	B
Card column	10	20
Format type	F	F

Item no. 6: 1 card

Variable	MO	Q0	
Card column	10	20	
Format type	F	F	

Item no. 7: K cards,  $K = 1 + \text{INT}(N/8)$ , 8 values per card

Variable	XBASE(1)	XBASE(2)	XBASE(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F

Item no. 8: K cards, K as above, 8 values per card

Variable	YBASE(1)	YBASE(2)	YBASE(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F

Item no. 9: 1 card

Variable	M1	Q1	
Card column	10	20	
Format type	F	F	

Item no. 10: K cards, K as above, 8 values per card

Variable	XCALB(1)	XCALB(2)	XCALB(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F

Item no. 11: K cards, K as above, 8 values per card

Variable	YCALB(1)	YCALB(2)	YCALB(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F



Item no. 12: 1 card

Variable	M2	Q2	
Card column	10	20	
Format type	F	F	

Item no. 13: K cards,  $K = 1 + \text{INT}(N/8)$ , 8 values per card

Variable	XCHEK(1)	XCHEK(2)	XCHEK(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F

Item no. 14: K cards, K as above, 8 values per card

Variable	YCHEK(1)	YCHEK(2)	YCHEK(3)	----				
Card column	10	20	30	40	50	60	70	80
Format type	F	F	F	F	F	F	F	F

## A.6 DESCRIPTION OF OUTPUT

The first output item consists of a banner page, and the card images of the input data, the latter only if LECHO = 1.

The second item is a page headed by the job title, listing:

1. the input parameters relevant to the actual calculation;
2. the critical values of the dependent variable;
3. the locations of the minimum, maximum, and critical points found by the program;
4. the straining points selected;
5. the invariant points.

Results for unit straining of XBASE, and unit perturbation of the dependent variable are the third item output; this is done only if LUNIT = 1.

The fourth item (repeated for each case computed) summarizes the results of the calculation. The Mach number, the value of the perturbation parameter, and the critical value of the dependent variable are printed first, followed by the locations of the minimum, maximum, and critical points in the perturbation solution and comparison solution (if any). Then follows a table listing XBASE, YBASE, XCALB, YCALB, XPERT (the strained coordinate), and YPERT (the computed value of the dependent variable). If LCHEK = 1, three additional columns list XCHEK, YCHEK, and YPERT(INT), the latter being interpolated values of YPERT (the computed solution) at the points given by XCHEK. This allows direct numerical comparison of YPERT with YCHEK, since the values of XPERT and XCHEK do not coincide in general.

## A.7 ERROR MESSAGES

NUMBER OF CRITICAL POINTS IN  
BASE AND CALIBRATION SOLUTIONS  
ARE UNEQUAL - CALCULATION ENDED

This message will be printed if critical points are specified in straining (LSPEC = 0) and the number of critical points in base and calibration solutions are unequal. The remedy is to avoid use

of critical points in straining, or to use base and calibration solutions having equal numbers of critical points.

NUMBER OF CRITICAL POINTS  
SELECTED EXCEEDS NUMBER  
ACTUALLY LOCATED - CALCULATION  
ENDED

This message will be printed if more critical points are specified in straining (LSPEC = 0) than the number located by the program. The remedy is to specify a number of points less than or equal to the actual number.

ORDER OF SPECIFIED POINTS IN  
BASE AND CALIBRATION SOLUTIONS  
DOES NOT CORRESPOND

This message will be printed if the fixed points specified (LSPEC = 0) occur in a different sequence in the base and calibration solutions. The remedy is to use base and calibration solutions having the same qualitative features.

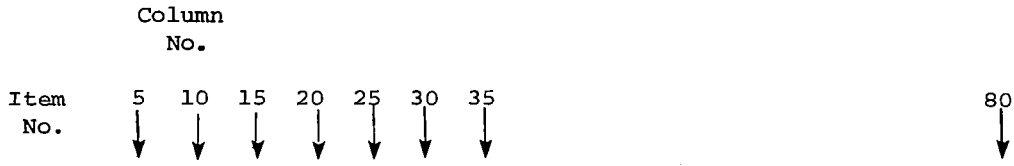
## A.8 SAMPLE CASE

The sample case presented in this section provides results (6 perturbation calculations and comparisons with 'exact' nonlinear solutions) for a multiple-shock flow for which partial results were provided in figure 14 of the main text. The calculation is for angle-of-attack perturbations of full potential flows past an isolated NACA 0012 airfoil at  $M_\infty = 0.80$ . The base and calibration angles-of-attack are  $\alpha_b = 0.500^\circ$  and  $\alpha_c = 0.200^\circ$ . Perturbation results are determined at  $\alpha = 0.00^\circ, 0.10^\circ, 0.30^\circ, 0.40^\circ, 0.60^\circ, \text{ and } 0.70^\circ$  and are compared with previously-calculated 'exact' nonlinear flows at those angles.

The input data is tabulated in figure A.1, with item numbers corresponding to those identified in Section A.5.1 and A.5.2. The first card, item 0, indicates that there are 149 points ( $N=149$ ) at which data will be input for the base, calibration, and comparison solutions; that there will be 5 cases ( $NCASE=6$ ) for which perturbation solutions are to be computed, that the invariant points will be located by the program ( $LSEPC=0$ ), that the input card deck will not be printed ( $LECHO=0$ ), that the information regarding the unit perturbation will be printed ( $LUNIT=1$ ), that there will be a comparison of the perturbation results with the exact solution ( $LCHEK=1$ ), and that the plot output will denote an angle-of-attack perturbation ( $LPERT=2$ ). The second card, item 2a, indicates that there will be three invariant points ( $NSELECT=3$ ) in addition to the end points; and that those points will be (1) where the maximum occurs ( $LSELECT(1)=2$ ) i.e. the stagnation point, (2) the first critical point ( $LSELECT(2)=3$ ) i.e. the 1st shock point found when moving forward on the bottom surface from the trailing edge, and (3) the second critical point ( $LSELECT(3)=4$ ) i.e. the 2nd shock point. The next card, item 3, contains the identifying title. On the next card, item 4, the 2 length character string indicates that the dependent variable for print output will be symbolized by a 'CP' denoting pressure coefficient. Item 5 indicates that the coordinates of the data points to be read in will start at  $x=1.0$  on the upper surface (refer to descriptions in A.4). The next card, item 6, indicates that the base flow values of Mach number and perturbation parameter (angle-of-attack in this case) are  $M_0 = 0.80$  and  $Q_0 = 0.50$ , respectively. The following 19 cards, item 7, provide the 149 base flow values of the surface coordinates, while the next 19 cards, item 8, provide the 149 base flow values of the dependent variable (pressure coefficient). Items 9, 10, and 11 indicate for the calibration flow the corresponding information given by the items 6,7, and 8 for the base flow. Items 12, 13, and 14, of which there are six sets corresponding to the 6 cases to be studied, provide analogous information as items 6,7, and 8, but now refer to the 'exact' nonlinear results. These, of course have been previously computed at the indicated

values of angle-of-attack ( $Q_2$ ) given in Item 12, and are included here for comparative purposed to enable assessment of the perturbation results.

Figure A.2 provides an abbreviated print output for the sample case, while figure A.3 provides the plot output of the results for the six cases, and display the base (.....), calibration (----), perturbation (\*\*\*\*), and 'exact' nonlinear (——) flow solutions.



1	149	6	0	0	1	1	2
2	3	2	3	4			
3	SAMPLE CASE - ALPHA PERTURBATION FOR MULTI-SHOCK ISOLATED AIRFOIL FLOW						
4	CP						
5	-1.0	1.0					
6	.800000	.500000					
7	.999495	.997758	.995132	.991597	.987091	.981656	.975290
	.959827	.950749	.940799	.930018	.918425	.906020	.892863
	.864393	.849140	.833235	.816758	.799690	.782110	.764038
	.726603	.707279	.687646	.667702	.647510	.627108	.606518
	.564955	.544062	.523144	.502239	.481369	.460595	.439937
	.395095	.378972	.359110	.339530	.320253	.301322	.282758
	.246838	.229547	.212731	.196413	.180616	.165342	.150655
	.123054	.110166	.097919	.086338	.075427	.065211	.055678
	.038773	.031422	.024818	.018980	.013897	.009586	.006089
	.001483	.000362	.000000	.000362	.001483	.003379	.006089
	.013897	.018980	.024818	.031422	.038773	.046895	.055678
	.075427	.086338	.097919	.110166	.123054	.136558	.150655
	.180616	.196413	.212731	.229547	.246838	.264582	.282758
	.320253	.339530	.359110	.378972	.399095	.419417	.439937
	.481369	.502239	.523144	.544062	.564955	.585801	.606518
	.647510	.667702	.687646	.707279	.726603	.745536	.764038
	.799690	.816758	.833235	.849140	.864393	.878994	.892863
	.918425	.930018	.940799	.950749	.959827	.968014	.975290
	.987091	.991597	.995132	.997758	.999495		
8	.454705	.494256	.414440	.401123	.370901	.335399	.303686
	.247635	.222849	.198568	.174301	.153539	.133375	.110559
	.072839	.053774	.034598	.016276	-.000859	-.019323	-.037707
	-.071935	-.090813	-.109358	-.126685	-.146033	-.164340	-.183752
	-.221343	-.241733	-.260834	-.278560	-.295525	-.311898	-.326422
	-.403291	-.497623	-.600835	-.668220	-.700287	-.711940	-.712103
	-.699075	-.687197	-.672616	-.656143	-.635506	-.613957	-.591617
	-.530317	-.496814	-.460480	-.418394	-.372256	-.319692	-.268845
	-.131694	-.056739	.037329	.148457	.285927	.443489	.629985
	1.019305	1.145394	1.161699	1.053470	.849325	.606187	.371494
	.005374	-.133558	-.243041	-.335771	-.410394	-.483396	-.543403
	-.638836	-.679361	-.716673	-.749187	-.779281	-.807499	-.834131
	-.876312	-.896774	-.914886	-.932147	-.947948	-.961781	-.974530
	-.997187	-1.006749	-1.016031	-1.022717	-1.029553	-1.034871	-1.038277
	-1.041697	-1.039425	-1.033151	-1.019966	-.993560	-.947606	-.857172
	-.047182	.060158	.067065	.059619	.053938	.049536	.050777
	.061931	.069406	.079745	.092114	.105449	.119529	.134216
	.170576	.188732	.210704	.233000	.256263	.281298	.309273
	.374275	.403616	.416177	.425153	.4294705		
6	.800000	.200000					
7	.998133	.994611	.990123	.984707	.978365	.971116	.962959
	.944001	.933260	.921710	.909350	.896242	.882424	.867877
	.836834	.820417	.803411	.785896	.767891	.749457	.730593
	.691779	.671909	.651791	.631465	.610951	.590310	.569541
	.527234	.507056	.486263	.465566	.444984	.424539	.404292
	.364455	.344946	.325740	.306870	.288381	.270271	.252590
	.218602	.202340	.186596	.171371	.156730	.142674	.129207
	.104130	.092568	.081668	.071455	.061912	.053105	.044940
	.030821	.024841	.019557	.014959	.011060	.007786	.005125
	.001399	.000360	.000000	.000360	.001399	.002995	.005125
	.011060	.014959	.019557	.024841	.030821	.037523	.044940
	.061912	.071455	.081668	.092568	.104130	.116351	.129207
	.156730	.171371	.186596	.202340	.218602	.235359	.252590
	.288381	.306878	.325740	.344946	.364455	.384243	.404292
	.444984	.465566	.486263	.507056	.527884	.548725	.569541

Figure A.1- Card input for sample case

Item  
No.

7	.610951	.631465	.651791	.671909	.691779	.711341	.730593	.749457
	.767891	.785896	.803411	.820417	.836834	.852680	.867877	.882424
	.896242	.909350	.921710	.933260	.944001	.953914	.962959	.971116
8	.978365	.984707	.990123	.994611	.998133			
	.449164	.448979	.368724	.356677	.323122	.286384	.256968	.228251
	.201234	.174994	.152770	.131478	.107828	.088414	.069020	.049604
	.030206	.011790	-.005341	-.023660	-.041742	-.059474	-.074743	-.092628
	-.109767	-.125161	-.141841	-.156432	-.170621	-.181949	-.190490	-.196776
	-.198872	-.203294	-.235719	-.367492	-.586069	-.746810	-.804612	-.822557
	-.830010	-.830430	-.828029	-.823224	-.815167	-.805731	-.794610	-.780953
	-.765371	-.748500	-.728200	-.707375	-.685933	-.658961	-.629807	-.598986
	-.555587	-.528007	-.488403	-.441165	-.394459	-.336198	-.271604	-.211304
	-.136609	-.051358	.052134	.163559	.296591	.449001	.621890	.810551
	.994877	1.129638	1.169160	1.096662	.933160	.727769	.526208	.346353
	.190742	.056809	-.053975	-.156994	-.241464	-.316023	-.377381	-.443985
	-.500232	-.542507	-.588047	-.626762	-.662591	-.693996	-.723334	-.751171
	-.777519	-.798919	-.819680	-.839982	-.857221	-.873286	-.888000	-.901015
	-.912830	-.923922	-.932634	-.940086	-.946688	-.949963	-.952335	-.951561
	-.946329	-.937133	-.917745	-.885770	-.780090	-.450920	-.135657	-.055363
	-.048879	-.054497	-.059285	-.058713	-.056033	-.048960	-.039039	-.030023
	-.017360	-.003396	.011570	.025968	.042113	.059613	.077446	.095514
	.113310	.136492	.156974	.178507	.204143	.230640	.258908	.289933
	.324329	.357590	.369390	.449350	.449164			
6	.800000	.000000						
7	.999495	.997758	.995132	.991597	.987091	.981656	.975290	.968014
	.959827	.950749	.940799	.930018	.918425	.906020	.892863	.878994
	.864393	.849140	.833235	.816753	.799690	.782110	.764038	.745536
	.726603	.707279	.687646	.667702	.647510	.627108	.606518	.585801
	.564955	.544062	.523144	.502239	.481369	.460595	.439937	.419417
	.399095	.378972	.359110	.339530	.320253	.301322	.282758	.264582
	.246838	.229547	.212731	.196413	.180616	.165342	.150655	.136558
	.123054	.110166	.097919	.086338	.075427	.065211	.055678	.046895
	.038773	.031422	.024818	.018980	.013897	.009586	.006089	.003379
	.001403	.000362	.000000	.000362	.001483	.003379	.006089	.009586
	.013897	.018990	.024818	.031422	.038773	.046895	.055678	.065211
	.075427	.086338	.097919	.110166	.123054	.136558	.150655	.165342
	.180616	.196413	.212731	.229547	.246838	.264582	.282758	.301322
	.320253	.339530	.359110	.378972	.399095	.419417	.439937	.460595
	.481369	.502239	.523144	.544062	.564955	.585801	.606518	.627108
	.647510	.667702	.687646	.707279	.726603	.745536	.764038	.782110
	.799690	.816753	.833235	.849140	.864393	.878994	.892863	.906020
	.918425	.930018	.940799	.950749	.959827	.968014	.975290	.981656
	.987091	.991597	.995132	.997758	.999495			
	8	.436539	.485539	.406037	.392728	.362350	.326615	.294714
.238623		.213577	.189264	.164999	.144313	.124278	.101626	.083098
.064554		.045974	.027430	.009934	-.006143	-.023238	-.039876	-.055862
-.068951		-.084043	-.097616	-.108363	-.118872	-.125296	-.128789	-.127030
-.123011		-.131088	-.139768	-.1450296	-.149836	-.1539648	-.1566610	-.1582942
-.1889736		-.1891484	-.1891050	-.1886414	-.1880317	-.1872670	-.1862403	-.1851047
-.1838181		-.1822989	-.1805990	-.1787772	-.1766224	-.1744070	-.1721154	-.1692722
-.1662052		-.1629341	-.1593876	-.1553475	-.1510373	-.1459592	-.1407829	-.1341769
-.1267177		-.1193761	-.1101301	.008967	.147792	.309038	.505174	.724631
.941542		1.107686	1.170402	1.107686	.941542	.724631	.505174	.309038
.147792		.008967	-.101302	-.193761	-.267177	-.341769	-.407829	-.459592
-.510373		-.553475	-.593876	-.629341	-.662053	-.692722	-.721154	-.744070
-.766224		-.787772	-.805991	-.822989	-.838182	-.851047	-.862403	-.872670
-.880317		-.886415	-.891050	-.891484	-.889736	-.882942	-.868610	-.839649
-.729837		-.450298	-.199788	-.131088	-.123011	-.127030	-.128789	-.125296
-.118872		-.108362	-.097616	-.064043	-.068951	-.055862	-.039376	-.023238
-.006143		.009934	.027430	.045974	.064554	.083098	.101626	.124277

Figure A.1- Continued

Item  
No.

8	.144313	.164999	.189264	.213576	.238623	.265263	.294714	.326615
	.362350	.392727	.406037	.486539	.486539			
6	.800000	.100000						
7	.998831	.996261	.992734	.988238	.982815	.976464	.969203	.961034
	.951976	.942049	.931291	.919724	.907346	.894219	.880381	.865812
	.850593	.834722	.818282	.801251	.783710	.765678	.747217	.728326
	.709045	.689455	.669555	.649407	.629051	.608507	.587835	.567035
	.546189	.525317	.504458	.483635	.462907	.442295	.421820	.401543
	.381464	.361646	.342110	.322876	.303986	.285462	.267326	.249620
	.232366	.215585	.199302	.183537	.168293	.153635	.139563	.126082
	.113216	.100986	.089419	.078517	.068305	.058770	.049977	.041834
	.034450	.027794	.021880	.016686	.012213	.008485	.005447	.003104
	.001403	.000361	.000000	.000361	.001403	.003104	.005447	.008485
	.012213	.016686	.021880	.027794	.034450	.041834	.049977	.058770
	.068305	.078517	.089419	.100986	.113216	.126082	.139563	.153635
	.168293	.183537	.199302	.215585	.232366	.249620	.267326	.285462
	.303986	.322876	.342110	.361646	.381464	.401543	.421820	.442295
	.462907	.483635	.504458	.525317	.546189	.567035	.587835	.608507
	.629051	.649407	.669555	.689455	.709045	.728326	.747217	.765812
	.783710	.801251	.818282	.834722	.850593	.865812	.880381	.894219
	.907346	.919724	.931291	.942049	.951976	.961034	.969203	.976464
	.982815	.988238	.992734	.996261	.998831			
8	.458893	.458791	.378081	.371267	.338882	.305476	.273846	.245302
	.218790	.193380	.168300	.146953	.126374	.103259	.084303	.065328
	.046302	.027281	.009255	-.007440	-.025268	-.042763	-.059836	-.074259
	-.091057	-.106811	-.120384	-.134618	-.145906	-.155618	-.161043	-.162445
	-.162500	-.168324	-.216625	-.409216	-.665613	-.798270	-.838444	-.852963
	-.858704	-.860938	-.858258	-.853377	-.847175	-.837775	-.827215	-.815116
	-.800639	-.784352	-.766850	-.748036	-.724693	-.702694	-.675207	-.645547
	-.614010	-.580003	-.541413	-.500299	-.451419	-.402259	-.340192	-.270822
	-.204426	-.121671	-.025336	.093513	.226491	.387549	.572665	.774910
	.971786	1.120096	1.170101	1.103351	.940608	.733454	.524863	.336346
	.173240	.040723	-.078005	-.173334	-.256245	-.322903	-.393605	-.455428
	-.501918	-.549224	-.589936	-.627684	-.660684	-.691319	-.720232	-.747349
	-.762284	-.790569	-.811391	-.829099	-.845728	-.860684	-.873407	-.884928
	-.895825	-.904302	-.911372	-.917277	-.919415	-.920039	-.916639	-.907461
	-.892168	-.860794	-.748397	-.440787	-.163682	-.093649	-.088298	-.093776
	-.096178	-.095105	-.088954	-.081614	-.070719	-.057721	-.046263	-.031594
	-.016004	.000265	.015695	.032676	.050827	.069131	.087505	.105956
	.123637	.142566	.169896	.194712	.219897	.246217	.274591	.306084
	.339369	.371654	.378391	.458995	.456893			
6	.800000	.300000						
7	.999455	.997755	.995132	.991597	.987091	.981656	.975290	.968014
	.959527	.950749	.940799	.930018	.918425	.906020	.892863	.878994
	.864393	.849140	.833235	.816758	.799690	.782110	.764038	.745536
	.726603	.707279	.687646	.667702	.647510	.627108	.606518	.585801
	.564955	.544062	.523144	.502239	.481369	.460595	.439937	.419417
	.399095	.378972	.359110	.339530	.320253	.301322	.282758	.264502
	.246833	.229547	.212731	.196413	.180616	.165342	.150655	.136558
	.123054	.110166	.097919	.086338	.075427	.065211	.056878	.046895
	.038773	.031422	.024818	.018980	.013897	.009586	.006089	.003379
	.001483	.000362	.000000	.000362	.001483	.003379	.006089	.009586
	.013897	.018980	.024818	.031422	.038773	.046895	.056878	.065211
	.075427	.086338	.097919	.110166	.123054	.136553	.150655	.165342
	.180616	.196413	.212731	.229547	.246838	.264502	.282758	.301322
	.320253	.339530	.359110	.378972	.399095	.419417	.439937	.460595
	.481369	.502239	.523144	.544062	.564955	.585801	.606518	.627108
	.647510	.667702	.687646	.707279	.726603	.745536	.764038	.782110
	.799690	.816758	.833235	.849140	.864393	.878994	.892863	.906020
	.918425	.930018	.940799	.950749	.959827	.968014	.975290	.981656

Figure A.1- Continued



Item  
No.

7	.987091	.991597	.995132	.997758	.999495			
8	.489108	.488886	.408558	.395215	.364845	.329132	.297230	.267752
	.241058	.215931	.191510	.167103	.146228	.125956	.103014	.084120
	.065123	.045986	.026752	.008404	-.008728	-.027140	-.045429	-.063491
	-.079201	-.097685	-.115657	-.132155	-.150327	-.166937	-.183882	-.198952
	-.212493	-.225064	-.233421	-.238540	-.247696	-.284055	-.404125	-.589304
	-.720706	-.772628	-.790788	-.795047	-.794404	-.790553	-.782963	-.773702
	-.762569	-.748667	-.732627	-.715095	-.693840	-.671857	-.649105	-.620354
	-.582698	-.554654	-.516249	-.477942	-.432722	-.378478	-.326477	-.262355
	-.188943	-.114700	-.021243	.089693	.227968	.387372	.579277	.787676
	.908193	1.131475	1.167507	1.078406	.890077	.657697	.429117	.228601
	.066214	-.072675	-.182169	-.274184	-.342235	-.424712	-.488345	-.535603
	-.585089	-.627328	-.666135	-.699952	-.731234	-.760614	-.788219	-.810703
	-.832370	-.853387	-.871153	-.888087	-.903869	-.917736	-.930370	-.942268
	-.951914	-.960503	-.968506	-.973536	-.978179	-.980470	-.979557	-.976684
	-.966756	-.949158	-.918589	-.860190	-.600836	-.220297	-.035550	-.008065
	-.014402	-.021252	-.025646	-.024423	-.019218	-.013954	-.004358	.007052
	.019861	.032428	.046981	.063042	.079533	.096314	.113339	.134662
	.153582	.173302	.196701	.220258	.244641	.270691	.299605	.331015
	.366298	.396308	.409349	.489330	.489108			
6	.800000	.400000						
7	.998883	.996261	.992734	.988238	.982815	.976464	.969203	.961034
	.951976	.942049	.931291	.919724	.907346	.894219	.880381	.865812
	.850593	.834722	.818282	.801251	.783710	.765678	.747217	.728326
	.709045	.689455	.669555	.649407	.629051	.608507	.587835	.567035
	.546189	.525317	.504458	.483635	.462907	.442295	.421820	.401543
	.381464	.361646	.342110	.322876	.303936	.285462	.267326	.249620
	.232366	.215585	.199302	.183537	.168293	.153635	.139563	.126082
	.113216	.100986	.089419	.078517	.068305	.058770	.049977	.041834
	.034450	.027794	.021680	.016066	.0112213	.006485	.002447	.001304
	.001403	.000361	.000000	.000361	.001403	.003104	.005447	.008485
	.012213	.016686	.021680	.027794	.034450	.041834	.049977	.058770
	.068305	.078517	.089419	.100986	.113216	.126082	.139563	.153635
	.168293	.183537	.199302	.215585	.232366	.249620	.267326	.285462
	.303936	.322876	.342110	.361646	.381464	.401543	.421820	.442295
	.462907	.483635	.504458	.525317	.546189	.567035	.587835	.608507
	.629051	.649407	.669555	.689455	.709045	.728326	.747217	.765678
	.783710	.801251	.818282	.834722	.850593	.865812	.880381	.894219
	.907346	.919724	.931291	.942049	.951976	.961034	.969203	.976464
	.982815	.988238	.992734	.996261	.998883			
8	.463906	.463542	.383210	.376335	.344041	.310753	.279223	.250763
	.224307	.198933	.173878	.152520	.131821	.108699	.089614	.070459
	.051192	.031852	.013404	-.003844	-.022380	-.040809	-.059045	-.075003
	-.093807	-.112203	-.129293	-.146288	-.166067	-.184700	-.202242	-.219409
	-.237169	-.252367	-.264630	-.275050	-.287088	-.310622	-.384371	-.520500
	-.651464	-.723049	-.749049	-.757014	-.757358	-.752284	-.744825	-.735153
	-.722462	-.707479	-.699900	-.670467	-.649353	-.627514	-.599640	-.568836
	-.535928	-.501070	-.461913	-.418310	-.367155	-.318745	-.258751	-.190444
	-.123351	-.039674	.057341	.175989	.308200	.466303	.645304	.836752
	1.016671	1.141682	1.165213	1.072227	.857264	.664673	.446682	.253724
	.089560	-.043323	-.161762	-.257318	-.340367	-.408038	-.476755	-.534029
	-.576586	-.625309	-.664559	-.700900	-.732626	-.762251	-.790133	-.816546
	-.838046	-.850656	-.878883	-.896170	-.913119	-.928684	-.942283	-.954687
	-.936554	-.976330	-.985336	-.993910	-.999744	-1.005513	-1.009427	-1.010893
	-1.011613	-1.007216	-.993251	-.981310	-.951188	-.898942	-.675924	-.267956
	-.014439	.020454	.026006	.017221	.012851	.013018	.014088	.020132
	.028577	.038929	.049469	.062311	.076941	.092261	.108099	.124370
	.145106	.163688	.183282	.206792	.230823	.256119	.283563	.314198
	.346700	.378313	.384611	.464269	.463906			
6	.800000	.600000						

Figure A.1- Continued

Item  
No.

7	.998133	.994611	.990123	.984707	.978365	.971116	.962959	.953914	
	.944001	.933260	.921710	.909350	.896242	.882424	.867877	.852680	
	.836234	.820417	.803411	.785896	.767891	.749457	.730593	.711341	
	.691779	.671909	.651791	.631465	.610951	.590310	.569541	.548725	
	.527884	.507056	.486263	.465566	.444984	.424539	.404292	.384243	
	.364455	.344946	.325740	.306878	.288381	.270271	.252590	.235359	
	.218602	.202340	.186596	.171371	.156730	.142674	.129207	.116351	
	.104130	.092568	.081668	.071455	.061912	.053105	.044940	.037523	
	.030821	.024841	.019557	.014959	.011060	.007786	.005125	.002995	
	.001399	.000360	.000000	.000360	.001399	.002995	.005125	.007786	
	.011060	.014959	.019557	.024841	.030821	.037523	.044940	.053105	
	.061912	.071455	.081668	.092568	.104130	.116351	.129207	.142674	
	.156730	.171371	.186596	.202340	.218602	.235359	.252590	.270271	
	.288381	.306878	.325740	.344946	.364455	.384243	.404292	.424539	
	.444984	.465566	.486263	.507056	.527884	.548725	.569541	.590310	
	.610951	.631465	.651791	.671909	.691779	.711341	.730593	.749457	
	.767891	.785896	.803411	.820417	.836834	.852680	.867877	.882424	
	.896242	.909350	.921710	.933260	.944001	.953914	.962959	.971116	
	.978365	.984707	.990123	.994611	.998133				
	8	.462483	.461570	.382264	.370167	.336918	.302577	.271523	.243132
		.216423	.190478	.168483	.147409	.124007	.104765	.085520	.066233
		.046930	.028541	.011355	-.007085	-.025401	-.043535	-.059445	-.078202
		-.096599	-.113795	-.133011	-.151223	-.170603	-.189406	-.208636	-.229710
		-.250142	-.270238	-.291004	-.312540	-.331581	-.351660	-.373582	-.403607
		-.454352	-.519560	-.581649	-.625410	-.647867	-.656544	-.656523	-.649851
		-.638908	-.625234	-.606912	-.587538	-.567393	-.540859	-.511204	-.479994
		-.446358	-.407620	-.366183	-.319089	-.274597	-.218665	-.155032	-.093624
		-.017795	.068219	.171463	.282458	.413040	.560212	.723428	.895624
		1.054411	1.155458	1.157109	1.047605	.852882	.625219	.410777	.224683
		.066900	-.066937	-.176675	-.280304	-.366630	-.442649	-.499694	-.558762
		-.611902	-.653198	-.695634	-.731009	-.764265	-.793388	-.820540	-.846013
		-.870094	-.899922	-.909713	-.929431	-.946512	-.962757	-.977730	-.990991
		-1.003343	-1.015392	-1.025594	-1.035091	-1.044424	-1.051292	-1.058472	-1.066473
-1.062684		-1.073640	-1.075761	-1.077063	-1.076656	-1.073327	-1.063527	-1.045402	
-1.011530		-.950877	-.666516	-.212324	.066583	.111356	.110028	.101566	
.097627		.077297	.100041	.104126	.111381	.121190	.132408	.144782	
.158122		.176327	.192981	.211133	.233658	.257373	.283129	.311842	
.344092		.375500	.385985	.463397	.462483				
6		.800000	.700000						
7	.997307	.992822	.987410	.981072	.973828	.965677	.956638	.946732	
	.935998	.924455	.912104	.899005	.885196	.870659	.855472	.839636	
	.823231	.806236	.788732	.770739	.752317	.733466	.714226	.694678	
	.674821	.654716	.634403	.613903	.593275	.572520	.551718	.530890	
	.510077	.489297	.468614	.448046	.427614	.407381	.387345	.367569	
	.348073	.328880	.310029	.291544	.273445	.255774	.238553	.221805	
	.205551	.189815	.174597	.159960	.145909	.132444	.119590	.107368	
	.095902	.084896	.074673	.065117	.056291	.048101	.040650	.033903	
	.027865	.022504	.017806	.013774	.010327	.007438	.005001	.002992	
	.001397	.000359	.000000	.000359	.001397	.002992	.005001	.007438	
	.010327	.013774	.017806	.022504	.027865	.033903	.040650	.048101	
	.056291	.065117	.074673	.084896	.095802	.107368	.119590	.132444	
	.145909	.159960	.174597	.189815	.205551	.221805	.238553	.255774	
	.273445	.291544	.310029	.328880	.348073	.367569	.387345	.407381	
	.427614	.448046	.468614	.489297	.510077	.530890	.551718	.572520	
	.593275	.613903	.634403	.654716	.674821	.694678	.714226	.733466	
	.752317	.770739	.788732	.806236	.823231	.839636	.855472	.870659	
	.885196	.899005	.912104	.924455	.935998	.946732	.956638	.965677	
	.973828	.981072	.987410	.992822	.997307				
	8	.668728	.668307	.522300	.441943	.369762	.314595	.273042	.239321
		.209654	.185547	.163233	.139135	.119414	.099830	.080298	.060812

Figure A.1- Continued

8

.042266	.024934	.006399	-.011958	-.030082	-.045955	-.064578	-.082762
-.099687	-.118538	-.136324	-.155210	-.173483	-.192150	-.212669	-.232627
-.252466	-.273361	-.295668	-.316378	-.338889	-.361301	-.383949	-.409232
-.437231	-.471741	-.508264	-.539430	-.562879	-.577161	-.581808	-.579150
-.571404	-.557257	-.540964	-.523283	-.498949	-.471497	-.441655	-.409089
-.372779	-.334533	-.290600	-.249155	-.196686	-.137268	-.081064	-.012045
.064911	.156050	.250797	.360544	.482506	.618924	.770906	.932586
1.076892	1.162004	1.149382	1.027122	.829499	.598496	.380563	.198798
.049066	-.077266	-.187047	-.280101	-.373953	-.451564	-.515760	-.564158
-.619415	-.668576	-.705960	-.744879	-.777275	-.808043	-.834895	-.859770
-.883137	-.906457	-.925930	-.944910	-.963698	-.979917	-.995537	-1.010093
-1.023092	-1.035272	-1.047226	-1.057403	-1.066934	-1.076385	-1.083463	-1.090957
-1.097319	-1.102275	-1.108152	-1.111634	-1.114988	-1.117839	-1.119957	-1.119530
-1.117793	-.112354	-1.098409	-1.070169	-1.009848	-.890889	-.417645	.062612
.154654	.163569	.160082	.157352	.156120	.158521	.164013	.171487
.180676	.191367	.207425	.222816	.240828	.265001	.293844	.331050
.382321	.450996	.528020	.669150	.668728			

Figure A.1- Concluded

```

*****
*
*          PROGRAM PERTURB
*
*    CALCULATES NONLINEAR SINGLE-PARAMETER
*
*          CONTINUOUS OR DISCONTINUOUS
*
*          PERTURBATION SOLUTIONS
*
*    WHICH REPRESENT A CHANGE IN EITHER
*
*          GEOMETRY OR FLOW CONDITIONS
*
*    BY EMPLOYING A STRAINED-COORDINATE PROCEDURE
*
*    UTILIZING A UNIT PERTURBATION DETERMINED FROM
*
*          TWO PREVIOUSLY CALCULATED
*
*          'BASE' AND 'CALIBRATION' SOLUTIONS
*
*    DISPLACED FROM ONE ANOTHER BY SOME REASONABLE
*
*          CHANGE IN GEOMETRY OR FLOW CONDITION
*
*
*          WRITTEN BY
*
*    JAMES P. ELLIOTT AND STEPHEN S. STAHARA
*
*
*    NIELSEN ENGINEERING AND RESEARCH, INC.
*
*          MOUNTAIN VIEW, CALIFORNIA
*
*****

```

```

*****
*          SAMPLE CASE - ALPHA PERTURBATION FOR MULTI-SHOCK ISOLATED AIRFOIL FLOW
*
*****

```

<<<<<<<< LIST OF INPUT PARAMETERS >>>>>>>>

```

N = 149
A = -1.0   B = 1.0
BASE SOLUTION:      M0 = 0.8000   Q0 = 0.5000
CALIBRATION SOLN:  M1 = 0.8000   Q1 = 0.2000

```

Figure A.2- Abbreviated print output for sample case

<<<<<<<<< CRITICAL VALUES OF CP >>>>>>>>>

BASE SOLUTION:           CPCRT = -0.4346

CALIBRATION SOLN:       CPCRT = -0.4346

<<<<< LOCATIONS OF MIN., MAX., AND CRITICAL PTS. >>>>>

(\* DENOTES POINT ON LOWER SURFACE)

BASE SOLUTION:

MINIMUM AT X = 0.4606     (PPOINT #112)  
MAXIMUM AT X = 0.0000     (PPOINT # 75)  
2 CRITICAL POINT(S):  
  1ST AT X = 0.3924\*     (AFTER POINT # 41)  
  2ND AT X = 0.6288     (AFTER POINT #120)

CALIBRATION SOLN:

MINIMUM AT X = 0.4043     (PPOINT #111)  
MAXIMUM AT X = 0.0000     (PPOINT # 75)  
2 CRITICAL POINT(S):  
  1ST AT X = 0.4592\*     (AFTER POINT # 36)  
  2ND AT X = 0.5498     (AFTER POINT #118)

<<<<<<<<< STRAINING POINTS SELECTED >>>>>>>>>

NUMBER OF FIXED POINTS : 5

FIXED POINTS SELECTED (IN ADDITION TO END POINTS) :

POINT OF MAXIMUM CP  
CPCRT (1ST OF 2)  
CPCRT (2ND OF 2)

<<<<<<<<< LOCATION OF FIXED POINTS >>>>>>>>>

(\* DENOTES POINT ON LOWER SURFACE)

BASE SOLUTION:

XFIX(1) = 1.0000\*  
XFIX(2) = 0.3924\*  
XFIX(3) = 0.0000  
XFIX(4) = 0.6288  
XFIX(5) = 1.0000

CALIBRATION SOLN:

XFIX(1) = 1.0000\*  
XFIX(2) = 0.4592\*  
XFIX(3) = 0.0000  
XFIX(4) = 0.5498  
XFIX(5) = 1.0000

```

*****
* UNIT PERTURBATION OF CP *
*           C           *
* UNIT STRAINING OF XBASE *
*****

```

POINT	XBASE	XSTRUNIT	CPUNIT
1	0.9995	0.9996	0.1516
2	0.9978	0.9980	0.1503
3	0.9951	0.9957	-0.1153
4	0.9916	0.9925	-0.0350
5	0.9871	0.9885	0.0192
6	0.9817	0.9837	-0.0527
7	0.9733	0.9780	-0.0591
8	0.9680	0.9715	-0.0533
9	0.9619	0.9642	-0.0470
10	0.9507	0.9562	-0.0418
11	0.9408	0.9473	-0.0390
12	0.9300	0.9377	-0.0386
13	0.9184	0.9274	-0.0339
14	0.9060	0.9164	-0.0339
15	0.8929	0.9046	-0.0415
16	0.8790	0.8923	-0.0351
17	0.8644	0.8793	-0.0381
18	0.8491	0.8657	-0.0417
19	0.8332	0.8516	-0.0455
20	0.8168	0.8369	-0.0468
21	0.7997	0.8217	-0.0471
22	0.7821	0.8061	-0.0556
23	0.7640	0.7900	-0.0611
24	0.7455	0.7735	-0.0662
25	0.7266	0.7567	-0.0647
26	0.7073	0.7395	-0.0775
27	0.6876	0.7220	-0.0888
28	0.6677	0.7043	-0.0928
29	0.6475	0.6863	-0.1067
30	0.6271	0.6681	-0.1201
31	0.6065	0.6498	-0.1349
32	0.5858	0.6314	-0.1532
33	0.5650	0.6128	-0.1734
34	0.5441	0.5942	-0.2064
35	0.5231	0.5756	-0.2428
36	0.5022	0.5570	-0.2809
37	0.4814	0.5384	-0.3257
38	0.4606	0.5199	-0.3711
39	0.4399	0.5015	-0.3818
40	0.4194	0.4833	-0.3191
41	0.3991	0.4652	-0.1193
42	0.3790	0.4435	0.2948
43	0.3591	0.4203	0.5270
44	0.3395	0.3974	0.4753
45	0.3203	0.3748	0.4194
46	0.3013	0.3526	0.3944
47	0.2828	0.3309	0.3886
48	0.2646	0.3096	0.3889
49	0.2468	0.2889	0.3877
50	0.2295	0.2686	0.3917
51	0.2127	0.2490	0.3971

Figure A.2- Continued

52	0.1964	0.2299	0.3990
53	0.1806	0.2114	0.4079
54	0.1653	0.1935	0.4105
55	0.1507	0.1763	0.4084
56	0.1366	0.1598	0.4252
57	0.1231	0.1440	0.4357
58	0.1102	0.1289	0.4411
59	0.0979	0.1146	0.4454
60	0.0863	0.1010	0.4572
61	0.0754	0.0883	0.4672
62	0.0652	0.0763	0.4799
63	0.0557	0.0652	0.4718
64	0.0469	0.0549	0.4757
65	0.0388	0.0454	0.4779
66	0.0314	0.0368	0.4874
67	0.0248	0.0290	0.4954
68	0.0190	0.0222	0.4945
69	0.0139	0.0163	0.5134
70	0.0096	0.0112	0.5078
71	0.0061	0.0071	0.4605
72	0.0034	0.0040	0.3529
73	0.0015	0.0017	0.2112
74	0.0004	0.0004	0.0803
75	0.0000	0.0000	-0.0247
76	0.0004	0.0003	-0.1734
77	0.0015	0.0013	-0.3329
78	0.0034	0.0030	-0.4227
79	0.0061	0.0053	-0.4710
80	0.0096	0.0084	-0.4994
81	0.0139	0.0121	-0.4931
82	0.0190	0.0166	-0.5033
83	0.0248	0.0217	-0.4911
84	0.0314	0.0275	-0.4720
85	0.0388	0.0339	-0.4489
86	0.0469	0.0410	-0.4620
87	0.0557	0.0487	-0.4517
88	0.0652	0.0570	-0.4046
89	0.0754	0.0659	-0.4024
90	0.0863	0.0755	-0.3963
91	0.0979	0.0856	-0.3821
92	0.1102	0.0963	-0.3693
93	0.1231	0.1076	-0.3594
94	0.1366	0.1194	-0.3552
95	0.1507	0.1317	-0.3520
96	0.1653	0.1446	-0.3368
97	0.1806	0.1579	-0.3235
98	0.1964	0.1717	-0.3246
99	0.2127	0.1860	-0.3201
100	0.2295	0.2007	-0.3143
101	0.2468	0.2158	-0.3123
102	0.2646	0.2313	-0.3079
103	0.2828	0.2472	-0.3037
104	0.3013	0.2634	-0.3028
105	0.3203	0.2800	-0.2994
106	0.3395	0.2969	-0.2961
107	0.3591	0.3140	-0.2961
108	0.3790	0.3313	-0.2930
109	0.3991	0.3489	-0.2937
110	0.4194	0.3667	-0.2927
111	0.4399	0.3846	-0.2942

Figure A.2- Continued

112	0.4606	0.4027	-0.2989
113	0.4814	0.4209	-0.3000
114	0.5022	0.4391	-0.3053
115	0.5231	0.4574	-0.3079
116	0.5441	0.4757	-0.3077
117	0.5650	0.4939	-0.2921
118	0.5858	0.5122	-0.2933
119	0.6065	0.5303	-0.3833
120	0.6271	0.5483	-0.0654
121	0.6475	0.5724	0.2949
122	0.6677	0.5969	0.3781
123	0.6876	0.6211	0.3958
124	0.7073	0.6449	0.3910
125	0.7266	0.6684	0.3758
126	0.7455	0.6913	0.3521
127	0.7640	0.7138	0.3283
128	0.7821	0.7357	0.3056
129	0.7997	0.7570	0.2892
130	0.8168	0.7777	0.2640
131	0.8332	0.7977	0.2435
132	0.8491	0.8170	0.2301
133	0.8644	0.8355	0.2154
134	0.8790	0.8532	0.1976
135	0.8929	0.8700	0.1802
136	0.9060	0.8860	0.1773
137	0.9184	0.9011	0.1615
138	0.9300	0.9151	0.1423
139	0.9408	0.9282	0.1398
140	0.9507	0.9403	0.1259
141	0.9598	0.9513	0.1090
142	0.9680	0.9612	0.0929
143	0.9753	0.9700	0.0783
144	0.9817	0.9777	0.0614
145	0.9871	0.9843	0.0620
146	0.9916	0.9898	0.1164
147	0.9951	0.9941	-0.0799
148	0.9978	0.9973	0.1531
149	0.9995	0.9994	0.1520

Figure A.2- Continued



\*\*\*\*\*  
 \* OUTPUT FOR CASE #1 OF 6 \*  
 \*\*\*\*\*

M2 = 0.6000  
 Q2 = 0.0000  
 CPCRIT = -0.4346

<<<< LOCATIONS OF MIN., MAX., AND CRITICAL PTS. >>>>  
 (\* DENOTES POINT ON LOWER SURFACE)

PERTURBATION SOLN:

MINIMUM AT X = 0.4112\* (POINT # 45)  
 MAXIMUM AT X = 0.0000\* (POINT # 75)  
 2 CRITICAL POINT(S):  
 1ST AT X = 0.5024\* (AFTER POINT # 41)  
 2ND AT X = 0.4961 (AFTER POINT #120)

COMPARISON SOLN:

MINIMUM AT X = 0.3790\* (POINT # 42)  
 MAXIMUM AT X = 0.0000 (POINT # 75)  
 2 CRITICAL POINT(S):  
 1ST AT X = 0.5035\* (AFTER POINT # 35)  
 2ND AT X = 0.5035 (AFTER POINT #114)

POINT	XBASE	CPBASE	XCALB	CPCALB	XPERT	CPPERT	XCHEK	CPCHEK	CPPERT(INT)
1	0.9995	0.4947	0.9981	0.4492	0.9996	0.4189	0.9995	0.4865	0.4189
2	0.9978	0.4943	0.9946	0.4490	0.9982	0.4191	0.9978	0.4865	0.4293
3	0.9951	0.4144	0.9901	0.3687	0.9960	0.4721	0.9951	0.4060	0.4556
4	0.9916	0.4011	0.9847	0.3567	0.9931	0.4186	0.9916	0.3927	0.3946
5	0.9871	0.3709	0.9784	0.3231	0.9895	0.3613	0.9871	0.3623	0.3615
6	0.9817	0.3354	0.9711	0.2884	0.9850	0.3618	0.9817	0.3266	0.3433
7	0.9753	0.3037	0.9630	0.2570	0.9798	0.3332	0.9753	0.2947	0.3087
8	0.9680	0.2744	0.9539	0.2283	0.9739	0.3010	0.9680	0.2653	0.2750
9	0.9598	0.2478	0.9440	0.2012	0.9672	0.2713	0.9598	0.2386	0.2440
10	0.9507	0.2228	0.9333	0.1750	0.9598	0.2438	0.9507	0.2136	0.2155
11	0.9408	0.1986	0.9217	0.1528	0.9517	0.2180	0.9408	0.1893	0.1886
12	0.9300	0.1743	0.9093	0.1315	0.9428	0.1936	0.9300	0.1650	0.1638
13	0.9184	0.1535	0.8962	0.1078	0.9334	0.1705	0.9184	0.1443	0.1418
14	0.9060	0.1334	0.8824	0.0884	0.9232	0.1503	0.9060	0.1243	0.1187
15	0.8929	0.1106	0.8679	0.0690	0.9125	0.1313	0.8929	0.1016	0.0972
16	0.8790	0.0918	0.8527	0.0496	0.9012	0.1093	0.8790	0.0831	0.0777
17	0.8644	0.0728	0.8368	0.0302	0.8893	0.0919	0.8644	0.0646	0.0581
18	0.8491	0.0538	0.8204	0.0118	0.8768	0.0746	0.8491	0.0460	0.0382
19	0.8332	0.0346	0.8034	-0.0053	0.8638	0.0574	0.8332	0.0274	0.0195
20	0.8168	0.0163	0.7859	-0.0237	0.8504	0.0397	0.8168	0.0099	0.0029
21	0.7997	-0.0009	0.7679	-0.0417	0.8364	0.0227	0.7997	-0.0061	-0.0151
22	0.7821	-0.0193	0.7495	-0.0595	0.8221	0.0085	0.7821	-0.0232	-0.0338
23	0.7640	-0.0377	0.7306	-0.0747	0.8073	-0.0072	0.7640	-0.0399	-0.0496
24	0.7455	-0.0559	0.7113	-0.0926	0.7922	-0.0228	0.7455	-0.0559	-0.0645
25	0.7266	-0.0719	0.6918	-0.1098	0.7767	-0.0396	0.7266	-0.0690	-0.0818

Figure A.2- Continued

26	0.7073	-0.0908	0.6719	-0.1252	0.7609	-0.0520	0.7073	-0.0840	-0.0961
27	0.6876	-0.1094	0.6518	-0.1418	0.7449	-0.0650	0.6876	-0.0976	-0.1099
28	0.6677	-0.1267	0.6315	-0.1564	0.7286	-0.0803	0.6677	-0.1084	-0.1225
29	0.6475	-0.1460	0.6110	-0.1706	0.7121	-0.0927	0.6475	-0.1189	-0.1332
30	0.6271	-0.1643	0.5903	-0.1819	0.6955	-0.1043	0.6271	-0.1253	-0.1385
31	0.6065	-0.1838	0.5695	-0.1905	0.6787	-0.1163	0.6065	-0.1288	-0.1391
32	0.5858	-0.2024	0.5487	-0.1968	0.6617	-0.1259	0.5858	-0.1270	-0.1356
33	0.5650	-0.2213	0.5279	-0.1989	0.6447	-0.1347	0.5650	-0.1230	-0.1284
34	0.5441	-0.2417	0.5071	-0.2033	0.6277	-0.1385	0.5441	-0.1311	-0.1347
35	0.5231	-0.2608	0.4863	-0.2357	0.6106	-0.1394	0.5231	-0.1398	-0.2159
36	0.5022	-0.2786	0.4656	-0.3675	0.5935	-0.1381	0.5022	-0.4503	-0.4368
37	0.4814	-0.2955	0.4450	-0.5861	0.5765	-0.1327	0.4814	-0.7298	-0.6896
38	0.4606	-0.3119	0.4245	-0.7468	0.5595	-0.1263	0.4606	-0.8396	-0.8651
39	0.4399	-0.3264	0.4043	-0.8046	0.5426	-0.1355	0.4399	-0.8686	-0.8992
40	0.4194	-0.3504	0.3842	-0.8226	0.5259	-0.1909	0.4194	-0.8829	-0.9086
41	0.3991	-0.4033	0.3645	-0.8300	0.5093	-0.3436	0.3991	-0.8897	-0.9096
42	0.3790	-0.4976	0.3449	-0.8304	0.4866	-0.6450	0.3790	-0.8915	-0.9082
43	0.3591	-0.6009	0.3257	-0.8280	0.4611	-0.8644	0.3591	-0.8910	-0.9056
44	0.3395	-0.6682	0.3069	-0.8232	0.4359	-0.9059	0.3395	-0.8864	-0.9016
45	0.3203	-0.7003	0.2884	-0.8152	0.4112	-0.9100	0.3203	-0.8803	-0.8942
46	0.3013	-0.7119	0.2703	-0.8057	0.3869	-0.9091	0.3013	-0.8727	-0.8860
47	0.2828	-0.7121	0.2526	-0.7996	0.3630	-0.9064	0.2828	-0.8624	-0.8765
48	0.2646	-0.7073	0.2354	-0.7810	0.3397	-0.9017	0.2646	-0.8510	-0.8648
49	0.2458	-0.6991	0.2186	-0.7654	0.3169	-0.8929	0.2458	-0.8382	-0.8514
50	0.2295	-0.6872	0.2023	-0.7485	0.2947	-0.8830	0.2295	-0.8230	-0.8370
51	0.2127	-0.6726	0.1866	-0.7282	0.2731	-0.8711	0.2127	-0.8060	-0.8197
52	0.1964	-0.6561	0.1714	-0.7074	0.2522	-0.8556	0.1964	-0.7878	-0.7995
53	0.1806	-0.6355	0.1567	-0.6859	0.2319	-0.8395	0.1806	-0.7662	-0.7814
54	0.1653	-0.6140	0.1427	-0.6590	0.2123	-0.8192	0.1653	-0.7441	-0.7601
55	0.1507	-0.5916	0.1292	-0.6298	0.1934	-0.7958	0.1507	-0.7212	-0.7348
56	0.1366	-0.5629	0.1164	-0.5989	0.1753	-0.7755	0.1366	-0.6927	-0.7068
57	0.1231	-0.5308	0.1041	-0.5656	0.1580	-0.7487	0.1231	-0.6621	-0.6767
58	0.1102	-0.4968	0.0926	-0.5280	0.1414	-0.7174	0.1102	-0.6293	-0.6450
59	0.0979	-0.4605	0.0817	-0.4884	0.1257	-0.6832	0.0979	-0.5939	-0.6090
60	0.0863	-0.4184	0.0715	-0.4412	0.1108	-0.6470	0.0863	-0.5535	-0.5689
61	0.0754	-0.3723	0.0619	-0.3945	0.0968	-0.6058	0.0754	-0.5104	-0.5224
62	0.0652	-0.3197	0.0531	-0.3362	0.0837	-0.5597	0.0652	-0.4596	-0.4704
63	0.0557	-0.2688	0.0449	-0.2716	0.0715	-0.5047	0.0557	-0.4078	-0.4116
64	0.0469	-0.2053	0.0375	-0.2113	0.0602	-0.4431	0.0469	-0.3418	-0.3492
65	0.0388	-0.1317	0.0308	-0.1366	0.0498	-0.3706	0.0388	-0.2672	-0.2838
66	0.0314	-0.0567	0.0248	-0.0514	0.0403	-0.3004	0.0314	-0.1938	-0.2038
67	0.0248	0.0373	0.0196	0.0521	0.0319	-0.2104	0.0248	-0.1013	-0.1055
68	0.0190	0.1485	0.0150	0.1636	0.0244	-0.0988	0.0190	0.0090	0.0069
69	0.0139	0.2859	0.0111	0.2966	0.0178	0.0292	0.0139	0.1478	0.1436
70	0.0096	0.4435	0.0078	0.4490	0.0123	0.1896	0.0096	0.3090	0.3170
71	0.0061	0.6300	0.0051	0.6219	0.0078	0.3998	0.0061	0.5052	0.5266
72	0.0034	0.8314	0.0030	0.8106	0.0043	0.6549	0.0034	0.7246	0.7570
73	0.0015	1.0193	0.0014	0.9949	0.0019	0.9137	0.0015	0.9415	0.9699
74	0.0004	1.1454	0.0004	1.1296	0.0005	1.1052	0.0004	1.1077	1.1206
75	0.0000	1.1617	0.0000	1.1692	0.0000	1.1740	0.0000	1.1704	1.1739
76	0.0004	1.0535	0.0004	1.0967	0.0003	1.1402	0.0004	1.1077	1.1294
77	0.0015	0.8493	0.0014	0.9332	0.0012	1.0158	0.0015	0.9415	0.9745
78	0.0034	0.6062	0.0030	0.7278	0.0027	0.8175	0.0034	0.7246	0.7479
79	0.0061	0.3715	0.0051	0.5262	0.0048	0.6070	0.0061	0.5052	0.5197
80	0.0096	0.1683	0.0078	0.3464	0.0076	0.4180	0.0096	0.3090	0.3201
81	0.0139	0.0054	0.0111	0.1907	0.0110	0.2519	0.0139	0.1478	0.1549
82	0.0190	-0.1336	0.0150	0.0568	0.0150	0.1181	0.0190	0.0090	0.0185
83	0.0248	-0.2430	0.0196	-0.0540	0.0196	0.0025	0.0248	-0.1013	-0.0993
84	0.0314	-0.3358	0.0248	-0.1570	0.0248	-0.0998	0.0314	-0.1938	-0.1939
85	0.0388	-0.4104	0.0308	-0.2415	0.0306	-0.1859	0.0388	-0.2672	-0.2684

Figure A.2-Continued

86	0.0469	-0.4834	0.0375	-0.3160	0.0371	-0.2524	0.0469	-0.3418	-0.3445
87	0.0557	-0.5434	0.0449	-0.3774	0.0440	-0.3175	0.0557	-0.4078	-0.4134
88	0.0652	-0.5903	0.0531	-0.4440	0.0515	-0.3880	0.0652	-0.4596	-0.4659
89	0.0754	-0.6388	0.0619	-0.5002	0.0596	-0.4376	0.0754	-0.5104	-0.5160
90	0.0863	-0.6794	0.0715	-0.5425	0.0682	-0.4812	0.0863	-0.5535	-0.5615
91	0.0979	-0.7167	0.0817	-0.5880	0.0774	-0.5256	0.0979	-0.5939	-0.6014
92	0.1102	-0.7492	0.0926	-0.6268	0.0871	-0.5645	0.1102	-0.6293	-0.6355
93	0.1231	-0.7793	0.1041	-0.6626	0.0973	-0.5996	0.1231	-0.6621	-0.6681
94	0.1366	-0.8075	0.1164	-0.6940	0.1080	-0.6299	0.1366	-0.6927	-0.7005
95	0.1507	-0.8341	0.1292	-0.7233	0.1191	-0.6581	0.1507	-0.7212	-0.7271
96	0.1653	-0.8557	0.1427	-0.7512	0.1307	-0.6874	0.1653	-0.7441	-0.7504
97	0.1806	-0.8763	0.1567	-0.7775	0.1428	-0.7145	0.1806	-0.7662	-0.7737
98	0.1964	-0.8968	0.1714	-0.7989	0.1553	-0.7345	0.1964	-0.7878	-0.7933
99	0.2127	-0.9149	0.1866	-0.8197	0.1682	-0.7548	0.2127	-0.8060	-0.8115
100	0.2295	-0.9321	0.2023	-0.8400	0.1815	-0.7750	0.2295	-0.8230	-0.8279
101	0.2468	-0.9479	0.2186	-0.8572	0.1951	-0.7918	0.2468	-0.8382	-0.8424
102	0.2646	-0.9618	0.2354	-0.8733	0.2092	-0.8078	0.2646	-0.8510	-0.8559
103	0.2828	-0.9745	0.2526	-0.8880	0.2235	-0.8227	0.2828	-0.8624	-0.8673
104	0.3013	-0.9868	0.2703	-0.9010	0.2382	-0.8354	0.3013	-0.8727	-0.8769
105	0.3203	-0.9972	0.2884	-0.9128	0.2532	-0.8475	0.3203	-0.8803	-0.8844
106	0.3395	-1.0067	0.3069	-0.9239	0.2684	-0.8587	0.3395	-0.8864	-0.8898
107	0.3591	-1.0160	0.3257	-0.9326	0.2839	-0.8680	0.3591	-0.8910	-0.8920
108	0.3790	-1.0227	0.3449	-0.9401	0.2996	-0.8762	0.3790	-0.8915	-0.8918
109	0.3991	-1.0296	0.3645	-0.9467	0.3155	-0.8827	0.3991	-0.8897	-0.8858
110	0.4194	-1.0349	0.3842	-0.9500	0.3316	-0.8885	0.4194	-0.8829	-0.8746
111	0.4399	-1.0383	0.4043	-0.9523	0.3478	-0.8912	0.4399	-0.8686	-0.8550
112	0.4606	-1.0418	0.4245	-0.9516	0.3641	-0.8924	0.4606	-0.8396	-0.8082
113	0.4814	-1.0417	0.4450	-0.9463	0.3805	-0.8917	0.4814	-0.7298	-0.6388
114	0.5022	-1.0394	0.4656	-0.9371	0.3970	-0.8868	0.5022	-0.4503	-0.3785
115	0.5231	-1.0332	0.4863	-0.9177	0.4136	-0.8792	0.5231	-0.1998	-0.1929
116	0.5441	-1.0200	0.5071	-0.8858	0.4301	-0.8661	0.5441	-0.1311	-0.1426
117	0.5650	-0.9936	0.5279	-0.7801	0.4466	-0.8475	0.5650	-0.1230	-0.1300
118	0.5858	-0.9478	0.5487	-0.4509	0.4631	-0.8012	0.5858	-0.1270	-0.1325
119	0.6065	-0.8572	0.5695	-0.1357	0.4795	-0.6655	0.6065	-0.1288	-0.1356
120	0.6271	-0.4706	0.5903	-0.0554	0.4957	-0.4378	0.6271	-0.1253	-0.1342
121	0.6475	-0.0472	0.6110	-0.0489	0.5224	-0.1946	0.6475	-0.1189	-0.1288
122	0.6677	0.0602	0.6315	-0.0545	0.5498	-0.1289	0.6677	-0.1084	-0.1200
123	0.6876	0.0671	0.6518	-0.0593	0.5768	-0.1308	0.6876	-0.0976	-0.1087
124	0.7073	0.0596	0.6719	-0.0587	0.6034	-0.1359	0.7073	-0.0840	-0.0961
125	0.7266	0.0539	0.6918	-0.0560	0.6296	-0.1340	0.7266	-0.0690	-0.0839
126	0.7455	0.0495	0.7113	-0.0490	0.6552	-0.1265	0.7455	-0.0559	-0.0679
127	0.7640	0.0508	0.7306	-0.0390	0.6803	-0.1133	0.7640	-0.0399	-0.0512
128	0.7821	0.0551	0.7495	-0.0300	0.7048	-0.0977	0.7821	-0.0232	-0.0349
129	0.7997	0.0619	0.7679	-0.0174	0.7286	-0.0827	0.7997	-0.0061	-0.0188
130	0.8168	0.0695	0.7859	-0.0034	0.7517	-0.0625	0.8168	0.0099	-0.0017
131	0.8332	0.0797	0.8034	0.0116	0.7740	-0.0420	0.8332	0.0274	0.0175
132	0.8491	0.0921	0.8204	0.0260	0.7956	-0.0229	0.8491	0.0460	0.0370
133	0.8644	0.1054	0.8368	0.0421	0.8163	-0.0023	0.8644	0.0646	0.0552
134	0.8790	0.1195	0.8527	0.0596	0.8360	0.0207	0.8790	0.0831	0.0742
135	0.8929	0.1342	0.8679	0.0774	0.8548	0.0441	0.8929	0.1016	0.0958
136	0.9060	0.1534	0.8824	0.0955	0.8727	0.0648	0.9060	0.1243	0.1190
137	0.9184	0.1706	0.8962	0.1138	0.8895	0.0898	0.9184	0.1443	0.1391
138	0.9300	0.1887	0.9093	0.1365	0.9052	0.1176	0.9300	0.1650	0.1631
139	0.9408	0.2107	0.9217	0.1570	0.9198	0.1413	0.9408	0.1893	0.1895
140	0.9507	0.2330	0.9333	0.1785	0.9333	0.1700	0.9507	0.2136	0.2172
141	0.9598	0.2563	0.9440	0.2041	0.9456	0.2018	0.9598	0.2386	0.2462
142	0.9680	0.2813	0.9539	0.2306	0.9567	0.2348	0.9680	0.2653	0.2769
143	0.9753	0.3093	0.9630	0.2589	0.9665	0.2701	0.9753	0.2947	0.3098
144	0.9817	0.3398	0.9711	0.2899	0.9751	0.3091	0.9817	0.3266	0.3393
145	0.9871	0.3743	0.9784	0.3243	0.9825	0.3433	0.9871	0.3623	0.3449

Figure A.2- Continued

146	0.9916	0.4036	0.9847	0.3576	0.9886	0.3454	0.9916	0.3927	0.4144
147	0.9951	0.4162	0.9901	0.3694	0.9934	0.4561	0.9951	0.4060	0.4379
148	0.9978	0.4952	0.9946	0.4493	0.9970	0.4186	0.9978	0.4865	0.4186
149	0.9995	0.4947	0.9981	0.4492	0.9993	0.4187	0.9995	0.4865	0.4187

\*\*\*\*\*  
 \* OUTPUT FOR CASE #2 OF 6 \*  
 \*\*\*\*\*

M2 = 0.8000

Q2 = 0.1000

CPCRIT = -0.4346

<<<< LOCATIONS OF MIN., MAX., AND CRITICAL PTS. >>>>  
 (\* DENOTES POINT ON LOWER SURFACE)

PERTURBATION SOLN:

MINIMUM AT X = 0.3834 (POINT #112)  
 MAXIMUM AT X = 0.0000\* (POINT # 75)  
 2 CRITICAL POINT(S):  
 1ST AT X = 0.4805\* (AFTER POINT # 41)  
 2ND AT X = 0.5229 (AFTER POINT #120)

COMPARISON SOLN:

MINIMUM AT X = 0.4015 (POINT #110)  
 MAXIMUM AT X = 0.0000 (POINT # 75)  
 2 CRITICAL POINT(S):  
 1ST AT X = 0.4816\* (AFTER POINT # 36)  
 2ND AT X = 0.5258 (AFTER POINT #116)

POINT	XBASE	CPBASE	XCALB	CPCALB	XPERT	CPPERT	XCHEK	CPCHEK	CPPERT(INT)
1	0.9995	0.4947	0.9981	0.4492	0.9996	0.4341	0.9989	0.4589	0.4341
2	0.9978	0.4943	0.9946	0.4490	0.9981	0.4341	0.9963	0.4588	0.4557
3	0.9951	0.4144	0.9901	0.3687	0.9958	0.4606	0.9927	0.3781	0.4138
4	0.9916	0.4011	0.9847	0.3567	0.9928	0.4151	0.9882	0.3713	0.3621
5	0.9871	0.3709	0.9784	0.3231	0.9890	0.3632	0.9828	0.3389	0.3483
6	0.9817	0.3354	0.9711	0.2884	0.9843	0.3565	0.9765	0.3055	0.3148
7	0.9753	0.3037	0.9630	0.2570	0.9789	0.3273	0.9692	0.2738	0.2811
8	0.9680	0.2744	0.9539	0.2283	0.9727	0.2957	0.9610	0.2453	0.2503
9	0.9598	0.2478	0.9440	0.2012	0.9657	0.2666	0.9520	0.2188	0.2216
10	0.9507	0.2228	0.9333	0.1750	0.9580	0.2396	0.9420	0.1934	0.1944
11	0.9408	0.1986	0.9217	0.1528	0.9495	0.2142	0.9313	0.1683	0.1692
12	0.9300	0.1743	0.9093	0.1315	0.9403	0.1897	0.9197	0.1470	0.1468
13	0.9 84	0.1535	0.8962	0.1078	0.9304	0.1671	0.9073	0.1264	0.1249
14	0.9060	0.1334	0.8824	0.0884	0.9198	0.1469	0.8942	0.1033	0.1022
15	0.8929	0.1106	0.8679	0.0690	0.9086	0.1271	0.8804	0.0843	0.0828
16	0.8790	0.0918	0.8527	0.0496	0.8967	0.1058	0.8658	0.0653	0.0634
17	0.8644	0.0728	0.8368	0.0302	0.8843	0.0881	0.8506	0.0463	0.0438
18	0.8491	0.0538	0.8204	0.0118	0.8713	0.0705	0.8347	0.0273	0.0246
19	0.8332	0.0346	0.8034	-0.0053	0.8577	0.0528	0.8183	0.0093	0.0071
20	0.8168	0.0163	0.7859	-0.0237	0.8436	0.0350	0.8013	-0.0074	-0.0105
21	0.7997	-0.0009	0.7679	-0.0417	0.8291	0.0180	0.7837	-0.0253	-0.0286
22	0.7821	-0.0193	0.7495	-0.0595	0.8141	0.0029	0.7657	-0.0428	-0.0469
23	0.7640	-0.0377	0.7306	-0.0747	0.7986	-0.0133	0.7472	-0.0598	-0.0623
24	0.7455	-0.0559	0.7113	-0.0926	0.7829	-0.0295	0.7283	-0.0743	-0.0786
25	0.7266	-0.0719	0.6918	-0.1098	0.7667	-0.0461	0.7090	-0.0911	-0.0955

Figure A.2- Continued

26	0.7073	-0.0908	0.6719	-0.1252	0.7502	-0.0598	0.6895	-0.1068	-0.1106
27	0.6876	-0.1094	0.6518	-0.1418	0.7335	-0.0738	0.6696	-0.1204	-0.1257
28	0.6677	-0.1267	0.6315	-0.1564	0.7164	-0.0896	0.6494	-0.1346	-0.1393
29	0.6475	-0.1460	0.6110	-0.1706	0.6992	-0.1034	0.6291	-0.1459	-0.1518
30	0.6271	-0.1643	0.5903	-0.1819	0.6818	-0.1163	0.6085	-0.1556	-0.1598
31	0.6065	-0.1838	0.5695	-0.1905	0.6642	-0.1298	0.5878	-0.1610	-0.1644
32	0.5858	-0.2024	0.5487	-0.1968	0.6466	-0.1412	0.5670	-0.1624	-0.1658
33	0.5650	-0.2213	0.5279	-0.1989	0.6288	-0.1520	0.5462	-0.1625	-0.1641
34	0.5441	-0.2417	0.5071	-0.2033	0.6109	-0.1592	0.5253	-0.1683	-0.1718
35	0.5231	-0.2608	0.4863	-0.2357	0.5931	-0.1637	0.5045	-0.2166	-0.2237
36	0.5022	-0.2786	0.4656	-0.3675	0.5752	-0.1662	0.4836	-0.4092	-0.3977
37	0.4814	-0.2955	0.4450	-0.5861	0.5574	-0.1652	0.4629	-0.6656	-0.6327
38	0.4606	-0.3119	0.4245	-0.7468	0.5397	-0.1634	0.4423	-0.7983	-0.7985
39	0.4399	-0.3264	0.4043	-0.8046	0.5221	-0.1737	0.4218	-0.8384	-0.8483
40	0.4194	-0.3504	0.3842	-0.8226	0.5046	-0.2228	0.4015	-0.8530	-0.8645
41	0.3991	-0.4033	0.3645	-0.8300	0.4872	-0.3556	0.3815	-0.8587	-0.8689
42	0.3790	-0.4976	0.3449	-0.8304	0.4650	-0.6156	0.3616	-0.8609	-0.8689
43	0.3591	-0.6009	0.3257	-0.8280	0.4407	-0.8117	0.3421	-0.8583	-0.8665
44	0.3395	-0.6682	0.3069	-0.8232	0.4166	-0.8583	0.3229	-0.8537	-0.8621
45	0.3203	-0.7003	0.2884	-0.8152	0.3930	-0.8681	0.3040	-0.8472	-0.8546
46	0.3013	-0.7119	0.2703	-0.8057	0.3698	-0.8697	0.2855	-0.8378	-0.8457
47	0.2828	-0.7121	0.2526	-0.7946	0.3470	-0.8675	0.2673	-0.8272	-0.8352
48	0.2646	-0.7073	0.2354	-0.7810	0.3247	-0.8628	0.2496	-0.8151	-0.8225
49	0.2468	-0.6991	0.2186	-0.7654	0.3029	-0.8542	0.2324	-0.8006	-0.8081
50	0.2295	-0.6872	0.2023	-0.7485	0.2817	-0.8439	0.2156	-0.7844	-0.7920
51	0.2127	-0.6726	0.1866	-0.7282	0.2610	-0.8314	0.1993	-0.7669	-0.7735
52	0.1964	-0.6561	0.1714	-0.7074	0.2410	-0.8157	0.1835	-0.7460	-0.7533
53	0.1806	-0.6355	0.1567	-0.6859	0.2216	-0.7987	0.1683	-0.7247	-0.7339
54	0.1653	-0.6140	0.1427	-0.6590	0.2029	-0.7782	0.1536	-0.7027	-0.7095
55	0.1507	-0.5916	0.1292	-0.6298	0.1849	-0.7550	0.1396	-0.6752	-0.6821
56	0.1366	-0.5629	0.1164	-0.5989	0.1676	-0.7330	0.1261	-0.6455	-0.6523
57	0.1231	-0.5308	0.1041	-0.5656	0.1510	-0.7051	0.1132	-0.6140	-0.6204
58	0.1102	-0.4968	0.0926	-0.5280	0.1352	-0.6733	0.1010	-0.5800	-0.5857
59	0.0979	-0.4605	0.0817	-0.4884	0.1202	-0.6386	0.0894	-0.5414	-0.5472
60	0.0863	-0.4184	0.0715	-0.4412	0.1059	-0.6013	0.0785	-0.5003	-0.5047
61	0.0754	-0.3723	0.0619	-0.3945	0.0926	-0.5591	0.0683	-0.4514	-0.4574
62	0.0652	-0.3197	0.0531	-0.3362	0.0800	-0.5117	0.0588	-0.4023	-0.4026
63	0.0557	-0.2688	0.0449	-0.2716	0.0683	-0.4575	0.0500	-0.3402	-0.3403
64	0.0469	-0.2053	0.0375	-0.2113	0.0575	-0.3955	0.0418	-0.2708	-0.2775
65	0.0388	-0.1317	0.0308	-0.1366	0.0476	-0.3229	0.0345	-0.2044	-0.2056
66	0.0314	-0.0567	0.0248	-0.0514	0.0386	-0.2517	0.0278	-0.1217	-0.1194
67	0.0248	0.0373	0.0196	0.0521	0.0305	-0.1608	0.0219	-0.0253	-0.0199
68	0.0190	0.1485	0.0150	0.1636	0.0233	-0.0493	0.0167	0.0935	0.0917
69	0.0139	0.2859	0.0111	0.2966	0.0171	0.0806	0.0122	0.2265	0.2268
70	0.0096	0.4435	0.0078	0.4490	0.0118	0.2404	0.0085	0.3875	0.3973
71	0.0061	0.6300	0.0051	0.6219	0.0075	0.4458	0.0054	0.5727	0.5947
72	0.0034	0.8314	0.0030	0.8106	0.0041	0.6902	0.0031	0.7749	0.7999
73	0.0015	1.0193	0.0014	0.9949	0.0018	0.9348	0.0014	0.9718	0.9890
74	0.0004	1.1454	0.0004	1.1296	0.0004	1.1133	0.0004	1.1201	1.1243
75	0.0090	1.1617	0.0000	1.1692	0.0000	1.1716	0.0000	1.1701	1.1715
76	0.0004	1.0535	0.0004	1.0967	0.0003	1.1228	0.0004	1.1034	1.1138
77	0.0015	0.8493	0.0014	0.9332	0.0012	0.9825	0.0014	0.9406	0.9603
78	0.0034	0.6062	0.0030	0.7278	0.0028	0.7753	0.0031	0.7335	0.7474
79	0.0061	0.3715	0.0051	0.5262	0.0051	0.5599	0.0054	0.5249	0.5349
80	0.0096	0.1683	0.0078	0.3464	0.0080	0.3680	0.0085	0.3363	0.3447
81	0.0139	0.0054	0.0111	0.1907	0.0116	0.2026	0.0122	0.1738	0.1820
82	0.0190	-0.1336	0.0150	0.0568	0.0158	0.0677	0.0167	0.0407	0.0469
83	0.0248	-0.2430	0.0196	-0.0540	0.0207	-0.0466	0.0219	-0.0780	-0.0689
84	0.0314	-0.3358	0.0248	-0.1570	0.0262	-0.1470	0.0278	-0.1738	-0.1694
85	0.0388	-0.4104	0.0308	-0.2415	0.0323	-0.2308	0.0344	-0.2562	-0.2526

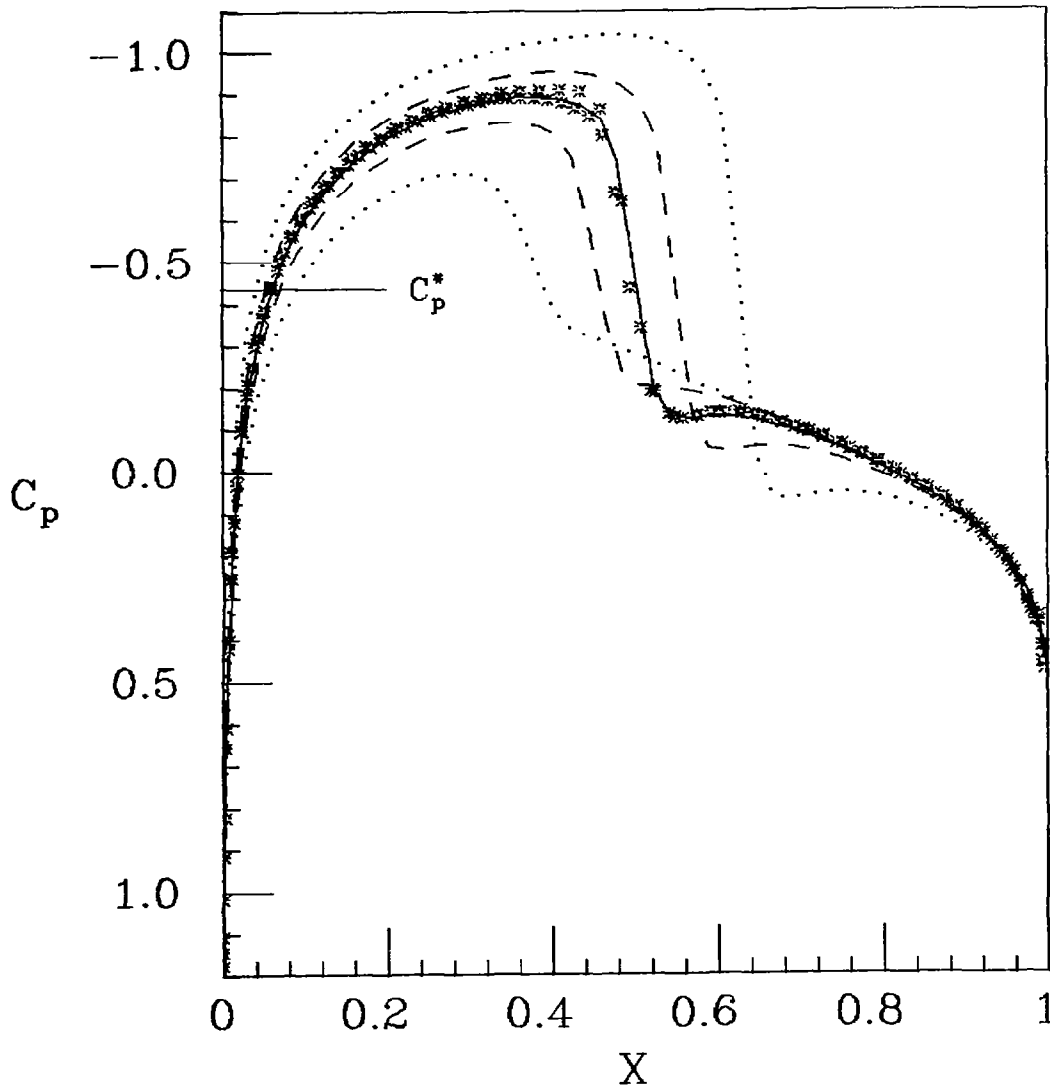
Figure A.2- Continued

86	0.0469	-0.4834	0.0375	-0.3160	0.0390	-0.2986	0.0418	-0.3229	-0.3231
87	0.0557	-0.5434	0.0449	-0.3774	0.0463	-0.3627	0.0500	-0.3936	-0.3928
88	0.0652	-0.5903	0.0531	-0.4440	0.0543	-0.4285	0.0588	-0.4554	-0.4545
89	0.0754	-0.6388	0.0619	-0.5002	0.0628	-0.4779	0.0683	-0.5019	-0.5040
90	0.0863	-0.6794	0.0715	-0.5425	0.0719	-0.5208	0.0785	-0.5492	-0.5505
91	0.0979	-0.7167	0.0817	-0.5880	0.0815	-0.5638	0.0894	-0.5899	-0.5930
92	0.1102	-0.7492	0.0926	-0.6268	0.0917	-0.6015	0.1010	-0.6277	-0.6309
93	0.1231	-0.7793	0.1041	-0.6626	0.1024	-0.6355	0.1132	-0.6607	-0.6642
94	0.1366	-0.8075	0.1164	-0.6940	0.1137	-0.6654	0.1261	-0.6913	-0.6949
95	0.1507	-0.8341	0.1292	-0.7233	0.1254	-0.6933	0.1396	-0.7202	-0.7250
96	0.1633	-0.8557	0.1427	-0.7512	0.1376	-0.7210	0.1536	-0.7473	-0.7519
97	0.1806	-0.8763	0.1567	-0.7775	0.1503	-0.7469	0.1683	-0.7693	-0.7740
98	0.1 64	-0.8968	0.1714	-0.7989	0.1635	-0.7669	0.1835	-0.7906	-0.7959
99	0.2127	-0.9149	0.1866	-0.8197	0.1771	-0.7868	0.1993	-0.8114	-0.8159
100	0.2295	-0.9321	0.2023	-0.8400	0.1911	-0.8064	0.2156	-0.8291	-0.8337
101	0.2468	-0.9479	0.2186	-0.8572	0.2055	-0.8230	0.2324	-0.8457	-0.8502
102	0.2646	-0.9618	0.2354	-0.8733	0.2202	-0.8386	0.2496	-0.8607	-0.8647
103	0.2828	-0.9745	0.2526	-0.8880	0.2354	-0.8530	0.2673	-0.8734	-0.8779
104	0.3013	-0.9868	0.2703	-0.9010	0.2508	-0.8657	0.2855	-0.8849	-0.8899
105	0.3203	-0.9972	0.2884	-0.9128	0.2666	-0.8774	0.3040	-0.8958	-0.9000
106	0.3395	-1.0067	0.3069	-0.9239	0.2826	-0.8883	0.3229	-0.9043	-0.9084
107	0.3591	-1.0160	0.3257	-0.9326	0.2989	-0.8976	0.3421	-0.9114	-0.9154
108	0.3790	-1.0227	0.3449	-0.9401	0.3155	-0.9055	0.3616	-0.9173	-0.9198
109	0.3991	-1.0296	0.3645	-0.9467	0.3322	-0.9121	0.3815	-0.9194	-0.9221
110	0.4194	-1.0349	0.3842	-0.9500	0.3491	-0.9178	0.4015	-0.9200	-0.9215
111	0.4399	-1.0383	0.4043	-0.9523	0.3662	-0.9206	0.4218	-0.9166	-0.9157
112	0.4606	-1.0418	0.4245	-0.9516	0.3834	-0.9223	0.4423	-0.9075	-0.9049
113	0.4814	-1.0417	0.4450	-0.9463	0.4007	-0.9217	0.4629	-0.8922	-0.8853
114	0.5022	-1.0394	0.4656	-0.9371	0.4181	-0.9173	0.4836	-0.8608	-0.8411
115	0.5231	-1.0332	0.4863	-0.9177	0.4355	-0.9100	0.5045	-0.7489	-0.7069
116	0.5441	-1.0200	0.5071	-0.8858	0.4529	-0.8969	0.5253	-0.4408	-0.4081
117	0.5650	-0.9936	0.5279	-0.7801	0.4703	-0.8767	0.5462	-0.1637	-0.1787
118	0.5858	-0.9478	0.5487	-0.4509	0.4876	-0.8305	0.5670	-0.0936	-0.1091
119	0.6065	-0.8572	0.5695	-0.1357	0.5049	-0.7038	0.5878	-0.0883	-0.0912
120	0.6271	-0.4706	0.5903	-0.0554	0.5220	-0.4444	0.6085	-0.0938	-0.0933
121	0.6475	-0.0472	0.6110	-0.0489	0.5474	-0.1651	0.6291	-0.0962	-0.0967
122	0.6677	0.0602	0.6315	-0.0545	0.5733	-0.0911	0.6494	-0.0951	-0.0963
123	0.6876	0.0671	0.6518	-0.0593	0.5990	-0.0912	0.6696	-0.0890	-0.0921
124	0.7073	0.0596	0.6719	-0.0587	0.6242	-0.0968	0.6895	-0.0816	-0.0840
125	0.7266	0.0539	0.6918	-0.0560	0.6490	-0.0964	0.7090	-0.0707	-0.0736
126	0.7455	0.0495	0.7113	-0.0490	0.6733	-0.0913	0.7283	-0.0577	-0.0624
127	0.7640	0.0508	0.7306	-0.0390	0.6970	-0.0805	0.7472	-0.0463	-0.0502
128	0.7821	0.0551	0.7495	-0.0300	0.7202	-0.0672	0.7657	-0.0316	-0.0353
129	0.7997	0.0619	0.7679	-0.0174	0.7428	-0.0537	0.7837	-0.0160	-0.0195
130	0.8168	0.0695	0.7859	-0.0034	0.7647	-0.0361	0.8013	0.0003	-0.0043
131	0.8332	0.0797	0.8034	0.0116	0.7859	-0.0176	0.8183	0.0157	0.0118
132	0.8491	0.0921	0.8204	0.0260	0.8063	0.0001	0.8347	0.0327	0.0293
133	0.8644	0.1054	0.8368	0.0421	0.8259	0.0193	0.8506	0.0508	0.0477
134	0.8790	0.1195	0.8527	0.0596	0.8446	0.0405	0.8658	0.0691	0.0662
135	0.8929	0.1342	0.8679	0.0774	0.8624	0.0621	0.8804	0.0875	0.0841
136	0.9060	0.1534	0.8824	0.0955	0.8793	0.0825	0.8942	0.1060	0.1045
137	0.9184	0.1706	0.8962	0.1138	0.8953	0.1060	0.9073	0.1286	0.1270
138	0.9300	0.1887	0.9093	0.1365	0.9101	0.1318	0.9197	0.1489	0.1480
139	0.9408	0.2107	0.9217	0.1570	0.9240	0.1552	0.9313	0.1699	0.1709
140	0.9507	0.2330	0.9333	0.1785	0.9368	0.1826	0.9420	0.1947	0.1963
141	0.9598	0.2563	0.9440	0.2041	0.9484	0.2127	0.9520	0.2199	0.2233
142	0.9680	0.2813	0.9539	0.2306	0.9589	0.2441	0.9610	0.2462	0.2517
143	0.9753	0.3093	0.9630	0.2589	0.9683	0.2780	0.9692	0.2746	0.2822
144	0.9817	0.3398	0.9711	0.2899	0.9764	0.3153	0.9765	0.3061	0.3154
145	0.9871	0.3743	0.9784	0.3243	0.9834	0.3495	0.9828	0.3394	0.3465
146	0.9916	0.4036	0.9847	0.3576	0.9892	0.3571	0.9882	0.3717	0.3558
147	0.9951	0.4162	0.9901	0.3694	0.9937	0.4481	0.9927	0.3784	0.4278
148	0.9978	0.4952	0.9946	0.4493	0.9971	0.4339	0.9963	0.4590	0.4375
149	0.9995	0.4947	0.9981	0.4492	0.9994	0.4339	0.9989	0.4589	0.4339

Figure A.2- Concluded

# Plot of $C_p$

Full  $\phi_{\text{Holst}}$  Nonl.  $\alpha$  Pert.  $\alpha_b = 0.500$   $\alpha_c = 0.200$

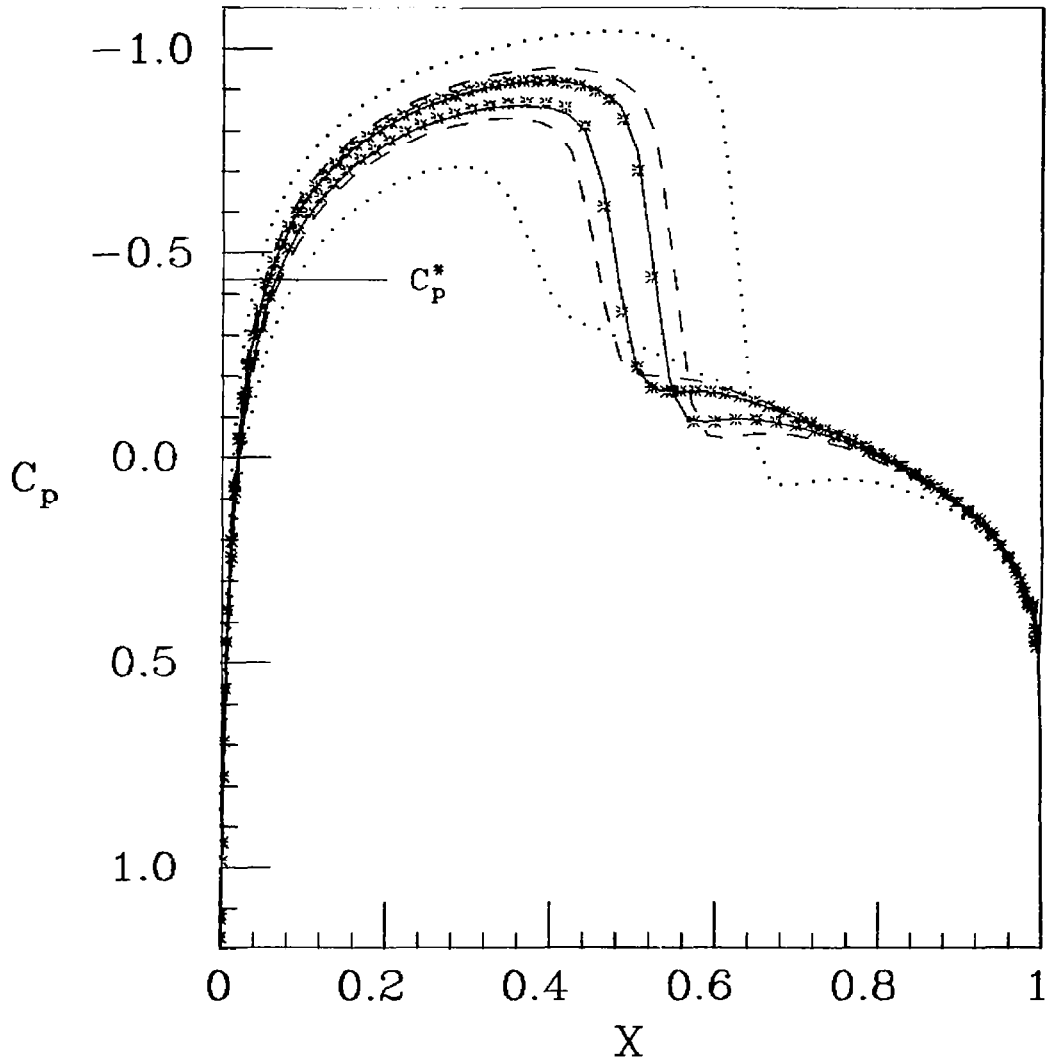


$M_\infty = 0.800$   $\alpha = 0.000$

Figure A.3- Plot output for sample case

# Plot of $C_p$

Full  $\phi_{\text{Holst}}$  Nonl.  $\alpha$  Pert.  $\alpha_b = 0.500$   $\alpha_c = 0.200$



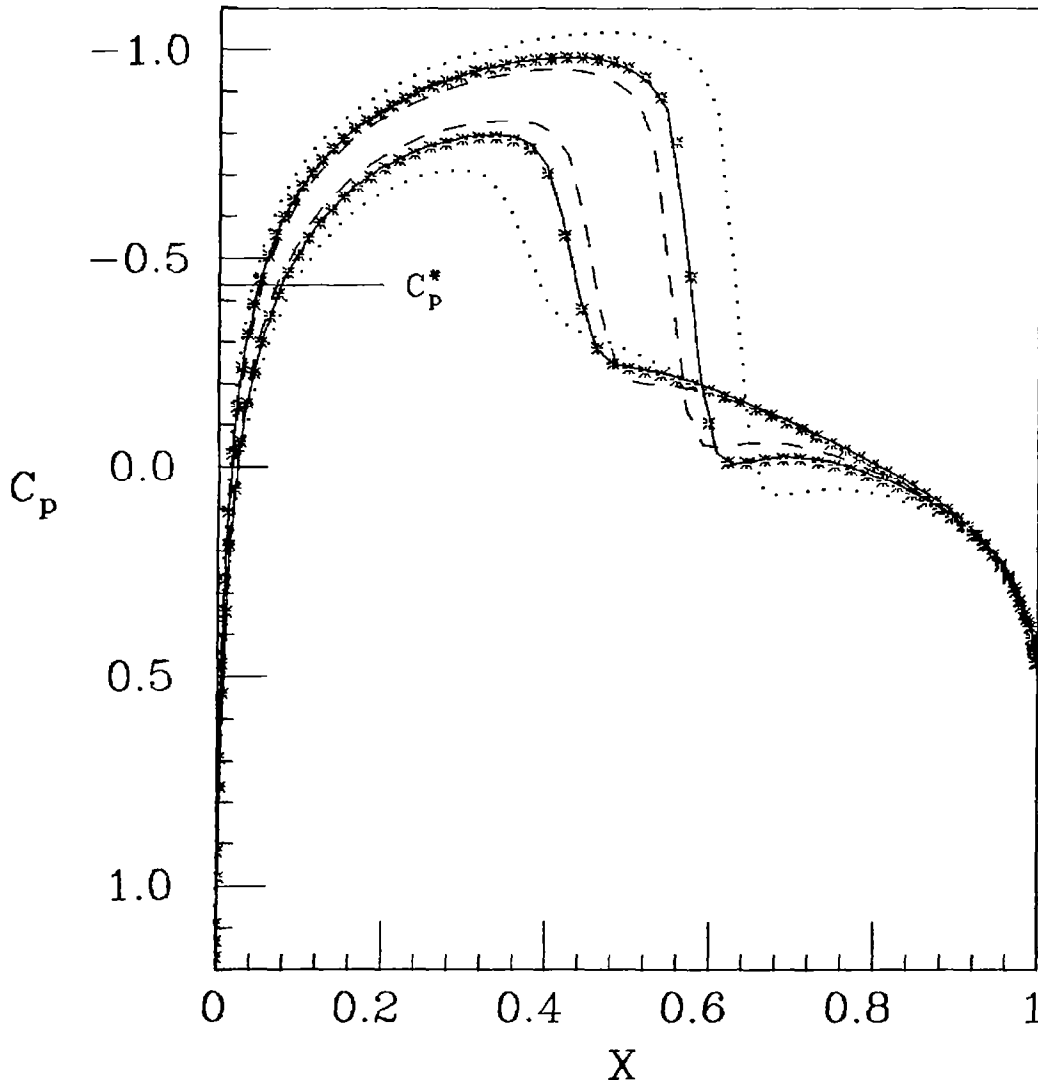
$M_\infty = 0.800$   $\alpha = 0.100$

Figure A.3- Continued



# Plot of $C_p$

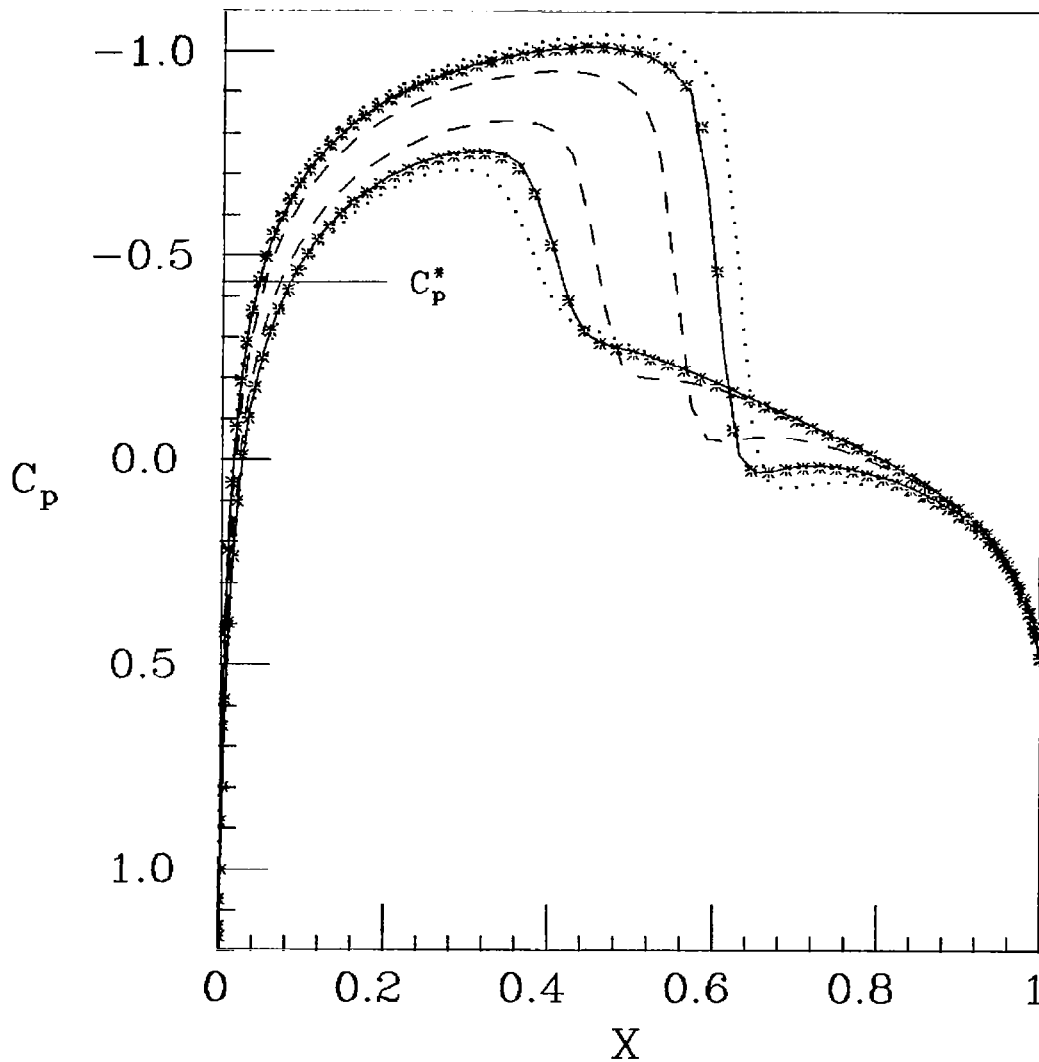
Full  $\phi_{Holst}$    Nonl.  $\alpha$  Pert.    $\alpha_b = 0.500$     $\alpha_c = 0.200$



$M_\infty = 0.800$     $\alpha = 0.300$

# Plot of $C_p$

Full  $\phi_{\text{Holst}}$  Nonl.  $\alpha$  Pert.  $\alpha_b = 0.500$   $\alpha_c = 0.200$

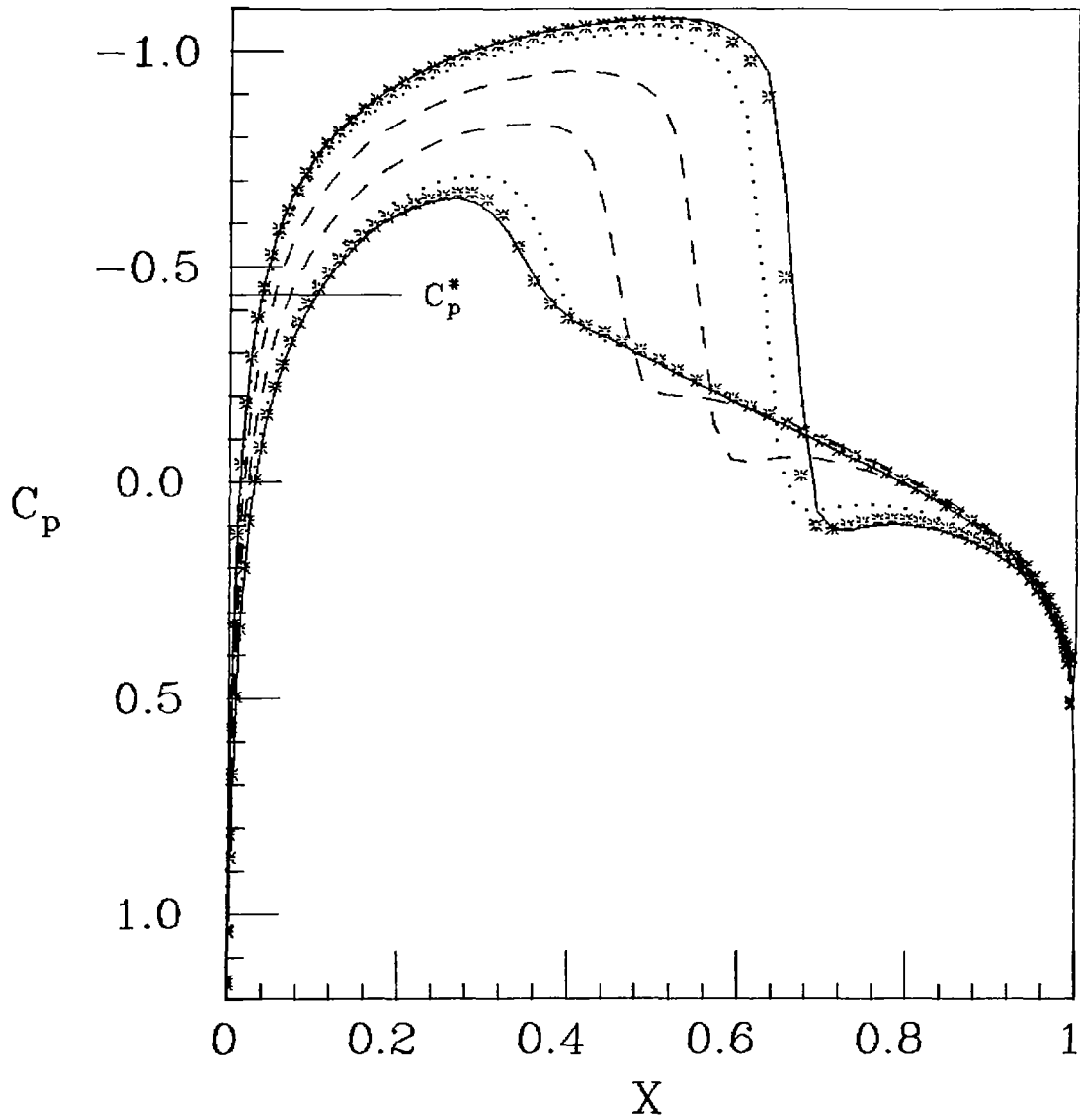


$M_\infty = 0.800$   $\alpha = 0.400$

Figure A.3- Continued

# Plot of $C_p$

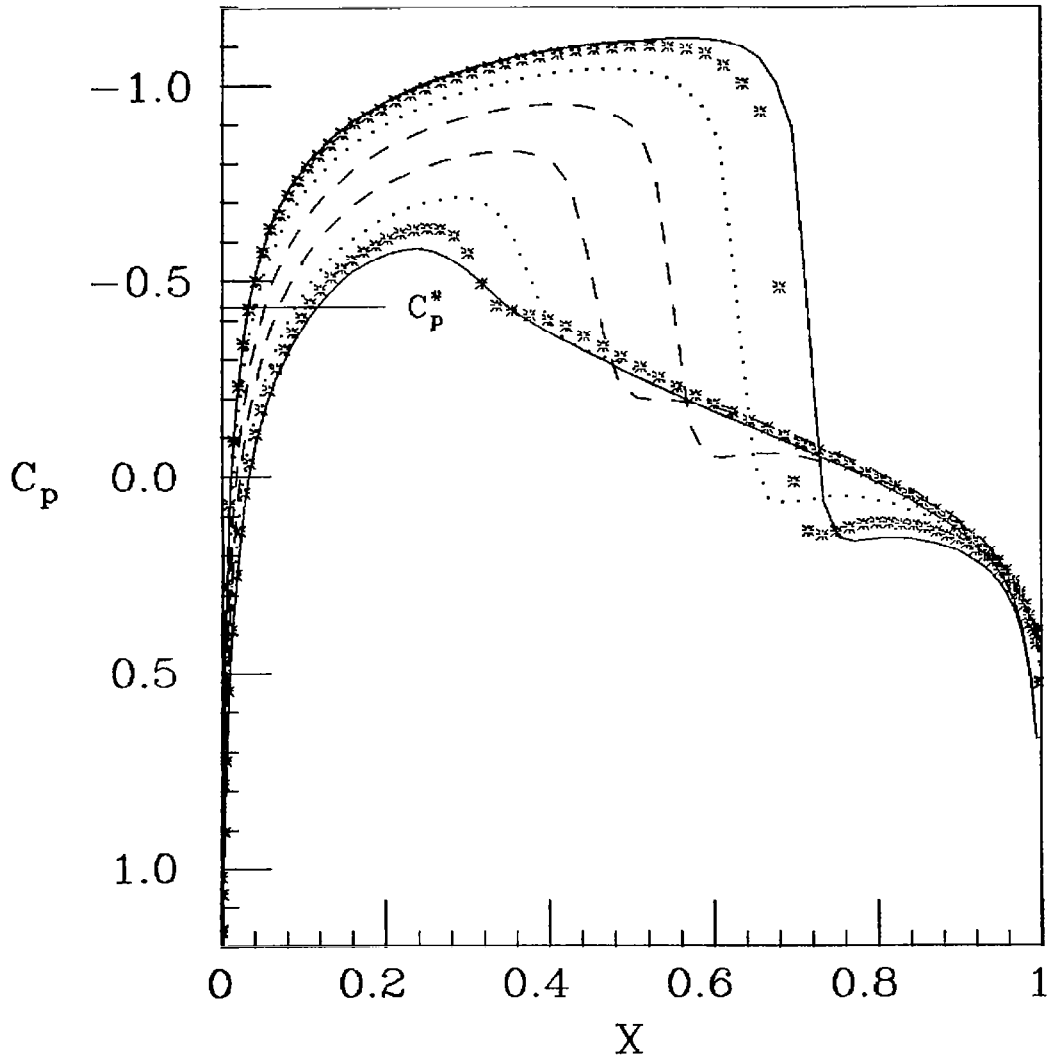
Full  $\phi_{\text{Holst}}$  Nonl.  $\alpha$  Pert.  $\alpha_b = 0.500$   $\alpha_c = 0.200$



$M_\infty = 0.800$   $\alpha = 0.600$

# Plot of $C_p$

Full  $\phi_{\text{Holst}}$  Nonl.  $\alpha$  Pert.  $\alpha_b = 0.500$   $\alpha_c = 0.200$



$M_\infty = 0.800$   $\alpha = 0.700$

Figure A.3- Concluded



C	DEL1=Q1-Q0	MAIN121	XFIX1(I+1)=XCALB(LOC1(I))	MAIN181
	YCR0=YCRIT(M0)	MAIN122	GO TO 10	MAIN182
	YCR1=YCRIT(M1)	MAIN123	5 CONTINUE	MAIN183
	WRITE (6,2000) TITLE	MAIN124	XFIX0(I+1)=XLOC0(LSELCT(I))	MAIN184
	WRITE (6,2010) N,A,B,HEAD0,M0,Q0,HEAD1,M1,Q1	MAIN125	XFIX1(I+1)=XLOC1(LSELCT(I))	MAIN185
	WRITE (6,2020) NAME,HEAD0,NAME,YCR0,HEAD1,NAME,YCR1	MAIN126	10 CONTINUE	MAIN186
C		MAIN127	WRITE (6,2050) NFIX	MAIN187
C....	NORMALIZE X COORDINATES AND LOCATE MINIMUM, MAXIMUM, AND CRITICAL	MAIN128	IF (LSPEC .EQ. 0) GO TO 14	MAIN188
C	POINTS FOR BASE AND CALIBRATION SOLUTIONS.	MAIN129	WRITE (6,2046) HEAD0,HEAD1	MAIN189
C		MAIN130	DO 13 I=1,NSELC	MAIN190
C	CALL SCALE (N,XBASE,1,A,B)	MAIN131	WRITE (6,2047) LOC0(I),LOC1(I)	MAIN191
	CALL LOCATE (N,XBASE,YBASE,YCR0,IGRAD,LMN0,LMX0,NCRO,LCRO,XLOC0)	MAIN132	13 CONTINUE	MAIN192
	CALL SCALE (N,XCALB,1,A,B)	MAIN133	GO TO 18	MAIN193
	CALL LOCATE (N,XCALB,YCALB,YCR1,IGRAD,LMN1,LMX1,NCR1,LCR1,XLOC1)	MAIN134	14 CONTINUE	MAIN194
	WRITE (6,2030)	MAIN135	DO 15 I=1,NSELCT	MAIN195
	WRITE (6,2035)	MAIN136	IF (LSELCT(I) .EQ. 1) WRITE (6,2060) NAME	MAIN196
	CALL UPLOW (A,B,XLOC0,6,NCRO+2,XOUT,FLAG)	MAIN137	IF (LSELCT(I) .EQ. 2) WRITE (6,2070) NAME	MAIN197
	WRITE (6,2040) HEAD0,XOUT(1),FLAG(1),LMN0,XOUT(2),FLAG(2),LMX0	MAIN138	IF (LSELCT(I) .LE. 2) GO TO 15	MAIN198
	IF (NCRO .GT. 0) WRITE (6,2045) NCRO,	MAIN139	LCORH=1	MAIN199
	% (ORD(I),XOUT(I+2),FLAG(I+2),LCRO(I),I=1,NCRO)	MAIN140	LPR=LSELCT(I)-2	MAIN200
	CALL UPLOW (A,B,XLOC1,6,NCR1+2,XOUT,FLAG)	MAIN141	WRITE (6,2080) NAME,ORD(LPR),NCRO	MAIN201
	WRITE (6,2040) HEAD1,XOUT(1),FLAG(1),LMN1,XOUT(2),FLAG(2),LMX1	MAIN142	15 CONTINUE	MAIN202
	IF (NCR1 .GT. 0) WRITE (6,2045) NCR1,	MAIN143	16 CONTINUE	MAIN203
	% (ORD(I),XOUT(I+2),FLAG(I+2),LCR1(I),I=1,NCR1)	MAIN144		MAIN204
C		MAIN145	C....ARRANGE SELECTED FIXED POINTS IN A MONOTONE SEQUENCE.	MAIN205
C....	CHECK FOR INVALID STRAINING SPECIFICATION IF LSPEC = 0.	MAIN146	C	MAIN206
C		MAIN147	CALL SORT (NFIX,XFIX0,ISEQ0)	MAIN207
	IF (LSPEC .EQ. 1) GO TO 4	MAIN148	CALL SORT (NFIX,XFIX1,ISEQ1)	MAIN208
	ICOUNT=0	MAIN149	WRITE (6,2090)	MAIN209
	DO 2 I=1,NSELCT	MAIN150	WRITE (6,2035)	MAIN210
	IF (LSELCT(I) .LE. 2) GO TO 2	MAIN151	CALL UPLOW (A,B,XFIX0,8,NFIX,XOUT,FLAG)	MAIN211
	ICOUNT=ICOUNT+1	MAIN152	WRITE (6,2100) HEAD0,(I,XOUT(I),FLAG(I),I=1,NFIX)	MAIN212
	IF (NCRO .NE. NCR1) LTERM=1	MAIN153	CALL UPLOW (A,B,XFIX1,8,NFIX,XOUT,FLAG)	MAIN213
2	CONTINUE	MAIN154	WRITE (6,2100) HEAD1,(I,XOUT(I),FLAG(I),I=1,NFIX)	MAIN214
C		MAIN155	C	MAIN215
C....	STOP EXECUTION IF CRITICAL POINTS ARE TO BE USED IN STRAINING AND	MAIN156	C....STOP EXECUTION IF ORDER OF OCCURRENCE OF CRITICAL POINTS IN BASE	MAIN216
C	NUMBER OF CRITICAL POINTS IN BASE AND CALIBRATION SOLUTIONS ARE	MAIN157	C AND CALIBRATION SOLUTIONS DOES NOT CORRESPOND.	MAIN217
C	UNEQUAL.	MAIN158	C	MAIN218
C		MAIN159	IF (LSPEC .EQ. 1) GO TO 25	MAIN219
	IF (LTERM .EQ. 1) GO TO 900	MAIN160	DO 20 I=1,NFIX	MAIN220
C		MAIN161	IF (ISEQ0(I) .NE. ISEQ1(I)) GO TO 910	MAIN221
C....	STOP EXECUTION IF NUMBER OF CRITICAL POINTS SELECTED EXCEEDS	MAIN162	20 CONTINUE	MAIN222
C	NUMBER ACTUALLY LOCATED.	MAIN163	25 CONTINUE	MAIN223
C		MAIN164	C	MAIN224
	IF (ICOUNT .GT. NCRO) GO TO 905	MAIN165	C....COMPUTE COEFFICIENTS IN UNIT STRAINING OF XBASE:	MAIN225
C		MAIN166	C	MAIN226
	4 CONTINUE	MAIN167	XSTR = C(I) + D(I)*XBASE, I=1,2, ... ,NSEG,	MAIN227
C		MAIN168	C	MAIN228
C....	LOAD SELECTED STRAINING POINTS INTO FIXED-POINT ARRAYS FOR BASE	MAIN169	C WHERE NSEG IS THE NUMBER OF LINEAR SEGMENTS.	MAIN229
C	AND CALIBRATION SOLUTIONS.	MAIN170	C	MAIN230
C		MAIN171	NSEG=NFIX-1	MAIN231
	NFIX=NSELCT+2	MAIN172	DO 30 I=1,NSEG	MAIN232
	XFIX0(I)=0.0	MAIN173	CNUM=XFIX1(I)*XFIX0(I+1)-XFIX1(I+1)*XFIX0(I)	MAIN233
	XFIX1(I)=0.0	MAIN174	DNUM=XFIX1(I+1)-XFIX1(I)	MAIN234
	XFIX0(NFIX)=1.0	MAIN175	DENOM=XFIX0(I+1)-XFIX0(I)	MAIN235
	XFIX1(NFIX)=1.0	MAIN176	C(I)=CNUM/DENOM	MAIN236
	DO 10 I=1,NSELCT	MAIN177	D(I)=DNUM/DENOM	MAIN237
	IF (LSPEC .EQ. 0) GO TO 5	MAIN178	30 CONTINUE	MAIN238
	XFIX0(I+1)=XBASE(LOC0(I))	MAIN179	C	MAIN239
		MAIN180	C....DETERMINE UNIT STRAINING OF XBASE	MAIN240

C		MAIN241	C		MAIN301
	CALL STRAIN (N,NSEG,XFIX0,XBASE,1.0,XUNIT)	MAIN242	C	....LOCATE MINIMUM, MAXIMUM, AND CRITICAL POINTS IN PERTURBATION.	MAIN302
C		MAIN243	C	SOLUTION.	MAIN303
C	....INTERPOLATE CALIBRATION SOLUTION TO BASE FLOW POINTS CORRESPONDING	MAIN244	C		MAIN304
C	TO UNIT STRAINING.	MAIN245		CALL SCALE (N,XPERT,2,A,B)	MAIN305
C		MAIN246		CALL SCALE (N,XPERT,1,A,B)	MAIN306
	CALL INTERP (N,XCALB,YCALB,XUNIT,YINTP)	MAIN247		CALL LOCATE (N,XPERT,YPERT,YCR2,IGRAD,LMN2,LMX2,NCR2,LCR2,XLOC2)	MAIN307
C		MAIN248		WRITE (6,2130) ICASE,NCASE,M2,Q2,NAME,YCR2	MAIN308
C	....CORRECT VALUES ON EITHER SIDE OF CRITICAL POINTS, IF THESE ARE	MAIN249		WRITE (6,2030)	MAIN309
C	USED IN STRAINING.	MAIN250		WRITE (6,2035)	MAIN310
		MAIN251		CALL UPLOW (A,B,XLOC2,6,NCR2+2,XOUT,FLAG)	MAIN311
	IF (LCORR .EQ. 0) GO TO 36	MAIN252		WRITE (6,2040) HEAD2,XOUT(1),FLAG(1),LMN2,XOUT(2),FLAG(2),LMX2	MAIN312
	DO 35 I=1,NCRI	MAIN253		IF (NCR2 .GT. 0) WRITE (6,2045) NCR2,	MAIN313
	YINTP(LCR0(I))=YCALB(LCR1(I))	MAIN254		% (ORD(I),XOUT(I+2),FLAG(I+2),LCR2(I),I=1,NCR2)	MAIN314
	YINTP(LCR0(I+1))=YCALB(LCR1(I+1))	MAIN255		CALL SCALE (N,XBASE,2,A,B)	MAIN315
	35 CONTINUE	MAIN256			MAIN316
	36 CONTINUE	MAIN257	C	....IF LCHEK .NE. 0 READ IN DATA FOR COMPARISON SOLUTION AND LOCATE	MAIN317
C		MAIN258	C	MINIMUM, MAXIMUM, AND CRITICAL POINTS.	MAIN318
C	....RESTORE PHYSICAL X IN CALIBRATION SOLUTION SINCE IT WILL NOT BE	MAIN259	C		MAIN319
C	FURTHER USED.	MAIN260		IF (LCHEK .EQ. 0) GO TO 90	MAIN320
C		MAIN261		CALL INPUT(3)	MAIN321
	CALL SCALE (N,XCALB,2,A,B)	MAIN262		CALL SCALE (N,XCHEK,1,A,B)	MAIN322
C		MAIN263		CALL LOCATE (N,XCHEK,YCHEK,YCR3,IGRAD,LMN3,LMX3,NCR3,LCR3,XLOC3)	MAIN323
C	....DETERMINE THE UNIT PERTURBATION.	MAIN264		CALL UPLOW (A,B,XLOC3,6,NCR3+2,XOUT,FLAG)	MAIN324
C		MAIN265		WRITE (6,2040) HEAD3,XOUT(1),FLAG(1),LMN3,XOUT(2),FLAG(2),LMX3	MAIN325
	DO 40 I=1,N	MAIN266		IF (NCR3 .GT. 0) WRITE (6,2045) NCR3,	MAIN326
	40 YUNIT(I)=(YINTP(I)-YBASE(I))/DEL1	MAIN267		% (ORD(I),XOUT(I+2),FLAG(I+2),LCR3(I),I=1,NCR3)	MAIN327
C		MAIN268		CALL INTERP (N,XPERT,YPERT,XCHEK,YPRTI)	MAIN328
C	....PRINT UNIT PERTURBATION AND UNIT STRAINING IF LUNIT .NE. 0.	MAIN269		CALL SCALE (N,XPERT,2,A,B)	MAIN329
C		MAIN270		CALL SCALE (N,XCHEK,2,A,B)	MAIN330
		MAIN271		WRITE (6,2135) NAME,NAME,NAME,NAME,NAME	MAIN331
	IF (LUNIT .EQ. 0) GO TO 50	MAIN272		WRITE (6,2140) (I,XBASE(I),YBASE(I),XCALB(I),YCALB(I),	MAIN332
	CALL SCALE (N,XBASE,2,A,B)	MAIN273		% XPERT(I),YPERT(I),XCHEK(I),YCHEK(I),YPRTI(I),	MAIN333
	CALL SCALE (N,XUNIT,2,A,B)	MAIN274		% I=1,N)	MAIN334
	WRITE (6,2110) NAME,NAME	MAIN275		GO TO 100	MAIN335
	WRITE (6,2120) (I,XBASE(I),XUNIT(I),YUNIT(I),I=1,N)	MAIN276		90 CONTINUE	MAIN336
	CALL SCALE (N,XBASE,1,A,B)	MAIN277	C		MAIN337
	50 CONTINUE	MAIN278		CALL SCALE (N,XPERT,2,A,B)	MAIN338
C		MAIN279		WRITE (6,2145) NAME,NAME,NAME	MAIN339
C	....CONSTRUCT PERTURBATION SOLUTIONS FOR TEST CASES (AND COMPARE WITH	MAIN280		WRITE (6,2150) (I,XBASE(I),YBASE(I),XCALB(I),YCALB(I),	MAIN340
C	EXACT SOLUTION, IF AVAILABLE).	MAIN281		% XPERT(I),YPERT(I),I=1,N)	MAIN341
		MAIN282		100 CONTINUE	MAIN342
	DO 200 ICASE=1,NCASE	MAIN283	C		MAIN343
	CALL INPUT(2)	MAIN284		....IF LCHEK .NE. 0 PLOT PERTURBATION AND COMPARISON SOLUTIONS.	MAIN344
	DEL2=Q2-Q0	MAIN285	C		MAIN345
	DEL21=DEL2/DEL1	MAIN286		IF (LCHEK .EQ. 1) CALL PLOT (N,LPERT)	MAIN346
	YCR2=YCRIT(M2)	MAIN287		CALL SCALE (N,XBASE,1,A,B)	MAIN347
	YCR3=YCR2	MAIN288		200 CONTINUE	MAIN348
C		MAIN289		GO TO 999	MAIN349
C	....DETERMINE STRAINED COORDINATE FOR GIVEN PERTURBATION.	MAIN290	C		MAIN350
C		MAIN291		....ABNORMAL TERMINATION OF COMPUTATION.	MAIN351
	CALL STRAIN (N,NSEG,XFIX0,XBASE,DEL21,XPERT)	MAIN292	C		MAIN352
C		MAIN293		900 WRITE (6,9000)	MAIN353
C	....DETERMINE PERTURBATION SOLUTION.	MAIN294		GO TO 999	MAIN354
C		MAIN295		905 WRITE (6,9050)	MAIN355
	DO 60 I=1,N	MAIN296		GO TO 99	MAIN356
	60 YPERT(I)=YBASE(I)+DEL2*YUNIT(I)	MAIN297		910 WRITE (6,9100)	MAIN357
C		MAIN298	C		MAIN358
C	....ADJUST VALUES NEAR CRITICAL POINT FOR MONOTONE BEHAVIOR.	MAIN299		999 WRITE (6,9500)	MAIN359
C		MAIN300		STOP	MAIN360
	IF (LCORR .EQ. 1) CALL MONO (NCR0,LCR0,XPERT,YPERT)				

C  
C.....I/O FORMAT STATEMENTS FOLLOW.

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2000 FORMAT (1H1,132(1H*//
Z 1X,1H*,25X,20A4,25X,1H*//
Z 1X,132(1H*//)
2010 FORMAT (1H ,10(1H<),1X,24HLIST OF INPUT PARAMETERS,1X,10(1H>)//
Z 6X,3HN =,1X,I3//
Z 6X,3HA =,1X,F4.1,4X,3HB =,1X,F4.1//
Z 6X,5A4,1X,4HMO =,1X,F6.4,4X,4HQ0 =,1X,F6.4//
Z 6X,5A4,1X,4HM1 =,1X,F6.4,4X,4HQ1 =,1X,F6.4//)
2020 FORMAT (1H ,10(1H<),1X,18HCritical VALUES OF,1X,A2,1X,10(1H>)//
Z 2(6X,5A4,1X,A2,6HCrit =,1X,F7.4//)
2030 FORMAT (1H ,5(1H<),1X,37HLOCATIONS OF MIN., MAX., AND CRITICAL,
Z 1X,4HPTS.,1X,5(1H>)
2035 FORMAT (1H ,2X,34H(* DENOTES POINT ON LOWER SURFACE))
2040 FORMAT (/6X,5A4//
Z 11X,14HMINIMUM AT X =,1X,F6.4,A1,3X,8H(POINT #,I3,1H)/
Z 11X,14HMAXIMUM AT X =,1X,F6.4,A1,3X,8H(POINT #,I3,1H))
2045 FORMAT (1H ,10X,I1,1X,18HCritical POINT(S)//
Z (15X,A4,6HAT X =,1X,F6.4,A1,3X,
Z 14H(AFTER POINT #,I3,1H))
2046 FORMAT (1H ,10X,2(5A4)/)
2047 FORMAT (1H ,14X,2HX(I,3,1H),15X,2HX(I,3,1H))
2050 FORMAT (///1X,10(1H<),1X,25HSTRAINING POINTS SELECTED,1X,
Z 10(1H>)//
Z 6X,24HNUMBER OF FIXED POINTS ,1X,I1//
Z 6X,5HFIXED POINTS SELECTED (IN ADDITION TO END POINTS) :
Z /)
2060 FORMAT (1H ,10X,16HPOINT OF MINIMUM,1X,A2)
2070 FORMAT (1H ,10X,16HPOINT OF MAXIMUM,1X,A2)
2080 FORMAT (1H ,10X,A2,6HCrit (,A4,3HOF ,I1,1H))
2090 FORMAT (///1X,10(1H<),1X,24HLOCATION OF FIXED POINTS,1X,10(1H>))
2100 FORMAT (/6X,5A4//
Z (1H ,10X,5HXFIX(I,1,3H) =,1X,F6.4,A1))
2110 FORMAT (1H1,27(1H*//
Z 1X,1H*,1X,20HUNIT PERTURBATION OF,1X,A2,1X,1H*//
Z 1X,1H*,12X,1HC,12X,1H*//
Z 1X,1H*,1X,23HUNIT STRAINING OF XBASE,1X,1H*//
Z 1X,27(1H*//)
Z 1X,5HPOINT,4X,5HXBASE,4X,8HXSTRUNIT,3X,A2,4HUNIT//)
2120 FORMAT (1H ,1X,I3,1X,3F10.4)
2130 FORMAT (1H1,27(1H*//
Z 1X,19H* OUTPUT FOR CASE #,I1,4H OF ,I1,2H #/
Z 1X,27(1H*//)
Z 6X,4HM2 =,1X,F6.4//
Z 6X,4HQ2 =,1X,F6.4//
Z 6X,A2,6HCrit =,1X,F7.4//)
2135 FORMAT (///1X,5HPOINT,4X,5HXBASE,5X,A2,4HBASE,
Z 4X,5HXCALB,5X,A2,4HCALB,
Z 4X,5HXPert,5X,A2,4HPert,
Z 4X,5HXCHEK,5X,A2,4HCHEK,
Z 2X,A2,9HPert(INT//)
2140 FORMAT (1H ,1X,I3,1X,9F10.4)
2145 FORMAT (///1X,5HPOINT,4X,5HXBASE,5X,A2,4HBASE,
Z 4X,5HXCALB,5X,A2,4HCALB,
Z 4X,5HXPert,5X,A2,4HPert//)
2150 FORMAT (1H ,1X,I3,1X,6F10.4)
9000 FORMAT (///1X,28HNUMBER OF CRITICAL POINTS IN/
Z 1X,30HBASE AND CALIBRATION SOLUTIONS/

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|MAIN361 % 1X,31HARE UNEQUAL - CALCULATION ENDED) |MAIN421
|MAIN362 9050 FORMAT (///1X,25HNUMBER OF CRITICAL POINTS/ |MAIN422
|MAIN363 % 1X,23HSELECTED EXCEEDS NUMBER/ |MAI423
|MAIN364 % 1X,30HACTUALLY LOCATED - CALCULATION/ |MAI424
|MAIN365 % 1X,5HENDED) |MAI425
|MAIN366 9100 FORMAT (///1X,28HORDER OF SPECIFIED POINTS IN/ |MAI426
|MAIN367 % 1X,30HBASE AND CALIBRATION SOLUTIONS/ |MAI427
|MAIN368 % 1X,39HDOES NOT CORRESPOND - CALCULATION ENDED) |MAI428
|MAIN369 9500 FORMAT (1H1) |MAI429
|MAIN370 END |MAI430
|MAIN371 SUBROUTINE INPUT (ICALL) |INPU001
|MAIN372 C |INPU002
|MAIN373 C ALL INPUT FOR PROGRAM PERTURB IS READ BY THIS SUBROUTINE, AND IS |INPU003
|MAIN374 C REQUIRED IN THE FOLLOWING ORDER (FOR DETAILS, REFER TO |INPU004
|MAIN375 C ACCOMPANYING MANUAL). |INPU005
|MAIN376 C |INPU006
|MAIN377 C**** CARD #1 (1615) ***** |INPU007
|MAIN378 C |INPU008
|MAIN379 C N NUMBER OF DATA POINTS IN BASE AND CALIBRATION |INPU009
|MAIN380 C SOLUTIONS. |INPU010
|MAIN381 C |INPU011
|MAIN382 C NCASE NUMBER OF CASES FOR WHICH PERTURBATION SOLUTIONS ARE |INPU012
|MAIN383 C TO BE COMPUTED. |INPU013
|MAIN384 C |INPU014
|MAIN385 C LSPEC CONTROLS HOW INVARIANT POINTS IN STRAINING ARE |INPU015
|MAIN386 C SPECIFIED (SEE CARD #2). |INPU016
|MAIN387 C |INPU017
|MAIN388 C LSPEC = 0 ... INVARIANT POINTS SELECTED FROM AMONG |INPU018
|MAIN389 C THOSE LOCATED BY THE PROGRAM |INPU019
|MAIN390 C LSPEC = 1 ... INVARIANT POINTS PRESELECTED BY USER |INPU020
|MAIN391 C |INPU021
|MAIN392 C LECHO CONTROLS WHETHER OR NOT INPUT DECK IS PRINTED. |INPU022
|MAIN393 C |INPU023
|MAIN394 C LECHO = 0 ... NO OUTPUT |INPU024
|MAIN395 C LECHO = 1 ... OUTPUT |INPU025
|MAIN396 C |INPU026
|MAIN397 C LUNIT CONTROLS WHETHER OR NOT UNIT COORDINATE STRAINING |INPU027
|MAIN398 C AND UNIT PERTURBATION ARE PRINTED. |INPU028
|MAIN399 C |INPU029
|MAIN400 C LUNIT = 0 ... NO OUTPUT |INPU030
|MAIN401 C LUNIT = 1 ... OUTPUT |INPU031
|MAIN402 C |INPU032
|MAIN403 C LCHEK SPECIFIES WHETHER OR NOT PERTURBATION SOLUTION IS TO |INPU033
|MAIN404 C BE CHECKED WITH AND PLOTTED AGAINST AN EXACT |INPU034
|MAIN405 C COMPARISON SOLUTION. |INPU035
|MAIN406 C |INPU036
|MAIN407 C LCHEK = 0 ... NO COMPARISON, NO PLOT |INPU037
|MAIN408 C LCHEK = 1 ... COMPARISON, PLOT |INPU038
|MAIN409 C |INPU039
|MAIN410 C LPERT SPECIFIES TYPE OF PERTURBATION. OPERATIONAL ONLY |INPU040
|MAIN411 C WHEN LCHEK = 1 AND AFFECTS ONLY OUTPUT FROM PLOT |INPU041
|MAIN412 C SUBROUTINE. |INPU042
|MAIN413 C |INPU043
|MAIN414 C LPERT = 1 ... THICKNESS-RATIO PERTURBATION |INPU044
|MAIN415 C LPERT = 2 ... ANGLE-OF-ATTACK PERTURBATION |INPU045
|MAIN416 C LPERT = 3 ... MACH-NUMBER PERTURBATION |INPU046
|MAIN417 C |INPU047
|MAIN418 C**** CARD #2 (1615) ***** |INPU048
|MAIN419 C |INPU049
|MAIN420 C REQUIRED INPUT DEPENDS ON VALUE OF LSPEC: |INPU050

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C |INPU051
C |INPU052
C |INPU053
C |INPU054
C |INPU055
C |INPU056
C |INPU057
C |INPU058
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C |INPU060
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C |INPU168
C |INPU169
C |INPU170

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INPUT FOR LSPEC = 0

NSELCT NUMBER OF POINTS (IN ADDITION TO END POINTS) TO BE HELD INVARIANT IN STRAINING. NOTE: 1 <= NSELCT <= 6.

LSELCT(I) ... ARRAY OF LENGTH 6 OF WHICH NSELCT ELEMENTS ARE READ IN. SPECIFIES NATURE OF POINTS TO BE HELD INVARIANT ACCORDING TO THE CODE:

1 ... MINIMUM PT. HELD INVARIANT  
 2 ... MAXIMUM PT. HELD INVARIANT  
 3 ... 1ST CRITICAL PT. HELD INVARIANT  
 4 ... 2ND CRITICAL PT. HELD INVARIANT  
 5 ... 3RD CRITICAL PT. HELD INVARIANT  
 6 ... 4TH CRITICAL PT. HELD INVARIANT

NOTE THAT THE CODE NUMBERS CAN BE ASSIGNED IN ANY ORDER, E.G.

LSELCT(1) = 1  
 LSELCT(2) = 3  
 LSELCT(3) = 4

IS EQUIVALENT TO

LSELCT(1) = 4  
 LSELCT(2) = 1  
 LSELCT(3) = 3

BOTH CORRESPONDING TO NSELCT = 3 WITH THE MINIMUM, AND FIRST AND SECOND CRITICAL POINTS HELD INVARIANT.

INPUT FOR LSPEC = 1

NSELCT NUMBER OF POINTS (IN ADDITION TO END POINTS) TO BE HELD INVARIANT IN STRAINING. NOTE: 1 <= NSELCT <= 6.

LOC0(I) ARRAY OF LENGTH 6 OF WHICH NSELCT ELEMENTS ARE READ IN. SPECIFIES SUBSCRIPTS OF THOSE BASE FLOW POINTS WHICH ARE TO BE HELD INVARIANT.

LOC1(I) ARRAY OF LENGTH 6 OF WHICH NSELCT ELEMENTS ARE READ IN. SPECIFIES SUBSCRIPTS OF THOSE CALIBRATION FLOW POINTS WHICH ARE TO BE HELD INVARIANT.

CARD #3 (20A4)

TITLE IDENTIFIES JOB - PRINTED AS HEADLINE ON FIRST PAGE OF OUTPUT.

CARD #4 (A2)

NAME CHARACTER STRING OF LENGTH 2 WHICH SYMBOLIZES DEPENDENT VARIABLE, E.G. 'CP' FOR PRESSURE

COEFFICIENT.

CARD #5 (8F10.6)

A SCALING PARAMETER (A = -X(1), WHERE X(1) IS FIRST DATA POINT ON LOWER SURFACE ... SEE MANUAL).

B SCALING PARAMETER (B = X(N), WHERE X(N) IS LAST DATA POINT ON UPPER SURFACE ... SEE MANUAL).

CARD #6 (8F10.6)

H0 ONCOMING MACH NUMBER IN BASE SOLUTION.

Q0 VALUE OF PERTURBATION PARAMETER IN BASE SOLUTION.

ONE SET OF K CARDS (8F10.6), WHERE K = 1 + INT(N/8)

XBASE(I), I=1,N ... X COORDINATE IN BASE SOLUTION.

ONE SET OF K CARDS (8F10.6), K AS ABOVE

YBASE(I), I=1,N ... DEPENDENT VARIABLE IN BASE SOLUTION.

NEXT CARD (8F10.6)

M1 ONCOMING MACH NUMBER IN CALIBRATION SOLUTION.

Q1 VALUE OF PERTURBATION PARAMETER IN CALIBRATION SOLUTION.

ONE SET OF K CARDS (8F10.6), K AS ABOVE

XCALB(I), I= 1,N ... X COORDINATE IN CALIBRATION SOLUTION.

ONE SET OF K CARDS (8F10.6), K AS ABOVE

YCALB(I), I=1,N ... DEPENDENT VARIABLE IN CALIBRATION SOLUTION..

NCASE SETS OF CARDS, EACH SET COMPRISED AS FOLLOWS:

FIRST CARD (8F10.6)

M2 ONCOMING MACH NUMBER IN SOLUTION TO BE COMPUTED.

Q2 VALUE OF PERTURBATION PARAMETER IN SOLUTION TO BE COMPUTED.

ONE SET OF K CARDS (8F10.6), K AS ABOVE

XCHEK(I), I= 1,N ... X COORDINATE IN COMPARISON SOLUTION.

ONE SET OF K CARDS (8F10.6), K AS ABOVE

C  
C  
C  
C  
C  
C  
C  
C

```

YCHEK(I), I=1,N ...
DEPENDENT VARIABLE IN COMPARISON SOLUTION.

***** THE LATTER TWO SETS OF K CARDS ARE
* NOTE * OMITTED WHEN LCHEK = 0 (NO COMPARISON
***** SOLUTION AVAILABLE).

DIMENSION LOC0(6),LOC1(6),LSELCT(6),TITLE(20)
DIMENSION XBASE(200),XCALB(200),XPERT(200),XCHEK(200),
% YBASE(200),YCALB(200),YPERT(200),YCHEK(200)
REAL M0,M1,M2
INTEGER*2 NAME
COMMON /PARAM/ TITLE,LOC0,LOC1,LSELCT,N,NCASE,LSPEC,LECHO,LUNIT,
% LCHEK,LPERT,NSELCT,A,B,NAME
COMMON /PERT/ M0,M1,M2,Q0,Q1,Q2,YCRO,YCR1,YCR2
COMMON /XY/ XBASE,XCALB,XPERT,XCHEK,YBASE,YCALB,YPERT,YCHEK
GO TO (100,200,300), ICALL
100 READ (5,1000) N,NCASE,LSPEC,LECHO,LUNIT,LCHEK,LPERT
IF (LSPEC.EQ.0) READ (5,1000) NSELCT,(LSELCT(I),I=1,NSELCT)
IF (LSPEC.EQ.1) READ (5,1000) NSELCT,(LOC0(I),I=1,NSELCT),
% (LOC1(I),I=1,NSELCT)
READ (5,1050) TITLE
READ (5,1100) NAME
READ (5,1200) A,B
READ (5,1200) M0,Q0
READ (5,1200) (XBASE(I),I=1,N)
READ (5,1200) (YBASE(I),I=1,N)
READ (5,1200) M1,Q1
READ (5,1200) (XCALB(I),I=1,N)
READ (5,1200) (YCALB(I),I=1,N)
RETURN
200 READ (5,1200) M2,Q2
RETURN
300 READ (5,1200) (XCHEK(I),I=1,N)
READ (5,1200) (YCHEK(I),I=1,N)
RETURN
1000 FORMAT (16I5)
1050 FORMAT (20A4)
1100 FORMAT (A2)
1200 FORMAT (8F10.6)
END
SUBROUTINE ECHINP
DIMENSION LOC0(6),LOC1(6),LSELCT(6),TITLE(20)
DIMENSION XBASE(200),XCALB(200),XPERT(200),XCHEK(200),
% YBASE(200),YCALB(200),YPERT(200),YCHEK(200)
REAL M0,M1,M2
INTEGER*2 NAME
COMMON /PARAM/ TITLE,LOC0,LOC1,LSELCT,N,NCASE,LSPEC,LECHO,LUNIT,
% LCHEK,LPERT,NSELCT,A,B,NAME
COMMON /PERT/ M0,M1,M2,Q0,Q1,Q2,YCRO,YCR1,YCR2
COMMON /XY/ XBASE,XCALB,XPERT,XCHEK,YBASE,YCALB,YPERT,YCHEK
WRITE (6,1400)
WRITE (6,1500) N,NCASE,LSPEC,LECHO,LUNIT,LCHEK,LPERT
IF (LSPEC.EQ.0) WRITE (6,1500) NSELCT,(LSELCT(I),I=1,NSELCT)
IF (LSPEC.EQ.1) WRITE (6,1500) NSELCT,(LOC0(I),I=1,NSELCT),
% (LOC1(I),I=1,NSELCT)
%
WRITE (6,1550) TITLE
WRITE (6,1600) NAME
WRITE (6,1700) A,B
WRITE (6,1700) M0,Q0

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INPU171  
INPU172  
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INPU178  
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INPU210  
INPU211  
ECHI001  
ECHI002  
ECHI003  
ECHI004  
ECHI005  
ECHI006  
ECHI007  
ECHI008  
ECHI009  
ECHI010  
ECHI011  
ECHI012  
ECHI013  
ECHI014  
ECHI015  
ECHI016  
ECHI017  
ECHI018  
ECHI019

```

WRITE (6,1700) (XBASE(I),I=1,N)
WRITE (6,1700) (YBASE(I),I=1,N)
WRITE (6,1700) M1,Q1
WRITE (6,1700) (XCALB(I),I=1,N)
WRITE (6,1700) (YCALB(I),I=1,N)
RETURN
1400 FORMAT (1H1,25(1H*)/
% 1X,1H*,1X,21HLISTING OF INPUT DECK,1X,1H*/
% 1X,25(1H*)//)
1500 FORMAT (1X,16I5)
1550 FORMAT (1X,20A4)
1600 FORMAT (1X,A2)
1700 FORMAT (1X,8F10.6)
END
SUBROUTINE BANNER
WRITE (6,1300)
WRITE (6,1310)
1300 FORMAT (1H1,10(/),49X,55(1H*)/49X,1H*,53X,1H*/
% 49X,1H*,19X,15HPROGRAM PERTURB,19X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,8X,37HCALCULATES NONLINEAR SINGLE-PARAMETER,
% 8X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,13X,27HCONTINUOUS OR DISCONTINUOUS,
% 13X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,15X,22HPERTURBATION SOLUTIONS,16X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,9X,34HWHICH REPRESENT A CHANGE IN EITHER,
% 10X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,13X,27HGEOMETRY OR FLOW CONDITIONS,
% 13X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,4X,44HBY EMPLOYING A STRAINED-COORDINATE PROCEDURE,
% 5X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,4X,45HUTILIZING A UNIT PERTURBATION DETERMINED FROM,
% 4X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,14X,25HTWO PREVIOUSLY CALCULATED,
% 14X,1H*/49X,1H*,53X,1H*)
1310 FORMAT ( 49X,1H*,9X,34H'BASE' AND 'CALIBRATION' SOLUTIONS,
% 10X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,4X,45HDISPLACED FROM ONE ANOTHER BY SOME REASONABLE,
% 4X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,8X,36HCHANGE IN GEOMETRY OR FLOW CONDITION,
% 9X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,21X,10HWRITTEN BY,22X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,7X,39HJAMES P. ELLIOTT AND STEPHEN S. STAHARA,
% 7X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,7X,38HNIELSEN ENGINEERING AND RESEARCH, INC.,
% 8X,1H*/49X,1H*,53X,1H*/
% 49X,1H*,14X,25HMOUNTAIN VIEW, CALIFORNIA,14X,1H*/49X,1H*,53X,1H*/
% 49X,55(1H*))
RETURN
END
SUBROUTINE SCALE (N,X,M,A,B)

```

ECHI020  
ECHI021  
ECHI022  
ECHI023  
ECHI024  
ECHI025  
ECHI026  
ECHI027  
ECHI028  
ECHI029  
ECHI030  
ECHI031  
ECHI032  
ECHI033  
BANN001  
BANN002  
BANN003  
BANN004  
BANN005  
BANN006  
BANN007  
BANN008  
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BANN033  
BANN034  
BANN035  
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BANN037  
SCAL001  
SCAL002  
SCAL003  
SCAL004  
SCAL005  
SCAL006  
SCAL007  
SCAL008  
SCAL009

C  
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GO TO (1,6),M |SCAL010
1 CONTINUE |SCAL011
  NZ=0 |SCAL012
  DO 2 I=2,N |SCAL013
  IF (X(I) .LT. X(I-1)) NZ=I |SCAL014
2 CONTINUE |SCAL015
  DO 5 I=1,N |SCAL016
  IF (I .LE. NZ) T=-X(I) |SCAL017
  IF (I .GT. NZ) T=X(I) |SCAL018
  X(I)=(T-A)/(B-A) |SCAL019
5 CONTINUE |SCAL020
  RETURN |SCAL021
6 DO 7 I=1,N |SCAL022
  X(I)=ABS((B-A)*X(I)+A) |SCAL023
7 CONTINUE |SCAL024
  RETURN |SCAL025
  END |SCAL026
  SUBROUTINE LOCATE (N,X,Y,YCRIT,IGRAD,LMIN,LMAX,NCRIT,LCRIT,XLOC) |LOCA001
C |LOCA002
C .....OPERATES ON THE INPUT ARRAY Y, LOCATING MINIMUM AND MAXIMUM |LOCA003
C VALUES, AND ALL CRITICAL POINTS (Y=YCRIT) FOR WHICH DY/DX (IN |LOCA004
C PHYSICAL COORDINATES) HAS ALGEBRAIC SIGN GIVEN BY IGRAD. NCRIT |LOCA005
C IS NUMBER OF CRITICAL POINTS. POINTS FOUND ARE STORED IN THE ARRAY |LOCA006
C XLOC AS FOLLOWS: |LOCA007
C |LOCA008
C |LOCA009
C |LOCA010
C |LOCA011
C |LOCA012
C |LOCA013
C |LOCA014
C |LOCA015
C |LOCA016
C |LOCA017
C |LOCA018
C |LOCA019
C |LOCA020
C |LOCA021
C |LOCA022
C |LOCA023
C |LOCA024
C |LOCA025
C |LOCA026
C |LOCA027
C |LOCA028
C |LOCA029
C |LOCA030
C |LOCA031
C |LOCA032
C |LOCA033
C |LOCA034
C |LOCA035
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C |LOCA037
C |LOCA038
C |LOCA039
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C |LOCA041
C |LOCA042
C |LOCA043
C |LOCA044
C |LOCA045
C |LOCA046
C |LOCA047
C |SORT001
C |SORT002
C |SORT003
C |SORT004
C |SORT005
C |SORT006
C |SORT007
C |SORT008
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C |SORT011
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C |SORT020
C |SORT021
C |SORT022
C |SORT023
C |INTE001
C |INTE002
C |INTE003
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C |INTE022
C |INTE023
C |INTE024
C |INTE025
C |INTE026
C |INTE027
C |INTE028
C |INTE029
C |STRA001
C |STRA002
C |STRA003
C |STRA004
XLOC(1) = MINIMUM PT. |LOCA009
XLOC(2) = MAXIMUM PT. |LOCA010
XLOC(3) = CRITICAL PT. #1 |LOCA011
... = ... |LOCA012
XLOC(4) = CRITICAL PT. #4 |LOCA013
DIMENSION X(200),Y(200),LCRIT(4),XCRIT(4),XLOC(6) |LOCA015
COMMON /FLOREV/ IREV |LOCA016
IFLOW=-1 |LOCA017
LMIN=1 |LOCA018
LMAX=1 |LOCA019
ISTART=2 |LOCA020
IF (IREV .EQ. 0) GO TO 5 |LOCA021
LMIN=2 |LOCA022
LMAX=2 |LOCA023
ISTART=3 |LOCA024
5 CONTINUE |LOCA025
  NCRIT=0 |LOCA026
  DO 100 I=ISTART,N |LOCA027
  IF (IREV .NE. 0 .AND. I .EQ. N) GO TO 10 |LOCA028
  IF (Y(I) .GT. Y(LMAX)) LMAX=I |LOCA029
  IF (Y(I) .LT. Y(LMIN)) LMIN=I |LOCA030
10 CONTINUE |LOCA031
  IF ((Y(I) .GT. YCRIT .AND. Y(I-1) .GT. YCRIT) .OR. |LOCA032
  % (Y(I) .LT. YCRIT .AND. Y(I-1) .LT. YCRIT)) GO TO 100 |LOCA033
  IF (I .GT. IREV) IFLOW=1 |LOCA034
  IF ((Y(I)-Y(I-1))*FLOAT(IFLOW*IGRAD) .LT. 0.0) GO TO 100 |LOCA035
  NCRIT=NCRIT+1 |LOCA036
  LCRIT(NCRIT)=I-1 |LOCA037
  SLOPE=(X(I)-X(I-1))/(Y(I)-Y(I-1)) |LOCA038
  XCRIT(NCRIT)=X(I-1)+SLOPE*(YCRIT-Y(I-1)) |LOCA039
100 CONTINUE |LOCA040
  XLOC(1)=X(LMIN) |LOCA041
  XLOC(2)=X(LMAX) |LOCA042
  IF (NCRIT .EQ. 0) RETURN |LOCA043
DO 200 I=1,NCRIT |LOCA044
200 XLOC(I+2)=XCRIT(I) |LOCA045
  RETURN |LOCA046
  END |LOCA047
  SUBROUTINE SORT (N,X,ISEQ) |SORT001
C |SORT002
C .....ARRANGES THE SET X(1), X(2), ... , X(N) IN A MONOTONE INCREASING |SORT003
C SEQUENCE. ISEQ GIVES ORDER OF SUBSCRIPTS IN REARRANGED SEQUENCE. |SORT004
C |SORT005
C |SORT006
C |SORT007
C |SORT008
C |SORT009
C |SORT010
C |SORT011
C |SORT012
C |SORT013
C |SORT014
C |SORT015
C |SORT016
C |SORT017
C |SORT018
C |SORT019
C |SORT020
C |SORT021
C |SORT022
C |SORT023
C |INTE001
C |INTE002
C |INTE003
C |INTE004
C |INTE005
C |INTE006
C |INTE007
C |INTE008
C |INTE009
C |INTE010
C |INTE011
C |INTE012
C |INTE013
C |INTE014
C |INTE015
C |INTE016
C |INTE017
C |INTE018
C |INTE019
C |INTE020
C |INTE021
C |INTE022
C |INTE023
C |INTE024
C |INTE025
C |INTE026
C |INTE027
C |INTE028
C |INTE029
C |STRA001
C |STRA002
C |STRA003
C |STRA004
  DIMENSION X(6),ISEQ(6) |SORT006
  NM1=N-1 |SORT007
  DO 1 I=1,N |SORT008
  ISEQ(I)=I |SORT009
  10 ITEST=0 |SORT010
  DO 100 I=1,NM1 |SORT011
  IF (X(I) .LE. X(I+1)) GO TO 100 |SORT012
  XSAVE=X(I) |SORT013
  X(I)=X(I+1) |SORT014
  X(I+1)=XSAVE |SORT015
  ISAVE=ISEQ(I) |SORT016
  ISEQ(I)=ISEQ(I+1) |SORT017
  ISEQ(I+1)=ISAVE |SORT018
  ITEST=1 |SORT019
  100 CONTINUE |SORT020
  IF (ITEST .EQ. 1) GO TO 10 |SORT021
  RETURN |SORT022
  END |SORT023
  SUBROUTINE INTERP (N,X,Y,XI,YI) |INTE001
C |INTE002
C .....GIVEN THE SET OF POINTS X(I), Y(I), I=1,N, AND THE SET XI(J), |INTE003
C J=1,N, USES LINEAR INTERPOLATION TO COMPUTE THE SET YI(J), J=1,N. |INTE004
C |INTE005
C |INTE006
C |INTE007
C |INTE008
C |INTE009
C |INTE010
C |INTE011
C |INTE012
C |INTE013
C |INTE014
C |INTE015
C |INTE016
C |INTE017
C |INTE018
C |INTE019
C |INTE020
C |INTE021
C |INTE022
C |INTE023
C |INTE024
C |INTE025
C |INTE026
C |INTE027
C |INTE028
C |INTE029
C |STRA001
C |STRA002
C |STRA003
C |STRA004
  DIMENSION X(200),Y(200),XI(200),YI(200) |INTE006
  NM1=N-1 |INTE007
  JSTART=1 |INTE008
  DO 100 I=1,N |INTE009
  IF (XI(I) .LE. X(1)) GO TO 10 |INTE010
  IF (XI(I) .GE. X(N)) GO TO 20 |INTE011
  GO TO 30 |INTE012
  10 J=1 |INTE013
  GO TO 95 |INTE014
  20 J=N-1 |INTE015
  GO TO 95 |INTE016
  30 CONTINUE |INTE017
  DO 90 J=JSTART,NM1 |INTE018
  IF (XI(I) .NE. X(J)) GO TO 40 |INTE019
  YI(I)=Y(J) |INTE020
  GO TO 100 |INTE021
  40 IF (XI(I) .GT. X(J) .AND. XI(I) .LT. X(J+1)) GO TO 95 |INTE022
  90 CONTINUE |INTE023
  95 SLOPE=(Y(J+1)-Y(J))/(X(J+1)-X(J)) |INTE024
  YI(I)=Y(J)+SLOPE*(XI(I)-X(J)) |INTE025
  JSTART=J |INTE026
  100 CONTINUE |INTE027
  RETURN |INTE028
  END |INTE029
  SUBROUTINE STRAIN (N,NSEG,XFIX,XIN,PARM,XOUT) |STRA001
C |STRA002
C .....COMPUTES STRAINED COORDINATE FROM INPUT ARRAY XIN, USING PIECEWISE |STRA003
C LINEAR STRAINING WITH NSEG LINEAR SEGMENTS. FOR UNIT STRAINING, |STRA004

```

C	INPUT VALUE OF PARM IS 1.0; FOR GENERAL CASE,	STPA005	WRITE (4,1300)	PLOT022
C	PARM = (Q2-Q0)/(Q1-Q0).	STPA006	WRITE (4,1350) SYM(LPRT),SYM(LPRT),Q0,SYM(LPRT),Q1	PLOT023
C		STPA007	WRITE (4,1360)	PLOT024
	DIMENSION XFIX(8),XIN(200),XOUT(200)	STPA008	WRITE (4,1370) SUB(LPRT),SUB(LPRT),SUB(LPRT)	PLOT025
	COMMON /COEFF/ C(7),D(7)	STPA009	WRITE (4,1400)	PLOT026
	JSTART=1	STPA010	IF (LPRT .NE. 3) WRITE (4,1500) M0,SYM(LPRT),Q2	PLOT027
	DO 50 I=1,N	STPA011	IF (LPRT .EQ. 3) WRITE (4,1505) Q2	PLOT028
	DO 40 J=JSTART,NSEG	STPA012	IF (LPRT .HE. 3) GO TO 10	PLOT029
	IF (XIN(I) .GE. XFIX(J) .AND. XIN(I) .LE. XFIX(J+1)) GO TO 45	STPA013	WRITE (4,5000) XL,YCR0,XR0,YCR0	PLOT030
40	CONTINUE	STPA014	WRITE (4,1700)	PLOT031
45	XOUT(I)=XIN(I)+PARM*(C(J)+D(J)-1.0)*XIN(I)	STPA015	WRITE (4,1530) XT0,YCR0,S0	PLOT032
	JSTART=J	STPA016	WRITE (4,5000) XL,YCR1,XR1,YCR1	PLOT033
50	CONTINUE	STPA017	WRITE (4,1800)	PLOT034
	RETURN	STPA018	WRITE (4,1530) XT1,YCR1,S1	PLOT035
	END	STPA019	10 CONTINUE	PLOT036
	SUBROUTINE MONO (N,L,X,Y)	MOHO001	WRITE (4,5000) XL,YCR2,XR2,YCR2	PLOT037
C		MOHO002	WRITE (4,1510)	PLOT038
C	.....CHECKS POINTS IN VICINITY OF A CRITICAL POINT FOR MONOTONE	MOHO003	WRITE (4,1550) XT2,YCR2	PLOT039
C	BEHAVIOR, AND ADJUSTS VALUES TO GIVE A LINEAR PROFILE.	MOHO004	WRITE (4,5000) (X0(I),Y0(I),I=1,N)	PLOT040
C		MOHO005	WRITE (4,1700)	PLOT041
	DIMENSION L(4),X(200),Y(200)	MOHO006	WRITE (4,5000) (X1(I),Y1(I),I=1,N)	PLOT042
	DO 100 I=1,N	MOHO007	WRITE (4,1800)	PLOT043
	LS=L(I)	MOHO008	WRITE (4,5000) (X2(I),Y2(I),I=1,N)	PLOT044
	Y1=Y(LS-1)	MOHO009	WRITE (4,1900)	PLOT045
	Y2=Y(LS)	MOHO010	WRITE (4,5000) (X3(I),Y3(I),I=1,N)	PLOT046
	Y3=Y(LS+1)	MOHO011	WRITE (4,1510)	PLOT047
	Y4=Y(LS+2)	MOHO012	RETURN	PLOT048
	IF ((Y1 .LT. Y2) .AND. (Y2 .LT. Y3) .AND. (Y3 .LT. Y4)) GO TO 100	MOHO013	1000 FORMAT (45H//PERTPLOT JOB , ' JIM ELLIOTT',REGION=512K/	PLOT049
	IF ((Y1 .GT. Y2) .AND. (Y2 .GT. Y3) .AND. (Y3 .GT. Y4)) GO TO 100	MOHO014	%	PLOT050
	X1=X(LS-1)	MOHO015	% 40HSET DEVICE VERSATEC CONTINUOUS INTENSITY/	PLOT051
	X2=X(LS)	MOHO016	% 18HSET CARD LENGTH 80/	PLOT052
	X3=X(LS+1)	MOHO017	% 15HSET FONT DUPLEX)	PLOT053
	X4=X(LS+2)	MOHO018	1050 FORMAT (9HNEW FRANE/	PLOT054
	SLOPE=(Y4-Y1)/(X4-X1)	MOHO019	% 27HSET TICKS TOP OFF RIGHT OFF/	PLOT055
	Y(LS)=Y1+SLOPE*(X2-X1)	MOHO020	% 22HSET WINDOW X 2 7 Y 2 8/	PLOT056
	Y(LS+1)=Y1+SLOPE*(X3-X1)	MOHO021	% 20HSET SYMBOL 8P SIZE 1)	PLOT057
100	CONTINUE	MOHO022	1100 FORMAT (18HSET LIMITS X 0 1 Y,2F5.1)	PLOT058
	RETURN	MOHO023	1200 FORMAT (51HTITLE 4.5 9.5 CENTER SIZE 3 SPACES 7 'PLOT OF C2P3'/	PLOT059
	END	MOHO024	% 19HCASE ' LLL LL CLC')	PLOT060
	SUBROUTINE PLOT (N,LPERT)	PLOT001	1300 FORMAT (39HTITLE 4.5 8.7 CENTER SIZE 1.5 SPACES 38,1X,	PLOT061
C		PLOT002	% 24H'FULL F2HOLST3 NONL.')	PLOT062
C	.....CREATES FILE OF COMMANDS FOR PROGRAM 'TOPDRAW' AT STANFORD CENTER	PLOT003	1350 FORMAT (4HNDRE,1X,1H',A1,1X,5HPERT.,3X,A1,5H2B3 =,F6.3,X,A1,	PLOT063
C	FOR INFORMATION PROCESSING (S.C.I.P.). CALLED ONLY ONCE IN MAIN	PLOT004	% 5H2C3 =,F6.3,1H')	PLOT064
C	PROGRAM AND MAY BE DELETED OR REPLACED.	PLOT005	1360 FORMAT (4HCASE,1X,24H' LLL GC LLLC LLL ')	PLOT065
C		PLOT006	1370 FORMAT (4HCASE,1X,1H',A1,2X,3HLL,4X,A1,3HCLC,11X,A1,3HCLC,8X,1H')	PLOT066
	DIMENSION X0(200),X1(200),X2(200),X3(200),	PLOT007	1400 FORMAT (25HTITLE 0.8 5 SIZE 2 'C2P3'/	PLOT067
	Y0(200),Y1(200),Y2(200),Y3(200)	PLOT008	% 4HCASE,1X,6H' CLC'/	PLOT068
	COMMON /PERT/ M0,M1,M2,Q0,Q1,Q2,YCR0,YCR1,YCR2	PLOT009	% 23HTITLE BOTTOM SIZE 2 'X')	PLOT069
	COMMON /XY/ X0,X1,X2,X3,Y0,Y1,Y2,Y3	PLOT010	1500 FORMAT (47HTITLE 4.5 0.5 CENTER SIZE 1.5 SPACES 17 'M203 =,	PLOT070
	LOGICAL*1 SYM(3) /IHT,1HA,1HM/, SUB(3) /IHG,1HG,1H /,	PLOT011	% F6.3,3X,A1,2H =,F6.3,1H'/	PLOT071
	S0 /IHB/, S1 /IHC/	PLOT012	% 10HCASE ' CSC,11X,1HG,8X,1H')	PLOT072
	REAL MJ,M1,M2	PLOT013	1505 FORMAT (34HTITLE 4.5 0.5 CENTER SIZE 1.5 'M =,F6.3,1H')	PLOT073
	DATA ICALL /0/, XL,XR0,XT0,XR1,XT1,XR2,XT2	PLOT014	1510 FORMAT (6HJOIN 1)	PLOT074
	/0.0, 0.12, 0.14, 0.16, 0.18, 0.2, 0.24/	PLOT015	1530 FORMAT (5HTITLE,1X,2(F6.3,1X),	PLOT075
	IF (ICALL .EQ. 0) WRITE (4,1000)	PLOT016	% 27HDATA SIZE 1.0 'C42P350*1)2,A1,2H3'/	PLOT076
	ICALL=1	PLOT017	% 4HCASE,1X,16H' CCLCCC C CLC')	PLOT077
	CALL LIMITS (N,YMIN,YMAX)	PLOT018	1550 FORMAT (5HTITLE,1X,2(F6.3,1X),25HDATA SIZE 1.5 'C42P350*1'/	PLOT078
	WRITE (4,1050)	PLOT019	% 16HCASE ' CCLCCC C')	PLOT079
	WRITE (4,1100) YMIN,YMAX	PLOT020	1700 FORMAT (15HSET INTENSITY 3/	PLOT080
	WRITE (4,1200)	PLOT021	% 11HJOIN 1 DOTS/	PLOT081

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%          15HSET INTENSITY 2)
1800 FORMAT (13HJOIN 1 DASHES)
1900 FORMAT (4HPLOT)
5000 FORMAT (5(2(F6.3,1X),1H;))
END
SUBROUTINE LIMITS (N,YMIN,YMAX)
C
C.....SEARCHES FOUR DATA ARRAYS Y0,Y1,Y2,Y3 FOR MINIMUM AND MAXIMUM.
C      CALLED ONLY BY PLOT SUBROUTINE.
C
C      DIMENSION X0(200),X1(200),X2(200),X3(200),
%          Y0(200),Y1(200),Y2(200),Y3(200)
C      DIMENSION Z(800)
C      COMMON /XY/ X0,X1,X2,X3,Y0,Y1,Y2,Y3
C      EQUIVALENCE (Y0(1),Z(1))
C      YMIN=Z(1)
C      YMAX=Z(1)
C      DO 10 I=1,4
C      JSTART=200*(I-1)+1
C      JSTOP=JSTART+N-1
C      DO 10 J=JSTART,JSTOP
C      IF (Z(J).GT.YMAX) YMAX=Z(J)
C      IF (Z(J).LT.YMIN) YMIN=Z(J)
10 CONTINUE
C      YSAVE=YMAX
C      YMAX=YMIN
C      YMIN=YSAVE
C      CALL ROUND (YMIN)
C      CALL ROUND (YMAX)
C      RETURN
C      END
C      SUBROUTINE ROUND (Y)
C
C.....ROUNDS Y LIMITS FOR OUTPUT IN F5.1 FORMAT. CALLED ONLY BY
C      SUBROUTINE LIMITS.
C
C      Z=ABS(Y)
C      IF (10.*Z-INT(10.*Z).LT.5) Z=Z+.05
C      IF (Y.GT.0.) GO TO 1
C      Y=-Z
C      RETURN
1 Y=Z
C      RETURN
C      EHD
C      SUBROUTINE UPLOW (A,B,XIN,K,N,XOUT,FLAG)
C
C.....CONVERTS NORMALIZED ARRAY XIN TO PHYSICAL ARRAY XOUT AND FLAGS
C      POINTS ON LOWER SURFACE WITH A '*'.
C
C      DIMENSION XIN(K),XOUT(8)
C      LOGICAL*1 FLAG(8), BLANK/1H /, STAR/1H*/
C      XNOSE=-A/(B-A)
C      DO 1 I=1,N
C      FLAG(I)=BLANK
C      IF (XIN(I) .LT. XNOSE) FLAG(I)=STAR
C      XOUT(I)=ABS((B-A)*XIN(I)+A)
1 CONTINUE
C      RETURN
C      END

```

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|PLOT082
|PLOT083
|PLOT084
|PLOT085
|PLOT086
|LIMI001
|LIMI002
|LIMI003
|LIMI004
|LIMI005
|LIMI006
|LIMI007
|LIMI008
|LIMI009
|LIMI010
|LIMI011
|LIMI012
|LIMI013
|LIMI014
|LIMI015
|LIMI016
|LIMI017
|LIMI018
|LIMI019
|LIMI020
|LIMI021
|LIMI022
|LIMI023
|LIMI024
|LIMI025
|LIMI026
|ROUN001
|ROUN002
|ROUN003
|ROUN004
|ROUN005
|ROUN006
|ROUN007
|ROUN008
|ROUN009
|ROUN010
|ROUN011
|ROUN012
|ROUN013
|UPL0001
|UPL0002
|UPL0003
|UPL0004
|UPL0005
|UPL0006
|UPL0007
|UPL0008
|UPL0009
|UPL0010
|UPL0011
|UPL0012
|UPL0013
|UPL0014
|UPL0015

```

APPENDIX C  
LIST OF SYMBOLS

C	blade chord, m
H	blade spacing for nonstaggered cascades, m
i	invariant point index; eq. (6); also, index for Lagrangian coefficients; eq. (22)
k	dummy index; eq. (20)
L	two-dimensional full potential operator; eq. (2)
$L_1$	linear operator representing first-order perturbation of two-dimensional full potential equation; eq. (4)
$L_2$	linear operator representing first-order perturbation terms arising from coordinate straining; eq. (9)
$L_i$	Lagrangian coefficients; eq. (22)
n	total number of shock points and high-gradient maxima points; eq. (24)
N	total number of invariant points, equal to $n+2$ ; eq. (24)
q	arbitrary geometric or flow parameter to be perturbed; eq. (13)
$q_c$	calibration flow value of q; eq. (9)
$q_o$	base flow value of q; eq. (3)
Q	approximate flow solution for arbitrary flow quantity; eq. (1)
$Q_c$	calibration flow solution for value $q_c$ of arbitrary parameter; eq. (8)
$Q_o$	base flow solution for value $q_o$ of arbitrary parameter; eq. (1)
$Q_p$	linearized perturbation solution per unit change of perturbed parameter; eq. (1)
(s,t)	strained (x,y) coordinates; eq. (5)

- $(x,y)$  nondimensional blade-fixed orthogonal coordinates; eq. (5), normalized by C
- $(\bar{x},\bar{y})$  nondimensional blade-fixed orthogonal coordinates related to calibration solution; eq. (9)
- $(x_1,y_1)$  straining functions associated with  $(x,y)$  coordinates; eq. (6)
- $(x_{1_i},y_{1_i})$  straining functions associated with  $i$ th invariant point; eq. (6)
- $(\delta x_i, \delta y_i)$  unit displacements in  $(x,y)$  directions associated with  $i$ th invariant point; eq. (6)
- $\delta x_i^c$  unit displacement in  $x$  direction between base and calibration flows of the  $i$ th invariant point; eq. (18)
- $\epsilon$  perturbation change of geometric or flow parameter; eq. (17)
- $\epsilon_c$  perturbation of geometric or flow parameter between base and calibration flows; eq. (18)
- $\Phi$  nondimensional total velocity potential; eq. (2), normalized by  $CV_\infty$
- $\Phi_o$  nondimensional base flow velocity potential; eq. (3), normalized by  $CV_\infty$
- $\Phi_1$  nondimensional perturbation velocity potential; eq. (3), normalized by  $CV_\infty$

Subscripts

- o denotes base flow quantities
- 1 denotes perturbation quantities
- c denotes quantities associated with calibration flow

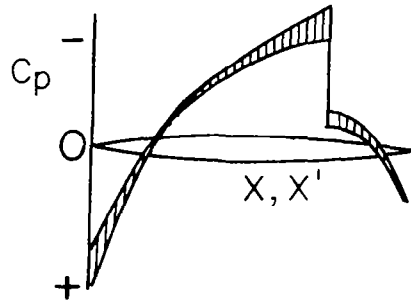
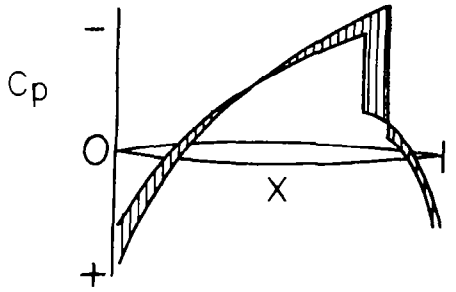
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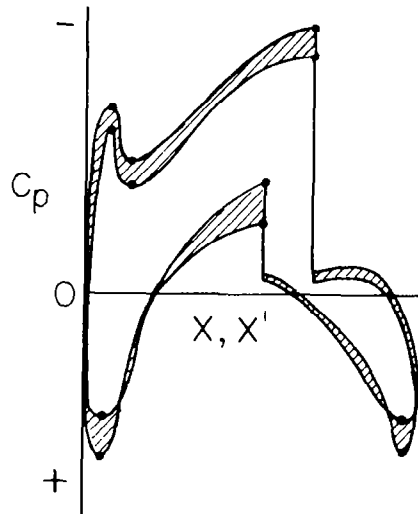
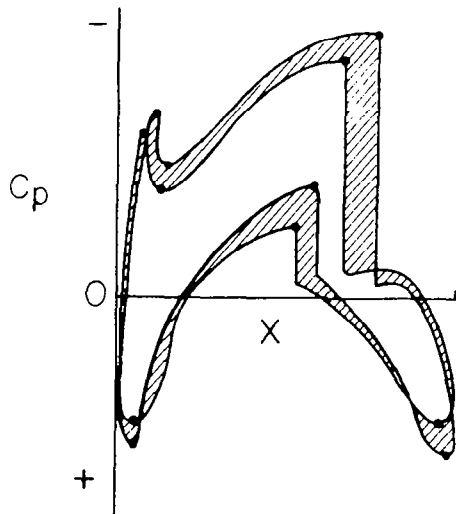


Perturbation for  
calibration solution  
in physical coordinates

Perturbation for  
calibration solution  
in strained coordinates



(a) Single shock.



(b) Multiple shock and high-gradient locations.

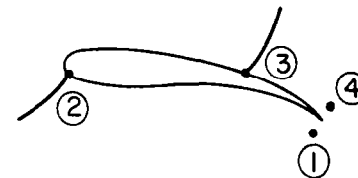
Figure 1.- Illustration of perturbation solution  
for calibration solution in physical and  
strained coordinates

FLOW TYPE:

SUPERCRITICAL  
(SYMMETRIC)

SUBCRITICAL

SUPERCRITICAL



STRAINING  
FUNCTION:

PARABOLIC,  
PIECEWISE CONTINUOUS

PARABOLIC,  
PIECEWISE CONTINUOUS

CUBIC, PIECEWISE  
CONTINUOUS

POINTS HELD  
INVARIANT:

L.E., SHOCK, T.E.

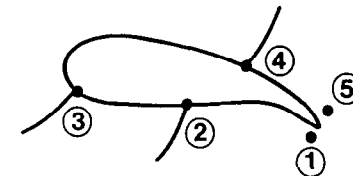
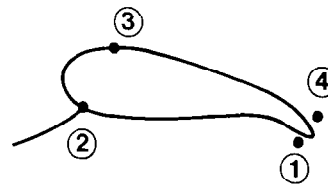
T.E., STAG. PT., T.E.

T.E., STAG. PT.,  
SHOCK PT., T.E.

FLOW TYPE:

SUBCRITICAL

SUPERCRITICAL



STRAINING  
FUNCTION:

CUBIC,  
PIECEWISE CONTINUOUS

QUARTIC,  
PIECEWISE CONTINUOUS

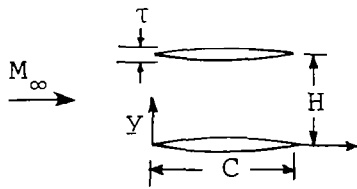
POINTS HELD  
INVARIANT:

T.E., STAG. PT.,  
MAX SUC. PRES. PT., T.E.

T.E., SHOCK PT., STAG. PT.,  
SHOCK PT., T.E.

Figure 2.- Summary of various flows and straining functions considered

BICONVEX PROFILES



- ..... BASE
- CALIBRATION
- oooooo PERTURBATION
- EXACT NONLINEAR

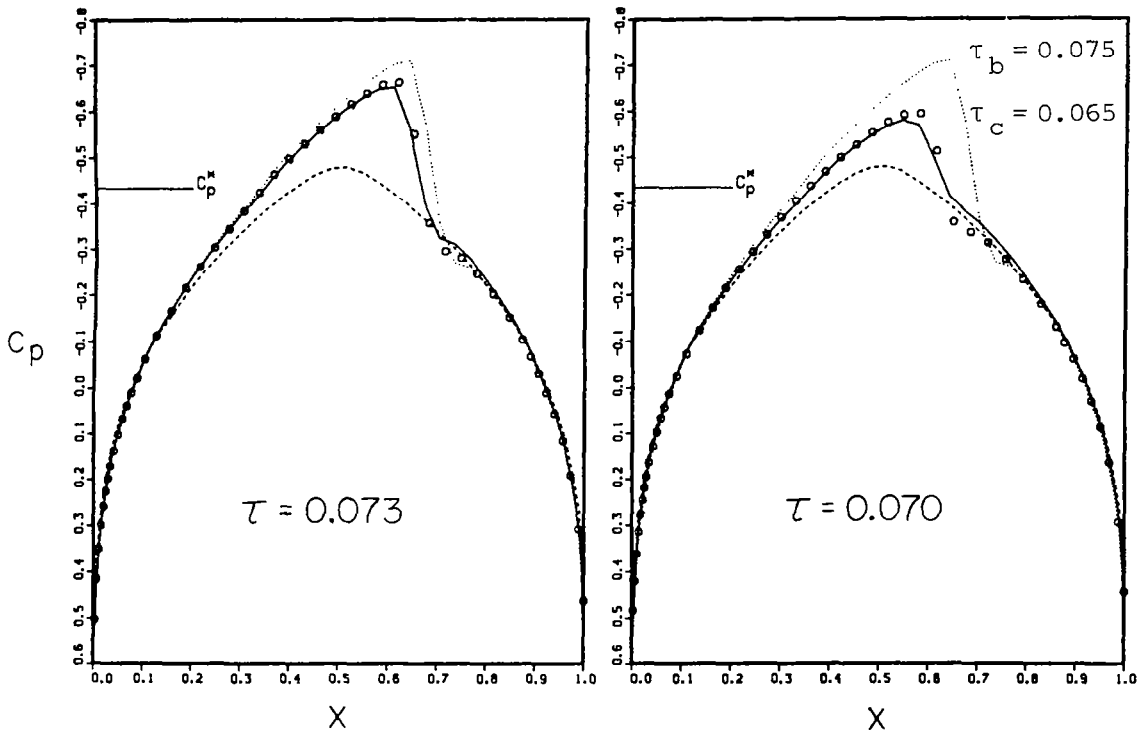


Figure 3.- Comparison of perturbation (O) and non-linear (—) surface pressures for a thickness-ratio perturbation of a nonlifting cascade of biconvex profiles with  $H/C = 1.0$  at  $M_\infty = 0.80$

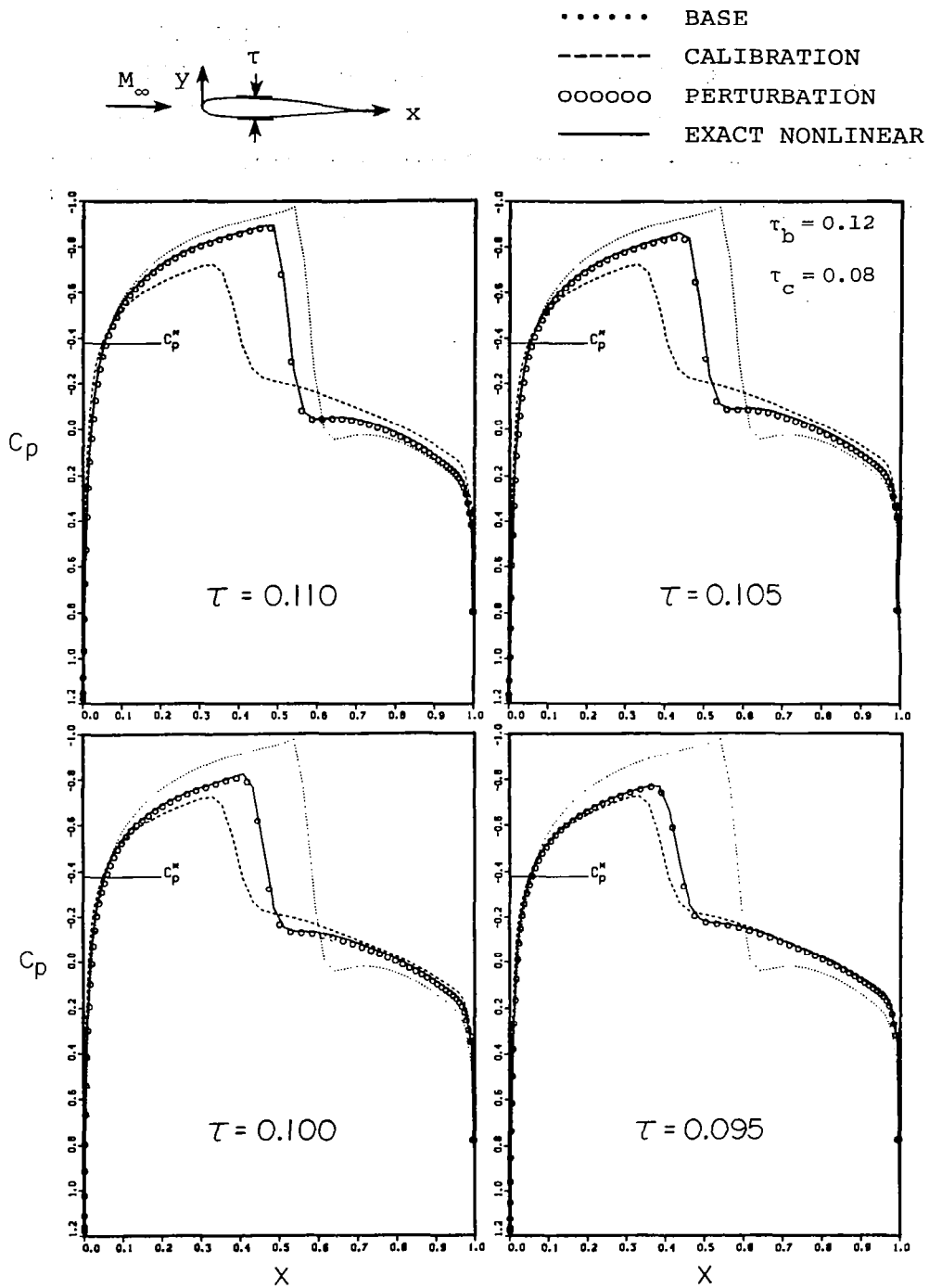


Figure 4.- Comparison of perturbation (o) and non-linear (—) surface pressures for a thickness-ratio perturbation for an isolated NACA 00XX airfoil at  $M_\infty = 0.820$  and  $\alpha = 0^\circ$  for solution interpolation

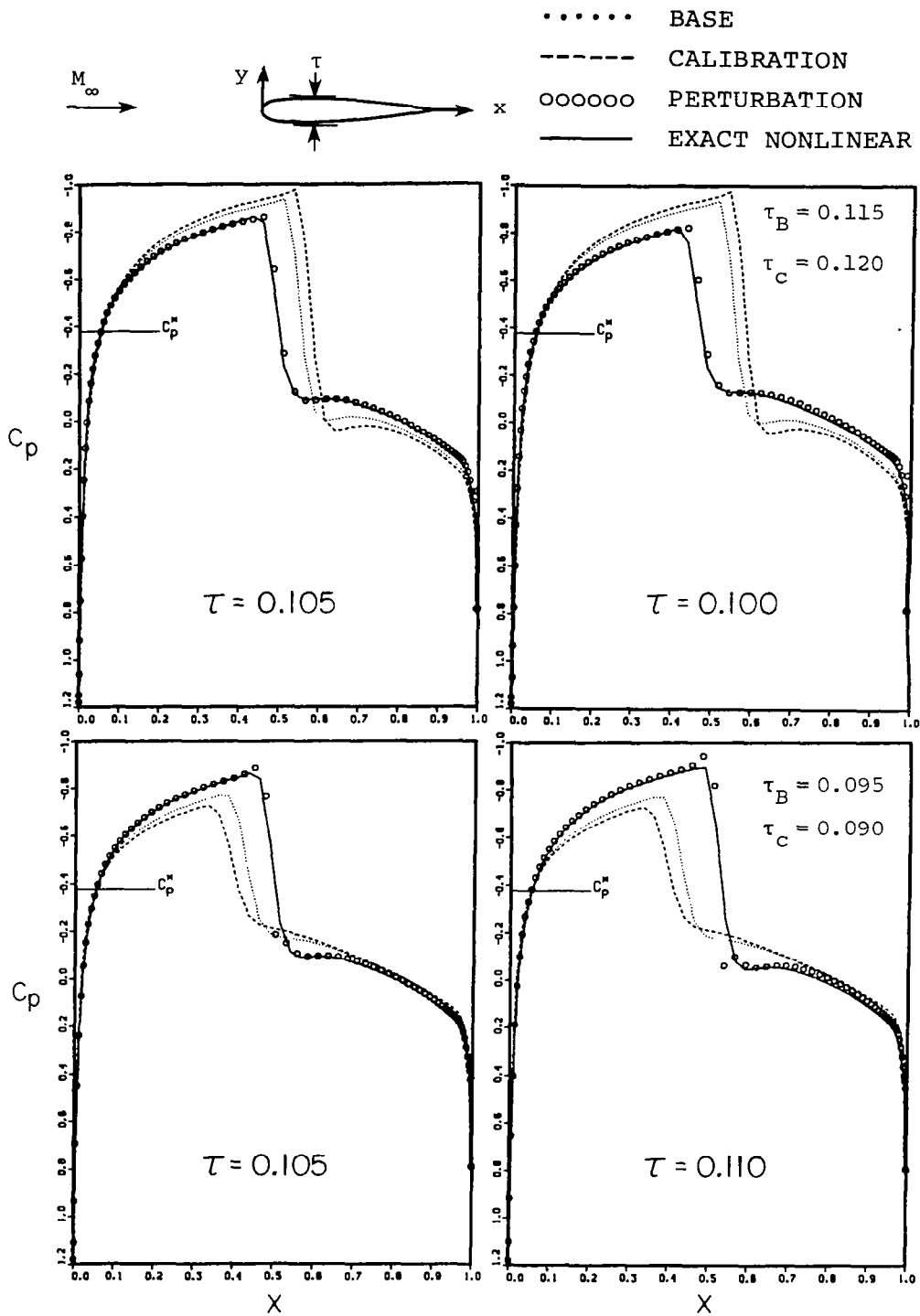


Figure 5.- Comparison of perturbation (o) and non-linear (—) surface pressures for a thickness-ratio perturbation for an isolated NACA 00XX airfoil at  $M_\infty = 0.820$  and  $\alpha = 0^\circ$  for extreme solution extrapolation

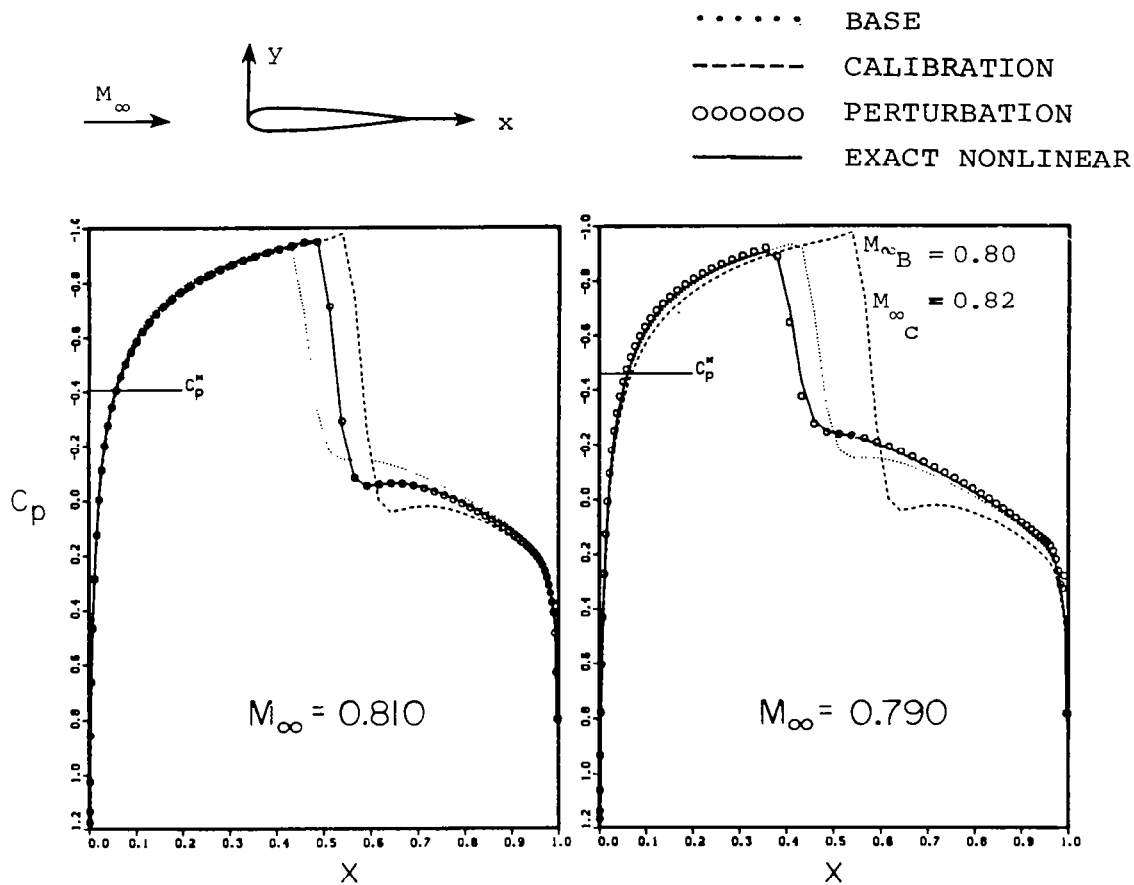


Figure 6.- Comparison of perturbation (○) and nonlinear (—) surface pressures for a Mach number perturbation of an isolated NACA 0012 airfoil at  $\alpha = 0^\circ$

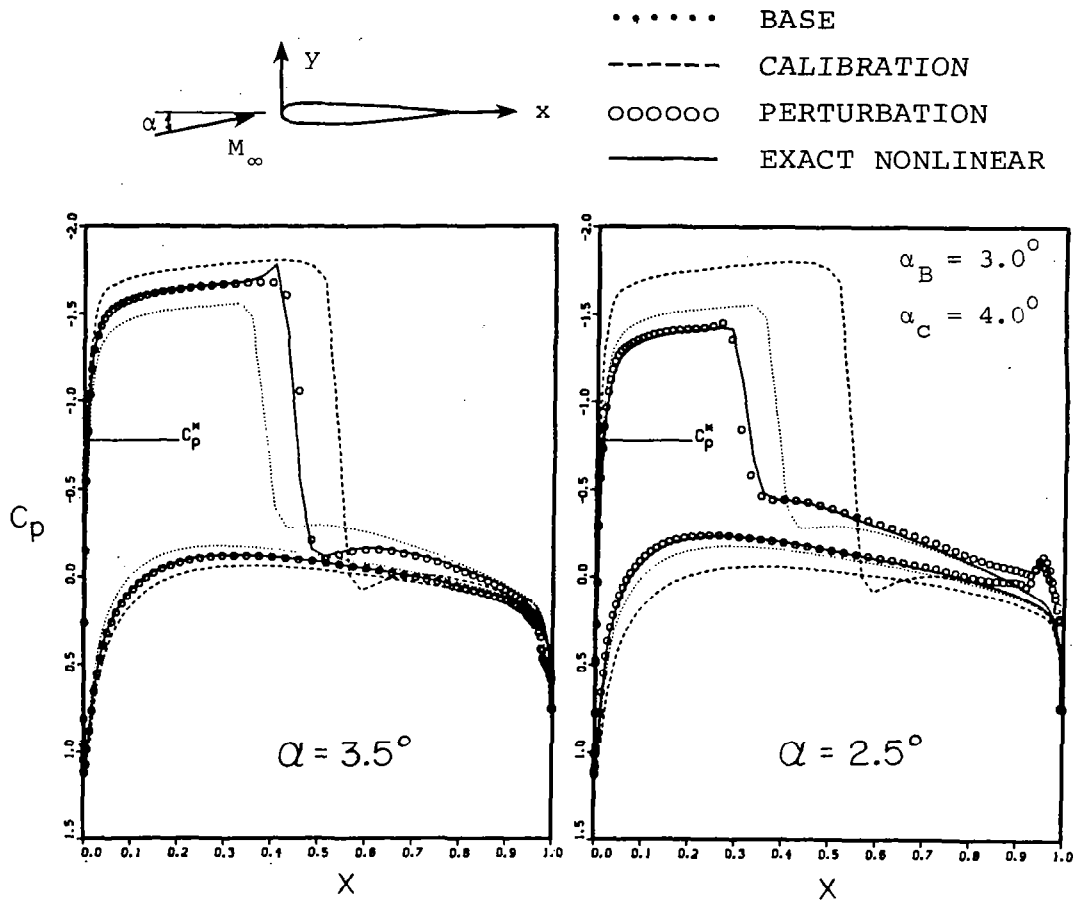


Figure 7.- Comparison of perturbation (O) and nonlinear (—) surface pressures for an angle-of-attack perturbation of an isolated NACA 0012 airfoil at  $M_\infty = 0.70$

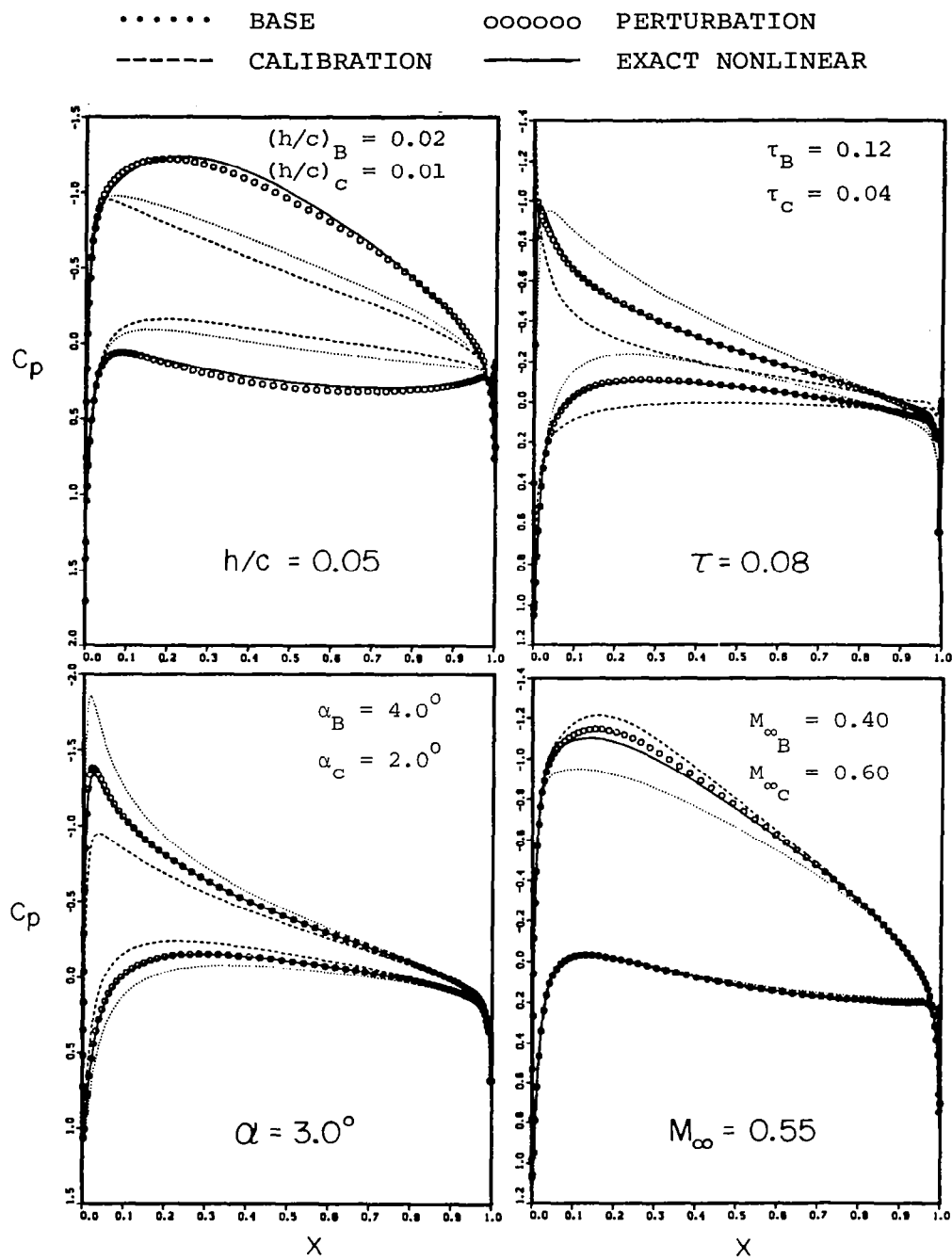


Figure 8.- Comparison of perturbation (o) and nonlinear (—) surface pressures for various geometry and flow parameter perturbations of isolated airfoils at subcritical speeds



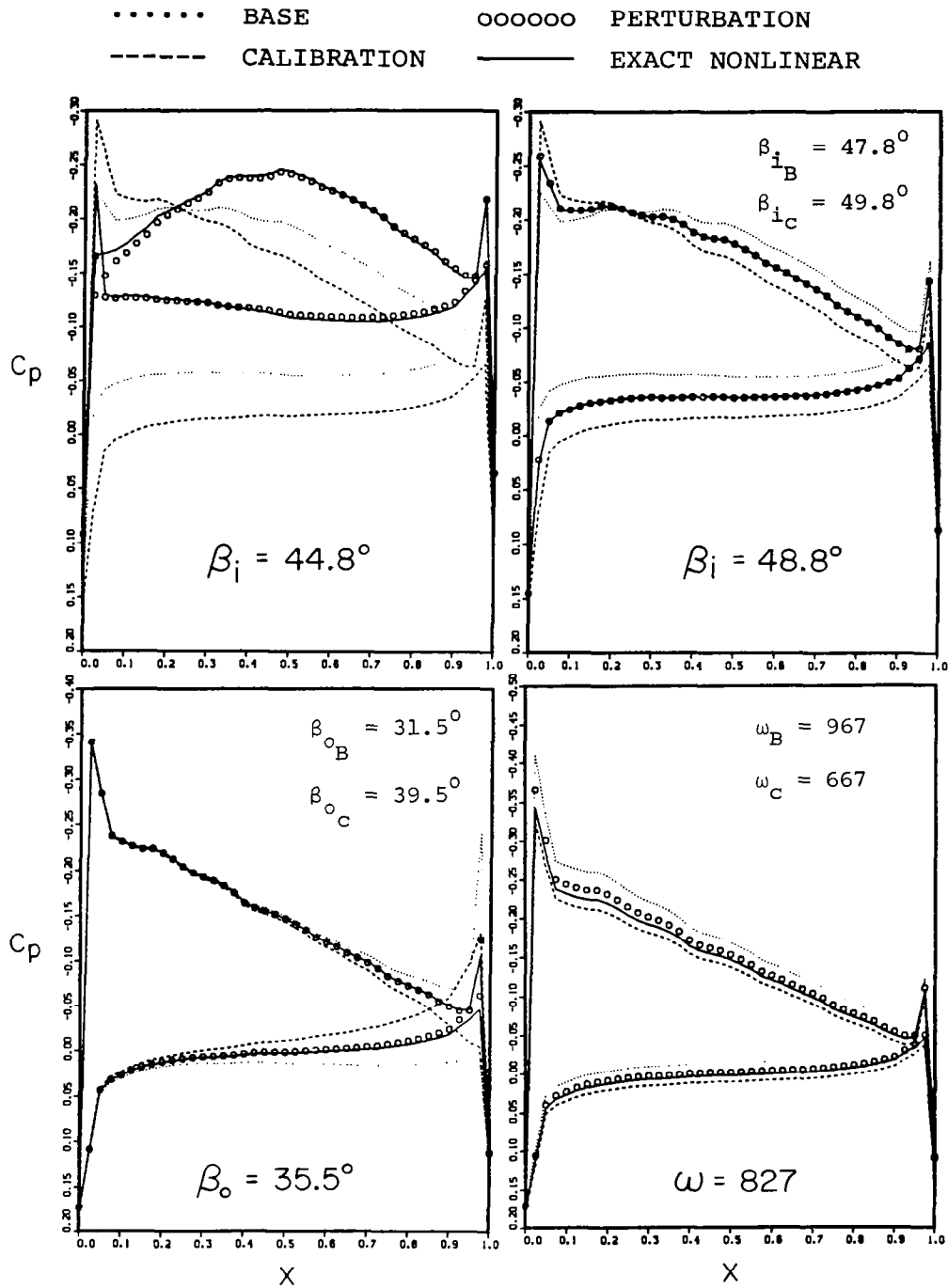


Figure 9.- Comparison of perturbation (o) and nonlinear (—) surface pressures for various flow parameter perturbations of a compressor cascade at subcritical speeds

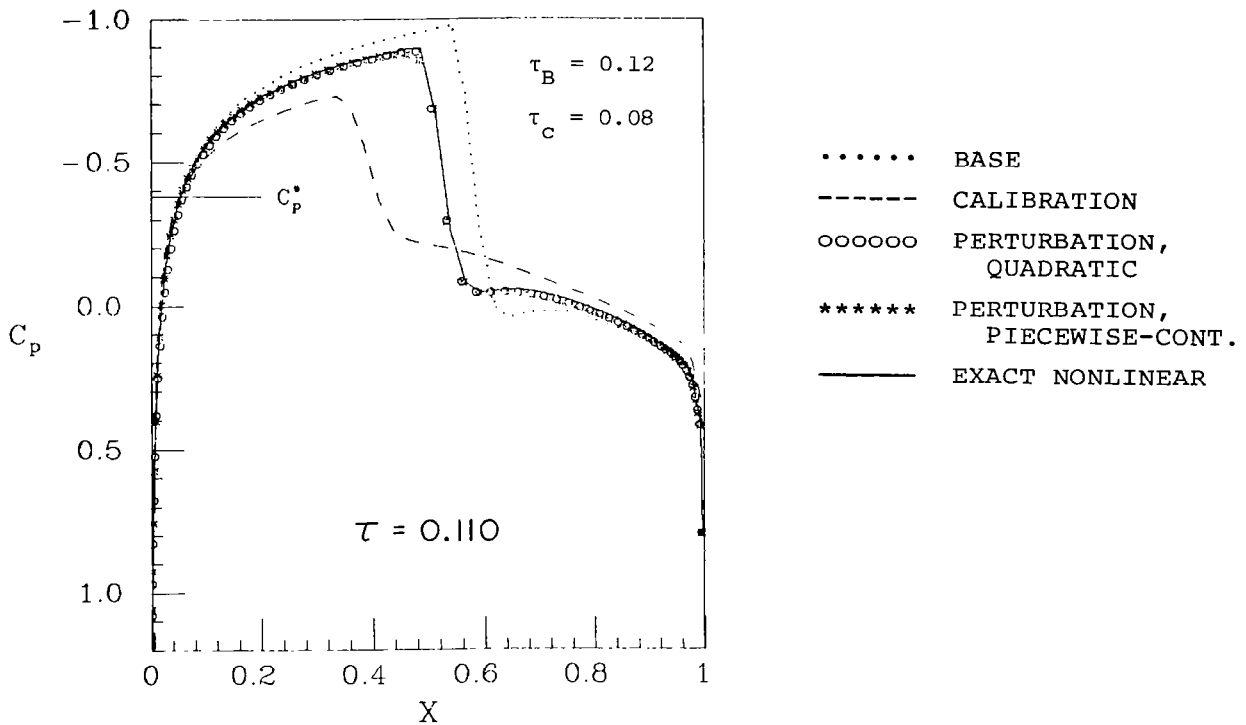


Figure 10.- Comparison of nonlinear (—) surface pressures with perturbation results using quadratic (O) and linear piecewise-continuous (\*) straining functions for a thickness-ratio perturbation of an isolated NACA 00XX airfoil at  $M_\infty = 0.820$  and  $\alpha = 0^\circ$

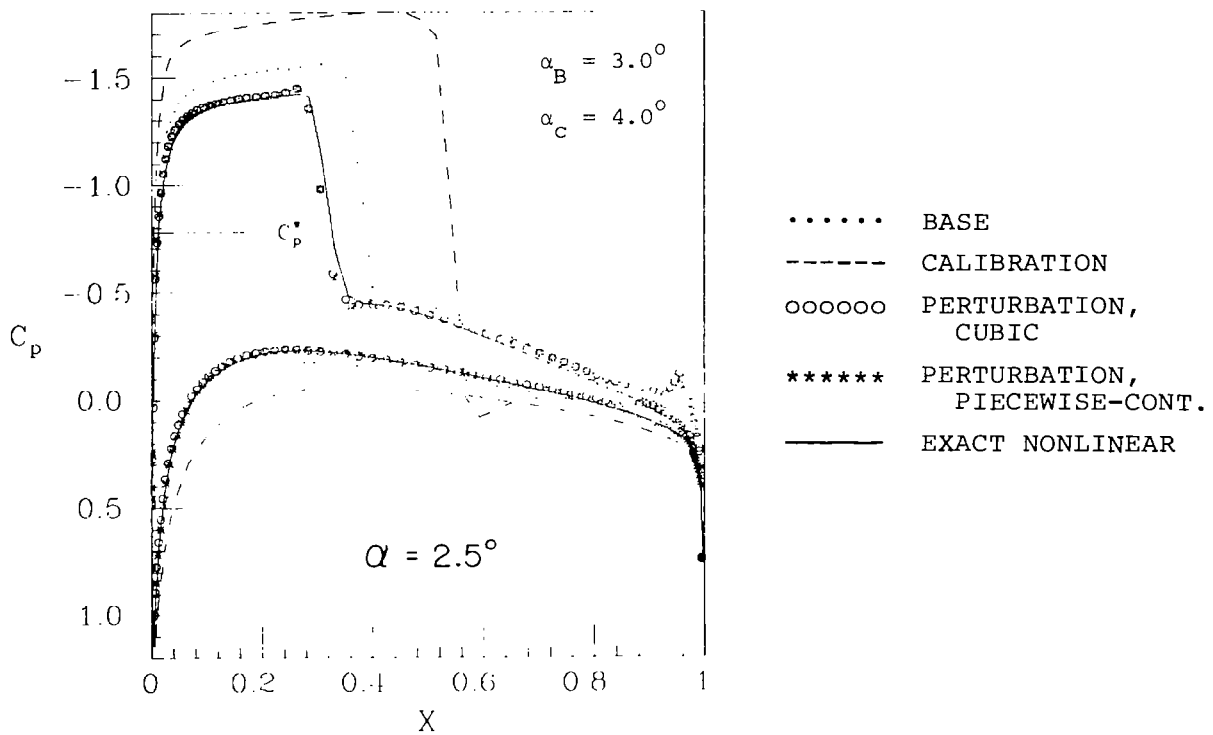


Figure 11.- Comparison of nonlinear (—) surface pressures with perturbation results using cubic (O) and linear piecewise-continuous (\*) straining functions for an angle-of-attack perturbation of an isolated NACA 0012 airfoil at  $M_\infty = 0.70$

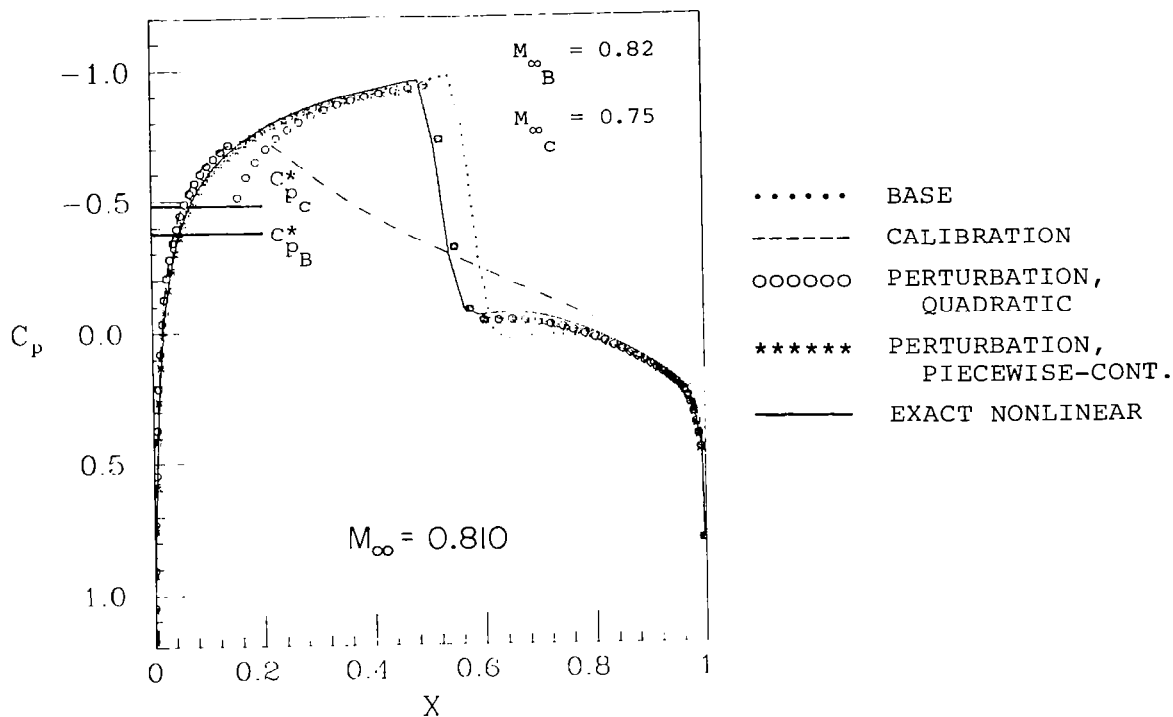


Figure 12.- Comparison of nonlinear (—) surface pressures with perturbation results using quadratic (O) and linear piecewise-continuous (\*) straining functions for a Mach number perturbation of an isolated NACA 0012 airfoil at  $\alpha = 0^\circ$

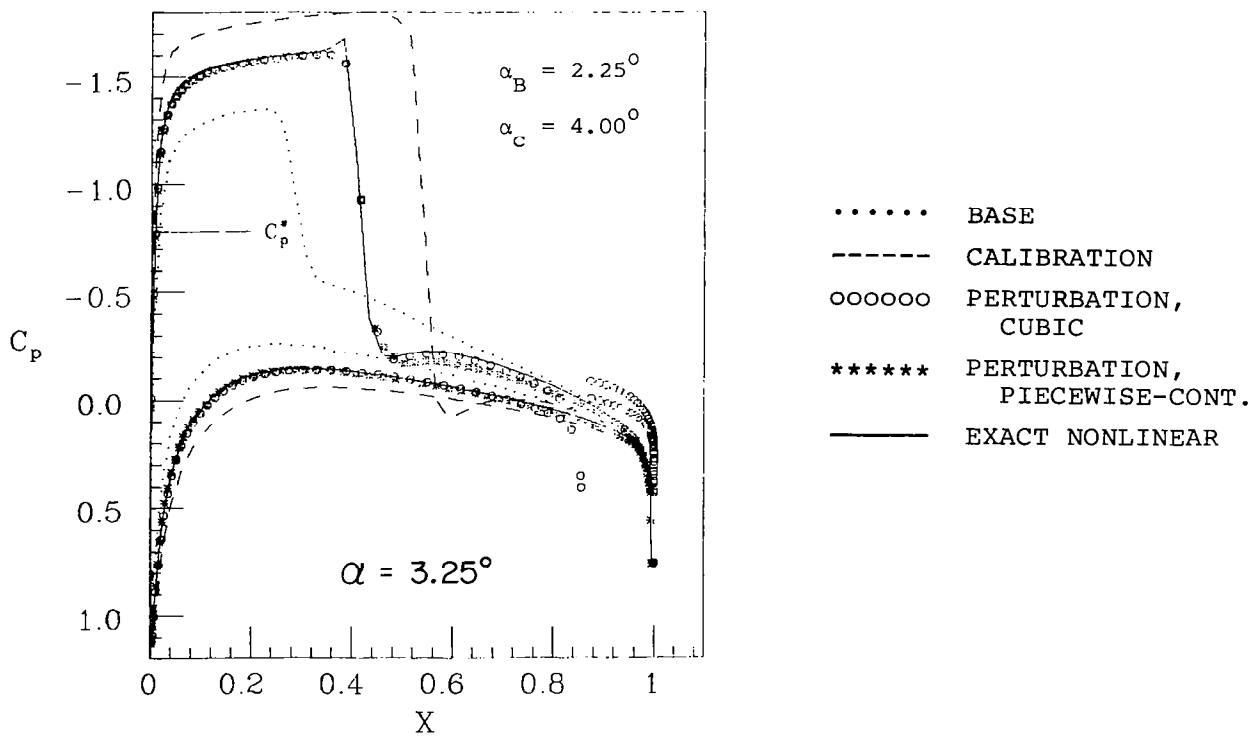
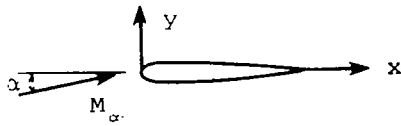


Figure 13.- Comparison of nonlinear (—) surface pressures with perturbation results using cubic (O) and linear piecewise-continuous (\*) straining functions for an angle-of-attack perturbation of an isolated NACA 0012 airfoil at  $M_\infty = 0.70$



- ..... BASE
- CALIBRATION
- \*\*\*\*\* PERTURBATION
- EXACT NONLINEAR

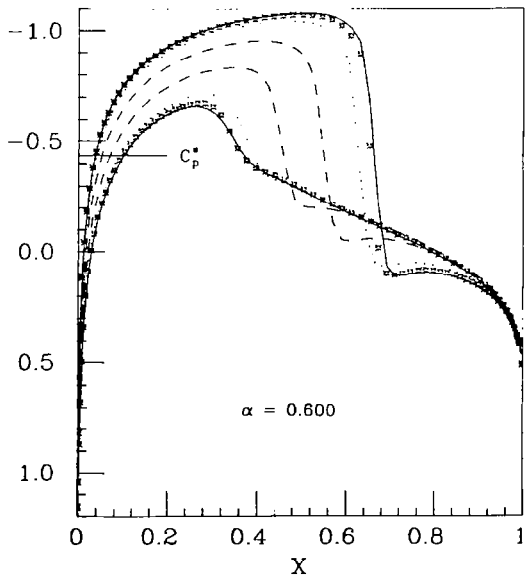
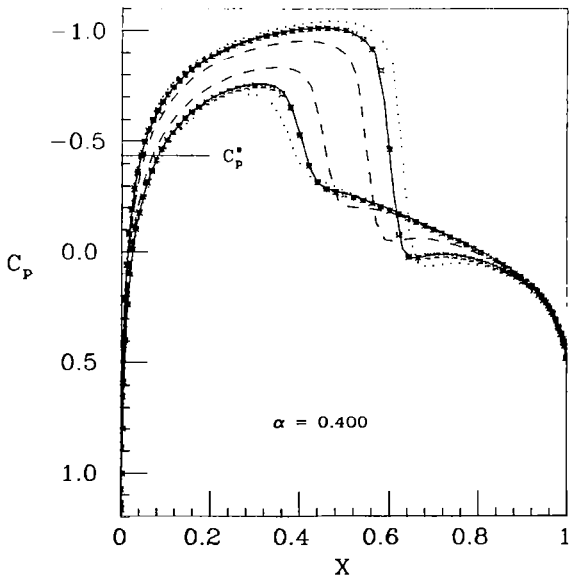
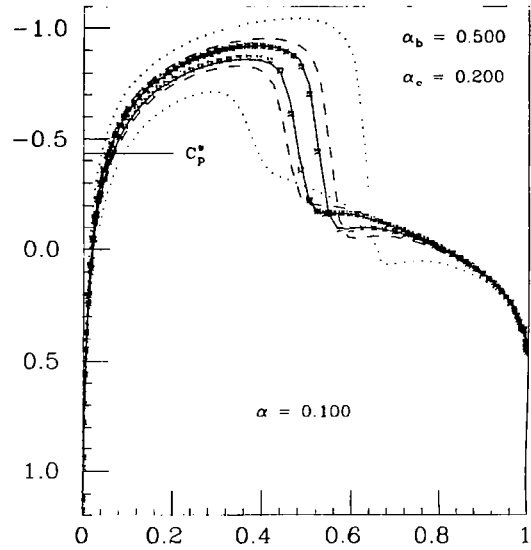
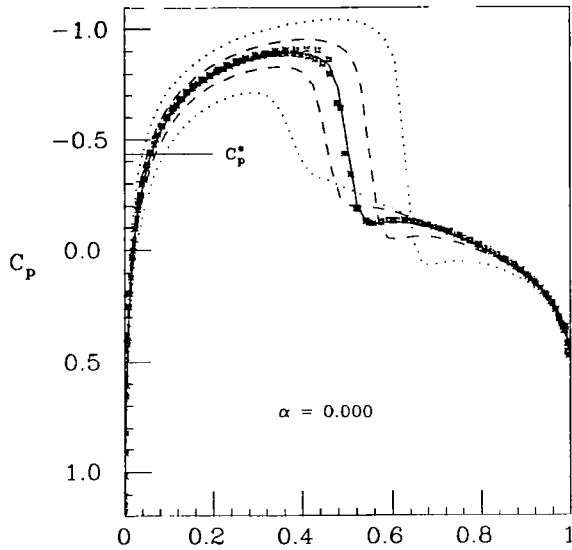


Figure 14.- Comparison of perturbation (\*) and nonlinear (—) surface pressures for an angle-of-attack perturbation of an isolated NACA 0012 airfoil at  $M_\infty = 0.80$  having multiple shocks

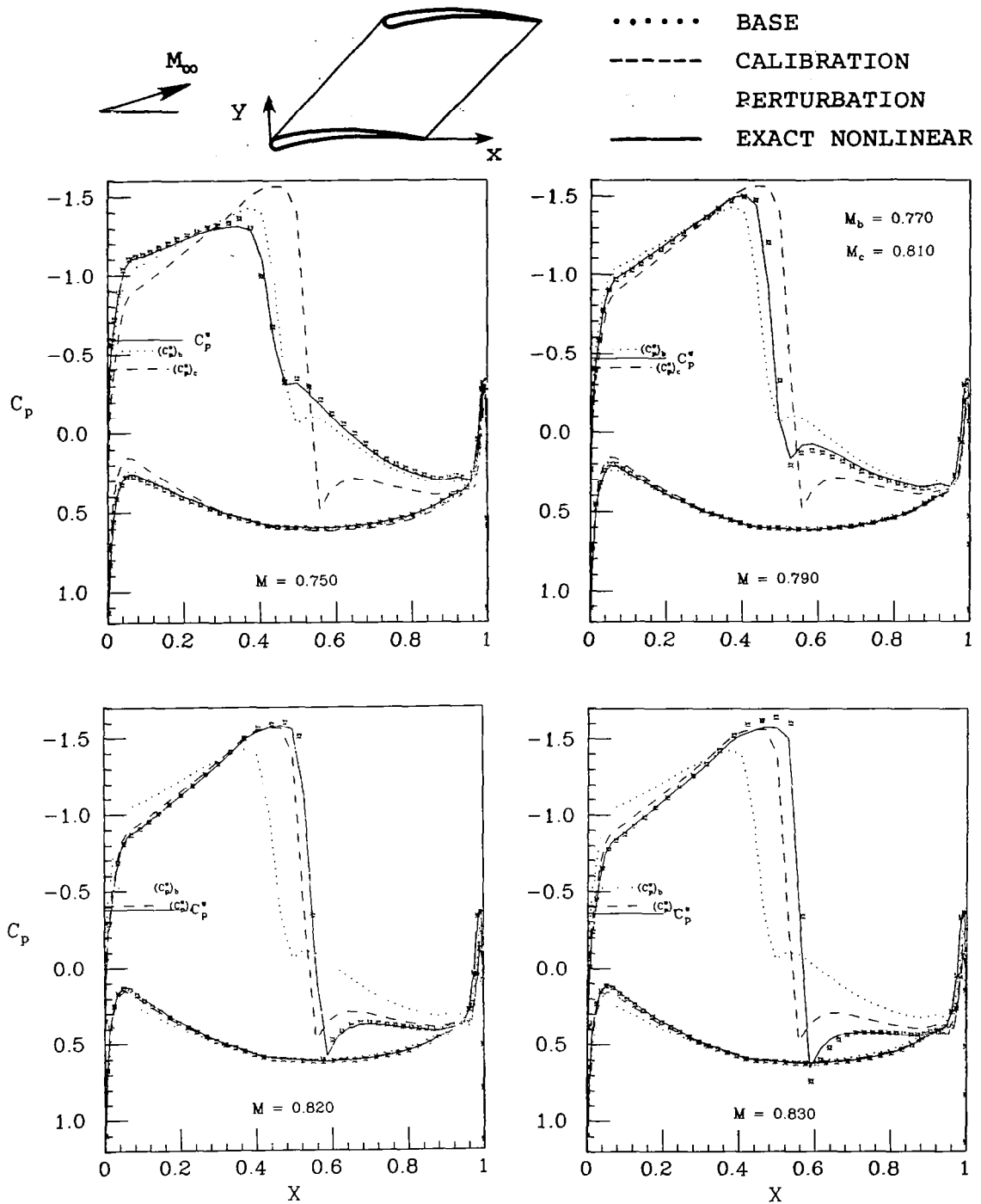


Figure 15.- Comparison of perturbation (\*) and nonlinear (—) surface pressures for an oncoming Mach number perturbation of supercritical flow past a cascade of Jose Sanz blade profiles

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16. Abstract An investigation was conducted to develop perturbation procedures and associated computational codes for determining nonlinear flow solutions, with the objective of establishing a method for minimizing computational requirements associated with parametric studies of transonic flows in turbomachines. The procedure that was developed and evaluated was found to be capable of determining highly accurate approximations to families of strongly nonlinear solutions which are either continuous or discontinuous, and which represent variations in some arbitrary parameter. Coordinate straining is employed to account for the movement of discontinuities and maxima of high-gradient regions due to the perturbation. Although simultaneous multiple-parameter perturbations can be treated, the development and results reported here are for the single-parameter perturbation problem. Flows past both isolated airfoils and compressor cascades involving a wide variety of flow and geometry parameter changes are reported. Attention is focused in particular on transonic flows which are strongly supercritical and exhibit large surface shock movement over the parametric range studied; and on subsonic flows which display large pressure variations in the stagnation and peak suction pressure regions. Comparisons with the corresponding 'exact' nonlinear solutions indicate a remarkable accuracy and range of validity of such a procedure. Computational time of the method, beyond the determination of the base solutions, is trivial.					
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