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Semi-Actuator Disk Theory for Compressor Choke Flutter

J. Micklow and J. Jeffers

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Semi-Actuator Disk Theory for Compressor Choke Flutter

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West Palm Beach, Florida

Prepared for
Lewis Research Center
under Contract NAS3-20060

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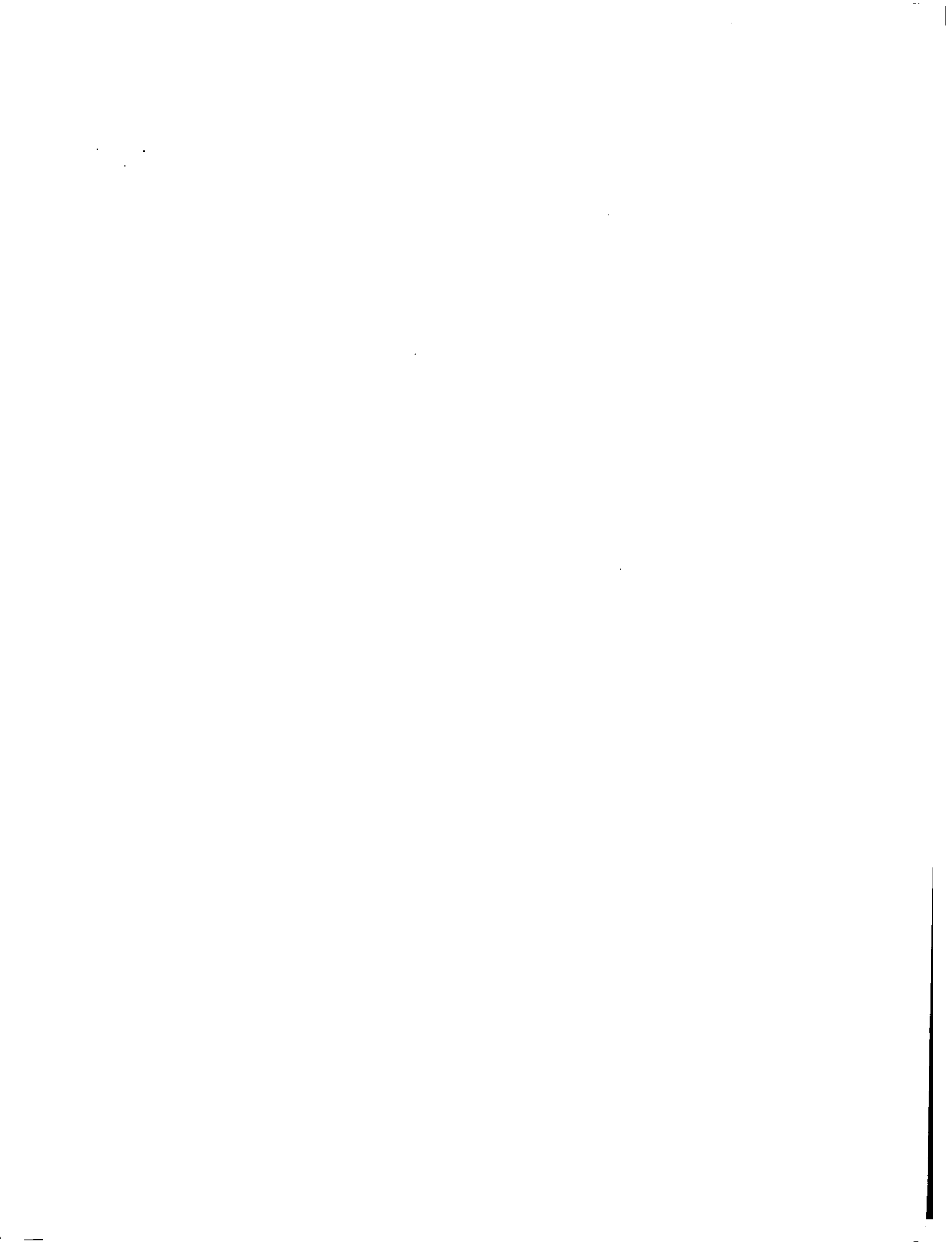
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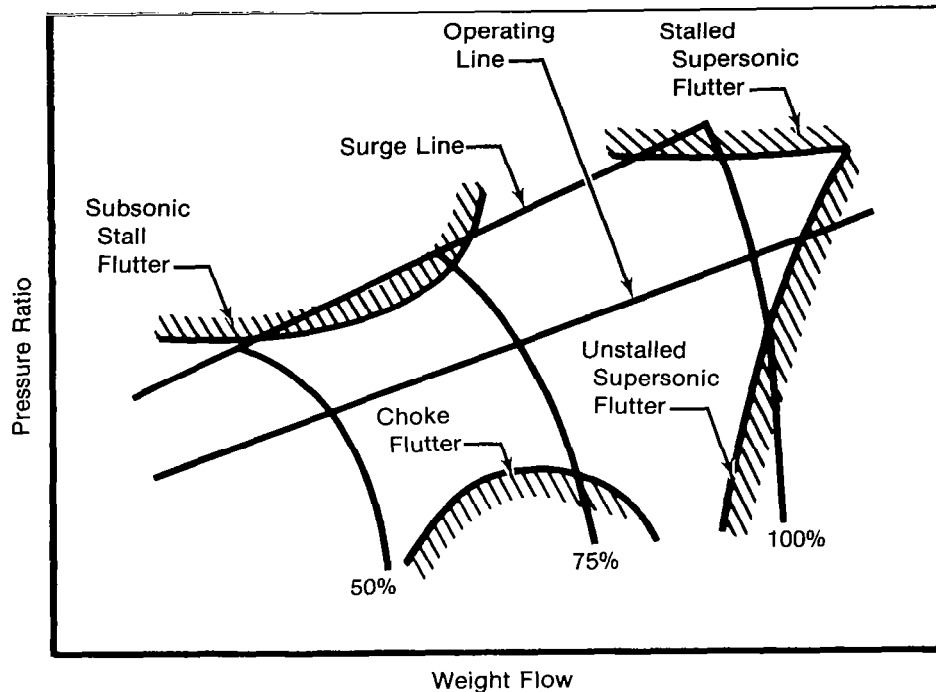
SUMMARY

Utilizing semi-actuator disk theory, a mathematical analysis was developed to predict the unsteady aerodynamic environment for a cascade of airfoils harmonically oscillating in choked flow. In the model, a normal shock is located in the blade passage, its position depending on the time dependent geometry and pressure perturbations of the system. In addition to shock dynamics, the model includes the effect of compressibility, interblade phase lag, and an unsteady flow field upstream and downstream of the cascade.

Calculated unsteady aerodynamic forces using the semiactuator disk model were compared to experimental data from isolated airfoil wind tunnel tests. The wind tunnel data simulate the special cascade condition of 180 deg interblade phase. Agreement between experimental and theory was reasonable. The semiactuator theory was also evaluated using compressor airfoil choke flutter data from single-spool tests of the F100 turbofan engine. The model was incorporated into a flutter prediction program in which calculated aerodynamic damping is correlated to construct flutter onset boundaries. The calculated flutter boundaries compared well with the measured flutter boundaries. Based on these evaluations, it was concluded that a conservative choke flutter design system could be established based on the semiactuator disk model.

INTRODUCTION

Compressor airfoil flutter remains a continuing problem in the design and development of advanced aircraft gas turbine engines. Flutter occurs over a wide range of operating conditions, but can be categorized into four regions: (1) subsonic/transonic stall, (2) subsonic/transonic choke, (3) supersonic unstalled, and (4) supersonic stalled, as shown in Figure 1.



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Figure 1. Possible Flutter Boundaries

The subsonic stall flutter problem has been investigated by a number of authors. Jeffers (Reference 1) devised a semi-empirical unsteady aerodynamic theory based on combining the unsteady unstalled aerodynamic forces from Smith's theory (Reference 2) with correction from theory and experimental data of isolated airfoils operating at high incidence in incompressible flows. Sisto (Reference 3) used steady aerodynamic data to treat the unsteady flow problem in a quasi-steady manner. Perumal (Reference 4) developed an essentially "Helmholtz flow" model, while Yashima and Tanaka (Reference 5) adapted a rigid wake model to obtain reasonably good correlation with linear cascade experimental data. Most recently, Chi (Reference 6) used a small perturbation technique to model flow separation.

The supersonic flow region has also been discussed by a number of authors. For the unstalled regime, a finite difference method was first used by Verdon (Reference 7) and Brix and Platzler (Reference 8) to model the unsteady supersonic aerodynamics. Other approaches by Kurosaka (Reference 9) and Verdon and McCune (Reference 10) extend a velocity potential method first developed by Miles (References 11 and 12) for simple supersonic cascade configurations.

Recently, an unsteady actuator disk model was developed by Adamczyk (Reference 13) with encouraging results for supersonic stall bending flutter. The supersonic stalled region was also investigated by Goldstein, Braun and Adamczyk (Reference 14) in which the small perturbation analysis included the presence of a strong in passage shock.

The choke flutter problem that has arisen in advanced gas turbine engines with variable inlet guide vanes poses a very serious problem and no analytical model exists at present to predict the unsteady aerodynamic environment. The complex nature of this environment has thus far resisted rigorous mathematical formulation, but a "simplified" model has been undertaken herein based on a modified semi-actuator disk approach with one-dimensional channel flow. The channel flow approach originally used by NASA-NACA to analyze inlet diffusers of ramjet and turbojet engines was selected because airfoil cascades can exhibit flow characteristics similar to those of inlet diffusers. The flow in an inlet diffuser and a choked blade passage both contain a shock wave whose position strongly affects the pressure forces on the channel walls or blade surfaces. The position of the shock depends upon channel geometry and, therefore, in the case of the airfoil cascade, can be related to the vibratory motion of the airfoils in flutter. A preliminary analysis was completed in the initial phase of this effort which produced promising results when used in a stability prediction of a compressor rotor that experienced choke flutter at off-schedule operating conditions. However, concern for certain aspects of the preliminary model led to the present approach which includes a modified semi-actuator disk method to describe the upstream and downstream flow fields.

The section following contains the analytical derivation and definition of the mathematical model, including a steady-state interblade analysis and a linearized small perturbation analysis. The next section details the results obtained using this channel flow model. In order to rid the text of this report of complex, cumbersome, and lengthy mathematical manipulations and assumptions, numerous appendices are included herewith which allow model development in a straight forward manner.

ANALYTICAL MODEL

Model Definition

The semi-actuator disk model consists of two solutions: a steady-state intrablade analysis and an unsteady linearized small perturbation analysis. The steady-state analysis utilizes steady isentropic one-dimensional relations to define the intrablade conditions. An iterative procedure, ending in the match of the known static pressure ratio across the blade, locates the steady-state normal shock position. The procedure appears in Appendix H. The flow entering and leaving the cascade is defined externally by a streamline analysis. The unsteady solution consists of three basic flow fields: (1) upstream flow field, (2) intrablade flow analysis, and (3) a downstream flow field, as shown in Figure 2.

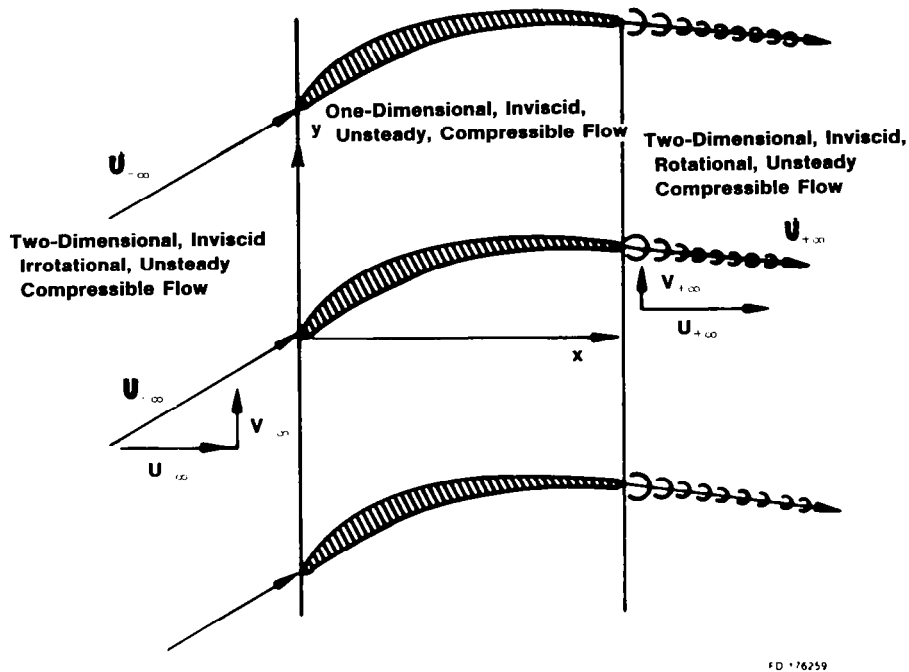


Figure 2. Unsteady Flow Field Description

Assumptions and Boundary Conditions

Assumptions

The following assumptions were made relative to the flow within the three basic flow fields of the unsteady solution:

1. *Upstream Flow Field* — The flow is assumed to be two-dimensional, inviscid, irrotational, unsteady, and compressible.

2. *Intrablade Flow Analysis* — The flow is assumed to be one-dimensional, inviscid, unsteady, and compressible. Figure 3 details the division of the flow field into three sections: a subsonic section from blade leading edge to the blade throat or $M = 1$, a supersonic section from blade throat to shock location, and a subsonic section from shock location to blade trailing edge.

3. *Downstream Flow Field* — The flow is assumed to be two-dimensional, rotational, inviscid, compressible, and constructed of the sum of two basic solutions: an irrotational part similar to the upstream flow field and a rotational part due to the vortices being shed off the blade trailing edge.

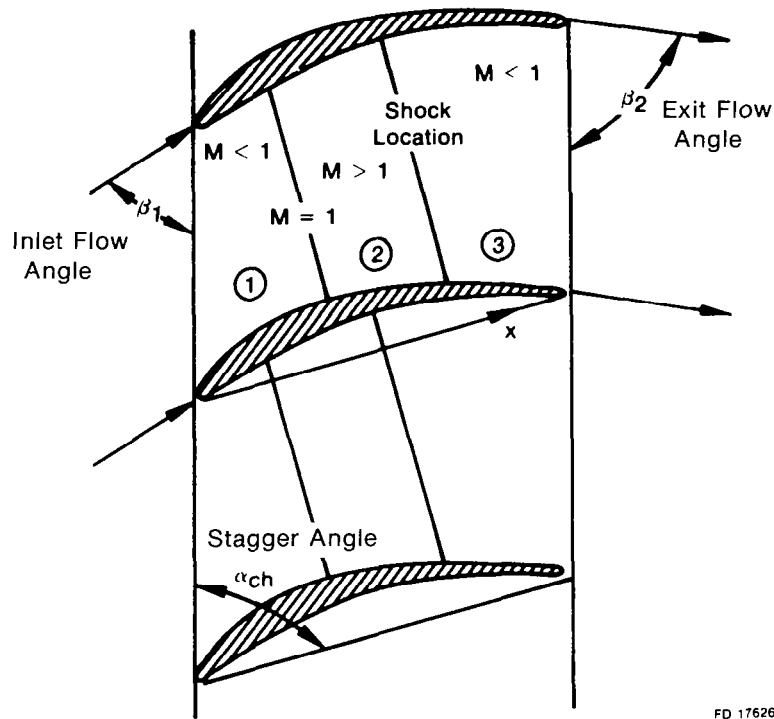


Figure 3. Intrablade Flow Field

Boundary Conditions

Boundary conditions for the unsteady solutions consist of the following:

1. The mass flow is continuous at the leading- and trailing-edge lines.
2. Conservation of mass, energy, and momentum was observed within each section of the blade channel.
3. The Kutta condition at the trailing edge is satisfied by specifying the exit air angle.

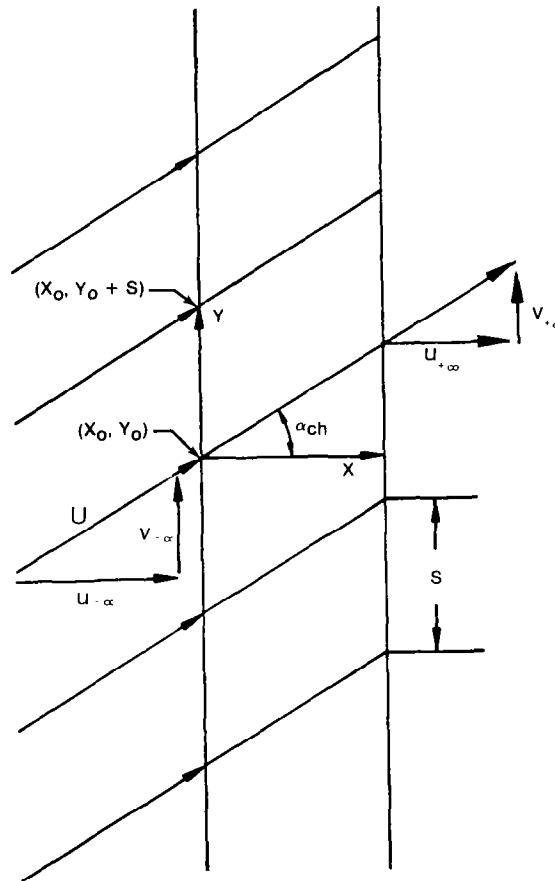
Derivation of the Unsteady Model

Upstream and Downstream Irrotational Flow Equations

Continuity Equation

Establishing a coordinate system, as shown in Figure 4 produces the following form for the continuity equation:

$$\frac{\partial(\rho u A)_{z_\infty}}{\partial x} + \frac{\partial(\rho v A)_{z_\infty}}{\partial y} + \frac{\partial(\rho A)_{z_\infty}}{\partial t} = 0 \quad (1)$$



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Figure 4. Cascade Geometry

Small perturbations of the flow variables are assumed as follows:

$$\rho = \bar{\rho} + \rho' \quad u = \bar{u} + u' \quad v = \bar{v} + v'$$

Substituting the above relationships into Equation 1, subtracting out the steady-state equation, and neglecting higher order terms gives the small perturbation form of the continuity equation, as follows:

$$\frac{\partial \rho'_{\pm\infty}}{\partial t} + \bar{u}_{\pm\infty} \left(\frac{\partial \rho'_{\pm\infty}}{\partial x} \right) + \bar{v}_{\pm\infty} \left(\frac{\partial \rho'_{\pm\infty}}{\partial y} \right) + \bar{p}_{\pm\infty} \left(\frac{\partial u'_{\pm\infty}}{\partial x} + \frac{\partial v'_{\pm\infty}}{\partial y} \right) = 0 \quad (2)$$

Momentum Equation

In tensor notation, the two-dimensional form of the momentum equation is as follows:

$$\frac{\partial (\rho u_i u_j)}{\partial x_j} + \frac{\partial (\rho u_i)}{\partial t} = - \frac{\partial p}{\partial x_i} + \rho f_i$$

which becomes

x-direction

$$\frac{\partial (\rho u_{\pm\infty} u_{\pm\infty})}{\partial x} + \frac{\partial (\rho u_{\pm\infty} v_{\pm\infty})}{\partial y} + \frac{\partial (\rho u_{\pm\infty})}{\partial t} = - \frac{\partial p_{\pm\infty}}{\partial x} \quad (3a)$$

y-direction

$$\frac{\partial (\rho u_{\pm\infty} v_{\pm\infty})}{\partial x} + \frac{\partial (\rho v_{\pm\infty} v_{\pm\infty})}{\partial y} + \frac{\partial (\rho v_{\pm\infty})}{\partial t} = - \frac{\partial p_{\pm\infty}}{\partial y} \quad (3b)$$

Assuming small perturbations of the flow variables in Equations 3a and 3b, neglecting higher order terms and subtracting out the mean flow equation, gives the small perturbation forms of the momentum equation, as shown below:

x-direction

$$\frac{\partial u'_{\pm\infty}}{\partial t} + \bar{u}_{\pm\infty} \frac{\partial u'_{\pm\infty}}{\partial x} + \bar{v}_{\pm\infty} \frac{\partial u'_{\pm\infty}}{\partial y} = \frac{-1}{\bar{\rho}_{\pm\infty}} \left(\frac{\partial p'_{\pm\infty}}{\partial x} \right) \quad (4a)$$

y-direction

$$\frac{\partial v'_{\pm\infty}}{\partial t} + \bar{u}_{\pm\infty} \frac{\partial v'_{\pm\infty}}{\partial x} + \bar{v}_{\pm\infty} \frac{\partial v'_{\pm\infty}}{\partial y} = \frac{-1}{\bar{\rho}_{\pm\infty}} \left(\frac{\partial p'_{\pm\infty}}{\partial y} \right) \quad (4b)$$

Nondimensionalized Wave Equation

Flow entering the cascade is assumed to be irrotational and can therefore be represented by a potential function as can the irrotational portion of the flow leaving the cascade. The velocity perturbations can then be represented in the following forms:

$$u'_{\pm\infty} = \frac{\partial\Phi'_{\pm\infty}}{\partial x} \quad v'_{\pm\infty} = \frac{\partial\Phi'_{\pm\infty}}{\partial y} \quad (5)$$

Substituting these relationships into the continuity and momentum equations and combining these equations gives the small perturbation form of the nondimensionalized wave equation (after some manipulation), as shown:

$$\begin{aligned} \bar{M}_{\pm\infty}^2 \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial t^{*2}} \right) + 2\bar{M}_{\pm\infty}\bar{M}_{x\pm\infty} \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial x^*\partial t^*} \right) + 2\bar{M}_{\pm\infty}\bar{M}_{y\pm\infty} \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial y^*\partial t^*} \right) + \\ 2\bar{M}_{x\pm\infty}\bar{M}_{y\pm\infty} \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial x^*\partial y^*} \right) + (\bar{M}_{x\pm\infty}^2 - 1) \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial x^{*2}} \right) + \\ (\bar{M}_{y\pm\infty}^2 - 1) \left(\frac{\partial^2\Phi'_{\pm\infty}}{\partial y^{*2}} \right) = 0 \end{aligned} \quad (6)$$

where x and y are nondimensionalized by semichord b , and time is nondimensionalized by the quantity U/b . The derivation of the wave equation appears in Reference 15.

Because Equation 6 is linear, a solution can be obtained by superposition of fundamental solutions, taking the following form:

$$\Phi'_{\pm\infty} = A_{\pm\infty} \exp i(Bx + Cy + kt) \quad (7)$$

where A , B , and C are unknown constants and k represents the reduced frequency based on semichord $k = b\omega/U$. Assuming the blades vibrate with a constant interblade phase angle σ , the tangential wave constant C is controlled by an unsteady periodicity condition. Any perturbation velocity potential at $(x_0, y_0 + s)$ leads or lags the same potential at (x_0, y_0) by σ at all times. This may be expressed as follows:

$$\Phi'(x_0, y_0 + s, t) = \Phi'(x_0, y_0, t) e^{i\sigma} \quad (8)$$

where s defines the blade gap-to-semichord ratio. Substituting Equation 8 into Equation 7 gives

$$C = \frac{\sigma}{s}$$

Substituting Equation 7 into Equation 6, dividing by $A_{\pm\infty} \exp i(Bx + Cy + kt)$ and solving for B yields

$$B_{1,2} = \frac{-D_1 \pm \sqrt{D_1^2 + \beta_x^2 D_2}}{-\beta_x^2} \quad (9)$$

where,

$$\begin{aligned}
D_1 &= \bar{M}_{x \pm \infty} \bar{M}_{\pm \infty} k + \bar{M}_{x \pm \infty} \bar{M}_{y \pm \infty} C \\
D_2 &= k^2 \bar{M}_{\pm \infty}^2 + 2 \bar{M}_{\pm \infty} \bar{M}_{y \pm \infty} C k - C^2 \beta_y^2 \\
\beta_x^2 &= 1 - \bar{M}_x^2 \\
\beta_y^2 &= 1 - \bar{M}_y^2
\end{aligned}$$

Thus, the solution takes the form:

$$\Phi'_{\pm \infty}(x, y, t) = [A_{1 \pm \infty} e^{iB_1 x} + A_{2 \pm \infty} e^{iB_2 x}] [e^{i(Cy + kt)}] \quad (10)$$

The velocity perturbation must approach zero in the far field. Now if the quantity $D_1^2 + \beta_x^2 D_2$ < 0 , B_1 and B_2 are complex conjugates. In order to satisfy the far field condition, the solution must take the following form:

$$\Phi'_{\pm \infty}(x, y, t) = A_{\pm \infty} \exp(iB_1 x) \exp i(Cy + kt) \quad (11)$$

Now consider the case where $D_1^2 + \beta_x^2 D_2 > 0$. Here B_1 and B_2 will be real numbers and the boundedness condition cannot be applied in the far field. This problem can be solved by representing the flow field as containing a transient part. The assumed solution becomes form

$$\Phi'_{\pm \infty}(x, y, t) = A_{\pm \infty} \exp i[(Bx + Cy + kt) + k_1 t] \quad (12)$$

Substituting Equation 12 into Equation 6 and dividing by $A_{\pm \infty} \exp[i(Bx + Cy + kt) + K_1 t]$ gives

$$B_{1, 2 \text{ real}} = \frac{D_1^2 \pm \sqrt{D_1^2 + \beta_x^2 D_2^2}}{\beta_x^2} \quad (13)$$

and

$$B_{\text{imag}} = \frac{-i(\bar{M}_{y \pm \infty} C + \bar{M}_{\pm \infty} k)}{\bar{M}_{x \pm \infty}} \quad (14)$$

Equation 12 physically implies an unsteady solution under the influence of a slowly divergent, oscillating source with the earlier emitted wave being stronger than the one following it. Since a finite period of time is required for a disturbance to propagate to the far field, the amplitude of the response should decrease with increasing distance from the source. Therefore, the proper upstream solution can be chosen. By letting k , approach zero, the correct wave solutions on either side of the harmonically vibrating source can be recovered, as shown in Equation 15.

$$\Phi'_{\pm\infty} = A_{\pm\infty} \exp i[x (B_1 + B_{imag}) + Cy + kt] \quad (15)$$

A summary of the correct solutions for the irrotational upstream and downstream flow fields is presented in Table 1.

Intrablade Flow Equations and Solutions

Perturbation Equations and Solutions for Region 1

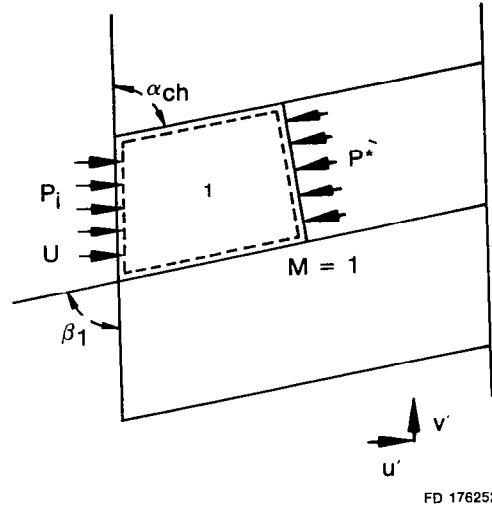


Figure 5. Region 1 from Blade Leading-Edge Line to the Channel Throat, or $M = 1$

The first region, as shown in Figure 5, defines a control volume from the blade leading-edge line to the channel throat, or $M = 1$, with the coordinate system aligned with the chord line. In order to define the unsteady flow field, three unknowns must be obtained: (1) the complex constant describing the upstream flow field, (2) the density perturbation at the throat, and (3) the velocity perturbation at the throat. This requires the use of the unsteady form of the mass, momentum, and energy equations.

Equation 16 expresses the conservation of mass for a control volume (Reference 16):

$$\frac{dm}{dt} = 0 = \iint_{cs} \rho U dA + \frac{\partial}{\partial t} \iiint_{vol} \rho \partial (vol) \quad (16)$$

For the control volume shown in the above figure, this can be expressed as follows:

$$\dot{w}^* - \dot{w}_{inlet} = - \frac{\partial(\rho_1 V_1)}{\partial t} \quad (17)$$

Assuming small perturbations on the mean flow variables, neglecting higher order terms and subtracting out the mean flow equation, the following result is obtained after expanding Equation 17:

TABLE 1. IRROTATIONAL FLOW FIELD SOLUTIONS

Upstream Flow Fields		
$D_{1-\infty}^2 + \beta_{x-\infty}^2 D_{2-\infty} < 0$	$D_{1-\infty}^2 + \beta_{x-\infty}^2 D_{2-\infty} > 0$	
$B_{1R} = D_{1-\infty}/\beta_{x-\infty}^2$ $B_{1I} = -\sqrt{D_{1-\infty}^2 + \beta_{x-\infty}^2 D_{2-\infty}}/\beta_{x+\infty}^2$	$\bar{M}_{-\infty} k + \bar{M}_{y-\infty} C > 0$	$\bar{M}_{-\infty} k + \bar{M}_{y-\infty} C < 0$
	$B_{1R} = \frac{D_{1-\infty} + \sqrt{D_{1-\infty}^2 + \beta_{x-\infty}^2 D_{2-\infty}}}{\beta_{x-\infty}^2}$ $B_{1I} = \frac{-(\bar{M}_{y-\infty} C + \bar{M}_{-\infty} k)}{\bar{M}_{x-\infty}}$	$B_{1R} = \frac{D_{1-\infty} - \sqrt{D_{1-\infty}^2 + \beta_{x-\infty}^2 D_{2-\infty}}}{\beta_{x-\infty}^2}$ $B_{1I} = \frac{+(\bar{M}_{y-\infty} C + \bar{M}_{-\infty} k)}{\bar{M}_{x-\infty}}$
Downstream Flow Fields		
$D_{1+\infty}^2 + \beta_{x+\infty}^2 D_{2+\infty} < 0$	$D_{1+\infty}^2 + \beta_{x+\infty}^2 D_{2+\infty} > 0$	
$B_{1R} = D_{1+\infty}/\beta_{x+\infty}^2$ $B_{1I} = +\sqrt{D_{1+\infty}^2 + \beta_{x+\infty}^2 D_{2+\infty}}/\beta_{x+\infty}^2$	$\bar{M}_{+\infty} k + \bar{M}_{y+\infty} C > 0$	$\bar{M}_{+\infty} k + \bar{M}_{y+\infty} C < 0$
	$B_{2R} = \frac{D_{1+\infty} - \sqrt{D_{1+\infty}^2 + \beta_{x+\infty}^2 D_{2+\infty}}}{\beta_{x+\infty}^2}$ $B_{2I} = \frac{+(\bar{M}_{y+\infty} C + \bar{M}_{+\infty} k)}{\bar{M}_{x+\infty}}$	$B_{2R} = \frac{D_{1+\infty} + \sqrt{D_{1+\infty}^2 + \beta_{x+\infty}^2 D_{2+\infty}}}{\beta_{x+\infty}^2}$ $B_{2I} = \frac{-(\bar{M}_{y+\infty} C + \bar{M}_{+\infty} k)}{\bar{M}_{x+\infty}}$

where,

$$D_{1\pm\infty} = (\bar{M}_{x\pm\infty} \bar{M}_{\pm\infty} k + \bar{M}_{y\pm\infty} \bar{M}_{\pm\infty} C)$$

$$\beta_{x\pm\infty}^2 = 1 - \bar{M}_{x\pm\infty}^2$$

$$D_{2\pm\infty} = (k^2 \bar{M}_{\pm\infty}^2 + 2 \bar{M}_{\pm\infty} \bar{M}_{y\pm\infty} C k - \beta_{y\pm\infty}^2 C^2)$$

$$\beta_{y\pm\infty}^2 = 1 - \bar{M}_{y\pm\infty}^2$$

$$\begin{aligned}
& [\rho'^* \bar{A}^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + \bar{\rho}^* \bar{A}^* a'^*] - [\bar{\rho}_1 \bar{A}_1 (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& \quad \bar{\rho}_1 A'_1 \bar{U}_1 \cos (\alpha_{ch} - \beta_1) + \rho'_1 \bar{A}_1 \bar{U}_1 \cos (\alpha_{ch} - \beta_1)] \\
& = - \left(\bar{\rho}_1 \frac{\partial V'_1}{\partial t} + \nabla'_1 \frac{\partial \rho'_1}{\partial t} \right) \quad (18)
\end{aligned}$$

The derivation of the small perturbation forms of the mass, momentum, and energy equations is shown in Appendix A.

The next step involves the derivation of the unsteady control volume form of the momentum equation. The equation takes the form:

$$\mathbf{F} = \frac{d(m\mathbf{U})}{dt} = \iint_{cs} \mathbf{U}(\rho\mathbf{U})dA + \frac{\partial}{\partial t} \iiint_{vol} \mathbf{U}\rho d(\text{Vol}) \quad (19)$$

Expanding Equation 19 and assuming small perturbations yields for Region 1:

$$\begin{aligned}
& \rho^* \bar{A}_1 + \bar{\rho} A'_1 - \bar{\rho}^* A'^* - \rho^* \bar{A}^* = [a'^* \bar{w}^* + \bar{a}^* w'^*] - \\
& \quad [\bar{w}_1 (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \bar{U}_1 w'_1 \cos (\alpha_{ch} - \beta_1)] + \\
& \quad U_{ave} \bar{\nabla}_1 \frac{\partial \rho'_{ave}}{\partial t} + \bar{\rho}_{ave} \bar{U}_{ave} \frac{\partial V'_1}{\partial t} + \rho_{ave} \bar{\nabla}_1 \frac{\partial U'_{ave}}{\partial t} \quad (20)
\end{aligned}$$

The next step requires the derivation of the small perturbation form of the energy equation for an inviscid fluid with no external heat addition. The control volume for Region 1 is therefore (Reference 16):

$$\frac{\partial}{\partial t} \iiint_{vol} \rho e d(\text{Vol}) = - \iint_{cs} (\rho \mathbf{U} \cdot \mathbf{n}) e ds \quad (21)$$

where, \mathbf{n} = Unit vector normal to the surface and,

$$e = \frac{a^2}{\gamma(\gamma - 1)} + \frac{U^2}{2}$$

Substituting the value for e into Equation 21 gives:

$$\begin{aligned}
& \frac{\partial}{\partial t} \left[\rho_1 V_1 \left(\frac{a^2}{\gamma(\gamma - 1)} + \frac{U^2}{2} \right) \right] \\
& = - \left[w^* \left(\frac{a^2 (\gamma^2 - \gamma + 2)}{2\gamma (\gamma - 1)} \right) - w_1 \left(\frac{a_1^2}{\gamma (\gamma - 1)} + \frac{U_1^2}{2} \right) \right] \quad (22)
\end{aligned}$$

Expanding Equation 22 by assuming small perturbations and subtracting out the mean flow equation produces the small perturbation form of the energy equation, as shown below:

$$\begin{aligned}
& \frac{1}{\gamma(\gamma-1)} [2\bar{a}_i a'_i \bar{w}_i + \bar{a}_i^2 w'_i] + \frac{1}{2} [\bar{w}_i (2\bar{U}_i \cos (\alpha_{ch} - \beta_i)) (u'_i \sin \alpha_{ch} + \\
& \quad v'_i \cos \alpha_{ch}) + w'_i (\bar{U}_i \cos (\alpha_{ch} - \beta_i))^2] - \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} [2\bar{a}^* a'^* \bar{w}^* + \bar{a}^{*2} w'^*] \\
& = \frac{1}{\gamma(\gamma-1)} \left[\bar{a}^* \bar{V}_1 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'_1}{\partial t} + 2\bar{a} \bar{\rho} \bar{V}_1 \frac{\partial a'}{\partial t} \right]_1 + \\
& \quad \frac{1}{2} \left[\bar{U}^2 \bar{V}_1 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{U}^2 \frac{\partial V'_1}{\partial t} + 2 \bar{U} \bar{\rho} \bar{V}_1 \frac{\partial U'}{\partial t} \right]_1, \tag{23}
\end{aligned}$$

Before solving Equations 18, 20, and 23, densities are nondimensionalized by ρ_i , pressures by $\rho_i U_i^2$, velocities by U_{IRE} , lengths by semichord b , time by b/U_{IRE} , areas by A^* , and volumes by $A^* b$. The nondimensionalized form of the equations of motion for Region 1 are as follows:

Continuity Equation

$$\begin{aligned}
& [\rho'^* \bar{a}^* + \bar{\rho}^* \bar{a}^* A'^* + \bar{\rho}^* a'^*] - [\bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& \quad \bar{A}'_i \cos (\alpha_{ch} - \beta_i) + \rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i)] \\
& = \left[\bar{\rho} \frac{\partial V'}{\partial t} + \bar{V} \frac{\partial \rho'}{\partial t} \right]_1, \tag{24}
\end{aligned}$$

Momentum Equation

$$\begin{aligned}
& [\rho'_i \bar{A}_i + \bar{\rho}_i A'_i - p'^* - \bar{p}^* A'^*] = \bar{a}^* [2\bar{\rho}^* a'^* + \bar{\rho}^* A'^* \bar{a}^* + \rho'^* \bar{a}^*] - \\
& \quad \cos (\alpha_{ch} - \beta_i) [\rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i) + A'_i \cos (\alpha_{ch} - \beta_i) + 2\bar{A}_i (u'_i \sin \alpha_{ch} + \\
& \quad v'_i \cos \alpha_{ch})] + \bar{U}_1 \bar{V}_1 \frac{\partial \rho'_1}{\partial t} + \bar{\rho}_1 \bar{U}_1 \frac{\partial V'_1}{\partial t} + \bar{\rho}_1 \bar{V}_1 \frac{\partial U'_1}{\partial t} \tag{25}
\end{aligned}$$

Energy Equation

$$\begin{aligned}
 & \left[\frac{\bar{a}_i}{\gamma(\gamma-1)} + \frac{\cos^2(\alpha_{ch} - \beta_1)}{2} \right] \left[\rho'_i \bar{A}_i \cos(\alpha_{ch} - \beta_1) + A'_i \cos(\alpha_{ch} - \beta_1) + \right. \\
 & \left. \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \right] + \frac{2\bar{A}_i \cos(\alpha_{ch} - \beta_1) \bar{a}_i a'_i}{\gamma(\gamma-1)} + \\
 & \bar{A}_i \cos^2(\alpha_{ch} - \beta_1) (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) - \\
 & \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} \left[\bar{a}^{*2} (\rho'^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + 3\bar{\rho}^* a'^*) \right] \\
 = & \frac{1}{\gamma(\gamma-1)} \left[\bar{V} \bar{a}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'}{\partial t} + 2\bar{a} \bar{\rho} \bar{V} \frac{\partial a'}{\partial t} \right] + \\
 & \frac{1}{2} \left[\bar{V} \bar{U}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{U}^2 \frac{\partial V'}{\partial t} + 2\bar{U} \bar{V} \bar{\rho} \frac{\partial U'}{\partial t} \right], \quad (26)
 \end{aligned}$$

In Equations 24, 25, and 26, all quantities are nondimensionalized. The solution to these equations will next be presented.

After separating the equations into real and imaginary parts, a method of substitution is used to solve for the unsteady, complex flow field. All flow parameters within the blade channel are assumed to vary harmonically with time, as follows:

$$f' = \bar{f}' e^{ik t} \quad (27)$$

where f is any flow parameter within the blade channel. The two-dimensional upstream flow field is converted to a one-dimensional flow field utilizing a technique commonly known in analyzing the turbulent channel flows (Reference 17):

$$\bar{f}' = \frac{1}{\bar{y}} \int_0^{\bar{y}} f' dy \quad (28)$$

The energy equation is utilized to solve for the complex constant ($A_{-\infty}$) describing the upstream flowfield. The result is as follows:

$$\begin{aligned}
 A_{-\infty R} = & E_{16} \bar{A}'_{iR} + E_{17} \bar{A}'_{iI} + E_{18} \bar{\rho}'_{R*} + E_{19} \bar{\rho}'_{I*} + E_{20} \bar{A}'_{R*} + E_{21} \bar{A}'_{I*} + \\
 & E_{22} \bar{a}'_{R*} + E_{23} \bar{a}'_{I*} + E_{24} \bar{V}'_{iI} + E_{25} \bar{V}'_{iR} \quad (29)
 \end{aligned}$$

$$\begin{aligned}
 A_{-\infty I} = & E_{26} \bar{A}'_{iR} + E_{27} \bar{A}'_{iI} + E_{28} \bar{\rho}'_{R*} + E_{29} \bar{\rho}'_{I*} + E_{30} \bar{A}'_{R*} + E_{31} \bar{A}'_{I*} + \\
 & E_{32} \bar{a}'_{R*} + E_{33} \bar{a}'_{I*} + E_{34} \bar{V}'_{iI} + E_{35} \bar{V}'_{iR} \quad (30)
 \end{aligned}$$

The derivations of Equations 29 and 30 are presented in Appendix B. The E constants in the above equations are a function of the steady flow parameter and are presented in detail in Appendix E. Calculation of the steady flow parameters appears in Appendix F. Area and volume perturbations are a function of the known vibrational mode shapes. The calculation procedure appears in Appendix G. The two complex unknowns in Equations 29 and 30, $\bar{\rho}'^*$ and \bar{a}'^* , are obtained through the use of the momentum and continuity equations.

Using the momentum equation to solve for the complex density perturbation at the throat gives:

$$\begin{aligned} \bar{\rho}'_R^* = & M_{32}\bar{A}'_{iR} + M_{33}\bar{A}'_{iI} + M_{34}\bar{A}'_{R^*} + M_{35}\bar{A}'_{I^*} + \\ & M_{36}\bar{a}'_{R^*} + M_{37}\bar{a}'_{I^*} + M_{38}\bar{V}'_{iI} + M_{39}\bar{V}'_{iR} \end{aligned} \quad (31)$$

$$\begin{aligned} \bar{\rho}'_I^* = & M_{40}\bar{A}'_{iR} + M_{41}\bar{A}'_{iI} + M_{42}\bar{A}'_{R^*} + M_{43}\bar{A}'_{I^*} + \\ & M_{44}\bar{a}'_{R^*} + M_{45}\bar{a}'_{I^*} + M_{46}\bar{V}'_{iI} + M_{47}\bar{V}'_{iR} \end{aligned} \quad (32)$$

The M constants represent a function of the steady flow parameters. The derivation of Equations 31 and 32 appear in Appendix B. The final unknown, \bar{a}'^* , is obtained through the use of the continuity Equation (24). The result is:

$$\begin{aligned} \bar{a}'_{R^*} = & C_{46}\bar{A}'_{iR} + C_{47}\bar{A}'_{iI} + C_{48}\bar{A}'_{R^*} + C_{49}\bar{A}'_{I^*} + \\ & C_{50}\bar{V}'_{iI} + C_{51}\bar{V}'_{iR} \end{aligned} \quad (33)$$

$$\begin{aligned} \bar{a}'_{I^*} = & C_{52}\bar{A}'_{iR} + C_{53}\bar{A}'_{iI} + C_{54}\bar{A}'_{R^*} + C_{55}\bar{A}'_{I^*} + \\ & C_{56}\bar{V}'_{iI} + C_{57}\bar{V}'_{iR} \end{aligned} \quad (34)$$

The C constants also represent functions of the steady flow parameters and are presented in Appendix E. The derivations of Equations 33 and 34 also appear in Appendix B. This completes the analysis for Region 1.

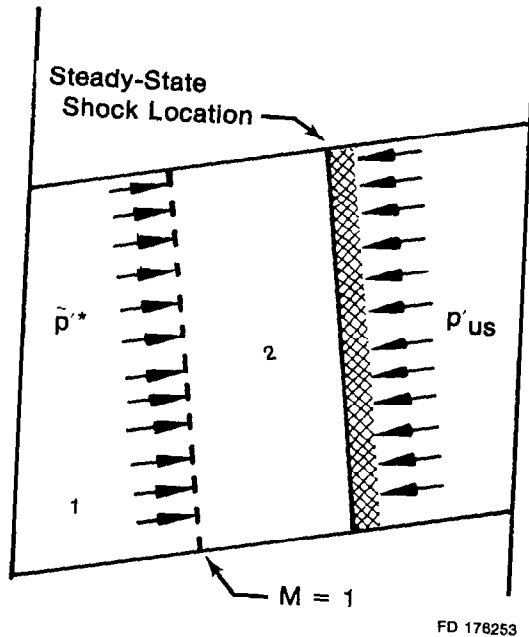
Perturbation Equations for Region 2

Region 2 in Figure 6 represents the supersonic region from the blade throat ($M=1$) to the steady-state shock position. Since all the inlet flow parameters are known, only two flow parameters need to be found: (1) perturbation velocity (U_{us}) and (2) density (ρ'_{us}) on the upstream side of the shock. Thus, only the continuity and momentum equations are needed to solve for the density and velocity perturbations. Dealing first with the continuity equation, recall that:

$$\frac{dm}{dt} = 0 = \iint \rho U dA + \frac{\partial}{\partial t} \iiint \rho d(\text{Vol}) \quad (16)$$

which becomes for Region 2:

$$\rho_{us} U_{us} A_b - \rho^* A^* a^* = \frac{-\partial \rho V_2}{\partial t} \quad (35)$$



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Figure 6. Region 2 from Channel Throat ($M = 1$) to the Steady-State Shock Location

Assuming small perturbations produces:

$$\begin{aligned}
& (\rho'_{us} \bar{U}_{us} \bar{A}'_s + \bar{\rho}_{us} U'_{us} \bar{A}'_s + \bar{\rho}_{us} \bar{U}_{us} \bar{A}'_s) - (\rho'^* \bar{A}^* \bar{a}^* + \bar{\rho}^* \bar{A}^* \bar{a}^* + \\
& \bar{\rho}^* \bar{A}^* \bar{a}^*) = - \left(\bar{V}_2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \frac{\partial V'_2}{\partial t} \right)_2
\end{aligned} \tag{36}$$

Following similar lines to the derivation of the momentum equation in Region 1, the resultant equation becomes:

$$\begin{aligned}
& (-p'_{us} \bar{A}'_s - \bar{p}_{us} \bar{A}'_s + p'^* \bar{A}^* + \bar{p}^* \bar{A}^*) \\
& = (U'_{us} \bar{w}_{us} + \bar{U}_{us} w'_{us}) - (a'^* \bar{w}^* + \bar{a}^* w'^*) + \\
& \bar{U}_2 \bar{V}_2 \frac{\partial \rho'_2}{\partial t} + \bar{\rho}_2 \bar{U}_2 \frac{\partial V'_2}{\partial t} + \bar{\rho}_2 \bar{V}_2 \frac{\partial U'_2}{\partial t}
\end{aligned} \tag{37}$$

These equations are nondimensionalized in the same manner as in Region 1 to yield the following equations for continuity and momentum:

Continuity Equation

$$\begin{aligned}
& (\rho'_{us} \bar{U}_{us} \bar{A}'_s + \bar{\rho}_{us} U'_{us} \bar{A}'_s + \bar{\rho}_{us} \bar{U}_{us} \bar{A}'_s) - (\rho'^* \bar{a}^* + \bar{\rho}^* \bar{a}^* + \bar{\rho}^* \bar{A}^* \bar{a}^*) \\
& = - \left(\bar{V}_2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \frac{\partial V'_2}{\partial t} \right)_2
\end{aligned} \tag{38}$$

Momentum Equation

$$\begin{aligned}
& (-p'_{us} \bar{A}'_s - \bar{p}_{us} \bar{A}'_s + p'^* + \bar{p}^* \bar{A}^*) \\
& = (U'_{us} \bar{w}_{us} + \bar{U}_{us} w'_{us}) - [\bar{\rho}^* \bar{a}^* \bar{a}^* + \bar{a}^* (\rho'^* \bar{a}^* + \bar{\rho}^* \bar{A}^* \bar{a}^* + \bar{\rho}^* \bar{a}^*)] + \\
& \bar{U}_2 \bar{V}_2 \frac{\partial \rho'_2}{\partial t} + \bar{\rho}_2 \bar{U}_2 \frac{\partial V'_2}{\partial t} + \bar{\rho}_2 \bar{V}_2 \frac{\partial U'_2}{\partial t}
\end{aligned} \tag{39}$$

The momentum equation is used to obtain the expression for the density perturbation upstream of the shock to yield:

$$\begin{aligned}
\bar{\rho}'_{USI} = & M_{52} \bar{A}'_{SR} + M_{53} \bar{A}'_{SI} + M_{54} \bar{p}'_R + M_{55} \bar{p}'_I + M_{56} \bar{A}'_R + \\
& M_{57} \bar{A}'_I + M_{58} \bar{U}'_{USR} + M_{59} \bar{U}'_{USI} + M_{60} \bar{a}'_I + M_{61} \bar{a}'_R + \\
& M_{62} \bar{\rho}'_I + M_{63} \bar{\rho}'_R + M_{64} \bar{V}'_{2I} + M_{65} \bar{V}'_{2R}
\end{aligned} \tag{40}$$

$$\begin{aligned}
\bar{\rho}'_{USR} = & M_{66} \bar{A}'_{SR} + M_{67} \bar{A}'_{SI} + M_{68} \bar{p}'_R + M_{69} \bar{p}'_I + M_{70} \bar{A}'_R + \\
& M_{71} \bar{A}'_I + M_{72} \bar{U}'_{USR} + M_{73} \bar{U}'_{USI} + M_{74} \bar{a}'_I + M_{75} \bar{a}'_R + \\
& M_{76} \bar{\rho}'_I + M_{77} \bar{\rho}'_R + M_{78} \bar{V}'_I + M_{79} \bar{V}'_R
\end{aligned} \tag{41}$$

The derivation of Equations 40 and 41 appears in Appendix C. The expressions for the density perturbations contain one unknown: the complex velocity perturbation upstream of the shock (\bar{U}'_{US}). The continuity equation (37) gives this relationship:

$$\begin{aligned} \bar{U}'_{USR} = & C_{73} \bar{A}'_{SR} + C_{74} \bar{A}'_{SI} + C_{75} \bar{\rho}'_{R*} + C_{76} \bar{\rho}'_{I*} + C_{77} \bar{a}'_{R*} + C_{78} \bar{a}'_{I*} + C_{79} \bar{p}'_{R*} + \\ & C_{80} \bar{p}'_{I*} + C_{81} \bar{A}'_{R*} + C_{82} \bar{A}'_{I*} + C_{83} \bar{V}'_1 + C_{84} \bar{V}'_R \end{aligned} \quad (42)$$

$$\begin{aligned} \bar{U}'_{USI} = & C_{85} \bar{A}'_{SR} + C_{86} \bar{A}'_{SI} + C_{87} \bar{\rho}'_{R*} + C_{88} \bar{\rho}'_{I*} + C_{89} \bar{a}'_{R*} + \\ & C_{90} \bar{a}'_{I*} + C_{91} \bar{p}'_{R*} + C_{92} \bar{p}'_{I*} + C_{93} \bar{A}'_{R*} + C_{94} \bar{A}'_{I*} + C_{95} \bar{V}'_R + C_{96} \bar{V}'_1 \end{aligned} \quad (43)$$

The derivation of Equations 42 and 43 is presented in Appendix C. The M and C coefficients are found in Appendix E.

This completes the analysis for Region 2. The unsteady shock wave movements are next defined along with equations describing flow discontinuities across the shock.

Unsteady Shock Movement

For a normal shock wave moving at a velocity, U_s with respect to the channel, the pressures on the two sides of the wave relate in the following manner (Reference 18):

$$p_{ds} = p_{us} \left[\frac{2\gamma \left(M_{us} - \frac{U_{us}}{a_{us}} \right) - (\gamma - 1)}{\gamma + 1} \right] \quad (44)$$

Assuming small perturbations and nondimensionalizing Equation 44 yields:

$$\begin{aligned} \frac{2\gamma}{\gamma + 1} \bar{p}_{us} U'_s = & \left[\bar{p}_{us} \bar{M}_{us} - \bar{p}_{ds} - \left(\frac{\gamma - 1}{\gamma + 1} \right) \bar{p}_{us} \right] a'_{us} + \\ & \left[\frac{2\gamma}{\gamma + 1} \bar{a}_{us} \bar{M}_{us} - U'_s - \left(\frac{\gamma - 1}{\gamma + 1} \right) \bar{a}_{us} \right] p'_{us} + \\ & \left[\frac{2\gamma}{\gamma + 1} \bar{a}_{us} \bar{p}_{us} M'_{us} - \bar{a}_{us} p'_{ds} \right] \end{aligned} \quad (45)$$

Again, all flow parameters are assumed to vary harmonically with time. The change in shock position can be related to the shock velocity perturbation by:

$$U'_s = \bar{U}'_s \exp ikt = \frac{\partial \bar{x}'_s}{\partial t} \quad (46)$$

Integrating Equation 46 gives:

$$\begin{aligned} \bar{x}'_{SR} = & \frac{\bar{U}'_{SI}}{k} \\ \bar{x}'_{SI} = & - \frac{\bar{U}'_{SR}}{k} \end{aligned} \quad (47)$$

Substituting Equation 46 into Equation 45, dividing by e^{ikt} , and solving for \bar{U}'_s produces:

$$\begin{aligned}\bar{U}'_{sR} &= \frac{1}{S_1} \left[S_2 \bar{a}'_{usR} + S_3 \bar{p}'_{usR} + S_4 \frac{U'_{usR}}{\bar{a}_{us}} - \bar{a}_{us} \bar{p}'_{dsR} \right] \\ \bar{U}'_{sI} &= \frac{1}{S_1} \left[S_2 \bar{a}'_{usI} + S_3 \bar{p}'_{usI} + S_4 \frac{U'_{usI}}{\bar{a}_{us}} - \bar{a}_{us} \bar{p}'_{dsI} \right]\end{aligned}\quad (48)$$

where the S coefficients are presented in Appendix E.

The pressure perturbation upstream of the shock can now be related to the density perturbation in the following manner:

$$\bar{p}'_{us} = \bar{a}_{us}^2 \bar{\rho}'_{us}$$

and to the speed of sound perturbation as

$$\bar{a}'_{us} = \left(\frac{\gamma - 1}{2} \right) \left(\frac{\bar{a}_{us}}{\bar{\rho}_{us}} \right) \bar{\rho}'_{us}$$

Thus, there exists only one unknown in Equation 48: \bar{p}'_{ds} . The downstream pressure perturbation is found by expanding the following equation in small perturbation form. The density discontinuity across the shock is

$$\rho_{ds} [(\gamma - 1) M_{us}^2 + 2] = \rho_{us} [(\gamma + 1) M_{us}^2] \quad (49)$$

This equation in small perturbation form can be expressed as

$$\begin{aligned}\bar{\rho}'_{dsR} &= S_8 \bar{U}'_{usR} + S_9 \bar{\rho}'_{usR} \\ \bar{\rho}'_{dsI} &= S_8 \bar{U}'_{usI} + S_9 \bar{\rho}'_{usI}\end{aligned}\quad (50)$$

Substituting the relationships of Equation 50 into Equation 48 produces the shock perturbation velocity.

The velocity perturbation downstream of the shock is found by satisfying the continuity across the shock, as noted in the following equation:

$$\rho_{us} U_{us} A_s = \rho_{ds} U_{ds} A_s \quad (51)$$

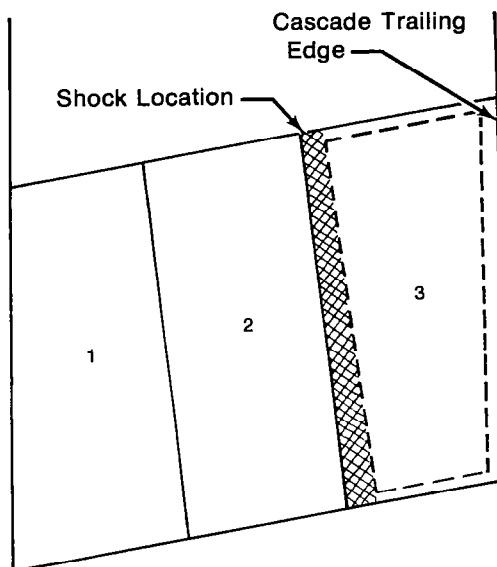
Expanding Equation 51 in small perturbation form and solving for \bar{U}'_{ds} gives

$$\bar{U}'_{dsR} = \frac{1}{\bar{\rho}_{ds}} \left[\bar{\rho}'_{usR} \bar{U}_{us} + \bar{\rho}_{us} \bar{U}'_{usR} - \bar{\rho}'_{dsR} \bar{U}_{ds} \right] \quad (52)$$

$$\bar{U}'_{dsI} = \frac{1}{\bar{\rho}_{ds}} \left[\bar{\rho}'_{usI} \bar{U}_{us} + \bar{\rho}_{us} \bar{U}'_{usI} - \bar{\rho}'_{dsI} \bar{U}_{ds} \right] \quad (53)$$

This completes the relationships describing the flow perturbations across the normal shock.

Perturbation Equations for Region 3



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Figure 7. Region 3 from the Steady-State Shock Location to the Cascade Trailing Edge

Region 3 is the subsonic intrablade region from the shock location to the cascade trailing edge, as shown in Figure 7. The unsteady flow field is derived using the perturbation relationships for the flow entering the region across the normal shock and exiting the cascade into the downstream flow field. The unsteady flow field exiting Region 3 is represented as the sum of two flow fields: (1) the irrotational part derived previously in the section on upstream and downstream irrotational flow, and (2) a rotational part related to the vorticities being shed off the trailing edge of the blades due to the unsteady vibratory motion.

Rotational Downstream Flowfield

Lacking viscosity, the equation of motion in vorticity form for two-dimensional flow is (Reference 19):

$$\rho \frac{D\zeta}{Dt} - \zeta \frac{D\rho}{Dt} = 0 \quad (54)$$

In small perturbation form this equation becomes

$$\frac{D\zeta'}{Dt} - \frac{\bar{\zeta}}{\bar{\rho}} \frac{D\rho'}{Dt} = 0 \quad (55)$$

where, D/Dt represents the substantial derivative.

The vorticity of a nonviscous field remains zero, if at the beginning the vorticity was equal to zero and the fluid is only subjected to the forces which have a potential associated with them. A shock wave does not fit into this category. However, according to Crocco's Theorem (Reference 20), if the fluid passes through a stationary shock wave, the flow can conserve its irrotational character only if the entropy rise is uniform across the shock. This is the case for the normal shock in a one-dimensional channel flow. Thus, the mean flow vorticity is zero. Using this fact in Equation 55 produces

$$\frac{D\zeta'}{Dt} = 0 \quad (56)$$

A solution to Equation 56 is assumed.

$$\zeta(x,y,t) = Z_{+r} \exp i(Rx + Cy + kt) \quad (57)$$

where C is defined in Equation 9 and k is the reduced frequency. Substituting Equation 57 into Equation 56 and solving for R results in

$$R = \frac{-(k + \bar{v}_E C)}{\bar{u}_E} \quad (58)$$

The vorticity can be related to the stream function as follows (Reference 21):

$$\nabla^2 \psi' = -\zeta'$$

where,

$$u' = \frac{\partial \psi'}{\partial y}$$

$$v' = -\frac{\partial \psi'}{\partial x}$$

A solution for the stream function ψ exists in the following form:

$$\psi' = \frac{Z_{+\infty}}{R^2 + C^2} \exp i (Rx + Cy + kt) \quad (59)$$

which will give the rotational velocity perturbation at the exit.

Continuity and Momentum Equations for Region 3

Because there are only two unknowns in the downstream flow field, the complex irrotational and rotational constants, the momentum and continuity equations are all that are required. The nondimensionalized small perturbation forms of these equations are

Continuity

$$\begin{aligned} & [(\rho'_{\text{E}} \bar{A}_{\text{E}} \bar{U}_{\text{E}} + \bar{\rho}_{\text{E}} A'_{\text{E}} \bar{U}_{\text{E}} + \bar{\rho}_{\text{E}} \bar{A}_{\text{E}} U'_{\text{E}}) - (\bar{\rho}_{\text{ds}} \bar{A}_{\text{E}} U'_{\text{ds}} + \bar{\rho}_{\text{ds}} A'_{\text{s}} \bar{U}_{\text{ds}} + \\ & \rho'_{\text{ds}} \bar{A}_{\text{s}} \bar{U}_{\text{ds}})] = \left[\bar{\rho} \frac{\partial V'}{\partial t} + \bar{V} \frac{\partial \rho'}{\partial t} \right]_3 \end{aligned} \quad (60)$$

Momentum

$$\begin{aligned} & (\rho'_{\text{ds}} \bar{A}_{\text{s}} + \bar{\rho}_{\text{ds}} A'_{\text{s}} - \rho'_{\text{E}} \bar{A}_{\text{E}} - \bar{\rho}_{\text{E}} A'_{\text{E}}) = (U'_{\text{E}} \bar{w}_{\text{E}} + \bar{U}_{\text{E}} w'_{\text{E}}) - \\ & (U'_{\text{ds}} \bar{w}_{\text{ds}} + \bar{U}_{\text{ds}} w'_{\text{ds}}) + \bar{U}_{\text{ave}} \bar{V}_3 \frac{\partial \rho'_{\text{3}}}{\partial t} + \bar{\rho}_{\text{3}} \bar{U}_3 \frac{\partial V'_{\text{3}}}{\partial t} + \\ & \bar{\rho}_{\text{3}} \bar{V}_3 \frac{\partial U'_{\text{3}}}{\partial t} \end{aligned} \quad (61)$$

The momentum equation is used to solve for the complex constant $A_{+\infty}$ to produce

$$A_{+\infty R} = Z_{+\infty} M_{82} + Z_{+\infty I} M_{83} + RC M_{84} + IC M_{85} \quad (62)$$

$$A_{+\infty I} = Z_{+\infty R} M_{86} + Z_{+\infty I} M_{87} + RC M_{88} + IC M_{89} \quad (63)$$

The M constants, along with RC and IC, appear in Appendix E.

The continuity equation is used to solve for the complex constant describing the rotational downstream flow field resulting in

$$Z_{+\infty R} = RC C_{112} + IC C_{113} + C_{114} \quad (64)$$

$$Z_{+\infty I} = RC C_{115} + IC C_{116} + C_{117} \quad (65)$$

where the C constants are found in Appendix E.

Equations 64 and 65 can then be substituted into Equations 62 and 63 to solve for $A_{+\infty}$ giving a complete description of the flow field in Region 3. The derivation of Equations 62 through 65 is presented in Appendix D.

Thus, with the solution of the flow field in Region 3, a complete description of the cascade flow field is obtained. A compilation of the computer code for the semiactuator disk model is presented in Appendix J.

RESULTS

The semi-actuator disk theory was evaluated in two ways. The first method consisted of a comparison with test data presented by Tanida and Saito (Reference 22) for an isolated airfoil oscillating in choked flow in a wind tunnel. The second method involved a flutter analysis of the F100(3) 6th stage of the high-pressure compressor of the F100(3) turbofan engine, which encountered choke flutter while operating at off-design conditions in a core engine (no low rotor) at the Arnold Engineering Development Center (AEDC).

Wind Tunnel Test Data

Tanida and Saito oscillated an airfoil in a wind tunnel at constant amplitude for various combinations of inlet Mach number, back pressure, reduced frequency twist axis location, and tunnel wall separation, and recorded both steady and unsteady aerodynamic characteristics. Because reflections from the tunnel walls create a special case in a cascade in which adjacent blades are exactly out of phase, the experimental test conditions were simulated with an interblade phase angle of 180 deg. After calculating unsteady pressures through the blade channel, the unsteady lift and momentum coefficients were calculated in the manner presented in Appendix I.

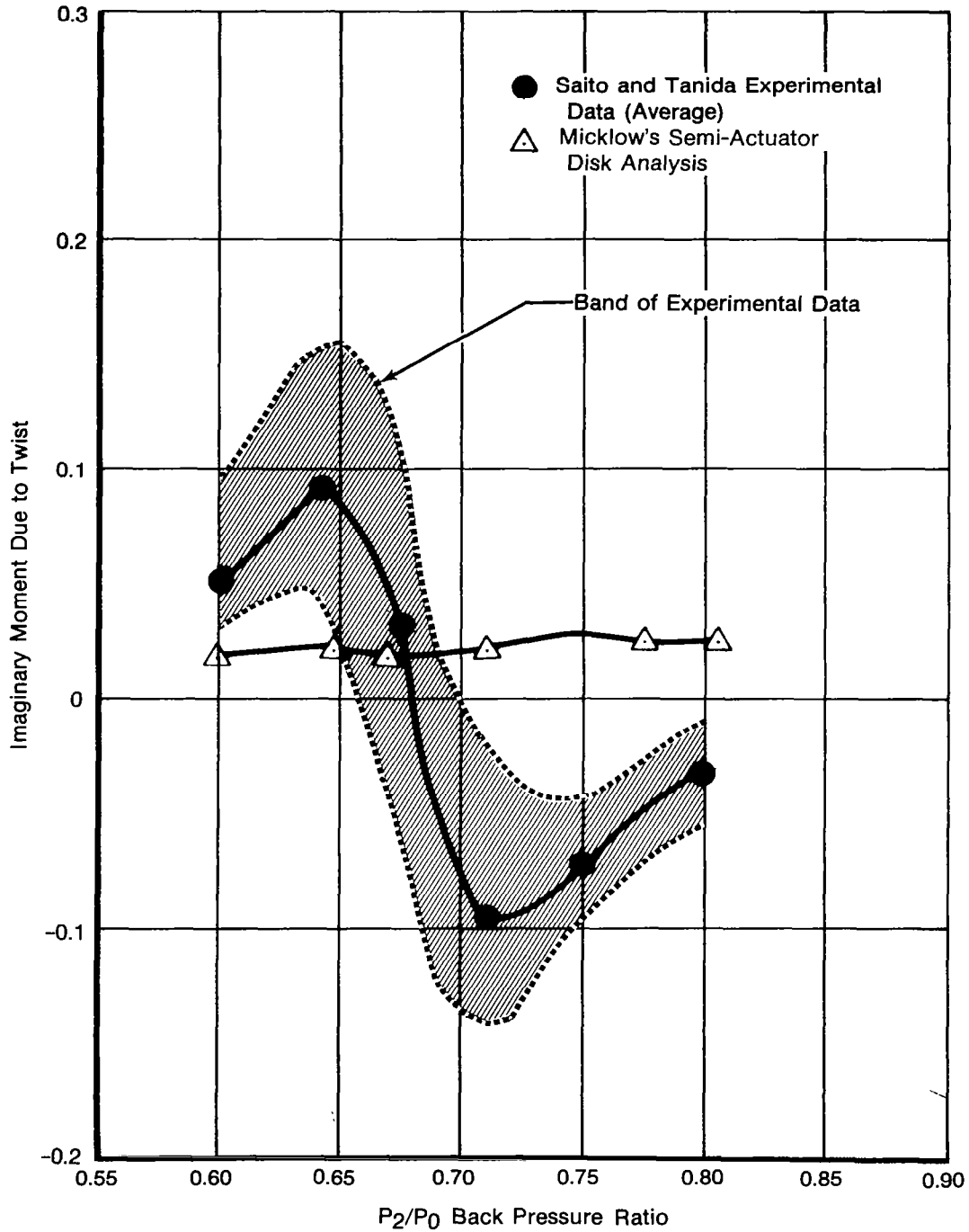
General agreement was achieved for the imaginary moment due to pitch for back-pressure ratios less than 0.7, as shown in Figure 8. At higher back-pressure ratios, flow is not fully transonic during the full cycle of operation; i.e., a weak shock appears and disappears on the airfoil surface. A basic assumption of the analytical model stipulates that flow is fully transonic, and a strong normal shock exists in the blade passage throughout the oscillatory cycle. Because this was not the case at high back-pressure ratios, good correlation was not anticipated.

Flutter Analysis

Computational Method

Investigations also involved an assessment of the semi-actuator disk unsteady aerodynamics by combining the model with existing P&WA cyclic work and aerodynamic damping calculations, and then performing a flutter analysis for the F100(3) sixth compressor stage, which experienced choke flutter at off-design conditions. The P&WA approach to flutter prediction stems from a cyclic energy method in which total system damping is calculated. The system becomes unstable when total damping, which is comprised of aerodynamic and mechanical damping components, is less than zero; i.e.,

$$\delta_{TOT} = (\delta_{aero} + \delta_{mech}) < 0 \Rightarrow \text{flutter}$$



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Figure 8. Comparison Between Wind Tunnel Data for Isolated Airfoil and Semi-Actuator Disk Analysis

No method currently exists to determine mechanical damping. Thus, stability predictions are based on correlations of aerodynamic damping for an assumed level of mechanical damping. To calculate δ_{aero} , a quasi-three-dimensional analysis is used in which two-dimensional unsteady aerodynamic work (W) is calculated for strips along the airfoil span. As shown by Carta (Reference 23), the two-dimensional work can be expressed as

$$W = \rho\pi b^2 U^2 k^2 \left\{ L_{HI} \bar{h}^2 + \bar{\alpha} \bar{h} [(L_{oR} - M_{HR}) \sin \theta + (L_{oI} + M_{HI}) \cos \theta] + M_{oI} \bar{\alpha}^2 \right\}$$

where,

- L = unsteady lift coefficient
- M = unsteady moment coefficient
- $\bar{\alpha}$ = normalized mode shape deflection of maximum twist
- \bar{h} = normalized mode shape deflection of maximum bending
- θ = phase relationship between α and h .

Numerically integrating the two-dimensional work along the span produces total unsteady cyclic work for one blade (W_{TOT}). The logarithmic decrement, or aerodynamic damping, is then calculated as

$$\delta_{aero} = - \frac{n W_{TOT}}{4KE}$$

where,

- n = number of blades in the system
- KE = normalized average kinetic energy of the system vibrating in the normalized mode.

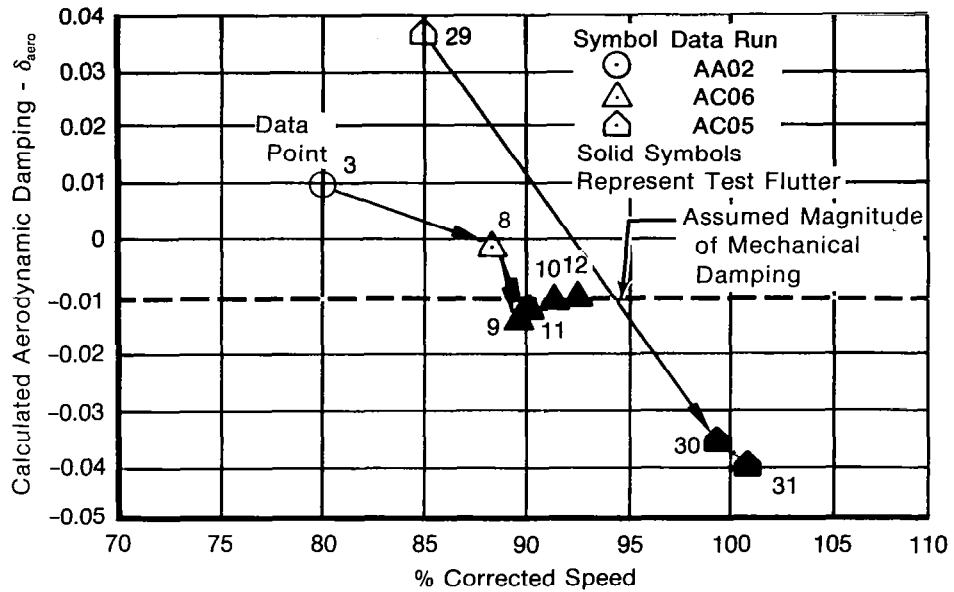
Flutter Prediction Results

A flutter analysis of the F100 sixth stage compressor operating at sea level conditions was performed, and a summary of the data points, operating conditions and predicted damping values is presented in Table 2. The model predicted the first bending mode to be least stable, which is consistent with the results observed by Lubomski (Ref. 24). However, test data showed Rotor 6 encountered negative incidence flutter in the second coupled mode, with some secondary first-mode response.

For both the first bending and second coupled modes, the model shows a distinct difference in damping level between flutter and non-flutter points, with all of the flutter points analyzed having a negative damping value, indicating an unstable condition as seen in Figure 9 and 10. In general, the model predicts the correct trend with increasing speed, i.e., a decrease in stability with increasing speed.

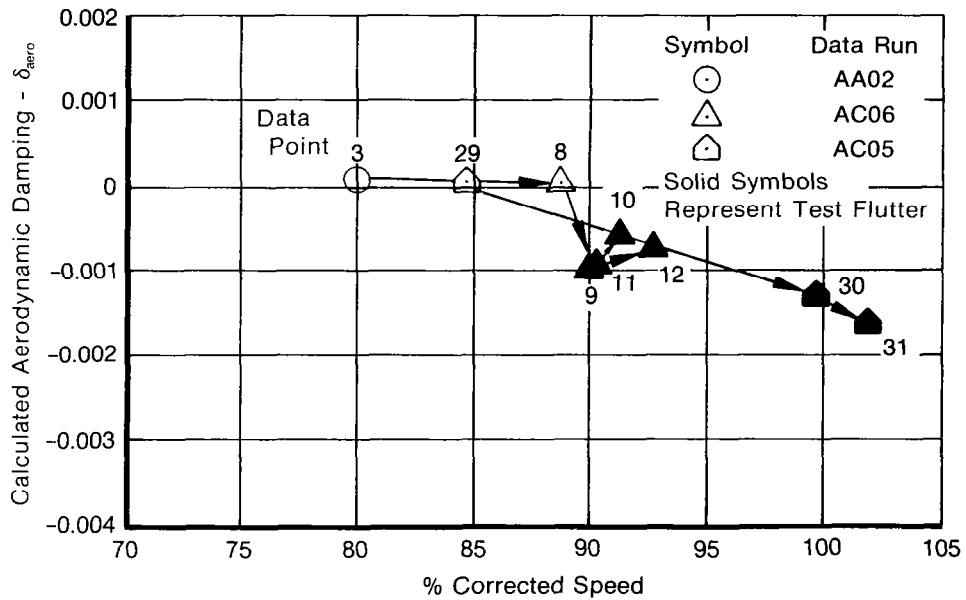
TABLE 2. SUMMARY OF LEAST-STABLE NODAL DIAMETERS AND DAMPING VALUES FOR SEA LEVEL CONDITIONS

<i>Test Data</i>				<i>Stability Calculations</i>		
<i>Data Point</i>	<i>Test Stability</i>	<i>% N_{cor}</i>	<i>Rear Compressor Variable Vane, RCVV</i>	<i>Least Stable Nodal-Dia.</i>	<i>Aerodynamic Damping δ_{aer}</i>	
					<i>Smith</i>	<i>Semi-Actuator Disk Theory</i>
<i>F100(3) Rotor 6 1st Bending Mode</i>						
AA02PT3	NF	80.0	-30.0	2	0.0164	0.0102
AC06PT8	NF	88.7	-30.0	2	0.02978	-0.00104
AC06PT9	FB	89.75	-33.0	4	0.0558	-0.01394
AC06PT10	F	91.4	-37.2	3	0.04605	-0.0110
AC06PT11	FB	90.21	-29.2	3	0.05793	-0.01313
AC06PT12	F	92.4	-29.0	3	0.05516	-0.0106
AC05PT29	NF	84.76	-20.0	2	0.03718	0.03718
AC05PT30	NB	99.68	-20.8	4	0.0795	-0.0356
AC05PT31	F	102.00	-21.0	4	0.06492	-0.03990
<i>F100(3) Rotor 6 2nd Coupled Mode</i>						
AA02PT3	NF	80.0	-30.0	3	0.000125	0.000101
AC06PT8	NF	88.7	-30.0	3	0.000275	0.000097
AC06PT9	FB	89.75	-33.0	6	0.01894	-0.00101
AC06PT10	F	91.4	-37.2	6	0.01334	-0.00053
AC06PT11	FB	90.21	-29.2	6	0.0182	-0.00097
AC06PT12	F	92.4	-29.0	4	0.00170	-0.00077
AC05PT29	NF	84.76	-20.0	2	0.000354	0.000354
AC05PT30	FB	99.68	-20.8	6	0.0299	-0.00130
AC05PT31	F	102.0	-21.0	6	0.0243	-0.00166
where,	NF	= No Flutter				
	FB	= Flutter Boundary				
	F	= Flutter				



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Figure 9. Comparison of Calculated Aerodynamic Damping and Observed Stability for the First Vibratory Mode of the F100 6th-Stage Compressor



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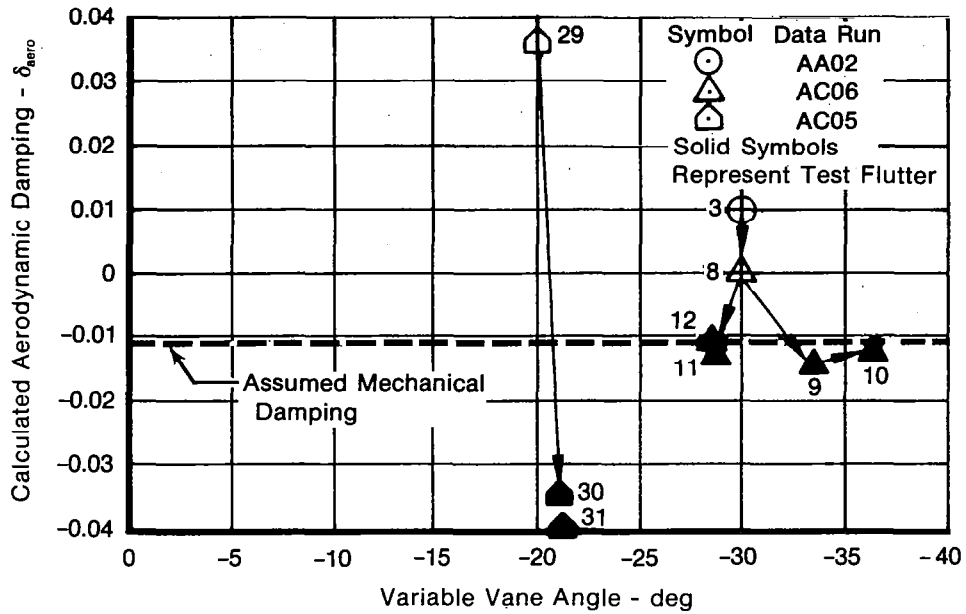
Figure 10. Comparison of Calculated Aerodynamic Damping and Observed Stability for the Second Vibratory Mode of the F100 6th-Stage Compressor

The slight discrepancy in this trend exhibited by the calculated damping values for flutter data points AC06-10 and -12 may be attributed to inaccuracies in the steady streamline aerodynamic input. The inaccuracies arise due to the difficulty in predicting the steady aerodynamic environment while the compressor is operating well off the design condition in flutter. The model is very sensitive to inlet air angle, relative inlet and exit Mach numbers, and static pressure ratio across the stage. The inlet Mach number and air angle are used to check for choked flow in the blade passage, and the static pressure ratio is used to locate the steady-state shock location. For data points AC06-9 and -11, the static pressure ratio across the stage was matched within 2 percent, while for AC06-10, the error was as high as 15 percent. This magnitude of discrepancy in matching the stage static pressure ratio could lead to inaccuracies in locating the steady-state normal shock. In the analysis, the chordwise location of the shock greatly affects the magnitude and sign of the unsteady aerodynamic coefficients, thus greatly affecting the damping calculation. It was noted that data points AC06-9 and -11 were calculated to be choked along the entire span, while AC06-10 and -12, operating deeper into the flutter boundary, were not. This has a large effect on the stability calculation because if the flow is determined to be unchoked, the unsteady, zero-incidence Smith coefficients (Reference 2) are used in place of the semi-actuator disk coefficients. The Smith coefficients always predicted the rotor to be stable.

The magnitude of the inlet and exit Mach numbers proved to be important parameters in describing the upstream and downstream flow fields. With increasing inlet Mach number, the unsteady lift due to flap increased in a destabilizing manner and with increasing exit Mach number. The unsteady moment due to twist showed a similar trend. Data points AC05-30 and -31 are approximately double the value of the negative damping of the other flutter points. This is due in part to the greater inlet and exit Mach numbers. The magnitude of the inlet air angles are related to the magnitude of the inlet Mach number through the continuity equation in the steady streamline analysis. It is noted that the streamline analysis indicated data point AC06-9 was operating above β_{\max} at the tip. The air angle is related to the flow area; therefore, to satisfy continuity the inlet Mach number must decrease at this station. This is particularly significant because the maximum unsteady work occurs at the tip. A similar trend was noted at other sections. AC05-30 had similar trends, but not as severe, giving a more accurate inlet Mach number. It is speculated that a more accurate representation of AC06-9 would give a larger inlet Mach number and increase the damping in a negative sense, bringing the flutter points closer to the same damping level.

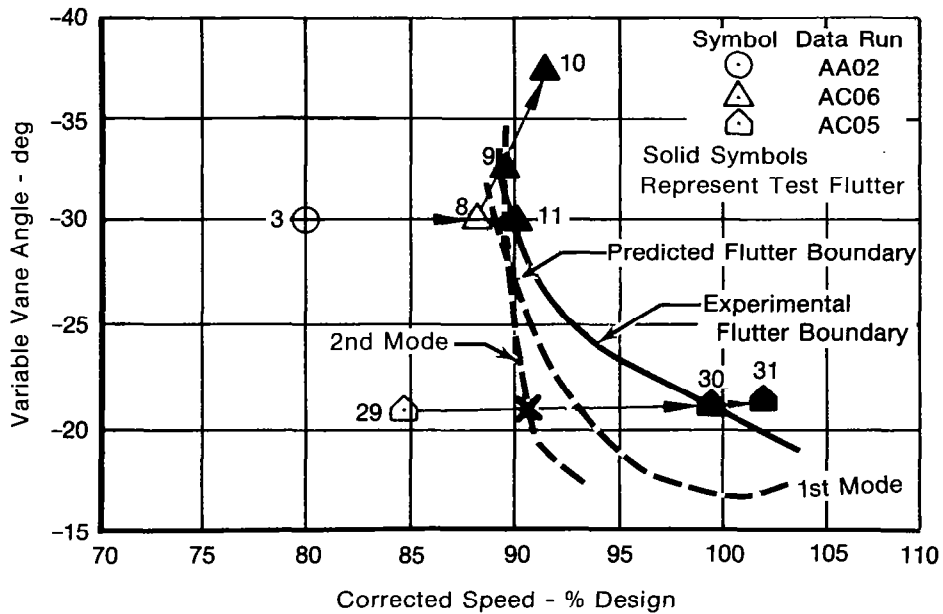
By plotting the aerodynamic damping versus corrected speed, as in Figures 9 and 10, or versus variable inlet guide vane setting, as in Figure 11, a predicted flutter boundary can be obtained by assuming a value for mechanical damping. As shown in Figure 12, the results of the flutter analysis prove to be conservative, particularly at higher corrected speeds. This conservatism is believed to be due to the inlet Mach number and air angle discrepancy discussed above resulting in two negative damping levels. A summary of the least stable nodal diameters and damping values is presented in Table 2 for the 1st bending and 2nd coupled modes of vibration. All the predicted least stable nodal diameters are relatively small because the model encountered numerical difficulties at large interblade phase angles as shown in Figure 13.

The General Electric annular cascade data (Reference 25), which was originally to be analyzed, had to be abandoned because the published data gave an inadequate description of the steady flow field. No exit data were published and attempts to calculate aerodynamic damping using exit conditions based on PWA 2-dimensional cascade test data produced no meaningful results.



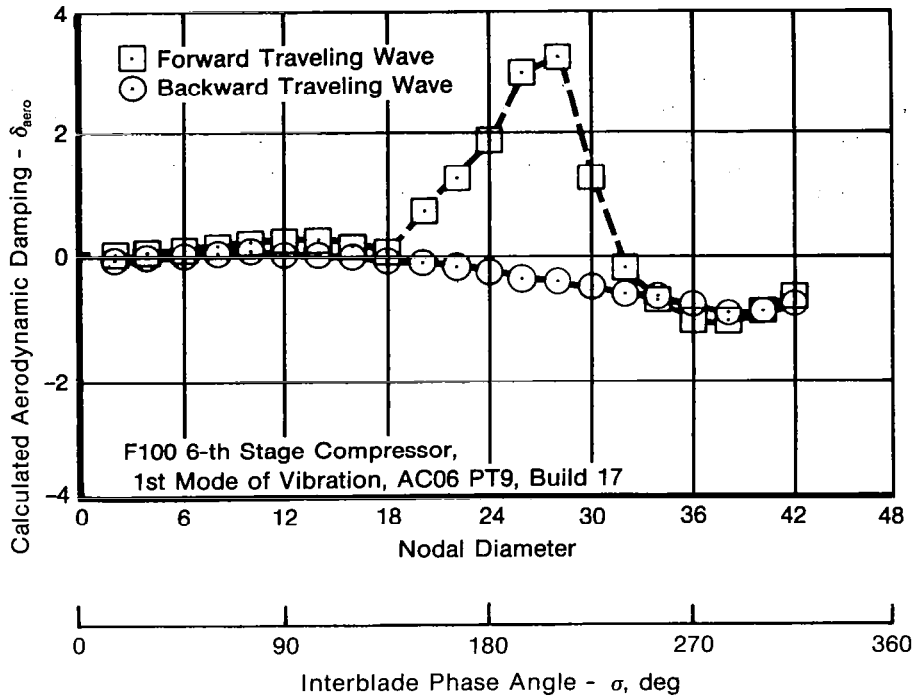
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Figure 11. Calculated Aerodynamic Damping Versus Variable Vane Angle With 1st Mode of Vibration, Backward Traveling Wave



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Figure 12. Comparison of Predicted and Experimental Flutter Boundaries



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Figure 13. Effect of Nodal Diameter and Interblade Phase Angle on Aerodynamic Damping

Summary of Results

The following is a summarization of the results of the flutter prediction study.

1. The model showed a distinct difference in damping level between flutter and non-flutter points indicating the importance of the unsteady shock in the blade passage.
2. The model shows the correct trend with increasing corrected engine speed and vane angle; i.e., with increasing choked conditions, the model predicts the aerodynamic damping to become less stable.
3. The model was very sensitive to the steady aerodynamic input. In order to perform an accurate flutter analysis, an accurate steady aerodynamic description of the flow field must be obtained.
4. The model was conservative in predicting the flutter boundary at high corrected engine speeds.
5. Due to limitations on interblade phase angle inherent in semiactuator disk theory, the model is limited to interblade phase angles of ± 90 deg.

Conclusions

Based on the results of the F100(3) 6th-stage compressor flutter analysis, it is concluded that the model is useful as a conservative choke flutter design system. The model is sensitive to the steady-state aerodynamic input, particularly inlet and exit relative Mach number, inlet air angle, and static pressure ratio across the stage. In order to perform an accurate flutter analysis, an accurate description of the steady flowfield must be obtained.

APPENDIX A
DERIVATION OF THE SMALL PERTURBATION FORM
OF THE EQUATIONS OF MOTION IN REGION 1

CONTINUITY EQUATION

The continuity equation for Region 1 takes the form

$$\dot{w}^* - \dot{w}_{inlet} = - \frac{\partial(\rho_1 V_1)}{\partial t} \quad (A1)$$

Mean Flow

Assuming small perturbations about the mean flow gives the following expressions for the flow rates, neglecting higher order terms.

$$\begin{aligned} \dot{w}_{inlet} &= \bar{w}_{inlet} + \dot{w}'_{inlet} = (\bar{\rho}_1 + \rho'_1) (\bar{A}_1 + A'_1) (\bar{U}_{inlet} + U'_1) \\ &= \bar{\rho}_1 \bar{A}_1 \bar{U}_1 + \rho'_1 \bar{A}_1 \bar{U}_1 + \bar{\rho}_1 A'_1 \bar{U}_1 + \bar{\rho}_1 \bar{A}_1 U'_1 \end{aligned}$$

Let U'_1 be the perturbation velocity aligned with chord line. Then

$$U'_1 = u'_1 \sin \alpha_{ch} + v'_1 \cos \alpha_{ch}$$

The inlet velocity is $U_{IRE} \cos(\alpha_{ch} - \beta_1)$. The inlet flowrate is then defined as

$$\begin{aligned} \dot{w}_{inlet} &= \bar{w}_{inlet} + \dot{w}'_{inlet} = \bar{\rho}_1 \bar{A}_1 \bar{U}_1 \cos(\alpha_{ch} - \beta_1) + \\ &\quad \bar{\rho}_1 \bar{A}_1 (v'_1 \cos \alpha_{ch} + u'_1 \sin \alpha_{ch}) + \bar{\rho}_1 \bar{U}_1 \cos(\alpha_{ch} - \beta_1) A'_1 + \\ &\quad \rho'_1 \bar{A}_1 \bar{U}_1 \cos(\alpha_{ch} - \beta_1) \end{aligned} \quad (A2)$$

The flowrate at the throat is represented by

$$\dot{w}^* = \bar{w}^* + \dot{w}^* = \bar{\rho}^* \bar{A}^* \bar{a}^* + \rho'^* \bar{A}^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + \bar{\rho}^* \bar{A}^* a'^* \quad (A3)$$

Next, expand the term $\frac{\partial(\rho_1 V_1)}{\partial t}$

$$\frac{\partial(\rho_1 V_1)}{\partial t} = \rho_1 \frac{\partial V_1}{\partial t} + V_1 \frac{\partial \rho_1}{\partial t}$$

By assuming small perturbations, the following relationships emerge:

$$\rho_1 \frac{\partial V_1}{\partial t} = (\bar{\rho}_1 + \rho'_1) \left[\frac{\partial V_1}{\partial t} + \frac{\partial V'_1}{\partial t} \right] = \bar{\rho}_1 \frac{\partial V_1}{\partial t}$$

$$V_1 \frac{\partial \rho_1}{\partial t} = (V_1 + V'_1) \left[\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial \rho'_1}{\partial t} \right] = V_1 \frac{\partial \rho'_1}{\partial t}$$

$$\frac{\partial(\rho V)}{\partial t} = \bar{\rho}_1 \frac{\partial V_1}{\partial t} + V_1 \frac{\partial \rho'_1}{\partial t} \quad (A4)$$

The small perturbation form of the continuity equation for Region 1 is then:

$$\begin{aligned}
& [\rho'^* \bar{A}^* \bar{a}^* + \bar{\rho}^* \bar{A}'^* \bar{a}^* + \bar{\rho}^* \bar{A}^* \bar{a}'^*] - [\bar{\rho}_i \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& \quad \bar{\rho}_i A'_i \bar{U}_i \cos (\alpha_{ch} - \beta_i) + \rho'_i \bar{A}_i \bar{U}_i \cos (\alpha_{ch} - \beta_i)] \\
& = - \left(\bar{\rho}_i \frac{\partial V'_i}{\partial t} + \bar{V}_i \frac{\partial \rho'_i}{\partial t} \right) \tag{A5}
\end{aligned}$$

Unsteady Control Volume Form of Momentum Equation

The next step involves the derivation of the unsteady control volume form of the momentum equation. The equation has the form

$$F = \frac{d(mU)}{dt} = \iint_{cs} U(\rho U) dA + \frac{\delta}{\delta t} \iiint_{vol} U \rho d(Vol) \tag{A6}$$

with

$$\sum F = p_i A_i - p^* A^* = (\bar{p}_i + p'_i)(\bar{A}_i + A'_i) - (\bar{p}^* + p'^*)(\bar{A}^* + A'^*)$$

Neglecting higher order terms and subtracting out the mean flow forces gives

$$\sum F' = (p'_i \bar{A}_i + \bar{p}_i A'_i) - (\bar{p}^* A'^* + p'^* \bar{A}^*) \tag{A7}$$

Now to evaluate the first part of Equation A6,

$$\iint_{cs} U (\rho U \cdot dA)$$

which is essentially

$$a^* \dot{w}^* - U_i (\cos (\alpha_{ch} - \beta_i)) \dot{w}_i$$

In terms of small perturbations, this expands to

$$\begin{aligned}
& (\bar{a}^* \bar{w}^* + a'^* \bar{w}^* + \bar{a}^* w'^*) - (\bar{U}_i \bar{w}_i \cos (\alpha_{ch} - \beta_i) + U'_i \bar{w}_i + \\
& \quad \bar{U}_i \cos (\alpha_{ch} - \beta_i) w'_i)
\end{aligned}$$

Subtracting out the mean flow quantities gives

$$(a'^* \bar{w}^* + \bar{a}^* w'^*) - (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \bar{w}_i + \bar{U}_i \cos (\alpha_{ch} - \beta_i) w'_i \tag{A8}$$

To evaluate the second term of Equation A6

$$\begin{aligned}
& \frac{\partial}{\partial t} \iiint U \rho d(Vol) = \frac{\partial}{\partial t} (\rho_1 U_1 V_1) \\
& = U_1 V_1 \frac{\partial \rho_1}{\partial t} + \rho_1 U_1 \frac{\partial V_1}{\partial t} + \rho_1 V_1 \frac{\partial U_1}{\partial t}
\end{aligned}$$

Working with the first term and neglecting higher order terms produces

$$\begin{aligned} U_1 V_1 \frac{\partial \rho_1}{\partial t} &= (\bar{U}_1 \bar{V}_1 + U_1' \bar{V}_1 + \bar{U}_1 V_1') \left(\frac{\partial \bar{\rho}}{\partial t} + \frac{\partial \rho'}{\partial t} \right) \\ &= \bar{U}_1 \bar{V}_1 \frac{\partial \rho_1'}{\partial t} \end{aligned}$$

Similarly, for the second and third terms

$$\begin{aligned} \rho_1 U_1 \frac{\partial V_1}{\partial t} &= \bar{\rho}_1 \bar{U}_1 \frac{\partial V_1'}{\partial t} \\ \rho_1 V_1 \frac{\partial U_1}{\partial t} &= \bar{\rho}_1 \bar{V}_1 \frac{\partial U_1'}{\partial t} \end{aligned} \quad (\text{A9})$$

Combining Equations A7, A8, and A9 gives the small perturbation form of the momentum equation, as shown in Equation A10.

$$\begin{aligned} p' \bar{A}_1 + \bar{p} \bar{A}'_1 - \bar{p}^* A'^* - p'^* \bar{A}^* &= [a'^* \bar{w}^* + \bar{a}^* w'^*] - \\ &[\bar{w}_1 (u_1' \sin \alpha_{ch} + v_1' \cos \alpha_{ch}) + \bar{U}_1 w_1' \cos (\alpha_{ch} - \beta_1)] + \\ \bar{U}_1 \bar{V}_1 \frac{\partial \rho_1'}{\partial t} + \bar{\rho}_1 \bar{U}_1 \frac{\partial V_1'}{\partial t} + \bar{\rho}_1 \bar{V}_1 \frac{\partial U_1'}{\partial t} & \end{aligned} \quad (\text{A10})$$

Small Perturbation Form of Energy Equation

The next step requires the derivation of the small perturbation form of the energy equation for an inviscid fluid with no external heat addition. The control volume form for Region 1 is therefore

$$\frac{\partial}{\partial t} \iiint_{\text{Vol}} \rho e d(\text{Vol}) = - \iint_{\text{cn}} (\rho \mathbf{U} \cdot \mathbf{n}) e ds \quad (\text{A11})$$

where,

\mathbf{n} = unit vector normal to surface

$$e = \frac{a^2}{\gamma(\gamma - 1)} + \frac{U^2}{2}$$

Substituting this relationship into Equation A11 produces

$$\begin{aligned} \frac{\partial}{\partial t} \iiint_{\text{Vol}} \rho \left(\frac{a^2}{\gamma(\gamma - 1)} + \frac{U^2}{2} \right) d(\text{Vol}) \\ = - \iint (\rho \mathbf{U} \cdot \mathbf{n}) \left(\frac{a^2}{\gamma(\gamma - 1)} + \frac{U^2}{2} \right) ds \end{aligned}$$

which becomes

$$\begin{aligned}
& \underbrace{\frac{\partial}{\partial t} \left[\rho V \left(\frac{a^2}{\gamma(\gamma-1)} + \frac{U^2}{2} \right) \right]}_3 \\
&= - \underbrace{\left[w^* \left(\frac{a^{*2}(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} \right) \right]}_2 \quad - \underbrace{w_i \left(\frac{a_i^2}{\gamma(\gamma-1)} + \frac{U_i^2}{2} \right)}_1
\end{aligned} \tag{A12}$$

Now, in order to expand by assuming small perturbations.

Term 1 (neglecting higher order terms)

$$- \frac{\gamma^2 - \gamma + 2}{2\gamma(\gamma-1)} \left[\bar{a}^{*2} \bar{w}^* + 2\bar{a}^* a' \bar{w}^* + \bar{a}^{*2} w'^* \right] \tag{A13}$$

Term 2 (neglecting higher order terms)

$$\begin{aligned}
& (\bar{w}_i + w'_i) \left[\frac{\bar{a}_i^2 + 2\bar{a}_i a'_i}{\gamma(\gamma-1)} + \frac{1}{2} \left[(\bar{U}_i \cos(\alpha_{ch} - \beta_i))^2 + 2\bar{U}_i \cos(\alpha_{ch} - \beta_i) \cdot \right. \right. \\
& \quad \left. \left. (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \right] \right] \\
&= \frac{1}{\gamma(\gamma-1)} \left[\bar{a}_i^2 \bar{w}_i + 2\bar{a}_i a'_i \bar{w}_i + \bar{a}_i^2 w'_i \right] + \frac{1}{2} \left[\bar{w}_i (\bar{U}_i \cos(\alpha_{ch} - \beta_i))^2 + \right. \\
& \quad \left. 2\bar{U}_i \cos(\alpha_{ch} - \beta_i) (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \bar{w}_i + w'_i (\bar{U}_i \cos(\alpha_{ch} - \beta_i))^2 \right] \tag{A14}
\end{aligned}$$

Term 3

$$\begin{aligned}
& \frac{\partial}{\partial t} \left[\rho_1 V_1 \left(\frac{a_1^2}{\gamma(\gamma-1)} + \frac{U_1^2}{2} \right) \right] = \underbrace{\frac{1}{\gamma(\gamma-1)} \frac{\partial}{\partial t} (\rho_1 V_1 a_1^2)}_a + \\
& \quad \underbrace{\frac{1}{2} \frac{\partial}{\partial t} (\rho_1 V_1 U_1^2)}_b
\end{aligned}$$

Term 3a

$$\begin{aligned}
& \frac{1}{\gamma(\gamma-1)} \left[\frac{\partial(\bar{\rho}_1 \bar{V}_1 \bar{a}_1^2)}{\partial t} + \frac{\partial(\rho'_1 \bar{V}_1 \bar{a}_1^2)}{\partial t} + \frac{\partial(\bar{\rho}'_1 V'_1 \bar{a}_1^2)}{\partial t} + \frac{\partial(2a'_i \bar{a}_i \bar{\rho}_i \bar{V}_i)}{\partial t} \right] \\
&= \frac{1}{\gamma(\gamma-1)} \left[\bar{V}_i \bar{a}_i^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'_1}{\partial t} + 2\bar{a} \bar{\rho} \bar{V}_1 \frac{\partial a'}{\partial t} \right], \tag{A15}
\end{aligned}$$

Term 3b

$$\frac{1}{2} \frac{\partial(\bar{\rho}\bar{V}\bar{U}_i^2)}{\partial t} = \frac{1}{2} \left[\bar{V}_i\bar{U}_i^2 \frac{\partial\rho'}{\partial t} + \rho\bar{U}_i^2 \frac{\partial\bar{V}'_i}{\partial t} + 2\bar{U}_i\bar{\rho}\bar{V}_i \frac{\partial\bar{U}'_i}{\partial t} \right], \quad (\text{A16})$$

Combining Equations A13 through A16 and subtracting out the mean flow equation yields the small perturbation form of the energy equation:

$$\begin{aligned} & \frac{1}{\gamma(\gamma-1)} [2\bar{a}_i\bar{a}'_i\bar{w}_i + \bar{a}_i^2\bar{w}'_i] + \frac{1}{2} [\bar{w}_i (2\bar{U}_i \cos(\alpha_{ch} - \beta_i)) (u'_i \sin \alpha_{ch} + \\ & v'_i \cos \alpha_{ch}) + w'_i (\bar{U}_i \cos(\alpha_{ch} - \beta_i))^2] - \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} [2\bar{a}^*\bar{a}'^*\bar{w}^* + \bar{a}^{*2}\bar{w}'^*] \\ & = \frac{1}{\gamma(\gamma-1)} \left[\bar{a}^2\bar{V}_i \frac{\partial\rho'}{\partial t} + \bar{\rho}\bar{a}^2 \frac{\partial\bar{V}'_i}{\partial t} + 2\bar{a}\bar{\rho}\bar{V}_i \frac{\partial\bar{a}'_i}{\partial t} \right] + \\ & \frac{1}{2} \left[\bar{U}_i^2\bar{V}_i \frac{\partial\rho'}{\partial t} + \bar{\rho}\bar{U}_i^2 \frac{\partial\bar{V}'_i}{\partial t} + 2\bar{U}_i\bar{\rho}\bar{V}_i \frac{\partial\bar{U}'_i}{\partial t} \right], \quad (\text{A17}) \end{aligned}$$

In order to nondimensionalize Equations A14, A15, and A17, densities will be nondimensionalized by $\bar{\rho}_{inlet}$, pressure by $\bar{\rho}_i\bar{U}_i^2$, velocities by \bar{U}_{IRE} , lengths by semichord, time by b/\bar{U}_i , areas by A^* , and volumes by A^*b . Starting with the continuity equation:

$$\begin{aligned} & \left[\bar{\rho}_i \left(\frac{\rho'^*}{\rho_i} \right) A^* \left(\frac{\bar{A}^*}{\bar{A}^*} \right) \bar{U}_i \left(\frac{\bar{a}^*}{\bar{U}_i} \right) + \bar{\rho}_i \left(\frac{\bar{\rho}^*}{\bar{\rho}_i} \right) A^* \left(\frac{A'^*}{\bar{A}^*} \right) \bar{U}_i \left(\frac{\bar{a}^*}{\bar{U}_i} \right) + \right. \\ & \left. \bar{\rho}_i \left(\frac{\bar{\rho}^*}{\bar{\rho}_i} \right) A^* \left(\frac{\bar{A}^*}{\bar{A}^*} \right) \bar{U}_i \left(\frac{a'^*}{\bar{U}_i} \right) \right] - \left[\bar{\rho}_i \left(\frac{\bar{\rho}_i}{\bar{\rho}_i} \right) A^* \left(\frac{\bar{A}_i}{\bar{A}^*} \right) \bar{U}_i \left(\frac{u'_i}{\bar{U}_i} \sin \alpha_{ch} + \right. \right. \\ & \left. \left. \frac{v'_i}{\bar{U}_i} \cos \alpha_{ch} \right) + \bar{\rho}_i \left(\frac{\bar{\rho}_i}{\bar{\rho}_i} \right) A^* \left(\frac{\bar{A}_i}{\bar{A}^*} \right) \bar{U}_i \left(\frac{\bar{U}_i}{\bar{U}_i} \right) \cos(\alpha_{ch} - \beta_i) + \bar{\rho}_i \left(\frac{\rho'_i}{\bar{\rho}_i} \right) A^* \left(\frac{\bar{A}_i}{\bar{A}^*} \right) \bar{U}_i \left(\frac{\bar{U}_i}{\bar{U}_i} \right) \cos(\alpha_{ch} - \beta_i) \right] \\ & = - \left(\bar{\rho}_i \left(\frac{\bar{\rho}_{ave}}{\bar{\rho}_i} \right) \left(A^* b \frac{\bar{U}_i}{b} \right) \frac{b_i \partial \bar{V}'_i}{A^* b \bar{U}_i \partial t} + \frac{\bar{A}^* b \bar{V}_i}{\bar{A}^* b} \bar{\rho}_i \frac{\bar{U}_i}{b} \frac{b_i \partial \rho' V}{\bar{U}_i \partial t \partial \bar{\rho}_i} \right) \end{aligned}$$

Dividing through by $\bar{\rho}_i \bar{A}^* \bar{U}_i$ gives

$$\begin{aligned}
& [\rho'^* \bar{a}^* + \bar{\rho}^* A' \bar{a}^* + \bar{\rho}^* a'^*] - [\bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& \quad \bar{A}_i \cos (\alpha_{ch} - \beta_i) + \rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i)] \\
& = - \left[\bar{\rho} \frac{\partial V'_1}{\partial t} + \bar{V}_1 \frac{\partial \rho'}{\partial t} \right]_i, \tag{A18}
\end{aligned}$$

Consider next the momentum equation (Equation A10). The result is, after dividing by $\bar{\rho} \bar{U}_i^2 \bar{A}_i^*$

$$\begin{aligned}
& (p'_i \bar{A}_i + \bar{p}_i A'_i - p'^* - \bar{p}^* A'^*) \\
& = [a'^* \bar{\rho}^* \bar{a}^* + \bar{a}^* (\bar{\rho}^* a'^* + \bar{\rho}^* A' \bar{a}^* + \rho'^* \bar{a}^*)] - \\
& \quad [(u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \bar{A}_i \cos (\alpha_{ch} - \beta_i) + \\
& \quad \cos (\alpha_{ch} - \beta_i) (\rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i) + A'_i \cos (\alpha_{ch} - \beta_i) + \\
& \quad \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}))] + \left(\bar{U} \bar{V}_1 \frac{\partial \rho'}{\partial t} \right)_i + \\
& \quad \left(\bar{\rho} \bar{U} \frac{\partial V'_1}{\partial t} \right)_i + \left(\bar{\rho} \bar{V}_1 \frac{\partial U'}{\partial t} \right)_i, \tag{A19}
\end{aligned}$$

Proceeding next to the energy equation (Equation A17):

$$\begin{aligned}
& \frac{\bar{\rho}_i \bar{A}^* \bar{U}_i^3}{\gamma(\gamma-1)} [2\bar{A}_i \cos (\alpha_{ch} - \beta_i) \bar{a}_i a'_i + (\rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i) + \\
& \quad A'_i \cos (\alpha_{ch} - \beta_i) + \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch})) \bar{a}_i^2] + \\
& \quad \frac{\bar{\rho}_i \bar{A}^* \bar{U}_i^3}{2} [2\bar{A}_i \cos^2 (\alpha_{ch} - \beta_i) (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& \quad \cos^2 (\alpha_{ch} - \beta_i) (\rho'_i \bar{A}_i \cos (\alpha_{ch} - \beta_i) + A'_i \cos (\alpha_{ch} - \beta_i) + \\
& \quad \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}))] - \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma - 1)} \bar{\rho}_i \bar{A}^* \bar{U}_i^3 [2\bar{a}^{*2} a'^* \bar{\rho}^* + \\
& \quad \bar{a}^{*2} (\rho'^* \bar{a}^* + \bar{\rho}^* A' \bar{a}^* + \bar{\rho}^* a'^*)] \\
& = \frac{\bar{\rho}_i \bar{A}^* \bar{U}_i^3}{\gamma(\gamma-1)} \left[\bar{V} \bar{a}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'_1}{\partial t} + 2\bar{\rho} \bar{a} \bar{V}_1 \frac{\partial a'}{\partial t} \right]_i + \\
& \quad \frac{\rho'_i \bar{A}^* \bar{U}_i^3}{2} \left[\bar{V}_1 \bar{U}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{U}^2 \frac{\partial V'_1}{\partial t} + 2\bar{U} \bar{\rho} \bar{V}_1 \frac{\partial U'}{\partial t} \right]_i,
\end{aligned}$$

Dividing the above equation through by $\bar{\rho}_i \bar{A}^* \bar{U}_i^3$ and rearranging the terms yields

$$\begin{aligned}
& \left[\frac{\bar{a}_i^2}{\gamma(\gamma-1)} + \frac{\cos^2(\alpha_{ch} - \beta_i)}{2} \right] \left[\rho'_i \bar{A}_i \cos(\alpha_{ch} - \beta_i) + A'_i \cos(\alpha_{ch} - \beta_i) + \right. \\
& \left. \bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) \right] + \frac{2\bar{A}_i \cos(\alpha_{ch} - \beta_i) \bar{a}_i a'_i}{\gamma(\gamma-1)} + \\
& \bar{A}_i \cos^2(\alpha_{ch} - \beta_i) (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) - \\
& \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} \left[\bar{a}^{*2} (\rho'^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + 3\bar{\rho}^* a'^*) \right] \\
& = \frac{1}{\gamma(\gamma-1)} \left[\bar{V}_i \bar{a}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'_i}{\partial t} + 2\bar{a} \bar{\rho} \bar{V}_i \frac{\partial a'}{\partial t} \right]_i + \\
& \frac{1}{2} \left[\bar{V}_i \bar{U}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{U}^2 \frac{\partial V'}{\partial t} + 2\bar{U} \bar{V}_i \frac{\partial U'}{\partial t} \right]_i, \tag{A20}
\end{aligned}$$

The nondimensional equations of motion are as follows

Continuity

$$\begin{aligned}
& [\rho'^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + \bar{\rho}^* a'^*] - [\bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
& A'_i \cos(\alpha_{ch} - \beta_i) + \rho'_i \bar{A}_i \cos(\alpha_{ch} - \beta_i)] \\
& = - \left[\bar{\rho} \frac{\partial V'_i}{\partial t} + \bar{V}_i \frac{\partial \rho'}{\partial t} \right]_i, \tag{A18}
\end{aligned}$$

Momentum

$$\begin{aligned}
& [\rho'_i \bar{A}_i + \bar{\rho}_i A'_i - \rho'^* - \bar{\rho}^* A'^*] = \bar{a}^* [2\bar{\rho}^* a'^* + \bar{\rho}^* A'^* \bar{a}^* + \rho'^* \bar{a}^*] - \\
& \cos(\alpha_{ch} - \beta_i) [\rho'_i \bar{A}_i \cos(\alpha_{ch} - \beta_i) + A'_i \cos(\alpha_{ch} - \beta_i) + 2\bar{A}_i (u'_i \sin \alpha_{ch} + \\
& v'_i \cos \alpha_{ch})] + \bar{U}_i \bar{V}_i \frac{\partial \rho'_i}{\partial t} + \bar{\rho}_i \bar{U}_i \frac{\partial V'_i}{\partial t} + \bar{\rho}_i \bar{V}_i \frac{\partial U'_i}{\partial t} \tag{A19}
\end{aligned}$$

Energy

$$\begin{aligned}
& \left[\frac{\bar{a}_1^2}{\gamma(\gamma-1)} + \frac{\cos^2(\alpha_{ch} - \beta_1)}{2} \right] \left[\rho'_1 \bar{A}_1 \cos(\alpha_{ch} - \beta_1) + A'_1 \cos(\alpha_{ch} - \beta_1) + \right. \\
& \left. \bar{A}_1 (u'_1 \sin \alpha_{ch} + v'_1 \cos \alpha_{ch}) \right] + \frac{2\bar{A}_1 \cos(\alpha_{ch} - \beta_1) \bar{a}_1 a'_1}{\gamma(\gamma-1)} + \\
& \bar{A}_1 \cos^2(\alpha_{ch} - \beta_1) (u'_1 \sin \alpha_{ch} + v'_1 \cos \alpha_{ch}) - \\
& \frac{(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)} \left[\bar{a}^{*2} (\rho'^* \bar{a}^* + \bar{\rho}^* A'^* \bar{a}^* + 3\bar{\rho}^* a'^*) \right] \\
= & \frac{1}{\gamma(\gamma-1)} \left[\bar{V}_1 \bar{a}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{a}^2 \frac{\partial V'_1}{\partial t} + 2\bar{a} \bar{\rho} \bar{V}_1 \frac{\partial a'}{\partial t} \right]_1 + \\
& \frac{1}{2} \left[\bar{V}_1 \bar{U}^2 \frac{\partial \rho'}{\partial t} + \bar{\rho} \bar{U}^2 \frac{\partial V'}{\partial t} + 2\bar{U} \bar{V}_1 \frac{\partial U'}{\partial t} \right]_1 \tag{A20}
\end{aligned}$$

The small perturbation forms of the equations of motion for Region 2 and 3 are derived in a similar manner.

APPENDIX B
SOLUTIONS TO THE EQUATIONS OF MOTION FOR REGION 1

The nondimensionalized linear equations of motion for Region 1 are represented by Equations 24, 25, and 26, and are solved to obtain the three unknowns: $A_{-\infty}$, ρ^* , and a^* . Starting with the inlet flow parameters

$$u'_i = \frac{\partial \Phi'}{\partial x^*} \quad v'_i = \frac{\partial \Phi'}{\partial y^*}$$

The expression for Φ is given in Equation 11. The inlet velocity perturbations become

$$u'_i = \frac{i B_1 A_{-\infty}}{b U_1} \exp i(B_1 x + C y + k t) \quad (\text{nondimensional}) \quad (\text{B1})$$

$$v'_i = \frac{i C A_{-\infty}}{b U_1} \exp i(B_1 x + C y + k t) \quad (\text{nondimensional}) \quad (\text{B2})$$

The inlet pressure perturbation can be obtained from Bernoulli's relationship in the following manner:

$$\frac{\partial p_i}{\partial t} = -\bar{\rho}_i \left[\underbrace{\frac{\partial^2 \Phi'_{-\infty}}{\partial t^2}}_1 + \underbrace{\bar{u}_i \frac{\partial u'_i}{\partial t}}_2 + \underbrace{\bar{v}_i \frac{\partial v'_i}{\partial t}}_3 \right]$$

or

$$\int \partial p'_i = -\bar{\rho}_i \int \left(\frac{\partial^2 \Phi'_{-\infty}}{\partial t^2} + \bar{u}_i \frac{\partial u'_i}{\partial t} + \bar{v}_i \frac{\partial v'_i}{\partial t} \right) \partial t \quad (\text{B3})$$

Substituting for $\Phi'_{-\infty}$, u'_i , and v'_i , and integrating produces

$$p'_i = -i [A_{-\infty} \exp i(B_1 x + C y + k t)] [k + \bar{u}_i B_1 + \bar{v}_i C] \quad (\text{B4})$$

where all quantities are nondimensional. The inlet density perturbation can be related to the inlet pressure for isentropic flow

$$\partial \rho'_i = \frac{1}{\bar{a}^2} p'_i \quad (\text{B5})$$

Substituting Equation B4 into Equation B5 gives

$$\rho'_i = \frac{-i}{\bar{a}_i^2} [A_{-\infty} \exp i(B_1 x + C y + k t)] [k + \bar{u}_i B_1 + \bar{v}_i C] \quad (\text{B6})$$

The speed of sound perturbation at the inlet is related to p_i by

$$a'_i = \frac{\gamma - 1}{2} \left(\frac{p'_i}{\rho_i \bar{a}_i} \right) \quad (\text{B7})$$

Substituting Equation B4 into Equation B7 gives

$$a'_i = \frac{i(\gamma - 1)}{2\bar{a}_i} [A_{-\infty} \exp i(B_1x + Cy + kt)] [k + \bar{u}_i B_1 + \bar{v}_i C]$$

The inlet flow parameters represented by a two-dimensional flow field which is converted to a one-dimensional flow field by a technique commonly used in turbulent channel flows. Here the relationship holds.

$$\bar{f}'(x) = \frac{1}{\bar{y}_i} \int_0^{\bar{y}_i} f(x, y) dy \quad (\text{B8})$$

where f is any flow parameter. Starting with the velocities

$$\begin{aligned} u'(x, y) &= i B_1 A_{-\infty} \exp i(B_1x + Cy + kt) \\ \bar{u}'(x) &= \frac{1}{\bar{y}_i} \int_0^{\bar{y}_i} i B_1 A_{-\infty} \exp i(B_1x + Cy + kt) dy \\ \Rightarrow \bar{u}'(x=0) &= \frac{B_1 A_{-\infty}}{C\bar{y}_i} e^{ikt} (e^{iC\bar{y}_i} - 1) \end{aligned} \quad (\text{B9})$$

The following steps are used to evaluate the y component of the velocity:

$$\begin{aligned} v'(x, y) &= i C A_{-\infty} \exp i(B_1x + Cy + kt) \\ \bar{v}'(x=0) &= \frac{1}{\bar{y}_i} \int_0^{\bar{y}_i} i C A_{-\infty} \exp i(B_1x + Cy + kt) dy \\ \Rightarrow \bar{v}'_i &= \frac{A_{-\infty} e^{ikt} (e^{iC\bar{y}_i} - 1)}{\bar{y}_i} \end{aligned} \quad (\text{B10})$$

Returning to the inlet pressure expression noted in Equation B4, the following results:

$$\bar{p}'_i = - \frac{(k + \bar{u}_i B_1 + \bar{v}_i C)}{C\bar{y}_i} [A_{-\infty} e^{ikt} (e^{iC\bar{y}_i} - 1)] \quad (\text{B11})$$

Returning to the inlet density expression, the following results:

$$\bar{\rho}'_i = - \frac{(k + \bar{u}_i B_1 + \bar{v}_i C)}{C\bar{y}_i \bar{a}_i^2} [A_{-\infty} e^{ikt} (e^{iC\bar{y}_i} - 1)] \quad (\text{B12})$$

Finally, the speed of sound expression becomes

$$\bar{a}'_i = - \left(\frac{\gamma - 1}{2\bar{a}_i} \right) \left(\frac{1}{C\bar{y}_i} \right) [A_{-\infty} e^{ikt} (e^{iC\bar{y}_i} - 1)] [k + \bar{u}_i B_1 + \bar{v}_i C] \quad (\text{B13})$$

The next step assumes that all flow parameters within the blade channel are harmonically varying with time. As such, the following equations apply:

$$\begin{aligned}
p'^* &= \bar{p}'^* \exp ikt \\
\rho'^* &= \bar{\rho}'^* \exp ikt \\
a'^* &= \bar{a}'^* \exp ikt \\
A'^* &= \bar{A}'^* \exp ikt \\
V' &= \bar{V}' \exp ikt
\end{aligned} \tag{B14}$$

In addition to these expressions, the time rate of change of density, velocity, and speed of sound are also needed. Starting with the velocity at the inlet from Equation B9, the sequence progresses as follows:

$$\bar{u}'_1 = \frac{B_1 A_{-\infty}}{C \bar{y}_1} \exp ikt (e^{iC\bar{y}_1} - 1) \tag{B9}$$

$$\frac{\partial \bar{u}'_1}{\partial t} = \frac{ik B_1 A_{-\infty}}{C \bar{y}_1} \exp ikt (e^{iC\bar{y}_1} - 1) \tag{B14}$$

$$\bar{v}'_1 = \frac{A_{-\infty} \exp ikt (e^{iC\bar{y}_1} - 1)}{\bar{y}_1} \tag{B10}$$

$$\frac{\partial \bar{v}'_1}{\partial t} = \frac{ik A_{-\infty}}{\bar{y}_1} \exp ikt (e^{iC\bar{y}_1} - 1) \tag{B15}$$

The time rate of change of density at the inlet equals

$$\frac{\partial \bar{\rho}'_1}{\partial t} = \frac{-ik (k + \bar{u}_1 B_1 + \bar{v}_1 C)}{C \bar{y}_1 \bar{a}_1^2} [A_{-\infty} \exp ikt (e^{iC\bar{y}_1} - 1)] \tag{B16}$$

and at the throat

$$\frac{\partial \rho'^*}{\partial t} = ik \bar{\rho}'^* \exp ikt \tag{B17}$$

The time rate of change of the speed of sound at the inlet equals

$$\frac{\partial \bar{a}'_1}{\partial t} = -ik \left(\frac{\gamma - 1}{2 \bar{a}'_1} \right) \left(\frac{1}{C \bar{y}_1} \right) [A_{-\infty} \exp ikt (e^{iC\bar{y}_1} - 1)] [k + \bar{u}_1 B_1 + \bar{v}_1 C] \tag{B18}$$

and at the throat

$$\frac{\partial a'^*}{\partial t} = ik \bar{a}'^* \exp ikt \tag{B19}$$

Working first with the energy equation, let

$$E_1 = \left[\frac{\bar{a}_1^2}{\gamma(\gamma-1)} + \frac{\cos^2(\alpha_{ch} - \beta_1)}{2} \right]$$

$$I = \cos(\alpha_{ch} - \beta_1)$$

$$E_2 = \frac{\bar{a}^{*2}(\gamma^2 - \gamma + 2)}{2\gamma(\gamma-1)}$$

Inserting Equations B9 through B19 into Equation 26 and dividing through by e^{ikt} produces the following expression:

$$\begin{aligned} & E_1 \left[-I\bar{A}_1 \left(\frac{\mathbf{k} + \bar{u}_1\mathbf{B}_1 + \bar{v}_1\mathbf{C}}{C\bar{y}_1\bar{a}_1^2} \right) A_{-\infty}(e^{iC\bar{y}_1} - 1) + \bar{A}'_1 I + \right. \\ & \left. \bar{A}_1 \left(\sin\alpha_{ch} \left(\frac{B_1 A_{-\infty}}{C\bar{y}_1} \right) (e^{iC\bar{y}_1} - 1) + \left(\frac{A_{-\infty}}{\bar{y}_1} \right) \cos\alpha_{ch}(e^{iC\bar{y}_1} - 1) \right) \right] - \\ & \frac{2\bar{A}_1 I \bar{a}_1}{\gamma(\gamma-1)} \left(\frac{\gamma-1}{2\bar{a}_1} \right) \left(\frac{1}{C\bar{y}_1} \right) (\mathbf{k} + \bar{u}_1\mathbf{B}_1 + \bar{v}_1\mathbf{C}) A_{-\infty}(e^{iC\bar{y}_1} - 1) + \\ & \bar{A}_1 I^2 \left[\sin\alpha_{ch} \left(\frac{B_1 A_{-\infty}}{C\bar{y}_1} \right) (e^{iC\bar{y}_1} - 1) + \left(\frac{A_{-\infty}}{\bar{y}_1} \right) \cos\alpha_{ch}(e^{iC\bar{y}_1} - 1) \right] - \\ & E_2 [\bar{\rho}'^* \bar{a}^* + \bar{\rho}^* \bar{A}'^* \bar{a}^* + 3\bar{\rho}^* \bar{a}^*] \\ & = \frac{1}{\gamma(\gamma-1)} \left\{ \frac{\bar{V}_1 \bar{a}^2}{2} \left[ik \left(\bar{\rho}'^* - \left(\frac{\mathbf{k} + \bar{u}_1\mathbf{B}_1 + \bar{v}_1\mathbf{C}}{C\bar{y}_1\bar{a}_1^2} \right) A_{-\infty}(e^{iC\bar{y}_1} - 1) \right) \right] + \right. \\ & ik\bar{\rho}\bar{a}^2\bar{V}'_1 + \bar{\rho}\bar{a}\bar{V}_1 \left[ik \left(\bar{a}'^* - \left(\frac{\gamma-1}{2\bar{a}_1} \right) \left(\frac{1}{C\bar{y}_1} \right) A_{-\infty}(e^{iC\bar{y}_1} - 1) \right) \right. \\ & \left. \left. (\mathbf{k} + \bar{u}_1\mathbf{B}_1 + \bar{v}_1\mathbf{C}) \right) \right] \Bigg\}_1 + \frac{1}{2} \left\{ \frac{\bar{U}^2 \bar{V}_1}{2} \left[ik (\bar{\rho}'^* - (\mathbf{k} + \bar{u}_1\mathbf{B}_1 + \bar{v}_1\mathbf{C}) \right. \right. \\ & \left. \left. \left(\frac{A_{-\infty}}{C\bar{y}_1\bar{a}_1^2} \right) (e^{iC\bar{y}_1} - 1) \right) \right] + ik\bar{\rho}\bar{U}^2\bar{V}'_1 + \bar{\rho}\bar{U}\bar{V}_1 \left[ik \left(\bar{a}'^* + \right. \right. \\ & \left. \left. \left(\frac{B_1 A_{-\infty}}{C\bar{y}_1} \right) (e^{iC\bar{y}_1} - 1) \sin\alpha_{ch} + \left(\frac{A_{-\infty}}{\bar{y}_1} \right) (e^{iC\bar{y}_1} - 1) \cos\alpha_{ch} \right) \right] \Bigg\}_1, \end{aligned} \quad (B20)$$

The E, M, and C constants presented in this section are found in Appendix E.

The next step requires that Equation 220 be broken down into its real and imaginary parts. The solution deals first with the four terms containing the unknown complex constant $A_{-\infty}$. These are broken into real and imaginary parts as follows:

$$1. \quad B_1 A_{-\infty} (e^{iC\bar{y}_i} - 1)$$

$$B_1 = B_{1R} + iB_{1i}$$

$$A_{-\infty} = A_{-\infty R} + iA_{-\infty i}$$

$$e^{iC\bar{y}_i} = \cos(C\bar{y}_i) + i \sin(C\bar{y}_i)$$

Thus, the real part of the expression is

$$B_{1R} A_{-\infty R} (\cos(C\bar{y}_i) - 1) - B_{1i} A_{-\infty i} (\cos(C\bar{y}_i) - 1) -$$

$$B_{1R} A_{-\infty i} \sin(C\bar{y}_i) - B_{1i} A_{-\infty R} \sin(C\bar{y}_i)$$

and the imaginary part of the expression is

$$B_{1R} A_{-\infty i} (\cos(C\bar{y}_i) - 1) + B_{1i} A_{-\infty R} (\cos(C\bar{y}_i) - 1) +$$

$$B_{1R} A_{-\infty R} \sin(C\bar{y}_i) - B_{1i} A_{-\infty i} \sin(C\bar{y}_i)$$

$$2. \quad A_{-\infty} (e^{iC\bar{y}_i} - 1)$$

The real part of the expression is

$$A_{-\infty R} (\cos(C\bar{y}_i) - 1) - A_{-\infty i} \sin(C\bar{y}_i)$$

The imaginary part of the expression is

$$A_{-\infty i} (\cos(C\bar{y}_i) - 1) + A_{-\infty R} \sin(C\bar{y}_i)$$

$$3. \quad ikB_1 A_{-\infty} (e^{iC\bar{y}_i} - 1)$$

The real part of the expression is

$$-k[B_{1R} A_{-\infty i} (\cos(C\bar{y}_i) - 1) + B_{1i} A_{-\infty R} (\cos(C\bar{y}_i) - 1) +$$

$$B_{1R} A_{-\infty R} \sin(C\bar{y}_i) - B_{1i} A_{-\infty i} \sin(C\bar{y}_i)]$$

The imaginary part of the expression is

$$k[B_{1R} A_{-\infty R} (\cos(C\bar{y}_i) - 1) - B_{1i} A_{-\infty i} (\cos(C\bar{y}_i) - 1) -$$

$$B_{1R} A_{-\infty i} \sin(C\bar{y}_i) - B_{1i} A_{-\infty R} \sin(C\bar{y}_i)]$$

$$4. \quad ikA_{-\infty} (e^{iC\bar{y}_i} - 1)$$

The real part of the expression is

$$-k[A_{-\infty I} (\cos(C\bar{y}_i) - 1) + A_{-\infty R} \sin(C\bar{y}_i)]$$

The imaginary part of the expression is

$$k[A_{-\infty R} (\cos(C\bar{y}_i) - 1) - A_{-\infty I} \sin(C\bar{y}_i)]$$

Now let,

$$E_3 = - \frac{E_1 \bar{I} \bar{A}_i \bar{u}_i}{C\bar{y}_i \bar{a}_i^2} + \frac{E_1 \bar{A}_i \sin \alpha_{ch}}{C\bar{y}_i} - \frac{2\bar{A}_i \bar{I} \bar{u}_i}{2\gamma C\bar{y}_i} + \frac{\bar{A}_i \bar{I}^2 \sin \alpha_{ch}}{C\bar{y}_i} \quad (B21)$$

$$E_4 = - \frac{E_1 \bar{I} \bar{A}_i (k + \bar{v}_i C)}{C\bar{y}_i \bar{a}_i^2} + \frac{E_1 \bar{A}_i \cos \alpha_{ch}}{\bar{y}_i} - \frac{2\bar{A}_i \bar{I} (k + \bar{v}_i C)}{2\gamma C\bar{y}_i} + \frac{\bar{A}_i \bar{I}^2 \cos \alpha_{ch}}{\bar{y}_i} \quad (B22)$$

$$E_5 = \frac{-\bar{a}_i^2 \bar{V}_i \bar{u}_i}{2\gamma(\gamma-1)C\bar{y}_i \bar{a}_i^2} - \frac{\bar{\rho}_i \bar{a}_i \bar{V}_i \bar{u}_i}{\gamma(\gamma-1)} \left(\frac{\gamma-1}{2\bar{a}_i C\bar{y}_i} \right) - \frac{\bar{U}_i^2 \bar{V}_i \bar{u}_i}{4C\bar{y}_i \bar{a}_i^2} + \frac{\bar{\rho}_i \bar{U}_i \bar{V}_i \sin \alpha_{ch}}{2C\bar{y}_i} \quad (B23)$$

$$E_6 = \frac{-\bar{a}_i^2 \bar{V}_i (k + \bar{v}_i C)}{2\gamma(\gamma-1)C\bar{y}_i \bar{a}_i^2} - \bar{\rho}_i \bar{a}_i \bar{V}_i \left(\frac{\gamma-1}{2C\bar{y}_i \bar{a}_i} \right) \left(\frac{k + \bar{v}_i C}{\gamma(\gamma-1)} \right) + \frac{\bar{U}_i^2 \bar{V}_i (k + \bar{v}_i C)}{4C\bar{y}_i \bar{a}_i^2} + \frac{\bar{\rho}_i \bar{U}_i \bar{V}_i \cos \alpha_{ch}}{2\bar{y}_i} \quad (B24)$$

Substituting the relationships for the real part of the A_{∞} terms, along with Equations B21 through B24 gives

$$\begin{aligned} & E_3 [B_{IR} A_{\infty R} (\cos(C\bar{y}_i) - 1) - B_{II} A_{\infty I} (\cos(C\bar{y}_i) - 1) - \\ & B_{IR} A_{\infty I} \sin(C\bar{y}_i) - B_{II} A_{\infty R} \sin(C\bar{y}_i)] + \\ & E_4 [A_{\infty R} (\cos(C\bar{y}_i) - 1) - A_{\infty I} \sin(C\bar{y}_i)] + \\ & E_5 \bar{A}'_{IR} \bar{I} - E_2 (\bar{\rho}'_R \bar{a}'^* + \bar{\rho}^* \bar{A}'_R \bar{a}'^* + 3\bar{\rho}^* \bar{a}'_R) \\ & = -kE_5 [B_{IR} A_{\infty I} (\cos(C\bar{y}_i) - 1) + B_{II} A_{\infty R} (\cos(C\bar{y}_i) - 1) + \\ & B_{IR} A_{\infty R} \sin(C\bar{y}_i) - B_{II} A_{\infty I} \sin(C\bar{y}_i)] + \\ & E_6 [-k(A_{\infty I} (\cos(C\bar{y}_i) - 1) + A_{\infty R} \sin(C\bar{y}_i))] \\ & - \left(\frac{\bar{a}_i^2 \bar{V}}{2\gamma(\gamma-1)} \right)_I (k\bar{\rho}'_I)^* - \bar{V}'_{II} k \left(\frac{\bar{\rho} \bar{a}_i^2}{\gamma(\gamma-1)} \right)_I - k\bar{a}'_I \left(\frac{\bar{\rho} \bar{a}_i \bar{V}}{\gamma(\gamma-1)} \right)_I \\ & - k\bar{\rho}'_I \left(\frac{\bar{U}^2 \bar{V}}{4} \right)_I - k\bar{V}'_{II} \left(\frac{\bar{\rho} \bar{U}^2}{2} \right)_I - k\bar{a}'_I \left(\frac{\bar{\rho} \bar{U} \bar{V}}{2} \right)_I \end{aligned} \quad (B25)$$

Collecting terms and rearranging Equation B25 produces:

$$\begin{aligned}
E_9 A_{-\infty R} + E_9 A_{-\infty I} &= -E_1 \bar{A}'_{IR} I + \dots \\
&\left(\frac{\bar{a}^2 \bar{V}_1}{2\gamma(\gamma-1)} + \frac{U^2 \bar{V}_1}{4} \right)_1 k \bar{\rho}'_1^* - \\
&k \bar{a}'_1^* \left(\frac{\bar{\rho} \bar{a} \bar{V}_1}{\gamma(\gamma-1)} + \frac{\rho U \bar{V}_1}{2} \right)_1 - \\
&k \bar{V}'_{1I} \left(\frac{\bar{\rho} U^2}{2} + \frac{\bar{\rho} \bar{a}^2}{\gamma(\gamma-1)} \right)_1 \tag{B26} \\
&+ E_2 (\bar{a}^* \bar{\rho}'_R^* + \bar{\rho}^* \bar{A}'_R^* \bar{a}^* + 3 \bar{a}'_R^* \bar{\rho}^*)
\end{aligned}$$

with

$$\begin{aligned}
E_z &= \cos(C\bar{y}_i) - 1 \\
E_x &= E_t (E_r B_{IR} - B_{II} \sin(C\bar{y}_i)) + E_s E_r + E_r k (B_{II} E_r + \\
&\quad B_{IR} \sin(C\bar{y}_i)) + E_r k \sin(C\bar{y}_i) \\
E_y &= -E_t (B_{II} E_r + B_{IR} \sin(C\bar{y}_i)) - E_s \sin(C\bar{y}_i) + \\
&\quad E_r k (B_{IR} E_r - B_{II} \sin(C\bar{y}_i)) + E_r k E_r
\end{aligned}$$

Equation B26 represents the real component of Equation B20. The imaginary component is

$$\begin{aligned}
&E_1 [B_{IR} A_{-I} E_r + B_{II} A_{-OI} E_r + B_{IR} A_{-OR} \sin(C\bar{y}_i) - \\
&\quad B_{II} A_{-OI} \sin(C\bar{y}_i)] + E_1 [A_{-I} E_r + A_{-OR} \sin(C\bar{y}_i)] + \\
&\quad E_1 \bar{A}'_{II} I - E_2 [\bar{\rho}'_1^* \bar{a}^* + \bar{\rho}^* \bar{A}'_1^* \bar{a}^* + 3 \bar{\rho}^* \bar{a}'_1^*] \\
&= E_3 k [B_{IR} A_{-OR} E_r - B_{II} A_{-OI} E_r - B_{IR} A_{-OI} \sin(C\bar{y}_i) - \\
&\quad B_{II} A_{-OR} \sin(Cy)] + E_r k [A_{-OR} E_r - A_{-OI} \sin(C\bar{y}_i)] + \\
&\quad k \bar{\rho}'_R^* \left(\frac{\bar{a}^2 \bar{V}_1}{2\gamma(\gamma-1)} + \frac{U^2 \bar{V}_1}{4} \right)_1 + \\
&\quad k \bar{a}'_R^* \left(\frac{\bar{\rho} \bar{a} \bar{V}_1}{\gamma(\gamma-1)} + \frac{\bar{\rho} U \bar{V}_1}{2} \right)_1 \tag{B27} \\
&+ k \left[\frac{\bar{\rho} \bar{a}^2}{\gamma(\gamma-1)} + \frac{\bar{\rho} U^2}{2} \right]_1 \bar{V}'_{IR}
\end{aligned}$$

Again, collecting terms for $A_{-\infty l}$ and $A_{-\infty R}$, let

$$\begin{aligned} E_{10} &= E_3 [B_{1l}E_7 + B_{1R} \sin (C\bar{y}_l)] + E_4 \sin (C\bar{y}_l) - \\ &E_5 k [B_{1R}E_7 - B_{1l} \sin (C\bar{y}_l)] - E_6 k E_7 \\ E_{11} &= E_3 [B_{1R}E_7 - B_{1l} \sin (C\bar{y}_l)] + E_4 E_7 + \\ &E_5 k [B_{1l}E_7 + B_{1R} \sin (C\bar{y}_l)] + E_6 k \sin (C\bar{y}_l) \end{aligned}$$

Substituting the above relationships into Equation B27 gives

$$\begin{aligned} E_{10}A_{-\infty R} + E_{11}A_{-\infty l} &= -E_l \bar{A}'_{il} I + E_2 [\bar{\rho}'_l \bar{a}^* + \bar{\rho}^* \bar{A}'_l \bar{a}^* + 3\bar{\rho}^* \bar{a}'_l] + \\ k\bar{\rho}'_R & \left(\frac{\bar{a}^2 \bar{V}_l}{2\gamma(\gamma-1)} + \frac{\bar{U}^2 \bar{V}_l}{4} \right)_l + k\bar{V}'_{1R} \left(\frac{\bar{\rho} \bar{a}^2}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{U}^2}{2} \right)_l + \\ k\bar{a}'_R & \left(\frac{\bar{\rho} \bar{a} \bar{V}_l}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{U} \bar{V}_l}{2} \right)_l \end{aligned} \quad (B28)$$

The next step involves the combination of Equation B26 with Equation B28 and solving for $A_{-\infty R}$ and $A_{-\infty l}$. The sequence requires that Equation B26 be divided by E_8 and Equation B28 be divided by E_{10} , as shown in Equations B29 and B30.

$$\begin{aligned} A_{-\infty R} + A_{-\infty l} \left(\frac{E_9}{E_8} \right) &= -\bar{A}'_{iR} I \left(\frac{E_1}{E_8} \right) + (\bar{\rho}'_R \bar{a}^* + \bar{\rho}^* \bar{A}'_R \bar{a}^* + \\ 3\bar{\rho}^* \bar{a}'_R) & \left(\frac{E_2}{E_8} \right) - \bar{\rho}'_l \left(\frac{E_{12}}{E_8} \right) - \bar{V}'_{1l} \left(\frac{E_{13}}{E_8} \right) - \bar{a}'_l \left(\frac{E_{14}}{E_8} \right) \end{aligned} \quad (B29)$$

$$\begin{aligned} A_{-\infty R} + A_{-\infty l} \left(\frac{E_1}{E_{10}} \right) &= -\bar{A}'_{il} I \left(\frac{E_1}{E_{10}} \right) + \\ (\bar{\rho}'_l \bar{a}^* + \bar{\rho}^* \bar{A}'_l \bar{a}^* + 3\bar{\rho}^* \bar{a}'_l) & \left(\frac{E_2}{E_{10}} \right) + \\ \bar{\rho}'_R \left(\frac{E_{12}}{E_{10}} \right) + \bar{V}'_{1R} \left(\frac{E_{13}}{E_{10}} \right) + \bar{a}'_R \left(\frac{E_{14}}{E_{10}} \right) & \end{aligned} \quad (B30)$$

where,

$$\begin{aligned} E_{12} &= \left(\frac{\bar{a}^2 \bar{V} k}{2\gamma(\gamma-1)} + \frac{\bar{U}^2 \bar{V} k}{4} \right)_l \\ E_{13} &= \left(\frac{\bar{\rho} \bar{a}^2 k}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{U}^2 k}{2} \right)_l \\ E_{14} &= \left(\frac{\bar{\rho} \bar{a} \bar{V} k}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{U} \bar{V} k}{2} \right)_l \end{aligned}$$

Subtracting Equation B30 from Equation B29 yields

$$\begin{aligned}
A_{-\infty I} = & \left[-E_1 I \left(\frac{\bar{A}'_{iR}}{E_8} - \frac{\bar{A}'_{ii}}{E_{10}} \right) + E_2 \left[\bar{a}^* \left(\frac{\bar{\rho}'_R}{E_8} - \frac{\bar{\rho}'_I}{E_{10}} \right) + \right. \right. \\
& \left. \left. \bar{\rho}^* \bar{a}^* \left(\frac{\bar{A}'_{R^*}}{E_8} - \frac{\bar{A}'_{I^*}}{E_{10}} \right) + 3\bar{\rho}^* \left(\frac{\bar{a}'_R}{E_8} - \frac{\bar{a}'_I}{E_{10}} \right) \right] - \right. \\
& E_{12} \left(\frac{\bar{\rho}'_I}{E_8} - \frac{\bar{\rho}'_R}{E_{10}} \right) - E_{13} \left(\frac{\bar{V}'_{II}}{E_8} - \frac{\bar{V}'_{IR}}{E_{10}} \right) - \\
& \left. E_{14} \left(\frac{\bar{a}'_I}{E_8} - \frac{\bar{a}'_R}{E_{10}} \right) \right] \div \left[\frac{E_9}{E_8} - \frac{E_{11}}{E_{10}} \right] \tag{B31}
\end{aligned}$$

Substituting Equation B31 into Equation B29 and solving for $A_{-\infty R}$ produces

$$\begin{aligned}
A_{-\infty R} = & -\bar{A}'_{iR} I \left(\frac{E_1}{E_8} \right) + (\bar{\rho}'_R \bar{a}^* + \bar{\rho}^* \bar{A}'_{R^*} \bar{a}^* + 3\bar{\rho}^* \bar{a}'_R) \left(\frac{E_2}{E_8} \right) - \\
& \bar{\rho}'_I \left(\frac{E_{12}}{E_8} \right) - \bar{V}'_{II} \left(\frac{E_{13}}{E_8} \right) - \bar{a}'_I \left(\frac{E_{14}}{E_8} \right) - A_{-\infty I} \left(\frac{E_9}{E_8} \right) \tag{B32}
\end{aligned}$$

Before proceeding further, the following relationships must be given:

$$\begin{aligned}
E_{15} &= \left(\frac{E_9}{E_8} - \frac{E_{11}}{E_{10}} \right) \\
E_{16} &= I \left(\frac{E_1}{E_8} \right) \left(\frac{E_9}{E_{15}E_8} - 1 \right) \\
E_{17} &= I \left(\frac{-E_1}{E_{10}E_{15}} \right) \left(\frac{E_8}{E_9} \right) \\
E_{18} &= \left(\frac{E_2}{E_8} \right) \left(\bar{a}^* - \frac{\bar{a}^* E_9}{E_{15}E_8} \right) - \left(\frac{E_{12}E_9}{E_8E_{10}E_{15}} \right) \\
E_{19} &= - \frac{E_{12}}{E_8} + \left(\frac{\bar{a}^* E_9}{E_{10}E_{15}} \right) \left(\frac{E_9}{E_8} \right) + \left(\frac{E_{12}}{E_8E_{15}} \right) \left(\frac{E_9}{E_8} \right) \\
E_{20} &= \bar{\rho}^* \bar{a}^* \left[\frac{E_2}{E_8} - \left(\frac{E_2}{E_8E_{15}} \right) \left(\frac{E_9}{E_8} \right) \right] \\
E_{21} &= \bar{\rho}^* \bar{a}^* \left(\frac{E_2}{E_{10}E_{15}} \right) \left(\frac{E_9}{E_8} \right) \\
E_{22} &= 3\bar{\rho}^* \left(\frac{E_2}{E_8} \right) \left(1 - \frac{E_9}{E_8E_{15}} \right) - \left(\frac{E_9E_{14}}{E_8E_{10}E_{15}} \right) \\
E_{23} &= - \frac{E_{14}}{E_8} + 3\bar{\rho}^* \left(\frac{E_2E_9}{E_8E_{10}E_{15}} \right) + \left(\frac{E_{14}}{E_8E_{15}} \right) \left(\frac{E_9}{E_8} \right) \\
E_{24} &= - \frac{E_{13}}{E_8} + \left(\frac{E_{13}}{E_8E_{15}} \right) \left(\frac{E_9}{E_8} \right) \\
E_{25} &= - \left(\frac{E_{13}}{E_{10}E_{15}} \right) \left(\frac{E_9}{E_8} \right) \tag{B33}
\end{aligned}$$

Substituting Equation B31 into Equation B32 and using the relationships detailed in Equation B33 yields

$$\begin{aligned}
A_{-\infty R} = & (E_{16}\bar{A}'_{iR} + E_{17}\bar{A}'_{iI}) + (E_{18}\bar{\rho}'_{R^*} + E_{19}\bar{\rho}'_{I^*}) + \\
& (E_{20}\bar{A}'_{R^*} + E_{21}\bar{A}'_{I^*}) + (E_{22}\bar{a}'_{R^*} + E_{23}\bar{a}'_{I^*}) + \\
& (E_{24}\bar{V}'_{iI} + E_{25}\bar{V}'_{iR})
\end{aligned} \tag{B34}$$

The nondimensional momentum equation has the form

$$\begin{aligned}
\bar{p}'_i\bar{A}_i + \bar{p}_i\bar{A}'_i - \bar{p}^*\bar{A}'^* - \bar{p}'^* &= \bar{a}^* (2\bar{\rho}^*\bar{a}'^* + \bar{\rho}^*\bar{A}'^*\bar{a}^* + \bar{\rho}'^*\bar{a}^*) - \\
& I [\bar{\rho}_i\bar{A}_i I + \bar{A}'_i I + 2\bar{A}_i (\bar{u}'_i \sin \alpha_{ch} + \bar{v}'_i \cos \alpha_{ch})] + \\
\bar{U}_i\bar{V}_i \frac{\partial \bar{\rho}'_i}{\partial t} + \bar{\rho}_i\bar{U}_i \frac{\partial \bar{V}'_i}{\partial t} + \bar{\rho}_i\bar{V}_i \frac{\partial \bar{U}'_i}{\partial t}
\end{aligned}$$

Inserting the relationships given in Equations B9 through B19 into Equation B25 and dividing through by $e^{iC\bar{y}_i}$ produces

$$\begin{aligned}
& \left[- \left(\frac{k + \bar{u}_i B_i + \bar{v}_i C}{C\bar{y}_i} \right) \right. \\
& \left. (A_{-\infty} (e^{iC\bar{y}_i} - 1)) \bar{A}_i + \bar{p}_i\bar{A}'_i - \bar{p}^*\bar{A}'^* - \bar{\rho}'^*\bar{a}^* \right] \\
& - \bar{a}^* (2\bar{\rho}^*\bar{a}'^* + \bar{\rho}^*\bar{A}'^*\bar{a}^* + \bar{\rho}'^*\bar{a}^*) - I \left\{ - \right. \\
& \left. \left(\frac{k + \bar{u}_i B_i + \bar{v}_i C}{C\bar{y}_i\bar{a}_i^2} \right) \bar{A}_i I A_{-\infty} (e^{iC\bar{y}_i} - 1) + \bar{A}'_i I + \right. \\
& 2\bar{A}_i \left[\left(\frac{B_i A_{-\infty}}{C\bar{y}_i} \sin \alpha_{ch} \right) (e^{iC\bar{y}_i} - 1) + \right. \\
& \left. \left. \left(\frac{A_{-\infty}}{\bar{y}_i} \cos \alpha_{ch} \right) (e^{iC\bar{y}_i} - 1) \right] \right\} + \\
& \frac{\bar{U}_i\bar{V}_i}{2} \left\{ ik \left[\bar{\rho}'^* - \left(\frac{k + \bar{u}_i B_i + \bar{v}_i C}{C\bar{y}_i\bar{a}_i^2} \right) \right. \right. \\
& \left. \left. A_{-\infty} (e^{iC\bar{y}_i} - 1) \right] \right\} + \bar{\rho}_i\bar{U}_i ik\bar{V}'_i + \frac{\bar{\rho}_i\bar{V}_i}{2} \left\{ ik\bar{a}'^* + \right. \\
& \left. ik \left[\frac{B_i A_{-\infty}}{C\bar{y}_i} (e^{iC\bar{y}_i} - 1) \sin \alpha_{ch} + \frac{A_{-\infty}}{\bar{y}_i} (e^{iC\bar{y}_i} - 1) \cos \alpha_{ch} \right] \right\}
\end{aligned} \tag{B35}$$

Rearrange the Equation B35 to collect $A_{-\infty}$ on right side of expression. Assume the following relationships:

$$\begin{aligned}
M_1 &= -\frac{\bar{u}_i \bar{A}_i}{C\bar{y}_i} - \frac{\bar{u}_i \bar{A}_i I^2}{C\bar{y}_i \bar{a}_i^2} + \frac{2\bar{A}_i I}{C\bar{y}_i} \sin \alpha_{ch} \\
M_2 &= -\frac{(k + \bar{v}_i C) \bar{A}_i}{C\bar{y}_i} - \frac{(k + \bar{v}_i C) \bar{A}_i I^2}{C\bar{y}_i \bar{a}_i^2} + \frac{2\bar{A}_i I^2}{\bar{y}_i} \cos \alpha_{ch} \\
M_3 &= \frac{\bar{U}_i \bar{V}_i}{2C\bar{y}_i \bar{a}_i^2} - \frac{\bar{\rho}_i \bar{V}_i \sin \alpha_{ch}}{2C\bar{y}_i} \\
M_4 &= \left(\frac{k + \bar{v}_i C}{C\bar{y}_i \bar{a}_i^2} \right) \left(\frac{\bar{U}_i \bar{V}_i}{2} \right) - \frac{\bar{\rho}_i \bar{V}_i \cos \alpha_{ch}}{2\bar{y}_i}
\end{aligned} \tag{B36}$$

Substitute the relationships of Equation B36 into Equation B35.

$$\begin{aligned}
&M_1 [B_1 A_{-\infty} (e^{iC\bar{y}_i} - 1)] + M_2 [A_{-\infty} (e^{iC\bar{y}_i} - 1)] + \\
&M_3 [ikB_1 A_{-\infty} (e^{iC\bar{y}_i} - 1)] + M_4 [ikA_{-\infty} (e^{iC\bar{y}_i} - 1)] \\
&= (-\bar{p}_i - I^2) \bar{A}_i + (\bar{p}^* + \bar{a}^2 \bar{\rho}^*) \bar{A}'^* + 2\bar{a}^{*2} \bar{\rho}'^* + \frac{ik\bar{\rho}'^* \bar{U}_i \bar{V}_i}{2} + \\
&ik\bar{\rho}_i \bar{U}_i \bar{V}_i + \frac{ik\bar{\rho}_i \bar{a}'^* \bar{V}_i}{2} + 2\bar{\rho}^* \bar{a}^* \bar{a}'^*
\end{aligned} \tag{B37}$$

The next step in the sequence involves the separation of Equation B37 into real and imaginary parts. First, assume

$$\begin{aligned}
M_5 &= -\bar{p}_i - I^2 \\
M_6 &= \bar{p}^* + \bar{\rho}^* \bar{a}^{*2}
\end{aligned} \tag{B38}$$

and

$$\begin{aligned}
M_7 &= M_1 [B_{1R} E_7 - B_{1i} \sin (C\bar{y}_i)] + M_2 E_7 - \\
&M_3 [kB_{1i} E_7 - kB_{1R} \sin (C\bar{y}_i)] - M_4 k \sin (C\bar{y}_i) \\
M_8 &= -M_1 [B_{1i} E_7 + B_{1R} \sin (C\bar{y}_i)] - M_2 \sin (C\bar{y}_i) + \\
&kM_3 [-B_{1R} E_7 + B_{1i} \sin (C\bar{y}_i)] - kM_4 E_7
\end{aligned} \tag{B39}$$

Real Part

$$\begin{aligned}
& M_1 [B_{1R}A_{-\infty R}E_7 - B_{1I}A_{-\infty I}E_7 - B_{1R}A_{-\infty I} \sin(C\bar{y}_i) - \\
& \quad B_{1I}A_{-\infty R} \sin(C\bar{y}_i)] + M_2 [A_{-\infty R}E_7 - A_{-\infty I} \sin(C\bar{y}_i)] - \\
& \quad kM_3 [B_{1R}A_{-\infty I}E_7 + B_{1I}A_{-\infty R}E_7 + B_{1R}A_{-\infty R} \sin(C\bar{y}_i) - \\
& \quad B_{1I}A_{-\infty I} \sin(C\bar{y}_i)] - kM_4 [A_{-\infty I}E_7 + A_{-\infty R} \sin(C\bar{y}_i)] \\
& = M_5 \bar{A}'_{iR} + M_6 \bar{A}'_{iR}^* + 2\bar{\rho}'_R \bar{a}^{*2} - \frac{k\bar{\rho}'_I \bar{U}_1 \bar{V}_1}{2} - \\
& \quad k\bar{\rho}_I \bar{U}_1 \bar{V}'_{iI} - \frac{k\bar{\rho}_I \bar{a}'_I \bar{V}_1}{2} + 2\bar{\rho}^* \bar{a}^* \bar{a}'_R^* \tag{B40}
\end{aligned}$$

Substituting Equation B39 into Equation B40 yields

$$\begin{aligned}
M_7 A_{-\infty R} + M_8 A_{-\infty I} & = M_5 \bar{A}'_{iR} + M_6 \bar{A}'_{iR}^* + 2\bar{\rho}'_R \bar{a}^{*2} - \\
& \quad \frac{k\bar{\rho}'_I \bar{U}_1 \bar{V}_1}{2} - k\bar{\rho}_I \bar{U}_1 \bar{V}'_{iI} - \frac{k\bar{\rho}_I \bar{a}'_I \bar{V}_1}{2} + 2\bar{\rho}^* \bar{a}^* \bar{a}'_R^* \tag{B41}
\end{aligned}$$

Imaginary Part

$$\begin{aligned}
& M_1 [B_{1R}A_{-\infty I}E_7 + B_{1I}A_{-\infty R}E_7 + B_{1R}A_{-\infty R} \sin(C\bar{y}_i) - B_{1I}A_{-\infty I} \sin(C\bar{y}_i)] + \\
& \quad M_2 [A_{-\infty I}E_7 + A_{-\infty R} \sin(C\bar{y}_i)] + kM_3 [B_{1R}A_{-\infty R}E_7 - B_{1I}A_{-\infty I}E_7 - \\
& \quad B_{1R}A_{-\infty I} \sin(C\bar{y}_i) - B_{1I}A_{-\infty R} \sin(C\bar{y}_i)] + kM_4 [A_{-\infty R}E_7 - A_{-\infty I} \sin(C\bar{y}_i)] \\
& = (-\bar{\rho}_i - I^2) \bar{A}'_{iI} + (\bar{\rho}^* + \bar{\rho}^* \bar{a}^{*2}) \bar{A}'_{iI}^* + 2\bar{\rho}'_I \bar{a}^{*2} + 2\bar{\rho}^* \bar{a}^* \bar{a}'_{iI}^* + \\
& \quad \frac{k\bar{\rho}'_R \bar{U}_1 \bar{V}_1}{2} + k\bar{\rho}_I \bar{U}_1 \bar{V}'_{iR} + \frac{k\bar{\rho}_I \bar{a}'_I \bar{V}_1}{2} \tag{B42}
\end{aligned}$$

Assume

$$\begin{aligned}
M_9 & = M_1 [B_{1I}E_7 + B_{1R} \sin(C\bar{y}_i)] + M_2 \sin(C\bar{y}_i) + \\
& \quad kM_3 [B_{1R}E_7 - B_{1I} \sin(C\bar{y}_i)] + M_4 E_7 k \\
M_{10} & = M_1 [B_{1R}E_7 - B_{1I} \sin(C\bar{y}_i)] + M_2 E_7 - \\
& \quad kM_3 [B_{1I}E_7 + B_{1R} \sin(C\bar{y}_i)] - kM_4 \sin(C\bar{y}_i) \tag{B43}
\end{aligned}$$

Substituting the relationships of Equation B43 into Equation B42 yields

$$\begin{aligned}
M_9 A_{-\infty R} + M_{10} A_{-\infty I} & = M_5 \bar{A}'_{iI} + M_6 \bar{A}'_{iI}^* + 2\bar{\rho}'_I \bar{a}^{*2} + \\
& \quad 2\bar{\rho}^* \bar{a}^* \bar{a}'_{iI}^* + \frac{k\bar{\rho}'_R \bar{U}_1 \bar{V}_1}{2} + k\bar{\rho}_I \bar{U}_1 \bar{V}'_{iR} + \frac{k\bar{\rho}_I \bar{a}'_I \bar{V}_1}{2} \tag{B44}
\end{aligned}$$

In order to solve for $A_{-\infty R}$ and $A_{-\infty I}$, make the following assumptions.

$$E_{26} = -I \left(\frac{E_1}{E_8 E_{15}} \right)$$

$$E_{27} = \frac{E_1}{E_{10} E_{15}}$$

$$E_{28} = \frac{\bar{a}^* E_2}{E_8 E_{15}} + \frac{E_{12}}{E_{10} E_{15}}$$

$$E_{29} = - \frac{\bar{a}^* E_2}{E_{10} E_{15}} - \frac{E_{12}}{E_8 E_{15}}$$

$$E_{30} = \frac{\bar{\rho}^* \bar{a}^* E_2}{E_8 E_{15}}$$

$$E_{31} = - \frac{\bar{\rho}^* \bar{a}^* E_2}{E_{10} E_{15}}$$

$$E_{32} = \frac{3\bar{\rho}^* E_2}{E_8 E_{15}} + \frac{E_{14}}{E_{10} E_{15}}$$

$$E_{33} = - \frac{3\bar{\rho}^* E_2}{E_{10} E_{15}} - \frac{E_{14}}{E_8 E_{15}}$$

$$E_{34} = - \frac{E_{13}}{E_8 E_{15}}$$

$$E_{35} = \frac{E_{13}}{E_{10} E_{15}}$$

(B45)

Substitute the relationships of Equation B45 into Equations B41 and B44 to solve for $A_{-\infty R}$ and $A_{-\infty I}$ such that

$$A_{-\infty R} = E_{16}\bar{A}'_{iR} + E_{17}\bar{A}'_{iI} + E_{18}\bar{\rho}'_{R*} + E_{19}\bar{\rho}'_{I*} + E_{20}\bar{A}'_{R*} + E_{21}\bar{A}'_{I*} + E_{22}\bar{a}'_{R*} + E_{23}\bar{a}'_{I*} + E_{24}\bar{V}'_{iI} + E_{25}\bar{V}'_{iR} \quad (B46)$$

$$A_{-\infty I} = E_{26}\bar{A}'_{iR} + E_{27}\bar{A}'_{iI} + E_{28}\bar{\rho}'_{R*} + E_{29}\bar{\rho}'_{I*} + E_{30}\bar{A}'_{R*} + E_{31}\bar{A}'_{I*} + E_{32}\bar{a}'_{R*} + E_{33}\bar{a}'_{I*} + E_{34}\bar{V}'_{iI} + E_{35}\bar{V}'_{iR} \quad (B47)$$

Substitute Equations B46 and B47 into Equation B41 to give

$$\begin{aligned} M_7 (E_{16}\bar{A}'_{iR} + E_{17}\bar{A}'_{iI} + E_{18}\bar{\rho}'_{R*} + E_{19}\bar{\rho}'_{I*} + E_{20}\bar{A}'_{R*} + E_{21}\bar{A}'_{I*} + E_{22}\bar{a}'_{R*} + E_{23}\bar{a}'_{I*} + E_{24}\bar{V}'_{iI} + E_{25}\bar{V}'_{iR}) + M_8 (E_{26}\bar{A}'_{iR} + E_{27}\bar{A}'_{iI} + E_{28}\bar{\rho}'_{R*} + E_{29}\bar{\rho}'_{I*} + E_{30}\bar{A}'_{R*} + E_{31}\bar{A}'_{I*} + E_{32}\bar{a}'_{R*} + E_{33}\bar{a}'_{I*} + E_{34}\bar{V}'_{iI} + E_{35}\bar{V}'_{iR}) \\ = M_5\bar{A}'_{iR} + M_6\bar{A}'_{R*} + 2\bar{\rho}'_{R*}\bar{a}'_{I*} - \frac{k\bar{\rho}'_{I*}\bar{U}_i\bar{V}_i}{2} - k\bar{\rho}_i\bar{U}_i\bar{V}'_{iI} - \frac{k\bar{\rho}_i\bar{a}'_{I*}\bar{V}_i}{2} + 2\bar{\rho}^*\bar{a}^*\bar{a}'_{R*} \end{aligned} \quad (B48)$$

Let;

$$\begin{aligned} M_{11} &= E_{18}M_7 + E_{28}M_8 - 2\bar{a}^*{}^2 \\ M_{12} &= E_{19}M_7 + E_{29}M_8 + \frac{k\bar{U}_i\bar{V}_i}{2} \\ M_{13} &= -E_{16}M_7 - E_{26}M_8 + M_5 \\ M_{14} &= -E_{17}M_7 - E_{27}M_8 \\ M_{15} &= -E_{20}M_7 - E_{30}M_8 + M_6 \\ M_{16} &= -E_{21}M_7 - E_{31}M_8 \\ M_{17} &= -E_{22}M_7 - E_{32}M_8 + 2\bar{\rho}^*\bar{a}^* \\ M_{18} &= E_{23}M_7 - E_{33}M_8 - \frac{k\bar{\rho}_i\bar{V}_i}{2} \\ M_{19} &= -E_{24}M_7 - E_{34}M_8 - k\bar{\rho}_i\bar{U}_i \\ M_{20} &= -E_{25}M_7 - E_{35}M_8 \end{aligned} \quad (B49)$$

Rearrange Equation B48 to collect $\bar{\rho}'_{R*}$ and $\bar{\rho}'_{I*}$ on the left side of the equation and all other terms on the right side. Substitute the relationships of Equation B49 into Equation B48.

$$\begin{aligned} M_{11}\bar{\rho}'_{R*} + M_{12}\bar{\rho}'_{I*} &= M_{13}\bar{A}'_{iR} + M_{14}\bar{A}'_{iI} + M_{15}\bar{A}'_{R*} + M_{16}\bar{A}'_{I*} + M_{17}\bar{a}'_{R*} + M_{18}\bar{a}'_{I*} + M_{19}\bar{V}'_{iI} + M_{20}\bar{V}'_{iR} \end{aligned} \quad (B50)$$

Working with the imaginary part of the momentum equation, substitute Equation B46 and B47 into Equation B44 to give

$$\begin{aligned}
M_9 [E_{16}\bar{A}'_{iR} + E_{17}A'_{ii} + E_{18}\bar{\rho}'_{R*} + E_{19}\bar{\rho}'_{i*} + E_{20}\bar{A}'_{R*} + E_{21}\bar{A}'_{i*} + \\
E_{22}\bar{a}'_{R*} + E_{23}\bar{a}'_{i*} + E_{24}\bar{V}'_{ii} + E_{25}\bar{V}'_{iR}] + M_{10} [E_{26}\bar{A}'_{iR} + \\
E_{27}\bar{A}'_{ii} + E_{28}\bar{\rho}'_{R*} + E_{29}\bar{\rho}'_{i*} + E_{30}A'_{R*} + E_{31}\bar{A}'_{i*} + E_{32}\bar{a}'_{R*} + E_{33}\bar{a}'_{i*} + \\
E_{34}\bar{V}'_{ii} + E_{35}\bar{V}'_{iR}] = M_5\bar{A}'_{ii} + M_6\bar{A}'_{i*} + 2\bar{\rho}'_{i*}\bar{a}^{*2} + 2\bar{\rho}^*\bar{a}^*\bar{a}'_{i*} + \\
\frac{k\bar{\rho}'_{R*}\bar{U}_1\bar{V}_1}{2} + k\bar{\rho}_1\bar{U}_1\bar{V}'_{iR} + \frac{k\bar{\rho}_1\bar{a}^*\bar{V}_1}{2}
\end{aligned} \tag{B51}$$

Further, assume the following relationships.

$$\begin{aligned}
M_{21} &= E_{18}M_9 + E_{28}M_{10} - \frac{k\bar{U}_1\bar{V}_1}{2} \\
M_{22} &= E_{19}M_9 + E_{29}M_{10} - 2\bar{a}^{*2} \\
M_{23} &= -E_{16}M_9 - E_{26}M_{10} + M_5 \\
M_{24} &= -E_{17}M_9 - E_{27}M_{10} + M_6 \\
M_{25} &= -E_{20}M_9 - E_{30}M_{10} \\
M_{26} &= -E_{21}M_9 - E_{31}M_{10} \\
M_{27} &= -E_{22}M_9 - E_{32}M_{10} + \frac{k\bar{\rho}_1\bar{V}_1}{2} \\
M_{28} &= -E_{23}M_9 - E_{33}M_{10} + 2\bar{\rho}^*\bar{a}^* \\
M_{29} &= -E_{24}M_9 - E_{34}M_{10} \\
M_{30} &= -E_{25}M_9 - E_{35}M_{10} + k\bar{\rho}_1\bar{U}_1
\end{aligned} \tag{B52}$$

Substituting the relationships of Equation B52 into Equation B51 and rearranging the resultant to isolate $\bar{\rho}'_{R*}$ and $\bar{\rho}'_{i*}$ gives

$$\begin{aligned}
M_{21}\bar{\rho}'_{R*} + M_{22}\bar{\rho}'_{i*} = M_{23}\bar{A}'_{iR} + M_{24}\bar{A}'_{ii} + M_{25}\bar{A}'_{R*} + M_{26}\bar{A}'_{i*} + \\
M_{27}\bar{a}'_{R*} + M_{28}\bar{a}'_{i*} + M_{29}\bar{V}'_{ii} + M_{30}\bar{V}'_{iR}
\end{aligned} \tag{B53}$$

Solve for $\bar{\rho}'_{R*}$ using Equations B50 and B53. First multiply Equation B53 by M_{12}/M_{22} to give

$$\begin{aligned}
\left(\frac{M_{21}M_{12}}{M_{22}} \right) \bar{\rho}'_{R*} + M_{12}\bar{\rho}'_{i*} = \frac{M_{12}}{M_{22}} [M_{23}\bar{A}'_{iR} + \\
M_{24}\bar{A}'_{ii} + M_{25}\bar{A}'_{R*} + M_{26}\bar{A}'_{i*} + \\
M_{27}\bar{a}'_{R*} + M_{28}\bar{a}'_{i*} + \\
M_{29}\bar{V}'_{ii} + M_{30}\bar{V}'_{iR}]
\end{aligned}$$

Next, subtract Equation B53 from Equation B50.

$$\begin{aligned}
\bar{\rho}'_{R^*} = & \left\{ \left(M_{13} - \frac{M_{12}M_{23}}{M_{22}} \right) \bar{A}'_{iR} + \right. \\
& \left(M_{14} - \frac{M_{12}M_{24}}{M_{22}} \right) \bar{A}'_{iI} + \left(M_{15} - \frac{M_{12}M_{25}}{M_{22}} \right) \bar{A}'_{R^*} + \\
& \left(M_{16} - \frac{M_{12}M_{26}}{M_{22}} \right) \bar{A}'_{I^*} + \left(M_{17} - \frac{M_{12}M_{27}}{M_{22}} \right) \bar{a}'_{R^*} + \\
& \left(M_{18} - \frac{M_{12}M_{28}}{M_{22}} \right) \bar{a}'_{I^*} + \left(M_{19} - \frac{M_{12}M_{29}}{M_{22}} \right) \bar{V}'_{iI} + \\
& \left. \left(M_{20} - \frac{M_{12}M_{30}}{M_{22}} \right) \bar{V}'_{iR} \right\} \div \left\{ M_{11} - \frac{M_{12}M_{21}}{M_{22}} \right\} \quad (B54)
\end{aligned}$$

Now, let

$$\begin{aligned}
M_{31} &= M_{11} - \frac{M_{12}M_{21}}{M_{22}} \\
M_{32} &= \left(M_{13} - \frac{M_{12}M_{23}}{M_{22}} \right) \div M_{31} \\
M_{33} &= \left(M_{14} - \frac{M_{12}M_{24}}{M_{22}} \right) \div M_{31} \\
M_{34} &= \left(M_{15} - \frac{M_{12}M_{25}}{M_{22}} \right) \div M_{31} \\
M_{35} &= \left(M_{16} - \frac{M_{12}M_{26}}{M_{22}} \right) \div M_{31} \\
M_{36} &= \left(M_{17} - \frac{M_{12}M_{27}}{M_{22}} \right) \div M_{31} \\
M_{37} &= \left(M_{18} - \frac{M_{12}M_{28}}{M_{22}} \right) \div M_{31} \\
M_{38} &= \left(M_{19} - \frac{M_{12}M_{29}}{M_{22}} \right) \div M_{31} \\
M_{39} &= \left(M_{20} - \frac{M_{12}M_{30}}{M_{22}} \right) \div M_{31} \quad (B55)
\end{aligned}$$

Substituting the relationships in Equation B55 into Equation B54 results in the following:

$$\begin{aligned}
\bar{\rho}'_{R^*} = & M_{32}\bar{A}'_{iR} + M_{33}\bar{A}'_{iI} + M_{34}\bar{A}'_{R^*} + \\
& M_{35}\bar{A}'_{I^*} + M_{36}\bar{a}'_{R^*} + M_{37}\bar{a}'_{I^*} + \\
& M_{38}\bar{V}'_{iI} + M_{39}\bar{V}'_{iR} \quad (B56)
\end{aligned}$$

$$\begin{aligned}
\bar{p}'_1^* = & \frac{1}{M_{12}} [M_{13} - M_{32}M_{11}] \bar{A}'_{iR} + (M_{14} - M_{33}M_{11}) \bar{A}'_{iI} + \\
& (M_{15} - M_{34}M_{11}) \bar{A}'_{R^*} + (M_{16} - M_{35}M_{11}) \bar{A}'_{I^*} + \\
& (M_{17} - M_{36}M_{11}) \bar{a}'_{R^*} + (M_{18} - M_{37}M_{11}) \bar{a}'_{I^*} + \\
& (M_{19} - M_{38}M_{11}) \bar{V}'_{iI} + (M_{20} - M_{39}M_{11}) \bar{V}'_{iR}]
\end{aligned}$$

Assume the following relationships:

$$\begin{aligned}
M_{40} &= \frac{M_{13} - M_{32}M_{11}}{M_{12}} \\
M_{41} &= \frac{M_{14} - M_{33}M_{11}}{M_{12}} \\
M_{42} &= \frac{M_{15} - M_{34}M_{11}}{M_{12}} \\
M_{43} &= \frac{M_{16} - M_{35}M_{11}}{M_{12}} \\
M_{44} &= \frac{M_{17} - M_{36}M_{11}}{M_{12}} \\
M_{45} &= \frac{M_{18} - M_{37}M_{11}}{M_{12}} \\
M_{46} &= \frac{M_{19} - M_{38}M_{11}}{M_{12}} \\
M_{47} &= \frac{M_{20} - M_{39}M_{11}}{M_{12}}
\end{aligned}$$

Substituting the relationships in Equation B51 into Equation B57 results in the value of \bar{p}'_1^* in the momentum equation.

$$\begin{aligned}
\bar{p}'_1^* = & M_{40}\bar{A}'_{iR} + M_{41}\bar{A}'_{iI} + M_{42}\bar{A}'_{R^*} + \\
& M_{43}\bar{A}'_{I^*} + M_{44}\bar{a}'_{R^*} + M_{45}\bar{a}'_{I^*} + \\
& M_{46}\bar{V}'_{iI} + M_{47}\bar{V}'_{iR}
\end{aligned} \tag{B59}$$

Working with the continuity equation, recall that

$$\begin{aligned}
[\rho'^* \bar{a}^* + \bar{p}^* A'^* \bar{a}^* + \bar{p}^* a'^*] - [\bar{A}_i (u'_i \sin \alpha_{ch} + v'_i \cos \alpha_{ch}) + \\
A'_i I + \rho'_i \bar{A}_i I] = - \left[\bar{\rho}_i \frac{\partial V'_i}{\partial t} + \bar{V}_i \frac{\partial \rho'_i}{\partial t} \right]
\end{aligned} \tag{B26}$$

Substituting for the perturbation quantities and dividing through by e^{ikt} gives

$$\begin{aligned}
& (\bar{\rho}'^* \bar{a}^* + \bar{\rho}^* \bar{A}'^* \bar{a}^* + \bar{\rho}^* \bar{a}'^*) - \left\{ \bar{A}_i \left[\frac{B_i A_{-\infty}}{C \bar{y}_i} (e^{iC \bar{y}_i} - 1) \sin \alpha_{ch} + \right. \right. \\
& \quad \left. \left. \frac{A_{-\infty}}{\bar{y}_i} (e^{iC \bar{y}_i} - 1) \cos \alpha_{ch} \right] + \bar{A}'_i I - \left(\frac{k + \bar{u}_i B_i + \bar{v}_i C}{C \bar{y}_i \bar{a}_i^2} \right) \bar{A}_i A_{-\infty} I (e^{iC \bar{y}_i} - 1) \right\} \\
& = - \left[ik \bar{\rho}_1 \bar{V}'_1 + \frac{ik \bar{V}_1}{2} \left(\bar{\rho}'^* - \left(\frac{k + \bar{u}_i B_i + \bar{v}_i C}{C \bar{y}_i \bar{a}_i^2} \right) A_{-\infty} (e^{iC \bar{y}_i} - 1) \right) \right] \quad (B60)
\end{aligned}$$

Assume the following relationships:

$$\begin{aligned}
C_1 &= - \frac{\bar{A}_i}{C \bar{y}_i} \sin \alpha_{ch} + \frac{\bar{u}_i \bar{A}_i I}{C \bar{y}_i \bar{a}_i^2} \\
C_2 &= - \frac{\bar{A}_i}{\bar{y}_i} \cos \alpha_{ch} + \left(\frac{k + \bar{v}_i C}{C \bar{y}_i \bar{a}_i^2} \right) \bar{A}_i I \\
C_3 &= \frac{\bar{u}_i \bar{V}_1}{2 C \bar{y}_i \bar{a}_i^2} \\
C_4 &= \left(\frac{k + \bar{v}_i C}{2 C \bar{y}_i \bar{a}_i^2} \right) V_1 \quad (B61)
\end{aligned}$$

Substituting the relationships in Equation B61 into Equation B60 and separating the real parts of the resultant equation yields

$$\begin{aligned}
& C_1 [B_{1R} A_{-\infty R} E_7 - B_{1I} A_{-\infty I} E_7 - B_{1R} A_{-\infty I} \sin (C \bar{y}_i) - \\
& \quad B_{1I} A_{-\infty R} \sin (C \bar{y}_i)] + C_2 [A_{-\infty R} E_7 - A_{-\infty I} \sin (C \bar{y}_i)] + \\
& \quad [\bar{\rho}'_R \bar{a}^* + \bar{\rho}^* \bar{A}'_R \bar{a}^* + \bar{\rho}^* \bar{a}'_R - \bar{A}'_{IR} I] \\
& = -k C_3 [B_{1R} A_{-\infty I} E_7 + B_{1I} A_{-\infty R} E_7 + B_{1R} A_{-\infty R} \sin (C \bar{y}_i) - \\
& \quad B_{1I} A_{-\infty I} \sin (C \bar{y}_i)] - k C_4 [A_{-\infty I} E_7 + A_{-\infty R} \sin (C \bar{y}_i)] + \\
& \quad k \bar{\rho}_1 \bar{V}'_1 + \frac{k \bar{\rho}'_1 \bar{V}_1}{2} \quad (B62)
\end{aligned}$$

Again, assume

$$\begin{aligned}
C_5 &= C_1 B_{1R} E_7 - C_1 B_{1I} \sin (C \bar{y}_i) + C_2 E_7 + \\
& \quad C_3 k B_{1I} E_7 + C_3 k B_{1R} \sin (C \bar{y}_i) + C_4 k \sin (C \bar{y}_i) \\
C_6 &= - C_1 B_{1I} E_7 - C_1 B_{1R} \sin (C \bar{y}_i) - C_2 \sin (C \bar{y}_i) + \\
& \quad C_3 k B_{1R} E_7 - C_3 k B_{1I} \sin (C \bar{y}_i) + C_4 k E_7 \quad (B63)
\end{aligned}$$

Substituting the relationships in Equation B63 into Equation B62 yields

$$C_5 A_{-\infty R} + C_6 A_{-\infty I} + (\bar{\rho}'_R \bar{a}^* + \bar{\rho}^* \bar{A}'_R \bar{a}^* + \bar{\rho}^* \bar{a}'_R - I \bar{A}'_{II}) - k \bar{\rho}'_I \bar{V}'_{II} - \frac{k \bar{\rho}'^* \bar{V}'_I}{2} = 0 \quad (\text{B64})$$

Next, substitute the values for $A_{-\infty R}$ and $A_{-\infty I}$ into Equation B64

$$C_5 (E_{16} \bar{A}'_{iR} + E_{17} \bar{A}'_{iI} + E_{18} \bar{\rho}'_R + E_{19} \bar{\rho}'_I + E_{20} \bar{A}'_R + E_{21} \bar{A}'_I + E_{22} \bar{a}'_R + E_{23} \bar{a}'_I + E_{24} \bar{V}'_{II} + E_{25} \bar{V}'_{iR})$$

Combining like terms produces

$$\begin{aligned} & (C_5 E_{16} + C_6 E_{26} - I) \bar{A}'_{iR} + (C_5 E_{17} + C_6 E_{27}) \bar{A}'_{iI} + \\ & (C_5 E_{18} + C_6 E_{28} + \bar{a}^*) \bar{\rho}'_R + (C_5 E_{19} + C_6 E_{29} - k \bar{V}_I / 2) \bar{\rho}'_I + \\ & (C_5 E_{20} + C_6 E_{30} + \bar{\rho}^* \bar{a}^*) \bar{A}'_R + (C_5 E_{21} + C_6 E_{31}) \bar{A}'_I + \\ & (C_5 E_{22} + C_6 E_{32} + \bar{\rho}^*) \bar{a}'_R + (C_5 E_{23} + C_6 E_{33}) \bar{a}'_I + \\ & (C_5 E_{24} + C_6 E_{34} - k \bar{\rho}'_I) \bar{V}'_{II} + (C_5 E_{25} + C_6 E_{35}) \bar{V}'_{iR} = 0 \end{aligned}$$

Again, assume

$$\begin{aligned} C_7 &= C_5 E_{16} + C_6 E_{26} - I \\ C_8 &= C_5 E_{17} + C_6 E_{27} \\ C_9 &= C_5 E_{18} + C_6 E_{28} + \bar{a}^* \\ C_{10} &= C_5 E_{19} + C_6 E_{29} - k \bar{V}_I / 2 \\ C_{11} &= C_5 E_{20} + C_6 E_{30} + \bar{\rho}^* \bar{a}^* \\ C_{12} &= C_5 E_{21} + C_6 E_{31} \\ C_{13} &= C_5 E_{22} + C_6 E_{32} + \bar{\rho}^* \\ C_{14} &= C_5 E_{23} + C_6 E_{33} \\ C_{15} &= C_5 E_{24} + C_6 E_{34} - k \bar{\rho}'_I \\ C_{16} &= C_5 E_{25} + C_6 E_{35} \end{aligned} \quad (\text{B65})$$

Substituting Equation B65 into Equation B64 gives

$$C_7 \bar{A}'_{iR} + C_8 \bar{A}'_{iI} + C_9 \bar{\rho}'_R + C_{10} \bar{\rho}'_I + C_{11} \bar{A}'_R + C_{12} \bar{A}'_I + C_{13} \bar{a}'_R + C_{14} \bar{a}'_I + C_{15} \bar{V}'_{II} + C_{16} \bar{V}'_{iR} = 0 \quad (\text{B66})$$

Substituting for $\bar{\rho}'_R^*$ and $\bar{\rho}'_I^*$ results in

$$\begin{aligned} C_9 (M_{32}\bar{A}'_{iR} + M_{35}\bar{A}'_{iI} + M_{34}\bar{A}'_{R^*} + M_{35}\bar{A}'_{I^*} + M_{36}\bar{a}'_{R^*} + M_{37}\bar{a}'_{I^*} + \\ M_{38}\bar{V}'_{i1} + M_{39}\bar{V}'_{iR}) + C_{10} (M_{40}\bar{A}'_{iR} + M_{41}\bar{A}'_{iI} + M_{42}\bar{A}'_{R^*} + M_{43}\bar{A}'_{I^*} + \\ M_{44}\bar{a}'_{R^*} + M_{45}\bar{a}'_{I^*} + M_{46}\bar{V}'_{i1} + M_{47}\bar{V}'_{iR}) - C_7\bar{A}'_{iR} + C_8\bar{A}'_{iI} + \\ C_{11}\bar{A}'_{R^*} + C_{12}\bar{A}'_{I^*} + C_{13}\bar{a}'_{R^*} + C_{14}\bar{a}'_{I^*} + C_{15}\bar{V}'_{i1} + C_{16}\bar{V}'_{iR} = 0 \end{aligned}$$

Combining like terms gives

$$\begin{aligned} (C_9M_{32} + C_{10}M_{40} + C_7)\bar{A}'_{iR} + (C_9M_{33} + C_{10}M_{41} + C_8)\bar{A}'_{iI} + \\ (C_9M_{34} + C_{10}M_{42} + C_{11})\bar{A}'_{R^*} + (C_9M_{35} + C_{10}M_{43} + C_{12})\bar{A}'_{I^*} + \\ (C_9M_{36} + C_{10}M_{44} + C_{13})\bar{a}'_{R^*} + (C_9M_{37} + C_{10}M_{45} + C_{14})\bar{a}'_{I^*} + \\ (C_9M_{38} + C_{10}M_{46} + C_{15})\bar{V}'_{i1} + (C_9M_{39} + C_{10}M_{47} + C_{16})\bar{V}'_{iR} = 0 \end{aligned} \quad (B67)$$

Assume

$$\begin{aligned} C_{17} &= C_9M_{32} + C_{10}M_{40} + C_7 \\ C_{18} &= C_9M_{33} + C_{10}M_{41} + C_8 \\ C_{19} &= C_9M_{34} + C_{10}M_{42} + C_{11} \\ C_{20} &= C_9M_{35} + C_{10}M_{43} + C_{12} \\ C_{21} &= - (C_9M_{36} + C_{10}M_{44} + C_{13}) \\ C_{22} &= - (C_9M_{37} + C_{10}M_{45} + C_{14}) \\ C_{23} &= C_9M_{38} + C_{10}M_{46} + C_{15} \\ C_{24} &= C_9M_{39} + C_{10}M_{47} + C_{16} \end{aligned} \quad (B68)$$

Substituting the relationships in Equation B68 into Equation B67 and solving for \bar{a}'_{R^*} and \bar{a}'_{I^*} gives

$$C_{21}\bar{a}'_{R^*} + C_{22}\bar{a}'_{I^*} = C_{17}\bar{A}'_{iR} + C_{18}\bar{A}'_{iI} + C_{19}\bar{A}'_{R^*} + C_{20}\bar{A}'_{I^*} + C_{23}\bar{V}'_{i1} + C_{24}\bar{V}'_{iR} \quad (B69)$$

The imaginary part of Equation 24 is

$$\begin{aligned} (\bar{\rho}'_I^*\bar{a}^* + \bar{\rho}\bar{A}'_{I^*}\bar{a}^* + \bar{\rho}^*\bar{a}'_{I^*}) - I\bar{A}'_{iI} + C_1 [B_{iR}A_{-\infty I}E_7 + \\ B_{iI}A_{-\infty R}E_7 + B_{iR}A_{-\infty R} \sin(C\bar{y}_i) - B_{iI}A_{-\infty I} \sin(Cy_i)] + \\ C_2 [A_{-\infty I}E_7 + A_{-\infty R} \sin(C\bar{y}_i)] - kC_3 [B_{iR}A_{-\infty R}E_7 - \\ B_{iI}A_{-\infty I}E_7 - B_{iR}A_{-\infty I} \sin(C\bar{y}_i) - B_{iI}A_{-\infty R} \sin(C\bar{y}_i)] - \\ kC_4 [A_{-\infty R}E_7 - A_{-\infty I} \sin(C\bar{y}_i)] + k\bar{\rho}_I \bar{V}'_{iR} + \frac{k\bar{\rho}'_R^*\bar{V}'_{i1}}{2} = 0 \end{aligned} \quad (B70)$$

Again, let

$$\begin{aligned}
C_{25} &= C_1 [B_{11}E_7 + B_{1R} \sin(C\bar{y}_i)] + C_2 \sin(C\bar{y}_i) - kC_3 [B_{1R}FE_7 - \\
&\quad B_{11} \sin(C\bar{y}_i)] - kC_4E_7 \\
C_{26} &= C_1 [B_{1R}E_7 - B_{11} \sin(C\bar{y}_i)] + C_2E_7 + kC_3 [B_{11}E_7 + \\
&\quad B_{1R} \sin(C\bar{y}_i)] + kC_4 \sin(C\bar{y}_i)
\end{aligned} \tag{B71}$$

Substituting the relationships of Equation B71 into Equation B70 yields

$$\begin{aligned}
C_{25}A_{-\infty R} + C_{26}A_{-\infty I} + (\bar{\rho}'_1 \bar{a}^* + \bar{\rho}^* \bar{A}'_1 \bar{a}^* + \bar{\rho}^* \bar{a}'_1) - I \bar{A}'_{11} + k\bar{\rho}_1 \bar{V}'_{1R} + \\
\frac{k\bar{\rho}'_R \bar{V}'_1}{2} = 0
\end{aligned} \tag{B72}$$

Substituting for $A_{-\infty R}$ and $A_{-\infty I}$ into Equation B72 gives

$$\begin{aligned}
C_{26} [E_{26} \bar{A}'_{1R} + E_{27} \bar{A}'_{1I} + E_{28} \bar{\rho}'_R + E_{29} \bar{\rho}'_1 + E_{30} \bar{A}'_R + E_{31} \bar{A}'_1 + \\
E_{32} \bar{a}'_R + E_{33} \bar{a}'_1 + E_{34} \bar{V}'_{11} + E_{35} \bar{V}'_{1R}] + C_{25} [E_{16} \bar{A}'_{1R} + E_{17} \bar{A}'_{1I} + \\
E_{18} \bar{\rho}'_R + E_{19} \bar{\rho}'_1 + E_{20} \bar{A}'_R + E_{21} \bar{A}'_1 + E_{22} \bar{a}'_R + E_{23} \bar{a}'_1 + \\
E_{24} \bar{V}'_{11} + E_{25} \bar{V}'_{1R}] + [\bar{\rho}'_1 \bar{a}^* + \bar{\rho}^* \bar{A}'_1 \bar{a}^* + \bar{\rho}^* \bar{a}'_1] - I \bar{A}'_{11} + k\bar{\rho}_1 \bar{V}'_{1R} + \frac{kV_1}{2} \bar{\rho}'_R = 0
\end{aligned}$$

Collecting like terms, as before, produces

$$\begin{aligned}
(C_{26} E_{26} + C_{25} E_{16}) \bar{A}'_{1R} + (C_{26} E_{27} + C_{25} E_{17} - I) \bar{A}'_{1I} + \\
(C_{26} E_{28} + C_{25} E_{18} + \frac{k\bar{V}'_1}{2}) \bar{\rho}'_R + (C_{26} E_{29} + C_{25} E_{19} + \bar{a}^*) \\
\bar{\rho}'_1 + (C_{26} E_{30} + C_{25} E_{20}) \bar{A}'_R + (C_{26} E_{31} + C_{25} E_{21} + \bar{\rho}^* \bar{a}^*) \\
\bar{A}'_1 + (C_{26} E_{32} + C_{25} E_{22}) \bar{a}'_R + (C_{26} E_{33} + C_{25} E_{23} + \bar{\rho}^*) \\
\bar{a}'_1 + (C_{26} E_{34} + C_{25} E_{24}) \bar{V}'_{11} + (C_{26} E_{35} + C_{25} E_{25} + k\bar{\rho}_1) \bar{V}'_{1R} = 0
\end{aligned} \tag{B73}$$

Again, assume

$$\begin{aligned}
C_{27} &= C_{26} E_{26} + C_{25} E_{16} \\
C_{28} &= C_{26} E_{27} + C_{25} E_{17} - I \\
C_{29} &= C_{26} E_{28} + C_{25} E_{18} + \frac{k\bar{V}'_1}{2} \\
C_{30} &= C_{26} E_{29} + C_{25} E_{19} + \bar{a}^* \\
C_{31} &= C_{26} E_{30} + C_{25} E_{20} \\
C_{32} &= C_{26} E_{31} + C_{25} E_{21} + \bar{\rho}^* \bar{a}^*
\end{aligned}$$

$$\begin{aligned}
C_{33} &= C_{26} E_{32} + C_{25} E_{22} \\
C_{34} &= C_{26} E_{33} + C_{25} E_{23} + \bar{\rho}^* \\
C_{35} &= C_{26} E_{34} + C_{25} E_{24} \\
C_{36} &= C_{26} E_{35} + C_{25} E_{25} + k \bar{\rho}_1
\end{aligned} \tag{B74}$$

Substituting the relationships of Equation B74 and the relationships for $\bar{\rho}'_R^*$ and $\bar{\rho}'_I^*$, and collecting terms gives

$$\begin{aligned}
&(C_{29}M_{32} + C_{30}M_{40} + C_{27}) \bar{A}'_{iR} + (C_{29}M_{33} + C_{30}M_{41} + C_{28}) \bar{A}'_{iI} + \\
&(C_{29}M_{34} + C_{30}M_{42} + C_{31}) \bar{A}'_R^* + (C_{29}M_{35} + C_{30}M_{43} + C_{32}) \bar{A}'_I^* + \\
&(C_{29}M_{36} + C_{30}M_{44} + C_{33}) \bar{a}'_R^* + (C_{29}M_{37} + C_{30}M_{45} + C_{34}) \bar{a}'_I^* + \\
&(C_{29}M_{38} + C_{30}M_{46} + C_{35}) \bar{V}'_{iI} + (C_{29}M_{39} + C_{30}M_{47} + C_{36}) \bar{V}'_{iR} = 0
\end{aligned} \tag{B75}$$

Further, assume

$$\begin{aligned}
C_{37} &= C_{29}M_{32} + C_{30}M_{40} + C_{27} \\
C_{38} &= C_{29}M_{33} + C_{30}M_{41} + C_{28} \\
C_{39} &= C_{29}M_{34} + C_{30}M_{42} + C_{31} \\
C_{40} &= C_{29}M_{35} + C_{30}M_{43} + C_{32} \\
C_{41} &= - (C_{29}M_{36} + C_{30}M_{44} + C_{33}) \\
C_{42} &= - (C_{29}M_{37} + C_{30}M_{45} + C_{34}) \\
C_{43} &= C_{29}M_{38} + C_{30}M_{46} + C_{35} \\
C_{44} &= C_{29}M_{39} + C_{30}M_{47} + C_{36}
\end{aligned} \tag{B76}$$

Substituting the relationships of Equation B76 into Equation B75 gives

$$\begin{aligned}
C_{41} \bar{a}'_R^* + C_{42} \bar{a}'_I^* &= C_{37} \bar{A}'_{iR} + C_{38} \bar{A}'_{iI} + \\
&C_{39} \bar{A}'_R^* + C_{40} \bar{A}'_I^* + C_{43} \bar{V}'_{iI} + C_{44} \bar{V}'_{iR}
\end{aligned} \tag{B77}$$

Recall that

$$\begin{aligned}
C_{21} \bar{a}'_R^* + C_{22} \bar{a}'_I^* &= C_{17} \bar{A}'_{iR} + C_{18} \bar{A}'_{iI} + C_{19} \bar{A}'_R^* + \\
&C_{20} \bar{A}'_I^* + C_{23} \bar{V}'_{iI} + C_{24} \bar{V}'_{iR}
\end{aligned} \tag{B69}$$

Multiplying Equation B78 by C_{42}/C_{22} and subtracting the resultant from Equation B77 yields

$$\begin{aligned} \bar{a}'_R * &= \left[\left(C_{37} - \frac{C_{17}C_{42}}{C_{22}} \right) \bar{A}'_{iR} + \left(C_{38} - \frac{C_{18}C_{42}}{C_{22}} \right) \bar{A}'_{ii} + \right. \\ &\quad \left(C_{39} - \frac{C_{19}C_{42}}{C_{22}} \right) \bar{A}'_R * + \left(C_{40} - \frac{C_{20}C_{42}}{C_{22}} \right) \bar{A}'_i * + \\ &\quad \left(C_{43} - \frac{C_{23}C_{42}}{C_{22}} \right) \bar{V}'_{iI} + \left(C_{44} - \frac{C_{24}C_{42}}{C_{22}} \right) \bar{V}'_{iR} \left. \right] \div \\ &\quad \left[C_{41} - \frac{C_{21}C_{42}}{C_{22}} \right] \end{aligned}$$

Further, assure

$$\begin{aligned} C_{45} &= C_{41} - \frac{C_{21}C_{42}}{C_{22}} \\ C_{46} &= \frac{1}{C_{45}} \left(C_{37} - \frac{C_{17}C_{42}}{C_{22}} \right) \\ C_{47} &= \frac{1}{C_{45}} \left(C_{38} - \frac{C_{18}C_{42}}{C_{22}} \right) \\ C_{48} &= \frac{1}{C_{45}} \left(C_{39} - \frac{C_{19}C_{42}}{C_{22}} \right) \\ C_{49} &= \frac{1}{C_{45}} \left(C_{40} - \frac{C_{20}C_{42}}{C_{22}} \right) \\ C_{50} &= \frac{1}{C_{45}} \left(C_{43} - \frac{C_{23}C_{42}}{C_{22}} \right) \\ C_{51} &= \frac{1}{C_{45}} \left(C_{44} - \frac{C_{24}C_{42}}{C_{22}} \right) \end{aligned} \tag{B78}$$

Thus,

$$\begin{aligned} \Rightarrow \bar{a}'_R * &= C_{46} \bar{A}'_{iR} + C_{47} \bar{A}'_{ii} + C_{48} \bar{A}'_R * + C_{49} \bar{A}'_i * + \\ &\quad C_{50} \bar{V}'_{iI} + C_{51} \bar{V}'_{iR} \end{aligned} \tag{B79}$$

Substituting Equation B79 into Equation B77 and solving for $a'_i *$ and then combining like terms gives

$$\begin{aligned} \bar{a}'_i * &= \frac{1}{C_{42}} [(C_{37} - C_{41}C_{46}) \bar{A}'_{iR} + (C_{38} - C_{41}C_{47}) \bar{A}'_{ii} + \\ &\quad (C_{39} - C_{41}C_{48}) \bar{A}'_R * + (C_{40} - C_{41}C_{49}) \bar{A}'_i * + \\ &\quad (C_{43} - C_{41}C_{50}) \bar{V}'_{iI} + (C_{44} - C_{41}C_{51}) \bar{V}'_{iR}] \end{aligned}$$

Further, assume

$$\begin{aligned}
C_{52} &= (C_{37} - C_{41}C_{46}) \div C_{42} \\
C_{53} &= (C_{38} - C_{41}C_{47}) \div C_{42} \\
C_{54} &= (C_{39} - C_{41}C_{48}) \div C_{42} \\
C_{55} &= (C_{40} - C_{41}C_{49}) \div C_{42} \\
C_{56} &= (C_{43} - C_{41}C_{50}) \div C_{42} \\
C_{57} &= (C_{44} - C_{41}C_{51}) \div C_{42}
\end{aligned} \tag{B80}$$

Substituting the relationships of Equation 136 into the relationship for a'_1^* gives

$$\begin{aligned}
\Rightarrow \bar{a}'_1^* &= C_{52}\bar{A}'_{IR} + C_{53}\bar{A}'_{II} + C_{54}\bar{A}'_{R^*} + C_{55}\bar{A}'_{I^*} + \\
&C_{56}\bar{V}'_{II} + C_{57}\bar{V}'_{IR}
\end{aligned} \tag{B81}$$

This completes the analysis for Region 1.

APPENDIX C
SOLUTIONS TO THE EQUATIONS OF MOTION FOR REGION 2

The nondimensionalized equations of motion are represented by Equations 38 and 39. Only two equations are needed because all the inlet parameters to Region 2 are known. This leaves two unknowns: the density and velocity perturbations just upstream of the shock. To obtain a solution to these equations, again make the assumption that all flow parameters vary harmonically in time. Thus,

$$\begin{aligned}
 p'_{us} &= \bar{p}'_{us} \exp ikt \\
 A'_s &= \bar{A}'_s \exp ikt \\
 \rho'_{us} &= \bar{\rho}'_{us} \exp ikt \\
 U'_{us} &= \bar{U}'_{us} \exp ikt
 \end{aligned} \tag{C1}$$

making these substitutions into the momentum equation (39) and dividing by e^{ikt} gives

$$\begin{aligned}
 &(\bar{p}'^* \bar{A}'^* + \bar{p}'^* - \bar{p}'_{us} \bar{A}'_{us} - \bar{p}'_{us} \bar{A}'_s) \\
 &= [\bar{U}'_{us} \bar{w}'_{us} + \bar{U}'_{us} (\bar{\rho}'_{us} \bar{U}'_{us} \bar{A}'_s + \bar{\rho}'_{us} \bar{U}'_{us} \bar{A}'_s + \bar{\rho}'_{us} \bar{U}'_{us} \bar{A}'_s)] - \\
 &\quad [\bar{\rho}'^* \bar{a}'^* \bar{a}'^* + \bar{a}'^* (\rho' \bar{a}'^* + \bar{\rho}'^* \bar{A}'^* \bar{a}'^* + \bar{\rho}'^* \bar{a}'^*)] + \\
 &\quad \frac{ik \bar{U}'_2 \bar{V}'_2}{2} (\bar{\rho}'^* + \bar{\rho}'_{us}) + ik \bar{\rho}'_2 \bar{U}'_2 \bar{V}'_2 + ik \frac{\bar{\rho}'_2 \bar{V}'_2}{2} (\bar{a}'^* + \bar{U}'_{us})
 \end{aligned} \tag{C2}$$

Making the substitution $p' = \bar{a}'^2 \rho'$ and collecting terms produces

$$\begin{aligned}
 &(\bar{p}'^* \bar{A}'^* + \bar{p}'^* - \bar{p}'_{us} \bar{A}'_s - \bar{\rho}'_{us} \bar{A}'_s \bar{a}'^2_{us}) \\
 &= (2\bar{U}'_{us} \bar{w}'_{us} + \bar{\rho}'_{us} \bar{U}'^2_{us} \bar{A}'_s + \bar{\rho}'_{us} \bar{U}'^2_{us} \bar{A}'_s) - (2\bar{\rho}'^* \bar{a}'^* \bar{a}'^* + \rho' \bar{a}'^2 + \bar{\rho}'^* \bar{A}'^* \bar{a}'^2) + \\
 &\quad \frac{ik \bar{U}'_2 \bar{V}'_2}{2} (\bar{\rho}'^* + \bar{\rho}'_{us}) + ik \bar{\rho}'_2 \bar{U}'_2 \bar{V}'_2 + ik \bar{\rho}'_2 \bar{V}'_2 (\bar{a}'^* + \bar{U}'_{us})
 \end{aligned} \tag{C3}$$

Real Part

$$\begin{aligned}
 &\bar{p}'^*_R + \bar{p}'^* \bar{A}'^*_R - \bar{\rho}'_{USR} \bar{A}'_s \bar{a}'^2_{US} - \bar{p}'_{US} \bar{A}'_{SR} \\
 &= \bar{\rho}'_{USR} \bar{U}'^2_{US} \bar{A}'_s - \frac{k \bar{\rho}'_{USI} \bar{U}'_2 \bar{V}'_2}{2} + 2\bar{U}'_{USR} \bar{w}'_{US} - \frac{k \bar{\rho}'_2 \bar{U}'_{USI} \bar{V}'_2}{2} - \\
 &\quad 2\bar{\rho}'^* \bar{a}'^*_R \bar{a}'^* - \frac{k \bar{\rho}'_2 \bar{a}'^*_I \bar{V}'_2}{2} - \bar{\rho}'^*_R \bar{a}'^2 - \frac{k \bar{\rho}'_I \bar{U}'_2 \bar{V}'_2}{2} - k \bar{\rho}'_2 \bar{U}'_2 \bar{V}'_{21} + \\
 &\quad \bar{\rho}'_{US} \bar{U}'^2_{US} \bar{A}'_{SR} - \bar{\rho}'^* \bar{A}'^*_R \bar{a}'^2
 \end{aligned} \tag{C4}$$

Rearranging terms produces

$$\begin{aligned}
& (-\bar{a}_{US}^2 \bar{A}_S - \bar{U}_{US}^2 \bar{A}_S) \bar{\rho}'_{USR} + \left(\frac{k \bar{U}_2 \bar{V}_2}{2} \right) \bar{\rho}'_{USI} \\
& = (\bar{\rho}_{US} \bar{U}_{US}^2 + \bar{p}_{US}) \bar{A}'_{SR} - 2\bar{p}'_R - (\bar{p}^* + \bar{\rho}^* \bar{a}^{*2}) \bar{A}'_R + 2\bar{U}'_{USR} \bar{w}_{US} - \\
& \quad \frac{k \bar{\rho}_2 \bar{U}'_{USI} \bar{V}_2}{2} - \frac{k \bar{\rho}'_1 \bar{U}_2 \bar{V}_2}{2} - k \bar{\rho}_2 \bar{U}_2 \bar{V}'_{21} - 2\bar{\rho}^* \bar{a}'^*_{R} \bar{a}^* - \frac{k \bar{\rho}_2 \bar{a}'_1 \bar{V}_2}{2}
\end{aligned} \tag{C5}$$

Next, assume

$$\begin{aligned}
M_{48} & = (-\bar{a}_{US}^2 - \bar{U}_{US}^2) \bar{A}_S \\
M_{49} & = \frac{k \bar{U}_2 \bar{V}_2}{2} \\
M_{50} & = \bar{p}_{US} + \bar{\rho}_{US} \bar{U}_{US}^2 \\
M_{50A} & = \bar{p}^* + \bar{\rho}^* \bar{a}^{*2}
\end{aligned} \tag{C6}$$

Substituting Equation C6 into Equation C5 gives

$$\begin{aligned}
M_{48} \bar{\rho}'_{USR} + M_{49} \bar{\rho}'_{USI} & = M_{50} \bar{A}'_{SR} - 2\bar{p}'_R + M_{50A} \bar{A}'_R + 2\bar{U}'_{USR} \bar{w}_{US} - \\
& \quad \frac{k \bar{\rho}_2 \bar{U}'_{USI} \bar{V}_2}{2} - \frac{k \bar{\rho}'_1 \bar{U}_2 \bar{V}_2}{2} - k \bar{\rho}_2 \bar{U}_2 \bar{V}'_{21} - 2\bar{\rho}^* \bar{a}'^*_{R} \bar{a}^* - \frac{k \bar{\rho}_2 \bar{a}'_1 \bar{V}_2}{2}
\end{aligned} \tag{C7}$$

Imaginary Part

$$\begin{aligned}
M_{48} \bar{\rho}'_{USI} - M_{49} \bar{\rho}'_{USR} & = M_{50} \bar{A}'_{SI} - 2\bar{p}'_1 - M_{50A} \bar{A}'_1 + 2\bar{U}'_{USI} \bar{w}_{US} + \\
& \quad \frac{k \bar{\rho}_2 \bar{U}'_{USR} \bar{V}_2}{2} + \frac{k \bar{\rho}'_R \bar{U}_2 \bar{V}_2}{2} + k \bar{\rho}_2 \bar{U}_2 \bar{V}'_{2R} - 2\bar{\rho}'^* \bar{a}'^*_{1} \bar{a}^* + \frac{k \bar{\rho}_2 \bar{a}'_R \bar{V}_2}{2}
\end{aligned} \tag{C8}$$

Multiply Equation C8 by M_{48}/M_{49} and add to Equation C7 to produce

$$\begin{aligned}
& \left(M_{49} + \frac{M_{48}^2}{M_{49}} \right) \bar{\rho}'_{USI} = M_{50} \bar{A}'_{SR} + \left(\frac{M_{48} M_{50}}{M_{49}} \right) \bar{A}'_{SI} - 2 \bar{p}'_{R} - 2 \left(\frac{M_{48}}{M_{49}} \right) \bar{p}'_{I^*} + \\
& M_{50A} \bar{A}'_{R^*} + \left(\frac{M_{48} M_{50A}}{M_{49}} \right) \bar{A}'_{I^*} + 2 \bar{U}'_{USR} \bar{w}_{US} + 2 \left(\frac{M_{48} \bar{w}_{US}}{M_{49}} \right) \bar{U}'_{USI} - \\
& \frac{k \bar{\rho}_2 \bar{U}'_{USI} \bar{V}_2}{2} + \frac{k \bar{\rho}_2 \bar{U}'_{USR} \bar{V}_2}{2} \left(\frac{M_{48}}{M_{49}} \right) - \frac{k \bar{\rho}'_1 \bar{U}_2 \bar{V}_2}{2} + \frac{k \bar{\rho}'_R \bar{U}_2 \bar{V}_2}{2} \left(\frac{M_{48}}{M_{49}} \right) \\
& - k \bar{\rho}_2 \bar{U}_2 \bar{V}'_{2I} + \frac{k \bar{\rho}_2 \bar{U}_2 \bar{V}'_{2R}}{2} \left(\frac{M_{48}}{M_{49}} \right) - 2 \bar{\rho}^* \bar{a}'_{R^*} \bar{a}^* - 2 \bar{\rho}^* \bar{a}'_{I^*} \left(\frac{M_{48}}{M_{49}} \right) \bar{a}^* - \\
& \frac{k \bar{\rho}_2 \bar{a}'_1 \bar{V}_2}{2} + \frac{k \bar{\rho}_2 \bar{a}'_R \bar{V}_2}{2} \left(\frac{M_{48}}{M_{49}} \right)
\end{aligned} \tag{C9}$$

Assume the following relationships

$$\begin{aligned}
M_{51} &= M_{49} + \frac{M_{48}^2}{M_{49}} \\
M_{52} &= \frac{M_{50}}{M_{51}} \\
M_{53} &= \frac{M_{48} M_{50}}{M_{49} M_{51}} \\
M_{54} &= \frac{-2}{M_{51}} \\
M_{55} &= \frac{-2 M_{48}}{M_{49} M_{51}} \\
M_{56} &= \frac{M_{50A}}{M_{51}} \\
M_{57} &= \frac{M_{48} M_{50A}}{M_{49} M_{51}} \\
M_{58} &= \left[2\bar{w}_{US} + \frac{k\bar{\rho}_2\bar{V}_2}{2} \left(\frac{M_{49}}{M_{49}} \right) \right] \div M_{51} \\
M_{59} &= \frac{1}{M_{51}} \left(\frac{2\bar{w}_{US}M_{48}}{M_{49}} - \frac{k\bar{\rho}_2\bar{V}_2}{2} \right) \\
M_{60} &= \frac{1}{M_{51}} \left(-\frac{\bar{a}^*2\bar{\rho}^* M_{48}}{M_{49}} - \frac{k\bar{\rho}_2\bar{V}_2}{2} \right) \\
M_{61} &= \frac{1}{M_{51}} \left(-2\bar{\rho}^*\bar{a}^* + \frac{k\bar{\rho}_2\bar{V}_2 M_{48}}{2M_{49}} \right) \\
M_{62} &= \frac{1}{M_{51}} \left(\frac{-k\bar{U}_2\bar{V}_2}{2} \right) \\
M_{63} &= \frac{1}{M_{51}} \left(\frac{k\bar{U}_2\bar{V}_2 M_{48}}{2M_{49}} \right) \\
M_{64} &= \frac{-k(\bar{\rho} \bar{U})_2}{M_{51}} \\
M_{65} &= k(\bar{\rho} \bar{U})_2 \left(\frac{M_{48}}{M_{49} M_{51}} \right)
\end{aligned} \tag{C10}$$

Substituting these values into Equation 3.7 and solving for ρ'_{USI} gives

$$\begin{aligned}\bar{\rho}'_{USI} = & M_{52}\bar{A}'_{SR} + M_{53}\bar{A}'_{SI} + M_{54}\bar{\rho}'_{R^*} + M_{55}\bar{\rho}'_{I^*} + M_{56}\bar{A}'_{R^*} + \\ & M_{57}\bar{A}'_{I^*} + M_{58}\bar{U}'_{USR} + M_{59}\bar{U}'_{USI} + M_{60}\bar{a}'_{I^*} + M_{61}\bar{a}'_{R^*} + \\ & M_{62}\bar{\rho}'_{I^*} + M_{63}\bar{\rho}'_{R^*} + M_{64}\bar{V}'_{2I} + M_{65}\bar{V}'_{2R}\end{aligned}\quad (C11)$$

Substituting the relationships of Equation C11 into C7 and solving for $\bar{\rho}'_{USR}$ gives

$$\begin{aligned}\bar{\rho}'_{USR} = & \frac{1}{M_{48}} \left[(M_{50} - M_{49}M_{52}) \bar{A}'_{SR} + (M_{50A} - M_{49}M_{53}) \bar{A}'_{SI} + \right. \\ & (-2 - M_{49}M_{54}) \bar{\rho}'_{R^*} + (-M_{49}M_{55}) \bar{\rho}'_{I^*} + \\ & (M_{50A} - M_{49}M_{56}) \bar{A}'_{R^*} + (-M_{49}M_{57}) \bar{A}'_{I^*} + (2\bar{w}_{US} - M_{49}M_{58}) \bar{U}'_{USR} + \\ & \left(\frac{-k\bar{\rho}_2\bar{V}_2}{2} - M_{49}M_{59} \right) \bar{U}'_{USI} + \left(\frac{-k\rho_2\bar{V}'_2}{2} - M_{49}M_{60} \right) \bar{a}'_{I^*} + \\ & (-2\bar{\rho}^* - M_{49}M_{61}) \bar{a}'_{R^*} + \left(\frac{-kU_2\bar{V}_2}{2} - M_{49}M_{62} \right) \bar{\rho}'_{I^*} + (-M_{49}M_{63}) \bar{\rho}'_{R^*} + \\ & \left. (-k\bar{\rho}_2\bar{U}_2 - M_{49}M_{64}) \bar{V}'_{2I} + (-M_{49}M_{65}) \bar{V}'_{2R} \right]\end{aligned}\quad (C12)$$

Further, let

$$M_{66} = \frac{M_{50} - M_{49}M_{52}}{M_{48}}$$

$$M_{67} = -\frac{M_{49}M_{54}}{M_{48}}$$

$$M_{68} = \frac{-2 - M_{49}M_{54}}{M_{48}}$$

$$M_{69} = \frac{-M_{49}M_{55}}{M_{48}}$$

$$M_{70} = (M_{50A} - M_{49}M_{56})/M_{48}$$

$$M_{71} = -M_{49}M_{57}/M_{48}$$

$$M_{72} = (2\bar{w}_{US} - M_{49}M_{58})/M_{48}$$

$$M_{73} = (-\frac{1}{2}k\bar{\rho}_2\bar{V}_2 - M_{49}M_{59})/M_{48}$$

$$M_{74} = (-\frac{1}{2}k\rho_2\bar{V}'_2 - M_{49}M_{60})/M_{48}$$

$$M_{75} = (-2\bar{\rho}^*\bar{a}^* - M_{49}M_{61})/M_{48}$$

$$\begin{aligned}
M_{76} &= (-\frac{1}{2}k\bar{U}_2\bar{V}_2 - M_{49}M_{62})/M_{48} \\
M_{77} &= -M_{49}M_{63}/M_{48} \\
M_{78} &= (-k(\bar{\rho}\bar{U})_2 - M_{49}M_{63})/M_{48} \\
M_{79} &= -M_{49}M_{64}/M_{48}
\end{aligned} \tag{C13}$$

Substituting Equation C13 into Equation C12 gives

$$\begin{aligned}
\rho'_{USR} &= M_{76}\bar{A}'_{SR} + M_{77}\bar{A}'_{SI} + M_{78}\bar{P}'_{R^*} + M_{79}\bar{P}'_{I^*} + M_{80}\bar{A}'_{R^*} + M_{81}\bar{A}'_{I^*} + \\
&M_{82}\bar{U}'_{USR} + M_{83}\bar{U}'_{USI} + M_{84}\bar{a}'_{I^*} + M_{85}\bar{a}'_{R^*} + M_{86}\bar{\rho}'_{I^*} + M_{87}\bar{\rho}'_{R^*} + \\
&M_{88}\bar{V}'_{2I} + M_{89}\bar{V}'_{2R}
\end{aligned} \tag{C14}$$

Next, the continuity equation is used to solve for the velocity perturbation upstream of the shock. Separating Equation 38 into real and imaginary parts gives

Real Part

$$\begin{aligned}
\bar{\rho}_{US}\bar{U}'_{USR}\bar{A}_s &= -(\bar{\rho}'_{USR}\bar{U}_{US}\bar{A}_s + \bar{\rho}_{US}\bar{U}_{US}\bar{A}'_{SR}) + (\bar{\rho}'_{R^*}\bar{a}^* + \bar{\rho}^*\bar{a}'_{R^*} + \bar{\rho}^*\bar{A}'_{R^*}\bar{a}^*) + \\
k \left[\frac{V_2}{2} (\bar{\rho}'_{I^*} + \bar{\rho}'_{USI}) + \bar{\rho}_2\bar{V}'_{2I} \right]
\end{aligned} \tag{C15}$$

Imaginary Part

$$\begin{aligned}
\bar{\rho}_{US}\bar{U}'_{USI}\bar{A}_s &= -(\bar{\rho}'_{USI}\bar{U}_{US}\bar{A}_s + \bar{\rho}_{US}\bar{U}_{US}\bar{A}'_{SI}) + (\bar{\rho}'_{I^*}\bar{a}^* + \bar{\rho}^*\bar{a}'_{I^*} + \bar{\rho}^*\bar{A}'_{I^*}\bar{a}^*) - \\
k \left[\frac{\bar{V}_2}{2} (\bar{\rho}'_{R^*} + \bar{\rho}'_{USR}) + \bar{\rho}_2\bar{V}'_{2R} \right]
\end{aligned} \tag{C16}$$

Substituting for $\bar{\rho}'_{USR}$ and $\bar{\rho}'_{USI}$ in Equation C15 produces

$$\begin{aligned}
\bar{\rho}_{US}\bar{U}'_{USR}\bar{A}_s &= -\bar{U}_{US}\bar{A}_s (M_{66}\bar{A}'_{SR} + M_{67}\bar{A}'_{SI} + M_{68}\bar{P}'_{R^*} + M_{69}\bar{P}'_{I^*} + M_{70}\bar{A}'_{R^*} + \\
&M_{71}\bar{A}'_{I^*} + M_{72}\bar{U}'_{USR} + M_{73}\bar{U}'_{USI} + M_{74}\bar{a}'_{I^*} + M_{75}\bar{a}'_{R^*} + M_{76}\bar{\rho}'_{I^*} + M_{77}\bar{\rho}'_{R^*} + \\
&M_{78}\bar{V}'_{2I} + M_{79}\bar{V}'_{2R}) - \bar{\rho}_{US}\bar{U}_{US}\bar{A}'_{SR} + (\bar{\rho}'_{R^*}\bar{a}^* + \bar{\rho}^*\bar{a}'_{R^*} + \bar{\rho}^*\bar{A}'_{R^*}\bar{a}^*) + \\
&\frac{k\bar{V}_2}{2}\bar{\rho}'_{I^*} + \frac{k\bar{V}_2}{2}(M_{52}\bar{A}'_{SR} + M_{53}\bar{A}'_{SI} + M_{54}\bar{P}'_{R^*} + M_{55}\bar{P}'_{I^*} + M_{56}\bar{A}'_{R^*} + \\
&M_{57}\bar{A}'_{I^*} + M_{58}\bar{U}'_{USR} + M_{59}\bar{U}'_{USI} + M_{60}\bar{a}'_{I^*} + M_{61}\bar{a}'_{R^*} + M_{62}\bar{\rho}'_{I^*} + M_{63}\bar{\rho}'_{R^*} + \\
&M_{64}\bar{V}'_{2I} + M_{65}\bar{V}'_{2R}) + k\bar{\rho}_2\bar{V}'_{2I}
\end{aligned}$$

Next, combine like terms and rearrange so that $(\bar{U}'_{\text{USR}}$ and $\bar{U}'_{\text{USI}})$ is on the left-hand side of the equation and let:

$$\begin{aligned}
M_{80} &= \bar{\rho}_{\text{US}} \bar{A}_S + \bar{U}_{\text{US}} \bar{A}_S M_{72} - \frac{k\bar{V}_2}{2} M_{58} \\
M_{81} &= \bar{U}_{\text{US}} \bar{A}_S M_{73} - \frac{k\bar{V}_2}{2} M_{59} \\
M_{82} &= -\bar{U}_{\text{US}} \bar{A}_S M_{66} - \bar{\rho}_{\text{US}} \bar{U}_{\text{US}} + \frac{k\bar{V}_2}{2} M_{52} \\
M_{83} &= -\bar{U}_{\text{US}} \bar{A}_S M_{67} + \frac{k\bar{V}_2}{2} M_{53} \\
M_{84} &= -\bar{U}_{\text{US}} \bar{A}_S M_{77} + \bar{a}^* + \frac{k\bar{V}_2}{2} M_{63} \\
M_{85} &= -\bar{U}_{\text{US}} \bar{A}_S M_{76} + \frac{k\bar{V}_2}{2} M_{62} \\
M_{86} &= -\bar{U}_{\text{US}} \bar{A}_S M_{68} + \frac{k\bar{V}_2}{2} M_{54} \\
M_{87} &= -\bar{U}_{\text{US}} \bar{A}_S M_{69} + \frac{k\bar{V}_2}{2} M_{55} \\
M_{88} &= -\bar{U}_{\text{US}} \bar{A}_S M_{70} + \bar{\rho}^* \bar{a}^* + \frac{k\bar{V}_2}{2} M_{56} \\
M_{89} &= -\bar{U}_{\text{US}} \bar{A}_S M_{71} + \frac{k\bar{V}_2}{2} M_{57} \\
M_{90} &= -\bar{U}_{\text{US}} \bar{A}_S M_{74} + \frac{k\bar{V}_2}{2} M_{60} \\
M_{91} &= -\bar{U}_{\text{US}} \bar{A}_S M_{75} + \bar{\rho}^* + \frac{k\bar{V}_2}{2} M_{61} \\
M_{92} &= -\bar{U}_{\text{US}} \bar{A}_S M_{78} + \frac{k\bar{V}_2}{2} M_{64} + k\bar{\rho}_2 \\
M_{93} &= -\bar{U}_{\text{US}} \bar{A}_S M_{79} + \frac{k\bar{V}_2}{2} M_{65}
\end{aligned} \tag{C17}$$

Using these relationships, the real part becomes

$$\begin{aligned} M_{80}\bar{U}'_{USR} + M_{81}\bar{U}'_{USI} &= M_{82}\bar{A}'_{SR} + M_{83}\bar{A}'_{SI} + M_{84}\bar{\rho}'_R^* + M_{85}\bar{\rho}'_I^* + M_{86}\bar{\rho}'_R^* + \\ &M_{87}\bar{\rho}'_I^* + M_{88}\bar{A}'_R^* + M_{89}\bar{A}'_I^* + M_{90}\bar{a}'_I^* + M_{91}\bar{a}'_R^* + M_{92}\bar{V}'_{2I} + M_{93}\bar{V}'_{2R} \end{aligned} \quad (C18)$$

Dealing next with the imaginary part, substituting for $\bar{\rho}'_{USI}$ and $\bar{\rho}'_{USR}$ gives

$$\begin{aligned} \bar{\rho}_{US}\bar{U}'_{USI}\bar{A}_S &= -\bar{\rho}_{US}\bar{U}_{US}\bar{A}_{SI} + (\bar{\rho}'_I^*\bar{a}^* + \bar{\rho}^*\bar{a}'_I^* + \bar{\rho}^*\bar{A}'_I^*\bar{a}^*) - \\ &k \left(\frac{V_2}{2} \bar{\rho}'_R^* + \bar{\rho}\bar{V}'_{2R} \right) - \bar{U}_{US}\bar{A}_S (M_{52}\bar{A}'_{SR} + M_{53}\bar{A}'_{SI} + M_{54}\bar{\rho}'_R^* + M_{55}\bar{\rho}'_I^* + \\ &M_{56}\bar{A}'_R^* + M_{57}\bar{A}'_I^* + M_{58}\bar{U}'_{USR} + M_{59}\bar{U}'_{USI} + M_{60}\bar{a}'_I^* + M_{61}\bar{a}'_R^* + \\ &M_{62}\bar{\rho}'_I^* + M_{63}\bar{\rho}'_R^* + M_{64}\bar{V}'_{2I} + M_{65}\bar{V}'_{2R}) - \frac{k\bar{V}_2}{2} (M_{66}\bar{A}'_{SR} + M_{67}\bar{A}'_{SI} + \\ &M_{68}\bar{\rho}'_R^* + M_{69}\bar{\rho}'_I^* + M_{70}\bar{A}'_R^* + M_{71}\bar{A}'_I^* + M_{72}\bar{U}'_{USR} + M_{73}\bar{U}_{USI} + M_{74}\bar{a}'_I^* + \\ &M_{75}\bar{a}'_R^* + M_{76}\bar{\rho}'_I^* + M_{77}\bar{\rho}'_R^* + M_{78}\bar{V}'_{2I} + M_{79}\bar{V}'_{2R}) \end{aligned}$$

Now, let

$$\begin{aligned} C_{58} &= \bar{\rho}_{US}\bar{A}_S + \bar{U}_{US}\bar{A}_S M_{59} + \frac{k\bar{V}_2}{2} M_{73} \\ C_{59} &= \bar{U}_{US}\bar{A}_S M_{58} + \frac{k\bar{V}_2}{2} M_{72} \\ C_{60} &= -\bar{U}_{US}\bar{A}_S M_{52} + \frac{k\bar{V}_2}{2} M_{66} \\ C_{61} &= -\bar{\rho}_{US}\bar{U}_{US} - \bar{U}_{US}\bar{A}_S M_{53} - \frac{k\bar{V}_2}{2} M_{67} \\ C_{62} &= -\bar{U}_{US}\bar{A}_S M_{63} - \frac{k\bar{V}_2}{2} (1 + M_{77}) \\ C_{63} &= \bar{a}^* - \bar{U}_{US}\bar{A}_S M_{62} - \frac{k\bar{V}_2}{2} M_{76} \\ C_{64} &= -\bar{U}_{US}\bar{A}_S M_{61} - \frac{k\bar{V}_2}{2} M_{75} \\ C_{65} &= \bar{\rho}^* - \bar{U}_{US}\bar{A}_S M_{60} - \frac{k\bar{V}_2}{2} M_{74} \\ C_{66} &= -\bar{U}_{US}\bar{A}_S M_{54} - \frac{k\bar{V}_2}{2} M_{68} \\ C_{67} &= -\bar{U}_{US}\bar{A}_S M_{55} - \frac{k\bar{V}_2}{2} M_{69} \end{aligned}$$

$$\begin{aligned}
C_{68} &= -\bar{U}_{US}\bar{A}_S M_{56} - \frac{k\bar{V}_2}{2} M_{70} \\
C_{69} &= \bar{\rho}^* \bar{a}^* - \bar{U}_{US}\bar{A}_S M_{57} - \frac{k\bar{V}_2}{2} M_{71} \\
C_{70} &= -\bar{U}_{US}\bar{A}_S M_{65} - \frac{k\bar{V}_2}{2} (2\bar{\rho}_2 + M_{79}) \\
C_{71} &= -\bar{U}_{US}\bar{A}_S M_{64} - \frac{k\bar{V}_2}{2} M_{78}
\end{aligned} \tag{C19}$$

Substituting these relationships into the imaginary part of the continuity equation yields

$$\begin{aligned}
C_{59}\bar{U}'_{USR} + C_{98}\bar{U}'_{US1} &= C_{60}\bar{A}'_{SR} + C_{61}\bar{A}'_{S1} + C_{62}\bar{\rho}'_{R^*} + C_{63}\bar{\rho}'_{I^*} + C_{64}\bar{a}'_{R^*} + C_{65}\bar{a}'_{I^*} + \\
&C_{66}\bar{\rho}'_{R^*} + C_{67}\bar{\rho}'_{I^*} + C_{68}\bar{A}'_{R^*} + C_{69}\bar{A}'_{I^*} + C_{70}\bar{V}'_{2I} + C_{71}\bar{V}'_{2R}
\end{aligned} \tag{C20}$$

Multiplying Equation C18 by C_{58}/M_{81} and subtracting the resultant from Equation C20 produces, after solving for \bar{U}'_{USR} :

$$\begin{aligned}
\bar{U}'_{USR} &= C_{73}\bar{A}'_{SR} + C_{74}\bar{A}'_{S1} + C_{75}\bar{\rho}'_{R^*} + C_{76}\bar{\rho}'_{I^*} + C_{77}\bar{a}'_{R^*} + C_{78}\bar{a}'_{I^*} + \\
&C_{79}\bar{\rho}'_{R^*} + C_{80}\bar{\rho}'_{I^*} + C_{81}\bar{A}'_{R^*} + C_{82}\bar{A}'_{I^*} + C_{83}\bar{V}'_{2I} + C_{84}\bar{V}'_{2R}
\end{aligned} \tag{C21}$$

where,

$$\begin{aligned}
C_{72} &= C_{69} - \frac{C_{58}M_{80}}{M_{81}} \\
C_{73} &= \frac{1}{C_{72}} \left(C_{60} - \frac{C_{58}M_{82}}{M_{81}} \right) \\
C_{74} &= \frac{1}{C_{72}} \left(C_{61} - \frac{C_{58}M_{83}}{M_{81}} \right) \\
C_{75} &= \frac{1}{C_{72}} \left(C_{62} - \frac{C_{58}M_{84}}{M_{81}} \right) \\
C_{76} &= \frac{1}{C_{72}} \left(C_{63} - \frac{C_{58}M_{85}}{M_{81}} \right) \\
C_{77} &= \frac{1}{C_{72}} \left(C_{64} - \frac{C_{58}M_{86}}{M_{81}} \right) \\
C_{78} &= \frac{1}{C_{72}} \left(C_{65} - \frac{C_{58}M_{87}}{M_{81}} \right) \\
C_{79} &= \frac{1}{C_{72}} \left(C_{66} - \frac{C_{58}M_{88}}{M_{81}} \right) \\
C_{80} &= \frac{1}{C_{72}} \left(C_{67} - \frac{C_{58}M_{89}}{M_{81}} \right)
\end{aligned}$$

$$\begin{aligned}
C_{81} &= \frac{1}{C_{72}} \left(C_{68} - \frac{C_{58}M_{88}}{M_{81}} \right) \\
C_{82} &= \frac{1}{C_{72}} \left(C_{69} - \frac{C_{58}M_{89}}{M_{81}} \right) \\
C_{83} &= \frac{1}{C_{72}} \left(C_{70} - \frac{C_{58}M_{93}}{M_{81}} \right) \\
C_{84} &= \frac{1}{C_{72}} \left(C_{71} - \frac{C_{58}M_{92}}{M_{81}} \right)
\end{aligned}$$

Substituting Equation C21 into C20 and solving for \bar{U}'_{USI} produces

$$\begin{aligned}
\bar{U}'_{USI} &= C_{85}\bar{A}'_{SR} + C_{86}\bar{A}'_{SI} + C_{87}\bar{\rho}_R^{*'} + C_{88}\bar{\rho}_1^{*'} + C_{89}\bar{a}_R^{*'} + C_{90}\bar{a}_1^{*'} + \\
&C_{91}\bar{p}_R^{*'} + C_{92}\bar{p}_1^{*'} + C_{93}\bar{A}_R^{*'} + C_{94}\bar{A}_1^{*'} + C_{95}\bar{V}_{21}' + C_{96}\bar{V}_{2R}'
\end{aligned} \tag{C22}$$

where,

$$\begin{aligned}
C_{85} &= \frac{1}{C_{58}} (C_{60} - C_{59}C_{73}) \\
C_{86} &= \frac{1}{C_{58}} (C_{61} - C_{59}C_{74}) \\
C_{87} &= \frac{1}{C_{58}} (C_{62} - C_{59}C_{75}) \\
C_{88} &= \frac{1}{C_{58}} (C_{63} - C_{59}C_{76}) \\
C_{89} &= \frac{1}{C_{58}} (C_{64} - C_{59}C_{77}) \\
C_{90} &= \frac{1}{C_{58}} (C_{65} - C_{59}C_{78}) \\
C_{91} &= \frac{1}{C_{58}} (C_{66} - C_{59}C_{79}) \\
C_{92} &= \frac{1}{C_{58}} (C_{67} - C_{59}C_{80}) \\
C_{93} &= \frac{1}{C_{58}} (C_{68} - C_{59}C_{81}) \\
C_{94} &= \frac{1}{C_{58}} (C_{69} - C_{59}C_{82}) \\
C_{95} &= \frac{1}{C_{58}} (C_{70} - C_{59}C_{83}) \\
C_{96} &= \frac{1}{C_{58}} (C_{71} - C_{59}C_{84})
\end{aligned} \tag{C23}$$

The value for the velocity perturbation upstream of the shock given in Equations C21 and C22 can then be inserted into Equations C11 and C14, given the density perturbation upstream of the shock, for a complete description of the flowfield in Region 2.

APPENDIX D
SOLUTION TO THE EQUATIONS OF MOTION FOR REGION 3

The nondimensionalized equations of motion appear as Equations 60 and 61. Again, only the momentum and continuity equations are needed because these are only two unknowns, $A_{+\infty}$ and $Z_{+\infty}$. These parameters describe the irrotational and rotational flowfields.

Exit flow perturbations are defined first. The velocity is represented as the sum of the rotational and irrotational fields, while pressure is a function of the irrotational flowfield only¹. Conversion of the two-dimensional flowfield into a one-dimensional field proceeds as follows:

$$p'_E = \frac{1}{\bar{y}_E} \int_0^{y_E} p'_E dy$$

which produces

$$p'_E = -\bar{\rho}_E \left[\frac{k + \bar{u}_E B_2 + \bar{v}_E C}{C \bar{y}_E} \right] \left[A_{+\infty} e^{ikt} (e^{iC\bar{y}_E} - 1) \right] \quad (D1)$$

The density perturbation can be related through isentropic form relationships to the pressure perturbations as

$$\bar{\rho}'_E = -\bar{\rho}_E \left[\frac{k + \bar{u}_E B_2 + \bar{v}_E C}{C \bar{y}_E \bar{a}_E^2} \right] \left[A_{+\infty} e^{ikt} (e^{iC\bar{y}_E} - 1) \right] \quad (D2)$$

The expressions for the velocity component are listed below:

Irrotational Field in the x-direction

$$\bar{u}'_{IE} = \frac{B_2 A_{+\infty}}{C \bar{y}_E} e^{ikt} (e^{iC\bar{y}_E} - 1) \quad (D3)$$

$$\bar{v}'_{IE} = \frac{A_{+\infty}}{\bar{y}_E} e^{ikt} (e^{iC\bar{y}_E} - 1) \quad (D4)$$

Rotational Field

$$\bar{u}'_{RE} = \frac{Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} e^{ikt} (e^{iC\bar{y}_E} - 1) \quad (D5)$$

$$\bar{v}'_{RE} = \frac{-R Z_{+\infty}}{C (R^2 + C^2) \bar{y}_E} e^{ikt} (e^{iC\bar{y}_E} - 1) \quad (D6)$$

Next, expressions for the time derivatives of the exit density and velocity perturbations are obtained in the following manner:

Density

$$\frac{\partial \bar{\rho}'_E}{\partial t} = -\bar{\rho}_E \left[\frac{i k (k + \bar{u}_E B_2 + \bar{v}_E C)}{C \bar{a}_E^2 \bar{y}_E} \right] \left[A_{+\infty} e^{ikt} (e^{iC\bar{y}_E} - 1) \right] \quad (D7)$$

¹Goldstein, M. E., *Aeroacoustics*, pp. 220, 221.

Velocity

Irrotational Component

$$\frac{\partial \bar{u}'_{1E}}{\partial t} = \left(\frac{i k B_2 A_{+\infty}}{C \bar{y}_E} \right) e^{ikt} (e^{iC\bar{y}_k} - 1) \quad (D8)$$

$$\frac{\partial \bar{v}'_{1E}}{\partial t} = \left(\frac{i k A_{+\infty}}{\bar{y}_E} \right) e^{ikt} (e^{iC\bar{y}_k} - 1) \quad (D9)$$

Rotational Component

$$\frac{\partial \bar{u}'_{RE}}{\partial t} = \left[\frac{i k Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} e^{ikt} (e^{iC\bar{y}_k} - 1) \right] \quad (D10)$$

$$\frac{\partial \bar{v}'_{RE}}{\partial t} = \left[\frac{-i k R Z_{+\infty}}{C (C^2 + R^2) \bar{y}_E} e^{ikt} (e^{iC\bar{y}_k} - 1) \right] \quad (D11)$$

Assuming the flow parameters vary harmonically with time, substituting the relationships of Equations D3 through D11 into the momentum equation, and dividing through by e^{ikt} produces

$$\begin{aligned} & \bar{\rho}'_{ds} \bar{A}_s + \bar{\rho}_{ds} \bar{A}'_s + \bar{\rho}_E \left[\left(\frac{k + \bar{u}_E B_2 + \bar{v}_E C}{C \bar{y}_E} \right) \bar{A}_E A_{+\infty} (e^{iC\bar{y}_k} - 1) \right] - \bar{\rho}_E \bar{A}'_E \\ & = \left\{ \left[\left(\frac{B_2 A_{+\infty}}{C \bar{y}_E} + \frac{Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \sin \alpha_{ch} + \right. \right. \\ & \quad \left. \left(\frac{A_{+\infty}}{\bar{y}_E} - \frac{R Z_{+\infty}}{C (R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \cos \alpha_{ch} \right] \bar{w}_E + \\ & \quad U_{2REL} \cos (\alpha_{ch} - \beta_2) \left[-\bar{\rho}_E \left(\frac{k + \bar{u}_E B_2 + \bar{v}_E C}{C \bar{y}_E \bar{a}_E^2} \right) \bar{A}_E A_{+\infty} (e^{iC\bar{y}_k} - \right. \\ & \quad \left. 1) \cos (\alpha_{ch} - \beta_2) \bar{U}_{2REL} + \bar{\rho}_E \bar{A}_E \left(\left(\frac{B_2 A_{+\infty}}{C \bar{y}_E} + \frac{Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \sin \alpha_{ch} + \right. \right. \\ & \quad \left. \left. \left(\frac{A_{+\infty}}{\bar{y}_E} - \frac{R Z_{+\infty}}{C (R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \cos \alpha_{ch} \right) \right] + \\ & \quad \left. \bar{\rho}_E \bar{U}_E \bar{A}'_E \cos (\alpha_{ch} - \beta_2) \right\} - \{ \bar{U}'_{ds} \bar{w}_{ds} + \bar{U}_{ds} \bar{w}'_{ds} \} + \\ & \quad \frac{\bar{U}_3 \bar{V}_3}{2} \left\{ ik \bar{\rho}'_{ds} - ik \bar{\rho}_E \left(\frac{k + \bar{u}_E B_2 + \bar{v}_E C}{C \bar{y}_E \bar{a}_E^2} \right) A_{+\infty} (e^{iC\bar{y}_k} - 1) \right\} + \\ & \quad + ik (\bar{\rho} \bar{U} \bar{V})_3 + \frac{ik \bar{\rho}_3 \bar{V}_3}{2} \left\{ \left(\frac{B_2 A_{+\infty}}{C \bar{y}_E} + \frac{Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \sin \alpha_{ch} + \right. \\ & \quad \left. \left(\frac{A_{+\infty}}{\bar{y}_E} - \frac{R Z_{+\infty}}{C (R^2 + C^2) \bar{y}_E} \right) (e^{iC\bar{y}_k} - 1) \cos \alpha_{ch} + \bar{U}'_{ds} \right\} \quad (D12) \end{aligned}$$

The next step requires that Equation D12 be divided into its real and imaginary parts. The momentum equation is used to solve for Z_{+coR} , Z_{+coI} and the continuity equation is used to solve for A_{+coR} , A_{+coI} . Combine like terms and make the following assumptions:

$$\begin{aligned}
M_{94} &= \frac{\bar{u}_E \bar{A}_E}{C \bar{y}_E} \bar{\rho}_E - \frac{\bar{w}_E \sin \alpha_{ch}}{C \bar{y}_E} + \\
&\quad \left[U_{2REL} \cos (\alpha_{ch} - \beta_2) \right]^2 \left[\frac{U_E \bar{A}_E}{C \bar{y}_E \bar{a}_E^2} \bar{\rho}_E \right] - \\
&\quad \left[\left(\frac{\bar{\rho}_E U_{2REL} \bar{A}_E}{C \bar{y}_E} \right) (\sin \alpha_{ch}) \right] \cos (\alpha_{ch} - \beta_2) \\
M_{95} &= \bar{\rho}_E \left(\frac{k + \bar{v}_E C}{C \bar{y}_E} \right) \bar{A}_E - \frac{\bar{w}_E \cos \alpha_{ch}}{\bar{y}_E} + \\
&\quad \left[U_{2REL} \cos (\alpha_{ch} - \beta_2) \right]^2 \left[\left(\frac{k + \bar{v}_E C}{C \bar{y}_E \bar{a}_E^2} \right) \bar{A}_E \bar{\rho}_E \right] - \\
&\quad \bar{\rho}_E \bar{U}_{2REL} \bar{A}_E \cos \alpha_{ch} \cos (\alpha_{ch} - \beta_2) \\
M_{96} &= \frac{\bar{\rho}_E \bar{u}_E \bar{U}_3 \bar{V}_3}{2 C \bar{y}_E \bar{a}_E^2} - \frac{\rho_3 \bar{V}_3 \sin \alpha_{ch}}{2 C \bar{y}_E} \\
M_{97} &= \frac{\bar{\rho}_E (k + \bar{v}_E C) \bar{U}_3 \bar{V}_3}{2 C \bar{y}_E \bar{a}_E^2} - \frac{\rho_3 \bar{V}_3 \cos \alpha_{ch}}{2 \bar{y}_E} \\
M_{98} &= \left[\frac{\bar{w}_E \sin \alpha_{ch}}{(R^2 + C^2) \bar{y}_E} - \frac{R \bar{w}_E \cos \alpha_{ch}}{C(R^2 + C^2) \bar{y}_E} \right] + \\
&\quad \left[\frac{\bar{\rho}_E U_{2REL} \bar{A}_E \cos (\alpha_{ch} - \beta_2)}{(R^2 + C^2) \bar{y}_E} \right] \left[\sin \alpha_{ch} - \frac{R \cos \alpha_{ch}}{C} \right] \\
M_{99} &= \frac{\bar{\rho}_3 \bar{V}_3}{2 (R^2 + C^2) \bar{y}_E} \left(\sin \alpha_{ch} - \frac{R \cos \alpha_{ch}}{C} \right) \\
M_{100} &= \cos (C \bar{y}_E) - 1
\end{aligned} \tag{D13}$$

Substituting these relationships into equation D12 and extracting the real part yields

$$\begin{aligned}
& \mathbf{M}_{94} [\mathbf{B}_{2R} \mathbf{A}_{+\infty R} \mathbf{M}_{100} - \mathbf{B}_{2I} \mathbf{A}_{+\infty I} \mathbf{M}_{100} - \mathbf{B}_{2R} \mathbf{A}_{+\infty} \sin(C\bar{y}_E) - \\
& \quad \mathbf{B}_{2I} \mathbf{A}_{+\infty R} \sin(C\bar{y}_E)] + \mathbf{M}_{95} [\mathbf{A}_{+\infty R} \mathbf{M}_{100} - \mathbf{A}_{+\infty I} \sin(C\bar{y}_E)] - \\
& \quad k \mathbf{M}_{96} [\mathbf{B}_{2R} \mathbf{A}_{+\infty I} \mathbf{M}_{100} + \mathbf{B}_{2I} \mathbf{A}_{+\infty R} \mathbf{M}_{100} + \mathbf{B}_{2R} \mathbf{A}_{+\infty R} \sin(C\bar{y}_E) - \\
& \quad \mathbf{B}_{2I} \mathbf{A}_{+\infty I} \sin(C\bar{y}_E)] - k \mathbf{M}_{97} [\mathbf{A}_{+\infty I} \mathbf{M}_{100} + \mathbf{A}_{+\infty R} \sin(C\bar{y}_E)] \\
& = -k \mathbf{M}_{99} [\mathbf{Z}_{+\infty R} \mathbf{M}_{100} + \mathbf{Z}_{+\infty I} \sin(C\bar{y}_E)] + \mathbf{M}_{98} [\mathbf{Z}_{+\infty I} \mathbf{M}_{100} - \\
& \quad \mathbf{Z}_{+\infty R} \sin(C\bar{y}_E)] - \bar{p}'_{daR} \bar{\mathbf{A}}_s - \bar{p}_{da} \bar{\mathbf{A}}'_{sR} + [\bar{p}_E + \bar{\rho}_E (U_{2RE} \cos(\alpha_{ch} - \beta_2))^2] \bar{\mathbf{A}}'_{ER} - \\
& \quad \bar{\mathbf{U}}'_{daR} \bar{\mathbf{W}}_{da} - \bar{\mathbf{U}}_{da} \bar{\mathbf{W}}'_{daR} - \frac{k \bar{\rho}'_{daI} \bar{\mathbf{U}}_3 \bar{\mathbf{V}}_3}{2} - k \bar{\rho}_3 \bar{\mathbf{U}}_3 \bar{\mathbf{V}}'_{3I} - \frac{k \bar{\rho}_3 \bar{\mathbf{U}}'_{daI} \bar{\mathbf{V}}_3}{2} \tag{D14}
\end{aligned}$$

Collect like terms in the following manner and let

$$\begin{aligned}
\mathbf{M}_{101} &= \mathbf{M}_{94} (\mathbf{M}_{100} \mathbf{B}_{2R} - \mathbf{B}_{2I} \sin(C\bar{y}_E)) + \mathbf{M}_{95} \mathbf{M}_{100} \\
&\quad - k \mathbf{M}_{96} (\mathbf{B}_{2I} \mathbf{M}_{100} + \mathbf{B}_{2R} \sin(C\bar{y}_E)) - k \mathbf{M}_{97} \sin(C\bar{y}_E) \\
\mathbf{M}_{102} &= -\mathbf{M}_{94} (\mathbf{M}_{100} \mathbf{B}_{2I} + \mathbf{B}_{2R} \sin(C\bar{y}_E)) - \mathbf{M}_{95} \sin(C\bar{y}_E) \\
&\quad - k \mathbf{M}_{96} (\mathbf{M}_{100} \mathbf{B}_{2R} - \mathbf{B}_{2I} \sin(C\bar{y}_E)) - k \mathbf{M}_{97} \mathbf{M}_{100} \\
\mathbf{M}_{103} &= -k \mathbf{M}_{99} \sin(C\bar{y}_E) + \mathbf{M}_{98} \mathbf{M}_{100} \\
\mathbf{M}_{104} &= -k \mathbf{M}_{99} \mathbf{M}_{100} - \mathbf{M}_{98} \sin(C\bar{y}_E) \tag{D15}
\end{aligned}$$

Substituting Equation D15 into D14 yields

$$\mathbf{M}_{101} \mathbf{A}_{+\infty R} + \mathbf{M}_{102} \mathbf{A}_{+\infty I} = \mathbf{M}_{103} \mathbf{Z}_{+\infty R} + \mathbf{M}_{104} \mathbf{Z}_{+\infty I} + \text{RC} \tag{D16}$$

where,

$$\begin{aligned}
\text{RC} &= -\bar{p}'_{daR} \bar{\mathbf{A}}_s - \bar{p}_{da} \bar{\mathbf{A}}'_{sR} + [\bar{p}_E + \bar{\rho}_E (U_{2RE} \cos(\alpha_{ch} - \beta_2))^2] \bar{\mathbf{A}}'_{ER} - \\
&\quad \bar{\mathbf{U}}'_{daR} \bar{\mathbf{W}}_{da} - \bar{\mathbf{U}}_{da} \bar{\mathbf{W}}'_{daR} - \frac{k \bar{\rho}'_{daI} \bar{\mathbf{U}}_3 \bar{\mathbf{V}}_3}{2} - k \bar{\rho}_3 \bar{\mathbf{U}}_3 \bar{\mathbf{V}}'_{3I} - \frac{k \bar{\rho}_3 \bar{\mathbf{U}}'_{daI} \bar{\mathbf{V}}_3}{2} \tag{D17}
\end{aligned}$$

with

$$\bar{\mathbf{W}}'_{daR} = (\bar{p}'_{daR} \bar{\mathbf{U}}_{da} \bar{\mathbf{A}}_s + \bar{p}_{da} \bar{\mathbf{U}}'_{daR} \bar{\mathbf{A}}_s + \bar{p}_{da} \bar{\mathbf{U}}_{da} \bar{\mathbf{A}}'_{sR}) \tag{D18}$$

The imaginary part of the momentum equation is expressed as

$$M_{105}A_{+\infty R} + M_{106}A_{+\infty I} = M_{107}Z_{+\infty R} + M_{108}Z_{+\infty I} + IC \quad (D19)$$

where,

$$\begin{aligned} M_{105} &= M_{94} [B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)] + M_{95} \sin (C\bar{y}_E) \\ &\quad + kM_{96} [B_{2R}M_{100} - B_{2I} \sin (C\bar{y}_E)] + kM_{97}M_{100} \\ M_{106} &= M_{94} [B_{2R}M_{100} - B_{2I} \sin (C\bar{y}_E)] + M_{95}M_{100} \\ &\quad - kM_{96} [B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)] - kM_{97} \sin (C\bar{y}_E) \\ M_{107} &= M_{98} \sin (C\bar{y}_E) + kM_{99}M_{100} \\ M_{108} &= M_{70}M_{72} - kM_{71} \sin (C\bar{y}_E) \end{aligned} \quad (D20)$$

and

$$\begin{aligned} IC &= -\rho_{ds} \bar{A}_S - \bar{\rho}_{ds} \bar{A}'_{S1} + [\bar{p}_E + \bar{\rho}_E (\bar{U}_{2REL} \cos (\alpha_{ch} - \beta_2))^2] \bar{A}'_{E1} - \\ &\quad \bar{U}'_{ds1} \bar{w}_{ds} - \bar{U}_{ds} \bar{w}'_{ds1} + \frac{k\bar{\rho}'_{dsR} \bar{U}_3 \bar{V}_3}{2} + k\bar{\rho}_3 \bar{U}_3 \bar{V}'_{3R} + \frac{k\bar{\rho}_3 \bar{U}'_{dsR} \bar{V}_3}{2} \end{aligned} \quad (D21)$$

with

$$\bar{w}'_{ds1} = (\bar{\rho}'_{ds1} \bar{U}_{ds} \bar{A}_S + \bar{\rho}_{ds} \bar{U}'_{ds1} \bar{A}_S + \bar{\rho}_{ds} \bar{U}_{ds} \bar{A}'_{S1})$$

In order to solve for $A_{+\infty R}$, multiply Equation D16 by M_{78}/M_{74} and subtract the resultant from Equation D19 to produce

$$A_{+\infty R} = M_{110}Z_{+\infty R} + M_{111}Z_{+\infty I} + M_{112}IC + M_{113}RC \quad (D22)$$

where,

$$\begin{aligned} M_{109} &= M_{105} - M_{101}M_{106}/M_{102} \\ M_{110} &= (M_{107} - M_{103}M_{106}/M_{102})/M_{109} \\ M_{111} &= (M_{108} - M_{104}M_{106}/M_{102})/M_{109} \\ M_{112} &= 1/M_{109} \\ M_{113} &= -M_{106}/M_{102}M_{109} \end{aligned}$$

Substitute Equation D22 into Equation D19 and solve for $A_{+\infty I}$ to give

$$A_{+\infty I} = M_{114}Z_{+\infty R} + M_{115}Z_{+\infty I} + M_{116}RC + M_{117}IC \quad (D23)$$

where,

$$\begin{aligned}
M_{114} &= (M_{107} - M_{105}M_{110})/M_{106} \\
M_{115} &= (M_{108} - M_{105}M_{109})/M_{106} \\
M_{116} &= -M_{105}M_{113}/M_{106} \\
M_{117} &= (1 - M_{105}M_{113})/M_{106}
\end{aligned} \tag{D24}$$

The continuity equation is used to obtain $Z_{+\infty}$. The equation takes the following form after substituting Equation D2 through D7 into Equation 38 and dividing through by $e^{i\mathbf{k}\cdot\mathbf{r}}$.

$$\begin{aligned}
& \left[-\bar{\rho}_E \left(\frac{\mathbf{k} + \bar{u}_E \mathbf{B}_2 + \bar{v}_E \mathbf{C}}{C \bar{y}_E \bar{a}_E^2} \right) \bar{U}_E \bar{A}_E \bar{A}_{+\infty} (e^{iC\bar{y}_E} - 1) \cos(\alpha_{ch} - \beta_2) \right] + \\
& \bar{\rho}_E \bar{U}_E \bar{A}'_E + \bar{\rho}_E \bar{A}_E (e^{iC\bar{y}_E} - 1) \left[\left(\frac{B_2 A_{+\infty}}{C \bar{y}_E} + \frac{Z_{+\infty}}{(R^2 + C^2) \bar{y}_E} \right) \sin \alpha_{ch} + \right. \\
& \left. \left(\frac{A_{+\infty}}{\bar{y}_E} - \frac{R Z_{+\infty}}{C (R^2 + C^2) \bar{y}_E} \right) \cos \alpha_{ch} \right] - \bar{w}'_{ds} \\
& = i \bar{\rho}_3 \bar{V}'_3 - \frac{i \bar{k} \bar{V}_3}{2} \left[-\bar{\rho}_E \left(\frac{\mathbf{k} + \bar{u}_E \mathbf{B}_2 + \bar{v}_E \mathbf{C}}{C \bar{y}_E \bar{a}_E^2} \right) \bar{A}_{+\infty} (e^{iC\bar{y}_E} - 1) + \bar{\rho}'_{ds} \right]
\end{aligned} \tag{D25}$$

The following steps separate out and solve for the real and imaginary components of $\bar{A}_{+\infty}$ and $Z_{+\infty}$. First, make the assumptions noted in Equation D26 below.

$$\begin{aligned}
C_{97} &= -\frac{\bar{\rho}_E \bar{u}_E \bar{U}_{2REL} \bar{A}_E \cos(\alpha_{ch} - \beta_2)}{C \bar{y}_E \bar{a}_E^2} + \frac{\bar{\rho}_E \bar{A}_E \sin \alpha_{ch}}{C \bar{y}_E} \\
C_{98} &= -\frac{\bar{\rho}_E (\mathbf{k} + \bar{v}_E \mathbf{C}) \bar{U}_{2REL} \bar{A}_E \cos(\alpha_{ch} - \beta_2)}{C \bar{y}_E \bar{a}_E^2} + \frac{\bar{\rho}_E \bar{A}_E \cos \alpha_{ch}}{\bar{y}_E} \\
C_{99} &= \frac{\bar{\rho}_E \bar{A}_E}{(R^2 + C^2) \bar{y}_E} \left(\sin \alpha_{ch} - \frac{R}{C} \cos \alpha_{ch} \right) \\
C_{100} &= \frac{\bar{u}_E \bar{V}_3 \bar{\rho}_E}{2 C \bar{y}_E \bar{a}_E^2} \\
C_{101} &= \frac{(\mathbf{k} + \bar{v}_E \mathbf{C}) \bar{V}_3 \bar{\rho}_E}{2 C \bar{y}_E \bar{a}_E^2}
\end{aligned} \tag{D26}$$

Substituting Equation D26 into Equation D25 and rearranging term yields

$$C_{102}A_{+\infty R} + C_{103}A_{+\infty I} + C_{104}Z_{+\infty R} + C_{105}Z_{+\infty I} = CR \quad (D27)$$

where,

$$\begin{aligned} C_{102} &= C_{97}B_{2R}M_{100} - C_{97}B_{2I} \sin (C\bar{y}_E) + C_{98}M_{100} \\ &\quad + kC_{100} (B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)) + kC_{101} \sin (C\bar{y}_E) \\ C_{103} &= -C_{97} (B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)) - C_{98} \sin (C\bar{y}_E) \\ &\quad + kC_{100}(B_{2R}M_{100} - B_{2I} \sin (C\bar{y}_E)) + kC_{101}M_{100} \\ C_{104} &= C_{99}M_{100} \\ C_{105} &= -C_{99} \sin (C\bar{y}_E) \end{aligned} \quad (D28)$$

and

$$CR = -\bar{\rho}_E \bar{A}'_{ER} \bar{U}_E + w'_{dr} + \bar{\rho}_3 k \bar{V}'_{3I} + \frac{1}{2} k \bar{V}'_{3R} \bar{\rho}'_{dr} \quad (D29)$$

Substituting Equation D22 and D23 into Equation D27 and combining the terms gives

$$C_{106}Z_{+\infty R} + C_{107}Z_{+\infty I} + C_{108}RC + C_{109}IC = CR \quad (D30)$$

where,

$$\begin{aligned} C_{106} &= C_{102}M_{110} + C_{103}M_{114} + C_{104} \\ C_{107} &= C_{102}M_{111} + C_{103}M_{115} + C_{105} \\ C_{108} &= C_{102}M_{113} + C_{103}M_{116} \\ C_{109} &= C_{102}M_{112} + C_{103}M_{117} \end{aligned} \quad (D31)$$

Next, the imaginary part of the continuity equation is dealt with. Let,

$$\begin{aligned} C_{110} &= C_{97} [B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)] + C_{98} \sin C\bar{y}_E - \\ &\quad kC_{100} [B_{2R}M_{100} - B_{2I} \sin (C\bar{y}_E)] - k C_{101}M_{100} \\ C_{111} &= C_{97} [B_{2R}M_{100} - B_{2I} \sin (C\bar{y}_E)] + C_{98}M_{100} + \\ &\quad kC_{100} [B_{2I}M_{100} + B_{2R} \sin (C\bar{y}_E)] + k C_{101} \sin (C\bar{y}_E) \\ C_{112} &= C_{99} \sin (C\bar{y}_E) \\ C_{113} &= C_{99}M_{100} \end{aligned} \quad (D32)$$

This yields

$$C_{110}A_{+\infty R} + C_{111}A_{+\infty I} + C_{112}Z_{+\infty R} + C_{113}Z_{+\infty I} = CI \quad (D33)$$

where,

$$CI = -\bar{p}_E \bar{U}_{2REL} \cos(\alpha_{ch} - \beta_2) \bar{A}'_{EI} + \bar{w}'_{dal} - \bar{\rho}_3 k \bar{V}'_{3R} - \frac{1}{2} k \bar{V}_3 \bar{\rho}'_{daR} \quad (D34)$$

Next, substitute the relationships for $A_{+\infty}$ presented in Equations D22 and D23 into D33. Let,

$$\begin{aligned} C_{114} &= C_{110}M_{110} + C_{111}M_{114} + C_{112} \\ C_{115} &= C_{110}M_{111} + C_{111}M_{115} + C_{113} \\ C_{116} &= C_{110}M_{113} + C_{111}M_{116} \\ C_{117} &= C_{110}M_{112} + C_{111}M_{117} \end{aligned} \quad (D35)$$

This produces

$$C_{114}Z_{+\infty R} + C_{115}Z_{+\infty I} + C_{116}RC + C_{117}IC = CI \quad (D36)$$

Now, multiply Equation D30 by C_{115}/C_{107} and subtract the resultant from Equation D30.

Let,

$$\begin{aligned} C_{118} &= C_{114} - C_{115}C_{106}/C_{107} \\ C_{119} &= -(C_{116} - C_{115}C_{108}/C_{107})/C_{118} \\ C_{120} &= -(C_{117} - C_{115}C_{109}/C_{107})/C_{118} \\ C_{121} &= (CI - C_{115}CR/C_{107})/C_{118} \end{aligned} \quad (D37)$$

This gives

$$Z_{+\infty R} = C_{119}RC + C_{120}IC + C_{121} \quad (D38)$$

Substituting Equation D38 into Equation D36 and solving for $Z_{+\infty I}$ gives

$$Z_{+\infty I} = C_{122}RC + C_{123}IC + C_{124} \quad (D39)$$

where,

$$\begin{aligned} C_{122} &= (-C_{116} - C_{114}C_{119})/C_{115} \\ C_{123} &= (-C_{117} - C_{114}C_{120})/C_{115} \\ C_{124} &= (CI - C_{121}C_{114})/C_{115} \end{aligned} \quad (D40)$$

The relationships given for $Z_{+\infty}$ in Equations D38 and D39 are then substituted into Equations D22 and D23 to yield $A_{+\infty}$. This, then, produces a complete description of the flowfield for Region 3.

APPENDIX E
STEADY-STATE FLOW COEFFICIENTS FOR REGIONS 1, 2, AND 3

REGION 1

Energy

$$E_1 = \left[\frac{\bar{a}_1^2}{\gamma(\gamma - 1)} + \frac{\cos^2(\alpha_{ch} - \beta_1)}{2} \right]$$

$$E_2 = \frac{\bar{a}^{*2}(\gamma^2 - \gamma + 2)}{2\gamma(\gamma - 1)}$$

$$E_3 = \frac{-E_1 \bar{I} \bar{A}_1 \bar{u}_1}{C \bar{y}_i \bar{a}_1^2} + \frac{E_1 \bar{A}_1 \sin \alpha_{ch}}{C \bar{y}_i} - \frac{2 \bar{A}_1 \bar{I} \bar{u}_1}{2\gamma C \bar{y}_i} + \frac{\bar{A}_1 \bar{I}^2 \sin \alpha_{ch}}{C \bar{y}_i}$$

$$E_4 = \frac{-E_1 \bar{I} \bar{A}_1 (k + \bar{v}_1 C)}{C \bar{y}_i \bar{a}_1^2} + \frac{E_1 \bar{A}_1 \cos \alpha_{ch}}{\bar{y}_i} - \frac{2 \bar{A}_1 \bar{I} (k + \bar{v}_1 C)}{2\gamma C \bar{y}_i} + \frac{\bar{A}_1 \bar{I}^2 \cos \alpha_{ch}}{\bar{y}_i}$$

$$E_5 = \frac{-\bar{a}_1^2 \bar{V}_1 \bar{u}_1}{2\gamma(\gamma - 1) C \bar{y}_i \bar{a}_1^2} - \frac{\bar{\rho}_1 \bar{a}_1 \bar{V}_1 \bar{u}_1}{\gamma(\gamma - 1)} \left(\frac{\gamma - 1}{2 \bar{a}_1 C \bar{y}_i} \right) - \frac{\bar{U}_1^2 \bar{V}_1 \bar{u}_1}{4 C \bar{y}_i \bar{a}_1^2} + \frac{\bar{\rho}_1 \bar{U}_1 \bar{V}_1 \sin \alpha_{ch}}{2 C \bar{y}_i}$$

$$E_6 = \frac{-\bar{a}_1^2 \bar{V}_1 (k + \bar{v}_1 C)}{2\gamma(\gamma - 1) C \bar{y}_i \bar{a}_1^2} - \bar{\rho}_1 \bar{a}_1 \bar{V}_1 \left(\frac{\gamma - 1}{2 C \bar{y}_i \bar{a}_1} \right) \left(\frac{k + \bar{v}_1 C}{\gamma(\gamma - 1)} \right) + \frac{\bar{U}_1^2 \bar{V}_1 (k + \bar{v}_1 C)}{4 C \bar{y}_i \bar{a}_1^2} + \frac{\bar{\rho}_1 \bar{U}_1 \bar{V}_1 \cos \alpha_{ch}}{2 \bar{y}_i}$$

$$E_7 = \cos(C \bar{y}_i) - 1$$

$$E_8 = E_3 (E_7 B_{1R} - B_{1I} \sin(C \bar{y}_i)) + E_4 E_7 + E_5 k (B_{1I} E_7 + B_{1R} \sin(C \bar{y}_i)) + E_6 k \sin(C \bar{y}_i)$$

$$E_9 = -E_3 (B_{1I} E_7 + B_{1R} \sin(C \bar{y}_i)) - E_4 \sin(C \bar{y}_i) + E_5 k (B_{1R} E_7 - B_{1I} \sin(C \bar{y}_i)) + E_6 k E_7$$

$$E_{10} = E_3 [B_{1I} E_7 + B_{1R} \sin(C \bar{y}_i)] + E_4 \sin(C \bar{y}_i) - E_5 k [B_{1R} E_7 - B_{1I} \sin(C \bar{y}_i)] - E_6 k E_7$$

$$E_{11} = E_3 [B_{1R} E_7 - B_{1I} \sin(C \bar{y}_i)] + E_4 E_7 + E_5 k [B_{1I} E_7 + B_{1R} \sin(C \bar{y}_i)] + E_6 k \sin(C \bar{y}_i)$$

$$\mathbf{E}_{12} = \left(\frac{\bar{a}^2 \bar{\mathbf{V}}\mathbf{k}}{2\gamma(\gamma-1)} + \frac{\bar{\mathbf{U}}^2 \bar{\mathbf{V}}\mathbf{k}}{4} \right),$$

$$\mathbf{E}_{13} = \left(\frac{\bar{\rho} \bar{a}^2 \mathbf{k}}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{\mathbf{U}}^2 \mathbf{k}}{2} \right),$$

$$\mathbf{E}_{14} = \left(\frac{\bar{\rho} \bar{a} \bar{\mathbf{V}}\mathbf{k}}{\gamma(\gamma-1)} + \frac{\bar{\rho} \bar{\mathbf{U}} \bar{\mathbf{V}}\mathbf{k}}{2} \right),$$

$$\mathbf{E}_{15} = \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} - \frac{\mathbf{E}_{11}}{\mathbf{E}_{10}} \right)$$

$$\mathbf{E}_{16} = \mathbf{I} \left(\frac{\mathbf{E}_1}{\mathbf{E}_6} \right) \left(-\frac{\mathbf{E}_9}{\mathbf{E}_{15} \mathbf{E}_8} - 1 \right)$$

$$\mathbf{E}_{17} = \mathbf{I} \left(\frac{-\mathbf{E}_1}{\mathbf{E}_{10} \mathbf{E}_{15}} \right) \left(-\frac{\mathbf{E}_8}{\mathbf{E}_9} \right)$$

$$\mathbf{E}_{18} = \left(\frac{\mathbf{E}_2}{\mathbf{E}_8} \right) \left(\bar{a}^* - \frac{\bar{a}^* \mathbf{E}_9}{\mathbf{E}_{15} \mathbf{E}_8} \right) - \left(\frac{\mathbf{E}_{12} \mathbf{E}_9}{\mathbf{E}_8 \mathbf{E}_{10} \mathbf{E}_{15}} \right)$$

$$\mathbf{E}_{19} = -\frac{\mathbf{E}_{12}}{\mathbf{E}_8} + \left(\frac{\bar{a}^* \mathbf{E}_2}{\mathbf{E}_{10} \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right) + \left(\frac{\mathbf{E}_{12}}{\mathbf{E}_8 \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right)$$

$$\mathbf{E}_{20} = \bar{\rho}^* \bar{a}^* \left[\frac{\mathbf{E}_2}{\mathbf{E}_8} - \left(\frac{\mathbf{E}_2}{\mathbf{E}_8 \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right) \right]$$

$$\mathbf{E}_{21} = \bar{\rho}^* \bar{a}^* \left(\frac{\mathbf{E}_2}{\mathbf{E}_{10} \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right)$$

$$\mathbf{E}_{22} = 3\bar{\rho}^* \left(\frac{\mathbf{E}_2}{\mathbf{E}_8} \right) \left(1 - \frac{\mathbf{E}_9}{\mathbf{E}_8 \mathbf{E}_{15}} \right) - \left(\frac{\mathbf{E}_8 \mathbf{E}_{14}}{\mathbf{E}_8 \mathbf{E}_{10} \mathbf{E}_{15}} \right)$$

$$\mathbf{E}_{23} = -\frac{\mathbf{E}_{14}}{\mathbf{E}_8} + 3\bar{\rho}^* \left(\frac{\mathbf{E}_2 \mathbf{E}_9}{\mathbf{E}_8 \mathbf{E}_{10} \mathbf{E}_{15}} \right) + \left(\frac{\mathbf{E}_{14}}{\mathbf{E}_8 \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right)$$

$$\mathbf{E}_{24} = -\frac{\mathbf{E}_{13}}{\mathbf{E}_8} + \left(\frac{\mathbf{E}_{13}}{\mathbf{E}_8 \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right)$$

$$\mathbf{E}_{25} = -\left(\frac{\mathbf{E}_{13}}{\mathbf{E}_{10} \mathbf{E}_{15}} \right) \left(\frac{\mathbf{E}_9}{\mathbf{E}_8} \right)$$

$$\mathbf{E}_{26} = -\mathbf{I} \left(\frac{\mathbf{E}_1}{\mathbf{E}_8 \mathbf{E}_{15}} \right)$$

$$\mathbf{E}_{27} = \frac{\mathbf{E}_1}{\mathbf{E}_{10} \mathbf{E}_{15}}$$

$$\mathbf{E}_{28} = \frac{\bar{a}^* \mathbf{E}_2}{\mathbf{E}_8 \mathbf{E}_{15}} + \frac{\mathbf{E}_{12}}{\mathbf{E}_{10} \mathbf{E}_{15}}$$

$$\mathbf{E}_{29} = -\frac{\bar{a}^* \mathbf{E}_2}{\mathbf{E}_{10} \mathbf{E}_{15}} - \frac{\mathbf{E}_{12}}{\mathbf{E}_8 \mathbf{E}_{15}}$$

$$E_{30} = \frac{\bar{\rho}^* \bar{a}^* E_2}{E_8 E_{15}}$$

$$E_{31} = - \frac{\bar{\rho}^* \bar{a}^* E_2}{E_{10} E_{15}}$$

$$E_{32} = \frac{3\bar{\rho}^* E_2}{E_8 E_{15}} + \frac{E_{14}}{E_{10} E_{15}}$$

$$E_{33} = - \frac{3\bar{\rho}^* E_2}{E_{10} E_{15}} - \frac{E_{14}}{E_8 E_{15}}$$

$$E_{34} = - \frac{E_{13}}{E_8 E_{15}}$$

$$E_{35} = \frac{E_{13}}{E_{10} E_{15}}$$

Momentum

$$M_1 = - \frac{\bar{u}_i \bar{A}_i}{C\bar{y}_i} - \frac{\bar{u}_i \bar{A}_i I^2}{C\bar{y}_i \bar{a}_i^2} + \frac{2\bar{A}_i I}{C\bar{y}_i} \sin \alpha_{ch}$$

$$M_2 = - \frac{(k + \bar{v}_i C) \bar{A}_i}{C\bar{y}_i} - \frac{(k + \bar{v}_i C) \bar{A}_i I^2}{C\bar{y}_i \bar{a}_i^2} + \frac{2\bar{A}_i I}{\bar{y}_i} \cos \alpha_{ch}$$

$$M_3 = \frac{\bar{U}_i \bar{V}_i}{2C\bar{y}_i \bar{a}_i^2} - \frac{\bar{\rho}_i \bar{V}_i \sin \alpha_{ch}}{2C\bar{y}_i}$$

$$M_4 = \left(\frac{k + \bar{v}_i C}{C\bar{y}_i \bar{a}_i^2} \right) \left(\frac{\bar{U}_i \bar{V}_i}{2} \right) - \frac{\bar{\rho}_i \bar{V}_i \cos \alpha_{ch}}{2\bar{y}_i}$$

$$M_5 = -\bar{p}_i - I^2$$

$$M_6 = \bar{p}^* + \bar{\rho}^* \bar{a}^{*2}$$

$$M_7 = M_1 [B_{1R} E_7 - B_{1l} \sin (C\bar{y}_i)] + M_2 E_7 -$$

$$M_3 [kB_{1l} E_7 - kB_{1R} \sin (C\bar{y}_i)] - M_2 k \sin (C\bar{y}_i)$$

$$M_8 = -M_1 [B_{1l} E_7 + B_{1R} \sin (C\bar{y}_i)] - M_2 \sin (C\bar{y}_i) +$$

$$kM_3 [-B_{1R} E_7 + B_{1l} \sin (C\bar{y}_i)] - kM_4 E_7$$

$$M_9 = M_1 [B_{1l} E_7 + B_{1R} \sin (C\bar{y}_i)] + M_2 \sin (C\bar{y}_i) +$$

$$kM_3 [B_{1R} E_7 - B_{1l} \sin (C\bar{y}_i)] + M_4 E_7 k$$

$$M_{10} = M_1 [B_{1R} E_7 - B_{1l} \sin (C\bar{y}_i)] + M_2 E_7 -$$

$$kM_3 [B_{1l} E_7 + B_{1R} \sin (C\bar{y}_i)] - kM_4 \sin (C\bar{y}_i)$$

$$M_{11} = E_{16} M_7 + E_{26} M_8 - 2\bar{a}^{*2}$$

$$M_{12} = E_{19}M_7 + E_{29}M_8 + \frac{k\bar{U}_1\bar{V}_1}{2}$$

$$M_{13} = -E_{16}M_7 - E_{26}M_8 + M_5$$

$$M_{14} = -E_{17}M_7 - E_{27}M_8$$

$$M_{15} = -E_{20}M_7 - E_{30}M_8 + M_6$$

$$M_{16} = -E_{21}M_7 - E_{31}M_8$$

$$M_{17} = -E_{22}M_7 - E_{32}M_8 + 2\bar{\rho}^*\bar{a}^*$$

$$M_{18} = E_{23}M_7 - E_{33}M_8 - \frac{k\bar{\rho}_1\bar{V}_1}{2}$$

$$M_{19} = -E_{24}M_7 - E_{34}M_8 - k\bar{\rho}_1\bar{U}_1$$

$$M_{20} = -E_{25}M_7 - E_{35}M_8$$

$$M_{21} = E_{18}M_9 + E_{28}M_{10} - \frac{k\bar{U}_1\bar{V}_1}{2}$$

$$M_{22} = E_{19}M_9 + E_{29}M_{10} - 2\bar{a}^{*2}$$

$$M_{23} = -E_{16}M_9 - E_{26}M_{10} + M_5$$

$$M_{24} = -E_{17}M_9 - E_{27}M_{10} + M_6$$

$$M_{25} = -E_{20}M_9 - E_{30}M_{10}$$

$$M_{26} = -E_{21}M_9 - E_{31}M_{10}$$

$$M_{27} = -E_{22}M_9 - E_{32}M_{10} + \frac{k\bar{\rho}_1\bar{V}_1}{2}$$

$$M_{28} = -E_{23}M_9 - E_{33}M_{10} + 2\bar{\rho}^*\bar{a}^*$$

$$M_{29} = -E_{24}M_9 - E_{34}M_{10}$$

$$M_{30} = -E_{25}M_9 - E_{35}M_{10} + k\bar{\rho}_1\bar{U}_1$$

$$M_{31} = M_{11} - \frac{M_{12}M_{21}}{M_{22}}$$

$$M_{32} = \left(M_{13} - \frac{M_{12}M_{23}}{M_{22}} \right) \div M_{31}$$

$$M_{33} = \left(M_{14} - \frac{M_{12}M_{24}}{M_{22}} \right) \div M_{31}$$

$$M_{34} = \left(M_{15} - \frac{M_{12}M_{25}}{M_{22}} \right) \div M_{31}$$

$$M_{35} = \left(M_{16} - \frac{M_{12}M_{26}}{M_{22}} \right) \div M_{31}$$

$$M_{36} = \left(M_{17} - \frac{M_{12}M_{27}}{M_{22}} \right) \div M_{31}$$

$$M_{37} = \left(M_{18} - \frac{M_{12}M_{28}}{M_{22}} \right) \div M_{31}$$

$$M_{38} = \left(M_{19} - \frac{M_{12}M_{29}}{M_{22}} \right) \div M_{31}$$

$$M_{39} = \left(M_{20} - \frac{M_{12}M_{30}}{M_{22}} \right) \div M_{31}$$

$$M_{40} = \frac{M_{13} - M_{32}M_{11}}{M_{12}}$$

$$M_{41} = \frac{M_{14} - M_{33}M_{11}}{M_{12}}$$

$$M_{42} = \frac{M_{15} - M_{34}M_{11}}{M_{12}}$$

$$M_{43} = \frac{M_{16} - M_{35}M_{11}}{M_{12}}$$

$$M_{44} = \frac{M_{17} - M_{36}M_{11}}{M_{12}}$$

$$M_{45} = \frac{M_{18} - M_{37}M_{11}}{M_{12}}$$

$$M_{46} = \frac{M_{19} - M_{38}M_{11}}{M_{12}}$$

$$M_{47} = \frac{M_{20} - M_{39}M_{11}}{M_{12}}$$

Continuity

$$C_1 = - \frac{\bar{A}_i}{C\bar{y}_i} \sin \alpha_{ch} + \frac{\bar{u}_i \bar{A}_i I}{C\bar{y}_i \bar{a}_i^2}$$

$$C_2 = - \frac{\bar{A}_i}{C\bar{y}_i} \cos \alpha_{ch} + \left(\frac{k + \bar{v}_i C}{C\bar{y}_i \bar{a}_i^2} \right) \bar{A}_i I$$

$$C_3 = \frac{\bar{u}_i \bar{V}_1}{2C\bar{y}_i \bar{a}_i^2}$$

$$C_4 = \left(\frac{k + \bar{v}_i C}{2C\bar{y}_i \bar{a}_i^2} \right) \bar{V}_1$$

$$\begin{aligned}
C_5 &= C_1 B_{1R} E_7 - C_1 B_{11} \sin(C\bar{y}_i) + C_2 E_7 + \\
&\quad C_3 k B_{11} E_7 + C_3 k B_{1R} \sin(C\bar{y}_i) + C_4 k \sin(C\bar{y}_i) \\
C_6 &= C_1 B_{11} E_7 - C_1 B_{1R} \sin(C\bar{y}_i) - C_2 \sin(C\bar{y}_i) + \\
&\quad C_3 k B_{1R} E_7 - C_3 k B_{11} \sin(C\bar{y}_i) + C_4 k E_7 \\
C_7 &= C_5 E_{16} + C_6 E_{26} - I \\
C_8 &= C_5 E_{17} + C_6 E_{27} \\
C_9 &= C_5 E_{18} + C_6 E_{28} + \bar{a}^* \\
C_{10} &= C_5 E_{19} + C_6 E_{29} - k\bar{V}_1/2 \\
C_{11} &= C_5 E_{20} + C_6 E_{30} + \bar{\rho}^* \bar{a}^* \\
C_{12} &= C_5 E_{21} + C_6 E_{31} \\
C_{13} &= C_5 E_{22} + C_6 E_{32} + \bar{\rho}^* \\
C_{14} &= C_5 E_{23} + C_6 E_{33} \\
C_{15} &= C_5 E_{24} + C_6 E_{34} - k\bar{\rho}_1 \\
C_{16} &= C_5 E_{25} + C_6 E_{35} \\
C_{17} &= C_9 M_{32} + C_{10} M_{40} + C_7 \\
C_{18} &= C_9 M_{33} + C_{10} M_{41} + C_8 \\
C_{19} &= C_9 M_{34} + C_{10} M_{42} + C_{11} \\
C_{20} &= C_9 M_{35} + C_{10} M_{43} + C_{12} \\
C_{21} &= - (C_9 M_{36} + C_{10} M_{44} + C_{13}) \\
C_{22} &= - (C_9 M_{37} + C_{10} M_{45} + C_{14}) \\
C_{23} &= C_9 M_{38} + C_{10} M_{46} + C_{15} \\
C_{24} &= C_9 M_{39} + C_{10} M_{47} + C_{16} \\
C_{25} &= C_1 [B_{11} E_7 + B_{1R} \sin(C\bar{y}_i)] + C_2 \sin(C\bar{y}_i) - kC_3 [B_{1R} E_7 - \\
&\quad B_{11} \sin(C\bar{y}_i)] - kC_4 E_7 \\
C_{26} &= C_1 [B_{1R} E_7 - B_{11} \sin(C\bar{y}_i)] + C_2 E_7 + kC_3 [B_{11} E_7 + \\
&\quad B_{1R} \sin(C\bar{y}_i)] + kC_4 \sin(C\bar{y}_i) \\
C_{27} &= C_{26} E_{26} + C_{25} E_{16} \\
C_{28} &= C_{26} E_{27} + C_{25} E_{17} - I
\end{aligned}$$

$$C_{29} = C_{26}E_{28} + C_{25}E_{18} + \frac{k\bar{V}_1}{2}$$

$$C_{30} = C_{26}E_{29} + C_{25}E_{19} + \bar{a}^*$$

$$C_{31} = C_{26}E_{30} + C_{25}E_{20}$$

$$C_{32} = C_{26}E_{31} + C_{25}E_{21} + \bar{\rho}^*\bar{a}^*$$

$$C_{33} = C_{26}E_{32} + C_{25}E_{22}$$

$$C_{34} = C_{26}E_{33} + C_{25}E_{23} + \bar{\rho}^*$$

$$C_{35} = C_{26}E_{34} + C_{25}E_{24}$$

$$C_{36} = C_{26}E_{35} + C_{25}E_{25} + k\bar{\rho}_1$$

$$C_{37} = C_{29}M_{32} + C_{30}M_{40} + C_{27}$$

$$C_{38} = C_{29}M_{33} + C_{30}M_{41} + C_{28}$$

$$C_{39} = C_{29}M_{34} + C_{30}M_{42} + C_{31}$$

$$C_{40} = C_{29}M_{35} + C_{30}M_{43} + C_{32}$$

$$C_{41} = - (C_{29}M_{36} + C_{30}M_{44} + C_{33})$$

$$C_{42} = - (C_{29}M_{37} + C_{30}M_{45} + C_{34})$$

$$C_{43} = C_{29}M_{38} + C_{30}M_{46} + C_{35}$$

$$C_{44} = C_{29}M_{39} + C_{30}M_{47} + C_{36}$$

$$C_{45} = C_{41} - \frac{C_{21}C_{42}}{C_{22}}$$

$$C_{46} = \frac{1}{C_{45}} \left(C_{37} - \frac{C_{17}C_{42}}{C_{22}} \right)$$

$$C_{47} = \frac{1}{C_{45}} \left(C_{38} - \frac{C_{18}C_{42}}{C_{22}} \right)$$

$$C_{48} = \frac{1}{C_{45}} \left(C_{39} - \frac{C_{19}C_{42}}{C_{22}} \right)$$

$$C_{49} = \frac{1}{C_{45}} \left(C_{40} - \frac{C_{20}C_{42}}{C_{22}} \right)$$

$$C_{50} = \frac{1}{C_{45}} \left(C_{43} - \frac{C_{23}C_{42}}{C_{22}} \right)$$

$$C_{51} = \frac{1}{C_{45}} \left(C_{44} - \frac{C_{24}C_{42}}{C_{22}} \right)$$

$$C_{52} = (C_{37} - C_{41}C_{46}) \div C_{42}$$

$$C_{53} = (C_{38} - C_{41}C_{47}) \div C_{42}$$

$$C_{54} = (C_{39} - C_{41}C_{48}) \div C_{42}$$

$$C_{55} = (C_{40} - C_{41}C_{49}) \div C_{42}$$

$$C_{56} = (C_{43} - C_{41}C_{50}) \div C_{42}$$

$$C_{57} = (C_{44} - C_{41}C_{51}) \div C_{42}$$

REGION 2

Momentum

$$M_{48} = (-\bar{a}_{US}^2 - \bar{U}_{US}^2) \bar{A}_S$$

$$M_{49} = \frac{k\bar{U}_2\bar{V}_2}{2}$$

$$M_{50} = \bar{p}_{US} + \bar{\rho}_{US}\bar{U}_{US}^2$$

$$M_{50A} = \bar{p}^* + \bar{\rho}^*\bar{a}^{*2}$$

$$M_{51} = M_{49} + \frac{M_{48}^2}{M_{49}}$$

$$M_{52} = \frac{M_{50}}{M_{51}}$$

$$M_{53} = \frac{M_{48}M_{50}}{M_{49}M_{51}}$$

$$M_{54} = \frac{-2}{M_{51}}$$

$$M_{55} = \frac{-2M_{48}}{M_{49}M_{51}}$$

$$M_{56} = \frac{M_{50A}}{M_{51}}$$

$$M_{57} = \frac{M_{48}M_{50A}}{M_{49}M_{51}}$$

$$M_{58} = \left[2\bar{w}_{US} + \frac{k\bar{\rho}_2\bar{V}_2}{2} \left(\frac{M_{48}}{M_{49}} \right) \right] \div M_{51}$$

$$M_{59} = \frac{1}{M_{51}} \left(\frac{2\bar{w}_{US}M_{48}}{M_{49}} - \frac{k\bar{\rho}_2\bar{V}_2}{2} \right)$$

$$\begin{aligned}
M_{60} &= \frac{1}{M_{51}} \left(-\frac{\bar{a}^* 2\bar{\rho}^* M_{48}}{M_{49}} - \frac{k\bar{\rho}_2 \bar{V}_2}{2} \right) \\
M_{61} &= \frac{1}{M_{51}} \left(-2\bar{\rho}^* \bar{a}^* + \frac{k\bar{\rho}_2 \bar{V}_2 M_{48}}{2M_{49}} \right) \\
M_{62} &= \frac{1}{M_{51}} \left(\frac{-k\bar{U}_2 \bar{V}_2}{2} \right) \\
M_{63} &= \frac{1}{M_{51}} \left(\frac{k\bar{U}_2 \bar{V}_2 M_{48}}{2M_{49}} \right) \\
M_{64} &= \frac{-k(\bar{\rho}\bar{U})_2}{M_{51}} \\
M_{65} &= k(\bar{\rho}\bar{U})_2 \left(\frac{M_{48}}{M_{49} M_{51}} \right) \\
M_{66} &= \frac{M_{50} - M_{49} M_{52}}{M_{48}} \\
M_{67} &= -\frac{M_{49} M_{53}}{M_{48}} \\
M_{68} &= \frac{-2 - M_{49} M_{54}}{M_{48}} \\
M_{69} &= \frac{-M_{49} M_{55}}{M_{48}} \\
M_{70} &= (M_{50A} - M_{49} M_{56}) / M_{48} \\
M_{71} &= -M_{49} M_{57} / M_{48} \\
M_{72} &= (2\bar{w}_{US} - M_{49} M_{58}) / M_{48} \\
M_{73} &= (-\frac{1}{2} k\bar{\rho}_2 \bar{V}_2 - M_{49} M_{59}) / M_{48} \\
M_{74} &= (-\frac{1}{2} k\bar{\rho}_2 \bar{V}_2 - M_{49} M_{60}) / M_{48} \\
M_{75} &= (-2\bar{\rho}^* \bar{a}^* - M_{49} M_{61}) / M_{48} \\
M_{76} &= (-\frac{1}{2} k\bar{U}_2 \bar{V}_2 - M_{49} M_{62}) / M_{48} \\
M_{77} &= -M_{49} M_{63} / M_{48} \\
M_{78} &= (-k(\bar{\rho}\bar{U})_2 - M_{49} M_{64}) / M_{48} \\
M_{79} &= -M_{49} M_{65} / M_{48} \\
M_{80} &= \bar{\rho}_{US} \bar{A}_S + \bar{U}_{US} \bar{A}_S M_{72} - \frac{k\bar{V}_2}{2} M_{58} \\
M_{81} &= \bar{U}_{US} \bar{A}_S M_{73} - \frac{k\bar{V}_2}{2} M_{59}
\end{aligned}$$

$$\begin{aligned}
M_{62} &= -\bar{U}_{US}\bar{A}_S M_{66} - \bar{\rho}_{US}\bar{U}_{US} + \frac{k\bar{V}_2}{2} M_{52} \\
M_{63} &= -\bar{U}_{US}\bar{A}_S M_{67} + \frac{k\bar{V}_2}{2} M_{53} \\
M_{64} &= -\bar{U}_{US}\bar{A}_S M_{71} + \bar{a}^* + \frac{k\bar{V}_2}{2} M_{63} \\
M_{65} &= -\bar{U}_{US}\bar{A}_S M_{76} + \frac{k\bar{V}_2}{2} M_{62} \\
M_{66} &= -\bar{U}_{US}\bar{A}_S M_{68} + \frac{k\bar{V}_2}{2} M_{54} \\
M_{67} &= -\bar{U}_{US}\bar{A}_S M_{69} + \frac{k\bar{V}_2}{2} M_{55} \\
M_{68} &= -\bar{U}_{US}\bar{A}_S M_{70} + \bar{\rho}^*\bar{a}^* + \frac{k\bar{V}_2}{2} M_{56} \\
M_{69} &= -\bar{U}_{US}\bar{A}_S M_{71} + \frac{k\bar{V}_2}{2} M_{57} \\
M_{90} &= -\bar{U}_{US}\bar{A}_S M_{74} + \frac{k\bar{V}_2}{2} M_{60} \\
M_{91} &= -\bar{U}_{US}\bar{A}_S M_{75} + \bar{\rho}^* + \frac{k\bar{V}_2}{2} M_{61} \\
M_{92} &= -\bar{U}_{US}\bar{A}_S M_{78} + \frac{k\bar{V}_2}{2} M_{64} + k\bar{\rho}_1 \\
M_{93} &= -\bar{U}_{US}\bar{A}_S M_{79} + \frac{k\bar{V}_2}{2} M_{60}
\end{aligned}$$

Continuity

$$\begin{aligned}
C_{58} &= \bar{\rho}_{US}\bar{A}_S + \bar{U}_{US}\bar{A}_S M_{59} + \frac{k\bar{V}_2}{2} M_{71} \\
C_{59} &= \bar{U}_{US}\bar{A}_S M_{58} + \frac{k\bar{V}_2}{2} M_{72} \\
C_{60} &= -\bar{U}_{US}\bar{A}_S M_{92} - \frac{k\bar{V}_2}{2} M_{65} \\
C_{61} &= -\bar{\rho}_{US}\bar{U}_{US} - \bar{U}_{US}\bar{A}_S M_{63} - \frac{k\bar{V}_2}{2} M_{67} \\
C_{62} &= -\bar{U}_{US}\bar{A}_S M_{63} - \frac{k\bar{V}_2}{2} (1 + M_{71}) \\
C_{63} &= \bar{a}^* - \bar{U}_{US}\bar{A}_S M_{62} - \frac{k\bar{V}_2}{2} M_{76}
\end{aligned}$$

$$\begin{aligned}
C_{64} &= -\bar{U}_{US}\bar{A}_{US}M_{61} - \frac{k\bar{V}_2}{2} M_{75} \\
C_{65} &= \rho^* - \bar{U}_{US}\bar{A}_S M_{60} - \frac{k\bar{V}_2}{2} M_{74} \\
C_{66} &= -\bar{U}_{US}\bar{A}_S M_{54} - \frac{k\bar{V}_2}{2} M_{68} \\
C_{67} &= -\bar{U}_{US}\bar{A}_S M_{55} - \frac{k\bar{V}_2}{2} M_{69} \\
C_{68} &= -\bar{U}_{US}\bar{A}_S M_{56} - \frac{k\bar{V}_2}{2} M_{70} \\
C_{69} &= \bar{\rho}^* \bar{a}^* - \bar{U}_{US}\bar{A}_S M_{57} - \frac{k\bar{V}_2}{2} M_{71} \\
C_{70} &= -\bar{U}_{US}\bar{A}_S M_{65} - \frac{k\bar{V}_2}{2} (2\bar{\rho}_2 + M_{79}) \\
C_{71} &= -\bar{U}_{US}\bar{A}_S M_{64} - \frac{k\bar{V}_2}{2} M_{78} \\
C_{72} &= C_{59} - \frac{C_{58}M_{80}}{M_{81}} \\
C_{73} &= \frac{1}{C_{72}} \left(C_{60} - \frac{C_{58}M_{82}}{M_{81}} \right) \\
C_{74} &= \frac{1}{C_{72}} \left(C_{61} - \frac{C_{58}M_{81}}{M_{81}} \right) \\
C_{75} &= \frac{1}{C_{72}} \left(C_{62} - \frac{C_{58}M_{84}}{M_{81}} \right) \\
C_{76} &= \frac{1}{C_{72}} \left(C_{63} - \frac{C_{58}M_{85}}{M_{81}} \right) \\
C_{77} &= \frac{1}{C_{72}} \left(C_{64} - \frac{C_{58}M_{81}}{M_{81}} \right) \\
C_{78} &= \frac{1}{C_{72}} \left(C_{65} - \frac{C_{58}M_{90}}{M_{81}} \right) \\
C_{79} &= \frac{1}{C_{72}} \left(C_{66} - \frac{C_{58}M_{86}}{M_{81}} \right) \\
C_{80} &= \frac{1}{C_{72}} \left(C_{67} - \frac{C_{58}M_{87}}{M_{81}} \right) \\
C_{81} &= \frac{1}{C_{72}} \left(C_{68} - \frac{C_{58}M_{88}}{M_{81}} \right) \\
C_{82} &= \frac{1}{C_{72}} \left(C_{69} - \frac{C_{58}M_{89}}{M_{81}} \right)
\end{aligned}$$

$$C_{83} = \frac{1}{C_{72}} \left(C_{70} - \frac{C_{58}M_{93}}{M_{81}} \right)$$

$$C_{84} = \frac{1}{C_{72}} \left(C_{71} - \frac{C_{58}M_{92}}{M_{81}} \right)$$

$$C_{85} = \frac{1}{C_{58}} (C_{60} - C_{59}C_{73})$$

$$C_{86} = \frac{1}{C_{58}} (C_{61} - C_{59}C_{74})$$

$$C_{87} = \frac{1}{C_{58}} (C_{62} - C_{59}C_{75})$$

$$C_{88} = \frac{1}{C_{58}} (C_{63} - C_{59}C_{76})$$

$$C_{89} = \frac{1}{C_{58}} (C_{64} - C_{59}C_{77})$$

$$C_{90} = \frac{1}{C_{58}} (C_{65} - C_{59}C_{78})$$

$$C_{91} = \frac{1}{C_{58}} (C_{66} - C_{59}C_{79})$$

$$C_{92} = \frac{1}{C_{58}} (C_{67} - C_{59}C_{80})$$

$$C_{93} = \frac{1}{C_{58}} (C_{68} - C_{59}C_{81})$$

$$C_{94} = \frac{1}{C_{58}} (C_{69} - C_{59}C_{82})$$

$$C_{95} = \frac{1}{C_{58}} (C_{70} - C_{59}C_{83})$$

$$C_{96} = \frac{1}{C_{58}} (C_{71} - C_{59}C_{84})$$

REGION 3

Momentum

$$M_{94} = \frac{\bar{u}_E \bar{A}_E}{C \bar{y}_E} - \frac{\bar{w}_E \sin \alpha_{ch}}{C y_E} +$$

$$[U_{2REL} \cos (\alpha_{ch} - \beta_2)]^2 \left[\frac{\bar{u}_E \bar{A}_E}{C \bar{y}_E \bar{a}_E^2} \bar{p}_E \right] -$$

$$\left[\left(\frac{\bar{p}_E \bar{U}_{2REL} \bar{A}_E}{C \bar{y}_E} \right) (\sin \alpha_{ch}) \right] \cos (\alpha_{ch} - \beta_2)$$

$$\begin{aligned}
M_{95} &= \bar{\rho}_E \left(\frac{k + \bar{v}_E C}{C \bar{y}_E} \right) \bar{A}_E - \frac{\bar{w}_E \cos \alpha_{ch}}{\bar{y}_E} + \\
& [U_{2REL} \cos (\alpha_{ch} - \beta_2)]^2 \left[\left(\frac{k + \bar{v}_E C}{C \bar{y}_E \bar{a}_E^2} \right) \bar{A}_E \bar{\rho}_E \right] - \\
& \bar{\rho}_E \bar{U}_{2REL} \bar{A}_E \cos \alpha_{ch} \cos (\alpha_{ch} - \beta_2) \\
M_{96} &= \bar{\rho}_E \frac{\bar{u}_E \bar{U}_3 \bar{V}_3}{2 C \bar{y}_E \bar{a}_E^2} - \frac{\rho_3 \bar{V}_3 \sin \alpha_{ch}}{2 C \bar{y}_E} \\
M_{97} &= \bar{\rho}_E \frac{(k + \bar{v}_E C) \bar{U}_3 \bar{V}_3}{2 C \bar{y}_E \bar{a}_E^2} - \frac{\rho_3 \bar{V}_3 \cos \alpha_{ch}}{2 \bar{y}_E} \\
M_{98} &= \left[\frac{\bar{w}_E \sin \alpha_{ch}}{(R^2 + C^2) \bar{y}_E} - \frac{R \bar{w}_E \cos \alpha_{ch}}{C (R^2 + C^2) \bar{y}_E} \right] + \\
& \left[\frac{\bar{\rho}_E U_{2REL} \bar{A}_E \cos (\alpha_{ch} - \beta_2)}{(R^2 + C^2) \bar{y}_E} \right] \left[\sin \alpha_{ch} - \frac{R \cos \alpha_{ch}}{C} \right] \\
M_{99} &= \frac{\bar{\rho}_3 \bar{V}_3}{2 (R^2 + C^2) \bar{y}_E} \left(\sin \alpha_{ch} - \frac{R \cos \alpha_{ch}}{C} \right) \\
M_{100} &= \cos (C \bar{y}_E) - 1 \\
M_{101} &= M_{94} (M_{100} B_{2R} - B_{2l} \sin (C \bar{y}_E)) + M_{95} M_{100} - \\
& k M_{96} (B_{2l} M_{100} + B_{2R} \sin (C \bar{y}_E)) - k M_{97} \sin (C \bar{y}_E) \\
M_{102} &= M_{94} (M_{100} B_{2l} + B_{2R} \sin (C \bar{y}_E)) - M_{95} \sin (C \bar{y}_E) - \\
& k M_{96} (M_{100} B_{2R} - B_{2l} \sin (C \bar{y}_E)) - k M_{97} M_{100} \\
M_{103} &= -k M_{99} \sin (C \bar{y}_E) + M_{98} M_{100} \\
M_{104} &= -k M_{99} M_{100} - M_{98} \sin (C \bar{y}_E) \\
M_{105} &= M_{94} [B_{2l} M_{100} + B_{2R} \sin (C \bar{y}_E)] + M_{95} \sin (C \bar{y}_E) + \\
& k M_{96} [B_{2R} M_{100} - B_{2l} \sin (C \bar{y}_E)] + k M_{97} M_{100} \\
M_{106} &= M_{94} [B_{2R} M_{100} - B_{2l} \sin (C \bar{y}_E)] + M_{95} M_{100} - \\
& k M_{96} [B_{2l} M_{100} + B_{2R} \sin (C \bar{y}_E)] - k M_{97} \sin (C \bar{y}_E) \\
M_{107} &= M_{98} \sin (C \bar{y}_E) + k M_{99} M_{100} \\
M_{108} &= M_{98} M_{100} - k M_{99} \sin (C \bar{y}_E) \\
M_{109} &= M_{105} - M_{101} M_{106} / M_{102} \\
M_{110} &= (M_{107} - M_{103} M_{106} / M_{102}) / M_{109} \\
M_{111} &= (M_{108} - M_{104} M_{106} / M_{102}) / M_{109} \\
M_{112} &= 1 / M_{109}
\end{aligned}$$

$$M_{113} = -M_{106}/M_{102}M_{109}$$

$$M_{114} = (M_{107} - M_{105}M_{110})/M_{106}$$

$$M_{115} = (M_{108} - M_{105}M_{109})/M_{106}$$

$$M_{116} = M_{105} M_{113}/M_{106}$$

$$M_{117} = (1 - M_{105}M_{112})/M_{106}$$

Continuity

$$C_{97} = - \frac{\bar{\rho}_E \bar{u}_E \bar{U}_{2REL} \bar{A}_E \cos(\alpha_{ch} - \beta_2)}{C \bar{y}_E \bar{a}_E^2} + \frac{\bar{\rho}_E \bar{A}_E \sin \alpha_{ch}}{C \bar{y}_E}$$

$$C_{98} = - \frac{\bar{\rho}_E (k + \bar{v}_E C) \bar{U}_{2REL} \bar{A}_E \cos(\alpha_{ch} - \beta_2)}{C \bar{y}_E \bar{a}_E^2} + \frac{\bar{\rho}_E \bar{A}_E \cos \alpha_{ch}}{\bar{y}_E}$$

$$C_{99} = \frac{\bar{\rho}_E \bar{A}_E}{(R^2 + C^2) \bar{y}_E} \left(\sin \alpha_{ch} - \frac{R}{C} \cos \alpha_{ch} \right)$$

$$C_{100} = \frac{\bar{u}_E \bar{V}_3 \bar{\rho}_E}{2 C \bar{y}_E \bar{a}_E^2}$$

$$C_{101} = \frac{(k + \bar{v}_E C) \bar{V}_3 \bar{\rho}_E}{2 C \bar{y}_E \bar{a}_E^2}$$

$$C_{102} = C_{47} B_{2R} M_{100} - C_{97} B_{21} \sin(C \bar{y}_E) + C_{98} M_{100} +$$

$$k C_{100} (B_{21} M_{100} + B_{2R} \sin(C \bar{y}_E)) + k C_{101} \sin(C \bar{y}_E)$$

$$C_{103} = -C_{97} (B_{21} M_{100} + B_{2R} \sin(C \bar{y}_E)) - C_{98} \sin(C \bar{y}_E) +$$

$$k C_{100} (B_{2R} M_{100} - B_{21} \sin(C \bar{y}_E)) + k C_{101} M_{100}$$

$$C_{104} = C_{99} M_{100}$$

$$C_{105} = -C_{99} \sin(C \bar{y}_E)$$

$$C_{106} = C_{102} M_{110} + C_{103} M_{114} + C_{104}$$

$$C_{107} = C_{102} M_{111} + C_{103} M_{115} + C_{105}$$

$$C_{108} = C_{102} M_{113} + C_{103} M_{116}$$

$$C_{109} = C_{102} M_{112} + C_{103} M_{117}$$

$$C_{110} = C_{97} [B_{21} M_{100} + B_{2R} \sin(C \bar{y}_E)] + C_{98} \sin(C \bar{y}_E) -$$

$$k C_{100} [B_{2R} M_{100} - B_{21} \sin(C \bar{y}_E)] - k C_{101} M_{100}$$

$$C_{111} = C_{97} [B_{2R} M_{100} - B_{21} \sin(C \bar{y}_E)] + C_{98} M_{100} +$$

$$k C_{100} [B_{21} M_{100} + B_{2R} \sin(C \bar{y}_E)] + k C_{101} \sin(C \bar{y}_E)$$

$$\begin{aligned}
C_{112} &= C_{99} \sin (C\bar{y}_E) \\
C_{113} &= C_{99}M_{100} \\
C_{114} &= C_{110}M_{110} + C_{111}M_{114} + C_{112} \\
C_{115} &= C_{110}M_{111} + C_{111}M_{115} + C_{113} \\
C_{116} &= C_{110}M_{113} + C_{111}M_{116} \\
C_{117} &= C_{110}M_{112} + C_{111}M_{117} \\
C_{118} &= C_{114} - C_{115}C_{106}/C_{107} \\
C_{119} &= -(C_{116} - C_{115}C_{108}/C_{107})/C_{118} \\
C_{120} &= -(C_{117} - C_{115}C_{109}/C_{107})/C_{118} \\
C_{121} &= (CI - C_{115}CR/C_{107})/C_{118} \\
C_{122} &= (-C_{116} - C_{114}C_{119})/C_{115} \\
C_{123} &= (-C_{117} - C_{114}C_{120})/C_{115} \\
C_{124} &= (CI - C_{121}C_{114})/C_{115}
\end{aligned}$$

S Coefficients

$$\begin{aligned}
S_1 &= \frac{2\gamma}{\gamma + 1} \bar{p}_{us} \\
S_2 &= \frac{2\gamma}{\gamma + 1} \bar{p}_{us}\bar{M}_{us} - \bar{p}_{ds} - \frac{(\gamma - 1)}{\gamma + 1} \bar{p}_{us} \\
S_3 &= \frac{2\gamma}{\gamma + 1} \bar{a}_{us}\bar{M}_{us} - \bar{U}_e - \frac{(\gamma - 1)}{\gamma + 1} \bar{a}_{us} \\
S_4 &= \frac{2\gamma}{\gamma + 1} \bar{a}_{us}\bar{p}_{us} \\
S_5 &= (\gamma - 1) \bar{M}_{us}^2 + 2 \\
S_6 &= (2 - 2\gamma) \bar{p}_{ds} + (2\gamma + 2) \bar{p}_{us} \\
S_7 &= (\gamma + 1) \bar{M}_{us}^2 \\
S_8 &= S_6/(\bar{a}_{us}S_5) \\
S_9 &= S_7/S_5
\end{aligned}$$

Miscellaneous

$$RC = -\bar{p}'_{daR} \bar{A}_S - \bar{p}_{da} \bar{A}'_{SR} + [\bar{p}_E + \rho_E (U_{2REL} \cos(\alpha_{ch} - \beta_2))^2] \bar{A}'_{ER} -$$

$$\bar{U}'_{daR} \bar{w}_{da} - \bar{U}_{da} \bar{w}'_{daR} - \frac{k \bar{\rho}'_{daR} \bar{U}_3 \bar{V}_3}{2} - k \bar{\rho}_3 \bar{U}_3 \bar{V}'_{3I} - \frac{k \bar{\rho}_3 \bar{U}'_{daI} \bar{V}_3}{2}$$

with

$$\bar{w}'_{daR} = (\bar{\rho}'_{daR} \bar{U}_{da} \bar{A}_S + \bar{\rho}_{da} \bar{U}'_{daR} \bar{A}_S + \bar{\rho}_{da} \bar{U}_{da} \bar{A}'_{SR})$$

$$IC = -\bar{p}_{daI} \bar{A}_S - \bar{p}_{da} \bar{A}'_{SI} + [\bar{p}_E + \bar{\rho}_E (\bar{U}_{2REL} \cos(\alpha_{ch} - \beta_2))^2] \bar{A}'_{EI} -$$

$$\bar{U}'_{daI} \bar{w}_{da} - \bar{U}_{da} \bar{w}'_{daI} + \frac{k \bar{\rho}'_{daR} \bar{U}_3 \bar{V}_3}{2} + k \bar{\rho}_3 \bar{U}_3 \bar{V}'_{3R} + \frac{k \bar{\rho}_3 \bar{U}'_{daR} \bar{V}_3}{2}$$

with

$$\bar{w}'_{daI} = (\bar{\rho}'_{daI} \bar{U}_{da} \bar{A}_S + \bar{\rho}_{da} \bar{U}'_{daI} \bar{A}_S + \bar{\rho}_{da} \bar{U}_{da} \bar{A}'_{SI})$$

$$CR = -\bar{\rho}_E \bar{A}'_{ER} \bar{U}_E + \bar{w}'_{daR} + \bar{\rho}_3 k \bar{V}'_{3I} + 1/2 k \bar{V}'_{3I} \bar{\rho}'_{daI}$$

$$CI = -\bar{\rho}_E \bar{U}_{2REL} \cos(\alpha_{ch} - \beta_2) \bar{A}'_{EI} + \bar{w}'_{daI} - \bar{\rho}_3 k \bar{V}'_{3R} - 1/2 k \bar{V}_3 \bar{\rho}'_{daR}$$

APPENDIX F
CALCULATION OF MEAN FLOW AERODYNAMICS

Pressure jump across the shock (steady state)

$$\frac{\bar{\rho}_{ds}}{\bar{\rho}_{us}} = 1 + \frac{2\gamma}{\gamma+1} (\bar{M}_{us}^2 - 1) \quad (F1)$$

Steady-state pressure as a function of chordwise location

$$\bar{p}(x) = \bar{p}_{TOT} \left[1 + \frac{\gamma-1}{2} (\bar{M}_{(x)}^2) \right]^{\left(\frac{-\gamma}{\gamma-1} \right)} \quad (F2)$$

Steady-state density as a function of chordwise location

$$\bar{\rho}(x) = \bar{\rho}_{TOT} \left[1 + \frac{\gamma-1}{2} (\bar{M}_{(x)}^2) \right]^{\left(\frac{-1}{\gamma-1} \right)} \quad (F3)$$

where,

$$\rho_{TOT} = \frac{P_{TOT}}{RT_{TOT}}$$

Steady-state temperature as a function of chordwise position

$$\bar{T}(x) = \bar{T}_{TOT} \left[1 + \frac{\gamma-1}{2} (\bar{M}_{(x)}^2) \right]^{-1} \quad (F4)$$

Local steady-state speed of sound as a function of chordwise position.

$$\bar{a}(x) = \left[\frac{\gamma RT_{TOT}}{1 + \frac{\gamma-1}{2} (\bar{M}_{(x)}^2)} \right]^{1/2} \quad (F5)$$

Total pressure, temperature, and density downstream of the shock

$$p_{TOT_{ds}} = p_{TOT_{us}} \left[1 + \frac{2\gamma}{\gamma+1} (\bar{M}_{us}^2 - 1) \right]^{\left(\frac{-1}{\gamma-1} \right)}$$

$$\left[\frac{(\gamma+1) \bar{M}_{us}^2}{(\gamma-1) \bar{M}_{us}^2 + 2} \right]^{\left(\frac{\gamma}{\gamma-1} \right)} \quad (F6)$$

$$T_{TOT_{ds}} = T_{TOT_{us}}$$

Mach number across the shock

$$\bar{M}_{ds}^2 = \frac{\bar{M}_{ds}^2 (\gamma - 1) + 2}{2\gamma \bar{M}_{ds}^2 - (\gamma - 1)} \quad (\text{F7})$$

APPENDIX G

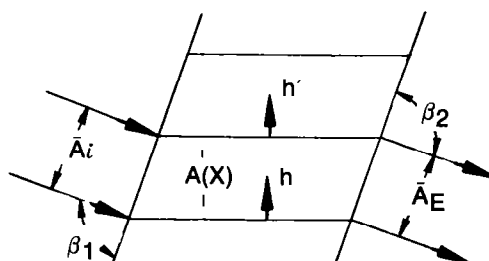
DEFINITION OF AREA PERTURBATIONS

Bending Mode Perturbations

Figure G1 describes the inlet area to be considered. The blades will be considered to vibrate harmonically with a constant interblade phase lag σ . Assuming this, then the deflections can be represented as follows:

$$h = \bar{h} e^{ikt} \quad (\text{reference blade})$$

$$h' = \bar{h} e^{ikt} e^{i\sigma} \quad (\text{blade adjacent to reference blade})$$



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Figure G1. Bending Mode Area Perturbations

Thus, the area perturbation can be represented as shown below:

$$\begin{aligned} A'_i &= (h' - h) \frac{\sin \beta_1}{\sin \alpha_{ch}} \\ &= (\bar{h}^{ikt} e^{i\sigma} - \bar{h} e^{ikt}) \frac{\sin \beta_1}{\sin \alpha_{ch}} \end{aligned} \quad (G1)$$

where all quantities are nondimensional. Simplifying this expression yields

$$\bar{A}'_i = \bar{h} e^{ikt} (e^{i\sigma} - 1) \frac{\sin \beta_1}{\sin \alpha_{ch}} \quad (G2)$$

Dividing through by e^{ikt} gives

$$\bar{A}'_i = \bar{h} (e^{i\sigma} - 1) \frac{\sin \beta_1}{\sin \alpha_{ch}} \quad (G3)$$

Separating the expression into its real and imaginary components produces

$$\bar{A}'_{iR} = \bar{h} (\cos \sigma - 1) \frac{\sin \beta_1}{\sin \alpha_{ch}} \quad (G4)$$

$$\bar{A}'_{iI} = \bar{h} \sin \sigma \left(\frac{\sin \beta_1}{\sin \alpha_{ch}} \right) \quad (G5)$$

Similarly, the other area perturbations for the bending mode will be

$$\bar{A}'_{*R} = \bar{h} (\cos \sigma - 1) \quad (G6)$$

$$\bar{A}'_{*I} = \bar{h} \sin \sigma \quad (G7)$$

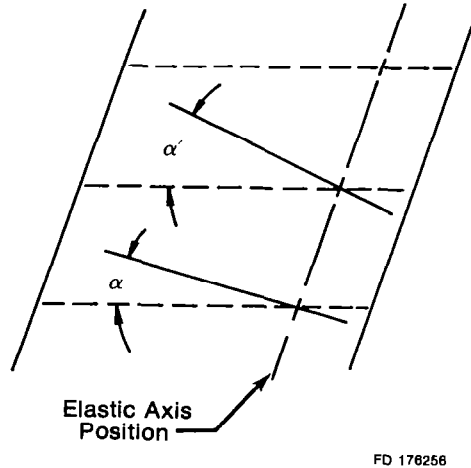
$$\bar{A}'_{SR} = \bar{h} (\cos \sigma - 1) \quad (G8)$$

$$\bar{A}'_{SI} = \bar{h} \sin \sigma \quad (G9)$$

$$\bar{A}'_{ER} = \bar{h} (\cos \sigma - 1) \frac{\sin \beta_2}{\sin \alpha_{ch}} \quad (G10)$$

$$\bar{A}'_{EI} = \bar{h} \sin \sigma \left(-\frac{\sin \beta_2}{\sin \alpha_{ch}} \right) \quad (G11)$$

Torsional Mode Area Perturbations



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Figure G2. Torsional Mode Area Perturbation

In the torsional mode area perturbations, the airfoils will be assumed to be undergoing rigid body torsional deflections about an arbitrary elastic axis position a . Consider first a cascade of flat plate airfoils oscillating out of phase, as shown in Figure G2. The area perturbation can be represented by

$$A'(x) = c \left(Z - \frac{x}{c} \right) (\alpha' - \alpha) \quad (G12)$$

where,

$$Z = \frac{b(1 + a)}{c}$$

a = Elastic axis position referenced to midchord and nondimensionalized by b .

At the blade leading edge

$$A'_i = Zc (\alpha' - \alpha) \left(\frac{\sin \beta_1}{\sin \alpha_{ch}} \right) \quad (G13)$$

Next, assume

$$\begin{aligned} \alpha &= \bar{\alpha} e^{ikt} \\ \alpha' &= \bar{\alpha} e^{ikt} e^{i\sigma} \end{aligned} \quad (G14)$$

Thus,

$$\alpha' - \alpha = \bar{\alpha} e^{ikt} (e^{i\sigma} - 1)$$

Dividing this by e^{ikt} and substituting the resultant into Equation G13 gives

$$\bar{A}'_i = cZ\bar{\alpha}(e^{i\sigma} - 1) \left(\frac{\sin \beta_1}{\sin \alpha_{ch}} \right)$$

Separating this expression into its real and imaginary components yields

$$\bar{A}'_{ir} = cZ\bar{\alpha}(\cos \sigma - 1) \left(\frac{\sin \beta_1}{\sin \alpha_{ch}} \right) \quad (G15)$$

$$\bar{A}'_{ii} = cZ\bar{\alpha} \sin \sigma \left(\frac{\sin \beta_1}{\sin \alpha_{ch}} \right) \quad (G16)$$

Similarly, the other area perturbations will be

$$\bar{A}'_{*R} = \left(Z - \frac{x^*}{c} \right) c\bar{\alpha} (\cos \sigma - 1) \quad (G17)$$

$$\bar{A}'_{*I} = \left(Z - \frac{x^*}{c} \right) c\bar{\alpha} \sin \sigma \quad (G18)$$

At the shock wave, the area perturbations will be

$$\bar{A}'_{sR} = \left(Z - \frac{x_s}{c} \right) c\bar{\alpha} (\cos \sigma - 1) \quad (G19)$$

$$\bar{A}'_{sI} = \left(Z - \frac{x_s}{c} \right) c\bar{\alpha} \sin \sigma \quad (G20)$$

At the exit, the area perturbations will be

$$\bar{A}'_{ER} = (Z - 1) (\cos \sigma - 1) c\bar{\alpha} \left(\frac{\sin \beta_2}{\sin \alpha_{ch}} \right) \quad (G21)$$

$$\bar{A}'_{EI} = (Z - 1) c\bar{\alpha} \sin \sigma \left(\frac{\sin \beta_2}{\sin \alpha_{ch}} \right)$$

DEFINITION OF THE VOLUME PERTURBATIONS

Bending Mode Perturbations

Figure G3 defines and divides the volumes into three sections. Each section is dealt with separately in this Appendix.

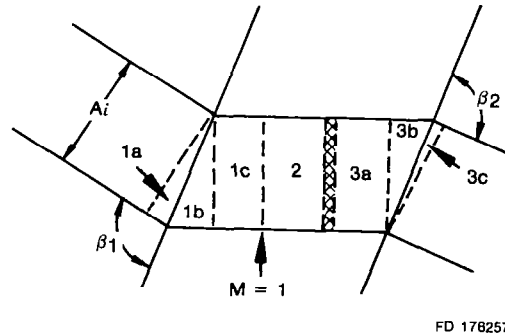


Figure G3. Bending Mode and Torsional Mode Volume Perturbations

Section 1

For ease of determination, Section 1 is divided into three subsections, such that

$$V_1 = V_{1a} + V_{1b} + V_{1c} \quad (G22)$$

Subsection 1a

$$V_{1a} = \frac{A_1^2}{2 \cotan (90 - \beta_1)}$$

Expanding this expression in small perturbations and neglecting higher order terms gives

$$V_{1a} = \frac{(\bar{A}_1^2 + 2\bar{A}_1 A'_1)}{2 \cotan (90 - \beta_1)} \quad (G23)$$

which produces for the perturbation volume in terms of real and imaginary components

$$\bar{V}'_{iR} = \frac{\bar{A}_i \bar{A}'_{iR}}{\cotan(90 - \beta_i)}$$

and

$$\bar{V}'_{iI} = \frac{\bar{A}_i \bar{A}'_{iI}}{\cotan(90 - \beta_i)} \quad (G24)$$

where \bar{A}'_{iR} and \bar{A}'_{iI} are given in Equations A4 and A5, respectively, in Appendix A.

Subsection 1b

$$V_{ib} = \frac{1}{2} (\delta \cos \alpha_{ch}) (\delta \sin \alpha_{ch})$$

where δ represents the gap between the blades. Expanding this expression in small perturbations gives

$$V_{ib} = \frac{1}{2} (\bar{\delta}^2 + 2 \bar{\delta} \delta') \sin \alpha_{ch} \cos \alpha_{ch}$$

or

$$V_{ib} = \bar{\delta} \left(\frac{h' - h}{\sin \alpha_{ch}} \right) \sin \alpha_{ch} \cos \alpha_{ch}$$

Breaking this expression down into its real and imaginary components gives

$$\bar{V}'_{iR} = \bar{\delta} \bar{h} (\cos \sigma - 1) \cos \alpha_{ch} \quad (G25)$$

$$\bar{V}'_{iI} = \bar{\delta} \bar{h} \sin \sigma \cos \alpha_{ch} \quad (G26)$$

Subsection 1c

$$V_{ic} = A^* (x^* - \delta \cos \alpha_{ch})$$

Expanding this expression in small perturbations and neglecting higher order terms produces

$$V'_{ic} = (\bar{A}^* - A'^*) x^* - (\bar{\delta} \bar{A}^* + \bar{\delta}' A'^* + \delta' \bar{A}^*) \cos \alpha_{ch}$$

or

$$V'_{ic} = A'^* (x^* - \bar{\delta} \cos \alpha_{ch}) + \delta' \bar{A}^* \cos \alpha_{ch} \quad (G27)$$

or

$$\bar{V}'_{iR} = \bar{h} (x^* - \bar{\delta} \cos \alpha_{ch}) (\cos \sigma - 1) + \bar{A}^* \bar{h} (\cos \sigma - 1) \cot \alpha_{ch}$$

$$\bar{V}'_{iI} = \bar{h} (x^* - \bar{\delta} \sin \alpha_{ch}) \sin \sigma + \bar{A}^* \bar{h} \sin \sigma \cot \alpha_{ch} \quad (G28)$$

Thus, the real and imaginary components for Section 1 represent a summation of the three subsections, as shown in Equation G29.

$$\begin{aligned}\bar{V}'_{1R} &= \bar{V}'_{1aR} + \bar{V}'_{1bR} + \bar{V}'_{1cR} \\ \bar{V}'_{1I} &= \bar{V}'_{1aI} + \bar{V}'_{1bI} + \bar{V}'_{1cI}\end{aligned}\tag{G29}$$

Section 2

Since this section represents supersonic flow, downstream disturbances cannot propagate upstream. Therefore, the volume perturbation can be represented as being bounded on the downstream side by a steady-state shock location, as shown in Equation G30.

$$V_2 = \frac{1}{2} (A^* + A_g) (\bar{x}_s - x^*)\tag{G30}$$

Expanding in small perturbations produces

$$V'_2 = \frac{1}{2} (A'^* + A'_g) (\bar{x}_s - x^*)\tag{G31}$$

Breaking the expression into its real and imaginary components gives

$$\begin{aligned}\bar{V}'_{2R} &= \bar{h} (\cos \sigma - 1) (\bar{x}_s - x^*) \\ \bar{V}'_{2I} &= \bar{h} \sin \sigma (\bar{x}_s - x^*)\end{aligned}\tag{G32}$$

Section 3

Section 3 also has been divided into three subsections. Each subsection is dealt with separately with the volume perturbations of Section 3 equal to the summation of the subsections.

Subsection 3a

$$\begin{aligned}V_{3a} &= \frac{1}{2} (A_s + A_{\frac{x}{c} - 1}) (c - x_s) \\ V'_{3a} &= \frac{1}{2} [(A'_s + A'_{\frac{x}{c} - 1}) (c - \bar{x}_s) - (\bar{A}_s + \bar{A}_{\frac{x}{c} - 1}) x'_s] \\ \bar{V}'_{3aR} &= \bar{h} (\cos \sigma - 1) (c - \bar{x}_s) - \frac{1}{2} \bar{x}'_{SR} (\bar{A}_s + \bar{A}_{\frac{x}{c} - 1})\end{aligned}\tag{G33}$$

$$\bar{V}'_{3aI} = \bar{h} \sin \sigma (c - \bar{x}_s) - \frac{1}{2} x'_{SI} (\bar{A}_s + \bar{A}_{\frac{x}{c} - 1})\tag{G34}$$

Subsection 3b

$$V_{3b} = \frac{1}{2} (\delta \cos \alpha_{ch})(\delta \sin \alpha_{ch})$$

$$V'_{3b} = \delta\delta' \cos \alpha_{ch} \sin \alpha_{ch}$$

$$\bar{V}'_{3bR} = \frac{1}{2} \bar{h} (\cos \sigma - 1)(\delta \sin 2\alpha_{ch}) \quad (G35)$$

$$\bar{V}'_{3bI} = \frac{1}{2} \bar{h} \sin \sigma (\delta \sin 2\alpha_{ch}) \quad (G36)$$

Subsection 3c

$$V_{3c} = \frac{A_E^2}{2 \tan \beta_2}$$

$$V'_{3c} = \frac{\bar{A}_E A'_E}{\tan \beta_2}$$

$$\bar{V}'_{3cR} = \frac{\bar{A}_E \bar{h} (\cos \sigma - 1) \cos \beta_2}{\sin \alpha_{ch}} \quad (G37)$$

$$V'_{3cI} = \frac{\bar{A}_E \bar{h} \sin \sigma \cos \beta_2}{\sin \alpha_{ch}} \quad (G38)$$

As stated, the real and imaginary components of Section 3 represent a summation of the three subsections, as shown below:

$$\begin{aligned} \bar{V}'_{3R} &= \bar{V}'_{3aR} + \bar{V}'_{3bR} + \bar{V}'_{3cR} \\ \bar{V}'_{3I} &= \bar{V}'_{3aI} + \bar{V}'_{3bI} + \bar{V}'_{3cI} \end{aligned} \quad (G39)$$

Torsional Mode Perturbations

As before, the individual sections will be dealt with separately in the following paragraphs.

Section 1

Section 1 is again divided into three subsections for ease of determination.

Subsection 1a

$$V'_{1a} = \frac{\bar{A}_i A'_i}{\cotan(90 - \beta_i)}$$

$$\bar{V}'_{1aR} = \frac{\bar{A}_i c Z \bar{\alpha} (\cos \sigma - 1) \sin \beta_i}{\sin \alpha_{ch} \cotan(90 - \beta_i)} \quad (G40)$$

$$\bar{V}'_{1aI} = \frac{\bar{A}_i c Z \bar{\alpha} \sin \sigma \sin \beta_i}{\sin \alpha \cotan(90 - \beta_i)} \quad (G41)$$

Subsection 1b

$$V'_{1b} = \frac{1}{2} \bar{\delta} \delta' \sin(2 \alpha_{ch})$$

$$\bar{V}'_{1bR} = c Z \bar{\alpha} \bar{\delta} (\cos \sigma - 1) \cos \alpha_{ch} \quad (G42)$$

$$\bar{V}'_{1bI} = c Z \bar{\alpha} \bar{\delta} \sin \sigma \cos \alpha_{ch} \quad (G43)$$

Section 1c

$$V'_{1c} = \int_{x_{ENT}}^{x^*} A'(x) dx \quad A'(x) = c \left(Z - \frac{x}{c} \right) (\alpha' - \alpha)$$

$$= c Z (x^* - \bar{\delta} \cos \alpha_{ch}) (\alpha' - \alpha) - \frac{c [x^{*2} - (\bar{\delta} \cos \alpha_{ch})^2] (\alpha' - \alpha)}{2c}$$

$$\bar{V}'_{1cR} = c \bar{\alpha} (\cos \sigma - 1) \left[Z (x^* - \bar{\delta} \cos \alpha_{ch}) - \frac{x^{*2} - (\bar{\delta} \cos \alpha_{ch})^2}{2c} \right] \quad (G44)$$

$$\bar{V}'_{1cI} = c \alpha \sin \sigma \left[Z (x^* - \bar{\delta} \cos \alpha_{ch}) - \frac{x^{*2} - (\bar{\delta} \cos \alpha_{ch})^2}{2c} \right] \quad (G45)$$

The summation of the real and imaginary components of the subsections is shown in Equation G46.

$$\bar{V}'_{1R} = \bar{V}'_{1aR} + \bar{V}'_{1bR} + \bar{V}'_{1cR}$$

$$\bar{V}'_{1I} = \bar{V}'_{1aI} + \bar{V}'_{1bI} + \bar{V}'_{1cI} \quad (G46)$$

Section 2

$$V'_2 = \int_{x^*}^{x_s} A'(x) dx$$

$$\bar{V}'_{2R} = c \bar{\alpha} (\cos \sigma - 1) \left[Z (x_s - x^*) - \left(\frac{x_s^2 - x^{*2}}{2c} \right) \right] \quad (G47)$$

$$\bar{V}'_{2I} = c \bar{\alpha} \sin \sigma \left[Z (x_s - x^*) - \left(\frac{x_s^2 - x^{*2}}{2c} \right) \right] \quad (G48)$$

Section 3

As before, Section 3 will be divided into three subsections.

Subsection 3a

$$\begin{aligned}
 V'_{3a} &= \int_{x_s}^{x_E} A'(x) dx - \bar{A}_s x'_s \\
 &= c (\alpha' - \alpha) \left[Z (x_E - x_s) - \left(\frac{x_E^2 - x_s^2}{2c} \right) \right] - \bar{A}_s x'_s \\
 \bar{V}'_{3aR} &= c \alpha (\cos \sigma - 1) \left[Z (c - \bar{x}_s) - \left(\frac{c^2 - \bar{x}_s^2}{2c} \right) \right] - \bar{A}_s \bar{x}'_{sR} \tag{G49}
 \end{aligned}$$

$$\bar{V}'_{3aI} = c \bar{\alpha} \sin \sigma \left[Z (c - \bar{x}_s) - \left(\frac{c^2 - \bar{x}_s^2}{2c} \right) \right] - \bar{A}_s \bar{x}'_{sI} \tag{G50}$$

Subsection 3b

$$V_{3b} = \frac{1}{2} (\bar{\delta}^2 + 2 \bar{\delta} \delta') \sin \alpha_{ch} \cos \alpha_{ch}$$

$$V'_{3b} = \bar{\delta} \delta' \sin \alpha_{ch} \cos \alpha_{ch}$$

where,

$$\delta' = \frac{c(Z-1)(\bar{\alpha}-\alpha)}{\sin \alpha_{ch}}$$

$$\bar{V}'_{3bR} = c \bar{\alpha} (Z-1) (\cos \sigma - 1) \cos \alpha_{ch} \cdot \bar{\delta} \tag{G51}$$

$$\bar{V}'_{3bI} = c \bar{\alpha} (Z-1) \sin \sigma \cos \alpha_{ch} \cdot \bar{\delta} \tag{G52}$$

Subsection 3c

$$V_{3c} = \frac{A_E^2}{2 \tan \beta_2}$$

$$V'_{3c} = \frac{\bar{A}_E A'_E}{\tan \beta_2}$$

$$\bar{V}'_{3cR} = \frac{\bar{A}_E c \bar{\alpha} (Z-1) (\cos \sigma - 1) \cos \beta_2}{\sin \alpha_{ch}} \tag{G53}$$

$$\bar{V}'_{3cI} = \frac{\bar{A}_E c \bar{\alpha} (Z-1) \sin \sigma \cos \beta_2}{\sin \alpha_{ch}} \tag{G54}$$

The resulting summations for Section 3 are given in Equation G55.

$$\bar{V}'_{3R} = \bar{V}'_{3aR} + \bar{V}'_{3bR} + \bar{V}'_{3cR}$$

$$\bar{V}'_{3I} = \bar{V}'_{3aI} + \bar{V}'_{3bI} + \bar{V}'_{3cI} \tag{G55}$$

APPENDIX H
CALCULATION OF THE STEADY-STATE SHOCK
LOCATION AND TEST FOR CHOKED FLOW

The first step in the analysis requires checking for choked flow. Equation H1 is obtained from the isentropic flow relationships.

$$\left(\frac{A_c}{A^*} \right)^2 = \frac{1}{M_c^2} \left[\frac{2}{\gamma + 1} \left(1 + \frac{\gamma - 1}{2} M_c^2 \right) \right] \left(\frac{\gamma + 1}{\gamma - 1} \right) \quad (H1)$$

Rearranging this relationship and solving for A_c gives

$$A_c = \left\{ \frac{1}{M_c^2} \left[\frac{2}{\gamma + 1} \left(1 + \frac{\gamma - 1}{2} M_c^2 \right) \right] \left(\frac{\gamma + 1}{\gamma - 1} \right) \right\}^{1/2} A^* \quad (H2)$$

For a choked flow condition, A^* will be equal to or greater than the minimum area between the blades (A_{min}). Thus, for choked flow Equation H2 becomes

$$A_c \geq \left\{ \frac{1}{M_c^2} \left[\frac{2}{\gamma + 1} \left(1 + \frac{\gamma - 1}{2} M_c^2 \right) \right] \left(\frac{\gamma + 1}{\gamma - 1} \right) \right\}^{1/2} A^* \quad (H3)$$

During steady-state conditions

$$A_c = \delta \sin \beta_1 \quad (H4)$$

and

$$A_{min} = (\delta \sin \alpha_{cn}) + y_L x_2^* - y_u x_1^* \quad (H5)$$

Substituting the relationships of Equations H4 and H5 into Equation H3 and solving for β_1 yields

$$\beta_1 = \sin^{-1} \left[\frac{1}{\delta} \left\{ \frac{1}{M_c^2} \left[\frac{2}{\gamma + 1} \left(1 + \frac{\gamma - 1}{2} M_c^2 \right) \right] \left(\frac{\gamma + 1}{\gamma - 1} \right) \right\}^{1/2} \left\{ \delta \sin \alpha_{cn} + y_L(x_2^*) - y_u(x_1^*) \right\} \right] \quad (H6)$$

Thus, the flow will be choked if the input β_1 is equal to or greater than the β_1 calculated in Equation H6. If the flow is choked, the areas at various cross-sections will be calculated along with the exit area, as shown.

$$A_{exit} = \delta \sin \beta_2 \quad (H7)$$

The next step in the analysis requires that the supersonic Mach numbers downstream of the throat be calculated by solving the following equation for M at the specific chordwise locations:

$$\frac{A(x)}{A_{min}} = \left\{ \frac{1}{M^2(x)} \left[\frac{2}{\gamma + 1} \left(1 + \frac{\gamma - 1}{2} M^2(x) \right) \right] \left(\frac{\gamma + 1}{\gamma - 1} \right) \right\}^{1/2} \quad (H8)$$

Once the area relationships and Mach numbers as a function of x have been determined, the analysis proceeds with the calculation of the steady-state shock position. The derivation requires only the Mach numbers downstream of the throat since it can be shown that the shock position is stable only in a diverging channel. With the pressure ratio across the stage (PR) input, the capture Mach number fixes the pressure at the entrance, as shown in Equation H9.

$$\frac{P_c}{P_t} = \left(1 + \frac{\gamma-1}{2} M_c^2 \right)^{\left(\frac{-\gamma}{\gamma-1} \right)} \quad (\text{H9})$$

The pressure ratio from the throat to the exit can be defined as

$$\text{PR}_{\text{corr}} = \frac{P_c}{P_t} \left(\frac{2}{\gamma+1} \right)^{\frac{-\gamma}{\gamma+1}} \text{PR} \quad (\text{H10})$$

The use of an iteration procedure finds the shock location by first assuming the shock is located at the throat and then incrementally moving downstream until the obtainable pressure ratio (PR_{obt}) at the specific shock location matches the PR_{corr} within some tolerance ϵ . PR_{obt} is calculated in the following manner:

1. Calculate the ratio of pressure entering the shock to the pressure at the throat as

$$\frac{P_{\text{us}}}{P^*} = \frac{\left(1 + \frac{\gamma-1}{2} M_s^2 \right)^{\left(\frac{-\gamma}{\gamma-1} \right)}}{\left(\frac{\gamma+1}{2} \right)^{\left(\frac{-\gamma}{\gamma-1} \right)}} \quad (\text{H11})$$

2. Calculate the pressure rise across the throat as

$$\frac{P_{\text{ds}}}{P_{\text{us}}} = 1 + \frac{2\gamma}{\gamma+1} (M_s^2 - 1) \quad (\text{H12})$$

3. Calculate the Mach number exiting the shock to determine the nature of the flow (subsonic or supersonic) as

$$M_{\text{ds}}^2 = \frac{1 + \frac{\gamma-1}{2} M_{\text{us}}^2}{\gamma M_{\text{us}}^2 - \frac{\gamma-1}{2}} \quad (\text{H13})$$

where,

$$M_{\text{us}} = M_s$$

Knowing the nature of the flow, M_{exit} can be calculated from the area relationship in Equation H8. Knowing M_{exit} , the pressure ratio from the shock location to the exit can be calculated by

$$\frac{P_{exit}}{P_{ds}} \left(\frac{1 + \frac{\gamma-1}{2} M_{exit}^2}{1 + \frac{\gamma-1}{2} M_{ds}^2} \right)^{\left(\frac{-\gamma}{\gamma-1} \right)} \quad (H14)$$

From this point, PR_{obt} can be defined from Equations H11, H12, and H13 as

$$PR_{obt} = \left(\frac{P_{us}}{P^*} \right) \left(\frac{P_{ds}}{P_{us}} \right) \left(\frac{P_{exit}}{P_{ds}} \right) \quad (H15)$$

Equation H14 is valid for all values for M_s less than or equal to 1.1. If $M_s > 1.1$, there will be a loss in total pressure. Calculate total pressure in this situation in the following manner:

$$\frac{P_{totc}}{P_{totexit}} = \left[1 + \frac{2\gamma}{\gamma+1} (M_s^2 - 1) \right]^{\left(\frac{1}{\gamma-1} \right)} \left[\frac{(\gamma-1) M_s^2 + 2}{(\gamma+1) M_s^2} \right]^{\left(\frac{\gamma}{\gamma+1} \right)} \quad (H16)$$

Then, Equation H15 becomes

$$PR_{obt} = \left(\frac{P_{us}}{P^*} \right) \left(\frac{P_{ds}}{P_{us}} \right) \left(\frac{P_{exit}}{P_{ds}} \right) \left(\frac{P_{totexit}}{P_{totc}} \right) \quad (H17)$$

The shock will be located at the point where

$$| PR_{corr} - PR_{obt} | \leq \epsilon \quad (H18)$$

If this relationship is not valid, move to the next discrete section downstream and repeat the calculations. If the shock position is not located in this manner before reaching the end of the channel, the shock location is downstream of the blades.

APPENDIX I
LIFT AND MOMENT COEFFICIENT CALCULATION

With pressure perturbations known at the inlet and outlet of each section, the mean section pressure perturbations can be defined as follows:

Section 1

$$\bar{p}'_{1R} = (\bar{p}'_{iR} + \bar{p}'_{*R}) \div 2 \qquad \bar{p}'_{1I} = (\bar{p}'_{iI} + \bar{p}'_{*I}) \div 2 \qquad (I-1)$$

Section 2

$$\bar{p}'_{2R} = (\bar{p}'_{*R} + \bar{p}'_{USR}) \div 2 \qquad \bar{p}'_{2I} = (\bar{p}'_{*I} + \bar{p}'_{USI}) \div 2 \qquad (I-2)$$

Section 3

$$\bar{p}'_{3R} = (\bar{p}'_{dR} + \bar{p}'_{ER}) \div 2 \qquad \bar{p}'_{3I} = (\bar{p}'_{dI} + \bar{p}'_{EI}) \div 2 \qquad (I-3)$$

Equations I-1 through I-3 give the unsteady pressure distribution on the reference airfoil suction surface. The pressure perturbation for the "channel" below the reference channel can be described in the following manner:

$$p_L = p_u \exp(-i \sigma) \qquad (I-4)$$

where, L denotes the lower airfoil surface and U denotes the upper airfoil surface,

or

$$p_{LR} + ip_{LI} = (p_{UR} + ip_{UI}) (\cos \sigma - i \sin \sigma) \qquad (I-5)$$

which gives

$$\begin{aligned} p_{LR} &= p_{UR} \cos \sigma + p_{UI} \sin \sigma \\ p_{LI} &= p_{UI} \cos \sigma - p_{UR} \sin \sigma \end{aligned} \qquad (I-6)$$

Then, the pressure perturbations for each section can be calculated as follows:

Section 1

$$\begin{aligned} \bar{p}'_{1RL} &= \bar{p}'_{1RU} \cos \sigma + \bar{p}'_{1IU} \sin \sigma \\ \bar{p}'_{1IL} &= \bar{p}'_{1IU} \cos \sigma - \bar{p}'_{1RU} \sin \sigma \end{aligned} \qquad (I-7)$$

Section 2

$$\begin{aligned} \bar{p}'_{2RL} &= \bar{p}'_{2RU} \cos \sigma + \bar{p}'_{2IU} \sin \sigma \\ \bar{p}'_{2IL} &= \bar{p}'_{2IU} \cos \sigma - \bar{p}'_{2RU} \sin \sigma \end{aligned} \qquad (I-8)$$

Section 3

$$\begin{aligned} \bar{p}'_{3RL} &= \bar{p}'_{3RU} \cos \sigma + \bar{p}'_{3IU} \sin \sigma \\ \bar{p}'_{3IL} &= \bar{p}'_{3IU} \cos \sigma - \bar{p}'_{3RU} \sin \sigma \end{aligned} \qquad (I-9)$$

With the unsteady pressures defined on the suction and pressure sides of the airfoil, the unsteady lift and moment coefficients can be defined as follows for each mode. The real part of the unsteady lift coefficient is

$$C_{LR} = C_{LUR} + C_{LLR} \quad (I-10)$$

where,

$$C_{LUR} = -(\bar{p}'_{1UR} x^* + \bar{p}'_{2UR} (x_s - x^*) + \bar{p}'_{3UR} (c - x_s)) \div b \quad (I-11)$$

$$C_{LLR} = +(\bar{p}'_{1LR} (x^* - \delta \cos \alpha_{ch}) + \bar{p}'_{2LR} (x_s - x^*) + \bar{p}'_{3LR} (c - x_s + \delta \cos \alpha_{ch})) \div b \quad (I-12)$$

The imaginary part of the unsteady lift coefficient is

$$C_{Li} = C_{LUI} + C_{LLI} \quad (I-13)$$

where, C_{LUI} and C_{LLI} are defined in the same manner as in Equations I-11 and I-12 with the imaginary pressure perturbations used. The real part of the moment coefficient can be defined for each mode as

$$C_{MR} = C_{MUR} + C_{MLR} \quad (I-14)$$

where,

$$C_{MUR} = - \left[\bar{p}'_{1UR} x^* \left(Zc - \frac{x^*}{2} \right) + \bar{p}'_{2UR} (x_s + x^*) \left(Zc - \frac{(x^* + x_s)}{2} \right) + \bar{p}'_{3UR} (c - x_s) \left(Zc - \frac{(c + x_s)}{2} \right) \right] \div b^2 \quad (I-15)$$

$$C_{MLR} = + \left[\bar{p}'_{1LR} (x^* - \delta \cos \alpha_{ch}) \left(Zc \left(\frac{x^* - \delta \cos \alpha_{ch}}{2} \right) \right) + \bar{p}'_{2LR} (x_s - x^*) (Zc - \frac{1}{2} (x_s + x^* - 2\delta \cos \alpha_{ch})) + \bar{p}'_{3LR} (c - x_s + \delta \cos \alpha_{ch}) (Zc - \frac{1}{2} (c + x_s - \delta \cos \alpha_{ch})) \right] \div b^2 \quad (I-16)$$

The imaginary part of the moment coefficient for each mode is

$$C_{MI} = C_{MUI} + C_{MLI} \quad (I-17)$$

where, C_{MUI} and C_{MLI} are defined as in Equations I-15 and I-16 with the imaginary pressure perturbations used.

*VERSION 1.3.0 (01 MAY 80) SYSTEM/370 FORTRAN H EXTENDED (ENHANCED) DATE 20.353/15.02.47 PAGE 1
 REQUESTED OPTIONS: EBDCIC,MAP,HOLIST,NOBDECK,XREF,OPT=2
 OPTIONS IN EFFECT: H/ME(MAIN) OPTIMIZE(2) LINECOUNT(60) SIZE(MAX) AUTOBOL(NONE)
 SOURCE EBDCIC HOLIST NOBDECK OBJECT MAP NOFOPMAT GCSHTM XREF NOALC NOANSF TERM IBM_FLAG(I)

C		DATA SET 9066MAIN	AT LEVEL 001 AS OF 07/30/80	00000000				
C		DATA SET 9066MAIN	AT LEVEL 001 AS OF 04/21/80	00000010				
C		DECK 9066	REVISED VERSION OF DECK 8303	00000020				
C			NON TO BE USED AS THE CHANNEL FLOW CHOKE FLUTTER	00000030				
C			UNSTEADY AERODYNAMIC MODEL	00000040				
C				00000050				
C		ANALYST	JEFFREY F. SIMPSON	00000060				
C		DATE	JANUARY, 1980	00000070				
C		EXT	4315	00000080				
C		MAIL LOC	R-47, BLDG 32, PEG	00000090				
C				00000100				
C	*	ALPBAR	- MEAN TORSIONAL DEFLECTION THRU CYCLE (DEG.)	00000110				
C	*	ALPCH	- STAGGER OF BLADE FOM (DEG.)	00000120				
C	*	BETA1	- INLET AIR ANGLE (DEG.)	00000130				
C	*	BETA2	- EXIT AIR ANGLE (DEG.)	00000140				
C	*	C	- CHORD (IN.)	00000150				
C	*	DELTA	- GAP BETWEEN BLADES (IN.)	00000160				
C	*	E	- NONDIMENSIONAL ELASTIC AXIS LOCATION REF. TO MIDSPAN	00000170				
C	*	EPS	- TOLERANCE FOR PRESSURE RATIO	00000180				
C	*	GAM	- SPECIFIC HEAT RATIO	00000190				
C	*	HBAR	- MEAN FLAPPING DEFLECTION OF BLADE THRU THE CYCLE (IN.)	00000200				
C	*	MI	- INLET MACH NUMBER	00000210				
C	*	M1	- L.E. MACH NUMBER	00000220				
C	*	NP	- NUMBER OF AIRFOIL COORDINATES	00000230				
C	*	NSECT	- NUMBER OF SEGMENTS FOR CHANNEL AREA, MACH NUMBER, ETC.	00000240				
C	*	NTIME	- NUMBER OF TIME INCREMENTS	00000250				
C	*	OMEGA	- FREQUENCY OF VIBRATION (RAD. / SEC.)	00000260				
C	*	P?	- STATIC PRESSURE RATIO ACROSS STAGE	00000270				
C	*	PT	- TOTAL PRESSURE ENTERING STAGE (PSI)	00000280				
C	*	RGAS	- GAS CONSTANT (FT-LBF / DEG. R-LE(MCL))	00000290				
C	*	TOLMIT	- TOLERANCE FOR MACH NUMBER ITERATION	00000300				
C	*	TT	- TOTAL TEMPERATURE ENTERING STAGE (DEG. R)	00000310				
C	*	XIN	- ARRAY OF AIRFOIL X COORDINATES (IN.)	00000320				
C	*	YLIN	- ARRAY OF AIRFOIL Y LOWER COORDINATES (IN.)	00000330				
C	*	YUIN	- ARRAY OF AIRFOIL Y UPPER COORDINATES (IN.)	00000340				
C				00000350				
ISH 0002	COMMON	VOL1,	VOL2,	VOL3,	AE,	ASTAR,	E,	00000360
?		AC,	AC,	MO,	MSE,	MOI,	TT,	00000370
?		XS,	XSTAR,	P1CS,	P2SS,	P3SS,	C,	00000380
?		MI,	PT,	BETA1,	BETA2,			00000390
?		ALPCH,	PHIBB,	FHIBT,	HBAR,	DELTA,		00000400
?		OMEGA,	ALPBAR,	V				00000410
C								00000420
ISH 0003	REAL	LAMBDA,	LSS,	LLSS,	MI,	M1,	ME,	00000430
?		MO,	MSS,	MCS,	MOI,	MC,	MOSS,	00000440
?		LSS,	MUSS,	HLSS				00000450
C								00000460
ISH 0004	DIMENSION	ASS(51),	HSS(51),	TITLE(20),	XIN(50),	YLIN(50),		00000470
?		YUIN(50),	DMACH(25),	XIP(51),	XMACH(25),	YMACH(25)		00000480
ISH 0005		FLOF(XM) =	1. + .5 * (GAM - 1.) * XM * XM					00000500
C								00000510
C		READ INPUT						00000520
C								00000530
ISH 0006		RGAS =	1716.26					00000540
ISH 0007	10	READ(5,20,END=500)	TITLE					00000550

ISN 0008	80	FORMAT(20A4)	00000560
ISN 0009		READ(5,85) NP,NS,IMAVE,NSECT	00000570
ISN 0010	85	FORMAT(4I5)	00000580
ISN 0011		READ(5,86) ALPCH,BETA1,BETA2,C,DELTA,E,EPS,GAM,MI,M1,OMEGA,PR,PT, TT,V,DIAM,TOLMIT,ALPBAR,HBAR	00000590
ISN 0012	86	FORMAT(8F10.0)	00000600
ISN 0013		READ(5,86) (XIN(I),I=1,NP)	00000620
ISN 0014		READ(5,86) (YLIN(I),I=1,NP)	00000630
ISN 0015		READ(5,86) (YUIN(I),I=1,NP)	00000640
ISN 0016		IF (EPS .EQ. 0.) EPS = .05	00000650
ISN 0018		IF (GAM .EQ. 0.) GAM = 1.4	00000660
ISN 0020		IF (M1 .EQ. 0.) M1 = MI	00000670
ISN 0022		IF (RGAS .EQ. 0.) RGAS = 1716.26	00000680
ISN 0024		IF (TOLMIT .EQ. 0.) TOLMIT = .00001	00000690
ISN 0026		TAU = DELTA	00000700
ISN 0027		PHIIB3 = 6.2832 * DIAM/NS	00000710
ISN 0028		IF (IMAVE.LT.0) PHIIB5 = -1.0 * PHIIB3	00000720
ISN 0030		PHIIB7 = PHIIB3	00000730
	C		00000740
	C	PRINT INPUT	00000750
	C		00000760
ISN 0031		WRITE(6,5) TITLE,	00000770
	?	ALPBAR,ALPCH,BETA1,BETA2,C,DELTA,DIAM,E,EPS,GAM,HBAR,	00000780
	?	IMAVE,MI,M1,NS,NP,NSECT,HTIME,OMEGA,PHIIB3,PHIIB7,PR,V	00000790
ISN 0032	5	FORMAT('ICH'//MACH FLOW CHOICE FLUTTER ANALYSIS DECK 9066'// 20A4'//	00000800
	1'	ALPBAR',F10.5,5X,'MEAN TORSIONAL DEFLECTION THRU CYCLE (DEG.)'//	00000810
	2'	ALPCH',F10.3,5X,'STAGGER ANGLE OF BLADE ROW (DEG.)'//	00000820
	3'	BETA1',F10.3,5X,'INLET AIR ANGLE (DEG.)'//	00000830
	4'	BETA2',F10.3,5X,'EXIT AIR ANGLE (DEG.)'//	00000840
	5'	C',F10.5,5X,'CHORD (IN.)'//	00000850
	6'	DELTA',F10.5,5X,'GAP BETWEEN BLADES (IN.)'//	00000860
	7'	DIAM',F10.3,5X,'NOZLE DIAMETER'//	00000870
	8'	E',F10.6,5X,'ELASTIC AXIS LOCATION REF. TO MIDSPAN'//	00000880
	9'	EPS',F10.5,5X,'TOLERANCE FOR PRESSURE RATIO'//	00000890
	X'	GAM',F10.5,5X,'SPECIFIC HEAT RATIO'//	00000900
	1'	HBAR',F10.6,5X,'MEAN FLAPPING DEFLECTION OF BLADE (IN.)'//	00000910
	2'	IMAVE',I10,5X,'HAVE MOTION'//	00000920
	3'	MI',F10.5,5X,'INLET MACH NUMBER'//	00000930
	4'	M1',F10.5,5X,'L. E. MACH NUMBER'//	00000940
	5'	NS',I10,5X,'NUMBER OF BLADES'//	00000950
	6'	NP',I10,5X,'NUMBER OF AIRFOIL COORDINATES'//	00000960
	7'	NSECT',I10,5X,'NUMBER OF SECTIONS FROM L.E. TO T.E.'//	00000970
	8'	HTIME',I10,5X,'NUMBER OF TIME INCREMENTS'//	00000980
	9'	OMEGA',F10.2,5X,'FREQUENCY OF VIBRATION (CYCLES. / SEC.)'//	00000990
	X'	PHIIB3',F10.4,5X,'BENDING MODE INTERBLADE PHASE ANGLE (RAD.)'//	00001000
	1'	PHIIB7',F10.4,5X,'TORSIONAL MODE INTERBLADE PHASE ANGLE (RAD.)'//	00001010
	2'	PR',F10.5,5X,'STATIC PRESSURE RATIO ACROSS STAGE'//	00001020
	3'	V',F10.3,5X,'RELATIVE INLET VELOCITY (FT/SEC)'//	00001030
		ERRC? DETECTED - SCAN POINTER = 0	
ISN 0033		WRITE (6,6) PT,RGAS,TAU,TOLMIT,TT	00001040
ISN 0034	6	FORMAT(' PT',F10.4,5X,'TOTAL PRESSURE ENTERING STAGE (PSI)'//	00001050
	1'	RGAS',F10.2,5X,'GAS CONSTANT (FT**2 / (SEC**2 DEG. R))'//	00001060
	2'	TAU',F10.5,5X,'GAP BETWEEN AIRFOILS WHILE OSCILLATING (IN.)'//	00001070
	3'	TOLMIT',F10.6,5X,'TOLERANCE FOR MACH NUMBER ITERATION'//	00001080
	4'	TT',F10.2,5X,'TOTAL TEMPERATURE ENTERING STAGE (DEG. R)'//	00001090
ISN 0035		WRITE (6,7) (XIN(I),YLIN(I),YUIN(I),I=1,NP)	00001100
ISN 0036	7	FORMAT('0AIRFOIL COORDINATES (IN.)'//10X,'X Y LOWER Y UPPER'//	00001110
		1(1X,3F10.5))	00001120
	C		00001130

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ISN 0037      DCR = 57.2057795100      00001140
ISN 0038      ALPBAR = ALPBAR / DCR      00001150
ISN 0039      ALPCH = ALPCH / DCR      00001160
ISN 0040      BETA1 = BETA1 / DCR      00001170
ISN 0041      BETA2 = BETA2 / DCR      00001180
C      00001190
C      DEFINE QUANTITIES USED FREQUENTLY IN CALCULATIONS      00001200
C      00001210
ISN 0042      PI = 3.141592653589800      00001220
ISN 0043      OMEGA = OMEGA      00001230
ISN 0044      OMEGA = OMEGA * 2. * PI      00001240
ISN 0045      NSECT1 = NSECT + 1      00001250
ISN 0046      NTIME1 = NTIME + 1      00001260
C      00001270
ISN 0047      B = .5 * C      00001280
ISN 0048      DIPE = B * (1. - E)      00001290
ISN 0049      SIPE = B * (1. + E)      00001300
C      00001310
ISN 0050      RHO1 = PI / (TT * RGAS * 12.)      00001320
ISN 0051      ZP = DIPE / C      00001330
C      00001340
ISN 0052      CASCADE STACGER ANGLE      00001350
ISN 0053      LAMDA = .5 * FI - ALPCH      00001360
ISN 0054      COSLAM = COS(LAMDA)      00001370
ISN 0055      SINLAM = SIN(LAMDA)      00001380
ISN 0056      TSLAM = TAU * SINLAM      00001390
ISN 0057      GAM = GAM / (1. - GAM)      00001400
ISN 0058      THROPL = 2. / (GAM + 1.)      00001410
ISN 0059      COSACH = COS(ALPCH)      00001420
ISN 0060      SINACH = SIN(ALPCH)      00001430
ISN 0061      SINB1 = SIN(BETA1)      00001440
ISN 0062      COSB2 = COS(BETA2)      00001450
ISN 0063      SINB2 = SIN(BETA2)      00001460
C      00001470
ISN 0064      EXIT STEADY STATE AREA      00001480
C      ASSE = TAU * SINB2      00001490
C      X INCREMENT FOR CHANNEL AREA, MACH NUMBER, ETC.      00001500
ISN 0065      DXIP = (C - TSLAM) / FLOAT(NSECT)      00001510
C      00001520
C      CALCULATE STEADY STATE AREAS THRU CHANNEL & LOCATE THROAT      00001530
C      00001540
ISN 0066      ASTAR = 1.E20      00001550
ISN 0067      DO 20 NS = 1, NSECT1      00001560
C      X COORDINATE OF REFERENCE BLADE      00001570
ISN 0068      XIP(NS) = FLOAT(NS-1) * DXIP + TSLAM      00001580
C      X COORDINATE OF BLADE ABOVE REFERENCE BLADE      00001590
ISN 0069      X2P = XIP(NS) - TSLAM      00001600
C      Y COORDINATE OF UPPER SURFACE OF REFERENCE BLADE      00001610
ISN 0070      CALL LINE (NP, XIN, YUIN, XIP(NS), YU, DUM, IDUM)      00001620
C      Y COORDINATE OF LOWER SURFACE OF UPPER BLADE      00001630
ISN 0071      CALL LINE (NP, XIN, YLIN, X2P, YL, DUM, IDUM)      00001640
C      ARRAY OF STEADY STATE AREAS THRU CHANNEL ABOVE REFERENCE BLADE      00001650
ISN 0072      ASS(NS) = TAU * COSLAM - YU + YL      00001660
ISN 0073      WRITE(6,16) NS,ASS(NS)      00001670
ISN 0074      16      FORMAT(2X,I2,4X,E15.6)      00001680
ISN 0075      IF (ASS(NS) .GT. ASTAR) GO TO 20      00001690
C      THROAT AREA      00001700
ISN 0077      ASTAR = ASS(NS)      00001710
C      DISTANCE FROM REFERENCE BLADE L.E. TO THROAT      00001720
ISN 0078      XSTAR = XIP(NS)

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ISH 0079		YUSTAR = YU	00001730
ICN 0080		YLSTAR = YL	00001740
	C	SUBSCRIPT OF THROAT LOCATION	00001750
ISH 0081		IT = HS	00001760
ISH 0092		20 CONTINUE	00001770
	C		00001780
	C	CALCULATE SUPERSONIC MACH NUMBERS FROM THROAT TO EXIT	00001790
	C		00001800
ISH 0083		CALL ZERO (MSS(1), MSS(NSECT1), 1.)	00001810
ISH 0084		DO 30 HS = 1, NSECT1	00001820
ISH 0085		IF (XIP(NS) .LT. XSTAR) GO TO 30	00001830
	C	RATIO OF AREA TO THROAT AREA	00001840
ISH 0087		AOAMIN = ASS(NS) / ASTAR	00001850
ISH 0088		ISS = 1	00001860
ISH 0089		CALL MACHIT (AOAMIN, MSS(NS), GAM, TOLMIT, ISS, KILL)	00001870
	C	IF (KILL .NE. 0) WRITE (6,25) AOAMIN, XIP(NS)	00001880
ISH 0090		25 FORMAT('0A/ASTAR = ',1PE15.5,' X = ',E15.5)	00001890
ISH 0091		30 CONTINUE	00001900
	C		00001910
	C	TEST FOR CHOKED FLOW AND LOCATE STEADY STATE SHOCK	00001920
	C		00001930
ISH 0092		BETAIC = ARSIN(SORT(((THOSP1 * FIOF(M1))**GPIGM1) / M1**2) / TAU * COSLAM + YLSTAR - YUSTAR))	00001940
	1	(TAU * COSLAM + YLSTAR - YUSTAR))	00001950
ISH 0093		IF (BETA1 .GE. BETAIC) GO TO 41	00001960
ISH 0095		WRITE (6,40)	00001970
ISH 0096		40 FORMAT(' ***** FLOW IS NOT CHOKED')	00001980
ISH 0097		ICHCKE = 0	00001990
ISH 0098		ASTAR = TAU*SINB1*MI*(2./(GAM+1.) * (1.+(GAM-1.)/2. * MI**2)) **	00002000
	?	(-(GAM+1.)/(2.*(GAM-1.)))	00002010
ISH 0099		XSTAR = DELTA * COSACH + 0.1	00002020
ISH 0100		IF(C-DELTA+COSACH .GT. 0.1) GOTO 330	00002030
ISH 0102		XSTAR = C - DELTA*COSACH	00002040
	C	WRITE(6,340) XSTAR	00002050
ISH 0103		340 FORMAT('/',' *** XSTAR RESET TO',E15.6)	00002060
ISH 0104		330 X0 = XSTAR + 0.15	00002070
ISH 0105		IF(C-X0 .GT. 0.) GOTO 61	00002080
ISH 0107		X0 = XSTAR + 0.05	00002090
	C	WRITE(6,360) X0	00002100
ISH 0108		360 FORMAT('/',' *** X0 RESET TO',E15.6)	00002110
ISH 0109		GO TO 61	00002120
	C	INLET STATIC TO TOTAL PRESSURE RATIO	00002130
ISH 0110		41 AC = TAU * SINB1	00002140
ISH 0111		AOAMIN = AC / ASTAR	00002150
ISH 0112		ISS = 0	00002160
ISH 0113		IF (MI .GT. 1.) ISS = 1	00002170
ISH 0115		CALL MACHIT (AOAMIN, MC, GAM, TOLMIT, ISS, KILL)	00002180
ISH 0116		PIPT = FIOF(MC)**GDIIMG	00002190
ISH 0117		BETA1 = BETAIC	00002200
ISH 0118		BETAPR = BETA1 * DCR	00002210
ISH 0119		WRITE(6,117) BETAPR1	00002220
ISH 0120		117 FORMAT('/',' BETA1',E15.6)	00002230
ISH 0121		ICHCKE = 1	00002240
	C	EXIT AREA	00002250
ISH 0122		AE = DELTA * SINB2	00002260
	C	RATIO OF EXIT AREA TO THROAT AREA	00002270
ISH 0123		AOAMIN = AE / ASTAR	00002280
ISH 0124		ISS = 0	00002290
ISH 0125		IF (AE .LT. ASTAR) ISS = 1	00002300
ISH 0127		CALL MACHIT (AOAMIN, ME, GAM, TOLMIT, ISS, KILL)	00002310

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*VERSION 1.3.0 (01 MAY 80)      MAIN      SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)      DATE 80.353/15.02.47      PAGE 5
C      IF (KILL.NE. 0) WRITE (6,45) AOAMIN, XIP(INS)      00002300
ISN 0128      45 FORMAT('0AE/ASTAR = ',IPE13.5,' X = ',E13.5)      00002330
ISN 0129      DO 60 NS = IT, NSECT1      00002340
C      SHOCK MACH NUMBER      00002350
ISN 0130      MC = MSS(NS)      00002360
ISN 0131      IF(MO.GT.0.999 .AND. MO.LT.1.001) MO=1.01      00002370
C      TOTAL INLET TO EXIT PRESSURE RATIO      00002380
ISN 0133      PIIPTI = 1.      00002390
ISN 0134      IF (MO .GT. 1.1) PIIPTI =
1          (1. + GAM * TNGSP1 * (MO**2 - 1.))**(.5 / (GAM - 1.)) *      00002410
2          (2. * FIOF(MO) / ((GAM + 1.) * MO**2))**(-GOIMG)      00002420
C      STATIC PRESSURE RATIO ACROSS STAGE      00002430
ISN 0136      PEPI = PR * PIIPTI      00002440
C      THROAT TO EXIT PRESSURE RATIO      00002450
ISN 0137      PROCCR = PIPT * TNGSP1**GOIMG * PR      00002460
C      SHOCK INLET MACH NUMBER      00002470
ISN 0138      MOI = MO      00002480
C      SHOCK EXIT MACH NUMBER      00002490
ISN 0139      MOE = SQRT(FIOF(MOI) / (GAM * MOI**2 - (GAM - 1.) * .5))      00002500
C      SHOCK INLET TO THROAT STATIC PRESSURE RATIO      00002510
ISN 0140      50 POIPTH = FIOF(MO)**GOIMG / ((GAM + 1.) * .5)**GOIMG      00002520
C      SHOCK EXIT TO INLET STATIC PRESSURE RATIO      00002530
ISN 0141      POEPOI = 1. + GAM * TNGSP1 * (MO**2 - 1.)      00002540
C      STAGE EXIT TO SHOCK EXIT STATIC PRESSURE RATIO      00002550
ISN 0142      PEPE = (FIOF(ME) / FIOF(MOE))**GOIMG / PIIPTI      00002560
C      OBTAINABLE PRESSURE RISE (STAGE EXIT TO THROAT STATIC PRESS. RAT.)      00002570
ISN 0143      PROCT = POIPTH * POEPOI * PEPE      00002580
ISN 0144      IF (ABS(PROCCR/PROCT-1.) .GT. EPS) GO TO 60      00002590
C      DISTANCE FROM REFERENCE BLADE L.E. TO STEADY STATE SHOCK      00002600
ISN 0146      X0 = XIP(INS)      00002610
ISN 0147      XS=X0      00002620
ISN 0148      IF(XSTAR.NE.X0) GOTO 71      00002630
ISN 0150      WRITE(6,75)      00002640
ISN 0151      75 FORMAT(' ', 'THE SHOCK IS LOCATED AT THE THROAT')      00002650
ISN 0152      GOTO 500      00002660
ISN 0153      71 X2P = X0 - TSLAM      00002670
ISN 0154      CALL LINE (NP, XIN, YUIN, X0, YU0, YUOX0, IDUM)      00002680
ISN 0155      CALL LINE (NP, XIN, YLIN, X2P, YL0, YLOX0, IDUM)      00002690
ISN 0156      EO = ZP - X0 / C      00002700
C      STEADY STATE SHOCK AREA      00002710
ISN 0157      AO = ASS(NS)      00002720
ISN 0158      FO = FIOF(MO)**GOIMG * PT      00002730
C      CHECKPOINT OF SHOCK LOCATION      00002740
ISN 0159      IS = NS      00002750
ISN 0160      GO TO 65      00002760
ISN 0161      60 CONTINUE      00002770
ISN 0162      61 IS = 0      00002780
ISN 0163      WRITE (6,62)      00002790
ISN 0164      62 FORMAT(' ***** STEADY STATE SHOCK POSITION NOT FOUND')      00002800
ISN 0165      GO TO 500      00002810
C      00002820
C      CALCULATE CAPTURE MACH NUMBER      00002830
C      00002840
ISN 0166      65 AC = TAU * SIN(BETAIC)      00002850
C      RATIO OF CAPTURE AREA TO THROAT AREA      00002860
ISN 0167      AOAMIN = AC / ASTAR      00002870
ISN 0168      ISS = 0      00002880
ISN 0169      IF (MI .GT. 1.) ISS = 1      00002890
ISN 0171      CALL MACHIT (AOAMIN, MC, GAM, TOLMIT, ISS, KILL)      00002900

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ISN 0172	C	IF (KILL .NE. 0) WRITE (6,66) AOAMIN	00002910
	C	66 FORMAT('0AC/ASTAR = ',IPE13.5)	00002920
	C		00002930
	C	CALCULATE STEADY STATE PRESSURES	00002940
	C		00002950
ISN 0173		PCPT = FIOF(MC)**GOIMG	00002960
ISN 0174		PC = PCPT * PT	00002970
ISN 0175		IF(ICH0KE.EQ.1) GOTO 370	00002980
ISN 0177		CALL LINE (NP,XIN,YUIN,XSTAR,YU,DUM,IDUM)	00002990
ISN 0178		CALL LINE (NP,XIN,YLIN,XSTAR-TSLAM,YL,DUM,IDUM)	00003000
ISN 0179		AREA = TAU * COSLAM - YU + YL	00003010
ISN 0180		AOAMIN = AREA / ASTAR	00003020
ISN 0181		CALL MACHIT (AOAMIN,XMS,GAM,TOLMIT,ISS,KILL)	00003030
ISN 0182		PHIPT = FIOF(XMS) ** COIMG	00003040
	C	WRITE(6,201) XSTAR,TSLAM,YU,YL,AREA,AOAMIN,XMS,PHIPT	00003050
ISN 0183	201	FORMAT('/' XSTAR,TSLAM,YU,YL,AREA,AOAMIN,XMS,PHIPT'/8E15.6)	00003060
ISN 0184		GOTO 380	00003070
ISN 0185	370	PHIPT = ((GAM + 1.) * .5)**COIMG	00003080
ISN 0186	380	PHI = PHIPT * PT	00003090
ISN 0187		PIPT = FIOF(M0)**GOIMG	00003100
ISN 0188		PENT = PIPT * PT	00003110
ISN 0189		PCEPI = 1. + GAM * THCGPI * (M0**2 - 1.)	00003120
	C		00003130
ISN 0190		IF(ICH0KE.EQ.1) GOTO 375	00003140
ISN 0192		CALL LINE (NP,XIN,YUIN,X0,YU,DUM,IDUM)	00003150
ISN 0193		CALL LINE (NP,XIN,YLIN,X0-TSLAM,YL,DUM,IDUM)	00003160
ISN 0194		AREA = TAU * COSLAM - YU + YL	00003170
ISN 0195		AOAMIN = AREA / ASTAR	00003180
ISN 0196		CALL MACHIT (AOAMIN,XM0,GAM,TOLMIT,ISS,KILL)	00003190
ISN 0197		P0E = FIOF(XM0) ** GOIMG * PT	00003200
	C	WRITE(6,202) X0,TSLAM,YU,YL,AREA,AOAMIN,XM0,P0E	00003210
ISN 0198	202	FORMAT('/' X0,TSLAM,YU,YL,AREA,AOAMIN,XM0,P0E'/8E15.6)	00003220
ISN 0199		GOTO 385	00003230
ISN 0200	375	P0E = PCEPI * PENT	00003240
ISN 0201	385	PEPT = FIOF(ME)**GOIMG	00003250
ISN 0202		PE = PEPT * PT	00003260
ISN 0203		P1SS = (PHI + PE) * .5	00003270
ISN 0204		P2SS = (PENT + PH1) * .5	00003280
ISN 0205		IF(ICH0KE.EQ.0) P2SS=(P0E+PM1) * 0.5	00003290
ISN 0207		P3SS = (P0E + PE) * .5	00003300
	C		00003310
	C	CALCULATE STEADY STATE MOMENT	00003320
	C		00003330
ISN 0208		MUSS = -P1SS * XSTAR * (B1PE - XSTAR * .5) - P2SS * (X0 - XSTAR) * (B1PE - (XSTAR + (X0 - XSTAR) * .5)) - P3SS * (C - X0) * (B1PE - (X0 + (C - X0) * .5))	00003340
	1		00003350
	2		00003360
ISN 0209		DCACH = DELTA * COSACH	00003370
ISN 0210		MLSS = P1SS * (XSTAR - DCACH) * (B1PE - (XSTAR - DCACH) * .5) + P2SS * (X0 - XSTAR) * (B1PE - (XSTAR - DCACH + (X0 - XSTAR) * .5)) + P3SS * (C - (X0 - DCACH)) * (B1PE - (C + (X0 - DCACH)) * .5)	00003380
	1		00003390
	2		00003400
	3		00003410
	4		00003420
ISN 0211		MOSS = (MUSS + MLSS) / 12.	00003430
	C		00003440
	C	CALCULATE STEADY STATE LIFT	00003450
	C		00003460
ISN 0212		LUSS = -P1SS * XSTAR - P2SS * (X0 - XSTAR) - P3SS * (C - X0)	00003470
ISN 0213		LLSS = P1SS * (XSTAR - DCACH) + P2SS * (X0 - XSTAR) + P3SS * (C - (X0 - DCACH))	00003480
	1		00003490

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*VERSICH 1.3.0 (01 MAY 80)      MAIN      SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)      DATE 80.353/15.02.47      PAGE 7
ISN 0214      LSS =(LUSS + LLSS)/12.      00003500
ISN 0215      WRITE (6,100) PISS,PCSS,P3SS,MOSS,LSS      00003510
ISN 0216      100 FORMAT(// ' STEADY STATE VALUES'//11X,'P1',10X,'P2',10X,'P3',6X,      00003520
              1 'MOMENT',EX,'LIFT'//1X,3F12.4,F12.5,F12.6)      00003530
              C      00003540
ISN 0217      DUM = PROBT/PRCORR      00003550
ISN 0218      WRITE(6,200) XSTAR,X0,ASTAR,M0,DUM,PTIPTE      00003560
ISN 0219      200 FORMAT( /T5,'XSTAR',T22,'X0',T36,'ASTAR',T52,'M0',T63,      00003570
              ?      'PRCB/PRCORR',T80,'PRLCSS'/6E15.6)      00003580
ISN 0220      WRITE(6,600) MSS      00003590
ISN 0221      600 FORMAT(// ' MSS'/(6E15.6))      00003600
              C      00003610
              C DEFINE STEADY STATE VOLUMES      00003620
              C      00003630
ISN 0222      AC = DELTA*SINACH      00003640
ISN 0223      AE = DELTA*SINACH      00003650
ISN 0224      VIA = 0.0      00003660
ISN 0225      XENT = DELTA*COSACH      00003670
ISN 0226      CALL LINE (NP, XIN, YUIN, XENT, YUENT, DUM, IDUM)      00003680
ISN 0227      VIB = 0.5*DELTA**2*COSACH*SINACH-.5*DELTA*COSACH*YUENT      00003690
ISN 0228      X2P = 0      00003700
ISN 0229      CALL LINE (NP, XIN, YLIN, X2P, YLENT, DUM, IDUM)      00003710
ISN 0230      VIC = (DELTA*SINACH-(YUSTAR+YUENT)/2.0+(YLSTAR+YLENT)/2.0)*      00003720
              ?      (XSTAR-XENT)      00003730
ISN 0231      VOL1 = VIA+VIB+VIC      00003740
ISN 0232      VOL2 = (DELTA*SINACH-(YU0+YUSTAR)/2.0+(YLO+YLSTAR)/2.0)*      00003750
              ?      (X0-XSTAR)      00003760
ISN 0233      YUE = YUIN(NP)      00003770
ISN 0234      V3A = (DELTA*SINACH-(YUE+YU0)/2.0+(YLO+YL)/2.0)*      00003780
              ?      (C-X0)      00003790
ISN 0235      XCD = C-DELTA*COSACH      00003800
ISN 0236      CALL LINE (NP, XIN, YLIN, XCD, YLW, DUM, IDUM)      00003810
ISN 0237      V3B = 0.5*DELTA**2*SINACH*COSACH-.5*DELTA*COSACH*YLW      00003820
ISN 0238      V3C = 0.0      00003830
ISN 0239      VOL3 = V3A+V3B+V3C      00003840
ISN 0240      WRITE(6,605)      00003850
ISN 0241      605 FORMAT(' ',T3,'VIA',T15,'VIB',T27,'VIC',T39,'YUENT',T51,'YUSTAR',      00003860
              ?      T63,'YLSTAR',T75,'YLENT')      00003870
ISN 0242      WRITE(6,610) VIA,VIB,VIC,YUENT,YUSTAR,YLSTAR,YLENT      00003880
ISN 0243      610 FORMAT(' ',8(E11.4,1X))      00003890
              C      00003900
ISN 0244      CALL UNST      00003910
ISN 0245      500 STOP      00003920
ISN 0246      END      00003930

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***** FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL STATEMENT NUMBERS
B	0047 0043 0049
C	0002 0011 0031 0047 0051 0065 0100 0102 0105 0156 0208 0208 0210 0210 0212 0213 0234 0235
E	0002 0011 0031 0048 0049
I	0013 0013 0013 0014 0014 0015 0015 0015 0035 0035 0035 0035 0035
V	0002 0011 0031
AC	0002 0110 0111 0166 0167 0222
AE	0002 0122 0123 0125 0223
A0	0002 0157
EO	0156
IS	0159 0162
IT	0081 0129
MC	0003 0115 0116 0171 0173

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL	STATEMENT	NUMBERS																			
ME	0005	0127	0142	0201																		
MI	0002	0003	0011	0020	0031	0098	0098	0113	0169													
MO	0002	0003	0130	0131	0131	0131	0134	0134	0134	0138	0140	0141	0158	0187	0189	0218						
M1	0003	0011	0020	0020	0031	0092	0092															
MS	0009	0027	0031																			
MP	0009	0013	0014	0015	0031	0035	0070	0071	0154	0155	0177	0178	0192	0193	0226	0229	0233	0236				
MS	0037	0038	0058	0059	0070	0072	0073	0073	0075	0077	0078	0081	0084	0085	0087	0089	0129	0130	0146			
PC	0174	0203																				
PE	0002	0207																				
PI	0042	0044	0052																			
PR	0011	0051	0136	0137																		
PT	0002	0011	0033	0050	0158	0174	0186	0188	0197	0202												
PO	0150																					
TT	0002	0011	0033	0050																		
XS	0002	0147																				
XO	0104	0105	0107	0146	0147	0148	0153	0154	0156	0192	0193	0208	0208	0208	0208	0208	0210	0210	0210			
	0210	0212	0212	0213	0213	0218	0232	0234														
YL	0071	0072	0030	0178	0179	0193	0194	0234														
YU	0070	0072	0079	0177	0179	0192	0194															
ZP	0051	0156																				
ABS	0144																					
ASS	0004	0072	0073	0075	0077	0007	0157															
COS	0053	0059	0062																			
OCR	0037	0030	0037	0040	0041	0110																
DUM	0070	0071	0177	0178	0192	0193	0217	0218	0226	0229	0236											
EPS	0011	0016	0016	0031	0144																	
GAM	0005	0011	0018	0018	0031	0056	0056	0057	0057	0058	0069	0098	0098	0098	0115	0127	0134	0134				
	0134	0139	0139	0140	0141	0171	0181	0185	0189	0196												
ISS	0033	0009	0112	0113	0115	0124	0125	0127	0160	0169	0171	0181	0196									
LSS	0003	0214	0215																			
MSS	0003	0004	0083	0083	0089	0130	0220															
NJE	0002	0003	0139	0142																		
MJI	0002	0003	0139	0139	0139																	
PHI	0186	0203	0204	0205																		
PQE	0177	0200	0205	0207																		
SIN	0054	0050	0061	0063	0166																	
TAU	0026	0033	0055	0064	0072	0092	0092	0098	0110	0166	0179	0194										
VIA	0224	0231	0242																			
VIB	0227	0231	0242																			
VIC	0230	0231	0242																			
VJA	0234	0239																				
VJB	0237	0239																				
VJC	0233	0239																				
XCD	0235	0236																				
XIH	0004	0015	0035	0070	0071	0154	0155	0177	0178	0192	0193	0226	0229	0236								
XMS	0181	0182																				
XND	0196	0197																				
XIP	0004	0058	0069	0070	0078	0085	0146															
XCP	0069	0071	0153	0155	0228	0229																
YLI	0236	0237																				
YLO	0155	0232	0234																			
YUE	0233	0234																				
YUG	0154	0232	0234																			
AREA	0179	0180	0194	0195																		
ASSE	0064																					

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL STATEMENT NUMBERS
BIME	0048
BIFE	0049 0051 0208 0208 0208 0210 0210 0210
DIAM	0011 0027 0031
DXLP	0065 0068
FIOF	0005 0092 0116 0134 0139 0140 0142 0142 0158 0173 0182 0187 0197 0201
HEAR	0002 0011 0031
IDUM	0070 0071 0154 0155 0177 0178 0192 0193 0226 0229 0236
KILL	0009 0115 0127 0171 0181 0196
LINE	0070 0071 0154 0155 0177 0178 0192 0193 0226 0229 0236
LLSS	0003 0213 0214
LUSS	0003 0212 0214
MLSS	0003 0210 0211
MORS	0003 0211 0215
MUES	0003 0208 0211
CMES	0043
FCPT	0173 0174
FENT	0183 0000 0204
FEP1	0136
FEP2	0201 0202
PIPT	0116 0137 0187 0188
PISS	0002 0203 0208 0210 0212 0213 0215
PGSS	0002 0204 0205 0208 0210 0212 0213 0215
PGSS	0002 0207 0208 0210 0212 0213 0215
RGAS	0005 0022 0022 0033 0050
RNOT	0050
SRRT	0092 0139
UNST	0044
VCL1	0002 0231
VCL2	0002 0232
VCL3	0002 0239
XENT	0025 0026 0230
YLIN	0004 0016 0035 0071 0155 0178 0193 0229 0236
YUIN	0004 0015 0035 0070 0154 0177 0192 0226 0233
ZERO	0003
ALPCH	0002 0011 0031 0039 0039 0052 0059 0060
ARSDN	0092
ASTAR	0002 0066 0075 0077 0087 0090 0111 0123 0125 0167 0180 0195 0218
BETA1	0002 0011 0031 0040 0040 0061 0093 0117 0118
BETA2	0002 0011 0031 0041 0041 0062 0063
COSE2	0062
DCACH	0009 0210 0210 0210 0210 0210 0213 0213
DELTA	0002 0011 0026 0031 0099 0100 0102 0122 0209 0222 0223 0225 0227 0227 0230 0232 0234 0235 0237
DNACH	0004
FLOAT	0035 0068
GOING	0056 0116 0134 0137 0140 0140 0142 0158 0173 0182 0185 0187 0197 0201
INAVE	0009 0028 0031
NSECT	0009 0031 0045 0065
NYINE	0031 0046
OMEGA	0002 0011 0031 0043 0044 0044
PEPDE	0142 0143
PHLPT	0182 0185 0186
PROBT	0143 0144 0217
POEPI	0189 0200
SINB1	0061 0098 0110
SINB2	0063 0064 0122

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL	STATEMENT	NUMBERS																
TITLE	0004	0007	0031																
TSLAM	0055	0065	0063	0069	0153	0178	0193												
XIMACH	0004																		
XSTAR	0202	0078	0085	0099	0102	0104	0107	0148	0177	0178	0208	0208	0208	0208	0210	0210	0210		
	0210	0212	0212	0213	0213	0218	0230	0232											
YLENT	0229	0230	0242																
YLOX0	0155																		
YMECH	0004																		
YUENT	0226	0227	0230	0242															
YUDX0	0154																		
ALPDAR	0032	0011	0031	0038	0038														
ADAMIN	0087	0089	0111	0115	0123	0167	0171	0180	0181	0195	0196								
BETAFR	0118	0119																	
BETAIC	0092	0093	0117	0166															
COSACH	0059	0099	0100	0102	0209	0225	0227	0227	0235	0237	0237								
COSLAM	0053	0072	0092	0179	0194														
GPIGH1	0057	0082																	
ICHCKE	0097	0121	0175	0190	0205														
LAWDA	0003	0052	0053	0054															
MACHIT	0089	0115	0127	0171	0181	0196													
MSECT1	0045	0067	0083	0084	0129														
NTIME1	0046																		
PHICD	0002	0027	0028	0028	0030	0031													
PHIBT	0002	0030	0031																
PCCER	0137	0144	0217																
PIIPE	0133	0134	0136	0142	0218														
PSEFGI	0141	0143																	
POIPTH	0140	0143																	
SINACH	0060	0222	0223	0227	0230	0232	0234	0237											
SINLAM	0054	0055																	
TOLMIT	0011	0024	0024	0033	0089	0115	0127	0171	0181	0196									
TNOSP1	0053	0092	0134	0137	0141	0189													
YLSTAR	0080	0092	0230	0232	0242														
YUSTAR	0079	0092	0230	0232	0242														

*****FORTRAN CROSS REFERENCE LISTING*****

LABEL	DEFINED	REFERENCES														
5	0032	0031														
6	0034	0033														
7	0036	0035														
10	0007															
16	0074	0073														
20	0082	0067	0075													
25	0090															
30	0091	0084	0085													
40	0095	0095														
41	0110	0093														
48	0128															
50	0140															
60	0161	0129	0144													
61	0162	0105	0109													
62	0164	0163														
65	0166	0160														
66	0172															
71	0153	0148														
75	0151	0150														

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*****FORTRAN CROSS REFERENCE LISTING*****

LABEL	DEFINED	REFERENCES
80	0033	0007
85	0010	0009
96	0012	0011 0013 0014 0015
109	0216	0215
117	0100	0119
200	0219	0218
201	0133	
202	0193	
330	0104	0100
340	0103	
360	0100	
370	0185	0175
375	0200	0190
390	0185	0184
395	0201	0199
500	0245	0007 0152 0165
600	0221	0220
605	0241	0240
610	0243	0242

MAIN / SIZE OF PROGRAM 00247C HEXADECIMAL BYTES

NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.				
B	SF	R*4	000958	C	SF	C	R*4	000044	E	SF	C	R*4	000014	I	F	I*4	00095C		
V	SF	C	R*4	000074	AC	SF	C	R*4	000018	AE	SF	C	R*4	00000C	AG	S	C	R*4	00001C
EO	S	R*4	000060	IS	S	I*4	000084	IT	SF	I*4	000068	MC	SFA	R*4	000B5C				
HE	SFA	R*4	000070	NI	SF	C	R*4	000043	HO	SFA	C	R*4	000020	MI	SFA	R*4	000B74		
HS	SF	I*4	000073	NP	SFA	I*4	00007C	NS	SFA	I*4	000080	PC	SF	R*4	000B84				
PE	SF	R*4	000003	PI	SF	R*4	00000C	FR	SF	R*4	000050	PT	SF	C	R*4	00004C			
FO	S	R*4	000094	TT	SF	C	R*4	00000C	XS	S	C	R*4	000030	XO	SFA	R*4	000B98		
YL	SFA	R*4	00009C	YU	SFA	R*4	000040	ZP	SF	R*4	000BA4	ASS	SF	R*4	000CF0				
CO3	F	XF	R*4	000000	DCR	SF	R*4	0000A8	DUM	SFA	R*4	000BAC	EPS	SF	R*4	000B80			
GAM	SFA	R*4	000004	ICS	SFA	I*4	000008	LSS	SF	R*4	000B0C	MSS	SFA	R*4	000B9C				
NCE	SFA	C	R*4	000004	MOI	SFA	C	R*4	000008	PMI	SF	R*4	000C00	POE	SF	R*4	000B04		
SIN	F	XF	R*4	000000	TJU	SFA	R*4	000008	VIA	SF	R*4	000B0C	VIB	SF	R*4	000B00			
VIC	SF	R*4	000004	V3A	SF	R*4	000000	V3B	SF	R*4	000B0C	V3C	SF	R*4	000B00				
X00	SFA	R*4	0000E4	XIN	SFA	R*4	0000E8	XIS	SFA	R*4	0000E8	XMO	SFA	R*4	000B0C				
XLP	SFA	R*4	0000F0	X2P	SFA	R*4	0000F0	YLI	SFA	R*4	000EF4	YLO	SFA	R*4	000BF8				
YUF	SF	R*4	0000FC	YUJ	SFA	R*4	000000	APEA	SF	R*4	000004	ASSE	S	R*4	000C08				
DINE	S	R*4	00000C	DIFE	SF	R*4	000010	DIAM	SF	R*4	000014	DXIP	SF	R*4	000C18				
FIOF	AA3F	R*4	000000	FJAR	SF	C	R*4	000064	ICUM	SFA	I*4	00001C	KILL	SFA	I*4	000C20			
LIME	SF	XF	R*4	000000	LLS3	SF	R*4	000004	LUSS	SF	R*4	000028	MLSS	SF	R*4	000C2C			
M003	SF	R*4	000030	M005	SF	R*4	000034	OMCG	S	R*4	000038	FCPT	SF	R*4	000C3C				
PEPT	SF	R*4	000040	PEPI	S	R*4	000044	PEPT	SF	R*4	000048	PIPT	SF	R*4	000C4C				
PL03	SF	C	R*4	000030	PCSS	SF	C	R*4	00003C	P3SS	SF	C	R*4	000040	R6AS	SF	R*4	000C50	
PHCT	S	R*4	000034	SRCT	FA	XF	R*4	000000	UNST	SF	XF	000000	VOLI	S	C	R*4	000C00		
V012	S	C	R*4	000004	V013	S	C	R*4	000008	XENT	SFA	R*4	000058	YLIH	SFA	R*4	00101C		
YUEN	SFA	R*4	0010E4	ZERO	SF	XF	000000	ALPCH	SFA	C	R*4	000058	A3SIN	F	XF	R*4	000000		
ASTAR	SF	C	R*4	000010	BETA1	SFA	C	R*4	000050	BETA2	SFA	C	R*4	000054	COE32	S	R*4	000C5C	
OCACH	SF	R*4	000060	DELTA	SF	C	R*4	000060	DMACH	R*4	NR	R*4	NR	G01MG	SF	R*4	000C64		
INAVE	SF	I*4	000068	NSECT	SFA	I*4	00006C	NTIME	F	I*4	000070	ONEGA	SF	C	R*4	00006C			
PEPOE	SF	R*4	000074	PHIPT	SF	R*4	000078	PPCST	SFA	R*4	00007C	POEPI	SF	R*4	000C80				
SIN31	SF	R*4	000084	SIN32	SF	R*4	00008C	TITLE	SF	R*4	0011AC	TSLAH	SFA	R*4	000C8C				
XHACH	R*4	NR	XRSTAR	SFA	C	R*4	000034	YLENT	SFA	R*4	000090	YLOXO	SFA	R*4	000C94				
YHACH	R*4	NR	YUENT	SFA	R*4	000098	YUOXO	SFA	R*4	00009C	FRXFR#	XF	R*4	000000					
ALPBAR	SF	C	R*4	000070	AOADMIN	SFA	R*4	000CA0	BETAPR	SF	R*4	000CA4	BETA1C	SFA	R*4	000CA8			

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COSACH SF	R*4	00CCAC	COCLAM SFA	R*4	000C80	GPLG11 SFA	R*4	000C84	IBCOM# F XF	I*4	000C00
ICHICE S	I*4	000C23	LAMDA SFA	R*4	000C6C	MACHIT SF XF		000000	NSECT1 SFA	I*4	000C00
ITIME1 S	I*4	000C04	PHIIB3 SF	R*4	00005C	PHIIBT SF C	R*4	000060	FRCORR SFA	R*4	000C08
PIIPIE SF	R*4	000CCC	PEEPCI SF	R*4	000C00	POIPIH SF	R*4	000C04	SINACH SF	R*4	000C08
SINLAM SF	R*4	000C0C	TOLMIT SFA	R*4	000CE0	THGSP1 SFA	R*4	000CE4	YLSTAR SFA	R*4	000C08
YUSTAR SFA	R*4	000CEC									

***** COMMON INFORMATION *****

NAME OF COMMON BLOCK * * SIZE OF BLOCK 000078 HEXADECIMAL BYTES

VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.
VCL1	R*4	000000	VOL2	R*4	000C04	VOL3	R*4	000C08	AE	R*4	00000C
ASTAR	R*4	00C010	E	R*4	000014	AC	R*4	000018	A0	R*4	00001C
MO	R*4	000020	MOE	R*4	000024	MOI	R*4	000028	TT	R*4	00002C
XS	R*4	000030	XSTAR	R*4	000034	PIES	R*4	000038	P2SS	R*4	00003C
PISS	R*4	000040	C	R*4	000044	MI	R*4	000048	PT	R*4	00004C
BETA1	R*4	000050	BETA2	R*4	000054	ALPCH	R*4	000058	PHIIB3	R*4	00005C
PHIIBT	R*4	000060	HBAR	R*4	000064	DELTA	R*4	000068	OMEGA	R*4	00006C
ALFBAR	R*4	000070	V	R*4	000074						

SOURCE STATEMENT LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
10	7	001298 NR	20	82	001C44	30	91	001924	330	104	001A42
41	110	001A64	50	140	001C00 NR	71	153	001D70	60	161	001E0C
61	162	001E16	65	166	001E38	370	185	001F32	380	186	001F54
375	200	002036	335	201	002042	500	245	00244C			

COMPILER GENERATED LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
100000	1	001204	100001	7	0012AC	200001	17	00142C	100008	17	001436
100009	18	00143E	100010	19	001448	100011	20	001450	100012	21	00145A
100013	22	001462	100014	23	00146C	100015	24	001474	100016	25	00147E
100017	26	001485	100018	29	00148A	100019	30	0014C4	100020	35	0015FE
100021	35	00161A	100022	68	0017F4	100023	77	001830	100024	83	0018DC
100025	85	0018F2	100026	87	0018FE	100027	92	001928	100028	95	001945
100029	102	001A3E	100030	107	001A54	100031	114	001A88	100032	115	001A8C
100033	123	001D1E	100034	127	001D22	100035	130	001DDE	200002	132	001E5E
100036	132	001DF6	100037	133	001DFE	100038	135	001C12	100039	136	001C76
100040	146	001D52	100041	150	001D76	100043	170	001E66	100044	171	001E6A
100045	177	001EE6	100046	192	001FE6	100047	206	0020A2	100048	207	0020B2
200003	217	0021F8	200004	235	00237E						

FORMAT STATEMENT LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
60	8	000C28	85	10	000C2E	86	12	000034	5	32	00003B
6	34	00046F	7	35	000537	16	74	0005CB	25	90	0005D6 NR
40	55	0005F4	340	103	000511 NR	360	108	00052D NR	117	120	000646
45	123	000656 NR	75	151	000675	62	164	00069E	66	172	0006CE NR
201	183	0006E3 NR	202	193	000716 NR	100	216	000744	200	219	000793
600	221	0007D3	605	241	0007E6	610	243	000826			

NUMBER LEVEL FORTRAN H EXTENDED ERROR MESSAGES

IFE029I 8(E) ISN 0032 THE NUMBER OF CONTINUATION CARDS EXCEEDS 19. COMPILER PROCESSING OF THE STATEMENT CONTINUES.

IFE0226I 8(E) ISN 0119 THE STATEMENT HAS A VARIABLE WITH MORE THAN SIX CHARACTERS. THE RIGHTMOST CHARACTERS ARE TRUNCATED.

*OPTIONS IN EFFECT*NAME(MAIN) OPTIMIZE(2) LINECOUNT(60) SIZE(MAX) AUTOCDBL(NONE)
 *OPTIONS IN EFFECT*SOURCE EBCDIC NOLIST NOCHECK OBJECT MAP NOFORMAT GOSTMT XREF NOALC NOANSF TERM IBM FLAG(I)

C	DATA SET 936GUNST	AT LEVEL 002 AS OF 08/15/80	00000000
C	DATA SET 936GUNST	AT LEVEL 001 AS OF 07/30/80	00000010
C	DATA SET 936GUNST	AT LEVEL 001 AS OF 04/21/80	00000020
ISN 0002	SUBROUTINE UNST		00000030
C	DECK 9066	CHANNEL FLOW CHOKE FLUTTER UNSTEADY AERODYNAMIC MODEL	00000040
C			00000050
C			00000060
C	ANALYST	J. F. SIMPSON	00000070
C	DATE	JANUARY, 1980	00000080
C	EXT	4315	00000090
C	MAIL LOC	R-47, BLDG 32, P5G	00000100
C			00000110
C			00000120
C			00000130
C		-----DESCRIPTIONS-----	00000140
C			00000150
C	B	SENTECHORD OF THE AIRFOIL	00000160
C	DI	DENSITY AT THE INLET	00000170
C	DD	DENSITY DOWNSTREAM	00000180
C	DTOT	TOTAL DENSITY	00000190
C	DU	DENSITY UPSTREAM	00000200
C	E	ELASTIC AXIS POSITION REFERENCED TO MIDCHORD	00000210
C	FLOW	S.S FLOW RATE ENTERING THE SHOCK	00000220
C	FRE	S.S FLOW RATE AT THE EXIT	00000230
C	GAMMA	SPECIFIC HEAT RATIO	00000240
C	MEXIT	MACH NUMBER AT THE EXIT	00000250
C	MDS	MACH NUMBER DOWNSTREAM	00000260
C	MI	MACH NUMBER AT THE INLET	00000270
C	MSHOCK	MACH NUMBER AT THE SHOCK	00000280
C	MT	MACH NUMBER AT THE THROAT	00000290
C	MUS	MACH NUMBER UPSTREAM	00000300
C	OMEGA	FREQUENCY OF VIBRATIONS	00000310
C	PE	PRESSURE AT THE EXIT	00000320
C	PD	PRESSURE DOWNSTREAM	00000330
C	PIN	PRESSURE AT THE INLET	00000340
C	PIOT	TOTAL PRESSURE	00000350
C	PU	PRESSURE UPSTREAM	00000360
C	PGAS	GAS CONSTANT	00000370
C	SIGMAB	BENDING MODE INTERBLADE PHASE ANGLE	00000380
C	SIGMAT	TORSIONAL MODE INTERBLADE PHASE ANGLE	00000390
C	SOEDS	SPEED OF SOUND DOWNSTREAM	00000400
C	SOSEX	SPEED OF SOUND AT THE EXIT	00000410
C	SSI	SPEED OF SOUND AT THE INLET	00000420
C	SST	SPEED OF SOUND AT THE THROAT	00000430
C	SEU	SPEED OF SOUND UPSTREAM	00000440
C	TDS	TEMPERATURE DOWNSTREAM	00000450
C	TE	TEMPERATURE AT THE EXIT	00000460
C	TI	TEMPERATURE AT THE INLET	00000470
C	UE	AXIAL VELOCITY AT THE EXIT	00000480
C	VE	TANGENTIAL VELOCITY AT THE EXIT	00000490
C	VU	VELOCITY UPSTREAM	00000500
C			00000510
C			00000520
C			00000530
ISN 0003	DIMENSION FE(35), FM(47), FC(57), FM2(18), M(28),		00000540

*VERSION 1.3.0 (01 MAY 80)		UNST SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)					DATE 80.353/15.02.89	PAGE 2
	?	CL(39),	S(9),	LH(24),	LC(28),	AAR(2),	00000550	
	?	AAI(2),	BAR(2),	BAI(2),	AHR(2),	AHI(2),	00000560	
	?	BHR(2),	CHR(2)				00000570	
	C						00000580	
ISN 0004	REAL	K,	IC,	LC,	LM,	MU,	MD,	
	?	MDS,	MUS,	MXD,	MXU,	MYD,	MYU,	
	?	IMAP,	IKDP,	INPP,	IMSP,	IRFF,	ISSE,	
	?	ICDS,	IVFD,	MCPT,	MIRE,	M2RE,	IMAPI,	
	?	IMAPT,	IMDFU,	IMPPU,	INVPU,	ISOSU,	MEXIT,	
	?	MSHOCK,	MT,	MI,	M,	I,	IACE,	
	?	IVPU,	IDDS,	IFDS,	MXSTAR,			
	?	MYU,	MYD,	IMPPI,	IMDPI,	IMSSPI,	IITVPI,	
	?	IIAVPI,	IMPPE,	IMSPE,	IMSSPE,	IIAVPE,	IITVPE,	
	?	IRAVPE,	IRTVPE,	IP1U,	IPC1U,	IP3U,	IP1L,	
	?	IP2L,	IP3L,	ICLU,	ICLL,	ICL,	ICPU,	
	?	ICML,	ICM				00000700	
	C						00000710	
ISN 0095	COMMON	VOL1,	VOL2,	VOL3,	ACE,	ASTAR,	E,	
	?	AC,	AO,	MSHOCK,	MDS,	MUS,	TTOT,	
	?	XS,	XSTAR,	PRES1,	PRES2,	PRES3,	C,	
	?	MIRE,	PTOT,	BETA1,	BETA2,			
	?	ALPCH,	SIGMAB,	SIGMAT,	H,	DELTA,		
	?	OMEGA,	ALPBAR,	UIIRE			00000770	
	C						00000780	
ISN 0006	READ(5,1000)	M2RE					00000790	
ISN 0007	MEXIT =	M2RE*COS(ALPCH - BETA2)					00000800	
ISN 0008	1000	FORMAT(F10.0)					00000810	
	C						00000820	
	C						00000830	
ISN 0009	B	=	C/2.0				00000840	
	C						00000850	
	C	UNDIMENSIONALIZE	TAU				00000860	
ISN 0010	TAU	=	DELTA				00000870	
ISN 0011	TAU	=	TAU/B				00000880	
	C						00000890	
	C	DEFINE	QUANTITIES	USED	FREQUENTLY	IN	CALCULATIONS	
	C						00000900	
ISN 0012	PI	=	3.141592653589800				00000920	
ISN 0013	MT	=	1.0				00000930	
ISN 0014	MI	=	MIRE				00000940	
ISN 0015	GAMMA	=	1.4				00000950	
ISN 0016	MXSTAR	=	1.0				00000960	
ISN 0017	Z	=	B*(1+E)/C				00000970	
ISN 0018	RGAS	=	1716.26				00000980	
ISN 0019	DTOT	=	PTOT/(RGAS*TTOT)*144.0				00000990	
ISN 0020	I	=	COS(ALPCH-BETA1)				00001000	
ISN 0021	K	=	C/2*OMEGA/UIIRE/12.0				00001010	
ISN 0022	WRITE(6,95)	K					00001020	
ISN 0023	95	FORMAT(//	K'/E15.6)				00001030	
ISN 0024	FU	=	PTOT*(1+(GAMMA-1)/2*MSHOCK**2)**(-GAMMA/(GAMMA-1))				00001040	
ISN 0025	PT	=	PTOT*(1+(GAMMA-1)/2*MT**2)**(-GAMMA/(GAMMA-1))				00001050	
ISN 0026	PIN	=	PTOT*(1+(GAMMA-1)/2*MI**2)**(-GAMMA/(GAMMA-1))				00001060	
ISN 0027	TIN	=	TTOT/(1+(GAMMA-1)/2*MI**2)				00001070	
ISN 0028	DI	=	PIN/(RGAS*TIN)*144				00001080	
ISN 0029	PD	=	PU*(1+2*GAMMA/(GAMMA+1))*(MUS**2-1)				00001090	
ISN 0030	TDS	=	TTOT/(1+(GAMMA-1)/2*MDS**2)				00001100	
ISN 0031	DD	=	PD/(RGAS*TDS)*144				00001110	
ISN 0032	DU	=	DTOT*(1+(GAMMA-1)/2*MSHOCK**2)**(-1/(GAMMA-1))				00001120	
	C						00001130	

ISN 0033	C	CO	=	SIGMA8/TAU	00001140
ISN 0034		MD	=	M2RE*COS(ALPCH-BETA2)	00001150
ISN 0035		MXU	=	M1RE*COS(ALPCH-BETA1)	00001160
ISN 0036		MXD	=	M2RE*COS(ALPCH-BETA2)*SIN(ALPCH)	00001170
ISN 0037		MYU	=	M1RE*COS(ALPCH-BETA1)*SIN(ALPCH)	00001180
ISN 0038		MYD	=	M2RE*COS(ALPCH-BETA2)*COS(ALPCH)	00001190
ISN 0039		MYU	=	M1RE*COS(ALPCH-BETA2)*COS(ALPCH)	00001200
ISN 0040		BXU	=	1-MXU**2	00001210
ISN 0041		BXD	=	1-MXD**2	00001220
ISN 0042		BYU	=	1-MYU**2	00001230
ISN 0043		BYD	=	1-MYD**2	00001240
ISN 0044		D1U	=	MXU*MYU*K+MXU*MYU*CO	00001250
ISN 0045		D1D	=	MXD*MYD*K+MXD*MYD*CO	00001260
ISN 0046		D2U	=	K**2*MXU**2+2*MXU*MYU*CO*K-CO**2*BYU	00001270
ISN 0047		D2D	=	K**2*MXD**2+2*MXD*MYD*CO*K-CO**2*BYD	00001280
ISN 0048	C	QUANT	=	D1U**2+BXU*D2U	00001290
ISN 0049		IF(QUANT.GT.0)	GOTO	80	00001300
ISN 0051		WRITE(6,71)			00001310
ISN 0052	71	FORMAT(' ', 'UPSTREAM SOLUTION NUMBER 1')			00001320
ISN 0053		B1I	=	-SQRT(-(D1U**2+BXU*D2U))/BXU	00001330
ISN 0054		B1R	=	D1U/BXU	00001340
ISN 0055		GOTO	90		00001350
ISN 0056	80	WRITE(6,75)			00001360
ISN 0057	75	FORMAT(' ', 'UPSTREAM SOLUTION NUMBER 2')			00001370
ISN 0058		QUANT	=	MU*K+MYU*CO	00001380
ISN 0059		IF(QUANT.GT.0)	GOTO	84	00001390
ISN 0061		B1R	=	(D1U-SQRT(D1U**2+BXU*D2U))/BXU	00001400
ISN 0062		B1I	=	(MYU*CO+MU*K)/MXU	00001410
ISN 0063		GOTO	90		00001420
ISN 0064	84	B1R	=	(D1U+SQRT(D1U**2+BXU*D2U))/BXU	00001430
ISN 0065		B1I	=	-(MYU*CO+MU*K)/MXU	00001440
ISN 0066	C	QUANT	=	D1D**2+BXD*D2D	00001450
ISN 0067		IF(QUANT.GT.0)	GOTO	91	00001460
ISN 0069		WRITE(6,76)			00001470
ISN 0070	76	FORMAT(' ', 'DOWNSTREAM SOLUTION NUMBER 1')			00001480
ISN 0071		B2I	=	SQRT(-(D1D**2+BXD*D2D))/BXD	00001490
ISN 0072		B2R	=	D1D/BXD	00001500
ISN 0073		GOTO	99		00001510
ISN 0074	91	QUANT	=	MD*K+MYD*CO	00001520
ISN 0075		WRITE(6,78)			00001530
ISN 0076	78	FORMAT(' ', 'DOWNSTREAM SOLUTION NUMBER 2')			00001540
ISN 0077		IF(QUANT.GT.0)	GOTO	94	00001550
ISN 0079		B2R	=	(D1D+SQRT(D1D**2+BXD*D2D))/BXD	00001560
ISN 0080		B2I	=	-(MYD*CO+MD*K)/MYD	00001570
ISN 0081		GOTO	99		00001580
ISN 0082	94	B2R	=	(D1D-SQRT(D1D**2+BXD*D2D))/BXD	00001590
ISN 0083		B2I	=	(MYD*CO+MD*K)/MYD	00001600
ISN 0084	C	K	=	K * 12	00001610
ISN 0085		PE	=	PTOT*(1+(GAMMA-1)/2*MEXIT**2)**(-GAMMA/(GAMMA-1))	00001620
ISN 0086		TE	=	TTOT/(1+(GAMMA-1)/2*MEXIT**2)	00001630
ISN 0087		DE	=	PE/(RGAS*TE)*144	00001640
ISN 0088		SOSEX	=	SQRT(GAMMA*RGAS*TTOT/(1+(GAMMA-1)/2*MEXIT**2))	00001650
ISN 0089		REV	=	M2RE * SOSEX	00001660
ISN 0090		VE	=	REV*COS(ALPCH-BETA2)*COS(ALPCH)	00001670
ISN 0091		UE	=	REV*COS(ALPCH-BETA2)*SIN(ALPCH)	00001680

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ISN 0092 ULIRE = REV 00001730
ISN 0093 US = 0.0 00001740
ISN 0094 SSI = SQRT(GAMMA*RGAS*TTOT/(1+(GAMMA-1)/2*MI**2)) 00001750
ISN 0095 SST = SQRT(GAMMA*RGAS*TTOT/(1+(GAMMA-1)/2*MXSTAR**2)) 00001760
ISN 0096 SOGDS = SQRT(GAMMA*RGAS*TTOT/(1+(GAMMA-1)/2*MDS**2)) 00001770
ISN 0097 SSU = SQRT(GAMMA*RGAS*TTOT/(1+(GAMMA-1)/2*NUS**2)) 00001780
ISN 0098 VU = MUS*SSU 00001790
ISN 0099 VD = MDS*SOGDS 00001800
ISN 0100 WDS = VD*AD*DD/(DI*ULIRE*ASTAR) 00001810
ISN 0101 AVGSOS = (SSI+SST)/2.0 00001820
ISN 0102 DT = DTOT*(1+(GAMMA-1)/2*MXSTAR**2)**(-1/(GAMMA-1)) 00001830
ISN 0103 AVGD1 = (DI+DT)/2.0 00001840
ISN 0104 AVGD2 = (DT+DU)/2.0 00001850
ISN 0105 AVGD3 = (DD+DE)/2.0 00001860
ISN 0106 AVGV1 = (ULIRE*COSS(ALPCH-BETA1)+SST)/2.0 00001870
ISN 0107 AVGV2 = (SST+VU)/2.0 00001880
ISN 0108 AVGV3 = (VD+REV*COSS(ALPCH-BETA2))/2.0 00001890
ISN 0109 REV = REV/ULIRE 00001900
ISN 0110 VELAX = ULIRE*COSS(ALPCH-BETA1)*SIN(ALPCH) 00001910
ISN 0111 VEL = ULIRE*COSS(ALPCH-BETA1)*COS(ALPCH) 00001920
C 00001930
ISN 0112 WRITE(6,100) 00001940
ISN 0113 100 FORMAT('1',T37,'INLET',T51,'THROAT',T66,'UPSTREAM',T81,
? 'DOWNSTREAM',T96,'EXIT') 00001950
ISN 0114 WRITE(6,105) MI,NT,MUS,MDS,MEXIT 00001960
ISN 0115 WRITE(6,115) PIN,PT,FU,FD,FE 00001970
ISN 0116 WRITE(6,110) DI,DT,DU,DD,DE 00001980
ISN 0117 WRITE(6,120) SSI,SST,SSU,SOGDS,SOSEX 00002000
ISN 0118 105 FORMAT('0','MACH NUMBER',T32,G14.7,T47,F14.7,T62,F14.7,T77,F14.7,
? T92,F14.7) 00002010
ISN 0119 110 FORMAT('0','DENSITY',T32,F14.7,T47,F14.7,T62,F14.7,T77,F14.7,
? T92,F14.7) 00002030
ISN 0120 115 FORMAT('0','PRESSURE',T32,F14.7,T50,F14.7,T62,F14.7,T77,F14.7,
? T92,F14.7) 00002050
ISN 0121 120 FORMAT('0','SPEED OF SOUND',T32,F14.7,T47,F14.7,T62,F14.7,T77,
? F14.7,T92,F14.7) 00002070
C 00002080
C 00002090
C NONDIMENSIONALIZE ALL QUANTITIES 00002100
C 00002110
C DIVIDE ALL THE DENSITIES BY THE INLET DENSITY 00002120
C 00002130
ISN 0122 DU = DU/DI 00002140
ISN 0123 DD = DD/DI 00002150
ISN 0124 DE = DE/DI 00002160
ISN 0125 DT = DT/DI 00002170
ISN 0126 DTOT = DTOT/DI 00002180
ISN 0127 AVGD1 = AVGD1/DI 00002190
ISN 0128 AVGD2 = AVGD2/DI 00002200
ISN 0129 AVGD3 = AVGD3/DI 00002210
C 00002220
C PUT PRESSURE IN P-S-F UNITS AND DIVIDE BY (INLET DENSITY*INLET
C VELOCITY) 00002230
C 00002240
C 00002250
ISN 0130 PIN = PIN*144.0/(DI*ULIRE**2) 00002260
ISN 0131 PU = PU*144.0/(DI*ULIRE**2) 00002270
ISN 0132 PD = PD*144.0/(DI*ULIRE**2) 00002280
ISN 0133 PE = PE*144.0/(DI*ULIRE**2) 00002290
ISN 0134 PT = PT*144.0/(DI*ULIRE**2) 00002300
ISN 0135 PTOT = PTOT*144.0/(DI*ULIRE**2) 00002310

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ISH 0136	PRES1 =	FRES1*144.0/(DI*UIRE**2)	00002320
ISH 0137	PRES2 =	FRES2*144.0/(DI*UIRE**2)	00002330
ISH 0138	PRES3 =	FRES3*144.0/(DI*UIRE**2)	00002340
	C		00002350
	C	DIVIDE ALL VELOCITIES AND THE SPEED OF SOUNDS BY THE INLET VELOCITY	00002360
	C		00002370
ISH 0139	SSI =	SSI/UIRE	00002380
ISH 0140	SST =	SST/UIRE	00002390
ISH 0141	SOSEX =	SOSEX/UIRE	00002400
ISH 0142	SCSDS =	SCSDS/UIRE	00002410
ISH 0143	SSU =	SSU/UIRE	00002420
ISH 0144	AVGSOS =	AVGSOS/UIRE	00002430
ISH 0145	VE =	VE/UIRE	00002440
ISH 0146	UE =	UE/UIRE	00002450
ISH 0147	VU =	VU/UIRE	00002460
ISH 0148	VD =	VD/UIRE	00002470
ISH 0149	AVGV1 =	AVGV1/UIRE	00002480
ISH 0150	AVGV2 =	AVGV2/UIRE	00002490
ISH 0151	AVGV3 =	AVGV3/UIRE	00002500
ISH 0152	VELAX =	VELAX/UIRE	00002510
ISH 0153	VEL =	VEL/UIRE	00002520
ISH 0154	U2IRE =	U2IRE/UIRE	00002530
	C		00002540
	C	UNDIMENSIONALIZE THE VOLUMES	00002550
	C		00002560
ISH 0155	VOL1 =	VOL1/(ASTAR*B)	00002570
ISH 0156	VOL2 =	VOL2/(ASTAR*B)	00002580
ISH 0157	VOL3 =	VOL3/(ASTAR*B)	00002590
	C		00002600
	C	UNDIMENSIONALIZE THE AREAS	00002610
	C		00002620
ISH 0158	ACE =	ACE/ASTAR	00002630
ISH 0159	AC =	AC/ASTAR	00002640
ISH 0160	AO =	AO/ASTAR	00002650
	C		00002660
	C		00002670
	C	INTERBLADE ANALYSIS OF SECTION 1	00002680
	C		00002690
	C		00002700
ISH 0161	FLOW =	DU*AO*VU	00002710
ISH 0162	FRE =	DE*ACE*U2IRE	00002720
ISH 0163	FE(1) =	SSI**2/(GAMMA*(GAMMA-1))+((COS(ALPCH-BETA1))**2)/2	00002730
	C		00002740
ISH 0164	FE(2) =	(GAMMA**2-GAMMA+2)/(2*(GAMMA*(GAMMA-1)))*SST**2	00002750
	C		00002760
ISH 0165	FE(3) =	-FE(1)*I*AC*VELAX/(CO*TAU+SSI**2)+FE(1)*AC	00002770
	?	*SIN(ALPCH)/(CO*TAU)-2*AC*I*VELAX/(2*GAMMA*CO*	00002780
	?	TAU)+AC*I**2*SIN(ALPCH)/(CO*TAU)	00002790
	C		00002800
ISH 0166	FE(4) =	-FE(1)*I*AC*(K+CO*VEL)/(CO*TAU*SST**2)+(FE(1)*AC	00002810
	?	*COS(ALPCH))/TAU-2*AC*I*(K+CO*VEL)/(2*GAMMA	00002820
	?	*CO*TAU)+(AC*I**2*COS(ALPCH))/TAU	00002830
	C		00002840
ISH 0167	FE(5) =	(-VOL1*AVGSOS**2*VELAX)/(2*GAMMA*(GAMMA-1)*CO*TAU	00002850
	?	*SSI**2)-AVGSOS*AVGD1*VOL1/GAMMA/(2*SSI*CO*TAU)	00002860
	?	*VELAX-(VOL1*AVGV1**2*VELAX)/(4*CO*TAU*SSI**2)	00002870
	?	+(AVGV1*AVGD1*VOL1*SIN(ALPCH))/(2*CO*TAU)	00002880
	C		00002890
ISH 0168	FE(6) =	(-VOL1*AVGSOS**2*(K+CO*VEL))/(2*GAMMA*(GAMMA-1)	00002900

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		?	*CO*TAU*SSI**2)-AVGSS05*AVGD1*VOLL/GAMMA/(2*SSI	00002910	
		?	*CO*TAU)*(K+CO*VEL)+(VOLL*AVGV1**2)/4*(K+CO*VEL)/	00002920	
		?	(CO*TAU*SSI**2)+AVGV1*AVGD1*VOLL*COS(ALPCH)/(2*TAU)	00002930	
ISN 0169	C	FE(7) =	COS(CO*TAU)-1	00002940	
ISN 0170	C	FE(8) =	FE(3)*(FE(7)*B1R-B1I*SIN(CO*TAU))+FE(4)*FE(7)+FE(5)	00002950	
		?	+K*(B1I+FE(7)+B1R*SIN(CO*TAU))+FE(6)*K*SIN(CO*TAU)	00002960	
ISN 0171	C	FE(9) =	-FE(3)*(FE(7)*B1I+B1P*STN(CO*TAU))-FE(4)*SIN(CO*TAU)+	00002990	
		?	FE(5)*K*(B1R*FE(7)-B1I*SIN(CO*TAU))+FE(6)*K*FE(7)	00003000	
ISN 0172	C	FE(10) =	FE(3)*(B1I*FE(7)+B1R*SIN(CO*TAU))+FE(4)*SIN(CO*	00003010	
		?	TAU)-FE(5)*(K*B1R*FE(7)-K*B1I*SIN(CO*TAU))	00003020	
		?	-FE(6)*K*FE(7)	00003030	
ISN 0173	C	FE(11) =	FE(3)*(B1R*FE(7)-B1I*SIN(CO*TAU))+FE(4)*FE(7)	00003040	
		?	+FE(5)*(K*B1I*FE(7)+K*B1R*SIN(CO*TAU))+FE(6)	00003050	
		?	*K*SIN(CO*TAU)	00003060	
ISN 0174	C	FE(12) =	VOLL * AVGSS05**2* K/12 / (2*GAMMA*(GAMMA-1)) +	00003070	
		?	VOLL * AVGV1**2 * K/12 / 4	00003080	
ISN 0175	C	FE(13) =	K/12 * AVGD1 * AVGSS05**2 / (GAMMA * (GAMMA-1)) +	00003090	
		?	K/12 * AVGD1 * AVGV1**2 / 2	00003100	
ISN 0176	C	FE(14) =	AVGSS05 * AVGD1 * AVGV1 / (GAMMA * (GAMMA - 1)) *	00003110	
		?	K/12 + AVGV1 * AVGD1 * VOLL * K/12 / 2	00003120	
ISN 0177	C	FE(15) =	FE(9)/FE(8)-FE(11)/FE(10)	00003130	
ISN 0178	C	FE(16) =	FE(1)/FE(8)*I/(FE(9)/FE(8)/FE(15)-1)	00003140	
ISN 0179	C	FE(17) =	-FE(8)/FE(9)*FE(1)/FE(10)*I/FE(15)	00003150	
ISN 0180	C	FE(18) =	FE(2)/FE(8)*(SST-SST/FE(15)*FE(9)/FE(8))-FE(12)/	00003160	
		?	(FE(10)*FE(15))*FE(9)/FE(8)	00003170	
ISN 0181	C	FE(19) =	-FE(12)/FE(8)+FE(2)*SST/(FE(10)*FE(15))*FE(9)/FE(8)+	00003180	
		?	FE(12)/(FE(8)*FE(15))*FE(9)/FE(8)	00003190	
ISN 0182	C	FE(20) =	FE(2)/FE(8)*DT*SST-FE(2)*DT*SST/(FE(8)*FE(15))*	00003200	
		?	FE(9)/FE(8)	00003210	
ISN 0183	C	FE(21) =	FE(2)*DT*SST/(FE(10)*FE(15))*FE(9)/FE(8)	00003220	
ISN 0184	C	FE(22) =	3*FE(2)/FE(8)*(DT-DT/FE(15))*FE(9)/FE(8)-FE(14)/	00003230	
		?	(FE(10)*FE(15))-FE(9)/FE(8)	00003240	
ISN 0185	C	FE(23) =	-FE(14)/FE(8)+(FE(2)*3*DT*FE(9)/(FE(10)*FE(15)*FE(8)))+	00003250	
		?	FE(14)*FE(9)/(FE(8)**2*FE(15))	00003260	
ISN 0186	C	FE(24) =	-FE(13)/FE(8)+FE(13)/(FE(8)*FE(15))*FE(9)/FE(8)	00003270	
ISN 0187	C	FE(25) =	-FE(13)/(FE(10)*FE(15))*FE(9)/FE(8)	00003280	
ISN 0188	C	FE(26) =	-FE(11)*I/(FE(8)*FE(15))	00003290	
ISN 0189	C	FE(27) =	FE(1)/(FE(10)*FE(15))	00003300	
				00003310	
				00003320	
				00003330	
				00003340	
				00003350	
				00003360	
				00003370	
				00003380	
				00003390	
				00003400	
				00003410	
				00003420	
				00003430	
				00003440	
				00003450	
				00003460	
				00003470	
				00003480	
				00003490	

ISH 0190	C	FE(28) =	FE(2)*SST/(FE(8)*FE(15))+FE(12)/(FE(10)*FE(15))	00003500 00003510
ISN 0191	C	FE(29) =	-FE(2)*SST/(FE(10)*FE(15))-FE(12)/(FE(8)*FE(15))	00003520 00003530
ISH 0192	C	FE(30) =	FE(2)*SST*DT/(FE(8)*FE(15))	00003540 00003550
ISN 0193	C	FE(31) =	-FE(2)*SST*DT/(FE(10)*FE(15))	00003560 00003570
ISH 0194	C	FE(32) =	3*DT*FE(2)/(FE(8)*FE(15))+FE(14)/(FE(10)*FE(15))	00003580 00003590
ISN 0195	C	FE(33) =	-3*DT*FE(2)/(FE(10)*FE(15))-FE(14)/(FE(8)*FE(15))	00003600 00003610
ISH 0196	C	FE(34) =	-FE(13)/(FE(8)*FE(15))	00003620 00003630
ISN 0197	C	FE(35) =	FE(13)/(FE(10)*FE(15))	00003640 00003650
ISH 0198	C	FM(1) =	-VELAX*AC/(CO*TAU)-VELAX*AC*I**2/(CO*TAU*SSI**2)+ 2*AC*SIN(ALPCH)*I/(CO*TAU)	00003660 00003670 00003680
ISN 0199	C	FM(2) =	-(K+CO*VEL)*AC/(CO*TAU)-(K+CO*VEL)*AC*I**2/(CO*TAU* SSI**2)+2*AC*COS(ALPCH)*I/TAU	00003690 00003700 00003710
ISH 0200	C	FM(3) =	AVGV1*VOLL*VELAX/(2*CO*TAU*SSI**2)-AVGD1*VOLL* SIN(ALPCH)/(2*CO*TAU)	00003720 00003730 00003740
ISN 0201	C	FM(4) =	(K+CO*VEL)/(CO*TAU*SSI**2)*AVGV1*VOLL/2-AVGD1*VOLL/2* COS(ALPCH)/TAU	00003750 00003760 00003770
ISH 0202	C	FM(5) =	-PIN-I**2	00003780 00003790
ISN 0203	C	FM(6) =	PT+SST**2*DT	00003800 00003810
ISH 0204	C	FM(7) =	FM(1)*(B1I*FE(7)-B1I*SIN(CO*TAU))+FM(2)*FE(7)-FM(3)* (K*B1I*FE(7)-K*B1R*SIN(CO*TAU))-FM(4)*K*SIN(CO*TAU)	00003820 00003830 00003840
ISN 0205	C	FM(8) =	-FM(1)*(B1I*FE(7)+B1R*SIN(CO*TAU))-FM(2)*SIN(CO*TAU)+ FM(3)*K*(-B1R*FE(7)+B1I*SIN(CO*TAU))-FM(4)*K*FE(7)	00003850 00003860 00003870 00003880
ISH 0206	C	FM(9) =	FM(1)*(B1I*FE(7)+B1R*SIN(CO*TAU))+FM(2)*SIN(CO*TAU)+ FM(3)*K*(B1R*FE(7)-B1I*SIN(CO*TAU))+FM(4)*FE(7)*K	00003890 00003900 00003910
ISN 0207	C	FM(10) =	FM(1)*(B1R*FE(7)-B1I*SIN(CO*TAU))+FM(2)*FE(7)- FM(3)*K*(B1I*FE(7)+B1R*SIN(CO*TAU))-FM(4)*K* SIN(CO*TAU)	00003920 00003930 00003940
ISH 0208	C	FM(11) =	FE(18)*FM(7)+FM(8)*FE(28)-2*SST**2	00003950 00003960
ISN 0209	C	FM(12) =	FM(7) * FE(19) + FM(8) * FE(29) + AVGV1 * VOLL * K/12 /2	00003970 00003980 00003990
ISH 0210	C	FM(13) =	-FM(7)*FE(16)-FM(8)*FE(26)+FM(5)	00004000 00004010
ISN 0211	C	FM(14) =	-FM(7)*FE(17)-FM(8)*FE(27)	00004020 00004030
ISH 0212	C	FM(15) =	-FM(7)*FE(20)-FM(8)*FE(30)+FM(6)	00004040 00004050
ISN 0213	C	FM(16) =	-FM(7)*FE(21)-FM(8)*FE(31)	00004060 00004070 00004080

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ISN 0214	C	FM(17) =	-FM(7)*FE(22)-FM(8)*FE(32)+2*SST*DT	00004090	
ISN 0215	C	FM(18) =	-FM(7) * FE(23) - FM(8) * FE(33) - AVGD1 * VOL1 * K/12 / 2	00004100 00004110 00004120 00004130	
ISN 0216	C	FM(19) =	-FM(7) * FE(24) - FM(8) * FE(34) - AVGD1 * AVGV1 * K/12	00004140 00004150	
ISN 0217	C	FM(20) =	-FM(7)*FE(25)-FM(8)*FE(35)	00004160 00004170 00004180	
ISN 0218	C	FM(21) =	FM(9) * FE(18) + FM(10) * FE(28) - AVGV1 * VOL1 * K/12 / 2	00004190 00004200 00004210	
ISN 0219	C	FM(22) =	FM(9)*FE(19)+FM(10)*FE(29)-2*SST**2	00004220 00004230	
ISN 0220	C	FM(23) =	-FM(9)*FE(16)-FM(10)*FE(26)+FM(5)	00004240 00004250	
ISN 0221	C	FM(24) =	-FE(17)*FM(9)-FM(10)*FE(27)+FM(6)	00004260 00004270	
ISN 0222	C	FM(25) =	-FM(9)*FE(20)-FM(10)*FE(30)	00004280 00004290	
ISN 0223	C	FM(26) =	-FM(9)*FE(21)-FM(10)*FE(31)	00004300 00004310	
ISN 0224	C	FM(27) =	-FM(9) * FE(22) - FM(10) * FE(32) + AVGD1 * VOL1 * K/12 / 2	00004320 00004330 00004340	
ISN 0225	C	FM(28) =	-FM(9)*FE(23)-FM(10)*FE(33)+2*SST*DT	00004350 00004360	
ISN 0226	C	FM(29) =	-FM(9)*FE(24)-FM(10)*FE(34)	00004370 00004380	
ISN 0227	C	FM(30) =	-FM(9) * FE(25) - FM(10) * FE(35) + AVGD1 * AVGV1 * K/12	00004390 00004400 00004410	
ISN 0228	C	FM(31) =	FM(11)-FM(12)*FM(21)/FM(22)	00004420 00004430	
ISN 0229	C	QUOT =	FM(12)/FM(22)	00004440	
ISN 0230		DO 10	J=32,37	00004450	
ISN 0231		FM(J) =	(FM(J-19)-QUOT*FM(J-9))/FM(31)	00004460	
ISN 0232	10	CONTINUE		00004470 00004480	
ISN 0233	C	FM(33) =	(FM(19)-QUOT*FM(29))/FM(31)	00004490 00004500	
ISN 0234	C	FM(39) =	(FM(20)-QUOT*FM(30))/FM(31)	00004510 00004520	
ISN 0235	C	DO 20	J=40,47	00004530	
ISN 0236		FM(J) =	(FM(J-27)-FM(J-8)*FM(11))/FM(12)	00004540	
ISN 0237	20	CONTINUE		00004550 00004560 00004570	
ISN 0238	C	FC(1) =	-AC*SIN(ALPCH)/(CO*TAU)+VELAX*AC*I/(CO*TAU*SSI**2)	00004580 00004590	
ISN 0239	C	FC(2) =	-AC*COS(ALPCH)/TAU+(K+CO*VEL)*AC*I/(CO*TAU*SSI**2)	00004600 00004610	
ISN 0240	C	FC(3) =	VOL1/2*VELAX/(CO*TAU*SSI**2)	00004620 00004630	
ISN 0241	C	FC(4) =	VOL1*(K+CO*VEL)/(CO*TAU*SSI**2)/2	00004640 00004650	
ISN 0242	C	FC(5) =	FC(1)*B1R*FE(7)-FC(1)*B1I*SIN(CO*TAU)+FC(2)*FE(7) +FC(3)*K+B1I*FE(7)+FC(3)*B1R*SIN(CO*TAU)*K+FC(4)*	00004660 00004670	

		?	K*SIN(CO*TAU)	00004680
	C			00004690
ISN 0243	FC(6) =	-FC(1)*B1I*FE(7)-FC(1)*B1R*SIN(CO*TAU)-FC(2)*	SIN(CO*TAU)+FC(3)*K*B1R*FE(7)-FC(3)*K*B1I*	00004700
		?	SIN(CO*TAU)+FC(4)*K*FE(7)	00004710
	C			00004720
ISN 0244	FC(7) =	FC(5)*FE(16)+FC(6)*FE(26)-I		00004730
ISN 0245	FC(8) =	FC(5)*FE(17)+FE(6)*FE(27)		00004740
ISN 0246	FC(9) =	FC(5)*FE(18)+FC(6)*FE(28)+SST		00004750
ISN 0247	FC(10) =	FC(5) * FE(19) + FC(6) * FE(29)- K/12 * VOL1 /2		00004760
ISN 0248	FC(11) =	FC(5)*FE(20)+FC(6)*FE(30)+DT*SST		00004770
ISN 0249	FC(12) =	FC(5)*FE(21)+FC(6)*FE(31)		00004780
ISN 0250	FC(13) =	FC(5)*FE(22)+FE(6)*FE(32)+DT		00004790
ISN 0251	FC(14) =	FC(5)*FE(23)+FC(6)*FE(33)		00004800
ISN 0252	FC(15) =	FC(5) * FE(24) + FC(6) * FE(34)- K/12 * AVGD1		00004810
ISN 0253	FC(16) =	FC(5)*FE(25)+FC(6)*FE(35)		00004820
ISN 0254	FC(17) =	FC(9)*FM(32)+FC(10)*FM(40)+FC(7)		00004830
ISN 0255	FC(18) =	FC(9)*FM(33)+FC(10)*FM(41)+FC(6)		00004840
	C			00004850
ISN 0256	DO 30 J=	19,24		00004860
ISN 0257	FC(J) =	FC(9)*FM(J+15)+FC(10)*FM(J+23)+FC(J-8)		00004870
ISN 0258	30	CONTINUE		00004880
	C			00004890
ISN 0259	FC(21) =	-FC(21)		00004900
ISN 0260	FC(22) =	-FC(22)		00004910
	C			00004920
ISN 0261	FC(25) =	FC(1)*(B1I*FE(7)+B1R*SIN(CO*TAU))+FC(2)*SIN(CO*TAU)-		00004930
		?	FC(3)*K*(B1R*FE(7)-B1I*SIN(CO*TAU))-FC(4)*K*FE(7)	00004940
	C			00004950
ISN 0262	FC(26) =	FC(1)*(B1R*FE(7)-B1I*SIN(CO*TAU))+FC(2)*FE(7)+FC(3)*		00004960
		?	K*(B1I+FE(7)+B1R*SIN(CO*TAU))+FC(4)*K*SIN(CO*TAU)	00004970
	C			00004980
ISN 0263	FC(27) =	FC(26)*FE(26)+FC(25)*FE(16)		00004990
ISN 0264	FC(28) =	FC(26)*FE(27)+FC(25)*FE(17)-I		00005000
ISN 0265	FC(29) =	FC(26) * FE(28) + FC(25) * FE(18) + VOL1 * K/12 / 2		00005010
ISN 0266	FC(30) =	FC(26)*FE(29)+FC(25)*FE(19)+SST		00005020
ISN 0267	FC(31) =	FC(26)*FE(30)+FC(25)*FE(20)		00005030
ISN 0268	FC(32) =	FC(26)*FE(31)+FC(25)*FE(21)+DT*SSI		00005040
ISN 0269	FC(33) =	FC(26)*FE(32)+FC(25)*FE(22)		00005050
ISN 0270	FC(34) =	FC(26)*FE(33)+FC(25)*FE(23)+DT		00005060
ISN 0271	FC(35) =	FC(26)*FE(34)+FC(25)*FE(24)		00005070
ISN 0272	FC(36) =	FC(26) * FE(35) + FC(25) * FE(25) + K/12 * AVGD1		00005080
ISN 0273	FC(37) =	FC(29)*FM(32)+FC(30)*FM(40)+FC(27)		00005090
ISN 0274	FC(38) =	FC(29)*FM(33)+FC(30)*FM(41)+FC(28)		00005100
	C			00005110
ISN 0275	DO 40 J=	39,44		00005120
ISN 0276	FC(J) =	FC(29)*FM(J-5)+FC(30)*FM(J+3)+FC(J-8)		00005130
ISN 0277	40	CONTINUE		00005140
	C			00005150
ISN 0278	FC(41) =	-FC(41)		00005160
ISN 0279	FC(42) =	-FC(42)		00005170
	C			00005180
ISN 0280	QUOT =	FC(42)/FC(22)		00005190
ISN 0281	FC(45) =	FC(41)-QUOT*FC(21)		00005200
ISN 0282	FC(46) =	1/FC(45)*(FC(37)-QUOT*FC(17))		00005210
ISN 0283	FC(47) =	1/FC(45)*(FC(39)-QUOT*FC(18))		00005220
ISN 0284	FC(48) =	1/FC(45)*(FC(39)-QUOT*FC(19))		00005230
ISN 0285	FC(49) =	1/FC(45)*(FC(40)-QUOT*FC(20))		00005240
ISN 0286	FC(50) =	1/FC(45)*(FC(43)-QUOT*FC(23))		00005250

ISN	Code	Statement	Address
*VERSION 1.3.0 (01 MAY 80) UNST SYSTEM/376 FORTRAN H EXTENDED (ENHANCED) DATE 80.353/15.02.49 PAGE 10			
ISN 0207	FC(51) =	1/FC(45)*(FC(44)-QUOT*FC(24))	00005270
ISN 0208	FC(56) =	1/FC(42)*(FC(43)-FC(41)*FC(50))	00005280
ISN 0209	FC(57) =	1/FC(42)*(FC(44)-FC(41)*FC(51))	00005290
C			
ISN 0290	DO 50 J=	52,55	00005310
ISN 0291	FC(J) =	1/FC(42)*(FC(J-15)-FC(41)*FC(J-6))	00005320
ISN 0292	50 CONTINUE		00005330
C			
C			
C			
C INTERBLADE ANALYSIS OF SECTION 2			
C			
ISN 0293	FM2(1) =	AO*(-SSU**2-VU**2)	00005370
ISN 0294	FM2(2) =	AVGV2 * VOL2 * K/12 / 2	00005410
ISN 0295	FM2(3) =	FU*DU**VU**2	00005420
ISN 0296	FM2(3A) =	PT*DT*SST**2	00005430
ISN 0297	FM2(4) =	FM2(2)+FM2(1)**2/FM2(2)	00005440
ISN 0298	FM2(5) =	FM2(3)/FM2(4)	00005450
ISN 0299	FM2(6) =	FM2(3)*FM2(1)/(FM2(2)*FM2(4))	00005460
ISN 0300	FM2(7) =	-2/FM2(4)	00005470
ISN 0301	FM2(8) =	-2*FM2(1)/(FM2(2)*FM2(4))	00005480
ISN 0302	FM2(9) =	FM2(3A)/FM2(4)	00005490
ISN 0303	FM2(10) =	FM2(1)/FM2(2)*FM2(3A)/FM2(4)	00005500
ISN 0304	FM2(11) =	1 / FM2(4) * (2 * FLOW + FM2(1) / FM2(2) * AVGD2 * VOL2 * K/12 / 2)	00005520
ISN 0305	FM2(12) =	1 / FM2(4) * (FM2(1) / FM2(2) * 2 * FLOW - AVGD2 * VOL2 * K/12 / 2)	00005530
ISN 0306	FM2(13) =	1 / FM2(4) * (-2 * FM2(1) / FM2(2) * DT * SST - AVGD2 * VOL2 * K/12 / 2)	00005550
ISN 0307	FM2(14) =	1 / FM2(4) * (-2 * DT * SST + K/12 * FM2(1) / 2 / FM2(2) * VOL2 * AVGD2)	00005570
ISN 0308	FM2(15) =	-AVGV2 * VOL2 * K/12 / (2 * FM2(4))	00005590
ISN 0309	FM2(16) =	FM2(1) / FM2(2) * AVGV2 * VOL2 * K/12 / (2 * FM2(4))	00005600
ISN 0310	FM2(17) =	-AVGV2 * AVGD2 * K/12 / FM2(4)	00005610
ISN 0311	FM2(18) =	AVGD2 * AVGV2 * K/12 * FM2(1) / (FM2(2) * FM2(4))	00005620
C			
ISN 0312	M(1) =	(FM2(3)-FM2(2)*FM2(5))/FM2(1)	00005650
ISN 0313	M(2) =	-FM2(2)*FM2(6)/FM2(1)	00005660
ISN 0314	M(3) =	(-2-FM2(7)*FM2(8))/FM2(1)	00005670
ISN 0315	M(4) =	-FM2(2)*FM2(8)/FM2(1)	00005680
ISN 0316	M(5) =	(FM2(3A)-FM2(2)*FM2(9))/FM2(1)	00005690
ISN 0317	M(6) =	-FM2(2)*FM2(10)/FM2(1)	00005700
ISN 0318	M(7) =	(2*FLOW-FM2(2)*FM2(11))/FM2(1)	00005710
ISN 0319	M(8) =	(-K/12 * AVGD2 * VOL2 / 2 - FM2(2) * FM2(12)) / FM2(1)	00005720
ISN 0320	M(9) =	(-AVGD2 * VOL2 * K/12 / 2 + FM2(2) * FM2(13)) / FM2(1)	00005730
ISN 0321	M(10) =	(-2*DT*SST-FM2(2)*FM2(14))/FM2(1)	00005740
ISN 0322	M(11) =	(-AVGV2 * VOL2 * K/12 / 2 + FM2(2) * FM2(15)) / FM2(1)	00005750
ISN 0323	M(12) =	(-FM2(2)*FM2(16))/FM2(1)	00005760
ISN 0324	M(13) =	(-K/12 * AVGD2 * AVGV2 - FM2(2) * FM2(17)) / FM2(1)	00005770
ISN 0325	M(14) =	-FM2(2)*FM2(18)/FM2(1)	00005780
ISN 0326	M(15) =	DU * AO + VU * AO * M(7) - K/12 * VOL2 / 2 * FM2(11)	00005790
ISN 0327	M(16) =	VU * AO * M(8) - K/12 * VOL2 / 2 * FM2(12)	00005800
ISN 0328	M(17) =	-VU * AO * M(1) - DU * VU + K/12 * VOL2 / 2 * FM2(5)	00005810
ISN 0329	M(18) =	-VU * AO * M(2) + K/12 * VOL2 / 2 * FM2(6)	00005820
ISN 0330	M(19) =	-VU * AO * M(12) + SST + K/12 * VOL2 / 2 * FM2(16)	00005830
ISN 0331	M(20) =	-VU * AO * M(11) + K/12 * VOL2 / 2 * FM2(15)	00005840
ISN 0332	M(21) =	-VU * AO * M(3) + K/12 * VOL2 / 2 * FM2(7)	00005850

ISN 0333	M(22) =	-VU * AO * M(4) + K/12 * VOL2 / 2 * FM2(8)	00005060
ISN 0334	M(23) =	-VU * AO * M(5) + DT * SST + K/12 * VOL2 / 2 * FM2(9)	00005070
ISN 0335	M(24) =	-VU * AO * M(6) + K/12 * VOL2 / 2 * FM2(10)	00005080
ISN 0336	M(25) =	-VU * AO * M(9) + K/12 * VOL2 / 2 * FM2(13)	00005090
ISN 0337	M(26) =	-VU * AO * M(10) + K/12 * VOL2 / 2 * FM2(14) + DT	00005100
ISN 0338	M(27) =	-VU * AO * M(13) + K/12 * VOL2 / 2 * FM2(17) + K/12 * AVGD2	00005110
ISN 0339	M(28) =	-VU * AO * M(14) + K/12 * VOL2 * FM2(18) / 2	00005120
	C		00005130
	C		00005140
ISN 0340	CL(1) =	DU * AO + VU * AO * FM2(12) + K/12 * VOL2 / 2 * M(8)	00005160
ISN 0341	CL(2) =	VU * AO * FM2(11) + K/12 * VOL2 / 2 * M(7)	00005170
ISN 0342	CL(3) =	-VU * AO * FM2(5) - K/12 * VOL2 / 2 * M(1)	00005180
ISN 0343	CL(4) =	-DU * VU - VU * AO * FM2(6) - K/12 * VOL2 / 2 * M(2)	00005190
ISN 0344	CL(5) =	-K/12 * VOL2 / 2 - VU * AO * FM2(16) - K/12 * VOL2 / 2 * M(12)	00006010
ISN 0345	CL(6) =	SST - VU * AO * FM2(15) - K/12 * VOL2 / 2 * M(11)	00006020
ISN 0346	CL(7) =	-VU * AO * FM2(14) - K/12 * VOL2 / 2 * M(10)	00006030
ISN 0347	CL(8) =	DT - VU * AO * FM2(13) - K/12 * VOL2 / 2 * M(9)	00006040
ISN 0348	CL(9) =	-VU * AO * FM2(7) - K/12 * VOL2 / 2 * M(3)	00006050
ISN 0349	CL(10) =	-VU * AO * FM2(8) - K/12 * VOL2 / 2 * M(4)	00006060
ISN 0350	CL(11) =	-VU * AO * FM2(9) - K/12 * VOL2 / 2 * M(5)	00006070
ISN 0351	CL(12) =	-VU * AO * FM2(10) - K/12 * VOL2 / 2 * M(6) + AVGD2*SST	00006080
ISN 0352	CL(13) =	-VU * AO * FM2(16) - K/12 * VOL2 / 2 * (2 * DT + M(14))	00006090
ISN 0353	CL(14) =	-VU * AO * FM2(17) - K/12 * VOL2 / 2 * M(13)	00006100
ISN 0354	CL(15) =	CL(2)-CL(1)/M(16)*M(15)	00006110
ISN 0355	DO 60 J=	16,19	00006120
ISN 0356	CL(J) =	1/CL(15)*(CL(J-13)-CL(1)/M(16)*M(J+1))	00006130
ISN 0357	60 CONTINUE		00006140
	C		00006150
ISN 0358	CL(20) =	1/CL(15)*(CL(7)-CL(1)/M(16)*M(26))	00006170
ISN 0359	CL(21) =	1/CL(15)*(CL(8)-CL(1)/M(16)*M(25))	00006180
ISN 0360	CL(22) =	1/CL(15)*(CL(9)-CL(1)/M(16)*M(21))	00006190
ISN 0361	CL(23) =	1/CL(15)*(CL(10)-CL(1)/M(16)*M(22))	00006200
ISN 0362	CL(24) =	1/CL(15)*(CL(11)-CL(1)/M(16)*M(23))	00006210
ISN 0363	CL(25) =	1/CL(15)*(CL(12)-CL(1)/M(16)*M(24))	00006220
ISN 0364	CL(26) =	1/CL(15)*(CL(13)-CL(1)/M(16)*M(28))	00006230
ISN 0365	CL(27) =	1/CL(15)*(CL(14)-CL(1)/M(16)*M(27))	00006240
	C		00006250
ISN 0366	DO 70 J =	26,39	00006260
ISN 0367	CL(J) =	(CL(J-25)-CL(2)*CL(J-12))/CL(1)	00006270
ISN 0368	70 CONTINUE		00006280
	C		00006290
	C		00006300
	C	INTERBLADE ANALYSIS OF SECTION 3	00006310
	C		00006320
	C		00006330
ISN 0369	S(1) =	2*GAMMA*PU/(GAMMA+1)	00006340
ISN 0370	S(2) =	(2*GAMMA/(GAMMA+1))*PUNUS-PO-(GAMMA-1)/(GAMMA+1)*PU	00006350
ISN 0371	S(3) =	2*GAMMA/(GAMMA+1)*SSU*NUS+US-(GAMMA-1)/(GAMMA+1)*SSU	00006360
	C		00006370
ISN 0372	S(4) =	2*GAMMA/(GAMMA+1)*SSU*PU	00006380
ISN 0373	S(5) =	(GAMMA-1)*NUS**2+2	00006390
ISN 0374	S(6) =	(2-2*GAMMA)*PO+(2*GAMMA+2)*DU	00006400
ISN 0375	S(7) =	(GAMMA+1)*NUS**2	00006410
ISN 0376	S(8) =	S(6)/(SSU*S(5))	00006420
ISN 0377	S(9) =	S(7)/S(5)	00006430
	C		00006440

*VERSION 1.3.0 (01 MAY 80)	CNST	SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)	DATE 80.353/15.02.49	PAGE 12
ISN 0378	C	LM(1) = DE*UE*ACE/(CO*TAU)-FRE*SIN(ALPCH)/(CO*TAU)+(REV*	00006450	
	?	COS(ALPCH-BETA2))*2*ACE*UE*DE/(CO*TAU*SOSEX**2)-DE*	00006460	
	?	ACE/(CO*TAU)*SIN(ALPCH)*REV/COS(ALPCH-BETA2)	00006490	
ISN 0379	C	LM(2) = DE*(K+CO*VE)*ACE/(CO*TAU)-COS(ALPCH)*FRE/TAU+DE*(REV*	00006500	
	?	COS(ALPCH-BETA2))*2*ACE*(K+CO*VE)/(CO*TAU*SOSEX**2)-	00006510	
	?	DE*ACE*COO(ALPCH)*REV*COO(ALPCH-BETA2)	00006520	
ISN 0380	C	LM(3) = AVGD3*VOL3/2*UE*DE/(CO*TAU*SOSEX**2)-AVGD3*VOL3*	00006530	
	?	SIN(ALPCH)/(2*CO*TAU)	00006540	
			00006550	
ISN 0381	C	LM(4) = AVGD3*VOL3/(2*CO*TAU*SOSEX**2)*(K+CO*VE)*DE-	00006560	
	?	AVGD3*VOL3*COO(ALPCH)/(2*TAU)	00006570	
			00006580	
ISN 0382	C	R = -(K+VE*CO)/UE	00006600	
ISN 0383	C	LM(5) = FRE*SIN(ALPCH)/(TAU*(R**2+CO**2))-R*COO(ALPCH)*FRE/(CO*	00006610	
	?	TAU*(R**2+CO**2))+REV*COO(ALPCH-BETA2)*DE*ACE*	00006620	
	?	(SIN(ALPCH)-R/CO*COO(ALPCH))/(TAU*(R**2+CO**2))	00006630	
ISN 0384	C	LM(6) = AVGD3*VOL3/(2*TAU*(R**2+CO**2))*(SIN(ALPCH)-R/CO*	00006640	
	?	COO(ALPCH))	00006650	
			00006660	
ISN 0385	C	LM(7) = COS(CO*TAU)-1	00006670	
			00006680	
ISN 0386	C	LM(8) = LM(1)*LM(7)*B2R-B2I*SIN(CO*TAU))+LM(2)*LM(7)-K*	00006690	
	?	LM(3)*(B2I*LM(7)+B2R*SIN(CO*TAU))-K*LM(4)*SIN(CO*TAU)	00006700	
			00006710	
ISN 0387	C	LM(9) = -LM(1)*(LM(7)*B2I+B2R*SIN(CO*TAU))-LM(2)*SIN(CO*TAU)-	00006720	
	?	K*LM(3)*(B2R*LM(7)-B2I*SIN(CO*TAU))-K*LM(4)*LM(7)	00006730	
			00006740	
ISN 0388	C	LM(10) = -K*LM(6)*SIN(CO*TAU)+LM(5)*LM(7)	00006750	
			00006760	
ISN 0389	C	LM(11) = -K*LM(6)*LM(7)-LM(5)*SIN(CO*TAU)	00006770	
			00006780	
ISN 0390	C	LM(12) = LM(1)*(B2I*LM(7)+B2R*SIN(CO*TAU))+LM(2)*SIN(CO*TAU)+	00006790	
	?	LM(3)*K*(B2R*LM(7)-B2I*SIN(CO*TAU))+LM(4)*K*LM(7)	00006800	
			00006810	
ISN 0391	C	LM(13) = LM(1)*(B2R*LM(7)-B2I*SIN(CO*TAU))+LM(2)*LM(7)-LM(3)*	00006820	
	?	K*(B2I*LM(7)+B2R*SIN(CO*TAU))-LM(4)*K*SIN(CO*TAU)	00006830	
			00006840	
ISN 0392	C	LM(14) = LM(5)*SIN(CO*TAU)+LM(6)*K*LM(7)	00006850	
			00006860	
ISN 0393	C	LM(15) = LM(5)*LM(7)-LM(6)*K*SIN(CO*TAU)	00006870	
			00006880	
ISN 0394	C	LM(16) = LM(12)-LM(13)/LM(9)*LM(8)	00006890	
			00006900	
ISN 0395	C	LM(17) = 1/LM(16)*(LM(14)-LM(13)*LM(10)/LM(9))	00006910	
			00006920	
ISN 0396	C	LM(18) = 1/LM(16)*(LM(15)-LM(13)*LM(11)/LM(9))	00006930	
			00006940	
ISN 0397	C	LM(19) = 1/LM(16)	00006950	
			00006960	
ISN 0398	C	LM(20) = -LM(13)/(LM(9)*LM(16))	00006970	
			00006980	
			00006990	
ISN 0399	C	LM(21) = (LM(14)-LM(12)*LM(17))/LM(13)	00007000	
ISN 0400	C	LM(22) = (LM(15)-LM(12)*LM(18))/LM(13)	00007010	
ISN 0401	C	LM(23) = -LM(12)*LM(20)/LM(13)	00007020	
ISN 0402	C	LM(24) = (1-LM(12)*LM(19))/LM(13)	00007030	

ISN 0403		WRITE(6,333) B1R,B1I	00007040
ISN 0404	333	FORMAT(// ' B1'/(2E15.6))	00007050
	C		00007050
ISN 0405		WRITE(6,6000) FE	00007070
ISN 0406		WRITE(6,6100) AVGV1	00007080
ISN 0407		WRITE(6,6200) VOL1	00007090
ISN 0408		WRITE(6,6001) FM	00007100
ISN 0409		WRITE(6,6002) FC	00007110
ISN 0410		WRITE(6,6003) FM2	00007120
ISN 0411		WRITE(6,6004) M	00007130
ISN 0412		WRITE(6,6005) CL	00007140
ISN 0413		WRITE(6,6006) S	00007150
ISN 0414		WRITE(6,6007) LM	00007160
ISN 0415		WRITE(6,6008) LC	00007170
ISN 0416	6000	FORMAT('1', ' FE',/,9(E12.4,2X))	00007180
ISN 0417	6100	FORMAT('-', ' AVGV1',/,1X,E12.4,2X)	00007190
ISN 0418	6200	FORMAT('-', ' VOL1',/,1X,E12.4,2X)	00007200
ISN 0419	6001	FORMAT('-', ' FM',/,9(E12.4,2X))	00007210
ISN 0420	6002	FORMAT('-', ' FC',/,9(E12.4,2X))	00007220
ISN 0421	6003	FORMAT('-', ' FM2',/,9(E12.4,2X))	00007230
ISN 0422	6004	FORMAT('-', ' M',/,9(E12.4,2X))	00007240
ISN 0423	6005	FORMAT('-', ' CL',/,9(E12.4,2X))	00007250
ISN 0424	6006	FORMAT('-', ' S',/,9(E12.4,2X))	00007260
ISN 0425	6007	FORMAT('-', ' LM',/,9(E12.4,2X))	00007270
ISN 0426	6008	FORMAT('-', ' LC',/,9(E12.4,2X))	00007280
	C		00007290
	C	CALCULATE_EITHER IN_TORSIONAL_OR_BENDING_MODE	00007300
	C		00007310
	C		00007320
ISN 0427		DO 15 II = 1,2	00007330
ISN 0428		IF (II.EQ.1) GOTO 15	00007340
	C		00007350
	C	CALCULATES TORSIONAL MODES AREAS	00007360
	C		00007370
ISN 0430		RMAPI = C*Z*ALFAP*(COS(SIGMA)-1)/ASTAR	00007380
ISN 0431		IMAPI = C*Z*ALFAP*(SIN(SIGMA))/ASTAR	00007390
ISN 0432		RMAPT = (Z-XSTAR/C)*(COS(SIGMA)-1)*C*ALFAP/ASTAR	00007400
ISN 0433		IMAPT = (Z-XSTAR/C)*SIN(SIGMA)*C*ALFAP/ASTAR	00007410
ISN 0434		RMAP = (Z-XS/C)*(COS(SIGMA)-1)*C*ALFAP/ASTAR	00007420
ISN 0435		IMAP = (Z-XS/C)*SIN(SIGMA)*C*ALFAP/ASTAR	00007430
ISN 0436		RACE = (Z-1)*(COS(SIGMA)-1)*C*ALFAP/ASTAR	00007440
ISN 0437		IACE = (Z-1)*SIN(SIGMA)*C*ALFAP/ASTAR	00007450
	C		00007460
	C	CALCULATES TORSIONAL MODE VOLUMES FOR THE FIRST TWO SECTIONS	00007470
	C		00007480
ISN 0438		V1AR = 0.0	00007490
ISN 0439		V1AI = 0.0	00007500
	C		00007510
ISN 0440		V1BR = C*DELTA*Z*ALFAP*(COS(SIGMA)-1)*COS(ALFCH)/(ASTAR*B)	00007520
ISN 0441		V1BI = C*DELTA*Z*ALFAP*(SIN(SIGMA))*COS(ALFCH)/(ASTAR*B)	00007530
	C		00007540
ISN 0442		V1CR = C*ALFAP*(COS(SIGMA)-1)*(Z*(XSTAR-DELTA*COS(ALFCH))-	00007550
	?	(XSTAR**2-(DELTA*COS(ALFCH))**2)/(2*C))/(ASTAR*B)	00007560
ISN 0443		V1CI = C*ALFAP*(SIN(SIGMA))*(Z*(XSTAR-DELTA*COS(ALFCH))-	00007570
	?	(XSTAR**2-(DELTA*COS(ALFCH))**2)/(2*C))/(ASTAR*B)	00007580
	C		00007590
ISN 0444		V1R = V1AR+V1BR+V1CR	00007600
ISN 0445		V1I = V1AI+V1BI+V1CI	00007610
	C		00007620

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  ISN 0446 V2R = C*ALPBARY(COS(SIGMAB)-1)*(Z*(XS-XSTAR)-((XS**2- 00007630
    ? XSTAR**2)/(2*C)))/(ASTAR*B) 00007640
  ISN 0447 V2I = C*ALPBARY(SIN(SIGMAB))*(Z*(XS-XSTAR)-((XS**2- 00007650
    ? XSTAR**2)/(2*C)))/(ASTAR*B) 00007660
  C 00007670
  ISN 0448 GOTO 17 00007680
  C 00007690
  C CALCULATES BENDING MODE AREAS 00007700
  C 00007710
  ISN 0449 13 RACE = H*(COS(SIGMAB)-1)/ASTAR 00007720
  ISN 0450 IACE = H*SIN(SIGMAB)/ASTAR 00007730
  ISN 0451 IMAP = H*GIN(SIGMAB)/ASTAR 00007740
  ISN 0452 RMAP = H*(COS(SIGMAB)-1)/ASTAR 00007750
  ISN 0453 RMAPI = H*(COS(SIGMAB)-1)/ASTAR 00007760
  ISN 0454 IMAPT = H*SIN(SIGMAB)/ASTAR 00007770
  ISN 0455 IMAPI = H*SIN(SIGMAB)/ASTAR 00007780
  ISN 0456 RMAPT = H*(COS(SIGMAB)-1)/ASTAR 00007790
  C 00007800
  C CALCULATES BENDING MODE VOLUMES FOR THE FIRST TWO SECTIONS 00007810
  C 00007820
  ISN 0457 VIAR = 0.0 00007830
  ISN 0458 VIAI = 0.0 00007840
  C 00007850
  ISN 0459 VISR = DELTA*H*(COS(SIGMAB)-1)*COS(ALFCH)/(ASTAR*B) 00007860
  ISN 0460 VIBI = DELTA*H*(SIN(SIGMAB))*COS(ALFCH)/(ASTAR*B) 00007870
  C 00007880
  ISN 0461 VICR = (H*(XSTAR-DELTA*COS(ALFCH))*(COS(SIGMAB)-1)+ASTAR*H* 00007890
    ? (COS(SIGMAB)-1)*COTAN(ALFCH))/(ASTAR*B) 00007900
  ISN 0462 VICI = (H*(XSTAR-DELTA*COS(ALFCH))*(SIN(SIGMAB))+ASTAR*H* 00007910
    ? (SIN(SIGMAB))*COTAN(ALFCH))/(ASTAR*B) 00007920
  C 00007930
  ISN 0463 VIR = VIAR+VIBI+VICR 00007940
  ISN 0464 VII = VIAI+VIBI+VICI 00007950
  C 00007960
  ISN 0465 V2R = H*(COS(SIGMAB)-1)*(XS-XSTAR)/(ASTAR*B) 00007970
  ISN 0466 V2I = H*(SIN(SIGMAB))*(XS-XSTAR)/(ASTAR*B) 00007980
  C 00007990
  C 00008000
  C 00008010
  ISN 0467 17 ISSP = FC(52)*RMAPI+FC(53)*IMAPI+FC(54)*RMAPT+FC(55)*IMAPT+ 00008020
    ? FC(56)*VII+FC(57)*VIR 00008030
  C 00008040
  ISN 0468 RSSP = FC(46)*RMAPI+FC(47)*IMAPI+FC(48)*RMAPT+FC(49)*IMAPT+ 00008050
    ? FC(50)*VII+FC(51)*VIR 00008060
  C 00008070
  ISN 0469 IMDP = FM(40)*RMAPI+FM(41)*IMAPI+FM(42)*RMAPT+FM(43)*IMAPT+ 00008080
    ? FM(44)*RSSP+FM(45)*ISSP+FM(46)*VII+FM(47)*VIR 00008090
  C 00008100
  ISN 0470 RMDP = FM(32)*RMAPI+FM(33)*IMAPI+FM(34)*RMAPT+FM(35)*IMAPT+ 00008110
    ? FM(36)*RSSP+FM(37)*ISSP+FM(38)*VII+FM(39)*VIR 00008120
  C 00008130
  ISN 0471 IVPU = FE(26)*RMAPI+FE(27)*IMAPI+FE(28)*RMDP+FE(29)*IMDP+ 00008140
    ? FE(30)*RMAPT+FE(31)*IMAPT+FE(32)*RSSP+FE(33)*ISSP+ 00008150
    ? FE(34)*VII+FE(35)*VIR 00008160
  C 00008170
  ISN 0472 RVPU = FE(16)*RMAPI+FE(17)*IMAPI+FE(18)*RMDP+FE(19)*IMDP+ 00008180
    ? FE(20)*RMAPT+FE(21)*IMAPT+FE(22)*RSSP+FE(23)*ISSP+ 00008190
    ? FE(24)*VII+FE(25)*VIR 00008200
  C 00008210

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ISN 0473	IMFP =	IMDP*SSST**2	00008220
ISN 0474	RMFP =	RMDP*SSST**2	00008230
ISN 0475	RMVPU =	CL(16)*RMFP+CL(17)*IMFP+CL(18)*RMDP+CL(19)*IMDP+CL(20)*RSDP+CL(21)*ISSP+CL(22)*RMFP+CL(23)*IMFP+CL(24)*RMFP+CL(25)*IMFP+CL(26)*V2R+CL(27)*V2I	00008250 00008260 00008270 00008280
ISN 0476	IMVPU =	CL(28)*RMFP+CL(29)*IMFP+CL(30)*RMDP+CL(31)*IMDP+CL(32)*RSDP+CL(33)*ISSP+CL(34)*RMFP+CL(35)*IMFP+CL(36)*RMFP+CL(37)*IMFP+CL(38)*V2R+CL(39)*V2I	00008290 00008300 00008310 00008320
ISN 0477	RMDPU =	M(1)*RMFP+M(2)*IMFP+M(3)*RMFP+M(4)*IMFP+M(5)*RMFP+M(6)*IMFP+M(7)*RMFP+M(8)*IMFP+M(9)*ISSP+M(10)*RSDP+M(11)*IMDP+M(12)*RMDP+M(13)*V2I+M(14)*V2R	00008330 00008340 00008350 00008360
ISN 0478	IMDPU =	FM2(5)*RMFP+FM2(6)*IMFP+FM2(7)*RMFP+FM2(8)*IMFP+FM2(9)*RMFP+FM2(10)*IMFP+FM2(11)*RMVPU+FM2(12)*IMVPU+FM2(13)*ISSP+FM2(14)*RSDP+FM2(15)*IMDP+FM2(16)*RMDP+FM2(17)*V2I+FM2(18)*V2R	00008370 00008380 00008390 00008400 00008410
ISN 0479	RDDS =	S(8)*RMVPU+S(9)*RMDPU	00008420
ISN 0480	IDDS =	S(3)*IMVPU+S(9)*IMDPU	00008430
ISN 0481	RMPPU =	SSU**2*RMDPU	00008440
ISN 0482	IMPPU =	SSU**2*IMDPU	00008450
ISN 0483	RSSU =	(GAMMA-1)/2*RMPPU/(DU*SSU)	00008460
ISN 0484	ISSU =	(GAMMA-1)/2*IMPPU/(DU*SSU)	00008470
ISN 0485	RFDS =	SOEDS**2*RDDS	00008480
ISN 0486	IFDS =	SOEDS**2*IDDS	00008490
ISN 0487	USR =	1/S(1)*(S(2)*RSDSU+S(3)*RMPPU+S(4)*SSU*RMVPU-SSU*RPDS)	00008500
ISN 0488	USI =	1/S(1)*(S(2)*ISOSU+S(3)*IMPPU+S(4)*SSU*IMVPU-SSU*IPDS)	00008510
ISN 0489	IMSP =	-USR/K	00008520
ISN 0490	RMSP =	USI/K	00008530 00008540 00008550
ISN 0491	RMFPI =	-VELX/(CO*TAU)*(B1R*RVPU*(COS(CO*TAU)-1)-B1I*IVPU*(COS(CO*TAU)-1)-IVPU*B1R*SIN(CO*TAU)-B1I*RVPU*SIN(CO*TAU))-(K+CO*VEL)/(CO*TAU)*(RVPU*(COS(CO*TAU)-1)-IVPU*SIN(CO*TAU))	00008560 00008570 00008580 00008590
ISN 0492	IMFPI =	-VELX/(CO*TAU)*(IVPU*B1R*(COS(CO*TAU)-1)+B1I*RVPU*(COS(CO*TAU)-1)+B1R*RVPU*SIN(CO*TAU)-B1I*IVPU*SIN(CO*TAU))-(K+CO*VEL)/(CO*TAU)*(IVPU*(COS(CO*TAU)-1)+RVPU*SIN(CO*TAU))	00008600 00008610 00008620 00008630 00008640
ISN 0493	RMFPI =	RMFPI/SSI**2	00008650
ISN 0494	IMFPI =	IMFPI/SSI**2	00008660 00008670
ISN 0495	RMSSPI =	(GAMMA-1)/(2*SSI)*RMFPI	00008680
ISN 0496	IMSSPI =	(GAMMA-1)/(2*SSI)*IMFPI	00008690 00008700
ISN 0497	RITVPI =	1/TAU*(RVPU*(COS(CO*TAU)-1)-IVPU*SIN(CO*TAU))	00008710
ISN 0498	ITVPI =	1/TAU*(IVPU*(COS(CO*TAU)-1)+RVPU*SIN(CO*TAU))	00008720 00008730
ISN 0499	RIAVPI =	1/(CO*TAU)*(B1R*RVPU*(COS(CO*TAU)-1)-B1I*IVPU*(COS(CO*TAU)-1)-IVPU*B1R*SIN(CO*TAU)-B1I*RVPU*SIN(CO*TAU))	00008740 00008750 00008760 00008770
ISN 0500	IIAVPI =	1/(CO*TAU)*(B1R*IVPU*(COS(CO*TAU)-1)+B1I*RVPU*(COS(CO*TAU)-1)+B1R*RVPU*SIN(CO*TAU)-B1I*IVPU*SIN(CO*TAU))	00008780 00008790 00008800

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ISN 0501	C	IF(II.EQ.2) GOTO 24	00003310
	C		00003320
	C	CALCULATES THE BENDING MODE VOLUMES FOR THE THIRD SECTION	00003330
	C		00003340
ISN 0503		V3AR = (H*(COS(SIGMAB)-1)*(C-XS)-(AO+ACE)*ASTAR/2*RHSP*B)/(ASTAR*B)	00003350
ISN 0504		V3BR = DELTA/2.0*SIN(2*ALPCH)*H*(COS(SIGMAB)-1)/(ASTAR*B)	00003360
ISN 0505		V3CR = 0.0	00003370
ISN 0506		V3AI = (H * SIN(SIGMAD) * (C - XS) - (AO + ACE) * ASTAR / 2 * RHSP * B) / (ASTAR * B)	00003380
ISN 0507		V3BI = DELTA/2.0*SIN(2*ALPCH)*H*(SIN(SIGMAB))/(ASTAR*B)	00003390
ISN 0508		V3CI = 0.0	00003400
ISN 0509		WRITE(6,410)	00003410
ISN 0510	410	FORMAT('1',' BENDING')	00003420
ISN 0511		GOTO 25	00003430
	C		00003440
	C	CALCULATES THE TORSIONAL MODE VOLUMES FOR THE THIRD SECTION	00003450
	C		00003460
ISN 0512	24	V3AR = (C*ALFBAR*(COS(SIGMAT)-1)*(Z*(C-XS)-(C**2-XS**2)/(2*C))-AO*ASTAR*RHSP*B)/(ASTAR*B)	00003470
ISN 0513		V3AI = (C*ALFBAR*(SIN(SIGMAT))*(Z*(C-XS)-(C**2-XS**2)/(2*C))-AO*ASTAR*RHSP*B)/(ASTAR*B)	00003480
ISN 0514		V3BR = ALFBAR*(Z-1)*(COS(SIGMAT)-1)*C*COS(ALPCH)*DELTA/(ASTAR*B)	00003490
ISN 0515		V3BI = ALFBAR*(Z-1)*SIN(SIGMAT)*C*COS(ALPCH)*DELTA/(ASTAR*B)	00003500
ISN 0516		V3CR = 0.0	00003510
ISN 0517		V3CI = 0.0	00003520
ISN 0518		WRITE(6,420)	00003530
ISN 0519	420	FORMAT('1',' TORSIONAL')	00003540
ISN 0520	25	V3I = V3AI+V3BI+V3CI	00003550
ISN 0521		V3R = V3AR+V3BR+V3CR	00003560
	C		00003570
	C		00003580
	C		00003590
	C		00003600
ISN 0522		RUDS = 1/DD*(RNDPU*VU+DU*RMVPU-RDDS*VD)	00003610
ISN 0523		IUDS = 1/DD*(INDFU*VU+DU*IMVPU-IDDS*VD)	00003620
ISN 0524		WDSI = IDDS*VD*AO+DD*IUDS*AO+DD*VD*IMAP	00003630
ISN 0525		WDSR = RDDS*VD*AO+DD*RUDS*AO+DD*VD*RMAP	00003640
ISN 0526		CI = -DE*REV*COS(ALPCH-BETA2)*IACE+WDSI-AVG03 * K/12 * V3R - K/12 *VOL3/2*RDDS	00003650
ISN 0527		CR = -DE*RACE*U2IRE+WDSR+AVG03 * K/12 * V3I+K*VOL3/2*IDDS	00003660
ISN 0528		IC = -IPDS*AO-PD*IMAP+IACE*(PE+(REV*COS(ALPCH-BETA2)))**2*DE) - WDS * IUDS + AVGV3 * VOL3 * K/12 * RDDS / 2 + AVGD3 * AVGV3 * K/12 * V3R + K/12 *AVGD3 * VOL3 * RUDS / 2 - WDSI * VD	00003670
ISN 0529		RC = -RPDS*AO-PD*RMAP+RACE*(PE+(REV*COS(ALPCH-BETA2)))**2*DE) - WDS * RUDS - AVGV3 * VOL3 * K/12 * IDDS / 2 - AVGD3 * AVGV3 * K/12 * V3I - K/12 * AVGD3 * VOL3 * IUDS / 2 - WDSR * VD	00003680
	C		00003690
	C		00003700
ISN 0530		LC(1) = -UE*DE*ACE*REV*COS(ALPCH-BETA2)/(CO*TAU*SOSEX**2)+DE*	00003710

	C		ACE*SIN(ALPCH)/(CO*TAU)	00009400
ISN 0531	C	LC(2) =	-DE*(K+CO*VE)*ACE*REV*COS(ALPCH-BETA2)/ (CO*TAU*SOSEX**2)+DE*ACE*COS(ALPCH)/TAU	00009410 00009420 00009430
ISN 0532	C	LC(3) =	DE*ACE/(TAU*(R**2+CC**2))*(SIN(ALPCH)-R/CO*COS(ALPCH))	00009440 00009450 00009460
ISN 0533	C	LC(4) =	VOL3*UE*DE/(2*CO*TAU*SOSEX**2)	00009470 00009480
ISN 0534	C	LC(5) =	VOL3/2*(K+CO*VE)*DE/(CO*TAU*SOSEX**2)	00009490 00009500
ISN 0535	C	LC(6) =	LC(1)*B2R*LM(7)-LC(1)*B2I*SIN(CO*TAU)+LC(2)*LM(7)+ K*LC(4)*(B2I*LM(7)+B2R*SIN(CO*TAU))+K*LC(5)* SIN(CO*TAU)	00009510 00009520 00009530 00009540
ISN 0536	C	LC(7) =	-LC(1)*(B2I*LM(7)+B2R*SIN(CO*TAU))-LC(2)*SIN(CO*TAU)+ K*LC(4)*(LM(7)*B2R-B2I*SIN(CO*TAU))+K*LC(5)*LM(7)	00009550 00009560 00009570
ISN 0537	C	LC(8) =	LC(3)*LM(7)	00009580 00009590
ISN 0538	C	LC(9) =	-LC(3)*SIN(CO*TAU)	00009600 00009610
ISN 0539	C	LC(10) =	LC(6)*LM(17)+LC(7)*LM(21)+LC(8)	00009620 00009630
ISN 0540	C	LC(11) =	LC(6)*LM(18)+LC(7)*LM(22)+LC(9)	00009640 00009650
ISN 0541	C	LC(12) =	LC(6)*LM(20)+LC(7)*LM(23)	00009660 00009670
ISN 0542	C	LC(13) =	LC(6)*LM(19)+LC(7)*LM(24)	00009680 00009690
ISN 0543	C	LC(14) =	LC(1)*(B2I*LM(7)+B2R*SIN(CO*TAU))+LC(2)*SIN(CO*TAU)- LC(4)*K*(B2R*LM(7)-B2I*SIN(CO*TAU))-LC(5)*K*LM(7)	00009700 00009710 00009720
ISN 0544	C	LC(15) =	LC(1)*(B2R*LM(7)-B2I*SIN(CO*TAU))+LC(2)*LM(7)+ LC(4)*K*(B2I*LM(7)+B2R*SIN(CO*TAU))+LC(5)*K* SIN(CO*TAU)	00009730 00009740 00009750 00009760
ISN 0545	C	LC(16) =	LC(3)*SIN(CO*TAU)	00009770 00009780
ISN 0546	C	LC(17) =	LC(3)*LM(7)	00009790 00009800
ISN 0547	C	LC(18) =	LC(14)*LM(17)+LC(15)*LM(21)+LC(16)	00009810 00009820
ISN 0548	C	LC(19) =	LC(14)*LM(18)+LC(15)*LM(22)+LC(17)	00009830 00009840
ISN 0549	C	LC(20) =	LC(14)*LM(20)+LC(15)*LM(23)	00009850 00009860
ISN 0550	C	LC(21) =	LC(14)*LM(19)+LC(15)*LM(24)	00009870 00009880
ISN 0551	C	LC(22) =	LC(18)-LC(19)/LC(11)*LC(10)	00009890 00009900
ISN 0552	C	LC(23) =	-(LC(20)-LC(19)/LC(11)*LC(12))/LC(22)	00009910 00009920
ISN 0553	C	LC(24) =	-(LC(21)-LC(19)/LC(11)*LC(13))/LC(22)	00009930 00009940
ISN 0554	C	LC(25) =	(CI-CR*LC(19)/LC(11))/LC(22)	00009950 00009960
ISN 0555	C	RRFF =	LC(23)*RC+LC(24)*IC+LC(25)	00009970 00009980

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ISN 0536	C	LC(26) =	-(LC(20)+LC(18)*LC(23))/LC(19)	00009990	
ISN 0557	C	LC(27) =	-(LC(21)+LC(19)*LC(24))/LC(19)	00010000	
ISN 0553	C	LC(28) =	(CI-LC(25)*LC(18))/LC(19)	00010020	
ISN 0559	C	IRFF =	LC(26)*RC+LC(27)*IC+LC(28)	00010030	
ISN 0560	C	IVPD =	LM(21)*RRFF+LM(22)*IRFF+LM(23)*RC+LM(24)*IC	00010050	
ISN 0561	C	RVPD =	LM(17)*RRFF+LM(18)*IRFF+LM(19)*IC+LM(20)*RC	00010060	
ISN 0562	C	RMPPE =	-UE*DE/(CO*TAU)*(B2R*RVPD*(COS(CO*TAU)-1)-B2I*IVPD*	00010070	
	?		(COS(CO*TAU)-1)-IVPD*B2R*SIN(CO*TAU)-B2I*RVPD*SIN(CO*	00010080	
	?		TAU))-(K+CO*VE)/(CO*TAU)*(RVPD*(COS(CO*TAU)-1)-IVPD*	00010090	
	?		SIN(CO*TAU))	00010100	
ISN 0563	C	IMPPE =	-UE*DE/(CO*TAU)*(IVPD*B2R*(COS(CO*TAU)-1)+B2I*RVPD*	00010110	
	?		(COS(CO*TAU)-1)+B2R*RVPD*SIN(CO*TAU)-B2I*IVPD*SIN(CO*	00010120	
	?		TAU))-(K+CO*VE)/(CO*TAU)*(IVPD*(COS(CO*TAU)-1)+RVPD*	00010130	
	?		SIN(CO*TAU))	00010140	
ISN 0564	C	RMPPE =	RMPPE/SOSEX**2	00010150	
ISN 0565	C	IMPPE =	IMPPE/SOSEX**2	00010160	
ISN 0566	C	RMSPE =	(GAMMA-1)/(2*SOSEX*DE)*RMPPE	00010170	
ISN 0567	C	IMSPE =	(GAMMA-1)/(2*SOSEX*DE)*IMPPE	00010180	
	C	IRROTATIONAL VELOCITY PERTURBATIONS AT THE EXIT			00010190
	C			00010200	
ISN 0568	C	RITVPE =	1/TAU*(RVPD*(COS(CO*TAU)-1)-IVPD*SIN(CO*TAU))	00010210	
ISN 0569	C	IITVPE =	1/TAU*(IVPD*(COS(CO*TAU)-1)+RVPD*SIN(CO*TAU))	00010220	
ISN 0570	C	RIAVPE =	1/(CO*TAU)*(B2R*RVPD*(COS(CO*TAU)-1)-B2I*IVPD*	00010230	
	?		(COS(CO*TAU)-1)-IVPD*B2R*SIN(CO*TAU)-B2I*RVPD*	00010240	
	?		SIN(CO*TAU))	00010250	
ISN 0571	C	IIAVPE =	1/(CO*TAU)*(B2R*IVPD*(COS(CO*TAU)-1)+B2I*RVPD*	00010260	
	?		(COS(CO*TAU)-1)+RVPD*B2R*SIN(CO*TAU)-B2I*IVPD*	00010270	
	?		SIN(CO*TAU))	00010280	
	C	ROTATIONAL VELOCITY PERTURBATIONS AT THE EXIT			00010290
	C			00010300	
ISN 0572	C	RRAVPE =	1/((R**2+CO**2)*TAU)*(RRFF*(COS(CO*TAU)-1)-IRFF*	00010310	
	?		SIN(CO*TAU))	00010320	
ISN 0573	C	IRAVPE =	1/((R**2+CO**2)*TAU)*(IRFF*(COS(CO*TAU)-1)-RRFF*	00010330	
	?		SIN(CO*TAU))	00010340	
ISN 0574	C	RRTVPE =	-R/CO*RAVPE	00010350	
ISN 0575	C	IRTVPE =	-R/CO*IRAVPE	00010360	
ISN 0576	C	IP1U =	(IMPPI+IMPPI)/2.0	00010370	
ISN 0577	C	RPIU =	(RMPPI+RMPPI)/2.0	00010380	
ISN 0578	C	IP2U =	(IMPPI+IMPPI)/2.0	00010390	
ISN 0579	C	RP2U =	(RMPPI+RMPPI)/2.0	00010400	
ISN 0580	C	IP3U =	(IPDS+IMPPE)/2.0	00010410	
ISN 0581	C	RP3U =	(RPDS+RMPPE)/2.0	00010420	
ISN 0582	C	RP1L =	RP1U*COS(SIGMAT)+IPIU*SIN(SIGMAT)	00010430	
				00010440	
				00010450	
				00010460	
				00010470	
				00010480	
				00010490	
				00010500	
				00010510	
				00010520	
				00010530	
				00010540	
				00010550	
				00010560	
				00010570	

ISH 0583	IP1L	=	IP1U*COS(SIGMAT)-RP1U*SIN(SIGMAT)	00010530	
ISH 0584	RP2L	=	RP2U*COS(SIGMAT)+IP2U*SIN(SIGMAT)	00010590	
ISH 0595	IP2L	=	IP2U*CCS(SIGMAT)-RP2U*SIN(SIGMAT)	00010600	
ISH 0586	RP3L	=	RP2U*CCS(SIGMAT)+IP3U*SIN(SIGMAT)	00010610	
ISH 0587	IP3L	=	IP3U*CCS(SIGMAT)-RP3U*SIN(SIGMAT)	00010620	
	C				
ISH 0588	RCLU	=	(RP1U*XSTAR+RP2U*(XS-XSTAR)+RP3U*(C-XS))/B	00010630	
ISH 0589	ICLU	=	(IP1U*XSTAR+IP2U*(XS-XSTAR)+IP3U*(C-XS))/B	00010640	
	C				
ISH 0590	RCLL	=	(RP1L*(XSTAR-DELTA*COS(ALPCH))+RP2L*(XS-XSTAR)+ RP3L*(C-XS+DELTA*COS(ALPCH)))/B	00010660	
ISH 0591	ICLL	=	(IP1L*(XSTAR-DELTA*COS(ALPCH))+IP2L*(XS-XSTAR)+ IP3L*(C-XS+DELTA*COS(ALPCH)))/B	00010670	
	C				
ISH 0592	RCL	=	-RCLU+PCLL	00010680	
ISH 0593	ICL	=	-ICLU+ICLL	00010690	
	C				
ISH 0594	RCMU	=	-(RP1U*XSTAR*(Z*C-XSTAR/2)+RP2U*(XS-XSTAR)*(Z*C- (XSTAR+XS)/2)+RP3U*((C-XS)*(Z*C-(C+XS)/2)))/(B**2)	00010700	
ISH 0595	ICMU	=	-(IP1U*XSTAR*(Z*C-XSTAR/2)+IP2U*(XS-XSTAR)*(Z*C- (XSTAR+XS)/2)+IP3U*((C-XS)*(Z*C-(C+XS)/2)))/(B**2)	00010710	
	C				
ISH 0596	RCML	=	(RP1L*(XSTAR-DELTA*COS(ALPCH))*(Z*C-(XSTAR-DELTA* COS(ALPCH))/2)+RP2L*(XS-XSTAR)*(Z*C-(XS+XSTAR-2*DELTA* COS(ALPCH))/2)+RP3L*(C-XS+DELTA*COS(ALPCH))*(Z*C- (XS+C-DELTA*COS(ALPCH))/2)))/(B**2)	00010720	
ISH 0597	ICML	=	(IP1L*(XSTAR-DELTA*COS(ALPCH))*(Z*C-(XSTAR-DELTA* COS(ALPCH))/2)+IP2L*(XS-XSTAR)*(Z*C-(XS+XSTAR-2*DELTA* COS(ALPCH))/2)+IP3L*(C-XS+DELTA*COS(ALPCH))*(Z*C- (XS+C-DELTA*COS(ALPCH))/2)))/(B**2)	00010730	
	C				
ISH 0598	RCM	=	RCMU+RCML	00010740	
ISH 0599	ICM	=	ICMU+ICML	00010750	
	C				
ISH 0600			IF (II.EQ.1) GOTO 33	00010760	
ISH 0602	AAR(2)	=	(-RCL/(PI*K**2*ALPBAR))/12	00010770	
ISH 0603	AAI(2)	=	(-ICL/(PI*K**2*ALPBAR))/12	00010780	
ISH 0604	DAR(2)	=	RCM/(PI*K**2*ALPBAR)/12	00010790	
ISH 0605	BAI(2)	=	ICM/(PI*K**2*ALPBAR)/12	00010800	
ISH 0606	AARK	=	(AAR(2)*K**2)	00010810	
ISH 0607	AAIK	=	(AAI(2)*K**2)	00010820	
ISH 0608	BARK	=	(BAR(2)*K**2)	00010830	
ISH 0609	BAIK	=	(BAI(2)*K**2)	00010840	
ISH 0610	AHRK	=	(AHR(2)*K**2)	00010850	
ISH 0611	AHIK	=	(AHI(2)*K**2)	00010860	
ISH 0612	BHRK	=	(BHR(2)*K**2)	00010870	
ISH 0613	BHIK	=	(BHI(2)*K**2)	00010880	
ISH 0614			GOTO 34	00010890	
	C				
ISH 0615	33	AHR(2)	=	(S*RCL/(PI*K**2*H))/12	00010900
ISH 0616		AHI(2)	=	(S*ICL/(PI*K**2*H))/12	00010910
ISH 0617		BHR(2)	=	-B*RCM/(PI*K**2*H)/12	00010920
ISH 0618		BHI(2)	=	-B*ICM/(PI*K**2*H)/12	00010930
	C				
ISH 0619	34		WRITE(6,6666) LC	00010940	
ISH 0620	6666		FORMAT('-', ' LC', 10(E8.2,3X))	00010950	
ISH 0621			WRITE(6,440) RMAP1,IMAPI,RHAPT,IMAPT,RMAP,IMAP	00010960	

ISN 0657	570	FORMAT('-',T11,'RHVPU',T28,'IMVPU',T45,'RHITVPI',T62,'IMITVPI', T79,'RMIAPVI',T96,'IMIAVPI'/4X,6(F15.6,2X))	00011760 00011770 00011780
ISN 0658	580	FORMAT('I',T11,'RMIAPVE',T28,'IMIAVPE',T45,'RHITVPE',T62, 'IMITVPE',T79,'RMIAPVE',T96,'IMIAVPE'/4X,6(F15.6,2X))	00011790 00011800 00011810
ISN 0659	590	FORMAT('-',T11,'RHRTVPE',T28,'IMRTVPE'/4X,2(F15.6,2X))	00011820 00011830
ISN 0660	600	FORMAT('-',T11,'RUDS',T28,'IUDS',T45,'WDSR',T62,'WDSI' /4X,4(F15.6,2X))	00011840 00011850 00011860
ISN 0661	610	FORMAT('-',T11,'CI',T28,'CR',T45,'IC',T62,'RC' /4X,4(F15.6,2X))	00011870 00011880 00011890
ISN 0662	620	FORMAT('-',T11,'IRFF',T28,'IRFF'/4X,2(F15.6,2X))	00011900 00011910
ISN 0663	630	FORMAT('-',T11,'RPIU',T28,'IPIU',T45,'RPU',T62,'IPU', T79,'RPSU',T96,'IPSU'/4X,6(F15.6,2X))	00011920 00011930 00011940
ISN 0664	640	FORMAT('-',T11,'RPIL',T28,'IPI',T45,'RP2L',T62,'IP2L', T79,'RP3L',T96,'IP3L'/4X,6(F15.6,2X))	00011950 00011960 00011970
ISN 0665	650	FORMAT('-',T11,'RCLU',T28,'ICLU',T45,'RCLL',T62,'ICLL', T79,'RCL',T96,'ICL'/4X,6(F15.6,2X))	00011980 00011990
ISN 0666	660	FORMAT('-',T11,'RCMU',T28,'ICMU',T45,'RCML',T62,'ICML', T79,'RCM',T96,'ICM'/4X,6(F15.6,2X))	00012000 00012010 00012020
ISN 0667	15	CONTINUE	00012030
ISN 0668		WRITE(6,670) AAR(2), AAI(2), BAR(2), BAI(2)	00012040
ISN 0669		WRITE(6,680) AHR(2), AHI(2), BHR(2), BHI(2)	00012050
ISN 0670		WRITE(6,690) AARK, AAIK, BARK, BAIK	00012060
ISN 0671		WRITE(6,700) AHRK, AHK, BHRK, BHK	00012070 00012080
ISN 0672	670	FORMAT('-',T11,'AAR',T28,'AAI',T45,'BAR',T62,'BAI' /4X,4(F15.6,2X))	00012090 00012100
ISN 0673	680	FORMAT('-',T11,'AHR',T28,'AHI',T45,'BHR',T62,'BHI' /4X,4(F15.6,2X))	00012110 00012120
ISN 0674	690	FORMAT('-',T11,'AARK',T28,'AAIK',T45,'BARK',T62,'BAIK' /4X,4(F15.6,2X))	00012130 00012140
ISN 0675	700	FORMAT('-',T11,'AHRK',T28,'AHK',T45,'BHRK',T62,'BHK' /4X,4(F15.6,2X))	00012150 00012160 00012170
ISN 0676		RETURN	00012180
ISN 0677		END	00012190

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL	STATEMENT	CROSS	REFERENCE	LISTING	NUMBERS													
B	0039	0011	0017	0155	0157	0440	0441	0442	0443	0446	0447	0459	0460	0461	0462	0465	0466	0503	
	0503	0504	0506	0506	0507	0512	0512	0513	0513	0514	0515	0588	0589	0590	0591	0594	0595	0596	0597
	0515	0616	0617	0618															
C	0005	0009	0017	0021	0430	0431	0432	0432	0433	0433	0434	0434	0435	0435	0436	0437	0440	0441	0442
	0442	0443	0443	0446	0446	0447	0447	0503	0506	0512	0512	0512	0512	0513	0513	0513	0513	0514	0515
	0508	0539	0590	0591	0594	0594	0594	0594	0594	0595	0595	0595	0595	0595	0596	0596	0596	0596	0596
	0597	0597	0597	0597	0597														
	0005	0017																	
H	0005	0449	0450	0451	0452	0453	0454	0455	0456	0459	0460	0461	0461	0462	0462	0465	0466	0503	0504
	0506	0507	0615	0616	0617	0618													
I	0004	0020	0165	0165	0165	0166	0166	0166	0178	0179	0188	0198	0198	0199	0199	0202	0238	0239	0244
	0264																		

*****F O R T R A N C R O S S R E F E R E N C E L I S T I N G*****

SYMBOL		INTERNAL STATEMENT NUMBERS																	
		0251	0251	0251	0252	0252	0252	0253	0253	0253	0254	0254	0254	0255	0255	0255	0255	0257	0257
		0257	0257	0259	0259	0260	0260	0261	0261	0261	0261	0262	0262	0262	0262	0263	0263	0263	0263
		0264	0264	0264	0265	0265	0265	0266	0266	0266	0267	0267	0268	0268	0269	0269	0269	0269	0270
		0270	0270	0271	0271	0271	0272	0272	0272	0273	0273	0273	0274	0274	0274	0274	0276	0276	0276
		0276	0278	0278	0279	0279	0280	0281	0281	0281	0282	0282	0282	0282	0283	0283	0283	0283	0284
		0284	0284	0284	0285	0285	0285	0286	0286	0286	0286	0287	0287	0287	0287	0287	0288	0288	0288
		0288	0289	0289	0289	0289	0289	0291	0291	0291	0291	0291	0409	0467	0467	0467	0467	0467	0467
		0468	0468	0468	0468	0468													
FE		0003	0163	0164	0165	0165	0165	0166	0166	0166	0167	0168	0169	0170	0170	0170	0170	0170	0170
		0170	0171	0171	0171	0171	0171	0171	0171	0171	0172	0172	0172	0172	0172	0172	0172	0172	0172
		0173	0173	0173	0173	0173	0173	0174	0175	0176	0177	0177	0177	0177	0178	0178	0178	0178	0178
		0178	0179	0179	0179	0179	0179	0180	0180	0180	0180	0180	0180	0180	0180	0180	0180	0181	0181
		0181	0181	0181	0181	0181	0181	0181	0181	0181	0181	0181	0182	0182	0182	0182	0182	0182	0182
		0182	0183	0183	0183	0183	0183	0184	0184	0184	0184	0184	0184	0184	0184	0184	0184	0184	0185
		0185	0185	0185	0185	0185	0185	0185	0185	0185	0185	0185	0185	0186	0186	0186	0186	0186	0186
		0187	0187	0187	0187	0187	0187	0188	0188	0188	0188	0188	0188	0189	0189	0190	0190	0190	0190
		0190	0190	0191	0191	0191	0191	0191	0191	0191	0192	0192	0192	0192	0193	0193	0193	0193	0194
		0194	0194	0194	0194	0194	0195	0195	0195	0195	0195	0195	0195	0196	0196	0196	0197	0197	0197
		0197	0204	0204	0204	0205	0205	0205	0206	0206	0206	0207	0207	0207	0208	0208	0209	0209	0210
		0211	0211	0212	0212	0213	0213	0214	0214	0215	0215	0216	0216	0217	0218	0218	0219	0219	0220
		0220	0221	0222	0222	0222	0223	0224	0224	0225	0225	0226	0226	0227	0227	0227	0227	0227	0227
		0243	0243	0244	0244	0245	0245	0246	0246	0247	0247	0248	0248	0249	0249	0250	0250	0250	0251
		0251	0252	0252	0253	0253	0261	0261	0261	0262	0262	0262	0263	0263	0264	0264	0265	0265	0266
		0267	0267	0268	0268	0269	0269	0270	0270	0271	0271	0272	0272	0405	0471	0471	0471	0471	0471
FM		0471	0471	0471	0471	0472	0472	0472	0472	0472	0472	0472	0472	0472	0472				
		0003	0158	0199	0200	0201	0202	0203	0204	0204	0204	0204	0204	0205	0205	0205	0205	0206	0206
		0206	0206	0206	0207	0207	0207	0207	0207	0208	0208	0208	0209	0209	0209	0210	0210	0210	0210
		0211	0211	0212	0212	0212	0212	0213	0213	0214	0214	0214	0214	0215	0215	0216	0216	0216	0217
		0217	0217	0218	0218	0218	0219	0219	0219	0220	0220	0220	0220	0221	0221	0221	0222	0222	0222
		0223	0223	0223	0224	0224	0224	0225	0225	0225	0226	0226	0227	0227	0227	0228	0228	0228	0228
		0228	0229	0229	0231	0231	0231	0231	0233	0233	0233	0233	0234	0234	0234	0235	0236	0236	0236
		0236	0254	0254	0255	0255	0257	0257	0273	0273	0274	0274	0276	0276	0408	0469	0469	0469	0469
		0469	0469	0469	0470	0470	0470	0470	0470	0470	0470	0470							
IC		0004	0528	0555	0559	0560	0561	0638											
II		0427	0428	0501	0600														
LC		0003	0004	0415	0530	0531	0532	0533	0534	0535	0535	0535	0535	0535	0535	0536	0536	0536	0536
		0537	0537	0538	0538	0539	0539	0539	0539	0540	0540	0540	0540	0541	0541	0541	0542	0542	0542
		0543	0543	0543	0543	0544	0544	0544	0544	0544	0545	0545	0545	0546	0546	0547	0547	0547	0548
		0548	0548	0549	0549	0549	0550	0550	0550	0551	0551	0551	0551	0551	0552	0552	0552	0552	0552
		0553	0553	0553	0553	0553	0553	0554	0554	0554	0555	0555	0555	0555	0556	0556	0556	0556	0557
		0557	0557	0557	0557	0558	0558	0558	0558	0559	0559	0559	0619						
LH		0003	0004	0378	0379	0380	0381	0383	0384	0385	0386	0386	0386	0386	0386	0386	0386	0387	0387
		0387	0387	0387	0387	0387	0388	0388	0388	0389	0389	0389	0389	0389	0389	0390	0390	0390	0390
		0390	0390	0390	0391	0391	0391	0391	0391	0391	0391	0392	0392	0392	0392	0393	0393	0393	0393
		0394	0394	0394	0394	0395	0395	0395	0395	0395	0395	0396	0396	0396	0396	0396	0396	0397	0397
		0398	0398	0398	0398	0399	0399	0399	0399	0400	0400	0400	0400	0400	0401	0401	0401	0401	0402
		0402	0402	0402	0414	0535	0535	0535	0536	0536	0536	0537	0539	0539	0540	0541	0541	0542	0542
		0543	0543	0543	0544	0544	0544	0546	0547	0547	0548	0548	0549	0549	0550	0550	0560	0560	0560
		0561	0561	0561	0561														
MD		0004	0004	0045	0047	0047	0074	0080	0083										
MI		0004	0014	0026	0027	0054	0114												
MT		0004	0013	0025	0114														
MU		0004	0035	0044	0046	0046	0058	0062	0065										
PD		0029	0031	0115	0132	0132	0370	0528	0529										
PE		0085	0087	0115	0133	0133	0528	0529											
PI		0012	0602	0603	0604	0605	0615	0616	0618										

		*****FORTRAN CROSS REFERENCE LISTING*****																			
SYMBOL	INTERNAL STATEMENT NUMBERS																				
PT	0025 0115 0134 0134 0203 0256																				
FU	0024 0029 0115 0131 0131 0255 0369 0370 0370 0372																				
PC	0329 0555 0559 0530 0561 0633																				
TE	0036 0057																				
UE	0091 0146 0146 0378 0378 0330 0302 0530 0533 0562 0563																				
US	0093 0371																				
VD	0099 0100 0103 0140 0148 0522 0523 0524 0524 0525 0525 0528 0529																				
VE	0090 0145 0145 0379 0379 0301 0332 0331 0534 0562 0563																				
VU	0093 0107 0147 0147 0161 0293 0295 0326 0327 0328 0328 0329 0330 0331 0332 0333 0334 0335 0336																				
	0337 0338 0339 0340 0341 0342 0343 0343 0344 0345 0346 0347 0348 0349 0350 0351 0352 0353 0522																				
	0523																				
XS	0005 0434 0435 0446 0446 0447 0447 0465 0466 0503 0506 0512 0512 0513 0513 0588 0588 0589 0589																				
	0520 0590 0591 0591 0594 0594 0594 0594 0595 0595 0595 0595 0596 0596 0596 0596 0597 0597 0597																				
	0597																				
AAI	0003 0603 0607 0668																				
AAR	0003 0602 0606 0668																				
ACE	0005 0158 0158 0162 0378 0378 0378 0379 0379 0379 0303 0503 0506 0530 0530 0531 0531 0532																				
AHI	0003 0611 0616 0669																				
AHR	0003 0610 0615 0669																				
BAI	0003 0605 0609 0668																				
BAR	0003 0604 0603 0668																				
BHI	0003 0613 0616 0669																				
BHR	0003 0612 0617 0669																				
END	0041 0066 0071 0071 0072 0079 0079 0082 0082																				
ENU	0040 0048 0053 0053 0054 0061 0061 0064 0064																				
EYD	0043 0047																				
BYU	0042 0046																				
BII	0053 0062 0065 0170 0170 0171 0171 0172 0172 0173 0173 0204 0204 0205 0205 0206 0206 0207 0207																				
	0242 0242 0243 0243 0261 0261 0262 0262 0403 0491 0491 0492 0492 0499 0499 0500 0500																				
BIR	0054 0051 0054 0170 0170 0171 0171 0172 0172 0173 0173 0204 0204 0205 0205 0206 0206 0207 0207																				
	0242 0242 0243 0243 0261 0261 0262 0262 0403 0491 0491 0492 0492 0499 0499 0500 0500																				
BZI	0071 0080 0083 0326 0326 0327 0327 0390 0390 0391 0391 0535 0535 0536 0536 0543 0543 0544 0544																				
	0562 0562 0563 0563 0570 0570 0571 0571																				
BZR	0072 0079 0082 0326 0326 0327 0327 0390 0390 0391 0391 0535 0535 0536 0536 0543 0543 0544 0544																				
	0562 0562 0563 0563 0570 0570 0571 0571																				
COS	0087 0020 0034 0035 0036 0037 0038 0038 0039 0039 0090 0090 0091 0106 0106 0110 0111 0111 0163																				
	0166 0166 0168 0169 0199 0201 0239 0378 0378 0379 0379 0379 0381 0383 0383 0383 0384 0385																				
	0430 0432 0434 0436 0440 0440 0441 0442 0442 0442 0443 0443 0446 0449 0452 0453 0456 0459 0459																				
	0450 0451 0461 0461 0462 0465 0491 0491 0491 0492 0492 0497 0498 0499 0499 0500 0500 0503																				
	0504 0512 0514 0514 0515 0516 0528 0529 0530 0531 0531 0532 0562 0562 0562 0563 0563 0563 0568																				
	0539 0570 0573 0571 0571 0572 0573 0582 0583 0584 0585 0586 0587 0590 0590 0591 0591 0596 0596																				
	0566 0596 0596 0597 0597 0597 0597 0597																				
DID	0045 0056 0071 0072 0079 0079 0082 0082																				
DIU	0044 0048 0053 0054 0051 0061 0064 0064																				
DSD	0047 0066 0071 0079 0082																				
DZU	0046 0048 0053 0051 0054																				
FM2	0003 0293 0294 0295 0297 0297 0297 0297 0298 0298 0298 0299 0299 0299 0299 0299 0300 0300 0301																				
	0301 0301 0301 0302 0302 0303 0303 0303 0303 0304 0304 0304 0304 0305 0305 0305 0305 0306 0306																				
	0306 0306 0307 0307 0307 0308 0308 0309 0309 0309 0310 0310 0311 0311 0311 0311 0311 0312																				
	0312 0312 0312 0313 0313 0313 0314 0314 0314 0315 0315 0316 0316 0316 0316 0317 0317 0318																				
	0318 0318 0319 0319 0319 0320 0320 0320 0321 0321 0321 0322 0322 0322 0323 0323 0324 0324																				
	0324 0325 0325 0325 0326 0327 0328 0329 0330 0331 0332 0333 0334 0335 0336 0337 0338 0339 0340																				
	0341 0342 0343 0344 0345 0346 0347 0348 0349 0350 0351 0352 0353 0410 0478 0478 0478 0478 0478																				
	0478 0478 0478 0478 0478 0478 0478 0478 0478																				
FRE	0162 0378 0379 0383 0383																				
ZCL	0004 0593 0603 0616 0642																				

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL STATEMENT NUMBERS									
ICH	0004	0599	0605	0618	0543					
MDS	0004	0005	0030	0055	0099	0114				
MUS	0204	0005	0329	0097	0098	0114	0370	0371	0373	0375
MXD	0004	0036	0041	0045	0045	0080	0083			
MXU	0004	0037	0040	0044	0044	0062	0065			
MYD	0004	0004	0033	0043	0045	0047	0074	0080	0083	
MYU	0004	0004	0039	0042	0044	0046	0058	0062	0065	
PIN	0026	0028	0115	0130	0130	0202				
RCL	0592	0602	0615	0642						
RCH	0598	0604	0617	0543						
REV	0089	0090	0091	0092	0108	0109	0109	0378	0378	0379
SIN	0036	0037	0091	0110	0165	0165	0167	0170	0170	0170
	0193	0200	0204	0204	0204	0205	0205	0205	0206	0206
	0243	0243	0261	0261	0261	0262	0262	0262	0370	0378
	0307	0308	0309	0390	0390	0390	0391	0391	0392	0393
	0451	0454	0455	0460	0462	0462	0456	0491	0491	0491
	0504	0506	0507	0507	0513	0515	0530	0532	0535	0535
	0544	0544	0545	0562	0562	0562	0563	0563	0563	0568
	0584	0595	0595	0587						
SSI	0094	0101	0117	0139	0139	0163	0165	0167	0167	0167
	0240	0241	0268	0493	0494	0495	0495			
SST	0095	0101	0156	0107	0117	0140	0140	0164	0166	0180
	0203	0208	0214	0219	0225	0246	0248	0266	0296	0307
SSU	0097	0098	0117	0143	0143	0293	0371	0371	0372	0373
TAU	0010	0011	0011	0033	0165	0165	0165	0165	0166	0166
	0168	0169	0170	0170	0170	0171	0171	0171	0172	0172
	0199	0200	0200	0201	0201	0204	0204	0204	0205	0205
	0239	0239	0240	0241	0242	0242	0242	0243	0261	0261
	0378	0379	0379	0379	0320	0320	0351	0381	0383	0383
	0383	0389	0390	0390	0390	0391	0391	0391	0392	0393
	0492	0492	0492	0492	0492	0492	0492	0497	0497	0498
	0500	0500	0500	0500	0530	0530	0531	0531	0532	0533
	0543	0543	0544	0544	0544	0545	0562	0562	0562	0562
	0563	0563	0563	0568	0568	0568	0569	0569	0569	0570
	0572	0572	0572	0573	0573	0573				
TDS	0030	0031								
TIN	0027	0028								
USI	0483	0490	0629							
UCR	0487	0489	0629							
VEL	0111	0153	0153	0166	0166	0168	0168	0168	0199	0201
VII	0445	0464	0467	0468	0469	0470	0471	0472	0631	
VIR	0444	0463	0467	0468	0469	0470	0471	0472	0631	
V2I	0447	0466	0475	0476	0477	0478	0631			
V2R	0446	0465	0475	0476	0477	0478	0631			
V3I	0520	0527	0529	0632						
V3R	0521	0526	0528	0632						
WDS	0100	0528	0529							
AAIK	0607	0670								
AARK	0606	0670								
AHIK	0611	0671								
AMRK	0610	0671								
BAIK	0609	0670								
BASK	0608	0670								
BHIK	0613	0671								
BHRK	0612	0671								
DTOT	0019	0032	0102	0126	0126					

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL STATEMENT NUMBERS																		
VOL2	0005	0156	0156	0294	0304	0305	0306	0307	0308	0309	0319	0320	0322	0326	0327	0328	0329	0330	0331
	0332	0333	0334	0335	0336	0337	0338	0339	0340	0341	0342	0343	0344	0344	0345	0346	0347	0348	0349
	0350	0351	0352	0353															
VOL3	0005	0157	0157	0303	0380	0381	0381	0384	0526	0527	0528	0528	0529	0529	0533	0534			
VIAI	0439	0445	0453	0464	0630														
VIAI	0439	0444	0457	0463	0630														
VIBI	0441	0445	0450	0464	0630														
VIER	0440	0444	0459	0463	0630														
VICI	0443	0445	0462	0464	0630														
VICR	0442	0444	0451	0453	0630														
V3AI	0506	0513	0520	0631															
V3AR	0503	0512	0521	0631															
V3BI	0507	0515	0520	0632															
V3BR	0504	0514	0521	0632															
V3CI	0500	0517	0520	0632															
V3CR	0505	0516	0521	0632															
WDSI	0524	0526	0523	0637															
WDSR	0525	0527	0529	0637															
ALPCH	0005	0007	0020	0034	0035	0036	0036	0037	0037	0038	0038	0039	0039	0090	0090	0091	0091	0106	0108
	0110	0110	0111	0111	0163	0165	0165	0166	0166	0167	0168	0198	0199	0200	0201	0238	0239	0378	0378
	0378	0373	0379	0379	0379	0379	0330	0381	0393	0393	0393	0383	0383	0394	0384	0440	0441	0442	0442
	0443	0443	0459	0460	0461	0461	0462	0462	0504	0507	0514	0515	0526	0528	0529	0530	0530	0531	0531
	0532	0532	0590	0590	0591	0591	0596	0596	0596	0596	0596	0597	0597	0597	0597				
ASTAR	0005	0100	0155	0156	0157	0153	0159	0160	0430	0431	0432	0433	0434	0435	0436	0437	0440	0441	0442
	0443	0445	0447	0449	0450	0451	0452	0453	0454	0455	0456	0459	0460	0461	0461	0462	0462	0465	0466
	0503	0503	0504	0506	0506	0507	0512	0512	0513	0513	0514	0515							
AVGD1	0103	0127	0127	0167	0167	0168	0168	0175	0175	0176	0176	0200	0201	0215	0216	0224	0227	0252	0272
AVGD2	0104	0128	0128	0304	0305	0306	0307	0310	0311	0319	0320	0324	0338	0351					
AVGD3	0105	0129	0129	0300	0381	0384	0526	0527	0528	0528	0529	0529							
AVGV1	0103	0149	0149	0167	0167	0168	0168	0174	0175	0176	0176	0200	0201	0209	0216	0218	0227	0406	
AVGV2	0107	0150	0150	0294	0308	0309	0310	0311	0322	0324									
AVGV3	0108	0151	0151	0300	0381	0528	0528	0529	0529										
BETA1	0005	0020	0035	0037	0106	0110	0111	0163											
BETA2	0005	0007	0034	0036	0038	0039	0090	0091	0108	0378	0378	0379	0379	0383	0526	0528	0529	0530	0531
COTAN	0451	0452																	
DELTA	0005	0010	0440	0441	0442	0442	0443	0443	0459	0460	0461	0462	0504	0507	0514	0515	0590	0590	0591
	0591	0596	0596	0596	0596	0596	0597	0597	0597	0597	0597								
FM23A	0296	0302	0303	0316															
GAMMA	0015	0024	0024	0024	0025	0025	0025	0026	0026	0026	0027	0029	0029	0030	0032	0032	0065	0085	0085
	0026	0033	0033	0094	0094	0095	0095	0096	0096	0097	0097	0102	0102	0163	0163	0164	0164	0164	0164
	0165	0166	0167	0167	0167	0168	0168	0168	0174	0174	0175	0175	0176	0176	0369	0369	0370	0370	0370
	0370	0371	0371	0371	0371	0372	0372	0373	0374	0374	0375	0483	0484	0495	0496	0566	0567		
IMAPI	0004	0431	0455	0467	0468	0469	0470	0471	0472	0621									
IMAPT	0004	0433	0454	0467	0468	0469	0470	0471	0472	0475	0476	0477	0478	0621					
IMPPE	0004	0565	0628																
IMPPI	0004	0494	0628																
IMPPU	0004	0478	0480	0482	0523	0627													
IMPPE	0004	0563	0565	0567	0590	0626													
IMPPI	0004	0492	0494	0496	0576	0626													
IMPPU	0004	0482	0484	0488	0578	0625													
IMPPI	0004	0476	0477	0478	0480	0488	0523	0634											
IOSSU	0004	0484	0488	0623															
MEXIT	0004	0007	0085	0086	0088	0114													
OMEGA	0005	0021																	
PRES1	0005	0136	0136																
PRES2	0005	0137	0137																

*****FORTRAN CROSS REFERENCE LISTING*****

SYMBOL	INTERNAL STATEMENT NUMBERS
FRES3	0005 0130 0130
QUANT	0048 0049 0058 0059 0066 0057 0074 0077
FMAPI	0430 0453 0467 0468 0469 0470 0471 0472 0621
FMAPT	0432 0456 0467 0468 0469 0470 0471 0472 0475 0476 0477 0478 0621
ENDPE	0534 0503
ENDPI	0493 0528
ENDFU	0477 0479 0481 0522 0627
RMPE	0562 0564 0566 0581 0626
RMPI	0491 0493 0495 0577 0626
RMFU	0431 0483 0487 0579 0625
RMVPU	0475 0477 0478 0479 0487 0522 0634
RSOSU	0403 0487 0623
SOSES	0096 0099 0117 0142 0142 0435 0486
SOSEX	0023 0039 0117 0141 0141 0378 0379 0380 0381 0530 0531 0533 0534 0564 0565 0566 0567
UIIRE	0005 0921 0100 0106 0109 0110 0111 0130 0131 0132 0133 0134 0135 0136 0137 0138 0139 0140 0141
	0142 0143 0144 0145 0146 0147 0148 0149 0150 0151 0152 0153 0154
UIIFE	0072 0154 0134 0162 0527
VELAX	0110 0152 0152 0165 0165 0167 0167 0167 0198 0198 0200 0238 0240 0491 0492
XSTAR	0305 0432 0433 0442 0442 0443 0443 0446 0446 0447 0447 0461 0462 0462 0588 0588 0589 0589
	0590 0590 0591 0591 0594 0594 0594 0595 0595 0595 0595 0596 0596 0596 0596 0597 0597 0597
	0597
ALFAR	0205 0430 0431 0432 0433 0434 0435 0436 0437 0440 0441 0442 0443 0446 0447 0512 0513 0514 0515
	0402 0603 0604 0605
AVGOS	0101 0144 0144 0167 0167 0168 0168 0174 0175 0176
IIAVPE	0304 0571 0635
IIAVPI	0004 0500 0634
IIIVPE	0004 0569 0635
IIIVPI	0004 0493 0634
INSCPE	0004 0567 0624
INSCPI	0004 0496 0623
IRAVPE	0004 0573 0575 0635
IRIVPE	0004 0575 0635
IMOCK	0004 0095 0032
IMSTAR	0004 0016 0095 0102
RIAVPE	0570 0635
RIAVPI	0499 0434
RIIVPE	0568 0635
RITVPI	0497 0634
RMSPE	0566 0624
RMSPI	0495 0623
RRAVPE	0572 0574 0635
RRTVPE	0574 0636
SIGMAB	0305 0033 0449 0450 0451 0452 0453 0454 0455 0456 0459 0460 0461 0461 0462 0462 0465 0466 0503
	0504 0506 0507
SIGHAT	0005 0430 0431 0432 0433 0434 0435 0436 0437 0440 0441 0442 0443 0446 0447 0512 0513 0514 0515
	0582 0582 0583 0583 0584 0584 0585 0585 0586 0586 0587 0587

*****FORTRAN CROSS REFERENCE LISTING*****

LABEL	DEFINED	REFERENCES
10	0232	0230
13	0449	0428
15	0667	0427
17	0467	0448
20	0237	0235
24	0512	0501
25	0510	0511

*****FORTRAN CROSS REFERENCE LISTING*****

LABEL	DEFINED	REFERENCES
30	0258	0256
33	0515	0600
34	0619	0614
40	0277	0275
50	0292	0290
60	0357	0355
70	0368	0366
71	0052	0051
75	0057	0056
76	0070	0069
78	0076	0075
80	0056	0049
84	0064	0359
90	0065	0355 0063
91	0074	0057
94	0032	0077
95	0023	0022
99	0034	0073 0081
100	0113	0112
105	0118	0114
110	0119	0116
115	0120	0115
120	0121	0117
333	0404	0403
410	0510	0509
420	0519	0518
440	0644	0621
450	0645	0622
460	0646	0623
470	0647	0624
480	0648	0625
490	0649	0626
500	0650	0627
510	0651	0628
520	0652	0629
530	0653	0630
540	0654	0631
550	0655	0632
560	0656	0633
570	0657	0634
580	0658	0635
590	0659	0636
600	0660	0637
610	0661	0638
620	0662	0639
630	0663	0640
640	0664	0641
650	0665	0642
660	0666	0643
670	0672	0668
680	0673	0669
690	0674	0670
700	0675	0671
1000	0008	0006
6000	0416	0405
6001	0419	0408

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***** FORTRAN CROSS REFERENCE LISTING*****

LABEL	DEFINED	REFERENCES
6002	0420	0409
6003	0421	0410
6004	0422	0411
6005	0423	0412
6006	0424	0413
6007	0425	0414
6008	0426	0415
6100	0417	0406
6200	0418	0407
6666	0620	0619

/ UNST / SIZE_OF_PROGRAM,006999,HEXADECIMAL,BYTES

NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.	NAME	TAG	TYPE	ADD.
B*SF		R*4	000D50	C F C	R*4	000044	E F C	R*4	000014	H F C	R*4	H F C	R*4	R*4	000064
I SF		R*4	000D54	J SF		I*4	000053	K SF		R*4	000D5C	M SF		R*4	00105C
R SF		R*4	000D60	S SF		R*4	0010CC	Z SF		R*4	000D64	AC SF	C	R*4	000018
AO SF	C	R*4	00001C	SI SF		R*4	000D68	CL SF		R*4	0010F0	CO SFA		R*4	000D6C
CR SF		R*4	000D70	DD SF		R*4	000D74	CE SF		R*4	000D78	DI SF		R*4	000D7C
DT SF		R*4	000D80	DU SF		R*4	000D84	FC SF		R*4	00118C	FE SF		R*4	001270
FM SF		R*4	0012FC	IC SF		R*4	000D88	II SF		R*4	000D8C	LC SF		R*4	001388
LM SF		R*4	001428	IO SF		R*4	000D90	MI SFA		R*4	00CD94	MT SF		R*4	000D98
LU SF		R*4	000D9C	FD SF		R*4	000DA0	PE SF		R*4	00CDA4	PI SF		R*4	000DA8
PT SF		R*4	000DAC	PU SF		R*4	000DB0	RC SF		R*4	000DB4	TE SF		R*4	000DB8
UE SF		R*4	000DBC	US SF		R*4	000DC0	VD SF		R*4	000DC4	VE SF		R*4	000DC8
VU SF		R*4	000DCC	XS F	C	R*4	000D30	AAI SF		R*4	001488	AAR SF		R*4	001490
ACE SF	C	R*4	000D0C	AHI SF		R*4	001493	AHR SF		R*4	0014A0	BAI SF		R*4	0014A8
BAR SF		R*4	0014D0	EHI SF		R*4	0014B3	BHR SF		R*4	0014C0	BXD SFA		R*4	000D00
BYU SFA		R*4	000DD4	EYD SF		R*4	000DD3	BYU SF		R*4	000DDC	BLI SF		R*4	000DE0
BLR SF		R*4	000DE4	E2I SF		R*4	000DE8	B2R SF		R*4	000DEC	COS F	XF	R*4	000000
D1D SFA		R*4	000DF0	D1U SFA		R*4	000DF4	D2D SFA		R*4	000DF8	DCU SFA		R*4	000DFC
FM2 SF		R*4	0014C8	FRE SF		R*4	000E00	ICL SF		R*4	000E04	ICH SF		R*4	000E08
HDS FA	C	R*4	000E24	HXS FA	C	R*4	000E28	HXD SF		R*4	000E0C	HXU SF		R*4	000E10
HYD SF		R*4	000E14	MYU SF		R*4	000E18	PIN SF		R*4	000E1C	RCL SF		R*4	000E20
RCM SF		R*4	000E24	REV SF		R*4	000E28	SIH F	XF	R*4	000E00	SGI SF		R*4	000E2C
SST SF		R*4	000E30	SSU SF		R*4	000E34	TAU SFA		R*4	000E38	TOS SF		R*4	000E3C
TIN SF		R*4	000E40	USI SF		R*4	000E44	USR SF		R*4	000E48	VEL SF		R*4	000E4C
V1I SF		R*4	000E50	V1R SF		R*4	000E54	V2I SF		R*4	000E58	V2R SF		R*4	000E5C
V3I SF		R*4	000E60	V3R SF		R*4	000E64	WDS SF		R*4	000E68	AAIK SF		R*4	000E6C
AARK SF		R*4	000E70	AHIK SF		R*4	000E74	AHRK SF		R*4	000E78	BAIK SF		R*4	000E7C
BACK SF		R*4	000E80	BHIK SF		R*4	000E84	BHRK SF		R*4	000E88	DTOT SF		R*4	000E8C
FLCN SF		R*4	000E90	IACE SF		R*4	000E94	ICLL SF		R*4	000E98	ICLU SF		R*4	000E9C
ICML SF		R*4	000EA0	ICNU SF		R*4	000EA4	IODS SF		R*4	000EA8	IMAP SF		R*4	000EAC
INDP SF		R*4	000EA0	IHPF SF		R*4	000EE4	IMSP SF		R*4	000EB8	IPDS SF		R*4	000EDC
IP1L SF		R*4	000EC0	IFIU SF		R*4	000EC4	IP2L SF		R*4	000EC8	IP2U SF		R*4	000ECC
IP3L SF		R*4	000ED0	IP2U SF		R*4	000ED4	IRFF SF		R*4	000ED8	ISSP SF		R*4	000EDC
IUDS SF		R*4	000EE0	IVFD SF		R*4	000EE4	IVFU SF		R*4	000EE8	NCPT		R*4	NR
M1RE F	C	R*4	000E48	MCRE SF		R*4	000EEC	PTOT SF	C	R*4	00004C	QUOT SF		R*4	000EF0
RACE SF		R*4	000EF4	RCLL SF		R*4	000EF8	RCLU SF		R*4	000EFC	RCHL SF		R*4	000EF0
RCMU SF		R*4	000F04	RDOS SF		R*4	000F08	RGAS SFA		R*4	000F0C	RMAP SF		R*4	000F10
RNDP SF		R*4	000F14	RKPP SF		R*4	000F18	RKSP SF		R*4	000F1C	RPDS SF		R*4	000F20
RP1L SF		R*4	000F24	RP1U SF		R*4	000F28	RP2L SF		R*4	000F2C	RP2U SF		R*4	000F30
RP3L SF		R*4	000F34	RP3U SF		R*4	000F38	RFFP SF		R*4	000F3C	RSSP SF		R*4	000F40
RUDS SF		R*4	000F44	RVPD SF		R*4	000F48	RVFU SF		R*4	000F4C	SQRT F	XF	R*4	000000
TTOT FA	C	R*4	00002C	UHST		R*4	000F50	VOLL SF	C	R*4	000000	VOL2 SF	C	R*4	000004
VOL3 SF	C	R*4	000008	V1AI SF		R*4	000F54	V1AR SF		R*4	000F58	V1BI SF		R*4	000F5C

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*VERSION 1.3.0 (01 MAY 80)      UNST      SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)      DATE 80.353/15.02.49      PAGE 31
VICR SF R*4 000F60      VICI SF R*4 000F64      VICR SF R*4 000F68      V3AI SF R*4 000F6C
V3AR SF R*4 000F70      V3DI SF R*4 000F74      VICR SF R*4 000F78      V3CI SF R*4 000F7C
V3CR SF R*4 000F80      WDSI SF R*4 000F84      WDS3 SF R*4 000F88      ALRCH FA C R*4 000053
ASTAR F C R*4 000010      AVGD1 SF R*4 000F8C      AVGD2 SF R*4 000F90      AVGD3 SF R*4 000F94
AVGV1 SF R*4 000F98      AVGV2 SF R*4 000F9C      AVGV3 SF R*4 000FA0      BETAL FA C R*4 000050
BETA2 FA C R*4 000054      COTAN F XF R*4 000000      DELTA F C R*4 000068      FHC3A SF R*4 000F44
GAMMA SFA R*4 000FAS      IMAPI SF R*4 000FAC      IMAPT SF R*4 000FD0      IMPPE SF R*4 000FD4
INDPI SF R*4 000FD3      INDPU SF R*4 000FDC      IMPPE SF R*4 000F00      IMPPI SF R*4 000FC4
INPPI SF R*4 000FC8      INVPU SF R*4 000FCC      ISCSI SF R*4 000FD0      MEXIT SFA R*4 000FD4
OMEGA F C R*4 00003C      PRES1 SF C R*4 000038      PRES2 SF C R*4 00003C      PRES3 SF C R*4 000040
QUANT S R*4 000FD8      RHAPI SF R*4 000FDC      RHAPT SF R*4 000FE0      RNDPE SF R*4 000FE4
RNDPI SF R*4 000FE8      RNDPU SF R*4 000FEC      RMPPE SF R*4 000FF0      RMPPI SF R*4 000FF4
RHPFU SF R*4 000FF8      RHPFU SF R*4 000FFC      RPSU SF R*4 001000      S3SDS SF R*4 001004
SOSEX SF R*4 001008      ULIRE F C R*4 000074      UZIRE SF R*4 00100C      VELAX SF R*4 001010
XSTAR F C R*4 000034      FRXCF* XF R*4 000000      ALFBAR F C R*4 000070      AVGSOS SF R*4 001014
ISCCM# F XF R*4 000000      IIAVPE SF R*4 001018      IIAVPI SF R*4 00101C      IITVPE SF R*4 001020
IITVPI SF R*4 001024      IHSSPE SF R*4 001028      IHSSPI SF R*4 00102C      IRAVPE SF R*4 001030
IRTVPE SF R*4 001034      MSHOCK F C R*4 000020      MXSTAR SFA R*4 001038      RIAVPE SF R*4 00103C
RIAVPI SF R*4 001040      RITVPE SF R*4 001044      RITVPI SF R*4 001048      RHSSPE SF R*4 00104C
RHSSPI SF R*4 001050      RRAVPE SF R*4 001054      RRTVPE SF R*4 001058      SIGMAB FA C R*4 00005C
SIGNAT FA C R*4 000060

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***** COMMON INFORMATION *****

NAME OF COMMON BLOCK * * SIZE OF BLOCK 000078 HEXADECIMAL BYTES											
VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.	VAR. NAME	TYPE	REL. ADDR.
VOL1	R*4	000000	VOL2	R*4	000004	VOL3	R*4	000008	ACE	R*4	00000C
ASTAR	R*4	000010	E	R*4	000014	AC	R*4	000018	AO	R*4	00001C
MSHOCK	R*4	000020	MDS	R*4	000024	MUS	R*4	000028	TTOT	R*4	00002C
XS	R*4	000030	XSTAR	R*4	000034	PRES1	R*4	000038	PRES2	R*4	00003C
PRES3	R*4	000040	C	R*4	000044	MIRE	R*4	000048	PTOT	R*4	00004C
BETA1	R*4	000050	BETA2	R*4	000054	ALPCH	R*4	000058	SIGMAB	R*4	00005C
SIGNAT	R*4	000060	H	R*4	000064	DELTA	R*4	000068	OMEGA	R*4	00006C
ALFBAR	R*4	000070	ULIRE	R*4	000074						

SOURCE STATEMENT LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
80	56	001ED2	84	64	001F2E	90	66	001F6E	91	74	001FCC
94	82	00204E	99	84	00208E	10	232	0031DE	20	237	00322E
30	253	003402	40	277	00358A	50	292	0037C6	60	357	003EE6
70	358	00400A	13	449	005166	17	467	0051F6	24	512	00539A
25	520	005900	33	615	0062CE	34	619	006322	15	667	006800

COMPILER GENERATED LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
100001	2	001B04	200001	27	001C76	200002	40	001D78	100002	51	001E6A
100003	61	001EEA	100004	69	001F83	100005	79	00200A	200003	93	00218A
200004	112	0022CA	200005	136	002492	200006	166	0026EE	200007	169	00284A
200008	173	002972	200009	177	002A58	200010	182	002B1A	200011	186	002CF2
200012	194	002CE0	200013	200	002D23	200014	205	002E96	200015	203	002F42
200016	216	003052	200017	224	003130	100006	231	0031C8	100007	233	0031E2
100008	236	003218	100009	238	003232	200018	243	0032FE	200019	248	0033C0
200020	256	00340E	100010	257	003493	100011	259	0034B6	200021	264	003540
200022	272	003618	100012	276	003670	100013	278	00369E	200023	290	00379E
100014	291	0037AC	100015	293	0037CA	200024	305	003CD4	200025	312	0034D6
200026	321	003AF2	200027	329	003C1C	200028	336	003CC2	200029	343	003D74
200030	350	003E2C	100016	356	003E08	100017	358	003EEA	200031	365	003FCA
100018	367	003FF4	100019	369	00400E	200032	379	0041A8	200033	382	004293

```

*VERSION 1.3.0 (01 MAY 80)      UNST      SYSTEM/370 FORTRAN H EXTENDED (ENHANCED)      DATE 80.353/15.02.49      PAGE 32
200034 337 0043D6      200035 391 004476      200036 396 004522      200037 408 004634
100020 423 0050DE      100021 430 0050CC      200038 441 00511A      200039 447 00515A
200040 452 0051DE      200041 471 00537A      200042 476 005468      200043 479 0055FE
200044 492 00577C      200045 499 00580A      100022 503 00582C      200046 529 005A7A
200047 533 00582E      200048 537 0058CC      200049 544 005874      200050 550 005808
200051 559 0058DA      200052 553 0058CC      200053 569 005F5C      200054 573 005F92
200055 587 005866      200056 595 006194      200057 593 006292      100023 602 006224
200053 603 0064D0      200059 636 0065C0      100024 668 006890

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FORMAT STATEMENT LABELS

LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR	LABEL	ISN	ADDR
1000	8	000020	95	23	000020	71	52	00003A	75	57	00005B
76	70	00007C	78	76	00009F	100	113	0000C2	105	118	0000FB
110	119	000126	115	120	00014D	120	121	000175	333	404	0001A3
6000	416	000135	6100	417	0001C3	6200	418	0001DD	6001	419	0001F1
6002	420	000204	6003	421	000217	6004	422	00022B	6005	423	00023E
6006	424	000251	6007	425	000264	6008	426	000277	410	510	00028A
420	519	000279	6666	620	0002AA	440	644	0002DC	450	645	0003C0
440	646	000320	470	647	000366	430	648	00038A	490	649	0003CC
500	650	000400	510	651	000442	500	652	000476	530	653	0004A4
540	654	0004E4	550	655	000520	560	656	00055E	570	657	00058E
520	658	0005DC	590	659	00062E	600	660	000654	610	661	000684
620	662	0006AC	630	663	0006CC	640	664	00070C	650	665	00074C
640	666	00073A	670	672	0007C8	680	673	0007F4	690	674	000820
700	675	000850									

```
*OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(60) SIZE(MAX) AUTODBL(NONE)
```

```
*OPTIONS IN EFFECT: SOURCE EBCDIC HOLIST NODECK OBJECT MAP NOFORNAT GOSTHT XREF NOALC NOANSF TERM IBM FLAG(I)
```

```
*STATISTICS* SOURCE STATEMENTS = 676, PROGRAM SIZE = 27026, SUBPROGRAM NAME = UNST
```

```
*STATISTICS* NO DIAGNOSTICS GENERATED
```

```
***** END OF COMPILATION *****
```

```
280K BYTES OF CORE NOT USED
```

```
*STATISTICS* 2 DIAGNOSTICS THIS STEP, HIGHEST SEVERITY CODE IS 8
```

FORTRAN H EXTENDED (ENHANCED)

*** FORTRAN H EXTENDED ERROR MESSAGES ***

```

LINE 00000300 5 FORMAT('CHANNEL FLOW CHOKE FLUTTER ANALYSIS DECK 9066'// 20A4//
LINE 00000310 1' ALPDAR',F10.5,5X,'MEAN TORSIONAL DEFLECTION THRU CYCLE (DEG.)'/
LINE 00000320 2' ALPCH ',F10.3,5X,'STAGGER ANGLE OF BLADE ROW (DEG.)'/
LINE 00000330 3' BETA1 ',F10.3,5X,'INLET AIR ANGLE (DEG.)'/
LINE 00000340 4' BETA2 ',F10.3,5X,'EXIT AIR ANGLE (DEG.)'/
LINE 00000350 5' C ',F10.5,5X,'CHORD (IN.)'/
LINE 00000360 6' DELTA ',F10.5,5X,'GAP BETWEEN BLADES (IN.)'/
LINE 00000370 7' DIAM ',F10.3,5X,'NOODAL DIAMETER'/
LINE 00000380 8' E ',F10.6,5X,'ELASTIC AXIS LOCATION REF. TO MIDSPAN'/
LINE 00000390 9' EPS ',F10.5,5X,'TOLERANCE FOR PRESSURE RATIO'/
LINE 00000900 X' GAM ',F10.5,5X,'SPECIFIC HEAT RATIO'/
LINE 00000910 1' HBAR ',F10.6,5X,'MEAN FLAPPING DEFLECTION OF BLADE (IN.)'/
LINE 00000920 2' HAVR ',I10 ,5X,'RAVE MOTION'/
LINE 00000930 3' NI ',F10.5,5X,'INLET MACH NUMBER'/
LINE 00000940 4' MI ',F10.5,5X,'L. E. MACH NUMBER'/
LINE 00000950 5' NB ',I10 ,5X,'NUMBER OF BLADES'/
LINE 00000960 6' NP ',I10 ,5X,'NUMBER OF AIRFOIL COORDINATES'/
LINE 00000970 7' NSECT ',I10 ,5X,'NUMBER OF SEGMENTS FROM L.E. TO T.E.'/
LINE 00000980 8' NTIME ',I10 ,5X,'NUMBER OF TIME INCREMENTS'/
LINE 00000990 9' OMEGA ',F10.2,5X,'FREQUENCY OF VIBRATION (CYCLES. / SEC.)'/
LINE 00001000 X' PHIIB',F10.4,5X,'BENDING MODE INTERBLADE PHASE ANGLE (RAD.)'/
LINE 00001010 1' PHIIBT',F10.4,5X,'TORSIONAL MODE INTERBLADE PHASE ANGLE (RAD.)'/
LINE 00001020 2' PR ',F10.5,5X,'STATIC PRESSURE RATIO ACROSS STAGE'/
LINE 00001030 3' V ',F10.3,5X,'RELATIVE INLET VELOCITY (FT/SEC)'/

```

IFE029I SEVERITY 8(E) ISN 0032 THE NUMBER OF CONTINUATION CARDS EXCEEDS 19. COMPILER PROCESSING OF THE STATEMENT CONTINUES.

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LINE 00002220 WRITE(6,117) BETAPRI
IFE226I SEVERITY 8(E) ISN 0119 THE STATEMENT HAS A VARIABLE WITH MORE THAN SIX CHARACTERS.
THE RIGHTMOST CHARACTERS ARE TRUNCATED.

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SOURCE STATEMENTS = 245, PROGRAM SIZE = 9340, SUBPROGRAM NAME = MAIN

SOURCE STATEMENTS = 676, PROGRAM SIZE = 27026, SUBPROGRAM NAME = UNST

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NOMENCLATURE

<i>Symbol</i>	<i>Definition</i>
A	Area between the blades
$A_{z\infty}$	Complex constant describing irrotational flow field
a	Speed of sound
B	Axial wave number
b	Semi-chord of the airfoil
C	Tangential wave number
c_p	Specific heat at constant pressure
c_v	Specific heat at constant volume
$\frac{D}{Dt}$	Substantial derivative
e	Specific internal energy
F	Force on the control volume
h	Enthalpy
\bar{h}	Blade deflection in bending mode
k	Reduced frequency based on semi-chord
M	Mach number
m	Mass
n	Unit vector normal to surface
p	Pressure
R	Universal gas constant
S	Nondimensional interblade spacing
T	Temperature
U	Velocity along airfoil chord
u	Axial velocity component

NOMENCLATURE (Continued)

<i>Symbol</i>	<i>Definition</i>
\tilde{u}	Specific intrinsic energy
V	Volume
v	Tangential velocity component
\dot{w}	Mass flowrate
Z	Nondimensional elastic axis position
$Z_{+\infty}$	Complex constant describing rotational flow field
$\bar{\alpha}$	Mean torsional deflection
α_{ch}	Cascade stagger angle
β_1	Inlet air angle
β_2	Exit air angle
δ	Gap between the blades
γ	Ratio of specific heats
$\Phi'_{z\infty}$	Perturbation velocity potential
ρ	Density
τ	Gap between the blades
σ	Interblade phase lag
ψ	Stream function
ζ	Vorticity

Subscripts

$-\infty$	Far upstream of blade row
$+\infty$	Far downstream of blade row
x	Axial component
y	Tangential component
ave	Average flow parameter

NOMENCLATURE

Subscripts

Definition

IRE	Relative inlet quantity
i	Inlet to blade row
E	Outlet from blade row
1,2,3	Average quantity in the control volume
us	Upstream of the shock
ds	Downstream of the shock
R	Real part
I	Imaginary part
IR	Irrrotional component
S	Shock
E	Cascade exit

Superscripts

-	Steady-state quantity
'	Perturbation quantity
-'	Mean perturbation quantity
*	Blade throat

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16. Abstract Utilizing semi-actuator disk theory, a mathematical analysis was developed to predict the unsteady aerodynamic environment for a cascade of airfoils harmonically oscillating in choked flow. In the model, a normal shock is located in the blade passage, its position depending on the time dependent geometry and pressure perturbations of the system. In addition to shock dynamics, the model includes the effect of compressibility, interblade phase lag, and an unsteady flow field upstream and downstream of the cascade. The theory was evaluated by comparing calculated unsteady aerodynamics with isolated airfoil wind tunnel data and predicted choke flutter onset boundaries with data from testing of an F100 high pressure compressor stage.			
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