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Semiconductor Technology Program Progress Briefs

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SEMICONDUCTOR TECHNOLOGY PROGRAM

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ABSTRACT - This report provides information on the current status of NBS work on measurement technology for semiconductor materials, process control, and devices. Results of both in-house and contract research are covered. Highlighted activities include: modeling of diffusion processes, analysis of model spreading resistance data, and studies of resonance ionization spectroscopy, resistivity-dopant density relationships in p-type silicon, deep level measurements, photoresist sensitometry, random fault measurements, power MOSFET thermal characteristics, power transistor switching characteristics, and gross leak testing. In addition, brief descriptions of new and selected on-going projects are given. The report is not meant to be exhaustive; contacts for obtaining further information are listed. Compilations of recent publications and publications in press are also included.

KEY WORDS - Electronics; integrated circuits; measurement technology; microelectronics; semiconductor devices; semiconductor materials; semiconductor process control; silicon.

Preface

κ.,

This report covers results of work during the forty-fifth quarter of the NBS Semiconductor Technology Program. This Program serves to focus NBS research on improved measurement technology for the use of the semiconductor device community in specifying materials, equipment, and devices in national and international commerce, and in monitoring and controlling device fabrication and assembly. This research leads to carefully evaluated, well-documented test procedures and associated technology which, when applied by the industry, are expected to contribute to higher yields, lower cost, and higher reliability of semiconductor devices and to provide a basis for controlled improvements in fabrication processes and device performance. By providing a common basis for the purchase specifications of government agencies, improved measurement technology also leads to greater economy in government procurement. Financial support of the Program is provided by a variety of Federal agencies. The sponsor of each technical project is identified at the end of each entry in accordance with the following code: i. The Defense Advanced Research Projects Agency; 2. The National Bureau of Standards; 3. The Division of Electric Energy Systems, Department of Energy; 4. The Division of Distributed Solar Technology, Department of Energy; 5. The Defense Nuclear Agency; 6. The C. S. Draper Laboratory; 7. The Army Electronics R&D Command; 8. The Air Force Avionics Laboratory; 9. The Naval Material Command; 10. The Naval Weapons Support Center; 11. The Solar Energy Research Institute; 12. The Naval Avionics Center; 13. The Lewis Research Center, National Aeronautics and Space Administration; and 14. The Office of Naval Research.

This report is provided to disseminate results rapidly to the semiconductor community. It is not meant to be complete; in particular, references to prior work either at NBS or elsewhere are omitted. The Program is a continuing one; the results and conclusions reported here are subject to modification and refinement. Further information may be obtained by referring to more formal technical publications or directly from responsible staff members, telephone: (301) 921-listed extension. General information, past issues of progress briefs, and a list of publications may be obtained from the Electron Devices Division, National Bureau of Standards, Washington, D.C. 20234, telephone: (301) 921-3786.



Semiconductor Technology Program Progress Briefs



Modeling of Diffusion Processes

Fick's Law treats the diffusion coefficient as the factor of proportionality between the flux and the spatial gradient of a diffusing species. A recently reported model for impurity diffusion in semiconductors computes the flux in a different way by taking the spatial gradient of the product of the diffusion coefficient and the density of diffusing It was suggested that this species. modified form of Fick's Law implicitly accounts for electric-field effect. The significance of the difference between the two forms of the diffusion equation was explored by viewing diffusion as a Markov process, that is, as a stochastic process in which the diffusing species is undergoing a random walk wherein each step is independent of all other steps.

It was concluded from this analysis that the additional term in the modified form of Fick's Law, which contains the gradient of the diffusion coefficient, is a force term which may or may not be appropriate in the total flux equation. For the case of singly charged vacancy diffusion, and only for this case, there is a mathematical equivalence between this term and the force term associated with the self electric field. However, there is no physical relationship between this additional term and terms resulting from the effects of self or applied electric fields. It was therefore concluded that the conventional form of Fick's Law should be used to express the diffusion component of flux to which any force-driven components should be added explicitly to give the total flux. [Sponsor: 2] (J. R. Lowney, x3625)

Model Spreading Resistance Data

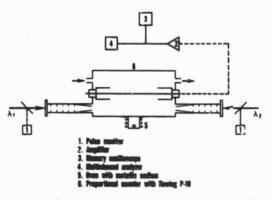
For uniform or shallow (<5 µm) diffused layers, the spreading resistance depends linearly on the logarithm of the probe spacing with a slope which is propertional to the sheet resistance of the layer and an intercept which is equal to the effective radius of the probas. Previous study of the relationship between the spreading resistance and probe spacing showed that for deeper diffusions the intercept cannot be interpreted in terms of the probe radius but that the slope is approximately proportional to the sheet resistance. To extend the study to ion-implanted layers, model spreading resistance data were generated from Gaussian resistivity profiles corresponding to implants at several beam energies and fluences into substrates with resistivity of 100 or 108 R. cm. The lower value was used to simulate the case of implantation into a substrate of the same conductivity type, while the higher was used to simulate the isolation which arises from the junction formed by implantation into substrates of opposite conductivity type.

In each case, data were generated for several probe spacings; the same effective electrical probe radius was used in all cases. At various points along the profile from the surface to the substrate, curves of spreading resistance as a function of the logarithm of the probe spacing were analyzed by linear regression to ascertain whether the slope was proportional to the incremental sheet resistance calculated directly from the model registivity profile and whether the intercept was related to the probe radius used in the simulation. For the $10^8 - \Omega \cdot cm$ substrate simulations, good agreement was obtained between the sheet resistance calculated from the slope and that calculated from the resistivity profile. For the 100-Q.cm substrate simulations, fluences of the order of 10^{13} cm⁻² were required to achieve comparable agreement. For both cases, the agreement between the intercepts and the probe radius used in the simulations was generally not satisfactory which suggests that the probespacing experiment is not an accurate method for determining the probe radius. However, the results of these simulations and those previously reported for deep diffusions suggest that probespacing experiments may be effectively used to determine the sheet resistance of small geometry specimens which contain diffusions or implants into substrates of the opposite conductivity or, if the substrate resistivity is high enough, into substrates of the same conductivity type. This result is especially important where specimen size precludes accurate sheet resistance measurement by the four-probe metho but where measurements of this quantity are important for process control and process modeling. [Sponsor: 1]

(J. H. Albers, x3625) '

Resonance Ionization Spectroscopy

resonance ionization spectroscopy (RIS) The pulse energy was typically 0.5 J; in a proportional counter filled with the asymmetric triangular pulse has a sodium-contaminated P-10 (90% argon plus rise time of about 100 ns and a base 10% methane) gas to a pressure of 94 width of about 700 ns. The counter was Torr. The apparatus used for this ex- calibrated for single-electron detection periment is shown schematically in the with uv light from a small mercury lamp accompanying figure. A laser of wave- appropriately collimated to illuminate a length λ_1 = 588.995 nm induces the small area of the counter cathode. The 3s-3p transition (ground state to one of normalized pulse-height distribution was



Schematic diagram of apparatus for resonance ionization spectroscopy experiment.

the "D" levels), and a second laser with wavelength $\lambda_2 = 568.820$ nm induces the 3p-4d transition leading to a level only 0.855 eV below the ionization potential in sodium. From this highly excited level, either laser λ_1 (2.104 eV) or λ_2 (2.179 eV) can induce photoionization, releasing an atomic electron with very low kinetic energy. This photoelectron is accelerated by the electric field in the counter, drifting towards the anode to develop a Townsend avalanche in the gas.

The pulsed dys lasers are pumped by a coaxial menon flashlamp and tuned to the appropriate wavelength by a dispersive system in the laser cavity. The radiation is collimated to a beam 3 mm in diameter by a reducing telescope provided with a spatial filter. The intersection of the two beams in the counter gener-A single atom of sodium was detected by ates an active volumo of about 100 mm³.

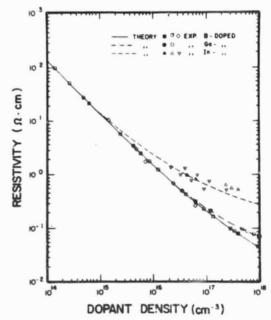
found to fit the expected probability distribution for a series of single-Calibration for sevelectron events. eral primary electrons (up to 350) was done by use of x-ray fluorescence lines from a variety of metals. The average number of primary electrons generated by each line, which is proportional to the line energy, was a linear functic. of the peak amplitude of the signal.

In the single-atom detection experiment, a controlled amount of sodium vapor was injected into the counter from The the small oven attached to it. sodium-RIS signal appeared only when both lasers were tuned to the proper wavelengths; it disappeared when either laser was detuned or blocked. Singleatom detection was confirmed by decreasing the amount of injected sodium until the RIS signal disappeared; at sodium levels just above extinction, the signal amplitude distribution was consistent with the single-electron pulse height Jistribution. [Sponsors: 2, 14]

> (S. Mayo, x3625, and T. B. Lucator:o,* x2031)

Pesistivity-Dopant Density Evaluation

Theoretical expressions were derived at into account the effect of hole-hole the University of Florida to compute scattering on both lattice and ionized hole mobility and resistivity as func- impurity scattering relaxation times and tions of dopant density and temperature the effect of interband transitions on for silicon doped with boron, gallium, or indium. The valence band of silicon tion time. Resistivity and dopant denwas represented by a three-band model sity measurements were made as a funcwhich takes into account the nonpara- tion of temperature over the range from bolic nature of the bands. This attrib- 100° to 400°C on silicon wafers doped ute of the valence band is included in with boron, gallium, or indium. Agreethe effective mass calculations. Con- ment between the experimental and theotributions from scattering by acoustical retical results was within 10% over the and optical phonons, ionized impurities, entire temperature range. and neutral impurities were considered panying figure shows the calculated in the calculation of average relaxation resistivity-dopant density curves for a time. In addition, the model also takes temperature of 300 K together with ex-* NBS Radiation Physics Division.



Resistivity as a function of dopant density for boron-, gallium-, and indiumdoped silicon at 300 K. Theoretical curves are shown as lines, experimental data obtained in the present work are shown as solid symbols, and data obtained from the literature are shown as open symbols.

the acoustic phonon scattering relaxa-The accomperimontal data from this study and from the literature. The differences which thus would obscure the proper identifiarise because of differences in ioniza- cation. tion energy are clearly evident. [Spon- study experimental procedures for detersor: 1]

Deep Level Measurements

In making doep level measurements to identify impurities in a test specimen, Photoresist Sensitometry one determines the activation energy and prefactor associated with carrier emis- Experimental work was completed which sion from one or more centers and re- verifies the applicability of Van Krevlates these to values which have been eld's additivity law to exposures of established from measurements on inten- positive photoresists and therefore also tionally contaminated specimens. Calcu- confirms the suitability of the techlations were made to determine the deep niques for determining exposure characlevel transient spectroscopy (DLTS) re- teristics of photoresists which had b.en sponse which might be obtained from a postulated some years ago. The law had specimen which contains two impurity been previously shown to apply whe centers relatively closely spaced in an- films of a commonly used positive resist ergy. pairs of levels were calculated for a radiation of 356, 405, and 436 nm. In range of sampling times typical of those these experiments, exposures were made used in experiments. From these curves, both for equal times and for times adeach of which had a single maximum (al- justed to compensate for variations in though some appeared broader and some- the transmittance of the narrow-pass what distorted when compared with sig- filters so that the exposures matched nals computed for single isolated lev- the relative irradiance of a mercury arc els), the apparent activation energies lamp at each of the three wavelengths. (AE) and prefactors (B) were determined When these experiments were repeated and compared with the originally assumed with exposure first to the 436-nm irraparameters. The results of two repre- diation followed successively by exposentative cases are shown in the accom- sure to the 405- and 365-nm irradiation, panying table. The parameters of the the experimental exposure vs. exposed composite signal differ from those of depth curve also agreed well with calcueither of the two assumed levels and lations based on Van Krevela's Law.

Parameters for Model DLTS Calculations

	Case 1	Case 2
∆E ₁ , eV	0.2018	0.2036
B ₁ , s ⁻¹ .K ⁻²	3.115 × 10 ⁶	2.48 × 10 ⁶
ΔĒ ₂ , eV	0.175	0.187
B ₂ , s ⁻¹ ·K ⁻²	3.115×10^{6}	2.48 × 10 ⁶
ΔĒ, eV	0.212	0.2018
B, s ⁻¹ .K ⁻²	9.71×10^{6}	3.115×10^{6}

Note - In both cases, level 2 is assumed to be present with half the density of level 1. . NBS Contact

Further work is underway to (W. R. Thurber, * x3625) mining that interference of this type are present and ways of suparating peaks associated with closely spaced levels. [Sponsor: 3] (R. Y. Koyama, x3625)

The DLTS signals for several were successively exposed to filtered

Another set of experiments was carried out in which resist films were exposed simultaneously to filtered radiation of wavelengths of 365, 405, and 436 nm from the mercury arc lamp. A specially designed jig was used for these experi-Three mirrors were mounted in ments. the plane bisecting and normal to the mercury arc lamp at azimuthal angles of 120° from each other. These mirrors, in

through respective lenses and narrow cannot be detected by single-fault-type band filters onto a resist-coated glass plate approximately 76 cm above the lamp and on the same center line. This created a superposition of impinging light of three wavelengths about 4° from the normal on the resist-coated plate. The actual measured intensities used for the respective exposures were within 5% of those calculated from Van Kreveld's additivity law and monochromatic exposure data except for the smallest exposures where differences of up to 20% were found. For these small exposures, the experimental errors are greatest.

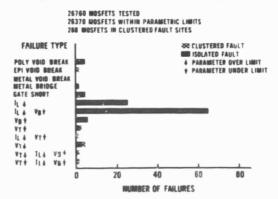
Measurements of the power density spectrum of the mercury arc lamp confirmed the hypothesis that unaccountedfor background radiation (radiation from the continuum existing between the principal mercury arc lines) was responsible for the previously reported discrepancy between the measured and calculated exposure data when the resist was exposed by unfiltered radiation from the lamp. It was found that (1) the principal lines at 365, 405, and 436 nm were only 3 to 4 nm wide (FWHM), one-third or less than the published spectral width, and (2) the energy radiated at wavelengths between the principal lines was 35% of the energy contained in these lines. [Sponsor: 1] (D. B. Novotny, x3621)

Random Fault Measurements - I

Test results from a random access fault structure can be used to detect, identify, and analyze a variety of fault mechanisms introduced during the fabrication of integrated circuits. The structure allows one to determine the relative density of different faults which limit process yield without previous knowledge of what fault types might be expected.

turn, refl-cted light vertically upward In some cases, it detects faults which random fault test structures. A prototype structure which consists of a 10 by 10 array of electrically isolated n- or p-channel MOSFETs has been implemented on test pattern NBS-16. This test pattern is being used to develop methods to assess the electrical performance and yield potential of radiation-hardened, silicon-gate CMOS/SOS LSI circuits. The structures are repeated at intervals across a wafer so that the threshold voltage (V_T) at a drain current of 1 µA with gate tich to drain, source-drain breakdown voltage (V_R) at a drain current of 10 µA with gate tied to source, and source-to-drain leakage current (I_C) at a drain voltage of 10 V with gate tied to source can be measured on a statistically significant number of individual MOSFET's. These results are then compared with failurs criteria which define the parametric limits required for satisfactory circuit operation. In addition, open and short circuits in the structure can be detected.

> Typical results from a wafer representative of a lot with relatively low device yield are shown in the accompanying



Test results from random access fault structures on a wafer representative of a lot with relatively low device yield.

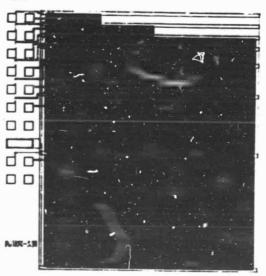
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devices to meet a given criterion is lization in each array are each contreated as a single clustered fault. In nected to a probe pad; all of the epi this example, 40 of the MOSFETs could lines in each group of arrays are connot be tested because of opens or shorts nected to a single probe pad. There are in the common connections between the 95 such croups of seven arrays on a waf-MOSFET and the probe pad. The measured er fabricated with test pattern NBS-16. parameters of 26370 of the 26760 tested Preliminary analysis of initial measure-MOSFETs fell within the dsfined limits; ments of leakage current between pairs there were 102 isolated faults, and 288 of conducting layers on wafers taken MOSFETs were contained in the clustered- from three processed lots suggests that fault sites. From these results, it can the array sizes are appropriate for be seen that the dominant failure type evaluating the process under study. is an excess source-to-drain leakage However, in some cases, the origin of current, usually coupled with low break- the leakage current cannot be estabdown voltage. Analysis of test results lished unambiguously; work is underway suggests that this failure mode, which to determine the causes of the various had not been anticipated at the time the test pattern was designed, was the primary contributor to yield degradation of the accompanying product wafer lot. This failure mode would not have been detected from test results of other test structures on the pattern. [Sponsor: (L. W. Linholm, x3541) 81

Random Fault Measurements - II

One of the two random fault test structures on test pattern NBS-16 is designed to measure gate and field oxide integrity in a CMDS-SOS process. A composite drawing of this structure is shown in the accompanying figure. The structure is comprised of five layers on a sapphire substrate: n epitaxial silicon lines, 80-nm thick gate oxide, p⁺ polysilicon lines running at right angles to the epi linas, 600-nm thick field oxide, and aluminum metallization. At the crossovers, the poly and epi lines are separated by gate oxide. The poly lines are grouped to form seven arrays with 57, 114, 225, 375, 750, 1500, and 3000 crossovers. Each array is covered with the metallization layer which is separated from the poly or epi lines by

figure. Failure of two or more adjacent field oxide. The polysilicon and metal-



Composite drawing of random fault test structure for evaluating gate and field oxide integrity in a CMDS-SOS process. The n - epi lines which run vertically are connected together by the wide horizontal bars and brought out to the double probe pad E. The poly lines run horizontally; each array is connected to a separate probe pad P. Each array is completely covered by metallization connected to a separate probe pad M.

ambiguities, to investigate methods of perature, at least for drain currents up be removed, and to evaluate design modifications to provide better separation of leakage current sources. [Sponsors: 1,8] (M. A. Mitchell and L. W. Linholm, x3541)

Power MOSFET Characterization

Three techniques for measuring the temperature of silicon power MOS fieldeffect transistors (MOSFETs) are being investigated. Two of the techniques em-, loy temperature-sensitive electrical parameters (TSPs) of the device, the forward voltage of the drain-body diode (V_{DB}) at constant low current, and the applied gate voltage (V_{C}) at constant drain current. The third is a direct measurement of the surface temperature using an automated infrared microradiometer (IRM).

The temperature variation of VDB is both linear and well characterized. Measurement of MOSFET temperature with this TSP is similar to a method used to measure the temperature of bipolar transistors using the forward voltage of the collector-base diode as the TSP. However as in the bipolar case, the diode senses only an average chip temperature because it occupies a large area of the chip, whereas the heat source (defined by the channel region) occupies a very small area. Also, the large electrical capacitance of the forward-biased diode makes the require. electrical switching extremely slow.

Much less is known about the temperature depandence of VG. It is known that this voltage can increase, decrease, or remain constant with temperature depending upon the value of the applied gate voltage and the channel characteristics. Preliminary results suggest that VG varies linearly with tem-

analysis by which the ambiguities might to about 3 5; the temperature coefficient varies from about -6 mV/°C at 10 mA to about -4 mV/°C at 1 A to about -2 mV/°C at 3 A. At higher drain currents, typical of those encountered in device operation, the temperature coefficient would be expected to approach or pass through zero; therefore, to obtain adequate sensitivity, it may be desirable to use a switching method, similar in concept to the standard method employed for measuring the junction temperature of bipolar transistors with the emitterbase voltage (V_{EB}) as the TSP. In this method, the device is operated at the conditions for which the temperature is desired and then rapidly switched to a low value of drain current for measurement of the Vc.

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If the temperature measured by the IRM is to be accurate, the power MOSFET surface must be coated with a uniform, high emissivity film because the very fine metallization patterns on these devices cause the spatial variation of surface emissivity to be extremely rapid and large. In the present work, a flatblack, commercial spray paint with low carbon content (to reduce electrical conductivity) is used.

Preliminary comparative measurements of the temperatures determined by the three methods confirm that the temperature indicated with VDB as the TSP is very nearly equal to the average chip temperature as determined by an unweighted spatial averaging of the IRM results. However, contrary to results found with the use of VER in bipolar devices, use of VG as the TSP results in a device temperature in excess of the peak temperature as determined by the IRM. This may occur because the spatial resolution of the IRM is not sufficient to resolve the peak gate temperature or because of errors in the gate voltage

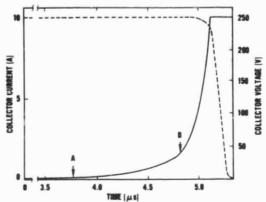
measurement method. These part bilities are being investigated furth a [Sponsor: 2] (D. L. Blackburn, x3621)

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Transistor Switching Characteristics

Work was undertaken to develop improvea methods for characterizing the switching properties, particularly the reversebias safe operating limits, of power transistors. Initially, the effect of the magnitude of the base current, IBR, on transistor switching characteristics is being studied. Previous work had shown that the magnitude of IBR has a strong influence on the reverse-bias safe operating limits of silicon $n^+ - p - v - n^+$ power transistors. Generally, as IBR is increased, the voltage at which second breakdown occurs decreases. It was also observed that for the turnoff of these devices, a time interval exists which occurs during, but is distinct from, the traditional storage time. During this newly observed interval, designated the dynamic saturation time, the collector voltage slowly of the collector voltage can be attribrises as the collector current remains uted to an increased resistive voltage nearly constant. It was found that as drop as stored charge is being removed much as 30% of the total energy dissi- from the collector, narrowing the width pated during turnoff may occur during of conduct vity-modulated portion of the this interval. It is therefore impor- v-region of the collector. tant to understand its cause and effects.

the base-emitter voltage during tarroff tion of the dependence of the secondshows that current focusing to the cen- breakdown voltage on base current. For ter of the emitter fingers does not be- reverse-bias second breakdown to occur, gin to occur until after the device has a high current density is required; the entered dynamic saturation. consistent with the idea that the effec- the current focusing. Thus, the magnitive base width (i.e., the mecallurgical tude of IBR during dynamic saturabase width plus the width of the tion should affect the second breakdown conductivity-modulated portion of the characteristics, but the magnitude belightly doped (v) collector region) be- fore dynamic saturation (i.e., during gins to decrease when the device enters classical saturation) should have no efdynamic saturation. Thus the slow rise fect. This was verified experimentally;



Current (dashed curve) and voltage (solid curve) transients of an n⁺-p-v-n⁺ power transistor following switching of the base current from +4 A to -0.5 A at time zero. The transistor is operating in a dynamic saturation condition between the points labeled A and B on the voltage trace. In the dynamic saturation region, the collector voltage increases slowly and the collector current is essentially constant as discussed in the text.

The finding that current focusing does not occur until the device enters dynam-Observation of the time variation of ic saturation also provides an explana-This is high density is realized as a result of

the voltage at which second oreakdown chamber for both pressurization and deoccurred was observed to be exactly the tection, while the second mode separates same if the value of IBR to which these functions. Present test times for the device was switched after dynamic saturation began was applied instead the single chamber mode, the process ocduring the entire turnoff time. Thus, turnoff time may be reduced with minimum risk of second breakdown if IBR is large at the start of the turnoff cycle but is reduced to a small value near the to the atmosphere and the remaining exend. [Sponsor: 13]

(D. L. Blackburn, x3621)

Gross Leak Testing

Although the correlation between gross leaks in hermetically packaged semiconductor devices and their failure rates is clear cut, the dry gas hermetic test methods currently in use are far from satisfactory. The commonly used helium mass spectrometer and radioisotope leak tests are based on back pressurization techniques. At present the limitation in applying these procedures to the gross leak range is the rapid depletion of the gas, from the package interior which leads to inability to detect large leaks. A new noncontaminating, quantitative dry gas test method has been developed to eliminate these problems. This method is intended for use in the glass capsules fabricated with capillary leak size range from about 10⁻⁵ to about leaks which were measured initially by 1 atm·cm³/s. It employs a rapid gas cy- means of an absolute procedure. cling technique to extend the upper sizes determined from the experimentally range of the helium mass spectrometer measured leak rate appear to be within leak detector by reducing the delay be- ± 50 % of the true (initial) value over tween the pressurization phase and the the leak size range from 10^{-5} to 0.5 detection of the tracer gas which has atm.cm³/s. Closer agreement is expected penetrated the package. the environmental conditions during the tector response time. [Sponsor: 2] test are tailored so that a simple quantitative relationship based on viscous flow theory exists between the true and measured leak values.

Two modes of operation have been em- The third in the scries of measurement

either mode are of the order of 8 s. In curs in four rapid steps. The package is pressurized for about 2 s in an atmosphere of mitrogen which contains about 1% helium. The chamber is then vented cess gas exhausted within a fraction of a second. After the chamber is refilled with a neutral gas such as nitrogen or dry air to further dilute the helium concentration in the chamber, it is exhausted a second time and pumped to low enough pressure for the helium loak detector to be valved in. This process, which is carried out under automatic control, reduced the partial pressure of helium in the ambient around the test package by a factor of about 107 within a few seconds. The second mode of operation is inherently simpler and less subject to sorption effects, but it requires rapid transfe of the test specimen between the pressurizing and detection cnambers; in this mode, the pressurization chamber is exhausted only once.

Measurements were made on a group of Leak In addition, after corrections are made for leak de-

(S. Ruthberg, x3621)

Linewidth Seminar Rescheduled

ployed. One mode utilizes a single test seminars on integrated circuit linewidth

measurements has been rescheduled for contrast. The procedure includes defi-July 15-18, 1980 at NBS/Gaithersburg. r tions of appropriate terms and speci-The primary emphasis in this seminar fications for the necessary equipment; will be on the measurement of widths of it describes techniques for setting up features in the range 0.5 to 10 µm on the microscope with either transmitted silicon wafers. The seminar will con- or reflected illumination. In addition, sist of lectures, laboratory demonstrations, hands-on training with optical the ratio of condenser numerical apermicroscope systems, and small group dis- ture to objective numerical aperture cussions of industry measurement problems. [Sponsor: 2] Cohen, x3786, and J. M. Jerke, x3621)

New Topic . . .

Carbon in Silicon -Average carbon densities typically present in silicon crystals grown by the Czochralski process range from 10^{16} to 10^{17} cm⁻³. Although as a neutral substitutional impurity carbon has no direct effects on device performance, it has recently been implicated in a wide range of deleterious indirect effects such as the nucleation of intrinsic defects, oxide precipitation, the "X" centers in indium-doped silicon, gold diffusion, and dopant distribution. A new project to study carbon in silicon was initiated. The first step in this project is to develop methods for measuring the carbon distribution across production wafers. These methods will then be applied in a thorough study of the role of carbon in silicon, from its initial distribution to the effects of processing. [Sponsor: (A. Baghdadi, x3625) 21

Work in Progress . . .

optical microscope to operate in Kohler determined whether the difference occurs illumination is being documented in the because the same center is not being obformat of an ASTM Practice. This oper- served in the two experiments or because ating mode provides uniform illumination the center responds differently to the of the specimen, reduces stray light, two measurement techniques. and results in bright images with good 2]

the procedure tells the user how to set equal to two thirds in order to provide (E. C. images with steep edge gradients. This procedure is an essential part of the instructions currently being developed for using the photomask-like chromiumon-glass artifacts (soon to be made available by NBS as standard reference material SRM 474) to calibrate an optical microscope for linewidth measurements in the 0.5- to 10-pm regime. These instructions will be applicable to most types of measurement systems, such as filar, image-shearing, and videomicrometer systems, used throughout the microelectronics industry to measure linewidths on photomasks. [Sponsor: 1] (J. M. Jerke, x3621)

The optical spectra of the deepest known sulfur center in silicon are being studied. Specimens were prepared by diffusing natural sulfur (95% 32S) or isotopically enriched sulfur (90% 34S) into ptype silicon doped with boron to a density of about 3×10^{16} cm⁻³. The diffusions were carried out at 1350°C for about 200 h in an evacuated, sealed quartz tube. In contrast to earlier studies of this center by isothermal transient capacitance and thermally stimulated current measurements, no significant isotopic shift of the energy A procedure developed for adjusting an level was observed. It remains to be (Sponsor: (R. A. Forman, x3625)

Recent Publications . . .

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