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Measurements in the Wake of an 'Infinite' Swept Airfoil

Iowa Inst. of Hydraulic Re3earch, Iowa City

Prepared for
National Aeronautics and Space Administration
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by

C. J. Novak and B. R. Ramaprian

Sponsored by
National Aeronautics and Space Administration
Grants No. NSG-1200 and NAG 2-110


IIHR Report No. 240
Iowa Institute of Hydraulic Research
The University of Iowa
Iowa City, Iowa 52242
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## LIST OF SYMBOLS

| $a_{i}, b_{i}, c_{i}$ | Senscr direction cosines $\mathbf{i}=1,2,3$ |
| :---: | :---: |
| $A_{2 d}$ | Bradshaw structure parameter as calculated in twodimensional case |
| $A_{3 d}$ | Bradshaw structure parameter as calculated in threedimensionai case |
| $A^{\prime}, B^{\prime}$ | Constants in King's Law |
| $b_{x}$ | Half width of longitudinal velocity in the wake |
| $b_{z}$ | Half width of spanwise velocity in the wake |
| C, L | Airfoil chord length $=922 \mathrm{~mm}$ |
| $C_{p}$ | Pressure coefficient |
| $A_{i}, B_{i}, C_{i}$ | Variables used in hot wire data analysis, $\boldsymbol{i}=1,2,3$ |
| $\begin{aligned} & a_{4}, b_{4}, c_{6} \\ & d_{4}, e_{4}, f_{4} \end{aligned}$ | Constants used in hot wire data analysis |
| E | Hot wire mean voltage output |
| H | Shape factor ( $\delta_{1} / \theta_{f}$ ) |
| k | Sensitivity correction factor $=.1$ |
| $n$ | Exponent in King's Law |
| $\bar{Q}_{1}, \bar{Q}_{2}, \bar{Q}_{3}$ | Effective mean velocities of sensors 1, 2, and 3, respectively |
| $q_{1}, q_{2}, q_{3}$ | Fluctuating effective velocities of sensors 1, 2, and 3 |
| $Q_{1}, Q_{2}, Q_{3}$ | Instantanecus effective velocities of sensors 1, 2, and 3 |
| $\bar{U}, \bar{V}, \bar{W}$ | Mean velocities in the $x, y$, and $z$ directions |
| U,V,W | Instantaneous velocities in $x, y$, and $z$ directions |
| $u^{\prime}, v^{\prime}, w^{\prime}$ | Root mean squares of $u, v$, and $w$ |
| $\mathrm{U}_{\mathrm{n}}, \mathrm{J}_{\infty}$ | Constant freestream velocity |


| $W_{C l}$ | W at the centerline |
| :---: | :---: |
| $W_{x}$ | Defect velocity in the longitudinal or chordwise direction |
| $W_{Z}$ | Defect velocity in the spanwise or crossflow direction |
| ${ }^{\text {wo }}$ | $w_{x}$ at the centerline of the wake |
| $w_{20}$ | $w_{z}$ at the centerline of the wake |
| e | Fluctuating hot-wire voltage output |
| $u_{\tau}$ | Friction velocity |
| 5 | Boundary layer thickness |
| $x, y, z$ | Streamwise (longitudinal or chordwise), normal and spanwise (crossflow) coordinates |
| $\varepsilon_{x}$ | Eddy viscosity, as calculated from the gradient of $U$ |
| $\varepsilon_{z}$ | Eddy viscosity, as calculated from the gradient of $w$ |
| $\theta$ | Sensor orientation in the $x-z$ plane |
| 4 | Sensor orientation in the $y-z$ plane |
| ${ }^{\theta} \mathrm{f}$ | Final station momentum thickness (2-0 definition) |
| $\theta_{t}$ | Trailing edge momentum thickness (2-D definition) |
| ${ }^{7} \times$ | $y / b_{x}$ |
| $\eta^{\prime}$ | ${ }^{y /} / b_{x}$ |
| ${ }^{5}$ | $\int_{-\infty}^{\infty}\left(1-\frac{U}{U_{\infty}}\right) d y$ |
| $\delta_{2}$ | $\int_{-\infty}^{\infty}\left(1-\frac{W}{U_{\infty}}\right) d y$ |
| ${ }^{8} 11$ | $\int_{-\infty}^{\infty} \frac{U}{U_{\infty}}\left(1-\frac{U}{U_{\infty}}\right) d y$ |
| ${ }^{*} 22$ | $\int \frac{W}{U_{\infty}}\left(1-\frac{W}{U}\right) d y$ |

$$
\begin{array}{ll}
\theta_{12} & \int \frac{U}{U_{\infty}}\left(1-\frac{W}{U_{\infty}}\right) d y \\
{ }^{\theta} 21 & \int \frac{W}{U_{\infty}}\left(1-\frac{W}{U_{\infty}}\right) d y
\end{array}
$$

$r \quad$ Angle between the velocity vector and the $x$-direction in the case of the swept air foil
$\phi$
Sweep angle of the airfoil

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## CHAPTER I

introduction

## 1. The Probiem Introduced

There is a growing interest, in recent years, in the development of calculation methods of three dimensional turbulent shear flows such as boundary layers and wakes. This is largely because of their application in aerodynamics and ship hydrodynamics. Calculation of the drag on the swept wing of an aircraft is a typical example. This problem presents many complexities (especially in the trailing edge and near-wake regions) such as viscid-inviscid interaction, three dimensionality and the modeling of turbulence. Computational techniques capable of handling these complexities are currently under development at various research centers. In order to rerify the accuracy of such methods and also to aid the development of new methods, it is essential to have a comprehensive data base on three-dimensional flows in general and three dimensional wake flows with particular reference to the problem cited above. Such data should include information on velocities, pressure and the details of the turbulence properties. Measurement of turbulence properties in three dimensional flows is still a challenging task since the required measurement techniques are only in their early stages of developmeni. Consequently, there is not much data available on the turbulence properties in three dimensional flows. The present research has been aimed at obtaining a set of detailed measurements, in a
relatively simple three-dimensional wake, namely the developing wake behind an infinite swept airfoil. Also as a part of this research program, significant effort has been directed towards the development and validation of a suitable technique for turbulence measurements in three-dimensional flows.

Previous work done in the wake of streamlined bodies have been largely restricted to two-dimensional flows. These include the work of Chevray and Kovasznay (1969), Pot (1979), Andreopoulos and Bradshaw (1980) as well as that of Sastry (1981) in the symmetric and asymmetric wakes of flat plates and airfoils. In these studies, the relevant components of the Reynolds stress tensor have been measured. Computational models have also been developed and used to describe these experimental findings. Of the above, Sastry's studies included the twodimensional wake of a Korn-Garabedian airfoil at incidence. The development of various wake parameters were studied and compared with those in the case of the flat piate wake. It should be noted that most of the data reported by the previous investigators have been restricted to the near or intermediate wake regions.

It is only very recently that experimental studies of wakes 'exhibiting three-dimensionality have been performed. One study, that of Cousteix and Pailhaus (1980) explored the wake behind a swept 'ONERA D' airfoil at incidence. This study contains hot-wire measurement of all the components of the Reynolds stress tensor in addition to data on mean velocities and flow inclination. Similar hot-wire data in the threedimensional wakes of turbomachinery blade have been obtained at Penn State University. These data, however, were obtained under very
difficult experimental conditions, with many uncertainties introduced by such factors as small physical size, probe orientation, probe dimensions, data reduction, etc. The data of Cousteix and Pailhaus (1980) were obtained under more favorable conditions. Yet, even in these experiments, (the airfoil studied had a chord of 200 mm and a span of 300 mm ) tunnel wall interference on the wake development could have been significant.

In the present study some of the limitations of the atove experiments have been removed. Present experiments were performed in the wake behind an "infinite" NACA 0012 airfoil swept at 30 degrees, and at zero incidence. Such a flow, being significantly three dimensional, yet relatively simple, is expected to provide useful information on growth rates, viscosity models and other properties of three-dimensional wakes. Also, these flows can be studied with regard to the manner in which they approach two-dimensionality at a large dirtance from the trailing edge.

A complete set of measurements are obtained both in the boundarylayer in the trailing edge region and in the developing wake behind the airfoil. These measurements include pressure distribution, mean velocity components, in the boundary layer and wake. Additionally, the six components of the Reynn?ds shear stress were also measured in the wake. The measurements extend up to two chord lengths downstream from the trailing edge. The turbulence measurements have been made using triplesensor hot-wire anemometry. A significant contribution of this study is the development and validation of a suitable technique for the measurement of turbulence properties using the triple-sensor probe and
a comparative study of alternative techniques. In fact, development of the present technique forms a part of a larger program of study on general three-dimensional boundary layer flows, currently in progress at the Institute of Hydraulic Research.

Chapter II will be devoted to the development of the triple-sensor anemometry. A complete description of the method is presented. Particulars concerning the airfoil experiments such as calibration procedures for hot-wire instrumentation, details of the airfoil fabrication, instrumentation, and mounting are discussed in Chapter III. The results and discussion are presented and discussed in Chapter IV. The effects of three-dimensionality are compared with previous two dimensional experiments and theory. Further, Chapter IV also contains the concluding remarks as well as suggestions toward further work.

## CHAPTER II

DEVELOPMENT OF THREE-DIMENSIONAL
HOT-WIRE INSTRUMENTATION
$\frac{\text { 1. Problem Description and Quantities }}{\text { Requiring Measurement }}$
In studies of three-dimensional turbulent flows, it is necessary to measure the instantaneous values of three mutually orthogonal velocities at a given point in space. The instantaneous values will contain three mean velocities as well as three time-dependent or fluctuating velocities. It. is these three instantaneous fluctuating velocity components, in addition to the mean components that require retrieval in this study. This retrieval must be general enough, such that it allows one to study flow fields of arbitrary geometry.

## 2. Available Techniques and Relative Assessment

Hot-wire anemometry has been the most commonly used instrumentation for velocity measurements especially in three-dimensional flows. Laser velocimetry, which is currently gaining popularity, is complex and not well tested in three-dimensional flows. Two techniques, using hot-wire anemometry have been tried in three-dimensional flows. One of these uses a triple-sensor probe in which there are three independent sensing elements oriented in certain specific relative directions, and the other uses a single sensor or $x$-wire probe oriented sequentially at different
angles to the flow. The first method, used by Gorton and Lakshminarayana (1976) and others, requires that the predominant direction of flow is known and that the probe axis coincides with it. The three instantaneous outputs $u, v, w$, from the sensors are then processed by an analog computer which solves for the mean values $\bar{U}, \bar{V}$, and $\bar{W}$ as well as for the six Reynolds stresses $u^{2}, v^{\prime}, w^{\prime}, \overline{u v}, \overline{v w}$ and $\overline{u w}$ via the soiution of a set of linearized simultaneous equations. One of the drawbacks of this nethod is that, often, there may not be a predominant direction in the flow. Also, nonlinear effects may affect the accuracy of this method.

The second method, used by Elsenaar and Soelsma (1974), consists of 1 single inclined wire probe rotating about it's roll axis. The probe outputs at successive orientations of the probe are used to obtain solutions for the mean velocities and turbulent stresses. Here again, the major difficulties are the need for actual physical rotation of the probe itself and the interference of the probe support system with the flow. Elsenaar and Boelsma have noted this and the associated uncertainties.

Apart from the shortcornings mentioned above, information is lost in these methods during the analog processing for time-averaged results. Thus instantaneous velocities, their triple correlations, and higher order moments cannot be determined.

The method selected for the present study use a triple sensor probe oriented in an arbitrary direction to the flow. The instantaneous outputs from the three sensors are used to obtain not only the three mean velocity components and the six Reynolds stresses but also the instantaneous fluctuating velocity components. This is done through
digital da :a acquisition software. Two alternate schemes were developed for comparison purposes. Each of these is described below in detail.

### 2.1 Method A

Figure 1 shows the probe and flow coordinate system. Consider hot wire sensor 1 with respect to the $X, Y, Z$ (flow) coordinate system. This sensor has an orthogonal coordinate system $\left(\zeta_{1}, \zeta_{2}, \zeta_{3}\right)$ associated with its orientation. We, then, have

$$
\begin{align*}
& \left(Q_{1}\right)=a_{1} U+b_{1} V+c_{1} W  \tag{1}\\
& \left(Q_{2}\right)=a_{2} U+b_{2} V+c_{2} W  \tag{2}\\
& \left(Q_{3}\right)=a_{3} U+b_{3} V+c_{3} W \tag{3}
\end{align*}
$$

where $Q_{1}, Q_{2}$ and $Q_{3}$ are the instantaneous velocitics in the $\zeta_{i}, \zeta_{2}, \zeta_{3}$ coordinate system, and $U, V$, and $W$ are the instantaneous velocities in the $X, Y$, and $Z$ coordinate system. The coefficients $a_{j}, b_{p}, c_{1}, \ldots$, are the appropriate direction cosines of the $X, Y$, and $Z$ axis with respect to the $\left(\zeta_{1}, \zeta_{2}, \zeta_{3}\right)$ coordinate system as shown in figure 1.

The effective velocity sensed by hot wire sensor 1 is given by [see for example Friehe and Schwarz (1968)]

$$
\begin{equation*}
Q_{1}=\left[\left(Q_{2}\right)^{2}+\left(Q_{3}\right)^{2}+k_{1}^{2}\left(Q_{1}\right)^{2}\right]^{1 / 2} \tag{4}
\end{equation*}
$$

with $k_{1}$ being the cross-sensitivity factor associated with deviation from the cosine law. $Q_{1}$ is obtained from the voltage output $E_{1}$ of the sensor 1, via its calibration curve. Now, if we substitute Eq. (1,2,3) Into Eq. (4), the resulting expression becomes

$$
\begin{equation*}
Q_{1}=\left[a_{4} U^{2}+b_{4} v^{2}+c_{4} W^{2}+d_{4} U V+e_{4} U W+f_{4} W W\right]^{1 / 2} \tag{5}
\end{equation*}
$$

The coefficients in the above equations are:

$$
\begin{align*}
& a_{4}=a_{2}^{2}+a_{3}^{2}+k_{1}^{2} a_{1}^{2}  \tag{6}\\
& b_{4}=b_{2}^{2}+b_{3}^{2}+k_{1}^{2} b_{1}^{2}  \tag{7}\\
& c_{4}=c_{2}^{2}+b_{3}^{2}+k_{1}^{2} c_{1}^{2}  \tag{8}\\
& d_{4}=2\left(a_{2} b_{2}+a_{3} b_{3}+k_{1}^{2} a_{1} b_{1}\right)  \tag{9}\\
& e_{4}=2\left(a_{2} c_{2}+a_{2} c_{3}+k_{1}^{2} a_{1} c_{1}\right)  \tag{10}\\
& f_{4}=2\left(b_{2} c_{2}+b_{3} c_{3}+k_{1}^{2} b_{1} c_{1}\right) \tag{11}
\end{align*}
$$

Similar expressions can be derived for sensors 2 and 3. One, then gets a set of three nonlinear algebraic equations for the three unknowns $U, V$, and W. These require solution by an iterative scheme. The method used is a Newton-Raphson scheme, which requires an initial guess to start the algorithm.

In practice, the instantaneous outputs $E_{1}, E_{2}$ and $E_{3}$ of the three sensors are sampled, digitized and stored. These data are later recalled, calibrated and the velocity components $U, V, W$ are computed on an instant to instant basis using the above procedure. After the entire set of instantaneous velocity components are evaluated, they are time averaged to obtain $\bar{U}, \bar{V}$, and $\bar{W}$. The turbulent fluctuations, $u, v$, and $w$ are then recovered and stored for subsequent procezsing. Since this method involves the solution of nonlinear equations, for each set of instantaneously sampled data, it is computationally very time-consuming and expensive. However, it does not involve any linearization and can therefore be used even for very large fluctuations.

### 2.2 Method B

In this method the effective velocities and resultant velocities are both decomposed into their mean and fluctuating components. Eqs. (1-3), thei,, become:

$$
\begin{align*}
& (\bar{Q}+q)_{1}=a_{1}(\bar{U}+u)+b_{1}(\bar{V}+v)+c_{1}(\bar{W}+w)  \tag{12}\\
& (\bar{Q}+q)_{2}=a_{2}(\bar{U}+u)+b_{2}(\bar{V}+v)+c_{2}(\bar{W}+w)  \tag{13}\\
& (\bar{Q}+q)_{3}+a_{3}(\bar{U}+u)+b_{3}(\bar{V}+v)+c_{3}(\bar{W}+w) \tag{14}
\end{align*}
$$

Now, if they are substituted into Eq. (4), we get

$$
\begin{align*}
& {\left[\left(a_{4} \bar{J}^{2}+b_{4} \bar{v}^{2}+c_{4} \bar{W}^{2}+d_{4} \overline{u v}+e_{4} \overline{U W}+f_{4} \overline{V W}\right)+\right.} \\
& \quad+\left(2 a_{4} \bar{u} u+2 b_{4} \overline{V_{v}}+2 c_{4} \bar{W} w+d_{4}(\bar{u} v+\bar{V} u)+e_{4}(\bar{u} w+\bar{W} u)+f_{4}\left(\overline{W_{w}} \bar{W} \bar{W}\right)\right. \\
& \left.\quad+\left(a_{4} u^{2}+b_{4} v^{2}+c_{4} w^{2}+d_{4} u v+c_{4} u w+d_{4} u v+c_{4} u w+f_{4} u w\right)\right] \\
& \quad=(\bar{Q}+q)_{1} \tag{15}
\end{align*}
$$

which can be written (collecting terms from Eq. 15 such that $A_{7}$ contains $\bar{W}$ terms, $B_{1}$ contains terms as $\overline{\mathrm{V}}$, and $\mathrm{C}_{1}$ contains uv terms):

$$
\begin{equation*}
\bar{Q}_{1}+q_{1}=\left[r_{1}+B_{1}+c_{1}\right]^{1 / 2} \tag{16}
\end{equation*}
$$

If we assume that $B_{1}+C_{1} \leq A_{1}$, we can linearize Eq. (16) as

$$
\begin{equation*}
\bar{Q}_{1}+G_{1}=A_{1}^{1 / 2}\left[1+\frac{B_{1}+C_{1}}{2 A_{1}}-\frac{1}{8} \frac{\left(B_{1}+C_{1}\right)^{2}}{A_{1}}+\ldots\right] \tag{17}
\end{equation*}
$$

Squaring Eq. (16), time-āveraging and neglecting terms of higher order, the resulting expressions for hot-wire sensor 1 in terms of thie mean and fluctuating velocity components are

$$
\begin{align*}
\bar{Q}_{1}^{2} & =a_{4} \bar{U}^{2}+b_{4} \bar{V}^{2}+c_{4} \bar{W}^{2}+d_{4} \overline{U V}+e_{4} \overline{U W}+f_{4} \overline{V U}  \tag{18}\\
q_{1} & =\frac{1}{2 A_{1} 1 / 2}\left[u\left(2 a_{4} \bar{U}+d_{4} \bar{V}+e_{4} \bar{W}\right)+v\left(2 b_{4} \bar{V}+d_{4} \bar{U}+f_{4} \bar{W}\right)+\right. \\
& \left.+w\left(2 c_{4} \bar{W}+e_{4} \bar{U}+f_{4} \bar{V}\right)\right] . \tag{19}
\end{align*}
$$

$$
\begin{equation*}
=\frac{1}{2 A_{1} / / 2}\left(u P_{1}+v S_{1}+w F_{1}\right) \tag{20}
\end{equation*}
$$

where $P_{1}$, $R_{1}$ and $S_{1}$ denote the terms in parentheses in Eq. (19). Simultaneous so'ution of the tiree nonlinear algebraic equation like Eq. (: 8 ) will directiy yie!d the mean velocity components $\bar{U}, \bar{V}$, and $\bar{W}$. Now, looking at (20) we find that it contains, as constants, the oredetermined mean velocities and those of geometry. The variables remaining are the instantancous fluctuations about the mean. Eq. (20) and the corresponding equations for the other two sensors will again yield three Tinear equations for the instantaneous turbulence velocities. Thus, the instantanecus details can be recovered in this methed aiso, as was done in Hethod $A$. Since the equations for $u, v, w$ are linear in nature, this procedure requ res much less computer time (less by an order of 10) than method $A$, and also ensures absoluce convergence. However, it should be reminded that this method involves the process of linearization and is therefore acceptable only for flow situations wher the linearization assumption is valid.

### 2.3 Method C

If we consider Eq. (20) again, square it and time-average, we get

$$
\begin{align*}
\bar{q}_{1}^{2} & =\frac{1}{4 A_{1}}\left[u^{2} p_{1}^{2}+\bar{v}^{2} S_{1}^{2}+\bar{w}^{2} R_{1}^{2}+2 \overline{u v} p_{1} \xi_{1}+2 \overline{u w} p_{1} R_{1}+\right. \\
& +2 \bar{v} v S_{1} R_{1} ; \tag{21}
\end{align*}
$$

Also for the effective shear stresses, the resulting expression is

$$
\begin{gather*}
\bar{q}_{1} q_{2}=\frac{1}{4 A_{1} 1 / 2 A_{2}^{1 / 2}}\left[\bar{u}^{2} P_{1} P_{2}+\bar{v}^{2} S_{1} S_{2}+\bar{w}^{2} R_{1} R_{2}+\overline{u v}\left(P_{1} S_{2}+\right.\right. \\
\left.\left.P_{2} S_{1}\right)+\overline{u w}\left(P_{1} R_{2}+P_{2} R_{1}\right)+\overline{v w}\left(S_{1} R_{2}+S_{2} R_{1}\right)\right] \tag{22}
\end{gather*}
$$

Similarly, one can obtain two such equations for each of the other two sensors. There would be a set of six linear algebraic equations with known coefficients, which can be solved once for all, for the six Reynold's stress components $\overline{u^{2}}, \overline{v^{2}}, \overline{w^{2}}, \overline{u v}, \overline{v w}$, and $\overline{u w}$ in terms of the corresponding effective components as sensed by the three sensors. This method does not allow the instantaneous $u, v$, and $w$ values to be recovered. It is, however, very economical computationally. This procedure is very similar to that used by Gorton and Lakshminarayana (1976), and has been used in the present studies chiefly to provide a comparison with methods $A$ and $B$. In fact, the linearization assumption implied in this is less restrictive than in the method of Gorton and Lakshminayama.

### 2.4 Description of the Probe

Two prubes were used for the test of the methodology and softriare. One of these (Probe 1) was developed in house and the other (Probe 2) was a commercial probe, Disa triaxial hot-wire probe. Probe 1 has dimensions and sensor orientation, as shown in Fig. 2a. It can be seen that the ratio of probe diameter to sensor support length is such that significant probe interference may be present. This probe was a first attempt at constructing a 3 -wire probe at the Institute. The Disa
probe, Fig. 2b, has a much longer sensor support as well as a small head diameter. Sensing elements in the Disa probe sense over the central 1.25 mm , whereas the in-house probe senses over its entire length (approximately 2.1 mm ). This is due to the fact that the Disa sensor is a gold plated tungsten wire of 5 micron diameter, with the central portion etched to give a 1.25 mm sensing length. The in-house sensor is a bare platinum-rhodium wire of the same diameter. It is obvious that the spatial resolution of the Disa probe is much better. However, for flow measurement where spatial resolution was not crucial, the inhouse probe proved to be quite adequate.

### 2.5 Signal Processing: Hardware and Software

Each hot-wire sensor of the triaxiai probe was operated by a Disa model 55 M 10 constant temperature anemometer. These instantaneous output voltages from the three anemometers were filtered using a high pass filter of $36 \mathrm{db} / o c t a v e$ roll-off, set at 0.1 Hz . The voltage signals were then amplified by 3 Preston SW 8300 amplifiers and sampled at a rate of 50 samples per second for 50 seconds by a Preston GMAD-3 analog to digital converter (with a 3-channel'sample and hold' front end). The data were stored on disk file on an HP-1000 minicomputer. At the end of the experiment the data were recalled from the disk file and processed to obtain the velocities. A block diagram of the instrumentation is shown in Fig. 3.

The computer program used for data acquisition records the time and date when the samples were taken. This information is used to apply corrections for sma!l drifts in probe calibration caused by temperature
changes in the tunnel and probe contamination. This is done by calibrating the probe at the beginning and end of the experiment and applying a simple linear correction to the calibration coefficients with respect to time. The corrected calibration constants are then used to convert the voltages to effective velocities in each case. Time-associated changes in the calibration for the hot-wire anemometry are therefore eliminated.

Double precision arithmeti, was used in calculating accumulated sums while computing average valuse where the numerics were suspected to exceed the single precision register size. This way, round-off and truncation errors were held to a minimum. Any errors still remaining have to be associated with the uncertainties in the overall process (i.e., probe orientation, tunnel speed, ... etc).

### 2.6 Accuracy of the Software

The stability of the Newton-Raphson algorithm used in Method A depends on the "closeness" of the initial velocity quesses as well as the convergence criteria specified. Tests on the program itself were performed to provide values of accuracy and reliability. In these tests, dummy values for the effective velocities, computed (independently) for a set of known velocities, were input to the program. The resultant velocities output by the program were compared with the correct values. It was found, that, in each case, the dlgorithm converged to the correct value in approximately 5 iterations with a convergence criterion of one
part in one thousand. The initial guess needed to be only of the same order of magnitude as the final converged correct value. In order to make the scheme converge quicker, the equations were analyzed and certain characteristics were noted. The predominant term in the three equations (Eqs. 5), were those that involved the $U^{2}$ terms and if the effective velocities were averaged and divided by the sum of the coefficients the resultant velocity was found to be within $2 \%$ of the final value. For the other two velocities the relative magnitudes of the effective velocities were noted and the appropriate sign for one of the two could be recovered. The remaining velocity was guessed as the mean of all the velocities in that direction by averaging first the effective velocities and then using that solution as the starting point for the discrete solution in question. Furthermore, a relaxation factor was applied to $U$ such that instabilities were damped. These techniques helped in arriving at the final result for the particular sample set. However, if convergence was not achieved, the set would be discarded and calculation would proceed to the next sample set. This dropout rate would be recorded and later examined to decide whether or not the results should be accepted. Typical dropout rates were .05 pir cent.

### 2.7 Triaxial Hoi Wire Probe Calibration

Each sensor of the probe was calibrated at the beginning and end of an experiment. The probe was traversed to the freestream in the wind tunnel with the probe axis aligned with the flow direction. This condition represents a limiting case for the three simultaneous equations, namely the equivalent velocities sensed by each sensor must be equal.

- This is true, of course, if the sensors are all symmetrically oriented with respect to the probe axis. A computer program was developed exclusively for probe calibration.

For the hot-wire anemometer, the output voltage across the sensor varies as a function of velocity, this variation being given by the well-known King's law. This law most commonly takes the form

$$
\begin{equation*}
E^{2}=A+B V^{n} \tag{23}
\end{equation*}
$$

where n is generally between . 45 and .55 .
The computer program developed would sample the sensor output at 50 samples per second for 50 seconds. These vaiues were time averaged to obtain $E$. The calibration was carried out at several wind tunnel velocities. The computer program would fit Eq. (23) to these data by a least square procedure. This method of calibration required a minimum amount of time due to the fact that the three sensors were calibrated simultaneously.

### 2.8 Validation Tests of Hot Wire Anemometry

Two of the well known two-dimensional flow fields, namely, a flat plate boundary layer and a wake were used to validate the performance of the triple-sensor probe and the related software and hardware. The results obtained by the methods $A, B$, and $C$, in each case were compared with those obtained earlier in the same flow with a conventional $x$-wire probe by Sastry (1981). In the case of the flat plate boundary layer
measurements, the triaxial probe results were also compared with the data of Klebannof (1955). In this situation, however, only the in-house probe was tested.

The flat plate used by Sastry (1981) was mounted as in the earlier experiments. Velocity traverses were made in the boundary layer (at -76.3 mm from the leading edge), and in the wake (at 177.8 mm from the trailing edge). Various tests were then performed on the probe and the software to determine uncertainties, accuracies, as well as the sensitivity to errors in nrientation of the probes with respect to the flow field. Effect of sample sizes and rates on drop out rates (in Method A) and accuracy was also studied.

Sensor orientation angles were measured to be within $\pm 1^{\circ}$ in the $\theta$ plane and $\pm 3^{\circ}$ in the $\psi$ plane (see Fig. 1). Effects of uncertainties in probe orientation were estimated by sirulating these uncertainties in probe yaw, pitch and roll in the processing algorithm. Values determined from the adjusted geometric inputs deviated from the best known inputs substantially enough (approximately 5 to $10 \%$ ) to verify the uncertainty ranges and their centers. Sample size and rates were varied also to find an efficient rate and size. The values used by Sastry (1981) were rechecked to see if such rates and si es were adequate for the three dimensional probe, and were, in fact, found to be so. The convergence criterion specified for the Newton-Raphson scheme was $0 . i \%$ of the freestream velocity, and further constraining (reducing the value) resulted in negligible changes (less than $2 \%$ ) in the shear stresses (the most sensitive of the results).

### 2.9 Test F.zsults

It is seen from Fig. 4 that at $y / \delta=.2$ in the boundary layer, the predominant turbulent shear stress $\overline{\mathrm{UV}}$ (obtained by Method A) is $25 \%$ below the data of Klebanoff (1955), whereas the results obtained by Method $C$, developed by the authors and very similar to that of Gorton and Lakshimaryana (1976) appear to agree very well with the data of Klebanoff. Intensities $u^{\prime}$ and $v^{\prime}$ determined by Method A or $C$ show smaller departures from the data of Klebanoff than in the case of the shear stress $\overline{\mathrm{uv}}$ (see Fig. 5). Hoivever, the most disturbing observation was that the nonprincipal shear stresses $\overline{u v}$ and $\overline{v w}$ (which should be zero) were found to be as large as the principal component $\overline{\mathrm{uv}}$ at some y locations. Results in the wake (shown in Table Al) for the in-house probe were found to be even worse in this regard, even though values for the principal stress $\overline{u v}$, were found to be of proper magnitude and to deviate by less than $20 \%$ from those of Sastry (1981).

The Disa triaxial probe 55 Pg 1 was traversed only in the wake because of probe support and mounting problems encountered in the boundary layer traverse. Principal shear stress component $\overline{u v}$, shown in Fig. 6 is in good agreement with the previous measurements. However, it was found that even though the nonprincipal stresses had decreased substantially in comparison with the inhouse probe, they are still as high as $30 \%$ of $\overline{u v}_{\max }$ in the worst case. Also shown for this case is a comparison of the three methods of data reduction presented earlier. It is seen that methods $A$ and $B$ agree well with each other and also with the data of Sastry (1981), especially in the evaluation of $\overline{u v}$. On the other hand, the time-averaged method C , while still giving good results
for $u^{\prime}, v^{\prime}$, and $w^{\prime}$, estimates the shear stresses very poorly (see Figs: $6,7,8,9)$. Since method B takes about $1 / 10$ the computer time of method A and shows comparable accuracy, it is concluded that this is the most useful procedure for data reduction.

Looking at the spatial resolution of the two probes, Disa 55P91 and the in-house probe, and the corresponding values of the $\overline{v w}$ and $\overline{u w} o b-$ tained using them provides some insight into the requirements for threedimensional flow measurement. The in-house probe with a spatial resolution of approximately 5 mm has failed to determine turbulent quantities in the moderately high shear regions. On the other hand, the DISA 55P91 probe with a resolution of 3.2 mm has given better results. It is therefore recommended that a probe of even smaller size (such as the custom made Disa subminiature triaxial probe with a resolution of 1.2 mm ) be used in future measurements. It is felt that such a probe is needed to measure $\overline{u w}$ and $\overline{v w}$ to a better level of accuracy.

### 2.10 Accuracies

An uncertainty of $\pm .007$ meters per second, due to micromanometer error, is inherent in the calibration procedure. Other uncertainties are associated with the software as well as with probe size and orientation. The convergence criteria for velocities in the software is less than . 007 U . This is the maximum uncertainty associated with software, hence, remaining inaccuracies are best described as being well within the uncertainty associated with geometry. However, in high shear regions, the errors extend beyond this and are due to size and displacement effects. Those errors due to misalignment and sensor angle error are approximately $.5 \%$
in the mean velocity components and $5 \%$ and $15 \%$ for the intensitites and shear stresses respectively. However, because of the finite size, the accuracies are reduced even further. While the probe size effect is not significant on the intensities $u^{\prime}, v^{\prime}, w^{\prime}$, the estimated values of $\overline{v w}$ and $\overline{u w}$ may be in erior by as much $3 \pm 35 \%$ of $\overline{u v}_{\text {max }}$.

## CHAPTER III

## DESCRIPTION OF AIRFOIL EXPERIMENTS

In selecting a suitable airfoil for the study, several design requirements had to be satisfied before construction could begin. The test section of the wind tunnel used is octogonal throughout, with a distance of 1.67 m between the flats. This dimension provided the constraint for the span of the airfoil. The requirement of 'infinite' spail condition; restricted the maximum chord. By combining these two constraints with the range of velocities possible in the wind tunnel used, the sweep angle, span and chord could be specified.

The profile selected is the NACA 0012 symmetric airfoil - a well studied profile. For the purpose of fabrication the coordinates describing the profile geometry were taker, fror Abott and Van Joenhoff (1059).

A sweep angle of $30^{\circ}$ was chosen, so that a sizeable crossflow component would be present at the trailing edge. The construction of the airfoil is shown in Fig. 10. Blocks of aluminum, milled to size, were secured to a central frame 12.5 mm thick aluminum plate, to form spanwise ribs. A wooden nose cone was attached to the central frame with pressure taps on both sides. Two aluminum sheets ( 0.75 mm ) used as skins, were wrapped around the aluminum blocks on each side of the frame and were feathered to give a nominal trailing edge thickness of 1 mm . Pressure taps of 1.2 mm diameter (totaling 92) were provided on both sides
of the airfoil along midspan, and along two cff-mid span planes 15 cm away on either side of the midspan plane.

The airfoil was mounted in the wind tunnel with its spanwise direction located $30^{\circ}$ from vertical. A false floor and boundary layer fence were used to remove end-wa?l boundary effects. This provided conditions close to the desired infinite airfoil configuration (see Fig. 11). Tubing and other mounting fixtures were concealed so as not to affect the flow under measurement. Flow visualization studies made, using wool tufts did not indicate any areas of separation or tunrel-wall effects. These tests also showed that the fence had no appreciable effect on the flow and hence was not used during the main experiments.

## 1. Coordinate System

The coordinate system used is a right handed orthogonal system, the freestream velocity direction corresponding to the $x$-direction. The $y$ direction is normal to the airfoil surface and towards the near wall of the tunnel. The z-direction is perpendicular to the tunnel ceiling. This system is shown in Fig. 10 and it should be noted that the sensor's direction cosines are referenced to this system.

Traversing the $\dot{\text { u }}$ oundary layer and wake in the $y$-direction was accomplished by a servo-controlled stepping notor with a range of 20 cm and ar accuracy of .025 mm . The traverse is movi.ble in the $x$-direction for the longitudinal repositioning. A list of traverse stations and their locations are shown in Table 1.

## 2. Pressure Measurements on the Airfoil

Measurements of static pressure were used to insure that the airfoil was mounted at zero incidence, and more importantly to provide the ą:ual pressure distribution on the airfoil. A total of 92 pressure taps were used for this purpose. In the midpları of measurement, 22 taps were located on each side of the airfoil with additional 11 taps $\pm 15 \mathrm{~cm}$ from the midplane. These were also located on both sides of the airfoil.

To perform the measurements a 48-port-selecting scani-valve was used in conjunction with an alcohol micromanometer. The micromanometer had an accuracy of .025 mm and was used for monitoring tunnel speed as well as measuring the static pressure.

Static pressure taps on the two sides of the airfoil and directly opposite to each other were used for initial alignment of the airfoil. For a further and more accurate check on alignment complete scans of all the pressure taps were made at a tunnel speed of 21.84 meters per second. The coefficients of pressure, $C_{p}$, defined as

$$
\begin{equation*}
C_{p}=\frac{P-P_{\infty}}{\frac{1}{2} \rho U_{\infty}^{2}} \tag{24}
\end{equation*}
$$

are shown in Figs. 12 and 13. It is seen that the static pressure distributions at the three spanwise locations coincide reasonably well. This is an indication that the airfoil can be considered to be nearly 'infinite' in span. However, it should be noted that this test will be
pressure, which, according to Bussman and Ulrich (see Schlichting (1974)) is the region where natural transition occurs on an NACA 0012 airfoil. Past experience at the Institute has shown that sand paper strips perform better than trip wires often used for this purpose. With this tripping, the boundary layer on the airfoil was found to be fully turbulent and the Reynolds number based on the momentum thickness at 2.5 cm upstream of the trailing ed! :vas 7280.

## 5. Experimental Procedure

The wind tunnel was started and allowed to warm up before starting the experiment. Typically, the temperature rise over the duration of an experiment was of the order of $1^{\circ} \mathrm{C}$. First, yaw probe traverses were made in the bcundary layer and wake. These included profiles on either side of the airfoil to check for symmetry and profiles at spanwise locations 15 cm above and below the midspan plane to check for infinite conditions. These locations are indicated in Table 1, referred to eariier. The yaw probe traverses were followed by traverses of the Disa 55p9i triaxial hot-wire probe in the wake. Location of the wake centerline was inferred from the minimum in :he velocity distribution across the wake.

Uncertainty in the velocity measurements is estimated to be $2 \%$. Static pressure is considered to be accurate within 15-20\% and the flow angle within $0.5^{\circ}$. Hot wire uncertainties are primarily due to errors in probe alignment and resolution. Preset yaw is expected to be less than $+0.5^{\circ}$ and : freset error in pitch is likely to be less than $+2^{\circ}$. Inaccuracies in probe manufacturing were of the order of $2^{\circ}$ as determined by direct measurement through a steroscope with a graduated reticie. Roll
inaccuracies were eiminated through the Disa type mounting that fixed the roll axis to the traversing mechanism. The overall uncertainties in the turbulence measurement were $u^{\prime}, v^{\prime}, w^{\prime}: 5 \%, \overline{\mathrm{Lv}}: 15 \%, \overline{\mathrm{vw}}: 30 \%$. The uncertainties in the $\overline{u w}$ were too large to be acceptable. Hence, these data are not discussed in this report. Further miniaturization of the probe is necessary in order to improve these results. Also not used in the discussion are the data on the various triple correlations such as $\overline{u^{3}}, \overline{u^{2} v}$, etc., since their accuracy has not been established so far. All the data are, however, reported in the Appendix, for the purpose of documentation.
followed by boundary layer traverses in "the midplane" and "off-planes" for this aspect to be further verified.

## 3. Flow Pigle Measurement and Calibration

Magnitude and cirection of the velocity vector were determined by a directionally sensitive thee-hole yaw probe. This probe of similar type and size to thrt used by Ramaprian, Patel, and Choi (1978) and was calibrated in, a similar manner. However, several changes were made to make the calioration more accurate. The probe was also caiibrator to yield the local static pressure, in addition to the magnitude and direction of the velocity. Details of this calibration, inciuding probe dimensions are shown in Fig. 14. The yaw probe outputs we read via a set of three STATHAM pressure transducers, amplified, scanned, digitized and reco:-ded by the HP/ 1000 computer. The results were averaged over 20 seconds to obtain the mean pressure for each tube. The transducer calibration was repeated several, times each day to see if the instrumentation and methodology used ensured repeatability.

## 4. Experimental Conditions

The experiments were performed at a tunnel velocity of about 22 meters per second corresponding to a trailing edge Reynol is number of approximately $1.36 \times 10^{6}$. To fix transition te t.jrbulence, a boundary layer 'trip' consisting of a strip of 20 -grid sand japer, 15 cm wide and extending over the entire span, was glued to the surface on boch sides of the airfoil at a distance of 20 cm from the leading edge. The location of the sand paper corresponds roughly to the region of lowest

# CHAPTER IV <br> RESULTS AND DISCUSSION 

## 1. General

The results of the experiments are presented and discussed in this chapter. A complete set of tabulated experimental data is provided in Appendix A.
2. Mear: Flow Measurements

The variation of static pressure across the tr: ? ing edge boundary layer and wake is shown in Fig. 15 . It is interesting to see that the static pressure varies across the wake and reaches a maximum value at the wake centerline. This trend persists mildly even at the last measuring s+ation in the wake. This observation confirms similar otservations made by Sastry (1981) in the developing two-dimensional wake of an airfoil. The static pressure variation across the wake can be viewed as the result of the interaction among the boundary layer, wake and the external inviscid flow. These data should, therefore, provide a good test case for interactive calculation methods.

Figure 16 shows the distributions of the longitudinal velocity component $U$ across the boundary layer in the trailing edge region. Specifically, the velocity profiles at $x / L=-0.220$ and $x / L=-0.014$ are shown. It is seen that these profiles, especially the latter, resemble a typical distribution in a moderate adverse pressure gradient. The
profile does not show any evidence of flow separation. The variation of the crossflow angle $\gamma$ in the boundary layer at the same locations are also shown in Fig. 16. It is seen that the crossflow increases towards the trailing edge. Also, at the trailing edge, the crossflow is strong ! $\gamma=15^{\circ}$ ) pmough to introduce significant three-dimensional effects into the wake flow. Furthermore, at $x / L=-.124$, spanwise variations in the boundary layer are seen to be small (see Fig. 17).

Figures 18 and 19 show the distributions across the wake of the chordwise velocity component. A close study of these profiies shows that there is slight asymmetry in the flow. The reason for this is not known. Also, the profile at the last station has been shown only for half of the wake since the traverse could not be extended to cover the other half of the wake at this station. The crossflow angle, $(\gamma)$ profiles across the wake shown in Fig. 20 exhibit similar properties as the chordwise velocity profiles, namely decay of the crossflow angle and the increase in spread with distance downstream. It is also seen that at the last measuring station, the crossflow is verv small, indicating that the mean flow is virtually two-dimensional beyond this point.

## 3. Integral Parameters

We now define the foliowing integral parameters for the wake:

$$
\begin{align*}
& \text { displacement thickness } \delta_{1}=\int_{-\infty}^{\infty}\left(1-\frac{U}{U_{\infty}}\right) d y  \tag{25}\\
& \text { momentum thickness } \theta_{11}=\int_{-\infty}^{\infty} \frac{U}{U_{\infty}}\left(1-\frac{U}{U_{\infty}}\right) d y \tag{26}
\end{align*}
$$

momentum thickness $\theta_{21}=\int_{-\infty}^{\infty} \frac{W}{U_{\infty}}\left(1-\frac{U}{U_{\infty}}\right) d y$
shape factor $\quad H=\frac{\delta_{1}}{\theta_{11}}$
The distributions of these parameters with dcwnstream distance are shown in Figs 19-22. First, Fig. 21 shows the variation of $\theta_{11}$. In this figure the reference length scale used is the momentum thickness ${ }_{f}$ at the far wake, defined as

$$
{ }^{\theta_{\mathrm{f}}}={ }^{\theta} 11 \text { at the last measurement location }
$$

From Fig. 21 certain aspects concerning the flow may be noted. Clearly, the momentum thickness increases in the boundary layer due to the adverse pressure gradient. There is a small increase in ${ }^{9} 11$ again in the very near wake due to the finite thickness of the trailing edge of the airfoil. Continuing into the wake we see a gradual decrease of $\theta_{11}$ brought about by the combined contributions from the favorable longitudinal pressure gradient and the rotation of the velocity vector towards the streamwise direction. The effect of the gradual weakening of three-dimensionality in the flow is also seen from Fig. 22 which shows that the momentum deficit thickness $\theta_{21}$ decreases nearly to zero at the last station. Unfortunately, measurements could not continue beyond $x / L=1.928$. It is assumed that the flow will be two-dimensional beyond this distance and hence $\theta_{11}$ will remain constant at its value $\theta_{f}$ at the last station. This value is, therefore, used as the reference length scale in the rest of the figures.

Figure 23 shows the variation of the displacement thickness along the streamwise direction. The initial increase in the boundary layer is caused again by the adverse pressure gradient in the trailing edge region. Subsequent reduction in $\delta_{i}$ is primarily due to the evolution oi the wakelike profile from the original boundary-layer-lik? profile. This trend is, of course, also influenced by factors such as the prevailing favorable pressure gradient and gradual decay of three-dimensionality. The changes in profile shape are more clearly seen from Fig. 24 , winch shows the variation of the shape factor $H$ with $\theta_{f}$. It is seen that the shape factor, at the trailing edge, has a value of about 1.5 , indicating that tine boundary layer is only under a moderate adverse pressure gradient. The shape factor is seen to drop quick', and at $x / \theta_{f}=326$, is only slightly greater than 1, its asymptotic value at very large distances.

## 4. Growth and Decay of the Velocity Defect and Wake Width

The decay of the maximum longitudinal yelocity defect $w_{x 0}$ is shown in Fig. 25 in the usual coordinates used for two-dimensional wakes. The distance downstream is normalized with $\theta_{f}$. Also shown in the figure is the asymptotic decay law [See Sastry (1981)].

$$
\begin{equation*}
\left(\frac{U_{\infty}}{w_{x 0}}\right)^{2}=.4\left(\frac{x}{\theta_{f}}\right) \tag{29}
\end{equation*}
$$

for two-dimensional far-wake. The data seem to indicate that the decay rate approaches this law near the farthest downstream station $\left(x / \theta_{f}=326\right)$. Likewise, Fig. 26 shows the half-width of the wake again plotted in the
usual two-dimensional coordinates. The theoretical two-dimensional far-wake behavior given by [Sastry (1981)].

$$
\begin{equation*}
\left(\frac{b_{x}}{\theta_{f}}\right)^{2}=.355\left(x / \theta_{f}\right) \tag{30}
\end{equation*}
$$

is also shown in the figure. The wake width does not show any clear signs of approach to two-dinensional far-wake state, though changes in slope qualitatively resemble two-dimensicnal flow results. In Figure 25, if the dotted line can be assumed to repiresent the decay process at large distances, it intersects the x-axis, at a virtual origin corresponding to $x / \theta_{f}=70$. If, this point is joined to the last data point in Fig. 26 the line is approximately parallel to the asymptotic growth line. This is a very rough indication that the wake has perhaps reached very nearly the two-dimensional far-wake state at $x / \theta_{f}=326$. This observation is in conformity with the findings of Sastry (1981), who inferred that the two-dimensional wake behind a flat plate reaches an asymptotic state around $x / \theta_{f}=350$. The present data, are inadequate to confirm this definitively. More closely spaced data as well as data extending further downstream are needed to substantiate the present observations.

Figures 27 and 28 show the corresponding decay and growth rates of parameters $w_{z 0}$ and $b_{z}$, the maximum defect velocity in the z-direction and half-width of this profile. The velocity defect in this case is normalized by $U_{\infty} \sin \phi$ (where $\phi$ is the swend angle), since this is the maximum value that $w_{z}$ can attain. It is seen that the crossflow defect velocity decays very rapidly compared to the longitudinal velocity. The wake half-width behaves qualitatively like $b_{x}$ but is seen to be much larger.

## 5. Self Similar Velocity Profiles

The velocity data plotted in the conventional self-similar coordinates are shown in Figs. 29 and 30. In Fig. 29 showing the streamwise velocities, it is seen that beyond $\quad x / L=0.138\left(x / \theta_{f}=23.4\right)$ the profiles become very nearly self-similar and that the profiles follow closely the profile for asymptotic two-dimensional far-wakes. The spanwise velocity profiles shown in Fig. 30 also appear to approach selfsimilarity. Compared to the $w_{x o}$ profile, however, the half-width of this self-similar profile is approximately $25 \%$ larger.

## 6. Turbulent Stress Measurement

The turbulence quantities measured using the three-dimensional hot wire probe can be analyzed in different ways. First of all, the three intensities $u^{\prime}, v^{\prime}, w^{\prime}$ and the three shear stresses $\overline{u v}, \overline{u w}$ and $\overline{v w}$ are normalized with respect to the freestream velocity so tha - general trends may be observed. These results are shown in Figs. 31 thru 36. In general, the intensities are of similar magnitude and behave in a similar way. This indicates that the turbulence is not too far from isotropy. Furthermore, in Figs. 31, 32, and 33 it is seen that the profiles exhibit a minimum in the center and are fairly symmetric about the wake centerline. The shear stresses $\overline{u v}, \overline{v w}$ and $\overline{u w}$ are shown in Figs. 34, 35 and 36 . It is seen that $\overline{u v}$ and $\overline{\mathrm{w}} \mathrm{W}$ arenearly anti-symmetric about the wake centerline changing sign as expected. In contrast, however $\overline{u w}$ measurements are considerably in error. This stress has almost the same magnitude as $\overline{\mathrm{vw}}$. This trend is similar to that observed during the validation tests in the flat plate wake mentioned earlier. The ratio of $\overline{v w}$ to $\overline{u v}$ is shown in

Fig. 37. This ratio is a good measure of the three-dimensionality of the shear flow field. It is seen that at the last station $\left(x / \theta_{f}=326\right), \overline{v w}$ is of the order of $15 \%$ of $\overline{u w}$, indicating the near two-dimensionality of the flow field. Next, the decay rates for $u_{\text {max }}^{\prime} \overline{u v}_{\text {max }}$ and $\overline{v w}_{\text {max }}$ are shown in Fig. 38. The centerline turbulent intensity is seen to increase continuously with distance downstream. The normalized Reynolds shear stress $\frac{\overline{u v}_{\text {max }}}{w_{x 0}}$ is also increasing, but at the last measuring station its value is only slightiy higher than the value of 0.48 for the asymptotic two-dimensional far-wake (Sastry 1981). The shear stress $\frac{\nabla W_{\max }}{2}$ on the other hand is found to be slow?y decreasing at large values of $x / \theta_{f}$. It is difficult to predict from the present measurements whether it would continue to decrease or reach an asymptotic value farther downstream. The developmient of $\overline{u v}_{\text {max }}$ is faster then in the case of the flat plate wake of Sestry or Pot (1979), but otherwise qualitatively similar in that it approaches the asymptotic vaiue from "below". This is in contrast with the behevior of either the wake behind a cylinder or that behind a two-dimensional dirfoil at incidence [Sastry (1981)]. In both these cases, $\overline{u v}_{\text {max }}$ decreases 0 the asymptotic value from 'above". The trends in ali these cases are consistent with the relative influence of the wall on the velocity profile just before the flow detaches from the wake generator.

## 7. Self Preservation of the Turbulence Profiles

Figures 39-41 show the profiles of $u^{\prime}, \overline{u v}$ and $\overline{v w}$ in self preserving coordinates. It is seen tha: the profiles are evolving continuously.

It appears that the turbulence properties have not attained self-similarity even at the last location. However, since $\overline{u v}_{\max }$ at $x / \theta_{f}=326$ is nearly equal to the dsymptotic value, further data in the far wake are needed to detemmine whether the turbulence profiles indeed have attained the asymptotic self-preserving form.

## 8. Eddy Viscosity Results

The preseni data can be used to calculate the eddy viscosity in both $x$ and $z$ directions defined by

$$
\begin{equation*}
\varepsilon_{x}=\frac{\overline{u v}}{\partial U / \partial y} \tag{31}
\end{equation*}
$$

and

$$
\begin{equation*}
\varepsilon_{z}=\frac{\overline{\bar{l} w}}{\partial w / \partial y} \tag{32}
\end{equation*}
$$

Plots in self-preserving coordinates of both $\varepsilon_{x}$ and $\varepsilon_{z}$ are shown in Figs. 42 and 43. From this it is seen that both viscosities are evolving continuously. The eddy viscosity in the $z$ direction exhibits considerable scatter, but is still of similar magnitude as $\varepsilon_{x}$. As to whether or not the flow has reached asymptotic behavior it is difficult to confirm because of the lack of adequate data points between $x / \theta_{f}=$ 180 and 326. However, the value of $\varepsilon_{x}$ obtained from the present experiments can be compared with the asymptotic values given by Schlicnting (:-3) and Sastry (1981). At $x / \theta_{f}=326$ the present value for $\varepsilon_{x} / \mathrm{U} \mathrm{\theta}_{\mathrm{m}}$ is .039. This compares with .032 quoted by Sastry and a value of .044 by

Schlichting. Therefore, it seems reasonable to say that at $x / \theta_{f}=326$ the streamwise flow is sufficiently close to the two-dimensional far-wake.

If we now assume an average eddy viscosity from Figs. 42 and 43 for each $x / L$ location, we should in turr be able to recalculate the stresses. Typical results for $\overline{u v}$ and $\overline{v W}$ are shown in Figs. 44-47. These figures allow us to compare the extent to which it is realistic to use a constant eday viscosity model to describe the shear stress in the wake. For the shear stress $\overline{u v}$ it is seen that the agreement with experiment is good. The plots for $\overline{\mathrm{VW}}$ show the difference in result obtained by using $\varepsilon_{z}$ or $\varepsilon_{x}$ (scalar eddy viscosity assumption). Average values of the viscosity for the given $x$-iocation were used in eacn case. The comparison (considering experimencal scatter) shows little difference and hence we conclude, that for calcuiation purposes, $\varepsilon_{x}$ can be used as a scalar eddy viscosity at least for mild crossflows.

## 2. The Structure Parameter

Nash proposed for three-dimensional turbulent shear flow, the following relations (see Nash and Patel 1972):

$$
\begin{align*}
& |\overline{u v}|=a \bar{q}^{2} \frac{\partial H / \partial y}{\left[\left(\frac{\partial U}{\partial y}\right)^{2}+\left(\frac{\partial \dot{W}}{\partial y}\right)^{2}\right]}  \tag{33}\\
& |\overline{v W}|=a \bar{q}^{2}\left(\frac{\overline{\partial W}}{\partial y}\right) /\left[\left(\frac{\overline{\partial U}}{\partial y}\right)^{2}+\left(\frac{\partial \bar{W}^{\partial y}}{\partial y}\right)^{2}\right]
\end{align*}
$$

where $q^{2}=\left(u^{\prime 2}+v^{\prime 2}+w^{\prime 2}\right) / 2$ and

# "a" is often called the structure parameter. For two-dimensional flows, the proposal takes the form 

$$
\begin{equation*}
|\overline{u v}|=a \bar{q}^{2} \tag{35}
\end{equation*}
$$

Results of calculating 'a' from Eq. (34) are shown in Fig. 48. It is Seen that except near the centerline and in the outer intermittent region, 'a' ras a nearly constant value. Typical comparisons of the value of 'a' ubtained using the alternate definitions Eq. (34) and Eq. (35) (three-dimensional and two-dimensional definitions) are shown in Fig. 49. It is seen that tho-dimensional definition is adequate to evaluate 'a' in this mildly three-dimensional fiow. Again, to assess the validity of the model for practical use in three-dimensioral flows an average value for a (= .15) was used in Eq. (34) to calculate $\overline{\mathrm{vw}}$ for a few nearwake locations. These results are compared with measurements in Fig. 50. The agreement is seen to be moderate considering the uncertainties in measurement. Therefore, it can be concluded that the structure parameter model in the for of Eos. (33) and (34) is reasonably satisfactory for describing dimensional wakes.

## 10. Conciusions

The following conclusions can be arrived at from the study reported in this thesis.

1. A triaxial hot-wire probe can be used, with one of the techniques developed in the present study, for turbulence measurements in three-dimensional shear flows. The accuracy of such measurements,
critically depends on probe orientation, probe size and correct knowledge of sensor angles.
2. The present techniques $A$ and $B$ are superior to that used by Gorton and Lakshminarayana (1976) both in accuracy and versatility.
3. At present, the measurement of $\overline{u w}$ is unsatisfactory. Further miniaturization of probe is necessary to get better results.
4. The study of the developing three-dimensional wake behind an infinite swept airfoil shows that the static pressure varies across the near-wake with the maximum occurring at the wake centerline.
5. The longitudinal and crossflow wake defect components reach near self-similar distributions within about 10 momentum thicknesses downstream of the trailing edge.
6. Three-dimensionality of the flow becomes negligible within about 325 momentum thicknesses. Also, at this distance, the wake begins to exhibit many of the mean and turbulent flow properties of two-dimensional far-wakes.
7. A scalar eddy viscosity, constant across the wake, can be used to describe both the shear stress components $\overline{u v}$ and $\overline{v w}$ reasonably well.
8. The turbulence in the wake exhibits structural similarity in the manner proposed by Nash.

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## Table 1

Boundary Layer and Wake Traverse Stations

| Station \# | $x / L$ | $x / \theta_{f}$ |
| :--- | :--- | :--- |
|  |  |  |
|  | -.220 | -36.7 |
| 2 | -.124 | -21.0 |
| 3 | -.014 | -.37 |
| 4 | +.014 | 3.37 |
| 5 | .027 | 4.57 |
| 6 | .041 | 6.94 |
| 7 | .055 | 9.32 |
| 8 | .138 | 23.4 |
| 9 | .275 | 46.6 |
| 10 | . .669 | 67.6 |
| 11 | . .992 | 112.0 |
| 12 | 1.928 | 168.1 |
| 13 |  | 326.7 |



Figure 1. Sensor coordinate system and sensor direction cosines

a. IIHR constructed three-dimensional hot-wire probe

b. Disa 55P91 three-dimensional hot-wire probe

Figure 2. Hot-Wire Probes

Figure 3. Instrumentation block diagram


Figure 4. Distribution of $\overline{u v}$ in the flat plate boundary layer at -76.3 mm


Figure 5. Distribution of $u^{\prime}$ and $v^{\prime}$ in flat plate boundary layer at -76.3 mm


Figure 6. Distribution of $\overline{u v}$ in the flat plate wake at 177.8 mm


Figure 7. Distribution of $u^{\prime}$ in the flat plate wake at 177.8 mm


Figure 8. Distribution of $v^{\prime}$ in the fiat plate wake at 177.8 mm


Figure 9. Distribution of $w$ in the flat plate wake at 177.8 mm

Figure i0. Airfoil cross-section (dimensions in mm)

Figure 11. Airfoil in tunnel with measurement station locations


Figure 12. Coefficient of pressure variation on both sides of wing


Figure 13. Coefficient of -ressure variation on both sides of wing on à s .nwise basis


$$
\begin{aligned}
& \text { AKP }=\frac{n_{1}^{-D} 3}{\left.P_{2}-\frac{P_{1}+P_{3}}{2}\right]} \\
& \text { AK23 }=P_{2}-P_{3} \\
& 1 / 2 \rho Q^{2}=A K 23 /\left[.51885+.023779 * \gamma-.0000142 * \gamma^{2}-\right. \\
& \left..0000035 * \gamma^{3}-.0000001 * \gamma^{4}\right] \\
& \gamma=-1.4345+10.28219 * \text { AKP }+.04959500 * \text { AKP }^{2} \\
& -.0246590 \text { * AKP }{ }^{3}-.0148545 \text { * AKP }{ }^{4} \\
& P_{\text {static }}=P_{2}-\left[1.0088296-.0011767 * \gamma-.0002621 \gamma^{2}\right] * 1 / 2 \rho Q^{2}
\end{aligned}
$$

Figure 14. Yaw probe description and calibration









$$
\begin{aligned}
& \text { (2.8: } \\
& \text { (W) } 2^{01} \times{ }^{\circ}
\end{aligned}
$$








Figure 30. Self-similar spanwise velocities














$\left(2^{s / 2} 2^{(1) \wedge n}\right.$
(1.8


$$
\left(2^{s /} 2^{W}\right) \underline{M n}
$$





APPENDIX
TABLES OF EXPERIMENTAL DATA

## KEY TO CYMBOLS USED IN DATA TABLES

| Beta | angle between mean flow vectors and the chordline |
| :---: | :---: |
| Gama | angle between mean flow vectors and the free streamline |
| Delta*sub x | $\delta_{1}$ meters |
| Delta*sub z | $\delta_{2}$ meters |
| Theta sub x | $\delta_{11}$ meters |
| Theta sub $z$ | $\delta_{22}$ meters |
| Theta sub xz | ${ }^{\text {o }} 12$ meters |
| Theta sub $\mathrm{z} \mathrm{\lambda}$ | $\delta_{21}$ meters |
| Pstatic | static pressure in inches of alcohol |
| $u /$ Uinf | loca? chordwise velocity/reference velocity |
| w/Jinf | locai spanwise velocity/reference velocity |
| U bar 1 | U/U from 3-D hot wire probe |
| $V$ bar | normal velocity component/ $U_{\infty}$ from 3-D hot wire probe |
| W bar | W/U from 3-D hot wire probe |
| uu bar | $\sqrt{\bar{u}^{2}} / U_{\infty} \times 100$ |
| IV bar | $\sqrt{\mathrm{V}^{2}} / U_{\infty} \times 100$ |
| ww bar | $\sqrt{W^{*}} / U_{\infty} \times 100$ |
| uv bar | $\frac{\overline{u v}}{U_{\infty}{ }^{2}} \times 10^{4}$ |
| $Y(m m)$ | distance from wake centerline or wall in millimeters |
| uw bar | $\frac{\overline{u w}}{U_{\alpha}^{2}} \times 10^{4}$ |

vw bar
uuu bar
vov bar
www bar
uuv bar
uuw bar
vvu bar
vvw bar
wwu bar
liwv bar
uvw bar
$\frac{\overline{V W}}{U_{\alpha}{ }^{2}} \times 10^{4}$
$\frac{\bar{u}^{3}}{U_{\alpha}^{3}} \times 10^{5}$
$\frac{\bar{v}^{3}}{U_{\alpha}^{3}} \times 10^{5}$
$\frac{\bar{w}^{3}}{U_{\infty}} \times 10^{5}$
$\frac{\overline{u^{2} v}}{U_{\alpha}^{3}} \times 10^{5}$
$\overline{\frac{u^{2} w}{U_{a}^{2}}} \times 10^{5}$

- $\frac{u v^{2}}{U_{x}^{2}} \times 10^{5}$
$\frac{\overline{v^{2} w}}{\overline{U_{\infty}^{3}}} \times 10^{5}$
$\overline{\frac{U w^{2}}{U_{\alpha}^{3}}} \times 10^{5}$

$$
\frac{v w^{2}}{w_{x}^{3}} \times 10^{5}
$$

$$
\frac{\overline{U V W}}{U_{\sigma}{ }^{3}} \times 10^{5}
$$

## Table Ai

Symmetric Flat Plate Wake Data $(x=177.8)$ using IIHR probe

| $\bar{Y}_{\text {mm }}$ | $u / U_{\text {inf }}$ | ${ }^{\text {u }}$ bar | $\mathrm{Vv}_{\text {bar }}$ | ${ }^{W} w_{\text {bar }}$ | $\mathrm{uv}_{\text {bar }}$ | ${ }^{\mathbf{V W}}$ bar | ${ }^{\text {uw }}$ bar |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 9.62 | . 81 | 6.273 | 6.460 | 3.476 | +10.879 | -14.358 | -7.370 |
| 5.81 | . 76 | 6:219 | 6.764 | 3.755 | 10.091 | -13.694 | -0.821 |
| 2.00 | . 72 | 5.381 | 6.709 | . 317 | 5.551 | -6.361 | -15.538 |
| . 73 | . 709 | 5.098 | 6.422 | 4.713 | 4.056 | -1.497 | -17.972 |
| -. 54 | . 710 | 4.933 | 6.009 | 5.075 | -2.966 | +4.693 | -17.587 |
| -1.30 | . 711 | 5.131 | 5.758 | 5.349 | -4.775 | 6.742 | -17.594 |
| 2.06 | . 714 | 5.292 | 5.645 | 5.613 | -7.217 | 9.636 | -17.105 |
| 2.32 | . 717 | 5.530 | 5.572 | 5.724 | -12.607 | 10.014 | -15.674 |
| 4.09 | . 729 | 5.799 | 5.084 | 5.873 | -12.315 | 12.802 | -14.390 |
| 5.36 | . 742 | 5.086 | 4.965 | 5.926 | -15.098 | 14.263 | -13.412 |
| 7.90 | . 769 | 6.409 | 4.931 | 5.829 | -15.113 | 14.137 | -11.317 |
| 12.98 | . 839 | 6.259 | 4.819 | 5.279 | -13.286 | 12.468 | -9.424 |
| 18.06 | . 882 | 5.397 | 4.504 | 4.570 | -9.85i | 8.262 | -6.942 |
| 23.84 | . 930 | 4.648 | 3.660 | 3.646 | -6.184 | 5.515 | -4.753 |
| 33.0 | . 922 | 2.199 | 2.063 | 1.837 | -1.058 | 1.273 | -2.113 |
| 48.24 | 1.00 | . 477 | . 792 | . 603 | -. 002 | . 099 | -. 339 |

# Airfoil Boundary Layer Data at $x / L=-.220$ (central plane) 



Integral Parameters
Melta*sut $x=4776659 E-02$
Deltaxsub
Theta sut
$=.34150475 E-02$
Theta sub $z=.6247743 \mathrm{E}-03$
Theta sut. $2=2=2697000 \mathrm{E}-01$
Theta sut $2 x=.1681640 \mathrm{E}-03$
Shape factar $H=1.298$
Foundaru Laver Profile

| $\because \mathrm{Mm}$. | u'linf | Gama | w/Uinf | Heta | Pstatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 51 | 454 | 3. 37 | $0 こ 7$ | 33.37 | $-.0016$ |
| . 76 | 5010 | 4.8 | $03 \%$ | 34.19 | 0007 |
| 1.0 | 52 | \& 57 | 042 | 34.5 | 0043 |
| 1.27 | 535 | 4.51 | 042 | 34.51 | 0062 |
| 1.53 | 5 | 4.43 | 043 | 34. $0^{4}$ | 0093 |
| 2.03 | 官74 | 4.47 | 045 | 34.47 | 0021 |
| 3.54 | 596 | 4.35 | 045 | 34.35 | 0.0000 |
| 5.05 | E16 | 4.20 | 045 | 34.20 | $-0012$ |
| 3.5 | 639 | 4.10 | 640 | $\because 4.10$ | -0095 |
| 4.83 | 676 | 3.65 | 144 | 37.69 | -.0132 |
| 6.10 | 716 | 3.36 | 041 | 33.26 | - 0252 |
| \%. 37 | 739 | 3.07 | 039 | 쿤.7 | -. 0161 |
| 8.64 | 76 | 2. 68 | 036 | 32.63 | - 0206 |
| 5.91 | 795 | 2.31 | 032 | 3231 | -.0301 |
| 12.45 | 842 | 1.77 | 026 | 31.77 | $-0.094$ |
| 14.95 | 859 | 1.19 | 016 | 31.10 | -0303 |
| it. Ef | $\bigcirc$ | . 55 | 019 | 30.55 | - 16A 4 |
| 26.61 | 984 | 46 | 003 | 30.18 | - 045 |
| 26. 0 | $\pm 0110$ | 0.65 | 0.010 | 30.00 | - 0380 |
| 30.23 | 1. 000 | 0.00 | 0.000 | 30.00 | $-039 \%$ |

TABLE A3
Airfoil Boundary Layer Data at $x / L=-.124$ (central plane)
RUNT OF RUN 29381.

Urefernce (m/5) = $21 . \therefore$

Uinfinity ( $\mathrm{m} / \mathrm{s}$ ) $=22.31$
Integral Parameters
Dejta*sut $x=6346047 \mathrm{E}-02$
Dejta*sut z $=.7913063 E-012$
Theta sut $z=.7212564 \mathrm{E}-03$
Theta sut $x==7305337 E-01$
Shape Factor $H=1.443$

Eoundary Laver Frofile

| Y' MM | u'Uinf | Gama | w/Uinf | Hota | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| . 51 | 416 | 4.80 | 035 | 34.80 | 0428 |
| 1.27 | 400 | 6.09 | 051 | 36.09 | 0612 |
| 2.03 | 526 | 6. 27 | 058 | 36.27 | 0635 |
| 3.30 | 578 | 5.87 | 059 | 35.87 | 0604 |
| 5.08 | 630 | 5.20 | 057 | 35.20 | 0583 |
| 7.11 | 678 | 4.36 | 059 | 34.36 | 0520 |
| 10.16 | 742 | 3.27 | 042 | 33.27 | $045 \%$ |
| 1245 | 796 | 2.44 | 034 | 32.44 | 0407 |
| 16.29 | 887 | 1.20 | 015 | 31.20 | 0235 |
| 25.91 | 977 | 1.04 | 001 | 30.04 | 0173 |
| 33.53 | $9{ }_{9} 9$ | $-13$ | -. 002 | 29.67 | 0 28 |
| 43.69 | 997 | -. 28 | -. 005 | 29.72 | 0214 |
| 58.53 | ¢¢8 | $-.15$ | -. 003 | 29.65 | 0182 |
| 74.17 | 1.000 | 0.00 | 0.000 | 30.00 | 0145 |

TABLE A4
Airfoil Boundary Layer Data at $x / L=-.124$ ( $-6^{\prime \prime}$ plane)


TABLE A5
Airfoil Boundary Layer Data at $x / L=-.124$ ( $+6^{\prime \prime}$ plane)

```
RUN # OF RÜN 29884
Urefernce (m/s) = 21.84
Uinfinitu (m/E) = 22.65
Integral Farameters
Delta*sut x = 5613346E-02
DElta*sut z =.4455480E-01
Theta sut x = . 3972160E-02
Theta sut z = .9642757E-03
The:a sut xz = 3918079E-01
Theta sut :x = .2592422E-03
Shape factor H = 1.413
```

Eoundary Laver Profile

| $Y \mathrm{~mm}$ | u/Uinf | Gama | W/Uinf | Feta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 51 | 41.8 | 5.36 | 035 | 75.36 | 1560 |
| . 76 | $40^{\circ}$ | 5.58 | 045 | 35.58 | 9640 |
| 1.03 | 496 | 5.91 | 051 | 35.91 | 1640 |
| 1.27 | 520 | 5.84 | 053 | 35.84 | 1000 |
| $\frac{1}{2} 9$ | $54 \%$ | 5.6 | 056 | 35.80 | 1620 |
| 3.05 | 605 | 5.36 | 05 | 35.36 | 1540 |
| 4.32 | 635 | 4.91 | 055 | 34.91 | 1520 |
| 5.59 | 6.1 | 4.46 | 052 | 34.46, | 1450 |
| 6.86 | 691 | 4.19 | 051 | 34.18 | 1560 |
| 8.13 | 795 | 3.67 | 040 | 33.63 | 1450 |
| 10.67 | 776 | 2.89 | $0{ }^{\circ}$ | 32.88 | 1380 |
| 13.21 | 825 | 2. 19 | 032 | 32.19 | 1260 |
| $\frac{95}{19} 95$ | 884 | 1.65 | 025 | 31.65 | 1260 |
|  | 963 |  | 009 | 30.49 | 1160 |
| 20.75 | 987 | 10 | 003 | 30.10 | 1910 |
| 36.07 | 1.001 | 02 | 000 | 30.09 | 1130 |
| 42.42 | 1.000 | -0\% | 01.1 | 29.97 | 1070 |

TABLE A6

## Airfoil Roundary Layer Data at $x / L=-.014$ (far side)



```
\refernce (m/0) = 214 = 21.84
Uinfinitv(m!́)=22.23
```

Integral Parametere
Delta*sut $x=7026338 E-0:$
Delta*sut $z=524734$ PE-01
Trieta sut $x=.4873316 \mathrm{E}-02$
Theta sut $z=1921229 E-02$
Theta sut $x z=.4601236 E-01$
Theta sut $x=.565845 \mathrm{E}-0 \frac{1}{3}$
Shade factor $H=1.442$

Houndary Layer Profije

| Y Mm. | u'Uirif | Gamä | wiUirf | FEta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Tif | 475 | 12.62 | 07 | A2. 6 | 205 |
| 76 | 496 | 13.53 | 109 | 42.53 | 2296 |
| 1.02 | 455 | 12.55 | 110 | $A 2.56$ | E.29 |
| 1. 27 | 507 | 12.23 | 110 | 42.23 | 2206 |
| 1.52 | 517 | 15.01 | 130 | 42.01 | 人ator |
| 1.72 | $5 こ 3$ | 11.93 | 110 | 41.95 | 2270 |
| 2.03 | 534 | 1 1. 44 | 108 | A1. 44 | 930 |
| 2. 54 | 551 | 10.85 | 105 | 4 (1) 8 ? | 2204 |
| 3.05 | 56 | 10.72 | 106 | 40.72 | 203? |
| -5. 56 | 573 | 10.33 | 104 | 40.33 | 2016 |
| 4.06 | 584 | 10.05 | 103 | 40.03 | 628 |
| 5.08 | 607 | 9.46 | 1.01 | 39.46 | 5201 |
| 6.10 | 637 | 8.55 | 096 | 33.55 | 21.40 |
| 6.13 | 678 | 7.36 | 088 | 37.36 | 2074 |
| 19.16 | 702 | 6.72 | 083 | $36.7 \%$ | 2188 |
| 12.70 | $76 \%$ | 526 | 070 | 35.26 | 4979 |
| 45.24 | 815 | - 17 | 059 | 34.17 | 1005 |
| 20.30 | 98' | 2.60 | 140 | 32.60 | 1695 |
| 20.40 | 940 | 1.43 | 024 | 31.83 | 49\% |
| 30.48 | \%78 | 79 | 013 | 30.79 | 1597 |
| 35.56 | 992 | 47 | 000 | 30.47 | 1. 98 |
| $4 \% \cdot 13$ | $99 \%$ | . 11 | 002 | 30.11 | 1314 |
| 50.80 | 1. 000 | 0.00 | 0.000 | 30.00 | 1. O 1.6 |



TABLE A8

## Airfoil Wake Data at $x / L=.014$ (central plane)



## TABLE A8（Continued）

| $Y \mathrm{Mm}$ | U／UArif | Gumer | w／Un的 | He¢a | $\mathrm{F}=$ ¢atar |
| :---: | :---: | :---: | :---: | :---: | :---: |
| －$\overline{5}$ | $4 \% 9$ | 14.16 | ¢11 | 44.16 | 2154 |
| －-76 | 442 | 14.45 | 112 | 44.15 | ¢10 |
| － 5.27 | 465 | 13.30 | 110 | 43.34 | ご96 |
| －1．78 | 46 | 13．86 | 196 | $43 \cdot 66$ | \％280 |
| － 54 | $4{ }^{4} 9$ |  | 194 | 42.8 | ¢180 |
| －3．30 | 5？ | 12.20 | ${ }^{1} 1414$ | $4{ }^{4} 4$ | 3154 |
| －A． 32 | 594 | 11.29 | 1910 | 41.29 |  |
| －5， 5 | 500 | 10.47 | 107 | 40.6 |  |
| －6．80 | 614 | 9.69 | 103 1097 | 39.60 | 6164 |
| $-7.97$ | 836 | $\frac{8}{7} 80$ | 097 093 | 37.87 | ¢08\％ |
| －9．65 | 703 7 | 6.86 | 0¢ | 36.8 | 2074 |
| －1 $\frac{1}{3} .97$ | 747 | 5.76 | 075 | 35.7 | 2024 |
| $-17.76$ | 610 | 4.23 | 060 | 34.23 | 196 |
| －31．59 | ．86\％ | 2.91 | 044 | 39.91 | 1050 |
| $-26.67$ | 6 | 1.57 | 026 | 31.57 | 164 |
| $-31.76$ | 976 | 89 | 016 | 30.89 | 1545 |
| $-39.37$ | －9\％ | ． 41 | 007 |  | 4 ${ }^{4} 0$ |
| －46．99 | 1．0100 | 0.10 | $0.00 \%$ | 30.12 | 125 |
| $-54.63$ | 1．000 | 0.00 | 0.000 | $\cdots 0.00$ | 1．… |

TABLE A9
Airfoil Wake Data at $x / L=.027$ (central plane)

| ROTE OF Rut eqeat |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & \text { Urefern } \\ & x \text { loca. } \end{aligned}$ | $\begin{gathered} (m / 0) \\ .027 \end{gathered}$ | 21.84 |  |  |  |
| Uirifinity (m/s) $=22.26$ |  |  |  |  |  |
| Integraj. Farameters |  |  |  |  |  |
|  |  |  |  |  |  |
|  |  |  |  |  |  |
| Shape Factor $H=1.453$ |  |  |  |  |  |
| Wake Frofile |  |  |  |  |  |
| $Y \mathrm{MM}$. | u/Uirif | Gama | W/Uinf | Heta | Pstatic |
| -54.56 | 1.005 | 35 | 006 | 30.35 | 1. 1317 |
| -44.78 | 98 | 09 | 905. | 30.30 | 1.453 |
| - 56 | 96 | 86 | 015 | 30.36 | 1508 |
| -21.84 | $8 \%$ | $\bigcirc$ | 05 | 83.9 | 1275 |
| -56.76 | 887 | \% 64 | 075 | \%E.93 | 9.769 |
| -11.68 | 675 | 7.05 | 08 | 8.05 | 1980 |
| $\cdots 6.60$ | 63. | 6. 3 | 05. | 38.27 | 1970 |
| -5.08 | $59 \%$ | 9.39 | $09 \%$ | 35.36 | 2049 |
| $\cdots 4.06$ | 558 | 19.76 | 106 | 40.64 | 2047 |
| -2.05 | 545 | 10.84 | \%s? | 40.74 | 2013 |
| -3.0E | 505 | 11.13 | 099 | 41.13 | 2009 |
| 0.00 | 500 | 11.15 | 697 | 41.12 | 2015 |
|  | 509 | 19.70 |  | 39.94 | 151. |
|  | 536 |  | 093 | 39.48 | 1947 |
| 3.98 | 59 | 8.89 | 09 | 38.89 | 1 c |
|  |  | Q. 37 | 188 | $3 \mathrm{3C} .17$ | 1971 |
| \% $=7$ | 6 m | 7.3 | 164 | 37.38 | 1048 |
| 51 | $6 \%$ | ¢. 0 | 878 | $36.9 \%$ | ${ }^{18 \%}$ |
| 12.5 | 746 | 4.96 | 05 | $\square 4.06$ | 377 |
| 14.99 | 788 |  | 056 | 3\%.39 | 3505 |
| 2! | 98 | 1. 20 | 019 | 31.20 | 1401 |
| 注•者, | 96\% | -43 | 0107 | 30.44 | 1304 |
| 4.39 | 996 | $1 \%$ | 003 | 30.10 | 1370 |
| AC. 6 | . 090 | (1.00 | 0.000 | 30.00 | 1113 |

## Airfoil Wake Data at $x / L=.027$ ( $-6^{\prime \prime}$ plane)

## RUNTE OF RUN 3981

Urefernce $(\mathrm{m} / 5)$
X joca $=21.84$
Uinfinity ( $\mathrm{m} / \mathrm{s}$ ) $=2194$

Inteoral Farameters
yeltawsub $x=\{524869 E-01$
Dejta*sut $z=.1011579 E+00$
Theta sut $y=.104263 E-01$


Shade factor $H=1.463$
Wake Profile

| $Y \mathrm{MM}$ | u'Uirif | Gama | w'Uinf | Eeta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | 1.003 | 00 | 000 | 30.00 | 1.243 |
| 44.20 | 1.009 | 14 | 002 | 30.14 | 1367 |
| 36.55 | 998 | 26 | 005 | -1.68 | 1375 |
| 29.96 | $87 \%$ | 2.84 | 035 | 32.27 | 1.46 |
| 56.26 | 806 | 3.70 | 05 | 33.70 | 1854 |
| 13.2 | 758 | 4.69 | 062 | 34.69 | 18988 |
| 11. 18 | 717 | 5.59 | 070 | 35.58 | 1939 |
|  |  | 789 | 085 | 37.84 | 1950 |
| 4.3 | ¢ 7 | 6.63 | 085 | 38.63 | 2031 |
| 3.05 | 55 | 8.98 | 087 | 38.98 | 2009 |
| 2.29 | 539 | \%. 50 | 090 | 39.50 | 2050 |
| 1. 52 | 526 | 40.88 | 089 | 40.88 | 20¢ |
| 0.00 | 516 | $10 . t 5$ | 097 | 40.69 | 2057 |
| -26 | 54 | 10.67 | 095 | $40 \cdot 6$ |  |
| -3. 56 | $55 \%$ | $10.8 \%$ | 109 | 40.81 | 205 |
| -3.56 | 563 | 16.19 | 101 | 41.19 | 2109 |
| -5. 53 | 6.4 | E.56 | 152 | 38.56 | 2014 |
| -7.11 | 642 | 7.85 | 088 | 37.85 | 906 |
| -9.6. 9 | 695 738 | 6.478 | 070 | 35.43 | 196 |
| -1 ${ }^{-1} 97$ | 926, | 3.5 | 05 |  | 1860 |
| -20.35 | 94.5 | $2.2 \%$ | 035 | 32.23 | 1740 |
| -29.97 | 974 | 1.11 | 119 | 31.11 | 1.630 |
| -37.59 | . 997 | 6 | 011 | 30.80 | 1435 |
| -45.at | 1.000 | 54 | cor | $\bigcirc 4$ | 13. |


| $\stackrel{10}{\text { RUN }}$ |
| :---: |
|  |  |



Uinfinjtu (mis) = 2215

Integial Parametere
Deltax 5 ub: = 1577514E-01
Deitawsut i =. $1282576 \mathrm{E}+00$
Theta sub $x=. \hat{0} 6 E 428 E-0$ S.
Theta sut $z=.3429951 E-02$
Theta sub $x=1135999 E+09$
Theta sub $2 x=107508 \mathrm{E}=02$
Shape Factor $H=\frac{2}{4} 47^{\circ}$
Wate Profile

| $\bar{Y} \mathrm{MM}$. | u/Uinf | Gama | n/Uinf | Feta | Pstatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 60.96 | 1050 | -. 18 | -. 003 | 29.82 | 1106 |
| 5. ${ }^{4}$ | 1.009 | - 18 | -. 002 | 39.89 | 1207 |
| 45.72 | 1.006 | -. 03 | -. 000 | 29.97 | 129 |
| 3 Sa 10 | 1.000 | 07 | 001 | 30.07 | 1421 |
| 30.43 | 954 | . 36 | 006 | 30.36 | 1463 |
| 20.80 | 507 | 1.39 | 025 | 31.39 | 1695 |
| 12.70 | 743 | 4. | 054 | 4 | 2012 |
| 4 C .16 | 70 | 5.0 | 065 | 35.07 | 2054 |
| 7.6 | $65 i$ | G.06 | 069 | 30.09 | 2117 |
| 5.84 | $\bigcirc \mathrm{C}$ | 713 | 076 | 37.12 | 2205 |
| 4.06 | 589 | 0.01 | 080 | 38.01 | 2244 |
| $\cdots \mathrm{F}$ | 54.3 | 8. 61 | $0 ¢$ | 38.61 | 2957 |
| 2. 03 | 527 | 9.11 | 095 | 37.11 | 2250 |
| 1.52 | 590 | 9.29 | $0 ¢ 5$ | 30.29 | 2248 |
| 1.02 | 513 | 9.74 | 083 | 39.74 | 2271 |
|  | 509 | 9.9 | 109 | 39.97 | 2957 |
| 0.00 | 510 | 10.13 | 051 | 40.13 | 2275 |
| -. 76 | 513 | 46.52 | 095 | 40.52 | 2349 |
| -1.78 | 523 | 18.81 | 098 | 40.61 | 2300 |
| $-3.05$ | 509 | 40.24 | $0 \%$ | 40.2A | 256.1 |
| -4.32 | 574 | 9.75 | 098 |  |  |
| -6. 10 | 614 | 8.5 | 093 | 38.59 | 236 |
| - 96.64 | 658 | 7.20 | 084 | 37.63 | 2298 |
| -19.53 | 0.9 | 3 \% |  | 35.5 | 5 |
| -29.61 | Est | 2.3 | 035 | 気: ${ }^{3}$ | 4975 |
|  | 96: | \% | 116 | 30.9 | 1866 |
| - 4.6 | 095 | 24 | 004 | 3084 | 1694 |
| -53 if | 1. 000 | 0: | 01 | 30.05 | 1470 |
| -60.74 | 1.000 | 0.011 | 0.000 | 30.00 | 1369 |

Airfoil Wake Data at $x / L=.041$ (central plane)


```
RUNM*OF 12
DATE OF RUN 3981
Urefernce (m/E) = 055 = 21.84
Uinfinity (m/s) = 22.14
Integral Parameters
Delta*sub x= 14509999E-01
DElta*sut z =.1879447E+00
Theta sut x = .1045328E-01
Theta sut z = . 2475021E-02
Theta sut xz= = 17423005+00
Shape Factar H = 1.3%7
Wake Profile
```

| $Y \mathrm{~mm}$. | u/Uinf | Gama | w.Uinf | Eet | Stat |
| :---: | :---: | :---: | :---: | :---: | :---: |
| -33.31 | 1.005 | -. 44 | -. 008 | 29.56 | 0964 |
| -73.15 | 1.002 | -. 39 | -. 007 | 29.62 | 0961 |
| -62.99 | 1.000 | -. 23 | -. 005 | 29.72 | 1130 |
| -52.83 | 995 | -. 10 | -. 002 | 29.90 | 1259 |
| -42. 67 | . 997 | 00 | 000 | 30.00 | 1278 |
| -33.76 | 992 | $2 i$ | 004 | 30.21 | 1357 |
| -2. 16 | 969 | 70 | 012 | 30.70 | 1353 |
| -19.84 | 989 | 4.95 | 130 | 31.96 | 1659 |
| -14.73 | 819 | 3.26 | 047 | 33.26 | 1710 |
| -10.92 | 752 | 4.55 | 060 | 34.55 | 1898 |
| -7. 11 | 696 | 5.65 | 068 | 35.65 | 1707 |
| -4.57 | 635 | 6.98 | 078 |  | 1869 |
| - 2.54 | .602 | 7.69 | 081 | 37.69 | 1872 |
| -1.02 | . 573 |  | 084 |  |  |
| 0.00 | 565 | 8.30 | 08, | 38.35 | 1897 |
| 1. 51 | . 575 | 8.83 | $08 \frac{2}{1}$ | $38.20{ }^{38}$ | 1885 |
| 1.52 | .569 | 3.21 | 082 | 38.21 | 1855 |
| 2.54 | 575 | 9.14 | 085 | 39.14 | 1833 |
| 4.06 | 587 | 7.65 | 081 | 37.85 | 1870 |
| ¢. 10 | 6 67 | 6.88 | 076 | 36.89 | 1776 |
| a. 64 | 672 | 5.93 | 070 | 35.93 | 1725 |
| 16.26 | 798 | 3.31 | 046 | 33.31 | 1669 |
| 23.88 | 910 | 1.40 | 022 | 31.40 | 1449 |
| 34.14 | 978 | 42 | 007 | 30.42 | 1514 |
| 44.20 | 989 | 19 | 003 | 30.19 | 1371 |
| 54.36 |  | 16 | 003 | 30.16 |  |
| 64.52 | .996 | 9 | 003 | 30.17 30 | 1178 |
| 56.90 | 1.000 | -2.97 | -1.147 | 27.03 | 0995 |

## TABLE Al3 (Continued)

| Turbulerice Data |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y / L$ Location $=.055$ <br> Reference Velocity $=21.8 \mathrm{~m} / \mathrm{s}$ |  |  |  |  |  |  |
| Wake Profile |  |  |  |  |  |  |
| $\begin{array}{r} Y \\ 43 M \\ 33.69 \\ 25.93 \\ 20.93 \\ 18.29 \\ 15.75 \end{array}$ | Ubar $1.000$ | Ubar 000 | Wtar | yubar | $\begin{aligned} & \text { unbar } \\ & \hline 78 \end{aligned}$ | Whbar |
|  | . .994 | . 001 | 003 | 1.440 | 4.703 | 1.772 |
|  | 952 | -. 000 | 061 | 4.244 | 3.817 | 4.210 |
|  | 895 | -. 002 | . 015 | 5.674 | 5.120 | 5.715 |
|  | 859 | -. 002 | 019 | 6.303 | 5.701 | 6. 126 |
|  | 825 | 004 | . 026 | 6.788 | 6.052 | 6. 570 |
| 10.67 | 740 | 011 | 038 | 7.724 | 6. 646 | 7.241 |
| 8.13 | 694 | 015 | 045 | 7.897 | 6.794 | 7.426 |
| 5.59 | . 644 | 018 | 046 | 7.723 | 6.993 | 7.435 |
| 3.05 | . 602 | 023 | 052 | 7.037 | 8.802 | 7.022 |
| - 51 | -572 | 033 | 066 | 6.032 | 6.534 | 7.261 |
| -5.73 | . 579 | 040 | 081 | ¢. 036 | 6.564 | 7.598 |
| -4.57 | 613 | 051 | 095 | \%.876 | 6.709 6.894 | - 955 |
| -7.1i | 662 | 058 | 065 | 8.313 | 7.037 | 8.485 |
| -5.65 | 751 | 062 | 077 | 8.085 | 6.848 | S. 242 |
| - 12.19 | . 755 | 064 | 063 | 7.538 | 6.677 | 7.601 |
| -14.73 | \% 93 | 065 | $0{ }_{0} 9$ | 7.335 | 6.555 | 7.592 |
| - 19.81 | 876 | 067 | 034 | 6.087 | 5.70\% | 6.063 |
| -22.35 | 905 | 066 | 034 | 5.335 | 5190 | 5.314 |
| -27.43 | 962 | 067 | 007 | 3.637 | 3.728 | 3.361 |
| - 35.05 | .993 | 067 | -. 011 | 1.072 | 1.617 | 1.191 |
| -45.21 | . 995 | 688 | -. 015 | 481 | . 817 | 663 |

## TABLE Al3 (Continued)

| Profile | conti | a* | $=$. |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $Y \mathrm{Mm}$ | uruar | untorar | vebar | uuubar | uuubar | whwbar |
| 43.69 | -. 049 | 027 | 03 | 002 | -. 001 |  |
| 33.53 | -. 547 | 627 | 14 | 630 | 400 |  |
| 25.91 | -5 893 | 2. 112 | 49 | 9.619 | 1.883 | -4.075 |
| 20.83 | -1 5.54 | - 805 | 2.67 | 10.649 | $\frac{1}{5} .743$ | -4.621 |
| 18.29 | -14.57 | 1.907 | 3.59 | 11.487 | 5.278 | -6. 527 |
| 15.75 | -38.74 | 1. 053 | 5.15 | 9.449 | -. 708 | -5.240 |
| 13.24 | -21.14 | -. 557 | 6.56 | 6.700 | 6.456 | 6.142 |
| 10.6 | -25.83 | - -3.301 | 5.85 | 6.253 | 8. 880 | -2.970 |
| 5.59 | -25.30 | -2.125 | 5.79 | -1204 | 10.199 | 203 |
| 3.05 | - 59.99 | -6. 697 | 4.23 | - 15.932 | 4.814 | 2. 375 |
| -. 76 | -5.256 | -4.799 | -2. 2.67 | -8.657 | $-5.607$ | - $\mathrm{S}^{\text {. }}$ S 97 |
| -2.03 | 13.694 | -6.051 | -4.59 | -16.48 | -3 931 | -4. 271 |
| -4.57 | 23.008 | -1i. 20 | -9.72 | -19.39 | -9.6.99 | -3.926 |
| -7. 11 | 25. 356 | - 17.82 | -13.46 | -7.130 | -5.490 | $\frac{1}{2} .066$ |
| -19.65 | 24.408 | -13.18 | -10.5\% | 3.687 | -4.966 -7.206 | 2.195 -.188 |
| -14.73 | 20.784 | -15.07 | -10.83 | 11.366 | -10.72 | 3.759 |
| -17.27 | 16.665 | -9.671 | -5.9\% | 12.142 | -5.797 | 2.727 |
| -19.81 | 15.484 | -8. 233 | -6 43 | 10.567 | -5.905 | 8.045 |
| -22. ${ }^{-25}$ | 12.033 | -6.786 | -5 63 | 10.977 | -4.757 | 6. 872 |
| -35.05 | -. 226 | - -3.97 | - -36 | , 14\% | -. 4 - | . 609 |
| -45.21 | -. 021 | -. 043 | -. 04 | -. 001 | -. 002 | 011 |

## TABLE Al3 (Continued)

| Pr | continued | et $X / L$ | 055 |  |
| :---: | :---: | :---: | :---: | :---: |
| $Y$ mm | uuvbar | uuntar | uoubar | uuwbar |
| 43.69 | -. 001 | 001 | -. 001 | 004 |
| 33.53 | - .188 | 210 | . 460 | 134 |
| 25.91 | -3.774 | 384 | 4.183 |  |
| 20.83 | -4.480 | 343 | 5.362 | - $-\frac{1}{3} .671$ |
| 15.75 | -4.451 | - 488 | 5. \% $^{\text {¢ }}$ | - 2.328 |
| 13.21 | -4.200 | -3.980 | 6. 057 | -2.786 |
| 10.67 | -3.003 | -2.856 | 3.313 | - 2.144 |
| 8. 13 | -. 308 | -1. 178 | . 474 | -2. 308 |
| 5.59 3.05 | 2.917 | .802 -.454 | -5.411 | -1.817 |
| . 51 | 3.282 | 1.600 | -9.519 | 829 |
| -. 76 | -3.500 | ${ }^{2} .164$ | - 10.83 | 2.869 |
| -2.03 | -8.344 | 3.249 | - 40.45 | 1.491 |
| -4. 57 | -7.127 | 8.861 | -4.762 | -4. 309 |
| -9.65 | 2.665 | 542 | 2.788 | $-5.753$ |
| -12.15 | 3.205 | -3.025 | 4.332 | -5.779 |
| -14.73 | 5.442 | -2.941 | 5.843 | -4.322 |
| -19.27 | 4. 7171 | - 2.349 | 6.175 | - 5.445 |
| -22. 35 | 6.059 | -1.892 | 6.331 | -4.016 |
| -27. 43 | 3.053 | -1.522 | 4.503 | -2. 555 |
| - 45.05 | 06 | - 1180 | - 1505 | - 0001 |

## TABLE Al3 (Continued)


table al4

## Airfoil Wake Data at $x / L=.138$ (central plane)



## TABLE A14 (Continued)

| Turbulence Data |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $x /$ L Location $=.130$ <br> Refererice Velocity $=2: .8 \mathrm{~m} / \mathrm{s}$ |  |  |  |  |  |  |
| Wake Profile |  |  |  |  |  |  |
| $Y$ MM. | Ubar1.000 | Ubar | Whar | uubar | $\begin{aligned} & \text { uybar. } \\ & 772 \end{aligned}$ | wwbar |
| 44.45 |  | -. 001 | .001 |  |  |  |
| 36.83 | 997 | -.003 | $00 \frac{1}{5}$ | 970 | 1.227 | 1. 340 |
| 29.21 | . 976 | $-007$ | 003 | 2.938 | 2.937 | 3.186 |
| 24.13 | 932 | $-.007$ | 006 | 4.383 | 4.357 | 4.868 |
| 19.05 | . 875 | -. 004 | 015 | 6.142 | 5.577 | 6.109 |
| 16.51 | .843 | -.004 | 020 | 6.668 | 5.919 | 6.576 |
| 13.97 | . 802 | . 002 | 096 | 7.347 | 6.256 | 6.784 |
| 11.43 | 768 | 602 | 030 | 7.510 | 6.466 | 6.886 |
| 8.89 | 725 | 005 | 035 | 7.530 | 6.727 | 7.038 |
| 6.35 | 695 | 006 | 039 | 7.210 | 6. 454 | 6.703 |
| 3.81 | . 666 | 012 | 049 | 6.460 | 6.348 | 6.300 |
| 1.27 | . 648 | 017 | 05 | 5.804 | 6.238 | 6.215 |
| 0.00 | . 648 | 8 C | 06 | 5.801 | 6.225 | 6.513 |
| $-1.27$ | . 651 | QEi | 065 | 5.947 | 6.089 | 6.854 |
| -2.54 | . 655 | 054 | 065 | 6.191 | $6.14 \%$ | 6.783 |
| -5.08 | . 684 | 030 | 070 | 7.562 | 6.276 | 7.113 |
| -7.62 | 718 | 031 | 066 | 7.483 | 6.612 | 7.546 |
| $-10.16$ | 757 | $0 \cdot 5$ | 16.4 | 7.835 | 6.538 | 7.769 |
| -12.70 | 797 | 037 | 055 | 7.84 | 6.619 | 7.492 |
| -15.24 | 836 | 830 | 046 | 7.416 | 6.223 | 6.950 |
| -17.78 | 874 | . 040 | 077 | 6.554 | 5.751 | 6.604 |
| $-20.32$ | 906 | $0<0$ | 030 | 6.009 | 5.295 | 5.694 |
| -55.40 | 958 | 037 | $0 \pm 4$ | 4.278 | 3.985 | 4.020 |
| -30.46 | 990 | $0 \frac{3}{5}$ | -.001 | 2.306 | 2.614 | 玉. 210 |
| $-38.10$ | 959 | 035 | - 009 | . 728 | 1.119 | . 868 |
| -45.72 | 1. 000 | 035 | -.011 | . 521 | . 789 | 700 |

TABLE A14 (Continued)

|  | $\alpha$ 3 3 | 0 0 0 + |
| :---: | :---: | :---: |
|  |  | $\stackrel{\sim}{\sim}$ |
|  | E | $\square$ <br> 3 |
| -ovbonwifmo | + | - |
|  | 3 | $\stackrel{H}{3}$ |
|  |  | m |
|  | $E$ | $\square$ |
| Enunco | cr | 0 |
|  | a | $\rightarrow$ |
|  <br>  | - | $x$ |
|  |  |  |
|  |  |  |
|  | $\Sigma$ | $!$ |
| - | ${ }^{1}$ |  |
|  | 9 |  |
| ) |  |  |
|  |  |  |
|  | ${ }_{5}^{5}$ |  |
|  | c |  |
|  | - |  |
|  | ${ }^{\text {a }}$ |  |
|  | $\rightarrow$ |  |
| 1111111 |  |  |
|  | $c$ |  |
|  |  |  |
|  | c |  |
|  | $\square$ |  |
|  | $\stackrel{\square}{3}$ |  |
|  |  |  |
|  | $\Sigma$ |  |
|  | $\Sigma$ |  |
|  | $\Sigma_{0}$ |  |
|  |  |  |
|  | $\underset{\sim}{w}$ |  |

## TABLE A'4 (Continued)

| Profile | continued | X/L | 138 |  |
| :---: | :---: | :---: | :---: | :---: |
| $Y \mathrm{~mm}$ | uuubar | uuwbar | uvubar | uuwbar |
| 44.45 | -. 001 | 001 | -. 003 | 003 |
| 56.83 | - 0.088 | 024 | . 0.05 | 028 |
| 27.293 | - -4.735 | - 588 | 2. 5.435 | -1.095 |
| 19.05 | -5.021 | -2. 110 | 7.200 | -1.891 |
| $\frac{16.51}{46}$ | -3.191 | -5.483 | 4.848 | - 1.570 |
| 13.97 | -3.605 | -5.825 | 4.730 5 | $-3.189$ |
| 18.89 | 2.601 | -1.033 | 2. 253 | -1.436 |
| 6.35 | 6.625 | -. 707 | -3.610 | 453 |
| 3.81 | 7.794 | 713 | -6.815 | 1.028 |
| 1.27 | 3.210 | 1.364 | -9.432 | 2.138 |
| -0.00 | - $\frac{1}{2} \cdot 102$ | $\frac{1}{2} \cdot 284$ | -9.934 | 2. 031 |
| -1. 274 | - 2.470 | 2.913 | -8.629 | . 920 |
| -5. 08 | -7.042 | 5.685 | -8. 434 | - ${ }^{4} \cdot 408$ |
| -7.62 | -3.55 | 3.402 | -3. 33.9 | -1. 153 |
| $\cdots 10.16$ | 2.776 | 2.670 | 3.660 | -3.459 |
| -12.70 | 6.000 | -1.368 | 5.990 | -6.357 |
| - 15.34 | 8. 467 | -3.561 | 6.901 | -6.042 |
| - 20.32 | 6.543 | - 1.810 | 6.287 | -4.105 |
| -25.40 | $4.30 \%$ | -2.791 | 5.008 | - $\because .777$ |
| -30.48 | 985 | -. 975 | 1.287 | -. 802 |
| -48.72 | 001 | -. 0001 | -084 -.004 | 012 |

TABLE A14 (Continued)

| Prof | continued | $x / L$ | 1.38 |
| :---: | :---: | :---: | :---: |
| $Y$ MM. | wwutar | wwubar | unwbar |
| 44.45 | -. 005 | -. 002 | -. 007 |
| 36.83 | 2.093 | -. ${ }^{1707}$ | - 1115 |
| $22^{29} \cdot 2 \frac{1}{1}$ | 2.186 | - .807 | $\overline{5} 3 \mathrm{Cl}$ |
| 99.05 | 5.461 | -1.946 | 8.473 |
| 46.51 | 2. 466 | -1.156 | 6.856 |
| 13.97 | 2. 021 | 301 | 6.356 |
| 14.43 | 581 | 189 | 8.123 |
| 8.89 | -4.450 | 1.396 | 3.524 |
| $\frac{8.35}{} 3.8$ | - -5.411 | 1.869 2.704 | 4. 682 |
| 1.27 | -5.440 | . 310 | -3.292 |
| 0.00 | -3.704 | -3. 243 | 171 |
| - $-\frac{1}{2} \cdot 54$ | -7.0418 | -3. 5 - 8.1 | 6.080 |
| -5.08 | -6.139 | -2.059 | 6.899 |
| -7.62 | -4.458 | -. 590 | 6.842 |
| -10.16 | -1.203 | 1.012 | -1.012 |
| -12.70 | 3.147 | 4.042 | -8.345 |
| - 15.38 | $4 . \frac{1}{5} 46$ | 3.455 | -8.901 |
| -20.32 | 7.090 | 2.998 | -6.711 |
| -25.40 | 4.589 | 2. 750 | -5. 718 |
| -30.48 | 1.397 | 595 | -1.432 |
| -3a. 10 | -. 010 | 026 | -. 032 |
| -45.72 | -. 003 | 001 | 005 |

TABLE A15
Airfoil Wake Data at $x / L=.225$ (central plane)

|  |
| :---: |
| $\begin{aligned} & \text { Unefernce } \\ & \times \text { loca }=(5) \\ & .275 \end{aligned}=21.84$ |
| Uirifinity (mis) = 25.3i |
| Integral Parameters |
| De1ta*sub $x=1050085 E-01$. <br> DE1ta*sub z =.8172131E-01 <br> Theta sut $x=.8346590[-02$ |
| Theta sut $z^{2}=.2319180 \mathrm{EE} 02$ <br> Theta sub $\times 2=.7163487 E-01$ Theta sut $2 x=41984 E-03$ |
| 58 |

Wake frofile

| $Y$ MM. | u/Ujnf | Gama | w/uinf | Eeta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 30.10 | 1.000 | 87 | 012 | 30.67 | 0865 |
| 30.48 | . 987 | 89 | 015 | 30.89 | $034 \%$ |
| 22.85 | 936 | 1.41 | 023 | 31.41 | 0847 |
| 15.24 | $84 \%$ | 2.46 | 036 | 35.46 | 0947 |
| 10.16 | 794 | 3.06 | 042 | 33.06 | 0954 |
| 6.3 "' | 755 | 3.56 | 047 | 35.56 | 1015 |
| 3.91 | 735 | 3.80 | 047 | 33.7\% | 1040 |
| 2.54 | 726 | 3.82 | 049 | 33.82 | 1037 |
| 1. ${ }^{2}$ | 722 | 3.86 | 049 | 33.86 | 1047 |
| 0.00 | 75 | 3.89 | 049 | 33.89 | 1057 |
| -1.3\% | 719 | 3.80 | 149 | 33.80 | 1057 |
| -5. 08 | 734 | 3.85 | 048 | 33.81 | 1056 |
| -3.68 | 754 | 3. ${ }^{2} 3$ | 043 | 3, ${ }_{3} \cdot 6$ | - 030 |
| -11.43 | 790 | 2.76 | 039 | 䢒. 70 | 102 |
| -16.51. | 847 | 2.10 | 631 | 32.10 | 0973 |
| -23.86 | 920 | 1.19 | 019 | -31.19 | 0876 |
| -30.40 | 934 | 36 | 006 | 30.36 | 0812 |
| -39.37 | 1.000 | 0.00 | 000 | 30.00 | 0797 |

## TABLE A15 (Continued)

Turbulence Data
Xeferenceton Velocity ${ }^{275}=21.8 \mathrm{~m} / \mathrm{s}$
Wake Prafile

| $5^{\text {Mm }}{ }^{\text {m }}$ |
| :---: |
| 52.07 |
| 31.75 |
| 24.13 |
| 20.32 |
| 12.70 |
| 10.16 |
| 7.82 |
| 5.5 |
| 0.00 |
| 2. 54 |
| - 0.9 |
| -1 |
| 7 |
| 7.78 |
| 9 |
|  |
| 49. 5 |






[^0]
## TABLE Al5 (Continued)

Frafile continued at $X / L=.275$

| Y MM. | uvtiar | Uwbar | voter | uuutar | vuutar | Wwwtar |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 52.07 | 006 | -. 023 | -. 01 | 000 | -. 001 | 003 |
| 41.91 | -.051 | -. 003 | -. 01 | -. 045 | . 021 | -. 060 |
| 31.75 | $-2.453$ | 1.639 | . 19 | 4.557 | 1.699 | -1.926 |
| 24.13 | -9.112 | 1.762 | 1.95 | 13.305 | 2. 161 | -5.148 |
| 20.32 | -13. 27 | -.933 | 4.34 | 13.430 | 2.915 | -6.466 |
| 16.51 | -17.63 | -2.174 | 7.40 | 9.217 | 3.760 | -5.441 |
| 12.70 | -17.67 | -2. 539 | 7.45 | -1.671 | 4.802 | -2.570 |
| 10.10 | $-18.2$ | -2.730 | 6.44 | -8. 135 | 5.111 | 811 |
| 7.62 | -15.91 | - 5.501 | 5.60 | -7.595 | . 298 | 1.789 |
| 5.08 | $-1040$ | -3.607 | 3.25 | -9.080 | 5.573 | 1.602 |
| 2.54 | -4.087 | -4.305 | 1.30 | -6.399 | 3.501 | 1.296 |
| 0.01 | 2.821 | -2. 569 | $2 \pm$ | -5.828 | 1.926 | 075 |
| -2.54 | 7.942 | - 5.368 | - 84 | -8.859 | 2.509 | -1.802 |
| -5.08 | 43.833 | -5. 263 | -2. 26 | $-10.26$ | 4.494 | -2.376 |
| -762 | 15.025 | -7.119 | $-3.17$ | $-6.043$ | -. 355 | -3.026 |
| -10.16 | 17.619 | -6.591 | -3.37 | -. 154 | . 88 | $-353$ |
| $-13.97$ | 10.744 | $-\frac{1}{1} 0.2$ | -3.39 | 7.688 | -6.1800 | 3.368 |
| -17.70 | 15.170 | -7.134 | -1.42 | 16.395 | -4.547 | 6.841 |
| -31.59 | 10.690 | -7.574 | -1.66 | 15.791 | -4.276 | 5.619 |
| -29.21 | 2.750 | -2.022 | -1.05 | 5.076 | $-3.106$ | 3.515 |
| -39.37 | 040 | -. 147 | 02 | -. 010 | -. 001 | 150 |
| -49.53 | -.018 | . 110 | 04 | 0.000 | $-.005$ | 000 |

## TABLE A15 (Continued)

| Profile | ntinued | at X /L | 275 |  |
| :---: | :---: | :---: | :---: | :---: |
| Y MM | uubtar | uuwbar | vuubar | vuwbar |
| 52.07 | -. 000 | 0.000 | -. 002 | 004 |
| $4{ }^{4} .91$ | - 009 | $-.002$ | -.013 | -. 005 |
| 31.75 | -1.674 | . 415 | 2.249 | -. 318 |
| 24.13 | -5. 566 | - -3.305 | 6.052 | - 1.843 |
| 20.32 | -5. 5.363 | - -1.169 | 6.557 | -2. 3 - 35 |
| 12.70 | 2. 163 | -1.086 | . 918 | -1.222 |
| 10.16 | 4.336 | - 264 | -2.435 | -. 210 |
| 7.62 | 5.680 | -. 056 | -4.862 | 119 |
| 5.08 | 4.192 | 1.309 | -5.271 | 748 |
| 2. 04 | - 2.2814 | 1.249 | -5. 534 C | 705 |
| -3. 54 | -4.360 | 1.479 | -6.622 | 028 |
| -5.08 | -5.463 | 2.413 | -4.165 | $-.657$ |
| -7.62 | -2.350 | 4.173 | -2.225 | -1.487 |
| - 10.16 | 3. 877 | 1.773 | 4.186 | -1.496 |
| -17.78 | S. 736 | -1.662 | 8.204 | -3.2488 |
| -21. 59 | 5.501 | -2.744 | 5.813 | $\cdots 3.393$ |
| -29.21 | 2.270 | -. 958 | 2. 214 | -1.582 |
| - -39.37 | a. 025 | - 005 | - 030 | -006 |
| -49.53 | 0.000 | -. 002 | -.003 | -.003 |

## TABLE A15 (Continued)

| Profile | continued | at $X / L=$ | 275 |
| :---: | :---: | :---: | :---: |
| $Y$ MM. | wwutar | wwutar | uuwtar |
| 52.07 | -. 004 | 001 | -. 004 |
| 41.91 | $-024$ | -. 017 | $-.001$ |
| 31.75 | 1.956 | $-.694$ | . 737 |
| 24.15 | 5.259 | -1.785 | 4.886 |
| 20.32 | 5.305 | -. 446 | 7. 127 |
| 13.51 | 2. 491 | -. 898 | 7.050 |
| 15.70 | -1.851 | 1.469 | 4.444 |
| 10.16 | -3.131 | 3.119 | . 710 |
| 7.62 | -1.994 | . 391 | -2.272 |
| 5.08 | -3.984 | 1.529 | - 429 |
| 3.54 | $-\frac{3}{3} .053$ | 1. 405 | $\because 095$ |
| 0.00 | -3.449 | $-307$ | -283 |
| -2. 54 | -3.157 | -1 114 | $-638$ |
| -5.08 | -3. 346 | -1.386 | $\therefore .834$ |
| -7.62 | -2.758 | -.920 | 4.353 |
|  |  |  |  |
| $-13.97$ | $\frac{1}{3} .586$ | 2.436 | $-3 \cdot \frac{2}{2}$ |
| $-7.78$ | $6 \cdot 2$ | $5.805$ | $-1 .$ |
| $-2 i$ | $5.825$ | $5.400$ | $-\frac{2}{3} .6$ |
| - 59.21 | 1.807 | . 686 | -i. 644 |
| $-393$ | . 005 | -.002 | 012 |
| -47. 53 | 0.009 | $-.002$ | 004 |

Airfoil Wake Data at $x / L=.399$ (central plane)
RUN \# OF KUTE KI 3981


Uinfinjty (m/s) = 22.26
Integr. Parameters
Dfi:is*sut $x=9473079 E-02$
Delta*sut $z=1359371 E+00$
T.うta sub $x=.7767224 E-02$
-neta sut $z=.17513645-02$
The ta sub $x z=-1267331 E+00$
Shape Factor $H=1.220$

Wake Profile

| $Y \mathrm{~mm}$. | u/Uinf | Gama | wruinf | Eeta | Pstatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 64.77 | 1.008 | 25 | 004 | 30.25 | 0736 |
| 54.81 | 1.005 | 19 | 003 | 30.19 | -. 0718 |
| 44.45 | 1.004 | 14 | 002 | 30.14 | -. 0743 |
| 34.29 | 946 | $3{ }^{3}$ | 006 | $30.3{ }^{3}$ | -. 0801 |
| 24.13 | 925 | 1.03 | 017 | 31.03 | -. 0524 |
| 15.97 | 835 | 2.07 | 030 | 32.67 | -. 0218 |
| 6.35 | 773 | 2. 52 | $0{ }^{0} 4$ | 52. 52 | - 0013 |
| 1.27 | 769 | 2. 56 | $00_{04}$ | 弪2. 59 | 0040 |
| -1.27 | 761 | 2. 61 | 035 | $\frac{3}{32} \cdot 61$ | 0047 |
| -3.81 | 771 | 2.47 | 035 | 32.47 | 0025 |
| -6.35 | 786 | 2.35 | 032 | 32.32 | -.001\% |
| -11.43 | 825 | 1.97 | 028 | 51.97 | - 0127 |
| -16.51 | 873 | 1.54 | 023 | 31.54 | -. 0250 |
| -24.13 | 954 | 74 | ${ }^{12}$ | 30.74 | -. 0860 |
| -31.75 | 985 | 30 | 005 | 30.30 | -. 0547 |
| -41.91 | 977 | 08 | 001 | 30.08 | -.0618 |
| -52.07 | 998 | 01 | 000 | 30.01 | -. 0643 |
| -62. 23 | 1.000 | 0.00 | 0.000 | 30.00 | -. 3609 |

TABLE Al7
Airfoil Waike Data at $x / L=.66 i$ (central plane)


| $Y$ MM. | u/tinf | Gama | w'Uinf | Eeta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 63.25 | 1. 1008 | 7 | 013 | 30.72 | -. 0766 |
| 46.99 | 4. 006 | 4.70 | 012 | 30.70 | -. 0757 |
| 49.05 | 885 | 1.31 | 028 | 31.81 | -. 0387 |
| 11.43 | $0 \cdot 57$ | -20 | 033 | 32.29 | -. 0207 |
| 8.89 | 826 | E. 3. | 034 | 32.37 | -. 0161 |
| 7.62 | 95i | 2. 411 | 034 | 32.40 | -. 045 |
| ¢. 5.08 | 816 810 | 5.4. | $0{ }^{1} 34$ | 32.42 | -.0176 |
| 3.81 | 807 | 5.42 | 034 | 35.4? | -.0123 |
| 2.54 | 307 | 2.48 | 035 | 32.48 | -. 0117 |
| 1.27 | 806 | 2. 47 | 035 | 3.49 | -. 0092 |
| 0.00 | 803 | 2.47 | 035 | 32.47 | -.0104 |
| -1.27 | 805 | 2.41 | $\bigcirc 34$ | 32.41 | -. 0119 |
| - -4.05 | 8818 | 2. 40 | ? 34 | 32. $\frac{40}{77}$ | -. 0125 |
| -7. ${ }^{-3}$ | 8 | 2. 25 | 32 | 起: 25 | -. 0147 |
| \%919 | 539 |  | V39 | 32. 13 | - 0211 |
|  | 884 | 1. 61 | 029 | $\frac{31}{3} \cdot 915$ | - 0.028 |
| -5.15 | 946 | 3.16 | 019 | 31.16 | - 055 |
| -32.77 | 975 | . 74 | 013 | 30.74 | -.05\%2 |
| - 61.06 | 1.954 | 46 | 018 | 30.46 | -. 0413 |
| -51.02 | 1.003 | 69 | 234 | 30.60 | -. 0609 |

## TABLE AI7 (Continued)

## Turbulence Data

X/L Location $=$
Reference Velocit ${ }^{661}=21.8 \mathrm{~m} / \mathrm{s}$
Wake Profile

| $Y_{53}{ }^{\text {MM }} \mathbf{3} 4$ | utar | Utiar | Whar | uubar | yutiar | unbar |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 43.18 | 996 | 001 | . 000 | 1. $5 \frac{1}{3} 2$ | 1.783 | 1. 550 |
| 33.02 | 965 | -. 000 | 001 | 3.760 | 3.460 | 3.385 |
| 22.86 | 903 | -. 005 | 008 | 5.414 | 4.94 | 4.89 |
| 15.24 | 847 | -. 001 | 014 | 5.617 | 5.469 | 4.909 |
| 7.62 | 904 | 004 | 026 | 5.046 | 5.204 | 4.656 |
| 2. 54 | \%9 | 007 | 027 | 4.392 | 5.054 | 4.4 |
| -2. 54 | 793 | 008 | 038 | 4.607 | 5.206 | 4. 547 |
| -12.70 | 843 | 015 | 026 | 5.609 | 5.386 | 5.112 |
| -17.78 | 875 | 015 | 035 | 5.722 | 5.046 | 5.008 |
| -22.86 | 908 | 014 | 020 | 5.442 | 4.699 | 4.659 |
| -37.94 | 941 | 013 | 013 | 4.514 | 4.076 | 3.840 |
| -35.56 | 972 | 009 | 000 | 2. 552 | 2.835 | 2.308 |
| -43.18 | 981 | 008 | -. 003 | 1.137 | 1.683 | 1.203 |
| -50.80 | 982 | 005 | -.003 | 749 | 1.049 | 780 |
| Profile | contin | d at $X$ | $=.6$ |  |  |  |
| $Y \mathrm{~mm}$ | uubar. | untiar | untiar | uuubar | vuutar | wwwb a |
|  | -. 002 | C12 | . 05 | -. 008 |  | -. 070 |
| 43.18 | - 5.847 | -138 |  | - ${ }^{5}{ }^{\text {a }}$ | 1.36 | - -4.054 |
| 53.02 | -5.382 | 1. 373 | . 27 | 6.543 | 1.359 | -4.095 |
| 122.86 | -13.39 | -. 0150 | 3.98 | 6.881 | 1.480 | -1.893 |
| 7.62 | -7.482 | -1.279 | 1.03 | -4.980 | 1.646 | 007 |
| 2.54 | -1.641 | -1. 578 | 1.25 | -3.183 | 803 | 606 |
| -2. ${ }^{-7} 4$ | 5.110 | -1.7i6 | - 22 | -2.7 41 | 593 | - 1.288 |
| -12.70 | $1{ }^{1} 8.85$ | -4.007 | -1.95 | -4.41\% | - 097 | 1.957 |
| -17.78 | 13.984 | -4.990 | -1.36 | 3.543 | .334 | 2. 347 |
| -22.86 | 12.680 | -5. 607 | -1.40 | 10.473 | -3.069 | 4.595 |
| -27.94 | 7.794 | -4.646 | - 81 | 10.951 | -2.931 | 5.573 |
| -35.56 | 2.648 | -2.037 | -. 69 | 3.880 | -1.895 | 3.100 |
| -43. 18 | 546 | -. 340 | -. 10 | 199 | -. 288 | 836 |
| -50.80 | 150 | -. 028 | -. 05 | 114 | -. 014 | 204 |

## TABLE Al7 (Continued)

| Profile | ontinued | at X / | 661 |  |
| :---: | :---: | :---: | :---: | :---: |
| Y MM. | vuutiar | Uuwbar | vuubar | vowbar |
| 53.34 | -. 012 | -. 003 | 010 | -. 013 |
| 43.18 | -. 585 | -. 003 | 575 | -. 145 |
| 33.02 | -2.751 | . 037 | 3.647 | - 492 |
| 23. 86 | -2.923 | - 686 | 4.345 | -1. 524 |
| 15.24 | 3.355 | -. 5198 | -2.461 | -. 309 |
| $2 \cdot 54$ | 1. 131 | 341 | -2. -1.167 | -. 228 |
| -2. 54 | -2.075 | 517 | -4. 330 | $-.954$ |
| -7.62 | - 2.915 | 687 | -2. 521 | - -2.47 |
| -12.76 | こ. 287 | 494 152 | 2. 188 | -5.281 |
| -32.86 | 4.802 | - 928 | 4.124 | - 1.070 |
| -27.94 | 4636 | -2.095 | 4.746 | -1.639 |
| - 35.56 | 1.984 | -1.014 | 2.293 | -1. 158 |
| - 50.18 | . 285 | -. 01418 | 141 | -. 1019 |

Profile continued at $x / L=.661$

| $Y \mathrm{~mm}$ | wwutar | wwotar | uuntiar |
| :---: | :---: | :---: | :---: |
| 53.34 | -. 001 | -. 014 | 025 |
| 43.18 | . 185 | - 33 | 315 |
| 37.00 | 2.586 | 1.343 | . 956 |
| $\frac{22}{5} .84$ | - 2.046 | $-1.353$ | 5.694 |
| 17.6 | -1.781 | $1 \cdot 755$ | 530 |
| 2. 54 | -1. 591 | - 298 | 1.260 |
| -2. 54 | -1.619 | -. 790 | 1.612 |
| - -6.6 | -1.405 | -. 570 | 4.469 |
| $-17.78$ | 1. 8.96 | 1. 469 | - 1.721 |
| -25.96 | 3.989 | 3.350 | -2.953 |
| -27.94 | 3. 6.75 | 2.300 | -3.658 |
| -35.56 | 1.247 | c\%0 | -1.939 |
| $-43.18$ | 126 | 194 | -. 151 |
| -50.80 | 005 | 00 | -. 1146 |

TABLE A18

## Airfoil Wake Data at $x / L=.661$ ( $+6^{\prime \prime}$ plane)

| RUY ${ }^{\text {DATE OF }}$ R ${ }^{\text {RUN }}$ ( 4981 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{array}{r} \text { Urefernce }(\mathrm{m} / \mathrm{s}) \\ \mathrm{n} \text { loca }= \\ .661 \end{array}=21.84$ |  |  |  |  |  |
| Uinfinitu (m/s) $=22.50$ |  |  |  |  |  |
| Integral farameters |  |  |  |  |  |
| Delta*sut $x=9693721 E-02$Deltaxsub $z=9094906 E-7425145 E-02$Theta sub $x=.74254$ |  |  |  |  |  |
| Theta sub z $=.21108515-02$Theta sut $x z=.850125 E-01$Theta sub $x x=.245953 E-03$ |  |  |  |  |  |
| Shape factor $H=1.171$ |  |  |  |  |  |
| Wake Profile |  |  |  |  |  |
| $Y \mathrm{MM}$ | u/Uinf | Gama | w/Uinf | Eeta | Fstatic |
| 48.26. | $1.04{ }^{1}$ | $\underline{9}$ | 015 | 30.86 | 0495 |
| 27.94 | $1.95 \%$ | 1.25 | 021 | 31.24 | 0381 |
| 20.32 | 915 | 1.5 | 024 | 31.50 | 0334 |
| 12.70 | 861 | 1.96 | 0 | 31.86 | 0350 |
| 7.62 | 833 | 2.07 | 031 | 32.07 | 0359 |
| 5.88 | 818 |  | 036 | $32 \cdot 29$ | 0376 |
| 2. 54 | 807 | 2. ${ }^{1}$ | 033 | 32.31 | 0396 |
| 1.27 | 800 | 2.36 | 035 |  | 0445 |
| 0.00 | 796 | 2. 37 | 033 | 32.36 | 0438 |
| - $-2 \cdot 54$ | 814 | 2. 3.9 | ${ }_{0}^{03}$ | 32.19 | 0154 |
| -4. 5.7 | 816 | 2. 19 | 031 | $3{ }^{3}$ 2 19 | 0140 |
| -7. 916 | 825 | 2.29 | 030 | 32. ${ }^{3}$ 2 | 0198 |
| -14.73 | 854 | 1.86 | 028 | 31.66 | 0153 |
| -21.08 | 888 | 1.57 | 024 | 31.56 | 0192 |
| - 37.5 | 1. 000 | 44 | 175 | 30.40 | 0037 |

TABLE A19
Airfoil Wake Data at $x / L=.661$ ( $-6^{\prime \prime}$ plane)


TABLE A2O
Airfoil Wake Data at $x / L=.992$ (central plane)


## TASLE A21

## Airfoil Wake Data at $x / L=1.928$ (central plane)

## RUN \# 18 <br> DATE OF RUN 4981

```
Urefernce. (m/5) \(=21.84\)
Uinfinitり (m/5) = 33.34
```

Integral Parametere
Delta*sub $x=3385820 E-02$
De1ta*sub $z=.2070857 E+00$
Theta sub $x=2963738 \mathrm{E}-02$
Theta sub $z=.3562933 E-03$
Theta sut $\times 2=.1039173 \mathrm{E}+00$
Theta sut $z x=.2139227 E-04$
Stape Facior $H=1.075$

Wake Profi?e

| $Y \mathrm{Mm}$ | u/Uinf | Gama | w/ binf | Eeta | Petatic |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.00 | 899 | 77 | 012 | 30.77 | 0116 |
| 3.81 | 901 | 51 | 008 | 30.51 | 0103 |
| 7.6 | $90 \%$ | 50 | 008 | 30.50 | 0084 |
| 12.70 | 9 | 47 | 0084 | 30.47 | 0001 |
| 26.67 | 947 | 55 | 0.19 | 30.55 | 0029 |
| 25.56 | 968 | 20 | 003 | 30.15 | $-.0034$ |
| 45.72 | 981 | 09 | 002 | 30.09 | 0039 |
| 57.15 | 994 | 13 | 002 | 30.93 | 0037 |
| 69.8 | 999 | 05 | 001 | 30.05 | 0035 |
| O. 06 | 1.000 | 00 | 001 | 30.00 | 0072 |

TABLE A21 (Continued)

Turbulence Data

```
X/L Locatign = {
```

Half Wake Profile

| $Y$ MM | Utär | Vtar | Whar | uutar | vツbap | Wubam |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -6. 35 | . 902 | . 004 | . 015 | 3. 378 | 4.385 | ㄷ.59 |
| -5.08 | 900 | 003 | 016 | 3.376 | 4.485 | 3.656 |
| -2.54 | 905 | 005 | 017 | 3.358 | 4.351 | 7. 674 |
| 0.00 | 900 | 003 | 014 | 3.439 | 4.545 | 3.637 |
| 2.54 | 903 | 004 | 016 | 3.482 | 4.308 | $\because .654$ |
| 5.08 | 904 | 003 | 016 | 3.454 | 4.293 | 5.726 |
| 7.62 | 908 | 004 | 015 | 3.584 | 4.315 | 3.755 |
| 10.16 | 911 | 000 | 014 | 3.650 | 4.127 | 3.841 |
| 12.70 | 914 | 002 | 012 | 3.537 | 4.196 | 3.714 |
| 17.78 | 923 | 005 | 011 | 3.669 | 3.973 | 3.710 |
| 22.86 | 934 | 001 | 010 | 3.735 | 3.802 | 3. 635 |
| 27.94 | 946 | 005 | 010 | 3.504 | 3.730 | 3.553 |
| 33.02 | 955 | 002 | 009 | 3.379 | 3.580 | 3.414 |
| 38.10 | 965 | 001 | . 000 | 3.139 | 3.257 | ㅍ.006 |
| 43.15 | 975 | 001 | 007 | 2.756 | 2.988 | 2.853 |
| 5 (1. 8 (1) | 987 | $-000$ | 004 | 2.310 | 2.453 | \%.326 |
| 58.42 | 992 | 001 | 006 | 1.766 | 1.967 | 1. 909 |
| 66.04 | 998 | 002 | 005 | 1.290 | 1.691 | 1. 473 |
| 76.20 | 1.00\% | 001 | 005 | 883 | 1.162 | 1. 044 |
| 86.36 | 1.000 | $-.000$ | $-000$ | 62 c | . 930 | . 640 |

## TABLE A21 (Coniinued)

Profile continued at $X / L=1.928$

| $Y \mathrm{MN}$ | untiar | UWtar | UWtat | vuutar | quubar | wwwter |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $-6.35$ | 801 | $\text { - } 184$ | $-69$ |  | $906$ | $\begin{array}{r} 389 \\ 159 \end{array}$ |
| $-5.06$ | $961$ | $-979$ | $-70$ | $\text { - } 331$ | $.590$ | $345$ |
| $-2.54$ | -. 018 | - 425 | - 60 | -. ${ }^{2} 7$ | 656 | - 455 |
| 0.00 | - -740 | -. 149 | - 39 | -. 167 | 175 | $-1.085$ |
| 2.54 | -2.947 | 342 | $\cdots .08$ | . 065 | 849 | -1.309 |
| 7.6 | $-3.60{ }^{-3}$ | 479 | .37 | 249 | $\frac{2}{9} \frac{1}{7}$ | 7 |
| 10.16 | $-3.806$ | 700 | 46 | 114 |  |  |
| 12.70 | -4.57\% | . 828 | 66 | 081. | 756 | -1 74 |
| 17.78 | -4.812 | 1. 428 | 40 | 1.66? |  | - 2480 |
| 22.86 | -5.369 | 967 | $\frac{15}{5}$ | 1.808 | 1.029 | - 5.189 |
| 27.94 | -5.393 | 1.383 | 9 | $\frac{2}{2} .484$ | 1.682 | -2.807 |
| 33.02 | -5.111 | 1.249 | 82 | 2.360 | 482 | -5.46 |
| 38.10 | -4.299 | 1. 004 | 62 | 2.803 | 447 | $\bigcirc 2045$ |
| 43.16 | -3.451 | $7 \mathrm{7a}$ | 60 | $1.90 \%$ | 94 | -1. 69 |
| 50.80 | -2.359 | 769 | 14 | 95 | 68 | -1.067 |
| 58.42 | -1.412 | $5 \% 7$ | 5 | - 0 | 方i | -1.097 |
| 60.04 | -. 608 | - 150 | 1. ${ }_{4}$ | $05 \frac{1}{5}$ | -00\% | . 225 |
| 76.20 | -. 188 | -. 029 | 0 | - $00 \frac{3}{4}$ | -. 016 | $-.053$ |
| 86.36 | -. 041 | -.012 | 6 | -.003 | 010 | . |

## TABLE A21 (Continued)

| Y mm | uuubar | uuwtiar. | quubar | vuwtar. |
| :---: | :---: | :---: | :---: | :---: |
| -6.35 | -. 160 | 6 | -1.831. | - 402 |
| -5.08 | - ${ }^{3}$ | .004 | -1. 554 | -. 303 |
| -2.54 | 306 | -. 043 | -1.494 | -. 401 |
| 0.00 | 543 | -. 016 | - 5.635 | -. 151 |
| 2.54 | 455 | -. 264 | - -1.694 | -. 210 |
| 7.6 2 | 492 | - 550 | -1.066 | -. 325 |
| 10.16 | 326 | -. 307 | $-.584$ | 035 |
| 42.70 | 383 | -. 005 | -. 404 | - 077 |
| 17.78 | . 240 | -. 068 | 038 | -. 0.38 |
| 25.88 | -. 325 | -.063 | 1. 236 | -. 187 |
| 3'02 | -1.149 | 169 | 1. 508 | - 5 |
| 38.10 | -1.401 | 056 | 1. 461 | - 408 |
| 43. 180 | -1. 0174 | 288 | 1.309 | -. 088 |
| 58.42 | -. 577 | 194 | 1.657 | -. $18 \%$ |
| 66.14 | -. 217 | 104 | 172 | - 480 |
| 76.20 | -. 888 | 009 | 068 | -.026 |
| 86.36 | -. 012 | 004 | 005 | © 4 |

## TABLE A21 (Continued)

| Y MM | wwubar | wwubar | uuwbar |
| :---: | :---: | :---: | :---: |
| -6. 35 | -. 036 | -. 104 | 1.675 |
| $-5.09$ | . 063 | -. 089 | 1.560 |
| -2. 54 | -. 228 | . 103 | 1.470 |
| 0.00 | . 061 | $\underline{137}$ | 2. 025 |
| 5.08 | -. 2345 | - . 2 2를 | 2. 0851 |
| 7.62 | . 114 | - 2 ¢ 3 | 1.335 |
| 10.18 | 139 | 085 | 1.378 |
| 17.78 | 429 | 054 | 868 |
| 29.86 | 650 | - 511 | 456 |
| 27.94 | 1.009 | -. 610 | 244 |
| 33.02 | . 828 | -. 686 | 718 |
| 38.10 | 1.020 | -1.089 | 765 |
| 43.18 | 896 | -. 607 | 711 |
| 50.30 | 893 | - 550 | 253 |
| 66.04 | 403 | - 473 | - 060 |
| 76.20 | -. 019 | -. 089 | 017 |
| 86.36 | -. 002 | -. 015 | 020 |


[^0]:    Arusinulguninginininutoronarue
    vet
    762
    034
    686
    018
    321
    103
    069
    102
    905
    861
    515
    614
    804
    896
    876
    989
    895
    961
    104
    
    Whbar

