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(NASA-CR-167873) ATMOSPHEREIC RADIATION
                                    N83-21677
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GODEL FOR WATEh SURFACES Final Technical
Report, 4 Mar. 1981 - 15 Jui. 1982 (science
Applications. IHC.) $67 \mathrm{pHCAO4/MFAOI}$ Unclas


# ATMOSPHERIC RADIATION MODEL FOR WATER SURFACES 

## FINAL TECHiviICAL REPORT

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\text { MAY } 1982
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Prepared For
NATIONAL AERONAUTICS AND SPACE ADMINISTRATION
Lewis Research Center Cleveland, Ohio 44135 CONTRACT NO. NAS3-22495

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| 1. REPORT NUMGEA <br> NASA CR-167873 2. COVT ACCESSION NO | 2. MECIPIEMT'S CATALOC MUMEER |
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| 7. AUTMOR(a) <br> Robert E. Turner <br> Daniel W. Gaskill <br> James R. Lierzer | a. Comtract on ghawt mumoeng NAS3-22495 |
| 9. PENFORMIMG ORGANIzATIOM NAME AND adDRESS <br> Science Applications, Incorporated 15 Research Drive Ann Arbor, Michigan 48103 | 10. PROGRAM ELEMEMT. PNOSECT, TASK AnEA A WONK UNIT WUMEES |
| 11. COMTMOLLIMG OFFICE MAME AND ADONESS <br> NASA - Lewis Research Center <br> 21000 Brookpark Road | $\qquad$ |
| 21000 Brookpark Road <br> Cleveland, Ohio 44135 | $\begin{aligned} & \text { 13. Mumeer or Pages } \\ & \text { ix }+56 \\ & \hline \end{aligned}$ |
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additional components such as the atmospheric path radiance which results from singly-scattered sky radiation specularly reflected by the water surface. It also considers a component which is referred to as the virtual sun path radiance, i.e. the singly-scattered path radiance which results from the solar radiation which is specularly reflected by the water surface.

These atmospheric radiation components are coded into a computer program for the analysis of multispectral remote sensor data over the Great Lakes of the United States. The user must know certain parameters, such as the visibility or spectral optical thickness of the atmosphere and the geometry of the sensor with respect to the sun and the target elements under investigation.

Suggestions and recommendations are given for further investigation of the problem of the remote sensing of water surfaces. If all of these extrinsic radiation components are properly accounted for, then the intrinsic water radiance can be found by applying the algorithm or an adaptation of the algorithm in this report. As a result, one would then be able to know the actual surface water spectral radiation field independent of the atmosphere.

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## PREFACE

This report describes part of the work done on a research program in the remote sensing of the Great Lakes using a multispectral scanner aboard an aircraft. The research has been conducted for the NASA-Lewis Research Center, Cleveland, Ohio, by Science Applications, Inc., at the Dayton, Ohio, office. The primary objective of this program is to develop remote sensing as a practical technique for the analysis of the Great Lakes.

Remote sensing of the environment involves the transfer of radiation from the Earth's surfa' 2 through the atmosphere to a sensor which is located at some point within the atmosphere. For water surfaces with their inherently low reflectances, the atmospheric scattering of solar radiation acts as a significant noise factor. In this report we have extended an existing model to include various atrospheric radiation components so that the resulting mathematical algorithm will allow one to extract a radiance value which is more nearly representative of the actual radiance of the water, independent of atmospheric effects.

This research was performed under contract NRS3-22495 and covers the period from 4 March 1981 through 15 July 1982. Mr. Thom Coney served as Technical Monitor of the contrace and Dr. Robert E. Turner of Science Applications, Inc., was the program manager and principal investigator.

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## SYMBOLS

| $E_{0}$ | solar irradiance at the top of the atmosphere |
| :---: | :---: |
| $\mathrm{E}_{0}{ }^{\text {- }}$ | solar irradiance of the virtual sun |
| $L_{0}(i)$ | radiance incident at angle i |
| $L_{s}(\mathrm{i})$ | surface radiance at angle i |
| $L_{\text {SKYRE }}$ | reflected sky radiance |
| $L_{\text {SOLRE }}$ | reflected solar radiance |
| $L_{\text {SKY }}$ | sky radiance |
| $L_{\text {SOL }}$ | direct beam radiance |
| $L_{\text {PMS }}$ | multiply scattered path radiance |
| $L_{\text {PSS }}$ | singly scattered path radiance |
| $L_{U}$ | upward scattered radiance beneath water surface |
| $L_{\text {INIR }}$ | radiance scattered from beneath water surface |
| $L_{\text {SKYSS }}$ | singly scattered sky radiance |
| $L_{\text {PTOT }}$ | total radiance |
| LVSOL | radiance from virtual sun |
| $L_{\text {VP }}$ | virtual sun path radiance |
| n | index of refraction of water relative to air |
| p | slore probability |
| $p(x)$ | scattering phase function |
| T | transmittance |
|  | (GREEK SYMBOLS) |
| ${ }^{\text {i }}$ i | incident angle |
| ${ }^{\theta} \mathrm{r}$ | reflected angle |
| ${ }^{\circ} \mathrm{s}$ | Fresnel reflectance |


| $\rho_{f}(\mathrm{i})$ | Fresnel reflectance at angle i |
| :---: | :---: |
| $\theta$ | nadir view angle |
| ${ }_{0}$ | solar zenith angle |
| $\phi$ | azimuth view angle |
| $\phi_{0}$ | solar azimuth angle |
| ${ }^{\mu}$ | cosine of nadir view angle |
| $\mu_{0}$ | cosine of solar zenith angle |
| $\chi_{\text {SKY }}$ | scatterirg angle |
| $\tau$ | optical depth |
| ${ }^{\tau} 0$ | optical thickness |
| $n$ | forward scattering parameter |
| $\rho$ | surface albedo |
| $\omega_{0}$ | single scattering albedo |

In the analysis of remotely sensed data on bodies of water, the atmosphere obscures the inherent surface features as a result of the scattering and absorption of solar radiation. In the case of multispectral data acquired by aircraft or spacecraft sensors, one can preprocess the data by applying mathematical models and algorithms to the digitized data. The mathematical model developed in this investigation is specifically designed to account for various components of the visible and infrared radiation in the atmosphere which interfere with the inherent signal from a water surface. If the atmospheric parameters are known, then when the algorithm is applied to the multispectral data sets, an improved or corrected data set will result.

This improved atmospheric correction model allows for the path radiance in the atmosphere as a result of singly-scattered solar radiation and also singly-scattered solar-reflected radiation. In addition, the model includes a singly-scattered sky radjation component for the radiation which is reflected by the water surface. Comparisons are made among the relative magnitudes of these radiation components in terms of the geomet:ic and environmertal factors. Recommendations are presented for a more advanced model whicin would include the corresponding radiation components for multiple scattering.

## INTRODUCTION

Multispectral scanner data obtained by sensors aboard aircraft and spacecraft allow a Lser to examine the detailed physical properties of a surface. 'hese properties are of interest to many investjgations in various disciplines such as land use studies, agriculture, hycrology, forestry, and oceanography. In all of these investigations, however, the scattering of visible and infrared radiation by the atmospheric constituents will reduce the inherent surface radiance and add a palh radiance to the attenuated radiance from the target. For many cases of the remote sensing of bright land areas on relatively clear days the attenuated radiation from the surface is rather large as compared to the atmospheric path radiance. For water bodies, with inherently low reflectances, this is no longer true and the path radiance can be a major effect in the total radiance at the sensor.

The purpose of this investigation was to extend an existing atmospheric radiative transfer model to include other radiation components which did not exist in the previous model. These additional atmospheric radiation components include specific effects for the remote sensing of water surfaces. The model is used in conjunction with an algorithm specifically designed for the analysis of multispectral data.

The determination of the atmospheric radjation components is important for the analysis of the probability of misclassification of various classes of surface materials. To first order one may consider the so called linear transfer problem in which the path radiance is constant over varying surface reflectances for a horizontally spatially uniform haze. For a non-uniform haze, however, the path radiance can vary, thereby resulting in a higher
probability of misclassification of objects if the degree of nonunifirmity is unknown. A second-order effect, but one which can become quite important for the remote sensing of high-contrast targets is the adjacency effect. This is when radiation from a bright target causes an increase in the path radiance with respect to the radiance from a neighboring dark target. This problem would exist, for exampla, in the remote sensing of water bodies near bright sandy beacheg. The results should be evident in the brightened image of the wat:er near the shoreline, provided the effects of waves and whitecaps are eliminated. This second-order adjacency effect is not included in the model or algorithm in this investigation but the effect can be accounted for if the investigator has sufficiently detailed atmospheric data on the horizontal stratification of aerosols.

Multiple scattering is particularly important if the sky is hazy. These effects are considered in the model for path radiance which results from the sun as a source. We have not included the multifle scattering effects for solar-reflected radiation.

## OPTICAL PROPERTIES OF WATER

The interpretation of remote sensing data collected over water surfaces requires a detailed knowledge of the optical properties of water and the air-water interface. Water is unusual as a natural surface because it is a specular reflector and because in the visible regions sensible data can be obtained from well below the water surface [1]. Also of importance is the phenomonon of refraction, which occurs when radiation passes through the air-water boundary. These properties of water and their significance in terms of remote sensing are discussed in detail in this chapter.
3.1 REFLECTION, REFLECTANCE AND REFRACTION

Most natural surfaces are approximately Lambertian--reflecting incident radiance equally in all directions. A smooth water surface, however, is a specular reflector, and reflection of radiation from it follows the geometrical law of reflection. This geometric law requires that the angle with respect to the normal to the surface of the reflected ray equal the angle of incidence of the incident ray and that the reflected ray be in the same plane as the incident ray. Specular reflection is depicted in Figure 1. Reflectance, $\rho_{s}$, of the water surface is given by the Fresnel equation

$$
\begin{equation*}
\rho_{s}=\frac{1}{2}\left[\frac{\sin ^{2}\left(\theta_{i}-\theta_{r}\right)}{\sin ^{2}\left(\theta_{i}+\theta_{r}\right)}+\frac{\tan ^{2}\left(\theta_{i}-\theta_{r}\right)}{\tan ^{2}\left(\theta_{i}+\theta_{r}\right)}\right] \tag{1}
\end{equation*}
$$

where $\theta_{i}$ is the incident angle and $\theta_{r}$ is the angle of refraction.
The transmitted part of the incident ray experiences refraction at the water surface, as shown in Figure 2. Snell's law, given by

$$
\begin{equation*}
\frac{\sin \left(\theta_{i}\right)}{\sin \left(\theta_{r}\right)}=n, \tag{2}
\end{equation*}
$$

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FIGURE 1. SPECULAR REFLECTION FROM A SMOOTH WATER SURFACE. THE INCIDENT AND REFLECTED RAYS ARE IN THE SAME PLANE.

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FIGURE 2. RE: RACTION OF THE TRANSMITTED BEAM AT THE WATER SURFACE AND THE LAW OF REFRACTION $\mathrm{n}=\mathrm{REFRACTIVE}$ INDEX OF WATER RELATIVE TO AIR.
describes the relationship between the angle of incidence of the incoming beam, $\theta_{i}$, the angle with respect to the normal to the surface of the refracted beam, $\theta_{r}$, and the refractive index of air relative to water. The refracted beam lies in the same plane as the incident beam. The phenomenon of refraction is depicted in Figure 2. The law of refraction requires that for a water surface, downward sky radiation and direct sunlight enter the water within $48.5^{\circ}$ of the vertical. Only when the water surface is roughened by wind or another disturbance can direct sunlight or sky radiation penetrate the water surface outside this range of angles. Back-scattered radiation from beneath the water-air inte::face also experiences refraction on reaching the water surface. When the water surface is calm, upward radiation incident at angles greater than $48.5^{\circ}$ with tre vertical is totally internally reflected [2]. Thus, downward radiation beneath the water surface at angles with the vertical greater than $48.5^{\circ}$ is upward radiation in the water which has been totally internally reflected [3]. Equations 1 and 2 show that the reflectance and transmittance of the water surface are dependent on the refractive index of water. The refractive index is influenced by changes in temperature and by the concentration of various solutes in the water. Figure 3 shows how the reflectance function, equation 1 , varies as a function of the angle of incidence for refractive indices of $1.20,1.33,1.40,1.45$. This range of refractive indices encompasses the range of natural variability in the refractive index for water; and Figure 3 shows that, over this range, variation in the refractive index is of little importance in determining the surface reflectance, For all of the calculations shown in this repurt, a refractive index of $4 / 3$ is used. Figure 4 shows transmittance, $T$, and reflectance, $\rho_{s}$, as a function of angle of incidence of the incoming radiation for a refractive index of 1.33 . Since many applications of remote sensing over water require

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FIGURE 3. FRESNEL REFLECTANCE FOR UNPOLARIZED LIGHT AS A function of angle of incidence of incoming beam.

examining the signature from beneath the water surface, Fresnel reflectance is of paramount importance. Figure 4 shows that transmittance is highest when the incident beam is perpendicular to the water surface, while the reflectance is lowest at this angle. Thus, from a consideration of the Fresnel formulas, one would expect the return from beneath the water surface relative to the total, to be greatest when viewing that surface at the nadir. Contributing to this effort is the fact that most scattering phase functions for polydispersions in water have a secondary peak at $180^{\circ}$ [2]. Figure 4 shows that while the Fresnel reflectance function is at a minimum at the nadir view angle, this function changes very little out to viewing angles as great as $40^{\circ}$, at which point it begins to rise steeply. Internal reflection at the water surface of upwelling radiation is also at a minimum normal to the water surface. Thus, for remote sensing work where the return from beneath the water surface is of greatest interest, scan angles should be maintained within $40^{\circ}$ to $45^{\circ}$ of the vertical. Beyond an angle of $48.5^{\circ}$ no radiation from beneath the water surface will reach the sensor when the water surface is calm. Scanning in the solar plane is also problematic in this regard since specularly reflected light on the solar side of the scan plane would saturate the sensor. For some purposes, such as viewing the glitter pattern on the water surface, scanning in the solar plane may be desirable.

### 3.2 ABSORPTION AND SCATTERING

Water is a good absorber of electromagnetic radiation. Only in the relatively narrow spectral region from about 400 to 600 nanometers is the transparency of water such that radiation can penetrate more than a few meters in depth below the water surface. Both at wavelengths shorter than 400 nm and longer than 600 nm , absorption increases rapidly and only very small amounts of radiation are scattered back out of the water
into the atmosphere. At the very short wavelengths this radiation is further strongly attenuated in the atmosphere.

Scattering of radiation in water is caused by water molecules, by dissolved salts and by particles in suspension. These effects are usually assumed to be additive [4].

Scattering by water molecules is described by fluctuation theory which predicts scattering of radiation as a result of molecular movements which cause fluctuations in the density of the medium. As in Rayleigh scattering, this type of scattering is proportional to $\lambda^{-4}$, where $\lambda$ is the wavelength of the radiation being scattered. The effect of dissolved salts on the molecular scattering phase function is usually small enough to be neglected.

In general, most scattering in water is accomplished by particles in suspension [4]. Particulate matter in water derives from runoff from land, deposition from the atmosphere, and organic processes within the water. Thus, particles may be quite irregular in shape and particle size distributions are difficult to sharacterize precisely [2]. Because some sources of particles such as runoff of organic processes may be highly localized in space, size distributions may vary greatly in space and time. Although particle shapes vary considerably from the spherical ideal of scattering theory, it has been shown $[5,6]$ that systems of irregularly shaped particles can be adequately approximated by systems of polydisperse systems of spherical particles. The major observed features of phase functions of particulate suspensions in water are a strong forward scattering peak, a broad minimum around $100^{\circ}-130^{\circ}$ and a small secondary peak in the back scattering direction [4].
3.3 OPTICAL PROPERTIES OF A WIND ROUGHENDED WATER SURFACE Roughness of the water surface caused by wind presents an additional problem in the calculation of the optical properties of the surface. Waves increase the angle of incidence of direct
radiation for high solar elevations. The effect on the Fresnel reflectance, however, is of little consequence since the reflectance does not vary much with solar zenith angle for zenith angles less than $40^{\circ}$ (see Figure 3). Waves reduce the angle of incidence of direct radiation from a low sun, greatly reducing the reflectance of the water surface. Cox and Munk [7] have shown that wave action becomes a significant factor for solar elevations below $20^{\circ}$. At these low sun elevations, reduced reflection, shadowing and multiple reflections greatly reduce the reflected radiance.

The reflection of diffuse radiation by the water surface is little affected by surface roughness [2] , although complete agreement on this matter is lacking [8]. Burt [9] found that the albedo of a wind roughened water surface was slightly less than the albedo of a smooth water surface--a decrease from 6.68 to 5.78 for the roughened surface. Cox and Munk [7] measured a small increase in the albedo of a smooth water surface of 58 to $5.5 \%$ for a water surface roughened by waves. Kondratyev [8] on the other hand, calculates that where the solar zenith angle is $0^{0}$ the albedo of calm water surface of 2.18 will increase to 13.18. When the solar zenith angle is $30^{\circ}$, the increase will be from $2.2 \%$ to $3.8 \%$, and for a solar zenith angle of $60^{\circ}$ there will be a decrease from $6.2 \%$ down to 2.48 for a roughened surface. Plass et al. [10], using a Monte Carlo model of the atmosphere ocean system, demonstrate that the downward flux just below the surface always increases with wind speed, even at high sun elevations. They attribute this result to the fact that more sky radiance near the horizon enters the water when waves are present.

The effect of waves on the radiance of the water surface can be calculated if the probability distribution of surface slopes is known. For an observer looking down on a water surface, the specular angle will vary from place to place over the surface of the water. Since in most remote sensing applications
the light source (Sun) and observer (sensor) are high enough above the surface and the region viewed sufficiently small that variation in the specular angle can be neglected. The radiance of the surface is then directly proportional to the probability of finding a surface element with slope, $S_{0}$, at the specular angle [11]. If $p$ is this probability, the radiance of the surface, $L_{s}$, at vertical angle $i$ is given by

$$
\begin{equation*}
L_{s}(i)=L_{0}(i) \rho_{f}(i) p \tag{3}
\end{equation*}
$$

where
$L_{0}(i)$ is radiance incident at the surface at vertical angle $i$, and
$\rho_{f}(i)$ is the Fresnel reflectance at vertical angle i.
Duntley [12] and Cox and Munk [7,13] have studied the statistical distribution of wave slope as a function of wind speed. Observations of the effect of wind speed on spatially or temporally averaged reflectance of the water surface indicate that it is not significant for view angles less than $70^{\circ}$ from vertical. Angles in excess of $50^{\circ}$ from the vertical are seldom used in remote sensing systems because of the large optical air mass at these angles.

COMPONENTS OF REFLECTED AND PATH RADIANCE IN REMOTE SENSING OVER WATER

In many applications of oceanographic remote sensing the quantity of greatest interest is the radiance information transmitted from below the water surface to a sensor, sometimes called the intrinsic radiance. To determine this quantity from raw remote sensing data we must not only estimate atmospheric path radiance but also the magnitude of radiance reflected off the water surface and transmitted to the sensor. In this chapter we describe in detail analytical modela appropriate for estimating the following quantities: reflected sky radiance, LSKYRE; singly scattered reflected solar path radiance, LVP; singly scattered path radiance, LPSS; and multiply scattered path radiance, LPMS. The first two of these quantities are radiances resulting from specular reflection off the water surface, the latter two are atmospheric path radiances. Each of these radiances augments the radiance detected by a sensor, masking the radiance signal from beneath the water surface, as shown in Figure 5.

### 4.1 REFLECTED SKY RADIANCE, LSKYRE

The geometry for sky radiance reflected into the line of sight of the sensor is depicted in Figure 6 , where $\theta$ represents the nadir view angle of the sensor, $\phi$ the azimuth of the sensor scanning plane. We assume a plane parallel uniform atmosphere. Sky radiance downwelling in the scan plane and incident at an angle $\theta$ with the normal to the surface is reflected in the direction of the sensor, and attenuated by the atmosphere as it travels to the sensor. The surface reflectance is given by the Fresnel formula for unpolarized light described in Chapter 3.

We consider in this model only singly scattered sky radiance, generated by scattering of solar beam radiation along the straight



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line path from the top of the acnosphere to the water surface. In Figure 6, this straight lins path has direction $(-\mu, \phi)$, where $\mu=\cos \theta$. In order to define the angle of scattering we muat first define 'wo vectors, one defininq the direction of a photon leaving the Sun and the other the direction of the singly scattered sky radiation. If $\left(-\mu_{0}, \phi_{0}\right)$ is the direction of the photon leaving the Sun (where $\mu_{0}$ is the cosine of the solar zenith angle, $\theta_{0}$, and $\phi_{0}$ is the photon azimuth), the vector direction of the photon leaving the Sun is

$$
\hat{L}_{S O L}=\left[\begin{array}{l}
\sin \left(\pi-\theta_{0}\right) \cos \phi_{0}  \tag{4}\\
\sin \left(\pi-\theta_{0}\right) \sin \phi_{0} \\
\cos \left(\pi-\theta_{0}\right)
\end{array}\right]=\left[\begin{array}{ll}
\sqrt{1-\mu_{0}^{2}} & \cos \phi_{0} \\
\sqrt{1-\mu_{0}^{2}} & \sin \phi_{0} \\
-\mu_{0}
\end{array}\right]
$$

and

$$
\dot{L}_{S K Y}=\left[\begin{array}{l}
\sin (\pi-\theta) \cos \phi  \tag{5}\\
\sin (\pi-\theta) \sin \phi \\
\cos (\pi-\theta)
\end{array}\right]=\left[\begin{array}{cc}
\sqrt{1-\mu^{2}} & \cos \phi \\
\sqrt{1-\mu^{2}} & \sin \phi \\
-\mu
\end{array}\right]
$$

The cosine of the scattering angle, $\lambda$ SKY, is given by the dot product I,SOL - LSKY, i.e.

$$
\begin{equation*}
\cos X_{\text {SKX }}=\mu \mu_{0}+\sqrt{1-\mu^{2}} \sqrt{1-\mu_{0}^{2}} \cos \left(\phi-\phi_{0}\right) . \tag{6}
\end{equation*}
$$

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The equation for singly-scattered aky radiance, $L_{\text {SKYSS }}$, is the well-known formula [14]

$$
\begin{equation*}
L_{S K Y S S}=\frac{\omega_{0} \mu_{0} E_{0} P\left(X_{S K Y}\right)}{4 \pi\left(\mu_{0}-\mu\right)}\left[e^{-\tau 0 / \mu_{0}}-e^{-\tau} 0 / \mu\right] \tag{7}
\end{equation*}
$$

$$
\text { where } \begin{aligned}
E_{0} & =\text { solar irradiance at the top of the atmosphere; } \\
\omega_{0} & =\text { atmospheric single scattering albedo; } \\
{ }_{0} \quad & =\text { optical thickness of atmosphere: } \\
P\left(\chi_{S K Y}\right)= & \text { scattering phase function for scattering } \\
& \text { angle } X_{\text {SKY }} .
\end{aligned}
$$

When the sky radiance is reflected off the surface of the water it is diminished by the fresnel reflectance of the water surface, $\rho_{F}$, and further attenuated by the atmosphere on its way to the sensor. Thus, the complete formula for the iflected sky radiance is

$$
\begin{equation*}
L_{\text {SKYRE }}=\rho_{F} e^{-\left(\tau_{0}-\tau\right) / \mu} L_{\text {SKYSS }} \tag{8}
\end{equation*}
$$

where $\mu$ is the cosine of the scan angle, $t$ is the optical depth of the sensor, and $e^{-\left(\tau_{0}-T\right) / H}$ is the transmittance of the atmosphere between the water surface and the sensor. At this point we take note of the fact that in the above formula for the singly scattered reflected sky radiance, the reflectance of the water surface and the transmittance of the atmosphere are opposing effects. Assuming the refractive index of water to be 1.33 , the

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Fresnel reflectance of water at the nadir vier: angle reaches a minimum of 0.021 and attains a maximum of $i .0$ at the grazing angle. Typical values of $\rho_{F}$ for angles commonly used in remote sensing range between 2.1 at the nadir to 3.0 at a scan angle of $46^{\circ}$. The transmittance, on the other hand, reaches a maximum at the nadir and becomes increasingly small as the scan angle increases. These effects will be discussed further in the following chapter.

### 4.2 SINGLY SCATTERED REFLECTED SOLAR RADIANCE, $L_{V P}$ •

The phenomenon of specular reflection produces an image of the radiation source on the surface of the water. We refer to the image of the Sun on the water surface as the virtual sun. If the scan plane is coincident with the solar plane and the sensor is scanning on the solar side of the scan plane at a view angle equal to the solar zenith angle, the field of view becomes saturated with the radiance of the Sun's image. Radiance from the virtual Sun is also scattered into the line of sight of the sensor. In the terminology of this report, we refer to singly-scattered path radiance from the Sun as virtual Sun path radiance, $\mathrm{I}_{\mathrm{VP}}$.

To find the scattering angle for the computation of singlyscattered virtual Sun path radiance, we note that the zenith angle of a photon leaving the virtual Sun is $\theta_{\text {SUN }}$ and the azimuth angle is $\phi_{0}=\phi_{S U N}+\pi$ (see Figure 7). The vector direction of a photon leaving the virtual Sun is

$$
\hat{\mathrm{L}}_{\text {VSOL }}=\left[\begin{array}{ll}
\sin \theta_{\text {SUN }} & \cos \phi_{0}  \tag{9}\\
\sin \theta_{\text {SUN }} & \sin \phi_{0} \\
\cos \theta_{\text {SUN }}
\end{array}\right]=\left[\begin{array}{cc}
\sqrt{1-\mu_{0}^{2}} & \cos \phi_{0} \\
\sqrt{1-\mu_{0}^{2}} & \sin \phi_{0} \\
\mu_{0} &
\end{array}\right]
$$

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FIGURE 7. GEOMETRY FOR VIRTUAL SUN PATH RADIANCE
CALCULATION.
and the vector direction into which photons from the virtual Sun are scattered, creating virtual Sun path radiance is

$$
L_{V P}=\left[\begin{array}{l}
\sin \theta \cos \phi  \tag{10}\\
\sin \theta \sin \phi \\
\cos \theta
\end{array}\right]=\left[\begin{array}{cc}
\sqrt{1-\mu^{2}} & \cos \phi \\
\sqrt{1-\mu^{2}} & \sin \phi \\
\mu &
\end{array}\right]
$$

The scattering angle is the dot product,

$$
\begin{equation*}
L_{\text {VSOL }} \cdot L_{V P}=\mu \mu_{0}+\sqrt{1-\mu^{2}} \sqrt{1-\mu_{0}^{2}} \cos \left(\phi-\phi_{0}\right) . \tag{11}
\end{equation*}
$$

As in the case of singly scattered sky radiance, the same types of physical interactions which generate singly-scattered sky radiance from direct solar radiation also scatter radiation from the virtual Sun to generate virtual Sun path radiance. Thus, we may use the same equation for singly-scattered sky radiance, with some modifications, to find the singly-scattered virtual Sun path radiance, $\mathrm{I}_{\mathrm{vp}}$. One difference is in that the computation of $L_{V P}$ we will now sum the scattered radiation over a path beginning at $\tau=\tau_{0}$ (the optical depth of the scene viewed by the sensor) and ending at the optical depth of the sensor, $\tau$. If we denote the irradiance of the virtual Sun by $E_{0}{ }^{\prime}$, we obtain the following formula for singly-scattered path radiance from the virtual Sun:

$$
\begin{gather*}
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L_{V P}=\frac{\omega_{0} \mu_{0} E_{0} \cdot \mathrm{p}\left(\cos \chi_{V P}\right)}{4 \pi\left(\mu_{0}-\mu\right)}\left[e^{-\left(\tau_{0}-\tau\right) / \mu_{0}}-e^{-\left(\tau_{0}-\tau\right) / \mu}\right] \cdot \tag{12}
\end{gather*}
$$

We define $E_{0}{ }^{\text {- }}$ by noting that the $E_{0}{ }^{-}$is the image of the Sun reflected in the water surface. Hence, the irradiance of the virtual Sun is the irradiance of the true Sun at the top of the atmosphere attenuated by the atmosphere and the reflectance of the water surface, i.e.,

$$
\begin{equation*}
E_{0}^{\prime}=\rho_{F} e^{-\tau_{0} / \mu_{0}} E_{0} \tag{13}
\end{equation*}
$$

where $\rho_{F}$ is the Fresnel reflectance and $e^{-\tau 0 / \mu 0}$ is the transmittance of the atmosphere.

### 4.3 SINGLY SCATTERED PATI! RADIANCE, $L_{\text {PSS }}$.

The geometry for singly-scattered path radiance is shown in Figure 8. The formula for singly-scattered path radiance is similar to that for singly-scattered path radiance from the Sun--the same straight line path from $\tau_{0}$ to $\tau$ is used, but $E_{0}$ is substituted for $E_{0}{ }^{\prime}$ in the formula. The cosine of the scattering angle for singly-scattered path radiance, $\cos X_{P S S}$ is also the negative of the cosine of the scattering angle used to compute the phase function for $I_{\text {VP }}$. Thus, $\cos X$ is the dot product, $\hat{\mathrm{L}}_{\text {SOL }}$. $\hat{\mathrm{I}}_{\text {VP' }}$, vector directions which have already been defined. The formula used to compute singly-scattered solar path radiance is

$$
\begin{equation*}
L_{\text {PSS }}=\frac{\omega_{0} \mu_{0} E_{0} p\left(\cos \chi_{\text {PSS }}\right)}{4 \pi\left(\mu+\mu_{0}\right)} e^{-\tau_{0} / \mu_{0}} \quad\left[e^{\left(\tau_{0}-\tau\right) / \mu_{0}}-e^{-\left(\tau_{0}-\tau\right) / \mu}\right] \tag{14}
\end{equation*}
$$



FIGURE 8 . GEOMETRY FOR SINGLY SCATTERED PATH RADIANCE, $L_{\text {PSS }}$, CALCULATION.

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where the variables in the above equation are as previously defined.
4.4 MULTIPLY SCATTERED PATH RADIANCE, $L_{\text {PMS }}$ •

To compute multiply scattered path radiance we use an analytical approximation described in detail in an earlier report [15]. The formula for the computation of $L_{p M S}$ is

$$
\begin{align*}
L_{P M S} & =\frac{E_{0}}{4 \pi\left[\mu_{0}+(1-\eta) \tau_{0}\right]} \\
& \left(\left\{(1-\eta) \tau_{0}\left[p\left(\mu, \phi_{,} \mu_{0}, \pi+\phi_{0}\right)+p\left(\mu, \phi_{1}-\mu_{0} \phi_{0}\right)\right]+\mu_{0} p\left(\mu, \phi_{1}-\mu_{0} \phi_{0}\right)\right.\right. \\
& \left.+\frac{2 \mu_{0}{ }^{2} \rho}{1+2(1-\eta)(1-\rho) \tau_{0}}\right\}\left[1-e^{-(\tau 0-\tau) / \mu}\right]+\left\{(1-\eta) p\left(\mu, \phi_{, ~} \mu_{0}, \pi+\phi_{0}\right)\right. \\
& \left.\left.+p\left(\mu, \phi_{1}-\mu_{0}, \phi_{0}\right)\right]-\frac{8(1-\eta) \mu_{0}^{2} \rho}{1+2(1-\eta)(1-\rho) \tau_{0}}\right\}\left[\left(\tau_{0}+\mu\right) e^{\left.\left.-\left(\tau_{0}-\tau\right) / \mu_{-}(\tau+\mu)\right]\right)}\right. \tag{15}
\end{align*}
$$

The single-scattering phase functions are given by:

$$
\begin{aligned}
& p\left(\mu, \phi, \mu_{0}, \pi+\phi_{0}\right)=p\left[\mu \mu_{0}-\sqrt{\left(1-1^{2}\right)\left(1-\mu_{0}{ }^{2}\right)} \cos \left(\phi-\phi_{0}\right)\right] \\
& p\left(\mu, \phi,-\mu_{0}, \phi_{0}\right)=p\left[-\mu \mu_{0}+\sqrt{\left(1-\mu^{2}\right)\left(1-\mu_{0}{ }^{2}\right)} \cos \left(\phi-\phi_{0}\right)\right] .
\end{aligned}
$$

In this chapter we describe implementation of the formulas discussed in Chapter 4 in the computer program called ATCOR. Many of the details of ATCOR have been presented in a previous report [16], so only a brief description of the entire program will be given here. Our discussion will focus primarily on the subroutine ATMSFR, in which the formulas presented in Chapter 4 have been introduced.

### 5.1 SPECIFICATION OF SOLAR AND SENSOR GEOMETRY

The geometric relationship of the sencur to the environment is shown in Figure 9. The geographic coordinates of the sensor locate the center of a spherical coordinate system used to define angles needed in model calculations. In the diagram the scanner scans along a path from $P_{1}$ (the first pixel) to $P_{n}$ (the last pixel). The azimuth of the scan plane is measured in the counterclockwise direction from north to the first pixel and is read into the program by the routine ATMSFR. The first pixel is always $90^{\circ}$ in a clockwise direction from a vector pointing in the direction of the flight.

The solar zenith and azimuth angles are computed automatically once the latitude, longitude, date, time of day (standard time), and zone number are specified. The extraterrestrial solar irradiance is also computed based on these inputs.

### 5.2 ATMOSPHERIC CORRECTION OPTIONS

Two input parameters set by the user determine which calculations are performed in routine ATMSFR. These parameters are SCATT and OPTION. SCATT may assume the value of either 0 or $1 ;$ OPTION can take on the values of 1,2 , or 3 . If SCATT is 0 , only multiply scattered path radiance is calculated and the value of OPTION can be any integer and will be ignored since only LPMS is then calculated.


FIGURE 9. SCAN PLANE AND SOLAR GEOMETRY. SSCAN IS MEASURED COUNTERCLOCKWISE FROM NORTH TO PIXEL \#l, \$SUN IS MEASURED COUNTERCLOCKWISE FROII SOUTH. POSITIVE SCAN ANGLES ARE MEASURED FROM THE FIRST PIXEL TO THE NADIR. NEGATIVE SCAN ANGLES ARE MEASURED FROM THE NADIR TO THE LAST PIXEL. $\mathrm{P}_{1}$ AND $\mathrm{P}_{\mathrm{n}}$ ARE THE FIRST AND LAST PIXELS, RESPECTIVELY.

If SCATT is 1 , single scattering computations are performed and the value of OPTION is used to determine which values to calculate. If OPTION $=1$, singly scattered path radiance, LPSS, and reflected sky radiance, LSKYRE, are computed. If OPTION $=2$, LPSS and virtual sun path radiance, LVP, are computed. If OPTION $=3$, LPSS, LSKYRE and LVP are calculated.

OPTION and SCATT are read into the data file on logical unit 4. The format for this record is (5I5), and the variables read in are:

FSTP LSTP PTINC SCATT OPTION
where
FSTP $=$ the number of the first pixel to be processed, LSTP $=$ the number of the last pixel to be processed, PTINC = the pixel increment to use in the processing, and SCATT and OPTION are as previously defined.

The output file which is used by ATMSFR is given in Table 1 and the new subroutine ATMSFR2 is given in Table 2.

### 5.3 MODEL INPUT PARAMETERS

In addition to the geometric paramete\&s, we must specify parameters characterizing the medium and the measurement system.

The model makes use of several "altitude" values which must be input by the user. First, one must know the actual altitude (km) of the sensor above the surface. Second, one must know the pressure (millibars) of the atmosphere at the surfare, and third, one must know the atmospheric pressure (in millibars) at flight altitude. If only the altitudes are known, one can use the tables relating pressure to altitude as given by the U.S. Standard Atmosphere [17].

TABLE 1

| LINE M | READ OCCURS IN ROUTINE | INPUT VARIABLES | FORMAT |
| :---: | :---: | :---: | :---: |
| 1 | ATMSFR | FSTP,LSTP,PTINC, SCATT, OPTION | (515) |
| 2 | DATE | MNTH, DAY , YEAR | (3I4) |
| 3 | ANGLES | HOUR,MIN,SEC | (2I5,F6.3) |
| 4 | ANGLES | NZ | (I5) |
| 5 | ANGLES | LATD, LATM, LATS, LONDG, LONGK, LONGS | $\left(\begin{array}{c} (2 I 5, F 6.3,2 I 5, \\ F 6.3) \end{array}\right.$ |
| 6 | ATMSFR | ZSCAN, ZGRND, LSW | (2F8.5,15) |
| 7 | ATMSFR | (WAVE (I), I= 1, QNCHAN) | (10F8.5) |
| 8 | ATMSFR | (RHO(I), I=I.QNCHAN) | (I0F8.5) |
| 9 | PHASE | R,IM | (2F8.6) |
| $\left\lvert\, \begin{gathered} 10 \\ 11-14 \end{gathered}\right.$ | PHASE | NWT, NANG, (C, ( $), I=1, N A N G)$ | (2I5/(10F8.6) ) |
| 15-79 | PHASE | $\begin{aligned} & \text { (WTAB(I), (PF (I,J), } \\ & J=1, N A N G), I=1, N W T) \end{aligned}$ | $\begin{gathered} (F 10.6 /(10 \\ F 8.4)) \end{gathered}$ |
| 80 | RAYLEI | PRESO, PRES 2 | (2F10.4) |
| 81 | OZONE | NO2, NPROF, NO3W1,NO3W2 | (4I5) |
| 82 | OZONE | WIVCL, WAVC2 | (2F8.4) |
| 83-153 | OZONE | $\left(\begin{array}{l} ((Z O Z(I Z), 03 I N T(I Z, I P) \\ I Z=1, N O Z), I P=1, N P R O F) \end{array}\right.$ | $\begin{aligned} & \text { (F7.0, } \\ & \text { E11.4) } \end{aligned}$ |
| 154 | OZONE | (03MAX (IP), IP=1, NPROF) | (10E13.6) |
| 155-169 | OZONE | (WAVO3 (I), A(I), I=1,NO3W) <br> (NOTE:NO3W=NO3W1 + NO3W2) | (F7.0,E11.4) |

TABLE 1 (Cont.)
INPUT FILE ON LOGICAL UNIT FOR USE BY ATMSFR (CONTINUED)

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| LINE NO. | READ OCCURS IN ROUTINE | INPUT VARIABLES | FORMAT |
| :---: | :---: | :---: | :---: |
| 170 | OZONE | NOP | (4I5) |
| 171 | THICK | NTEX | (I2) |
| 177 | THICK | ( $\operatorname{WAVEX}(\mathrm{I}), \operatorname{TAUEX}(\mathrm{I}), \mathrm{I}=1, \mathrm{NTEX})$ | (2F8.4) |
| 178 | PARAMS | FSCAT | (10F8.6) |
| 179 | AERO | NAER,MPROF, NUZ, MAXG | (515) |
| 180 | AERO | IPROF | (I5) |
| 181-184 | AERO | (WAVAER (I) , RIN ( I ) , I=1,NAER) | (10F8.4) |
| 185-230 | AERO | $\begin{gathered} ((Z U N(I Z), \\ N U Z), I P=1, M P(I Z, I P), I Z=1, \end{gathered}$ | (F7.0,E11.4) |
| 231-258 | AERO | ( $\mathrm{X}(\mathrm{I}), \mathrm{Z}(\mathrm{I}), \mathrm{I}=1, \mathrm{MAXG})$ | (F7.0,E11.4) |
| 259 | ATMSFR | PHID, PHIM, PHIS | (2I3,F6.3) |
|  |  | \$ENDFILE |  |

TABLE 2
PROGRAM LISTING


INIEGER FUNCTION ATMSFR(L. CONTRL, MCHAN. MSS)

VERSION VER 3.0 SAAPRTI -- 200/24 CHANNEL VERSION ALL PUT JOGETHER

- VER 3.oa guUnti -- finish making fortran compatible. ogreal ver. 3.ob tuunt7 -- caang, doang

VER. 3. OD 3OUUNT7 -- QIESQL; /GANCIL/ IN ALL ROUTINES


- VWITCH VER B.OA GJUnT7 -- FINISH making fortran compatible. obreal
 VER 3 OD 3OJUNTT -- QIESQL: /QANCIL/ IN ALL ROUTINES

lpath - path radiance generated by the direct solar radiation. trsky - reflected sky radiance observed at sensor.
iviep - path radiance generaite by the virtual image of the sun ob-
modifications programed by jim lienzer MNEGER FUNCIION ATMSFR(L, CONTRL MCHAN, MSS)
 uvuリuv


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TABLE 2 (cont.)



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$253 \begin{aligned} & \text { WRITE(OUPRNT, 253)2SC } \\ & z \text { ORMAI(', 19X, 'SENSOR ALTITUDE', FIO. } 1,{ }^{\prime} \text {, FEET') }\end{aligned}$
Z = ZSCAN-ZGRND READ (OUDATA, 264) (WAVE(i), I=1, ONCHAN)
READ (OUDATA, 264) (RHO(i), $1=1$ ONCHAN)



$$
\text { TABLE } 2 \text { (Cont.) }
$$



[^2]


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[^3]ORIGINAL PAQE IJ
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TABLE 2 (Cont.)
SINGLE SCATIERING CALCULATIONS
SIG(I) = TAU(I) - TAUZ(I)
EFACIO = EXP(SIG(I)/MUQ)
CONST = CON•OM(I) EOO(I)
EFACTS = EXP(-TAU(I)/MUO)
ETOTF = (I. EETA(I) * VAU(I)
E(I) = (MUO*MUO*EO(I)/(MUO+EIOTF))*(I+(2.*RHO(I)*ETOTF)/


## (1.42. ©ETOTF))

FE(I) = MUO*EO(I)/E(I)
COMPUTE DIRECT IRRADIANCE AT SURFACE
ED( I) = MUO*EO( I) ${ }^{\circ}$ EFACTS
GO TO 15
GO TO IS
CONT INE
COMPUTE WAVELENGTH DEPENDENT QUANTITIES NEEDED FOR
MULTI SCATTERING CALCULATIONS MULT SCATTERINO CALCULATIONS

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SF = O-ETA(I)
CPII
IHS = (1) vHdTV
$G(1)=4 O^{\circ} \mathrm{CHI}$
SIG(i)


\begin{tabular}{|c|c|}
\hline 58

59
310 \& ```
IF (THETA) 58. 59. 58
THETTA=ABS(THETA)
REFRAC=ARSIN(SIN(THETTA)*I./INDEX)
AX = THETTA-REFRAC
BX= THETTA REFRAC
SINA=SIN(AX)
SINB=SIN(7x)
TANA = TAN(AX)
TANB=TAN(BX)
FRES(IP)-FRESO
CONTINUE
TABLE 2 (Cont.)
FRES(IP)=0.5*((SINA*SINA )/(SINE'SINB ) \& (TANA -TANA )/(TANGPTANE ))
GO to 3to

``` \\
\hline c
c
c
c & calculate the other scan angle dependent quantities needed TO FIND THE REFLECTED SKY RADIANCE. \\
\hline C & ```
ARGSKY = FMU-FPHI -ROOT
PFSKY=PF(ARGSKY.I.TR(I).TAS(I))
EXPO2-TAU(I)/ms
MUF AC2-(muO-mu)/muo
CONO2-EXPO2.mUFAC2
IF (R-)(CONO2).LT COND) 60 to 503
``` \\
\hline c & Calculate the reflecteo sky radiance \\
\hline C 1 & ```
LRSKY(I.IP) = FRES(IP)*(CONST/MUFAC2)*PFSKY•EFACT(I.IP)*
G0 t0 505
GO (EFACTS-EXP(-EXPO2))
``` \\
\hline c
c
c
c & EITHER mu and muo or tauif) ano tauz(I) are too close together SO THE FOURTH ORDER EXPANSION OF THE REFLECTED SKY RADIANCE must be USED. \\
\hline \[
\begin{array}{r}
C \\
503 \\
1 \\
2
\end{array}
\] & \begin{tabular}{l}
LRSKY(I,IP)=FRES(IP) ©CONST•PFSKY ©EFACT(I, IP) ©EFACTS•EXPO2*(I. \\
-(COND2/2 ) \(\cdot(1 .-(C O N D 2 / 3.) \cdot(1 .-(C O N O 2 / 4))\). \\
))
\end{tabular} \\
\hline & IF (OPTION.LT. 2) Go to 2 \\
\hline C & Calculate the scan angle dependent quantities needed to find \\
\hline c & \begin{tabular}{l}
the path radiance generated by the virtual sun. \\
ARGVIR-TMU-FPHI •ROOT \\
PFVIR \(=\) PF (ARGVIR.I.TR(I), TAS(I)) \\
MUFAC2 \(=(\) MNO-MU)/ MUO \\
CONO3-EXPOI•MUFAC2 \\
IF (ABS(COND3).LT COND) 60 to 506
\end{tabular} \\
\hline \(\begin{array}{ll}\text { c } & \\ \text { c } & \\ & \\ & \\ & \text { 1 }\end{array}\) & Calculate the path radiance generated by the virtual sun.
```

LVIRP(I.IP)=FRESUN* (CONST/MUFAC2)PPFVIR•EFACTS*(I./EFACTO -
EFACT(I.IP))

```
GO 102019 \\
\hline & EITHER MU ANO MUO OR TAU(I) aNO TAUZ(I) are too close together SO we use the expansion of the virtual path radiance. \\
\hline 506 &  \\
\hline
\end{tabular}

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\text { TABLE } 2 \text { (Cont.) }
\]


The user must also define the center wavelengths ( \(\mu \mathrm{m}\) ) of the multispectral scanner and the corresponding values of the surface background albedo (values from zero to one). Also, the center wavelength 3 ( \(\mu \mathrm{m}\) ) of the surface radiometers and the corresponding optical thicknesses must be known. It should be noted that the opticai thicknesses used should be those measured as closely as possible in time with the multispectral data.

\subsection*{5.4 MODEL CALCULATIONS}

In this section we present several examples of the radiances for the various components. Because our main interest is in the radiance components as a function of scan angle and visibility, we will presert the results of the calculations in terms of these parameters.

Figure 30 depicts the variation in the singly-scattered reflected sky radiance at the sensor as a function of the nadir scan angle and visibility. In this example the solar zenith angle is \(45^{\circ}\) and the scan plane is perpendicular to the solar plane. The curves which result are a combination of the variation of the sky radiance, the transmittance from the surface to the sensor, and the Fresnel reflectance of the water surface. For a practical scanner with a maximum scan angle of about \(45^{\circ}\) the curves indicate that one would not observe the large radiance peaks at the large angles.

In Figure 11 we display the corresponding path radiance as a result of singly-scattered radiation from the reflection of the sun in the water. In this case the radiance peaks do not exist at the large scan angles.

In Figure 12 we illustrate the relative magnitudes of the various radiation components as a function of scan angle for a moderately hazy atmosphere. The virtual sun path radiance is the smallest value and the multiply scattered sky radiance is the largest islue.
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\[
W^{H-} \frac{\cdot y \exists \perp S z^{W} \partial}{M W}
\]

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Figure 12. SURFACE AND PATH RADIANCE COMPONENTS DETECTED BY SENSOR AT OPTICAL DEPTH OF \(\tau=0.346\) AS A FUNCTION OF SCAN ANGLE, \(\operatorname{OSCAN}\). SCAN PLANE 1 SOLAR PLANE, USUN \(=30^{\circ}, \lambda=0.55 \mu M\), PHASE FUNCTION \(=\) CONTINENTAL REFRACTIVE INDEX \(1.5-0.01 i, \operatorname{VISIBILITY}=10 \mathrm{KM}\).

In Figure 13 we indicate the variation in the ratio of the singly-scattared sky radiance to the aingly-scattered path radiance as function of the optical depth \(r\) of the sensor. As the curves illustrate, the reflected sky radiance component is relatively more important for the larger optical depths.

Figure 14 illuatrates the variation in the ratio of the virtual sun path radiance to the singly-scattered path radiance with scan angle for four optical depths.

Figure 15 depicts the large ratio of the multiply-scattered component to the singly-scattered path radiance component as a function of optical depth and scan angle.

Because optical thickness or visibility is of major importance in remote sensing investigations, we want to consider the variation of the radiance components with respect to visibility. This effect is illustrated in Figure 16 for three different atmospheres. We chose the continental aerosol because it more nearly represents the type which would be found over the Great Lakes. The three refractive indices are: 1.5-0.0i which corresponds to a "clean" haze, i.e., one where there is no absorption; 1.5-0.01i which corresponds to a haze with some aerosol absorption; and 1.5-0.1i, a complex index of refraction which corresponds to a haze with more absorption. As the curves indicate, an absorbing haze or one which corresponds \(t\), considerable air pollution gives rise to a large rario of reflec ed sky radiance relative to the singly-scattered path radiance.

The effect of the complex index of refraction is also evident in the ratio of the virtual sun path radiance ta the singly-scattered path radiance as indicated in Figure 17.

Finally, we illustrate in figure 18 the variation of various combinations of ratios in terms of the visibility for a refractive index of 1.5-0.01i.


FIGURE 13. RATIO OF REFLECTED SINGLY SCATTERED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE AS A FUNCTION OF SCAN ANGIE, OSCAN, FOR OPTICAL DEPTH, \(\tau\), OF THE SENSOR OF \(0.132,0.175,0.225\) and 0.346 . SCAN PLANE \(\mid\) SOLAR PLANE, PHASE FUNCTION \(=\) CONTINENTAL REFRACTIVE INDEX 1.5 - 0.01i, VISIBILITY \(=10 \mathrm{KM}, \lambda=\) \(0.55 \mu \mathrm{M}, \theta\) SUN \(=30^{\circ}\).


FIGURE 14. RATIO OF SINGLY SCATTERED PATH RADIANCE FFOM THE VIRTUAL SUN, LVP, TO SINGLY SCATTERED PATH RADIF"CE, LPSS, AS A FUNCTION OF SCAN ANGLE, OSCAN, FOR OPTICAL DEPTHS, \(\tau\), OF THE SENSOR OF A) \(0.132, ~ B) ~ 0.175, ~ C) ~ 0.225 ~ A N D ~ D) ~ 0.346 . ~\) SCAN PLANE SOLAR PLANE, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX 1.5 - 0.01i, VISIBILITY \(=10 \mathrm{KM}, \lambda=\) \(0.55 \mu \mathrm{M}, \theta_{\text {SUN }}=30^{\circ}\).

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FIGURE 15. RATIO OF MULTIPLY SCATTERED PATH RADIANCE, LPMS, TO SINGLY SCATTERED PATH RADIANCE, LPSS, AS A FUNCTION OF SCAN ANGLE, \({ }^{\circ}\) SCAN, FOR OPTICAL DEPTHS, \(\tau\), OF THE SENSOR OF \(0.132,0.175,0.229,0.346\). SCAN PLANE SOLAR PLANE, VISIBILITY \(=10 \mathrm{KM}\), PHASE FUNCTION \(=\) CONTINENTAL REFRACTIVE INDEX \(1.5-0.01 i, \lambda=0.55 \mu \mathrm{M}\), \(\theta\) SUN \(=30^{\circ}\).


Figure 16. RATIO OF REFLECTED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE FOR THREE CONTINENTAL AEROSOL MODELS: a) 1.5 \(0.0 i\), b) \(1.5-0.01 i\), and c) \(1.5-0.1 i\) AS A FUNCTION OF ATMOSPHERIC VISIBILITY. SCAN PLANE \(\perp\) SOLAR PLANE, \(\theta_{\text {SUN }}=30^{\circ}, \lambda=0.55 \mu \mathrm{M}\).
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FIGURE 17. RATIO OF SINGLY SCATTERED PATH RADIANCE FROII THI VIRTUAL SUN, LVP, TO SIINGLY SCATTERED PATH RADIANCE, LPSS, AS A FUNCTIOI OF ATMOSPHERIC VISIBILITY FOR THREE CONTINENTAL AEROSOL HODELS WITH REFRACTIVE INDICES OF a) 1.5 - 0.0i, b) 1.5 - 0.0li, and c) 1.5 - 0.1i. \(\lambda=0.55 \mu \mathrm{M}, \theta_{\text {SUN }}=30^{\circ}\), \(\theta_{\text {SCAN }}=0^{0}\).


LVP \(=\) virtual sun path radiance;
LPSS \(=\) singly scattered path radiance;
\(L_{S K Y R E}=\) reflected sky radiance.
FIGURE 18. RATIOS OF (a) LVP/LPSS, (b) LSKYRE/LPSS, (c) LSKYRE/ LVP AS A FUNCTION OF VISIBILITY (KM). \(\lambda=0.55 \mu \mathrm{M}\), \(\theta_{\text {SUN }}=30^{\circ}, \theta_{\text {SCAN }}=0^{\circ}\), AEROSOL MODEL \(=\) CONTINENTAL REFRACTIVE INDEX 1.5 - 0.01i.

The problem of developing an atmospheric correction algorithm for remote sensing is an old and difficult one. The main difficulty lies in not being able to have available sufficient data which can be used to specify the values of the relevant atmospheric parameters The problem is all the more difficult in the case of the remote sensing of water bodies because of the low signal-to-noise ratio involved. In this investigation we have extended an existing computer algorithm so as to include additional radiation components. The original algorithm included the path radiance which arises from the singly-scattered solar radiation in the atmosphere. We have now included the radiance which arises from the sky radiation which is reflected by the water surface and is then attenuated as it propagates from the surface to the sensor. In addition, we have included the path radiance component which arises from the single scattering of radiation as a result of a virtual sun, i.e., of the sun's reflection in the water. It should be realized that this component is always present regardless of the scan plane, i.e., it does not only occur when the scanner is looking at the specular angle. In addition to these components, we have also included a multiple-scattering approximation. It should be realized, however, that the multiple scatterirg applies only to an atmosphere with the sun as a source. Enother multiple scattering calculation should be performed to include the effect due to the virtual sun.

The general result of all these calculations indicates that the various components are all about equal in magnitude but that there is considerable variation with respect to scan angle and visibility. Also, it appears that the multiply-scattered component is of major significance.

It must be pointed out that the objective of this investigation is to provide an algorithm for the correction of remotely sensed data for atmospheric effects so that one can extract from the multispectral data the radiance which is characteristic of the water itself. In our investigation we have dealt with the water surface as a flat, specular reflector, which in general is not true. A wind-roughened surface will be characterized by a complex wave structure which leads to a more complicated representation of the reflected and virtual sun radiances than presented in this report. A further investigation should be conducted to model the water surface in terms of wind speed and a stochastic representatior. of the reflecting facets of the water surface. If this is done, then a more realistic model could be developed which should provide better values for the sky-reflected and the virtual sun path radiances. It may even be possible to establish a method for the determination of wind speed by observing the average radiance as a function of the instantaneous field of view.

A further recommendation is to improve the accuracy of the algorithm by including a more detailed calculation of the multiplyscattered path radiance, both for the direct sun as a source and for the virtual \(s \mathrm{n}\) as a source.

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[^0]:    FIGURE ס. GEOMATRY FOR REFLECTED SKY RADIANCE COMPUTATION.

[^1]:    :

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    30 FORMAT('O'. IOX. 'ALL RADIANCES ARE IN UNITS DF MILLIMATTS PER CEN'.

[^3]:    

[^4]:    SH4 = SINH(ARG4)
    CHA = COSH(ARG4)
    COMPUTE LAMRDA DOUELE PRIME
    LAMPP(I) - CON*NUSQ(I) - OM(I) *EO(I) /(B(I)*NU*SH2 * MHO*CH2)
    
    
    ARG4 = ARG3*TAUZ (I)/TAU(I)
    SH4 = SINH(ARG4)
    CHA = COSH(ARG4)
    COMPUTE LAMBDA DOUELE PRIME
    COMPUTE PHI PRIME
    (OM(I).G1
    A(I) $=001$
    $\mathbf{g}(I)=1.0$
    G(i) $=1.0$
    NUS = MUO
    MUSQ(I) = MU'NU
    SIG(I) = TAU(I) - TAUZ(I)
    ARGI $=$ SIG(I)/NU
    SHI
    

    ARG2 = TAU(I)/MU
    SH2 $=$ SINH(ARG2)
    CH2 $=$ COSH(ARG2)
    ARG3 $=$ CO2•ARG2
    SH3 $=$ SINH(ARG3)
    CH3 $=$ COSH(ARG3)
    ARG4 $=$ ARG3*TAUZ (I)/TAU(I)
    -

