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ATMOSPHERIC RADIATION MODEL FOR WATER SURFACES

FINAL TECHNICAL REPORT

ROBERT E. TURNER DANIEL W. GASKILL JAMES R. LIERZER

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MAY 1982

PREPARED FOR NATIONAL AERONAUTICS AND SPACE ADMINISTRATION LEWIS RESEARCH CENTER CLEVELAND, OHIO 44135 CONTRACT NO. NAS3-22495

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additional components such as the atmospheric path radiance which results from singly-scattered sky radiation specularly reflected by the water surface. It also considers a component which is referred to as the virtual sun path radiance, i.e. the singly-scattered path radiance which results from the solar radiation which is specularly reflected by the water surface.

These atmospheric radiation components are coded into a computer program for the analysis of multispectral remote sensor data over the Great Lakes of the United States. The user must know certain parameters, such as the visibility or spectral optical thickness of the atmosphere and the geometry of the sensor with respect to the sun and the target elements under investigation.

Suggestions and recommendations are given for further investigation of the problem of the remote sensing of water surfaces. If all of these extrinsic radiation components are properly accounted for, then the intrinsic water radiance can be found by applying the algorithm or an adaptation of the algorithm in this report. As a result, one would then be able to know the actual surface water spectral radiation field independent of the atmosphere.

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#### PREFACE

This report describes part of the work done on a research program in the remote sensing of the Great Lakes using a multispectral scanner aboard an aircraft. The research has been conducted for the NASA-Lewis Research Center, Cleveland, Ohio, by Science Applications, Inc., at the Dayton, Ohio, office. The primary objective of this program is to develop remote sensing as a practical technique for the analysis of the Great Lakes.

Remote sensing of the environment involves the transfer of radiation from the Earth's surfare through the atmosphere to a sensor which is located at some point within the atmosphere. For water surfaces with their inherently low reflectances, the atmospheric scattering of solar radiation acts as a significant noise factor. In this report we have extended an existing model to include various atmospheric radiation components so that the resulting mathematical algorithm will allow one to extract a radiance value which is more nearly representative of the actual radiance of the water, independent of atmospheric effects.

This research was performed under contract NAS3-22495 and covers the period from 4 March 1981 through 15 July 1982. Mr. Thom Coney served as Technical Monitor of the contract and Dr. Robert E. Turner of Science Applications, Inc., was the program manager and principal investigator.

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## SYMBOLS

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<sup>е</sup> о	solar irradiance at the top of the atmosphere
E <sub>0</sub> ´	solar irradiance of the virtual sun
L <sub>0</sub> (1)	radiance incident at angle i
L <sub>s</sub> (i)	surface radiance at angle i
LSKYRE	reflected sky radiance
LSOLRE	reflected solar radiance
LSKY	sky radiance
LSOL	direct beam radiance
L PMS	multiply scattered path radiance
<sup>L</sup> PSS	singly scattered path radiance
тu	upward scattered radiance beneath water surface
L <sub>INIR</sub>	radiance scattered from beneath water surface
LSKYSS	singly scattered sky radiance
<sup>L</sup> PTOT	total radiance
LVSOL	radiance from virtual sun
<sup>L</sup> VP	virtual sun path radiance
n	index of refraction of water relative to air
Р	slore probability
p(X)	scattering phase function
Т	transmittance
	(GREEK SYMBOLS)
θi	incident angle
<sup>θ</sup> r	reflected angle
۹ <b>s</b>	Fresnel reflectance

<sup>ρ</sup> f <sup>(i)</sup>	Fresnel reflectance at angle i
θ	nadir view angle
θo	solar zenith angle
φ	azimuth view angle
φ	solar azimuth angle
μ	cosine of nadir view angle
μo	cosine of solar zenith angle
XSKY	scattering angle
τ	optical depth
<sup>τ</sup> 0	optical thickness
η	forward scattering parameter
ρ	surface albedo
<sup>ω</sup> 0	single scattering albedo

**...** (1)

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## 1 SUMMARY

In the analysis of remotely sensed data on bodies of water, the atmosphere obscures the inherent surface features as a result of the scattering and absorption of solar radiation. In the case of multispectral data acquired by aircraft or spacecraft sensors, one can preprocess the data by applying mathematical models and algorithms to the digitized data. The mathematical model developed in this investigation is specifically designed to account for various components of the visible and infrared radiation in the atmosphere which interfere with the inherent signal from a water surface. If the atmospheric parameters are known, then when the algorithm is applied to the multispectral data sets, an improved or corrected data set will result.

This improved atmospheric correction model allows for the path radiance in the atmosphere as a result of singly-scattered solar radiation and also singly-scattered solar-reflected radiation. In addition, the model includes a singly-scattered sky radiation component for the radiation which is reflected by the water surface. Comparisons are made among the relative magnitudes of these radiation components in terms of the geometric and environmental factors. Recommendations are presented for a more advanced model which would include the corresponding radiation components for multiple scattering.

#### INTRODUCTION

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Multispectral scanner data obtained by sensors aboard aircraft and spacecraft allow a user to examine the detailed physical properties of a surface. These properties are of interest to many investigations in various disciplines such as land use studies, agriculture, hydrology, forestry, and oceanography. In all of these investigations, however, the scattering of visible and infrared radiation by the atmospheric constituents will reduce the inherent surface radiance and add a path radiance to the attenuated radiance from the target. For many cases of the remote sensing of bright land areas on relatively clear days the attenuated radiation from the surface is rather large as compared to the atmospheric path radiance. For water bodies, with inherently low reflectances, this is no longer true and the path radiance can be a major effect in the total radiance at the sensor.

The purpose of this investigation was to extend an existing atmospheric radiative transfer model to include other radiation components which did not exist in the previous model. These additional atmospheric radiation components include specific effects for the remote sensing of water surfaces. The model is used in conjunction with an algorithm specifically designed for the analysis of multispectral data.

The determination of the atmospheric radiation components is important for the analysis of the probability of misclassification of various classes of surface materials. To first order one may consider the so called linear transfer problem in which the path radiance is constant over varying surface reflectances for a horizontally spatially uniform haze. For a non-uniform haze, however, the path radiance can vary, thereby resulting in a higher

probability of misclassification of objects if the degree of nonunifermity is unknown. A second-order effect, but one which can become quite important for the remote sensing of high-contrast targets is the adjacency effect. This is when radiation from a bright target causes an increase in the path radiance with respect to the radiance from a neighboring dark target. This problem would exist, for example, in the remote sensing of water bodies near bright sandy beaches. The results should be evident in the brightened image of the water near the shoreline, provided the effects of waves and whitecaps are eliminated. This second-order adjacency effect is not included in the model or algorithm in this investigation but the effect can be accounted for if the investigator has sufficiently detailed atmospheric data on the horizontal stratification of aerosols.

Multiple scattering is particularly important if the sky is hazy. These effects are considered in the model for path radiance which results from the sun as a source. We have not included the multiple scattering effects for solar-reflected radiation.

## 3 OPTICAL PROPERTIES OF WATER

The interpretation of remote sensing data collected over water surfaces requires a detailed knowledge of the optical properties of water and the air-water interface. Water is unusual as a natural surface because it is a specular reflector and because in the visible regions sensible data can be obtained from well below the water surface [1]. Also of importance is the phenomenon of refraction, which occurs when radiation passes through the air-water boundary. These properties of water and their significance in terms of remote sensing are discussed in detail in this chapter.

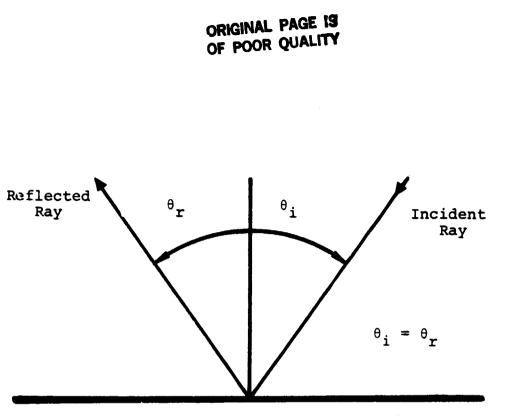
#### 3.1 REFLECTION, REFLECTANCE AND REFRACTION

Most natural surfaces are approximately Lambertian--reflecting incident radiance equally in all directions. A smooth water surface, however, is a specular reflector, and reflection of radiation from it follows the geometrical law of reflection. This geometric law requires that the angle with respect to the normal to the surface of the reflected ray equal the angle of incidence of the incident ray and that the reflected ray be in the same plane as the incident ray. Specular reflection is depicted in Figure 1. Reflectance,  $\rho_s$ , of the water surface is given by the Fresnel equation

$$\rho_{s} = \frac{1}{2} \left[ \frac{\sin^{2}(\theta_{i} - \theta_{r})}{\sin^{2}(\theta_{i} + \theta_{r})} + \frac{\tan^{2}(\theta_{i} - \theta_{r})}{\tan^{2}(\theta_{i} + \theta_{r})} \right], \qquad (1)$$

where  $\theta_i$  is the incident angle and  $\theta_r$  is the angle of refraction. The transmitted part of the incident ray experiences refraction at the water surface, as shown in Figure 2. Snell's law, given by

$$\frac{\sin(\theta_{i})}{\sin(\theta_{r})} = n, \qquad (2)$$



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Water Surface

FIGURE 1. SPECULAR REFLECTION FROM A SMOOTH WATER SURFACE. THE INCIDENT AND REFLECTED RAYS ARE IN THE SAME PLANE.



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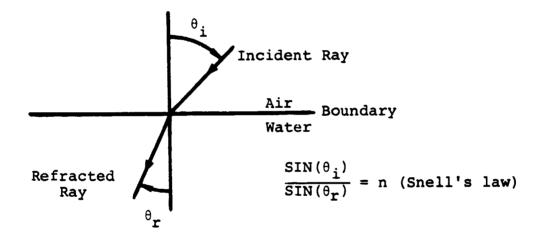
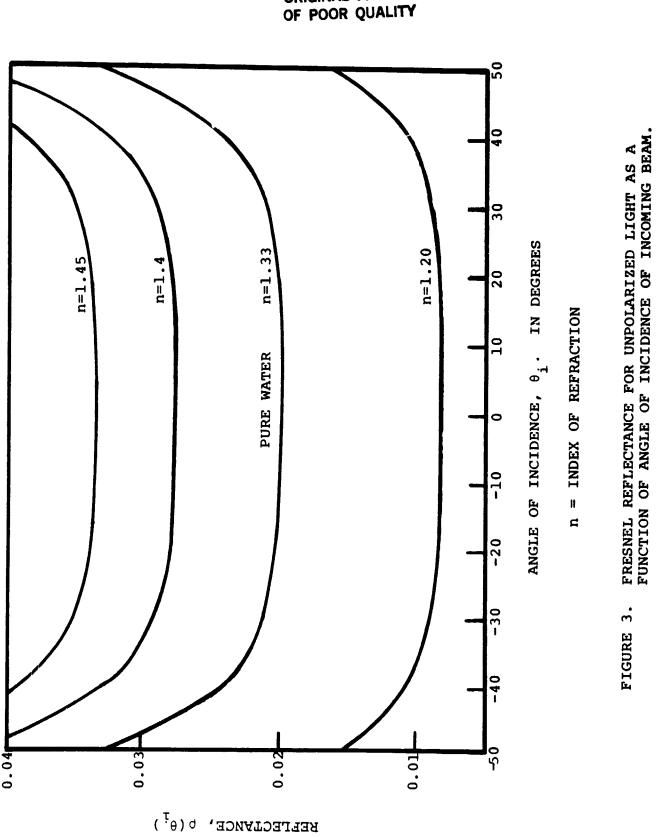


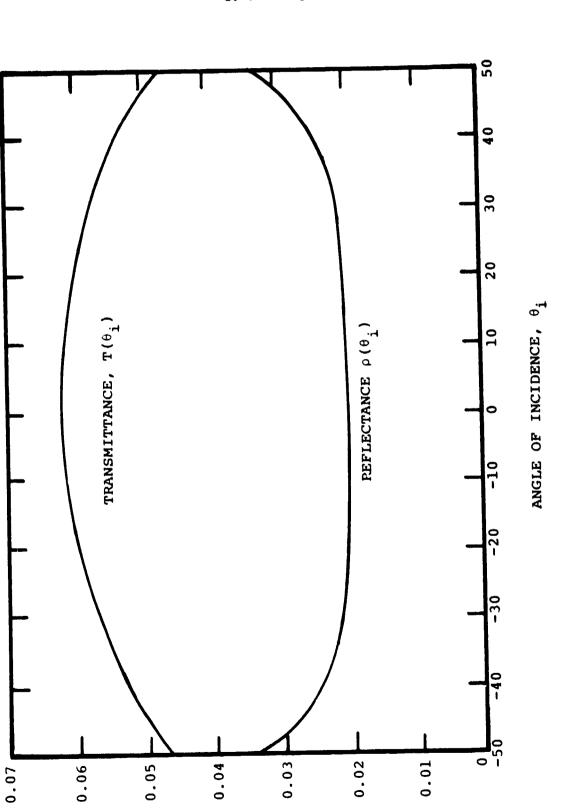
FIGURE 2. REFRACTION OF THE TRANSMITTED BEAM AT THE WATER SURFACE AND THE LAW OF REFRACTION n = REFRACTIVE INDEX OF WATER RELATIVE TO AIR.

describes the relationship between the angle of incidence of the incoming beam,  $\theta_i$ , the angle with respect to the normal to the surface of the refracted beam,  $\theta_r$ , and the refractive index of air relative to water. The refracted beam lies in the same plane as the incident beam. The phenomenon of refraction is depicted in Figure 2. The law of refraction requires that for a water surface, downward sky radiation and direct sunlight enter the water within 48.5° of the vertical. Only when the water surface is roughened by wind or another disturbance can direct sunlight or sky radiation penetrate the water surface outside this range of angles. Back-scattered radiation from beneath the water-air intexface also experiences refraction on reaching the water surface. When the water surface is calm, upward radiation incident at angles greater than 48.5° with the vertical is totally internally reflected [2]. Thus, downward radiation beneath the water surface at angles with the vertical greater than 48.5° is upward radiation in the water which has been totally internally reflected [3].

Equations 1 and 2 show that the reflectance and transmittance of the water surface are dependent on the refractive index of water. The refractive index is influenced by changes in temperature and by the concentration of various solutes in the water. Figure 3 shows how the reflectance function, equation 1, varies as a function of the angle of incidence for refractive indices of 1.20, 1.33, 1.40, 1.45. This range of refractive indices encompasses the range of natural variability in the refractive index for water; and Figure 3 shows that, over this range, variation in the refractive index is of little importance in determining the surface reflectance. For all of the calculations shown in this report, a refractive index of 4/3 is used. Figure 4 shows transmittance, T, and reflectance,  $\rho_s$ , as a function of angle of incidence of the incoming radiation for a refractive index of 1.33.

Since many applications of remote sensing over water require





FRESNEL REFLECTANCE AND TRANSMITTANCE AS A FUNCTION OF ANGLE OF INCIDENCE OF INCOMING BEAM. REFRACTIVE

OF ANGLE OF INCIDENCE OF INCOMING BEAM.

iNDEX = 1.33.

FIGURE 4.

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examining the signature from beneath the water surface, Fresnel reflectance is of paramount importance. Figure 4 shows that transmittance is highest when the incident beam is perpendicular to the water surface, while the reflectance is lowest at this angle. Thus, from a consideration of the Fresnel formulas, one would expect the return from beneath the water surface relative to the total, to be greatest when viewing that surface at the nadir. Contributing to this effort is the fact that most scattering phase functions for polydispersions in water have a secondary peak at 180° [2]. Figure 4 shows that while the Fresnel reflectance function is at a minimum at the nadir view angle, this function changes very little out to viewing angles as great as 40°, at which point it begins to rise steeply. Internal reflection at the water surface of upwelling radiation is also at a minimum normal to the water surface. Thus, for remote sensing work where the return from beneath the water surface is of greatest interest, scan angles should be maintained within 40° to 45° of the vertical. Beyond an angle of 48.5° no radiation from beneath the water surface will reach the sensor when the water surface is calm. Scanning in the solar plane is also problematic in this regard since specularly reflected light on the solar side of the scan plane would saturate the sensor. For some purposes, such as viewing the glitter pattern on the water surface, scanning in the solar plane may be desirable.

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### 3.2 ABSORPTION AND SCATTERING

Water is a good absorber of electromagnetic radiation. Only in the relatively narrow spectral region from about 400 to 600 nanometers is the transparency of water such that radiation can penetrate more than a few meters in depth below the water surface. Both at wavelengths shorter than 400 nm and longer than 600 nm, absorption increases rapidly and only very small amounts of radiation are scattered back out of the water

into the atmosphere. At the very short wavelengths this radiation is further strongly attenuated in the atmosphere.

Scattering of radiation in water is caused by water molecules, by dissolved salts and by particles in suspension. These effects are usually assumed to be additive [4].

Scattering by water molecules is described by fluctuation theory which predicts scattering of radiation as a result of molecular movements which cause fluctuations in the density of the medium. As in Rayleigh scattering, this type of scattering is proportional to  $\lambda^{-4}$ , where  $\lambda$  is the wavelength of the radiation being scattered. The effect of dissolved salts on the molecular scattering phase function is usually small enough to be neglected.

In general, most scattering in water is accomplished by particles in suspension [4]. Particulate matter in water derives from runoff from land, deposition from the atmosphere, and organic processes within the water. Thus, particles may be quite irregular in shape and particle size distributions are difficult to characterize precisely [2]. Because some sources of particles such as runoff of organic processes may be highly localized in space, size distributions may vary greatly in space and time. Although particle shapes vary considerably from the spherical ideal of scattering theory, it has been shown [5,6] that systems of irregularly shaped particles can be adequately approximated by systems of polydisperse systems of spherical particles. The major observed features of phase functions of particulate suspensions in water are a strong forward scattering peak, a broad minimum around 100°-130° and a small secondary peak in the back scattering direction [4].

3.3 OPTICAL PROPERTIES OF A WIND ROUGHENDED WATER SURFACE

Roughness of the water surface caused by wind presents an additional problem in the calculation of the optical properties of the surface. Waves increase the angle of incidence of direct

radiation for high solar elevations. The effect on the Fresnel reflectance, however, is of little consequence since the reflectance does not vary much with solar zenith angle for zenith angles less than  $40^{\circ}$  (see Figure 3). Waves reduce the angle of incidence of direct radiation from a low sun, greatly reducing the reflectance of the water surface. Cox and Munk [7] have shown that wave action becomes a significant factor for solar elevations below  $20^{\circ}$ . At these low sun elevations, reduced reflection, shadowing and multiple reflections greatly reduce the reflected radiance.

The reflection of diffuse radiation by the water surface is little affected by surface roughness [2], although complete agreement on this matter is lacking [8]. Burt [9] found that the albedo of a wind roughened water surface was slightly less than the albedo of a smooth water surface -- a decrease from 6.6% to 5.7% for the roughened surface. Cox and Munk [7] measured a small increase in the albedo of a smooth water surface of 5% to 5.5% for a water surface roughened by waves. Kondratyev [8] on the other hand, calculates that where the solar zenith angle is 0° the albedo of calm water surface of 2.1% will increase to 13.1%. When the solar zenith angle is  $30^\circ$ , the increase will be from 2.2% to 3.8%, and for a solar zenith angle of  $60^{\circ}$ there will be a decrease from 6.2% down to 2.4% for a roughened surface. Plass et al. [10], using a Monte Carlo model of the atmosphere ocean system, demonstrate that the downward flux just below the surface always increases with wind speed, even at high sun elevations. They attribute this result to the fact that more sky radiance near the horizon enters the water when waves are present.

The effect of waves on the radiance of the water surface can be calculated if the probability distribution of surface slopes is known. For an observer looking down on a water surface, the specular angle will vary from place to place over the surface of the water. Since in most remote sensing applications

the light source (Sun) and observer (sensor) are high enough above the surface and the region viewed sufficiently small that variation in the specular angle can be neglected. The radiance of the surface is then directly proportional to the probability of finding a surface element with slope,  $S_0$ , at the specular angle [11]. If p is this probability, the radiance of the surface,  $L_e$ , at vertical angle i is given by

$$L_{e}(i) = L_{0}(i)\rho_{f}(i)p$$
(3)

where

 $L_0(i)$  is radiance incident at the surface at vertical angle i, and

 $\rho_f(i)$  is the Fresnel reflectance at vertical angle i.

Duntley [12] and Cox and Munk [7,13] have studied the statistical distribution of wave slope as a function of wind speed. Observations of the effect of wind speed on spatially or temporally averaged reflectance of the water surface indicate that it is not significant for view angles less than  $70^{\circ}$  from vertical. Angles in excess of  $50^{\circ}$  from the vertical are seldom used in remote sensing systems because of the large optical air mass at these angles.

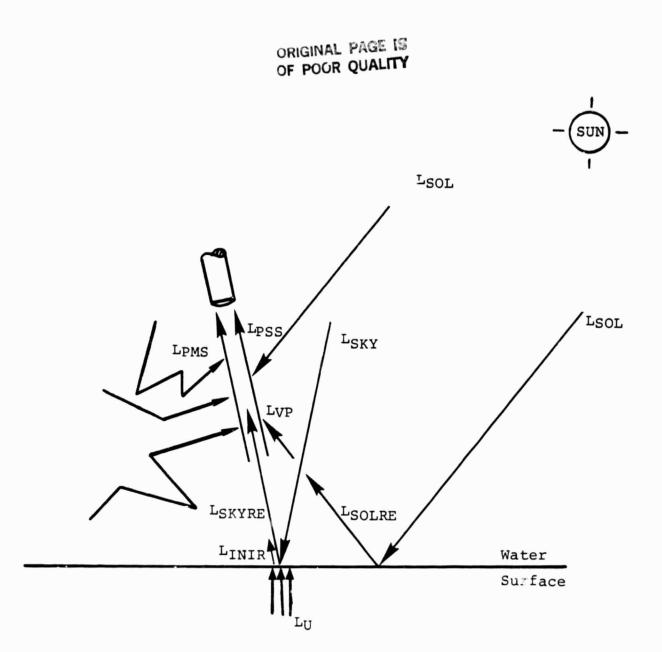
# COMPONENTS OF REFLECTED AND PATH RADIANCE IN REMOTE SENSING OVER WATER

In many applications of oceanographic remote sensing the quantity of greatest interest is the radiance information transmitted from below the water surface to a sensor, sometimes called the intrinsic radiance. To determine this guantity from raw remote sensing data we must not only estimate atmospheric path radiance but also the magnitude of radiance reflected off the water surface and transmitted to the sensor. In this chapter we describe in detail analytical models appropriate for estimating the following quantities: reflected sky radiance, LSKYRE; singly scattered reflected solar path radiance, Lyp; singly scattered path radiance, Lpss; and multiply scattered path radiance, Lpms. The first two of these quantities are radiances resulting from specular reflection off the water surface, the latter two are atmospheric path radiances. Each of these radiances augments the radiance detected by a sensor, masking the radiance signal from beneath the water surface, as shown in Figure 5.

#### 4.1 REFLECTED SKY RADIANCE, LSKYRE

The geometry for sky radiance reflected into the line of sight of the sensor is depicted in Figure 6, where  $\theta$  represents the nadir view angle of the sensor,  $\phi$  the azimuth of the sensor scanning plane. We assume a plane parallel uniform atmosphere. Sky radiance downwelling in the scan plane and incident at an angle  $\theta$  with the normal to the surface is reflected in the direction of the sensor, and attenuated by the atmosphere as it travels to the sensor. The surface reflectance is given by the Fresnel formula for unpolarized light described in Chapter 3.

We consider in this model only singly scattered sky radiance, generated by scattering of solar beam radiation along the straight



LSOL	=	direct beam radiance;
LPMS	=	multiply scattered path radiance;
L <sub>PSS</sub>	=	singly scattered path radiance;
LSKY	5	sky radiance;
LSKYRE	=	reflected sky radiance;
LSOLRE	=	reflected solar radiance;
LU	=	upward scattered radiance beneath water surface;
L <sub>INIR</sub>	=	radiance scattered from beneath water surface.
ETCUDE	F	COMPONENTS OF TOTAL PARTANCE DETERMED BY CENCOR
FIGURE	5.	COMPONENTS OF TOTAL RADIANCE DETECTED BY SENSOR

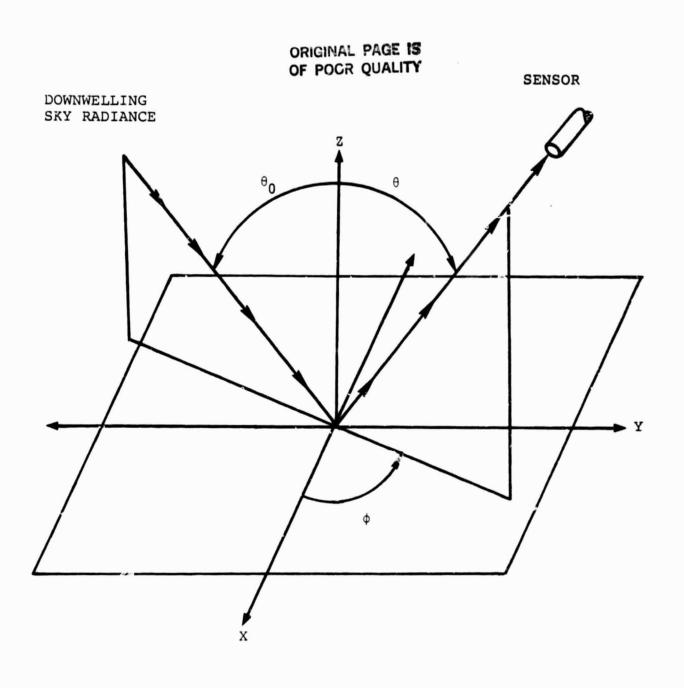


FIGURE 6. GEOMETRY FOR REFLECTED SKY RADIANCE COMPUTATION.

line path from the top of the atmosphere to the water surface. In Figure 6, this straight line path has direction  $(-\mu, \phi)$ , where  $\mu = \cos\theta$ . In order to define the angle of scattering we must first define we vectors, one defining the direction of a photon leaving the Sun and the other the direction of the singly scattered sky radiation. If  $(-\mu_0, \phi_0)$  is the direction of the photon leaving the Sun (where  $\mu_0$  is the cosine of the solar zenith angle,  $\theta_0$ , and  $\phi_0$  is the photon azimuth), the vector direction of the photon leaving the Sun is

$$\hat{\mathbf{L}}_{\text{SOL}} = \begin{bmatrix} \sin(\pi - \theta_0) \cos\phi_0 \\ \sin(\pi - \theta_0) \sin\phi_0 \\ \cos(\pi - \theta_0) \end{bmatrix} = \begin{bmatrix} \sqrt{1 - \mu_0^2} \cos\phi_0 \\ \sqrt{1 - \mu_0^2} \sin\phi_0 \\ -\mu_0 \end{bmatrix}, \quad (4)$$

and

$$\hat{\mathbf{L}}_{SKY} = \begin{bmatrix} \sin(\pi-\theta)\cos\phi \\ \sin(\pi-\theta)\sin\phi \\ \cos(\pi-\theta) \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu^2}\cos\phi \\ \sqrt{1-\mu^2}\sin\phi \\ -\mu \end{bmatrix}.$$
 (5)

The cosine of the scattering angle,  $\lambda_{SKY}$ , is given by the dot product  $I_{iSOL}$  ·  $L_{SKY}$ , i.e.

$$\cos \chi_{SKY} = \mu \mu_0 + \sqrt{1-\mu^2} \sqrt{1-\mu_0^2} \cos(\phi-\phi_0).$$
 (6)

The equation for singly-scattered sky radiance, L<sub>SKYSS</sub>, is the well-known formula [14]

$$L_{SKYSS} = \frac{\omega_0^{\mu} 0^E 0^{P(\chi_{SKY})}}{4\pi (\mu_0 - \mu)} \left[ e^{-\tau_0/\mu_0} - e^{-\tau_0/\mu} \right]$$
(7)

where 
$$E_0$$
 = solar irradiance at the top of the atmosphere;  
 $\omega_0$  = atmospheric single scattering albedo;  
 $\tau_0$  = optical thickness of atmosphere;  
 $P(\chi_{SKY})$  = scattering phase function for scattering  
angle  $\chi_{SKY}$ .

When the sky radiance is reflected off the surface of the water it is diminished by the Fresnel reflectance of the water surface,  $\rho_{\rm F}$ , and further attenuated by the atmosphere on its way to the sensor. Thus, the complete formula for the reflected sky radiance is

$$L_{SKYRE} = \rho_F e^{-(\tau_0 - \tau)/\mu} L_{SKYSS}$$
(8)

where  $\mu$  is the cosine of the scan angle,  $\tau$  is the optical depth of the sensor, and  $e^{-(\tau_0 - \tau)/\mu}$  is the transmittance of the atmosphere between the water surface and the sensor. At this point we take note of the fact that in the above formula for the singly scattered reflected sky radiance, the reflectance of the water surface and the transmittance of the atmosphere are opposing effects. Assuming the refractive index of water to be 1.33, the

Fresnel reflectance of water at the nadir view angle reaches a minimum of 0.021 and attains a maximum of 1.0 at the grazing angle. Typical values of  $\rho_F$  for angles commonly used in remote sensing range between 2.1 at the nadir to 3.0 at a scan angle of 46°. The transmittance, on the other hand, reaches a maximum at the nadir and becomes increasingly small as the scan angle increases. These effects will be discussed further in the following chapter.

4.2 SINGLY SCATTERED REFLECTED SOLAR RADIANCE, Lyp.

The phenomenon of specular reflection produces an image of the radiation source on the surface of the water. We refer to the image of the Sun on the water surface as the virtual Sun. If the scan plane is coincident with the solar plane and the sensor is scanning on the solar side of the scan plane at a view angle equal to the solar zenith angle, the field of view becomes saturated with the radiance of the Sun's image. Radiance from the virtual Sun is also scattered into the line of sight of the sensor. In the terminology of this report, we refer to singly-scattered path radiance from the Sun as virtual Sun path radiance,  $L_{\rm Up}$ .

To find the scattering angle for the computation of singlyscattered virtual Sun path radiance, we note that the zenith angle of a photon leaving the virtual Sun is  $\theta_{SUN}$  and the azimuth angle is  $\phi_0 = \phi_{SUN} + \pi$  (see Figure 7). The vector direction of a photon leaving the virtual Sun is

$$\hat{\mathbf{L}}_{\text{VSOL}} = \begin{bmatrix} \sin\theta_{\text{SUN}} & \cos\phi_{0} \\ \sin\theta_{\text{SUN}} & \sin\phi_{0} \\ \cos\theta_{\text{SUN}} \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu_{0}^{2}} & \cos\phi_{0} \\ \sqrt{1-\mu_{0}^{2}} & \sin\phi_{0} \\ \mu_{0} \end{bmatrix}$$
(9)

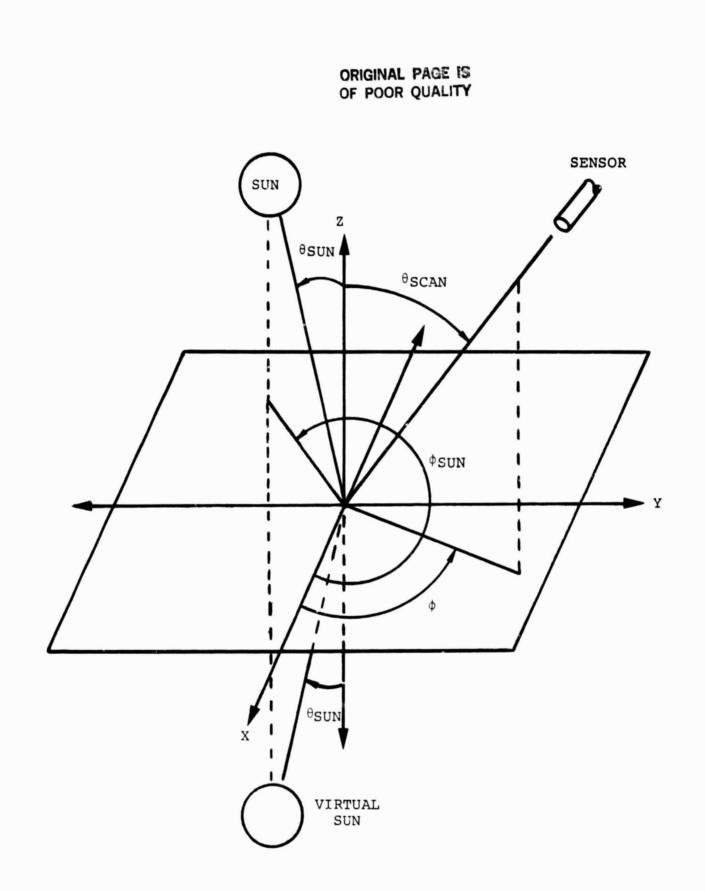


FIGURE 7. GEOMETRY FOR VIRTUAL SUN PATH RADIANCE CALCULATION.

and the vector direction into which photons from the virtual Sun are scattered, creating virtual Sun path radiance is

$$L_{VP} = \begin{bmatrix} \sin\theta\cos\phi\\ \sin\theta\sin\phi\\ \cos\theta \end{bmatrix} = \begin{bmatrix} \sqrt{1-\mu^2} & \cos\phi\\ \sqrt{1-\mu^2} & \sin\phi\\ \mu \end{bmatrix} .$$
(10)

The scattering angle is the dot product,

$$L_{VSOL} \cdot L_{VP} = \mu \mu_0 + \sqrt{1 - \mu^2} \sqrt{1 - \mu_0^2} \cos(\phi - \phi_0).$$
(11)

As in the case of singly scattered sky radiance, the same types of physical interactions which generate singly-scattered sky radiance from direct solar radiation also scatter radiation from the virtual Sun to generate virtual Sun path radiance. Thus, we may use the same equation for singly-scattered sky radiance, with some modifications, to find the singly-scattered virtual Sun path radiance,  $L_{VP}$ . One difference is in that the computation of  $L_{VP}$  we will now sum the scattered radiation over a path beginning at  $\tau = \tau_0$  (the optical depth of the scene viewed by the sensor) and ending at the optical depth of the sensor,  $\tau$ . If we denote the irradiance of the virtual Sun by  $E_0$ , we obtain the following formula for singly-scattered path radiance from the virtual Sun:

$$L_{VP} = \frac{\omega_0 \mu_0 E_0 p(\cos \chi_{VP})}{4\pi (\mu_0 - \mu)} \left[ e^{-(\tau_0 - \tau)/\mu_0} - e^{-(\tau_0 - \tau)/\mu} \right].$$
(12)

We define  $E_0$  by noting that the  $E_0$  is the image of the Sun reflected in the water surface. Hence, the irradiance of the virtual Sun is the irradiance of the true Sun at the top of the atmosphere attenuated by the atmosphere and the reflectance of the water surface, i.e.,

$$E_0' = \rho_F e^{-\tau_0/\mu_0} E_0$$
 (13)

where  $\rho_F$  is the Fresnel reflectance and  $e^{-\tau_0/\mu_0}$  is the transmittance of the atmosphere.

4.3 SINGLY SCATTERED PATH RADIANCE, LPSS.

The geometry for singly-scattered path radiance is shown in Figure 8. The formula for singly-scattered path radiance is similar to that for singly-scattered path radiance from the Sun--the same straight line path from  $\tau_0$  to  $\tau$  is used, but  $E_0$ is substituted for  $E_0$  in the formula. The cosine of the scattering angle for singly-scattered path radiance,  $\cos\chi_{PSS}$  is also the negative of the cosine of the scattering angle used to compute the phase function for  $L_{VP}$ . Thus,  $\cos\chi$  is the dot product,  $\hat{L}_{SOL} \cdot \hat{L}_{VP}$ , vector directions which have already been defined. The formula used to compute singly-scattered solar path radiance is

$$L_{PSS} = \frac{\omega_0^{\mu} 0^E 0^{\mu} 0^E (\cos \chi_{PSS})}{4\pi (\mu + \mu_0)} e^{-\tau_0 / \mu_0} \left[ e^{(\tau_0 - \tau) / \mu_0} - e^{-(\tau_0 - \tau) / \mu} \right]$$

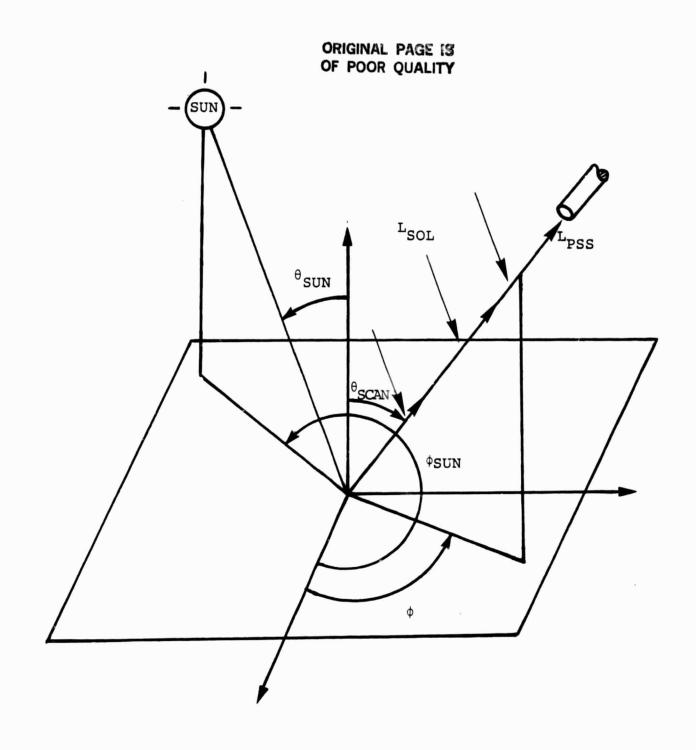


FIGURE 8. GEOMETRY FOR SINGLY SCATTERED PATH RADIANCE, L<sub>PSS</sub>, CALCULATION.

where the variables in the above equation are as previously defined.

## 4.4 MULTIPLY SCATTERED PATH RADIANCE, L<sub>PMS</sub>.

To compute multiply scattered path radiance we use an analytical approximation described in detail in an earlier report [15]. The formula for the computation of  $L_{PMS}$  is

$$\begin{split} \mathbf{L}_{\mathbf{PMS}} &= \frac{\mathbf{E}_{\mathbf{0}}}{4\pi \left[ \mu_{\mathbf{0}} + (\mathbf{1} - \mathbf{\eta}) \tau_{\mathbf{0}} \right]} \\ & \left( \left\{ (1 - \eta) \tau_{\mathbf{0}} \left[ \mathbf{p}(\mu, \phi, \mu_{\mathbf{0}}, \pi + \phi_{\mathbf{0}}) + \mathbf{p}(\mu, \phi, -\mu_{\mathbf{0}} \phi_{\mathbf{0}}) \right] + \mu_{\mathbf{0}} \mathbf{p}(\mu, \phi, -\mu_{\mathbf{0}} \phi_{\mathbf{0}}) \right. \\ & + \frac{2\mu_{\mathbf{0}}^{2} \rho}{1 + 2 (1 - \eta) (1 - \rho) \tau_{\mathbf{0}}} \right\} \left[ 1 - \mathbf{e}^{-(\tau_{\mathbf{0}} - \tau) / \mu} \right] + \left\{ (1 - \eta) \mathbf{p}(\mu, \phi, \mu_{\mathbf{0}}, \pi + \phi_{\mathbf{0}}) \right. \\ & + \mathbf{p}(\mu, \phi, -\mu_{\mathbf{0}}, \phi_{\mathbf{0}}) \right] - \frac{8 (1 - \eta) \mu_{\mathbf{0}}^{2} \rho}{1 + 2 (1 - \eta) (1 - \rho) \tau_{\mathbf{0}}} \right\} \left[ (\tau_{\mathbf{0}} + \mu) \mathbf{e}^{-(\tau_{\mathbf{0}} - \tau) / \mu} - (\tau + \mu) \right] \right) \end{split}$$

$$\tag{15}$$

The single-scattering phase functions are given by:

$$p(\mu,\phi,\mu_{0},\pi+\phi_{0}) = p\left[\mu \mu_{0} - \sqrt{(1-\mu^{2})(1-\mu_{0}^{2})} \cos(\phi-\phi_{0})\right]$$

$$p(\mu,\phi,-\mu_{0},\phi_{0}) = p\left[-\mu \ \mu_{0} + \sqrt{(1-\mu^{2})(1-\mu_{0}^{2})} \ \cos{(\phi-\phi_{0})}\right].$$

#### COMPUTER MODEL AND RESULTS

In this chapter we describe implementation of the formulas discussed in Chapter 4 in the computer program called ATCOR. Many of the details of ATCOR have been presented in a previous report [16], so only a brief description of the entire program will be given here. Our discussion will focus primarily on the subroutine ATMSFR, in which the formulas presented in Chapter 4 have been introduced.

#### 5.1 SPECIFICATION OF SOLAR AND SENSOR GEOMETRY

The geometric relationship of the sensor to the environment is shown in Figure 9. The geographic coordinates of the sensor locate the center of a spherical coordinate system used to define angles needed in model calculations. In the diagram the scanner scans along a path from  $P_1$  (the first pixel) to  $P_n$  (the last pixel). The azimuth of the scan plane is measured in the counterclockwise direction from north to the first pixel and is read into the program by the routine ATMSFR. The first pixel is always  $90^{\circ}$  in a clockwise direction from a vector pointing in the direction of the flight.

The solar zenith and azimuth angles are computed automatically once the latitude, longitude, date, time of day (standard time), and zone number are specified. The extraterrestrial solar irradiance is also computed based on these inputs.

## 5.2 ATMOSPHERIC CORRECTION OPTIONS

Two input parameters set by the user determine which calculations are performed in routine ATMSFR. These parameters are SCATT and OPTION. SCATT may assume the value of either 0 or 1; OPTION can take on the values of 1, 2, or 3. If SCATT is 0, only multiply scattered path radiance is calculated and the value of OPTION can be any integer and will be ignored since only LPMS is then calculated.

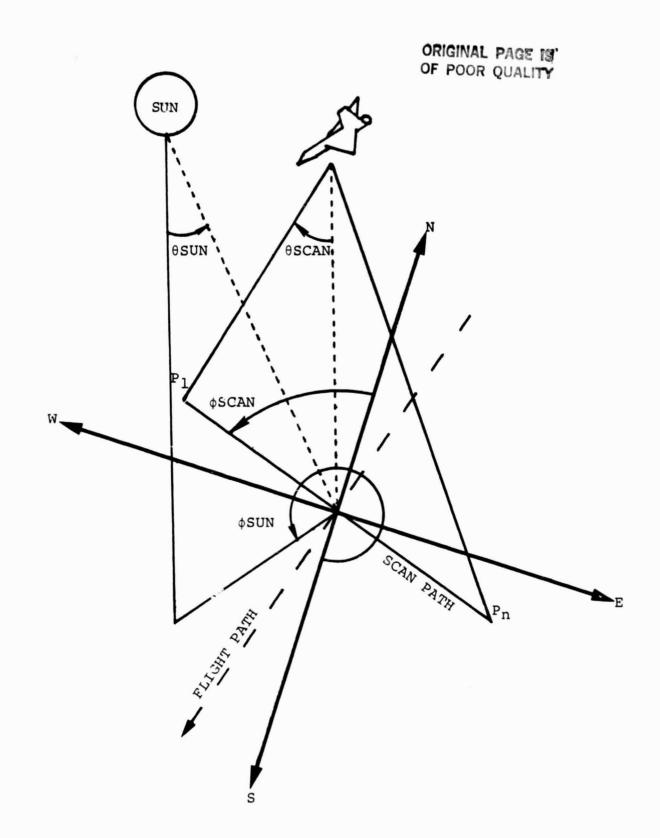


FIGURE 9. SCAN PLANE AND SOLAR GEOMETRY.  $\phi$ SCAN IS MEASURED COUNTERCLOCKWISE FROM NORTH TO PIXEL #1,  $\phi$ SUN IS MEASURED COUNTERCLOCKWISE FROM SOUTH. POSITIVE SCAN ANGLES ARE MEASURED FROM THE FIRST PIXEL TO THE NADIR. NEGATIVE SCAN ANGLES ARE MEASURED FROM THE NADIR TO THE LAST PIXEL. P1 AND Pn ARE THE FIRST AND LAST PIXELS, RESPECTIVELY.

If SCATT is 1, single scattering computations are performed and the value of OPTION is used to determine which values to calculate. If OPTION = 1, singly scattered path radiance, LPSS, and reflected sky radiance, LSKYRE, are computed. If OPTION = 2, LPSS and virtual sun path radiance, LVP, are computed. If OPTION = 3, LPSS, LSKYRE and LVP are calculated.

OPTION and SCATT are read into the data file on logical unit 4. The format for this record is (515), and the variables read in are:

FSTP LSTP PTINC SCATT OPTION

where

FSTP = the number of the first pixel to be processed, LSTP = the number of the last pixel to be processed, PTINC = the pixel increment to use in the processing, and SCATT and OPTION are as previously defined.

The output file which is used by ATMSFR is given in Table 1 and the new subroutine ATMSFR2 is given in Table 2.

# 5.3 MODEL INPUT PARAMETERS

In addition to the geometric parameters, we must specify parameters characterizing the medium and the measurement system.

The model makes use of several "altitude" values which must be input by the user. First, one must know the actual altitude (km) of the sensor above the surface. Second, one must know the pressure (millibars) of the atmosphere at the surface, and third, one must know the atmospheric pressure (in millibars) at flight altitude. If only the altitudes are known, one can use the tables relating pressure to altitude as given by the U.S. Standard Atmosphere [17].

# TABLE 1INPUT FILE ON LOGICAL UNIT 4FOR USE BY ATMSFR

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LINE NO	READ OCCURS IN ROUTINE	INPUT VARIABLES	FORMAT
1	ATMSFR	FSTP,LSTP,PTINC, SCATT,OPTION	(515)
2	DATE	MNTH, DAY, YEAR	(314)
3	ANGLES	HOUR, MIN, SEC	(2I5,F6.3)
4	ANGLES	NZ	(15)
5	ANGLES	LATD, LATM, LATS, LONDG, LONGM, LONGS	(215,F6.3,215, F6.3)
6	ATMSFR	ZSCAN, ZGRND, LSW	(2F8.5,15)
7	ATMSFR	(WAVE(I),1=1,QNCHAN)	(10F8.5)
8	ATMSFR	(RHO(I),I=1.QNCHAN)	(10F8.5)
9	PHASE	R,IM	(2F8.6)
10 11-14	PHASE	NWT, NANG, (C(I), I=1, NANG)	(215/(10F8.6))
15-79	PHASE	(WTAB(I),(PF(I,J), J=1,NANG),I=1,NWT)	(F10.6/(10 F8.4))
80	RAYLEI	PRESO, PRESZ	(2F10.4)
81	OZONE	NOZ, NPROF, NO3W1, NO3W2	(415)
82	OZONE	WAVC1,WAVC2	(2F8.4)
83-153	OZONE	((ZOZ(IZ),O3INT(IZ,IP), IZ=1,NOZ),IP=1,NPROF)	(F7.0, E11.4)
154	OZONE	(O3MAX(IP),IP=1,NPROF)	(10E13.6)
155-169	OZONE	(WAVO3(I),A(I),I=1,NO3W) (NOTE:NO3W=NO3W1 + NO3W2	

# INPUT FILE ON LOGICAL UNIT 4 FOR USE BY ATMSFR

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# (CONTINUED)

LINE NO.	READ OCCURS IN ROUTINE	INPUT VARIABLES	FORMAT
170	OZONE	NOP	(415)
171	THICK	NTEX	(12)
177	THICK	(WAVEX(I), TAUEX(I), I=1, NTEX)	(2F8.4)
178	PARAMS	FSCAT	(10F8.6)
179	AERO	NAER, MPROF, NUZ, MAXG	(515)
180	AERO	IPROF	(15)
181-184	AERO	(WAVAER(I),RIN(I),I=1,NAER)	(10F8.4)
185-230	AERO	((ZUN(IZ),UNIZ(IZ,IP),IZ=1, NUZ),IP=1,MPROF)	(F7.0,E11.4)
231-258	AERO	(X(I),Z(I),I=1,MAXG)	(F7.0,Ell.4)
259	ATMSFR	PHID, PHIM, PHIS	(213,F6.3)
		\$ENDFILE	

2 TABLE

PROGRAM LISTING

5 9 8 C:

CORRECT FOR ATMOSPHERIC EFFECTS UNIQUE TO LARGE BODIES OF WATER. IT Will Perform Identically as 115 previous versions when the value zero ( or blanks ) are read into the integer variable "scatt." Thus any VER. MA3. 3 5JANB2 -- ADDED SINGLE SCATTER, LPATH, LRSKY, LVIRP ...... ROUTINES AS ITS PREDECESSOR VERSIONS AND IS CALLED IN THE SAME MANNER +VERSION IILINE/QLINE COMMON BLOCKS VEP. NAS. 2 15AUG77 RHH--NASA UNIVAC WHEN SCATT-1, HOWEVER, ATMSFR2 WILL CALCULATE ONLY SINGLY SCATTERED QUANTITIES. ANOTHER NEW VARIABLE, "OPTION," IS READ IN. WHEN SCATT-O 0P110N-1 IMAGE" ) OF THE SUN FROM THE WATE. IS SUBTRACTED. WHEN OPTION-3, ALL THREE QUANTITIES ARE SUBTRACTED. ATMSFR2 CALLS ALL AND OMLY THE SAME 6-UN77 -- FINISH MAKING FORTRAN COMPATIBLE, QBREAL OPTION=2. THE DIRECT PATH RADIANCE AS WELL AS THE PATH RADIANCE PRO-FLECTED FROM THE WATER SURFACE INTO THE SENSOR ARE SUBTRACTED. WHEN VER. 3.0 14APR77 -- 200/24 CHANNEL VERSION ALL PUT TOGETHER MODIFICATIONS PROGRAMMED BY JIM LIERZER UN'S QANCIL, QIOPAR COMMON BLOCKS ADDED JAN. 5, 1982 PREVIOUS DATA SET MAY BE USED WITH THIS ROUTINE TO PRODUCE THE SAME LVIRP - PATH RADIANCE GENERATED BY THE VIRTUAL IMAGE OF THE SUN OB-DUCEU BY THE SPECIALARLY REFLECTED > STUAL IMAGE ( I.E. THE "MIRROR THIS IS A MODIFIED VERSION OF THE ROUTINE "ATMSFR" TO BE USED TO RADIATION IS SUBTRACTED FROM THE EXPERIMENTAL RADIANCES. IF OPTION Both the direct path radiance and the SKY radiance specularly re-ONLY THE PATH RADIANCE PRODUCED BY THE DIRECT SOLAR OPTION HAS NO EFFECT. WHEN SCATT=1, OPTION CHOOSES WHICH RADIANCE VER. 3.2 12AUG77 -- MO02, NOT 1.0; QPOLY; QM1,2; Q. . 889; LPATH - PATH RADIANCE GENERATED BY THE DIRECT SOLAR RADIATION VER. 3. OD 30JUN77 -- QIBSQL; /QANCIL/ IN ALL ROUTINES 15AUG77 -- DIMENSIONS FOR NASA-LEWIS UNIVAC LRSKY - REFLECTED SKY RADIANCE OBSERVED AT SENSOR INTEGER FUNCTION ATMSFR(L, CONTRL, MCHAN, MSS) MULTIPLY SCATTERED PATH RADIANCE COMRECTIONS. QUANG, QUANG ATMSFR2 WE DEFINE THE FOLLOWING VARIABLES: SERVED AT THE SENSOR CONTRABUTIONS TO CORRECT FOR ; ł VER 3. OC 10JUN77 TUUUTT VER. NA3.2 VER. 3.06 VER. 3. 0A IF OPTION=0. EIC VERSION. OBSEQL. SWI TCH ٠ C U U c 5 2222 26 239 42 53 

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QCVERS, OLINE, QSLINE, OMSAS, OMSA, OMSE, OKS, OMA, OMB, OKP, OMSS OFRSTL, QLASTL, QSTEPL, OFRSTP, QLASTP, QSTEPP QSTEP, OMCHAN, QMSSS, QMCHAN, OMDACH, QFILE, QSPARE, QSCHAN QBNOME, QBEMD, QBPAST, QBPOLY, QBLSEQ, QBTRSE, QBSEQL QBNDNE, QBEND, QBPAST, QBPOLY, QBTRSE, QFLDID(24), QSPARE (B3), CBLSEG, QBEEQB, QBSECL, QBANG, QDANG, QBREAL, QBFACT, 0111LE (120), 0111L2 (120), 0L1ST (6, 20), QSCHAN(200), 0FACTM (20C), QFAC, A (200), QCMEDED QFACTM, QFACTA, QBANG, QDANG "MAIN COMMON BLOCK FOR USE BY MONITOR, IMPUT FSR, MUDULES, AND DUTPUT LOGICAL OFLDID.OTITLE.OTITL2.OLIST EQUIVALENCE (ONSAS(1).OFRSTL.ONSA), (ONSAS(2).OLASTL.ONSB) EQUIVALENCE (ONSAS(2).OSTEPL.OKS), (ONSAS(4).OFRSTP.ONA) EQUIVALENCE (ONSAS(5).QLASTP.OMB), (ONSAS(6).OSTEPP.OKP) 00CHAN, 00CHAN, 00HSS, 00HSS2, 00I BUF, 000BUF 00DILN, 00MX2, 00MX4, 00LDTA, 00LDT2, 00LCTL 0UDATA, 0UPTRY, 0UPRNT, 0UERR, 0UPPNCH, 0UPPNT, 0UMAP C+ANCILLARY DATA COMMON ALOCK COMMON /QANCIL/ QHDNUM,QANEAD(200),QNCIL(180),QAN999 / QCOM/ QCVERS(4), QMSAS(6), QMSS, QLINE, QSLINE, QSTEP, QMCHAN, QMSSS, QMCHAN, QMDACH, QFILE, C+PARAMETERS -- MACHINE DEFENDENT LENGTH ATTRIBUTES, ETC C+PARAMETERS -- DECLARATION SIZES FOR DATA ARRAYS C ... PANAMETERS . FOR Use IN FORTRAN PROGRAMS GOVRTX, GOMREG, GOMSEG CHEMAIM, GAHE AD, CINCIL QANOOO, QANIB99 METERS -- STANDARD 1/0 UNITS EQUIVALENCE (QCMODO, QCVERS) EQUIVALENCE (QANOOO, QHDMSH) QCM000, QCM999 QOCVER(4) QONSS2/824/ QOIBUF/22293/ 0008UF / 22293/ 00VR1 X / 101/ QONSS /3264/ DATA QUCHAN/200/ 00CH042/24/ DATA QUNSEG/24/ QOMRE G/24/ 00LDTA/4/ 00LDT2/4/ DATA QODILN/4/ QONK2 /2/ QONIXA /4/ DATA QOLCTL/4/ QUDATA/4/ QUPRMT/6/ QUERR /6/ QUIRTRY/5 CUPNCH/7 QUPRNT/8 INTEGER INTEGER INTEGER NIEGER LOGICAL LOGICAL INTEGER INTEGER INTEGER INTEGER INTEGER DGICAL COMMON • ..... ----\*\*\*\* DATA DATAD DATA DATA DATA DATA DATA REAL E Su C-PAR • U U U υ U U 0100 28 11 60 -15 8 20 5

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C QSTEP=1 -- CONSTRUCT QLINE LINKAGE TO PRIOR MODULE SUPPLYING DATA: 100 FDRMAT('OATWSFR: ENTER FIRST PIXEL, LAST PIXEL, PIXEL INCREMENT,' 1 .' Scatt, and option (515)') CALL ANGLES(QUDATA,QUPRNJ,MDAY, PHIO,THETAO) Read height of scannes, height of ground, print/write tape suitch CALL QOVERS(VERMES, QOCVER) DON'T FORGET TO REMOVE THE "C" IN COLUMN 1 OF THE ABOVE COMMENT WHEN OSTEP-3 -- INITIALIZE; PRINT MODULE IDENTIFICATION, READ CONTROL SO THAT TAS(I) IS THE AEROSOL OPTICAL THICKNESS DUE ONLY TO SCAT-AND TO CHECK THAT THE COMMON BLOCK VERSION "QOCVER" IS CURRENT CALL QOVERS TO PRINT MODULE NAME, DESCRIPTION, VERSION, DATE, NOTE THAT WHEN YOU ENCOUNTER THE CALLS TO FUNCTION PF. VIZ I HAVE DELETED THE CALL TO QQVERS IN ORDER TO RUN ATMSFR Pf(ARGUMENT.I.TR(I).TAS(I))
THE FOURTH ARGUMENT SHOULD BE TAS(I). WHERE
TAS(I)-FSCAT(I)+TA(I) OPTION-3 FOR LPATH+LRSKY+LVIRP . ') SCATT-0 FOR MULTISCATTERING, SCATT-1 FOR SINGLE SCATTERING, OPTION-1 FOR LPATH+LRSKY. READ ( QUDATA, 105) FSTP, LSTP, PTINC, SCATT, OPTION GO TO (1.99.3.99.99.6.99.8.99.99.99), QSTEP OPTION-2 FOR LPATH+LVIRP CALL QQCALD(PRIDR, DATA, CONTRL, QMCHAN, QMSS) OPTION-O FOR LPATH ONLY. JIM LIERZER COMPUTE WAVELENGTH DEPENDENT CONSTANTS READ(QUDATA, 263) ZSCAN, ZGRND, LSW CALL DATE (QUDATA, QUPRNT, NDAY) • OSTEP=6 -- BEGINNING OF REGION IF (QBNOME OR QBPASI)RETURN INTERFACING WITH "ATCOR ALONE FOR DEBUGGING PURPOSES. CON - 7.9577472E-02 WRITE (OUPRMT, 100) ZSC-ZSCAN+3280.833 WHERE WRITE (QUPRMT, 101) FORMAT(2F8.5,15) PRIOR - L(1,1) 105 FORMAT(515) ATMSFR - 0 CONT INUE CONT INUE IOI FORMAT( RETURN RETURN **PROGRAMMER** RETURN TERING C DATA: 263 -• 66 G J U J c ........ 000 U c c S U J U C 200 221 193 202 204 210 181 89 6 561 196 197 66 206 208 209 225 229 530 233 239 68 184 185 186 181 188 161 192 207 212 213 214 215 216 217 218 219 222 223 224 226 227 228 162 232 234 235 236 237 238 Ξ

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(Cont.) 2 TABLE

PHI = 1.74532936-02\*PHID + 2.90888216-04\*PHIM + 4.84813686-06\*PHIS FPH = SQRT(1.0 - MU0\*MU0)\*COS(PHI-PHIO) FORMAT( '0', 10X, 'ALL RADIANCES ARE IN UNITS OF MILLIWATTS PER CEN'. 561 FORMAT(' ', 10X, 'SWITCH POSITION', 3X, 11, 5X, 'SCATTERING MODE', 3X, 11, 552, 'OPTION', 3X, 11) READ (QUDATA,265) PHID,PHIM,PHIS WFITE (QUPRNT,270) PHID,PHIM,PHIS 270 F'JRMAT('O',19X,'AZIMUTH ANGLE MEASURED COUNTERCLOCKWISE FROM MORT' ', 'H TO FIRST PIXEL'/'',20X,'OF SCAN PLANE IS ',13, 2 ', DEGREES,',13,' MIMJTES,',F5.1,' SECONDS.') READ THE AZIMUTH ANGLE COUNTERCLOCKWISE FROM NORTH TO ( THE FIRST CALL RAYLEI(QUDATA,ONCHAN,WAVE, TR.TK2) Call Dzone(Qudata,Ouprnf,Onchan,Wave,Z, Tauda,Tauzda) Call Thick(Qudata,Quprnf,Onchan,Wave, Tau) Call Params(Qudata,Quprnf,Tr,Tauda,Tau,QNCHAN, ETA,OM,TAS,FSCAT) Call Aero(Qudata,Quprnf,QNCHAN,Wave,Z,TAU,TR,TAUDA, TAZ,TA) CALCULATE ANGULAR QUANTITIES (DEPENDENT ON AZIMUTH ANGLE ONLY) ECHO A LOT OF DATA TO TELL THE USER THE INPUT DATA WHICH GIVE WRITE(QUPRNT,300)FRESUN FORMAT(' ',19X,'THE FRESNEL REFLECTANCE FOR THE DIRECT SOLAR 'RADIATION IS',F10.5) FRESUN=0.5+((SINA+SINA)/(SINB+SINB)+(TANA+TANA)/(TANB+TANB)) READ SCANNER WAVELENGTHS, ESTIMATED BACKGROUND REFLECTANCES READ (QUDATA.264) (WAVE(1),1-1,0NCHAN) READ (QUDATA.264) (RHO(1),1-1,0NCHAN) FIND THE FRESNEL REFLECTANCE OF THE DIRECT SOLAR RADIATION FORMAT(' ', 19X,'SENSOR ALTITUDE',F10.1,' FEET') Z = ZSCAN-ZGRND IF (SCATT.EQ.O.OR.OPTION.LT.2) GO TO 301 CALL PHASE (QUDATA, QUPRNT, QNCHAN, WAVE ) REFRAC-APSIN(SIN(THETTA)+1./INDEX) WRITE (OUPRNT, 561)LSW, SCATT, OPTION CALL SOLAR(WAVE, QNCHAN, NDAY, EO) PIXEL OF THE ) SCAN PLANE WHENEVER LVIRP IS DESIRED HIM THE REST OF THE NUMBERS WRITE (QUPRNT, 253) 25C IF (THETAO)39.43.39 THETTA-ABS(THETAO) MUO - COS(THETAO) FORMAT(213, F6.3) AX=THETTA-REFRAC BX=THETTA+REFRAC WRITE (OUPRNT, 30) OUM-OUM - OSOUM FORMAT( 10F8.5) FRE SUN=FRE SO 2. • MUO SINA-SIN(AX) SINB=SIN(BX) TANA - TAN(AX) TANB-TAN(BX GO VO 299 300 FORMAT ( 301 CONTINUE c02 = 299 253 43 90 66 265 264 υ 00000 0000 υu υu 242 243 245 245 245 245 249 255 255 255 255 268 269 270 2552 2554 2555 2555 2556 2556 2557 2558 2559 2559 275 281 282 283 284 285 285 285 285 287 260 262 263 264 265 266 267 272 273 274 276 278 279 288 289 241 261 271 277

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ANGULAR RE 36 WRITE(QUPRNT,37) 37 FORMAT(' ',10X,'WE SHALL CALCULATE MULTIPLY SCATTERED RADIANCES ', 'TIMETER SQUARED'/' '14X,'PER MICROMETER PER STERADIAN '/ 11X,'ALL IRRADIANCES ARE IN UNITS OF MILLIWATTS PER CENTIMETER' '' SQUARED'/15X,'PER MICROMETER '/ IF (OPTION.EQ.O) WRITE(QUPRNT.32) 32 FORMAT(' ',10X,'ONLY THE DIRECT PATH RADIANCE WILL BE CONSIDERED 33 FORMAT(' ',10X,'THE DIRECT PATH RADIANCE AND THE REFLECTED SKY 31 FORMAT(' ', 10X, 'WE SHALL CALCULATE SINGLY SCATTERED RADIANCES 34 FORMAT(' ',10X, THE DIRECT PATH RADIANCE AND THE VIRTUAL PATH 1 'RADIANCE WILL BE CONSIDERED.') 1 (OPTION.EQ.3) WRITE(OUPRNT,35) 35 FORMAT(' ', 10X, 'THE DIRECT PATH RADIANCE AND BOTH THE VIRTUAL 1 'Path Radiance and the reflected Sky Radiance ' 2 'Will be considered.') COMPUTE WAVELENGTH DEPENDENT QUANTITIES NEEDED FOR WRITE(QUPRNT, 275) QBANG, QDANG 275 FORMAT(20X, MAXIMUM SCAN ANGLE-', F9.7, ' RADIANS. 150Lution of Scanner IS ', F9.7, ' Radians.') 'RADIANCE WILL BE CONSIDERED. ') 11X, 'ALL WAVELENGTHS ARE IN MICROMETERS ') , 10F 10.5) WRITE(QUPRNT, 279)(TAU(I), I-1, QNCHAN) FORMAT('- INTERPOLATED OPTICAL', 10F10.5) ', 10F 10.5) (, 10F 10.5) FORMAT( '- SCATTERING PARAMETER', 10F 10.5) WRITE(OUPRNT, 295)(TR(I), I-1, ONCHAN) FORMAT('- RAYLEIGH OPTICAL', 4X, 10F10.5) TAUZ(I) - TAUZO3(I) + TRZ(I) + TAZ(I) WRITE (QUPRNT, 290) (TAU03(I), I-1, QNCHAN) QUPRNT, 285)(FSCAT(1), I-1, QNCHAN) WRITE(QUPRNT.277)(WAVE(I),I-1,QUCHAN) FORMAT('- WAVELENGTH ',10F10 WRITE (QUPRNT, 281) (RH0(I), I-1, QNCHAN) FORMAT ( '- DZONE OPTICAL ', 7X, 10F 10. 5) WRITE(QUPRNT, 283)(OM(I), I-1, QNCHAN) (01 101 . IF (OPTION EQ.2) WRITE (QUPRNT, 34) IF (OPTION.EQ.1) WRITE(QUPRNT, 33) WRITE (QUPRNI, 276) (I. I-1, QNCHAN) IF (SCATT.EQ.0) GO TO 500 FORMAT('- BACKGROUND ALBEDO THICKNESS') FORMAT( '- SINGLE SCATTERING THICKNESS') (FSCAT)') THICKNESS ') IF (SCATT.EQ.0) GO TO 36 WRITE(QUPRNT,31) ALBEDO') (, , NULY . ) (, 'ATNO, WRITE (QUPRNT, 284) (QUPRNT, 287) WRITE (QUPRNT, 292) WRITE (OUPRNT, 296) FORMAT('1 CHANNEL WRITE (QUPRNT. 280) DG 10 I-1. ONCHAN FORMAT( ' FORMAT( ' GO TO 38 38 CONTINUE FORMAT FORMAT FORMAT WRITE WRITE 280 290 284 287 292 295 296 281 283 285 276 279 277 υu 316 345 346 01 E 313 866 301 303 303 303 306 **806** E 312 31E 916 916 320 322 323 324 325 326 328 329 OEE IEE **332** EEE 334 **9**336 LEE 340 341 342 343 344 347 348 349 **35**0 1351 352 ESE 354 355 356 1357 358 359 321 327

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COMPUTE LAMBDA DOUBLE PRIME LAMPP(I) = CON\*NUSQ(I) \* OM(I)\*EO(I) /(B(I)\*NU\*SH2 \* MU0\*CH2) E(I) = (MUO+MUO+EO(I)/(MUO+ETOTF))+(I+(2,+RHO(I)+ETOTF)/ (1,+2,+ETOTF)) DTPP(I) = 2.0+C\* ((1.0-0M(I))•NU+SH3 + MU0+CH3) compute total irradiance at surface E(I) = 6.28318531\*PHIP(I) \* (NU+8(I)+SH3 +MU0+CH3)/ COMPUTE WAVELENGTH DEPENDENT QUANTITIES NEEDED FOR Multi scattering calculations 2.0+C+((1.0-DM(I))+NU+SH4 + MU0+CH4) CALCULATE SIMILAR CONSTANTS FOR DMEGA-1 COMPUTE DIRECT IRRADIANCE AT SURFACE COMPUTE PHI PRIME PHIP(1) = 2 0\*RHD(1)\*MUOSQ\*LAMPP(1)/ (MUO\*CH3+NU\*(B(1)-A(1)\*RHD(1))\*SH3) COMPUTE DTHER CONSTANTS FE(1)-MUO+EO(1)/E(1) COMPUTE DIRECT IRRADIANCE AT SURFACE ED(1)-MUO+EO(1)+EFACTS ED(1)-WUO+EO(1)+EXP(-TAU(1)/MUO) IF (OM(1).GT. OML) GD TO 20 A(1) = OM(1)\* (1.0-ETA(1)) B(1) = 1.0 + A(1) = OM(1) C = A(1) + B(1) Nu = MUO / SORY(C\*(B(1)-A(1))) SINGLE SCATTERING CALCULATIONS - B(I)+CH1 + MUO+SH1/NU ATP(1) - B(1) • NU • SH1 + MU0+CH1 - C+NU+SH3 + MUO+CH3 - C+NU+SHA + MU0+CH4 ARG4 - ARG3+TAU2(I)/TAU(I) ETOTF - (1.-ETA(1)) • TAU(1) SIG(I) - TAU(I) - TAUZ(I) EFACTO - EXP(SIG(I)/MUQ) CONST - CON+OM(I)+EO(I) EFACTS - EXP(-TAU(I)/MUO) SIG(1) - TAU(1) - TAU2(1) (1)0HN+(1)W0+(1)0SNW) FE(1) - MUO'EO(1)/E(1) ARG1 - SIG(1)/NU SH4 - SINH(ARG4) SHI - SINH(ARGI) (EBAR)HNIS - EHS CH4 - COSH( ARG4 ) CHI = COSH(ARGI) ARG2 - TAU(I)/NU CH2 - COSH( ARG2 ) CH3 - COSH( ARG3) NUSQ(1) - NU-NU SH2 - SINH(ARG2) ARG3 - C02+ARG2 AT(1) - NU-SHI BI(1) - CHI G0 T0 15 GO TO 15 CONTINUE BIP(1) ATPP(1 BIPP(I CIPP(I CONTINUE 500 20 U 0000 J U U c C 00 C 376 409 372 SLE 380 404 **378** 395 397 398 398 398 410 162 364 368 369 370 116 LLE 979 381 386 89 90 102 103 101 108 = 412 EI 1.4 115 116 120 365 366 367 E86 385 186 988 392 E 6E 394 118 419 384 6 1 96

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LPATH(I,IP)-CONST+PF2+EFACTS+EFACT(I,IP)+EXPO1+(I,+(COND1/2.)+ E(1) - 6.28318531+CH1+FP(1)+(1.0 + 2.0+SF+TAU(1))/(RH0(1)+SF) NOW CALCULATE THOSE QUANTITIES WHICH ARE DEPENDENT ON SCAN ANGLE. EITHER MUO AND -MU OR TAU(I) AND TAUZ(I) ARE TOO CLOSE TO-Gether, so we must expand the exponentials in the Path Radi-Ance formula to fourth order terns and cancel terns. LPATH(I,IP)=(CONST/MUFAC1)+PF2+EFACTS+(EFACTO-EFACT(I,IP)) FP(I) = CON+SF+EO(I) / (MUO + SF+TAU(I)) CHI = CO2+MUO+RHO(I) / (1.0 + 2.0+SF+(1.0-RHO(I))+TAU(I)) ALPHA(I) = CHI+ (4.0 + 1.0/SF) G(I) = 4.0+CHI G(I) = 4.0+CHI IN THIS PROGRAM REPRESENTS THE SOLAR AZIMUTH ANGLE WHILE IN THE LITERATURE PHIO USUALLY REPRESENTS THE AZIMUTH ANGLE OF THE DIR-ECTION OF THE SOLAR RADIATION. THE TWO ANGLES ARE SHIFTED BY 90 NOW CALCULATE THE FRESNEL REFLECTANCE AT THIS SCAN ANOLE THE FOLLOWING EQUATION FOR ARG IS CORRECT. REMEMBER THAT PHIO BLOCK OF COMPUTATIONS FOR SINGLY SCATTERED R.DIANCES. NEEDED TO FIND THE REFLECTED SKY RADIANCE. IF (OPTION EQ. 0. OR . OPTION EQ. 2) GO TO 505 CALCULATE SINGLY SCATTERED PATH RADIANCE THESE QUANTITIES DEPEND ON SCAN ANGLE COMPUTE DIRECT IRRADIANCE AT SURFACE SIG(1) - TAU(1) - TAU2(1) Compute Total Irradiance at Surface ED(1) - MUO+EO(1)+EXP(-TAU(1)/MUO) IF (ABS(COND1).LT.COND) GO TO 501 (1.+(COMD1/3.)+(1.+(COMD1/4.))) IF (THETA .LE. 0.0) SPHI - -FPH DEGREES, CAUSING FPHI TO CHANGE SIGN. DO 2019 IP-FSTP,LSTP,PTIMC THETA - QBANG - (IP-1)+QDANG PF2-PF(-ARG, I, TR(I), TAS(I)) IF (SCATT EQ. 0) 60 TO 510 EFACT(I, IP)-EXP(-EXP01) FE(1) - MUO·EO(1)/E(1) ROOT - SQRT( 1.0-MUSQ) ARG - FMU + FPHI + ROOT NUFAC1=(MU+MU0)/MU0 COND 1-EXPO 1 - MUFAC 1 EXPOI-51G(1)/MU MU - COS(THETA) = 1.0-ETA(1) D(I) - MUO/SF FMU - MU·MUO IM-IM - DSIM Hdd - IHd 30 10 502 CONT INUE ŝ 502 201 5 υ υu J ........ c 0000 000 00000 υu υu 421 422 423 425 425 426 428 428 429 431 450 451 432 433 434 435 435 436 437 438 439 440 443 444 445 446 448 452 153 455 455 457 457 457 460 164 469 470 471 441 442 447 454 159 462 463 166 179 461 167 168 173 174 175 176 178 111

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IF (THETA) 58, 58, 58 TABLE 2 (Cont.)	THETTA - ABS(THETA)	MEFMAG-AMSIN(SIN(IHETTA)*1./INDEX) AX=THETTA-DFFDAC	BX=1HETTA+REFRAC	SINA-SIN(AX)	SIMB-SIN(7X)	TANA - TAN(AX)	TANB-TAN(BX)	FRES(IP)=0.5+((SINa+SINA)/(SINE+SINB)+(TANA+TANA)/(TANB+TANB))	G0 T0 310	FRES(IP)-FRESO	CONTINUE		CALCULATE THE DIHER SCAN ANGLE DEPENDENT QUANTITIES MEEDED	TO FIND THE REFLECTED SKY RADIANCE.	ARGSKY=FMJ-FPHI+ROOT	PFSKY-PF(argsky, I, Ta(I), Tas(I))	EXP02-TAU(I)/MJ	MJF AC2 = ( MJO - MJ ) / MJO	COND2 - EXPO2 - MUFAC2	IF (# .(COND2).LT.COND) GD TD 503	CALCULATE THE RETLECTED SKY NAULANCE			G0 10 505			SO THE FOURTH ORDER EXPANSION OF THE REFLECTED SKY RADIANCE		LRSKY(1.1P)+FRES(1P)•CONST+PFSKY+FFACT(1.1P)•FFACT<+F	- (COMD3/2 ) ( COMD3/2 ) ( ) ( - (COMD3/2 ) )		IF (0PTION.LT.2) G0 T0 2019			THE PATH RAUTANCE GENERATED BY THE VIRTUAL SUN.	ARGV   R = [ M1 - F PH ] + R00 [	PFVIR=PF(ARGVIR,I,TR(I),TAS(I))	MUF AC2=(MUO-MU)/MUO	CONDJEEXPOINMUFAC2 If (ARS(COND)) IT COND) SO TO ESS		CALCULATE THE PATH RADIANCE GENERATED BY THE VIRTUAL SUN		LVCRP(I,IP)=FRESUN•(CONST/MUFAC2)•PFVIR•EFACT5•(1,/EFACT0 -	GD TD 3019		EITHER MU AND MUO OR TAU(I) AND TAUZ(I) ARE TOD CLOSE TOGETHED	JSE THE EXPANSION OF THE VIRTUAL PATH RADIANCE.		LVINP(1,1P)=FMESON-CONSI PFVIN*EFAC(1S*EFAC(1,1P)*EXPO1+(1,+)	G0 T0 2019	
	58									59	310	0	0 0	50	,							1	-		c	C	00		503	-	2	505	0	50	5 0	,				0	0	U		-	0	, u	U O		90c	-	
481	482	484	485	486	487	488	489	490	491	492	493	494	495	494	498	499	500	501	502	505	506	507	508	509	510	511	212	115	515	516	517	518	519	075	522	523	524	525	976	528	529	063	165	511	534	535	536	537	970	540	

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SINGI = NUSQ(I) - MUSQ SING2 = NUSQ(I) - CO2\*CO2\*MUSQ IF (ABS(SINGI).GT.EPS .AND. ABS(SING2).GT.EPS) GG TO 55 WRITE (QUERR,380) FORMAT('-ATMSFR: \*\*\*\*\*ERROR-- SINGULARITY EXISTS IN EQUATI ONS\*\*\*\*\*) LPATH(I.TP) = LAMPP(I) • (A(I)•(AT(I)-BT(I)•MU + MFACT) •PF1 +(ATP(I) -BTP(I)•MU -MUO•EFACT(I.IP) +B(I)•MFACT)•PF2)/SIM01 CALCULATE THE SCAN ANGLE DEPENDENT QUANTITIES NEEDED TO FIND (ATPP(I) +BTPP(I)+MU -CTPP(I)+EFACT(I,IP) -DTPP(I)+MFACT) LPATH(I,IP) = FP(I) • ((SIG(I) -MU+CFACT)+PF1 + (SIG(I) + (D(I)-MU)+CFACT)+PF SINGULAR AT THETA-' ,F7.4) ECHO MORE QUANTITIES FOR THE USER'S CONVENIENCE WRITE (QUPRNT, 555) (ETA(I), I = 1, QMCHAN) FORMAT('- ANISOTROPY PARAMETER', 10F 10.5) (, 10F 10.5) (, 10F 10.5) 282 FORMAT('- OPTICAL DEPTH ', 10F10.5 WRITE(OUPRN1,297)(E(1),1-1,ONCHAN) 297 FORMAT('- TOTAL IRRADIANCE',4X,10F10.5) WITE ( OUPRNT , 282) ( TAUZ ( I ) , I - 1, ONCHAN) + (ALPHA(I) + G(I) • MU) • CFACT) WRITE( QUPRNT.370)(ED(1).1-1.QNCHAN) FDRMAT('- DIRECT IRRADIANCE '.10F1( WRITE(QUPRNT.371) PF1 - PF(ARG.I.TR(I). TAS(I)) PF2 - PF(-ARG.I.TR(I). TAS(I)) EFACT(I.IP) - EXP(-SIG(I)/WU) CALCULATE PHASE FUNCTIONS PF1 - PF(ARG.1.TR(1).TAS(1)) PF2 - PF(-ARG.1.TR(1).TAS(1)) FF2 - PF(-ARG.1.TR(1).TAS(1)) CFACT(1.1P) - EXP(-SIG(1)/MU) CFACT - 1.0 - EFACT(1.1P) THE MULTIPLY SCATTERED RADIANCE. IF (OM(I) .GT. OML) GO TO 52 ATHET - 57 29578•THETA WRITE (QUERR,385) ATHET Format(' Atmsfr: GO TO 2019 MFACT - MU · EFACT(1,1P) (ETA) ') CALCULATE PATH RADIANCE CALCULATE PATH RADIANCE UN SURFACE ') ON SURFACE ') DEFINE SINGULARITIES WRITE (QUPRNT, 291) WRITE (OUPRNT, 298) GO TO 2019 +(I)dIHd+ CONT INUE CONTINUE /SING2 CONT INUE CONTINUE FORMAT( ' FORMAT ( FORMAT( ' CONTINUE **RETURN** 167 01E 525 116 510 2019 298 55 380 385 9 52 J J 000 c υu J U C 589 541 544 544 545 545 545 546 547 5549 5549 5549 5553 5554 5555 5555 5555 5555 558 562 565 566 569 570 575 576 578 579 580 582 583 584 585 586 587 588 590 591 592 593 595 595 596 599 560 561 567 568 572 573 574 577 185 598

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CONTINUE IF (LSW.NE.6) GO TO 3019 PRINT RADIOMETRIC QUANTITIES PRINT RADIOMETRIC QUANTITIES IF(SCATT.EQ.0.OR.OPTION.EQ.0.OR.OPTION.EQ.2)WRITE(QUPRNT.410 IF(SCATT.EQ.0.OR.OPTION.EQ.0.OR.OPTION.EQ.2)WRITE(QUPRNT.410 10 IF (SCATT.EQ.O) GO 10 520 IF (OPTION.EQ.1.OR.OPTION.EQ.3)LBEAM(IV)-LBEAM(IV)-LRSKY(IV. FORMAT(5x,'LINE MUMBER',16,5x,'PIXEL MUMBER',14,5x, 'SCAN ANGLE-',F7.3,'DEGREES FRESNEL DEGLETCE' HIS PIXEL. IF NON-ZERO, DON'T PROCESS 'SCAN ANGLE-',FT.3,'DEGREES') If(SCATT.NE.O.AND.(OPTION.EQ. OR.OPTION.EQ.3))WRITE(QUPRNT 8 FORMAT('O', 14X,'1',9X,'2',9X,'3',9X,'4',9X,'5',9X,'6' 9X,'7',9X,'8',9X,'9',9X,'10') WRITE (QUPRNT,425) (L(IW,IP),IW-1,ONCHAN) IF WE ARE PRINTING RADIOMETRIC QUANTITIES, START EACH LINE FORMAT(5X, 'LINE NUMBER', 16, 5X, 'PIXEL NUMBER', 14, 5X IF (OPTION.GE.2)LBEAM(IW)-LBEAM(IW)-LVIRP(IW, IP) FORMAT( 'O LTGT ', 10F10.5) WRITE (QUPRNT, 430) (LPATH(IW.IP), IW-1, QNCHAN) TOTAL RADIANCE (EXPERIMENTAL) - L(IV.IP) LBEAM(IV) - L(IV.IP) - LPATH(IV.IP) LOOP FOR EACH PIXEL - - QNSS-MUMBER OF PIXELS DO 1009 IP-FSTP, LSTP, PTINC THETA = (QBANG-QDANG\*(IP-1)) • 180/PI - LBEAM( IV)/EFACT( IV. IP) IF (L(IW, IP).EQ.0) LBEAM(IW) - 0 IF (CONTRL(IP) .NE. 0) GO TO 2009 CALCULATE INTRINSIC RADIANCE LINTR(IW) - FE(IW)\*LSURF(IW) CALCULATE SURFACE REFLECTANCE GO TO (71.72.73.74.75.76).LSW CALCULATE SURFACE RADIANCE RHOS(IV) - LSURF(IV)/E(IV) LPATH ', 10F 10.5) IF (LSW.EQ.6) WRITE (QUPRNT, 400) L(IV. IP) - LPATH(IV. IP) TEST "BAD DATA" FLAG FOR L(IV, IP) - LBEAM(IV) L(IV, IP) - LSURF(IV) L(IV. IP) - LINIR(IW) L(IV. IP) - RHOS(IV) ( OUPRNT , 420) DO 3009 14-1, QNCHAN LSURF (IV) GO TO 60 GO TO 60 GO TO 60 G0 10 60 GO TO 60 CONT INUE CONT / NUE FORMAT ( A NEW PAGE CONT I NUE CONT INUE ARI TE FDRMAT( '1') CONT INUE -520 410 411 3009 400 420 425 430 90 1 72 51 14 15 76 • J U U U J υ U U U 610 604 605 609 613 614 615 615 615 615 619 620 622 623 624 625 626 627 628 629 663 635 636 638 638 638 640 643 644 645 645 650 653 655 658 659 660 603 612 642 549 501 607 617 621 169 632 1034 647 648 223 656 657

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IF (SCATT EQ.0)G0 T0 433	IF (OPTION EQ. 1. OR. OPTION EQ. 3) WRITE (QUPRNT 431) (LRSKY (IV. IP)	LV-1 ONCHAN)	FORMAT( ' LRSKY ', 10F 10.5)	IF (0PT10N, GE, 2) WRITE (OUPRNT, 432) (LVIRP (IV, IP), IV-1, ONCHAN)	-	-	-	WRITE (QUPRNT, 435) LBEAM	-	WRITE (QUPRNT, 440) LSURF	-	WRITE (QUPRNT, 445) LINTR		WRITE (QUPRNT, 450) RHOS		SET QBNONE TRUE TO TELL ANY FOLLOWING MODULES THAT NO OUTPUT	DATA ARE PRESENT.	QBNONE - TRUE	3	09 CONTINUE	09 CONTINUE	RETURN	END	
		-	164		432	664	434		435		440		445		450	J	U		3019	2009	1009 C	œ		LE L
661	662	663	664	665	666	667	668	699	670	671	672	673	674	675	676	677	678	619	680	681	682	683	684	END OF FILE

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The user must also define the center wavelengths ( $\mu$ m) of the multispectral scanner and the corresponding values of the surface background albedo (values from zero to one). Also, the center wavelength<sub>3</sub> ( $\mu$ m) of the surface radiometers and the corresponding optical thicknesses must be known. It should be noted that the optical thicknesses used should be those measured as closely as possible in time with the multispectral data.

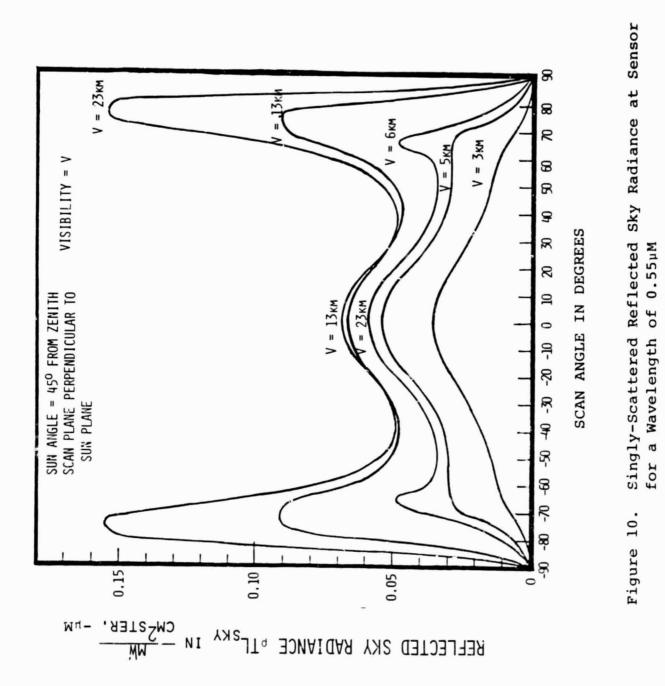
#### 5.4 MODEL CALCULATIONS

In this section we present several examples of the radiances for the various components. Because our main interest is in the radiance components as a function of scan angle and visibility, we will presert the results of the calculations in terms of these parameters.

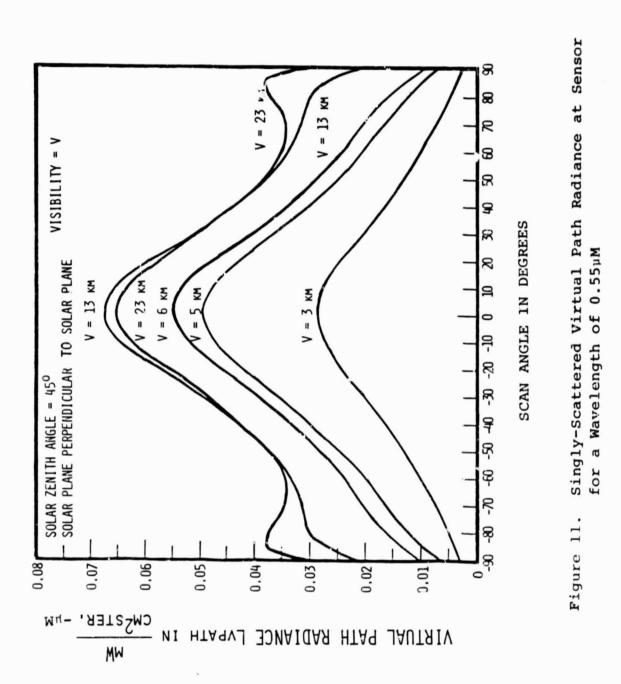
Figure 39 depicts the variation in the singly-scattered reflected sky radiance at the sensor as a function of the nadir scan angle and visibility. In this example the solar zenith angle is  $45^{\circ}$  and the scan plane is perpendicular to the solar plane. The curves which result are a combination of the variation of the sky radiance, the transmittance from the surface to the rensor, and the Fresnel reflectance of the water surface. For a practical scanner with a maximum scan angle of about  $45^{\circ}$  the curves indicate that one would not observe the large radiance peaks at the large angles.

In Figure 11 we display the corresponding path radiance as a result of singly-scattered radiation from the reflection of the sun in the water. In this case the radiance peaks do not exist at the large scan angles.

In Figure 12 we illustrate the relative magnitudes of the various radiation components as a function of scan angle for a moderately hazy atmosphere. The virtual sun path radiance is the smallest value and the multiply scattered sky radiance is the largest value.



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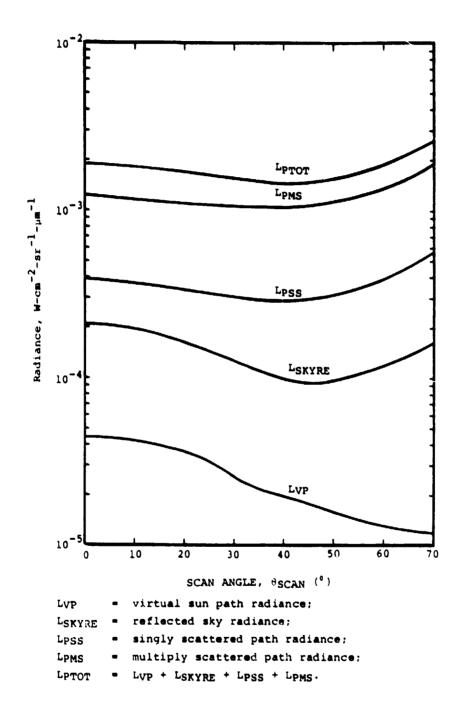


Figure 12. SURFACE AND PATH RADIANCE COMPONENTS DETECTED BY SENSOR AT OPTICAL DEPTH OF  $\tau = 0.346$  AS A FUNCTION OF SCAN ANGLE,  $\theta$ SCAN. SCAN PLANE  $\perp$  SOLAR PLANE,  $\theta$ SUN = 30°,  $\lambda = 0.55 \mu$ M, PHASE FUNCTION = CONTI-NENTAL REFRACTIVE INDEX 1.5 - 0.01i, VISIBILITY = 10 KM. In Figure 13 we indicate the variation in the ratio of the singly-scattered sky radiance to the singly-scattered path radiance as a function of the optical depth  $\tau$  of the sensor. As the curves illustrate, the reflected sky radiance component is relatively more important for the larger optical depths.

Figure 14 illustrates the variation in the ratio of the virtual sun path radiance to the singly-scattered path radiance with scan angle for four optical depths.

Figure 15 depicts the large ratio of the multiply-scattered component to the singly-scattered path radiance component as a function of optical depth and scan angle.

Because optical thickness or visibility is of major importance in remote sensing investigations, we want to consider the variation of the radiance components with respect to visibility. This effect is illustrated in Figure 16 for three different atmospheres. We chose the continental aerosol because it more nearly represents the type which would be found over the Great Lakes. The three refractive indices are: 1.5-0.0i which corresponds to a "clean" haze, i.e., one where there is no absorption; 1.5-0.0li which corresponds to a haze with some aerosol absorption; and 1.5-0.1i, a complex index of refraction which corresponds to a haze with more absorption. As the curves indicate, an absorbing haze or one which corresponds to considerable air pollution gives rise to a large ratio of reflected sky radiance relative to the singly-scattered path radiance.

The effect of the complex index of refraction is also evident in the ratio of the virtual sun path radiance to the singly-scattered path radiance as indicated in Figure 17.

Finally, we illustrate in Figure 18 the variation of various combinations of ratios in terms of the visibility for a refractive index of 1.5-0.01i.

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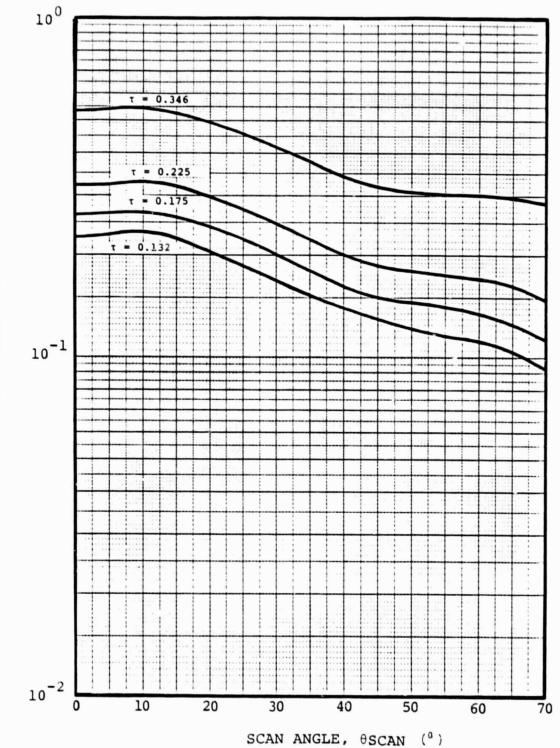


FIGURE 13. RATIO OF REFLECTED SINGLY SCATTERED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE AS A FUNCTION OF SCAN ANGLE,  $\theta$ SCAN, FOR OPTICAL DEPTH,  $\tau$ , OF THE SENSOR OF 0.132, 0.175, 0.225 and 0.346. SCAN PLANE  $\perp$  SOLAR PLANE, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX 1.5 - 0.01i, VISIBILITY = 10 KM,  $\lambda$  = 0.55  $\mu$ M,  $\theta$ SUN = 30°.

RATIO, LSKYRE/LPSS

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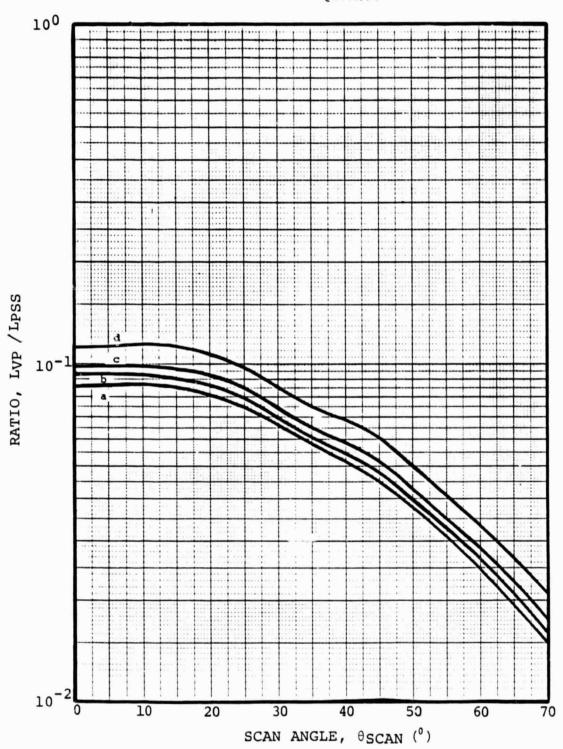
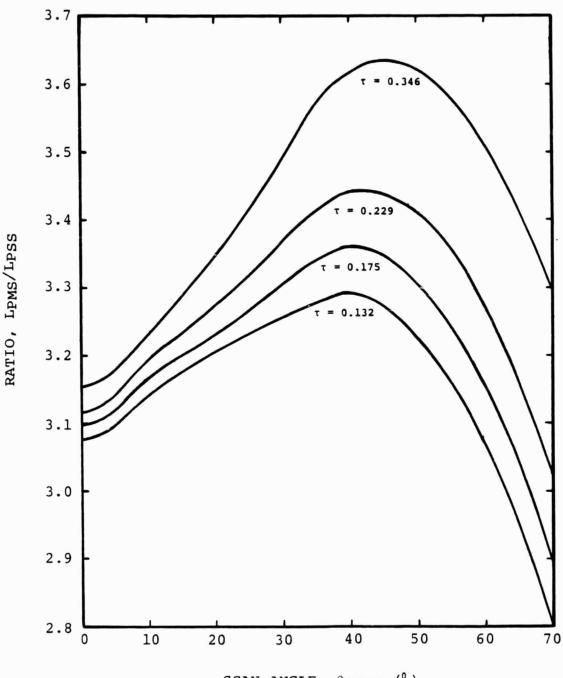


FIGURE 14. RATIO OF SINGLY SCATTERED PATH RADIANCE FROM THE VIRTUAL SUN, LVP, TO SINGLY SCATTERED PATH RADIANCE, LPSS, AS A FUNCTION OF SCAN ANGLE,  $\theta$ SCAN, FOR OPTICAL DEPTHS,  $\tau$ , OF THE SENSOR OF A) 0.132, B) 0.175, C) 0.225 AND D) 0.346. SCAN PLANE  $\perp$  SOLAR PLANE, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX 1.5 - 0.011, VISIBILITY = 10 KM,  $\lambda = 0.55 \mu$ M,  $\theta$ SUN = 30°.

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SCAN ANGLE,  $\theta_{SCAN}$  (°)

FIGURE 15. RATIO OF MULTIPLY SCATTERED PATH RADIANCE, LPMS, TO SINGLY SCATTERED PATH RADIANCE, LPSS, AS A FUNCTION OF SCAN ANGLE,  $\theta$ SCAN, FOR OPTICAL DEPTHS,  $\tau$ , OF THE SENSOR OF 0.132, 0.175, 0.229, 0.346. SCAN PLANE  $\perp$ SOLAR PLANE, VISIBILITY = 10 KM, PHASE FUNCTION = CONTINENTAL REFRACTIVE INDEX 1.5 - 0.01i,  $\lambda$  = 0.55  $\mu$ M,  $\theta$ SUN = 30<sup>0</sup>.

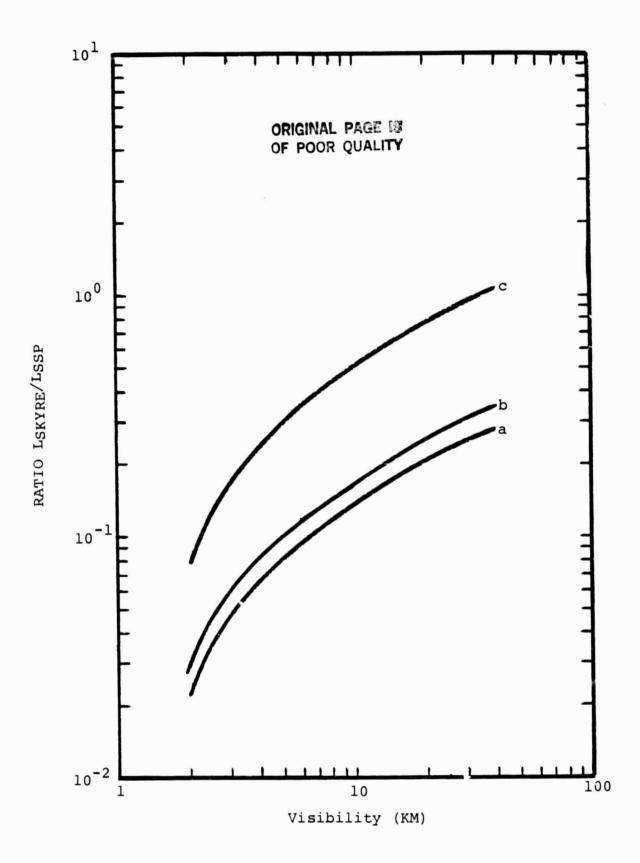
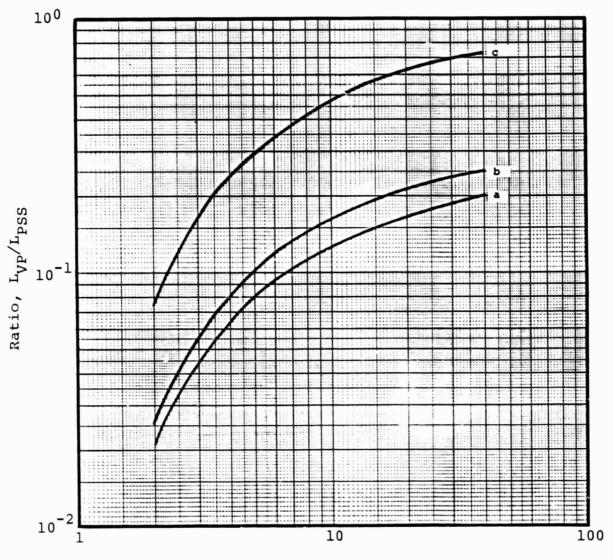


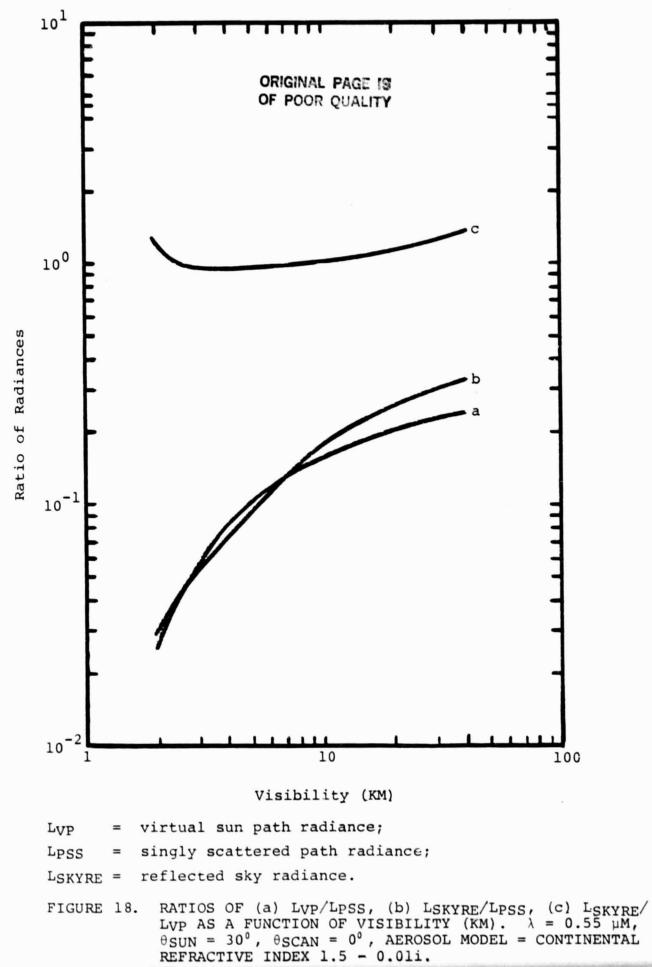
Figure 16. RATIO OF REFLECTED SKY RADIANCE TO SINGLY SCATTERED PATH RADIANCE FOR THREE CONTINENTAL AEROSOL MODELS: a) 1.5 – 0.0i, b) 1.5 – 0.0li, and c) 1.5 – 0.1i AS A FUNCTION OF ATMOSPHERIC VISIBILITY. SCAN PLANE  $\perp$  SOLAR PLANE,  $\theta$ SUN = 30°,  $\lambda$  = 0.55  $\mu$ M.

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Visibility (KM)

FIGURE 17. RATIO OF SINGLY SCATTERED PATH RADIANCE FROM THE VIRTUAL SUN, LVP, TO SINGLY SCATTERED PATH RADI-ANCE, LPSS, AS A FUNCTION OF ATMOSPHERIC VISIBIL-ITY FOR THREE CONTINENTAL AEROSOL MODELS WITH REFRACTIVE INDICES OF a) 1.5 - 0.0i, b) 1.5 - 0.0li, and c) 1.5 - 0.1i.  $\lambda = 0.55 \mu$ M,  $\theta$ SUN = 30°,  $\theta$ SCAN = 0°.



#### CONCLUSIONS AND RECOMMENDATIONS

6

The problem of developing an atmospheric correction algorithm for remote sensing is an old and difficult one. The main difficulty lies in not being able to have available sufficient data which can be used to specify the values of the relevant atmospheric para-The problem is all the more difficult in the case of the meters remote sensing of water bodies because of the low signal-to-noise ratio involved. In this investigation we have extended an existing computer algorithm so as to include additional radiation components. The original algorithm included the path radiance which arises from the singly-scattered solar radiation in the atmosphere. We have now included the radiance which arises from the sky radiation which is reflected by the water surface and is then attenuated as it propagates from the surface to the sensor. In addition, we have included the path radiance component which arises from the single scattering of radiation as a result of a virtual sun, i.e., of the sun's reflection in the water. It should be realized that this component is always present regardless of the scan plane, i.e., it does not only occur when the scanner is looking at the specular angle. In addition to these components, we have also included a multiple-scattering approximation. It should be realized, however, that the multiple scattering applies only to an atmosphere with the sun as a source. Another multiple scattering calculation should be performed to include the effect due to the virtual sun.

The general result of all these calculations indicates that the various components are all about equal in magnitude but that there is considerable variation with respect to scan angle and visibility. Also, it appears that the multiply-scattered component is of major significance.

It must be pointed out that the objective of this investigation is to provide an algorithm for the correction of remotely sensed data for atmospheric effects so that one can extract from the multispectral data the radiance which is characteristic of the water In our investigation we have dealt with the water surface itself. as a flat, specular reflector, which in general is not true. Α wind-roughened surface will be characterized by a complex wave structure which leads to a more complicated representation of the reflected and virtual sun radiances than presented in this report. A further investigation should be conducted to model the water surface in terms of wind speed and a stochastic representation. of the reflecting facets of the water surface. If this is done, then a more realistic model could be developed which should provide better values for the sky-reflected and the virtual sun path radiances. It may even be possible to establish a method for the determination of wind speed by observing the average radiance as a function of the instantaneous field of view.

A further recommendation is to improve the accuracy of the algorithm by including a more detailed calculation of the multiplyscattered path radiance, both for the direct sun as a source and for the virtual sun as a source.

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