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A METHOD FOR ESTIMATING THE ROLLING MOMENT DUE TO SPIN RATE FOR ARBITRARY PLANFORM WINGS

FOR REFERENCE

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NOT TO BE TAKEN FROM THIS ROOM

Inchestratement (KPJd)

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Langley Research Center Hampton, Virginia 23665

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A ME**TH**O**DF**O**E** E**STIMATINGTH**E **R**O**LL**I**NGM**O**MENT.DU**E **T**O **SPINRAT**E**F**O**E ARB**I**TR**A**RYPI**A**NF**O**RMWINGS**

Willia**m** A. Poppen, Jr.

ABSTRA**CT**

The application o**f** aer**o**dynamic theory for estimating the force and moments acting upon spinning airplanes is of interest. For example, strip theory has been used to generate estimates of the aerodynamic characteristics as a function of spin
rate for wing-dominated configurations for angles of attack up to 90 degrees. This rate for wing-dominated configurations for angles of attack up to 90 degrees. work, which had been limited to constant chord wings, is extended here to wings comprised of tapered segments. Comparison of the analytical predictions with rotary balance wind tunnel results shows that large discrepancies remain, particularly for those angles-of-attack greater than 40 degrees.

NOME**NCLATURE**

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The use of parameter estimation in **m**odeling aircraft dynamics has been quite su**cc**essful for many mathemati**c**al models of flight. Parameter estimation is most read**i**ly applied when linear models representing smal**l** perturbations from straight equilibrium paths are appropriate. Flight data is most a**cc**urate in this reg**i**me and the mat**h**emati**c**al model is the slmplest. 1

Par**a**mete**r** estimation be**c**omes m**o**re **c**omp**l**ex in app**li**cati**o**n to spinning air**cr**aft. Modeling nonlinear aerodynamics, including rotational flow effects, is much more difficult and many more unknown parameters are introduced. **2** In order to reduce the large number of unknowns it is helpful to apply strip theory of referen**c**e 3. Strip theory "**l**inks the **w**ing ai**r**foi**l** section **c**hara**c**teristi**c**s to the ro**ll**ing and yawing moment of the wing in spinning f**l**ight. ''I

In referen**c**e I, **s**trip theory provided a mathematical model that was used to determine the rolling moment of a wing in spinning flight. Cal**c**ulated rolling moment for**c**es due to the wing were about 50 percent larger than the experimental rotary balance spin-tunnel measurements of a wing-dominated aircraft. It is the purpose of this paper to expand the existing mathemati**c**a**l** mode**l** of a spinning wing in order to more closely repre**s**ent an aircraft in spinning flight, and to further explore the limitations and poss**i**bilit**i**es of the more general model. Spec**i**fically, the str**i**p theory te**c**hnique of reference I w**i**ll be extended to wings comprised of tapered segments. The same **l**im**i**tat**i**on of reference 1 will be used **i**n that the flow angle at ea**c**h strip lo**c**ation is independent of the incremental lift at other locations.

DI**SCUSS**I**O**N

In order to decrease the complex**i**ty of estimating the rolling moment due to spinning, the authors in referen**c**e 1 restri**c**ted their analysis to the rolling moment produced by an untapered wing of a wing-dominated aircraft. In this paper the approach is extended to wings of arbitr**a**ry planform hy considerelng **a** wing to be made up of se**c**tions of differing taper.

Let us first **c**onsider the lo**c**al flow **c**hara**c**teristi**c**s for the general sp**a**nwise lo**c**ation y, shown in figure I.

$$
v_{\ell}^{2} = (u - ry)^{2} + (w + py)^{2}
$$
\n
$$
\alpha_{\ell} = \arctan\left(\frac{w + py}{u - ry}\right) = \arcsin\left(\frac{w + py}{\sqrt{(u - ry)^{2} + (w + wy)^{2}}}\right)
$$
\n
$$
q_{\ell} = \frac{\rho}{2} \left[(u - ry)^{2} + (w + py)^{2} \right]
$$

and

$$
q_{\ell} = \frac{\rho}{2} \left[(u - ry)^{2} + (w + py)^{2} \right]
$$

For wings having a constant taper*,* the wing chord can be represented by a linear equation:

$$
c = c_0 - hy \quad \text{for} \quad y > 0
$$

$$
c = c_0 + hy \quad \text{for} \quad y < 0
$$

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The equation for rolling moment for a single strip would be:

$$
dL = -\frac{\rho}{2} \left[(u - ry)^2 + (w + py)^2 \right] C_N \left(c_0 \pm hy y \, dy \right)
$$

For the entire wing the rolling moment becomes:

$$
L = -\frac{\rho}{2} \int_{-\frac{b}{2}}^{\frac{b}{2}} \left[(u - ry)^{2} + (w + py)^{2} \right] c_{N}(y) \left(c_{o} \pm hy \right) y \ dy
$$

In order to easily represent aerodynamic data at high angles of attacks, the normal force coefficient (fig. 2) is given the form:

$$
C_N(\alpha) = C_N + C_{N_{\text{sin}\alpha}}
$$

It follows then that a single wing section over which the normal force equation is applicable will have the following contribution to rolling moment:

$$
\Delta L = -\frac{{}^{6}C_{0}C_{N}}{2} \int_{y_{lower}}^{y_{upper}} [(u - ry)^{2} + (w + py)^{2}] y dy =
$$

$$
\frac{{^{h\rho C}}N_0}{2}\int_{y_{lower}}^{y_{upper}} [(u - ry)^2 + (w + py)^2]y^2 dy -
$$

$$
\frac{\rho c_o C_{N_{\text{sin}\alpha}}}{2} \int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{(u - ry)^2 + (w + py)^2} (w + py)y \, dy =
$$

$$
\frac{{^{h\rho C}}_{N}}{2} \int_{y_{lower}}^{y_{upper}} \sqrt{{(u - ry)}^{2} + (w + py)}^{2}} (w + py)y^{2} dy
$$

After integrating,

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$$
\Delta L = - \frac{\rho c_o C_{N_o}}{2} \frac{1}{4} \left[A \left(y_{upper}^4 - y_{lower}^4 \right) + \frac{1}{3} B \left(y_{upper}^3 - y_{lower}^3 \right) + \frac{1}{2} C \left(y_{upper}^2 - y_{lower}^2 \right) \right] + \frac{h \rho C}{2} \left[\frac{1}{5} A \left(y_{upper}^5 - y_{lower}^5 \right) - 3 - \frac{1}{5} C \left(y_{upper}^2 - y_{lower}^2 \right) \right]
$$

$$
+\frac{\frac{1}{4} B \left(y_{\text{upper}}^{4} - y_{\text{lower}}^{4}\right) + \frac{1}{3} C \left(y_{\text{upper}}^{3} - y_{\text{lower}}^{3}\right)}{\frac{WPC_{0} C_{N_{\text{sing}}}}{2} \int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{\phi} y \, dy - \frac{C_{N_{\text{sing}}}}{2} \left(\text{pc}_{0} \pm \text{wh}\right) \int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{\phi} y^{2} dy}
$$

$$
\pm \qquad \frac{\text{pphC}_{N_{\text{sing}}}}{2} \int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{\phi} y^{3} dy
$$

where:

 \overline{a}

$$
A = p2 + r2 = \Omega2
$$

\n
$$
B = -2ur + 2wp = 0
$$

\n
$$
C = u2 + w2 = v2
$$

\n
$$
\Phi = Ay2 + By + C
$$

$$
\int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{\Phi} \ dy = \frac{1}{4A} \left[\left(2Ay_{\text{upper}} \right) + B \sqrt{Ay_{\text{upper}}^2 + By_{\text{upper}} + C} \right]
$$

$$
-\left(2Ay_{lower}\right)+B\sqrt{Ay_{lower}^2+By_{lower}+C}
$$

+ $\frac{4AC-B^2}{8A\sqrt{A}}$ log $\left[\frac{2Ay_{upper}+B+2\sqrt{A^2y_{upper}^2+ABy_{upper}+AC}}{2Ay_{lower}+B+2\sqrt{A^2y_{lower}^2+ABy_{lower}+AC}}\right]$

$$
\int_{y_{\text{lower}}}^{y_{\text{upper}}} \sqrt{\Phi} \, y \, dy = \frac{1}{3A} \left[\left(A y_{\text{upper}}^2 + B y_{\text{upper}} + c \right)^{3/2} \right]
$$

$$
-\left(Ay_{lower}^2 + By_{lower} + c\right)^{3/2}\right]
$$

$$
-\frac{B}{2A} \int_{y_{lower}}^{y_{upper}} \sqrt{\phi} dy
$$

$$
\int_{y_{lower}}^{y_{upper}} \sqrt{\phi} y^{2} dy = \frac{6Ay_{upper} - 5B}{24A^{2}} \left(Ay_{upper}^{2} + By_{upper} + c \right) \frac{3}{2}
$$

$$
-\frac{6Ay_{lower} - 5B}{24A^{2}} \left(Ay_{lower}^{2} + By_{lower} + c \right) \frac{3}{2}
$$

$$
-\frac{4AC - 5B^{2}}{16A^{2}} \int_{y_{lower}}^{y_{upper}} \sqrt{\phi} dy
$$

$$
\int_{y_{lower}}^{y_{upper}} \sqrt{\phi} y^{3} dy = \left(\frac{y_{upper}^{2}}{5A} - \frac{7By_{upper}}{40A^{2}} + \frac{7B^{2}}{48A^{3}} - \frac{2C}{15A^{2}} \right) \left(Ay_{upper}^{2} + By_{upper} + c \right) \frac{3}{2}
$$

$$
-\left(\frac{y_{lower}^{2}}{5A} - \frac{7By_{lower}}{40A^{2}} + \frac{7B^{2}}{48A^{3}} - \frac{2C}{15A^{2}} \right) \left(Ay_{upper}^{2} + By_{lower} + c \right) \frac{3}{2}
$$

$$
-\left(\frac{7B^{3}}{32A^{3}} - \frac{3CB}{8A^{2}} \right) \left(\int_{y_{lower}}^{y_{upper}} \sqrt{\phi} dy \right)
$$

 T_{max} in the normal force equation, $\frac{N_{\text{O}}}{N_{\text{O}}}$ and $\frac{N_{\text{S}}}{N_{\text{O}}}$ the lo**c**al angle-of-attack ranges (see fig. 2) listed in Table 1 from reference I:

-5-

The wing span locations having local angles of attack of -16 , -10.5 , 10.5 and 16 degrees are determined by:

/

$$
y_{\text{boundary}} = \frac{w - u \tan(\alpha \text{ boundary})}{-p - r \tan(\alpha \text{ boundary})}
$$

These will serve as limits of integration in the above equations if they fall in the confines of the panel being considered. If they do not, the boundaries of the panel will be used as limits.

The program used to calculate the rolling moment of the wing using the above equations is listed in the appendix. It is a series of subroutines that will calculate the rolling moment coefficient of any tapered section of a flat wing given the following data: the two boundary chord lengths of each panel; the distance of these chords from the origin; the air density; the velocity of the aircraft; the wingspan; and the area of the wing. There is an option to calculate the rolling moment coefficient of a single panel, or both symmetrical panels having the given dimensions.

The spin subroutine accepts the dimensions of the panel and calculates the slope of the linear equation describing wing taper (h). It then computes the limits of integration along the panel. These limits are sent to the intermediate tests integration along the panel. These limits are sent to the intermediate tests subroutine. Tests classifies the limits and sends only those that are within the bounds of the desired panel(s) to the panel subroutine. The panel subroutine does the actual rolling moment calculation of the panel between the limits using the above
equations. The split subroutine is an optional subroutine which given the The split subroutine is an optional subroutine which, given the dimensions of the wing, will split a wing into its component panels and send each panel in succession to the spin subroutine.

With this program, a wing comprised of tapered panels can be modeled, panel by panel. Through a simple modification, the program can accumulate the total rolling moment of an aircraft wing due to each panel at a selected angle-of-attack. For the airplane shown in figure 3, this was done at an angle-of-attack of 14 degrees in order to obtain figure 4. Figure 4 is a plot of the total rolling moment coefficient of the wing of the aircraft, as well as the rolling moment coefficient of each of the wing's component panels as a function of nondimensional spin rate. The bottom curve of figure 4 represents the total rolling moment coefficient for the airplane of figure 3 at 14 degrees angle-of-attack.

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In a typical light, wing-dominated aircraft such as the one illustrated in figure 3, the panels that cause the greatest moment are the outer panels as is shown
in figure 4 and in table 2. The upper curve in figure 4 represents the rolling The upper curve in figure 4 represents the rolling moment contribution of the inner panels, the next curve represents the contribution of the middle panels and the third curve from the top is the contribution of the
outer panels. This figure is for a fixed angle-of-attack while the rotation This figure is for a fixed angle-of-attack while the rotation rate varies. On the other hand, table I shows the relationship between the panels when the rotation rate is fixed and the angle-of-attack is varied. The data of table I and Figure 4 clearly show that the outer panels contribute from 78% to 97% of the total rolling moment. Of course, this is expected since these panels are larger than the others, have the longest moment arm, and experience the greatest variation in dynamic pressure.

Figure 5 shows the improvement caused by taking into account wing taper as
red to the values obtained with a constant chord. There is significant compared to the values obtained with a constant chord. improvement in the data, particularly at higher rates of rotation. The upper curves are the spin-tunnel test data. Obviously, improvements in the model must be made
before the method can be considered acceptable. It is interesting to note (see before the method can be considered acceptable. figures 5 and 6) that there is little difference when the wing of the aircraft in figure 3 is simplified in the calculations to two large trapezoidal panels instead of However, the multi-panel approach is more accurate and is applicable to the more general case.

- 7-

In reference I, it was noted that at angles-of-attack around 50 degrees the experimental rolling moments were autorotative at low rotation rates. The calculated data of reference I did not represent this phenomenon. The plot of 30 and 50 degrees angle-of-attack in figure 7 shows that the new calculated data does not show autorotative moments either. With the theory being used here, it would be impossible to obtain autorotative moments except over an angle-of-attack range of 10.5 to 16 degrees since the slope of the line of normal force coefficient vs. angle-of-attack (fig. 2) is always positive except over this range. Note that figures 5 and 6 show an autorotative moment at low rates of rotation both in the test data and in the calculated data for 14 degrees angle-of-attack. However, an extension must be made to this simplified aerodynamic theory for higher angles-of-attack.l

The amount of error in the mathematical representation of a spinning wing has been decreased by describing the wing as a set of tapered panels. However, the been decreased by describing the wing as a set of tapered panels. However, the errors are still large. The next step might be to consider the contribution of the tail section to the rolling moment. Since the program calculates the rolling moment of any tapered panel, the three tail panels could be input in order to determine the
tail effects. The present method will compute the rolling moment for swept-wing tail effects. The present method will compute the rolling moment for swept-wing configurations since rolling moment is independent of sweep. However, an extension of the model should also incorporate pitching moment. Of course, this method will
not hold for aircraft where body effects cannot be neglected. The effects of the not hold for aircraft where body effects cannot be neglected. body would have to be considered by some other method such as the strip theory of reference 5. Improved estimates of aerodynamic moments would be expected if the Improved estimates of aerodynamic moments would be expected if the induced flow effects on the flow angles were included in the formulation. Past results and future extentions promise further improvements in predicting the aerodynamic forces and moments of spinning airplanes.

CON**C**L**U**DIN**G**RE**M**ARK**S**

Mathematical representations of nonlinear phenomena such as the aerodynamics of a spinning aircraft are characterized by having large numbers of unknown parameters. Analytical methods such as strip theory can be used to reduce the number of unknown parameters. In this paper, strip theory is applied to compute aerodynamic forces for a wing composed of several variable taper trapezoidal panels in order to obtain a model structure which requires only the unknowns of the normal force equations. Although the error is decreased significantly by using strip theory in this manner as compared to approximating the wing as untapered, there is still much more to be done in order to analytically predict aerodynamic force of spinning aircraft. In order to extend the model further, many new parameters would have to be added. Also, it is clear that aerodynamic theory for angles-of-attack greater than 40 degrees must be improved since it is impossible to predict the results of spin-tunnel rotary balance tests with strip theory methods.

Since the program that calculates the data is general enough to accept any wing panel of an aircraft, the revised model is currently useful in comparing the effects on a panel of changing parameters such as rotation rate, angle-of-attack, velocity, taper, etc. It is also useful for comparing aircraft components. However, the error between anlytical predictions and the experimental data is still too large to consider the strip theory representation to be an effective model of a spinning aircraft.

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APPEN**DIX**

Listed in the f**o**llowing pages is the pr**o**gra**m** used t**o** calc**u**late the rolling moment due to spin rate for an arbitrary planform wing. Inputs to the program are set in a sh**o**rt main program which calls the s**u**br**o**utines necessary f**o**r the calculations.

The first page shows an example of the simplest case where a single panel is input t**o** the pr**o**gra**m**. Va**r**iable CHORD is the inb**o**ard cho**r**d length and **C**HTIP is the outboard chord length. D1 and D2 are spanwise distances from the center line of the fuselage to CHORD and CHTIP respectively. AREA refers to the area of the entire wing containing the panel, and SPAN is the wing span. RHO is air density and VEL is the velocity of the aircraft. The last integer tells the program whether to compute the rolling moment for panel with the given dimensions on the positive side of the aircraft (0), the negative side of the aircraft (1) or both (2). Normally, this aircraft (0), the negative side of the aircraft (1) or both (2). value will be 2, except when isolation a single wing panel is desired.

The second page shows a case where the panels of an entire wing will be input to the program. In this case, variable CHORD is the root chord CHI and CHI are the the progra**m**. In this case, variable CHORD is the root chord, CH**I** and CH2 are the chord lengths at the point of wing taper change, and CHTIP is the chord length at the wing tip. D! and D2 are the spanwise distances from the center line to CH**I** and CH2 respectively. FUSE is the width of the fuselage.

As was mentioned in the text, the program can be modified to accu**m**ulate the total rolling moment of a multi-paneled wing. To do this, a one-dimenslonal array can be defined and placed inside the main loop of the spin subroutine such that each time spin is called with a new panel's dimensions, the nine values of the panel's rolling moment coefficient array (CLW) are added to the new array for a selected angle-of-attack (corresponding to an iteration of the main loop). For more than one angle-of-attack, a two-dimensional array would be necessary.

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SYMBOLIC REFERENCE MAP (R=1)

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SYNSOLIC REFERENCE MAP (R+1)

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SYNODLIC REFERENCE NAP (R=1) $\ddot{}$

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SYMBOLIC REFERENCE HAP (R=1)

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Table 2

Rolling moment coefficient for the contributing trapezoidal panels on the main wing of the airplane shown in figure 3 at angles of attack from 0 to 24 degrees and $\Omega_0/2V = 0.5$.

Figure 1. Schematic Sketch of a Spinning Wing

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Figure 3. Drawing of 1/16-scale typical light airplane of reference 4. Sote
the 6 distinct panels on the main wing.

Figure 5. Effects of rotation rate on rolling moment coefficient for 8 and 14 degrees
angle of attack.

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Figure 7. Effect of rotation rate on rolling moment coefficient for 30 and
50 degrees angle of attack.

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 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ $\label{eq:2.1} \mathcal{L}(\mathcal{L}) = \mathcal{L}(\mathcal{L}) \mathcal{L}(\mathcal{L}) = \mathcal{L}(\mathcal{L}) \mathcal{L}(\mathcal{L})$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$ \mathcal{L}_{max} and \mathcal{L}_{max} $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}}),\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{$ $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$ and $\mathcal{L}^{\mathcal{L}}(\mathcal{L}^{\mathcal{L}})$. Then the contribution of $\mathcal{L}^{\mathcal{L}}$ $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$ fl $\mathcal{L}^{\text{max}}_{\text{max}}$

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 $\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{0}^{\infty}\frac{dx}{\sqrt{2\pi}}\,dx\leq \frac{1}{\sqrt{2\pi}}\int_{0}^{\infty}\frac{dx}{\sqrt{2\pi}}\,dx$

 $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2\frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^2.$

 $\label{eq:2.1} \begin{split} \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) & = \frac{1}{2} \sum_{i=1}^{N} \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \\ & = \frac{1}{2} \sum_{i=1}^{N} \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r}) \mathcal{L}_{$