

ON A POSSIBILITY OF INELASTICITY PARTIAL COEFFICIENT  $K_{\gamma}$   
 DETERMINATION IN  $\pi C$  AND  $\pi Pb$  INTERACTIONS AT  $10^{14}$  EV  
 (EXPERIMENT "PAMIR" I)

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The investigation of hadron-nuclear interactions in  
 "Pamir" experiment is carried out by means of X-ray emulsion  
 chambers of two types: carbon (C) and lead (Pb) ones [1].

While comparing the results from the chambers of both  
 types [2] it was found a discrepancy in  $\langle n_h \rangle$  and  $\langle E_h^{(0)} \rangle$   
 values.

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\* This article only

Here  $\langle n_h \rangle$  is an average hadron multiplicity normalized to one  $\gamma$ -h-family fitted the following selection criterion A:  $\sum E_\gamma \geq 100$  Tev,  $n_h \geq 0$ ,  $E_h^{(v)} \geq 4$  Tev;  $\langle E_h^{(v)} R \rangle$  is a space-energetic characteristics of hadrons in families fitted the criterion B:  $\sum E_\gamma \geq 100$  Tev,  $n_h > 3$ ,  $E_h^{(v)} \geq 4$  Tev;  $E_h^{(v)}$  is the energy registered in hadron block.  $n_h$  values are corrected on interaction probability in each chamber.

Tables I and II represent our experimental results, where  $\nu$  and  $\chi$  parametres defined as

$$\nu = \frac{\langle n_h \rangle_{Pb}}{\langle n_h \rangle_c} \quad (1) \quad ; \quad \chi = \frac{\langle E_h^{(v)} R \rangle_{Pb}}{\langle E_h^{(v)} R \rangle_c} \quad (2)$$

Table I

	$N_{fam}$	$\langle n_h \rangle$	$\beta - 1$
C	169	$3.0 \pm 0.2$	$1.0 \pm 0.1$
Pb	41	$4.2 \pm 0.3$	$1.1 \pm 0.1$
$\nu$	-	$1.4 \pm 0.2$	

Table II

	$N_{fam}$	$E_h^{(v)} R$ Tev mm
C	33	$350 \pm 30$
Pb	22	$455 \pm 70$
$\chi$	-	$1,3 \pm 0,2$

$N_{fam}$  is a number of families in C and Pb chambers fitted the above mentioned criteria,  $\beta - 1$  is a slope of integral energy spectrum of hadrons from families selected by A-criterion

In the present paper we connect the observed discrepancy between  $\langle n_h \rangle$  and  $\langle E_h^{(v)} R \rangle$  in C and Pb chambers with the difference in values of effective coefficients of energy transfer to the soft component  $K_{eff}$  for C and Pb chambers. The following considerations can be a ground of this suggestion:

1) It is known that a probability of hadron registration (in case of power like spectrum) is equal to  $\langle K_\gamma^{\beta-1} \rangle$ . Thus, the ratio of multiplicities

$$\nu = \frac{\langle n_h \rangle_{Pb}}{\langle n_h \rangle_c} = \frac{\langle K_\gamma^{\beta-1} \rangle_{Pb}}{\langle K_\gamma^{\beta-1} \rangle_c}$$

As  $\beta - 1 \approx 1$  (see Table I) we have:

$$\nu = \langle K_\gamma \rangle_{Pb} / \langle K_\gamma \rangle_c \quad (3)$$

2)  $\langle E_h^{(v)} R \rangle$  value can be written down as  $\langle K_\gamma \rangle \langle E_h R \rangle$  (in case of  $K_\gamma$  independent of  $E_h$ ). Hence:  $\chi = \langle K_\gamma \rangle_{Pb} / \langle K_\gamma \rangle_c$  (4).

To test this considerations some simulations on nuclear electromagnetic cascades in the atmosphere were made on the basis of fireball scaling model (S-model) [3]. Hadrons passage through the chamber was imitated in the simulations. For each hadron the value of  $K_\gamma$  coefficient was found by means of  $f(K_\gamma)$  distribution function taken in the form of  $f(K_\gamma) \sim K_\gamma^\alpha \cdot \exp(-K_\gamma/\beta)$  [4]. The parametres  $\alpha$  and  $\beta$  define the momenta of this distribution. Two series of simulation with various  $\langle K_\gamma \rangle$  values were carried out. One of them with  $\langle K_\gamma \rangle = 0.17$  imitated the situation in C-chamber, the other with  $\langle K_\gamma \rangle = 0.30$  in Pb-chamber. The latter value is greater than

the real one in Pb-chamber, was chosen especially to emphasize the effects connected with  $K_\gamma$  variations. The ratio of simulation values of  $\langle K_\gamma \rangle$  is equal to

$$\langle K_\gamma \rangle_{\text{Pb}} / \langle K_\gamma \rangle_{\text{C}} = 1.76 \quad (5)$$

A hadron was considered registered one in case of  $E_h^\gamma = K_\gamma E_h \geq 4 \text{ Tev}$ .

For  $\langle E_h^\gamma R \rangle$  determination we selected the families with number of registered hadrons  $n_h > 3$ . Table III represents the obtained results.

Table III

	$N_{\text{fam}}$	$E_h K$ Tev mm
C	126	285 $\pm$ 20
Pb	275	510 $\pm$ 15
$\chi$		1,80 $\pm$ 0.1

Table IV

	$N_{\text{fam}}$	$\langle n \rangle$	$\beta - 1$
C	505	2.3 $\pm$ 0.1	0.9 $\pm$ 0.05
Pb	505	3.9 $\pm$ 0.1	1.0 $\pm$ 0.05
$\gamma$		1.7 $\pm$ 0.1	

As it is seen from Table III

$$\chi = \langle E_h^{(\gamma)} R \rangle_{\text{Pb}} / \langle E_h^{(\gamma)} R \rangle_{\text{C}} = 1.80 \pm 0.10 \quad (6)$$

and that is in good agreement with value (5).

The incoming to the installation hadron families with  $n_h > 3$  and  $\sum E_h \geq 250 \text{ Tev}$  were selected to calculate  $\gamma$  parametre. The last condition provides registration of the same events in both sets of simulation (in "both chambers").

Table IV represents the obtained values of mean multiplicities and energy spectra slope.  $\gamma$  value which is equal to  $1.70 \pm 0.10$  is also in agreement with the given ratio of  $\langle K_\gamma \rangle$  (5). Thus, the experimental estimation of  $\chi$  and  $\gamma$  values in the two types of chambers gives an opportunity to get a ratio of  $\langle K_\gamma \rangle$  coefficients for lead and carbon nuclei.

Returning to preliminary experimental evaluations of  $\chi$  and  $\gamma$  parametres (Table 1) one can see that the coefficients ratio turned to be roughly equal to  $1.2 \pm 1.5$ . This ratio corresponds to  $K_\gamma$  coefficients for pion-nuclear interactions as the majority of hadrons in families are  $\pi$ -mesons.

Experiments carried out in the low energy range yield that  $K_{\pi^0}$  partial inelastic coefficient dependence on atomic number of target nucleus is described as  $K_{\pi^0} \sim A^\alpha$ . Using this dependence and  $\langle K_\gamma \rangle_{\text{Pb}} / \langle K_\gamma \rangle_{\text{C}}$  value we can estimate  $\alpha$ -index at energy  $\sim 10^4 \text{ ev}$ . Here we should bear in mind, that  $K_\gamma^{\text{Pb}} \approx K_{\pi^0}^{\text{Pb}}$ , while  $K_\gamma^{\text{C}}$  can considerably differ from  $K_{\pi^0}^{\text{C}}$ . However, it was shown in simulations [5]  $K_\gamma$  value in carbon chamber in  $E_h^{(\gamma)}$  energy range from 3 up to 30 Tev is similar to  $K_{\pi^0}$  value in pion-nucleon interactions. In such case

$$\frac{\langle K_\gamma \rangle_{\text{Pb}}}{\langle K_\gamma \rangle_{\text{C}}} \approx A^{\alpha} \text{ Pb} \quad \text{and} \quad \alpha \approx 0.06 \pm 0.02. \text{ This value is}$$

not in contradiction to  $\alpha$ -measureings at lower energies

and confirms a possibility of the suggested method.

To determine  $\alpha$ -value with 20% accuracy it is necessary to have approximately 100 families with  $\sum E_{\gamma} \geq 100$  Tev in each chamber as well as to make precise simulations of connection between partial  $K_{\pi^0}$  and effective  $K_{\gamma}$  inelasticity coefficients in chambers of both types.

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