

MODELING CURRENTS AT SATELLITE ALTITUDES

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In a conducting medium, the magnetic field cannot be written $\vec{B} = -\vec{\nabla}\psi$ because $\mu_0 \vec{J} = \vec{\nabla} \times \vec{B} \neq \vec{0}$. However, \vec{B} is still solenoidal, and any solenoidal field can be written $\vec{B} = \vec{\nabla} \times \vec{\Lambda} p + \vec{\Lambda} q$ where $\vec{\Lambda} = \vec{r} \times \vec{\nabla}$. Let $S(r)$ be the spherical surface of radius r centered on the origin, and let $\langle f \rangle_r$ be the average of the function f on $S(r)$. For each r the poloidal and toroidal fields, $\vec{\nabla} \times \vec{\Lambda} p$ and $\vec{\Lambda} q$, are uniquely determined on $S(r)$ by \vec{B} on $S(r)$, and p and q are determined up to arbitrary additive constants. These can always be chosen uniquely by demanding $\langle p \rangle_r = \langle q \rangle_r = 0$. A toroidal field is a solenoidal field without radial component. In the solid spherical shell between $S(a)$ and $S(b)$, $\vec{J} = \vec{0}$ if and only if $q = 0$ and $\nabla^2 p = 0$. Thus a vacuum field is poloidal, and its poloidal scalar is harmonic. In a source-free shell, the poloidal scalar p represents \vec{B} as economically as does the magnetic potential ψ ; and p has the advantage that it continues to be physically meaningful where $\vec{J} \neq \vec{0}$.

Being solenoidal, the current density \vec{J} can be analyzed into poloidal and toroidal parts, and in fact $\mu_0 \vec{J} = \vec{\nabla} \times \vec{B} = \vec{\nabla} \times \vec{\Lambda} q - \vec{\nabla}(\nabla^2 p)$. Thus the toroidal currents are the source of the poloidal magnetic field and the poloidal currents are the source of the toroidal magnetic field. For each r , the radial component of \vec{J} on $S(r)$ determines the toroidal part of \vec{B} on $S(r)$. If \vec{B} is known on $S(r)$ when J_r is determined there, as are the poloidal magnetic fields produced by the toroidal currents inside $S(r)$ and by those outside $S(r)$. The sources of the poloidal magnetic field on $S(r)$ are inside and outside $S(r)$, while the sources of the toroidal magnetic field on $S(r)$ are on $S(r)$ itself. If \vec{B} is known on $S(a)$ and $S(b)$ then the radial averages in $a < r < b$ of the toroidal current and the tangential component of the poloidal current can be determined.

Analysis of p and q into surface spherical harmonics can replace the conventional Gaussian analysis of the vacuum field. However, the radial dependence of the spherical harmonic coefficients for p and q is arbitrary in current-carrying regions unless some further physical hypothesis is introduced. At MAGSAT altitudes, a reasonable hypothesis is $\vec{J} \times \vec{B} = \vec{0}$ (field-aligned currents, or a force-free plasma). This hypothesis greatly reduces the space of field-models to be considered, and at MAGSAT altitudes it can be implemented by linear iteration with a vacuum field as the first step. One must recognize, however, that even with $\vec{J} \times \vec{B} = \vec{0}$ the magnetic effects of \vec{J} are non-local. Polar currents produce equatorial magnetic fields.