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# The Use of Interleaving for Reducing Radio Loss in Convolutionally Coded Systems 

D. Divsalar<br>Communications Systems Research Section<br>M. K. Simon<br>Telecommunications Systems Section<br>J. H. Yuan<br>Telecommunications Science and Engineering Division


#### Abstract

In-this-article the use of interleaving after convolutional coding and deinterleaving before Viterbi decoding is proposed. This effectively reduces radio loss at low-loop SNIt by several decibels and at high-loop SNRs by a few tenths of a decibel. Performance of the coded system can further be enhanced if the modulation index is optimized for this system. This will correspond to a reduction of bit SNR at a certain bit error rate for the overall system. The introduction of interleaving/deinterleaving into communication systems designed for future deep space missions does not substantially complicate their hardware design or increase their system cost.




## I. Introduction

In analyzing the performance of coherent receivers of convolutionally encoded phase modulation, one often makes the simplifying assumption that the reference signal used for demodulation is perfectly phase synchronized to the transmitted signal, ie., one assumes ideal coherent detection [1] .

In a practical receiver, such as that used by NASA's Deep Space Network, the coherent demodulation reference is derived from a carrier synchronization subsystem, e.g., a type of phase-locked loop or Costa loop, resulting in a performance degradation due to the phase error between the received signal and the locally generated reference. Since this subsystem
forms its demodulation reference from a noise-perturbed version of the transmitted signal, the phase error is a random process. The manner in which this process degrades the system error probability performance depends on the ratio of data rate to loop bandwidth, ie., the rate of variation of the phase error over the data symbol interval $[2,3]$.

In this article, the above degradation in signal-to-noise ratio (SNR) performance (often referred to as radio or noisy referene loss) is discussed, and it is shown how it can be reduced by employing interleaving/deinterleaving. Specific closed form results are derived for both discrete and suppressed carrier systems and the differences between the two are discussed and numerically illustrated.

## II. Upper Bound on the Average Bit Error Probability Performance

## A. Perfect Carrier Phase Synchronization

For a convolutionally encoded binary phase-shift-keyed (BPSK) modulation transmitted over a perfectly phasesynchronized additive white Gaussian noise (AWGN) channel (Fig. 1), an upper bound on the average bit error probability is given as [1]

$$
\begin{equation*}
P_{b} \leqslant \sum_{\underline{x}, \underline{\hat{x}} \in C} \sum a(\underline{x}, \underline{\hat{x}}) p(\underline{x}) P(\underline{x} \rightarrow \underline{\hat{x}}) \tag{1}
\end{equation*}
$$

where $a(\underline{x}, \underline{\hat{x}})$ is the number of bit errors that occurs when the sequence $\underline{x}$ is transmitted and the sequence $\underline{\hat{x}} \neq \underline{x}$ is chosen by the decoder, ${ }^{1} p(\underline{x})$ is the a priori probability of transmitting $\underline{x}$, and $C$ is the set of all coded sequences. Also, in Eq. (1), $\underline{P}(\underline{x} \rightarrow \underline{\hat{x}})$ represents the pairwise error probability, i.e., the probability that the decoder chooses $\underline{\widehat{x}}$ when indeed $\underline{x}$ was transmitted. The upper bound of Eq. (1) is efficiently evaluated using the transfer function bound approach [1].

In general, evaluation of the pairwise error probability depends on the proposed decoding metric, the presence or absence of channel state information (CSI), and the type of detection used, i.e., coherent versus differentially coherent. For the case of interest here, namely, coherent detection with no CSI and a Gaussian metric (optimum for the AWGN channel), it is well known [1] that the pairwise error probability is given by

$$
\begin{equation*}
P(\underline{x} \rightarrow \underline{\hat{x}}) \leqslant \exp \left\{-\frac{E_{s}}{N_{0}} d_{H}(\underline{x}, \underline{\hat{x}})\right\} \tag{2}
\end{equation*}
$$

where

$$
\begin{equation*}
d_{H}(\underline{x}, \underline{\hat{x}})=\frac{1}{4} \sum_{n}\left|x_{n}-\hat{x}_{n}\right|^{2} \tag{3}
\end{equation*}
$$

represents the Hamming distance between $\underline{x}$ and $\underline{\hat{x}}$, i.e., the number of symbols in which the sequences $\underline{x}$ (the correct one) and $\underline{\hat{x}}$ (the incorrect one) differ. In Eq. (2), $E_{s}$ is the energy-

[^0]per-output coded symbol and $N_{0}$ is the single-sided noise spectral density.

## B. Imperfect Carrier Phase Synchronization

1. Discrete Carrier (No Interleaving). When a carrier phase error $\phi(t)$ exists between the received signal and the locally generated demodulation reference, then the result in Eq. (2) is modified as follows.

Assuming the case in which the data symbol rate $1 / T_{s}$ is high compared to the loop bandwidth $B_{L}$, then $\phi(t)$ can be assumed constant (independent of time) over a number of symbols on the order of $1 / B_{L} T_{s}$. In this case, let $\phi(t)=\phi$. Since the decoder has no knowledge of $\phi$, the decoding metric can make no use of this information and as such is mismatched to the channel. Under these conditions, it can be shown (see the Appendix) that using the maximum-likelihood metric for a perfectly phase-synchronized AWGN, one obtains

$$
\begin{align*}
& P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda) \leqslant \\
& \left\{\begin{array}{l}
\frac{1}{2} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda(\cos \phi-\lambda) d_{H}(\underline{x}, \underline{\hat{x}})\right\} ; \quad 0 \leqslant|\phi| \leqslant \frac{\pi}{2} \\
1 ; \quad \frac{\pi}{2} \leqslant|\phi| \leqslant \pi
\end{array}\right. \tag{4}
\end{align*}
$$

where $\lambda \geqslant 0$ is a parameter to be optimized. Note that for $\phi=$ 0 , the expression $4 \lambda(1-\lambda)$ is maximized by the value $\lambda=1 / 2$, which when substituted in Eq. (4) yields Eq. (2) as it should.

Letting $p(\phi)$ denote the probability density function (p.d.f.) of the phase error $\phi$, the average bit error probability is upper bounded by ${ }^{2}$
where $E_{\phi}\{\cdot\}$ denotes statistical averaging over the p.d.f. $p(\phi)$. Using Eq. (4), the statistical average required in Eq. (5) becomes

[^1]\[

$$
\begin{align*}
& E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\}= \\
& \quad \int_{-\pi / 2}^{\pi / 2} \frac{1}{2} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda(\cos \phi-\lambda) d_{H}(\underline{x}, \underline{\hat{x}})\right\} p(\phi) d \phi \\
& \quad+2 \int_{\pi / 2}^{\pi} p(\phi) d \phi \tag{6}
\end{align*}
$$
\]

2. Discrete Carrier (With Interleaving). Ordinarily, one thinks of using interleaving/deinterleaving to break up the effects of error bursts in coded communication systems. One can gain an intuitive notion of how it may be applied in systems with noisy carrier phase reference by considering the $\cos \phi$ degradation factor as an "amplitude fade" whose duration is on the order of $1 / B_{L} T_{s}$ symbols. Thus, if we break up this "fade" by interleaving to a depth on the order of $1 / B_{L} T_{s}$, then, after deinterleaving, the degradation due to $\cos \phi$ will be essentially independent from symbol to symbol. From a mathematical standpoint, this is equivalent to replacing Eq. (4) by ${ }^{3}$
$P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda) \leqslant$

$$
\left.\left\{\begin{array}{l}
\frac{1}{2} \exp \left\{-\frac{E_{s}}{N_{0}} \sum_{n=1}^{d_{H}} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right. \tag{7}
\end{array}\right\} ; \sum_{n=1}^{d_{H}} \cos \phi_{n}>0\right\}
$$

where the $\phi_{n}$ variables are independent identically distributed (i.i.d.) random variables with p.d.f. $p(\phi)$, and $\phi$ refers to the vector whose components are $\phi_{n}$ s. The derivation of Eq. (7) is given in the Appendix. The expectation required in Eq. (5) now involves computation of multidimensional integrals over regions of $\phi$ corresponding to the inequalities in Eq. (7). In these regions, since the intervals of integration per dimension are dependent on one another, the expectation required in Eq. (5) is extremely difficult to compute.

Thus, instead we turn to a looser upper bound on conditional pairwise error probability, which has the advantage of not having to separate the multidimensional integration required in Eq. (5) into two disjoint regions. Indeed, it is straight-

[^2]forward to see that the right-hand side of Eq. (7) is upper bounded by the exponential in its first line (without the factor of $1 / 2$ ) over the entire domain of $\phi$, i.e., $\left\{\phi_{n} \in(-\pi, \pi) ; n \in \eta\right\}$. Hence,
\[

$$
\begin{aligned}
P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda) & \leqslant \exp \left\{-\frac{E_{s}}{N_{0}} \sum_{n=1}^{d_{H}} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right\} \\
& =\prod_{n=1}^{d_{H}} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right\}
\end{aligned}
$$
\]

which is identically equal to the Chernoff bound. Now, substituting Eq. (8) into Eq. (5) gives the much simpler result

$$
\begin{align*}
E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\}= & \prod_{n=1}^{d_{H}} \int_{-\pi}^{\pi} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right\} \\
& \times p\left(\phi_{n}\right) d \phi_{n} \tag{9}
\end{align*}
$$

3. Suppressed Carrier (No Interleaving). When the carrier synchronization loop used to track the input phase is of the suppressed carrier type (e.g., a Costas loop), then the results of Section B. 1 have to be somewhat modified since the appropriate domain for $\phi$ is no longer ( $-\pi, \pi$ ). In fact, for suppressed carrier tracking of BPSK with a Costas-type loop, and assuming perfect phase ambiguity resolution, $\phi$ takes on values only in the interval $(-\pi / 2, \pi / 2)$ [2]. Thus, the interval of integration for the first integral in Eq. (6) becomes $(-\pi / 2, \pi / 2)$ and the second integral in Eq. (6) disappears, i.e.,

$$
\begin{align*}
& E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\}= \\
&  \tag{10}\\
& \quad \int_{-\pi / M}^{\pi / M} \frac{1}{2} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda(\cos \phi-\lambda) d_{H}(\underline{x}, \underline{\hat{x}})\right\} p(\phi) d \phi
\end{align*}
$$

The significance of the second integral in Eq. (6) being equal to zero will be mentioned shortly relative to a discussion of irreducible error probability.
4. Suppressed Carrier (With Interleaving). Once again assuming suppressed carrier tracking of BPSK with a Costas-type
loop, and perfect phase ambiguity resolution, one obtains, analogous to Eq. (9), ${ }^{4}$

$$
\begin{align*}
& E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\}= \\
& \quad \frac{1}{2} \prod_{n=1}^{d_{H}} \int_{-\pi / 2}^{\pi / 2} \exp \left\{-\frac{E_{s}}{N_{0}} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right\} p\left(\phi_{n}\right) d \phi_{n} \tag{11}
\end{align*}
$$

## III. Carrier Synchronization Loop Statistical Model and Average Pairwise Error Probability Evaluation

To evaluate Eq. (5), using Eqs. (6), (9), (10), or (11), one must specify the functional form of the probability density function (p.d.f.) $p(\phi)$ of the modulo $2 \pi$ reduced phase error $\phi$. For a discrete carrier synchronization loop of the phaselocked type, $p(\phi)$ is given by the Tikhonov p.d.f. [2]

$$
p(\phi)= \begin{cases}\frac{\exp (\rho \cos \phi)}{2 \pi I_{0}(\rho)} ; & |\phi| \leqslant \pi  \tag{12}\\ 0 ; \quad \text { otherwise }\end{cases}
$$

where $\rho$ is the SNR in the loop bandwidth.
In order to allow evaluation of Eq. (5) in closed form, one, must recognize that for the case of no interleaving, Eq. (6) can be further upper bounded by using $(-\pi, \pi)$ instead of $(-\pi / 2$, $\pi / 2$ ) in the first integral. Then, making this replacement

$$
\begin{align*}
& \min _{\lambda} E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\} \leqslant \\
& \min _{\lambda}\left(\frac{1}{2} \exp \left\{4 d_{H} \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{I_{0}\left(\rho_{A}\right)}{I_{0}(\rho)}+I\right)  \tag{13}\\
& \rho_{A} \triangleq \rho-4 d_{H} \lambda \frac{E_{s}}{N_{0}} \\
& I=\frac{1}{\pi I_{0}(\rho)}\left[\int_{\pi / 2}^{\pi} \exp (\rho \cos \phi) d \phi\right]
\end{align*}
$$

[^3]When Eq. (13) is substituted into Eq. (5), the term $I$ will contribute an irreducible error probability, i.e., the system will exhibit a finite error probability when $\rho$ is held fixed and $E_{s} / N_{0}$ approaches infinity. An example of such a system is one which apportions a fixed amount of the total available input power to a discrete carrier component for the purpose of deriving a coherent carrier reference at the receiver.

When interleaving is employed, Eq. (9) (minimized over $\lambda$ ) together with Eq. (12) become

$$
\begin{align*}
& \min _{\lambda} E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \Phi ; \lambda)\} \leqslant \\
& \min _{\lambda}\left\{\prod_{n \in \eta} \exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{I_{0}\left(\rho_{B}\right)}{I_{0}(\rho)}\right\}  \tag{14}\\
&=\min _{\lambda}\left\{\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{I_{0}\left(\rho_{B}\right)}{I_{0}(\rho)}\right)^{d_{H}}\right\} \\
& \rho_{B} \triangleq \rho-4 \lambda \frac{E_{s}}{N_{0}}
\end{align*}
$$

For suppressed carrier tracking with a biphase Costas loop, $p(\phi)$ again has a Tikhonov-type p.d.f., which is given by [2] as

$$
p(\phi)= \begin{cases}\frac{\exp (\rho \cos 2 \phi)}{\pi I_{0}(\rho)} ; & |\phi| \leqslant \frac{\pi}{2}  \tag{15}\\ 0 ; \quad \text { otherwise }\end{cases}
$$

Here $\rho$ is the "effective" loop SNR which includes the effects of $S \times S, S \times N$, and $N \times N$ degradations commonly referred to as "squaring loss." Since suppressed carrier systems of this type derive their carrier demodulation reference from the databearing signal, the loop $\operatorname{SNR}, \rho$, is directly proportional to $E_{s} / N_{0}$; thus there can be no irreducible error probability since $\rho \rightarrow \infty$ when $E_{s} / N_{0} \rightarrow \infty$. Furthermore, for perfect phase ambiguity resolution, it has been shown previously that for no interleaving, the term $I$ is identically zero since the p.d.f. is zero in the region $(\pi / 2, \pi)$. Thus, the average pairwise error probability results become

$$
\begin{gather*}
\min _{\lambda} E_{\phi}\{P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\} \leqslant \min _{\lambda}\left\{\frac { 1 } { 2 } \operatorname { e x p } \left\{4 d_{H} \lambda^{2} \frac{E_{s}}{N_{0}}\left\{\frac{f_{A}(\rho)}{I_{0}(\rho)}\right\}\right.\right. \\
f_{A}(\rho) \triangleq \frac{1}{\pi} \int_{-\pi / 2}^{\pi / 2} \exp \left\{\rho \cos 2 \phi-4 d_{H} \lambda \frac{E_{s}}{N_{0}} \cos \phi\right\} d \phi \tag{16}
\end{gather*}
$$

for no interleaving and

$$
\begin{align*}
& \min _{\lambda} E_{\phi}\{P(\underline{x} \rightarrow \underline{\widehat{x}} \mid \underline{\phi} ; \lambda)\} \leqslant \min _{\lambda}\left\{\prod_{n \in \eta} \exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{f_{B}(\rho)}{I_{0}(\rho)}\right\} \\
&=\min _{\lambda}\left\{\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\left(\frac{f_{B}(\rho)}{I_{0}(\rho)}\right)^{d_{H}}\right\}\right.\right. \\
& f_{B}(\rho) \triangleq \frac{1}{\pi} \int_{-\pi / 2}^{\pi / 2} \exp \left\{\rho \cos 2 \phi-4 \lambda \frac{E_{s}}{N_{0}} \cos \phi\right\} d \phi \tag{17}
\end{align*}
$$

for the case of interleaving.
In arriving at Eqs. (13), (14), (16), and (17), we have assumed the "same type" of Chernoff bound in the sense that in all cases, the minimization over $\lambda$ was performed after the averaging over $\phi$. The principal reason for doing this is to allow comparison of performance with and without interleaving using bounds with "similar degrees of looseness." For the case of no interleaving, one can actually achieve a tighter bound than that given above by performing the minimization over $\lambda$ on the conditional pairwise probability in Eq. (4). When this is done, one obtains

$$
\begin{equation*}
\lambda_{\mathrm{opt}}=\frac{1}{2} \cos \phi \tag{18}
\end{equation*}
$$

and Eq. (4) becomes
$P(\underline{x} \rightarrow \underline{\widehat{x}} \mid \phi) \leqslant\left\{\begin{array}{l}\frac{1}{2} \exp \left\{-\frac{E_{s}}{4 N_{0}} d_{H} \cos ^{2} \phi\right\} ; \quad 0 \leqslant|\phi| \leqslant \frac{\pi}{2} \\ 1 ; \quad \frac{\pi}{2} \leqslant|\phi| \leqslant \pi\end{array}\right.$

Unfortunately, the integral of Eq. (19) over the p.d.f.s of Eqs. (12) and (15) cannot be obtained in closed form. Defining the integral

$$
\begin{equation*}
L(J)=\int_{-\pi / J}^{\pi / J} \exp \left\{-\frac{E_{s} d_{H}}{N_{0}} \cos ^{2} \phi\right\} \frac{\exp (\rho \cos J \phi)}{\left(\frac{2 \pi}{J}\right) I_{0}(\rho)} d \phi \tag{20}
\end{equation*}
$$

then, the average pairwise error probabilities are now as follows:

## Discrete Carrier

$$
\begin{equation*}
E_{\phi}\left\{\min _{\lambda} P(\underline{x} \rightarrow \underline{\hat{x}} \mid \Phi ; \lambda)\right\} \leqslant \frac{1}{2} L(1)+1 \tag{21}
\end{equation*}
$$

where $I$ is defined in Eq. (13).

## Suppressed Carrier

$$
\begin{equation*}
E_{\phi}\left\{\min _{\lambda} P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda)\right\} \leqslant \frac{1}{2} L(2) \tag{22}
\end{equation*}
$$

Using Eqs. (21) and (22) (rather than Eqs. 13 and 16) will result in a smaller improvement in performance due to interleaving/deinterleaving since Eqs. (21) and (22) result in a tighter bound on $P_{b}$ (no interleaving).

## IV. Evaluation of Bit Error Probability

## A. Discrete Carrier Tracking

To evaluate the upper bound on bit error probability, e.g., Eq. (5), we use the transfer function bound approach [1], which, for the ideal case of perfect carrier synchronization, gives

$$
\begin{equation*}
P_{b} \leqslant\left.\frac{1}{2} \frac{d}{d z} T(D, z)\right|_{z=1} \tag{23}
\end{equation*}
$$

where $T(D, z)$ is the transfer function of the pair-state diagram associated with the trellis diagram of the code. When noisy carrier synchronization references are present, the appropriate upper bound on bit error probability for the case of no interleaving becomes

$$
P_{b} \leqslant \frac{1}{2}\left[\left.2 \int_{0}^{\pi / 2} \frac{d}{d z} T(D(\phi), z)\right|_{z=1} p(\phi) d \phi\right]
$$

$$
+2 \int_{\pi / 2}^{\pi} p(\phi) d \phi
$$

$$
=\left.\int_{0}^{\pi / 2} \frac{d}{d z} T(D(\phi), z)\right|_{z=1} p(\phi) d \phi+I
$$

where $I$ is defined in Eq. (13) and from Eq. (19), the equivalent Bhattacharyya parameter becomes

$$
\begin{equation*}
D(\phi)=\exp \left\{-\frac{E_{s}}{N_{0}} \cos ^{2} \phi\right\} ; \quad 0 \leqslant|\phi| \leqslant \frac{\pi}{2} \tag{25}
\end{equation*}
$$

For the case of interleaving, we use Eq. (23) without the factor $1 / 2$ and with $D$ defined in accordance with Eq. (14), namely,

$$
\begin{equation*}
D=\min _{\lambda}\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{I_{0}\left(\rho_{B}\right)}{I_{0}(\rho)}\right) \tag{26}
\end{equation*}
$$

In arriving at Eq. (26), we have made use of the fact that, for any $d$,

$$
\begin{align*}
& \min _{\lambda}\left\{\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}} \frac{I_{0}\left(\rho_{B}\right)}{I_{0}(\rho)}\right)^{d}\right\}=\right. \\
&\left\{\min _{\lambda}\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{I_{0}\left(\rho_{B}\right)}{I_{0}(\rho)}\right)\right\}^{d} \tag{27}
\end{align*}
$$

Figures 2 through 10 illustrate the upper bound on bit error probability versus bit energy-to-noise ratio $E_{b} / N_{0}$ ( $E_{b}$ is related to $E_{s}$ by $E_{s}=r E_{b}$ where $r$ is the code rate) for the rate $1 / 2$, constraint length 7 Voyager code, the rate $1 / 4$, constraint length 15 Galileo experimental code [4], the rate $1 / 6$, constraint length 15 " 2 dB " code [5], and various values of loop SNR, $\rho$. In these curves we have assumed carrier synchronization with a PLL, i.e., the discrete carrier case. ${ }^{5}$ On each figure are plotted the results for the case of no interleaving, the case of interleaving, and the case of no radio loss, i.e., ideal carrier synchronization. The transfer function bounds (truncated to 15 terms) on $P_{b}$ for the above three codes are given in Table 1.

One may observe from the results in Figs. 2 through 10 that even for large values of $\rho$, e.g., 13 dB , a substantial reduction of bit error rate is possible by using interleaving. Also, since the tighter bound was used for the no-interleaving case and the looser bound for the interleaving case, the performance improve-

[^4]ment illustrated is, as previously mentioned, somewhat pessimistic, i.e., in reality, one will do even better than shown.

## B. Suppressed Carrier Tracking

The upper bound on bit error probability for the no-interleaving case is analogous to Eq. (23), and is given by

$$
\begin{equation*}
P_{b} \leqslant\left.\int_{0}^{\pi / 2} \frac{d}{d z} T(D(\phi), z)\right|_{z=1} p(\phi) d \phi \tag{28}
\end{equation*}
$$

with $p(\phi)$ as in Eq. (15) and $D(\phi)$ as in Eq. (25). For the case of interleaving and the tighter upper bound, we again use Eq. (23) with, however, $D$ now defined analogous to Eq. (26) by

$$
\begin{equation*}
D=\min _{\lambda}\left(\exp \left\{4 \lambda^{2} \frac{E_{s}}{N_{0}}\right\} \frac{f_{B}(\rho)}{I_{0}(\rho)}\right) \tag{29}
\end{equation*}
$$

where $f_{B}(\rho)$ is given in Eq. (17).

Assuming a Costas loop with integrate-and-dump arm filters (matched filters), the equivalent loop SNR is given by [2, 3]

$$
\rho=\frac{1}{4}\left(\frac{S}{N_{0} B_{L}}\right) S_{L}=\frac{1}{4}\left(\frac{E_{b}}{N_{0}}\right)\left(\frac{1}{B_{L} T_{b}}\right) S_{L} ; \quad S_{L}=\frac{2 \frac{E_{s}}{N_{0}}}{1+2 \frac{E_{s}}{N_{0}}}
$$

where $S_{L}$ denotes the "squaring loss" associated with BPSK modulation. Figures 11 through 18 illustrate results analogous to Figs. 2 through 10 for the case of suppressed carrier tracking and various values of $B_{L} T_{b}$. Since, as already mentioned, in suppressed carrier systems there is no irreducible error (since, from Eq. (30), $\rho \rightarrow \infty$ as $E_{b} / N_{0} \rightarrow \infty$ ), the noisy reference losses are much smaller to begin with (i.e., no interleaving) than for the discrete carrier case. Thus, for sufficiently small $B_{L} T_{b}$, interleaving does not provide significant improvement. We also observe from the above figures that the noisy reference losses are much larger for the rate $1 / 4$ and rate $1 / 6$ codes than for the rate $1 / 2$ code. The reason for this is that for a given $E_{b} / N_{0}$, the value of $E_{s} / N_{0}$ is 3 dB smaller for the rate $1 / 4$ code and 4.77 dB smaller for the rate $1 / 6$ code than for the rate $1 / 2$ code and thus, from Eq. (30), for the same value of $B_{L} T_{h}$, the equivalent loop SNR is smaller because of the increased squaring loss.

## V. Concluding Remarks

It has been shown that by interleaving the transmitted coded bits in convolutionally coded systems the radio loss can be significantly reduced. The amount of this reduction depends on
the particular convolutional code used and the region of operation of the system as characterized by such parameters as bit error rate and loop SNR. In some cases, the performance of the interleaved system is close to that of the ideal coherent detection case.

## Acknowledgment

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Table 1. Transfer function bounds of $r=1 / 2, K=7, r=1 / 4$, $K=15$, and $r=1 / 6, K=15$ convolutional codes

${ }^{\text {a }}$ This code has been incorporated as an experiment for the Galileo mission. It is a good code for a concatenated coding scheme but not the optimum code from the standpoint of maximizing $d_{f}$.


Fig. 1. System block diagram.


Fig. 2. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 2$, constraint length 7 convolutional code; loop SNR = 7 dB; discrete carrier.


Fig. 3. Upper bound on average bit error probablity versus bit energy-to-noise ratio for rate $1 / 2$, constraint length 7 convolutional code; loop SNR = 10 dB ; discrete carrier.


Fig. 4. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 2$, constraint length 7 convolutional code; loop SNR = 13 dB ; discrete carrier.


Fig. 5. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 4$, constraint length 15 convolutional code; loop SNR = 7 dB ; discrete carrier.


Fig. 6. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 4$, constraint length 15 convolutional code; loop SNR = 10 dB ; discrete carrier.


Fig. 7. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 4$, constraint length 15 convolutional code; loop SNR = 13 dB ; discrete carrier.


Fig. 8. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 6$, constraint length 15 convolutional code; loop SNR = 7 dB ; discrete carrier.


Fig. 9. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 6$, constraint length 15 convolutional code; loop SNR = $\mathbf{1 0} \mathbf{~ d B}$; discrete carrier.


Fig. 10. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate 1/6, constraint length 15 convolutional code; loop SNR = 13 dB ; discrete carrier.


Fig. 11. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 2$, constraint length 7 convolutional code; suppressed carrier; $\mathbf{1} / B_{L} T_{b}=10$.


Fig. 12. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 2$, constraint length 7 convolutional code; suppressed carrier; $1 / B_{L} T_{b}=20$.


Fig. 13. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 4$, constraint length 15 convolutional code; suppressed carrier; $1 / B_{L} \boldsymbol{T}_{b}=10$.


Fig. 14. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 4$, constraint length 15 convolutional code; suppressed carrier; $1 / B_{L} T_{b}=20$.


Fig. 15. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate 1/4, constraint length 15 convolutional code; suppressed carrier; $1 / B_{L} T_{b}=40$.


Fig. 16. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 6$, constraint length 15 convolutional code; suppressed carrier; $1 / B_{L} T_{b}=10$.


Fig. 17. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 6$, constraint length 15 convolutional code; suppressed carrier; $\mathbf{1 / B} \boldsymbol{B}_{\boldsymbol{L}}=\mathbf{2 0}$.


Fig. 18. Upper bound on average bit error probability versus bit energy-to-noise ratio for rate $1 / 6$, constraint length 15 convolutional code; suppressed carrier; $1 / B_{L} T_{b}=60$.

## Appendix

## Derivation of an Upper Bound on the Painwise Error Probability for Convolutionally Coded BPSK With Imperfect Carrier Phase Reference

Let $\underline{y}=\left(y_{1}, y_{2}, \ldots, y_{N}\right)$ denote the received sequence when the normalized (to unit power) sequence of MPSK symbols $\underline{x}=\left(x_{1}, x_{2}, \ldots, x_{N}\right)$ is transmitted. A pairwise error occurs if $\underline{\hat{x}}=\left(\hat{x}_{1}, \hat{x}_{2}, \ldots, \hat{x}_{N}\right) \neq \underline{x}$ is chosen by the receiver, which, if the receiver uses a distance metric to make this decision, implies $\underline{y}$ is closer to $\hat{x}$ than to $\underline{x}$. Assuming that distance metric which is maximum-likelihood for ideal coherent detection (perfect carrier phase reference), then such an error occurs whenever

$$
\begin{equation*}
\sum_{n=1}^{N}\left|y_{n}-\hat{x}_{n}\right|^{2}<\sum_{n=1}^{N}\left|y_{n}-x_{n}\right|^{2} \tag{A-1}
\end{equation*}
$$

Since BPSK is a constant envelope signaling set, we have $\left|x_{n}\right|^{2}=\left|\widehat{x}_{n}\right|^{2}=1$ (assuming a normalized signal) and Eq. (A-1) reduces to

$$
\begin{equation*}
\sum_{n=1}^{N} \operatorname{Re}\left\{y_{n} \hat{x}_{n}^{*}\right\}>\sum_{n=1}^{N} \operatorname{Re}\left\{y_{n} x_{n}^{*}\right\} \tag{A-2}
\end{equation*}
$$

Letting $n_{n}$ represent the additive noise in the $n$th signaling interval, and $\phi_{n}$ the phase shift introduced by imperfect carrier demodulation in that same interval, then $y_{n}$ and $x_{n}$ are related by

$$
\begin{equation*}
y_{n}=x_{n} e^{j \phi_{n}}+n_{n} ; \quad n=1,2, \ldots, N \tag{A-3}
\end{equation*}
$$

Substituting Eq. (A-3) into Eq. (A-2) and simplifying gives

$$
\begin{equation*}
\operatorname{Re}\left\{\sum_{n \in \eta}\left(\hat{x}_{n}-x_{n}\right)^{*} n_{n}\right\}>\operatorname{Re}\left\{\sum_{n \in \eta} x_{n}\left(x_{n}-\hat{x}_{n}\right)^{*} e^{j \phi_{n}}\right\} \tag{A-4}
\end{equation*}
$$

where $\eta$ is the set of all $n$ such that $x_{n} \neq \hat{x}_{n}$.

Since for an AWGN channel, $n_{n}$ is a complex Gaussian random variable whose real and imaginary components have variance

$$
\begin{equation*}
E\left\{\left[\operatorname{Re}\left(n_{n}\right)\right]^{2}\right\}=E\left\{\left[\operatorname{Im}\left(n_{n}\right)\right]^{2}\right\} \triangleq \sigma^{2} \tag{A-5}
\end{equation*}
$$

then

$$
\begin{equation*}
\operatorname{var}\left\{\operatorname{Re}\left\{\sum_{n \in \eta}\left(\hat{x}_{n}-x_{n}\right)^{*} n_{n}\right\}\right\}=\sigma^{2} \sum_{n \in \eta}\left|x_{n}-\hat{x}_{n}\right|^{2} \tag{A-6}
\end{equation*}
$$

and the conditional pairwise error probability $P(\underline{x} \rightarrow \underline{\hat{x}} \mid \underline{\phi})$ is given by

$$
\begin{align*}
P(\underline{x} \rightarrow \underline{\hat{x}} \mid \underline{\phi}) & =\operatorname{Pr}\left\{\operatorname{Re}\left\{\sum_{n \in \eta}\left(\hat{x}_{n}-x_{n}\right)^{*} n_{n}\right\}\right. \\
& \left.>\operatorname{Re}\left\{\sum_{n \in \eta} x_{n}\left(x_{n}-\hat{x}_{n}\right)^{*} e^{j \phi_{n}}\right\}\right\} \\
& =Q\left\{\frac{\operatorname{Re}\left\{\sum_{n \in \eta} x_{n}\left(x_{n}-\hat{x}_{n}\right)^{*} e^{j \phi_{n}}\right\}}{\sigma \sqrt{\sum_{n \in \eta}\left|x_{n}-\hat{x}_{n}\right|^{2}}}\right\} \tag{A-7}
\end{align*}
$$

where $\phi=\left(\phi_{1}, \phi_{2}, \ldots, \phi_{N}\right)$ is the sequence of carrier phase errors and $Q(x)$ is the Gaussian integral defined by

$$
\begin{equation*}
Q(x)=\frac{1}{\sqrt{2 \pi}} \int_{x}^{\infty} \exp \left(-\frac{y^{2}}{2}\right) d y \tag{A-8}
\end{equation*}
$$

To simplify Eq. (A-7), proceed as follows. For BPSK modulation

$$
\begin{equation*}
\left|x_{n}-\hat{x}_{n}\right|^{2}=4 ; \quad n \in \eta \tag{A-9}
\end{equation*}
$$

and

$$
\begin{equation*}
\sum_{n \in \eta} \operatorname{Re}\left\{x_{n}\left(x_{n}-\hat{x}_{n}\right)^{*} e^{j \phi_{n}}\right\}=2 \sum_{n \in \eta} \cos \phi_{n} \tag{A-10}
\end{equation*}
$$

Thus, Eq. (A-7) can be written as

$$
\begin{equation*}
P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi)=Q\left(\frac{\sum_{n \in \eta} \cos \phi_{n}}{\sqrt{\sigma^{2} d_{H}(\underline{x}, \underline{\hat{x}})}}\right) \tag{A-11}
\end{equation*}
$$

where $d_{H}(\underline{x}, \underline{\hat{x}})$ is the Hamming distance between $\underline{x}$ and $\underline{\hat{x}}$, i.e., the number of elements in the set $\eta$ (see Eq. 3).

The argument of the Gaussian integral in Eq. (A-11) is in the form $a / \sqrt{b}$. For $a>0$, we can upper bound this integral by ${ }^{1}$

$$
\begin{equation*}
Q\left(\frac{a}{\sqrt{b}}\right) \leqslant \frac{1}{2} e^{-a^{2} / 2 b} \tag{A-12}
\end{equation*}
$$

Since for any $\lambda$, we have $(a-2 \lambda b)^{2}>0$, rearranging this inequality gives the equivalent form

$$
\begin{equation*}
\frac{a^{2}}{b} \geqslant 4 \lambda a-4 \lambda^{2} b \tag{A-13}
\end{equation*}
$$

[^5]Thus, for $a>0$,

$$
\begin{equation*}
Q\left(\frac{a}{\sqrt{b}}\right) \leqslant \frac{1}{2} \exp \{-2 \lambda[a-\lambda b]\} \tag{A-14}
\end{equation*}
$$

For $a<0$, the loose upper bound must be used

$$
\begin{equation*}
Q\left(\frac{a}{\sqrt{b}}\right)=Q\left(-\frac{|a|}{\sqrt{b}}\right)=1-Q\left(\frac{|a|}{\sqrt{b}}\right) \leqslant 1 \tag{A-15}
\end{equation*}
$$

Finally, using Eqs. (A-14) and (A-15) in Eq. (A-11) gives the desired upper bound on pairwise error probability as

$$
\begin{align*}
& P(\underline{x} \rightarrow \underline{\hat{x}} \mid \phi ; \lambda) \leqslant \\
& \left\{\begin{array}{l}
\frac{1}{2} \exp \left\{-\frac{E_{s}}{N_{0}} \sum_{n \in \eta} 4 \lambda\left(\cos \phi_{n}-\lambda\right)\right\} ; \quad \sum_{n \in \eta} \cos \phi_{n}>0 \\
1 ; \quad \sum_{n \in \eta} \cos \phi_{n} \leqslant 0
\end{array}\right. \tag{A-16}
\end{align*}
$$

In Eq. (A-16), use is made of the fact that for the unnormalized system, $1 / 2 \sigma^{2}=E_{s} / N_{0}$ where $E_{s}$ is the symbol energy and $N_{0}$ the noise spectral density, and $\lambda$ is replaced by the normalized quantity $\lambda \sigma^{2}$. Also, note that if Eq. (A-16) is minimized over $\lambda$, then it is identically in the form of a Chernoff bound.


[^0]:    ${ }^{1}$ For simplicity of notation, the sequences $\underline{x}$ and $\underline{\hat{x}}$ are assumed to be normalized such that their elements take on values $\pm 1$.

[^1]:    ${ }^{2}$ Later on a tighter bound for this case is presented by optimizing on $\lambda$ prior to performing the expectation over $\phi$.

[^2]:    ${ }^{3}$ Herein, for simplicity of notation, we drop the dependence of $d_{H}$ on $\underline{x}$ and $\underline{\hat{x}}$.

[^3]:    ${ }^{4}$ Note that the factor of $1 / 2$ can be included here since for $0 \leqslant\left|\phi_{\boldsymbol{n}}\right|$ $\leqslant \pi / 2 ; n \in \eta$, the condition on the first line of Eq. (7) is always satisfied and thus we need not use the looser upper bound of Eq. (8).

[^4]:    ${ }^{5}$ Also, in the computation of Eq. (24), we have set the value of onehalf the derivative of the transfer function evaluated at $z=1$ to one whenever it would normally exceed one. This is allowable since the conditional error probability cannot exceed one. Doing so results in a tighter bound.

[^5]:    ${ }^{1}$ Note that for perfect carrier demodulation, i.e., $\phi=0$, we always have $a>0$.

