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Nimbus-7 Stratospheric and Mesospheric Sounder (SAMS) Experiment Data User's Guide

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National Aeronautics and Space Administration Office of Management Scientific and Technical Information Division

#### FOREWORD

3-7 Stratospheric and Mesospheric Sounder (SAMS) Experiment Data User's Guide is intended to provide immunity with the background information necessary for understanding and using data products on SAMS ved Temperature Tapes (GRID-T), and Zonal Mean Methane and Nitrous Oxide Tapes (ZMT-G). The ent was flown aboard the Nimbus-7 spacecraft and collected data from October 26, 1978 through June 9, ument provides users with information concerning the operational principles of the SAMS instrument and the retrieval of temperature and atmospheric constituents, and the scientific validity of SAMS data. nts that influence the quality of data are included along with the mission history. Data formats of the MT-G tape products and descriptions of SAMS data are also given.

al discussion in this document was prepared originally by the SAMS processing team at Oxford, England. nvestigator was Dr. F. W. Taylor and the NET chairman was Dr. C. D. Rodgers. All questions of a should be addressed to these individuals, at the address given in Section 1.5. The description of the tape ed upon the data tapes provided for conversion from the DEC format available from Oxford to the IBM by the Nimbus project. The text provided by the SAMS processing team and the information gained from the data tapes were compiled into this document for NASA by S. T. Nutter and N. Oslik of ST Systems TX) under contract NAS5-28063.

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### TABLE OF CONTENTS

- - - - -

Section	Page										
	FOREWORD										
1	PURPOSE, INTRODUCTION, AND OBJECTIVES										
	1.1Purpose of the Data User's Guide11.2Introduction11.3Objectives21.4Data Products21.5Principal Investigators3										
2	INSTRUMENT DESCRIPTION										
3	DATA PROCESSING										
	3.1       Calibration										
4	VALIDATION										
5	MISSION HISTORY AND CHANNEL STATUS										
6	DATA FORMATS										
	<ul> <li>6.1 Format of SAMS Gridded Retrieved Temperature Tapes (GRID-T)</li></ul>										
REFERENCE	ES										
LIST OF AC	RONYMS, INITIALS, AND ABBREVIATIONS										
Appendixes											
A B C D	SEQUENCE NUMBERS OF SAMS DATA TAPES       71         FORTRAN SOURCE LISTING OF SAMS TAPE DUMP       73         STRATOSPHERIC AND MESOSPHERIC SOUNDER ON NIMBUS-7       89         RETRIEVAL OF TEMPERATURE AND COMPOSITION FROM       113         NIMBUS-7 SAMS MEASUREMENTS       113										
Е	TEMPERATURE COMPARISONS BETWEEN THE NIMBUS-7 SAMS, ROCKET/RADIOSONDES AND THE NOAA 6 SSU										
F	OBSERVATIONS OF CH4 AND N2O BY THE NIMBUS-7 SAMS: A COMPARISON WITH IN SITU DATA AND TWO-DIMENSIONAL										
G	NUMERICAL MODEL CALCULATIONS										

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#### **SECTION 1**

### PURPOSE, INTRODUCTION, AND OBJECTIVES

### 1.1 Purpose of the Data User's Guide

The SAMS Data User's Guide is intended to provide the scientific community with the background information needed to access data on SAMS GRID-T and ZMT-G tapes. The document describes the SAMS instrument, the method of temperature and constituent retrieval, data calibration, and data validation.

Section 1 of this document explains the scientific and technical objectives of the SAMS instrument. Section 2 briefly discusses the literature from which a description of the instrument can be obtained. Reprints of two relevant papers appear as Appendixes C and G to this User's Guide. Section 3 outlines the data processing system, explains instrument calibration, briefly describes the temperature and constituent retrieval process, and explains the temperature gridding procedure. Section 4, like Section 2, provides an introduction to literature that discusses the accuracy of the retrievals and compares the SAMS results with those of other measurements. Reprints of two relevant papers are included as Appendixes E and F of this User's Guide. Important dates in the history of the SAMS instrument and information concerning the status of SAMS channels are given in Section 5. The data formats of the SAMS GRID-T and ZMT-G tapes are described in detail in Section 6. A FORTRAN source listing of a SAMS tape dump program is included in Appendix B. This program enables users to access data on GRID-T and ZMT-G tapes.

### **1.2 Introduction**

The Stratospheric and Mesospheric Sounder (SAMS) instrument flown on Nimbus-7 is the fourth in a series of multi-channel infrared radiometers designed to measure emission from the upper atmosphere, for which conventional spectral filtering techniques do not give adequate performance. The technique used in these radiometers is known as gas correlation spectroscopy and is based on the use of gas cells to select emission from chosen spectral lines or from particular parts of spectral lines.

In the Selective Chopper Radiometer (SCR) on Nimbus-4, a beam-chopping technique was employed to switch the scene (at 10 Hz) between the atmosphere and space view in a differential manner between two gas cells containing different amounts of CO<sub>2</sub>. The "difference" signal was then detected by a thermistor bolometer. The chief limitation in performance of the Nimbus-4 SCR was the difficulty in maintaining a balance condition due to stray thermal emission from within the instrument.

Nimbus-5 employed a variation of this technique, in which cells containing different amounts of CO<sub>2</sub> were switched in sequence (one each second) into the optical path to a pyroelectric detector. The difference signals were then extracted on the ground. Performance in this case was limited by gas leakage with time, by uncertainties in the effects of degradation, and by contamination of the cell windows (which give spurious difference signals).

The Pressure Modulator Radiometer (PMR) on Nimbus-6 overcame earlier difficulties by employing a single gas cell and no moving parts in front of the detector. The gas (CO<sub>2</sub>) amount in the cell is modulated at approximately 35 Hz by an oscillating piston, and the oscillatory component of the signal arriving at the detector is related directly to the radiance of the scene, but only at the frequencies corresponding to the variation in absorption of the spectral lines of the gas in the modulator cell.

The SAMS instrument views the limb of the atmosphere, rather than employing vertical sounding as in the earlier radiometers. The technique has been extended to gases other than CO<sub>2</sub>.

### **1.3 Objectives**

The SAMS is a 12-channel infrared radiometer observing thermal emission and solar resonance fluorescence from the atmospheric limb. Global measurements are made of radiation from the molecular species listed in Table 1-1. These measurements, when interpreted together with results from the LIMS and SBUV/TOMS instruments, provide extensive data for chemical and dynamical models of the stratosphere and mesosphere.

Specific objectives of the SAMS experiment were to derive the following quantities:

- a) Temperature for altitudes from 15 km to 80 km, from emission in the 15  $\mu$ m CO<sub>2</sub> band.
- b) Vibrational temperature of CO<sub>2</sub> bands where they depart from local thermodynamic equilibrium (LTE), between 50 km and 140 km.
- c) Distributions of CO, NO, CH4, N2O, and H2O from 15 km to 60 km.
- d) Distributions of CO<sub>2</sub> (4.3  $\mu$ m) from 100 km to 140 km and H<sub>2</sub>O from 60 km to 100 km, to study dissociation in the lower thermosphere.

The study of planetary waves that used data from the Nimbus-5 SCR and Nimbus-6 PMR can be continued and extended by making use of these new data. In addition, the measurements allow calculation of the transfer of momentum, energy, and trace gases by mean motions and eddies.

The SAMS instrument is designed to exploit the selectivity, energy grasp, and tuning capability of the pressure modulation technique proved earlier for CO<sub>2</sub> emission measurements in the Nimbus-6 PMR (Curtis *et al.*,1974). The main technical innovations in SAMS are as follows:

- a) The extension of the pressure modulation technique to other gases.
- b) The simultaneous use of conventional chopping and pressure modulation to
  - 1) Extend the range of heights that can be sounded
  - 2) Determine the pressure at the viewing level
  - 3) Enable some interfering radiance signals to be eliminated
  - 4) Provide additional calibration information and confidence checks
- c) The use of a programmable scan system with two independent axes to optimize usage of the observing time and to accommodate uncertainties in spacecraft attitude.

### **1.4 Data Products**

Retrieved temperature and constituent information from the SAMS instrument is stored on two data products, GRID-T and ZMT-G tapes, respectively, which are archived at the National Space Science Data Center (NSSDC). The temperature data on the GRID-T tapes cover the period December 24, 1978 to June 9, 1983. The ZMT-G tape consists of methane and nitrous oxide mixing ratios and covers the period January 1, 1979 to December 30, 1981.

The tapes are 9-track, 6250 bpi. Data records are of variable length. The data formats of the SAMS tapes are described in detail in Section 6.

#### Table 1-1

#### Molecular Species and Spectral Bands

<u>Constituent</u>	Spectral Band
Carbon dioxide	4.3 μm and 15 μm
Water vapor	2.7 $\mu m$ , and 25 $\mu m$ to 100 $\mu m$
Carbon monoxide	4.7 μm
Nitrous oxide	7.7 μm
Methane	7.7 μm
Nitric oxide	5.3 μm

### **1.5 Principal Investigators**

The principal investigator for the Nimbus-7 SAMS experiment at the University of Oxford is F. W. Taylor. Clive D. Rodgers is the NET chairman. Users should contact these individuals if questions arise concerning SAMS data or scientific algorithms: Their address is

> Department of Atmospheric Physics Clarendon Laboratory University of Oxford Oxford OX1 3PU Great Britain

Telephone: Oxford (0865) 272905 Telex 83650 ATMOXF G

The NASA contact regarding SAMS is

Nimbus-7 Manager NASA/GSFC, Code 636 Greenbelt, MD 20771

Telephone: (301) 286-9846

### **SECTION 2**

### **INSTRUMENT DESCRIPTION**

A complete discussion of the instrument, its principles of operation, calibration, and measurement procedures is given by Drummond *et al.* (1980). For convenience, this article is reproduced in its entirety from the *Philosophical Transactions of the Royal Society of London* as Appendix C of this report. Sections 2-5 (pp. 220-237) of that paper present relevant information about the instrument's behavior. Additional material on the design and behavior of the SAMS is contained in a paper by Wale and Peskett (1984), which is reproduced from the Journal of Geophysical Research and presented here as Appendix G. The articles are reproduced here with the permission of the authors and the journals in which they were originally published.

A somewhat less detailed description of the SAMS instrument is found in Section 6 of the Nimbus-7 User's Guide (Drummond *et al.*, 1978).

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#### **SECTION 3**

### DATA PROCESSING

Continuous observations from all eight experiments were recorded aboard the spacecraft. In normal operation, one tape recorder was recording while another was playing back data to a receiving station, resulting in no loss of data. Initial processing of all data took place at Goddard Space Flight Center (GSFC), after which SAMS data were transmitted by data lines to Oxford for all further processing.

Approximately 13 orbits per day requiring about 2 hours of transmission time were received at Oxford for further processing. The data were received by a PDP 11/70 computer, where they were simultaneously written directly to 9-track, 1600 bpi phase-encoded magnetic tape in raw format and processed to a "quick-look" plotting stage. The processing was performed by a suite of programs that ran under a controlling program requiring operator intervention only when problems occurred. The data were quality checked and calibrated. A "quick-look" estimate of spacecraft attitude was calculated using data from two channels of the SAMS instrument. Various data products were plotted for each orbit. These plots were regularly checked to determine instrument health. When all data had been received, the software produced data products required on a daily basis rather than by orbit, as well as additional plots, e.g., retrieved temperature and listings of housekeeping functions. The temperature plots enabled a watch to be kept for sudden warmings and the housekeeping information assisted in defining operating modes and data quality of the instrument.

A stage of final processing took place when all the raw data had been sorted into time order. The same suite of programs was used but a better calibration and other data, e.g., ILT, were included.

Figure 3-1 provides a schematic overview of the data processing.

### 3.1 Calibration

The calibration of SAMS can be divided into two broad areas:

- a) Conversion of signal channel telemetry counts into a radiance, expressed as a fraction of the signal obtained from a black body at a specified temperature, with sufficient additional information to define the spectral properties of the instrument and its field of view.
- b) Calculation of the spectroscopic properties of each channel and the physical properties of the gases being investigated in the atmosphere, in order to retrieve useful quantities (temperature, composition) from the calibrated radiances.

This section deals with the first area only; the second area is discussed in Section 3.2. The Nimbus-7 User's Guide contains an outline of the various preflight calibration measurements made, including the following:

- a) Measurement of the response of the instrument to black bodies at temperatures between 77 K and 320 K. The temperature variations experienced in orbit were simulated during this test, so that temperature coefficients could be determined dynamically and likely operating conditions estimated. The basic measurement verifies the linearity of the instrument response and that the spectral properties of each channel are known.
- b) Measurements of the field of view of each channel, and calibration of the scan mirror telemetry system.

The response of each channel is linear with respect to the incoming radiance, and may, therefore, be characterized by a slope (the instrument gain) and an offset, corresponding to the signals observed when observing space (zero radiance input).

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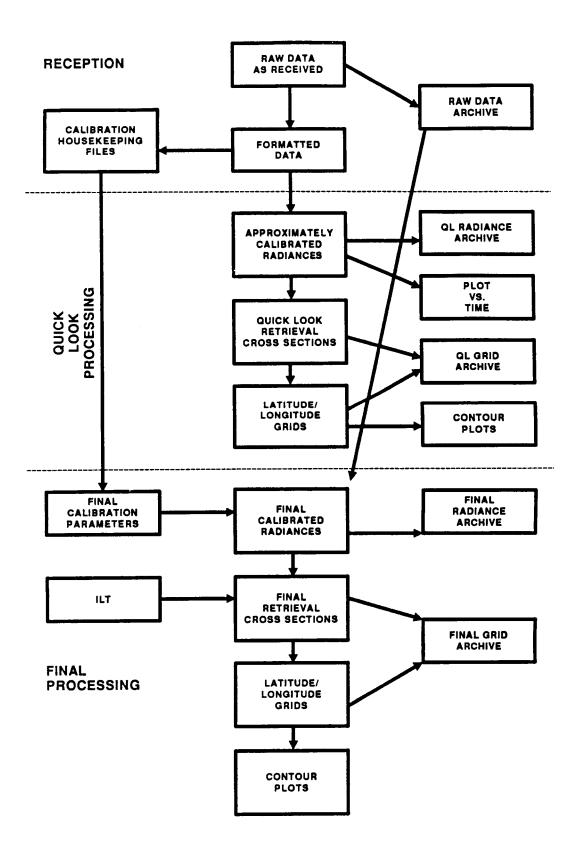


Figure 3-1. Overall view of data processing.

The following are some of the most important factors that affect the gain of a *narrow band* pressure-modulated channel (PMC):

a) <u>Cell pressure</u>

The equivalent width of the pressure modulation depends on the pressure in the cell, as a result of changes in the absorber amount and line shape. An additional influence is the variation of the frequency of the modulation with the mean cell pressure, as the response of the detector and signal channel is not uniform with frequency. There is also a phase difference between the cycling of the cell pressure and the piston position, which is used to derive the phase sensitive detector (PSD) reference wave forms. This results from the flow of gas to and from the PMC head and is also pressure dependent; it may be significant in some cases at low cell pressures.

#### b) <u>Cell temperature</u>

The molecular sieve gas containment system ensures that the mean pressure of gas in the PMC is held sensibly constant as the temperature T varies. The absorber amount, therefore, changes as 1/T, with a corresponding change in the equivalent width. This variation is quite distinct from the very small changes that result from variation of the line width and strength with temperature.

#### c) <u>Detector temperature</u>

The responsivity of all of the detectors varies with temperature; in the case of the PbS unit employed in the B1 channel, the phase shift is also temperature dependent.

d) <u>PMC amplitude</u>

There is a small variation in the amplitude of the oscillation of the pressure modulator with temperature; this is significant not only because of its direct effect on the gain, but also because it causes a small shift in resonant frequency.

For the *wideband channels*, the situation is very much simpler; the gain depends principally on the detector temperature and the chopper amplitude.

The offset is the sum of a large number of components, the principal contributions being from

- a) The emission from the modulator or chopper against which the incoming radiation is being chopped.
- b) Temperature cycling within the pressure modulator cell, which results in an additional emission signal, with amplitude and phase that are strong functions of cell pressure but not significantly dependent on temperature.
- c) The emission from optical components in front of the modulator or chopper, including the system aperture stop (in the case of the PMR channels, the mask of the ellipsoid mirror M3) and the front optics components.
- d) Stray radiation detected as a result of aberrations and imperfections in the optical system.

e) Synchronous pick-up. Pyroelectric detectors such as the triglycine sulfate (TGS) detectors employed in SAMS also display a piezoelectric effect and show a significant microphonic response. The pressure modulators and the chopper cause a small level of vibration in the instrument structure. This vibration is detected directly through the detector piezoelectric response. The level of vibration is, unfortunately, not very stable, and varies with operating modes; there is also variation around the orbit, which is of a constant form over a period of a day or so. Such a response is especially noticeable in the detector fitted to the C2/3 channel. This particular detector has been noted in other ways to exhibit an unusually high level of microphonic response. This problem is treated by fitting the residual space view variation to terms periodic in the orbital position, after normal calibration.

Most of the above factors depend on the instrument gain and are, therefore, subject to all the same variations with temperature, etc.

Because of the very large number of factors that contribute to the gain and offset of each channel, overall calibration of the instrument must be performed frequently during flight. The zero radiance signal is checked by moving the scan mirror so that the tangent height of the path viewed by the lowest of the three fields of view is at least 150 km (space view); a second calibration point is found by inserting a room temperature black body into the optical path at the focus of the paraboloid, M2. Since SAMS measures signals that are very close to zero, the zero offset of the system must be known extremely precisely. Every scan program, therefore, includes frequent space view calibration periods. One of the commonly used programs has two samples of space view data every 32 seconds. Black body calibrations are required less frequently.

It is not possible to find the magnitude of all the relevant factors by calibration in flight, as many of them show a high correlation, and the variations in operating conditions experienced in a typical period of operation in a given mode are not sufficiently wide. Taken in isolation, the calibration sequences do not have good enough signal to noise ratios to permit their direct use. Many of the factors (e.g., the effect of temperature on the PMR equivalent width, the emission from the cell or chopper and the emission from optical components) can be calculated explicitly; the remaining coefficients may then be found by regression analysis with respect to the temperatures, etc. involved. Normally, SAMS is kept in a given operating mode for one 24-hour period at a time, the first orbit of the day being used to allow the cell pressures to stabilize after the mode change. During this period, continuous black body and space view information is needed.

Calibration of the instrument is performed on a day-by-day basis, and the effects of the variations (e.g., cell pressure) are treated as perturbations from a nominal value. The cell pressures show long-term trends, however, and for this reason the calibration coefficients cannot be fixed in a single test, for instance as part of the preflight calibration tests. Various special tests have been run in orbit to provide additional calibration information, including one 3-day period of continuous calibration sequences. During this test, auxiliary heaters were employed to provide a wider variation in temperature than that found in normal operation. Tests have also been run to investigate the signal levels observed at all possible positions of the scan mirror. These tests have shown that the variation of the stray signal (including microphonic terms) with mirror position is negligible.

### 3.2 Retrieval of Temperature and Constituents

The retrieval problem has two aspects, temperature and composition. The temperature profile is determined from the two 15  $\mu$ m CO<sub>2</sub> pressure-modulated (PM) channels and the two corresponding wideband (WB) channels in the A and C fields of view, which are separated by 0.56 degrees (about 35 km). The composition retrieval is determined from radiances in the other channels, but requires the temperature profile as input.

As with remote sounding in general, the problem is underdetermined and cannot be solved explicitly without recourse to *a priori* information. The temperature retrieval problem has some complications that make it less nearly linear than its nadir sounding counterpart. The most serious complication is that the attitude of the spacecraft is not adequately measured or controlled, so that the tangent height of the line of sight is not known accurately enough to compute the corresponding transmittance function. However, there is enough information in the radiances to determine both temperature profile and tangent height. The composition retrieval is nonlinear because the unknown enters the equation of transfer as an argument of the transmittance function.

Our approach to both retrieval problems is to linearize the direct equation about an *a priori* profile, and to use a statistically optimum estimator (Rodgers, 1976) to find the most likely solution. Iteration is avoided by ensuring

that the *a priori* profile is sufficiently close to the solution that the linearization is adequate. In the case of temperatures this is achieved by using a sequential maximum-likelihood estimator, in which the retrieval corresponding to time t and its error covariance matrix are used to construct the *a priori* profile and its error covariance at time  $t + \Delta t$ , where  $\Delta t$  is the time interval between measurements. For composition retrieval, it has proved necessary to average the data, for example zonally, in order to improve the signal-to-noise ratio and to reduce computing time. In this case, the previous day's retrieval for the same averaging box is used to construct the *a priori* profile. Full details of the methods adopted have been described by Rodgers *et al.* (1984), reprinted as Appendix D of this user's guide.

### 3.3 SAMS Temperature Gridding

The sequential estimator used to retrieve temperature and attitude obtains estimated retrieval coefficients and the corresponding covariance matrix every 2 seconds around the orbit (about every 14 km). These values are ultimately interpolated onto a regular latitude-longitude grid at latitudes 50° S, 47.5° S, 45° S, ....., 67.5° N and longitude 0°, 10°, 20°, 30°, ....., 350° E. As a first step, the sequential estimator selects the 2-second retrieval nearest to each 2.5° latitude line (but no farther than 1.25° away), and arranges for a combined forward + backward estimate of the temperature retrieval and error information to be written to an output file.

These results are combined into two "orbit" grids per day, one for the northbound (ascending) portions of each orbit (all between about 1400 and 1800 hours local time) and one for the southbound (descending) portions (1800 to 2200 hours local time). Each grid has 48 rows (latitudes) and 14 columns (orbits). Some grid points will be empty because of missing orbits. For example, the first orbit each day is generally spent calibrating while the PMC pressure settles down, so no retrievals are obtained; other orbits or parts of orbits were never received. Note that the data here are divided into "orbits" starting and ending when the satellite is at 80° N or S; they bear no relation to an 'orbit' as received from NASA, which is normally one data dump of the satellite tape recorder. Note also that the orbit number has only approximate meaning, because those given by NASA can be the orbit on which the recorder was replayed, which may be several orbits after the observations were made.

Given a northbound and southbound grid for each day, gaps of up to two points along an orbit are filled by linear interpolation between values at the north and south ends of the gap for each grid. The values are then interpolated in longitude using linear interpolation, to obtain still separate northbound and southbound latitude/longitude grids. Where there is a missing orbit, no attempt is made to interpolate into the gap, except that the last good observation on either side of the gap is used for points up to a quarter of the orbital separation into the gap on that side. This procedure is carried out for the following quantities:

- a) Eigenfunction coefficients of Planck deviation from the climatological first guess.
- b) Standard deviations of (a).
- c) Temperature at 10 special pressure levels.
- d) Standard deviations of (c).

The off-diagonal elements of the eigenfunction covariance matrix are not gridded because there are so many of them; however, the standard deviations of temperature at the selected levels are obtained in the retrieval program using the full covariance matrix.

The northbound and southbound grids are now combined by taking both values where possible to form a combined daily average grid. In the case of temperatures and eigenfunction coefficients, the values are weighted inversely by the reciprocal of their variances, while standard deviations are combined by taking the square root of the sum of reciprocals of the variances. If data are available from only one of the grids (northbound or southbound, not both) the retrieved values from that grid are used unchanged in the final mean analysis.

The combined mean analyses are supplied on the data tape in a format where each field is given as a separate tape block, with eigenfunction coefficient or temperature fields followed by the corresponding error grid. In addition, the grids of eigenfunction coefficients are used to obtain temperature grids with a vertical resolution of 0.2 pressure scale heights (62 levels in all). These temperatures are stored on the tape in profiles together with the first guess climatological profile used and the zonal mean for each latitude.

#### **SECTION 4**

### VALIDATION

Two papers concerning the accuracy of the retrieved SAMS temperatures and the methane and nitrous oxide distributions are reprinted from the *Journal of Geophysical Research* (Barnett and Corney, 1984; Jones and Pyle, 1984) as Appendixes E and F of this User's Guide. The articles are included with the permission of the authors and the Journal. Independent comparisons of several stratospheric data sources, including SAMS and LIMS, may be found in the Middle Atmosphere Program [MAP] Handbook, (Rodgers, 1984*a*; Grose and Rodgers, 1986), published by SCOSTEP, and available from the SCOSTEP secretariat, University of Illinois, 1406 W. Green Street, Urbana, IL 61801.

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#### **SECTION 5**

### **MISSION HISTORY AND CHANNEL STATUS**

Table 5-1 lists the significant dates in the history of SAMS. The first 2 months of instrument operation were an experimental period, during which many changes were made to the scan pattern and sieve setting programs. Tables 5-2 and 5-3 provide full details of the settings used, along with comments on some problems.

The following points concerning the instrument have significant consequences for the quality of the data.

- a) SAMS operated on a 3-day-on, 1-day-off duty cycle; atmospheric phenomena with a 4-day period will not be well represented in the data.
- b) The first orbit of any day is almost always a calibration orbit.
- c) The passive cooler did not reach the design temperature; as a result, the shortwave channels (NO, CO, and CO<sub>2</sub> 4.3 microns) all had a poorer signal-to-noise ratio than expected.
- d) A chemical problem with the NO molecular sieve rendered the channel unusable. The NO modulator failed by driving at too high an amplitude on day 91 of 1980.
- e) PMC frequencies depend on cell pressure, and hence on sieve setting. For some combinations, there are "beats" between different modulators at similar frequencies. These occasions are noted in Table 5-2. One significant consequence of the possible beating is that the upper CO<sub>2</sub> 15 micron channel (A1) is sometimes in a low pressure setting to avoid beats with the CO channel. There is a systematic difference between temperatures retrieved with A1 sieve 3 and those retrieved with other sieve settings, particularly in the mid and upper mesosphere. We suggest that temperatures from days with A1 settings other than sieve 3 be treated with caution.
- f) A problem with the housekeeping analog/digital (A/D) converter caused some instrument temperature measurements to be made with low accuracy. This problem was largely circumvented in the processing.
- g) Changes were made in the water PMC drive circuits after the final pressure calibration runs were made. Consequently, the water cell pressure measurements are uncertain.
- h) The instrument failed partially during 1982 when the scan mechanism stuck. After a period of recovery, final failure occurred in June 1983.
- i) The eruption of El Chichon in April 1982 injected a considerable amount of aerosol into the lower stratosphere. This event is reflected in significant changes to the temperature retrievals, but it is uncertain whether these changes are real or are the result of unmodeled emission from the aerosol affecting the retrieval process. Before September 1982, the effects were confined to the region equatorwards of about 30° N and below about 30 km but, after this date, the dust cloud quickly moved to higher latitudes but with lower concentration. We have not carried out nitrous oxide and methane retrievals for this period because of these uncertainties.

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#### Table 5-1

#### Significant Dates in the History of SAMS

Day	Year	Event
297	1978	Launch.
299	1978	SAMS turned on.
317	1978	Cooler door opened.
358	1978	First usable data.
48	1979	Cooler decontamination heaters on.
91	1980	NO modulator failed.
100	1982	El Chichon effects first seen in data.
132	1982	Cooler decontamination heaters on.
277	1982	Scan problem - first appearance.
301	1982	Cooler decontamination heaters on.
44	1983	Scan problem - apparent recovery.
164	1983	Scan problem - final failure.

Following is the status of the channels:

a) A1 and C1; CO<sub>2</sub> 15 micron temperature sounding channels.

Temperature retrievals have been carried out for the whole data set, except for a few occasions when the data quality was poor, or when the instrument was not scanning due to scan problem or azimuth scan tests.

b) A2; 4.3 micron CO<sub>2</sub> channel. Non LTE studies.

This channel suffered from the poor signal-to-noise (S/N) problem of all the shortwave channels.

c) A3; carbon monoxide.

The SAMS measured CO during the daytime using a pressure modulator to measure resonant fluorescent scattering of sunlight near 4.7 microns. The noise level was high, and it was necessary to average over time periods of 6 months and  $35-50^{\circ}$  wide latitude bands. Consequently, the total number of profiles obtained was small, and they are all given in Figure 5-1. It should be noted that,

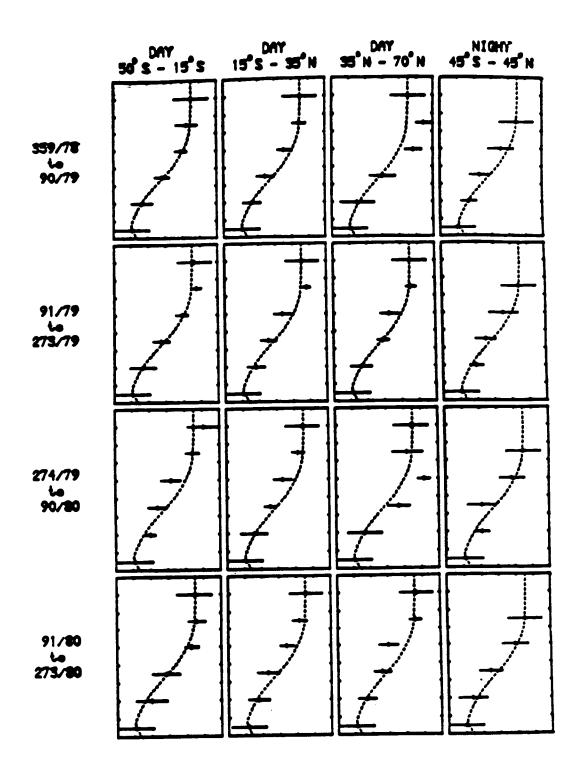


Figure 5-1. Comparison of the retrieved carbon monoxide profiles. Dotted line is the *a priori* profile (identical in all cases). Vertical scale:  $-\ln$  (pressure/atmospheres) = 2 to 20. Horizontal scales:  $\log_{10}$  (volume mixing ratio) = -9 to -3.

because of the geometry of the orbit, the density of observations is greatest at 50° S and 70° N. Consequently, averages for bands that include these latitudes are biased toward them.

The precision (systematic error) of the retrieval is approximately  $\pm 25$  percent, but the accuracy (random error) is worse, about -57 to +130 percent. The profile is valid over the range 6 to 14 pressure scale heights (about 45 to 100 km).

Although the random errors are high, marked variations can be seen. The most significant feature is that mixing ratios are very high in the mesosphere during the northern hemisphere winter. This effect is clearly seen in the retrievals for the 1978-79 and 1979-80 winters. There is a difference of well over a factor of 10 between summer and winter mixing ratios at some levels. This difference is far too large to be attributed to errors in the measurements or the retrieval process. There is some evidence for a similar effect in the Southern Hemisphere, at about 85 km (12 pressure scale heights), but the most southerly latitude zone extends only to 50° S. The variations between the remaining profiles are generally comparable with the level of errors on the retrievals. The central latitude band shows a profile that exhibits little variation between the data periods, and the mixing ratios are comparable with summertime values in the northern and southern hemisphere zones.

d) A4: nitric oxide.

The nitric oxide in the PMC unexpectedly suffered from a catalytic disproportion reaction on the molecular sieve material, so that the PMC actually contained a mixture of nitrous oxide and a small amount of nitric oxide. Consequently, the data from this channel are of little value. The channel failed mechanically after 520 days in orbit.

e) B1: water vapor 2.7 micron resonance fluorescence.

This channel should provide water concentration in the 65-90 km region, but suffers from poor S/N, as do the other shortwave channels. So far only a limited number of retrievals have been performed for monthly zonal means, in the latitude range of  $45^{\circ}$  S to  $45^{\circ}$  N for the months of January, April, July, and October 1979. Processing is in hand to produce similar products for the whole of 1979 and for smaller latitude boxes.

f) B2: water vapor rotation band.

This is the primary water vapor channel, which should produce both zonal means and global distributions of water from 15 to 60 km. However, our retrieved profiles to date have failed to satisfy our validation criteria. The problem appears to be associated with uncertainties in the spectroscopy of water vapor in pressure modulators, and the symptom is that the retrieved profiles show an unreasonably large amount of water at high levels. Laboratory studies are currently being undertaken in an attempt to solve this problem.

Preliminary processing of zonal mean fields for 1979 and 1980 has been completed; similar processing of the remainder of the data and the production of global fields is proceeding. These data can be used only to illustrate the qualitative behavior of the water distribution, however, and will not be archived with NSSDC.

g) C2 and C3: nitrous oxide and methane.

Nitrous oxide and methane have been retrieved from the SAMS radiances for the period January 1, 1979 to December 31, 1981. The data are too sparse during the experimental period before January 1979 to be useful, and 1982 contains both the El Chichon eruption and the scan failure. Some parts of the 1982 data will be retrieved and placed in the archive. Note that because C2 and C3 both use the same detector, only one of the two channels can be used at any one time.

For channels A2/3/4 and C2/3, the wideband measurements are of little use since they were the result of emission from several gases. They were not used in the retrievals. They were, however, affected by the El Chichon dust, and they could be used to study the wavelength dependence of aerosol absorption.

The pattern of use for the instrument is given in Table 5-2. The 'SCAN PROG' column is the internal reference number for the scan program being employed on that day. It is included so that different days when the same program was being used can be identified. Detailed scan patterns can be supplied by the principal investigators at Oxford (Section 1.5) for individual cases on request. The nominal scan range is given in km for the 'D' field of view. Relative to this direction, the A field of view (FOV) is 0.32° above, B is 0.16° above, and C is 0.24° below (approximately 18, 9, and -7 km respectively). The nominal scan range is only an approximate guide to the actual scan range obtained in the atmosphere because of the attitude variations of the spacecraft relative to the horizon. The "sieve setting" number indicates the choice of PMC cell pressure (see Table 5-3), and the "E" indicates which modulator is energized, in those cases where there is more than one modulator associated with a detector.

Table 5-2 Mode Table 1: Launch - day 239, 1982

		664 <b>8</b>		) SIEVE	SETTING		
DATE Year day	ORBITS (APPROX)	PROG	(APPROX)	A1 A2 A3	BI CI C	2 C3	
			<b></b>		82		
1978 299	32:35	1.1		JE OE O	32 32 0		WIDE SCAN - LIND SEARCH
1978 299	35:37	2.2	15:90	IE O O			
1978 300	37:52	2.2	15:90		3E 3E 1		
1978 301	52:58	2.7			3E 3E 0		NO SCAN
1978 301	58:64	2.3	30:105		2E 3E 1		BEAT B12/CH MEAT C1/C3
1978 302	64:78	2.8	10:85		32 32 1		BEAT C1/C3 BEAT C1/C3
1978 303 1978 303	78:89 89:92	2.8 2.10	10:85 10:85	3200	312 312 1 312 312 1		BERT CIVES
1978 304	92:106	2.10	10:85	-	32 32 1		
1978 305	106:119	2.12	30:105		IE SE O		
1978 305	119:120	2.17	501105		3E 3E 1		NO SCAN BRODY/SPACE CAL
1978 306	120:130	2.11	29:95		2E 3E 1		BEAT B12/CR
1978 306	130:134	2.9	10:85	372 0 0	3E 3E 1	2 1	
1978 307	134:148	2.13	30:105	3E 0 0	1E 3E 1	12	
1978 308	148:163	2.11	20:95		28 38 1		BFAT B12/CH
1978 309	163:175	2.16	25:100		3E 3E 1		
1978 310	175:189	2.16	25:100		3E 3E 1		
1978 311	189:203	2.16	25:100		3E 3E I	12	
1978 312	203:210	2.16	25:100	3E 0 0	38 38 1	12	
1978 312	210:213	5.24	10.85		32 32 1		MOON SCAN ATTEMPT (PCL PROG ERROR)
1978 312 1978 313	214:216 217:231	2.9 3.18	10:85 15:35		312 312 11 212 312 1		BEAT BIZ/CH
1978 314	231:245	3.19	25:45		12 32 1		BERT BIL/ON
1978 315	245:258	3.21	-5:30	320 0			
1978 316	258:261	3.21	-5:30		3E 3E 1		
1978 316	261:272	4.20		1E 0 0			NO SCAN
1978 317	272:282	4.22			OE SE O		NO SCAN
1978 317	282:286	4.23		12 0 0	18 38 1	IE	NO SCAN COOLER DOOR OPEN 282 ETC 17002
1978 318	286:300	4.23		1E 0 0	1E 3E 1	12	NO SCAN
1978 319	300:315	3.25	5:25	1E OE 1	3e 3e 1	<b>E</b> 1	
1978 320	315:325	3.28	5:40		3E 3E 1	12	
1978 321	337:341	3.28	5:40		3E 3E 1	12	BALLOON FLIGHT
1978 322	341:355	8.29	10:45		3E 3E 1	12	
1978 323	355:368	6.30		IE OE I		IE	NO SCAN
1978 324	368:383	6.32		0211			NO SCAN
1978 326	397:411	6.31	20.48		05, 37, 1		NO SCAN
1978 327 1978 328	411:424	8.33 8.33	30:65		18 38 1		
1978 330	424:438 452:465	8.34	30:65 -5:30		: 18 38 11 : 38 38 1		DATUM MAY BE TOO LOW FOR GOOD SPACE VIEW
1978 331	465:480	8.34	-5:30		38 38 1		DATUM MAY BE TOO LOW FOR GOOD SPACE VIEW
1978 332	480:494	8.35	-5:30	3E 1 1			DATUM MAY BE TOO LOW FOR GOOD SPACE VIEW
1978 334	508:521	11.36	-5:50	3E 1 0E			DRIVE WE DE TOU EUW FOR OUDD STREE VIEW
1978 335	521:535	11.36	-5:50		3E 3E 1		
1978 336	535:549	11.37	-5:50	32 1 02			
1978 338		11.37		3E I OE			
1978 339	578 <b>: 585</b>	5.38			3E 3E 1		BAD SCAN PROGRAM
1978 339	586:590	11.39	-5:50	3E 1E 1	3E 3E 1	IE	BEAT A12/A3
1978 340	590:604	11.39	-5:50		3E 3E 1		BEAT A12/A3
1978 342	618:627	12.42		3E 1 0E			NO SCAN
1978 342	628:633	5.38			3E 3E 1		BAD DATA
1978 343	634:645	11.40	-5:50		3E 3E 1		BEAT A12/A3
1978 344	645:660	11.40	-5:50	BELEI			BEAT A12/A3
1978 346	673:688	8.41	25:60	35 1 05			BAD SCAN DROOMAN
1978 347 1978 347	688:690 696:702	0.0 3.28	5:40	JE I E I	2E 3E 1		BAD SCAN PROGRAM
1978 347	702:715	3.28	25:60	321 02			BFAT B12/CH
17/0 340	/04:/13	0.41	23:00		46 36 1	15	

DA 779			ON) SIEVE SETTINGS	
DATE Year day	ORBITS	PROG (APPROD	() AI A2 A3 BI CI CZ C3	
TEAK ONL	(APPROX)	PRUG (APPRUG	B2	
1978 350	738:743	15.56 -20:35	312 I IE 312 313 112 1	
1978 351	743:756	15.44 0:55	3E 1 OE 3E 3E 1E 1	
1978 352	757:770	8.46 25:60	3E 1 OE 2E 3E 1E 1	POSS BEAT B12/CR
1978 354	784:798	108.50 35:70	1E 1 1 1E 3E OE 1	
1978 355	798:812	108.158 10:45	1g 1 0 1g 2g 1 0g	
1978 356	812:825	108.147 20:55	32 OE 0 22 32 1 12	POSS BEAT B12/CR
1978 358	839:853	108.148 20:55	312 I O 212 312 112 1	POSS BEAT B12/CH
1978 359	854:867 867:881	115.161 -5:50 115.162 -25:30	2E 1E 0 3E 3E 1 1E	
1978 360 1978 362	895:908	114.153 50:115		
1978 363	906:922	114.159 50:115		
1978 364	922:936		011 1 01 31 01 1	
1979 1	950:964	115.156 -25:30	3E 1 OE 3E 3E 1E 1	
1979 2	964:977	115.144 -5:50	3E 1E 0 3E 3E 1E 1	POSS BEAT A12/A3
1979 3	97 <b>7:982</b>	115.144 -5:50	3E 1E 0 3E 3E 1E 1	POSS BEAT A12/A3
1979 3¥	982:991	5.43	3E 0 0 3E 3E 1 1E	MOON SCAN
1979 5	1005:1019		1E 1 0 1E 3E 0E 1	
1979 6	1019:1037		IE I O IE 2E I OE	
1979 7 1979 7	1033:1037	108.158 35:70 5.43	1E 1 0 1E 2E 1 0E 3E 0 0 3E 3E 1 1E	NOON SCAN
1979 9	1057:1047		BEI O BEBEI IE	Moon Scan Poss Beat B12/Ch
1979 10	1075:1088		2E 1E 0 3E 3E 1 1E	FUSS BEAT STEFCI
1979 11		115.162 -25:30	2E 1E 0 3E 3E 1 1E	
1979 13		114.153 50:115		
1979 14		114.159 50:115		
1979 15	1143:1157	114.155 50:115	OFE 1 0 OFE 372 OFE 1	
1979 17		115.156 -25:30	3E 1 OE 3E 3E 1E 1	
1979 18		115.144 -5:50	3E 1E 0 3E 3E 1E 1	POSS BEAT A12/A3
1979 19		115.144 -5:50	3E 1E 0 3E 3E 1E 1	POSS BEAT A12/A3
1979 21	1226:1240		1E 1 0 1E 3E 0E 1	
1979 22 1979 23		115.145 -5:50 108.147 20:55	32 1 02 32 32 1 12 32 02 0 22 32 1 12	POSS BEAT B12/CH
	1263:1268		3E IE 0 3E 3E IE 1	POSS BEAT A12/A3
		108.147 20:55	3E OE O 2E 3E 1 1E	STRAT WARMING POSS BEAT B12/CR
-	1282:1297			STRAT WARMING POSS BEAT A12/A3
1979 26		115.144 -5:50	3E IE 0 3E 3E 1E 1	STRAT WARMING POSS BEAT A12/A3
1979 27	1310:1323	115.246 10:65	3E 1 OE 2E 3E 1E 1	STRAT WARMING POSS BEAT A12/A3
		115.250 25:80	IE I O IE JE OE I	STRAT WARMING POSS BEAT B12/CH
		115.145 -5:50	3E 1 OE 3E 3E 1 1E	STRAT WARMING
		115.247 10:65	3E OE O 2E 3E 1 1E	STRAT WARHING POSS BEAT B12/CH
	1365:1379	115.251 25:80 115.261 -5:50	1E 1 0 1E 3E 1 0E 2E 1E 0 3E 3E 1E 1	STRAT WARMING POSS BEAT CI/C3 STRAT WARMING
	1393:1406		3E 1 OE 2E 3E 1E 1	STRAT WARNING POSS BEAT B12/CH
1979 34		115.145 -5:50	3E 1 OE 3E 3E 1 1E	Sticki waking 1000 april 512/ ch
			3E 1 OE 3E 3E 1 1E	
		115.145 -5:50	3E 1 OE 3E 3E 1 1E	
1979 37	1448:1461	115.145 -5:50	3E 1 OE 3E 3E 1 1E	
	1462:1475	115.145 -5:50	3E 1 OE 3E 3E 1 1E	
		115.145 -5:50	3E 1 OE 3E 3E 1 1E	
	1504:1516	115.145 -5:50	32 1 02 32 32 1 12	
		115.145 -5:50	3E 1 OE 3E 3E 1 1E	
		115.145 -5:50	3E 1 OE 3E 3E 1 1E	
		114.253 40:105 114.254 40:105		POSS BEAT CI/C3
		116.262 -15:50	1E OFE O OFE 3E 1 OFE 2E 1E O 3E 3E 1 1E	POSS BEAT CI/C3
		116.262 -15:50	2E IE O 3E 3E 1 IE	NO SCANNING

DATE	ORBITS		I) SIEVE SETTINGS	
YEAR DAY		PROG (APPROX)	A1 A2 A3 B1 C1 C2 C3	
			82	
		_		
1979 49		116.262 -15:50	2E 1E 0 3E 3E 1 1E 3E 1 0E 2E 3E 1E 1	"NO" ROCKET COVERAGE
1979 50		115.246 10:65	3E 1 OE 2E 3E 1E 1	"NO" ROCKET COVERAGE
1979 51 1979 53		115.246 10:65 118.348 -10:90	372 1 0 272 372 112 1	POSS BEAT B12/CR
1979 54		118.345 -25:75		
1979 55		118.362 -45:55		``
1979 56		118.346 -10:90		POSS BEAT B12/CH
1979 57		118.346 -10:90		POSS BEAT BI2/CR
1979 58		118.347 -10:90		NO SCANNING, POSS BEAT B12/CH
1979 58 1979 59		118.347 -10:90 119.355 40:125	312 OFE 0 212 312 1 112 OFE 1 0 OFE 312 OFE 1	POSS BEAT B12/CH FREQ SPACE CAL
1979 61		118.361 -25:75		FREQ SPACE CAL
1979 62		118.361 -25:75	28 1E 0 3E 3E 1 1E	
	1802:1807		38 0 0 3E 3E 1 1E	MOON SCAN
		118.346 -10:90	31E 1 OE 21E 31E 1	POSS BEAT B12/CH
•	1835:1848			
	1648:1855			LOOM COALE
1979 66 1979 67	1855:1863	5.143 118.347 -10:90	3E 0 0 3E 3E 1 1E 3E 0E 0 2E 3E 1 1E	HOON SCAN POSS BEAT B12/CH
1979 69		118.348 -10:90		POSS BEAT B12/CR
1979 70	1904:1918	118.345 -25:75	3E 1 OE 3E 3E 1E 1	
		118.362 -45:55		
1979 73		119.453 30:115		PREQ SPACE CAL
1979 73 1979 74	1958:1959		1E 0 0 3E 3E 1 0E 1E 0E 0 0E 3E 1 0E	SCAN CAL CHECK FRED SPACE CAL
1979 75			08.1 0 08.38.08.1	FRED SPACE CAL
1979 77		118.356 -45:55	2E IE 0 3E 3E IE 1	THEY STARL CAL
1979 78		118.361 -25:75	2E 1E 0 3E 3E 1 1E	
1979 79		118.346 -10:90	3E 1 OE 2E 3E 1E 1	POSS BEAT B12/CH
1979 81		118.350 5:105		
1979 82	2070:2083		1E 1 0 1E 3E 1 0E	
1979 83 1979 85		118.347 -10:90 118.348 -10:90	3E OE O 2E 3E 1 1E 3E 1 0 2E 3E 1E 1	POSS BEAT B12/CH POSS BEAT B12/CH
1979 86		118.345 -25:75		PUSS BEAL BIZ/CA
1979 87		118.362 -45:55		
1979 87		118.364 -25:75	3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
1979 89	2167:2179	119.453 30:115	1E 1 OE OE 3E OE 1	FREQ SPACE CAL
1979 90	2180:2194		3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
1979 91		118.364 -25:75	32 1 02 32 32 1 12	NORMAL BALLOON MODE
1979 92 1979 92	2208:2213	118.364 -25:75	3E 1 OE 3E 3E 1 1E 3E 0 0 3E 3E 1 1E	NORMAL BALLOON MODE
1979 92	2219:2218			1 DOPPLER SCAN
1979 93		118.356 -45:55		
1979 94	2236:2249		3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
1979 95	2250:2263	118.364 -25:75	3E 1 OE 3E 3E 1 18	NORMAL BALLOON MODE
1979 96		118.364 -25:75	3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
1979 96		118.364 ~25:75		NORMAL BALLOON MODE
1979 97 1979 98		118.364 -25:75 118.364 -25:75	3E 1 OE 3E 3E 1 IE	NORMAL BALLOON MODE
1979 98 1979 99		118.364 -25:75	3E 1 OE 3E 3E 1 1E 3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
		118.364 -25:75	3E 1 OE 3E 3E 1 1E	NORMAL BALLOON MODE
	2332:2344		OE 1 0 0E 3E 0E 1	1 DOPPLER SCAN/ORBIT
			3E 1 OE 3E 3E 1E 1	
1979 103	2360:2374	118.362 -45:55	2E 1E 0 3E 3E 1 1E	
1979 105			IE 1 OF OF 3E OF 1	FREQ SPACE CAL
1979 107	2415:2429	119.455 30:115	OF I O OF BE OF I	FRED SPACE CAL

DATE	ORBITS	SCAN	RANCE (KH)	) !	SIE	ve :	SET	TN	55		
YEAR DAY	(APPROX)	PROG	(APPROX)	AL	A2	A3	81	C1	C2	đ	
		•	•				82				
1979 109	2443:2452	118.356	-45:55	Æ	t			3E			
1979 109	2453:2457	118.364	-25:75	3E				3E			NORMAL BALLOON MODE
1979 110	2457:2471	118.361	-25:75		12			Æ		1E	
1979 111	2471:2484			38				3E			POSS BEAT B12/CH
1979 113	2498:2512	118.364		32				JE		12	NORMAL BALLOON MODE
1979 114								38		Œ	
1979 114		118.364	-	æ				3E		12	NORMAL BALLOON MODE
1979 115 1979 117			-23:/3			OE OE				IE	ALT DOPPLER AND LINE SCANS
1979 119			-25.75			OE					NORMAL BALLOON MODE
1979 120		118.364	_			Œ					NORMAL BALLOON MODE
1979 121	2610:2623	118.350				õ					
1979 122	2624:2630	5.143				0					HOON SCAN
1979 122	2631:2636	119.453	30:115	12	1	0E	œ	3 <b>e</b>	Œ	1	FREQ SPACE CAL
1979 123	2637:2650	119.453	30:115	12		0E					FREQ SPACE CAL
1979 125	2678:2685	5.143				0					MOON SCAN
1979 125		119.454				0					FRED SPACE CAL
1979 126		119.455				0					FREQ SPACE CAL
1979 127 1979 127		119.454				0 05					FRED SPACE CAL Normal Balloon Hode
1979 127		118.356				OE					
1979 130	2736:2747					õ					
1979 131		118.346				Œ		_			POSS BEAT B12/CH
1979 132		118.364				OE				IE	NORMAL BALLOON MODE
1979 133	2775:2788	118.364	-25:75			0E		_		12	NORMAL BALLOON MODE
1979 135	2802:2816	118.347	-10:90	3E	Œ	0	2 <b>E</b>	3E	1	12	Poss Beat B12/CR
1979 136		118.364				0E				12	NORMAL BALLOON HODE
1979 137		118.364			1			3E			NORMAL BALLOON MODE
1979 138		118.345				OE					
1979 139	2857:2875		-43:55			0				12	
1979 141 1979 141	2885:2886 2887:2899	119.100	30:115			0 OE					scan cal check Freq space cal
1979 142	2900:2913	119.454				0					FREQ SPACE CAL
1979 143	2913:2926	119.455	-		1			3E			FREQ SPACE CAL
1979 145	2941:2953	118.363			iE			3E			
1979 146	2954:2967	118.361			12			32			
1979 147	2968:2981	118.366			Œ			3E			
1979 149	2996:3009	118.350	5:105	IE	1	0	12	3E	Œ	1	
1979 150	3010:3023	118.351	5:105		1			3E		OE.	
1979 151	3024:3037	118.347	-10:90		œ			3E		12	POSS BEAT B12/CH
1979 152	3049:3052	5.143			0			32		IE	MOON SCAN
1979 153	3053:3055	5.143	-10.00		0			32		IE	MOON SCAN POSS BEAT B12/CH
1979 153 1979 154	3055:3064 3065:3078	118.348			1	0 02		32 72			PUSS BEAL BIZ/CH
1979 154	3079:3089	118.362				0					
1979 155	3090:3092	5.143	-43.33		0			3E		12	MOON SCAN
1979 156	3093:3097	5.143		3E		ŏ		3E		12	MOON SCAN
1979 157	3106:3120	300,200		32				32			ALT DOPPLER AND LIMB SCANS
1979 158	3121:3130	300.200		32				32			ALT DOPPLER AND LIMB SCANS
979 158	3131:3134	118.346	-10:90	3E	L	Œ	2E	3E	12	1	POSS BEAT B12/CH
1979 159	3135:3147	300.201		3E	1	0E	Z	3 <b>e</b>	12	1	DOPPLER SCANS
1979 161	3161:3175	25.101			Œ			32		12	
1979 162	3175:3188	118.361			12			3E			
1979 163	3189:3203	118.366			OE		-	3E			
1979 165		118.364	-25:75			OE					NORMAL BALLOON MODE
1979 166	3231:3244	301.202		2 <b>E</b>	1	UE	2 <b>E</b>	3 <b>E</b>	UE	L	ALT DOPPLER AND LIMB SCANS

DA' Year i		ORBITS	SCAN PROG	RANGE (10) (APPROX)	)	SIE	VE :	SET		3	~	
TEAK		(APTROX)	PROG	(APTRUA)		~*	~	82				
	147	3944 9959			-	~	•	10		,	1 🕶	
1979 1979		3244:3258	26.101			100 100			運			SPACE VIEW/BB. MIRROR AT AZ BACKWARD LIMIT DOPPLER SCANS IN SPACE VIEW
1979		3272:3285 3286:3299	26.102	-25:75					32			MIRROR AT AZDAUTH BACKWARD LDATT
1979		3300:3313	118.362						32			MIRROR AT AZIMUTH BACKHARD LIMIT
1979		3327:3339		20:125			-					FRED SPACE CAL
1979		3340:3341	129.100						32			SCAN CAL CHECK
1979	174	3342:3355	129.475	20:125	Œ	1	0	Œ	32	Œ	1	FREQ SPACE CAL
1979		3355:3368	129.473						Æ			PREQ SPACE CAL
1979		3383:3396						-	Æ		-	
1979   1979		3397:3410 3411:3423	118.361						浭			
1979		3438:3451	118.350			1			312 312			
1979		3452:3465		5:105								
1979 1		3466:3479							32		-	POSS BEAT B12/CH
1979 1	185	3493:3506	118.348	-10:90	3E	1	0	2E	3E	12	1	POSS BEAT BIZ/CH
1979 1		3507:3520	118.345	-25:75					3E			
1979 1		3521:3535							Æ		-	
1979 1		3549:3562							32			FREQ SPACE CAL
1979 I 1979 I		3563:3575 3576:3589							JE JE			FRED SPACE CAL FRED SPACE CAL
1979 1		3606:3617	118.365						32			FREU SPACE CAL
1979 1		3618:3632	118.361						32			
1979 1	95	3632:3646	118.366		-				3E			
1979 1		3669:3672	118.350		12	1	0	12	32	œ	1	
1979 1		3673:3687	118.351			1			32			
1979 1		3687:3700	118.365			12			32			
1979 2 1979 2		3701:3704 3714:3728	118.365						3E 3E	_	-	NORMAL BALLOON MODE
1979 2		3728:3741							32	-		NUMPER BELLOUN HUDE
1979 2		3742:3758	118.362						3E	_		
1979 2	05	3769:3782	129.473	20:125	12	1	Œ	Œ	3E	Œ	1	FREO SPACE CAL
1979 2		3783:3784	129.100						3E			SCAN CAL CHECK
1979 2			129.474						JE			FREO SPACE CAL
1979 2		3797:3811	129.475						32			FREO SPACE CAL
1979 2 1979 2		3828:3838 3839:3852	118.365 118.364						3E 3E			
1979 2		3852:3867							3E			Normal Balloon Mode Normal Balloon Mode
1979 2		3880:3894	118:350						32			
1979 2			118:351			i			3E			
1979 2	15		118.347		3E	œ			3E			POSS BEAT B12/CH
1979 2	-	3936:3949	118.103			0			Œ	. –	-	
1979 2			118,104	-					12			SPECTAL CI PMC CAL MODE
1979 2		3963:3967		-					32		12	NORMAL BALLOON MODE
1979 2 1979 2		3968:3977 3991:4004	118:105 · 118.475						212 312			SPECIAL C1 PMC CAL MODE FREO SPACE CAL
1979 2				20:125								FREQ SPACE CAL
1979 2		4018:4032	118.475		0E				3E			FREO SPACE CAL
1979 2		4047:4060	118.476			i			3E			FREQ SPACE CAL
1979 2		4060:4073	118.476		12				3E			FREQ SPACE CAL
1979 2		4073:4087	118.476	20:125	lZ	1			3E		0E	FREQ SPACE CAL
1979 2		4101:4115	27.106			0			3E		Œ	WIDE LINE SCAN FOR WE STRAYS CHECK
1979 2		4115:4129	118.351			1			3E		OE	
1979 2 1979 2		4129:4143	118.347			OE .			3E			POSS REAT B12/CR
1979 2		4156:4171 4171:4184	118.348 · 118.345 ·			1			312 312			POSS BEAT B12/CH
1979 2			118.362						3E			
••••					-	-		_	-	÷		

<b></b>			RANCE (ICH)		e T <b>F</b>	<b>.</b>	51.11				
DATE	ORBITS	SCAN	(APPROX)	, : • • •	A7	۲۹	81	C1	<u>.</u>	C	
YEAR DAY	(APPROX)	PROG	(APPROX)	AI	74	2	82	Ç1		3	
1979 236	4211:4212	129.100		12	0	0	3E	3E	1	0E	SCAN CAL CHECK
1979 237	4212:4225	129.473	20:125					ĴĒ	Œ	1	FREQ SPACE CAL
1979 238	4225:4238	129.474	20:125	IZ	Œ	0	œ	32	1	œ	PREQ SPACE CAL
1979 239	4238:4252	129.475		ÔE	1	0	Œ	3E	œ	1	FREQ SPACE CAL
1979 241	4267:4281	118.365	-25:75	æ	12	0	32	Æ	12	1	
1979 242		118.361	-25:75		12			3E			
1979 243		118.366			Œ			3E			
1979 245		118.350	5:105		1			32			
1979 246 1979 247		118.351	5:105		l Ore			3E 3E			POSS BEAT B12/CH
1979 247		118.347			1			3E			POSS BEAT BI2/CR
1979 250		118.345		3E				32			1000 Barr Distan
1979 251		118.362			12			32			
1979 253		129.473		12	1	ÖE.	0E	3E			FREQ SPACE CAL
1979 254		129.474			œ			3E			FREO SPACE CAL
1979 255	4461:4474	129.475	20:125	Œ	1	0	0Ľ	3 <b>E</b>	Œ	1	FREQ SPACE CAL
1979 257		118.365			12			36			
1979 258		118.361			IE			3E			
1979 259	4515:4529				Œ			3E			
1979 261		118.350	5:105	12		0		3E			
1979 262		118.351	5:105	12		0		312 312			
1979 263 1979 265	4571:4585 4599:4602	118.348		32 32		0		3E			
1979 265		118.348		3E				32			
1979 266		118.345			i			3E			
1979 267	4626:4640	118.362			iE			3E			
1979 269		129.100			0			3E			SCAN CAL CHECK
1979 269	4654:4668	129.473	20:125	12	1	Œ	Œ	3E	Œ	1	SPACE VIEW EACH SCAN CYCLE
1979 270	4668:4682	129.474	20:125					3E			SPACE VIEW EACH SCAN CYCLE
1979 271	4682:4695	129.475			1			3 <b>E</b>			SPACE VIEW EACH SCAN CYCLE
1979 273	4709:4723	118.365			IE			32			
1979 274	4723:4737	118.346						3E			
1979 275	4737:4751	118.366			OE .			3E			
1979 277 1979 278	4765:4778 4778:4792	118.350	5:105 5:105		1			3e 3e			
1979 279	4792:4806	118.347			ÓE.			3E		IE	
1979 281	4820:4834	118.348			1			32			
1979 282	4834:4847	118.345			-	-		3E			
1979 283	4847:4861	118.362			12			3E			
1979 285	4875:4889	129.473	20:125	1E	1	0E	0E	3E	Œ	1	SPACE VIEW EACH SCAN CYCLE
1979 286	4889:4903	129.474			œ			3E			SPACE VIEW EACH SCAN CYCLE
1979 287	4903:4917	129.475			1			32			SPACE VIEW EACH SCAN CYCLE
	4930:4944	118.365			IE			38			
1979 290	4944:4958	118.361			IE			38			
	4958:4971							3E 3E			
	4986:4999 5000:5013	118.350	5:10 <b>5</b> 5:10 <b>5</b>		1			32			
1979 295	5014:5027	118.347			oe.			死			
1979 297	5041:5055	118.348		3E		ŏ		32			
1979 298	5055:5068	118.345		3E				3E			
1979 299	5068:5082	118.345		3E				32			
1979 300	5095:5096	129.100		12	0			3E			scan cal check
1979 301	5096:5110	129.473	20:125	1E	1	Œ		3E			SPACE VIEW EACH SCAN CYCLE
1979 302	5110:5124	129.474			Œ			3 <b>e</b>			SPACE VIEW EACH SCAN CYCLE
1979 303	5124:5138	129.475			1			3E			SPACE VIEW EACH SCAN CYCLE
1979 305	5152:5165	118.365	-25:75	ZE	12	0	3E	3E	12	1	

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	DATE	ORBITS	SCAN	RANGE ( KH	) SI	ÉVE	SET	TIN	3		
v		(APPROX)	PROC	(APPROX)	AI A	2 13	81	C1	C2	C3	
Ľ				(			82				
1	979 306	5165:5179	118.347	-10:90		20					
		5179:5193				۳ ٥					
		5208:5220				Œ					
		5220:5234				OZ					
1	979 311	5234:5248	118.364	-25:75		OE					
		5262:5276				0					
		5276:5290				_ 02					
		5290:5303				E 0			-		SPACE VIEW EACH SCAN CYCLE
		5317:5331				OE					ROCKET PLIGHT COVERAGE
		5331:5345 5345:5359							-		SPACE VIEW EACH SCAN CYCLE
		5373:5386				εŎ		-			SPACE VIEW EARLY SCHEV CICLE
		5386:5400				ĒÖ					
		5400:5414				ĒÖ					
		5428:5442			12 1						
		5442:5455		5:105	12 1						
		5455:5469				E O					
		5483:5497				Ō					
			118.345			Ű.					SYNC SLIP
		5511:5525	118.362	-45:55	2E 1	EO	32	32	t	12	SEVERAL SCAN LOGIC JUMPS
19	79 333	5538:5552	118.364	-25:75	3E 1	OZ	æ	3E	1	IE	NO SCANNENG
19	79 334	5552:5566	118.364	-25:75	3TE 1	Œ	3E	3E	1	12	
		5566:5579		-25:75	3 <b>E</b> 1	Œ	Æ	32	1	lE	NORMAL BALLOON MODE
-		5594:5596				0					GSPC MESSING ABOUT WITH PCL
			118.365			e o					
			118.361			EO					
			118.366			60					
		5635:5649				_ 02					Poss beat Biz/CH Rocket Flight Mode
			118.366	-		20					
			118.351		IE I						
			118.347			<b>E</b> 0					POSS BEAT BI2/CH
		-	118.348			0					POSS BEAT BI2/CH
			118.345	-23:73		012 012					NO. COMMITTE BOLL OF THE BOLL MAL PROPERTY
			118.345	-75.75	3E 1						NO SCANNING FOLLOWING PCL MALFUNCTION
	-		118.362			E O					
			129.100			0					SCAN CAL CHECK
			129.473	20+125	12 1						FREQ SPACE CAL
		5773:5787			IE O						FREQ SPACE CAL
			118.346		3E 1						POSS BEAT BI2/CH ROCKET FLIGHT MODE
			118.346			0E					POSS BEAT BIZ/CH ROCKET FLIGHT MODE
			118.365		2E 1						
				-25:75		E O					
			118.366		-	EO					
		5870:5884									
		5884:5898		5:105							
		5898:5912									POSS BEAT B12/CH
		592 <b>5:5939</b>				0	2 <b>E</b>	3E	12	1	POSS BEAT BI2/CH
		5939:5953	118.345	-25:75							
19	79 363	5953:5966	118.362	-45:55	2 <b>E</b> 1	ΕO	3 <b>E</b>	3E	1	1 <b>E</b>	
19	79 365			20:125	12-1	Œ	0E	3E	Œ	1	FREO SPACE CAL
19	80 1	5994:6008	129.474	20:125	1 <b>E</b> 0	E 0	0 <b>E</b>	3E	1	Œ	FREO SPACE CAL
	80 2			20:125							FREQ SPACE CAL
	80 3		129.100			0					SCAN CAL CHECK
	80 4	6036:6049									
19	80 5	6050:6063	118.351	5:105	1 <b>E</b> 1	0	12	3E	I	OF.	

DATE	ORBITS	SCAN B	ANCE (IN)	STE	TE S	EIII	NCS		
YEAR DAY	(APPROX)	PROG (	APPROX)	AL A2	N	BI C	1 C2	C	
	(ALL WAY					82			
				_		_			
1980 6		118.347 -		JE OE					POSS BEAT BI2/CH
1980 8		118.348 -		312 1					POSS BEAT BI2/CH
1980 9		118.345 -		第1	_		-		
1980 10	6119:6133	118.362 -	45:55						FREO SPACE CAL
1980 12	6146:6160	129.473	20:125		012	ທະ ສ ທ	E UE. 191		FREO SPACE CAL
1980 13		129.474		012 1	Å i	012 JU 012 JU		1	FREQ SPACE CAL
1980-14 1980-16		129.475 118.365 ~		22 12					
1980 17		118.361 -		ZEIE					
1980 18		118.366 -		38 08					
1980 20	6257:6271		5:105						
1980 21	6271:6284	118.351		1 <b>E</b> 1					
1980 22		118.347 -		3E OE	0	2E 31	21	12	POSS BEAT BI2/CH
1980 24	6312:6326			312 I					POSS BEAT BI2/CH
1980 25	6329:6340			38 1					
1980 26	6340:6354			ZE IE					FREQ SPACE CAL
1980-28 1980-29	6367:6381 6381:6395								FREQ SPACE CAL
1980 30	6396:6409								FREQ SPACE CAL
1980 32	6423:6437			ZEIE					FREQ SPACE CAL. SEE NOTE 1.
1980 33		129.361 -		22 12					FRED SPACE CAL. SEE NOTE 1
1980 34		129.366 -		38 08					FRED SPACE CAL. SEE NOTE !
1980 36		129.350	5:105						FREQ SPACE CAL. SEE NOTE 1
1980 37	6492:6497	129.351	5:105	1 <b>E</b> 1	0	12 J	21	Œ	FRED SPACE CAL. NO SCANNING
1980 37	6497:6505	129.351	5:105	1 <b>E</b> 1	0	12 J	<b>2</b> 1	Œ	FREQ SPACE CAL. SEE NOTES 1 AND 2
1980-38	6505:6519	129.347 -		Je oe					FRED SPACE CAL. SEE NOTES 1 AND 2 AND 3
1980 40	6533:6547			3 <b>E</b> 1					FREO SPACE CAL. SEE NOTES 1, 2 AND 3.
1980 41		129.345 -		32 1					FREQ SPACE CAL. SEE NOTE 1
1980 42		129.362 -		ZE IE					FREQ SPACE CAL. SEE NOTE 1
1980 44 1980 45	6589:6603 6603:6617	129.484 -		12 1 12 02					FREQ SPACE CAL. Freq space cal.
1980 46	6617:6629			OEL					FREO SPACE CAL.
1980 46		129.100		IEO					SCAN CAL CHECK.
1980 48		118.365 -		22 12					
1980 49		118.361 -		2E 1E		3E 31			
1980 50		118.366 -		JE OE					
1980 52	6699:6713	118.350	5:105	1E 1	0	1E 31	e ne	1	
1980 53		118.351	5:105						
1980 54		118.347 -		32 02					POSS BEAT B12/CH.
1980 56		118.348 -		312 I					POSS BEAT B12/CH.
1980 57 1980 58		118.345 - 118.362 -		3E 1 2E 1E					
1980 60		129.473							FREO SPACE CAL
1980 61	6823:6837	129.484 -		IE OE					FREO SPACE CAL
1980 62	6837:6851	129.475							FREO SPACE CAL
1980 64	6865:6879	118.365 -		ZE IE		3E 31			
1980 65	6879:6892	118.361 -		2E 1E		3E 31			
1980 66	6893:6907	118.366 -		JE OE		2E 38			
1980 68	6920:6934	118.350	5:105	1E 1	0	1E 31	2 O 2	1	
1980 69	6934:6948	118.351		ie i		1E 31			
1980 70	6948:6962	118.347 -		je oe		2E 31			POSS BEAT BI2/CH
1980 72	6975:6990	118.348 -		3E 1					POSS BEAT BI2/CR
1980 73	6990:7003	118.345 -				3E 31			
1980 74	7003:7017	118.362 -		2E IE					BAD SCAN PATTERN ORB. 7010
1980 76	7031:7045	129.483 -				3E 31			FRED SPACE CAL
1980 77	7045:7058	129.484 -	10:90	IE OE	υ.	אניכו	5 1	0E	FREQ SPACE CAL

DA 77	CRATER		DI) SIEVE SETTINGS	
DATE YEAR DAT	ORBITS (APPROX)	PROG (APPRO	() AI AZ AJ BI CI CZ CJ	
			82	
	_			
1980 78	7058:7071	129.485 -10:90	012 1 0 012 312 012 1	FREQ SPACE CAL
1980 78	7071:7072		1E 0 0 3E 3E 1 0E	scan cal check
1980 80	7086:7094	26.101	32 02 0 32 32 1 12	
1980 80 1980 81	7095:7100 7100:7114	26.101 26.101	377. OPE 0 372. 372. 1 12 372. 072. 0 372. 372. 1 12	
1980 82	7114:7127			
1980 84	7141:7155			
1980 85	7155:7168	118.351 5:105		
1980 86	7168:7182	118.347 -10:90	3E OE O 2E 3E 1 1E	POSS BEAT BI2/CH
1980 88		118.356 -45:55	3E 1 OF 3E 3E 1E 1	
1980 89		118.356 -45:55	3E 1 OE 3E 3E 1E 1	
1980 90		118.356 -45:55	3E 1 OE 3E 3E 1E 1	
1980 92 1980 93	7251:7265		3E 1 1E 3E 3E 1E 1 3E 1 1E 3E 3E 1E 1	
1980 94	7279:7293	118.368	3E 1 1E 3E 3E 1E 1	
1980 95	7293:7307		JE 1E 0 JE JE 1E 1	
1980 96	7307:7321	118.365 -25:75		
1980 97	7321:7335	118.361 -25:75	2E 1E 0 3E 3E 1 1E	
1980 98	7335:7349	118.366 -10:90	3E 0E 0 2E 3E 1E 1	
1980 100	7362:7376			
1980 101	7376:7390			
1980 102 1980 104	7418:7431	118.347 -10:90	3E OE O 2E 3E 1 1E 3E 1 O 2E 3E 1E 1	POSS BEAT BIZ/CH
		118.348 -10:90 118.345 -25:75	3E 1 OF 3E 3E 1E 1	Poss reat Bi2/CR
1980 106	7445:7459	118.362 -45:55	2E 1E 0 3E 3E 1 1E	
1980 108	7473:7487	36.101	3E OE O 3E 3E 1 1E	BB/SPACE CAL AUX HEATERS ON
1980 109	7487:7500	36,101	3E OE O 3E 3E 1 1E	BB/SPACE CAL
1980 110	7500:7513	129.485 20:125		FREQ SPACE CAL
1980 110	7513:7514	129.100	1E 0 0 3E 3E 1 0E	SCAN CAL CHECK
1980 112		118.365 -25:75	2E 1E 0 3E 3E 1E 1	
1980 113		118.361 -25:75	2E 1E 0 3E 3E 1 1E	
1980 114	_	118.366 -10:90	32 02 0 22 32 12 1	
1980 116 1980 117	7584:7597 7597:7611	118.350 5:105 118.351 5:105		
1980 118		118.347 -10:90	3E OE O 2E 3E 1 1E	BAD ROLL COMP PHASING POSS BEAT BI2/CH
1980 120	7639:7653	118.348 -10:90	3E 1 0 2E 3E 1E 1	POSS BEAT BI2/CH
1980 121		118.345 -25:75	3E 1 OF 3E 3E 1E 1	FUSS REAL BLE/CR
1980 122		118.362 -45:55	22 1E 0 3E 3E 1 1E	
1980 124	7694:7708	129.484 20:125	1E OE O 3E 3E 1 OE	FREQ SPACE CAL
1980 125	7708:7722	129.484 20:125	1E OE 0 3E 3E 1 OE	FREQ SPACE CAL
1980 126		129.485 20:125		FREQ SPACE CAL
1980 128	7750:7763	118.365 -25:75	2E 1E 0 3E 3E 1E 1	
1980 129	7763:7777	118.361 -25:75	2E 1E 0 3E 3E 1 1E	
1980 130 1980 132	7777:7791	118.366 +10:90	3E OE O 2E 3E 1E 1 1E 1 O 1E 3E OE 1	
1980 132	7818:7832		1E 1 0 1E 3E 1 0E	
1980 134		118.347 -10:90	3E OE O 2E 3E 1 1E	POSS BEAT BI2/CH
1980 136	7860:7874	118.348 -10:90	3E 1 0 2E 3E 1E 1	POSS BEAT BIZ/CH
1980 137		118.348 -10:90	31E 1 0 21E 31E 1	POSS BEAT BIZ/CH
1980 138		118.362 -45:55	2E 1E 0 3E 3E 1 1E	
1980 140		129.484 20:125		FREO SPACE CAL
1980 141		129.485 20:125		FREO SPACE CAL
1980 142		129.485 20:125		FREQ SPACE CAL
1980 142		129.100	1E O O 3E 3E 1 OE	SCAN CAL CHECK
1980 144		118.365 -25:75	22 12 0 32 32 12 1 27 12 0 32 37 1 12	
1980 145	/ 704 : / 998	118.361 -25:75	2E IE O 3E 3E 1 1E	

			RANGE (RM)								
DATE	ORBITS	SCAN PROG	(APPROX)	<b>^</b>	A2 /		611 B1	CI	~2 C2	<b>C3</b>	
YEAR DAY	(APPROX)	FRUG					82				
1980 146		118,366	-10:90		02 (				12		
1960 148			5:105						0E		
1980 149		118.351	5:105						1		POSS BEAT BI2/CH
1980 150		118.347	-10:90 5:110		02 (				1 05		FREQ SPACE CAL
1980 152			5:110						1		FRED SPACE CAL
1980 153		129.508			0 0				iz		FREQ SPACE CAL
1980 156	8136:8150				IE C				iE.		FREQ SPACE CAL
1980 157	8150:8164	129.513	20:125	Z	0 0	) 2	2E	3E	1	12	FREO SPACE CAL
	8164:8178				072 C	) 3	3e i	3E	IE	1	FREQ SPACE CAL
		118,500			0 0				œ		
	8206:8219				IE C						
	8219:8233				0 0				1		
	8247:8261 8261:8274				1E 0				l.		
	8274:8288				OE C				ie		
	8302:8316										FREQ SPACE CAL
	8316:8330				OE C				1		FREQ SPACE CAL
1980 170	8330:8343	129.508	5:110	3E	0 0	<b>)</b> 3	3E	3E	12	1	FREQ SPACE CAL
			40:145								FRED SPACE CAL
	8371:8385										FREQ SPACE CAL
	8385:8397		40:145	IZ							FREQ SPACE CAL. SOME BAD DATA
	8398:8399		40.148		0 0						scan cal check freq space cal
		129.511	40:145		1 0		_		12	-	NO SCANNING
	8413:8427		-10:90		0 0						
	8427:8440				12 0				ĩ		
	8440:8454				0 0				i		
1980 180	8468:8482	118.503	-10:90	22	1E 0	) 1	te :	3E	Œ	1	
1980 181	8482:8496	118,504	-10:90	3E	0 0	) 2	ZE. I	3E	1	12	ROLL COMP SLIP. SOME DATA LOST
	8496:8509				OF				IZ		
1980 184	8523:8537		5:110		0 0				œ.		FRED SPACE CAL
		129.507			020				1		FREQ SPACE CAL
1980 186 1980 188		129.508	5:110 20:125		00				12 12		FREQ SPACE CAL FREQ SPACE CAL. SEE NOTE 5.
	8593:8606	129.513			0 0				1		FRED SPACE CAL. SEE NOTE 6
		129.514			OE C				iz		FREQ SPACE CAL
1980 192	8634:8648				0 0					-	
1980 193		118.501		22	12 (	) 2	2E	3E	1	œ	
1980 194		118.502			0 0				1		
	8689:8703				12 (				12		BEAT IN CHANNEL A3P
1980 197	8703:8717				0 0				1 12		
1980 198 1980 200	8717:8731 8745:8758	118.505	5:110		0 0	-		_	15		FREQ SPACE CAL
	8758:8772										FREO SPACE CAL
	8772:8786										FREO SPACE CAL
	8799:8813										FREQ SPACE CAL
1980 205	8813:8827										FREQ SPACE CAL
1980-20 <b>6</b>	8827:8840	129.511		12	1 0	) (	Œ	3E	12	1	FREQ SPACE CAL
1980 207	8840:8841	129.100			0 0						SCAN CAL CHECK
	8855:8869	118.500			0 0				œ		
		118.501			IE C				1		SEE NOTE 4
	8882:8896 8910:8924	118.502			0 0				1		
1980 212		118,504			0 0				1		
	8938:8952				ÕE C				iE		
						•			_		

DATE	ORBITS	SCAN	RANCE ( RM	)	SIE	VE	SET	TIN	3										
YEAR DAY		PROG	(APTROX)	้งเ	A2	. 13	81	CI	2	C									
							82												
1980 216	8965: <b>8979</b>	129.506	5:110						0E			SPACE							
1980 217	8979 <b>: 8993</b>	129.507			Œ				1		,	SPACE							
1980 218	8993:9007	129.508			0		Æ		-			SPACE							
1980 220	9021:9034	129.512					12					SPACE							
1980 221	9034:9048	129.513			0		22					SPACE	-						
1980 222	9048:9062		20:125				22				FREQ	SPACE	CAL						
	9076:9090	118.500			0		)E												
1980 225 1980 226	9090:9104 9104:9118				1E 0				l										
	9131:9145						12		1										
	9145:9159				0				1										
1980 230	9159:9173			-	ÖE				ie										
	9186:9200		5:110				22				FREO	SPACE	CAL						
1980 233	9200:9214	129.507					32					SPACE							
1980 234	9214:9228	129.508					3E					SPACE							
1980 236	9242:9254	129.509	40:145	0E	1	0	12	3E	02	1		SPACE							
1980 237	9254:9269			12			12				FREO	SPACE	CAL						
1980 238	9269:9281	129.511	40:145	12	1	0	02	X	1E	1	FREQ	SPACE	CAL						
1980 238	9281:9283			12	0	0	32	æ	1	œ		SCAN	CAL	CRECK					
	9297:9311	118.500			0				œ										
1980 241		118.501			12		Z												
1980 242					0				1										
	9352:9366 9366:9380						22				-	SPACE							
		129.506			0				œ		•	SPACE							
1980 248	9408:9421	129.506			0			-	30			SPACE		AZIM	9 06	G PW			
		129.506	-		ŏ				OE OE		. – 4	SPACE		AZIM	1 11				
	9436:9449	129.506			õ				OE		•	SPACE		AZIM					
	9463:9477	129.506			ŏ				0E		-	SPACE		AGEN	0.06	A3 (W			
1980 253	9477:9490	129.506			ŏ				Œ		•	SPACE		AZTH	3 0				
		129.506			Ō				ne.			SPACE							
1980 256	9518:9532	129.506			0				Œ			SPACE							
1980 257	9532:9545	129.506			0				Œ			SPACE		AZIM	12 0	EG BK			
1980 258	9546:9558	129.506	5:110	3E	0	0	22	32	0E	1	-	SPACE						NOTE	6
1980 260	9572:9586	129.514	5:110	3E	œ	0	Æ	32	12	1	FREQ	SPACE	CAL						
1980 261	9586:9600	129.514	5:110	Æ	Œ	0	3E	3E	IE	1	FREQ	SPACE	CAL	AZDH	12 0	EG FW			
		129.514	5:110	Æ	OE,	0	3E	3E	ìΕ	1	FREQ	SPACE	CAL	AZIN	9 DE	GFW			
1980 264		129.514	5:110						IE		FREQ	SPACE	CAL						
1980 265		129.514			Œ				1E			SPACE		AZTH					
		129.514			90				12			SPACE		AZTM	6 DE	G FW.	SEE	NOTE	4
1980 268		129.514			<b>0</b> 2				12			SPACE							
		129.514			Œ				1E			SPACE		AZTH	-				
1980 270		129.514			Œ				12			SPACE		AZ IM	JDE	GBK			
1980 272		129.514			OE.				IE.			SPACE							
1980 273 1980 274	9753:9767 9767:9781	129.514	5:110								•	SPACE							
1980 276		129.514	5:110 5:110		0E				ie Ie			SPACE		AZ IM	4 05	G 5K			
1980 277		129.514	5:110						IE			SPACE		AZIN	6 100	C 17			
1980 278		129.514			OE				12			SPACE		AZIM					
1980 280		129.506			õ				OE			SPACE		1944 ANT					
1980 281		129.507			œ.				1			SPACE							
1980 282		129.508			õ				iz			SPACE							
		129.100			Õ				ī					CHECK					
1980 284	9917:9919	29.513			ñ		2E				FREO	SPACE							
1980 285	9919:9933	29.513			0				1			SPACE							
1980 286	9933:9946	29.514			OE.				١Ľ			SPACE							

<b>D</b> .	ATE	ORBITS	SCAN	RANCE ( KM	) 5	IE	7E :	SET	TN	3	
YEAR			PROG	(APPROX)	Â	A2	AJ	81	C1	C2	C
				(				82			
						_	_				
1980		9960:9974	118.500		32		0	12		œ	
1980			118.501		22 31	31	0	涩涩		1	oe 1e
1980			118.502	-10:90		12		312		-	1
1980 1980		9996:10002 10016:10030	118.363	-10.90		12		12			i
1980		10030:10043			3E	-	ō	22	38		12
1980		10043:10057	118.505		3E	Œ	0	32	3E	12	1
1980	-	10071:10085	129.506	5:110	3E	0	0	2E	3E	œ	1
1980	297	10085:10098		5:110			0	32		1	0E
1980			129.508	5:110		0	0	3E	3E		1
1980		10126:10140		20:125		1	30	OE.		œ	1
1980		10140:10154		20:125 20:125		01E	0 0	OFE.		1	0E
1980		10154:10168 10181:10195			32		ŏ	OE 1E	范		1
1980		10195:10209			-	iz.		ZE		1	OE.
1980	<u>د</u>	10209:10223			32		ō	æ		i	12
		10237:10251			2E	12		12		-	1
1980	309	10251:10265	218.504	-10:90	3E	0	0	ZE	æ	1	12
1980	310	10265:10279	218.505	-10:90	3E	œ	0	38	3E	12	1
		10293:10306		5:110	3E		0	2E	-	œ	-
1980		10306:10320		5:110		0E		3E		1	0E
		10320:10334		5:110		0	0	3E		IZ	
1980 1980		10347:10361		20:125	212 212	12	0	IE 2E	32 32	1E 1	1 12
1980	• -	10375:10389		5:110		0E		3E	光	1	1
		10402:10417			32		ŏ	IE	3E		i
1980		10417:10430			22	1E		Z		1	œ.
1980		10430:10444	218.502	-10:90		0	Ō	Æ		ī	12
1980	324	10458:10472	218.503	-10:90	2 <b>E</b>	iz	0	IE	3E	Œ	1
1980	325	10472:10486	218.504	-10:90	3E	0	0	2E	3E	1	12
1980		10486:10499				Œ	0	3E	3E	12	1
1980		10513:10527	229.506	5:110		0	0	2E		œ	1
1980		10527:10541	229.507	5:110		Œ	0	32		1	OE .
1980		10541:10555	229.508	5:110		0	0	32	32	1E	1
1980 1980		10568:10582	229.509	40:145 40:145		1	0	IE IE	涯	90 1	1 072
1980		10597:10610		40:145		ì	ŏ	0E	32	12	1
1980		10624:10638				ò	ŏ	12	3E		i
1980	337		218.501	-		lE	Ō	22	32	1	OF.
1980	338	10652:10665	218.502	-10:90	3E	0	0	3E	3E	1	12
1980		10679:10693			22	lE	0	IE		Œ	1
1980		10693:10707				0	0	22	32	1	IE
1980		10707:10720					0	32		IE	1
1980	-	10734:10748		5:110	32		0	2E	JE JE	OE.	1 0e
1980		10748:10762 10762:10776		5:110 5:110	32 32		0	3E 3E		12	-
		10790:10803		20:125		12		12			1
		10803:10817		20:125	22		ŏ	22	3E	1	ie
		10817:10831		5:110		Ŭ.		32		12	1
		10845:10859			32		Ō	IE		0E	1
		10859:10872		-10:90		12	0	2E	3E	1	0E
		10872:10886			32	0	0	3E	3E	1	12
		10900:10914			22	12		12		Œ	1
		10914:10928			32		0	ZE	32	1	1E
		10928:10941				œ		32	32	12	1
1980	360	10955:10969	229.506	5:110	32	U	0	2E	JE	07.	L

DATE	ORBITS	SCAN	RANGE (KH	) :	SIE	Æ:	seri	TN	3	
TEAR DAT	(APTROE)	PROG	(APPROX)							C
			<b>,</b>				82			
					~	^		-		~
1980 361	10969:10983		5:110		02 0	0	32	32 32	1 12	0E
1980 362 1980 364	10963:10997 11011:11025		5:110 40:145	3E OE	1	õ	3E 1E	運	UE OE	1
1980 365	11025:11025		40:145		i	ŏ	IE	32	1	OE.
1980 366	11039:11051	229.511	40:145	IE	i	ŏ	OE.	ĴE.	iz.	1
1980 366	11051:11052			12		ō	32		i	ÖE .
1981 2	11066:11080	218.500	-10:90	Æ		Ō	12	32	-	1
1981 3	11080:11094	218.501	-10:90	2E	12	0	2E		1	Œ
1981 4	11094:11108	218.502		32	0	0	3E	3E	1	12
1981 6	11122:11135			22	IĽ	-	1E	æ		1
1981 7	11135:11148			32	0	0	2E		1	12
1981 8 1981 10	11148:11163				OE		32			1
1981 11	11191:11205	229.506	5:110 5:110	32 32	0	0	že Je	3E 3E	082. 1	1 012
1981 12	11205:11218		5:110	-	0	ŏ	3E 3E		-	1
1981 14	11232:11246		40:145	Œ	ĭ	ŏ	12	32		i
1981 15	11246:11259		40:145	IE	i	ŏ	12	3E	1	OZ
1981 16	11259:11274	229.511	40:145	lĽ	1	0	Œ	3E	12	1
1981-18	11288:11301	218.500	-10:90	3E	0	0	12	æ	Œ	1
1981 19	11301:11315			2E	12	_	Z	Æ	1	0E
1981 20	11315:11329				0	0	32	3Ľ	1	IE
1981 22 1981 23	11342:11356 11356:11370			2E	IE			3E		1
1981 24	11370:11384				0 OZ	0	22	3E	1	12
1981 24	11398:11412		5:110			õ	312 272	_	ie Oe	1
1981 27	11412:11425		5:110		-	ŏ	32	32	1	0E
	11425:11439		5:110	2.7	0	õ		JE	iz	1
	11453:11467		40:145	OE		Ō		3E		i
1 <b>981 31</b>	11467:11481	229.510	40:145	1E	1	0	12	3E	1	0E
1981 32	11481:11494	229.511	40:145	12	1	0	œ	3E	12	I
	11494:11495				0	0	3 <b>E</b>	3E	1	Œ
	11508:11522					0		3E		I
			-10:90	22		0	22	3E	1	œ
	11536:11550					0		32	1	1E
	11564:11577 11577:11590		-10:90	狸		0	1E 2E		072 1	1 12
	11591:11605				OE .		3E		-	1
	11619:11632		5:110	JE JE		ŏ		32		1
	11632:11646		5:110			ŏ	3E	3E	1	œ.
	11646:11660		5:110	32		Ō			12	1
1981 46	11674:11688	229.509	40:145	0E	1	0	1E	32	Œ	1
1981 47	11688:11702	229.510	40:145	12	1	0	1E	3E	1	0E
	11702:11716		40:145	IE	1	0	Œ	3E	lE	1
1981 50	11730:11743	218.500	-10:90			0	lE	3E	Œ	1
	11743:11757				12			3E		Œ
	11757:11771			3E		0	3E			18
	11785:11798 11798:11812				1Z			3E		1
	11812:11826			3E 18	U De	0		32 32		12 1
	11833:11838			3E		0			15	I IE
	11840:11854			3E		ŏ		JE		12
	11854:11867		5:110	3E				312		OE
	11867:11881		5:110	3E		ŏ		3E		
	11895:11909		20:125		IE			3E		
1981 63	11909:11923	229.513	20:125	22		0		3E		12
1981 64	11923:11937	229.514	5:110	3 <b>E</b>	Œ	0	3 <b>E</b>	3E	IE	L

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•	ATE	ORBITS	SCAN	RANGE (RM	) 9	IE	Æ S	ET	TIN	3	
YEAR			PROG	(APPROX)							C3
								82			
1981	66	11951:11964	218.500	-10:90	3E	0	0	12	3E	Œ	1
1981	67	11964:11978	218.501	-10:90	2E	12	0	22	T	1	0E
1981	68	11978:11992	218.502	-10:90	Æ	0	0	Æ		1	12
1981	70	12006:12019	218.503	-10:90	22	i E	0	ΙĒ	3E	0E	1
1981	71	12019:12033		-10:90	3 <b>E</b>	0	0	22	32	1	12
1981	72	12033:12047		-10:90	3E	œ	0	Æ	æ	lE	1
1981	74	12061:12075		5:110		0	0	2 <b>E</b>	32	0E	1
1981	75	12075:12068		5:110		0E	0	32	32	1	OE.
1981	76	12088:12102		5:110		0	0	32	312	IE.	1
1981	78	12116:12130		40:145	OE	1	0	12		œ	1
1981	79		229.510	40:145	12	1	0	1E	32	1	0E
1981	80	12144:12157	229.511	40:145	12	1	0	OE.	32	12	1
1981	80		229.107	10.00	IE	0	0	32	32	1	0E
1981 1981	82 83	12172:12186			3E	O IE	õ	ie Ze		02 1	i Orz
1981	84		218.501 218.502		2E 3E	0	ŏ	3E	312	1	12
1981	86			-10:90	22	1E	ŏ	IE	32	-	1
1981	87	12241:12255			32	0	ŏ	22	3E	1	12
1981	88	12255:12269		-10:90		ÖE	ŏ	32	3E	iE	l
1981	90	12282:12296	229.506	5:110	32	0	ŏ	22		OE	i
1981	91		229.507	5:110	3E	ŐE.	ŏ	3E	æ	1	ÖE -
1981	92	12309:12323	229.508	5:110	32	õ	ŏ	3E	32	ie	1
1981	94	12337:12351	229.512	20:125	22	1E	ŏ	IE	312	IE	i
1981	95	12351:12365	229.513	20:125	ZE	0	ō	22	3E	1	IE
1981	96	12365:12378	229.514	5:110	32	Œ	Õ	32	3E	iz	i
1981	98	12392:12406		-10:90	32	0	0	12		Œ	1
1981	99	12406:12420	218.501	-10:90	2E	lE	0	Z	3E	1	ÖE
1981	100	12420:12434		-10:90	3E	0	0	32	JE	1	12
1981	102	12448:12462	218.503	-10:90	ZĽ	iE	0	IE	X	Œ	1
1981	103	12462:12475	218.504	-10:90	3E	0	0	2E	Æ	1	12
1981	104	12475:12489	218.505	-10:90	3 <b>E</b>	0E	0	32	3E	12	1
1981	106	12503:12517	229.506	5:110	3E	0	0	2E	32	Œ	1
1981	107	12517:12531	229.507	5:110	3E	Œ	0	3E	3E	1	Œ
1981	108	12531:12545	229.508	5:110	3 <b>E</b>	0	0	3E	32	1E	1
1981	110	12559:12573	229.509	40:145	Œ	1	0	١Z	3E	Œ	1
1981	ш	12573:12587	229.510	40:145	12	1	0	IE	3E	1	0E
1981	112	12587:12600	229.511	40:145	12	1	0	Œ	3E	1E	1
1981	112	12600:12601	229.107			0	0	32	32	l	0E
1981	114	12614:12627	218.500		32	0	0	12	32	Œ	1
1981	115	12627:12641		-10:90	22	12	0	22	32	1	30
1981	116	12641:12655		-10:90	32	0	0	32	32	1	12
1981	118	12669:12682		-10:90	2	12	0	12	32		1
1981	119	12682:12696		-10:90		0	0	22	32	1	12
1981	120	12696:12710			32	OE.	0	32	32	IE	1
1981	122	12724:12738					0	3E	32	IE	1
		12738:12752				QE.				12	
1981		12752:12766				0E				IE	
		12779:12793				0E		38		1E	
1981		12793:12807				30		32	32		1
-		12807:12821				OE .		32	32	1E	1
1981		12835:12848 12848:12862			3E 79		0	1E 2E	泥	012	0E
		12862:12876				ie Oe		2E 3E		1 1E	1
1981 1981		12890:12904				12		12 12	JE JE		l
1981		12904:12918			2E 3E		0	22	3E 3E		12
1981		12918:12931				0E		32	32		1
1701	1.70	12710:12731	210,303	10.70	75		•	16	76	1 6	•

<u>م</u>	ATE	ORBITS	SCAN	RANCE (RM	) 9	SIE	ne s	SETT	TN	3	
TEAR		(APTROX)	PROG	(APPROX)							C3
				(1211000)				82			
1 <b>981</b>	-	12945:12959		5:110	3E		0	22		OE	
1961		12959:12973		5:110		œ	0	3E	Æ	1	œ
1981		12973:12987		5:110		0	0	32	X		1
1981		13001:13015		40:145	OE	1	0	12	32		1
1981 1981	143	13015:13029		40:145 40:145	IE IE	1	0	ie Oe	32 32	1 12	018 1
1981	144	13041:13042	229.107	40:143	1E	0	ŏ	32	32	1	ŬE.
1981	146	13056:13070		-10:90	32	-	ŏ	12	3E	-	1
1981	147		218.501		22	-	ō	Z	32	1	ÖE -
1981	148	13064:13097	218.502	-10:90	32	0	0	3E	3E	1	IE
1981		13111:13125			2E	IE	0	12	3E	œ	1
1981	151	13125:13139			3 <b>E</b>		0	212	3E	t	12
1981	152	13139:13153					0	32	叉	12	1
1981	154	13167:13180		5:110	32		0	2E	æ	OE.	1
1981 1981	155	13180:13194		5:110 5:110	312 312		0	32 32	范范	1 12	012 1
1981	158	13222:13236		20:125	ле 212	IZ	0	IE	笼	1E 1E	1
1981		13236:13250		20:125	22		ŏ	22	JE.	1	iz
1981		13250:13264		5:110	æ	-	ŏ	32	π	ie	i
1981		13278:13291			32		0	12		Œ	1
1 <b>981</b>	163	13291:13305	218.501	-10:90	2E	12	0	25	Æ	1	Œ
1981		13305:13318			3E	0	0	3E	3E	1	12
1981		13333:13346			22	1E		IE	3E	02	1
1981	167	13346:13360			32	-	0	22	æ	1	12
1981 1981		13360:13374		5:110	312 312	012 0	0	38	312	12	1
1981	171	13402:13416		5:110	JE JE	ÔE	0	2E 3E	32 32	1	l OE
1981	172	13416:13430		5:110	32	. –	ŏ	3E	38	iz	1
1981	174	13444:13457		40:145	Œ	ĩ	ŏ	12	_		i
1981	175	13457:13471		40:145	IE	i	0	12	32	1	0E
1981	176	13471:13483	229.511	40:145	IE	1	0	Œ	3E	12	1
1981	176	13483:13484	229.107		12	0	0	3E	3 <b>E</b>	1	Œ
1 <b>981</b>	178	13498:13512			3E	0	0	12	32	œ	1
1981		13512:13526			2 <b>E</b>	12		22	32	1	Œ
1981	180	13526:13540			32	0	0	32	32	1	12
1981		13554:13567				IE		12		OR	1
1981		13567:13581 13581:13595			32		0	22	TE TE	1	12
1981 1981	184 186	13609:13622		5:110	312	05	0	JE ZE	3E 3E	1E OE	1
1981	187	13622:13636		5:110		ÖE	ŏ	3E	32	1	ŐE.
1981	188	13636:13650		5:110	32	0	ŏ	32	3E	iE	ĩ
1981	190	13664:13678		20:125	22	IE		12	3E	12	i
1981	191	13678:13692		20:125	2E	0	0	22	3E	1	12
1981	192	13692:13706	229.514	5:110	3E	œ	0	3E	3E	12	t
		13720:13733			3E	0	0		3Z		
		13733:13747				12	0		3E		
		13747:13761			32		0		32		12
		13775:13788				12		12	32		1
1981		13788:13802			32		0	22	3E		12
		13802:13816				OE		3E		12	
		13830:13843		5:110	3E	0 De	0	22 312	3E 3E		l OE
		13843:13857 13857:13871		5:110 5:110	JE JE		0		3E		
1981		13885:13899		40:145	3C 02		0	12		02	1
		13899:13913		40:145	IE		ŏ	IE	æ		OE
1981		13913:13925			IE		ŏ		3E		1
						-	-				-

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	ATE	(1) 1 TTT	SCAN	RANGE ( KH	) STE	re se	TTINGS
	DAT	ORBITS (APPROX)	PROG	(APPROX)	A1 A2	A3 8	1 C1 C2 C3
							2
		13925:13926			IE O	-	BE 3E 1 OE
		13940:13954			32 0		E 3E OE 1
	211	13954:13968			272 112 372 0		12 312 1 012 12 312 1 12
	212	13996:13702			2E 1E		E SE OE 1
	215	14009:14023			32 0		232112
	216	14023:14037			312 012		E JE IE I
	218	14051:14064		5:110	3E 0		E SE OE I
1981	219	14064:14078	229.507	5:110	3E 0E	0 3	E 3E 1 OE
		14078:14092		5:110	3E 0		e je ie i
		14105:14120			ZE IE		E 3E 1E 1
				20:125	222.0		E 3E 1 1E
		14134:14148		5:110	32 02		E 3E IE I
1981		14162:14175 14175:14189			32 0 22 12	-	E 31E 01E 1 1E 31E 1 01E
1981		14189:14203			32 0		E 3E 1 1E
		14217:14231			22 12		E 3E OE 1
1981		14231:14245			32 0		2 32 1 12
	232	14245:14258			JE OE		E JE IE I
1981	234	14272:14286	229.506	5:110	37E 0	0 Z	E 3E OE 1
1981	235	14286:14299	229.507	5:110	JE OE	0 3	E 3E 1 OE
		14299:14313		5:110	3E 0	0 3	e 3e 1e 1
		14327:14341			072 1		E 3E 0E 1
1981		14341:14355		40:145	1E 1	-	E 3E 1 OE
1981		14355:14368		40:145	12 1		E 3E 1E 1
		14368:14369 14383:14396		40:145	12 1 32 0	-	12 312 112 1 E 312 012 1
1981 1981		14396:14410			37E 0 27E 1		E 3E 0E 1 1E 3E 1 0E
1981		14410:14424			320		E 3E 1 1E
1981		14438:14452			ZE IE		E 372 072 1
1981	247	14452:14465			32 0		Z 3Z 1 1Z
		14465:14479			SE OE		E SE IE I
1981	-	14493:14507		5:110	3E 0	-	2 32 02 1
1981	251	14507:14521	229.507	5:110	3E OE	0 3	E 3E 1 OE
1981	252	14521:14535	229.508	5:110	3E 0	0 3	E 3E 1E 1
1981	254	14549:14562	229.512	20:125	2E 1E	0 1	E 37E 1 E 1
19 <b>81</b>		14562:14576		20:125	2E 0	0 2	E 3E 1 1E
1981		14576:14590		5:110	3E 0E		E 3E 1E 1
		14604:14618			320		
		14618:14632			ZE IE		EJEIOE
1981 1981		14632:14645 14659:14673			32 0 22 12		E 3E 1 1E E 3E 0E 1
-		14673:14687			320		E 3E 1 1E
1981		14687:14701			32 02	_	
		14715:14728			3E 0	-	E SE OF 1
		14728:14742		5:110	SE OE		EBEIOE
		14742:14756		5:110	3E 0		E 3E 1E 1
		14770:14783		40:145	OE I		E JE OE I
1981	271	14783:14797	229.510	40:145	1E 1	0 1	E 3E 1 0E
1981		14797:14810		40:145	1E 1		e 3e 1e 1
		14810:14811			IE O		E 3E 1 OE
1981		14825:14839			JE O		E 3E OE I
		14839:14853			2E 1E		E 3E 1 OE
1981		14853:14867			3E 0		
		14880:14894					E 3E 1E 1
1981	2/9	14894:14908	210.000	-10:90	3E 0E	0 3	EGELEL

DA 77		664 <b>M</b>	RANCE (RH	•	IE	æ s	se n		3	
DATE Year Da		SCAN PROG	(APPROX)							C3
	A (APPRUA)	PROV	(ne : parts)				82			
						_		_		_
1981 28	0 14908:14922	218,505	-10:90		Œ		32	30		1
	2 14936:14950				02	-	3E	32	IE.	1
	3 14950:14963				OE .		32	32	12	1
	4 14963:14977				OE OE	0	32 32	JE Je	1E 1E	1 1
1981 28 1981 28					0E		3E		12	1
	8 15018:15032					ŏ	32	32	12	i
	0 15046:15060				OE		3Z		12	i
1981 29				3E	OE.	0	3E	32	12	1
1981 29	2 15074:15088			3E	Œ	0	3E	3E	12	1
1981 294	4 15101:15115	218.505	+10:90	32	œ	0	32	3E	lE	1
1981 29					œ		32		IE	1
1981 29					OE.		32	32	12	1
1981 299			-10:90	3E 3E	OE. OE.	0	312 312		IE IE	t 1
-	9 15171:15185 0 15185:15198					0	3E	3E 3E	12	1
1981 302		-	40:145	OE	1	ŏ	IE		OE.	i
1981 30		329.510	40:145	IE	i	ŏ	iE		1	ÔE.
1981 304		329.511	40:145	12	1	Õ	OE.			1
1981 30	15253:15254	329.107		١Ľ	0	0	3E	3E	1	0E
1981 306		318.500		3E	0	0	IE		œ	L
	7 15281:15295			2E		0	2E	3E	1	OE
1981 308		318.502		3E	0	0	32	3E	1	12
1981 310 1981 311	) 15323:15337 1 15337:15350			212 312	1E 0	0	18 28	3e 3e	072 1	1 12
1981 312			-		-	0	2E 3E	3E	1 12	1
1981 314		329.506	5:110		0	ŏ	22		Œ	i
1981 315			5:110			õ	32	32	1	ÔE -
1981 316	5 15406:15420	329.508	5:110	3E	0	0	3E	3E	12	1
1981 318	3 15434:15448	329.512	20:125	2E	12	0	12	3E	12	1
1981 319		329.513	20:125	22	0	0	22	3E	1	12
	0 15461:15475		5:110	312		0	3E	3E	12	1
1981 322		318.500		32	0	0	IE		œ	1
1981 323				22	12		2E	JE	1	OE
1981 324 1981 326			-10:90 -10:90	3E ZE	0 12	0	3E	3E 3E	1 072	1 <b>E</b> 1
1981 327				2E 3E		0	1E 2E	32	1	1
1981 328				32		ŏ	3E	32	ie	1
1981 330			5:110	3E		ŏ	2E		Œ	i
1981 331	15613:15626	329.507	5:110	32	0E	0	3E	3E	1	0E
1981 332	2 15626:15640	329.508	5:110	3E	0	0	3E	3F.	1E	1
1981 334	15654:15668	329.509	40:145	œ	1	0	1E	3E	Œ	1
1981 335			40:145	12	1	0	IE	3E	1	OE
1981 336			40:145	lE	1	0	OE	3E	12	1
	5 15695:15696				0			32		OF.
	15709:15723			3E		0	1E	3E		1
	15723:15737 15737:15751			2E 3E	IE 0	0	2E		1 1	OE
	15765:15779				IE.	-	3E 1E	3E 3E		1 <b>E</b> 1
	15779:15793			32		0	212		1	1 1E
	15793:15806				ñe.		38		ie	
	15820:15834		5:110	3E		ŏ	2E	JE		i
	15834:15848		5:110		07	-	32		1	ÔE.
	15848:15861		5:110	3E		0	3E	32		1
	15875:15889		20:125		12	0	12	3E	18	1
1981 351	15889:15903	329.513	20:125	2E	0	0	2E	3 <b>E</b>	1	12

			SCAN	RANGE ( RH	) SIE	VIE S		TNES	
YZAR		ORBITS	PROG	(APPROX)					3
12.44	UNI	(APPROX)					B2		
1961	352	15903:15917	329.514	5:110	3E 0E	0	3E	3E 1E 1	
	354				3E 0	0	12	372 OF 1	
1961		15944:15958			2E 1E	0	ŻĒ	3E I 0E	!
1981	356	15958:15972	318.502	-10:90	3E 0	0	3 <b>E</b>	3E 1 1F	
1981	358	15986:16000	318.503	-10:90	22 IE	0	lE	372, OFE 1	
1981	359	16000:16013	318.504	-10:90	3E 0	0	2E	32 1 1E	:
1981	360	16013:16027	318.505	-10:90	JE OE	0	3E	32 IE I	
1 <b>981</b>	362	16041:16055	329.506	5:110	3E 0	0	2E	3E OE 1	
1981	363	16055:16069		5:110	3e 0e		3E	3E 1 OE	2
1981	364	16069:16083		5:110	3E 0		3E	3E 1E 1	
1982		16097:16111		40:145	OEI		IE	312 OF 1	
	2	16111:16125		40:145	12 1		12	3E 1 0E	
1982		16125:16137		40:145	1E 1		œ	3E 1E 1	
	3	16137:16138		10.00	IE O		32	3E 1 OE	
1982	-	16151:16166			32 0		12	32 02 1	
1982		16166:16180			ZE IE		ZE.	3E 1 0E 3E 1 1E	
1982 1982		16180:16194 16207:16221			32 0 22 12		3E 1E	312 1 112 312 012 1	
	10	16221:16235			32 0		21	312 1 12	,
	11	16235:16248					32	JE IE I	
1982		16262:16276		5:110	32 0	-	2E	32 02 1	
1982		16276:16290		5:110	JE OE		32	37E 1 07E	
1982	-	16290:16303		5:110	320		32	SE LE L	•
1982		16303:16304	-	5.110	IEO		3E	3E 1 0E	,
1982		16318:16332		20:125	22 12		12	JE IE I	•
1982	-	16332:16345		20:125	220		22	3E 1 1E	•
1982		16345:16359		5:110	SE OE	-	32	3E 1E 1	•
1982		16373:16387			32 0	-	12	SE OE 1	
1982		16387:16400			28 18		22	32 1 02	2
1982		16400:16414			SE O	-	3E	3E 1 1E	
1982		16428:16442			2E 1E		IE.	3E OE 1	
1982		16442:16456			JE O		2E	3E 1 1E	
1982	-	16456:16470	-				32	3E 1E 1	
1982	29	16483:16497		5:110	BE O	0	2E	BE OF 1	
1982	30	16497:16511	329.507	5:110	3E OE	0	3E	3E 1 OE	!
1982	31	16511:16525	329.508	5:110	JE O	0	3E	3E 1E 1	
1982	33	16539:16553	329.509	40:145	01E 1	0	12	3E OE 1	
1982	34	16553:16567	329.510	40:145	1E 1	0	ie	3E 1 OE	:
19 <b>82</b>	35	16567:16580	329.511	40:145	1E 1	0	œ	JE IE I	
1982	37	16594:16608	318.500	-10:90	3 <b>E</b> O	0	1 E	3e oe 1	
	38	16608:16622			2E 1E	0	2 <b>E</b>	312 I 012	:
	39	16622:16636			3E 0	0	3E	3E 1 1E	:
1982		16649:16663					IE	JE OE I	
1982		16663:16677			3E 0		2E	3E 1 1E	
1982		16677:16690			3e oe			3E 1E 1	
1982		16705:16719		5:110	3E 0			JE OE I	
1982		16719:16733		5:110			32	3E 1 OE	
1982		16733:16746		5:110	3E 0		3E	SE LE 1	
1982		16760:16774		20:125	2E 1E		IE	3E 1E 1	
1982		16774:16788		20:125	2E 0		2E	32 1 12	•
1982		16788:16801		5:110	3E OE		3E	3E 1E 1	
1982		16816:16829			3E 0		IE Te	3E OE 1	,
1982	-	16829:16843			2E 1E		2E	JE 1 OE	
1982		16843:16857			32 0		3E	3E 1 1E	•
1982		16870:16885			ZE IE		12	BEOEL	,
1982	79	16885:16898	210,204	-10:50	JE O	0	2E	3E 1 1E	•

			RANGE (RM)	۱ c	1		SETI			
DATE	ORBITS	SCAN	(APPROX)							G
YEAR DAY	(APPROX)	PROG	(APPROX)	~		~	B2	<b>.</b> .		
							~			
1982 59	16898:16912	318.505	-10:90	æ	0E	0	3E	æ	12	1
1982 61	16926:16940		5:110	3E (	0	0	2E	æ	Œ	1
1982 62	16940:16954	329.507	5:110	3E (	œ	0	3E	X	1	Œ
1982 63	16954:16968	329.508	5:110	3E (	0	0	3E	3E	IE	1
1982 65	16981:16995	329.509	40:145	œ	1	0	12	Ľ	œ	1
1982 66	16995:17009	329.510	40:145			0		ĴĒ	1	0E
1982 67	17009:17023	329.511	40:145		-	0	Œ	æ	12	t
19 <b>82 69</b>	17036:17050			3 <b>e</b> (		0	12	3E		1
1982 70	17050:17064			32	-	0	12			1
1982 71	17064:17078			32 (		0		3E		t
1982 73		318.503				0	IE	38		1
1982 74	17106:17119			32 (		0	22		1	12
1982 75	17119:17135					0	32	32	12	1
1982 77 1982 78	17147:17161	329.506	5:110 5:110	3E ( 3E (	-	0	2E JE	3E 3E		1
1982 78 1982 79	17174:17188		5:110	312 (		õ	JE JE	JE T	1 12	02
1982 81	17202:17216		20:125		IE.	-	12	光		1
1982 82	17216:17230		20:125	22 (		ŏ	2E	æ	1	iz
1982 83	17230:17244		5:110	3E (	-	õ	32	3E	iE	i
1982 85	17257:17271			3E (	_	ŏ			OE.	i
1982 86	17271:17285				12	-	22	-	ī	ÔE -
1982 87	17285:17298			3E (	0	0	3E		i	18
1982 89	17312:17326	318.503	-10:90	22.	lE	0	12	32	Œ	1
1982 90	17326:17340	318.504	-10:90	3E (	0	0	2 <b>2</b>	X	1	12
1 <b>982 91</b>	17340:17354	318.505	-10:90	3E (	Œ	0	3E	3E	15	1
19 <b>82 93</b>	17368:17382		5:110	3E (	0	0	2E	32	0E	1
1982 94	17382:17396	329.507	5:110	3E	Œ	0	3E	3E	1	Œ
1982 95	17396:17410		5:110		0	0	3E	3E		1
1982 97	17424:17438		40:145		1	0	12		œ	-
1982 98	-	329.510	40:145		ł	0	12	32	1	œ
1982 99	17451:17465		40:145		1	0	œ	38	IE	1
1982 101	17479:17493			32 (		0	12	32		1
1982 102	17493:17506			_	IE.	-	2E	32	1	OE
1982 103	17506:17519				0	0	32	32	1	12
1982 105					12		1E		0E	1
	17548:17562			32		0	22		1	1E
1982 107		318.505				0	32	32		1
1982 109 1982 110	17589:17603	329.507	5:110 5:110	3E ( 3E (		0	272 372	范辺	1	1 012
1982 111	17617:17629		5:110	32 (		0	3E			1
1982 111	17630:17631	329.107	3.110		0	0	3E	312 312	1	0E
1982 113	17645:17658		20:125		IE	-	1E	32	12	1
1982 114	17658:17672		20:125	22		õ	22	3E	1	iz
	17672:17686		5:110	32 (			32	JE		1
	17700:17714			3E (	-	-	IE	_	_	
	17714:17728			28			22			ÔE.
	17728:17742			32 (		õ		32	i	12
	17756:17770			2E			12	32		
	17770:17783			38		Ō		32		12
	17783:17797			32			32		12	
	17811:17825		5:110	3E (		0		3E		
	17825:17838		5:110	3E (	Œ	0	3E	Æ	1	Œ
1982 127	17838:17852	329.508	5:110	3E (	0	0	3E	3E	IE	1
	17866:17879		40:145	œ	1	0	12	3E	Œ	1
	17879:17893		40:145	12	1	0		3E		OE.
	17893:17907		40:145	12	1	0	œ	3E	IE	1

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		SCAN	RANGE ( KH	) 9	SIE	ne :	SET.	TIN	3	
DATE YEAR DAT	ORBITS (APPROX)	PROG	(APTROX)							C3
		-					B2			
					_	-			_	• •
1982 133	17921:17935	318.500	-10:90		0				0E	
1962 134	17935:17949	318.301	-10:90		12		2世 3世	32	1	012 112
1962 135	17949:17963 17977:17990	318.302	-10:90	3E	U IE	0		JE JE		1
1982 137	17990:18004	318 504	-10:90	24 312		0	22		1	12
1982 138	18004:18018	318.505	-10:50		0E		32		iz	
	18032:18046			32	-	ŏ	ZE		OE.	
	18046:18060		5:110		õe		32	Æ		Œ
	18060:18073			32		0	3E	3E	IE	1
1982 145	18087:18100	329.512	20:125	2E	12	0	IE	Æ	1E	1
	18100:18114		20:125	Z	0	0	22		1	IE
	18114:18126		5:110		œ		3E		IE	
	18126:18129		5:110		œ				IE.	
	18143:18156			32		0			œ.	
	18156:18170 18170:18184				12			3E 3E		012 112
	18198:18211			3E 2E	IE	0			0E	
	18211:18225			32		ŏ		JE		iE
	18225:18239				ÖE				iE	
	18253:18266		5:110	32		ŏ			Œ	
	18266:18280			32		ō	-		1	ÖE
1982 159	18280:18294	329.508	5:110	3E	0	0	32	3E	18	1
	18308:18322			3E	0	0	3E	3E	12	1
	18322:18336			3E	0	0	3E	3E	12	1
	18336:18350			3E		0			12	-
	18363:18377			32		0	22	32	1	12
	18377:18391			3E		0	22	32	1	12
	18391:18404 18418:18432			3E		0	2E	32	1	1E
	18432:18446			3E 3E		0	范	JE JE	ie Ie	1
	18446:18460			3E		ŏ		_		i
	18474:18487			22	12		32	3E	1	ie -
	18487:18501			22	IE				i	IE
	18501:18515			2E	12		32	32	1	12
1982 177	18529:18543	333.631		3E	0	0	3E	38	1E	1
1982 178	18543:18557	333.631		3E	0	0	3 <b>E</b>	3E	12	1
1982 179	18557:18569	333.631		3 <b>E</b>	0	0	3E	3E	IE	1
	18569:18571			3E		0	3E	3E	12	1
	18585:18596			32		0	22	32	1	12
	18598:18612			32		0	2E	32	1	IE
	18612:18626			32		0	22	3E.	1	12
	18640:18654 18654:18668			3E 3E		0 0	3E 3E	JE JE	02. (77)	1
	18668:18681			эе. Зе		0		JE JE		1
	18695:18709			3E		ŏ		3E		0E
	18709:18723			3E		ŏ	IE		i	OE
	18723:18737			3E		ŏ	IE	3E		OE.
	18751:18765			3E		Ō	3 <b>E</b>			1
1982 194	18765:18778	333.631		3E		0	3 <b>E</b>	3E	12	1
	18778:18792			3E	0	0	3E	3 <b>E</b>	12	1
	18806:18820			3E	0	0	2E		1	12
	18820:18834			3E		0	2E	3E	1	IE
	18834:18847			3E		0	2E	3 <b>E</b>	1	12
	18861:18875			3E		0	32	3E	1E	1
	18875:18889			3E		0	32	32		1
1962 203	18889:18902	179.07		3 <b>E</b>	U	0	3 <b>E</b>	3 <b>E</b>	IE	1

D	TZ	ORBITS	SCAN	RANGE (ICH)							
YEAR	DAY	(APPROX)	PROG	(APPROX)	A1	A2	<b>A3</b>	81 82	CI	C2	ය
1962	205	18916:18930	333.632		38	0	0	22	X	1	12
1982	206	18930:18944	333.632		3E	-	0	22	32	1	18
		18944:18957			32	0	0	28	38	1	12
1982	208	18957:18971	333.632		3E	-	0	22	ЗE	1	12
1982	209	18971:18985	333.631		32	-	0	32	Z	IZ	1
1982	210	18985:18998	333.631		3E	0	0	3E	3E	12	1
1982	211	18998:19012	333.631		3E	-	0	38	3E	12	1
		19012:19013			3E	-	0	3 <b>E</b>	X	12	1
		19013:19027			3 <b>E</b>	-	0	3E	3E	12	1
		19027:19042			3E		0	Z		1	12
		19042:19055			3E	-	0			1	12
		19055:19068			3E		0	22		1	12
		19068:19082			3 <b>e</b>	-	0	ZE		1	12
		19082:19096			3E	-	0		3E		1
		19096:19110			32	-	0	3Z	Æ		1
		19110:19124			3E	-	0	3 <b>E</b>	æ		1
		19124:19138			3E	-	0	T	Æ		1
		19138:19151			3E	-	0	12	72	-	0E
		19151:19165			3E	-	0	12	Æ	1	œ
		19165:19179			3E		0	IE	_	1	OE
		19179:19193			3E	-	0	12	32	-	0E
		19193:19207			3E	-	0	32		12	-
		19207:19221			3E	-	0	Æ	3E		1
		19221:19234			32	-	0	3E	32		-
		19234:19248			3E	-	0	æ		18	-
		19248:19262			3E	-	0	22	32	-	18
		19262:19276			3E	-	0	Z	JE	-	12
		19276:19289			3E	-	0	Ž	38	-	12
		19303:19317			32	-	0	32		12	-
		19318:19331			32	-	0	32	æ		1
•••==		19331:19344			3E	-	0	38	X		1
		19358:19372					0	3E		1	12
		19372:19386			ZZ	lE	-	3E	38	1	12
1982	2 <b>39</b>	19386:19400	333.633		2E	12	0	3E	3E	1	IE

NOTE 1: Wrong scan pattern (GSFC error). Space views suspect.

NOTE 2: Bad roll compensation phasing. Some data lost.

NOTE 3: Possible beat (PMC B1/2 and chopper).

NOTE 4: Spikes on data in some orbits.

NOTE 5: C1 signal channel spikes evident on some orbits.

NOTE 6: Large number C1 signal channel spikes. Also some on A23W.

#### Table 5-3

### Mode Table 2: Day 286, 1981 to 160, 1983

This table also contains the start and end times of daily data sets (year, day of year, second of day) and the number of major frames of data in the data set, as well as the sieve setting/-PMC mode as before.

NOMINAL	START TIME	END TIME 16 SEC	PMC MODE	PROCESSING
DATE	YR DAY SECS	YR DAY SECS FRAMES	A1 A2 A3 B1 C1 C2 C3	DATE
YEAR DAY			82	YEAR DAY
1981 286	1981 286 1668	1981 286 86388 5259	3E OE O 3E 3E 1E 1	1983 146
1981 287	1981 287 4	1981 287 86388 5181	3E OE O 3E 3E 1E 1	1983 146
1981 288	1981 288 4	1981 289 3412 5474	3E OE O 3E 3E 1E 1	1983 146
1981 290	1981 289 86276	1981 290 86388 5357	3E OE O 3E 3E 1E 1	1983 146
1981 291	1981 291 4	1981 291 86388 4991	3E OE O 3E 3E 1E 1	1983 146
1981 292 1981 294	1981 29 <b>2</b> 4 1981 294 4356	1981 293 1684 5417 1981 294 86388 4906	3E OE O 3E 3E 1E 1 3E OE O 3E 3E 1E 1	1983 146 1983 171
1981 295	1981 294 4356 1981 295 4	1981 295 86388 5338	3E OE O 3E 3E 1E 1	1983 171
1981 296	1981 295 4	1981 296 86356 5358	3E OE O 3E 3E 1E 1	1983 171
1981 298	1981 298 2564	1981 298 86388 5210	3E OE O 3E 3E 1E 1	1983 171
1981 299	1981 299 4	1981 299 86388 5204	3E OE O 3E 3E 1E 1	1983 171
1981 300	1981 300 4	1981 301 4532 5064	3E OE O 3E 3E 1E 1	1983 171
1981 302	1981 302 788	1981 303 388 5326	OE 1 0 1E 3E 0E 1	1983 172
1981 303	1981 303 1908	1981 304 1508 5354	IEI O IE JEI OE	1983 172
1981 304	1981 304 3028	1981 304 81508 4881	1E 1 0 0E 3E 1E 1	1983 172
1981 306	1981 306 5252	1981 306 85140 4963	3E 0 0 1E 3E 0E 1	1983 172
1981 307	1981 307 116	1981 308 596 5274		1983 172
1981 308	1981 308 1236	1981 309 900 529 <b>3</b>		1983 172
1981 310	1981 310 3460	1981 311 2980 5326	2E 1E 0 1E 3E 0E 1	1983 216
1981 311	1981 311 4580	1981 311 84500 4956		1983 216
1981 312	1981 311 85844	1981 312 85572 5301	3E OE O 3E 3E 1E 1	1983 216
1981 314	1981 314 1684	1981 314 37220 2207	3E 0 0 2E 3E 0E 1	1983 216
1981 314	1981 314 43460	1981 315 1220 2744 1981 316 2452 5299	3E 0 0 2E 3E 0E 1	1983 216
1981 315 1981 316	1981 315 2804 1981 316 3908	1981 316 2452 5299 1981 317 3476 5327		2 1983 216 1983 216
1981 318	1981 317 86292	1981 318 85988 5323	3E 0 0 3E 3E 1E 1 2E 1E 0 1E 3E 1E 1	1983 227
1981 319	1981 319 1012	1981 320 676 5138		1983 227
1981 320	1981 320 2132	1981 321 1716 5181	3E OE O 3E 3E 1E 1	1983 227
1981 322	1981 322 4356	1981 323 4212 5131	3E 0 0 1E 3E 0E 1	1983 227
1981 323	1981 323 5476	1981 323 85220 4958		1983 227
1981 324	1981 324 340	1981 324 86260 5309		1983 227
1981 326	1981 326 2596	1981 327 2260 5351	2E 1E 0 1E 3E 0E 1	1983 228
1981 327	1981 327 3684	1981 328 3284 5326		1983 228
1981 328	1981 328 4804	1981 328 84580 4967	3E OE O 3E 3E 1E I	1983 228
1981 330	1981 330 788	1981 331 500 5258	3E 0 0 2E 3E 0E 1	1983 228
1981 331	1981 331 1924	1981 332 1444 5293	3E OE O 3E 3E 1 OE	1983 228
1981 332	1981 332 3012	1981 333 2564 5271	3E 0 0 3E 3E 1E 1	1983 228
1981 334	1981 334 5380	1981 334 85140 4852	OE 1 0 1E 3E 0E 1	1983 231
1981 335	1981 335 116	1981 335 86340 5224		2 1983 231
1981 336	1981 336 1236	1981 336 79796 4805	IE 1 0 DE 3E IE 1	1983 231
1981 338	1981 338 3460	1981 339 3252 5113	3E 0 0 1E 3E 0E 1	1983 231
1981 339 1981 340	1981 339 6420 1981 339 85844	1981 339 84596 4824 1981 340 85524 4861		E 1983 231
1981 340	1981 342 1668	1981 343 1204 5281	3E 0 0 3E 3E 1 11 2E 1E 0 1E 3E 0E 1	E 1983 231 1983 234
		1981 344 2356 5289		
1981 344	1981 344 3908	1981 345 3380 5318	3E OF 0 3E 3E 1E 1	
1981 346	1981 346 6132	1981 346 85988 4907	3E 0 0 2E 3E 0E 1	1983 234
1981 347	1981 347 996	1981 348 724 5294		1983 234
1981 348	1981 348 2132	1981 349 1636 5279	3E 0 0 3E 3E 1E 1	1983 234
1981 350	1981 350 4356	1981 351 3892 5185	2E IE 0 IE 3E IE 1	1983 235
1981 351	1981 351 5460	1981 352 5220 5329		1983 235
1981 352	1981 352 6580	1981 353 20 4931	3E OE 0 3E 3E 1E 1	1983 235
1981 354	1981 354 2564	1981 355 2164 5286	3E 0 0 1E 3E 0E 1	1983 235
1981 355	1981 355 3668	1981 356 3316 5298		1983 235
1981 356	1981 356 4788	1981 357 4324 5295		1983 235
1981 358	1981 358 772	1981 359 1220 5380	2E 1E 0 1E 3E 0E 1	1983 238

NOMI	NAL	ST	ART 1	TIME	E!	T DR	IME	16 SEC				с мо				PROCE	SSING
DAT	Z	TR	DAT	SECS	YR	DAT	SECS	FRAMES	A1	A2	A3	B1	Cl	C2	C3	DA1	TE
YEAR	DAT											82				TEAR	DAT
1981		1981			1981			5271	3E		0		38			1983	
1981		1981	-	2996	1981		-	5182		OE				12		1983	
1981			362	5236			84964	_	3E		0			0E		1983	
1981		1981		100			86036			OE			3E			1983	
1981		1981		1220	1981			5255	32		0			IE		1983	
1982	1	1982	-	11771	1982	2	-	4706	OE		0			0E		1983	
1982	2	1982	2	4571	1982		84331		12		0		3E			1983	
1982	2	1982		21979	1982		84331		12		0		32			1983	
1982	3	1982	3	3499	1982		85563		32		0			OE	-	1983	
1982	5	1982	5	7899 2763	1982	6		4898	38	12	0			OE		1983	
1982	6	1982	6		1982	7		5297					38			1983	
1982	7	1982	7	3883 86251	1982		83803		32		0		32	0E		1983	
1982	.9	1982			1982		85851			12						1983	
1982	10	1982	10	971 2091	1982	11		5198	32		0		32	12		1983	
1982 1982	11 13	1982 1982	11	4331	1982 1982	12		5317		0E	0			0E		1983	
1982	14	1982		85611	1982		84219 85355		32	OE			32			1983 1983	
1982	15	1982	15	299	1982		78827		38 38		ŏ			1		1983	
1982	17	1982	17	2539	1982	18		5142		12				IE		1983	
1982	18	1982	18	3643	1982	19		5192	22		ŏ		32			1983	
1982	19	1982	19	4763	1982		84635			0E				ie		1983	
1982	21	1982	21	747	1982	22		5296	32		ŏ			OE		1983	
1982	22	1982	22	1851	1982	23		5328		12		22	32			1983	
1982	23	1982	23	2971	1982	24		5307	35		ŏ		32			1983	
1982	25	1982	25	5195	1982	26		5314		12				ÖE		1983	
1982	26	1982	26	6315	1982		86251		38		Ō		32			1983	
1982	27	1982	27	1179	1982	28	715	5287		0E		32	3E	12	1	1983	308
1982	29	1982	29	3403	1982	30		5279	3E		0			OE		1983	
1982	30	1982	30	4523	1982	30	84411	4880	3E	OE	0	32	3E	1	0E	1983	308
1982	31	1982	30	85787	1982	31	85563	5308	3E	0	0	3E	3E	IE	1	1983	308
1982	33	1982	33	1611	1982	34	1227	5275	OE	1	0	12	3 <b>E</b>	0E	1	1983	311
1982	34	1982	34	2731	1982	35	2363	5240	12	1	0	ΙĒ	3E	1	0E	1983	311
1982	35	1982	35	3835	1982	- 35	83643	4835	lE	1	0	0E	3E	12	1	1983	311
1982	37	1982	36	86219	1982	37	85787	5307	3E	0	0	12	3E	0E	1	1983	311
1982	38	1982	- 38	939	1982	- 39	475	5303	2E	12	0	2E	3E	1		1983	
1982	39	1982	39	2107	1982	40	1627	5249	3E	0	0	3E	3E	1	lΕ	1983	311
1982	41	1982	41	4283	1982	42	3707	5286	2E	12	0	18	3E	0E	l	1983	312
1982	42	1982	42	540 <b>3</b>	1982	43		5241	3e		0		3E			1983	
1982	43	1982	43	6571	1982	43	86171		_	0E	0			18		1983	312
1982	45	1982	45	2491	1982	46		5529	38		0			OE		1983	
1982	46	1982	46	9851	1982	47		4907		0E		-	3E	-		1983	
1982	47	1982	47	4715	1982		83355		3E		0	-		12		1983	
1982	48	1982		23483	1982	48	23531	4		0E		_		OE		1983	
1982	49	1982	49	827	1982	50		5294		1E				12		1983	
1982		1982		1803			1435									1983	
1982	51	1982	51	2923	1982	52		5154		0E				IE		1983	
1982	53	1982	53	5147	1982		84923			0				0E		1983	
1982	54	1982	54	11	1982	-	86107			IE				1		1983	
1982	55	1982	55	1131	1982	56		5194		0			3E			1983	
1982	57	1982 1982	57	<b>3355</b> 10715	1982	58	84187	5425	2 K 3 E	IE				OE		1983	
1982	58 50		-		1982						0			1		1983	
1982	59	1982		85739	1982		85451			0E				1E		1983	
1982	61	1982	61	1547	1982			5540			0			0E		1983	
1982 1982	62 63	1982	62	8907	1982	63 64		492 <b>4</b> 5257		0E 0						1983	
1982	63 65	1982	63		1982	66		5734		1				IE OE			315
1997	00	1982	04	86155	1982	00	5//1	31.34	UE,	1	U	16	שנ	UE	ı	1983	214

1000						D TI	INTE	IS SEC			PM(	СМО	DE			PROCE	SSING
NOMI DAT		ST/ YR	ART " DAY	SECS	YR		SECS I		<b>A</b> 1	A2		B1		C2	с3	DAT	
YEAR		1 R	URI	3203	14	541						B2	``		••	YEAR	
	-																
1982	66	1982	66	7147	1982	67	491	4946	12	1	0		3E			1983	319
1982	67	1982	67	1979	1982	68		5376	12		0	OE			-	1983	
1982	69	1982	69	2811	1982		86395		3E	-	0		3E	-	-	1983	-
1982	70	1982	70	11	1982		86395		3E		0		3E			1983	
1982	71	1982	71	11	1982		86219		3E		0	3E	3E	1		1983 1983	
1982 1982	73 74	1982 1982	73 74	2427 3531	1982	74	83595	5334	2E 3E	IE	ŏ		32			1983	
1982	75	1982		84795	1982	76	4107			ÖE			3E			1983	
1982	77	1982	77	619	1982	78		5015	32		ō		32			1983	-
1982	78	1982	78	1739	1982	79		5305			0	3E	3E	1	0E	1983	321
1982	79	1982	79	2843	1982	79	81723	4876	3E	0	0	3E	3 <b>E</b>	12	1	1983	321
1982	81	1 <b>982</b>	80	85227	1982	81	85083	5314		12	0		3E			1983	
1982	82	1982		86347	1982		84555		2E		0		3 <b>E</b>			1983	
1982	83	1982	83	1035	1982	84		5531		OE			38			1983	
1982	85	1982	85	3275	1982	86		5201	38		0		32			1983	
1982	86	1982	86	4379	1982		84251		2E 3E	12	0		3E 3E			1983 1983	
1982 1982	87 89	1982 1982	89	85643 1483	1982 1982	88 90		5711 5308		IE			3E			1983	
1982	90	1982	90	2571	1982	91		5313	32		ŏ		3E			1983	
1982	91	1982	91	3691	1982	92		5307		ÖE	-		38		1	1983	-
1982	93	1982	93	5915	1982	94		4860	3E		Õ		3E		1	1983	175
1982	94	1982	94	779	1982	95	811	5354	3E	OE	0	3 <b>E</b>	3E	1	0E	1983	175
1982	95	1982	95	18 <b>99</b>	1982	96	1563	5323	3E	0	0		3E			1983	175
1982	97	1982	97	4123	1982	98		5317	0e	t	0		3 <b>E</b>			1984	60
1982	98	1982	98	5227	1982		85211	-	IE		0		32		-	1984	60
1982	99	1982	99	6347	1982		86203		1E		0		3E			1984	60
1982		1982		8571	1982			4889	3E.		0		3E			1984	60
1982		1982 1982		3435 4539	1982		2955 84363	5289	2E 3E	IE	0	-	3E 3E			1984 1984	60 60
1982 1982		1982		523	1982			5322		12	-		35			1983	
1982		1982		1627	1982	-		5295	3E		õ		3E			1983	
1982		1982		2747	1982			5316		ŐE			3E			1983	
1982		1982		4971			84843		32		Ō		3E			1983	
1982	110			86235	1982	110	85643	5294	3E	0E	0	3E	3E	i	0E	1983	179
1982	111	1982	111	939	1982	111	79403	4885	3E	0	0	3E	3E	1 E	1	1983	179
1982	113	1982	113	3163	1982			5317		1 E	0		3 <b>e</b>			1983	
1982		1982		4283	1982			5297	2E		0	2F	3e			1983	180
1982		1982		5403			85163			0E			3E			1983	
1982 1982		1982 1982		1371 24 <b>75</b>	1982 1982			5325 546 <b>9</b>	32	0 1E	0		3E 3E		1	1983 1983	180
1982		1982		9835	1982			4953	38		0		JE JE			1983	
1982				85963			85627			1E		IE		0 <b>E</b>		1983	
1982		1982		683	1982			5265	3E		ŏ		3E				
1982		1982		1787	1982	124		5303		0E	0		3E		1	1983	181
1982	125	1982	125	4011	1982	126	3723	5313	3E	0	0	2E	3E	0E	1	1983	181
1982	126	1982	126	5131			85131		3E	0 <b>e</b>		3E	3E	1	0E	1983	
1982				86395			85995			0	0		3 <b>e</b>			1983	
1982		1982		2219	1982			5316	0E		0		3E			1983	
1982		1982		3323	1982			5282	12		0					1983	
1982		1982		4443			84283		15		0		3E			1983	
1982		1982		411			86299	-		0	0		3E			1983	
1982 1982		1982 1982		1515 2651	1982 1982			5304 5308		1E 0			3E 3E			1983	
1982		1982		4875			84651			IE			35	-		1983	
1982				86123	1982			5549		0			3E			1983	
1982		1982		7083	1982			4939		ÖE			3E			1983	
												-	. –		-		

NOMINAL	START TIME	END TIME 16 SEC	PMC MODE PROCESSING
DATE	YR DAY SECS	YR DAY SECS FRAMES	A1 A2 A3 BI C1 C2 C3 DATE
YEAR DAY			BZ YEAR DAY
1982 141	1982 141 3051	1982 142 3099 5350	3E 0 0 2E 3E 0E 1 1983 185
1982 142	1982 142 4155	1982 143 3739 5320	3E OE O 3E 3E 1 OE 1983 185
1982 143	1982 143 5275	1982 143 77483 448 <b>7</b>	3E 0 0 3E 3E 1E 1 1983 185
1982 145	1982 145 1243	1982 146 859 5305	2E 1E 0 1E 3E 1E 1 1983 328
1982 146	1982 146 2363	1982 147 1963 5282	2E 0 0 2E 3E 1 1E 1983 328
1982 147	1982 147 3467	1982 148 3019 5330	3E OE O 3E 3E 1E 1 1983 328
1982 149	1982 148 85851	1982 149 85467 5296	3E 0 0 1E 3E 0E 1 1983 328
1982 150	1982 150 555	1982 151 219 5300 1982 152 1419 5358	2E 1E 0 2E 3E 1 0E 1983 328
1982 151	1982 151 1675 1982 153 3899	1982 152 1419 5358 1982 154 3483 5295	3E 0 0 3E 3E 1 1E 1983 328 2E 1E 0 1E 3E 0E 1 1983 329
1982 153 1982 154	1982 153 3899 1982 154 5003	1982 154 84955 4953	2E 1E 0 1E 3E 0E 1 1983 329 3E 0 0 2E 3E 1 1E 1983 329
1982 155	1982 154 86267	1982 155 86139 5336	3E OE O 3E 3E 1E 1 1983 329
1982 157	1982 157 2091	1982 158 1755 5317	3E 0 0 2E 3E 0E 1 1983 329
1982 158	1982 158 3195	1982 159 2811 5305	3E OE O 3E 3E 1 OE 1983 329
1982 159	1982 159 4315	1982 159 84107 4937	3E 0 0 3E 3E 1E 1 1983 329
1982 161	1982 161 283	1982 161 86395 5328	3E 0 0 3E 3E 1E 1 1983 332
1982 162	1982 162 11	1982 162 86395 5358	3E 0 0 3E 3E 1E 1 1983 332
1982 163	1982 163 11	1982 164 1979 5445	3E 0 0 3E 3E 1E 1 1983 332
1982 165	1982 165 4731	1982 165 86395 5012	3E 0 0 2E 3E 1 1E 1983 332
1982 166	1982 166 11	1982 166 86395 5362	3E 0 0 2E 3E 1 1E 1983 332
1982 167	1982 167 11	1982 168 315 5376	3E 0 0 2E 3E 1 1E 1983 332
1982 169	1982 169 2923	1982 169 86395 5152	3E 0 0 3E 3E 1E 1 1983 334
1982 170	1982 170 11	1982 170 86395 5272	3E 0 0 3E 3E 1E 1 1983 334
1982 171	1982 171 11	1982 172 4587 5594	3E 0 0 3E 3E 1E 1 1983 334
1982 173	1982 173 1115	1982 173 86395 5249	2E 1E 0 3E 3E 1 1E 1983 334
1982 174	1982 174 11	1982 174 86395 5278	2E 1E 0 3E 3E 1 1E 1983 334
1982 175	1982 175 11	1982 176 2859 5487	2E 1E 0 3E 3E 1 1E 1983 334
1982 177 1982 178	1982 177 5563 1982 178 11	1982 177 86395 5009	3E 0 0 3E 3E 1E 1 1983 336
1982 178	1982 179 11	1982 178 86395 5353 1982 180 1067 5389	3E 0 0 3E 3E 1E 1 1983 336 3E 0 0 3E 3E 1E 1 1983 336
1982 181	1982 181 3755	1982 181 86395 5051	3E 0 0 3E 3E 1E 1 1983 336 3E 0 0 2E 3E 1 1E 1983 336
1982 182	1982 182 11	1982 182 86395 5342	3E 0 0 2E 3E 1 1E 1983 336
1982 183	1982 183 11	1982 183 85883 5291	3E 0 0 2E 3E 1 1E 1983 336
1982 185	1982 185 1947	1982 185 86395 5145	3E 0 0 3E 3E 0E 1 1983 339
1982 186	1982 186 11	1982 186 86395 5343	3E 0 0 3E 3E 0E 1 1983 339
1982 187	1982 187 11	1982 188 3627 5560	3E 0 0 3E 3E 0E 1 1983 339
1982 189	1982 189 6395	1982 189 86395 4919	3E 0 0 1E 3E 1 0E 1983 339
1982 190	1982 190 11	1982 190 86395 5327	3E 0 0 1E 3E 1 0E 1983 339
1982 191	1982 191 11	1982 192 1995 5451	3E 0 0 1E 3E 1 0E 1983 339
1982 193	1982 193 4587	1982 193 86395 5051	3E 0 0 3E 3E 1E 1 1983 343
1982 194	1982 194 11	1982 194 86395 5310	3E 0 0 3E 3E 1E 1 1983 343
1982 195	1982 195 11	1 <b>982</b> 196 251 5369	3E 0 0 3E 3E 1E 1 1983 343
1982 197	1982 197 2779	1982 197 86395 5142	3E 0 0 2E 3E 1 1E 1983 343
1982 198	1982 198 11	1982 198 86395 5321	3E 0 0 2E 3E 1 1E 1983 343
			3E 0 0 2E 3E 1 1E 1983 343
1982 201	1982 201 971	1982 201 86395 5269	3E 0 0 3E 3E 1F 1 1983 347
1982 202	1982 202 11	1982 202 86395 5350	3E 0 0 3E 3E 1E 1 1983 347
1982 203	1982 203 11	1982 204 2667 5266	3E 0 0 3E 3E 1E 1 1983 347
1982 205	1982 204 85563	1982 205 86395 5365	3E 0 0 2E 3E 1 1E 1983 347
1982 206	1982 206 11	1982 206 86395 5310	3E 0 0 2E 3E 1 1E 1983 347
1982 207	1982 207 11	1982 207 86395 5322	3E 0 0 2E 3E 1 1E 1983 347
1982 208	1982 208 11	1982 208 86363 5352	3E 0 0 2E 3E 1 1E 1983 347
1982 209	1982 209 3611 1982 210 11	1982 209 86395 -1	3E 0 0 3E 3E 1E 1 1983 61
1982 210 1982 211	1982         210         11           1982         211         11	1982 210 86395 -1 1982 211 85579 -1	<b>3E 0 0 3E 3E 1E 1 1983 61</b> <b>3E 0 0 3E 3E 1E 1 1983 61</b>
1982 211	1982 210 683	1982 211 85579 -1 1982 212 60795 -1	
1704 414	1704 210 003	1704 212 00/93 -1	3E 0 0 3E 3E 1E 1 1983 61

				ת ה	1178	16 SEC			<b>9</b> 14(	5 MO	77			2200	SSING
NOMINAL		T TIME				FRAMES	A1	A2		B1		<b>C</b> 2	<b>C</b> 3		
DATE	TR DA	AT SECS	IR	UAI	3603	r NAULS		~-	~~	B2		~	0,0	YEAR	
YEAR DAY														LONK	URI
1982 213	1982 2	13 1899	1982	213	86395	-1	3E	0	0	2 <b>E</b>	3E	ì	12	1983	61
1982 214	1982 21				86395		3E		Ō	2E				1983	61
1982 215	1982 2				86395	-	32		ō	22				1983	61
1982 216	1982 21				85003		3E		Ó	2E			12	1983	61
1982 217		16 86395		-	86395		3E		0			0E	1	1983	62
1982 218	1982 2				86395		3 <b>E</b>	0	0	3E	3E	0E	1	1983	62
1982 219	1982 2				86395		3E	0	0	3E	3E	0E	1	1983	62
1982 220	1982 22		1982		3019		3E	0	0	3E	3E	0Z	1	1983	62
1982 221	1982 22	21 10667	1982	221	86395	-1	3E	0	0	1 E	3E	1	0E	1983	62
1982 222	1982 22	22 11	1982	222	86389	-1	3E	0	0	12	3E	1	0E	1983	62
1982 223	1982 22	23 5	1982	223	86389	-1	3E	0	0	I E	3E	1	0E	1983	62
1982 224	1982 22	24 5	- 1982		357		3e	0	0	12	3 <b>E</b>	1	OE	1983	62
1982 225	1982 22	25 2725	1982	225	86389	5154	3E	0	0			12		1983	116
1982 226	1982 22		1982	2 <b>26</b>	86389	5337	3E	0	0	3E	3 <b>E</b>	12	1	1983	
1982 227	1982 22	27 5				4818	3E		0			12		1983	
1982 228	1982 22				85717		3E	0	0			12		1983	
1982 229	1982 22		-			5172	3E		0	2E				1983	
1982 230	1982 23					4803	3E		0	2E				1983	
1982 231	1982 23		1982			5496	3E		0	2E				1983	
1982 233		33 22661				3895	3E		0			IE	-	1983	
1982 234	1982 23					5358	3E		0			12		1983	
1982 235 1982 237	1982 23	-	1982			5178	3E		0			12		1983 1983	
	1982 23				86389	5122		12		3E 3E				1983	
1982 238	1982 21 1982 21					5271		1E			32			1983	
1982 239 1982 241	1982 24					5196	3E		ŏ			IE		1983	
1982 242	1982 24					5309	32		õ			12		1983	
1982 243	1982 24		1982			5566	35		ñ			IE		1983	
1982 245		4 86229	_			5342	3E		ő		3E			1983	
1982 246	1982 24					5349	32		ŏ		32			1983	
1982 247	1982 24		1982			5421	3E		õ	_	3E			1983	
1982 241	1982 24					5196	3E		0			IE		1984	6
1982 242	1982 24					5309	3E	0	0	3E	3E	12	1	1984	6
1982 243	1982 24	13 5	1982	244	3893	5566	3E	0	0	3E	3E	lΕ	1	1984	6
1982 245	1982 24	44 86229	1982	245	86389	5342	3E	0	0	2E	3E	1	IE	1984	6
1982 246	1982 24	6 5	1982	246	86389	5349	3E	0	0		3E		1 E	1984	6
1982 247	1982 24	47 5	1982	248	1621	5421	3 <b>e</b>	0	0	2 <b>E</b>	3E	1	1 E	1984	6
1982 241	1982 24	41 1637	1982	241	86389	5196	3E	0	0			ΙE		1984	9
1982 242	1982 24					5309	3E		0			ΪE	-	1984	9
1982 243	1982 24		1982			5566	3E		0			IE		1984	9
1982 245		4 86229				5342	3E		0		3E		_	1984	9
1982 246	1982 24					5349	38		0		3E			1984	9
1982 247	1982 24		1982			5421 5029	3e 3e		0		3E	เ 0ย		1984	9 10
1982 249	1982 24 1982 25	44 4 4 2 0 1				5295			0					1984	
1982 250	1982 2					5219									
1982 251 1982 253	1982 2					5191	3e 3e		0		3E		1	1984	
1982 255	1982 2					5314	32		õ		32			1984	
1982 255	1982 2					5 5600	35		0		3E			1984	
1982 257	1982 25					5305	32	-	ŏ			ίε		1984	11
1982 258	1982 2					5354	3E		ő			10		1984	
1982 259	1982 2		1982			5193	35		ő			IE		1984	11
1982 261		51 13109				4542	3E		ŏ		ĴĒ			1984	11
1982 262	1982 26					5349	32		ō		3E			1984	11
1982 263	1982 20		1982			5 391	3E		ñ		3E			1984	11
1982 265	1982 20					5136	3E		0		38			1984	12
			_									-	_		

NOR	IRAL	67	ART	TTME	E	T DR	ime	16 SEC			PM	с м	DDE			PROCI	ESSING
DA		78	DAY	SECS	TR			FRAMES	<b>A</b> 1	A2				C2	C3	DA	
YEAR		•••	••••									B2				YEAR	DAT
1982	266		266	5			86389		3 <b>E</b>		0	-	3E			1984	12
1982			267	5			85781		3E		0		3E			1984	12
1982			269	6181			86389		3E		0	22		1		1984	12
1982	-		270	5	1982		86389		3E 3E		0 0		3E 3E			1984 1984	12 12
1982 Gap		1982	fai	lure'	ofsc	an	mech	5466 anism.	35	U	U	25	JE	•	15	1704	14
1983	44	1983		331	1983		86395			IE	0	3E	3E	1	12	1983	110
1983	45	1983		11	1983		86395			IE	-		38			1983	
1983	46	1983		11	1983	47	2347	5438	2E	12		3E	3E	1	12	1983	110
1983	48	1983	48	4763	1983		86395		3₽	0	0			IE		1983	110
1983	49	1983		11	1983		85051		3E		0			IE	-	1983	
1983	50	1983		59	1983		66187	-	32		0		-	12		1983	
1983	52	1983		2939	1983		86395		32		0		32	1		1983	
1983 1983	53 54	1983	53 54	11 59	1983		86395 84923		38 38		0		3E 3E	1		1983 1983	
1983	56	1983	56	1115	1983	-	86395		32		0			OE		1983	
1983	57	1983	57	11	1983		86395		32		ŏ			OE		1983	
1983	58	1983	58	- 11	1983	59		5512	32		ŏ		-	OE		1983	
1983	60	1983	59		1983		86395		32		ŏ		32	1		1984	13
1983	61	1983	61	11	1983		86395		32		Ō	IE	32	1	OE	1984	13
1983	62	1983	62	11	1983	63		5391	3E	0	0	12	3E	1		1984	13
1983	64	1983	64	9963	1983		86395		3E		0			12		1984	13
1983	65	1983	65	11	1983		86395		32		0	JE		12		1984	13
1983	66	1983	66	11	1983		85707		3E		0			12		1984	13
1983 1983	6 <b>8</b> 6 <b>9</b>	1983 1983	68 69	1899	1983 1983	68 69	86395 86395	-	32 32		0	22	32 32			1984 1984	16
1983	70	1983	70	11	1983	71	-	5557	35		0		32			1984	16 16
1983	72	1983	72	75	1983	-	86395	-	32		õ			12	. –	1984	16
1983	73	1983	73	11	1983		86395		32		ŏ			IE		1984	16
1983	74	1983	74	11	1983	75		5417	32		õ			IE		1984	16
1983	76	1983	76	4491	1983	76	86395				0		32			1984	17
1983	77	1983	77	11	1983	77	86395	5320	2E	15	0	3E	32	t	12	1984	17
1983	78	1983	78	11	1983	79	91	5246	2 <b>E</b>	1 E	0	3E	3E	1	12	1984	17
1983	80	1983	80	2667	1983		86395		3E	0	0	3E	3E	12	l	1984	17
1983	81	1983	81	11	1983		86395		3E		0			IE		1984	17
1983	82	1983	82	11	1983		84651		38		0	3E	35	IE	1	1984	17
1993	84	1983	84	843	1983		R6395		3E		0		38	1		1984	19
1983 1983	85 96	1983 1983	85 86	11	1983 1983	75 87	86395	5334	3e 3e		0 0	2E 2E	3E 3E	1		1984	19 19
1983	88	1983	88	5275	1983		86395	-	36		0			OE	12	1984	19
1983	<b>29</b>	1983	80	-11	1983	-	86395		35		n			0E	-	1984	19
1983	90	1983	90	11	1983	91		5388	32		0	32		OE	i	1984	19
1983	92	1983	92	3451	1983		86395		3E	-	ŏ			1	-	1984	20
1983	93	1983	93	11	1983	93	86395	5354	3E	0	0	12	32	l	0E	1984	20
1983	94	1983	94	11	1983	- 94	85499	5280	32	0	0	12	3 <b>E</b>	1	٥Z	1984	20
1983		1983	96	7867	1983		86395		3 <b>e</b>		0			12		1984	20
1983	97	1983	97	11	1983		86395		38		0			12		1984	20
1983	98	1983	98	!1	1983		83787		32		0			12		1984	20
1983		1983		4523			86395		32		0		32			1984	23
1983		1983		11			86395		38		0		32			1984	23
1983		1983		11	1983				32		0		32			1984	23
1983 1983		1983		421 <b>9</b> 11			86395 86395		3e 3e		0			12		1984	23
1983		1983 1983		11			86251		3E		0			1E 1E		1984 1984	23 23
1983		1983		8635			86395			IE			3E			1984	23
1983		1983		11			86395			IE			32			1984	24
				• •							-		ربا در	•		1704	• •

NONTHAL	ST	ART 1	TIME	E	T D	DE	16 SEC				C MC				PROCE	SSING
DATE	TR	DAT	SECS	TR	DAT	SECS	PRAMES	AI	A2	Α3	<b>B</b> 1	C1	C2	C3	DAT	E
YEAR DAT											B2				YEAR	DAT
													ς.			
1983 110	1983	110	11	1983	110	84363	5212	2E	18	0	3E	3E	1	12	1984	24
1983 112	1983	112	571	1983	112	86395	4965	3E	0	0	3E	35	12	1	1984	24
1983 113	1983	113	11				5349	3E	0	0	3E	3E	12	1	1984	24
1983 114	1983	114	11	1983	115	5659	5689	3E	0	0	3E	3E	12	1	1984	24
1983 116	1983	115	85131	1983	116	86395	5428	32	0	0	2E	38	1	12	1984	25
1983 117	1983		11				5165	32	0	0	2E	3E	1	12	1984	25
1983 118	1983	118	11	1983	119	571	5369	3E	0	0	2E	3E	1	I E	1984	25
1983 120	1983	120	3163	1983	120	86395	5144	3E	0	0	3E	3E	0E	1	1984	25
1983 121	1983	121	11	1983	121	86395	5327	3E	0	0	3 <b>E</b>	3E	0E	1	1984	25
1983 122	1983	122	11				5268	3E	0	0	3E	3E	0E	1	1984	25
1983 116	1983	115	85131	1983	116	86395	5428	3E	0	0	2E	3E	1	12	1984	31
1983 117	1983	117	11	1983	117	86395	5165	3E	0	0	2E	3E	1	12	1984	31
1983 118	1983	118	11	1983	119	571	5369	3E	0	0	2E	3E	1	1 E	1984	31
1983 120	1983	120	3163	1983	120	86395	5144	3E	0	0	3 <b>Z</b>	3 <b>E</b>	0e	1	1984	31
1983 121	1983	121	11	1983	121	86395	5327	3E	0	0	3E	3E	0E	1	1984	31
1983 122	1983	122	11	1983	122	85179	5268	3E	0	0	3E	32	0E	I	1984	31
1983 124	1983	124	1339	1983	124	86395	5211	3B	0	0	12	3E	1	0E	1984	32
1983 125	1983	125	11	1983	125	86395	5289	- 3B	0	0	12	3E	1	0E	1984	32
1983 126	1983	126	11	1983	127	12571	6093	3E	0	0	I Z	3 <b>E</b>	-		1984	32
1983 128	1983	127	85915	1983	128	86395	5327	3B	0	0	3E	3E	12	1	1984	32
1983 129	1983	129	11				5 5034	3E	0	0	3 <b>E</b>		IE	1	1984	32
1983 130	1983	130	11	1983			5453	3E		0	38	3E	12	1	1984	32
1983 132	1983	132	3931				5 5094	3E		0	2E	3E	-		1984	33
1983 133		133	11				5279	3E		0	2E	3E	1		1984	33
1983 134	1983		11				5303	3E		0	2E	3E	-		1984	33
1983 136	1983		20 <b>9</b> 1				5 5189	3E		0	32	3E	12	-	1984	33
1983 137	1983		11	- · · · •			5 5340	38	-	0			12		1984	33
1983 138	1983		11				5 5220	3E	-	0		3E		-	1984	33
1983 140	1983	140	12763				5 4543	2E			3E	3E			1984	39
1983 141	1983	141	11				5 5357		12	-	3E	3E	-	12		39
1983 142	1983		11	1983			3 5411		12	-	3E	3E		12		39
1983 144	1983	144	4683	1983	-		5 50 <b>59</b>	3E	-	0	3E		12		1984	39
1983 145	1983	145	11	1983			5 5301	3E		0	3E	3E			1984	39
1983 146	1983		11	1983	-		5346	3e	-	0	3E		12		1984	39
1983 148	1983	148	90 <b>99</b>	1983	148	8639	5 4774	3E	0	0	2E	3E	-	ΙE		40
1983 149	1983	149	11	1983	149	86395	5 4925	3E	0	0	2E	3E	1	12	1984	40
1983 150	1983		11				5208		0	0	2E	32	-		1984	40
1983 152	1983		1019				5 5264	3E	-	0	3E		0E	-	1984	40
1983 153	1983		11				5 5332	3E		0	32		0E	-	1984	40
1983 154	1983		11	1983			5694	3E		0			OE		1984	40
1983 156	1983		11691				5 4557	32		0	12		-		1984	41
1983 157	1983		11				5 4983	3E		0	-	32	_	0E		41
1983 158	1983		11				5210	32		0	IE	3E			1984	41
1983 160	1983	160	9867	1983	161	1631	5 5722	3 <b>e</b>	0	0	3E	3E	ie	1	1984	41

#### **SECTION 6**

## **DATA FORMATS**

# 6.1 Format of SAMS Gridded Retrieved Temperature Tapes (GRID-T)

#### **Overall Structure**

SAMS GRID-T tapes contain latitude-longitude grids of temperature retrievals as a function of pressure, averaged over northbound and southbound orbits. Zonal mean and climatology values are given. Error bars are also included for some grids.

The data are written on 9-track, half-inch wide magnetic tapes at 6250 bpi. Tapes contain one or more files separated by an end of file mark and are terminated by a double end of file mark. The first file on a tape is a header file consisting of two 630-byte physical records written in EBCDIC. These two records are identical and serve to identify the data type (temperature or composition), data start and end dates, generation date, and software version. The overall file structure is shown in Figure 6-1.

Data files are divided into records, each giving a convenient set of results such as an entire grid or a set of profiles. There is one record per physical tape block. Record lengths (hence block lengths) vary according to the type of data they contain. The maximum block length (for a type 7402 record) is 4882 bytes.

All data are written as 16-bit, twos complement integers (CDC users beware: CDC integers are ones complement). Hence, one word occupies exactly 2 bytes.

The byte ordering within words follows the IBM convention of most significant byte first; note that is the reverse of some machines, (e.g., DEC, which puts the least significant byte first).

Each physical block has the following structure:

Word 1	Length of record in bytes (=2N where N is record length in words).
Word 2	Record serial number (this does not always increment by 1).
Word 3	Record type identifier (uniquely defines format of remainder of record).
Word 4 to N-1	Data section of record.
Word N	Least significant byte is the checksum of words 3 to N-1.
Word N+1	Zero.

Hence, the total length of a physical block is two bytes larger than the record length, i.e., 2N+2 bytes, except for record type 7400, which has a record length of 22 bytes and a block length of 40 bytes. The checksum (the low-order byte of word N of the physical record) is the least significant 8 bits of the sum of bytes of words 3 to N-1; i.e., checksum is between 0 and 255, it is put into byte 2N, and it is the sum of bytes 5 to 2N-2. Note that byte 2N-1 should be ignored.

#### **Example**

The following type 7400 block is at the beginning of one of the tapes. It is 12 words (24 bytes) long. Those words in decimal are as follows: 22, 1, 7400, 1, 1979, 281, 7401, 7402, 7403, 0, -13, 0. The bytes (in the order that they are on the tape) are: 0, 22, 0, 1, 28, 232, 0, 1, 7, 187, 1, 25, 28, 233, 28, 234, 28, 235, 0, 0, 255, 243, 0, 0. To verify the checksum we take

28+232+0+1+7+1+25+28+234+28+235+0+0=1267; 1267=4\*256+243,

so the checksum should be 243, and byte 22, indeed, contains this value.

### PRECEDING PAGE BLANK NOT FILMED

FI	LE	#1

|

FILE #2, #3,....

HEADER HEADER O DATA O RECORD RECORD F F	E O F DATA	E O DATA F		E O F
---	---------------------	------------------	--	-------------

Figure 6-1. Overall structure of SAMS tapes.

#### **Header File Structure**

The header file consists of two identical 630-byte physical records that identify a tape as a SAMS product. The format of each of the records is given here. The character "b" indicates a blank, and the character "x" indicates characters that vary from tape to tape.

#### **Characters**

bNIMBUS-7bSAMS	(identifies tape as SAMS product)	1 - 14
bXXXXXXXXXXX	(identifies tape type)	15 - 26
bSQbNOb	(sequence number)	27 - 33
XXXXX	(5-digit sequence number)	34 - 38
XX	(redo character and copy number 1 or 2)	39 - 40
<sub>b</sub> START <sub>b</sub> 19XX <sub>b</sub> XXX	(data start year and day of year)	41 - 55
bTOb19XXbXXX	(data end year and day of year)	56 - 67
bGENb19XXbXXXbXXXXX	(year, day of year, and hours, minutes, seconds tape was generated	68 - 87
bPROGRAMbSAMSbXXXXXXX	(SAMS software version name)	88 - 109
bXX/XX/XX	(software version month, day, year)	110-118
bbb	(blanks to fill out the record)	119-630

The following is an example of a SAMS temperature tape header record.

#### **Example**

A SAMS temperature tape starting in December 1978 and ending in December 1979 was generated on December 27, 1984 at 19:10:15 using software version VERVS02A.

bNIMBUS-7bSAMSbTEMPERATUREbSQbNOb83581-2bSTARTb1978b358bTOb1979b365bGENb1984b 362b191015bPROGRAMbSAMSbVERVS02Ab12/24/84 followed by 512 blanks.

The SAMS header record tape type, sequence number, redo character, and copy number fields are explained here. The tape type can be one of two values and starts with character 16.

CHARACTERS 16 - 26: This field distinguishes the type of SAMS tapes. TEMPERATURE identifies a tape as a Gridded Retrieved Temperature Tape. COMPOSITION identifies a tape as a Zonal Mean Methane and Nitrous Oxide Composition tape.

The sequence number consists of five digits and starts with character 34 of the record.

CHARACTER 34: The last digit of the year in which data were acquired.

CHARACTERS 35 - 37:	Day of the year on which data were acquired.
CHARACTER 38:	Value is 1 unless necessary to remove ambiguities in the year (character 34), as may occur if data are acquired on or after December 24, 1988.
CHARACTER 39:	Redo character. It will remain as a hyphen unless there is a remake of the tape. In this case, an ascending alpha character will replace the hyphen.
CHARACTER 40:	Tape copy number. The original tape will have a value of 1; the archived copy will have a value of 2.

#### **Data File Structure**

GRID-T tape data files have the following types of records:

- a) File header record (type 7400).
- b) Data records (type 7402 or 7403).
- c) End of file mark.

Table 6-1 shows the data file structure.

There is one file per "data day" of data. A "data day" is a period when the instrument is in just one operating mode. Mode changes generally occur near midnight, in which case one data day corresponds to a calendar day. However, this is not always the case and sometimes there will be more than one data day in one calendar day.

Format of File Header Record

Table 6-2 shows the structure of file header records.

#### Format of latitude-longitude-ln(pressure) retrieved temperatures, block type 7402

This grid contains temperature averaged over day and night at 62 pressure levels. The error bars are not given in this format, but can be obtained for fewer levels (which should be adequate) from 7403 blocks. Table 6-3 shows the block type 7402 format. Detailed comments on selected words follow:

Word	Description
8 + 64(n-1)	Latitude of grid point in degrees N*100 (-5000 to 6750).
9 + 64(n-1)	Longitude of grid point in degrees $E^{*100}$ (-18000 to 20000) (19000 and 20000 indicate zonal mean and climatological first guess).
10 + 64(n-1) to 71 + 64 (n-1)	Temperature profile; 62 values in K*100 starting at $\ln(p_0/p) = 1.4$ , up every 0.2 in $\ln(p_0/p)$ up to 13.6 (p_0=1000).

The value of n ranges from 1 to 38 for each record. There are 64 words for each value of n: latitude, longitude, and temperature for n=1 are for 180°W, those for n=2 are for 170°W, and continuing at 10° intervals to 170°E for n=36. The last two groups of 64 words contain zonal means (n=37) and climatology (n=38). The first record in a type 7402 block contains data for latitude 50°S, the second contains data for 47.5°S, the following records are spaced at 2.5° intervals, and the last record contains data for 67.5°N.

## GRID-T Temperature Tape Data File Structure

## Record #

1	Туре 7400	Header Record
2	Type 7402	Retrieved Temperature Grid at - 50.0N
3	Туре 7402	Retrieved Temperature Grid at - 47.5N
		:
49	Туре 7402	Retrieved Temperature Grid at + 67.5N
50	Туре 7403	Temperature Grid at Level 2303
51	Туре 7403	Error Bars at Level 2303
52	Туре 7403	Temperature Grid at Level 3507
53	Туре 7403	Error Bars at Level 3507
÷		:
68	Туре 7403	Temperature Grid at Level 12717
69	Туре 7403	Error Bars at Level 12717
		EOF

NOTE:

.4

Not all these records will necessarily be present.

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## GRID-T Temperature Tape Data File Header Record Format

16-Bit Word #	MSB 16		LSB 1
1	Record length	(22 bytes)	
2	Serial number		
3	Record type 7400		
4	Data file number		
5	Data year		
6	Data day		
7	7401		
8	7402		
9	7403		
10	Zero fill		
11		Checksum	
12-20	Zero fill		

## GRID-T Temperature Tape Latitude-Longitude-1n (Pressure) Record Format Type 7402

16-Bit Word #	MSB 16	LSB 1
1	Record length 4880 bytes	
2	Serial number	
3	Record type 7402	
4	Data day	
5	Data year	
6	Processing day	
7	Processing year	
8	Grid point latitude (degrees *100)	
9	Grid point longitude (degrees *100)	-180 E
10	Temperature (degrees K *100) for 62 pressure levels	
:		
71		
72	Grid point latitude degrees *100	
73	Grid point longitude degrees *100	-170 E

16-Bit <u>Word #</u>	MSB 16	LSB 1
74	Temperature degrees K *100 for 62 pressure levels	
÷		
135		
136 :	:	
2248	Grid point latitude degrees *100	
2249	Grid point longitude degrees *100	+170 E
2250 	Temperature degrees K *100 for 62 pressure levels	
2311		
2312	Grid point latitude degrees *100	
2313	Grid point longitude degrees *100	+190 E
2314 :	Zonal mean	
2375		

GRID-T Temperature Tape Latitude-Longitude-1n (Pressure) Record Format Type 7402 (continued)

GRID-1 Temperatu	ie rape Latitude-Longitude-In (rie	ssure) Record Pormat Type 7402 (com	inucci)
16-Bit <u>Word #</u>	MSB 16		LSB 1
2376	Grid point latitude degrees *100	)	
2377	Grid point longitude degrees *1	00	+200 E
2378	Climatology, i.e., first guess pro	ofile	
:			
2439			
2440		Checksum	
2441	Zero fill		

GRID-T Temperature Tape Latitude-Longitude-1n (Pressure) Record Format Type 7402 (continued)

-----

16-Bit Word #	MSB 16		LSB 1
ì	Record length	3504 bytes	·····
2	Serial number		
3	Record type 7403		
4	Measurement type	(3)	
5	Data day		
6	Data year		
7	Processing day		
8	Processing year		
9	Ignore		
10	Scale factor	(divide by this integer)	
11	Data type 2 or 102		
12	Pressure level		
13	Ignore		
14	Ignore		
15	Ignore		

## Grid-T Temperature Tape Latitude-Longitude Grid Record Format Type 7403

58

# Grid-T Temperature Tape Latitude-Longitude Grid Record Format Type 7403 (continued)

16-Bit <u>Word #</u>	MSB 16	LSB 1
16 :	Zero fill	
23		
24 :	Temperatures or error bars for Word 12 level at -50N, -180 E to+170 E	
59		
60 :	Temperatures or error bars for Word 12 level at -47.5N, -180 E to +170	E
95 :		
1680	Temperatures or error bars for Word 12 level at +65N, -180 E to +170 E	,
1715		
1716	Temperatures or error bars for Word 12 level at +67.5N, -180E to +170	E
1751		
1752	Checksum	
1753	Zero fill	

The climatology should always contain valid data. Bad or missing data are replaced with a fill value of -32768.

#### Format of latitude-longitude grid block type 7403

These contain grids that extend from 50°S to 67.5°N (every 2.5°) and 180°W to 170°E (every 10°). Each quantity gridded is put into a separate record, which contains a header and then an array of 16-bit integers. Grids of an item and its error count as separate quantities and accordingly are put into different records (adjacent on the tape). All grids are for the mean of northbound and southbound measurements. Table 6-4 shows the format for type 7403 records. Detailed descriptions of selected words follow. In these descriptions, the quantities in brackets (words 9, 13, 14, 15) should be ignored.

Detailed descriptions of block type 7403 words

Word	Description			
3	7403 identifier.			
4	1, 2, or 3 meaning northbound orbits, southbound orbits or combined average (only type 3 present on these tapes).			
5	Day of data.	Day of data.		
6	Year of data.			
9	[address in orbit g	[address in orbit grid records of quantity gridded].		
10	Scale factor: divide data by this integer.			
11	Data type:	1-eigenfunction coefficient; 2-temperature; 101-error in eigenfunction coefficient; 102-temperature error (only types 2 and 102 present on these tapes).		
12	Level:	for type 1 or 101, this is eigenfunction number (1, 2, etc.), for 2 or 102, this is $1000*\ln(\text{pressure/p}_0)$ where $p_0$ is $1000$ mb.		
13	[Format/version r	[Format/version number of CAL, RETRIEVE programs, etc.].		
14	[Format/version number of OGRID program].			
15	[Format/version number of LGRID program].			
16-23	Spares, set to 0.			
24-23+36*48	Data array A dimensioned (36,48) in order A(1,1), A(2,1), A(3,1)A(36,48). -32768 means no information. A(1,1) is for (50°S, 180°W), A(2,1) is (50°S, 170°W),etc.; A(I,J) is for (J-21)*2.5°N, (I-19)*10°E.			

Temperatures are given at the following pressures in mbar (figures in parentheses are the corresponding values of word 12): 100 (2303), 30 (3507), 10 (4605), 3 (5809), 1 (6908), 0.3 (8112), 0.1 (9210), 0.03 (10414), 0.01 (11513), 0.003 (12717).

At the very top pressure levels, the values are primarily climatology. Significant measurements extend up to about 0.03 mb to 0.01 mb (70 to 80 km).

### 6.2 Format of SAMS Zonal Mean Methane and Nitrous Oxide Composition Tape (ZMT-G)

#### **Overall Structure**

The ZMT-G tape contains zonal mean values of nitrous-oxide and methane retrievals and errors expressed as a function of pressure and averaged over 24-hour periods.

The data are written on 9-track, half-inch wide magnetic tapes at 6250 bpi. Tapes contain one or more data files and are terminated by a double end of file mark. Like the SAMS temperature tape, the SAMS composition tape has a tape header file as the first file on the tape. It contains two identical 630-byte physical records written in EBCDIC. The description of header records in Section 6.1 applies to the ZMT-G tapes, as well as to the temperature tapes.

Data files are divided into records, each giving values for 1 day. There is one record per physical tape block. Refer to Table 6-5 for the ZMT-G tape file structure.

All data are written as 16-bit twos complement integers (CDC users beware: CDC integers are ones complement). Hence, one word occupies exactly 2 bytes.

The byte ordering within words follows the IBM convention of most significant byte first; note that is the reverse of some machines, (e.g., DEC, which puts the least significant byte first).

Each physical block has the following structure:

Word 1	Length of record in bytes (=2N where N is record length in words) (2N=5986 in this case).
Word 2	Record serial number.
Word 3	Record type identifier (uniquely defines format of remainder of record) (7405 or 7406 in this case).
Word 4 to N-2	Data section of record.
Word N-1	Least significant byte is the checksum of words 3 to N-2.
Word N	Zero.

Hence, the total length of a physical block is 2N bytes, the same as the record length.

The checksum (the low-order byte of word N-1 of the physical record) is the least significant 8 bits of the sum of bytes of words 3 to N-2; i.e., the checksum is between 0 and 255, it is put into byte 2N-2, and is the sum of bytes 5 to 2N-4. Note that byte 2N-3 should be ignored. An example of checksum computation is given in the description of the overall structure of the temperature tape.

#### Format of latitude-In(pressure) zonal mean composition blocks 7405 and 7406

This grid contains zonal 24-hour average (day and night) mean mixing ratios of methane and nitrous oxide at 31 pressure levels. The SAMS cannot measure methane and nitrous oxide simultaneously. Consequently, measurements of only one of the above gases are made during each 24-hour period. Table 6-6 gives the format of block types 7405 and 7406.

Detailed descriptions of selected words follow:

Word	Description
3	Identifier (7405-nitrous oxide; 7406-methane).
6	Enabled channel (8-nitrous oxide; 9-methane).
7	Sieve setting of enabled PMC.
8	Sieve setting of clamped PMC.
9	Sieve setting of A1 PMC (temperature sounding).
10	Sieve setting of channel C1 PMC (temperature sounding).
11	Creation date (day).
12	Creation date (year).
13	Number of elements in each profile.
14	Bottom level $(\ln(P_0/P)*10)$ .
15	Top level (units as word 14).
15 + 31 (n-1) to 45 + 62 (n-1)	Mixing ratio profile: 31 values (in ppbv*50 for nitrous oxide or ppmv*10000 for methane), starting at $\ln(P_0/P)=3.0$ , every 0.2 in $\ln(P_0/P)$ up to 9.0 (P_0=1013.25 mb); n varies from 1 to 48; n=1 corresponds to 50°S, n=2 to 47.5°S, continuing at 2.5° intervals to 67.5°N for n=31.

Bad or missing data are set to -32768 (note that in 16 bit twos complement arithmetic -32768=32767+1=32768, although +32768 is generally taken to be an invalid number).

There is one block per "data day" of data. A "data day" is a period when the instrument is in just one operating mode. Mode changes generally occur near midnight, in which case one data day corresponds to a calendar day. However, this is not always the case and sometimes there will be more than one data day in one calendar day.

## ZMT-G Composition Tape Data File Structure

### Record #

1	Nitrous oxide (7405) or methane (7406) for one data day
2	Nitrous oxide (7405) or methane (7406) for one data day
3	Nitrous oxide (7405) or methane (7406) for one data day
÷	:
N-1	Nitrous oxide (7405) or methane (7406) for one data day
N	Nitrous oxide (7405) or methane (7406) for one data day
	End of file mark

ZMT-G Composition Ta	pe Latitude-In (Press	re) Zonal Mcan Record	Format Types 7405 and 7406

16-Bit <u>Word #</u>	MSB 16	LSB 1
1	Record length	
2	Serial number	
3	Record type: 7405 or 7406	
4	Data day	
5	Data year	
6	Enabled channel	
7	Sieve setting enabled PMC	
8	Sieve setting clamped PMC	
9	Sieve setting A1 PMC	
10	Sieve setting C1 PMC	
11	Processing day	
12	Processing year	
13	Number of elements in profile	
14	Bottom level	

ZMT-G Composition Tape Latitude-In (Pressure) Zonal Mean Record Format Types 7405 and 7406 (continued)

16-Bit Word #	MSB 16	LSB 1
15	Top level	
16 :	Retrieved mixing ratio at - 50N 31 values	
46		
47 :	Retrieved mixing ratio at - 47.5N 31 values	
77 :	:	
1442 	Retrieved mixing ratio at + 65N 31 values	
1447		
1473 	Retrieved mixing ratio at +67.5N 31 values	
1503		

ZMT-G Composition Tape Latitude-In (Pressure) Zonal Mean Record Format Types 7405 and 7406 (continued)

16-Bit <u>Word #</u>	MSB 16		LSB 1
1504 :	Mixing ratio errors at - 50N 31 values		
1534			
1535 	Mixing ratio errors at -47.5N 31 values		
1565 :			
2930 :	Mixing ratio errors at + 65 N 31 values		
2960			
2961 :	Mixing ratio errors at + 67.5N 31 values		
2991			
2992		Checksum	
2993	Zero fill	· · · · · · · · · · · · · · · · · · ·	

66

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## LIST OF ACRONYMS, INITIALS, AND ABBREVIATIONS

A/D	Analog/digital
bpi	Bits per inch (9-track tapes)
CDC	Control Data Corporation
Composition Tape	SAMS Zonal Mean Methane and Nitrous Oxide Composition Tape. Also known as a Trace Species, Constituent or ZMT-G tape.
DEC	Digital Equipment Corporation
EBCDIC	Extended Binary Coded Decimal Interchange Code
EOF	End of file mark
FOV	Field of view
GRID-T	Gridded Retrieved Temperature Tape
GSFC	Goddard Space Flight Center
IBM	International Business Machines
ILT	Image Location Tape
JCL	Job Control Language
JGR	Journal of Geophysical Research
LIMS	Limb Infrared Monitor of the Stratosphere
LSB	Least Significant Bit
LTE	Local Thermodynamic Equilibrium
MSB	Most Significant Bit
МАР	Middle Atmosphere Program
NASA	National Aeronautics and Space Administration
NET	Nimbus Experiment Team
NSSDC	National Space Science Data Center
РМС	Pressure-modulated Channel
РМ	Pressure-modulated
PSD	Phase Sensitive Detector

PM Signal	Pressure Modulated Signal (the PM and WB signals are the two radiation components measured by each of the SAMS detectors)
PMR	Pressure Modulator Radiometer
QL	Quick look
SAMS	Stratospheric and Mesospheric Sounder
S/N	Signal-to-Noise
SCOSTEP	Scientific Committee on Solar-Terrestrial Physics
SCR	Selective Chopper Radiometer
STX	ST Systems Corporation
Temperature Tape	A SAMS Gridded Retrieved Temperature Tape (GRID-T)
TGS	Triglycine sulfate
WB	Wideband
WB Signal	Wideband Signal
ZMT-G	Zonal Means Tape - Gas. Same as a Composition Tape

#### APPENDIX A

1

# SEQUENCE NUMBERS OF SAMS DATA TAPES

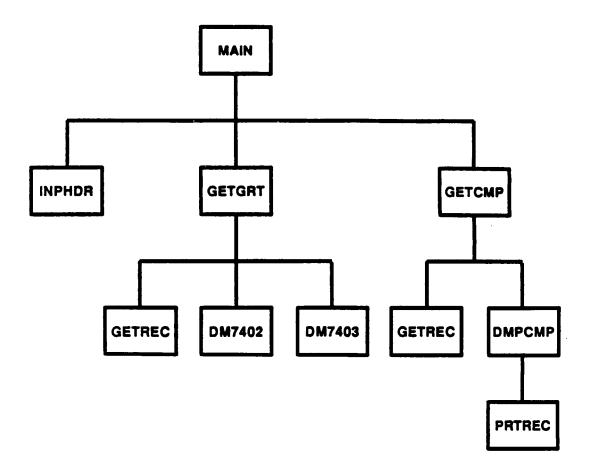
This table shows the sequence numbers and start and end dates of SAMS temperature and composition tapes that are archived at the NSSDC. The file column includes tape header and data files.

Type	Sequence	Start	End	<u>Files</u>
Temperature	83581-2	1078/358	1979/365	273
Temperature	00011-2	1980/001	1980/366	273
Temperature	10021-2	1981/002	1981/364	273
Temperature	20021-2	1982/002	1983/160	298
Composition	90011-2	1979/001	1981/364	2

#### **APPENDIX B**

## FORTRAN SOURCE LISTING OF SAMS TAPE DUMP

This appendix contains the FORTRAN source listing of a program that will enable the user to produce formatted dumps of retrieved temperature data, and methane and nitrous oxide mixing ratios on SAMS GRID-T and ZMT-G tapes, respectively. Figure B-1 is a hierarchy chart showing the relation between dump program modules. The program is written in VS FORTRAN for use on an IBM 3081 system operating under MVS.



B-1. Hierarachy Chart for SAMS tape dump program.

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#### APPENDIX B

00000020 C PROGRAM - MATH 00000030 C PURPOSE - THIS PROGRAM GENERATES A FORMATTED DUMP OF A SAMS GRIDDED 00000040 00000050 RETRIEVED TEMPERATURE TAPE (GRID-T) OR A SAMS ZONAL MEAN METHANE C AND NITROUS-OXIDE COMPOSITION TAPE (ZMT-G). THE BYTE ORDERING C 00000060 WITHIN THE 16-BIT WORDS MUST FOLLOW THE IBM CONVENTION. 00000070 C 00000080 C <<< MARNING >>> LARGE VOLUMES OF PRINTOUT ARE GENERATED BY DUMPING 00000090 C SAMS TAPES UNLESS DUMPING IS RESTRICTED BY CAREFULLY USING THE DUMP DATES AND DUMP OPTIONS. 00000100 C 00000110 Ĉ 00000120 ĉ THE PROGRAM READS ONE INPUT CARD WHICH DEFINES THE TAPE VOLSER, 00000130 C DUMP START AND STOP DATES, AND EIGHT DUMP OPTIONS. THE DATES 00000140 HAVE FORMAT YYDDD WHERE YY ARE THE TWO LOW ORDER DIGITS OF THE YEAR 00000150 AND DDD IS THE DAY OF YEAR. THE EIGHT DUMP OPTIONS CORRESPOND TO 00000160 r C C DIFFERENT RECORD TYPES: 7400, 7402, 7403 TEMPERATURE GRID, C 00000170 7403 TEMPERATURE GRID ERRORS, 7405 N20 MIXING RATIOS, 7405 N20 MIXING RATIO ERRORS, 7406 CH4 MIXING RATIOS, 7406 CH4 MIXING RATIO c 00000180 C 00000190 C ERRORS. OPTIONS CAN HAVE ONE OF THO VALUES, O FOR NO DUMP, 00000200 1 FOR DUMP. 00000210 C Ĉ 00000220 A TYPICAL INPUT CARD WITH THE TAPE VOLSER NUMBER STARTING IN 00000230 C COLUMN ONE HOULD APPEAR AS FOLLOWS: 00000240 C TAP001 79358 80004 0 0 0 0 1 0 1 0 00000250 C THIS HOULD CAUSE THE CH4 AND N20 MIXING RATIOS FOR DAY 358 1979 C 00000260 TO DAY 004 1980 TO BE DUMPED. 00000270 c 00000280 THE PROGRAM DETERMINES WHAT TYPE OF TAPE IS BEING PROCESSED: 00000290 EITHER A GRIDDED RETRIEVED TEMPERATURE TAPE OR AN CHA/N2O ZONAL MEAN TAPE. IT SEARCHES A TAPE FOR RECORDS THAT ARE MITHIN THE 00000300 Ċ Ċ 00000310 START AND STOP DATES GIVEN ON THE INPUT CARD, AND DUMPS RECORDS 00000320 C FOR MHICH OPTIONS ARE SET TO 1. 00000330 00000340 THE START AND STOP DATES HUST BE ON OR AFTER 78297 (THE SATELLITE 00000350 LAUNCH DATE, DAY 297 1978). THE STOP DATE MUST NOT BE LESS THAN 00000360 THE START DATE, AND MUST NOT BE GREATER THAN 99365, DAY 365 00000370 Ċ (DEC. 31) 1999. 00000380 C 00000390 LANGUAGE - VS FORTRAN 00000400 С 00000410 C C COMPUTER - IBM 3081 HVS 00000420 00000430 с C DATA SETS -00000440 FT05F001 - INPUT DATA CARD FT06F001 - HARDCOPY PRINTOUT 00000450 C 00000460 ٢ FT10F001 - INPUT SAMS TAPE 0000670 C 00000480 C SUBPROGRAM DEFINITIONS -00000490 INPHOR = READ AND CHECK INPUTS, MOUNT TAPE, CHECK HEADER 00000500 GETGRT = READ RECORDS FROM A TEMPERATURE TAPE GETCMP = READ RECORDS FROM A COMPOSITION TAPE 00000510 00000520 DM7402 = DUMP TYPE 7402 RECORD DM7403 = DUMP TYPE 7403 RECORD (GRID AND GRID ERRORS) С 00000530 C 00000540 DMPCMP = DETERMINE WHICH PART OF TYPE 7405 AND 7406 RECORDS TO DUMP 00000550 С PRTREC = DUMP TYPE 7405 AND 7406 RECORDS С 00000560 GETREC = READ RECORDS FROM SAMS TAPE 00000570 C 00000580 C C LIBRARY SUBPROGRAM DEFINITIONS -00000590 MOUNT = FTIO MOUNT & TAPE (FTIO = FORTRAN INPUT/OUTPUT PACKAGE) 00000600 C = FTIO POSITION TAPE TO A FILE POSN 00000610 C FREAD . FTIO READ TAPE RECORDS C 00000620 # PERFORMS LOGICAL COMPARISON OF BYTES С KCLC 00000630 FMOVE # STORAGE-TO-STORAGE MOVEMENT OF BYTES C 00000640 ZTIME = GETS CURRENT DATE AND TIME 00000650 С 00000660 C SUBROUTINES CALLED - ZTIME.INPHOR.GETGRT.GETCMP 00000670 00000680 C PROGRAMMER - TOM NUTTER 03/05/85 SASC TECHNOLOGIES, INC. 00000690 

APPENDIX B

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		**1	·.*8
ISN	1	CHARACTER*8 VRNAME, VRDATE	00000710
ISN	2	CHARACTER*1 RUNDAT(28)	00000720
ISN	3	INTEGER TAPTYP, ISTART, ISTOP, FILNUM, OPTION(8), IRETN	00000730
ISN	4	DATA TAPTYP/0/, IRETN/0/, VRNAME/'SAMDMP1A'/, VRDATE/'03/15/85'/,	00000740
	•	FILNUN/1/, OPTION/8+0/, ISTART, ISTOP/2+-99999.	00000750
		CHAN GET RUN DATE AND TIME, PRINT JOB HEADING AND VERSION NAME	00000760
ISN	5	CALL ZTIME (RUNDAT, 3)	00000770
ISN	6	WRITE(6,5010) VRNAME, VRDATE, RUNDAT	00000780
		CHAN READ AND CHECK CARD INPUTS, MOUNT TAPE, CHECK TAPE HEADER,	00000790
		CHAN DETERMINE TAPE TYPE	00000800
ISN	7	CALL INPHDR(TAPTYP,ISTART,ISTOP,FILNUM,OPTION,IRETN)	00000810
ISN	8	IF (IRETN.GT.O) THEN	00000820
ISN	9	WRITE(6,5020) IRETN	00000830
ISN	10	STOP 999	00000840
ISN	11	ENDIF	00000850
		C*** DUMP A TEMPERATURE TAPE OR A COMPOSITION TAPE	00000860
ISN	12	IF (TAPTYP.EQ.1) THEN	00000870
ISN	13	CALL GETGRT(ISTART, ISTOP, FILNUM, OPTION, IRETN)	00000880
ISN	14	IF (IRETN.GT.O) THEN	00000890
ISN	15	WRITE(6,5030) IRETN	00000900
ISN	16	STOP 999	00000910
ISN	17	ENDIF	00000920
ISN	18	ELSE	00000930
ISN	19	CALL GETCMP(ISTART, ISTOP, FILNUM, OPTION, IRETN)	00000940
ISN	20	IF (IRETN.GT.O) THEN	00000950
ISN	21	WRITE(6,5040) IRETN	00000960
ISN	22	STOP 999	00000970
ISN	23	ENDIF	00000980
ISN	24	ENDIF	00000990
		C*** PRINT JOB COMPLETION MESSAGE	00001000
ISN	25	WRITE(6,5050)	00001010 00001020
ISN	26	STOP	00001030
ISN	27	5010 FORMAT(//1X,5('<'),' NIMBUS-7 SAMS TAPE DUMP PROGRAM ', 5('>')/' PROGRAM VERSION: ',A8,5X,'DATE LAST MODIFIED: ',A8	00001030
			00001050
TCN	28	<pre>. /' RUN DATE ',28A1//) 5020 FORMAT(/' *** MAIN *** ERROR CODE OF ',13,' RETURNED FROM ',</pre>	00001060
ISN	20	. '"INPHDR" ABORT RUN. '/)	00001070
ISN	29	5030 FORMAT(/' *** MAIN *** ERROR CODE OF ',13,' RETURNED FROM ',	00001080
TON	27	. "GETGRT" ABORT RUN. '/)	00001090
ISN	30	5040 FORMAT(/' *** MAIN *** ERROR CODE OF ',I3,' RETURNED FROM ',	00001100
1.94	20	. '"GETCMP" ABORT RUN. '/)	00001110
ISN	31	5050 FORMAT(//' END OF DUMP')	00001120
ISN	32	END	00001130
1.914	26	LITE	

		**1	*8
ISN	1	SUBROUTINE INPHOR(TAPTYP, ISTART, ISTOP, FILNUH, OPTION, IRETN)	00000010
		C <del>x#xxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxx</del>	+00000020
		C SUBROUTINE - INPHOR	00000030
		c	00000040
		C PURPOSE - READ AND ECHO PRINT CARD INPUT, CHECK INPUT VALUES, MOUNT	00000050
		C TAPE AND CHECK TAPE HEADER, AND DETERMINE TAPE TYPE.	00000060
			00000070
		C ARGUMENTS PASSED -	00000080
		C TAPTYP = TAPE TYPE (1=TEMPERATURE TAPE, 2=CH4-N2O ZONAL MEAN) C ISTART = START YEAR, DAY OF YEAR (YYDDD) OF DATA TO BE DUMPED	00000090 00000100
		C ISTOP = STOP YEAR, DAY OF YEAR (YYODD) OF DATA TO BE DUMPED	00000110
		C FILNUM = TAPE FILE NUMBER	00000120
		C OPTION = DUMP OPTIONS FROM INPUT CARD (1=DUMP, 0=NO DUMP)	00000130
		C IRETN = ERROR RETURN CODE (IRETN=0 NO ERROR, >0 ERROR)	00000140
		C	00000150
		C SUBPROGRAMS CALLED - MOUNT, FREAD, KCLC	00000160
		C	00000170
		C CALLED BY - MAIN	00000180
		c	00000190
		C PROGRAMMER - TOM NUTTER 03/05/85	00000200
	•		
ISN	2	INTEGER TAPTYP, ISTART, ISTOP, FILNUM, OPTION(8), IRETN	00000220
ISN ISN	3	CHARACTER*1 INLINE(80),HEADER(630) CHARACTER*4 DUMP	00000230 00000240
ISN	5	CHARACTER*3 DASH	00000250
ISN	- 6	CHARACTER*6 VOLSER	00000260
ISN	7	CHARACTER*11 TEMP, COMP, ACTION(8)	00000270
ISN	8	CHARACTER*15 SAMSID	00000280
		C <del>XFF</del> INITIALIZE VARIABLES	00000290
ISN	9	DATA DASH/''/,DUMP/'DUMP'/,HEADER/630*' '/	00000300
ISN	10	DATA TEMP/'TEMPERATURE'/,COMP/'COMPOSITION'/	00000310
ISN	11	DATA SAMSID/' NIMBUS-7 SAMS '/,ACTION/8*'DO NOT DUMP'/	00000320
ISN	12	C*** READ AND ECHO PRINT INPUT CARDS WRITE(6,5010)	00000330 00000340
ISN	13	READ(5,5020,END=900,ERR=910) INLINE	00000350
ISN	14	WRITE(6,5030) INLINE	00000360
ISN	15	REWIND 5	00000370
		C*** READ INPUTS AGAIN AND CHECK START AND STOP DATES	00000380
ISN	16	READ(5,5060,END=900,ERR=910) VOLSER,ISTART,ISTOP,OPTION	00000390
ISN	17	IF ((ISTART.LT.78297.OR.ISTOP.LT.78297).OR.(ISTART.GT.ISTOP)) THE	N00000400
ISN	18	IRETN=1	00000410
ISN	19	WRITE(6,5065) ISTART,ISTOP,IRETN	00000420
ISN	20	RETURN ENDIF	00000430
ISN	21	CHAR CHECK DUMP OPTIONS FOR UNREASONABLE VALUES	00000440 00000450
ISN	22	DO 200 I=1,8	00000460
ISN	23	IF (OPTION(I).GT.1) THEN	00000470
ISN	24	IRETN=9	00000480
ISN	25	WRITE(6,5070) I,OPTION(I)	00000490
ISN	26	ENDIF	00000500
ISN	27	200 CONTINUE	00000510
ISN	28	IF (IRETN.GT.0) RETURN	00000520
		C*** MOUNT AND POSITION TAPE TO HEADER FILE, READ AND PRINT HEADER, AND	
701	70	C*** CHECK RECORD LENGTH	00000540
ISN ISN	30 31	CALL MOUNT(1,10,VOLSER,FILNUM) CALL FREAD(HEADER,10,LENGTH,*920,*930)	00000550 00000560
ISN	32	WRITE(6,5075) HEADER	00000570
ISN	33	IF (LENGTH.NE.630) THEN	00000580
ISN	34	IRETN=2	00000590
ISN	35	WRITE(6,5080) LENGTH, IRETN	00000600
ISN	36	RETURN	00000610
ISN	37	ENDIF	00000620
		C*** CHECK IF TAPE IS A SAMS TAPE	00000630
ISN	38	IF (KCLC(HEADER,1,SAMSID,1,15).NE.0) THEN	00000640
ISN	39		00000650
ISN	40	WRITE(6,5090) IRETN	00000660
ISN ISN	41 42	RETURN ENDIF	00000670 00000680
2	76	C*** DETERMINE WHETHER TAPE IS A TEMPERATURE OR COMPOSITION TAPE, AND	00000690
		C*** PRINT OUT DUMP OPTION TABLE	00000700

		*	*1234567.1						
ISN	43		IYRSTR=ISTART/1000	00000710					
ISN	44		IDYSTR=ISTART-IYRSTR#1000	00000720					
ISN	45		IYRSTP=ISTOP/1000	00000730 00000740					
ISN	46		IDIDIF-IDIGE-XIKDIFA1000						
ISN	47								
ISN	48		I (GFIAGHLI):EQII) ACTIONLI/ DONN						
ISN	50	300							
ISN	51		IF (KCLC(HEADER,16,TEMP,1,11).EQ.0) THEN TAPTYP=1	00000780					
ISN	52 53		WRITE(6,5110) TEMP,IDYSTR,IYRSTR,IDYSTP,IYRSTP	00000800					
ISN ISN	55 54		WRITE(6,5120) TEMP,(DASH,J=1,7)	00000810					
ISN	55		WRITE(6,5130) (OPTION(I),ACTION(I),I=1,4)	00000820					
ISN	56		ELSE	00000830					
ISN	57		IF (KCLC(HEADER, 16, COMP, 1, 11). EQ.0) THEN	00000840					
ISN	58		TAPTYP=2	00000850					
ISN	59		WRITE(6,5110) COMP, IDYSTR, IYRSTR, IDYSTP, IYRSTP	00000860					
ISN	60		WRITE(6,5120) COMP,(DASH, I=1,7)	00000870					
ISN	61		WRITE(6,5140) (OPTION(I),ACTION(I),I=5,8)	00000880					
ISN	62		ELSE	00000 <b>890</b>					
ISN	63		IRETN=4	00000900					
ISN	64		WRITE(6,5100) IRETN	00000910					
ISN	65		RETURN	00000920					
ISN	66		ENDIF	00000930					
ISN	67		ENDIF	00000950					
ISN	68		RETURN	00000960					
TCN	69	000	INPUT CARD END-OF-FILE CONTINUE	00000970					
ISN ISN	70	700	IRETN=5	00000980					
ISN	71		WRITE(6,5050) IRETN	00000990					
ISN	72		RETURN	00001000					
2011		C×××	INPUT CARD READ ERROR	00001010					
ISN	73	910	CONTINUE	00001020					
ISN	74		IRETN=6	00001030					
ISN	75		WRITE(6,5040) IRETN	00001040					
ISN	76		RETURN	00001050					
		C***	TAPE END-OF-FILE	00001060					
ISN	77	920	CONTINUE	00001070					
ISN	78		IRETN=7	00001080					
ISN	79		WRITE(6,5150) IRETN	00001090					
ISN	80		RETURN	00001100 00001110					
			TAPE READ ERROR	00001110					
ISN	81	930		00001120					
ISN	82 83		IRETN=8 WRITE(6,5160) IRETN	00001140					
ISN ISN	84		RETURN	00001150					
ISN	85	5010	FORMAT(' INPUT DATA')	00001160					
ISN	86	5020	FORMAT( BOA1 )	00001170					
ISN	87	5030	FORMAT(1X,80A1)	00001180					
ISN	88	5040	FORMATI/' *** INPHDR *** READ ERROR ON INPUT DATA CARDS ',	00001190					
			. ' IRETN= ',I3/)	00001200					
ISN	89	5050	FORMAT(/' *** INPHOR *** EOF ON INPUT DATA CARDS ',	00001210					
		5	. ' IRETN= ',I3/)	00001220					
ISN	90	5060	FORMAT(A6,1X,2(15,1X),8(11,1X))	00001230					
ISN	91	5065		00001240					
			. 'ISTART= ',18,' ISTOP= ',18,' IRETN= ',13/)	00001250					
ISN	92	5070	FORMAT(/' *** INPHOR *** DUMP OPTION IS NOT VALID. ',	00001260					
			. 'OPTION(',11,')= ',12)	00001270					
ISN	93	5075		00001290					
ISN	94	5080	. 'LENGTH= ',16,' IRETN= ',13/)	00001300					
TCN	95	5000	FORMAT(/' *** INPHDR *** TAPE DOES NOT HAVE A SAMS HEADER ',	00001310					
ISN	75	5070	. 'RECORD IRETN= ',I3/)	00001320					
ISN	96	E100	FORMAT(/' *** INPHDR *** TAPE IS NOT A TEMPERATURE OR ',	00001330					
TOM	70	5100	. 'COMPOSITION TAPE IRETN= ',I3/)	00001340					
ISN	97	5110	FORMAT(//' A SAMS ',A11,' TAPE WILL BE DUMPED STARTING WITH ',	00001350					
2000			. 'DAY ',13,' 19',12,' AND ENDING WITH DAY ',13,' 19',12)	00001360					
ISN	98	5120	FORMAT( ' THE FOLLOWING ', All, ' TAPE DUMP OPTIONS WERE ',	00001370					
	~		'SELECTED: '/' RECORD', 18X, 'DUMP', 5X, 'ACTION TO'/' TYPE',	00001380					
			18X, 'OPTION',4X, 'BE TAKEN'/1X,2A3,17X,2A3,4X,3A3)	00001390					
		E170	FORMAT(2X, '7400 HEADER RECORDS', 6X, I1, 6X, A11/	00001400					
ISN	99	2720							
ISN	99	9130	. 2X, '7402 RETRIEVED TEMP', 6X, 11, 6X, 411/	00001410					
ISN	99	9130							

5140 FORMAT(2X,'7405 N20 MIXING RATIO',4X,I1,6X,A11/ . 2X,'7405 N20 RATIO ERRORS',4X,I1,6X,A11/ . 2X,'7406 CH4 MIXING RATIO',4X,I1,6X,A11/ . 2X,'7406 CH4 MIXING RATIO',4X,I1,6X,A11/ . 2X,'7406 CH4 RATIO ERRORS',4X,I1,6X,A11) 5150 FORMAT(/' #MM INPHOR ### UNEXPECTED EOF ON TAPE -- IRETN= ',I3/) 5160 FORMAT(/' #MM INPHOR ### READ ERROR ON TAPE -- IRETN= ',I3/) ISN ISN ISN END ISN 

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ISN	1	SUBROUTINE GETGRT(ISTART, ISTOP, FILNUM, OPTION, IRETN)	000000
		C SUBROUTINE - GETGRT	000000 <del>×</del> 000000
			000000
		C C PURPOSE - READ RECORDS FROM A SAMS TEMPERATURE TAPE AND DUMP THOSE	000000
			000000
		C THAT FALL BETWEEN THE START AND STOP DATES.	000000
			000000
		C ARGUMENTS PASSED -	000000
		C ISTART = START DATE (YYDDD) FROM INPUT CARD	
		C ISTOP = STOP DATE (YYDDD) FROM INPUT CARD	000001
		C FILNUM = TAPE FILE NUMBER	000001
		C OPTION = DUMP OPTIONS FROM INPUT CARD (1=DUMP, 0=NO DUMP)	000001
		C IRETN = ERROR RETURN CODE (IRETN=0 NO ERROR, >0 ERROR)	000001
		C	000001
		C SUBPROGRAMS CALLED - POSN,GETREC,DM7402,DM7403	000001
		C	000001
		C CALLED BY - MAIN	000001
		C	000001
		C PROGRAMMER - TOM NUTTER 03/05/85	000001
		C*************************************	<b>*000</b> 002
ISN	2	INTEGER ISTART, ISTOP, FILNUM, OPTION(8), IRETN, FILDAT	000002
ISN	3	INTEGER*2 IBLOCK(2994),IYEAR,IDAY	000002
ISN	4	LOGICAL EOF,EOV	000002
		C*** STATEMENT FUNCTION	000002
ISN	5	FILDAT(IYEAR, IDAY)=(IYEAR-(IYEAR/100)*100)*1000+IDAY	000002
		C*** INITIALIZE END-DF-VOLUME FLAG	000002
[SN	6	EQY=.FALSE.	000002
	•	C*** SEARCH THROUGH DATA FILES UNTIL FILE START DATE MATCHES	000002
		C*** THE REQUESTED DUMP START DATE	000002
(SN	7	100 CONTINUE	000003
	'	C*** RESET RECORD COUNTER, END-OF-FILE FLAG, INCREMENT FILE COUNT,	000003
		CHAR RESET RECORD COUNTER, END-OF-FILE FLAS, INCREMENT FILE COUNTY	000003
ISN	•	KNTREC=0	000003
	8 9		000003
ISN	-	EOF=.FALSE.	
ISN	10	FILNUM=FILNUM+1	000003
SN	11	CALL POSN(1,10,FILNUM)	000003
ISN	12	CALL GETREC(10,KNTREC, IBLOCK, LEN, NHRD, EOV, EOF, IRETN)	000003
ISN	13	IF (IRETN.GT.O) THEN	000003
ISN	14	IRETN=10	000003
(SN	15	WRITE(6,5010) IRETN	000004
ISN	16	RETURN	000004
(SN	17	ENDIF	000004
		CHAN CHECK IF END-OF-VOLUME OCCURRED BEFORE ANY RECORDS WERE DUMPED	000004
ISN	18	IF (EOV) THEN	000004
(SN	19	IRETN=11	000004
ISN	20	WRITE(6,5020) IRETN	000004
ISN	21	RETURN	000004
ISN	22	ENDIF	000004
		C*** VERIFY THAT FIRST RECORD IS A HEADER RECORD (TYPE 7400)	000004
ISN	23	IF (IBLOCK(3).NE.7400) THEN	000005
SN	24	IRETN=12	000005
SN	25	WRITE(6,5030) IBLOCK(3), IRETN	000005
SN	26	RETURN	000005
ISN	27	ENDIF	000005
	<b>~</b> '	CHECK IF FILE START DATE MEETS REQUESTED DUMP LIMITS	000005
SN	28	IF (FILDAT(IBLOCK(5),IBLOCK(6)).GT.ISTOP) THEN	000005
			000005
SN	29	IRETN=15 WEITE(4 EDED) TETART TETOR ETIDAT(TRIDEV(E), TRIDEV(A)), TRETN	
SN	30	WRITE(6,5050) ISTART, ISTOP, FILDAT(IBLOCK(5), IBLOCK(6)), IRET	
SN	31	RETURN	000005
SN	32	ENDIF	000006
		C*** GET NEXT FILE HEADER RECORD	000006
SN	33	IF (FILDAT(IBLOCK(5), IBLOCK(6)).LT.ISTART) GO TO 100	000006
		C*** THE FILE WITH THE REQUESTED START DATE HAS BEEN FOUND. DUMP THE	000006
		C*** HEADER RECORD	000006
SN	34	WRITE(6,5040) IBLOCK(3),IBLOCK(0),IBLOCK(5)	000006
		C*** READ THROUGH RECORDS AND DUMP THOSE FOR WHICH THE DUMP	000006
		C*** OPTION IS SET TO 1	000006
SN	35	300 CONTINUE	000006
SN	36	CALL GETREC(10,KNTREC, IBLOCK, LEN, NARD, EOV, EOF, IRETN)	000006

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		<b>*</b>	.*1	.*8		
ISN	38		IRETN=13	00000710		
ISN	39		WRITE(6,5010) IRETN	00000720		
ISN	40		RETLIN	00000730		
ISN	40		ENDIF			
1.34	74	CXXX		00000740 00000750		
ISN	42	•	IF (EOF) THEN			
ISN	43		KNTREC=0	00000760 00000770		
ISN	44		EOF=.FALSE.	00000780		
ISN	45		FILNUM=FILNUM+1	00000790		
ISN	46		CALL POSN(1,10,FILNUM)	00000800		
ISN	47		CALL GETREC(10,KNTREC, IBLOCK, LEN, NWRD, EOV, EOF, IRETN)	00000810		
ISN	48		IF (IRETN.GT.O) THEN	00000820		
ISN	49		IRETN=14	00000830		
ISN	50		WRITE(6,5010) IRETN	00000840		
ISN	51		RETURN	00000850		
ISN	52		ENDIF	00000860		
ISN	53		ENDIF	00000870		
		C×××	AT END-OF-VOLUME ALL DATA HAS BEEN DUMPED	00000880		
ISN	54		IF (EOV) RETURN	00000890		
		CXXX	GET FILE DATE, DUMP HEADER RECORD OR RETURN IF STOP DATE	00000900		
		C <del>xxx</del>		00000910		
ISN	56		IF (IBLOCK(3).EQ.7400) THEN	00000920		
ISN	57		IF (FILDAT(IBLOCK(5), IBLOCK(6)).GT.ISTOP) THEN	00000930		
ISN	58		RETURN	00000940		
ISN	59		ELSE	00000950		
ISN	60		WRITE(6,5040) IBLOCK(3),IBLOCK(6),IBLOCK(5)	00000960		
ISN	61		ENDIF	00000970		
ISN	62		ENDIF	00000980		
		CXXX	DUMP A TYPE 7402 RECORD	00000990		
ISN	63		IF (IBLOCK(3).EQ.7402) THEN	00001000		
ISN	64		IF (OPTION(2).EQ.1) CALL DM7402(IBLOCK)	00001010		
ISN	66		ENDIF	00001020		
		C×××	DUMP A TYPE 7403 RECORD	00001030		
ISN	67		IF (IBLOCK(3).EQ.7403) THEN	00001040		
ISN	68		IF (OPTION(3).EQ.1.AND.IBLOCK(11).EQ.2) CALL DM7403(IBLOCK,2			
ISN	70		<pre>IF (OPTION(4).EQ.1.AND.IBLOCK(11).EQ.102) CALL DM7403(IBLOCK</pre>	,00001060		
	_		102)	00001070		
ISN	72		ENDIF	00001080		
		C×××	GET NEXT RECORD	00001090		
ISN	73		GO TO 300	00001100		
ISN	74	5010	FORMAT(/' *** GETGRT *** ERROR FOUND IN GETREC IRETN= ',I3)	00001110		
ISN	75	5020	FORMAT(/' *** GETGRT *** EOV OCCURRED BEFORE ANY RECORDS ', . 'WERE DUMPED IRETN= ',I3)	00001120 00001130		
ISN	76	5030	<pre>FORMAT(/' *** GETGRT *** FIRST RECORD IN DATA FILE IS ICORRECT ',</pre>	00001140		
ISN	77	5040	FORMAT(///' RECORD TYPE ',14,' DATA DAY/YEAR ',13,1X,14)	00001160		
ISN	78	5050	FORMAT(/' *** GETGRT *** THE TAPE START DATE EXCEEDS THE ',	00001170		
			. 'REQUESTED DUMP STOP DATE'/' CHECK DUMP DATES AGAINST TAPE ',			
			'HEADER DATES. ISTART= ',16,' ISTOP= ',16,' TAPE DATE= ',16,	00001190		
			. 'IRETN= ',I3)	00001200		
ISN	79		END	00001210		

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ISN	1	SUBROUTINE GETCHP(ISTART, ISTOP, FILNUM, OPTION, IRETN)	000000
		Cx <del>xxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxxx</del>	+000000
		C SUBROUTINE - GETCHP	000000
		C	000000
		C PURPOSE - READ RECORDS FROM A SAMS COMPOSITION TAPE AND DUMP THOSE	000000
		C THAT FALL BETWEEN THE START AND STOP DATES.	000000
			000000
			000000
		C ARGUMENTS PASSED -	
		C ISTART = START DATE (YYDDD) FROM INPUT CARD	00000
		C ISTOP = STOP DATE (YYDDD) FROM INPUT CARD	00000
		C FILNUM = TAPE FILE NUMBER	00000
		C OPTION = DUMP OPTIONS FROM INPUT CARD (1=DUMP, 0=NO DUMP)	00000
		C IRETN = ERROR RETURN CODE (IRETN=0 NO EPROR, >0 ERROR)	00000
			00000
		C SUBPROGRAMS CALLED - POSN, GETREC, DMPCMP	00000
			00000
		C	
		C CALLED BY - MAIN	00000
		C	00000
		C PROGRAMMER - TOM NUTTER 03/05/85	00000
		C*************************************	+00000;
TCM	2	INTEGER ISTART, ISTOP, FILNUM, OPTION(8), IRETN, RECDAT	00000
ISN	2		00000
ISN	3	INTEGER*2 IBLOCK(2994),IYEAR,IDAY	
ISN	4	LOGICAL EDF, EOV	00000
		C*** STATEMENT FUNCTION	00000
ISN	5	RECDAT(IYEAR,IDAY)=(IYEAR-(IYEAR/100)*100)*1000+IDAY	00000
		C*** INITIALIZE RECORD COUNTER, END-OF-FILE FLAG, END-OF-VOLUME FLAG,	00000
		C*** INCREMENT FILE COUNT, POSITION TO NEXT FILE	00000
ISN	6	KNTREC=0	00000
ISN	7	EOF=.FALSE.	00000
ISN	8	EOV=.FALSE.	00000
ISN	9	FILNUM=FILNUM+1	00000
ISN	10	CALL POSN(1,10,FILNUM)	00000
		C*** SEARCH THROUGH DATA RECORDS UNTIL RECORD START DATE MATCHES	00000
		C*** THE REQUESTED INPUT START DATE	00000
ISN	11	100 CONTINUE	00000
			00000
ISN	12	CALL GETREC(10,KNTREC, IBLOCK, LEN, NWRD, EOV, EOF, IRETN)	
ISN	13	IF (IRETN.GT.O) THEN	00000
ISN	14	IRETN=50	00000
ISN	15	WRITE(6,5010) IRETN	00000
ISN	16	RETURN	00000
ISN	17	ENDIF	00000
		C*** CHECK IF END-OF-VOLUME OCCURRED BEFORE ANY RECORDS WERE DUMPED	00000
TCN	10		00000
ISN	18	IF (EOV) THEN	
ISN	19	IRETN=51	00000
ISN	20	WRITE(6,5020) IRETN	00000
ISN	21	RETURN	00000
ISN	22	ENDIF	00000
		C*** POSITION TO NEXT FILE IF AN END-OF-FILE IS ENCOUNTERED	00000
ISN	23	IF (EOF) THEN	00000
			00000
ISN	24	KNTREC=0	
ISN	25	EOF=.FALSE.	00000
ISN	26	FILNUM=FILNUM+1	00000
ISN	27	CALL POSN(1,10,FILNUM)	00000
ISN	28	ENDIF	00000
		C*** CHECK IF FILE START DATE MEETS REQUESTED DUMP LIMITS	00000
TCN	29	IF (RECDAT(IBLOCK(5),IBLOCK(4)).GT.ISTOP) THEN	00000
ISN			00000
ISN	30		
ISN	31	WRITE(6,5030) ISTART,ISTOP,RECDAT(IBLOCK(5),IBLOCK(4)),IRETN	
ISN	32	RETURN	00000
ISN	33	ENDIF	00000
		C*** GET NEXT RECORD	00000
TCN	34	IF (RECDAT(IBLOCK(5),IBLOCK(4)).LT.ISTART) GO TO 100	00000
ISN	24		00000
		C*** THE RECORD START DATE NOW FALLS WITHIN THE REQUESTED DUMP LIMITS.	
		C*** READ AND DUMP RECORDS UNTIL REQUESTED STOP DATE IS EXCEEDED	00000
ISN	35	CALL DMPCMP(IBLOCK, OPTION)	00000
ISN	36	200 CONTINUE	00000
ISN	37	CALL GETREC(10,KNTREC, IBLOCK, LEN, NWRD, EOV, EOF, IRETN)	00000
ISN	38	IF (IRETN.GT.O) THEN	00000
4 J I I		IRETN=52	00000
7.0.1			
ISN ISN	39 40	WRITE(6,5010) IRETN	00000

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		**1	*8
ISN	41	RETURN	00000710
ISN	42	ENDIF	00000720
1.041		CHAN POSITION TO NEXT FILE IF AN END-OF-FILE IS ENCOUNTERED,	00000730
		CHAR READ THE NEXT RECORD	00000740
ISN	43	IF (EOF) THEN	00000750
ISN	44	KNTREC=0	00000760
ISN	45	EOF=.FALSE.	00000770
ISN	46	FILNUM=FILNUM+1	00000780
ISN	47	CALL POSN(1,10,FILNUM)	00000790
ISN	48	CALL GETREC(10,KNTREC, IBLOCK, LEN, NMRD, EOV, EOF, IRETN)	00800000
ISN	49	IF (IRETN.GT.0) THEN	00000810
ISN	50	IRETN=53	00000820
ISN	51	WRITE(6,5010) IRETN	00000830
ISN	52	RETURN	00000840
ISN	53	ENDIF	00000850
ISN	54	ENDIF	00000860
		CXXX AT END-OF-VOLUME, OR WHEN RECORD DATE EXCEEDS DUMP STOP DATE,	00000870
		CXXX ALL DATA HAS BEEN DUMPED	00000880
ISN	55	IF (EOV) RETURN	00000890
ISN	57	IF (RECDAT(IBLOCK(5),IBLOCK(4)).GT.ISTOP) RETURN	00000900
		CXXX DUMP A RECORD	00000910
ISN	59	CALL DMPCMP(IBLOCK, OPTION)	00000920
		CHAH, GET NEXT RECORD	00000930
ISN	60	GO TO 200	00000940
ISN	61	5010 FORMAT(/' *** GETCMP *** ERROR FOUND IN GETREC IRETN= ',13)	00000950
ISN	62	5020 FORMATI' *** GETCMP *** EOV OCCURRED BEFORE ANY RECORDS ',	00000960
		. 'WERE DUMPED IRETN# ',I3)	00000970
ISN	63	5030 FORMAT(/' *** GETCMP *** THE TYPE START DATE EXCEEDS THE ',	00000980
		. 'REQUESTED DUMP STOP DATE'/' CHECK DUMP DATES AGAINST TAPE ',	00000990
		'HEADER DATES. ISTART= ',16,' ISTOP= ',16,' TAPE DATE= ',16,	00001000
		. ' IRETN= ',I3)	00001010
ISN	64	END	00001020

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		**1	·. <b>#</b> 8
			00001220
ISN	1	SUBROUTINE WETRECIJUNIT, KNTREC, IBLOCK, LENGTH, NHORDS, EOV, EOF,	00000010
	-	. IRETN)	00000020
		C <del>************************************</del>	H##00000030
		C SUBROUTINE - GETREC	00000040
		C	00000050
		C PURPOSE - READ A RECORD FROM A SPECIFIED LOGICAL UNIT, INCREMENT	00000060
		C RECORD COUNTER, AND CALCULATE THE NUMBER OF 32-BIT WORDS IN THIS	00000070 00000080
		C RECORD. DETERMINE END-OF-FILE AND END-OF-VOLUME CONDITIONS C AS THEY OCCUR.	00000090
		C AS THET OCCOR.	00000100
		C ARGUMENTS PASSED -	00000110
		C IUNIT = LOGICAL UNIT NUMBER	00000120
		C KNTREC = RECORDS PER FILE	00000130
		C IBLOCK = BUFFER INTO WHICH RECORD IS READ	00000140
		C LENGTH = LENGTH OF RECORD IN BYTFS	00000150
		C NHORDS = NUMBER OF 32-BIT HORDS IN RECORD	00000160
		C EOV = END-OF-VOLUME FLAG	00000170
		C EOF = END-OF-FILE FLAG	00000180
		C IRETN * ERROR RETURN CODE (IRETN=0 NO ERROR, >0 ERROR)	00000190
			00000200
		C SUBPROGRAMS CALLED - FMOVE, FREAD	00000210
		C CALLER BY A CETCHE CETCHE	00000220 00000230
		C CALLED BY - GETGRT,GETCMP C	00000240
		C PROGRAMMER - TOM NUTTER 03/05/85	00000250
ISN	2	LOGICAL EOV,EOF	00000270
ISN	3	INTEGER IBLOCK(1497)	00000280
	-	CHAN INITIALIZE BUFFER TO O	00000290
ISN	4	IBLOCK(1)=0	00000300
ISN	5	CALL FMOVE(IBLOCK(2),5984,IBLOCK(1))	00000310
		CHAN READ A RECORD, INCREMENT RECORD COUNTER AND CALCULATE THE	00000320
		CANN NUMBER OF 32-BIT HORDS IN THIS RECORD	00000330
ISN	6	CALL FREAD(IBLOCK(1), IUNIT, LENGTH, *800, *930)	00000340
ISN	7	KNTREC=KNTREC+1	00000350
ISN	8	IF (LENGTH.EQ.40 .OR. LENGTH.EQ.3506 .OR. LENGTH.EQ.4882 .OR.	00000360
TON	•	LENGTH.EQ.5986) THEN	00000370
ISN	9 10	NHORDS=LENGTH/4 IF (MOD(LENGTH,4).GT.0) NHORDS=NHORDS+1	00000380 00000390
ISN ISN	10	RETURN	00000400
ISN	12	ELSE	00000410
ISN	14	IRETN=40	00000420
ISN	15	WRITE(6,5010) LENGTH, IRETN	00000430
ISN	16	RETURN	00000440
ISN	17	ENDIF	00000450
		C*** SET END-OF-FILE FLAG AND END-OF-VOLUME FLAG	00000460
ISN	18	800 CONTINUE	00000470
ISN	19	EOF=.TRUE.	00000480
ISN	20	IF (KNTREC.LT.1) EOV=.TRUE.	00000490
ISN	22	RETURN	00000500
		C*** READ ERROR	00000510
ISN	23	930 CONTINUE	00000520
ISN	24	IRETN=41	00000530
ISN	25	WRITE(6,5020) IUNIT, IRETN	00000540
ISN	26	RETURN	00000550
ISN	27	5010 FORMAT(/' *** GETREC *** INVALID RECORD LENGTH OF ',	00000560
TON		. 16, 'BYTES. IRETN= ',13/)	00000570
ISN	28	5020 FORMAT(/' *** GETREC *** READ ERROR ON UNIT ',13,	00000580
ISN	29	. ' IRETN= ',I3/)	00000590 00000600
194	47	END	00000000

ISN         1         SLERGUTTME (BTFAG2) ISLOCK )         00000120           C++++++++++++++++++++++++++++++++++++			**1	′. <b>*8</b>
SUBRCUTINE - 947402         00000050           C         00000060           C         PURPOSE - FORM AND PRINT OUTPUT LINES FOR SAMS GRIDDED RETRIEVED         00000050           C         TEMPERATURE TYPE 7402 RECORDS.         00000050           C         RARUMENTS PASSED -         00000050           C         IBLOCK = ARRAY CONSISTING OF GRIDDED RETRIEVED TEMPERATURE DATA         00000120           C         SUBPROGRAMS CALLED - NONE         00000120           C         CALLED BY - GETGRT         00000120           C         CALLED BY - GETGRT         00000120           C         CHARACTERMA F - TOM NUTTER 05/02/85         00000150           C         CHARACTERMA F A TSH.UK(2994),LINE(19)         00000180           ISN         S         DATA LONGY'-180','-170','-160','-150','-140','-130','-120', 00000210           .         '-10','-10','', '0','-40','-0',' 40', '00000210         .'-10','-10','','',''10','+10','+10','+10','+120', 00000220           .         '+30','+10','',''10','+1	ISN	1		
C PURPOSE - FORM AND PRINT OUTPUT LINES FOR SAMS GRIDDED RETRIEVED 00000050 C TEMPERATURE TYPE 7402 RECORDS. 00000050 C ARGUMENTS PASSED - 00000050 C ARGUMENTS PASSED - 00000010 C SUBPROGRAMS CALLED - NONE 0000010 C SUBPROGRAMS CALLED - NONE 0000010 C CALLED BY - GETGRT 00000150 C HARACTER*4 DASM./LINE19) 00000210 . '-10', '10', '10', '10', '10', '10', '10', '10', '10', '00', '10',			÷	
C PURPOSE - FORM AND PRINT CUTPUT LINES FOR SAMS GRIDDED RETRIEVED 00000050 C THEPERATURE TYPE 7-02 RECORDS. 00000070 C ARCUMENTS PASSED - 00000070 C IBLOCK = ARRAY CONSISTING OF GRIDDED RETRIEVED TEMPERATURE DATA 0000010 C SUBPROGRAMS CALLED - NONE 0000010 C SUBPROGRAMS CALLED - NONE 0000010 C CALLED BY - GETGRT 00000120 C C CALLED BY - GETGRT 00000120 C C CALLED BY - GETGRT 00000150 C POGRAMMER - TOM NUTTER 05/02/25 00000150 C HAMASTANAMANANANANANANANANANANANANANANANANANA				
C TEMPERATURE TYPE 7402 RECORDS. 00000070 C ARGUMENTS PASSED - 00000070 C TBLOCK = ARRAY CONSISTING OF GRIDDED RETRIEVED TEMPERATURE DATA 0000010 C SUBPROGRAMS CALLED - NONE 0000010 C CALLED BY - GETGRT 0000010 C CHARACTER*6 DASH, LCNG(38), HEDLIN(19) 00000170 ISN 5 DATA LCNG/'-150', '-120', '-120', '-120', '-120', 00000210 - '-110', '-100', '-90', '-30', '-70', '-40', '-130', '-120', 00000210 - '-110', '-10', '-90', '-90', '-10', '-40', '-130', '-120', 00000220 - '-130', '-20', '-10', '-90', '-10', '-40', '-130', '-120', 00000220 - '-130', '-20', '-10', '+90', '+10', '+10', '+10', '+10', '00000220 - '-130', '-10', '+00', '+10', '+10', '+10', '+10', '+10', '00000220 - '-130', '-10', '-90', '-90', '-90', '-90', '-10', '-0', 00000220 - '-130', '-10', '+00', '+10', '+10', '+10', '+10', '+10', '00000250 - CHAR HATTE (6,500) IBLOCK(I),I,I-1,8) 00000250 C CHAR HATTE (6,500) IBLOCK(I),I,I-1,8) 00000350 C CHAR HOND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND RRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND RRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUTPUT LINES FOR THO LCNGTUDE GROUPS -180E TO OE S C CHAR FORM AND PRINT QUTPUT LINES FOR THO LONGTUDE GROUPS -180E TO OE C CHAR FORM AND PRINT QUT				
C ARGUMENTS PASSED - 0000000 C ARGUMENTS PASSED - 00000010 C IBLOCK = ARRAY CONSISTING OF GRIDDED RETRIEVED TEMPERATURE DATA 0000010 C SUBPROGRAMS CALLED - NONE 00000110 C CALLED BY - GETGRT 00000110 C CALLED BY - GETGRT 00000130 C CREATER - TOM NUTTER 03/02/85 00000130 C PROGRAMMER - TOM NUTTER 03/02/85 00000130 C PROGRAMMER - TOM NUTTER 03/02/85 00000130 C NO000130 C THITEGER*2 IBLOCK(29%).LINE(19) 00000140 OU000100 ISN 3 CHARACTER*4 DASH.LONG(33).HEDLIN(19) 00000180 OU000100 . '-100','-100','-10','-10','-10','-10','-10','-10','-10', 00000210 . '-10','-10','-0','-10','-0','-6','-5','-40', 00000220 . '-10','-10','-10','-10','-10','-40','-10','+40', 00000220 . '-10','-10','-10','-10','-10','-10','+10','+20',' 40', 00000220 . '-10','-10','-10','-10','-10','-10','+10','+20',' 40', 00000220 . '-10','-10','-10','-10','+10','+20','+30','+40', 00000220 C CHEME FORM AND FRIMT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E 00000320 ISN 6 MRITE(6,500) IBLOCK(1),I=1,8) 00000250 C CHEME FORM AND FRIMT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E 00000320 ISN 9 IL0G=9 100G=9 00000320 ISN 10 IF (K.EQ.2) IL0G=1225 00000310 ISN 12 CALL FMOVE(HEDLIN(1),7,6,LONC'((K-1)+19+1)) 00000320 C CHEME FORM AND FRIMT OUTPUT LINES FOR 62 PRESSURES 00000320 ISN 14 D0 100 I=1,19 0000510 COMOS50 ISN 15 D0 200 J=1,62 00000510 ISN 16 IDFF=JAEDCK(IDFF) 0000040 ISN 17 D0 100 I=1,19 00000510 ISN 18 ID COMTINUE 0000050 ISN 18 ID COMTINUE 00000510 ISN 18 ID COMTINUE 00000510 ISN 18 ID COMTINUE 00000510 ISN 18 ID COMTINUE 00000510 ISN 18 ID COMTINUE 110/CK/IDFF) 0000040 ISN 22 00 COMTINUE 00000510 ISN 23 200 COMTINUE 122,5,%, TECE', 15,2X, 10000510 ISN 24 5000 FORMATI/' RETRIEVED TEMPERATURE GRID') 00000540 ISN 25 5000 FORMATI/' RETRIEVED TEMPERATURE GRID') 00000540 ISN 24 5000 FORMATI/' RETRIEVED TEMPERATURE GRID') 00000540 ISN 25 5000 FORMATI/' RETRIEVED TEMPERATURE GRID') 00000540 ISN 24 5000 FORMATI/' RETRIEVED TEMPERATURE GRI				
C ARGUMENTS PASSED - 00000090 C IBLOCK = ARAY CONSISTING OF GRIDED RETRIEVED TEMPERATURE DATA 00000100 C SUBPROGRAMS CALLED = NONE C SUBPROGRAMS CALLED = NONE C CALLED BY - GETGRT C 00000120 C CALLED BY - GETGRT C PROGRAMS CALLED - NONE C PROGRAMS CALLED - NONE C PROGRAMS CALLED - NONE C CALLED BY - GETGRT C 00000120 C CALLED BY - GETGRT C PROGRAMS CALLED - NONE C CHARACTERS LOCK(29%)LLINE(19) ISN 3 C CHARACTERS LOCK(29%)LLINE(19) ISN 4 DATA LONG'-180','-170','-160','-130','-120', 00000210 . '-130','-20','-00','-70','-40','-10','-10','-10','-10', '00000210 . '-150','+40','100','-10','-10','-10','+10','+10','+40', 00000220 C - '+50','+40','+70','+60','+170','+10','+10','+10','+120', 00000250 C - '+50','+40','+70','+60','+170','+10','+10','+10','+120', 00000250 C - '+50','+40','+70','+60','+170','+10','+10','+10','+120', 00000250 C C+++ FORM AND PRINT OUTPUT LINES FOR THO LONGTUDE CROUPS -180E TO 0E 00000280 C C+++ FORM AND PRINT OUTPUT LINES FOR THO LONGTUDE CROUPS -180E TO 0E 00000230 ISN 4 D 00 300 K=1,2 D 0000510 ISN 12 C CALL FROVE(HEDLIN(1),76,LONC'(K-1)+19+1)) 00000510 ISN 12 C CALL FROVE(HEDLIN(1),76,LONC'(K-1)+19+1)) 00000510 ISN 12 C CALL FROVE(HEDLIN(1),76,LONC'(K-1)+19+1)) 00000510 ISN 12 C CALL FROVE(HEDLIN(1),76,LONC'(K-1)+19+1)) 00000510 ISN 12 C CONTINUE 15N 12 D 0 100 I=1,19 00000510 ISN 12 C CONTINUE 15N 12 D 0 100 I=1,19 00000510 ISN 13 C D 0 00 J=1,62 00000510 ISN 14 PLEVEL=1.2 00000510 ISN 15 D 0 00 J=1,62 00000510 ISN 15 D 0 00 J=1,62 00000510 ISN 15 D 0 00 J=1,62 00000510 ISN 15 D 0 00 CONTINUE 15N 25 D00 FORMATIC' ERTISERDE TEMPERATURE GRID') 00000640 IS			· · · · · ·	
C IBLOCK = ARRAY CONSISTING OF GRIDDED RETRIEVED TEMPERATURE DATA 00000090 C SUBPROGRAMS CALLED - NONE 00000110 C CALLED BY - GETGRT 00000130 C CALLED BY - GETGRT 00000130 C CALLED BY - GETGRT 00000130 C CHARACTER*6 DASH, LONG (30, HED LIN (19) 00000140 C PROGRAMMER - TOM MUTTER 03/02/85 00000150 C PROGRAMMER - TOM MUTTER 03/02/85 00000150 ISN 3 CHARACTER*6 DASH, LONG (30, HED LIN (19) 00000100 ISN 4 DATA DASH, '0', '0', '60', '-130', '-120', 00000210 . '-110', '-100', '-10', '-160', '-150', '-10', '-120', 00000210 . '-10', '-10', '-0', '-0', '-60', '-50', '-0', '00000210 . '-10', '-10', '-10', '-10', '-10', '-10', '+10', '+20', '+00', 00000220 . '-10', '+10', '+10', '+0', '+0', '+0', '+0', '+0', 00000220 . '+150', '+60', '+10', '+10', '+10', '+10', '+120', 00000250 C H## FITE (6,500) ISN 6 MRITE(6,500) (IBLOCK(I),I=1,8) 00000250 C H## FORM AND PRINT OUTPUT LINES FOR THO LONGTUDE GROUPS -180E TO 00 C H## FORM AND PRINT OUTPUT LINES FOR THO LONGTUDE GROUPS -180E TO 00 C H## FORM AND PRINT OUTPUT LINES FOR THO LONGTUDE GROUPS -180E TO 00 C CH## PRINT HEADING 00000320 ISN 8 DO 300 K-1,2 ISN 10 IF (K.EQ.2) ILOG=1225 00000310 ISN 12 C LALL FMOVE (HEDLIN (1),76,LONC'(K-1)+19+1)) 00000350 C CH## FORM AND PRINT OUTPUT LINES FOR C6 2 PRESSURES 00000350 ISN 13 MRITE(6,5020) HEDLIN, (DASH,I=1,21) 00000350 C H## FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES 00000350 ISN 14 PLEVEL=1,2 ISN 15 DO 100 I=1,19 100 CONTINUE 0000040 ISN 16 I DFF=J+ILOG 15N 17 DO 100 I=1,19 100 CONTINUE 0000040 ISN 16 I OFF=J+ILOK 0000041 ISN 17 DO 100 I=1,19 10000440 ISN 16 I DFF=J+EUEVEL,LINE 00000450 ISN 16 I DFF=J+EUEVEL,LINE 00000450 ISN 16 I DFF=J+ILOK 0000450 ISN 17 DO 100 I=1,19 15N 22 OR CONTINUE 0000040 ISN 22 DO CONTINUE 0000040 ISN 23 DO CONTINUE 0000040 ISN 24 S000 CONTINUE 0000450 ISN 25 OR FORMATI/' RETRIEVE TEMPERATURE GRID') 00000450 ISN 25 OR FORMATI/' RETRIEVE TEMPERATURE GRID') 00000450 ISN 25 OFFORMATI/' RETRIEVE TEMPERATURE GRID') 00000450 ISN 26 S000 FORMATI/' RETRIEVE TEM				0800000
C SUBPROGRAMS CALLED - NONE 00000100 C SUBPROGRAMS CALLED - NONE 00000120 C CALLED BY - GETGRT 00000120 C CALLED BY - GETGRT 00000130 C PROGRAMMER - TOM MUTTER 03/02/05 C MARACTER* DASH.LAWELTSTATESTATESTATESTATESTATESTATESTATEST				00000090
C SUBPROGRAMS CALLED - NONE 00000120 C CALLED BY - GETGRT 00000130 C CALLED BY - GETGRT 00000130 C PROGRAFMER - TOM NUTTER 03/02/85 00000150 C************************************				00000100
C CALLED BY - GETGRT         00000130           C         0000140           C         PROGRAIMER - TOM NUTTER 03/02/85         00000150           C************************************				00000110
C         D0000140         D0000150           C         PROGRAPHER - TOM NUTTER 03/02/85         00000150           INTEGER*2         INTEGER*2         D0001150           ISN         2         INTEGER*2         D0001160           ISN         3         CHARACTER*4         DASSN/LONG(39),LINE(19)         00000180           ISN         4         DATA DASN/''         00000120			c	00000120
$\begin{tabular}{ c c c c c c c c c c c c c c c c c c c$			C CALLED BY - GETGRT	00000130
CHARGE CONSTRUCT           ISN 2         INTEGER*2 IBLOCK(2994).LIN(19)         00000120           ISN 3         CHARACTER*4 DASH,LONG(39).HEDLIN(19)         00000120           ISN 4         DATA DASH//         00000120           ISN 5         DATA LONG('180','-170','-160','-150','-140','-130','-120', 00000220         .'-110','-100','-90','-40', '-40', '-40', 00000220           .'-10','-10','-0','-10','-10','+0','+20','+20','+40', 00000220         .'+50','+60','+70','+80','+90','+100','+110','+120', 00000230           .'+130','+140','+150','+160','+170','+190','+200'/         00000250           ISN 6         MRITE(6,5000)         00000220           ISN 7         MRITE(6,5001) (IBLOCK(1),I=1,8)         00000220           CH+# FORM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E         00000230           ISN 7         MRITE(6,5020) IGD=1225         00000310           ISN 9         ILOG=9         00000320           ISN 10         IF (K.GG.2) ILOG=1225         00000320           CH+# FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000350           ISN 12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)+19+1))         00000350           ISN 13         HRITE(6,5020) HEDLIN,(DASH,I=1,21)         00000350           ISN 14         PLEVEL+1.2         00000350           ISN 15         DO 200			c	00000140
ISN         2         INTEGER#2 IBLOCK (29%), LINE [19)         0000170           ISN         3         CHARACTER#4 DASH, LONG(38), HEDLIN(19)         0000180           ISN         4         DATA DASH/''/         00000190           ISN         5         DATA LONG'-180', '-170', '-160', '-150', '-140', '-120', 00000200         .           .         .'10', '-100', '-00', '-0', '-40', '-50', '-40', 00000220         .'-50', '-20', '-10', '-0', '+10', '+20', '+30', '+40', 00000220           .         .'+50', '+60', '+70', '+40', '+90', '+100', '+20', '+30', '+40', 00000220         .'+150', '+40', '+150', '+10', '+19', '+20', '           .         .'+50', '+60', '+70', '+40', '+10', '+10', '+20', '         00000220           .         .'+150', '+160', '+150', '+10', '+19', '+20', '         00000220           .         .'+150', '+160', '+150', '+10', '+19', '10', '+20', '         00000220           .         .'+150', '+160', '+150', '+170', '+190', '+200'/         00000220           .         .'+150', '+160', '+150', '+170', '+190', '+200'/         00000220           .         .         .'+150', '+160', '+170', '+190', '+200'/         00000220           .         .         .'+150', '+150', '+160', '+170', '+190', '+200'/         00000220           .         .         .'+150', '-150', '-160', '+170', '+190', '+200'/         00000220			÷ • • • • • • • • • • • • • • • • • • •	
ISN         G         CHARACTER*& DASH, LONG(38), HEDLIN(19)         00000180           ISN         4         DATA DASH/''/         00000190           ISN         5         DATA LONG/'-180', '-170', '-160', '-130', '-120', 00000210         .           .         '-110', '-100', '-90', '-60', '-70', '-60', '-50', '-40', 00000220         .         .           .         '-50', '-20', '-10', '', '0', '+10', '+20', '+30', '+40', 00000230         .         .           .         '+50', '+60', '+70', '+40', '+100', '+10', '+20', '         00000230           .         '+50', '+60', '+70', '+60', '+10', '+10', '+20', '         00000230           .         '+510', '+160', '+160', '+170', '+190', '+200'/         00000220           .         '+510', '+160', '+150', '+10', '+190', '+200'/         00000220           .         WRITE(6,5010) (IBLOCK(I),I=1,8)         00000220           .         MRITE(6,5010) IBLOCK(I),I=1,8)         00000220           .         S         00000220         00000230           .         ILGE=9         00000220         00000220           .         MRITE(6,5020) IEDL LINES FOR THO LONGTUDE GROUPS -180E TO 0E         00000220           .         ILGE=9         00000310         0000320           .         ISN         10         IF			-	
ISN         4         DATA DASH/''/         0000190           ISN         5         DATA LONG/'-180', '-170', '-160', '-130', '-120', 0000020         0000190           . '-110', '-100', '-90', '-80', '-70', '-60', '-50', '-40', 00000210         . '-30', '-20', '-10', '0', '+10', '+20', '40', 00000220           . '-50', '-60', '-70', '+80', '+70', '+20', '40', 00000220         . '+50', '+60', '+70', '+90', '+20', '40', '+120', 00000240           CH## HTTE RECORD HEADING         00000260           ISN         6         MRITE(6,5010) (IBLOCK(I),I=1,8)           CH## FORM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E         00000280           CH## AND +10E TO +170E, PLUS ZONAL MEAN AND CLIMATOLOGY         00000320           ISN         8         D0 300 K=1,2         00000330           ISN         9         IL0G=9         00000330           ISN         10         IF (K.EQ.2) IL0G=1225         00000330           CH## PRINT HEADING         00000340         00000340           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)+19+1))         00000350           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)+19+1)         00000350           ISN         14         PLEVEL=1,2         00000350           ISN         15         DO 200 J=1,62         000000350				
ISN         5         DATA LONG/'-180','-170','-160','-150','-140','-120',         00000200           '-110','-100','-90','-90','-70','-60','-50','-40',         00000210           '-30','-20','-10','','','10','+60','+40',         00000220           '-50','+60','+70','+80','+100','+210','+40',         00000220           '+50','+60','+70','+80','+100','+110','+120',         00000240           CHMM HRITE RECORD HEADING         00000220           ISN         6         HRITE(6,5000)           ISN         7         HRITE(6,5000)           ISN         7         HRITE(6,5000)           ISN         7         HRITE(6,5000)           ISN         7         HRITE(6,5000)           ISN         8         D0 300 K=1,2           00000220         CHMM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E           ISN         9         IL0G=9           ISN         8         D0 300 K=1,2           00000320         GO000320           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)*19+1))           00000320         GO000320           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)*19+1))           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)*19+1))           ISN				
<ul> <li>'-110','-100','-90','-80','-70','-60','-50','-60',' 0000210</li> <li>'-30','-20','-10',' 0','+10','+20','+30','+40', 0000220</li> <li>'+50','+60','+70','+80','+20','+20','+30','+40', 0000220</li> <li>'+130','+140','+150','+150','+190','+20',' 0000230</li> <li>'+130','+140','+150','+150','+190','+200',' 0000250</li> <li>CHH HAITE RECORD HEADING</li> <li>WRITE(6,500) (IBLOCK(1),I=1,8)</li> <li>O0000260</li> <li>CHH HAITE RECORD HAND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E</li> <li>CHH AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E</li> <li>CHH AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E</li> <li>CHH AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E</li> <li>CHH AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E</li> <li>CHH FORM AND PRINT OUTPUT LINES FOR CA PRESSURES</li> <li>O0000350</li> <li>ISN 12</li> <li>CALL FMOVE(HEDLIN(1),76,LONC'(K-1)+19+1))</li> <li>O0000350</li> <li>ISN 14</li> <li>PLEVEL=1.2</li> <li>O0000350</li> <li>ISN 15</li> <li>DO 200 J=1,62</li> <li>O0000360</li> <li>ISN 16</li> <li>IDFF=JOFF+64</li> <li>O0000420</li> <li>ISN 18</li> <li>LINE(1)=IBLOCK(IDFF)</li> <li>O0000420</li> <li>ISN 20</li> <li>IDO CONTINUE</li> <li>O0000420</li> <li>ISN 23</li> <li>CONTINUE</li> <li>O0000420</li> <li>ISN 24</li> <li>SO0 FORMATI'/ RETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 25</li> <li>FORMATI'/ RETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 24</li> <li>SO10 FORMATI'/ NETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 24</li> <li>SO10 FORMATI'/ RETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 24</li> <li>SO10 FORMATI'/ RETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 24</li> <li>SO10 FORMATI'/ NETRIEVED TEMPERATURE GRID')</li> <li>O0000420</li> <li>ISN 28<td></td><td></td><td></td><td></td></li></ul>				
<pre>. ' -30',' -20',' -10',' 0',' +10',' +20',' +30',' +40', . ' +50',' +60',' +70',' +80',' +90',' +10','+120', . '1130','+120',' +120','+120',' +200'/ . 0000220 CHHH HRITE RECORD HEADING OCOND260 ISN 6 MRITE(6,500) CHHE HRITE RECORD MEADING OCOND260 CHHE AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO DE OCOND270 CHHE AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO DE OCOND280 CHHE AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO DE OCOND280 ISN 8 DO 300 K=1,2 00000320 ISN 9 ILOG=9 ISN 8 DO 300 K=1,2 00000320 CHHE PRINT HEADING CALL FHOVE(HEDLIN(1),76,LONC'(K-1):419+1)) 00000320 CHHE PRINT HEADING CALL FHOVE(HEDLIN(1),76,LONC'(K-1):419+1)) 00000350 CHHE FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES 00000350 ISN 12 CALL FHOVE(HEDLIN(1),76,LONC'(K-1):419+1)) 00000350 ISN 14 PLEVEL=1.2 00000370 ISN 15 DO 200 J=1,62 00000380 ISN 16 IOFF=J-FILOG ISN 18 LINE(I)=IBLOCK(IOFF) 00000420 ISN 20 LOO CONTINUE 00000420 ISN 21 PLEVEL=0.2 00000430 ISN 22 MRITE(6,5030) PLEVEL,LINE 00000420 ISN 23 200 CONTINUE 00000450 ISN 24 300 CONTINUE 00000450 ISN 25 RETURN 00000450 ISN 25 RETURN 00000450 ISN 26 5000 FORMATI// RETRIEVE TEMPERATURE GRID') ISN 25 RETURN 00000450 ISN 26 5000 FORMATI// RETRIEVE TEMPERATURE GRID') ISN 28 5000 FORMATI// NETRIEVE TEMPERATURE GRID') ISN 29 5000 FORMATI// NETRIEVE TEMPERATURE GRID') ISN 28 5000 FORMATI// NETRIEVE TEMPERATURE GRID') ISN 28 5000 FORMATI// NETRIEVED TEMPERATURE GRID') ISN 28 5000 FORMATI// NETRIEVED TEMPERATURE GRID') ISN 28 5000 FORMATI// NETRIEVENTEME TEMPERATURE GRI</pre>	ISN	5		
<pre></pre>				
'+130','+140','+150','+160','+170','+190','+200'/         00000240           CH## WRITE RECORD HEADING         00000250           ISN         6         MRITE(6,500)         00000260           ISN         7         MRITE(6,5010) (IBLOCK(I),I=1,8)         00000270           C### FORM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO 0E         00000280           C### AND +10E TO +170E, PLUS ZONAL MEAN AND CLIMATOLOGY         00000310           ISN         8         DO 300 K=1,2         00000320           ISN         9         ILOG=9         00000320           ISN         10         IF (K.EQ.2) ILOG=1225         00000330           CH##         PRINT HEADING         00000350         00000350           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1):19+1))         00000350           ISN         13         WRITE(6,5020) HEDLIN(10ASH,I=1,21)         00000350           ISN         14         PLEVEL=1.2         00000350           ISN         15         DO 200 J=1,62         00000370           ISN         16         IOFF=3.142         00000400           ISN         18         LINE(1)=IBLOCK(IOFF)         00000450           ISN         18         LINE(1)=IBLOCK(IOFF)         00000450 </td <td></td> <td></td> <td></td> <td></td>				
CHHH HRITE RECORD HEADING         00000250           ISN         6         HRITE(6,5000)         00000270           ISN         7         HRITE(6,5000)         00000270           CHHH FORM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO DE         00000270           CHHH FORM AND PRINT OUTPUT LINES FOR THO LONGITUDE GROUPS -180E TO DE         00000270           ISN         8         DO 300 K=1,2         00000310           ISN         9         ILOG=9         00000310           ISN         10         IF (K.EQ.2) ILOG=1225         00000320           CHHH PRINT HEADING         000000350         00000350           CHHH FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000350           CHHH FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000370           ISN         13         HRITE(6,5020) HEDLIN,(DASH,I=1,21)         00000370           ISN         14         PLEVEL=1.2         00000370           ISN         15         D0 200 J=1,62         00000370           ISN         16         IDFF=J+ILOG         00000420           ISN         17         D0 100 I=1,19         00000420           ISN         18         LINE(I)=IBLOCK(IOFF)         00000420           ISN         20				
ISN         6         HRITE(6,500)         00000260           ISN         7         HRITE(6,5010) (IBLOCK(I),I=1,8)         00000270           C#H# FORM AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO DE         00000280           C#H# FORM AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO DE         00000280           ISN         8         DO 300 K=1,2         00000310           ISN         9         ILOG=9         00000320           ISN         10         IF (K.EQ.2) ILOG=1225         00000320           C#H# PRINT HEADING         00000350         00000350           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1):19+1))         00000350           ISN         13         WRITE(6,5020) HEDLIN,(DASH,I=1,21)         00000370           ISN         14         PLEVEL=1.2         00000380           ISN         15         DO 200 J=1,62         00000380           ISN         16         IOFF=J+ILOG         00000400           ISN         16         IOFF=JOFF+64         00000420           ISN         19         IOFF=IOFF+64         00000420           ISN         22         WRITE(6,5030) PLEVEL,LINE         000000420           ISN         23         200         C				
ISN         7         WRITE(6,5010) (IBLOCK(I),I=1,8)         00000270           C### FORM AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO 0E         00000280         00000280           ISN         8         DO 300 K=1,2         00000310           ISN         9         IL0G=9         00000320           ISN         9         IL0G=1225         00000320           C### AND PRINT MEADING         00000330         00000330           ISN         12         CALL FMOVE(HEDLIN(1),76,L0NC'(K-1):*19+1))         00000350           ISN         12         CALL FMOVE(HEDLIN(1),76,L0NC'(K-1):*19+1))         00000350           ISN         12         CALL FMOVE(HEDLIN(1),76,L0NC'(K-1):*19+1))         00000350           ISN         14         PLEVEL=1.2         00000370           ISN         14         PLEVEL=1.2         00000370           ISN         16         IOFF=J+IL0G         00000420           ISN         16         IOFF=IOFF+64         00000420           ISN         18         LINE(1)=IBLOCK(IOFF)         00000420           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000440           ISN         23         200         CONTINUE         00000440           I	TCN	4		
C###         FORM         AND         PRINT         OUTPUT         LINES         FOR         THO         LONGITUDE         GROUPS         -180E         TO         OE           ISN         8         DO         300         K=1,2         00000300         00000310           ISN         9         ILOG=9         00000310         00000310           ISN         10         IF         (K.Eq.2)         ILOG=1225         00000320           C###         PRINT         HEADING         00000330         00000330           ISN         12         CALL         FMOVE(HEDLIN(1),76,LONC'(K-1):419+1))         00000330           ISN         13         WRITE(6,5020)         HEDLIN,(DASH,I=1,21)         00000350           ISN         13         WRITE(1,5020)         HEDLIN,(DASH,I=1,21)         00000350           ISN         14         PLEVEL=1.2         00000350         00000350           ISN         15         DO         200         J=1,62         00000350           ISN         16         IOFF=J+ILOG         00000450         00000450           ISN         16         IDF=IOFF+64         00000450         00000450           ISN         23         200				
CNAME AND +10E TO +170E, PLUS ZONAL MEAN AND CLIMATOLOGY         00000290           ISN         8         D0 300 K=1,2         00000310           ISN         9         ILOG=9         00000320           ISN         10         IF (K.EQ.2) ILOG=1225         00000320           CNAME         PRINT HEADING         00000330           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1):19+1))         00000330           ISN         12         CALL FMOVE(HEDLIN:(DASH,I=1,21)         00000350           ISN         13         MRITEI(6,5020) HEDLIN:(DASH,I=1,21)         00000350           ISN         14         PLEVEL=1.2         00000370           ISN         15         D0 200 J=1,62         00000380           ISN         16         IOFF=JAILOG         00000400           ISN         18         LINE(I)=IBLOCK(IOFF)         00000420           ISN         20         100         CONTINUE         00000420           ISN         21         PLEVEL=PLEVEL+0.2         00000420           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000460           ISN         23         200         CONTINUE         00000460           ISN         24         300	1.34	'		
ISN         8         D0 300 K=1,2         00000300           ISN         9         ILOG=9         00000310           ISN         10         IF (K.Eq.2) ILOG=1225         00000330           CMMW         PRINT HEADING         00000330           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1)*19+1))         00000330           ISN         13         WRITE(6,5020) HEDLIN,(DASH,I=1,21)         00000360           CMMW         FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000370           ISN         14         PLEVEL=1.2         00000380           ISN         15         D0 200 J=1,62         00000380           ISN         16         IOFF=J+ILOG         00000400           ISN         17         D0 100 I=1,19         00000410           ISN         17         D0 100 I=1,19         00000410           ISN         17         D0 100 CONTINUE         00000430           ISN         18         LINE(I)=IBLOCK(IOFF)         00000450           ISN         21         PLEVEL=PLEVEL+0.2         00000450           ISN         22         CONTINUE         00000450           ISN         23         200         CONTINUE         000000460     <				
ISN       10       IF (K.EQ.2) ILOG=1225       00000320         C++++       PRINT HEADING       00000330         ISN       12       CALL FMOVE(HEDLIN(1),76,LONC'(K-1)+(19+1))       00000330         ISN       13       WRITE(6,5020) HEDLIN,(DASH,J=1,21)       00000350         C++++       FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES       00000360         ISN       14       PLEVEL=1.2       00000370         ISN       15       D0 200 J=1,62       00000380         ISN       16       IOFF=J+ILOG       00000400         ISN       16       IOFF=J+ILOG       00000400         ISN       18       LINE(1)=IBLOCK(IOFF)       00000420         ISN       20       100       CONTINUE       00000420         ISN       21       PLEVEL=PLEVEL+0.2       00000450         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000460         ISN       25       RETURN       00000460       00000460         ISN       25       RETURN       00000460       00000460         ISN	ISN	8	DO 300 K=1,2	00000300
CHHH         PRINT HEADING         00000330           ISN         12         CALL FMOVE(HEDLIN(1),76,LONC'(K-1):19+1))         00000340           ISN         13         WRITE(6,5020) HEDLIN,(DASH,I=1,21)         00000350           CHHH         FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000360           ISN         14         PLEVEL=1.2         00000370           ISN         15         DO 200 J=1,62         00000380           ISN         16         IOFF=J+ILOG         00000400           ISN         17         DO 100 I=1,19         00000410           ISN         18         LINE(I)=IBLOCK(IOFF)         00000420           ISN         19         IOFF=IOFF+64         00000440           ISN         21         PLEVEL=PLEVEL+0.2         00000440           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000440           ISN         23         200         CONTINUE         00000440           ISN         24         300         CONTINUE         00000440           ISN         25         RETURN         00000440         00000440           ISN         25         RETURN         00000440         000000440           ISN	ISN	9	ILOG=9	00000310
ISN       12       CALL FMOVE(HEDLIN(1),76,LONC'(K-1)*19+1))       00000340         ISN       13       WRITE(6,5020) HEDLIN,(DASH,I=1,21)       00000350         C###       FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES       00000360         ISN       14       PLEVEL=1.2       00000380         ISN       15       D0 200 J=1,62       00000380         ISN       16       IOFF=J+ILOG       00000380         ISN       16       IOFF=J+ILOG       00000400         ISN       17       D0 100 I=1,19       00000410         ISN       18       LINE(I)=IBLOCK(IOFF)       00000420         ISN       19       IOFF=IOFF+64       00000430         ISN       20       CONTINUE       00000430         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000440         ISN       25       RETURN       00000440       00000440         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       000000460 <t< td=""><td>ISN</td><td>10</td><td>IF (K.EQ.2) ILOG=1225</td><td>00000320</td></t<>	ISN	10	IF (K.EQ.2) ILOG=1225	00000320
ISN       13       WRITE(6,5020) HEDLIN,(DASH,I=1,21)       00000350         CMMM       FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES       00000360         ISN       14       PLEVEL=1.2       00000370         ISN       15       DO 200 J=1,62       00000380         ISN       16       IOFF=J+ILOG       00000390         ISN       17       DO 100 I=1,19       00000400         ISN       18       LINE(I)=IBLOCK(IOFF)       00000420         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000440         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000460         ISN       25       RETURN       00000460       00000460         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000460         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000500         ISN       27       5010       FORMAT(/' RET			CXXX PRINT HEADING	00000330
C****         FORM AND PRINT OUTPUT LINES FOR 62 PRESSURES         00000360           ISN         14         PLEVEL=1.2         00000370           ISN         15         DO 200 J=1,62         00000380           ISN         16         IOFF=J+ILOG         00000400           ISN         17         DO 100 I=1,19         00000400           ISN         18         LINE(I)=IBLOCK(IOFF)         00000420           ISN         19         IOFF=IOFF+64         00000420           ISN         20         100         CONTINUE         00000420           ISN         21         PLEVEL=PLEVEL+0.2         00000440           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000440           ISN         23         200         CONTINUE         00000450           ISN         23         200         CONTINUE         00000460           ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000490         00000490           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(/' RETRIEVED TEMPERATURE GRID')	ISN	12	CALL FMOVE(HEDLIN(1),76,LONC((K-1):+19+1))	00000340
ISN       14       PLEVEL=1.2       00000370         ISN       15       DO 200 J=1,62       00000380         ISN       16       IOFF=J+ILOG       00000390         ISN       16       IOFF=J+ILOG       00000400         ISN       17       DO 100 I=1,19       00000410         ISN       18       LINE(I)=IBLOCK(IOFF)       00000420         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000440         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000470         ISN       24       300       CONTINUE       00000470         ISN       25       RETURN       00000470       00000470         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000470         ISN       26       5000       FORMAT(/' RECSIZE',IS,2X,'REC\$*',IS,2X,'TYPE',IS,2X,       00000520         ISN       27       5010       FORMAT(/' RECSIZE',IS,2X,'REC\$*',IS,2X,'TYPE',IS,2X,       00000520         ISN       28       <	ISN	13		000003 <b>50</b>
ISN       15       D0 200 J=1,62       00000380         ISN       16       IOFF=J+ILOG       00000390         ISN       17       D0 100 I=1,19       00000400         ISN       17       D0 100 I=1,19       00000400         ISN       18       LINE(I)=IBLOCK(IOFF)       00000420         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000420         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000460         ISN       25       RETURN       00000460       00000460         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000490         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000500         .       'DATA DAY/YEAR',1X,XI,31X,14,42X, 'GEN DATE',1X,I3,1X,14,2X,       00000500         .       'DATA DAY/YEAR',1X,XI,51X,244,3X,19(A4,2X)/1X,2A4,3X,19(A4,2X))       00000520         ISN       28       5020 <td></td> <td></td> <td></td> <td>00000360</td>				00000360
ISN       16       IOFF=J+ILOG       00000390         ISN       17       D0 100 I=1,19       00000400         ISN       18       LINE(I)=IBLOCK(IOFF)       00000410         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000430         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000450         ISN       23       200       CONTINUE       00000460         ISN       23       200       CONTINUE       00000450         ISN       24       300       CONTINUE       00000460         ISN       25       RETURN       00000480       00000480         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000480         ISN       27       5010       FORMAT(/' RECSIZE',15,2X,'REC\$',15,2X,'TYPE',15,2X,       00000500         .       'DATA DAY/YEAR',1X,13,1X,14,2X,'GEN DATE',1X,13,1X,14,2X,       00000520       .       .'DATA DAY/YEAR',1X,13,1X,14,2X,'GEN DATE',1X,13,1X,14,2X,       00000520         .       'GRID LAT(X100				
ISN       17       D0 100 I=1,19       00000400         ISN       18       LINE(I)=IBLOCK(IOFF)       00000410         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000430         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000470         ISN       25       RETURN       00000470       00000470         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000470         ISN       27       5010       FORMAT(/' RECSIZE',I5,2X,'REC\$',I5,2X,'TYPE',I5,2X,       00000500         .       'DATA DAY/YEAR',1X,I3,IX,I4,2X,'GEN DATE',1X,I3,IX,I4,2X,       00000510         .       'GRID LAT(X100)',1X,I5)       00000520       00000520         ISN       28       5020       FORMAT(/' LNO/P')',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X))       00000520         ISN       29       5030       FORMAT(2X,F4.1,4X,1916)				
ISN       18       LINE(I)=IBLOCK(IOFF)       00000410         ISN       19       IOFF=IOFF+64       00000420         ISN       20       100       CONTINUE       00000430         ISN       21       PLEVEL=PLEVEL+0.2       00000440         ISN       22       WRITE(6,5030) PLEVEL,LINE       00000450         ISN       23       200       CONTINUE       00000460         ISN       23       200       CONTINUE       00000460         ISN       23       200       CONTINUE       00000460         ISN       24       300       CONTINUE       00000470         ISN       25       RETURN       00000480       00000480         ISN       26       5000       FORMAT(/' RETRIEVED TEMPERATURE GRID')       00000490         ISN       27       5010       FORMAT(/' RECRIZE',15,2X, 'TYPE',15,2X, 'TYPE',15,2X, 'O0000510				
ISN         19         IOFF=IOFF+64         00000420           ISN         20         100         CONTINUE         00000430           ISN         21         PLEVEL=PLEVEL+0.2         00000440           ISN         21         WRITE(6,5030) PLEVEL,LINE         00000450           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000450           ISN         23         200         CONTINUE         00000460           ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000480         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(' RECSIZE',I5,2X,'REC\$',I5,2X,'TYPE',I5,2X,         00000500           ISN         27         5010         FORMAT(' RECSIZE',I5,2X,'REC\$',I5,2X,'TYPE',I5,2X,         00000520           ISN         27         5010         FORMAT(' RECSIZE',I5,2X,'REC\$',I5,2X,'TYPE',I5,2X,         00000520           ISN         28         5020         FORMAT(/' LNIPO/P)',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X))         00000520           ISN         29         5030         FORMAT(2X,F4.1,4X,1916)         000000540				
ISN         20         100         CONTINUE         00000430           ISN         21         PLEVEL=PLEVEL+0.2         00000440           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000450           ISN         23         200         CONTINUE         00000450           ISN         24         300         CONTINUE         00000460           ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000480         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(/' RECSIZE', J5,2X, 'REC\$', J5,2X, 'TYPE', J5,2X, '00000510         '00000500           ISN         27         5010         FORMAT(/' LNIPO/P)', JX, J9, (A4,2X, 'GEN DATE', JX, J3, JX, J4, 2X, '00000510         '00000520           ISN         28         5020         FORMAT(/' LNIPO/P)', JX, J9, (A4,2X)/1X, 2A4,3X, J9(A4,2X))         00000520           ISN         29         5030         FORMAT(2X,F4.1,4X, J916)         00000540				
ISN         21         PLEVEL=PLEVEL+0.2         00000440           ISN         22         WRITE(6,5030) PLEVEL,LINE         00000450           ISN         23         200         CONTINUE         00000460           ISN         24         300         CONTINUE         00000460           ISN         25         RETURN         00000480         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000480           ISN         26         5010         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000500           ISN         27         5010         FORMAT(' RECSIZE', J5,2X, 'REC\$', J5,2X, 'TYPE', J5,2X,         00000500           .         'DATA DAY/YEAR', 1X, J3, 1X, I4,2X, 'GEN DATE', 1X, J3, 1X, J4,2X,         00000510           .         'GRID LAT(X100)', 1X, J5)         00000520         00000520           ISN         28         5020         FORMAT(/' LNPO/P)', 3X, 19(A4,2X)/1X, 2A4, 3X, 19(A4,2X))         00000530           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540				
ISN         22         WRITE(6,5030) PLEVEL,LINE         00000450           ISN         23         200         CONTINUE         00000460           ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(/' RECSIZE',IS,2X,'REC#',IS,2X,'TYPE',IS,2X,         00000520           .         'DATA DAY/YEAR',IX,I3,1X,I4,2X,'GEN DATE',1X,I3,1X,I4,2X,         00000520         00000520           ISN         28         5020         FORMAT(/' LN(PO/P)',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X))         00000530           ISN         29         5030         FORMAT(2X,F4.1,4X,1916)         00000540				
ISN         23         200         CONTINUE         00000460           ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(' RECSIZE', I5,2X, 'REC\$', I5,2X, 'TYPE', I5,2X, 00000500         00000520           .         'DATA DAY/YEAR', 1X, I3, 1X, I4,2X, 'GEN DATE', 1X, I3, 1X, I4,2X, 00000520         00000520           ISN         28         5020         FORMAT(/' LN(PO/P)', 3X, 19(A4,2X)/1X, 2A4, 3X, 19(A4,2X))         00000530           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540				
ISN         24         300         CONTINUE         00000470           ISN         25         RETURN         00000480           ISN         26         5000         FORMAT(/' RETRIÉVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(' RECSIZE', I5,2X, 'REC\$', I5,2X, 'TYPE', I5,2X, 'O0000500         00000500           .         'DATA DAY/YEAR', IX, I3, IX, I4,2X, 'GEN DATE', IX, I3, IX, I4,2X, '00000510         00000520           ISN         28         5020         FORMAT(/' LNIPO/P)', 3X, 19(A4,2X)/1X, 2A4, 3X, 19(A4,2X))         00000530           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540		-		
ISN         25         RETURN         00000480           ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(' RECSIZE', J5,2X, 'REC\$', I5,2X, 'TYPE', I5,2X, 'TYPE', I5,2X, 'O0000500         00000500           'DATA DAY/YEAR', 1X, I3, 1X, I4,2X, 'GEN DATE', 1X, I3, 1X, I4,2X, 'O0000510         'GRID LAT(X100)', 1X, I5)         00000520           ISN         28         5020         FORMAT(/' LNIP0/P)', 3X, 19(A4,2X)/1X, 2A4, 3X, 19(A4,2X))         00000520           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540				
ISN         26         5000         FORMAT(/' RETRIEVED TEMPERATURE GRID')         00000490           ISN         27         5010         FORMAT(' RECSIZE', I5,2X, 'REC\$', I5,2X, 'TYPE', I5,2X, 00000500         00000500           'DATA DAY/YEAR', IX, I3, IX, I4,2X, 'GEN DATE', IX, I3, IX, I4,2X, 00000510         'GRID LAT(X100)', IX, I5)         00000520           ISN         28         5020         FORMAT(/' LNPO/P)', 3X, 19(A4,2X)/1X, 2A4, 3X, 19(A4,2X))         00000520           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540				
ISN         27         5010         FORMAT(' RECSIZE', I5,2X, 'REC\$', I5,2X, 'TYPE', I5,2X, 00000500           'DATA DAY/YEAR', 1X, I3, 1X, I4,2X, 'GEN DATE', 1X, I3, 1X, I4,2X, 00000510         'GRID LAT(X100)', 1X, I5)         00000520           ISN         28         5020         FORMAT(' LN(P0/P)', 3X, 19(A4, 2X)/1X, 2A4, 3X, 19(A4, 2X))         00000530           ISN         29         5030         FORMAT(2X, F4.1, 4X, 1916)         00000540				
. 'DATA DAY/YEAR',1X,I3,1X,I4,2X, 'GEN DATE',1X,I3,1X,I4,2X, 00000510 'GRID LAT(X100)',1X,I5) 00000520 ISN 28 5020 FORMAT(/' LN(P0/P)',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X)) 00000530 ISN 29 5030 FORMAT(2X,F4.1,4X,1916) 00000540				
'GRID         LAT(X100)',1X,15)         00000520           ISN         28         5020         FORMAT(/' LN(P0/P)',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X))         00000530           ISN         29         5030         FORMAT(2X,F4.1,4X,1916)         00000540				
ISN 28 5020 FORMAT(/'LN(P0/P)',3X,19(A4,2X)/1X,2A4,3X,19(A4,2X)) 00000530 ISN 29 5030 FORMAT(2X,F4.1,4X,1916) 00000540				
ISN 29 5030 FORMAT(2X,F4.1,4X,1916) 00000540	ISN	28		
	ISN	30		

ł

		· · · · · · · · · · · · · · · · · · ·	00000010
(SN	1	SUBROUTINE DM7403(IBLOCK,DATTYP) C#************************************	
			00000030
		C SUBROUTINE - DH7403	00000040
		C C PURPOSE - FORM AND PRINT OUTPUT LINES FOR SAMS TYPE 7403 RECORDS -	00000050
		C TEMPERATURE GRIDS AND TEMPERATURE GRID ERRORS.	00000060
			00000070
		C ADCIMENTS DASSED -	00000080
		C TRUCCK = ARRAY CONSISTING OF TEMPERATURE GRID DATA OR ERRORS	00000090
		C DATTYP = DATA TYPE DESIGNATING TEMPERATURE GRID OR GRID ERRORS	00000100
		C DATE DATE TO DESCRIPTION OF T	00000110
		C SUBPROGRAMS CALLED ~ NONE	00000120
		c	00000130
		C CALLED BY - GETGRT	00000140
		C	00000150
		C PROGRAMMER - TOM NUTTER 03/01/85	00000160
		C PROBRIER - FOR HOLLER CLOSE - CONTROL - CONT	+00000170
ISN	2	INTEGER DATTYP	00000180
(SN	3	INTEGER#2 IBLOCK(2994)	00000200
ISN	4	CHARACTER#4 DASH,LONG(36),HEDLIN(18)	00000210
[SN	5	CHARACTER#5 LAT(48)	00000220
(SN	6	DATA DASH/''/	00000230
ISN	7	DATA LONG/'-180','-170','-160','-150','-140','-130','-120',	0000024
		'-110','-100',' -90',' -80',' -70',' -60',' -50',' -40', ' -30',' -20',' -10',' 0',' +10',' +20',' +30',' +40',	0000025
		. ' _30',' _20',' _10',' 0',' +10',' +20',' +30',' +40', . ' +50',' +60',' +70',' +80',' +90','+100','+110','+120',	0000026
			0000027
		'+130','+140','+150','+160','+170'/ DATA LAT/'-50.0','-47.5','-45.0','-42.5','-40.0','-37.5','-35.0',	
ISN	8	DATA LATY -50.0', -47.5', -45.0', -42.5', -40.0', 57.5', -15.0', -32.5', -30.0', -27.5', -25.0', -22.5', -20.0', -17.5', -15.0',	0000029
		· -12.5', '-10.0', ' -7.5', ' -5.0', ' -2.5', ' 0.0', ' +2.5', ' +5.0',	0000030
		· +7.5', '+10.0', '+12.5', '+15.0', '+17.5', '+20.0', '+22.5', '+25.0',	0000031
		+27.5', '+30.0', '+32.5', '+35.0', '+37.5', '+40.0', '+42.5', '+45.0'	0000032
		· '+47.5', '+50.0', '+52.5', '+55.0', '+57.5', '+60.0', '+62.5', '+65.0',	0000033
		. '+67.5'/	0000034
		CHAR WRITE RECORD IDENTIFICATION INFORMATION	0000035
	9	WRITE(6,5000) (IBLOCK(I),I=1,8),IBLOCK(10),IBLOCK(11)	0000036
ISN	10	IF (DATTYP.EQ.2) WRITE(6,5010)	0000037
ISN	12	TE (DATTYP EQ 102) WRITE(6.5020)	0000038
ISN	16	C*** FORM AND PRINT OUTPUT LINES FOR TWO LONGITUDE GROUPS -180E TO -100	0000039
		C*** AND OE TO +170E	0000040
ISN	14	DO 200 J=1,2	0000041
ISN	15	ISTRT=-12	0000042
ISN	16	IF (J.EQ.2) ISTRT=6	0000043
		CHAN PRINT HEADING	0000044
ISN	18	CALL FMOVE(HEDLIN(1),72,LONG((J-1)*18+1))	0000045
ISN	19	WRITE(6,5030) IBLOCK(12), HEDLIN, (DASH, M=1, 18)	0000046
1011	- /	C**** FORM AND PRINT OUTPUT LINES FOR LATITUDES -50 N TO +67.5 N	0000047
ISN	20	DO 100 LOOP=1,48	0000048
ISN	21	ISTRT=ISTRT+36	000004
ISN	22	ISTOP=ISTRT+17	0000050
ISN	23	WRITE(6,5040) LAT(LOOP),(IBLOCK(I),I=ISTRT,ISTOP)	000005
ISN	24	100 CONTINUE	000005
ISN	25	200 CONTINUE	000005
ISN	26	RETURN	000005
ISN	27	5000 FORMAT(/' RECSIZE', 15, 2X, 'REC#', 15, 2X, 'TYPE', 15, 2X, 'ORBIT TYPE',	
		13,2X, 'DATA DAY/YEAR',1X,13,1X,14,2X, 'GEN DATE',1X,13,1X,14,2	000005
		'SCALE FACTOR', I3, 2X, 'DATA TYPE', I4)	
ISN	28	5010 FORMAT(/' TEMPERATURE GRID DATA')	000005
ISN	29	5020 FORMAT(/' TEMPERATURE GRID ERRORS')	000005
ISN	30	5030 FORMAT( ' PRESSURE LEVEL= ', 15/9X, 18(A4, 2X)/9X, 18(A4, 2X))	
ISN	31	5040 FORMAT(1X,A5,1X,1816)	000006
1044			

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		**1	7.*8
ISN	1	SUBROUTINE DMPCHP(IBLOCK, OPTION)	00000010
1.314	•		***00000020
		C SUBROUTINE - DHPCHP	00000030
		C	00000040
		C PURPOSE - PRINT A FORMATTED DUMP OF A SAMS ZONAL MEAN METHANE OR	00000050
		C NITROUS-OXIDE COMPOSITION RECORD.	00000060
		C C	00000070
		C ARGUMENTS PASSED -	00000080
		C IBLOCK = ARRAY CONSISTING OF ZONAL MEAN DATA	00000090
		C OPTION = DUMP OPTIONS FROM INPUT CARD (1=DUMP, 0=NO DUMP)	00000100
		c	00000110
		C SUBPROGRAMS CALLED - PRTREC	00000120
		c	00000130
		C CALLED BY - GETCMP	00000140
		C	00000150
		C PROGRAMMER - TOM NUTTER 03/05/85	00000160
		C <del>************************************</del>	+**00000170
ISN	2	INTEGER OPTION(8)	00000180
ISN	3	INTEGER*2 IBLOCK(2994)	00000190
		C*** PRINT RECORD IDENTIFYING INFORMATION	00000200
ISN	4	WRITE(6,5010) (IBLOCK(I),I=1,15)	00000210
		CHHH DUMP NITROUS-OXIDE RECORD	00000220
ISN	5	IF (IBLOCK(3).EQ.7405) THEN	00000230
		CXXXX DUMP NITROUS-OXIDE MIXING RATIOS	00000240
ISN	6	IF (OPTION(5).EQ.1) THEN	00000250
ISN	7	WRITE(6,5020)	00000260
ISN	8	CALL PRTREC(IBLOCK, -480)	00000270
ISN	9	ENDIF	00000280
		CXXXX DUMP NITROUS-OXIDE MIXING RATIO ERRORS	00000290
ISN	10	IF (OPTION(6).EQ.1) THEN ~	00000300
ISN	11	WRITE(6,5030)	00000310
ISN	12	CALL PRTREC(IBLOCK, 1008)	00000320
ISN	13	ENDIF	00000330
ISN	14	ENDIF	00000340
		C*** DUMP METHANE RECORD	00000350
ISN	15	IF (IBLOCK(3).EQ.7406) THEN	00000360
		C*** DUMP METHANE MIXING RATIOS	00000370
ISN	16	IF (OPTION(7).EQ.1) THEN	00000380
ISN	17	WRITE(6,5040)	00000390
ISN	18	CALL PRTREC(IBLOCK,-480)	00000400
ISN	19	ENDIF	00000410
		C*** DUMP METHANE MIXING RATIO ERRORS	00000420
ISN	20	IF (OPTION(8).EQ.1) THEN	00000430
ISN	21	WRITE(6,5050)	00000440
ISN	22	CALL PRTREC(IBLOCK, 1008)	00000450
ISN	23	ENDIF	00000460
ISN	24	ENDIF	00000470
ISN	25	RETURN	00000480
ISN	26	5010 FORMAT(/' RECSIZE', 15,2X, 'REC#', 15,2X, 'TYPE', 15,2X, 'DATA DATE',	00000490
		. 1X,I3,1X,I4,2X,'CHAN',I2,2X,'SIEVE SET ENPMC',I2,2X,'CLPMC',	00000500
		. I2,2X, 'A1PMC',I2,2X, 'C1PMC',I2,2X, 'GEN DATE',1X,I3,1X,I4,2X,	00000510
7011		. /' ELEMENTS IN PROFILE',14,2X,'BOT/TOP LEVELS',13,1X,13)	00000520
ISN	27	5020 FORMAT(/' NITROUS-OXIDE MIXING RATIOS')	00000530
ISN	28	5030 FORMAT(/' NITROUS-OXIDE MIXING RATIO ERRORS')	00000540
ISN	29	5040 FORMAT(/' METHANE MIXING RATIOS')	00000550
ISN	30	5050 FORMAT(/' METHANE MIXING RATIO ERRORS')	00000560
ISN	31	END	00000570

.

ISN	1		SUBROUTINE PRTREC(IBLOCK, INDEX)	0000001
		CXXXX	***************************************	H#0000002
		C SUB	ROUTINE - PRTREC	0000003
		С		0000004
		C PUR	POSE - FORM AND PRINT OUTPUT LINES FOR ZONAL MEAN METHANE OR	0000005
		CN	ITROUS-OXIDE COMPOSITION RECORDS.	0000006
		С		0000007
		C ARG	UMENTS PASSED -	0000008
			BLOCK = ARRAY CONSISTING OF ZONAL MEAN DATA	000009
		C I C	NDEX = INDEX THAT POSITIONS TO MIXING RATIOS OR MIXING RATIO ERRO	0000010 000010 0000010
		-		0000012
		C	PROGRAMS CALLED - NONE	0000012
		+		0000014
		C	LED BY - DMPCMP	0000015
		-	GRAMMER - TOM NUTTER 03/02/85	0000016
			GRAMMER - 100 MG112R 05/02/05 ************************************	
ISN	2	CARAA	INTEGER*2 IBLOCK(2994),LINE(16),IGROUP(496)	0000018
ISN	3		CHARACTER*4 DASH	0000019
ISN	4		CHARACTER*5 HEDLIN(16),LAT(48)	0000020
ISN	5		DATA DASH/''/	000002
ISN	6		DATA LAT/'-50.0','-47.5','-45.0','-42.5','-40.0','-37.5','-35.0',	
	•		. '-32.5', '-30.0', '-27.5', '-25.0', '-22.5', '-20.0', '-17.5', '-15.0',	
			· '-12.5', '-10.0', ' -7.5', ' -5.0', ' -2.5', ' 0.0', ' 2.5', ' 5.0',	
			. ' 7.5', ' 10.0', ' 12.5', ' 15.0', ' 17.5', ' 20.0', ' 22.5', ' 25.0',	
			. ' 27.5', ' 30.0', ' 32.5', ' 35.0', ' 37.5', ' 40.0', ' 42.5', ' 45.0',	
			. ' 47.5', ' 50.0', ' 52.5', ' 55.0', ' 57.5', ' 60.0', ' 62.5', ' 65.0';	
			. ' 67.5'/	0000028
		CHHH	FORM AND PRINT OUTPUT LINES FOR THREE LATITUDE GROUPS -50.0N TO	000002
			-12.5N, -10.0N TO +27.5N, +30.0N TO +67.5N	0000031
ISN	7		IOFF=INDEX	0000033
ISN	8		DO 300 K=1,3	000003
ISN	9		IOFF=IOFF+496	000003
ISN	10		CALL FMOVE(IGROUP(1),992,IBLOCK(IOFF))	0000034
		CHHH	PRINT HEADING	000003
ISN	11		CALL FMOVE(HEDLIN(1),80,LAT((K-1)*16+1))	000003
ISN	12		WRITE(6,5000) HEDLIN,(DASH,I=1,18)	000003
ISN	13		PLEVEL=2.8	000003
ISN	14		DO 200 J=1,31	000003
ISN	15		DO 100 I=1,16	000004
ISN	16		L=(I-1)*31+J	000004
ISN	17		LINE(I)=IGROUP(L)	000004
ISN	18	100	CONTINUE	000004
		CXXX	PRINT OUTPUT LINE	0000044
ISN	19		PLEVEL=PLEVEL+0.2	000004
ISN	20		WRITE(6,5010) PLEVEL,LINE	000004
ISN	21	200	CONTINUE	000004
ISN	22	300	CONTINUE	000004
ISN	23		RETURN	000004
ISN	24	5000	FORMAT(/' LN(PO/P)',2X,16(A5,1X)/1X,2A4,3X,16(A4,2X))	000005
ISN	25	5010	FORMAT(2X,F4.1,4X,1616)	000005
ISN	26		END	000005

## THE STRATOSPHERIC AND MESOSPHERIC SOUNDER ON NIMBUS-7

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The stratospheric and mesospheric sounder on Nimbus 7

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The stratospheric and mesospheric sounder (s.a.m.s.) instrument was launched on the Nimbus G satellite on 24 October 1978. It is designed to measure temperature and concentration profiles of various gases in the height range 20-100 km by detecting either their thermal emission or, in some cases, resonant scattering of sunlight. The gases selected, CO<sub>2</sub>, CO, CH<sub>4</sub>, NO, N<sub>2</sub>O and H<sub>2</sub>O, significantly affect the

The gases selected,  $CO_2$ , CO,  $CH_4$ , NO,  $N_2O$  and  $H_2O$ , significantly affect the upper atmosphere energy budget by their influence on the concentration of the primary sunlight absorber, ozone. This influence is disproportionate to their own concentration because of the existence of 'catalytic cycles' which destroy ozone while regenerating the catalyst.

A description of the instrument, its principles of operation and some of the methods of retrieval used is presented, together with some preliminary results from the first 3 months of operations.

## 1. INTRODUCTION

As has been explained earlier in this meeting by Professor Labitzke and by Dr Barnett, during the last 10 years remote sounding observations from satellites of atmospheric temperature have been extended through the stratosphere into the mesosphere. Much extra insight into atmospheric phenomena has arisen through the provision of these observations substantially continuous in time and with near global coverage. The purpose of the stratospheric and mesospheric sounder (s.a.m.s.) on Nimbus 7 is to provide similar coverage in time and space for observations of the concentration of a number of important minor constituents in the stratosphere and mesosphere, all of which are directly or indirectly involved in the chemistry of ozone (cf. Thrush, this symposium). Like previous remote sounding instruments built for Nimbus satellites in which the Oxford University Atmospheric Physics Department has been involved, s.a.m.s. observes infrared radiation emitted by the atmospheric constituents it is required to measure. In this paper, the principles of the techniques employed will first be outlined, then the different measurements being made by s.a.m.s. will be discussed together with a description of the instrument and the method of calibration. Finally, some early results will be presented.

# 2. PRINCIPLES OF OPERATION

## 2.1. Limb-sounding

The advantages of the limb view (see figure 7) are that emission observations can be made from as long a path through the atmosphere as possible, and that behind the atmospheric path is the cold radiation background of space. Such a limb path possesses an air mass approximately 70 times that in a vertical path above the lowest point of the path. The atmosphere's pressure profile and the geometry of the path both weight the material in the path strongly towards the tangent height giving good vertical resolution. For emission from spectral lines that are

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## 220 J. R. DRUMMOND AND OTHERS

not saturated at their centres, 60% of the emission from a limb path originates from altitudes within 3 km of the lowest altitude of the path (see figure 4). By contrast, the horizontal resolution of a limb observation is poor, the length of the limb path being of the order of 400 km.

Radiometers carried by balloons or aircraft have, for a long time, observed emission from the atmosphere's limb. Results from the limb radiance inversion radiometer (l.r.i.r.) on Nimbus 6 (Gille *et al.* 1975), the first satellite radiometer to observe radiance from the atmosphere's limb, have demonstrated the vertical resolution which can be achieved in measurements of temperature and ozone (Gille, this symposium).

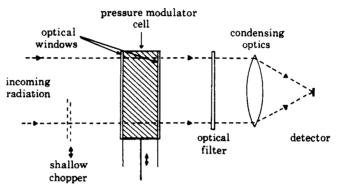


FIGURE 1. Simplified diagram of a p.m.r. optical system employing a 'shallow chopper' which only modulates part of the incoming radiation beam.

#### 2.2. The pressure modulator

The pressure modulator is a device which selectively modulates the emission from a gas by using the absorption lines of the same gas as an optical filter. A cell of gas is included in the optical path of the radiometer (figure 1) and its pressure varied cyclically. The transmission of the cell is modulated only at optical frequencies that lie within the absorption lines of the gas, thus selectively modulating emission from the same gas incident on the radiometer. After detection, the signal at the frequency of the modulation may be recovered by phase-sensitive detection. A broad optical filter limits the spectral bandwidth to the appropriate emission band.

The pressure modulator technique is a development from the selective chopper technique employed for temperature sounding instruments on the Nimbus 4 and 5 satellites (Houghton & Smith 1970; Ellis et al. 1973). It has previously been flown in a temperature sounder on the Nimbus 6 satellite (Taylor et al. 1972; Curtis et al. 1974). The advantages of the technique for composition sounding have been described by Chaloner et al. (1978), who employed it in a balloon-borne radiometer for the measurement of NO, NO<sub>2</sub> and H<sub>2</sub>O in the stratosphere. Further details of the technique are also given in Drummond & Jarnot (1978).

### 2.3. Double chopping

The gas pressure in the pressure modulator cell (p.m.c.) is modulated at the resonant frequency (20-40 Hz) of the mechanical system (see § 4.3). It is also useful to include a conventional chopper at a much higher frequency (ca. 240 Hz) in the optical system. This modulates all 'incoming' radiation with the optical filter profile and produces a 'wideband' signal

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

at the detector. The 'wideband' signal, in the absence of emission from other species, provides information for lower in the atmosphere than the p.m.r. signal (see figure 2). The wideband signal may also be used for the elimination of signals from atmospheric constituents whose absorption bands overlap that of the wanted species (see Drummond & Jarnot 1978), and in the determination of the atmospheric level being observed (§ 3.3). In s.a.m.s. a shallow chopper is employed (see § 4), which chops only ca. 2.5% of the aperture of the system. The shallow chopper is preferred over a conventional 100% chopper because it does not obstruct the beam for the pressure modulator further back in the optical chain.

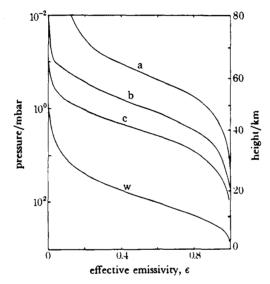


FIGURE 2. The effective emissivity, ε (see equation (5)) for water vapour evaluated by using a climatological equatorial temperature profile and a constant volume mixing ratio of 5 × 10<sup>-6</sup>. Curves (a), (b) and (c) are for the s.a.m.s. B2 p.m.r. channel with mean cell pressures of 0.87 mbar, 4.48 mbar and 15.3 mbar respectively. Curve (w) is for the corresponding wideband channel and is virtually independent of modulator pressure.

#### 3. MEASUREMENTS WITH S.A.M.S.

The radiant power, W, emitted by a limb path at frequency  $\nu$  and received by one of the s.a.m.s. detectors is given by

$$W = A\Omega \int_0^\infty \int_0^\infty J_\nu(x) \tau_{0\nu} \tau_{m\nu} \frac{\mathrm{d}\tau_\nu(x)}{\mathrm{d}x} \,\mathrm{d}x \,\mathrm{d}\nu, \tag{1}$$

where A is the aperture of the telescope,  $\Omega$  the field of view,  $\tau_{\nu}(x)$  the transmission between the satellite and the position along the path described by the coordinate x,  $\tau_{0\nu}$  the static component of the optical transmission in the appropriate channel,  $\tau_{m\nu}$  the modulated component of the optical transmission, and  $J_{\nu}(x)$  is the source function appropriate to the emission band in question at the location described by the coordinate x. Under conditions of local thermodynamic equilibrium (l.t.e.), the source function is equal to the Planck function  $B_{\nu}(T)$  at the temperature T of the atmosphere at location x.

Integration is along the limb path and over the frequencies where the optical components have any significant transmission.

222

## J. R. DRUMMOND AND OTHERS

Calibration of the instrument is carried out by inserting a black-body at temperature  $T_b$  in the path when the radiant power  $W_b$  received is

$$W_b = A\Omega \int_0^\infty B_\nu(T_b) \tau_{0\nu} \tau_{m\nu} \,\mathrm{d}\nu.$$
<sup>(2)</sup>

The radiance, L, of an atmospheric path is then defined as

$$L = \frac{W}{W_b} \frac{\int_0^\infty B_{\nu}(T_b) \tau_{0\nu} \tau_{m\nu} \, \mathrm{d}\nu}{\int_0^\infty \tau_{0\nu} \tau_{m\nu} \, \mathrm{d}\nu}.$$
 (3)

For channels covering a reasonably narrow frequency interval the ratio of the integrals in equation (3) is  $B_{\bar{\nu}}(T_b)$ , where  $\bar{\nu}$  is the centre frequency of the interval.

To illustrate the radiometric measurements being made it is useful when l.t.e. applies to write a simplified expression for the radiance L measured by a particular channel, namely:

1

$$L = \epsilon B_{\bar{\nu}}(\bar{T}), \tag{4}$$

$$\epsilon = \frac{\int_0^\infty \int_0^\infty \tau_{0\nu} \tau_{m\nu} \frac{\mathrm{d}\tau_{\nu}(x)}{\mathrm{d}x} \,\mathrm{d}x \,\mathrm{d}\nu}{\int_0^\infty \tau_{0\nu} \tau_{m\nu} \,\mathrm{d}x}$$
(5)

where

may be defined as the *effective emissivity* of the path and where  $\overline{T}$  is a mean temperature for the path.

#### 3.1. Temperature measurements

If the radiance is measured from the emission band of an atmospheric constituent whose distribution is accurately known, for any given path the emissivity  $\epsilon$  in equation (4) may be calculated and hence a mean temperature  $\overline{T}$  for that path may be deduced. A set of radiances received from different levels as the instrument scans over the limb may be retrieved to give a temperature profile (see Rodgers (1976) for a review of retrieval methods and § 6.2). In s.a.m.s. two wideband and two pressure modulator channels are included observing emission from the 15  $\mu$ m band of CO<sub>2</sub>, a gas that is substantially uniformly mixed up to at least 90 km altitude. Above about 80 km, local thermodynamic equilibrium no longer applies for this band; the quantity that is then measured is the vibrational temperature of the band, which is of considerable importance because emission from it is the major energy sink in the upper mesosphere and lower thermosphere.

#### 3.2. Composition measurements

Given that the atmosphere's temperature structure has been determined by the method of § 3.1, from radiance observations on any other band an effective emissivity  $\epsilon$  may be determined from equation (4), and hence the distribution of the emitting constituent observed. In figure 2 are curves showing the variation of effective emissivity for emission from the long wave water vapour band as measured by a wide band channel and by channels containing pressure modulator cells with water vapour at different pressures, demonstrating that the water vapour distribution can be measured from the tropopause through into the upper mesosphere. Some results for H<sub>2</sub>O from an early orbit of s.a.m.a. are given in § 6.

## APPENDIX C

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

The method by which a detailed distribution is derived is as follows. First of all, through theoretical calculations on the basis of spectral information about the band in question, together with a programme of laboratory measurements as described in § 5, algorithms describing the transmission of any atmospheric path are derived. A 'first guess' profile for the distribution of the gas in question is then assumed and, on the basis of the temperature profile derived by the method of § 3.1, an expected set of radiances is computed. These computed radiances are then compared with the measured radiances, and by an iterative process a profile of the gas distribution is found for which computed and measured radiances agree.

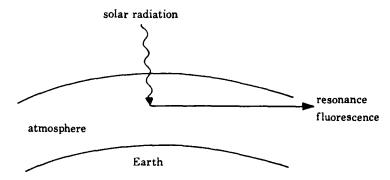


FIGURE 3. The generation of resonance fluorescence signals by scattering of sunlight within the atmosphere.

To acquire information about the distribution of some constituents in the mesosphere and lower thermosphere it is possible to use observations of resonance fluorescence of solar radiation from those molecules which have strong bands in the near infrared (figure 3) (cf. Table 1 for list of bands included in s.a.m.s.). Provided that absorption by solar radiation is the only process of excitation of the band in question, and provided that collisional quenching is negligible, the radiance observed is a function of the solar flux incident on the layer and the number of molecules along the path. In practice, an elaborate radiative transfer calculation has to be carried out for each situation because multiple scattering, excitation by other non-thermal processes (e.g. by particles), and details of quenching processes have also to be considered. Interpretation of measurements from fluorescence channels, therefore, is extremely complex.

### 3.3. Reference pressure determination

Because s.a.m.s. is a limb viewing instrument it is necessary that accurate information is available regarding the pressure of the atmosphere at the levels being viewed at any given time. Consider information such as that contained in figure 2. When, for instance, the effective emissivity of an atmospheric path is about 0.5, a change in signal corresponding to a change in emissivity of 0.01 (equivalent to a change in mixing ratio of *ca*. 6%) could also result from a change in the level being viewed of 0.15 km, which is equivalent to a change in pressure at the tangent point of *ca*. 2% or 0.003° in viewing direction.

The attitude of the spacecraft is controlled only to  $ca. 0.5^{\circ}$  in all three axes; it is therefore necessary from the measurements made by s.a.m.s. itself to determine the appropriate information about the pressure at the part of the atmosphere being observed.

A method for providing this information arises from comparing the signals from a wideband and a pressure modulator channel in the 15  $\mu$ m CO<sub>2</sub> band. Under the approximation of equation

## J. R. DRUMMOND AND OTHERS

(4) the two signals will be  $\epsilon_w B(\overline{T}_w)$  and  $\epsilon_p B(\overline{T}_p)$  where the subscripts w and p respectively denote the wideband and pressure modulator channels. Now it is possible to choose the mean cell pressure in the pressure modulator channel such that over a certain range of altitudes the relative contributions of different segments of the atmospheric path being observed are very similar for the two channels, even though the effective emissivity of the atmosphere for the two channels is very different (figure 4). This means that  $\overline{T}_w$  and  $\overline{T}_p$  will be almost equal. We then find that the ratio  $\epsilon_p/\epsilon_w$  is almost independent of temperature but strongly dependent on the pressure at the level of observation. Figure 5 illustrates that, for a particular cell pressure, atmospheric pressure over the range between 0.8 mbar and 20 mbar may be derived from the ratio of signals in the two channels. A first order correction for the atmospheric temperature profile can be included, enabling pressure measurement of adequate accuracy (ca. 2%) to be achieved. In § 6 an example of the use of the method with s.a.m.s. data is presented.

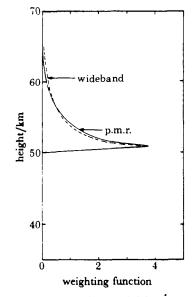


FIGURE 4. A typical pair of CO<sub>3</sub> 15  $\mu$ m weighting functions for the s.a.m.s. attitude-determining channels before the finite field-of-view is accounted for. The tangent height of the line-of-sight is 50 km.

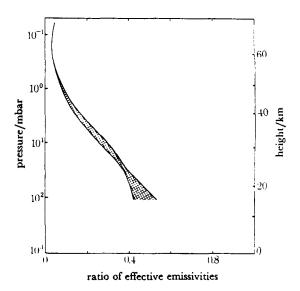


FIGURE 5. Ratio of p.m.r. to wideband emissivities for the 15  $\mu$ m CO<sub>2</sub> channels C1 at 15 mbar mean cell pressures. The envelope shown is the extreme of a series of curves derived using climatological temperature profiles at 10° latitude intervals from Newell (1977). The effect of the finite instrument field-of-view has also been included in the calculation.

## 3.4. Wind determination

Through the incorporation in the instrument of cells containing the gases whose emission is being detected, the possibility exists of measuring atmospheric motion. If the emitting gas in the atmosphere possesses a significant velocity along the line of sight relative to the gas in the absorbing cells, the emission lines will be shifted relative to the absorbing lines. Figure 6 shows the variation of signal with azimuth angle for a particular case. The direction along which the signal is a maximum is that along which there is zero Doppler shift and hence zero relative velocity of the atmosphere with respect to the spacecraft. If the position of this maximum is

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

determined, knowing the spacecraft velocity vector, and making allowance for the rotation of the Earth, the wind speed along the line of sight can be found. For directions nearly perpendicular to the spacecraft's velocity vector, a change in azimuth direction of 0.1° introduces ca. 10 ms<sup>-1</sup> of the spacecraft's velocity along the line of sight.

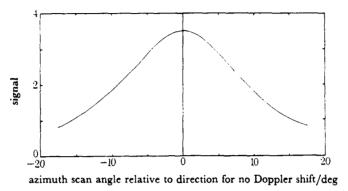


FIGURE 6. Variation of the signal from the 15  $\mu$ m CO<sub>2</sub> p.m.r. channel with azimuth scan angle relative to the central point where there is no relative motion along the line-of-sight between the atmosphere and the instrument.

So that the influence of atmospheric variations is a minimum, the azimuth scan is carried out at such a rate that during the scan the same part of the atmosphere is kept under observation. A scan over  $30^{\circ}$  (figure 6) then takes about 240 s. The main problem in carrying out the measurement is that of the maintenance of adequate attitude stability during the period of scan so that the slope of the curve of signal against scan angle is not disturbed by attitude changes. With the Nimbus 7 spacecraft the attitude stability is not adequate for accurate wind determination, although it is hoped that the stability will sometimes be sufficient for the lower mesospheric jet at *ca*. 60 km altitude to be detected.

field of view	channel	constitu <del>e</del> nt (gas in modu- lator)	mean cell pressures mbar	spectral band µm	derived quantities and altitude range
A	A1 A2	CO	17, 2.4, 0.69, 0.25	15	kinetic temperature 15–80 km; attitude $\nu_2$ vibrational temperature 80–100 km
	A3	CO	14.8, 4.5	4–5	distribution 15-60 km
	A4	NO	45, 20		distribution 15-00 km
В	B1 B2	H <sub>1</sub> O	16, 4.5, 0.8, 0.5	$2.7 \cdot 25-100$	distribution 80–100 km distribution 15–100 km
С	C1 C2 C3	CO2 N2O CH4	36, 11.2, 3.25, 0.87 24.4, 7.15 47.8, 22.5	15 7.7	kinetic temperature 15–80 km; attitude distribution 15–60 km

### TABLE 1. S.A.M.S. RADIOMETRIC CHANNELS

#### 3.5. Instrument performance

S.a.m.s. contains a number of channels for various measurements. Some of these channels may be operated simultaneously and some are mutually exclusive. A list of all channels is shown in table 1. Each entry in the constituent column corresponds to a modulator and each entry in

Vol. 296. A

15

## 226 J. R. DRUMMOND AND OTHERS

the spectral band column to a detector. Since it is only possible to pass one pair of signals (p.m.r. and wideband) through a detector at any one time only one modulator signal using a particular detector may be processed at any time. Therefore, only one of the gases  $N_2O$  and  $CH_4$  may be measured at a particular time and only one of CO, NO and 4.3  $\mu$ m CO<sub>2</sub>. A detailed description of how the channels are fitted together in the instrument is given in § 4.

gas	CO3	NO	N <sub>3</sub> O/CH <sub>4</sub>	CO3	CO2	$H_2O$
band/µm	4.3	5.3	7.7	15	15	25-100
altitude of observation/km	120	40	40	50	100	50
approximate integration time for n.e.e. = 0.03† as measured by s.a.m.s. on Nimbus 7 /s	10000	<b>400</b>	10	4‡	2000	5

TABLE 2.	INSTRUMENT	PERFORMA	NCE FOR	TEMPERATU	RE AND COMPO	SITION
:	MEASUREMEN	TS FROM P	RESSURE	MODULATOR	CHANNELS	

 $\dagger$  n.e.e. = noise equivalent emissivity. An n.e.e. = 0.03 implies a S/N of 33 when instrument observing black body at atmospheric temperature.

 $\pm$  n.e.e. = 0.01 as this is a temperature sounding channel required to measure to  $\pm 1$  K.

It can be seen from figure 2 that over a large part of the sensitive region for composition sounding a 10% change in composition produces a change of about 3% in the observed signal. In table 2 some typical integration times required to detect such a change are tabulated for various channels in the instrument at typical heights in the atmosphere.

### 4. THE S.A.M.S. INSTRUMENT

The s.a.m.s. radiometer instrument consists of two modules: the sensor housing containing the optics, mechanisms and detectors, and a separate electronics module. The overall size of the sensor (see figure 8) is  $55 \text{ cm} \times 30 \text{ cm} \times 55 \text{ cm}$  and it weighs 23.6 kg. The electronics module weighs 6.7 kg. The average power consumption is about 20 W.

There are a total of twelve radiometric channels: six pressure modulator and six wideband channels distributed between three fields-of-view denoted A, B and C. The relation between the three fields is fixed by the geometry of the instrument optics and is illustrated in figure 7. Some of the channels may be used to sense one of several constituents selected by programmable control logic within the instrument (see § 4.7). Figure 7 also shows the scanning range of the instrument for limb views and space calibration.

#### 4.1. Optical system

The optical system may be conveniently subdivided into two sections:

(a) the front optics comprising scan mirror, telescope, calibration system and wideband chopper;

(b) the rear optics comprising field separator mirror, pressure modulator cells, detectors and optical relay components. The front optics is shown in figure 9. Incident radiation is first reflected from the scanning mirror (M1), which can rotate about two axes for both limb

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

(vertical) and azimuth (horizontal) scanning (see § 4.5). It is then reflected from the main paraboloid (M2) to the primary focus forming a telescope of  $177 \text{ cm}^2$  area and 22 cm focal length.

During in-flight calibration a small black body is rotated into the beam at the primary focus to give a known input radiance to the rest of the optical system.

After the primary focus the beam is directed by ellipsoid (M3) and plane (M4) mirrors through the fast, shallow chopper (see § 4.4) to the secondary focus at the field separation mirror

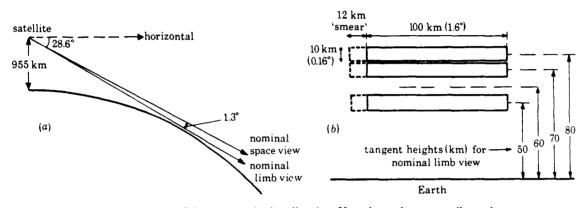


FIGURE 7. (a) The geometry of the s.a.m.s. viewing direction. Note that only very small angular moments are necessary to scan the entire atmosphere. As drawn the satellite velocity is into the paper. (b) The relative positions of the three fields-of-view A, B and C viewed along the direction of the line of sight in the 'nominal limbview' position. The horizontal smear is caused by spacecraft motion during a 2 s measurement period.

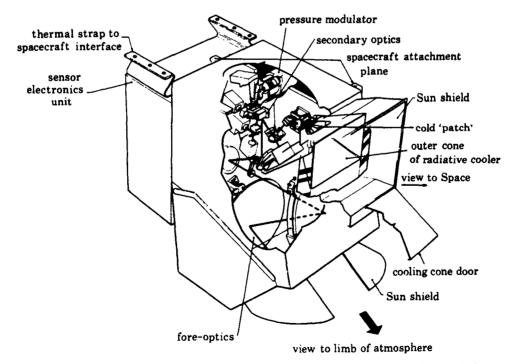


FIGURE 8. Cutaway diagram of the s.a.m.s. sensor unit showing the relative orientation of the optical components, pressure modulators and radiative cooler. The whole unit is located on the sensor ring below the satellite.

15-2

228

## J. R. DRUMMOND AND OTHERS

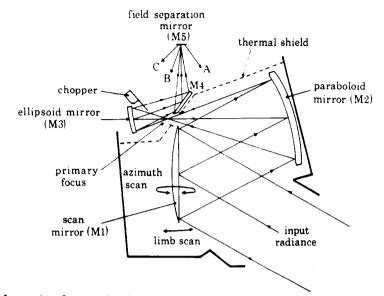


FIGURE 9. The fore-optics of s.a.m.s. in plan view. The calibration black body (not shown) is inserted into the beam at the primary focus. The thermal shield isolates the secondary optics compartment at 25 °C from the front optics compartment at 0 °C.

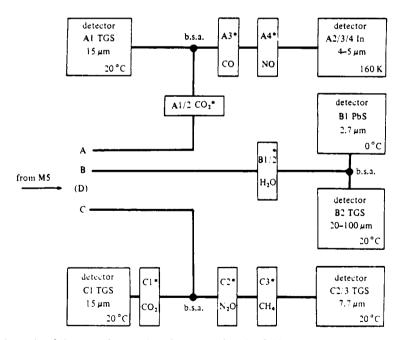


FIGURE 10. Schematic of the secondary optics of s.a.m.s. after the field-separation mirror (M5). The dichroic beamsplitters are marked b.s.A, b.s.B and b.s.C. Secondary mirrors relay lenses and other optical components have been omitted for clarity. The beam marked (D) corresponds to a facet of M5 that reflects the input beam back through the front optics and is used for optical alignment. The pressure modulator cells are identified by an asterisk (\*).

## APPENDIX C

### STRATOSPHERIC AND MESOSPHERIC SOUNDER

(M5). The use of an ellipsoid in the M3 position cancels some of the aberrations of the off-axis paraboloid (M2).

Thermal isolation of the compartment containing the first two mirrors, i.e. those before the primary focus and calibration point, allows these mirrors to be run at a somewhat lower temperature  $(0 \, ^{\circ}C)$  than the rest of the instrument, thus reducing the thermal strays from this area and increasing the reliability of the calibration.

The rear optics are shown in the correct spatial orientation in figure 8 and in schematic form in figure 10. Field separation occurs at three of the four facets of the mirror M5 which are angled to direct parts of the beam into one of three separate optical chains, and curved to image M3 on the detector aperture stop. Each of the fields so defined is rectangular and subtends an angle in the instrument line-of-sight  $0.16^{\circ} \times 1.6^{\circ}$ , a solid angle of  $7.8 \times 10^{-5}$  sr (figure 7).

Broad spectral selection for the different molecular bands is achieved by dichroic beamsplitters (b.s.A, b.s.B, b.s.C) and filter packs of up to three separate components on the various detectors. The pressure modulator windows are either anti-reflexion coated or are made of low refractive index materials. In all there are about 40 individual optical components in the system after M5. High optical transmission in the system is maintained by the use of anti-reflexion coatings. As an example, the overall transmission of the C2/3 chain with 19 surfaces is still 67 % at band centre.

The rear optics, modulators and detectors are mechanically and, except for the 2.7  $\mu$ m PbS and 4-5  $\mu$ m InSb detectors (B1 and A2/3/4), thermally attached to the main optics plate made of aluminium 1 cm thick giving good thermal uniformity and capacity as well as mechanical stability.

### 4.2. Detectors

There are six detectors in s.a.m.s., four triglycine sulphate (TGS), one lead sulphide (PbS) and one indium antimonide (InSb). Each detector flake is imaged by its optical system on to one of the facets of M5 and, therefore, the aspect ratio of the detectors is the aspect ratio of the field of view. Each detector is  $3.12 \text{ mm} \times 0.312 \text{ mm}$ , giving an area of approximately  $1 \text{ mm}^2$ .

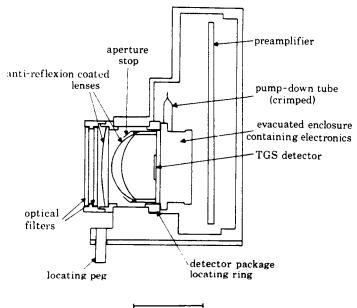
The detectors are each assembled and tested as sub-assemblies consisting of detector, preamplifier, condensing optical system and filters before being integrated with the rest of the instrument. A TGS sub-assembly is shown in figure 11. The condensing optical doublet, consisting of a meniscus and a plano-convex lens (see Murray 1962), is made of anti-reflexion coated germanium (except in the case of the 25-100  $\mu$ m H<sub>2</sub>O channel (B2) where silicon is used). The package consisting of the meniscus lens detector and mounting flange is evacuated to protect the detectors from moisture during ground operations.

The sub-assemblies for the cooled detectors are more complicated than those for the TGS because of the need to isolate the detector thermally while mechanically attaching it to the optics plate. The B1 and A2/3/4 detectors are cooled radiatively, the B1 to 0 °C by a radiator in the fore optics compartment and the A2/3/4 to 160 K by a separate radiative cooler (see  $\S$  4.6).

### 4.3. Pressure modulator assemblies

The pressure modulators used in s.a.m.s. are a development of those used in the Nimbus 6 p.m.r. instrument (see Curtis *et al.* 1974). The pressure of the gas in the optical cell is altered by means of a piston oscillating at between 25 and 40 Hz in a connecting volume (figure 12). A pressure variation over the cycle of the order 2:1 is achieved in this design, the exact value depending on the dimensions of the particular optical cell and the mean cell pressure. The principal

## J. R. DRUMMOND AND OTHERS



20 mm

FIGURE 11. S.a.m.s. detector sub-assembly (TGS). Because the impedance of a TGS detector is very high and the signal small, part of the preamplifier is built into the back of the detector package within the evacuated enclosure. Precise optical alignment is essential, both of the package to the sub-assembly and of the sub-assembly to the secondary optics mounting plate. In the former case a close-fitting ring is used and in the latter a precision locating peg.

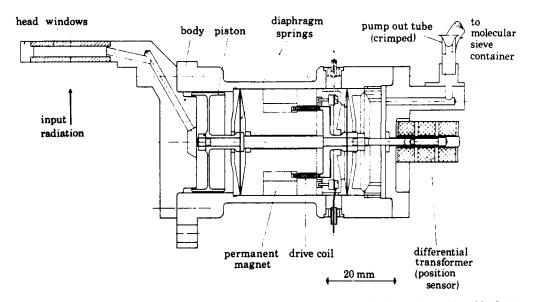


FIGURE 12. Cross section of a p.m.r. assembly. The body of the PMR is made of titanium with gold 'O' ring seals at each end. The support springs and coil lead-outs are made of beryllium-copper. The thermostatted container of molecular sieve is not shown.

#### APPENDIX C

### STRATOSPHERIC AND MESOSPHERIC SOUNDER 231

changes from the devices used in the Nimbus 6 instrument are a reduction of size by a factor of approximately 2 and the use of a moving coil drive leading to better amplitude control and electrical efficiency. Other improvements have made the modulators easier to manufacture and assemble. It is necessary, as was observed in § 3.2, to alter the pressure of gas in the cell in flight, the pressure selected depending on the altitude range being scanned. This is achieved in the s.a.m.s. modulators as follows. The gas in the cell is in equilibrium with a much larger amount adsorbed on a zeolite molecular sieve material in a reservoir (figure 12). In this system the mean pressure of gas is determined by the temperature of the sieve, and provision has therefore been made for programmable temperature control of the sieve container. Each modulator is provided with two or four pressure settings, which may encompass a range of over 20:1, corresponding to a thermostat temperature range of 30-120 °C. These settings are selected by the program control logic. The adsorption and desorption of gas from the sieve is entirely reversible and the system provides good pressure stabilization. Measurement of the oscillation frequency, which changes rapidly with pressure, is used as the primary method of determining cell pressure since slow variations in gas amount, such as might be caused by a leak in a modulator, could not be detected by monitoring the sieve temperature.

#### 4.4. Fast chopper assembly

The fast chopper used on s.a.m.s. only modulates 2.5 % of the beam for reasons described above (see § 2.3). It consists of two etched, black-painted copper grids, mounted one above the other. One grid is fixed to the mount but the other is free to move on the end of a flat cantilever leaf spring. Oscillation, at the mechanical resonant frequency of the spring-grid assembly, is maintained by an electronic loop using redundant pairs of piezoelectric ceramic plates as both drive and position sensors. Oscillation is at a frequency of approximately 245 Hz with a peakto-peak amplitude of 2 mm.

### 4.5. Scan mirror assembly

The scan mirror must be rotated with high positional accuracy in two mutually perpendicular directions to produce limb and azimuth scanning. This is realized by supporting the mirror in two frames, one within the other, for the two scan directions. Each mirror drive consists of a lead-screw directly driven by a 45° stepping motor. Coupling to the mirror is by means of a recirculating-ball nut, which is lubricated by a lead coating. The axis of rotation of the mirror is defined by a pair of cross-leaf springs (flexipivots). The step sizes of the drive system, expressed in terms of motion of the line-of-sight, are 4.8 arc min in the limb direction and 7.3 arc min in azimuth.

Mirror position is independently sensed in both directions by a linear variable differential transformer (l.v.d.t.), actuated directly by the mirror, and associated electronics. The output of the limb l.v.d.t. is digitized with a resolution of 0.045 arc min and the azimuth l.v.d.t. to 1.8 arc min. It is therefore possible to monitor the limb position of the mirror to a very much higher accuracy than the step size. The limb position is under full control of the program control logic (p.c.l.) (see § 4.7), whereas the azimuth scan can only be programmed as a continuous scan or stepped.

Limit detectors, both optical and mechanical, are provided for resynchronizing the control logic and preventing damage under fault conditions.



### J. R. DRUMMOND AND OTHERS

#### 4.6. Radiation cooler

The indium antimonide detector is cooled by a two stage radiation cooler, the first highcapacity stage forming a cone around the second stage as shown in figure 13. The first stage cools the front lens of the optical doublet (see § 4.2) as well as pre-cooling the rear lens and detector. In orbit, the inner and outer cooler temperatures are typically 160 K and 190 K.

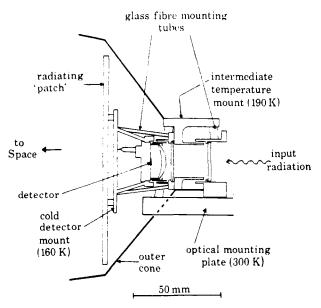


FIGURE 13. The s.a.m.s. cold detector assembly (InSb). The mount is in three parts: the detector mount at 160 K, the intermediate lens mount at 190 K and the mechanical interface at 298 K. The parts are interconnected by 8 glass fibre tubes, some of which may be seen in the diagram. The relative sizes of the 'radiating patch', which is black-painted, and the outer cone, which is coated with vapour deposited aluminium, can be seen in figure 8.

The cooler has been designed to minimize the effects of contamination on the optics as the spacecraft outgases; however, provision has been made to heat the detector and optics electrically to remove the debris as necessary. The cooler surfaces were protected during launch and the first 3 weeks of orbit by a door which was opened on 13 November 1978. The door cannot be closed again and will remain open for the duration of the mission.

#### 4.7. Electronics

The electronics of s.a.m.s. is considerably more complicated than that of previous instruments built at Oxford for three reasons:

(a) there are more signal channels (twelve in all) and modulator drives (seven);

(b) the increased number of signal channels and drives means that more care must be taken to ensure that coupling and cross-talk, both electrical and mechanical, is minimized.

(c) The scan and calibration control system is considerably more flexible than previous instruments allowing in-flight programming.

The major changes in construction details since previous instruments are the use of integrated circuit amplifiers in the linear electronics and CMOS in the digital electronics. A full block diagram of the system is shown in figure 14.

APPENDIX C

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

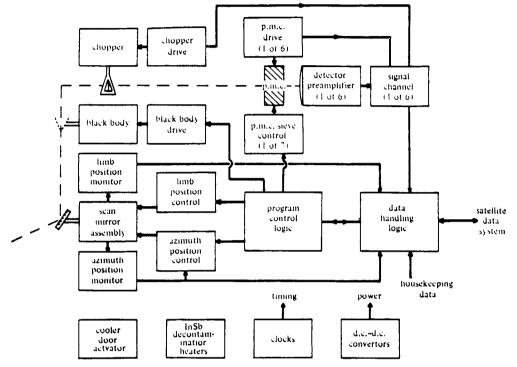


FIGURE 14. Block diagram of the s.a.m.s. electronics subsystems.

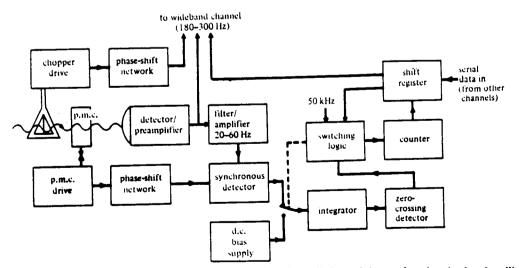


FIGURE 15. Block diagram of one half of a signal channel. Data from all channels is transferred to the data handling system in a single serial stream.

234

### J. R. DRUMMOND AND OTHERS

### (a) Signal channels

The signal processing electronics follows the pattern of the system used for a previous balloon radiometer (see Chaloner *et al.* 1978). The high and low frequencies (wideband and pressure modulated signals respectively) present at the detector output are processed by separate channels as shown in the block diagram in figure 15. The output integrator of each channel is also made to act as a dual-slope analogue-to-digital converter. This is achieved by integrating the signal, plus a small offset to ensure that the result is positive, for 1.8 s, and then integrating a fixed negative signal until the result is zero. By timing the second integration period with the use of the 50 kHz spacecraft clock, a resolution of 1 part/10<sup>4</sup> is obtained in the remaining 200 ms of the 2 s period. Only 12 bits are actually telemetred to the ground but the 13th can be inferred, if necessary, by context. The 200 ms period while the signals are evaluated is also used to step the scan mirror.

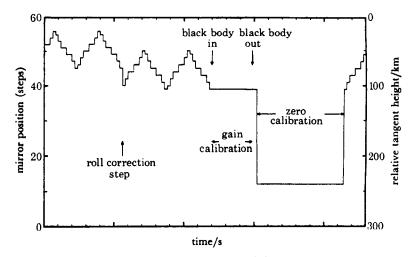


FIGURE 16. A typical scan program. There are 18 scans (up and down ramps) between calibration sequences leading to an interval between calibrations which is not synchronous with the orbital period. Roll correction is applied independently, synchronous with the orbital period.

## (b) Pressure modulator drives

The drive circuits for the pressure modulators are considerably more complex than those used for the Nimbus 6 p.m.r. or the balloon instrument. The complexity is necessary because of the need to maintain a constant amplitude despite large internal pressure variations, which in turn lead to large changes in frequency and power requirements. The use of a moving coil drive, with its inherently tighter coupling than the previous moving magnet drive, assists with amplitude control and also allows active damping to be employed when a modulator is inhibited. Active damping is useful in multiple channels with only one detector (e.g. A2/3/4) since unwanted mechanical energy could otherwise be coupled into the 'off' modulators and appear as unwanted modulation in the detector output. This can be prevented by switching the drive from positive (oscillating) to negative (damping) feedback when the modulator is inhibited.

## APPENDIX C

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

### (c) Program control logic (p.c.l.)

The extremely large number of possible operating modes for s.a.m.s. using various combinations of modulators, scanning ranges and mean cell pressures, and the need to correct the scanning sequences for any predictable orbital roll of the spacecraft, made a programmable operating mode very desirable. The system used operates the scan mirror, modulators, molecular sieves, thermostats and calibration black body by reference to a stored 'program' in a  $128 \times 10$  bit memory which may be reloaded from the ground. The instruction set of the p.c.l. also includes jump, wait, synchronizing and timing instructions to allow program loops to be constructed which operate in a fixed relationship to the spacecraft data system. The variations of scanning and calibration are, therefore, infinite and it is possible to adapt the scanning sequence to partially compensate for orbital roll. It is also possible to scan outside the atmosphere, for example, to locate the moon for calibration purposes.

A typical scan sequence is shown in figure 16. This program is designed to scan a height range of 55 km in 10 km steps with regular calibration sequences not related to the orbital period. A roll correction step can also be seen.

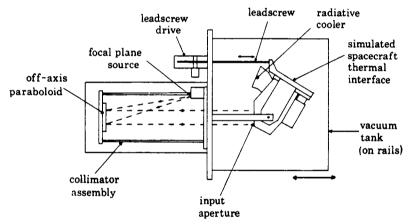


FIGURE 17. S.a.m.s. environmental test chamber at Oxford. Baffles, cooled by liquid nitrogen, surround the instrument to simulate radiative loss to space but these have been omitted for clarity. Access is provided by rolling the main vacuum tank away on rails. The collimator assembly is shown in position. The black body assembly (see text) may be substituted for the collimator when required.

## 5. VERIFICATION AND CALIBRATION PERFORMANCE

In addition to the tests that were necessary to prove the design of the instrument and the environmental tests specified by N.A.S.A., s.a.m.s. has undergone the following series of measurements intended to facilitate the interpretation of data when in orbit.

(a) The response of the instrument to known black body radiance inputs has been measured, and a calibration scheme devised to take account of variations in operating conditions.

(b) The response of the instrument to emission from gases in the atmosphere has been determined, using a simulated atmospheric path, in order to verify the spectroscopic calculations by which atmospheric quantities are deduced.

(c) The field of view profiles of all the channels and their interrelation has been measured.

(d) The scan mirror position telemetry (see § 4.5) has been accurately calibrated in terms of the

236

#### J. R. DRUMMOND AND OTHERS

instrument line-of-sight. A combined vacuum and thermal environment chamber at Oxford provides correct simulation of the spacecraft interface and orbital conditions. Various additional radiation sources are also available within the chamber for calibration purposes (see figure 17).

#### 5.1. Radiometric calibration

The response of the instrument is determined for a range of accurately known black body radiances in order to

(a) verify that the instrument response is linear and has the correct gain and noise levels;

(b) check that the spectral selection of the optics is correct and that there are no unknown spectral 'leaks' in the overall response;

(c) provide a basis for the model of the instrument used in the in-flight calibration scheme;

(d) determine the magnitude of any stray responses, particularly those which depend on either azimuth or limb scan angle. The source used for this test is a large, baffled black-body cone, completely filling the instrument aperture, whose temperature may be held constant and uniform at any temperature between 77 K and 320 K. The input radiance is, therefore, very precisely calculable. There are in fact two sources, one on the instrument axis viewed with the azimuth scan centralized and one off-axis viewed at the limit of the azimuth scan. By using both targets, the variations of strays with azimuth scan angle may be established.

Variations with limb-scan angle are measured by tilting the whole instrument using a precision lead-screw, thus varying the relative position of the on-axis target and the instrument.

The sensor temperature is held constant during measurements of stray radiation but may be varied or cycled over a considerable range (-5 to +40 °C), in order to dynamically evaluate the temperature coefficients of the response. By simulating the orbital temperature variations of the spacecraft experience is gained of the likely instrument temperature variations in orbit – a verification of the sensor thermal design.

### 5.2. Response of p.m. channels to atmospheric paths

The response of the channels to paths of gas is checked by measurment of the transmission of simulated atmospheric paths in the laboratory, using either the actual sensor modulators, filters and detectors before final assembly of devices identical in design and construction. The apparatus consists of a multiple-pass absorption white cell which provides paths of up to 10 m, which may be used at temperatures down to 240 K (see Houghton & Taylor 1973). Transmissions of isothermal paths of the required absorber or absorber/nitrogen mixtures measured with this equipment as a function of pressure are then compared with the response calculated from line data, and information from both sources is combined in the transmission model used in the final attitude, temperature and composition retrievals.

## 5.3. Field of view

The attitude of the sensor line of sight is determined in orbit from the ratio of the pressure modulator and wideband signals in the  $CO_2$  channels (see § 3.3). The line-of-sight of the other channels, therefore, needs to be measured with respect to these. The details of the field profiles, especially the shape of the lower edge, also critically affect the weighting function of the channels and must also be determined.

For field of view determination the black cones are replaced by a collimator (see figure 17. This consists of an off-axis paraboloid mirror, with a narrow slit illuminated by a hot ceramic

### STRATOSPHERIC AND MESOSPHERIC SOUNDER

filament in the focal plane. The lead screw system is used to rotate the s.a.m.s. sensor through the parallel beam of radiation from the collimator at a precisely known rate, and the field response is thus determined.

The collimator assembly and leadscrew are also used for various optical checks and in particular for the calibration of the l.v.d.t. system used for limb scan angle telemetry. The lead screw is driven slowly and continuously while the p.c.l. is programmed to step the s.a.m.s. scan mirror so that the edge of the field of view of a selected channel passes through the collimated beam once per step. By matching the respective points on the field response, the l.v.d.t. system is calibrated directly in terms of the line-of-sight. This test and the primary field of view determination are carried out at several temperatures over the desired operating range.

#### 5.4. In-flight calibration

A two point calibration is provided in flight as follows.

(a) The scan mirror may be programmed to move so that the tangent height viewed in the atmosphere is sufficiently high that the atmospheric radiance is insignificant ('space' view). This is arranged in operational programs such that the minimum tangent height for the lowest field of view (the C channels) is at least 150 km.

(b) The black body (see § 4.1) may be placed in the optical path to provide a standard radiance at approximately 295 K. This is also under the control of the p.c.l.

In practice the estimate of the space view signal is much more important than the measurement of gain, as many of the measurements made are of small differences from the zero radiance level. Operational p.c.l. programs are therefore arranged to have a higher proportion of space view calibration than of black body view as shown in figure 16.

In flight the zero offset and radiometric gain are calculated from instrument housekeeping data by means of a model, the parameters of which have been obtained by combining the results of ground tests and in-flight calibration measurements.

### 6. PRELIMINARY RESULTS

Calibrated radiance data from s.a.m.s. is processed in the following sequence. A preliminary estimate is obtained of the tangent pressure of the line-of-sight of all three fields-of-view at all times using the method described in § 3.3. Then the atmospheric temperature profile and a better estimate of the tangent pressure are derived from the preliminary estimate of tangent pressure and the  $CO_3$  channel radiances. Finally, this information is used, together with radiances from the other channels, to retrieve constituent profiles. These three steps are described in some detail below.

#### 6.1. Attitude determination

As has already been shown in § 3.3 the tangent pressure of the line-of-sight of the  $CO_2$  channels can be determined from the ratio of the wideband to p.m.r. radiances assuming that the tangent height is within a suitable range (approximately 20-50 km). Since the space-craft roll rate is reasonably low and the orbital period long compared to the rate of sounding (1 every 2 s), smoothing along the orbit path, i.e. in time, is applied using a linear sequential estimator (Kalman-Bucy filtering, see Rodgers 1976). To achieve this the tangent pressure is converted, using a climatological temperature profile and the scan mirror position, to the tangent height of a reference scan mirror position. This tangent height estimate is then smoothed.

238

#### J. R. DRUMMOND AND OTHERS

The estimator is stepped forwards along the orbit and for each new observation a smoothed estimate of the tangent height and its uncertainty are calculated given the observation and its error, the previous smoothed estimate and its error, the variation in scan mirror position measured by the l.v.d.t., the time between observations and a knowledge of the likely maximum roll rate. The process is repeated backwards along the orbit and the two results at each observation point are combined to derive a best estimate of the roll angle at any point. The smoothed roll-angle is then reconverted to tangent pressure for each field of view. The results for a typical orbit are shown in figure 18.

The smoothing process is capable of dealing automatically with missing data caused either by bad soundings or a tangent height out of the range where attitude determination is possible. Prolonged loss of information produces a steadily increasing uncertainty in the attitude.

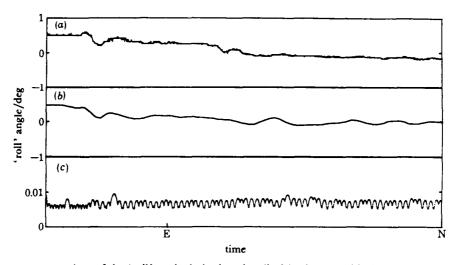


FIGURE 18. A comparison of the 'roll' angle derived as described in the text with the output of the spacecraft attitude control system (for part of orbit 135 of Nimbus 7). For reasons described in the text these are not necessarily comparable but short timescale features should be reproduced. The error in the derived roll error is in a scale magnified  $100 \times$  relative to the other two plots. E = equator, N = most northerly latitude of orbit. (a) Spacecraft attitude control system roll output; (b) s.a.m.s. derived 'roll'; (c) s.a.m.s. derived 'roll' error.

#### 6.2. Temperature retrieval

The short wave channels are particularly sensitive to atmospheric temperature (5%) per degree for the 5 µm channel) so accurate temperature retrievals are of particular importance for making good constituent measurements. It is also possible to improve further upon the attitude determination described above which assumed climatological temperature profiles. Temperature profiles and an improved attitude are, therefore, retrieved simultaneously.

Since there are a large number of possible scan patterns in use, the retrieval method must be able to handle any pattern selected and must, therefore, be very flexible. Were the instrument quickly scanning up and down through the whole atmosphere, complete sets of radiances for a scan could be assembled and the retrieval performed by simpler methods. But there is a need to incorporate observations when and where they occur, essentially at random, and a sequential estimator is again used. For economy of computation the retrieval solution is represented by coefficients of 12 eigenfunctions, giving the deviation of source function at 15 µm from the

## STRATOSPHERIC AND MESOSPHERIC SOUNDER

climatological mean (interpolated linearly from monthly means at 10° latitude intervals). They are eigenfunctions of the covariance matrix

$$C_{ij} = \exp\left(-(z_i - z_j)^2/\alpha^2\right) \frac{\mathrm{d}B(z_i)}{\mathrm{d}T} \frac{\mathrm{d}B(z_i)}{\mathrm{d}T},$$

where  $z_i$  is the scale height of level *i* and  $dB(z_i)/dT$  is the differential with respect to temperature at level  $z_i$ . This is the Planck function covariance resulting from a set of wave perturbations in

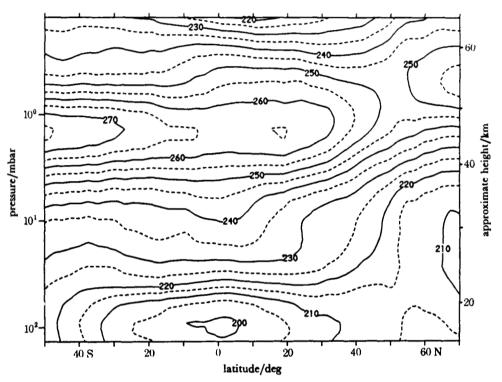


FIGURE 19. Zonal mean temperature cross section for 6 November 1978 derived entirely from Nimbus 7 s.a.m.s. data.

the vertical with a Gaussian spectrum and temperature amplitude independent of height. Using the estimate of attitude derived as above, the deviation of observed radiance from its climatological value can be represented linearly in terms of the solution which consists of the eigenfunction coefficients and the attitude. By these means a linear sequential estimator can be employed in a similar manner to the attitude determination described above. In the absence of input data the solution is made to decay exponentially with time towards zero (i.e. towards its climatological value) so that when there is no information about some or all levels, any deviations for those levels are slowly forgotten. When run forwards along the orbit the estimator is recursive and one-sided, using just past and present data to give the solution. However, an independent set of estimates is obtained by running the estimator backwards along the orbit. At  $2\frac{1}{2}$ ° latitude intervals along the orbit both estimates are combined in accordance with their uncertainties and the result retained for derivation of temperatures and analysis on latitudelongitude grids. A typical zonal mean cross section is shown in figure 19.

#### J. R. DRUMMOND AND OTHERS

#### 6.3. Concentration retrievals

Using the derived tangent heights and temperatures it is now possible to obtain concentration profiles from the data. Figure 20 shows some typical results for the  $H_2O$  rotation band (channel B2) in both the p.m.r. and wideband cases.

The solid lines were calculated using the temperature data derived using the methods described above, water vapour band data from McClatchey *et al.* (1973) and Chaloner (1976) and a computer program to simulate atmospheric signals due to Eyre (1978).

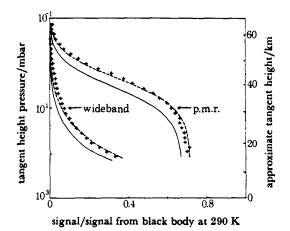


FIGURE 20. Water vapour signal from s.a.m.s. (on 6 November 1978 for latitudes between 20° N and 20° S) compared with signals calculated for constant volume mixing ratio profiles. It can be seen that the p.m.r. signals indicate a slight fall in mixing ratio from 50 to 30 km which becomes pronounced in the wideband signal between 30 and 18 km where there are uncertainties due to clouds and other effects associated with the tropopause. +, S.a.m.s. data; --, 5 parts/10<sup>6</sup> by volume; -, 2 parts/10<sup>6</sup> by volume.

Because the data is of a preliminary nature, a profile of concentration versus altitude is not presented. However, it can be seen that the results are inconsistent with a constant mixing ratio profile but would be consistent with a slow rise starting at 20 km levelling off above 35 km. This type of behaviour would be expected if the photolytic destruction of  $CH_4$  were a source of  $H_2O$  in the middle stratosphere (McConnell *et al.* 1971).

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## STRATOSPHERIC AND MESOSPHERIC SOUNDER

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Vol. 296. A

241

**16** 

## **RETRIEVAL OF TEMPERATURE AND COMPOSITION FROM NIMBUS-7 SAMS MEASUREMENTS**

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## Retrieval of Temperature and Composition From NIMBUS 7 SAMS Measurements

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A general purpose stochastic sequential estimator is described, which allows the inverse problem of remote sounding to be solved in an optimum manner. The use of a sequential estimator allows continuity in the horizontial or in time to be included, making the a priori estimate close to the final solution, so that iteration is not required. The scan pattern in which measurements are made is not used and need not be known, so that the problem of spacecraft attitude stability is minimized. The method gives both a solution profile and its error covariance. It is applied to both temperature and composition sounding from a limb scanning instrument, the SAMS on NIMBUS 7.

#### 1. INTRODUCTION

The Stratospheric and Mesospheric Sounder (SAMS) on NIMBUS 7 is described in detail by *Drummond et al.* [1980]. It is designed to measure temperature and concentration profiles of various gases in the height range of 15-100 km by detecting thermal emission or, in some cases, resonantly scattered sunlight. It obtains information about vertical distribution by scanning the limb and uses the technique of gas correlation by pressure modulation to select the required spectral signal from the gases being measured. A summary of the spectral channels used and the characteristics of these channels is given in Table 1. There are three fields of view, A, B, and C, separated in the vertical, with A highest in the atmosphere. All the fields of view are approximately Gaussian, with a vertical width of about 8 km.

In the technique of pressure modulation, a cell of the gas whose emission is to be measured is used as a filter, in conjunction with an interference filter to select the appropriate spectral band. The pressure in the cell is modulated, and the transmitted radiation is detected at the modulation frequency to give a measurement of radiation originating near the line centers, integrated over the whole spectral interval. By chopping the beam with a black chopper at a different frequency, a parallel measurement of emission over the whole band (excluding the line centers) is obtained. These measurements are known as pressure modulated (PM) and wide band (WB), respectively.

The radiance  $L_{i}(h)$  received at the instrument from the limb (Figure 1) can be written

$$L_{v}(h) = \int_{-\infty}^{\infty} B[v, T(x)] \frac{d\tau}{dx} (v, x) dx$$

where h is the tangent height of the line of sight, B[v, T] is the Planck function at wave number v and temperature T, and  $\tau(v, x)$  is the transmittance from position x to the instrument.

The retrieval (inverse) problem divides into two parts, temperature and composition. The temperature profile is determined from the two 15  $\mu$ m CO<sub>2</sub> PM channels and the two corresponding WB channels in the A and C fields of view which are separated by 0.56<sup>2</sup> (about 35 km). The composition profile is determined from radiances in the remaining channels but requires the temperature profile as input.

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The limb scanning temperature retrieval problem has some complications that make it less linear than its nadir sounding counterpart. The most serious complication is that the attitude of the spacecraft is not adequately measured or controlled, so that the tangent height of the line of sight is not known accurately enough to compute the corresponding transmittance function. However, there is enough information in the radiances to determine both temperature profile and tangent height. The composition retrieval is nonlinear because the unknown enters the equation of transfer nonlinearly as an argument of the transmittance function.

Our approach to both retrieval problems is to linearize the direct equation about an a priori profile and to use a statistically optimum estimator [Rodgers, 1976] to find the most likely solution. The question of iteration is avoided by ensuring that the a priori profile is sufficiently close to the solution that the linearization is adequate. In the case of temperature this is achieved by using a sequential maximumlikelihood estimator, in which the retrieval corresponding to time t, and its error covariance matrix, are used to construct the a priori and its error covariance at time  $t + \Delta t$ , where  $\Delta t$ (typically 2 s) is the time interval between measurements. For composition retrieval it has proved necessary to average the data, for example, zonally, in order to improve the signal to noise ratio and to reduce computing time. In this case, the previous day's retrieval for the same averaging box is used to construct the a priori.

#### 2. RADIANCE EMITTED FROM A LIMB PATH

To simplify the equations, we will consider first the case of an infinitesimal field of view. The radiant flux  $\Phi(h)$  emitted by a limb path with a tangent height h and reaching the detector of the radiometer may be written as

$$\Phi(h) = A\Omega \iint F(v)B[v, T(x)] \frac{d\tau_k}{dx}(v, x) dv dx \qquad (1)$$

where A is the aperture of the telescope.  $\Omega$  the (infinitesimal) instrument field of view, F(v) the spectral response of the channel in question (including the modulated component of the cell transmittance in the case of PMR channels), and  $\tau_{s}(v, x)$  is the monochromatic transmittance of the atmosphere at the frequency v from the satellite to the location described by the coordinate x. Provided that the molecular bands being observed are in local thermodynamic equilibrium at the atmosphere at the transmittance of the satellite to the location described by the coordinate x.

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Channel	Gas	Spectral Region. cm <sup>-1</sup>	Ригрозе
A1	CO <sub>2</sub>	600-800	temperature, attitude
A2	CO <sub>2</sub>	1700-2700	non-LTE
A3	CO	1700-2700	composition
A4	NO	1700-2700	composition
B1	Н2О	3650-4250	composition
B2	Н2О	170-330	composition
C1	CO2	600-800	temperature, attitude
C2	N2O	1180-1350	composition
C3	CH4	1180-1350	composition

TABLE 1. SAMS Spectral Channels

spheric levels of interest, the source function in equation (1) is the Planck function B(v, T).

Explicit integrations along the limb path and over the instrument pass band both involve substantial amounts of computation so that if a direct calculation is to be made sufficiently fast, some more rapid and hence necessarily approximate means of evaluating  $\Phi(h)$  must be found.

We first simplify equation (1) by dividing by  $A\Omega \int F(v) dv$ , obtaining

$$L(h) = \iint \tilde{F}(v)B[v, T(x)] \frac{d\tau_h}{dx}(v, x) dv dx$$

where L(h) has the units of spectral radiance and  $\vec{F}(v)$  is the normalized spectral response of the instrument. As the internal calibration is made by viewing a black body at a known temperature of about 290 K, we find it convenient to express the measured radiance as a fraction of that from a black body at T = 290 K. The radiance seen by the instrument from a black body at temperature T is

$$L_{\mathbf{g}}(T) = \int B(v, T) \vec{F}(v) \, dv$$

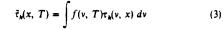
so that the normalized radiance I(h) is

$$I(h) = L(h)/L_{B}(T_{0}) = \iint \frac{L_{B}[T(x)]}{L_{B}[T_{0}]} f(v, T(x)) \frac{d\tau_{b}}{dx}(v, x) dv dx$$

where  $f(v, T) = B(v, T)F(v)/L_{p}(T)$  is a weight which depends mainly on the instrumental spectral response, but also varies slightly with the temperature of the atmospheric path. We can now formally remove the frequency integration by writing

$$l(h) = \int \vec{B}(T(x)) \frac{d\vec{\tau}_h}{dh}(x, T) dx \qquad (2$$

where  $\tilde{B}(T) = L_{g}(T)/L_{g}(T_{0})$  is a known function of temperature and



This spectrally integrated transmittance must be precomputed and parameterized for  $CO_2$ , so that equation (2) can be solved for the temperature profile. The functional form of the transmittance must be parameterized for the trace gases so that equation (2) can be solved for their concentration.

#### 2.1. Integration Along the Path

RODGERS ET AL., SAMS TEMPERATURE AND COMPOSITION RETRIEVALS

The integration along the line of sight is carried out by dividing the atmosphere into a number of concentric shells, with  $z = \ln (p_{0/P})$  as the vertical coordinate (p is pressure and  $p_0$  is one atmosphere) and equal shell thicknesses of  $\Delta z = 0.2$ (about 1.4 km). Each limb path thus falls naturally into a number of segments, each bounded by the shell boundaries (see Figure 1). We now replace the integration along the limb path by a summation over the N segments, with the gradient of path transmittance being evaluated by finite differences so that the emission from the limb path may be written

$$l(h) = \sum_{1}^{N} \tilde{B}(\tilde{T}_{i})(\tilde{\tau}_{i} - \tilde{\tau}_{i-1})$$

where T is the mean (mass weighted) temperature of the *i*th segment.

The next step is to use a version of the Curtis-Godson approximation (C.G.) to replace each inhomogeneous path with a homogeneous path that has an equivalent absorption. The required equivalence is obtained by matching the absorber amounts and the mass weighted mean pressures and temperatures of each pair of paths. Details are not given here, but it may be shown, by considering the spectroscopic domains where the C.G. approximation is exact, the spectroscopy of the bands concerned, and the expected atmospheric distributions of the gases, that errors in the transmittance of paths arising from the use of the C.G. approximation are negligible in this instance. Typical errors are found to be  $\ll 0.2\%$  for methane and  $\ll 0.5\%$  for nitrous oxide at 50% transmittance [Jones, 1983].

Using this approximation, the mean atmospheric path transmittance  $\tilde{\tau}_i$  (equation (3)) may be written as a function of four variables, the absorber mass  $\tilde{u}_i$ , mean pressure  $\tilde{p}_i$ , mean temperature  $\tilde{T}_i$ , and the local temperature  $T_i$ ; that is,

$$\tilde{\tau}_i = \tilde{\tau}(T_i, \, \bar{u}_i, \, \bar{p}_i, \, \bar{T}_i) \tag{4}$$

The mean path temperature and the local temperature are in general not equal, but, in practice, the mean path transmittance depends only weakly on the local temperature  $T_i$ . We have therefore taken  $T_i$  to be equal to  $\tilde{T}_i$  in equation (4), thereby reducing the number of independent variables describing the state of any path to three. Such an assumption only has the effect of weighting the mean path transmittance by the frequency dependence of the Planck function of a slightly wrong temperature. The error due to this effect depends on a number of factors, principally the gradient of the Planck func-

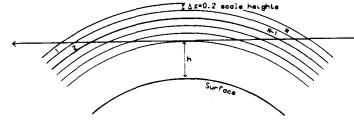


Fig. 1. The geomary of limb sounding.

TABLE 2.	The Effect of	<b>Planck Function</b>	Weighting on Calculated
Pressu	are Modulated	Transmittance in	Two Sample Cases

	Pressure S Transn			
Local temperature. T	CH <sub>4</sub> 21.2 mbar	N <sub>2</sub> O 23.9 mbar	$\frac{B(1150, T)}{B(1350, T)}$	
195 K	0.6103	0.6731	2.57	
220 K	0.6083	0.6730	2.19	
245 K	0.6067	0.6729	1.93	
270 K	0.6045	0.6728	1.74	

The C.G. parameters for the absorbing path are  $u = 10^{18}$  mols/cm, p = 10 mbar. T = 195 K. The modulator length is 3 mm, its temperature is 296 K, and its mean pressures are as in the table.

tion across the passband and the distribution of the spectral lines. For the pressure modulated  $CH_4$  and  $N_2O$  sensing channels in SAMS, the resulting errors in transmittance are shown in Table 2 for an arbitrarily chosen path. Despite the substantial change of the Planck function gradient with local temperature (column 4), the calculated  $N_2O$  channel transmittances (column 3) remain almost constant (due to the high degree of symmetry of the *P* and *R* branches about the band center) while for the CH<sub>4</sub> channel (column 2), the changes in transmittance, although larger, are small (<0.5% in transmittance).

Precomputing a data base of homogeneous transmittances for suitable ranges of the C.G. equivalent u, p, and T is practical, and the costly line-by-line integration across the instrument passband is thus avoided by parameterizing the data base in terms of simple, readily evaluated functions (called the "transmittance functions") of the three independent variables.

#### 2.2. Transmittance Function Parameterization

The path transmittances are found to depend only weakly on the mean path temperature T, allowing their temperature dependence to be treated in a simple manner. A suitable functional form for a transmittance function is as a number of bivariate polynomial surfaces in log (u) and log  $(\bar{p})$  at a number of preselected temperatures  $T_m$ ,  $m = 1, \dots, M$ , each of the form

$$\ln\left(-\ln\left(\tilde{\tau}_{i}\right)\right) = \sum_{ij} A_{ij} \left(\ln \bar{p}\right)^{i} (\ln \bar{u})^{j-1}$$

with linear interpolation between the temperatures. The coefficients  $A_{ij}^{m}$  were obtained by fitting a bivariate Chebyshev series to  $\ln(-\ln(\tau))$  at values of  $\log_{10}(\vec{u})$  and  $\log_{10}(\vec{p})$  in intervals of 0.5 (i.e.,  $\vec{u}$  and  $\vec{p}$  spaced by factors of  $\sqrt{10}$ ), and path temperatures between 170 K and 320 K in intervals of 25 K. The degrees of the polynomials needed depend on the application and intended coverage, a seventh order fit in  $\log(\vec{v})$  and between a third and seventh order fit in  $\log(\vec{p})$  being used for the pressure modulated channels.

The homogeneous path transmittance computations use a line-by-line integration across the instrument passband. The spectral data used in these calculations were based on those available from the spectral data compilations of *McClatchey* et al. [1973] and Rothman [1981], with modifications required to fit a series of laboratory measurements on simulated atmospheric paths in the case of methane and nitrous oxide [Jones, 1983].

#### 2.3. Incorporation of the FOV

The discussion has up to this point been restricted to radiances from a single tangent height. SAMS in fact possesses an approximately Gaussian field of view with a width at the tangent point of about 8 km at half height (measurement of the instrument's fields of view forms an important part of the prelaunch calibration program and is described by Wale and Peskett [this issue]). The effects of the finite field of view may be accounted for by convolving the field of view with the radiance profile. This is done to good precision in the case of constituent retrieval by computing the infinitesimal field of view radiances at each shell boundary (i.e., every 0.2 scale heights or 1.4 km) and performing a weighted average. In the case of temperature retrieval, weighting functions are used in the linearization. The infinitesimal field of view weighting function is calculated and then convolved with the actual field of view.

#### 3. TEMPERATURE RETRIEVAL

The SAMS views at right angles to the direction of satellite travel, thus the variation of the roll angle causes corresponding changes of the observation tangent level. As a consequence of the uncertainty in satellite roll angle, the tangent height corresponding to any measurement is not known a priori with adequate accuracy. Therefore, the temperature retrieval process must estimate roll as well as the temperature profile.

The retrieval is a sequential maximum likelihood estimator. At each observation time (i.e., each point in the scan) several pieces of information are known about the unknown quantities, roll angle, and temperature profile. These are (1) the wide band and PMR radiances at two different altitudes (channels A1 and C1), (2) the look angle relative to the space-craft, (3) the roll angle and rate derived from the spacecraft attitude control system (for NIMBUS 7 these are 2 or 3 orders of magnitude less accurate than needed for specifying the tangent level), (4) the retrieved profile, roll angle, and roll rate at the previous observation time, 2 s earlier, (5) the climatological temperature distribution.

If we make some assumptions about the statistics of the horizontal variation of the temperature profile, and the time evolution of roll, then item 4 can be used to provide an a priori estimate of these quantities at the observation time. All these pieces of information have error estimates, and they are all combined by using a statistically optimum maximum likelihood estimator [Rodgers, 1976] to provide the best estimate

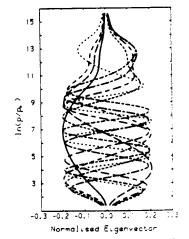


Fig. 2. Eigenfunctions of the solution covariance that have been used as an efficient representation for the temperature prohie

## RODGERS ET AL.: SAMS TEMPERATURE AND COMPOSITION RETRIEVALS

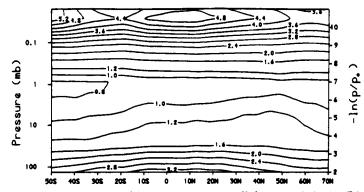


Fig. 3. An estimate of the rms precision of the retrieved temperature (K) for one particular day, February 5, 1981, as a function of latitude and altitude. This estimate is obtained by using equation (7b) and includes both instrumental noise and "unseen structure."

of the unknown profile and roll at the observation time, together with the error statistics of this estimate.

#### 3.1. The Temperature Profile Representation

To minimize computing time, it is important to use the most efficient representation possible for the temperature profile. This has been done here by using a set of empirical orthogonal functions which represent an ensemble of retrieved profiles most efficiently, rather than functions representing an ensemble of actual profiles. If the covariance of an ensemble of actual profiles is S, and the instrument has weighting function matrix K and noise covariance S<sub>e</sub>, then it is straightforward to show that the covariance of the ensemble of maximum likelihood retrieved profiles S<sup>R</sup> is

$$\mathbf{S}^{R} = \mathbf{S}\mathbf{K}^{T}(\mathbf{K}\mathbf{S}\mathbf{K}^{T} + \mathbf{S}_{*})^{-1}\mathbf{K}\mathbf{S}$$

The representation with least error of a member of this ensemble, for a given number N of coefficients, is as a linear combination of those eigenvectors of  $S^R$  which correspond to the N largest eigenvalues.

SAMS does not have a fixed set of weighting functions, nor is the atmospheric covariance S known adequately for our range of heights. Nevertheless, we have used the above representation as a reasonably efficient one for our purpose. For the elements of S we have assumed

$$S_{ii} = S_0 \exp(-(i-j)^2/a^2)$$

where *i* and *j* are level numbers, with levels spaced at intervals of 0.2 scale height,  $a^2$  is 50, and  $S_0$  is 0.0625 in units of the square the radiance of a 290 K black body. This assumes that the atmosphere is correlated over a range of about 3-4 km and ignores longer-range correlation. For the matrix K in this calculation we have used a set of weighting functions corresponding to tangent heights from 2.3 to 8.3 scale heights for the wide band channels and from 4.7 to 13.5 scale heights for the PMR, both in steps of 0.4 scale heights. The resulting set of eigenfunctions are shown in Figure 2.

#### 3.2. Stochastic Equation for Roll

The instrument datum used is the direction of the D field of view when the vertical scan mirror is at step 41, its central position. The fields of view of each channel were determined relative to channel 1 during prelaunch tests, and the mirror angle relative to step 41 is known continuously to high accuracy by telemetry. Roll angle  $\theta_t$  at time t is defined for the purpose of the retrieval as the angle of the datum above the direction which has a tangent point at pressure level p where  $\ln (p_0/p) = 7$  (about 0.92 mbar). A roll rate r is also defined such that  $r_t = \theta_t - \theta_{t-1}$ . Time t is expressed in units of observation intervals of 2 s. Given estimates  $\theta_{t-1}$  and  $\hat{r}_{t-1}$  at time t - 1, we assume an a priori estimate at time t given by

$$\theta_t^0 = \hat{\theta}_{t-1} + \hat{r}_{t-1} + \delta_r$$
$$r_t^0 = \beta \hat{r}_{t-1} + \delta_r$$

where  $\delta$ , is a gaussian random change of roll rate and  $\beta$  is a constant between 0 and 1 which represents the long-term control of mean roll rate to be zero. This stochastic equation is equivalent to assuming that the torque on the spacecraft is random and uncorrelated on a time scale of 2 s. This is a conservative assumption, but it appears to be adequate for our purposes.

#### 3.3. Stochastic Equation for Temperature

Temperature is represented by a vector x of coefficients of eight eigenfunctions of the deviation of profile of Planck function at 675 cm<sup>-1</sup> from the climatology interpolated to the appropriate latitude and day. This wave number was chosen because the Planck function at that frequency most closely matched the spectrally weighted black body radiances  $L_g(T)$  for the two PMR CO<sub>2</sub> channels.

We assume that the coefficients obey a first-order stochastic process thus:

$$\mathbf{x}_r = \alpha \mathbf{x}_{r-1} + \mathbf{\delta}_r \tag{5}$$

where  $\delta_x$  is a gaussian random vector and  $\alpha$  represents the fact that in the long term  $\bar{x} \rightarrow 0$ .

A relationship between the coefficient  $\alpha$  and  $S_{\sigma}$ , the covariance of  $\delta_{\pi}$  can be found by taking the covariances of both sides of (5) to be equal to  $S^{R}$ :

$$\mathbf{S}^{\mathbf{R}} = E\{\mathbf{x}_{t}\mathbf{x}_{t}^{T}\} = E\{(\alpha\mathbf{x}_{t-1} + \mathbf{\delta}_{x})(\alpha\mathbf{x}_{t-1} + \mathbf{\delta}_{x})^{T}\}$$

therefore

hence

$$S^{R} = \alpha^{2}S^{R} + S$$

$$\mathbf{S}_{\mathbf{A}} = (1 - \alpha^2) \mathbf{S}^{\mathbf{R}}$$

5284

A second (or higher) order stochastic process would clearly be more representative, but leads to excessively large computing requirements.

# 3.4. Construction of the A Priori Information, u,<sup>0</sup>

We define the state vector **u** as  $(\mathbf{x}, \theta, r)$ . The stochastic equations for roll and temperature can be written jointly as  $\mathbf{u}_r^0 = \mathbf{A}\hat{\mathbf{u}}_{r-1} + \mathbf{\Delta}$ 

where

and

$$\mathbf{A} = \begin{pmatrix} \mathbf{1} & \mathbf{0} \\ & \mathbf{1} & & \mathbf{1} \\ \mathbf{0} & & & \\ & & \mathbf{0} & & \boldsymbol{\beta} \end{pmatrix}$$

$$\Delta = (\delta_r, \delta_r, \delta_s)$$

If the covariance matrix of  $\hat{u}_{t-1}$  is  $\hat{S}_{t-1}$  then the covariance matrix of the a priori  $u_t^o$  is given by

$$\mathbf{S}_{t}^{0} = \mathbf{A}\mathbf{\hat{S}}_{t-1}\mathbf{A}^{T} + \Delta\Delta^{T} = \mathbf{A}\mathbf{\hat{S}}_{t-1}\mathbf{A}^{T} + \mathbf{S}_{\Delta}$$

# 3.5. Linearization of the Direct Model

When the a priori profile is close to the retrieved profile, as is the case with this sequential estimator, linearization is likely to be an accurate approximation to the direct model. This allows the usual linear theory to be used without iteration. The linearization for temperature retrieval takes place in two stages. Equation (2) is sufficiently linear in  $\tilde{B}$  that it can be adequately linearized once and for all with respect to Planck function about a climatological (time and space dependent) distribution. However the nonlinearity due to roll is so large that this linearization must be done about the a priori at each observation time.

Radiances are calculated as functions of pressure for climatological mean profiles taken every  $10^{\circ}$  latitude and every month and are interpolated to the exact latitude and day. Differential weighting functions W(h) relating radiance from tangent height h to the coefficient of the representation are calculated numerically by finding the changes of radiance resulting from separate small changes in each of the coefficients and averaging over all climatological profiles.

$$W_i(h) = \{I(h, x_i + \Delta x_i) - I(h, x_i)\}/\Delta x_i$$

The radiance profile is therefore related to tangent height h and eigenvector coefficients by

$$I(h) = I_c(h) + \mathbf{W}^{T}(h)\mathbf{x}$$
 (6)

where  $I_c(h)$  is the climatological radiance profile. If we linearize (6) with respect to h about a nominal value  $h_0$  we obtain

$$I(h) = I_c(h_0) + \mathbf{W}^T(h_0)\mathbf{x} + \left\{\frac{\partial I_c}{\partial h}(h_0) + \frac{\partial \mathbf{W}^T}{\partial h}(h_0)\mathbf{x}\right\}(h - h_0) + O(h - h_0)^2$$

This can be put in the form

$$\Delta I = \mathbf{M}^{\mathsf{T}} \Delta \mathbf{u}$$

where  $\Delta I = I - I_c(h_0) - \mathbf{W}^T(h_0)\mathbf{x}$  is the deviation of the observation from that expected using  $\mathbf{x}_0$  and  $\theta_0$ ,  $\Delta \mathbf{u}$  is the vector  $(\mathbf{x} - \mathbf{x}_0, \theta - \theta_0, r - r_0)$  and **M** is a vector of coefficients, the last of which is zero as radiance does not depend on roll rate.

Given an a priori estimate  $\mathbf{u}^{\circ}$  with covariance  $\mathbf{S}^{\circ}$ , we can update this to find an improved estimate  $\hat{\mathbf{u}}$  with covariance  $\hat{\mathbf{S}}$ by using the observed  $\Delta I$  and its variance  $\sigma^2$  with equations of *Rodgers* [1976]

$$\hat{\mathbf{u}} = \mathbf{u}^{0} + \mathbf{S}^{0}\mathbf{M}(\mathbf{M}^{T} - \mathbf{S}^{0}\mathbf{M} + \sigma^{2})^{-1}\Delta I$$
(7a)

$$\hat{\mathbf{S}} = \mathbf{S}^{0} - \mathbf{S}^{0}\mathbf{M}(\mathbf{M}^{T} - \mathbf{S}^{0}\mathbf{M} + \sigma^{2})^{-1}\mathbf{M}^{T}\mathbf{S}^{0}$$
(7b)

For this application we have up to four observations at each time (two WB/PM pairs), but because the number varies (mainly because one or more of the channels may be viewing outside its useful height range) we choose to apply equations (7a) and (7b) separately for each channel as a scalar, using the new  $\hat{\mathbf{u}}$  and  $\hat{\mathbf{S}}$  from one update as the  $\mathbf{u}^0$  and  $\mathbf{S}^0$  of the next, including using the revised roll angle to provide an improved tangent height linearization point.

#### 3.6. Practical details

We have initially available for each observation time an approximate estimate of the roll angle determined by a sequential estimation process which uses the ratio of wide band to PM radiances and the output of the satellite attitude control system. This method gives roll angle sufficiently accurately to be a starting point from which the full retrieval can converge. It is only necessary at the start of sets of atmospheric data (e.g., after data gaps).

Data are normally processed in sets of about 24 hours. The retrieval algorithm is applied to a whole 24 hour period which is available in sequential order. The estimator is started at a point 10 min from the start of data and run back to the start. the results being discarded except for the estimates of x and  $\theta$ for the first observation. The estimator is then run forward through the data set,  $\theta$  and its variance being stored every observation time (every 2 s), and x and  $\theta$  and their covariances being stored at every crossing of a 2.5° latitude line. At the end of the data set the final x and  $\theta$  are used as starting values and the estimator run backward along the orbit to the start. This time the backward estimate of  $\theta$  is combined with the forward estimate previously stored, weighted by the reciprocals of their variances. Similarly at 2.5° latitude intervals combined values of x are also produced in accordance with the covariances. At the start of both the forward and backward runs S<sup>0</sup> is set to a large value representing no data. When forward and backward are combined no attempt is made to combine roll rate (since it is not required) or to use its covariance since, for a given time for which x and  $\theta$  are applicable, the forward and backward values of r correspond to times separated by one time interval.

The approximate first estimate roll angle is incorporated into the full estimate as a pseudo observation (with relatively large standard deviation of  $0.4^{\circ}$ ) every time slot. It is only effective during long periods when there is no atmospheric data sufficient to constrain the solution.

#### 3.7. Sequential Estimation of Space-Radiance Offset

The precision of calibration of radiances near zero limits the lowest values of atmospheric emissivity, and thus the highest altitude of measurement, that can be used. For example, an emissivity of 0.1 when viewing an atmosphere at 240 K would give a radiance 0.05 of that of the 290 K internal black body. Since with the SAMS implementation of pressure modulation the channels chop against references at instrument temperatures (about 290 K), changes of gain and instrument temperature cause apparent changes of radiance when viewing space. Instrument scan patterns are chosen to give views of space

# APPENDIX D

5285

well above the atmosphere at about 10 min intervals. In addition, every scan cycle (32 s) the channels view levels sufficiently high in the atmosphere for the radiance to be regarded as zero. Only the retrieval algorithm has knowledge of whether these are high enough to be usable as a zero calibration. Therefore, an estimator of space view radiance offset is run in parallel with the attitude and temperature estimator, using the roll angle estimate to provide estimates of radiance correction for each channel.

#### 4. COMPOSITION RETRIEVAL

The retrieval of the distribution of trace gases could in principle be carried out in the same way as the temperature retrieval, by linearizing the direct model at every observation time and by updating the profile estimate and its covariance with the measured radiance. However, because the direct model is much more complicated in this case, the computing requirements of this method are prohibitive. Furthermore, the signal to noise in the composition channels is relatively low, so that this approach would be rather like taking a sledgehammer to crack a nut. We have therefore chosen to analyze the radiances in a relatively coarse grid before retrieving profiles, thus considerably reducing the number of profiles retrieved. The details of the grid vary with the gas, depending on the signal to noise ratio.

#### 4.1. Use of Sequential Estimation

The starting point for a composition retrieval is a radiance profile sampled at a fixed interval in  $-\ln(p)$ , usually 0.2 scale heights. The profile to be retrieved is modeled by straight line segments in the logarithm of mixing ratio against  $-\ln(p)$ , using a relatively small number, *n*, of points to determine the profile. Thus the problem is over constrained, and does not, in principle, need an a priori estimate. In practice, of course, a rather good first guess is needed as a linearization point, and to prevent the solution behaving badly at those altitudes where there is little information in the measurements.

The profile is retrieved, starting at the top, by sequentially updating the coefficients x representing the profile and its co-variance S with each measured radiance  $I_i$ (and its variance  $\sigma_i$ ) in turn, thus

$$\mathbf{x}_{i} = \mathbf{x}_{i-1} + \mathbf{S}_{i-1} \mathbf{W}_{i} \Delta I_{i}^{m} / (\mathbf{W}^{T} \mathbf{S}_{i-1} \mathbf{W}_{i} + \sigma_{i}^{2})$$
  
$$\mathbf{S}_{i} = \mathbf{S}_{i-} - \mathbf{S}_{i-1} \mathbf{W}_{i} \mathbf{W}_{i}^{T} \mathbf{S}_{i-1} / (\mathbf{W}_{i}^{T} \mathbf{S}_{i-1} \mathbf{W}_{i} + \sigma_{i}^{2})$$

where subscript *i* refers to the radiance level being used for the update, i - 1 is the result of the previous stage,  $W_i$  is a vector of weighting functions, and  $\Delta I_i^{m}$  is the departure of the measured radiance from that calculated for level *i* using the retrieval of stage i - 1:

$$\Delta I_i^{\mathbf{m}} = I_i^{\mathbf{m}} - I_i^{i-1} (\text{calc})$$

The weighting functions  $\mathbf{W}_i$  are calculated by finite differences at each stage

$$\mathbf{W}_i = \frac{\Delta \mathbf{I}_i}{\Delta \mathbf{x}}$$

 $W_i$ , depends on the temperature structure, the a priori concentration profile, and the pressure of gas in the PMC.

The final solution of this iteration  $\hat{\mathbf{x}}$  and  $\hat{\mathbf{S}}$  are equal to  $\hat{\mathbf{x}}_n$ and  $\hat{\mathbf{S}}_n$ , respectively. This solution has second-order convergence and is usually found to converge adequately in one step. If further iterations are required, the covariance  $\mathbf{S}^0$  is reinitia lized to the first guess, the a priori absorber distribution  $x^0$  is set equal to  $\hat{x}$ , and the weighting function matrix is reevaluated by linearizing about the new initial guess  $x^0$ .

# 4.2. Practical Details

The CH<sub>4</sub> and N<sub>2</sub>O sensing channels on SAMS share the same detector and thus cannot both be enabled at any one time. To maximize the coverage obtained, the CH<sub>4</sub> and N<sub>2</sub>O PM channels are thus generally arranged to operate on alternate 24 hour periods. It is thus convenient, and in fact appropriate from signal to noise considerations, to treat each 24 hour period independently. Although the longitudinal variations of CH<sub>4</sub> and N<sub>2</sub>O are being investigated, the quantities processed for archiving are zonal mean mixing ratio cross sections. After calibration and removal of space view offset, the radiances are zonally averaged into 10° latitude bands and in 0.2 scale height intervals using the retrieved attitude and the tangent point latitude. Radiances with estimated attitude errors greater than 0.15 scale heights are rejected at this point. At the same time the corresponding zonal mean temperature profiles are obtained by averaging the Planck function at 1290 cm<sup>-1</sup> (the center of the CH<sub>4</sub> and N<sub>2</sub>O channels) and retrieved temperatures over the same bands. The zonal mean radiances and temperatures are then archived as an intermediate product.

To retrieve a given data set, latitude dependent a priori estimate of the CH<sub>4</sub> and N<sub>2</sub>O zonal mean mixing ratios are determined using the Oxford two-dimensional model [Harwood and Pyle, 1975] assuming large covariances (implying uncertainties of 25 and 100% in the CH<sub>4</sub> and N<sub>2</sub>O mixing ratios, respectively). Starting a month into the data set, and working backward in time, each day's data is retrieved by using a single iteration. At the completion of each day, the a priori estimate for the species being measured that day (either CH<sub>4</sub> or N<sub>2</sub>O) is then replaced with the retrieved zonal mean mixing ratio cross section, and the covariance is reinitialized. This process is repeated until the start of the data set is reached, whereupon the retrieval is run forward through the data set in the same manner with the zonal mean cross sections and their covariances now being stored.

#### 5. ERROR ANALYSIS

Random errors arise mainly from the contribution of detector noise to the measured radiances. One reason for using the optimal estimation equations is to treat this component of the total error correctly and automatically, obtaining the covariance matrix of the random error of the solution (equation (7b)) as one of the outputs of the system. Thus every retrieved profile has its own error estimate attached. An example of this error estimate is given in Figure 3, which shows contours of temperature error in a zonal cross section. Note that this error includes contributions from both instrument noise and "unseen structure" (i.e., fine structure in the temperature profile (as expressed in its covariance matrix) which is invisible at the vertical resolution of the instrument).

Systematic errors arise from many sources, some of which have already been touched on:

1. Calibration error: gain, space offset, temperature dependences, etc.

2. Instrumental: PMC pressure, field of view profile. spectral filter profile.

3. Location error: tangent height (for composition), lati-

# APPENDIX D

\$286

tude, time. Winds and spacecraft yaw stability (via Doppler shift).

4. Spectroscopy: line strengths, widths, temperature dependences, air broadening, and self broadening. C.G. approximation. Transmittance parameterization. Overlap with other gases.

5. Retrieval: zonal averaging (for composition). Stochastic prediction equation. Use of eigenvectors as a representation. Nonlinearity. Temperature errors in the case of constituent retrieval.

Error analysis may in principle be carried out by perturbing the retrieval equations in a way appropriate to the error source, to determine the effect on the resulting profile. The details are straightforward but tedious. A discussion of these errors in the case of temperature is given by *Barnett and Corney* [this issue] and in the case of methane and nitrous oxide by *Jones and Pyle* [this issue].

# 6. DISCUSSION

Retrieval of temperature and composition profiles from SAMS presents both theoretical and logistical problems. The inverse problem for temperature is nonlinear both because of the temperature dependence of the transmittance and because the scan is geometrical, while the atmospheric mass distribution with height varies according to the hydrostatic equation. The inverse problem for composition is more nonlinear because the absorber distribution to be measured appears as an argument of the transmittance function. Logistical problems arise because the scan pattern of SAMS is programable and therefore arbitrary and because the attitude control of the spacecraft is inadequate. Furthermore, owing to the satellite's horizontal motion during the vertical scan, the radiance profile is along a slant path, not a vertical path, so that vertical profiles are not measured directly. The problems are further compounded by the low signal/noise ratio in some channels.

In the case of temperature, the retrieval method presented here retrieves all the unknowns, including roll angle and rate, at every observation time, using a single radiance and the optimal estimation equations. This may appear unlikely but is no different in principle from retrieving a continuous vertical profile from perhaps 10 simultaneous nadir sounding channels. In practice, up to four channels are used at any one observation time, two PM and two WB  $CO_2$  15  $\mu$ m radiances. Continuity along the tangent track is provided by way of a stochastic equation which provides the a priori profile using that retrieved at the previous observation time. Forward and backward estimates are combined so that continuity is used in both directions.

Apart from the question of statistical optimality, the method has a number of practical advantages, which follow:

1. The a priori is always close to the solution, so that problems of nonlinearity and iteration are minimized.

2. Horizontal continuity is incorporated via the stochastic equation.

3. The scan pattern need not be pre-specified.

4. The first stage of global gridding is done automatically.

5. Objective interpolation is automatically carried out over data gaps.

We have not included in this paper examples of retrieved cross sections of temperature and composition and their accuracy and precision. For these we refer the reader to the two companion papers, *Barnett and Corney* [this issue] and *Jones* and *Pyle* [this issue]. Barnett and Corney discuss the validation of the temperature profiles in comparison with Tiros-N SSU and rocket measurements, while Jones and Pyle present methane and nitrous oxide cross sections in comparison with those computed by the Oxford two-dimensional model.

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# TEMPERATURE COMPARISONS BETWEEN THE NIMBUS-7 SAMS, ROCKET/RADIOSONDES, AND THE NOAA-6 SSU

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Temperature Comparisons Between the NIMBUS 7 SAMS, Rocket/Radiosondes and the NOAA 6 SSU

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The stratospheric and mesospheric sounder (SAMS) on the NIMBUS 7 satellite is a limb sounder measuring infrared thermal emission from gases in the earth's atmosphere. Carbon dioxide measurements near  $15 \,\mu$ m are used to determine the temperature profile and satellite roll angle. These temperature fields are compared with measurements by the NOAA 6 SSU and radio/rocketsondes over a 13-month period in 1980/81. Consistent patterns of bias between the instruments are found but are generally less than 2 K. Some of these are attributed to atmospheric tides, but others appear to indicate errors, including a problem caused by atmospheric ozone affecting lower-stratosphere SAMS measurements. Remarkably good agreement is found between variations measured by the three sensors, with standard deviations of difference typically of 1 K and correlation coefficients of up to 0.994 over the whole period.

#### 1. INTRODUCTION

The stratospheric and mesospheric sounder (SAMS) is a limb-scanning infrared radiometer carried on board the NIMBUS 7 satellite, which was launched in October 1978. The SAMS worked well until October 1982, when the scan mechanism began to malfunction, thus giving 4 years of continuous data; several months of good data were also obtained in 1983.

This study was undertaken to see how well the SAMS temperature measurements agreed with those of rockets, radiosondes, and the TIROS-N series stratospheric sounder units. The temperature field is an important product of the SAMS measurements for the study of atmospheric thermal structure, radiation, and dynamics, but it is also an essential input for the retrieval of the SAMS minor constituent channels.

#### 1.1 The SAMS

The SAMS hardware was described in detail by Drummond et al. [1980] and by Wale and Peskett [this issue], who also discuss the SAMS calibration and its operation in space.

The SAMS radiometer uses the technique of pressure modulation to select emission from a specified part of the band for a given gas. The principals are outlined by Taylor et al. [1972] and Curtis et al. [1974], and its advantages for limb composition sounding are described by Chaloner et al. [1978] and Drummond and Jarnot [1978]. The pressure modulator selectively modulates the emission from a gas in the atmosphere by using the absorption lines of the same gas as an optical filter. A cell containing the gas to be measured is included in the optical path of the radiometer. The gas pressure is varied so that the transmission of the cell is modulated at frequencies in the range 20-40 Hz. Pressure modulation is sensitive to emission at higher altitudes. By using a conventional black chopper in series with the pressure modulator at a much higher frequency (approximately 240 Hz), it is possible to obtain information from as low in the atmosphere as does a simple wideband radiometer. By combining the two methods of modulation, known as pressure modulation (PM) and wideband (WB) modulation, data have been obtained over altitudes from 10 to 100 km.

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Paper number 3D1991 0148-0227/84/003D-1991\$05.00 Pressure level and temperature sounding is accomplished by two such optical paths in the SAMS: C1, which uses  $\sim 30$ mbar of CO<sub>2</sub> in a 10-mm long cell; A1, which uses  $\sim 12$  mbar of CO<sub>2</sub> in a 3-mm long cell. Both pairs measure emission from the whole 15-µm  $v_2$  band with similar filters, hence the wideband measurements are nearly identical. However the A1 pair are offset to view 0.56° (approximately 30 km) above the C1 pair. The A1 cell pressure is sometimes set to much lower values than above; in this case the signal-to-noise ratio of the pressure-modulated measurement is too poor for it to be useful other than in the zonal mean, and only the C1 measurements are used.

The weighting functions (i.e., the Frèchet derivatives of the measured radiance with respect to the vertical radiance coordinate) for these channels are shown in Figure 1. The retrieval method is described in detail by Rodgers et al. [this issue]. It uses a sequential estimator in which the retrieved profile at each observation time is used as the a priori profile for the retrieval at the next observation time, 2 s later (an observation is a single tangent height, not a complete scan-the scan takes about 32 s). This allows continuity in the horizontal to be included and makes the a priori profile so close to the solution at each stage that linearization is adequate and iteration is not required. The profiles are stored at every 2.5° of latitude along the tangent track, thus automatically doing the first stage of gridding, the "orbit grid." At a later stage, values are interpolated in longitude onto a regular grid with 10° longitude and 2.5° latitude intervals. The northbound and southbound orbits are gridded separately but are then combined to give an average field for the day derived from all values between 0000 and 2400 Z.

#### 1.2 The SSU

5294

The stratospheric sounder unit (SSU) is a temperature sounder provided by the British Meteorological Office and carried on the NOAA operational TIROS-N series satellites. It measures thermal emission from carbon dioxide by using three pressure modulator channels. Although scanning sideways by  $\pm 35^{\circ}$ , it is essentially a nadir sounder, providing measurements of a fundamentally different type from the SAMS. The weighting functions peak near 28 km, 36 km, and 43 km for channels 25, 26, and 27, respectively, and are about 15-km broad (see Figure 1). The data were processed by the Meteorological Office and were supplied in the form of analyses on a

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# APPENDIX E

#### BARNETT AND CORNEY: SAMS TEMPERATURES VS. SONDES AND SSU

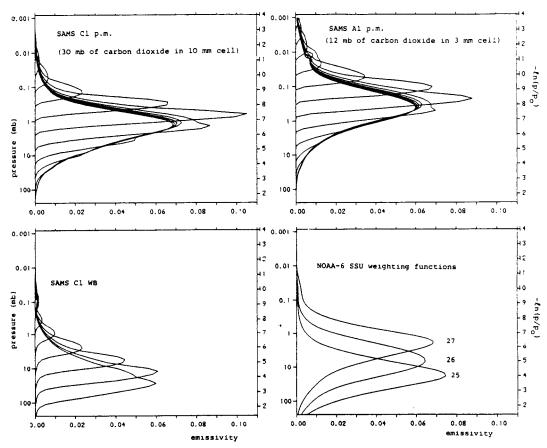


Fig. 1. Weighting functions for SAMS C1 PM, C1 WB, and A1 PM channels and for the NOAA 6 SSU. SAMS weighting functions are given for tangent pressures of the center of the field of view at intervals of 1.0 in  $\ln(p/p_0)$ .

5° by 5° latitude/longitude grid interpolated in time to 1200 Z for each day. Two types of data were used: (1) radiances for each SSU channel (corrected to exact nadir view); (2) thickness analyses from 100 mbar to various levels, derived by the Meteorological Office by regression, and incorporating, to some extent, measurements by the HIRS-II infrared and MSU microwave vertical sounders on the same satellite. For this study, data from the NOAA 6 satellite were used. This has equatorial crossing times of approximately 0730 and 1930 local mean time (see Table 1). Of the SSUs launched so far, this is probably the best understood and the most stable. The quality of the measurements were discussed in detail by Nash and Brownscombe [1983]. They show the overall radiometric precision for channels 25, 26, and 27 to be 0.15, 0.15, and 0.2 mw  $m^{-2}$  ster<sup>-1</sup> (cm<sup>-1</sup>)<sup>-1</sup>, respectively, with an additional component of 0.06, 0.1, and 0.3, respectively, arising from the uncertainty of cell pressure, implying a weighting function mean height uncertainty. At this wavelength these units are approximately equivalent to K for atmospheric temperatures. Pick and Brownscombe [1981], Brownscombe and Schlapp [1983], and Miller et al. [1980] give additional information about the SSU's and their data.

#### 1.3 Rocket/Radiosonde Measurements

These were prepared by the Meteorological Office as part of their SSU comparison program and supplied on magnetic tape. The rocket sensors were Super Loki datasondes, which measure temperature to a precision of 1.0-2.5 K over the 20-60 km height range. The rocket data were originally taken from teletype Rocobs messages and corrected where necessary for dynamic and radiative heating. They were combined with nearly coincident radiosonde profiles and interpolated to give temperature, geopotential height, and wind as functions of pressure. The radiosonde provided both temperature data below 50-mbar pressure and a pressure and height near that level to enable the rocket temperature measurements, which are measured as a function of height, to be obtained on a pressure scale that is the fundamental vertical coordinate for the satellite measurements. Rocket temperature was used exclusively above 10 mbar and radiosonde below 50 mbar, a smooth transition being made in any overlap region between these pressures. In this region the standard levels are 50, 30, 20, and 10 mbar. Where there was no overlap such that one or more levels was missing between the two parts of the profile, no further attempt was made to use the profile. For this study, data from three rocketsonde stations were used. In every case the radiosonde ascent was from the same station. At Ascension Island (8°S, 14°W) the 1200-Z radiosonde was always used, and rocket launches were typically within 2 hours and always within 4 hours. At Kwajalein Island (9°N, 168°E) the 0000-Z radiosonde was always used with rocket measurements typically within 2 hours but up to 5 hours away. At Primrose

5295

5296

BARNETT AND CORNEY: SAMS TEMPERATURES VS. SONDES AND SSU

	SAMS			S		
Latitude	Ascending	Descending	Bias, %	Ascending	Descending	Bias, %
50°S	1657	1842	97	2015	0644	-7
40°S	1517	2021	88	2002	0658	~ 5
20°S	1414	2125	59	1944	0716	-2
0°	1349	2150	50	1930	0730	0
20°S	1342	2157	47	1917	0743	2
40°N	1353	2146	51	1859	0801	5
50°N	1412	2126	58	1842	0817	8
60°N	1451	2048	71	1825	0834	11
67°N	1604	1934	90	1758	0902	15

TABLE 1. Approximate Local Times of SAMS and SSU Observations

SAMS times are for the tangent point location, whereas for the SSU they relate to the NOAA 6 subsatellite point (times in hours and minutes). The bias is explained in section 3.

Lake (55°N, 110°W) the rocket launches were all at about 1800 Z, and either the 0000-Z or 1200-Z radiosonde was used, whichever was more appropriate; the biggest time difference was 8 hours.

## 2. RADIANCE AND THICKNESS COMPARISONS

#### 2.1 Comparisons at Rocket Stations

For this comparison we have studied in detail three rocket stations for which there were particularly frequent launches. Satellite data were interpolated to these locations from the gridpoint analyses. A 406-day period was used for the comparison to give a complete annual cycle plus some overlap. SSU data were available for most days, but the SAMS is routinely switched off for one complete day every fourth, giving 300 days of data. For rocket/radiosonde data the number of profiles available at Primrose Lake and Kwajalein and Ascension Islands were 103, 75, and 83, respectively. Thicknesses of the 100-20 mbar, 20-5 mbar, 5-1 mbar layers were calculated for all sensors at the three rocket stations. In Figure 2 we show a time sequence plot for the three layers at Primrose Lake. Agreement between all sensors appears best in the layer 20-5 mbar, with SAMS being slightly warmer than

SSU at 5-1 mbar but cooler at 100-20 mbar. With almost continuous data the SAMS and SSU values appear to track each other extremely well, both traces exhibiting the same short-term variations. It is more difficult to follow these variations with the rocket data, but when they are sufficiently frequent, e.g., during days 1-80 and 300-350, 1980, there is good correlation.

Figure 3(a-c) and Table 2 summarize the comparisons. Here the three data types have been compared for the whole period in pairs, using days when both types of data were present. Using days when all three types of data were available would in some ways have been preferable but would have substantially reduced the number of comparisons. Thickness differences and standard deviations are plotted as equivalent mean temperatures by using lines that span the entire layer. Temperature profiles are compared at the standard rocket/radiosonde levels, but no SSU profiles were available. SSU radiances were compared with simulated radiances calculated from SAMS and rocket/radiosonde profiles. Weighting functions appropriate to the NOAA 6 SSU were used; they were not varied in time, since this radiometer is sufficiently stable for this to be unnecessary. They were appropriate to a global annual mean temperature profile, and no account was taken

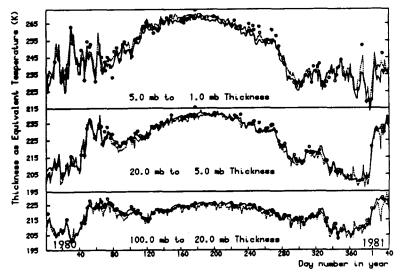


Fig. 2. SSU (solid line), SAMS (dashed line), and rocket (open circles) thicknesses expressed as equivalent layer mean temperature over Primrose Lake rocket station (55°N, 110°W).

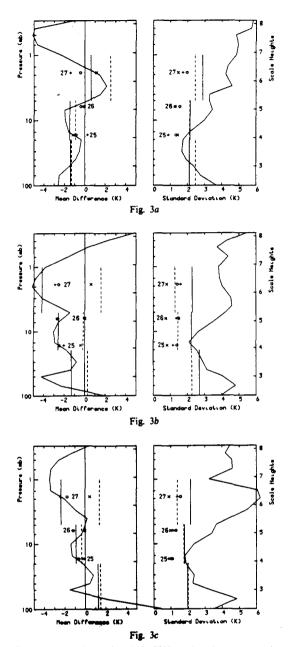


Fig. 3. Comparison of SAMS, SSU, and rockets expressed as equivalent temperatures. The profile is SAMS-rocket given at the standard levels of 100, 70, 50, 40, 30, 20, 15, 10, 7, 5, 4, 3, 2, 1.5, 1, 0.7, 0.5, 0.4, 0.2, 0.15, and 0.1 mbar (the ticks on the pressure scale); the bars are mean layer temperature differences (100-20 mbar, 20-5 mbar, 5-1 mbar) of SAMS-rocket (solid line) and SAMS-SSU (dashed line); the spot points are at SSU weighting function maxima SAMS-rocket (open circles), SAMS-SSU (×), and SSU-rocket (+) for (a) Primrose Lake, (b) Ascension Island, and (c) Kwajalein Island.

of temperature dependence, an approximation that may lead to errors of up to about 1 K. Radiances were converted to brightness temperature each day before differencing and averaging. The differences are plotted at the levels of weighting function maxima.

The mean differences for radiances, temperatures, and thick-

ness are consistent with each other. Standard deviations are also consistent: the smallest variations occur with radiances, which average over the broadest layers. The largest variations occur with temperature profiles, where the SAMS is unable to resolve the fine structure, which itself may change rapidly within the time scale of a day and which would be reduced by horizontal averaging.

Table 2 shows that at Primrose Lake the standard deviation of temperature over the period was 7 to 14 K, depending on the level and the layer. Correlation coefficients are all larger than 0.97, except for the SAMS components of the 100-20 mbar thickness. Correlations are generally higher for radiance than for thickness, being 0.988-0.993 for SAMS vs. SSU and 0.985-0.994 for SAMS vs. rocket (becoming worse with increasing height). The gradients of least squares fit straight lines are given, in each case using first one then the other variable of a pair as the independent (assumed noise-free) variable; this was done because the relative errors of each were unknown. With perfect correlation the two gradients would be equal, while for zero correlation the first would be zero and the second infinite. Ideally, they would both be 1, indicating that both sensors of a pair show the same magnitude of response to a change. At Kwajalein and Ascension Islands the standard deviation of temperature is 1-2 K over the period. This is little more than the standard deviation between the instruments, so correlation coefficients are much poorer, being as low as 0.39. For the same reason the gradients at these places are of little value and are given only for completeness.

The mean differences given in Table 2 show interesting patterns, which will be discussed in section 3. The three differences for a given channel or layer at a single station do not add to zero because different sets of days were used for each comparison, as already discussed. In the middle and upper stratosphere, standard deviations of difference are generally smallest for SAMS-SSU and largest for SAMS-rocket, indicating that the SSU shows the smallest and the rocket the largest variation about some compromise "truth." This interpretation must be treated with caution, since errors may be correlated between sensors, e.g., this could result from error-free rocket measurements but with SAMS and SSU both having offset errors that vary seasonally together. Estimates of the absolute standard deviation of each sensor can be obtained from the variances and covariances between differences (preferably using only days when all three measurements are present), but too few days were available for this to be done reliably.

Figure 4 shows the scatter between channel 27 SSU and SAMS at Primrose Lake; days with and without A1 PM available are shown separately, but there is no suggestion that these two classes of measurement fall on separate lines. The line of points tends to be slightly curved, falling closer to the SAMS-SSU line in the middle 230-250 K range than at higher or lower temperatures. From Table 2a the mean and standard deviation of SAMS-SSU for this set of measurements is 1.14 K and 1.46 K, respectively.

# 2.2 SAMS and SSU Comparisons as a Function of Latitude

Zonal mean SAMS-SSU brightness temperatures have been calculated for each day. They show consistent seasonal variations, changing smoothly from month to month. Figure 5 gives monthly means at 3-month intervals as being representative of the behavior. Differences of up to 3 K occur, usually peaking toward either pole and being smallest in the

# APPENDIX E

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5298

# BARNETT AND CORNEY: SAMS TEMPERATURES VS. SONDES AND SSU

	Overall Standard				Standard		dy	/dx
Channel Layer	Deviation, K	Difference, x y	Number of Pairs	Mean, K	Deviation, K	Correlation Coefficient	x independent	y independen
	~	~ ,					macpendent	mdependen
Channel 25	8.0	SAMS-SSU	262	imrose Lake – 1.16	1.42	0.988	1.002	1.027
Channel 25	6.0	SAMS-SSO SAMS-Rocket	74	-0.86	1.33	0.988	1.002	1.027 1.031
		SSU-Rocket	94	0.26	0.88	0.994	1.005	1.031
Channei 26	9.5	SAMS-SSU	262	-0.26	1.27	0.993	1.003	1.017
	5.5	SAMS-Rocket	74	-0.42	1.54	0.988	1.002	1.026
		SSU-Rocket	94	-0.07	1.27	0.991	1.005	1.024
Channel 27	9.8	SAMS-SSU	262	1.14	1.46	0.990	0.988	1.008
		SAMS-Rocket	74	0.44	1.97	0.982	1.007	1.043
		SSU-Rocket	94	- 1.42	1.73	0.985	1.010	1.042
100–20 mbar	7.0	SAMS-SSU	261	-1.32	2.47	0.936	1.027	1.172
		SAMS-Rocket	74	-1.43	2.13	0.941	1.018	1.149
		SSU-Rocket	94	-0.22	1.44	0.971	0.963	1.021
20–5 mbar	12.2	SAMS-SSU	261	-0.92	2.13	0.985	0.991	1.022
		SAMS-Rocket	74	-1.49	2.18	0.980	0.978	1.019
e i subse	14.0	SSU-Rocket	94	-0.60	1.55	0.989	1.011	1.034
5-1 mbar	14.2	SAMS-SSU	261	2.53	2.45	0.985	1.021	1.052
		SAMS-Rocket	74 94	0.61 1.73	2.91 2.97	0.979 0.976	1.030 1.001	1.075
		SSU-Rocket				0.970	1.001	1.050
			Asc	ension Islan				
Channel 25	1.4	SAMS-SSU	268	-0.40	0.81	0.939	0.80	1.14
		SAMS-Rocket	66	., -2.49	1.39	0.677	0.79	1.73
		SSU-Rocket	76	- 2.02	1.12	0.810	1.13	1.72
Channel 26	1.8	SAMS-SSU	268	0.01	0.73	0.894	0.96	1.21
		SAMS-Rocket	66	-2.71	1.48	0.706	0.95	1.90
		SSU-Rocket	76	- 2.61	1.34	0.785	0.99	1.62
Channel 27	1.9	SAMS-SSU	268	0.60	0.79	0.878	1.00	1.30
		SAMS-Rocket	66	- 2.58	1.40	0.698	0.86	1.77
		SSU-Rocket	76	-2.94	1.60	0.680	0.73	1.58
100-20 mbar	1.8	SAMS-SSU	270	0.29	2.26	0.403	0.32	1.99
		SAMS-Rocket	65	-1.30	2.72	0.389	0.31	2.06
		SSU-Rocket	73	-1.77	1.56	0.695	0.86	1.79
20–5 mbar	2.1	SAMS-SSU	270	-0.15	1.45	0.788	0.71	1.14
		SAMS-Rocket	66	-2.59	2.24	0.657	0.82	1.89
		SSU-Rocket	74	-2.29	1.87	0.768	1.07	1.81
5-1 mbar	2.0	SAMS-SSU	270	1.59	1.28	0.842	0.71	1.00
		SAMS-Rocket	66	-4.21	2.30	0.676	0.81	1.76
		SSU-Rocket	74	- 5.55	2.07	0.740	1.02	1.86
			Kwa	ijalein Islan	d			
Channel 25	1.8	SAMS-SSU	267	-0.20	0.93	0.839	0.83	1.18
		SAMS-Rocket	48	-0.73	1.08	0.824	0.82	1.20
		SSU-Rocket	69	-0.59	0.88	0.876	0.91	1.19
Channei 26	2.0	SAMS-SSU	267	-0.02	0.98	0.842	0.91	1.53
		SAMS-Rocket	48	-1.16	1.30	0.823	0.98	1.44
		SSU-Rocket	69	-1.24	1.13	0.873	1.01	1.32
Channel 27	2.1	SAMS-SSU	267	0.46	0.90	0.872	0.92	1.21
		SAMS-Rocket	48	-1.75	1.58	0.799	0.96	1.51
100 00		SSU-Rocket	69 267	-2.42	1.37	0.845	1.03 0.54	1.44
100–20 mbar	2.1	SAMS-SSU	267 48	1.51	2.03 2.00	0.626 0.673	0.82	1.36 1.81
		SAMS-Rocket	40 69	1.28 0.53	1.49	0.829	1.02	1.48
20-5 mbar	2.3	SSU-Rocket SAMS-SSU	267	-0.35	1.76	0.752	0.66	1.17
20-5 moar	23	SAMS-SSC SAMS-Rocket	48	-0.91	1.81	0.796	0.76	1.20
		SAMS-Rocket	40 69	-0.42	1.61	0.820	0.88	1.20
5-1 mbar	2.2	SAMS-SSU	267	-0.42	1.38	0.823	0.73	1.08
2-1 moan	<b>4</b> -4	SAMS-SSO	48	- 2.36	2.19	0.809	1.00	1.53
		SSU-Rocket	69	-4.12	1.94	0.880	1.31	1.55
			Kwajalein (]			0.000		1.07
Channel 27	1.8	SAMS-SSU	112	0.69	1.10	0.817	0.89	1.33
	1.0	SAMS-SSC SAMS-Rocket	24	- 1.79	1.99	0.722	0.91	1.55
		GUILD ROOM				19.7 da da	v. / 1	1.10
			Kwajalein (C	mly A1 12-1	nbar days)			
Channel 27	1.8	SAMS-SSU	155	0.29	0.68	0.924	0.97	1.14
		SAMS-Rocket	24	-1.71	1.06	0.893	1.01	1.26

TABLE 2. Radiance and Thickness Comparisons at Primrose Lake, Ascension Island, and Kwajalein Island

125

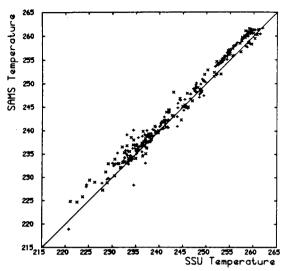


Fig. 4. SAMS and SSU brightness temperatures for SSU channel 27 over Primrose Lake rocket station ( $55^{\circ}N$ ,  $110^{\circ}W$ ) for January 1, 1980 to February 11, 1981; + and × denote days with and without A1 PM data for 12 mbar available.

tropics. The locations of Primrose Lake, Kwajalein Island, and Ascension Island are marked, and it can be seen that the overall means for these places are approximately consistent with mean SAMS-SSU values given in Table 2. Channels 25 and 27 appear to show quite different patterns, with channel 26 behaving in an intermediate way. Channel 27 shows large differences toward 50°S; these are shown in more detail in Figure 6 togeher with 50°N and 67.5°N for comparison. The differences at 50°S are seen to be 2.5-3 K from August to January, with no counterpart in the northern hemisphere. Conversely, for Channel 25 the differences are more symmetric about the equator, with maximum difference occurring in autumn in each hemisphere.

#### 3. DISCUSSION

Since this paper is primarily intended to discuss the SAMSderived temperatures, possible problems with the SSU or rocket/radiosonde will not be discussed. There are several potential problems with the SAMS retrievals: (1) motional chopping in the wideband channels and field of view uncertainties, (2) horizontal gradient effects in pressure-modulated channels at levels where the emissivity is near unity, (3) Doppler shifts caused by the earth's rotation affecting pressuremodulated channels, (4) inadequate allowance for ozone effects in the wideband channels, (5) solar tides, given the consistent differences between measurement times.

Table 3 gives values for the magnitudes of some of these effects. For Doppler shift they are the errors introduced in the tropics by neglecting the earth's rotation, and they were found by analyzing at the equator a special sequence of data when the SAMS was viewing with azimuth angles up to  $12^{\circ}$  away from the normal giving Doppler Shifts of up to  $1600 \text{ m s}^{-1}$ . To quantify other problems a single day (8 February 1981) was taken and processed several times with different assumptions.

#### 3.1. Field of View Uncertainty and Motional Chopping

The SAMS fields of view are approximately Gaussian in the vertical, with a thickness at half peak value of about 8 km.

They were measured during prelaunch tests and, for CO<sub>2</sub>, are subject to an error of the level of the center of  $\pm 0.003^{\circ}$  (0.2 km) [Wale and Peskett, this issue]. However, there is also a potential error caused by motional chopping. This is the result of vibration of the scan mirror (or some other part of the optics) in synchrony with the radiometric chopper. In the case of wideband channels, motional chopping has the effect of changing the field of view. The components were measured during test and have an effect similar to a downward shift of the field of view by 0.02°. The temperature retrieval uses wideband fields of view that include motional chopping on the assumption that it was unchanged between prelaunch test and operation in space. Wale [1981] has shown that for the pressure-modulated channels motional chopping effects should be much smaller than for the wideband. Table 3 gives the effect of changing wideband motional chopping by 100%; it is reasonable to allow possible errors of half of this, giving errors of up to 1 K. Motional chopping is discussed further by Wale and Peskett.

# 3.2 Horizontal Weighting Function Shifts

The weighting functions are primarily functions of log-(pressure). The emission originates within about 200 km horizontally of the tangent point in cases where the total absorption along the tangent path is less than about 0.5 (the optically thin case). However, for absorption near unity the emission is centered about a point nearer the satellite than the tangent point; for the range of data used in the temperature retrieval this horizontal displacement can be up to 300 km. Temperature is assumed constant on pressure surfaces along the line of sight. The net effects of this approximation are (1) the temperature derived for the upper stratosphere and lower mesosphere corresponds to a point roughly 100 km nearer the satellite than the tangent point; (2) the pressure registration

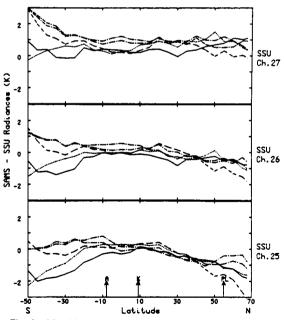


Fig. 5. Monthly zonal mean SAMS-SSU brightness temperatures for SSU channels 25, 26, and 27: (dashed dotted line) January 1980, 1981; (dotted line) April 1980; (solid line) July 1980; (dashed line) October 1980.

will be in error where there is a strong horizontal temperature gradient in the atmosphere. For zonal means the effects cancel at latitudes where the SAMS views primarily east-west, errors only arising north of about 60°N and south of 40°S. In this case a temperature error corresponding to the change in about 100 km is to be expected, thus approximately  $\pm 0.5$  to  $\pm 1$  K, depending on season and pressure, with corresponding tangent height errors of 125 to 250 m (0.002° to 0.004°).

#### 3.3 Doppler Shift

The SAMS views perpendicular to the direction of motion of the satellite, so the radiation suffers no Doppler shift as a result of satellite velocity, except for relatively minor errors caused by the orientation of the satellite. However, the earth's rotation causes a Doppler shift of about  $\pm 400 \text{ m s}^{-1}$  in the tropics, falling to zero at extreme northerly and southerly latitudes. In the case of the CO<sub>2</sub> channels the effects are small, and no attempt is made to allow for them.

#### 3.4 Ozone Contamination

The mid-latitude annual mean ozone profile given by Krueger and Minzner [1976] was used to correct the wideband channels. Table 3 shows the effect of reducing the ozone mixing ratio in the model by 30% at all levels. Consequently, the retrieved temperature increases in the stratosphere to account for the observed emission. Minor decreases occur in the mesosphere because of compensation associated with the pressure-modulated channels being unaffected by ozone. These changes in temperature are to be expected from 30% ozone mixing ratio increases in the atmosphere (more ozone giving a higher apparent temperature). Ozone mixing ratios given by Heath [1980] show maximum deviations of this magnitude about the mean model, with a lower and mid-stratospheric maximum in the tropics and minima at the poles. This would adequately explain the behavior of SAMS-SSU as a function

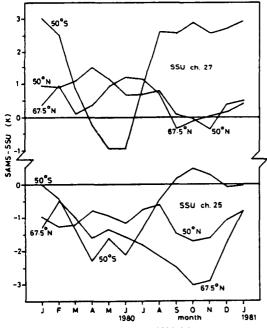


Fig. 6. Monthly zonal mean SAMS-SSU brightness temperature for SSU channels 25 and 27 at 50°S, 50°N, and 67.5°N.

 
 TABLE 3.
 Changes in Temperature Retrieval (K) Resulting From Changes in Assumed Conditions

	Level $(\ln (p_0/p)$				SSU Channel		
	2-4	4-6	6-8	8-10	25	26	27
Ozone less 30%	1.1	1.9	-0.4	-0.3	1.3	1.2	0.4
C1 WB shifted up 0.02°	-0.1	2.2	-0.6	- 1.0	1.0	1.5	0.7
A1 PM not used	0.2	0.5	-0.3	6.3	0.4	0.3	0.0
Doppler shift corrected at equator	-0.4	-0.3	0.4	0.9			

of latitude (Figure 5) for channel 25, although only if the curve is moved upward by about 1 K.

#### 3.5 The Effects of Tides

SSU measurements from two satellites measuring both temperature and its east-west gradient every 6 hours have shown solar tides of up to 3 K amplitude (J. L. Brownscombe and J. Nash, private communication, 1983) for channel 27. Model calculations by C. F. Rogers et al. (private communication, 1983) indicate that in the tropics the semidiurnal tides averaged over the SSU weighting function have amplitudes of  $\sim 1$  K, whereas the diurnal oscillation, because its phase changes rapidly with height, should have a much smaller amplitude. However, at mid-latitudes the diurnal wave would be expected to dominate, with a maximum temperature at about 1800 hours local time. These calculations appear to be confirmed quite well by SSU measurements. Table 1 shows the biases of SAMS and SSU zonal means produced by a diurnal sine wave peaking at 1800 hours. For each satellite the northbound and southbound measurements are assumed to be averaged together, and deviations from the true mean arise because, in general, the two measurements for a given satellite at a given latitude are not 12 hours apart. At 50°S the bias of the SAMS relative to the SSU daily mean is expected to be 104% of the diurnal amplitude, whereas at 50°N, where comparable tidal amplitudes can be expected (in the corresponding season), the bias is only 50%-consistent with the much smaller channel 27 SAMS-SSU differences found there. Thus it is reasonable to ascribe these differences to tides. At 67.5°N, where the bias would be 75% of the tidal amplitude, smaller tidal amplitudes are expected, so the smaller discrepancies there are reasonable. Tidal amplitudes are believed to increase rapidly with height in the stratosphere, so although tides may explain discrepancies in channel 27, their effects should be very much smaller at channel 25 levels, and much less than the observed discrepancies. The SAMS measurements show little difference  $(\sim 1 \ K)$  between northbound and southbound measurements, consistent with a diurnal tide peaking at 1800 hours local time, since the SAMS measurements are symmetrically placed about that time.

Brownscombe and Schlapp [1983] have shown that the lunar semidiurnal tide is detectable in SSU channel 27 data, with a maximum amplitude of about 0.2 K in January and February and a mean of 0.15 K averaged over the year investigated (1980/81). The amplitude is smaller for lower channels. This would have the effect of modulating the differences with a period of 14 days and a zero mean. A simple calculation, using the times given in Table 1, shows that this would cause an rms SAMS-SSU component of 0.15 K in the tropics (averaged over the year). Compared with the smallest observed standard deviation for channel 27, which is 0.79 K at Ascension Island, the lunar tide contributes only 4% of the variance. Hence it has been ignored in this study.

#### 3.6 Comparisons Between Satellites and Rockets

Comparing mean differences between the three sensors at the three rocket stations shows that at Primrose Lake the differences are considerably smaller than at Kwajalein or Ascension Islands. Nash and Brownscombe [1983] found a similar effect, comparing the TIROS-N and the NOAA 6 SSU's with rockets for different periods. They found good agreement to  $\sim 1$  K at mid and high latitudes of the northern hemisphere, but going further south, the rocket-SSU difference progressively increased, being larger at Ascension Island than at Kwajalein Island. The results presented here (see Table 2) support those findings, with the SAMS and SSU generally agreeing with each other but not with the rocket. At Kwajalein Island the difference is only marked for channel 27, but at Ascension Island, differences are evident at all levels. Several potential problems with the SAMS have been outlined, but the offsets they can introduce should all be identical within a fraction of a degree at both tropical stations. This is also true for the SSU, but also no doubt for the rocket measurements, both rocket stations measuring at about the same local time.

#### 4. CONCLUSIONS

There are various, quite different uses to which temperature fields derived from the SAMS measurements will be put, and they determine the types of error that are acceptable:

(1) The temperature fields are used directly as the source of data for study of the dynamical and radiative properties of the atmosphere, e.g., calculation of momentum fluxes. For this application, slow changes of bias with height and latitude of the magnitude found are probably not too serious.

(2) They may be used to monitor long-term trends related to climate change or the solar cycle; here biases do not matter, provided they do not change. Some of the sources of error discussed here can, in principle, drift (notably motional chopping), but there are also others, such as mirror surface degradation and second-order effects associated with  $CO_2$  cell pressure and satellite temperature drifts, that are discussed by *Wale and Peskett* [this issue]. Thus long-term changes of bias must be expected, with a probable magnitude of 1–2 K.

(3) For many studies it is advantageous to combine SAMS with SSU and other data because each measures over different time and space domains with different resolution. However, this is difficult where they do not agree, whether because of temporal effects or instrument errors, because sharp gradients far more serious than the absolute error of either alone arise at the interfaces.

(4) The pressure-scale of the SAMS scan is determined from  $CO_2$  meaurements by a process intimately linked with the temperature sounding such that they both depend strongly on each other. The SAMS minor constituent measurements rely on correct registration of the scan, and since there is no way to verify this, a level of confidence can only be obtained indirectly from the instrument's performance as a temperature sounder. Based on discrepancies found in this comparison, it appears that the pressure-scale registration may have a bias of up to  $\pm 0.006^{\circ}$  in satellite roll angle, corresponding to about 360 m at the limb. The implication of pressure-scale errors on the constituent sounding are outlined by *Jones and Pyle* [this issue].

(5) Temperature measurements are a necessary input for constituent sounding, since thermal emission measured by the minor constituent channels is proportional to the Planck function. The effect of temperature errors is discussed by *Jones and Pyle* [this issue] for the 7- $\mu$ m CH<sub>4</sub>/N<sub>2</sub>O channel. At this wavelength, at 240 K the Planck function changes by 3% K<sup>-1</sup>, and the temperature uncertainties are a limiting factor.

The SAMS has been compared with other sensors in the stratosphere, and differences with random components of 1-2 K have been bound. Mean differences show marked patterns that, to some extent, are understood. They indicate that SAMS temperature measurements are probably subject to biases (excluding those due to tides) of up to 2 K in the low stratosphere and 1 K in the upper stratosphere. The difference is perhaps to be expected, since the upper stratosphere is the middle of the SAMS temperature retrieval range and is the lower level for which the temperature is measured by the pressure modulator channels.

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# APPENDIX E

5302

BARNETT AND CORNEY: SAMS TEMPERATURES VS. SONDES AND SSU

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# OBSERVATIONS OF CH4 AND N2O BY THE NIMBUS-7 SAMS: A COMPARISON WITH IN SITU DATA AND TWO-DIMENSIONAL NUMERICAL MODEL CALCULATIONS

[Reproduced from Journal of Geophysical Research, 89, 5263-5279 (1984).]

# Observations of CH<sub>4</sub> and N<sub>2</sub>O by the NIMBUS 7 SAMS: A Comparison With In Situ Data and Two-Dimensional Numerical Model Calculations

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Monthly mean, zonal mean measurements of  $CH_4$  and  $N_3O$  for 1979 made by the stratospheric and mesospheric sounder (SAMS) on the NIMBUS 7 satellite are presented for the first time. Comparison with in situ and other available data confirms a general pattern of mixing ratios decreasing with height and of maxima, at a given pressure level, in low latitudes. The SAMS data, by virtue of its extensive coverage, reveals new features. These include during certain months a "double peak" when, along a constant pressure surface, mixing ratio maxima are found in low latitudes of both hemispheres with a local minimum at the equator. In the upper stratosphere are found regions of weak and strong honzon-tal gradients. Comparison with a two-dimensional model shows many areas of agreement, especially for  $CH_4$ , and the model is used to interpret atmospheric behavior. The model overestimates the observed  $N_2O$  in the upper stratosphere. This is investigated in terms of a possible underestimation of the photochemical sink. We believe strongly, however, that uncertainties in transport representation cannot be ruled out.

#### 1. INTRODUCTION

The stratospheric and mesospheric sounder (SAMS) instrument launched on the NIMBUS 7 satellite in October 1978 is a multi-channel infrared limb scanning radiometer employing conventional chopping and pressure modulation techniques to measure atmospheric temperature and the abundances of a number of minor constituents by detecting either their thermal emission or, in some cases, resonantly scattering sunlight. In this paper we present monthly mean cross sections of CH4 and N<sub>2</sub>O for 1979 derived from SAMS measurements. The satellite data provide a great increase in spatial and temporal coverage compared with previous measurements and many features which could not be resolved by limited in situ measurements become accessible to investigation. Since the production of CH<sub>4</sub> and N<sub>2</sub>O is at the ground and the photochemical sink is weak in the stratosphere, the satellite data should prove particularly useful for tracing atmospheric motions. Comparison of the zonal mean data with a twodimensional model places a strong constraint on the transport and photochemical schemes used by the model. In areas of good agreement the model can be used diagnostically to make inferences about atmospheric processes. For the comparison we have used the two-dimensional circulation model developed initially at the University of Oxford.

In the following sections, the  $CH_4$  and  $N_2O$  channels of the SAMS instrument are briefly described. An error analysis of the retrieval of the radiance data to give mixing ratio profiles is summarized. The data for 1979 are then presented and a comparison with other measurements made. This is followed by a brief description of the numerical model. After a comparison with the model results, a number of particular features shown by the data are discussed in more detail.

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Paper number 3D1644. 0148-0227/84/003D-1644505.00 2. THE CH4 AND N2O SENSING CHANNEL ON SAMS

A detailed description of the SAMS instrument has been given by Drummond et al. [1980]. Broad spectral selection for the CH<sub>4</sub> and N<sub>2</sub>O (C2/3) sensing channel is provided by a dichroic beam splitter and a filter consisting of three separate components, which together define a spectral passband with maximum transmission at 7.8  $\mu$ m (1280 cm<sup>-1</sup>) and half heights at 7.5 µm (1340 cm<sup>-1</sup>) and 8.3 µm (1200 cm<sup>-1</sup>). Atmospheric emission in this region of the spectrum arises primarily from the P and Q branches of the  $v_{\pm}$  band of CH<sub>4</sub> and the  $v_{\pm}$ band system of N2O, with minor additional contributions from CO<sub>2</sub>, H<sub>2</sub>O, SO<sub>2</sub>, O<sub>3</sub>, and HNO<sub>3</sub>. Atmospheric radiant energy directed into the C2/3 channel passes through two pressure modulated cells (PMCs), the first containing N2O gas and the second CH<sub>4</sub> gas, which act as optical filters allowing emission from each of these species to be detected selectively [see Drummond et al., 1980]. The mean pressure of gas in each modulator can be set to either of two preselected pressures (see Table 1) extending the useful vertical coverage that may be obtained with each channel. To illustrate this effect, in Figure 1 are curves showing the variation with height of the emissivity of the CH<sub>4</sub> and N<sub>2</sub>O channels for the two available mean PMC pressures. The best measurements are obtained over the height range where the emissivity gradient is large. with little useful information being obtained from the levels where the emissivity is close to zero or unity. By altering the mean PMC pressure, the spectral response of the pressure modulated channel is moved closer to or further from the centers of emission lines, shifting the region of maximum emissivity gradient in the vertical. The effect is evident in the curves shown in Figure 1 (compare curves a and b and curves c and d). The emissivity curves of the N2O channel tend to zero at lower levels than those of the CH, channel because of the much more rapid drop-off of N2O with height.

Because the CH<sub>4</sub> and N<sub>2</sub>O PMCs share a common detector, only one channel can be enabled at any one time so that

5263

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5264

TABLE 1.	Nominal Mean	PMC	Pressures	in	the	CH₄	and	N <sub>2</sub> O
Pressure Modulators							•	

Channel	Mode	Mean PMC Pressure and Estimated Error, mbar
C2 (N <sub>2</sub> O)	C2S0 C2S1	7.15 (±3.5%) 24.4 (±2.0%)
C3 (CH <sub>4</sub> )	C3S0 C3S1	22.5 (±1.0%) 47.8 (±1.0%)

Cell length 3 mm, approximate compression ratio: 1.7.

simultaneous measurements of both gases are not possible. Normally, therefore, measurements of CH<sub>4</sub> and N<sub>2</sub>O are made on alternate days and this, together with the SAMS duty cycle of 3 days on and 1 day off, means that measurements of each species are possible for  $\sim 40\%$  of the time ( $\sim 12$  days a month).

#### 3. RETRIEVALS

The signal-to-noise ratio of the measured radiances was not in general sufficient to obtain individual profiles of CH<sub>4</sub> and N<sub>2</sub>O. Normally, therefore, we derived zonal mean mixing ratio profiles directly from zonally averaged radiance profiles and effective zonal mean temperature profiles. To account for the nonlinearity of the Planck function at 7.8  $\mu$ m, at each level we first average zonally the Planck function at this wavelength using the SAMS derived temperatures [see *Barnett and Corney*, this issue]. Each zonal mean Planck function profile for each latitude band. Some error does result through averaging zonally rather than by using smaller longitudinal segments and then averaging the retrieved profiles but, as is discussed further in section 4.3, we found this to be minor when compared with other systematic errors.

Normally, we averaged radiances over  $10^{\circ}$  latitude bands and 0.2 scale height (~1.4 km) intervals. The choice of these values is somewhat subjective, with these representing a suitable compromise between latitudinal resolution and the computation time required for the retrievals. The signal-to-noise ratios of the zonal mean radiance profiles depend greatly on the atmospheric temperature structure (see Figure 2) and the instrument channel used. Signal-to-noise ratios reached typically ~30 in the low stratosphere. The restrictions that this places on the accuracy of retrieved solutions are discussed in section 4.5.

The retrieval method employed a sequential estimator to update an a priori profile and covariance. To ensure rapid convergence, the most recent previous measurement at each latitude was used for the a priori profile with an overestimated variance. This method is described in detail elsewhere in this issue [Rodgers et al., this issue].

#### 4. Accuracy and Precision of the SAMS $CH_4$ and $N_2O$ Measurements

An important feature of atmospheric measurements made from a satellite-mounted instrument such as SAMS is that spatial and temporal variations in atmospheric state become accessible to measurement. It is particularly important, therefore, that those errors that may introduce spurious features of these kinds are well understood and their magnitudes known. In the following discussion of error sources we have attempted, where possible, to distinguish between the errors of this type and those which to first order only lead to constant

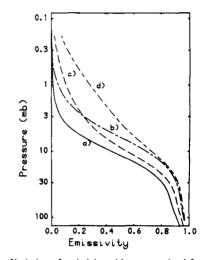


Fig. 1. Variation of emissivity with pressure level for the SAMS  $N_2O$  and  $CH_4$  channels assuming typical equatorial mixing ratio and temperature profiles. Curves a and b are for the  $N_2O$  channel for mean PMC pressures of 24.4 and 7.15 mbar, respectively, and curves c and d are for the CH<sub>4</sub> channel, 47.8 and 22.5 mbar, respectively.

biases in the retrieved fields. Where the distinction cannot be made reliably, we have, if anything, erred on the conservative side.

Systematic errors in the retrieved constituent fields fall broadly into four categories: uncertainties in the spectroscopy of the absorbing gases; instrumental uncertainties; limitations and simplifications in the retrieval method and algorithm (e.g., uncompensated Doppler shifts and errors associated with the zonal averaging of radiances and temperatures); and inaccurate knowledge of the atmospheric state (primarily of the temperature field). These various error sources and their impacts

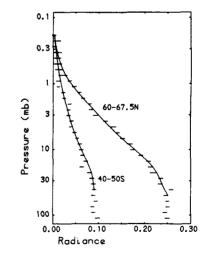


Fig. 2. Examples of zonal mean radiance profiles for July 2, 1979 (expressed as a fraction of the signal observed from a black body at 290 K). The horizontal bars denote measurements and their  $(\pm 1\sigma)$  errors, and the solid curves are synthetic profiles computed from the retrieved mixing ratio profile. The difference between the profiles is mainly due to the large interhemispheric temperature gradient present at the time.

#### JONES AND PYLE: SAMS OBSERVATIONS VERSUS MODEL CH4 AND N2O

Pressure Level	20 mbar	7 mbar	2 mbar	0.6 mbar	0.2 mbar
1. Spectroscopy	±7%	±7%	± 5%	±5%	± 5%
2. Field of view	±2%	± 1%	± 3%	+6%	+ 3%
3. Mean PMC pressure	± 3%	±2%	±6%	±6%	± 3%
<ol> <li>Uncompensated Doppler shifts</li> </ol>	±1.5%	±,1%	±1.5%	±6%	±7%
5. Zonal averaging	±8%	±8%	±8%	±8%	±8%
6. Interference of N <sub>2</sub> O	± 5%	±2%	±0.5%		-
7. Temperature $(\pm 2 K)$	± 50%	±15%	±10%	±10%	±10%
8. Line-of-sight attitude (±0.006°)	±4%	±2%	±6%	±10%	±6%
9. RSS of latitude and time dependent errors	±51%	±15%	±12%	±17%	±14%
0. RSS of bias errors	±9%	±6%	±8%	±10%	±7%
1. Net RSS accuracy	$\pm 52\%$	$\pm 18\%$	$\pm 17\%$	$\frac{1}{\pm}21\%$	±17%
12. Precision of monthly mean cross section	±5-15%	$\pm 3 - 10\%$	±3-13%	±4-13%	±20-40%

TABLE 2a. Summary of Error Budget for the C3 (CH<sub>4</sub>) Sensing Channel on SAMS

on the retrieved constituent fields are described below. Other more minor error sources also exist, and, although these are not discussed individually below, their effects have been included in the error budget summaries given in Tables 2a and 2b.

To compute the magnitudes of systematic errors, a synthetic radiance profile was computed for typical mixing ratio and temperature profiles, with all the uncertain parameters set to their nominal values. This data set was then retrieved with each uncertain parameter offset in turn within its limit of uncertainty, and the retrieval is then compared with the original. Several tests were performed that demonstrated that the errors were sufficiently insensitive to the temperature and mixing ratio profiles used for the results presented below to be of general applicability.

#### 4.1. Spectroscopic Errors

Uncertainty in the atmospheric absorptions of  $CH_4$  and  $N_2O$  was minimized by measuring in the laboratory the response of sensors, identical in design and construction to those used in the flight instrument, to simulated atmospheric limb paths containing these gases. Corresponding line-by-line transmission calculations were then performed, using spectral data from the compilations of *McClatchey et al.* [1973] and *Rothman* [1981], which demonstrated that minor improve-

ment and modifications to the spectral data were required [Jones, 1983]. Using the modified spectral data, residual uncertainty in the absorptions of CH<sub>4</sub> and N<sub>2</sub>O is estimated to be equivalent to bias error of  $\sim \pm 5-7\%$  in the retrieved mixing ratios of both CH<sub>4</sub> and N<sub>2</sub>O (Tables 2a and 2b, row 1).

#### 4.2. Instrumental Uncertainties

The  $CH_4/N_2O$  channel field of view was measured as an integral part of the prelaunch testing to be approximately Gaussian in the vertical with a half width at half height of ~4 km. Some uncertainty is present in this measurement which may be thought of as being approximately equivalent to an uncertainty in view angle of  $\pm 0.003^{\circ}$  or approximately  $\pm 0.2$  km at the tangent point. The effect of such an error is to change the vertical gradient of the retrieved profiles, leading to biases of up to  $\pm 5\%$  in the retrieved mixing ratios of  $CH_4$  and up to  $\pm 7\%$  in those of the N<sub>2</sub>O (see Tables 2a and 2b, row 2).

The mean pressure of gas in a PMC is critical in defining the spectral response of that channel (refer for example to Figure 1); any error leads to both a vertical gradient change and an offset in the derived fields. The gas pressures in the CH<sub>4</sub> modulator are known to  $\pm 1\%$  and those in the N<sub>2</sub>O modulator to  $\pm 3.5\%$  and  $\pm 2\%$  for the low and high mean cell pressures, respectively (see Table 1). These lead to bias

TABLE 2b. Summary of Error Budget for the C2 (N<sub>2</sub>O) Sensing Channel on SAMS

Pressure Level	20 mbar	7 mbar	2 mbar	0.6 mbar
1. Spectroscopy	±7%	±7%	± 5%	±5%
2. Field of view	$\pm 2\%$	±4%	±7%	±8%
3. Mean PMC pressure	± 3%	±2%	±17%	±21%
4. Uncompensated Doppler shifts	±3%	± 2%	±45%	±15%
5. Zonal averaging	+8%	±8%	±8%	±8%
6. Interference of CH <sub>4</sub>	+8.5%	±2%	± 5%	<u>+</u> 25%
7. Temperature $(\pm 2 \text{ K})$	± 50%	$\pm 15\%$	±10%	<u>+</u> 20%
8. Line-of-sight attitude (+0.006°)	±4%	±8%	±12%	±3%
9. RSS of latitude and time dependent errors	± 51%	±17%	±48%	±35%
0. RSS of bias errors	+9%	±7%	±20%	±24%
1. Net RSS accuracy	$\frac{-}{\pm}$ 52%	$\pm 20\%$	± 50%	±43%
2. Precision of monthly mean cross section	$\pm 12 - 20\%$	$\pm 6 - 17\%$	$\pm \overline{10}-24\%$	± 50-100%

5266

errors of up to  $\pm 6\%$  in the retrieved CH<sub>4</sub> mixing ratios (Table 2a, row c) and up to  $\pm 25\%$  in those of N<sub>2</sub>O (Table 2b, row 3).

#### 4.3. Limitations in the Retrieval Method

A problem inherent in using a gas correlated device such as a pressure modulator is the Doppler shifting of atmospheric emission lines relative to those in the PMC on the orbiting spacecraft. In normal operation the instrument line-of-sight is nominally perpendicular to the direction of spacecraft motion so that the shift is due to the combined effects of the earth's rotation and atmospheric winds. In practice, an additional component arises because of yaw angle errors in the spacecraft orientation (estimated to be  $\pm 0.6^{\circ}$ ). The dominant component-that due to the earth's rotation-can be precomputed from a knowledge of the instrument viewing geometry and is therefore included explicitly in the model for calculating atmospheric radiances. The other sources cannot be dealt with reliably and have not been allowed for. This results in uncertainties of up to ~5% in the CH<sub>4</sub> abundances and up to  $\sim \pm 30\%$  at higher levels in those of N<sub>2</sub>O (see Tables 2a and 2b, row 4).

As was stated above, zonal mean mixing ratio cross sections have been derived directly from zonal mean radiance and temperature and cross sections. While this approach is likely to be satisfactory when the temperature and constituent fields are zonally symmetric (frequently the case in the summer hemisphere), significant errors might arise during more active winter periods when large waves in temperature and constituent concentration may be present around a latitude circle. A direct test of the accuracy of this approach was therefore made, using as an example a single days data (February 9, 1981). At this time a large temperature wave was present at high northern latitudes with an amplitude in excess of 30 K peak-to-peak between 30 and 1 mbar. In contrast, the temperature field in the summer (southern) hemisphere was almost zonally symmetric. Radiances in the latitude bands from 60°N to 67.5°N and from 50°S to 40°S were retrieved by using different numbers of boxes around each latitude circle, namely, (1) one  $360^{\circ}$  longitude box (the standard), (2) four  $90^{\circ}$  longitude boxes, and (3) twelve 30° longitude boxes.

A mixing ratio profile was retrieved for each box. A zonal mean mixing ratio profile was then obtained for each test by averaging the individual boxes together with due regard for their uncertainties. As expected, the zonal mean profiles for the band between 50°S and 40°S were almost indistinguishable, while those for the  $60^{\circ}N-67.5^{\circ}N$  band showed minor differences, with an rms scatter of between 3.5 and 8% at different heights. However, in spite of these differences the important features were faithfully reproduced in all three profiles. We have taken an error of  $\pm 8\%$  as being the maximum error incurred through the zonal averaging of radiances (Tables 2a and 2b, row 5). This, of course, represents a worst case, appropriate only when large deviations from a zonally symmetric atmospheric state are present.

#### 4.4. Uncertainties in the Atmospheric State

Atmospheric thermal emission in the 7.8  $\mu$ m region of the spectrum arises predominantly from CH<sub>4</sub> and N<sub>2</sub>O, with minor additional contributions from CO<sub>2</sub>, H<sub>2</sub>O, O<sub>3</sub>, SO<sub>2</sub>, and HNO<sub>3</sub>. The use of a gas correlated technique like pressure modulation virtually excludes emission from the minor contributors, although the mutual overlap of the CH<sub>4</sub> and N<sub>2</sub>O

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bands is such that a small correction caused by absorption by the other must be applied in the retrieval of both species. As the atmospheric distributions of both species are of course measured by the SAMS, we have applied a correction using the most recent previous measurement of the abundance of the interfering gas at that latitude. Uncertainties in abundances of the contaminating gas lead to uncertainties of  $\pm 5\%$ in the retrieved mixing ratios of CH<sub>4</sub> at low levels, decreasing rapidly with height, and <6% over most of the stratosphere for N<sub>2</sub>O, increasing to ~25% at the stratopause level (see Tables 2*a* and 2*b*, row 6).

An error in the assumed atmospheric temperature structure leads to an anti-correlated error in the retrieved minor constituent fields. The atmospheric temperatures used in the retrievals are derived from SAMS measurements of thermal emission from the 15- $\mu$ m band of CO<sub>2</sub> in a coupled temperature/line-of-sight attitude retrieval. Uncertainty in the retrieved temperature field is estimated to be  $\pm 2$  K [Barnett and Corney, this issue] and  $\pm 0.006^{\circ}$  in attitude, equivalent to  $\pm 0.37$  km at the tangent point (J. J. Barnett, private communication, 1983), both with possibly significant latitudinal and seasonal components. An uncertainty in temperature of  $\pm 2$  K translates to a mixing ratio error of  $\gtrsim \pm 15\%$  (See Tables 2a and 2b, row 7), while an uncertainty in attitude of  $\pm 0.006^{\circ}$  gives rise to an error of  $\lesssim 10\%$  (Tables 2a and 2b, row 8).

#### 4.5. Random Errors

Random errors on the retrieved solutions are predominantly due to radiance noise. We can determine the effects of radiance noise on the retrieved mixing ratio profiles straightforwardly from the diagonal elements of the error covariance matrix obtained from the retrieval process [see *Rodgers et al.*, this issue]. Typical values for the random errors on a monthly mean cross section obtained in this way are shown in Tables 2a and 2b, row 12. A precise value cannot be given as this would depend on the atmospheric temperature structure (which as can seen from Figure 2 largely determines the signal-to-noise ratios of the radiance measurements) and to a lesser extent on the temporal coverage obtained during each month. Errors obtained in this way are overestimates in that they take no account of the correlations that are present between the retrieved mixing ratios at the various levels.

#### 4.6. Summary of Errors

For both channels, insufficiently accurate knowledge of the atmospheric temperature structure and to a lesser extent of the line-of-sight attitude are potentially significant sources of spurious latitudinal or seasonal features. For the N<sub>2</sub>O channel, uncompensated Doppler shifts are another potentially serious problem, particularly in the upper stratosphere (see Table 2b row 4) where the emission lines of N<sub>2</sub>O at 7.8  $\mu$ m are Doppler broadened and hence narrow. Comparison of Tables 2a and 2b, rows 9 and 10 shows that purely bias errors are generally small in comparison.

The net RSS accuracies of the CH<sub>4</sub> and N<sub>2</sub>O measurements made by the SAMS are shown in percentage terms at various heights in Tables 2a and 2b row 11. Over much of the stratosphere the CH<sub>4</sub> measurements are seen to be superior to those of N<sub>2</sub>O with an RSS accuracy of  $\leq 20\%$  compared with 25-50% for the latter. This is a consequence of the greater sensitivity of the N<sub>2</sub>O channel to unwanted Doppler shifts and to interference by CH<sub>4</sub> in the upper stratosphere and is offset to some extent by the more rapid vertical drop-off of  $N_2O$  mixing ratio with height and the consequential increase, in percentage terms, of the features observed (e.g., latitudinal gradients) in the cross sections. These values are to be compared with the precisions at various heights obtained on a monthly mean cross section (Tables 2a and 2b row 12).

The vertical extent of useful  $CH_4$  and  $N_2O$  measurements made by the SAMS is influenced by a number of factors. The spectrocopy and the atmospheric distributions of each species, together with the spectral response of each channel, couple to define the height range over which information may in principle be obtained (refer for example to Figure 1). In practice, however, the finite signal-to-noise ratio of the radiance measurements and the sensitivity of the retrievals to systematic errors reduce the useful coverage. We feel that useful measurements of monthly zonal mean cross sections of  $CH_4$ are made between 20 mbar (~30 km) and 0.2 mbar (~60 km) and of  $N_2O$  between 20 mbar (~30 km) and 0.6 mbar (~53 km). The data presented below have been restricted to within these confidence limits.

#### 5. THE DATA SET

Since launch in October 1978, over 4 years of CH4 and N<sub>2</sub>O radiance measurements have been accumulated. Data from the whole of 1979 and a 3 month period in early 1981 have been retrieved and studied extensively. In the following discussion we will concentrate our attention on the first full year of data. In Figures 3a-3l and 4a-4l are shown monthly mean, zonal mean cross sections of N<sub>2</sub>O and CH<sub>4</sub> for January-December 1979. Mixing ratios are plotted only between the confidence limits discussed in the previous section. Note that the coverage is asymmetric about the equator, extending from 50°S to 70°N. The gross structure suggested by these cross sections is of a low latitude, low stratosphere maximum, with the mixing ratios of both CH<sub>4</sub> and N<sub>2</sub>O decreasing with height throughout the stratosphere, the former by a factor of  $\sim 6$ , the latter much more strongly, by a factor of  $\sim 60$ . Superimposed on this basic structure are large, apparently seasonal, changes. Of particular note is the marked correlation between the changes observed in the CH<sub>4</sub> and N<sub>2</sub>O cross sections throughout the year (simple inspection of Figures 3 and 4 is sufficient to reveal this).

During January (refer to Figures 3a and 4a) the distributions of CH<sub>4</sub> and N<sub>2</sub>O showed a marked asymmetry about the equator, with a region of elevated mixing ratios extending throughout the stratosphere, tilting from the equator at  $\sim 20$ mbar into the summer (southern) hemisphere, reaching a latitude of  $\sim 20^{\circ}$ S at the stratopause level. By way of contrast, at this time a more rapid drop-off of mixing ratio was evident in low northern (winter) latitudes with a region of weaker gradients above. This structure persisted through February (Figures 3b and 4b)), but during March and April (Figures 3c and 3dand 4c and 4d) the southern hemisphere maximum began to subside, and a second similar feature began to develop in the northern hemisphere, giving by May (Figures 3e and 4e) an almost symmetric pattern. This "double peak" structure in which two low latitude maxima, on a fixed pressure surface, are separated by a local mixing ratio minimum near the equator will be discussed in detail below.

By July, the pattern was a reversal of that of January, with an area of elevated mixing ratio now extending into the northern hemisphere with again an area of weaker vertical gradients in the winter mid-stratosphere (Figures 3g and 4g). The structure during the second half of 1979 differed in many respects from that of the first. During August and September throughout the middle stratosphere there existed a pronounced maximum at low latitudes. At high northern latitudes very little height variation of the mixing ratio is evident in the upper stratosphere (Figures 3h and 3i and 4h and 4i). Coverage does not extend sufficiently far into the southern hemisphere to establish the existence of a corresponding feature 6 months earlier. By October, the maximum had subsided somewhat, although the region of uniform mixing in the vertical still persisted. The "double peak" structure evident during May was not observed. In the final months of 1979, the maximum continued to subside giving, by December, an almost symmetric distribution (Figures 3l and 4l).

The high degree of correlation between the distributions of  $CH_4$  and  $N_2O$  is expected because both originate in the troposphere, and their long photochemical lifetimes (of the order of many months for both gases in the mid-stratosphere) mean that the distributions of both species are determined in large part by the same dynamical processes. It is important to note that because of the more rapid vertical decrease of  $N_2O$ , the errors required to introduce similar spurious features into the distributions of both species are radically different. However, the sensitivities of both species to known errors are of comparable size or do not influence the  $CH_4$  distribution significantly (see Tables 2a and 2b), and we therefore have considerable confidence that the features evident in Figures 3 and 4 are substantially real and are not artifacts of the instrument calibration or the retrieval algorithm.

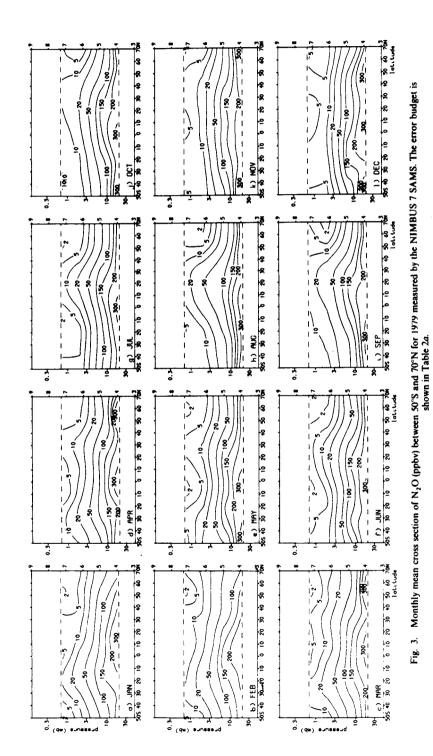
#### 6. COMPARISON WITH OTHER MEASUREMENTS

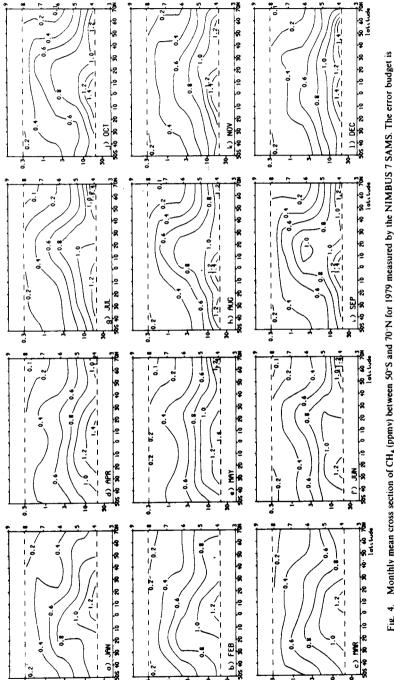
Vertical profiles of CH<sub>4</sub> and N<sub>2</sub>O have been measured on numerous occasions over the last decade, the majority by local sampling using either evacuated grab samplers or cryosamplers although several remote measurements using infrared absorption techniques have also been made [see World Meteorological Organization, 1982] (WMO). These data do not provide uniform spatial coverage, being strongly weighted to latitudes of 32°N, 44°N, and 52°N for CH<sub>4</sub> and to equatorial latitudes and 44°N for N<sub>2</sub>O.

Because of the large variability evident in the profiles of  $CH_4$  and  $N_2O$  measured by other workers (much larger than the precisions of 2-5% quoted by the individual experimenters) and the absence of any clearly defined seasonal trends, we feel that a comparison of SAMS zonal mean  $CH_4$  and  $N_2O$  profiles with the few profiles measured during 1979 would be of dubious value. Instead, in Figures 5-10 we compare all the available measurements made using in situ and other remote techniques (as presented in WMO [1982]) with annual mean profiles of  $CH_4$  and  $N_2O$  for 1979 measured by SAMS at various latitudes and leave a discussion of the observed seasonal changes to section 8.

As the SAMS measurements are made with respect to pressure coordinates while those of other workers are generally quoted in relation to geometric height, we have used a climatological annual mean, zonal mean temperature cross section kindly provided by J. J. Barnett (private communication, 1983) to relate the two.

For CH<sub>4</sub>, SAMS and other measurements are consistent where overlap occurs (the SAMS measurements extend to much higher levels). At low latitudes ( $< 10^{\circ}$ ), there are too few measurements to draw firm conclusions, but the SAMS profile and the one available in situ measurement are consistent





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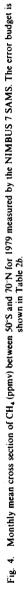
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JONES AND PYLE: SAMS OBSERVATIONS VERSUS MODEL CH4 AND N2O

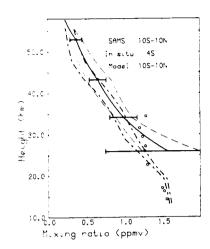


Fig. 5. Comparison of the SAMS CH<sub>4</sub> annual mean profile (solid line) for  $10^{\circ}$ S to  $10^{\circ}$ N with other measurements (circle, *WMO* [1982]). The dashed envelope (dashed line) shows the standard deviation of the monthly mean profiles for the latitude band and the horizontal bars the estimated accuracy. The heavy dot-dash envelope shows the annual variation of the model profiles over the indicated latitude band.

(Figure 5). Between 25°N and 35°N, many more measurements have been performed. These show an almost linear drop-off of mixing ratio with height from 1.6 ppmv at 15 km to 0.5 ppmv at 40 km, although with much greater variability than at other latitudes. A similar drop-off is evident in the SAMS measurements, reaching a somewhat higher mixing ratio of 0.7 ppmv at 40 km, with a region of weaker vertical gradients above (Figure 6). Between 40°N and 60°N (Figure 7), the in situ measurements show a rather more rapid, although still approximately linear drop-off of mixing ratio with height up to 35 km. The SAMS measurements reproduce well the vertical gradient, showing the drop-off extending to above 40 km with, again, a region of weak vertical gradient above. The SAMS measurements are also consistent with the one profile measured at 65°N (Figure 8).

For  $N_2O$ , at equatorial latitudes (Figure 9) above 30 km the

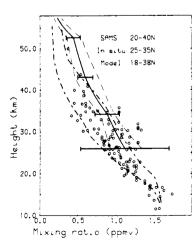


Fig. 6. As Figure 5, except at 20°-40°N.

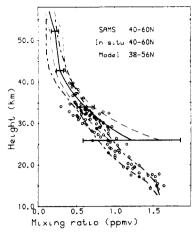


Fig. 7. As Figure 5, except at 40°-60°N.

SAMS measurements agree well with other measurements both in terms of vertical gradient and absolute amount. Below 30 km the SAMS measurements are biased high, but nevertheless the two data sets are consistent to within their estimated accuracies. At 45°N (Figure 10), although the SAMS measurements reproduce fairly well the vertical gradient suggested by the in situ measurements, they appear biased high relative to the latter, by 20–30% at 30 km and rather more below. It is conceivable that this discrepancy is exacerbated by a combination of the more rapid vertical drop-off of N<sub>2</sub>O compared with that of CH<sub>4</sub> and inaccuracies in relating the geometric height and pressure scales of the two data sets.

Overall, the SAMS measurements reproduce well the gross features shown by other measurements, namely, the low stratosphere, low latitude maxima, the almost linear decrease of  $CH_4$  with height at mid- and high latitudes, and the much more rapid vertical drop-off of N<sub>2</sub>O. There is some indication of a bias (SAMS higher than other measurements), particularly in the low stratosphere at high latitudes, although the two data sets are consistent to within their estimated accuracies.

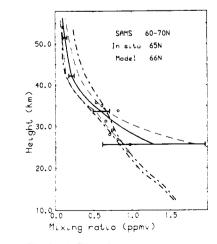


Fig. 8. As Figure 5, except at 60°-70°N.

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## JONES AND PYLE: SAMS OBSERVATIONS VERSUS MODEL $CH_4$ and $N_2O$

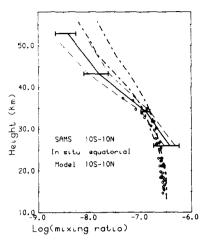


Fig. 9. As Figure 5, except the SAMS  $N_2O$  is plotted for the 10°S-10°N latitude band.

## 7. MODEL

The model used in this study is the two-dimensional circulation model described by Harwood and Pyle [1975]. The model calculates the zonal mean values of temperature, wind components, and chemical constituent mixing ratios with a resolution of  $\pi/19$  in latitude, 0.5 in  $\ln (p_0/p)$  (approximately 3.5 km) in the vertical and with a 6 hour time step. A secondorder partial differential equation is solved for the meridional stream function given the forcing by radiative and other diabatic heating and eddy heat and momentum fluxes. The dynamical and radiative formulation employed is that described by Haigh and Pyle [1982].

The horizontal eddy momentum fluxes are derived from data from the selective chopper radiometer flown on NIMBUS 5. Monthly means of the 1973 values are used. The eddy fluxes of heat and tracers are calculated by using the Kcoefficients derived by Luther [1973] and based on atmospheric statistics. It is clear that by using these data we should not expect the model to reproduce the behavior of any particular month's synoptic situation, still less that it should reproduce behavior on a shorter time scale. Neither will the model necessarily reproduce long-term average behavior since, for example, the momentum fluxes were obtained from just one year of data. However, experiments using different K coefficients and momentum fluxes [Harwood and Pyle, 1980] still preserve the main features of the month-to-month changes in the model and maintain the broad agreement found on comparison with observations. Thus, we are confident that the model development is at least representative of some mean atmospheric behavior. This should be borne in mind when considering the comparisons presented later.

The photochemical scheme is based on that used by Haigh and Pyle [1982] in which family groupings are employed. The kinetic data have been updated by the recommendations of WMO [1982] and Baulch et al. [1980]. HO<sub>2</sub>NO<sub>2</sub> and HOCI are now included in the photochemical scheme, and in this case continuity equations are also solved for CH<sub>4</sub> and N<sub>2</sub>O with sinks calculated using the reactions in Table 3. At the bottom boundary the mixing ratios of CH<sub>4</sub> and N<sub>2</sub>O are fixed at their tropospheric values. The upper boundary at approximately 60 km is a rigid lid to tracers. Unlike in Haigh and

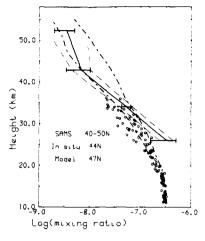


Fig. 10. As Figure 9, except at 40°-50°N.

Pyle, the photochemical model used here contains a source of chlorine which simulates the natural background.

We have discussed above some of the problems related to comparing model and observations arising from the representativeness of the model transport fields. For nonconservative tracers, like  $CH_4$  and  $N_2O$ , the comparison will also depend, for example, on how well the model reproduces the fields of  $O(^1D)$ , Cl, and OH, with which  $CH_4$  and  $N_2O$ react. A detailed analysis of the behavior of all the model constituent fields and their budgets is beyond the scope of this paper, and, indeed, for the important radicals the data base is not adequate to allow latitudinal and seasonal comparisons. Nevertheless, it can be stated that for OH and Cl there is reasonable agreement between model profiles and the very limited number of observations.

#### 8. COMPARISON WITH A TWO-DIMENSIONAL MODEL

## 8.1. Model Behavior

In Figures 11 and 12 are shown cross sections of the  $CH_4$ and  $N_2O$  volume mixing ratios for equinox and solstice from a model run, run A. Notice that the model domain covers the entire globe. There is a broad similarity in the two distributions, as would be expected for gases that originate at the surface and are transported into the stratosphere where their lifetimes against photochemical destruction are long.

The model exhibits much less variability than the observations on "short" spatial ( $< 20^{\circ}$  latitude) and temporal (< month) scales. This is to be expected since the model is driven by, for example, radiative heating rates and eddy fluxes which are either smoothly varying or monthly averaged.

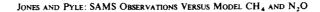
TABLE 3. Model Photochemical Sinks for CH<sub>4</sub> and N<sub>2</sub>O

	Rate Constants
$N_2O + h\nu \rightarrow N_2 + O$ $N_2O + O(^1D) \rightarrow 2NO$ $N_2O + O(^1D) \rightarrow N_2 + O_2$ $CH_4 + O(^1D) \rightarrow CH_3 + OH$ $CH_4 + CI \rightarrow CH_3 + HCI$ $CH_4 + OH \rightarrow CH_3 + H_2O$	7.2 × $10^{-11}$ 4.4 × $10^{-11}$ 1.4 × $10^{-10}$ 2.4 × $10^{-12} \exp(-1710/T)$ 9.6 × $10^{-12} \exp(-1350 T)$

Rate constants are in units of cm<sup>3</sup> molec<sup>-1</sup> s<sup>-1</sup>.

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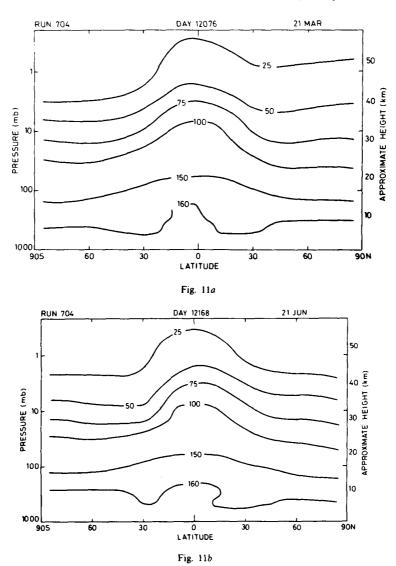


Fig. 11. Model distributions of CH<sub>4</sub> mixing ratio (ppmv  $\times$  100) for (a) March, (b) June, (c) September, (d) December.

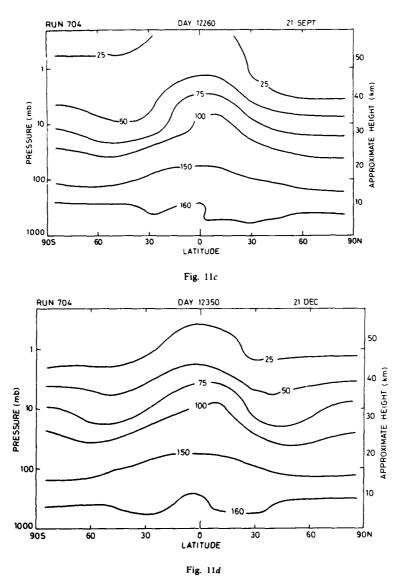
Short-term or small-scale features, which might have a significant influence on the monthly mean observations for a particular month, should not be expected to be reproduced by the model. It should be noted, for example, that there is much larger variation in the monthly observations than found between successive months of the model. The fact of atmospheric variability must be borne in mind when considering the comparisons discussed below.

The model behavior is quite simple. For any constant pressure surface the maximum mixing ratios are generally found throughout the year in equatorial latitudes. The latitudinal excursions of this maximum are small with the maximum confined to within 10° of the equator. The model does not produce a "double peak," a point discussed in more detail below.

The distributions for the solstices (Figures 11b, 11d, 12b, and 12d) are quite symmetric. In contrast, there are significant hemispheric differences in the upper stratosphere at the equinoxes (Figures 11a, 11c, 12a, and 12c). For example, in March

(Figure 12a) the upper stratospheric  $N_2O$  mixing ratio decreases from about 25 ppbv at the equator to less then 5 ppbv at the southern pole. The photochemical lifetime of  $N_2O$  decreases in the upper stratosphere, being a few weeks at 1 mbar, and the decrease of  $N_2O$  in the high latitudes of the sunlit upper stratosphere is, at least in part, due to the integrated photochemical destruction there throughout the polar summer. Six months later the pattern is reversed with a mixing ratio at the southern hemisphere stratopause of about 10 ppbv, produced by transport from the summer hemisphere and the much reduced loss by photochemical processes.

An interesting feature of the model behavior (see, e.g., Figure 11d) is the presence of a mixing ratio minimum in middle latitudes of the midstratosphere. This appears to be produced when the reverse mean meridional circulation is established in high latitudes, carrying high mixing ratios aloft at the pole and bringing down air relatively poor in  $CH_4$  and  $N_2O$  into middle latitudes. This Eulerian circulation is op-



posed by eddy transports. Nevertheless, the net effect in the model is to produce the mid-latitude minimum.

#### 8.2. Comparison With Data

When the model and data are compared, there is satisfactory agreement with respect to the gross features. For both gases there is a decrease of mixing ratio with altitude, this decrease being most pronounced in high latitudes. Both model and data have equatorial or low latitude maxima associated with upward transport in the stratospheric extension of the Hadley cell.

In both observations and calculations the decrease with height is more pronounced for  $N_2O$  which has a much shorter photochemical lifetime than  $CH_4$  in the upper stratosphere. This also results in a greater dynamically produced variability of  $N_2O$  in the summer upper stratosphere, as discussed above in connection with the model fields and confirmed by the observations. Thus, transport in the meridional plane converts the large  $N_2O$  vertical gradients into correspondingly large horizontal gradients.  $CH_4$  is more nearly mixed. For example, there is a factor of 20 difference between N<sub>2</sub>O mixing ratios in low and high northern latitudes at about 2 mbar in September, compared with a factor of only about 2 for  $CH_4$  (Figures 11c and 12c).

As discussed above, the model, at certain times of the year, produces mid-latitude minima in  $N_2O$  and  $CH_4$  in the middle stratosphere. As a consequence of the more limited latitudinal coverage provided by the data it is not possible to confirm this feature with confidence, although there are suggestions of such a feature in, for example, Figure 4c.

To consider more quantitative aspects of the comparison, the range of  $CH_4$  and  $N_2O$  mixing ratios over a model year has been plotted at a number of latitudes in Figures 5-10 along with the SAMS data and the other available observations. For  $CH_4$  the agreement between model and SAMS is generally extremely good (Figures 5, 7, 8). In the low stratosphere the model gives mixing ratios below those observed, but these appear to be biased high when compared with the in

# Jones and Pyle: SAMS Observations Versus Model $CH_4$ and $N_2O$

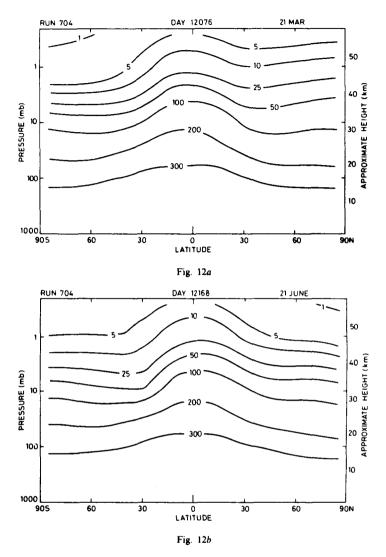


Fig. 12. Model distributions of N<sub>2</sub>O mixing ratio (ppbv) for (a) March, (b) June, (c) September, (d) December.

situ measurements. The upper stratospheric values agree well. Notice that the model reproduces the change in the vertical gradient found between 40 and 45 km. We have also found this feature in a model whose vertical domain extended to 80 km, thus ruling out spurious boundary effects.

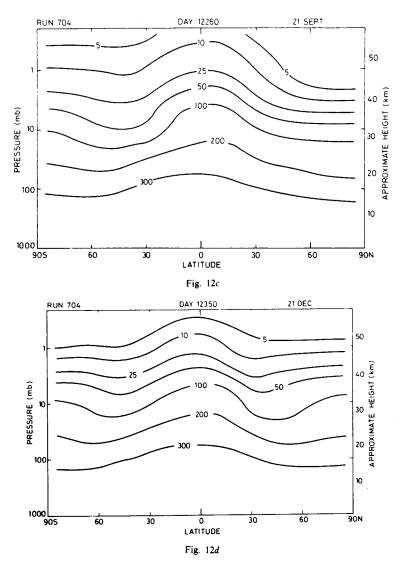
Agreement is least satisfactory between 20°N and 40°N (Figure 6) where the SAMS observations excede the model results throughout the middle and upper stratosphere. At these latitudes the model shows a minimum, compared with higher and lower latitudes, throughout much of the stratosphere. While this feature, as stated above, appears possibly to be present in some of the observations it cannot be identified unambiguously and is clearly not so pronounced as in the model. Data from other years need to be analyzed before it can be said that this represents a serious discrepancy between the model and data. In the model, the minimum is associated with downward transport, which shows only slow latitudinal variation with time. It is possible that, on the average, down-

ward motion in the atmosphere covers a greater latitudinal extent, thus producing a less strong minimum.

For  $N_2O$  the agreement is less satisfactory than in the case of CH<sub>4</sub>. At the equator (Figure 9), where for CH<sub>4</sub> there is excellent agreement between model and observations, the decrease of the mixing ratio with height is significantly underestimated in the model. On the other hand, the comparison at  $47^\circ N$  (Figure 10) is somewhat more satisfactory.

The variability of the data makes comparison with an averaged model difficult, even when the annual range is plotted. By choosing specific latitudes it is possible that agreement will appear more satisfactory or disagreements more serious than is the case. A comparison of the two-dimensional distributions should provide a good complementary test, and this has been attempted in Figures 13 and 14 for the month of February.

For  $CH_4$  (Figure 13) there are many satisfactory areas of agreement. The data exhibit a low latitude maximum, dis-



placed toward the summer hemisphere compared with the model. The slopes of the contours in middle latitude of the model northern hemisphere compare well with the observations, and there is a suggestion in the data of a mid-latitude minimum in the midstratosphere, although this is not a strong feature of the model in this month. Around the stratopause the data have less strong latitudinal gradients, suggesting perhaps a stronger horizontal transport than found in the model. On the whole, the agreement is good.

Although the basic distributions are similar, there are more areas of discrepancy for  $N_2O$  (Figure 14). For example, the observed mixing ratios decrease more rapidly with height than the model would indicate, although a comparison at about 25°S would suggest reasonably good agreement up to 1 mbar. The data show a minimum in low latitudes in the upper stratosphere and the double peak discussed earlier. As with CH<sub>4</sub>, the horizontal gradients in the upper stratosphere are weaker in the observations than those found in the model.

While the gross features of the observations are satisfactorily reproduced by the model, there are important differences of detail. The variability exhibited by the satellite data emphasizes the difficulty of comparing averaged models with observations. In this regard, notice the sharp horizontal gradients which sometimes exist. For example, at about  $30^{\circ}$ N where many observations are made there are often large gradients in CH<sub>4</sub> in the middle statosphere (see Figure 4*i*). Furthermore, this is the latitude at which many one-dimensional models are run. These variations indicate the problem of using a limited number of vertical profiles to infer coefficients for vertical transfer.

In one-dimensional models the problem of satisfactorily modeling  $CH_4$  and  $N_2O$  (and, for example,  $CFCl_3$ ) simultaneously with a single eddy diffusion profile has been apparent for some time. It has been suggested that a major source of the discrepancy for  $N_2O$ , which tends to be over estimated in the upper stratosphere when reasonable  $CH_4$  mixing ratios are reproduced, may be an underestimation of the photochemical sink. Froidevaux and Yung [1982] have recently considered model sensitivity to the absorption cross sections of molecular oxygen around 200 nm. By using average data in agreement

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#### JONES AND PYLE: SAMS OBSERVATIONS VERSUS MODEL CH4 AND N2O

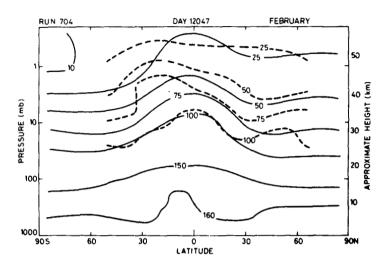


Fig. 13. A comparison of model (solid line) and SAMS (dashed line) CH<sub>4</sub> cross sections for February 1979.

with the observations of *Herman and Mentall* [1982] and more in line with those suggested by *Brewer and Wilson* [1965], they were able to improve their model representation of various trace gases.

We have performed a similar study within the twodimensional model. Between 196 and 225 nm the oxygen absorption cross sections were reduced by a factor of 0.55. This gives values a little lower than those calculated by *Brewer and Wilson* [1965]. Since our primary aim is to investigate the model sensitivity, the use of this constant factor is justifiable, particularly as the value of the cross sections in still the subject of debate.

Figure 15 shows equatorial profiles for two model runs, run A and a run including the modified oxygen cross section, run B. Both profiles are for the northern hemisphere winter solstice. Also plotted are SAMS data. To emphasize again the variability shown by the data, equatorial profiles for the December and January means are plotted. It is clear that the calculations in run A excede the observations in the upper stratosphere. The discrepancy is reduced in run B where the agreement with the January observed profile is good. On the other hand, the December observations are still exceeded by the model.

Changes in the oxygen absorption cross sections do not simply influence  $N_2O$ . The varying penetration of ultraviolet around 200 nm influences many other species, as well as affecting the radiation balance. The impact of all these processes needs to be considered. Such detailed investigation is beyond the scope of this paper. However, to summarize the comparison of runs A and B, it appears that changing the  $O_2$  absorption cross section has both a positive and negative influence. For example, the ozone distribution in run B is less satisfac-

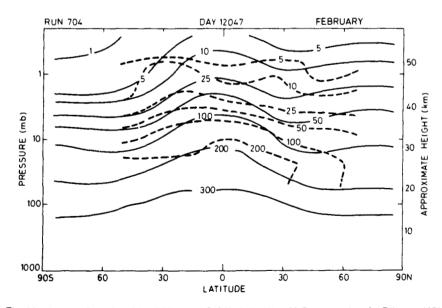


Fig. 14. A comparison of model (solid line) and SAMS (dashed line) N<sub>2</sub>O cross sections for February 1979.

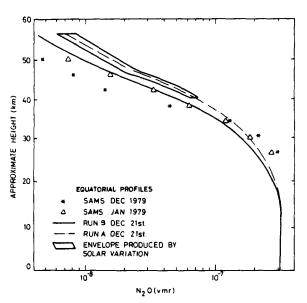


Fig. 15. SAMS equatorial N<sub>2</sub>O profiles compared with various model runs.

tory. The total column amounts are increased by 20-30 Dobson Units, depending on the latitude, leading to a significantly overestimated ozone column in low latitudes. The changes in the ozone column amount arise from increases in ozone in the low stratosphere which more than compensates for decreases in the upper stratosphere. These changes will also affect the penetration of ultraviolet radiation. In fact, the changes in ultraviolet heating lead to a somewhat improved model temperature and zonal wind fields. Following the changes in ozone, the temperatures are reduced in the upper stratosphere and increased in the lower stratosphere. Latitudinal variations in these changes lead to reductions in the stratospheric jets which tend to be overestimated in the model.

Since the N<sub>2</sub>O profile is sensitive to the penetration of radiation around 200 nm, any solar cycle variations at these wavelengths may be significant. To study to what extent the solar cycle variations, or simply uncertainty in the solar flux, might affect the comparison of model and observations we have imposed on the model the solar variability suggested by *Brasseur* and Simon [1981] for the 11 year sunspot cycle, and these results are also plotted in Figure 15, the envelope indicating the ensuing variations in N<sub>2</sub>O.

In reasonable agreement with Brasseur and Simon we find a variation of  $N_2O$  during the solar cycle of more than 30% in the upper stratosphere. The decrease in  $N_2O$  when the flux is maximum does not appear to be sufficient to explain the observed profiles.

In some respects, model experiments including the lower oxygen cross sections and solar flux variations can lead to an improvement in modeled  $N_2O$ ; they clearly indicate the necessity of establishing these parameters accurately before  $N_2O$  can be modeled with confidence. Nevertheless, this should not be taken to mean that the discrepancy is necessarily photochemical in origin. The lower chemical activity of methane results in much weaker gradients than for  $N_2O$ . Its distribution is therefore much less sensitive to transport processes. The modeled methane will be more robust to the representation of transport (just as an inert, uniformly mixed gas with no net sources would be completely insensitive to transport). In consequence, while the overestimate of the  $N_2O$  mixing ratio in the model could plausibly be ascribed to an underestimation of the sink, we believe that this is yet to be conclusively demonstrated.

#### 9. DISCUSSION

The year of results presented here provide an enormous increase in our knowledge of the distributions of  $CH_4$  and  $N_2O$  and will clearly be the subject of much detailed analysis in the future. For the present we comment on just a small number of the more obvious features that require explanation. These are the maxima found in low latitudes, the double peak structure found in the midstratosphere during a few months of the data and the uniform region sometimes found in mid-latitudes in the upper stratosphere.

The low latitude peak we have attributed to the stratospheric extension of the Hadley circulation. It is interesting to ask what velocities are implied by the changes in the monthly mean mixing ratios. For example, if we compare the methane cross sections for January and February, the 1 ppmv contour of CH<sub>4</sub> at the equator, situated near 10 mbar, ascends by approximately 0.5 pressure scale heights between January and February. This corresponds to a vertical velocity of about 1.4 mm s<sup>-1</sup> which is comparable with that calculated in the twodimensional model. Of course, this can only be regarded as a very rough estimate since horizontal motion and photochemistry, both of which would be important, have been ignored. We believe that the data are sufficiently good that we will be able to use them to trace stratospheric motions, and this work is already underway.

The double peak structure is most strongly evident in the upper stratosphere between February and May. It does not appear to be present during the southern hemisphere spring. It is not clear yet whether it would be present in a climatology of  $CH_4$  and  $N_2O$ ; analysis of further data is clearly required. Although the model does not produce a double peak, we believe that some insight into its development can be gained with the help of the model. Figure 16 overlays the model  $N_2O$ cross sections for January and July. While the peak mixing ratios are always found in equatorial latitudes, it is evident that there are significant seasonal variations at around 30° latitude. The increase in January at 30°S is consistent with the fluxes marked by the arrows. Such a flux could well give rise to the maximum seen in the data which characteristically tilt with height. Of course the model results are also consistent with a predominantly horizontal flow in the middle atmosphere. The model mean circulation, however, would not support this interpretation. If we accept the view shown in Figure 16 then the double peak will arise when the sun crosses the equator, and the upward motion is predominantly in the northern hemisphere.

It is unclear why there should be a minimum at the equator or, alternatively, why the southern hemisphere ascending branch should preserve its identity. It would appear from the data that if this is a plausible explanation, the stratosphere is more stable in the horizontal than the model indicates.

At these latitudes the rising motion in the model is driven by radiative heating which varies only slowly with latitude. It is likely that the lower stratosphere is much more variable than in the model, and, consequently, the radiative heating and associated motions will also show greater structure which might be needed to explain the observations.

5277

5278

JONES AND PYLE: SAMS OBSERVATIONS VERSUS MODEL CH4 AND NO

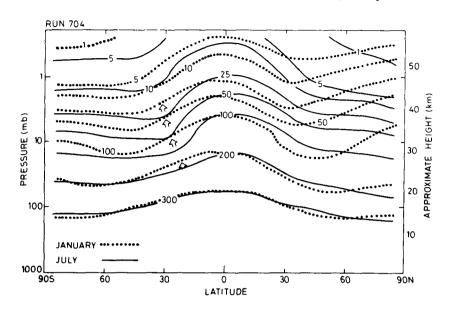


Fig. 16. Model cross-section of N<sub>2</sub>O mixing ratio (ppbv) for January (dotted line) and July (solid line). Tentative fluxes are shown by the arrows.

It is interesting to note another feature in Figure 16 commented on above. In the southern hemisphere middle and high latitude mid-stratosphere the contours of constant mixing ratio are at a greater height in January than in July. This is due to the predominantly rising motion found throughout the summer stratosphere transporting high mixing ratios upward. However, at about 1 mbar the pattern changes with the summer hemisphere having less  $N_2O$ , a reflection of the fact that the photochemical time constant is decreasing with height.

A number of the observed cross sections show a region in the upper stratosphere of the summer hemisphere across which there is a strong horizontal gradient, while in the winter hemisphere the horizontal gradient is small. The area of marked horizontal gradient is also the region, commented on earlier, where the vertical gradient becomes very small (see, e.g., Figures 3h and 4h at  $\sim 40^{\circ}$ N and 2 mbar).

The behavior involves both transport and photochemistry. We identify the area of very small vertical gradient as a transition region. Equatorward of this the high mixing ratios are produced by strong upward transport confined to low latitudes. The horizontal flow in the upper stratosphere is from the summer to the winter hemisphere. The weak horizontal gradients in the winter hemisphere are thus produced by strong horizontal transport from low latitudes. Poleward of the transition region the air has low mixing ratios of N<sub>2</sub>O and CH<sub>4</sub>. This can be thought of as old stratospheric air, having been at high altitudes for sufficient time for the slow photochemistry of N<sub>2</sub>O and CH<sub>4</sub> to reduced the mixing ratios. This air is transport equatorward. The transition region arises when the vertical transport becomes of comparable importance to the horizontal.

The model produces these features, although less markedly than does the data. This was discussed in section 8.2. Figure 17 shows profiles of N<sub>2</sub>O at 47°N for February and October. In October, representative of late summer/early autumn, the features found in the data are reproduced, namely, a sharpening of the gradient in the low stratosphere and a much weaker vertical gradient found in the upper stratosphere.

#### 10. CONCLUSIONS

In this paper, monthly mean, zonal mean mixing ratios of  $CH_4$  and  $N_2O$  for 1979 obtained from the SAMS instrument on NIMBUS 7 have been presented for the first time, with a coverage from 67.5°N to 50°S and including most of the stratosphere. These data represent a vast increase in the knowledge of the behavior of the two gases in the upper atmosphere.

The data have been compared with the small number of available measurements. At low altitudes, there appears to be a consistent difference between the two data sets with the satellite data higher by  $\sim 10-40\%$ . Nevertheless, the two sets of data are consistent within their estimated errors.

The SAMS data confirm the broad features revealed by in situ and other remote measurements. In particular, the decrease of the mixing ratios with height is more rapid in middle latitudes then in the tropics. Many new features are revealed by the satellite data. The large variability from month to month and with latitude is of interest and emphasizes the problems of comparing theoretical models with single observed profiles for a particular day.

For any pressure level in the stratosphere, maximum mixing ratios, as we have stated, are found in low latitudes and are associated with upward transport. Between February and May 1979 a double peak is found with maxima in the sub tropics in the middle stratosphere and a local minimum at the equator. This feature is not found in the model, but the seasonal change in the model fields does suggest a somewhat similar behavior with upward transport following the sun. In the model, horizontal transport prevents the occurrence of an equatorial minimum. It will be of great interest to see if the double peak is found in subsequent years. These data are yet to be analyzed. A more detailed investigation of the double

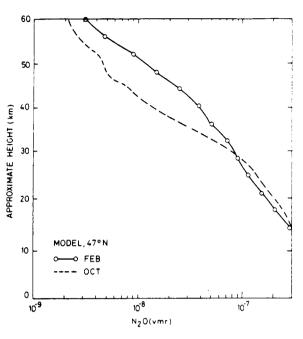


Fig. 17. Model N<sub>2</sub>O profile at 47°N for February and October.

peak in which the velocity fields are diagnosed is also important.

Another feature discovered in the data is a region of very strong horizontal gradients, and weak vertical gradients in the mixing ratios, found in the upper stratosphere summer hemisphere. A similar, if less pronounced feature, is found in the model. This transition region marks a loose boundary between air of high mixing ratio, which is being transported upward and into the winter hemisphere, and low mixing ratio air which has been in the stratosphere for a sufficiently long time for the slow photochemistry to produce low mixing ratios of CH<sub>4</sub> and N<sub>2</sub>O.

Comparison of the data with the model has proved a very useful exercise. The broad features agree well, suggesting that the gross behavior of the model is satisfactory. One feature found in the model is a middle latitude minimum. Although suggestions of this are found in the data, the coverage is not sufficient for an unambiguous identification. As discussed above, investigation of the model has helped in understanding the double peak and transition regions found in certain months.

The quantitative comparison for  $CH_4$  is extremely good; the model agrees with the data very well. For  $N_2O$  the agreement is not so satisfactory with the model results tending to exceed the observed values in the upper stratosphere. Such a disagreement has been found in other models and for other species (e.g., CFCl<sub>3</sub>). It could be suggested that because the agreement of CH<sub>4</sub> is generally good, the model's failure to reproduce the N<sub>2</sub>O profile may be due to an underestimation of the stratosphere sink. We have considered two possibilities, a reduced oxygen absorption cross section around 200 nm and solar flux variations. While both cases can help to reduce the discrepancy, the results are not conclusive. Since the CH<sub>4</sub> gradients are weak, it is less sensitive to transport than N<sub>2</sub>O. We believe that inaccuracies in the representation of the transport cannot be ruled out as a cause of the discrepancy.

Finally, we stress again the advantages provided by this large data set. Although the absolute accuracy is not high compared with some in situ measurements, this disadvantage is more than offset by the spatial and temporal coverage. We believe much useful analysis on the interaction of transport and photochemistry is possible by using the data.

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# SOME ASPECTS OF THE DESIGN AND BEHAVIOR OF THE STRATOSPHERIC AND MESOSPHERIC SOUNDER

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# Some Aspects of the Design and Behavior of the Stratospheric and Mesospheric Sounder

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The measuring task, the engineering constraints, and the resulting design of the stratospheric and mesospheric sounder are described in broad terms. The calibration system and its performance are discussed in more detail. One significant side effect in the design and its impact on the fields of view is also discussed. The major lessons for future pressure modulator radiometer designs are summarized.

#### 1. INTRODUCTION

The rapid development of remote sounding techniques has meant that new instruments are rarely simple developments of older designs. Radiometers such as the stratospheric and mesospheric sounder (SAMS) consist of optical, electronic, mechanical, and thermal systems that are related to one another in complex ways. Consequently, the designer who, after having optimized the various trade-offs to produce the best overall performance, does not find when the instrument has been built that his decisions have led to unexpected side effects may count himself as being very fortunate. Particularly troublesome are those side effects whose magnitudes may be different in orbit (e.g., owing to the change to a zero g environment). Such effects cannot be characterized satisfactorily in the laboratory, and ideally they should be eliminated by redesign, but this may not be possible because of constraints of finance or of schedule. We feel that a discussion of some of these aspects of the SAMS radiometer may be of interest to geophysicists engaged on the analysis of SAMS data as well as to others involved with the design of instruments. This paper should be read in conjunction with the description of the SAMS instrument given by Drummond et al. [1980].

#### 2. BASIC DESIGN AIMS AND THEIR REALIZATION

The basic task of the SAMS radiometer is to measure the thermal emission from chosen molecular species present in atmospheric limb paths whose tangent point altitudes range upward from 20 km. Since the radiometer is itself a significant source of radiation, it is necessary to label the incoming radiation. In some of the channels in SAMS this function is combined with the spectral selection of the part of the incoming radiation which has been emitted by the desired constituent.

The problem of spectral selection has two main aspects: the separation of the wanted lines from other atmospheric emission and the minimizing of the contribution due to components of the instrument's own optical system. (Variations in the instrumental emission can often decide the minimum detectable radiance.) Both objectives would be achieved if a spectral response profile limited to just those very narrow regions embracing the wanted lines could be realized.

A good approximation to the ideal spectral response is obtained in SAMS by employing multilayer dielectric filters to

Paper number 3D1831.

5287

isolate the desired parts of the chosen emission bands, combined with pressure modulators [*Curtis et al.*, 1974] which provide the necessary detailed spectral selectivity. A pressure modulator is an optical absorption cell which is filled with the gas whose emission is to be measured. The pressure of the gas is modulated at a few tens of Hertz, resulting in cyclic transmission variations which are confined to spectral regions within the absorption lines of the molecule in question. The desired output from the detector is extracted by "lock-in" type electronic signal processing referenced to the frequency and phase of the pressure cycling. In the channels where pressure modulation is not used (the so-called wideband channels) a vibrating black vane chopper is used to label the incoming radiation, and lock-in techniques are again employed.

As the tangent height of the atmospheric limb path is increased, the signal observed in a radiometer channel falls toward zero and is eventually lost in instrumental noise. In order to sound the widest possible altitude range, therefore, it is necessary to use the best available detectors and to maximize the energy gathered by the telescope by employing the largest practical aperture and field of view.

The choice of detectors for SAMS was constrained by the limited cooling technology then available for long missions, and only two of them were cooled. These were an indium antimonide photovoltaic detector covering the 4 to 5-µm region, cooled by a two-stage radiation cooler, and a lead sulphide photoconductor for 2.5-2.7  $\mu$ m, cooled by a simpler radiator. The remaining detectors were triglycine sulphate (TGS) pyroelectric devices operated at room temperature. These offered the highest radiometric sensitivity which could be achieved without cooling, but they also showed unwelcome microphonic behavior, which will be discussed later. The costs of increasing the aperture of the telescope become disproportionately high once a mirror diameter of about 20 cm is exceeded, bearing in mind that the signal to noise ratios obtained are proportional to the diameter (not the area) of the collecting optics when, as in SAMS, the detector area is matched to the collecting area. The field of view in the vertical direction is determined largely by the requirements of vertical resolution and is usually of the same order as the natural height range of a limb path (about 4 km), while for a number of reasons the horizontal width of the field of view is best limited to at most 10 times the vertical height. The telescope in SAMS had an aperture of 177 cm<sup>2</sup> and a field of view of  $0.16^{\circ} \times 1.6^{\circ}$  (nominally 10 km  $\times$  100 km at the tangent point).

Since the objective of a limb emission sensor such as SAMS is to measure radiances that range down to those that are barely distinguishable from zero radiance (space view) it is

PAGE 148 INTENTIONALER BEAM

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5288

essential to measure the atmospheric radiances as differences from the zero radiance level, so that low frequency noise (drifts) in the instrument response can be eliminated. Difference measurements could in principle be made either by rotating the whole instrument between the chosen limb path orientation and a clear view to space or by using an additional tilting mirror. The former approach would require a high precision mechanism capable of rotating the whole instrument sufficiently rapidly to provide adequate coverage with minimal loss of observing time, a difficult engineering requirement. In addition, the associated angular momentum changes would have had a significant effect on the roll attitude of the spacecraft, to the detriment of other limb viewing instruments. Consequently, a scanning mirror system was chosen, in spite of its disadvantages, which are that the optical system must contain an additional component and that the system is unavoidably changed as the mirror is tilted. The mirror tilt required to achieve a clear space view is, however, only of the order of 2°, and, as will be shown later, this does not cause significant changes in instrument response.

Standard calibration radiances for the various channels are provided by introducing a black-body source into the optical path. The most straightforward way of achieving this would be to move a large black body cavity into place across the viewing aperture. This approach would, however, have given rise to difficulties in achieving adequate temperature stability and uniformity, and, indeed, accommodation of such a system on the spacecraft would not have been at all straightforward. An alternative scheme of introducing a small black body at the first focal point of the telescope was therefore chosen. The implications of this are discussed in section 4.

#### 3. ZERO RADIANCE (SPACE VIEW) SIGNAL

There are a number of contributions to the zero-radiance signal. The main ones are listed below.

1. Emission from the optical components situated in front of the modulators makes a contribution which depends on the temperatures of the components and their emissivities.

2. Radiation enters the beam in front of the modulators as a result of scattering from imperfections in the optical surfaces and through optical aberrations and diffraction.

3. When the pressure modulator (or chopper) in any given SAMS channel is at the point in its modulation cycle where the transmission is low (or zero), it acts as a room temperature radiation source of high emissivity. Consequently, the modulated (or chopped) radiance at the detector has the largest amplitude when the incoming atmospheric radiance is zero, and this amplitude will change as the temperature of the modulator (chopper) changes.

4. The pressure oscillation in the pressure modulator cell is unavoidably accompanied by an oscillation in the gas temperature. These temperature variations lead to an oscillating radiance at the detector which is synchronous with the pressure modulation, yielding an offset in the signal channel output which is dependent on the modulation depth, on the mean pressure in the modulator, and on the temperature of the cell.

Finally, a system employing lock-in type signal processing will respond to any synchronous input, including those not related to the modulated radiance. For example, vibrations in the instrument structure excited by the mechanical modulators may be picked up by microphonic detectors. The TGS detectors in SAMS have a large microphonic response and significant offsets due to this occur in both the pressure modulated and wideband channels. Another source of offsets of this type is electrical pick-up of phase sensitive detector (PSD) reference or modulator drive waveforms.

The difference method of measuring small signals discussed in section 2 relies on all the contributions to the zero radiance signal mentioned above (and also the gain) remaining constant when the scan mirror is moved between space view and the atmospheric limb view. (In addition, the contributions must drift sufficiently slowly that the zero radiance signal can be estimated with the necessary accuracy between space views.)

In general, the instrument is designed such that the temperature changes that affect the instrument response occur only slowly; similarly, the signal offsets due to temperature cycling in the modulator cell and electrical synchronous pickup would also be expected to be stable in the short term. Our attention should therefore be focused on the stray radiation reaching the detector as a result of optical system imperfections and on the microphonic component of the synchronous pickup. Both of these effects may vary as the scan mirror is tilted, the variation in the microphonic case being due to changes in mechanical couplings between the various parts of the instrument. The variations likely to be encountered in orbit are not easily assessed in the laboratory before launch, because it is difficult to arrange for the radiation environment of the instrument to be sufficiently well simulated and because the mechanical couplings responsible for the microphonic component may well change when the sensor is in zero-g. A check obtained in orbit that the zero-radiance offset does not depend on scan mirror position is therefore of special value.

Ideally, we should like to be able to tilt the spacecraft about its roll axis, so that the instrument has a clear view to space throughout the entire scan range: we could then investigate the variation in offset with scan angle directly. It has unfortunately not been possible to do this with NIMBUS 7, but a clear view to space is nonetheless available over a scan range of about 2.5° in the line of sight.

Figure 1 shows the variation in space view offset with scan mirror position in channel A1 WB (a temperature sounding wideband channel operating near 15  $\mu$ m in the v<sub>2</sub> band of CO<sub>2</sub>) during a typical orbit. The data were obtained as follows. The instrument was programed to perform a uniform up and down scan of 2.9° amplitude in the nominal space view region of the scan range, with occasional excursions to view the middle atmosphere so that a check on the attitude of the spacecraft could be obtained by using the CO<sub>2</sub> channel radiances [see Barnett et al. this issue]. The data are presented in Figure 1 in terms of the tangent height of the atmospheric path viewed at any given time, determined from these attitude checks. The dependence of the observed radiance on scan angle is very weak, typically less than 0.1% of the radiance, which would be received from a 290 K black body (or about 0.2% of the radiance received from an atmospheric path of low transmission and stratospheric temperatures (240 K)). It should be noted that this is much less than the noise level in a normal measurement sequence.

#### 4. GAIN CALIBRATION

As mentioned earlier, gain calibration is achieved in SAMS by introducing a small black-body source at the first focal point of the telescope. The temperature of this source is very close to the temperature of the secondary optics components which accept its emitted radiation. Consequently, the emiss-

## WALE AND PESKETT: SAMS DESIGN AND BEHAVIOR

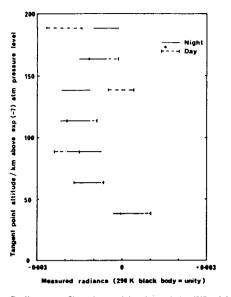


Fig. 1. Radiance profiles observed in channel A1 WB (CO<sub>2</sub>, 15  $\mu$ m band) in the space view region of the scan range. The data shown are taken from the portion between latitudes 50°N and 90°N of a single orbit of NIMBUS 7. The bars indicate 1 $\sigma$  error. "Day" and "night" indicate the ascending and descending parts of the orbit, respectively.

ivity of the source does not have to be very high, since any radiation which is reflected by it is of the same temperature as the radiation emitted by the source itself. *Wale* [1981] has shown that the gain error resulting from the assumption that the SAMS internal black body is of unity emissivity should be less than 0.1% under all normal operating conditions.

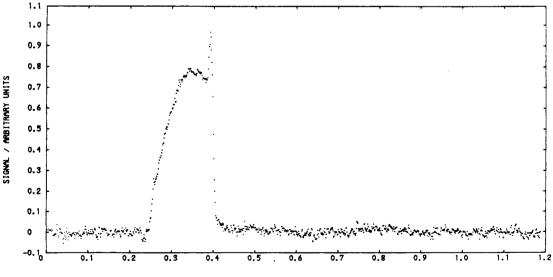
The main drawback to the use of a calibration black body that is not at the front of the whole optical system is that corrections need to be made for emission from the scan mirror and objective mirror, which are excluded from the system when the black body is in the beam. It should be noted, however, that if the mirror temperatures are close to the black body temperature the corrections are small, tending to zero in the case of an isothermal instrument. In SAMS the error in the correction may be as large as 30% before the error in the temperature measured by the  $15-\mu$ m band carbon dioxide channels exceeds 0.1 K.

It is possible that the level of microphonic synchronous pickup is dependent upon whether the black body is in or out of the beam. The black-body mechanism is a small one, and the only change in configuration involved in moving the source into the beam is slight, but other small perturbations (e.g., the opening of the solar door on the ERB instrument) did have an observable effect in ground test. Since, however, we have no way of quantifying the effect, we have no choice but to ignore it. The ultimate test is comparison with results from other instruments [see Barnett and Corney, this issue].

# 5. MOTIONAL CHOPPING

One of the sets of measurements carried out during ground testing was designed to determine the shape of the field of view of each channel. This was achieved by rotating the whole instrument in such a way that the line of sight elevated slowly through a nearly parallel beam of radiation from a collimator.

The measurements revealed unexpected features: The profiles were found to be asymmetric, and they often showed spikes at the edges. A typical profile is shown in Figure 2. The general features are similar in all of the channels, both wide band and pressure modulated, although the positions of the negative and positive peaks are reversed in some cases. The spikes arise from vibrations at the edge of the scanning or telescope mirrors at modulator or chopper frequency, resulting in unwanted synchronous scanning about the roll axis. Estimates of the effective reaction mass of the instrument and its mountings and a knowledge of the uncompensated momentum of the pressure modulators and chopper indicate that motions of the required amplitude (of the order of 1  $\mu$ m) are to



ANGULAR POSITION OF LINE OF SIGHT / DEGREES

Fig. 2. Field of view response profile of channel C1 PMR, measured during ground testing.

5290

be expected. The question arises as to what are the "real" fields of view.

Consider the situation shown in Figure 3, where a source of very limited angular extent (the collimator beam) is filling the aperture of the sensor. If the field of view is described by the function  $F(\theta)$ , a small movement  $\Delta \theta$  in the angular position of the source will result in a change in the radiance of  $(d/d\theta)F(\theta)\Delta\theta$ . If  $\Delta\theta$  takes the form of an oscillatory function synchronous with the modulation in the channel under consideration, an additional signal proportional to  $(d/d\theta)F(\theta)$  will be observed at the output of the channel. We call this effect motional chopping of the second kind, to distinguish it from amplitude modulation of a beam resulting from unwanted vibration of a nonuniform optical component near an aperture plane, long known to us as motional chopping.

It should be noted that in a pressure-modulated channel, the motionally chopped contribution is a wide band one derived from the whole channel bandwidth defined by the band pass filter, not just the pressure modulated fraction of that bandwidth. It is for this reason that the effect can be significant, even though the motions involved are very small. In SAMS an enhanced effect is also present in the wide band channels, since only a small part of the optical system aperture is modulated by the vibrating chopper, but the whole aperture contributes to the motionally chopped signal.

We describe the signals obtained as follows. In the absence of motional chopping the signal consists of a single term

$$S(\theta') = \alpha \left\{ \int L_1(\theta) h_1(\theta - \theta') \ d\theta \right\}$$

where  $L_1(\theta)$  describes the angular distribution of the source radiance which is selected spectrally by the channel,  $\theta$  is an elevation angle referred to the instrument mountings,  $h_1(\theta)$  $-\theta'$ ) is the field of view profile when the line of sight is directed at  $\theta'$  in this coordinate system, and  $\alpha$  is a constant. If the line of sight is now allowed to move at the modulation frequency,

$$\theta' = \theta'_0 + \Delta \theta' \sin(\omega t + \phi),$$

θ we obtain a second term, of the form

$$\alpha \Delta \theta' \cos \phi \cdot \frac{d}{d\theta'} \int L_2(\theta) h_2(\theta - \theta') d\theta$$

It should be noted that  $L_2$  and  $h_2$  are new functions which may differ from  $L_1$  and  $h_1$ . For example, in the SAMS

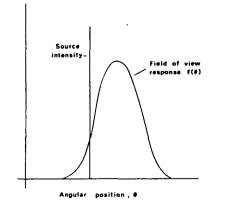


Fig. 3. Schematic illustration of field of view measurement.

pressure-modulated channels,  $L_1$  is the pressure modulated (narrow band) radiance, while  $L_2$  is a wide band radiance. We would, however, expect  $h_1$  to be the same as  $h_2$  in this case. In the wide band channels,  $h_1$  is the field of view resulting from the modulation of the small part of the aperture occupied by the chopper blade, while  $h_2$  is the field of view resulting from the use of the whole aperture. The main difference between  $h_1$ and  $h_2$  is a scaling factor, although there are minor differences in the shape of the profile caused by the off-axis siting of the chopper.  $L_1$  and  $L_2$  are in this case essentially the same function. but the off-axis nature of the wide band chopping has an effect on the spectral selection of the source radiance, owing to angle of incidence effects on dielectric interference filter coatings.

In the laboratory we approximate the shape of  $L_1$  and  $L_2$ by a delta function (the collimated beam); that is,  $L_1 = k_1 \delta(\theta)$ and  $L_2 = k_2 \delta(\theta)$ , where  $k_1$  and  $k_2$  are constants. In this case the signal observed is

$$S_{L}(\theta') = \alpha k_{1} \cdot h_{1}(\theta') + \alpha \Delta \theta' \cos \phi \cdot k_{2} \frac{d}{d\theta'} h_{2}(\theta')$$

When observing the atmosphere, the signal obtained is

$$S_{A}(\theta') = \alpha \int L_{1}^{A}(\theta)h_{1}(\theta - \theta') d\theta + \alpha \Delta \theta' \cos \phi \int L_{2}^{A}(\theta) \frac{d}{d\theta'} h_{2}(\theta - \theta') d\theta$$

where we have used the fact that  $L_2(\theta)$  does not depend on  $\theta'$ .

Considering first the wide band channels and assuming that the differences between  $L_1$  and  $L_2$  (and, hence, between  $k_1$  and  $k_2$ ) can be neglected,

$$S_{L}(\theta') = \alpha k_{1} \{ h_{1}(\theta') + \Delta \theta' \cos \phi \, \frac{d}{d\theta'} \, h_{2}(\theta') \}$$

and

$$S_{A}(\theta') = \alpha \int L_{1}^{A}(\theta) \left\{ h_{1}(\theta - \theta') + \Delta \theta' \cos \phi \cdot \frac{d}{d\theta'} h_{2}(\theta - \theta') \right\} d\theta$$

Combining the two equations,

$$S_{\mathcal{A}}(\theta') = \frac{1}{k_1} \int L_1^{\mathcal{A}}(\theta) S_L(\theta - \theta') \ d\theta$$

We conclude that under the assumption that the effect is the same in zero-g, the motional chopping is taken into account simply by employing the field of view profile in the raw form measured in the laboratory (i.e., including the spikes). The consequences of this assumption being incorrect have been studied by Barnett and Corney [this issue].

The treatment of motional chopping in the pressuremodulated channels is not so straightforward. The reader will perhaps have realized that the conclusions in the wide band case do not in fact depend on the precise functional form of the additional term so long as it depends only on  $\theta'$ . In the pressure-modulated case, however, we must establish the correctness of the model in order to proceed. The first step is to derive the underlying pressure-modulated field of view profile from the laboratory measurements. Equation (1) may be integrated numerically to obtain an estimate of  $h_1(\theta)$  for trial values of the constant  $\Delta \theta' \cos \phi \cdot k_2/k_1$ . The result of such an integration using a value for  $\Delta\theta \cos \phi \cdot k_2/k_1$  of  $-0.01^\circ$  is shown in Figure 4. It should be noted that the integration has had a useful smoothing effect, reducing the noise level. Also

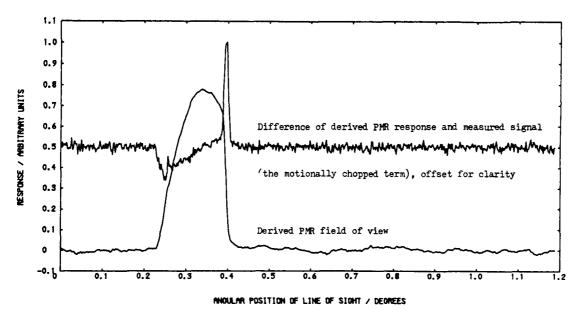


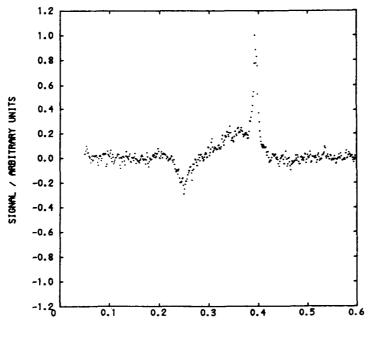
Fig. 4. Underlying pressure modulated field of view in channel C1 PMR, derived from the profile shown in Figure 2, using a value for the constant  $\Delta\theta \cos k_2/k_1$  of  $-0.01^\circ$ .

shown is the difference between the derived profile and the raw measurement.

sensor. The amount of gas in the path was sufficiently large for the pressure-modulated radiance  $L_1$  to be strongly attenuated, but  $L_2$  was scarcely affected.

Compelling evidence of the correctness of the model may be adduced by comparing the difference profile with the profile shown in Figure 5. The latter was obtained with an absorbing path of gas interposed between the collimator source and the

When observing the atmosphere, the ratio of the magnitudes of  $L_1$  and  $L_2$  in the pressure-modulated channels is much smaller than the ratio of  $k_1$  and  $k_2$  in the laboratory.



ANGULAR POSITION OF LINE OF SIGHT / DEGREES

Fig. 5. Field of view profile observed in channel C1 PMR with a path of carbon dioxide interposed between the source and the SAMS instrument.

5292

#### WALE AND PESKETT: SAMS DESIGN AND BEHAVIOR

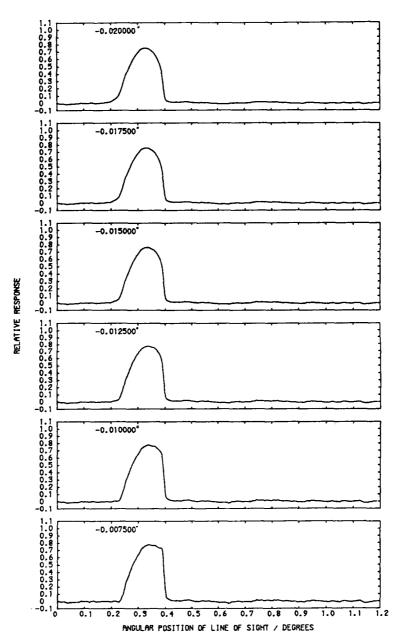


Fig. 6. Field of view profiles for channel C1 PMR derived by using different values of the constant  $\Delta\theta \cos k_2/k_1$ .

This is because a spectrally uniform source (a hot filament) was used in the latter case, and the motionally chopped term contains in consequence a large contribution from spectral regions between the absorption lines of the gas. In the atmospheric case, however, the emission lines are usually narrow and well separated, and the contribution from spectral regions between the wanted lines is much smaller than in the laboratory case. The best course in the pressure-modulated case is to ignore the motional chopping term and to use a derived field of view profile such as that shown in Figure 4.

The criteria for choosing the optimum value of the constant  $\Delta\theta \cos \phi \cdot k_2/k_1$  are less clear than one might have hoped. A

typical set of trial derivations is shown in Figure 6. The bottom profile may be rejected on the basis that spikes are still visible and the top two look suspiciously oversmoothed, but the remaining profiles are all plausible. The uncertainty in the value of the constant leads to an uncertainty of about  $\pm 0.003^{\circ}$  in the position of the center of the field of view in this case.

# 6. CONCLUDING REMARKS

We believe that the basic approaches employed in the design of SAMS were sound, but two changes of detail would be desirable if another instrument were to be built within the

#### WALE AND PESKETT: SAMS DESIGN AND BEHAVIOR

same constraints. First, the linear momentum of the pressuremodulator pistons should be compensated in some way to eliminate the motional chopping (of the second kind) which results from vibrations excited in the instrument structure. Second, in the wide band channels, the use of a balanced rotating chopper instead of an unbalanced vibrating one would reduce the motional chopping effect to negligible levels and also eliminate the microphonic component of synchronous pickup. (This chopper should modulate the entire aperture, so that the spectral response profiles applicable to the pressure-modulated and wide band channels are identical.)

New spacecraft such as the Upper Atmospheric Research Satellite (UARS) offer a relaxation in the power and weight constraints, and closed cycle detector coolers suitable for use on long missions (e.g., 2 years) are now available. The much lower noise levels and absence of microphony offered by quantum detectors operating at medium temperatures (70–80 K) could improve the performance achieved in SAMS by 2 orders of magnitude. Such an enormous improvement, welcome though it is, will of course mean that all the effects mentioned in this paper and indeed all the design trade-offs will need to be re-examined very carefully if the additional signal to noise ratio is to be effectively utilized.

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<ul> <li>16. Abstract The Stratospheric and Mesospheric Sounder (SAMS) aboard Nimbus-7 observes infrared radiation from the atmospheric limb. Global upper atmosphere temperature profiles and vertical concentrations of H<sub>2</sub>O, NO, N<sub>2</sub>O, CH<sub>4</sub> and CO<sub>2</sub> are derived from these measurements. The status of all channels was carefully monitored. Temperature and composition were retrieved from the measurements by linearizing the direct equation about an a priori profile and using an optimum statistical estimator to find the most likely solution. The derived temperature and composition profiles are archived on two tape products whose file structure and record formats are described in detail. The Gridded Retrieved Temperature Tape (GRID-T) contains daily day and night average temperatures at 62 pressure levels in a 2.5° latitude by 10° longitude grid extending from 67.5°N to 50°S. The Zonal Mean Methane and Nitrous Oxide Composition Tape (ZMT-G) contains zonal mean day and night average CH<sub>4</sub> and N<sub>2</sub>O mixing ratios at 31 pressure levels for 2.5° latitude zones extending from 67.5°N to 50°S.</li> <li>17. Key Words (Suggested by Author(s)) Stratospheric and Mesospheric Sounder Temperature Methane</li> <li>18. Distribution Statement</li> </ul>						
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