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NACA TN 3809 8	NATIONAL ADVISORY COMMITTEE
	TECHNICAL NOTE 3809
	A METHOD FOR CALCULATION OF FREE-SPACE SOUND PRESSURES
	NEAR A PROPELLER IN FLIGHT INCLUDING CONSIDERATIONS
	OF THE CHORDWISE BLADE LOADING
	By Charles E. Watkins and Barbara J. Durling
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	NACA
	Washington
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NATIONAL ADVISORY COMMITTEE FOR AERONAUT.



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A METHOD FOR CALCULATION OF FREE-SPACE SOUND PRESSURES

NEAR A PROPELLER IN FLIGHT INCLUDING CONSIDERATIONS

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SUMMARY

The general expressions of NACA Report 1198 for determining the sound-pressure field for a rotating propeller in uniform subsonic flight are reviewed for the case of a propeller with uniform chordwise forces. Consideration is given to effects of nonuniform chordwise and radial blade loading. Tabulated values of certain definite integrals that are involved in the calculation of a near-field propeller noise regardless of the form of the chordwise forces are presented. These tabulations cover a wide range of operating conditions and are useful for estimating propeller noise when either the concept of an effective radius or radial distributions of forces are used in conjunction with the concept of an effective radius to make evaluations for some specific propellers. Results of these evaluations are presented and discussed.

INTRODUCTION

In references 1 and 2 the sound-pressure field of an "on-stand" propeller is analyzed by replacing the normal-pressure distributions associated with thrust and torque over the propeller by distributions of pressure doublets acting at the propeller disk. In reference 3 the analyses of references 1 and 2 are extended to the case of an in-flight propeller by considering the pressure doublets that represent the thrust and torque to be subjected to a uniform rectilinear motion.

In references 1 and 3 the equations for the sound pressure are derived in exact form, that is, within the realm of linearized potential theory. For convenience in calculation, however, these equations are ultimately simplified so that they pertain only to the first few harmonics of a propeller of low solidity. These simplifications involve the assumptions that the radial load distribution on a propeller blade can be considered as concentrated at some effective radial position and that the

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propeller blade chord is so small that the chordwise load at any radial station can be considered as having the form of an impulse or Dirac delta function. Calculations in references 1 to 3 are based on these assumptions.

The present report begins with the general equations that are derived in reference 3 for the sound field of an in-flight propeller and reviews the specialization of these equations to the case of a uniform or rectangular-type chordwise loading. It is deduced from results for this case that, as long as the chordwise loading remains unchanged throughout the propeller cycle, certain definite integrals involved in the expressions for the sound field are the same regardless of the form of the chordwise loading. In the case of the "far field," or for field points located at distances of several propeller diameters from the propeller, the definite integrals can be evaluated in terms of Bessel functions of the first kind, but, in the case of the "near field" or for field points located within a distance of two or three propeller diameters from the propeller, the integrals must be evaluated by numerical procedures. These numerical evaluations for the near-field case can be made once and for all, however, so as to apply to any chordwise loading that remains unchanged throughout the propeller cycle. A main purpose of this report is therefore to present tabulated values of these definite integrals for a range of parameters that correspond to a large range of operating conditions so that the sound field of a given propeller can be fairly rapidly estimated.

Other purposes of the report are: (a) to consider the effects on the sound field of nonuniform types of chordwise blade loading, (b) to consider a method of evaluating expressions for the sound field when radial distributions of loading are taken into account, and (c) to present and discuss results of some calculations, based on the tabulated integrals, of the sound pressures for some specific propellers. The reduction of definite integrals pertaining to the far-field case to Bessel functions of the first kind and a Fourier series development of some simple shapes that are useful for representing nonuniform chordwise force distributions are presented in the appendixes.

SYMBOLS

B number of propeller blades

b(r) width of propeller blades or blade chord

c velocity of sound

cn section normal-force coefficient

÷,

D propeller diameter

 $F(r, \theta, t)$ chordwise distribution of forces perpendicular to thrust of propeller blade and giving rise to torque

 G_1, G_2 functions defined in equations (31) and (32)

 $I_{\nu}, I_{\nu}, I_{\nu}', I_{\nu}''$ functions defined in equations (13), (15), (18), and (19)

J_n Bessel function of first kind with index n

$$K_{1} = \frac{2\beta^{2}ry}{x^{2} + \beta^{2}y^{2} + \beta^{2}r^{2}}$$

$$K_2 = \frac{\Omega}{\beta^2 c} \sqrt{x^2 + \beta^2 y^2 + \beta^2 r^2}$$

 $k = \omega/c$

М	Mach number, V/c
М _Е ,	flight Mach number
м _R	rotational Mach number
M _t	tip Mach number
m	order of harmonic
n	propeller rotational speed
ъ ^щ	pressure due to thrust
р _Q	pressure due to torque
p _{rms}	root-mean-square pressure
ନ୍	torque
R	length (radius) of propeller blades
Re	effective (radius) length of propeller blades

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R.P.	real part of any expression
r	radius to blade element
r,0	polar coordinates in yz-plane
$S = \sqrt{x^2}$	$\beta^2 + \beta^2 y^2 + \beta^2 r^2 - 2\beta^2 ry \cos \theta$
T	thrust
t	time
ť	dummy variable
V	forward velocity
$V_{\mathbf{F}}$	forward flight velocity
vt	tip velocity
x,y,z,y _l	,z ₁ Cartesian coordinates
$\beta = \sqrt{1 - 1}$	M2
^β 0.75R	blade angle at 0.75 tip radius
μ ₁ ,μ ₂	constants
η	propeller efficiency
$\sigma = \frac{Mx + \beta^2}{\beta^2}$	<u></u> .8
$\tau = b/r\Omega$	
τ 0	period, $2\pi/B\Omega$
Ω	angular velocity
ω	frequency of mth harmonic, $mB\Omega$
Subscrip	ts:
Ĵ	jth element

It left triangle

Ú,

r rectangle

rt right triangle

t triangle

 $\nu = 1, 2, \text{ and } 3$

ANALYSIS

The sound-pressure equations for chordwise rectangular loading are derived and reduced to expressions involving certain definite integrals denoted by I_{ν} ($\nu = 1, 2, \text{ and } 3$). It is deduced that the integrals I_{ν} are common to propeller-noise theory; hence, they are evaluated and tabulated over useful ranges of certain key parameters. This tabulation is discussed briefly and then consideration is given to effects of other types of chordwise loading and to a method of approximating a radial integration involved in the sound-pressure equations.

Sound-Pressure Equations for Chordwise Rectangular

Blade Loading

Although the main purpose of this report is to present tables of certain definite integrals that arise in expressions for the near sound field of a propeller, it is desirable, for the sake of completeness and discussion, to derive the integrals by specializing the general expression for the sound field to some particular form of chordwise loading. For this purpose the general equations derived in reference 3 are specialized to the case of rectangular chordwise loading. More general types of loading are discussed in a subsequent section.

From equations (6) to (12) of reference 3, the equations for the sound pressure $p_{\rm T}$ associated with a thrust distribution $F_1(t,r,\theta)$ and the sound pressure $p_{\rm Q}$ associated with a torque distribution $F_2(t,r,\theta)$, F_1 and F_2 being periodic with respect to time, can be written as (see fig. 1 for illustration of coordinate system):

$$p_{\rm T} = -\frac{1}{4\pi} \int_0^R \int_0^{2\pi} F_{\rm L}(t,r,\theta) \frac{\partial}{\partial x} \frac{e^{-ik\sigma}}{s} dr d\theta \qquad (1)$$

$$p_{Q} = \frac{1}{4\pi} \int_{0}^{R} \int_{0}^{2\pi} F_{2}(t,r,\theta) \frac{1}{r} \frac{\partial}{\partial \theta} \frac{e^{-ik\sigma}}{S} dr d\theta$$
(2)

In these equations R is radius, r is the blade element radius, t is time, θ is a polar angle in the yz-plane, $\beta = \sqrt{1 - M_F^2}$, $S = \sqrt{x^2 + \beta^2 y^2 + \beta^2 r^2 - \beta^2 r y \cos \theta}$, and $\sigma = \frac{Mx + S}{\beta^2}$; $k = \frac{\omega}{c} = \frac{mB\Omega}{c}$ is a

frequency parameter involving order of harmonic m, number of blades B, angular velocity Ω , and velocity of sound c. It is to be noted that the expressions for $p_{\rm T}$ and $p_{\rm Q}$ (eqs. (1) and (2)) have been made independent of the coordinate z. This is a convenient simplification that follows from rotational symmetry when the loading is assumed to remain unchanged throughout the propeller cycle.

When the periodic functions F_1 and F_2 are known, they can be expressed in the form of Fourier series that enables one to treat each harmonic of the associated noise spectra separately. When, for simplicity, the forces acting on the propeller are assumed to be uniformly distributed at each radial station of the propeller blade, that is, when the chordwise distributions of forces are assumed to be rectangular, the functions $F_1(t,r,\theta)$ and $F_2(t,r,\theta)$ may be expressed, first for $\theta = 0$, as (see ref. 3)

$$F_{1}(t,r,0) = \frac{1}{B} \frac{dT}{dr} \frac{r}{b} \frac{d\theta}{dr} dr \qquad (0 < t < \tau)$$

$$F_{1}(t,r,0) = 0 \qquad (\tau < t < \tau_{0})$$

$$(3)$$

$$F_{2}(t,r,0) = \frac{1}{B} \frac{1}{r} \frac{dQ}{dr} \frac{r d\theta}{b} dr \quad (0 < t < \tau)$$

$$F_{2}(t,r,0) = 0 \qquad (\tau < t < \tau_{0})$$

$$(4)$$

where T is thrust, b is blade chord, Q is torque, $\tau = \frac{b}{r\Omega}$ is the time elapsed for a blade chord with angular velocity Ω to pass a given point, and $\tau_0 = \frac{2\pi}{B\Omega}$ is time elapsed for corresponding parts (for example, the leading edges) of two consecutive blades to pass a given point.

and

The expressions for $F_1(t,r,0)$ and $F_2(t,r,0)$ in equations (3) and (4), respectively, may be harmonically analyzed to yield

$$F_{1}(t,r,0) = \frac{r}{bBr_{0}} \frac{dT}{dr} d\theta dr \sum_{m=-\infty}^{\infty} \int_{0}^{T} e^{\frac{2m\pi i}{T_{0}}(t-t')} dt'$$
$$= \frac{1}{\pi} \sum_{-\infty}^{\infty} \frac{r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{b}{2\Omega r} \right)} \frac{dT}{dr} dr d\theta$$
(5)

and

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$$F_{2}(t,r,0) = \frac{1}{\pi} \sum_{-\infty}^{\infty} \frac{r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{b}{2\Omega r} \right)} \frac{1}{r} \frac{dQ}{dr} dr d\theta$$
(6)

In these expression the zeroeth harmonic, m = 0, corresponds to the instantaneous average thrust and torque over the blade element which does not give rise to the rotational-type noise under consideration and need not be considered further in the present analysis. For all other harmonics the expressions in equations (5) and (6) may be written as

$$F_{1}(t,r,0) = \frac{1}{\pi} R.P. \sum_{m=1}^{\infty} \frac{2r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{b}{2\Omega r} \right)} \frac{dT}{dr} dr d\theta \qquad (7)$$

and

$$F_{2}(t,r,0) = \frac{1}{\pi} R.P. \sum_{m=1}^{\infty} \frac{2r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{b}{2\Omega r} \right)} \frac{1}{r} \frac{dQ}{dr} dr d\theta \qquad (8)$$

Although only the real parts of the summations in equations (7) and (8) are necessary in the formulation of the forces, it is convenient to deal with expressions of the forces in complex form throughout the analysis and then, for final results, extract only real parts. Expressions for F_1 and F_2 that pertain to any value of θ may be obtained directly from the expressions in equations (7) and (8) for $\theta = 0$, because the forces on a second blade element that is shifted with respect to the first by an angle θ in the rotational direction are the same as the forces on the first but are retarded by the time θ/Ω . The corresponding Fourier developments are therefore real parts of

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$$F_{1}(t,r,\theta) = \frac{1}{\pi} \sum_{m=1}^{\infty} \frac{2r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{\theta}{\Omega} - \frac{b}{2\Omega r} \right)} \frac{dT}{dr} dr d\theta$$
(9)

and

$$F_{2}(t,r,\theta) = \frac{1}{\pi} \sum_{m=1}^{\infty} \frac{2r}{mbB} \left(\sin \frac{mbB}{2r} \right) e^{imB\Omega \left(t - \frac{\theta}{\Omega} - \frac{b}{2\pi r} \right)} \frac{1}{r} \frac{dQ}{dr} dr d\theta \qquad (10)$$

It can be seen from this development that, regardless of the form of the chordwise loading employed, that is, as long as the loading remains unchanged throughout the propeller cycle, the variable θ will always enter into the Fourier development of the forces in the same manner. As will be discussed subsequently, this relationship leads to a fixed set of integrals in the expressions for the sound field.

With the use of equations (9) and (10) for $F_1(t,r,\theta)$ and $F_2(t,r,\theta)$ the equations for the sound pressure associated with thrust and torque (eqs. (1) and (2)) corresponding to rectangular chordwise loading can be written for any harmonic different from zero as

$$p_{\rm T} = -\frac{e^{\rm i}\omega t}{\mu_{\pi}^2} \int_0^{\rm R} \frac{2r}{\rm mbB} \left(\sin \frac{\rm mbB}{2r}\right) \frac{dT}{dr} e^{-\frac{\rm i}mBb}{2r} dr \int_0^{2\pi} e^{-\rm i}mB\theta} \frac{\partial}{\partial x} \frac{e^{-\rm i}k\sigma}{S} d\theta$$
(11a)

 \mathbf{or}

$$P_{T} = -\frac{e^{i\omega t}}{\mu_{\pi}^{2}} \int_{0}^{R} \frac{2r}{mBb} \left(\sin \frac{mBb}{2r} \right) \frac{dT}{dr} e^{-\frac{imBb}{2r}} \frac{\mathbf{I}_{k}M_{F}}{\beta^{2}} \mathbf{I}_{1}(r) + \frac{\mathbf{i}_{kx}}{\beta^{2}} \mathbf{I}_{2}(r) + \mathbf{x}\mathbf{I}_{3}(r) dr$$
(11b)

and

$$p_{Q} = \frac{e^{i\omega t}}{\mu_{\pi}^{2}} \int_{0}^{R} \frac{2r}{mBb} \left(\sin \frac{mBb}{2r} \right) \frac{1}{r} \frac{dQ}{dr} e^{-\frac{imB}{2r}} dr \int_{0}^{2\pi} e^{-imB\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{e^{-ik\sigma}}{s} d\theta$$
(12a)

or

$$p_{Q} = \frac{imBe^{i\omega t}}{4\pi^{2}} \int_{0}^{R} \frac{2r}{mBb} \left(\sin \frac{mBb}{2r} \right) \frac{1}{r^{2}} \frac{dQ}{dr} e^{-\frac{imBb}{2r}} \Pi_{1}(r) dr \qquad (12b)$$

where

$$\Pi_{\nu}(r) = \int_{0}^{2\pi} \frac{e^{-i(mB\theta + k\sigma)}}{s^{\nu}} d\theta \qquad (\nu = 1, 2, \text{ and } 3) \qquad (13)$$

In the limit b = 0, these expressions for p_T and p_Q are equivalent to those given in equations (21) and (23) of reference 3. This equivalency may be easily verified by use of the limiting value of

$$\lim_{b \to 0} \left[\frac{2r}{mBb} \left(\sin \frac{mBb}{2r} \right) e^{\frac{-imBb}{2r}} = 1 \right]$$

The integrals in equations (11) to (13) must, in general, be evaluated by approximate or numerical procedures. The integrals II_v (for v = 1, 2, and 3) in equation (13) occur in expressions for propeller noise regardless of the form of chordwise loading employed, that is, provided the loading remains unchanged throughout the propeller cycle. These integrals therefore play an important role in the analytical determination of propeller noise and are the integrals for which the treatment and evaluations, pertaining to the near field, constitute the main purpose of this report. In the case of the far field the integrals II_v can, as was done in references 1 and 3, be reduced to Bessel functions of the first kind. This reduction is, for the sake of completeness, carried out in appendix A in a slightly different and more extended form than was done in references 1 and 3.

The Integrals Π_{ν} in Terms of Lumped Parameters mb, K_1 ,

and Ko and Equivalent Integrals I,

The integrals Π_{ν} are functions of six parameters (x, y, r, $M_{\rm F}$, mB, and k). These parameters can be lumped, however, so that the integrals can be conveniently expressed in terms of related integrals that are functions of only three parameters (mB, K₁, and K₂). This lumping is accomplished as follows:

$$II_{\nu} = \int_{0}^{2\pi} \frac{e^{-i(mB\theta + k\sigma)}}{s^{\nu}} d\theta = e^{-\frac{ikx}{\beta^{2}}} \int_{0}^{2\pi} \frac{e^{-imB(\theta + \sqrt{x^{2} + \beta^{2}y^{2} + \beta^{2}r^{2} - 2\beta^{2}ry \cos \theta)}}{(x^{2} + \beta^{2}y^{2} + \beta^{2}r^{2} - 2\beta^{2}ry \cos \theta)^{\nu/2}} d\theta$$

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$$\Pi_{\nu}(\mathbf{r}) = \left(\frac{\Omega}{\beta^{2}c}\right)^{\nu} \frac{e^{\frac{ikM_{F}x}{\beta^{2}}}}{\left(K_{2}\right)^{\nu}} I_{\nu}(\mathbf{r}) = \left(\frac{M_{R}}{\beta^{2}RK_{2}}\right)^{\nu} e^{\frac{ikM_{F}x}{\beta^{2}}} I_{\nu}(\mathbf{r})$$
(14)

where

$$I_{\nu}(\mathbf{r}) = \int_{0}^{2\pi} \frac{e^{-i\mathbf{m}B\left(\theta + K_{2}\sqrt{1 - K_{1}\cos\theta}\right)}}{\left(\sqrt{1 - K_{1}\cos\theta}\right)^{\nu}} d\theta = I_{\nu}' - iI_{\nu}'' \qquad (15)$$

$$K_{1} = \frac{2\beta^{2}ry}{x^{2} + \beta^{2}y^{2} + \beta^{2}r^{2}} = \frac{2\beta^{2}\frac{r}{R}\frac{y}{R}}{\left(\frac{x}{R}\right)^{2} + \beta^{2}\left(\frac{y}{R}\right)^{2} + \beta^{2}\left(\frac{r}{R}\right)^{2}}$$
(16)

$$K_{2} = \frac{\Omega}{\beta^{2}c} \sqrt{x^{2} + \beta^{2}y^{2} + \beta^{2}r^{2}}$$
$$= \frac{\Omega R}{\beta^{2}c} \sqrt{\left(\frac{x}{R}\right)^{2} + \beta^{2}\left(\frac{y}{R}\right)^{2} + \beta^{2}\left(\frac{r}{R}\right)^{2}}$$
$$= \frac{M_{R}}{\beta^{2}} \sqrt{\left(\frac{x}{R}\right)^{2} + \beta^{2}\left(\frac{y}{R}\right)^{2} + \beta^{2}\left(\frac{r}{R}\right)^{2}}$$
(17)

and $\,M_{\rm R}\,$ denotes the Mach number associated with rotational velocity. The integrals $\,I_{\nu}{\,}^{\prime}\,$ and $\,I_{\nu}{\,}^{\prime\prime}\,$ may be expressed as

$$I_{\nu}' = \int_{0}^{2\pi} \frac{\cos mB\left(\theta + K_{2}\sqrt{1 - K_{1} \cos \theta}\right)}{\left(\sqrt{1 - K_{1} \cos \theta}\right)^{\nu}} d\theta \qquad (18a)$$

or

$$I_{\nu}' = 2e^{-imB\pi} \int_{0}^{\pi} \frac{\cos mB\theta \cos mBK_{2}\sqrt{1 + K_{1} \cos \theta}}{\left(\sqrt{1 + K_{1} \cos \theta}\right)^{\nu}} d\theta \qquad (18b)$$

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or

and

$$I_{\nu}'' = \int_{0}^{2\pi} \frac{\sin mB(\theta + K_{2}\sqrt{1 - K_{1} \cos \theta})}{\left(\sqrt{1 - K_{1} \cos \theta}\right)^{\nu}} d\theta \qquad (19a)$$

 or

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$$I_{\nu}'' = 2e^{-imB\pi} \int_{0}^{\pi} \frac{\cos mB\theta \sin mBK_{2}\sqrt{1 + K_{1} \cos \theta}}{\left(\sqrt{1 + K_{1} \cos \theta}\right)^{\nu}} d\theta \qquad (19b)$$

Expressions for the pressures ${\rm p}_{\rm T}$ and ${\rm p}_{\rm Q}$ in terms of the integrals ${\rm I}_{\nu}$ may be written as

$$p_{T} = -\frac{M_{R}^{2}e^{imB\left(\Omega t - \frac{M_{R}M_{F}}{\beta^{2}}\frac{x}{R}\right)}}{\mu_{\pi}^{2}\beta^{\mu}R^{2}K_{2}}\int_{0}^{R}\frac{2r}{mBb}\left(\sin\frac{mBb}{2r}\right)\frac{dT}{dr}e^{-\frac{imBb}{2r}}\left[\left(mBM_{F}I_{1}^{"} + \frac{x}{R}\frac{mBM_{R}}{\beta^{2}K_{2}}I_{2}^{"} + \frac{x}{R}\frac{M_{R}}{\beta^{2}K_{2}^{2}}I_{3}^{"}\right) + i\left(mBM_{F}I_{1}^{"} + \frac{x}{R}\frac{mbM_{R}}{\beta^{2}K_{2}}I_{2}^{"} - \frac{x}{R}\frac{M_{R}}{\beta^{2}K_{2}^{2}}I_{3}^{"}\right)\right]dr$$
(20)

and

$$p_{Q} = \frac{mBM_{R}e}{\frac{mBM_{R}e}{\mu\pi^{2}\beta^{2}RK_{2}}} \int_{0}^{R} \frac{2r}{mBb} \left(\sin \frac{mBb}{2r}\right) \frac{1}{r^{2}} \frac{dQ}{dr} e^{-\frac{imBb}{2r}} \left(I_{1}" + iI_{1}'\right) dr$$
(21)

Tables of I_{ν} Corresponding to the Near Field

Since the integrals I_v are expressed in terms of lumped parameters K_l (eq. (16)) and K_2 (eq. (17)), it is necessary to determine the range of values of these parameters that pertain to the near field and to the far field. Regardless of the values that x, y, r, and β may have, numerical values of K_1 will always fall in the range

 $0 \leq K_1 \leq 1 \tag{22}$

but numerical values of K, fall in the semi-infinite range

$$0 \leq K_{2} \leq \infty$$
(23)

Inasmuch as the present development is inherently restricted to propeller tip speeds V_t that correspond to propeller-tip Mach numbers M_t less than about unity, the factor $\frac{M_R}{\beta^2} = \frac{M_R}{1 - M_F^2}$ of K_2 does not become large even when M approaches 1. This condition exists because

$$M_{t} = \sqrt{\left(\frac{V_{F}}{c}\right)^{2} + \left(\frac{\Omega R}{c}\right)^{2}} = \sqrt{M_{F}^{2} + M_{R}^{2}}$$
(24)

thus, when $M_F \approx 1$, $M_R \approx 0$ for $M_t \approx 1$. It therefore follows that large numerical values of K2 are associated only with large values of $\sqrt{x^2 + \beta^2 y^2}$ and these large values imply large distances from the propeller disk. Furthermore, large distances from the propeller disk imply small numerical values of K_1 ($K_1 \ll 1$). Hence, small values of K_1 and large values of K_2 ($K_2 \gg 1$) apply to the far field and large values of K_1 and small values of K_2 pertain to the near field. The specific ranges of values of K_1 and K_2 considered in the present report for the near field are $0.4 \leq K_1 \leq 0.985$ and $0.75 \leq K_2 \leq 6.0$, respectively. These ranges correspond to distances less than about 2.5 propeller diameters from the propeller tip. Tables of values of the integrals I_{ν} that cover these ranges of K_1 and K_2 for a number of values of the parameter mB are presented in tables I to VII. Evaluations are presented for values of mB of 2, 3, 4, 6, 8, 9, and 12; values of \bar{K}_1 of 0.4, 0.5, 0.6, 0.7, 0.75, 0.8, 0.85, 0.9, 0.925, 0.95, and 0.985; and values of K_2 at intervals of 0.25 in the range $0.75 \leq K_{2} \leq 6.0.$

The evaluations were made with the use of a card-programed electronic calculator and are presented, as calculated, to eight significant figures. They can be considered accurate throughout, however, to no more than four significant figures. In order to maintain this accuracy throughout the calculations, it was necessary to decrease the θ -increments at which the

integrands were evaluated in making the numerical integrations as the parameter mB was increased. The specific increments employed are for values of mB of 2, 3, and 4, $\Delta \theta = 6^{\circ}$; for a value of mB of 6, $\Delta \theta = 4^{\circ}$; and for values of mB of 8, 9, and 12, $\Delta \theta = 3^{\circ}$. With regard to interpolation in the tables, the increments in K₁ and K₂ are not small enough to yield good results by linear interpolation. Reasonably good interpolated results can be obtained, however, by graphical procedures or by a nonlinear scheme such as Aitken's method.

Consideration of Effects of Other Types of Chordwise Loading

It is recalled that the equations for the sound field of propellers derived in foregoing sections of this report are based on the assumption of a uniform chordwise loading. With real propellers, however, the actual form of the chordwise loading, which is, of course, not uniform, may be very significant with regard to the generation of rotational noise. For example, a propeller may be loaded, as in "turn ups", so as to produce no net thrust, yet it may produce considerable noise, mainly of a higher harmonic content. The fact is that negative forces produce noise as well as positive forces and it is only when the separate forces are in one direction that the assumption of uniform loading over the chord should be expected to yield realistic values for the calculated noise. Some examples that bear out this statement and demonstrate the significance of nonuniform chordwise loading are presented in the following paragraphs.

Examples based on square-wave type chordwise loading.- Consider first a propeller to have a constant loading F_0 on the front half-chord and constant loading $-F_0$ on the rear half-chord of a section of each blade at some radial station r. That is, let a force acting on each blade be defined by

$$f(t,r,\theta) = \frac{r}{bB} F_{0} \qquad \left(0 < t < \frac{T}{2}\right)$$

$$f(t,r,\theta) = -\frac{r}{bB} F_{0} \qquad \left(\frac{T}{2} < t < \tau\right)$$

$$f(t,r,\theta) = 0 \qquad \left(\tau < t < \tau_{0}\right)$$

$$f(t+\tau_{0},r,\theta) = f(t,r,\theta) \qquad (25)$$

This function corresponds to a zero net force and may be written as a Fourier series, with the zeroeth harmonic neglected, as

$$f(t,r,\theta) = R.P. \frac{4F_0}{\pi} \sum_{l}^{\infty} \frac{ir}{mBb} \left(\sin^2 \frac{mB}{4} \frac{b}{r} \right) e^{imB\left(\Omega t - \theta - \frac{b}{2r}\right)}$$
(26)

Next, for the sake of comparison, consider the full-wave rectification of $f(t,r,\theta)$. That is,

$$\begin{vmatrix} f(t,r,\theta) &= \frac{r}{bB} F_0 \qquad (0 < t < \tau) \\ \left| f(t,r,\theta) &= 0 \qquad (\tau < t < \tau_0) \right| \\ f(t + \tau_0,r,\theta) &= \left| f(t,r,\theta) \right| \qquad (27)$$

The Fourier expansion of this function is found to be, the zeroeth harmonic being neglected

$$|f(t,r,\theta)| = R.P. \frac{2F_0}{\pi} \sum_{m=1}^{\infty} \frac{r}{mBb} \left(\sin \frac{mBb}{2r} \right) e^{imB\left(\Omega t - \theta - \frac{b}{2r}\right)}$$
 (28)

It may be noted that there is a phase difference of 90°, indicated by the pure imaginary term in equation (26), between the harmonics of $f(t,r,\theta)$ and $|f(t,r,\theta)|$. The main thing of interest in the present discussion, however, is the relative values of the amplitude functions $\frac{hr}{mBb} \sin^2 \frac{mBb}{4r}$ of $f(t,r,\theta)$ and $\frac{2r}{mBb} \sin \frac{mBb}{2r}$ of $|f(t,r,\theta)|$. These functions are shown plotted as a function of $\frac{mBb}{2r}$ in figure 2. Examination of this figure shows that, for a range of values of $\frac{mBb}{2r} > \frac{\pi}{2}$, the amplitude function for $f(t,r,\theta)$ is greater than that for $|f(t,r,\theta)|$ even though the function $f(t,r,\theta)$ corresponds to no net force.

Examples based on measured distributions of normal force .- Under actual operating conditions the chordwise distribution of force is usually such that it can be closely approximated with the use of a few simple shapes such as rectangles and triangles. Furthermore, the Fourier development for such shapes can be made once and for all in such a way that, when a chordwise force distribution is known, its Fourier development can be obtained by simple superposition. In order to illustrate this point and to demonstrate further the significance of negative loading with regard to sound pressures, two cases of measured distribution of normal-force coefficients chosen from reference 4 are plotted in figure 3 for the operating conditions shown. Each of the two chosen distributions is for a blade section at r = 0.78R; however, one distribution (fig. 3(a)) corresponds to a small net loading with negative forces acting over part of the chord, and the other distribution (fig. 3(b)) corresponds to a large net loading with no negative forces. The area under each curve is divided into rectangles and triangles by the dashed lines shown in figure 3. (The right triangle facing to the right is herein designated right triangle, and the right triangle facing to the left is herein designated left triangle. Their elements are designated right triangular and left triangular, respectively.) The Fourier developments of arbitrarily placed rectangular- and right and left triangularshaped elements of loading are given in appendix B and have been employed to obtain the Fourier expansions of the distributions in figure 3. The amplitudes so obtained are plotted as a function of mB in figure 4 and compared with amplitude functions for uniform chordwise loading that yield the same net values as do the distributions in figure 3 and with amplitude functions associated with Dirac delta-type loadings. The results in figure 4(a) are for the case of low net loading and these results show that taking into account the negative forces as well as the positive forces can lead to considerably higher amplitude of sound in the higher harmonics. The results in figure 4(b) for the case of no negative forces show that about the same level of noise is obtained from the uniform distribution of forces as from the detailed distribution of forces. The amplitude functions for the Dirac delta-type loadings agree, as would be expected, with those for rectangular-type loading for lower values of mB.

The Integrations in the Radial Direction

A brief consideration is now given to a method of approximating the integrations with respect to r indicated in equations (11) and (12). These integrations can be accomplished by a generalization of the concept of effective radius discussed in references 1 to 3. As discussed in these references, the concept of effective radius is based on the consideration that all the sound-producing forces acting on the propeller can be assumed to be concentrated on an annular ring that roughly coincides with the circumference along which the resultant force acts.

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In the present case the assumption of an effective radius $r = R_e$ gives, from equations (20) and (21),

$$p_{T} = -\frac{M_{R}^{2} e^{imB\left(\Omega t - \frac{M_{R}M_{F}}{\beta^{2}} \frac{x}{R}\right)}}{4\pi^{2}\beta^{4}R^{2}K_{2}(R_{e})} G_{1}(R_{e}) \int_{0}^{R} \frac{dT}{dr} dr$$
$$= -\frac{M_{R}^{2} e^{imB\left(\Omega t - \frac{M_{R}M_{F}}{\beta^{2}} \frac{x}{R}\right)}}{4\pi^{2}\beta^{4}R^{2}K_{2}(R_{e})} TG_{1}(R_{e})$$
(29)

and

$$p_{Q} = \frac{mEM_{R}e}{4\pi^{2}\beta^{2}RK_{2}(R_{e})} G_{2}(R_{e}) \int_{0}^{R} \frac{dQ}{dr} dr$$

$$= \frac{\operatorname{imB}\left(\Omega t - \frac{M_{\mathrm{F}}M_{\mathrm{R}}}{\beta^{2}}\frac{x}{\mathrm{R}}\right)}{4\pi^{2}\beta^{2}RK_{2}(\mathrm{R}_{\mathrm{e}})} QG_{2}(\mathrm{R}_{\mathrm{e}})$$
(30)

where

$$G_{1}(R_{e}) = F_{m}\left(\frac{mBb}{R_{e}}\right) \left(\left\{ mBM_{F}I_{1}"(R_{e}) + \frac{x}{R_{e}} \frac{mBM_{R}}{\beta^{2}K_{2}(R_{e})} I_{2}"(R_{e}) + \frac{x}{R} \frac{M_{R}}{\beta^{2}\left[K_{2}(R_{e})\right]^{2}} I_{3}"(R_{e}) \right\} + \left. i \left\{ mBM_{F}I_{1}"(R_{e}) + \frac{x}{R} \frac{mBM_{R}}{\beta^{2}K_{2}(R_{e})} I_{2}"(R_{e}) - \frac{x}{R} \frac{M_{R}}{\beta^{2}\left[K_{2}(R_{e})\right]^{2}} I_{3}"(R_{e}) \right\} \right)$$
(31)

$$G_{2}(R_{e}) = F_{m}\left(\frac{mBb}{R_{e}}\right) \frac{1}{R_{e}^{2}} \left[I_{l}''(R_{e}) + iI_{l}'(R_{e})\right]$$
(32)

and $F_m\left(\frac{mBb}{R_e}\right)$ is an amplitude function that depends on the form of the chordwise blade loading. For uniform or rectangular-type chordwise loading (see eqs. (20) and (21)),

 $F_{m} = \frac{2R_{e}}{mBb} \left(\sin \frac{mBb}{2R_{e}} \right) e^{-\frac{imBb}{2R_{e}}}$

In many cases, especially when the point at which the sound pressure is to be calculated is outside and not too near the cylinder formed by the propeller disk and its wake, the total sound pressure is evidently not sensitive to the details of the radial distribution of forces; thus, the assumption of an effective radius serves as a means of approximating the total sound pressure. See, for example, comparisons of calculated and measured results in references 1 and 2. For cases where details of the radial loading are important, as would be true, for example, at near-field points inside the cylinder formed by the propeller disk and its wake, the effective-radius concept can be generalized as follows:

Suppose the distributions of thrust and torque are graphically represented as indicated in the sketch in figure 5. Divide the areas under the curve into n parts by dividing the blade length R into n parts. The parts into which R is divided need not be equal but they should be fairly small to insure a reasonably accurate approximation to the integrations. Let $R_{e,j}$ denote the abscissa of the center of area of the jth element of area under the curve. Approximate expressions for p_T and p_0 (eqs. (20) and (21)) may then be written as

$$p_{T} = -\frac{M_{R}^{2}e^{\frac{imB\left(\Omega t - \frac{M_{R}M_{F}}{\beta^{2}} \frac{x}{R}\right)}}{\mu_{\pi}^{2}\beta^{4}R^{2}} \sum_{j=1}^{n} \frac{G_{l}(R_{e,j})}{K_{2}(R_{e,j})} \Delta T_{j}$$
(33)

and

$$p_{Q} = \frac{\frac{\text{imB}\left(\Omega t - \frac{M_{R}M_{F}}{\beta^{2}}\frac{x}{R}\right)}{4\pi^{2}\beta^{2}R}}{\sum_{j=1}^{n} \frac{G_{2}(R_{e,j})}{K_{2}(R_{e,j})} \Delta Q_{j}$$
(34)

where $G_1(R_{e,j})$ and $G_2(R_{e,j})$ are defined in equations (31) and (32), respectively, and

$$\Delta \mathbf{T}_{j} = \int_{\mathbf{r}_{j}}^{\mathbf{r}_{j}+1} \frac{d\mathbf{T}}{d\mathbf{r}} d\mathbf{r}$$
(35)

and

$$\Delta Q_{j} = \int_{r_{j}}^{r_{j+1}} \frac{dQ}{dr} dr \qquad (36)$$

are the jth elements of area under the radial distributions of thrust and torque curves.

Root-Mean-Square Value of Pressure

Before some specific examples are discussed, it is appropriate to point out that, in applications of propeller noise theory, one is generally interested in the magnitude of the root mean square of the total pressure $p_{\rm rms}$ for each harmonic of noise. In order to obtain this quantity, it is convenient to treat the propeller forces as complex quantities. As may be noted in equations (20) and (21) or equations (33) and (34) the pressures $p_{\rm T}$ and $p_{\rm Q}$ are of forms

$$p_{\pi} = (\cos \omega t + i \sin \omega t) (A_1 + iA_2)$$
(37)

and

$$P_{Q} = (\cos \omega t + i \sin \omega t) (B_{1} + iB_{2})$$
(38)

The form of an expression for the total sound pressure p associated with thrust and torque is obtained by adding these expressions and retaining only the real parts of the sum. That is,

$$p = R.P.(p_T + p_Q) = (A_1 + B_1)\cos \omega t - (A_2 + B_2)\sin \omega t$$
(39)

The root-mean-square value of this expression is

$$P_{\rm rms} = \left\{ \frac{\omega}{2\pi} \int_{0}^{2\pi/\omega} \left[\left[A_{\rm l} + B_{\rm l} \right] \cos \omega t - \left(A_{\rm 2} + B_{\rm 2} \right) \sin \omega t \right]^2 dt \right\}^{1/2}$$
$$= \frac{1}{\sqrt{2}} \sqrt{\left(A_{\rm l} + B_{\rm l} \right)^2 + \left(A_{\rm 2} + B_{\rm 2} \right)^2} = \frac{1}{\sqrt{2}} \left| P_{\rm T} + P_{\rm Q} \right| \qquad (40)$$

Thus $p_{\rm rms}$ is obtained by simply multiplying the absolute value of the complex sum $p_{\rm T} + p_{\rm Q}$ by $\frac{1}{\sqrt{2}}$.

SOME APPLICATIONS TO SPECIFIC PROPELLERS

The propellers chosen for the applications of the tables are threeand four-blade 16-foot-diameter propellers with solidity factors ranging from 0.12 to 0.16. The efficiency, total thrust, and total torque for such propellers have been determined from charts of references 5 and 6 and are summarized in figure 6. The total thrust for three different altitudes and the efficiency are presented as functions of flight Mach number $M_{\rm F}$ in figures 6(a) and 6(b), respectively. The total torque is presented as a function of rotational Mach number $M_{\rm R}$ in figure 6(c).

The present examples are based on the assumption of uniform chordwise loading and, except for one case, on the concept of an effective radius (eqs. (29) and (30)) with $R_e = 0.8R$. When the calculated results were obtained, the necessary interpolations in the tables of I_V were performed by graphical procedures.

The first results are presented in figure 7 and are intended to give some indication as to the differences that might be expected between calculations based on the concept of an effective radius and calculations based on radial distributions of forces. Shown in the figure are values of the root-mean-square pressure, associated with the first harmonic of a four-blade 16-foot-diameter propeller, calculated with the use of equations (33) and (34) (by using five steps), for radial distributions of forces shown in the upper plot and calculated with the use of equations (29) and (30) for the forces assumed to be concentrated on an annulus of radius 0.8R. Calculations by each procedure were made for operating conditions indicated in figure 7 at various values of x in the range $-0.3D \leq x \leq 0.3D$ along a line y = 0.7D (11.2 feet); that is, along a line two-fifths of the propeller radius from the tip and extending a distance three-fifths of the radius behind to three-fifths of the radius ahead of the plane of rotation. For the conditions considered, the pressure or sound levels based on distributed forces are below those

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based on effective radius or concentrated forces at each point considered. The significant point is that the shapes of the two curves, considered as functions of x/D, are about the same; this fact indicates that calculations based on an effective radius can at least be used for studying trends.

In figure 8 calculated values of the root-mean-square pressure associated with the first harmonic of a four-blade 16-foot-diameter propeller are plotted as functions of x/D for two altitude conditions, three values of rotational Mach number M_R , and a fixed value (0.6) of flight Mach number M_F .

As can be noted in figure 8 for the set of conditions under consideration, the peaks or highest amplitudes of pressure occur behind the plane of rotation at values of x/D in the range from -0.05 to -0.125. A somewhat surprising feature of the results is the rather sharp rearward trend of the position of peak amplitude as M_R is increased. As noted in reference 3, the trend of position of peak amplitude of pressure with increasing M_F is from some rearward position toward the plane of rotation or just the opposite of that indicated in figure 7 for increasing M_R . The results also indicate a rather sharp attenuation of sound-pressure level with altitude and with distance away from the location of peak amplitude. The apparent attenuation with altitude, however, can be directly associated with the decrease in thrust and torque (see fig. 6) required for the higher altitude.

In figure 9 the calculated root-mean-square pressures associated with the first harmonic of a four-blade, 16-foot-diameter propeller for a point x = -0.15D, y = 11.2 feet (or 0.7D) are shown plotted as a function of flight Mach number for three altitude conditions and a fixed rotational Mach number $(M_R = 0.8)$. These results show, as do the examples in reference 3, that, for a fixed altitude, the amplitudes of sound pressure decrease substantially with flight Mach number for Mach numbers up to about 0.4 or 0.5 and then increase rapidly with flight Mach number.

Figure 10 shows the effect of flight Mach number on the calculated root-mean-square pressures for various values of the parameter mB for two altitude conditions. The calculations are for a point x = 0, y = 11.2 feet (or 0.7D) and for a fixed value of the rotational Mach number ($M_R = 0.7$). A change in the parameter mB may correspond to either a change in number of propeller blades B or a change in harmonic m. For example, the curves in figure 10 for mB = 8 may be interpreted as the calculated sound pressure corresponding to either the first harmonic of a 8-blade propeller, the second harmonic of a 4-blade propeller.

All the curves in figure 10 show the falling-off or diminishing characteristic as the Mach number $M_{\rm F}$ is increased from zero up to 0.2 or 0.3 and then the rapid increase with Mach number when $M_{\overline{W}}$ exceeds about 0.4 or 0.5. The total amount that the calculated pressures decrease from the $M_{T} = 0$ condition is about the same percentage (about 20) in each case. It may be implied therefore that, in the neighborhood of take-off conditions, the Mach number effect is about the same, percentagewise, regardless of the value of mB. The increase with Mach number, after a Mach number of 0.4 or 0.5 has been attained, however, depends strongly on the value of mB and becomes very large as mB increases beyond 6 or 8. For these large values of mB the present simplified theory, especially with the assumption of an effective radius, could only serve to indicate the trends roughly. In order to yield realistic magnitudes, it is likely that the details of the radial loading and, perhaps, details of the chordwise loadings would have to be taken into account.

The effect of distance from the propeller on the sound level along a line of constant y is implied in figures 7 and 8. Figure 11 shows the corresponding effect along a line of constant x (x = -0.15D). The results shown are the calculated root-mean-square pressures corresponding to the first harmonic of a four-blade 16-foot-diameter propeller for $M_{\rm F}$ = 0 and $M_{\rm R}$ = 0.7 at sea level plotted as a function of d/D where d is the distance from the propeller tip. Figure 11 shows that, for the conditions considered, the amplitude of pressure drops off very rapidly as the distance from the propeller tip is increased. This type of dropoff is typical of that expected for other values of the parameters $M_{\rm F}$, $M_{\rm R}$, and mB.

CONCLUDING REMARKS

A main purpose of this report is to present tabulated values of certain definite integrals that are involved in the calculation of nearfield propeller noise. These tabulations cover a wide range of operating conditions and are useful for estimating lower harmonic, near-field propeller noise when either the concept of an effective radius or radial distributions of forces are considered.

Consideration is given to the general forms of chordwise and radial load distributions and it is shown that propellers operating even with zero net forces can give rise to considerable noise, especially in the

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higher harmonics. Results of applications of noise theory to some specific 16-foot-diameter propellers are presented and discussed.

Langley Aeronautical Laboratory, National Advisory Committee for Aeronautics, Langley Field, Va., June 26, 1956.

APPENDIX A

EVALUATION OF THE INTEGRAIS I, FOR THE FAR FIELD

As pointed out in the text, the integrals I_{ν} corresponding to the far field or small values of the parameter K_{l} ($K_{l} \ll l$) can be expressed in terms of Bessel functions of the first kind. The purpose of this appendix is to reduce the integrals to such expressions.

For this purpose it is convenient to consider the integral expression for I_{ν} given in equation (15); that is,

$$I_{\nu}(r) = \int_{0}^{2\pi} \frac{e^{-imB\left(\theta + K_{2}\sqrt{1 - K_{1} \cos \theta}\right)}}{\left(\sqrt{1 - K_{1} \cos \theta}\right)^{\nu}} d\theta = I_{\nu}' - iI_{\nu}''$$

The value of the radical $\sqrt{1 - K_1 \cos \theta}$ appearing in the exponential term and in the denominator of the integrand of this expression is, for $K_1 \ll 1$, closely approximated by

$$\sqrt{1 - K_{\perp} \cos \theta} \approx 1 - \frac{K_{\perp}}{2} \cos \theta$$
 (A1)

Correspondingly, the integral I_{ν} is closely approximated by

$$I_{\nu} \approx \int_{0}^{2\pi} \frac{e^{-imB\left(\theta + K_{2} - \frac{K_{1}K_{2}}{2}\cos\theta\right)}}{\left(1 - \frac{K_{1}}{2}\cos\theta\right)^{\nu}} d\theta$$

 \mathbf{or}

$$I_{\nu} \approx e^{-imBK_{2}} \int_{0}^{2\pi} e^{-imB\left(\theta - \frac{K_{1}K_{2}}{2}\cos\theta\right)} \left(1 + \frac{\nu K_{1}}{2}\cos\theta\right) d\theta \qquad (A2)$$

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From the known integral representation of Bessel functions,

$$\int_{0}^{2\pi} e^{-i(n\theta - \lambda \cos \theta)} d\theta = 2\pi i^{n} J_{n}(\lambda)$$
 (A3)

and

$$\cos \theta_{e} - i(n\theta - \lambda \cos \theta)_{d\theta} = -i \frac{\partial}{\partial \lambda} \int_{0}^{2\pi} e^{-i(n\theta - \lambda \cos \theta)}_{d\theta}$$
$$= -\pi i^{n+1} [J_{n-1}(\lambda) - J_{n+1}(\lambda)] \qquad (A4)$$

The expression for I_{ν} given in equation (A2) may be written as

$$I_{v} \approx \pi i^{mB} e^{-imBK_{2}} \left\{ 2J_{mB} \left(\frac{mBK_{1}K_{2}}{2} \right) - \frac{ivK_{1}}{2} \left[J_{mB-1} \left(\frac{mBK_{1}K_{2}}{2} \right) - J_{mB+1} \left(\frac{mBK_{1}K_{2}}{2} \right) \right] \right\}$$
(A5)

or, by substituting the relations for $\ensuremath{\,\mathrm{K}_{\mathrm{l}}}$ and $\ensuremath{\,\mathrm{K}_{\mathrm{2}}},$

$$I_{v} \approx \pi i^{mB} e^{-\frac{ik}{\beta^{2}} \sqrt{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}}} \left\{ 2J_{mB} \left(\frac{kry}{\sqrt{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}}} \right) - \frac{i \frac{\nu \beta^{2} ry}{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}}}{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}} \left[J_{mB-1} \left(\frac{kry}{\sqrt{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}}} \right) - J_{mB+1} \left(\frac{kry}{\sqrt{x^{2} + \beta^{2} y^{2} + \beta^{2} r^{2}}} \right) \right] \right\}$$
(A6)

or, since $\beta r \ll \sqrt{x^2 + \beta^2 y^2}$ for large distances from the propeller disk,

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$$\begin{split} \mathbf{I}_{\nu} &\approx \pi \mathbf{1}^{mB} e^{-\frac{\mathbf{1}\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}}} \begin{cases} 2 J_{mB} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) - \\ \frac{\mathbf{1} \nu \beta^{2} \mathbf{ry}}{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \left[J_{mB-1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}}} \right) - J_{mB+1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}}} \right) \right] \end{cases} \\ &= \pi \mathbf{1}^{mB} \left(\left(2 \cos \frac{\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) J_{mB} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}}} \right) - \\ \frac{\nu \beta^{2} \mathbf{ry}}{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \left(\sin \frac{\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) J_{mB} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}}} \right) - \\ \frac{1}{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \left(\sin \frac{\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) \left[J_{mB-1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) - J_{mB+1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) \right] - \\ \frac{1}{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \left(\cos \frac{\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) J_{mB} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) + \\ \frac{\nu \beta^{2} \mathbf{ry}}{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \left(\cos \frac{\mathbf{k}}{\beta^{2}} \sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) \left[J_{mB-1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) - J_{mB+1} \left(\frac{\mathbf{kry}}{\sqrt{\mathbf{x}^{2} + \beta^{2} \mathbf{y}^{2}} \right) \right] \right\} \right) \\ (A7) \end{split}$$

Thus, for mB even,

$$I_{\nu}' = \pi (-1)^{mB/2} \Biggl\{ \Biggl(2 \cos \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \Biggr) J_{mB} \Biggl(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \Biggr) - \frac{\nu \beta^2 ry}{x^2 + \beta^2 y^2} \Biggl(\sin \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \Biggr) \Biggl[J_{mB-1} \Biggl(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \Biggr) - J_{mB+1} \Biggl(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \Biggr) \Biggr] \Biggr\}$$
(A8)

and

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$$I_{\nu}'' = \pi (-1)^{mB/2} \left\{ \left(2 \sin \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) J_{mB} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) + \frac{\nu \beta^2 ry}{x^2 + \beta^2 y^2} \left(\cos \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) \left[J_{mB-1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) - J_{mB+1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) \right] \right\}$$
(A9)

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For mB odd,

$$I_{\nu}' = \pi(-1)^{\frac{mB-1}{2}} \left\{ \left(2 \sin \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) J_{mB} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) + \frac{\nu \beta^2 ry}{x^2 + \beta^2 y^2} \left(\cos \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) \left[J_{mB-1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) - J_{mB+1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) \right] \right\}$$
(A10)

and

$$I_{\nu}'' = \pi(-1)^{\frac{mB+1}{2}} \left\{ \left(2 \cos \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) J_{mB} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) - \frac{\nu \beta^2 ry}{x^2 + \beta^2 y^2} \left(\sin \frac{k}{\beta^2} \sqrt{x^2 + \beta^2 y^2} \right) \left[J_{mB-1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) - J_{mB+1} \left(\frac{kry}{\sqrt{x^2 + \beta^2 y^2}} \right) \right] \right\}$$
(A11)

These approximations of I_{ν} for the far field go a step beyond that given in reference 3 in that the terms herein involving ν as a factor were not considered in reference 3. It is not deemed necessary to tabulate I_{ν} for the far-field case since tabulations of the Bessel functions J_n are so abundant in the literature.

APPENDIX B

THE FOURIER DEVELOPMENT OF SOME SIMPLE SHAPES USEFUL FOR

REPRESENTING CHORDWISE FORCE DISTRIBUTIONS

FOR PROPELLER BLADES

It is assumed that any chordwise distribution of force acting on a propeller-blade section can be satisfactorily represented by superimposing rectangular- and triangular-shaped elements of force as indicated in figure 3. This, at least, is a convenient assumption because such elements can be expanded into Fourier series in such a way as to make the Fourier development of the force distribution obtainable by simply adding the Fourier series representations of the different components into which the force distribution is divided.

Rectangular Element of Loading

Consider first a rectangular element of periodic force $f_r(t,r,\theta)$ that is located on some element of chord in a force-time system.



If the origin of coordinates is chosen to be the leading edge of the blade section, f_r may be expressed as follows:

$$f_{r}(t,r,\theta) = 0 \qquad 0 < t < \tau_{1}$$

$$f_{r}(t,r,\theta) = F_{r} d\theta dr \qquad \tau_{1} < t < \tau_{2}$$

$$f_{r}(t,r,\theta) = 0 \qquad \tau_{2} < t < \tau_{0}$$

$$f_{r}(t+\tau_{0},r,\theta) = f_{r}(t,r,\theta) \qquad (B1)$$

where

$$\tau_{1} = \mu_{1}\tau = \mu_{1}\frac{b}{r\Omega}$$
$$\tau_{2} = \mu_{2}\tau = \mu_{2}\frac{b}{r\Omega}$$

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The terms μ_1 and μ_2 are constants such that $0 \leq \mu_1 \leq 1$ and $0 \leq \mu_2 \leq 1$. Equations (B1) may be expressed as a Fourier series, the term corresponding to the zeroeth harmonic being neglected, as

$$f_{\mathbf{r}}(\mathbf{t},\mathbf{r},\theta) = \frac{2F_{\mathbf{r}}}{\tau_{0}} \frac{d\theta}{dr} \sum_{l}^{\infty} \left[\cos \frac{2m\pi}{\tau_{0}} \left(\mathbf{t} - \frac{\theta}{\Omega} \right) \int_{\tau_{1}}^{\tau_{2}} \cos \frac{2m\pi}{\tau_{0}} \mathbf{t}' \, d\mathbf{t}' + \frac{\sin \frac{2m\pi}{\tau_{0}} \left(\mathbf{t} - \frac{\theta}{\Omega} \right) \int_{\tau_{1}}^{\tau_{2}} \sin \frac{2m\pi}{\tau_{0}} \mathbf{t}' \, d\mathbf{t}' \right] \\ = \frac{2F_{\mathbf{r}}}{\pi} \frac{d\theta}{\tau} \frac{d\mathbf{r}}{\tau_{0}} \sum_{l}^{\infty} \frac{1}{m} \left[\cos \frac{2m\pi}{\tau_{0}} \left(\mathbf{t} - \frac{\theta}{\Omega} \right) \cos \frac{m\pi}{\tau_{0}} (\tau_{2} + \tau_{1}) \sin \frac{m\pi}{\tau_{0}} (\tau_{2} - \tau_{1}) + \frac{\sin \frac{2m\pi}{\tau_{0}} \left(\mathbf{t} - \frac{\theta}{\Omega} \right) \sin \frac{m\pi}{\tau_{0}} (\tau_{2} + \tau_{1}) \sin \frac{m\pi}{\tau_{0}} (\tau_{2} - \tau_{1}) \right] \\ = \frac{2F_{\mathbf{r}}}{\pi} \frac{d\theta}{\tau} \frac{d\mathbf{r}}{\tau_{0}} \sum_{l}^{\infty} \frac{1}{m} \left[\cos \frac{2m\pi}{\tau_{0}} \left(\mathbf{t} - \frac{\theta}{\Omega} - \frac{\tau_{1}}{2} + \frac{\tau_{2}}{2} \right) \sin \frac{m\pi}{\tau_{0}} (\tau_{2} - \tau_{1}) \right]$$
(B2)

This expression may be written in complex form as

$$f_{r}(t,r,\theta) = R.P. \frac{2F_{r}}{\pi} \sum_{l}^{\infty} \frac{1}{m} e^{\frac{2im\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega}\right)} \left[\sin \frac{m\pi}{\tau_{0}} \left(\tau_{2} - \tau_{1}\right)\right] e^{-\frac{im\pi}{\tau_{0}} \left(\tau_{1} + \tau_{2}\right)} d\theta dr$$
(B3)

 \mathbf{or}

$$f_{r}(t,r,\theta) = R.P. \frac{2F_{r}}{\pi} \sum_{l}^{\infty} \frac{e^{imB(\Omega t - \theta)}}{m} \left[\sin \frac{mBb}{2r} (\mu_{2} - \mu_{1}) \right] e^{-\frac{imBb}{2r} (\mu_{2} + \mu_{1})} d\theta dr$$
(B4)

The total force represented by $f_r(t,r,\theta)$ that acts on a blade section may be seen from equation (B1) to be

$$f_r(t,r,\theta) = F_r(\mu_2 - \mu_1)\tau \, d\theta \, dr \tag{B5}$$

If this force is to represent some preassigned percentage h_r of a total force of magnitude F, such as total thrust or total torque, that acts on the section, then

$$F_{r}(\mu_{2} - \mu_{1})\tau = \frac{h_{r}r}{bB}F\tau$$
 (B6)

or

$$\mathbf{F}_{\mathbf{r}} = \frac{\mathbf{h}_{\mathbf{r}}}{\mu_2 - \mu_1} \frac{\mathbf{r}}{\mathbf{Bb}} \mathbf{F}$$
(B7)

Substituting this expression for F_r into equation (B4) gives

$$f_{r}(t,r,\theta) = R.P. \frac{h_{r}F}{\pi(\mu_{2} - \mu_{1})} \sum_{l}^{\infty} \frac{2r}{mBb} e^{imB(\Omega t - \theta)} \left[sin \frac{mBb}{2r}(\mu_{2} - \mu_{1}) \right] e^{-\frac{imBb}{2r}(\mu_{2} + \mu_{1})} d\theta dr$$
(B8)

Note that in this equation, if $\mu_1 = 0$, $\mu_2 = 1$, and $h_r F = \frac{dT}{dr}$, the results given in equation (9) are obtained. If $\mu_1 = 0$, $\mu_2 = 1$, and $h_r F = \frac{1}{r} \frac{dQ}{dr}$, the results given in equation (10) are obtained.

Left-Triangular Element of Loading

Consider next a left-triangular element of loading $f_{lt}(t,r,\theta)$ on the chord element $\overline{\tau_1 \tau_2}$ in a time-coordinate system (see sketch).



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The equation for this case may be expressed as

$$f_{lt}(t,r,\theta) = 0 \qquad (0 < t < \tau_{1})$$

$$f_{lt}(t,r,\theta) = \frac{F_{t}(t - \tau_{1})}{\tau_{2} - \tau_{1}} d\theta dr \qquad (\tau_{1} < t < \tau_{2})$$

$$f_{lt}(t,r,\theta) = 0 \qquad (\tau_{2} < t < \tau_{0})$$

$$f_{lt}(t+\tau_{0},r,\theta) = f_{lt}(t,r,\theta) \qquad (B9)$$

The Fourier expansion for this function is

$$f_{lt}(t,r,\theta) = \frac{2F_{t} d\theta dr}{\tau_{0}(\tau_{2} - \tau_{1})} \sum_{l}^{\infty} \left[\cos \frac{2m\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega} \right) \int_{\tau_{1}}^{\tau_{2}} \left(t' - \tau_{1} \right) \cos \frac{2m\pi}{\tau_{0}} t' dt' + \frac{\sin \frac{2m\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega} \right) \int_{\tau_{1}}^{\tau_{2}} \left(t' - \tau_{1} \right) \sin \frac{2m\pi}{\tau_{0}} t' dt' \right]$$

$$= \frac{F_{t} dr d\theta}{\pi(\tau_{2} - \tau_{1})} \sum_{l}^{\infty} \left[-\frac{\tau_{2} - \tau_{1}}{m} \sin \frac{2m\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega} - \tau_{2} \right) + \frac{\tau_{0}}{m^{2}\pi} \sin \frac{m\pi}{\tau_{0}} \left(\tau_{2} - \tau_{1} \right) \sin \frac{2m\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega} - \frac{\tau_{1} + \tau_{2}}{2} \right) \right]$$
(B10)

or in complex form

$$f_{lt}(t,r,\theta) = R.P. \frac{F_{t} dr d\theta}{\pi \langle \tau_{2} - \tau_{1} \rangle} \sum ie^{\frac{2im\pi}{\tau_{0}} \left(t - \frac{\theta}{\Omega}\right)} \left\{ \frac{\tau_{2} - \tau_{1}}{m} e^{\frac{2im\pi\tau_{2}}{\tau_{0}}} - \frac{\tau_{1}}{m} e^{\frac{2im\pi\tau_{2}}{\tau_{0}}} - \frac{\tau_{1}}{m} e^{\frac{\tau_{1}}{\tau_{0}}} \right\}$$
(B11)

or

$$f_{lt}(t,r,\theta) \approx R.P. \frac{F_{t} dr' d\theta}{\pi} \sum_{l}^{\infty} i \frac{e^{imB(\Omega t - \theta)}}{m} \left\{ e^{-\frac{imBb}{r}\mu_{2}} - \frac{2r}{mBb(\mu_{2} - \mu_{1})} \left[sin \frac{mBb}{2r}(\mu_{2} - \mu_{1}) \right] e^{-\frac{imBb}{2r}(\mu_{2} + \mu_{1})} \right\}$$
(B12)

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As in the preceding paragraph, if the total force acting on a blade section is to represent some preassigned percentage $\,h_{\rm t}\,$ of a total force F, then

$$\mathbf{F}_{t} = \frac{2h_{lt}r}{\left(\mu_{2} - \mu_{1}\right)} \frac{r}{Bb} \mathbf{F}$$
(B13)

and

$$f_{lt}(t,r,\theta) = \text{R.P.} \frac{h_{lt}F}{\pi(\mu_2 - \mu_1)} \sum_{l}^{\infty} \frac{2ire^{imB(\Omega t - \theta)}}{mBb} \left\{ e^{-\frac{imBb}{r}\mu_2} - \frac{2r}{mBb(\mu_2 - \mu_1)} \left[sin \frac{mBb}{2r}(\mu_2 - \mu_1) \right] e^{-\frac{imBb}{2r}(\mu_2 + \mu_1)} \right\} dr d\theta \quad (B14)$$

This result is easily checked by considering the limit

$$\lim_{\mu_1 \to \mu_2} f_{lt}(t,r,\theta)$$
(B15)

This limit leads, as it should, to the Fourier expansion of the Dirac delta function of strength $h_{lt}F$.

Right-Triangular Element of Loading

A right-triangular element of loading f_{rt} in a time-coordinate system is shown in the following sketch:



The equation for this element of loading may be written as

$$f_{rt}(t,r,\theta) = 0 \qquad (0 < t < \tau_1)$$

$$f_{rt}(t,r,\theta) = -\frac{F_t(t - \tau_2)}{\tau_2 - \tau_1} \qquad (\tau_1 < t < \tau_2)$$

$$f_{rt}(t,r,\theta) = 0 \qquad (\tau_2 < t < \tau)$$

$$f_{rt}(t+\tau_0,r,\theta) = f_{rt}(t,r,\theta)$$
(B16)

The Fourier expansion of this equation can be obtained from equation (B14) by simply interchanging μ_1 and μ_2 . That is,

$$f_{rt}(t,r,\theta) = -R.P. \frac{h_{rt}F}{\pi(\mu_2 - \mu_1)} \sum_{l}^{\infty} \frac{2ire^{imB(\Omega t - \theta)}}{mBb} \left\{ e^{-\frac{imBb}{r}\mu_l} - \frac{2r}{mBb(\mu_2 - \mu_1)} \left[sin \frac{mBb}{2r}(\mu_2 - \mu_1) \right] e^{-\frac{imBb}{2r}(\mu_2 + \mu_1)} \right\} dr d\theta \quad (B17)$$

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REFERENCES

- 1. Gutin, L.: On the Sound Field of a Rotating Propeller. NACA TM 1195, 1948. (From Physik. Zeitschr. der Sowjetunion, Bd 9, Heft 1, 1936, pp. 57-71.)
- 2. Hubbard, Harvey H., and Regier, Arthur A.: Free-Space Oscillating Pressures Near the Tips of Rotating Propellers. NACA Rep. 996, 1950. (Supersedes NACA TN 1870.)
- 3. Garrick, I. E., and Watkins, Charles E.: A Theoretical Study of the Effect of Forward Speed on the Free-Space Sound-Pressure Field Around Propellers. NACA Rep. 1198, 1954. (Supersedes NACA TN 3018.)
- 4. Maynard, Julian D., and Murphy, Maurice P.: Pressure Distributions on the Blade Sections of the NACA 10-(3)(066)-033 Propeller Under Operating Conditions. NACA RM L9LL2, 1950.
- 5. Crigler, John L., and Talkin, Herbert W.: Charts for Determining Propeller Efficiency. NACA WR L-144, 1944. (Formerly NACA ACR L4129.)
- 6. Biermann, David, Gray, W. H., and Maynard, Julian D.: Wind-Tunnel Tests of Single- and Dual-Rotating Tractor Propellers of Large Blade Width. NACA WR L-385, 1942. (Formerly NACA ARR, Sept. 1942.)

TABLE I. - VALUES OF THE INTEGRAL $\mathbf{1}_{p}$ FOR all = 2

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(a) I₁'

	,4	.,	,6	.7	.75	.8	.85	.9	.925	-95	.985
0.1.1.1.2.2.4.5.5.0.4.4.4.5.5.5.0 0.1.1.1.2.2.4.5.5.0.4.5.5.5.0 0.1.1.1.2.2.4.4.5.5.5.5.0 0.1.1.1.2.2.4.4.5.5.5.5.0	0.1571;01,77 .21129075 .21305216 .25229673 .100507h3 .00707379 24232753 51107301 90958072 -1.1139484 1059263 60051938 20921306 .57139147 1.3781786 2.213822 1.9285725 1.1373823 01932500 -1.2871545 -2.35710058	0.26185519 .33250404 .h0105844 .h2075502 .33229132 .06785937 330568999 83992925 -1.7037759 -1.7037759 -1.7037759 -1.7037759 -1.7037759 -1.7037597 -2.55450668 .55460668 .55460668 .53460696 .53460696 .53460696 .53460696 .53460696	0.1100003 .5160076 .62203162 .6606921 .51633111 .21721289 -31666473 -1.0733165 -2.3135130 -2.3135130 -2.3135130 -2.413536 -2.0716911 -1.166639 .15520710 1.5562380 2.7732202 3.3276358	0.63076358 .75228543 .93714078 1.0081107 .86276332 .46953152 25273869 -1.2001658 -2.9128572 -3.2399538 -2.9128572 -3.2399538 -2.99202607 -1.9901114 62651362 .86342967	0.78344608 .96305196 1.2140751 1.1157200 .67154751 12180917 12180917 -1.1802554 -2.8880475 -3.1734767 -3.57576666	0.9826282h 1.1946901 1.6138008 1.6138008 1.6260419 .95430731 .9745277 -1.064584 -2.293330 -3.30396h -3.8223117 -3.6213621 -2.9280483	1.2573318 1.5775615 1.7768553 1.9120805 1.8162901 1.359205 .14138211 79711107 -2.133676 -3.269261 -3.26982677	1.6713358 1.5707001 2.5963038 2.5151841 2.5564360 1.67756860 1.67756860 1.6755833 26962642 -2.5061465 -3.6479637	1.9894793 2.3132179 2.6723681 2.9260391 2.9260391 2.4289153 1.48682171 -1.6282171 -1.6282171 -2.520504 -3.2997286	2.4560452 2.41114529 3.5102675 3.5059026 3.5115576 3.07118939 2.1106715 .82676565 64876565 1.877553 -2.67111858	3.9553356 k.3631590 k.840004 5.2166602 5.301723k k.9322010 k.9595045 2.7957927 1.397633k .1856520k 57126580

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TABLE I.- VALUES OF THE INTERAL I, FOR = 2 - Continued

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0.0 1.0 2.5 5.0 5.0 5.0 5.0 5.0 5.0 5.0 5.0 5.0 5	0.29511042 .28066053 .23819112 .145445 02211564 57937951 57937951 57937951 59762634 -1.1386299 -1.200374 -1.0283461 56420387 .11308239 .7471389 2.2579079 2.3337759 1.8877974 .95576261 29659990 -1.66558472	0.51539787 .152201379 .12539760 .00159736 -37322017 -86695005 -1.37122017 -86695005 -1.371269 -1.7735279 -1.1592506 -1.7551169 1.02071226 2.1833015 3.0351027 3.3399963 2.9722565 1.9616098 .15971801 .10791801 -1.108939 -2.5013784	0.961.71044 .83220931 .73644402 .52644958 .14497036 ~422448010 -1.1505315 -1.9289534 -2.59331173 -2.948007 -2.848007 -2.848007 .53866131 2.0772258 3.3794222 3.98660235	1.4564637 1.4120775 1.781604 .97201426 .47201547 -1.3250160 -2.4462791 -3.4555922 -4.1227686 -4.2280267 -3.4631154 -2.4666927 83353448 .92995030	1.9233477 1.8724745 1.7225193 1.3895692 .75023455 -11884207 -2.6213255 -3.8536950 -4.7343089 -5.0309606	2.6036168 2.5159134 2.3752146 1.9956853 1.3109753 .66816813 1.0095525 2.5158234 -1.0985525 2.5158234 -5.8168722 -5.8168355 -5.67683552 -4.8159139	3.6809421 3.0.55577 3.1.226151 2.9932116 2.2169011 1.0295650 5.3809950 -2.3406428 -4.1256926 -5.598323 -6.50151.94	5-0120098 5-5686813 5-3557042 1-24671246 3-9965740 2-67521276 -1-2226336 -3-31465956 -5-225756 -6-5562645	7.4109188 7.333459 7.1056250 6.5994426 5.6686080 4.210319 2.336587 .09658710 -2.2201409 -5.9165575	10.525224 10.44,3491 10.202012 9.6632493 8.684,2970 7.1733719 5.14,1900 2.74,1967 - 2.111,5646 - 2.111,5646 - 2.111,564	2.538569 25.150876 25.190891 31.610711 12.5751572 21.918991 19.709073 17.061390 11.256326 11.518777 9.15550055

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(d) I₂"

L ₂	.4	.5	,6	.7	.75	.8	.85	.9	.925	.95	.985
505550555055505550 0111112880550055505505550 011111288055005550550555055550	0.18339268 .27271003 .381,42362 .51039267 .624,8038 .6339333 .4296689 .051,4464 1,4602,827 .021,4464 1,271740 1,2721740 52171740 52171740 52171740 521737 34751823 34751855555555555555555	0.30630511 .45560791 .66051028 .81922376 1.0431363 1.1552980 .26612926 .5101971 5101971 5101971 5101971 2.657621 2.657621 2.13705786 3.5912186 3.5912186 3.5912186 3.5912186	0.60032970 .72134548 1.0772356 1.3322645 1.3225645 1.6225340 1.4425433 1.84255340 1.4427539 .74549531 1.303225 2.509308 -3.533125 -3.535125 -2.549308 -3.5161519 -3.7415072 -3.34645185 73208484	0.76613847 1.1179336 1.5593419 2.4415250 2.5246764 2.5246764 2.5246764 2.5246764 2.5846146 1.7758821 -1.1296985 -3.9477951 -1.1296985 -3.9477954 -2.6966406 -3.9477242 -4.037774 -4.5525006	0.96234789 1.3972517 1.9268792 2.5317469 3.1343089 3.5966754 3.14293685 2.5579965 1.1774527 ~.53436265	1.2231707 1.7654034 2.1420236 3.1679259 3.9237915 4.5357846 4.5125992 4.5776593 3.7754791 2.3220511 .511476307 -1.3922100 -3.0702968	1.5906480 2.2793165 3.1017699 4.0390573 5.002111 5.8195913 6.2879271 6.2087075 5.4737462 4.1101212 2.2956450	2.1621804 3.0703066 h.1388567 5.388467 6.6778145 7.7341760 8.55778893 8.57772218 8.2170199 7.0524050 5.3915693	2.6027921 3.6749691 4.9223502 6.3300157 7.8007213 9.1516915 10.573613 10.296652 9.334006 7:8552272	3.265 3184 4.5783072 6.0014699 7.7716298 9.5418700 12.530936 13.884396 13.335361 12.657027 11.535792	5.4458802 7.5221264 9.8243854 12.346530 14.592268 17.571830 19.839662 21.565318 22.612661 22.954280 22.875077

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TABLE I. - VALUES OF THE INTRUBAL I, FOR mB = 2 - Concluded

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0.1.25,5750 25,5250 25,5750 25,5750 25,5750 25,5750 25,5750 25,5750 25,5750 25,5750 25,5750 25,555,550	0.46312106 .35001001, .186608701, 32196821 45253002, 965253002, 965253002, 9759159 -1.2519650 91910850 91910850 91910850 1.4550745 1.4550745 2.58601751, 2.1959229 2.58601751, 2.1959229 618266762 618266762 5013039 3010399 -3.01461670	0.85314925 .66352265 .39078019 .01679228 -1.5614253 -1.5014253 -1.501423 -2.3554681 -2.3554681 -2.3554681 -2.3554681 -2.3554681 -2.3554681 1.6586373 2.8861738 3.703822 3.8917756 3.3502961 2.1371109 .1.35502961 2.1371109	1.5397235 1.2397189 .8097189 .82102651 -52102651 -52102651 -2.327186 -3.1651692 -3.1651692 -3.1951651 -3.33958635 -3.33958635 1.0653661 2.8516309 8.3213724 5.0847085	2.8703i23 2.103i369 1.7369075 .86270258 301k2817 -1.6601785 -3.1297716 -1.5313171 -5.6321727 -6.3921803 -6.0257716 -5.0625326 -3.3821186 -1.2133618 1.1093867	4.0532208 3.4670790 2.6397713 1.527329 -1.5238798 -3.4365995 -5.2525992 -5.772268 -7.77268 -7.7247720 -7.7437720	5.9666950 5.2239880 4.1826096 2.788232 1.01292b0 -1.1121083 -3.466922h -5.8388793 -7.9126277 -9.4781236 -10.197714 -9.9755727 -8.8442114	9.1160684 8.15381.09 7.114.1932 5.3324838 3.0707213 -2.6501851 -2.6501851 -2.6801851 -2.6801851 -3.8801851 -11.208248 -12.823040	16,88283 15,58235 13,808995 11,14995 8,459955 1,154595 -7,83509552 -7,835095 -7,835095 -7,835095 -7,835095 -11,478687 -11,786807	21.352313 22.296536 21.154333 15.317095 11.517009 10.566020 5.7721316 5.7221316 -1.3737860 -1.47347142 -9.6685200 -1.3.962138	k1.501668 39.548196 33.43058 29.104865 23.908465 17.953066 11.478609 4.7877470 -1.7460730 -7.8218220	164.18367 160.55029 156.62355 151.08873 114.25586 136.10717 126.73608 116.35810 105.28467 93.857569 82.375422

(1) I3"

41	.4	.5	.6	.7	-75	.8	.85	.9	.925	.95	.985
Bo Wa Ka Wa Ko Wa Ko Wa Ko Wa Ko	0.4h717227 .557.871h .7232077h .8255.8645 .785171h .8555.8645 .785171h .3051695 .2015695h .3051695h .3051695h .3051695h .1.8860187 .1.8764003 .1.8860187 .1.8764023 .1.186921 2.2273039 2.5553342 3.1251000 2.6557106	0.750138b0 1.0333730 1.2657151 1.bb6358b 1.5259779 1.b1422535 1.1351872 22259713 22559713 2136527 2136597 -2.136527 -2.136597 -2.136597 -2.136597 -2.2181695 22181695 35350612 b.266658 3.1136252	1.3072519 1.7331848 2.1250229 2.15902568 2.5253939 2.5538176 2.1713948 1.1036212 .264313427 -3643334276 -3.8735217 -3.8735217 -4.677576 -3.8735217 -4.677576 -3.8735217 -4.973803 -4.1201744 -2.90510593	2.1984602 2.9117451 3.5785883 4.4730304 4.538382 4.5093427 3.2533040 1.768350 14657815 2609684 -4.282140 -5.8179973 -6.6555378 -6.4296553	2.5008368 3.852857 4.7569772 5.5545233 6.1013487 6.28260754 4.9577903 3.3299340 1.1629026 -1.3066002	3.9238036 5.2112727 6.4512763 7.5518845 8.3958621 8.48054353 8.6091412 7.6761985 5.57732043 3.5389766 .79026230 ~2.1209518 ~1.209518	5.5217967 7.3693369 9.1362955 10.752571 12.071152 12.30013 13.034613 12.322050 10.696656 8.2265783 5.1956965	8.4658116 11.292322 14.036765 16.50020 18.466347 21.465312 21.455312 21.19953 20.19953 20.19953 20.19953 20.19953 20.19953 20.19953 20.19953	11.110550 11.80590 18.149000 21.888151 21.973905 21.1473187 29.135010 29.753525 29.23550 21.571251 21.590086	15.813365 21.059813 26.259813 26.259813 35.82080 39.807357 42.506522 44.886683 45.688246 45.111128 45.587515	38.335520 51.037107 63.757103 76.221.973 89.282304 110.11136 119.31826 127.11494 133.57331 138.72353

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NACA TIN 3809

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TABLE II. - VALUES OF THE INTEGRAL I, FOR mB = 3

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(a) I_l'

12	, k	.5	.6	.7	.75	.8	.85	.9	.925	.95	.985
8.4 8 5 15 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6 8 5 16 6	0.032154.37 .0/796502 .06133654 .0522954A 0619714/0 31928896 326896 326896 326896 32767772 .03391791 .56921102 1.0975931 1.0310378 .4406101A 55119789 -1.5210639 -1.5210639 -1.5216639 -1.5215855 2.16789591 2.5175855	0.00408847 .1028861 .1347500 .01115321 ~.2747100 .01115321 ~.27483233 74589587 1.4763005 1.850897 1.4763005 1.855325 1.855325 1.4855325 1.9612102 5508875 1.4165250 2.6771148 2.6101263	0.131148425 .10029672 .55795999 .5575555 .00923903 365595988 95669751 -3.102903 1.3173938 1.3173938 .0138280 2.2263142 2.7338281 2.732621 1.9996893 .55369551 -2.8097385	0.24,376978 .34,42631, .46087835 .4797307 .2227053 -1.34,33065 -2.314,2944 -2.1930332 -1.94837 .65069003 2.229040 1.9671066	0.33188772 .46118671 .65147631 .36655477 35535772 -1.5003829 -2.4794389 -2.4794389 -2.4794389 -2.17062102 -1.8160255 03181474	0-15657674 .%223214 .42326119 .49789012 .98765114 *86014769 -1.*896976 -2.*9777893 -3.2205692 -2.*5568185 7170704 1.1708306 2.*715952	0.64352178 .85537169 1.1145110 1.2161157 .91106223 -0.5057805 -1.5367891 -2.9175652 -3.6476470 -3.0628462 -1.5332101	0.95267785 1.7251790 1.7781319 1.7785422 1.1923571 .15192817 -1.2053237 -2.8853787 -3.7996806 -3.5793019 -2.2765710	1.2015940 1.5131042 1.9138936 2.1757183 1.9228475 .57480637 8442513 -2.6177250 -3.5018438 -2.4647169	1.5893701 1.9521007 2.4149323 2.7487550 2.5167759 1.5140236 23205956 -2.0700601 -3.1730555 -2.3314097	2.93k7915 3.9664965 3.9632810 4.4674512 h.1069080 3.44708066 1.97726857 04776399 -1.1823850 -1.210188 56343880

(ъ) I₁"

K -1	.4	.5	.6	.7	.75	.8	,85	.9	.925	.95	.985
ຬຬຬຆຬຬຬຑຆຬຬຬຬຆຬຬຬໟຏຏຬຬຬຑຑຬຬ	0.00217913 .0129259h .ctild39936 .10550132 .17731701 .00506649 .11506738 .11506738 .12853100 46153668 16223663 .4609257 1.2993378 1.4638121 .8069257 1.2993378 1.4638121 .8069257 1.2993378 1.4638121 .8069257 1.2993378 1.4638121 .8069257 1.2993378 1.4638121 .8069257 1.29530262 1.55577566	0.00124061 .00528257 .06779356 .5577734 .12234130 .71112212 -14310469 -78377067 -1.2774692 -1.2612131 -53350044 .7847768 1.9339619 2.3517761 1.4012696 -2.85726121 -2.5471690 -2.55726121 -2.5471690 -1.0229364 .03329225	0.0073/619 (1376202 1525778 6325678 6325678 6325678 55197219 071%/110 1.010992 101992 101992 101992 101992 101992 1012053 2.1320547 3.0009437 2.5385269 90909091	0.01170869 .05963064 .71360055 .88130116 1.0251111 1.3219271 1.3219271 1.3219271 1.3219271 1.3219271 1.3219271 2.2614018 -2.56140188 -2.56140188 -2.56140188 -2.56140188 -2.5792811	0.014/0641 .00573668 .30050093 .72380173 1.2817707 1.6376561 2.5360458 .57542200 -1.0113887 -2.5685075 -3.2857582	0.01749127 .10117167 .35586373 .6550754 1.55821132 2.1290275 2.0559322 1.0660591 -2.0110269 -3.3479525 -3.0048005 -1.6522509	0.04098767 -12510992 -11031907 1.0702011 1.9324524 2.6579127 2.7151840 1.7606381 02115399 -1.69841149 -2.9995364	0.072192336 .11,870750 .52851,788 1.26512861 2.3371,820 3.3276572 3.55126321 2.7138416 .99984,648 862,28770 -2.0019293	0.02706381 .16155827 .57054154 1.3770567 2.5562943 3.6449585 4.0296443 3.3161521 1.6915655 08466026	0.02932352 .17513612 .61923892 2.8035706 h.0332171 h.5715611 h.0054560 2.5272395 .90397666 ~0301640h	0.03269291 .19511333 .69212751 1.705563k 3.1666687 4.632892k 5.1141277 3.97167882 2.711535 2.1567689

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TABLE II.- VALUES OF THE INTERAL I, FOR mB = 3 - Continued

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27	.4	.,	.6	.7	.13	_8	.85	.9	.925	.95	.985
702557025570255702557025570 7111122225750255702557025570	0.~C61.6771.0 	0.13830183 .12897922 .008993033 03763323 22906191 52931110 52931110 50979025 3659821 .27769314 1.9793913 1.9362598 .913328400 77427911 -2.579592 -3.0670035 360428 56260602 1.5539667 2.94642839	0.288638077 .2725779F .274863552 .07751303 15437786 45572710 -1.15271913 -1.6464161 -1.2164205 05331857 1.5560971 2.8470939 3.1882959 2.27571477 -219512570 -3.28514513	0.59567078 .57002133 .46212267 .16193115 12407892 -2.3810388 -2.300219 -2.8599995 -2.5315467 -1.1386893 .5831376 3.0211396 3.0211396 3.7118321 2.1375417	0.86951923 .87791747 .70125803 .3395838 -111637679 -1.8571880 -2.8072307 -3.6495193 -3.5094779 -2.1387163 .13015204	1, 3037751 1, 2651280 1, 1036666 .65024647 27080592 -1, 656507h -3, 3080120 -4, 5366680 -4, 1177087 -3, 5394662 -1, 3055461 1, 1565360 2, 9444504	2.0518675 2.0058310 1.8713336 1.8750780 .11530225 -1.6010225 -3.6729606 -5.1160885 -6.1020586 -5.3571151 -3.1521150	3.51,201,25 3.4,8734,59 3. 2757058 1. 2562097 - 8,7800390 -3.4,977566 -5.91701.51 -7.3485713 -7.3485713 -7.3485713	1	7.6248195 7.5604289 7.2875013 6.515377 4.9082023 2.3213824 95689360 -h.2513802 -6.742708 -8.0082378 -8.2750198	21.129987 21.358163 21.053355 20.186633 18.376776 15.432252 11.577535 7.5617961 1.1065777 1.625777 1.625777

(a) 12"

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12	.4	.5	.6	.7	.73	.8	.85	.9	.925	.97	.985
2.0 25 5 0 25 5 0 25 5 7 0 25 7 0 25	0.0501.7170 .0801.2192 .12272871 .17119948 .19576852 .14457746 -029711597 317404.11 617501.39 745312552 52327961 .08530669 .88509217 1.46205296 67775399 24118055 24118055 67775399 24118055 721.3005 	*.10860988 .17093011 .55992565 .36291482 .42367133 .3223645 -52296858 -1.1882086 -1.1882086 -1.1892047 -1.189217 -1.189217 -1.18951365 -2.10840513 -2.1084051 -3.1954571 -2.1084051 -3.1954571 -2.1084051 -3.1954571 -2.1084051 -3.1954571 -2.1084051 -3.1954571 -2.1084051 -3.1954571	0.211222800 .33336201 .50145022 .70051866 .83635652 .7105581 .2100531 .21704531 .21704531 .21004531 .2004531 .2	0.k08L1651 .62622877 .92958315 1.\$37716 1.\$798083 1.\$59703 .87180075 k5136665 -2.1216530 -3.k336132 -3.8315042 \$6665130 3.1554097	0.56579647 .85970133 1.7265754 1.7513280 2.1648038 2.1648038 2.1637706 1.48220539 328147967 -2.0448640 3.813246 3.813246 3.813246	0.79463387 1.1946962 1.7379866 2.39967th 2.9880205 3.1367919 2.44677557 .771447689 -1.55005 -1.752005 -1.55005 -1.5500571 -3.7525005 -1.59800671 -3.1152112	1.1477424 1.7039619 2.44459663 3.3486789 4.2018847 4.5779383 4.0058232 2.2825966 287505372 -2.8765303 -k.5477964	1.7528035 2.5620393 3.6111183 4.8866501 6.1553995 6.9341307 6.67714982 5.1193911 2.5439252 ~.27226050 ~2.41782637	2.25b0958 3.2631512 4.5177418 6.1010285 7.46819503 8.7828956 8.8218822 7.5012401 5.0843404 2.3289568 .05933550	3.0533840 4.3699899 6.0013482 7.950785 9.3995699 11.578937 12.084996 9.1616984 6.705985 4.6131054	5.9305802 8.2060942 11.04,8549 14.237196 17.612805 20.624,680 22.623622 23.251447 22.777186 21.753790 21.04,9327

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TABLE 11.- VALUES OF THE INTEGRAL I, FOR all = 3 - Concluded.

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(a) Iz'

12	<u>, ,4</u>	.5	.6	•7	.75	.8	.85	.9	.985	.95	.985
ໞ຺຺ຨຏຑຏຌຏຑຏຬ຺ຨຏຑຏຬ຺ຨຏຑຏຬ຺ ຉຏຏຏຏຬຌຏຑຏຬ຺ຨຏຑຏຬ຺ຨຏຑຏຬ຺ຨ	0.09212288 .0(,389362 -,03138911 -,1118607 -,2821263h -,4105665 .01161705 .53071326 1.0407366 1.2258906 .84511710 -,05101231 -1.2028218 -2.0456516 2.2728893 2.4951509 2.3251050	0.221/2139 .11811110 0k180519 77566943 57566943 57566973 87563577 85369180 28569180 28569180 28569180 2.0327786 .66854092 2.0327786 .66854092 -2.25097863 -3.5271006 -2.6509161 58223033 1.7805466 3.3306483 3.5358970	0.50671731 .30k13977 0557322 k558126 -1.0k17572 -1.6538786 -1.0k17572 -1.6538786 -1.0213267 .58065169 2.1277780 3.48977735 2.53465340 .12333469 -2.45346759 -2.45346759 -2.45346759	1.1826519 .79938575 .22151129 61120057 -1.7007467 -2.8596674 -3.6399913 -1.0342528 -3.0722187 92528483 1.8929230 4.11690016 5.1222036 5.1272519 3.5724489	1.8677359 1.3392544 5695140 5595140 -2.0631109 -3.1296895 -5.1567519 -5.1567519 -5.7145325 -4.97135279 -2.7340829 .46350650	3.0825310 2.314,3153 1.223,1009 -2.3253,080 -2.3280330 -4.65539316 -7.6316,717 -8.0871798 -7.6316,717 -2.40570537 -2.4055537 1.3171,886 4.11477585	5_1828331 4_h233907 2,0800163 -72,1350330 -2,1325620 -8,706102 -8,706102 -11,1869h1 -11,971085 -10,762571 -7,8676177	11.215339 9.6099853 7.3114325 4.1347043 01925100 -1.976468 -10.159752 -11.662400 -17.579383 -18.427377 -18.427377	17,727785 15,673330 12,763094 8,762935 3,6021718 -8,6213545 -9,3042312 -15,509217 -23,650707 -23,650707 -23,6897331	32.086045 29.319118 35.450519 30.229899 13.1,67979 5.3530060 -3.601.8930 -12.1,60104 -26.057653 -30.261974	116.1.12361 133.90055 128.1.6117 112.1.9516 98.1.12250 61.613692 64.524731 107.229239 30.516402 114.526806

(f) I3"

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12	.4	.,	.6	.7	.75	8	.85	.9	.925	.95	.985
5085550855085508550855085508550	0.11006206 .18502516 .22117437 .2117937 .18107584 .03134086 55779134 55779134 55779134 45768922 35818313 .362539999 1.2002833 1.65794624 .14056034 1592565 .14050034 -2.258334 -2.5935653 -1.7044308	0.11370295 .11149852 .49950089 .5266816 .44707196 .16602022 .320339 99096657 -1.5307713 -1.6531460 -1.0696197 .17190611 2.65408795 1.7190611 2.65408795 1.7032706 11747860	0.65391345 	1.3471969 1.7869691 2.1821488 2.43221637 2.3483784 1.6912741 1.6912743 -3.7465378 -3.7465378 -3.7465378 -3.7318687 -3.7318687 9405765 2.0308349 4.2786075	1.9429525 2.6005252 3.1971205 3.602098 2.8832817 1.216543 -1.2219433 -3.9846400 -6.1763955 -6.9636262	2.9438552 3.9112796 4.8116793 5.4912142 5.6696613 4.961706 3.0655989 0.09727170 -3.4544732 -6.4393255 -8.4550068 -8.5506814 -6.93227116	h.6291250 6.1555383 7.6016566 8.7010936 9.3345049 8.8k30631 6.6663058 3.h150273 \$88640800 -5.3750804 -8.7334000	7.5843342 10.685031 13.172248 15.105043 16.902282 17.05914 15.159773 11.889705 6.8339215 1.2227485 -3.9390505	11,203138 14,5125260 21,79396 24,256006 25,28635h 24,256736 21,093420 16,066978 10,100396 4,1453342	17.173123 22.574279 28.562925 33.661.267 38.025598 10.754864 11.336392 39.343763 39.34376553 39.3437655555555555555555555555555555555555	64.236430 64.209475 80.21.6698 95.790158 110.26198 123.01024 133.00024 133.000024 133.00000000000000000000000000000000000

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TABLE ICT. - VALUES OF THE DESIGNAL I, FOR = 4

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(e) I₁'

12	.4	.5	.6	.7	<u>.</u> 75	.8	.85	.9	.925	.95	.985
0.1.1.1.1.0.45.5.75 0.45.5.55 0.45.5.55 0.450000000000	0.00682117 .01190926 .0176381h .0176381h .0176580 .033601h8 11038726 20183873h 20183873h 20183873h 2018587 2015165b 6672587 2015165b 6771587 09781645 097816759 1.16716759 1.16716759 1.16716759 1.16717591a 22808109 28908109 	0.01838656 .07119520 .01698145520 .0316997 -07016378 -3809971 -3809971 -3809971 -380997135 -38097135 -38097135 -38097135 -38097135 -3939510 -393951	0.041368909 .(7288285 .10928373 .08617382 .3005482 .8005482 .8005482 .8005482 .45966921 1.660486 .4560486 .527048 .52704	0.09785398 .15731575 .23577920 .21515902 .11793613 81581555 .12511559 .18672507 2.0650262 2.8118943 1.55876476 8179712 -2.3369458 .1.9680071	0. Ustolski 77 .2291971k .3209662 .3336688k 07571059 9793317k -1.8783325 -1.9790190k 236397136 1.959685k 3.9511851	0.220h7526 .33558001 .19785597 .51557880 .0355317h -1.0966215 -2.3991791 -2.439013 1.4707215 2.0193218 2.3673762 .71747165	0.34215192 .50255039 .7328305 .80201693 .261114778 -1.41026511 -2.6645027 -3.4321578 -1.7715411 .55112393 2.0232126	0.56293125 .78783773 1.118389 1.2806386 .72242393 85429732 -2.7932762 -3.6788767 -2.6595492 -3.6788767 -2.6595492 785162	0.75364012 1.0224427 1.1216636 1.6558772 1.1181426 - 2.638400 - 2.638400 - 3.7548962 - 3.0276153 - 1.3541439 14853110	1.0675639 1.3928923 1.4805289 2.2137968 1.7311478 .03553487 -2.2115695 -3.5169658 -3.5169658 -3.5169658 -3.525358 -1.2363477	2.251,3744 2.70091,90 3.3738767 3.9320764 3.9327764 3.9327764 7.952782 5622122 562014 -1.4621626 60970505 70962078

(b) I₁"

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1,1	, k	.7	.6	.7	51.	.8	.85	• •9	.925	.95	.985
ogwy.ogwy.ogwy.ogywy.og ogwy.ogwy.ogywy.ogywy.og	0.00035723 .00336277 .01516684, .01141491 .04,33925 .03517547 036510165 .2377547 .00623245 .14154063 .13360780 .31462711 .39697467 .1.3212920 5669746 1.7953165 1.795216 1.79524 19511660 -1.5064988	0.00073/9 .008231/2 .03810556 .103079/2 .1055763 10905763 55207222 6598/9,665 11957700 .0511/725 1.52009/1 .50779530 66568078 -2.160692 06565377 2.0639631 .593750.6 -1.2586/21 -2.1596321	0.00181283 .01712242 .07970549 .218571714 .36513076 .238571000 .28005495 .295755697 37555697 37555697 31555697 3165667 2.3487311 2.0038133 05067699 2.309578 2.309578 2.309578	0.00336276 .03183267 .11977861 .12972581 .5972581 .6977514 00101466 -1.16018712 -2.27144687 -1.6149881 .699496 2.7310283 1.0689009 81911479	0.001413396 .012034100 .355610503 .955610503 .95121467 1.0218072 .13446194 -1.1628775 -2.7052850 -2.8412285 13075718	0.00574328 .05453171 .25718775 .7302030 1.3336396 1.360327 -2.5718562 -2.5718562 926312968 .1.3065928 1.3065928 1.7181005	0.00732482 .30595710 .91680599 1.7689178 2.1048078 1.1655311 95127522 -2:8712562 -3.1078560 -1.6461167	0,00921279 .03776056 1.2203685 2.3168420 2.9113102 2.1572560 .00317643 -2.1573688 -2.6647351 -1.1742304	0.01028309 .09904917 4.6795644 1.3624550 2.44m1375 3.462457 2.8266055 .75277927 -1.337370 ~1.9509310 -1.2569314	0.01144697 .10972804. 3.52251417 1.5254651 3.60955330 1.7365969 211464763 77236330 17316603	0.01323485 .12617094 60772019 1.7904895 3.5712813 5.035605 5.0455153 3.537001 2.1218969 1.9982907 2.1218969

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TABLE III.- VALUES OF THE INTEGRAL, $\mathbf{L}_{\mathbf{y}}$ FOR $\mathbf{nB}=\mathbf{k}$ - Continued

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¥1 42	.6	.5	.6	•7	•75	ā,	. ð5	. 9	.925	• ઝ	.985
50 255 550 250 2	0.01.289313 .011.6125 .0131.3706 073.9695 13576519 13576519 3311.0019 3311.0019 3311.0019 3321.2952 .212.2952 .212.2952 .212.2019 99795680 8521.2410 .12596814 1.3681277 1.7311831 .55599115 -1.261.651 -2.4319615 -1.4789901 .26521592	C.03711696 	0.09637728 .08979693 .01641744 -0968004 39548289 75480400 83446330 23734828 1.9554822 .17825163 -2.1677939 -2.1677939 -2.9303788 -2.0300941 .57065736 2.53708014	0.24,3591,39 .23007,292 .15048561, 199684, -1.6260718 -1.6260736 .1.0996426 .83117619 2.9243985 3.55590015 1.977.6052 82845149 -2.8440397 -2.9183542	0.39302607 37518008 26667826 09017015 87257426 -1.9388590 -2.5019389 -1.943920 2.8677011 2.8672899 1.1609254	0.65264533 .62250841 .h2260551 .02115302 -1.0182797 -2.4453439 -3.6001573 -3.7094877 -9.3386808 2.1377732 4.1063875 3.9436189 2.3867526	1.11.94540 1.11.97193 .99865062 .331.0905 -1.0305577 -3.0465888 -1.305577 -2.9654902 -2.9654902 2.6675105	2.2217908 2.1846071 1.9633522 1.1963700 57056880 3.286085 -5.975577 -7.1259042 -5.9944027 -3.4503924 -1.1779688	3. 3348144 3. 2932863 3.0456748 2.1782176 .18395126 -2.9533423 -6.2821150 -8.1086735 -7.74.10530 -5.9548284 -4.2634158	5.5219009 5.1756556 5.1994607 1.9760421 -1.4380923 -5.6327503 -8.3991140 -9.0880197 -8.1593815 -7.9470969	17.968301 17.914787 17.554402 16.461536 13.802637 9.4163919 4.2259250 17692700 -2.9608683 -4.8924088 -7.2968770

(a) I₂

12 1	-4	.5	•6	.7	.75	.8	.85	.9	.925	95	. 985
0.1.4 5750 555750 555750 555750 555750 555750 555750 555750 5555750 5555750 555555	0.01296395 .02207119 .03713679 .05306056 .04505321 05573756 16148141 28005303 20314888 .15503795 .01169778 .15705213 87007208 41169718 87102530 .66366590 1.98642011 1.786638 25545046 -2.511221/6	0.03570001 .055971k7 .059612959 .114017h7 .13369530 02706851 362817h3 5956h19 1.3239089 1.17565h0 1.651358h .17565h0 1.651358h .17565h0 1.651358h .17565h0 -1.0160350 -2.1279763 .02h86905 2.34469431 2.5929931 1.2202953 -1.2029003 -2.1279503	0.05725200 .11404033 .2303291 .34351313 .35912591 .05261691 -1.3560259 -1.3560259 -1.3567065 -32068141 1.5552531 2.8192469 2.1931562 18366826 -2.5816550 -3.18366826	0.22285359 .32731118 .52508292 .76732917 .84970311 .10249377 77528485 -2.2328484 -2.4513583 -1.7017249 .89140116 3.2292813 3.6035774 1.8888630 1503723	0.30980866 .19309125 .77975122 1.1387718 1.315217 .81515780 65014953 -2.5999267 -3.7233624 -2.8199535 19773429	0.16052149 .75272939 1.1698150 2.0317908 1.5128519 1397212 -2.6959009 -1.1996395 -1.1996395 -1.8901045 .828026330 2.3388005	0.77033211 1.1813659 1.8012251 2.5593093 3.1996059 2.8356700 .9167351 -2.1215977 -4.7157036 -5.3312907 -3.9093703	1. 3200605 1. 9838949 2. 9310496 1. 1085078 5. 2527635 5. 2638916 3. \$119561 03347070 -3. 3013444 -5. 0230430 -5. 0263425	1.0123046 2.6668141 3.9041970 5.4795809 6.9563051 7.3310940 5.719286 2.3766343 -1.1569794 -3.358930 -4.158930 -4.1375423	2.6476075 3.8616306 5.4911966 7.575716 9.6413919 10.620850 9.5202828 6.5195130 3.0990897 66221230	5.5636804 8.4202508 11.44548 15.12548 19.023052 21.956229 22.851528 21.723994 20.0131.74 18.968437 18.450548

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TABLE III. - VALUES OF THE DITERAL L, FOR mB = 4 - Concluded.

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K2 K1	-14	.5	.6	•7	.75	.8	.85	.9	.925	.95	.985
50 25 50 20 20 20 20 20 20 20 20 20 20 20 20 20	0.01651833 07057349 07268221 0757349 12805125 12805125 12805125 12805125 17819348 42670328 62063985 12120915 805039500 805039500 80500000	0,05248302 .(0571,270 07782118 19593710 19593710 1925518 1925518 1925518 2812,2730 .30987063 .99751281 1.2798901 .62012717 88055068 -2.1177289 -33131770 2.1168259 3.3000173 2.0373758 78010128 29681916 -2.96255668	0.151,23967 .01,22539 -11,57364,5 8373161, 80010683 +1.0337092 77335500 .2395664, 1,7135350 21322623 -2.6987797 -3.59371,658 -2.5267688 .53371578 3.039921,7	0.45772183 .19758351 2219308 166951829 -1.7056516 -2.3300851 -2.2592306 73828700 1.9263100 1.9263100 1.9263100 2.5648663 2.56486632 3.8260648 -4.5268359	0.31115150 .11963260 20717910 -1.1667076 -2.1258096 -3.5652120 -3.7236792 -2.0926079 1.1996998 1.4196998 1.4396998	1.5169372 .90993350 03764510 -1.4715153 -3.3764180 -5.8206452 -6.0371419 -4.6049305 -88205082 7.3082207 5.5912712	306280277 2	7.1906959 5.5571540 3.1263871 - 5.1051140 -5.1671154 -15.057784 -15.057784 -16.856335 -15.139441 -10.810041 -5.6836857	12.269533 10.042300 6.7804592 2.1090222 -4.1530615 -18.133069 -22.280568 -22.828398 -22.828398 -22.628398 -16.610565	81.1900h6 20.962b91 16.327116 9.8223950 1.17276b0 -9.1021130 -19.362286 -27.16232h -72.261192 -71.053028 -31.034117	128,58088 121,52597 111,52698 81,26598 60,576703 37,997123 15,593651 -5,2557720 -28,681,39 -13,108858

(f) I3"

K2 1	.4	•5	•6	•7	.75	.8	.85	•9	.925	.95	.985
01.025550255502555025550255550255550255550255550255550255550255550255550	0.03899175 .05110966 .05970665 .0512611 .0114098 1063364 21232566 2713309 10163309 10163309 1016337 0688867 5136537 5136537 5136537 5136537 5136537 5136537 5136537 5136537 217638 1.8013804 217638 1.8013804 2173215 -2.7239276	0.112164 .11606951 .17332586 .15853769 .03986201 23165222 45281792 .5005962 1.7517751 .60845118 -1.1638329 -2.8817766 -2.2818129 .18612675 2.7281067 3.3321292 1.5592550 -2.28216732 -2.7921725	0.29071.793 .384,68170 .4576.320 .21041875 3703040 -1.2071167 -1.8076781 -1.4578688 .1556.384 .1556.384 2.3237276 3.5029194 2.47755127 39728755 -3.2137785 -3.5707412 -2.2954036	0.73372619 .97283907 1.1727846 1.2103952 .82718750 25251337 -1.553480571 -3.53480571 -3.53480571 -3.51484057 1.55489459 5.1359429 3.1603037 .11764158	1.1829557 1.5700690 1.9054175 2.0252777 1.5806836 .17855880 -2.1655880 -2.1655880 -2.1655880 -2.1655880 -2.1655880 -2.1655880 -2.1655880 -2.22329220	1.9629577 2.6079407 3.1857035 3.1827035 3.1827035 3.1827035 1.2931673 -1.6366455 -5.4030918 -7.6088762 -5.9588762 -3.7117108 .48538690 3.7011819	3.4387638 4.5704384 5.6181663 6.3016819 6.0245927 4.0152937 -5.0514767 -9.1783745 -10.578722 -9.0202828	6.6734225 8.8829048 10.963990 12.62730 13.030566 11.156045 6.4557439 2731300 -7.0120694 -11.780809 -14.015582	13.33499 13.33499 14.534391 19.22005 20.509806 19.195270 14.534801 7.1942789 - 30715680 - 7.4043714 - 7.4043714 - 7.2654725	16.575692 22.08291 27.153825 32.251659 35.177677 35.736400 32.066015 24.089503 15.980993 7.1627030 ~.98650500	53.516411 71.866978 89.661078 106.78773 122.06984 133.80307 140.62588 142.52756 140.83452 136.89820 130.87657

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TABLE IV.- VALUES OF THE INTEGRAL I, FOR #8 = 6

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(a) I₁'

21	•4	.5	.6	.7	.75	•8	.85	.9	.925	•95	.985
5025550 85550 85570 85570 85570 85570 85570 85570 85570 85570 85570 85570 85570 85570 85570 85570	0.00032613 .000780714 .00131983 00161147 01288490 01140976 .03150473 .03150473 21910796 06073951 .10051189 .5397673 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53975081 53958851 2.1795087 8.155764	0.0011/2502 .0012/2503 .00589191 001/13108 06758194 .001781094 06758194 .02102286 64557275 45572753 .01160057 1.2178239 80600222 20607214 .00023245 -1.9665712 -1.9665712 -1.0691913 1.0017404	0.00512868 .01133611 .02084267 -00578510 -1352357 -22509647 .13957553 .36151206 -1.2183700 -1.2183700 -1.2183700 -1.2183700 -2.255650 -89098000 1.0555050	0.01378064 .03176388 .06509589 .0133080 30192130 30192130 30192130 30192130 30192130 .12135851 2.5310116 2.151068 1.91207494	0.0300342 .05961618 .11216362 .3038259 41829609 29132153 1.7703121 2.0677891 60903480 -2.5150803	0.05473900 .10336450 .19385712 .1285645 52856594 82956570 1.0855207 2.0052709 .30037587 -1.7417236 97186882 53490927	0.10307132 .18277192 .33356707 .29430845 53488578 16896579 .1219932 2.8250371 1.2102785 21394026	0.21947636 .34251329 .59967831 .64107598 -14119976 -2.4667285 -2.6167558 -19212518 1.7996876 1.0644752 .93079619	0.3159021 .42068276 .83002275 .95930320 -219711514 -2.5790961 -3.3522575 -1.0015610 .72014573 .17678267 .63755106	0.51292989 .7/1778559 1.2025148 1.4655863 .22 450188 -2.3581680 -3.6893835 -1.5975369 5994940 -1.3139691 6020200	1.4171019 1.8063998 2.5305233 3.1710204 2.2172507 -2.4111329 -2.4111329 -2.4399274 -1.2599274 -1.2411508 -2.8391194 -2.8391194 -2.7218211

(b) I₁"

						(8) 1	,				
12	.4	.5	.6	.7	.75	.8	.85	.9	.925	.95	.965
๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	0.0000965 .00021174 .00181123 .00558125 02258203 02258203 03137636 .14044131 .0117943 31202801 19286390 19286390 19286390 19286390 1.1055135 .6450395 6550395 0755029 -2.0985207 5014028	0.0003785 .0078825 .00702805 .015k7765 -07k84820 .115k7765 -07k84820 .06340060 .kG14452 -30593162 -30593162 -30593162 -30593162 -30593162 -30593162 -356225677 -2.271058 2.1979570 2.273058 -1.51534b0 -1.0619873 -1.0619873	0.00011293 ,00266805 .02112279 .07116066 .06782729 -1.1768503 -0299967h 1.0129912 .0299967h 1.0129912 .02999677 -1.95677218 -1.35697218 2.14596411 .539028135 -1.7530813 -353908435 -1.7530843 -353908435	0.00028536 .00570162 .05533668 .1390036 .23915215 27763367 58052928 1.5288450 2.0730024 58052973 56250373 56250373 56250373 .0369397 .59759171	0.00113332 .01030116 .06143310 .0261499 .42208965 -1.502705 -1.502705 1.4116331 2.6212215 .27593112	0.00063936 .0152405 .12672588 .12672588 -12672589 -1.981978 -1.981978 2.6419778 2.6419579 2.6419579 2.6419579 2.6419103 ~16550120 ~165119103	0.00092150 .02201/757 .18506001, .699310425 1.1959154 .30911346 -2:0450713 -2:7760475 -2:7760475 1.7735059 1.0039543	0.00130076 .03123126 .26486825 1.0227188 1.1932315 -1.4026593 -3.1216102 -1.4026593 -0.6120430 -557731653	0.00153123 .03691278 .31476415 1.\$111803 2.1294606 1.9049640 95864580 -2.78715500 -1.6236706 859706 1.7590286	0.000.80176 .01313934 .3721.6148 1.1600327 3.0126858 2.3668330 .1256496 -1.0714569 -1.0714569 -1.0714569 -2.1873641	0.00223999 .0511910h .56832375 1.938057h 4.1301282 2.6916125 1.2567630 2.1176855 1.7801598 07536343

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TABLE IV .- VALUES OF THE INTEGRAL I, FOR mB = 6 - Continued

F2 F2	عام	•5	_6	۰،	.15	_8	.85	.9	.925	.95	•9 85
๛๚๚๚๚๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	0,00056369 .000kl752 -,0022269 -,0022269 -,00221667 .0500k356 -,25067569 .00223873 -,27067569 .00223873 -,27067569 .00223873 -,2607569 .00223873 -,2607569 .00223873 -,2607569 .00223873 -,2607569 -,2006423 -,2006423 -,09933082 2,3282231 .85171220	0.00267235 .00222868 ~.00256030 0730313 05952021 02611369 166395392 155395392 13056085 75302367 30058572 1.1581323 1.188080 -1.0095296 -2.1356316 354638662 2.6171682 729312650 -2.1333977 -1.b171899 889690115 1.1394111	0.c1073420 .00940457 -00510592 -07199103 -119517643 -11950704 .87622140 -1.089569 -1.43997031 -1.4390888 1.35997031 -1.4390888 2.7115594 .03906069 -2.6113959 -1.4609299 -395527100	0.04070362 .03733079 .00010671 -1.7690306 -51423628 -60060470 .b7200477 1.9505527 1.3372170 -1.601119 -3.2058536 -4.9030322 2.3336338 1.8938723 .49017958	0.08022835 	0.16339915 .15583318 .07124272 34511653 -1.3021637 -1.3021637 -1.49507888 45911813 2.8523235 3.8735859 3.8735859 -2.2158503 -2.2158503 -2.3900973 -1.9132330	0.35481375 .3438754 .22087072 -35538554 -1.8924570 -3.2725118 2.0277253 4.5761718 3.1260323 .86443866	0.87357810 .05812277 .68280400 21298189 -2.5062509 -5.1640560 -1.07765783 -1.0174650 2.6676628 3.16955919 3.9264405	1.1593664 1.151321 1.2738680 -2.6131360 -6.2116156 -7.0728280 -3.6773779 -3.3370050 1.8593561 3.2160796	2.8961402 2.37145744 2.6292661 1.3761920 -2.0995335 -6.8913105 -9.1163965 -7.6695110 -5.3312842 -1.3762612	12.779596 12.753118 12.145585 10.820228 6.2739707 75171060 6.513061 9.1800951 -11.589798 -11.589798 -16.071200

(c) I₂'

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K2 1	•4	.5	.6	.7	.75	.8	.85	.9	.925	.95	.985
01111120000000000000000000000000000000	0.00078032 .00155392 .00323949 .00394846 ~.0034453 03066667 02480591 .06802082 .13638275 05655071 34701514 .05053221 .4854603 76368283 10837602 .67056968 2.47659605 23346006 -2.4765955 23346006	0.00351754 .00582604 .00337097 -01718465 -1148465 -1148465 -12397278 .18309694 .02439769 .02439769 -06785069 1.2551303 1.6947694 78593323 -2.48993377 18374710 2.1990062 1.1814685 -1.5692631 .2006553	0.01319068 .02176365 .01548397 .070511666 -02273776 -32285528 .05139508 1.00233 .65139508 -1.66551138 -2.1693376 -1.55876732 2.8899561 .834123165 -1.5585615	0.01553337 .06197189 .15753901 .23915195 .05706405 69537121 3391490 24281777 2.214570 2.5786531 76969139 -3.1160220 -1.5370266 .937888355 1.3698524	0.084144826 .11426796 .27747507 .5311831. .21362901 -2.0764131396 -2.07641336 -1.0122547 2.203547 2.5038419 .69177230	0.15958700 .27233426 .16311551 .78021377 .57009682 -2.95032671 1.3050351 3.8502776 2.1427784 .14275060 73011134	0.31474801 .51962432 .50300605 1.4537860 1.3530576 66423681 -3.752077 -k.3411243 36547834 2.3614589 2.8047761	0.67777530 1.0768080 1.7732163 2.7681980 3.1281338 .92257260 -3.4566381 -5.9396710 -4.576585 -2.0385585 -2.0385585 -2.03778582	1.0593331 1.6452756 2.6202637 4.0475915 4.0395768 2.7796375 -5.6407086 -5.5860475 -4.7793075 -k.7793075	1.7965822 2.7176105 4.1573224 6.531067 7.8282574 6.3192067 1.396650 -3.0107144 -4.5820524 -5.9840527 -7.7233127	5.130306 7.6192862 11.037324 15.11395 19.716979 21.029360 18.346921 15.189657 13.454694 10.17897 5.9009733

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TABLE IV. - VALUES OF THE DETERRAL IN FOR MS = 6 - Concluded

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(a) I5'

12	.4	•5	.6	.7	.75	_ 8	.85	•9	.925	.95	.985
0.75 1.25 1.55 1.55 2.25 2.75 3.25 3.25 3.25 3.25 3.25 3.25 5.55 5.5	0.00038280 00136342 00136342 001463166 0076319 01125353 .01581776 05632575 04359376 12513672 22146550 0532524 21551680 056326250 21550400 .02382871 2.5512665 .87601070	0,00186912 00553997 02053120 04675189 055584357 .03912714 .03912714 .03912714 .03912714 34304453 3432452 3432452 09108801 1.1228231 -1.14139565 2.9590641 .88854601 -2.3742113 -1.9053652 .77529166 1.67718470	0.01045657 07065333 16519733 2550438 .01513388 .01513388 .0715875 .95035140 28609726 -2.0869278 -2.0869278 -2.0869278 -2.1158295 1.9082184 3.2366062 .04666480 -3.1466668 -2.1038376 .88031862	0.05557340 03671832 51685571 75926270 381,63930 1.22382890 2.7275820 1.2239239 -2.7765787 -4.2981,091, 460,36871 3.09391,9 3.31,809,7342	0.13078345 03821186 38795110 887951102 -1.4623052 -1.461962 1.95557 4.0221890 3.2229420 -1.5849937 -5.0973468	0.32169750 .00677540 51969690 -1.5151099 -2.6515992 -2.5306454 55076430 5-0744878 6.1093757 -3.3224321 -5.1021784 -4.5913364	0.86333750 .25218780 78596790 -2.5149285 4.800.0650 5.6212933 2.8855957 4.3257201 8.63(3113 7.28(5780 3.1508679	2.71.09055 1.8496009 6389008 -8.6796836 -32.108351 -10.87555 -2.7219928 5.3969786 10.056801 12.167739	5.5168733 3.5208157 - 3.5623530 - 4.5586376 - 11.467012 - 17.6377878 - 12.093687 - 3.3277663 13277663 1356022	13,076384 9,7343094 4,6643351 -3,0675192 -13,865084 -24,990550 -11,020915 -29,469973 -23,523710 -15,778391 -5,3947515	96_5181.64 88.517471 74.543295 54.831345 28.285600 - 22.720404 -57.681319 -79.111038 -79.31387 -109.54601

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K2 K1	.4	.5	.6	•7	.75	•8	.85	.9	.925	.95	.985
0.75 1.25 1.25 1.25 2.25 2.75 3.25 3.25 3.25 3.25 3.25 3.25 3.25 3.2	0.0025h876 .00334563 .00334558 11707624 03120309 0014039h .10144522 .117971174 13146661 36917833 .01013280 .89917833 .01013280 .89912836 .39280291 94379580 1.8239063 1.8239063 1.8239063 1.8239063 1.8239063 2.14668880	0.01207837 .01590103 .01677786 00021316 00530402 1453006 05565412 .32284267 .31673996 1551542 1515142 1515142 1.5187392 1.5187392 15768456 2.632640 1.5319386 15508450 2.632640 1.5319386 16507470 19565401 19565401 19565401 19565401 19565401 19565401 19565401 19565401 19565401 19565401	0.04846710 .05400396 .07056544 17782491 46196139 37233306 .64555200 1.5583395 -2.334265 -2.2392455 -2.44443213 1.0343630 3.5327177 1.2249376 -2.2259249 -2.4023451	0.18358125 .2k322265 .28016784 .173816512 35166512 -1.5645320 .37157718 3.2762818 3.0417101 -1.2268162 -1.5715055 .87555557 2.5732190	0.361660kp .1796533 .56317957 .2629150 30224585 -1.9135587 -2.8700.892 68655390 3.6385825 5.0557612 1.3138220	0.73623020 .97815370 1.1687791 1.0277010 -2.7505121 -14.0939868 -3.1314232 2.427929 6.3271476 4.59552931 1.2505513 -1.8925112	1.5980088 2.125656 2.5780955 2.5%19056 .92691550 -7.559555 -7.65031715 -2.2265734 3.939587 6.7774732	3.9330062 5.2560090 6.1350776 6.9315378 5.0675394 820195010 -8.9595383 -13.712721 -7.2815082 -1.0285693	6.7404792 8.9909265 11.107199 12.387082 10.804213 h.0671212 -6.4256667 -14.96926h -17.914757 -12.944757 -13.877263	13.035251 17.370495 21.550182 21.758966 21.400053 17.713213 5.0356900 -8.0913064 -17.701847 -21.732718 -22.732718	57.512699 76.671108 95.615577 113.36561 126.6009 130.87251 129.9970 112.93006 97.529403 77.529403 77.529403 58.070827

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TABLE V.- VALUES OF THE DIVERSAL I, FOR NB = 8

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F 2 1	-4	•5	-6	.7	.75	.8	.85	.9	.925	.95	.985
50 1.25 1.1.2 2.25 1.1.2 2.25 5.0 2.5 5.0 5.0 5.0 5.0 5.0 5.0 5.0 5	0.0003/016 .0005263 .00059280 00159280 .00155280 .00155816 .01151253 .02158159 .10812730 .02571728 .25912991 .7813509 .31339839 .791392782 .09211123 1.075793 81559036 1.9757196 1.9757196	0.00011571 .00072281 .00072281 00321045 01275367 .01275367 .01275367 .01275367 .110901655 110901655 110901697 1.8103695 87167786 2.1530113 82336328 23356380 .27589809	0.00263271 .701837127 .00309718 .01140168 75772810 .01227271: .221935555 222935812 75141003 .755481137 1.0645642 2.1822120 32813052 2.1822120 3293780 0826252 .04167511	0.00372219 .00799168 .m84b909 02756115 2756115 2756125 .018595433 .018595433 .018795436 188861666 .33766647 2.2957071 75634668 1.21551848 (27072046) 17218821	0.00652132 .01422539 .03803402 ~.37376394 18512265 1.241552265 2.2741285 35995325 1.58995325 1.58991244	0.01428270 .03%72795 	0.(132639320 .04899000 .158373bl .01911625 90113743 -1.2720149 2.1107376 2.1107376 2.1107376 2.1107376 745891420 765891420 765891420	C.0819302 .15148341 -33542557 -26403522 -1.1527491 -2.3735257 -37723748 2.5043316 .49131508 .96324207 -02096971	0.13912320 .24427250 .5249507 .51251431 -1.519691 60273141 1.7180204 .70375508 1.8646536 1.4686536	0.25870569 .h1611835 .79959886 .95170650 -,90016973 -3.4h11601 -1.7581782 .29087812 .29087812 .5697832 1.5898363	0.972995562 1.2169382 1.9601684 2.6116653 .83613280 -2.1263927 -3.1277588 -1.658602 -1.9610851

(b) I1[®]

K2 K1	" L	.5	•6	.7	.75	.8	.85	.9	.925	.95	.985
0.00 1.2555 1.120 2.2555 2.0255 2.05555 2.05555 2.05555 2.05555 2.05555 2.05555 2.05555 2.05555 2.05555 2.05555 2.055555 2.055555 2.055555 2.055555 2.055555 2.055555 2.0555555 2.05555555555	0.0000048 .0001607 .0020658 .00057751. -00150939 -00150939 -00150939 .01765376 -01755781 -01305174 .17939217 -11664222 -31332504 .51210065 .1553216 .51210055 -1.6532205 -1.6552205 -1.6552205 -1.6552205 -1.6552205	0.00000220 .0009760 .00127280 .00765775 -03205136 .0471097h .117581,00 -23056595 -1.9783880 .70561,90 -25012,077 -1.27577802 .1.27577802 3.15385768 -70295778 .32236327 .16077187	0.0000002 .0042412 .00528043 .00976285 .0211943 -14269577 .00610219 .50103518 - 36564716 1.97553160 1.97553160 -2.1926961 05875013 1.9735156 -39513967 -39513967	0.000021451 .0011/302 .07972510 01615686 46430538 1.0598581 1.759809 26536113 .62345580 1.4938500 51211039 36090512 80010513	0.00001309 .00257665 .03612865 .03379163 379901388 37991251 1.8123863 .65771257 2.37846095 36289561	0.0007256 .00835152 .008217515 .278482519 .17759950 -L.0540966 -L.050161 1.9767855 1.6311292 3775519 3775519 3775426 58900741 15081717	0.00011835 .00712592 .103461603 .53213738 -1.45039716 -2.05148669 1.4655336 2.1102961 00495167 .3224,3059	0.00018612 .01135623 .16330838 .83538879 1.23637899 99275208 -2.99313064 3333560 .93255979 .52883329 2.0673199	0.00023376 .01120270 .21253932 1.0063379 1.6306012 -3.06165082 -3.06165082 -3.06165082 -3.06165082 -3.0616508 1.1515251 -3.0610018 1.8117862	0,00029017 .0176623k .26690716 1.40217459 2.651x2012 .58756120 -2.41331255 -1.36040706 -1.51689911 -1.485711k0 .21764067	0.00038859 .02375049 .35394796 1.9399453 4.3365544 3.3135311 .72476514 1.4925547 .33135568 -955388700 -25871261

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TABLE V.- VALUES OF THE INTEGRAL I, FOR mB = 8 - Continued

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						(c) I ₂ '					
K2 K1	•4	.5	•6	.7	.75	.8	.85	.9	.925	.95	.965
50 255 50 250 50 250 20 250 20 250 20 250 20 250 20 250 20 250 20 250 20 250 25	0.0002117 .0070503 000259 00055119 .00695719 .00695719 .00695719 .00695719 .00695719 .00695719 .00695719 .00695719 .3386046 85975230 .23821881 1.209345 1230153 75660006 2.1001024 17196623 -1.7553159	0,00019232 -00013464 -00073641 -00073647 -00073647 -00073647 -16119190 -160790023 -51507104 -00995062 -1.0770103 -6455550 2.3645516 -20750761 -1.7661630 -51433167 -71909900 -34863723	0,00119526 ,10096470 01860209 02922462 017465133 .12460166 71764776 1,1495348 1.0085761 58620037 2,5173315 069864271 -1.5108312 33978173	0.00679561 .00599960 ~.007610364 ~.007610364 .2353977 ~.36803658 ~2,313932 .63734213 ~.313997 ~.313997 ~.313997 ~.0136948 ~.0136968	0.01.636473 0.01497308 01295065 17510919 .17629197 1.7424172 .30060015 -3.06881h3 83228016 2.6438972	0.04068060 .0343524 07912289 32446310 92712487 15637705 1.7728558 -2.5575705 1.7728558 -2.557380 .30458856 1.2903898 2.6725066	0.11002992 .10619820 51148636 -1.7260393 -1.853709 2.73191.7h 3.8714565 17631645 -2.3571838 -2.5160163	0.34330792 .33721332 .21003313 6%682512 -3.071282 -3.7136316 .894639955 4.7278510 3.3683939 2.2765408 14700181	0.67384292 .66623040 .50622520 65666600 -3.8421879 -5.8118577 -1.7112507 3.0182206 3.8764108 5.1304032 3.5149663	1.5130970 1.5086392 1.308k612 -1.7157k70 -8.k191k56 -6.0408kk8 -1.8065799 .k63k9907 1.4k938621 5.9931927	9.0281201; 9.0151220 8.7143203 6.67713203 6.67713203 -8.2589892 -11.874590 -13.408513 -15.032967 -14.0977744

(d) 12⁷

1.2	.4	.5	•6	•7	•75	-3	•85	.9	.925	.95	.985
0.1.1.1.2.2.2.5.5.0.25.5.5.5.	0.00004464 .000205769 .0002769723 002858724 .002858724 .002858726 06567028 .00958720 .178841452 17884120 .178841554 275449554 275449554 -1.70938511 .69350337 -1.2756554 -1.71670556 -1.1352263	0.00032778 .00074627 .00092841 -00092841 -00207996 .009958556 .00107170 -30172960 -09277606 -79817172 -24990163 -1.3137066 1.1761827 -2.185217h -2.185217h -33996696 2.32032h -5.54939785 -1.1966744 .1966744 .356296469	0.001.880.5 .00407989 .007989 .007989 .007989 .007989 .0012877 .01301303 .1301303 .13157592 .171387768 .072203390 1.7237768 .024680 2.163260 2.163260 2.163260 2.163260 2.163260	0.00962217 .0195526 .0590182 17449872 55511882 .255511882 .255678 1.670786 4877095 -2.7916257 -72936352 2.5821343 16308037 -1.6225807 -1.5013712	0.02165209 .01232715 .09655313 -1394330 -23807258 -1.0227272 -01537277 2.5509396 .559467615 -3.1703677 -1.0154220	0.016779935 .09322866 .20205561 .22267049 ~23693196 -1.721236 9250090 2.9609705 2.4080093 -1.907695 2.4080093 -1.907690 -2.1300185 -1.3595356 .30942063	0.12067706 .21509594 .43652228 .73567310 .02367685 -2.7869651 1.9920129 3.8009037 1.871736 ~.06758949	0.32600076 .54832967 1.0223494 1.7570929 1.1350267 -2.4326235 -5.436978 -1.7461666 1.5537287 2.625720 h.4219320	0.57933953 .93406290 1.6674131 2.0345624 2.535521 -2.0345624 -6.3139736 -2.0311620 .20494933 3.9591941	1.1387782 1.787645 2.5429366 4.0830661 5.3769657 .66224200 -5.2827110 -6.1730679 -6.4597853 -5.6152687 -1.7673625	4.5989519 6.7332384 9.6739515 13.626475 18.651013 16.730160 11.388837 8.3003507 3.3452142 -1.7411846 -4.5148548

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TABLE V.- VALUES OF THE INTEVENT I_{ψ} FOR mB = 8 - Concluded.

<u> </u>	_lı	.5	.6	•1	.75	.8	.85	-9	.925	•95	.985
2 50 55 50 85 70 8	-0.00141372 -00077735 -0012822 -0012803 -00761-63 -00761-00761-63 -00761-00760-000-000	-0.10008458 001408 00362343 00390872 .05252789 .01299492 20821166 01104473 592291492 20821166 01104473 200701403 -1.1611183 2091403 159380 1598459 2198459 .77540911 1.219456	-0.40005907 0562761 0513219 01333769 0612362 27153011 68231886 51149881 1.6130113 .86958450 -2.71195875 62370119 3.1025958 .21145751 81406011	0.00333739 	0.01388015 -0.01691711 17813128 43651405 43651405 43651405 43691480 .86005326 2.3992718 25513777 -4.15027176 -1.2166825 3.8691970	0.05398804 -28390459 -34373471 -012498864 -76135286 1.1098864 2.1056111 -2.1056111 -2.1056111 -2.1056111 -2.1056111 -4.7553512 -021287955 3.2850405 5.8320345	0.208371.70 115568160 73647380 -3.0218856 6369700 5.5363114 6.7793128 -5.2668105 -5.2668105 -5.2668105	0.95182600 .10055800 -1.1091273 -4.2025784 -7.5752007 -4.3171215 2.9002532 10.825149 10.3258149 10.3258149 10.3258149 -1.5601382	2. 3108331 .82175050 -1. 7015975 -6. 1903905 -12. 123965 -13. 287506 - h.1211559 7. Juli0629 13. 6034.64 15. 857975 11. 453315	6.7046050 3.33277830 79182600 -8.5731134 -19.153372 -26.380030 -21.153904 -9.0116121 3.3583778 16.138977 23.512726	72.276875 61.019804 44.683772 20.411547 -13.456116 -50.000215 -77.958848 -97.443363 -109.65852 -110.73052 -104.29281

(e) 15,'

(1) I3"

1.2	. h	.5	.6	.7	•75	-8	.85	.9	.925	.95	.985
	0,C(1221.69 -,C(1221.69 -,C(1025)6 -,C(100698 -,C(123399 -,C(123399 -,C(123399 -,C(123399 -,C(123399 -,C(123399 -,C(1233997 -,C(123395 -,C(123397 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395 -,C(123395) -,C(123395 -,C(123395) -,	0.00115736 .00115786 .00115967 00506781 002506781 .0021508 .002506781 .0011232 .0026070 36668847 .01599581 .87303270 1,909381 1.3733270 1,909381 1.3733270 1,909381 1.3733270 1,909381 1.3735287 4839774 48396287 .04875298	0.00718137 .00860713 10878131 10878191 07584385 .00148016 .39361211 -1.0870305 0676713500 2.2057691 .83366090 -3.0919536 37136840 2.71166840 2.71166840 2.71166840 37136640 2.71166840 37136640 2.71166840 37136640 2.71166840 37136640 2.71166840 37136640 2.71166895 38501730	0.04082302 .05103597 .05603597 -06003905 -10977751 -56387477 .81,855302 2.0291631 -1.11,51243 -3,5992502 .96949182 3,7084351 3,7084351 -2,0341270 -2,0341270 -2,5361849	0.09827055 .13033696 .1303596 02153876 7233311 -1.800874 .7306785 3.5519077 .39836035 -1.5667351 -2.0159627	0.24,51,2061 .32606591 .31373121,27 .11296529 -1.1624561 -2.7003483 .2,7003483 .2,7005660 1.,5095660 1.,5095630 -2.7262130 -4.6123598 -2.3965397 1.1213789	0.460k1230 .8766k900 1.0k26957 -685267k0 -1.5318089 -5.1222053 -3.979800k7 3.828010k 7.5117967 1.922319 79263760	2.0602275 2.7139500 3.3070631 49202250 -8.0222817 -11.577324 -5.0161908 3.5071524 8.9737152 11.371389	4.ch35233 5.967599h 6.6225366 6.7658260 2.5580372 -7.6895317 -16.328319 -11.333913 -7.16071/76 1.7819630 11.218760	9,1091553 12,110663 15,036647 14,551206 12,338368 -1,280782 -16,817916 -24,817916 -25,682133 -20,887113 -9,5513238	54.169392 72.213629 90.063399 105.05020 105.365726 105.365714 84.1677562 58.757140 88.297859 -3.4657340 -32.515404

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TABLE VI.- VALUES OF THE DECEMBRAL L, FOR mB = 9

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F 2	•h	. 5	.6	.7	.75	.8	.85	.9	.92 5	.95	.985
0.75 1.25 1.55 1.25 1.55 1.0 2.25 2.55 3.0 2.25 3.0 3.55 3.75 4.55 5.55 4.55 5.55 5.55 5.55 5.55 5	0.00001551 .00001551 .00001551 .00001551 .00001551 .00001551 .0002399 .00016861 .0012518 .0012518 .0012518 .0012518 .0012518 .002519 .002519 .002519 .002519 .002519 .002195 .002195 .002195	0.00003315 .(C012157 .C012157 .C012157 .C0194556 .(C11575h .C0465342 .13511729 .13511729 .1351172 .135157 .5399205 .3991205 .71282951 1.352157h .71282951 .352157h .2.0420333 .653110059 .86705640 	0.00022358 .007/1786 .00173221 00500513 03051036 .08106937 .10141359 16505105 .17013715 1.0101174 14505105 1.0101174 46505108 2.0855984 5932668 .71253564 .71253564 .31147537	0.00129953 .00386511 .00983017 -02862159 .11711940 19121987 -63759866 2.2332213 64570916 -1.1649012 .63347250 .83347250 .83347250	0.00306892 .00858288 .02221759 010817671 29158350 1.1139575 -1.0609006 -1.\$261637 2.0869668 .38560068	0.00738099 .0182532 .04948115 04918769 53570496 .01565997 1.79554736 53537266 -2.2531227 1.0591800 27508442	0.01858534 .01271131 .110759 0183845 90612720 2.3185098 .63181659 -1.4301848 43108671 96099688	0.05184854 .1045289 .5326661 .1182293 -1.3325343 -1.7105104 1.9083428 .106452 .10651876 .29101555 -1.9102727	0.09341538 .173814178 .39759995 .34783284 -2.5308206 1.0153187 1.6033512 .99856893 1.5646555 77582238	0.18585218 .31349710 .557677777 .75155120 -1.5912910 -3.3315549 98120288 .33010277 .71361324 2.1700501 .76342779	0.76110193 1.0149164 1.7471506 2.376781 1.6787060 -3.0019673 -1.6839355 -2.3715604 -2.4819986 -1.3059126 -1.3059126

(ъ) I₁*

12	5	. 4	•5	.6	•7	.75	.8	.85	.9	.925	.%	•985
		-0.0000151 .0000510 .0000510 .0000515 .0006372 .00066372 .00066375 .00066375 .00066375 .00066375 .00066375 .00066370 .00066370 .15556121 .000657072 .15556121 .76382059 1.76718019 -1.1285140 743451130 743451140 743451140	-0.00000126 .00003214 .00053263 .00146738 -00671661 -00971661 -00971661 -00971661 -009716141 -18489730 -35098187 -02718343 -85702523 1.0366838 .22220040 -1.7931152 1.6366838 .22220040 -1.7931152 1.63668316 -30659116 -1.6907078 1.0301778 .19011071 -,71021138	0.0000056 .0021266 .0021266 .0058899 02801362 06887583 .21860192 .0650622 180505270 .65115969 .88700273 .240137716 .69595612 1.1690081 -1.1057627 00613559	0_00000681 .00230701 .01230701 .0125070 .0125075 -33557915 .18690893 .7647502 -1.720615 -0599651 2.1698671 -1.705281 -35581173 .70954118	0.0001370 .0025331, .02359317 .10165313 07746516 6093392 .k9363481 1.5096806 -1.7370602 -1.042145 1.6906462	0.00002510 .01233605 .04351286 .2014.0903 -1.0963165 .16071517 2.3200517 -9.1174930 -1.5205316 .121725316 .12172556	0.0001521 .00107851 .9753975 .9553908 -1.6155265 76019930 2.5888619 .1233659 -1235869 .1233659 1233659 .1233659	0.0007678 .00689179 .13827753 .%157073 .%3556860 -1.8286298 -2.3986718 1.2657877 .72308291 1.558667 .72308291 1.558667 .65560083	0.000972k .008367k .17490845 1.0043662 1.1306857 -1.5187200 -2.7492545 .0599697k 29457288 1.48004k 1.3732470	0.00012416 .01133979 .22656122 1.3500866 2.3225502 -2.6371318 90977244 .0164532 .36670538	0.00017501 .01583912 .32166723 2.0239001 h.3050h75 2.3766603 .h0966762 1.1732803 -1.1223046 54127935 -1.7122966h

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TABLE VI VALU	18 OF 1906 DEL	RORAL L, FOR	2 ¥8 ≓ '	9 - Continued
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(c) I₂'

K2 K1	.4	.5	.6	.7	.75	.8	.85	.9	.925	-95	.985
5.1.1.1.1.2.2.5.5.5.5.5.5.5.5.5.5.5.5.5.	0.0000577 .00025777 -00036977 -00036977 -00125377 -00197970 -011327523 -011327523 -011327523 -011327625 -12128016 -01137655 -12128016 -011333360 -1.1217723 1.5442215 1.5442215 -1.7669517 -1.6411283 2.0217719 6442558	0.00005135 .00003527 00022221 00300285 00002576 .0262998h 1121149 .30453807 .10951473 99456669 .59184555 -1.7963263 .89776930 1.9847552 -2.2405250 .55696135 1.193491h 73216042 52407712	0.20039715 .00030510 00211979 01281047 12737786 .01735131 51027522 .h1387018 .9985640 -1.7107691 118577790 2.1.120405 -1.1873759 -1.1875759 -1.1855759 -1.1875759 -1.1975759 -1.1975759 -1.1975759 -1.1975759 -1.1975759 -1.1975759 -1.1	0.00277236 .00239223 -00770431 -076464 -11444107 .50050403 -14760486 45392012 2.7864984 70206516 -2.1577602 .1577602 .1577602 1.0610955 1.3106752	0.00736132 .00665721 01253095 11976106 2657137 .58070470 1.2201220 -1.7375404 -1.8217371 2.7736979 1.1192863	0.02013796 	0.06125758 .05900203 003005051 1.5301945 -1.153576 .07.619201 3.7277125 -2.9772269 -2.1115172 -1.3090372	0.21519182 .21138267 .171666336 77864159 -2.9388099 -2.9711412 3.5725323 4.3229436 1.5102363 45839012 -3.5765369	0.451701169 .54605544 .30670760 806217508 -4.02714542 -4.5777521 1.4883974 1.36214762 4.2797888 3.5247323 63321945	1.0990559 1.0990643 .9167610 63736390 -7.633736390 -7.633728 -3.0071111 .7312260 3.7223951 6.166668 5.2820797	7.5836291 7.5745065 7.3237126 5.0539732 -2.1151788 -10.828547 -13.071788 -15.260850 -15.360852 -12.985981 -11.043181

(d) 1₂

R ₂	ન્ધ	•5	•6	.7	.75	.8	.85	.9	.925	•95	.965
0.1.0 225 150 255 150	0CCCTORE 39 CCCCR 39 CCCCR 39 CCCCR 39 CCCCR 39 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCCR 30 CCCR 30 C	0.00019645 .00024636 .000395 .000395 .003955 .00221965 .6030552 .003552 .003552 .003552 .003552 .003552 .003552 .003522 .003522 .003525 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0033552 .0035552 .005552 .005552 .005552 .005552 .005552 .005552 .0055552 .005552 .0055	0.00069956 .00164459 .00164459 .002744060 .28521353 -09376356 -83304604 1.0023771 .8023771 .77566866 1.8645530 56516203 .579890 56516203 .29619247	0.00136899 .00523365 .02523365 .2312355 1837905 285571515 .32396710 .62067108 -2.2385680 .766206318 2.7708859 95308280 50596775 1.3303642	0.01082766 .02211855 .05700655 .06998979 30596935 68783611 1.0510167 1.6684576 -2.4312719 -1.5379610 2.2216034	0.027h6559 .05k01218 .12926922 .15205791 136158k4 -1.162631 3.1672637 -1.198k143 -2.7005251 -2.7005251 k2973315 .96022250 2.7121137	0.07371493 .13681332 .3085326 .5085124 42532127 -2.6705964 579941390 3.9875735 1.4682030 890730 850730 850730	0.22301245 .38710568 .77153450 1.5726128 .28300529 -3.9244422 -3.94171269 1.4639670 2.914901 3.5225617 1.7881473	0.4225678 .70577124 1.3189311 2.3430683 1.4422743 -3.8703776 -5.9843962 -1.9337501 .28766805 3.8400009 h.5228162	0.49361600 1.1299290 2.4711953 h.2707095 h.0631692 -2.013161 -5.4305113 -5.4305139 -5.4310162 -1.4661257 1.5897560	4,1630867 6,1513702 9,1963017 11,055140 13,860969 8,0792751 4,3212070 -2,0621116 -5,61366906 -9,7613532

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TABLE VI.- VALUES OF THE INTEGRAL I, FOR ME = 9 - Concluded

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2	. 4	۶.	.6	.7	.75	.8	•æ	.9	.925	•¥	.985
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(e) I3,

(f) I,

F 1 F 2	- k	•2	•6	.7	-75	•8	.85	.9	.925	.95	.985
៰៓៲៲៲៲៲៵៵៵៵៳៷៰៰ៜ៴៶៵៰ៜ៴៵៵៰	0.11172104 011/12203 0012924 .00136138 01261159 .00136138 01261159 .0013775 02392201 .0013677 12085493 0694113677 12085493 06941136 303201659 1.10287718 63493018 1.972956 3397488 8014200 2.2691089	C_CCC014552 .00045509 .00022597 0031600 01001957 .01766853 .0556.1386 12837999 095552925 .49857007 29231829 82837707 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 32337817 3237825 32378577857857857857857857857777785785777777	0.00250951 .00353927 .00353927 .01592005 -04352047 .01391949 .33365225 -31067506 -31067506 -31067507 -2.7900027 .77638446 2.3925852 -1.77743587 -1.2181466 .69261724	0.01075271 .02860295 .02386763 -05878237 -31762973 -11503521 1.2612863 31512850 -2.9806155 .3621113 3.7368317 98203366 -2.5506650 78185230 2.3890963	C.04950728 .06604501 .06844764 6798228 64347920 57316310 1.9399486 1.8418609 -2.9135782 3.2266886	0.1395551 .15339541 .20330542 -2053064226 -1.2133516 -1.7736222 2.1407062 14.6026570 -1.5296310 -1.5296316 2.5576077 1.3144549	0.41364770 .55035570 .64234860 .2228630 -2.0290746 -4.4541294 .1.070430 7.1334303 4.1340572 -1.929747 -5.9456852	1.4527999 1.9350868 2.3410992 1.8433449 -9.22718034 -9.2230307 -7.9296810 2.4426392 8.66944165 10.270926	3de22845 40533813 4.7754950 55563950 -11.440652 -7.4932362 2.0016519 11.4573320 11.4573320 11.457372011	7.1130301 9.8888012 12.231933 13.021007 6.9063601 -9.0415357 -22.230966 -24.075027 -7.2311311 6.0503198	51.150329 68.22936 95.129960 99.77283 103.75358 88.000534 60.21,3756 28.81,2766 -6.600,68226 -65.762724

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TABLE VII. -- VALUES OF THE INTEGRAL I_{μ} FOR nB = 12

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(a) I₁'

L 2	•4	.5	.6	•7	•75	.8	.85	.9	.925	•95	.985
0.10 1.255 150 255 150	0.00000123 0000140 .00007151 00007151 00007151 00077102 0006400 .00377381 01211006 0.0677352 01977159 .06677352 01977159 .06677352 011252288 03078183 011457736 11457736 11457376 1145032 -1.7881343 1.4673072 -1.28090545	0,0000205 .0000205 .0000205 .000571.79 .0125907 .0197206 .0197200 .116093 .1990530 .02027271 .0395550 -1.014617 1.014617 1.014538 .1.955777 1.6856211 .1.955777 1.6856211 .1.955728 .1.955728 .1.955728	0.99001182 .07005380 .00011770 -00953118 .00112132 .019651141 -11258200 .29538911 -112500565 .25106257 .38412811 -1.2700216 1.8700216 1.8700216 .25103029 -1.5311310 .554134303 .1349997 61191029	0.0000621 .00015373 .00157351 -00128568 .38658264 -55316638 .06581 -002553586 -1.380700 1.7145702 53350307 -45131620 .77703125 .02861045	0.00032760 .00127033 .00143871 -03387994 -0148671 -14159158 -87045700 -1205700 -1205700 -1205700 -120570053 .76190995	0_00104587 _00356799 _00356799 _056714 _056714 _0769235 _9230143 _56240027 1_2036192 _72548191 _8755296 _56992401	0.00353305 .0036820 .0377845 1018662 53455266 1.3332142 1485266 .3332142 14856264 .65165504 .65165504 66578465	0.0135458 .03325191 .11170161 06612771 -1.2667201 1.37647792 1.5761365 -1.130664 -1.0775531 32673300 .41396388	0.02911599 .061.3919 .19682175 .01108693 1.0137119 .113897h7 2.1238525 .9722730 42571461 -1.075002 -1.22501257	0.07098013 .1373002 .3750548 .28813995 -2.3047336 97539920 1.5235903 1.5235903 1.5235903 1.5235903 1.5235903 1.5255870 -1.3785603	0.12877715 .43536978 1.2141251 1.7803569 -1.6353328 -2.7853886 -1.6373328 -2.7853886 -1.637397070 -1.7899872 -1.0313979 .28635916 .011855998

(b)	Ľ1
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R.2 1	-4	.5	•5	•7	.75	.8	.85	.9	.925	.95	.985
0.10 1.255 1.255 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.257 1.20 2.255 1.20 2.555 1.50 2.555 1.50 2.555 1.50 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.5555 2.55555 2.55555 2.5555 2.55555 2.55555 2.55555 2.555	0,0000226 .000011,4 -,0000202 -,000008 -,0000284 -,00026449 .00021284 -,0013085 ,00271288 -,00270606 -,00725693 .03678848 -,10021013 .21062086 -,36014781 .51925569 -,62917031 .61330465 -,12971916 .05669238 .13403656 -,05569238		-0.0000845 .0001635 .00037268 -0330028 -04824707 -07152004 1.4214132 -1.0232455 1.567665 -1.2148695 -1.2148695 -1.2148695 -1.2148695 -1.218695		-0.0000638 .0016306 .00650232 .00356377 -15613692 .11456110 .5249080 -1.66538120 75106499 37434542	-0.00001481 .00036010 .00147045 .07141711 2900783 03811421 1.3587761 -2.1737959 .977554707 .h1927214 .36772943 - 1.6204617 .61286380	0.0000119 .0076466 .03553155 1280273 12537599 2537599 2537599 2537599 12577595 14546047 16346047 1.7339149	0.00011136 .00156063 .06812689 80647853 30166805 -2.00123285 2.0511564 .83713971 .56912072 -1.8102371 324155342	0.00001711 .00219559 .09751501 .06571367 - 2.1188136 .9695013 1.8676973 512662 52352361	0.0002776 .01305520 1.1505781 .95278150 5701765 8701765 1.6570166 .85512552 1.26778473	0.60005019 .0048143; .22388449 2.0778400 3.5937100 11514787 .11324885 1.9309199 -1.0367273 -2.2109988 -1.5289023

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TABLE VII. - VALUES OF THE INTEGRAL I, FOR nB = 12 - Continued

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- (a)	Lo'
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F 2	•4	-5	.6	.7	.75	.8	.55	.9	.925	.95	-985
0.10 1.25 1.575 2.25 2.75 3.25 3.25 3.25 5.05 5.55	0,0000555 -,00000548 -,00000913 ,00010563 -,00010913 ,00010563 -,00029339 ,00063057 -,0063056 ,030152169 -,0566308 ,03219177 ,11550356 -30207300 -1,3765551 1,7556366 -1,256359	0.0000866 0000327 000327 0003297 0003186 .01151625 00673816 .04618302 13228208 13228208 13228208 13228208 1587515182 99558986 1.7212268 99558986 1.7212268 99558986 1.8306030 98933386 01645281 .62945281 .52979167 1391678	0_0002181, _00001246 -00036981, _00020691, _00020091, _10507619 _35025704, _55025704, _55025704, _55025704, _55025704, _1050723, _1050723, _117833773, _112051, _117833773, _112051, _1178357, _112055, _91279655	0,00018950 .00015918 00236836 .01178956 .11262961 1.0637303 322113653 1.003735 2.3350469 932869 932869 2.3350469 .2566912 .28453778	0.00067276 .00059589 ~.00597179 .06221105 .k5696830 -1.2776672 .90757611 1.1636291 ~2.511195 .93196k77	0.00254007 .00234830 1111679 11641029 00922628 1.1552136 6697517 56917904 2.7817797 56917904 2.7817797 98700374 552721146 1.65463427	0.01055608 .01010995 01893531 31395312 2.2551240 66034108 -3.0611167 1.27147711 2.366699 .10400979	0.05295900 .05203997 00628086 7553950 -1.6213099 2.7217356 3.066287 -2.1123186 -3.2376274 -1.167928 1.8090019	0.13601290 .13471243 .04002759 -1.0567880 -3.3364296 1.4805926 5.017344 1.6522638 -1.9751008 -3.884693 -3.884693 -2.9735653	0.k1698850 .k1516120 .29199420 -1.3566911 -5.7275255 -2.3282653 3.9650612 5.4759120 4.87219211 .666095519 3.69005559	h.1936134 h.19001(7 h.2982,77 1.1,872552 3.4002610 3.598632 13.718894 13.71638 7.3376638 7.3376638 7.3376638 7.3376638 7.3376638 7.3376638

(d) I₂"

¥2	•4	.5	.6	•7	.75	.8	.85	.9	, 925	.95	.985
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TABLE VII. - VALUES OF THE INFRIENT I, FOR nB = 12 - Concluded

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(a) I.5'

1.2 L	-h	.5	.6	.7	.75	.8	.85	.9	.925	•95	•985
500 1.255 1.2023 2.570 2.575 2.222 2.575 2.225 2.575 2.225 2.575 2.225 2.575 2.225 2.575 2.225 2.575 2.255 2.5755 2.575 2.575 2.575 2.575 2.575 2.575 2.5755 2.575 2.575 2.575	0.00150211. 01187747 00193758 00193758 00190787 00193751 00196771 00196771 00196771 06967775 11826617 11826017 11860017 118600	9(1000688 (100800- (100800- (100800- (100800- (100800- (100800- (100800- (10080	-0.00013785 00022902 00057385 00069023 01549417 02737621 07131660 .81754297 703796674 .5057557 -1.5057557 -2.1641889 351759285 -2.1641889 .351759885 -1.4149548	- 0.00013850 00237917 00228678 0129291 .11709875 .02824878 501875286 5785286 5785286 78599795 -1.6890336 3.218*644 -1.45(12213) -1.6512213 67191637	-0.00116686 00737862 02752227 06176420 .28136933 .35822134 -1.8099512 1.7015127 -3.754.0366 1.2674961	- 0.002h3h09 02280195 06207(686 207(66180 39636390 1.39903h3 - 2.5795422 43785270 h.5592164 - 3.476536h8 .89319308 3.6445408	0.0015630 07110260 5380360 .2219610 3.789003 -1.6931689 -5.9956990 2.278793h 5.1111336 .h2721030	0.07151210 21608620 2503020 -2.3631837 -2.0075381 6.372631 -1.6398831 -3.9670186 -3.3924589 5.5777822	0.31257540 32942620 -1.6061998 -1.5043531 6.0068759 5.1402409 13.232364 5.1990945 -6.3180202 -2.2.76754 -8.9586036	1.5230303 11738420 - 3.1208991 - 9.4270711 - 15.960321 - 5.0153981 13.394711 21.056331 14.370063 2.1557811 - 13.187817	36.589500 25.505754 8.1090100 - 20.079508 - 61.121294 - 98.269229 - 98.167711 - 81.569588 - 53.167715 - 116.730129 20.381027

(f) I3"

¥2 1	. lı	-5	.6	.7	.75	•8	. 85	•9	.925	.95	•985
0.75 1.25 1.75 2.0 2.25 2.75 3.0 3.55 3.5 5.25 4.5 5.25 5.5,75 60	0.c0025279 00057106 00057106 00057106 00057177 00350134 0014505 01145276 01145276 01145276 01145276 01145276 01145276 01145276 059277 22400977 22400977 22400977 22400977 2529780 0044096 0058251	0.0000197 .0002188 0012785 0012785 .00051844 01453619 0233824 054531404 054531404 055517804 3592632 1.3713109 30159827 74425359 1.320712 -1.0036132 05955314 0595952 -1.8329635	0.00012161 .00018750 0009187 0.0053873665 105387361 .153518333 .264381551 -1.1719209 2.0607857 -2.1250818 .50669322 .80310610 -1.5314133 .13368220 1.1390602	0.001.67876 .00224931 .0007h731 -0075568990 .30083175 1530314 72651.655 2.584.0512 2.591.54.14 .571.228336 1.931.1278 1.201.7791 1.84.92310 2.71.3060h	0.00609266 .0080911h .00528738 07148610 20948811 .77709330 0586068 2.1607686 3.6483766 7070139h 2.5706113	0.02298114 .03051957 .02721702 ~.66722710 1.1855179 1.66722710 1.65012779 ~ b.032262 1.3653015 3.45301840 ~.53340259 ~ b.3466344 .86498473	0.0951.6060 .12657270 .33474900 2795870 3795870 3.1377470 5.5227023 3.13774870 5.5227023 3.1377489 48439517 33824060 6.0384738	0.17705170 .63512010 .715511870 07222110 - 1.14575291 - 1.4607315 7.1575590 7.1575590 7.1575590 - 23682360 - 8.9160060 - 8.9160060	1.228.6394 1.6319507 1.9711995 .95229210 - 6.0913481 -11.238222 1.557117k 11.057547 11.425315 .50865240 - 9.1252885	3, 7535959 5, 0135151 6, 1383368 5, 3437057 - 5, 3391347 - 20, 553731 - 15, 064872 - 2, 3780143 14, 371915 20, 751650 16, 742148	10.146386 53.526217 67.255203 77.31548 68.302223 31.573977 - 9.325860 - b9.636381 - 77.338006 - 91.88111 - 89.289720

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Figure 1.- Propeller disk and coordinate system.



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(a) $c_n = 0.029$; J = 2.55; $M_r = 0.8$; n = 1,500 rpm; $\beta_{0.75R} = 45^{\circ}$.

Figure 3.- Chordwise normal-force distribution at 0.78R for two sets of operating conditions (ref. 4). (Numbers indicate rectangles and triangles used for the analysis.)







(a) $c_n = 0.029$; J = 2.55; $M_r = 0.8$; n = 1,500 rpm; $\beta_{0.75R} = 45^\circ$.

Figure 4.- Amplitude functions associated with the normal-force distributions in figure 3 plotted as functions of mB compared with amplitude functions associated with uniform distributions of normal force. b/r = 0.1715.



(b) $c_n = 0.9529$; J = 1.009; $M_r = 0.789$; n = 2,000 rpm; $\beta_{0.75R} = 30^{\circ}$. Figure 4.- Concluded.



Figure 5.- Sketch of radial thrust or torque distribution.







(b) Efficiency.

(c) Torque.





Radial distribution of thrust and torque (per blade)



Figure 7.- Comparison of sound-pressure calculations for forces concentrated at an effective radius with calculations for forces distributed along the propeller radius.



Figure 8.- Distribution of the root-mean-square pressures for the first harmonic of a four-blade 16-foot-diameter propeller at two altitudes and several tip Mach numbers.

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Figure 9.- Calculated root-mean-square pressures associated with first harmonic of a four-blade 16-foot-diameter propeller at a point 0.15D behind the plane of rotation as functions of flight Mach number for three different altitudes.

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(a) Sea-level altitude.

Figure 10.- Calculated root-mean-square pressures for various values of of mB in the plane of rotation for 16-foot-diameter propellers as functions of flight Mach number for two different altitudes.

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(b) 35,000-foot altitude.

Figure 10. - Concluded.

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Figure 11.- Calculated root-mean-square pressure as a function of distance from propeller tip along a line 0.15D behind the plane of rotation for a four-blade 16-foot-diameter propeller operating at sea-level conditions.