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WIND-TUNNEL INVESTIGATION OF EFFECT OF SWEEP ON ROLLING DERIVATIVES AT ANGLES OF ATTACK UP TO $13^{\circ}$ AND AT HIGH SUBSONIC MACH NUMBERS, INCLUDING A SEMIEMPIRICAL METHOD OF ESTIMATING THE ROLLING DERIVATIVES

By James W. Wiggins
Langley Aeronautical Laboratory Langley Field, Va.


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# WIND-TUNNEL INVESTIGATION OF EFFECT OF SWEEP ON ROLLING DERIVATIVES AT ANGLES OF ATTACK UP TO $13^{\circ}$ AND AT HIGH <br> SUBSONIC MACH NUMBERS, INCLUDING A SEMIEMPIRICAL METHOD OF ESTIMATING THE ROLLING DERIVATIVES ${ }^{1}$ 

By James W. Wiggins

SUMMARY

An investigation was performed in the Langley high-speed 7- by l0-foot tunnel in order to determine the rolling derivatives for swept-wing-body configurations at angles of attack from $0^{\circ}$ to $13^{\circ}$ and at high subsonic Mach numbers. The wings had sweep angles of $3.60,32.6^{\circ}$, $45^{\circ}$, and $60^{\circ}$ at the quarter-chord line, an aspect ratio of 4 , a taper ratio of 0.6 , and an NACA 65 A 006 airfoil section parallel to the free stream. The results indicate a reduction in the damping-in-roll derivative ${ }^{C_{l p}}$ at the higher test angles of attack. Of the wings tested, instability of the damping-in-roll derivative ${ }^{C} l_{p}$ was experienced over the largest ranges of angle of attack and Mach number for the $32.6^{\circ}$ sweptback wing.

In general, the variation of the damping-in-roll derivative $C_{l_{p}}$ with sweep angle at zero angle of attack was only in fair agreement with the predicted variation, inasmuch as the $32.6^{\circ}$ sweptback wing showed more damping in roll at zero angle of attack in the Mach number range from 0.85 to 0.93 than any of the other plan forms. The predicted variation of $C_{l_{p}}$ at zero angle of attack with Mach number was in good agreement with the experimental trend to the critical Mach number. Contrary to predictions based on potential-flow theory, the yawing moment due to rolling was positive and the lateral force due to rolling was negative at the higher angles of attack throughout the range of Mach number for all configurations of the investigation. Presented herein is a method of estimating yawing moment due to rolling and lateral force due to rolling through the angle-of-attack range. The method is shown to be applicable over large ranges of leading-edge radii, wing thickness, and Mach number. The results indicate a loss of wing-tip suction within the ranges of Mach number and angle of attack investigated.
$l_{\text {Supersedes }}$ recently declassified NACA Research Memorandum L54C26 by James W. Wiggins, 1954.

## INIRODUCTION

The present investigation is a continuation of a program being conducted in the Langley high-speed 7- by l0-foot tunnel in order to determine the effects of wing geometry and angle of attack on rolling stability derivatives at high subsonic speeds. Reported herein are results on the effect of sweep angle and angle of attack on the rolling derivatives for a body in combination with various wings. The wings tested had sweep angles of $3.6^{\circ}, 32.6^{\circ}, 45^{\circ}$, and $60^{\circ}$ at the quarter-chord line and had a taper ratio of 0.6 , an aspect ratio of 4 , and an NACA 65 AO 06 airfoil section parallel to the free stream. Tests were also conducted on the $45^{\circ}$ sweptback wing with wing fences located at the 65 -percent-semispan station, since an appreciable loss of damping in roll $C_{l_{p}}$ was noted at the higher test angles of attack for the clean-wing configuration.

The longitudinal and lateral stability characteristics of the wingbody and body-alone configurations are presented in references 1 to 3 . The wing geometry is designated as in reference 2. For example, the designation 3.6-4-.6-006 denotes a wing with the quarter-chord line swept back 3.60 with an aspect ratio of 4 , a taper ratio of 0.6 , and a 6 -percentthick airfoil section with zero camber.

## COEF'FICIENTS AND SYMBOLS

The stability system of axes used for the presentation of the results, together with an indication of the positive forces, moments, velocities, and angles, is presented in figure l. All moments are referred to the projection of the quarter-chord point of the wing mean aerodynamic chord on the fuselage center line.
$C_{2}$ rolling-moment coefficient, $\frac{\text { Rolling moment }}{q S b}$
$\mathrm{C}_{\mathrm{n}}$ yawing-moment coefficient, Yawing moment
qSb
$C_{Y} \quad$ lateral-force coefficient, $\frac{\text { Lateral force }}{q S}$
$C_{D}$ drag coefficient, Drag/qS
$C_{L} \quad$ lift coefficient, Lift/qS
$q$ dynamic pressure, $\rho V^{2} / 2, l b / s q$ ft
$\rho \quad$ mass density of air, slugs/cu ft
$p$ rate of roll, radians/sec
$\Lambda$ sweep angle at quarter-chord line, deg
V free-stream velocity, ft/sec
M Mach number
$R$ Reynolds number
S wing area, sq ft
b wing span, ft
y lateral distance from plane of symmetry, ft
$\lambda$ wing taper ratio, $\frac{\text { Tip chord }}{\text { Root chord }}$
2 body length
d body diameter
A aspect ratio, $b^{2} / \mathrm{s}$
a angle of attack, deg
$\frac{\mathrm{pb}}{2 \mathrm{~V}}$ wing-tip helix angle, radians
$C_{q_{p}}=\frac{\partial C_{\imath}}{\partial \frac{p b}{2 V}}$ per radian
$C_{n p}=\frac{\partial C_{n}}{\partial \frac{p b}{2 V}}$ per radian
$C_{Y_{p}}=\frac{\partial C_{Y}}{\partial \frac{p b}{2 V}}$ per radian

MODEL AND APPARATUS

A sketch of the models investigated and details of the fence are shown in figure 2. All wings except the $45^{\circ}$ swept wing were constructed of $245-T$ aluminum alloy. The $45^{\circ}$ swept wing had a steel core with a bismuth-tin covering. The wings had a taper ratio of 0.6 , an aspect ratio of 4 , and an NACA 65 A006 airfoil section and were attached to the fuselage in a midwing position. The geometric characteristics of the body are presented in table I.

The models were tested on the forced-roll sting-support system shown in figure 3. Angles of attack were obtained by means of offset sting adapters in the sting behind the model (fig. 3). The forces and moments were measured on an internally mounted electrical strain-gage balance.

## TESTS AND CORRECTIONS

The forced-roll tests were made in the Langley high-speed 7-by 10 -foot tunnel through a Mach number range from 0.50 to 0.95 and through a maximum angle-of-attack range from $0^{\circ}$ to $13^{\circ}$. Tests were also conducted on the $45^{\circ}$ swept wing with wing fences located at the 65 -percent-semispan station. The variation with Mach number of the mean Reynolds number (based on the mean aerodynamic chord of the wing) and the maximum values of $\mathrm{pb} / 2 \mathrm{~V}$ are presented in figures 4 and 5, respectively.

The blocking corrections applied to the dynamic pressure and Mach number were determined by the velocity-ratio method of reference 4. Drag and angle of attack were corrected for jet-boundary effects by the method of reference 5, but an investigation of the jet-boundary corrections to the rolling derivatives by methods similar to those used in reference 6 indicated that these corrections were negligible. Tare tests were made at zero angle of attack with and without a simulated offset sting adapter behind a similar model and the effects were found to be negligible.

The data presented have been corrected for inertia forces and moments that were introduced as the model was rotated, with consideration being given to deflections of the entire support system under aerodynamic loads.

In order to evaluate the aeroelastic corrections to $C_{Z_{p}}$ at zero angle of attack, the wings were statically loaded in accordance with theoretical load distributions obtained from reference 7. The resulting incremental changes in wing-section angle of attack have been interpreted
in terms of equivalent rotational velocities and the results are presented in figure 6. Values of the equivalent linear twist distribution indicated at $\frac{y}{b / 2}=1.0$ can be interpreted as correction increments $\frac{\Delta\left(\frac{p b}{2 V}\right)}{q_{2}}$, which can be applied to the measured values of $\mathrm{pb} / 2 \mathrm{~V}$ according to the equation

$$
C_{q_{p}}=\frac{\left(C_{q}\right)_{\text {meas }}}{\frac{p b}{2 V}-\Delta \frac{p b}{2 V}}=\frac{\left(C_{z_{p}}\right)_{\text {meas }}}{I-\frac{\Delta(p b / 2 V)_{q}}{q C_{2}}\left(\frac{C_{q}}{p b / 2 V}\right)_{\text {meas }}}
$$

or

$$
c_{l_{p}}=\frac{\left(c_{l_{p}}\right)_{\text {meas }}}{1-k\left(c_{l_{p}}\right)_{\text {meas }}}
$$

where

$$
\mathrm{k}=\frac{\Delta(\mathrm{pb} / \mathrm{V}) \mathrm{q}_{\mathrm{q}}}{\mathrm{qC}_{2}}
$$

and is presented as a function of Mach number in figure 7.
Distortion effects on $C_{n_{p}}$ and $C_{Y_{p}}$ have been roughly estimated and, since these effects appeared to be small over the angle-of-attack range for all wings, they have been neglected.

The angle of attack at the plane of symmetry has been corrected for the deflection of the model and support system under load. All data are referred to the stability axes system.

RESULTS AND DISCUSSION

The results of the present investigation are presented as outlined in the following table:

Results
Figure
$\mathrm{C}_{2}$ against $\mathrm{pb} / 2 \mathrm{~V}$. . . . . . . . . . . . . . . . . . . . . . . . 8
$\mathrm{C}_{\mathrm{n}}$ against $\mathrm{pb} / 2 \mathrm{~V}$9

${ }^{C_{2}}$ against $\alpha$ ..... 10


$\mathrm{C}_{\chi_{p}}$ against $M$ (compared with calculations) . . . . . . . . . . 14
$C_{\eta_{p}}$ against $\Lambda$ (compared with calculations) . . . . . . . . . . . 15
$\mathrm{C}_{\eta_{p}}$ against $a$ (compared with calculations) . . . . . . . . . . 16
${ }^{C_{p}}$ against $\alpha$ (compared with calculations)17
$C_{D}$ against ..... $\alpha$
18 and 19
$\mathrm{C}_{\mathrm{n}_{\mathrm{p}}}$ against $\alpha$ (compared with calculations) . . . . . . . . 20 to 22 ${ }^{C_{Y}}$ against $\alpha$ (compared with calculations) ....... . 23 and 24

Results of $C_{2}$ and $C_{n}$ plotted against $\mathrm{pb} / 2 \mathrm{~V}$ are presented in figures 8 and 9 only for those angle-of-attack conditions for which pronounced nonlinearities, with respect to $\mathrm{pb} / 2 \mathrm{~V}$, were indicated. For all conditions not covered in figures 8 and 9, the data were sufficiently linear to permit adequate representations of the results by means of derivatives with respect to $\mathrm{pb} / 2 \mathrm{~V}$.

## Experimental Rolling Derivatives

Rolling moment due to rolling.- In general, the damping due to roll $\mathrm{C}_{i_{p}}$ decreased above an angle of attack of about $6^{\circ}$ (fig. 10), and in the region of low damping or positive values of $C_{q_{p}}$, the variation of rolling moment with rolling velocity was rather erratic (fig. 8). These nonlinearities and ranges of uncertainties of $C_{l}$ against $\mathrm{pb} / 2 \mathrm{~V}$ are difficult to analyze in quantitative terms of $C_{q_{p}}$ (see shaded areas
of figs. 10,11 , and 16 , which are approximately average values of $C_{l_{p}}$ ) and, therefore, the application of the data with respect to controllability, rolling, and dynamic stability should be carefully considered. Some additional remarks on the damping for these wings are given in reference 8.

At an angle of attack of $11^{\circ}$ and a Mach number of 0.85 , the $3.6^{\circ}$ swept-wing data (fig. 8(a)) show a pronounced hysteresis. The data were obtained by rolling the model from the extreme negative values of $\mathrm{pb} / 2 \mathrm{~V}$ to the extreme positive values of $\mathrm{pb} / 2 \mathrm{~V}$, then back through the $\mathrm{pb} / 2 \mathrm{~V}$ range. It should be pointed out that these nonlinearities and hystereses occur above the angle of attack at which this wing indicates a peak in the lift curve (ref. l).

The results for the $32.6^{\circ}$ swept wing (figs. 8(b) and $10(\mathrm{~b})$ ) show that, in general, at the higher Mach numbers and angles of attack, an unstable condition (positive values of $C_{l_{p}}$ ) is apparent over a wide range of rolling velocities; whereas a stable condition is indicated only at the extreme rates of roll. The results presented in figure ll show that, of the wings investigated, this plan form showed unfavorable damping-in-roll characteristics (indicated by $C_{\eta_{p}}$ ) over the largest test ranges of Mach number and angle of attack.

The $45^{\circ}$ sweptback wing shows only a small region of zero or reverse damping which occurred at a Mach number of 0.91 , an angle of attack of about $10.8^{\circ}$, and at low values of $\mathrm{pb} / 2 \mathrm{~V}$. (See figs. 8(c) and 10(c).) Considerable loss in damping did exist, however, at the higher test angles of attack throughout the test range of Mach number. The configuration with fences installed did not exhibit as much loss of damping at higher angles of attack as did the clean configuration. As was shown in reference 9, the fences improve the lifting capabilities of the airfoil sections near the wing tips on the $45^{\circ}$ swept wing and, therefore, improve the damping in roll as well as the high-lift longitudinal stability. (See refs. 8 and 10.)

The decrease in damping noted for all wings at the higher test angles of attack is probably associated with tip-stalling as shown for the $45^{\circ}$ sweptback wing in references 9 and 11.

Yawing moment and lateral force due to rolling.- The variation of the yawing-moment-due-to-rolling derivative $C_{n_{p}}$ with angle of attack is presented in figure 12. The general trend of $C_{n_{p}}$ with angle of attack is similar for all wings; that is, zero or slightly negative values are obtained at the lower angles of attack and positive values are obtained at the higher test angles of attack.

In figure 13, the variation of $C_{Y p}$ with angle of attack is presented. Negative values of $C_{Y_{p}}$ (lateral force due to rolling) are indicated at the higher test angles of attack and at most Mach numbers for all configurations tested. A discussion and possible explanation of the behavior of $C_{n_{p}}$ and $C_{Y_{p}}$ through the test angle-of-attack and Mach number range and a method for estimating these derivatives are presented in the following section.

## Estimations of Rolling Derivatives

Rolling moment due to rolling.- A comparison of the experimental and calculated variations of $C q_{p}$ at zero angle of attack with Mach number is presented in figure 14. The calculated variations were evaluated by the methods described in references 7 and 12. The predicted trend, at least up to the critical Mach number, is in fairly good agreement with experiment. The aeroelastic effects are seen to be of appreciable magnitude for the $45^{\circ}$ and $60^{\circ}$ sweptback wings.

The variation of $C_{l_{p}}$ at zero angle of attack with sweep angle presented in figure 15 shows that, at the higher Mach numbers ( 0.85 to 0.93 ), the 32.60 sweptback wing has higher values of damping, corrected for aeroelasticity, than the other test wings, although the calculations of references 7 and 12 predict a decrease in $C_{l_{p}}$ with increasing sweep angle.

A comparison of the variation of $C_{q_{p}}$ with angle of attack, for all configurations tested, determined by experiment and calculations, is shown in figure 16. Values of $\mathrm{C}_{l_{p}}$ at zero angle of attack were determined
from reference 7 , and compressibility effects were evaluated from reference 12. Angle-of-attack effects were determined by the procedure of reference 13 by using the experimental lift-curve slopes of references 1 and 2. It has been shown in reference 14 that root-bending-moment data would be more appropriate than lift data in evaluating angle-of-attack effects on ${ }^{C} Z_{p}$; however, bending-moment data were not available at these angles of attack for the wings of the present investigation. The quantitative agreement shown in figure 16 is not very good; however, the experimental and predicted values show similar trends.

Yawing moment due to rolling.- In references 14 and 15 , methods are presented for the prediction of $C_{n_{p}}$ through an angle-of-attack range using corresponding experimental drag data. The method in reference 14 has been shown to predict $C_{n_{p}}$ through the test angle-of-attack range
with considerably better results than the method described in reference 15. This result is probably due to the fact that the method of reference 15 predicts $C_{n_{p}}$ by extrapolating from the potential-flow theory by use of experimental drag data and an empirically determined factor that is proportional to a drag-center moment arm; whereas the procedure of reference 14 predicts $C_{n_{p}}$ by using the experimental drag data to proportion $C_{n_{p}}$ relative to two known conditions. Briefly, the method of reference 14 consists of proportioning $C_{n_{p}}$ relative to the condition of potential flow where the resultant force is normal to the relative wind and to the condition for which the resultant force is normal to the wing chord. An equation is presented therein for evaluating $C_{n_{p}}$ for triangular wings. However, for other wing plan forms, this equation must be modified as follows: the potential-flow value of $C_{n_{p}}$ for triangular plan forms (ref. 16) must be replaced by the value for wings of taper ratios other than 0, and an additional term, shown in reference 15 for low-aspect-ratio wings to be a result of wing-tip suction, must be considered. With these considerations applied, the following equation can be written:

$$
\begin{equation*}
C_{n_{p}}=-C_{l_{p}} \tan \alpha-K\left(-C_{l_{p}} \tan \alpha-\frac{C_{n_{p}}}{C_{L}} C_{L}\right)+\left(C_{n_{p}}\right)_{\text {tip suction }} \tag{1}
\end{equation*}
$$

where the potential-flow values of $C_{n_{p}} / C_{L}$ can be determined from reference 17 and either experimental or calculated values of $C_{l_{p}}$ can be used. The factor $K$ is a dimensionless factor that relates $C_{n_{p}}$ to any intermediate flow condition that exists between the conditions where the resultant force is normal to the relative wind and where the resultant force is normal to the wing chord and can be determined from the drag data of figures 18 and 19 as follows:

$$
K=\frac{\frac{\partial}{\partial \alpha}\left(C_{L} \tan \alpha\right)-\frac{\partial}{\partial \alpha}\left(C_{D}-C_{D_{C_{L}}=0}\right)_{\exp }}{\frac{\partial}{\partial \alpha}\left(C_{L} \tan \alpha\right)-\frac{\partial}{\partial \alpha} \frac{C_{L}{ }^{2}}{\pi A}}
$$

The tip-suction contribution is expressed in reference 15 as

$$
\left(\mathrm{C}_{\mathrm{n}_{\mathrm{p}}}\right)_{\text {tip suction }}=\left(-\mathrm{C}_{\mathrm{Y}}\right)_{\Lambda=0^{\circ}} \frac{\mathrm{d}}{\mathrm{~b}}
$$

where $d$ is the longitudinal distance from the midchord point at the wing tip to the coordinate origin and is defined as

$$
d=\frac{b}{2(1+\lambda)}\left(\frac{2+\lambda}{3} \tan \Lambda+\frac{\lambda}{A}\right)+X^{\prime}
$$

where $X^{\prime}$ is the longitudinal distance rearward from the coordinate origin to the wing aerodynamic center. An attempt was made in reference 15 to account for tip suction by correlating data obtained at low speeds on 12 -percent-thick unswept wings with the slender-triangular-body theory of reference 16. Shown in reference 15 is an empirical expression, determined from a limited amount of data, that expresses the tip-suction contribution as

$$
\left(\mathrm{C}_{\mathrm{Y}}\right)_{\Lambda=0^{\circ}}=\frac{\mathrm{C}_{\mathrm{L}}}{\mathrm{~A}}
$$

The tip-suction contribution to $C_{n_{p}}$ can be written now in terms of the empirically determined suction force and moment arm as

$$
\left(C_{n_{p}}\right)_{\text {tip suction }}=-\frac{C_{L}}{2 A(1+\lambda)}\left(\frac{2+\lambda}{3} \tan \Lambda+\frac{\lambda}{A}\right)-\frac{C_{L}}{A} \frac{X^{\prime}}{b}
$$

For the wings of the present investigation, $X^{\prime}$ is equal to 0 . The results presented in figure 20 at a Mach number of 0.70 show that values of $C_{n_{p}}$ evaluated by use of equation (l), using calculated values
of ${ }^{C} l_{p}$, are in better agreement with experiment than values determined independently of the consideration of tip suction by the method of reference 15. It should be pointed out, although not shown, that a similar comparison of the two methods was obtained for the other wings tested. Better agreement is indicated when the tip suction in equation (1) is assumed to be 0; however, this result is not surprising inasmuch as the experimental values of $\mathrm{C}_{\mathrm{Y}_{\mathrm{p}}}$ for the unswept wing (which are due primarily to tip suction) presented in figure 17 indicate that, within the test ranges of Mach number and angle of attack, the tip-suction contribution can be neglected. The data for a Mach number of 0.15 (fig. 17) were obtained from reference 18.

The leading-edge contribution to $C_{n_{p}}$ would be expected to vary considerably with leading-edge radii; consequently, in figure 21 low-speed
results are presented for 12-percent-thick wings (ref. 19) whose leadingedge radii vary from a very sharp one to a very blunt one. The wings reported in reference 19 had an aspect ratio of 2.61 , a taper ratio of 1.0 , and 450 sweepback. A comparison of experimental $C_{n_{p}}$ with calculated values evaluated by use of equation (l) is presented in figure 22 for the wings of the present investigation at various Mach numbers where both experimental and calculated values of $C_{q_{p}}$ were used; the tipsuction contribution to $C_{n_{p}}$ was assumed to be 0 . The agreement is considered good for all wings tested where either experimental or calculated $C_{l_{p}}$ are used in equation (1), and the results presented in figures 21 and 22 indicate that the present method of estimating $C_{n_{p}}$ (eq. (l), without tip-suction effects) is applicable over large ranges of leading-edge radii, wing thickness, and subsonic Mach numbers.

Lateral force due to rolling.- An expression for determining $\mathrm{C}_{Y_{p}}$ for the potential-flow case can be obtained from reference 17 and written as

$$
C_{Y_{p}}=C_{L} \frac{A+\cos \Lambda}{A+4 \cos \Lambda} \tan \Lambda
$$

and, for the nonpotential case, when the resultant force is normal to the wing chord plane, ${ }^{C} Y_{p}$ would be equal to 0 . By considering the tip-suction contribution and by applying the factor $K$, the equation can be written

$$
\begin{equation*}
C_{Y_{p}}=K\left(C_{L} \frac{A+\cos \Lambda}{A+4 \cos \Lambda} \tan \Lambda\right)+\left(C_{Y_{p}}\right)_{\text {tip suction }} \tag{2}
\end{equation*}
$$

where the tip contribution to $\mathrm{C}_{\mathrm{Y}_{\mathrm{p}}}$ can be expressed as

$$
\left({ }^{C_{p}}\right)_{\text {tip suction }}=\frac{C_{L}}{A}
$$

In figure 23, values of $C_{Y_{p}}$ evaluated by equation (2) and values determined by the potential-flow method of reference 15 are compared with low-speed data of reference 19. In the application of equation (2), the contribution of the tip was assumed to be zero; however, the order of magnitude of the tip-suction contribution is indicated
in the comparison of the values determined by the procedures described in references 15 and 17. Better agreement is indicated when equation (2) of the present paper is used; however, this agreement is not surprising inasmuch as the methods of references 15 and 17 do not account for any nonpotential-flow effects on the leading-edge contribution. The results of the present investigation are compared with values evaluated by use of equation (2), without tip-suction effects, in figure 24. The agreement shown is reasonably good, particularly the negative trend indicated at the higher test angles of attack.

## CONCLUSIONS

An investigation conducted to determine the effects of sweep angle on the rolling derivatives at high subsonic Mach numbers and high angles of attack for a series of swept wings with aspect ratio of 4 , taper ratio of 0.6 , and NACA 65A006 airfoil sections indicates the following conclusions:

1. The results show large reductions in the damping-in-roll derivative $C_{l_{p}}$ at the higher test angles of attack for all wings tested. Of the wings investigated, the results for the $32.6^{\circ}$ sweptback wing showed unfavorable damping-in-roll characteristics (indicated by $\mathrm{Clp}_{\mathrm{p}}$ ) over the largest ranges of angle of attack and Mach number.
2. Wing fences on the $45^{\circ}$ sweptback wing at the 65 -percent-semispan station are shown to improve the damping-in-roll derivative $C_{l_{p}}$ at the higher test angles of attack relative to the clean-wing configuration.
3. In general, the variation of the damping-in-roll derivative $C_{l_{p}}$ at zero angle of attack with sweep angle was only in fair agreement with the predicted variation with sweep angle, inasmuch as the $32.6^{\circ}$ sweptback wing showed more damping in roll in the Mach number range from 0.85 to 0.93 than any of the other plan forms.
4. The predicted variation of $C_{\eta_{p}}$ at zero angle of attack with Mach number was in good agreement with the experimental trend up to the critical Mach number.
5. Contrary to predictions based on potential-flow theory, the yawing moment due to rolling $C_{n_{p}}$ was positive and the lateral force due to rolling $C_{Y_{p}}$ was negative at the higher test angles of attack for all wings tested.
6. The results indicate a loss in wing-tip suction within the test ranges of Mach number and angle of attack investigated.
7. Presented herein is a method of estimating $C_{n_{p}}$ (yawing moment due to rolling) and $C_{Y_{p}}$ (lateral force due to rolling) through the test angle-of-attack range. This method is shown to be applicable over large ranges of leading-edge radii, wing thickness, and subsonic Mach numbers.

Langley Aeronautical Laboratory,
National Advisory Committee for Aeronautics, Langley Field, Va., March ll, 1954.

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## TABLE I.- FUSELAGE ORDINATES

[Basic fineness ratio 12, actual fineness ratio 9.8 achieved by cutting off rear portion of body


| Ordinates, percent length |  |
| :---: | :---: |
| Station | Radius |
| 0 | 0 |
| .61 | .28 |
| .91 | .36 |
| 1.52 | .52 |
| 3.05 | .88 |
| 6.10 | 1.47 |
| 9.15 | 1.97 |
| 12.20 | 2.40 |
| 18.29 | 3.16 |
| 24.39 | 3.77 |
| 30.49 | 4.23 |
| 36.59 | 4.56 |
| 42.68 | 4.80 |
| 48.78 | 4.95 |
| 54.88 | 5.05 |
| 60.98 | 5.08 |
| 67.07 | 5.04 |
| 73.17 | 4.91 |
| 79.27 | 4.69 |
| 85.37 | 4.34 |
| 91.46 | 3.81 |
| 100.00 | 3.35 |
| Leading-edge radius $=0.00062$ |  |



Figure l.- System of axes used showing positive direction of forces, moments, angles, and velocities.

## Fuselage:

$\begin{array}{ll}\text { Length } & 49.2 \mathrm{in} \\ \text { Max. diam. } & 5 \mathrm{in} .\end{array}$

Wing:

| Area | 2.25 sqft |
| :--- | :--- |
| Span | 3.0 ft |
| Chord |  |
| $\quad$ Tip | 6.75 in. |
| Root | 11.25 in. |
| Mean aerodynamic chord | .765 ft |
| Aspect ratio | 4.0 |
| Taper ratio | 0.6 |
| Incidence | 0 |
| Dihedral | 0 |

Airfoil section
parallel to fuselage \& NACA 65A006
2.25 sqft
6.75 in .
11.25 in. .765 f
0.6

0
,



Figure 3.- General arrangement of forced-roll sting-support system.


Figure 4.- Variation of mean test Reynolds number with Mach number for wings tested.


Figure 5.- Variation of maximum test values of $\mathrm{pb} / 2 \mathrm{~V}$ with Mach number for wings tested.


Figure 6.- Aeroelastic characteristics for wings.


Figure 7.- Aeroelastic correction factors for wings tested. $\alpha=0^{\circ}$.

$$
\begin{aligned}
& \text { From }-\frac{p b}{2 V} \text { to }+\frac{p b}{2 V} \\
& \text {----From }+\frac{p b}{2 V} \text { to }-\frac{p b}{2 V}
\end{aligned}
$$



(a) 3.6-4-.6-006 wing.

Figure 8.- Variation of $\mathrm{C}_{2}$ with $\mathrm{pb} / 2 \mathrm{~V}$ in high test angle-of-attack range for three of wings investigated.


Unflagged symbols, clean wing

(c) 45-4-.6-006 wing.

Figure 8.- Concluded.



Figure 9.- Variation of yawing moment with wing-tip helix angle $\mathrm{pb} / 2 \mathrm{~V}$ in high test angle-of-attack range. 3.6-4-.6-006 wing.

(a) 3.6-4-.6-006 wing.


Figure 10.- Concluded.

$$
\begin{aligned}
& \text { पाTVP } C_{l_{p}}=\text { o or pos itive } \\
& --C_{l_{p}}=\frac{1}{2}\left(C_{l_{p}}\right)_{C_{L}}=0 \\
& --- \text { Test limits }
\end{aligned}
$$



Figure 11.- Boundaries for $C_{l_{p}}$ for test ranges of lift coefficient and Mach number.


> -(Unflagged symbols) clean wing
> --(Flagged symbols) fence on


Figure l2.- Concluded.

_(Unflagged symbols) clean wing
---(Flagged symbols) fence on

(c) 45-4-.6-006 wing. (d) 60-4-.6-006 wing.

Figure 13.- Concluded.


Figure 14.- Comparison of experimental and calculated variation of ${ }^{C} l_{p}$ with Mach number. $\alpha=0^{\circ}$.


Figure 15.- Comparison of experimental and predicted variation of $C_{l_{p}}$ with sweep angle. $\alpha=0^{\circ}$.


Figure 16.- Comparison of experimental and calculated variation of $C_{q_{p}}$ with angle of attack for wings tested. Data not corrected for aeroelastic distortion.


Figure 16.- Concluded.


Figure 17.- Comparison of experimental $C_{Y_{p}}$ with calculated values for 3.6-4-.6-006 wing.

(a) 3.6-4-.6-006 wing.
(b) 32.6-4-.6-006 wing.

(c) 45-4-.6-006 wing.
(d) 60-4-.6-006 wing.
Figure 18.- Concluded.


Angle of attack, a, deg
Figure 19.- Variation of low-speed drag due to lift for untapered $45^{\circ}$ swept wings of aspect ratio 2.61 .

$$
\begin{array}{cc} 
& M=.70 \\
\text { Experiment } & \text { calculated } \\
\square & -- \text { Without tip suction } \\
& - \text { With tip suction }
\end{array}
$$



Figure 20.- Comparison of experimental values of $C_{l_{p}}$ with various methods of calculation for 32.6-4-.6-006 wing.


Figure 21.- Comparison of low-speed experimental $C_{n_{p}}$ with calculated values for untapered $45^{\circ}$ swept wings of aspect ratio 2.61 .


Figure 22.- Comparison of experimental and calculated variation of $C_{n_{p}}$ with angle of attack.

$$
\begin{aligned}
& 0 \square \diamond \text { Data from ref. } 19, M=.17 \\
& - \text { Withouttip suction } \\
& -- \text { With tip suction } 15 \\
& -- \text { Without tip suction, equation (2) }
\end{aligned}
$$



Figure 23.- Comparison of low-speed experimental $C_{Y_{p}}$ with calculated values for untapered $45^{\circ}$ swept wings of aspect ratio 2.61 .
-_ Equation (2), without tip suction. - $\diamond \Delta$ Experiment



Figure 24.- Comparison of experimental and calculated variation of $C_{Y_{p}}$
with angle of attack.

