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RESEARCH MEMORANDUM

AERODYNAMIC DAMPING AT MACH NUMBERS OF 1.3 AND 1.6 OF

A CONTROL SURFACE ON A TWO-DIMENSIONAL WING

BY THE FREE-OSCILLATION METHOD

By W. J. Tuovila and Robert W. Hess

Langley Aeronautical Laboratory Langley Field, Va.

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NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS

> WASHINGTON May 1, 1956

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RESEARCH MEMORANDUM

AERODYNAMIC DAMPING AT MACH NUMBERS OF 1.3 AND 1.6 OF

A CONTROL SURFACE ON A TWO-DIMENSIONAL WING

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SUMMARY

Tests have been made at two supersonic speeds to obtain experimentally the aerodynamic damping characteristics of a control surface on a two-dimensional wing. The control surface had a chord of 1.67 inches (1/3 of the wing chord) and a span of 7.25 inches and was supplied in three materials (steel, aluminum, and magnesium) having different mass, inertia, and stiffness properties. Two wing sections were tested, one being of 65A004 section and the other a 5-percent-thick hexagonal section. The test results are compared with results calculated by two- and three-dimensional oscillating air-force theories. At a Mach number of 1.6, both theories are in fairly good agreement with the experimental results. At a Mach number of 1.3, both theories predict negative (unstable) damping, whereas the tests indicate that the damping is slightly positive (stable). The in-phase or aerodynamic stiffness coefficients predicted by both theories are slightly higher than the experimentally determined coefficients.

INTRODUCTION

Theoretical studies have indicated that at low supersonic speeds control surfaces with a single degree of torsional freedom can encounter unstable aerodynamic damping at some values of reduced frequencies. Since existing theories do not account for many flow effects which may influence the problem, tests were made to obtain some experimentally determined aerodynamic damping coefficients for comparison with wheoretical values. Aerodynamic in-phase or stiffness coefficients and outof-phase or damping coefficients were determined for a 1/3-chord control surface attached to a two-dimensional wing at zero angle of attack. Wings with hexagonal and 65A004 section shape were used. The tests were

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made at Mach numbers of 1.3 and 1.6 over a reduced frequency range from 0.029 to 0.074. This paper presents the test results and compares them with results calculated using two- and three-dimensional theories for oscillating air forces. The test results are also compared with the results of some damping tests made on a control surface attached to a triangular wing (ref. 1).

SYMBOLS

Aa	aspect ratio of control surface, $l_a/2b_a$
b _a	semichord of control surface, ft
δ	control-surface deflection, radians
f _o	natural frequency of rotation of control surface about hinge line at zero airspeed, cps
ft	natural frequency of rotation of control surface about hinge line at test Mach number, cps
go	damping coefficient associated with fo
St	demping coefficient associated with ft
la	length of control surface, ft
ma	mass of control surface, slugs/ft of span
N ₅	in-phase aerodynamic coefficient per foot of span
<u>n</u> 6	out-of-phase or damping coefficient per foot of span
М	Mach number
v	airspeed, fps
ρ.	air density, slugs/ft ³
ωο ιτ. '	2af ₀
യപ്പ	21ft
k	reduced frequency $b_{\alpha}\omega_t/V$

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k_a spring constant, ft-lb/radian

Ché -2kN6

Ia mass moment of inertia about hinge line, slug-ft²/ft of span

 r_a^2 . $I_a/m_a b_a^2$

 $\mu = \frac{m_a}{4\rho b_a^2}$

Subscripts:

8.	refers to control surface
0	refers to conditions at zero wind velocity
t	refers to conditions at test Mach number

MODELS AND TEST METHODS

Wing, control surface, and hinge details are given in figure 1. Control surfaces made of steel, aluminum, and magnesium were tested on two steel wing models which differed only in section. One wing model had a 65A004 section and the other had a 5-percent-thick hexagonal section. Each wing had a 5-inch chord and spanned the tunnel test section with one end clamped in the sidewall and the other end pinned in the sidewall. The control-surface chord was 1/3 of the wing chord. Steel hinges of various stiffnesses were used to attach the control surfaces to the wings at three points. There was a gap of about 0.02 inch between the wing and the control surface. Table 1 lists some of the physical parameters of the models. The masses and inertias were determined experimentally and include the contribution of the hinges.

The tests were made in the 9- by 18-inch Langley supersonic flutter apparatus which is an intermittent-flow blow-down tunnel operated at atmospheric stagnation pressure. The testing technique used was first to obtain "no-wind" damping decrements with the wing in the testing configuration by flicking the control surface. The control surface was then deflected, the tunnel was brought up to speed, and the control surface was released and a "wind-on" damping decrement was obtained. The air flow was then stopped and the process was repeated using hinges of different stiffness.

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The initial amplitude of both the "no-wind" and the "wind-on" oscillations was not controlled precisely. It was judged by eye to range from about $\pm 1^{\circ}$ to $\pm 2\frac{1^{\circ}}{2}$, the larger amplitudes occurring at the lowest frequencies.

The system for deflecting the aileron is illustrated in figure 1 and consisted of a wire with an eye on the end which was inserted through a small hole at the trailing edge of the control surface. A straight release wire was then inserted through the eye of the cocking wire. The control surface was cocked by pulling the cocking wire until the desired deflection was obtained. The control surface was released by pulling the release wire out of the eye of the cocking wire.

Damping decrements were obtained from a strain gage glued to a thin metal strip fastened to the wing and control surface. This metal strip followed the control-surface motion and the strain-gage output was amplified and fed into a recording oscillograph.

REDUCTION OF DATA

The experimental decay decrements were reduced to average total supersonic aerodynamic coefficients \overline{N}_5 and \overline{N}_6 as was done in reference 2 for subsonic flow. All damping terms are assumed proportional to amplitude and in phase with velocity. The following equation of equilibrium,

$$I_{a}\ddot{\delta} + k_{a}\left(1 + ig_{0}\right)\delta = -4\rho b_{a}^{2}V^{2}k^{2}\delta\left(\overline{N}_{5} + i\overline{N}_{6}\right)$$
(1)

leads to the following results for the in-phase component,

$$\bar{N}_{5} = \mu r_{a}^{2} \left[1 - \left(\frac{\omega_{0}}{\omega_{t}} \right)^{2} \right]$$
(2)

and, for the out-of-phase or damping component,

$$\overline{N}_{6} = \mu r_{a}^{2} \left[g_{t} - g_{0} \left(\frac{\omega_{0}}{\omega_{t}} \right)^{2} \right]$$
(3)

The details of the analysis are given in the appendix.

It may be noted that the damping component is not obtained from just the difference in the damping coefficients of the "wind-on" and "no-wind" decrements. Instead, the "no-wind" damping coefficient is reduced by the factor $(\omega_0/\omega_t)^2$, which accounts for the difference in the structural damping coefficient due to the difference in frequency between "wind-off" and "wind-on" conditions. It is of interest to note that at M = 1.3 the "no-wind" damping coefficient g_t but the factor $(\omega_0/\omega_t)^2$ made the aerodynamic damping coefficient \overline{N}_6 slightly positive.

The experimentally determined \overline{N}_5 and \overline{N}_6 are compared with twoand three-dimensional air-force coefficients obtained from references 3 and 4. For comparison with the results obtained in reference 1, the damping coefficient \overline{N}_6 is expressed in stability notation using viscoustype damping terms as follows:

$$C_{h\delta} = \frac{\partial C_{h}}{\partial \frac{b\delta}{V}} = -2k\overline{N}_{6}$$
(4)

RESULTS AND DISCUSSION

Presentation of Data and Comparison With Theory

The control surfaces were attached to two-dimensional wings set at zero angle of attack. The aerodynamic in-phase and damping coefficients were obtained from the decay records and frequencies obtained in both still air and at the test Mach numbers of 1.3 and 1.6 and the data are presented in table 2. Sample "wind-off" and "wind-on" decrements are shown in figures 2(a) and 2(b). The hinge axis was so near the leading edge of the control surface that it was assumed to be there. The aerodynamic damping coefficients \overline{N}_6 are presented in figure 3 and the inphase coefficients \overline{N}_5 are presented in figure 4. The aerodynamic coefficients are plotted against the reduced frequency, based on the control-surface semichord.

The experimental results are compared with the two-dimensional theory of reference 3 by assuming the control surface to be a wing oscillating about its leading edge and with the three-dimensional theory of reference 4, assuming a sealed gap between the wing and the control surface. The theoretical results are also plotted on figures 3 and 4. Both theories predict negative aerodynamic damping at M = 1.3; however, the three-dimensional theory predicts only about 1/2 the damping of the



two-dimensional theory. The experimental aerodynamic damping at M = 1.3 is slightly positive and both theories approach it as k increases. At M = 1.6 both theories are in good agreement with the experimental damping results, the three-dimensional theory giving slightly higher values than the two-dimensional theory.

The experimental in-phase aerodynamic coefficients N_5 presented in figure 4 are fairly consistent and both theories predict the trends well. The two-dimensional theory gives slightly higher values than the three-dimensional theory does and both theories yield values that are higher than the experimental.

It appears that linearized flow theory, when applied to flow around trailing-edge control surfaces, begins to break down at low Mach numbers in the neighborhood of 1.3 or less. Adding an aspect-ratio correction to the two-dimensional-flow theory improves the results; however, some basic differences between the actual and the idealized flow appears to affect the results. Wing thickness, boundary layer, and the gap between the wing and control surface are some factors whose effects are not included in the theory. Also, the experimental results were obtained from decaying oscillations, whereas the theory assumes constant-amplitude oscillations. At M = 1.6 the theory seems to compensate for these effects and the agreement is good.

Comparison With Control-Surface Data for a Triangular Wing

The results of the present tests are compared in figure 5 with those of reference 1 through the Mach number range. Results for an amplitude of $\pm 3^{\circ}$ at a maximum k value of 0.03 from reference 4 are compared with the results of the present tests for amplitudes of about $\pm 2^{\circ}$ at k values of 0.045. The damping coefficients are expressed in stability notation as C_{hs} . The difference in the present results and those of reference 1

may be the result of differences in flow caused by the wings. It may be noted that in reference 1 the control surface is attached to an aspectratio-2 triangular wing and not to a two-dimensional wing. In reference 1 the damping varied from a small degree of instability at M = 1.3to neutral stability at M = 1.9, whereas the present tests indicate slight stability at M = 1.3 and considerable stability at M = 1.6. The two- and three-dimensional theory results are also presented in figure 5.

GENERAL OBSERVATIONS

At M = 1.3 there is considerable scatter in the results but the damping coefficients in all but one case are positive. This scatter is

due to the sensitivity of the equation for \overline{N}_6 to small changes in measured damping between the "wind-off" and "wind-on" conditions when the aerodynamic damping is low. No flutter was observed during these tests which indicates that the total damping was positive and shows that the aerodynamic damping could have been, at most, only slightly negative since the structural damping was small. At M = 1.6, where the aerodynamic damping is higher, the scatter is considerably reduced. Any effects due to wing-profile or control-surface material is lost within the scatter of the results.

The structural damping g_0 was principally in the range 0.006 to 0.01 with a few extreme values of $g_0 = 0.004$ on the low end and $g_0 = 0.034$ on the high end. This spread in the structural damping coefficient is believed to be due to variations in the hinge clamping force. Also, the structural damping coefficient generally decreased with decrease in amplitude and some unusually large changes are noted in table 2(a) for $\mu r_a^2 = 650$. The damping coefficients recorded in table 2(a) were measured near the maximum amplitude of oscillation.

The aerodynamic damping may also be affected by amplitude; however, since the present tests were made without amplitude control, no such effect can be determined. No appreciable amplitude effect is indicated in reference 1 at Mach numbers from 1.3 to 1.9 while reference 5 shows considerable effect for amplitudes up to $\pm 5^{\circ}$ at Mach numbers near 1.0.

Wing bending motion may also affect the results by introducing a translation degree of freedom to the control surface. Although the wing motion was not measured, it is believed to have been very slight since the wing was clamped at one end and pinned at the other. As the control-surface frequency approached the wing resonant frequency, the wing amplitude would increase rapidly and any bending effect should become evident. At M = 1.6 the NACA 65A004 wing with control surface $\mu r_{\rm g}^2 = 378$ reached the wing resonant frequency at k = 0.069 and yielded essentially the same results as the hexagonal wing with control surface $\mu r_{\rm g}^2 = 427$ where the control-surface frequency was 85 percent of the wing resonant frequency. The NACA 65A004 wing would have had about 5 times the amplitude of the hexagonal wing at this k value which indicates that the wing bending amplitude had no apparent effect on the damping results.

CONCLUDING REMARKS

The results of the tests of a control surface attached to a twodimensional wing at zero angle of attack indicate that at a Mach number of 1.3, a slight amount of aerodynamic damping exists on the control

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surface, whereas both two- and three-dimensional theories predict negative damping. At a Mach number of 1.6 the control surface has considerable aerodynamic damping which both two- and three-dimensional theories predict quite well. Both theories predict the trends of the in-phase aerodynamic coefficients, but they yield results which are slightly higher than experimental values. These results were obtained at reduced frequencies from 0.029 to 0.074.

Langley Aeronautical Laboratory, National Advisory Committee for Aeronautics, Langley Field, Va., January 9, 1956.

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APPENDIX A

Derivation of Aerodynamic Coefficients \overline{N}_5 and \overline{N}_6

The supersonic aerodynamic coefficients \overline{N}_5 and \overline{N}_6 are derived from the following equations of equilibrium, where the damping is assumed proportional to the displacement and in phase with the velocity.

 $\pi g = Logarithmic decrement$

"Wind-on" condition (aerodynamic and structural)

$$I_{a}\ddot{\delta} + k_{t}(1 + ig_{t})\delta = 0$$
 (A1)

"No-wind" condition (structural only)

$$I_{a}\ddot{\delta} + k_{a} \left(l + ig_{0} \right) \delta = 0$$
 (A2)

where

then,

$$k_{t}\delta - k_{a}\delta = Aerodynamic spring force$$
 (A3)

and,

$$k_{t}g_{t}\delta - k_{a}g_{o}\delta = Aerodynamic damping force$$
 (A4)

Equation (A4) implies that the structural damping force is independent of frequency.

By definition

$$\overline{N}_{5} = \frac{\text{Aerodynamic spring force}}{4\rho b_{a}^{2} V^{2} k^{2} \delta}$$

$$= \frac{(k_{t} - k_{a}) \delta}{4\rho b_{a}^{2} V^{2} k^{2} \delta}$$
(A5)

and

$$\overline{N}_{6} = \frac{\text{Aerodynamic damping force}}{4\rho b_{a}^{2} V^{2} k^{2} \delta}$$

$$= \frac{(k_{t}g_{t} - k_{a}g_{o})\delta}{4\rho b_{a}^{2} V^{2} k^{2} \delta}$$
(A6)

for small values of damping

$$\omega_{\rm t}^2 = k_{\rm t}/I_{\rm cl}, \quad \omega_{\rm o}^2 = k_{\rm a}/I_{\rm a} \tag{A7}$$

reduced frequency (by definition),

$$k = b_{a}\omega_{t}/V$$
 (A8)

Substituting equations (A7) and (A8) into (A5) and (A6)

$$\overline{N}_{5} = \frac{I_{a} \left(\omega_{t}^{2} - \omega_{0}^{2} \right)}{\mu_{\rho} b_{a}^{2} V^{2} \frac{b_{a}^{2} \omega_{t}^{2}}{V^{2}}} = \frac{I_{a}}{\mu_{\rho} b_{a}^{4}} \left[1 - \left(\frac{\omega_{0}}{\omega_{t}} \right)^{2} \right]$$
(A9)

$$\overline{N}_{6} = \frac{I_{a} \left(\omega t^{2} g_{t} - \omega_{0}^{2} g_{0}\right)}{\mu_{\rho} b_{a}^{2} v^{2} \frac{b_{a}^{2} \omega t^{2}}{v^{2}}}$$

$$= \frac{I_{a}}{\mu_{\rho} b_{a}^{4}} \left[g_{t} - g_{0} \left(\frac{\omega_{0}}{\omega_{t}}\right)^{2} \right]$$
(A10)

finally, substituting

$$\mu r_{a}^{2} = \frac{m_{a}}{\mu \rho b_{a}^{2}} \times \frac{T_{a}}{m_{a} b_{a}^{2}} = \frac{T_{a}}{\mu \rho b_{a}^{\mu}}$$

$$\overline{N}_{5} = \mu r_{a}^{2} \left[1 - \left(\frac{\omega_{0}}{\omega_{t}} \right)^{2} \right]$$
(A11)
$$\overline{N}_{6} = \mu r_{a}^{2} \left[g_{t} - g_{0} \left(\frac{\omega_{0}}{\omega_{t}} \right)^{2} \right]$$
(A12)

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Wing section		Hexagonal		
Control-surface material	Steel Aluminum Magnesium		Magnesium	
l_{a}	0.606 0.0696 0.0145 5.5 × 10 ⁻⁵ 0.782 1133	0.506 0.0692 0.00593 2.29 × 10 ⁻⁵ 0.806 469	0.606 0.070 0.00357 1.24 × 10 ⁻⁵ 0.71 276	0.600 0.0679 0.00679 2.39 × 10 ⁻⁵ 0.766 559

TABLE 1.- SOME CONTROL-SURFACE PHYSICAL PARAMETERS

The first natural wing frequency for the NACA 65A004 wing was about 260 cps and for the 5-percent hexagonal wing it was about 300 cps.

CONDITION



TABLE 2.- CONTROL-SURFACE DATA

(a) M = 1.3

fo	ft	ଞ _୦	8 _t	k	N5	₩ ₆	-Chi
Hexagonal wing; $\mu r_a^2 = 313$							
68 70. 68 72 89 101 101 133 133 145 146 178 180 185	182 182 184 183 188 193 191 215 213 222 217 240 242 245	0.0115 .0105 .0095 .0085 .0095 .011 .019 .034 .013 .014 .018 .014 .0165	0.0023 .0022 .007 .011 .0082 .0092 .007 .012 .015 .012 .010 .0115 .0074 .014	0.0550 .0555 .0552 .0568 .0585 .0578 .065 .0645 .0645 .0672 .0658 .0725 .073 .074	269 267 270 264 242 228 228 193 191 180 171 141 140 135	0.22 .19 1.75 2.98 1.97 2.07 1.25 1.47 .53 2.04 1.16 .50 09 1.44	0.024 .021 .194 .330 .224 .242 .144 .191 .068 .274 .153 .073 013 .213
		65A0	04 wing;	$\mu r_a^2 = 6$	50		
66 66 81 81 89 89 89 120 120 120 138 132 147 160 160	135 136 142 142 142 146 147 145 164 165 175 172 186 192 190	0.0085 .0070 .0109 .0097 .0091 .010 .010 .010 a.032 b.018 a.018 b.018 a.018 b.018 a.020 b.016 a.024 b.016	0.0068 .0030 .0085 .0075 .0052 .0093 .0070 .0050 .0055 .0072 .021 .021 .021 .021 .021 .021 .021 .02	0.0419 .0422 .0440 .0440 .0453 .0457 .0457 .0508 .0513 .0544 .0535 .0578 .0578 .0595 .0589	495 497 439 439 439 408 412 405 302 306 246 267 244 198 189	3.12 .85 3.25 2.80 1.43 3.64 2.80 1.04 .06 1.24 .65 10.40 6.50 10.40 6.50 4.03 9.10 3.90 8.05	0.262 .072 .286 .246 .126 .330 .256 .094 .006 .127 .071 .706 1.112 .695 .465 1.082 .465 1.082

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^aHigh amplitude. ^bLower amplitude.

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TABLE 2.- CONTROL-SURFACE DATA - Continued

(b) M = 1.6

fo	ft	ಕಂ	8 _t	k	₹	™ ₆	$-c_{h_{\delta}}$
	<u> </u>	Hexa	igonal wir	ıg; μr _a 2	= 427	-	
$\begin{array}{c} 70\\ 72\\ 71\\ 90\\ 90\\ 92\\ 92\\ 101\\ 102\\ 148\\ 146\\ 147\\ 189\\ 187\\ 213\\ 213\\ 233\\ 232\\ 232\end{array}$	160 163 160 170 171 170 170 172 176 210 201 202 202 233 231 233 250 251 265 260	0.014 .011 .0107 .0114 .0103 .0104 .010 .0103 .0103 .0103 .0103 .0115 .0097 .010 .0115 .0078 .0060 .0058 .0054 .0054 .0078 .0088	0.021 .020 .023 .021 .0205 .024 .0275 .0256 .0259 .0245 .024 .023 .0188 .0155 .018 .023 .018 .023 .019 .016 .0214 .020	0.0415 .0424 .0415 .0442 .0445 .0442 .0442 .0446 .0456 .0547 .0520 .0524 .0524 .0524 .0606 .0606 .0606 .065 .0654 .069 .0676	345 344 343 307 309 302 280 284 202 284 201 204 201 148 152 120 97 88	7.80 7.65 8.90 7.60 7.51 8.90 9.40 9.40 7.95 7.25 6.275 8.15 6.575	0.647 .648 .739 .671 .669 .786 .929 .838 .856 .920 .826 .759 .656 .600 .710 .989 .827 .675 .907 .750

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TABLE 2.- CONTROL-SURFACE DATA - Concluded

fo	ft	₿ ₀	8t	k	- N5	<u>N</u> 6	-C _h *
1	$65A004$ wing; $\mu r_a^2 = 886$						
52 52 53 53 66 65 58 80 80 90 90 88 153 153	$111 \\ 112 \\ 110 \\ 113 \\ 120 \\ 120 \\ 120 \\ 120 \\ 120 \\ 120 \\ 120 \\ 130 \\ 129 \\ 128 \\ 129 \\ 133 \\ 133 \\ 133 \\ 133 \\ 131 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 \\ 176 $	0.0078 .0080 .0073 .0067 .0063 .0085 .0085 .0083 .0076 .0079 .0060 .0063 .0062 .0063 .0055 .0069 .0047 .0040	0.016 .0154 .0172 .0147 .015 .0164 .0168 .0164 .0154 .0164 .0158 .0165 .0151 .0162 .0151 .0162 .0151 .0144 .0146	0.0296 .0298 .0294 .0301 .0321 .0318 .0321 .0321 .0347 .0344 .0354 .0354 .0354 .0354 .0354 .035 .047 .047	693 696 699 619 614 626 55 54 50 54 80 80 75 75 488 80 75 75 488 80 75 75 75 488 80 75 75 75 75 75 75 75 75 75 75 75 75 75	12.7 12.1 13.7 11.6 12.8 12.7 12.0 11.1 11.9 10.9 12.1 10.8 12.1 10.6 9.7 10.3	0.752 .723 .805 .705 .745 .815 .815 .770 .770 .770 .819 .744 .819 .856 .765 .856 .742 .910 .968
		65A00	04 wing;	$\mu r_{a}^{2} = 3$	78		
83 83 105 105 128 148 144 142 142 193 194 227 230	176 176 186 187 200 203 210 210 210 240 240 240 259	0.009 .009 .010 .010 .0073 .008 .0085 .008 .0078 .0065 .0064 .0058 .0063	0.026 .0223 .0224 .0252 .0215 .0213 .0224 .021 .0225 .021 .020 .0207 .0222	0.0467 .0493 .0495 .0532 .0558 .0559 .0559 .0552 .0637 .0637 .069 .0685	294 258 259 228 201 206 2034 131 91 80	9.13 7.72 7.26 8.33 7.00 6.85 6.96 6.55 7.15 6.36 5.98 6.17 6.50	0.852 .711 .717 .825 .745 .737 .737 .739 .732 .790 .811 .761 .851 .890
	65A004 wing; $\mu r_a^2 = 196$						
144 142 114 114 114	256 256 242 246	0.0075 .0076 .0065 .0079	0.030 .023 .0286 .032	0.0688 .0688 .065 .0658	134 135 152 154	5.40 4.05 5.32 5.93	0.743 •557 •691 •780

⁽b) M = 1.6

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Figure 1.- Sketch of wing and control surface used in tests.

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15 п 10 10 Q Q 5 Stable \diamond ₽₽₽ de la construction de la constru þog 0 Ν_θ <u>_</u>_____ Ñs ol ٥V Unstable - 5 -5 Hex. wing-mg control 0 65A004 wing-steel control 66A004 wing - mg control \diamond -Ю -10 65A004 wing - al. control Δ 3-D Theory (Ref.4) 2-D Theory (Ref. 3) ___ _ -15 -15 -204 -20 0 /03 .05 .06 .**0**7 ,04 .05 .08 .06 .07 k k

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Figure 3.- Variation of aerodynamic damping coefficient \overline{N}_6 with reduced frequency k at M = 1.3 and 1.6.

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Figure 5.- Variation of damping coefficient $C_{h_{\delta}^{\bullet}}$ with Mach number and comparison with results of reference 1 at k values of about 0.03.

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