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## AERODYNAMIC CHARACTERISTICS OF THIN EMPIRICAL PROFILES AND

THEIR APPLICATION TO THE TAIL SURFACES AND AILERONS OF AIRPLANES By A. Toussaint and E. Carafoli

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AERODYMAIC CHARACTERISTICS OF THIN EMPIRICAL PROFILES AND THEIR APPLICATION TO THE TAIL SURFACPS AND AILERONS OF AIRPLANES.* By A. Toussaint and ㅍ. Carafoli.

The increasing use of airplane wings equipped with ailerons and the importance of knowing the aerodynamic characteristics of tail units (empennages) provided with movable parts (rudders and elevators) impart some interest to the so-called "empirical profiles." This term is anplied to wing profilos which are not dravn according to any theoretical method. Generally these profiles are of no particular interest, since they can alweys be replaced by theoretical profiles winch satisly the same aerodynamic and structural conditions. For the abovementioned applications, however, the profiles, modified by the deflection of the aileron, elevator or rudder, necessarily become empirical profiles, of which it is important to know the aerodynamic characteristics and, above all, to kow how these characteristics are affected by the magnitude of the deflections.

For this puryose we thought best to employ the method proposed by Turk for the approximate theoretical study of thin, slightly curved profiles, assimilable, from an aerodynamic view*From L'aérophile, June, 1928, pp. 179-183.
point, to their mean camber line.

> Principle of the Method

Given a profile whose nean camber line $A S B$ is moderately curved and whose chord $A B=l$.

It is well known that, in the conformal transformation

$$
\begin{equation*}
z=\zeta+\frac{\mathrm{a}^{2}}{\zeta} \tag{1}
\end{equation*}
$$

a circle $C$, of radius $a=\frac{l}{4}$ and centered at 0 , 亏́ves $a$ straicht line $A B$ of length $l=4 a$.

On adopting the above transformation function (1), the line $A S B$, very close to the chord $A E$, will have for its antecodent in the $\zeta$ plane a curve $s^{\prime}$ which will lie very close to tho circle $C$, and which will out this circle at the points $A^{2}$ and $5^{1}$, antececients of $A$ and $B$.

The curve St may be defined by the expression $^{\text {m }}$

$$
\begin{equation*}
\xi^{\prime}=a(I+r) e^{i \theta} \tag{2}
\end{equation*}
$$

in which $r$ is the radial segment, as a function of $\theta$, measuring the smail distance between the circle $C$ and the curve S' along the radius vector of "argunent" $\theta$.

The meon line $A S B$ will be obtained by the transformation

$$
\begin{equation*}
z=\xi^{\prime}+\frac{a^{2}}{\xi^{\prime}} \tag{1'}
\end{equation*}
$$

which, on neglecting the torms in $r^{2}$, gives

$$
z=2 a \cos \theta+2 i \operatorname{ar} \sin \theta
$$

For the abscissa $X=2$ a $\cos \theta \quad\}$
the ordinate will be $Y=2$ a $r \sin \theta$
To transform the circle c into the curve $S^{\prime}$ so that a point $\zeta^{\prime}=a(1+r) e^{i \theta}$ of the curve $S^{\prime}$ will correspond to a point $\zeta=a e^{i \varphi}$ of the circle $c$, let

$$
\begin{equation*}
\varphi=\theta+\beta \tag{4}
\end{equation*}
$$

Under these conditions, the trailing edge $B$ of the profile $A S B$ will correspond to the point $B^{\prime},(\theta=\pi)$ of the curve $S^{\prime}$ and to the point $B^{\prime} \circ\left(\varphi=\pi+\beta_{0}\right)$ of the circle $C$. According to Joukowski's hypothesis, the point $B^{\prime} \circ$ must be a point of zero velocity. Hence, the circulation will have the value

$$
\begin{equation*}
\Gamma=4 \pi a V \sin \left(\alpha+\beta_{0}\right) \tag{5}
\end{equation*}
$$

The transformation function under the general form

$$
\begin{gather*}
\zeta^{\prime}=\zeta+a\left(A_{0}+i B_{0}\right)+a^{2} \frac{\left(A^{\prime}+i B^{\prime}\right)}{\zeta}+\ldots \\
+a \frac{A_{n}+i B_{n}}{\zeta_{n}} \tag{6}
\end{gather*}
$$

will enable us to pass from the circle $C$ to the curve $S^{\prime}$. on replacing $\zeta^{\prime}$ and $\zeta$ by their respective values,

$$
\zeta^{\prime}=a(I+r) e^{\dot{i} \theta} \text { and } \zeta=a e^{\dot{j} \varphi}
$$

and, since $\beta$ and $r$ are very small,

$$
\begin{aligned}
& r \cong \sum A_{n} \cos (n+1) \theta+\sum B_{n} \sin (n+1) \theta \\
& \beta \cong \sum A_{n} \sin (n+1) \theta-\sum B_{n} \cos (n+1) \theta
\end{aligned}
$$

Since the profile $A S B$ has no thickness, the ordinate

$$
y=2 \text { ar } \sin \theta
$$

takes the same value for $\pm \theta$. Hence, $r$ changes $\operatorname{sign}$ with $\theta$ and.

$$
A_{0}=A_{1}=A_{3}=\ldots \ldots=A_{n}=0
$$

This gives

$$
\left.\begin{array}{l}
r \cong \Sigma B_{n} \sin (n+I) \theta  \tag{?}\\
\beta \cong-\Sigma B_{n} \cos (n+I) \theta
\end{array}\right\}
$$

The coefficients $B_{n}$ are therefore defined by the formula

$$
\begin{equation*}
B_{n}=\frac{I}{\pi} \int_{0}^{2 \pi} r \sin (n+I) \theta d \theta \tag{8}
\end{equation*}
$$

The aerodynamic characteristics of the profile ASB are then given by the following formulas:

1. Lift $P=\rho V \Gamma=4 \pi$ a $\rho V^{2} \sin (\alpha+\beta)$
whence

$$
\begin{equation*}
C_{z} \cong 2 \pi\left(\alpha+\beta_{0}\right) \tag{9}
\end{equation*}
$$

2. Moment $M_{A}=M_{O}-P 2$ a $\cos \alpha$.

Leet $M_{A} \cong 4 \pi \rho V^{2}\left|X_{1}\right|(\alpha+\gamma)-8 \pi a^{2} \rho V^{2}\left(\alpha+\beta_{0}\right)$
whence

$$
\begin{equation*}
c_{m}(A) \cong \frac{\pi}{2} \frac{\left|x_{1}\right|}{a^{2}}(\alpha+\gamma)-\pi\left(\alpha+\beta_{0}\right) \tag{10}
\end{equation*}
$$

$X_{1}=\mathrm{C}^{2} i^{2 i \gamma}$ being the term in $\frac{1}{\zeta}$ of the transformation fundtion (6). Thus, we have
N.A.C.A. Technical Memorandum No. 493

$$
\begin{equation*}
X_{1}=C^{2} e^{2 i \gamma}=a^{2}\left(I+i B_{1}\right) \tag{II}
\end{equation*}
$$

The computation of $C_{z}$ and $C_{m}(A)$ leads therefore to the calculation of $\beta_{o},\left|x_{1}\right|$ and cf $\gamma$.

$$
\text { Calculation of } \beta_{0}
$$

This angle corresponds to the point $B_{0}{ }^{3}$, that is, to . $\theta=\pi$. Under these conditions the formulas (7) give

$$
\beta_{0}=-\sum B_{n} \cos (n+1) \pi=B_{0}-B_{1}+B_{2}-B_{3}+\ldots
$$

Replacing $B_{0}, B_{1}, B_{2} \ldots$ etc., by their values derived from formula (8), we obtain

$$
\beta_{0}=\frac{1}{\pi} \int_{e}^{2 \pi} r(\sin \theta-\sin 2 \theta+\sin 3 \theta \ldots) d \theta
$$

cr

$$
\beta_{0}=p \cdot i m \cdot \frac{I}{\pi} \int_{\theta}^{2 \pi} \frac{r e^{i \theta}}{1+e^{i \theta}} d \theta
$$

which reduces to

$$
\begin{equation*}
\beta_{0}=\frac{1}{\pi} \int_{\theta}^{2 \pi} \frac{r \sin \theta}{21+\cos \theta} d \theta=\frac{1}{2 \pi a} \int_{\theta}^{\pi} y \frac{d \theta}{1+\cos \theta} \tag{12}
\end{equation*}
$$

In practice the profile abscissas and ordinates are refired to the axes $x^{\prime} A y^{\prime}$ and the relative coordinates

$$
\left\{\begin{array}{l}
\xi=\frac{x^{\prime}}{4 a}=\frac{1}{2}(1-\cos \theta) \\
\eta=\frac{y^{t}}{4 a}=\frac{y}{4 a}
\end{array}\right.
$$

are preferably considered.
N.A.C.A. Technical Memorandum No. 493

Thus equation (12) gives for $\beta_{0}$. the value

$$
\begin{equation*}
\beta_{0}=\frac{1}{\pi} \int_{Q}^{1} \eta \frac{d \xi}{(1-\xi) \sqrt{\xi(1 \xi)}}=\int_{0}^{1} \eta f_{1}(\xi) d \xi \tag{13}
\end{equation*}
$$

by letting

$$
\begin{equation*}
f_{1}(\xi)=\frac{1}{\pi(1-\xi) \sqrt{\xi(1-\xi)}} \tag{}
\end{equation*}
$$

When the mean line $A S B$ of the profile is known, equation
(13) can be integrated by employing computed values of $f_{1}(\xi)$.

$$
\text { Calculation of }\left(X_{1}\right) \text { and of } \gamma
$$

Since $B_{1}$ is small, formula (II) gives

$$
C^{2}=\left|X_{1}\right|=a^{2}\left(1+B_{1}^{2}\right) \cong a^{2}
$$

Similarly

$$
\tan 2 \gamma \cong 2 \gamma=B_{1}
$$

The value of $B_{1}$ is found by equation (8), which gives

$$
B_{1}=\frac{1}{\pi} \int_{0}^{2 \pi} r \sin 2 \theta d \theta=\frac{2}{\pi a} \int_{0}^{\pi} y \cos \theta d \theta
$$

Letting $\mu_{0}=\frac{\pi}{4} B_{1}$ and employing the relative coordinates $\xi$ and $\eta$, we obtain
with

$$
\begin{equation*}
\mu_{0}=2 \int_{0}^{1} \eta \frac{(1-2 \xi) d \xi}{\sqrt{\xi(1-\xi)}}=2 \int_{0}^{1} \eta f_{2}(\xi) d \xi \tag{14}
\end{equation*}
$$

- As before, the values of $f_{z}(\xi)$ can be calculated for different abscissas $\xi$, and the integration of equation (14) will be accomplished according to the values of $\eta$ obtained on the mean camber line ASE. Expression (10) for $C_{m(A)}$ then becomes

$$
\begin{gather*}
C_{m(A)}=\frac{\pi}{2}\left(\alpha+\frac{B_{1}}{2}\right)-\pi\left(\alpha+\beta_{0}\right)=-0.25 C_{2}+C_{m_{0}}  \tag{1}\\
C_{m_{0}}=\frac{\pi}{4} B_{1}-\frac{\pi}{2} B_{0} \tag{10"}
\end{gather*}
$$

In short, the aerodynamic characteristics of empirical profiles, assimilable to their mean camber line, due to their relative thickness and curvature, can be calculated from the formulas

$$
\left\{\begin{array}{l}
\mathrm{C}_{\mathrm{z}} \cong 2 \pi\left(\alpha+\beta_{0}\right)=2 \pi\left[\alpha+\int_{0}^{1} \eta f_{1}(\xi) d \xi\right]  \tag{91}\\
\mathrm{C}_{\mathrm{m}} \cong-0.25 \mathrm{C}_{\mathrm{Z}}+\mathrm{C}_{\mathrm{m}_{0}}=-0.25 \mathrm{C}_{\mathrm{z}} \\
\quad+2 \int_{0}^{1} \eta f_{2}(\xi) d \xi-2 \int_{0}^{1} \eta f_{1}(\xi) d \xi
\end{array}\right.
$$

for which the function $f_{1}(\xi)$ and $f_{2}(\xi)$ are given in the table below as functions of $\xi$.

| $\xi$ |  | 0.025 | 0.05 | 0.10 | 0.20 | 0.30 | 0.40 | 0.50 |
| :--- | :--- | :--- | :--- | :---: | :---: | :---: | :---: | :---: |
| $f_{1}(\xi)$ | 2.09 | 1.54 | 1.18 | 1.00 | 0.99 | 1.08 | 1.27 |  |
| $f_{2}(\xi)$ | 6.10 | 4.13 | 2.67 | 1.50 | 0.87 | 0.41 | 0 |  |
| $\xi$ |  | 0.60 | 0.70 | 0.80 | 0.90 | 0.95 |  |  |
| $f_{1}(\xi)$ | 1.62 | 2.31 | 3.98 | 10.6 | 29.2 |  |  |  |
| $f_{2}(\xi)$ | -0.41 | -0.87 | -1.5 | -2.67 | -4.13 |  |  |  |

## Profiles Having Any Iiean Camber Line

If the nean camber line of the profile cannot be represented by a simple relation between $\eta$ and $\xi$, the integrations

$$
\int_{0}^{1} \eta f_{1}(\xi) d \xi \text { and } \int_{0}^{1} \eta f_{z}(\xi) d \xi
$$

can be made either graphically or mechanically, by making the products $\eta f_{1}$ and $\eta f_{2}$ according to the values of $\eta$ derived from the mean camber line. In general, the value of $\eta f_{1}$ ir this calculation increases indefinitely toward the trailing edge (i.e., for $\xi=1$ ). We can then evaluate the integral $\eta f_{1} d \xi$, from $\xi=0$ to $\xi=0.95$, and estimate the portion of the integral from $\xi=0.95$ to $\xi=1.0$. Assuming that this portion of the mean camber line is rectilinear, we find that the additional contribution ainounts to $2.9 \eta^{\prime}$, $\eta^{\prime}$ being the relative ordinate for $\xi=0.95$ (according to Glauert).

## Profiles with Particular Mean Camber Lines

When the form of the mean camber line can be expressed by a simple ratio between $\eta$ and $\xi$, the values of $\beta_{0}$ and $\mu_{0}$ can be calculated by direct integration. For example, the equation

$$
\eta=b \xi(I-\xi)(c-\xi)
$$

represents a profile whose mean camber line is of double curvature when $c$ varies between $\frac{1}{2}$ and 1 . Direct integration gives
N.A.C.A. Technical Memorandum No. 493

$$
\beta_{0}=\frac{\dot{b}}{8}(4 c-3)
$$

and

$$
\mu_{0}=\frac{\pi b}{32},
$$

from which

$$
c_{m_{0}}=\frac{\pi b}{32}(7-8 c) \text { is derived }
$$

For $\quad c=\frac{7}{8}$ we get $\mathrm{c}_{\mathrm{m}_{\mathrm{O}}}=0$.

Application to Tail Units with Movable Parts

For a stabilizer having a symmetrical biconvex profile, the deflection of the elevator $H B$ by a certain moderate angle $\beta$, forms a profile whose mean camber line is composed of two straight lines, AH and HB.

Let $\sigma$ be the ratio of the movable part $H B$ to the total surface area $A H+H B$, then:

$$
\sigma=\frac{H^{\prime} B}{A} ; \quad \xi(H) \cong \frac{A H^{\prime}}{A} \bar{B}=1-\sigma \text { and } h=\frac{H^{\prime} H}{A B}=\sigma(1-\sigma) \beta
$$

For $\xi$ varying from 0 to ( $1-\sigma$ ) we get

$$
\eta=\frac{h}{1-\sigma} \xi
$$

and for $\xi$ varying from ( $1-\sigma$ ) to 1 we get

$$
\eta=\frac{h}{\sigma}(1-\xi) .
$$

Under these conditions the values of $\beta_{0}$ and $\mu_{0}$ are calculable by direct integration of formulas (13) and (14).

Thus we get

$$
\begin{aligned}
& \beta_{0}= \frac{1}{\pi} \int_{0}^{1 \sigma} \frac{h}{1-\sigma} \xi \frac{d \xi}{(1-\xi) \sqrt{\xi(1-\xi)}} \\
&+\frac{1}{\pi} \int_{1-\sigma}^{1} \frac{h}{\sigma}(1-\xi) \frac{d \xi}{(1-\xi) \sqrt{\xi(1-\xi)}} \\
& \mu_{0}=2 \int_{0}^{1-\sigma} \frac{h}{1-\sigma} \xi \frac{(1-2 \xi) d \xi}{\sqrt{\xi(1-\xi)}} \\
&+2 \int_{1-\sigma}^{1}(1-\xi)-\frac{h(1-2 \xi) d \xi}{\xi(1-\xi)}
\end{aligned}
$$

and finally

$$
\begin{align*}
& \beta_{0}=\frac{2 h}{\pi}\left[\frac{1}{\sqrt{\sigma(1-\sigma)}}+\frac{\pi}{2 \sigma}-\frac{\operatorname{arc} \cos \sqrt{\sigma}}{\sigma(1-\sigma)}\right]  \tag{15}\\
& y_{0}=h\left\{\frac{2 \sigma-1}{\sqrt{\sigma(1-\sigma)}}+\frac{\pi}{2 \sigma}-\frac{\operatorname{arc} \cos \sqrt{\sigma}}{\sigma(1-\sigma)}\right\} \tag{15}
\end{align*}
$$

and

Aerodynamic Characteristics

Here we have:

$$
C_{z}=2 \pi\left(\alpha+\beta_{\dot{O}}\right)
$$

In practice, the lift coefficient is expressed as a fundtion of the deflection $\beta$ of the elevator and of the angle of attack of the relative wind on the stabilizer. AH. Therefore, we may write

$$
\begin{equation*}
c_{z}=2 \pi\left(\alpha_{1}+m_{0} \beta\right) \tag{17}
\end{equation*}
$$

N.A.C.A. Technical Memorandum No. 493

Since $h \cong \sigma(1-\sigma)$, and $\beta$ and $\alpha=\alpha_{1}+\gamma_{1}=\alpha_{1}+\sigma \beta$, we finally obtain

$$
C_{z}=2 \pi\left[\alpha_{1}+\left(1+\frac{2}{\pi} \sqrt{\sigma(1-\sigma)}-\frac{2}{\pi} \operatorname{arc} \cos \sqrt{\sigma)} \beta\right]\right.
$$

hence

$$
\begin{equation*}
m_{0}=1-\frac{2}{\pi}[\operatorname{arc} \cos \sqrt{\sigma}-\sqrt{\sigma(1-\sigma)}] \tag{18}
\end{equation*}
$$

Figure 3 represents the variation of $m_{0}$ with respect to $\sigma$ according to formula (18). The experimental results of different aerodynamic laboratories, for deflections $\beta$ varying between $\pm 5^{\circ}$ and $\pm 15^{\circ}$, are here shown. The locations of these representative points with reference to the theoretical curve show that the value of the coefficient $m$, applicable in the formula

$$
I_{z}=2 \pi(\alpha+m \beta),
$$

diminishes somewhat with the magnitude of the deflection $\beta$. As regards this analysis, we need only to call attention to the fact that the experimental points very close to the theoretical curve apply to elevators extending over the whole tail span without any central cut-out. When there is such a cut-out, the experimental values of $m$ are somewhat snaller than the theoretical values, and the empirical expression

$$
\begin{equation*}
m=\sqrt{\sigma} \tag{19}
\end{equation*}
$$

very closely represents the mean values applicable to moderate deflections $\left(\beta<10^{\circ}\right)$.

Similarly the moment coefficient $C_{m}(A)$ is finally obtained as

$$
\begin{equation*}
C_{m(A)}=-0.25 C_{z}-2 \sqrt{\sigma(1-\sigma)^{3}} \beta \tag{20}
\end{equation*}
$$

The experimental data available for verifying this theoretical formula are, however, very few.

Application to Wings Provided with Ailerons

The preceding results may be extended to the case of supporting wings equipped with ailerons for the purpose of increasinc the aaximum lift coefficient.

Let ASB represent the rean camber line of a wing profile, the rear portion $O f$ which $H B$ forms an aileron rotating about the hings $H$ (Fig. 4). The aerodynamic characteristics of the undisturbed nean camber line are given by formulas (9') and (10'), as previously demonstrated.

Let $\beta$ represent the angle of deflection, such that the mean camber line becomes AHB. As previously shown, the approximate theory is applicable to moderate values of deflection $\beta$, so that the chord $A B^{\prime}$ of the modified profile is very nearly equal to the chord $A B$ of the original profile.

The aerodynamic characteristics of the profile ASB: will be calculable by the application of formulas (13) and (14) to the equation of the line ASB' with respect to the axes ox' and oy', $\gamma$ being the angle which ox makes with ox'. Let
N.A.C.A. Technical Memorandum No. 493
$\sigma=\frac{H B}{A B}$ then $\gamma \cong \sigma \beta$. For the portion $A H$, characterized by the relative abscissa $\xi=(1-\sigma)$ we then get:

$$
\left\{\begin{array}{l}
\xi^{\prime}=\xi \cos \gamma-\eta \sin \gamma \cong \xi \\
\eta^{\prime}=\eta \cos \gamma+\xi \sin \gamma \cong \eta+\xi \sigma \beta
\end{array}\right.
$$

In an analogous manner for the portion $H B^{\prime}$ we get:

$$
\begin{aligned}
& \xi^{\prime}=\xi \\
& \eta^{\prime} \cong \eta+(I-\sigma) \beta(I-\xi)
\end{aligned}
$$

The calculation of $\beta_{O}^{\prime}$ and $\mu_{0}^{\prime}$ is in this case made by formulas (13) and (14) applied between 0 and ( $1-\sigma$ ) for the portion $A H$ and between $(I-\sigma)$ and $I$ for the portion HB'. We thus find
and

$$
\begin{align*}
\beta_{0}^{\prime} & =\int_{c}^{1} \eta f_{1}(\xi) \alpha \xi+\sigma \beta \int_{c}^{1-\sigma} \xi f_{1}(\xi) d \xi \\
& +(1-\sigma) \beta f_{1-0}^{1}(1-\xi) f_{1}(\xi) d \xi \tag{131}
\end{align*}
$$

$$
\mu_{0}=2 \int_{0}^{1} \eta f_{2}(\xi) d \xi+2\left[\sigma \beta \int_{0}^{2-\sigma} \xi f_{2}(\xi) d \xi\right.
$$

$$
\begin{equation*}
\left.+(I-\sigma) \beta \int_{1-\sigma}^{1}(1-\xi) f_{z}(\xi) d \xi\right] \tag{1}
\end{equation*}
$$

The first integral in $\beta_{0}{ }^{\prime}$ represents the angle of zero Iii $\beta_{0}$ of the profile ASB. The other two integrals ropesent the angle of zero lift of the deflection $\beta$ of the levator. This angle being $\Delta \beta_{O}$, we have

$$
c_{z(\Omega)}=2 \pi\left(\alpha+\beta_{0}+\Delta \beta_{0}\right)
$$

$\beta_{0}$ is calculable from the equation $\eta=f(\xi)$ of the initial mean camber line ASB.
$\Delta \beta_{0}$ can be calculated as already described for the movable elevator, i.e., by the formula
with

$$
\begin{aligned}
\Delta \beta_{0} & =\left(m_{0}-\sigma\right) \beta \\
m_{0} & =1-\frac{2}{\pi}[\operatorname{arc} \cos \sqrt{\sigma}-\sqrt{\sigma(1-\sigma)}]
\end{aligned}
$$

If the chord $A B$ of the original profile is taken as the reference line, the angle of attack is then $\alpha_{1}=\alpha-\gamma \cong \alpha-v \beta$ and the result is

$$
c_{z}=2 \pi\left(a_{1}+\beta_{0}+m \beta\right) .
$$

The first integral in the expression for $\mu_{0}{ }^{\prime}$ represents the value of $\mu_{0}$ for the original profile. The other two integrals represent the correction $\Delta \mu_{0}$ due to the deflection of the ailcron. $\Delta \mu_{0}$ will be calculated as in the case of the elevator, and we have

$$
c_{m_{0}(\beta)}=c_{m_{0}(\beta=0)}-\left(\frac{\pi}{2} \Delta \beta_{0}-\Delta \mu_{0}\right)
$$

or

Experiments with wings equipped with ailerons verify quite well the theoretical values of $\Delta \beta_{0}$, in particular for positive moderate deflections $\beta$. In general, the values of
N.A.C.A. Technical Memorandum $\mathbb{N Q} \cdot .493$.

$$
\Delta C_{m_{0}}=2 \sqrt{\sigma(1-\sigma)^{3}} \beta
$$

under the same conditions, are somewhat larger than those found experimentally. The deviation is similar to the one found in the case of wing profiles having an appreciable $\mathrm{C}_{\mathrm{m}_{\mathrm{O}}}$.

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Figs.1, 2, 3,4


Fig. 1


Fig. 4

