# A Study of DC-DC Converters With MCT's for Arcjet Power Supplies 

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Abstract<br>A Study of DC-DC Converters with MCT's<br>for Arcjet Power Supplies<br>by<br>Thomas A. Stuart

Many arcjet DC power supplies use PWM full bridge converters with large arrays of parallel FETs . This report investigates an alternative supply using a variable frequency series resonant converter with small arrays of parallel MCTs. The reasons for this approach are to: 1) Increase reliability by reducing the number of switching devices; 2) Decrease the surface mounting area of the switching arrays. The variable frequency series resonant approach is used because the relatively slow switching speed of the MCT precludes the use of PWM. The 10 kW converter operated satisfactorily with an efficiency of over $91 \%$. Test results indicate this efficiency could be increased further by additional optimization of the series resonant inductor.

## I. INTRODUCTION

Many of the present DC power supplies for arc jet thrusters utilize a full bridge DC-DC converter similar to those in [1-3]. These circuits employ four switches that usually consist of multiple FETs connected in parallel, and phase shift control is used to regulate the output. The operating frequency is usually limited to the $10-20 \mathrm{kHz}$ range where switching loss is very low because of the fast switching times of the FETs. Higher frequencies have been avoided because large ferrite transformer cores have not been considered satisfactory for a spacecraft environment. Most of these converters have been intended for lower power applications (e.g., 1-2 kW) which are supplied from a 28 Vdc source. At this voltage, no other switching device is capable of providing the efficiency and cost effectiveness of FETs.

However, for higher power levels such as 10 kW , it becomes desirable to increase the source voltage to $150-300 \mathrm{Vdc}$ in order to avoid the exceptionally high currents that would accompany a 28 Vdc source. At this higher voltage level, other switching devices such as the IGBT and the MCT become competitive with the FET. This is because the forward voltage drop of the FET climbs rapidly for breakdown voltages above 200 volts, and a very large number of devices must be operated in parallel to achieve the forward drop of an IGBT or an MCT. If the application is not cost sensitive, it is still feasible to operate FETs in this manner, but very large arrays of FETs are required. This is especially true when attempting to match the forward drop of an MCT, which promises to have a much lower forward drop than an IGBT with the same die size [4]. Because MCTs offer the prospect of a large reduction in the number of switching devices, it was decided to investigate the use of these devices in a 10 kW arcjet power supply.

## A. MCTs vs. FETs

The MCT module described in [4] consisted of 6 parallel die and had a nominal current rating of 300 amps . This device was a prototype MCT power module developed by the General Electric Corporate Research and Development Center.

A 10 kW converter with an input voltage of 150 Vdc (specified input voltage when research proposal was submitted) and an efficiency of $90 \%$ will draw an average input current of 74 Adc . At 74 amps and $\mathrm{T}_{\mathrm{J}}=125^{\circ} \mathrm{C}$, the GE MCT has a forward drop of about 0.75 volts. No breakdown voltage ratings were given in [4] for this device, but it is presumed to be about 600 volts since it was operated from a 270 Vdc bus and compared with a 600 volt IGBT.

To achieve this same drop at 74 amps with FETs, one might consider the IRFK4H450 module, which has an $\mathrm{R}_{\text {dson }}=0.19 \Omega$ at $\mathrm{T}_{\mathrm{J}}=125^{\circ} \mathrm{C}$ and consists of 4 parallel die. This device has a peak voltage rating of 500 volts. To get a drop of 75 V. © 74 A would require 19 parallel modules, or 76 parallel die. This means that a full bridge converter would need a total of $4 \times 76=304$ die. While it is feasible to construct such a circuit, it should be remembered that a short circuit failure of any single die will destroy the entire power supply unless each die has an individual fuse. Packaging restrictions may be another problem, and the drive circuits will need a high impulse current capability because of the high input capacitance of the FET arrays. All of these problems would be greatly alleviated by using MCTs, since a circuit with the same forward drop would contain only 6 x $4=24$ die .

MCTs are much slower than FETs however, so the MCT would not be competitive in a conventional 20 kHz phase shift circuit because of its higher switching loss. The GE MCT in [4], for example, has a current fall time of 1714 n.s. as compared to $70 \mathrm{n} . \mathrm{s}$. for the IRFK4H450. Therefore the MCT would have to
be used in some type of resonant or soft switching circuit where switching losses are reduced.

## B. Resonant vs Soft Switching Converters

When this project was first initiated, it appeared that one of the better candidate circuits might be the ZVS/ZCS (zero voltage switching, zero current switching) converter shown in Fig. I-1. Earlier versions of this circuit are described in [3, 5-8]. These earlier circuits are satisfactory for power levels in the $1-2 \mathrm{~kW}$ range, but they are difficult to implement at higher power levels. The more recent version in Fig. I-1 overcomes this limitation by using an active snubber which recycles part of the energy back to the source. To date, a 10 kW 20 kHz version of the circuit has achieved an efficiency of $91.3 \%$, even with high voltage drop IGBTs and a relatively low output voltage of 36 Vdc [9].

As seen from the timing diagram in Fig. I-1, on a typical half cycle for Q1 and Q4, both devices turn on simultaneously, but Q4 turns off first. When Q4 turns off, C 1 provides a path for the load current and $\mathrm{vCl}_{1}$ rises at a gradual rate. Thus C1 provides zero voltage switching (ZVS) for Q4. Assuming Q2 has reverse blocking capability, $\mathrm{V}_{\mathrm{C} 1}$ continues to increase beyond $\mathrm{V}_{\mathrm{s}}$, but when $\mathrm{v}_{\mathrm{C} 1}=\mathrm{V}_{\mathrm{B}}$, D 5 conducts and $v_{C 1}$ is clamped at this value. Since $V_{B 1} \gg V_{s}, i_{1}$ is rapidly driven to zero, and zero current switching (ZCS) is achieved for Q1 since $i_{1}=0$ before Q1 is turned off. After il reaches zero, a very slight reverse current flows through Q1 before it recovers. D6 is required to clamp the voltage across Q3 when Q1 recovers.

Although the circuit in Fig. I-1 has proven to be advantageous for higher input voltages, two developments occurred which indicated that it probably would not be competitive for the present application. First of all, the input voltage specification decreased from 150 to 120 Vdc , and there was some indication that it might drop even further. The second was that the available MCTs had lower than anticipated reverse breakdown voltages, and this would have required the use of

(a.) Power circuit with active snubbers.

(b.) timing diagram for gate drivers.

FigI-1. ZVS/ZCS $\underset{4}{\text { converter with MCT's }}$
series blocking diodes D1-4, as shown in Fig. I-2. The lower input voltage plus the extra drop of the diodes indicated that the ZVS/ZCS probably would have an excessive conduction loss as compared to other topologies. Therefore, it was decided to use the conventional series resonant circuit in Fig. I-3 instead of the ZVS/ZCS.

## C. Series Resonant Converters

The series resonant circuit (SRC) in Fig. I-3 is well known in the power electronics community, and it has been analyzed extensively [10-17]. If slower devices such as IGBTs or MCTs are used for Q1-Q4, the circuit should be controlled by varying the frequency (instead of PWM) in order to reduce the turnoff switching losses. D5 and D6 also have very low turn-off loss, but D1-D4 still have significant turn-off loss. All of the conduction losses will be higher than for conventional phase shift PWM because the current has a higher RMS/average ratio. This indicates an increased disadvantage for FETs in this circuit because of their $\mathrm{I}^{2} \mathrm{R}$ loss characteristic. In spite of the advantages for Q1-Q4 and D5-D6, it also should be noted that the higher AC currents are a disadvantage for all of the power level capacitors and magnetics because of their higher losses.

Although the SRC exhibits higher conduction losses, the 15 kHz MCT version achieved an efficiency of $91 \%$ at 10 kW . However, it should be noted that this was accomplished only after an extensive re-design of T 1 and $\mathrm{L}_{\mathrm{r}}$, and the temperature rise of $L_{r}$ in particular was still higher than expected. This proved to be the case even when foil conductors or cables with \#22 magnet wire were used (although the individual strands in this cable were insulated, it would not be proper to call this a Litz cable since the strands were not interwoven). Since these conductors appear to be satisfactory for the skin depth at 15 kHz , the excessive heating appears to be related to proximity effect losses and possibly heat transfer problems associated with larger cables. These heating problems did not appear to be caused


Fig I-2. ZVS/ZCS converter with MCT's and blocking diodes.


Fig I-3 Series resonant converter with MCTs.
by core losses however, since all calculations indicated that these should have remained well below the acceptable range. The final $L_{r}$ design operated satisfactorily, but this was achieved only after resorting to a three-turn three-layer winding. Naturally, this also produced a core that was larger than desirable. It appears the subjects of proximity effect losses and heat transfer for magnetic components in this power range deserve further attention, and additional research in this area is certainly warranted.

## II. POWER CIRCUIT OF THE 10-KW ARCJET POWER CONVERTER

## A. Series Resonant Converter (SRC) Description

Fig. II-1 shows the schematic of the SRC power circuit which was developed for this study. The variable frequency control circuit is shown in Fig. II-2, while the gate pulse logic is shown in Fig. II-3. The gate drivers for the MCTs are shown in Fig. II-4.

Since this converter is intended to power an arcjet thruster, it is designed to operate as a current regulator. A voltage limit loop is also required to prevent an excessive output voltage under light loading conditions. Variable frequency control is provided by the LM 566 voltage controlled oscillator (VCO) in Fig. II-2. The current regulation loop is controlled by the circuitry associated with op amps A1 and A2, while the voltage limit loop is controlled by the circuitry associated with A3.

The gate drive logic in Fig. II-3 can be explained with the aid of the waveforms in Fig. II-5. Fig. II-5 (a.) represents continuous current operation, while Fig. II-5 (b.) shows the discontinuous mode. Referring to Fig. 5 (a), as the switching frequency approaches the resonance frequency, it becomes necessary to decrease the width of the Q14 and Q23 drive signals while maintaining a constant deadtime, $\mathrm{t}_{\delta}$. However, in the discontinuous mode represented in Fig. II-5 (b.), Q14 and Q23 should be limited to some value, $\mathrm{t}_{\mathrm{n}}$, to avoid re-triggering the same switches at the end of each half cycle. The maximum $t_{n}$ is determined by the upper one-shot, while the minimum $\mathrm{t}_{\delta}$ is determined by the lower one-shot.

The isolated gate driver used with the MCTs is shown in Fig. II-4. A discrete design was used for this gate driver to allow increased flexibility over the available integrated gate drivers. It should be noted that the MCTs require a negative gateanode voltage (the MCT is characterized on its data sheet with -10 V ) for turn on. The manufacture recommends a positive gate-anode voltage (the MCT is


Fig II-1 Power circuit of $10-\mathrm{kW}, 100-120 \mathrm{~V}$ arcjet power converter.


Prevents Porositic Osc
"Sets 4.5 6 V5 6 6.1
$f_{0}=\frac{2\left(V^{+}-\sqrt{51}\right.}{\mathrm{RICIV}^{+}}$

Fig II-2 Variable frequency control circuit for SRC.


Fig II-3. Gate pulse logic for SRC.


Fig. II-4(a) MCT gate driver circut. One channel shown.


Fig. II-4(b) Isolated power supply for two channels of the MCT gate driver.

(a) Continuous operation

Fig II-5 Logic timing diagram

(b) Discontinuous operation

Fig II-5 Logic timing diagram (cont.)
characterized with +15 V ) to ensure turn off. Voltages of + and -15 V were used in the gate driver with good results (the MCT maximum gate-anode voltage is 20 V ). Isolated gate drive is supplied by a low-capacitance $60-\mathrm{Hz}$ transformer; the gate signals are transferred through a high speed opto-isolator (HP2602). The problem of coordinating the application of gate drive signals and the anode voltage to the MCTs was not addressed in this gate driver. However, this is an issue which would affect the reliable use of MCTs, because the MCT cannot be assumed to be off with zero gate voltage.

## B. Design of SRC

Fig. II-7 shows plots of the normalized continuous load current, $I_{\text {AN }}$, vs. the diode conduction angle, $\alpha$, for various ratios of $q$, where $q$ is the ratio of the reflected output voltage to the input voltage. These curves are based on the following equations which are derived in [14],

$$
\begin{equation*}
I_{A N}=\frac{2(1+q)(1-\cos \alpha)}{(\alpha+\beta)(q-\cos \alpha)} \tag{II-1}
\end{equation*}
$$

where $q=\frac{n V_{0}}{V}$
$\mathrm{n}=$ primary to secondary ratio
of T1 in Fig. I-3.
$\alpha=$ diode conduction angle in Fig. II-5 (a)
$\beta=$ MCT conduction angle in Fig. II-5 (a), which can be found from (II-3) below
$\alpha$ (also called the delay angle) is the independent variable which is varied to control $I_{A N}$. Actually this is done indirectly since the controller actually varies the angle, $\gamma$, where,

$$
\begin{equation*}
\gamma=\alpha+\beta \tag{II-2}
\end{equation*}
$$

Since $\beta$ is a function of $\alpha$ (see II- 3 below), controlling $\gamma$ also controls $\alpha$.
For continuous current,

$$
\begin{equation*}
\beta=\pi+\tan ^{-1}\left[\frac{\left(q^{2}-1\right) \sin \alpha}{2 q-\left(1+q^{2}\right) \cos \alpha}\right] \tag{II-3}
\end{equation*}
$$

rad.

To find the actual average current, $\mathrm{I}_{\text {avg }}$, use the equation,

$$
\begin{equation*}
\mathrm{I}_{\mathrm{avg}}=\mathrm{I}_{\mathrm{AN}} \times \mathrm{I}_{\mathrm{B}} \quad \mathrm{Amps} \tag{II-4}
\end{equation*}
$$

where $I_{B}=$ base current $=\frac{V_{s}}{Z}$ Amps

$$
\mathrm{Z}=\sqrt{\frac{\mathrm{L}_{\mathrm{r}}}{\mathrm{C}_{\mathrm{r}}}} \quad \Omega
$$

In this case,

$$
\mathrm{V}_{\mathrm{s}}=120 \mathrm{Vdc} \pm 10 \%(108-132 \mathrm{Vdc})
$$

Assume a range of $104-128 \mathrm{Vdc}$ due to series voltage drops. Also assume,

$$
\mathrm{q}=0.9
$$

Therefore, © low line, reflected $\mathrm{V}_{0}=.9 \times 104 \cong 94 \mathrm{~V}$.
Assume,

$$
\begin{aligned}
& \mathrm{V}_{\text {sec. }} \cong 200+2.4=202.4 \mathrm{~V} . \\
& (2.4 \mathrm{~V} . \cong \text { series voltage drop on secondary }) \\
\therefore \quad & \mathrm{n}=\frac{94}{202.4}=.46
\end{aligned}
$$

Assuming an efficiency of $90 \%$,

$$
P_{\mathrm{IN}}=\frac{10 \mathrm{~kW}}{.9}=11.1 \mathrm{~kW}
$$

Therefore,

$$
\mathrm{I}_{\text {avg }}=\frac{11.1 \mathrm{~kW}}{104 \mathrm{~V} .}=107 \mathrm{Adc}
$$

From Fig. II-7 for $\alpha=35^{\circ}$ and $\mathrm{q}=.9$ we get,

$$
\mathrm{I}_{\mathrm{AN}} \cong 2.5
$$

Therefore,

$$
\begin{aligned}
& \mathrm{I}_{\mathrm{B}}=\frac{\mathrm{I}_{\mathrm{avg}}}{\mathrm{I}_{\mathrm{AN}}}=\frac{107}{2.5}=42.8 \mathrm{~A} . \\
& \mathrm{Z}_{0}=\frac{\mathrm{V}_{\mathrm{S}}}{\mathrm{I}_{\mathrm{B}}}=\frac{104}{42.8}=2.43 \Omega
\end{aligned}
$$

Therefore,

$$
\frac{\mathrm{L}_{\mathrm{r}}}{\mathrm{C}_{\mathrm{r}}}=(2.43)^{2}=5.9
$$

We can now derive the following relationships between the switching frequency, $f_{s}$, and the natural resonance frequency, $f_{0}$. From Fig. II-5(a) for continuous current,

$$
\begin{equation*}
2 \gamma=\omega_{0} T_{s} \tag{II-5}
\end{equation*}
$$

where $T_{s}=1 / f_{s}$

$$
\omega_{0}=\frac{1}{\sqrt{L_{r} C_{r}}}
$$

Therefore,

$$
\begin{align*}
& \omega_{0}=2 \gamma f_{\mathrm{s}}  \tag{II-6}\\
& \text { or } \quad \mathrm{f}_{0}=\frac{\gamma}{\pi} \mathrm{f}_{\mathrm{s}} \tag{II-7}
\end{align*}
$$

Equations (II-2) and (II-3) can be used to plot $\gamma$ vs. $\alpha$ for various $q$ as in Fig. II-6. From Fig. II-6 for $q=.9$ and $\alpha=35^{\circ}$ we get,

$$
\gamma \equiv 196^{\circ}
$$

Therefore,

$$
\begin{aligned}
& \mathrm{f}_{\mathrm{o}}=\frac{196}{180} \mathrm{f}_{\mathrm{s}}=\frac{196}{180} \times 15 \times 10^{3}=16.33 \mathrm{kHz} \\
& \omega_{0}=2 \pi \mathrm{f}_{\mathrm{o}}=102.6 \times 10^{3} \mathrm{rad} .=\frac{1}{\sqrt{\mathrm{~L}_{\mathrm{r}} \mathrm{C}_{\mathrm{r}}}}
\end{aligned}
$$

Therefore,
$\mathrm{L}_{\mathrm{r}} \mathrm{C}_{\mathrm{r}}=94$
and since $\frac{\mathrm{L}_{\mathrm{r}}}{\mathrm{C}_{\mathrm{r}}}=5.9$,
we get

$$
\begin{aligned}
& \mathrm{C}_{\mathrm{r}}=4 \mu \mathrm{fd} \\
& \mathrm{~L}_{\mathrm{r}}=23.7 \mu \mathrm{H}
\end{aligned}
$$

We can now calculate several other important parameters for the circuit. Figs. II-8 through II-12 show the following normalized variables as functions of $q$ and $\alpha$. These curves are all based on equations derived in [14]. The base current, $I_{B}$, was previously defined and,
$\mathrm{V}_{\mathrm{B}}=$ base voltage $=\mathrm{V}_{\mathrm{s}}$
normalized RMS current $=I_{R N}$ (Fig. II-8)
normalized peak current $=I_{P N}$ (Fig. II-9)
normalized avg. MCT current $=\mathrm{I}_{\mathrm{QN}}$ (Fig. II-10)
normalized avg. diode current $=I_{D N}$ (Fig. II-11)
normalized peak capacitor voltage $=\mathrm{V}_{\mathrm{PN}}$ (Fig. II-12)
For $\mathrm{q}=.9$ and $\alpha=35^{\circ} \mathrm{min}$., we obtain the following values.

$$
\begin{array}{ll}
\mathrm{I}_{\mathrm{RN}}=2.75 & , \\
\mathrm{I}_{\mathrm{RMS}}=117.7 \mathrm{~A} . \mathrm{rms} \\
\mathrm{I}_{\mathrm{pN}}=4.15 & , \\
\mathrm{I}_{\mathrm{peak}}=177.6 \mathrm{~A} . \text { peak } \\
\mathrm{I}_{\mathrm{QN}}=1.125 & , \\
\mathrm{I}_{\text {Qavg }}=48.2 \mathrm{~A} . \mathrm{avg} \\
\mathrm{I}_{\mathrm{DN}}=.0625 & , \\
\mathrm{I}_{\text {Davg }}=2.7 \mathrm{~A} . \mathrm{avg} \\
\mathrm{~V}_{\mathrm{pN}}=4.15 & , \\
\mathrm{~V}_{\text {Cpeak }}=431.6 \mathrm{~V} . \mathrm{pk} .
\end{array}
$$

However, it should be noted that the above data does not represent the maximum values for all variables. This is illustrated by Figs. II-13 through II-18 where the track of $\alpha$ is plotted as $q$ approaches 0.1 (the lowest value plotted) while $I_{\text {avg }}$ is held constant. These values are obtained by plotting the $\alpha$ track for constant $I_{\text {AN }}$ in Fig. II-13 and then using each pair of $q$ and $\alpha$ values to find the corresponding track on Figs. II-14 through II-18. The IRN, IPN and IQN tracks either remain fairly flat or decrease in value, but IDN and VPN both increase to produce the following maximum values at $\mathrm{q}=0.1$,

$$
\begin{array}{ll}
I_{D N \max }=.563 & , \\
\mathrm{I}_{\mathrm{D} \max }=24.1 \mathrm{~A} . \mathrm{avg} . \\
\mathrm{VNmax}=4.7 & \mathrm{~V}_{\mathrm{CP} \max }=488.8 \mathrm{~V} . \mathrm{pk}
\end{array}
$$

To determine the $A C$ ripple currents for $C_{s}$ and $C_{0}$ we use,

$$
I_{\mathrm{rms}}=\left[I_{\mathrm{avg}}^{2}+\sum_{\mathrm{n}=1}^{\infty} I_{n}^{2}\right]^{1 / 2}
$$

where $I_{n}=$ RMS value of the $n$th harmonic
Therefore for $\mathrm{C}_{\mathrm{s}}$,

$$
\begin{aligned}
\mathrm{I}_{\mathrm{cs}} & =\left[\mathrm{I}_{\mathrm{rms}}^{2}-\mathrm{I}_{\mathrm{avg}}^{2}\right]^{1 / 2} \\
& =\left[(117.7)^{2}-(107)^{2}\right]^{1 / 2} \\
& =49 \mathrm{~A} \cdot \mathrm{rms} .
\end{aligned}
$$

The actual transformer used a turns ratio $\mathrm{n}=.412$, so for $\mathrm{C}_{0}$,

$$
\mathrm{I}_{\mathrm{co}}=20.2 \mathrm{~A} . \mathrm{rms} .
$$

These ripple currents are well within the ratings of polypropylene capacitors, but if electrolytics are used to provide lower ripple voltages, the current rating usually determines the capacitor size instead of the filtering requirements. This particular design used an electrolytic and polypropylene combination for $\mathrm{C}_{\mathbf{s}}$ and polypropylene only for $\mathrm{C}_{0}$.

The RMS load current values did not pose a particular problem for $\mathrm{C}_{\mathrm{r}}$, but it was necessary to use higher voltage SCR commutation type capacitors to achieve the necessary voltage ratings.

It was originally thought that the $23.7 \mu \mathrm{H}$ inductor could be implemented simply by increasing the leakage inductance of T 1 , which would eliminate one of the larger components. This was done using two ferrite " C " cores with the primary wound on one leg of the core. The secondary was then divided, with part of the winding on top of the primary and the remaining part on the opposite core leg. The two parts were divided experimentally to achieve a leakage close to $\mathbf{2 3 . 7}$
$\mu \mathrm{H}$. The resulting circuit performed well at smaller loads, but as the load approach $4-5 \mathrm{~kW}$, the conduction losses became excessive and the efficiency was only about $80 \%$. This heating appears to have been caused by eddy currents in the windings induced by the leakage flux and/or proximity effects.
$T 1$ and $L_{r}$ were then separated, and it was determined that it was necessary to limit the number of turns on each of these devices and to limit the number of winding layers. Various gapped core and winding combinations were evaluated for $L_{r}$, but the only one that was considered successful was a design that used only 3 turns on a rather large ferrite core. This inductor is discussed in greater detail in a following section.


Fig II-6. $\gamma$ vs $\alpha$


Fig II-7. IAN vs $\kappa$


Fig II-8. IRN vs $\alpha$


Fig II-9. IPN vs $\alpha$


Fig Il-10. IQN vs $\alpha$


Fig II-11. IDN vs $\alpha$


Fig II-12. VPN vs $\alpha$


Fig II-13. IAN vs $\alpha$ locus as $q$ decreases.


FIG II-14. IRN vs $\alpha$ locus as $q$ decreases.


Fig II-15. IPN vs $\alpha$ locus as $q$ decreases.


Fig ll-16. IQN vs $\alpha$ locus as $q$ decreases.


Fig II-17. IDN vs $\alpha$ locus as $q$ decreases.


Fig II-18. VPN vs $\alpha$ locus as $q$ decreases.

## III. POWER CIRCUIT PERFORMANCE AND EXPERIMENTAL DATA

## A. Power Circuit Description

The power circuit diagram is reproduced in Fig. III-1 for convenience. Each switch in the full-bridge series-resonant converter comprises 3 paralleled MCTs (Harris developmental number MCTA60P60) and 1 anti-parallel diode (Harris developmental number RURU15060). The MCTs in one switching pole of the inverter were later changed to the Harris commercial part number MCTA75P60, considered to be a close equivalent of the developmental part. Because peak currents in steady state operation approached 250 A , with switched currents of nearly 150 A under some conditions, a low-inductance layout was used for the MCTs and the dc-side bypass capacitors. These values proved to be significantly higher than those from the initial design calculations because a lower value of $\alpha$ was eventually used. The mechanical layout of the MCTs and diodes is shown in Fig. III-2.

The purpose of the low-inductance assembly shown in Fig. III-2 is two-fold: One purpose is the control of the layout inductance, the other is the facilitation of removal and replacement of individual TO-218 packages without damage to the other devices. The assembly is composed of a 15.2 -by- 20.3 -by- 0.48 cm (6-by- 8 -by-3/16-inch) aluminum carrier plate, which does not carry current, together with 4 sub-assemblies bolted to it and insulated from it using $0.15-\mathrm{mm}$ ( $6-\mathrm{mil}$ ) Wakefield type-175 Kapton thermally-conducting electrically-insulating material. Each subassembly is also built on $3 / 16$-inch aluminum plate. One carries 6 MCTs and 2 diodes, forming two of the bridge switching legs, and also serving as the negative dc bus. This sub-assembly is marked "-dc" in Fig. II-2. Two other sub-assemblies carry 3 MCTs and 1 diode each, forming two more bridge switching legs and the two ac output terminals of the bridge. These sub-assemblies are marked "ac 1" and "ac 2" in the figure. All MCTs and diodes are mounted directly to their


Fig. III-1 Power circuit of $10-\mathrm{kW} 100-120 \mathrm{~V}$ arcjet power converter.


Fig. III-2 Mechanical layout for the switching devices of the $10-\mathrm{kW}$ arcjet power converter.


Fig. III-3 Core arrangement used for the $10-\mathrm{kW}$ series-resonant inductor.
respective aluminum plates without any insulators. Heat flow from the switching devices is to their sub-assembly base plates, then through the Kapton insulating material to the aluminum carrier plate, and then to the cold plate to which the carrier plate is bolted.

Several comments can be made about this arrangement. One is that its heat flow properties can be very good if the problem of maintaining uniform pressure across the entire area of the sandwich of aluminum plates and insulators is properly solved. Space permitting, the sub-assembly base plates could be enlarged to spread the heat over a greater area before it passes through the insulator, thus lowering the net thermal resistance of this mounting system. Although it was not done on this project, EMI control could also be improved in some situations by adding a second layer of insulation between the aluminum carrier plate and the cold plate to which it is bolted. Connection to the carrier plate would then allow the redirection of the capacitive currents (displacement currents) through the insulating material away from the grounded cold plate.

It can also be noted that the arrangement described does not produce the lowest possible layout inductance because the commutating currents flow essentially side-by-side, although they have broad conducting paths. An alternative arrangement in which the commutating currents flow in facing layers of the structure would lower its inductance still farther, but was considered unnecessary here. The fairly low dc voltage of $100-120 \mathrm{~V}$ with the $600-\mathrm{V}$ ratings of the MCTs made the control of voltage overshoot much less critical. Individual removal and replacement of the switching devices was also more readily possible with this mounting arrangement.

Approximately 300 nH of commutating inductance, shown in Fig. III-1, was added to each switching pole. This controls the di/dt of the commutations involving turn-off of the diode, and turn-on of the incoming MCT. Turn-off of the

MCTs is essentially lossless because of the operation of the series resonant tank below resonance. The snubbers were the simple R-C networks shown in Fig. III1 , and produced negligible snubbing losses relative to the power levels processed.

The resonant inductor ( L 1 in Fig. III-1) was found to be a substantial component, both in terms of size and losses. Three different arrangements with a small number of turns were tried, each using an air-gapped ferrite core assembly. Each was still found to be lossy, with excessive temperature rise in the copper windings. (In each case the core losses were small; the ferrite did not become hot until heat was transferred to it from the windings.) The final version of the inductor core is shown in Fig. III-3. This assembly of 3C8 ferrite C-I cores was wound with 3 turns (in 3 layers) of copper strap, consisting of two paralleled pieces 3.2 cm wide and 0.89 mm thick ( 1.25 by 0.035 in ). Skin depth at 15 kHz is $0.54 \mathrm{~mm}(0.021 \mathrm{in})$, so no benefit from additional copper is expected. The magnetic path is air gapped in two places at approximately $1.9 \mathrm{~mm}(0.075 \mathrm{in})$, for a total gap of 3.8 mm . The inductance is $23 \mu \mathrm{H}$ at 250 A peak, for a stored energy of 0.72 J .

## B. Steady-State Performance

The steady-state performance data for the converter is summarized in Table III-1, which lists the input-output conditions for four steady state operating points. These four test conditions are numbered in the table for later reference. Measured efficiencies ranged from $93.7 \%$ at partial load to $91.6 \%$ at full load of 201 V and 51.4 A output ( 10.3 kW ).

Table III-1
Steady-State Performance Data for the 10-kW Arcjet Converter

|  | Input <br> Voltage <br> $\mathrm{V}_{\mathrm{s}}$ | Input <br> Current <br> $\mathrm{I}_{\mathrm{s}}$ | Input <br> Power <br> $\mathrm{P}_{\mathrm{s}}$ | Output <br> Voltage <br> $\mathrm{V}_{\mathrm{o}}$ | Output <br> Current <br> $\mathrm{I}_{\mathrm{o}}$ | Output <br> Power <br> $\mathrm{P}_{\mathrm{o}}$ | Power <br> Effic. | Power <br> Loss |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 120 V | 58.0 A | 6960 W | 182.6 V | 35.7 A | 6519 W | $93.7 \%$ | 441 W |
| 2 | 120 V | 70.8 A | 8496 W | 153.1 V | 49.5 A | 7579 W | $89.2 \%$ | 918 W |
| 3 | 120 V | 78.0 A | 9360 W | 183.0 V | 47.0 A | 8592 W | $91.8 \%$ | 768 W |
| 4 | 120 V | 94.0 A | 11280 W | 201 V | 51.4 A | 10330 W | $91.6 \%$ | 949 W |

Losses in the switching devices under the four conditions of Table III-1 were also measured. A Tektronix 11401 digital storage oscilloscope was used to sample the switch voltage and current waveforms for switch leg Q4/D4. These waveforms were event-averaged to reduce their noise levels, and then multiplied to form an instantaneous power waveform. This waveform was numerically integrated to determine the energy loss per switching period. The total Q1-Q4 and D1-D4 conduction and switching losses of the converter were then estimated by multiplying this result by 4. The data are summarized in Table III-2: The entries in this table correspond by row to those in Table III-1. Figs. III-4 through III-7 are also referenced in this table, and illustrate the switch loss-measurement data.

Table III-2
Measured Switching Device Losses, Based on the Q4/D4 Leg

|  | Switching <br> Period | Switching <br> Frequency | Energy | Power <br> (1 Leg) | Power <br> (4 Legs) | Total <br> Converter <br> Loss | See <br> Figure |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 | $72.0 \mu \mathrm{~S}$ | 13.89 kHz | 6.26 mJ | 87.0 W | 348 W | 441 W | III-4 |
| 3 | $70.1 \mu \mathrm{~S}$ | 14.27 kHz | 7.99 mJ | 114 W | 456 W | 918 W | III-5 |
| 3 | $69.9 \mu \mathrm{~S}$ | 14.37 kHz | 7.45 mJ | 107 W | 428 W | 768 W | III-6 |
| 4 | $67.7 \mu \mathrm{~S}$ | 14.78 kHz | 6.88 mJ | 102 W | 407 W | 949 W | III-7 |

Waveforms showing the details of the switching in the Q4/D4 leg are in Figs. III-8 and III-9. Fig. III-8 shows the turn-on transition. The upper window in this figure shows the MCT anode-cathode voltage (vaK, $100 \mathrm{~V} / \mathrm{div}$ ), anode current ( $i_{A}, 50 \mathrm{~A} / \mathrm{div}$ ) and gate-anode voltage ( $\mathrm{v}_{\mathrm{GA}}, 20 \mathrm{~V} / \mathrm{div}$ ). The operating conditions were an input voltage of 120 V , and output voltage of 137 V at 49.7 A , so the switched currents were close to those of full-power operation. (The turn-on current actually exceeds that of full-power by about 25 A.) The lower window of Fig. III-8 shows the turn-on waveforms for this switch at $500 \mathrm{nS} / \mathrm{div}$. It can be noted that turn-on of the MCT occurs with a negative gate-anode voltage, thus requiring gate drive polarity inverted from the usual arrangement. The manufacturer recommends a positive gate voltage to ensure turn-off of the MCTs. This raises questions of reliability in that the system must be designed to supply turn-off bias to the MCTs prior to the application of anode voltage. The turn-on transition in Fig. III-8 lasts about 500 nS and is lossy. A shelf of about 160 nS may be observed on the falling edge of the anode voltage waveform. Increased commutating inductance could be used to reduce this switching loss, at the expense of increased snubbing loss, but this was not done. Fig. III-7 shows that


Fig. III-4 Steady-state operation, case 1. Switch loss measurement (refer to Table III-2).
Upper window: Switch current ( 50 A /div) Switch voltage ( $100 \mathrm{~V} /$ div)
Lower window: Switch instantaneous power ( $200 \mathrm{~W} / \mathrm{div}$ )


Fig. III-5 Steady-state operation, case 2. Switch loss measurement (refer to Table III-2). Upper window: Switch current ( $50 \mathrm{~A} /$ div) Switch voltage ( $100 \mathrm{~V} / \mathrm{div}$ )
Lower window: Switch instantaneous power ( $500 \mathrm{~W} /$ div)


Fig. III-6 Steady-state operation, case 3. Switch loss measurement (refer to Table III-2).
Upper window: Switch current ( 50 A /div) Switch voltage ( $100 \mathrm{~V} / \mathrm{div}$ )
Lower window: Switch instantaneous power ( $500 \mathrm{~W} /$ div)


Fig. III-7 Steady-state operation, case 4. Switch loss measurement (refer to Table III-2). Upper window: Switch current ( $50 \mathrm{~A} /$ div) Switch voltage ( $100 \mathrm{~V} / \mathrm{div}$ )
Lower window: Switch instantaneous power ( $500 \mathrm{~W} /$ div)


Fig. III-8 Switch tum-on loss, operation with 137 V/50 A output ( 6850 W ).
Upper window: Switch voltage ( $100 \mathrm{~V} / \mathrm{div}$ )
Switch current ( 50 A /div)
Gate-anode voltage ( $20 \mathrm{~V} / \mathrm{div}$ )
at $10 \mu \mathrm{~S} / \mathrm{div}$. Baselines indicated.
Lower window: Same as upper window, except timebase is $500 \mathrm{nS} /$ div.


Fig. III-9 Switch turnoff loss, operation with 137 V/50 A output ( 6850 W).
Upper window: Switch voltage ( $100 \mathrm{~V} /$ div)
Switch current ( 50 A div)
Gate-anode voltage ( $20 \mathrm{~V} / \mathrm{div}$ )
at $10 \mu \mathrm{~S} / \mathrm{div}$. Baselines indicated.
Lower window: Same as upper window, except timebase is $500 \mathrm{nS} /$ div.
although there is an impulse of loss at turn-on exceeding 2.5 kW , its approximate energy content is only 0.75 mJ , thus its contribution relative to the conduction loss is not great.

The turn-off transition is detailed in Fig. III-9 for the same operating conditions as Fig. III-8. The upper window of this figure also shows the MCT anode-cathode voltage ( $\mathrm{v}_{\mathrm{AK}}, 100 \mathrm{~V} / \mathrm{div}$ ), anode current ( $\mathrm{i}_{\mathrm{A}}, 50 \mathrm{~A} / \mathrm{div}$ ) and gateanode voltage (vGA, $20 \mathrm{~V} / \mathrm{div}$ ). The lower window shows the turn-off waveforms at $500 \mathrm{nS} / \mathrm{div}$. This switching transition is lossless for the MCT due to the belowresonance operation of the series-resonant converter. It may be seen that the gate-anode voltage has made its transition to the off state well in advance of the current commutation, and also that turn-off losses for the anti-parallel diode are low. The diodes, Harris developmental device RURU15060 with a rated recovery time of 90 nS , did quite well in this regard compared to other $600 \mathrm{~V} / 150 \mathrm{~A}$ diodes. Below-resonance operation was selected for use with the MCTs because of the lengthy current fall time ( $1.4 \mu \mathrm{~S}$ ) specified for this device.

## C. Load Transient Tests

The transient response of the dc/dc converter was explored using stepped load changes. The converter was loaded with a resistive load bank, connected directly across the $20 \mu \mathrm{~F}$ output capacitor (the $200 \mu \mathrm{H}$ inductor in the ignition circuit was not included). Fig III-10 shows the results of a step increase in the load resistance. The output current (center of Fig. III-10, $50 \mathrm{~A} /$ div) is virtually unaffected by a doubling of the load resistance. The output voltage (top of Fig. III$10,100 \mathrm{~V} /$ div) nearly doubles. The transient in the resonant tank current can be seen at the bottom of the figure ( $200 \mathrm{~A} /$ div) to be very small.

Fig. III-11 shows the transient due to a step decrease in the load resistance. Here the load voltage (top of Fig. III-11, $100 \mathrm{~V} / \mathrm{div}$ ) halves following the transient. The load current (center, 50 A div) shows a momentary overshoot of about 38 A
and then returns to its set point of 50 A . There is a transient overcurrent in the resonant tank current (bottom, 200 A /div), but it is of quite manageable size. Flux capacity for this transient should be allowed in the design of the resonant inductor, but it is probable that the flux density needed for a low-loss design will ensure this.

The effects of the $200 \mu \mathrm{H}$ output inductor in the ignition circuit were next explored. Fig. III-12 shows the transient produced by a load resistance step increase with the output inductor in place. This waveform is virtually the same as that of Fig. III-10, showing that the $200 \mu \mathrm{H}$ inductor produces a insignificant time constant, relative to the others associated with the resonant converter. Fig. III-13 shows the transient produced by a load resistance step decrease. Again, the result is virtually the same as that of Fig. III-11. This transient is also shown on a different time scale ( $5 \mathrm{mS} / \mathrm{div}$ ) in Fig. III-14. Both the $20-\mu$ F capacitor voltage (top, $100 \mathrm{~V} / \mathrm{div}$ ) and the load voltage (next down, $100 \mathrm{~V} / \mathrm{div}$ ) are shown in Fig. III14. It can be seen that the inductor contributes negligible dyramics on this time scale, although it does serve to filter the load current ripple.


Fig. III-10 Load-resistance step-increase-transient with initial operation at $130 \mathrm{~V}, 50 \mathrm{~A}$ Ignition circuit not in place.

Upper: Load voltage ( $100 \mathrm{~V} /$ div).
Middle: Load current ( $50 \mathrm{~A} /$ div).
Lower: Resonant tank current ( 200 A /div).
Timebase: $10 \mathrm{mS} / \mathrm{div}$.


Fig. III-11 Load-resistance step-decrease-transient with initial operation at $220 \mathrm{~V}, 50 \mathrm{~A}$. Ignition circuit not in place.

Upper: Load voltage ( $100 \mathrm{~V} /$ div).
Middle: Load current ( 50 A div).
Lower: Resonant tank current ( 200 A /div).
Timebase: $10 \mathrm{mS} /$ div.


Fig. III-12 Load-resistance step-increase-transient with initial operation at $130 \mathrm{~V}, 50 \mathrm{~A}$. Ignition circuit in place.

Upper: Load voltage ( $100 \mathrm{~V} /$ div).
Middle: Load current ( $50 \mathrm{~A} /$ div).
Lower: Resonant tank current ( $200 \mathrm{~A} / \mathrm{div}$ ).
Timebase: $10 \mathrm{mS} / \mathrm{div}$.


Fig. III-13 Load-resistance step-decrease-transient with initial operation at $220 \mathrm{~V}, 50 \mathrm{~A}$. Ignition circuit in place.

Upper: Load voltage ( $100 \mathrm{~V} /$ div).
Middle: Load current ( $50 \mathrm{~A} /$ div).
Lower: Resonant tank current ( 200 A div).
Timebase: $10 \mathrm{mS} /$ div.


Fig. III-14 Load-resistance step-decrease-transient with initial operation at $220 \mathrm{~V}, 50 \mathrm{~A}$
(Same event as Fig. II-13, ignition circuit in place.)
Upper: Load voltage ( $100 \mathrm{~V} / \mathrm{div}$ ).
Middle: Load current ( $50 \mathrm{~A} /$ div).
Lower: Resonant tank current ( $200 \mathrm{~A} / \mathrm{div}$ ). Timebase: $5 \mathrm{mS} / \mathrm{div}$.

## IV. IGNITION CIRCUTT

## A. Power Circuit Design

The power circuit of the arcjet ignitor is shown in Fig. IV-1, and the gate driver is in Fig. IV-2. The design of Fig. IV-1 was motivated by [1]. Closure of the momentary contact switch in the control circuit sets the monostable to gate the IGBT Q1 for a timed pulse of about $24 \mu \mathrm{~S}$. Energy is stored in the primary of the coupled inductor $L_{P}$, with the IGBT interrupting about 400-500 A of primary current. At turn-off of the IGBT, resonance of $L_{P} / L_{S}$ and the IGBT snubber capacitor C1 $(1.1 \mu \mathrm{~F})$ together with CSEC $(0.005 \mu \mathrm{~F})$ begins, allowing the IGBT voltage to rise to about 550 V , assuming an open-circuit on the secondary of the ignition circuit. The secondary-side voltage will be ten times this, due to the turns ratio of the coupled inductor. After the IGBT voltage crests, the snubber diode D1 blocks, and resonance continues with $L_{P} / L_{S}$ and CSEC only. When the IGBT voltage reaches zero, its anti-parallel diode conducts, allowing recovery of the energy stored in $L_{p}$. If the arcjet ignites on the rising edge of the inductor voltage waveform, the energy stored in the inductor at that time establishes the initial current in the arcjet. In either case, some energy is transferred to C 1 ; the $2-\mathrm{K} \Omega$ resistor discharges this capacitor. There is an additional monostable in the control which prevents ignition pulses from exceeding a rate of about one per second, thus ensuring reset of the circuit on successive attempts. Winding data for the coupled inductor $L_{P} / L_{S}$ are given in Table IV-1.

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Fig. IV-1 Arcjet ignition power circuit.


## Table IV-1

Winding Data for the Coupled Inductor $\mathrm{L}_{\mathrm{p}} / \mathrm{L}_{\mathrm{S}}$

| Core: | Powdered iron |  | Micrometals T520-40D Relative permeability $=60$ |
| :--- | :---: | :---: | :---: |
| Primary | Wound over | $\mathrm{L}_{\mathrm{P}}=$ | 2 turns of \#14, repeated 20 times, spaced evenly |
| Winding: | secondary | 2 uH | around the circumference of the inductor, and all |
|  |  |  |  |
| paralleled. |  |  |  |

The inductor-charging portion of the ignition pulse is initiated by turn-on of Q1. Energy is transferred from the $100-u F$ capacitor $C_{B Y P}$ to $L_{P}$ in a resonant fashion. The size of $C_{B Y P}$ was limited to bound the collector current that would flow if the IGBT failed to turn off. The period of the inductor-charging resonance is given by:

$$
\begin{equation*}
\mathrm{T}_{\mathrm{s} 1}=2 \pi \sqrt{\mathrm{~L}_{\mathrm{P}} \mathrm{C}_{\mathrm{BYP}}}=89 \mu \mathrm{~S} \tag{IV-1}
\end{equation*}
$$

The amplitude of the inductor current, primary referred, which could occur if the IGBT failed to turn off is then given by (assuming a $120-\mathrm{V}$ supply):

$$
\begin{equation*}
\mathrm{I}_{\mathrm{PEAK}}=\frac{\mathrm{V}_{\mathrm{SUPPLY}}}{\sqrt{\frac{\mathrm{~L}_{\mathrm{P}}}{\mathrm{C}_{\mathrm{BYP}}}}}=848 \mathrm{~A} \tag{IV-2}
\end{equation*}
$$

Turn off of Q1 is followed by a rapidly rising inductor voltage as resonance involving the IGBT snubber and secondary-side capacitors occurs. The period of this resonance is given by:

$$
\begin{gather*}
\mathrm{T}_{\mathrm{S} 2}=2 \pi \sqrt{\mathrm{~L}_{\mathrm{P}} C_{E Q}}=11.2 \mu \mathrm{~S} \\
\text { where } \mathrm{C}_{\mathrm{EQ}}=\mathrm{C} 1+\mathrm{C}_{\mathrm{SEC}}\left(\frac{\mathrm{~N}_{\mathrm{SEC}}}{\mathrm{~N}_{\mathrm{PRI}}}\right)^{2}=1.6 \mu \mathrm{~F} \tag{IV-3}
\end{gather*}
$$

The peak value of the voltage, referred to the secondary of the coupled inductor is:

$$
\begin{equation*}
\mathrm{V}_{\mathrm{PEAK}}=\mathrm{I}_{\mathrm{INTTIAL}} \sqrt{\frac{\mathrm{~L}_{\mathrm{P}}}{\mathrm{C}_{\mathrm{EQ}}}} \cdot \frac{\mathrm{~N}_{\mathrm{SEC}}}{\mathrm{~N}_{\mathrm{PRI}}}=4.5 \mathrm{kV} \tag{IV-4}
\end{equation*}
$$

assuming an initial current of 400 A . If the arcjet fails to ignite, this open-circuit voltage will appear on the secondary side of the ignition circuit. The arcjet is expected to break over on the rising portion of this sinusoidal voltage, however. If breakover occurs at the voltage level $\mathrm{V}_{\mathrm{BO}}$, then the current transferred to the arc will be:

$$
\begin{equation*}
\mathrm{I}_{\mathrm{ARC}}=\frac{\mathrm{N}_{\mathrm{PRI}}}{\mathrm{~N}_{\mathrm{SEC}}} \mathrm{I}_{\mathrm{INTIAL}} \cos \theta \quad \text { where } \quad \sin \theta=\frac{\mathrm{V}_{\mathrm{BO}} \frac{\mathrm{~N}_{\mathrm{PRI}}}{\mathrm{~N}_{\mathrm{SEC}}}}{\mathrm{I}_{\mathrm{INTTAL}} \sqrt{\frac{\mathrm{~L}_{\mathrm{P}}}{\mathrm{C}_{\mathrm{EQ}}}}} \tag{IV-5}
\end{equation*}
$$

If the arcjet break over is assumed to occur at 1.6 kV , then $\theta=20.5 \mathrm{deg}$ and $\mathrm{I}_{\mathrm{ARC}}$ $=37.5 \mathrm{~A}$. Failure of the arcjet to ignite implies that at the peak of the voltage swing, the IGBT snubber capacitor will be isolated from the resonance by diode D1. The resonance will continue with only the secondary-side capacitor participating. When the IGBT collector voltage reaches zero, its internal antiparallel diode will conduct, permitting recovery of the energy stored in the coupled inductor back to CBYP.

It is expected that arcjet ignition will occur on the rising part of the inductor voltage waveform, with a substantial portion of the initial stored energy available to establish a secondary current of about 40 A upon arc-over of the arcjet. The closer this secondary current to the current set point in the arcjet power converter, the lesser the magnitude of the transient following ignition, and the less likely a blowout of the arc. Although it was not done here, it may be useful to use a current sense to cause IGBT turn-off at a specific initial current as done in [18]. This current setting could be coordinated with the current set point used in
the power converter; this added feature would be particularly useful if operation with widely varying supply voltage or current set point was expected.

## B. Experimental Results

Figs. IV-3 and IV-4 show the experimental waveforms obtained with an open-circuited secondary. (Because these waveforms were obtained as singleevent records, sufficient sampling rate was available for only two traces per figure, thus two figures are necessary to record the desired data.) Fig. IV-3 shows the coupled inductor primary current ( $50 \mathrm{~A} / \mathrm{div}$ ) and IGBT collector voltage ( $200 \mathrm{~V} / \mathrm{div}$ ) for operation with a $100-\mathrm{V}$ supply and no arcjet. This figure shows a peak IGBT collector current of 400 A , and a peak collector voltage of 550 V . A parasitic resonance involving the leakage inductance associated with the coupled inductor $L_{P} / L_{S}$ is evident in the current waveforms. The negative primary current after return of the collector voltage to zero corresponds to the energyrecovery phase of the event. It can also be noted in this figure that the $100-\mathrm{uF}$ capacitor, initially charged to 100 V , is left with about 40 V at the end of the pulse due to the losses in the circuit. Fig. IV-4 shows the inductor primary current (50 $\mathrm{A} / \mathrm{div}$ ) and secondary voltage ( $2 \mathrm{kV} /$ div) for the same open-circuit event. The opencircuit secondary voltage is 6 kV .

Figs. IV-5 through IV-8 were obtained using a spark gap to simulate the arcjet, with the gap adjusted to obtain various breakdown voltages. For each of these figures, the secondary voltage $v_{s}$ is shown at $2 \mathrm{kV} / \mathrm{div}$, and the secondary current is shown at 10 A /div. Table IV-2 records the breakdown voltage $\mathrm{V}_{\mathrm{BO}}$ observed, along with the current initially established in the arc.

## Table IV-2

Breakdown Voltage Versus Initial Arc Current for the Ignition Circuit

| Reference | Breakdown | Initial Arc | Theoretical |
| :---: | :---: | :---: | :---: |
| Figure. | Voltage | Current | Breakdown |
|  | $\mathrm{V}_{\text {BO }}$ |  | Voltage* |
| Fig. TV-5 | 1.0 kV | 45 A | 1.9 kV |
| Fig. IV-6 | 2.5 kV | 40 A | 2.7 kV |
| Fig. IV-7 | 3.5 kV | 30 A | 3.6 kV |
| Fig. IV-8 | 3.8 kV | 25 A | 3.9 kV |

*Based on 500 A initial current at IGBT turn off, and the observed initial arc current
It should be noted that the theoretical calculations are based.on a lossless resonance. However, the powdered-iron core and the IGBT forward voltage drop both contribute significant loss to the resonance; therefore there is some discrepancy between the expected and observed voltages and currents. The ignition circuit appears suitable to establish initial currents in the range of 30-40 A in the arcjet, and this can readily be increased to 50 A if necessary by increasing either the IGBT on-time or the size of $\mathrm{C}_{\mathrm{BYP}}$.


Fig. IV-3 Arcjet ignition circuit waveforms on open-circuit-load.
Upper: IGBT collector voltage $\mathrm{v}_{\mathrm{O} 1}(200 \mathrm{~V} /$ div)
Lower: Primary current $i_{p}$ ( 50 A /div)
Timebase: $10 \mu \mathrm{~S} / \mathrm{div}$.


Fig. IV-4 Arcjet ignition circuit waveforms on open-circuit-load.
Upper: Secondary voltage $\mathrm{v}_{\mathrm{s}}(2 \mathrm{kV} / \mathrm{div})$
Lower: Primary current $i_{p}$ ( $50 \mathrm{~A} /$ div)
Timebase: $10 \mu \mathrm{~S} / \mathrm{div}$


Fig. IV-5 Arcjet ignition circuit waveforms, spark gap adjusted for breakover at 1 kV . Upper: Secondary voltage $\mathrm{v}_{\mathrm{s}}(2 \mathrm{kV} /$ div) Lower: Secondary current $i_{s}$ ( $10 \mathrm{~A} /$ div)
Timebase: $20 \mu \mathrm{~S} /$ div.


Fig. IV-6 Arcjet ignition circuit waveforms, spark gap adjusted for breakover at 2.5 kV . Upper: Secondary voltage $\mathrm{v}_{\mathrm{s}}(2 \mathrm{kV} /$ div $)$ Lower: Secondary current $i_{s}$ ( $10 \mathrm{~A} /$ div)
Timebase: $20 \mu \mathrm{~S} / \mathrm{div}$.


Fig. IV-7 Arcjet ignition circuit waveforms, spark gap adjusted for breakover at 3.5 kV .
Upper: Secondary voltage $\mathrm{v}_{\mathrm{s}}(2 \mathrm{kV} /$ div)
Lower: Secondary current $i_{s}$ ( $10 \mathrm{~A} /$ div)
Timebase: $20 \mu \mathrm{~S} /$ div.


Fig. IV-8 Arcjet ignition circuit waveforms, spark gap adjusted for breakover at 3.8 kV Upper: Secondary voltage $\mathrm{v}_{\mathrm{s}}(2 \mathrm{kV} / \mathrm{div})$ Lower: Secondary current $i_{s}$ ( $10 \mathrm{~A} /$ div) Timebase: $20 \mu \mathrm{~S} /$ div.

## V. SUMMARY

The design of a $10-\mathrm{kW}$ dc/dc converter with a $120-\mathrm{V}$ source has been explored. At this voltage level, the switching device of choice is the MOSFET because of its low switching losses and low conduction losses if it is assumed that large numbers of dice can be satisfactorily paralleled. However, at the $10-\mathrm{kW}$ level the number of dice becomes so large that there may be a reliability problem. The use of the MCT as an alternative was therefore chosen for this project. Its high current density at reasonably low conduction losses allowed the use of only three paralleled devices in each switch of the $10-\mathrm{kW}$ full-bridge series-resonant converter.

The series-resonant converter, operating below resonance, was chosen to reduce the turn-off switching loss expected with the MCT. The use of a fast developmental diode (Harris RURU15060) controlled the MCT turn-on losses (and diode turn-off losses). This strategy produced good results with regard to the primary-side switches: The combined switching and conduction losses of the four switching legs was less than 460 W under all conditions. The main difficulty with the series-resonant approach was found to be the resonant inductor, which exhibited losses comparable to those of the switches, reducing the efficiency of the converter. It is believed that further work could significantly improve the design of the inductor, but it remains a large and heavy component.

The overall power efficiency achieved was $91.6 \%$ at 10.3 kW output ( 51.4 A at 201 V ). Under this condition, the primary-side switches contributed 407 W out of a total loss of 949 W ; the resonant inductor contributed the bulk of the remainder. The power efficiency was generally found to improve at reduced load currents; for example, it was $93.7 \%$ at 6.5 kW output ( 35.7 A at 182.6 V ). Under this condition, the primary-side switches contributed the bulk of the loss ( 348 W out of 441 W ).


The arcjet ignition circuit was readily able to produce open-circuit voltages of 6 kV , and initial currents in the arc in the range 30-40 A. The dependency of the initial arc current on the arc breakdown voltage was predicted theoretically, and confirmed experimentally. Because the arcjet breakdown voltage is not precisely repeatable, there will be some variation in the initial arc current established by the ignition circuit, and thus a current transient following ignition as the arcjet converter's current regulator takes control of the arc current. This is generally a large-signal transient, and is best explored experimentally using either an arcjet, or a load programmed to simulate an arcjet.

Transient testing of the converter using load resistance step changes was done with good results. The converter appears to be readily adapted to stabilizing the arcjet current under start-up and steady-state conditions.

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| 13. ABSTRACT (Maximum 200 words) Many arcjet DC power supp investigates an alternative s MCTs. The reasons for this Decrease the surface mount because the relatively slow satisfactorily with an efficie additional optimization of the | s) <br> plies use PWM full bridge conv upply using a variable frequency approach are to: 1) Increase re ting area of the switching array switching speed of the MCT pr ency of over $91 \%$. Test results is he series resonant inductor. | rters with large arrays of para series resonant converter with ability by reducing the numb The variable frequency series cludes the use of PWM. The dicate this efficiency could be | FETs. This report mall arrays of parallel of switching devices; 2) esonant approach is used kW converter operated creased further by |
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