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## FINAL REPORT

**DEFORMED ELLIPSOIDAL DIFFRACTION GRATING BLANK** 

Contract No. NAS5-31790

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## FINAL REPORT

## DEFORMED ELLIPSOIDAL DIFFRACTION GRATING BLANK

Contract No. NAS5-31790

Submitted by:

SPACE OPTICS RESEARCH LABS AN INTERGRAPH DIVISION 7 STUART ROAD CHELMSFORD, MA 01824

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## FINAL REPORT

## DEFORMED ELLIPSOIDAL DIFFRACTION GRATING BLANK

Contract No. NAS5-31790

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## 1.0 Introduction

The <u>D</u>eformed <u>Ellipsoidal Grating Blank</u> (DEGB) is the primary component in an ultraviolet spectrometer. Since one of the major concerns for these instruments is throughput, significant efforts are made to reduce the number of components and subsequently reflections. Each reflection results in losses through absorbtion, and scattering. It is these two sources of photon loss that dictated the requirements for the DEGB.

The first goal is to shape the DEGB in such a way that the energy at the entrance slit is focused as well as possible on the exit slit. This as a geometrical problem which must acount for the diffraction angle of the radiation. The result was the shape defined in the specification. This shape is a non-rotationally symetric general asphere. This unusual surface reguires an unusual test and unusual test equipment to verify the shape

The second goal is to produce a surface smooth enough to minimize the photon loss due to scattering. The production techniques to produce thiese surfaces are often held in the craftsmanship of the optician. The measurement is accomplished with modern scanning interferometric methods.

## 2.0 Applicable Documents

## 3.0 Program Summary

The program was accomplished in three phases. The first phase was the fabrication planning. The second phase was the actual fabrication and initial testing. The last phase was the final testing of the completed DEGB.

3.1 Phase 1 -- Fabrication planning

The project actually began at the proposal stage. At that time, SORL calculated the asphericity of the DEGB and used the results to verify the values presented in the RFP. These calculations and calculation methods are shown in Appendix A. This led to many of the conclusions regarding the test methods and equipment requirements.

First, we saw that the underlying ellipsoid for the DEGB is an optic which can be tested in a very conventional null test. The ellipsoid is generated by taking the ellipse and spinning it around its long axis. The section of the ellipsoid we are fabricating is taken from around the minor axis. A double pass test can be used by placing an interferometer at one focus and a retroreflecting sphere centered on the other focus. We also saw that the difference between the minor axis ellipse and a circle which best fits that ellipse is only  $0.11\mu$  ( $\lambda/5.75 @ \lambda=.633\mu$ ). This is a good approximation for estimating the wavefront over the entire piece if the part is generated as a "football torrus". This is where the longer radius is used as the generating surface, and the shorter radius is used to locate the generating axis. This method of generating a torrus closely approximates the method of making the ellipsoid of revoluition which can be tested with the previously mentioned null test. The same radii can also be used to generate a "bagel torrus", and as might be expected in this case, the shorter radius is used as the generating surface, and the longer radius is used to locate the generating axis.

The deformation factor is an addative expression combined with the standard equation for an ellipsoid. The deformation along the edge of the clear aperture as defined by the specification showed a maximum deveation of about 2.41 $\mu$  or 3.81 $\lambda$ . This would show up as 15.25 fringes in an interferogram. A simulated interferogram is presented in figure 1.

SORL prepared a fabrication and test plan. This is included in appendix B. The plan included both the fabrication steps and the testing steps which were intended to be used for the completion of the project. The plan includes drawings of the test configuration and the DEGB component. This plan was delivered 15 November 1991.

## 3.2 Phase 2 -- Fabrication and initial testing

The actual fabrication was carried out along the guidelines put forth in the fabrication plan. The primary goal of this phase was to create the optical figure. The geometry of this figure was worked out during the proposal and phase 1. This led to the conclusion that ordinary interferometry, even with phase shifted data would not be sufficient to meet the accuracy requirements of the surface. Therefore the holographic approach to the testing was adopted. The principles of operation of this instrument are described in the first paper in Appendix D.

## 3.2.1. The Holographic Interferometer

Negotiations for the procurement of the holographic interferometer as proposed began shortly after receipt of the contract. The purchase order to APA optics for the holographic interferometer and all the necessary holograms and accessories was dated 10 Oct 1991. The promised delivery date was 17 Jan 1992 Almost immediately the delivery of the interferometer began to slip. By November, the delivery was targeted for February. In order to ameliorate this slip, SORL sent our optician, Bob Scannell, who was to work on the figuring to have a hand on session with a demonstration unit which was at a trade show in California. The visit was successful with regard to the training. After this trip which occurred in January, we learned that the delivery of our interferometer would be 16 March 1994. In February Alan DeCew traveled to Florida to again observe the APA Interferometer for <u>Aspheric Testing (IAT) at a nother trade show. He was also able</u> to talk to APA's other customer for an IAT, Martin Marieta. By March the delivery hasd slipped to 16 April 1992.

On 22 May 1992 the IAT was installed at SORL. The hologram for the DEGB was not completed, and delivery was scheduled for mid-June. By June, the delivery of the hologram was slipped to July. Then other technical issues arose such as the calibration. The F/4 objective was not fully characterized and would have to be replaced. A new schedule for completing these tasks pushed the IAT qualification out to 24 Sep 1994. The new objective and holograms were actually ready for testing at SORL the first week in November. Unfortunately, the holograms were wrong. A correct hologram did not arive until 23 February, 1993.

Other problems with the IAT also surfaced. The fringes were almost impossible to see. After lengthy negotiations with APA, the instrument was returned to APA. It was to be replaced with a new unit. Delivery was promised for 1 June 1993. The unit was actually set up and working 26 August 1993, more than 19 months after promised delivery. It took APA about 7 times as long as they originally projected to produce a working IAT for SORL.

The Calibration of the unit was performed at APA. The results are presented in Appendix D along with a paper by Steve Arnold describing the calibration procedure.

## 3.2.2 The aspheric figuring

The Zerodur<sup>®</sup> blanks for the DEGB were ordered from Schott in November 1991. The delivery was scheduled for January 1992, and arrived on schedule. By February, the spherical shape had been generated into the blanks, concave on one, convex on the other. The blanks were then delivered into the hands of the optician, Bob Scannell for figuring. The first step was to generate a torroidal shape into the blanks. A number of preliminary forays into developing a process were attempted from March through June. The conclusion was that it was no advantage working on the positive and negative blanks together. Using a combination of full size tools and sub-aperture tools seemed necessary to make predictable progress.

The progress of the grinding was monitored with a programable 3-axis coordinate measurment machine. The sag of the part was measured in three positions. The part was rotated 90° and measured again. These results were studied by the optician to determine the next step of the grinding. The results also averaged and plotted to give an indication of the general direction of the progress. The charts are presented in Appendix C.

This was the state of the optic until the holgraphic interferometer was finally ready for use. At that point progress became rapaid and steady. The goal for the optician was to make the fringes straight. The progress over time can be seen in the interferograms shown in Appendic C.

One interesting development of the testing metrology which starts with the null configuration for the underlying ellipse is that the image of the ellipse as seen by the camera in the LUPI is foreshortened. The fringes which normally would be evenly spaced, appear to radiate from the vanishing point of the circumscribed rectangle.

The verification of this is to look at the null fringe case. The interferogram is difficult to read even in the original copy since it is a null. The visual inspection clearly showed that the the wavefront was adequate to show a complete null. This confirmed the foreshortening phenomenon. Reading the interferograms in the normal manner would indicate huge errors on the order of 1 wave. The reading also becomes worse the more tilt fringes are introduced into the interferogram.

## 3.2.3 The testing metrology

Obtaining the proper shape required the proper instrument, a trained optician, and finally, a careful set up of the equipment. The positions on the table and the procedure for placing the components is given in appendix E.

## 3.3 Phase 3 -- Final testing

The surface figure was established during Phase 2. In Phase 3 the object was to verify all the other requirements of the DEGB. This effort began immediately after establishing that we had met the surface figure requirement.

All the remaining requirements were relatively easy to establish with the exception of the midfrequency. To measure this we contracted with Bauer Associates to use their Model 100 scanner. This was a particularly good choice since Goddard SFC has a similar machine. We could be quite sure that measurements taken at Bauer would correspond closely with measurements taken at GSFC. A number of delays in the program occurred due to the lack of availability of this machine, including the fact that at one point the only available working sources had to go to GSFC to upgrade their Model 100 instrument. The final measurements were made 12 November 1994. This completed the tasks on the contract.

## 4.0 Specification Compliance

All the specification requirements are called out in Attachment 3 of the Statement of Work (Appendix H). These requirements were measured and the results are shown in the QA documentation which was provided with the part. A copy of the documentation is included as Appendix I.

In determining the mechanical requirements standard mechanical metrology tools were used to determine compliance. The blank diameter was measured with a vernier. The center thickness was measured with a drop gauge. The conjugate distances for the ellipsoid were measured with radius rods. The QA Documentation showing the results of these measurements is included in Appendix I.

## 4.1 Surface Figure

The actual shape produced depended on the correct set up of the optics and interferometer according to the set up shown in Appendix E. This was verified by independent verification of the set up by a second test technician. The positions of the conjugate points relative to the optic are within the nominal settings.

4.2 Figure error

The figure was toleranced in three spatial wavelength areas. On the overall figure level the requirement is for 1/8 wave. The midfrequency covers spatial wavelengths from .5mm to 10mm. Microroughness covers spatial wavelengths less than .5mm.

## 4.2.1 Overall Peak to Valley

As pointed out in a previous section, the fringes did not appear as evenly spaced straight fringes. As a result a manual reduction of the fringes was done. This construction is presented in the QA documentation. It was concluded that the surface figure met the 1/8 wave requirement.

## 4.2.2 Midfrequency roughness

Probably the second most important parameter after the surface figure is the midfrequency ripple. The specification for this is 2nm rms. The measurements were made on a Bauer Model 100 noncontact profilometer, and the results are shown in Appendix G. The data shown was taken from a scan from  $\{-25,50\}$  to  $\{+25,50\}$ in the  $\{Y,Z\}$  plane of the ellipsoid as described in the Statement of Work (Appendix H). It is appropriate to remove polynomial 0 through 6 since these polynomials represent spatial wavelengths greater than 10mm. A plot of exactly what each polynomial takes out of the data is included as part of Appendix G.

4.2.3 Microroughness

The microroughness was not directly measured. The results of the Model 100 instrument were below the microrougness requirement. The shorter spatial frequencies typically show smaller rms roughness values. This result is also presented in the QA Documentation.

## 5.0 Recommendations

5.1 Fabrication times and techniques

While considerable more time was spent on the first unit, the estimate of labor hours needed to complete a similar DEGB is 260 hours. There is also associated engineering of 50 hours needed to specify the testing and hologram from APA. Assuming there are quality assurance provisions commensurate with a space flight hardware, and applying industry rates, the selling price for other DEGBs will be about \$50,000.

## 5.2 Figure error metrology

The holographic approach to surface metrology proved quite adequate after the developmental nature of the interferometer was overcome. There are other ways of implementing holographic testing and one is presented in Appendix F.

# DEGB SURFACE FIGURE CALCULATIONS

- 1) Geometric properties of underlying ellipse
- 2) Comparison of ellipse to circle
- 3) Deformation from ellipse

1) Geometric properties of underlying ellipse

From the parameters supplied in the contract Statement of Work other descriptors of the ellipse were generated. In particular, the conic constant "K" which is used in the sag equation was calculated. The vertex curvature is "c" and the radial distance form the optical axis is "Y".

Sag = 
$$\frac{c\gamma^2}{1+\sqrt{1-(1+K)c^2 \gamma^2}}$$

The vertex radii were also calculated for both the minor axis and the major axis. The calculations match the SOW perfectly.

VARIABLE SHEET ======== Comment	****** ELLIPSE ****** semi-major axis semi-major axis semi-minor axis Sag of point on ellipsoid radial coordinate of point on ellipsoi radial coordinate of point on ellipsoi conic constant vertex curvature Vertex radius at major axis Vertex radius at minor axis Focus to Vertex distance Focus to Focus distance	
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Z = C \* X^2 /(1 + Sqrt(1 - (K + 1) \* C^2\*X^2))
K = ( b^2 - a^2) / a^2
C = 1/((K+1)\*a)
Rmajor=1/C
EFF = 2\* SQRT(a^2-b^2)
EVF = a-EFF/2
Rminor = a^2/b
Rminor = b^2/a

-

Vertical coordinate of point on ellips Filename: < CONIC2.TK > Rev. 21 Feb 94 Sag of point on ellipse (Horiz Coord) semi-major axis (Horizontal Axis) semi-minor axis (Vertical Axis) Comment------Vertex radius at major axis Vertex radius at minor axis ELLIPSE \*\*\*\*\*\*\* ------RULE SHEET ====== Focus to Focus distance Focus to Vertex distance vertex curvature conic constant \*\*\*\*\*\* 1111 Output--- Unit----VARIABLE SHEET 2.7235235 .20643616 .00054422 1386.6627 1384.0312 OBLATE 1837.5 St Input---- Name---Rmajor Rminor EFF EVF ANXX υ đ 1523.081 1672.92 100

+ Sqrt(1 - (K + 1) \* C^2\*X^2)) / a^2 Rminor = a^2/b Rmajor = b^2/a EFF =2\* SQRT(abs(a^2-b^2)) IF K>0 THEN EVF = 'OBLATE ELSE EVF = a-EFF/2 b^2 - a^2) = C \* X^2 /(1 C = 1/((K+1)\*a)Rmajor=1/C ;Ellipse S Rule----11 N м

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2) Comparison of ellipse to circle

cut which includes the major axis. This is essentially the difference between removed from his starting point which is a sphere. The third shows the difference between the ellipse and a circle which best fits the ellipse along a sags. This is valuable to the optician since it tells how much stock must be Three graphs are presented. The first shows the actual sags of the ellipsoid in two axes. The second shows the difference between the two an ellipsoid and a torrus.

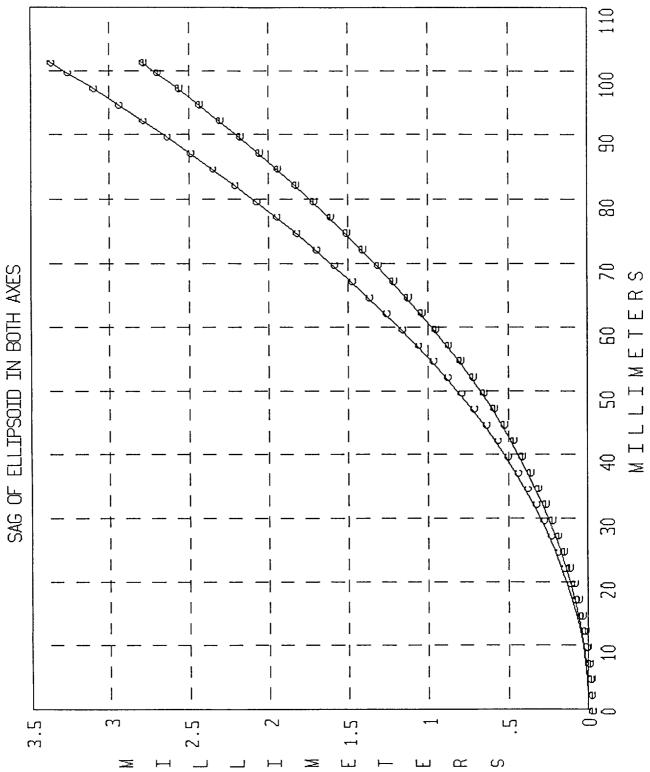
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st Input Name	.20643616 1837.5 203.2 72 72 0 1523.081	

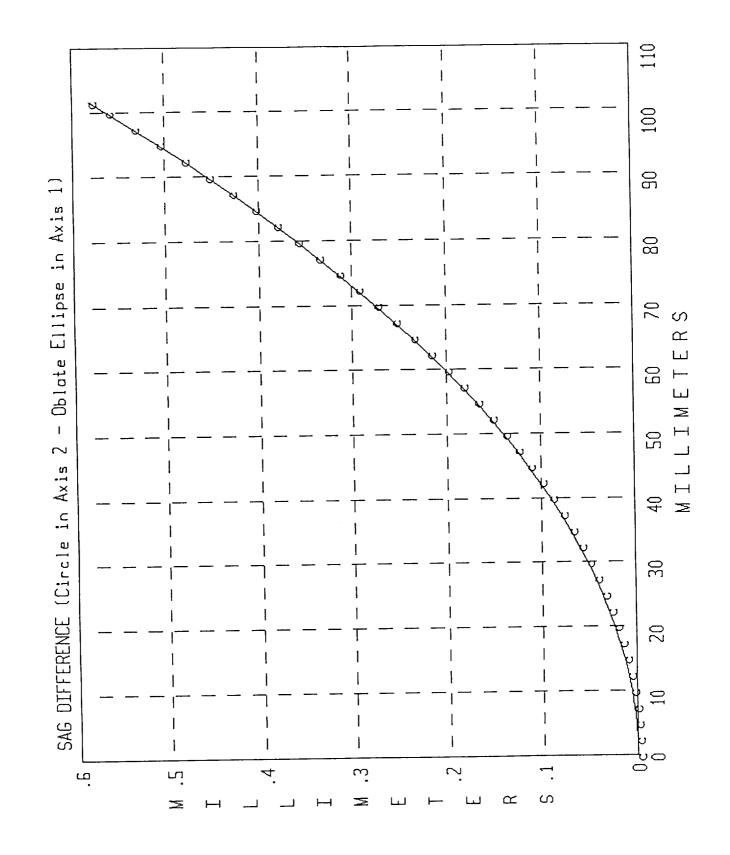
DZcirtan = Sgn(c)\*Ytan/SQRT(Rc^2-Ytan^2) dZdYtan = DZcirtan Zcirtan = 0 - Sgn(c)\*SQRT(Rc<sup>2</sup>-Ytan<sup>2</sup>) rx2tan = SQRT(1-(K2+1)\*c2^2\*Ytan^2) rxtan = SQRT(1-(K+1)\*c^2\*Ytan^2)  $Zcir = 0 - Sgn(c) * SQRT(Rc^2 - Y^2)$ D2sag = Z2 - Z; - (Z2tan - Ztan)  $DZcir = Sgn(c) * Y/SQRT(Rc^2-Y^2)$ =================== RULE SHEET rx2 = SQRT(1-(K2+1)\*c2^2\*Y^2) Z2 = c2\*Y^2/(1+rx2) Dsag2 = Z2 - Zcir  $Z2tan = c2*Ytan^2/(1+rx2tan)$  $CRL = Parfl+(ACD^2/(2*Rv))$ rx = SQRT(1-(K+1)\*c^2\*Y^2)
Z = c\*Y^2/(1+rx) Ztan = c\*Ytan<sup>2</sup>/(1+rxtan) Sind(OffAxAng) = ACD/CRL dZdYtan = c\*Ytan/rxtanDIAM/2 + OED = ACDDslope =dZdY-DZcir Rule-----Rule Parfl2 = 1/(c2\*2)dZdY2 = c2\*Y/rx2Parfl = 1/(c\*2)Ztan = Zcirtan dZdY = c\*Y/rxDsag = Z-ZcirRv2 = 1/c2RV = 1/CS

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	Y	2.5		7.5	0	12.5	15	17.5	0	22.5	പ	27.5	30	32.5	ഹ	37.5	0	42.5	പ	47.5	0	52.5	പ	57.5	0	62.5	പ	67.5	0	C.2/	ט / הינר	c.//	и С 0 С 0 С 0 С 0 С 0 С 0 С 0 С 0 С 0 С 0	c.28	85	87.5	0	92.5	0		100	$\sim$

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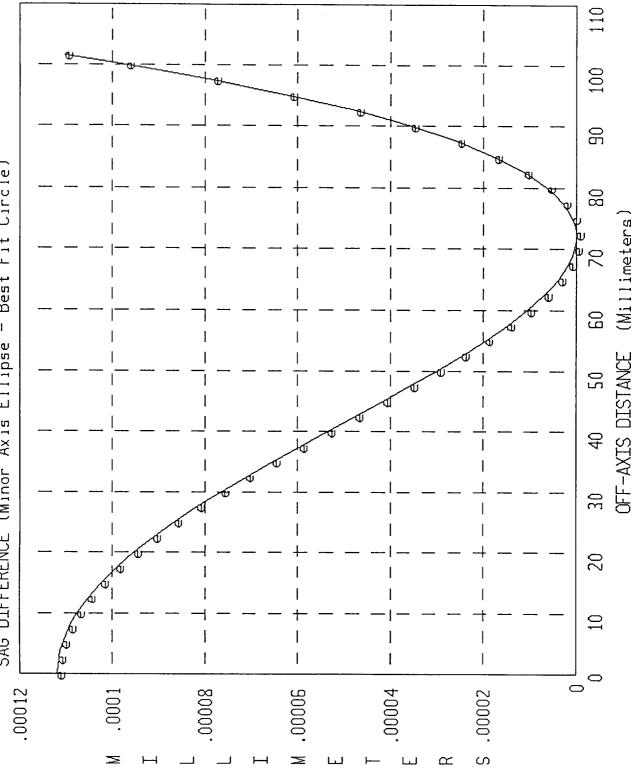




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SAG DIFFERENCE (Minor Axis Ellipse - Best Fit Circle)



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3) Deformation from ellipse

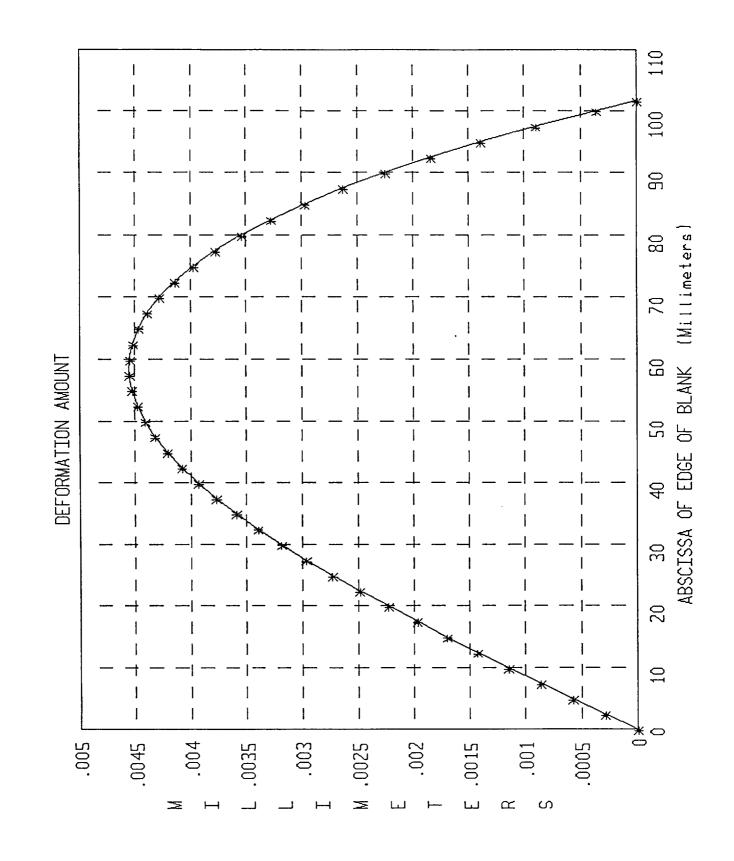
In these calculations, the departure from the ellipse is calculated. It is calculated at points around the clear aperture and then at points around the generate the points of equal deformation which were used to produce the edge of the blank. The graphs of this computation are used to show the maximum departure. The basic deformation equation was also used to topographical map of the surface and simulated fringe map.

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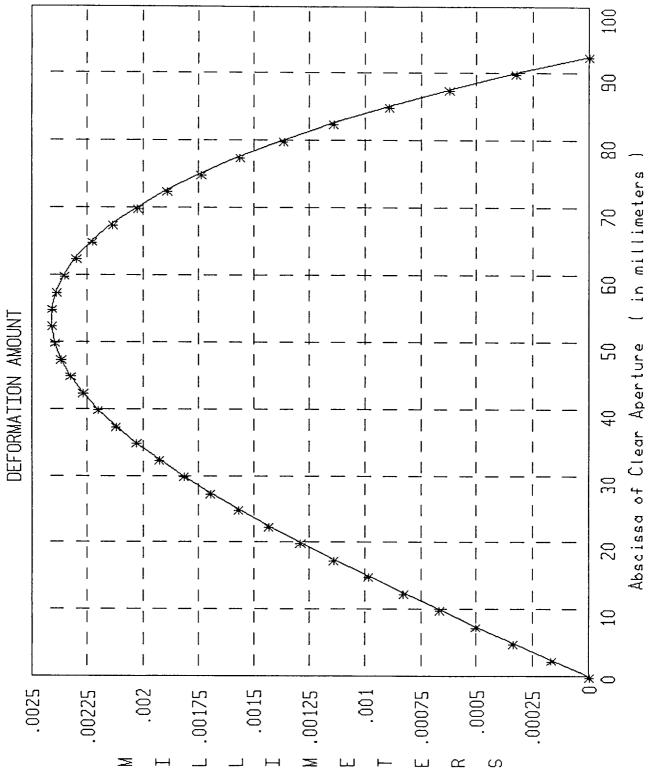
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37.5	0.845617	0212308	37.5	4.426214	0377159
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42.5	8.835402	022715	42.5	2.28385	0408271
45	7.710827	0232722	10	1.090943	0421184
47.5	6.501316	0236953	47.5	9.812638	0432192
0	5.202142	0239774	0	8.445237	0441192
52.5	3.807828	0241110	52.5	5.984538	0448076
ທ	2.311988	0240887	in	5.425757	0452740
57.5	0.707118	0239032	57.5	3.763416	0455077
0	3.984319	0235468	0	1.991219	0454982
62.5	7.132923	0230124	62.5	0.101872	0452349
0	5.139978	0222923	10	3.086874	0447073
67.5	2.989517	0213792	67.5	5.93622	00439046
0	0.661484	0202657	0	3.638033	0428165
72.5	3.130119	0189444	72.5	1.178016	00414322
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c.//	2.308736	0156484	77.5	5.698630	0377331
0	3.905367	0136589	~	2.630344	353971
82.5	5.049271	0114319	82.5	9.298482	327227
10	0.569601	089599	85	5.655727	296994
87.5	5.135135	062356	87.5	L.636324	263164
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92.5	.90161E-	123E-1	92.5	2.027491	184296
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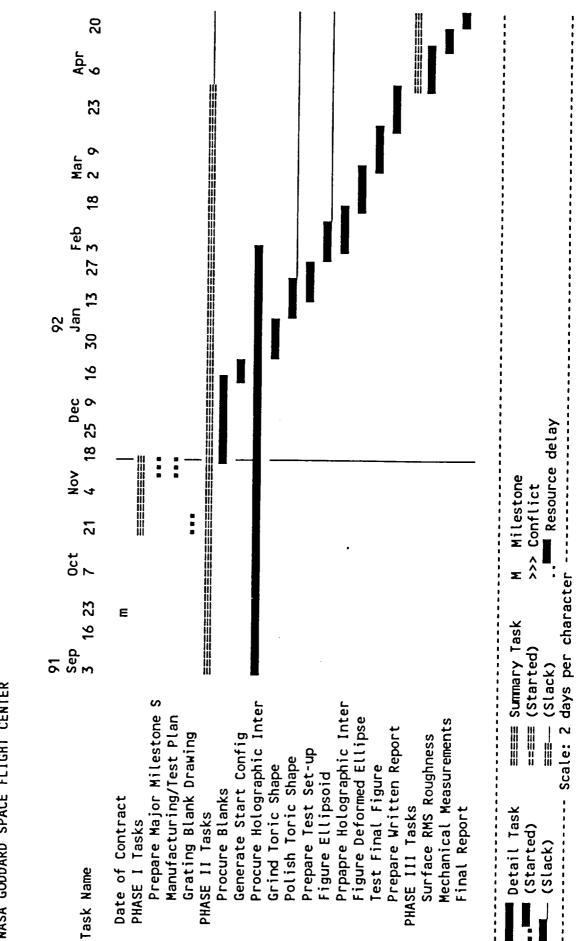


## APPENDIX B

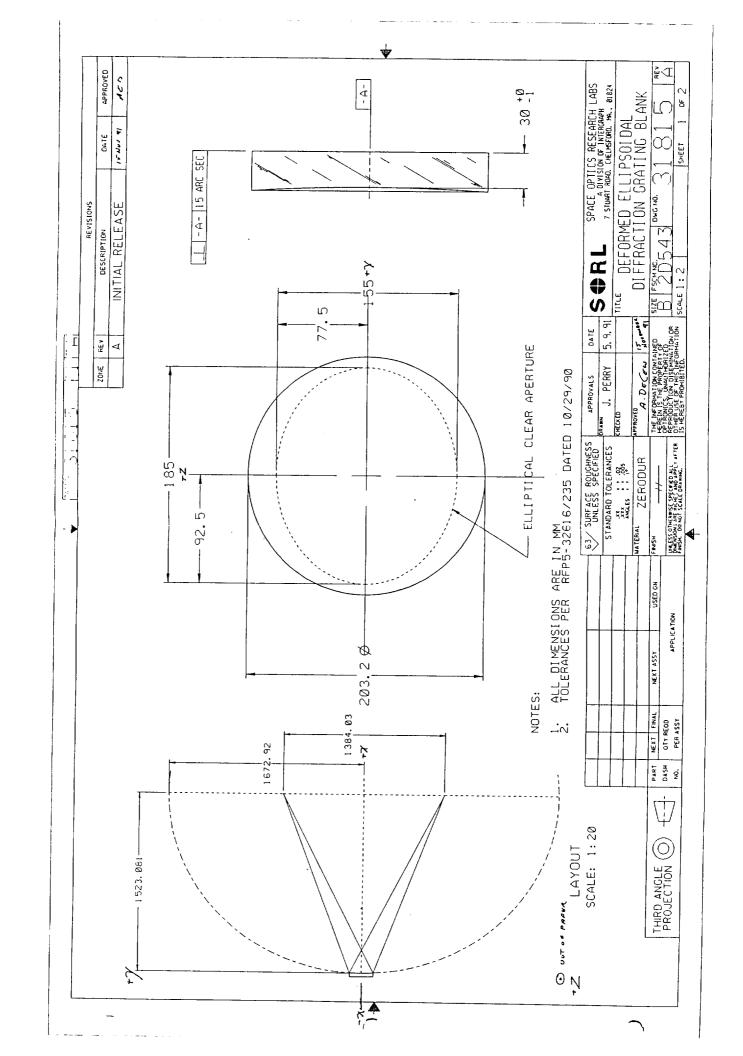
## MANUFACTURING AND TEST PLAN

Schedule File : C:\TL3\DATA\TEMPLATE Schedule Name : DEFORMED ELLIPSOIDAL DIFFRACTION GRATING BLANK : 15-Nov-91 9:00am : Alan DeCew, Jr. Responsible As-of Date

NASA GODDARD SPACE FLIGHT CENTER

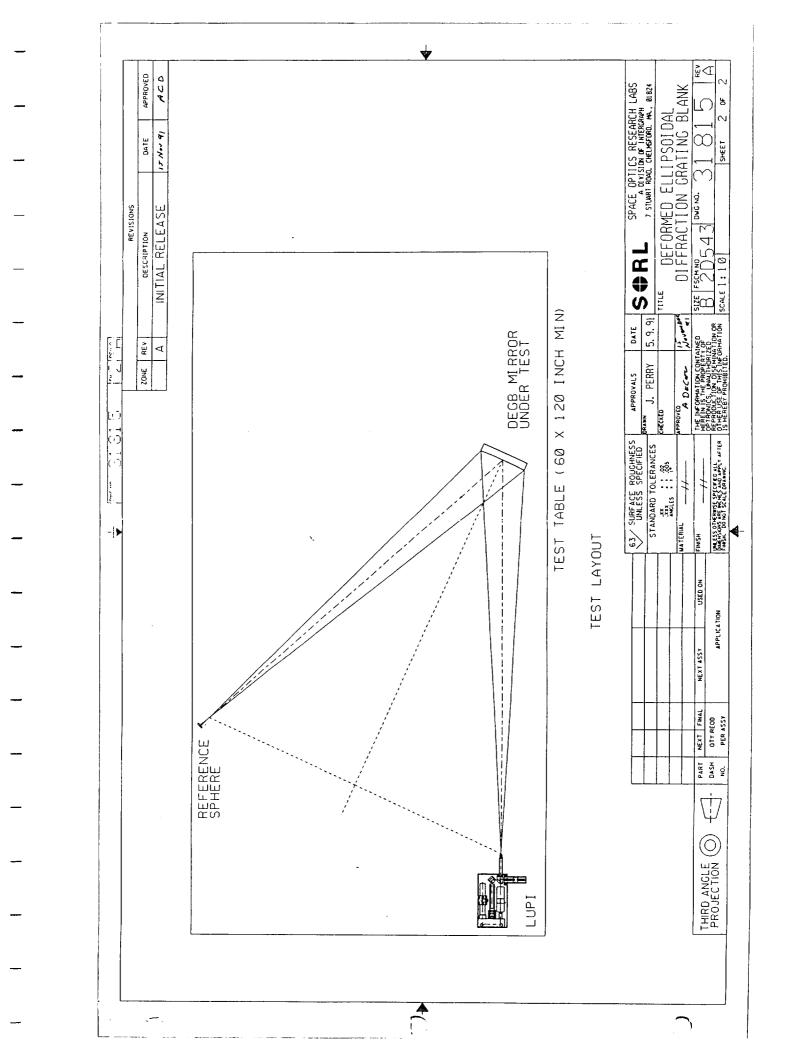


TIME LINE Gantt Chart Report, Strip



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## The Manufacturing/Test Plan

## 1M) PROCURING THE BLANKS

The Zerodur will be procured from Schott. The diameter of 8" leaves sufficient polishing area beyond the clear aperture. The blank will be ordered 8" Diameter with 33.4mm thickness. The thickness will provide for a 30mm vertex thickness after the figure is obtained.

## 2M) GENERATING THE STARTING CONFIGURATION

SORL will edge and generate the Zerodur blanks into the starting configuration. The generated blanks will then be a positive/negative spherical pair. The radius will be the shorter Rx = 59.964".

## 3M) GRINDING THE TORIC SHAPE

The spherical blanks will be formed into a toric pair using loose abrasive grinding. The parts are ground together using 40 micron Aluminum Oxide grit.

### 3T) TESTING THE TORIC SHAPE

SORL has a computer controlled 3 axis measurement machine which will be used to monitor the progress of the toric grinding. This will provide the information to the optician about the actual toric radii.

## 4M) POLISHING THE TORIC SHAPE

Having verified the basic toric radii, the next step is to polish the surface so interferometry can be used to test the surface. This begins with working down through the grit sizes from 40 to 30 to 15 to 10 and finally 5 micron.

The fine ground surface is then polished with a full size pitch lap and Cerium Oxide polishing compound. Both sides of the test plate pair are polished.

## 4T) OPTICALLY TESTING THE TORIC

The polished surfaces of the test plate pair allow the viewing of interference fringes between the surfaces. This is an accurate test of the toric surface since the test plates can be translated and flipped end for end. The optician uses this method to insure the smooth toric shape. The concave toric is then tested in the interferometric configuration shown in SORL DWG # 31815 Sheet 2 of 2. This test preparation is not performed by the optician. A member of the test engineering staff is always involved in preparing all new test set-ups. The difference between the toric and a perfectly imaging ellipsoid is very small, less than 1/5 wave. The toric is, therefore, an ideal starting point to use to verify the operation of the holographic interferometer.

## 5M) FIGURING THE TORIC INTO AN ELLIPSOID

The ellipsoid which will provide a perfect null test in the focus to focus configuration can be figured from the toroid by the removal of only 5 microinches of material. Using local polishing techniques, the optician will hand figure the toric into the ellipsoid.

## 5T) PREPARING THE HOLOGRAPHIC INTERFEROMETER

The holographic interferometer comes with a system hologram which nulls out any residual aberrations of the instrument. This can be readily checked against our reference sphere. Using this configuration, the figure of the toric can be measured. The starting point will show irregular fringes which identify high points corresponding to the difference between the toric and the ellipse. The polishing will continue until the fringe pattern in the interferometer shows straight fringes.

At this point, the hologram for nulling the effects of the deformation coefficients will be inserted into the interferometer. After alignment, the fringe pattern should match the calculated interferogram presented in the proposal as figure 1. This will not verify the hologram to 1/10 wave accuracy, but it will verify that no significant calculation or transcription error has occurred.

## 6M) FIGURING THE ELLIPSE INTO THE DEFORMED ELLIPSE

The figuring of the deformed ellipse requires the removal of as much as 8 microns of material. This is more than 30 times what was removed to change the toric to an ellipse. It is, however, still a small correction, and will be polished into the ellipse. The optician will again be using small laps and monitor the progress by viewing the straightness of the fringes in the interferometer.

## 6T) TESTING THE DEFORMED ELLIPSE

The holographic interferometer will be used throughout the manufacturing process to evaluate the progress of the figuring. The optician does not require additional engineering of test support until the part is felt to be ready. Hardcopy of the interferograms are then taken and reduced by the test engineering department. Meeting the 1/10 wave surface accuracy completes the requirements of Phase II.

## 7M) CREATING A 2.0nm RMS ROUGHNESS SURFACE

This step runs concurrently with the figuring. The optician, knowing that a smooth surface is a requirement, will approach the final figuring with a controlled polishing technique. As the final figure is approached, the stroking becomes lighter and the randomness of the direction of stroking is carefully maintained.

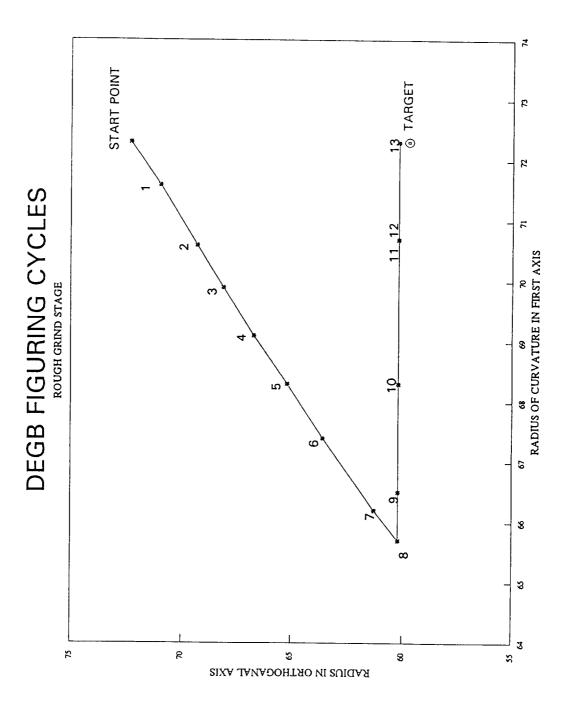
## 7T) TESTING THE RMS SURFACE ROUGHNESS

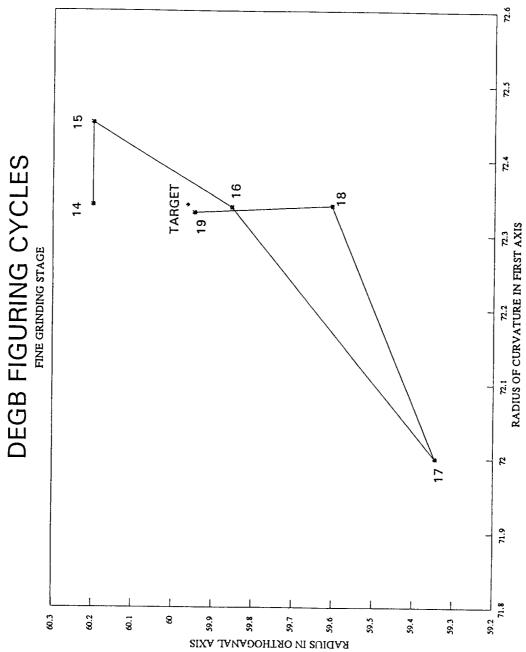
SORL does not have a profilometer capable of measuring surface roughness to this level. The part will be sent to an outside vendor for measuring. The data will be obtained from a Wyko or Chapman profilometer.

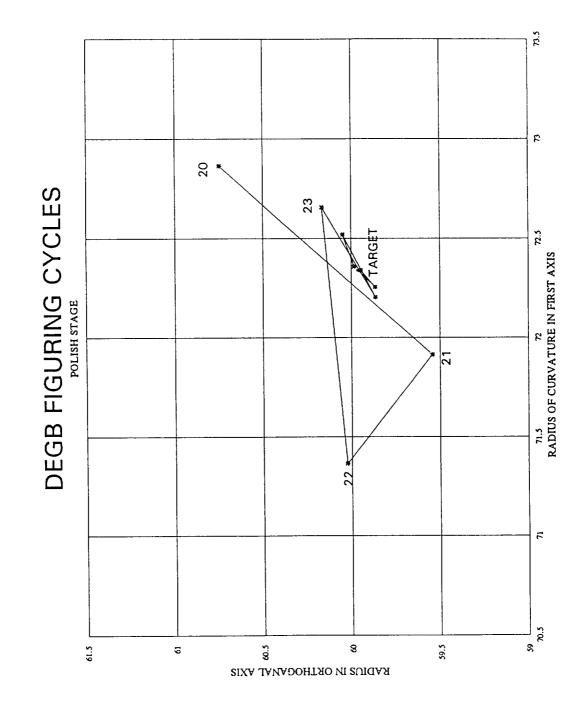
## APPENDIX C

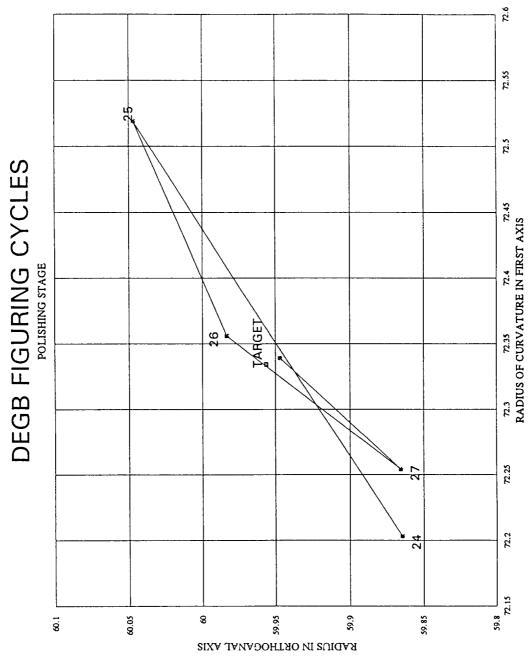
## FIGURING CYCLES TO FABRICATE DEGB

- Cycles required to take the generated spherical shape to the torrus which best approximates the ellipse.
- 2) Progess of the figure polishing of the deformation into the final DEGB configuration.







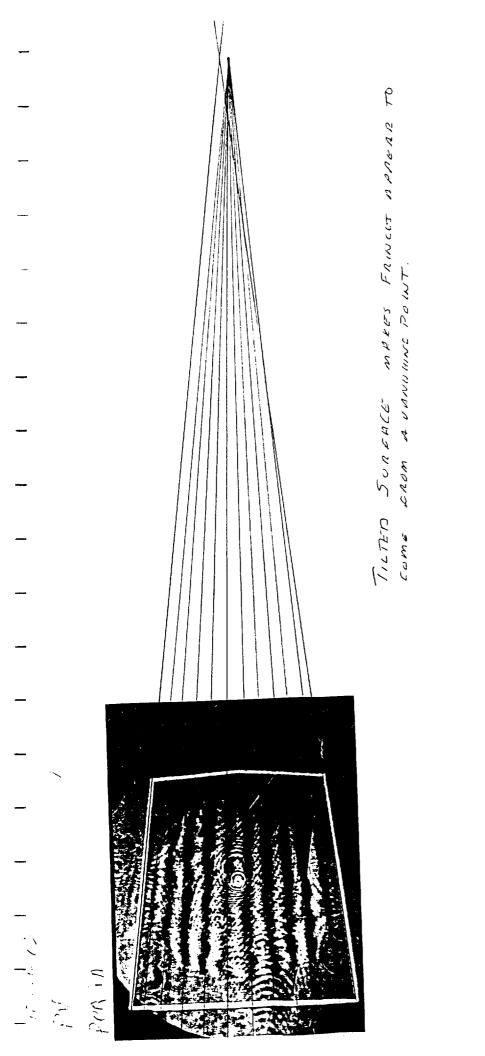


## Raw data used for figuring cycle graphs

13	72.3	60.2			
12	70.7	60.2			
11	70.7	60.2			
10	68.3	60.2			
თ	66.5	60.2			
ß	65.7	60.2			
7	66.2	61.3			
9	67.4	63.6			
ۍ	68.3	65.2	19	72.33	59.947
4	69.1	66.7	18	72.339	59.604
e	69.9	68.1	17	72	59.345
2	70.6	69.3	16	72.337	59.854
-	71.6	70.98	15	72.45	60.2
0	72.3	72.3	14	72.34	60.2

28	72.339	59.947
27	72.254	59.866
26	72.356	59.983
25	72.519	60.047
24	72.203	59.865
23	72.656	60.165
22	71.365	60.024
21	71.915	59.546
20	72.863	60.748

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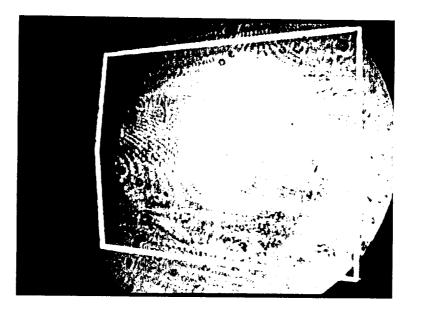
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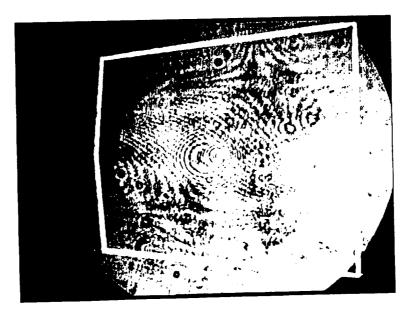
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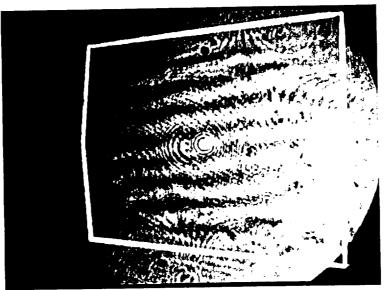
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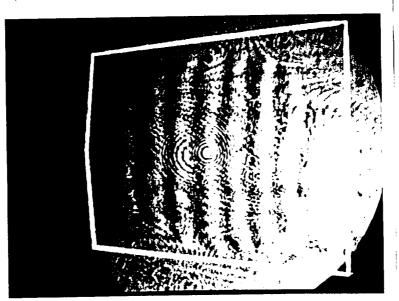
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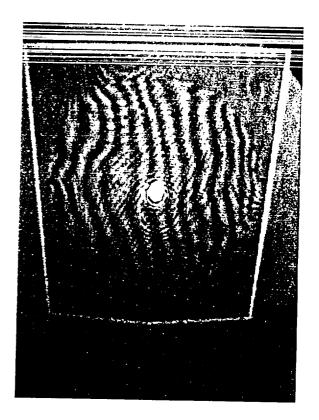




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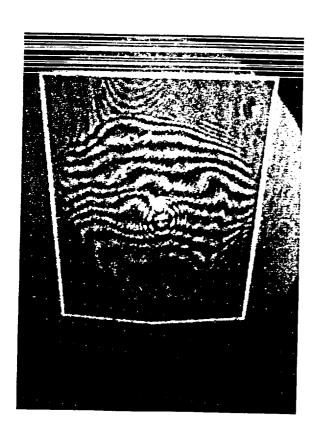
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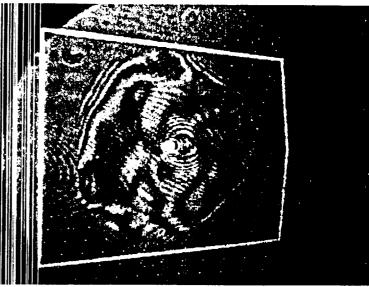


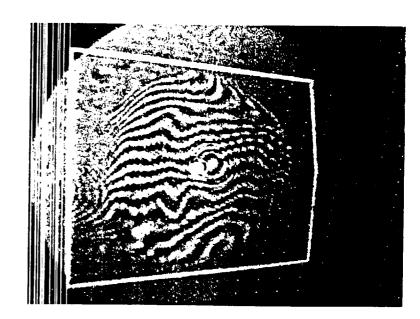
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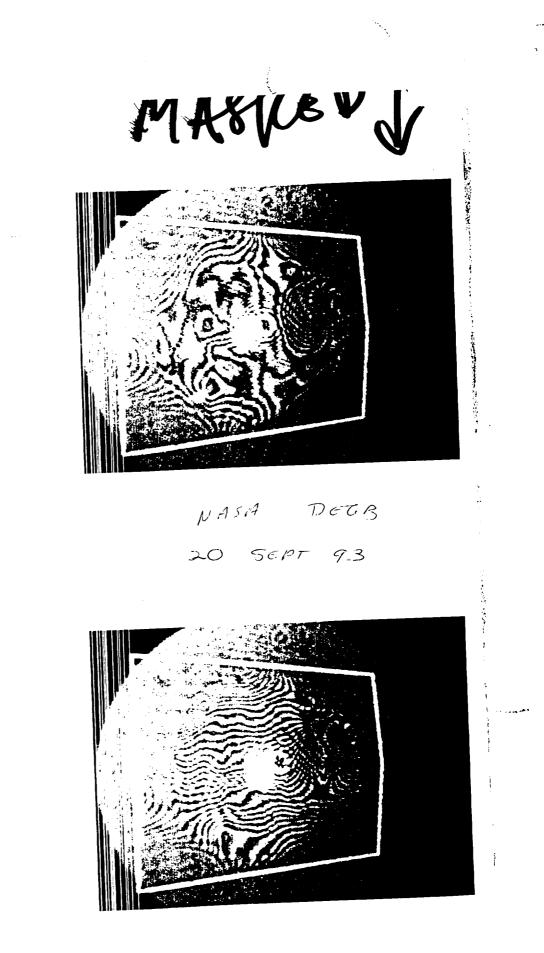


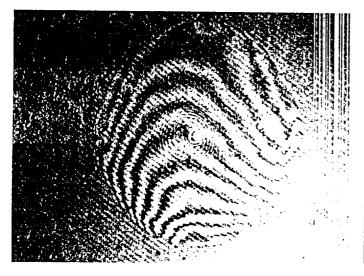


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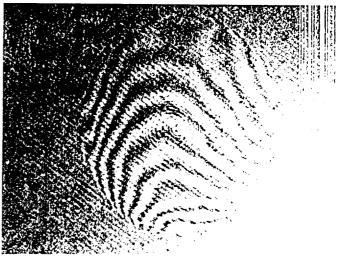
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### APPENDIX D

# INTERFEROMETER FOR ASPHERIC TESTING

- Arnold, Steven M., Jain, Anil K., "An interferometer for testing of general aspherics using computer generated holograms", SPIE Vol 1396, 1990.
- Arnold, Steven M., "An interferometer for aspheric testing: Calibration and error compensation", ASPE Proceedings, April 1992. 5
- 3) Calibration curves for IAT SN 003.

### An interferometer for testing of general aspherics using computer generated holograms

Steven M. Arnold and Anil K. Jain

APA Optics Inc. 2950 N.E. 84th Lane, Minneapolis, Minnesota 55434 PREPRINT SPIE Vol. 1396 27-28 September 1990 Chicago

### ABSTRACT

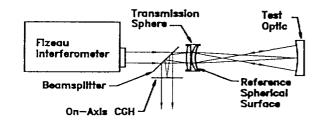
An interferometer for aspheric testing (IAT) is under development at APA Optics for testing of general aspherics using inexpensive electron-beam written computer generated holograms (CGHs) as null compensators. This 152-mm aperture Twyman-Green interferometer is compatible with standard transmission spheres, fringe analysis software, and phase measuring accessories. Aspheric departures of up to several hundred waves can be measured using only standard interferometer accessories. Deeper aspheres may be tested using simple auxiliary optics. The interferometer configuration, methods of operation, and performance specifications are presented.

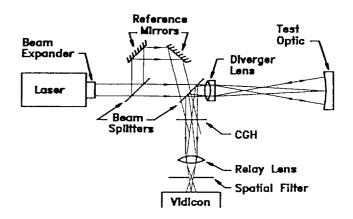
### **1. INTRODUCTION**

Modern optical systems demand the use of aspheric surfaces to improve system performance and reduce the number of optical components required. Fabrication of precision aspheres is becoming less expensive because of advances in such techniques as diamond turning and automated polishing. However, these aspheric components still require testing, and optical methods are preferred. Stylus methods are slow, usually require physical contact, and cannot test for index inhomogeneities or assembly errors. The standard interferometric method of testing an aspheric surface (or lens) is to use a Fizeau interferometer with auxilliary optics to produce a wavefront matching the (ideal) aspheric surface being tested. The problem with these refractive or reflective null compensators is that they are expensive, they require a long lead time, and a separate compensator must be made for each asphere design tested.

Much has been written on the use of CGHs as aspheric null compensators. A recent paper by the author reviews this work<sup>1</sup>. With the addition of a beamsplitter between the mainframe and transmission sphere, a commercial Fizeau-type interferometer can be adapted to test aspheres using CGHs<sup>2</sup>. Depicted in Figure 1, this testing technique has several limitations: convex test surfaces cannot be imaged onto the CGH, a necessary condition for strong aspheric wavefronts; the on-axis Moiré-type fringe pattern is difficult to interpret for complicated wavefronts where the slope reverses sign; and the technique does not lend itself to phase measuring. Off-axis CGHs have been used in the test arm of a Twyman-Green or laser unequal path (LUPI) interferometer. This offers the advantage of a spatially filtered (non-Moiré) interferogram but does not solve the pupil imaging limitation and requires an expensive high diffraction efficiency CGH on a high quality substrate.

Figure 2 depicts a preferred configuration: a Twyman-Green interferometer with a CGH in the viewing arm. Key features include imaging of the test surface onto the CGH, placement of the CGH in collimated test and reference beams so that both pass similarly through the substrate, and spatial filtering of diffracted orders. Although several interferometers of this general configuration have been built, none are currently available commercially.





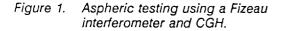


Figure 2. Twyman-Green (LUPI) interferometer with CGH in viewing arm.

### 2. REQUIREMENTS AND DESIGN GOALS

Our goal is a commercial instrument which can provide routine and economical interferometric testing of aspheric optics. To minimize the need for expensive null lenses, the instrument must be capable of testing aspheres which are much deeper than the few tens of fringes possible with currently available phase-measuring interferometers. Other design goals include  $\lambda/10$  system accuracy, multi-aperture capability, pupil imaging of both concave and convex aspheres, ease of use, phase measuring capability, and compatibility with industry standard accessories and interferogram analysis software.

Our approach makes maximum use of inexpensive e-beam written CGHs to null the test aspheres, to ease manufacturing tolerances by compensating known instrument imperfections, and to serve as alignment and calibration tools. The interferometer design therefore begins at the CGH. For unwanted diffraction orders to be blocked by spatial filtering, a CGH must include an off-axis carrier such that the minimum spatial frequency is more than half the maximum spatial frequency. Since the diffracted reference beam matches the test beam (or vice versa), test and reference beam ray slopes,  $\theta_t$  and  $\theta_r$ , are related to the local spatial frequency,  $1/\Lambda$ , through the grating equation

$$\sin \theta_{\rm t} = \sin \theta_{\rm r} - \lambda/\Lambda = \lambda/\Lambda_{\rm o} - \lambda/\Lambda. \tag{1}$$

Table 1 lists CGH bandwidths, linewidths ( $\Lambda/2$ ), ray slopes, and wavefront accuracies ( $\lambda/2 = 1$  fringe) achievable with standard (*i.e.* mask shop) e-beam resolutions. The collimated test or reference beam diameter will define the maximum CGH aperture. Other apertures in the system will depend on the CGH bandwidth (ray slopes) and pathlengths. We decided on a 20-mm CGH aperture and 38-mm cube beamsplitters as a reasonable compromise between functionality and cost. These choices will accomodate ray slopes of 100 to 150 lpmm (the practical limit of e-beam lithography) and allow a one-dimensional CGH space-bandwidth product of 2000 to 3000. This is adequate for testing of aspheric deformations (less tilt and focus) as large as several hundred fringes without the use of additional null optics.

### 3. OPTICAL SCHEMATIC

Figure 3 depicts the overall configuration of the APA Optics interferometer for aspheric testing (IAT). Its construction is modular, consisting of a Twyman-Green or LUPI mainframe of 20-mm beam diameter, detachable beam expanding telescopes (152-mm aperture shown), and a selection of accessory optics.

### 3.1 Interferometer Mainframe

The interferometer mainframe uses two polarizing beamsplitter cubes with attached  $\lambda/4$  plates to establish separate paths for orthogonally polarized test and reference beams. This approach is shared by an earlier holographic interferometer familiar to the author<sup>3</sup>. The He-Ne laser beam, after being spatially filtered and expanded to 20 mm diameter, is divided by the first cube into a transmitted reference beam and a reflected test beam. The test beam, which is circularly polarized as it exits the cube through a  $\lambda/4$  plate, is converted to the orthogonal linear polarization upon its return from the test arm and is therefore transmitted through the cube to the pupil relay arm. After being focused, retroreflected and recollimated by the pupil relay, the test beam is transformed back to its original state of linear polarization and is thus reflected by the beamsplitter cube to rejoin the reference beam. Test and reference beams then copropagate to the second cube with opposite linear polarizations; the reference beam is collimated whereas the test beam will have acquired up to several hundred waves of aberration from the test asphere.

On entering the second polarizing beamsplitter cube, the test beam is transmitted whereas the reference beam is diverted to a pair of reference mirrors which control its tilt and decentration. (For graphic purposes, both cubes are shown as oriented similarly; the second cube and its reference mirrors must be rotated 90 degrees about the optic axis in order that the test beam and not the reference beam be transmitted.)

CGH bandwidth	Minimum linewidth (µm)	Ray slopes (degrees)	Accuracy (waves) vs. e-beam resolution			tion
(lpmm)			$1.0\mu\mathrm{m}$	0.5 μm	0.25 µm	0.1 µm
30±10	12.5	±0.363	λ/50	λ/100	λ/200	λ/500
75±25	5.0	±0.906	λ/20	λ/40	λ/80	λ/200
$150 \pm 50$	2.5	±1.813	_	λ/20	λ/40	$\lambda/100$
$300 \pm 100$	1.25	±3.628	_	_	λ/20	λ/50
375±125	1.0	±4.537	_	_	_	λ/40
450±150	0.83	$\pm 5.447$	_	—		λ/33

Table 1. Capabilities of e-beam written CGH photomasks.

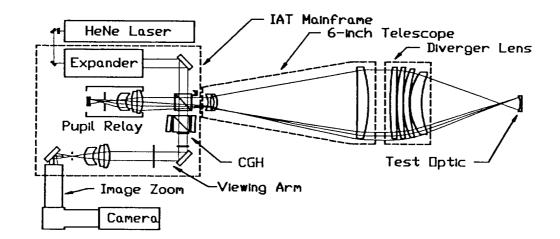


Figure 3. Overall configuration of proposed interferometer for aspheric testing (IAT) with attached 152-mm telescope and f/1.5 diverger lens.

On exiting the second beamsplitter cube, the test and reference beams would form an off-axis interferogram were it not for their opposite linear polarizations. Instead, the test and reference paths are raytraced from the expanded, collimated source to determine this expected interferogram which is then encoded as an e-beam written computer generated hologram. The CGH diffracts the reference beam, adding aberrations to match those expected in the test beam while simultaneously removing the reference beam tilt. The diffracted reference beam, in combination with the undiffracted (zero order) test beam, forms a nominally null interferogram. The viewing arm optics relay this interferogram to the video camera while also spatially filtering unwanted diffraction orders and converting the two beams to a like state of linear polarization by use of a diagonally oriented linear analyzer.

Like most Twyman-Green or LUPI interferometers, the proposed IAT can also operate in Fizeau mode. This requires blocking the Twyman-Green reference path (by rotating a waveplate and analyzer) and introducing a partially reflecting Fizeau transmission flat or sphere into the test arm, thus producing an on-axis reference beam of the same polarization as the test beam. Both beams then return via the pupil relay arm to the CGH.

### 3.2 Pupil Relay Arm

The beam expanding telescope and diverger lens will form a first image of the test surface somewhere in the vicinity of the mainframe-telescope junction. This 20-mm diameter first image is real if it lies between the telescope and the beamsplitter cube and virtual otherwise. The required image location will vary depending on the choice of diverger optics, the test surface radius, and whether the test surface is concave or convex. For concave (beyond focus) test surfaces, the first pupil image may lie outside the telescope, moving closer to and eventually into the telescope for larger test surface radii. For convex test surfaces, the image lies inside the telescope—further inside for smaller test surface radii. A generous range of image locations centered about the telescope/mainframe junction accomodates all but very small radius concave and convex parts.

The pupil relay arm transfers the first pupil image to the CGH. The relay optics consist of a multi-element focusing lens and a confocal spherical mirror for retroreflection. Both are mounted on a mechanical slide providing 150 mm of travel. From a pupil imaging standpoint, the spherical mirror acts as a field lens. The relay optics were designed by raytracing the test beam (backwards) from the CGH, through the beamsplitter cubes, out the relay arm, and back through the first cube to five different first pupil image locations, using the position of the mechanical slide as a zoom parameter. Three of the five zoom positions and their associated image locations are shown in Figure 4. An aperture stop at the CGH restricted design ray slopes to  $\pm 3.63$  degrees ( $\pm 100$  lpmm). In order to make the multi-element focusing lens separately testable, we required that it form a near perfect focus when presented with a 20-mm on-axis collimated beam. This required a sixth zoom configuration. The relay optics design did not attempt to compensate aberrations of the test arm optics (which were not included in the raytrace).

The focusing lens is a cemented pair of plano-convex elements plus a thick meniscus corrector lens. The pupil imaging performance of the relay optics design is better than 5  $\lambda$  for all five zoom positions and ray slopes up to  $\pm 100$  lpmm. The double pass de-collimation (*i.e.* zero ray slope performance) is better than  $\lambda/20$ . The approximate range of first pupil image locations is from 80 mm inside to 200 mm outside the test arm face of the first beamsplitter cube.

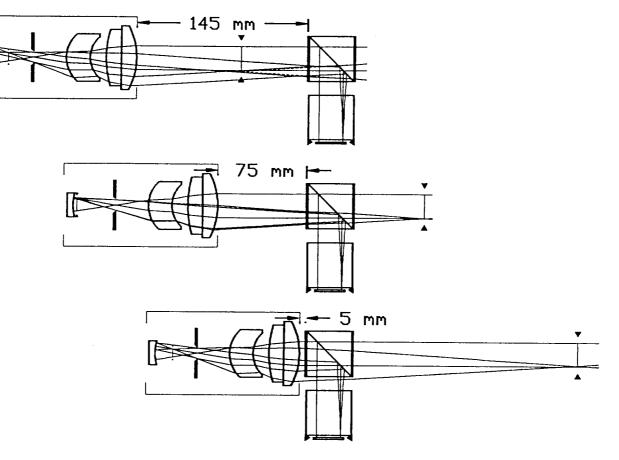


Figure 4. Pupil relay optics for three locations of first pupil image.

### 3.3 Test Arm

Test arm magnification should match the required test aperture. Excess magnification reduces the scale of the CGH, increasing its cost and/or forcing a reduction in space-bandwidth. The mainframe's unexpanded 20-mm beam, together with a diverger lens of suitable f/#, is adequate for testing most concave or small convex parts. A telescope is required only for testing of larger convex surfaces or for compatibility with standard Fizeau accessories.

We have designed or are designing test arm optics for a range of apertures and f/#'s. One completed design is for an f/3 Galilean telescope to expand the test beam from 20 mm to 152 mm. This telescope is atypical in that the negative lens aperture must be large enough to accomodate a returning aspheric wavefront. The design seeks to minimize coma for locations of the first pupil image which lie within the focus range of the pupil relay optics (cf. Figure 4). A desire to minimize vignetting without increasing the aperture of the objective lens beyond 152 mm drives the design to low f/#. The apertures of the two smaller lenses are sized to pass all such rays which are not vignetted by the 152-mm aperture objective. The expander requirement is for 7.5X beam expansion and  $\lambda/10$  collimation. The competing pupil imaging requirements can be expressed as a 20-mm field of view for a 280-mm range of image locations with an f/8 aperture stop at the image.

### 3.4 Reference Path

The function of the second beamsplitter cube and reference mirrors is to tilt the reference beam without decentration at the CGH plane. Figure 5 shows the ray paths and mirror configuration for selected CGH carrier frequencies. Note that the diffracted reference beam acquires aberrations to match those expected in the zero order (undiffracted) test beam. The 38-mm cube is decentered by 4 mm to accomodate larger reference path tilts without vignetting. The first reference mirror includes one axis of coarse tilt adjustment and a piezo transducer for phase measuring. The second reference mirror, which tilts in both axes, is used to fine adjust the number and orientation of interferogram fringes. When operating in Fizeau mode, an on-axis reference beam is provided by the final surface of a transmission sphere or flat and the orthogonally polarized Twyman-Green reference beam is blocked.

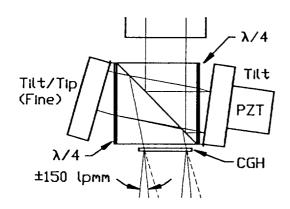


Figure 5. Reference path for CGH of 450±150 lpmm.

### 3.5 Viewing Arm

The viewing arm optics, shown in Figure 6, relay the 20-mm diameter interferogram at the CGH onto the 11-mm diagonal video camera while also spatially filtering unwanted diffraction orders and converting the two beams to a like polarization state by use of a diagonally oriented linear analyzer. The imaging requirements of the relay lens are rather unusual in that it must cover a full field of view with low distortion and no vignetting for an aperture stop located at the object (CGH). Because only two rays of interest (zero-order test and first-order reference) come simultaneously from any point on the CGH, and these rays are nominally copropagating, resolution is relevant only in the sense of an aperture dependance of distortion. Somewhat coincidentally, these requirements are well met by the multi-element lens previously designed for the pupil imaging relay. In Fizeau mode, an on-axis CGH collimates the first-order test beam to match the zero-order reference beam. A small aperture then blocks most light of other (spurious) diffraction orders.

Figure 3 included zoom optics between the first spatially filtered image and the camera. It appears that 5:1 image zooming can be provided by commercially available optics and a field lens, although this option has not been verified.

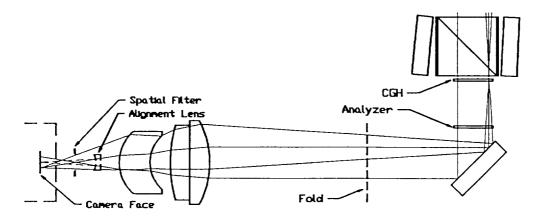
### **4. USER INTERFACE**

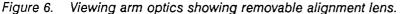
Although the Aspheric Interferometer is of greater optical complexity than the usual Fizeau interferometer, the user interface is similar in most respects. The primary differences relate to the CGH.

### 4.1 Test Arm Configuration and CGH Design

Planning begins with the specification of a test arm configuration to minimize ray slopes at the CGH. Typically this will involve selecting a combination of diverger lens and/or beam expander to produce a best-fit spherical wavefront. A partial null lens will be required for very deep aspheres. However, with a CGH to accomodate the higher order aberrations, a singlet can often suffice, both as diverger lens and to null the lower order aberrations.<sup>4,5</sup>

For each variety of asphere to be tested, a CGH must be designed and fabricated. The CGH design is performed using commercially available optical design software, most likely the same package as was used to design the asphere. The CGH is defined by its phase function which represents the expected optical path difference between the test and





reference beams as measured at the CGH plane. Design of the CGH requires a complete description of the interferometer and test arm optics. The raytrace begins with the 20-mm diameter collimated laser beam entering the first beamsplitter cube and ends at the CGH. The source arm and viewing arm optics need not be raytraced as they are common to both beams. The CGH design process includes specifying a test arm configuration to minimize ray slopes at the CGH and chosing a position for the relay arm slide to provide acceptable pupil imaging.

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Although initial setup of the raytrace model and CGH merit function is quite complicated, modification for subsequent asphere testing should be routine. Since the reference path varies only in the mirror tilts which are uniquely determined by the CGH carrier frequency, it is not necessary to explicitly raytrace this path for each CGH design; OPD maps for various carrier frequencies are instead included as part of the interferometer documentation and subtracted from CGH designs. Similarly, the Fizeau mode reference path can be documented for each beam expander and transmission sphere configuration. Given the baseline raytrace model for each expander/transmission sphere configuration, the designer need only modify the test optic description and, for very deep aspheres, insert a refractive null compensator.

The test arm is configured in accordance with the CGH design. Aspheric test surfaces will require 5-axis or 6-axis positioning. Similarly, since the test wavefront on returning through the diverger lens has acquired the aspheric aberrations of the test surface, it is necessary that the diverger lens be coaxial with the interferometer. These alignment complexities are characteristic of any interferometric testing of aspheres and have been addressed elsewhere<sup>6</sup>.

### 4.2 Interferometer Operation

Figure 7 identifies those parts of the interferometer which are available for manipulation by the user. The CGH is then inserted and positioned so its alignment fiducials match a corresponding pattern on the exit face of the second beamsplitter cube when viewed on the video monitor. The pupil relay slide and reference mirrors are set to values specified by the CGH design. These settings are inscribed onto the CGH and readable on the video monitor. The aperture turret is rotated to bring the negative alignment lens into the viewing path, thus focusing the test and reference beams onto the camera face. The test surface and second reference mirror are tilted to bring both focused spots to the center of the camera field. A small corner cube can be moved into the test arm path, providing a third focused spot to define the center. This corner cube is also available to verify collimation of the pupil relay optics which is adjustable via a focus screw on the retro-mirror.

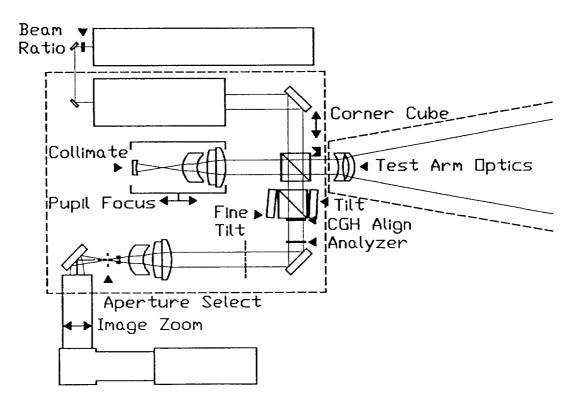


Figure 7. IAT user interface.

The corner cube is removed and the aperture turret is rotated to display the interferogram on the monitor. The smallest aperture which passes the entire interferogram is selected, blocking any unwanted diffraction orders. The apertures are labeled  $\pm 150$ ,  $\pm 100$ , etc. denoting the CGH bandwidth in lpmm. A  $\lambda/2$  plate prior to the beam expander adjusts the intensity ratio of the orthogonally polarized test and reference beams while the analyzer in the viewing arm determines the fraction of each beam which will reach the camera. By rotatation of the waveplate and analyzer, the user can manipulate both the intensity and contrast of the interferogram to suit any test surface reflectivity.

If the test surface includes alignment fiducials which are visible on the monitor, then it is translated and/or rotated to bring these fiducials into coincidence with corresponding marks on the CGH. It is then a matter of further manipulating the test surface to achieve straight, equally spaced fringes. For interferogram recording and analysis, both the video signal (RS-170) and the phase-shifting piezo transducer are plug compatible with commercially available fringe analysis and phase measuring attachments.

Flat surfaces can be tested in Twyman-Green mode with a generic CGH serving only to compensate aberrations of the interferometer optics. Flat or spherical surfaces (or shallow aspheres) can also be tested in Fizeau mode without the use of a CGH. One simply attaches the appropriate Fizeau transmission flat or sphere to produce a test arm reference beam and rotates the viewing arm analyzer to block the orthogonally polarized Twyman-Green reference beam.

### 5. TOLERANCING AND MANUFACTURABILITY

A complete sensitivity and tolerance analysis has been performed using methods similar to those outlined by Ginsberg<sup>7</sup>. The analysis clearly shows that the instrument cannot be economically built to provide  $\lambda/10$  testing capability without a high degree of software compensation. This should come as no surprise considering the optical complexity of the Aspheric Interferometer and that most components are unique to either the test or the reference path. The ability of a CGH to provide this compensation while also nulling the test asphere is what makes the Aspheric Interferometer a viable instrument for routine aspheric testing. Each CGH design is based on a full raytrace of the interferometer from expanded laser beam to CGH. Therefore, a precise interferometer is not necessarily required – providing we have a sufficiently precise optical description in a format acceptable to the optical design software.

Lens thicknesses, spacings, surface radii, and glass refractive indices can all be measured to much better accuracy than they can be manufactured. We therefore place fairly loose manufacturing tolerances on these parameters. If surface irregularities are describable by low order Zernike polynomials, they can be characterized to  $\lambda/50$  or better using phase measuring interferometry and modeled by several optical design codes. We therefore place fairly tight ( $\lambda/10$ ) tolerances on surface figures and further require that irregularities be well behaved. Since only the central zone of each element is ever used in double pass, we can double the tolerances for the larger aperture. In designing for fairly sturdy lens thicknesses and loose thickness tolerances, we reduce the difficulty of achieving these required irregularity tolerances.

Because refractive index inhomogeneities are difficult to characterize or model, they will be a major source of residual error. Several methods are used to minimize these error contributions. First, we specify Schlieren grades of glass wherever justified. Second, the beamsplitter cubes, pupil relay lens, and telescope are all designed as individually testable afocal subassemblies. This permits localization of residual errors and insertion of compensating dummy surfaces into the raytrace model. Third, by designing with identical components in both critical and non-critical locations, we have a better inventory of components from which to select.

This brief paper cannot detail the considerable opto-mechanical design involved in the aspheric interferometer. Primary mechanical considerations are component centration and thermal stability. A defocus of 26 microns in the pupil relay or 18 microns in the 152-mm telescope will produce 1 fringe of aberration at the interferogram. Both subassemblies are therefore designed for thermal stability through initial glass and metal choices and finally by fine-tuning the mount geometries. Manual focus adjustment of the pupil relay and telescope allow compensation of defocus errors due to humidity or barometric pressure. Plans are to assemble components using a transfer-spindle technique in which individual mechanical cells are elastomerically mounted to their centered optical components and final machining is performed without remounting. Subassembly tilt and decenter tolerances are therefore limited not by any mechanical tolerances by rather by our ability to optically detect tilt and decentration. Final assembly of the interferometer involves extensive self-testing using specific CGHs to bootstrap first the beam expander and viewing arm, then the beamsplitter cubes and reference mirrors, the pupil relay arm, and finally the telescope and other attachments.

We note that the 5-dimensional space of CGH pupil coordinates x and y (test pupil), sagittal and tangential ray slopes at the CGH (asphericity), and pupil relay slide position (test optic radius) includes all possible test beam ray paths through the interferometer. Consequently, the adequacy of the raytrace model in describing the interferometer

optics to within the desired system accuracy of  $\lambda/10$  can be verified by designing CGHs to test a flat mirror at selected tilt angles of up to 2.7 degrees (150 lpmm) and at various pupil relay slide positions. If sufficient accuracy does not result, the raytrace model must be improved. This will involve using the test interferograms to deduce mechanical misalignments and/or to insert dummy phase surfaces at key locations within the interferometer to compensate residual surface irregularities or glass inhomogeneities. By sufficient application of this test and calibration procedure, we can refine the raytrace model to reduce system errors to the limits imposed by interferogram measurability (<  $\lambda/50$ ), system stability, and the capabilites of the commercial raytrace software. We are confident that the interferometer stability and existing raytrace software (e.g. Code V<sup>®</sup>, Super Oslo<sup>®</sup> or GENII<sup>®</sup>) will be adequate to achieve the desired  $\lambda/10$  system accuracy.

### 6. SUMMARY

A laser unequal path interferometer (LUPI) is under development at APA Optics for testing of general aspherics. The interferometer design makes maximum use of inexpensive e-beam written CGHs to null test surface asphericity, to ease manufacturing tolerances by compensating known instrument imperfections, and to serve as alignment and calibration tools. The Interferometer for Aspheric Testing (IAT) is expected to be capable of testing aspheric departures (from base sphere) of several hundred waves with an accuracy of  $\lambda/10$  using only standard auxiliary optics. Deeper aspheres will be testable using simple optics as partial null compensators.

### 7. ACKNOWLEDGEMENT

The authors acknowledge the capable help of Mr. Lubomir Koudelka in accomplishing the opto-mechanical design and addressing the manufacturability of the aspheric interferometer.

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### AN INTERFEROMETER FOR ASPHERIC TESTING: CALIBRATION AND ERROR COMPENSATION

Steven M. Arnold

**PREPRINT** ASPE Spring Topical Meeting April 7-9, 1992, Tucson, AZ

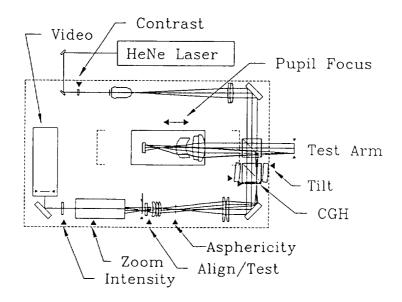
APA Optics, Inc. April 7-9, 19 2950 N.E. 84th Lane, Minneapolis, Minnesota 55434

### EXTENDED ABSTRACT

APA Optics has developed a phase measuring interferometer for testing of general aspherics using electron-beam written computer generated holograms as null compensators. Aspheric departures of up to several hundred waves are measured using only standard accessories; deeper aspheres may be tested using simple auxiliary optics.

The interferometer configuration is depicted in the figure below and has been described more completely elsewhere [1]. It is a Twyman-Green interferometer which uses polarization optics to separate the test and reference paths. The test beam is reflected by a first beamsplitter cube, enters the test arm, and returns through a 150-mm travel afocal telecentric relay system which images the test optic onto a 20-mm diameter computer generated hologram (CGH) located just after the second beamsplitter cube. Meanwhile, the collimated reference beam is transmitted by the first beamsplitter cube and undergoes four reflections by the second beamsplitter cube and the two reference mirrors before striking the CGH at an off-axis angle. The test beam and the diffracted reference beam form a nominally null interferogram which is spatially filtered and reimaged onto a diffuser glass by an afocal telescope. The image on the diffuser is then viewed by a video camera and zoom optics.

Because of 8 reflections and 8 double-pass air-glass interfaces unique to either the test or reference paths, the interferometer is not expected to achieve diffraction-limited performance without error compensation. Our approach has been to incorporate the system error compensation into the design of each CGH null compensator. In this manner, we can compensate known imperfections of the interferometer and accessory optics while presenting the user with a nominally null interferogram.

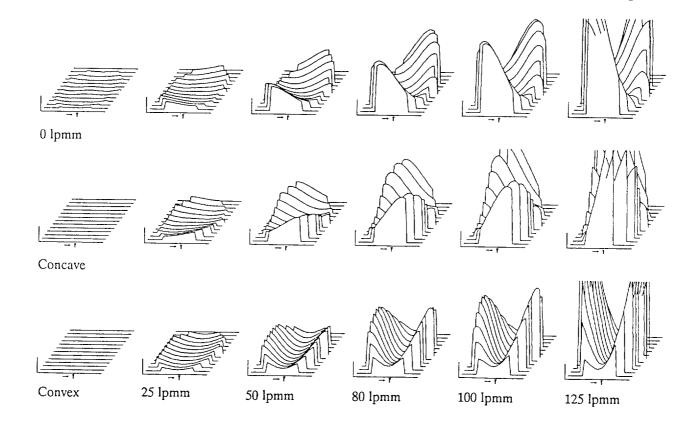


Several measures have been taken to assure that the interferometer alignment remains invariant from source arm pinhole to CGH. The front of the interferometer includes a kinematic bayonet mount for the attachment of various accessories. The removable CGH null compensator is also mounted kinematically. The laser collimator and the pupil relay optics are both designed to remain focused through a range of temperature. The positions of the pupil relay slide and the reference mirror tilt are displayed to the user and are to be set per the CGH design data. The optics following the CGH are common to both test and reference paths and therefore do not influence a null interferogram.

Calibration of the instrument is accomplished by attaching test fixtures to the accessory bayonet mount and inserting corresponding CGHs. The three test fixtures consist of a 500 mm radius concave mirror, a 500 mm radius convex mirror, and a variable tilt plano mirror. We find it useful to express the tilt angle of the plano mirror as an interferogram spatial frequency at the CGH (lpmm =  $2 \lambda^{-1} \sin \theta$ ). Each fixture can be mounted to the bayonet in three orientations differing by 120 degree increments. The plano fixture alone samples the entire range of possible rayslopes over the full CGH aperture. The concave and convex fixtures provide a sensitive measure of centration and some redundancy.

The 13 calibration CGHs are of three types. The concave and convex configurations use off-axis CGHs with a carrier frequency of 300 lpmm. The tilted plano configurations use triple crossed binary gratings comprised of equilateral triangles. These CGHs work for all three bayonet orientations and no off-axis carrier is required since the test beam is itself tilted. A third CGH group consists of 7 linear gratings used to map the reference beam path with the untilted plano fixture in the test beam.

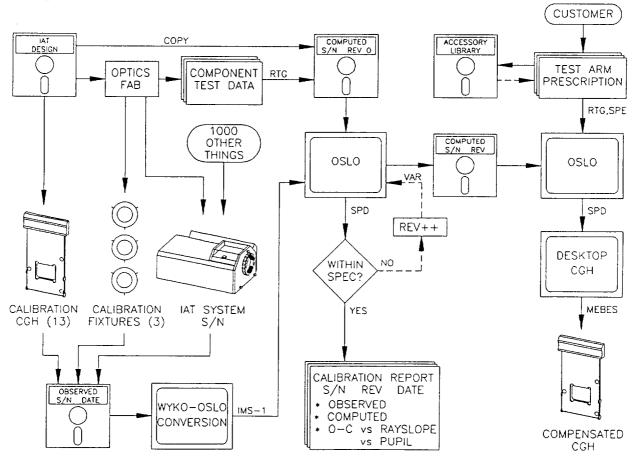
The figure below shows the expected interferograms resulting from the 18 test arm configurations. The leftmost column includes the concave, convex, and untilted plano



configurations. Additional columns represent increasing plano tilts of 25, 50, 80, 100 and 125 lpmm with each row being a different bayonet orientation (0, 120 and 240 degrees). The vertical scale is indicated by a bar representing two waves of OPD. The hardware calibration process consists of recording these 18 + 7 phase measured interferograms along with the system serial number and a date stamp. The process is simple enough to be performed in the field.

The flowchart diagram outlines the complete hardware and software calibration processes. After the 18 + 7 calibration interferograms are recorded to floppy disk, the data is ported to the OSLO raytrace program so that reverse optimization may be used to deduce the imperfections in the actual interferometer.

The WYKO DOS/RTI<sup>®</sup> software package is used to fit the interferogram data with 36 Zernike polynomials and a stand-alone program then converts the Zernike coefficients to an OSLO<sup>®</sup> diffractive surface (type DFX) described by a polynomial in x and y. This multi-configuration diffractive surface, located just after and coincident with the CGH phase function, acts to collimate the observed inteferogram wavefronts. Separate raytrace models are maintained for the test and reference paths. With the Zernike representation, we have effectively  $18 \times 36 = 648$  optimization targets for the test path model and  $7 \times 36 = 252$  optimization targets for the reference path model. Orthogonal Zernike polynomials are well suited to describing the expected decentering, tilting, focusing, and radius errors [2]. A 6-ring and 6-spoke rayset is now used, although we expect to identify a better rayset over which the Zernike polynomials are orthogonal [3]. The optimization operands then consist of OPD maps for each test configuration. Operands are inversely weighted by rayslope.



A question could be raised as to which optimization variables are legitimate. Clearly the 5 discrete tilt settings of the plano calibration fixture are valid optimization variables. Other key variables include the relative orientations of the bayonet, the test arm relay, and the CGH mount with respect to the optical axis (18 variables). Collimator and test relay focus are 2 additional variables. Decenter of the smaller half of the test relay doublet and tilt, decenter and airspace of the meniscus add another 7 variables. We therefore have at least 32 legitimate parameters. Some parameters will act as compensators for others, so the solution obtained may not be unique.

Another modeling approach is to sidestep questions of what constitutes a legitimate parameter and instead use whatever ad-hoc model best reproduces the observed interferogram data with a reasonable number of free parameters. This is what we have done with the reference path which has only a one-dimensional space of accessible configurations (*i.e.* tilt). For the test path, we have thus far opted for a strictly physical raytrace model in order to better understand the hardware. However, we would claim that the test arm space is sufficiently sampled by the 18 test configurations to justify the insertion of one or more fictitious phase surfaces if these are found to significantly reduce the residuals.

Accessory optics such as diverger lenses and an afocal beam expanding telescope are similarly modeled by recording their double-pass interferograms versus field angle. This requires no special CGHs or fixturing and can be accomplished using a conventional Fizeau interferometer.

Two principle outcomes of the interferometer calibration process are: 1) a report detailing the residual system error versus pupil coordinate and versus rayslope, and 2) a raytrace model which can be used in designing CGH null compensators.

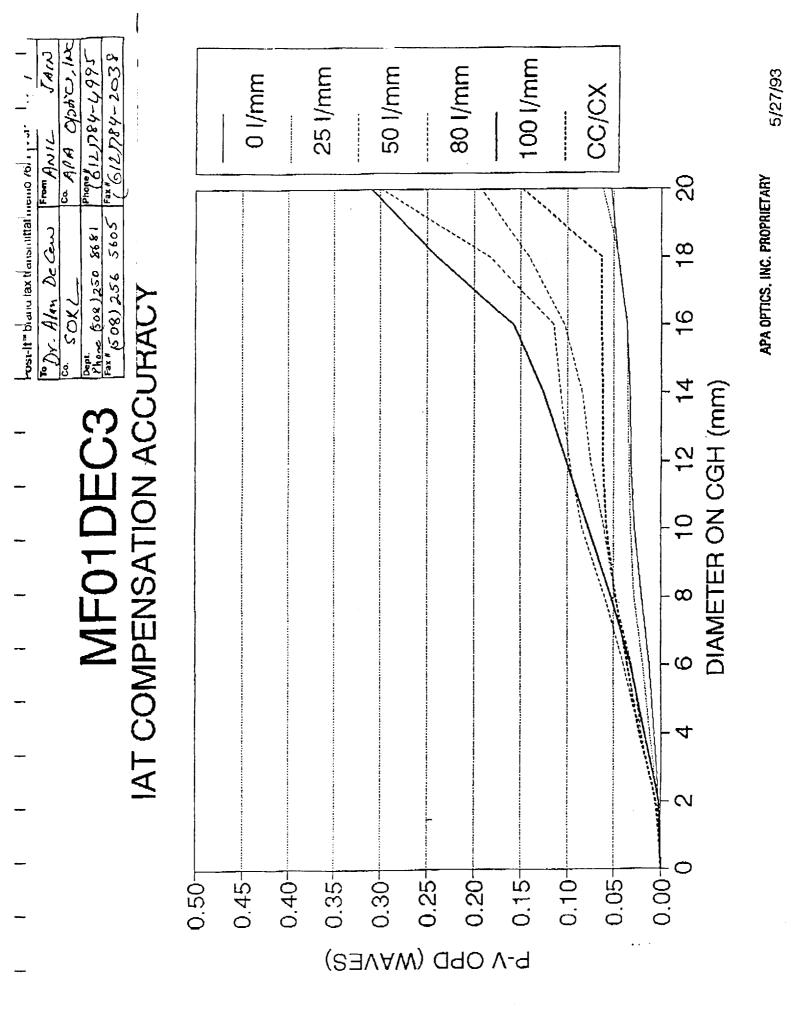
We plan to present initial calibration data from one of our interferometers.

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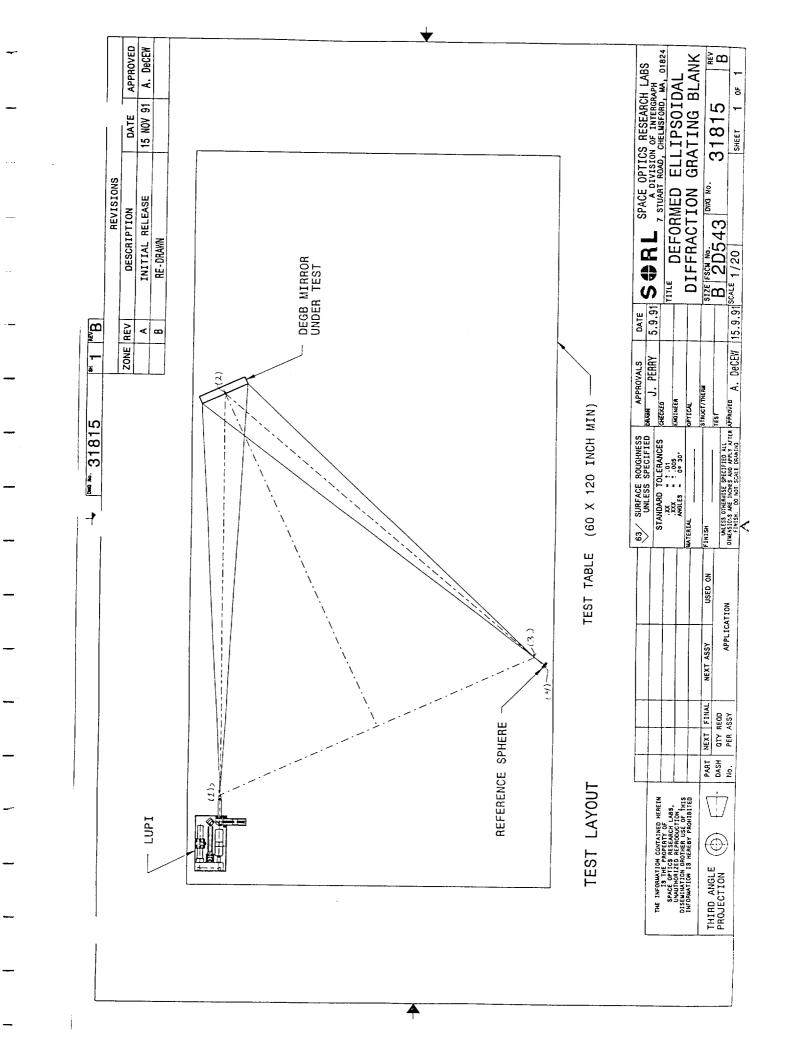
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### APPENDIX E

## OPTICAL TEST SET UP CONFIGURATION



Designated locations in Hologram design program:

{0,0}	{ 0 , 1672.918 }	{ 1260.067 , 572.512 }	{ 1448.371 , 408.069 }
LUPI-II focus	DEGB Vertex	Ellipsoid focus	Ref Sphere Vertex
(1)	(2)	(C)	(4)

Distances:

65.863 in	65.863 in	54.489 in	9.845 in
or	or	or	or
1672.918mm	1672.918mm	1384.030mm	250.063mm
(1) - (2)	(2) - (3)	(1) - (3)	(3) - (4)

### SET-UP PROCEDURE

- Set up components in rough positions. Place a refocusing sphere 0.125"R at the focus of the IAT. Locate center of DEGB with mask. ÷
- readingof radius rod should be 65.488 after taking into account the refocusing Using a radius rod position the center of the DEGB to be 65.863". Actual sphere radius and the diameter of the tangent balls on the radius rod ends. i
- 3. Using same radius rod. set up second ball on the arc of the second focus.
- With second radius rod, relocate second ball on both the arc from the DEGB vertex and the arc from the interferometer focus measuring 54.489". 4
- Extend radius rod from DEGB vertex to 75.708", and position the concave spherical return mirror. Ś.

### APPENDIX F

AN ALTERNATE TEST METHOD

**TECHNICAL REPORT** 

### CGH NULL TEST OF DEFORMED ELLIPSOIDAL GRATING BLANK

4 November 1993

Submitted to:

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Prepared by:

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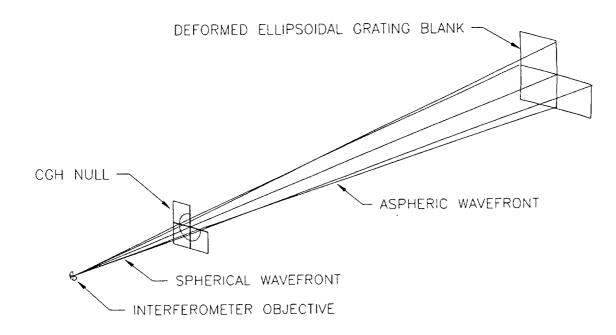
CGH & BINARY OPTICS - ASPHERIC METROLOGY - DESIGN ENGINEERING

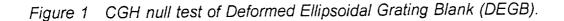
Diffraction International has recently developed a CGH null adapter to allow testing of highly aspheric optics using a conventional Fizeau or Twyman-Green interferometer. Based on this invention, we have designed and analyzed a null test for the Deformed Ellipsoidal Grating Blank (DEGB).

### Null Test Configuration

The test configuration is shown in Figure 1 below. The DEGB is tested from its center of curvature using a conventional interferometer such as SORL's LUPI-II. A CGH null positioned in the test beam between the focus and the DEGB transforms the spherical test beam into a highly astigmatic aspheric wavefront which is everywhere perpendicular to the DEGB surface. Peripheral portions of the CGH aperture, not required for the DEGB null, are devoted to a return sphere grating for aligning the CGH with respect to the interferometer. The proposed CGH is depicted schematically in Figure 2.

The CGH null must be positioned outside the caustic so that there will be a one-to-one correspondence between points on the CGH and points on the DEGB. Since CGH alignment tolerances, required resolution and aperture each scale with distance from test beam focus, the CGH should be located as far as practical from the focal position. For fabrication by e-beam lithography, the CGH aperture must be less than 150 mm. I have selected a CGH location 400 mm from the test beam focus as a reasonable compromise between performance and cost. Furthermore, to balance the x-axis and y-axis resolution requirements of the CGH, I have located the DEGB vertex 1670 mm from the test beam focus. The required CGH aperture for the DEGB null is then 62 x 28 mm.





-1-

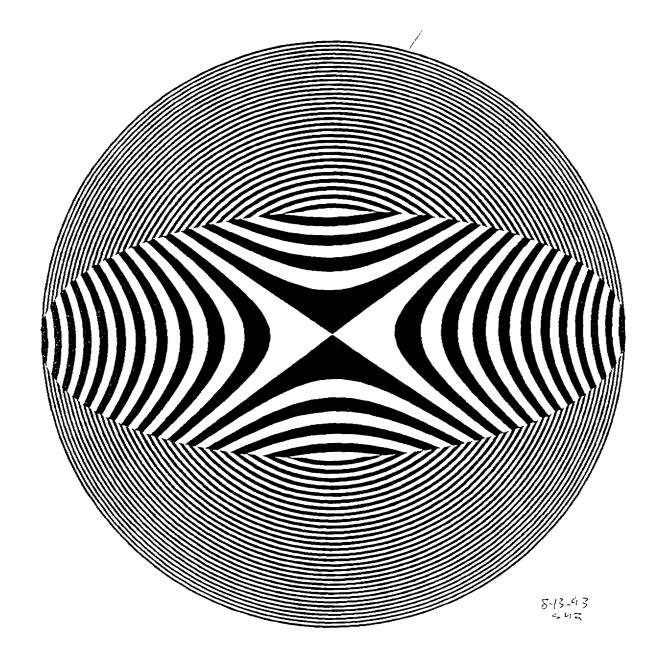


Figure 2 Proposed CGH including astigmatic DEGB null (center) and axisymmetric return sphere alignment grating. Optional fiducial projecting alignment gratings are not shown. The fringe spacing has been increased for purposes of illustration. Actual CGH diameter is 62 mm.

### Alignment

Since the DEGB lacks rotational symmetry, the proposed test requires 9 alignment degrees of freedom. Alignment of either the CGH null or the DEGB to the spherical test beam establishes an axis and consumes 3 degrees of freedom-1 of focus and 2 of tilt or decenter. Alignment of the remaining component (CGH or DEGB) then requires a full 6 degrees of freedom.

The alignment procedure is somewhat more intuitive if the test optic (DEGB) has 6 degrees of freedom, and the CGH null only 3. This is appropriate for an acceptance test. During the figuring process, however, the CGH will be aligned only once while the DEGB will be repeatedly removed and replaced. In this case, it is easier to use a 6-axis mount for the CGH and a simple tilt-tip-focus mount for the test optic. Mechanical repositioning of the test optic can be sufficiently repeatable that no decenter or clock adjustments are needed.

Locating the CGH null outside the region of the caustic results in a quite elliptical working aperture. This is perhaps fortunate as it leaves a large portion of the physical aperture available for alignment features. In particular, a reflection CGH is proposed for purposes of aligning the CGH to the spherical test beam. The reflection CGH will act like a 400 mm radius concave sphere and can be optically aligned by nulling its interferogram. It is also possible, as an aid in aligning the decenter and clock degrees of freedom, to incorporate various transmission CGHs which focus otherwise unused portions of the test beam onto two or more fiducial locations on the DEGB.

### **Optical Prescription**

The optical prescriptions of the DEGB surface and the CGH null have been approximated to a high degree of accuracy by cartesian polynomials. A raytrace model was constructed using Oslo Series 2 optical design software. The CGH wavefront coefficients were optimized to produce a null test to within 0.0002 fringe P-V and 0.00004 fringe RMS. The raytrace data is attached as Appendix A.

I do not know what tolerances have been assigned to the DEGB vertex radius and eccentricity. In testing from center of curvature with a null, the test wavefront is a good (deformed) ellipsoid only for the prescribed axial distance (vertex radius). This is not a problem for fabrication testing since the DEGB would be figured to match the null wavefront. However, for acceptance testing of an already completed DEGB, the vertex radius and eccentricity must be known with sufficient accuracy to assure a near null test.

### **Tolerance Analysis**

Error sources include CGH fabrication errors, interferometer errors, and alignment errors. A distinction is made between manufacturing tolerances and measurement tolerances. Errors which can be well measured and characterized may be removed during the test results analysis or compensated in the CGH design. Nonethless, it is best to minimize dependence on data analysis for error compensation, as the optician's job is much easier if the observed interferogram accurately represents only test optic errors.

The various error contributions are listed in the following table and discussed in the following paragraphs.

Parameter	Tolerance	Error (waves)	Compensation Method	Compensated Error (waves)
Interferometer	0.05-0.10 wave	0.100	Characterize	0.020
CGH Encoding	±0.01 wave	0.020	None	0.020
CGH Digitization	±0.125 micron	0.009	None	0.009
CGH Registration	0.15 micron	0.005	None	0.005
CGH Substrate	0.1 wave	0.100	Characterize	0.020
CGH Focus	25 micron	0.048	Phase Measure	0.005
CGH Tilt	l µrad	0.000	None	0.000
DEGB Focus	0.6 mm	2.750	System Alignment	0.007
Total (worst case)		3.032		0.086
Total (quadrature)		2.754		0.037
DEGB Error Budget		N/A		0.125

## Interferometer Errors

Interferometer system error depends on the focusing optics and test aperture and are typically 0.05 to 0.1 wave. System error is characterized and compensated by the standard technique of testing a return sphere in two clock orientations plus the cats-eye configuration. Compensation accuracy depends on the the test environment, but is typically better than 0.02 wave.

## **CGH Fabrication Errors**

CGH errors are comprised of pattern errors (encoding, digitization, registration) and substrate errors. Appendix B discusses these error sources in general terms.

The minimum fringe spacing of the proposed DEGB null is 28.3 microns, so pattern errors will be quite small in terms of phase. It is proposed that the CGH be encoded with an algorithm precision of  $\pm 0.01$  wave and a data grid (pixel size) of 0.25 microns. The guaranteed mean plus  $3\sigma$  registration accuracy of the MEBES e-beam system is 0.15 microns. Together, these error sources yield a CGH pattern error of 0.04 wave P-V. This could be reduced to 0.02 wave or less by tightening the algorithm precision and using a smaller pixel size.

Diffraction International offers photomask substrates with transmitted wavefront distortion of better than  $\lambda/10$  P-V. This error is measured and characterized interferometrically and may be verified by the end user.

## **CGH Alignment Errors**

By incorporating a 400 mm radius reflection CGH, the CGH null can be aligned with respect to the spherical test wavefront to better than 25 micron focus and 1  $\mu$ rad tilt (equivalently 1  $\mu$ m decenter). These values each correspond to 1/4 fringe in the alignment interferogram. A 25 micron CGH focus error contributes a significant 0.048 wave to the DEGB interferogram. In a good test environment, phase measuring interferometry can reduce this error by a factor of 10. The CGH tilt error of 1  $\mu$ rad is entirely negligible, contributing only 0.0001 fringe to the DEGB interferogram.

## **Test Optic Alignment Errors**

Assuming a null test is achieved, focus is the only DEGB degree of freedom which contributes to a surface figure error (*i.e.* vertex radius). The wavefront shape produced by the CGH null and the axial distance of the DEGB from the CGH completely define the DEGB surface figure. If a portion of the CGH periperal aperture is allotted to a transmission CGH which focuses onto the DEGB vertex, then the DEGB vertex radius can be optically determined to an accuracy of about 0.6 mm (1/2 fringe). This corresponds to a P-V figure error of 2.75 waves, but is it significant? Since it is almost entirely power and astigmatism, this figure error can be mostly compensated during final system alignment. To determine what fraction of the figure error caused by an incorrect CGH to DEGB spacing is but a minor change in vertex radius and eccentricity and what fraction is higher order aberrations, the  $x^2$  and  $y^2$  terms of the DEGB polynomial surface description were varied. It was found that only 0.007 wave is higher order aberrations.

The other 5 degrees of freedom are required only to establish that the backside of the blank is perpendicular to the optic axis to within  $\pm 15$  arc seconds and that the alignment fiducials coincide with the ellipsoid major and minor axes. Depending on how the available CGH aperture is partitioned, the size of fiducial spots focused onto the DEGB will be 20-50 microns. This should be sufficient to meet the DEGB mechanical centration specifications.

## **Pupil Distortion and Diffraction Effects**

There is significant pupil distortion in mapping the  $200 \times 160$  mm DEGB aperture onto the  $62 \times 28$  mm CGH aperture. Pupil distortion is of no consequence for a perfect null test, but must be considered when interpreting interferograms which depart from null.

Because the DEGB is not imaged onto the CGH null, there will be diffraction effects from the edges of the CGH aperture(s). For this reason, the proposed aperture of the DEGB null is made somewhat oversized so that the DEGB will act as the limiting aperture.

## Implementation

The proposed CGH null would be fabricated by e-beam lithography as a chrome pattern on a standard size 4-inch photomask blank. Figure 3 shows the expected fringe visibility for both chrome and phase CGH nulls when used in a Fizeau configuration. Visibility with a chrome CGH is adequate for visual inspection. For maximum accuracy with phase measuring Fizeau interferometry, the photomask should be ion etched to produce a phase-only CGH. It would be necessary to mask the alignment reflection CGH from the chrome etch. Alternatively, fringe visibility can be increased by reducing the reflectivity of the reference mirror of SORL's LUPI-II interferometer.

Diffraction International has recently developed a kinematic 3-axis mount for 4-inch photomask CGHs. The only additional hardware required to perform the proposed null test is a 6-axis mount for the DEGB.

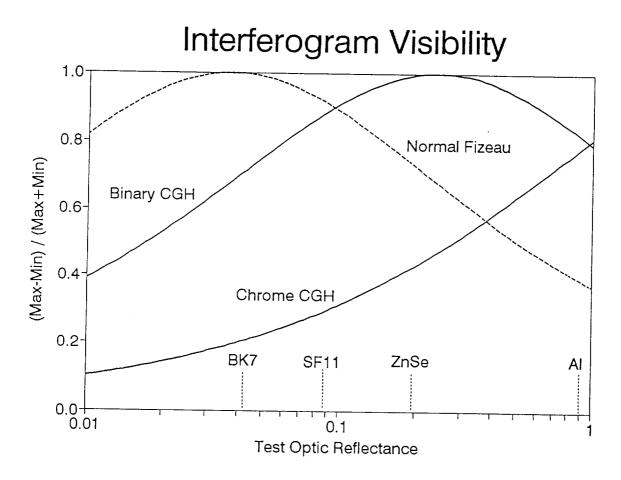


Figure 3 Fizeau interferogram visibility with chrome and phase CGH nulls.

## Conclusion

When combined with phase measuring interferometry to establish the correct interferometer to CGH spacing, the proposed CGH null test will measure the DEGB figure error to an expected accuracy of better than  $\lambda/20$  and the DEGB power to about  $3\lambda$ . The accuracy of the observed interferogram, before software compensation to correct interferometer system and CGH substrate errors, is about 0.15 waves.

\_LEN NEW "SORL DEGB CGH NULL TEST" AST 2 RFS 11 EPR 0.4 -OBY -0.001 THO -10 DES "sma" \_IDN 44 WW 1 0 0 0 0 0 WV 0.63282 0.48613 0.65627 0.43584 0.70652 0.40466 GTO 1 GTO 2 GLA BK7 RD 10 -AP 1 NOT "Micro-Objective" GTO 3 \_PK CV 2 TH 410 PK AP 2 NOT "Micro-Objective" GTO 4 NOT "Design CGH Face" GTO 5 -GLA SILICA AP 31 NOT "Actual CGH Face" ЗТО 6 GLA SILICA TH 2.286 ?Κ AP 5 -PDT PDR 2 SDT DFX 10 33 -3.700544119884E-04 \_35 7.200710687056E-04 S8 1.313929814807E-06 310 1.725699643609E-09 312 1.561993743483E-09 S14 -5.379144915847E-09 S17 -7.523552078414E-12 319 -3.387904665471E-11 -021 -6.397123329599E-15 S23 -7.42107296602E-15 325 4.745158907399E-14 327 3.616767921395E-14 NOT "CGH Phase Function" GTO 7 'H −2.286 TPK THM 6 SDT SAP 'LL 5 AY2 50 .X1 ~50 .X2 50 NOT "CGH Backside" GTO 8 K THM 4 TRE SDM 5 NOT "Restore Design CGH"

\_\_GTO 9  $\mathbf{TH}$ 1270 NOT "Nominal Test Optic " GTO 10 NOT "Test Optic Align" GTO 11 RFL — PK THM 10 PDT PDX 6 P3 -2.72108859282E-04 P5 -3.282819711021E-04 P8 -1.1283E-08 P10 -2.430705791785E-11 P12 -5.861123049922E-11 ---P14 -3.537863699515E-11 P21 -4.342622773699E-18 P23 -2.041886399956E-17 \_P25 -2.766986059897E-17 P27 -7.625444379005E-18 SDT SAP APT 2 AY1 -80 AY2 80 AX1 -100 -AX2 100 NOT "DEGB" GTO 12 -1270 \_\_TH PK THM 9 PKP SDM 10 NOT "Restore Design Axis" -GTO 13 PK TH 4 NOT "Design CGH Face" \_\_\_\_\_TO 14  $\mathbf{TH}$ 2.286 ?K TH 6 ?K AP 5 PKP SD 5 NOT "Actual CGH Face" JTO 15 -PKP GLA 5 TH -2.286 PK THM 6 \_YK AP 7 NOT "CGH Backside" GTO 16 үк тнм 13 ТРК АР 6 PKP SD 6 IOT "CGH Phase Function" -- TO 17  $\mathbf{TH}$ -400 'K THM 3 10 'KP SDM 14 NOT "Restore Design CGH" **TO** 18 LA BK7 'AP 10 NOT "Nominal Focus"

-GTO 19 PK AP 18 END 19 JAR NEW TLA 10 0 -1 1 1 1E-05 TLB 10 0 -1 1 1 1E-05 FLC 10 0 -1 1 1 0.001 -DCX 10 0 -0.5 0.5 1 0.001 0 -0.5 DCY 10 0.5 1 0.001 0 -5 5 1 0.01 ГН 10 END OPI F1 0 0 31 0 0 1 **R**2 0 0 0.5 R3 0 0 -0.5 R4 0 0 1 0 -1 0 R6 0 0 1.5 37 0 0 -1.5 \_R8 0 2 0 R9 0 0 -2 0.25 **R10** 0 0 0.25 R11 0 0.5 **-**R12 0 0.25 -0.5 0.25 R13 0 1 R14 0 0.25 - 11.5 \_\_R15 0 0.25 R16 0.25 -1.5 0 R17 0.25 2 0 R18 0 0.25 - 2R19 0 0.5 0 0.5 0.5 R20 0 R21 0.5 -0.5 0 0.5 1 -R22 0 0.5 - 1R23 0 R24 0 0.5 1.5 \_R25 0 0.5 -1.5 0.5 2 R26 0 0.5 - 2R27 0 0.75 0 R28 0 <sup>--</sup>R29 0 0.75 0.5 R30 0 0.75 -0.5 R31 0.75 1 0 \_R32 0 0.75 - 1R33 0.75 1.5 0 R34 0 0.75 -1.5 R35 0 0.75 2 R36 0.75 - 20 1 0 R37 0 R38 1 0.5 0 1 -0.5 **-**R39 0 R40 1 1 0 R41 1 -1 0 1 1.5 \_R42 0 R43 0 1 -1.5 1 2 R44 0 R45 0 1 -2 80 0 R46 1 R47 1 -80 0

----



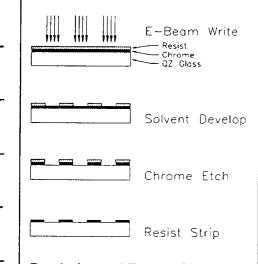
## CGH Photomasks

CGH & BINARY OPTICS . ASPHERIC METROLOGY . DESIGN ENGINEERING

At Diffraction International, we specialize in diffractive optics. Our staff has over 15 years experience in CGH and binary optics technology. Our goal is to become a leading supplier of custom precision diffractive optics. As a small business, we offer flexibility, attention and fast response.

## **General Description**

CGH Photomasks are the basis of our diffractive optic products and are also used directly as diffractive optical elements. CGH Photomasks are encoded using our HoloMask software and written using a MEBES e-beam lithography system. The MEBES is the standard photomask tool in the microelectronics industry. In resolution, accuracy and format size, it is unsurpassed for CGH writing. Unfortunately, the MEBES pattern vocabulary is limited to trapezoids. Our HoloMask software accurately encodes a diffractive optic element using a minimum number of these trapezoids.



## Resolution and Feature Size Patterns are written in a raster manner. The range of address resolution, or pixel size, is 0.10 to 1.1 micron. Common address sizes are 0.50, 0.25 and 0.10 micron.



 $\varphi(x,y) = \frac{Cr^2}{1 + \sqrt{1 - (\kappa + 1)C^2r^2}} + \sum_i A_i r^i + \sum_{ij} P_{ij} x^i y^j + \sum_i Z_i U_n^i(r,\theta) + custom$ 

Because of exposure proximity effects, the minimum feature size should be not less than 5 pixels. This is seldom a limitation for CGH patterns since a minimum of 10 pixels per grating period (5 pixels per line or space) is required to achieve a typical recording accuracy of 0.1 fringe. Feature sizes below 0.7 microns are extremely difficult to control in processing.

## Aperture

Within a rectangular data field, apertures of any shape may be encoded. Patterns may not exceed the MEBES physical reticle limit of 155×155 mm.

## **Optical Prescription**

The optical phase function is specified by an equation. Virtually any functional form  $\varphi(x,y)$  can be accomodated. Simple zone plate lenses may be specified by giving the focal length and wavelength. More general phase functions are described by radial or cartesian monomials or by Zernike polynomials. We will accept Sweatt model wavefront prescriptions. We can also write a custom wavefront evaluation subroutine or link to your code.

Most often, phase function coefficients are generated by an optical design program. Our preferred tool is Oslo Series 2, but we are also familiar with Code V and GENII-PC. If your phase function requires a large number of coefficients, we ask that you supply it on floppy disk.

## Units and Dimensions

We usually design in millimeters but can accomodate any unit of measure. Wavelength is in microns or nanometers. The MEBES works in microns.

## **Artwork Generation**

In addition to the phase function, other image elements such as fiducials, title blocks, etc. may be included. A conversion program for customer generated AutoCAD DXF files is available, or we can produce CAD layouts from your drawings. Points and zero width lines are ignored. Text is converted to a single monospace font.

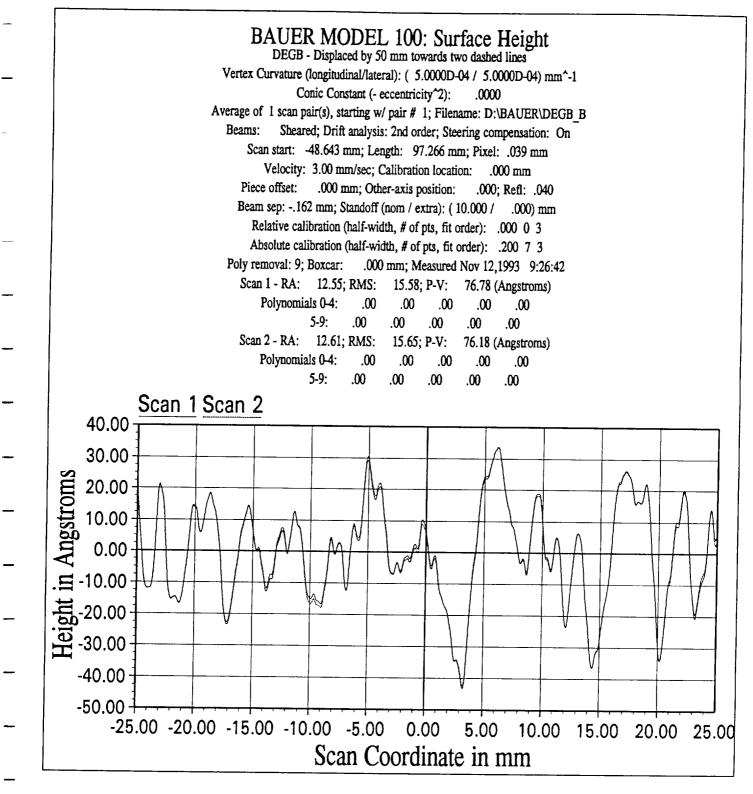
## **Pattern Composition**

Multiple images may be combined to form an overall pattern. Images may overlap and can be either positive (digitized area exposed/clear) or negative (digitized area unexposed/dark). It is difficult to add dark features within a positive pattern or clear features within a negative pattern. The MEBES writes most efficiently when patterns repeat in a rectangular array. This is often true of lenslet arrays or linear gratings.

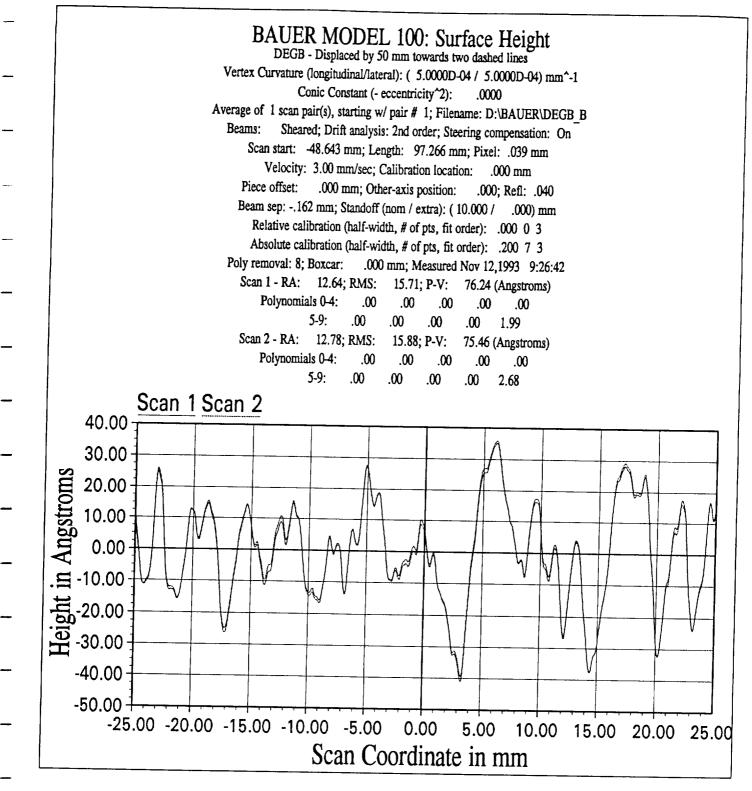
# APPENDIX G

# SURFACE SCAN DATA FOR MIDFREQUENCY

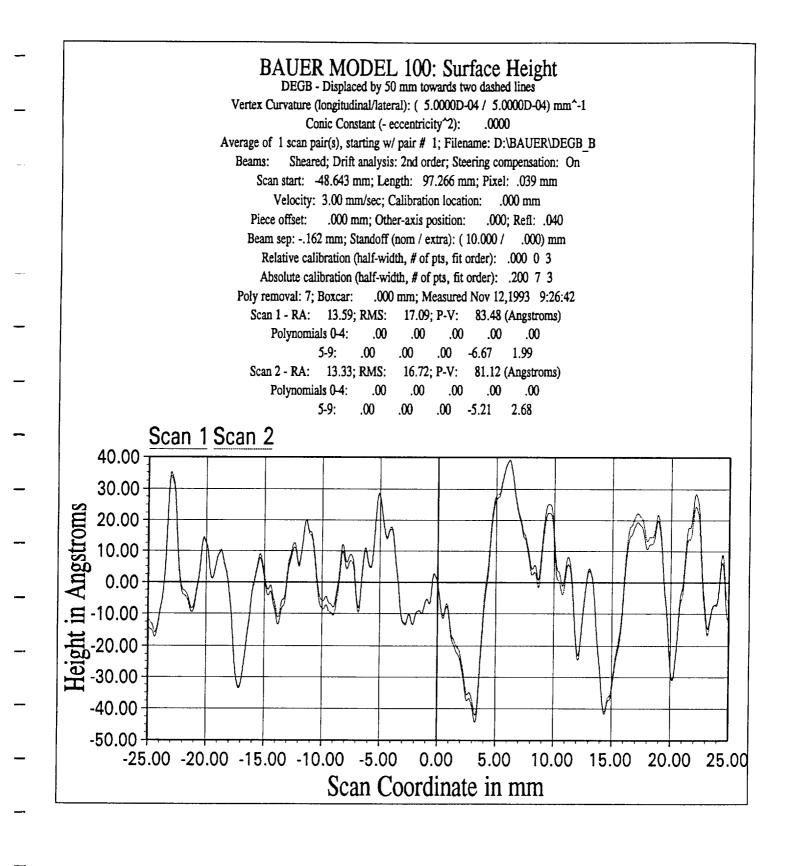
- 1.) Scan data from BAUER Model 100 Noncontact profilometer
- 2.) Legendre polynomials

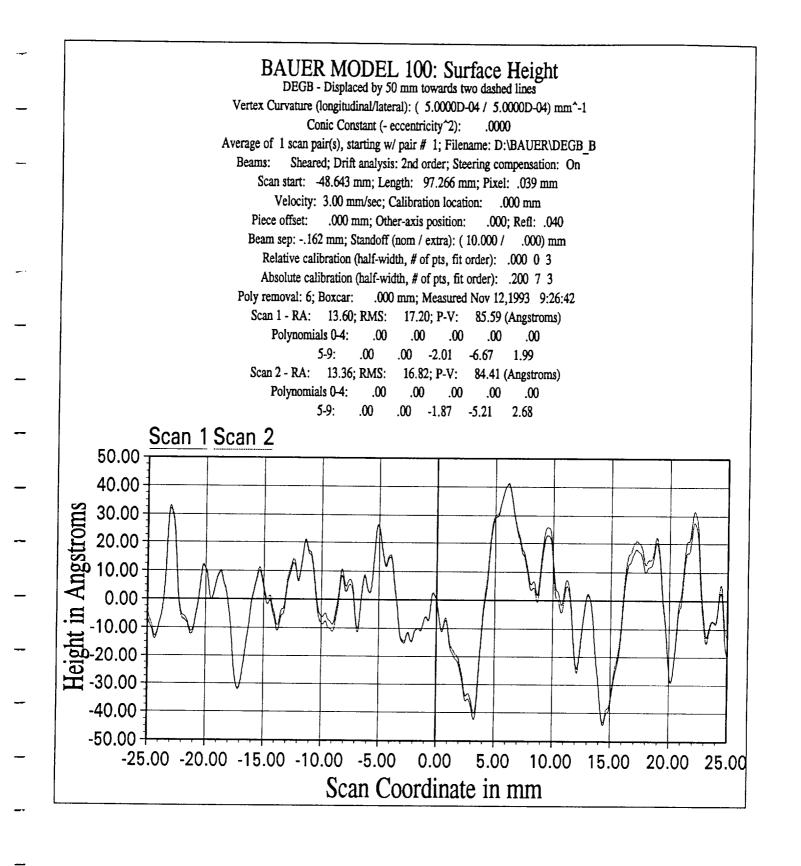


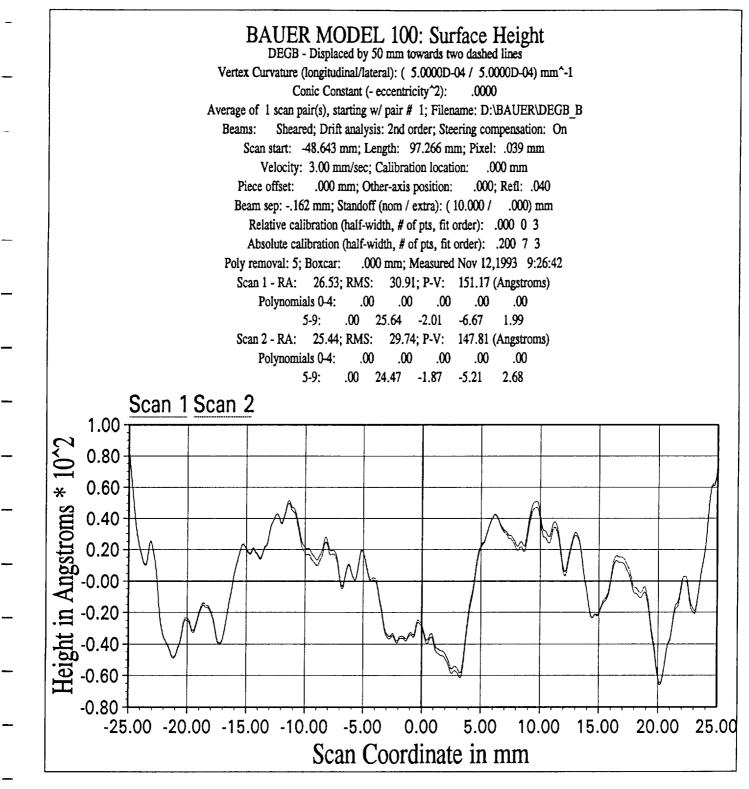
-. .



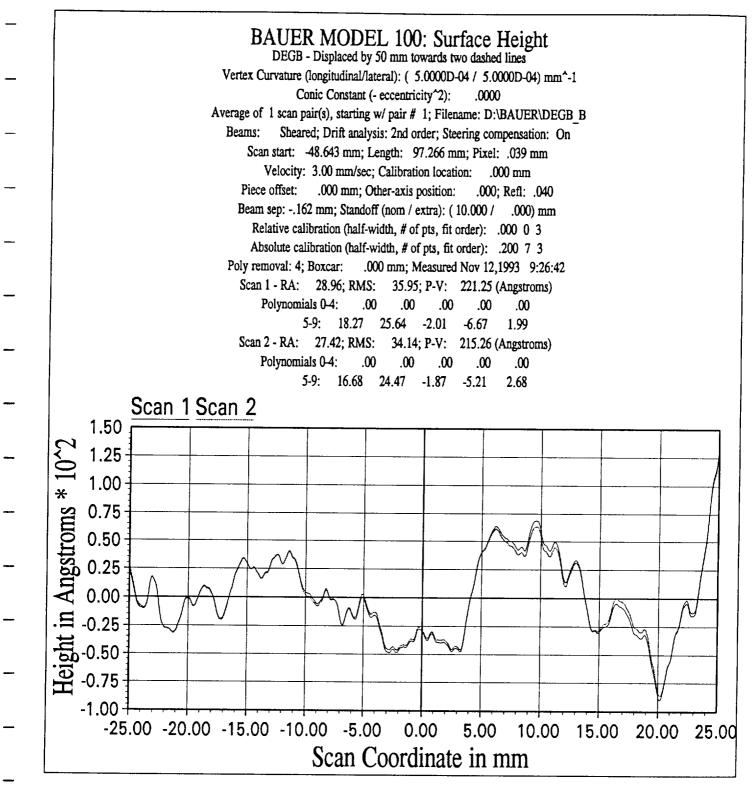
- -
- -

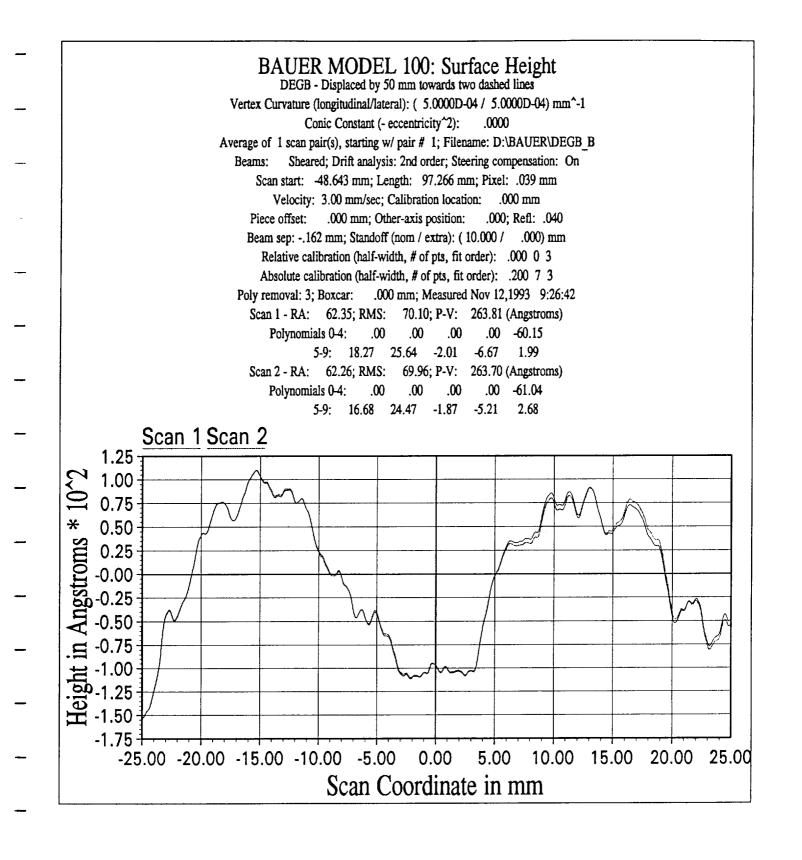


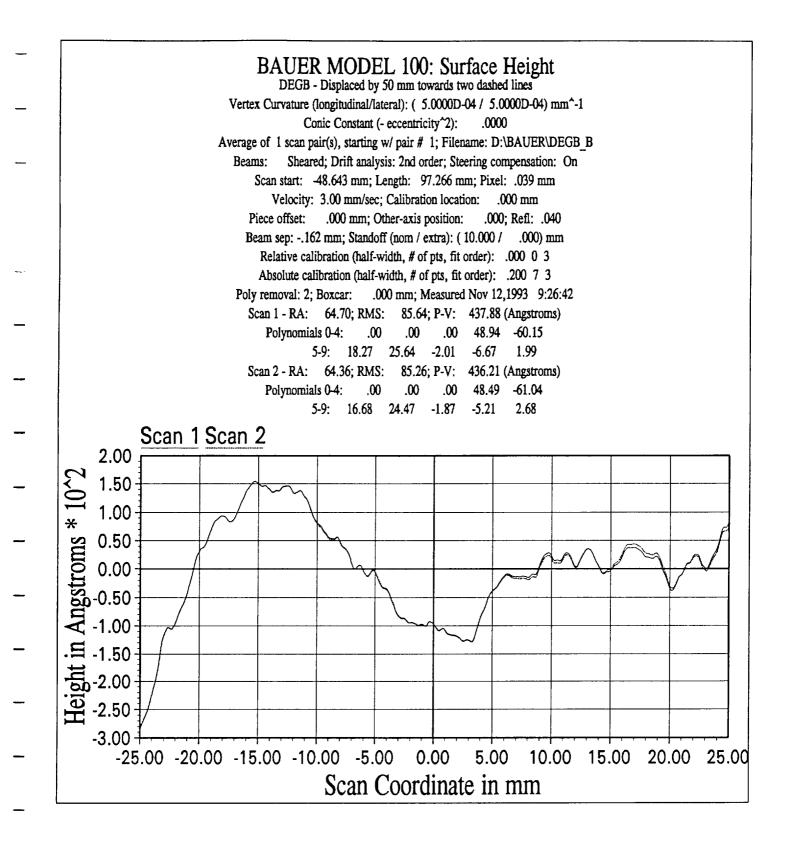


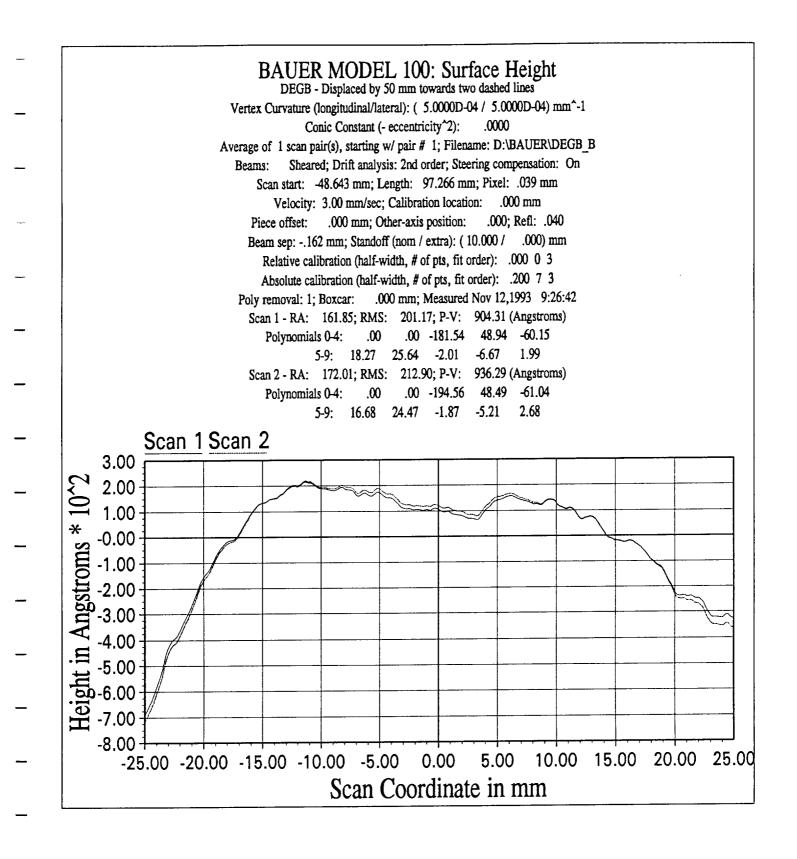


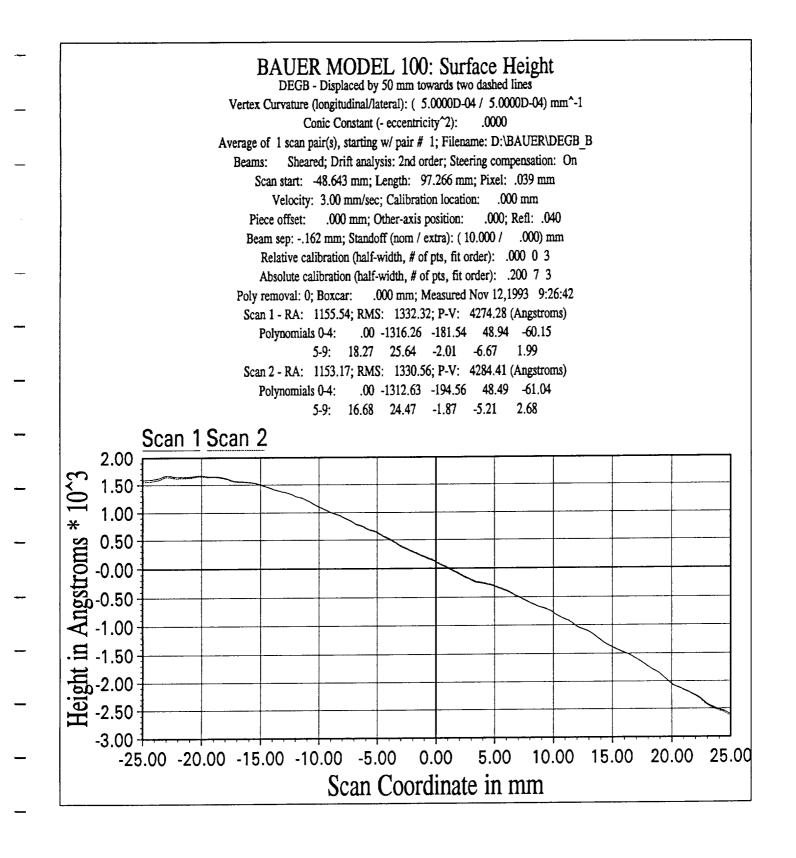
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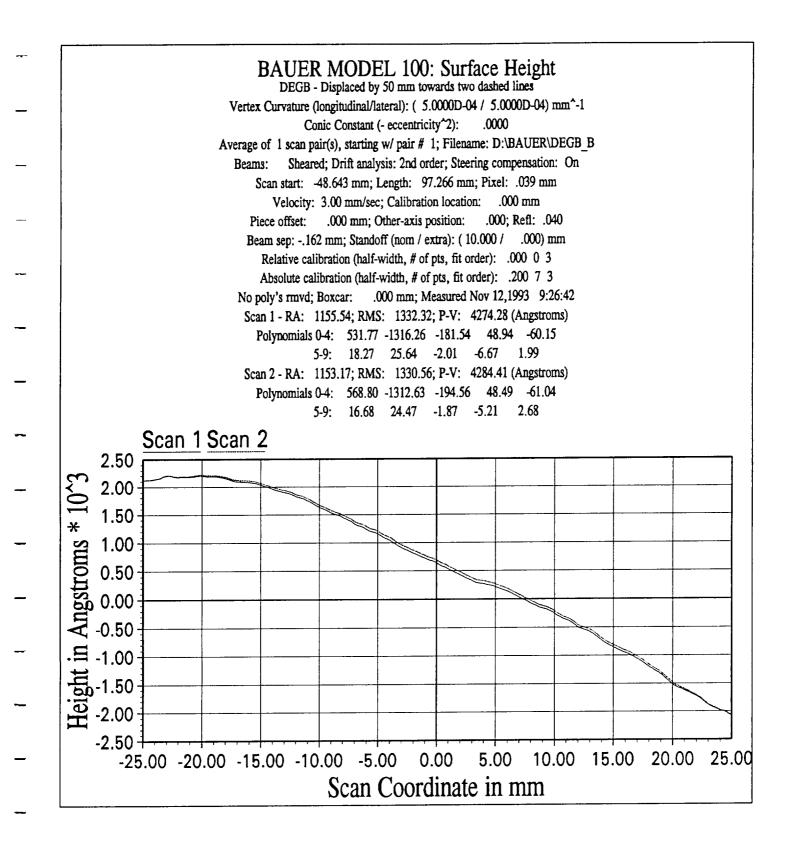








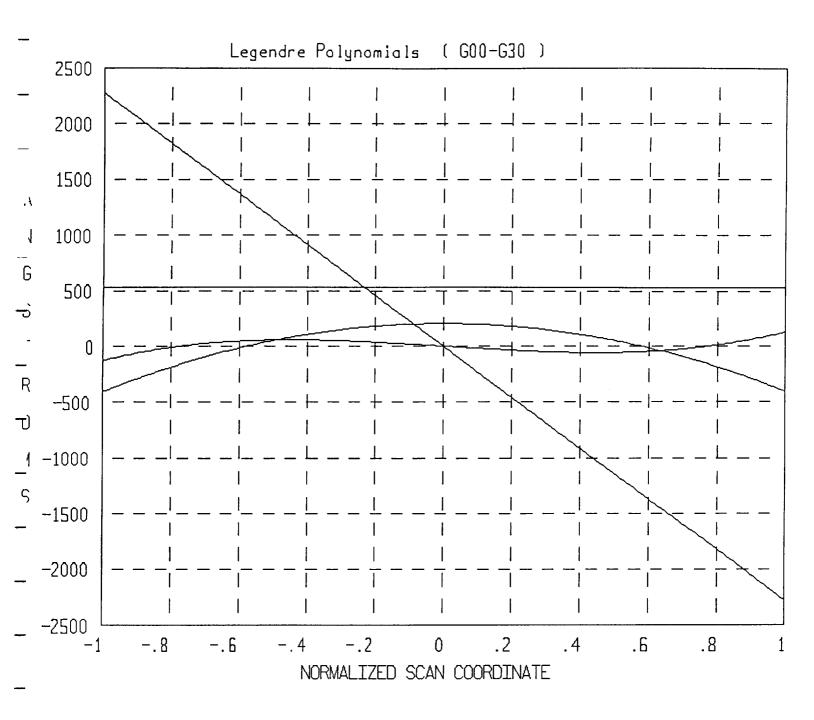




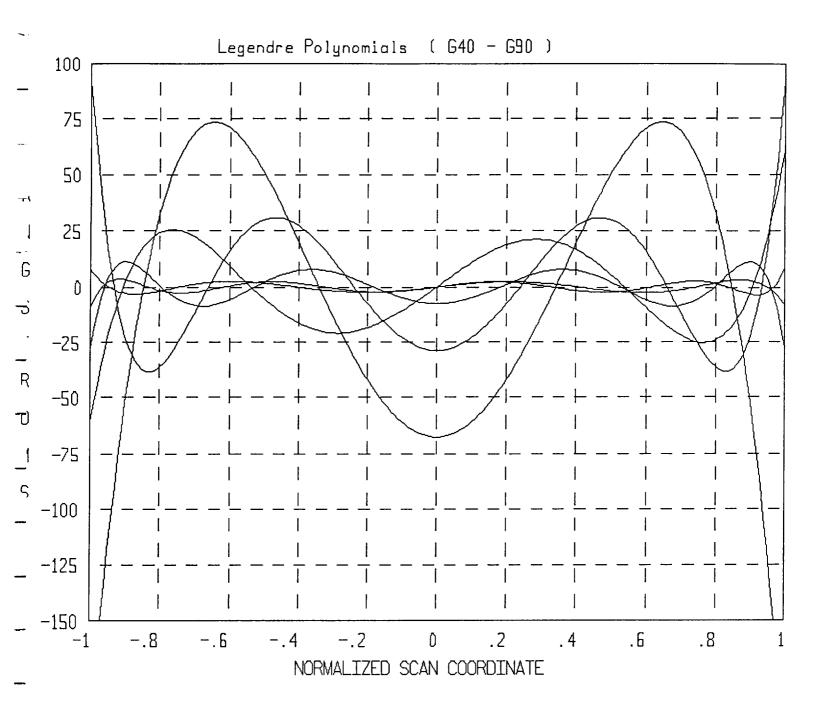
	===========		VARTABLE SH	EET =====	
	Input	Name	Output U	nit	Comment
	0	Z			
-	531.77	a0			
	-1316.26	a1			
	-181.54	a2			
_	48.94	a3			
	-60.15	a4			
	18.27	a5			
	25.64	a6			
	-2.01	a7			
	-6.67	a8			
	1.99	a9			
		<b>a</b>	F01 77		
Ц ·		G00	531.77		
ر -		G10	0		
		G20	202.96789		
ىل _		G30	0 -67.66875		
ر -		G40			
_ب		G50 G60	0 -28.88948		
T.		G70	0		
-		G80	-7.519836		
<u>, , , , , , , , , , , , , , , , , , , </u>		G90	0		
Ľ		0,0	•		
5		G9	630.65982		
· .		G8	630.65982		
L		G7	638.17966		
ייים ביי		G6	638.17966		
Ľ.		G5	667.06914		
ΞĹ		G4	667.06914		
L		G3	734.73789		
L L		G2	734.73789		
		G1	531.77		
$\mathbf{L}$		GO	531.77		
		G09	630.65982		
Ľ		G19	98.889825		
Ъ.		G29	98.889825		
- -		G39	-104.0781		
- <u>1</u>		G49	-104.0781		
Ŀ		G59	-36.40932		
		G69	-36.40932		
 L		G79	-7.519836		
Ц т		G89	-7.519836		
<b>*</b> _		G99	0		
_					

```
S Rule-----
  G00 = a0*1
  G10 = a1*SQRT(3)*z
  G20 = a2*SQRT(5)*(3*z^2 - 1)/2
  G30 = a3*SQRT(7)*(5*z^3 - 3*z)/2
 G40 = a4*SQRT(9)*(36*z^4 - 30*z^2 + 3)/8
  G50 = a5*SQRT(11)*(63*z^5 - 70*z^3 + 15*z)/8
  G60 = a6*SQRT(13)*(231*z^6 - 315*z^4 + 105*z^2 - 5)/16
  G70 = a7*SQRT(15)*(429*z^7 - 693*z^5 + 315*z^3 - 35*z)/16
  G80 = a8 \times SQRT(17) \times (6435 \times z^{8} - 12012 \times z^{6} + 6930 \times z^{4} - 1260 \times z^{2} + 35) / 128
  G90 = a9*SQRT(19)*(12155*z^9 - 25740*z^7 + 18018*z^5 - 4620*z^3 + 315*z)/128
- G9 = G00+G10+G20+G30+G40+G50+G60+G70+G80+G90
  G8 = G00+G10+G20+G30+G40+G50+G60+G70+G80
  G7 = G00+G10+G20+G30+G40+G50+G60+G70
  G6 = G00+G10+G20+G30+G40+G50+G60
  G5 = G00+G10+G20+G30+G40+G50
  G4 = G00+G10+G20+G30+G40
  G3 = G00+G10+G20+G30
 G2 = G00+G10+G20
  G1 = G00+G10
  G0 = G00
  G09 = G00+G10+G20+G30+G40+G50+G60+G70+G80+G90
  G19 = G10+G20+G30+G40+G50+G60+G70+G80+G90
  G29 = G20+G30+G40+G50+G60+G70+G80+G90
  G39 = G30+G40+G50+G60+G70+G80+G90
  G49 = G40+G50+G60+G70+G80+G90
  G59 = G50+G60+G70+G80+G90
- G69 = G60+G70+G80+G90
  G79 = G70 + G80 + G90
  G89 = G80+G90
```

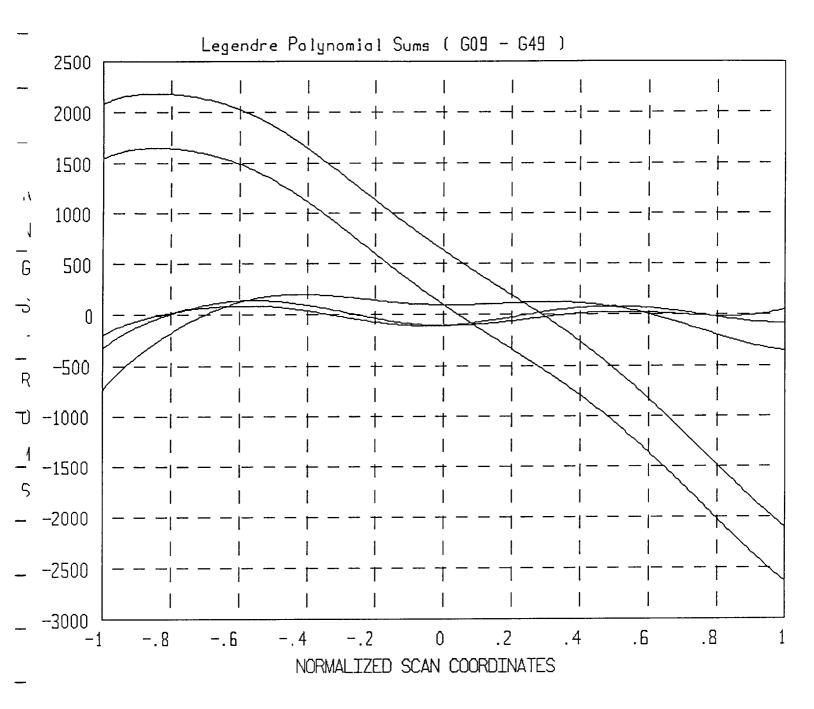
- G99 = G90

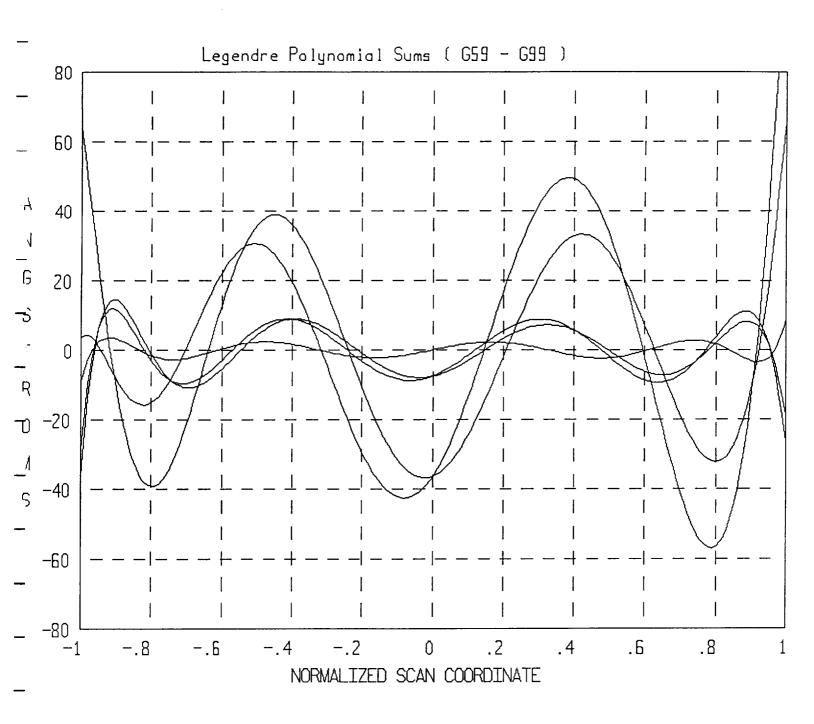


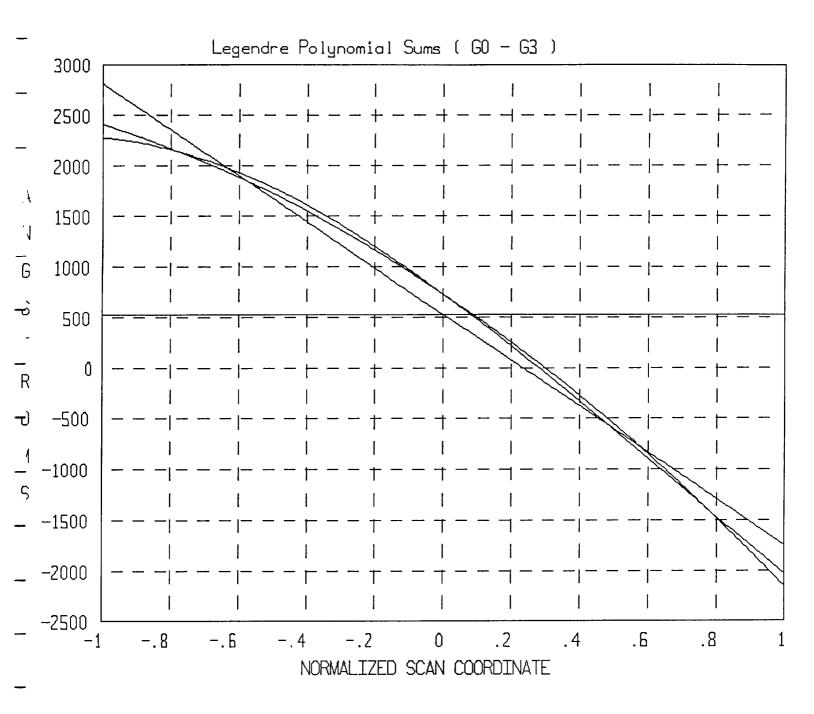
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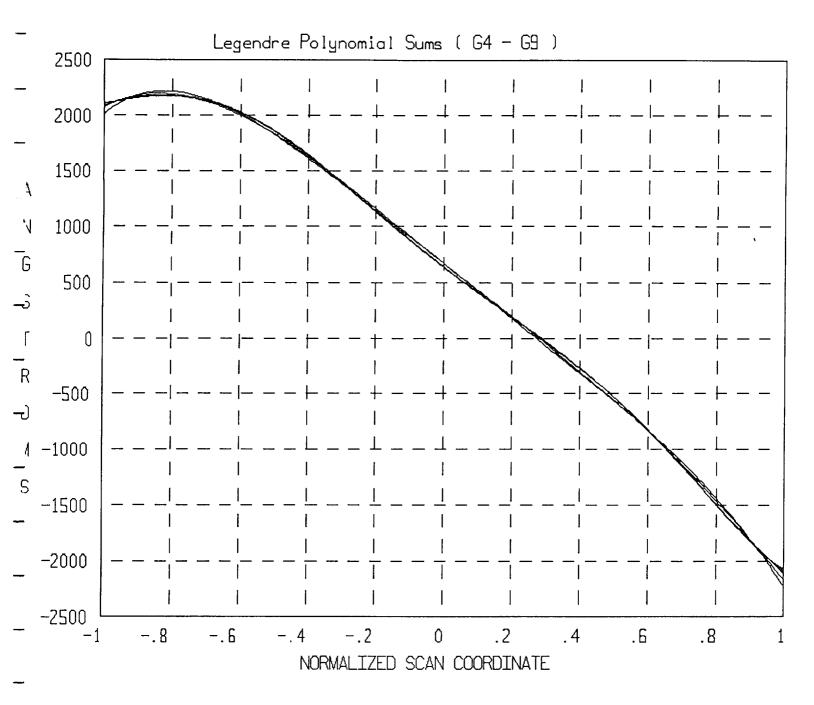
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# APPENDIX H

CONTRACT STATEMENT OF WORK

Attachment A

NAS5-31790 10/29/90 Page 1 of 7

## STATEMENT OF WORK

FABRICATION AND TESTING OF A DEFORMED ELLIPSOIDAL DIFFRACTION GRATING BLANK

1.0 Introduction

The grating blank to be fabricated is intended for use, when ruled as a high line density diffraction grating, as the sole optical element in a high spectral resolving power far ultraviolet (FUV) spectrograph. The ruling of this blank is not part of the current contract. The blank is a non-symmetric asphere (as specified below), designed to minimize optical aberrations when used in first diffracted order in a specific use geometry.

- The performance goal of a resolving power of 30000 when used as an - FUV diffraction grating leads to tight figuring requirements. The lack of a simple null lens interferometric test of the blank makes the manufacturing process challenging.
- This grating blank is intended as a technology demonstration item, and is not currently being considered for any flight mission. The successful demonstration of figuring and ruling such a grating would, however, be of strong interest to planned NASA ultraviolet spectroscopy missions.

## 2.0 Tasks\Specifications

- This statement of work consists of three phases. Phase 1 is the fabrication planning phase, Phase 2 is the fabrication and preliminary testing phase, and Phase 3 is the final testing phase. At the end of each Phase, a written report is required.
- RFP Note: Although there are several tests which shall be required on the diffraction grating blank, as listed below, a principal area
  of concern is that test or tests used to measure the surface figure. This may be an interferometric test, or otherwise.
  - 2.0.1

The contractor shall be responsible not only for the figuring and polishing of the diffraction grating blank, but also for designing, preparing, verifying, and implementing the testing procedure(s) required to verify the correct figuring of the blank to the accuracy required by the specifications below. No Government materials, designs, or test equipment would be supplied to the contractor as part of this contract.

## 2.0.2 Coordinate system

The coordinate system used to describe the diffraction grating blank is right handed, with  $\mathbf{x}$  axis normal to the piece at the vertex (oriented so as to emerge out of the illuminated face),  $\mathbf{y}$ axis parallel to the width, and  $\mathbf{z}$  axis parallel to the height. The diffraction rulings, when applied, would be parallel

to the z axis. Unless stated otherwise, all coordinates are given in this coordinate system.

## 2.0.3 Analytic expression for the optical surface

The analytic form of the surface can be expressed as a sag from the tangent plane x=0, with

$$x = a(1-sqrt(, 1-(y/b)^2 - (z/c)^2)) + \epsilon_0 yz^2$$

- The  $\epsilon_0$  term is the asymmetric deformation and represents the deformation from an ellipsoid. The maximum sag over the active area as compared to the undeformed ellipsoid is  $\pm 2.315$  microns. A sag table expressing the **x** coordinate of the surface on a rectangular grid of points (y,z) is attached for reference.
  - When ruled, the diffraction grating grooves shall be aligned with the z axis.

The specifications are enumerated below as tasks. For convenience, they also appear in attachment 3 to this Statement of Work. In the case of a conflict, the list of tasks below shall be used.

- 2.1 Mechanical specifications
  - 2.1.1 Surface Figure and shape

The surface of the diffraction grating blank shall be a deformed ellipsoid of the analytic form above in §2.0.4. The shape of the diffraction grating blank as delivered will be that of an 203.2 mm (8.00 inch) diameter circle. The thickness at the vertex shall be at least 20 mm, but no more than 30 mm. The backside of the diffraction grating blank shall be flat, with the normal to the backside aligned to the **-x** axis to within ±15 arcseconds.

- 2.1.2 Active area
- The active area of the diffraction grating blank shall be an ellipse with semi axes 92.5 mm (along the y axis) and 77.5 mm (along the z axis).

## 2.1.3 Alignment Fiducials

The contractor shall mark on the edge of the diffraction grating blank at the points where the positive y and z axes intersect the

blank. The two fiducial marks shall be distinguishable, permanent, and documented in project reports.

2.1.4 Material

The material to be used for the diffraction grating blank shall be an ultra low expansion glass (e.g. ULE or Zerodur).

2.1.5 Ellipsoidal radii

The ellipsoidal radii along the x,y, and z axes, respectively, shall be a =  $1523.081 \pm 0.15$  mm, b =  $1672.920 \pm 0.15$  mm, and c =  $1523.081 \pm 0.15$  mm.

2.1.6 Deformation coefficient

The deformation coefficient, described above in §2.0.4, shall have the value  $\epsilon_0 = 1.128E-8 \pm 0.01E-8 \text{ mm}^{-2}$ 

- 2.2 Figuring Tolerances
  - 2.2.1 Figure error
- The surface will be fabricated such that the peak to valley surface figure error will be 1/8 wavelength (wavelength = .6328  $\mu$ m).
  - 2.2.2 Roughness

The microroughness goal is 2.0nm rms or less. The midfrequency error (correlation length of 1 mm) is 2.0nm rms or less.

- 2.3 Phase 1 -- Fabrication planning
- The contractor shall prepare and submit a written fabrication and testing plan and schedule. The contractor shall provide mechanical drawings of the proposed diffraction grating blank.
  - 2.4 Phase 2 -- Fabrication and initial testing

In this Phase, the contractor shall fabricate the diffraction grating blank to the required surface figure tolerance as above in §2.2.1. Initial tests as required to verify that the surface figure tolerance has been met shall be performed. A written report establishing that this tolerance has been met shall be prepared and submitted. This report shall include test data demonstrating the surface figure quality that was achieved.

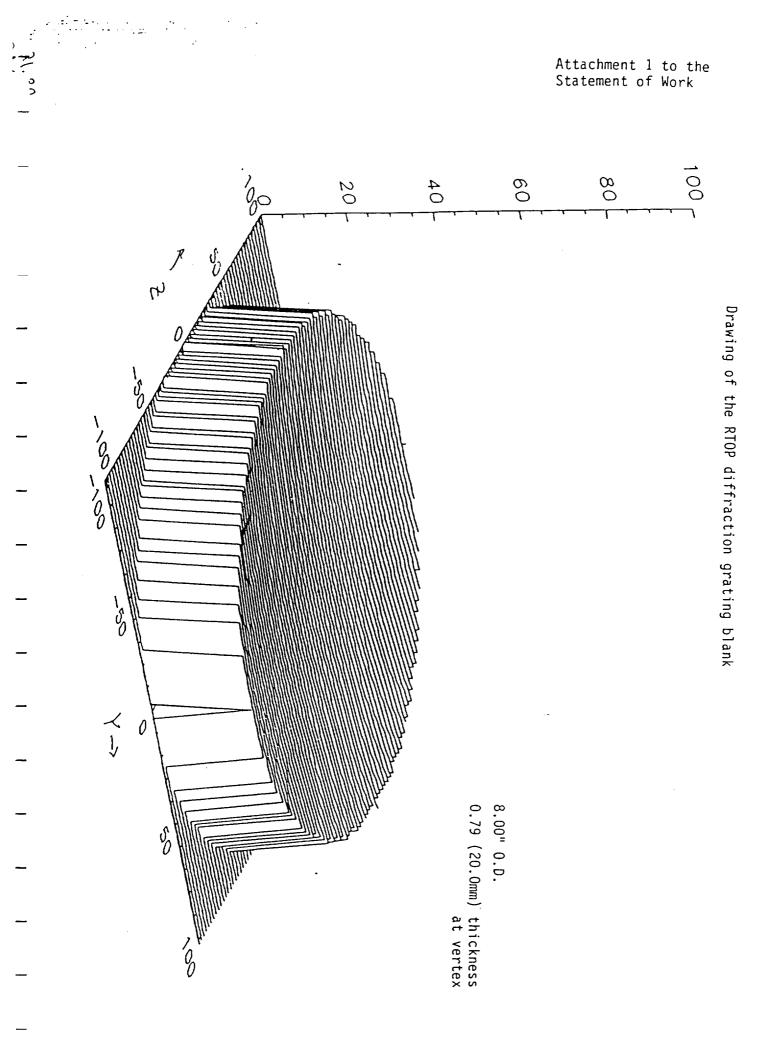
## 2.5 Phase 3 -- Final testing

In this phase, the completed diffraction grating blank will be tested for final surface figure, radii of curvature, flatness and parallelism of the back face, microroughness, midfrequency ripple, and mechanical dimensions. These tests shall measure these quantities to at least the precision required to ascertain whether they do or do not meet the specifications above.

## 2.5.1 Final testing requirements

The contractor is required to perform all necessary measurements to verify that the grating blank meets all specifications. The characterization shall include mechanical measurements, establishing overall surface figure and surface roughness, midfrequency ripple, radii of curvature, and mechanical dimensions. The contractor shall prepare and submit a written report on these tests. All test data shall be included in the report.

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Attachment 2 to the Statement of Work

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Sag table (in mm) for RTOP grating blank

-100 -50 0 50 100 У z -80 4.822531 2.780211 2.102456 2.787430 4.836970 3.248002 1.205098 0.525342 1.206903 3.251611 -40 2.723524 0.680424 0.000000 0.680424 2.723524 0 3.248002 1.205098 0.525342 1.206903 3.251611 40 4.822531 2.780211 2.102456 2.787430 4.836970 80

Attachment 3 to the Statement of Work

# Table of Specifications

	Wuantity	value	Units	Allowabie erro:	• •	
	Surface Figure Shape of blank		ellipsoid mm ter circle	±0.2 mm		
_	Material Active width Active height (active area is	Zerodur 185.00 155.00 s elliptic	mm mm al with the	±0.2 mm ±0.2 mm ese axes:		
	Thickness (@center)	20.00 ,	mai	<u>±1</u> mπ		
_	radius a (x) radius b (y) radius c (z)	1523.081 1672.920 1523.081	מש מומו מש	±0.15 mm ±0.15 mm ±0.15 mm		
	deformation coefficient ∈ <sub>0</sub>	1.128E-8	מומו <sup>2</sup> -	±1E-10mm <sup>-2</sup>		
	Figure error (P/V) microroughness (correlation length tbd)	1/8 wave 2 nm rms	.6328 µm goal			
-	roughness (1 mm correlation length)	2 nm rms	goal			
	Back surface shape	Flat (normal a	1/4 long -x ax 1 ±15 arcse	wave P/V (is, econds)	,*	
—	Mechanical mounting					
_	Notes: 1) The grooves wou is 155 mm long	uid be alo	ng the heig	nt, e.g. the lo	ndest groove	

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# **CERTIFICATE OF CONFORMANCE**

## Deformed Ellipsoidal Diffraction **MODEL NUMBER:** Grating Blank

**QUANTITY:** 1

COMPANY: NASA/Goddard Space Flight Center

SN4043 SORL SALES ORDER NUMBER:

NAS5-31790 **PURCHASE ORDER NUMBER:** 

November 19, 1993 DATE:

Space Optics Research Labs certifies that this product complies with the material, quality, and dimensional specifications as set forth in the enclosed Quality Assurance Documentation and the above mentioned order.

<u>Alan E. Lle Ceu.</u> Alan E. DeCew, Jr. General M

**General Manager** 

# **QUALITY ASSURANCE DOCUMENTATION**

NASA/Goddard Space Flight Center **CUSTOMER:** 

Deformed Ellipsoidal Diffraction Grating Blank **MODEL NUMBER:** 

NAS5-31790 **PURCHASE ORDER NUMBER:** 

SN 4043 SALES ORDER NUMBER:

		<b>SPECIFICATION</b>	ACTUAL
MATERIAL:		ZERODUR <sup>●</sup>	ZERODUR <sup>®</sup>
DIAMETER:		8.000" ± 0.005"	8.002"
VERTEX THIC	KNESS:	30mm +00 -10	25.45mm
CLEAR APERT	URE (Elliptcal): MAJOR AXIS: MINOR AXIS:	185m 155mm	185mm 155mm

ELLIPSOID OF REVOLUTION (ABOUT MAJOR AXIS) RADII:

MAJOR AXIS: MINOR AXIS:

FIGURE:

FLATNESS:

1672.92mm 1523.081 mm

Nominal <u>+</u> .1mm • Nominal <u>+</u> .1mm

SURFACE ACCURACY:

@ 632.8nm over 100% C.A. MICROROUGHNESS: MIDFREQUENCY:

20 A rms 20 A rms

1/8 Wave P-V

1/8 Wave P-V @ 632.8nm over 100% C.A. Not Measured 17 A rms

COATING:

None

**BACK SURFACE:** 

1/4 Wave P-V @ 632.8nm

4 Wave P-V Concave

<15 arc-sec

None

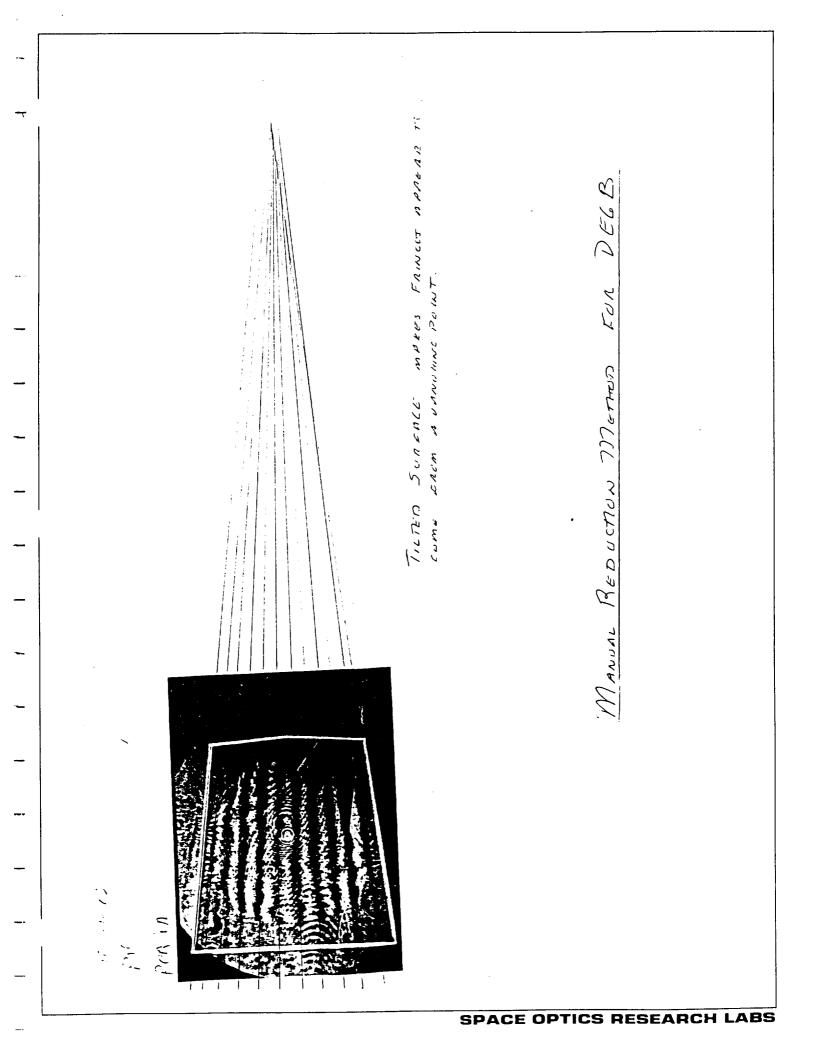
NORMAL TO +X AXIS: 15 arc-sec

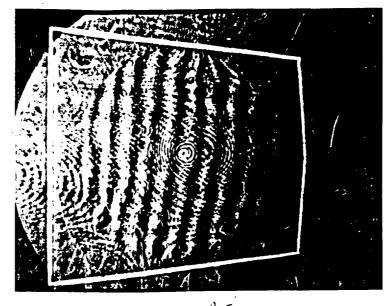
Alan E Alilia, J Alan E. DeCew, Jr.

General Manager

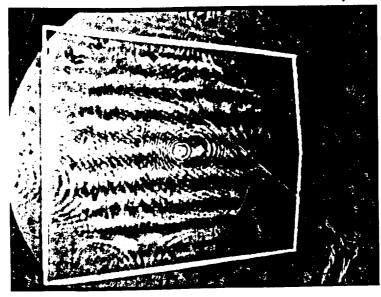
11 /18/93

Date





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<ul> <li>Subject Category 35</li> <li><b>13. ABSTRACT</b> (Maximum 200 words) The Deformed Ellipsoidal Grating Blank (DEGB) is the primary component in an ultraviolet spectrometer. Since one of the major concerns for these instruments is throughput, significant efforts are made to reduce the number of components and subsequently reflections. Each reflection results in losses through absorption, and scattering. It is these two sources of photon loss that dictated the requirements for the DEGB. The first goal is to shape the DEGB in such a way that the energy at the entrance slit is focused as well as possible on the exit slit. The second goal is to produce a surface smooth enough to minimize the photon loss due to scattering. </li> <li>The program was accomplished in three phases. The first phase was the fabrication planning. The second phase was the actual fabrication and initial testing. The last phase was the final testing of the completed DEGB. </li> </ul>					
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