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Pool and Flow Boiling in Variable and Microgravity

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INTRODUCTION

As is well known, boiling is an effective mode of heat transfer in that high heat **flux** levels **are** possible with relatively **small** temperature differences. Its optimal application **requires that** the process be adequately understood. A **measure** of the understanding of any physical event lies in the ability to predict its behavior in terms of the relevant parameters. Despite many years of research the predictability of boiling is currently possible only for quite **specialized** circumstances, e.g., the critical heat **flux** and film boiling for the pool boiling *case,* and then only with special geometries.

Variable gravity down to microgravity provides the opportunity **to** test this understanding, but possibly more important, by changing the dimensional and time scales involved permits more detailed observations of elements involved in the boiling process, and perhaps **discloses** phenomena heretofore unknown.

The focus here is on nucleate boiling although, **as** will be demonstrated below, under certain circumstances in microgravity it can take place concurrently with the dryout process. In the presence of earth gravity or forced convection effects, the latter process is usually referred **to as** film boiling. However, no vapor film as such forms with pool boiling in microgravity, only dryout. Initial results are presented here for pool boiling in microgravity, **and** were made possible at such an early date by the availability of the Get-Away-Specials (GAS).

Also presented here are some results of ground testing of a flow loop for the study of low velocity boiling, eventually to **take** place **also** in microgravity. In the interim, variable buoyancy normal **to** the heater surface is achieved by rotation of the entire loop relative to **earth** gravity. Of course, this is at the expense of varying the buoyancy patallel to the heater surface. Two questions which must be resolved early in the study of flow boiling in microgravity are **(1)** the lower limits of liquid flow velocity where buoyancy effects become significant to the boiling process **(2)** the effect of lower liquid flow velocities on the Critical Heat Flux when buoyancy is removed. Results of initial efforts in these directions **are** presented, albeit restricted currently to the ever present **earth** gravity.

POOL BOILING

Before a nucleate pool boiling system **qan** attain the steady periodic behavior normally observed in a gravity field, where buoyancy is the predominant vapor removal mechanism, the process must pass through a transient phase. Where the buoyancy is drastically reduced, **as** in microgravity, the process is anticipated to be inherently transient **unless** some special circumstances can be provided to maintain **a** subcooled bulk liquid domain. Absent this, the elements of transient boiling possible **are:**

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- (a) Conduction
(b) Onset of na Onset of natural convection
- (c) Nucleation
- (d) Vapor bubble growth/collapse

(e) Denarture
- (e) Departure **(f)** Motion for
- **(fl** Motion following **departwe**

In the transient pool boiling experiments conducted in the microgravity of space in the GAS on the shuttle **STS-47,** only elements (a), (c) and **(a)** were present and **are** considered here. The heater surface used is rectangular in shape, 1.91 cm by 3.81 cm $(0.75 \times 1.5$ inches), consisting of a 400 Angstrom thick semi-transparent gold film sputtexed on **a** quartz substrate, and serves simultaneously **as** a heater and **a** resistance thermometer. Viewing is provided simultaneously from the underside and side of the heater surface. Degassed R-113 with a normal boiling point of 47.2^oC (117^oF) was used as the working fluid. The subcooling level was controlled by varying the system pressure, which was maintained constant to within $+ 690 \text{ N/m}^2$ ($+ 0.1 \text{ psi}$) during any particular test. The experimental technique followed is described in ref. **(1).**

Measurement of the mean heater surface resistance, and hence temperature, was obtained at **10** *Hz,* and permitted the computation of heat transfer by conduction to the substrate, and to the fluid, from which a mean heat transfer coefficient was determined.

Figures **1** and **2** present the measured heater surface temperature and heat transfer coefficients for **2** of the **9 tests** conducted. It is noted that a quasi-steady process exists for this subcooled case during the some **75 seconds** of active boiling, although a tendency toward a **decrease exists at** the end. The 9 tests which constitute the matrix of test conditions *are* **shown** in Table **1.** The behavior **at** earth gravity obtained following the space flight **are** compared with that in microgravity. Boiling in microgravity under these conditions, with a relatively large flat heater surface, *appears* quite unstable.

The heater surface superheat at the moment of nucleation are plotted in Fig. **3 as** a function of heat flux with subcooling **as** a parameter for the identical system in microgravity, and post-and pre-flight. **A** rather anomalous behavior is noted in that the heater surface superheat required **to** nucleate the fluid is a maximum at the intermediate heat flux level. **A** similar behavior occurred with a different system in the **STS-57.** The maximum is also noted when the system is operated inverted in earth gravity. In ground based testing reported in ref. (2), nucleation at $a/g =$ -1 occurred at heater surface superheats of approximately 10° C at heat flux levels varying from 50 w/cm² at a Fig. **3.** Future space experiments are proposed to explore heat flux levels down to **0.5** w/cm2, which would extrapolate Fig. **3 to** the left. subcooling of 11° C to 22 w/cm² with a saturated liquid. These are consistent with extrapolations to the right in

It is obsewed in Fig. **2,** which applies for the lowest level of heat flux and virtually saturated liquid, that a correspondence exists between decreases in the mean surface temperature and increases in the mean heat transfer coefficient. This is not unexpected, since the latter is determined from the former. Upon viewing the motion pictures taken through the heating surface from the underside it was noted also **that** a decrease in the mean heat transfer coefficient corresponded **to** a distinct increase in the heater surface dryout area. Provision was made for the measurement of the fractional *dry* area of each frame of the **16** mm film, and examples **are** given in Fig. 4 for Run It ansfer coefficient corresponded to a distinct increase in the heater surface dryout area. Provision was made for the measurement of the fractional dry area of each frame of the 16 mm film, and examples are given in Fig. is for the region of the fractional dry area of each frame of the 16 mm film, and examples are given in Fig. 4 for Run No. 9 in Fig. 2. Fig. 4a is for the region of 61.5 — 67.5 seconds, in which the surface is rewetting, w random, and are the subject for further proposed microgravity experiments at lower levels of heat flux and higher levels of bulk liquid subcooling.

If the assumption is made that the heat transfer **to** the fluid in the *dry* portion of the heater surface can be neglected, and that the heat transfer coefficient over the remaining wetted portion of the heater surface is approximately uniform, then this latter quantity can be determined by dividing the overall mean heat transfer coefficient by the fractional wetted area. These results **are also** plotted in Eigs. 4a and 4b, and show that the microgravity boiling heat transfer coefficients are virtually the'same during the rewetting and dryout phases, **1250** and 1200 w/cm²C, respectively. Each data point in Fig. 4 corresponds to a single frame, taken at 10 pps, and the **oscillations** *are* a reality in the physical process. The rather large excursions in the boiling heat transfer coefficients early in Fig. 4a and late in Fig. 4b are a consequence of the inherent relatively large uncertainty in measurement of the **fractional** wet area at low levels of wetting, and should be disregarded.

FLOW BOILING

A schematic of the low velocity forced convection boiling loop for proposed studies with **R-113** in microgravity is shown in Fig. **5,** with a more detailed view of the test section in Fig. 6. Velocities can be varied from **0.5** cm/s to **32** cm/s by pump **speed** control combined with changes in test **section** height, from **2.54** cm **(1** inch) to **0.318** cm **(0.125** inch). The use of a flow loop permits the study of boiling under steady conditions **as** well **as** under transients, **as** in the **prior** described pool boiling studies, and thus *can* accommodate the use of more massive metallic surfaces which, while more representative of engineering surfaces, also introduce complications associated with heat capacity effects. The flow loop proper occupies a volume of about **1.22 x .61 x A6** m **(48 x** 24×18 inches).

Although up to 6 heater surfaces can be accommodated simultaneously, a maximum of 3 is anticipated, with the remaining ports used for visualization. In addition **to** the gold film on quartz substrate heaters described previously, two **flat** metallic substrate heaters *are* used, with the same dimensions **as** the gold film heaters, but with maximum heat flux capabilities to 15 w/cm² and 75 w/cm². The latter heater is used for Critical Heat Flux studies with R-113, and has a larger uncertainty in the heat flux measurement at the lower levels of heat flux, below about 20 w/cm2. System pressure, fluid temperature at the test section inlet, and flow rate controls *are* completely automated. More details on the flow loop and **heater surkes are** avaiiable in refs. (3) **and** (4).

Fig. 7 is a sample result from ref. (3, showing how buoyancy influences the boiling heat transfer behavior at a low level of veloeity of 4.1 cm/s. In single phase mixed convection heat transfer **both** buoyancy and **imposed** bulk liquid flow provide the mechanisms for fluid motion, where the relative significance of these is characterized by the Richardson number. For dealing with the combined effects of buoyancy and imposed bulk liquid flow with boiling, **a** "two-phase Richardson number" is developed in ref. (5) and (6). **and** is shown **as the** lower curve in Fig. 8 **as a** function of Nd, the product of the Weber and square of the Froude numbers. *Also* included **as** the upper curve **in** Fig. 8 is the **square** of the velocity ratio of Siegel (ref. 7). used to **desaibe** the **rise** velocity relative to a fluid flowing vertically. Tentative **bounds** are included in Fig. 8 for inertial and buoyant dominated domains, based **as yet on** limited experimental data It is anticipated that experimentation proposed for parabolic flights in **aircraft** will provide additional **data** for bracketing more closely **the** relative influences of buoyancy and inertial effects. limited experimental data. It is anticipated that experimentation proposed for parabolic flights in aircraft will provide
additional data for bracketing more closely the relative influences of buoyancy and inertial effects

The angle $\theta = 0$ applies to the horizontal upward facing orientation. The CHF is normalized relative to a pool boiling correlation of ref. (1), q_{co}, which includes the influence of bulk liquid subcooling. The curves labeled "model" in Figs. 9 and 10 are for pool boiling modified by multiplying by the square root of θ over the interval of **90 to** 270 deg., and **arise** from equating buoyancy and drag forces in the inverted positions where **the** vapor bubbles **are** held against the heater surface **as** they slide. The model is described in ref. (4).

The onset of forced convection effects become quite distinctive between Figs. 9 and 10, **and** it *can* be anticipated that the **CHF** will become independent of **8** at sufficiently large velocities. Plans *are* underway to increase the maximum velocity attainable **to** about *60* cm/s. It *can* also be anticipated that the behavior of the CHF will take on quite another character in microgravity, where buoyancy effects become totally negligible. The only forces remaining then **are** momentum (or inertia) of the bulk liquid flow and surface tension, both liquid-vapor and liquid-vapor-solid. These can be expected **to** influence not only the CHF *(or* drywt - more liely), but the departure *sizes* of **the** vapor bubbles **as** well.

REFERENCES

- 1. Ervin, J. **S.,** Merte, H., Jr., Keller, R. B., Kirk, K., "Transient Pool Boiling in Microgravity," Int. J. Heat **Mass** Trans., 35, March, 1992. pp. 659-674.
- 2. Ervin, J. **S.,** Merte, H. , Jr., "A Fundamental Study of Nucleate Pool Boiling under Microgravity," Final Report to NASA Lewis Research Center on NASA Grant NAG3-663, Report No. UM-MEAM-91-08. August, 1991.
- 3. Kirk, **K.** M., Merte, H., Jr., Keller, R. B.. "Low Velocity Nucleate Flow boiling at Various Orientations," ASME Symposium-Fluid Mechanics Phenomena in Microgravity, Ed. by D. A. Siginer and M. M. Weislogal, **AMD** - Vol. 154/FED - Vol. 142. 1992.
- 4. Brusstar, M. J., Merte, H., Jr., "Effects of Buoyancy on the Critical Heat Flux in Forced Convection," AIAA J. Thermophysics and Heat Transfer, 8, April - June, 1994, **pp.** 322-328.
- *5.* Kirk, **K.** M., Merte, H. Jr., **"A** Study of the Relative Effects of Buoyancy and Liquid Momentum **in** Forced Convection Nucleate Boiling," **Final** Report to NASA Lewis Research Center **on** NASA Grant NAG3- 1310, Report No. UM-MEAM-92-06. November, 1992.
- 6. Kirk, K. M., Merte, H., Jr., "A Mixed Natural/Forced Convection Nucleate Boiling Heat Transfer Criteria," proceedings of loth International heat Transfer Conference, Aug. 14-18,1994, Brighton, **U.K.**
- 7. Siegel, **R.,** "Effect of Reduced Gravity on Heat Transfer," Advances in Heat Transfer, *eds.* James P. Hartnett and Thomas F. bine, Vol. **4,** pp. 200-205.
- 8. Ivey, **H.** J., **Morris, D.** J., *"On* **the** Relevance of the Vapour-Liquid Exchange Mechanism for Subcooled Boiling Heat Transfer **at** High **Pressure,"** UKAEA, AEEW-R-137,1962.

Mean heater surface temperature and heat transfer coefficient. R-113. PBE Prototype. STS-47. Run No. 2. $q_T^* = 3.6$ Figure 1. w/cm². $\Delta T_{sub} = 11.9$ °C. $T_{sat} = 61.1$ °C.

Mean heater surface temperature and heat transfer coefficient. R-113. PBE Prototype. STS-47. Run No. 9. $q_T = 1.8$ w/cm². $\Delta T_{sub} = 0.2$ °C. T_{sat} = 49.4°C.

Computed from natural convection correlation: $Nu = 0.15 \times Ra^{1/3}$
Detailed results following.

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Table 1.

Comparison of measured mean heat transfer coefficients for PBE Prototype between $a/g = +1$ and STS-47 space flight.

Heater surface superheat at nucleation for PBE Prototype System for $\alpha/g = \pm 1$ and STS-47 space flight. Figure 3.

a. Rewetting poxtion of Figure 2. 61 *⁵*- **67.5 sec.**

b. Dryout portion of Figure 2. $80.5 - 85.5$ **sec.**

Figure 4. Fractional wetted area — nucleate boiling heat transfer coefficient — mean heat transfer coefficient with rewetting and dryout.
PBE Prototype. STS-47. Run No. 9. qT = 1.8 w/cm². $\Delta T_{sub} = 0.2$ °C. T_{sat} = 49.4°

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