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MEMORANDUM

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TO VIBRATIONS OF A 45° DELTA WING

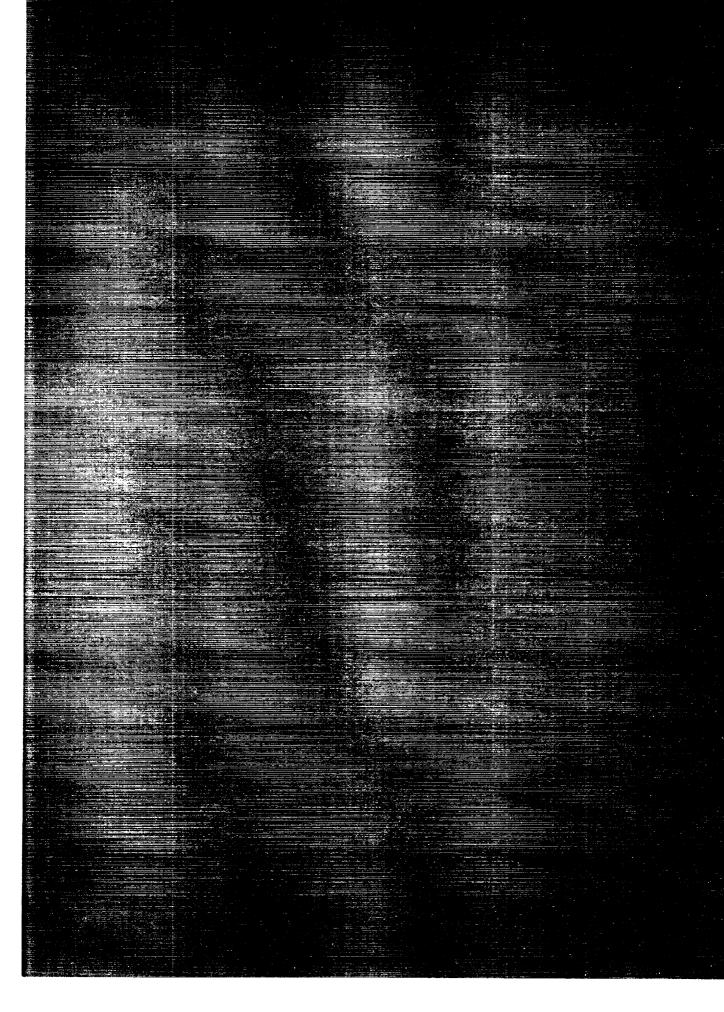
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EVALUATION OF THE LEVY METHOD AS APPLIED

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SUMMARY

The Levy method which deals with an idealized structure was used to obtain the natural modes and frequencies of a large-scale built-up 45° delta wing. The results from this approach, both with and without the effects of transverse shear, were compared with the results obtained experimentally and also with those calculated by the Stein-Sanders method. From these comparisons it was concluded that the method as proposed by Levy gives excellent results for thin-skin delta wings, provided that corrections are made for the effect of transverse shear.

INTRODUCTION

The literature contains many methods for obtaining the deflectional characteristics of low-aspect-ratio and delta wings. (See, for example, refs. 1 to 5.) Although these methods use a variety of approaches and assumptions, they can be classified into two categories: the method either deals with the actual structure and restricts the allowable deflection shape or deals with a simplified structure and allows arbitrary deflections. One analysis from the first category, the Stein-Sanders method, is described in reference 1. In this analysis, the actual structure was analyzed by assuming that its neutral surface was strain free, the effects of transverse shear were negligible, and its chordwise deformation was parabolic. An analysis from the second category, namely the Levy method, is described in reference 2. In this method an idealized structure consisting of interconnected beams and torque boxes whose deflections are unrestrained is analyzed.

Although methods of calculating the deflectional characteristics of low-aspect-ratio and delta wings do exist, there is available very little information concerning the application of the methods and the reliability of their results.

An experimental investigation of the stiffness and vibration characteristics of a large-scale built-up 45° delta wing has been discussed in

reference 6. Since the detailed stiffness and weight distributions of the specimen are presented therein, the results of the investigation can serve as a reliable basis for the evaluation of the analytical methods. These results have been used in reference 7 to evaluate the Stein-Sanders method. In the present paper the experimental results are used to evaluate the Levy method. A summary of some of the results of this investigation was presented in references 8 and 9.

The purpose of the present paper is threefold: First, to describe in detail the application of the Levy method to a 45° delta wing; second, to show how the Levy method can be easily modified to include approximately the influence of transverse shear; and third, to evaluate the method in the light of the results of the Stein-Sanders method and experimental results.

SYMBOLS

A_SG	shear stiffness of beam
D	constant defined in equation (A8)
E	modulus of elasticity
EI	bending stiffness of beam
GJ	torsional stiffness of torque box
h	depth of beam
i,j,n,N	integers
J	torsional constant
K_{ij}	constants defined in equation (A7)
ı	length of torque box
M, M_x, M_y M_{xx}, M_{yy}, M_z	constants defined in equation (A9)
P _i	concentrated load at station i
V	shear in beam web

```
deflection of ith station of free wing
w<sub>1</sub>
w, 3P
           deflection of ith station of wing on three-point support
           rigid-body translation
w_{o}
х,у
           coordinates of station
           distance of force from support
x_0
α,β
           a rigid body tipping about y- and x-axis, respectively
           influence coefficient of cantilevered beam
Δ
           stiffness coefficient of wing
\Delta^{3P}
           stiffness coefficient of wing on three-point support
\triangle_S^i, \triangle_R^i, \triangle_T^i stiffness coefficient of ith spar, rib, and torque box,
             respectively
           circular natural frequency
ω
\lceil F \rceil
           square matrix defined in equations (Al2) and (A24)
H
           square matrix defined in equations (All) and (A21)
unit matrix
           diagonal mass matrix for half-span
11
           row matrix of ones
11
           column matrix of ones
           rectangular matrix
           diagonal matrix
           row matrix
           column matrix
```

Subscripts:

C stations on center line

R stations on right side of center line

L stations on left side of center line

i,n integers

Superscripts:

s symmetrical

a antisymmetrical

ANALYSES

Specimen

The specimen used in the investigation discussed in reference 6 is a large-scale built-up 45° delta wing shown in figure 1. It has a span of 18 feet $11\frac{7}{8}$ inches, a midchord of 8 feet $1\frac{5}{8}$ inches and a uniform carrythrough bay of 2 feet 8 inches. The wing is uniform in depth in the chordwise direction but varies linearly in depth in the spanwise direction from $5\frac{1}{2}$ inches at the carrythrough section to $1\frac{3}{4}$ inches at the tip.

The top and bottom covers of the delta wing are of skin stringer construction with four light stringers between each spar. The interior construction consists of four straight spars spaced 24 inches apart, a bent leading-edge spar, and light streamwise bulkhead spaced 8 inches on centers. Detailed dimensions, section properties, and weight distribution of the specimen are given in reference 6. All parts were constructed of 2024-T6 aluminum alloy.

Idealization

In order to apply the Levy method, the actual structure in figure 1 was idealized as shown in figure 2 into an orthogonal set of crisscrossing beams with torque boxes attached at their four corners to the intersection of the beams. The locations of the idealized spars were chosen to coincide with the center line of the actual spars. The spacing of the ribs

in the idealized structure, however, was increased over that of the specimen in order to decrease the number of redundants in the analysis from 53 (if the actual rib locations are used) to 34.

All the spanwise normal-stress-carrying material of the spars, cover sheets, and stringers was concentrated into the spars of the idealized wing whereas all the chordwise bending ability of the actual ribs and covers was accounted for in the idealized ribs. The condition suggested by Levy (see ref. 2) of limiting the effectiveness of the sheets in the chordwise direction to 0.181 of the rib length to either side of the rib governed only in the last two outboard ribs of the actual structure. The stiffnesses of the idealized ribs were obtained by first distributing the moments of inertia of the actual ribs and then reconcentrating the inertias at the new stations. The moments of inertias of the idealized spars and ribs are given in tables I and II.

The shear-carrying capacity of the cover sheets is accounted for by the torque boxes in the spar-rib cells of the idealized structure. In the calculation of the torsional stiffness GJ of these boxes, the axis of twist was assumed to be in the spanwise direction. The values of J at the center section of each torque box are given in table III. Note that, when these values were obtained, the side walls of the torque boxes were considered to be rigid in shear as suggested by Levy in reference 2.

Application of Levy Method

The first step in the analysis of the idealized wing is to determine the loads carried by the individual components in terms of the deflection at the junctions of the spars and ribs. These loads can be expressed as follows:

$$|P| = E[\Delta_S^n]|w|$$
 $(n = 1,2,...5)$ (la)

$$|P| = E\left[\Delta_{R}^{n}\right]|W|$$
 (n = 1,2,...10) (1b)

$$|P| = E[\Delta_T^n]|w|$$
 (n = 1,2,...20) (1c)

where Δ_S^n , Δ_R^n , and Δ_T^n are the stiffness coefficients of the nth spar, rib, and torque box, respectively. In equation (lb), n = 10 refers to the swept portion of the leading-edge spar.

In these calculations the influence of shear deformation in the spar and rib webs along with the torque-carrying capacity of the triangular cells was neglected. Furthermore, no moment transfer was permitted to take place between the spars and ribs and between the straight and swept portion of the leading-edge spar. The stiffness coefficients of the nonuniform spars were obtained as described in reference 2 by inversion of the influence coefficients of cantilevered beams. influence coefficients were calculated by an approximate procedure described in reference 10, which was based on an assumption of a linear 1/EI variation between stations. An example of the resulting influencecoefficient matrix $\lceil \delta \rceil$ is shown in table IV(a) for the trailing-edge spar. When the stiffness coefficients of the spars were calculated, cognizance of the type of loading was taken. For the case of symmetrical loading the stiffness coefficients of the spars were obtained for the condition of zero slope at the center line, whereas for antisymmetrical loading the condition of zero deflection at the center line was maintained. The resulting stiffness coefficients for symmetrical loading for the trailing-edge spar $\begin{bmatrix} \triangle_S^1 \end{bmatrix}$ are shown in table V.

Inasmuch as the ribs and torque boxes were uniform, there were no complications involved in the calculations of their stiffness matrices. Typical examples of the stiffness coefficients are shown in table VI for rib number 4 and in table VII for torque boxes 15 and 16. The stiffness coefficients of the swept portion of the leading-edge spar were obtained by considering that the swept portion of the spar acts as a rib and that no moment is transferred at any point of attachment including the junction of the unswept and swept portion of the spar.

The loads carried by the idealized structure are considered to be the sum of the loads carried by the idealized spars in spanwise bending, by the ribs in chordwise bending, and by the torque boxes in torsion. Thus the stiffness coefficients of the composite structure were obtained by summing the stiffness coefficients of the components:

$$|P| = E[\Delta]|w|$$
 (2)

where

$$\left[\Delta\right] = \sum_{n=1}^{5} \left[\Delta_{S}^{n}\right] + \sum_{n=1}^{10} \left[\Delta_{R}^{n}\right] + \sum_{n=1}^{20} \left[\Delta_{T}^{n}\right]$$
(3)

The synthesis of a typical row of $\left[\Delta\right]$ for symmetrical loading is illustrated in table VIII for row 24. The elements of this row represent the contribution of the deflections at each station of the wing to

the load at station 24. As can be seen, elements associated with stations not on the spar, rib, or torque boxes common to station 24 are zero. The remaining elements of row 24 are the summations of the rows of $\begin{bmatrix} \Delta_S^1 \end{bmatrix}$, $\begin{bmatrix} \Delta_T^{4} \end{bmatrix}$, $\begin{bmatrix} \Delta_T^{15} \end{bmatrix}$ and $\begin{bmatrix} \Delta_T^{16} \end{bmatrix}$ associated with P_{24} and are shown in tables V to VII.

As yet there have not been any restraining or boundary conditions placed on the stiffness matrix $[\Delta]$. Thus the structure represented by this matrix is free to move with a rigid-body displacement. Obviously, the deflections of such a structure are not uniquely related to the loads and therefore the inverse of its stiffness matrix cannot exist, that is, the matrix $[\Delta]$ is singular. In order to obtain a structure whose stiffness matrix can be inverted, the wing was assumed to be simply supported at three points (stations 1 and 22) in figure 2. This particular support condition was used because the results from a three-point support can be converted to influence coefficients for most other support and loading conditions. The particular stations used were chosen to conform to the supporting condition used in the static tests of the delta wing described in reference 6.

The stiffness matrix of the wing on a three-point support was obtained by omitting from the $\left[\Delta\right]$ matrix the rows and columns associated with stations 1 and 22. The resulting stiffness matrix $\left[\Delta^{3P}\right]$ for a delta wing on a three-point support is shown in table IX for both symmetrical and antisymmetrical cases. The influence coefficients of the idealized structure on a three-point support were obtained by inverting the $\left[\Delta^{3P}\right]$ matrix

$$\left|\mathbf{w}\right| = \frac{1}{E} \left[\Delta^{3P} \right]^{-1} \left| \mathbf{P} \right| \tag{4}$$

The influence coefficient matrices $\left[\Delta^{3P}\right]^{-1}$ for symmetrical and antisymmetrical loading conditions are shown in table X.

Since the influence coefficients of the delta wing are known for a three-point support, the load deflection characteristics of the wing can be calculated for other support conditions. (See ref. 1.) Furthermore, the frequency equations necessary to determine the natural modes and frequencies can readily be obtained. A method for "freeing" a wing is discussed in the appendix. In this method the displacements of a free-free wing vibrating in a natural mode are described in terms of the influence coefficients of the wing on a three-point support. With the use of the results of the appendix, the frequency equation for a free-free wing can be written as follows (see eqs. (A23) and (A27)):

For symmetrical vibrations:

$$|w| = \frac{\mathbf{w}^2}{\mathbf{E}} \left[\mathbf{I} \right] + \left[\mathbf{F}^{\mathbf{S}} \right] \left[\mathbf{M}^{\mathbf{S}} \right] \left[\Delta^{\mathbf{S}} \right]^{-1} \left[\mathbf{M}^{\mathbf{S}} \right] |w|$$
 (5)

For antisymmetrical vibrations:

$$\left|\mathbf{w}\right| = \frac{\omega^2}{E} \left[\left[\mathbf{I}\right] + 2K_{22} \left|\mathbf{x}\right| \left[\mathbf{x}\right] \left[\Delta^{\mathbf{a}}\right]^{-1} \left[\mathbf{M}^{\mathbf{s}}\right] \left|\mathbf{w}\right|$$
 (6)

where

[FS] matrix defined in eq. (A24)

K₂₂ constant defined in eq (A7)

[I] unit matrix

x spanwise coordinate

 $\left[\triangle^{s}\right], \left[\triangle^{a}\right]$ stiffness matrix for wing on three-point support for symmetrical and antisymmetrical loading conditions, respectively

MS diagonal mass matrix for half-span

The elements of the diagonal mass matrix represent the mass that is considered to be concentrated at each station. In order to obtain these elements the components of the wing tabulated in reference 6 were divided into two groups. One group contained the cover sheets, stringers, spars, and spar-to-cover and stringer-to-cover rivets and the second group contained the weights of the ribs and the concentrated weights (such as those of the filler blocks, splice plates, pickup, and the moving elements of vibrators). The contribution of the components of the first group to the elements of the mass matrix was obtained by dividing the wing into regions (shown in fig. 3) and then allotting the weights of the portion of the components included in each region to the station associated with the region. The contribution of the components of the second group was obtained in such a way that the total and first and second moments about the wing center line of these contributions were the same as the total and first and second moments of the weight of the actual components in the second group. The sum of all the weights associated with the stations shown in figure 3 was within 0.1 percent of the actual weight of the wing.

Modification of the Levy Method to Include Transverse Shear

In the previous calculations the effects of transverse shear were neglected as suggested in reference 2. On the other hand in reference 9 it was shown that the influence of transverse shear could be of importance especially in the higher modes of vibration.

If the effects of transverse shear were to be included exactly in a consistent deformation analysis, such as that of reference 2, the slopes in both the spanwise and chordwise direction in addition to the deflections at each spar-rib intersection must be treated as unknowns. This requirement would, of course, cause a threefold increase in the number of redundants in the solution. The influence of transverse shear, however, can be included in the Levy method approximately with no increase in the number of redundants and with little additional labor.

In the previous calculations the stiffness coefficients of the spars and ribs were obtained by inversion of the influence coefficients of cantilever beams. These influence coefficients, however, contained only the deflections due to bending. The effects of shear deformation on the spars and the ribs can be included in the influence coefficients by super-imposing the deflections due to shear onto those due to bending. The influence coefficients including shear deformation can be obtained from the equation

$$w = \int_{0}^{x} \frac{P}{EI} (x_{O} - \eta) (x - \eta) d\eta + \int_{0}^{x} \frac{P}{A_{S}G} d\eta + \int_{0}^{x} \frac{P}{A_{S}G} \frac{h'}{h} (x_{O} - \eta) d\eta + \int_{0}^{x} \frac{P}{A_{S}G} \frac{h'}{h} (x_{O} - \eta) d\eta + \int_{0}^{x} \frac{P}{A_{S}G} (\frac{h'}{h})^{2} (x_{O} - \eta) (x - \eta) d\eta$$
 (7)

where w is the deflection of a cantilever beam at any point x (distance from the root) due to a load P at $x_{\rm O}$, h' is the derivative of h with respect to η , and EI and $A_{\rm S}G$ are the bending stiffness and effective shear stiffness, respectively. The first term on the right-hand side of equation (7) is the portion of the deflection due to bending stresses. The second term is the shear deformation that would occur if the beam was nontapered. The third term represents the deflection due to the effect of the normal stresses in the flanges of the tapered beams on the shear in the webs. The last two terms represent the deflections due to the effects of taper on the shear strain.

As an example, the influence coefficients with transverse-shear deformations included are shown in table IV(b) for the trailing-edge

spar. Comparisons of these coefficients with those in table IV(a) will give an indication of the magnitude of the transverse-shear deformation. In these calculations the effective shear areas of the spar and rib webs were taken to be the product of the web thickness and the depth of channel.

The set of influence coefficients for all spars and ribs resulting from the use of equation (7) was inverted to obtain the stiffness coefficients of the spars and ribs. The stiffness coefficients of the torque boxes were left unchanged.

The influence coefficients of the idealized delta wing were then obtained in the same manner as described in the previous section. The numerical values of the resulting influence coefficients including transverse shear are shown in table XI for the wing simply supported at three points and loaded both symmetrically and antisymmetrically.

RESULTS AND DISCUSSION

The first nine free-free modes (5 symmetrical and 4 antisymmetrical) of the delta wing were calculated with the use of equations (5) and (6) for both the case where transverse shear was neglected and the case where the influence of transverse shear was included.

In figure 4 the node lines and frequencies as obtained by the Levy method with transverse shear neglected are compared with the node lines and frequencies obtained by the Stein-Sanders method (ref. 7) and with the experimental node lines and frequencies (ref. 6).

Note that the frequencies given in figure 4 for the Levy method are smaller than those given in reference 8. This discrepancy was due to the fact that, in the calculations for the frequencies in reference 8, 12 inches of the cover sheet were included in the moments of inertia of the leading-edge spar whereas in the present calculation only 6.14 inches were included as suggested by the criteria of reference 2. Furthermore, in the calculations of the results in reference 8, moment transfer was allowed between the unswept and swept portions of the leading-edge spar whereas in the calculations of the present paper no moment transfer was allowed.

As can be seen in figure 4, the node-line patterns of both the Stein-Sanders and Levy methods agree fairly well with the ones obtained experimentally. The node lines obtained by the Levy method, however, are not as good as those obtained by the Stein-Sanders method, especially in the vicinity of the leading edge. Examination of the figure seems to indicate that the stiffness of the leading edge in the idealized structure is too great.

Although the Stein-Sanders method predicts the experimental nodeline pattern fairly well, the frequency agreement is poor. The errors range from 7 percent in the first mode to 38 percent in the fifth symmetrical mode. On the other hand, the frequency agreement in the Levy method is much better. The largest error in the first 8 modes occurs in the third antisymmetrical mode and is only $8\frac{1}{2}$ percent; the error in the fifth symmetrical mode is only 20 percent.

One of the principal sources of error in the Stein-Sanders method is the assumption of a parabolic chordwise variation of deformation. As this particular specimen had no extra chordwise stiffening in the center section such as would be furnished by a fuselage, for example, the errors due to this assumption may be large. Another source of error which is in both the Stein-Sanders and the Levy methods is that the results shown in figure 4 do not include the effects of transverse shear.

The results of the calculations of the frequencies of the first nine free-free modes of the delta wing by various methods are summarized in table XII. The frequencies that were obtained experimentally are given in the first row. The corresponding frequencies as calculated by the Stein-Sanders method and by the Levy method without shear are tabulated in the second and third rows, respectively. The frequencies obtained by the modified Levy method that includes transverse shear are given in the fourth row. The last row contains frequencies that were calculated from the experimentally determined influence coefficients of reference 6 by the method discussed in the appendix. This calculation was included because a popular method of obtaining frequencies is to measure influence coefficients on a model or full-scale structure and then use them in a vibrational analysis.

A comparison of the results tabulated in rows 1 and 4 of table XII shows that the frequencies calculated by the Levy method with shear are in excellent agreement with the experimental frequencies. The largest error occurs in the seventh (fourth symmetrical) mode and is slightly less than 4 percent. The effect of transverse shear on the calculated nodal-line patterns was slight. The changes that did occur, however, tended to improve the agreement between the calculated and experimental node lines.

Comparison of rows 3 and 4 of table XII indicates that the effect of transverse shear can be important. For instance, the inclusion of transverse shear caused an 18-percent reduction in the calculated frequencies of the fifth symmetrical mode. Also, a comparison of frequencies shown in rows 1, 4, and 5 shows that, for this particular specimen, the modified Levy method gave results which were as good as those obtained from experimental influence coefficients.

Although a comparison of experimental and calculated frequencies provides a test of the accuracy of calculated influence coefficients, a comparison of calculated to experimental deflections of a cantilever delta wing under static loading is of some interest. Therefore the deflections of a delta-wing specimen clamped along the center line under a uniform load of one pound per square inch were obtained from the influence coefficients shown in table XI and were compared with deflections obtained from the experimental influence coefficients shown in reference 6.

The results of these calculations are shown in figure 5. The deflections of the five spars as calculated by the Levy method with transverse shear are shown by the solid lines whereas the deflections as obtained from the experimental influence coefficients are shown as points. From figure 5 it can be seen that, with the exception of the tip, the deflections as given by the modified Levy method agree well with those obtained from experimental influence coefficients. The large discrepancy in the tip deflections can be attributed to the neglect of the torsional stiffness of triangular boxes in the analysis. As can be seen from figure 2, such an assumption in the idealized beam leaves only the leading- and trailing-edge spars to transfer the tip load to the inboard stations. In the actual structure, however, the triangular box contributed a large amount of the torsional stiffness.

CONCLUDING REMARKS

From a comparison of calculated and experimental frequencies it has been shown that a method which deals with an idealized structure, such as the method proposed by Levy, gives excellent results for thinskin wings, such as the 45° delta-wing specimen investigated, provided that corrections are made for the effects of transverse shear. Furthermore, the Stein-Sanders type of approach seems to be inapplicable to low-aspect-ratio wings with center sections which have not been stiffened against chordwise bending.

Langley Research Center,
National Aeronautics and Space Administration,
Langley Field, Va., October 20, 1958.

APPENDIX A

FREEING OF INFLUENCE-COEFFICIENT MATRIX

FOR GENERAL THREE-POINT SUPPORT

Asymmetrical Structure

The problem of obtaining influence coefficients for other loading and support conditions from the influence coefficients for a three-point support was discussed in reference 1. This appendix is concerned with the problem of obtaining a frequency determinant for a structure from its influence coefficients on an arbitrarily located three-point support.

It is assumed that a structure is simply supported at three arbitrary points and that the influence coefficients of this structure at N points (including the supports) are known. The coordinate system is chosen so that the x-axis goes through two of the supports and the y-axis through the third. The deflections of this structure at any of the points in terms of loading at the points i = 1 to N are given by the following matrix:

$$|\mathbf{w}^{3P}| = [\delta]|P| \tag{A1}$$

where the elements of the matrices are

$$w_i^{3P}$$
 deflection of point i when i = 1, 2, 3, ... N

$$P_i$$
 load at station i when i = 1, 2, 3, . . . N

$$\delta_{ij}$$
 deflection at point i due to a load at point j when i, j = 1, 2, 3, . . . N

If the system is permitted to be completely unrestrained, the deflection at any point can be written as

$$|\mathbf{w}| = |\mathbf{w}^{3P}| + \mathbf{w}_0 |\mathbf{1}| + \alpha |\mathbf{x}| + \beta |\mathbf{y}|$$
 (A2)

where

w_o rigid-body translation

α rigid-body rotation about y-axis

β rigid-body rotation about x-axis

x_i,y_i coordinates of point i

column matrix of ones

The loadings on this structure must then satisfy the following equilibrium conditions:

When a structure is vibrating in its natural mode, the inertial loading can be written as:

$$|P| = \omega^2 |M| |w|$$
 (A4)

where ω is the natural circular frequency and $M_{\dot{1}}$ is the effective concentrated mass of the structure at station i. With the use of equation (A2), equation (A4) can be written as:

$$|P| = \omega^2 \left[M \right] |w^{3P}| + w_0 |1| + \alpha |x| + \beta |y|$$
 (A5)

The values of α , β , and w_0 can be obtained in terms of $\left| w^{3P} \right|$ by substituting equation (A5) into equation (A3) and solving the resulting equations to yield

$$\mathbf{w}_{o} = \begin{bmatrix} \mathbf{K}_{11} \begin{bmatrix} 1 \end{bmatrix} + \mathbf{K}_{12} \begin{bmatrix} \mathbf{x} \end{bmatrix} + \mathbf{K}_{13} \begin{bmatrix} \mathbf{y} \end{bmatrix} \begin{bmatrix} \mathbf{M} \end{bmatrix} \begin{bmatrix} \mathbf{M} \end{bmatrix} \begin{bmatrix} \mathbf{w}^{3P} \end{bmatrix}$$

$$\alpha = \begin{bmatrix} \mathbf{K}_{12} \begin{bmatrix} 1 \end{bmatrix} + \mathbf{K}_{22} \begin{bmatrix} \mathbf{x} \end{bmatrix} + \mathbf{K}_{23} \begin{bmatrix} \mathbf{y} \end{bmatrix} \begin{bmatrix} \mathbf{M} \end{bmatrix} \begin{bmatrix} \mathbf{w}^{3P} \end{bmatrix}$$

$$\beta = \begin{bmatrix} \mathbf{K}_{13} \begin{bmatrix} 1 \end{bmatrix} + \mathbf{K}_{23} \begin{bmatrix} \mathbf{x} \end{bmatrix} + \mathbf{K}_{33} \begin{bmatrix} \mathbf{y} \end{bmatrix} \begin{bmatrix} \mathbf{M} \end{bmatrix} \begin{bmatrix} \mathbf{w}^{3P} \end{bmatrix}$$
(A6)

where

$$K_{11} = \frac{1}{D} (M_{xy}^{2} - M_{xx}M_{yy})$$

$$K_{12} = \frac{1}{D} (M_{x}M_{yy} - M_{y}M_{xy})$$

$$K_{13} = \frac{1}{D} (M_{xx}M_{y} - M_{x}M_{xy})$$

$$K_{22} = \frac{1}{D} (M_{y}^{2} - MM_{yy})$$

$$K_{23} = \frac{1}{D} (M_{xy} - M_{x}M_{y})$$

$$K_{33} = \frac{1}{D} (M_{x}^{2} - MM_{xx})$$
(A7)

$$D = MM_{xx}M_{yy} + 2M_xM_yM_{xy} - M_{xx}M_y^2 - M_{yy}M_x^2 - MM_{xy}^2$$
 (A8)

and

$$M = \lfloor 1 \rfloor \lceil M \rceil | 1 |$$

$$M_{\mathbf{x}} = \lfloor \mathbf{x} \rfloor \lceil M \rceil | 1 |$$

$$M_{\mathbf{y}} = \lfloor \mathbf{y} \rfloor \lceil M \rceil | 1 |$$

$$M_{\mathbf{xy}} = \lfloor \mathbf{x} \rfloor \lceil M \rceil | \mathbf{y} | = \lfloor \mathbf{y} \rfloor \lceil M \rceil | \mathbf{x} |$$

$$M_{\mathbf{xx}} = \lfloor \mathbf{x} \rfloor \lceil M \rceil | \mathbf{x} |$$

$$M_{\mathbf{yy}} = \lfloor \mathbf{y} \rfloor \lceil M \rceil | \mathbf{y} |$$

$$M_{\mathbf{yy}} = \lfloor \mathbf{y} \rfloor \lceil M \rceil | \mathbf{y} |$$

$$M_{\mathbf{yy}} = \lfloor \mathbf{y} \rfloor \lceil M \rceil | \mathbf{y} |$$

With equation (A6), equation (A2) becomes

$$|w| = |H| |w^{3P}|$$
 (AlO)

where

$$[H] = [I] + [F][M]$$
 (All)

[I] = unit matrix

and

Substitution of equations (Al) and (A4) into equation (AlO) yields the frequency equation:

$$|\mathbf{w}| = \omega^2 [\mathbf{H}] [\delta] [\mathbf{M}] |\mathbf{w}| \tag{A13}$$

From this frequency or characteristic equation, all modes and frequencies of the free-free asymmetrical structure can be calculated. However, much simplification of the calculation is possible if the structure is symmetrical.

Symmetrical Structure

For a symmetrical structure that is symmetrically supported and whose stations are symmetrically located, the stations can be arranged in three groups: The first group has stations on the center line $x_{C,i}$, $y_{C,i}$, the second group has stations on the right-hand side of the center line $x_{R,i},y_{R,i}$, and the third group has stations on the left-hand side $x_{L,i},y_{L,i}$. Furthermore, the stations of the last group should be numbered so that the ith station on the left is symmetrical with the

ith station on the right. Thus,

$$\begin{bmatrix} \mathbf{x}_{\mathbf{C}} \end{bmatrix} = 0$$

$$\begin{bmatrix} \mathbf{x}_{\mathbf{L}} \end{bmatrix} = - \begin{bmatrix} \mathbf{x}_{\mathbf{R}} \end{bmatrix}$$

$$\begin{bmatrix} \mathbf{y}_{\mathbf{L}} \end{bmatrix} = \begin{bmatrix} \mathbf{y}_{\mathbf{R}} \end{bmatrix}$$
(A14)

The characteristic or frequency equations (AlO) can now be partitioned as follows:

$$\begin{vmatrix}
|w_{C}| \\
|w_{R}| \\
|w_{L}|
\end{vmatrix} = \omega^{2} \begin{bmatrix} H_{11} & H_{12} & H_{13} \\ H_{21} & H_{22} & H_{23} \\ H_{31} & H_{32} & H_{33} \end{bmatrix} \begin{bmatrix} \delta_{11} & \delta_{12} & \delta_{13} \\ \delta_{21} & \delta_{22} & \delta_{23} \\ \delta_{31} & \delta_{32} & \delta_{33} \end{bmatrix} \begin{bmatrix} M_{C} \\ 0 & M_{R} \end{bmatrix} = 0 \quad |w_{C}| \\ |w_{R}| \\ |w_{L}| \end{bmatrix} |w_{L}|$$
(A15)

From consideration of the symmetry of the structure and the symmetry of the station location, the following relationships exist:

$$\delta_{12} = \delta_{21} = \delta_{13} = \delta_{31}$$
 (A17)

$$\delta_{23} = \delta_{32} \tag{A18}$$

From equations (A9) and (A14) it can be seen that for symmetrical structures

$$M_{X} = M_{XY} = 0 \tag{A19}$$

and therefore

$$K_{12} = K_{23} = 0 (A20)$$

Thus, the elements of [H] in equation (All) can be defined in terms of the locations of the stations on the center line and right-hand side of the structure as

$$\begin{aligned} & H_{11} = \begin{bmatrix} \mathbf{I} \end{bmatrix} + \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_C] + K_{13} | y_C | [1] + K_{33} | y_C | [y_C] \end{bmatrix} \begin{bmatrix} M_C \end{bmatrix} \\ & H_{12} = H_{13} = \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_R] + K_{13} | y_C | [1] + K_{33} | y_C | [y_R] \end{bmatrix} \begin{bmatrix} M_R \end{bmatrix} \\ & H_{21} = H_{31} = \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_C] + K_{13} | y_R | [1] + K_{33} | y_R | [y_C] \end{bmatrix} \begin{bmatrix} M_C \end{bmatrix} \\ & H_{22} = H_{33} = \begin{bmatrix} \mathbf{I} \end{bmatrix} + \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_R] + K_{22} | x_R | [x_R] + K_{13} | y_R | [1] + K_{33} | y_R | [y_R] \end{bmatrix} \begin{bmatrix} M_R \end{bmatrix} \\ & H_{23} = H_{32} = \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_R] - K_{22} | x_R | [x_R] + K_{13} | y_R | [x_R] \end{bmatrix} \end{bmatrix} \\ & H_{23} = H_{32} = \begin{bmatrix} K_{11} | 1 | [1] + K_{13} | 1 | [y_R] - K_{22} | x_R | [x_R] + K_{13} | y_R | [x_R] \end{bmatrix} \end{bmatrix} \end{aligned}$$

If the frequency equation (eq. (A13)) is used, both symmetrical and antisymmetrical modes are obtained. However, if the symmetrical and antisymmetrical vibrations are considered separately, the order of the frequency matrix can be considerably reduced.

Symmetrical modes. - For the symmetrical structure vibrating in a symmetrical mode,

$$|\mathbf{w}_{R}| = |\mathbf{w}_{L}|$$

Thus, only the deflections at the center line and on the right-hand side of the structure need to be considered and the frequency equation

(eq. (Al3)) reduces to

$$\begin{vmatrix} w_{C} \\ w_{R} \end{vmatrix} = \begin{bmatrix} H_{11} & H_{12} + H_{13} \\ H_{21} & H_{22} + H_{23} \end{bmatrix} \begin{bmatrix} 2\delta_{11} & \delta_{12} + \delta_{13} \\ 2\delta_{21} & \delta_{22} + \delta_{23} \end{bmatrix} \begin{bmatrix} M_{C} \\ 2 \end{bmatrix} = 0 \quad |w_{C}|$$

$$\begin{vmatrix} w_{C} \\ 0 \\ 0 \\ M_{R} \end{vmatrix} = |w_{R}|$$
(A22)

or

$$|\mathbf{w}| = \omega^2 \left[[\mathbf{I}] + [\mathbf{F}^{\mathbf{S}}] [\mathbf{M}^{\mathbf{S}}] [\mathbf{\delta}^{\mathbf{S}}] [\mathbf{M}^{\mathbf{S}}] |\mathbf{w}| \right]$$
 (A23)

where

$$\begin{bmatrix} \mathbf{F}^{\mathbf{E}} \end{bmatrix} = 2 \left[\mathbf{K}_{11} \middle| 1 \middle| 1 \middle| + \mathbf{K}_{13} \middle| 1 \middle| \mathbf{y} \middle| + \mathbf{K}_{13} \middle| \mathbf{y} \middle| 1 \middle| 1 \middle| + \mathbf{K}_{33} \middle| \mathbf{y} \middle| \mathbf{y} \right]$$
(A24)

Note that in equation (A21) only the properties of the stations on the center line and on the right-hand side of the structure are involved. Also note that the mass associated with the center-line stations in the $\begin{bmatrix} M^S \end{bmatrix}$ matrix is one-half of the total assigned mass. The matrix $\begin{bmatrix} \delta^S \end{bmatrix}$ is the influence coefficient of the structure on a three-point support under a symmetrical loading. When the coefficients K_{11} , K_{13} , and K_{33} as shown in equations (A7), (A8), and (A9) are calculated, the $\begin{bmatrix} M^S \end{bmatrix}$ matrix can be used instead of the total $\begin{bmatrix} M \end{bmatrix}$ matrix. In this case,

$$M = 2 \left[1 \right] \left[M^{S} \right] \left[1 \right]$$

$$M_{y} = 2 \left[y \right] \left[M^{S} \right] \left[1 \right]$$

$$M_{yy} = 2 \left[y \right] \left[M^{S} \right] \left[y \right]$$

$$M_{xx} = 2 \left[x \right] \left[M^{S} \right] \left[x \right]$$

$$M_{x} = M_{xy} = 0$$

$$(A25)$$

Antisymmetrical modes. - For a symmetrical structure vibrating in an antisymmetrical mode,

$$|\mathbf{w}_{\mathbf{C}}| = 0$$

and

$$|\mathbf{w}_{R}| = -|\mathbf{w}_{L}|$$

Thus, only the deflections on one side need to be considered. For this case, the frequency equation (Al5) reduces to

$$\left|\mathbf{w}_{R}\right| = \omega^{2} \left[\mathbf{H}_{22} - \mathbf{H}_{23}\right] \left[\delta_{22} - \delta_{23}\right] \left[\mathbf{M}_{R}\right] \left|\mathbf{w}_{R}\right| \tag{A26}$$

or

$$|\mathbf{w}| = \omega^2 \left[\mathbf{I} \right] + 2K_{22} |\mathbf{x}| |\mathbf{x}| |\mathbf{x}| |\mathbf{x}|$$
 (A27)

Note that in this equation only the properties of the stations on one side of the center line are involved. The influence-coefficient matrix $\left[\delta^a\right]$ is the influence coefficient matrix of the structure on a three-point support under an antisymmetrical loading.

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TABLE I.- MOMENTS OF INERTIA OF IDEALIZED SPARS

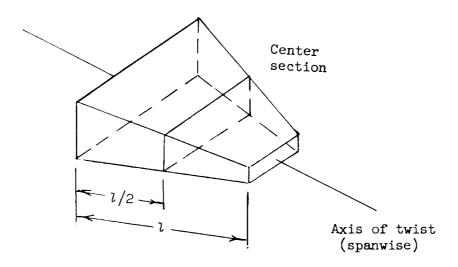
	Moments of inertia of -						
х	Spar 1	Spar 2	Spar 3	Spar 4	Spar 5	Swept leading edge	
0 16 28 40 52 64 76 88 100 112	26.123 26.123 21.748 17.785 14.230 11.080 8.331 5.980 4.022 2.456	43.546 43.546 36.374 29.858 23.995 18.781 14.213 10.287	45.643 45.643 31.131 31.306 25.163 19.700	40.300 40.300 33.720 27.738	20.635 20.635	8.624 7.119 5.785 4.599 3.559 2.659 1.897 1.268	

TABLE II.- MOMENT OF INERTIAS OF IDEALIZED RIBS

Rib	I (*)
123456789	11.588 15.647 11.550 9.402 7.509 5.816 4.369 3.115 1.255

*Ribs are assumed to be uniform.

TABLE III.- TORSIONAL CONSTANT OF TORQUE BOXES



Torque box	J
1, 3, 7, and 13 2 4, 8, and 14 5, 9, and 15 6 10 and 16 11 and 17 12 18 19 20	98.517 67.674 90.232 74.750 45.544 60.726 48.154 27.782 37.043 27.385 14.387

TABLE IV.- INFLUENCE COEFFICIENTS FOR SPAR 1 AS CANTILEVER BEAM

$$\left[\left[w \mid = \frac{1}{E} \left[\delta_{S} \right] \mid P \mid \right]$$

(a) Neglecting the effects of transverse shear

 $\left[\delta_{\rm S}\right]$

Station	22	23	24	25	26	27	28	29	30
22	52.266	111.064	169.863	228.662	287.461	346.259	405.058	463.857	522.656
23	111.064	281.220	463.508	645.797	828.086	1,010.375	1,192.664	1,374.953	1,557.242
24	169.863	463.508	834.758	1,220.725	1,606.692	1,992.659	2,378.626	2,764.593	3,150.560
25	228.662	645.797	1,220.725	1,890.308	2,578.308	3,266.207	3,954.107	4,642.006	5,329.905
26	287.461	828.086	1,606.692	2,5 7 8.308	3,668.163	4,781.134	5,894.105	7,007.076	8,120.047
27	346.259	1,010.375	1,992.659	3,266 .2 07	4,781.134	6,447.672	8,144.491	9,841.309	11,538.128
28	405.058	1,192.664	2,378.626	3,954.107	5,894.105	8,144.491	10,596.220	13,089.314	15,582.408
29	463.857	1,374.953	2,764.593	4,642.006	7,007.076	9,841.309	13,089.314	16,617.617	20,205.800
30	522.656	1,557.242	3,150.560	5,329.905	8,120.047	11,538.128	15,582.408	20,205.800	25,246.469

(b) Including the effects of transverse shear

 $\left[\delta_{\mathbb{S}}\right]$

Station	22	23	24	25	26	27	28	29	30
22	175.545	234.344	293.143	351.942	410.740	469.539	528.338	587.137	645.935
23	234.344	492.616	665.770	838.925	1,012.080	1,185.235	1,358.389	1,531.544	1,704.699
24	293.143	665.770	1,125.587	1,492.177	1,858.767	2,225.358	2,591.978	2,958.538	3,325.128
25	351.942	838.925	1,492.177	2,251.988	2,908.724	3,565.461	4,222.197	4,878.933	5,535.670
26	410.740	1,012.080	1,858.767	2,908.724	4,091.808	5,159.590	6,227.372	7,295.155	9,362.937
27	469.539	1,185.235	2,225.358	3,565.461	5,159.590	6,924.698	8,558.872	10,193.046	11,827.220
28	528.338	1,358.389	2,591.978	4,222.197	6,227.372	8,558.872	11,117.942	13,525.292	15,932.642
29	587.137	1,531.544	2,958.538	4,878.933	7,295.155	10,193.046	13,525.292	17,175.353	20,644.405
30	645.935	1,704.699	3,325.128	5,535.670	9,362.937	11,827.220	15,932.642	20,644.405	25,831.532

TABLE V.- SITFFNESS COEFFICIENTS FOR SPAR 1 UNDER SYMMETRICAL LOADING

_	1	Y	

_											$\overline{}$
	30	-0.000000933	949400000.	.000016929	945750000.	00019150	.00062716	00201,58	,0062964	0083585	.00359591
	29	-0.00056916 0.00013281 -0.00003010 0.000005979	.000032568	.00011823	0040124	.00013349	0043710	.0140343	0271777	.0248476	0083585
	28	-0.00003010	.00015617	0056472	.0019138	0063648	.0208393	0427307	,0476624	0271777	.0062964
	27	0.00013281	9689000.	.00249059	0084332	.0280409	0587813	.0679500	0427307	.0140345	.000627160020138
	26	-0.00056916	.0029607	010692	.036196	0776837	.0908740	0587813	.0208393	0043710	
	25	0.0024012	.012489	.045101	09788	650711.	0770837	.0280409	8495900-	.0013349	.00005754600019150
	77	0.0999941	.0519875	119798	.146262	097788	961960.	0084552	.0019138	00040124	
	23	0.0411729	131713	.173902	119798	.045101	010692	.00249059	00056472	.00011823	000016929
	22	-0.0822049	.172020	131713	.0519875	012489	.0029607	9689000:-	.00015617	000032568	000000933 .0000004646000016929
•	21	0.04908861 -0.0822049	0822049	.0411729	00999641	.0024012	00056916	.00013281	00003010	.000005979	000000933
	Station	21	22	23	77	25	56	27	28	59	30

TABLE VI.- STIFFNESS COEFFICIENTS FOR RIB 4

 $\left[\Delta_{
m R}^{
m l_4}
ight]$

Station	6	10	16	24
6	0.00108821	-0.0024847	0.00163231	-0.00027205
10	00244847	.00652924	00571307	.00163231
16	.00163231	00571308	.00652924	00244847
24	00027205	,00163231	00244847	.00108821

TABLE VII. - STIFFNESS COEFFICIENTS FOR TORQUE BOXES 15 AND 16

r L	נ	7	1	E-	
۰					

^16]	

Station	16	17	1 7.	25
54	0.004055431	0.004055431	0.004055431	0.004055431
23	0.004055431	0.004055431	0.004055431	0.004055431
16	0.004055431 0.004055431 0.004055431 0.004055431	0.004055431 0.004055431 0.004055431 0.004055431	0.004055431 0.004055431 0.004055431 0.004055431	0.004055431 0.004055431 0.004055431 0.004055431
15	0.004055431	0.004055431	0.004055431	0.004055431
Station	15	91	23	1 7

 Station	91	17	24	25
16	0.003294499	0.003294499 0.003294499 0.003294499 0.003294499	0.003294499	0.003294499
 17	0.003294499	0.003294499 0.003294499 0.003294499 0.003294499	0.003294499	0.003294499
†₹.	0.003294499	0.003294499 0.003294499 0.003294499 0.003294499	0.003294499	0.003294499
25	0.003294499	0.003294499 0.003294499 0.003294499 0.003294499	0.003294499	0.003294499

TABLE VIII. - ELEMENTS OF ROW 24 OF STIFFNESS COEFFICIENT OF DELTA WING UNDER SYMMETRICAL LOADING

$$P_{24} = E \left[\Delta_{24,n} \right] \left[w_i \right]$$

$$\left[\Delta_{24,n} \right]$$

n	Δ _{24,n}	n	∆ ₂₄ ,n
1 2 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17	0 0 0 0 0 0 0 00027205 0 0 .00163231 0 0 0 0 .004055431 009798400 .003294499	18 19 20 21 22 23 24 25 26 27 28 29 30 31 32 33	0 0 0 00999641 .0519875 1238534 .15470014 10108250 .036196 0084332 .019138 00040124 .00005746 0

TABLE IX.- STIFFNESS MATRIX FOR

(a) Cymmetrical deflections;

Sta- tion	2	,	4	5	6	7	8	9	10	11	12	13	14	15	16	3.1
2	0.02682143	0.0040087	-0.01180341	0.00734502	0.00363055	0	0.002910508	0	n	0	0.000202307	0	-0.000727625		1	
1 3 1	.0040067	.09182985	1358191	0595852	00933340	01191208			lo	ě	0	.00323320		10	'n	0
l úi	01180341	1558191	.28773245	1950508	.04860282	.00400870		.004895355	lo .	o .	0	0	.00456574	10	ñ	0
1 5 1	.00734302	.0595852	- 1950508	.2579946	07429795	0	.004895555	02154473	.004055451	o .	0	0	0	.00440182	0	n
6		00955540	.04860282	07429795	.05502345	o .	0	.004055431	008974781	.004941755	.004113883	c i	Ď	0	.00165251	l n
1 71		01191208	.00400870		0	.10530996	1517220	.0720054	0174452	.0039563	00061045		.00400870	0	0	ő
8	.002910508	.00400870	01957586	.004895355	o .	1517220	.3341983	24055251	.0907950	- 0205956	.00517793	.00400870	- 01957586	.00-895555	lo .	ŏ
9	c	a	.004895355	02154475	.004055451	.0720054	24055251	. 5551829	2171529	.0744876	0114950	0	.004895355	01714503	.004055431	o .
10	ō l	ō l	0	.004055453		0174432	.0907950	2171329	.2728925	1640481	.0389947	0	0	.004055451	0150b501	.003294499
	č l	6	0	G	.004941755		0205936	.0744876	- 1540481	.18031633	05584437	o l	0	0	.003294499	
12	.000202307	0	0	0	.00+115885	00061043	.00517795	0114950	.0389947	- 05584437	.05825592	lo l	0	l o	0	.002612571
13	0	.0032552	c l	0	0	01191208	.00400870	0	0	0	0	.09813152	- 1451224	.0687170	0167129	.00402300
	000727625		.00436574	0	e	.00400870	01957586	.004895355	0	0	0	1451224	. 5160310	- 2298577	0869980	0209390
15	0	0	0	.0044018	0	0	.004895555	01714303	.004055451	0	0	.0687170	2298377	31716106	2088919	.07573110
16	0	o	c l	0	.00163231	0	0	.004055451	- 01506501	.005294499	0	0167129	.0869980	2088919	2668151	1710800
17	0	0	0	0	0	0	0	0	.003294499	01331599	.002612571	.00402500	0209390	.0757311	1710800	.21488768
16	o i	0 :	0	0	0	0	0	0	0	.002612571	005381887	00095210	.0049575	0179286	.0608116	13-71724
19	0 !	0 .	0	0	0	0	0	0	0	0	.003014435	.00021018	00109551	.00596083		0-1759
20	.000010606	0 ,	0	0	.000216011	0	0	0	0	0	00234956	00005167	.000165125	000597105	00202525	00b7+ +
21	0	000538868	0	0	0	.00215547	0	0	0	0	0	00706228	.00400870	0	0	D
25	0 :	0	0	000753628	G	0	0	.00220089	0	0	0	0 .	.004895555	- 01200758	.004055451	ō
24	0	0	0	0	000272050	0	0	0	.00163231	0	0	0	0	.004055431	009798400	.003299499
29	0	0	0	0	0	0	0	0	0	,00177766	0	o !	C C	0	.005294499	
26	0	0	0	0	0	0	0	0	0	0	.00063105	0 !	o	0 .	0	.002612571
57	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0	0
28	0 .	0	0	0	0	0	0	0	0	0	0	0	0	0 :	0	0
29	0 .	0	0	0	0	0	0	0	0	0	0	0 '	0	0	0	0
30	.000000512	0	0	0	.000006587	0	0	0	0	0	.000069580		0	0	o	0
1 31 .	000002182	0	0	0	000044994		0	0	0	0	000489792		0	0	C	o
32	000046755	٥ -	0	0	000955418	0	0	0	0	0	00966743		0	0 ;	c	0
33 :	000856076	0	0	0	01605958	0	0	0	.004941755	01758671	006775619		0	0	0	.00355532
36	01248979	0	.00734502	02666847	01228-25	0	0	.005868998	0	0	001135430	0	0	00146/28	0	0

WING ON THREE-POINT CUPPORT

transverse stear neglested

18	19	20	21	23	24	25	26	27	26	29	5 0	31	52	35	34
	1	0.000010000			,	0	0	0	0	0	0.000000312	-0.000002182	-0.000046755	-0.000856076	0.01248979
100	10	0.000000	000538868	o .	0	o o	e .	ō	ō	o	0	0	0	0	0
l n	6	ñ	0	0	i .	ā	0	ō	0	0	0	0	0	0	.00754502
10		lő.	n	000755628	ő	o .	0	ō	6	٥	0	0	0	0	- 02666847
l n	ŏ	.000216011	n	0	000272050	0	0	0	0	a	.000006587	0000044994	000/55418	- 01605958	01228425
n n	ŏ	0	.00215547	ō	0	0	0	0	o	0	0	0	0	0	0
lő.	č	0	0	ō	ō i	ō	0	0	0	0	0	0	0	0	0
0	·ě	lõ	io ·	.002200F9	0	o .	0	0	0	0	0	0	0	0	.005868998
lō.	o	lo	o :	0	.00163251	0	0	0	0	0	0	0	0	.004941755	0
.002612	971 O	0	0	0	0	.00177760	0	0	0	0	0	0	0	01758671	10
005581	887 .00 5014435	.00234956	0	0	0	0	.00063105	0	0	0	.000069580	000489792	009607-3	008775619	001135430
0007/2		00005167	00706228	0	0	0	0	0	0	0	0	0	0	0	0 1
.004957	5 :00109551	.000169125	.00400870	.004895355	0	0	0	0	0	O	0	0	0	0	0
017928		000597105	0	01200758	.004055451		0	0	0	0	0	0	0	0	00146728
.060811			0	.004055431	00979840	.005294435		0	0	0	0	10	ļ ú		10
154717	14 .0447695		0	0	.003294499	007981009			0	0	:0	0	000001117	.00555552	10
.162512	18 1-10/582706	.0221608	0	0	0	.002612571				0	0	. 0	003014455		0
095827			0	0	0	0	.002009622	005591406	.001485664		0	0	009821085	000719575	0 000000000
.022160	805144626	.02058654	0	0	0	0	0	.001485664	002266188			00\582569		000(195)5	000059597
0	0	0	.05444448	.0411729	009996-1	.0024012	00056916	.00015261	00005010	.000005979	000000955	C	10	0	.000244542
0	0	0	.04117290	.184197771	12385543	.04)1010	0106920	.00249059	00056477	.00011825	000016929	1 0	10		
0	10	0	00999641	12585343	.15470014	10108250	.0361960	00845520	.00191380	00040124	.000057546	ů,	10	000592552	
.002612		0	.0024012	0411010	10108250	.1258549	07919627	.0280409	0063648	.0015549	00019150 .00062716	1 %	10	0	I,
	295 .002009622		00056916	0106920	.0561960	07969627	.09612724	06079092	.0208393	0045710	00201380	0	.00126408	Š	l č
.002009	6221005591406	.001+85664	.00013281	00249059	00845520	.0280409	06079092	.07207752	04421636	02873875	00201300	.001561045	0	ŏ	lŏ.
0	001+856b+	002266188	00005010	00056417	.0019156	0063648	.0208595	04421636	02875875	.02796970	00835650	003122095	l n	ň	Ič l
0	0	.001561045	.000005979	.00011825	- 00040124	.0015549	0045710	.0140343	.00629640	00855850	.00599202	+.000921125	00022262	0000212687	0000017616
0	0		000000955	000016929	.000057546	00019150	.00062716	00201380	.00029640			.005871415	.001567973	.000149948	.0000123806
0	0	00+5825tip		0	0	O .	10	.00126408	Characton.	2.00,122090	000921129	.001567973	.01622578	.00318224	.000265681
.003014	435009821085	-,00480245		0	0	0	10	.00126408	lo O	10	000222627	.000149948	.00118224	.02797735	000205001
0	0	000719573		0	0	000592552	l ^u	10	10	10	0000017616	.0000123806		004804766	055575975
0	0	.000059597	0	140445000	0	0	lc	lo.	U	10					1.033313413

TABLE IX. - STIFFNESS MATRIX FOR WING

(b) Antisymmetrical deflections;

Sta		4	5	6	8	9	10	11	1.2	14	15	16	17	18
24 55 66 8 9 10 11 12 14 15 16 17 18 19 20 25 26 27 28 29 20 31 20 20 20 20 20 20 20 20 20 20 20 20 20	0.01170746 01180541 .00754502 .00365035 .002910508 0 0 .0002025067 000727625	-0.01180/41 1.17812116 -1.17967939 -0.190779 -0.004895795 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0		0.00365035 0.430679 072513631 .05474396 .004955431 008974791 .004941755 .0041158834 0 .00165251 0	0.002910708 -0.0977695 -0.0977695 -0.04895959 -2.0486821 -0.04829596 -0.0482951 -0.0487595 -0.0487595 -0.0977695 -0.0977695 -0.0977695 -0.0977695 -0.0977695 -0.0977695	0 .004895355 021544726 .004055426 19741071 .32050486 21557686 .075681 01157046 .004895355	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0.0002023067 0.0041136834 .00281407 01137046 .0389646 055837571 .03823282	-0.00727625 .0045674 0 -0.019577855 .004895355 0 -1.186260 -1.18686471 .0770395 .0043917 -0.0014673 .004995355	0 .00440182 0 .004895355 -017143026 .004055431 0 0 -18886471 .303146662 -205483865 .0749101 -0177555 .0059184	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0

ON THREE-POINT SUPPORT - Concluded

transverse shear neglected

19	20	25	51*	25	26	27	28	29	30	51	32	55	34
)	0.000010608	0	c	0	0	0	0	0	0.000000312	-0.000002182	-0.000046755	-0.000856076	-0.01248979
3	0	0	0	0	0	0	0	0	0	0	0	0	.00754302
)	0	000735628	9	0	0	0	0	0	0	0	0	0	02666848
)	.0002160107	0	00027205	0	0	0	0	0	.000006387	0000449943	000955418	01605958	01228425
)	0	0	0	0	0	0	0	0	0	0	0	0	0
) .	0.	.00220089	0	0	0	0	0	0	10	0	0	0	00586899
)	0	0	.00163231	0	0	0	0	0	0	0	0	.004941755	0
)	C	0	0	.00177766	0	0	0	0	0	-0	0	017586705	0
.005014455	.002549564	0	0	10	.00063105	0	0	0	.000069580	0004897922	00966743	008775619	- 00113545
.000971205	.000146575	.004895355	0	0	C	0	0	0	0	0	0	0	0
.0059184	00059076	012007576	.004055431		0	0	0	0	0	0	0	0	00146726
0134257	.00202371	.004055431	009798400		0	0	0	0	0	0	0	0	0
.0447670	00674923	0	.003294499			0	0	0	0	0	0	.00355532	0
.09582646	.0221607	0	0	.002612571	005884292	.0020096217		ů .	0	0	.005014435		0
.10156481	051446264	0	0	0	.0020096217				0	0	009821085		0
.031446264	.020586341	0	0	0	0	.001485664	002266188		.000694609	+.004582568	004802426	000719573	- 00005959
)	0	.17581577	12181845	.0446126	0105760	.00246316	000558574		000016596	0	0	0	.00021455
)	0	12181845	.15420614	~.10096400	.0361678	00842654	.00191066	00040115	.000057465	0	0	0	0
) .	0	.0446126	10096400	.12580690	079689471	.0280393	0063644	.00153487	00019148	0	0	000592552	0
.0020096217		0105760	0561678	079689471	.09612564	06079052	.0208392	00437099	.00062716	0	0	0	0
.005391406	.001485664	.00246316	00842654	.0280393	06079052	.072077226	044216369		0020138	0	.00126408	0	0
.001485664	002266188	000558574	.00191066	0063644	.0208392	044216369	.049928588		.00629641	.001561045	0	0	0
	.0015610450	.00011776	00040113	.00133487	00437099	.0140345	028738745		0085585	~.0031220951	C	. 0	2000001
	.000694609	000016596	.000057465	00019148	.00062716	0020138	.00629641	0083585	.003992017	0009211247	000222820	0000212887	00000176
	- 004582569	0	0	0	0	0	.001561045	0031220951	0009211247	.00587142	.001567975	.000149948	.00001238
.009821085		0	0	0	0	.00126408	0	0	000222820	001567973	.016225776	.00518224	.00026560
	000719573		0	000592552	0	0	0	0	000021288	.000149948	.00318224	.02797735	
	0000595975	.000244542	0	0	0	0	0 .	0	0000017616	.0000123806	.000263681	.004804766	.03557398

TABLE X.- INFLUENCE COEFFICIENT MATRIX

(a) Symmetrical loading;

Sta- tion	2	3	l ₄	5	6	7	8	9	10	11	12	13	14	15	16	17
2	51.840556	25.292260	29.565920	35.316710	39.978140	19.903370	20.465700	21.979080	24.066920	26.079420	27.874650	10.371690	10.662470	11.151180	11.880520	12.859990
				225.04320				254.96158	241.53798	232.01163	225.45891	164.56462	163.44822	161.46636	158.78607	155.24781
			244.74328	253,15857			274.16204	271.45005	268.34342	268.80261	271.56951	167.60093	168.29999	170.05828	172.81753	175.78873
. 5!	35.316710		255.15837	301.02180		276.08520	286.89350	301.63130	319.61840	341.12420	363.51070	173.12440	177.60530	186.76790	200.24900	216.25470
6			262.86324	349.99570	464.82890	282.61720	302.58830	338.78200	386.00540	438.00640	488.87480	179.68910	189.38920		236.80540	270.80980
7			275.27832	278.08520	282.61720	400.21242	380.98263	356.70920	333.69741	316.57739	303.34514	262.04986	255.15322		232.30749	219.90623
			274.16204	286.89350	502.58830	380.98263	379.13397	368.82967	356.48657	349.05079	345.38343	254.06752		250.97131	248.46662	245.69783
			271 5005	301.63130	338.78200	356.70920	368.82967	386.51508	399.42266	413.65496			248.39250	262.01677	278.84214	296.30163
			268.34342	319.61840	386.00540	333.69741	356.48657	399.42266	455.56540	511.50330	567.69530			276.02460	521 20280	371.41560
			268.80261	341.12420	438.00640	316.57739	349.05079	413.65496	511.50330	635.76580	761.92350	210.82960		293.53450	372.83110	467.89410
			271.56951	363.51070	488.87480	303.34514	345.38543		567.69530	761.92350	996.95490	197.80460	236.17010	313.35490	129 32630	576.26070
13			167.60093	173.12440		262.04986	254.06752	240.75559	225.14600	210.82960	197.80460		209.72454	186.17612	164.32654	145.25174
			168.29999	177.60550		255.15522	252.97341	248.39250	242.68510	238.66860			208.85389	200.62860	190.55103	182.10886
	11.880520		170.05828 172.81753	186.76790	208.71050 236.80540	244.46439	250.97131 248.46662	262.01677	276.02460	293.53450	313.35490	186.17612	200.62860	222.92021	240.43412	255.52490
				216.25470		252.30749 219.90625	245.69783		321.20280	372.83110	129.32630	164.52654	190.55103	240.43412	304.29198	362.89558
	15.982550	151.74494	179.24462	233.83320			243.89594	296.30163 514.55785	371 41560 423 54160	\$67.89\$10 570.\$1270	576.26070	145.25174	182.10886	255.52490	362.89558	488.94566
			184.68959	255.22980			245.85556		479.52110	678.14670	742.74980 918.40940	126.92425 116.22396	176.28969	272.30171	419.03029	610.42877
				273.48250			250.38445		537.54060	787.05970	1,094,0155		174.33366 174.67950	295.22524	478.86119 542.12197	731.90409 857.10929
21	- 07296640			-5.8257790	-8.3152540	-3.8222471	-5.8823531	-10.127744	-16.442398						-28.981237	-42.554224
23			10.690521	14.348473		12.940574	17.179794		37.950270	53.822627			15.1207210		63.293266	91.419690
24		20.378383	24.953327	34.227550		29.228680	39.685375		91.042460	130.83680		9.0331190		82.388967		219.76164
25		51.571625	59.857275	56.602570	79.98220	43.866486		97 859950	152.01413	223.05955			51.792825	127,53538	236.35265	367.11114
			53.922929	79.759030	115.77058	54.995079	81.566313	134.81631	216.25918	324.71777	453.24553	14.441825	65.851401	168.74845	322.59439	518.77975
			66.449152	102.67839			97.686169	169.40080	280.51503	430.88413	611.19202	12.676594	76.394282	205.55581	403.71972	667.05769
			77.711786	125.16849			111.13500			538.88300	775.10536		84.456357	239.19491	480.78051	811.04487
29			86.706978			70.153652			407.50780	647.87559	941.53495		91.881388	272.02259	556.89418	954.43266
			99.703880			73.740310	136.76250		471.14950	757.09070	1,107.9447	13972000	99.266590	504.79918	632.97939	1,097.9474
			145.60106	221.96440	327.59470	134.95946			504.84210	773.14900	1,102.7935			310.67067	587.49599	978.21725
			235.99242						564.54620	791.76740	1,071.9111	157.87029		322.19852	495.88594	726.32478
			285.01765						507.40530	631.45210				274.29645	340.48261	421.36746
)4 <u> </u>	*). XXXXX	142.01974	171.27049	222.96610	279.80200	172.68259	181.55129	197.26304	217.33680	258.39220	258.73280	106.25322	110.01762	117.61472	128.65921	142.13633

FOR WING ON THREE-POINT SUPPORT

transverse shear neglected

18	19	20	21	25	24	25	26	27	28	29	30	31	32	33	34
15.982550	15,076700	16.109680	-0.07296640	0.17780080	0.47206800	0.88695100	1.4306880	2.1018470	2.8628750	5.6482160	4.4353400		21.931720	33.878620	45.386280
151.74494	150.32964	150.46747	-3.1444147	8.8867806	20.378383	31.571625	41.110207	48.470065		59.464670	64.753420	107.56165	192.34443	280.05895	142.07974
179.24462	184.68959	191.46258	-4.0419050		24.953327	39.857275	55.922929		77.711786	88.706978	99.703880		235.99242	285.01765	171.27049
253.85520	253.22980	273.48250	-5.8257790				79.759030	102.67839		147.64494	170.15070	221.96440	323.11560	378.63170	222.96610
508.14590	347.11230	386.25640	-8.5152540	19.454219	47.194640		115.77058	153.14557		229.74819	268.28680	327.59470	442.53220	505.83910	279.80200
208.69561	201.29282	196.49111	-3.8222471		29.228680		54.995079		66.530017	70.153652	73.740310	134.95946	256.92773	316.53644	172.68259
243.89594		250.38445	-5.8823531	17.179794			81.566515		111.13500	123.95811	136.76250	193.50242	505.52020	348.76665	181.55129
514.55785		360.90139	-10.127744	25.555828			134 81631		201.85722	233.91545	265.99583	313.57655	405.02772	412.07144	197.26304
423.54160		537.54060	-16.442398		91.042460		216.23918		343.96616	407.50780	471.14950	504.84210	564.54620	507.40550	217.55680
570.41270		787.05970		55.822627		223.05955	324 72777		538.88300	647.87559			791.76740	651.45210	238.39220
742.74980		1,094.0155	-33.371936	71.960922	176.85771		453.24553		775.10536	941.33495		1,102.7935			258.75280
128.92425	116.22396	105.59449	4.3631906				14.441825	12.676594	8.773488	4.3319400	13972000		157.87029	206.51744	106.25322
176.28969		174.67950	-3.5299065	15.1207210		51.792825	65.851401		84.456357	91.881388	99.266390		211.86575	229.08259	110.01762
272.30171		316.92398	-15 510915		82.588967	127.53538	168.74845		239.19491	272.02239		510.67067	522.19852	274.29645	117.61472
419.03029		542.12197	-28.981237	63.293266		236.35265	322 59439	403.71972	480.78051			587.49599	493.88594	340.48261	128.65921
610.42877		857.10929	-42.554224			367.11114	518.77975	667.05769			1,097.9474		726.32478	421.36746	142.15633
829.85701	1,050.7351		-55.921498	119.61066		508.42847	744 88680	687.92661	1,230.3223	1,474.2653	1,718.8660	1,499.2098	1,010.6292	510.42575	157.05819
1,050.7351				147.71831		654.54554	988.98264	1,354.9392	1,735.5502	2,125.0120	2,516.6172	2,170.0146	1,365.9733	601.87258	172.72000
1,273.2953			-82.424203	176.03564		803.82472	1,242.9649	1,749.6850	2,303.5498	2,882.7581	3,467.1308	2,967.7070	1,732.3414	692.86161	188.54721
-55.921498		-82.424203		-28.646445	-49.925225	-67.311928	-82.651004	-96.965158	-110.80212	-124.41642	-137.98376	-110.01505		-10.526759	-3.4812120
119.61066	147.71831	176.03564	-28.646445			143.24589	176.85787	207.83517	257.59495	266.79144	295.87045		120.53467	40.653402	8.5884820
294.36837	369.65915	445.86586	-49.925225				439.21412		605.45558	684.71425	763.59274		500.77097	99.250677	20.289094
508.42847		803.82472	-67.311928	143.24589		568.29663	764.81350	938.36308	1,101.2755	1,259.4146	1,416.5901	1,106.0620	552.42970	170.15355	34.145660
744.88680		1,242.9649		176.85787		764.81350	1,114.4713	1,432.2505	1,723.4281	2,003.0316	2,200.3699	1,751.7850	805.87976	249.08669	49.079710
687.92661		1,749 6830		207.83517	524.35111	938.36308	1,432.2505	1,963.1403	2,460.5436	2,930.4702	3,395.2474	2,549.7758	1,108.7179	352.02777	64.530150
1,230.3223	1,735.5302	2,303.5498	-110.80212	257 59495	605.45558	1,101.2755	1,723.4281	2,460.5436	3,266.4371	4,052.6611	4,627.6153	3,515.6047	1,426.3572	419.09675	80.200630
	2,125.0120				684.71425	1,259.4146	2,003.0516	2,930.4702	4,052.6611	5,541.8541	0,004.9163	4,657.7507	1,702.4211	500.36590	95.986690
	2,516.6172			295.87045	763.59274	1,416.5901	2,260.3699	5,595.2474	4,827.6153	0,004.9163	0,000.0505	5,880.4965	2,076.6729	295.79002	111.79745
	2,170.0146			255.46637			1,751.7830	2,549.7758	5,515.6047	4,657.7507	5,880.4965	4,519.7728	1,906.4817	044.13075	150.29541
	1,365.9735			120.33467			805 87976	1,108.7179	1,420.3372	11,752.4211	2,076.8729	1,906.4817	1,400.7911	755.8056	225.7527
510.42373	601.87258		-18.326739							506.36598		644.15873		706.99680	279.00310
157.03819	172.72000	188.54721	-3.4812120	8.3884820	20.289094	34.145660	49.079710	64.550150	100.200630	95.986690	TTT 79743	150.29541	225.7527	279.00310	215.88125

TABLE X.- INFLUENCE COEFFICIENT MATRIX FOR

(b) Antisymmetrical loading;

Sta- tion	2	4	5	6	8	9	10	11	12	14	15	16	17	18
2	575.7224	360.5361	601.7281	823.7979	232.9755	404.9876	571.0150	727.6955	876.5052	124.9455	226.5982	337.1404	454.6444	575.4706
14	360.5361	309.3417	518.1413	713.5195	220.3955	379.0422	529.5642	672.5078	809.4466	120.0355	217.4012	322.5558	433.2395	546.4655
5	601.7281	518.1413	894.3876	1250.551	379.7971	657.7621	924.7133	1179.133	1422.640	209.3141	379.4903	563.9921	758.9652	958.9030
6									2041.955					
- 8	232.9755	220.3955	379.7971	533.7412	194.2970	323.1545	441.2496	553.6178	662.5724	111.8606	198.2547	287.5461	379.1012	471.9543
9														835.8900
10	571.0150													
11	727.6955	672.5078	1179 133	1687.412	553.6178	976.0771	1408.627	1846.530	2271.723	334.4817	614.9803	928.2693	1265.833	1615.077
	876.5052	809.4466	1422.640	2041.955	662.5724	1172.523	1705.653	2271.723	2861.562	402.0729	744.8913	1136.470	1569.478	2027.530
14									402.0729					
15									744.8913					
16									1136.470					
17	454.6444	433.2395	758.9652	1085.073	379.1012	668.1963	693.9362	1265 833	1569.478	249.9188	470.0499	727.0367	1007.767	1286.898
18	575.4706	546.4655	958.9030	1373.601	471.9543	835 8900	1215.762	1615.077	2027.530	503.3133	578.6888	907.0376	1286.898	1698.932
	695.8834													
	814.5225													
									153.8919					
	105.4341													
	191.0006													
	288.6308													
27									1577.201					
	503.3678													
	614.3577													
30									2893.949					
31									2930.782					
	851.0952													
														1717.064
34	711.2854	559.0436	964.5270	1351.587	393.6721	683.6136	963.9637	1229.640	1482.730	214.7360	389.4617	579.0498	780.0244	986.4768

WING ON THREE-POINT SUPPORT - Concluded

transverse shear neglected

19	20	23	24	25	26	27	28	29	30	31	32	33	34
695.8834 659.7765 1159.022 1662.447 565.7222 1005.146 1469.988 1968.729 2494.570 362.3366 690.0625 1090.039 1565.981 2113.006 3323.006	814.5225 771.9030 1356.951 1596.951 1597.951 1659.3562 1174.044 1723.534 2320.541 2958.426 419.9036 802.5235 1274.774 1347.050 2526.946 3323.006 4226.304 4203.940 4226.304	58.35874 38.24327 67.09824 96.04945 36.60115 64.49765 93.29599 123.1929 123.1929 123.8919 28.58607 56.24736 88.30693 122.2403 125.7459	105.4341 104.7382 183.7961 265.2276 99.01192 175.0706 254.2079 337.1822 423.1287 423.1287 423.255.4857 332.7353 433.3575 535.3506 638.0764 86.37670	191.0006 188.963 351.7017 475.2769 176.1990 312.5988 455.8918 607.7461 766.5208 128.0436 253.2921 410.7206 592.4898 1786.9185 987.2658 1190.529 126.1979	288.6308 284.2654 499.2516 715.7050 261.5968 465.4293 681.5948 913.2406 1157.603 185.0113 365.7954 597.1513 874.5669 1184.964 11514.964 1184.964	393.8827 386.1788 678.5074 973.3144 351.1174 626.0849 920.0474 1238.396 1243.1567 480.2736 786.7557 1163.579 1602.011 2088.708 2088.708 2662.685 199.5231	503.3678 491.5430 864.0397 1240.165 442.5008 790.3360 1164.566 1573.427 301.6006 595.1023 976.6871 1453.633 2024.719 235.0022 684.338 3406.315 235.0022	614 - 3577 598 - 2163 1051 - 910 1510 - 1435 534 - 7234 996 - 1780 1411 - 678 1912 - 1448 2492 - 1473 360 - 3014 710 - 3657 1167 - 366 1745 - 137 2451 - 149 3291 - 864 14238 - 620 270 - 2040	725.5679 705.0981 1240.148 1781.231 627.1087 1122.324 1659.256 2893.949 419.0725 829.7551 1558.153 2037.103 2879.251 3901.957 5076.664 305.3432 881.4466	771.1284 7739.5294 1300.366 1867.218 644.0469 1149.697 1693.665 2289.752 22930.782 419.8628 814.8597 1317.677 1944.427 2708.292 3624.717 4666.654 265.5966	851. 0952' 797. 5192 1401. 582 2011. 412 668. 5964 1136. 548 1735. 363 2356. 621 2295. 807 1416. 5422 785. 7281 1220. 577 1728. 534 2302. 508 2931. 885. 7365 5577. 888 188. 7365 555. 1686 555. 1686 555. 1686 555. 1686 555.	866. 6755 782. 3820 1374. 774 1972. 989 617. 3092 1085. 215. 1560. 231 2040. 695 2512. 023 359. 8214 658. 7492 990. 2836 1347. 053 1717. 064 2088. 546 2456. 087 124. 2525 340. 5937	711. 2854 559.0456 964. 5270 1351.587 393.6721 683.6156 963.9657 1229.640 1482.735 214.7360 389.4617 759.0244 986.4768 1192.692 1396.264 67.4678 1185.1267
987.2658 1514.956 2088.708 2684.388 3291.864 3901.957 3624.717 2931.885 2088.546	1190.529 1854.487 2602.685 3406.315 4238.620 5076.664 4666.654 3577.888	126.1979	344.2830 458.3183 566.7128 672.5222 777.1148 881.4462 758.7331	601.1648 841.3429 1063.548 1277.445 1487.578 1696.936 1440.519	841.3429 1262.039 1657.844 2030.309 2392.611 2752.909 2295.124 1475.596	1063.548 1657.844 12298.178 12909.246 3494.658 4075.260 3317.985 2033.552	1277.445 2030.309 2909.246 3862.151 4797.624 5722.244 4516.403 2617.092	1487.578 2392.611 3494.658 4797.624 6270.249 7777.237 5894.344 3208.527	2752.909 4075.260 5722.244 7777.237 10187.43 7353.460 3800.904	2295.124 3317.985 4516.403 5894.344 7353.460 6108.619 3694.117	964, 2659 1475, 596 2033, 552 2617, 092 3208, 527 3800, 904 3694, 117 3370, 296 2506, 623 1450, 230	926.1097 1259.034 1603.616 1952.572 2302.291 2382.563 2506.623	504.6002 687.0963 876.4354 1068.264 1260.470 1330.225 1450.230

TABLE XI. - INFLUENCE COEFFICIENT

(a) Symmetrical loading;

Sta- tion	5	3	4	5	6	7	8	9	10	п	12	15	14	15	16	17
2 5 4 5 6 7 8 9 10 11 12 13 14 15 16	120. 35205 36. 734840 76. 193610 85. 271920 89. 270480 47. 999140 51. 984090 61. 451610 20. 759760 25. 944300 27. 867660	299.57711 265.03733 253.70413 249.02304 309.26130 289.69646 275.00126 259.89854 258.69112 257.57209 184.41614 176.35587 173.08019 169.12306	76.193610 265.03733 307.53218 306.90595 313.11701 302.30288 313.11302 308.10967 301.7563 308.89947 514.96415 187.65303 189.50129 191.70571 196.78247	253.70413 306.90595 3594.57370 432.09930 305.14420 325.09880 358.14680 374.36200 401.45120 424.6240 424.6240 189.3950 197.72210 215.60830 231.52290	89.270480 249.02504 313.17751 432.09930 578.44860 395.76650 460.05440 521.76270 572.98700 195.20400 208.69710 239.95970 279.66130	37.010830 309.26130 302.30288 303.11420 305.92510 464.77454 410.34276 391.07269 355.56917 345.87803 355.56917 299.32651. 271.67894 258.11004 245.030306	47.999140 269.69646 513.11302 525.09680 538.42500 410.34276 427.94241 406.63526 389.61244 406.63526 389.002688 389.002688 389.002688 273.99946 276.81015 273.07936	51.984090 275.00126 308.10967 358.14680 395.76650 381.07269 406.63526 453.16613 452.46764	56.100670 259.89634 301.75963 374.36200 460.05440 355.56917 389.61244 452.46764 540.29083 587.96176 460.17002 239.25271 261.87873 310.76598 373.70774 423.99431	58.890890 258.69112 308.89947 401.45120 521.76270 345.67803 349.002688 469.15185 587.96176 767.88230 584.33938 229.50153 562.21265 530.61334 427.75825 549.15978	61.451610 257.37209 314.96415 424.62490 572.98700 357.51787	20.759760	23.944300 176.35587 187.06303 197.72210 208.69710 271.67894 276.81015 269.41681 261.87873 262.21265 263.04370 227.80503 229.09940 221.19739 208.65229	25. 726650 173. 08019 189. 50129 215. 60850 239. 95970 258. 11004 273. 07956 359. 01494 310. 76958 350. 61354 550. 552457 221. 19739 274. 54093 279. 644538	27,867660 169,12506 199,70571 231,52290 279,66150 269,48204 487,75825 314,45566 373,70774 427,75825 483,27980 178,28642 208,65229 279,64538 577,55620 421,55907	29, 9164 50 167, 80 992 2 196, 7824 7 290, 835 10 321, 53590 269, 14111 532, 43914 423, 984 31 549, 13978 660, 57080 160, 931 51 201, 53023 290, 97712 421, 55907 788, 56288
18 19 20 21 23 24 25 26 27 28 29 30 31 32 33 34	52.033350 33.810090 55.537950 9904708 .7265680 1.7270130 2.890104 4.242889 5.656023 7.159676 8.622690 10.085890 10.085890 22.81697 48.55806	165.66927 169.25542 172.92590 4.8940853 11.9955484 23.198359 55.649807 42.141754 50.296616 57.189570 65.197020 75.116620 123.32579 219.75846 267.97473	200.92146 210.25077 219.46150 1.3998475 14.858489 29.807105 44.552239 57.738189 70.771349 82.697917 95.642390 108.49461 164.29627 271.79719 350.01652	269.06370 290.86860 512.13520 -1.881605 21.941970 44.680310 10.26675 115.20675 137.92692 137.92692 144.76460 248.76860 447.06050	361, 97740 403,16620 435,03620 -5,9751180 29,528450 64,846700 102,741310 141,44839 180,88158 20,20388 259,80150 571,48290 512,05020 601,96600	225, 27392 224, 94650 225, 10725 15, 859502 17, 631738 33, 400701 46, 600757 56, 099550 64, 657478 79, 112510 87, 197980 156, 64091 288, 99371 342, 36185	268.07614 277:117727 266.45244 4.4048441 24.638295 48.286027 69.893969 88.079581 120.7526 137.96675 137.96675 155.02138 221.28888 221.28888 24.28888 25.28888 26.28888 26.28888 26.28888 26.28888 26.28888 26.28888 26.28888	349.04786 376.33156 403.39847	471.59689 550.67243 5588.29190 -10.500781 55.662777 120.74555 183.79211 245.69146 308.75269 369.94624 433.03974 495.88755 542.94566 625.77449	651.17249 765.52273 871.35104 -19.525624 72.207728 164.53844 269.96626 374.42538 482.45680 589.54891 697.81623 805.78453 805.78453 809.60562 899.60562	858.45170 1051.8467 1252.6685	144.85740 136.15633 127.92834	195.69748 198.01456 201.11355 6.5459042 22.558272 42.562286 59.458217 72.718818 84.437318 94.437318 95.815174 105.15660 116.57356 159.17132 238.98590 249.77755	304.92057 327.97105 352.50732 -7.4569525 65.753774 111.76345 153.58342 190.88218 226.05044 258.41299 293.22075 327.87581 340.72594 359.30560 308.05854	469.82717 529.40613 591.93696	691.44124 809.24602 932.49011 -53.87989 117.28292 268.75058 459.08504 787.14597 864.76480 1,007.8759 1,150.4553 866.62977 498.91569 174.60405

MATRIX FOR WING ON THREE-POINT SUPPORT

transverse shear included

18	19	20	21	23	24	25	26	27	28	29	50	31	32	33	3 ² 4
32.033350	33.810090	35.537950	-0.9904708	0. 7265680	1.7270130	2.890104	4.242289	5.656023	7.159676	8.622690	10.085890	22.81697	48.35806	75.02467	104.10897
165.66927	169.23542	172.92590	4.8940853		25.198559	42.141734		57.189570	65.197020	73.116620	123.32579	219.73846	267.97473	160.71208	160.71208
200.92146	210.25057	219.46150			29.807105	44.552239	57.738189	70.771349	82.697917	95.642390	108.49461		271.79719	350.01652	215.01073
269.06370	290.86860	512.15520			44.680310	68.53420	91.765180	115.20675	137.92692	161.39040	184.76460		372.75020	447.06030	304.24550
561.97740	403.16620	443.03620	-5.9751180		64.846700	102.741310	141.44839	180.88158	220.20388	259.80150	299.31630	571.48290	512.05020	601.96600	354.26080
225.27392	224.94650	225.10725		17.651738	35.400701	46.600737	56.099550	64.657478	70.886545	79.112510	87.195980	156.64091	288.99371	342.36185	189.55261
268.07614		286.45244		24.638295	48.286027		88.079581	105.56544		137.96675		221.28888	346.49413	385 91400	210.88883
	576, 55156	405.39847	-2.5574521	41.030962	79.137810	156.41280	150.95999	185.40205	217.66388	251.95140		345.38947	455.27216	464.93350	244.59359
		588.29190	-10.500781		120.74555	185.79211	245.69146	308.75269	369.94624	455.03974			625.77449	578.11752	269.59819
	763.52273		-19.325624	72.207728	164.33844	269.96626	374.42338	482.45680	589.54891	697.81623			893.28144	757.76814	293.81508
	1,051.8467				212.26541	561.52591	526.89599	697.72878	869.13423			1,222.7202	180.63418		314.93210
	156.15635	127.92834	44.872860		14.801656	18.704249	19.418556	18.865657	16.195052	15.274590	14 . 240820	71.475320	258.98590	221.15924	116.96461
195.69748	198.01456	201.11535	6.5459042	22.558272	42.562286	59.458217	72.718818	84.437318	93.815174	105.15660		159.17132 340.72594	359.30560	249 77595	125.35542
304 92037			-7.4569525	65.753774	111.76345	153.58342	190.88218	226.05044	258.41299	293.22075	527.87581		545.91874	308.03854	
469.82717		591.93696	-20.852491	91.864732	199.44808	284.71994	364.25910	440.81864	513.81372	590.28029				590.80202 498.91569	157.57906
691.44124		932.49011	-33.87989		268.75058	139.08504	585.00631	727.14957	864.78480	1,007.0799	1,100.4000	1 50k 3683	1 152 0822	599.01743	191.42752
958.26547	1.164.2587	1,381.0311	-46.858128	142.58515		5/5.82010	857.88244	1,077.8267	1,312.5158	1,556.4907	1,799.2004	2 848 4504	1 568 9615	607 60606	208.84582
1.164.2587	1,620,6646	1,997.1294	-59.78820	167.85147	407.17793	713.94670	1,077.4731	1,485.9834	1,872.3459	2,272.9010	2,010.0297	x 273 8650	1 OhB hold	792.27889	225.77026
1.381.0311	1,997.1294	2,722.5282	-73.012755	193.66160	478.18059	855.70115	1,524.0896	1,887.2107	2,509.2269	5,145.4650	3, 109.4310	NE 656715	- 10 6L8730	-15.556155	
-46.858128	-19.78820	-73.012755	111.14219	-21.025117	-38.818408	-54.628194	-69.304289	-83.378962	-97.292852	-111.02125	-124.05274	253, 19317	139.23862	56,280553	14.052684
142.58515	167.83147	193.66160	-21.025117	107.97039	145.98901	175.60547		232.44035	259.81703	200.77407			537.42009	128.99924	29.388196
357.86859	407.17793	478.18059	-38.818408		319.78886	409.49785	491.16514	568.60192	644.73105	719.65171	795.10279			215.19818	45.775940
575.82010		855.70115	-54.628194			678.95479	847.49798	1,005.7842	1,160.5526	1,512.6409	1,400.0049	1 830 1063	804 86076	303.66028	62,457000
857.88244	1.077.4731	1,324.0896	-69.304289	204.68997	491.16514	847.49798	1,249.2974	1,535.2540	1,812.4621	2,004.5000	2,559.5150	2 600 6336	1 264 0141	305.00020	79.424060
1.077.8267	1,485.9834	1,887.2107	-85.378962	232.44035	568.60192	1,005.7842	1,535.2540	2,134.1227	2,604.4168	5,005,0102	5,752.0507	2 753 0088	1 503 0313	195 32786	96.300630
1.312.5158	1.872.3459	2,509.2269	-97.292852	259.81703		1,160.5526	1,812,4621	2,604,4168	3,510.4859	9,2(0.077)	7 117 6356	5 005 3811	1 030 0281	576 26305	113.3691
1.556,4907	2.272.9878	3.145.4850	-111.02123	286 . 77487	719.63171	1,312.6409	2,084.5686	3,065.6182	4,276.8773	7, 100.4179	0 626 7008	16 410 0561	2 285 4766	1667 Okara	130.3971
1,799,2554	2,670.8295	3,769.4370	-124.83234	313.91773	795 . 10279	1,466.0645	2,559.5150	3,532.6383	5,054.7197	1,11(.5555	6 112 0661	15 001 0441	2 121 5448	780 22235	178.23725
1.594.2682	2.343.4404	3,273.8650	-98.656715	253.19317	634.78426			2,690.6336	3,743.9088 1,593.0313	13,093.3011	2 295 1266	10 101 5448	1 722 8558	1840 23108	272.16214
1,132,9822	1,568.9615	1,948.4244	-49.648739	139.23862	337.42009	592.08790	894.86976		11,793.0513	1,939.9201	667 01:202	730 22235	849.23198	849.73590	343.92730
		792.27889	-15.536155	56.280553	128.99924	215.19818	303.66028	394.30871	467.32786	710.20307	667.04292 130.39714	178.23725	272.16214	343.92730	503.16737
		225.77026	-2.6479184	14.032684	29,588196	45.775940	62.457000	79.424060	96.300630	113.56913	120.39/14	110.27127	212110214	242.92130	305.1015

TABLE XI.- INFLUENCE COEFFICIENT MATRIX FOR

(b) Antisymmetrical loading;

Sta- tion	2	4	5	6	8	9	10	11	12	14	15	16	17	18
2	664.5865			612.6123	278.6191	472.3982	660.6795	833.0287	997.6331	147.2090	281.3849	422.8304	565.6384	708.4203
4	435.6276	380.6261	592.7688	798.8377	264.0352	440.5289	610.0023	772.0952	926.6229	141.6636	268.9312	401.7767	536.9514	671.5259
5	680.6178	592.7688	1,032.748	1,405.865	435.3010	766.0688	1,070.319	1,360.526	1,636.103	241.5507	470.6019			1,187.079
6	612.6123	798.8377	1,405.865	2,019.019	600.2362	1,065.821	1,528.929	1,958.564	2,363.037	338.7714	664.9336			1,713.716
8	278.6191	264.0352	435.3010	600.2362	241.1957	376.0705	508.6251	638.9806	766.0272	135.8592	243.7442			579.1817
.9	472.3982	440.5289	766.0688	1,065.821	376.0705			1,151.125						1.053.354
10	660.6795	610.0023	1,070.319	1,528.929	508.6251	913.7638	1,332.223	1,687.743	2,037.232	305.3881	607.9971	932.9471		1,557.214
11	833.0287	772.0952	1,360.526	1,958.564	638.9806	1,151.125	1,687.743	2,260.407	2,756.579	386.6501	777.0480		1,664.049	2,101.407
12	997.6331	926.6229	1,636.103	2,363.037	766.0272	1,383.316	2,037.232	2,756.579	3,507.773	465.9746	942.6906			2,663.030
14	147.2090			338.7714			305.3881	386.6501	465.9746	114.9871	174.8473	238.9946	305.9231	
15	281.3849			664.9336			607.9971	777.0480	942.6906	174.8473	359.3019	497.0443	640.1399	785.4992
16	422.8304	401.7767	706 8619	1,011.996	354.7006		932.9471	1,207.087	1,476.400	238.9946	497.0443	786.0877	1,021.877	1,261.737
17	565.6384	536.9514	947.5047	1,363.699	467.2502	847.6122	1,246.503	1,664.049	2,059.530	305.9231	640.1399	1,021.877	1,441.862	1,798.077
18	708.4203	671.5259	1,187.079	1,713.716	579.1817	1,053.354	1,557.214	2,101.407	2,663.030	373.9022	785.4992	1,261.737	1,798.077	2,379.692
19	846.5773	803.4266	1,421.402	2,054.819	692.0864	1,261.016	1,870.972	2,540.949	3,254.237	443.8083	234.6744	1,508.200	2,164.806	2,896.768
20	981.9168	932.6754	1,650.973	2,388.839	803.0764	1,465.055	2,178.733	2,970.331	3,826.232	513.2624	1,083.383	1,754.864	2,533.696	3,420.225
25	06.12594	04.65226	151 4270	217.4603	78.405401	146.5744	215.6564	283.4477	351.7224	56.47173	143.4252	214.2897	285.0582	355.7912
	190.0611		331.1825	478.8885	170.2997		472.7408	627.2312	783.7 252	119.8828	282.8905	466.7043	633.5189	
25	305.4440		530 1599	767.7709	270.4453	501.6768	748.8850	1,010.724	1,274.112	187.0861		711.4122		1,310.366
	528.3692		740.8234	1,073.765	375.1425	693.9846	1,037.821	1,409.856	1,800.193	256.0191	577.5021	959.8497	1,392.906	1,858.470
	554.9558			1,388.443	482.5512	891.3225	1,334.706	1,820.948	2,341.644	325.8487	727.6012	1,210.016	1,768.420	2,391.666
28	683.4830	662.3809	1,176.914	1,707.203	590.5572	1,089.779	1,633.408	2,235.168	2 ,888.08 9	395.4240	877.0961	1,459.004	2,142.037	2,922.448
29	811.8107	766.0634	1,396.354	2,025.800	699.2492	1,289.392	1,933.606	2,650.610	3,434.433	466.0090	1,028.603	1,711.512	2,521.628	3,463.496
30	939.8441	909.4441	1,615.272	2,343.665	807.6335	1,488.458	2,233.002	3,065.000	3,979.494	536.3698	1,179.693	1,963.338	2,900.094	4,002.557
31	961.9050	922.1010	1,634.921	2,368.778	806.4063	1,478.626	2,208.607	3,021.258	3,907.244	525.5784	1,132.961	1,861.426	2,720.708	3,718.099
32	994.1112	934.9840	1,652.729	2,388.806	791.0607	1,436.001	2,126.053	2,889.253	3,707.556	494.5318	1,021.882	1,628.866	2,314.906	3,066.618
33 I	965.75471	877.6160	1,548.466	2,235.464	697.1640	1,249,908	1,821.355	2,430,348	2,975,027	509.7573	816.1166	1,260,108	1,732.656	2,196.925
34	795 - 1994	632.0434	1,092.738	1,496.815	447.8560	787.8441	1,109.591	1,407.850	1,691.279	246.0945	479.4174	722.8857		1,215.342

WING ON THREE-POINT SUPPORT - Concluded

transverse shear included

19	20	23	24	25	26	27	28	29	30	31	32	33	34
846.5773	981.9168	86.12594	190.0611	305.4440						961.9050		965.7547	795.1994
803.4266	932.6754	84.65226	185.8989	297.9933		538.7776	662.3809	786.0634	909.4441	922.1010	1934.9640		632.0434
1,421.402	1,650.973	151.4270	331.1825	530.1599	740.8234	957.5168	1,176.914	1,396.354	1,615.272	1,654.921	1,652.729	1,540.400	1,092.75
2,054.819	2,388.839	217.4603	478.8885				1,707.203	2,025.000	2,545.005	2,500.770	2,388.806	697.1640	1,490.01
692.0864	803.0764	78.40340	170.2997			482.5512	590,5572	099.2492	807.6335	1 1000	1,436.001	1 010 008	787 8100
1,261.016	1,465.055		317.6516	501.6768	693.9846	091.3225	1,009.779	1,209.392	1,400.400	1,470.020	2,126.053	1,249.900	1 100 501
1,870.972	2,178.733		472.7408	748.8850	1,057.821	1,554.706	1,655.400	1,955.00	2,255.002	7 001 059	2,889.253	0 1.30 31.9	1 107 950
2,540.949	2,970.331	283.4477	627.2312	1,010.724	1,409.000	1,020.940	2,233.100	2 1 2 1 1 2 2	3,003.000	3,021.270	3,707.556	2 075 027	1 601 270
3,254.237	3,826.232	351.7224	783.7252	1,274.112	1,800.195	2,541.644	2,000.009	164 0000	15,919.494	5,901.244	494.5318	500 7573	246.0945
443.8083	513.2624	156.47173	119.0020	187.0861	256.0191	525.0407	999.4240	1 008 607	1 170 602	1 132 061	1,021.882	816 1166	
234.6744	1,083.383	143 4252	262.8905	428.7939	577.5021	127.6012	0 (1.0901	1,020.005	1,119.095	1 861 106	1,628.866	1 260 108	722 8857
1,508.200	1,754.864	214.2897	466.7045	711.4122	959.0497	1,210.010	2,459.004	2 503 508	12,909.990	2 720 708	2,314.906	1 732 656	960 3015
2,164.806	2,533.696	285.0582		1,019.566	1,592.900	1,700.420	2,142.057	7 1.67 106	1 000 557	2 718 000	3,066.618	2 106 025	1 215 342
2,896.768	3,420.225	355.7912	800.5205	1,510.566	1,000.470	7 077 602	z 816 52)	1 575 176	5 330 070	1 013 520	3,922.500	2 650.546	1.454.438
3,726.823	4,471.743	426.6003	967.6572	1,605.120	2,511.650	3,013.603	3,010.334	F 934 FO	6 870 561	6 281 682	4,715.509	3 003 606	1 688 660
	5,626.605	497.5801	1,135.697	1,097.947	2,709.574	175,000	4, 704.020	5,050.504	610 0058	677 5230	1,03 7588	283.8784	152.1718
426.6003	497.5801	142.6733	217.7445	291.0005	262 - 2977	455.2005	700.0420	3 70.5404	649.9258 1,515.471	3 30h 5hx	056 0355		334.0873
967.6372	1,135.697	217.7445		000.1000	201.0205	1 600 602	1,117,900	2 206 318	2 600 208	2 246 057	1,576.275	1 015 220	535.5112
1,603.128	1,897.947	291.0805		1,000.413	1,001.170	0 1.76 491	2 060 8)3	z liba zz8	3 ook 10k	3 330 307	2,263.881	1 422 653	749. 3836
2,311.630	2,769.574			1,701.110	1,904.701	7 770 180	1,071,006	1 80x 281	5 536 6ho	1 624 256	3,008.688	1.841.434	969,4498
3,073.603	3,745.672	435.2005	1,002.004	1,000.045	2,410.004	15,559.409	5 301 800	6 307 201	7 500 234	6 106 995	3,755.516	2.265.058	1.192.574
3,816.534	4,784.020	506.8428	1,175.900	1,993.520	z 1.1.1 zzg	h 807 081	6 307 200	8 266 250	10 007 76	7 800 180	4,502.317	2,688,553	1.415.512
4,575.176	15,836.504	578.3484	1,544.529	2,290.540	z ook 10k	5 536 600	7 500 234	10 007 76	12 050 40	638 002	5,246.729	3.111.087	1.637.936
5,330.070	6,879.561	649.9258	1 201 512	2 316 057	3 330 307	h 60h 256	6 106 995	7 800 480	9 638 002	8.144.210	4,989.119	3,105,691	1.665.079
4,913.520	6,284.682	575.5250	956.0355	1 576 275	2 263 881	3 008 688	3 755 516	4 502 317	5 246 729	4.989.119	4,361.136	3,054,714	1.698.327
3,922.500	4,715.509	425.7588	609 1660	1,015,220	1 102 653	1 841 434	265 058	2 688 553	3 111 087	3, 105, 691	3,054.714	2.731.569	1.617.700
	3,093.696		334.0873	535 5112	7ho 3836	060 1408	1 100 574	1 415 512	1 637 936	1.665.079	1,698.327	1.617.700	1,228.548
1,454.438	1,688.660	152.1718	354.0075	عندر دررر	149.5050	303.4490	12,226,717	1, 12, 12	12,001.900		1-,-,0,,		

TABLE XII. - COMPARISON OF EXPERIMENTAL AND CALCULATED

FREQUENCIES FOR FREE-FREE VIBRATION

			Freque	ency, (Frequency, cps, for	r -		
	Syı	Symmetrical	sal		4	ntisym	Antisymmetrical	
lst 2d mode mode	2d mode	3d mode	4th mode	5th mode	lst mode	2d mode	3d mode	4th mode
43.3	88.8	43.3 88.8 122.8 164.2 179.7 52.2	164.2	1.671	52.2	91.7	131.1	169.2
46.4	105.3	150.0	202.0	248.0	56.70	46.4 105.3 150.0 202.0 248.0 56.70 103.4	166.6	216.5
9.44	7.4	94.7 132.0 172.0 216.0 52.20	172.0	216.0	52.20		96.29 142.26 200.66	200.66
42.8	88.9	88.9 120.1 158.0 184.0 50.52	158.0	184.0	50.52	90.25	90.25 126.83 174.26	174.26
43.1	83.0	118.0	146.0	172.0	51.1	89.0	124.1	166.7
- 7	т. .	15.1 85.0	13.1 83.0 118.0	13.1 83.0 118.0 146.0	15.1 85.0 118.0 146.0 172.0	13.1 83.0 118.0 146.0 172.0 51.1	83.0 118.0 146.0 172.0 51.1 89.0	83.0 118.0 146.0 172.0 51.1

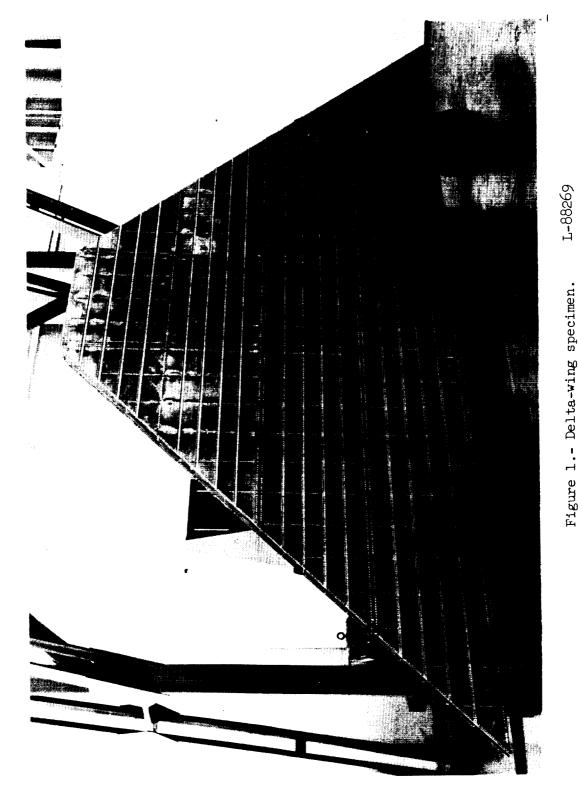
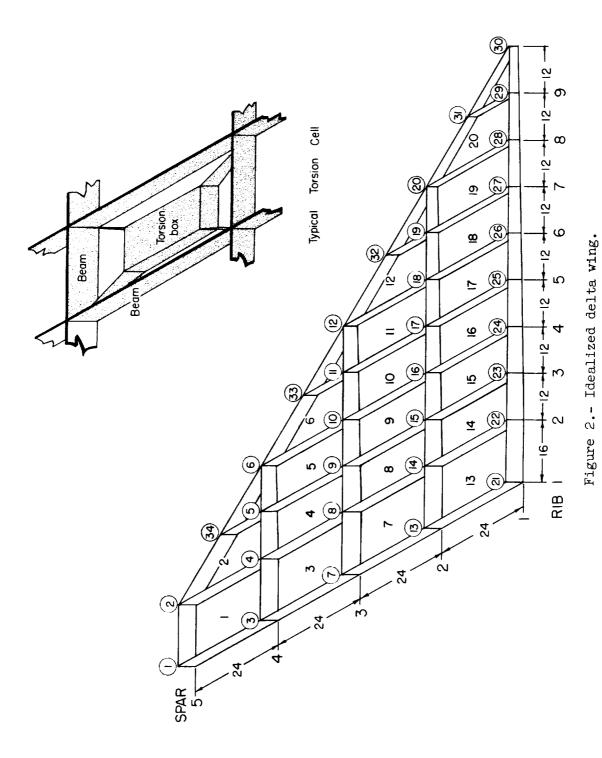
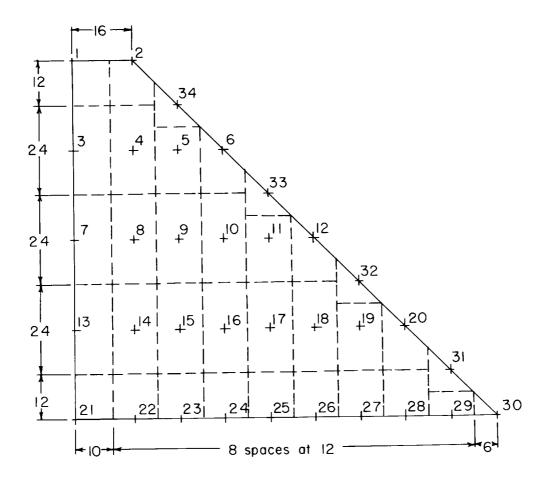


Figure 1.- Delta-wing specimen.

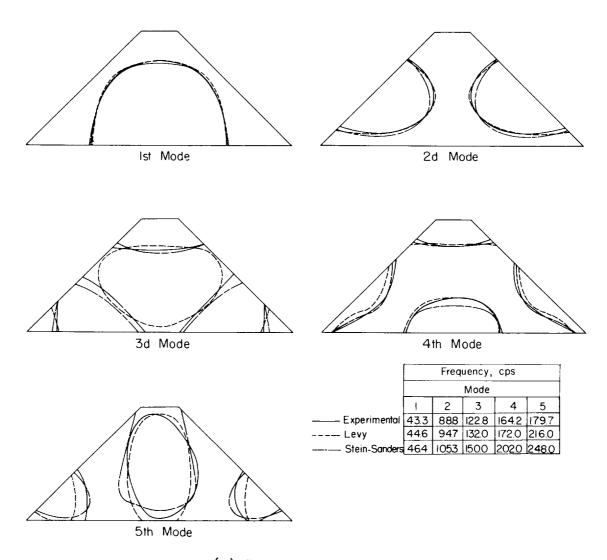




i	Wi	i	Wi	i	Wi	i	Wi
1	3.702	10	10726	19	6.649	28	6.488
2	5.975	11	7.2 32	20	5.458	29	2.486
3	6.575	12	5.840	21	4.457	30	1.717
4	10538	13	7.200	22	7.058	31	2358
5	6.404	14	11.316	23	5.496	32	2.535
6	5.553	15	8.963	24	5.8 84	33	2.705
7	7.477	16	9.095	25	52 94	34	2.877
8	11.649	17	8.652	26	5.729		
9	9.295	18	8874	27	5.095		

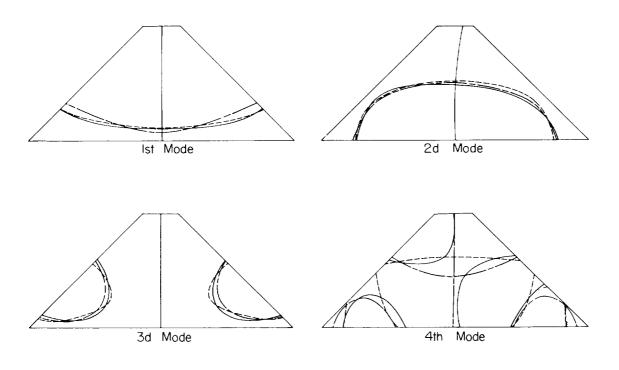
 W_i = Weight concentrated at i th station in pounds

Figure 3.- Mass distribution.



(a) Symmetrical modes.

Figure 4.- Calculated and experimental node lines and frequencies.



		Freque	ency, c	os
		М	ode	
		2	3	4
——Experimental	52.2	91.7	131.1	1692
Levy	522	96.3	142.3	200.7
Stein-Sanders	56.7	103.4	1666	2165

(b) Antisymmetrical modes.

Figure 4.- Concluded.

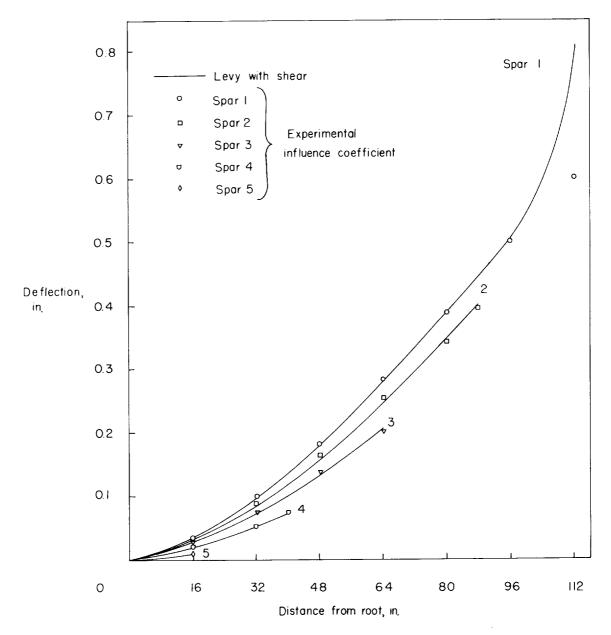


Figure 5.- Deflection of cantilevered wing under uniform load.

NASA MEMO 2-2-59L National Aeronautics and Space Administration. EVALUATION OF THE LEVY METHOD AS APPLIED TO VIBRATIONS OF A 45° DELTA WING. Edwin T. Kruszewski and Paul G. Waner, Jr. February 1959. 48p. diagrs., photo., tabs.	Vibrat Loads Struct Krusz	NASA MEMO 2-2-59L National Aeronautics and Space Administration. EVALUATION OF THE LEVY METHOD AS APPLIED TO VIBRATIONS OF A 45° DELTA WING. Edwin T. Kruszewski and Paul G. Waner, Jr. February 1959. 48p. diagrs., photo., tabs.	1. Vibration and Flutter (4.2) 2. Loads and Stresses. Structural (4.3.7) I. Kruszewski, Edwin T. II. Waner, Paul G., Jr.
(NASA MEMORANDUM 2-2-59L) The Levy method which deals with an idealized structure was used to obtain the natural modes and frequencies of a large-scale built-up 450 delta wing. The results from this approach, both with and without the effects of transverse shear, were compared with the results obtained experimentally and also with those calculated by the Stein-Sanders method. From these comparisons it was concluded that the method as proposed by Levy gives excellent results for thinskin delta wings, provided that corrections are made for the effect of transverse shear.	Ш. NASA MEMO 2-2-59L	(NASA MEMORANDUM 2-2-59L) The Levy method which deals with an idealized structure was used to obtain the natural modes and frequencies of a large-scale built-up 45° delta wing. The results from this approach, both with and without the effects of transverse shear, were compared with the results obtained experimentally and also with those calculated by the Stein-Sanders method. From these comparisons it was concluded that the method as proposed by Levy gives excellent results for thinskin delta wings, provided that corrections are made for the effect of transverse shear.	Ш. NASA MEMO 2-2-59L
	ASAN		ASAN
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	NASA		NASA