

# A 2D microphysical analysis of aerosol nucleation in the polar winter stratosphere: implications for H<sub>2</sub>SO<sub>4</sub> photolysis and nucleation mechanisms

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**Abstract.** Each spring a layer of small particles forms between 20 and 30 km in the polar regions. Results are presented from a 2D microphysical model of sulfate aerosol, which provide the first self-consistent explanation of the observed "CN layer." Photochemical conversion of sulfuric acid to SO<sub>2</sub> in the upper stratosphere and mesosphere is necessary for this layer to form. Recent laboratory measurements of H<sub>2</sub>SO<sub>4</sub> and SO<sub>3</sub> photolysis rates are consistent with such conversion, though an additional source of SO<sub>2</sub> may be required. Nucleation throughout the polar winter extends the top of the aerosol layer to higher altitudes, despite strong downward transport of ambient air. This finding may be important to heterogeneous chemistry at the top of the aerosol layer in polar winter and spring.

**Keywords:** Aerosols and particles, Middle atmosphere--composition and chemistry, Middle atmosphere--constituent transport and chemistry

## Introduction

*Rosen and Hofmann* [1983] first noted an anomalous enhancement in springtime particle concentration near 30 km altitude at mid-latitudes in air of recent Arctic origin. Similar enhancements were observed from September to November over Antarctica [*Hofmann and Rosen*, 1985; *Hofmann*, 1988; *Hofmann et al.*, 1988; *Hofmann et al.*, 1989; *Wilson et al.*, 1989; *Hofmann*, 1990; *Hofmann and Deshler*, 1991], and in January and February over the Arctic [*Hofmann*, 1990; *Wilson et al.*, 1990]. Because the concentration is dominated by the smallest particles, known as "condensation nuclei," the polar anomaly has become known as a "CN layer."

Antarctic in situ measurements indicate modest particle nucleation at the top of the sulfate layer before the arrival of sunlight, with dramatic increases producing a well-defined layer when sunlight arrives. Balloon-borne measurements from McMurdo Station in late August 1987 and 1988 show particle concentrations increasing to a maximum of ~20 particles/cm<sup>3</sup> near 25 km [*Hofmann et al.*, 1989; *Hofmann*, 1990]. Because particle concentrations should otherwise decrease with air density (as is observed in polar summer), these observations suggest that nucleation occurs between 20 and 30 km in polar night. Observations made in each austral spring from 1983 to 1989 are consistent with a well-defined layer of 100-150 cm<sup>-3</sup> forming each September between 22 and 28 km, with the altitude of the peak concentration descending through October to near 20-22 km in early November [*Hofmann and Rosen*, 1985; *Hofmann*, 1988; *Hofmann et al.*, 1988; *Hofmann et al.*, 1989; *Hofmann*, 1990; *Hofmann and Deshler*, 1991]. In 1993, the year following the eruption of Mt. Pinatubo, particle concentrations reached 50 cm<sup>-3</sup> by late August, and peaked at about 200 cm<sup>-3</sup> in early September [*Deshler et al.*, 1994].

The CN layer was first discovered in Northern Hemisphere measurements, and measurements clearly show well-defined CN events in the arctic, albeit with less consistency than the Antarctic. Observations over Kiruna, Sweden (68°N)

show a layer of  $40/\text{cm}^3$  between 22.5 and 26 km on January 30, 1989 [Hofmann, 1990]. But four profiles over Kiruna in January and February of 1992 show no definite layer, with CN concentrations constant or increasing slightly with altitude [Deshler, 1994]. The increasing particle mixing ratio with altitude in these profiles indicates that some nucleation occurred at the top of the layer, comparable to the weak nucleation indicated in Antarctic winter before sunrise. The variability in these arctic measurements are likely a product of the greater latitudinal excursions of vortex air in the Arctic as opposed to the Antarctic, as well as the lack of measurements well above the Arctic Circle. Trajectory analyses of the air sampled over Kiruna in 1992 indicated that it had recently been as far south as  $50^\circ\text{N}$ , and may have mixed with mid-latitude air. Such excursions and mixing explain the observation of the arctic CN layer over Laramie, Wyoming ( $41^\circ\text{N}$ ) in late winter and early spring of 1982, 1983, and 1986 [Rosen and Hofmann, 1983; Hofmann et al., 1985; Hofmann et al., 1988]. Because the Antarctic vortex is more circumpolar, leading to less variation in latitude and sunlit history of air parcels, and because Antarctic observations of the CN layer are more extensive and consistent than those in the Arctic, the focus of this paper is the Antarctic CN layer. It should be noted, however, that while the dynamical variability makes modeling the timing of the CN layer more complicated in the Arctic, the basic theory we present applies there as well.

## Theories

Rosen and Hofmann [1983] suggested two possible origins for the CN layer. First, sulfur-bearing gases might maximize in winter, due to lack of sunlight available to oxidize them to sulfuric acid. These gases might dissociate rapidly with the return of sunlight, OH, and O in spring. Their second theory relies on very rapid drops in temperature observed in the Arctic to create very large supersaturations. In the absence of large, existing particles on which to condense, the excess of ambient vapor would nucleate many new, small particles.

Building on the nucleation rate studies of Hamill et al. [1990], Zhao et al. [1995] investigated the CN layer with a 1-dimensional microphysical model. They questioned whether sufficient sulfuric acid is present in polar winter, given the lack of photochemical production there, to initiate the temperature-dependent mechanism of Rosen and Hofmann [1983]. They also calculated that OCS oxidation can not provide sufficient sulfuric acid vapor to allow new particles to grow to observable size, as suggested by Oppenheimer [1987]. They postulated that the “extra” sulfur required for the formation of the layer can be provided by  $\text{SO}_2$  descending from the mesosphere. The source of this  $\text{SO}_2$  was hypothesized to be  $\text{H}_2\text{SO}_4$  photolysis at lower latitudes. By setting upper boundary conditions of 100 pptv for both  $\text{SO}_2$  and  $\text{H}_2\text{SO}_4$  at 60 km and allowing for vertical subsidence, Zhao et al. were able to simulate the formation of a CN layer similar to that which is observed. The CN concentration was found to be much more sensitive to the  $\text{SO}_2$  concentration than the  $\text{H}_2\text{SO}_4$  concentration.  $\text{SO}_2$

concentrations increase during the polar night as advection carries SO<sub>2</sub>-rich air downward. During descent gas-phase sulfuric acid is lost on existing particles, and SO<sub>2</sub> remains as the dominant gaseous sulfur reservoir until the return of sunlight in spring. *Zhao et al.* did not model the production of SO<sub>2</sub> based on measured photolysis rates as we do here.

## Photochemistry of SO<sub>2</sub>

We present here calculations from a new microphysical model, which has been incorporated into the Garcia-Solomon two-dimensional dynamical/chemical model. The model spans 56 pressure levels from 2 to 112 km above sea level, and 36 latitude bins from 89.5°S to 89.5°N. Details of the photochemistry and dynamics of the model are described in *Garcia et al.* [1992] and *Garcia and Solomon* [1994]. We have added both sulfur oxidation chemistry and particle microphysics to this model.

OCS, tropopause SO<sub>2</sub>, and tropopause aerosol are the sources of stratospheric aerosol in the model. Oxidation of gas-phase sulfur species is calculated according to the recommendations of *DeMore et al.* [1994] except as noted below. The mixing ratio for OCS is maintained at 510 pptv at the lower boundary of the model (750 mbar). OCS is the dominant source of sulfate at the top of the stratospheric aerosol layer. In the lower stratosphere, most of the sulfate comes from tropospheric aerosol and SO<sub>2</sub>, which are maintained at constant mixing ratios of 0.27 ppbm of SO<sub>4</sub><sup>-</sup> and 80 pptv of SO<sub>2</sub> at the model's tropical tropopause. Tropospheric aerosol is represented by a bimodal lognormal distribution, with 30% of the mass in a 0.8- $\mu$ m radius mode, and the majority of particles distributed about 0.015 $\mu$ m. The details of these tropospheric sources of sulfur do not affect the calculation of the polar CN layer.

As temperatures increase with altitude in the stratosphere, sulfate aerosol evaporates, producing sulfuric acid vapor. This vapor is thought to photolyze to SO<sub>2</sub> at the short wavelengths of sunlight present in the upper stratosphere and mesosphere. *Rinsland et al.* [1995] concluded from observations of SO<sub>2</sub> in the upper stratosphere and mesosphere and from 2-dimensional model calculations, that sulfuric acid vapor should photolyze at rates roughly 0.3 times that of HCl. We present this assumption as model case A (see table 1). Recently, *Burkholder et al.* [1998] measured no photolytic cross sections for H<sub>2</sub>SO<sub>4</sub> vapor greater than their upper limit of 10<sup>-21</sup> cm<sup>2</sup>, for wavelengths greater than 195 nm. Our model case B sets the cross sections to that upper limit of 10<sup>-21</sup> cm<sup>2</sup> in the 179-226 nm region where HCl absorbs, representing maximum photolysis rates consistent with the findings of *Burkholder et al.* Case C assumes that the cross section is 10<sup>-21</sup> cm<sup>2</sup> at 195 nm, with the same wavelength dependence as HCl (decreasing with increasing wavelength). This assumption sets the cross sections to 0.011 times those of HCl. Case D assumes that H<sub>2</sub>SO<sub>4</sub> vapor does not photolyze at all.

*Burkholder and McKeen* [1997] recently measured  $\text{SO}_3$  photolysis cross sections, which we have included in the model. This reaction is an important source of  $\text{SO}_2$  near the top of the sulfate layer, where  $\text{SO}_3$  photolyzes faster than it oxidizes to  $\text{H}_2\text{SO}_4$ . The oxidation of  $\text{SO}_3$  proceeds through the production of an  $\text{SO}_3 \cdot \text{H}_2\text{O}$  adduct, resulting in a pressure-dependent effective reaction rate [*Lovejoy et al.*, 1996]. The timescale for oxidation increases exponentially with altitude, from a few seconds near the tropopause to a few days at 40 km. Photolysis becomes faster with altitude, overtaking oxidation as the dominant loss for  $\text{SO}_3$  near 40 km. Model cases A through D include  $\text{SO}_3$  photolysis. We test the importance of the new  $\text{SO}_3$  photolysis reaction to the formation of the CN layer with case E and F, which are the same as cases B and D, respectively, but without  $\text{SO}_3$  photolysis.

Calculated  $\text{SO}_2$  mixing ratios for cases B, D, E, and F are shown in Figure 1, revealing the effects of photolysis of  $\text{H}_2\text{SO}_4$ ,  $\text{SO}_3$ , and  $\text{SO}_2$ . The figures show  $\text{SO}_2$  on June 1 building up to more than 170 pptv in case B. Mixing ratios are greatest in polar winter, where destruction by photolysis is shut off, and SO thus formed reverts to  $\text{SO}_2$ .

## Particle Microphysics

The microphysical model includes classical binary homogeneous nucleation; coagulation; condensation and evaporation of water, sulfuric acid, and nitric acid; transport with ambient air parcels; and sedimentation [*Mills*, 1996]. The model simulates the formation of Type Ib (liquid ternary solutions) and Type II (water ice) polar stratospheric clouds, but does not deal with Type Ia (nitric acid hydrate) PSCs. Aerosol particles are divided into 45 size classes ranging from  $\sim 0.5$  nm to  $\sim 10$   $\mu\text{m}$ .

Calculation of binary homogeneous nucleation rates is highly uncertain. *Wyslouzil et al.* [1991] present laboratory measurements of nucleation rates, showing discrepancies of  $\pm 10$  orders of magnitude with classical theory. *Weber et al.* [1996] conclude from atmospheric measurements that particle number varies as the square of sulfuric acid concentrations, rather than the classically predicted exponential response. *Brock et al.* [1995] point out, however, that despite these uncertainties, classical nucleation theory correctly predicts the regions in the atmosphere in which large numbers of particles are observed. After experimenting with alternative nucleation parameterizations based on the findings of *Wyslouzil et al.* and *Weber et al.*, we conclude that classical nucleation theory provides the best correspondence with observed aerosol behavior in the stratosphere.

Figure 2 shows the calculated concentrations of particles larger than 0.01  $\mu\text{m}$ , the lower size limit of balloon-borne CN counters, from the 6 model cases described in Table 1. Cases A and B reproduce well the observed particle concentrations in both winter and spring. In these cases, particle production is evident throughout the winter, with dramatic

increases to 80-120 particles/cm<sup>3</sup> in September. The CN layer descends from 30 to 22 km as the austral spring progresses. Measurements made repeatedly in the spring of a given year show the CN layer to be descending in a similar manner [Hofmann *et al.*, 1988; Hofmann *et al.*, 1989; Deshler *et al.*, 1994]. Although the observed CN layer sometimes initiates as low as 26 km, these observations also indicate the presence of enhanced CN concentrations at 30 km. The calculations diverge from observations in the other model cases, when either H<sub>2</sub>SO<sub>4</sub> or SO<sub>3</sub> photolysis rates are decreased.

Figure 3 shows the seasonal cycle of temperature, nucleation, and particle concentration for case B. There is a small, but significant rate of nucleation near 30 km as early as April which continues throughout the winter, as cold temperatures (below 220 K) reach higher altitudes. Prior to July, at any given altitude nucleation returns to negligible rates within several weeks of its initiation, despite continued decreases in temperature, as aerosol formation and growth deplete ambient sulfuric acid vapor. The winter phase of nucleation produces small concentrations (10-50 /cm<sup>3</sup>) of particles, most of which grow to observable size at the 30 km maximum. Observations at McMurdo Station (78°S) show particle concentrations increasing slightly from about 10 /cm<sup>3</sup> at 16 km to about 20/cm<sup>3</sup> at 26 km in late August, before the well-defined CN layer forms [Hofmann *et al.*, 1989; Hofmann, 1990]. Our calculations suggest similar concentrations of observable particles at 26 km. Aerosol is not suppressed by the strong descent prevalent in polar winter, but instead is present at higher altitudes than in summer due to nucleation. If the only source of particles were at lower altitudes (i.e. the tropopause), subsidence would push down the top of the layer in polar stratospheric winter.

About a month after temperatures reach their minimum below 190K in late June and early July, particle nucleation increases dramatically at all altitudes below about 36 km. Particle concentrations respond rapidly, peaking at over 10<sup>4</sup>/cm<sup>3</sup> between 22 and 30 km. This event coincides with the return of sunlight in late July, and continues until temperatures increase to greater than 200 K at the beginning of October. A layer of observable-sized particles forms above 30 km in early September, descending in altitude and increasing in concentration until a maximum of about 80/cm<sup>3</sup> is found near 26 km in late October. This maximum in observable particles occurs nearly 2 months after the peak in total particle concentration.

The delay between the minimum in temperature and the maximum in calculated nucleation, and the subsequent lag between concentrations of total particles and observable particles depend critically on the descent of SO<sub>2</sub> from above. Figure 4 shows the annual cycles of various calculated gas-phase species and aerosol parameters at 26 km and 74°S for cases B, D, and E. The arrival of polar night at 74°S is evident from the dramatic drop in OH as photolytic production shuts off at the end of May. Air descends from the mesosphere to the stratosphere throughout polar winter. In case B, the air carries with it mixing ratios of SO<sub>2</sub> as high as 160 pptv at some altitudes (see Figure 1). Liquid sulfate mass, shown in Figure 4, decreases by sedimentation and descent throughout the austral winter in all three cases.

In mid-May, the disappearance of OH cuts off all production of  $\text{H}_2\text{SO}_4$  vapor, while the rate of loss due to condensation increases slightly on newly nucleated particles descending from above 30 km (see Figure 3). Loss of production and enhancement of the sink causes  $\text{H}_2\text{SO}_4$  mixing ratios to plummet by 2-4 orders of magnitude in cases B, D, and E. The effect is greatest in case D because more  $\text{H}_2\text{SO}_4$  vapor is available for nucleation above 30 km in May, resulting in greater surface areas and faster condensation.

Sulfuric acid vapor recovers in July, coincident with the return of sunlight and OH. At that time production of  $\text{H}_2\text{SO}_4$  vapor from  $\text{SO}_2$  resumes (quite rapidly in cases B and E) as  $\text{SO}_2$  has increased to more than 50 times its summertime minimum. The increase in sulfuric acid vapor concentration by two orders of magnitude in July, while cold temperatures continue to suppress equilibrium vapor pressures, generates significant nucleation rates, especially in cases B and E. Sustained over the months of August and September, the nucleation event produces particle concentrations in excess of  $10^4$  / $\text{cm}^3$  in these cases (see Figure 3).

By late September these particles at 26 km grow to observable size in case B. The growth in particle size coincides with a dramatic increase in liquid sulfate mass, indicating that this growth is not due merely to coagulation. The steady increase in OH with increasing sunlight as spring progresses hastens the conversion of  $\text{SO}_2$  to sulfuric acid, which condenses on the recently nucleated particles, rather than nucleating new ones. Nucleation shuts off in early October in favor of condensation due to rising temperatures. At 26 km in polar winter, homogeneous nucleation rates calculated from classical theory decrease by about 6 orders of magnitude as temperature increases from 180K in late September to 200K in early October. Condensation rates have a much lesser dependence on temperature.

Comparison of cases B, D, and E in Figure 4 reveals the importance of both  $\text{SO}_3$  and  $\text{H}_2\text{SO}_4$  photolysis to the production of a CN layer resembling that observed.  $\text{H}_2\text{SO}_4$  photolysis provides an essential mechanism for suppressing homogeneous nucleation in winter and conserving gas-phase sulfur as  $\text{SO}_2$ , which becomes available again as  $\text{H}_2\text{SO}_4$  in the spring to fuel nucleation. Although *Burkholder et al.*'s measurements reveal no evidence for  $\text{H}_2\text{SO}_4$  photolysis, their upper limit produces enough  $\text{SO}_2$ , in conjunction with  $\text{SO}_3$  photolysis, to create a CN layer similar to that observed.  $\text{SO}_3$  photolysis prevents  $\text{SO}_2$  from oxidizing back to  $\text{H}_2\text{SO}_4$ . At 40 km  $\text{SO}_3$  photolysis can increase descending  $\text{SO}_2$  mixing ratios by 50 pptv (see Figure 1), and at 26 km it increases the peak  $\text{SO}_2$  mixing ratio seen in September by a factor of 2.25 (see Figure 4).

Figure 5 shows the effect of the CN layer on aerosol surface areas at 74°S in each of the six model cases. In all cases, surface areas between 30 and 35 km are enhanced in the winter and spring relative to summer and autumn.  $\text{H}_2\text{SO}_4$  photolysis in cases A and B results in springtime surface areas enhanced by as much as a factor of 2 over cases D and F, in which  $\text{H}_2\text{SO}_4$  photolysis is turned off. Polar winter nucleation and the springtime CN layer may affect heterogeneous

reactions between sulfate aerosol and various gas-phase species on a seasonal basis. In all cases, surface areas of  $0.1 \mu\text{m}^2/\text{cm}^3$  or greater extend up to 36 km throughout winter and spring, and to less than 30 km in summer and autumn. Depending on reaction probabilities and ambient concentrations of gas-phase species, such surface areas could perturb gas-phase species via heterogeneous chemistry. In the absence of new particle nucleation above 30 km, dynamical descent in the polar winter would suppress the sulfate aerosol layer, reducing rather than increasing the altitude of the top of the layer relative to summer.

## Conclusions

The 2D microphysical model of sulfate aerosol that we have developed provides the first self-consistent explanation of the CN layer observed in Antarctic spring. Although dynamics complicate the Arctic situation, the same theory applies to enhanced CN observed there [Hofmann, 1990; Wilson *et al.*, 1990; Deshler, 1994].

Binary homogeneous nucleation and condensation following rapid production of sulfuric acid vapor from  $\text{SO}_2$  in spring is responsible for the appearance of a CN layer, as suggested by Zhao *et al.* [1995]. Photolysis of  $\text{SO}_3$  and  $\text{H}_2\text{SO}_4$  are proposed as the mechanism for producing  $\text{SO}_2$  in the upper stratosphere and mesosphere; downward transport in polar winter and spring then bring the  $\text{SO}_2$  into the lower stratosphere. The model demonstrates that photolysis of  $\text{SO}_3$  and  $\text{H}_2\text{SO}_4$  calculated within the constraints of recent laboratory measurements, can partition sulfur as  $\text{SO}_2$  in the upper stratosphere and mesosphere to produce a springtime CN layer much like that observed. However, when a wavelength dependence in the  $\text{H}_2\text{SO}_4$  cross sections matching that of HCl further constrains photolysis, the calculated CN layer has substantially smaller concentrations than that observed. An additional source for  $\text{SO}_2$  in the upper stratosphere and mesosphere may be indicated to explain the CN layer, as well as observations of  $\text{SO}_2$  at these altitudes.

The model also suggests that new particles nucleate and grow throughout the winter above 30 km, at a slower rate than that which produces springtime CN layer. Thus aerosol is present at higher altitudes in polar winter than in summer, despite the subsidence of polar air which would otherwise push the top of the aerosol layer downward. Concentrations of new particles thus formed are consistent with *in situ* measurements made in polar winter. These particles may be important to heterogeneous chemistry at the top of the aerosol layer in polar winter and spring.

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## Tables

Case	H <sub>2</sub> SO <sub>4</sub> cross sections	SO <sub>3</sub> photolysis	
A	$\sigma_{\text{H}_2\text{SO}_4} = 0.3 \cdot \sigma_{\text{HCl}}$	Rinsland fit to observations	On
B	$\sigma_{\text{H}_2\text{SO}_4} = 10^{-21} \text{ cm}^2$	<i>Burkholder et al.</i> upper limit	On
C	$\sigma_{\text{H}_2\text{SO}_4} = 0.011 \cdot \sigma_{\text{HCl}}$	Upper limit with HCl tail	On
D	$\sigma_{\text{H}_2\text{SO}_4} = 0.0$	No H <sub>2</sub> SO <sub>4</sub> photolysis	On
E	$\sigma_{\text{H}_2\text{SO}_4} = 10^{-21} \text{ cm}^2$	<i>Burkholder et al.</i> upper limit	Off
F	$\sigma_{\text{H}_2\text{SO}_4} = 0.0$	No H <sub>2</sub> SO <sub>4</sub> photolysis	Off

**Table 1.** Model cases, varying H<sub>2</sub>SO<sub>4</sub> and SO<sub>3</sub> photolysis.

## Figure Captions

**Figure 1.** SO<sub>2</sub> mixing ratios (pptv) for four model cases representing various assumptions about H<sub>2</sub>SO<sub>4</sub> and SO<sub>3</sub> photolysis. Calculations are shown for noon, June 1.

**Figure 2.** Calculated concentrations (/cm<sup>3</sup>) of particles large enough to be observed ( $r > 0.01 \mu\text{m}$ ) in six model cases, varying photolysis rates for H<sub>2</sub>SO<sub>4</sub> and SO<sub>3</sub>. Results are presented as annual cycles at 74°S, averaged over 5-day periods for each altitude gridpoint (2 km resolution). In cases A and B, CN layers are calculated which resemble those observed. SO<sub>3</sub> photolysis is included in cases A through D. Comparison of cases B and E reveals the importance of SO<sub>3</sub> photolysis to the concentration of particles in the CN layer.

**Figure 3.** Model calculations of (a) temperatures, (b) homogeneous binary nucleation rates, (c) total particle concentrations, and (d) particles of radius greater than 0.01 μm at 74°S in model case B. A delay of about one month is evident between the periods of minimum temperature and maximum nucleation. A similar delay occurs between the maximum in total particle concentration and that in particles large enough to be observed.

**Figure 4.** Variation in OH, SO<sub>2</sub>, sulfate mass, aerosol surface area, sulfuric acid vapor, sulfuric acid equilibrium vapor pressure, nucleation rate, and particles greater than 0.01 micron radius at 26 km, 74°S for model cases B, D, and E (see table 1 and text).

**Figure 5.** Calculated CN layer surface areas (μm<sup>2</sup>/cm<sup>3</sup>) in six model cases at 74°S.



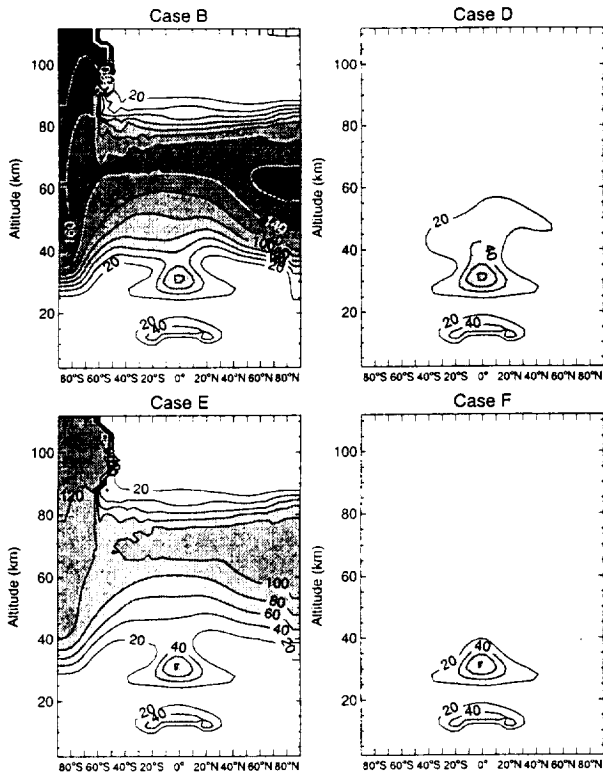


Figure 1

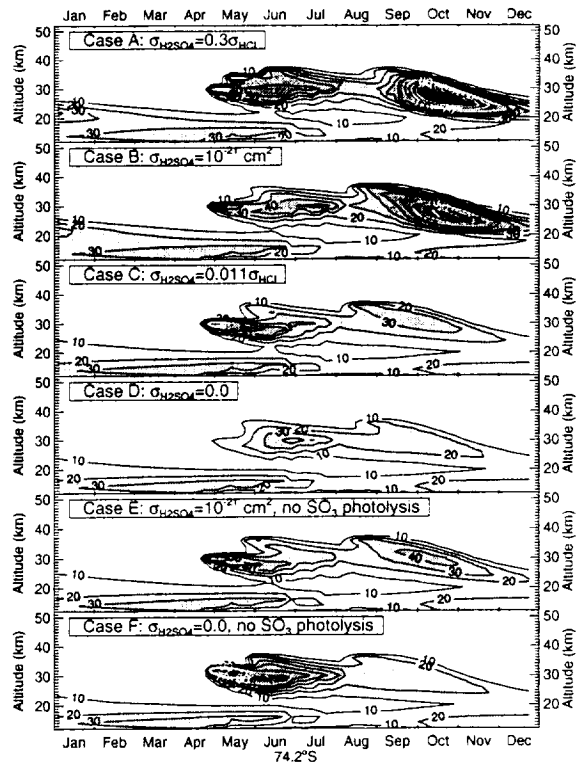


Figure 2

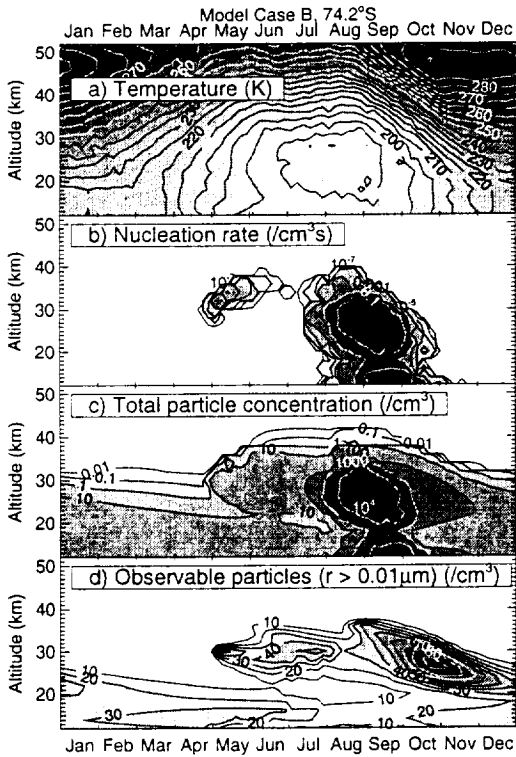


Figure 3

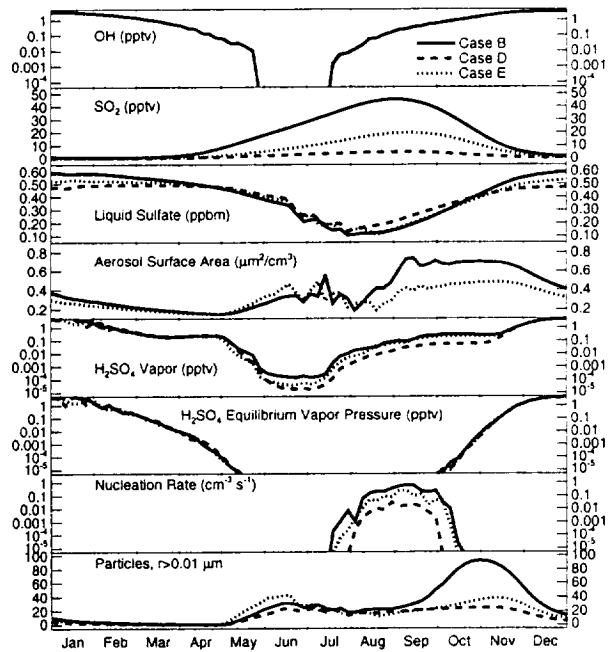


Figure 4

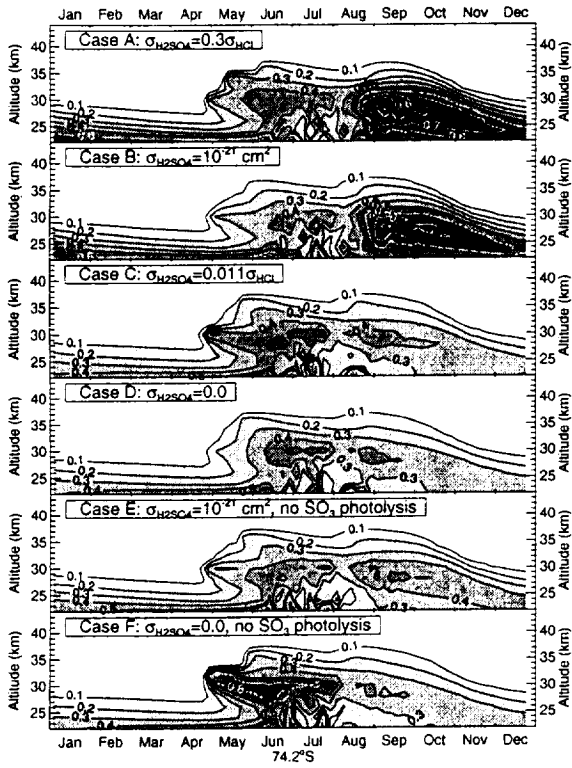


Figure 5

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