NASA LAW, February 14-16, 2006, UNLV, Las Vegas

# Laboratory experiments on interstellar ice analogs: The sticking and desorption of small physisorbed molecules

G. W. Fuchs<sup>1</sup>, K. Acharyya<sup>2</sup>, S. E. Bisschop<sup>1</sup>, K. I. Öberg<sup>1</sup>, F. A. van Broekhuizen<sup>1</sup>, H. J. Fraser<sup>3</sup>, S. Schlemmer<sup>4</sup>, E. F. van Dishoeck<sup>1</sup>, & H. Linnartz<sup>1</sup>

<sup>1</sup>Raymond and Beverly Sackler Laboratory for Astrophysics at Leiden Observatory, The Netherlands, <sup>2</sup>Center For Space Physics, Kolkata, India, <sup>3</sup> Department of Physics, University of Strathclyde, Scotland, <sup>4</sup> I.Physikalisches Institut, University of Cologne, Germany

fuchs@strw.leidenuniv.nl

## ABSTRACT

Molecular oxygen and nitrogen are difficult to observe since they are infrared inactive and radio quiet. The low  $O_2$  abundances found so far combined with general considerations of dense cloud conditions suggest molecular oxygen is frozen out at low temperatures (< 20 K) in the shielded inner regions of cloud cores. In solid form  $O_2$  and  $N_2$  can only be observed as adjuncts within other ice constituents, like CO. In this work we focus on fundamental properties of  $N_2$  and  $O_2$ in CO ice-gas systems, e.g. desorption characteristics and sticking probabilities at low temperatures for different ice morphologies.

## 1. Introduction

The understanding of dense cloud and pre-stellar core chemistry depends strongly on the available molecule budget and distribution. This is particularly true for oxygen and nitrogen. Since  $O_2$  nor  $N_2$  has an electric dipole moment direct observations of these species in cold environments are difficult and abundances are uncertain. Nevertheless, the low values that have been derived from SWAS and ODIN observations of the order  $N(O_2)/N(H_2) \approx$  $10^{-7} - 10^{-8}$  for gas phase  $O_2$  raise questions on the total oxygen budget, especially when compared with atomic oxygen abundances observed towards diffuse clouds. It is believed that at high enough densities  $(n(H_2)=10^6-10^8 \text{ cm}^{-3})$  and low enough temperatures (10-30

<sup>&</sup>lt;sup>1</sup>Correspondence to: G. W. Fuchs

K) even volatile molecules like CO,  $N_2$  and  $O_2$  are removed from the gas-phase by freezeout processes on grains, thus building interstellar ices. These freeze-out processes can also significantly influence the chemistry. For example, since CO is one of the main destroyers of  $N_2H^+$  a different freeze-out temperature of CO and  $N_2$  would alter any subsequent reaction of  $N_2H^+$ . An anti-correlation between  $C^{18}O$  and  $N_2H^+$  can be observed towards the B68 pre-stellar cloud and it has been suggested that different sticking probabilities or binding energies between CO and  $N_2$  (the parent molecule of  $N_2H^+$ ) are responsible for this effect (Flower *et al.* 2005; Bergin *et al.* 2002). Experiments in the laboratory help clarifying the situation by measuring key properties of ice-gas interactions, like the sticking probability, desorption kinetics and binding energies of the species.

#### 2. Experiment

Experiments have been performed using the CRYOPAD ultra-high vacuum set-up previously described in van Broekhuizen 2005 and Fuchs *et al.* 2006. CO, N<sub>2</sub> and O<sub>2</sub> have been deposited on a poly-crystalline gold surface at 14 K in mixed and layered structures forming amorphous ices. For astrophysical relevance, ices between 10 and 160 mono-layer thickness have been investigated using TPD (Temperature Programmed Desorption) techniques. The TPD spectra are used to determine the binding energies and desorption kinetics of these species under various conditions. They also give insight into the underlying intermolecular interactions. The TPD were taken by a controlled linear warm up (dT/dt = 0.1 K/min) of the ices between 14 and 80 K during which the evaporated molecules have been monitored by a quadrupole mass spectrometer. Simultaneously reflection absorption infrared (RAIR) spectra have been recorded using a Fourier Transform infrared spectrometer to investigate changes in the ice structure. The RAIR data have been discussed in Fuchs *et al.* 2006 and aid the modeling of the TPD data.

#### 3. Results

The experiments have been discussed in Fuchs *et al.* 2006 and Oberg *et al.* 2005, Bisschop *et al.* 2006, Acharyya *et al.* 2006. To elucidate the anti-correlation problem between CO and N<sub>2</sub>H<sup>+</sup> and also to confirm whether O<sub>2</sub> might be frozen out on grains the sticking coefficients S for CO sticking on CO, N<sub>2</sub> on CO, O<sub>2</sub> on O<sub>2</sub>, etc. have been determined It has been found that all sticking coefficients at 15 K have similar values and that their lower limit value is around 0.9, thus verifying the assumption S=1 made by most modelers. Hence, O<sub>2</sub> can easily freeze out below 20 K in dense cores and the anti-correlation between CO and N<sub>2</sub>H<sup>+</sup> is not due to different sticking coefficients of N<sub>2</sub> and CO.

TPD spectra have been recorded for pure  $O_2$ , CO and  $N_2$  and analyzed using Polanyi-Wigner type rate equations, e.g. to determine the desorption energies  $E_d$ . It has been found that

 $E_d(O_2) > E_d(CO) > E_d(N_2)$ , with  $E_d(O_2) \approx 925$  K,  $E_d(CO) \approx 855$  K and  $E_d(N_2) \approx 790$  K using 0<sup>th</sup>-order desorption kinetics. Layered and mixed CO-O<sub>2</sub> ices reveal similar TPD spectra with respect to pure ices but the desorption energies of the O<sub>2</sub> in these cases are lowered by 10 K whereas the  $E_d$  for CO is increased by 15 K. No significant co-desorption of CO with O<sub>2</sub> can be observed (Acharyya *et al.* 2006). In contrast, CO-N<sub>2</sub> ices show codesorption of N<sub>2</sub> with CO for layered and mixed ices. Thus, the CO-N<sub>2</sub> ice system reveals a two-step desorption process for the N<sub>2</sub> molecule involving 0<sup>th</sup>- and 1<sup>st</sup>-order kinetics due to diffusion and a sufficiently large  $E_d$  energy separation between the N<sub>2</sub>-N<sub>2</sub> binding sites and the N<sub>2</sub>-CO type binding sites (Bisschop *et al.* 2006).

Both N<sub>2</sub> and O<sub>2</sub> interact with CO mainly due to quadrupole interaction with quadrupole moments of  $Q(N_2) \leq Q(CO)$  and  $Q(O_2) \ll Q(N_2)$ . Furthermore, CO and N<sub>2</sub> have the same crystalline structure below 30 K whereas O<sub>2</sub> possesses a different crystalline architecture. Thus, N<sub>2</sub> can interact well with CO, i.e. stick, mix and co-desorb, whereas O<sub>2</sub> cannot. This follows the more general idea that desorption processes can be classified according to their intermolecular interaction pattern and crystalline structures. For example, CO, N<sub>2</sub> and O<sub>2</sub> in water ice behave differently (Collings *et al.* 2004, see their Fig. 3) causing CO to desorb at higher temperatures than O<sub>2</sub> (around 30-50 K)

# 4. Astrophysical implications



Fig. 1.— Desorption on astrophysical timescales  $(1 \text{ K}/10^3 \text{ yr})$  starting at 10 K and assuming 40 monolayer of pure, mixed or layered O<sub>2</sub>, CO, N<sub>2</sub> ice (i.e. no accretion has been considered).

The experimental data have been incorporated in an empirical kinetic model to predict desorption rates at a typical warm up rate  $(dT/dt = 1\text{K}/10^3 \text{ yr})$  for a newly born star (Fig. 1). N<sub>2</sub> desorbs first, then CO followed by O<sub>2</sub> at the highest temperatures. Independent of the ice morphology, oxygen molecules desorb mainly between 18.5 and 20 K. CO is more volatile and also has a wider desorption temperature range showing larger differences between pure, mixed and layered CO ices. Most remarkably is N<sub>2</sub> which in its layered CO-N<sub>2</sub> form desorbs in two steps covering a desorption range of 4 K between 15 and 19 K. Overall the desorption takes place within 5 K for all species, i.e. equivalent to only 5000 yr of the evolutionary stage of the star, and can thus hardly make up for the observed anti-correlation between CO and N<sub>2</sub>H<sup>+</sup>. Although O<sub>2</sub> can well be frozen out in dense cores it has not been seen in the later stages of star formation, e.g. in hot cores, probably due to beam dilution effects. Future deep searches for O<sub>2</sub> with the HERSCHEL Space Observatory will be able to put stringent limits on the gaseous O<sub>2</sub> abundance.

This work has been supported by NOVA and a NWO Spinoza grant.

# REFERENCES

Acharyya et al. 2006, Astron.&Astrophys., in prep. Bergin et al. 2002, Astrophys.J., 570, L101 Bisschop et al., 2006, Astron.&Astrophys., in press, (see astro-ph/0601082) Collings et al. 2004, Mon.Not.R.Astron.Soc., 354, 1133 Flower et al. 2005, Astron.&Astrophys., 436, 933 Fuchs et al. 2006, Faraday Discuss., in press (DOI:10.1039/B517262B) Öberg et al. 2005, Astrophys.J.Lett. 621, L33. van Broekhuizen 2005, PhD thesis, Leiden University