

### A cryogenic infrared calibration target

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## A cryogenic infrared calibration target

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A compact cryogenic calibration target is presented that has a peak diffuse reflectance,  $R \le 0.003$ , from 800 to 4800 cm<sup>-1</sup> (12 - 2  $\mu$ m). Upon expanding the spectral range under consideration to 400–10 000 cm<sup>-1</sup> (25 - 1  $\mu$ m) the observed performance gracefully degrades to  $R \le 0.02$  at the band edges. In the implementation described, a high-thermal-conductivity metallic substrate is textured with a pyramidal tiling and subsequently coated with a thin lossy dielectric coating that enables high absorption and thermal uniformity across the target. The resulting target assembly is lightweight, has a low-geometric profile, and has survived repeated thermal cycling from room temperature to  $\sim$ 4 K. Basic design considerations, governing equations, and test data for realizing the structure described are provided. The optical properties of selected absorptive materials—Acktar Fractal Black, Aeroglaze Z306, and Stycast 2850 FT epoxy loaded with stainless steel powder—are characterized and presented. [http://dx.doi.org/10.1063/1.4871108]

#### I. INTRODUCTION

Absorptive targets and light traps find widespread use in flux calibration, termination of residual reflections in optical systems, and establishing the zero point in reflectance spectrometry. Approximating a near-ideal absorber in a finite volume presents a challenge for such applications—the absence of reflected light or how "black" an object appears is not a unique material property or treatable as a boundary condition within the framework of electromagnetic theory<sup>2</sup>—but an object's reflectance and absorptance are intimately tied to the surface's material properties and the underlying geometry of the structure. Thus to implement a low reflectance calibrator design, the absorber material's dielectric and magnetic properties need to be either known or experimentally determined in order to suitably tailor the target's geometry.

At microwave through sub-millimeter wavelengths, precision absorptive standards find use in radiometric flux calibration and low reflectance targets for remote sensing, astrophysics, 5.6 and other metrology applications. Elements of the design techniques used in these microwave structures can be applied to realize compact, light-weight, and low-reflectance absorbers while maintaining the overall manufacturability for far infrared use at cryogenic temperatures. In this work, a calibration target for cryogenic applications is explored and its performance is described in detail. The design, manufacture, and optical characterization of the target structure and selected coatings of potential interest are summarized in Secs. II–IV, respectively.

#### II. TARGET DESIGN CONSIDERATIONS

Low reflectance targets have been realized from individual geometric structures<sup>5,6</sup> or tilings<sup>8</sup> of absorbing dielectric cones or pyramids. Wedges, rulings,<sup>4,7,9</sup> and Brewster angle light traps<sup>10–13</sup> have also been employed in instrumentation

settings where only a single polarization state is of interest. Through the microwave waveband, the absorptive coating's finish is typically smooth compared to the freespace wavelength,  $\lambda_o$ , and the resulting reflections from surfaces making up the target's facets are largely specular. In this limit, the effective reflectance for a pyramidal wedge with vertex angle,  $\alpha$ , can be estimated from geometric optics considerations,  $R \approx R_o^{\pi/\alpha}$ . Here, the single bounce reflection at the angle of incidence,  $R_o$ , is estimated from the Fresnel coefficients. <sup>14</sup> In effect, the geometry can be used to limit the impact of imperfect a priori knowledge of the coating properties on the final target absorptance at cryogenic temperatures. 15 As a result, the key geometric parameters that control the reflectance in this limit are the details of the pyramid's points, grooves at the structure base, and the number of reflections within the structure. 8 Ideally, these features would be perfectly sharp; however, in practice a finite length scale is achievable and set by fabrication tolerances and the coating morphology.

In the limit where the freespace wavelength,  $\lambda_o$ , is less than twice the tiling period, p, the target's response can be computed by treating the structure as an artificial dielectric adiabatic impedance taper. <sup>16,17</sup> As the wavelength is further increased beyond the length scale of the taper, h, the target's surface appears smooth and incident fields evanescently propagate into the structure. <sup>18</sup> As a result, the radiation is largely specularly reflected from the structure. More generally, these arguments can be extended to treat the diffuse reflectance from the surface roughness in the design process if sufficient details are available. <sup>19</sup> Here, the experimentally observed diffuse reflectance is taken as a property of the coating that in turn determines acceptable substrate geometry and the short wavelength limit of the structure's response.

The use of an appropriate coating with a controlled and sufficient thickness to absorb the incident radiation is also an important consideration. Here, the material is assumed to be non-magnetic and characterized as a function of frequency by a complex dielectric function,  $\epsilon_r = \epsilon'_r + i \cdot \epsilon''_r$ , relative to the permittivity of freespace. The thickness of the target's

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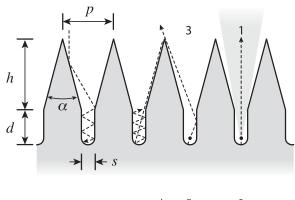
lossy dielectric layer employed must be greater than the material's field penetration length scale,  $\delta = \lambda_o/\pi \sqrt{2\epsilon_r''}$ , to suitably attenuate signals before reflection from the underlying substrate. By design, the substrate is optically thick and the absorptance is related to the reflectance by A=1-R. Consideration of the Fresnel coefficient magnitude at the surface of the lossy dielectric reveals that a lower index coating leads to lower insertion loss per bounce and thus has greater absorptance from the overall structure in the long wavelength limit. In the infrared, a reduction in coating thickness is desirable as long as the coating maintains structural integrity and minimal thermal gradients.

With appropriate attention to detail a variety of coatings could potentially be employed in this application<sup>20–22</sup> to achieve the desired response. The service temperature of epoxies, paints, and deposited films are typically limited by the material's adhesion and thermal match to the substrate material over the temperature range of interest. For high temperature applications, chemical breakdown, and aging may also require more careful consideration. At cryogenic temperatures typical anodized coating (thermal conductivity,  $\kappa \sim 1$  W/m/K) with thickness,  $\delta x < 100 \,\mu\text{m}$ , produces a thermal gradient across the absorber,  $\Delta T/q \simeq \delta x/\kappa \sim 0.1 \text{ K/mW}$ per unit incident flux, q. For our application this consideration is subdominant to the  $\sim 0.2$  K thermometry calibration uncertainty at 4 K. Cooling or achieving precise knowledge of the target temperature becomes a greater concern in going from bulk to low-volume-filling-fraction metal or semiconducting foams<sup>23,24</sup> and semi-metallic nanotubes<sup>25</sup> in the long wavelength limit due to such coatings' lower thermal conductivity. Additional considerations for space qualification and other low-outgassing applications can be found in the literature. 26-28 Based upon prior experience for the target implementation described here we use a custom specified thickness of Acktar Fractal Black<sup>29</sup> (Acktar FB) amenable to a variation on the vendor's standard (11 µm thickness coating) process for the absorptive coating.

#### III. SAMPLE PREPARATION AND FABRICATION

In implementing the absorptive calibrator design we concentrate on the use of simple manufacturing approaches to achieve the desired response in the infrared. The underlying approach adopted was to use a lossy dielectric on a highly conductive substrate to achieve high absorptance. See Figure 1 for target geometry and ray-trace summary. An aluminum alloy, 6061-T6, is chosen as the target's substrate material. It is durable, easy to machine, and its high thermal conductance minimizes thermal gradients and enables knowledge of the device temperature. The design described is driven by a need for the target to have a total thickness <6 mm. A total calibrator diameter,  $d_{cal} = 25$  mm, was chosen for compatibility with our existing integrating sphere configuration and for use as a low-reflectance reference in a cryogenic filter wheel assembly. A pyramidal tiling of the surface is used to limit the total required thickness of the target relative to a single pyramid.

The tips of the pyramids (radius  $<\!20\,\mu\text{m}$ ) used in the tiling were readily achievable by wire electrostatic discharge



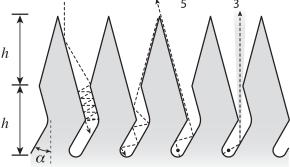


FIG. 1. The shaded gray regions in each figure represent the high thermal conductivity and optically thick patterned metallic substrate. The solid black lines are the lossy dielectric layer with thickness greater than the electrical penetration length,  $\delta$ . Representative geometric ray traces are indicated with dashed lines. The minimal number of bounces occur when rays go through the focus formed by the slot's radius (black dots, right). For a simple slot, the single bounce observation angle is controlled by the slot width and depth (shaded region, upper right). The overall absorptance can be improved through the addition of a "dog leg" in the slot as shown (lower right).

milling (EDM) with a full apex angle,  $\alpha=\pi/6$ , and are comparable to the surface roughness features of the coatings employed. A finite width slot was substituted for the sharp groove found at the base of the pyramid in a traditional microwave target design. Slots  $s\simeq 250\,\mu\mathrm{m}$  in width,  $d=940\,\mu\mathrm{m}$  deep, on a pitch  $p=1370\,\mu\mathrm{m}$  were defined by the EDM wire diameter. The finite slot width relative to the pitch limits the achievable reduction in the pyramidal tiling thickness by a factor of  $2d_{\rm cal}h/p^2\sim 50$ . The resulting patterned features in the aluminum substrate are less than half of the component's thickness and the remainder of the substrate is used to provide low lateral thermal conduction, mounting points for cooling, and thermometry.

A cross-section of the fabricated pyramidal target's pyramidal tiles is shown in the upper panel of Figure 1. In the specular limit, this geometry has >6 reflections for 80% of rays incident on the pyramidal structure. The ratio of the slot width over the depth limits number of rays with a single bounce and sets the target reflectance in this structure. By tipping the target to  $\sim 10^\circ$  from the normal the specular reflectance contribution from the bottom of the slot can be reduced. Alternatively, a "dog leg" at the bottom of the channel as shown in the lower panel of Figure 1 can be used to improve the absorptance by increasing the number of bounces ( $N \ge 3$ ) in the slot region and effectively reduces the angular extent of the caustic arising from the slot's finite radius. If greater volume

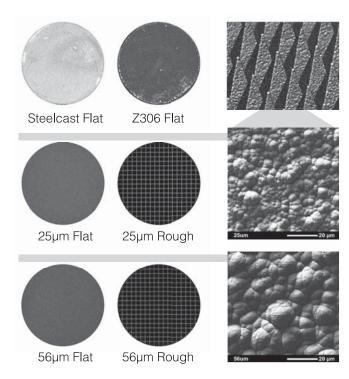


FIG. 2. Images of substrate samples: (top-left) 300  $\mu$ m Steelcast (Stycast 2850 FT epoxy loaded with 30% stainless steel powder by volume) sample on aluminum flat; (top-center) 40  $\mu$ m Aeroglaze Z306 black paint on Steelcast sample for enhanced long and short wavelength response; (top-right) Detail of Acktar FB coated calibrator on a substrate roughened with a pyramidal tiling viewed approximately  $\pi/4$  from vertical; (middle-left) 25  $\mu$ m Acktar FB witness coating on aluminum flat; (middle-center) 25 μm Acktar FB calibrator roughened with a pyramidal tiling; (middle-right) SEM image of 25  $\mu$ m Acktar FB coating; (lower-left) 56  $\mu$ m Acktar FB witness coating on aluminum flat; (lower-center) 56 µm Acktar FB calibrator roughened with a pyramidal tiling; and (lower-right) SEM image of  $56 \mu m$  Acktar FB coating. The images of the "flat" and "rough" samples (i.e., left- and centercolumns) were taken as a single exposure under an intense visible light source from approximately normal incidence. The logarithmic brightness scale has been uniformly adjusted in all digital images of the samples to convey the full range of the intensity recorded. The gray scale in the scanning electron microscope images (i.e., middle- and lower-right panels of figures) are uncorrected and convey the typical morphology of the Acktar FB samples investigated.

for the target is available, the influence of the finite slot width can be further reduced by using an adiabatically curved path for the slot to geometrically hide its termination from view.

The pyramidally tiled absorber structures and optical witness flats were subsequently coated with Acktar FB $^{29}$  via Physical Vapor Deposition (PVD) for spectral characterization. Two coating thicknesses,  ${\sim}25~\mu m$  and  ${\sim}50~\mu m$ , were specified in order to insure sufficient coating absorption. See Figure 2 for detailed digital images and a summary of the samples optically characterized in this study.

In addition, a flat sample was coated with "Steelcast," a dielectric mixture comprised of Emerson Cumming 2850 FT epoxy with a 30% volume filing fraction of 8  $\mu$ m stainless steel powder. The sample was then hand lapped to  $\sim 300~\mu$ m thickness with 600 grit sand paper. The relative dielectric permittivity and magnetic permeability of this dielectric mixture are  $\epsilon_r \approx 10.6 + 2.1i$  and  $\mu_r \approx 1$ , respectively. These properties were verified for the mixture employed by using a witness coupon as a waveguide Fabry-Perot resonator at 1 cm<sup>-1</sup>. To explore improving the Steelcast low frequency reflectance, a

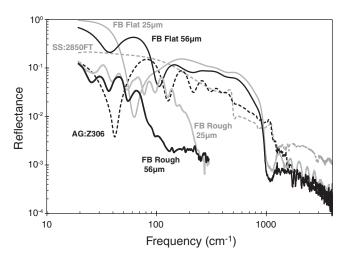


FIG. 3. Measured specular reflectance. Dashed lines denote the stainless steel loaded 2850 FT epoxy (gray) and the Aeroglaze Z306 paint (black); solid lines are the Acktar FB flat 25  $\mu m$  (gray) and 56  $\mu m$  (black) coatings; and the bold lines are the Acktar FB 25  $\mu m$  (gray) and 56  $\mu m$  (black) coatings on substrates roughened with a pyramidal tiling. The coatings are deposited on aluminum substrates.

sample was coated with an anti-reflection layer realized from a  ${\sim}40~\mu{\rm m}$  thick layer of Aeroglaze Z306 paint<sup>31</sup> with an RMS surface roughness,  $\sigma \simeq 2~\mu{\rm m}$ . In modeling this multi-layer dielectric coating<sup>20,32</sup> an effective relative permittivity,  $\epsilon_r \approx 2.6 + 0.6i$ , is estimated for the paint on the epoxy substrate using the measured sample parameters and the specular reflectance <500 cm<sup>-1</sup> (e.g., see the spectrum in Fig. 3 labeled "AG:Z306").

This combination of materials finds use in targets and baffles in cryogenic applications below  $\sim 400 \text{ cm}^{-1}$  and provides an informative point of reference to the microwave structures motivating the calibrator design presented. Loaded epoxy structures are typically fabricated by building up the absorber in layers via spraying, molding, or with great difficultly direct machining on a textured substrate. The achievable tolerances with these approaches have higher variability than is observed with PVD coating on an identical patterned substrate in the design's critical tip and slot regions. As a result, a significantly greater calibrator volume is required to insure identical optical performance in practice and the structure's thermal time constant and gradients are correspondingly increased. While relatively high index loaded epoxies provide a path in the microwave to minimize the coating thickness and mass, in going to the infrared, lower index lossy thin films, and conversion coatings more naturally serve this purpose. Realizing the desired substrate geometry presents the primary challenge to be address in this path to calibrator volume reduction.

#### IV. FTS CHARACTERIZATION

Infrared and sub-millimeter spectra were obtained using a Bruker IFS125hr Fourier transform spectrometer (FTS). This instrument is equipped with internal DTGS (Deuterated Triglycine Sulfate) detectors for use in the near-through far-infrared and an external LHe-cooled bolometer with two interchangeable filters for far-infrared and sub-millimeter

wavelengths (100–700 and 5–100 cm<sup>-1</sup>, respectively). An array of beamsplitters is used for the different spectral regions; the following were used here: Si on CaF2 (near-IR), Ge on KBr (mid-IR), 6  $\mu$ m multi-layer Mylar (Bruker T222/2), and 75  $\mu$ m Mylar (far-infrared and sub-millimeter). The FTS is equipped with three internal light sources: tungsten (near-IR), SiC Globar (mid-IR), and a mercury arc lamp (far-IR). Reflectance spectra for the samples were measured in each band and then merged together to form a composite spectrum. Spectra are recorded *in vacuo* ( $\sim$ 0.3 mbar).

Specular reflectance data were recorded over the frequency range of  $20\text{--}5000\,\mathrm{cm^{-1}}$  at a  $2\,\mathrm{cm^{-1}}$  resolution using a Pike Technologies accessory (Part No. SKU2084). Light from the spectrometer takes a W-path through the accessory, which has collimating optics and a fixed  $10^\circ$  angle of incidence on the sample. In this set-up, the sample's reflectance spectrum is obtained by taking the ratio of the light intensity detected from the sample being measured to the intensity of a gold coated flat provided with the Pike accessory. See Figure 3 for the observed specular reflectance spectra.

Diffuse reflectance spectra were recorded using a Bruker integrating sphere accessory (Part No. A562G/Q) over a frequency range of  $250-10\,000\,\mathrm{cm}^{-1}$ . This accessory is equipped with a side-mounted DTGS detector used primarily for mid- and near-infrared wavelengths and is supplied with a diffusely reflecting gold reference. Internal absorptive baffles are used to shield the detector from receiving light reflected directly from the samples. In order to mitigate the typical noise of measurements with the integrating sphere apparatus, and to help enhance the broad structure of the features (e.g., see the Acktar FB 25  $\mu$ m flat coating between  $\sim$ 1000– 3000 cm<sup>-1</sup>), spectra were recorded using spectral resolutions of 8 and 16 cm<sup>-1</sup>. In addition to the Bruker integrating sphere accessory, diffuse reflectance was measured for the roughened Acktar FB 56 μm coating using a custom cryogenic integrating sphere.<sup>33</sup> This apparatus is equipped with a Si bolometer element (4.2 K) and an internal temperature-controlled sample wheel. Far infrared spectra between 100 and 650 cm<sup>-1</sup> with 8 cm<sup>-1</sup> resolution were measured using this apparatus mounted to a Bruker IFS 113v FTS equipped with an internal SiC Globar source and 3.5 and 12  $\mu$ m Mylar beamsplitters. See Figure 4 for observed diffuse reflectance spectra.

#### V. DISCUSSION

From the observed diffuse reflectance of the witness samples in Figure 4 one notes both Acktar FB coatings have a slightly lower reflection than Aeroglaze Z306 paint in the spectral range of interest. The calibrator constructed with 56  $\mu$ m of Acktar FB was noted to have the best overall optical performance of the structures tested with an observed diffuse reflectance of R < 0.003 between 800 and 4800 cm<sup>-1</sup> (12–2  $\mu$ m). The minimum in the diffuse reflectance,  $R \simeq 0.002$ , occurs at 2000 cm<sup>-1</sup> (5  $\mu$ m). This represents a reduction in the diffuse reflectance by a factor of  $\sim$ 50 over a simple flat with an identical coating.

The low frequency response of the calibrator samples is set by the pyramidal tiling's pitch and finite coating thickness.

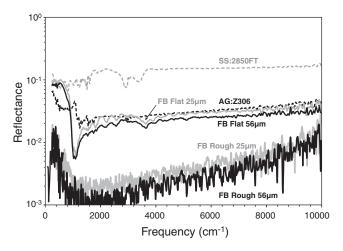


FIG. 4. Measured diffuse reflectance. Dashed lines denote the stainless steel loaded 2850 FT epoxy (gray) and the Aeroglaze Z306 paint (black); solid lines are the Acktar FB flat 25  $\mu m$  (gray) and 56  $\mu m$  (black) coatings; and the bold lines are the Acktar FB 25  $\mu m$  (gray) and 56  $\mu m$  (black) coatings on substrates roughened with a pyramidal tiling. All coatings are deposited on aluminum substrate disks. The room temperature Bruker integrating sphere (300 K, >250 cm^{-1}) and the OPASI-T cryogenic reflectometer (4.2 K, <650 cm^{-1}) data were calibrated with a polished aluminum reference. The measured data from the two instrument configurations are in agreement in the region of shared spectral overlap.

At greater than  $\sim \! 10^\circ$  from normal incidence the gray lines seen in Figure 2 that originate from spectral reflectance off the bottom of the calibrator's slots are no longer discernible. This artifact can be further mitigated by gradually turning the slot and removing its endpoint from view as indicated in Figure 1. Incorporation of this fabrication detail is recommended to improve the overall emittance. The calibrator is broadband, mechanically robust, and has survived multiple cool downs to  $\sim \! 4$  K. Due to the structure's low reflectance it has been adopted as the preferred absorptive light trap for general laboratory use.

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<sup>&</sup>lt;sup>1</sup>L. Palchetti, G. Bianchini, and F. Castagnoli, "Design and characterization of black-body sources for infrared wide-band Fourier transform spectroscopy," Infrared Phys. Technol. **51**, 207–215 (2008).

<sup>&</sup>lt;sup>2</sup>A. Sommerfeld, *Optik Vorlesungen über Theoretische Physik Band 4* (Dieterichsche Verlagsbuchhandlung, Munich, 1950), pp. 265–266.

<sup>&</sup>lt;sup>3</sup>H. V. de Hultst, *Light Scattering by Small Particles* (John Wiley and Sons, New York, 1957), p. 183.

<sup>&</sup>lt;sup>4</sup>M. C. Gaidis, M. S. Anderson, and D. G. Harding, "Calibration target for far-infrared spaceborne applications," Proc. SPIE 3795, 348–356 (1999).

<sup>&</sup>lt;sup>5</sup>H. P. Gush and M. Halpern, "Cooled submillimeter Fourier transform spectrometer flown on a rocket," Rev. Sci. Instrum. **63**, 3249–3260 (1992).

<sup>&</sup>lt;sup>6</sup>J. C. Mather, D. J. Fixsen, R. A. Shafer, C. Mosier, and D. T. Wilkinson, "Calibrator design for the COBE Far-Infrared Absolute Spectrophotometer (FIRAS)," Astrophys. J. **512**, 511–520 (1999).

- <sup>7</sup>S. Janz, D. A. Boyd, and R. F. Ellis, "Reflectance characteristics in the submillimeter and millimeter wavelength region of a vacuum compatible absorber," Int. J. Infrared Millimeter Waves **8**, 627–635 (1987).
- <sup>8</sup>D. J. Fixsen, E. J. Wollack, A. Kogut, M. Limon, P. Mirel, J. Singal, and S. M. Fixsen, "Compact radiometric microwave calibrator," Rev. Sci. Instrum. 77, 064905 (2006).
- <sup>9</sup>A. Kogut, E. Wollack, D. J. Fixsen, M. Limon, P. Mirel, S. Levin, M. Seiffert, and P. M. Lubin, "Design and calibration of a cryogenic blackbody calibrator at centimeter wavelengths," Rev. Sci. Instrum. 75, 5079–5083 (2004).
- <sup>10</sup>P. B. Elterman, "Brewster angle light trap," Appl. Opt. **16**, 2352–2352 (1977).
- <sup>11</sup>E. J. Breneman, "Light trap of high performance and simple construction," Appl. Opt. 20, 1118–1119 (1981).
- <sup>12</sup>F. J. J. Clarke and J. A. Compton, "Correction methods for integrating-sphere measurement of hemispherical reflectance," Color Res. Appl. 11, 253–262 (1986).
- <sup>13</sup> A. LaRocca, *The Infrared Handbook*, edited by W. L. Wolfe and G. J. Zissis (Infrared Information and Analysis Center, Environmental Research Institute University of Michigan, 1985), Chap. 2, pp. 2–9.
- <sup>14</sup>B. Carli, "Design of a blackbody reference standard for submillimeter region," IEEE Trans. Microwave Theory Tech. 22, 1094–1099 (1974).
- <sup>15</sup>E. J. Wollack, D. J. Fixsen, A. Kogut, M. Limon, P. Mirel, and J. Singal, "Radiometric-waveguide calibrators," IEEE Trans. Instrum. Meas. 56, 2073–2078 (2007).
- <sup>16</sup>E. Meyer and H. Severin, "Absorption devices for electromagnetic centimeter waves and their acoustic analogs," Z. Angew. Phys. 8, 105–114 (1956).
- <sup>17</sup>E. F. Kuester and C. Holloway, "A low-frequency model for wedge or pyramid absorber arrays-I: theory," IEEE Trans. Electromagn. Compat. 36, 300–306 (1994).
- <sup>18</sup>B. Baekelandt, F. Olyslager, and D. De Zutter, "Low-frequency reflectivity approximation for two- and three-dimensional EM absorbers," IEEE Trans. Electromagn. Compat. 41, 354–360 (1999).
- <sup>19</sup>R. E. Bedford and C. K. Ma, "Emissivities of diffuse cavities: Isothermal and nonisothermal cones and cylinders," J. Opt. Soc. Am. **64**, 339–349 (1974).
- <sup>20</sup>S. Smith, "Specular reflectance of optical-black coatings in the far infrared," Appl. Opt. 23, 2311–2326 (1984).

- <sup>21</sup>D. Betts, F. Clarke, L. Cox, and J. Larkin, "Infrared reflection properties of five types of black coating for radiometric detectors," J. Phys. E: Sci. Instrum. 18, 689–696 (1985).
- <sup>22</sup>M. Persky, "Review of black surfaces for space-borne infrared systems," Rev. Sci. Instrum. **70**, 2193–2217 (1999).
- <sup>23</sup>D. Advena, V. Bly, and J. Cox, "Deposition and characterization of farinfrared absorbing gold black films," Appl. Opt. 32, 1136–1144 (1993).
- <sup>24</sup>W. Becker, R. Fettig, and W. Ruppel, "Optical and electrical properties of black gold layers in the far infrared," Infrared Phys. Technol. 40, 431–445 (1999).
- <sup>25</sup>M. A. Quijada, J. G. Hagopian, S. Getty, R. E. Kinzer, Jr., and E. J. Wollack, "Hemispherical reflectance and emittance properties of carbon nanotubes coatings at infrared wavelengths," in *Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series* (SPIE, 2011), Vol. 8150.
- <sup>26</sup>L. Kauder, "Spacecraft thermal control coatings references," Technical Report TP-2005-212792 (NASA Goddard Space Flight Center, 2006).
- <sup>27</sup>T. Kralik and D. Katsir, "Black surfaces for infrared, aerospace, and cryogenic applications," in *Society of Photo-Optical Instrumentation Engineers* (SPIE) Conference Series (SPIE, 2009), Vol. 7298.
- <sup>28</sup>Y. Salomon, N. Sternberg, I. Gouzman, G. Lempert, E. Grossman, D. Katsir, R. Cotostiano, and T. Minton, "Qualification of Acktar Black coatings for space application," in *Proceedings of the International Symposium on Materials in a Space Environment, Aix-En-Provence, France*, 2009.
- <sup>29</sup> Acktar Ltd., "Acktar Fractal Black," Leshem St., Entrance A, P.O.B. 8643, Kirvat-Gat 82000.
- <sup>30</sup>E. J. Wollack, D. J. Fixsen, R. Henry, A. Kogut, M. Limon, and P. Mirel, "Electromagnetic and thermal properties of a conductively loaded epoxy," Int. J. Infrared and Millimeter Waves 29, 51–61 (2008).
- <sup>31</sup>Lord Corporation Chemical Products, "Aeroglaze Z306 Flat Black Absorptive Polyurethane Low-Outgassing Paint," 2000 West Grandview Blvd., P.O. Box 10038, Erie, PA 16514-0038.
- <sup>32</sup>P. Yeh, Optical Waves in Layered Media (John Wiley & Sons, New York, 1988), Chap. 5, pp. 102–117.
- <sup>33</sup>S. A. Rinehart, D. J. Benford, G. Cataldo, E. Dwek, R. Henry, J. R. E. Kinzer, J. Nuth, R. Silverberg, C. Wheeler, and E. Wollack, "Measuring the optical properties of astrophysical dust analogues: instrumentation and methods," Appl. Opt. 50, 4115–4123 (2011).