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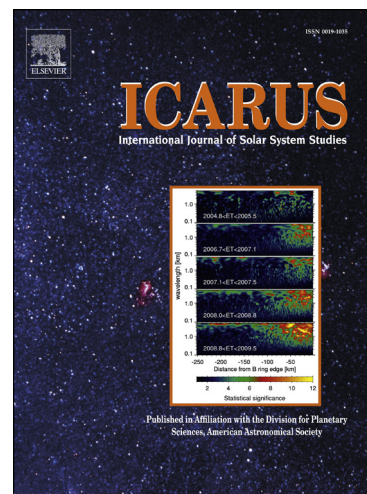
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1 Abstract

2 The empirical ‘volcano-scan’ atmospheric correction is widely applied to Martian near
3 infrared CRISM and OMEGA spectra between ~1000 and ~2600 nm to remove prominent
4 atmospheric gas absorptions with minimal computational investment. This correction method
5 employs division by a scaled empirically-derived atmospheric transmission spectrum that is
6 generated from observations of the Martian surface in which different path lengths through the
7 atmosphere were measured and transmission calculated using the Beer-Lambert Law.
8 Identifying and characterizing both artifacts and residual atmospheric features left by the
9 volcano-scan correction is important for robust interpretation of CRISM and OMEGA volcano-
10 scan corrected spectra. In order to identify and determine the cause of spectral artifacts
11 introduced by the volcano-scan correction, we simulated this correction using a multiple
12 scattering radiative transfer algorithm (DISORT). Simulated transmission spectra that are
13 similar to actual CRISM- and OMEGA-derived transmission spectra were generated from
14 modeled Olympus Mons base and summit spectra. Results from the simulations were used to
15 investigate the validity of assumptions inherent in the volcano-scan correction and to identify
16 artifacts introduced by this method of atmospheric correction. We found that the most prominent
17 artifact, a bowl-shaped feature centered near 2000 nm, is caused by the inaccurate assumption
18 that absorption coefficients of CO₂ in the Martian atmosphere are independent of column
19 density. In addition, spectral albedo and slope are modified by atmospheric aerosols. Residual
20 atmospheric contributions that are caused by variable amounts of dust aerosols, ice aerosols, and

21 water vapor are characterized by the analysis of CRISM volcano-scan corrected spectra from the
22 same location acquired at different times under variable atmospheric conditions.

23 **1 Introduction**

24 Images of the Martian surface acquired by the NASA Mars Reconnaissance Orbiter
25 (MRO) Compact Reconnaissance Imaging Spectrometer for Mars (CRISM) [Murchie *et al.*,
26 2007] and the Mars Express Observatoire pour la Minéralogie, l'Eau, les Glaces et l'Activité
27 (OMEGA) [Bibring *et al.*, 2004; 2005] measure solar light that was attenuated and scattered as it
28 traversed down through the Martian atmosphere, interacted with the surface, and traversed up
29 through the atmosphere. Therefore, each spectrum from standard CRISM and OMEGA
30 observations contains contributions from atmospheric gases (e.g., CO₂, CO, and H₂O),
31 atmospheric aerosols (e.g., dust and water ice), and the surface. Atmospheric gas contributions
32 dominate the spectrum at wavelengths that CO₂ absorbs (Figure 1).

33 The empirical 'volcano-scan' correction [Bibring *et al.*, 1989; Langevin *et al.*, 2005;
34 McGuire *et al.*, 2009] is widely applied to CRISM and OMEGA near infrared (NIR) spectra
35 between ~1000 and ~2600 nm to remove prominent atmospheric gas absorptions with minimal
36 computational investment. The volcano-scan correction method employs division by a scaled
37 atmospheric transmission spectrum that is generated from observations of the Martian surface in
38 which different path lengths through the atmosphere were measured. Low and high altitude
39 spectra acquired over the base and summit of the Olympus Mons volcano were used, giving the
40 correction its name. Transmission is calculated empirically using the Beer-Lambert Law, as
41 detailed in section 2.2.

42 Although the volcano-scan correction removes prominent gas absorptions to first order,
43 closer inspection reveals that corrected spectra exhibit spurious features in areas of strong gas
44 absorption. The most prominent artifact is a bowl-shaped feature that overlaps with the CO₂
45 triplet centered near 2000 nm (Figure 1). The volcano-scan correction can be applied
46 automatically to CRISM images using publically available CRISM Analysis Tools (CAT)
47 software released through the Planetary Data System (PDS). The current version of CAT at this
48 time (version 7.2.1) includes an option to empirically correct the 2000 nm bowl-shaped artifact
49 that is evident in volcano-scan corrected images on a pixel by pixel basis.

50 The occurrence of the apparent bowl-shaped artifact at 2000 nm could have several
51 causes. There are three major assumptions implicit in deriving empirical transmission spectra: 1)
52 surface contributions to the low and high altitude spectra used to create transmission spectra are
53 equivalent and therefore cancel out, 2) aerosol contributions to low and high altitude spectra can
54 be ignored both within and outside of gas absorption lines and therefore empirical transmission
55 spectra contain molecular absorption only, 3) absorption coefficients of CO₂ in the Martian
56 atmosphere are independent of column density and absorption strength therefore scales
57 exponentially with column density.

58 In order to identify and determine the cause of artifacts introduced by the volcano-scan
59 correction method, we simulated this correction using a multiple scattering radiative transfer
60 algorithm. Discrete Ordinate Radiative Transfer (DISORT) modeling allows for the explicit
61 treatment of aerosol, gas, and surface contributions simultaneously [*Stamnes et al.*, 1988;
62 *Thomas and Stamnes*, 2002]. The DISORT radiative transfer model was used to calculate
63 modeled high and low altitude spectra similar to Martian spectra of Olympus Mons. Simulated
64 transmission spectra were derived from these modeled spectra using the same method that was

65 used to produce CRISM- and OMEGA-derived transmission spectra. Generating simulated
66 transmission spectra in this manner and applying them to modeled Martian spectra using the
67 volcano-scan correction method allows the variables related to the three major assumptions
68 described above to be controlled and analyzed.

69 Because the volcano-scan correction is designed to remove absorptions from atmospheric
70 CO₂, this correction does not specifically address contribution from other atmospheric species,
71 including dust and ice aerosols and gaseous water vapor. CO absorption is minor (Figure 1, gray
72 arrow) and will not be considered. Aerosol and water vapor atmospheric contributions are
73 particularly important because they are spatially and temporally variable [e. g., *Smith et al.*,
74 2008] and produce noticeable features in CRISM and OMEGA spectra [e. g., *Smith et al.*, 2009].
75 Dust aerosol contributions affect spectral slope and amplitude and ice aerosols also have distinct
76 absorption features in the NIR. In addition, aerosol scattering within gas absorption lines alters
77 absorption features. Atmospheric water vapor absorptions that occur in empirically derived
78 transmission spectra cannot be scaled separately from CO₂ absorptions in the transmission
79 spectrum that is scaled during the volcano-scan correction. Failure to explicitly address
80 atmospheric water vapor can cause under- and over-corrected water vapor features in volcano-
81 scan corrected spectra.

82 Identifying and characterizing both artifacts and residual atmospheric features left by the
83 volcano-scan correction is important for proper interpretation of CRISM and OMEGA volcano-
84 scan corrected spectra. Results from the simulated volcano-scan correction are used to
85 investigate the validity of assumptions inherent in deriving and applying empirical transmission
86 spectra and to identify and determine the causes of artifacts (e.g., 2000 nm bowl-shape)
87 introduced by this atmospheric correction method. In addition, we characterize spectral features

88 that result from dust aerosols, ice aerosols, and water vapor that are not specifically addressed by
89 the volcano-scan correction. Variations in these species cause volcano-scan corrected spectra to
90 exhibit variable residual features depending on atmospheric conditions at the time of image
91 acquisition. These residual features and artifacts are highlighted by comparison of CRISM
92 volcano-scan corrected spectra from the same location on Mars but acquired at different times
93 under variable atmospheric conditions.

94 **2 Background**

95 **2.1 CRISM and OMEGA Datasets**

96 The CRISM spectrometer has short (S) and long (L) wavelength detectors that operate
97 between 360 to 1053 nm and 1003 to 3920 nm, respectively [Murchie *et al.*, 2007; 2009].
98 CRISM is capable of acquiring hyperspectral images with 544 channels spaced ~7 nm apart.
99 CRISM is mounted on a gimbal platform that allows ground tracking of a target to acquire full
100 resolution targeted (FRT) images with up to 20 m/pixel spatial resolution. The CRISM
101 spectrometer utilizes 2-dimensional detector arrays that cause slight column dependencies in
102 CRISM spectral data, including a smooth variation in the central wavelength position of each
103 channel [Murchie *et al.*, 2007]. Image rows are built up as a result of along track spacecraft
104 motion. CRISM I/F (where I is the radiance at sensor divided by the solar irradiance, F, at the
105 top of the Martian atmosphere divided by π) spectra available from the PDS ([http://pds-
106 geoscience.wustl.edu](http://pds-geoscience.wustl.edu)) were calibrated as described by Murchie *et al.* [2007, 2009].

107 Atmospheric transmission files derived from CRISM hyperspectral, non-gimbaled FFC
108 images acquired at the base and summit of Olympus Mons are contained in Atmospheric
109 Transmission (AT) Calibration Data Record (CDR) files that are available through the PDS.

110 Because a relatively small (< 1.5 nm) temperature dependent wavelength shift is observed in
111 CRISM L detector spectra [*Smith et al.*, 2009], AT CDRs for several different wavelength shifts
112 were derived. A slope correction was applied to transmission spectra in the AT CDR files in an
113 effort to normalize the shapes of the transmission spectra and ensure that continuum values are
114 near unity. The most current version of CAT at this time (version 7.2.1) includes an option to
115 empirically correct the 2000 nm bowl-shaped artifact that is evident in volcano-scan corrected
116 images on a pixel by pixel basis (algorithm developed by Frank Morgan and formulation present
117 in CAT v7.2.1 code). The 2000 nm bowl-shaped artifact is defined relative to a straight line
118 continuum derived from a reference observation. The same low altitude Olympus Mons spectra
119 used to generate empirical transmission spectra are used to define the artifact (to minimize
120 temperature dependent wavelength shift and surface variations) for wavelengths between ~ 1760
121 and 2240 nm. The artifact correction spectrum is determined by subtracting the volcano-scan
122 corrected reference spectrum from the continuum fit to that spectrum. After the volcano-scan
123 correction is applied, the artifact correction spectrum is multiplicatively scaled to match the
124 artifact expressed in each pixel. The 2000 nm bowl-shaped artifact is then corrected by adding
125 the scaled artifact correction to the volcano-scan corrected spectrum, for each pixel in the image.

126 The OMEGA spectrometer [*Bibring et al.*, 2004] covers the wavelength range from 350
127 to 5100 nm in 352 channels using VNIR (350 to 1070 nm), C (930 to 2700 nm), and L (2530 to
128 5100 nm) detectors. C detector channels have a wavelength spacing of ~ 14 nm [*Bibring et al.*
129 2004, 2005]. Images vary in spatial resolution depending on spacecraft altitude and are built up
130 by across track scanning and along track spacecraft motion. OMEGA data lack the column
131 dependent issues inherent in CRISM data. OMEGA images as well as software to calibrate and
132 process the images to units of I/F as described by *Bibring et al.* [2004, 2005], are available

133 through the PDS at <http://pds-geoscience.wustl.edu>. An OMEGA-derived atmospheric
 134 transmission spectrum for use with the empirical volcano-scan correction is not publically
 135 available but can be calculated from Olympus Mons observations using methods described in the
 136 following section.

137 2.2 Empirical Volcano-Scan Correction

138 The empirical volcano-scan correction method relies on empirically derived transmission
 139 spectra to remove CO₂ absorption features from spectra of the Martian surface [*Bibring et al.*,
 140 1989; *Langevin et al.*, 2005; *McGuire et al.*, 2009]. Transmission spectra are generated from
 141 observations in which different path lengths through the atmosphere were measured. High
 142 altitude spectra were acquired at the summit of Olympus Mons (~20 km above datum) and low
 143 altitude spectra at its base (~0 km) (Figure 2a). Transmittance between the base and summit of
 144 Olympus Mons is determined using the Beer-Lambert Law (equation 1) integrated with the
 145 assumption that the absorption coefficient, k , is constant with height in the atmosphere (equation
 146 2).

$$147 \quad T(\lambda) = I_{\lambda}(s_{0 \rightarrow 2}) / I_{\lambda}(s_{0 \rightarrow 1}) = e^{-\int_{s_1}^{s_2} k ds} \quad (1)$$

$$148 \quad T(\lambda) = I_{\lambda}(s_{0 \rightarrow 2}) / I_{\lambda}(s_{0 \rightarrow 1}) = e^{-k s_{1 \rightarrow 2}} = e^{-\tau} \quad (2)$$

$$149 \quad -\ln(T(\lambda)) = k s_{1 \rightarrow 2} = \tau \quad (3)$$

150
 151 In equations 1 through 3, T is transmittance, I_{λ} is equivalent to CRISM or OMEGA I/F at
 152 wavelength (λ), $s_{0 \rightarrow 1}$ is the sum of the incoming and outgoing path lengths from the top of the
 153 atmosphere (s_0) to height s_1 within the atmosphere, $s_{0 \rightarrow 2}$ is the sum of the incoming and outgoing

154 path lengths from the top of the atmosphere (s_0) to height s_2 within the atmosphere, and τ is the
155 total extinction. Note, for nonzero solar incidence (i) angles or detector emission (e) angles, the
156 incoming path length is equal to the height (s_0 or s_1) divided by $\cos(i)$ and the outgoing path
157 length is the height divided by $\cos(e)$. Implicit in equation (2) is that the surface contribution to
158 $I_\lambda(s_{0 \rightarrow 1})$ and $I_\lambda(s_{0 \rightarrow 2})$ is identical and that k is independent of path length (s). The effects of
159 aerosol scattering within wavelengths of gas absorption are ignored for the purposes of deriving
160 empirical transmission spectra and τ is assumed to result from molecular absorption only.
161 Although T at wavelengths outside of gas bands should equal unity, in practice, the ratio of the
162 low and high altitude spectra must be multiplicatively scaled to ~ 1.0 at wavelengths outside of
163 gas absorptions to account for differences in surface spectral features and/or aerosol opacities in
164 high and low altitude spectra.

165 To perform the volcano-scan correction for an arbitrary CRISM or OMEGA I/F
166 spectrum, the empirically derived transmission spectrum is scaled using an exponential factor
167 determined so that the most prominent CO_2 absorption feature in the transmission spectrum has
168 an equivalent depth to the same CO_2 feature in the I/F spectrum that is to be corrected, as
169 described below. For simplicity, the 'I/F spectrum that is to be corrected' will be referred to as
170 the target I/F spectrum. That is, division of the surface spectrum by the scaled transmission
171 spectrum, in principle, eliminates CO_2 absorption features from the target I/F spectrum. During
172 application of the volcano-scan correction, the transmission spectrum must be scaled separately
173 for each target I/F spectrum because the surface pressure, which affects the depth of CO_2
174 absorptions, varies as a function of surface elevation and solar longitude (L_s) (e. g., *Smith et al.*,
175 2004). Calculating the exponential scaling factor that is needed for the most prominent CO_2
176 feature in the transmission spectrum to match up with the depth of the same CO_2 feature in the

177 target I/F spectrum relies on band depth estimates. The band depth of an absorption feature in
 178 the target I/F spectrum and the band depth of the same feature in the empirical transmission
 179 spectrum are proxies for the transmittance that occurred over the atmospheric path lengths
 180 traversed for the I/F target spectrum and the empirical transmission spectrum, respectively. In
 181 the equations below, the arbitrary CRISM or OMEGA I/F spectrum that is to be corrected is
 182 labeled ‘target’, the transmission spectrum is labeled ‘trans’, I_{feature} is the I/F value of a CO₂
 183 absorption feature at some wavelength, T_{feature} is the transmission value of a CO₂ absorption
 184 feature at the same wavelength, and BD is the band depth of that CO₂ feature. The target I/F
 185 spectrum has an arbitrary atmospheric path length of $s_{0 \rightarrow 3}$.

$$186 \quad BD_{\text{targ}} = I_{\text{feature}}(\text{target})/I_{\text{continuum}}(\text{target}), \quad (4)$$

187 assuming $T=1$ on the continuum,

$$188 \quad BD_{\text{targ}} = I_{\text{feature}}(s_{0 \rightarrow 3})/I_{\text{continuum}} = e^{-ks_{0 \rightarrow 3}}, \quad (5)$$

189 Similarly,

$$190 \quad BD_{\text{trans}} = T_{\text{feature}}(\text{trans})/T_{\text{continuum}}(\text{trans}), \quad (6)$$

$$191 \quad BD_{\text{trans}} = T_{\text{feature}}(s_{1 \rightarrow 2})/T_{\text{continuum}} = e^{-ks_{1 \rightarrow 2}}. \quad (7)$$

192 The exponential scaling factor is determined from

$$193 \quad \ln(BD_{\text{targ}})/\ln(BD_{\text{trans}}) = \ln(e^{-ks_{0 \rightarrow 3}})/\ln(e^{-ks_{1 \rightarrow 2}}) = s_{0 \rightarrow 3}/s_{1 \rightarrow 2}. \quad (8)$$

194 Because the exponential scaling factor represents the difference in path length between
 195 the target spectrum and the transmission spectrum, its value should be the same for all

196 wavelengths. However, uncertainties are associated with determining the band depth from
 197 CRISM spectra. An estimate of the band depth for the strongest CO₂ band in the CO₂ triplet near
 198 2000 nm is commonly calculated as the ratio of the slope corrected I/F value along the
 199 continuum at ~1900 nm to the I/F value in the CO₂ absorption feature at ~2010 nm. OMEGA
 200 measures only a single channel near 2010 nm, which is located at 2012 nm, whereas CRISM has
 201 channels at 2007 and 2014 nm. *McGuire et al.*, [2009] note that variations in the surface spectral
 202 contribution between the wavelength along the continuum and the wavelength of the CO₂
 203 absorption feature can cause error in the calculated band depth and instead suggests using the
 204 ratio of the I/F continuum value at 1980 nm to the I/F value of the CO₂ absorption feature at
 205 2007 nm. These two band depth estimates will be referred to as the BD1 and BD2, respectively.
 206 All wavelengths of the transmission spectrum are scaled by the single (wavelength independent)
 207 exponential factor calculated using the band depth estimate.

208 Applying the volcano-scan correction to the target spectrum of interest, remembering that

$$209 \quad I_{\lambda}(s_{0 \rightarrow 3}) = I_{\lambda}(s_0) * e^{-ks_{0 \rightarrow 3}}, \quad (9)$$

210 results in

$$211 \quad I_{feature}(s_0) e^{-ks_{0 \rightarrow 3}} / (e^{-ks_{1 \rightarrow 2}})^{s_{0 \rightarrow 3}/s_{1 \rightarrow 2}} = I_{feature}(s_0), \quad (10)$$

212 where $I_{feature}(s_0)$ is the CRISM or OMEGA I/F value corrected for gas absorption, provided that
 213 all assumptions inherent in the volcano-scan correction are valid.

214 2.3 Radiative Transfer Modeling

215 Our radiative transfer modeling utilized a publically available general purpose Fortran
216 program for discrete-ordinate-method radiative transfer in scattering and emitting layered media,
217 Discrete Ordinate Radiative Transfer (DISORT) [*Stamnes et al.*, 1988; *Thomas and Stamnes*,
218 2002]. The atmosphere is treated as a plane-parallel medium in which individual layers are
219 homogenous but interlayer properties can be varied. Attenuation and scattering of the solar
220 beam down through the atmosphere, interaction with the surface, and attenuation and scattering
221 up through the atmosphere are modeled. The numerical implementation is discussed in *Stamnes*
222 *et al.*, [1999] and *Thomas and Stamnes* [2002]. Gas absorptions from CO₂, CO, and H₂O were
223 defined using correlated-k distributions with parameters derived from the HITRAN database
224 [*Rothman et al.*, 2005].

225 We utilized ‘front-end’ routines optimized for study of the Martian atmosphere,
226 DISORT_multi, to compute modeled CRISM and OMEGA I/F spectra [*Wolff et al.*, 2009].
227 Atmospheric temperature profiles relevant to atmospheric conditions at the latitude, longitude,
228 and Ls of interest were derived from spatially binned Mars Global Surveyor Thermal Emission
229 Spectrometer (TES) observations of Mars acquired between 1998 and 2000 [*Smith*, 2004]. Dust
230 aerosol single scattering albedos for 1.5 μm particles and wavelength dependent phase functions
231 utilized in our modeling were derived from analysis of CRISM hyperspectral data [*Wolff et al.*,
232 2009]. DISORT modeled I/F spectra presented in this paper are similar to those that would be
233 observed at the top of the Martian atmosphere by CRISM or OMEGA. Similar procedures were
234 described by *Arvidson et al.* [2006]. Absorption, emission, and multiple scattering from
235 atmospheric gases and aerosols were modeled.

236 3. Volcano-Scan Correction Simulation Results

237 To evaluate the validity of assumptions implicit in deriving empirical transmission
238 spectra and to identify and determine causes of artifacts introduced by application of the
239 volcano-scan correction, we simulated the volcano-scan correction using DISORT radiative
240 transfer modeling. Comparing volcano-scan corrected spectra generated using simulated
241 transmission spectra applied to modeled surface spectra enables identification and
242 characterization of artifacts introduced by the volcano-scan correction without the ambiguity
243 inherent in real data because all variables can be controlled and systematically varied. In this
244 section, we assessed potential artifacts that could arise from variation in surface spectral features
245 and atmospheric dust aerosols present in high and low altitude spectra used to generate empirical
246 transmission spectra as well as from ranges of exponential scaling factors potentially used during
247 application of the volcano-scan correction.

248 **3.1 Simulation of CRISM and OMEGA Empirical Transmission Spectra**

249 In order to demonstrate that simulated transmission spectra are similar to actual CRISM-
250 and OMEGA-derived transmission spectra, our initial simulation results were aimed at
251 approximating measured high and low altitude Olympus Mons spectra. Although we modeled
252 empirical transmission spectra that appear similar to both CRISM and OMEGA empirical
253 transmission spectra (Figure 3), we focus on CRISM results. The higher spectral resolution of
254 CRISM data (Figure 1) allows for more detailed characterization of artifacts. In addition, the
255 volcano-scan correction as applied to CRISM spectra has the added complication of a
256 temperature dependent wavelength shift, which we explore.

257 We chose to model CRISM high and low altitude spectra of Olympus Mons from
258 FFC000061C4. Figure 2a shows that the high and low altitude CRISM spectra do not overlap

259 and that the summit spectrum is slightly darker than the base spectrum. The dust aerosol
260 contribution is expected to be much less at the summit than at the base of Olympus Mons
261 because, in a uniformly mixed atmosphere, the amount of aerosol particles is proportional to the
262 atmospheric column density, which is much less at 20 km than 0 km altitude. Given the
263 expected small dust aerosol contribution at the summit of Olympus Mons, we attempted to
264 reproduce the low and high altitude spectra by modeling equivalent surface spectral properties in
265 the high and low altitude spectra, modest dust opacity for the low altitude spectrum, and zero
266 dust opacity for the summit spectrum. The surface spectrum was modeled as a straight line with
267 a positive slope. It was necessary to lower the surface albedo of the high altitude spectrum by
268 5% relative to the low altitude spectrum (shown in Figure 4b) to produce modeled base and
269 summit I/F spectra that are similar to measured base and summit I/F spectra (Figure 3). The
270 simulated transmission spectrum has values near 1.1 on the continuum (Figure 3b, gray line)
271 rather than expected values of 1.0, similar to the CRISM-derived transmission spectrum (Figure
272 2b, gray). Therefore, it was multiplicatively scaled to equal unity at wavelengths lacking gas
273 absorption (Figure 3b, black line) as was done for the actual CRISM-derived transmission
274 spectrum (Figure 2b, black line).

275 We produced simulated transmission spectra that appear broadly similar to actual
276 CRISM- and OMEGA-derived transmission spectra (Figure 3c, d). Volcano-scan corrections
277 using the simulated transmission spectra shown in Figure 3 were applied to modeled low altitude
278 Olympus Mons spectra in order to demonstrate that our simulated volcano-scan correction results
279 are similar to CRISM and OMEGA volcano-scan corrected spectra (Figure 4). Gaussian noise
280 with a standard deviation of 1% was added to the modeled I/F spectrum prior to application of
281 the volcano-scan correction to make the results more realistic. The resulting simulated volcano-

282 scan corrected spectra shown in Figure 4 are similar to actual CRISM and OMEGA volcano-scan
283 corrected spectra shown in Figure 1. The bowl-shape evident at 2000 nm (Figure 4) is similar to
284 the bowl-shaped artifact observed in CRISM and OMEGA volcano-scan corrected spectra
285 (Figure 1) even though the modeled input spectrum was linear (Figure 4b, c gray line). This
286 indicates that our simulation results are sufficient to characterize artifacts introduced by the
287 volcano-scan correction.

288

289 **3. 2 Impact of Aerosols and Surface Contributions in Transmission Spectra on the 2000 nm** 290 **Bowl-Shaped Artifact**

291 Two assumptions in empirically derived transmission spectra are that surface
292 contributions to low and high altitude spectra used to create transmission spectra are equivalent
293 and aerosol contributions to low and high altitude spectra can be ignored. If these assumptions
294 are valid, then spectral contributions should be equivalent outside wavelengths of gas absorption
295 in the low and high altitude spectra, and the base and summit spectra should overlap at these
296 wavelengths, resulting in empirical transmission spectra having values of unity outside of gas
297 absorptions. This is not the case (Figure 2a) and CRISM empirical transmission spectra must be
298 multiplicatively scaled so that continuum values are near unity (Figure 2b). The transmission
299 spectrum simulation results shown in Figure 4, in which the 2000 nm bowl-shaped artifact is
300 evident, were generated to appear similar to CRISM and OMEGA data and had variable surface
301 albedos (but no absorption feature at 2000 nm) and aerosol opacities between base and summit
302 spectra.

303 In order to determine whether or not surface spectral features and atmospheric dust
304 aerosols present in the high and low altitude spectra used to generate empirical transmission

305 spectra are responsible for the 2000 nm bowl-shaped artifact in the corrected spectra, we show
306 simulation results in which these two variables are equivalent in the low and high altitude
307 spectra. Figure 5a shows simulation results in which the base and summit spectra were identical
308 (assumption 1 satisfied) and aerosol opacities were identical. The 2000 nm bowl-shaped artifact
309 is evident. Figure 5b shows a similar simulation in which no aerosols were modeled (both
310 assumption 1 and 2 satisfied). This simulation demonstrates that a 2000 nm bowl-shaped
311 artifact occurs independently of the presence of aerosols and surface variability in the high and
312 low altitude spectra used to generate empirical transmission spectra. Although a surface
313 absorption feature near 2000 nm present in the base or summit spectrum could contribute to a
314 2000 nm artifact, such a feature is not required to explain the observed 2000 nm bowl-shaped
315 artifact in CRISM volcano-scan corrected spectra (Figure 1). However, it is important to note
316 that variable aerosol contributions between the empirical transmission spectrum and the
317 spectrum to be corrected causes a mismatch in shape between gas bands that results from aerosol
318 scattering into gas lines. This mismatch in shape between gas bands in the transmission
319 spectrum and the spectrum to be corrected causes additional artifacts near 2000 nm (section 4).

320

321 **3. 3 Impact of Exponential Scaling Factor on the 2000 nm Artifact**

322 The gas absorption coefficient (k) is assumed to be constant (equation 2) in the
323 calculation of the empirical transmission spectra. The volcano-scan correction is enabled by this
324 assumption (equations 4-8). During the application of the volcano-scan correction, the empirical
325 transmission spectrum is scaled exponentially by a single (wavelength independent) exponential
326 scaling factor that is calculated so that the strength of its strong CO₂ absorption near 2010 nm is
327 similar in strength to the same CO₂ feature in the spectrum to be corrected so that CO₂

328 absorptions can be divided out (see section 2.2). Small changes to the value of the exponential
329 scaling factor alter the expressions of the artifacts observed near 2000 nm in volcano-scan
330 corrected spectra (Figure 6). In the absence of artifacts, too small of an exponential scaling
331 factor would result in a residual CO₂ triplet near 2000 nm and too large of a scaling factor would
332 result in an inverse triplet. Applying the scaling factor calculated using the BD1 estimate results
333 in ‘hash’ superimposed on the bowl-shaped artifact near 2000 nm caused by over correction of
334 some parts of the CO₂ triplet, even though under correction of weaker CO₂ bands is evident at
335 other wavelengths (Figure 6). Applying the scaling factor calculated using the BD2 estimate
336 results in under correction of the CO₂ triplet near 2000 nm that has the appearance of a deeper
337 bowl-shape (Figure 6). This is because the BD2 estimate results in a consistently smaller
338 calculated exponential scaling factor. The addition of noise to the simulation (Figure 6b) reduces
339 the structured appearance of the residual gas bands but does not mask them. Similar artifacts are
340 evident in volcano-scan corrected CRISM spectra that were corrected using variable scaling
341 factors (Figure 6c), and there does not appear to be a single exponential scaling factor that
342 mitigates both the ‘hash’ and bowl-shape near 2000 nm (Figures 6).

343 The variable appearance of the artifact at 2000 nm in simulated and actual volcano-scan
344 corrected spectra caused by using different exponential scaling factors to scale the transmission
345 spectrum prior to its division demonstrates that assumption 3 is invalid and causes a prominent
346 artifact at 2000 nm and smaller artifacts at other wavelengths of CO₂ absorption. Gas bands that
347 appear broad at CRISM and OMEGA spectral resolution contain many individual absorption
348 lines which have variable strengths (Figure 7a). Larger absorption coefficients and/or more
349 molecules (higher column density) in the path length result in stronger and broader lines. A
350 detailed discussion of gas line broadening and the curve of growth in the context of radiative

351 transfer modeling can be found in *Thomas and Stamnes*, [2002]. In the weak line limit ($\tau \ll 1$),
352 or linear regime, wings of individual lines do not completely overlap and the absorptance ($1 - T$)
353 is proportional to the amount of absorbing molecules. In the strong line limit ($\tau \gg 1$), or
354 saturated regime, there is a square root dependence on column density. In between these two
355 limits, the wings of the gas lines overlap but the gas lines do not saturate, the absorption
356 coefficient (k) is constant with column density, and an exponential relationship between
357 absorptance and column density is observed. This regime is described by the integrated form of
358 the Beer-Lambert Law in which k is assumed to be constant for each wavelength (equation 2).
359 According to equation 3, $-\ln(T)$ should be linearly proportional to the path length, s , with the
360 slope equal to k . A plot of T calculated for different path lengths through the atmosphere using
361 DISORT reveals that T is not necessarily linearly proportional to s for all s for all wavelengths
362 (Figure 7b). This result is expected because multiple regimes are experienced for the range of
363 pressures in the Martian atmosphere and occur over different path lengths depending on the
364 strength of CO_2 absorption at a given wavelength.

365

366 **4. Artifacts and Residual Atmospheric Spectral Features in CRISM Volcano-Scan**

367 **Corrected Spectra**

368 The volcano-scan correction does not specifically address potentially large spectral
369 contributions from spatially and temporally variable dust aerosols, ice aerosols, and water vapor
370 that are present in the Martian atmosphere. In addition, the appearance of the 2000 nm bowl-
371 shaped artifact has variable expression depending on options selected in CAT during application
372 of the volcano-scan correction and the temperature dependent wavelength shift of the
373 transmission spectrum relative to the corrected spectrum. We assessed sensitivity to atmospheric

374 dust aerosols, ice aerosols, and water vapor and also present examples showing how the
375 temperature dependent wavelength shift, band depth estimate (BD1 or BD2), and empirical
376 artifact correction options present in CAT further impact the 2000 nm bowl-shaped artifact in
377 CRISM spectra.

378

379 **4.1 Residual Atmospheric Dust Aerosol Features**

380 Dust and ice aerosol abundances are spatially and temporally variable on Mars. Higher
381 dust opacities are typically observed in the Martian perihelion season and there is regional
382 variation in dust loading in the atmosphere [e.g., *Smith et al.*, 2004]. In the NIR, dust aerosols
383 generally impart a negative slope, affect overall spectral brightness, and impact the depth and
384 shape of the CO₂ triplet (Figure 8). Because multiple scattering from dust aerosols brightens
385 dark spectra and darkens bright spectra (Figure 8), with the magnitude of the effect depending on
386 dust opacity, the overall brightness of volcano-scan corrected spectra is not equivalent to surface
387 albedo. This is especially evident in simulated volcano-scan corrected spectra in which the
388 modeled surface spectrum is compared with the volcano-scan corrected spectrum (Figure 5a).
389 The simulated volcano-scan corrected spectrum has a lower albedo and a more negative slope as
390 a result of uncorrected dust aerosol contributions.

391 Variable aerosol contributions between the empirically-derived transmission spectrum
392 and the spectrum to be corrected cause a mismatch in shape between gas bands, which results
393 from aerosol scattering into gas lines. This mismatch in shape between gas bands in the
394 transmission spectrum and the spectrum to be corrected causes additional artifacts near 2000 nm.
395 The effect of dust aerosols on the CO₂ bands is most pronounced in higher opacity images
396 because empirical transmission spectra are derived from images with low dust opacities.

397 Volcano-scan corrected high opacity spectra exhibit additional ‘hash’ near 2000 nm because of
398 the larger mismatch in shape between gas bands in the transmission spectrum and the spectrum
399 to be corrected. This is evident in Figure 9 in which volcano-scan corrected CRISM spectra
400 from the same location on Mars but acquired at different times with variable atmospheric
401 conditions are compared. For this comparison, we chose images covering the dusty plains a few
402 km to the south of the Mars Exploration Rover Spirit landing site at Gusev crater. Surface
403 spectra of dusty surfaces are relatively featureless in the NIR, which allows atmospheric residual
404 features to stand out in volcano-scan corrected spectra. Multiple images were acquired of this
405 area at different times of year that capture a variety of atmospheric dust aerosol optical depths.
406 In addition, aerosol optical depth measurements derived from Spirit Pancam observations have
407 been reported [Lemmon *et al.*, 2004]. The 2000 nm bowl-shaped artifact has a variable
408 appearance in example spectra from all three images (Figure 9). Variations are evident that
409 result from different atmospheric conditions and different volcano-scan correction options. All
410 spectra shown in Figure 9 were corrected using transmission spectra with the closest match to the
411 wavelength shift. Both the BD1 and BD2 estimates were used to scale the transmission spectra
412 and the CAT empirical 2000 nm bowl-shaped artifact correction was applied (Figure 9).
413 Volcano-scan corrected spectra from the dustiest image (FRT000553B) exhibit more ‘hash’ in
414 the 2000 nm artifact (regardless of CAT volcano scan options) that results from changes in gas
415 band shape caused by aerosol scattering.

416

417 **4.2 Residual Atmospheric Ice Aerosol Features**

418 Water ice aerosols also effect brightness and spectral slope but are particularly
419 problematic because they introduce absorption features at 1500, 2000, and 2400 nm (Figure 10).

420 The depth of these features depends on the ice aerosol optical depth and grain size (e. g., *Clancy*
421 *et al.*, 2003). Water ice aerosol opacities are highly variable and these aerosols are commonly
422 present at high latitudes during the Martian winter but occur near the equator in the aphelion
423 season [e.g., *Clancy et al.*, 2003].

424 In order to illustrate the appearance of uncorrected water ice aerosols, we compared
425 volcano-scan corrected spectra from two different locations within the same image.
426 FFC0000A3F6 was acquired over a relatively featureless, dusty area of Mars. The differences
427 between the two spectra shown in Figure 10 result from under corrected water ice aerosols.
428 Absorption features at 1500 and 2000 nm as well as the pronounced negative slope at
429 wavelengths > 2400 nm are evident in the lower spectrum. These water ice aerosol absorptions
430 overlap with some features in hydrated minerals, particularly hydrated sulfates [e.g., *Cloutis et*
431 *al.*, 2006], and can interfere with mineral identifications when present.

432

433 **4.3 Residual Atmospheric Water Vapor Features**

434 The volcano-scan correction scales all gases (e.g., CO₂, H₂O and CO) present in the
435 transmission spectrum by the exponential scaling factor calculated using the ~ 2100 nm CO₂ band
436 depth estimate. Water vapor concentrations in the Martian atmosphere vary by L_s , latitude, and
437 elevation with typical values between ~ 5 and 20 precipitable μm [e.g., *Smith et al.*, 2008].

438 Under correction of water vapor leaves atmospheric absorption features at ~ 1400 , 1900, and
439 2600 nm that can impact diagnostic spectral features exhibited by alteration phases that occur
440 near these wavelengths. Figure 11 shows a relatively featureless dusty surface with water
441 absorption features present. Water vapor spectral contributions to surface spectra with alteration
442 phases are less obviously attributable to water vapor because they overlap with mineral

443 absorption features. This is well illustrated by Figure 12 which shows volcano-scan corrected
444 CRISM spectra from the same carbonate-bearing [Ehlmann *et al.*, 2009] location on Mars but
445 acquired at different times under variable atmospheric conditions. The 2600 nm water vapor
446 feature overlaps with the ~2500 nm carbonate feature and contributes to mismatches between the
447 two different spectra of the same location. In addition, this example shows that the 2000 nm
448 bowl-shaped artifact has a variable appearance and complicates interpretation of the ~1910 nm
449 surface hydration feature. Spectra were corrected using transmission spectra with the closest
450 match to the wavelength shift. Both the BD1 and BD2 estimates were used to scale the
451 transmission spectrum and the CAT empirical 2000 nm bowl-shaped artifact correction was
452 applied (Figure 12a). A ratio spectrum is shown to highlight the fact that apparent differences
453 between spectra taken from the two different images are related to artifacts introduced by the
454 volcano-scan correction as well as residual atmospheric water vapor (Figure 12b). Both
455 FRT00003FB9 and FRT0000A09C have aerosol opacities of ~0.4 calculated using methods
456 described by Wolff *et al.* [2009].

457

458 **5 Summary and Implications**

459 The empirical ‘volcano-scan’ correction [Bibring *et al.*, 1989; Langevin *et al.*, 2005;
460 McGuire *et al.*, 2009] is widely applied to CRISM and OMEGA NIR spectra between ~1000 to
461 ~2600 nm to remove prominent atmospheric CO₂ absorptions with minimal computational
462 investment. However, detailed examination of volcano-scan corrected spectra reveals a bowl-
463 shaped artifact that overlaps with prominent CO₂ features near 2000 nm. The identification and
464 characterization of both artifacts and residual atmospheric features left by the volcano-scan

465 correction is important for the proper interpretation of CRISM and OMEGA volcano-scan
466 corrected spectra.

467 Simulation of the volcano-scan correction with radiative transfer modeling enabled
468 assessment of assumptions underlying this empirical method of correction. We found that the
469 most prominent artifact, a bowl-shape centered near 2000 nm, is caused by the inaccurate
470 assumption that absorption coefficients of CO₂ in the Martian atmosphere are independent of
471 column density (Figure 7). This means that transmission is not accurately derived by the
472 division of high and low altitude Martian spectra.

473 Expression of the 2000 nm artifact varies depending on the choice of the exponential
474 scaling factor used to scale the empirical transmission spectrum during application of the
475 volcano-scan correction (Figure 6). In addition, differences in aerosol scattering into gas bands
476 between the transmission spectrum and the spectrum to be corrected also causes a mismatch in
477 shape between gas bands that results in ‘hash’ in the 2000 nm region (Figure 9). CAT software
478 version 7.2.1 provides an option to empirically correct the bowl-shape at 2000 nm that largely
479 removes this feature (Figures 9 and 12). However, the performance of this correction on
480 different types of surface spectra acquired under variable atmospheric conditions has not been
481 quantified in the literature and variations between spectra of the same surface acquired under
482 different atmospheric conditions are evident (Figure 12).

483 Residual atmospheric contributions caused by dust aerosols, ice aerosols, and water vapor
484 are also observed in volcano-scan corrected spectra. Because concentrations of these
485 atmospheric species are temporally and spatially variable, features resulting from these
486 atmospheric contributions in volcano-scan corrected spectra are also variable. Uncorrected dust
487 aerosols modulate the spectral slope and albedo (Figure 8) causing volcano-scan corrected

488 spectra to have inaccurate low frequency shapes (Figure 5a). Uncorrected ice aerosols, when
489 present, result in absorptions centered at 1500 and 2000 nm and a negative slope at 2400 nm
490 (Figure 10). Under corrected water vapor in some volcano-scan corrected spectra results in
491 absorption features centered at 1400, 1870, and 2650 nm (Figure 11). These features are
492 particularly problematic in spectra of alteration minerals (Figure 12).

493 The presence of artifacts and residual atmospheric features left by the volcano-scan
494 correction can impact interpretation of volcano-scan corrected CRISM and OMEGA spectra.
495 Because spurious features introduced by the volcano-scan correction overlap with and obscure
496 diagnostic spectral absorptions that occur between ~1900 and ~2150 nm (e.g., Figure 12),
497 spectral ratios are commonly performed to remove these residual features (as well as calibration
498 artifacts). Comparison of volcano-scan corrected spectra to ratioed volcano-scan corrected
499 spectra highlights artifacts that are evident in unratioed spectra (Figure 12). Spectral ratios of
500 volcano-scan corrected spectra are useful for confirming the presence of narrow absorption
501 features caused by alteration phases. However, although spectral ratios largely remove the
502 volcano-scan artifacts and residual atmospheric contributions, they impact low frequency
503 spectral shape, i.e. continuum shape, and can produce spurious features that mimic spectral
504 absorptions if an inappropriate denominator is utilized. Multiple scattering from atmospheric
505 aerosols is both additive and multiplicative; therefore, aerosol spectral contributions are not
506 completely removed by ratioing. In addition, the amplitude of ratio spectra is arbitrary and
507 cannot be related to physically meaningful units. Quantitative spectral analyses that address
508 band shape, band depth, and mineral abundances are strongly dependent on the quality of
509 atmospherically corrected spectra. Because of the presence of spurious features and/or
510 uncorrected atmospheric contributions in volcano-scan and ratio spectra, radiative transfer

- 511 modeling is necessary for quantitative analyses [e.g., *Arvidson et al.*, 2006; *McGuire et al.*, 2008;
- 512 *Poulet et al.*, 2009a; *Poulet et al.*, 2009b; *Cull et al.*, 2010a; *Cull et al.*, 2010b; *Liu et al.*, 2012].

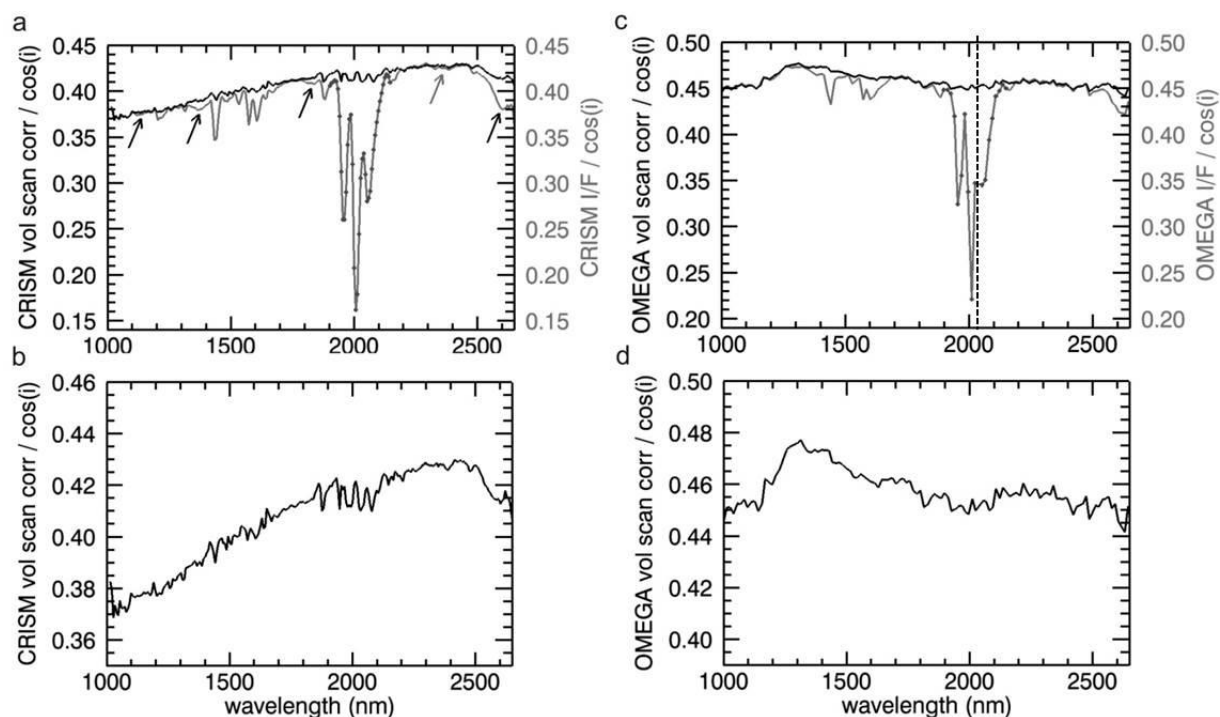
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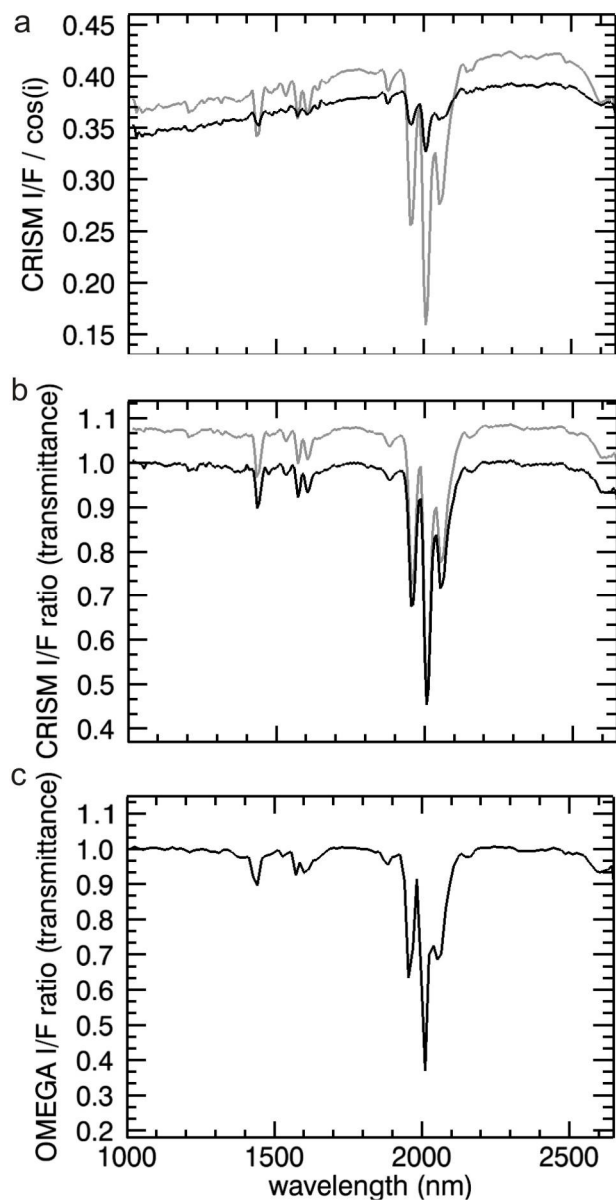
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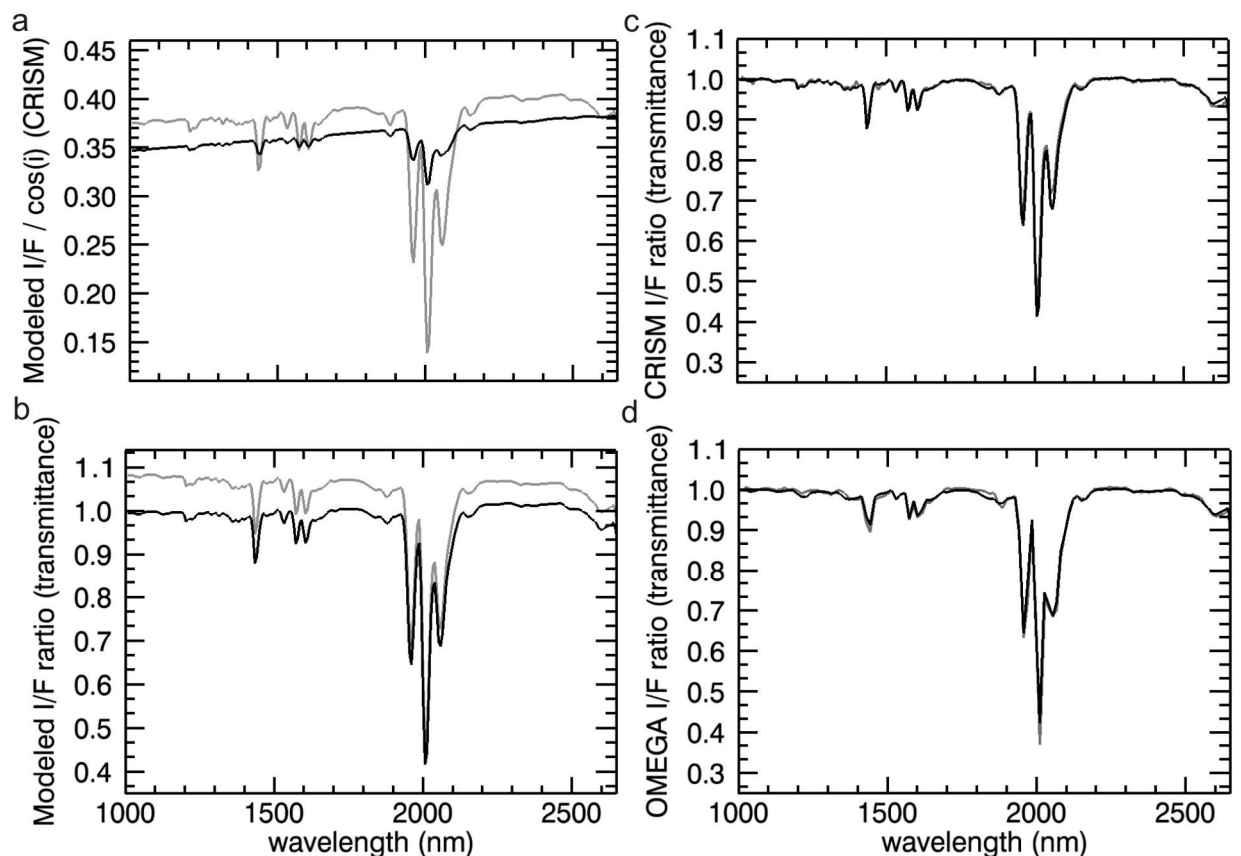
FIGURES



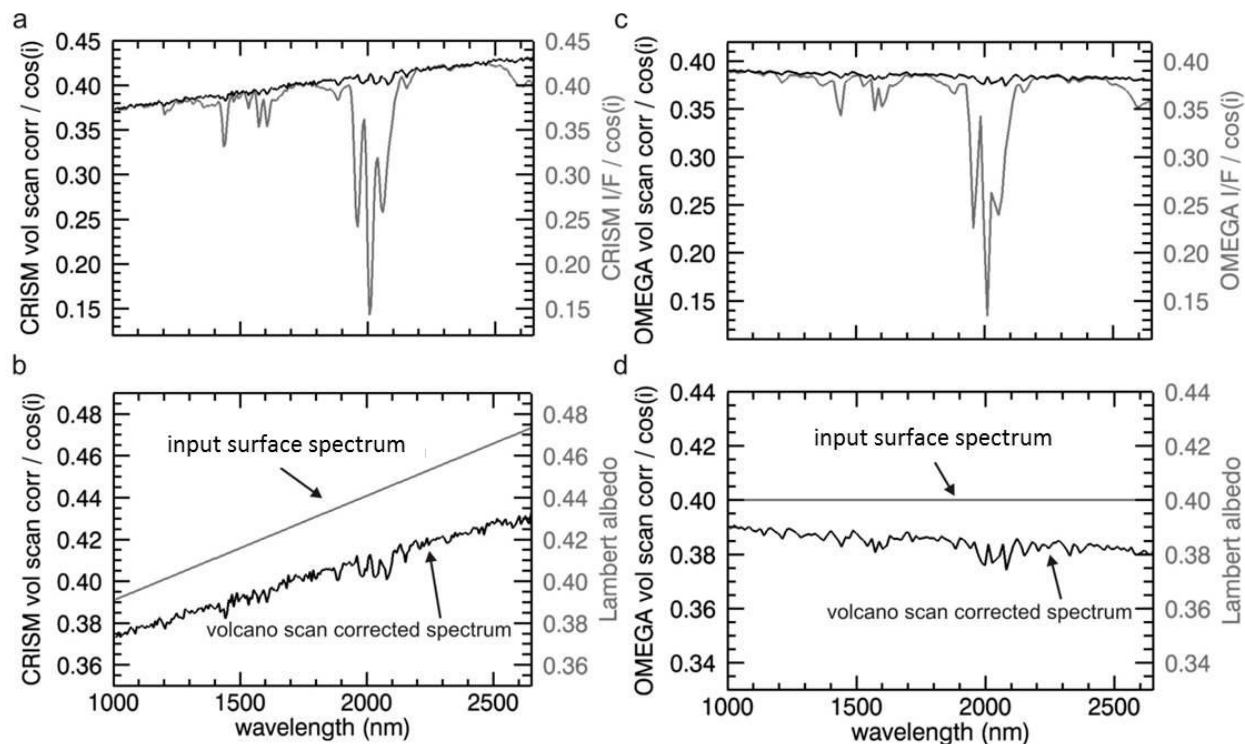
589 Figure 1. a) CRISM spectrum extracted from FFC000061C4 before (gray) and after (black)
 590 volcano-scan correction. Because the transmission spectrum and corrected spectrum were both
 591 derived from the same image, the temperature dependent wavelength shift is minimized. Gray
 592 dots show CRISM channel positions located between 1900 and 2150 nm (7 nm intervals). Black
 593 arrows indicate atmospheric H₂O vapor absorptions centered near 1130, 1380, 1880, and 2590
 594 nm. A shallow CO feature near 2350 nm is also present. Unlabeled absorptions result from
 595 atmospheric CO₂. b) Close up of the black CRISM spectrum shown in part (a). Note the bowl-
 596 shape and 'hash' in the 2000 nm region. c) OMEGA spectrum extracted from ORB0037_2
 597 before (gray) and after (black) volcano-scan correction. Gray dots indicate OMEGA channel
 598 positions located between 1900 and 2150 nm. OMEGA acquires data at 14 nm intervals;
 599 however the channel at 2040 nm (dotted line) is dead. Lack of data at this wavelength causes the
 600 longest wavelength minimum in the CO₂ triplet to appear less well defined in OMEGA spectra
 601 than CRISM spectra. d) Close up of the black OMEGA spectrum shown in part (c). Both
 602 CRISM and OMEGA spectra were extracted from dust covered surfaces located to the south of
 603 Olympus Mons. Differences between the CRISM and OMEGA spectra shown in parts (b) and
 604 (d) could result from calibration characteristics as well as differences in viewing geometries and
 605 atmospheric conditions.



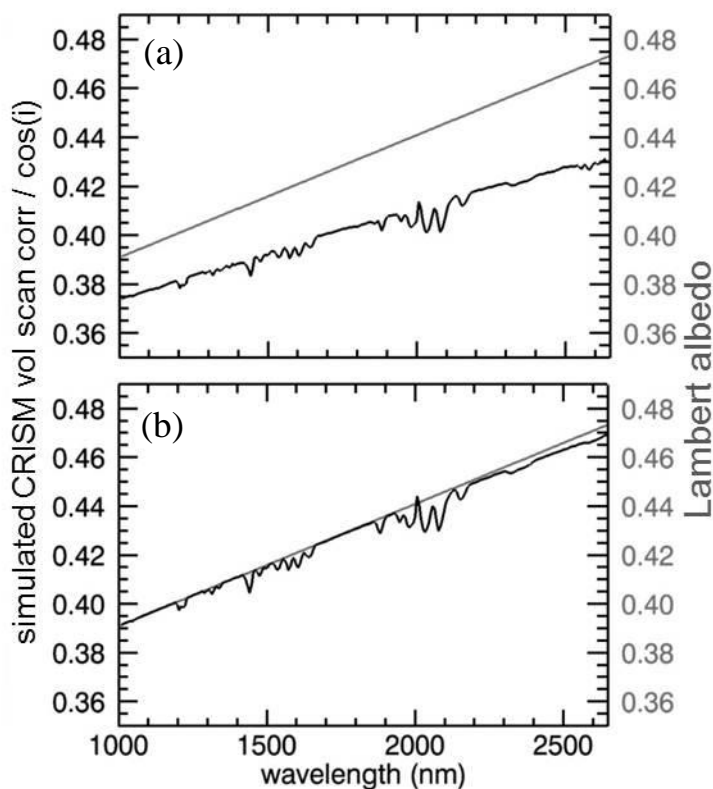
606
607 Figure 2. a) CRISM high (black) and low (gray) altitude spectra acquired at the base and summit
608 of Olympus Mons, respectively, extracted from FFC000061C4. b) CRISM scaled (black) and
609 unscaled (gray) transmission spectra generated from low and high altitude spectra shown in part
610 (a). c) OMEGA scaled transmission spectrum.



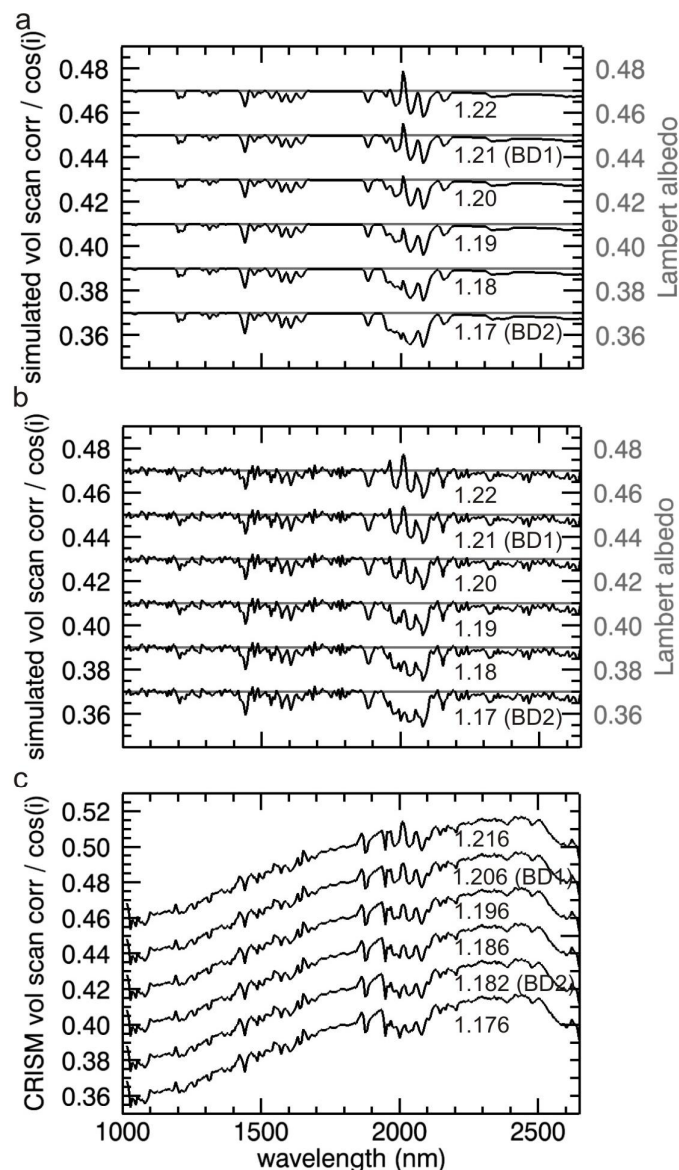
611 Figure 3. a) DISORT modeled high (black) and low (gray) altitude spectra convolved to CRISM
 612 spectral resolution computed using pressure/temperature profiles appropriate for the base and
 613 summit of Olympus Mons, respectively. A dust opacity of 0.3 and 5 precipitable μm of water
 614 vapor were modeled for the low altitude spectrum and no dust or water vapor were modeled for
 615 the high altitude spectrum. b) Scaled (black) and unscaled (gray) transmission spectra generated
 616 from spectra shown in part (a). c) Modeled transmission spectrum (black) overlain on an actual
 617 CRISM transmission spectrum (gray). d) Modeled transmission spectrum (black) overlain on an
 618 actual OMEGA transmission spectrum (gray).



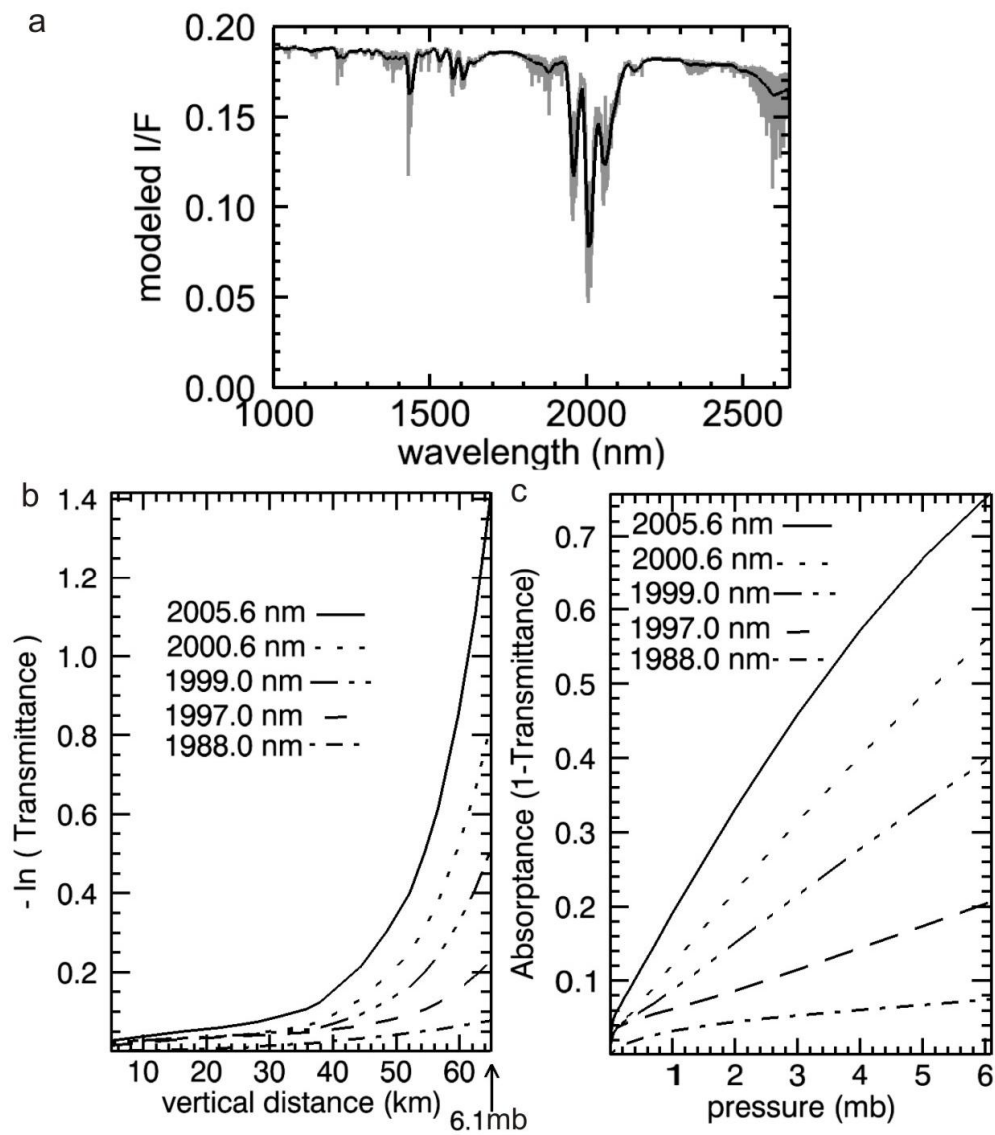
619 Figure 4. a) Modeled CRISM spectrum before (gray) and after (black) simulated volcano-scan
 620 correction. b) Close up of the black spectrum shown in part (a) but with the input surface
 621 spectrum over-plotted in gray. A linear function was used to model the surface spectrum to
 622 highlight artifacts in the volcano-scan correction. Perfect atmospheric correction would result in
 623 the black and gray spectra overlapping. c-d) Same as parts (a-b) but for OMEGA spectral
 624 resolution. Although CRISM and OMEGA spectra shown in parts (a) and (c) were both acquired
 625 at Olympus Mons, the CRISM spectrum has a positive slope whereas the OMEGA spectrum is
 626 near horizontal. This mismatch in continuum shape between the two datasets, as released in the
 627 PDS, is likely related to specifics of the instrument calibrations and is beyond the scope of this
 628 paper to address.
 629



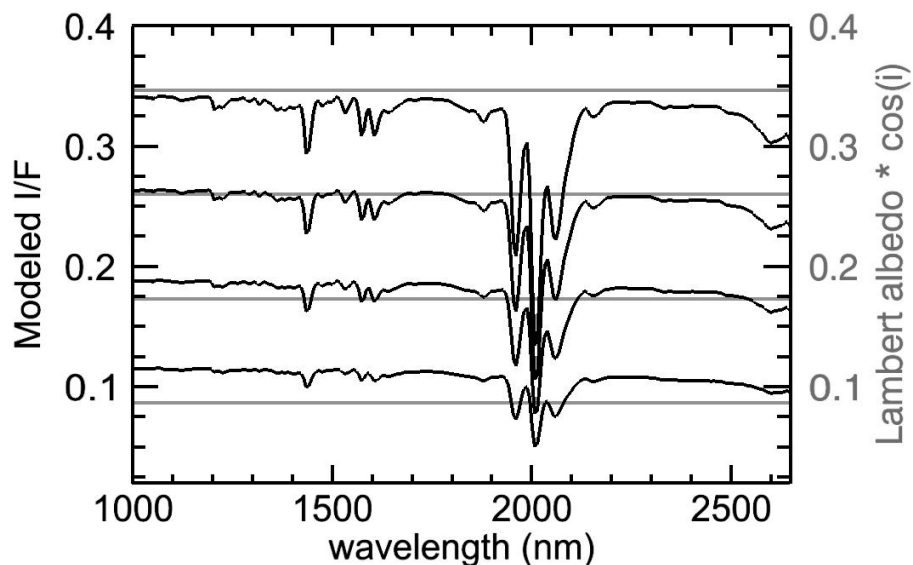
630 Figure 5. (a) Simulated CRISM volcano-scan corrected spectrum (black) and input surface
 631 Lambert albedo spectrum (gray). Aerosols were present in both the simulated transmission
 632 spectrum (Figure 3b, black) and the modeled spectrum that was corrected. Noise was not added
 633 to the simulation to highlight artifacts. Perfect atmospheric correction would result in the black
 634 spectrum overlapping with the gray spectrum. (b) Same as part (a), but with no aerosols
 635 modeled in either the simulated transmission spectrum or the corrected spectrum. Perfect
 636 atmospheric correction would result in the black spectrum overlapping with the gray spectrum.



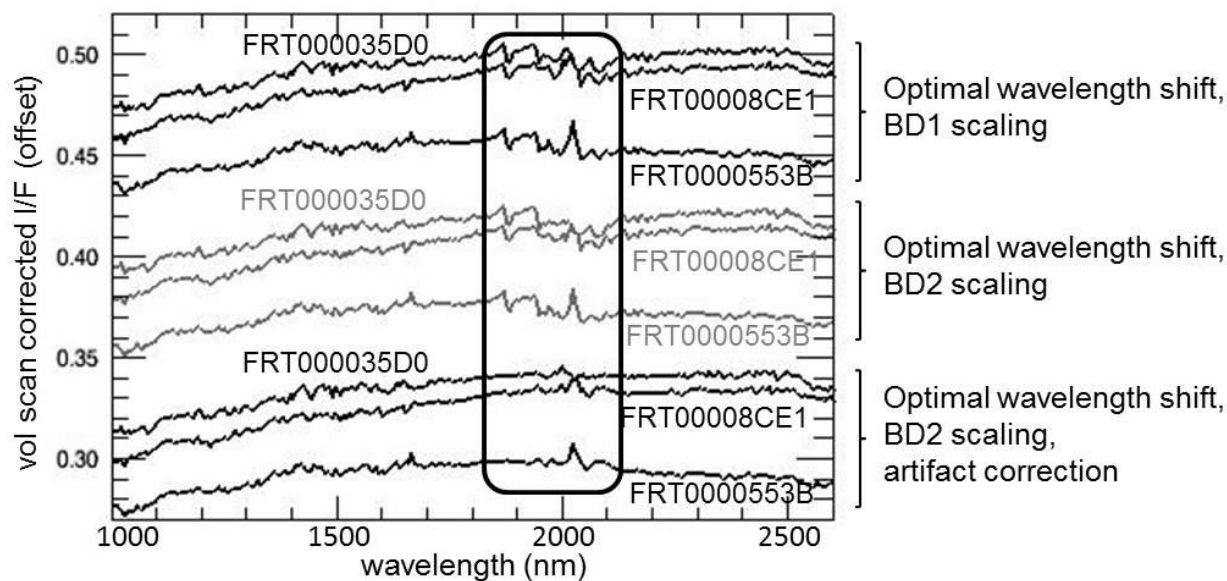
637 Figure 6. a) Black lines: Simulated volcano-scan corrected spectra using simulations with no
 638 aerosols modeled, offset for clarity. The exponential scaling factors (see Equation 8) used to
 639 create the simulated volcano-scan corrected spectra were systematically varied. The exponential
 640 scaling factor used to scale the simulated transmission spectrum prior to division (labeled on
 641 plot) was the only variable changed among the plotted spectra. Note the change in character of
 642 the artifacts near 2000 nm as the scaling factor value is increased. The gray lines show actual
 643 surface spectra (perfect atmospheric correction would result in the black spectra overlapping
 644 with the gray spectra). b) Same as part (a) but with noise added. c) Actual CRISM volcano-
 645 scan corrected spectra with systematically varied exponential scaling factors, offset for clarity.
 646 The exponential scaling factor used to scale the transmission spectrum prior to division (labeled
 647 on plot) was the only variable changed among the plotted spectra. Note the change in character
 648 of the artifacts near 2000 nm as the scaling factor value is increased.



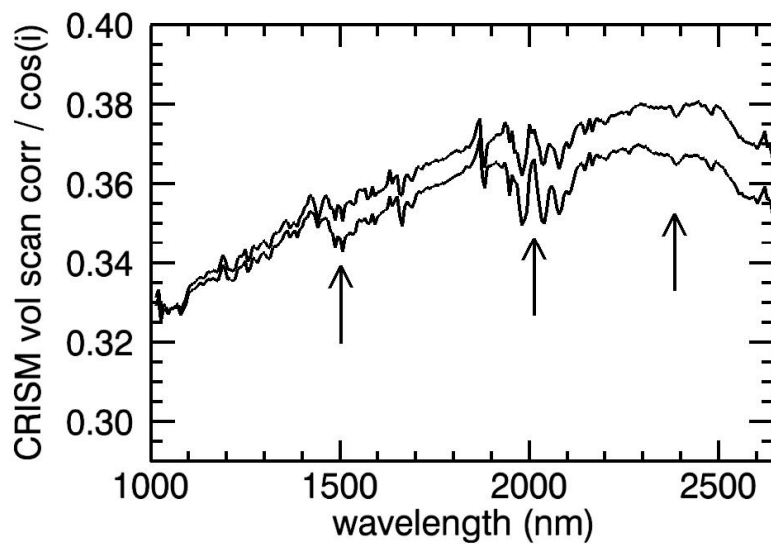
649 Figure 7. a) Shapes of gas absorption features measured by CRISM are a convolution of very
 650 narrow gas bands and the CRISM instrument response function. Gray line: DISORT generated
 651 I/F spectrum with I/F values calculated every 0.4 nm. Black line: High spectral resolution gray
 652 spectrum convolved to CRISM spectral resolution. b) Plots of negative natural logarithm
 653 transmittance calculated using DISORT versus vertical distance at several different wavelengths.
 654 Vertical distance times 2 is equivalent to the total path length, s . According to equation 3, there
 655 should be a linear relationship between the negative natural logarithm of transmittance and
 656 distance, s . If this were the case, then the curves would be linear with slopes of k . For this plot,
 657 the top of the atmosphere (vertical distance of 0) was set to 65 km above the surface, a distance
 658 at which atmospheric pressure is negligible. At the surface (vertical distance of 65 km), the
 659 pressure is 6.1 mb, which is Mar's average atmospheric pressure. c) Plots of absorbance versus
 660 pressure. Pressure is exponentially proportional to vertical distance (assuming constant
 661 temperature). Therefore, according to equation 3, these curves should also appear linear.



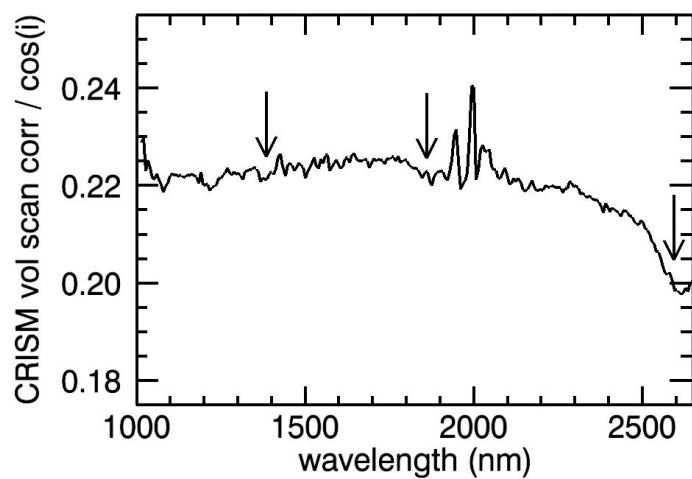
662 Figure 8. DISORT generated I/F spectra (black lines) for surface Lambert albedo spectra shown
 663 by gray lines. Long wavelength I/F spectral shape is affected by atmospheric dust, which
 664 imparts a negative slope. The surface Lambert albedo spectra with no atmospheric contributions
 665 (gray lines) have slopes of 0. Multiple scattering caused by dust aerosol particles brightens I/F
 666 spectra of dark surfaces and darkens spectra of bright surfaces. I/F spectra were modeled with
 667 DISORT for a modest dust opacity of 0.6 at 900 nm and a viewing geometry of $i = g = 30^\circ$, $e =$
 668 0° and surface Lambert albedos of 0.1, 0.2, 0.3, or 0.4 at all wavelengths.



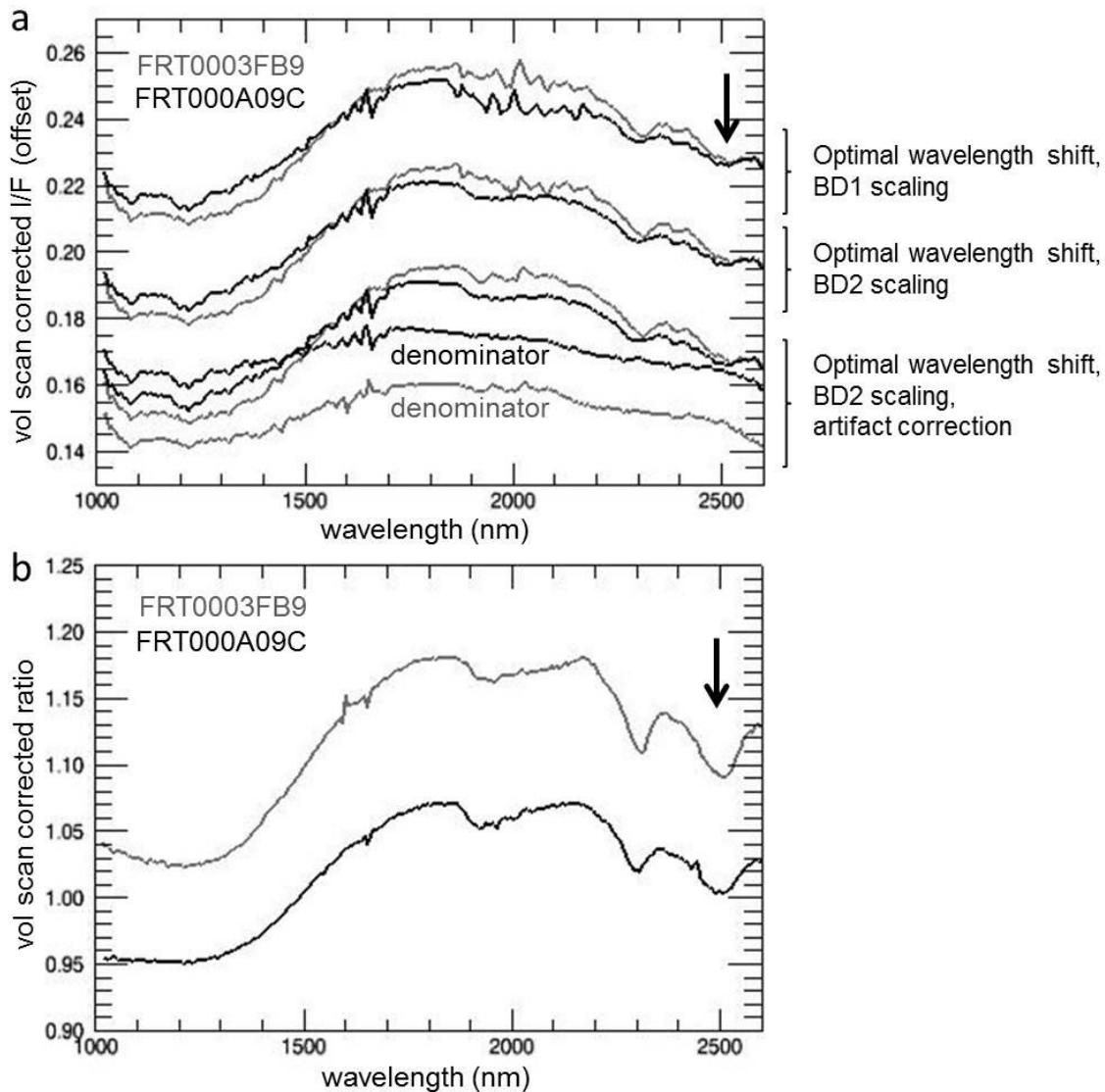
669 Figure 9. CRISM volcano-scan corrected spectra from the same location (-14.656N, 175.463E)
 670 in Gusev crater images FRT000035D0 ($\tau = 0.4$), FRT0000553B ($\tau = 0.6$), and FRT00008CE1 (τ
 671 = 0.9) (10x10 pixel average). FRT0000553B has the highest dust opacity and the most hash in
 672 the 2000 nm bowl-shaped artifact. The 2000 nm bowl-shaped artifact has a different appearance
 673 depending on the transmission spectrum and exponential scaling factor applied during the
 674 volcano-scan correction.



675 Figure 10. Volcano-scan corrected CRISM I/F spectra extracted from FFC0000A3F6. The
676 lower spectrum has distinct water ice aerosol features, which include absorptions centered at
677 1500 and 2000 nm and a decrease in slope starting at 2300 nm. The 2000 nm ice absorption is
678 partially obscured by artifacts induced by the volcano-scan correction.



679 Figure 11. Volcano-scan corrected CRISM I/F spectrum extracted from FRT0000CAB3 in which
680 water vapor is under-corrected. The water vapor absorption features near 1380, 1880, and 2590
681 nm are indicated with black arrows. Note, the CO₂ triplet near 2000 nm is poorly corrected as a
682 result of high atmospheric dust opacity (> 1.0) at the time of image acquisition.



683 Figure 12. (a) CRISM volcano-scan corrected spectra from the same location (21.24N, 78.6E)
 684 covering a carbonate- and phyllosilicate- bearing outcrop (10x10 pixel average) in Nili Fossae.
 685 Spectra from FRT00003FB9 were extracted from column 220, line 21 and spectra from
 686 FRT0000A09C from column 52, line 100. Although spectra were extracted from roughly the
 687 same location, they are not expected to have identical brightnesses because of differences in
 688 illumination resulting from different solar incidence angles and measurement emission angles.
 689 The 2000 nm bowl-shaped artifact, which overlaps with a ~1910 nm surface hydration feature,
 690 has a different appearance depending on the transmission spectrum and exponential scaling
 691 factor applied during volcano-scan correction. Variable amounts of residual water vapor alter
 692 the expression of the 2500 nm carbonate feature (black arrow). This feature looks more similar
 693 in ratio spectra shown in part (b). (b) Spectral ratios calculated using the lower set of spectra
 694 (optimal wavelength shift, BD2 scaling, artifact correction) shown in in part (a). Denominator
 695 spectra used to generate these ratio spectra were extracted from the same columns as the

696 corresponding numerator spectra rather than the same surface location to minimize column
697 dependent artifacts. Note, it was not possible to extract spectra from both the same surface
698 location and corresponding columns as numerator spectra due to CRISM measurement geometry.
699 Ratio spectra remove volcano-scan correction artifacts and highlight narrow absorption features.
700 However, low frequency spectral shape is variable between numerator spectra shown in part (a)
701 and ratioed spectra shown in part (b). In addition, the artifact correction as currently
702 implemented in CAT v7.2.1 impacts wavelengths between ~1760 and 2240 nm and contributed
703 to the apparent peak at ~ 2200 nm in the gray ratioed spectrum.

The empirical volcano scan correction leaves a prominent artifact near 2000nm.

Spectral albedo and slope are modified by uncorrected atmospheric aerosols.

Residual water vapor features overlap with mineral hydration features.

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