

# Aeroelasticity of Axially Loaded Aerodynamic Structures for Truss-Braced Wing Aircraft

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**This paper presents an aeroelastic finite-element formulation for axially loaded aerodynamic structures. The presence of axial loading causes the bending and torsional stiffnesses to change. For aircraft with axially loaded structures such as the truss-braced wing aircraft, the aeroelastic behaviors of such structures are nonlinear and depend on the aerodynamic loading exerted on these structures. Under axial strain, a tensile force is created which can influence the stiffness of the overall aircraft structure. This tension stiffening is a geometric nonlinear effect that needs to be captured in aeroelastic analyses to better understand the behaviors of these types of aircraft structures. A frequency analysis of a rotating blade structure is performed to demonstrate the analytical method. A flutter analysis of a truss-braced wing aircraft is performed to analyze the effect of geometric nonlinear effect of tension stiffening on the flutter speed. The results show that the geometric nonlinear tension stiffening effect can have a significant impact on the flutter speed prediction. In general, increased wing loading results in an increase in the flutter speed. The study illustrates the importance of accounting for the geometric nonlinear tension stiffening effect in analyzing the truss-braced wing aircraft.**

## I. Introduction

Future aircraft design trends are employing high-aspect ratio wing technology to improve aerodynamic efficiency for reduced fuel burn. As a wing aspect ratio increases, the need for maintaining sufficient structural load carrying capacity is becoming more important. Cantilever wing design can only accommodate up to a certain aspect ratio beyond which the wing root bending moment can become too large that will impose structural limitations on a wing design. At the same time, wing deflection can also become too excessive that could cause aerodynamic penalty. Truss-braced wing aircraft concepts provide a structural solution to high aspect ratio wing aircraft designs. The long slender wing would employ structural bracing via the use of axially loaded strut members to provide intermediate span supports in addition to the wing root attachment. These struts generally support a portion of the spanload carried by the wing and are loaded in tension. Under a negative-g flight condition such as during a dive, a load reversal could occur

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that could put the struts in compression. The compressive loading would require design considerations for buckling strength.

Under aerodynamic loading, an axially loaded member also experiences the normal bending and torsion generated by aerodynamic lift force and pitching moment. Aeroelasticity of an axially loaded structure undergoing transverse bending can be significantly different from that with transverse bending alone.

The SUGAR Truss-Braced Wing (TBW) aircraft concept is a Boeing-developed N+3 aircraft configuration funded by NASA ARMD Fixed Wing project.<sup>1,2</sup> The TBW aircraft concept is designed to be aerodynamically efficient by employing a high-aspect ratio wing design. The aspect ratio of the TBW is in the order of 14 which is significantly greater than those of conventional aircraft wings. As a result, intermediate structural supports are required. The main wings are braced at approximated at mid-span by two main struts. In addition, two jury struts; one on each wing, provide additional reinforcement. Figure 1 is an illustration of the TBW aircraft.



Fig. 1 - Boeing SUGAR Truss-Braced Wing (TBW) Aircraft Concept

The main struts and jury struts can impact aerodynamics of the main wings since they create aerodynamic interference drag that can offset the aerodynamic benefit of high-aspect ratio wing design. Thus, the TBW is designed to strive a balance between aerodynamic efficiency and structural efficiency.

In this paper, we will develop an aeroelastic analysis for axially loaded aerodynamic structures such as the TBW main struts.<sup>3</sup> The presence of axial loading causes the bending and torsional stiffnesses to change. A tensile loading will result in an increase in bending and torsional stiffnesses. This tension stiffening is a geometric nonlinear effect which should be captured in aeroelastic analyses of aircraft structures that carry significant axial loads. A finite-element formulation is developed to model the axially loaded structures in the presence of the bending-torsion motion of a typical aircraft wing. A flutter analysis is performed to analyze the geometric nonlinear effect of tension stiffening on the flutter speed.

## II. Elastic Analysis

In the subsequent analysis, the combined motion of the left wing is considered. The motion of the right wing is a mirror image of that of the left wing for symmetric flight. The wing has a varying pre-twist angle  $\gamma(x)$  commonly designed in many aircraft. Typically, the wing pre-twist angle varies from being nose-up at the wing root to nose-down at the wing tip. The nose-down pre-twist at the wing tip is designed to delay stall onsets. This is called a wash-out twist distribution. Under aerodynamic forces and moments, the aeroelastic deflections of a wing introduce stresses and strains into the wing structure. The internal structure of a wing typically comprises a complex arrangement of load carrying spars and wing boxes. Nonetheless, the elastic behavior of a wing can be captured by the use of equivalent stiffness properties. These properties can be derived from structural certification testing that yields information about wing deflections as a function of loading. For high aspect ratio wings, an equivalent beam approach can be used to analyze aeroelastic deflections with good accuracy.<sup>4</sup>

Consider an airfoil section on the left wing as shown in Fig. 2 undergoing bending and torsional deflections. Let  $(x, y, z)$  be the undeformed coordinates of point Q on a wing airfoil section in the reference frame D. Let  $\mathbf{p}_0 = x\mathbf{d}_1$  be a position vector along the elastic axis. Then, point Q is defined by a position vector  $\mathbf{p} = \mathbf{p}_0 + \mathbf{q}$  where  $\mathbf{q} = y\mathbf{d}_2 + z\mathbf{d}_3$  defines point Q in the  $y-z$  plane from the elastic axis. Then the undeformed local airfoil coordinates of point Q are

$$\begin{bmatrix} y \\ z \end{bmatrix} = \begin{bmatrix} \cos \gamma & -\sin \gamma \\ \sin \gamma & \cos \gamma \end{bmatrix} \begin{bmatrix} \eta \\ \xi \end{bmatrix} \quad (1)$$

where  $\eta$  and  $\xi$  are local airfoil coordinates, and  $\gamma$  is the wing section pre-twist angle, positive nose-down.<sup>5</sup>

Differentiating  $y$  and  $z$  with respect to  $x$  gives

$$\begin{bmatrix} y_x \\ z_x \end{bmatrix} = \gamma' \begin{bmatrix} -\sin \gamma & -\cos \gamma \\ \cos \gamma & -\sin \gamma \end{bmatrix} \begin{bmatrix} \eta \\ \xi \end{bmatrix} = \begin{bmatrix} -z\gamma' \\ y\gamma' \end{bmatrix} \quad (2)$$

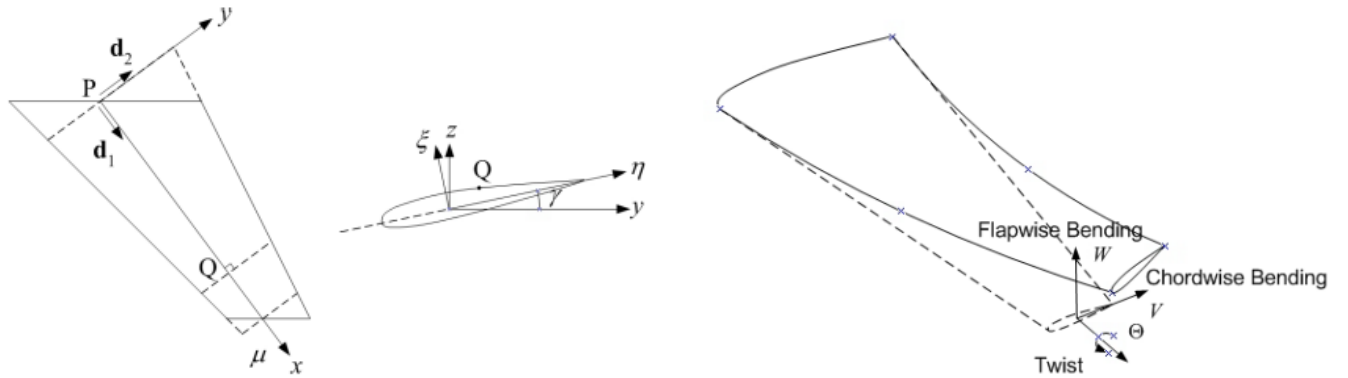


Fig. 2 - Left Wing Reference Frame of Wing in Combined Bending-Torsion

Let  $\Theta$  be a torsional twist angle about the  $x$ -axis, positive nose-down. Let  $W$  and  $V$  be flapwise and chordwise bending deflections of point Q, respectively. Let  $U$  be the axial displacement of point Q. Then, the displacement and rotation vectors due to the elastic deformation can be expressed as

$$\mathbf{r} = U\mathbf{d}_1 + V\mathbf{d}_2 + W\mathbf{d}_3 \quad (3)$$

$$\phi = \Theta \mathbf{d}_1 - W_x \mathbf{d}_2 + V_x \mathbf{d}_3 \quad (4)$$

where the subscripts  $x$  and  $t$  denote the partial derivatives of  $\Theta$ ,  $W$ , and  $V$ .

Let  $(x_1, y_1, z_1)$  be the deformed coordinates of point Q on the airfoil in the reference frame D and  $\mathbf{p}_1 = x_1 \mathbf{d}_1 + y_1 \mathbf{d}_2 + z_1 \mathbf{d}_3$  be its position vector. Then the coordinates  $(x_1, y_1, z_1)$  are computed as

$$\mathbf{p}_1 = \mathbf{p} + \mathbf{r} + \phi \times \mathbf{q} \quad (5)$$

where

$$\begin{bmatrix} x_1 \\ y_1 \\ z_1 \end{bmatrix} = \begin{bmatrix} x + U - yV_x - zW_x \\ y + V - z\Theta \\ z + W + y\Theta \end{bmatrix} \quad (6)$$

Differentiating  $x_1$ ,  $y_1$ , and  $z_1$  with respect to  $x$  yields

$$\begin{bmatrix} x_{1,x} \\ y_{1,x} \\ z_{1,x} \end{bmatrix} = \begin{bmatrix} 1 + U_x - yV_{xx} + z\gamma'V_x - zW_{xx} - y\gamma'W_x \\ -z\gamma' + V_x - z\Theta_x - y\gamma'\Theta \\ y\gamma' + W_x + y\Theta_x - z\gamma'\Theta \end{bmatrix} \quad (7)$$

Neglecting the transverse shear effect, the longitudinal strain is computed as<sup>6</sup>

$$\varepsilon = \frac{ds_1 - ds}{ds} = \frac{s_{1,x}}{s_x} - 1 \quad (8)$$

where

$$s_x = \sqrt{1 + y_x^2 + z_x^2} = \sqrt{1 + (y^2 + z^2)(\gamma')^2} \quad (9)$$

$$\begin{aligned} s_{1,x} &= \sqrt{x_{1,x}^2 + y_{1,x}^2 + z_{1,x}^2} \\ &= \sqrt{s_x^2 + 2U_x - 2yV_{xx} - 2zW_{xx} + 2(y^2 + z^2)\gamma'\Theta_x + (x_{1,x} - 1)^2 + (y_{1,x} + z\gamma')^2 + (z_{1,x} - y\gamma')^2} \end{aligned} \quad (10)$$

Using the Taylor series expansion,  $s_{1,x}$  is approximated as

$$s_{1,x} \approx s_x + \frac{U_x - yV_{xx} - zW_{xx} + (y^2 + z^2)\gamma'\Theta_x}{s_x} + \frac{(x_{1,x} - 1)^2 + (y_{1,x} + z\gamma')^2 + (z_{1,x} - y\gamma')^2}{2s_x} \quad (11)$$

The longitudinal strain is then obtained as

$$\varepsilon = \frac{U_x - yV_{xx} - zW_{xx} + (y^2 + z^2)\gamma'\Theta_x}{s_x^2} + \frac{(x_{1,x} - 1)^2 + (y_{1,x} + z\gamma')^2 + (z_{1,x} - y\gamma')^2}{2s_x^2} \quad (12)$$

For a small wing twist angle  $\gamma$ ,  $(\gamma')^2 \approx 0$ . Then

$$\varepsilon = U_x - yV_{xx} - zW_{xx} + (y^2 + z^2)\gamma'\Theta_x + \frac{(x_{1,x} - 1)^2 + (y_{1,x} + z\gamma')^2 + (z_{1,x} - y\gamma')^2}{2} \quad (13)$$

The axial force and moments acting on a wing are then obtained as<sup>6</sup>

$$\begin{aligned}
 \begin{bmatrix} P_x \\ M_x \\ M_y \\ M_z \end{bmatrix} &= \begin{bmatrix} \iint E \varepsilon dydz \\ GJ\Theta_x + \iint E \varepsilon (y^2 + z^2) (\gamma' + \Theta_x) dydz \\ -\iint E \varepsilon z dydz \\ -\iint E \varepsilon y dydz \end{bmatrix} \\
 &= \begin{bmatrix} EA & EI_{xx}\gamma' & -EAe_z & -EAe_y \\ EI_{xx}\gamma' & \left[ GJ + EB_1 (\gamma')^2 + EI_{xx}U_x \right] & -EB_2\gamma' & -EB_3\gamma' \\ -EAe_z & -EB_2\gamma' & EI_{yy} & -EI_{yz} \\ -EAe_y & -EB_3\gamma' & -EI_{yz} & EI_{zz} \end{bmatrix} \begin{bmatrix} U_x \\ \Theta_x \\ W_{xx} \\ V_{xx} \end{bmatrix} + \begin{bmatrix} \Delta P_x \\ \Delta M_x \\ \Delta M_y \\ \Delta M_z \end{bmatrix} \quad (14)
 \end{aligned}$$

where  $\Delta P_x$  and  $\Delta M_{(x,y,z)}$  are the nonlinear terms;  $E$  is the Young's modulus;  $G$  is the shear modulus;  $A$  is the tensile area;  $e_y$  and  $e_z$  are the offsets of the centroid from the elastic axis;  $I_{yy}$ ,  $I_{yz}$ , and  $I_{zz}$  are the section area moments of inertia about the flapwise axis;  $J$  is the torsional constant; and  $B_1$ ,  $B_2$ , and  $B_3$  are the bending-torsion coupling constants which are defined as

$$\begin{bmatrix} B_1 \\ B_2 \\ B_3 \end{bmatrix} = \int \int \begin{bmatrix} y^2 + z^2 \\ z \\ y \end{bmatrix} (y^2 + z^2) dydz \quad (15)$$

The slope of the twist angle  $\gamma'$  can play a significant role in structures with large twists such as turbomachinery blades. For aircraft wings, this effect is negligible and therefore can be neglected. Also, in general, the offset of the centroid from the elastic axis in the  $z$ -direction,  $e_z$ , is small compared to that in the  $y$ -direction  $e_y$ . Therefore, it may be neglected. Then, the moments acting a wing are reduced to

$$\begin{bmatrix} P_x \\ M_x \\ M_y \\ M_z \end{bmatrix} = \begin{bmatrix} EAU_x - EAe_y V_{xx} \\ (GJ + EI_{xx}U_x) \Theta_x \\ EI_{yy}W_{xx} - EI_{yz}V_{xx} \\ -EAe_y U_x - EI_{yz}W_{xx} + EI_{zz}V_{xx} \end{bmatrix} + \begin{bmatrix} \Delta P_x \\ \Delta M_x \\ \Delta M_y \\ \Delta M_z \end{bmatrix} \quad (16)$$

Note that the signs of the moments are defined in the positive deflection sense such that

$$\mathbf{M} = M_x \mathbf{d}_1 - M_y \mathbf{d}_2 + M_z \mathbf{d}_3 \quad (17)$$

The nonlinear terms are given by

$$\begin{aligned}
\begin{bmatrix} \Delta P_x \\ \Delta M_x \\ \Delta M_y \\ \Delta M_z \end{bmatrix} &= \frac{1}{2} E \int \int \left[ (U_x - yV_{xx} - zW_{xx})^2 + (V_x - z\Theta_x)^2 + (W_x + y\Theta_x)^2 \right] \begin{bmatrix} 1 \\ (y^2 + z^2)\Theta_x \\ -z \\ -y \end{bmatrix} dydz \\
&= \frac{1}{2} (U_x^2 + V_x^2 + W_x^2) \begin{bmatrix} EA \\ EI_{xx}\Theta_x \\ -EAe_z \\ -EAe_y \end{bmatrix} + \frac{1}{2} V_{xx}^2 \begin{bmatrix} EI_{zz} \\ EB_4\Theta_x \\ -EB_7 \\ -EB_{10} \end{bmatrix} + \frac{1}{2} W_{xx}^2 \begin{bmatrix} EI_{yy} \\ EB_5\Theta_x \\ -EB_8 \\ -EB_9 \end{bmatrix} + \frac{1}{2} \Theta_x^2 \begin{bmatrix} EI_{xx} \\ EB_1\Theta_x \\ -EB_2 \\ -EB_3 \end{bmatrix} \\
&+ U_x V_{xx} \begin{bmatrix} -EAe_y \\ -EB_3\Theta_x \\ -EI_{yz} \\ EI_{zz} \end{bmatrix} + U_x W_{xx} \begin{bmatrix} -EAe_z \\ -EB_2\Theta_x \\ EI_{yy} \\ -EI_{yz} \end{bmatrix} + V_{xx} W_{xx} \begin{bmatrix} -EI_{yz} \\ EB_6\Theta_x \\ -EB_9 \\ -EB_7 \end{bmatrix} + V_x \Theta_x \begin{bmatrix} -EAe_z \\ -EB_2\Theta_x \\ EI_{yy} \\ -EI_{yz} \end{bmatrix} \\
&+ W_x \Theta_x \begin{bmatrix} EAe_y \\ EB_3\Theta_x \\ EI_{yz} \\ -EI_{zz} \end{bmatrix} \tag{18}
\end{aligned}$$

where

$$\begin{bmatrix} B_4 \\ B_5 \\ B_6 \\ B_7 \\ B_8 \\ B_9 \\ B_{10} \end{bmatrix} = \int \int \begin{bmatrix} y^2 (y^2 + z^2) \\ z^2 (y^2 + z^2) \\ yz (y^2 + z^2) \\ y^2 z \\ z^3 \\ yz^2 \\ y^3 \end{bmatrix} dydz \tag{19}$$

Consider two cases:

1. If the nonlinear strain component due to axial displacement  $(x_{1,x} - 1)^2 \approx 0$ , then the nonlinear terms are simplified to

$$\begin{bmatrix} \Delta P_x \\ \Delta M_x \\ \Delta M_y \\ \Delta M_z \end{bmatrix} = \frac{1}{2} (V_x^2 + W_x^2) \begin{bmatrix} EA \\ EI_{xx}\Theta_x \\ -EAe_z \\ -EAe_y \end{bmatrix} + \frac{1}{2} \Theta_x^2 \begin{bmatrix} EI_{xx} \\ EB_1\Theta_x \\ -EB_2 \\ -EB_3 \end{bmatrix} + V_x \Theta_x \begin{bmatrix} -EAe_z \\ -EB_2\Theta_x \\ EI_{yy} \\ -EI_{yz} \end{bmatrix} + W_x \Theta_x \begin{bmatrix} EAe_y \\ EB_3\Theta_x \\ EI_{yz} \\ -EI_{zz} \end{bmatrix} \tag{20}$$

2. If, in addition, the nonlinear strain component due to chordwise deflection  $(y_{1,x} + z\gamma)^2 \approx 0$ , then the nonlinear

terms are simplified to

$$\begin{bmatrix} \Delta P_x \\ \Delta M_x \\ \Delta M_y \\ \Delta M_z \end{bmatrix} = \frac{1}{2} W_x^2 \begin{bmatrix} EA \\ EI_{xx} \Theta_x \\ -EAe_z \\ -EAe_y \end{bmatrix} + \frac{1}{2} \Theta_x^2 \begin{bmatrix} EI_{yy} \\ EB_4 \Theta_x \\ -EB_7 \\ -EB_{10} \end{bmatrix} + W_x \Theta_x \begin{bmatrix} EAe_y \\ EB_3 \Theta_x \\ EI_{yz} \\ -EI_{zz} \end{bmatrix} \quad (21)$$

### III. Aeroelastic Angle of Attack

The relative velocity of the air approaching a wing section includes the contribution from the wing elastic deflection that results in changes in the local angle of attack. Since aerodynamic forces and moments are dependent on the local angle of attack, the wing aeroelastic deflections will generate additional elastic forces and moments. The local angle of attack depends on the relative approaching air velocity as well as the rotation angle  $\phi$  from Eq. (4). The relative air velocity in turn also depends on the deflection-induced velocity. The velocity at point Q due to the aircraft velocity and angular velocity in the reference frame D is computed as

$$\begin{aligned} \mathbf{v}_Q &= \bar{\mathbf{v}} + \boldsymbol{\omega} \times \mathbf{r} = (u\mathbf{b}_1 + v\mathbf{b}_2 + w\mathbf{b}_3) + (p\mathbf{b}_1 + q\mathbf{b}_2 + r\mathbf{b}_3) \times (-x_a\mathbf{b}_1 - y_a\mathbf{b}_2 - z_a\mathbf{b}_3) \\ &= (u + ry_a - qz_a)\mathbf{b}_1 + (v - rx_a + pz_a)\mathbf{b}_2 + (w + qx_a - py_a)\mathbf{b}_3 \\ &= x_t\mathbf{d}_1 + y_t\mathbf{d}_2 + z_t\mathbf{d}_3 \end{aligned} \quad (22)$$

where

$$\begin{bmatrix} \mathbf{b}_1 \\ \mathbf{b}_2 \\ \mathbf{b}_3 \end{bmatrix} = \begin{bmatrix} -\sin\Lambda \cos\Gamma & -\cos\Lambda & \sin\Lambda \sin\Gamma \\ -\cos\Lambda \cos\Gamma & \sin\Lambda & \cos\Lambda \sin\Gamma \\ -\sin\Gamma & 0 & -\cos\Gamma \end{bmatrix} \begin{bmatrix} \mathbf{d}_1 \\ \mathbf{d}_2 \\ \mathbf{d}_3 \end{bmatrix} \quad (23)$$

$$\begin{bmatrix} x_t \\ y_t \\ z_t \end{bmatrix} = \begin{bmatrix} -(u + ry_a - qz_a) \sin\Lambda \cos\Gamma - (v - rx_a + pz_a) \cos\Lambda \cos\Gamma - (w + qx_a - py_a) \sin\Gamma \\ -(u + ry_a - qz_a) \cos\Lambda + (v - rx_a + pz_a) \sin\Lambda \\ (u + ry_a - qz_a) \sin\Lambda \sin\Gamma + (v - rx_a + pz_a) \cos\Lambda \sin\Gamma - (w + qx_a - py_a) \cos\Gamma \end{bmatrix} \quad (24)$$

and  $(p, q, r)$  are aircraft angular velocity components in the roll, pitch, and yaw axes, and  $(x_a, y_a, z_a)$  is the coordinate of point Q in the aircraft body-fixed reference frame B relative to the aircraft C.G. (center of gravity) such that  $x_a$  is positive when point Q is aft of the aircraft CG,  $y_a$  is positive when point Q is toward the left wing from the aircraft C.G., and  $z_a$  is positive when point Q is above the aircraft C.G.

Consider a trim problem when  $\beta = 0$ ,  $p = q = r = 0$ . Then

$$\begin{bmatrix} x_t \\ y_t \\ z_t \end{bmatrix} = \begin{bmatrix} -u \sin\Lambda \cos\Gamma - w \sin\Gamma \\ -u \cos\Lambda \\ u \sin\Lambda \sin\Gamma - w \cos\Gamma \end{bmatrix} \quad (25)$$

The local velocity at point Q due to aircraft rigid-body dynamics and aeroelastic deflections in the reference frame D is obtained as<sup>7</sup>

$$\mathbf{v} = \mathbf{v}_Q + \frac{\partial \mathbf{r}}{\partial t} + \frac{\partial \phi}{\partial t} \times (\mathbf{r} + \phi \times \mathbf{q}) = v_x\mathbf{d}_1 + v_y\mathbf{d}_2 + v_z\mathbf{d}_3 \quad (26)$$

where

$$\begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix} = \begin{bmatrix} x_t + U_t - (z + W + y\Theta)W_{xt} - (y + V - z\Theta)V_{xt} \\ y_t + V_t + (U - yV_x - zW_x)V_{xt} - (z + W + y\Theta)\Theta_t \\ z_t + W_t + (U - yV_x - zW_x)W_{xt} + (y + V - z\Theta)\Theta_t \end{bmatrix} \quad (27)$$

In order to compute the aeroelastic forces and moments, the velocity must be transformed from the reference frame D to the airfoil local coordinate reference frame defined by  $(\mu, \eta, \xi)$  as shown in Figure 2. Then the transformation can be performed using successive rotation matrix multiplication operations as

$$\begin{aligned} \begin{bmatrix} v_\mu \\ v_\eta \\ v_\xi \end{bmatrix} &= \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos(\Theta + \gamma) & \sin(\Theta + \gamma) \\ 0 & -\sin(\Theta + \gamma) & \cos(\Theta + \gamma) \end{bmatrix} \begin{bmatrix} \cos V_x & \sin V_x & 0 \\ -\sin V_x & \cos V_x & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \cos W_x & 0 & \sin W_x \\ 0 & 1 & 0 \\ -\sin W_x & 0 & \cos W_x \end{bmatrix} \begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix} \\ &= \begin{bmatrix} \cos V_x (v_x \cos W_x + v_z \sin W_x) + v_y \sin V_x \\ \cos(\Theta + \gamma) [-\sin V_x (v_x \cos W_x + v_z \sin W_x) + v_y \cos V_x] + \sin(\Theta + \gamma) (-v_x \sin W_x + v_z \cos W_x) \\ -\sin(\Theta + \gamma) [-\sin V_x (v_x \cos W_x + v_z \sin W_x) + v_y \cos V_x] + \cos(\Theta + \gamma) (-v_x \sin W_x + v_z \cos W_x) \end{bmatrix} \\ &\approx \begin{bmatrix} v_x + v_y V_x + v_z W_x \\ -v_x [V_x + W_x (\Theta + \gamma)] + v_y + v_z [(\Theta + \gamma) - V_x W_x] \\ v_x [-W_x + V_x (\Theta + \gamma)] - v_y (\Theta + \gamma) + v_z [1 + V_x W_x (\Theta + \gamma)] \end{bmatrix} \end{aligned} \quad (28)$$

for small deflections.

The local aeroelastic angle of attack on the airfoil section due to the velocity components  $v_\eta$  and  $v_\xi$  in the reference frame D, as shown in Fig. 8, is computed as

$$\alpha_c = \frac{v_\xi + w_i}{v_\eta} = \frac{\bar{v}_\xi + \Delta v_\xi + w_i}{\bar{v}_\eta + \Delta v_\eta} = \frac{v_\xi + w_i}{\bar{v}_\eta} - \frac{(\bar{v}_\xi + w_i) \Delta v_\eta}{\bar{v}_\eta^2} \quad (29)$$

where  $w_i$  is the downwash due to the three-dimensional lift distribution over a finite-aspect ratio wing and

$$\bar{v}_\xi = u \sin \Lambda \sin \Gamma - w \cos \Gamma \quad (30)$$

$$\Delta v_\xi = W_t + (U - yV_x - zW_x)W_{xt} + (y + V - z\Theta)\Theta_t + v_x [-W_x + V_x (\Theta + \gamma)] - v_y (\Theta + \gamma) + v_z V_x W_x (\Theta + \gamma) \quad (31)$$

$$\bar{v}_\eta = -u \cos \Lambda \quad (32)$$

$$\Delta v_\eta = V_t + (U - yV_x - zW_x)V_{xt} - (z + W + y\Theta)\Theta_t - v_x [V_x + W_x (\Theta + \gamma)] + v_z [(\Theta + \gamma) - V_x W_x] \quad (33)$$

Assuming that  $z \approx 0$  and neglecting any third-order terms or higher, the local aeroelastic angle of attack is evaluated



as

$$\begin{aligned}
\alpha_c = & -\frac{u \sin \Lambda \sin \Gamma - w \cos \Gamma + w_i}{u \cos \Lambda} - \frac{W_t + (U - yV_x)W_{xt} + (y + V)\Theta_t}{u \cos \Lambda} \\
& - \frac{[-u \sin \Lambda \cos \Gamma - w \sin \Gamma + U_t - (W + y\Theta)W_{xt} - (y + V)V_{xt}][ -W_x + V_x(\Theta + \gamma)]}{u \cos \Lambda} \\
& + \frac{[-u \cos \Lambda + V_t + (U - yV_x)V_{xt} - (W + y\Theta)\Theta_t](\Theta + \gamma)}{u \cos \Lambda} \\
& - \frac{u \sin \Lambda \sin \Gamma - w \cos \Gamma + w_i}{u^2 \cos^2 \Lambda} \left\{ V_t + (U - yV_x)V_{xt} - (W + y\Theta)\Theta_t \right. \\
& - [-u \sin \Lambda \cos \Gamma - w \sin \Gamma + U_t - (W + y\Theta)W_{xt} - (y + V)V_{xt}][V_x + W_x(\Theta + \gamma)] \\
& \left. + [u \sin \Lambda \sin \Gamma - w \cos \Gamma + W_t + (U - yV_x)W_{xt} + (y + V)\Theta_t][(\Theta + \gamma) - V_x W_x] \right\} \quad (34)
\end{aligned}$$

Let  $u \approx V_\infty$ ,  $w \approx V_\infty \alpha$ , and  $w_i = V_\infty \alpha_i$ , where  $\alpha_i$  is the induced angle of attack along the wing span. Then the nonlinear partial derivatives of the local aeroelastic angle of attack are evaluated as

$$\alpha_0 = -\gamma - \frac{(\sin \Lambda \sin \Gamma + \alpha_i)(1 + \gamma \tan \Lambda \sin \Gamma)}{\cos \Lambda} \quad (35)$$

$$\frac{\partial \alpha_c}{\partial \alpha} = \frac{\cos \Gamma}{\cos \Lambda} \left[ 1 + \frac{\gamma(2 \sin \Lambda \sin \Gamma + \alpha_i)}{\cos \Lambda} \right] \quad (36)$$

$$\frac{\partial \alpha_c}{\partial \alpha^2} = -\frac{\gamma \cos^2 \Gamma}{\cos^2 \Lambda} \quad (37)$$

$$\frac{\partial \alpha_c}{\partial \Theta} = -1 - \frac{(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma)(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{\cos^2 \Lambda} \quad (38)$$

$$\frac{\partial \alpha_c}{\partial W_x} = -\frac{\sin \Lambda \cos \Gamma + \alpha \sin \Gamma}{\cos \Lambda} \left[ 1 + \frac{\gamma(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{\cos \Lambda} \right] \quad (39)$$

$$\frac{\partial \alpha_c}{\partial V_x} = \frac{\sin \Lambda \cos \Gamma + \alpha \sin \Gamma}{\cos \Lambda} \left( \gamma - \frac{\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i}{\cos \Lambda} \right) \quad (40)$$

$$\frac{\partial \alpha_c}{\partial U_t} = \frac{W_x}{V_\infty \cos \Lambda} + \frac{V_x(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{V_\infty \cos^2 \Lambda} \quad (41)$$

$$\frac{\partial \alpha_c}{\partial \Theta_t} = \frac{y + V}{V_\infty \cos \Lambda} - \frac{(y\gamma - W)(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{V_\infty \cos^2 \Lambda} \quad (42)$$

$$\frac{\partial \alpha_c}{\partial W_t} = \frac{1}{V_\infty \cos \Lambda} - \frac{(\Theta + \gamma)(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{V_\infty \cos^2 \Lambda} \quad (43)$$

$$\frac{\partial \alpha_c}{\partial V_t} = \frac{\Theta + \gamma}{V_\infty \cos \Lambda} - \frac{\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i}{V_\infty \cos^2 \Lambda} \quad (44)$$

$$\frac{\partial \alpha_c}{\partial \Theta W_x} = -\frac{(\sin \Lambda \cos \Gamma + \alpha \sin \Gamma)(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{\cos^2 \Lambda} \quad (45)$$

$$\frac{\partial \alpha_c}{\partial \Theta V_x} = \frac{\sin \Lambda \cos \Gamma + \alpha \sin \Gamma}{\cos \Lambda} \quad (46)$$

$$\frac{\partial \alpha_c}{\partial W_x V_x} = \frac{(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma)(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{\cos^2 \Lambda} \quad (47)$$

$$\frac{\partial \alpha_c}{\partial W_{xt}} = -\frac{U - yV_x}{V_\infty \cos \Lambda} \quad (48)$$

$$\frac{\partial \alpha_c}{\partial V_{xt}} = -\frac{yW_x}{V_\infty \cos \Lambda} - \frac{U(\sin \Lambda \sin \Gamma - \alpha \cos \Gamma + \alpha_i)}{V_\infty \cos^2 \Lambda} \quad (49)$$

We note that the time derivative of the axial displacement contributes to the aeroelastic angle of attack as a damping

source. Otherwise, the axial displacement  $U$  is weakly coupled to the wing aeroelasticity.

The nonlinear angle of attack is then expressed as

$$\alpha_c(x, y) = \alpha_r(x) + \alpha_e(x, y) \quad (50)$$

where  $\alpha_r(x)$  is the rigid angle of attack and  $\alpha_e(x, y)$  is the elastic angle of attack

$$\alpha_r(x, y) = \alpha_0 + \frac{\partial \alpha_c}{\partial \alpha} \alpha + \frac{\partial \alpha_c}{\partial \alpha^2} \alpha^2 \quad (51)$$

$$\begin{aligned} \alpha_e(x, y) = & \frac{\partial \alpha_c}{\partial \Theta} \Theta + \frac{\partial \alpha_c}{\partial W_x} W_x + \frac{\partial \alpha_c}{\partial V_x} V_x + \frac{\partial \alpha_c}{\partial \Theta W_x} \Theta W_x + \frac{\partial \alpha_c}{\partial \Theta V_x} \Theta V_x + \frac{\partial \alpha_c}{\partial W_x V_x} W_x V_x \\ & + \frac{\partial \alpha_c}{\partial U_t} U_t + \frac{\partial \alpha_c}{\partial \Theta_t} \Theta_t + \frac{\partial \alpha_c}{\partial W_t} W_t + \frac{\partial \alpha_c}{\partial V_t} V_t + \frac{\partial \alpha_c}{\partial W_{xt}} W_{xt} + \frac{\partial \alpha_c}{\partial V_{xt}} V_{xt} \end{aligned} \quad (52)$$

The elastic angle of attack can be linearized as a function of the static deflection as

$$\begin{aligned} \alpha_e(x, y) = & \left( \frac{\partial \alpha_c}{\partial \Theta} + \frac{\partial \alpha_c}{\partial \Theta W_x} \bar{W}_x + \frac{\partial \alpha_c}{\partial \Theta V_x} \bar{V}_x \right) \Theta + \left( \frac{\partial \alpha_c}{\partial W_x} + \frac{\partial \alpha_c}{\partial \Theta W_x} \bar{\Theta} + \frac{\partial \alpha_c}{\partial W_x V_x} \bar{V}_x \right) W_x \\ & + \left( \frac{\partial \alpha_c}{\partial V_x} + \frac{\partial \alpha_c}{\partial \Theta V_x} \bar{\Theta} + \frac{\partial \alpha_c}{\partial W_x V_x} \bar{W}_x \right) V_x + \frac{\partial \alpha_c}{\partial U_t} U_t + \frac{\partial \alpha_c}{\partial \Theta_t} \Theta_t + \frac{\partial \alpha_c}{\partial V_t} V_t + \frac{\partial \alpha_c}{\partial W_t} W_t + \frac{\partial \alpha_c}{\partial V_{xt}} V_{xt} + \frac{\partial \alpha_c}{\partial W_{xt}} W_{xt} \end{aligned} \quad (53)$$

Finally, if we ignore all the deflection-dependent terms, then the linearized elastic angle of attack becomes

$$\alpha_e(x, y) = \frac{\partial \alpha_c}{\partial \Theta} \Theta + \frac{\partial \alpha_c}{\partial W_x} W_x + \frac{\partial \alpha_c}{\partial V_x} V_x + \frac{\partial \alpha_c}{\partial \Theta_t} \Theta_t + \frac{\partial \alpha_c}{\partial W_t} W_t + \frac{\partial \alpha_c}{\partial V_t} V_t \quad (54)$$

#### IV. Aerodynamic Forces and Moments

The wing section unsteady lift coefficient is given by

$$c_L = c_{L\alpha} [\alpha_r + C(k) \alpha_{ac}] \cos \Lambda + \frac{\pi \dot{\alpha}_{ac} \cos \Lambda c}{2V_\infty} \quad (55)$$

where  $k = \frac{\omega c}{2V_\infty}$  is the reduced frequency parameter,  $\omega$  is the frequency of wing oscillations,  $c$  is the section chord in the streamwise direction,  $c_{L\alpha}$  is the section lift curve slope,  $\alpha_{ac}$  is the effective elastic angle of attack at the three-quarter-chord point,<sup>5</sup> and  $C(k)$  is the Theodorsen's complex-valued function.<sup>8</sup>

The section pitching moment coefficient about the aircraft pitch axis is evaluated as

$$c_m = c_{m_{ac}} + \frac{e}{c} c_{L\alpha} [\alpha_r + C(k) \alpha_{ac}] - \frac{e_u}{c} \frac{\pi \dot{\alpha}_{ac} c}{2V_\infty} \quad (56)$$

where  $c_{m_{ac}}$  is the section pitching moment coefficient at the aerodynamic center, positive nose up,  $e$  is the offset of the aerodynamic center forward of the elastic axis, and  $e_u$  is the offset of the three-quarter chord point aft of the elastic axis.

The section drag coefficient is expressed in a parabolic drag polar form as

$$c_D = c_{D_0} + K c_L^2 \quad (57)$$

where  $c_{D_0}$  is the section parasitic drag coefficient and  $K$  is a parabolic drag parameter.

The force and moment coefficients can be written as

$$c_L = c_{L_r} + \frac{\partial c_L}{\partial \Theta} \Theta + \frac{\partial c_L}{\partial W_x} W_x + \frac{\partial c_L}{\partial V_x} V_x + \frac{\partial c_L}{\partial \Theta_t} \Theta_t + \frac{\partial c_L}{\partial W_t} W_t + \frac{\partial c_L}{\partial V_t} V_t + \frac{\partial c_L}{\partial W_{xt}} W_{xt} + \frac{\partial c_L}{\partial V_{xt}} V_{xt} + \frac{\partial c_L}{\partial \Theta_{tt}} \Theta_{tt} + \frac{\partial c_L}{\partial W_{tt}} W_{tt} + \frac{\partial c_L}{\partial V_{tt}} V_{tt} \quad (58)$$

$$c_m = c_{m_r} + \frac{\partial c_m}{\partial \Theta} \Theta + \frac{\partial c_m}{\partial W_x} W_x + \frac{\partial c_m}{\partial V_x} V_x + \frac{\partial c_m}{\partial \Theta_t} \Theta_t + \frac{\partial c_m}{\partial W_t} W_t + \frac{\partial c_m}{\partial V_t} V_t + \frac{\partial c_m}{\partial W_{xt}} W_{xt} + \frac{\partial c_m}{\partial V_{xt}} V_{xt} + \frac{\partial c_m}{\partial \Theta_{tt}} \Theta_{tt} + \frac{\partial c_m}{\partial W_{tt}} W_{tt} + \frac{\partial c_m}{\partial V_{tt}} V_{tt} \quad (59)$$

$$c_D = c_{D_r} + \frac{\partial c_D}{\partial c_L} (c_L - c_{L_r}) = c_{D_r} + 2Kc_{L_r} (c_L - c_{L_r}) \quad (60)$$

where  $c_{L_r}$ ,  $c_{m_r}$ , and  $c_{D_r}$  are the rigid lift, pitching moment, and drag coefficient contributions.

The lift force, drag force, and pitching moment in the aircraft coordinate reference frame along the elastic axis are expressed as

$$l = c_L q_\infty \cos \Lambda c \quad (61)$$

$$d = c_D q_\infty \cos \Lambda c \quad (62)$$

$$m = c_m q_\infty \cos \Lambda c^2 \quad (63)$$

The axial and normal force coefficients are defined as

$$c_n = c_L \cos \alpha + c_D \sin \alpha \approx c_L \quad (64)$$

$$c_a = -c_L \sin \alpha + c_D \cos \alpha \quad (65)$$

The forces and moments in the local coordinate reference frame are obtained as

$$f_x^a = c_n q_\infty c \cos \Lambda \sin \Gamma + c_a q_\infty c \cos \Lambda \sin \Lambda \cos \Gamma \approx 0 \quad (66)$$

$$f_y^a = c_a q_\infty c \cos^2 \Lambda \quad (67)$$

$$f_z^a \approx c_n q_\infty c \cos \Lambda \cos \Gamma - c_a q_\infty c \cos \Lambda \sin \Lambda \sin \Gamma \approx c_L q_\infty c \cos \Lambda \cos \Gamma \quad (68)$$

$$m_x^a = -c_m q_\infty c^2 \cos^2 \Lambda \cos \Gamma \quad (69)$$

$$m_y^a = c_m q_\infty c^2 \sin \Lambda \cos \Lambda \quad (70)$$

$$m_z^a = c_m q_\infty c^2 \cos^2 \Lambda \sin \Gamma \quad (71)$$

## A. Propulsive Forces and Moments

In addition, the propulsive effects of the aircraft engines must be accounted for in the analysis. Both the engine mass and thrust can contribute to the aeroelasticity.<sup>7</sup> The propulsive force and moment vector are expressed in the reference frame D as

$$\begin{aligned} \mathbf{f}^e &= \delta (x - x_e) (F_x^e \mathbf{d}_1 + F_y^e \mathbf{d}_2 + F_z^e \mathbf{d}_3) \\ &= \delta (x - x_e) [(-T \sin \Lambda \cos \Gamma - m_e g \sin \Gamma) \mathbf{d}_1 - T \cos \Lambda \mathbf{d}_2 + (T \sin \Lambda \sin \Gamma - m_e g \cos \Gamma) \mathbf{d}_3] \end{aligned} \quad (72)$$

$$\mathbf{m}^e = \mathbf{r}_e \times \mathbf{f}_e = (x_e \mathbf{d}_1 - y_e \mathbf{d}_2 - z_e \mathbf{d}_3) \times \mathbf{f}_e \quad (73)$$

where  $T$  is the engine thrust,  $m_e$  is the engine mass,  $(x_e, y_e, z_e)$  is the coordinate of the engine thrust center such that  $y_e > 0$  forward of the elastic axis and  $z_e > 0$  below the elastic axis, and  $\delta(x - x_e)$  is the Dirac delta function such that

$$\int \delta(x - x_e) f(x) dx = f(x_e) \quad (74)$$

Transforming into the local coordinate reference frame and neglecting nonlinear contributions, the propulsive forces and moments are given by

$$f_x^e = \delta(x - x_e) (F_x^e + F_z^e W_x + F_y^e V_x) \quad (75)$$

$$f_y^e = \delta(x - x_e) [F_y^e + F_z^e (\Theta + \gamma) - F_x^e V_x] \quad (76)$$

$$f_z^e = \delta(x - x_e) [F_z^e - F_y^e (\Theta + \gamma) - F_x^e W_x] \quad (77)$$

$$m_x^e = \delta(x - x_e) [F_y^e z_e - F_z^e y_e + (F_x y_e + F_y x_e) W_x - (F_x z_e + F_z x_e) V_x] \quad (78)$$

$$m_y^e = \delta(x - x_e) [-F_x^e z_e - F_z^e x_e + (F_x y_e + F_y x_e) (\Theta + \gamma) - (F_y z_e - F_z y_e) V_x] \quad (79)$$

$$m_z^e = \delta(x - x_e) [F_x y_e + F_y x_e + (F_x z_e + F_z x_e) (\Theta + \gamma) - (F_y z_e - F_z y_e) W_x] \quad (80)$$

## B. Inertial Forces and Moments

The inertial forces and moments are due to acceleration experienced by the wing, the air volume surrounding it which gives rise to the apparent mass acting at mid-chord, as well as the engine mass and half of the mass of the fuselage and tail empennage. The acceleration results in the inertial forces and moments as follows:

$$f_x^i = -mg \sin \Gamma - mU_{tt} - \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda U_{tt} - \delta(x) \frac{1}{2} m_f g \sin \Gamma \quad (81)$$

$$f_y^i = -mV_{tt} - \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda V_{tt} + \delta(x - x_e) (-m_e z_e \Theta_{tt} - m_e V_{tt}) + \delta(x) \left( \frac{1}{2} m_f z_f \Theta_{tt} - \frac{1}{2} m_f V_{tt} \right) \quad (82)$$

$$f_z^i = -mg \cos \Gamma + m e_{cg} \Theta_{tt} + \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda e_m \Theta_{tt} - mW_{tt} - \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda W_{tt} + \delta(x - x_e) (m_e y_e \Theta_{tt} - m_e W_{tt}) + \delta(x) \left( -\frac{1}{2} m_f g \cos \Gamma - \frac{1}{2} m_f y_f \Theta_{tt} - \frac{1}{2} m_f W_{tt} \right) \quad (83)$$

$$m_x^i = m g e_{cg} \cos \Gamma - m r_k^2 \Theta_{tt} - \frac{\pi}{4} \rho_\infty c^2 e_m^2 \cos \Lambda \Theta_{tt} - \frac{\pi}{128} \rho_\infty c^4 \cos \Lambda \Theta_{tt} + m e_{cg} W_{tt} + \frac{\pi}{4} \rho_\infty c^2 e_m \cos \Lambda W_{tt} + \delta(x - x_e) (-m_e r_e^2 \Theta_{tt} + m_e y_e W_{tt} - m_e z_e V_{tt}) + \delta(x) \left( -\frac{1}{2} m_f g y_f \cos \Gamma - \frac{1}{2} I_{xx}^f \Theta_{tt} - \frac{1}{2} m_f y_f W_{tt} + \frac{1}{2} m_f z_f V_{tt} \right) \quad (84)$$

$$m_y^i = \delta(x) \left( -\frac{1}{2} m_f g z_f \sin \Gamma + \frac{1}{2} I_{yy}^f W_{xtt} \right) \quad (85)$$

$$m_z^i = -m g e_{cg} \sin \Gamma + \delta(x) \left( \frac{1}{2} m_f g y_f \sin \Gamma - \frac{1}{2} I_{zz}^f V_{xtt} \right) \quad (86)$$

where  $e_m$  is the offset of the mid-chord aft of the elastic axis;  $e_{cg}$  is the offset of the center of mass forward of the elastic axis;  $(0, y_f, z_f)$  is the coordinate of the aircraft CG such that  $y_f > 0$  aft of elastic axis and  $z_f > 0$  above the elastic axis;  $m_f$  is the mass of the fuselage and tail empennage; and  $I_{xx}^f$ ,  $I_{yy}^f$ , and  $I_{zz}^f$  are the inertias about the left wing reference frame of the fuselage and tail empennage.

The aircraft CG is located in the aircraft body-fixed reference frame B at a distance  $x_B$  aft of the elastic axis and  $z_B$

above the elastic axis at  $x = 0$ . Then by coordinate transformation into the reference frame D

$$y_f = x_B \cos \Lambda \quad (87)$$

$$z_f = -x_B \sin \Lambda \sin \Gamma + z_B \cos \Gamma \quad (88)$$

The roll, pitch, yaw, and roll-yaw inertias of the fuselage and the tail empennage at the elastic axis at  $x = 0$  are computed as

$$I_{XX}^f = \bar{I}_{XX}^f \quad (89)$$

$$I_{YY}^f = \bar{I}_{YY}^f + m_f (x_B^2 + z_B^2) \quad (90)$$

$$I_{ZZ}^f = \bar{I}_{ZZ}^f + m_f x_B^2 \quad (91)$$

$$I_{XZ}^f = \bar{I}_{XZ}^f + m_f x_B z_B \quad (92)$$

where  $\bar{I}_{XX}^f$ ,  $\bar{I}_{YY}^f$ ,  $\bar{I}_{ZZ}^f$ , and  $\bar{I}_{XZ}^f$  are the roll, pitch, yaw, and roll-yaw inertias at the aircraft CG.

The inertias can be transformed into the wing reference frame D as

$$\begin{aligned} I_{xx}^f &= \begin{bmatrix} -\sin \Lambda \cos \Gamma & -\cos \Lambda \cos \Gamma & -\sin \Gamma \end{bmatrix} \begin{bmatrix} I_{XX}^f & 0 & -I_{XZ}^f \\ 0 & I_{YY}^f & 0 \\ -I_{XZ}^f & 0 & I_{ZZ}^f \end{bmatrix} \begin{bmatrix} -\sin \Lambda \cos \Gamma \\ -\cos \Lambda \cos \Gamma \\ -\sin \Gamma \end{bmatrix} \\ &= I_{XX}^f \sin^2 \Lambda \cos^2 \Gamma + I_{YY}^f \cos^2 \Lambda \cos^2 \Gamma + I_{ZZ}^f \sin^2 \Gamma - I_{XZ}^f \sin \Lambda \sin 2\Gamma \end{aligned} \quad (93)$$

$$I_{yy}^f = \begin{bmatrix} -\cos \Lambda & \sin \Lambda & 0 \end{bmatrix} \begin{bmatrix} I_{XX}^f & 0 & -I_{XZ}^f \\ 0 & I_{YY}^f & 0 \\ -I_{XZ}^f & 0 & I_{ZZ}^f \end{bmatrix} \begin{bmatrix} -\cos \Lambda \\ \sin \Lambda \\ 0 \end{bmatrix} = I_{XX}^f \cos^2 \Lambda + I_{YY}^f \sin^2 \Lambda \quad (94)$$

$$\begin{aligned} I_{zz}^f &= \begin{bmatrix} \sin \Lambda \sin \Gamma & \cos \Lambda \sin \Gamma & -\cos \Gamma \end{bmatrix} \begin{bmatrix} I_{XX}^f & 0 & -I_{XZ}^f \\ 0 & I_{YY}^f & 0 \\ -I_{XZ}^f & 0 & I_{ZZ}^f \end{bmatrix} \begin{bmatrix} \sin \Lambda \sin \Gamma \\ \cos \Lambda \sin \Gamma \\ -\cos \Gamma \end{bmatrix} \\ &= I_{XX}^f \sin^2 \Lambda \sin^2 \Gamma + I_{YY}^f \cos^2 \Lambda \sin^2 \Gamma + I_{ZZ}^f \cos^2 \Gamma + I_{XZ}^f \sin \Lambda \sin 2\Gamma \end{aligned} \quad (95)$$

## V. Aeroelastic Equations for Combined Extensional, Bending, and Torsion Motion

### A. Equilibrium Conditions

Consider the equilibrium of forces and moments acting on a wing section. The resulting equilibrium equations are given by<sup>6</sup>

$$\frac{\partial P_x}{\partial x} + f_x = 0 \quad (96)$$

$$\frac{\partial P_y}{\partial x} + f_y = 0 \quad (97)$$

$$\frac{\partial P_z}{\partial x} + f_z = 0 \quad (98)$$

$$\frac{\partial M_x}{\partial x} - P_y W_x + P_z V_x + m_x = 0 \quad (99)$$

$$\frac{\partial M_y}{\partial x} - P_x W_x + P_z - m_y = 0 \quad (100)$$

$$\frac{\partial M_z}{\partial x} - P_x V_x + P_y + m_z = 0 \quad (101)$$

These equations become

$$\frac{\partial M_x}{\partial x} + \left( \frac{\partial M_z}{\partial x} - P_x V_x + m_z \right) W_x - \left( \frac{\partial M_y}{\partial x} - P_x W_x - m_y \right) V_x + m_x = 0 \quad (102)$$

$$\frac{\partial^2 M_y}{\partial x^2} - \frac{\partial (P_x W_x)}{\partial x} - \frac{\partial m_y}{\partial x} - f_z = 0 \quad (103)$$

$$\frac{\partial^2 M_z}{\partial x^2} - \frac{\partial (P_x V_x)}{\partial x} + \frac{\partial m_z}{\partial x} - f_y = 0 \quad (104)$$

Let  $U(x, t) = \bar{U}(x) + u(x, t)$ ,  $\Theta(x, t) = \bar{\Theta}(x) + \theta(x, t)$ ,  $W(x, t) = \bar{W}(x) + w(x, t)$ , and  $V(x, t) = \bar{V}(x) + v(x, t)$  where the bar symbols denote the static deflection. The static tensile force acting on the wing section is defined as  $T = EA\bar{U}_x$ . Then, considering the linear deflection terms, we have

$$\begin{bmatrix} P_x \\ M_x \\ M_y \\ M_z \end{bmatrix} = \begin{bmatrix} EAu_x - EAe_y v_{xx} \\ (GJ + Tk^2) \theta_x \\ EI_{yy} w_{xx} - EI_{yz} v_{xx} \\ -EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx} \end{bmatrix} \quad (105)$$

$$\frac{\partial (EAu_x - EAe_y v_{xx})}{\partial x} = -f_x \quad (106)$$

$$\begin{aligned} \frac{\partial [(GJ + Tk^2) \theta_x]}{\partial x} &= -m_x - \left( \frac{d\bar{M}_z}{dx} - \bar{P}_x \bar{V}_x + \bar{m}_z \right) w_x - \left( \frac{\partial \bar{M}_z}{\partial x} - \bar{P}_x v_x - P_x \bar{V}_x + m_z \right) \bar{W}_x \\ &+ \left( \frac{d\bar{M}_y}{dx} - \bar{P}_x \bar{W}_x - \bar{m}_y \right) v_x + \left( \frac{\partial \bar{M}_y}{\partial x} - \bar{P}_x w_x - P_x \bar{W}_x - m_y \right) P \bar{V}_x \end{aligned} \quad (107)$$

$$\frac{\partial^2 (EI_{yy} w_{xx} - EI_{yz} v_{xx})}{\partial x^2} = f_z + \frac{\partial m_y}{\partial x} + \frac{\partial (\bar{P}_x w_x + P_x \bar{W}_x)}{\partial x} \quad (108)$$

$$\frac{\partial^2 (-EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx})}{\partial x^2} = f_y - \frac{\partial m_z}{\partial x} + \frac{\partial (\bar{P}_x v_x + P_x \bar{V}_x)}{\partial x} \quad (109)$$

where  $I_{xx} = Ak^2$ .

Since  $\bar{V}_x$  and  $\bar{V}_{xx}$  are generally small, so they can be neglected. Then,  $\bar{P}_x = T$  so that the equations above become

$$\begin{aligned} \frac{\partial [(GJ + Tk^2) \theta_x]}{\partial x} &= -m_x - \left[ \frac{d(-Te_y - EI_{yz} \bar{W}_{xx})}{dx} + \bar{m}_z \right] w_x \\ &- \left[ \frac{\partial (-EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx})}{\partial x} - T v_x + m_z \right] \bar{W}_x + \left[ \frac{d(EI_{yy} \bar{W}_{xx})}{dx} - T \bar{W}_x - \bar{m}_y \right] v_x \end{aligned} \quad (110)$$

$$\frac{\partial^2 (EI_{yy} w_{xx} - EI_{yz} v_{xx})}{\partial x^2} = f_z + \frac{\partial m_y}{\partial x} + \frac{\partial [T w_x + (EAu_x - EAe_y v_{xx}) \bar{W}_x]}{\partial x} \quad (111)$$

$$\frac{\partial^2 (-EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx})}{\partial x^2} = f_y - \frac{\partial m_z}{\partial x} + \frac{\partial (T v_x)}{\partial x} \quad (112)$$

Finally, if the vertical bending deflection is not too large so that the nonlinear effect is not dominant, then in the simplest form these equations are expressed as

$$\frac{\partial [(GJ + Tk^2) \theta_x]}{\partial x} = -m_x \quad (113)$$

$$\frac{\partial^2 (EI_{yy} w_{xx} - EI_{yz} v_{xx})}{\partial x^2} = f_z + \frac{\partial m_y}{\partial x} + \frac{\partial (T w_x)}{\partial x} \quad (114)$$

$$\frac{\partial^2 (-EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx})}{\partial x^2} = f_y - \frac{\partial m_z}{\partial x} + \frac{\partial (T v_x)}{\partial x} \quad (115)$$

The aeroelastic equations of motion for axially loaded structures with no deflection-dependent terms are given by

$$\frac{\partial (EAu_x - EAe_y v_{xx})}{\partial x} = \left( m + \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda \right) u_{tt} - \delta(x - x_e) (F_z^e w_x + F_y^e v_x) \quad (116)$$

$$\begin{aligned} \frac{\partial [(GJ + Tk^2) \theta_x]}{\partial x} &= c_m q_\infty c^2 \cos^2 \Lambda \cos \Gamma + \left( mr_k^2 + \frac{\pi}{4} \rho_\infty c^2 e_m^2 \cos \Lambda + \frac{\pi}{128} \rho_\infty c^4 \cos \Lambda \right) \theta_{tt} \\ &- \left( me_{cg} + \frac{\pi}{4} \rho_\infty c^2 e_m \cos \Lambda \right) w_{tt} - \delta(x - x_e) [(F_x y_e + F_y x_e) w_x - (F_x z_e + F_z x_e) v_x] \\ &+ \delta(x - x_e) (m_e r_e^2 \theta_{tt} - m_e y_e w_{tt} + m_e z_e v_{tt}) + \delta(x) \left( \frac{1}{2} m_f y_f w_{tt} + \frac{1}{2} I_{xx}^f \theta_{tt} - \frac{1}{2} m_f z_f v_{tt} \right) \end{aligned} \quad (117)$$

$$\begin{aligned} \frac{\partial^2 (EI_{yy} w_{xx} - EI_{yz} v_{xx})}{\partial x^2} &= c_L q_\infty c \cos \Lambda + \frac{\partial}{\partial x} (c_m q_\infty c^2 \sin \Lambda \cos \Lambda) + \left( me_{cg} + \frac{\pi}{4} \rho_\infty c^2 e_m \cos \Lambda \right) \theta_{tt} \\ &- \left( m + \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda \right) w_{tt} + \delta(x - x_e) (-F_y^e \theta - F_x^e w_x) \\ &+ \frac{\partial}{\partial x} \{ \delta(x - x_e) [(F_x y_e + F_y x_e) \theta - (F_y z_e - F_z y_e) v_x] \} + \delta(x - x_e) (m_e y_e \theta_{tt} - m_e w_{tt}) \\ &+ \delta(x) \left( -\frac{1}{2} m_f y_f \theta_{tt} - \frac{1}{2} m_f w_{tt} \right) + \frac{\partial}{\partial x} \left[ \delta(x) \left( \frac{1}{2} I_{yy}^f w_{xtt} \right) \right] + \frac{\partial (T w_x)}{\partial x} \end{aligned} \quad (118)$$

$$\begin{aligned} \frac{\partial^2 (-EAe_y u_x - EI_{yz} w_{xx} + EI_{zz} v_{xx})}{\partial x^2} &= c_a q_\infty c \cos^2 \Lambda - \frac{\partial}{\partial x} (c_m q_\infty c^2 \cos^2 \Lambda \sin \Gamma) - \left( m + \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda \right) v_{tt} \\ &+ \delta(x - x_e) (F_z^e \theta - F_x^e v_x) - \frac{\partial}{\partial x} \{ \delta(x - x_e) [(F_x z_e + F_z x_e) \theta - (F_y z_e - F_z y_e) w_x] \} \\ &+ \delta(x - x_e) (-m_e z_e \theta_{tt} - m_e v_{tt}) + \delta(x) \left( \frac{1}{2} m_f z_f \theta_{tt} - \frac{1}{2} m_f v_{tt} \right) - \frac{\partial}{\partial x} \left[ \delta(x) \left( -\frac{1}{2} I_{zz}^f v_{xtt} \right) \right] + \frac{\partial (T v_x)}{\partial x} \end{aligned} \quad (119)$$

## B. Discretization

The partial differential equations can be solved by weak-form solutions such as the finite-element method (FEM). Using interpolation functions, the displacements can be expressed as

$$u(x,t) = \sum_{i=1}^n N_u(x) u_i(t) \quad (120)$$

$$\theta(x,t) = \sum_{i=1}^n N_\theta(x) \theta_i(t) \quad (121)$$

$$w(x,t) = \sum_{i=1}^n N_w(x) \phi_i(t) \quad (122)$$

$$v(x,t) = \sum_{i=1}^n N_v(x) \psi_i(t) \quad (123)$$

where  $N_u(x)$ ,  $N_\theta(x)$ ,  $N_w(x)$ , and  $N_v(x)$  are the interpolation functions,  $\phi_i = \begin{bmatrix} w_i & w_{x,i} \end{bmatrix}^\top$ ,  $\psi_i = \begin{bmatrix} v_i & v_{x,i} \end{bmatrix}^\top$ , and  $n$  is the number of elements.

The discretization results in a spring-mass-damper system in the form

$$\sum_{i=1}^n M_i \ddot{x}_i + \sum_{i=1}^n C_i \dot{x}_i + \sum_{i=1}^n K_i x_i = 0 \quad (124)$$

where  $M_i$  is the elemental mass matrix,  $C_i$  is the elemental damping matrix,  $K_i$  is the elemental stiffness matrix, and  $x_i = \begin{bmatrix} u_i & \theta_i & w_i & w_{x,i} & v_i & v_{x,i} \end{bmatrix}^\top$  is the nodal displacement vector.

The mass matrix due to the structural mass, aerodynamic mass, engine mass, and fuselage mass is given by

$$M_i = M_i^s + M_i^a + M_i^e + M_i^f \quad (125)$$



where

$$M_i^s = \int_0^{l_i} m \begin{bmatrix} N_u^\top N_u & 0 & 0 & 0 \\ 0 & N_\theta^\top r_k^2 N_\theta & -N_\theta^\top e_{cg} N_w & 0 \\ 0 & -N_w^\top e_{cg} N_\theta & N_w^\top N_w & 0 \\ 0 & 0 & 0 & N_v^\top N_v \end{bmatrix} dx \quad (126)$$

$$M_i^a = \int_0^{l_i} \frac{\pi}{4} \rho_\infty c^2 \cos \Lambda \begin{bmatrix} N_u^\top N_u & 0 & 0 & 0 \\ 0 & N_\theta^\top \left( e_m^2 + \frac{c^2}{32} \right) N_\theta & -N_\theta^\top e_m N_w & 0 \\ 0 & -N_w^\top e_m N_\theta & N_w^\top N_w & 0 \\ 0 & 0 & 0 & N_v^\top N_v \end{bmatrix} dx$$

$$+ \int_0^{l_i} q_\infty c \cos^2 \Lambda \left\{ \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial \theta_{it}} N_\theta & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial w_{it}} N_w & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial v_{it}} N_v \\ 0 & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial \theta_{it}} N_\theta & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial w_{it}} N_w & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial v_{it}} N_v \\ 0 & -N_v^\top \frac{\partial c_a}{\partial \theta_{it}} N_\theta & -N_v^\top \frac{\partial c_a}{\partial w_{it}} N_w & -N_v^\top \frac{\partial c_a}{\partial v_{it}} N_v \end{bmatrix} \right.$$

$$\left. + \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial \theta_{it}} N_\theta & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial w_{it}} N_w & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial v_{it}} N_v \\ 0 & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial \theta_{it}} N_\theta & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial w_{it}} N_w & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial v_{it}} N_v \end{bmatrix} \right\} dx \quad (127)$$

$$M_i^e = m_e \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & N_\theta^\top r_e^2 N_\theta & -N_\theta^\top y_e N_w & N_\theta^\top z_e N_v \\ 0 & -N_w^\top y_e N_\theta & N_w^\top N_w & 0 \\ 0 & N_v^\top z_e N_\theta & 0 & N_v^\top N_v \end{bmatrix}_{x=x_e} \quad (128)$$

$$M_i^f = \frac{1}{2} \begin{bmatrix} 0 & 0 & 0 & 0 \\ 0 & N_\theta^\top I_{xx}^f N_\theta \ddot{\Theta}_i & N_\theta^\top m_f y_f N_w & -N_\theta^\top m_f z_f N_v \\ 0 & N_w^\top m_f y_f N_\theta & N_w^\top m_f N_w + N_w'^\top I_{yy}^f N_w' & 0 \\ 0 & -N_v^\top m_f z_f N_\theta & 0 & N_v^\top m_f N_v + N_v'^\top I_{zz}^f N_v' \end{bmatrix}_{x=0} \quad (129)$$

The damping matrix due to structural damping and aerodynamic damping is given by

$$C_i = C_i^s + C_i^a \quad (130)$$

where

$$C_i^s = 2M_i^s \Phi \zeta \Omega \Phi^{-1} \quad (131)$$

$$C_i^a = \int_0^{l_i} q_\infty c \cos^2 \Lambda \left\{ \begin{array}{l} \left[ \begin{array}{cccc} 0 & 0 & 0 & 0 \\ 0 & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial \theta} N_\theta & N_\theta^\top c \cos \Gamma \left( \frac{\partial c_m}{\partial w_t} N_w + \frac{\partial c_m}{\partial w_{xt}} N'_w \right) & N_\theta^\top c \cos \Gamma \left( \frac{\partial c_m}{\partial v_t} N_v + \frac{\partial c_m}{\partial v_{xt}} N'_v \right) \\ 0 & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial \theta} N_\theta & -N_w^\top \frac{1}{\cos \Lambda} \left( \frac{\partial c_L}{\partial w_t} N_w + \frac{\partial c_L}{\partial w_{xt}} N'_w \right) & -N_w^\top \frac{1}{\cos \Lambda} \left( \frac{\partial c_L}{\partial v_t} N_v + \frac{\partial c_L}{\partial v_{xt}} N'_v \right) \\ 0 & -N_v^\top \frac{\partial c_a}{\partial \theta} N_\theta & -N_v^\top \left( \frac{\partial c_a}{\partial w_t} N_w + \frac{\partial c_a}{\partial w_{xt}} N'_w \right) & -N_v^\top \left( \frac{\partial c_a}{\partial v_t} N_v + \frac{\partial c_a}{\partial v_{xt}} N'_v \right) \end{array} \right] dx \\ + \left[ \begin{array}{cccc} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial \theta} N_\theta & N_w'^\top c \tan \Lambda \left( \frac{\partial c_m}{\partial w_t} N_w + \frac{\partial c_m}{\partial w_{xt}} N'_w \right) & N_w'^\top c \tan \Lambda \left( \frac{\partial c_m}{\partial v_t} N_v + \frac{\partial c_m}{\partial v_{xt}} N'_v \right) \\ 0 & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial \theta} N_\theta & -N_v'^\top c \sin \Gamma \left( \frac{\partial c_m}{\partial w_t} N_w + \frac{\partial c_m}{\partial w_{xt}} N'_w \right) & -N_v'^\top c \sin \Gamma \left( \frac{\partial c_m}{\partial v_t} N_v + \frac{\partial c_m}{\partial v_{xt}} N'_v \right) \end{array} \right] dx \end{array} \right. \quad (132)$$

where  $\zeta = \text{diag}(\zeta_1, \zeta_2, \dots, \zeta_n)$  is the diagonal viscous damping ratio matrix,  $\Phi$  is the structural dynamic eigenvector matrix, and  $\Omega = \text{diag}(\omega_1, \omega_2, \dots, \omega_n)$  is a diagonal matrix of the structural frequencies.

The stiffness matrix due to structural stiffness including the nonlinear tension-induced stiffness, aerodynamic stiffness, and engine thrust-induced stiffness is given by

$$K_i = K_i^s + K_i^a + K_i^e \quad (133)$$

where

$$K_i^s = \int_0^{l_i} \left[ \begin{array}{cccc} N_u'^\top E A N_u' & 0 & 0 & -N_u'^\top E A e_y N_v'' \\ 0 & N_\theta'^\top (GJ + Tk^2) N_\theta' & 0 & 0 \\ 0 & 0 & N_w''^\top E I_{yy} N_w'' + N_w'^\top T N_w' & -N_w''^\top E I_{yz} N_v'' \\ -N_v''^\top E A e_y N_u' & 0 & -N_v''^\top E I_{yz} N_w'' & N_v''^\top E I_{zz} N_v'' + N_v'^\top T N_v' \end{array} \right] dx \quad (134)$$

$$K_i^a = \int_0^{l_i} q_\infty c \cos^2 \Lambda \left\{ \begin{array}{l} \left[ \begin{array}{cccc} 0 & 0 & 0 & 0 \\ 0 & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial \theta} N_\theta & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial w_x} N'_w & N_\theta^\top c \cos \Gamma \frac{\partial c_m}{\partial v_x} N'_v \\ 0 & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial \theta} N_\theta & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial w_x} N'_w & -N_w^\top \frac{1}{\cos \Lambda} \frac{\partial c_L}{\partial v_x} N'_v \\ 0 & -N_v^\top \frac{\partial c_a}{\partial \theta} N_\theta & -N_v^\top \frac{\partial c_a}{\partial w_x} N'_w & -N_v^\top \frac{\partial c_a}{\partial v_x} N'_v \end{array} \right] dx \\ + \left[ \begin{array}{cccc} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial \theta} N_\theta & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial w_x} N'_w & N_w'^\top c \tan \Lambda \frac{\partial c_m}{\partial v_x} N'_v \\ 0 & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial \theta} N_\theta & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial w_x} N'_w & -N_v'^\top c \sin \Gamma \frac{\partial c_m}{\partial v_x} N'_v \end{array} \right] dx \end{array} \right. \quad (135)$$

$$K_i^e = \left[ \begin{array}{cccc} 0 & 0 & -N_u^\top F_z^e N_w' & -N_u^\top F_y^e N_v' \\ 0 & 0 & -N_\theta^\top (F_x y_e + F_y x_e) N_w' & N_\theta^\top (F_x z_e + F_z x_e) N_v' \\ 0 & N_w^\top F_y^e N_\theta + N_w'^\top (F_x y_e + F_y x_e) N_\theta & N_w^\top F_x^e N_w' & -N_w'^\top (F_y z_e - F_z y_e) N_v' \\ 0 & -N_v^\top F_z^e N_\theta - N_v'^\top (F_x z_e + F_z x_e) N_\theta & N_v^\top (F_y z_e - F_z y_e) N_w' & N_v^\top F_x^e N_v' \end{array} \right] \quad (136)$$

The effect of the tensile force can be clearly seen in the structural stiffness as the additive stiffness terms in both the

bending and torsion stiffness matrices. The presence of a tensile force causes a geometric nonlinear tension stiffening effect which increases the natural frequencies of the structural dynamic modes. Conversely, a compressive force causes the natural frequencies to decrease. Static instability due to compressive loading could result when the compressive force-induced stiffness is large enough to offset the torsion or bending stiffness.

## VI. Applications

### A. Rotary Wings

Rotary wings are axially loaded structures that experience tensile loading due to the centrifugal acceleration. The structural stiffness of a rotary wing is given by

$$\begin{aligned}
K^s = \int_0^{l_i} & \begin{bmatrix} N'_\theta{}^\top \left[ GJ + EB_1 (\gamma')^2 \right] N'_\theta & -N'_\theta{}^\top EB_2 \gamma' N''_w & -N'_\theta{}^\top EB_3 \gamma' N''_v \\ -N''_w{}^\top EB_2 \gamma' N'_\theta & N''_w{}^\top EI_{yy} N''_w & -N''_w{}^\top EI_{yz} N''_v \\ -N''_v{}^\top EB_3 \gamma' N'_\theta & -N''_v{}^\top EI_{yz} N''_w & N''_v{}^\top EI_{zz} N''_v \end{bmatrix} dx \\
+ \Omega^2 \int_0^{l_i} & \begin{bmatrix} N'_\theta{}^\top T k^2 N'_\theta & 0 & 0 \\ 0 & N'_w{}^\top T N'_w & 0 \\ 0 & 0 & N'_v{}^\top T N'_v \end{bmatrix} dx
\end{aligned} \tag{137}$$

where  $T(x) = \int_x^L mr\Omega^2 dr$  is the centrifugal force,  $r$  is the radius, and  $\Omega$  is the angular speed.

Therefore, the system matrix equations at the angular speed  $\Omega$  and zero speed can be expressed as

$$-\omega_n^2 x + (M^{-1}K_0 + \Omega^2 M^{-1}F)x = 0 \tag{138}$$

and

$$-\omega_0^2 x + M^{-1}K_0 x = 0 \tag{139}$$

Thus, the natural frequencies of a rotary wing are dependent on the angular speed of rotation according to

$$\omega_n^2 = \omega_0^2 + (f\Omega)^2 \tag{140}$$

where  $\omega_n$  is the natural frequency at the angular speed  $\Omega$  and  $\omega_0$  is the zero-speed natural frequency, and  $f$  is the centrifugal stiffening factor.

As an example, the natural frequencies of compressor rotor blades in the NASA Ames 11-Ft Unitary Plan 11-Foot By 11-Foot Transonic Wind Tunnel three-stage compressor<sup>9</sup> are computed. The experimental natural frequencies were obtained from modal analysis measurements. In addition, wind-on strain gauge measurements were conducted at various compressor speeds to obtain the variation in the natural frequencies with the compressor speed. The natural frequencies were computed from the Fast Fourier Transform (FFT) of the time signatures of the strain gauges at various compressor speeds. This allows the centrifugal stiffening factor to be computed. Table 1 shows the computed natural frequencies of the three-stage compressor rotor blades by the finite-element method presented herein as compared to the experimental results. The agreement between the computed natural frequencies and the experimental results is excellent.

Mode	FEM $\omega_0$ (Hz)	Bench Test $\omega_0$ (Hz)	FEM $\omega_n$ at 700 rpm (Hz)	FEM $f$	Wind-On $f$
1B	54.43	54.18	61.67	2.488	2.375
2B	167.32	160.92	174.51	4.252	4.842
1T	302.37	305.56	309.25	5.560	4.661
3B	370.29	374.61	375.25	5.24	9.110
2T	549.98	554.51	561.91	9.87	8.764

Table 1 - Natural Frequencies of Rotor Blades in NASA Ames Unitary Plan 11-Foot By 11-Foot Transonic Wind Tunnel Three-Stage Compressor

### B. Truss-Braced Wing Aircraft

A recent flutter wind tunnel test of a dynamically scaled TBW was conducted in NASA Langley Transonic Dynamic Wind Tunnel (TDT).<sup>3</sup> The model is 16% scaled, semi-span, wall-mounted model as shown in Fig. 3. It was reported that the flutter speed is dependent on the angle of attack at  $\pm 1$  deg.<sup>3</sup> While there may be many nonlinear factors that could contribute to this observation, one possible exploration could be the geometric nonlinear tension stiffening effect of the main struts at a positive angle of attack, and conversely the softening effect at a negative angle of attack, assuming the axial force is below the critical buckling load.

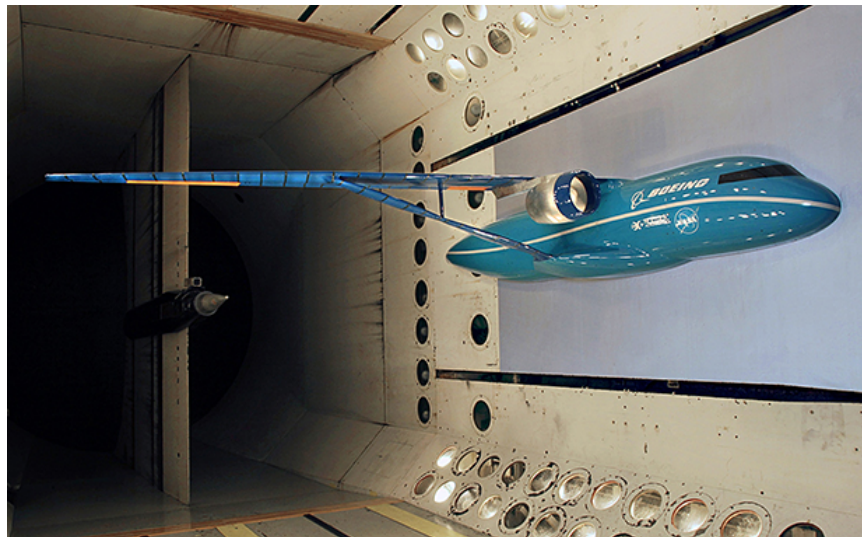


Fig. 3 - Truss-Braced Wing (TBW) Model in NASA Langley Transonic Dynamic Tunnel (TDT)

A static aeroelastic analysis is performed to compute the static deflections and the tension carried by the struts. The finite-element model (FEM) of the TBW utilizes an equivalent beam/frame configuration that uses a total of five components. Only the right side of the aircraft is modeled, and the wing configuration is assumed to be cantilevered to the side of the fuselage. The three main physical components of the wing configuration are the wing, the strut, and the jury strut. The FEM uses 30 elements for the wing, 28 for the strut, and 5 for the jury strut. Two rigid elements are used to connect the strut and the jury strut to the wing. Figure 4 shows the FEM of the TBW.

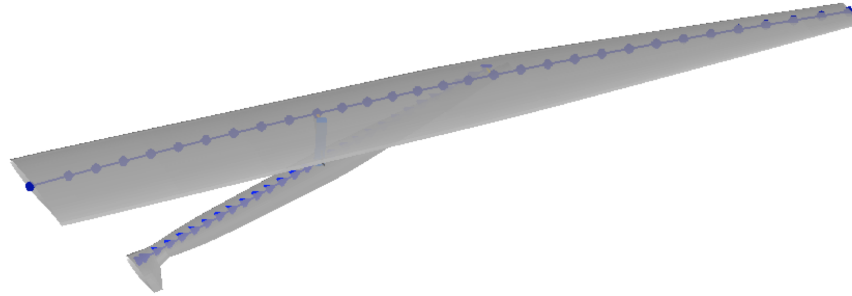


Fig. 4 - Finite-Element Model of Truss-Braced Wing (TBW)

The FEM uses lumped mass and inertia properties provided by Boeing Research & Technology.<sup>10</sup> The lumped mass data is converted into the running mass. Extensional stiffness of the wing is not provided and instead is estimated. The engine-pylon mass is not available and therefore is not modeled. This initial model is aimed at understanding the geometric nonlinear effect of the tension stiffening, and is not intended to provide an accurate representation of the TBW model. Therefore, the results are used only for the purpose of illustrating the geometric nonlinear tension stiffening, but are not to be compared to the established TBW data.

Aerodynamic loads are computed by an aerodynamic analysis code. Figures 5 and 6 are the plots of lift and pitching moment distributions for the wing and the strut at the design cruise  $C_L = 0.81$  at Mach 0.7 and an altitude of 42,000 ft, corresponding to the 1g flight load or a cruise gross weight of 170,250 lbs.

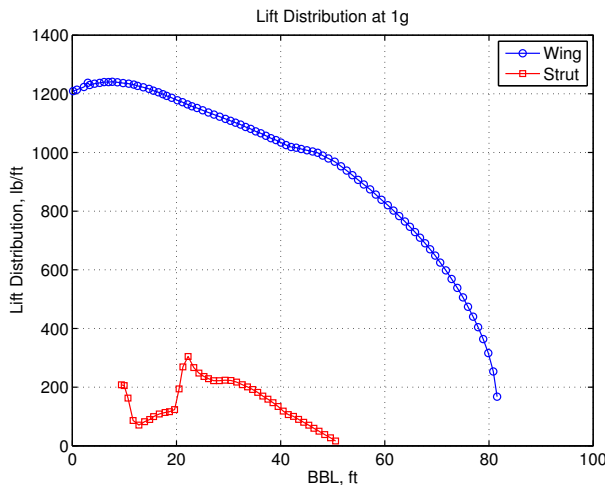


Fig. 5 - Lift Distribution

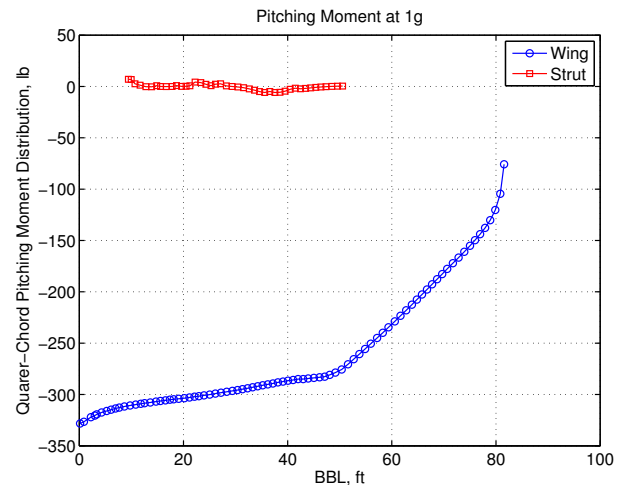


Fig. 6 - Pitching Moment Distribution

A nonlinear static analysis is performed to include the geometric nonlinear effect of tension stiffening. The solution is initialized with zero initial tension. After the first iteration, the tension is computed and the stiffness matrix is updated. The solution is iterated until the wing tip solution converges. The wing tip deflection from the linear solution with no tension stiffening is 1.5434 ft which indicates that the combined wing-strut structure is quite stiff. The nonlinear solution predicts a wing tip deflection of 1.5248 ft, which only reduces by about 1%. So the tension stiffening effect does not appear to be significant at 1g load. For comparison, the nonlinear wing tip deflections at 2.5g and -1g are computed to be 3.0622 ft and -1.8885 ft, respectively. These are compared to the corresponding linear wing tip deflections of 3.1009 ft and -1.9194 ft. The computed wing static deflections are shown in Fig. 7.

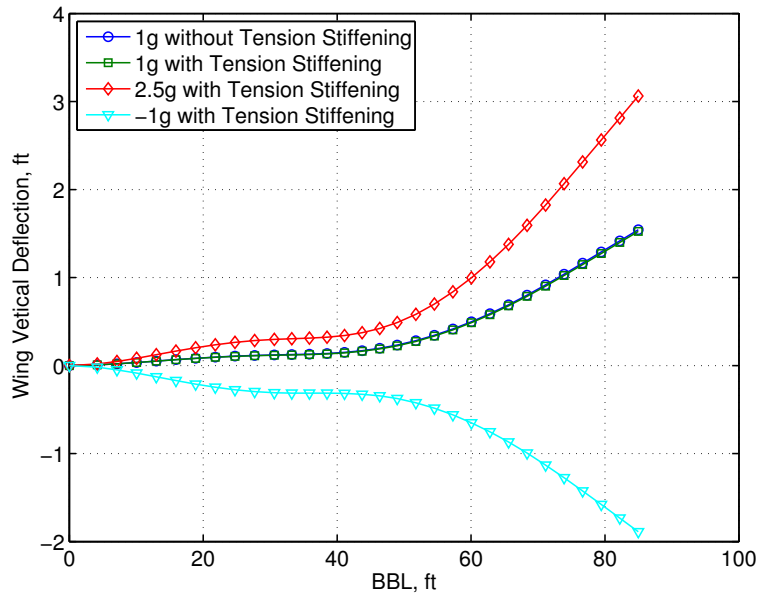


Fig. 7 - Wing Static Deflection

The axial extension of the strut for 1g, 2.5g, and -1g flight loads is plotted in Fig. 8. The tension carried by the strut is then computed from the axial strain and is plotted in Fig. 9. The static aeroelastic analysis shows that the strut on the average carries about 153,000 lbs to support a take-off gross weight of 170,250 lbs.

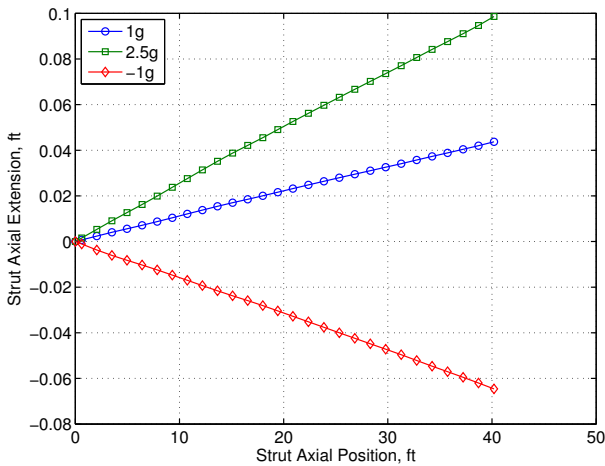


Fig. 8 - Axial Extension of Strut

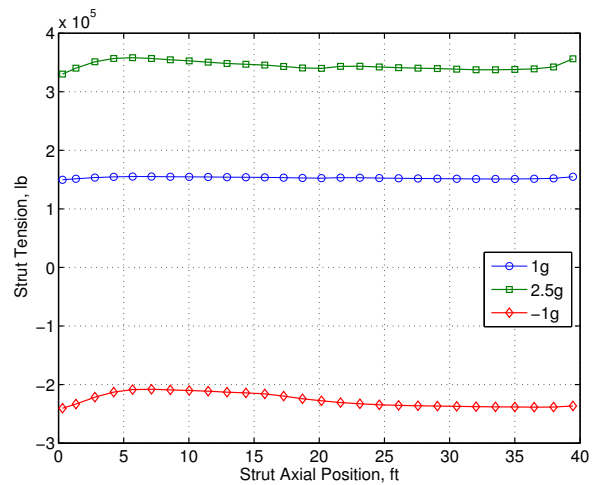


Fig. 9 - Tension in Strut

An eigenvalue analysis is performed to compute the no-load structural dynamic natural frequencies of the TBW and the structural dynamic natural frequencies with tension stiffening due to the loaded structure at 1g, 2.5g, and -1g loads. The natural frequencies of the first ten modes are shown in Table 2. It can be seen that the natural frequencies with tension stiffening at 1g and 2.5g are larger than the no-load natural frequencies. The increase in the natural

frequencies from the no-load natural frequencies ranges from about 1% to 8% for the 1g flight load case and from about 1% to 15% for the 2.5g flight load case. The -1g flight load case exhibits a significant reduction in the natural frequencies ranging from about 1% to 24%, particularly at higher modes.

Mode	No-Load Frequency (Hz)	Frequency at 1g (Hz)	Frequency at 2.5g (Hz)	Frequency at -1g (Hz)
1	1.9413	2.0052	2.0718	1.8062
2	2.3662	2.4541	2.5483	2.2090
3	4.0659	4.1832	4.3246	3.8891
4	6.6962	6.7525	6.8226	6.6082
5	8.5328	8.6784	8.8275	8.2343
6	10.0313	10.0461	10.0611	9.9881
7	11.4553	11.4804	11.5082	11.4106
8	17.1444	17.2365	17.3415	15.5022
9	18.2639	19.8070	21.3600	16.6709
10	22.1617	22.2387	22.5115	17.9237

Table 2 - Natural Frequencies of TBW Aircraft

A flutter analysis is performed for the no-load case and the three load cases. A 2% structural damping is assumed. The no-load case shows a flutter occurring at modes 4 and 5 as these two modes coalesce together, as shown in Fig. 10. The flutter speeds are 589.00 knots equivalent airspeed (KEAS) for mode 4 and 574.85 KEAS for mode 5, as shown in Fig. 11. Figures 12 to 17 are the plots of natural frequencies and damping of the TBW aircraft for the three flight loads at 1g, 2.5g, and -1g. The flutter speeds for the three flight load cases are shown in Table 3. As can be seen, the flutter speeds are dependent on the flight loads due to the geometric nonlinear effect of tension stiffening. The difference in the flutter speed is as much as 35 KEAS between the 1g flight load and -1g flight load and 15 KEAS between the 1g flight load and 2.5g flight load.

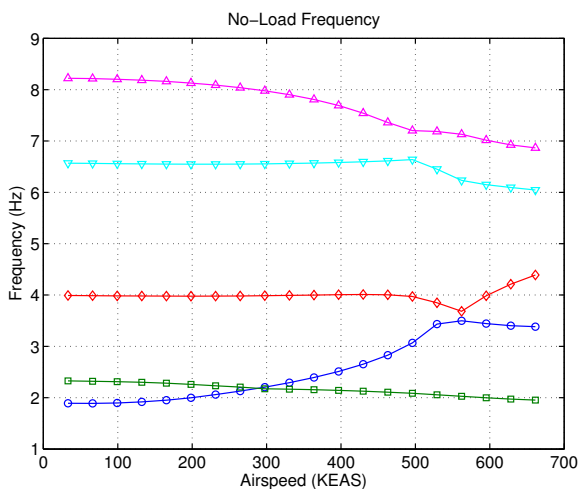


Fig. 10 - No-Load Frequency

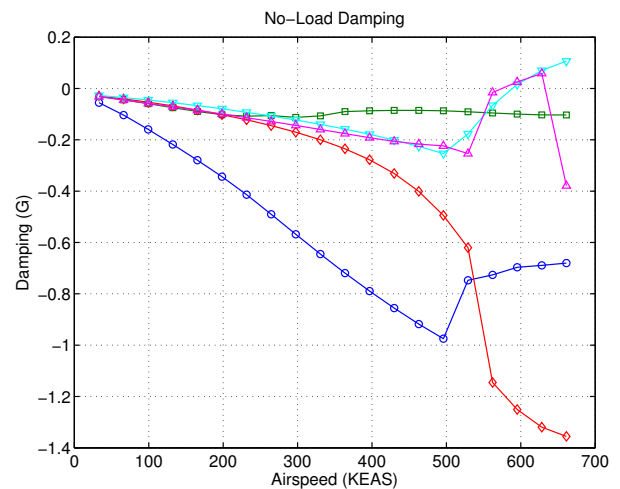


Fig. 11 - No-Load Damping

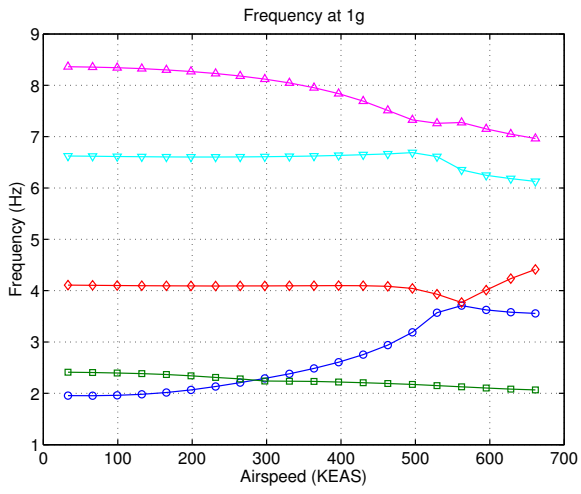


Fig. 142 - Frequency at 1g

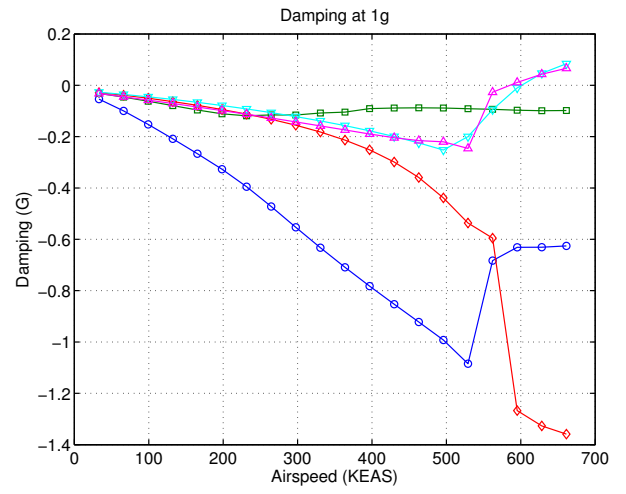


Fig. 13 - Damping at 1g

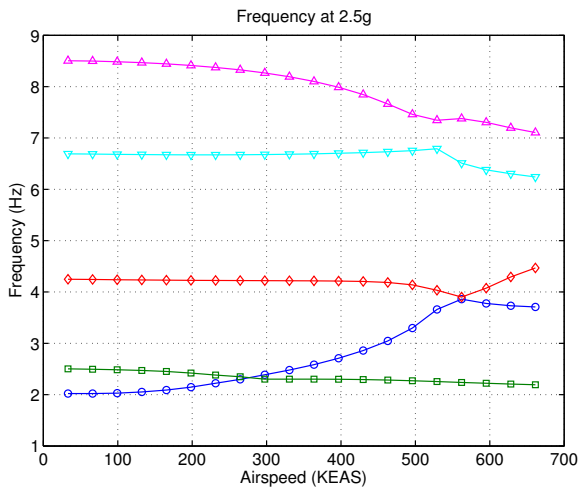


Fig. 14 - Frequency at 2.5g

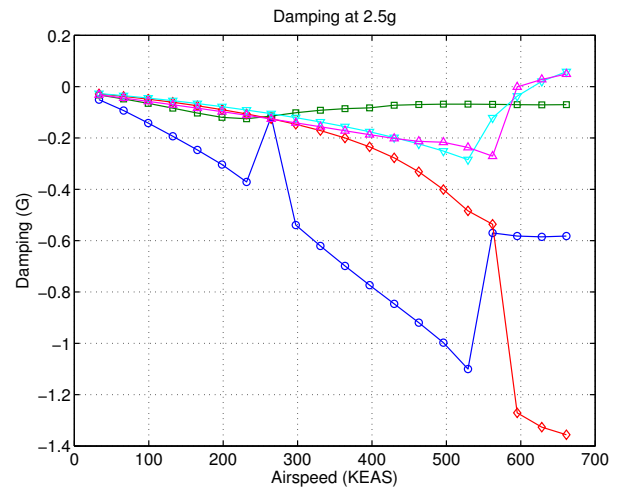


Fig. 15 - Damping at 2.5g



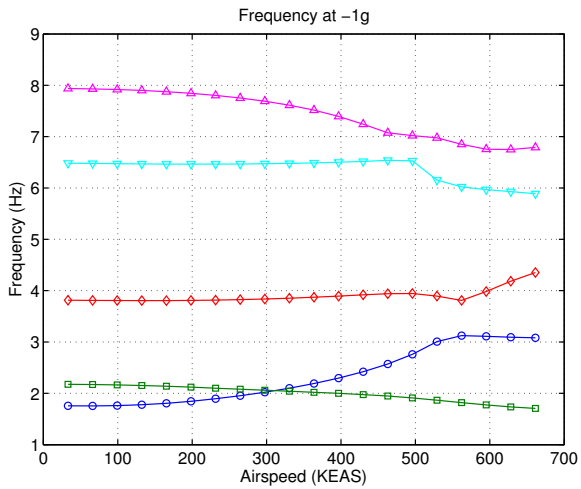


Fig. 16 - Frequency at -1g

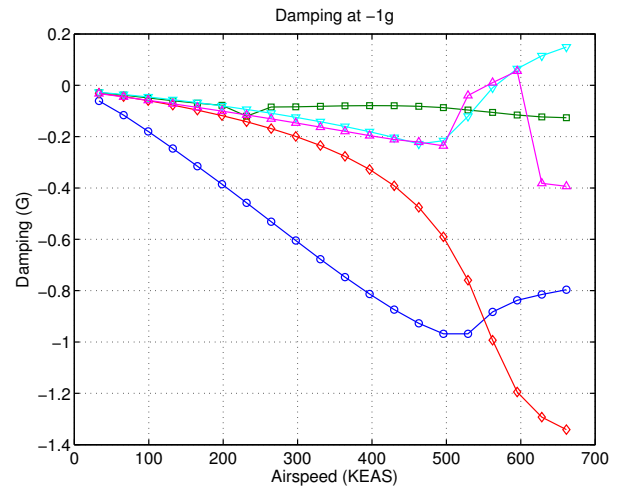


Fig. 17 - Damping at -1g

Mode	No-Load Flutter Speed (KEAS)	Flutter Speed at 1g (KEAS)	Flutter Speed at 2.5g (KEAS)	Flutter Speed at-1g (KEAS)
4	589.00	601.15	616.54	566.48
5	574.85	585.62	596.83	555.19

Table 3 - Flutter Speed of TBW Aircraft

The flutter Mach numbers at an altitude of 42,000 ft are shown in Table 4. The Boeing Report indicates a dive speed of Mach 0.82. A flutter margin of 15% results in a minimum clearance flutter Mach 0.9430. Thus, the flutter Mach numbers are well above the clearance flutter Mach. It should be noted that this analysis does not include the engine-pylon mass as well as the correct fuel mass. So the computed flutter speeds in this study are optimistic.

Mode	No-Load Flutter Mach	Flutter Mach at 1g	Flutter Mach at 2.5g	Flutter Mach at-1g
4	1.2867	1.3133	1.3469	1.2375
5	1.2558	1.2793	1.3038	1.2129

Table 4 - Flutter Mach Number of TBW Aircraft at Altitude of 42,000 ft

In general, increased wing loading results in an increase in the flutter speed. This is consistent with the observation from the wind tunnel test of the TBW in NASA Langley’s Transonic Dynamic Tunnel (TDT) which shows that the flutter speed at a negative angle of attack is consistently lower than that at a positive angle of attack at the same dynamic pressure.<sup>3,10</sup>

## VII. Conclusion

This paper presents a study of the geometric nonlinear effect of axial loading on the aeroelasticity of a truss-braced wing aircraft structure. The main struts which restrain the high-aspect ratio wings generally carry a tensile force due

to the lift force generated by the wings. The tension causes a geometric nonlinear stiffening effect which effectively raises the stiffness of the structure. This geometric nonlinear effect therefore causes the natural frequencies to increase, hence changes to flutter speed solutions. A nonlinear finite-element model of the trussed-braced wing aircraft has been developed. The nonlinear static deflection analysis is performed and yields a slightly lower wing deflection than the linear analysis. Eigenvalue analysis is performed and shows that the natural frequencies at the 1g and 2.5g flight loads are higher than the no-load natural frequencies. Conversely, the natural frequencies at the -1g flight load is lower than the no-load natural frequencies. A flutter analysis is performed and shows that two flutter modes exist. Modes 4 and 5 flutter at different airspeeds depending on the three flight loads. This is caused by the geometric nonlinear effect of tension stiffening due to the strut which affects the flutter frequencies. In general, increased wing loading results in an increase in the flutter speed. The study illustrates the importance of accounting for the geometric nonlinear tension stiffening effect in analyzing the truss-braced wing aircraft.

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