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# BAKER-AKHIEZER FUNCTION AS ITERATED RESIDUE AND SELBERG-TYPE INTEGRAL

GIOVANNI FELDER AND ALEXANDER P. VESELOV

ABSTRACT. A simple integral formula as an iterated residue is presented for the Baker-Akhiezer function related to  $A_n$  type root system both in the rational and trigonometric cases. We present also a formula for the Baker-Akhiezer function as a Selberg-type integral and generalise it to the deformed  $A_{n,1}$ -case. These formulas can be interpreted as new cases of explicit evaluation of Selberg-type integrals.

## 1. INTRODUCTION

The notion of rational Baker-Akhiezer (BA) function related to a configuration of hyperplanes with multiplicities was introduced in [1, 2, 3] as a multi-dimensional version of Krichever's axiomatic approach [4]. Such function exists only for special configurations, in particular for all Coxeter configurations. For the configuration of type  $A_{n-1}$  with multiplicity  $m$  it has the form

$$(1) \quad \Psi_m^{(n)} = \frac{P_m^{(n)}(x, \lambda)}{A(x)^m A(\lambda)^m} \exp(\lambda, x),$$

where  $A(x) = \prod_{i < j}^n (x_i - x_j)$ ,  $(\lambda, x) = \sum_{i=1}^n \lambda_i x_i$  and  $P_m^{(n)}(x, \lambda)$  is a polynomial in both  $x = (x_1, \dots, x_n) \in \mathbb{C}^n$  and  $\lambda = (\lambda_1, \dots, \lambda_n) \in \mathbb{C}^n$  with the leading term  $A(x)^m A(\lambda)^m$ . It satisfies the Schrödinger equation

$$(2) \quad L_m^{(n)} \Psi_m^{(n)} = -(\lambda, \lambda) \Psi_m^{(n)},$$

where  $L_m$  is the  $n$  particle *Calogero-Moser operator*

$$(3) \quad L_m^{(n)} = -\Delta + \sum_{i < j}^n \frac{2m(m+1)}{(x_i - x_j)^2}.$$

The rational BA function  $\Psi_m^{(n)}$  is determined uniquely by these properties and has a remarkable symmetry:

$$\Psi_m^{(n)}(z, \lambda) = \Psi_m^{(n)}(\lambda, z)$$

(see [3]). It plays an important role in the theories of commutative rings of differential operators, Huygens principle and quasi-invariants of Coxeter groups [1, 3, 5, 6].

Three ways for computing this function are known. The first one due to Chalykh and one of the authors [1] is known only in the first non-trivial case  $m = 1$  and is a recursive formula in the number  $n$  of variables. The second one uses the iteration of the shift operator (increasing  $m$  by one) by Heckman and Opdam, which can be effectively described using the Dunkl operators [9]. The third one, based on the formula due to Berest [10], is the most general: it works for all locus configurations, see [3].

In this note we present a new formula for the BA function both in rational and trigonometric cases as a simple iterated residue. The structure of the integrand is the same as in other integral formulas known in the theory of Calogero–Moser system and related Jack polynomials (see [11, 12, 13, 14]), but the integration cycle is different and adapted only for the case of integer multiplicities.

We present also another formula for the Baker-Akhiezer function as a Selberg-type integral, which can be considered as analytic continuation of the residue formula from  $m$  to  $-m - 1$ , which is an obvious symmetry of the Calogero–Moser operator. The comparison of these two formulas with other known forms of the BA function can be interpreted as new explicit evaluation of the special Selberg-type integrals, which are probably new. We present similar results also in the deformed  $A(n, 1)$ -case discovered in [15, 16].

Our approach is based on a generalisation of the identity, which plays an important role in the theory of Jack polynomials [17, 18] and in various versions used in [11, 12, 19, 20]. In particular, Langmann [20] suggested a simple explanation of this identity within the theory of Calogero–Moser models with different types of particles [21, 22], which is very convenient for us and allows to extend it for the general  $A(n, m)$  deformation [23].

## 2. RATIONAL BA FUNCTION

The following result can be considered as a version of the “adding particle” approach from [1]. Let us introduce for any two set of variables  $u_1, \dots, u_k$  and  $v_1, \dots, v_l$  the function

$$(4) \quad A(u, v) = \prod_{i=1}^k \prod_{j=1}^l (u_i - v_j).$$

We will also use the notation

$$\bar{u} = u_1 + \dots + u_k.$$

For fixed distinct  $x_1, \dots, x_{k+1} \in \mathbb{C}$  let us choose the cycle of integration  $\sigma$  in  $k$  integration variables  $z_i$  as a product of small circles  $|z_i - x_i| = \epsilon$  around the first  $k$  points  $x_i, i = 1, \dots, k$  and denote by  $dz$  the differential form

$$dz = dz_1 \wedge dz_2 \wedge \dots \wedge dz_k.$$

**Theorem 2.1.** *The BA functions with  $k$  and  $k + 1$  particles are related by the following iterated residue formula*

$$(5) \quad \oint_{\sigma} \frac{A(x)^{m+1} A(z)^{m+1}}{A(z, x)^{m+1}} e^{\lambda_{k+1}(\bar{x} - \bar{z})} \Psi_m^{(k)}(z, \lambda_1, \dots, \lambda_k) dz = C_1 \Psi_m^{(k+1)}(x, \lambda_1, \dots, \lambda_{k+1})$$

where

$$C_1 = C_1(k, m, \lambda) = \frac{(2\pi i)^k}{(m!)^k} \prod_{i=1}^k (\lambda_i - \lambda_{k+1})^m.$$

Iterating this procedure we come to the following formula for the BA function. Note that for  $n = 1$  the Calogero–Moser operator (3) becomes simply the second derivative and

$$\Psi_m^{(1)}(x, \lambda) = e^{\lambda_1 x_1}.$$

By adding one integration variable at each step we will have  $\frac{n(n-1)}{2}$  integration variables, which we denote  $t_{i,j}$  with  $1 \leq i \leq j \leq n-1$ . It is convenient also to denote  $t_{i,n} = x_i, i = 1, \dots, n$ . The integrand has the following form (cf. [11, 12])

$$\omega_m = \prod_{i \leq j, l \leq j+1 \leq n} (t_{i,j} - t_{l,j+1})^{-m-1} \prod_{1 \leq i < l \leq j < n} (t_{i,j} - t_{l,j})^{2+2m} \prod_{l \leq j < n} e^{(\lambda_j - \lambda_{j+1})t_{l,j}} dt_{l,j}.$$

We assume that  $x_i$  are distinct and choose the cycle of integration  $\Sigma$  as the product of circles  $|t_{k,j} - x_k| = \epsilon(n-j)$  with  $\epsilon$  small enough.

**Corollary 2.2.** *For any given positive integer  $m$  the rational Baker-Akhiezer function (1) can be given by the following iterated residue formula*

$$(6) \quad \Psi_m^{(n)}(x, \lambda) = \left( \frac{m!}{2\pi i} \right)^{\frac{n(n-1)}{2}} e^{\lambda_n \bar{x}} A(x)^{1+m} A(\lambda)^{-m} \oint_{\Sigma} \omega_m.$$

To explain another integral representation of BA function note that the Calogero-Moser operator (3) is invariant under the change

$$m \rightarrow -1 - m.$$

This leads to the following formula for BA function as a Selberg-type integral [25].

Let us assume for convenience that  $x_i, i = 1, \dots, k+1$  have distinct imaginary parts and  $\lambda_i - \lambda_j$  have negative real parts for all  $i < j, i, j = 1, \dots, k+1$ . Choose the contour of integration  $\gamma$  such that  $z_i = x_i + \tau_i, i = 1, \dots, k$  with real variables  $\tau_i$ , changing from 0 to  $\infty$ .

**Theorem 2.3.** *The rational BA functions for  $k$  and  $k+1$  particles are related by the following Selberg-type integral formula*

$$(7) \quad \int_{\gamma} \frac{A(z, x)^m}{A(x)^m A(z)^m} e^{\lambda_{k+1}(\bar{x} - \bar{z})} \Psi_m^{(k)}(z, \lambda_1, \dots, \lambda_k) dz = C_2 \Psi_m^{(k+1)}(x, \lambda_1, \dots, \lambda_{k+1}),$$

where

$$C_2 = C_2(m, k, \lambda) = (m!)^k \prod_{i=1}^k (\lambda_{k+1} - \lambda_i)^{-m-1}.$$

Consider now in the same variables  $t_{i,j}$  the form  $\alpha_m = \omega_{-m-1}$  :

$$\alpha_m = \prod_{i \leq j, l \leq j+1 \leq n} (t_{i,j} - t_{l,j+1})^m \prod_{1 \leq i < l \leq j < n} (t_{i,j} - t_{l,j})^{-2m} \prod_{l \leq j < n} e^{(\lambda_j - \lambda_{j+1})t_{l,j}} dt_{l,j}$$

and choose the integration contour  $\Gamma$  by assuming that  $t_{i,j} = t_{i,j+1} + \tau_{i,j}$  with real variables  $\tau_{i,j}, 1 \leq i \leq j = 1, \dots, n-1$  changing from zero to infinity.

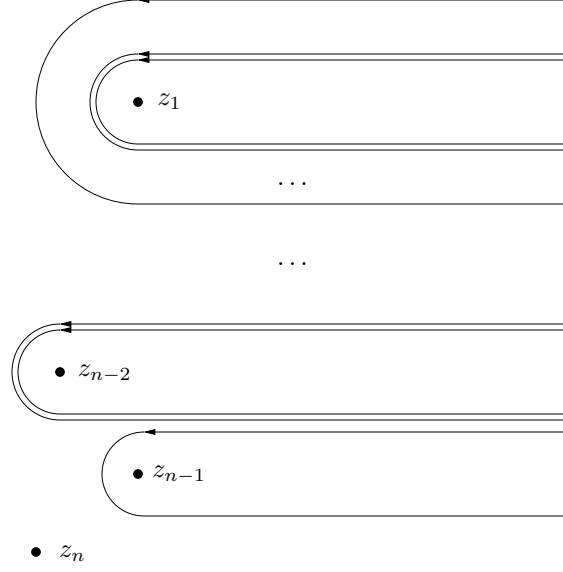
**Corollary 2.4.** *For any positive integer  $m$  the rational Baker-Akhiezer function (1) can be given by the following Selberg-type integral*

$$(8) \quad \Psi_m^{(n)}(x, \lambda) = ((-1)^{m+1} m!)^{-\frac{n(n-1)}{2}} e^{\lambda_n \bar{x}} A(x)^{-m} A(\lambda)^{m+1} \int_{\Gamma} \alpha_m.$$

These two different formulas are actually related by analytic continuation. To see this consider the same integral (8) but over Pochhammer contour  $\Pi$ :

$$I_P(m) = \int_{\Pi} \alpha_m$$

(see Figure 1). It converges for all  $m \in \mathbb{C}$  and is related for positive real  $m$  to the

FIGURE 1. The Pochhammer contour  $\Pi$  in the rational case

Selberg-type integral

$$I(m) = \int_{\Gamma} \alpha_m$$

in a simple way:

$$I_P(m) = (e^{2\pi im} - 1)^{\frac{n(n-1)}{2}} I(m).$$

Now replace  $m!$  in (8) by  $\Gamma(m+1)$ , where  $\Gamma(x)$  is the classical Euler gamma-function and note that

$$\Gamma(m+1)(e^{2\pi im} - 1) = 2ie^{\pi im} \sin \pi m \Gamma(m+1) = (-1)^{m+1} 2\pi i \Gamma^{-1}(-m)$$

because of the reflection property of gamma-function:

$$\Gamma(x)\Gamma(1-x) = \frac{\pi}{\sin \pi x}.$$

We note here a remarkable similarity to Riemann's first proof [24] of the reflection property of the Riemann zeta function

$$\pi^{-\frac{s}{2}} \Gamma\left(\frac{s}{2}\right) \zeta(s) = \pi^{-\frac{1-s}{2}} \Gamma\left(\frac{1-s}{2}\right) \zeta(1-s).$$

### 3. TRIGONOMETRIC CASE

It is actually more convenient for us to use the hyperbolic rather than trigonometric functions, but all the results are automatically applied to both cases because of the algebraic nature of BA function.

The trigonometric version of the BA function satisfies the equation

$$(9) \quad \mathcal{L}_m^{(n)} \Phi_m^{(n)} = -4(\nu, \nu) \Phi_m^{(n)},$$

where

$$(10) \quad \mathcal{L}_m^{(n)} = -\Delta + \sum_{i < j}^n \frac{2m(m+1)}{\sinh^2(x_i - x_j)}$$

is the Sutherland operator. It has the form

$$(11) \quad \Phi_m^{(n)} = \frac{P(x, \nu)}{B(x)^m C_m(\nu)} \exp 2(\nu, x),$$

where  $P(x, \nu)$  is a trigonometric polynomial in  $x$  and a usual polynomial in  $\nu$  with the leading term  $B(x)^m A(\nu)^m$ , where

$$B(x) = \prod_{i < j}^n \sinh(x_i - x_j)$$

and

$$C_m(\nu) = \prod_{k=1}^m \prod_{i < j}^n (\nu_i - \nu_j - k).$$

The normalisation constant  $C_m(\nu)$  is chosen in such a way that

$$\lim_{x \rightarrow +\infty} \frac{P(x, \nu)}{B(x)^m C_m(\nu)} = 1$$

when  $x \rightarrow +\infty$  in the Weyl chamber  $x_1 > x_2 > \dots > x_n$  (see [7, 8]). In the exponential variables

$$u_i = \exp 2x_i, \quad i = 1, \dots, n$$

the BA function can be rewritten as

$$(12) \quad \Phi_m^{(n)} = \frac{Q(u, \nu)}{A(u)^m C_m(\nu)} u^\nu,$$

where  $u^\nu = u_1^{\nu_1} \dots u_n^{\nu_n}$  and  $Q(u, \nu)$  is polynomial in  $\nu$  with the leading term  $A(u)^m A(\nu)^m$ .

It is convenient to modify the definition of  $A(w, u)$  as follows

$$A^*(w, u) = \prod_{i \leq j} (w_i - u_j) \prod_{i > j} (u_j - w_i)$$

to include some sign factor.

Let  $w = (w_1, \dots, w_k)$ ,  $u = (u_1, \dots, u_{k+1})$  and the cycle  $\sigma$  similarly to the rational case be the product of circles  $|u_i - w_i| = \epsilon$ ,  $i = 1, \dots, k$  with small positive  $\epsilon$ .

**Theorem 3.1.** *The trigonometric BA functions  $\Phi_m^{(k)}(w, \nu_1, \dots, \nu_k)$  and  $\Phi_m^{(k+1)}(u, \nu_1, \dots, \nu_{k+1})$  are related by the following iterated residue formula*

$$(13) \quad \prod_{i=1}^{k+1} u_i^{\nu_{k+1}} \int_{\sigma} \frac{A(u)^{m+1} A(w)^{m+1}}{A^*(w, u)^{m+1}} \prod_{i=1}^k w_i^{m-\nu_{k+1}} \Phi_m^{(k)}(w, \nu_1, \dots, \nu_k) dw = C \Phi_m^{(k+1)}(u, \nu_1, \dots, \nu_{k+1})$$

where

$$C = (2\pi i)^k \prod_{i=1}^k \binom{\nu_i - \nu_{k+1} - 1}{m}$$

and

$$\binom{a}{m} = \frac{a(a-1)\dots(a-m+1)}{m!}.$$

As before we introduce  $\frac{n(n-1)}{2}$  integration variables  $t_{i,j}$  with  $1 \leq i < j \leq n$  with the convention that  $t_{i,n} = u_i$ ,  $i = 1, \dots, n$ . The integrand in the trigonometric case has the following form (cf. [11])

$$\omega_m^* = \prod_{j=1}^{n-1} A_{j,j+1}^{-m-1} \prod_{j=1}^n A_{j,j}^{2+2m} \prod_{l < j} \prod_{i \leq j}^{n-1} t_{i,j}^{\nu_j - \nu_{j+1} + m} dt_{i,j},$$

where

$$A_{j,j+1}(t) = \prod_{i \leq j, l \leq j+1, i \leq l}^n (t_{i,j} - t_{l,j+1}) \prod_{i \leq j, l \leq j+1, i > l}^n (t_{i,j} - t_{l,j+1})$$

and

$$A_{j,j}(t) = \prod_{1 \leq i < l \leq j \leq n-1} (t_{i,j} - t_{l,j}).$$

The cycle of integration  $\Sigma$  as before is the product of circles  $|t_{i,j} - u_i| = \epsilon(n-j)$  with  $\epsilon$  small enough. Note that in contrast to [11] the origin is outside of these circles, so we have no problems with the multi-valuedness of the integrand.

**Corollary 3.2.** *The trigonometric BA function (12) can be given as an iterated residue*

$$(14) \quad \Phi_m^{(n)}(u, \nu) = C(n, m, \nu) \prod_{i=1}^n u_i^{\nu_i} A(u)^{1+m} \int_{\Sigma} \omega_m^*$$

with

$$C(n, m, \nu)^{-1} = (2\pi i)^{\frac{n(n-1)}{2}} \prod_{i < j}^n \binom{\nu_i - \nu_j - 1}{m}.$$

Similarly to the rational case we have also the following Selberg-type representation of BA functions.

Assume for convenience that the complex numbers  $u_1, u_2, \dots, u_{k+1}$  have different arguments and consider the contour  $\gamma^*$  when  $w_i$  belongs to the segment joining 0 and  $u_i$  for  $i = 1, \dots, k$ . In other words, we assume that  $w_i = \tau_i u_i$ ,  $i = 1, \dots, k$  with real  $\tau_i$  between 0 and 1. We assume also that  $\nu_i - \nu_{k+1}$  have large positive real parts to guarantee the convergence of the following integral.

**Theorem 3.3.** *The BA functions  $\Phi_m^{(k)}$  and  $\Phi_m^{(k+1)}$  are related by the Selberg-type integral formula*

$$(15) \quad \prod_{i=1}^{k+1} u_i^{\nu_{k+1}} \int_{\gamma^*} \frac{A^*(w, u)^m}{A(u)^m A(w)^m} \prod_{i=1}^k w_i^{-m-1-\nu_{k+1}} \Phi_m^{(k)}(w, \nu_1, \dots, \nu_k) dw = C_3 \Phi_m^{(k+1)}(u, \nu_1, \dots, \nu_{k+1})$$

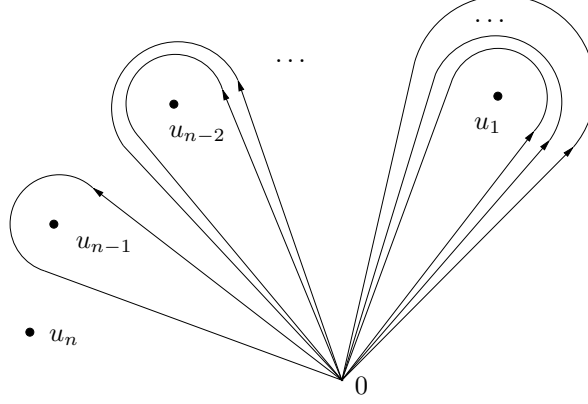
where

$$C_3 = (-1)^{km} \prod_{i=1}^k \frac{\Gamma(m+1)\Gamma(\nu_i - \nu_{k+1})}{\Gamma(\nu_i - \nu_{k+1} + m + 1)}.$$

Consider

$$\alpha_m^* = \omega_{-m-1}^* = \prod_{j=1}^{n-1} A_{j,j+1}^m \prod_{j=1}^n A_{j,j}^{-2m} \prod_{l < j} \prod_{i \leq j}^{n-1} t_{i,j}^{\nu_j - \nu_{j+1} - m - 1} dt_{i,j}$$

and choose the contour of integration  $\Gamma^*$  such that  $t_{i,j} = \tau_{i,j} t_{i,j+1}$  with  $\tau_{i,j} \in [0, 1]$ .


 FIGURE 2. The Pochhammer contour  $\Pi$  in the trigonometric case

**Corollary 3.4.** *The trigonometric BA function (12) can be given as a Selberg-type integral*

$$(16) \quad \Phi_m^{(n)}(u, \nu) = C_4 \prod_{i=1}^n u_i^{\nu_i} A(u)^{1+m} \int_{\Gamma^*} \alpha_m^*$$

with

$$C_4^{-1} = (-1)^{m \frac{n(n-1)}{2}} \prod_{i < j}^n \frac{\Gamma(m+1)\Gamma(\nu_i - \nu_j)}{\Gamma(\nu_i - \nu_j + m + 1)}.$$

The same calculation as in the rational case shows that these two integral representations are related by analytic continuation. The corresponding analogue of the Pochhammer contour is shown in Figure 2.

#### 4. DEFORMED CASE

Consider now the deformed Calogero–Moser operator [3, 16]

$$(17) \quad L_m^{(n,1)} = - \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} - m \frac{\partial^2}{\partial y^2} + \sum_{i < j}^n \frac{2m(m+1)}{(x_i - x_j)^2} + \sum_{i=1}^n \frac{2(m+1)}{(x_i - y)^2},$$

corresponding to an additional particle with mass  $\frac{1}{m}$  interacting with  $n$  Calogero–Moser particles in a special way. According to [16] for any integer  $m \in \mathbb{Z}$  it has the eigenfunction (which we will call deformed Baker–Akhiezer function) of the form

$$(18) \quad \Psi_m^{(n,1)} = \frac{P(x, y, \lambda, \mu)}{A(x)^{m^*} A(x, y) A(\lambda)^{m^*} A(\lambda, \mu)} \exp((\lambda, x) + \frac{1}{m} \mu y),$$

where  $m^* = \max(m, -m - 1)$ ,  $\lambda = (\lambda_1, \dots, \lambda_n) \in \mathbb{C}^n$ ,  $\mu \in \mathbb{C}$ ,

$$A(x) = \prod_{i < j}^n (x_i - x_j), \quad A(x, y) = \prod_{i=1}^n (x_i - y)$$



and  $P(x, y, \lambda, \mu)$  is a polynomial in all variables with the highest degree term  $A(x)^{m^*} A(x, y) A(\lambda)^{m^*} A(\lambda, \mu)$ .

Let  $\Psi_m^{(n)}(z_1, \dots, z_n; \lambda_1, \dots, \lambda_n)$  be the BA function (1) from the first section. For positive  $m$  we assume as before that all  $x_i$  have different imaginary parts, real part of  $\lambda_i - \mu$ ,  $i = 1, \dots, n$  are negative and choose the contour  $\gamma$  of integration  $z_i = x_i + \tau_i$  by considering  $\tau_i \in \mathbb{R}_+$ ,  $i = 1, \dots, n$ .

For negative  $m$  we assume simply that  $x_i$  are distinct and choose the cycle  $\sigma$  as a product of circles  $|z_i - x_i| = \epsilon$ ,  $i = 1, \dots, n$ .

**Theorem 4.1.** *The rational deformed Baker-Akhiezer function (18) for positive integer  $m$  can be given by the following Selberg-type integral*

$$(19) \quad \Psi_m^{(n,1)}(x, y, \lambda, \mu) = \frac{C_5}{A(x)^m A(x, y)} \int_{\gamma} \frac{A(z, x)^m A(z, y)}{A(z)^m} e^{\mu(\bar{x} - \bar{z} + \frac{y}{m})} \Psi_m^{(n)}(z, \lambda) dz$$

with

$$C_5 = \frac{\prod_{i=1}^n (\mu - \lambda_i)^{m+1}}{(m!)^n}.$$

For negative  $m = -m^* - 1$ ,  $m^* \in \mathbb{Z}_+$  it can be represented as an iterated residue

$$(20) \quad \Psi_m^{(n,1)} = C_6 \frac{A(x)^{m^*+1}}{A(x, y)} \int_{\sigma} \frac{A(z)^{m^*+1} A(z, y)}{A^*(z, x)^{m^*+1}} e^{\mu(\bar{x} - \bar{z} + \frac{y}{m})} \Psi_m^{(n)}(z, \lambda) dz,$$

where

$$C_6 = \left( \frac{m^*!}{2\pi i} \right)^n \prod_{i=1}^n (\lambda_i - \mu)^{-m^*}.$$

In the trigonometric case (see [16]) we have the operator

$$(21) \quad \mathcal{L}_m^{(n,1)} = - \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} - m \frac{\partial^2}{\partial y^2} + \sum_{i < j}^n \frac{2m(m+1)}{\sinh^2(x_i - x_j)} + \sum_{i=1}^n \frac{2(m+1)}{\sinh^2(x_i - y)},$$

and the BA function of the form

$$(22) \quad \Phi_m^{(n,1)} = \frac{Q(u, v, \nu, \mu)}{A(u)^{m^*} A(u, v) C_{m^*}(\nu) C(\nu, \mu)} u^\nu v^{\frac{\mu}{m}},$$

where  $u_i = e^{2x_i}$ ,  $i = 1, \dots, n$ ,  $v = e^{2y}$ ,  $Q(u, v, \nu, \mu)$  is polynomial in  $\nu \in \mathbb{C}^n$ ,  $\mu \in \mathbb{C}$  with the leading term  $A(u)^{m^*} A(u, v) A(\nu)^{m^*} A(\nu, \mu)$  and

$$C(\nu, \mu) = \prod_{i=1}^n \left( \nu_i - \mu - \frac{1+m}{2} \right) = \prod_{i=1}^n \left( \nu_i - \mu + \frac{m^*}{2} \right).$$

As in the non-deformed case, the form of  $C(\nu, \mu)$  is uniquely determined by the property

$$\lim \frac{Q(u, v, \nu, \mu)}{A(u)^{m^*} A(u, v) C_{m^*}(\nu) C(\nu, \mu)} = 1$$

when  $(u, v) \rightarrow +\infty$  in the chamber  $u_1 > u_2 > \dots > u_n > v$  (see [7, 8]).

Let  $\Phi_m^{(n)}(u_1, \dots, u_n, \nu_1, \dots, \nu_n)$  be the non-deformed  $n$  particle BA function (12). For positive  $m$  we assume that all  $u_i$  have different arguments, all  $\lambda_i - \nu$  are real negative and choose the contour  $\gamma^*$  of integration  $w_i = \tau_i u_i$  with real  $\tau_i \in [0, 1]$ . For negative  $m = -1 - m^*$  we assume simply that  $u_i$  are distinct and choose the cycle  $\sigma^*$  as a product of circles  $|w_i - u_i| = \epsilon$ ,  $i = 1, \dots, n$ .

**Theorem 4.2.** *The deformed trigonometric Baker-Akhiezer function  $\Phi_m^{(n,1)}(u, v, \nu, \mu)$  for positive integer  $m$  can be given by the following Selberg-type integral (23)*

$$\Phi_m^{(n,1)} = C_7 \frac{\prod_{i=1}^n u_i^{\mu + \frac{1-m}{2}} v_i^{\frac{\mu}{m}}}{A(u)^m A(v)} \int_{\gamma^*} \frac{A(w, u)^m A(w, v)}{A(w)^m} \prod_{i=1}^n w_i^{-\mu - \frac{1+m}{2}} \Phi_m^{(n)}(w, \nu) dw$$

with

$$C_7^{-1} = \prod_{k=1}^n \frac{\Gamma(m+1) \Gamma(\nu_k - \mu + \frac{1-m}{2})}{\Gamma(\nu_k - \mu + \frac{m+3}{2})}$$

and for negative  $m = -m^* - 1$ ,  $m^* \in \mathbb{Z}_+$  as an iterated residue (24)

$$\Phi_m^{(n,1)} = C_8 \frac{\prod_{i=1}^n u_i^{\mu + \frac{m^*+2}{2}} v_i^{\frac{\mu}{m}} A(u)^{m^*+1}}{A(u, v)} \int_{\gamma^*} \frac{A(w)^{m^*+1} A(w, v)}{A(w, u)^{m^*+1}} \prod_{i=1}^n w_i^{-\mu + \frac{m^*}{2}} \Phi_m^{(n)}(w, \nu) dw$$

where

$$C_8^{-1} = (2\pi i)^n \prod_{i=1}^n \begin{pmatrix} \nu_i - \mu + \frac{m^*}{2} \\ m^* \end{pmatrix}.$$

## 5. PROOFS: MAIN IDENTITY

Let

$$\begin{aligned} L_m^{k,l}(x, y) &= - \left( \frac{\partial^2}{\partial x_1^2} + \cdots + \frac{\partial^2}{\partial x_k^2} \right) - m \left( \frac{\partial^2}{\partial y_1^2} + \cdots + \frac{\partial^2}{\partial y_l^2} \right) + \sum_{i < j}^k \frac{2m(m+1)}{\sinh^2(x_i - x_j)} \\ (25) \quad &+ \sum_{i < j}^l \frac{2(m^{-1} + 1)}{\sinh^2(y_i - y_j)} + \sum_{i=1}^l \sum_{j=1}^k \frac{2(m+1)}{\sinh^2(x_i - y_j)} \end{aligned}$$

and

$$\begin{aligned} L_m^{k,l}(x, y) &= - \left( \frac{\partial^2}{\partial x_1^2} + \cdots + \frac{\partial^2}{\partial x_k^2} \right) - m \left( \frac{\partial^2}{\partial y_1^2} + \cdots + \frac{\partial^2}{\partial y_l^2} \right) + \sum_{i < j}^k \frac{2m(m+1)}{(x_i - x_j)^2} \\ (26) \quad &+ \sum_{i < j}^l \frac{2(m^{-1} + 1)}{(y_i - y_j)^2} + \sum_{i=1}^l \sum_{j=1}^k \frac{2(m+1)}{(x_i - y_j)^2}, \end{aligned}$$

be respectively trigonometric and rational deformed Calogero–Moser–Sutherland (CMS) operators in  $x, y$  variables [23]. Let  $L_m^{p,q}$  be a similar operator in variables  $z_1, \dots, z_p, w_1, \dots, w_q$ . Let  $A(u, v)$  be given by (4),

$$B(u, v) = \prod_{i=1}^k \prod_{j=1}^l \sinh(u_i - v_j)$$

and  $\bar{x} = x_1 + \cdots + x_k$  as before.

The key observation <sup>1</sup> comes from the following result (cf. Langmann [20]).

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<sup>1</sup>As we have learnt from Martin Hallnas a similar result can be extracted from his paper with Edwin Langmann [28].

**Theorem 5.1.** *The following identity holds for the trigonometric deformed CMS operators:*

$$(27) \quad L_m^{k,l} K = L_m^{p,q} K + C_0 K,$$

where

$$K(x, y; z, w) = \frac{B(x, z)^m B(y, z) B(x, w) B(y, w)^{1/m}}{B(x)^m B(x, y) B(y)^{1/m} B(z)^m B(z, w) B(w)^{1/m}} e^{\mu(\bar{x} - \bar{z} + \frac{1}{m}(\bar{y} - \bar{w}))}$$

and

$$C_0 = \frac{1}{4} m^2 \left[ \left( k - p + \frac{1}{m}(l - q) \right)^3 - \left( k - p + \frac{1}{m^3}(l - q) \right) \right] + \mu^2 \left( p + \frac{q}{m} - k - \frac{l}{m} \right).$$

In the rational case we have the same identity (27) for

$$K(x, y; z, w) = \frac{A(x, z)^m A(y, z) A(x, w) A(y, w)^{1/m}}{A(x)^m A(x, y) A(y)^{1/m} A(z)^m A(z, w) A(w)^{1/m}} e^{\mu(\bar{x} - \bar{z} + \frac{1}{m}(\bar{y} - \bar{w}))}$$

and  $C_0 = \mu^2 \left( p + \frac{q}{m} - k - \frac{l}{m} \right)$ .

The identity (27) with  $l = q = 0$  goes back to Stanley and Macdonald [17, 18]. In the case  $q = 0$  and arbitrary  $l$  it appeared in [27] (see part (iii) in Lemma 3). To prove it in the general case we borrow the idea from the work of Langmann [20]. It is based on the following result by Sen [22] (see also [21]).

Consider the following generalised CMS operator describing the interacting particles of different masses on a line:

$$(28) \quad H = - \sum_{j=1}^N \frac{1}{m_j} \frac{\partial^2}{\partial x_j^2} + \sum_{j < k} \frac{\gamma_{jk}}{\sinh^2(x_j - x_k)}.$$

**Theorem 5.2** (Sen [22]). *For the coupling constants of the special form*

$$(29) \quad \gamma_{ij} = (m_i + m_j) \beta (m_i m_j \beta - 1),$$

where  $\beta$  is an arbitrary parameter, the operator (28) has the following eigenfunction

$$(30) \quad \Phi_0 = \prod_{i < j} \sinh^{\beta m_i m_j} (x_i - x_j) :$$

$$H \Phi_0 = E_0 \Phi_0$$

with the eigenvalue

$$(31) \quad E_0 = -\frac{\beta^2}{3} \left( \left( \sum_{j=1}^N m_j \right)^3 - \sum_{j=1}^N m_j^3 \right).$$

Now we note that  $\gamma_{ij} = 0$  if  $m_j = -m_i$  or  $m_j = (m_i \beta)^{-1}$ . Choosing

$$m_1 = \dots = m_k = 1, \quad m_{k+1} = \dots = m_{k+l} = m^{-1},$$

$$m_{k+l+1} = \dots = m_{k+l+p} = -1, \quad m_{k+l+p+1} = \dots = m_{k+l+p+q} = -m^{-1},$$

where  $m = -\beta$ , we see that the operator  $H$  reduces to the difference

$$H = L_m^{k,l} - L_m^{p,q}$$

of two decoupled deformed CMS operators and the relation  $H \Phi_0 = E_0 \Phi_0$  implies the identity (27).

We note that this decoupling can be used to characterise the deformed CMS operators among all generalised CMS operators with different masses (28). It does

not imply though the quantum integrability, which had to be proven by other means (see [23, 27]).

The identity (27) suggests that the function  $K(x, y, z, w)$  could be used as the kernel of the integral representation transforming the eigenfunctions of one of the deformed CMS operators to another, although to make this precise could be a non-trivial task (see e.g. [19], where a similar problem is discussed).

In particular, choosing in the rational case  $k = p + 1, l = q = 0$  we come to the integral of the type (5). Choosing a suitable cycle (contour) of integration we come to the formula for the BA function. For example, in Theorem 2.1 the cycle  $\sigma$  is chosen in such a way that each integration has only one non-zero residue to compute, which guarantee that the result will be of the required form.

The rest of the proof is based on the fact well-known in the theory of BA function (see e.g. [8]) that properly normalised trigonometric BA function can be characterised as the special eigenfunction of CMS operator, which is a particular case of the Heckman-Opdam asymptotic solution [29]. The rational case can be treated as the limit of the trigonometric one (see [1]).

## 6. EXAMPLES

In the simplest case  $n = 2$  the rational BA function is known to have the form (see e.g. [1]):

$$(32) \quad \Psi_m^{(2)} = (\lambda_1 - \lambda_2)^{-m} \left(D_{12} - \frac{2m}{x_1 - x_2}\right) \left(D_{12} - \frac{2(m-1)}{x_1 - x_2}\right) \dots \left(D_{12} - \frac{2}{x_1 - x_2}\right) \exp(\lambda_1 x_1 + \lambda_2 x_2),$$

where

$$D_{12} = \frac{\partial}{\partial x_1} - \frac{\partial}{\partial x_2}.$$

We have two different representations for it. The first one is as a residue

$$(33) \quad \Psi_m^{(2)} = \frac{m!(x_1 - x_2)^{m+1}}{(\lambda_1 - \lambda_2)^m} e^{\lambda_2(x_1+x_2)} \operatorname{Res}_{z=x_1} \frac{e^{(\lambda_1-\lambda_2)z}}{(z-x_1)^{m+1}(z-x_2)^{m+1}},$$

the second one is the integral

$$(34) \quad \Psi_m^{(2)} = \frac{(\lambda_2 - \lambda_1)^{m+1}}{m!(x_1 - x_2)^m} e^{\lambda_2(x_1+x_2)} \int_{x_1}^{+\infty} (z-x_1)^m (z-x_2)^m e^{(\lambda_1-\lambda_2)z} dz,$$

which in this case can be effectively computed using the  $\Gamma$ -integral

$$\Gamma(a) = \int_0^{+\infty} z^{a-1} e^{-z} dz = (a-1)!$$

for positive integer  $a$ .

For  $n = 3$  the corresponding BA function  $\Psi_m^{(3)}(x_1, x_2, x_3, \lambda_1, \lambda_2, \lambda_3)$  can be written respectively as follows:

$$(35) \quad \Psi_m^{(3)} = C \operatorname{Res}_{z_1=x_1} \operatorname{Res}_{z_2=x_2} \operatorname{Res}_{w=z_1} \frac{(z_1 - z_2)^{2m+2} e^{(\lambda_2-\lambda_3)\bar{z}} e^{(\lambda_1-\lambda_2)w}}{\prod_{i=1}^2 (w - z_i)^{m+1} \prod_{i=1}^2 \prod_{j=1}^3 (z_i - x_j)^{m+1}}$$

where

$$C = \frac{(m!)^3 \prod_{i<j}^3 (x_i - x_j)^{m+1}}{\prod_{i<j}^3 (\lambda_i - \lambda_j)^m} e^{\lambda_3 \bar{x}}$$

(36)

$$\Psi_m^{(3)} = D \int_{x_1}^{+\infty} \int_{x_2}^{+\infty} \int_{z_1}^{+\infty} \frac{\prod_{i=1}^2 (w - z_i)^m \prod_{i=1}^2 \prod_{j=1}^3 (z_i - x_j)^m e^{(\lambda_2 - \lambda_3)\bar{z}} e^{(\lambda_1 - \lambda_2)w}}{(z_1 - z_2)^{2m}} dz_1 dz_2 dw,$$

with

$$D = (-1)^{m+1} \frac{\prod_{i < j}^3 (\lambda_i - \lambda_j)^{m+1}}{(m!)^3 \prod_{i < j}^3 (x_i - x_j)^m} e^{\lambda_3 \bar{x}}$$

where as before  $\bar{x} = x_1 + x_2 + x_3$ ,  $\bar{z} = z_1 + z_2$ .

One can interpret these formulas either as a new way of representing of BA function, or as an explicit evaluation of the Selberg-type integral (36) in terms of the BA function, which can be computed by other methods as well (see [1, 2, 3]). The same of course is true for general  $n$  and in the deformed case.

We should mention here the work of Kazarnovski-Krol [26], who found explicit expression of certain generalised Selberg integrals using Opdam's results. Our formulas work only for integer values of parameter  $m$ , but depend on the additional variables  $x$ .

## 7. CONCLUDING REMARKS

It would be interesting to explore the possibilities of choosing different cycles to produce the integral formulas for the super Jack polynomials [27]. For the usual Jack polynomials such integral formulas were obtained in [12, 13, 14, 20].

Our approach can be also naturally extended to the (deformed)  $BC_n$ -case and related theory of (super) Jacobi polynomials [30]. We will discuss this in more detail elsewhere.

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## REFERENCES

- [1] O.A. Chalykh, A.P. Veselov *Commutative rings of partial differential operators and Lie algebras*. Comm. Math. Phys. 126 (1990), no. 3, 597–611.
- [2] A.P. Veselov, K.L. Styrkas, O.A. Chalykh *Algebraic integrability for Schrödinger equation and finite reflection groups*. Theor. Math. Phys., 1993, v.94(2), p.253–275.
- [3] O.A. Chalykh, M.V. Feigin, A.P. Veselov *Multidimensional Baker-Akhiezer functions and Huygens' principle*. Comm. Math. Phys. 206 (1999), no. 3, 533–566.
- [4] I.M. Krichever *Methods of algebraic geometry in the theory of nonlinear equations*. Uspekhi Mat. Nauk **32** (6) (1977), 183–208.
- [5] M. Feigin, A.P. Veselov *Quasi-invariants of Coxeter groups and  $m$ -harmonic polynomials*. Int. Math. Res. Not. 2002, no. 10, 521–545.
- [6] P. Etingof, V. Ginzburg *On  $m$ -quasi-invariants of a Coxeter group*. Mosc. Math. J. 2 (2002), no. 3, 555–566.
- [7] O.A. Chalykh *Bispectrality for the quantum Ruijsenaars model and its integrable deformation*. J. Math. Phys. 41 (2000), no. 8, 5139–5167.

- [8] M. Feigin *Bispectrality for deformed Calogero-Moser-Sutherland systems*. J. Nonlinear Math. Phys. 12 (2005), suppl. 2, 95–136.
- [9] G. Heckman *A remark on the Dunkl differential-difference operators*. In *Harmonic analysis on reductive groups* (Brunswick, ME, 1989), 181–191, Progr. Math., 101, Birkhuser Boston, Boston, MA, 1991.
- [10] Yu. Berest *Huygens' principle and the bispectral problem*. The bispectral problem (Montreal, PQ, 1997), 11–30, CRM Proc. Lecture Notes, 14, Amer. Math. Soc., Providence, RI, 1998.
- [11] A. Kazarnovski-Krol. *Cycles for asymptotic solutions and the Weyl group*. The Gelfand Mathematical Seminars, 1993–1995, 123–150, Gelfand Math. Sem., Birkhäuser Boston, Boston, MA, 1996.
- [12] H. Awata, Y. Matsuo, S. Odake, J. Shiraishi *Excited states of the Calogero-Sutherland model and singular vectors of the  $W_N$  algebra*. Nuclear Phys. B 449 (1995), no. 1-2, 347–374.
- [13] A. Okounkov and G. Olshanski *Shifted Jack polynomials, binomial formula, and applications* Math. Res. Letters, 4 (1997), 6978.
- [14] V.B. Kuznetsov, V.V. Mangazeev, E. K. Sklyanin *Q-operator and factorised separation chain for Jack polynomials*. *Indag. Mathem., N.S.*, 14 (3,4) (2003), 451–482.
- [15] A.P. Veselov, M.V. Feigin, O.A. Chalykh *New integrable deformations of quantum Calogero - Moser problem*. Russian Math. Surveys 51, no.3, 1996, 185–186.
- [16] O.A. Chalykh, M.V. Feigin, A.P. Veselov *New integrable generalizations of Calogero-Moser quantum problem*. J. Math. Phys. 39 (2), 1998, 695–703.
- [17] R. Stanley *Some combinatorial properties of Jack symmetric functions*. Adv. Math. 77 (1989), no. 1, 76–115.
- [18] I. G. Macdonald *Symmetric functions and Hall polynomials*. 2nd edition, Oxford Univ. Press, New York, 1995.
- [19] S.N.M. Ruijsenaars *Elliptic integrable systems of Calogero-Moser type: Some new results on joint eigenfunctions*. In: Proceedings of the 2004 Kyoto Workshop on Elliptic integrable systems, Noumi, M., Takasaki, K., eds., Rokko Lectures in Math., Dept. of Math., Kobe Univ., pp. 223240. Available at: [www.math.kobe-u.ac.jp/publications/rml18/18elliptic.html](http://www.math.kobe-u.ac.jp/publications/rml18/18elliptic.html)
- [20] E. Langmann *Singular eigenfunctions of Calogero-Sutherland type systems and how to transform them into regular ones*. SIGMA 3 (2007), Paper 031, 18 pp.
- [21] V.Ya. Krivnov and A.A. Ovchinnikov *An exactly solvable one-dimensional problem with several particle species*. Theor. Math. Phys. 50 (1982), 100–103.
- [22] D. Sen *A multispecies Calogero-Sutherland model*. Nuclear Phys. B 479 (1996), no. 3, 554–574.
- [23] A. N. Sergeev, A. P. Veselov *Deformed quantum Calogero-Moser systems and Lie superalgebras*, Comm. Math. Phys. 245 (2004), 249–278.
- [24] B. Riemann *Ueber die Anzahl der Primzahlen unter einer gegebenen Grösse*, Monatsberichte der Berliner Akademie, November 1859. In *B. Riemann Gesammelte mathematische Werke, wissenschaftlicher Nachlass und Nachtrge*. Springer-Verlag, Berlin; B. G. Teubner Verlagsgesellschaft, Leipzig, 1990.
- [25] A.N. Varchenko *Selberg integrals*. arXiv:math/0408308
- [26] A. Kazarnovski-Krol *Variation on a theme of Selberg integral*. arXiv:q-alg/9611012.
- [27] A.N. Sergeev, A.P. Veselov *Generalised discriminants, deformed Calogero-Moser-Sutherland operators and super-Jack polynomials*. Adv. Math. 192 (2005), no. 2, 341–375.
- [28] M. Hallnas, E. Langmann *Quantum Calogero-Sutherland type models and generalised classical polynomials*. arXiv:math-ph/0703090
- [29] G. Heckman, E. Opdam *Root systems and hypergeometric functions. I*, Comp. Math. 64 (1987), 329–352.
- [30] A.N. Sergeev, A.P. Veselov  *$BC_\infty$  Calogero-Moser operator and super Jacobi polynomials*. arXiv: 0807.3858.

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