

1 The VOL-CALPUFF Model for Atmospheric Ash
2 Dispersal: II. Application to the Weak Mt. Etna
3 Plume of July 2001

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5 **Abstract.** The application of the VOL-CALPUFF model (*Barsotti et al.*
6 [2007], this issue) to a weak plume erupted from Mt. Etna in July 2001 is
7 here presented and discussed. The reconstruction of the explosive event was
8 obtained by using high-resolution weather forecasts, produced by a mesoscale
9 non-hydrostatic model, and volcanic source data coming from observations
10 and analytical studies. The plume rise and atmospheric dispersal models were
11 investigated over five days of eruption mostly in terms of column height, aer-
12 ial ash concentration and ground deposition. Modeling results are shown as
13 a function of various source conditions and compared to independent obser-
14 vations derived from satellite images and deposit mapping. The application
15 of VOL-CALPUFF clearly highlights the crucial role played by meteorolog-
16 ical conditions in determining dispersal dynamics. Some of the most impor-
17 tant effects described by the model are: a) the large wind field influence on
18 the plume height determination and tilting, b) the contrasting dispersal pat-
19 terns of ash particles of different sizes, c) the complex and somehow non-intuitive
20 distribution of ash on the ground resulting in preferential directions of dis-
21 persal and quite irregular deposit patterns, d) the impossibility to reproduce
22 both the column height and the deposit accumulation pattern by adopting
23 a steady-state vent mass flow rate over the investigated four-day period due
24 to observed temporal changes in eruption dynamics. Modeling results also
25 suggest the need for further integration of simulation outcomes with remote
26 sensing and field reconstructions on ash dispersal processes in future.

1. Introduction

27 We present the first application of the VOL-CALPUFF model, fully described in the
28 companion paper (*Barsotti et al.* [2007], this issue), to a real event in order to highlight its
29 potentialities and weaknesses. The code originates from the U.S. non-commercial model
30 CALPUFF (*Scire et al.* [2000]), developed and used for air quality applications, modified
31 and developed to make it appropriate for volcanological applications. VOL-CALPUFF
32 is an hybrid code in which a part of the process, i.e. the plume rise phase, is described
33 in a Eulerian framework, whereas dispersal and fallout processes are solved with a La-
34 grangian approach. Both processes are modelled taking into account the influence of
35 time-dependent, 3D meteorological conditions throughout the duration of the eruption.
36 VOL-CALPUFF relates the computation of the column height, computed by solving the
37 plume theory equations, to the altitude of release of pyroclastic material. The Lagrangian
38 description of the dispersal process guarantees model reliability both close to vent and
39 far from it (*Nguyen et al.* [1997]). The original CALPUFF code was validated through
40 extensive comparison with other widely used codes for air-pollution modeling and through
41 experimental tests concerning passive and active tracer transport and estimation of ground
42 deposition and atmospheric concentrations (*EPA U.S.* [1998]; *Scire et al.* [2000]). Due
43 to its good performance it is currently proposed by the U. S. EPA as a guideline model
44 for environmental applications such as impact studies of proposed and existing pollutant
45 sources. Similarly, VOL-CALPUFF needs to be validated and tested against volcanolog-
46 ical applications, and this is a challenge due to the uncertainties when characterizing the
47 natural phenomena.

48 The paper presents quantitative comparisons between model results and experimental
49 data and observations from the 2001 Etna eruption. This event was characterized by
50 moderately intense explosive activity that produced a weak plume lasting for several
51 weeks. Several recent studies (*Taddeucci et al.* [2002]; *Metrich et al.* [2004]; *Taddeucci*
52 *et al.* [2004]; *Scollo et al.* [2007]) provide new data and measurements for this event, al-
53 lowing thorough comparison with VOL-CALPUFF modeling.

54 In the following section we will illustrate how application of VOL-CALPUFF to this event
55 provides new fundamental insights into the dynamics of the dispersal processes as well as
56 suggests future research needs for data collection and data intercomparison.

2. Mount Etna's 2001 Eruption

57 The explosive event we refer to started at 1900 LST on 19 July when two pit craters
58 opened on Piano del Lago, on the S flank of Mt.Etna, at about 2570 m asl (see Fig.
59 1) (*Calvari and Pinkerton* [2004]; *Scollo et al.* [2007]). The first few days of explosive
60 activity, which lasted until 6 August, were characterized by a black and dense ash and
61 lapilli column, of likely phreatomagmatic origin, rising up to 3 km above the vent (*Coltelli*
62 *et al.* [2001]; GVP/USGS Weekly Volcanic Activity Report - GVP/USGS web-site [2004];
63 *Scollo et al.* [2007]). During the first days the wind was blowing towards E-NE and the
64 city of Reggio Calabria in the Calabria Region was affected by ash fallout. Due to a
65 wind change on 21 July, the ash cloud started to move S-SE, blanketing Catania on 22
66 July and creating major problems for the "Fontanarossa" International Airport. The
67 finer portion of ash reached the Island of Malta and then the N Africa coast. The ash
68 cloud evolution was visible on several satellite images and also inferred through deposit
69 characterization studies, making this event suitable for modeling applications. VOL-

70 CALPUFF was applied to simulate the first five days (20-24 July) of this eruption (the
71 tephra fallout on 19-20 July was negligible, see *Scollo et al.* [2007]) in order to describe
72 the main features of atmospheric plume rise and ash dispersal. Modeling results are then
73 compared to available field and satellite data for a first semi-quantitative validation of
74 VOL-CALPUFF.

2.1. Simulation Input Data and Parameters

75 By using the geophysical pre-processors of the system (see *Barsotti et al.* [2007], this
76 issue) an area of 22,500 square km was defined, covering the E part of Sicily and S tip of
77 Calabria. Terrain elevation and land use data were downloaded from the USGS website
78 (<http://edc.usgs.gov>). Data spatial resolution was about 900 m, even though for visualiza-
79 tion purposes a 90 m DEM of Sicily extracted from the Eurasia SRTM3 archive (available
80 at <ftp://e0mss21u.ecs.nasa.gov/srtm/Eurasia>) was used. The computational domain was
81 discretized by 150*150 cells with a uniform grid size of 1 km. Meteorological input data
82 needed to run the processor CALMET were provided by the Servizio Idro-Meteorologico of
83 the Emilia Romagna Region (Italy), which runs the non-hydrostatic Limited Area Model
84 Italy (LAMI) (*Doms and Schättler* [1997]; COSMO's web-site [2004]) throughout the
85 Italian territory with a resolution of 0.0625 degrees, equivalent to about 7 km. CAL-
86 MET is then run to compute both the 3D wind field and 2D fields containing the values
87 of the micrometeorological variables used by VOL-CALPUFF to compute atmospheric
88 turbulence parameters (e.g. standard deviations of the horizontal and vertical wind com-
89 ponents). All meteorological fields were computed on 16 terrain-following levels varying
90 between 20 and 8800 m of altitude.

91 Finally, VOL-CALPUFF needs as input, in addition to the CALMET output, a file
92 containing data characterizing the volcanic source. These inputs include vent altitude and
93 radius, and temperature, pressure, velocity, gas (water vapor) mass fraction and grain size
94 distribution of the eruptive mixture. Table 1 summarizes all performed simulations and
95 the associated vent conditions.

96 Vent conditions were derived, as much as possible, from the existing published literature
97 and from direct observations described in technical reports. Vent altitude was relatively
98 well constrained. We assumed a value of 2570 m in all simulations performed (*INGV -*
99 *CT Staff* [2001]; *Calvari and Pinkerton* [2004]; *Scollo et al.* [2007]). Erupted mixture
100 temperature was derived from *Taddeucci et al.* [2004] and fixed at 1300 K, based on
101 residual glass composition, experimentally calibrated geothermometers and assuming no
102 significant cooling effects during magma ascent. Based on the absence of any observed
103 overpressure effects, vent flow pressure was assumed to be atmospheric in all simulations
104 (*Calvari and Pinkerton* [2004]; *Scollo et al.* [2007]). Vent radius, exit flow velocity and
105 water vapor mass fraction are more difficult to estimate through direct measurement, and
106 each parameter was varied within a specific range of values. Vent radius was varied from
107 a few meters up to about 20 m, exit velocity from 20 to 100 m/s and water content from
108 1.5 to 4.5 wt%. Adopted vent radius values (hereafter considered to approximate the
109 plume radius at the vent) are consistent with photo estimates (*Calvari and Pinkerton*
110 [2004]; *Scollo et al.* [2007]). Exit velocities in the 40 - 100 m/s range were estimated
111 by *Scollo et al.* [2007]. The adopted range of water contents (1.5 - 4.5 wt%) tries to
112 capture the significant uncertainty associated with this variable, also in light of the possible
113 interaction between the eruptive mixture and external water (the average magma water

114 content estimated by FTIR analyzes of scoriae and lapilli inclusions is about 3.4 wt%
115 (*Metrich et al.* [2004])). The relative effects of these variables, including the influence of
116 the associated mass flow rate, are discussed in the following section. Table 2 reports the
117 granulometric distribution as reconstructed by *Scollo et al.* [2007] on the basis of field
118 work carried out during the eruption and the corresponding μm -size particle distribution
119 assumed at the vent in the performed simulations. The reconstructed distribution is based
120 on the spatial integration of deposit data and does not account for the amount of fine
121 particles dispersed by wind into the far field and inaccessible to field sampling. It should
122 also be pointed out that, although VOL-CALPUFF can potentially treat time-dependent
123 flow conditions at the vent, all simulations were performed under steady-state conditions.
124 In the models, reported temporal variations of plume height and ash dispersal are only
125 caused by changing meteorological conditions.

2.2. Main Results

126 2.2.1. Plume Rise Phase

127 As already discussed in the companion paper (*Barsotti et al.* [2007], this issue) the rise
128 phase model implemented in VOL-CALPUFF reproduces the contrasting styles charac-
129 terizing an explosive eruption (collapsing, super-buoyant and buoyant) (*Woods* [1988];
130 *Bursik and Woods* [1991]). For the constant vent conditions reported in Table 1, the
131 plume is typically super-buoyant; i.e. the column accelerates upwards after an initial ve-
132 locity decrease. The simulations performed highlight the critical effect of the wind field on
133 column dynamics, not only in terms of eruptive style but also regarding the plume height
134 and dispersal direction. In the presence of a strong wind field the plume axis becomes
135 more tilted toward the ground, whereas a weak wind can generate a higher plume reaching

136 shorter downwind distances (*Sparks et al.* [1997]; *Bursik* [2001]; *Barsotti et al.* [2007], this
137 issue).

138 The effect of mass flow rate on column dynamics can be assessed from variations of plume
139 height during the simulated period taking into account the current meteorological con-
140 ditions. Fig. 2a shows the temporal variation of column height over five days for some
141 of the simulations performed. Due to the significant uncertainty affecting column height,
142 a range of mass flow rate values was chosen for sensitivity tests. The smaller mass flow
143 rates (i.e. *sim_etna1a* and *sim_etna2a*) approximately correspond to the averaged mass
144 flow rates estimated from the deposit volume (see *Scollo et al.* [2007]) when assuming a
145 constant eruption flux over 5 days. As Fig. 2a illustrates, VOL-CALPUFF, for such dis-
146 charge rates, produces buoyant columns able to rise up to only a few hundreds of meters.
147 By increasing flow rate, ascent behavior changes to super-buoyant and the column reaches
148 greater heights. For comparison, in figure 2a is also reported the approximate range of
149 observed column heights. Some chronicles mention that the column reached a peak height
150 of at least 5500 m above sea level (i.e. about 3000 m above the vent) on 23 July (*Coltelli*
151 *et al.* [2001]). This observation allows us to use our modeling results in an inverse mode,
152 and to quantify the expected mass flow rate from the column height estimation. It results
153 that an eruption characterized by a mass flow rate in the range $1.5 * 10^4$ (*sim_etna3a*)
154 - $2.5 * 10^4$ (*sim_etna4a*) kg/s can produce an eruptive column able to reach an altitude
155 comparable with that observed. Sensitivity analysis of column height as a function of
156 different velocity/radius values (at a constant mass flow rate and water content) was also
157 carried out (Fig. 2b). Simulation results clearly show how a variation in velocity of 25 -
158 100 m/s does not change the column height by more than 10%. We therefore fixed the

159 initial mixture velocity to 25 m/s in most simulations. Lastly, we investigated the effects
160 of variations in water vapor mass content (Fig. 2c). Variations up to 50% produced
161 a maximum change in column height of about 15 %, making this a significant variable
162 (*sim_etna1a-c*, *sim_etna4a-c*, *sim_etna6a-c*). In summary, the sensitivity study shows that
163 the plume mass flow rate is the main volcanological variable controlling column behavior,
164 with water content also playing a significant role. Variations in velocity/radius values
165 appear to play a secondary role.

166 2.2.2. Dispersion of the Ash Cloud and Ground Deposition

167 Once the plume height has been computed, the VOL-CALPUFF code begins to track
168 the movements of the puff centers and to calculate the puff diffusion in the defined domain.
169 The puff feeding rate corresponds to the mass flow rate feeding the rising plume at the vent
170 minus the amount of mass radially lost by the plume during its ascent. Simulation outputs
171 consist of the temporal and spatial distribution of particle concentration in air and at the
172 ground for the nine grain sizes considered (see Table 2). In this section just a few selected
173 ash dispersal simulation outputs are presented to illustrate the complex and unsteady
174 cloud dynamics. Some semi-quantitative comparisons with independent observations of
175 the phenomenon are also made. An initial quantitative description of the large-scale ash
176 cloud dispersal can be gained by plotting the concentration of the smaller particles in the
177 eruptive mixture, i.e. the 31 μm particles. These particles can be transported to large dis-
178 tances even for relatively weak events, such as the 2001 Etna eruption, (*Sparks et al.* [1997])
179 and, at the same time, they can be detected by satellite thus allowing a semi-quantitative
180 comparison of modeling results with independent satellite observations (*Rose et al.* [2001];
181 *Bluth and Rose* [2004]). Fig. 3 illustrates the evolution of spatial distribution for the total

182 mass of $31 \mu\text{m}$ particles during the four-day period for simulation *sim_etna4a* starting at
183 00UTC on 20 July (results illustrated every 12 hours). This quantity is obtained by inte-
184 grating the concentration of particles over a column of unitary base up to the maximum
185 extension of the ash plume. It allows dispersal mapping at a given time for a selected par-
186 ticle size. During the first hours of 20 July, fine ash was transported towards Calabria in
187 the NE direction as reported by direct observations. In the following hours the ash cloud
188 started moving towards the E and then S (clockwise rotation of ca. 90 degrees). During
189 22 July the ash cloud passed several times over Catania showing a series of fluctuations
190 in the S-SE directions. Finally, on 23-24 July the cloud remained mainly directed SE. An
191 initial quantitative comparison between modeling results and observations can be made
192 by using satellite images. A number of similar comparisons were carried out at different
193 times during the four days of dispersal for different mass flow rates and the agreement
194 between remote sensing observations and predictions for the finer ash distribution was
195 always a close one. Fig. 4 presents the comparison between the dispersal patterns of the
196 different granulometric classes of *sim_etna4a* (characterized by a mass flow rate of $2.5 \cdot 10^4$
197 kg/s), as computed by VOL-CALPUFF, and NOAA/AVHRR satellite image (available
198 at the web site http://www.nerc-pml.ac.uk/rsdas/projects/etna/pic_gallery.html). The
199 comparison refers to 1114UTC on 20 July. For comparison purposes, $3 \mu\text{m}$ particles were
200 also considered in the simulation to evaluate the behavior of micron-size particles. Figure
201 4 illustrates the spatial distribution of the total mass of each particle size obtained by
202 vertically integrating the concentration of each class over the entire plume height. It can
203 be seen that the 3 and $31 \mu\text{m}$ particles closely follow the dispersal pattern apparent in
204 the satellite image (E-NE), matching the main dispersal axis and the lateral extension of

205 the detected plume, whereas all the other particles are transported in different directions
206 not detected by satellite. In particular, particles of 62 and 125 μm seem to move firstly
207 towards E and then slightly deviate towards the S-SE direction. Similarly, the larger re-
208 ported particles (250 μm and 2 mm) show a totally independent movement, characterized,
209 from the beginning, by a northward displacement.

210 In order to better understand the reason for such a complex dispersal pattern some
211 features of the wind field around the volcano were investigated. Fig. 5 shows the wind
212 field streamlines at the levels of 400 and 2200 m a.g.l computed by the meteorological
213 pre-processor CALMET for the acquisition time of the satellite image shown in Fig. 4.
214 As expected, at levels closer to the ground the wind is strongly affected by the volcano's
215 topography thus generating a more irregular field, whereas at higher altitudes the wind
216 shows a quasi-geostrophic trend.

217 A more complete picture of the wind field spatial and temporal variability, as produced
218 by the CALMET code, can be gained by analysing the wind direction and speed at
219 different altitudes. Fig. 6 describes variations of wind direction and speed over four days
220 (21-24 July) above the vent area, 15 km E and 15 km S from the vent, at two different
221 altitudes above the ground (400 and 2200 m). A key result is the strongly irregular
222 patterns of the flow field close to the ground (400 m) (see Figs. 6a and 6c), above the
223 vent and 15 km S and E. This irregularity is particularly pronounced for wind direction.
224 At the vent and at a location 15 km E of it, wind direction varies by 90-130 degrees over
225 the four days (blowing towards W above the vent and towards E some 15 km E from it).
226 These variations become even more dramatic 15 km S of the vent where three inversions of
227 180 degrees, between the evening of the 23 and the evening of the 24 July, are predicted.

228 In contrast, at 2200 m a.g.l the wind is blowing from NW to SE on almost any day of
229 the investigated period with a speed on the first day about five times greater than that
230 at lower levels. Wind speed at high altitude remains twice near-ground wind speed on all
231 other days.

232 VOL-CALPUFF can be used to estimate ground deposition as a function of assumed
233 particle size distribution and mass flow rate. Fig. 7 illustrates particle deposition for six
234 of the nine granulometric sizes employed in the *sim_etna4a* simulation. Deposition refers
235 to total ash accumulation over four days, between 21 and 24 July. We considered this
236 shorter period because of the availability of deposit measurements (*Scollo et al.* [2007])
237 which enable the comparison of model predictions with deposit data (see hereafter). It
238 is evident from the figure how the various particles disperse and deposit following very
239 different dynamics. On figure 7, the smaller particles (e.g. $31\ \mu\text{m}$) are mostly transported
240 in the direction of the prevailing high-altitude winds and remain confined in the SE
241 quadrant. Intense wind quickly advects these low settling velocity particles out of the
242 domain, so that the deposit iso-contours show values up to several orders of magnitude
243 smaller than those of the other particles sizes. In contrast, for increasingly large particles
244 the location of maximum deposition moves towards the vent, as expected. $500\ \mu\text{m}$ and
245 $1\ \text{mm}$ particles are much less dispersed than finer particles and their deposit is more
246 uniform in the azimuthal direction. Similar patterns were obtained for the 2, 4 and 8 mm
247 particles.

248

249 By summing the partial deposits produced by the nine different particle sizes it is
250 possible to compute the total particle accumulation. Fig. 8a and 8c show the mass isolines

251 of ash deposited after four days of activity for the two mass flow rates (i.e. 1.5 and $2.5 \cdot 10^4$
252 kg/s) suggested as end-member values for the event investigated (*sim_etna3a*, *sim_etna4a*).
253 Fig. 8b and 8d show a comparison of the two computed deposition patterns with the
254 deposit measurements, in kg/m^2 , as reported by *Scollo et al.* [2007]. The total deposit
255 isolines show a quite irregular and non-symmetric pattern in both cases. Both results
256 present a deposit characterized by a preferred E-SE direction, due to higher columns that
257 inject material at atmospheric levels characterized by intense wind, which, as we have
258 already seen (Fig. 6b), blew most of the time towards this quadrant. The predicted
259 deposit distributions show that maximum deposition is at a few kilometers from the vent
260 in the S-SE direction, as in the measured deposit (*Scollo et al.* [2007]). The lower end-
261 member mass flow rate (Fig. 8a and Fig. 8b - *sim_etna3a*) produces a deposit which
262 fits the sampled data with a correlation coefficient of 0.77 and with more than 97% of
263 simulated data ranging from 1/5 to 5 times the observed values. For the larger end-
264 member mass flow rate (Fig. 8c and Fig. 8d - *sim_etna4a*) the numerical prediction
265 largely overestimates the amount of ash deposited. A comparison between computed and
266 measured values quantifies this discrepancy to be up to about a factor of 5.

3. Discussion

267 The above results show that VOL-CALPUFF is helpful to investigate ash dispersal
268 processes. The new model captures the strong influence of meteorological conditions on
269 plume rise dynamics, shows the necessity to use codes able to determine column height
270 using a plume rise model and the importance of using an accurate meteorological data
271 set. A useful result that can be inferred *a posteriori* from modeling of column height in
272 the presence of realistic meteorological conditions is the estimation of the mass flow rate

273 feeding the plume. For the 2001 Etna event, the results reported in Fig. 2 illustrate that
274 only mass flow rates between $1.5 * 10^4$ kg/s (*sim_etna3a*) and $2.5 * 10^4$ kg/s (*sim_etna4a*)
275 are likely able to produce a maximum column height of about 3000 m above the vent on
276 23 July. Based on the analysis reported in the companion paper (see Fig. 4 of *Barsotti*
277 *et al.* [2007], this issue) on the effect of the entrainment coefficients and the values here
278 assumed, these values should be considered as minimum values in order to reproduce the
279 observed column height. In contrast, using an *averaged* mass flow rate, obtained by di-
280 viding the total amount of mass erupted ($1.02 * 10^9$ and $2.31 * 10^9$ kg from *Scollo et al.*
281 [2007]) by event duration (*sim_etna1a* and *sim_etna2a*), it would not be possible to match
282 the observed column height. Only in a still environment, the plume could reach altitudes
283 comparable with observed ones if fed with averaged mass flow rates (as assumed in *Costa*
284 *et al.* [2006]). For the long-lasting weak plume considered here, the assumption of an aver-
285 age mass flow rate is inadequate to simulate the whole dispersal process, since it strongly
286 underestimates column height. However the range of mass flow rate estimated above from
287 the column height (*sim_etna3a* and *sim_etna4a*) is likely representative of the *peak* mass
288 flow rate. This could significantly overestimate the erupted volume and ash ground de-
289 position. If we consider the column height and deposit estimates, we see that each of the
290 two end-member values matches only one of the two observed variables. Plume feeding
291 with a steady mass flow rate of $1.5 * 10^4$ kg/s allows a match with the amount of material
292 deposited (see Fig. 8b), but the predicted column height only approximately matches the
293 minimum observed values (see Fig. 2a). By feeding the plume with the higher mass flow
294 rate (i.e. $2.5 * 10^4$ kg/s) it is possible to match the observed plume height but not deposit
295 estimates (Fig. 8d). A reasonable explanation for this discrepancy is that events lasting

296 several days, such as the 2001 Etna eruption, are characterized by unsteady, intermittent
297 plume feeding (*INGV - CT Staff* [2001]; *Taddeucci et al.* [2004]; *Scollo et al.* [2007], P. Al-
298 lard written communication (2007)). Plume-feeding is therefore quite difficult to quantify
299 and sometimes also hard to detect (for instance during nocturnal periods or bad weather
300 conditions).

301 We believe that a better match between modeling results and observations cannot be
302 achieved under the assumption of steady-state plume feeding over the investigated four-
303 day period. The few available measurements are affected by significant uncertainty and
304 do not allow more accurate comparison or the definition of time-dependent plume feeding
305 over the simulated period. In order to support this interpretation we performed addi-
306 tional simulations adopting two different mass flow rates for the plume and puffs. We ran
307 VOL-CALPUFF assuming a mass flow rate of $2.5 \cdot 10^4$ and $1.0 \cdot 10^4$ kg/s for the plume
308 and puff release, respectively. In this way we were able to match the observed column
309 height and produce a deposit distribution (not shown) that fits the measured values with
310 a correlation coefficient ($R=0.82$) higher than that for the simulations reported in Fig. 8.
311 Another variable that could significantly reduce the above discussed discrepancy is the
312 gas mass fraction of the eruptive mixture at the vent. Water contents greater than those
313 assumed here could favor higher columns for a given mass flow rate (Fig. 2c).

314 VOL-CALPUFF allows analysis of new and non-intuitive effects of ash dispersal in an
315 orographically complex area under the action of 3D meteorological conditions. These
316 effects are apparent in the dispersal behavior of contrasting particles sizes. The major
317 mass *splitting* effect we observe in Fig. 4 and Fig. 7 between particles of different sizes
318 is the result of two main factors: (a) the dependence of settling velocities on particle

319 dimension and density and (b) the great variability and complexity of wind field patterns
320 in the vertical and horizontal directions. These factors make larger particles cross the
321 atmosphere in a temporal interval up to two orders of magnitude faster than the smaller
322 particles (minutes vs. hours), travelling faster towards the ground where wind fields are
323 more irregular. Particles up to a few tens of microns remain much longer at high at-
324 mospheric levels, where more uniform winds take them far from the vent. These results
325 are in agreement with the very first attempt to investigate the control of wind-shear and
326 grain size distribution on ash dispersal by *Woods et al.* [1995]. In terms of validation of
327 the predicted particle dispersal behavior, an initial comparison of simulation outputs with
328 satellite images was carried out. Results suggest that VOL-CALPUFF is able to correctly
329 predict the dispersal of the smaller particles moving at high altitude above the ground.
330 The path observed in satellite images was closely followed by particles up to about $30\ \mu\text{m}$
331 in our simulations. In contrast, the AVHRR satellite images taken by the NOAA satel-
332 lite do not show the more irregular behavior that the simulations produce for particles
333 larger than about $30\ \mu\text{m}$. This result is consistent with the present limitations of satellite
334 remote sensing techniques in detecting suspended material with a radius greater than a
335 few tens of microns (*Wen and Rose* [1994]; *Yu and Rose* [2002]; *Bluth and Rose* [2004]).
336 Given the importance, in terms of mass dispersed in the atmosphere, of particles with
337 size in the range $62\ \mu\text{m}$ -1 mm (see Table 2), an important contribution to the detection
338 of these coarse particles should come from the use of other remote sensing techniques
339 such as weather radar observation (*Lacasse et al.* [2004]; *Marzano and Vulpiani* [2006]).
340 Unfortunately, such data are not available for the 2001 Etna event and cannot be used
341 to validate the numerical predictions. This type of investigation appears particularly im-

342 portant for particle sizes (e.g. $125 \mu\text{m}$) that can reach concentrations of a few g/m^3 far
343 away from the area indicated by the satellite images. Such concentrations are well above
344 those used by other ash forecasting codes as a threshold for ash cloud visibility (*Simpson*
345 *et al.* [2002]). Qualitative support for the modeling results here presented comes from
346 the observation that ash fallout of coarse particles was actually recorded several times,
347 during the 2001 event, well away from the main plume dispersal direction (Andronico D.,
348 personal communication, 2005).

349 From the above discussion the fundamental role played by meteorology in the determi-
350 nation of both plume rise and ash dispersal is evident. A better understanding of the
351 minimum resolution of meteorological data required to capture, with sufficient accuracy,
352 most of the results illustrated here, is necessary for future progress. The present modeling
353 results could also contribute towards a more quantitative interpretation of remote sensing
354 and field data.

4. Conclusions

355 The first application of VOL-CALPUFF, to the 2001 explosive event of Mt. Etna,
356 highlights new key features of the dynamics of weak volcanic plumes in the presence of
357 a complex wind field. Some of the most important aspects are: (a) the major effect
358 of the wind field in the determination of the column height and bending, as well as of
359 plume ascent style, (b) the great difference in the dispersal patterns followed by ashes
360 of different sizes producing a very complex and non-intuitive dispersal, (c) the complex
361 distribution of ash on the ground showing preferential directions of dispersal and, for
362 events characterized by long duration and variable winds, a quite irregular deposit pattern,
363 (d) the impossibility to reproduce both column height and deposit features by adopting

364 a vent steady-state mass flow rate due to temporal variations in plume feeding during
365 the four-day period investigated. Furthermore, this initial application of VOL-CALPUFF
366 to weak plumes poses new questions about the use of low-resolution meteorological data
367 for the accurate prediction of the complex ash dispersal and fallout dynamics. Finally, it
368 highlights, once more, the importance of comparing modeling results to accurate field data
369 and observations in order to reach a consistent understanding of the dispersal dynamics
370 and also prompts new developments in modeling and observational techniques.

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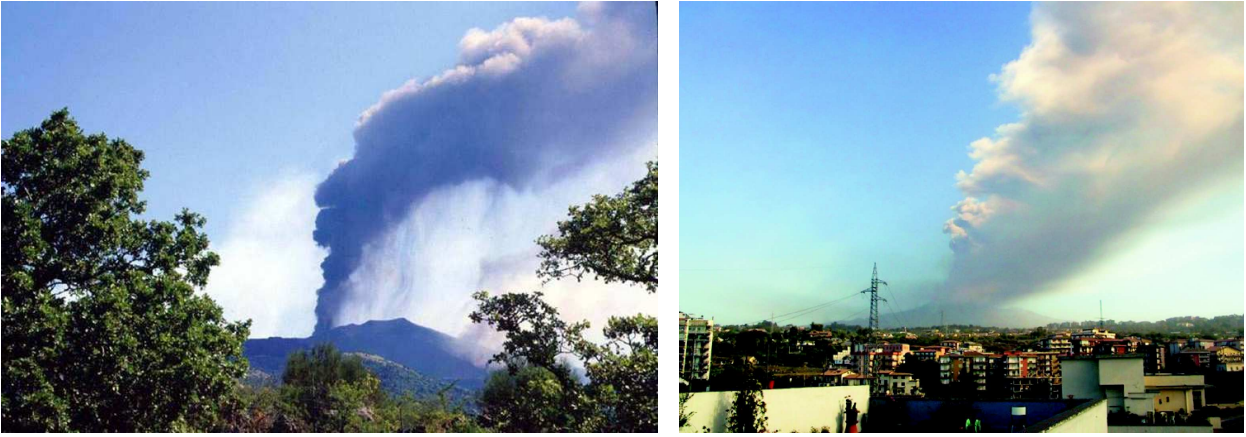


Figure 1. Photographs of the volcanic plume and ash fallout over Catania and nearby villages, during the recent eruptive events at Mt.Etna. Photos by P. Papale (INGV Pisa) and from www.larepubblica.it, respectively.

Table 1. List of performed simulations. The table contains the input data values assumed at the vent. V is velocity, R vent radius, H_2O the weight percent of water vapor and mfr the resulting mass flow rate. Vent radius was derived from the other three variables by using the mass conservation equation. Vent altitude, vent flow temperature and exit mixture pressure were kept constant in all simulations and are equal to 2570 m, 1300 K and atmospheric value, respectively. As explained in *Barsotti et al.* [2007] (this issue) all the simulations were performed using entrainment coefficients values (α and γ) equal to 0.09 and 0.6, respectively.

Name	V (m/s)	R (m)	H_2O (wt%)	mfr (kg/s)
<i>sim_etna1a</i>	20	2.7	3	$2.0 * 10^3$
<i>sim_etna1b</i>	30	2.7	4.5	$2.0 * 10^3$
<i>sim_etna1c</i>	25	1.7	1.5	$2.0 * 10^3$
<i>sim_etna2a</i>	20	4.8	3	$6.6 * 10^3$
<i>sim_etna3a</i>	25	6.7	3	$1.5 * 10^4$
<i>sim_etna4a</i>	25	8.7	3	$2.5 * 10^4$
<i>sim_etna4b</i>	30	9.7	4.5	$2.5 * 10^4$
<i>sim_etna4c</i>	25	6.2	1.5	$2.5 * 10^4$
<i>sim_etna4d</i>	50	6.2	3	$2.5 * 10^4$
<i>sim_etna4e</i>	75	5	3	$2.5 * 10^4$
<i>sim_etna4f</i>	100	4.3	3	$2.5 * 10^4$
<i>sim_etna5a</i>	25	11	3	$4.0 * 10^4$
<i>sim_etna6a</i>	25	15.1	3	$7.5 * 10^4$
<i>sim_etna6b</i>	30	16.7	4.5	$7.5 * 10^4$

Name	V (m/s)	R (m)	H_2O (wt%)	mfr (kg/s)
<i>sim_etna6c</i>	25	10.6	1.5	$7.5 * 10^4$
<i>sim_etna6d</i>	50	10.6	3	$7.5 * 10^4$
<i>sim_etna6e</i>	75	8.7	3	$7.5 * 10^4$
<i>sim_etna6f</i>	100	7.5	3	$7.5 * 10^4$
<i>sim_etna7a</i>	25	19.86	3	$1.3 * 10^5$

Table 2. The particle size distribution of the eruptive mixture assumed at the vent in the simulations performed as coming from the reconstructed total granulometric distribution (*Scollo et al.* [2007]).

Φ	Diameter (μm)	wt%
-3	8000	0.11
-2	4000	0.69
-1	2000	2.47
0	1000	7.53
1	500	20.36
2	250	40.2
3	125	24.06
4	62	3.89
5	31	0.69

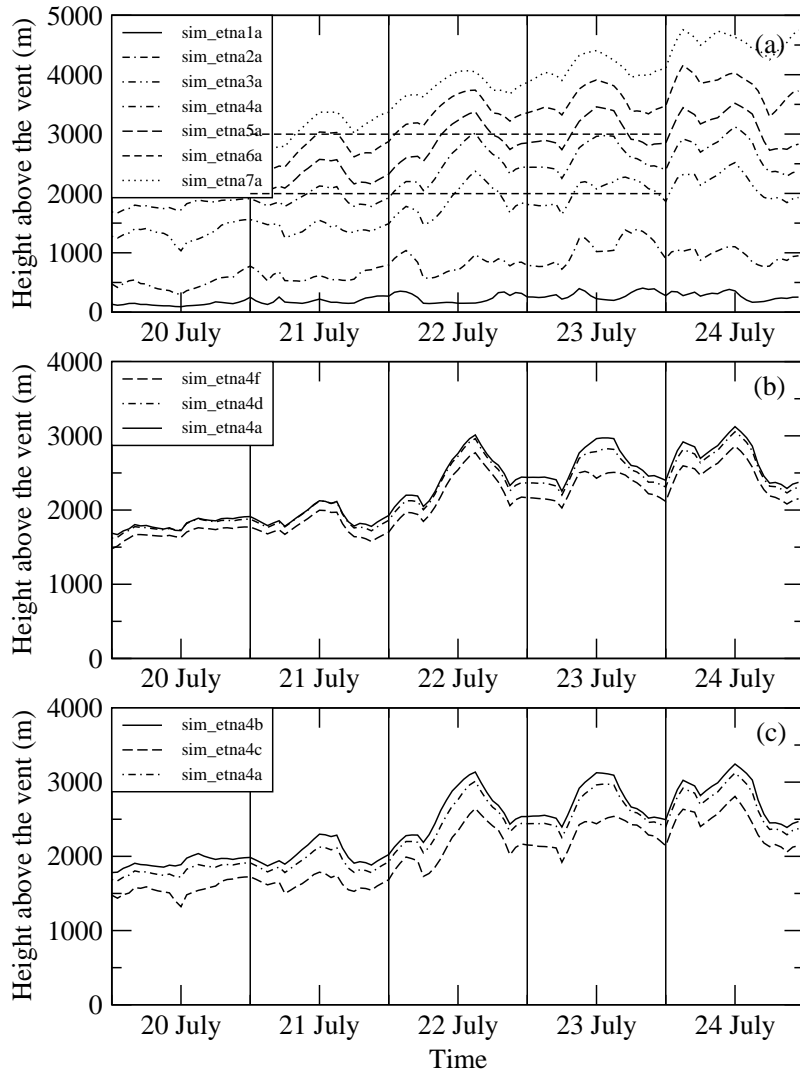


Figure 2. Temporal variations of column height as computed by VOL-CALPUFF for some of the simulations reported in Table 1. (a) Comparison between column heights computed for plumes with different mass flow rates. The vent mass flow rate is kept constant in time during each simulation. The superimposed horizontal lines delimit the observed range for the 21-23 July period (*Scollo et al.* [2007]); (b) Comparison between column heights computed for plumes with different mixture velocities (at constant mass flow rate and water content); (c) Comparison between column heights computed for plumes with different water vapor mass fractions.

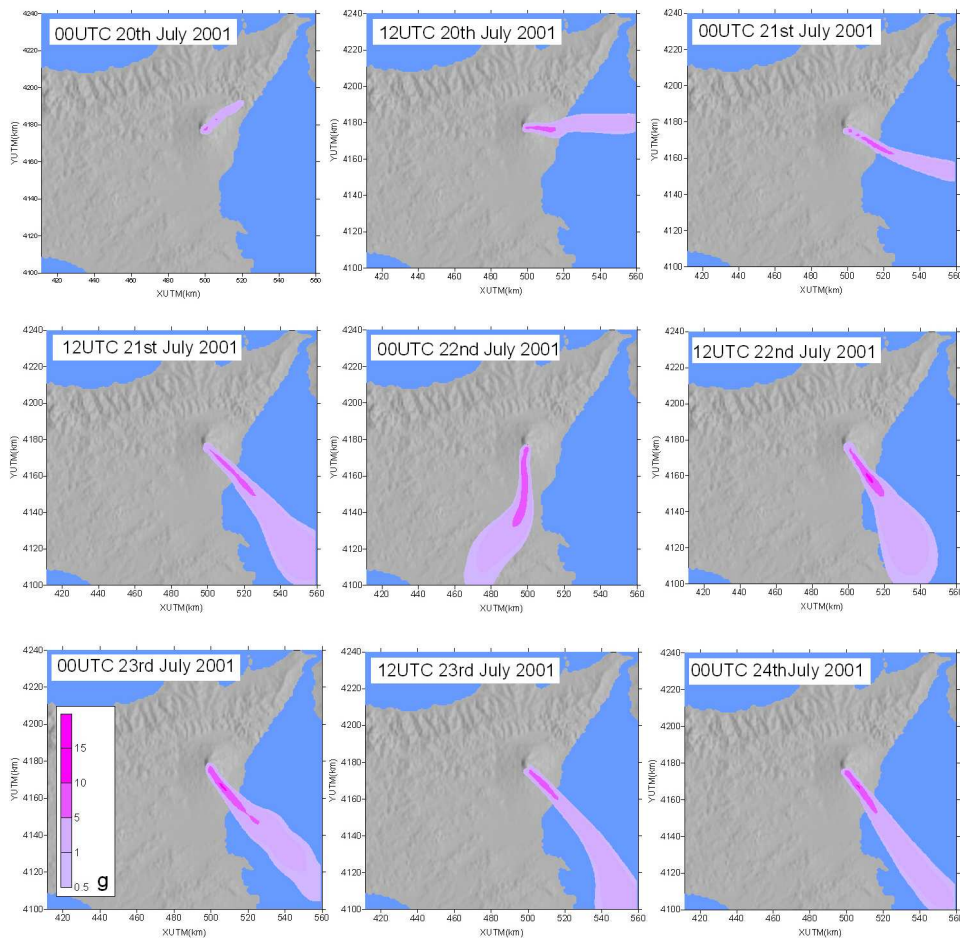


Figure 3. Cumulative mass of $31 \mu\text{m}$ particles, obtained by integration of its concentration over all vertical levels in the atmosphere. From the upper left-hand corner the concentration is mapped each 12 hours starting at 00UTC on 20 July. Isocontour values are reported on the bottom left-hand corner and refer to the total mass in grams. Results from simulation *sim_etna4a*.

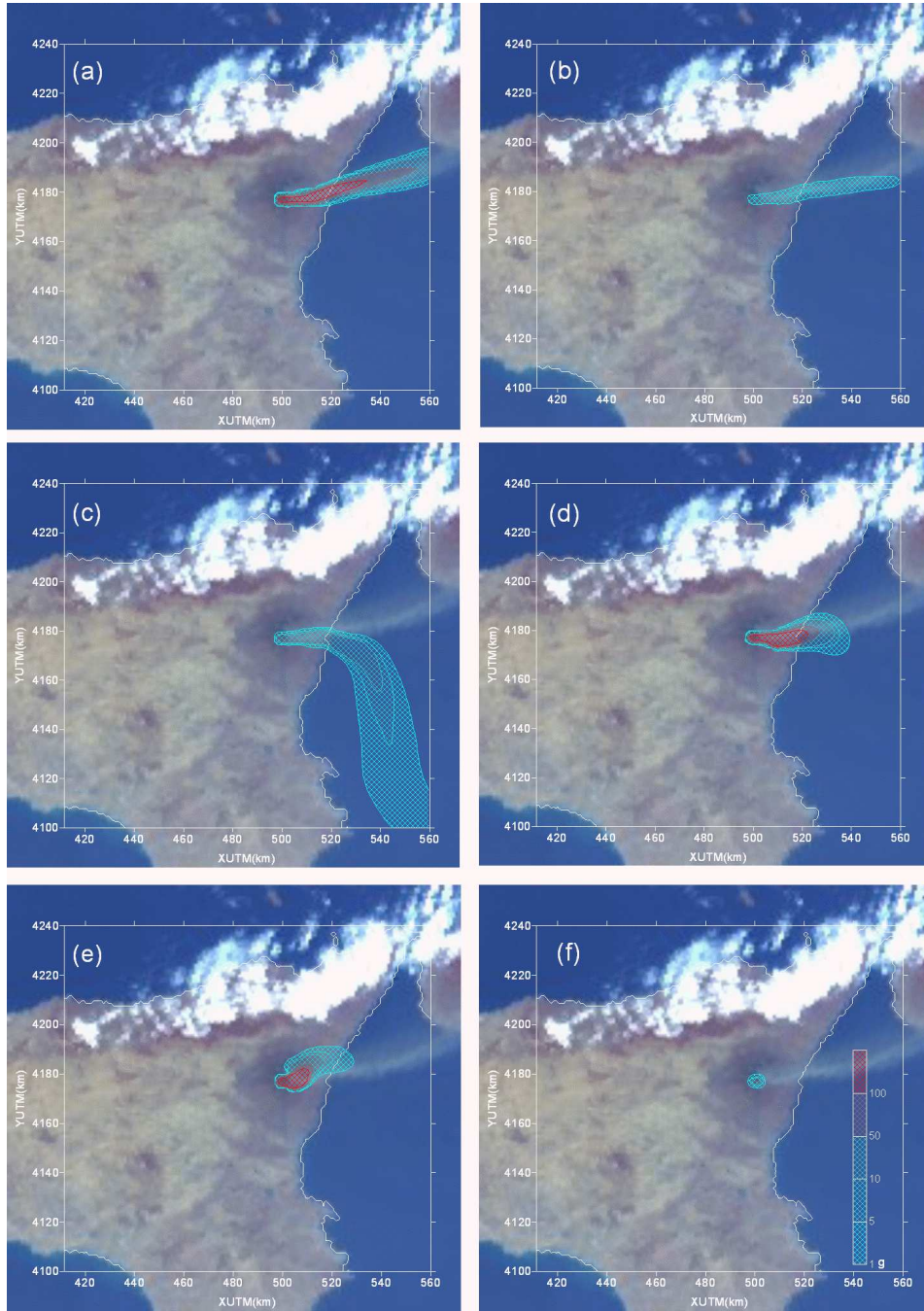


Figure 4. SeaStar satellite image at 1114UTC on 20 July compared with *sim_etna4a* results. The figures show the amount of ash integrated in the vertical extension of the domain for six different classes (3 (a), 31 (b), 62 (c), 125 (d), 250 (e) μm and 2 mm (f)). In all images the minimum reported isocontour is equal to 1 g. The comparison shows how the detection of the plume by the satellite image is limited to the detection of particles of a few- μm (*Bluth and Rose [2004]*).

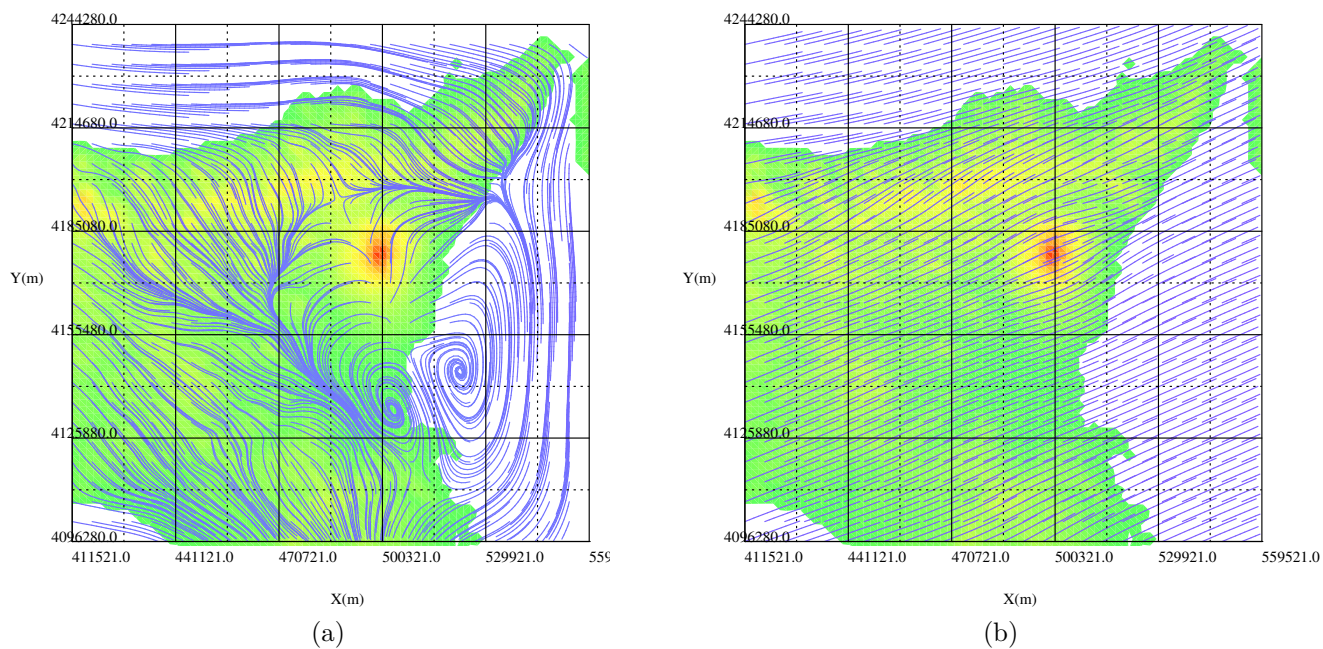


Figure 5. Wind field streamlines, as produced by CALMET model, at 400 m (a) and 2200 m (b) above ground level at 1100UTC on 20 July.

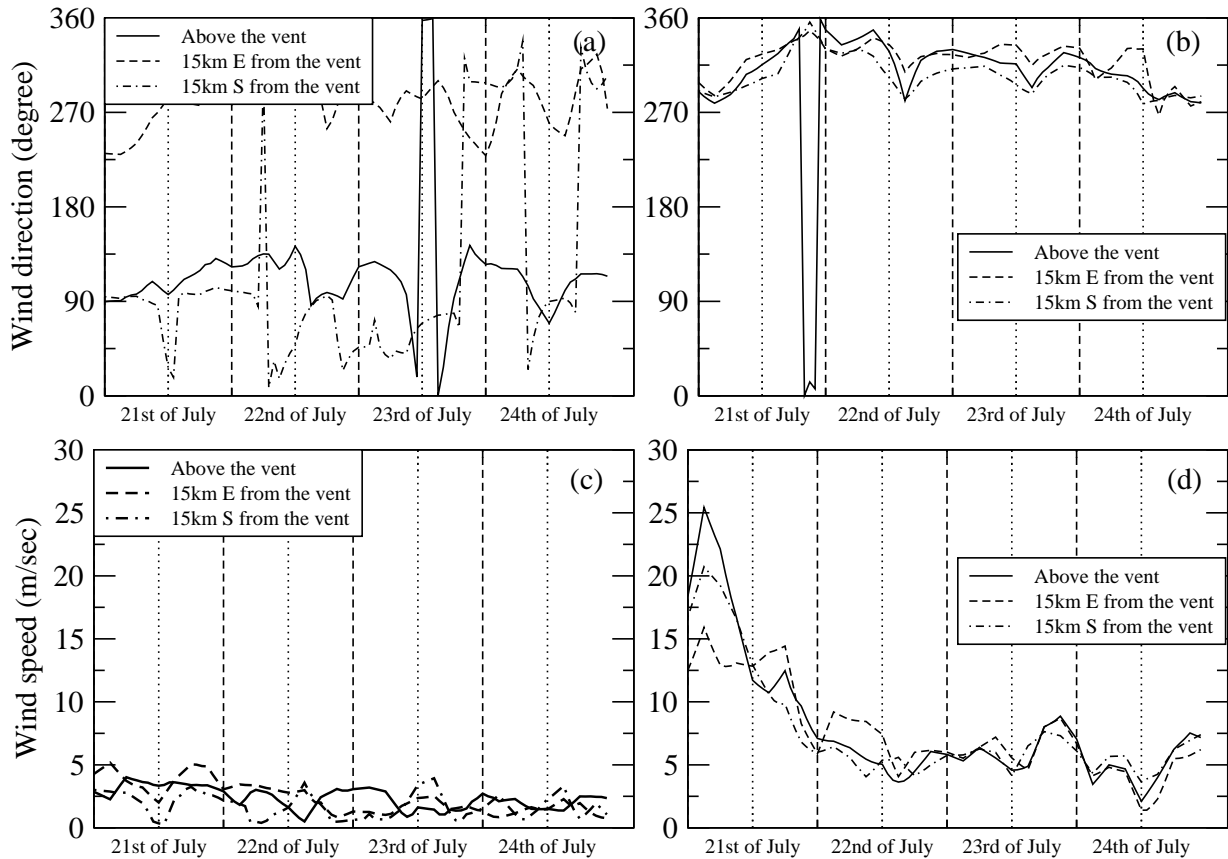


Figure 6. Temporal variations of wind direction (a and b) and wind speed (c and d) at different spatial locations. The wind information, produced by the CALMET model, is reported hourly for three different locations: above the vent, 15 km from the vent toward E or S. Wind information related to a terrain-following vertical level at 400 m agl (a and c) are compared with the same informations obtained for a 2200 m terrain-following level (b and d). The figures illustrate the effects on low atmosphere wind direction caused by complex orography (the effect of eruption dynamics on local atmospheric conditions is not considered by the model). At a higher level the wind field is more regular and faster than closer to the ground.

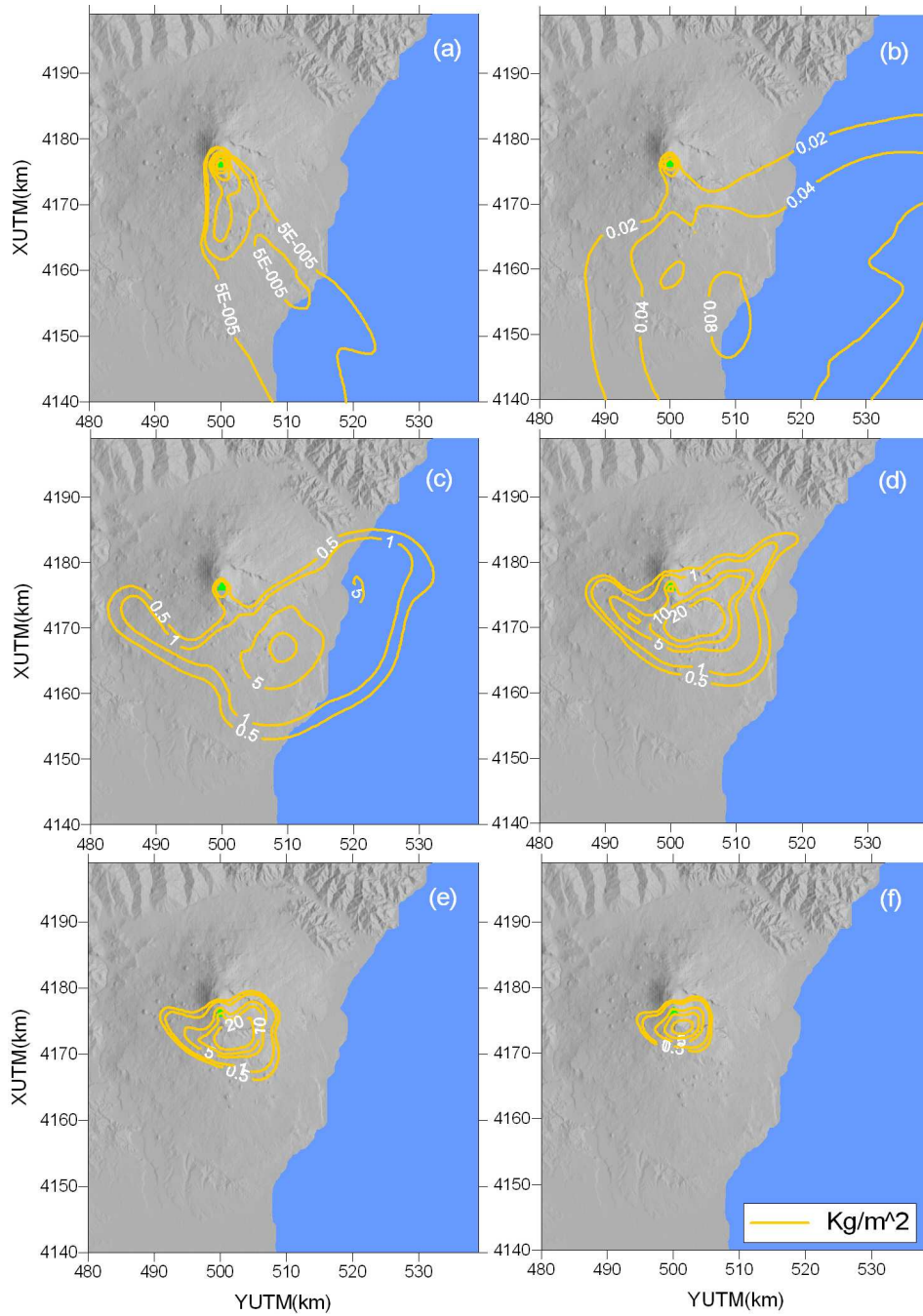


Figure 7. Ground deposit as a function of particle diameter for an emission rate of $2.5 \cdot 10^4$ kg/s (*sim_etna4a*) over four days (21-24 July). Six classes with diameter of (a) $31 \mu\text{m}$, (b) $62 \mu\text{m}$, (c) $125 \mu\text{m}$, (d) $250 \mu\text{m}$, (e) $500 \mu\text{m}$ and (f) 1 mm are reported. Particles with larger diameters show a behavior similar to the 1 mm particles.

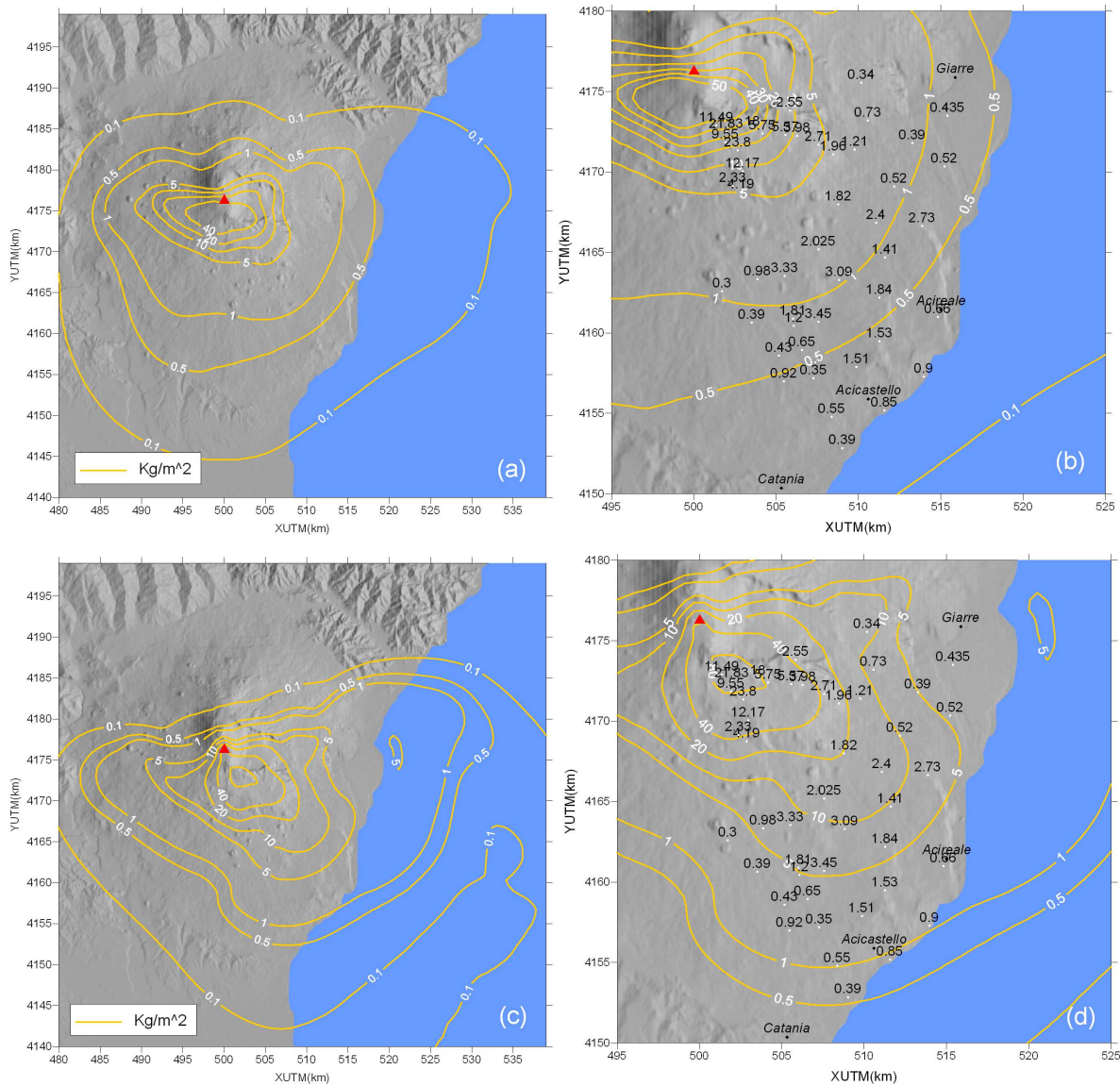


Figure 8. Total accumulation on the ground of all particle sizes after four days (21-24 July) of steady ash emission. A comparison between two different mass flow rates, i.e. $1.5 \times 10^4 \text{ kg/s}$ (*sim_etna3a*) (a, b) and $2.5 \times 10^4 \text{ kg/s}$ (*sim_etna4a*) (c, d), is shown. Both deposit patterns are characterized by an irregular shape with a peak at some kms from the vent ((a) and (c)). A zoomed comparison with measured field data (from Scollo *et al.* [2007]) highlights a good agreement for the case *sim_etna3a* (correlation coefficient $R = 0.77$) (b), and an overestimation up to about a factor of 5 for the case *sim_etna4a* (d).