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The VOL-CALPUFF Model for Atmospheric Ash Dispersal: II. Application to the Weak Mt. Etna Plume of July 2001

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Abstract. The application of the VOL-CALPUFF model (Barsotti et al. 5 [2007], this issue) to a weak plume erupted from Mt. Etna in July 2001 is 6 here presented and discussed. The reconstruction of the explosive event was 7 obtained by using high-resolution weather forecasts, produced by a mesoscale 8 non-hydrostatic model, and volcanic source data coming from observations 9 and analytical studies. The plume rise and atmospheric dispersal models were 10 investigated over five days of eruption mostly in terms of column height, aer-11 ial ash concentration and ground deposition. Modeling results are shown as 12 a function of various source conditions and compared to independent obser-13 vations derived from satellite images and deposit mapping. The application 14 of VOL-CALPUFF clearly highlights the crucial role played by meteorolog-15 ical conditions in determining dispersal dynamics. Some of the most impor-16 tant effects described by the model are: a) the large wind field influence on 17 the plume height determination and tilting, b) the contrasting dispersal pat-18 terns of ash particles of different sizes, c) the complex and somehow non-intuitive 19 distribution of ash on the ground resulting in preferential directions of dis-20 persal and quite irregular deposit patterns, d) the impossibility to reproduce 21 both the column height and the deposit accumulation pattern by adopting 22 a steady-state vent mass flow rate over the investigated four-day period due 23 to observed temporal changes in eruption dynamics. Modeling results also 24 suggest the need for further integration of simulation outcomes with remote 25 sensing and field reconstructions on ash dispersal processes in future. 26

1. Introduction

We present the first application of the VOL-CALPUFF model, fully described in the 27 companion paper (*Barsotti et al.* [2007], this issue), to a real event in order to highlight its 28 potentialities and weaknesses. The code originates from the U.S. non-commercial model 29 CALPUFF (Scire et al. [2000]), developed and used for air quality applications, modified 30 and developed to make it appropriate for volcanological applications. VOL-CALPUFF 31 is an hybrid code in which a part of the process, i.e. the plume rise phase, is described 32 in a Eulerian framework, whereas dispersal and fallout processes are solved with a La-33 grangian approach. Both processes are modelled taking into account the influence of 34 time-dependent, 3D meteorological conditions throughout the duration of the eruption. 35 VOL-CALPUFF relates the computation of the column height, computed by solving the 36 plume theory equations, to the altitude of release of pyroclastic material. The Lagrangian 37 description of the dispersal process guarantees model reliability both close to vent and 38 far from it (Nguyen et al. [1997]). The original CALPUFF code was validated through 39 extensive comparison with other widely used codes for air-pollution modeling and through experimental tests concerning passive and active tracer transport and estimation of ground 41 deposition and atmospheric concentrations (EPA U.S. [1998]; Scire et al. [2000]). Due 42 to its good performance it is currently proposed by the U. S. EPA as a guideline model 43 for environmental applications such as impact studies of proposed and existing pollutant 44 sources. Similarly, VOL-CALPUFF needs to be validated and tested against volcanolog-45 ical applications, and this is a challenge due to the uncertainties when characterizing the 46 natural phenomena. 47

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The paper presents quantitative comparisons between model results and experimental data and observations from the 2001 Etna eruption. This event was characterized by moderately intense explosive activity that produced a weak plume lasting for several weeks. Several recent studies (*Taddeucci et al.* [2002]; *Metrich et al.* [2004]; *Taddeucci et al.* [2004]; *Scollo et al.* [2007]) provide new data and measurements for this event, allowing thorough comparison with VOL-CALPUFF modeling.

In the following section we will illustrate how application of VOL-CALPUFF to this event provides new fundamental insights into the dynamics of the dispersal processes as well as suggests future research needs for data collection and data intercomparison.

2. Mount Etna's 2001 Eruption

The explosive event we refer to started at 1900 LST on 19 July when two pit craters 57 opened on Piano del Lago, on the S flank of Mt.Etna, at about 2570 m asl (see Fig. 58 1) (Calvari and Pinkerton [2004]; Scollo et al. [2007]). The first few days of explosive 59 activity, which lasted until 6 August, were characterized by a black and dense ash and 60 lapilli column, of likely phreatomagmatic origin, rising up to 3 km above the vent (Coltelli 61 et al. [2001]; GVP/USGS Weekly Volcanic Activity Report - GVP/USGS web-site [2004]; 62 Scollo et al. [2007]). During the first days the wind was blowing towards E-NE and the 63 city of Reggio Calabria in the Calabria Region was affected by ash fallout. Due to a 64 wind change on 21 July, the ash cloud started to move S-SE, blanketing Catania on 22 65 July and creating major problems for the "Fontanarossa" International Airport. The 66 finer portion of ash reached the Island of Malta and then the N Africa coast. The ash 67 cloud evolution was visible on several satellite images and also inferred through deposit 68 characterization studies, making this event suitable for modeling applications. VOL-69

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CALPUFF was applied to simulate the first five days (20-24 July) of this eruption (the tephra fallout on 19-20 July was negligible, see *Scollo et al.* [2007]) in order to describe the main features of atmospheric plume rise and ash dispersal. Modeling results are then compared to available field and satellite data for a first semi-quantitative validation of VOL-CALPUFF.

2.1. Simulation Input Data and Parameters

By using the geophysical pre-processors of the system (see *Barsotti et al.* [2007], this 75 issue) an area of 22,500 square km was defined, covering the E part of Sicily and S tip of 76 Calabria. Terrain elevation and land use data were downloaded from the USGS website 77 (http://edc.usgs.gov). Data spatial resolution was about 900 m, even though for visualiza-78 tion purposes a 90 m DEM of Sicily extracted from the Eurasia SRTM3 archive (available 79 at ftp://e0mss21u.ecs.nasa.gov/srtm/Eurasia) was used. The computational domain was 80 discretized by 150*150 cells with a uniform grid size of 1 km. Meteorological input data 81 needed to run the processor CALMET were provided by the Servizio Idro-Meteorologico of 82 the Emilia Romagna Region (Italy), which runs the non-hydrostatic Limited Area Model 83 Italy (LAMI) (Doms and Schättler [1997]; COSMO's web-site [2004]) throughout the 84 Italian territory with a resolution of 0.0625 degrees, equivalent to about 7 km. CAL-85 MET is then run to compute both the 3D wind field and 2D fields containing the values 86 of the micrometeorological variables used by VOL-CALPUFF to compute atmospheric 87 turbulence parameters (e.g. standard deviations of the horizontal and vertical wind com-88 ponents). All meteorological fields were computed on 16 terrain-following levels varying 89 between 20 and 8800 m of altitude. 90

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Finally, VOL-CALPUFF needs as input, in addition to the CALMET output, a file containing data characterizing the volcanic source. These inputs include vent altitude and radius, and temperature, pressure, velocity, gas (water vapor) mass fraction and grain size distribution of the eruptive mixture. Table 1 summarizes all performed simulations and the associated vent conditions.

Vent conditions were derived, as much as possible, from the existing published literature QF and from direct observations described in technical reports. Vent altitude was relatively 97 well constrained. We assumed a value of 2570 m in all simulations performed (INGV -98 CT Staff [2001]; Calvari and Pinkerton [2004]; Scollo et al. [2007]). Erupted mixture 90 temperature was derived from *Taddeucci et al.* [2004] and fixed at 1300 K, based on 100 residual glass composition, experimentally calibrated geothermometers and assuming no 101 significant cooling effects during magma ascent. Based on the absence of any observed 102 overpressure effects, vent flow pressure was assumed to be atmospheric in all simulations 103 (Calvari and Pinkerton [2004]; Scollo et al. [2007]). Vent radius, exit flow velocity and 104 water vapor mass fraction are more difficult to estimate through direct measurement, and 105 each parameter was varied within a specific range of values. Vent radius was varied from 106 a few meters up to about 20 m, exit velocity from 20 to 100 m/s and water content from 107 1.5 to 4.5 wt%. Adopted vent radius values (hereafter considered to approximate the 108 plume radius at the vent) are consistent with photo estimates (Calvari and Pinkerton 109 [2004]; Scollo et al. [2007]). Exit velocities in the 40 - 100 m/s range were estimated 110 by Scollo et al. [2007]. The adopted range of water contents (1.5 - 4.5 wt%) tries to 111 capture the significant uncertainty associated with this variable, also in light of the possible 112 interaction between the eruptive mixture and external water (the average magma water 113

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content estimated by FTIR analyzes of scoriae and lapilli inclusions is about 3.4 wt%114 (Metrich et al. [2004])). The relative effects of these variables, including the influence of 115 the associated mass flow rate, are discussed in the following section. Table 2 reports the 116 granulometric distribution as reconstructed by Scollo et al. [2007] on the basis of field 117 work carried out during the eruption and the corresponding μ m-size particle distribution 118 assumed at the vent in the performed simulations. The reconstructed distribution is based 119 on the spatial integration of deposit data and does not account for the amount of fine 120 particles dispersed by wind into the far field and inaccessible to field sampling. It should 121 also be pointed out that, although VOL-CALPUFF can potentially treat time-dependent 122 flow conditions at the vent, all simulations were performed under steady-state conditions. 123 In the models, reported temporal variations of plume height and ash dispersal are only 124 caused by changing meteorological conditions. 125

2.2. Main Results

¹²⁶ 2.2.1. Plume Rise Phase

As already discussed in the companion paper (*Barsotti et al.* [2007], this issue) the rise 127 phase model implemented in VOL-CALPUFF reproduces the contrasting styles charac-128 terizing an explosive eruption (collapsing, super-buoyant and buoyant) (Woods [1988]; 129 Bursik and Woods [1991]). For the constant vent conditions reported in Table 1, the 130 plume is typically super-buoyant; i.e. the column accelerates upwards after an initial ve-131 locity decrease. The simulations performed highlight the critical effect of the wind field on 132 column dynamics, not only in terms of eruptive style but also regarding the plume height 133 and dispersal direction. In the presence of a strong wind field the plume axis becomes 134 more tilted toward the ground, whereas a weak wind can generate a higher plume reaching 135

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¹³⁶ shorter downwind distances (Sparks et al. [1997]; Bursik [2001]; Barsotti et al. [2007], this
¹³⁷ issue).

The effect of mass flow rate on column dynamics can be assessed from variations of plume 138 height during the simulated period taking into account the current meteorological con-139 ditions. Fig. 2a shows the temporal variation of column height over five days for some 140 of the simulations performed. Due to the significant uncertainty affecting column height, 141 a range of mass flow rate values was chosen for sensitivity tests. The smaller mass flow 142 rates (i.e. sim_etna1a and sim_etna2a) approximately correspond to the averaged mass 143 flow rates estimated from the deposit volume (see Scollo et al. [2007]) when assuming a 144 constant eruption flux over 5 days. As Fig. 2a illustrates, VOL-CALPUFF, for such dis-145 charge rates, produces buoyant columns able to rise up to only a few hundreds of meters. 146 By increasing flow rate, ascent behavior changes to super-buoyant and the column reaches 147 greater heights. For comparison, in figure 2a is also reported the approximate range of 148 observed column heights. Some chronicles mention that the column reached a peak height 149 of at least 5500 m above sea level (i.e. about 3000 m above the vent) on 23 July (Coltelli 150 et al. [2001]). This observation allows us to use our modeling results in an inverse mode, 151 and to quantify the expected mass flow rate from the column height estimation. It results 152 that an eruption characterized by a mass flow rate in the range $1.5 * 10^4$ (sim_etna3a) 153 - $2.5 * 10^4$ (sim_etna4a) kg/s can produce an eruptive column able to reach an altitude 154 comparable with that observed. Sensitivity analysis of column height as a function of 155 different velocity/radius values (at a constant mass flow rate and water content) was also 156 carried out (Fig. 2b). Simulation results clearly show how a variation in velocity of 25 -157 100 m/s does not change the column height by more than 10%. We therefore fixed the 158

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¹⁵⁹ initial mixture velocity to 25 m/s in most simulations. Lastly, we investigated the effects ¹⁶⁰ of variations in water vapor mass content (Fig. 2c). Variations up to 50% produced ¹⁶¹ a maximum change in column height of about 15 %, making this a significant variable ¹⁶² $(sim_etna1a-c, sim_etna4a-c, sim_etna6a-c)$. In summary, the sensitivity study shows that ¹⁶³ the plume mass flow rate is the main volcanological variable controlling column behavior, ¹⁶⁴ with water content also playing a significant role. Variations in velocity/radius values ¹⁶⁵ appear to play a secondary role.

¹⁶⁶ 2.2.2. Dispersion of the Ash Cloud and Ground Deposition

Once the plume height has been computed, the VOL-CALPUFF code begins to track 167 the movements of the puff centers and to calculate the puff diffusion in the defined domain. 168 The puff feeding rate corresponds to the mass flow rate feeding the rising plume at the vent 169 minus the amount of mass radially lost by the plume during its ascent. Simulation outputs 170 consist of the temporal and spatial distribution of particle concentration in air and at the 171 ground for the nine grain sizes considered (see Table 2). In this section just a few selected 172 ash dispersal simulation outputs are presented to illustrate the complex and unsteady 173 cloud dynamics. Some semi-quantitative comparisons with independent observations of 174 the phenomenon are also made. An initial quantitative description of the large-scale ash 175 cloud dispersal can be gained by plotting the concentration of the smaller particles in the 176 eruptive mixture, i.e. the 31 μm particles. These particles can be transported to large dis-177 tances even for relatively weak events, such as the 2001 Etna eruption, (Sparks et al. [1997]) 178 and, at the same time, they can be detected by satellite thus allowing a semi-quantitative 179 comparison of modeling results with independent satellite observations (*Rose et al.* [2001]; 180 Bluth and Rose [2004]). Fig. 3 illustrates the evolution of spatial distribution for the total 181

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mass of 31 μm particles during the four-day period for simulation sim_etna4a starting at 182 00UTC on 20 July (results illustrated every 12 hours). This quantity is obtained by inte-183 grating the concentration of particles over a column of unitary base up to the maximum 184 extension of the ash plume. It allows dispersal mapping at a given time for a selected par-185 ticle size. During the first hours of 20 July, fine ash was transported towards Calabria in 186 the NE direction as reported by direct observations. In the following hours the ash cloud 187 started moving towards the E and then S (clockwise rotation of ca. 90 degrees). During 188 22 July the ash cloud passed several times over Catania showing a series of fluctuations 189 in the S-SE directions. Finally, on 23-24 July the cloud remained mainly directed SE. An 190 initial quantitative comparison between modeling results and observations can be made 191 by using satellite images. A number of similar comparisons were carried out at different 192 times during the four days of dispersal for different mass flow rates and the agreement 193 between remote sensing observations and predictions for the finer ash distribution was 194 always a close one. Fig. 4 presents the comparison between the dispersal patterns of the 195 different granulometric classes of $sim_etna 4a$ (characterized by a mass flow rate of $2.5*10^4$ 196 kg/s), as computed by VOL-CALPUFF, and NOAA/AVHRR satellite image (available 197 at the web site http://www.nerc-pml.ac.uk/rsdas/projects/etna/pic_gallery.html). The 198 comparison refers to 1114UTC on 20 July. For comparison purposes, 3 μ m particles were 199 also considered in the simulation to evaluate the behavior of micron-size particles. Figure 200 4 illustrates the spatial distribution of the total mass of each particle size obtained by 201 vertically integrating the concentration of each class over the entire plume height. It can 202 be seen that the 3 and 31 μ m particles closely follow the dispersal pattern apparent in 203 the satellite image (E-NE), matching the main dispersal axis and the lateral extension of 204

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²⁰⁵ the detected plume, whereas all the other particles are transported in different directions ²⁰⁶ not detected by satellite. In particular, particles of 62 and 125 μ m seem to move firstly ²⁰⁷ towards E and then slightly deviate towards the S-SE direction. Similarly, the larger re-²⁰⁸ ported particles (250 μ m and 2 mm) show a totally independent movement, characterized, ²⁰⁹ from the beginning, by a northward displacement.

In order to better understand the reason for such a complex dispersal pattern some features of the wind field around the volcano were investigated. Fig. 5 shows the wind field streamlines at the levels of 400 and 2200 m a.g.l computed by the meteorological pre-processor CALMET for the acquisition time of the satellite image shown in Fig. 4. As expected, at levels closer to the ground the wind is strongly affected by the volcano's topography thus generating a more irregular field, whereas at higher altitudes the wind shows a quasi-geostrophic trend.

A more complete picture of the wind field spatial and temporal variability, as produced 217 by the CALMET code, can be gained by analysing the wind direction and speed at 218 different altitudes. Fig. 6 describes variations of wind direction and speed over four days 219 (21-24 July) above the vent area, 15 km E and 15 km S from the vent, at two different 220 altitudes above the ground (400 and 2200 m). A key result is the strongly irregular 221 patterns of the flow field close to the ground (400 m) (see Figs. 6a and 6c), above the 222 vent and 15 km S and E. This irregularity is particularly pronounced for wind direction. 223 At the vent and at a location 15 km E of it, wind direction varies by 90-130 degrees over 224 the four days (blowing towards W above the vent and towards E some 15 km E from it). 225 These variations become even more dramatic 15 km S of the vent where three inversions of 226 180 degrees, between the evening of the 23 and the evening of the 24 July, are predicted. 227

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In contrast, at 2200 m a.g.l the wind is blowing from NW to SE on almost any day of the investigated period with a speed on the first day about five times greater than that at lower levels. Wind speed at high altitude remains twice near-ground wind speed on all other days.

VOL-CALPUFF can be used to estimate ground deposition as a function of assumed 232 particle size distribution and mass flow rate. Fig. 7 illustrates particle deposition for six 233 of the nine granulometric sizes employed in the *sim_etna4a* simulation. Deposition refers 234 to total ash accumulation over four days, between 21 and 24 July. We considered this 235 shorter period because of the availability of deposit measurements (Scollo et al. [2007]) 236 which enable the comparison of model predictions with deposit data (see hereafter). It 237 is evident from the figure how the various particles disperse and deposit following very 238 different dynamics. On figure 7, the smaller particles (e.g. $31 \ \mu m$) are mostly transported 239 in the direction of the prevailing high-altitude winds and remain confined in the SE 240 quadrant. Intense wind quickly advects these low settling velocity particles out of the 241 domain, so that the deposit iso-contours show values up to several orders of magnitude 242 smaller than those of the other particles sizes. In contrast, for increasingly large particles 243 the location of maximum deposition moves towards the vent, as expected. 500 μ m and 244 1 mm particles are much less dispersed than finer particles and their deposit is more 245 uniform in the azimuthal direction. Similar patterns were obtained for the 2, 4 and 8 mm 246 particles. 247

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²⁴⁹ By summing the partial deposits produced by the nine different particle sizes it is ²⁵⁰ possible to compute the total particle accumulation. Fig. 8a and 8c show the mass isolines

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of ash deposited after four days of activity for the two mass flow rates (i.e. 1.5 and $2.5 * 10^4$ 251 kg/s suggested as end-member values for the event investigated (sim_etna3a, sim_etna4a). 252 Fig. 8b and 8d show a comparison of the two computed deposition patterns with the 253 deposit measurements, in kg/m^2 , as reported by Scollo et al. [2007]. The total deposit 254 isolines show a quite irregular and non-symmetric pattern in both cases. Both results 255 present a deposit characterized by a preferred E-SE direction, due to higher columns that 256 inject material at atmospheric levels characterized by intense wind, which, as we have 257 already seen (Fig. 6b), blew most of the time towards this quadrant. The predicted 258 deposit distributions show that maximum deposition is at a few kilometers from the vent 259 in the S-SE direction, as in the measured deposit (Scollo et al. [2007]). The lower end-260 member mass flow rate (Fig. 8a and Fig. 8b - sim_etna3a) produces a deposit which 261 fits the sampled data with a correlation coefficient of 0.77 and with more than 97% of 262 simulated data ranging from 1/5 to 5 times the observed values. For the larger end-263 member mass flow rate (Fig. 8c and Fig. 8d - sim_etna4a) the numerical prediction 264 largely overestimates the amount of ash deposited. A comparison between computed and 265 measured values quantifies this discrepancy to be up to about a factor of 5. 266

3. Discussion

The above results show that VOL-CALPUFF is helpful to investigate ash dispersal processes. The new model captures the strong influence of meteorological conditions on plume rise dynamics, shows the necessity to use codes able to determine column height using a plume rise model and the importance of using an accurate meteorological data set. A useful result that can be inferred *a posteriori* from modeling of column height in the presence of realistic meteorological conditions is the estimation of the mass flow rate

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feeding the plume. For the 2001 Etna event, the results reported in Fig. 2 illustrate that 273 only mass flow rates between 1.5×10^4 kg/s (sim_etna3a) and 2.5×10^4 kg/s (sim_etna4a) 274 are likely able to produce a maximum column height of about 3000 m above the vent on 275 23 July. Based on the analysis reported in the companion paper (see Fig. 4 of *Barsotti* 276 et al. [2007], this issue) on the effect of the entrainment coefficients and the values here 277 assumed, these values should be considered as minimum values in order to reproduce the 278 observed column height. In contrast, using an *averaged* mass flow rate, obtained by di-279 viding the total amount of mass erupted $(1.02 * 10^9 \text{ and } 2.31 * 10^9 \text{ kg from Scollo et al.}$ 280 [2007]) by event duration (*sim_etna1a* and *sim_etna2a*), it would not be possible to match 281 the observed column height. Only in a still environment, the plume could reach altitudes 282 comparable with observed ones if fed with averaged mass flow rates (as assumed in Costa 283 et al. [2006]). For the long-lasting weak plume considered here, the assumption of an aver-284 age mass flow rate is inadequate to simulate the whole dispersal process, since it strongly 285 underestimates column height. However the range of mass flow rate estimated above from 286 the column height $(sim_etna3a \text{ and } sim_etna4a)$ is likely representative of the peak mass 287 flow rate. This could significantly overestimate the erupted volume and ash ground de-288 position. If we consider the column height and deposit estimates, we see that each of the 289 two end-member values matches only one of the two observed variables. Plume feeding 290 with a steady mass flow rate of 1.5×10^4 kg/s allows a match with the amount of material 291 deposited (see Fig. 8b), but the predicted column height only approximately matches the 292 minimum observed values (see Fig. 2a). By feeding the plume with the higher mass flow 293 rate (i.e. 2.5×10^4 kg/s) it is possible to match the observed plume height but not deposit 294 estimates (Fig. 8d). A reasonable explanation for this discrepancy is that events lasting 295

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²⁹⁶ several days, such as the 2001 Etna eruption, are characterized by unsteady, intermittent
²⁹⁷ plume feeding (*INGV* - *CT Staff* [2001]; *Taddeucci et al.* [2004]; *Scollo et al.* [2007], P. Al²⁹⁸ lard written communication (2007)). Plume-feeding is therefore quite difficult to quantify
²⁹⁹ and sometimes also hard to detect (for instance during nocturnal periods or bad weather
³⁰⁰ conditions).

We believe that a better match between modeling results and observations cannot be 301 achieved under the assumption of steady-state plume feeding over the investigated four-302 day period. The few available measurements are affected by significant uncertainty and 303 do not allow more accurate comparison or the definition of time-dependent plume feeding 304 over the simulated period. In order to support this interpretation we performed addi-305 tional simulations adopting two different mass flow rates for the plume and puffs. We ran 306 VOL-CALPUFF assuming a mass flow rate of 2.5^{*10^4} and 1.0^{*10^4} kg/s for the plume 307 and puff release, respectively. In this way we were able to match the observed column 308 height and produce a deposit distribution (not shown) that fits the measured values with 309 a correlation coefficient (R=0.82) higher than that for the simulations reported in Fig. 8. 310 Another variable that could significantly reduce the above discussed discrepancy is the 311 gas mass fraction of the eruptive mixture at the vent. Water contents greater than those 312 assumed here could favor higher columns for a given mass flow rate (Fig. 2c). 313

VOL-CALPUFF allows analysis of new and non-intuitive effects of ash dispersal in an orographically complex area under the action of 3D meteorological conditions. These effects are apparent in the dispersal behavior of contrasting particles sizes. The major mass *splitting* effect we observe in Fig. 4 and Fig. 7 between particles of different sizes is the result of two main factors: (a) the dependence of settling velocities on particle

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dimension and density and (b) the great variability and complexity of wind field patterns 319 in the vertical and horizontal directions. These factors make larger particles cross the 320 atmosphere in a temporal interval up to two orders of magnitude faster than the smaller 321 particles (minutes vs. hours), travelling faster towards the ground where wind fields are 322 more irregular. Particles up to a few tens of microns remain much longer at high at-323 mospheric levels, where more uniform winds take them far from the vent. These results 324 are in agreement with the very first attempt to investigate the control of wind-shear and 325 grain size distribution on ash dispersal by Woods et al. [1995]. In terms of validation of 326 the predicted particle dispersal behavior, an initial comparison of simulation outputs with 327 satellite images was carried out. Results suggest that VOL-CALPUFF is able to correctly 328 predict the dispersal of the smaller particles moving at high altitude above the ground. 329 The path observed in satellite images was closely followed by particles up to about 30 μ m 330 in our simulations. In contrast, the AVHRR satellite images taken by the NOAA satel-331 lite do not show the more irregular behavior that the simulations produce for particles 332 larger than about 30 μ m. This result is consistent with the present limitations of satellite 333 remote sensing techniques in detecting suspended material with a radius greater than a 334 few tens of microns (Wen and Rose [1994]; Yu and Rose [2002]; Bluth and Rose [2004]). 335 Given the importance, in terms of mass dispersed in the atmosphere, of particles with 336 size in the range 62 μ m-1 mm (see Table 2), an important contribution to the detection 337 of these coarse particles should come from the use of other remote sensing techniques 338 such as weather radar observation (Lacasse et al. [2004]; Marzano and Vulpiani [2006]). 339 Unfortunately, such data are not available for the 2001 Etna event and cannot be used 340 to validate the numerical predictions. This type of investigation appears particularly im-341

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³⁴² portant for particle sizes (e.g. 125 μ m) that can reach concentrations of a few g/m³ far ³⁴³ away from the area indicated by the satellite images. Such concentrations are well above ³⁴⁴ those used by other ash forecasting codes as a threshold for ash cloud visibility (*Simpson* ³⁴⁵ *et al.* [2002]). Qualitative support for the modeling results here presented comes from ³⁴⁶ the observation that ash fallout of coarse particles was actually recorded several times, ³⁴⁷ during the 2001 event, well away from the main plume dispersal direction (Andronico D., ³⁴⁸ personal communication, 2005).

From the above discussion the fundamental role played by meteorology in the determination of both plume rise and ash dispersal is evident. A better understanding of the minimum resolution of meteorological data required to capture, with sufficient accuracy, most of the results illustrated here, is necessary for future progress. The present modeling results could also contribute towards a more quantitative interpretation of remote sensing and field data.

4. Conclusions

The first application of VOL-CALPUFF, to the 2001 explosive event of Mt. Etna, 355 highlights new key features of the dynamics of weak volcanic plumes in the presence of 356 a complex wind field. Some of the most important aspects are: (a) the major effect 357 of the wind field in the determination of the column height and bending, as well as of 358 plume ascent style, (b) the great difference in the dispersal patterns followed by ashes 359 of different sizes producing a very complex and non-intuitive dispersal, (c) the complex 360 distribution of ash on the ground showing preferential directions of dispersal and, for 361 events characterized by long duration and variable winds, a quite irregular deposit pattern, 362 (d) the impossibility to reproduce both column height and deposit features by adopting 363

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³⁶⁴ a vent steady-state mass flow rate due to temporal variations in plume feeding during ³⁶⁵ the four-day period investigated. Furthermore, this initial application of VOL-CALPUFF ³⁶⁶ to weak plumes poses new questions about the use of low-resolution meteorological data ³⁶⁷ for the accurate prediction of the complex ash dispersal and fallout dynamics. Finally, it ³⁶⁸ highlights, once more, the importance of comparing modeling results to accurate field data ³⁶⁹ and observations in order to reach a consistent understanding of the dispersal dynamics ³⁷⁰ and also prompts new developments in modeling and observational techniques.

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Figure 1. Photographs of the volcanic plume and ash fallout over Catania and nearby villages, during the recent eruptive events at Mt.Etna. Photos by P. Papale (INGV Pisa) and from www.larepubblica.it, respectively.

Table 1. List of performed simulations. The table contains the input data values assumed at the vent. V is velocity, R vent radius, H_2O the weight percent of water vapor and mfr the resulting mass flow rate. Vent radius was derived from the other three variables by using the mass conservation equation. Vent altitude, vent flow temperature and exit mixture pressure were kept constant in all simulations and are equal to 2570 m, 1300 K and atmospheric value, respectively. As explained in *Barsotti et al.* [2007] (this issue) all the simulations were performed using entrainment coefficients values (α and γ) equal to 0.09 and 0.6, respectively.

Name	V (m/s)	R (m)	$H_2O (\mathrm{wt\%})$	mfr~(kg/s)
sim_etna1a	20	2.7	3	$2.0 * 10^3$
sim_etna1b	30	2.7	4.5	$2.0 * 10^3$
sim_etna1c	25	1.7	1.5	$2.0 * 10^3$
sim_etna2a	20	4.8	3	$6.6 * 10^3$
sim_etna3a	25	6.7	3	$1.5 * 10^4$
sim_etna4a	25	8.7	3	$2.5 * 10^4$
sim_etna4b	30	9.7	4.5	$2.5 * 10^4$
sim_etna4c	25	6.2	1.5	$2.5 * 10^4$
sim_etna4d	50	6.2	3	$2.5 * 10^4$
sim_etna4e	75	5	3	$2.5 * 10^4$
sim_etna4f	100	4.3	3	$2.5 * 10^4$
sim_etna5a	25	11	3	$4.0 * 10^4$
sim_etna6a	25	15.1	3	$7.5 * 10^4$
sim_etna6b	30	16.7	4.5	$7.5 * 10^4$

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Name	V (m/s)	R (m)	$H_2O (\mathrm{wt\%})$	mfr (kg/s)
sim_etna6c	25	10.6	1.5	$7.5 * 10^4$
sim_etna6d	50	10.6	3	$7.5 * 10^4$
sim_etna6e	75	8.7	3	$7.5 * 10^4$
sim_etna6f	100	7.5	3	$7.5 * 10^4$
sim_etna7a	25	19.86	3	$1.3 * 10^5$

Table 2. The particle size distribution of the eruptive mixture assumed at the vent in the simulations performed as coming from the reconstructed total granulometric distribution (*Scollo et al.* [2007]).

Φ	Diameter (μm)	wt%
-3	8000	0.11
-2	4000	0.69
-1	2000	2.47
0	1000	7.53
1	500	20.36
2	250	40.2
3	125	24.06
4	62	3.89
5	31	0.69

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Figure 2. Temporal variations of column height as computed by VOL-CALPUFF for some of the simulations reported in Table 1. (a) Comparison between column heights computed for plumes with different mass flow rates. The vent mass flow rate is kept constant in time during each simulation. The superimposed horizontal lines delimit the observed range for the 21-23 July period (*Scollo et al.* [2007]); (b) Comparison between column heights computed for plumes with different mixture velocities (at constant mass flow rate and water content); (c) Comparison between column heights computed for plumes with different water vapor mass fractions.



Figure 3. Cumulative mass of 31 μ m particles, obtained by integration of its concentration over all vertical levels in the atmosphere. From the upper left-hand corner the concentration is mapped each 12 hours starting at 00UTC on 20 July. Isocontour values are reported on the bottom left-hand corner and refer to the total mass in grams. Results from simulation sim_etna4a .



Figure 4. SeaStar satellite image at 1114UTC on 20 July compared with sim_etna4a results. The figures show the amount of ash integrated in the vertical extension of the domain for six different classes (3 (a), 31 (b), 62 (c), 125 (d), 250 (e) μ m and 2 mm (f)). In all images the minimum reported isocontour is equal to 1 g. The comparison shows how the detection of the plume by the satellite image is limited to the detection of particles of a few- μ m (*Bluth and Rose* [2004]).



Figure 5. Wind field streamlines, as produced by CALMET model, at 400 m (a) and 2200 m (b) above ground level at 1100UTC on 20 July.



Figure 6. Temporal variations of wind direction (a and b) and wind speed (c and d) at different spatial locations. The wind information, produced by the CALMET model, is reported hourly for three different locations: above the vent, 15 km from the vent toward E or S. Wind information related to a terrain-following vertical level at 400 m agl (a and c) are compared with the same informations obtained for a 2200 m terrain-following level (b and d). The figures illustrate the effects on low atmosphere wind direction caused by complex orography (the effect of eruption dynamics on local atmospheric conditions is not considered by the model). At a higher level the wind field is more regular and faster than closer to the ground.



Figure 7. Ground deposit as a function of particle diameter for an emission rate of $2.5 * 10^4$ kg/s (sim_etna_4a) over four days (21-24 July). Six classes with diameter of (a) 31 μ m, (b) 62 μ m, (c) 125 μ m, (d) 250 μ m, (e) 500 μ m and (f) 1 mm are reported. Particles with larger diameters show a behavior similar to the 1 mm particles.



Figure 8. Total accumulation on the ground of all particle sizes after four days (21-24 July) of steady ash emission. A comparison between two different mass flow rates, i.e. $1.5 * 10^4$ kg/s (sim_etna3a) (a, b) and $2.5 * 10^4$ kg/s (sim_etna4a) (c, d), is shown. Both deposit patterns are characterized by an irregular shape with a peak at some kms from the vent ((a) and (c)). A zoomed comparison with measured field data (from *Scollo et al.* [2007]) highlights a good agreement for the case sim_etna3a (correlation coefficient R = 0.77) (b), and an overestimation up to about a factor of 5 for the case sim_etna4a (d).

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