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<tr>
<td>Citation</td>
<td>PHYSICAL REVIEW LETTERS (2009), 103(16)</td>
</tr>
<tr>
<td>Issue Date</td>
<td>2009-10</td>
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<tr>
<td>URL</td>
<td><a href="http://hdl.handle.net/2433/109854">http://hdl.handle.net/2433/109854</a></td>
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<tr>
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<tr>
<td>Type</td>
<td>Journal Article</td>
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<td>Textversion</td>
<td>publisher</td>
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Multilamellar Structures Induced by Hydrophilic and Hydrophobic Ions Added to a Binary Mixture of D$_2$O and 3-Methylpyridine

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(Received 15 March 2009; published 14 October 2009)

A phase transition is observed between a one-phase disordered phase and an ordered phase with multilamellar (onion) structures in an off-critical mixture of D$_2$O and 3-methylpyridine (3MP), containing 85 mM of a salt in the absence of a surfactant. The salt consists of hydrophilic cations and hydrophobic anions that interact asymmetrically with the solvent composition fluctuations inducing mesophases. The structure factor of the composition distribution obtained from small-angle neutron scattering has a peak at an intermediate wave number in the disordered phase and multiple peaks in the ordered phase. Lamellar layers forming onions are composed of solvation-induced charged membranes swollen by D$_2$O. The onion phase is realized only for small volume fractions of 3MP (in D$_2$O-rich solvent).

Binomial mixtures of water and organic solvents have been used extensively to study the universal aspects of their critical behavior and phase separation dynamics. However, the unique ion effects in such mixtures, where preferential hydration around each ion should affect the composition fluctuations, has not yet been studied in detail [1,2]. The addition of even small concentrations of salts composed of small hydrophilic ions to the above-mentioned mixtures can drastically change their phase behavior [3–7]. Many groups exhibit long-lived heterogeneities (occasionally extending over a few micrometers) in the one-phase state [8,9] and a third phase visible as a thin plate forms at a liquid-liquid interface in the two-phase state [10]. These observations were reproduced in different experiments using high-purity solvents and salts, and therefore, they should be regarded as ion-induced supramolecular aggregates. Although it has not yet been understood clearly, the solvation (ion-dipole) interaction among ions and polar molecules should play a major role in these phenomena together with the Coulombic interaction among charges [1,2].

Recently, the solvation effect on phase behavior was theoretically studied in polar mixtures, including the case of antagonistic ion pairs composed of hydrophilic and hydrophobic ions [11]. Such cations and anions interact differently with composition fluctuations, leading to a charge-density-wave phase for sufficiently large solvation asymmetry even at low salt concentrations. In a small-angle neutron scattering (SANS) experiment, Sadakane

et al. [12] found a peak at an intermediate wave number $Q_m$ (~0.1 Å$^{-1}$) in a near-critical mixture of D$_2$O and 3-methylpyridine (3MP) containing 100 mM sodium tetraphenylborate (NaBPh$_4$). The volume fraction of 3MP, denoted by $\phi_{3MP}$, was chosen to be 0.35, 0.42, and 0.54. This salt dissociates into hydrophilic Na$^+$ and hydrophobic BPh$_4^-$ ion pairs. The BPh$_4^-$ acquires strong hydrophobicity such that the salt is more soluble in pure 3MP than in pure D$_2$O despite the hydrophilic nature of Na$^+$ [13]. Furthermore, the mixture exhibited colors that changed dramatically on approaching criticality at low salt concentrations (~10 mM), indicating the emergence of large-scale heterogeneities. No distinct mesophases appeared in the same solvent when hydrophilic salts such as NaCl were added [14].

In this Letter, we again use the system of D$_2$O + 3MP + NaBPh$_4$. We purchased D$_2$O with isotopic purity of 99.9% from EURISO-TOP and 99.5% purity 3MP and 99.5% purity NaBPh$_4$ from Aldrich. These chemicals were mixed without further purification. Hereafter, we report the observation of multilamellar (onion) structures for small $\phi_{3MP}$. Detailed measurements were carried out for $\phi_{3MP} = 0.08$, 0.09, 0.11, and 0.14 in the temperature range $283 \leq T \leq 343$ K. The concentration of NaBPh$_4$ was fixed at 85 mM. Without the salt, the mixtures in the measured composition and temperature ranges remained homogeneous in one-phase states with small composition fluctuation, because the lower critical solution temperature is at $T = 310$ K and $\phi_{3MP} = 0.30$. Upon the addition of
Thus, our result indicates that the onion structure is formed below $T_L$, where $T_L = 318$ K is the transition temperature.

SANS experiments were performed at SANS-J, JRR-3M at the Japan Atomic Energy Agency. A 6.0 Å incident neutron beam was mechanically selected with a wavelength resolution of 13%, and the scattered neutrons were detected at distances of 2.5 and 10 m from the sample position. Each sample was kept in 2 mm thick quartz cell and placed in a temperature-controlled chamber with an accuracy better than 0.1 K. The measured wave number $Q$ ranged between $6 \times 10^{-3}$ Å$^{-1}$ and $1.3 \times 10^{-1}$ Å$^{-1}$ and the observed two-dimensional data were azimuthally averaged. Figure 2 shows the SANS intensity vs $Q$ from the sample of $\phi_{3MP} = 0.09$. By varying $T$, it was found that the system was in the disordered phase for $T > T_L$ and in the ordered phase for $T < T_L$. At the lowest temperature 283 K, a pronounced Bragg peak is observed at $Q_m = 3.4 \times 10^{-2}$Å$^{-1}$, a second one is observed at 6.8 × 10$^{-2}$ Å$^{-1}$, and a slight shoulder is observed at around 1.0 × 10$^{-1}$ Å$^{-1}$ (although this is not clearly seen in the figure). The same behavior was found for all the SANS data below $T_L$, indicating the formation of aligned lamellae. This is further evidence of the onion formation below $T_L$.

Our SANS profile at $T = 283$ K can be well fitted to the formula for lyotropic lamellae given by Nallet et al. [20] in the range $Q > 3 \times 10^{-2}$ Å$^{-1}$, as shown by the solid curve of $T = 283$ K in Fig. 2(a). The spatial distribution of the scattering density profile is reconstructed as shown in Fig. 2(b), yielding the mean repeat distance $d = 174.8$ Å and the membrane thickness $\delta = 16.2$ Å. The scattering density difference in their formula is $\Delta \rho = (4.57 \pm 0.02) \times 10^{10}$ cm$^{-2}$, which is almost equal to that between 3MP and D$_2$O.

In the disordered phase above $T_L$, the intensity exhibits a broad peak at $Q_m \sim 0.07$ Å$^{-1}$, with $Q_m$ shifting slightly to lower values as $T \to T_L$. It may fairly be fitted to the theoretical mean-field intensity $I_{OK}(Q)$ [solid curves for $T = 343$ and 318 K in Fig. 2(a)] [11]. Its inverse is of the form

$$I_{OK}(0)/I_{OK}(Q) = 1 + \left[1 - \frac{\gamma_p^2}{1 + \lambda^2 Q^2}\right]\xi^2 Q^2, \quad (2)$$

where $\gamma_p$ indicates the degree of solvation asymmetry between cations and anions. For $\gamma_p = 0$ it follows the Ornstein-Zernike form. For $\gamma_p > 1$ the structure factor

FIG. 1 (color online). Optical microscopic images obtained from the sample of $\phi_{3MP} = 0.09$: (a) $T = 323$ K (disordered), (b) $T = 313$ K (onion), and (c) $T = 293$ K (onion). Scale bars in (b) and (c) are 100 μm. (b') A magnified view of (b) (left) and a birefringent image of the same region (right) with the scale bar being 20 μm.
has a peak at an intermediate wave number. In the present case \( C_{13} \) is treated as an adjustable parameter given by 2.32 at \( T = 333 \) K and by 1.96 at \( T = 320 \) K. Slightly below \( T_L \), the system was phase separated into an ordered phase in an upper region and a disordered phase in a lower region. The data of \( T < T_L \) in Fig. 2(a) are those from the upper region. The transition is, thus, a first-order one. The same \( T_L \) was obtained with decreasing and increasing \( T \) around the transition. The mass density of \( D_2O \) is 1.11 g/cm\(^3\) and that of 3MP is 0.96 g/cm\(^3\); therefore, the ordered phase contains more 3MP than the coexisting disordered phase. In this case the mass density difference between the two phases is less than 10 mg/cm\(^3\). As \( T \) was further lowered, the onion droplets increasingly swelled with \( D_2O \). As a result, the lamellar spacing increased as will be shown in Fig. 3. At low temperatures below 293 K, the entire cell was filled with swollen onions.

Figure 3 shows the temperature dependence of \( d = 2\pi / Q_m \). For \( \phi_{3MP} = 0.09 \), \( d \) increases considerably with decreasing \( T \) in the range 293 K < \( T < T_L \) because of the swelling of the onions taking place on a time scale of 1 min. For \( \phi_{3MP} = 0.14 \), \( d \) varies very weakly in the entire temperature range. Above 330 K, the dependence of \( d \) on \( \phi_{3MP} \) is weak. The transition temperature \( T_L \) decreases with increasing \( \phi_{3MP} \), and reaches 318 and 303 K for \( \phi_{3MP} = 0.09 \) and 0.11, respectively. However, the system remained in the disordered phase for \( \phi_{3MP} \leq 0.14 \). The figure shows that the transition occurs for \( d \sim 130 \AA \).

Figure 4 shows an ultra-small-angle neutron scattering (USANS) spectrum observed at PNO spectrometer in JRR-3M (\( 2 \times 10^{-5} < Q < 5 \times 10^{-4} \) \( \AA^{-1} \)) together with the data obtained at SANS-J in the focusing-SANS mode (the lowest \( Q \) is expanded down to \( 3 \times 10^{-4} \) \( \AA^{-1} \) [21]). The measurement was carried out at \( T = 298 \) K for the sample of \( \phi_{3MP} = 0.09 \). The low-\( Q \) profile obeys the Porod law in the range of \( 1 \times 10^{-4} < Q < 5 \times 10^{-3} \) \( \AA^{-1} \). It should arise from the interfaces of the onion droplets in the range \( R^{-1} \ll Q \ll Q_m \), where \( R \) is the lattice spacing.
droplet radius. Each onion is composed of approximately 1000 concentric lamellae provided the droplet interior is filled with layers. The estimated value of $R$ from this data analysis is consistent with the observation by optical microscopy, although a more precise data analysis is needed here accounting for the effects of the resolution function and the polydispersity.

Now, we discuss the physical processes involved. In aqueous solutions, each Na\(^+\) ion is known to be surrounded by a hydration shell composed of several water molecules [1,2]. This suggests that each BPh\(_4\) ion should be solvated by a certain number of 3MP molecules around it due to its strong hydrophobicity [13], although this solvation number $N_{3MP}^s$ is currently not known. Equation (1) gives the particle number densities in the disordered phase at $\phi_{3MP} = 0.09$. Then, the number density of the solvating 3MP molecules is $N_{3MP}^s n_{salt}$, while that of the free 3MP molecules is $n_{3MP} - N_{3MP}^s n_{salt}$. In the onion phase, each onion should be composed of concentric thin membranes comprising BPh\(_4\) ions and solvating 3MP molecules. If all the BPh\(_4\) ions are trapped on the membranes, the areal density of BPh\(_4\) is given by

$$\sigma_a = n_x d = 92 \, (\text{Å}^2).$$

Here, $d = 180 \, \text{Å}$ is the spacing in case (c) shown in Fig. 1. On both sides of each membrane, Na\(^+\) ions are localized within “counterion layers” with thickness of the order of the Debye screening length $\lambda = (4\pi n_{salt} \ell_B)^{-1/2} \approx 15 \, \text{Å}$, where $\ell_B$ is the Bjerrum length in D\(_2\)O ($\sim 7 \, \text{Å}$). The segregation of charge and composition in this manner should significantly lower the solvation free energy. However, the mechanism leading to the onion structure remains to be investigated in future, where the electrostatic energy due to the deformation of the charged membranes should be crucial as a new aspect.

We also argue why the lamellar structure is formed in the window composition range 0.05 < $\phi_{3MP}$ < 0.12. When lamellae are formed, trapped 3MP molecules can simultaneously interact with a number of BPh\(_4\) ions on the plane, thus achieving high solvation. If $\phi_{3MP}$ is less than the lower bound, the number density of 3MP should be too small to trigger the aggregation of BPh\(_4\) ions. If $\phi_{3MP}$ is increased above the upper bound, 3MP molecules would be sufficiently abundant outside the thin membrane regions. This would eventually lead to the destruction of the membranes with increasing $\phi_{3MP}$, where the delocalization of BPh\(_4\) ions can increase their translational entropy without an increase in the solvation free energy.

In summary, we have realized an onion structure in a D\(_2\)O-3MP mixture with a low concentration of 3MP by adding a small amount of an antagonistic salt composed of hydrophilic and hydrophobic ions without a surfactant. In previous experiments [3–7], hydrophilic salts have mostly been used, and many problems have remained unresolved. As demonstrated in this Letter, the effects of adding an antagonistic salt are quite significant in soft matter. We predict that its addition decreases the surface tension of a macroscopic liquid-liquid interface to zero [22], leading to emulsification; this also explains our previous experiment [12].

This work was supported by Grant-in-Aid for Scientific Research on Priority Area “Soft Matter Physics” from the Ministry of Education, Culture, Sports, Science and Technology of Japan. One of the authors (K. S.) was supported by the JSPS (19-3802). We would like to thank Dr. D. Yamaguchi at the JAEA for assisting us in the SANS experiment.

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[9] A. F. Kostko, M. A. Anisimov, and J. V. Sengers, Phys. Rev. E 70, 026118 (2004); M. Wagner et al., Phys. Chem. Chem. Phys. 6, 580 (2004). In these experiments, H\(_2\)O + 3MP + NaBr mixture was used. Large-scale heterogeneities were observed at high concentrations of NaBr (~17 mass%).
[13] Le Quoc Hung, J. Electroanal. Chem. 115, 159 (1980). There arises a difference in the solvation chemical potential of ions across an interface, called the Gibbs transfer free energy. For water nitrobenzene at room temperature, it was estimated to be 17.1 for Na\(^+\) and as ~18.0 for BPh\(_4\) in units of $k_B T$. The dielectric constant of nitrobenzene is 35 and that of 3MP is 9, and therefore, the preference of BPh\(_4\) for the less polar component (hydrophobicity) should not be weakened for D\(_2\)O + 3MP.