The initial value problem for Schrödinger equations on the torus

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This note is a summary of a paper [2]. We are concerned with the initial value problems for linear Schrödinger-type equations of the form

\[ Lu \equiv \partial_t u - i\Delta u + \tilde{b}(x) \cdot \nabla u + c(x)u = f(t, x) \quad \text{in} \quad \mathbb{R} \times \mathbb{T}^n, \quad (1) \]

\[ u(0, x) = u_0(x) \quad \text{in} \quad \mathbb{T}^n, \quad (2) \]

and for semilinear Schrödinger equations of the form

\[ \partial_t u - i\Delta u = F(u, \nabla u, \bar{u}, \nabla \bar{u}) \quad \text{in} \quad \mathbb{R} \times \mathbb{T}^n, \quad (3) \]

\[ u(0, x) = u_0(x) \quad \text{in} \quad \mathbb{T}^n, \quad (4) \]

where \( u(t, x) \) is a complex valued unknown function of \( (t, x) = (t, x_1, \ldots, x_n) \in \mathbb{R} \times \mathbb{T}^n, \)
\( \mathbb{T}^n = \mathbb{R}^n/2\pi\mathbb{Z}^n, \ i = \sqrt{-1}, \ \partial_t = \partial/\partial t, \ \partial_j = \partial/\partial x_j \ (j = 1, \ldots, n), \ \nabla = (\partial_1, \ldots, \partial_n), \)
\( \Delta = \nabla \cdot \nabla, \) and \( \tilde{b}(x) = (b_1(x), \ldots, b_n(x)), \ c(x), \ f(t, x) \) and \( u_0(x) \) are given functions. Suppose that \( b_1(x), \ldots, b_n(x) \) and \( c(x) \) are smooth functions on \( \mathbb{T}^n, \) and that \( F(u, v, \bar{u}, \bar{v}) \)
is a smooth function on \( \mathbb{R}^{2+2n}, \) and

\[ F(u, v, \bar{u}, \bar{v}) = O(|u|^2 + |v|^2) \quad \text{near} \quad (u, v) = 0. \]

In [7], Mizohata proved that, when \( x \in \mathbb{R}^n, \) if the initial value problem (1)-(2) is \( L^2 \)-well-posed, then it follows that

\[ \sup_{(t, x, \omega) \in \mathbb{R}^{1+n} \times \mathbb{S}^{n-1}} \left| \int_0^t \text{Im} \tilde{b}(x - \omega s) \cdot \omega ds \right| < +\infty, \quad (5) \]

where \( \tilde{b} \cdot \xi = b_1 \xi_1 + \cdots + b_n \xi_n. \) Moreover, he gave sufficient condition for \( L^2 \)-well-posedness which is slightly stronger than (5). In particular, (5) is also sufficient condition for \( L^2 \)
well-posedness when \( n = 1. \) Roughly speaking, (5) gives an upper bound of the strength of the real vector field \( (\text{Im} \tilde{b}(x)) \cdot \nabla. \) In other words, if \( (\text{Im} \tilde{b}(x)) \cdot \nabla \) can be dominated by so-called local smoothing effect of \( e^{it\Delta}, \) then (5) must holds. After his results, many authors investigated the necessary and sufficient condition, and some weaker sufficient conditions were discovered. Unfortunately, however, the characterization of \( L^2 \)-well-posedness for (1)-(2) remains open except for one-dimensional case. Such linear theories were applied to solving (1)-(2) in case \( x \in \mathbb{R}^n. \) See, e.g., [3] for linear equations, [1], for nonlinear equations, and references therein.

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On the other hand, the periodic case is completely different from the Euclidean case. The local smoothing effect of $e^{it\Delta}$ fails because the hamiltonian flow generated by the hamiltonian vector field $2\xi \cdot \nabla$ is completely trapped. See [4] for the relationship between the global behavior of the hamiltonian flow and the local smoothing effect.

The purpose of this note is to present the necessary and sufficient condition of $L^2$-well-posedness of (1)-(2), and apply this condition to (3)-(4). To state a definition and our results, we here introduce notation. Let $s \in \mathbb{R}$. $H^s(\mathbb{T}^n)$ denotes the set of all distributions on $\mathbb{T}^n$ satisfying
\[
\|u\|_s^2 = \int_{\mathbb{T}^n} |(1 - \Delta)^{s/2}u(x)|^2 dx < +\infty.
\]
Set $L^2(\mathbb{T}^n) = H^0(\mathbb{T}^n)$, and $\|\cdot\|_s = \|\cdot\|_0$ for short. Let $I$ be an interval in $\mathbb{R}$. $C(I; H^s(\mathbb{T}^n))$ denotes the set of all $H^s(\mathbb{T}^n)$-valued continuous functions on $I$. Similarly $L^1(I; H^s(\mathbb{T}^n))$ is the set of $H^s(\mathbb{T}^n)$-valued integrable functions on $I$.

We here give the definition of $L^2$-well-posedness.

**Definition 1.** The initial-boundary value problem (1)-(2) is said to be $L^2$-well-posed if for any $u_0 \in L^2(\mathbb{T}^n)$ and $f \in L^1_{\text{loc}}(\mathbb{R}; L^2(\mathbb{T}^n))$, (1)-(2) has a unique solution $u \in C(\mathbb{R}; L^2(\mathbb{T}^n))$.

It follows from Banach's closed graph theorem that the condition required in Definition 1 is equivalent to a seemingly stronger condition, that is, for any $u_0 \in L^2(\mathbb{T}^n)$ and for any $f \in L^1_{\text{loc}}(\mathbb{R}; L^2(\mathbb{T}^n))$, (1)-(2) has a unique solution $u \in C(\mathbb{R}; L^2(\mathbb{T}^n))$, and for any $T > 0$ there exists $C_T > 0$ such that
\[
\|u(t)\| \leq C_T \left( \|u_0\| + \int_0^t \|f(s)\| ds \right), \quad t \in [-T, T].
\]  
(6)

Firstly, we present $L^2$-well-posedness results for linear equations.

**Theorem 2.** The following conditions are mutually equivalent:

1. (1)-(2) is $L^2$-well-posed.

2. For $x \in \mathbb{T}^n$ and $\alpha \in \mathbb{Z}^n$
\[
\int_0^{2\pi} \text{Im} \tilde{b}(x - \alpha s) \cdot \alpha ds = 0.
\]  
(7)

3. There exists a scalar function $\phi(x) \in C^\infty(\mathbb{T}^n)$ such that $\nabla \phi(x) = \text{Im} \tilde{b}(x)$.
When $n = 1$, set $b(x) = b_1(x)$. The condition (7) is reduced to

$$\int_0^{2\pi} \text{Im} b(x) dx = 0. \quad (8)$$

The condition (7) is the natural torus version of (5). More precisely, (7) is a special case of Ichinose's necessary condition of $L^2$-well-posedness discovered in [5]. On the other hand, the condition 3 corresponds to Ichinose's sufficient condition of $L^2$-well-posedness discovered in [6]. Theorem 2 makes us expect analogous results for nonlinear equations. In fact, we have local existence and local ill-posedness results as follows.

**Theorem 3.** Let $s > n/2 + 2$. Suppose that there exists a smooth real-valued function $\Phi(u, \bar{u})$ on $\mathbb{R}^2$ such that for any $u \in C^1(\mathbb{T}^m)$

$$\nabla \Phi(u, \bar{u}) = \text{Im} \nabla_u F(u, \nabla u, \bar{u}, \nabla \bar{u}). \quad (9)$$

Then for any $u_0 \in H^s(\mathbb{T}^m)$, there exists $T > 0$ depending on $\|u_0\|_s$ such that (3)-(4) possesses a unique solution $u \in C([-T, T]; H^s(\mathbb{T}^m))$. Furthermore, Let $\{u_{0,k}\}$ be a sequence of initial data belonging to $H^s(\mathbb{T}^m)$, and let $\{u_k\}$ be a sequence of corresponding solutions. If

$$u_{0,k} \to u_0 \text{ in } H^s(\mathbb{T}^m) \text{ as } k \to \infty,$$

then for any $m < s$

$$u_k \to u \text{ in } C([-T, T]; H^m(\mathbb{T}^m)) \text{ as } k \to \infty. \quad (10)$$

**Theorem 4.** Suppose that there exists a holomorphic $n$-vector function

$$\bar{G}(u) = (G_1(u), \cdots, G_n(u)), \quad u \in \mathbb{C}$$

such that $G(u) \not\equiv 0$, and

$$F(u, \nabla u, \bar{u}, \nabla \bar{u}) = \nabla \cdot \bar{G}(u) \quad (11)$$

for any $u \in C^1(\mathbb{T}^m)$. Then (3)-(4) is not locally well-posed in the sense of Theorem 3.

It seems to be hard to show the continuous dependence of the solution on the initial data because the gain of derivative of $e^{i\Delta}f$ fails when $x \in \mathbb{T}^n$. To prove Theorem 4, we construct a sequence of solutions which are real-analytic in $x$ by using the idea of the abstract Cauchy-Kowalewski theorem. Hence it is essential that $G(u)$ is holomorphic.

In what follows we give the sketch the proofs of Theorems 2 and 4. We omit the sketch of the proof of Theorem 3.

**Proof of Theorem 2.** To prove $1 \Rightarrow 2$, we suppose that the condition 2 fails, and construct a sequence of approximate solutions $\{u(t, x)\}$ which break an energy inequality (6). Suppose that there exist $x_0 \in \mathbb{T}^n$ and $\alpha \in \mathbb{Z}^n \setminus \{0\}$ such that

$$\int_0^{2\pi} \text{Im} \bar{b}(x_0 - \alpha s) \cdot \alpha ds \equiv 4\pi b_0 \neq 0.$$
Without loss of generality, we can assume that $b_0 > 0$. It follows that there exists a small positive constant $\delta$ such that

$$\int_0^{2\pi} \text{Im} \vec{b}(x - s\alpha) \cdot \alpha ds \geq 2\pi b_0$$

(12)

for any $x \in D$, which is defined by

$$D = \bigcup_{\beta \in \mathbb{Z}^n} \{x \in \mathbb{R}^n | |x - x_0 - 2\pi\beta - \alpha a| \leq 2\delta\}.$$

Fix an arbitrary $T > 0$. We construct a sequence $\{u_l\}_{l=1,2,3,\ldots}$ by

$$u_l(t, x) = \exp(i\phi_l(t, x))\psi(x),$$

$$\phi_l(t, x) = -l^2 t \alpha \cdot \alpha + l\alpha \cdot x - \frac{1}{2} \int_0^{2l(t-T)} \vec{b}(x - \alpha s) \cdot \alpha ds,$$

where the amplitude function $\psi$ is a smooth function on $\mathbb{T}^n$ and supported on $D/2\pi\mathbb{Z}^n$. It is easy to see that

$$||u_l(T)|| = 1, \quad ||u_l(0)|| = O(\exp(-lb_0T)),$$

$$||Lu_l(t)|| = O(\exp(-lb_0(T-t)/2)),$$

which means that the energy inequality fails for $\{u_l\}$.

Next we give the sketch of the proof $2\Rightarrow 3$ in case $n \geq 2$. Suppose (7). Since $\text{Im} \vec{b} \in (C(\mathbb{T}^n))^n$, $\text{Im} \vec{b}(x)$ is represented by a Fourier series

$$\text{Im} \vec{b}(x) = \sum_{\beta \in \mathbb{Z}^n} \vec{b}_{\beta} e^{i\beta \cdot x}, \quad \vec{b}_{\beta} \in \mathbb{C}^n.$$  

(13)

The substitution of (13) into (7) gives

$$0 = \sum_{\beta \in \mathbb{Z}^n} \vec{b}_{\beta} \cdot \alpha e^{i\beta \cdot x} \int_0^{2\pi} e^{-i\alpha \cdot \beta s} ds = 2\pi \sum_{\beta \cdot \alpha = 0} \vec{b}_{\beta} \cdot \alpha e^{i\beta \cdot x}.$$

(14)

Then it follows that $\vec{b}_{\beta} \cdot \alpha = 0$ for any $\alpha \in \mathbb{Z}^n$. Since the orthogonal complement of $\beta \neq 0$ is spanned by some $\alpha^1, \ldots, \alpha^{n-1} \in \mathbb{Z}^n$, there exists $a_\beta \in \mathbb{C}$ such that $\vec{b}_{\beta} = a_\beta \beta$ for $\beta \neq 0$. On the other hand, (14) implies $\vec{b}_{0} = 0$ since $V_0 = \mathbb{R}^n$ is spanned by $e_1, \ldots, e_n \in \mathbb{Z}^n$. Then we have

$$\text{Im} \vec{b}(x) = \sum_{\beta \neq 0} a_\beta \beta e^{i\beta \cdot x}.$$  

If we set

$$\phi(x) = -i \sum_{\beta \neq 0} a_\beta e^{i\beta \cdot x},$$

then $\nabla \phi(x) = \text{Im} \vec{b}(x)$. 

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It is easy to prove $3 \Rightarrow 1$. Since $\exp(\pm \phi(x)/2)$ is a smooth function on $\mathbb{T}^n$, a mapping $u \mapsto v = \exp(-\phi(x)/2)u$ is automorphic on $L^2(\mathbb{T}^n)$. Multiplying $Lu = f$ by $\exp(\phi(x)/2)$, we have

$$(\partial_t - i\Delta + \text{Re} \tilde{b}(x) \cdot \nabla + \tilde{c}(x))v = g(t, x),$$

where $\tilde{c}(x) \in C^\infty(\mathbb{T})$ and $g(t, x) = \exp(-\phi(x)/2)f(t, x)$. It is easy to obtain forward and backward energy inequalities in $t$. The duality arguments proves that $(1)-(2)$ is $L^2$-well-posed.

Proof of Theorem 4. We will construct a sequence which fails to satisfy (19). It suffices to do it for one dimensional case since a one dimensional counter example is also an any dimensional counter example. Suppose that there exists a nonconstant holomorphic function $G(u)$ in $\mathbb{C}$ such that for $u \in C^1(\mathbb{T})$

$$F(u, \nabla u, \bar{u}, \nabla \bar{u}) = \frac{\partial}{\partial x}G(u) = G'(u)u_x.$$ 

Set $g = G'$ for short. If $u$ is a smooth solution to (3), then

$$\frac{d}{dt} \int_{\mathbb{T}} u(t, x)dx = \int_{\mathbb{T}} \partial_t u(t, x)dx = \int_{\mathbb{T}} \frac{\partial}{\partial x} \{u_x(t, x) + G(u(t, x))\}dx = 0. \quad (16)$$

We here express $u$ by a Fourier series

$$u(t, x) = \sum_{l \in \mathbb{Z}} u_l(t)e^{ilx}.$$ 

Then (16) implies $u_0(t) = u_0(0)$. Set $u_0(0) = z_0$ and $v(t, x) = u(t, x) - z_0$ for short. Since $g(0) = 0$ and $u_x = v_x$, there exists an appropriate complex constant $z_0$ such that

$$g(u)u_x = -(\mu + i\lambda)v_x + h(v)v_x,$$

where $\mu \in \mathbb{R}$, $\lambda > 0$, and $h$ is holomorphic in $\mathbb{C}$. Then, $v$ solves

$$v_t - iv_{xx} + (\mu + i\lambda)v_x = h(v)v_x.$$ 

In what follows, fix $z_0$. Note that $u(t, x) \equiv z_0$ is a solution to (3)-(4).

Suppose that the conclusion of Theorem 3 holds. Consider the initial value problem of the form $v^{(m)}$ solves the initial value problem of the form

$$v_t^{(m)} - iv_{xx}^{(m)} + (\mu + i\lambda)v_x^{(m)} = h(v^{(m)})v_x^{(m)} \quad \text{in} \quad (0, T) \times \mathbb{T}, \quad (17)$$

$$v^{(m)}(0, x) = \frac{e^{imx}}{(1 + m)^s} \quad \text{in} \quad \mathbb{T}, \quad (18)$$
where $s > 5/2$, $m = 1, 2, 3, \ldots$. Since $\{v^{(m)}(0, x)\}$ is bounded in $H^s(T)$ and

$$v^{(m)}(0, x) \to 0 \text{ in } H^s(T) \text{ as } m \to \infty$$

for any $\sigma < s$, it follows from the hypothesis that

$$v^{(m)} \to 0 \text{ in } C([0, T]; H^s(T)) \text{ as } m \to \infty$$

for any $\sigma < s$. We investigate a formal Fourier series solution to (17)-(18) of the form

$$w^{(m)}(t, x) = \sum_{l=1}^{\infty} w_{l}^{(m)}(t)e^{ilmx}. \quad (20)$$

The substitution of (20) into (17)-(18) gives

$$\frac{d}{dt}w_{l}^{(m)}(t) + (il^2m^2 + i\mu lm - \lambda lm)w_{l}^{(m)}(t)$$

$$= \sum_{p=1}^{\infty} h_p \sum_{l_0+\cdots+l_p=l, l_0, \ldots, l_p \geq 1} il_0 m \prod_{j=0}^{p} w_{l_j}^{(m)}(t), \quad (21)$$

$$w_{l}^{(m)}(0) = \begin{cases} (1 + m)^{-s} & \text{if } l = 1 \\ 0 & \text{otherwise} \end{cases} \quad (22)$$

For $l = 1$, (21)-(22) is concretely solved by

$$w_{1}^{(m)}(t) = (1 + m)^{-s} \exp(-i(m^2 + \mu m)t + \lambda mt). \quad (23)$$

For $l \geq 2$, we apply the idea of the abstract Cauchy-Kowalewski theorem to (21)-(22). We can show that there exists $T_m \in (0, T)$ such that the formal series (20) converges in $C([0, T_m]; H^s(T))$. Then it follows from the hypothesis that

$$v^{(m)} = w^{(m)} \text{ in } C([0, T_m]; H^s(T)).$$

Finally we can find $\delta > 0$, $\alpha \in (0, 1)$ and $t_m \in (0, T_m)$ such that

$$\sup_{t \in [0, T]} \|v^{(m)}(t)\|_{(1-\alpha)s} \geq \|v^{(m)}(t_m)\|_{(1-\alpha)s}$$

$$= \|w^{(m)}(t_m)\|_{(1-\alpha)s}$$

$$= \left(\sum_{l=1}^{\infty} (1 + lm)^{2(1-\alpha)s} |w_l^{(m)}(t_m)|^2\right)^{1/2}$$

$$\geq (1 + m)^{(1-\alpha)s}|w_1^{(m)}(t_m)|$$

$$= (1 + m)^{-\alpha} \exp(\lambda mt_m)$$

$$= \delta,$$

which contradicts (19). Here we omit the detail. \qed
References


