# NOTE ON SINGULAR PERTURBATION FOR ABSTRACT DIFFERENTIAL EQUATIONS 

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The present paper is concerned with abstract differential equations in a Banach space containing a small parameter in its coefficient

$$
\begin{equation*}
d u_{\varepsilon}(t) / d t+A_{\varepsilon}(t) u_{\varepsilon}(t)=f_{\varepsilon}(t) . \tag{0.1}
\end{equation*}
$$

As $\varepsilon \downarrow 0$ (0.1) degenerates to

$$
\begin{equation*}
d u_{0}(t) / d t+A_{0}(t) u_{0}(t)=f_{0}(t), \tag{0.2}
\end{equation*}
$$

where $A_{0}(t)$ is weaker than $A_{\varepsilon}(t), \varepsilon>0$, in the sense usually employed. We shall be interested in the behaviour of the solution $u_{\varepsilon}(t)$ of ( 0.1 ) as $\varepsilon \downarrow 0$, chiefly in the pointwise convergence of $u_{\mathrm{s}}(t)$ to the solution $u_{0}(t)$ of (0.2). The main theorem of section 2 is concerned with a sufficient condition in order that not only $u_{\varepsilon}(t)$ but also $A_{\varepsilon}(t) u_{\varepsilon}(t)$ and $d u_{\varepsilon}(t) / d t$ converge to their corresponding limits in the weak topology for each fixed $t$. It is almost essential that the limit equation (0.2) is well posed, which should be admitted to be a restrictive assumption.

In section 3 an example to which the above theorem can be applied is considered making frequent use of T. Kato's results on maximal accretive operators ([1], [2], [3]). This example is the initial-boundary value problem for the equation with coefficients having a singularity along $x=t$

$$
\frac{\partial u}{\partial t}-\varepsilon \frac{\partial^{2} u}{\partial x^{2}}+\frac{u}{(x-t)^{2}}=f, \quad a<x<b
$$

or

$$
\frac{\partial u}{\partial t}-\varepsilon \frac{\partial^{2} u}{\partial x^{2}}+\frac{\varepsilon}{x-t} \frac{\partial u}{\partial x}+\frac{u}{(x-t)^{2}}=f, \quad a<x<b
$$

with the boundary condition $u(t, a)=u(t, b)=0$, and was first motivated

[^0]by the construction of an example to which the main result of [4] on the initial value problem for the evolution equation
\[

$$
\begin{equation*}
d u(t) / d t+A(t) u(t)=f(t) \tag{0.3}
\end{equation*}
$$

\]

can be applied although $A(t)^{\alpha}$ has a variable domain whenever $\alpha>0$.
As a preparation a theorem on the unique solvability of the initial value problem for ( 0.3 ) is given in section 1 assuming among other things that

$$
\begin{equation*}
A(t)^{\rho} \cdot d A(t)^{-1} / d t \quad \text { is bounded and continuous in } t \tag{0.4}
\end{equation*}
$$

for some $\rho>0$. This hypothesis which implies

$$
\left\|\frac{\partial}{\partial t}(\lambda-A(t))^{-1}\right\| \leqq \frac{C}{|\lambda|^{p}},
$$

makes it possible to weaken the smoothness assumption of $A(t)$ as was made in [4], namely it enables us to remove the Hoelder continuity of $d A(t)^{-1} / d t$. It is a little interesting to note that ( 0.4 ) with $\rho=1$ implies the independence of the domain of $A(t)$ while if $\rho<1$ there exists a simple example for which $D\left(A(t)^{\alpha}\right)$ is not independent of $t$ whenever $\alpha>0$ although (0.4) is true.

1. We begin with a variant of the main theorem of [4]. By $D(A)$ and $R(A)$ we denote the domain and the range of an operator $A$.

Theorem 1.1. For each $t \in[0, T] A(t)$ is a densely defined, closed linear operator in a Banach space $X$. Let $A(t)$ satisfy the following assumptions:
(I) For each $t \in[0, T]$ the resolvent set of $A(t)$ contains a fixed closed angular domain

$$
\Sigma=\left\{\lambda: \arg \lambda \notin\left(-\theta_{0}, \theta_{0}\right)\right\}
$$

where $\theta_{0}$ is a positive number satisfying $0<\theta_{0}<\pi / 2$. The resolvent of A(t) satisfies

$$
\begin{equation*}
\left\|(\lambda-A(t))^{-1}\right\| \leqq M /|\lambda| \tag{1.1}
\end{equation*}
$$

for any $t \in[0, T]$ and $\lambda \in \sum$, where $M$ is a constant which is independent of $\lambda$ and $t$;
(II) $A(t)^{-1}$, which is bounded by (I), is continuously differentiable in $t$ in the uniform operator topology;
(III) There exists a positive number $\rho \leqq 1$ such that $R\left(d A(t)^{-1} / d t\right) \subset$ $D\left(A(t)^{\rho}\right)$ and $A(t)^{\rho} \cdot d A(t)^{-1} / d t$ is strongly continuous in $t \in[0, T]$. Hence
with some positive constant $N$ independent of $t$ we have

$$
\begin{equation*}
\left\|A(t)^{\rho} \frac{d}{d t} A(t)^{-1}\right\| \leqq N \tag{1.2}
\end{equation*}
$$

Then there exists a fundamental solution $U(t, s), 0 \leqq s \leqq t \leqq T$, to the equation

$$
\begin{equation*}
d u(t) / d t+A(t) u(t)=f(t) ; \tag{1.3}
\end{equation*}
$$

if $s<t, R(U(t, s)) \subset D(A(t))$ and $U(t, s)$ satisfies

$$
\begin{array}{r}
(\partial / \partial t) U(t, s)+A(t) U(t, s)=0, \quad 0 \leqq s<t \leqq T \\
U(s, s)=I . \tag{1.5}
\end{array}
$$

There exists a positive constant $C_{0}$ such that

$$
\begin{equation*}
\left\|\frac{\partial}{\partial t} U(t, s)\right\|=\|A(t) U(t, s)\| \leqq \frac{C_{0}}{t-s} \tag{1.6}
\end{equation*}
$$

If $f(t)$ is strongly Hoelder continuous, then the unique solution of (1.3) in $s<t \leqq T$ satisfying the initial condition $u(s)=u$ is given by

$$
\begin{equation*}
u(t)=U(t, s) u+\int_{s}^{t} U(t, \sigma) f(\sigma) d \sigma \tag{1.7}
\end{equation*}
$$

Proof. In what follows $C_{1}, C_{2}, \cdots, C_{8}$ denote constants which depend only on $\theta_{0}, M, \rho, N$ and $T$. First we note that (III) implies

$$
\begin{equation*}
\left\|\frac{\partial}{\partial t}(\lambda-A(t))^{-1}\right\| \leqq \frac{C_{1}}{|\lambda|^{\rho}}, \tag{1.8}
\end{equation*}
$$

which is a consequence of the formula

$$
\begin{align*}
& \quad(\partial / \partial t)(\lambda-A(t))^{-1} \\
& =-A(t)(\lambda-A(t))^{-1} \cdot d A(t)^{-1} / d t \cdot A(t)(\lambda-A(t))^{-1}  \tag{1.9}\\
& =-A(t)^{1-\rho}(\lambda-A(t))^{-1} \cdot A(t)^{\rho} d A(t)^{-1} / d t \cdot A(t)(\lambda-A(t))^{-1}
\end{align*}
$$

and the inequality

$$
\begin{equation*}
\left\|A(t)^{1-\rho}(\lambda-A(t))^{-1}\right\| \leqq C_{2} /|\lambda|^{\rho} . \tag{1.10}
\end{equation*}
$$

Hence just as in [4], it is possible to construct the fundamental solution by means of E. E. Levi's method:

$$
\begin{align*}
& U(t, s)=\exp (-(t-s) A(t))+\int_{s}^{t} \exp (-(t-\tau) A(t)) R(\tau, s) d \tau  \tag{1.11}\\
& R_{1}(t, s)=-(\partial / \partial t+\partial / \partial s) \exp (-(t-s) A(t))
\end{align*}
$$

$$
\begin{aligned}
& R(t, s)=\sum_{m=1}^{\infty} R_{m}(t, s) \\
& R_{m}(t, s)=\int_{s}^{t} R_{1}(t, \sigma) R_{m-1}(\sigma, s) d \sigma, \quad m=2,3, \cdots
\end{aligned}
$$

$R_{1}(t, s)$ may also be expressed as

$$
R_{1}(t, s)=-\frac{1}{2 \pi i} \int_{\Gamma} e^{-\lambda(t-s)} \frac{\partial}{\partial t}(\lambda-A(t))^{-1} d \lambda
$$

where $\Gamma$ is a smooth path running in $\sum$ from $\infty e^{-\theta_{0}{ }^{i}}$ to $\infty e^{\theta_{0}{ }^{i}}$. We have

$$
\begin{equation*}
\left\|R_{1}(t, s)\right\| \leqq \frac{C_{3}}{(t-s)^{1-\rho}}, \quad\|R(t, s)\| \leqq \frac{C_{4}}{(t-s)^{1-\rho}} \tag{1.12}
\end{equation*}
$$

Lemma 1.1. If $s<t$, we have

$$
\begin{equation*}
\|A(t) \exp (-(t-s) A(t))-A(s) \exp (-(t-s) A(s))\| \leqq \frac{C_{5}}{(t-s)^{1-\rho}} . \tag{1.13}
\end{equation*}
$$

Proof. By (1.8)

$$
\left\|(\lambda-A(t))^{-1}-(\lambda-A(s))^{-1}\right\| \leqq \frac{C_{1}(t-s)}{|\lambda|^{\rho}},
$$

hence the right member of (1.13) which is equal to

$$
\left\|\frac{-1}{2 \pi i} \int_{\Gamma} \lambda e^{-\lambda(t-s)}\left\{(\lambda-A(t))^{-1}-(\lambda-A(s))^{-1}\right\} d \lambda\right\|
$$

is dominated by

$$
\frac{C_{1}(t-s)}{2 \pi} \int_{\Gamma}|\lambda|^{1-\rho} e^{-(t-s) R e \lambda}|d \lambda| \leqq \frac{C_{5}}{(t-s)^{1-\rho}} .
$$

Lemma 1.2. If $0<\beta<\rho$ and $s<t$, then $A(t)^{\beta} R(t, s)$ is bounded and

$$
\begin{equation*}
\left\|A(t)^{\beta} R(t, s)\right\| \leqq \frac{C_{6}}{(t-s)^{1-\rho+\beta}} \tag{1.14}
\end{equation*}
$$

Proof. First we show that

$$
\begin{equation*}
\left\|A(t)^{\beta} R_{1}(t, s)\right\| \leqq \frac{C_{7}}{(t-s)^{1-\rho+\beta}} \tag{1.15}
\end{equation*}
$$

(1.15) is a consequence of

$$
A(t)^{\beta} R_{1}(t, s)=-\frac{1}{2 \pi i} \int_{\Gamma} e^{-\lambda(t-s)} A(t)^{\beta} \frac{\partial}{\partial t}(\lambda-A(t))^{-1} d \lambda
$$

$$
\begin{gathered}
\left\|A(t)^{\beta}(\partial / \partial t)(\lambda-A(t))^{-1}\right\| \\
=\left\|A(t)^{1+\beta-\rho}(\lambda-A(t))^{-1} \cdot A(t)^{\rho} d A(t)^{-1} / d t \cdot A(t)(\lambda-A(t))^{-1}\right\| \leqq C_{8} /|\lambda|^{\rho-\beta} .
\end{gathered}
$$

(1.14) follows from (1.15), (1.12) and

$$
A(t)^{\beta} R(t, s)=A(t)^{\beta} R_{1}(t, s)+\int_{s}^{t} A(t)^{\beta} R_{1}(t, \tau) R(\tau, s) d \tau
$$

According to the above two lemmas we may write

$$
\begin{gathered}
A(t) U(t, s)=A(t) \exp (-(t-s) A(t)) \\
+\int_{s}^{t}\{A(t) \exp (-(t-\tau) A(t))-A(\tau) \exp (-(t-\tau) A(\tau))\} R(\tau, s) d \tau \\
+\int_{s}^{t} A(\tau)^{1-\beta} \exp (-(t-\tau) A(\tau)) A(\tau)^{\beta} R(\tau, s) d \tau
\end{gathered}
$$

The inequality (1.6) is a simple consequence of (1.13), (1.14) as well as the above formula. The remaining part of the proof is the same as the argument of [4].

Remark. The assumption (III) enables us to remove the Hoelder continuity of $d A(t)^{-1} / d t$ in $t$ which was used in [4] when we proved that $U(t, s)$ satisfies (1.4) in the strict sense.
2. Singular perturbation. Letting $A(t), 0 \leqq t \leqq T$, be a family of linear closed operators in $X$,

Definition 1. $u(t)$ is called a strict solution of

$$
\begin{gather*}
d u(t) / d t+A(t) u(t)=f(t), \quad s<t \leqq T,  \tag{2.1}\\
u(s)=u \tag{2.2}
\end{gather*}
$$

in $(s, T]$ if
(1) $u(t)$ is strongly continuous in the closed interval $[s, T]$ and strongly continuously differentiable in the left open interval $(s, T]$,
(2) for each $t \in(s, T], u(t) \in D(A(t))$,
(3) $u(t)$ satisfies (2.1)-(2.2);

Definition 2. $u(t)$ is called a weak solution of (2.1)-(2.2) in $(s, T]$ if
(1) $u(t)$ is weakly continuous in $[s, T]$,
(2) $u(t)$ satisfies

$$
\int_{s}^{T}\left(u(t), \mathscr{P}^{\prime}(t)-A^{*}(t) \mathscr{P}(t)\right) d t+\int_{s}^{T}(f(t), \mathscr{P}(t)) d t+(u, \mathscr{P}(s))=0
$$

for any function $\varphi(t)$ with values in $X^{*}$ satisfying
(i) for each $t, \varphi(t) \in D\left(A^{*}(t)\right)$,
(ii) $\varphi(t), \varphi^{\prime}(t)=d \varphi(t) / d t$ and $A^{*}(t) \varphi(t)$ are strongly continuous in [ $s, T]$,
(iii) $\varphi(T)=0$.

The above definition of weak solution is slightly different from the one given in [4] where a weak solution was assumed to be strongly continuous.

Theorem 2.1. Suppose that $X$ be reflexive. Let $A_{\mathrm{s}}(t), 0 \leqq t \leqq T$, $0 \leqq \varepsilon \leqq \varepsilon_{0}$, be a family of closed linear operators in $X$. Suppose that the assumptions of Theorem 1.1 are satisfied by $A_{\mathrm{\varepsilon}}(t), 0 \leqq t \leqq T, 0<\varepsilon \leqq \varepsilon_{0}$, with constants $\theta_{0}, M, \rho$ and $N$ which are independent of $t$ and $\varepsilon$. We assume also that letting $A(t)$ stand for $A_{\mathrm{\varepsilon}_{\mathrm{J}}}(t)$
(a) $D\left(A_{\mathrm{\varepsilon}}(t)\right) \equiv D(A(t))$ and $D\left(A_{\mathrm{\varepsilon}}^{*}(t)\right) \equiv D\left(A^{*}(t)\right)$ do not depend on $\varepsilon$ if $0<\varepsilon \leqq \varepsilon_{0} ;$
(b) $D\left(A_{0}(t)\right) \supset D(A(t))$ and $D\left(A_{0}^{*}(t)\right) \supset D\left(A^{*}(t)\right)$;
(c) for each $\rho \in D\left(A^{*}(t)\right) A_{\varepsilon}^{*}(t) \varphi \rightarrow A_{0}^{*}(t) \mathcal{P}$ strongly in $X^{*}$ as $\varepsilon \downarrow 0$;
(d) $A_{\mathrm{\varepsilon}}(t) A(t)^{-1}, A_{0}(t) A_{\mathrm{s}}(t)^{-1}, A_{\varepsilon}^{*}(t) A^{*}(t)^{-1}$ and $A_{0}^{*}(t) A_{\varepsilon}^{*}(t)^{-1}$ are all uniformly bounded with respect to $\varepsilon$ and $t$ and are continuous in $t$ for each fixed $\varepsilon$ in the strong operator topology in $X$ or $X^{*}$.

If the initial value problem

$$
\begin{gather*}
d u(t) / d t+A_{0}(t) u(t)=f(t),  \tag{2.4}\\
u(t)=u_{0} \tag{2.5}
\end{gather*}
$$

has only one weak solution which is also a strict solution when $f(t)$ is strongly Hoelder continuous in $t$, then the solution $u_{z}(t)$ of the equation

$$
\begin{equation*}
d u_{\mathrm{e}}(t) / d t+A_{\mathrm{e}}(t) u_{\mathrm{e}}(t)=f_{\mathrm{z}}(t) \tag{2.6}
\end{equation*}
$$

converges to the solution of (2.4)-(2.5) in the following sense:

$$
\begin{aligned}
& \text { for each } t \in(0, T] u_{\varepsilon}(t) \rightarrow u(t), \quad A_{\varepsilon}(t) u_{\varepsilon}(t) \rightarrow A_{0}(t) u(t), \\
& d u_{\varepsilon}(t) / d t \rightarrow d u(t) / d t \text { all in the weak topology, }
\end{aligned}
$$

provided that
(i) $u_{\mathrm{e}}(0) \rightarrow u_{0}$ weakly,
(ii) $f_{\varepsilon}(t)$ is uniformly Hoelder continuous:

$$
\begin{equation*}
\left\|f_{\mathrm{e}}(t)-f_{\mathrm{z}}(s)\right\| \leqq F|t-s|^{\infty}, \quad F>0, \quad \alpha>0 \tag{2.7}
\end{equation*}
$$

where $F$ and $\alpha$ are independent of $\varepsilon$,
(iii) $f_{\mathrm{e}}(t)$ converges to a strongly Hoelder continuous function $f(t)$ uniformly in the weak topology.

In this section we denote by $C_{9}, C_{10}, \cdots$ constants which are dependent only on $\theta_{0}, M, \rho, N, T, F, \alpha, \sup _{\varepsilon}\left\|u_{\mathrm{e}}(0)\right\|$ and $\sup _{\varepsilon, t}\left\|f_{\mathrm{z}}(t)\right\|$.

Proof. If

$$
\begin{equation*}
U_{\mathrm{e}}(t, s)=\exp \left(-(t-s) A_{\mathrm{e}}(t)\right)+W_{\mathrm{e}}(t, s) \tag{2.8}
\end{equation*}
$$

is the fundamental solution of (2.6), then by Theorem 1.1 there exists a constant $C$ which is independent to $t, s$ and $\varepsilon$ such that

$$
\begin{align*}
& \left\|\frac{\partial}{\partial t} U_{\mathrm{e}}(t, s)\right\|=\left\|A_{\mathrm{e}}(t) U_{\mathrm{\varepsilon}}(t, s)\right\| \leqq \frac{C}{t-s},  \tag{2.9}\\
& \left\|\left(\frac{\partial}{\partial t}+\frac{\partial}{\partial s}\right) \exp \left(-(t-s) A_{\mathrm{e}}(t)\right)\right\| \leqq \frac{C}{(t-s)^{1-\rho}},  \tag{2.10}\\
& \left\|\frac{\partial}{\partial t} \exp \left(-(t-s) A_{\mathrm{e}}(t)\right)\right\| \leqq \frac{C}{t-s},  \tag{2.11}\\
& \left\|\frac{\partial}{\partial t} W_{\mathrm{e}}(t, s)\right\| \leqq \frac{C}{(t-s)^{1-\rho}}, \quad\left\|A_{\mathrm{e}}(t) W_{\mathrm{e}}(t, s)\right\| \leqq \frac{C}{(t-s)^{1-\rho}},  \tag{2.12}\\
& \left\|A_{\mathrm{e}}(t)^{\gamma} \exp \left(-(t-s) A_{\mathrm{e}}(t)\right)\right\| \leqq \frac{C}{(t-s)^{\gamma}}, \quad 0 \leqq \gamma \leqq 1 . \tag{2.13}
\end{align*}
$$

By the formula

$$
\begin{align*}
& u_{\mathrm{e}}(t)=U_{\mathrm{e}}(t, 0) u_{\mathrm{e}}(0)+\int_{0}^{t} U_{\mathrm{\varepsilon}}(t, \sigma) f_{\mathrm{z}}(\sigma) d \sigma,  \tag{2.14}\\
& \frac{\partial}{\partial t} u_{\mathrm{\varepsilon}}(t)=\frac{\partial}{\partial t} U_{\mathrm{e}}(t, 0) u_{\mathrm{e}}(0)+\int_{0}^{t} \frac{\partial}{\partial t} \exp \left(-(t-\sigma) A_{\mathrm{e}}(t)\right)\left(f_{\mathrm{\varepsilon}}(\sigma)-f_{\mathrm{e}}(t)\right) d \sigma \\
& +\int_{0}^{t}\left(\frac{\partial}{\partial t}+\frac{\partial}{\partial \sigma}\right) \exp \left(-(t-\sigma) A_{\mathrm{e}}(t)\right) d \sigma \cdot f_{\mathrm{e}}(t)  \tag{2.15}\\
& +\exp \left(-(t-s) A_{\mathrm{e}}(t)\right) f_{\mathrm{\varepsilon}}(t)+\int_{0}^{t} \frac{\partial}{\partial t} W_{\mathrm{e}}(t, \sigma) f_{\mathrm{\varepsilon}}(\sigma) d \sigma
\end{align*}
$$

as well as $(2.7) \sim(2.13)$ we immediately see that

$$
\begin{equation*}
\left\|u_{\mathrm{e}}(t)\right\| \leqq C_{9}, \quad\left\|\frac{\partial}{\partial t} u_{\mathrm{e}}(t)\right\| \leqq \frac{C_{9}}{t}, \quad\left\|A_{\mathrm{e}}(t) u_{\mathrm{e}}(t)\right\| \leqq \frac{C_{9}}{t} . \tag{2.16}
\end{equation*}
$$

If $\varphi(s)$ is an arbitrary function with values in $X^{*}$ such that $\varphi(s) \in D\left(A^{*}(s)\right)$ for each $s$ (recall that $A(s)=A_{\varepsilon_{0}}(s)$ ) and $\varphi(s), d \varphi(s) / d s=\varphi^{\prime}(s), A^{*}(s) \varphi(s)$ are all strongly continuous in $0 \leqq s \leqq T$, then $A_{\varepsilon}^{*}(s) \varphi(s)=A_{\mathrm{e}}^{*}(s) A^{*}(s)^{-1}$ $\times A^{*}(s) \varphi(s)$ is also strongly continuous and for each $t$

$$
\begin{align*}
& \left(u_{\mathrm{e}}(t), \varphi(t)\right)-\left(u_{\mathrm{e}}(0), \varphi(0)\right)-\int_{0}^{t}\left(u_{\mathrm{e}}(s), \varphi^{\prime}(s)\right) d s  \tag{2.17}\\
& +\int_{0}^{t}\left(u_{\mathrm{e}}(s), A_{\mathrm{\varepsilon}}^{*}(s) \varphi(s)\right) d s=\int_{0}^{t}\left(f_{\mathrm{z}}(s), \varphi(s)\right) d s .
\end{align*}
$$

For $1<p<\infty$ let $L^{p}(0, T ; X)$ be the space of all measurable functions with values in $X$ in $0<t<T$ for which $\|u(t)\| \in L^{p}(0, T)$. By Theorem 5.7 of $[6],\left(L^{p}(0, T ; X)\right)^{*}=L^{p^{\prime}}\left(0, T ; X^{*}\right)$ where $p^{-1}+{p^{\prime-1}}^{\prime}=1$, hence $L^{p}(0, T ; X)$ is reflexive. Since $\left\{u_{\varepsilon}\right\}$ is a bounded sequence in $L^{p}(0, T ; X)$ by (2.16), it contains a subsequence $\left\{u_{\varepsilon_{i}}\right\}$ which converges weakly to some function $u \in L^{p}(0, T ; X)$. Replacing $\varepsilon$ by $\varepsilon_{i}$ in (2.17) and then letting $i \rightarrow \infty$, we get

$$
\begin{align*}
& \lim _{i \rightarrow \infty}\left(u_{\Sigma_{i}}(t), \varphi(t)\right)-\left(u_{0}, \varphi(0)\right)-\int_{0}^{t}\left(u(s), \varphi^{\prime}(s)\right) d s \\
& \quad+\int_{0}^{t}\left(u(s), A_{0}^{*}(s) \varphi(s)\right) d s=\int_{0}^{t}(f(s), \varphi(s)) d s \tag{2.18}
\end{align*}
$$

Choosing $A^{*}(s)^{-1} \varphi$ as $\varphi(s)$ in (2.18) with an arbitrary $\varphi \in X^{*}$, which is possible by the assumptions, we conclude that $\lim _{i \rightarrow \infty}\left(u_{\varepsilon_{i}}(t), A^{*}(t)^{-1} \varphi\right)$ exists. Since $A^{*}(t)^{-1} \varphi$ is an arbitrary element of $D\left(A^{*}(t)\right)$ which is dense in $X^{*}$ and $u_{\varepsilon_{i}}(t)$ is bounded by (2.16), it follows that $u_{\varepsilon_{i}}(t)$ converges weakly to some element $v(t)$ satisfying

$$
\begin{gather*}
\|v(t)\| \leqq C_{9} \\
(v(t), \varphi(t))-\left(u_{0}, \varphi(0)\right)-\int_{0}^{t}\left(u(s), \varphi^{\prime}(s)\right) d s  \tag{2.19}\\
+\int_{0}^{t}\left(u(s), A_{0}^{*}(s) \varphi(s)\right) d s=\int_{0}^{t}(f(s), \varphi(s)) d s
\end{gather*}
$$

Clearly

$$
\begin{equation*}
(v(t), \mathscr{P}(t))-(v(\tau), \mathscr{P}(\tau)) \rightarrow 0 \quad \text { as } \quad t-\tau \rightarrow 0 \tag{2.20}
\end{equation*}
$$

Choosing again $\varphi(s)=A^{*}(s)^{-1} \varphi$ and using (2.20)

$$
\begin{gather*}
\left(v(t)-v(\tau), A^{*}(\tau)^{-1} \mathcal{P}\right)=\left(v(t), A^{*}(t)^{-1} \mathcal{P}\right) \\
-\left(v(\tau), \quad A^{*}(\tau)^{-1} \mathcal{P}\right)-\left(v(t), A^{*}(t)^{-1} \mathcal{P}-A^{*}(\tau)^{-1} \mathcal{P}\right) \rightarrow 0 \quad \text { as } \quad t \rightarrow \tau . \tag{2.21}
\end{gather*}
$$

Using (2.19) and (2.21) and noting that $D\left(A^{*}(\tau)\right)$ is dense we conclude that $v(t)$ is weakly continuous in $0 \leqq t \leqq T$. If $\mathcal{P}$ is an arbitrary element of $L^{p^{\prime}}\left(0, T ; X^{*}\right)$, then

$$
\int_{0}^{T}\left(u_{2_{i}}(t)-v(t), \varphi(t)\right) d t \rightarrow 0 \quad \text { as } \quad i \rightarrow \infty
$$

because of the measurability of the integrand and of the well known
theorem on dominated convergence sequences. Thus $u$ and $v$ are both a weak limit of $\left\{u_{\varepsilon_{i}}\right\}$, which implies that $u(t) \equiv v(t)$ is weakly continuous in $0 \leqq s \leqq T$ and $u(0)=u_{0}$. When $s>0, A_{\varepsilon}(s) u_{\mathrm{e}}(s)$ is bounded by (2.16), and so is $A_{0}(s) u_{\varepsilon}(s)$ due to the assumed uniform boundedness of $A_{0}(s) A_{\mathrm{s}}(s)^{-1}$. It follows consequently that $u(s) \in D\left(A_{0}(s)\right)$ and

$$
\begin{align*}
& A_{0}(s) u_{\varepsilon_{i}}(s) \rightarrow A_{0}(s) u(s)  \tag{2.22}\\
& A_{\varepsilon_{i}}(s) u_{\varepsilon_{i}}(s) \rightarrow A_{0}(s) u(s) \tag{2.23}
\end{align*}
$$

both in the weak topology. Next if $\varphi(s)$ is an arbitrary function with values in $X^{*}$ such that $\varphi(s) \in D\left(A_{0}^{*}(s)\right)$ and $\varphi(s), \varphi^{\prime}(s), A_{0}^{*}(s) \varphi(s)$ are all strongly continuous in $0 \leqq s \leqq T$, then for any $\delta>0$, we have

$$
\begin{aligned}
& \left(u_{\varepsilon_{i}}(t), \varphi(t)\right)-\left(u_{\varepsilon_{i}}(\delta), \varphi(\delta)\right)-\int_{\delta}^{t}\left(u_{\varepsilon_{i}}(s), \varphi^{\prime}(s)\right) d s \\
& \quad+\int_{\delta}^{t}\left(A_{\varepsilon_{i}}(s) u_{\varepsilon_{i}}(s), \varphi(s)\right) d s=\int_{\delta}^{t}\left(f_{\varepsilon_{i}}(s), \varphi(s)\right) d s .
\end{aligned}
$$

Letting $i \rightarrow \infty$ and then $\delta \downarrow 0$, we get noting (2.23)

$$
\begin{aligned}
& (u(t), \varphi(t))-\left(u_{0}, \mathscr{P}(0)\right)-\int_{0}^{t}\left(u(s), \varphi^{\prime}(s)\right) d s \\
& \quad+\int_{0}^{t}\left(u(s), A_{0}^{*}(s) \varphi(s)\right) d s=\int_{0}^{t}(f(s), \varphi(s)) d s
\end{aligned}
$$

which shows that $u(t)$ is a weak solution of (2.4)-(2.5), hence by the assumption it is the unique strict solution of the same problem. Therefore it follows that the original sequence $\left\{u_{\varepsilon}\right\}$ itself converges to $u$ weakly in $L^{p}(0, T ; X)$. We furthermore conclude that for each $t \in(0, T]$ $A_{\mathrm{e}}(t) u_{\mathrm{e}}(t) \rightarrow A_{0}(t) u(t)$ in the weak topology of $X$, and hence also that $d u_{\mathrm{s}}(t) / d t \rightarrow d u(t) / d t$ in the same sense.
3. Example. As an application of Theorem 2.1 we consider the following example. Let $-\infty<a<0<T<b<\infty$ and for each $t \in[0, T]$

$$
V(t)=\left\{u \in L^{2}(a, b): \frac{d u}{d x}, \frac{u}{x-t} \in L^{2}(a, b), u(a)=u(b)=0\right\}
$$

where the derivatives in the above as well as in what follows are interpreted in the distribution sense. $a_{\mathrm{\varepsilon}}(t ; u, v)$ denotes a family of sesquilinear forms on $V(t) \times V(t)$ defined by either of

$$
\begin{align*}
& a_{\mathrm{e}}(t ; u, v)=\int_{a}^{b}\left\{\varepsilon \frac{d u}{d x} \frac{\overline{d v}}{d x}+\frac{u \bar{v}}{(x-t)^{2}}\right\} d x  \tag{1}\\
& a_{\mathrm{e}}(t ; u, v)=\int_{a}^{b}\left\{\varepsilon \frac{d u}{d x} \frac{\overline{d v}}{d x}+\varepsilon \frac{d u}{d x} \frac{\bar{v}}{x-t}+\frac{u \bar{v}}{(x-t)^{2}}\right\} d x \tag{2}
\end{align*}
$$

$A_{\mathrm{e}}(t)$ is an operator corresponding to $a_{\mathrm{e}}(t ; u, v)$ which is defined in the following usual manner:

$$
\begin{aligned}
& u \in V(t) \text { belongs to } D\left(A_{\mathrm{z}}(t)\right) \text { and } A_{\mathrm{z}}(t) u=f \in L^{2}(a, b) \\
& \text { if } a_{\mathrm{z}}(t ; u, v)=(f, v) \text { for every } v \in V(t) .
\end{aligned}
$$

Thus $A_{\mathrm{z}}(t)$ is a differential operator

$$
\begin{aligned}
& \left(A_{\varepsilon}(t) u\right)(x)=-\varepsilon \frac{d^{2} u}{d x^{2}}+\frac{u}{(x-t)^{2}} \quad \text { or } \\
& \left(A_{\mathbf{\varepsilon}}(t) u\right)(x)=-\varepsilon \frac{d^{2} u}{d x^{2}}+\frac{\varepsilon}{x-t} \frac{d u}{d x}+\frac{u}{(x-t)^{2}}
\end{aligned}
$$

restricted to some class of functions satisfying $u(a)=u(b)=0$. In the first case $A_{\mathrm{e}}(t)$ is positive definite while in the second case $A_{\mathrm{e}}(t)$ is not selfadjoint although it is regularly accretive in the terminology of Kato [1]. These two cases can be treated quite similarly and we shall confine ourselves to the second case in what follows. The adjoint form $a_{\mathrm{a}}^{*}(t ; u, v)$ of $a_{\mathrm{e}}(t ; u, v)$ is

$$
a_{\mathrm{e}}^{*}(t ; u, v)=\int_{a}^{b}\left\{\varepsilon \frac{d u}{d x} \frac{d \bar{v}}{d x}-\varepsilon \frac{d u}{d x} \frac{\bar{v}}{x-t}+(1+\varepsilon) \frac{u \bar{v}}{(x-t)^{2}}\right\} d x
$$

which is also defined on $V(t) \times V(t)$. As is easily seen we have

$$
\left|\operatorname{Im} a_{\mathrm{e}}(t ; u, u)\right| \leqq \frac{1}{2} \operatorname{Re} a_{\mathrm{e}}(t ; u, u)
$$

namely the index (Kato [1]) of $a_{\varepsilon}(t ; u, v)$ does not exceed $\frac{1}{2}$. Hence by Theorem 2.2 of [1] any complex number $\lambda$ with $|\arg \lambda|>\tan ^{-1} \frac{1}{2}$ belongs to the resolvent set of $A_{\mathrm{e}}(t)$ and

$$
\|\left(\lambda-A_{\varepsilon}(t)\right)^{-1}| | \leqq\left\{\begin{array}{l}
\left(|\lambda| \sin \left(|\arg \lambda|-\tan ^{-1} \frac{1}{2}\right)\right)^{-1} \\
\tan ^{-1} \frac{1}{2}<|\arg \lambda| \leqq \pi / 2+\tan ^{-1} \frac{1}{2} \\
|\lambda|^{-1}, \quad|\arg \lambda|>\pi / 2+\tan ^{-1} \frac{1}{2}
\end{array}\right.
$$

Thus (I) of Theorem 1.1 is satisfied by $\left\{A_{\mathrm{e}}(t)\right\}$ uniformly with respect to $t$ and $\varepsilon$. The real part of $a_{8}(t ; u, v)$ is

$$
\operatorname{Re} a_{\mathrm{e}}(t ; u, v)=\int_{a}^{b}\left\{\varepsilon \frac{d u}{d x} \frac{\bar{v}}{d x}+\left(1+\frac{\varepsilon}{2}\right) \frac{u \bar{v}}{(x-t)^{2}}\right\} d x
$$

its corresponding operator being denoted by $H_{\mathrm{e}}(t)$. It is also possible to express the solutions of $A_{\mathrm{e}}(t) u_{\mathrm{e}}(t)=g, H_{\mathrm{e}}(t) v_{\mathrm{e}}(t)=g$ and $A_{\mathrm{e}}^{*}(t) w_{\mathrm{\varepsilon}}(t)=g$ for given $g \in L^{2}(a, b)$ explicitly all of which may be written below for the sake of convenience:

$$
\begin{align*}
& u_{\varepsilon}(t, x)=\frac{1}{2 \varepsilon \sqrt{1+\varepsilon^{-1}}}\left\{\int_{x}^{t}(t-y)^{\sqrt{1+\varepsilon^{-1}}} g(y) d y(t-x)^{1-\sqrt{1+\varepsilon^{-1}}}\right. \\
& -\int_{a}^{t}(t-y)^{\sqrt{1+\varepsilon^{-1}}} g(y) d y(t-a)^{-2 \sqrt{1+\varepsilon^{-1}}}(t-x)^{1+\sqrt{1+\varepsilon^{-1}}} \\
& \left.+\int_{a}^{x}(t-y)^{-\sqrt{1+\varepsilon^{-1}}} g(y) d y(t-x)^{1+\sqrt{1+\varepsilon^{-1}}}\right\} \quad \text { if } \quad a \leqq x<t, \\
& u_{\mathrm{e}}(t, x)=\frac{1}{2 \varepsilon \sqrt{1+\varepsilon^{-1}}}\left\{\int_{t}^{x}(y-t)^{\sqrt{1+\varepsilon^{-1}}} g(y) d y(x-t)^{1-\sqrt{1+\varepsilon^{-1}}}\right.  \tag{3.1}\\
& -\int_{t}^{b}(y-t)^{\sqrt{1+\varepsilon^{-1}}} g(y) d y(b-t)^{-2 \sqrt{1+\varepsilon^{-1}}}(x-t)^{1+\sqrt{1+\varepsilon^{-1}}} \\
& \left.+\int_{x}^{b}(y-t)^{-\sqrt{1+\varepsilon^{-1}}} g(y) d y(x-t)^{1+\sqrt{1+\varepsilon^{-1}}}\right\} \quad \text { if } \quad t<x \leqq b ; \\
& v_{\mathrm{e}}(t, x)=\frac{1}{\varepsilon \sqrt{3+4 \varepsilon^{-1}}}\left\{\int_{x}^{t}(t-y)^{\frac{\sqrt{3+4 \varepsilon^{-1}}+1}{2}} g(y) d y(t-x)^{\frac{1-\sqrt{3+4 \varepsilon^{-1}}}{2}}\right. \\
& -\int_{a}^{t}(t-y)^{\frac{\sqrt{3+4 e^{-1}}+1}{2}} g(y) d y(t-a)^{-\sqrt{3+4 \varepsilon^{-1}}}(t-x)^{\frac{1+\sqrt{3+4 \varepsilon^{-1}}}{2}} \\
& \left.+\int_{a}^{x}(t-y)^{\frac{1-\sqrt{3+4 \varepsilon^{-1}}}{2}} g(y) d y(t-x)^{\frac{\sqrt{3+4 \varepsilon^{-1}}+1}{2}}\right\} \quad \text { if } \quad a \leqq x<t \text {, }  \tag{3.2}\\
& v_{\mathrm{\varepsilon}}(t, x)=\frac{1}{\varepsilon \sqrt{ } 3+4 \varepsilon^{-1}}\left\{\int_{x}^{t}(y-t)^{\frac{\sqrt{3+4 \varepsilon^{-1}}+1}{2}} g(y) d y(x-t)^{\frac{1-\sqrt{3+4 \varepsilon^{-1}}}{2}}\right. \\
& -\int_{t}^{b}(y-t)^{\frac{\sqrt{3+4 e^{-1}}+1}{2}} g(y) d y(b-t)^{-\sqrt{3+4 e^{-1}}}(x-t)^{\frac{1+\sqrt{3+4 \mathrm{e}^{-1}}}{2}} \\
& \left.+\int_{x}^{b}(y-t)^{\frac{1-\sqrt{3+4 e^{-1}}}{2}} g(y) d y(x-t)^{\frac{\sqrt{3+4 e^{-1}}+1}{2}}\right\} \quad \text { if } \quad t<x \leqq b ; \\
& w_{\mathrm{e}}(t, x)=\frac{1}{2 \varepsilon \sqrt{1+\varepsilon^{-1}}}\left\{\int_{x}^{t}(t-y)^{\sqrt{1+\varepsilon^{-1}}+1} g(y) d y(t-x)^{-\sqrt{1+\varepsilon^{-1}}}\right. \\
& -\int_{a}^{t}(t-y)^{\sqrt{1+\mathrm{e}^{-1}}+1} g(y) d y(t-a)^{-2 \sqrt{1+\varepsilon^{-1}}}(t-x)^{\sqrt{1+\varepsilon^{-1}}} \\
& \left.+\int_{a}^{x}(t-y)^{1-\sqrt{1+\varepsilon^{-1}}} g(y) d y(t-x)^{\sqrt{1+\varepsilon^{-1}}}\right\} \quad \text { if } \quad a \leqq x<t,  \tag{3.3}\\
& w_{\mathrm{e}}(t, x)=\frac{1}{2 \varepsilon \sqrt{1+\varepsilon^{-1}}}\left\{\int_{t}^{x}(y-t)^{\sqrt{1+\varepsilon^{-1}}+1} g(y) d y(x-t)^{-\sqrt{1+\varepsilon^{-1}}}\right. \\
& -\int_{t}^{b}(y-t)^{\sqrt{1+\varepsilon^{-1}}+1} g(y) d y(b-t)^{-2 \sqrt{1+\varepsilon^{-1}}}(x-t)^{\sqrt{1+\varepsilon^{-1}}} \\
& \left.+\int_{x}^{b}(y-t)^{1-\sqrt{1+\varepsilon^{-1}}} g(y) d y(x-t)^{\sqrt{1+\varepsilon^{-1}}}\right\} \quad \text { if } \quad t<x \leqq b .
\end{align*}
$$

Using (3.1), (3.2) and (3.3) as well as the following two inequalities in p. 245 of [5]:

$$
\begin{align*}
& \int_{0}^{\infty} y^{p(\alpha-1)}\left(\int_{0}^{y} x^{-\infty} f(x) d x\right)^{p} d y<\left(\frac{p}{p-\alpha p-1}\right)^{p} \int_{0}^{\infty} f^{p} d x  \tag{3.4}\\
& \int_{0}^{\infty} x^{-\infty p^{\prime}}\left(\int_{x}^{\infty} y^{\alpha-1} g(y) d y\right)^{p^{\prime}} d x<\left(\frac{p^{\prime}}{1-\alpha p^{\prime}}\right)^{p^{\prime}} \int_{0}^{\infty} g^{p^{\prime}} d x \tag{3.5}
\end{align*}
$$

where $\alpha<1 / p^{\prime}$ we can prove that if $0<\varepsilon<4 / 5$

$$
\begin{gathered}
D\left(A_{\mathfrak{e}}(t)\right)=D\left(H_{\mathfrak{e}}(t)\right)=D\left(A_{\mathrm{e}}^{*}(t)\right) \\
=\left\{u \in L^{2}(a, b): \frac{d^{2} u}{d x^{2}}, \frac{1}{x-t} \frac{d u}{d x}, \frac{u}{(x-t)^{2}} \in L^{2}(a, b), u(a)=u(b)=0\right\},
\end{gathered}
$$

which shows that (a) and (b) of the assumptions of Theorem 2.1 are satisfied where $\left(A_{0}(t) u\right)(x)=u(x) /(x-t)^{2}$ in the present case. Similarly we can show that

$$
\begin{align*}
& \frac{d u_{e}(t)}{d t} \in V(t), \\
& \left.\varepsilon \sqrt{\int_{a}^{b} \left\lvert\, \frac{1}{x-t}\right.} \frac{\partial}{\partial x} u_{\mathrm{q}}(t, x)\right|^{2} d x ~ \leqq K \sqrt{\varepsilon}\|g\|,  \tag{3.6}\\
& \sqrt{\int_{a}^{b}\left|\frac{u_{\varepsilon}(t, x)}{(x-t)^{2}}\right|^{2} d x} \leqq K\|g\|,  \tag{3.7}\\
& \left\|A_{\mathrm{e}}(t) H_{\mathrm{e}}(t)^{-1}-I\right\| \leqq K \sqrt{\varepsilon}, \tag{3.8}
\end{align*}
$$

where $K$ is a constant which does not depend on $t$ and $\varepsilon$. By a general result on sesquilinear forms ([1]) we get after an integration by part

$$
\begin{align*}
& \left\|H_{\mathrm{z}}(t)^{\frac{1}{2}} \frac{d}{d t} u_{\mathrm{e}}(t)\right\|^{2}=\operatorname{Re} a_{\mathrm{z}}\left(t ; \frac{d}{d t} u_{\mathrm{z}}(t), \frac{d}{d t} u_{\mathrm{e}}(t)\right) \\
& \quad=\varepsilon \int_{a}^{b}\left|\frac{\partial}{\partial x} \frac{\partial u_{\mathrm{z}}}{\partial t}\right|^{2} d x+\left(1+\frac{\varepsilon}{2}\right) \int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathrm{q}}}{\partial t}\right|^{2} d x . \tag{3.9}
\end{align*}
$$

Differentiating both sides of

$$
-\varepsilon \frac{\partial^{2}}{\partial x^{2}} u_{\mathrm{z}}(t, x)+\frac{\varepsilon}{x-t} \frac{\partial}{\partial x} u_{\mathrm{z}}(t, x)+\frac{u_{\mathrm{e}}(t, x)}{(x-t)^{2}}=g(x)
$$

in $t$ we obtain

$$
-\varepsilon \frac{\partial^{2}}{\partial x^{2}} \frac{\partial u_{\varepsilon}}{\partial t}+\frac{\varepsilon}{x-t} \frac{\partial}{\partial x} \frac{\partial u_{\varepsilon}}{\partial t}+\frac{1}{(x-t)^{2}} \frac{\partial u_{\varepsilon}}{\partial t}=-\frac{\varepsilon}{(x-t)^{2}} \frac{\partial u_{\varepsilon}}{\partial x}-\frac{2 u_{\varepsilon}}{(x-t)^{3}}
$$

Multiplying both members of the above relation by $\overline{\partial u_{\mathrm{z}} / \partial t}$, and integrating the resulting equality by part over $(a, b)$, and using the formula

$$
\operatorname{Re} \int_{a}^{b} \frac{1}{x-t} \frac{d u}{d x} \bar{u} d x=\frac{1}{2} \int_{a}^{b}\left|\frac{u}{x-t}\right|^{2} d x
$$

which holds for $u \in V(t)$ and may be proved by integration by part, and finally comparing the real parts of both sides of the relation thus derived, we get

$$
\begin{gather*}
\varepsilon \int_{a}^{b}\left|\frac{\partial}{\partial x} \frac{\partial u_{\mathrm{\varepsilon}}}{\partial t}\right|^{2} d x+\left(1+\frac{\varepsilon}{2}\right) \int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathrm{\varepsilon}}}{\partial t}\right|^{2} d x \\
=-\varepsilon R e \int_{a}^{b} \frac{1}{(x-t)^{2}} \frac{\partial u_{\mathrm{z}}}{\partial x} \frac{\overline{\partial u_{\mathrm{e}}}}{\partial t} d x-2 \operatorname{Re} \int_{a}^{b} \frac{u_{\mathrm{e}}}{(x-t)^{3}} \frac{\overline{\partial u_{\mathrm{\varepsilon}}}}{\partial t} d x . \tag{3.10}
\end{gather*}
$$

It is not difficult to prove the above procedure rigourously noting for example that if $u \in V(t)$ we have

$$
|u(x)|=\left|\int_{t}^{x} \frac{d u}{d y} d y\right| \leqq \sqrt{|x-t|} \sqrt{\int_{t}^{x}\left|\frac{d u}{d y}\right|^{2} d y}
$$

Applying Schwarz inequality to the right of (3.10) and recalling (3.9) we obtain

$$
\begin{gathered}
\int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathrm{e}}}{\partial t}\right|^{2} d x \leqq\left\|H_{\mathrm{e}}(t)^{1 / 2} \frac{d}{d t} u_{\mathrm{e}}(t)\right\|^{2} \\
\leqq\left(\varepsilon \sqrt{\int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathrm{e}}}{\partial x}\right|^{2} d x}+2 \sqrt{\int_{a}^{b}\left|\frac{u_{\mathrm{e}}}{(x-t)^{2}}\right|^{2} d x}\right) \sqrt{\int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathrm{e}}}{\partial t}\right|^{2} d x}
\end{gathered}
$$

(3.6), (3.7) and (3.11) implies

$$
\begin{equation*}
\sqrt{\int_{a}^{b}\left|\frac{1}{x-t} \frac{\partial u_{\mathfrak{z}}}{\partial t}\right|^{2} d x} \leqq K(2+\sqrt{\varepsilon})\|g\| \tag{3.12}
\end{equation*}
$$

Combining (3.11), (3.12), (3.16) and (3.17) we get

$$
\begin{equation*}
\left\|H_{\mathrm{e}}(t)^{1 / 2} \frac{d}{d t} A_{\mathrm{e}}(t)^{-1}\right\| \leqq K(2+\sqrt{\varepsilon}) \tag{3.13}
\end{equation*}
$$

(3.7) implies $\left\|A_{\mathrm{e}}(t) H_{\mathrm{e}}(t)^{-1}\right\| \leqq 1+K \sqrt{\varepsilon}$, and hence by the generalization of Heinz inequality by Kato [3] we conclude

$$
\left\|A_{\mathrm{e}}(t)^{1 / 2} H_{\mathrm{e}}(t)^{-1 / 2}\right\| \leqq e^{\pi^{2 / 8}}(1+K \sqrt{\varepsilon})^{1 / 2}
$$

It follows from (3.13) and (3.14) that

$$
\begin{equation*}
\left\|A_{\mathrm{e}}(t)^{1 / 2} \frac{d}{d t} A_{\mathrm{e}}(t)^{-1}\right\| \leqq e^{\pi^{2 / 8}}(1+K \sqrt{\varepsilon})^{1 / 2} K(2+\sqrt{\varepsilon}) \tag{3.14}
\end{equation*}
$$

which states that (III) of Theorem 1.1 holds for $A_{8}(t), 0<\varepsilon \leqq 4 / 5$, with constants independent of $\varepsilon$ and $t$. It is not difficult to prove that the
remaining part of the assumptions of Theorem 2.1 is satisfied by $A_{\mathrm{z}}(t)$, $0<\varepsilon \leqq 4 / 5$.

REMARK. If $B(t)$ is the multiplication operator

$$
(B(t) u)(x)=\frac{u(x)}{(x-t)^{2}}
$$

then we have $D(A(t)) \subset D(B(t))$. Application of T. Kato's generalization of Heinz's theorem ([3]) shows that $D\left(A(t)^{\rho}\right) \subset D\left(B(t)^{\rho}\right)$ for any $\rho$ with $0<\rho<1$, therefore any function belonging to $D\left(A(t)^{\rho}\right)$ must vanish at $t$ in some sense. Thus we conclude that $D\left(A(t)^{\rho}\right)$ is not independent of $t$ whenever $\rho>0$. The same thing remains true in the first case as a consequence of Heinz's theorem itself.

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