

Open Archive TOULOUSE Archive Ouverte (OATAO)

OATAO is an open access repository that collects the work of Toulouse researchers and makes it freely available over the web where possible.

This is an author-deposited version published in : <u>http://oatao.univ-toulouse.fr/</u> Eprints ID : 13899

> To link to this article : DOI:10.1038/NMAT3601 URL : http://dx.doi.org/10.1038/NMAT3601

To cite this version :

Augustyn, Veronica and Come, Jérémy and Lowe, Michael A. and Kim, Jong Woung and Taberna, Pierre-Louis and Tolbert, Sarah H. and Abruña, Héctor D. and Simon, Patrice and Dunn, Bruce *Highrate electrochemical energy storage through Li+ intercalation pseudocapacitance*. (2013) Nature Materials, vol. 12 (n° 6). pp. 518-522. ISSN 1476-1122

Any correspondance concerning this service should be sent to the repository administrator: staff-oatao@listes-diff.inp-toulouse.fr

High-rate electrochemical energy storage through Li⁺ intercalation pseudocapacitance

Veronica Augustyn¹, Jérémy Come^{2,3}, Michael A. Lowe⁴, Jong Woung Kim¹, Pierre-Louis Taberna^{2,3}, Sarah H. Tolbert⁵, Héctor D. Abruña⁴, Patrice Simon^{2,3} and Bruce Dunn¹*

Pseudocapacitance is commonly associated with surface or near-surface reversible redox reactions, as observed with RuO₂ · xH₂O in an acidic electrolyte. However, we recently demonstrated that a pseudocapacitive mechanism occurs when lithium ions are inserted into mesoporous and nanocrystal films of orthorhombic Nb₂O₅ (T-Nb₂O₅; refs 1,2). Here, we quantify the kinetics of charge storage in T-Nb₂O₅: currents that vary inversely with time, charge-storage capacity that is mostly independent of rate, and redox peaks that exhibit small voltage offsets even at high rates. We also define the structural characteristics necessary for this process, termed intercalation pseudocapacitance, which are a crystalline network that offers two-dimensional transport pathways and little structural change on intercalation. The principal benefit realized from intercalation pseudocapacitance is that high levels of charge storage are achieved within short periods of time because there are no limitations from solid-state diffusion. Thick electrodes (up to 40 μ m thick) prepared with T-Nb₂O₅ offer the promise of exploiting intercalation pseudocapacitance to obtain high-rate charge-storage devices.

Pseudocapacitance occurs whenever the charge (Q) depends on the change in potential (dE), yielding a capacitance (dQ/dE)(ref. 3). The capacity can be due to monolayer adsorption of ions at an electrode surface, as in the underpotential deposition of metals⁴; surface redox reactions as in RuO₂; or ion intercalation that does not result in a phase change. Although these redox processes are Faradaic in nature, their phenomenological behaviour, and response to experimental variables such as sweep rate, are those typical of capacitors. All of these scenarios produce a relationship between the fractional extent of charge storage, *X*, and the potential of the form^{3,5}:

$$E \sim \left(\frac{RT}{nF}\right) \ln\left[\frac{X}{(1-X)}\right] \tag{1}$$

where *R* is the ideal gas constant (J mol⁻¹ K⁻¹), *T* is the temperature (K), *F* is Faraday's constant (96,485 As mol⁻¹) and *n* is the number of electrons involved in the reaction. In all of these cases, a constant-current experiment yields a potential *E* that changes with the extent of charge *Q* according to:

$$Q = C\Delta E$$

where Q is the charge passed (Coulombs), ΔE is the potential change (V) and C is the pseudocapacitance (F). This behaviour

is typical of capacitive charge/discharge, thus leading to the term pseudocapacitance³.

Of the three pseudocapacitive mechanisms mentioned above, underpotential deposition and surface redox reaction pseudocapacitance exhibit kinetics indicative of surface-controlled electrochemical processes⁶:

i = Cv

where *i* is the current (A) and ν is the sweep rate (mV s⁻¹) of a cyclic voltammetry experiment. However, in intercalation pseudocapacitance, as described herein, charge storage does not occur on the surface but in the bulk material. The kinetics are not diffusion-limited and instead are limited by surface processes so that the overall behaviour seems capacitive. Intercalation pseudocapacitance is rarely observed because in most intercalation materials charge storage (even in thin films, as in anatase TiO₂; refs 7,8) is limited by solid-state diffusion and therefore the peak currents scale with $\nu^{1/2}$.

Here we investigate the phenomenon of intercalation pseudocapacitance and the high-rate behaviour of T-Nb₂O₅ using two different electrode techniques that provide a wide variation in sweep rates. For timescales between ~3 h and 60 s (sweep rates of 0.1–20 mV s⁻¹ within a voltage window of 1.2 V), we used a thin-film electrode. For shorter timescales where ohmic polarization is significant (60–500 mV s⁻¹), we used a cavity microelectrode where the active material was mixed with a conductive carbon black to alleviate the loss of electrical transport (ohmic losses)⁹. To confirm the small ohmic drop of this electrode, we performed cyclic voltammetry from 100–500 mV s⁻¹ in a bulky-ion electrolyte (Supplementary Fig. S1). In addition, we prepared thick films (~40 µm) of *T*-Nb₂O₅ to determine whether the high-rate capability was limited to thin films.

Charge storage from the intercalation of lithium ions into Nb_2O_5 can be expressed as:

$$Nb_2O_5 + xLi^+ + xe^- \leftrightarrow Li_xNb_2O_5$$

where the maximum capacity is x = 2 (ref. 10). Figure la shows cyclic voltammograms from 100 to 500 mV s⁻¹ in a cavity microelectrode where it is evident that both anodic and cathodic peaks are broad, about 600 mV. There is also a noticeable peak shift (and increase in peak separation ΔE_p) as the sweep rate increases, but the capacity remains reversible. Figure 1b presents a plot of log(*i*) versus log(ν) from 0.1 to 500 mV s⁻¹ for both cathodic and anodic peaks. Assuming that the current obeys a power-law

¹Department of Materials Science and Engineering, University of California, Los Angeles, California 90095, USA, ²Department of Materials Science, Université Paul Sabatier, CIRIMAT UMR CNRS 5085, Toulouse 31062, France, ³Réseau sur le Stockage Electrochimique de l'Energie (RS2E), FR CNRS 3459, France, ⁴Department of Chemistry and Chemical Biology, Cornell University, Ithaca, New York 14853, USA, ⁵Department of Chemistry and Biochemistry, University of California, Los Angeles, California 90095, USA. *e-mail: bdunn@ucla.edu.



Figure 1 | **Kinetic analysis of the electrochemical behaviour of** T**-Nb**₂**O**₅**. a**, Cyclic voltammograms from 100 to 500 mV s⁻¹ demonstrate the high-rate capability of the material. **b**, *b*-value determination of the peak anodic and cathodic currents shows that this value is approximately 1 up to 50 mV s⁻¹. This indicates that even at the peak currents, charge storage is capacitive. **c**, Capacity versus $v^{-1/2}$ allows for the separation of diffusion-controlled capacity from capacitive-controlled capacity; two distinct kinetic regions emerge when the sweep rate is varied from 1 to 500 mV s⁻¹. The dashed diagonal line corresponds to the extrapolation of the infinite sweep rate capacitance using the capacity between 2 and 20 mV s⁻¹. **d**, The variation of the cathodic peak voltage with the sweep rate exhibits a region of small peak separation followed by increased separation at 20 mV s⁻¹, and represents another method of identifying systems with facile intercalation kinetics.

relationship with the sweep rate leads to⁸:

$$i = av^b$$

where *a* and *b* are adjustable values. Whereas a *b*-value of 0.5 would indicate that the current is controlled by semi-infinite linear diffusion, a value of 1 indicates that the current is surface-controlled. For sweep rates ranging from 0.1 to 20 mV s^{-1} , corresponding to charging times >60 s, the *b*-value for both the cathodic and anodic peaks is 1, indicating that the kinetics are surface-controlled, and thus fast. Figure 1b also exhibits a change in the slope of the anodic and cathodic peak currents at ~50 mV s⁻¹. This change in slope corresponds to a decrease in *b*-value to 0.8 and 0.7 for the cathodic and anodic currents, respectively, at sweep rates >50 mV s⁻¹ (charging times <20 s). This limitation to the rate capability can arise from numerous sources including an increase of the ohmic contribution (active material resistance, solid–electrolyte interphase resistance) and/or diffusion constraints/limitations¹¹. In the limit of slow diffusion, *b* would approach a value of 0.5 as described above.

The relationship between capacity and sweep rate can also establish the rate-limiting step of a charge-storage mechanism¹². In a plot of Q versus $v^{-1/2}$, regions that are linear represent capacity limited by semi-infinite linear diffusion whereas capacitive contributions are independent of the sweep rate. At sweep rates below 20 mV s^{-1} , the extrapolated *y*-intercept yields the infinite sweep rate capacitance¹³. Figure 1c shows the plot of normalized capacity versus $v^{-1/2}$ for \tilde{T} -Nb₂O₅ from 1 to 500 mV s⁻¹ (the gravimetric capacity for the thin-film electrode is shown in Supplementary Fig. S2). Analogous to the behaviour of the peak current in Fig. 1b, there are two distinct regions in Fig. 1c. In region 1, at sweep rates $<20 \text{ mV s}^{-1}$, the capacity is mostly independent of sweep rate. The magnitude of the capacity is ${\sim}130\,\text{mAh}\,\text{g}^{-1}$ or ${\sim}65\%$ of the theoretical value based on a two-electron redox reaction with Nb₂O₅. In this range, solidstate lithium-ion diffusion is not the rate-limiting step for charge storage. In region 2, from 50 to 500 mV s^{-1} , the capacity decreases linearly with $v^{-1/2}$. This indicates that charge storage is mainly diffusion-controlled at high sweep rates. That is, for charging times



Figure 2 | **Electrochemical cycling of a 40-\mum-thick T-Nb₂O₅ electrode. a**, Galvanostatic cycling of a thick Nb₂O₅ electrode at a 10C rate. **b**, Comparison of the rate capability of *T*-Nb₂O₅ with a high-rate lithium-ion anode, Li₄Ti₅O₁₂, at various C-rates (Li₄Ti₅O₁₂ data reproduced from ref. 16).

of <20 s, diffusion is rate-limiting, similar to most traditional battery electrodes. However, for charging times of 1 min (60C) or longer, there is no indication of diffusion limitations and this intercalation-based system behaves in a fully capacitive manner.

Another feature of T-Nb₂O₅ at sweep rates $< 20 \text{ mV s}^{-1}$ is that the peak voltage shifts with sweep rate are small (Fig. 1d). The cathodic peak shift is <0.1 V at sweep rates below 10 mV s⁻¹. As a result, the anodic and cathodic peaks overlap at 0.1 mV s^{-1} (Supplementary Fig. S3) and it is in this behaviour that the similarity to surface redox reactions is most apparent¹⁴. In many lithium-ion intercalation materials, the peak separation is significant even in thin films and at slow sweep rates (for example, $\Delta E_{\rm p} = 0.13 \, {\rm V}$ for LiCoO₂ at 0.1 mV s⁻¹; ref. 15). This type of behaviour is often associated with crystallographic phase changes during the Faradaic process, and contrasts with intercalation materials that form a solid solution, such as T-Nb₂O₅. Besides identifying facile intercalation, the peak voltage separation is related to the high-power capability of a material. As the charging time decreases, that is, at higher current densities, the peak separation in a battery material increases owing to polarization (reflecting the higher overpotentials necessary to deliver the higher currents), so that at higher rates the energy required to fully charge the material is significantly larger than the energy available on discharge.

The high-rate behaviour of T-Nb₂O₅ is not limited to thin films or to experiments with small amounts of active material. The constant-current charge/discharge of a 40-µm-thick T-Nb₂O₅ (1 mg cm^{-2}) electrode at a 10C rate is shown in Fig. 2a. At this rate, the capacity is 130 mAh g^{-1} and *E* varies linearly with *Q* as expected for a pseudocapacitive process from equation (1). This represents capacities typical of battery materials but at rates closer to those of supercapacitors. The rate capability of T-Nb₂O₅ from 1 to 1,000C is shown in Fig. 2b and compared with that of Li₄Ti₅O₁₂ (of comparable electrode dimensions)¹⁶. Li₄Ti₅O₁₂ is chosen as an example of a high-rate lithium-ion anode material. The rate capability for T-Nb₂O₅ is significantly better than Li₄Ti₅O₁₂ above 30C and even at a 1,000C rate the capacity of the thick T-Nb₂O₅ electrode is \sim 40 mAh g⁻¹. The thick electrode results verify that the intercalation pseudocapacitance mechanism is not due to thin-film or surface effects, such as vacancies or contributions of the first few atoms from the surface. Indeed, even for the thick electrode, the b-values for the anodic and cathodic peak currents are 1 from 1 to 10 mV s^{-1} (Supplementary Fig. S4).

The high-rate capability of T-Nb₂O₅ implies that the crystal structure permits exceptionally rapid ionic transport. As shown in Fig. 3a, the unit cell has sheets of edge- or corner-sharing distorted polyhedra lying parallel to the (001) plane, with each Nb⁵⁺ surrounded by either 6 or 7 O²⁻. The polyhedra are exclusively corner-sharing along the [001] direction with 5% of the Nb⁵⁺ ions randomly located in 9-coordinate sites between the (001) polyhedral planes¹⁷. The mostly empty octahedral sites between (001) planes provide natural tunnels for lithium-ion transport throughout the *a*-*b* plane. Calculations indicate that the (001) plane allows degenerate pathways with low energy barriers for ion transport¹⁸.

A previous in situ X-ray diffraction study showed that the insertion of lithium into T-Nb₂O₅ results in a solid solution with no apparent phase changes¹⁹, and negligible changes to lattice constants² and unit-cell volume²⁰ up to ~1.25 Li⁺/Nb₂O₅. In situ X-ray absorption spectroscopy (XAS) demonstrated that lithiation reduces Nb⁵⁺ to Nb⁴⁺ (ref. 19). The in situ XAS studies carried out here confirm that lithiation results in a continuous change in oxidation state (Fig. 3b) and the Fourier-transform of the extended X-ray absorption fine structure (EXAFS) indicates that the insertion reaction proceeds through two stages (Fig. 3c). From 2.5 to 1.75 V, the EXAFS signal from the various Nb-O bond lengths in Nb₂O₅ (1.40-1.85 Å) merges to a single peak at an intermediate bond length (1.75 Å), indicating that lithiation increases the Nb-centred symmetry. The lithiation is probably faster at low Li⁺ levels²⁰ owing to greater availability of sites and less interaction between cations. At lower potentials, the new EXAFS peak shifts to longer bond distances (1.85 Å) as a consequence of increased Li–O interactions at higher Li⁺ content. These structural studies emphasize the value of an open, layered structure to enable rapid ion transport within the active material.

The results presented in this study establish that T-Nb₂O₅ exhibits electrochemical features of a pseudocapacitive material despite charge storage occurring in the bulk. Such behaviour is consistent with intercalation pseudocapacitance. The electrochemical features indicative of this mechanism are currents that are linearly proportional to the sweep rate, capacity that does not vary significantly with charging time, and peak potentials that do not shift significantly with sweep rate. A key design rule for intercalation pseudocapacitance at the atomic scale is a structure that does not undergo phase transformations on intercalation. In addition, facile two-dimensional (2D) lithium-ion diffusion pathways are



Figure 3 | **Structural features of lithium intercalation in T-Nb₂O₅. a**, The structure of T-Nb₂O₅ stacked along the *c* axis demonstrates the layered arrangement of oxygen (red) and niobium (inside polyhedra) atoms along the *a*-*b* plane. **b**, Derivative of Nb K-edge X-ray absorption near-edge spectra at selected cell voltages, showing a systematic shift to lower energies as Nb⁵⁺ is reduced to Nb⁴⁺. **c**, k^2 -weighted Fourier-transformed Nb K-edge EXAFS at selected cell voltages.

Table 1 | Comparison between charge storage in two different pseudocapacitive materials.

	RuO₂ ∙ <i>x</i> H₂O	T-Nb ₂ O ₅
Electrolyte	H ⁺ , aqueous ²¹	Li ⁺ , non-aqueous
Structure	Amorphous/ nanocrystalline ^{22,23}	Crystalline
Structural water	Necessary for proton conduction ²²	Not necessary
Charge storage mechanism	Surface/near- surface ^{23,24}	Bulk
Pseudocapacitance type	Surface redox	Intercalation

 ${\rm RuO_2} \cdot x {\rm H_2O}$ and $T \cdot {\rm Nb_2O_5}$ both exhibit capacitive behaviour. The mechanism of charge storage, however, is different and results in different structural requirements for high capacitance.

important. Charge storage that behaves as a quasi-2D process exhibits similar behaviour to 2D surface adsorption reactions¹⁴. These features contrast with those of pseudocapacitive $RuO_2 \cdot xH_2O$ where charge storage occurs mainly on the surface or near-surface²¹⁻²⁴ as summarized in Table 1.

The results here are exciting because they demonstrate that for charging times as fast as 1 min (60C rate), there are no diffusion limitations in T-Nb₂O₅. As the high-rate capability is due to fast ion diffusion in the bulk, this mechanism may be very good for thick electrodes because surface exposure to the electrolyte is not critical. To achieve devices with high energy density, further engineering at the nanoscale and beyond will be necessary to preserve the atomic-scale behaviour observed in thin films and microelectrodes. In particular, maintaining proper electronic conduction pathways will be critical.

Methods

Synthesis. The synthesis of T-Nb₂O₅ nanocrystals was reported previously². Briefly, 2.56 mmol of NbCl₅ (Sigma-Aldrich) was dissolved in 2 ml of ethanol (Fisher Scientific). In a separate vial, 0.23 ml of deionized water was mixed with 2 ml of ethanol. Both vials were then chilled for 2 h. The two solutions were then mixed together while 1 ml of propylene oxide (Sigma-Aldrich) was slowly added, forming a transparent gel. This gel was aged for 1 day and then soaked in acetone (Fisher Scientific) for 5 days. After supercritical drying with CO₂, the gel was transformed into an amorphous Nb₂O₅ aerogel. Crystallization to the *T*-phase occurred by heat treatment at 600 °C for 2 h in air.

Characterization. Thin-film electrodes were fabricated by drop-casting a well-sonicated solution of T-Nb₂O₃ in ethanol onto an oxygen plasma etched stainless-steel foil (Alfa Aesar). For thin-film measurements, lithium foils (Sigma-Aldrich) served as the reference and counter electrodes, and the electrolyte was 1 M LiClO₄ (Sigma-Aldrich) in propylene carbonate (Sigma-Aldrich). The microelectrode preparation was described elsewhere⁹. Briefly, T-Nb₂O₅ active material and carbon black (Timcal Super C65) powders were mixed in a 1:1 weight ratio to ensure good electrical contact. The mixture was packed into the microcavity (30 µm depth, 30 µm width), and subsequently immersed in the electrolyte. A 1 cm² platinum foil (Alfa Aesar) was used as the counter electrode

and lithium foil served as the reference. Thick Nb₂O₅ electrodes were prepared by mixing the active material, carbon black (Timcal Super C65) and PVdF (Kynar) binder in a 80:10:10 weight ratio. The slurry was drop-cast onto 12-mm-diameter aluminium-disc current collectors and dried at 80 °C for 12 h. The electrode loading was between 1 and 1.5 mg cm⁻² with a thickness of 40 ± 5 µm. Measurements of the thick Nb₂O₅ electrodes were performed in three-electrode Swagelok cells with the Nb₂O₅ electrode as the working electrode, activated carbon as the counter electrode, a glass fibre separator (Whatman, GF/A) saturated with 1 M LiClO₄ in propylene carbonate, and a lithium foil reference. All measurements were carried out in argon-filled glove boxes between 1.2 and 3 V (versus Li/Li⁺) with oxygen and moisture levels of <1 ppm and using PAR EG&G 273 and Bio-Logic VMP3 potentiostats.

In situ X-ray absorption was performed using thin films of T-Nb₂O₅ cast from a well-sonicated ethanol solution onto carbon-coated aluminium. 2032-type coin cells were modified to include an X-ray window by drilling ~3-mm-diameter holes through the cathode casing and epoxying a 125-µm-thick Kapton window on the outside of the casing. Both sides of the window were then coated with 100-nm-thick aluminium layers using a vacuum evaporator. Cells were assembled with a lithium anode, a 1 M LiClO₄ in propylene carbonate electrolyte, and a Whatman glassy fibre separator.

References

- Brezesinski, K. *et al.* Pseudocapacitive contributions to charge storage in highly ordered mesoporous Group V transition metal oxides with iso-oriented layered nanocrystalline domains. *J. Am. Chem. Soc.* 132, 6982–6990 (2010).
- Kim, J. W., Augustyn, V. & Dunn, B. The effect of crystallinity on the rapid pseudocapacitive response of Nb₂O₅. Adv. Energy Mater. 2, 141–148 (2012).
- Conway, B. E. Electrochemical Supercapacitors: Scientific Fundamentals and Technological Applications (Kluwer-Academic, 1999).
- Herrero, E., Buller, L. J. & Abruña, H. D. Underpotential deposition at single crystal surfaces of Au, Pt, Ag and other materials. *Chem. Rev.* 101, 1897–1930 (2001).
- Huggins, R. A. Supercapacitors and electrochemical pulse sources. Solid State Ion. 134, 179–195 (2000).
- Angerstein-Kozlowska, H., Klinger, J. & Conway, B. E. Computer simulations of the kinetic behavior of surface reactions driven by a linear potential sweep. Part I: Model 1-electron reaction with a single adsorbed species. *J. Electroanal. Chem.* **75**, 45–60 (1977).
- Brezesinski, T., Wang, J., Polleux, J., Dunn, B. & Tolbert, S. H. Templated nanocrystal-based porous TiO₂ films for next-generation electrochemical capacitors. *J. Am. Chem. Soc.* 131, 1802–1809 (2009).
- 8. Lindström, H. *et al.* Li⁺ insertion in TiO₂ (anatase). 2. Voltammetry on nanoporous films. *J. Phys. Chem. B* **101**, 7717–7722 (1997).
- Come, J., Taberna, P-L., Hamelet, S., Masquelier, C. & Simon, P. Electrochemical kinetic study of LiFePO₄ using cavity microelectrode. *J. Electrochem. Soc.* 158, A1090–A1093 (2011).
- Ohzuku, T., Sawai, K. & Hirai, T. Electrochemistry of L-niobium pentoxide in a lithium/non-aqueous cell. J. Power Sources 19, 287–299 (1987).
- Park, M., Zhang, X., Chung, M., Less, G. B. & Sastry, A. M. A review of conduction phenomena in Li-ion batteries. *J. Power Sources* 195, 7904–7929 (2010).
- Ardizzone, S., Fregonara, G. & Trasatti, S. 'Inner' and 'outer' active surface of RuO₂ electrodes. *Electrochim. Acta* 35, 263–267 (1990).
- Baronetto, D., Krstajić, N. & Trasatti, S. Reply to Note on a method to interrelate inner and outer electrode areas by H. Vogt. *Electrochim. Acta* 39, 2359–2362 (1994).

- 14. Conway, B. E. Transition from 'supercapacitor' to 'battery' behavior in electrochemical energy storage. J. Electrochem. Soc. 138, 1539–1548 (1991).
- 15. Xia, H., Lu, L. & Ceder, G. Substrate effect on the microstructure and electrochemical properties of LiCoO₂ thin films grown by PLD. *J. Alloys Compd.* **417**, 304–310 (2006).
- Zhang, N. Q. *et al.* Facile preparation of nanocrystalline Li₄Ti₅O₁₂ and its high electrochemical performance as anode material for lithium-ion batteries. *Electrochem. Commun.* 13, 654–656 (2011).
- 17. Kato, K. & Tamura, S. The crystal structure of *T*-Nb₂O₅. *Acta Cryst.* B31, 673–677 (1975).
- Liu, C-P., Zhou, F. & Ozolins, V. First principles study for lithium intercalation and diffusion behavior in orthorhombic Nb₂O₅ electrochemical supercapacitor *American Physical Society March Meeting 2012* abstr. B26.00003. Accessed March 20, 2013. http://meetings.aps.org/link/BAPS.2012.MAR.B26.3.
- Kodama, R., Terada, Y., Nakai, I., Komaba, S. & Kumagai, N. Electrochemical and *in situ* XAFS-XRD investigation of Nb₂O₅ for rechargeable lithium batteries. *J. Electrochem. Soc.* 153, A583–A588 (2006).
- Kumagai, N., Koishikawa, Y., Komada, S. & Koshiba, N. Thermodynamics and kinetics of lithium intercalation into Nb₂O₅ electrodes for a 2 V rechargeable lithium battery. J. Electrochem. Soc. 146, 3203–3210 (1999).
- Long, J. W., Swider, K. E., Merzbacher, C. I. & Rolison, D. R. Voltammetric characterization of ruthenium oxide-based aerogels and other RuO₂ solids: The nature of capacitance in nanostructured materials. *Langmuir* 15, 780–785 (1999).
- Zheng, J. P., Cygan, P. J. & Jow, T. R. Hydrous ruthenium oxide as an electrode material for electrochemical capacitors. *J. Electrochem. Soc.* 142, 2699–2703 (1995).
- Dmowski, W., Egami, T., Swider-Lyons, K. E., Love, C. T. & Rolison, D. R. Local atomic structure and conduction mechanism of nanocrystalline hydrous RuO₂ from X-ray scattering. *J. Phys. Chem. B* 106, 12677–12683 (2002).
- Liu, Y., Zhou, F. & Ozolins, V. *Ab initio* study of the charge-storage mechanisms in RuO₂-based electrochemical ultracapacitors. *J. Phys. Chem. C* 116, 1450–1457 (2012).

Acknowledgements

This work was supported by the Molecularly Engineered Energy Materials and the Energy Materials Center at Cornell, Energy Frontiers Research Centers funded by the US DOE Office of Basic Energy Sciences (DE-SC001342 and DE-SC0001086, respectively). XAS was performed at the Cornell High Energy Synchrotron Source, supported by the NSF and NIH/NIGMS (DMR-0936384). M.A.L. acknowledges support from the US DOD National Defense Science and Engineering Fellowship. J.C. was supported by Delegation Generale pour l'Armement (DGA). P.S. and P-L.T. acknowledge the support from the European Research Council (ERC, Advanced Grant, ERC-2011-AdG, Project 291543—IONACES) and the Chair of Excellence 'Embedded multi-functional nanomaterials' from the EADS Foundation.

Author contributions

V.A., J.C., M.A.L. and J.W.K.: experimental work, data analysis. P-L.T., S.H.T., H.D.A., P.S., B.D.: project planning, data analysis.