

1                    **Spectral Analysis of the September 2017 Solar**  
2                    **Energetic Particle Events**

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8                    **Key Points:**

- 9                    • Extreme solar activity occurred in early September 2017; three solar energetic par-  
10                    ticle events, including a GLE, were registered at Earth  
11                    • A comprehensive investigation is made of the proton energy spectra based on ACE,  
12                    GOES and STEREO observations  
13                    • Spectral features are interpreted in terms of acceleration, transport and connec-  
14                    tivity. Results are compared with those of the previous GLE

## Abstract

An interval of exceptional solar activity was registered in early September 2017, late in the decay phase of solar cycle 24, involving the complex Active Region 12673 as it rotated across the western hemisphere with respect to Earth. A large number of eruptions occurred between 4–10 September, including four associated with X-class flares. The X9.3 flare on 6 September and the X8.2 flare on 10 September are currently the two largest during cycle 24. Both were accompanied by fast coronal mass ejections and gave rise to solar energetic particle (SEP) events measured by near-Earth spacecraft. In particular, the partially-occulted solar event on 10 September triggered a ground level enhancement (GLE), the second GLE of cycle 24. A ~~third~~ further, much less energetic SEP event was recorded on 4 September. In this work we analyze observations by the Advanced Composition Explorer (ACE) and the Geostationary Operational Environmental Satellites (GOES), estimating the SEP event-integrated spectra above 300 keV and carrying out a detailed study of the spectral shape temporal evolution. Derived spectra are characterized by a low-energy break at few/tens of MeV; the 10 September event spectrum, extending up to  $\sim 1$  GeV, exhibits an additional rollover at several hundred MeV. We discuss the spectral interpretation in the scenario of shock acceleration and in terms of other important external influences related to interplanetary transport and magnetic connectivity, taking advantage of multi-point observations from the Solar Terrestrial Relations Observatory (STEREO). Spectral results are also compared with those obtained for the 17 May 2012 GLE event.

## 1 Introduction

It is generally accepted that solar energetic particles (SEPs) are accelerated by a mixture of processes associated with flares and coronal mass ejections (CMEs) (see, e.g. Desai & Giacalone (2016)). Such mechanisms are predicted to leave distinct signatures in the energy spectrum, whose measurement thus provides important constraints on SEP origin. However, spectral features observed at different energies may arise from particle acceleration in different locations (e.g., the flare region, corona or interplanetary space), so the spectral shapes may exhibit the combined signatures of several dynamic processes that may be complex to disentangle. Furthermore, the morphology and the evolution of SEP events are strongly influenced by the magnetic connection to sources and by interplanetary transport effects and transient/recurrent solar wind (SW) disturbances which significantly complicate the interpretation of spectral measurements.

The early September 2017 solar events were well-observed by several space- and ground-based instruments, receiving noteworthy attention by a number of papers in the literature (see, e.g., Chertok (2018); Gary et al. (2018); Gopalswamy et al. (2018); Guo et al. (2018); Long et al. (2018); Luhmann et al. (2018); Omodei et al. (2018); Seaton & Darnel (2018); Sharykin & Kosovichev (2018); Shen al. (2018); Sun & Norton (2017); Warren et al. (2018)). In this work we focus on the SEP events that accompany these eruptions, taking advantage of multi-spacecraft data ~~by~~ from the Advanced Composition Explorer (ACE) and the Geostationary Operational Environmental Satellites (GOES) to provide an assessment of the SEP spectral shapes over a complete range of energies spanning from few hundreds of keV to a few GeV. We also illustrate the effects of SW structures on the SEP spectra. In addition, observations from the Solar Terrestrial Relations Observatory-Ahead (STEREO-A) are used to provide a more complete view of these SEP events near 1 AU. The paper is structured as follows: the September 2017 events are introduced in Section 2; in Section 3 we analyze the various SEP measurements and examine the relevant interplanetary data; Section 4 describes the reconstruction and analysis of SEP spectra; results are presented and discussed in Section 5; finally, Section 6 reports our summary and conclusions.

Date	Flare					CME			
	Class	Onset	Peak	End	Location	Speed	1 <sup>st</sup> -app. time	Width	Direction
<b>04 Sept.</b>	<b>M5.5</b>	<b>20:28</b>	<b>20:33</b>	<b>20:37</b>	<b>S11W16</b>	<b>1418/1114</b>	<b>20:12/20:36</b>	<b>360/92</b>	<b>S10W10</b>
06 Sept.	X2.2	08:57	09:10	09:17	S07W33	391/260	09:48/10:00	80/48	S08W83
<b>06 Sept.</b>	<b>X9.3</b>	<b>11:53</b>	<b>12:02</b>	<b>12:10</b>	<b>S08W33</b>	<b>1571/1238</b>	<b>12:24/12:24</b>	<b>360/88</b>	<b>S15W23</b>
07 Sept.	M7.3	10:11	10:15	10:18	S08W47	470/597	10:24/10:48	32/26	S13W51
07 Sept.	X1.3	14:20	14:36	14:55	S11W49	433/477	15:12/15:12	58/32	S16W53
08 Sept.	M8.1	07:40	07:49	07:58	S10W57	500/450	07:36/07:24	31/40	S03W54
<b>10 Sept.</b>	<b>X8.2</b>	<b>15:35</b>	<b>16:06</b>	<b>16:31</b>	<b>S08W88</b>	<b>3163/2650</b>	<b>16:00/16:09</b>	<b>360/108</b>	<b>S12W85</b>

**Table 1.** List of eruptions associated with major flares ( $>M5.0$ ) originated from AR NOAA 12673 during September 2017. Data in bold refer to the three SEP events registered at Earth. For each event, the flare class, onset/peak/~~stop~~end times (UT) and location (deg) are shown, based on the GOES-15 X-ray archive (<ftp://ftp.ngdc.noaa.gov/STP/space-weather/solar-data/solar-features/solar-flares/x-rays/goes/>), along with first appearance time (UT), ~~sky-plane~~ speed ( $\text{km s}^{-1}$ ), angular width (deg) and direction (deg) of the linked CME. The first and the second values reported for CMEs are from the CDAW ([https://cdaw.gsfc.nasa.gov/CME\\_list/](https://cdaw.gsfc.nasa.gov/CME_list/)) and the DONKI (<https://kauai.ccmc.gsfc.nasa.gov/DONKI/>) catalogs, respectively; CME directions are based on the latter. **Sky-plane (space) speeds are reported in case of CDAW (DONKI).**

## 2 The September 2017 solar events

The first half of September 2017 was characterized by extreme solar activity mostly related to the complex Active Region (AR) NOAA 12673, which rapidly developed on 4–5 September when near central meridian (e.g., Sun & Norton (2017)) and rotated over the west limb on 10 September. A large number of bright eruptions were registered between 4 and 10 September, including 27 associated with M-class flares and four with X-class flares. Table 1 lists the  $>M5$  flares during this period. That such large AR can emerge late in the declining phase of solar cycles is also demonstrated by the December 2006 events, involving four X-class flares including the powerful X9.0 flare on 5 December and the X3.4 flare on 13 December associated with the 70<sup>th</sup> ground level enhancement (GLE), linked to AR 10930 during the analogous period of the previous solar cycle (Adriani et al., 2011). ~~In addition, Richardson et al. (2016) noted that the solar minimum between cycles 23 and 24 was actually unusual compared to previous minima in having no substantial SEP events within two years of sunspot minimum.~~

Three of the major flares, indicated by bold type in Table 1 were associated with fast CMEs and gave rise to SEP events. A first, small SEP event ~~was~~ observed late on 4 September, originated from the moderately intense flare (M5.5) and the geo-effective, halo CME that erupted on the same day. The coordinated data analysis workshops (CDAW, [https://cdaw.gsfc.nasa.gov/CME\\_list/](https://cdaw.gsfc.nasa.gov/CME_list/)) catalog of the Large Angle and Spectrometric Coronagraph (LASCO) on board the Solar and Heliospheric Observatory (SOHO) indicates a linear speed of  $1418 \text{ km s}^{-1}$ ; the Database Of Notifications, Knowledge, Information (DONKI, <https://kauai.ccmc.gsfc.nasa.gov/DONKI/>) reports a **space** speed of  $1114 \text{ km s}^{-1}$  and direction of S10W10, based on the observations of the Sun Earth Connection Coronal and Heliospheric Investigation (SECCHI) instrument on board STEREO-A and of SOHO/LASCO. **Discrepancies in the CME speeds/widths between catalogs are attributable to the different methods used to estimate them including whether they are sky-plane (projected) or space (3-D) speeds based on single- or multiple-point corona-**

103 graph observations, and the helioradial distances at which they are calculated (see Richard-  
104 son et al. (2015) and references therein).

105 The subsequent SEP event was linked to the X9.3 flare peaking at 12:02 UT on 6  
106 September, the largest soft X-ray flare in more than ~~12~~ 10 years (since December 2006)  
107 and the most intense in cycle 24. It generated strong white-light emission and multiple  
108 helioseismic waves observed by the Helioseismic and Magnetic Imager (HMI) on board  
109 the Solar Dynamics Observatory (SDO) (Sharykin & Kosovichev, 2018). The explosion  
110 was associated with an Earth-directed, nearly symmetrical halo CME with an estimated  
111 sky-plane velocity of 1571 km s<sup>-1</sup> according to the CDAW catalog; DONKI indicates a  
112 1238 km s<sup>-1</sup> **space** speed and a S15W23 direction. It was also accompanied by an in-  
113 tense and complex radio emission with interplanetary Type II, III and IV bursts, and  
114 by long-duration  $\gamma$ -ray emission.

115 Finally, a third large SEP event originated following another exceptional flare (X8.2)  
116 occurring on 10 September and peaking at 16:06 UT, when the AR NOAA 12673 had  
117 just rotated over the western solar limb, so the X-ray intensity may be underestimated  
118 due to partial occultation by the limb. To date, it is the second largest soft X-ray flare  
119 of cycle 24, and was associated with a very fast (3136 km s<sup>-1</sup> linear speed) asymmet-  
120 ric halo CME in the CDAW catalog; DONKI indicates a ~~linear~~ **space** speed of 2650  
121 km s<sup>-1</sup> and direction of S12W85. The eruption was accompanied by long-duration emis-  
122 sions at different frequencies, ranging from radio waves (Type II, III and IV bursts) to  
123  $\gamma$ -rays (Gary et al., 2018; Omodei et al., 2018). Spectacular post-flare coronal loops were  
124 observed for nearly a full day. Furthermore, the Solar Ultraviolet Imager (SUVI) on GOES-  
125 16 showed evidence of an apparent current sheet associated with magnetic reconnection  
126 at the beginning of the eruption, and of an extreme-ultraviolet wave at some of the largest  
127 heights ever reported (Long et al., 2018; Seaton & Darnel, 2018; Warren et al., 2018).  
128 The resulting SEP event was energetic enough to give rise to a secondary particle shower  
129 in the Earth's atmosphere which was subsequently detected by neutron monitors (NMs)  
130 on ground as a GLE, the second of solar cycle 24 and the 72<sup>nd</sup> since NM measurements  
131 started in the 1940s (<https://gle.oulu.fi/>).

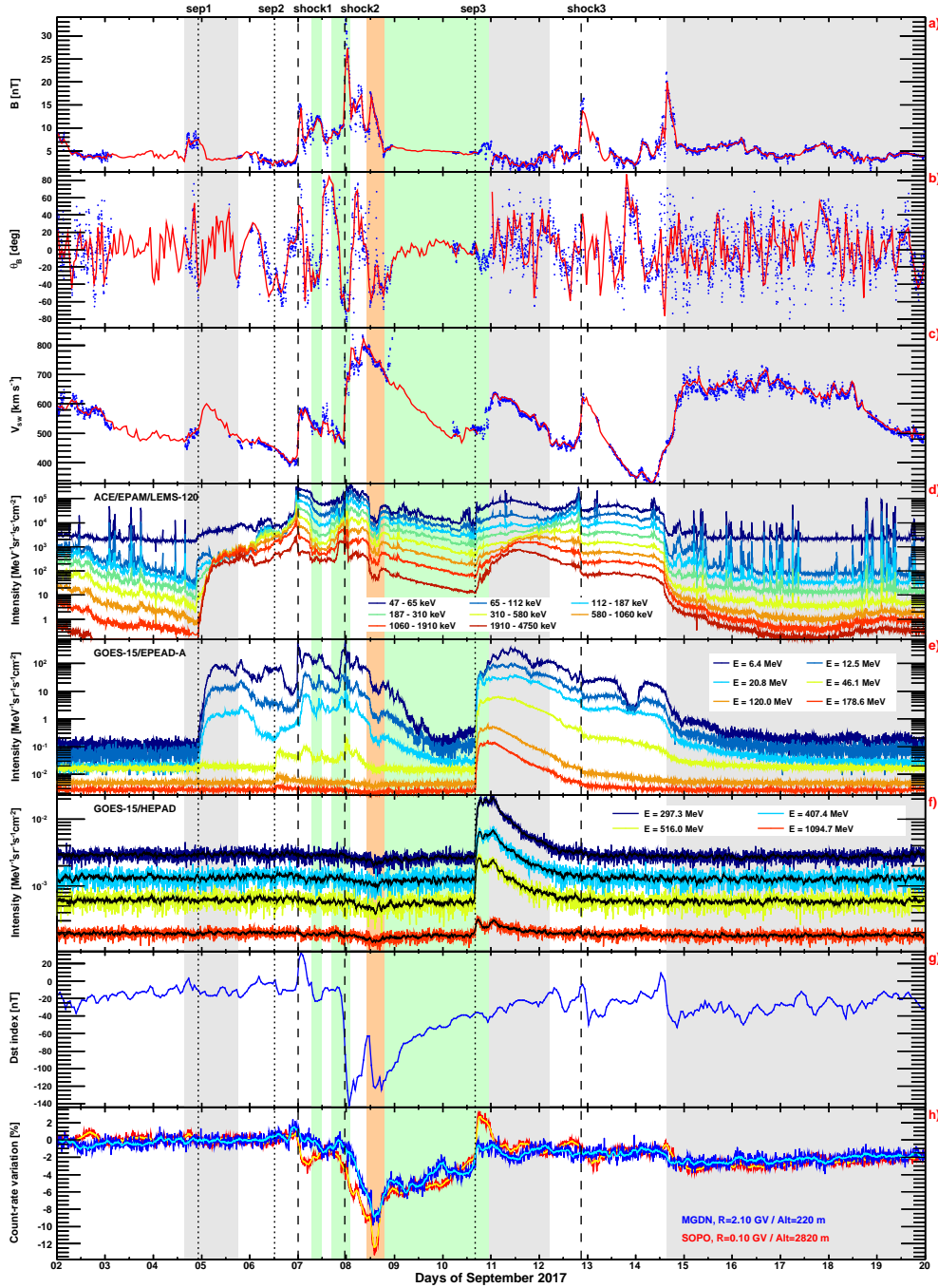
## 132 3 Data

### 142 3.1 SEP data

#### 143 3.1.1 Spacecraft observations

144 Figure 1 shows the temporal variation of the relevant interplanetary, geomagnetic  
145 and particle data between 2–19 September 2017. In particular, panels d), e) and f) dis-  
146 play the 5-min resolution proton intensities measured by near-Earth spacecraft. Specif-  
147 ically, panel d) reports the observations by the Low Energy Magnetic Spectrometer-120  
148 (LEMS-120) of the Electron, Proton, and Alpha Monitor (EPAM) on board ACE, for  
149 7 energy channels ranging from 47 keV to 4.75 MeV ([http://www.srl.caltech.edu/](http://www.srl.caltech.edu/ACE/)  
150 [ACE/](http://www.srl.caltech.edu/ACE/)). Panel e) shows the data from the westward-viewing Energetic Proton, Electron,  
151 and Alpha Detector (EPEAD) on board GOES-15; six energy channels (P2–P7) span-  
152 ning the nominal range 4.2–900 MeV are included. Finally, panel f) displays the intensi-  
153 ties measured by the four energy channels (P8–P11) of the High Energy Proton and  
154 Alpha Detector (HEPAD) on board GOES-15, with a 330–1500(?) MeV nominal energy  
155 interval; the black points correspond to the 1-hr running averages. In case of GOES ([https://](https://www.ngdc.noaa.gov/stp/satellite/goes/)  
156 [www.ngdc.noaa.gov/stp/satellite/goes/](https://www.ngdc.noaa.gov/stp/satellite/goes/)), reported mean energy values are based  
157 on the calibration schemes by Sandberg et al. (2014) and Bruno (2017), respectively be-  
158 low and above 80 MeV.

159 Vertical dotted lines indicate the onset times of the three SEP events introduced  
160 in the previous section, **based on a visual inspection of the intensity profile of the GOES**  
161 **highest-energy channel detecting the SEP arrival**. The first enhancement in the proton



133

134 **Figure 1.** From top to bottom: time profiles of IMF intensity (a), IMF latitude (b),  
 135 SW speed (c), proton intensities measured by ACE/EPAM (d), GOES/EPEAD (e) and  
 136 GOES/HEPAD (f), *Dst* index (g), count rate variations registered by SOPO and MGDN  
 137 NM stations (h). Combined ACE and Wind data (red, 1-hr resolution) are superimposed on  
 138 DSCOVR points (blue, 5-min resolution) in top three panels. The vertical dotted and dashed  
 139 lines mark the onset of the SEP events and the time of the shocks, respectively. The green, or-  
 140 ange and gray areas indicate the periods of the ICMEs, MC and HSSs, respectively. See the text  
 141 for details.

intensities, registered around 22:00 UT on 4 September and limited to energies below  $\sim 150$  MeV, originated from the M5.5 flare and the associated full halo CME reported by SOHO/LASCO at 20:12 UT (see Table 1). A new increase in the intensities of protons with energies up to a few hundreds of MeV was observed around 12:25 UT on 6 September, related to the X9.3 flare and the linked CME registered by SOHO/LASCO at 12:24 UT. The temporal evolution of the SEP event is complex and related to interplanetary structures described in Section 3.2.

A third, large SEP event was produced by the X8.2 flare and the associated very fast CME erupting on 10 September, with an onset around 16:05 UT, during the decaying phase of a Forbush decrease (FD). It was energetic enough to give rise to a GLE detected by high-latitude NM stations (see Section 3.1.2). ~~The sharp increase in proton intensities is indicative of a magnetic connection with the eastern flank of the shock (Cane et al. (1988)).~~ **The sharp increase in proton intensities is consistent with early connection to a shock following a western hemisphere event (Cane et al., 1988), though the W88 location of the event and W85 DONKI CME direction suggest that connection may have been to the eastern flank of the shock assuming nominal Parker spiral interplanetary magnetic field (IMF) lines. However, as pointed out below, the connectivity to the shock is uncertain because of the potential influence of transient SW structures between the Sun and the Earth.** Interestingly, a second peak can be observed in HEPAD profiles at the beginning of 11 September. The origin of this feature will be discussed ~~below~~ **in Section 3.2.**

As a final remark, we note that the EPAM/LEMS-120 low-energy channels ( $\lesssim 500$  keV) are affected by significant electron contamination, as suggested by the gradual enhancement observed apparently before the SEP event onsets. In addition, a number of approximately hour-long bursts can be noted, attributable to ions propagating upstream from the Earth's bow shock when the magnetic connectivity is favorable (see, e.g., Haggerty et al. (2000)).

### 3.1.2 Neutron monitor observations

Panel h) in Figure 1 shows the relative variation in the count-rates registered by the South Pole (SOPO, red points) and the Magadan (MGDN, blue points) NM stations, characterized by different values of geomagnetic cutoff rigidity  $R$  and altitude (see the legend; <http://www.nmdb.eu/>). For SOPO ~~the effective detection threshold is somewhat higher since the minimum particle rigidity is essentially controlled by the atmospheric absorption~~  **$R$  is negligible and the effective detection threshold is determined by the atmospheric cutoff ( $\sim 300$  MeV).**

The error bars refer to the statistical uncertainties. The yellow/cyan points denote the corresponding 1-hr running averages. The SEP event on 10 September gave rise to a GLE, the second of solar cycle 24, commencing at  $\sim 16:10$  UT during the decaying phase of a major FD, and lasting for several hours. It was a relatively small GLE event, as the maximum relative increase in the SOPO count-rates was  $\sim 6\%$ . The two-peak structure observed in the HEPAD profiles is also evident in the relatively high-cutoff stations, including MGDN.

## 3.2 Interplanetary and geomagnetic data

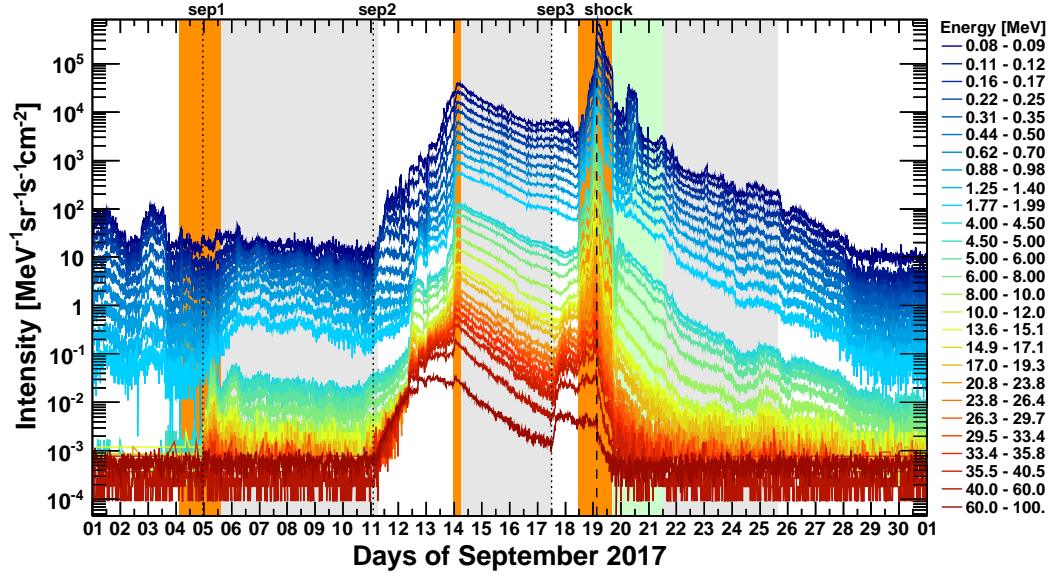
The aim of this section is to describe the SW structures influencing the near-Earth environment in early September 2017, and help to interpret the particle observations discussed in the previous sections. In particular, the profile of the IMF intensity, the IMF latitude in GSE coordinates and the SW speed are reported in panels a), b) and c) of Figure 1, respectively. Data are based on the OMNIWeb database (<http://OMNIWeb.gsfc.nasa.gov>), which provides in-situ observations time-shifted to the bow shock nose of

211 the Earth (King & Papitashvili, 2004). Specifically, combined ACE and Wind data (red,  
 212 1-hr resolution) are superimposed on DSCOVR points (blue, 5-min resolution). Gray shad-  
 213 ing indicates corotating high speed streams (HSSs), while the green regions are inter-  
 214 planetary CMEs (ICMEs; see, e.g., Kilpua et al. (2017); Zurbuchen & Richardson (2006)  
 215 and references therein); as discussed below, the orange shading emphasizes the presence  
 216 of a magnetic cloud (MC) structure.

217 Three interplanetary shocks passed by during this interval at the times indicated  
 218 by the vertical dashed lines. The first shock, marked by the commencement of a minor  
 219 geomagnetic storm at 23:43 UT on 6 September, as evident in the temporal profile of  
 220 the  $Dst$  index reported in panel g) of Figure 1, was driven by the interplanetary coun-  
 221 terpart of the CME observed by SOHO/LASCO on 4 September at  $\sim 19$  UT and asso-  
 222 ciated with the first SEP event which shows a local enhancement at low energies in the  
 223 vicinity of the shock. The first ICME interval indicated (shaded green) following the shock  
 224 was suggested by Shen et al. (2018), though the usual SW temperature ( $T_p$ ) decrease (Richard-  
 225 son & Cane, 1995) was not present, and it was associated with a decrease in the low-energy  
 226 particle intensity enhancement associated with this shock. The second ICME interval,  
 227 following this shock and commencing at  $\sim 19:40$  UT, did have a clear  $T_p$  relative reduc-  
 228 tion (and increase in the helium-proton ratio) and was present at Earth at the time of  
 229 arrival of the second shock, at 23:00 UT on 7 September (based on the storm sudden com-  
 230 mencement time). This shock was associated with the CME observed by SOHO/LASCO  
 231 on 6 September at 12:24 UT that was also associated with the second SEP event in Fig-  
 232 ure 1. Again there is a low-energy particle enhancement in the vicinity of this shock. An  
 233 intense geomagnetic storm occurred with  $Dst$  reaching  $-124$  nT early on 8 September,  
 234 as displayed in panel g) of Figure 1, following strong ( $\sim 30$  nT) southward (negative lat-  
 235 itude, see panel b) magnetic fields that were caused by the second shock compressing the  
 236 southward fields in the ICME through which it was propagating.

237 The ICME following this shock had two components. The first, marked by the or-  
 238 ange shading in Figure 1, exhibited many of the signatures of a magnetic cloud (MC)  
 239 MC (e.g., Klein & Burlaga (1982)), including ~~an~~ a distinct enhanced but declining IMF  
 240 intensity, declining SW speed, and low  $T_p$ , as well as enhanced He/proton ratio and oxy-  
 241 gen charge states, and bi-directional suprathermal electron beams. However, there was  
 242 no significant rotation of the IMF vector, so it may be termed a “MC-like” ICME (Wu  
 243 & Lepping, 2015); for brevity, we will refer to this region as the “MC” (shaded orange).  
 244 It was followed by a second, extended ICME structure (green shading) characterized by  
 245 a low variance, slightly enhanced, near-radial sunward magnetic field, depressed  $T_p$ , a  
 246 continuing decline in SW speed, and bidirectional suprathermal electrons. Following a  
 247 recovery as the field turned temporarily northward, a second peak in  $Dst$  ( $-109$  nT) was  
 248 driven by southward fields ( $\sim 17$  nT) inside the MC. Then, a recovery occurred as the  
 249 field returned northward in the following region of this ICME (shaded green). There is  
 250 a gap in the OMNIWeb data near the end of this region, but the DSCOVR data sug-  
 251 gest that it extended to  $\sim 00$  UT on 11 September based on the end of this region of low  
 252 variance, near-radial, magnetic field. This ICME was followed by a brief HSS (gray shad-  
 253 ing on 11–12 September) probably attributed to a weak influence from a negative po-  
 254 larity coronal hole. The SEP data show a local decrease during passage of the MC at  
 255 all energies from tens of keV to the peak of the FD observed by NMs.

256 A third shock on 12 September at  $\sim 20:02$  UT (storm commencement time) was likely  
 257 produced by the passage of the eastern flank of the shock associated with the 10 Septem-  
 258 ber event. This is consistent with the glancing blow with an arrival time of 13 Septem-  
 259 ber,  $\sim 02$  UT  $\pm 7$  hours based in ENLIL+CONE modeling indicated in the DONKI database.  
 260 However, closer examination of the SW data indicates that this was not a fully-steepened  
 261 shock. The subsequent lack of ICME-like signatures, in particular low  $T_p$ , indicates that  
 262 the associated ICME did not encounter Earth, consistent with the far western origin of  
 263 this event. Finally, a long-duration HSS was observed on 14 September, probably asso-



285

286 **Figure 2.** Temporal profiles of proton intensities measured by the SEPT, LET and HET in-  
 287 struments on board STEREO-A during September 2017. The vertical dotted and dashed lines  
 288 mark the onset of the SEP events and the time of the shock, respectively. The green and gray  
 289 areas indicate the periods of the ICMEs and HSSs, respectively. In this case, the orange shading  
 290 marks the CIRs.

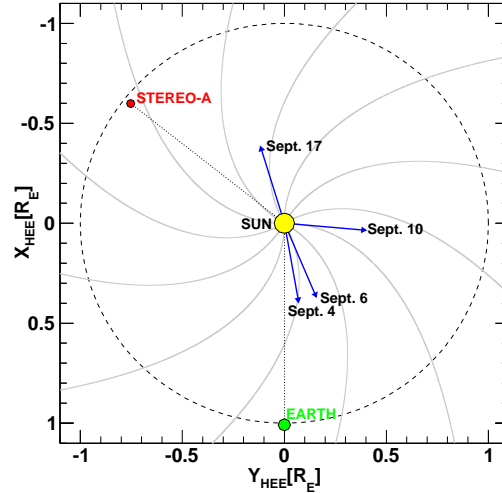
264 associated with the low-latitude extension of the northern polar coronal hole that passed cen-  
 265 tral meridian on 10 September. It carried an intermittent southward IMF and its effect  
 266 on the Earth endured for several days, triggering a moderate geomagnetic storm. The  
 267 SEP data show an enhancement at the lowest energies in the vicinity of the shock, and  
 268 also a rapid intensity decrease with the arrival of the HSS on September 14 which ter-  
 269 minated the event at low energies (below few hundreds of keV), while an extended decay,  
 270 already started before the HSS passage, can be observed at higher energies.

271 Returning to the onset of the 10 September event, this evidently occurred close to  
 272 the time when Earth was moving from an ICME to a HSS, so we suggest that the dou-  
 273 ble peak in the particle intensity at the highest energies may be associated with this tran-  
 274 sition, resulting in an improved connection to the particle source. This feature is less ev-  
 275 ident at lower energies. Possible reasons may be that the source of the high-energy par-  
 276 ticles was more spatially confined, and hence connectivity was more critical for the de-  
 277 tection of particles, and the low-energy particle intensities were still rising when Earth  
 278 exited the ICME whereas the highest energies had started to decay. Guo et al. (2018)  
 279 also proposed a second particle injection at the shock through merging of the ICME as-  
 280 sociated with the 10 September event with the two ICMEs that originated on 9 Septem-  
 281 ber from the same AR with similar directions. However, there does not appear to be ev-  
 282 idence of such a second particle injection in the available radio data from STEREO-A  
 283 or Wind, that clearly show only emissions associated with the original onset of the SEP  
 284 event.

### 296 3.3 Stereo observations

297 STEREO-A observations during this period made  $\sim 128$  deg east of Earth (see Fig-  
 298 ure 3) provide additional information on the SEP events discussed above and their lon-  
 299 gitudinal extent. Figure 2 displays the temporal profiles of proton intensities measured





291

292 **Figure 3.** Location of the Earth (ACE, GOES) and STEREO-A in Heliocentric Earth Ecliptic (HEE) coordinates during September 2017. The three arrows indicate the direction of the  
 293 parent flares of CMEs associated with the three SEP events observed at Earth/STEREO-A. The  
 294 nominal Parker-spiral IMF lines assuming  $V_{SW}=450 \text{ km s}^{-1}$  are also reported.  
 295

300 by the Solar Electron and Proton Telescope (SEPT; 0.084–6.5 MeV, 10-min resolution),  
 301 the Low Energy Telescope (LET; 4–12 MeV, 10-min resolution) and the High Energy  
 302 Telescope (HET; 13.6–100 MeV, 15-min resolution). In case of SEPT, only selected chan-  
 303 nels are shown for the sake of simplicity. As in Figure 1, the grey shading indicates HSSs  
 304 observed at STEREO-A, but here, orange shading indicates corotating interaction re-  
 305 gions (CIRs) at the stream leading edges, inferred from inspection of the STEREO-A  
 306 plasma and magnetic field data, not shown here.

307 The initial SEP enhancement in Figure 2 was associated with the 4 September event,  
 308 at  $\sim W143$  deg relative to the spacecraft longitude, while it was passing through the CIR  
 309 marking the arrival of a HSS. The prompt rise in the proton intensity suggests that par-  
 310 ticles propagated rapidly from the eastern flank off of the shock. There is a hint of an  
 311 increase from the 6 September event, but it is not compelling on the ongoing event. A  
 312 significant enhancement was registered early on 11 September, demonstrating that the  
 313 10 September event was very broad in longitude even at high energies, as the parent flare  
 314 was located at  $\sim E145$  deg relative to STEREO-A. In this case the magnetic footpoints  
 315 of STEREO-A were connected to the western flank of the shock, and measured inten-  
 316 sities exhibit a much more gradual increase. The delayed arrival ( $>10$  hours later than  
 317 the flare onset) may be attributed to cross-field diffusion in the SW. The event duration  
 318 can be inferred to be much longer with respect to near-Earth observations, well beyond  
 319 the onset of another high-energy event occurring on 17 September at  $\sim 12$  UT from the  
 320 same AR when at  $\sim W167$  ( $\sim E40$  of STEREO-A), that evidently was not observed at  
 321 Earth. This event was linked to a fast halo CME with a  $1385$  ( $1404$ )  $\text{km s}^{-1}$  speed ac-  
 322 cording to the CDAW (DONKI) catalog.

323 An interesting feature is the non-energy-dispersive increase in intensity early on  
 324 14 September which was associated with entry into – crossing of the stream interface –  
 325 a corotating HSS. This suggests that connection to the particle event and/or particle trans-  
 326 port in longitude was more favorable in the stream than in the preceding SW. In par-  
 327 ticular, a study based on the solar energetic particle event modeling (SEPMOD) of this  
 328 event (Luhmann et al., 2018) suggests that STEREO-A may have become connected to

329 the shock associated with the 10 September event beyond 1 AU at this time. Thus, the  
 330 observations suggest that field lines in the HSS were connected to this shock, but those  
 331 in the preceding slow SW were more poorly connected.

332 An interplanetary shock arrived on 19 September at 02:56 UT, when STEREO-  
 333 A was passing a CIR. At the same, the SW speed exceeded  $800 \text{ km s}^{-1}$  and a significant  
 334 enhancement of low-energy protons was observed. The CIR was followed by the arrival  
 335 of an ICME, as suggested by the drop in density and temperature, and an enhanced field  
 336 with a rotation, followed by a weaker, smoother field. The ICME caused a FD of proton  
 337 intensities. Then another HSS reached the spacecraft. Such interpretation is sup-  
 338 ported by the results of the ENLIL+CONE model in DONKI, with the flank of the ICME  
 339 passing STEREO-A at the time of a stream leading edge.

#### 340 4 SEP spectral analysis

341 In this section, the SEP observations introduced above will be used to construct  
 342 energy spectra over a wide energy range. The GOES data are affected by significant un-  
 343 certainties related to the poor resolution of the detector and high contamination by out-  
 344 of-acceptance particles (Bruno, 2017). In addition, the intensities measured by the HEPAD  
 345 channels and, to a lesser extent, the highest energy channels of the EPEADs, include a  
 346 high background associated with galactic cosmic rays (GCRs).

347 To improve the reliability of the EPEAD/HEPAD spectroscopic measurements, we  
 348 take advantage of two different cross-calibration schemes. For the data points below 80  
 349 MeV (P2–P5 channels), the mean energies by Sandberg et al. (2014) are used, based on  
 350 a calibration study of the Energetic Particle Sensors (EPSs) on board GOES-5, -7, -8,  
 351 and -11, using as reference the observations of the Goddard Medium Energy (GME) ex-  
 352 periment on board the Interplanetary Monitoring Platform-8 (IMP-8); the derived cross-  
 353 calibrated energies have been validated by Rodriguez et al. (2017) by comparison with  
 354 the STEREO data. Consistent with Sandberg et al. (2014), no background correction  
 355 is applied to the EPEAD intensities. This may result in an overestimate when SEP in-  
 356 tensities are low; conservatively, a 20% systematic uncertainty is assumed. To avoid east-  
 357 west effects (Rodriguez et al., 2010), more relevant at lower energies, only observations  
 358 from the westward viewing EPEADs are used.

359 The GOES data points above 80 MeV are based on Bruno (2017), who took ad-  
 360 vantage of the SEP measurements of the Payload for Antimatter Matter Exploration and  
 361 Light-nuclei Astrophysics (PAMELA) (Bruno et al., 2018) to calibrate the two most en-  
 362 ergetic channels (P6–P7) of the EPEADs and the four HEPAD channels (P8–P11), for  
 363 both GOES-13 and -15 units. As east-west effects are negligible at high energies, data  
 364 from both westward and eastward looking EPEADs are used in this range. A background  
 365 correction is applied by subtracting the average intensity measured during the 24-hr quiet  
 366 solar period prior to the SEP events. ~~Statistical uncertainties take into account the background~~  
 367 ~~subtraction.~~ It should be noted that derived “effective” mean energies represent aver-  
 368 age values and do not account for spectral index variations. A 20% (30%) systematic un-  
 369 certainty is assumed for the EPEAD (HEPAD) points, based on the comparison with  
 370 PAMELA measurements (Bruno, 2017).

371 In case of ACE and STEREO instruments, the background in each energy bin is  
 372 evaluated as the minimum intensity measured during a 20-day interval prior to the SEP  
 373 events, based on 1-hr resolution data. To a first approximation, the mean energy val-  
 374 ues are obtained by estimating the logarithmic center of each bin. **However, since the**  
 375 **two highest-energy channels of HET span a relatively much wider range (40–60 MeV and**  
 376 **60–100 MeV, respectively), the corresponding “true” mean energies are significantly af-**  
 377 **ected by spectral shape variations and, thus, the above assumption is no longer reason-**  
 378 **able. Consequently, a different approach based on Lafferty & Wyatt (1995) is used in**

379 this case:

$$380 \quad E_{mean} = \left[ \frac{E_{max}^{1-\gamma} - E_{min}^{1-\gamma}}{(E_{max} - E_{min})(1-\gamma)} \right]^{-\frac{1}{\gamma}}, \quad (1)$$

381 where  $E_{min}$  and  $E_{max}$  are the channel lower and upper energy limits, and  $\gamma$  is the spec-  
382 tral index derived by the power-law fit of HET spectral points between 30–40 MeV.

383 The “spikes” in the ACE temporal profiles of intensities, attributable to ions prop-  
384 agating upstream from the Earth’s bow shock (see Section 3.1.1), are removed. Since the  
385 lowest energy channels are affected by electron contamination, only the intensities above  
386 300 keV are considered; in addition, a 20% systematic uncertainty is associated with the  
387 data points.

388 In general, statistical errors are evaluated by accounting for the GCR background  
389 subtraction, by using 68.27% confidence level intervals for Poisson signal/background dis-  
390 tributions according to Feldman & Cousins (1998). Statistical and systematic uncertain-  
391 ties are summed in quadrature.

392 Event-integrated energy spectra are obtained by summing up the SEP intensities  
393 measured in each energy bin over the event duration. The integration interval is com-  
394 puted by identifying the event start/stop times in the intensity temporal profiles. When  
395 a new event commences while a preceding one was still in progress, the onset time of the  
396 second event is set as the end time of the first event. Consequently, the spectrum for the  
397 second event will include a contribution from the decay of the previous event. Finally,  
398 it should be noted that, since the background correction is based on pre-event intensi-  
399 ties, SEP event-integrated intensities are somewhat underestimated – especially above  
400 several tens of MeV – if FD periods are present, such as during the decaying phase of  
401 the 6 September event and the initial phase of the 10 September event.

#### 402 4.1 Spectral fits

403 In order to characterize the estimated event-integrated energy spectra, we fit them  
404 with several spectral shapes. A first, purely empirical model is given by the double power-  
405 law function by Band et al. (1993) (hereafter Band function):

$$406 \quad \Phi_{Band}(E) = \begin{cases} A E^{-\gamma_a} \exp(-E/E_0) & \text{for } E < (\gamma_b - \gamma_a) E_0, \\ A E^{-\gamma_b} [(\gamma_b - \gamma_a) E_0]^{(\gamma_b - \gamma_a)} \exp(\gamma_a - \gamma_b) & \text{for } E > (\gamma_b - \gamma_a) E_0, \end{cases} \quad (2)$$

407 originally developed to fit gamma-ray burst spectra. It is defined by four free param-  
408 eters ( $A$ ,  $\gamma_a$ ,  $\gamma_b$ ,  $E_0$ ), providing a smooth transition between two energy regions char-  
409 acterized by different spectral indices ( $\gamma_a$  and  $\gamma_b$ ); the transition energy is given by  $(\gamma_b -$   
410  $\gamma_a) E_0$ . While such spectral breaks, typically occurring at energies of few tens of MeV,  
411 have been often associated with the limits of shock acceleration (see, e.g., Desai et al.  
412 (2016) and references therein), they can be explained by accounting for interplanetary  
413 transport effects (Li & Lee, 2015; Zhao et al., 2016).

414 A second functional form is based on Ellison & Ramaty (1985) (hereafter referred  
415 as E-R), and consists of a power-law spectrum modulated by an exponential:

$$416 \quad \Phi_{E-R}(E) = A E^{-\gamma} \exp(-E/E_r), \quad (3)$$

417 where  $E_r$  is the cutoff or rollover energy. In the scenario of diffusive shock acceleration,  
418 the spectral rollover is attributed to particles escaping the shock region during accel-  
419 eration due to effects mostly related to the limited extension and lifetime of the shock (Lee,  
420 2005; Lee & Ryan, 1986). This function has been recently used by Bruno et al. (2018)  
421 to fit the time-integrated energy spectra of the high-energy (>80 MeV) SEP events ob-  
422 served by the PAMELA experiment.

In general, multiple spectral features can be present at different energies, and the above functional forms hardly reproduce the spectral shapes over the complete energy range of SEPs. In particular, the Band function reasonably describes the SEP spectra below several tens of MeV, but it reduces to a single power-law extending to infinity for energies much larger than the break energy; consequently, it can not be used to account for the high-energy (hundreds of MeV) spectral rollovers recently found in PAMELA observations (Bruno et al., 2018). In order to reproduce both the low-energy break and the high-energy rollover in the SEP spectra, Equations 2 and 3 can be combined into:

$$\Phi_{tot}(E) = \Phi_{Band}(E) \exp(-E/E_r), \quad (4)$$

i.e. a double-power law (Band) function multiplied by an (E-R) exponential cutoff. Hereafter we refer to the above functional form as the “combined” function.

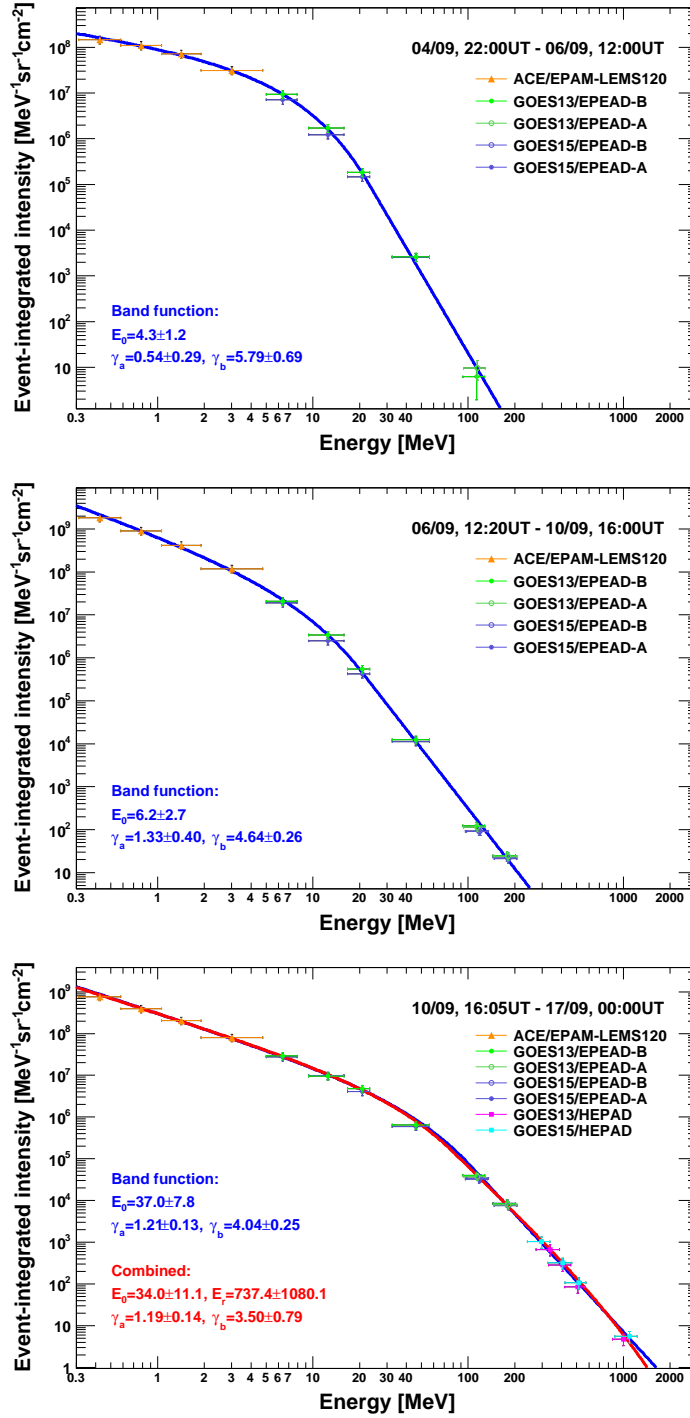
As a final remark we note that, overall, significant cross-correlations may exist between the fit parameters, in particular between the break/rollover energies and the spectral indices (Bruno et al., 2018; Desai et al., 2016), resulting in large parameter uncertainties. **Fit errors are evaluated with the MINOS technique (see, e.g., Ferbel (1993)).**

## 5 Results

The time-integrated energy spectra of the 4 and 6 September 2017 SEP events measured by ACE and GOES-13/15 above 300 keV are shown in top and middle panels of Figure 4, respectively. **The vertical error bars account for both statistical and systematic uncertainties.** The horizontal error bars denote the nominal energy ranges or, in the case of GOES, the “effective” energy ranges estimated by Sandberg et al. (2014) and Bruno (2017). The curves indicate the fits performed with the Band function; the fit parameters along with associated uncertainties are also reported. The Band function provides **a good fit of good fits to** the spectra, which are very soft ( $\gamma_b \approx 5.8$  and  $\gamma_b \approx 4.6$ , respectively) above the break energy (4.3 MeV and 6.2 MeV, respectively). In addition, the 4 September spectrum is almost flat below the break ( $\gamma_a \approx 0.5$ ). As reconstructed spectra are limited to energies below  $\sim 150$  MeV and  $\sim 200$  MeV, respectively, no reliable assumption can be made regarding an high-energy spectral rollover.

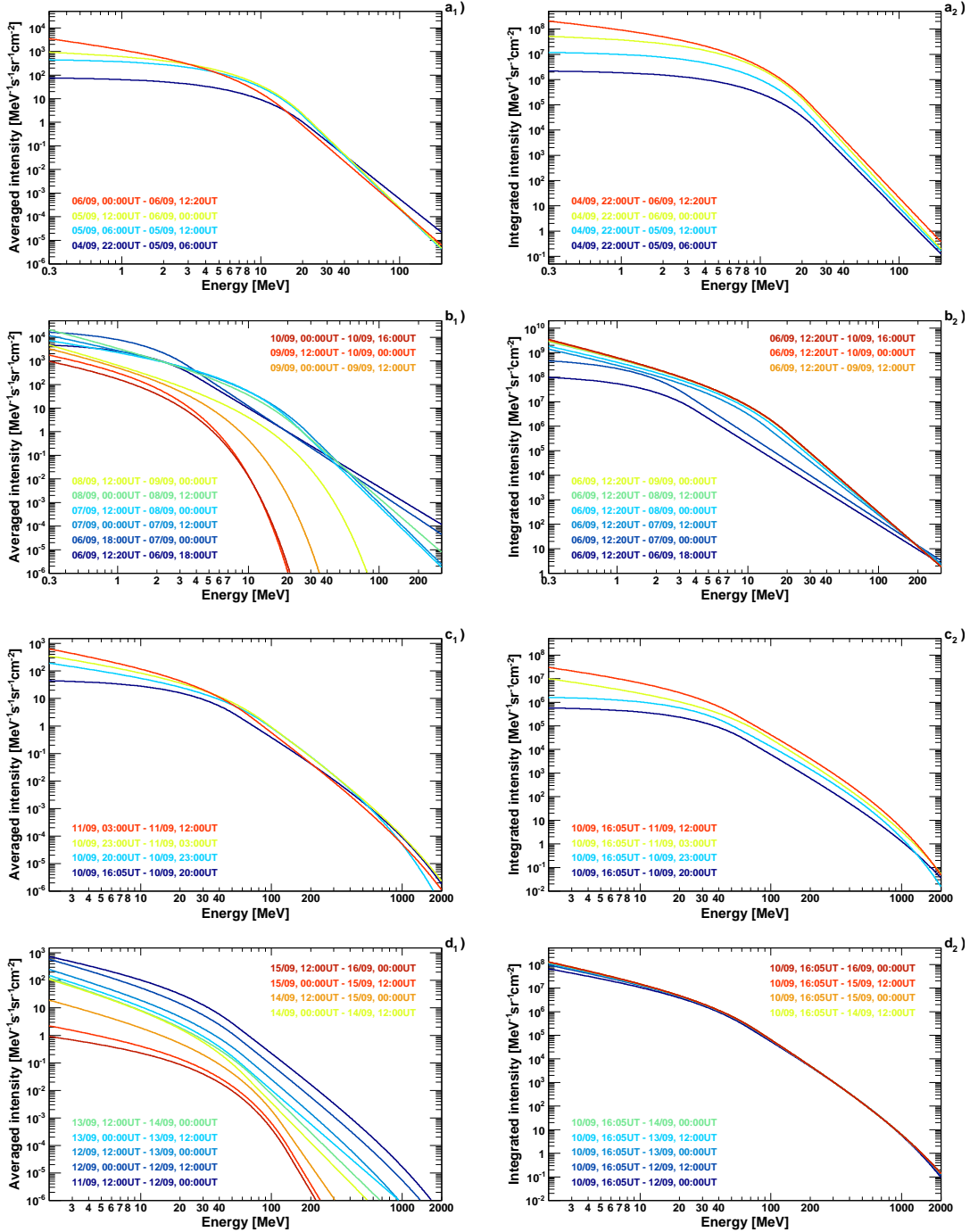
In contrast, as demonstrated in the bottom panel of Figure 4, the spectrum measured for the 10 September SEP event extends up to  $\sim 1$  GeV. Since faster shocks can accelerate particles to higher energies, the high energies reached in the 10 September event are consistent with the associated ultra-fast CME (see Table 1). In addition, **with respect in comparison** to 4 and 6 September events, a powerful radio emission at higher frequencies accompanied the event (Chertok, 2018), implying that SEPs were accelerated closer to the Sun, where the magnetic field is more intense **hence the maximum SEP energy is higher and hence the maximum energy to which SEPs can be accelerated is higher** (Gopalswamy et al., 2017; Zank et al., 2000). Gopalswamy et al. (2018) estimated a shock height of 1.4 Rs at Type II onset, in agreement with previous GLE observations. For comparison, the steeper radio spectrum with a peak at lower frequencies measured during the 4 September event is indicative of a post-eruption origin, while the 6 September event had intermediate features (Chertok, 2018).

The high-energy data in the spectrum of the 10 September event suggest the presence of a rollover – albeit with large uncertainties due to the few points and their error bars – similar to that found in the high-energy SEP observations reported by the PAMELA mission (Bruno et al., 2018), that may be consistent with the limits of diffusive shock acceleration (see Section 4.1). Comparing the fits performed with the Band (blue) and the combined (red curve) functions, we obtain a  $\sim 1.36$  value for the ratio of the corresponding reduced  $\chi^2$  ( $F$ -test). Therefore the spectral shape is better reproduced by the latter functional form, which provides a reasonable fit of the data points in the full energy range accounting for both the low-energy break (34 MeV) and the high-energy rollover



438

439 **Figure 4.** The time-integrated energy spectra of the 4, 6 and 10 September 2017 SEP events  
 440 (top, middle and bottom panel, respectively) measured by ACE and GOES-13/15. The vertical  
 441 error bars account for statistical and systematic uncertainties. The horizontal error bars denote  
 442 the channel nominal/effective energy ranges. The blue and the red curves denote the fits per-  
 443 formed by using the Band and the combined functions. The integration intervals, along with fit  
 444 parameters and associated uncertainties are also reported with the same color code.



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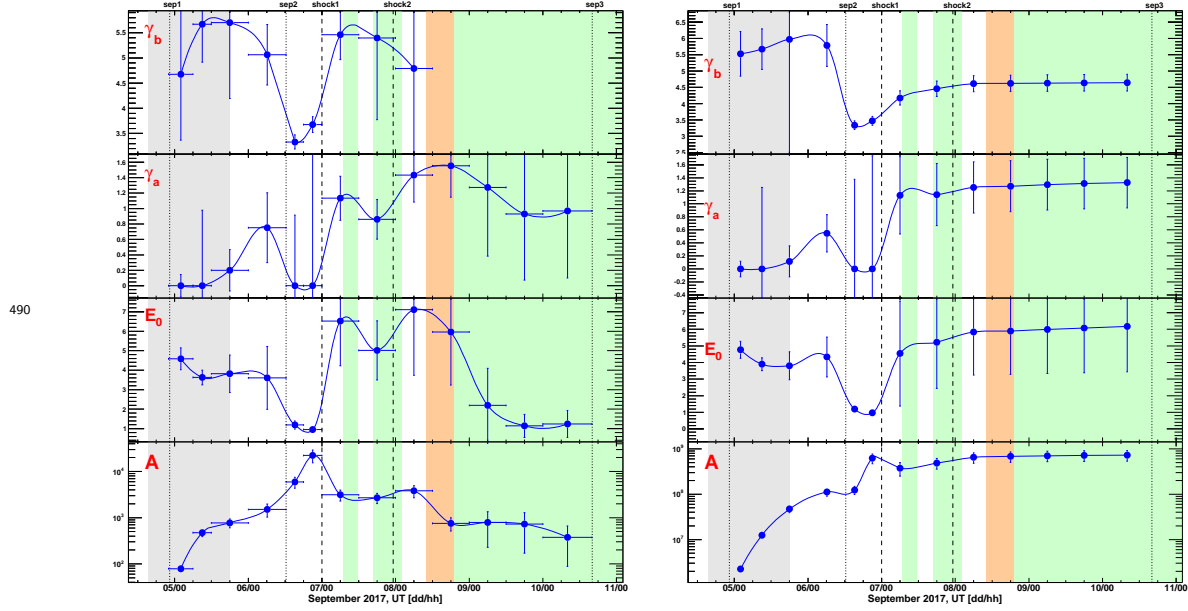
**Figure 5.** Spectral fits obtained for the 4 September event (a), the 6 September event (b) and the long-duration 10 September event (c and d). Left panels are based on the energy spectra averaged during successive time intervals, while right panels show the fits of the corresponding spectra integrated over cumulative intervals, with same color code (see labels).

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(737 MeV). However, the interpretation of spectra shapes is significantly complicated by a series of overlapping events and related interplanetary structures (local shocks, ICMEs and HSSs), as discussed in Section 3.2, influencing SEP intensities hence spectra. Con-

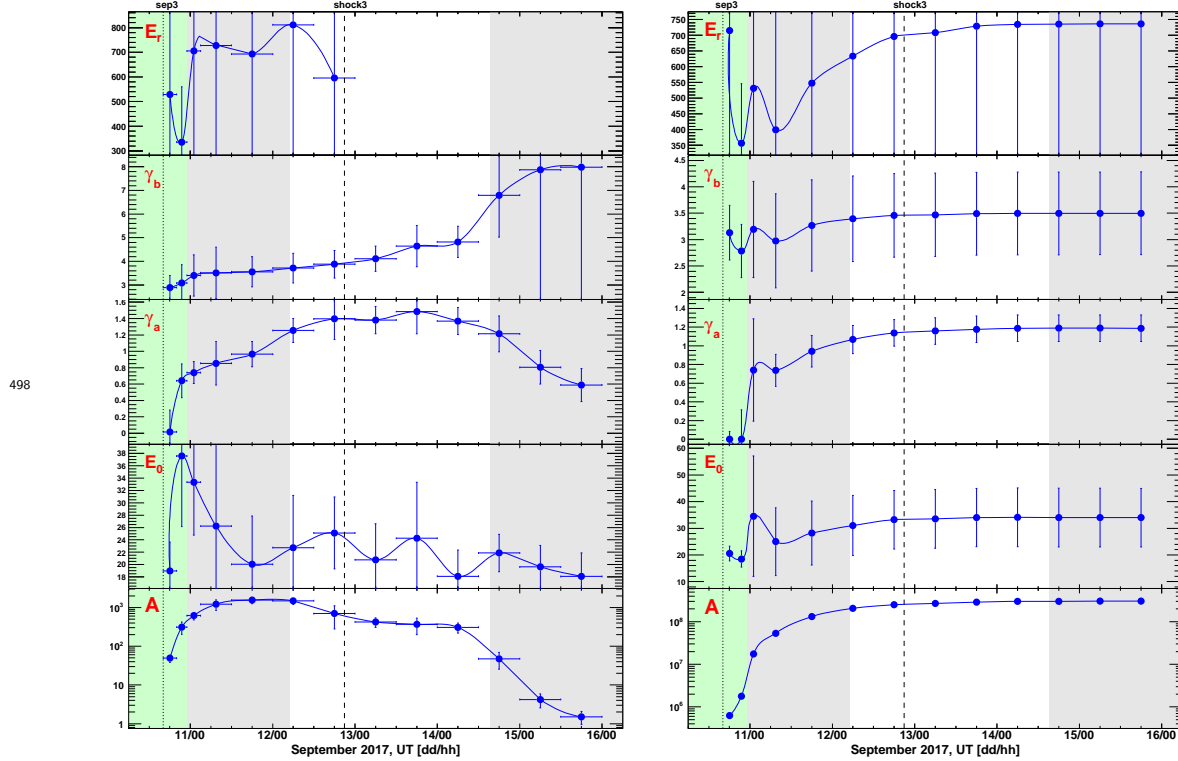


491 **Figure 6.** Left - Evolution of the Band fit parameters for the average spectra of the 4 and  
 492 6 September 2017 events reported in left panels of Figure 5. Right - Evolution of the Band fit  
 493 parameters for the cumulative spectra of the 4 and 6 September 2017 events reported in right  
 494 panels of Figure 5. The curves are to guide the eye. The vertical error bars account for fit param-  
 495 eter uncertainties. The vertical dotted and dashed lines mark the onset of the SEP events and  
 496 the time of the shocks, respectively. The green, orange and gray areas indicate the periods of the  
 497 ICMEs, MC and HSSs, respectively.

488 sequently, it is not realistic to account for the spectral features only in terms of parti-  
 489 cle acceleration.

### 506 5.1 Spectra temporal evolution

507 ~~Figure 5 displays the temporal evolution of the spectral shapes. Top panels show~~  
 508 ~~the results relating to the 4 and 6 September events, while bottom panels refer to the~~  
 509 ~~long-duration event on 10 September. The left panels in Figure 5 display the fits of the~~  
 510 ~~SEP spectra obtained in successive time intervals during the 4 September event (a), the~~  
 511 ~~6 September event (b) and the long-duration 10 September event (c and d). The fits for~~  
 512 ~~the 4 and 6 September events are based on the Band function, while the combined func-~~  
 513 ~~tional form was used for the 10 September event. Differential~~ The spectra are evaluated  
 514 by averaging intensities on a 12-hr timescale; a ~~smaller~~ **higher** time resolution (3–6 hours)  
 515 is used during the initial phase of the events (see labels). In addition, only data above  
 516 2 MeV are included for the 10 September event due to the difficulty in fitting the com-  
 517 plete energy spectrum, which exhibits a further softening at lower energies **in the early**  
 518 **phase attributable to a low energy component from the previous event.** The time vari-  
 519 ations of the fit parameters are summarized in **left panels of Figures 6 and 7.** ~~the results~~  
 520 ~~for the 4 and 6 September events, based on the Band function, and for the 10 September~~  
 521 ~~event, based on the combined functional form of Equation 4, are displayed in the left and~~  
 522 ~~the right panels, respectively.~~ It should be stressed that fit parameters are typically cor-  
 523 related. **The right-hand panels of Figure 5 show the cumulative spectra for each event**  
 524 **integrated up to the end time of each spectrum in the left panels and indicated with the**  
 525 **same color code. The corresponding fits to the cumulative spectra are shown in the right**  
 526 **panels of Figures 6 and 7.**

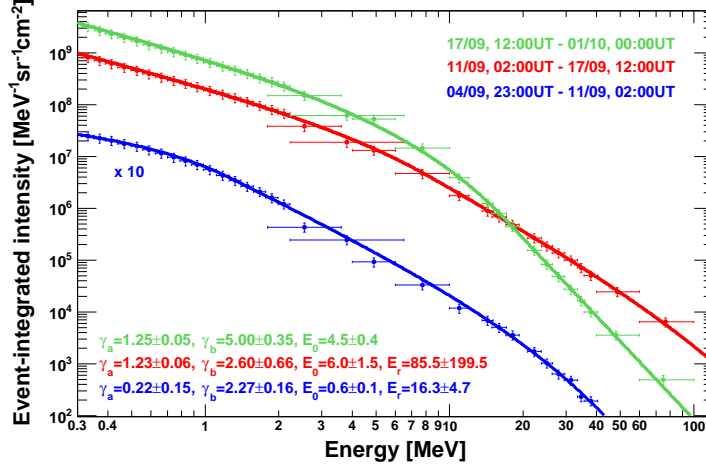


498 **Figure 7.** Left - Evolution of the combined fit parameters for the average spectra of the 10  
 499 September 2017 event reported in left panels of Figure 5. Right - Evolution of the combined fit  
 500 parameters for the cumulative spectra of the 10 September 2017 event reported in right panels of  
 501 Figure 5. The vertical error bars account for fit parameter uncertainties. The curves are to guide  
 502 the eye. The vertical dotted and dashed lines mark the onset of the SEP event and the time of  
 503 the shock, respectively. The green and gray areas indicate the periods of the ICMEs and HSSs,  
 504 respectively. The green and gray areas indicate the periods of the ICMEs and HSSs,  
 505 respectively.

527 The initial phase of the 4 September event – as well as the other events – was char-  
 528 acterized by velocity dispersion effects, with higher-energy particles arriving earlier, re-  
 529 sulting in relatively hard spectra. The spectra was almost flat at low-energies ( $\gamma_a \approx 0$ ).  
 530 In the subsequent three intervals the high-energy part of the spectrum did not change  
 531 significantly, in particular the break energy remained constant, while the low-energy spec-  
 532 trum became softer due to the increasing intensities.

533 The spectral evolution of the 6 September event can be divided into three phases.  
 534 During the first one (first two time bins), the break energy was very low ( $E_0 \approx 1$  MeV)  
 535 and the spectrum was flat ( $\gamma_a = 0$ ) and relatively hard ( $\gamma_b \approx 3.5$ ) in the energy ranges be-  
 536 low and above the spectral transition, respectively. Derived spectra, especially at low en-  
 537 ergies, include a particle component associated with the ongoing 4 September event, along  
 538 with the related shock. The second phase (subsequent three time bins) commenced af-  
 539 ter the arrival of the interplanetary shock at the end of 6 September: the break energy  
 540 increased (5–6 MeV) and the spectrum became softer ( $\gamma_a \approx 1$  and  $\gamma_b \approx 5$ ). The arrival  
 541 of the shock-ICME complex structure at the end of 7 September caused large FD effects,  
 542 inducing an enhancement of  $E_0$  and  $\gamma_a$ . The third phase (last four time bins) started  
 543 with arrival of the MC, corresponding to the peak of the FD, and extended over its de-  
 544 caying phase up the onset of the following SEP event. At the same time, intensities de-  
 545 creased significantly, especially at high-energy. As a consequence, the estimated spec-





568

569 **Figure 8.** Time-integrated energy spectra of the 4, 10 and 17 September 2017 SEP events  
 570 (blue, red and green points respectively) measured by STEREO-A. **The vertical error bars**  
 571 **account for statistical and systematic uncertainties.** The horizontal error bars show the nomi-  
 572 nal range of each energy channel. The curves represent the fits based on the Band (for the 17  
 573 September event) and the combined (for the 4 and 10 September events) functions. The integra-  
 574 tion intervals, along with the fit parameters and associated uncertainties are also reported with  
 575 the same color code.

546 trum is better reproduced by a truncated power-law (E-R function), i.e. without a tran-  
 547 sition to a high-energy spectral index, so no value of  $\gamma_b$  during this phase is reported in  
 548 Figure 6.

549 A complex temporal evolution characterized the initial phase of the 10 September  
 550 event. During the first three time bins, the spectrum was relatively hard with  $\gamma_a$  almost  
 551 constant ( $\sim 0.6$ ) and  $\gamma_b$  very slowly increasing. At the same time, two peaks were observed  
 552 in the intensity profiles of the HEPADs; a minimum of the rollover energy  $E_r$  and a maxi-  
 553 mum of the break energy  $E_0$  were found in the interval between the peaks (20–23 UT).  
 554 As discussed in section 3.1.1, there may be alternative interpretations of this feature. In  
 555 particular, the event commenced in the recovery phase of the FD, while the Earth was  
 556 in a ICME region, and the second peak occurred after the arrival of a HSS following the  
 557 trailing edge of the ICME. The SEP event lasted for several days, with a monotonic in-  
 558 crease of a  $\gamma_b$  and, hence, a gradual softening of the spectrum, as the intensities of the  
 559 higher energy particles accelerated earlier and closer to the Sun decline. The break en-  
 560 ergy remained relatively stable, within uncertainties, around a value of  $\sim 20$  MeV. Af-  
 561 ter 13 September the rollover energy was probably higher than the maximum explored  
 562 energy, and the spectra were better reproduced by the Band function. A significant sup-  
 563 pression of intensities was registered as a consequence of the arrival of a HSS on 14 Septem-  
 564 ber which terminated the event at low energies and caused an abrupt increase of  $\gamma_b$  from  
 565 5 to 7. Starting on 16 September the derived spectra between 2 and a few tens of MeV  
 566 can be described by a simple power-law gradually approaching the background inten-  
 567 sities, so results are not reported in Figure 7.

576

## 5.2 Comparison with STEREO-A spectra

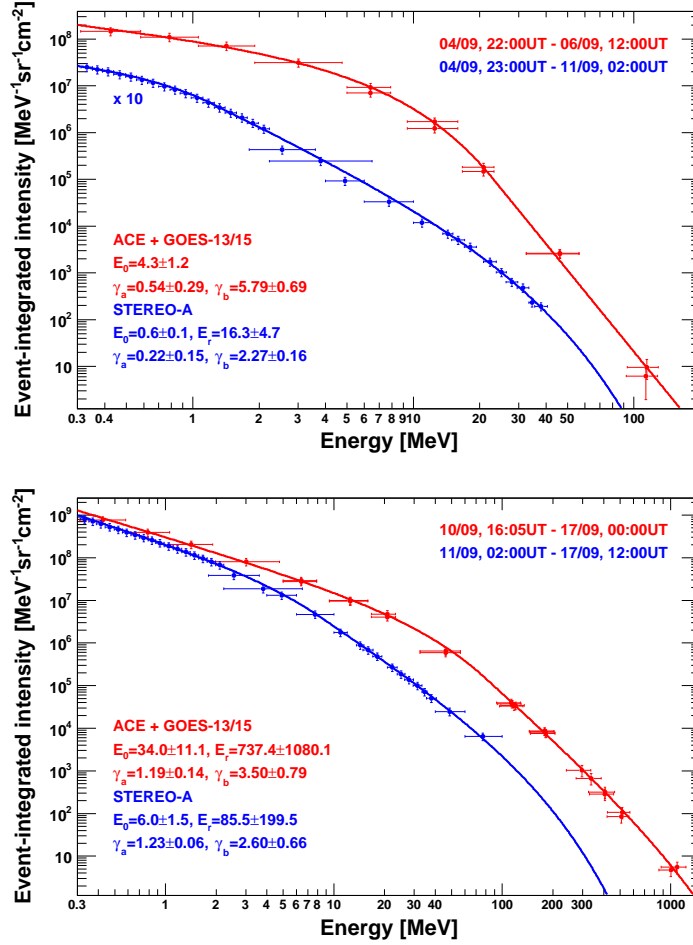
577 Figure 8 displays the time-integrated energy spectra of the 4, 10 and 17 Septem-  
 578 ber events measured by STEREO-A (see Section 3.3), denoted by blue, red and green  
 579 points respectively. The spectra extend over the full energy range (300 keV – 100 MeV)

580 covered by the SEPT, LET and HET instruments. The curves represent the fits based  
 581 on the Band (for 17 September event) and the combined (for the 4 and 10 September  
 582 events) functions. The integration intervals, along with the fit parameters and associ-  
 583 ated uncertainties are also reported with the same color code. The spectrum derived for  
 584 the 4 September event is much less intense, and was multiplied by 10 to improve the com-  
 585 parison. Albeit data points are limited to 40 MeV, it exhibits a break at very low en-  
 586 ergies ( $E_0 \approx 0.5$  MeV) along with a rollover at higher energies ( $E_r \approx 16$  MeV). In contrast,  
 587 the spectra of the other two events extend above 60 MeV. While the high-energy data  
 588 of the 10 September event spectrum suggest a rollover corresponding to  $E_r \approx 79$  MeV,  
 589 although affected by very large uncertainties due to the limited number of points, the  
 590 spectral shape of the 17 September event is significantly softer above the break energy  
 591 ( $\gamma_b \approx 5$ ); consequently, no rollover can be identified and the data are well reproduced by  
 592 the Band function. However, it should be noted that measured intensities include a con-  
 593 tribution from the previous event that is apparently larger at lower energies. In addi-  
 594 tion, a component of low-energy particles is associated with the interplanetary shock ar-  
 595 riving on 19 September (see Section 3.3). Finally, the spectrum is influenced by the FD  
 596 caused by the subsequent ICME, whose effects are not accounted for in the background  
 597 subtraction, as described in Section 4.

608 Figure 9 shows the comparison between ~~the fits of~~ the time-integrated energy spectra  
 609 measured by ACE and GOES-13/15 (red), and by STEREO-A (blue), during the 4  
 610 and 10 September SEP events (top and bottom panel, respectively). **The curves are the**  
 611 **fits based on combined functional form and, for the 4 September event spectrum measured by ACE and GOES, on the Band function. In case of STEREO-A, the fit are ex-**  
 612 **trapolated beyond the 100 MeV limit of the observations.** The integration intervals, along  
 613 with the fit parameters and associated uncertainties are also displayed with the same color  
 614 code. Overall, the spectra differ in both magnitude and shape. In particular, ~~the near-Earth~~  
 615 ~~intensities are more intense and the resulting spectra extend to higher energies~~ **the SEP**  
 616 **events are larger near the Earth and their spectra extend to higher energies.** Discrepan-  
 617 cies are emphasized during the 4 September event, with a  $\sim 100$  factor for the time-  
 618 integrated intensities at 1 MeV, while are less evident during the 10 September event.  
 619 Such differences can be mostly attributed to the different magnetic connection of the space-  
 620 craft: for both events, ACE and GOES footpoints were best connected to the solar event,  
 621 detecting higher ~~intensity magnitudes and more energetic, hence~~ **particle intensities and**  
 622 **harder spectra** (see, e.g., Hu et al. (2017)); ~~on~~ **On** the other hand, STEREO-A was con-  
 623 nected to the back side of the Sun (see Figure 3) and, as suggested by SEPMOD sim-  
 624 ulations (Luhmann et al., 2018), for the 10 September event it may have predominantly  
 625 detected particles streaming from the distant shock beyond 1 AU (see Section 3.3). STE-  
 626 REO observations demonstrate that this event was very broad in longitude ~~event~~ at high  
 627 energies. A major role was likely played by transport effects such as cross-field diffusion  
 628 and IMF corotation, possibly in combination with widespread particle sources associ-  
 629 ated with a CME-driven shock accelerating and injecting particles onto an extended re-  
 630 gion of the heliosphere (see, e.g., Lario et al. (2017); Richardson et al. (2014) and ref-  
 631 erences therein). Additional factors should be considered when comparing the two sets  
 632 of measurements, including the effects of SW structures. In particular, near-Earth ob-  
 633 servations of the 10 September event were influenced by the interplanetary counterpart  
 634 of the 6 September CME and the subsequent HSS (see Section 3.1.1). We also note that  
 635 measured SEP time-integrated spectra include a component from previous events and  
 636 that the used integration intervals are limited by the onset of the subsequent events, e.g.  
 637 the commencement of the 17 September event in case of STEREO-A.

### 639 5.3 Comparison with the 17 May 2012 GLE event

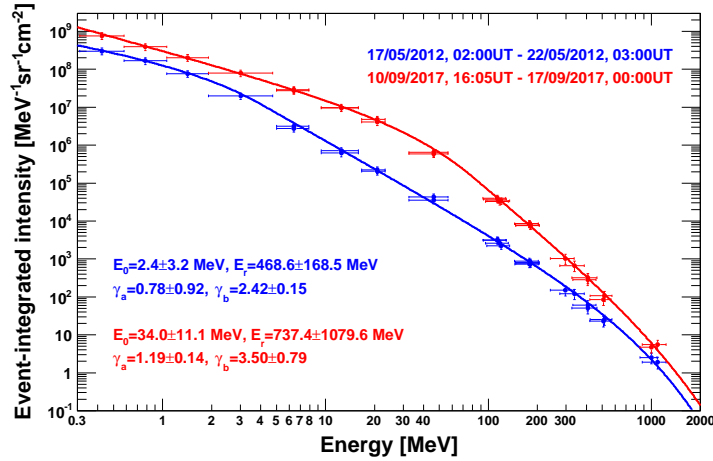
640 Figure 10 compares the time-integrated energy spectrum ~~fit~~ of the 10 September  
 641 2017 event (red) with that of the 17 May 2012 event (blue), associated with the previ-  
 642 ous GLE (n.71) of the solar cycle 24. Both spectral fits, **based on the combined func-**



598

599 **Figure 9.** Comparison between the fits of the time-integrated energy spectra measured by  
 600 ACE and GOES-13/15 (red), and by STEREO-A (blue), during the 4 and 10 September 2017  
 601 SEP events (top and bottom panel, respectively). The vertical error bars account for statistical  
 602 and systematic uncertainties; the horizontal error bars denote the channel nominal/effective  
 603 energy ranges. The curves are the fits based on the combined functional form (Equation 4) and,  
 604 in case of the 4 September event spectrum measured by ACE and GOES, on the Band function  
 605 (Equation 2). The integration intervals, along with the fit parameters and associated uncertain-  
 606 ties are also displayed with the same color code. The STEREO-A spectrum derived for the 4  
 607 September was multiplied by 10 to improve the comparison.

643 tional form, rely on ACE and GOES observations according to the procedure described  
 644 in Section 4. The used functional form is based on Equation 4. The integration inter-  
 645 vals along with derived fit parameters and related uncertainties are also shown with the  
 646 same color code. While the discrepancy in the absolute intensities reflects the much shorter  
 647 duration of the 17 May 2012 event, the two spectral shapes are quite different, with the  
 648 10 September 2017 event exhibiting a softer spectrum above several tens of MeV, with  
 649 higher break and rollover energies. This is consistent with PAMELA measurements (Bruno  
 650 et al., 2018), showing that higher energy rollovers tend to be associated with larger spec-  
 651 tral indices. Based on a simple power-law fit of the data points above the transition en-  
 652 ergies (78.4 MeV and 3.9 MeV, respectively), a spectral index value of  $4.05 \pm 0.03$  and



655

656 **Figure 10.** Fits of the Time-integrated energy spectra (based on Equation 4) of the 17 May  
 657 2012 (blue) and the 10 September 2017 (red) SEP events measured by ACE and GOES. The  
 658 vertical error bars account for statistical and systematic uncertainties. The horizontal error bars  
 659 denote the channel nominal/effective energy ranges. The curves are the fits based on the com-  
 660 bined functional form (Equation 4). The integration intervals, along with the fit parameters and  
 661 associated uncertainties are also reported with the same color code.

653  $2.97 \pm 0.20$  is obtained for the 10 September 2017 and the 17 May 2012 events, respec-  
 654 tively.

662 In contrast to NM observations, reporting a larger GLE during the 17 May 2012  
 663 event, a slightly larger SEP signal was measured by GOES/HEPAD during the 10 September  
 664 2017 event. Such disagreement can be explained by accounting for the harder peak spectrum  
 665 and for effects related to the much higher level of anisotropy measured during the 17 May  
 666 2012 event (Adriani et al. 2015; Bruno et al., 2016; Mishev et al. 2018). A minor contribution  
 667 can be attributed to the somewhat larger GCR background during the 10 September 2017  
 668 event as it occurred during a period of solar minimum.

669 In general, several concomitant factors potentially contribute to the differences in  
 670 the observed spectral shapes, such as the parent flare and CME parameters, the shock  
 671 morphology and evolution, the ambient conditions, the magnetic connection to Earth  
 672 and the interplanetary transport. The 17 May 2012 GLE event was peculiar because of  
 673 the moderately strong source: ~~a~~ an M5.6 flare linked to a  $1582 \text{ km s}^{-1}$  linear speed CME  
 674 in the CDAW catalog. Such values are significantly lower compared with those associ-  
 675 ated with the 10 September 2017 event ( $X8.2$  and  $3163 \text{ km s}^{-1}$ ). However, the former  
 676 event originated in a region characterized by a better longitudinal connectivity to Earth  
 677 (N11W76) than the latter event (S08W88), and the 10 September 2017 flare reached peak  
 678 intensity when the involved AR had just rotated over the western solar limb. In addi-  
 679 tion, Gopalswamy et al. (2018) proposed that the non-radial motion of the CME along  
 680 with the favorable B0 angle (the inclination of the solar equator to the ecliptic) rendered  
 681 the shock nose latitudinally well connected to Earth in case of the 17 May 2012 event,  
 682 while the opposite situation occurred during the 10 September 2017 event. Consequently,  
 683 it can be speculated that the protons detected near the Earth at highest energies were  
 684 accelerated mostly at the eastern flank of the shock, where acceleration is less efficient  
 685 and the SEP maximum energy is lower (Hu et al., 2017), resulting in a softer spectrum  
 686 with respect to better connected events such as the 17 May 2012 event.

687 However, the prevailing interplanetary conditions may significantly complicate such  
 688 arguments based on simple assumptions for the connectivity. For instance, according to  
 689 Rouillard et al. (2016) the magnetic connectivity between the 17 May 2012 solar event  
 690 and the near-Earth environment was established via a MC that erupted from the same  
 691 AR a few days before. ~~On the other hand,~~ **Similarly**, the 10 September 2017 event com-  
 692 menced while the Earth was in a ICME region, during the recovery phase of a FD caus-  
 693 ing a depression in observed intensities. ~~However,~~ **Since** the applied GCR background  
 694 correction does not account for such effects being based on the average intensities reg-  
 695 istered prior to the three SEP events (see Section 4), derived high-energy SEP intensi-  
 696 ties are somewhat underestimated. In addition, the double-peak feature exhibited by the  
 697 temporal profiles of high-energy intensities may be related to the influence of SW struc-  
 698 tures **on particle transport**. Finally, measured time-integrated intensities include a low-  
 699 energy contribution from the previous SEP event on September 6. Consequently, the “true”  
 700 SEP spectrum is supposed to be harder.

## 701 6 Summary and conclusions

702 Despite the near solar minimum conditions, an exceptional interval of solar activi-  
 703 ty occurred between 4–10 September 2017 during the late decay phase of solar cycle  
 704 24 that involved the complex AR NOAA 12673 located in the western solar hemisphere.  
 705 A large number of bright eruptions were observed, including four associated with X-class  
 706 flares. The X9.3 flare on 6 September and the X8.2 flare on 10 September are currently  
 707 the two strongest soft X-ray flares of solar cycle 24. Both were linked to fast CMEs, giv-  
 708 ing rise to SEP events measured by near-Earth spacecraft. In particular, the western limb  
 709 event on 10 September triggered a GLE recorded by several NM stations, the second GLE  
 710 (no.72) of the solar cycle. A further, smaller SEP event, detected late on 4 September,  
 711 originated from the M5.5 flare and the related CME that erupted on the same day.

712 In this work we analyzed the space-based proton measurements by ACE and GOES-  
 713 13/15 to study the time integrated spectra and spectral evolution of in a wide energy  
 714 range ( $\geq 300$  keV). The spectra show a low-energy spectral break at few/tens of MeV,  
 715 that is often attributed to the limits of diffusive shock acceleration, though interplan-  
 716 etary transport may also introduce such features in SEP spectra. In addition, the 10 Septem-  
 717 ber 2017 event spectrum, extending up to  $\sim 1$  GeV, exhibits a high-energy rollover simi-  
 718 lar to that reported in the recent SEP observations of the PAMELA experiment, that  
 719 may be ascribed to the limited extension and lifetime of the shock in the scenario of dif-  
 720 fusive shock acceleration. However, for the September 2017 period, the study of SEP fea-  
 721 tures, including the interpretation of spectra shapes, is significantly complicated by a se-  
 722 ries of overlapping events and interplanetary structures (local shocks, ICMEs and HSSs),  
 723 that influenced SEP intensities and hence the spectra. Consequently, it is not realistic  
 724 to account for the spectral features only in terms of particle acceleration. ~~In particular~~  
 725 **In addition**, a double peak in the high-energy proton intensity profile during the 10 Septem-  
 726 ber may have originated from a change in the connection conditions as the Earth moved  
 727 from an ICME into a HSS; available radio burst data disfavor the alternative interpre-  
 728 tation of a second particle injection.

729 Near-Earth SEP observations for these events have been compared with those re-  
 730 ported by STEREO-A. ~~In addition~~ **Furthermore**, we compared the spectrum for the 2017  
 731 September 10 event with that obtained for the 2012 May 17 event, associated with the  
 732 previous GLE in cycle 24. Differences in the spectra and their temporal evolution can  
 733 be mostly attributed to the different magnetic connection of the spacecraft with respect  
 734 to the shocks accelerating particles, but local interplanetary structures such as shocks,  
 735 ICMEs and HSSs also have a relevant impact. STEREO data demonstrate that the 10  
 736 September 2017 event was very broad even at high energies, suggesting significant trans-  
 737 port effects such as cross-field diffusion and IMF corotation in combination with the ex-  
 738 tended SEP source provided by the CME-driven shock.

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