**Atmospheric
Measurement
Techniques**

Near infrared nadir retrieval of vertical column densities: methodology and application to SCIAMACHY

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Abstract. Nadir observations with the shortwave infrared channels of SCIAMACHY on-board the ENVISAT satellite can be used to derive information on atmospheric gases such as CO, CH_4, N_2O, CO_2, and H_2O. For the operational level 1b-2 processing of SCIAMACHY data, a new retrieval code BIRRA (Beer InfraRed Retrieval Algorithm) has been developed. BIRRA performs a nonlinear or separable least squares fit (with bound constraints optional) of the measured radiance, where molecular concentration vertical profiles are scaled to fit the observed data. Here we present the forward modeling (radiative transfer) and inversion (least squares optimization) fundamentals of the code along with the further processing steps required to generate higher level products such as global distributions and time series. Moreover, various aspects of level 1 (observed spectra) and auxiliary input data relevant for successful retrievals are discussed. BIRRA is currently used for operational analysis of carbon monoxide vertical column densities from SCIAMACHY channel 8 observations, and is being prepared for methane retrievals using channel 6 spectra. A set of representative CO retrievals and first CH_4 results are presented to demonstrate BIRRA's capabilities.

1 Introduction

Nadir sounding of molecular column densities is well established in atmospheric remote sensing. For UV instruments observing the back-scattered sunlight, the analysis is traditionally based on a DOAS (Differential Optical

Absorption Spectroscopy) methodology that considers the logarithm of the measured signal – essentially the optical depth. This approach has also been successfully applied to SCIAMACHY's (Scanning Imaging Absorption Spectrometer for Atmospheric CHartographY) (Bovensmann et al., 1999; Gottwald and Bovensmann, 2011) near infrared (NIR) channels (Buchwitz et al., 2007; Frankenberg et al., 2005b). Alternatively, the measured radiance spectra can be directly analyzed. This method is common practice in thermal infrared atmospheric spectroscopy and constitutes the core of the Iterative Maximum Likelihood Method (IMLM) developed by SRON for analysis of SCIAMACHY's near infrared channels (Gloudemans et al., 2005). Recently, Reuter et al. (2010) also presented an optimum estimation analysis of sun normalized radiances for improved SCIAMACHY CO_2 retrievals and Lerot et al. (2010) applied a direct fitting to nadir viewing UV spectra.

The aim of SCIAMACHY nadir NIR observations is to retrieve information on atmospheric gases such as CO, CH_4, CO_2, N_2O, or H_2O. Ideally, profiles of volume mixing ratio $q_X(z)$ or number density $n_X(z) = q_X(z) \cdot n_{air}(z)$ of a given molecule X can be retrieved, where z represents altitude and n_{air} air number density. However, vertical sounding inversions are generally ill-posed problems, so for weakly absorbing gases it is customary to retrieve only vertical column densities (VCD),

$$N_X \equiv \int_{z_{srf}}^{z_{TOA}} n_X(z)\,dz \,, \tag{1}$$

where z_{srf} is the surface altitude and z_{TOA} the altitude of *top of the atmosphere* (TOA).

In order to have a flexible and robust inversion algorithm for the efficient processing of level-1b to level-2 data, the "Beer InfraRed Retrieval Algorithm" (BIRRA) has been developed at DLR. BIRRA performs a nonlinear least squares

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fit of the observed near infrared (sun-normalized) radiances. First results including a careful intercomparison with the WFM-DOAS algorithm (Weighting Function Modified DOAS, University of Bremen, Buchwitz et al., 2007) for orbit 8663 (27 October 2003) have been presented in Schreier et al. (2007). Recently, the BIRRA prototype has been successfully implemented in version 5 of the level 1b to 2 SCIAMACHY processor for operational data retrievals.

In this paper, the algorithmic basis of BIRRA and its application to carbon gas VCD retrievals are presented. The next section reviews the forward model and the inversion approach used in BIRRA as well as further processing steps. Aspects of input data (nb. level 1 issues) and retrieval settings are discussed in Sect. 3. Section 4 contains a survey of VCD retrieval results. Conclusions and an outlook are given in Sect. 5. The results presented here have been obtained with the scientific prototype version of BIRRA and may differ from the SCIAMACHY operational product. The emphasis of this paper is on methodology, where carbon monoxide and methane retrievals are presented for illustrations. For a discussion of spatial and temporal patterns we refer to, e.g. Buchwitz et al., 2007; Clerbaux et al., 2008; Gloudemans et al., 2006, 2009; de Laat et al., 2006, 2007; Worden et al., 2010; Yurganov et al., 2008. A comprehensive discussion of results for CO and CH$_4$, including validation, will be the subject of a forthcoming paper.

2   Theory and algorithm

2.1   The forward model – near infrared radiative transfer

For an arbitrary slant path, the intensity (radiance) I at wavenumber ν received by an instrument at $s = 0$ is described by the equation of radiative transfer (Liou, 1980; Zdunkowski et al., 2007)

$$I(\nu) = I_{\mathrm{b}}(\nu)\mathcal{T}(\nu) - \int_0^\infty \mathrm{d}s' J(\nu, s') \frac{\partial \mathcal{T}(\nu; s')}{\partial s'}, \qquad (2)$$

where I_{b} is an external contribution, $\mathcal{T}$ represents the transmission through the atmosphere, and J is the source function comprised of thermal emission and scattering. In the near infrared, thermal emission of the atmosphere and Earth's surface is negligible compared to the reflected and scattered sunlight. The contribution of molecular (Rayleigh) scattering is far below 1 % and will be neglected. Aerosol and cloud scattering may have a significant contribution to the intensity but their effects on the retrieval can be mitigated by proxy modeling (see Subsect. 2.3.1, a detailed justification can be found in Frankenberg et al., 2006; Gloudemans et al., 2008). Thus, neglecting scattering, Eq. (2) reduces to Beer's law for a double path through the atmosphere

$$I(\nu) = \frac{r(\nu)}{\pi} \mu_\odot I_{\mathrm{sun}}(\nu) \mathcal{T}_\uparrow(\nu) \mathcal{T}_\downarrow(\nu)$$

$$= \frac{r(\nu)}{\pi} \mu_\odot I_{\mathrm{sun}}(\nu)$$

$$\exp\left[-\int_{\mathrm{earth}}^{\mathrm{sun}} \mathrm{d}s' \sum_m n_m(s') k_m\big(\nu, p(s'), T(s')\big) \right] \cdot$$

$$\exp\left[-\int_{\mathrm{earth}}^{\mathrm{sat}} \mathrm{d}s'' \sum_m n_m(s'') k_m\big(\nu, p(s''), T(s'')\big) \right], \qquad (3)$$

where r is the surface albedo and $\mathcal{T}_\uparrow$ and $\mathcal{T}_\downarrow$ (with $\mathcal{T} = \mathcal{T}_\uparrow \mathcal{T}_\downarrow$) denote transmission between reflection point (e.g. Earth surface at altitude z_{srf}) and observer and between Sun and reflection point, respectively.

Assuming spherical symmetry, the path variable s is uniquely related to the altitude z (in a plane–parallel approximation $s' = z'/\mu$ with $\mu \equiv \cos\theta$ for an observer zenith angle θ, similarly $s'' = z''/\mu_\odot$ for a solar zenith angle (SZA) $\theta_\odot$). k_m and n_m are the (pressure p and temperature T dependent) absorption cross section and number density of molecule m. In infrared line-by-line models, the absorption cross section k_m of molecule m is obtained by summing up the contributions from many lines,

$$k_m(\nu, p, T) = \sum_l S_l^{(m)}(T)\, g\big(\nu; \hat{\nu}_l^{(m)}, \gamma_l^{(m)}(p, T)\big), \qquad (4)$$

where each individual line l is described by the product of the temperature-dependent line strength $S_l^{(m)}$ and a normalized line shape function g describing the broadening mechanism. $\hat{\nu}$ and γ are line position and half width at half maximum (HWHM), respectively. The combined effect of pressure broadening (corresponding to a Lorentzian line shape) and Doppler broadening (Gaussian line shape) is represented by a Voigt line profile (Schreier, 2011). Line mixing (LM) is not yet considered in BIRRA. Tran et al. (2010) discussed consequences of line mixing in the $2\nu_3$ band of methane for molecular spectroscopy and atmospheric retrievals, concluding that "from the point of view of atmospheric retrievals, neglecting LM with suitable effective line parameters is convenient and accurate (within current retrieval uncertainties)."

The measured spectrum is modeled by convolution of the monochromatic intensity spectrum (3) with a normalized instrument spectral response function $\mathcal{S}$

$$\widehat{I}(\nu) \equiv (I \otimes \mathcal{S})(\nu) = \int_{-\infty}^\infty I(\nu') \times \mathcal{S}(\nu - \nu')\, \mathrm{d}\nu',$$

For SCIAMACHY measurements, a Gaussian function is commonly used,

$$\mathcal{S}_{\mathrm{G}}(\nu, \gamma) = \frac{1}{\gamma}\left(\frac{\ln 2}{\pi}\right)^{1/2} \cdot \exp\left[-\ln 2\left(\frac{\nu}{\gamma}\right)^2 \right], \qquad (5)$$

where γ is the half width at half maximum (HWHM). The BIRRA forward model is based on GARLIC (Generic Atmospheric Radiation Line-by-line Infrared Code), the Fortran 90/2003 reimplementation of the Fortran 77 code MIRART (Schreier and Schimpf, 2001). GARLIC/MIRART is

a line-by-line code for arbitrary observation constellations (up-looking, down-looking or limb-viewing in a spherical geometry), and instrumental field-of-view and spectral response. GARLIC/MIRART has been verified in extensive intercomparisons (e.g. von Clarmann et al., 2002; Melsheimer et al., 2005); recently, Hedelt et al. (2011) used the code successfully to model observations of the Venus transit 2004 made by the Tenerife Vacuum Tower Telescope. Furthermore, a multiple scattering enhanced version of MIRART has been used to study effects of cirrus clouds on mid infrared limb MIPAS observations (Mendrok et al., 2007). Molecular spectroscopic parameters from HITRAN, GEISA, or JPL databases can be selected; additionally continuum corrections to the molecular absorption are considered, e.g. the CKD continuum of water (Clough et al., 1989). Jacobians, i.e. derivatives of transmission or radiance spectra with respect to the parameters to be retrieved, are obtained by means of automatic differentiation (Griewank, 2000; Hascoët and Pascual, 2004). For nadir modeling, refraction is only taken into account for the Sun – Earth path element whereas for the "down-looking" path segment (satellite – Earth's surface) with viewing angles $\leq 30°$ refraction is neglected.

2.2 The inverse problem – least squares

Denoting by α_m the molecular scale factor to be estimated and by $n_m^{\mathrm{prior}}(z)$ the a-priori (e.g. climatological) number density of molecule m, the total path transmission in Eq. (3) can be written as

$$\mathcal{T} \equiv \mathcal{T}_\uparrow \mathcal{T}_\downarrow = \exp\left[-\sum_m \tau_m\right] = \exp\left[-\sum_m \alpha_m \tau_m^{\mathrm{prior}}\right], \quad (6)$$

where τ_m is the "true" total optical depth of molecule m (note that for simplicity this equation is written for a non-refractive plane-parallel geometry),

$$\tau_m(\nu) = \int_{z_{\mathrm{srf}}}^{z_{\mathrm{TOA}}} \mathrm{d}z' \left(\frac{1}{\mu} + \frac{1}{|\mu_\odot|}\right) n_m(z') k_m(\nu, z'), \quad (7)$$

and τ_m^{prior} the a-priori total optical depth computed with n^{prior} (Note that for simplicity this equation is written for a plane-parallel non-refractive atmosphere). From Eq. (6), we have that $\tau_m(\nu) = \alpha_m \tau_m^{\mathrm{prior}}$.

If the observation geometry (i.e. z_{srf}, z_{TOA}, μ and $\mu_\odot$) is known, changes in the a-priori molecular cross sections k_m are neglected and it is assumed that the scale factors do not depend on altitude, then, recalling Eq. (1), the "true" VCD of the molecule m can be written as

$$N_m = \alpha_m N_m^{\mathrm{prior}}, \quad (8)$$

where $N_m^{\mathrm{prior}} = \int_{z_{\mathrm{srf}}}^{z_{\mathrm{TOA}}} n_m^{\mathrm{prior}}(z)\mathrm{d}z$ is the a-priori VCD of the molecule m.

In addition to the scaling factors α_m, the surface albedo r (generally modeled by a polynomial in wavenumber) is treated as an unknown; furthermore, instrumental parameters (i.e. the slit function half width) and an optional baseline correction b (again a polynomial) can be considered as fit parameters.

2.2.1 Nonlinear least squares

The standard approach to estimate unknown quantities $x \in \mathbb{R}^n$ from measurements $y \in \mathbb{R}^m$ relies on a (generally nonlinear) least squares fit

$$\min_x \|y - F(x)\|^2 . \quad (9)$$

Here, F denotes the forward model $F : \mathbb{R}^n \to \mathbb{R}^m$ essentially given by the radiative transfer and instrument model,

$$F(x) \equiv \widehat{I}(\nu) = \frac{r(\nu)}{\pi} \mu_\odot I_{\mathrm{sun}}(\nu)$$

$$\times \exp\left[-\sum_m \alpha_m \tau_m(\nu)\right] \otimes \mathcal{S}(\nu, \gamma) + b(\nu) , \quad (10)$$

where the state vector $x \equiv (\alpha, \gamma, r, b)$ is comprised of molecular scale parameters α, and optionally "auxiliary" parameters γ, r, b (the arrays r and b denote the coefficients of the polynomials in wavenumber).

For the solution of the nonlinear least squares problem, Eq. (9), BIRRA uses solvers provided in the PORT Optimization Library (Dennis, Jr. et al., 1981, available at http://www.netlib.org/port/) based on a scaled trust region strategy. Furthermore, BIRRA provides the option to use a nonlinear least squares with simple bounds to avoid unphysical results, e.g. non-negativity,

$$\min_{x>0} \|y - F(x)\|^2 . \quad (11)$$

2.2.2 Separable least squares

Note that the surface albedo r and the baseline correction(s) b enter the forward model F, Eq. (10), linearly and the least squares problems, Eqs. (9) and (11), can be reduced to a separable nonlinear least squares problem (Golub and Pereyra, 2003). Splitting the vector x of parameters to be fitted into a vector η of nonlinear parameters and a vector β of linear parameters, i.e.

$$x \longrightarrow (\eta, \beta), \text{with } x \in \mathbb{R}^n,$$
$$\eta \equiv (\alpha, \gamma) \in \mathbb{R}^p, \beta \equiv (r, b) \in \mathbb{R}^q \text{ and } n = p + q, \quad (12)$$

the forward model can be written as

$$F(x) = \sum_{l=1}^q \beta_l f_l(\eta) , \quad (13)$$

where $f_l : \mathbb{R}^p \to \mathbb{R}^m$ for $l = 1, \ldots, q$. Combining these functions in a matrix

$$\mathsf{A}(\eta) \equiv \left(f_1(\eta), f_2(\eta), \ldots, f_q(\eta)\right), \text{with } \mathsf{A} \in \mathbb{R}^{m \times q}, \quad (14)$$

Equation (9) can be rewritten as a linear least squares problem

$$\min_{\beta} \| \boldsymbol{y} - \mathsf{A}\boldsymbol{\beta} \|^2 \qquad (15)$$

for the vector $\boldsymbol{\beta}$, that is formally solved by

$$\boldsymbol{\beta} = \left(\mathsf{A}^{\mathrm{T}} \mathsf{A} \right)^{-1} \mathsf{A}^{\mathrm{T}} \boldsymbol{y} . \qquad (16)$$

Inserting this solution in Eq. (13), the original least squares problem Eq. (9) becomes

$$\min_{\boldsymbol{\eta}} \left\| \boldsymbol{y} - \sum_{l} \left(\left(\mathsf{A}^{\mathrm{T}} \mathsf{A} \right)^{-1} \mathsf{A}^{\mathrm{T}} \boldsymbol{y} \right)_{l} f_{l}(\boldsymbol{\eta}) \right\|^2 . \qquad (17)$$

This is a nonlinear least squares problem for $\boldsymbol{\eta}$ independent of $\boldsymbol{\beta}$ and can be solved in the usual way by means of Gauss–Newton or Levenberg–Marquardt algorithms. Once the optimum $\boldsymbol{\eta}$ is found, the linear parameter vector $\boldsymbol{\beta}$ is obtained from Eq. (16). The main advantages of this approach are: (i) the nonlinear least squares solver has to iterate only over a reduced fit vector $\boldsymbol{\eta}$; (ii) no initial guess is required for the linear parameters $\boldsymbol{\beta}$; (iii) the size of the Jacobian matrix is reduced.

2.2.3 Jacobians and altitude sensitivity

The Jacobian, $J \equiv \partial \boldsymbol{F} / \partial \boldsymbol{x}$, the partial derivatives of the model function with respect to the fit parameters, is an essential quantity required for Gauss–Newton type iterative solvers of nonlinear least squares problems. Although BIRRA performs a least squares fit of vertical column densities, derivatives with respect to the (discretized) molecular abundance profiles deliver insight into the sensitivity of the retrievals to different altitude regions in addition to useful hints to the appropriate spectral region selection.

According to Fig. 1, in SCIAMACHY channel 8 both branches of the first CO overtone band at 2.4 μm have about the same impact, however, water interference is less significant in the R-branch between 4250 and 4300 cm^{-1}. Furthermore, the plot indicates that the highest altitude sensitivity is located in the troposphere.

For an assessment of altitude sensitivity in the context of profile retrievals, specification of averaging kernels relating the true and estimated state vector (comprising the discretized profile) is customary. In case of column density retrievals as performed by BIRRA, the state vector is composed of the profile scaling factors α_m of all relevant molecules (and some additional auxiliary parameters). In order to estimate the altitude sensitivity, an approach suggested by Buchwitz et al. (2004) has been used, i.e. a series of BIRRA fits has been performed using synthetic spectra generated with perturbed CO profiles. Denoting $N^{\mathrm{r}}(z)$ the VCD retrieved from the spectrum generated with a profile perturbed at an altitude
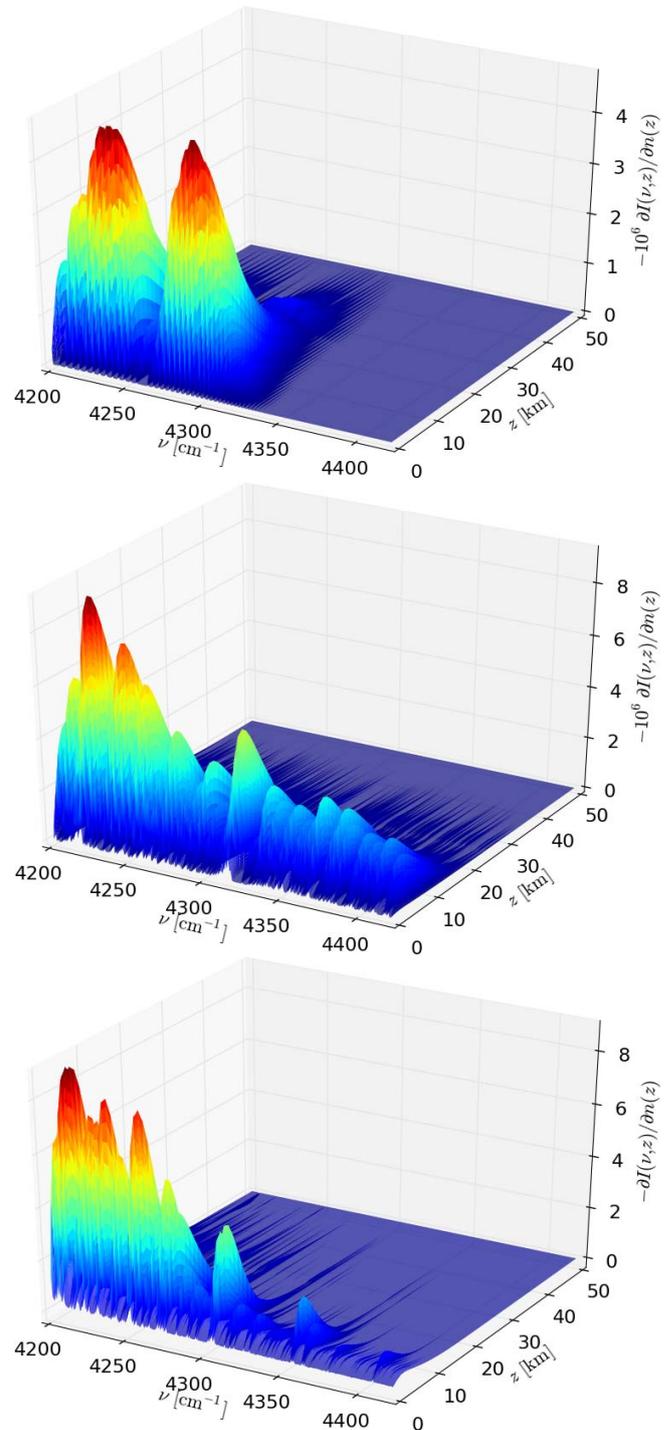

Fig. 1. Jacobians [erg/s/(cm^2 sr cm^{-1})/ppm] for molecular concentration profile retrieval in channel 8: CO (top), CH$_4$ (mid), and H$_2$O (bottom). Note the scaling by 10^6 of the CO and CH$_4$ Jacobians. The derivatives were calculated using GARLIC for a US standard atmosphere up to 50 km, a Gaussian slit function with $\gamma = 0.2$ cm^{-1}, and vertical downlooking observer at 800 km.

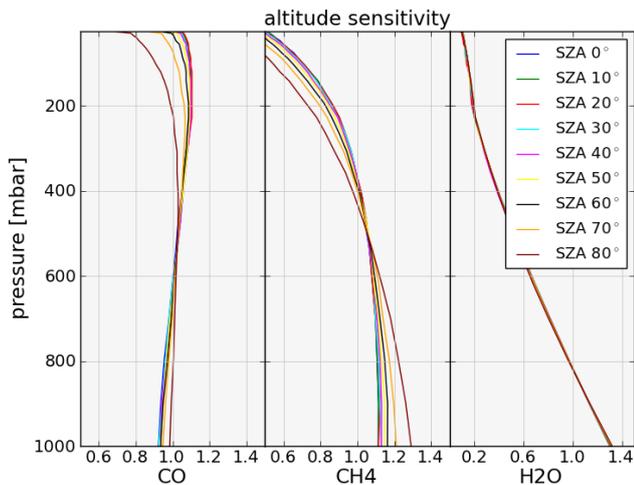**Fig. 2.** Altitude sensitivity of CO (left), CH_4 (mid), and H_2O in channel 8.

z, and $N^p(z)$ VCD corresponding to this perturbed profile, the altitude sensitivity is estimated by

$$A(z) = \frac{N^r(z) - N}{N^p(z) - N} ,$$

(18)

where N is the true VCD as in Eq. (1) (the subscript X is omitted for brevity). Fig. 2 confirms the high sensitivity of NIR retrievals in the lower troposphere; furthermore, as pointed out by Buchwitz et al. (2004) and Gloudemans et al. (2008), this sensitivity is depending on the solar zenith angle and, to lesser extent, on the observation angle.

2.3 Postprocessing

In the previous subsections, the basics of forward modeling (radiative transfer) and inversion (least squares optimization) as applied to a single observation have been presented. Here, we discuss further steps necessary to proceed from the fitted parameters during the inversion process to "higher level products" such as local and global spatial distributions and temporal evolutions.

2.3.1 Product definition

Since scattering is neglected in the BIRRA forward model, the photon path is considered to be the optical path that solar light travels from the top of the atmosphere to the Earth's surface, and reflected from the surface up to the observer. However, the measured radiance has a high probability of having also a (small or large, depending on the individual conditions) fraction coming from scattering events in the atmosphere, and consequently, having a photon path different from the pure geometrical one. In addition, the atmospheric conditions are set a-priori to climatological datasets but the actual meteorological conditions (e.g. pressure, temperature)

at the time of the observations are unknown. Moreover, the quality of the observed spectra has a significant impact on the fitted scaling factors of the atmospheric constituents and, therefore, calibration open issues and the known growth of an ice layer on channel 8 detector (Gloudemans et al., 2005) will affect the retrievals.

In order to account for these unconsidered effects, the results of the retrieved gases are presented as "proxy"-normalized vertical column densities instead of directly retrieved VCDs. As proxy, it is typical to select a fit parameter that contains information about the unaccounted issues. In case of CO retrievals from SCIAMACHY channel 8, methane is a reasonably good candidate to be used as proxy (Buchwitz et al., 2000). CH_4 has strong absorption lines across the CO spectral fitting window, so the amount of CH_4 can be determined with high accuracy. Furthermore, CH_4 is a well-mixed gas with long life time in the atmosphere and, consequently, it has quite homogeneous global distributions. The deviations of CH_4 concentrations from the a-priori due to sources and sinks are rarely bigger than 10 %, far smaller than the deviations of CO that can easily depart by hundreds percent from the reference value over emitting areas. Thus, in terms of CO variability, CH_4 can be considered as constant and variations in the retrieved CH_4 from the expected a-priori value can be interpreted as the effect of the unconsidered processes. So, the proxy-normalized CO vertical column density is defined as

$$xCO \equiv N_{CO}^{prior} \times \frac{\alpha_{CO}}{\alpha_{CH4}},$$

(19)

where N_{CO} represents the a-priori CO vertical column density, and α_{CO} and α_{CH4} are the fitted scaling factors of CO and CH_4, respectively (see Eq. 8). Note that both scaling factors are retrieved from channel 8 data.

Considering that the retrieval of CH_4 under the considerations of homogeneity exposed above is equivalent to the retrieval of the dry air mass, the ratioed quantity xCO is called *dry-air column density* (Wallace and Livingston, 1990; Yang et al., 2002). For channel 8 retrievals, however, the ice layer that grows over the detector has a different impact on CH_4 than on CO retrievals. Thus, the CH_4-proxy does not entirely account for the effect of the ice layer and special care has to be taken when correcting for this.

Analogously, using SCIAMACHY channel 6, one can define the *dry-air column mixing ratio* of CH_4 as

$$xCH4 \equiv q_{CH4}^{prior} \times \frac{\alpha_{CH4}}{\alpha_{CO2}},$$

(20)

where q_{CH4}^{prior} is the a-priori column mixing ratio of CH_4, and α_{CH4} and α_{CO2} are the scaling factors of CH_4 and CO_2, respectively. For the target species CH_4, the chosen dry-air proxy is CO_2 (Frankenberg et al., 2006). All conditions mentioned above for the CH_4-proxy also hold for the proxy CO_2. Further, CO_2 is far more homogeneous (vertical and horizontally) than CH_4, and the absorption signatures of target

and proxy species are of comparable magnitude in channel 6, which is desirable but not the case for the CO target with a CH_4 proxy.

2.3.2 Quality criteria

In order to use the fitted data, a variety of quality criteria have to be fulfilled. Convergence of the least squares fit is requested and only moderate solar zenith angles (typically < 80 dg) are selected. Retrievals with high fit errors are rejected (typically $\varepsilon(\alpha_{CO}) < 0.5$, $\varepsilon(\alpha_{CH4}) < 0.01$). Furthermore, for the acceptance or rejection of retrieved data a quantile analysis is performed, i.e. outliers far off the median are rejected. As suggested by Gloudemans et al. (2009), different cloud filters are used over land and over sea: in contrast to land observations, the NIR albedo of oceans is very low and, as a consequence, the signal-to-noise ratio of cloud free observations is very low. Thus, over oceans the presence of clouds with high albedo significantly enlarges the signal and enables reliable retrievals. Accordingly, only pixels with cloud fraction higher than 20 % are accepted. Over land, however, the presence of clouds is a source of uncertainties, since the radiative transfer model only accounts for one photon path and, in case of partial cloud cover, the observed radiances have contributions coming from reflections at the Earth's surface and from scattering at the cloud layer. Therefore, only observations with cloud fraction below 20 % are accepted.

Observations over ocean require special care. Cloud top height (CTH) is a crucial parameter for trace gas retrieval under cloudy conditions. This is specially true for retrievals over the ocean, since clouds are unavoidable for reliable retrievals. In these cases, the main contribution to the measured intensity comes from the cloud top region, so this information helps to understand and improve the retrievals. Since CH_4 is a well-mixed gas, variations in CH_4 VCDs can be related to "obstacles" along the photon path, mainly due to cloud shielding. Bounds on the CH_4 VCD are highly correlated with restrictions on cloud top height. This fact is exploited by Buchwitz et al. (2004, 2007). Gloudemans et al. (2009) used cloud information based on CH_4 retrieved from SCIAMACHY channel-8 for ocean CO retrievals.

High clouds introduce uncertainties and systematic errors to the CO retrievals especially by the scaling with an incorrectly retrieved CH_4 partial column. However, since retrievals in presence of high clouds translate, in most cases, to large CH_4 retrieval errors, these can be used for masking out high cloud observations. Indeed, comparisons of xCO using the quality criteria previously described and using an extra condition on cloud top height (namely, CTH < 2 km) show similar results, since most of the high cloud observations are already removed by the condition $\varepsilon(\alpha_{CH4}) < 0.01$.

For CO retrievals, the normalization by the CH_4 scaling factor accounts for deviations from geometrical photon paths. However, two systematic errors arise from this treatment. On the one hand, the ratio of the partial CO and CH_4 VCDs depends on the bottom boundary (surface or cloud top). Correcting for this, xCO retrieved only over clouds can be transferred to total vertical column densities, assuming that the retrieved-to-reference VCD ratio above and below the clouds is the same. On the other hand, CO and CH_4 have different altitude sensitivity, so deviations from the reference profiles at different altitudes will have a different impact on the retrieved CO and CH_4. This effect has not been considered. Nonetheless, below 500 mb the altitude sensitivity of both gases is similar and hence, the error introduced is far smaller than the other contributions (given that high-cloud observations have been filtered out).

3 Retrieval setup and input data: sensitivity studies

The quality of the input data greatly affects the accuracy of the retrievals. Since model parameters are optimally varied during the inversion process to mimic the measured values, errors in the input spectra will lead to wrong retrievals. In this section, sensitivity studies with respect to input data and retrieval settings are presented for carbon monoxide.

Level 1 data (spectra and geolocation) are taken from the SCIAMACHY level 1 product (SGP version 6.03). Unless otherwise noted, the dead and bad pixel mask (DBPM) used corresponds to the prototype of the new level 1 SGP version 7.03. For topographic information (surface elevation) (ETOPO4) data are used. Auxiliary data such as cloud information are taken from the SCIAMACHY level 2 product SGP version 3.01.

3.1 Level 1 spectra – trace gases fitting windows

SCIAMACHY spectra are spectrally and radiometrically calibrated and corrected for several effects, namely: leakage current, pixel-to-pixel gain, non-linear response, stray light, and polarization. Reflectances are calculated using in-flight sun diffuser spectra. Additionally, the degradation of the instrument is monitored and the quality of the individual spectral pixels is assessed. For the CO retrievals presented here, the spectral window in the middle of channel 8, ranging 4282–$4303\,cm^{-1}$ (equivalent to 2323–2336 nm) is used. For methane retrievals, two microwindows in channel 6 are selected: the 5986–$6139\,cm^{-1}$ (1629–1671 nm) interval with CH_4 as the strongest absorber, and the 6273–$6419\,cm^{-1}$ (1558–1594 nm) interval with CO_2 as the strongest absorber.

3.2 Dead & bad pixel mask (DBPM)

The DBPM contains pixels that are deemed unusable for the retrieval due to damage mostly caused by enhanced particle flux during the passing of the Southern Atlantic Anomaly (SAA). Consequences of the damage are, e.g. high noise values, disconnection or very high leakage currents. Dark, sun, and internal lamp measurements are used to determine bad
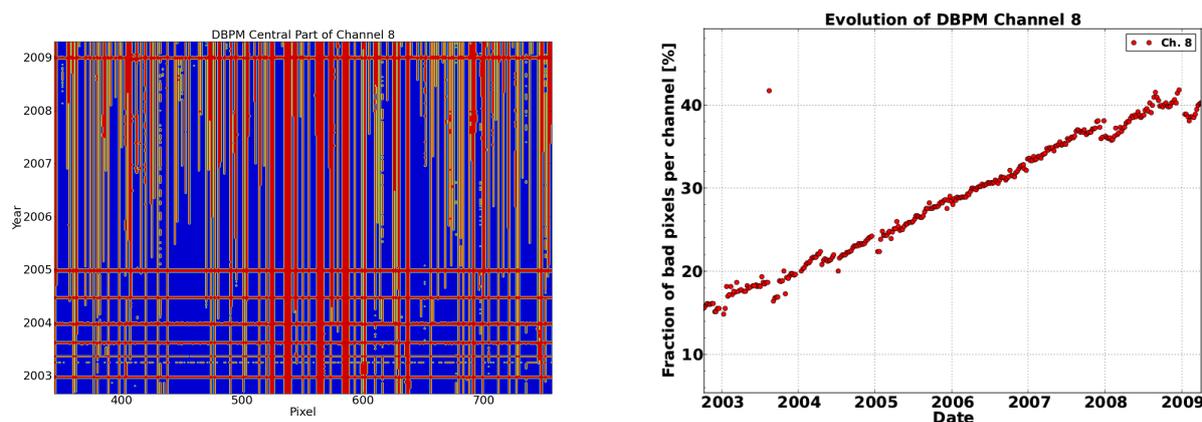
Fig. 3. Evolution of pixel mask from 2002 to 2009 (left) and percentage of flagged pixels (right). Good pixels are marked blue and bad pixels red. Several decontaminations rendering the detectors temporarily useless due to high temperatures and resulting noise are visible as horizontal red lines in the left diagram.

pixels (see, e.g. Lichtenberg et al., 2006), and parameters such as mean noise or error of dark parameters are calculated. The DBPM judges the quality of a pixel by setting channel wide thresholds for dark signal, dark signal noise, dark signal saturation, residual of dark correction, and sun and white light source signal. If one or more thresholds is/are violated in more than 40 % of the cases within one week, the pixel is masked as bad. The algorithm is based on an approach described by SRON (personal communication, October 2006). Furthermore, a pixel is always either bad or good, there are no intermediate values.

Since mission start, the number of dead/bad pixels has grown steadily, cf. Fig. 3a, and in June 2009 around 40 % of the pixels in channel 8 are marked as bad. Fig. 4 shows the effect on CO VCD retrievals of using a constant mask for one year vs. using a dynamic mask continuously updated Lichtenberg et al. (2010). Even though the total number of flagged pixels does not change dramatically between February and October 2004, cf. Fig. 3b, the retrieved VCDs differ significantly, i.e. CO fits depend on the presence of the individual pixels and their quality. This result was already found earlier by Gloudemans et al. (2005) for CO retrievals with the IMLM algorithm.

3.3 Sensitivity to signal changes in individual pixels

An inversion process aims at gaining information about model parameters from observed quantities. In the case of atmospheric gas retrievals, the observed quantities are the radiances measured at different wavelengths for a light beam that has traveled through the Earth's atmosphere. Since the absorption of light by atmospheric gas constituents is wavelength-dependent, the discrete measured radiance spectrum contains implicit information of gas concentrations: the higher the concentration of molecule m, the lower the radiance at wavelengths where gas m absorbs (see Sect. 2.1 for details).

In this subsection, the spectral sensitivity of the molecular scaling factors α_m is studied by sequential perturbations on individual pixels of SCIAMACHY channel 8. On the one hand, a response of a molecular scaling factor α_m to a perturbation on a given spectral pixel means that this pixel contains information about the atmospheric content of molecule m and that the inclusion of this pixel would be beneficial for the inversion. On the other hand, some radiative and spectroscopic aspects such as the interference of spectral lines of the target gas with strong lines of other gases, or insufficient knowledge of molecular absorption cross sections due to imprecise spectral line parameters can lead to errors during the inversion process. Consequently, those critical pixels where perturbations have large impacts on several molecules at the same time should be treated with care.

A synthetic SCIAMACHY spectrum covering the whole channel (1004 pixels ignoring the blinded 20 pixels at the left and right ends) with a representative viewing geometry was produced by means of GARLIC (the BIRRA forward model). An inversion of this unperturbed noise-free "reference" spectrum delivers scaling factors of unity for all gases. Gradually, the intensity spectrum was perturbed pixel by pixel by constant amounts:

$$I_{ij}^{\text{pert}} = (1 + a_j) \times I_i^{\text{ref}} \text{ with } a_j = 0.1 \times j$$
$$\text{for } i = 1, \ldots, 1004 \text{ and } j = -10, \ldots, 10 \tag{21}$$

where the index i denotes the spectral pixel number and j the perturbation. For instance, $I_{i,-10}$ represents a perturbation on the ith pixel of -100 % (i.e. $I_{i,-10} = 0$), $I_{i,0}$ is the unperturbed reference radiance I_i^{ref}, and $I_{i,10}$ a perturbation of 100 % (i.e. $2I_i^{\text{ref}}$).

Figure 5 shows the results of this study around the CO fitting window (4282–4303 cm^{-1}) in channel 8. The absorption cross sections of CO, CH$_4$, and H$_2$O are depicted in Fig. 5 (top left panel) for reference. Note that eight CO lines are shown, whereas only the central six are actually included
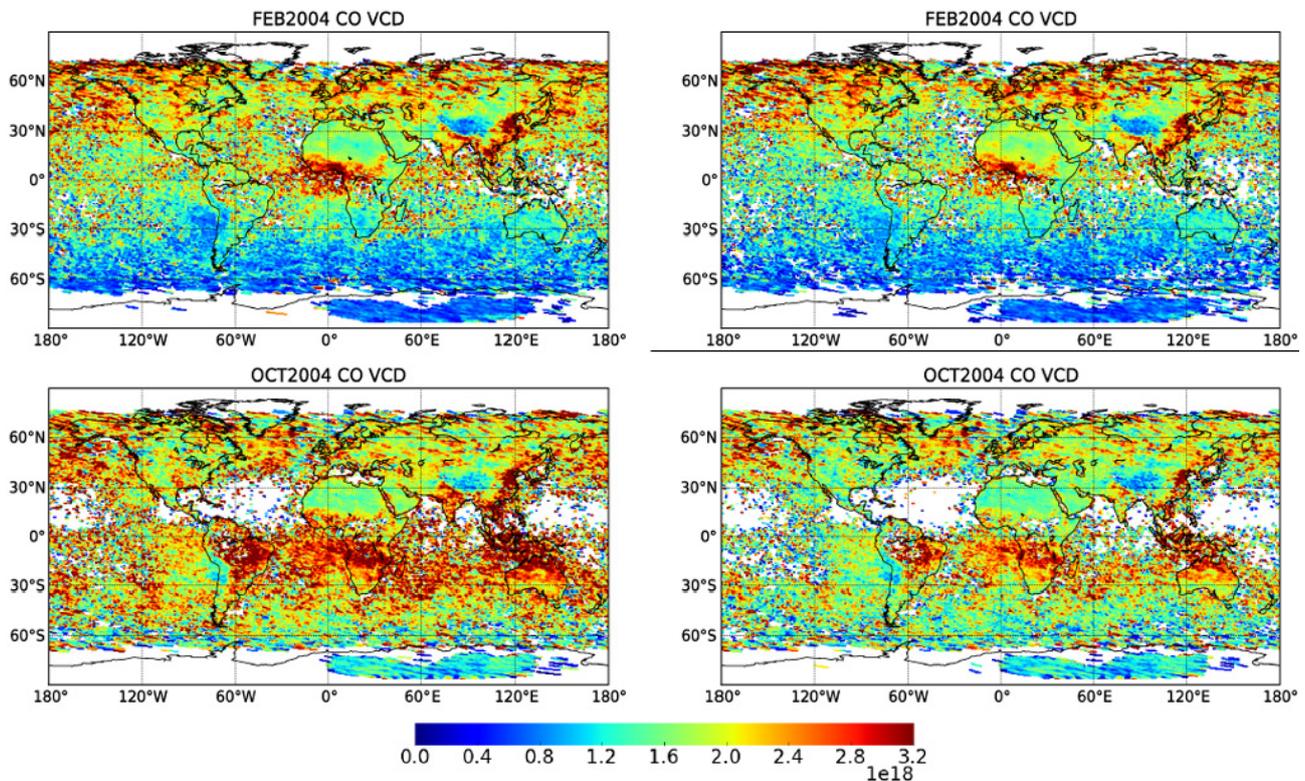
Fig. 4. Comparison of a retrieval using a constant mask flagging only pixels that are marked as "bad" for at least half of the cases of the year 2004 (left) and a retrieval using a dynamic mask appropriate for each measurement (right). In February (top) the results look similar while the result for October (bottom) is noisier for the constant mask that does not flag all bad pixels. CO VCD in units molec cm^{-2}.

in the aforementioned fitting window. The further panels show the impact of the perturbations on the fitted molecular scaling factors α_{CO}, α_{CH4} and α_{H2O}. The patterns of the effect of pixel perturbations on the different gas retrievals look completely different, reflecting the structure of the cross sections. Perturbations of pixels close to the center of strong absorption lines have mostly a large impact on the retrieved columns, whereas a perturbation on pixels far away in the wings does not alter significantly the retrievals. In absolute values, the effect is very different for CO (up to a factor 2) than for the other two gases (few percent). Both, water vapor and methane absorb substantially over the entire channel, so the perturbation of a single pixel is not critical. On the other hand, carbon monoxide has much weaker absorption lines and only in part of the channel, and consequently is more sensitive to the quality of the spectra. Indeed, some perturbations even resulted in negative CO scaling factors. Recall that for this sensitivity study, the measurement vector **y** comprises 1004 elements. In case of the CO retrieval window, the size of the measurement vector is considerably reduced and the effect of radiance perturbations on the α's increases significantly (orders of magnitude for CO), since individual pixels gain in relative weight. Note that for operational retrievals the fitting window cannot cover the whole channel due to the timeliness requirements on data availability.

According to Fig. 5, the effect on H_2O retrievals seems to be small. However, water vapor is highly variable spatially and temporally, and the water profile assumed in the model has a strong impact on the retrieval result, cf. Fig. 6. Furthermore, molecular spectroscopy of water is quite delicate and, according to Rothman et al. (2009), "the recommended line list for water remains in a state of constant evolution." In laboratory spectroscopy, an accurate determination of the amount of water in the absorption cell is difficult, thus any error in the number density n leads to a systematic error (over- or underestimate) of line strengths. For all 92 water lines in the 4280–4305 cm^{-1} wavenumber interval, HITRAN 2008 gives an uncertainty range between 5 and 10 % for line strengths. Note that optical depth τ and transmission $\mathcal{T}$ depend on the product of line strength and number density, Sn (see Eqs. 3 and 4); moreover, in the lower atmosphere the line center value of the molecular absorption is proportional to Sn/γ_L where γ_L is the Lorentz width. In conclusion, the uncertainty of both, water density profile and water line parameters suggests the omission of further pixels sensitive to water, i.e. those near strong H_2O absorption lines possibly not modeled sufficiently well.
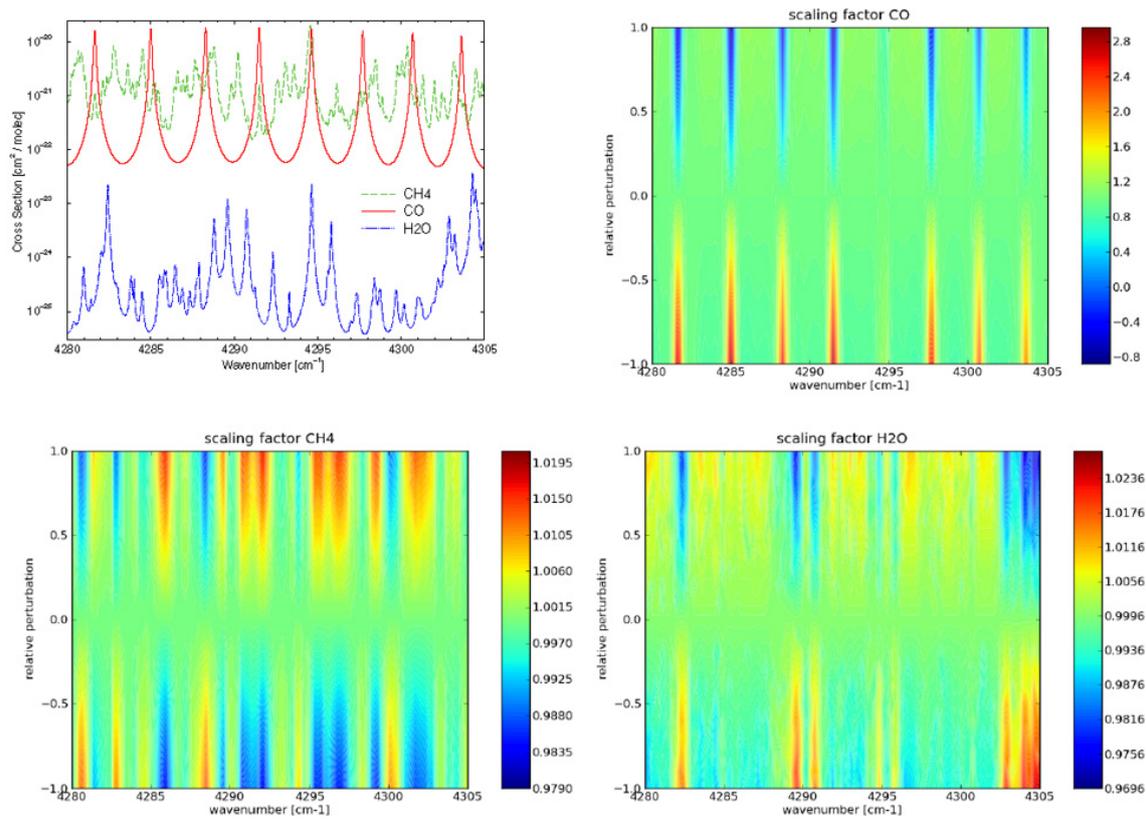
Fig. 5. Sensitivity of BIRRA CO VCDs with respect to perturbations of individual pixels. Top left: absorption cross sections of CO (red), CH$_4$ (green dashed), and H$_2$O (blue long dashed) in channel 8 CO fitting window. Other plots: molecular scaling factors as a function of individual pixel perturbations. Note the different range of the color bars.
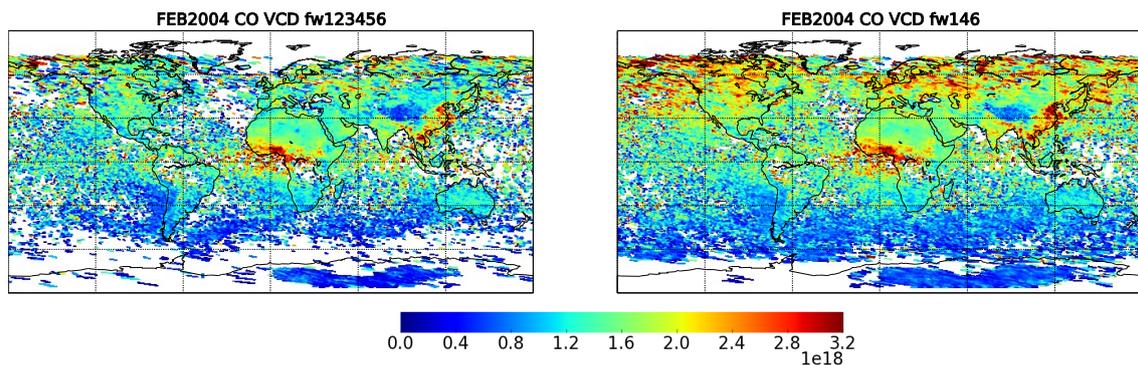

Fig. 6. Comparison of two VCD [molec cm^{-2}] retrievals for February 2004. Left: all good pixels in the retrieval window are used. Right: pixels over lines with strong water vapour interference are excluded. In the latter case, several features like enhanced CO values in South-East Asia and the North-South gradient are more clearly visible.

3.4 Ice layer and instrument transmission

The NIR detectors of SCIAMACHY are the coldest point of the instrument. Since not all water was removed from ENVISAT during the commissioning phase, an ice layer is deposited on the detector surface (this layer is regularly removed by heating the detectors). The ice reduces the trans-

mission in a wavelength dependent way; furthermore it scatters the incoming light and generally leads to a broadening of the spectrum (Lichtenberg et al., 2006; Gloudemans et al., 2005). To account for this effect in carbon monoxide retrievals from channel 8 spectra, BIRRA treats the slit function HWHM γ as an additional auxiliary fit parameter.
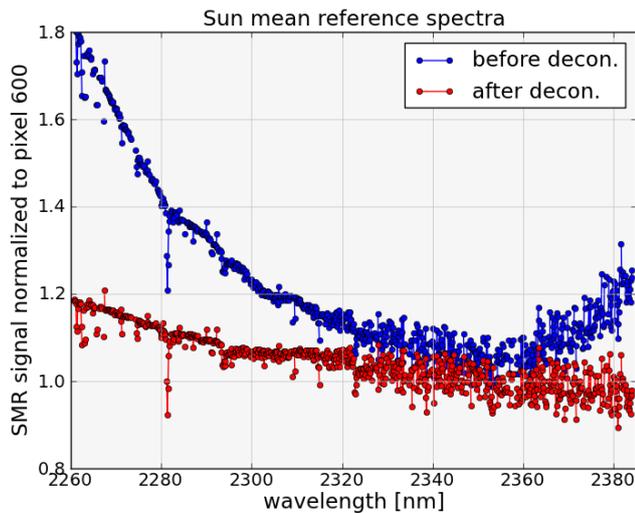
Fig. 7. Sun mean reference (SMR) spectrum in channel 8 with ice layer (blue) and with clean detector (red), normalised to the signal in an arbitrary pixel to illustrate the change in the spectral shape.

Figure 7 shows a typical sun mean reference (SMR) spectrum shortly after decontamination, i.e. with only a thin or no ice layer and a SMR spectrum several weeks later. The change of the spectral shape due to the ice extinction is clearly visible. See also Subsect. 4.1 for a discussion about the impact of the ice layer on mean instrument transmittance or throughput.

3.5 Solar spectrum

The solar irradiance spectrum used in the radiative transfer model, Eq. (10), also has some impact on the retrieval product. BIRRA can use the SMR spectrum measured daily (as an average of a series of measurements) by SCIAMACHY for calibration and scientific purposes. Alternatively, a variety of solar spectra models and measurements is available (e.g. Abrams et al., 1996; Hase et al., 2006), and BIRRA can read solar irradiance spectra of Kurucz (1995) (extracted from MODTRAN4, Berk et al., 1999).

Figure 8 illustrates relative differences of some selected fit parameters when using, on the one hand, the Kurucz model solar spectrum and the SCIAMACHY SMR spectrum, on the other hand. The first row shows relative differences for the CO scaling factor α_{CO}, the second row for the CH$_4$ scaling factor α_{CH4} and the third row for xCO column density. The left column shows the results of the month February 2004 whereas the second one presents those of July 2004. Reddish color (indicating higher values of the fit values for the Kurucz spectrum compared to the SMR spectrum) dominates in all plots of Fig. 8 with no exception, i.e. the use of the Kurucz spectrum biases α_{CO}, α_{CH4} and xCO to higher values. Note, however, that the color bar of α_{CH4} ranges from −0.1 to 0.1 (±10 %), whereas that of α_{CO} and xCO ranges from

−0.5 to 0.5 (±50 %). Thus, the impact of the solar spectrum is much higher on CO than on CH$_4$. The difference distributions of α_{CH4} are quite homogeneous, with the exception of the Tibetan Plateau and the Andes Cordillera, where the differences increase. The distributions of α_{CO} and xCO are very similar indicating that the effect of the solar spectrum on xCO is caused basically by the effect of the solar spectrum on α_{CO} rather than on α_{CH4}. The CO difference distributions show latitude (solar zenith angle) as well as seasonal dependency (the pattern in February and July differ considerably).

Since SMR and Earth's spectra are measured with the same detector, most artificial features are included in both spectra. This is an advantage of SCIAMACHY's SMR spectrum with respect to Kurucz or any other solar spectrum, since in the latter cases, artifacts in Earth's spectra would be attributed by the model to atmospheric effects. As a consequence SCIAMACHY's SMR spectra are used in the retrievals.

3.6 Spectral calibration

In order to ensure high spectral stability over the lifetime of the mission, SCIAMACHY has been equipped with a spectral calibration lamp, the "Spectral Line Source" (SLS), for determining the pixel-to-wavelength relationship. Whereas the SLS has proved to be suitable for a precise in-flight spectral calibration of channels 1 to 6, it is not sufficient for the calibration of channel (7 and) 8 due to the lack of enough SLS spectral lines within these channels. Because of this, the pixel-to-wavelength relationship of channel 8 in the level-1b product is set to the on-ground calibration. Although on-ground calibration was performed under representative flight conditions (temperature and vacuum), a similar spectral tuning as in the other channels should be applied to channel 8 for a precise spectral calibration.

Information from molecular spectroscopy (as provided by the HITRAN or GEISA database) can be exploited for spectral calibration, i.e. an in-flight spectral calibration for the SCIAMACHY channel 8 can be performed utilizing absorption signatures of atmospheric methane, water vapor and carbon monoxide. The spectral correction found has roughly a second-degree polynomial dependency and its value can be as large as 0.5 nm (more than 4 pixels) at the right edge of channel 8, see Fig. 9. Within the CO fitting window, the spectral correction is well approximated by a first-degree polynomial, i.e. shift and squeeze.

As shown in Fig. 10, the neglect of an appropriate wavelength calibration leads to negatively biased VCD distributions (note the different color bars) and unrealistic regional patterns (e.g. Himalayas).
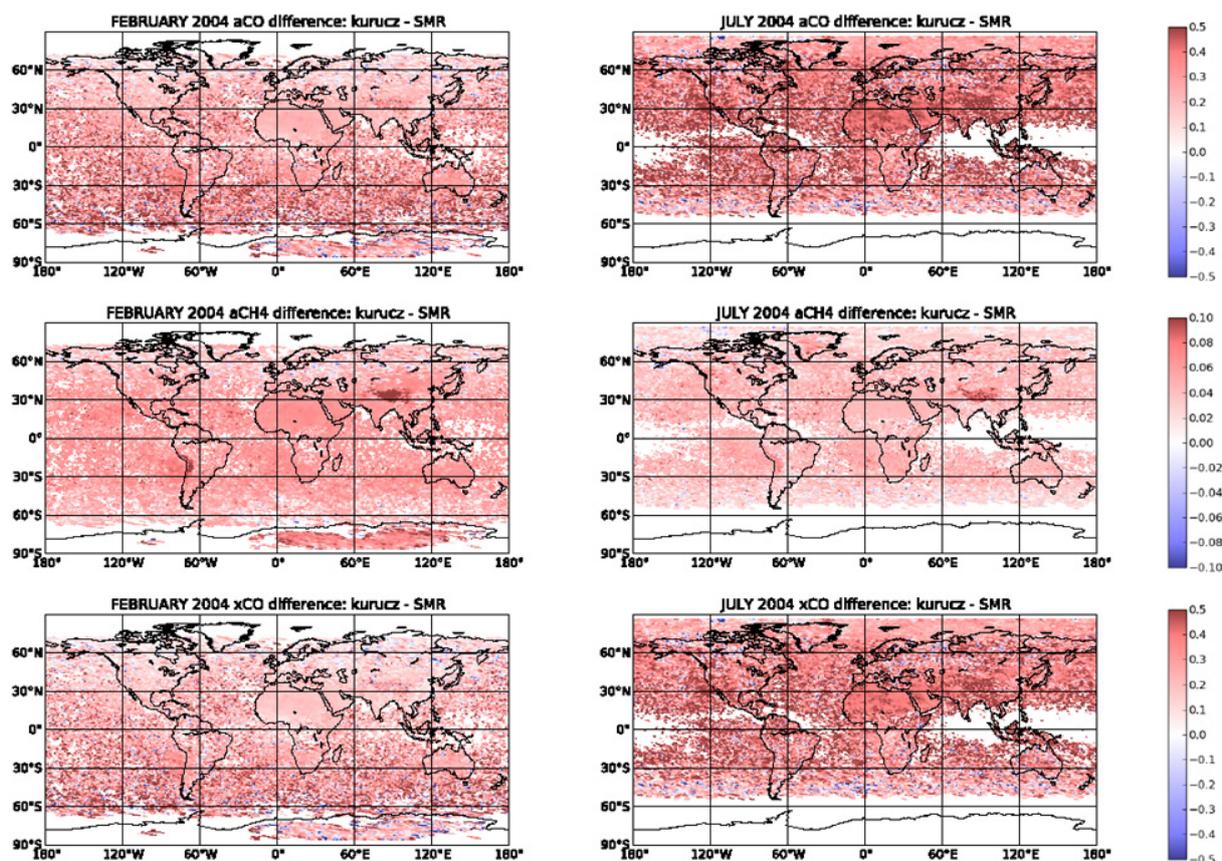
Fig. 8. Influence of the solar spectrum on monthly average CO retrievals. The plots show relative differences of α_{CO} (top), α_{CH4} (mid), and xCO (bottom) of retrievals with the Kurucz solar spectrum vs. the SCIAMACHY SMR spectrum for two months in 2004: February (left) and July (right).
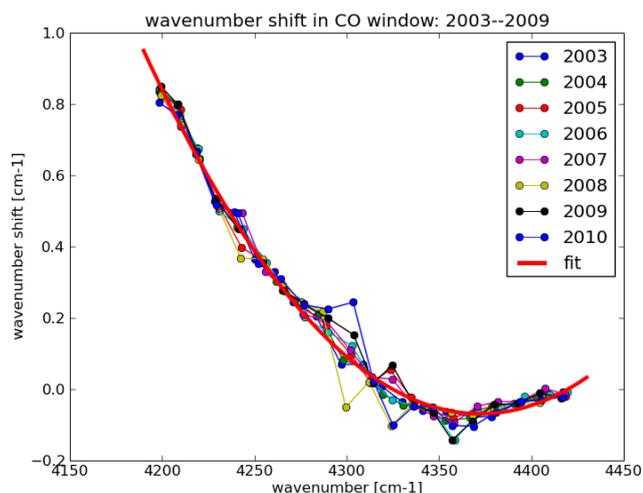

Fig. 9. Spectral correction to the on-ground pixel-to-wavelength relationship.

3.7 Spectroscopic input data: line parameters and continuum

Spectroscopic line parameter databases such as HITRAN (Rothman et al., 2009) or GEISA (Jacquinet-Husson et al., 2008) are an essential input for the computation of molecular cross sections (Eq. 4). Recently, Frankenberg et al. (2008a,b) have discussed the importance of accurate and complete line parameters for methane retrievals in SCIAMACHY channel 6. In the short–wave end of SCIAMACHY channel 8, water vapor spectroscopic data are significantly different in the recent version of these databases. Feng and Zhao (2009) discusses impacts of changes of the HITRAN database on near infrared transmittances, and in SCIAMACHY's channel 8 found the most noticeable changes for wavenumbers larger than $4350\,\mathrm{cm}^{-1}$. In the middle of this channel used for CO retrievals, there are only minor changes, and the retrievals do not show a strong impact.

The default setting for the retrieval of CO column densities from SCIAMACHY channel 8 considers three absorbers, i.e. CO, CH$_4$, and H$_2$O. Although there are no CO$_2$ lines in the center of channel 8, there is a contribution of a small, yet
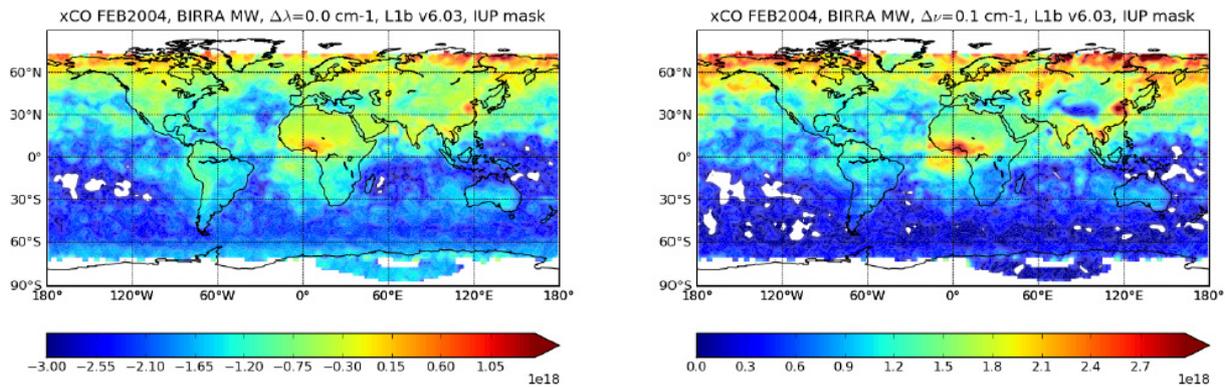
Fig. 10. Comparison of xCO retrievals without (left) and with wavenumber calibration. Note the different color bar with xCO values ranging down to $-3 \cdot 10^{18}$ molec cm^{-2} for the unshifted wavenumber retrievals (left). (DBPM provided by M. Buchwitz, personal communication, October 2006).

non-vanishing CO_2 continuum. However, adding CO_2 as an additional absorber has a negligible effect on the retrieved xCO product.

3.8 Atmospheric input data

VCDs are retrieved in BIRRA by fitting the scaling factors α_m of concentration profiles $n_m(z)$ (see Eqs. 1, 8). Accordingly, the quality of the retrieval clearly depends on the adequacy of the profiles used. Furthermore, pressure and temperature data are required to evaluate the molecular cross sections $k(\nu, p, T)$ in Eq. (4). Test retrievals for orbit 8663 using the six AFGL model atmospheres (Anderson et al., 1986) reveal that the fitted CO column densities are especially sensitive to temperature.

Clearly, the use of just a few atmospheric profiles (as provided by the AFGL data) does not really cover the full seasonal and spatial variability. There are different strategies to better represent the atmospheric state in the retrievals: Buchwitz et al. (2004) use a single profile of temperature, pressure, and trace gas mixing ratios from the US Standard atmosphere (with methane scaled to 1750 ppbv), and fits an additional temperature shift parameter to account for the temperature dependence of the molecular absorption cross sections. In Gloudemans et al. (2008), temperature and H_2O profiles are taken from the European Centre for Medium-Range Weather Forecasts (ECMWF).

Because molecular cross sections are computationally expensive and operational processing imposes timeliness constraints, BIRRA uses a compromise and takes a single pressure and temperature profile for each state selected from the CIRA-86 data base according to time and mean latitude. The Committee on Space Research (COSPAR) International Reference Atmosphere (CIRA) provides monthly mean profiles of pressure vs. temperature for the altitude range 0–120 km with almost global coverage (80° N–80° S) (Fleming et al.,

1990, http://badc.nerc.ac.uk/data/cira/). Trace gas profiles are taken from the US Standard Atmosphere.

3.9 Least squares settings

Solvers for nonlinear least squares problems usually offer several input parameters to control the iterations. In addition to termination code due to excessively large number of iterations, the PORT library delivers convergence codes for standard x–tolerance (relative change of the norm of the x state vector) and y–tolerance (relative change of the norm of the residual vector $y - F(x)$) (Gay, 1990). In our applications, "relative function convergence" was reached for most of the cases.

Ideally, the different least squares solvers provided by PORT should give identical retrieval results. Figure 11a shows that the retrieved carbon monoxide VCDs (averaged within a 1 dg latitude belt over all longitudes) of orbit 8663 (27 October 2003) are very similar for all methods. Differences can be expected for "difficult" observations, i.e. when the fit did not converge properly or some of the fitted parameters are exceptionally small or large. Especially in case where the iterative optimization algorithm suggests a step leading to one or several negative fit variables, differences show up in the constrained and unconstrained retrievals. In fact, the spikes around 20° N and 32° S do not show up in both constrained least squares results. Various strategies have been discussed, how an algorithm should proceed when some of the physical variables reach "forbidden" (usually negative) values; common approaches are to stop the iterative solution process, or to ignore and continue (and accounted for in the post-processing of the data). In case of unconstrained least squares fits, BIRRA stops the iteration when it encounters negative slit function widths or in case of too many negative scaling factors.

Separable least squares is slightly faster than the nonlinear least squares. The main advantage is the numerical

robustness and reliability of the separable least squares: using synthetic "observed spectra" with constant reflectivity ($r_0 > 0$, $r_1 = r_2 = 0$) and fitting with a second order reflectivity polynomial, the separable least squares delivers the "true" state vector (molecular scaling factors and auxiliary parameters) correctly in all cases, whereas the nonlinear least squares fails for about half of the spectra.

4 Carbon monoxide and methane vertical column densities – a survey

Despite the challenges of nadir NIR retrievals, SCIAMACHY observations can be used to estimate time averaged distributions of carbon gas vertical column densities. In the following subsections, a survey of BIRRA carbon monoxide retrieval results is given, followed by a brief presentation of first BIRRA methane retrievals.

4.1 Monitoring of fit parameters

In the previous subsections, the impact of the level 1b data quality on carbon monoxide retrievals has been discussed. Fig. 12 and 13 illustrate time series of the fit parameters for CO retrievals in channel 8. The panels of Fig. 12 show, from top to bottom, the mean channel transmission (throughput), the scaling factors of CO, CH_4 and H_2O, the HWHM of the instrument slit function, and the zeroth, first, and second degree coefficients of the albedo polynomial. Although the throughput is not a fit parameter, it has been included here for comparison. Because of the difficulty in modeling the ice layer on top of channel 8, it has an impact on the fitted parameters.

Since the ice layer reduces the mean intensity, the throughput (the mean instrument transmission) is a good indicator of its thickness. Table 1 shows the cross-correlation ρ of the different fit parameters with respect to the throughput. Note that these coefficients are calculated for the 14-days averaged time series, not for individual observations. They are normalized between -1 (full anti-correlation) and 1 (full correlation). A high correlation coefficient (in absolute value) means that the two curves follow a similar course and it will be taken here as evidence for a possible causal relationship. Table 1 illustrates that (the scaling factor of) CH_4 is the most affected parameter by the ice layer growth. The second degree polynomial coefficient r_2 has the second highest correlation, suggesting that the normalization of the observed spectra by the SMR spectrum does not completely eliminate the change in spectral shape due to scattering in the ice layer. These conclusions are backed by Fig. 12. According to Tab. 1, the first degree polynomial parameter (r_1), the HWHM and CO are less affected by the ice layer. A detailed examination of the curves in Fig. 12, however, shows that r_1 does have a strong anti-correlation with the throughput, but two outliers lower the value of ρ. Remarkably, the r_0

coefficient shows little correlation with the throughput. Using an external (constant) solar spectrum, a reduction of the instrument transmission would likely be interpreted as a reduction of albedo by the model. However, by using the measured SMR, the dependency of the albedo on the throughput is effectively reduced. H_2O shows almost no impact of the ice layer ($\rho \sim 0$) and the seasonal variation is clearly seen. Notice that the HWHM continuously increases with time, which may be an indication of general degradation of the instrument.

Since the ice layer affects CO and CH_4 differently, the use of xCO as estimator of the "true" CO VCD does not remove the impact of the ice layer and the results will be affected accordingly.

The time series in Fig. 13 show a general trend of increasing fit errors indicating that the results become continuously worse. Since the model remains the same, this can only be interpreted as a decrease of the measurement quality. The curves also show some time intervals with exceptionally high errors, e.g. June 2006 and 2007, most likely due to the incorporation of some bad pixels in the fit.

4.2 Carbon monoxide

Carbon monoxide is an important trace gas affecting air quality and climate. Although CO is not considered as a greenhouse gas, it is relevant as a precursor for carbon dioxide. CO is also one of the major precursors of tropospheric ozone. It is highly variable in space and time. In the troposphere, about half of the CO originates from anthropogenic sources (e.g. fossil fuel combustion), and further significant contributions are due to biomass burning. With its photochemical lifetime of one to three months, CO is a good tracer of transport in the troposphere as well as in the strato– and mesosphere.

With passive atmospheric remote sensing, carbon monoxide can be observed in several spectral regions from the microwave to the near infrared. CO is a target species of several spaceborne instruments, nb. AIRS (McMillan et al., 2005, 2008), MOPITT (Deeter et al., 2003, 2009), and TES (Rinsland et al., 2006) from NASA's nadir sounders aboard the EOS satellite series; MIPAS and SCIAMACHY on ESA's Envisat, and more recently it has also been observed by IASI on MetOp (Fortems-Cheiney et al., 2009; George et al., 2009; Illingworth et al., 2011).

4.2.1 Retrieval errors

The error of the proxy normalized VCD, Eq. (1), is estimated from the errors of the column scaling factors α_{CO} and α_{CH4}, that, in turn, are obtained from the diagonal elements of the least squares covariance matrix defined by

$$\Xi = \sigma^2 \left(J^T J \right)^{-1} \quad \text{with} \quad \sigma^2 = \| \boldsymbol{y} - \boldsymbol{F}(\boldsymbol{x}) \|^2 / (m - n)$$
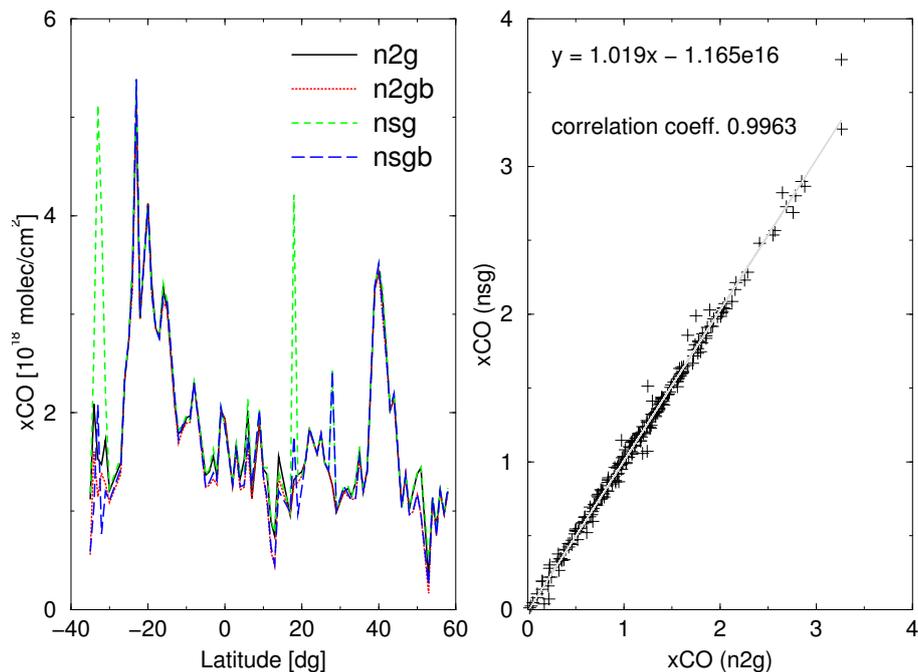
Fig. 11. Sensitivity of CO retrievals with respect to least squares algorithm for orbit 8663 (27 October 2003, covering Russia, the Arabic peninsula, and Eastern Africa): **(a)** comparison of CH_4-normalized CO vertical columns (xCO). "n2g" and "nsg" denote nonlinear and separable least squares, respectively; "b" indicates the bound constrained versions. **(b)** Scatter plot of xCO retrieved using nonlinear least squares ("n2g", horizontal axis) vs. separable least squares ("nsg", vertical axis).

Table 1. Cross correlation coefficients ρ between fit parameters and throughput.

parameter	CO	CH_4	H_2O	HWHM	refl0	refl1	refl2
correlation with throughput	0.374	0.701	0.044	-0.399	0.257	-0.401	0.526

where J denotes the Jacobian (Gay, 1990). Note that the scaled residual norm σ^2 contains both, errors due to instrumental noise and deficiencies of the forward model.

Since the spectral information on carbon monoxide in the observed spectrum is small, the retrievals are specially sensitive to instrumental noise. The signal-to-noise ratio depends on the surface (or cloud) albedo and solar irradiance (essentially SZA), hence dark areas and high latitudes are expected to have higher retrieval errors. Since in most of the cases the CO information is comparable or even lies under the noise level, the precision of the individual estimates is low and it is customary to deliver the CO retrievals as spatial and temporal averages on a regular longitude/latitude grid.

Deficiencies of the forward model introduce systematic errors to the CO estimates. One of the major assumptions in BIRRA is the neglect of scattering within the atmosphere. Homogeneous aerosol and cloud (e.g. low marine stratocumuli) layers are satisfactorily accounted for by proxy modeling. However, under highly convective conditions, cloud heterogeneity can cause large retrieval errors.

Figure 14 shows the spatial distribution of seasonally gridded xCO retrieval errors, averaged over the period 2003 to 2005. Note that the gridded errors represent the mean value of the individual retrieval errors and not the error associated with the spatial-temporal averaged xCO estimate. In this manner, it is possible to visualize the dependence of the retrieval errors on time (seasons) and location. Moreover, the contributions to retrieval errors can be identified taking into account extra information, like surface albedo and climatology. Regions where surface albedo is high (e.g. Sahara, Arabian Peninsula, Australia) exhibit low xCO errors (typically lower than $5 \cdot 10^{17}$ molec cm^{-2}). Analogously, regions with permanent homogeneous marine stratocumuli with high albedo (e.g. South-East Atlantic, Pacific west of Peru and California) also have associated low xCO errors. These findings are in accordance with the number of accepted observations depicted in Fig. 15, i.e. a large number of observations survive the quality screening under ideal retrieval conditions. On the other hand, systematic errors due to the presence of highly inhomogeneous convective clouds dominate the xCO errors. In particular, the highest retrieval errors are found
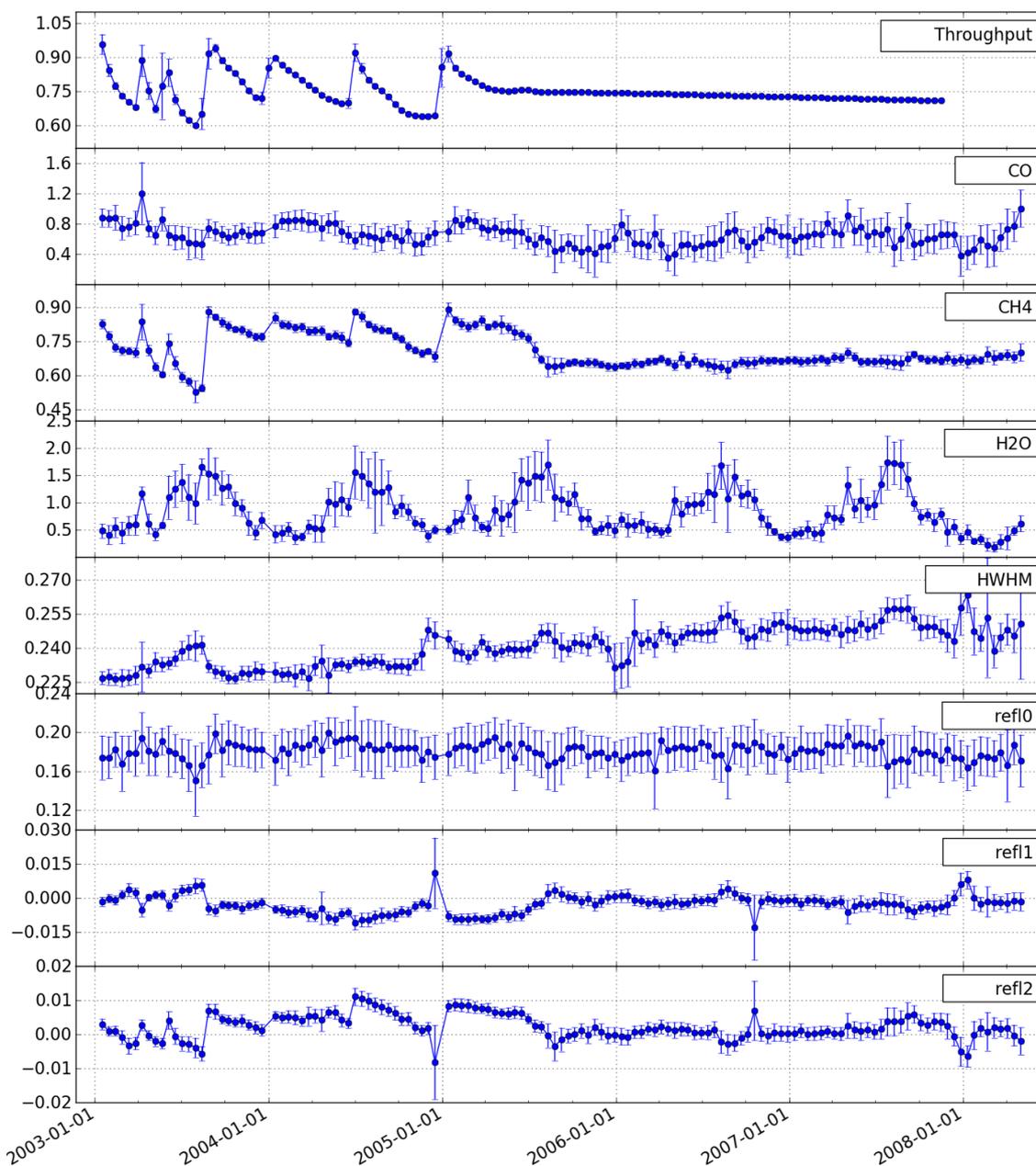
Fig. 12. Time series of 14-day averaged fit parameters included in CO retrievals from channel 8. From top down: the mean channel transmission (throughput), the scaling factors of CO, CH_4 and H_2O, the half width at half maximum of the instrument slit function, and the zeroth, first, and second degree coefficients of the albedo polynomial. See the dependency of some parameters on throughput.

over the Indian Ocean and West Equatorial Pacific during the Asian Monsoon period (from June to September, see bottom left summer panel in Fig. 14). Accordingly, only few observations pass the quality filter, see Fig. 15.

The histograms in Fig. 16 illustrate the occurrence frequency of xCO errors, i.e. the unnormalized error probability density function (PDF). The errors are slightly higher in summer (right panel) than in winter (left panel), which is reflected by the median values and in accordance with Fig. 14.

4.2.2 Spatial distributions

CO vertical columns have been processed for several years from 2003 to 2009 using the BIRRA algorithm and a dynamic bad & dead pixel mask (see Subsect. 3.2). Figure 17 shows the annual mean of CO vertical columns for the years 2003, 2004, and 2005. In addition to the selection criteria mentioned in Sect. 2.3, the relative errors of retrievals (errors of α_{CO} less than 50 %, error of α_{CH4} less than 1 %) are

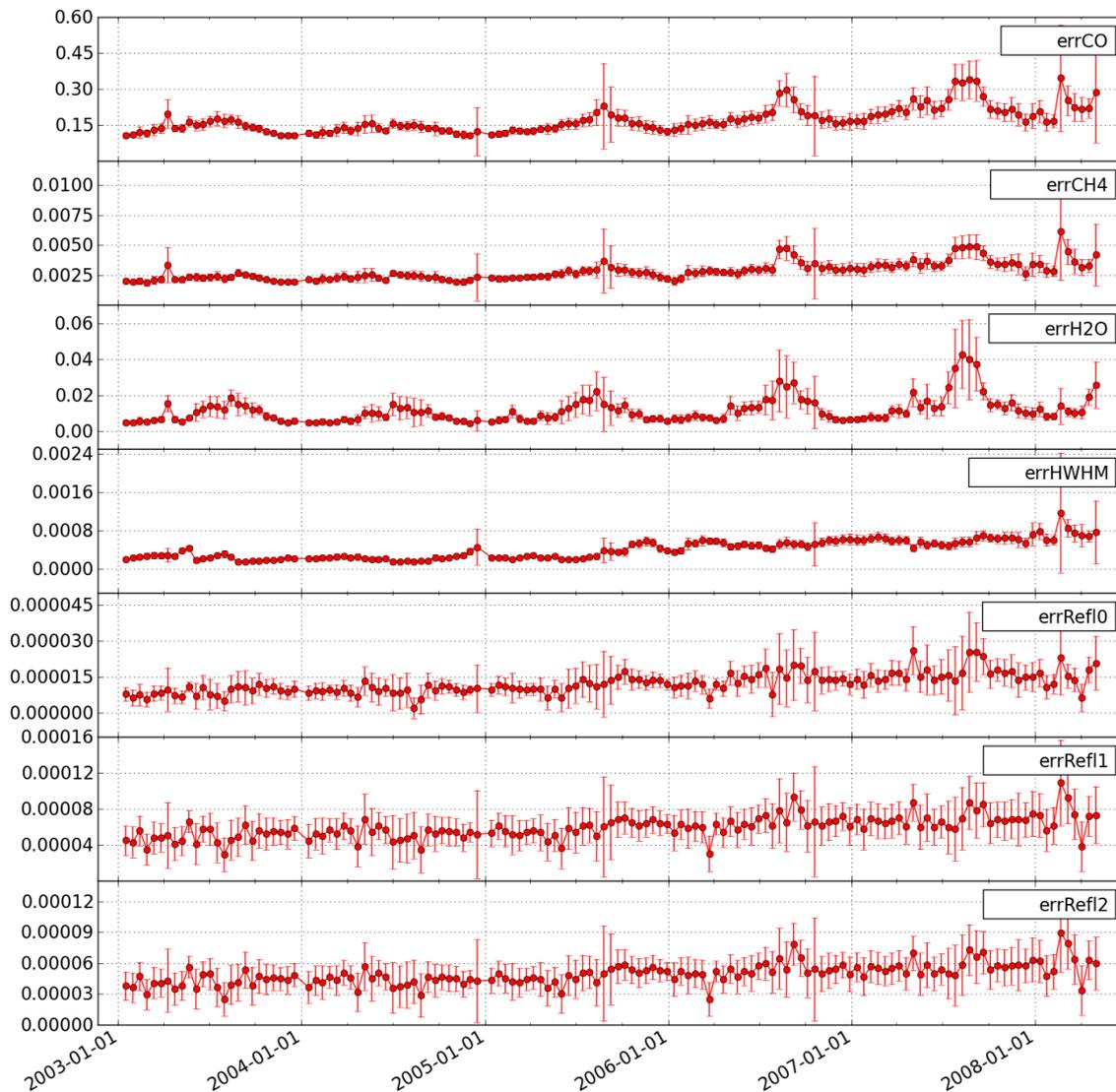**Fig. 13.** Time series of 14-day averaged errors of fit parameters (see Fig. 12). From top down: error of the scaling factors of CO, CH$_4$ and H$_2$O, the half width at half maximum of the instrument slit function, and the zeroth, first, and second degree coefficients of the albedo polynomial.

taken into account. All annual averages show high densities at South East Asia due to anthropogenic emissions and in Central Africa due to high density of biomass burnings during the dry seasons. It can be also noticed that the CO column densities were specially high during 2003. This CO increase in 2003 is related to the increase in biomass burnings during this year, and has also been reported by other studies (e.g. Buchwitz et al., 2007).

Figure 18 provides a closer look to the African continent. In this case, the results are presented as a three year average of the four seasons. Inter-tropical regions are the areas showing higher seasonality. The weather in the tropical region of the Earth is highly influenced by the tropical rain belt, which oscillates between the Northern and

Southern Hemisphere. In the Northern Hemisphere, the wet season comprises roughly the months from April to September, whereas the dry season lasts from October to March. Due to the rain bell oscillation, in the Southern Hemisphere the wet and dry seasons are reverted. During the dry season, biomass burning events are more likely to occur. The seasonality of the fires can be clearly seen in the carbon monoxide distributions: the inter-tropical regions present the highest CO VCDs at the end of the dry seasons (January-February-March in the northern and July-August-September in the Southern Hemisphere).

Figure 19 illustrates the three-year average of xCO over Southeastern Asia. Regions with a high population density such as the Sichuan Basin (Red Basin) in South-West
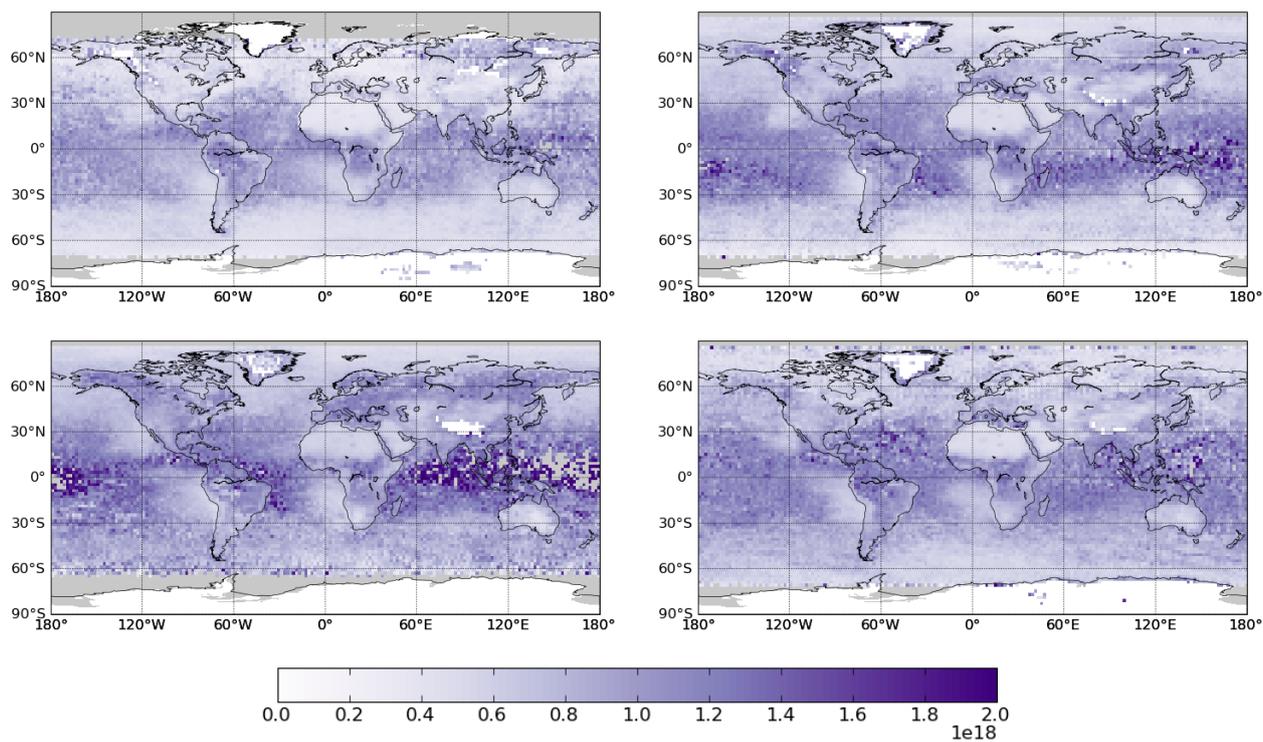
Fig. 14. xCO mean errors [molec cm^{-2}] averaged over the four seasons 2003–2005. Top-left: December-January-February, top-right: March-April-May, bottom-left: June-July-August, bottom-right: September-October-November.
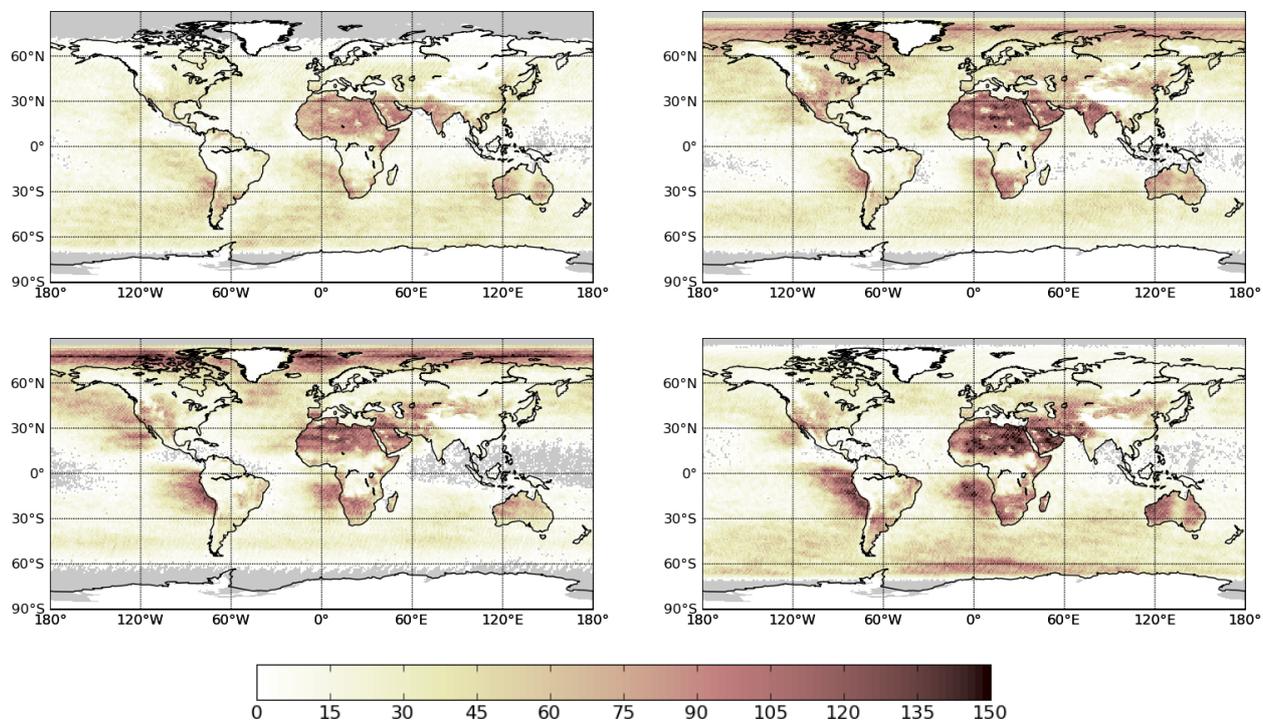

Fig. 15. Number of observations accepted for carbon monoxide VCD retrievals, seasonal three-year averages (arrangement as in Fig. 14).
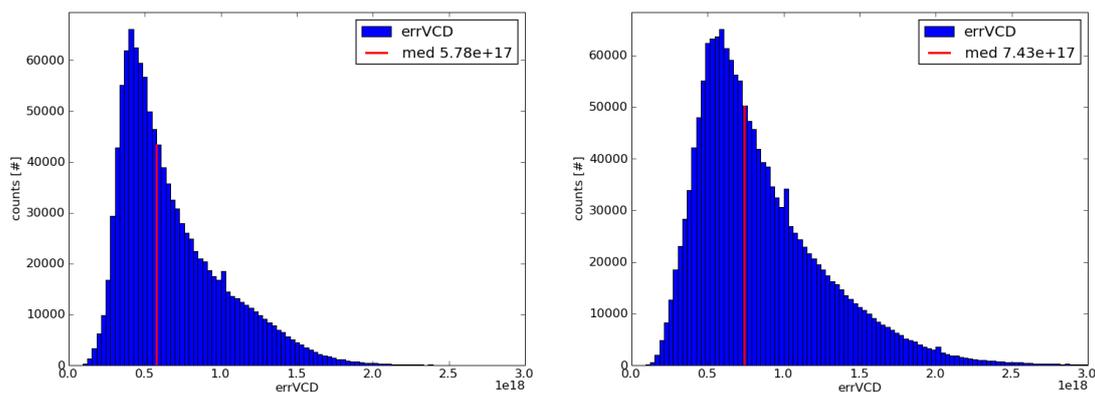
Fig. 16. Histograms of carbon monoxide vertical column density errors, three-year (2003–2005) seasonal averages. Left: winter; right: summer.
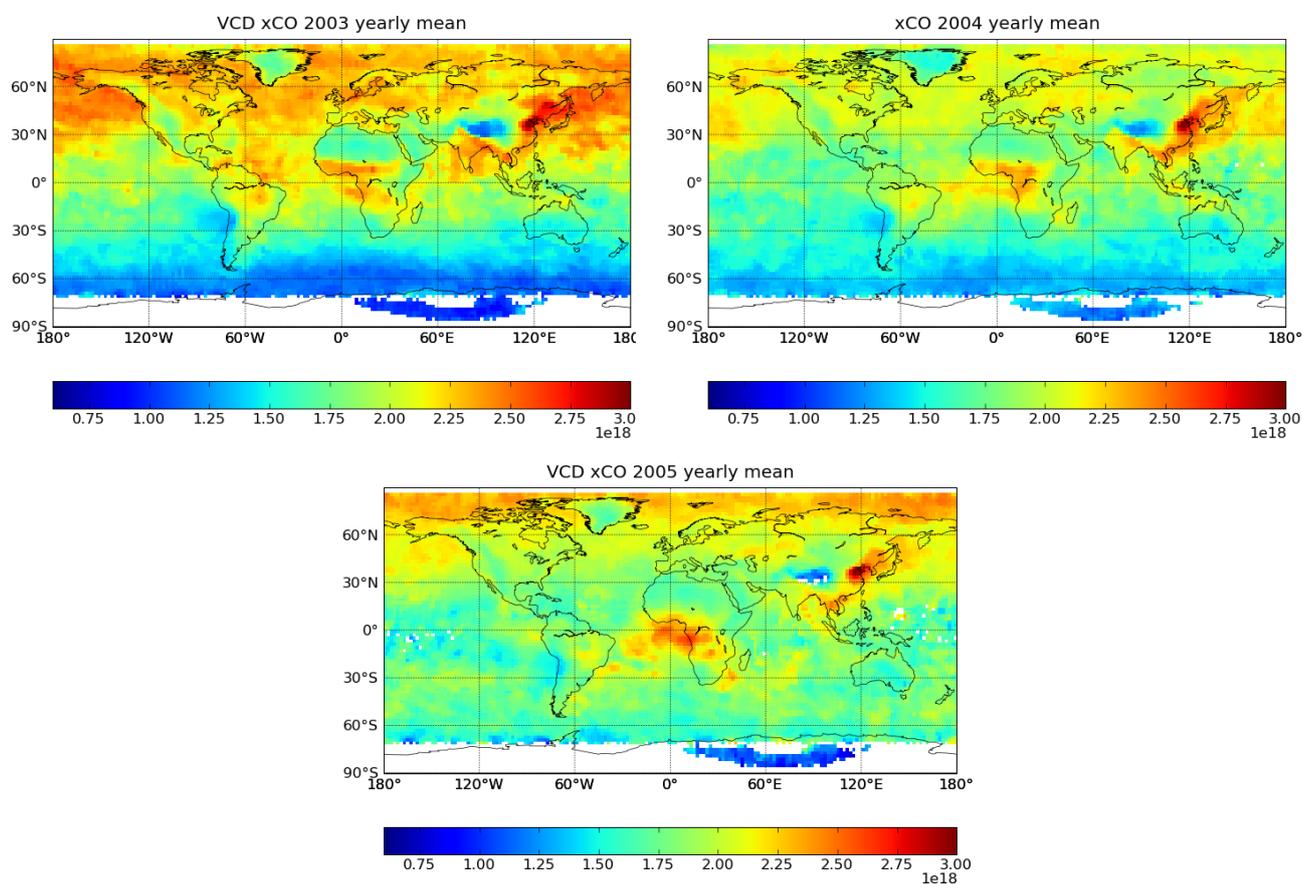

Fig. 17. Annual mean carbon monoxide vertical column densities [molec cm^{-2}] for 2003 to 2005. (1° × 1° grid with 3° × 3° median filter smoothing)

China or the Chinese eastern coast area are clearly visible with a high carbon monoxide abundance, as was already observed by, e.g. SCIAMACHY (Buchwitz et al., 2006; Gloudemans et al., 2009) and the MOPITT mission (Clerbaux et al., 2008).

4.2.3 Intercomparison with ground-based observations

Validation of satellite measurements is frequently performed by intercomparisons with ground-based observations. For example, Sussmann and Buchwitz (2005) have compared column CO measured by SCIAMACHY over Central Europe with Fourier Transform InfraRed (FTIR) profile retrievals at the NDACC Zugspitze station. In an extensive study,
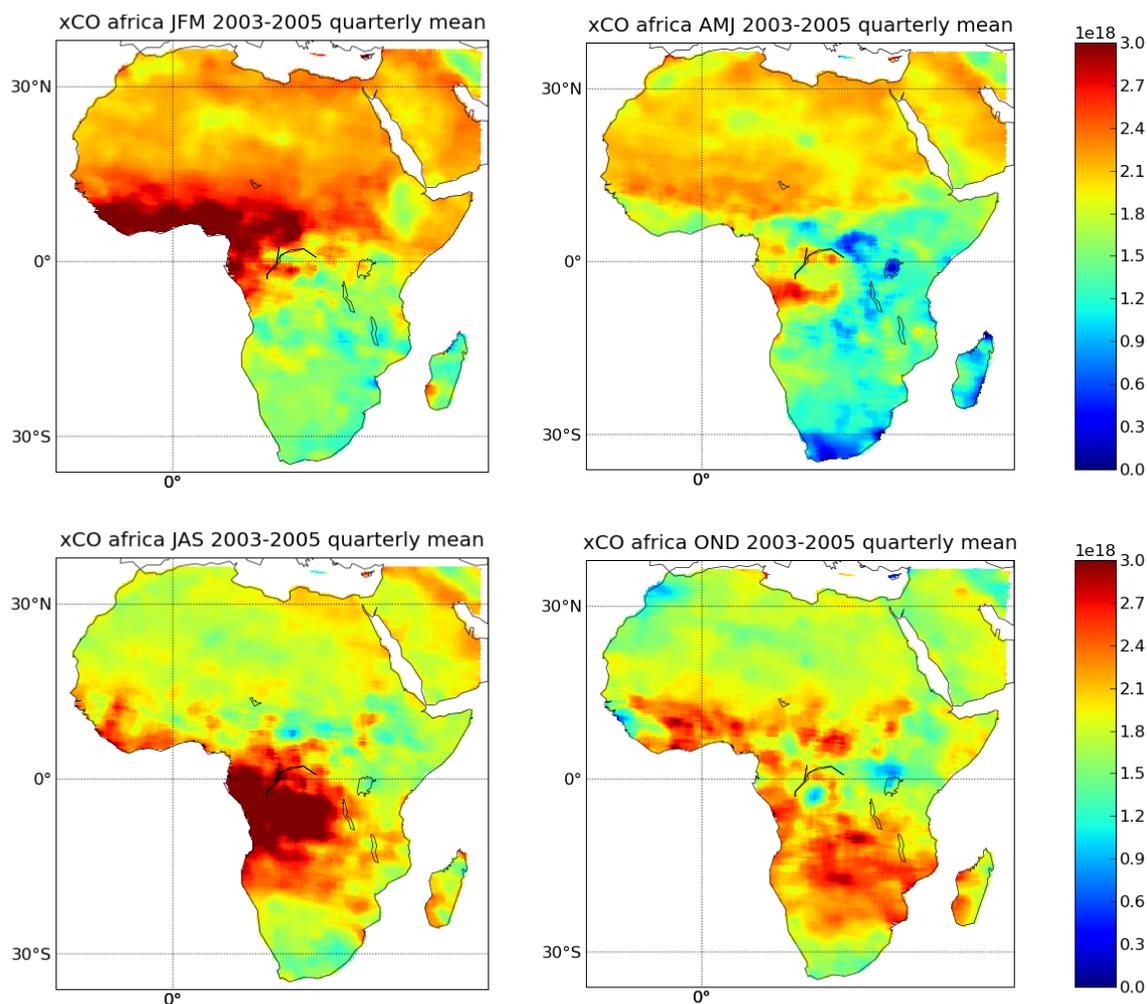
Fig. 18. Three year averages (2003–2005) of quarterly mean CO VCDs [molec cm^{-2}] over Africa. From top-left to bottom-right: January-February-March, April-May-June, July-August-September, and October-November-December. ($0.2° \times 0.2°$ grid with $1° \times 1°$ median filter smoothing)
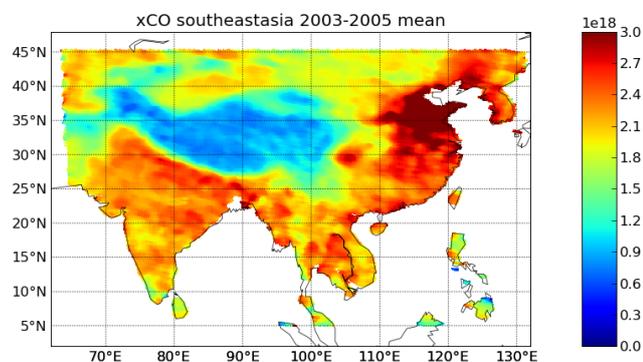

Fig. 19. Three-year average of CO vertical column densities [molec cm^{-2}] over South-East Asia. ($0.2° \times 0.2°$ grid with $1° \times 1°$ median filter smoothing).

Dils et al. (2006) described comparisons between SCIA-MACHY CO, CH$_4$, CO$_2$, and N$_2$O total columns retrieved by three different algorithms (WFM-DOAS, Buchwitz et al., 2004; IMAP-DOAS, Frankenberg et al., 2005a; and IMLM, de Laat et al., 2006) and ground-based FTIR data measured at eleven NDACC (then NDSC) stations. More recently, de Laat et al. (2010) reported good agreement between SCIAMACHY IMLM retrievals of carbon monoxide with twenty ground-based stations (mostly FTIR).

These intercomparisons are delicate since the instruments have different altitude sensitivity and horizontal resolution, the temporal match of the observations is not ideal and the location of the ground stations and the surrounding terrain may also be an impediment. Nevertheless, these studies allow an assessment of, e.g. relative biases or inter-annual variability and are an important element of the validation efforts.

Satellites observe large areas in contrast to the point-like view of uplooking ground-based spectrometers, so they see

in general a larger portion of the atmosphere. If comparisons were performed over a largely uniform surface terrain without emission sources, the problem of the horizontal resolution would be reduced. Accordingly, oceans are natural candidates but, in the absence of clouds, the signal received is rather weak because of the low albedo of water (see also Gloudemans et al., 2009). Another possibility is to use desert areas without significant emissions. Although deserts can exhibit altitude differences, they typically have high surface albedos providing high signal-to-noise ratios.

Here, SCIAMACHY CO BIRRA retrievals are compared to the ground-based measurements provided by the World Data Center for Greenhouse Gases (WDCGG) Assekrem Station (Novelli et al., 2003). It is important to note at this point that the SCIAMACHY CO retrievals represent (dry-air) *column* mixing ratios (i.e. considering the a-priori column mixing ratio q_{CO}^{prior} instead of the a-priori vertical column density N_{CO}^{prior} in Eq. (19)), whereas the Assekrem station measured *volume* mixing ratios at surface level. Such a comparison is only justified for gases with constant mixing ratio profiles (e.g. O_2, CO_2) and this is not the case of CO. This study in not intended to be a validation, since we are comparing two different quantities. However, both retrievals should show similar seasonality features and temporal evolution (and indeed they do) and this is our motivation here.

Figure 20 illustrates time series of 14-day averaged dry-air CO column mixing ratios as observed by SCIAMACHY and CO volume mixing ratios as measured at the Assekrem WDCGG ground station. In the Northern Hemisphere, higher column densities can be expected in winter, and this can be clearly seen in Fig. 20. Furthermore, the seasonal variation of carbon monoxide retrieved with BIRRA is also evident in the ground-based measurements. A good agreement of SCIAMACHY's and the Assekrem ground station CO is found for the years of 2003, 2004 and 2005. Afterward, due to instrument degradation, the SCIAMACHY CO shows a higher dispersion and the seasonal variation is worse represented (esp. in 2006).

4.3 Methane

Methane is the third (second anthropogenic) most important greenhouse gas representing one fifth of the whole radiative forcing of long-life well-mixed gases. Its concentration has increased by more than a factor of two since pre-industrial times with a growth rate of about 1 % per annum (until recently). Atmospheric methane results from anthropogenic (agriculture, fossil fuel combustion, . . .) as well as natural (e.g. wetlands, geological processes) sources. With a life time of about ten years, its spatial and temporal variation is considerably smaller than for CO. As a consequence, the required retrieval precision is much higher.

Atmospheric sounding of methane is performed in the near and thermal infrared, and it is observed by all sensors mentioned above (e.g. Buchwitz et al., 2005; Frankenberg
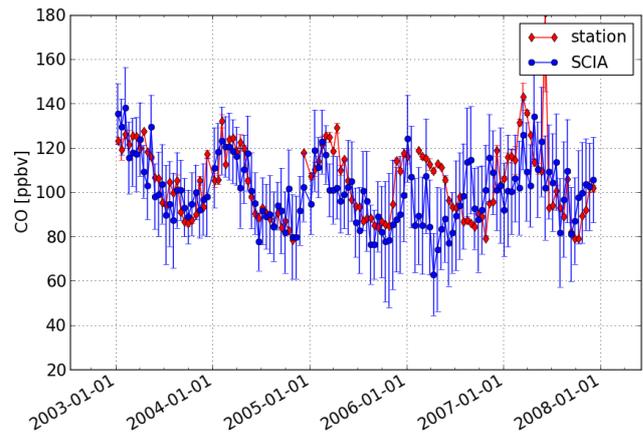

Fig. 20. Intercomparison of SCIAMACHY CO over Central Sahara with Assekrem WDCGG station (Ahaggar Mountains, 2710 m above sea level) data. Note that the SCIAMACHY CO retrievals are dry-air column mixing ratios and the Assekrem WDCGG CO data are volume mixing ratios at surface level. The time step is in both cases 14 days. The error bars represent the standard deviation of the data.

et al., 2006; Gloudemans et al., 2008). Furthermore, it is one of the two target gases of the TANSO Fourier transform spectrometer on-board the recently launched GOSAT satellite (Kuze et al., 2009).

For methane retrievals from SCIAMACHY observations, two microwindows in channel 6 are utilized: the 5986–6139 cm^{-1} interval with CH$_4$ as the strongest absorber, and the 6273–6419 cm^{-1} interval with CO$_2$ as the strongest absorber. For our retrievals, H$_2$O has been considered as additional absorber in both windows, the Gaussian slit function HWHM has been fixed to 2.45 and 2.64 cm^{-1} in the two windows, and albedo was modeled as a second degree polynomial.

Figure 21 gives an impression of three month averages for 2004. Regions of strong emissions, e.g. the northern South America, the equatorial region of Africa, and in Asia are clearly visible and reflect patterns found by, e.g. Frankenberg et al. (2006); Schneising et al. (2009). Seasonal variability such as increased emissions in South East Asia due to rice cultivation is evident especially for July-August-September. Furthermore, the plot shows the shift of methane emissions over wetlands from southern to northern Africa and back.

5 Summary and outlook

A new code BIRRA – Beer InfraRed Retrieval Algorithm – has been developed for level 1b → 2 processing of SCIAMACHY near infrared nadir observations. Recently, the BIRRA code has been successfully implemented into the operational SCIAMACHY processor for the retrieval of carbon monoxide vertical column densities. In view of its efficiency,
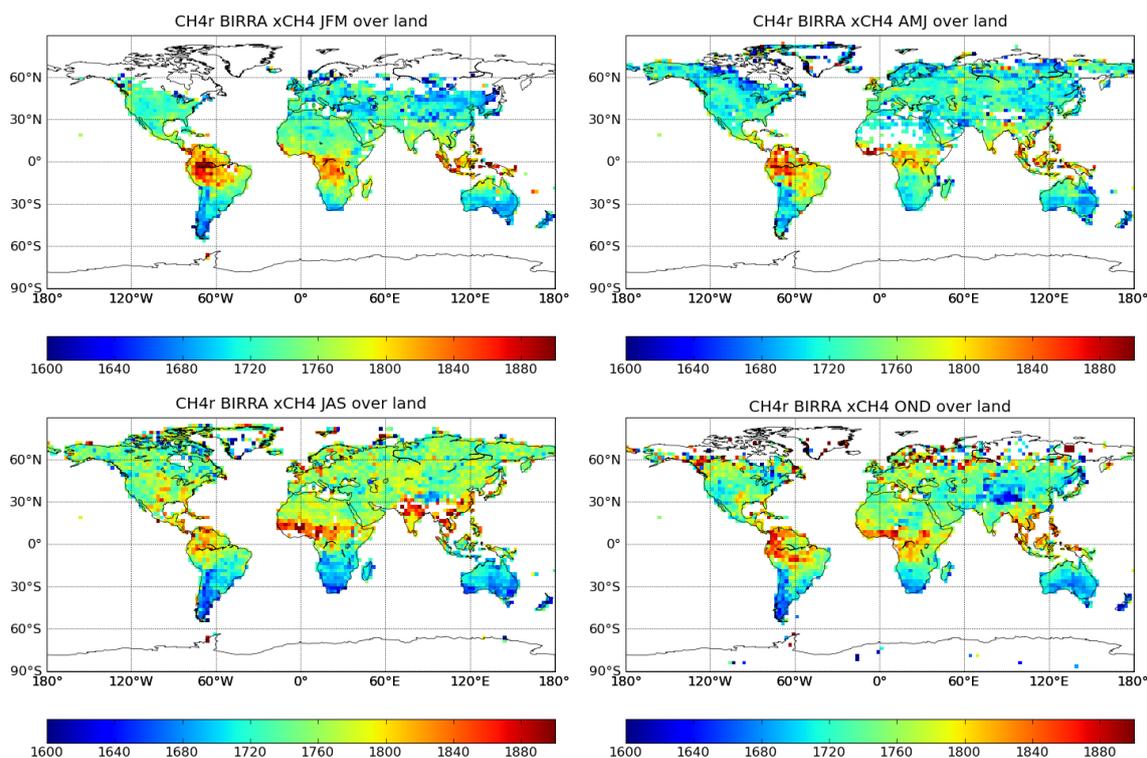
Fig. 21. Quarterly means of methane [ppm] for 2004. In view of the reduced signal over oceans, only CH_4 over land is plotted. ($2° \times 2°$ grid)

flexibility and precision, BIRRA is also a suitable tool for scientific investigations.

The fundamentals of the algorithm have been presented in Sect. 2. BIRRA analyzes radiance spectra instead of optical thicknesses (as in DOAS-like algorithms), since instruments measure the convolved radiance rather than the convolved optical thickness.

The inversion is done by standard nonlinear least squares or by separable least squares solvers, where linear and nonlinear parameters are treated separately. Further distinctive features of the code are the optional use of bound constrained least squares, exact analytical derivatives by automatic differentiation, and "on demand" line-by-line computation of molecular cross sections.

The results presented here are dry-air quantities: vertical column densities in case of CO, with CH_4 as proxy; and dry-air column mixing ratios in case of CH_4, with CO_2 as proxy. Whereas CO_2 is a good proxy for the dry air mass of the observations, CH_4 has some deficiencies (much stronger spectral signatures than CO, spatial variability, non-constant profile, etc). However, due to the much greater variability of CO, CH_4 can be used as an appropriate proxy.

Several aspects of the level 1 data have been investigated. It turned out that the effect of the dead/bad pixel mask (DBPM) has a major impact on the retrievals. In addition to masking spectral pixels with doubtful level 1 quality, one has

to consider further masking due to spectroscopic aspects; in particular, interference of spectral lines of the target gas with strong lines of other gases, or insufficient knowledge of the molecular cross sections due to imprecise spectral line parameters. The use of radiance spectra normalized by sun mean reference (SMR) spectra helps reducing the impact of the ice layer over the detector on the retrievals. However, in case of BIRRA, the ice layer affects CO and CH_4 differently and further treatment is needed. The pixel-to-wavelength relationship of channel 8 in SCIAMACHY level-1b product is set to the on-ground calibration and a spectral correction is needed. We found that the required spectral correction has roughly a second-degree polynomial shape that can be well approximated by a first-degree polynomial in the CO fitting window.

A survey of carbon monoxide and methane VCD retrievals has been presented with emphasis on the years 2003 to 2005. "Pre-" and "postprocessing" of the data turned out to be crucial, i.e. careful preparation of the level 1b data used as input to the least squares fitting and a meticulous examination of the fitted column density scaling factors for the generation of scientific products is mandatory. In addition, the final product is highly sensitive to the correct filtering of dubious retrievals and to the appropriate consideration of scene cloudiness. Moreover, careful analysis of the time series of all fit parameters with respect to the instrumental mean

transmittance provides valuable hints for the postprocessing. Preliminary work exploiting SCIAMACHY's channel 6 shows BIRRA's potential for methane (and carbon dioxide) retrievals.

The further development of BIRRA is motivated by its dual role as an operational processor and a scientific tool. For the "scientific prototype", we are currently working on an optimization and fine-tuning of the multiwindow fitting (required for methane and carbon dioxide retrievals). Furthermore, a better climatology esp. of temperature and water would be beneficial and likewise the treatment of clouds and aerosols can be improved in the forward model and/or in the post-processing. Finally, our verification and validation efforts will be intensified, e.g. by intercomparisons with ground-based observations (NDACC-TCCON) as well as thermal infrared sounders such as AIRS, GOSAT, IASI, MOPITT, and TES (Schreier et al., 2010). Clearly, the lessons learned from the scientific analysis will be a valuable guide for the ongoing upgrade of the operational processor. The further refinement will allow a better analysis of the years 2006 and beyond that are even more challenging due to the continuous channel degradation, nb. increasing number of bad or dead pixels.

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