

Spectral expansions of non-self-adjoint generalized Laguerre semigroups

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Abstract

We provide the spectral expansion in a weighted Hilbert space of a substantial class of invariant non-self-adjoint and non-local Markov operators which appear in limit theorems for positive-valued Markov processes. We show that this class is in bijection with a subset of negative definite functions and we name it the class of generalized Laguerre semigroups. Our approach, which goes beyond the framework of perturbation theory, is based on an in-depth and original analysis of an intertwining relation that we establish between this class and a self-adjoint Markov semigroup, whose spectral expansion is expressed in terms of the classical Laguerre polynomials. As a by-product, we derive smoothness properties for the solution to the associated Cauchy problem as well as for the heat kernel. Our methodology also reveals a variety of possible decays, including the hypocoercivity type phenomena, for the speed of convergence to equilibrium for this class and enables us to provide an interpretation of these in terms of the rate of growth of the weighted Hilbert space norms of the spectral projections. Depending on the analytic properties of the aforementioned negative definite functions, we are led to implement several strategies, which require new developments in a variety of contexts, to derive precise upper bounds for these norms.

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Introduction and main results

The importance of the spectral reduction of linear operators cannot be overemphasized as it is amply demonstrated by its diverse applications to far reaching fields of mathematics such as functional analysis, dynamical systems, topological groups, differential geometry, probability theory, harmonic and complex analysis and boundary value problems. Beyond the theoretical interests, these developments have proved to be fruitful in solving problems in the theory of infinite dimensional systems, in theoretical and experimental physics, economics, statistics and many other fields of applied mathematics.

Spectral theory at its best can be seen when one considers normal operators in Hilbert spaces. Although the class of non-self-adjoint (NSA) and non-local operators is central and generic in the study of linear operators, its spectral analysis is fragmentarily understood due to the fundamental technical difficulties arising when the properties of symmetry and locality are simultaneously relaxed. We refer the interested reader to the classical monographs [DS71] and [Dav07] and survey papers [Dun58] and [Dav02], [Sjö09] for a thorough and yet up-to-date account on the spectral reduction of (NSA) linear operators.

The main purpose of this work is to design an original and comprehensive theory to

- 1) provide the spectral representation in a weighted Hilbert space of the solution to the Cauchy problem associated to a class of NSA and non-local linear operators, which are central in the theory of Markov processes,
- 2) derive an eigenvalues representation of the heat kernel, i.e. the transition kernel of the underlying Markov process,
- 3) study the smoothness properties of the solution to the Cauchy problem and of the heat kernel,
- 4) obtain the speed of convergence to equilibrium,
- 5) develop a detailed analysis of the nature of the spectrum of this class of NSA and non-local operators.

Our approach is based on an in-depth study of an intertwining relation that we elaborate between the class of NSA and non-local operators and a specific self-adjoint differential operator whose spectral resolution is well understood. We already point out that it is flexible enough to study the substantial problems listed above in a unified way. We also believe that it is comprehensive and could be used to deal with the spectral expansion of different type of NSA operators. For these reasons,

a synthetic description of the methodology will be detailed in Chapter 2. These semigroups are defined below, where in the sequel $\mathbf{C}_0(\overline{\mathbb{R}}_+)$ (resp. $\mathbf{C}_0(\mathbb{R}_+)$) stands for the space of continuous functions on $\overline{\mathbb{R}}_+ = [0, \infty)$ (resp. $\mathbb{R}_+ = (0, \infty)$) both vanishing at infinity and endowed with the uniform topology and $\mathbf{C}_c^\infty(A)$ stands for the class of infinitely differentiable functions with compact support on $A \subseteq \mathbb{R}$.

DEFINITION 1.0.1. We say that a semigroup $P = (P_t)_{t \geq 0}$ is a **generalized Laguerre semigroup (for short gL semigroup)** of order $H > 0$ if

1. P is a Feller semigroup on $\mathbf{C}_0(\overline{\mathbb{R}}_+)$, i.e. for all $t \geq 0$ and $f \in \mathbf{C}_0(\overline{\mathbb{R}}_+)$ it holds that $P_t f \in \mathbf{C}_0(\overline{\mathbb{R}}_+)$, $P_t f \geq 0$ whenever $f \geq 0$, and, $\lim_{t \downarrow 0} P_t f(x) = f(x)$, $\forall x \geq 0$.
2. The family of operators $(K_t)_{t \geq 0}$ defined, for any $t \geq 0$, via

$$K_t f = P_{\ln(t+1)} d_{(t+1)^H} f,$$

with $d_c f(x) = f(cx)$, defines a Feller semigroup on $\mathbf{C}_0(\overline{\mathbb{R}}_+)$, which possesses the H -self-similarity property, i.e. the following relation holds, for all $t, x, c > 0$,

$$(1.1) \quad K_t f(cx) = K_{c^{-H}t} d_c f(x).$$

3. There exists a non-degenerate probability measure ϑ on \mathbb{R}_+ such that we have, for any $t \geq 0$,

$$(1.2) \quad \vartheta P_t f = \vartheta f,$$

where here and in the sequel $\vartheta f = \int_0^\infty f(x) \vartheta(dx)$.

4. For any $x \in \mathbb{R}_+$, there exists a positive measure $\Pi(x, dy)$, such that for any f with support included in $\mathbb{R}_+ \setminus \{x\}$ and $f \in \mathbf{C}_c^\infty(\mathbb{R}_+ \setminus \{x\})$, we have that

$$\lim_{t \rightarrow 0} \frac{1}{t} P_t f(x) = \int_{\mathbb{R} \setminus \{x\}} f(y) \Pi(x, dy)$$

and $\Pi(x, (x, \infty)) = 0$.

We say that a stochastic process $X = (X_t)_{t \geq 0}$ defined on a filtered probability space $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbb{P})$ is a generalized Laguerre process (for short gLp) if the family of linear operators $P = (P_t)_{t \geq 0}$ defined, for any $t \geq 0$ and $f \in \mathbf{C}_0(\overline{\mathbb{R}}_+)$, by

$$P_t f(x) = \mathbb{E}_x [f(X_t)]$$

is a gL semigroup, where \mathbb{E}_x stands for the expectation operator associated to $\mathbb{P}_x(X_0 = x) = 1$.

REMARK 1.0.2. Since there is no loss of generality considering only gL semigroups of order 1, we will assume from now on that $H = 1$. Indeed, note that if P is a gL semigroup of order $H > 0$ then, since $p_H(x) = x^H$ is a homeomorphism of \mathbb{R}_+ to \mathbb{R}_+ , then $\bar{K}_t f(x) = K_t f \circ p_H(x^{\frac{1}{H}})$ is 1-self-similar Feller semigroup on $\mathbf{C}_0(\overline{\mathbb{R}}_+)$ and $\bar{P}_t f(x) = P_t f \circ p_H(x^{\frac{1}{H}})$ is a gL semigroup of order precisely 1.

REMARK 1.0.3. The terminology generalized Laguerre semigroup is motivated by the well-known case when the invariant measure

$$\vartheta(dx) = \varepsilon_m(x) dx = \frac{x^m}{\Gamma(m+1)} e^{-x} dx, \quad x > 0,$$

is the probability measure of a Gamma random variable of parameter $m + 1 \geq 1$. Then, there exists a family of semigroups P , indexed by m , satisfying the conditions of Definition 1.0.1 above and each of which admits an eigenvalues expansion for its action on the Hilbert space of square integrable functions with respect to $\varepsilon_m(x)dx$, that is $L^2(\varepsilon_m)$, expressed in terms of the Laguerre polynomials of order m , which form an orthogonal basis in the Hilbert space $L^2(\varepsilon_m)$, see Chapter 3. In this case, P is self-adjoint in $L^2(\varepsilon_m)$, and we shall see that it is the only instance among the class of gL semigroups to be self-adjoint, see Theorem 7.0.3(4), which disproves the orthogonality of the eigenfunctions of these groups apart from the classical Laguerre case.

REMARK 1.0.4. Let us now indicate that the requirement $\Pi(x, (x, \infty)) = 0$ in condition 4, which has a nice pathwise interpretation as the imposition upon the corresponding generalized Laguerre process to have downward jumps only, is purely technical in the sense that one may define this class without this assumption. We choose to restrict our analysis to this situation since it requires already some substantial new developments which deserve to be detailed in one (long) paper and we postpone to a future paper the study of the general case. However, we emphasize that apart from some technical issues specific to the general framework, the main concepts, theories and methodologies needed to develop its spectral reduction are contained in this work. Indeed, from an operator viewpoint, the main mathematical difficulties stem from the non-self-adjointness combined with the non-local properties of the generators of these Markov processes and not necessarily on the support of the Lévy kernel, that is $\Pi(x, \cdot)$, see [PS18b] for the spectral expansion of two-sided self-similar Feller semigroups. As another illustration of this fact, one may extract from our main results Theorem 1.3.1, the spectral expansion of the semigroups constructed from the gL ones by performing a subordination in the sense of Bochner, which gives a class of non-self-adjoint semigroups whose associated Markov process may have jumps in both directions.

Beyond the theoretical interests of studying a class of NSA semigroups, there are several motivations underlying the investigation of this specific class of semigroups.

REMARK 1.0.5. On the one hand, the gL semigroups play a substantial role in probability theory. Indeed, from the celebrated transformation due to Lamperti [Lam62], which provides a bijection between stationary processes and self-similar stochastic processes, one can easily observe that the class of gL semigroups we have introduced above corresponds to the class of positive stationary Markov processes whose Lamperti transformation, see (1.1), leaves invariant the Markov property, that is the associated self-similar process remains a (homogeneous) Markov. The generalized Laguerre processes are intimately connected to the so-called generalized Ornstein-Uhlenbeck processes which have been introduced by Carmona et al. [CPY97]. Moreover, the aforementioned class of self-similar semigroups, whose heat kernel can easily be expressed in terms of the one of the associated gL semigroup via (1.1), appears naturally in limit theorems of Markov processes and have been intensively studied over the last two decades, see e.g. [Lam72], [VA94], [BS11], [BY02a], [Pat09], [Pat12] and [DZ15]. They also correspond to the class of positive Markov processes whose transition kernel satisfies a scale

invariance property similar to the classical Gaussian heat kernel whose index of self-similarity is $H = 2$. It is (one of) the aim of this work to make available additional explicit representations of Markov transition kernels satisfying this scale invariance property which has been observed in many physical phenomena.

We also emphasize that the class generators of the gL semigroups, see (1.6) below, encompasses a variety of substantial integro-differential operators such as the classical fractional derivatives, delay differential-operator of pantograph-type, differential Bessel operators perturbed by non-local operators, see e.g. [PS12b], [BP14] and [Die10], and, we refer to Chapter 3 for the description of some specific instances. Several classes of gL Markov processes have found applications in many fields of sciences, such as neurology, data transmission, economy, biology, epidemiology, see e.g. [DJG10], [YZLM14], [CKRT12], [Mal15] and the references therein.

1.1. Characterization and properties of gL semigroups

We shall provide a characterization of the class of gL semigroups in terms of their infinitesimal generators. To this end, let us introduce the function $\psi : i\mathbb{R} \mapsto \mathbb{C}$, which is defined, for any $z \in i\mathbb{R}$, by the following relation

$$(1.3) \quad \psi(z) = \sigma^2 z^2 + mz + \int_0^\infty (e^{-zy} - 1 + zy) \Pi(dy),$$

where $\sigma \geq 0$, $m \geq 0$ and Π is a σ -finite positive measure concentrated on $(0, \infty)$ and satisfying the integrability condition $\int_0^\infty (y^2 \wedge y) \Pi(dy) < \infty$. Note that the characteristic triplet (σ, m, Π) uniquely defines the function ψ .

We introduce the set

$$(1.4) \quad \mathcal{N} = \{\psi \text{ of the form (1.3)}\}.$$

\mathcal{N} is a subset of the negative definite functions or equivalently is a subset of the set of characteristic exponents of infinitely divisible distributions on \mathbb{R} , see e.g. [Jac01] for more details about Fourier transforms of infinitely divisible distributions. The class of infinitely divisible distributions in turn is the building block of Lévy processes via the fact that $(\xi_t)_{t \geq 0}$ is a Lévy process if and only if ξ_1 is infinitely divisible, see e.g. [Ber96, Chap. I] or for our specific case the discussion succeeding Theorem 1.1.1 below. We also write throughout

$$(1.5) \quad \bar{\Pi}(y) = \int_y^\infty \Pi(dr) \text{ and } \bar{\bar{\Pi}}(y) = \int_y^\infty \bar{\Pi}(r) dr$$

for the tails of the measure Π which due to the link to Lévy processes is usually called the Lévy measure. We say that $f \in \mathcal{C}^2([-\infty, \infty])$ if and only if f, f', f'' are continuous on $(-\infty, \infty)$ and have finite limits at $\pm\infty$. We are now ready to state our first result which provides a characterization of our class of generalized Laguerre semigroups of order 1, see Remark 1.0.2, together with some of their basic properties.

THEOREM 1.1.1. *(1) There exists a bijection between the class of generalized Laguerre semigroups and the subspace of negative definite functions \mathcal{N} . More specifically, for each $\psi \in \mathcal{N}$, the associated gL semigroup $P = (P_t)_{t \geq 0}$ is characterized by its infinitesimal generator \mathbf{G} which admits, at*

least, for any function $f \in \mathcal{D}_{\mathbf{G}} = \{f; x \mapsto f_e(x) = f(e^x) \in \mathcal{C}^2([-\infty, \infty])\}$, the representation

$$(1.6) \quad \begin{aligned} \mathbf{G}f(x) &= \sigma^2 x f''(x) + (m + \sigma^2 - x) f'(x) \\ &+ \int_0^\infty (f(e^{-y}x) - f(x) + yx f'(x)) \Pi(x, dy), \end{aligned}$$

where $\Pi(x, dy) = \frac{\Pi(dy)}{x}$ and (σ, m, Π) is the characteristic triplet of ψ .

- (2) For each $\psi \in \mathcal{N}$, P admits a unique invariant measure which is an absolutely continuous probability measure with a density denoted by ν . Moreover, ν has support on $(0, \mathfrak{r})$ and it is positive on this domain, where $\mathfrak{r} = \lim_{u \rightarrow \infty} \frac{\psi(u)}{u} = \phi(\infty)$, see (1.9) below for the justification of the notation, that is

$$(1.7) \quad \mathfrak{r} = \infty \text{ if } \sigma^2 > 0, \text{ and, } 0 < \mathfrak{r} = \overline{\overline{\Pi}}(0^+) + m \leq \infty \text{ otherwise.}$$

- (3) P can be extended uniquely to a strongly continuous contraction semigroup, still denoted by P , on the weighted Hilbert space

$$\mathbb{L}^2(\nu) = \left\{ f : (0, \mathfrak{r}) \rightarrow \mathbb{R} \text{ measurable; } \int_0^\mathfrak{r} f^2(x) \nu(x) dx < \infty \right\},$$

endowed with the norm $\|\cdot\|_{\mathbb{L}^2(\nu)} = \|\cdot\|_\nu$ and inner product $\langle \cdot, \cdot \rangle_\nu$. The algebra of polynomials, \mathbf{P} , is a core for its infinitesimal generator.

- (4) It admits a conservative standard Markov process on $(0, \mathfrak{r})$, see e.g. [BG68, Definition (9.2), p.45] for definition, as a (weak) dual with the measure $\nu(x)dx$ serving as reference measure. The corresponding semigroup $P^* = (P_t^*)_{t \geq 0}$ is Feller-Dynkin, i.e. $P_t^* \mathbf{C}_b(0, \mathfrak{r}) \subseteq \mathbf{C}_b(0, \mathfrak{r})$, where $\mathbf{C}_b(0, \mathfrak{r})$ is the space of continuous and bounded functions $(0, \mathfrak{r})$, when $m > 0$, and, $\mathfrak{r} = \infty$ or when $N_\mathfrak{r} \geq 1$, when $\mathfrak{r} < \infty$, where $N_\mathfrak{r}$ is defined by

$$(1.8) \quad N_\mathfrak{r} = \left\lceil \frac{\overline{\overline{\Pi}}(0^+)}{\mathfrak{r}} \right\rceil - 1 \in [0, \infty]$$

with $\lceil \cdot \rceil$ the ceiling function, i.e. it evaluates the smallest integer greater or equal to a positive real and we use the convention that $N_\mathfrak{r} = \infty$ when $\overline{\overline{\Pi}}(0^+) = \infty$. Moreover P^* admits a contraction semigroup extension in $\mathbb{L}^2(\nu)$ satisfying, for any $f, g \in \mathbb{L}^2(\nu)$,

$$\langle P_t f, g \rangle_\nu = \langle f, P_t^* g \rangle_\nu.$$

- (5) Moreover, for any $N_\mathfrak{r} > 2$, the infinitesimal generator \mathbf{G}^* of the (Feller-Dynkin) semigroup P^* takes the form, for at least any function f such that $f\nu \in \mathcal{D}_{\mathbf{G}}$, $x \in (0, \mathfrak{r})$, and, writing $\bar{f}(x, y) = f(e^{-y}x) - f(x) + yx f'(x)$,

$$\mathbf{G}^* f(x) = \sigma^2 x f''(x) + \left(\mathfrak{d}(x) + x \left(2 \frac{\nu'(x)}{\nu(x)} + 1 \right) \right) f'(x) + \int_0^\infty \bar{f}(x, y) \Pi_\nu^*(x, dy)$$

where $\mathfrak{d}(x) = \sigma^2 - m - \int_0^\infty \left(1 - \frac{\nu(x)}{\nu(e^{-y}x)} \right) xy \Pi_\nu^*(x, dy)$ and

$$\Pi_\nu^*(x, dy) = \frac{\nu(e^{-y}x)}{x\nu(x)} \Pi(dy).$$

- (6) If $\Pi \equiv 0$, then P is self-adjoint in $L^2(\nu)$. Otherwise, P is non-self-adjoint, i.e. $P \neq P^*$.

The proof of the item (1) is given in Section 2.2. The existence and the absolute continuity property of the invariant measure in item (2) along with the continuous extension of P in item (3) are justified in Section 2.3. The support and positivity properties of ν stated in item (2) are derived in Section 5.2.3 and are part of Theorem 5.0.2, whereas the expression of τ is justified in Proposition 4.1.1(8) and the proof of the uniqueness of the invariant measure is postponed to Section 7.4. The fact that the algebra \mathbf{P} is a core follows directly from Theorem 7.0.3(1) and (2). Finally, the remaining statements are proved in Section 5.7.

REMARK 1.1.2. Note that the bijection stated in (1) finds its root in a remarkable work from Lamperti [Lam72] which establishes a bijection between the class of positive self-similar Markov processes and the class of Lévy processes.

REMARK 1.1.3. We also point out that further fine distributional properties, including smoothness and small and large asymptotic behaviour, of the invariant density ν are stated in Chapter 5.

1.2. Definition and properties of subsets of \mathcal{N}

Our main results below regarding the spectral decomposition, regularity and speed of convergence for the $g\mathbb{L}$ semigroups are expressed in terms of some subsets of negative definite functions that we define and study here. To this end, we first point out that the set \mathcal{N} is also in bijection with the set of spectrally negative Lévy processes, i.e. processes that can jump downwards only, which have a non-negative mean, via the relation

$$\log \mathbb{E} [e^{z\xi_1}] = \psi(z), \quad z \in i\mathbb{R},$$

with $\xi = (\xi_t)_{t \geq 0}$ a Lévy process, see [Ber96, Chap. I]. We note that $\mathbb{E}[\xi_1] = \psi'(0^+) = m \geq 0$. We call ξ the Lévy process underlying the semigroup P or alternatively the Lévy process associated to ψ . It is a common fact that any $\psi \in \mathcal{N}$ has the Wiener-Hopf factorization given by

$$(1.9) \quad \psi(u) = u\phi(u), \quad u \in \mathbb{R}_+,$$

where, with $\bar{\Pi}(y) = \int_y^\infty \Pi(dr)$, see (1.5), for any $u \in \mathbb{R}_+$,

$$(1.10) \quad \phi(u) = m + \sigma^2 u + \int_0^\infty (1 - e^{-uy}) \bar{\Pi}(y) dy.$$

The function ϕ , which is a Bernstein function, is the Laplace exponent of the so-called descending ladder height process of ξ , say $\eta = (\eta_t)_{t \geq 0}$, that is, for any $u \in \mathbb{R}_+$,

$$(1.11) \quad \ln \mathbb{E} [e^{-u\eta_1}] = -\phi(u).$$

Note that η is a possibly killed subordinator (non-decreasing Lévy process) and we refer to [Ber96, Chap. VII] for more detail on spectrally negative Lévy processes including their Wiener-Hopf factorization. We point out that η is killed if and only if $m > 0$ as then $\phi(0) > 0$. Denote by

$$(1.12) \quad \mathcal{B}_{\mathcal{N}} = \{\phi \text{ is of the form (1.10)}\},$$

that is the convex sub-cone of the set of Bernstein functions which is in bijection with the set \mathcal{N} via the identity (1.9). Next, for any $\phi \in \mathcal{B}_{\mathcal{N}}$, we introduce the multiplicative Markov operator \mathcal{I}_{ϕ} , defined, for at least any $f \in \mathcal{C}_0(\overline{\mathbb{R}}_+)$, by

$$(1.13) \quad \mathcal{I}_{\phi}f(x) = \mathbb{E}[f(xI_{\phi})],$$

where I_{ϕ} is the positive random variable

$$(1.14) \quad I_{\phi} = \int_0^{\infty} e^{-nt} dt.$$

Write $\mathfrak{R} = \{(\alpha, \mathbf{m}); \alpha \in (0, 1] \text{ and } \mathbf{m} \geq 1 - \frac{1}{\alpha}\}$. Then, for any $(\alpha, \mathbf{m}) \in \mathfrak{R}$ we define for $u \geq 0$ the function

$$(1.15) \quad \phi_{\alpha, \mathbf{m}}^R(u) = \frac{\Gamma(\alpha u + \alpha \mathbf{m} + 1)}{\Gamma(\alpha u + \alpha \mathbf{m} + 1 - \alpha)},$$

and we simply write $\phi_{\mathbf{m}}^R(u) = \phi_{1, \mathbf{m}}^R(u) = u + \mathbf{m}$. We recall from [PS17], see also Lemma 4.2.1 below, that $\phi_{\alpha, \mathbf{m}}^R \in \mathcal{B}_{\mathcal{N}}$. We say that a function f is completely monotone if $f : \mathbb{R}_+ \rightarrow [0, \infty) \in \mathcal{C}^{\infty}(\mathbb{R}_+)$ and $(-1)^n f^{(n)}(u) \geq 0$, for all $n = 0, 1, \dots$ and $u > 0$. The space $\mathcal{C}^{\infty}(\mathbb{R}_+)$ consists of all infinitely differentiable functions on \mathbb{R}_+ .

We are now ready to introduce in Table 1 and Table 2 some substantial subclasses of \mathcal{N} and refer to Section 1.2 for more detailed information regarding these objects.

	σ^2	$\overline{\Pi}(0^+)$	$\overline{\overline{\Pi}}(0^+)$
\mathcal{N}_P	> 0		
\mathcal{N}_{∞}	> 0 0	$= \infty$	
\mathcal{N}_{∞}^c	0	$< \infty$	$< \infty$

TABLE 1. Definition of some classes in terms of the characteristic triplet.

$\mathcal{N}_{\underline{\Theta}}$	$\underline{\Theta}_{\phi} = \lim_{b \rightarrow \infty} \int_0^{\infty} \ln \left(\frac{ \phi(by+ib) }{\phi(by)} \right) dy > 0$
\mathcal{N}_{α}	$\psi(u) \underset{\infty}{\sim} C_{\alpha} u^{\alpha+1}, C_{\alpha} > 0, \alpha \in (0, 1)$
$\mathcal{N}_{\alpha, \mathbf{m}}$	$u \mapsto \frac{\phi_{\alpha, \mathbf{m}}^R(u)}{\phi(u)}$ is completely monotone
\mathcal{N}_R	$\bigcup_{(\alpha, \mathbf{m}) \in \mathfrak{R}} \mathcal{N}_{\alpha, \mathbf{m}}$

TABLE 2. Definition of other classes in terms of $\psi(u) = u\phi(u)$.

Although for most of the subsets of \mathcal{N} , such as $\mathcal{N}_\infty, \mathcal{N}_\infty^c, \mathcal{N}_P$, their definitions are given directly in terms of the triplet (σ, m, Π) , the other classes, i.e. $\mathcal{N}_\alpha, \mathcal{N}_{\underline{\Theta}}$ and \mathcal{N}_R , are more or less characterized through specific properties of the (associated ladder height) Laplace exponent ϕ . The aim of the next result is to provide for these latter subclasses sufficient conditions expressed in terms of (σ, m, Π) that allow to identify them.

THEOREM 1.2.1. (1) *If $\psi \in \mathcal{N}_{\underline{\Theta}}$ then $\tau = \phi(\infty) = \infty$. Moreover, we have that*

$$\mathcal{N}_G = \left\{ \psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \liminf_{u \rightarrow \infty} \frac{\overline{\Pi}(\frac{1}{u})}{\psi(u)} > 0 \right\} \subseteq \mathcal{N}_{\underline{\Theta}},$$

where if $\sigma^2 = 0$ we have the following equivalent criterion in terms of the characteristic triplet

$$(1.16) \quad \liminf_{u \rightarrow \infty} \frac{\overline{\Pi}(\frac{1}{u})}{\psi(u)} > 0 \iff \lim_{y \rightarrow 0} \frac{y^2 \overline{\Pi}(y)}{\int_0^y \overline{\Pi}(r) dr + my} > 0.$$

(2) *We also have*

$$(1.17) \quad \psi \in \mathcal{N}_\alpha \iff \overline{\Pi}(y) \stackrel{0}{\sim} \Gamma(\alpha + 1) C_\alpha y^{-\alpha} \iff \overline{\Pi}(y) \stackrel{0}{\sim} \Gamma(\alpha) C_\alpha y^{-\alpha-1},$$

where we recall that $C_\alpha > 0$.

(3) (a) $\mathcal{N}_P \subset \bigcap_{\alpha \in (0,1)} \bigcup_{\mathbf{m} \geq 1 - \frac{1}{\alpha}} \mathcal{N}_{\alpha, \mathbf{m}}$.

(b) *Let us write for $(\alpha, \mathbf{m}) \in \mathfrak{R} = \{(\alpha, \mathbf{m}); \alpha \in (0, 1] \text{ and } \mathbf{m} \geq 1 - \frac{1}{\alpha}\}$,*

$$U_{\alpha, \mathbf{m}}(y) = \frac{1}{\Gamma(\alpha + 1)} e^{-\mathbf{m}_\alpha y} \left(1 - e^{-\frac{y}{\alpha}}\right)^{\alpha-1}, \quad y > 0,$$

$\mathbf{m}_\alpha = (\alpha \mathbf{m} + 1 - \alpha)/\alpha \geq 0$. *Then, if $\psi \in \mathcal{N} \setminus \mathcal{N}_P$ with*

$$\sup_{y > 0} \inf_{A \in (0,1)} \left(\frac{\overline{\Pi}(y) + m}{\overline{\Pi}((1-A)y) + m} + \frac{U_{\alpha, \mathbf{m}}(y)}{U_{\alpha, \mathbf{m}}(Ay)} \right) \leq 1$$

for some $\alpha \in (0, 1), \mathbf{m} \geq 1 - \frac{1}{\alpha}$, then $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$.

(c) *For any $(\alpha, \mathbf{m}) \in \mathfrak{R}$, $\mathcal{N}_{\alpha, \mathbf{m}} \subset \mathcal{N}_{\underline{\Theta}}$ with when $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$, $\frac{\pi}{2}\alpha \leq \underline{\Theta}_\psi \leq \frac{\pi}{2}$.*

The proofs of the first item (1) are given in Section 6.7. The second one follows from augmentation of the classical estimate stated in Proposition 4.1.1 (7) to $\psi(u) = u\phi(u) \stackrel{\infty}{\sim} u^2 \int_0^{\frac{1}{u}} \overline{\Pi}(y) dy + mu$ in the regularly varying case, combined with a standard Tauberian theorem, i.e. $\int_0^y \frac{l(r)}{r^\alpha} dr \stackrel{0}{\sim} l(y)y^{\alpha-1}$, and the monotone density theorem. Items (3a) and (3b) are proved in Section 4.2, whereas item (3c) is part of the claim of Lemma 10.1.4

REMARK 1.2.2. Note that, in fact, under the conditions of the item (3b) we shall prove that the mapping $\Phi_{\alpha, \mathbf{m}}(u) = \frac{\psi(u)}{u\phi_{\alpha, \mathbf{m}}(u)} \in \mathcal{B}$, see (4.1) for the definition of general Bernstein functions, which is thanks to Proposition 4.1.1(4), a stronger statement than the requirement that $u \mapsto \frac{1}{\Phi_{\alpha, \mathbf{m}}(u)}$ is completely monotone.

1.3. Eigenvalue expansion and regularity of the gL semigroups

In order to state the next result regarding the spectral expansion and the regularity of the gL semigroups, we now introduce the following set which is the union of five sets of the form, for some linear space $L \subseteq L^2(\nu)$, set $\underline{\mathcal{N}} \subseteq \mathcal{N}$ and some $T > 0$ which may depend on $\psi \in \underline{\mathcal{N}}$,

$$\mathcal{D}_T^{\underline{\mathcal{N}}}(L) = \{(\psi, f, t); \psi \in \underline{\mathcal{N}}, f \in L, t > T = T(\psi) > 0\}$$

which shall serve as the domains of the spectral operator

$$\mathcal{D} = \mathcal{D}_0^{\mathcal{N}_P}(L^2(\nu)) \cup \mathcal{D}_{T_{\Theta_\phi}}^{\mathcal{N}_{\Theta_\phi}}(L^2(\mathbf{p}_\alpha)) \cup \mathcal{D}_{T_{\pi_\alpha}}^{\mathcal{N}_{\pi_\alpha}}(L^2(\bar{\nu}_\gamma)) \cup \mathcal{D}_{T_{\pi_\alpha, \rho_\alpha}}^{\mathcal{N}_{\pi_\alpha, \rho_\alpha}}(L^2(\nu)) \cup \mathcal{D}_{T_{\bar{\alpha}}}^{\mathcal{N}_{\bar{\alpha}}}(L^2(\nu))$$

where, for any fixed $\alpha \in (0, 1)$ and any fixed $\gamma > \alpha + 1$, we set

$$\mathbf{p}_\alpha(x) = x^{-\alpha}, \text{ and } \bar{\nu}_\gamma(x) = \max\left(\nu(x), e^{-x^{\frac{1}{\gamma}}}\right), x > 0,$$

and,

$$\begin{aligned} T_{\Theta_\phi} &= -\ln \sin \Theta_\phi, & T_{\pi_\alpha} &= -\ln \sin\left(\frac{\pi}{2}\alpha\right), \\ T_{\pi_\alpha, \rho_\alpha} &= \max\left(T_{\pi_\alpha}, 1 + \frac{1}{\rho_\alpha}\right), & T_{\bar{\alpha}} &= -\ln(2^{\bar{\alpha}} - 1), \end{aligned}$$

where, for each $\psi \in \mathcal{N}_R$, $\bar{\alpha} = \sup\{0 < \alpha \leq 1; \exists \mathbf{m} > \frac{1}{\alpha} - 1 \text{ and } \psi \in \mathcal{N}_{\alpha, \mathbf{m}}\}$ and ρ_α is the largest solution to the equation $(1 - \rho)^{\frac{1}{\alpha}} \cos\left(\frac{\arcsin(\rho)}{\alpha}\right) = \frac{1}{2}$. We mention that the notions introduced above and the ones entering in the next claim are reviewed and discussed at length in Chapters 4 and 5. It is also part of the ensuing statements that all functions are well defined. Next, we set, $W_\phi(1) = 1$ and, for $n \in \mathbb{N}$,

$$(1.18) \quad W_\phi(n+1) = \prod_{k=1}^n \phi(k),$$

and recall that $\mathfrak{r} = \infty$ if $\sigma^2 > 0$, and, $0 < \mathfrak{r} = \bar{\Pi}(0^+) + m \leq \infty$ otherwise, and, $N_{\mathfrak{r}} = \left\lceil \frac{\bar{\Pi}(0^+)}{\mathfrak{r}} \right\rceil - 1 \in [0, \infty]$ is the index of smoothness. Further, we set the quantity

$$(1.19) \quad d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\}$$

and note that if $\phi(0) = 0$ then necessarily $d_\phi = 0$ and always $d_\phi \in (-\infty, 0]$, see Proposition 4.1.1(3) for a justification. Finally, for two Banach spaces H_1, H_2 we denote by $\mathbf{B}(H_1, H_2)$ for the space of bounded linear operators from H_1 to H_2 . If $H_2 = H_1$ then we use $\mathbf{B}(H_1)$. We are now ready to state the main result of this work.

THEOREM 1.3.1. *Let $\psi \in \mathcal{N}$. Then, $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, where $\varepsilon(x) = \varepsilon_0(x) = e^{-x}, x > 0$ (ε_0 is defined in Remark 1.0.3), and \mathcal{I}_ϕ has a dense range, i.e. $\overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$.*

(1) For any $(\psi, f, t) \in \mathcal{D} \cup \mathcal{D}_0^{\mathcal{N}_\infty}(\text{Ran}(\mathcal{I}_\phi))$, the following holds.

(a) We have

$$(1.20) \quad P_t f = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n \in L^2(\nu)$$

where, for all $n \geq 0$,

$$(1.21) \quad \mathcal{P}_n(x) = \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{W_\phi(k+1)} x^k \in L^2(\nu),$$

and,

$$(1.22) \quad \mathcal{V}_n(x) = \frac{\mathcal{R}^{(n)}\nu(x)}{\nu(x)} = \frac{1}{n!} \frac{(x^n \nu(x))^{(n)}}{\nu(x)} \in L^2(\nu).$$

(b) If in addition $\overline{\Pi}(0^+) < \infty$, i.e. $\mathfrak{r} < \infty$, then, for any $f \in \text{Ran}(\mathcal{I}_\phi)$, $(t, x) \mapsto P_t f(x) \in \mathcal{C}^\infty(\mathfrak{C}_\mathfrak{r})$, with $\mathfrak{C}_\mathfrak{r} = \{(t, x) \in \mathbb{R}_+^2; x \leq \mathfrak{r}t\}$.

(c) Otherwise, $(t, x) \mapsto P_t f(x) \in \mathcal{C}^\infty((T, \infty) \times \mathbb{R}^+)$

(d) In both cases, for any non-negative integers k and p , we have

$$(1.23) \quad \frac{d^k}{dt^k} (P_t f)^{(p)}(x) = (-1)^k \sum_{n=p}^{\infty} n^k e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n^{(p)}(x),$$

where the series converges locally uniformly in (t, x) on the sets specified in (1b) and (1c).

(2) If $\psi \in \mathcal{N}_\infty^c$, then, we have for any $f \in \text{Ran}(\mathcal{I}_\phi)$, $(t, x) \mapsto P_t f(x) \in \mathcal{C}^\infty(\mathfrak{C}_\mathfrak{r})$. Moreover, we have, if $N_\mathfrak{r} \geq 0$ (resp. $N_\mathfrak{r} \geq 1$), that for any $0 \leq n < \frac{\overline{\Pi}(0^+)}{2\mathfrak{r}}$ (resp. $n > \frac{\overline{\Pi}(0^+)}{2\mathfrak{r}}$)

$$\mathcal{V}_n \in L^2(\nu) \quad (\text{resp. } \notin L^2(\nu)),$$

and, regardless of the value of $N_\mathfrak{r}$, for any $f \in \text{Ran}(\mathcal{I}_\phi)$ and $t \geq 0$,

$$(1.24) \quad P_t f = \sum_{n=0}^{\infty} e^{-nt} \langle \mathcal{I}_\phi^\dagger f, \mathcal{L}_n \rangle_\varepsilon \mathcal{P}_n \quad \text{in } L^2(\nu),$$

where \mathcal{I}_ϕ^\dagger is the pseudo-inverse of \mathcal{I}_ϕ , see [BIG02, p.234] for definition.

(3) **The heat kernel.** Let $\psi \in \mathcal{N}_{\underline{\Theta}}$ (resp. $\psi \in \mathcal{N}_R$). Then, for all $t > T_{\underline{\Theta}_\phi} = -\ln \sin \underline{\Theta}_\phi$ (resp. $t > \underline{T} = T_{\underline{\alpha}} \wedge T_{\underline{\Theta}_\phi}$) the heat kernel is absolutely continuous with a density $(t, x, y) \mapsto P_t(x, y) \in \mathcal{C}^\infty((T_{\underline{\Theta}_\phi}, \infty) \times \mathbb{R}_+^2)$ (resp. $\mathcal{C}^\infty((\underline{T}, \infty) \times \mathbb{R}_+^2)$) and, for any non-negative integers k, p, q ,

$$(1.25) \quad \frac{d^k}{dt^k} P_t^{(p,q)}(x, y) = \sum_{n=p}^{\infty} (-n)^k e^{-nt} \mathcal{P}_n^{(p)}(x) (\mathcal{V}_n(y) \nu(y))^{(q)},$$

where the series is locally uniformly convergent in $(t, x, y) \in (T_{\underline{\Theta}_\phi}, \infty) \times \mathbb{R}_+^2$ (resp. $(\underline{T}, \infty) \times \mathbb{R}_+^2$). If in addition $\nu \in \mathcal{A}(\frac{\pi}{2})$, that is ν is analytical on $\mathbb{C}(\frac{\pi}{2}) = \{z \in \mathbb{C}; |\arg z| < \frac{\pi}{2}\}$, then the expansion (1.25) is locally uniformly convergent even on $t > 0$. In particular, this is the case when $\psi \in \mathcal{N}_P$ or

$$\lim_{n \rightarrow \infty} \frac{\ln \left| \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{W_\phi(k+1)}{\Gamma(k+1)} \right|}{2\sqrt{n}} = -\infty.$$

This Theorem is proved in Chapter 12.

REMARK 1.3.2. This main result suggests some interesting and substantial differences between the spectral expansion of NSA and self-adjoint operators. The phenomenon that, for some classes, the expansion in the full Hilbert space holds only for t bigger than a positive constant has been observed in the framework of Schrödinger operator, see [Dav99], and is natural for non-normal operators. Indeed, in such a case, the spectral projections

$$(1.26) \quad P_n f = \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n$$

are not uniformly bounded as a sequence of operators. The projections are not orthogonal anymore and the sequence of eigenfunctions does not form a Riesz basis of the Hilbert space, see [Chr03] for definition and also the section 2.1 below. When compared to the spectral resolution of compact self-adjoint operators whose set of eigenfunctions forms an orthonormal basis of the Hilbert space, our study reveals that this requirement on the invariant subspace is, in general, too stringent. Indeed, even the range of the non-self adjoint operator, which is a linear subspace of the Hilbert space, can not be expanded into the invariant subspaces. These facts illustrate fundamental differences with the spectral reduction of self-adjoint Markov semigroups.

REMARK 1.3.3. In line with the previous remark, it is worth pointing out that the smoothness of the density of the heat kernel (or of $P_t f$, the solution to the associated Cauchy problem) remains the same for all t bigger than a constant. This fact may reveal that the eigenvalues expansion of the type (1.25) may be uniformly convergent only after the threshold time when the smoothness is established. Indeed, there are interesting and various examples of transition densities of Markov processes (with jumps) which are continuous of order $c(t)$ for some increasing functions c , see e.g. [Pic96].

REMARK 1.3.4. It is quite remarkable that our analysis allows to provide necessary and sufficient conditions for the existence of a sequence of co-eigenfunctions (at least for any $\psi \in \mathcal{N}_\infty$), a fact which requires not only smoothness properties of the invariant density but also very precise information about its small and large asymptotic behaviour along with its successive derivatives. Indeed, it is always difficult for non-local operators whose Lévy kernels span the whole set of Lévy measures to extract asymptotic properties beyond the classical regularly varying framework, or, the so-called stable-like case. We also indicate that the condition $\psi \in \mathcal{N}_\infty$ has a nice path interpretation since it means that the associated generalized Laguerre process has paths either of infinite variation when $\sigma^2 > 0$ or $\bar{\Pi}(0^+) = \infty$, or, when $\sigma^2 = 0$, of infinite activity in the sense that there are infinitely many jumps in any compact time interval.

REMARK 1.3.5. We mention that the fascinating and powerful techniques, such as Malliavin calculus, Hörmander analysis and PIDE techniques, see e.g. [IK06], [Pic96], [BRS14] and [CS09], that have proved successful for studying the smoothness properties of the transition kernel of diffusions semigroups or some Lévy type semigroups, are not general enough to be applied in our context. This is due to either a lack of symmetry and/or non-homogeneity of the Lévy kernel, or, unboundedness of the drift coefficient, or, the possible absence of diffusion part, or, simply the non-local feature of the generators. Our main results reveal that, in the context of

non-local Markov semigroups, the spectral expansion is a more flexible approach to derive these delicate regularity properties. Of course, our approach goes much beyond this issue as we also manage to obtain, for f in some various linear spaces, the smoothness of $P_t f$, that is the solution to the Cauchy problem. It would be interesting to characterize for each operator what is the maximal domain for which the stated regularity properties hold.

REMARK 1.3.6. We also observe from the space-time relationship between the self-similar semigroup $(K_t)_{t \geq 0}$ and the gL semigroup given in (2) of Definition 1.0.1, that, under the condition of the item (3) above, K has an absolutely continuous (smooth) kernel, $K_t(x, y)$, given by

$$K_t(x, y) = \sum_{n=0}^{\infty} (1+t)^{-n-1} \mathcal{P}_n(x) \mathcal{V}_n \left(\frac{y}{1+t} \right) \nu \left(\frac{y}{1+t} \right)$$

where the series is locally uniformly convergent in $(t, x, y) \in (e^T - 1, \infty) \times \mathbb{R}_+^2$, where T is either $T_{\underline{\Theta}_\phi}$ or \underline{T} .

The next result furnishes expansions for the adjoint semi-group P^* .

THEOREM 1.3.7. (1) If $\psi \in \mathcal{N}_{\underline{\Theta}}$ (resp. $\psi \in \mathcal{N}_R$), then for all $t > T_{\underline{\Theta}_\phi} = -\ln(\sin(\underline{\Theta}_\phi))$ (resp. $t > \underline{T} = T_{\bar{\alpha}} \wedge T_{\underline{\Theta}_\phi}$) and any $g \in L^2(\nu)$,

$$(1.27) \quad P_t^* g(y) = \sum_{n=0}^{\infty} e^{-nt} \langle g, \mathcal{P}_n \rangle_\nu \mathcal{V}_n(y),$$

where the series converges locally uniformly in $y > 0$.

(2) If $\psi \in \mathcal{N}_P$ with $0 < \bar{\Pi}(0^+) < \infty$, then for all $g \in L^2(\nu)$ and $t > 0$ (1.27) holds in $L^2(\nu)$.

This Theorem is proved in Chapter 12.

1.4. Convergence to equilibrium

There is a substantial and fascinating literature devoted to the study of the convergence to equilibrium of semigroups associated mainly to differential operators. In this framework, this problem has been investigated a lot under various coercive assumptions on the generator, such as spectral gap or logarithmic Sobolev inequalities, especially in the self-adjoint framework, see [BGL14] and the references therein. In this direction, we mention that recently Miclo [Mic15] has shown that a self-adjoint ergodic and hyperbounded Markov operator admits a spectral gap. Nevertheless, some recent works have identified asymptotic exponential convergence to equilibrium with bounds of the form $\mathfrak{m} e^{-ct}$, $\mathfrak{m}, c > 0$ and $t \geq 0$, when the generator satisfies some hypoelliptic type conditions. This phenomenon has been called hypocoercivity in Villani [Vil09], and, has recently attracted more and more attention, see e.g. Desvillettes and Villani [DV01], Eckmann and Hairer [EH03], Gadat and Miclo [GM13], Baudoin [Bau17] and Dolbeaut et al. [DMS15]. Note that in this literature the constants above are not necessarily optimal and are in general difficult to identify. At this stage, it is worth pointing out that among

the different rates of convergence that we obtain for the entire class of gL semi-groups, we also observe, see (1.28) below, for the (small) perturbation class the aforementioned hypocoercivity behaviour and the spectral decomposition enables us to explain it in terms of the spectral gap and projection norms which is a natural extension to the classical spectral gap observed in the self-adjoint case. We refer to the remarks following the next statement for further discussion on this and on other identified rates.

THEOREM 1.4.1. (1) *Let $\psi \in \mathcal{N}$. Then, for any $f \in \text{Ran}(\mathcal{I}_\phi) \subsetneq \mathbf{L}^2(\nu)$, we have for any $t > 0$,*

$$\|P_t f - \nu f\|_\nu \leq C_{\mathfrak{f},f}(\psi) e^{-t} \|f - \nu f\|_\nu,$$

where $\mathfrak{f} \in \mathbf{L}^2(\varepsilon)$ is such that $f = \mathcal{I}_\phi \mathfrak{f}$ and

$$C_{\mathfrak{f},f}(\psi) = \frac{\|\mathcal{I}_\phi\| \|\mathfrak{f} - \varepsilon \mathfrak{f}\|_\varepsilon}{\|f - \nu f\|_\nu} \geq 1$$

$$\text{with } \|\mathcal{I}_\phi\| = \sup_{g \in \mathbf{L}^2(\varepsilon), g \neq 0} \frac{\|\mathcal{I}_\phi g\|_\nu}{\|g\|_\varepsilon}.$$

(2) *For any $(\psi, f, t) \in \mathcal{D}$, there exist $C_L > 0$ and an integer $k \geq 0$ such that for any $t > T$,*

$$\|P_t f - \nu f\|_\nu \leq C_L 2^{\frac{k}{2}} \sqrt{\left(\frac{1}{e^{2(t-T)} - 1}\right)^{(k)}} \|f - \nu f\|_{\mathbf{L}},$$

where $\|\cdot\|_{\mathbf{L}}$ stands for the Hilbert space norm of

- $\mathbf{L} = \mathbf{L}^2(\mathbf{p}_\alpha)$ if $(\psi, f, t) \in \mathcal{D}_{T_{\Theta_\phi}}^{\mathcal{N}_{\Theta_\phi}}(\mathbf{L}^2(\mathbf{p}_\alpha))$;
- $\mathbf{L} = \mathbf{L}^2(\bar{\nu}_\gamma)$ if $(\psi, f, t) \in \mathcal{D}_{T_{\pi_\alpha}}^{\mathcal{N}_\alpha}(\mathbf{L}^2(\bar{\nu}_\gamma))$;
- $\mathbf{L} = \mathbf{L}^2(\nu)$ otherwise.

When $\psi \in \mathcal{N}_R$ then $k = 0$ and $\mathbf{L} = \mathbf{L}^2(\nu)$.

(3) *Finally, let $\psi \in \mathcal{N}_P$. Then, for any $\varepsilon > 0$, there exists $C_\varepsilon > 0$ such that for any $f \in \mathbf{L}^2(\nu)$ and $t > \varepsilon > 0$, we have that*

$$\|P_t f - \nu f\|_\nu \leq \frac{C_\varepsilon}{\sqrt{e^{2t} - 1}} \|f - \nu f\|_\nu.$$

Moreover, if $0 < \bar{\Pi}(0^+) < \infty$, then

$$\underline{\mathbf{m}} = \frac{m + \bar{\Pi}(0^+)}{\sigma^2} > \mathbf{d}_\varepsilon^+ = -(d_\phi + \varepsilon) \mathbb{I}_{\{-d_\phi > 0\}} \geq 0$$

for some $\varepsilon > 0$, and, for any $t > 0$ and any $f \in \mathbf{L}^2(\nu)$, the following bound

$$(1.28) \quad \|P_t f - \nu f\|_\nu \leq \sqrt{\frac{\underline{\mathbf{m}} + 1}{\mathbf{d}_\varepsilon^+ + 1}} e^{-t} \|f - \nu f\|_\nu$$

holds true with the inequality

$$\frac{\underline{\mathbf{m}} + 1}{\mathbf{d}_\varepsilon^+ + 1} \geq 1 + \frac{\bar{\Pi}(0^+)}{m + \sigma^2 + \varepsilon} \mathbb{I}_{\{-d_\phi > 0\}} + \frac{m + \bar{\Pi}(0^+)}{\sigma^2} \mathbb{I}_{\{\mathbf{d}_\varepsilon^+ = d_\phi = 0\}}.$$

REMARK 1.4.2. At this stage, it is worth pointing out that we also observe in (1.28), for the (small) perturbation class the aforementioned hypocoercivity behaviour. The intertwining approach that we develop enables us to characterize and interpret the role played by the two constants: the rate of decay e^{-t} corresponds to the second largest eigenvalue, as in the classical self-adjoint setting, whereas the second constant $\sqrt{\frac{m+1}{d_T^2+1}}$ may be interpreted as a measure of the quality of the decomposition of the Hilbert space in terms of the (co-)invariant subspaces naturally corresponding to the (co-)eigenfunctions. A particular justification for this is the estimate of the norm of the spectral projections as defined in (1.26), whose growth is in fact bounded by the order of the norms of the eigenfunctions and co-eigenfunctions, that is by $\|\mathcal{P}_n\|_\nu \|\mathcal{V}_n\|_\nu = O\left(\frac{n^m}{n^{d_T^2}}\right)$. The rate of growth of these norms seems to measure the departure from the classical orthonormal basis. We point out that we also identify in Theorem 9.0.1(2b), for some instances encompassing the non-perturbation case, an exponential rate of growth for the spectral projection norms of the adjoint operators. We refer to the recent papers [MSV17] where these rates have also been observed in the context of Schrödinger operators.

REMARK 1.4.3. In the general perturbation case, see Theorem 1.3.1 (3), the quality of the decomposition worsens and non-classical bounds on the speed of convergence to equilibrium of the type $C_{\nu,\epsilon}(e^{2t} - 1)^{-1/2}$, $t > \epsilon > 0$, appear, for some constant $C_{\nu,\epsilon} > 0$ which could be computed from Chapter 10 and the proof of [PS17, Proposition 2.3].

REMARK 1.4.4. Our results reveal additional interesting phenomena for this rate of convergence to equilibrium. Indeed, we identify more complex structure which involves different non-equivalent topologies. Upon substitution of the expression of $C_{f,f}(\psi)$ the estimate in Theorem 1.3.1(1) offers a classical spectral gap estimate but for functions in the range of the intertwining operator and against the topology of the reference self-adjoint semigroup which is not equivalent to the topology of $L^2(\nu)$, whereas item (2), when the quality of the decomposition of the Hilbert space improves, shows $C_L(e^{2t} - 1)^{-1/2}$, $C_L > 0, t > T$, speed of convergence against specific topologies. Thus, a general conclusion could be drawn to the effect that eventually the quality of decomposition of the Hilbert space is lost (item (2) and (3)), where this quality is measured by the rate of growth of the norms of the co-eigenfunctions, see (10.1), which shows slowest growth in the perturbation scenario and, (9.5) and (10.3), which reveal faster growth when we are strictly beyond perturbation. In turn we observe that in the latter scenario the speed of growth increases whenever the rate of convergence to infinity for small y of the tail of the Lévy measure $\bar{\Pi}(y) = \int_y^\infty \Pi(dr)$, slows down. We emphasize that, regardless of the latter, item (1) guarantees a classical speed of convergence against the topology of the self-adjoint semigroup.

1.5. Hilbert sequences and spectrum

We now present some substantial properties of the set of (co)-eigenfunctions and study in details the spectrum of the gL semigroups. This will follow from some general ideas, discussed in Chapter 11, that establish new and interesting connections

between three different concepts: intertwining relation, Hilbert sequences arising in non-harmonic analysis and spectrum of non self-adjoint operators.

We start by introducing some definitions related to Hilbert sequences which can be found in the monograph of Young [You01, Chap. 1]. Two sequences $(P_n)_{n \geq 0}$ and $(V_n)_{n \geq 0}$ are said to be biorthogonal in $L^2(\nu)$ if for any $n, m \in \mathbb{N}$,

$$(1.29) \quad \langle P_n, V_m \rangle_\nu = \delta_{nm}.$$

Moreover, a sequence that admits a biorthogonal sequence will be called *minimal* and a sequence that is both minimal and complete, in the sense that its linear span is dense in $L^2(\nu)$, will be called *exact*. It is easy to show that a sequence $(P_n)_{n \geq 0}$ is minimal if and only if none of its elements can be approximated by linear combinations of the others. If this is the case, then a biorthogonal sequence will be uniquely determined if and only if $(P_n)_{n \geq 0}$ is complete.

Next, we present some basic notions related to the concept of frames in Hilbert spaces. We point out that this generalization of orthogonal sequences has been introduced by Duffin and Schaeffer [DS52] in 1952 to study some deep problems in non-harmonic Fourier series and after the fundamental paper [DGM86] by Daubechies, Grossman and Meyer, frame theory began to be widely used, particularly in the more specialized context of wavelet frames and Gabor frames. A recent and thorough account on these Hilbert space sequences can be found in the book of Christensen [Chr03].

A sequence $(P_n)_{n \geq 0}$ in the Hilbert space $L^2(\nu)$ is a frame if there exist $A, B > 0$ such that the frame inequalities

$$(1.30) \quad A \|f\|_\nu^2 \leq \sum_{n=0}^{\infty} |\langle f, P_n \rangle_\nu|^2 \leq B \|f\|_\nu^2$$

hold, for all $f \in L^2(\nu)$. If only the upper bound exists, $(P_n)_{n \geq 0}$ is called a Bessel sequence. A frame sequence is always complete in the Hilbert space and when it is minimal, it is called a Riesz sequence. The latter are very useful objects as they share substantial properties with orthonormal sequences. Indeed, a Riesz sequence always admits a unique biorthogonal sequence $(V_n)_{n \geq 0}$ which is also a Riesz sequence and both together form the so-called Riesz basis. Moreover, the expansion in terms of the Riesz basis of any element of the Hilbert space is unique and convergent in the topology of the norm. When $(P_n)_{n \geq 0}$ is a Bessel sequence, that is only the upper frame condition in (1.30) is satisfied, then the so-called synthesis operator, that is the linear operator $\mathcal{S} : \ell^2(\mathbb{N}) \rightarrow L^2(\nu)$ defined by

$$(1.31) \quad \mathcal{S} : \underline{c} = (c_n)_{n \geq 0} \mapsto \mathcal{S}(\underline{c}) = \sum_{n=0}^{\infty} c_n P_n$$

is a bounded operator with norm $\|\mathcal{S}\|_\nu \leq \sqrt{B}$, that is, the series is norm convergent for any sequence in $\ell^2(\mathbb{N})$. However, \mathcal{S} is not in principle onto as the $(P_n)_{n \geq 0}$ does not form in general a basis of the Hilbert space.

Finally, we say that a sequence $(V_n)_{n \geq 0}$ in $L^2(\nu)$ is a Riesz-Fischer sequence if there exists a constant $B > 0$ such that $B \sum_{n=0}^{\infty} |c_n|^2 \leq \|\sum_{n=0}^{\infty} c_n V_n\|_\nu^2$ for all finite scalar sequences $(c_n)_{n \geq 0}$.

We proceed by recalling a few definitions concerning the spectrum of linear operators and we refer to [DS71, XV.8] for a thorough account on these objects. First, a complex number $\lambda \in S(P)$, the spectrum of the linear operator $P \in \mathbf{B}(L^2(\nu))$, if $P - \lambda \mathbf{I}$ does not have an inverse in $L^2(\nu)$ with the following three distinctions:

- $\lambda \in S_p(P)$, the point spectrum, if $\text{Ker}(P - \lambda \mathbf{I}) \neq \{0\}$.

- $\lambda \in S_c(P)$, the continuous spectrum, if

$$\text{Ker}(P - \lambda \mathbf{I}) = \{0\} \text{ and } \overline{\text{Ran}}(P - \lambda \mathbf{I}) = L^2(\nu) \text{ but } \text{Ran}(P - \lambda \mathbf{I}) \subsetneq L^2(\nu).$$

- $\lambda \in S_r(P)$, the residual spectrum, if

$$\text{Ker}(P - \lambda \mathbf{I}) = \{0\} \text{ and } \overline{\text{Ran}}(P - \lambda \mathbf{I}) \subsetneq L^2(\nu).$$

Clearly, $S(P) = S_p(P) \cup S_c(P) \cup S_r(P)$. Let $\lambda \in S_p(P)$ be an isolated eigenvalue. Then its geometric multiplicity, denoted by $\mathfrak{M}_g(\lambda, P)$ is computed as follows

$$(1.32) \quad \mathfrak{M}_g(\lambda, P) = \dim \text{Ker}(P - \lambda I),$$

that is the dimension of the corresponding eigenspace. Its algebraic multiplicity, denoted by $\mathfrak{M}_a(\lambda, P)$ is defined by

$$(1.33) \quad \mathfrak{M}_a(\lambda, P) = \dim \bigcup_{k=1}^{\infty} \text{Ker}(P - \lambda I)^k.$$

Note that always $\mathfrak{M}_g(\lambda, P) \leq \mathfrak{M}_a(\lambda, P)$. We are now ready to state both interesting properties satisfied by the sequence of eigenfunctions and co-eigenfunctions when viewed as sequences in Hilbert space and information on the structure of the spectrum of gL semigroups. These claims are proved in Chapter 11.

THEOREM 1.5.1. (1) Let $\psi \in \mathcal{N}$. Then, $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$ and $(\mathcal{P}_n)_{n \geq 0}$ is a Bessel sequence but it is not a Riesz basis in $L^2(\nu)$.

(2) Let us now assume that $\psi \in \mathcal{N}_\infty$. Then, the sequence $(\mathcal{V}_n)_{n \geq 0}$ is a minimal Riesz-Fischer sequence in $L^2(\nu)$ and $(\mathcal{P}_n)_{n \geq 0}$ is exact. Moreover, $(\mathcal{V}_n)_{n \geq 0}$ is the unique sequence biorthogonal to $(\mathcal{P}_n)_{n \geq 0}$ in $L^2(\nu)$.

(3) Assume that $\psi \in \mathcal{N}_P \cup \mathcal{N}_{\alpha, \mathbf{m}}^{d_\phi}$ where

$$(1.34) \quad \mathcal{N}_{\alpha, \mathbf{m}}^{d_\phi} = \left\{ \psi \in \mathcal{N}_{\alpha, \mathbf{m}}; d_\phi < 1 - \frac{\mathbf{m}}{2} - \frac{1}{2\alpha} \right\}$$

Then, $\overline{\text{Span}}(\mathcal{V}_n) = L^2(\nu)$ and hence $(\mathcal{V}_n)_{n \geq 0}$ is also exact and $(\mathcal{P}_n)_{n \geq 0}$ is its unique biorthogonal sequence in $L^2(\nu)$.

(4) Let $t \geq 0$ and write

$$\mathcal{E}_t = \{e^{-nt}, n \geq 0\}.$$

(a) Then, for any $\psi \in \mathcal{N}$, we have

$$\mathcal{E}_t \subseteq S_p(P_t) \text{ and } S_r(P_t) = \emptyset.$$

(b) Moreover, if $\psi \in \mathcal{N}_\infty$, then

$$\mathcal{E}_t \subseteq S_p(P_t^*) \text{ with } \mathfrak{M}_a(e^{-nt}, P_t^*) = \mathfrak{M}_g(e^{-nt}, P_t^*) = 1.$$

(c) If $\psi \in \mathcal{N}_P \cup \mathcal{N}_{\alpha, \mathfrak{m}}^{d_\phi}$, where we recall that the latter class is defined in (1.34), then

$$\mathcal{E}_t = S_p(P_t) = S_p(P_t^*),$$

and

$$\mathfrak{M}_a(e^{-nt}, P_t) = \mathfrak{M}_a(e^{-nt}, P_t^*) = 1 \text{ and } S_r(P_t) = S_r(P_t^*) = \emptyset.$$

(d) Finally, if $\psi \in \mathcal{N}_\infty^c$ with $N_\tau \geq 1$, then

$$\left\{ e^{-nt}, 0 \leq n < \frac{\bar{\Pi}(0^+)}{2\mathfrak{r}} \right\} \subseteq S_p(P_t^*)$$

and

$$\left\{ e^{-nt}, n > \frac{\bar{\Pi}(0^+)}{2\mathfrak{r}} \right\} \subseteq S_r(P_t^*).$$

This Theorem is proved in Chapter 12.

REMARK 1.5.2. Although we are able to provide precise information regarding the spectrum of the gL semigroups, the characterization of their full spectrum seems to be a delicate issue. We refer to Chapter 11 for interesting and detailed discussions regarding the spectrum of operators linked by an intertwining operator.

1.6. Plan of the paper

We now describe the contents of the remaining parts of the paper whose main purpose is merely to prove Theorem 1.3.1, that is to establish the statements related to the spectral expansions of gL semigroups. The length of the paper may be explained by the complexity of the problem that forces us to develop adequate mathematical tools almost from scratch. We emphasize that these new results may be of independent interests in a variety of contexts such as probability theory, the study of functional equations, asymptotic analysis, non-harmonic analysis, special functions and functional analysis. More specifically, we have the following organization.

- **Strategy of proof and auxiliary results.** Chapter 2 is a high-level description of our methodology whose main comprehensive ideas could be used to study the spectral representation of more general Markov semigroups. We also state a short version of the substantial results that are required to prove the main results. This part also includes the proof of Theorem 1.1.1(1) along with some preliminaries results on gL semigroups that are both useful for the remaining part of the paper and also easy to prove by merely adapting some results that can be found in the literature.
- **Examples.** Chapter 3 is devoted to the description of some specific examples which illustrate the variety of results we obtained ranging from the self-adjoint class, the perturbation cases to the pure compound Poisson case.
- **New developments in the theory of Bernstein functions.** In Chapter 4, we present some known and new results regarding the convex cone of Bernstein functions which are essential to develop a fine and detailed

analysis of gL semigroups. The new claims vary from new asymptotic estimates on the complex and (positive) real lines, to a new mapping leaving some subset of Bernstein functions invariant. Since this set of functions is central in a variety of contexts, ranging from potential theory, probability theory, operator theory to complex analysis, we gather and prove these results in one chapter.

- **Fine properties of the density of the invariant measure.** We start Chapter 5 by providing a simple but useful mapping relating the class of invariant densities to a subset of the substantial and well-studied class of positive self-decomposable variables. This connection enables us to use the work of Sato and Yamazato [SY78] to derive some fine distributional properties of the invariant densities. We proceed to deepen their study by providing the small exponential asymptotic decay of the densities together with their successive derivatives by showing that they satisfy the several and delicate conditions of a non-classical Tauberian theorem which was originally proved by Baalkema et al. [BKS93] and that we extend to fit to our framework. Moreover, by resorting to the theory of excursions of Lévy processes, we derive very precise bounds for the large asymptotic behaviour of these densities.
- **Bernstein-Weierstrass products and Mellin transforms.** In Chapter 6, we carry out an in-depth study of the functional equation $\mathcal{M}_{V_\psi}(z+1) = \phi(z)\mathcal{M}_{V_\psi}(z)$ valid in the right-half plane that is satisfied by the Mellin transform of the invariant measure. This part complements Webster's fascinating investigation in [Web97] of similar functional equations on the positive real line. We show that the solution of this equation is given by a class of generalized Weierstrass products in bijection with the convex cone of Bernstein functions. Then the study includes the analytical extension of the solution to the right-half plane as well as a description of its precise asymptotic behaviour along imaginary lines. The new special functions given by the generalized Weierstrass products and included in this study encompass many special functions that have appeared in different contexts in the literature and our approach provides a unified framework to their study and a common set of quantities describing their properties. We note that whilst this paper was reviewed we studied in even more detail the aforementioned special functions which we call now Bernstein-gamma functions. For more information, see [PS18a].
- **Intertwining relations and a set of eigenfunctions.** In Chapter 7, we first develop a factorization of Markov operators which turns out to be essential in the derivation of the intertwining relation between the gL semigroups and the classical Laguerre one. It also plays an important role in proving the continuity property of the intertwining operator between appropriate non-trivial weighted Hilbert spaces. From the intertwining relation we construct a sequence of polynomials which corresponds to a set of eigenfunctions for the gL semigroups. We also prove that the latter forms a Bessel sequence in $L^2(\nu)$, an object which has been introduced in non-harmonic analysis as a generalization of the concept of a basis in

Hilbert space. We also show that the polynomials are the Jensen polynomials of some entire function and exploit this connection to derive an upper bound for large orders of these polynomials in terms of the maximum modulus of the associated function.

- **Co-eigenfunctions: existence and characterization.** In Chapter 8, we first resort to the theory of distributions in the setting of the Mellin transform to characterize, in terms of the Rodrigues operator, the co-eigenfunctions. Using the precise information regarding the densities of the invariant measure obtained in the previous chapters, we provide (almost) necessary and sufficient conditions for the co-eigenfunctions to belong to the weighted Hilbert space $L^2(\nu)$.
- **Uniform and norms estimates of the co-eigenfunctions.** Chapter 9 contains the proof of two asymptotic estimates for large values of n for the norms of co-eigenfunctions considered in different topologies. The first one is based on a saddle-point approximation which applies for the class \mathcal{N}_{Θ} , see Table 2, whereas the second one relies on upper bounds for the co-eigenfunction that we derive by exploiting very precise information regarding the location of zeros of the successive derivatives of the invariant density.
- **The concept of reference semigroups: $L^2(\nu)$ -norm estimates and completeness of the set of co-eigenfunctions.** We develop in Chapter 10 the concept of reference semigroups. It consists in identifying gL semigroups \bar{P} which satisfy the following two criteria. First, their special structure permits to study their spectral reduction in detail. Furthermore, there exists a subclass of gL semigroups such that for each element in this class we have the adjoint intertwining relation $P_t^* \Lambda^* = \Lambda^* \bar{P}_t^*$ where Λ^* is the adjoint of a bounded operator between appropriate weighted Hilbert spaces. Although this approach may be extended to more general classes, we present in this part two different reference semigroups, which allow, in particular, to deal with the spectral expansion in the full Hilbert space of the perturbation class, that is when $\psi \in \mathcal{N}_P$. We describe two important applications of the reference concept regarding probably the two most technical difficulties of this work, namely the estimates of the $L^2(\nu)$ norm of the sequence of co-eigenfunctions \mathcal{V}_n and their completeness in $L^2(\nu)$.
- **Hilbert sequences, intertwining and spectrum.** In this Chapter, we gather new developments in the study of the spectrum of linear operators linked by an intertwining relation. In particular, we establish some interesting and new connections between the concepts of intertwining relation, Hilbert sequences arising in non-harmonic analysis and spectrum of non self-adjoint operators.
- **Proof of Theorems 1.3.1, 1.3.7 and 1.4.1.** The last Chapter contains the last arguments required to complete the lengthy proofs of Theorems 1.3.1, 1.3.7 and 1.4.1.

1.7. Notation, conventions and general facts

1.7.1. Functional spaces. Throughout, we denote by $L^2(\mathbb{R}_+)$ the Hilbert space of square integrable Lebesgue measurable functions on \mathbb{R}_+ endowed with the inner product $\langle f, g \rangle = \int_0^\infty f(x)g(x)dx$ and the associated norm $\|\cdot\|$. For any weight function ν defined on \mathbb{R}_+ , i.e. a non-negative Lebesgue measurable function, we denote by $L^2(\nu)$ the weighted Hilbert space endowed with the inner product $\langle f, g \rangle_\nu = \int_0^\infty f(x)g(x)\nu(x)dx$ and its corresponding norm $\|\cdot\|_\nu$. Similarly, we use the standard notation for the functional spaces $L^p(\nu)$, for $p \in [1, \infty]$. We preserve $L^p(\mathbb{R})$ and $L^p(\mathbb{R}_+)$ for the case when $\nu \equiv 1$.

For any $E \subseteq \mathbb{R}$, we use $\mathcal{C}^k(E)$, for $k \geq 0$, $k \in \mathbb{N} \cup \{\infty\}$, for functions with k continuous derivatives on E and write simply $\mathcal{C}(E) = \mathcal{C}^0(E)$. Additionally, we denote by $\mathcal{C}_0^k(\mathbb{R})$ (resp. $\mathcal{C}_0^k(\mathbb{R}_+)$) the subspaces of functions of $\mathcal{C}^k(\mathbb{R})$ (resp. $\mathcal{C}^k(\mathbb{R}_+)$), all of whose k derivatives and the functions themselves vanish at infinity. We also consider the spaces $\mathcal{C}_b^k(\mathbb{R})$ (resp. $\mathcal{C}_b^k(\mathbb{R}_+)$), i.e. the space of functions with k continuous, bounded derivatives. We reserve $\mathcal{C}_c^\infty(\mathbb{R})$ and $\mathcal{C}_c^\infty(\mathbb{R}_+)$ for all infinitely differentiable functions with compact support in \mathbb{R} and \mathbb{R}_+ . Also $\mathcal{C}^2([-\infty, \infty])$ stands for the space of functions $f \in \mathcal{C}^2(\mathbb{R})$ which along with their first and second derivatives admit a limit at $\pm\infty$. Furthermore we denote by $\mathbf{B}_b(\mathbb{R})$ and $\mathbf{B}_b(\mathbb{R}_+)$ the corresponding spaces of bounded measurable functions.

We also write for Banach spaces H_1, H_2

$$\mathbf{B}(H_1, H_2) = \{L : H_1 \rightarrow H_2 \text{ linear and continuous mapping}\}.$$

In the case of one Banach space H , the unital Banach algebra $\mathbf{B}(H, H)$ is denoted by $\mathbf{B}(H)$.

1.7.2. Complex plane, strips and analytic functions. We use \mathbb{C} for the complex plane and \mathbb{C}_\pm for the half planes with $\Re(z) \geq 0$ and $\Re(z) \leq 0$. For any $-\infty \leq \underline{a} < \bar{a} \leq \infty$ we define vertical strips of the complex plane by

$$\mathbb{C}_{(\underline{a}, \bar{a})} = \{z \in \mathbb{C}; \underline{a} < \Re(z) < \bar{a}\}$$

and, for any $a \in \mathbb{R}$, we denote by $\mathbb{C}_a = \{z \in \mathbb{C}; \Re(z) = a\}$ the vertical complex lines.

For $-\infty \leq \underline{a} < \bar{a} \leq \infty$, denote by

$$\mathcal{A}_{(\underline{a}, \bar{a})} \text{ the set of analytic functions on } \mathbb{C}_{(\underline{a}, \bar{a})}$$

and by $\mathcal{A}_{[\underline{a}, \bar{a}]} \subset \mathcal{A}_{(\underline{a}, \bar{a})}$ the set of analytic functions which have a continuous extension on $\mathbb{C}_{\underline{a}}$. Similarly, we define $\mathcal{A}_{[\underline{a}, \bar{a}]}$ and $\mathcal{A}_{(\underline{a}, \bar{a}]}$. Finally, for any $\theta \in (0, \pi)$, we denote by

$$\mathcal{A}(\theta) \text{ the set of analytic functions on } \mathbb{C}(\theta) = \{z \in \mathbb{C}; |\arg z| < \theta\}.$$

1.7.3. Asymptotic behaviour. We use the following notation

$$\begin{aligned} f &\asymp g \text{ means that } \exists c > 0 \text{ such that } c \leq \frac{f}{g} \leq c^{-1}, \\ f &\stackrel{a}{\sim} g \text{ means that } \lim_{x \rightarrow a} \frac{f(x)}{g(x)} = 1, \text{ for some } a \in \mathbb{R} \cup \{\pm\infty\}, \\ f &\stackrel{a}{=} O(g) \text{ means that } \overline{\lim}_{x \rightarrow a} \left| \frac{f(x)}{g(x)} \right| < \infty, \\ f &\stackrel{a}{=} o(g) \text{ means that } \lim_{x \rightarrow a} \left| \frac{f(x)}{g(x)} \right| = 0. \end{aligned}$$

We say that a function l is slowly varying at $a \in \{0\} \cup \{\infty\}$ if, for any $y > 0$,

$$\lim_{x \rightarrow a} \frac{l(yx)}{l(x)} = 1.$$

We then say that $f \in RV_\alpha(a)$, i.e. f is regularly varying of index $\alpha \in \mathbb{R}$ at a , if

$$(1.35) \quad f(x) = x^\alpha l(x)$$

where l is slowly varying at a . Note that the class of slowly varying functions at a coincides with $RV_0(a)$. We refer to [BGT87] for a complete account of the theory of regularly varying functions.

1.7.4. Random variables and Mellin transforms. Recall that a random variable is a measurable mapping $X : \Omega \rightarrow \mathbb{R}$ from a probability space (Ω, \mathbb{P}) endowed with a σ -algebra to \mathbb{R} endowed with the Borel σ -algebra. In this work we mostly consider mappings restricted to \mathbb{R}_+ , i.e. $X : \Omega \rightarrow \mathbb{R}_+$. In the latter case the Mellin transform is introduced via

$$\mathcal{M}_X(z) = \mathbb{E} [X^{z-1}]$$

and it is well defined at least for $z \in \mathbb{C}_1$, i.e. when $z = 1 + ib$, $b \in \mathbb{R}$. If $\mathcal{M}_X(z)$ is defined, absolutely integrable and uniformly decaying to zero along the lines of the strip $z \in \mathbb{C}_{(\underline{a}, \bar{a})}$, for $\underline{a} < 1 < \bar{a}$, the Mellin inversion theorem applies and gives, for $\underline{a} < a < \bar{a}$ and $x > 0$,

$$(1.36) \quad v(x) = \frac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} x^{-z} \mathcal{M}_X(z) dz,$$

where, by dominated convergence, $v \in \mathbf{C}_0(\mathbb{R}_+)$ and it is the density of X , i.e. $\mathbb{P}(X \in dx) = v(x)dx$. We write $\mathcal{M}_v = \mathcal{M}_X$. If for any fixed $a \in (\underline{a}, \bar{a})$ and some $\theta \in (0, \pi]$,

$$(1.37) \quad |\mathcal{M}_v(a + ib)| = O\left(e^{-\theta|b|}\right)$$

then a dominated convergence argument applied to (1.36) yields that $v \in \mathcal{A}(\theta)$.

Moreover if, for some $n \in \mathbb{N}$, \mathcal{M}_v is well-defined in $\mathbb{C}_{(\underline{a}, \bar{a})}$, then $\mathcal{M}_{v^{(n)}}(z) = (-1)^n (z-n)_n \mathcal{M}_v(z-n)$ is well-defined in $\mathbb{C}_{(\underline{a}+n, \bar{a}+n)}$, where

$$(z-n)_n = \frac{\Gamma(z)}{\Gamma(z-n)} = (z-1) \dots (z-n+1),$$

see [ML86, Chap. 11, (11.7.6)]. If, in addition, $z \mapsto z^n \mathcal{M}_v(z - n)$, for some $n \in \mathbb{N}$, is absolutely integrable and uniformly decaying along the complex lines of the strip $\mathbb{C}_{(\underline{a}+n, \bar{a}+n)}$, then we have, for $\underline{a} + n < a < \bar{a} + n$ and $x > 0$,

$$(1.38) \quad v^{(n)}(x) = \frac{(-1)^n}{2\pi i} \int_{a-i\infty}^{a+i\infty} x^{-z} (z - n)_n \mathcal{M}_v(z - n) dz$$

and by dominated convergence again $v^{(n)} \in \mathcal{C}_0(\mathbb{R}_+)$. Finally, we also use the L^2 -theory of Mellin transform and in particular the Parseval identity which reads off, for $v \in L^2(\mathbb{R}_+)$, as

$$(1.39) \quad \|v\|^2 = \int_{-\infty}^{\infty} \left| \mathcal{M}_v \left(\frac{1}{2} + ib \right) \right|^2 \frac{db}{2\pi}.$$

We refer to the monographs [ML86] and [PK01] for excellent accounts on Mellin transforms.

Strategy of proofs and auxiliary results

The aim of this chapter is two-fold. On the one hand, its primary purposes are to describe, in a synthetic way, the methodology we have developed to obtain the eigenvalues expansions stated in Theorem 1.3.1, to consider the main ideas underpinning the proof of the main results and to discuss the technical difficulties we had to overcome. On the other hand, both for the reader's convenience and to give flavour on the type of results we have derived, we also state, in general without proof, several interesting supplementary results. They could be of independent interest in a variety of contexts at the interplay between probability theory, the study of functional equations, asymptotic analysis, non-harmonic analysis, special functions and functional analysis. We wish to point out at this stage that the remaining parts of the paper contain additional interesting but more specific claims regarding the objects we are dealing with. We also mention them in this part.

2.1. Outline of our methodology

1. Intertwining relations. The first main idea is to identify a commutation relation between the NSA gL semigroups and the semigroup of a self-adjoint operator which admits a known spectral resolution. This is achieved by establishing an intertwining relation between the class of gL semigroups and the classical Laguerre semigroup $Q = (Q_t)_{t \geq 0}$, a self-adjoint operator in $L^2(\varepsilon)$ with $\varepsilon(x) = e^{-x}$, $x > 0$, whose main properties are reviewed in Example 3.0.1 below. More precisely, after recalling that for any $\psi(u) = u\phi(u) \in \mathcal{N}$, we define, from 1.13, the Markov operator $\mathcal{I}_\phi f(x) = \mathbb{E}[f(xI_\phi)]$ where $I_\phi = \int_0^\infty e^{-tu} dt$, we have the following statement which is a short version of Theorem 7.0.1

THEOREM 2.1.1. *Let $\psi \in \mathcal{N}$. Then, $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ and in $L^2(\varepsilon)$,*

$$(2.1) \quad P_t \mathcal{I}_\phi f = \mathcal{I}_\phi Q_t f.$$

Moreover, \mathcal{I}_ϕ is not bounded from below.

We prove this relation in a couple of steps utilizing that (2.1) can be reduced to the proof of a special multiplicative factorization linking the invariant measures of the involved semigroups, see [CPY97, Proposition 3.2], and the injectivity on appropriate Hilbert/Banach spaces of the multiplicative operator \mathcal{I}_ϕ . First, we derive the desired multiplicative factorization and then, by showing that the Mellin transform of ν , the density of the invariant measure, is zero-free on the imaginary line we deduce the injectivity of \mathcal{I}_ϕ by means of a Wiener Tauberian theorem combined with a standard approximation argument. This zero-free property is extracted from a generalized Weierstrass product representation of this Mellin transform, see (6.1), which is characterized as a solution to a Gamma type functional equation of the

form $\mathcal{M}(z+1) = \phi(z)\mathcal{M}(z)$, where ϕ stands for the descending ladder height exponent associated to ψ , that is a Bernstein function, see (1.9). This result has been announced in the note [PS13b] and discussed at much more depth in [PS18a], which succeeds this work in this direction. We mention that, for instance, when considering the trivial Bernstein function $\phi(u) = u$, i.e. $\psi(u) = u^2$, this Weierstrass product boils down to the infinite product representation of the Gamma function itself. We also point out that the proof of $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, that is the continuity property between weighted Hilbert spaces is in general a difficult problem. By means of the Marcinkiewicz multiplier theorem for Mellin transform, see [Ro082], one can show, from the asymptotic behaviour of its Mellin multiplier, see Theorem 6.0.2 below, that a Markov operator is bounded from $L^2(\mathbf{p}_\alpha)$, where $\mathbf{p}_\alpha(x) = x^{-\alpha}, x, \alpha > 0$, into itself. One classical approach is to consider weights which belong to the so-called class of Muchkenbaupt, the conditions of which are not satisfied by ε . Instead, the multiplicative factorization of Markov operators that we establish, allows us to derive by a simple application of the Jensen inequality the contraction property of the Markov operator \mathcal{I}_ϕ .

The idea of intertwining relation between Markov semigroups is not new and can be traced back to the works of Dynkin [Dyn65] and Rogers and Pitman [PR81] which yield such relation between a Brownian motion in \mathbb{R}^n and its radial part, the Bessel process of dimension n . This device, which is always difficult to identify, has also been used by Diaconis and Fill [DF90] in relation with strong uniform times, by Carmona, Petit and Yor [CPY98] in relation to the so-called self-similar saw tooth-processes, and, more recently by Fill [Fil09] for an elegant characterization of the distribution of the first passage time of some Markov chains, by Borodin and Corwin [BC14] in the context of Macdonald processes, by Pal and Shkolnikov [PS13a] for linking diffusion operators, and, by Patie and Simon [PS12b] to relate classical fractional operators.

On the other hand, this type of commutation relation between linear operators has been also intensively studied in functional analysis in the context of differential operators. This approach culminated in the work of Delsarte and Lions [DL57] who showed the existence of a transmutation operator between differential operators of the same order and acting on the space of entire functions. The transmutation operator, which plays the role of the intertwining operator, is in fact an isomorphism on this space. This property is very useful for the spectral reduction of these operators since it allows to transfer the spectral objects. We mention that Delsarte and Lions's development has been intensively used in scattering theory and in the theory of special functions, see e.g. Carroll and Gilbert [CG81].

In the context of this paper, the situation is more delicate, since on the one hand, the operators are non-local, and, on the other hand, the intertwining operator \mathcal{I}_ϕ is not in general an isomorphism. Perhaps, that is the price to pay in order to relate a non-local and non-symmetric operator to a self-adjoint and local (differential) operator. It is also worth mentioning that intertwining relations go beyond perturbation theory of self-adjoint operators. Indeed, it may relate the self-adjoint classical Laguerre semigroup, i.e. $\sigma^2 > 0, m = 0, \Pi(\mathbb{R}) = 0$ to a NSA semigroup without a diffusion part, that is when $\sigma^2 = 0, m > 0, \Pi(\mathbb{R}) > 0$ in (1.6). As far as the authors know, the intertwining theory has not been exploited for dealing with the spectral representations of Markov semigroups, or more generally of NSA linear operators,

which is rather surprising, as it seems to be a promising and natural technique as the following lines hope to illustrate convincingly.

2. A set of eigenfunctions. From the intertwining identity (2.1), we easily get that, for each $n \in \mathbb{N}$, $\mathcal{P}_n = \mathcal{I}_\phi \mathcal{L}_n$, with \mathcal{L}_n the classical Laguerre polynomials defined in (3.4), satisfy

$$(2.2) \quad P_t \mathcal{P}_n = P_t \mathcal{I}_\phi \mathcal{L}_n = \mathcal{I}_\phi Q_t \mathcal{L}_n = e^{-nt} \mathcal{I}_\phi \mathcal{L}_n = e^{-nt} \mathcal{P}_n$$

and since $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ but it is not bounded from below, we obtain with the notions introduced in Section 1.5 the following result.

THEOREM 2.1.2. *Let $\psi \in \mathcal{N}$. $(\mathcal{P}_n)_{n \geq 0}$ is a Bessel sequence in $L^2(\nu)$ of eigenfunctions for P_t but it is not a Riesz basis in $L^2(\nu)$. We also have, for any $n \in \mathbb{N}$,*

$$(2.3) \quad \|\mathcal{P}_n\|_\nu = O(1) \text{ and } \left| \mathcal{P}_n^{(p)}(x) \right| = O\left(n^{p+\frac{1}{2}} e^{\frac{x}{\tau} n}\right)$$

where the last bound holds for large n , any integer p and $x > 0$ and we recall that $\tau = \infty$ if $\sigma^2 > 0$, and, $0 < \tau = \overline{\Pi}(0^+) + m \leq \infty$ otherwise.

A more detailed version of this Theorem is Theorem 7.0.3.

3. A first spectral expansion. From the intertwining identity (2.1) and the expansion (3.2) of the classical Laguerre semigroup, one gets, by means of the continuity property of \mathcal{I}_ϕ and of the synthesis operator \mathcal{S} defined in (1.31), the following first expansion of P .

THEOREM 2.1.3. *Let $\psi \in \mathcal{N}$. Then, for any $f \in \text{Ran}(\mathcal{I}_\phi)$ and $t \geq 0$,*

$$(2.4) \quad P_t f = \sum_{n=0}^{\infty} e^{-nt} c_n^\dagger(f) \mathcal{P}_n \quad \text{in } L^2(\nu),$$

where $c_n^\dagger(f) = \langle \mathcal{I}_\phi^\dagger f, \mathcal{L}_n \rangle_\varepsilon$ and we recall that \mathcal{I}_ϕ^\dagger is the pseudo-inverse of \mathcal{I}_ϕ .

4. Existence and characterization of a set of co-eigenfunctions. With the aim of extending the domain of the spectral operator, i.e. the linear operator appearing on the right-hand side of (2.4), we proceed by first investigating the existence and the characterization of a set of co-eigenfunctions, that is, eigenfunctions for the adjoint semigroup P^* . More specifically, we say that, for some $t > 0$ and $n \geq 0$, \mathcal{V}_n is a co-eigenfunction for P_t , or equivalently, an eigenfunction for its adjoint P_t^* in $L^2(\nu)$, associated to the eigenvalue e^{-nt} if $\mathcal{V}_n \in L^2(\nu)$ and $P_t^* \mathcal{V}_n = e^{-nt} \mathcal{V}_n$, which can be rephrased as, for any $f \in L^2(\nu)$,

$$\langle f, P_t^* \mathcal{V}_n \rangle_\nu = \langle P_t f, \mathcal{V}_n \rangle_\nu = e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu.$$

We state the following theorem which summarizes the results that are presented in Chapter 8.

THEOREM 2.1.4. *Let $\psi \in \mathcal{N}$ and denote by \mathcal{I}_ϕ^* the adjoint of \mathcal{I}_ϕ . Then $\mathcal{I}_\phi^* \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ and in $L^2(\nu)$, for any $t \geq 0$,*

$$(2.5) \quad \mathcal{I}_\phi^* P_t^* f = Q_t \mathcal{I}_\phi^* f.$$

Moreover, for any $n \in \mathbb{N}$, the equation in $L^2(\nu)$

$$(2.6) \quad \mathcal{I}_\phi^* g = \mathcal{L}_n$$

has a unique solution (resp. no solution) given by

$$\mathcal{V}_n(x) = \frac{\mathcal{R}^{(n)}\nu(x)}{\nu(x)} = \frac{(x^n\nu(x))^{(n)}}{n!\nu(x)} = \frac{w_n(x)}{\nu(x)}$$

if $0 \leq n < \frac{\bar{\Pi}(0^+)}{2\tau}$ (resp. if $1 \leq N_\tau = \left\lceil \frac{\bar{\Pi}(0^+)}{\tau} \right\rceil - 1 < \infty$ and $n > \frac{\bar{\Pi}(0^+)}{2\tau}$), where by convention $\frac{\bar{\Pi}(0^+)}{2\tau} = \infty$ provided $\bar{\Pi}(0^+) = \infty$.

First, we observe that if \mathcal{V}_n is solution of (2.6) then, from the intertwining relation (2.5) it is indeed a co-eigenfunction of P , or an eigenfunction for its adjoint semigroup. Note that when such a solution does not exit for some n then e^{-nt} , the corresponding eigenvalue for P_t belongs to the residual spectrum of P_t^* .

The questions of existence and characterization of the solution to the equation (2.6) is extremely delicate as it requires the description of the range of \mathcal{I}_ϕ^* and of its unbounded inverse and also to deal with the complex structure of the weighted Hilbert space. As opposed to the self-adjoint framework where such a question is obvious, there do not seem to exist any results in the literature ensuring the (non)-existence of this set of co-eigenfunctions. To overcome this difficulty, we implement the following two-steps program. First, by considering the formal adjoint of \mathcal{I}_ϕ in $L^2(\mathbb{R}_+)$, we transfer equation (2.6) defined in $L^2(\nu)$ into a Mellin convolution equation that can be studied in $L^2(\mathbb{R}_+)$ or even in the sense of Mellin distribution, see (8.11) below for a precise statement. Then, by means of Mellin transform techniques we study this latter equation to obtain, in the distributional sense, necessary and sufficient conditions for existence, uniqueness and description of its solution. In particular, we get a characterization in terms of the Rodrigues operator acting on the density of the invariant measures ν . Then, applying the variety of results on ν (smoothness, positivity, small and large asymptotic equivalents or bounds) developed in Chapter 5, we obtain (almost) necessary and sufficient conditions for the existence of a unique solution to the original equation (2.6) considered in the Hilbert space $L^2(\nu)$. It is worth pointing out that the intensive study that we carry out on the density of the invariant measure includes some innovative techniques that could be used in a larger context. For instance, the large asymptotic behaviour of ν together with its successive derivatives of any order stems from a generalized version of a non-classical Tauberian theorem which was initially proved by Balkema et al. [BKS93] and that we state in Proposition 5.4.3. We simply mention that the establishment of the ultimate log-convexity property of the density and its derivatives is one of the several delicate Tauberian conditions. To the best of our knowledge, it seems that it is the first instance that can be found in the literature of a class of probability densities and its successive derivatives, for which such precise description of the (exponential) asymptotic decay is available at real rather than log-scale, i.e. for $\ln \nu(x, \infty)$. Finally, the small asymptotic behaviour of ν and its derivatives, is investigated by an appeal to the Itô's excursion theory for Lévy processes, which in its own right is a new and interesting result.

5. Extension of the domain of the spectral operator. Finally, we address the following three issues. The first one consists in characterizing the domain \mathcal{D} of the spectral operator S that is defined as

$$(2.7) \quad S_t f = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_{\nu} \mathcal{P}_n,$$

and \mathcal{D} is the union of domains of the form, for some $T > 0$, linear space $L \subseteq L^2(\nu)$ and $\underline{\mathcal{N}} \subseteq \mathcal{N}$,

$$(2.8) \quad \mathcal{D}_T^{\underline{\mathcal{N}}}(L) = \{(\psi, f, t); \psi \in \underline{\mathcal{N}}, f \in L \text{ and } t \in (T = T(\psi), \infty)\}.$$

Note that, since for any $(\psi, f, T) \in \mathcal{D}$, $\underline{c}_t(f) = (e^{-nt} \langle f, \mathcal{V}_n \rangle_{\nu})_{n \geq 0} \in \ell^2(\mathbb{N})$, the continuity property of the synthesis operator \mathcal{S} in (1.31) implies that for any $t > T$,

$$S_t f = \mathcal{S}(\underline{c}_t(f)) \in L^2(\nu).$$

The next step is to find conditions for the identity $S_t f = P_t f$ to hold in $L^2(\nu)$. Although the operators coincide on a dense domain of $L^2(\nu)$, it may not be obvious to show that any extension of S_t is bounded in $L^2(\nu)$. We shall design in Chapter 12 a specific methodology for each subdomain that deals with this issue. Finally, we are interested in establishing conditions under which the spectral operator (2.7) along with its space-time partial derivatives converge locally uniformly? Although the Hilbert space topology is natural for the spectral decomposition of the operator, the locally uniform convergence enables to derive regularity properties for the semigroups and the solution to the associated Cauchy problem. In the same vein, similar issues may be addressed for the heat kernel.

In order to characterize the domain \mathcal{D} , we resort to the Cauchy-Schwartz inequality that leads us to find upper bounds for the co-eigenfunctions \mathcal{V}_n which may be considered in different topologies. We emphasize that it is a very difficult problem to obtain these bounds as it requires precise information regarding the uniform asymptotic behaviour for large n of \mathcal{V}_n . We point out that there is a rich and fascinating literature on uniform asymptotic expansions of the Laguerre polynomials which reveals already that this issue, even in a simple case with explicit expression and several representations at hand, is far from being trivial, see e.g. [Sze75] and [Tem90] for a thorough description. We present below the different bounds we manage to extract for each of the subdomain defining \mathcal{D} . Recall the definition of various subspaces of \mathcal{N} given in Table 1 and Table 2 together with parameters such as $\underline{\Theta}_\phi$, the definition of d_ϕ in (1.19) and the relation $\mathcal{V}_n \nu = w_n$.

THEOREM 2.1.5. (1) Let $\psi \in \mathcal{N}_{\underline{\Theta}}$ and recall that $T_{\underline{\Theta}_\phi} = -\ln \sin \underline{\Theta}_\phi$ and $\mathbf{p}_\alpha(x) = x^{-\alpha}$, $x > 0$, $\alpha \in [0, 1)$. Then, we have, for any $\epsilon > 0$ and n large, uniformly on $x > 0$,

$$(2.9) \quad x^a |w_n(x)| = O\left(n^{\frac{1}{2}-a} e^{(T_{\underline{\Theta}_\phi} + \epsilon)n}\right), \quad a > d_\phi,$$

$$(2.10) \quad \left\| \frac{w_n}{\mathbf{p}_\alpha} \right\|_{\mathbf{p}_\alpha} = O\left(n^{\frac{1}{2}-a} e^{(T_{\underline{\Theta}_\phi} + \epsilon)n}\right), \quad d_\phi < a < \frac{\alpha + 1}{2}.$$

(2) Let $\psi \in \mathcal{N}_\alpha$ and recall that $T_{\pi\alpha} = -\ln \sin\left(\frac{\pi}{2}\alpha\right)$ and

$$\bar{\nu}_\gamma(x) = \max(\nu(x), e^{-x^{\frac{1}{\gamma}}}), \gamma > 1 + \alpha.$$

Then, for large n and for any $\epsilon > 0$, we have that

$$(2.11) \quad \left\| \frac{w_n}{\bar{\nu}_\gamma} \right\|_{\bar{\nu}_\gamma} = O\left(n^{\frac{1}{2}-a} e^{(T_{\pi_\alpha} + \epsilon)n}\right), \quad a > d_\phi,$$

$$(2.12) \quad \|\mathcal{V}_n\|_\nu = O\left(e^{(T_{\pi_\alpha, \rho_\alpha} + \epsilon)n}\right)$$

with $T_{\pi_\alpha, \rho_\alpha} = \max\left(T_{\pi_\alpha}, 1 + \frac{1}{\rho_\alpha}\right)$, where ρ_α is the largest root of

$$(1 - \rho)^\frac{1}{\alpha} \cos\left(\frac{\arcsin(\rho)}{\alpha}\right) = \frac{1}{2}.$$

(3) Let $\psi \in \mathcal{N}_P$. Then, for any $\epsilon > 0$ and large n ,

$$(2.13) \quad \|\mathcal{V}_n\|_\nu = O(e^{\epsilon n}).$$

If in addition $\bar{\Pi}(0^+) < \infty$ then, recalling that $\underline{\mathbf{m}} = \frac{m + \bar{\Pi}(0^+)}{\sigma^2}$, we have for large n ,

$$(2.14) \quad \|\mathcal{V}_n\|_\nu = O(n^{\underline{\mathbf{m}}}).$$

(4) Let $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$ with $(\alpha, \mathbf{m}) \in \mathfrak{R} = \{(\alpha, \mathbf{m}); \alpha \in (0, 1] \text{ and } \mathbf{m} \geq 1 - \frac{1}{\alpha}\}$. Then, for large n , with $T_\alpha = -\ln(2^\alpha - 1)$, we have that

$$(2.15) \quad \|\mathcal{V}_n\|_\nu = O(e^{T_\alpha n}).$$

To derive these bounds we develop several approaches which are of different nature.

For the first one, to obtain the bound (2.10) when $\psi \in \mathcal{N}_{\Theta}$, we are able to apply a classical saddle-point approximation to the Mellin-Barnes integral representation of $w_n = \mathcal{V}_n \nu$. This latter is expressed in terms of the Mellin transform of ν that we study in-depth in Chapter 4 which includes its characterization as an infinite product generalizing the classical Weierstrass product representation of the gamma function. We again emphasize that a more thorough study is now available in [PS18a].

A second path that we follow relies on a fine study of the locations of the real zeros of the successive derivatives of ν . This approach necessitates detailed information regarding the Itô's excursion measure of some Lévy processes, which forces us to specialize to the regularly varying case. Once the distribution of the real zeros of the derivatives of ν is approximated, one uses the basic theorem of calculus to first provide uniform estimates for $|\mathcal{V}_n(x)|$ and then deduce bounds for $\frac{\mathcal{V}_n \nu}{\bar{\nu}_\gamma}$ in the topology of the Hilbert space $L^2(\bar{\nu}_\gamma)$.

Still focussing on the regularly varying case, we develop a complex analytical approach based on the Phragmenn-Lindelöf principle to establish upper bounds on $|\mathcal{V}_n(x)|$ yielding to precise $L^2(\nu)$ -norm estimates for \mathcal{V}_n .

Finally, the last methodology is based on new structural ideas that we name the concept of reference semigroups. It consists on identifying gL semigroups P which satisfy the following two criteria. First, their special structure permits to study their spectral reduction in details, including information regarding the asymptotic behavior of the norm of co-eigenfunctions and the completeness of their sequence. On

the other hand, a reference semigroup intertwines with a subclass of gL semigroups in such way that these properties can be easily transferred to a priori intractable subclass of gL semigroups. For instance, there should exist a subclass of gL semigroups such that for each element in this class we have the adjoint intertwining relation $P_t^* \Lambda^* = \Lambda^* \overline{P}_t^*$, where Λ^* is the adjoint of a bounded operator between appropriate weighted Hilbert spaces.

We describe two important applications of the reference concept regarding probably the two most technical difficulties of this work, namely the estimates of the $L^2(\nu)$ -norm of \mathcal{V}_n and their completeness in $L^2(\nu)$. We manage to implement this approach for two different reference semigroups, which, in particular, enable us to deal with the spectral expansion, in a simple and optimal way as the expansion operator is proved to be bounded on $L^2(\nu)$ for all $t > 0$, for the perturbation class \mathcal{N}_P , that is when $\sigma^2 > 0$. It is also worth pointing out that when $\sigma^2 > 0$ and $\overline{\Pi}(0^+) = \infty$, the concept of reference semigroup has some unusual underlying mathematical idea. Indeed, it consists in approximating, in the sense of linking operators via intertwining relations, the class of perturbed operators (say the nice class) by families of operators for which the spectral operator is bounded in $L^2(\nu)$ only for $t > T_\alpha = -\ln(2^\alpha - 1)$ (say the non nice class). Finally, we mention that this approach goes well beyond this framework as it is characterized by the class \mathcal{N}_R which also encompasses gL semigroups whose infinitesimal generator does not have a diffusion component. It also enables to deal with the delicate issue of characterizing the class of co-eigenfunctions that form a complete sequence in the weighted Hilbert space.

2.2. Proof of Theorem 1.1.1(1)

According to Lamperti [**Lam72**], there is a bijection between the subspace of negative definite functions \mathcal{N} and the conservative Feller semigroups $(K_t)_{t \geq 0}$ on $(0, \infty)$ corresponding to processes that have only negative jumps and satisfying the following 1-self-similarity property, for any $c > 0$, $t \geq 0$,

$$K_t f(cx) = K_{c^{-1}t} d_c f(x),$$

where we recall that $d_c f(x) = f(cx)$. More specifically, Lamperti showed that for any $\psi \in \mathcal{N}$, the infinitesimal generator \mathbf{G}_0 of $(K_t)_{t \geq 0}$, takes, for any function f such that $x \mapsto f_e(x) = f(e^x) \in \mathcal{C}^2([-\infty, \infty])$, the form

$$(2.16) \quad \mathbf{G}_0 f(x) = \sigma^2 x f''(x) + (m + \sigma^2) f'(x) + \int_0^\infty f(e^{-y}x) - f(x) + yx f'(x) \frac{\Pi(dy)}{x},$$

where (σ, m, Π) is the characteristic triplet of ψ . Next, we define, for any $t \geq 0$,

$$(2.17) \quad P_t f(x) = K_{e^t - 1} d_{e^{-t}} f(x),$$

and note from the 1-self-similarity property of K that

$$(2.18) \quad P_t f(x) = K_{1 - e^{-t}} f(xe^{-t}).$$

Then for each $t \geq 0$, P_t is plainly linear, with $P_t \mathbf{C}_0(\overline{\mathbb{R}}_+) \subseteq \mathbf{C}_0(\overline{\mathbb{R}}_+)$. Moreover, from (2.18), we get that $\|P_t f\|_\infty \leq \|f\|_\infty$ and $\lim_{t \downarrow 0} P_t f = f$. Next, for any $t, s > 0$,

$$\begin{aligned} P_t P_s f(x) &= K_{1 - e^{-t}} K_{e^s - 1} d_{e^{-s}} f(xe^t) = K_{e^s - e^{-t}} d_{e^{-s}} f(xe^t) \\ &= K_{e^{t+s} - 1} d_{e^{-(t+s)}} f(x) = P_{t+s} f(x). \end{aligned}$$

It is easy to check now that the semigroup $(P_t)_{t \geq 0}$ satisfies the properties 1. and 2. of Definition 1.0.1. Moreover, from [BS11], we deduce that $\psi \in \mathcal{N}$ if and only if

$$\lim_{t \rightarrow \infty} P_t f(x) = \lim_{t \rightarrow \infty} K_{1-e^{-t}} f(e^{-t}x) = K_1 f(0) = \nu f,$$

where we have used (2.18) and for the last equality the identities (2.22) and (2.23), yielding to the condition 3. of the definition of the gL semigroup. Next, for f a smooth function, we have

$$\begin{aligned} \mathbf{G}f(x) &= \lim_{t \rightarrow 0} \frac{P_t f(x) - f(x)}{t} \\ &= \lim_{t \rightarrow 0} \frac{K_{1-e^{-t}} f(x) - f(x) + K_{1-e^{-t}} f(e^{-t}x) - K_{1-e^{-t}} f(x)}{1 - e^{-t}} \\ (2.19) \quad &= \mathbf{G}_0 f(x) - x f'(x), \end{aligned}$$

which combined with (2.16) gives the expression of \mathbf{G} . Moreover, following [Mey67] and performing a change of variables, we get that the Lévy kernel of P is characterized for any $f \in \mathcal{C}_c^\infty(\mathbb{R}_+ \setminus \{x\})$, by

$$\lim_{t \rightarrow 0} \frac{1}{t} P_t f(x) = \int_0^x f(r) \frac{\hat{\Pi}_l(x, dr)}{x},$$

where $\hat{\Pi}_l(x, \cdot)$ is the image of $\hat{\Pi}(\cdot)$ by the mapping $r \mapsto \ln(\frac{x}{r})$. Hence

$$\Pi(x, (x, \infty)) = \frac{\hat{\Pi}_l(x, (x, \infty))}{x} = 0$$

and the last condition 4. is also satisfied. By putting pieces together, we complete the proof this item.

2.3. Additional basic facts on gL semigroups

PROPOSITION 2.3.1. (1) *For any $\psi \in \mathcal{N}$, there exists a positive random variable V_ψ whose law is absolutely continuous with density ν which satisfies, for any $f \in \mathcal{C}_0(\overline{\mathbb{R}}_+)$ and $t \geq 0$,*

$$\nu P_t f = \nu f,$$

where here and hereafter $\nu f = \int_0^\infty f(x) \nu(x) dx$, that is, the measure $\nu(x) dx$ is an invariant measure of P . Moreover the law of V_ψ is determined by its entire moments given by

$$(2.20) \quad \mathcal{M}_{V_\psi}(n+1) = W_\phi(n+1), \quad n \in \mathbb{N},$$

where we recall from (1.18) that $W_\phi(n+1) = \prod_{k=1}^n \phi(k)$.

(2) *P can be extended uniquely to a strongly continuous contraction semigroup, still denoted by P , on the weighted Hilbert space $L^2(\nu)$.*

Proof. Item (1) can be deduced easily from [BY02b, Proposition 1(ii)], which states that for any $\psi \in \mathcal{N}$, there exists a positive variable V_ψ such that, for any $n \in \mathbb{N}$,

$$(2.21) \quad \mathcal{M}_{V_\psi}(n+1) = \mathbb{E}[V_\psi^n] = \frac{\prod_{k=1}^n \psi(k)}{n!} = W_\phi(n+1),$$

where the second identity follows from the definition of $\phi(u) = \frac{\psi(u)}{u}$. The absolute continuity of its law is proved in [Pat11, Proposition 2.4]. Then, write, for any $t > 0$,

$$(2.22) \quad t\nu_t(tx) = \nu(x)$$

i.e. $\nu f = \nu_t d_{1/t} f$ with d , the dilation operator. Then, from [BY02b, Proposition 1(ii)] augmented by a moment identification identifies $(\nu_t)_{t \geq 0}$ as the family of entrance laws for the semigroup K , that is, for any $t, s > 0$ and any $f \in \mathcal{C}_0(\overline{\mathbb{R}_+})$,

$$(2.23) \quad \nu_t K_s f = \nu_{t+s} f.$$

Next, using successively the definition of P , recalled in (2.17), the previous identity with $t = 1$ and $s = e^t - 1$, and the definition of ν_t above, we get, since plainly $d_{e^{-t}} f \in \mathcal{C}_0(\overline{\mathbb{R}_+})$, for any $t > 0$, that

$$\nu P_t f = \nu K_{e^t - 1} d_{e^{-t}} f = \nu_{e^t} d_{e^{-t}} f = \nu f,$$

which completes the proof of item (1). Turning now to (2), since $\nu(x)dx$ is an invariant measure, a standard result, see e.g. [DP06], provides the existence of a strongly continuous semigroup extension of P on $L^2(\nu)$. \square

CHAPTER 3

Examples

This part is devoted to the description of the eigenvalues expansions of specific instances of the gL semigroups, including the so-called reference semigroups that are exploited in Chapter 10. These examples illustrate the different situations that are treated in this work ranging from the self-adjoint case to perturbation of a self-adjoint differential operator through non-local operators without diffusion component.

EXAMPLE 3.0.1. *The self-adjoint diffusion case.* Let us consider, for any $m \geq 0$,

$$(3.1) \quad \psi(u) = u^2 + mu,$$

namely $\sigma^2 = 1$ and $\Pi \equiv 0$ in (1.3), that is the associated gL semigroups are the classical Laguerre semigroups which generate the class of squared radial Ornstein-Uhlenbeck processes of order m . We refer to the monographs [BS02] and [KT81] for a thorough account on these linear diffusions. We start by providing a detailed description of the Laguerre semigroup of order 0 as it plays a central role in this work and proceed with the essential elements which characterize the classical Laguerre semigroup of higher orders, which are used in Chapter 10. The Laguerre semigroup $Q = (Q_t)_{t \geq 0}$ of order 0 generates the so-called 2-dimensional squared radial Ornstein-Uhlenbeck process of parameter 1, denoted by $R = (R_t)_{t \geq 0}$. Its infinitesimal generator \mathbf{L} takes the form, for a function $f \in \mathcal{C}_0^2(\mathbb{R}_+)$,

$$\mathbf{L}f(x) = xf''(x) + (1-x)f'(x).$$

Note that the point 0 is an entrance-non exit boundary. The process R can be also realized as the solution to the stochastic differential equation

$$dR_t = (1 - R_t)dt + \sqrt{2R_t}dB_t, \quad R_0 = x \geq 0,$$

where $B = (B_t)_{t \geq 0}$ is a standard Brownian motion. The process R is a positively recurrent Feller diffusion on \mathbb{R}_+ with an absolutely continuous stationary measure, whose density is given by

$$\varepsilon(x) = e^{-x}, \quad x > 0,$$

that is the exponential distribution of parameter 1. The semigroup $Q = (Q_t)_{t \geq 0}$ is a strongly continuous contraction semigroup from the weighted Hilbert space $L^2(\varepsilon)$, endowed with inner product $\langle \cdot, \cdot \rangle_\varepsilon$, into itself. Moreover, it admits the eigenvalues expansions, valid for any $f \in L^2(\varepsilon)$ and $x \geq 0, t > 0$,

$$(3.2) \quad Q_t f(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{L}_n \rangle_\varepsilon \mathcal{L}_n(x) \quad \text{in } L^2(\varepsilon),$$

where, for any $n \geq 0$, \mathcal{L}_n is the Laguerre polynomial of order 0, defined either by means of the Rodrigues operator $\mathcal{R}^{(n)}$ as follows

$$(3.3) \quad \mathcal{L}_n(x) = \frac{\mathcal{R}^{(n)}\varepsilon(x)}{\varepsilon(x)} = \frac{1}{n!} \frac{(x^n \varepsilon(x))^{(n)}}{\varepsilon(x)},$$

or, through the polynomial representation

$$(3.4) \quad \mathcal{L}_n(x) = \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{x^k}{k!}.$$

The Laguerre polynomials also satisfy the well-known three terms recurrence relation, for any $n \geq 2$,

$$(3.5) \quad \mathcal{L}_n(x) = \left(2 + \frac{-1-x}{n}\right) \mathcal{L}_{n-1}(x) - \left(1 - \frac{1}{n}\right) \mathcal{L}_{n-2}(x).$$

The eigenvalues expansion is readily seen from the self-adjointness property of the semigroup Q in $L^2(\varepsilon)$, the facts that the sequence of normalized Laguerre polynomials $(\mathcal{L}_n)_{n \geq 0}$ forms an orthonormal basis of the Hilbert space $L^2(\varepsilon)$ and that it also corresponds to the sequence of eigenfunctions of Q_t associated to the set of eigenvalues $(e^{-nt})_{n \geq 0}$, that is

$$(3.6) \quad Q_t \mathcal{L}_n(x) = e^{-nt} \mathcal{L}_n(x).$$

Furthermore, the semigroup Q can be represented in terms of the semigroup $K^{(0)} = (K_t^{(0)})_{t \geq 0}$ of the 2-dimensional squared Bessel process as follows

$$(3.7) \quad Q_t f(x) = K_{e^{-t}-1}^{(0)} d_{e^{-t}} f(x).$$

More generally, for any $m > 0$, in (3.1), the infinitesimal generator of the associated Laguerre semigroup, which takes the form

$$\mathbf{L}_m f(x) = x f''(x) + (m+1-x) f'(x),$$

corresponds to the complete, up to a dilation, family (indexed by m) of second order differential operators included in our class of generators. It is the infinitesimal generator of a one-dimensional diffusion often referred to in the literature as the squared radial Ornstein-Uhlenbeck process, see e.g. [BS02, Appendix 1.26]. From $\psi(u) = u^2 + mu$ we get that $\phi(u) = u + m$, and, by moment identification a simple algebra yields, from (6.11) that $I_\phi \stackrel{(d)}{=} B(1, m)$ with $B(1, m)$ a Beta random variable of parameters 1 and m , and, from (2.20) that $V_\psi \stackrel{(d)}{=} G(m+1)$, where $G(m+1)$ is a Gamma random variable of parameter $m+1$ whose distribution admits a density $\varepsilon_m(x) = \frac{x^m e^{-x}}{\Gamma(m+1)}$, $x > 0$, which was already introduced in Remark 1.0.3. Moreover from (7.9), (7.24) and (1.21), the eigenfunctions are given, for any $n \geq 0$, by

$$(3.8) \quad \mathcal{P}_n(x) = \mathcal{I}_\phi \mathcal{L}_n(x) = \sum_{k=0}^n \binom{n}{k} \frac{k! \Gamma(m+1)}{\Gamma(k+m+1)} \frac{(-x)^k}{k!} = \mathbf{c}_n(m) \mathcal{L}_n^{(m)}(x)$$

where

$$\mathbf{c}_n(m) = \frac{\Gamma(n+1)\Gamma(m+1)}{\Gamma(n+m+1)},$$

$\mathcal{L}_n = \mathcal{L}_n^{(0)}$ and

$$\mathcal{L}_n^{(m)}(x) = \sum_{k=0}^n (-1)^k \binom{n+m}{n-k} \frac{x^k}{k!}$$

is the associated Laguerre polynomial of order m . From (1.22), we get that, for any $n \geq 0$,

$$\mathcal{V}_n(x) = \frac{\mathcal{R}^{(n)}\varepsilon_m(x)}{\varepsilon_m(x)} = \mathcal{L}_n^{(m)}(x),$$

where the second equality follows from a classical representation of the Laguerre polynomials, see [Leb72, (4.17.1)]. From these identities, we recover the well-known facts that the semigroup P is self-adjoint in $L^2(\varepsilon_m)$ and the sequence $(\sqrt{\mathbf{c}_n(m)}\mathcal{L}_n^{(m)})_{n \geq 0}$ forms an orthonormal basis in $L^2(\varepsilon_m)$. Finally, we get, for any $t > 0$, that

$$P_t(x, y) = \sum_{n=0}^{\infty} e^{-nt} \mathbf{c}_n(m) \mathcal{L}_n^{(m)}(y) \mathcal{L}_n^{(m)}(x) \varepsilon_m(y),$$

an expression which can be found, for instance, in [KT81, Chap. 15]. Note that, in this case since $\sigma^2 > 0$, the expansion is convergent in the Hilbert space topology and locally uniformly for all $f \in L^2(\varepsilon_m)$ and $t, x > 0$. The latter also follows from (1.25).

EXAMPLE 3.0.2. *Small perturbation of the Laguerre semigroup.* Let $\mathbf{m} \geq 1$ and consider, for any $u > 0$,

$$\phi_{\mathbf{m}}(u) = \frac{(u + \mathbf{m} + 1)(u + \mathbf{m} - 1)}{u + \mathbf{m}} = u + \frac{\mathbf{m}^2 - 1}{\mathbf{m}} + \int_0^{\infty} (1 - e^{-uy})e^{-my} dy.$$

Since plainly $\phi_{\mathbf{m}} \in \mathcal{B}_{\mathcal{N}}$, see Proposition 4.1.1(8), we have that

$$\psi(u) = u \frac{(u + \mathbf{m} + 1)(u + \mathbf{m} - 1)}{u + \mathbf{m}} \in \mathcal{N}_P,$$

as $\sigma^2 = 1$, $m = \frac{\mathbf{m}^2 - 1}{\mathbf{m}}$, $\bar{\Pi}(y) = e^{-my}$ in (1.3) and note that $\frac{\bar{\Pi}(0^+) + m}{\sigma^2} = \mathbf{m}$. The infinitesimal generator of the associated gL semigroup is the integro-differential operator

$$\mathbf{G}_{\mathbf{m}}f(x) = xf''(x) + (\mathbf{m} + 1 - x)f'(x) + \frac{\mathbf{m}}{x} \int_0^{\infty} (f(e^{-y}x) - f(x))e^{-my} dy.$$

Moreover, we get,

$$W_{\phi_{\mathbf{m}}}(n + 1) = \frac{\mathbf{m}}{n + \mathbf{m}} \frac{\Gamma(n + \mathbf{m} + 2)}{\Gamma(\mathbf{m} + 2)},$$

that is, by moment identification,

$$\nu(x) = \frac{1 + x}{\mathbf{m} + 1} \frac{x^{\mathbf{m}-1} e^{-x}}{\Gamma(\mathbf{m})} = \frac{(1 + x)}{\mathbf{m} + 1} \varepsilon_{\mathbf{m}-1}(x), \quad x > 0.$$

Thus, from (1.22) and (3.8), we have that, for $n \geq 1$, the ν_n 's can be expressed in terms of the Laguerre polynomials as follows,

$$\begin{aligned} \mathcal{V}_n(x) &= \frac{\mathcal{R}^{(n)}\nu(x)}{\nu(x)} = \frac{1}{x + 1} \mathcal{L}_n^{(\mathbf{m}-1)}(x) + \frac{x}{1 + x} \mathcal{L}_n^{(\mathbf{m})}(x) \\ &= \frac{1}{x + 1} \left((1 - n) \mathcal{L}_n^{(\mathbf{m}-1)}(x) + (n + \mathbf{m} - 1) \mathcal{L}_{n-1}^{(\mathbf{m}-1)}(x) \right) \end{aligned}$$

where we used the 3-point-rules of Laguerre polynomials. In this case, one gets, for any $n \geq 1$,

$$\begin{aligned}
\|\mathcal{V}_n\|_\nu^2 &= \int_0^\infty \left((1-n)\mathcal{L}_n^{(\mathbf{m}-1)}(x) + (n+\mathbf{m}-1)\mathcal{L}_{n-1}^{(\mathbf{m}-1)}(x) \right)^2 \frac{\varepsilon_{\mathbf{m}-1}(x)}{(\mathbf{m}+1)(1+x)} dx \\
&\leq 2 \left(\int_0^\infty \left(\frac{n-1}{\sqrt{\mathbf{m}+1}} \mathcal{L}_n^{(\mathbf{m}-1)}(x) \right)^2 + \left(\frac{n+\mathbf{m}-1}{\sqrt{\mathbf{m}+1}} \mathcal{L}_{n-1}^{(\mathbf{m}-1)}(x) \right)^2 \varepsilon_{\mathbf{m}-1}(x) dx \right) \\
&= 2 \frac{(n-1)^2}{\mathbf{m}+1} \|\mathcal{L}_n^{(\mathbf{m}-1)}\|_{\varepsilon_{\mathbf{m}-1}}^2 + 2 \frac{(n+\mathbf{m}-1)^2}{\mathbf{m}+1} \|\mathcal{L}_{n-1}^{(\mathbf{m}-1)}\|_{\varepsilon_{\mathbf{m}-1}}^2 \\
&= 2 \frac{(n-1)^2}{\mathbf{m}+1} \frac{\Gamma(n+\mathbf{m})}{\Gamma(\mathbf{m})n!} + 2 \frac{(n+\mathbf{m}-1)^2}{\mathbf{m}+1} \frac{\Gamma(n+\mathbf{m}-1)}{\Gamma(\mathbf{m})(n-1)!} \\
(3.9) \quad &= O(n^{\mathbf{m}+1}),
\end{aligned}$$

where we used for the second equality the fact that the sequence $(\sqrt{c_n(\mathbf{m})}\mathcal{L}_n^{(\mathbf{m}-1)})_{n \geq 0}$, with

$$c_n(\mathbf{m}) = \frac{\Gamma(n+1)\Gamma(\mathbf{m})}{\Gamma(n+\mathbf{m})},$$

forms an orthonormal sequence in $L^2(\varepsilon_{\mathbf{m}-1})$. On the other hand, we have from (1.21) applied with $\phi_{\mathbf{m}}$ identified with (3.8), for any $n \geq 0$,

$$\begin{aligned}
\mathcal{P}_n(x) &= \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{\Gamma(\mathbf{m}+2)}{\Gamma(\mathbf{m}+k+2)} \frac{\mathbf{m}+k}{\mathbf{m}} x^k \\
&= c_n(\mathbf{m}+1)\mathcal{L}_n^{(\mathbf{m}+1)}(x) - \frac{c_n(\mathbf{m}+1)}{\mathbf{m}} x \mathcal{L}_{n-1}^{(\mathbf{m}+1)}(x).
\end{aligned}$$

Note that from Theorem 1.5.1(2) the sequences $(\mathcal{P}_n(x))_{n \geq 0}$ and $(\mathcal{V}_n(x))_{n \geq 0}$ are biorthogonal in $L^2(\nu)$. Finally, as $(\mathcal{P}_n)_{n \geq 0}$ forms a Bessel sequence in $L^2(\nu)$, we have for all $f \in L^2(\nu)$ and all $t > 0$,

$$P_t f(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \frac{(n+\mathbf{m})}{\mathbf{m}} c_n(\mathbf{m}+1) \mathcal{L}_n^{(\mathbf{m}+1)}(x) \text{ in } L^2(\nu).$$

EXAMPLE 3.0.3. *The Gauss-Laguerre semigroup.* In [PS17], we introduce and study in depth the so-called Gauss-Laguerre semigroup, an instance of gL semigroup whose infinitesimal generator, for any $\alpha \in (0, 1)$ and $\mathbf{m} \in [1 - \frac{1}{\alpha}, \infty)$, and, for any given smooth function f , takes the form

$$\mathbf{G}_{\alpha, \mathbf{m}} f(x) = (\mathbf{m}_\alpha - x) f'(x) + \frac{\sin(\alpha\pi)}{\pi} x \int_0^1 f''(xy) g_{\alpha, \mathbf{m}}(y) dy, \quad x > 0,$$

where $\mathbf{m}_\alpha = \frac{\Gamma(\alpha\mathbf{m} + \alpha + 1)}{\Gamma(\alpha\mathbf{m} + 1)}$ and

$$g_{\alpha, \mathbf{m}}(y) = \frac{\Gamma(\alpha)}{\mathbf{m} + \frac{1}{\alpha} + 1} y^{\mathbf{m} + \frac{1}{\alpha} + 1} {}_2F_1(\alpha(\mathbf{m} + 1) + 1, \alpha + 1; \alpha(\mathbf{m} + 1) + 2; y^{\frac{1}{\alpha}}),$$

with ${}_2F_1$ the Gauss hypergeometric function. The terminology is motivated by the limit case $\alpha = 1$ which is proved to yield, writing simply $\mathbf{L}_{\mathbf{m}} = \mathbf{G}_{1, \mathbf{m}}$,

$$\mathbf{L}_{\mathbf{m}} f(x) = x f''(x) + (\mathbf{m} + 1 - x) f'(x),$$

that is the Laguerre differential operator of order \mathbf{m} , see Example 3.0.1 above. The algebra of polynomials \mathbf{P} is a core for $\mathbf{G}_{\alpha, \mathbf{m}}$ and the associated semigroup $P^{\alpha, \mathbf{m}} = (P_t^{\alpha, \mathbf{m}})_{t \geq 0}$ is a non-self-adjoint contraction semigroup in $L^2(\mathbf{e}_{\alpha, \mathbf{m}})$, where

$$(3.10) \quad \mathbf{e}_{\alpha, \mathbf{m}}(x) dx = \frac{x^{\mathbf{m} + \frac{1}{\alpha} - 1} e^{-x \frac{1}{\alpha}}}{\Gamma(\alpha \mathbf{m} + 1)} dx, \quad x > 0,$$

is its unique invariant measure. Observe that, for any $y > 0$, $z \mapsto \mathbf{e}_{\alpha, \mathbf{m}}(\frac{y}{1-z})$ is analytical in the interior of unit disc since $\mathbf{e}_{\alpha, \mathbf{m}}(z)$ is analytical on $\mathbb{C}_{(0, \infty)}$. Note that $P^{\alpha, \mathbf{m}}$ is the gL semigroup associated to $\psi_{\alpha, \mathbf{m}}^R \in \mathcal{N}$, where

$$\begin{aligned} \psi_{\alpha, \mathbf{m}}^R(u) &= u \phi_{\alpha, \mathbf{m}}^R(u) = u \frac{\Gamma(\alpha u + \alpha \mathbf{m} + 1)}{\Gamma(\alpha u + \alpha \mathbf{m} + 1 - \alpha)} \\ &= \frac{u}{\Gamma(1 - \alpha)} \int_0^\infty (1 - e^{-uy}) e^{-(\mathbf{m} + \frac{1}{\alpha})y} (1 - e^{-\frac{y}{\alpha}})^{-\alpha - 1} dy + u \frac{\Gamma(\alpha \mathbf{m} + 1)}{\Gamma(\alpha \mathbf{m} + 1 - \alpha)}, \end{aligned}$$

and we refer to Lemma 4.2.1 for more details on the computation. In order to recall some additional results, we proceed by setting some further notation. For any $x \geq 0$, we set $\mathcal{L}_0^{\alpha, \mathbf{m}}(x) = 1$ and for any $n \geq 1$, we introduce the polynomials

$$\mathcal{L}_n^{\alpha, \mathbf{m}}(x) = \Gamma(\alpha \mathbf{m} + 1) \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{\Gamma(\alpha k + \alpha \mathbf{m} + 1)} x^k.$$

Note that for $\alpha = 1$,

$$\mathcal{L}_n^{\alpha, \mathbf{m}}(x) = \mathbf{c}_n(\mathbf{m}) \mathcal{L}_n^{(\mathbf{m})}(x) = \Gamma(\mathbf{m} + 1) \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{\Gamma(k + \mathbf{m} + 1)} x^k$$

are the classical Laguerre polynomials of order $\mathbf{m} \geq 0$, see (3.8). Moreover, for any $x \geq 0$, we write $\mathcal{V}_n^{(\alpha, \mathbf{m})}(x) = \frac{\mathcal{R}^{(n)} \mathbf{e}_{\alpha, \mathbf{m}}(x)}{\mathbf{e}_{\alpha, \mathbf{m}}(x)}$, $n \in \mathbb{N}$, that is

$$\mathcal{V}_n^{(\alpha, \mathbf{m})}(x) = \frac{(-1)^n}{n! \mathbf{e}_{\alpha, \mathbf{m}}(x)} (x^n \mathbf{e}_{\alpha, \mathbf{m}}(x))^{(n)}.$$

From the Rodrigues representation of the Laguerre polynomials, we also get that for $\alpha = 1$,

$$\mathcal{V}_n^{(\alpha, \mathbf{m})}(x) = \mathcal{L}_n^{\mathbf{m}}(x).$$

We also define, for any $0 < \gamma < \alpha$ and $\bar{\eta}_\alpha > 0$ fixed,

$$\bar{\mathbf{e}}_{\gamma, \mathbf{m}}(x) = x^{\mathbf{m} + \frac{1}{\alpha} - 1} e^{\bar{\eta}_\alpha x^{\frac{1}{\gamma}}}, \quad x > 0,$$

and, recall that $T_\alpha = -\ln(2^\alpha - 1)$. By means of a delicate and non-classical saddle-point analysis, we obtain, in [PS17, Proposition 2.3], the following specific asymptotic estimates for the co-eigenfunctions for large values of the parameter n ,

$$(3.11) \quad \|\mathcal{V}_n^{(\alpha, \mathbf{m})}\|_{\mathbf{e}_{\alpha, \mathbf{m}}} = O(e^{T_\alpha n}),$$

and

$$(3.12) \quad \left\| \mathcal{V}_n^{(\alpha, \mathbf{m})} \frac{\mathbf{e}_{\alpha, \mathbf{m}}}{\bar{\mathbf{e}}_{\gamma, \mathbf{m}}} \right\|_{\bar{\mathbf{e}}_{\gamma, \mathbf{m}}} = O\left(n^{1 + \mathbf{m} + \frac{1}{\alpha} + \alpha} e^{\bar{\mathbf{t}}_\alpha n^{\frac{1}{\alpha+1}}}\right),$$

where

$$\bar{\mathbf{t}}_\alpha = (\alpha + 1) \alpha^{-\frac{\alpha}{\alpha+1}} \left(\frac{\alpha + 1}{\alpha} + \epsilon \right)^{\frac{1}{\alpha+1}},$$

for some small $\epsilon > 0$. From this asymptotic analysis, we deduce the following fine properties of the Gauss-Laguerre semigroup. For any $f \in L^2(\mathbf{e}_{\alpha, \mathbf{m}})$ (resp. $f \in L^2(\bar{\mathbf{e}}_{\gamma, \mathbf{m}})$) we have that

$$P_t^{\alpha, \mathbf{m}} f(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n^{(\alpha, \mathbf{m})} \rangle_{\mathbf{e}_{\alpha, \mathbf{m}}} \mathcal{L}_n^{\alpha, \mathbf{m}}(x),$$

where, for any $t > T_\alpha$ (resp. $t > 0$), the identity holds in $L^2(\mathbf{e}_{\alpha, \mathbf{m}})$. Moreover $P_t f \in \mathbf{C}^\infty((T_\alpha, \infty) \times \mathbb{R}_+)$ (resp. $\mathbf{C}^\infty(\mathbb{R}_+^2)$), and for any integers k, p ,

$$\frac{d^k}{dt^k} (P_t^{\alpha, \mathbf{m}} f)^{(p)}(x) = \sum_{n=p}^{\infty} (-n)^k e^{-nt} \langle f, \mathcal{V}_n^{(\alpha, \mathbf{m})} \rangle_{\mathbf{e}_{\alpha, \mathbf{m}}} (\mathcal{L}_n^{\alpha, \mathbf{m}})^{(p)}(x),$$

where, for any $t > T_\alpha$ (resp. $t > 0$), the series converges absolutely. Finally, the heat kernel is absolutely continuous with density $P_t(x, y) \in \mathbf{C}^\infty(\mathbb{R}_+^3)$, given for any $t, y > 0, x \geq 0$, and for any integers k, p, q , by

$$\frac{d^k}{dt^k} P_t^{(p, q)}(x, y) = \sum_{n=p}^{\infty} (-n)^k e^{-nt} (\mathcal{L}_n^{\alpha, \mathbf{m}})^{(p)}(x) (\mathcal{V}_n^{(\alpha, \mathbf{m})} \mathbf{e}_{\alpha, \mathbf{m}})^{(q)}(y),$$

where the series is absolutely convergent.

EXAMPLE 3.0.4. *The stationary saw-tooth semigroup.* Let $0 < a < 1 < a + b$ and

$$\psi(u) = u \frac{u+1-a}{u+b} = u \left(\frac{1-a}{b} + \int_0^\infty (1-e^{-uy}) (a+b-1) e^{-by} dy \right),$$

that is $m = \frac{1-a}{b}$ and $\bar{\Pi}(y) = (a+b-1)e^{-by}$. Hence

$$\mathfrak{r} = \bar{\Pi}(0^+) + m = \lim_{u \rightarrow \infty} \frac{u+1-a}{u+b} = 1,$$

see (1.7), and

$$N_{\mathfrak{r}} = \left\lceil \frac{\bar{\Pi}(0^+)}{\mathfrak{r}} \right\rceil - 1 = \lceil a+b-1 \rceil - 1 \geq 0,$$

see (1.8). We point out that the self-similar semigroup K , see Definition 1.0.1, associated to the gL semigroup was introduced and studied by Carmona et al. [CPY98]. They call the corresponding process the saw-tooth process due to the specific behaviour of its trajectories. Next, for any $n \geq 0$, we have from (1.18) with $\phi(u) = \frac{u+1-a}{u+b}$ that

$$W_\phi(n+1) = \frac{\Gamma(n+2-a)\Gamma(b+1)}{\Gamma(2-a)\Gamma(n+1+b)}$$

and thus from (1.21)

$$\mathcal{P}_n(x) = \frac{\Gamma(2-a)}{\Gamma(b+1)} \sum_{k=0}^n \frac{n! \Gamma(k+b+1)}{(n-k)! \Gamma(k+2-a)} \frac{(-x)^k}{k!}.$$

Moreover, by moment identification via (6.5), we easily observe that the invariant measure is the beta distribution, that is

$$\nu(x) = \frac{\Gamma(b+1)}{\Gamma(2-a)\Gamma(b+a-1)} x^{1-a} (1-x)^{b+a-2} \mathbb{I}_{\{0 < x < 1\}}.$$

From this expression, we get that indeed, for $N_{\mathfrak{r}} \geq 1$, $\nu \in \mathbf{C}^{N_{\mathfrak{r}}-1}(\mathbb{R}_+)$ and, in any case, the mapping $x \mapsto (x-1)\nu^{(N_{\mathfrak{r}})}(x)$ is continuous on \mathbb{R}_+ , as claimed in Theorem

5.0.2(2b). Moreover, for any $n \leq N_{\mathbf{r}}$, we have, for any $0 < x < 1$, and, for sake of simplicity assuming that $a + b \notin \mathbb{N}$,

$$\begin{aligned} \mathcal{V}_n(x) &= \frac{(-1)^n}{n! \nu(x)} (x^n \nu(x))^{(n)} \\ &= \frac{(-1)^n x^{a-1}}{n! (1-x)^{b+a-2}} \sum_{k=0}^n \binom{n}{k} \frac{(-1)^k \Gamma(n+2-a) \Gamma(b+a-1)}{\Gamma(k+2-a) \Gamma(b+a-1-k)} x^{k+1-a} (1-x)^{-k+b+a-2} \\ &= \frac{(-1)^n \Gamma(n+2-a) C_{a,b}}{n!} \sum_{k=0}^n \binom{n}{k} \frac{\Gamma(k+2-a-b)}{\Gamma(k+2-a)} \left(\frac{x}{1-x} \right)^k, \end{aligned}$$

where after using the reflection formula of the Gamma function, we have set

$$\pi C_{a,b} = \Gamma(b+a-1) \sin((2-a-b)\pi).$$

Thus, for any $n \geq 0$,

$$\mathcal{V}_n^2(x) \nu(x) \stackrel{1}{=} O((x-1)^{-2n+a+b-2})$$

and hence $\mathcal{V}_n \in L^2(\nu)$ if and only if $-2n+a+b-2 > -1$, that is $n < \frac{\overline{\Pi}(0^+)}{2\mathbf{r}}$, which is precisely the assertion of Theorem 1.5.1(4d).

New developments in the theory of Bernstein functions

The Bernstein functions play a central and recurrent role in this work. We review below some well known facts and present some new interesting developments which may find applications in the areas where this class of functions appears. An excellent reference on this topic is the monograph [SSV10] which is entirely devoted to their study.

4.1. Review and basic properties of Bernstein functions

We recall that the class of Bernstein functions is defined as follows

$$(4.1) \quad \mathcal{B} = \{\phi \in \mathcal{C}^\infty(\mathbb{R}_+); \phi \not\equiv 0, \phi \text{ is non-negative and } \phi' \text{ is completely monotone}\}$$

where a function f is called completely monotone if $f : \mathbb{R}_+ \rightarrow [0, \infty) \in \mathcal{C}^\infty(\mathbb{R}_+)$ and $(-1)^n f^{(n)} \geq 0$ on \mathbb{R}_+ , for all $n = 0, 1, \dots$. The next statement collects some well known and new properties of the Bernstein functions.

PROPOSITION 4.1.1. *Let $\phi \in \mathcal{B}$.*

(1) *We have that $\phi \in \mathcal{A}_{[0, \infty)}$ and ϕ has the form, for $z \in \mathbb{C}_{[0, \infty)}$,*

$$(4.2) \quad \phi(z) = m + \sigma^2 z + z \int_0^\infty e^{-zy} \bar{\mu}(y) dy = m + \sigma^2 z + \int_0^\infty (1 - e^{-zy}) \mu(dy),$$

where $m = \phi(0) \geq 0$, $\sigma^2 \geq 0$ and μ is a Lévy measure on $(0, \infty)$ such that $\int_0^\infty (1 \wedge y) \mu(dy) < \infty$ and as usual $\bar{\mu}(y) = \int_y^\infty \mu(du)$ is the tail of μ .

(2) *Moreover, ϕ is non-decreasing on \mathbb{R}_+ and*

$$(4.3) \quad \phi'(u) = \sigma^2 + \int_0^\infty e^{-uy} \bar{\mu}(y) dy - u \int_0^\infty e^{-uy} y \bar{\mu}(y) dy = \sigma^2 + \int_0^\infty e^{-uy} y \mu(dy)$$

as a completely monotone function is positive non-increasing on \mathbb{R}_+ . Hence, ϕ is strictly log-concave on \mathbb{R}_+ . Consequently, for any $u \in \mathbb{R}_+$,

$$(4.4) \quad 0 \leq u\phi'(u) = \phi(u) - m - u^2 \int_0^\infty e^{-uy} y \bar{\mu}(y) dy \leq \phi(u)$$

and

$$(4.5) \quad |\phi''(u)| \leq 2 \frac{\phi(u) - m}{u^2} \leq 2 \frac{\phi(u)}{u^2}.$$

(3) $\phi(u) \asymp \sigma^2 u + o(u)$ and $\phi'(u) \asymp \sigma^2 + o(1)$. Fix

$$a > d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\} \in (-\infty, 0],$$

then

$$|\phi(a + ib)| \cong \sigma^2 |a + ib| + o(|a + ib|),$$

as $|b| \rightarrow \infty$.

- (4) The mapping $u \mapsto \frac{1}{\phi(u)}$, $u \in \mathbb{R}_+$, is completely monotone, i.e. there exists a positive measure $\Upsilon(dy)$, whose support is contained in $[0, \infty)$, called the potential measure, such that the Laplace transform of Υ is given via the identity

$$\frac{1}{\phi(u)} = \int_0^\infty e^{-uy} \Upsilon(dy).$$

- (5) The mapping $u \mapsto \frac{u}{\phi(u)}$ is positive and log-concave on \mathbb{R}_+ .

(6) In any case,

$$(4.6) \quad \lim_{u \rightarrow \infty} \frac{\phi(u \pm \mathbf{a})}{\phi(u)} = 1 \quad \text{uniformly for } \mathbf{a}\text{-compact intervals on } \mathbb{R}_+.$$

(7) Uniformly, for $u \in \mathbb{R}_+$, we have that

$$(4.7) \quad \phi(u) \asymp u \int_0^{\frac{1}{u}} \bar{\mu}(y) dy + \sigma^2 u + m.$$

- (8) If $\psi \in \mathcal{N}$, then for $u \in \mathbb{R}_+$ put $\frac{\psi(u)}{u} = \phi(u)$. Then $\phi \in \mathcal{B}$. More precisely, $\psi'(0^+) = \phi(0) = m$, and $\mu(dy) = \bar{\Pi}(y)dy$, $y > 0$, where we recall that $\bar{\Pi}(y) = \int_y^\infty \Pi(dr)$ with Π the Lévy measure corresponding to ψ . Therefore,

$$\phi(u) = m + \sigma^2 u + \int_0^\infty (1 - e^{-uy}) \bar{\Pi}(y) dy$$

or equivalently $\phi \in \mathcal{B}_{\mathcal{N}}$, recall (1.12). Moreover, $\phi(\infty) < \infty$ if and only if $\sigma^2 = 0$ and $\bar{\Pi}(0^+) = \int_0^\infty \bar{\Pi}(y) dy < \infty$, and in fact then

$$(4.8) \quad \phi(\infty) = \mathbf{r} = m + \bar{\Pi}(0^+) < \infty.$$

- (9) Let $\beta > 0$. Then, for any $\phi \in \mathcal{B}$ (resp. $\phi \in \mathcal{B}_{\mathcal{N}}$ or $\psi \in \mathcal{N}$) $\mathcal{T}_\beta \phi(u) = \frac{u}{u+\beta} \phi(u+\beta) \in \mathcal{B}$ (resp. $\in \mathcal{B}_{\mathcal{N}}$ or $\mathcal{T}_\beta \psi(u) = \frac{u}{u+\beta} \psi(u+\beta) = u\phi(u+\beta) \in \mathcal{N}$). Moreover, $(\mathcal{T}_\beta \psi)'(0^+) = \phi(\beta)$ and the Lévy measure Π_β associated to $\mathcal{T}_\beta \psi \in \mathcal{N}$ is given by

$$(4.9) \quad \Pi_\beta(dy) = e^{-\beta y} (\Pi(dy) + \beta \bar{\Pi}(y) dy), \quad y > 0.$$

Proof. The proof of (1) is standard, see [Ber96, Chap. III]. Expression (4.2) gives (4.3). It together with (4.2) yield in turn the very first claim of item (2) and (4.4). Recall that $f > 0$ is strictly log-concave if $\log f$ is strictly concave and we verify that $(\log \phi)'' = \frac{\phi\phi'' - (\phi')^2}{\phi^2} < 0$ on \mathbb{R}_+ since ϕ' is a positive completely monotone function and thus ϕ'' is non-positive, see [Ber28] for more information. Differentiating the first expression of (4.3) we get that, for any $u > 0$,

$$\phi''(u) = -2 \int_0^\infty ye^{-uy} \bar{\mu}(y) dy + u \int_0^\infty y^2 e^{-uy} \bar{\mu}(y) dy.$$

However, since $\phi'' \leq 0$, then

$$|\phi''(u)| \leq 2 \int_0^\infty ye^{-uy} \bar{\mu}(y) dy.$$

Finally, using (4.3) we deduce that

$$|\phi''(u)| \leq 2 \int_0^\infty ye^{-uy} \bar{\mu}(y) dy = 2 \frac{\phi(u) - m}{u^2} - 2 \frac{\phi'(u)}{u},$$

and (4.5) follows from $\phi' \geq 0$ on \mathbb{R}_+ . The first claim of item (3) follows from (4.2) and (4.3), whereas the second one comes from the first expression for ϕ in (4.2) together with the Riemann-Lebesgue theorem applied to the integrable on \mathbb{R}_+ function $e^{-ay} \bar{\mu}(y)$. The fact that $d_\phi \in (-\infty, 0]$ is clear from $\lim_{u \rightarrow -\infty} \phi(u) = -\infty$, provided ϕ extends to an entire function. Item (4) can be found in [Ber96, Chap. III]. To prove item (5) we observe that the mapping $u \mapsto \frac{u}{\phi(u)}$ is positive on $(0, \infty)$ and its log-concavity property on $(0, \infty)$ is equivalent to

$$\begin{aligned} \frac{\phi(u)}{u} &= \frac{m}{u} + \sigma^2 + \int_0^\infty e^{-uy} \bar{\mu}(y) dy = \sigma^2 + \int_0^\infty e^{-uy} (\bar{\mu}(y) + m) dy \\ &= \sigma^2 + \mathcal{F}_{\bar{\mu}_m}^+(u) \end{aligned}$$

being log-convex on $(0, \infty)$, where for brevity $\bar{\mu}_m(y) = m + \bar{\mu}(y) \geq 0$, $y \in \mathbb{R}_+$, and $\mathcal{F}_{\bar{\mu}_m}^+$ stands for its Laplace transform. It therefore suffices to show that, for all $u > 0$,

$$(4.10) \quad \sigma^2 (\mathcal{F}_{\bar{\mu}_m}^+)'(u) + (\mathcal{F}_{\bar{\mu}_m}^+)'(u) \mathcal{F}_{\bar{\mu}_m}^+(u) - \left((\mathcal{F}_{\bar{\mu}_m}^+)'(u) \right)^2 > 0.$$

Noting that, for all $u > 0$, $\sigma^2 (\mathcal{F}_{\bar{\mu}_m}^+)'(u) \geq 0$, we deduce the inequality (4.10) from the Hölder's inequality which yields, for all $u > 0$, that

$$\begin{aligned} \left((\mathcal{F}_{\bar{\mu}_m}^+)'(u) \right)^2 &= \left(\int_0^\infty \left(\sqrt{\bar{\mu}_m(y)} e^{-\frac{u}{2}y} \right) \left(y \sqrt{\bar{\mu}_m(y)} e^{-\frac{u}{2}y} \right) dy \right)^2 \\ &< (\mathcal{F}_{\bar{\mu}_m}^+)'(u) \mathcal{F}_{\bar{\mu}_m}^+(u). \end{aligned}$$

Let us prove item (6), namely (4.6). It is obvious when $\sigma^2 > 0$ since item (3) holds. Let $\sigma^2 = 0$. We use the first relation in (4.2) and the monotonicity of ϕ to get, for $u > \mathfrak{a} > 0$,

$$\begin{aligned} 1 - \frac{\mathfrak{a}}{u} &\leq \frac{\left(1 - \frac{\mathfrak{a}}{u}\right) \int_0^\infty e^{-uy} \bar{\mu}(y) dy + \frac{m}{u}}{\int_0^\infty e^{-uy} \bar{\mu}(y) dy + \frac{m}{u}} = \frac{(u - \mathfrak{a}) \int_0^\infty e^{-uy} \bar{\mu}(y) dy + m}{\phi(u)} \\ &\leq \frac{\phi(u - \mathfrak{a})}{\phi(u)} \leq 1 \leq \frac{\phi(u + \mathfrak{a})}{\phi(u)} \\ &= \frac{(u + \mathfrak{a}) \int_0^\infty e^{-(u+\mathfrak{a})y} \bar{\mu}(y) dy + m}{\phi(u)} \leq \frac{\left(1 + \frac{\mathfrak{a}}{u}\right) \left(\int_0^\infty e^{-uy} \bar{\mu}(y) dy + \frac{m}{u}\right)}{\int_0^\infty e^{-uy} \bar{\mu}(y) dy + \frac{m}{u}} \\ &= 1 + \frac{\mathfrak{a}}{u}, \end{aligned}$$

and we deduce (4.6). Item (7) follows from [Ber96, Chapter III, Proposition 1]. Let us prove (4.8). Clearly, from item (3), $\phi(\infty) < \infty$ which implies that $\sigma^2 = 0$.

Then the substitution $z = \infty$ in the last expression in (4.2) and $\mu(dy) = \bar{\Pi}(y)dy$ yield that

$$\phi(\infty) = m + \mu(0, \infty) = m + \int_0^\infty \bar{\Pi}(y)dy = m + \bar{\Pi}(0^+).$$

Since from (1.3) $\psi(0) = 0$, and, $\psi(u) = u\phi(u)$ we get that $\psi'(0^+) = \phi(0^+) = m$. All other statements of item (8) follow from [Don07, p. 102, 9.4.7] and are standard for spectrally negative Lévy processes. A proof of item (9) can be found in [Pat11, Theorem 2.2] for $\beta = 1$, in [CKP10] or [PS12a, Proposition 2.1(2.2)] for $\beta > 0$. We only note that our definition (1.3) of ψ imposes $\Pi_+ \equiv 0$ and $\Pi_-(dy) = \Pi(-dy)$ for Π_\pm in [PS12a, Proposition 2.1(2.2)]. \square

The next claim focuses on further important for our work properties of the class $\mathcal{B}_{\mathcal{N}}$, see (1.12) for its definition. We recall that $\mathcal{N} = \mathcal{N}_\infty \cup \mathcal{N}_\infty^c$, $\mathcal{N}_\infty \cap \mathcal{N}_\infty^c = \emptyset$, where

$$\mathcal{N}_\infty = \{\psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \bar{\Pi}(0^+) = \infty\}$$

and \mathcal{N}_∞^c is its complement.

PROPOSITION 4.1.2. (1) Let $\phi \in \mathcal{B}_{\mathcal{N}}$ such that $\psi \in \mathcal{N}_\infty$ then

$$(4.11) \quad \lim_{u \rightarrow \infty} u^2 \frac{\phi'(u)}{\phi(u)} = \infty.$$

(2) Let $\phi \in \mathcal{B}_{\mathcal{N}}$ such that $\psi \in \mathcal{N}_\infty^c$ then

$$(4.12) \quad \lim_{u \rightarrow \infty} u^2 \frac{\phi'(u)}{\phi(u)} = \frac{\bar{\Pi}(0^+)}{\phi(\infty)} = \frac{\bar{\Pi}(0^+)}{m + \bar{\Pi}(0^+)} = \frac{\bar{\Pi}(0^+)}{\mathfrak{r}}.$$

Proof. When $\sigma^2 > 0$ (4.11) follows immediately from Proposition 4.1.1(3), i.e. $\phi(u) \asymp \sigma^2 u$ and $\phi'(u) \asymp \sigma^2$. Assume from now on that $\sigma^2 = 0$. Let $\underline{\lim}_{u \rightarrow \infty} \frac{u^2 \phi'(u)}{\phi(u)} < \infty$ and choose $C > 0$ and a sequence $(u_n)_{n \geq 1}$ tending to ∞ such that $\lim_{n \rightarrow \infty} \frac{u_n^2 \phi'(u_n)}{\phi(u_n)} = C^{-1}$. Then from this, the identity in (4.4) and the fact that when $\phi \in \mathcal{B}_{\mathcal{N}}$, $\mu(dy) = \bar{\Pi}(y)dy$ and $\bar{\mu}(y) = \bar{\bar{\Pi}}(y) = \int_y^\infty \bar{\Pi}(r)dr$, we deduce that

$$C u_n \phi'(u_n) \asymp \frac{\phi(u_n)}{u_n} = \phi'(u_n) + \frac{m}{u_n} + u_n \int_0^\infty e^{-y u_n} y \bar{\bar{\Pi}}(y) dy.$$

Therefore since trivially $\phi'(u_n) \asymp o(u_n \phi'(u_n))$ using the second formula for ϕ' in (4.3) we get from the last asymptotic relation that

$$\begin{aligned} C \phi'(u_n) &= C \int_0^\infty e^{-y u_n} y \bar{\bar{\Pi}}(y) dy \\ &\asymp \frac{m}{u_n^2} + \int_0^\infty e^{-y u_n} y \bar{\bar{\Pi}}(y) dy = \int_0^\infty e^{-y u_n} y (m + \bar{\bar{\Pi}}(y)) dy. \end{aligned}$$

However, this is impossible since $\bar{\Pi}$ and $\bar{\bar{\Pi}}$ being non-increasing on \mathbb{R}_+ determine the behaviour of the integrals above, as $u_n \rightarrow \infty$, solely via their local properties

at zero and when $\bar{\Pi}(0^+) = \infty$ then $m + \bar{\bar{\Pi}}(y) \stackrel{0}{=} o(\bar{\Pi}(y))$. Indeed, note that, for any $\epsilon > 0$,

$$\lim_{y \rightarrow 0} \frac{\bar{\bar{\Pi}}(y)}{\bar{\Pi}(y)} = \lim_{y \rightarrow 0} \frac{\int_y^\epsilon \bar{\Pi}(r) dr + \int_\epsilon^\infty \bar{\Pi}(r) dr}{\bar{\Pi}(y)} \leq \lim_{y \rightarrow 0} \frac{(\epsilon - y)\bar{\Pi}(y) + \int_\epsilon^\infty \bar{\Pi}(r) dr}{\bar{\Pi}(y)} \leq \epsilon,$$

and hence $\bar{\bar{\Pi}}(y) \stackrel{0}{=} o(\bar{\Pi}(y))$. Relation (4.12) follows immediately from the identity $u^2 \phi'(u) = u^2 \int_0^\infty e^{-uy} y \bar{\Pi}(y) dy$, which is the second relation in (4.3) with $\mu(dy) = \bar{\Pi}(y) dy$. \square

4.2. Products of Bernstein functions: new examples

There are many well known and fascinating mappings leaving invariant the set of Bernstein functions, that is \mathcal{B} , and we refer to [SSV10] for a nice account on these transformations. In this part, we show, through some substantial examples in our work, that products of some non-trivial subsets of Bernstein functions remain in the set of Bernstein functions. This simple transformation, which surprisingly does not seem to have been studied and used in the literature, plays a critical role in our development of the concept of reference semigroups. Indeed, this invariance allows us to identify a subset of gL semigroups which intertwines with a specific reference gL semigroup and whose intertwining operator is a bounded operator between weighted Hilbert spaces. Although we present this property for a two-parametric family of Bernstein functions, the approach can be easily extended to a more general framework and we believe that this idea of product factorization may be useful in a variety of contexts where the Bernstein functions appear. For example, since with each $\phi \in \mathcal{B}$ we associate a potential measure Υ on \mathbb{R}_+ via the identity

$$(4.13) \quad \frac{1}{\phi(u)} = \int_0^\infty e^{-uy} \Upsilon(dy), \quad u > 0,$$

see Proposition 4.1.1(4), if $\phi = \phi_1 \phi_2$ with $\phi_1, \phi_2 \in \mathcal{B}$, then $\Upsilon = \Upsilon_1 \star \Upsilon_2$, where \star stands for the additive convolution operator and Υ_1, Υ_2 are the potential measures associated to ϕ_1, ϕ_2 . Recall from (1.15) that, for any $\alpha \in (0, 1]$, $\mathbf{m} \geq 1 - \frac{1}{\alpha}$ and $u \geq 0$, we have

$$\phi_{\alpha, \mathbf{m}}^R(u) = \frac{\Gamma(\alpha u + \alpha \mathbf{m} + 1)}{\Gamma(\alpha u + \alpha \mathbf{m} + 1 - \alpha)} \quad \text{and} \quad \phi_{\mathbf{m}}^R(u) = \phi_{1, \mathbf{m}}^R(u) = u + \mathbf{m}.$$

We proceed with the following simple but useful result.

LEMMA 4.2.1. *Let $\alpha \in (0, 1)$ and $\mathbf{m} \geq 1 - \frac{1}{\alpha}$. Then*

$$(4.14) \quad \phi_{\alpha, \mathbf{m}}^R(u) = \int_0^\infty (1 - e^{-uy}) e^{-(\mathbf{m} + \frac{1}{\alpha})y} (1 - e^{-\frac{y}{\alpha}})^{-\alpha-1} \frac{dy}{\Gamma(1 - \alpha)} + \frac{\Gamma(\alpha \mathbf{m} + 1)}{\Gamma(\alpha \mathbf{m} + 1 - \alpha)} \in \mathcal{B}_{\mathcal{N}},$$

and its associated potential measure defined via (4.13) is absolutely continuous with a density given, for any $y > 0$, by

$$(4.15) \quad U_{\alpha, \mathbf{m}}(y) = \frac{e^{-\mathbf{m}_\alpha y} (1 - e^{-\frac{y}{\alpha}})^{\alpha-1}}{\Gamma(\alpha + 1)}$$

where $\mathbf{m}_\alpha = (\alpha\mathbf{m} + 1 - \alpha)/\alpha \geq 0$. Moreover, $U_{\alpha,\mathbf{m}}$ is non-increasing, convex on \mathbb{R}_+ and solves on \mathbb{R}_+ the differential equation

$$(4.16) \quad U'_{\alpha,\mathbf{m}}(y) = -U_{\alpha,\mathbf{m}}(y) \left(\mathbf{m}_\alpha + \frac{1-\alpha}{\alpha} (e^{\frac{y}{\alpha}} - 1)^{-1} \right).$$

Proof. First, from the integral representation of the Beta function \mathbf{B} , see [Leb72, (1.5.2) p.13], we get, for any $\alpha \in (0, 1)$ and $u > 0$,

$$\begin{aligned} \frac{\Gamma(u + \alpha)}{\Gamma(u)} &= \frac{u\mathbf{B}(u + \alpha, 1 - \alpha)}{\Gamma(1 - \alpha)} = \frac{u}{\Gamma(1 - \alpha)} \int_0^1 r^{u+\alpha-1} (1-r)^{-\alpha} dr \\ &= \frac{u}{\alpha\Gamma(1 - \alpha)} \int_0^\infty e^{-(\frac{u}{\alpha}+1)y} \left(1 - e^{-\frac{y}{\alpha}}\right)^{-\alpha} dy \\ &= \frac{1}{\Gamma(1 - \alpha)} \int_0^\infty e^{-y} \left(1 - e^{-\frac{y}{\alpha}}\right)^{-\alpha} d\left(1 - e^{-\frac{y}{\alpha}}\right) \\ &= \frac{1}{\Gamma(1 - \alpha)} \int_0^\infty \left(1 - e^{-\frac{y}{\alpha}}\right) e^{-y} \left(1 - e^{-\frac{y}{\alpha}}\right)^{-\alpha-1} dy. \end{aligned}$$

Using the above relation twice with $\alpha u + \alpha\mathbf{m} - \alpha + 1$ and $\alpha\mathbf{m} - \alpha + 1$ standing for u , we get, after some easy algebra, that

$$\frac{\Gamma(\alpha u + \alpha\mathbf{m} + 1)}{\Gamma(\alpha u + \alpha\mathbf{m} + 1 - \alpha)} - \frac{\Gamma(\alpha\mathbf{m} + 1)}{\Gamma(\alpha\mathbf{m} + 1 - \alpha)} = \frac{\int_0^\infty (1 - e^{-uy}) e^{-(\mathbf{m} + \frac{1}{\alpha})y} (1 - e^{-\frac{y}{\alpha}})^{-\alpha-1} dy}{\Gamma(1 - \alpha)}.$$

Thus we deduce the first claim after having easily checked that $y \mapsto e^{-(\mathbf{m} + \frac{1}{\alpha})y} (1 - e^{-\frac{y}{\alpha}})^{-\alpha-1}$ is non-increasing on \mathbb{R}_+ . Next, from the integral representation of the Beta function again, valid here for any $\alpha u + \alpha\mathbf{m} - \alpha + 1 > 0$, we get that

$$\begin{aligned} \int_0^\infty e^{-uy} e^{-\mathbf{m}_\alpha y} (1 - e^{-\frac{y}{\alpha}})^{\alpha-1} dy &= \alpha \int_0^1 v^{\alpha u + \alpha\mathbf{m} - \alpha} (1-v)^{\alpha-1} dv \\ &= \frac{\Gamma(\alpha + 1)\Gamma(\alpha u + \alpha\mathbf{m} + 1 - \alpha)}{\Gamma(\alpha u + \alpha\mathbf{m} + 1)} = \frac{\Gamma(\alpha + 1)}{\phi_{\alpha,\mathbf{m}}^R(u)} \end{aligned}$$

and (4.15) follows from (4.13). The other claims are obvious. \square

For any $\phi \in \mathcal{B}$ set

$$\Phi_{\alpha,\mathbf{m}}(u) = \frac{\phi(u)}{\phi_{\alpha,\mathbf{m}}^R(u)}, \quad u \geq 0.$$

The next statement furnishes a set of sufficient conditions on $\phi \in \mathcal{B}$ for which $\Phi_{\alpha,\mathbf{m}} \in \mathcal{B}$. Recall that $\mathcal{N}_P = \{\psi \in \mathcal{N}; \sigma^2 > 0\}$.

PROPOSITION 4.2.2. (1) Let $\psi \in \mathcal{N}_P$ with $\overline{\Pi}(0^+) < \infty$. Then for any $\mathbf{m} > \underline{\mathbf{m}} = \frac{m + \overline{\Pi}(0^+)}{\sigma^2}$ the mapping

$$\Phi_{\mathbf{m}}(u) = \frac{\phi(u)}{\phi_{\mathbf{m}}^R(u)} = \frac{\phi(u)}{u + \mathbf{m}} \in \mathcal{B}.$$

Moreover, if there exists $\mathbf{u}_0 > 0$ such that $\psi(-\mathbf{u}_0) = 0$ then $\mathbf{u}_0 \leq \frac{m}{\sigma^2}$ with identity if and only if $\overline{\Pi}(0^+) = 0$.

(2) If $\psi \in \mathcal{N}_P$ then, for any $\alpha \in (0, 1)$,

$$\underline{y}_\alpha = \inf \left\{ y \geq 0; (e^{\frac{y}{\alpha}} - 1) \overline{\Pi} \left(\frac{y}{2} \right) > \sigma^2 \frac{1 - \alpha}{\alpha} \right\} \in (0, \infty]$$

and $\Phi_{\alpha, \mathbf{m}} \in \mathcal{B}$ for any $\mathbf{m} > \frac{\overline{\Pi}(\frac{y_\alpha}{2}) + m}{\sigma^2} + 1 - \frac{1}{\alpha}$. Otherwise, if $\psi \in \mathcal{N} \setminus \mathcal{N}_P$ and there exist $\alpha \in (0, 1)$, $\mathbf{m} \geq 1 - \frac{1}{\alpha}$ such that for every $y > 0$ there is $A \in (0, 1)$

$$\frac{\overline{\Pi}(y) + m}{\overline{\Pi}((1 - A)y) + m} + \frac{U_{\alpha, \mathbf{m}}(y)}{U_{\alpha, \mathbf{m}}(Ay)} \leq 1$$

then $\Phi_{\alpha, \mathbf{m}} \in \mathcal{B}$.

Proof. From Proposition 4.1.1(3)

$$\lim_{u \rightarrow \infty} \Phi_{\mathbf{m}}(u) = \lim_{u \rightarrow \infty} \frac{\phi(u)}{\phi_{\mathbf{m}}^R(u)} = \lim_{u \rightarrow \infty} \frac{\phi(u)}{u + \mathbf{m}} = \sigma^2$$

and $\Phi_{\mathbf{m}}(0) = \frac{m}{\mathbf{m}} \geq 0$. Put $e_{\mathbf{m}}(y) = e^{-\mathbf{m}y}$, $y \geq 0$. Next, an integration by parts yields

$$\begin{aligned} \frac{\phi(u)}{u + \mathbf{m}} &= \frac{m}{u + \mathbf{m}} + \sigma^2 \frac{u}{u + \mathbf{m}} + \frac{u}{u + \mathbf{m}} \int_0^\infty e^{-uy} \overline{\Pi}(y) dy \\ &= \frac{m}{u + \mathbf{m}} + \sigma^2 \frac{u}{u + \mathbf{m}} + u \int_0^\infty e^{-uy} \overline{\Pi} \star e_{\mathbf{m}}(y) dy \\ (4.17) \quad &= \frac{m}{\mathbf{m}} + \int_0^\infty (1 - e^{-uy}) e^{-\mathbf{m}y} \left(\mathbf{m}\sigma^2 - m + \mathbf{m} \int_0^y e^{\mathbf{m}r} \overline{\Pi}(r) dr - e^{\mathbf{m}y} \overline{\Pi}(y) \right) dy. \end{aligned}$$

Since $\overline{\Pi}$ is non-increasing, we have, for all $y \geq 0$, $\mathbf{m} \int_0^y e^{\mathbf{m}r} \overline{\Pi}(r) dr \geq \overline{\Pi}(y) (e^{\mathbf{m}y} - 1)$ and thus

$$\mathbf{m}\sigma^2 - m + \mathbf{m} \int_0^y e^{\mathbf{m}r} \overline{\Pi}(r) dr - e^{\mathbf{m}y} \overline{\Pi}(y) \geq \mathbf{m}\sigma^2 - m - \overline{\Pi}(y) \geq \mathbf{m}\sigma^2 - m - \overline{\Pi}(0^+).$$

The latter expression is positive and the first multiplied by $e^{-\mathbf{m}y}$ is a density of a Lévy measure provided $\mathbf{m} > \underline{\mathbf{m}} = \frac{m + \overline{\Pi}(0^+)}{\sigma^2}$ and since then the last expression in (4.17) defines a Bernstein function the first claim of item (1) follows. Next, since $u_0 > 0$ and $\psi(-u_0) = -u_0 \phi(-u_0) = 0$, an application of the first identity in (4.2) with $\bar{\mu}(y) = \overline{\Pi}(y)$ since $\phi \in \mathcal{B}_{\mathcal{N}}$, yields that

$$\sigma^2 u_0 = \phi(-u_0) + \sigma^2 u_0 = m - u_0 \int_0^\infty e^{u_0 y} \overline{\Pi}(y) dy \leq m.$$

This completes the proof of item (1). Next, assume that $\sigma^2 > 0$. Our general assumption $\int_0^\infty (y^2 \wedge y) \Pi(dy) < \infty$ implies via twice integration by parts that

$$(4.18) \quad \int_0^1 \overline{\Pi}(y) dy = \int_0^1 y^2 \Pi(dy) + \overline{\Pi}(1) + \frac{1}{2} \overline{\Pi}(1) < \infty.$$

Then the fact that $\overline{\Pi}$ is non-increasing on \mathbb{R}^+ triggers from (4.18) that

$$(e^{\frac{y}{\alpha}} - 1) \overline{\Pi} \left(\frac{y}{2} \right) \stackrel{0}{\sim} \frac{y}{\alpha} \overline{\Pi} \left(\frac{y}{2} \right) \stackrel{0}{=} o(1)$$

and hence $\underline{y}_\alpha \in (0, \infty]$. Choose first $\mathbf{m}_\alpha = \alpha(\mathbf{m} + 1 - \alpha) > 0$ and observe from (4.13) and the definition of $U_{\alpha, \mathbf{m}}$ that

$$\begin{aligned} \Phi_{\alpha, \mathbf{m}}(u) &= \frac{\phi(u)}{\phi_{\alpha, \mathbf{m}}^R(z)} = \phi(u) \frac{\Gamma(\alpha u + \alpha \mathbf{m} + 1 - \alpha)}{\Gamma(\alpha u + \alpha \mathbf{m} + 1)} \\ &= \left(m + \sigma^2 u + u \int_0^\infty e^{-uy} \overline{\overline{\Pi}}(y) dy \right) \int_0^\infty e^{-uy} U_{\alpha, \mathbf{m}}(y) dy \\ &= m \int_0^\infty U_{\alpha, \mathbf{m}}(y) dy + \int_0^\infty (1 - e^{-uy}) (-m U_{\alpha, \mathbf{m}}(y) - \sigma^2 U'_{\alpha, \mathbf{m}}(y)) dy \\ &\quad + u \int_0^\infty e^{-uy} \overline{\overline{\Pi}} \star U_{\alpha, \mathbf{m}}(y) dy, \end{aligned}$$

where $0 < \int_0^\infty U_{\alpha, \mathbf{m}}(y) dy < \infty$ and the terms of the integration by parts vanish since from (4.15) we have that

$$U_{\alpha, \mathbf{m}}(y) \stackrel{0}{\sim} \frac{\alpha^{\alpha-1}}{\Gamma(1+\alpha)} y^{\alpha-1}, \quad U_{\alpha, \mathbf{m}}(y) \stackrel{\infty}{\sim} \frac{1}{\Gamma(\alpha+1)} e^{-\mathbf{m}_\alpha y}$$

and $\mathbf{m}_\alpha > 0$. Using (4.16) put, for $y > 0$,

$$\begin{aligned} -m U_{\alpha, \mathbf{m}}(y) - \sigma^2 U'_{\alpha, \mathbf{m}}(y) &= \left(\sigma^2 \mathbf{m}_\alpha - m + \sigma^2 \frac{1-\alpha}{\alpha} (e^{\frac{y}{\alpha}} - 1)^{-1} \right) U_{\alpha, \mathbf{m}}(y) \\ (4.19) \qquad \qquad \qquad &= \bar{u}_{\alpha, \mathbf{m}}(y) U_{\alpha, \mathbf{m}}(y). \end{aligned}$$

Then another integration by parts for the very last term in $\Phi_{\alpha, \mathbf{m}}$ above yields

$$\begin{aligned} \Phi_{\alpha, \mathbf{m}}(u) &= m \int_0^\infty U_{\alpha, \mathbf{m}}(y) dy \\ (4.20) \qquad \qquad \qquad &+ \int_0^\infty (1 - e^{-uy}) \left(\bar{u}_{\alpha, \mathbf{m}}(y) U_{\alpha, \mathbf{m}}(y) - \left(\overline{\overline{\Pi}} \star U_{\alpha, \mathbf{m}}(y) \right)' \right) dy. \end{aligned}$$

Indeed, the asymptotic relations for $U_{\alpha, \mathbf{m}}$ above allow to deduct that the boundary terms in this integration by parts do not contribute since from $1 - e^{-uy} \stackrel{0}{\sim} uy$, $1 - e^{-uy} \stackrel{\infty}{\sim} 1$,

$$\begin{aligned} &\lim_{y \rightarrow 0} y \int_0^y \overline{\overline{\Pi}}(y-r) U_{\alpha, \mathbf{m}}(r) dr \\ (4.21) \qquad \leq &\lim_{y \rightarrow 0} \left(y \overline{\overline{\Pi}}\left(\frac{y}{2}\right) + y \int_0^{\frac{y}{2}} \overline{\overline{\Pi}}(r) dr \sup_{r \in [\frac{y}{2}, y]} U_{\alpha, \mathbf{m}}(r) \right) = 0 \end{aligned}$$

and

$$\begin{aligned} &\lim_{y \rightarrow \infty} \int_0^y \overline{\overline{\Pi}}(r) U_{\alpha, \mathbf{m}}(y-r) dr \\ &\leq \lim_{y \rightarrow \infty} \left(\overline{\overline{\Pi}}\left(\frac{y}{2}\right) \int_{\frac{y}{2}}^y U_{\alpha, \mathbf{m}}(y-r) dr + e^{-\mathbf{m}_\alpha \frac{y}{2}} \int_0^{\frac{y}{2}} \overline{\overline{\Pi}}(r) dr \right) \\ (4.22) \qquad \leq &\int_0^\infty U_{\alpha, \mathbf{m}}(r) dr \lim_{y \rightarrow \infty} \overline{\overline{\Pi}}(y) + \lim_{y \rightarrow \infty} e^{-\mathbf{m}_\alpha \frac{y}{2}} \left(\int_0^1 \overline{\overline{\Pi}}(r) dr + \overline{\overline{\Pi}}(1)y \right) = 0. \end{aligned}$$

Next, choose $\mathbf{m} > 1 - \frac{1}{\alpha}$ so large such that $\mathbf{m}_\alpha = (\alpha \mathbf{m} + 1 - \alpha)/\alpha > \frac{\overline{\overline{\Pi}}(\frac{y}{2}) + m}{\sigma^2}$ and hence we get from (4.19) and the definition of \underline{y}_α that $\bar{u}_{\alpha, \mathbf{m}} > 0$ on \mathbb{R}_+ . Recalling the

expression of $\Phi_{\alpha, m}$ in (4.20) we aim to show that $\bar{u}_{\alpha, m}(y)U_{\alpha, m}(y) - \left(\bar{\Pi} \star U_{\alpha, m}(y)\right)'$ defines a density of a Lévy measure. For this purpose, for any $A \in (0, 1)$, put $y_A = Ay$ and $\bar{y}_A = (1 - A)y$. Then

$$\begin{aligned}
\left(\bar{\Pi} \star U_{\alpha, m}(y)\right)' &= \left(\int_0^{y_A} \bar{\Pi}(y-r)U_{\alpha, m}(r)dr + \int_{y_A}^y \bar{\Pi}(y-r)U_{\alpha, m}(r)dr\right)' \\
&= \left(\int_0^{y_A} \bar{\Pi}(y-r)U_{\alpha, m}(r)dr + \int_0^{\bar{y}_A} \bar{\Pi}(r)U_{\alpha, m}(y-r)dr\right)' \\
&= \bar{\Pi}(\bar{y}_A)U_{\alpha, m}(y_A) - \int_0^{y_A} \bar{\Pi}(y-r)U_{\alpha, m}(r)dr \\
(4.23) \quad &+ \int_0^{\bar{y}_A} \bar{\Pi}(r)U'_{\alpha, m}(y-r)dr.
\end{aligned}$$

Since $\bar{\Pi}(r) = \int_r^\infty \bar{\Pi}(v)dv$, $r > 0$, and as $U_{\alpha, m}$ is non-increasing, we get that

$$\begin{aligned}
&\bar{\Pi}(\bar{y}_A)U_{\alpha, m}(y_A) - \int_0^{y_A} \bar{\Pi}(y-r)U_{\alpha, m}(r)dr \\
&\leq \bar{\Pi}(\bar{y}_A)U_{\alpha, m}(y_A) - U_{\alpha, m}(y_A) \int_0^{y_A} \bar{\Pi}(y-r)dr \\
&= \bar{\Pi}(y)U_{\alpha, m}(y_A).
\end{aligned}$$

Thus, since $U'_{\alpha, m} < 0$ on \mathbb{R}_+ , see (4.16), from (4.23) we deduct that

$$\begin{aligned}
(4.24) \quad \left(\bar{\Pi} \star U_{\alpha, m}(y)\right)' &\leq \bar{\Pi}(y)U_{\alpha, m}(y_A) + \int_0^{\bar{y}_A} \bar{\Pi}(r)U'_{\alpha, m}(y-r)dr \\
&\leq \bar{\Pi}(y)U_{\alpha, m}(y_A) + \bar{\Pi}(\bar{y}_A)(U_{\alpha, m}(y) - U_{\alpha, m}(y_A)).
\end{aligned}$$

Next, from (4.24) we observe that, with

$$F_A(y) = \left(\bar{\Pi}(y) - \bar{\Pi}(\bar{y}_A)\right)U_{\alpha, m}(y_A) + \left(\bar{\Pi}(\bar{y}_A) - \bar{u}_{\alpha, m}(y)\right)U_{\alpha, m}(y),$$

we have that

$$(4.25) \quad \left(\bar{\Pi} \star U_{\alpha, m}(y)\right)' - \bar{u}_{\alpha, m}(y)U_{\alpha, m}(y) \leq F_A(y) \leq \left(\bar{\Pi}(\bar{y}_A) - \bar{u}_{\alpha, m}(y)\right)U_{\alpha, m}(y).$$

Choose $A = \frac{1}{2}$ and thus $\bar{y}_A = y_A = y/2$. As long as

$$m_\alpha = (\alpha m + 1 - \alpha)/\alpha > \frac{\bar{\Pi}\left(\frac{y}{2}\right) + m}{\sigma^2}$$

due to the definition of $\bar{u}_{\alpha, m}(y)$, see (4.19), and (4.25), we have, for all $y > 0$, that

$$\begin{aligned}
&\left(\bar{\Pi} \star U_{\alpha, m}(y)\right)' - \bar{u}_{\alpha, m}(y)U_{\alpha, m}(y) \\
&\leq F_{\frac{1}{2}}(y) \leq \left(\bar{\Pi}\left(\frac{y}{2}\right) - \bar{u}_{\alpha, m}(y)\right)U_{\alpha, m}(y) \\
&< \left(\left(\bar{\Pi}\left(\frac{y}{2}\right) - \bar{\Pi}\left(\frac{y}{2}\right)\right) - \sigma^2 \frac{1 - \alpha}{\alpha \left(e^{\frac{y}{\alpha}} - 1\right)}\right)U_{\alpha, m}(y).
\end{aligned}$$

Obviously from the fact that $\bar{\Pi}$ is non-increasing and $\alpha < 1$ the right-hand side is non-positive for $y \geq \underline{y}_\alpha$, whereas it is also non-positive for all $y < \underline{y}_\alpha$ thanks to the

definition of $\underline{y}_\alpha = \inf \left\{ y \geq 0; (e^{\frac{y}{\alpha}} - 1)\overline{\Pi}\left(\frac{y}{2}\right) > \sigma^2 \frac{1-\alpha}{\alpha} \right\} \in (0, \infty]$. Therefore,

$$f(y) = - \left(\overline{\Pi} \star U_{\alpha, m}(y) \right)' + \bar{u}_{\alpha, m}(y) U_{\alpha, m}(y) \geq 0.$$

To show that $f(y)dy$ defines a Lévy measure of a Bernstein function we first observe from (4.15) and (4.19) that

$$\lim_{y \rightarrow 0} y \bar{u}_{\alpha, m}(y) = \sigma^2(1 - \alpha), \quad \lim_{y \rightarrow \infty} \bar{u}_{\alpha, m}(y) = \sigma^2 \mathbf{m}_\alpha - m > 0, \quad U_{\alpha, m}(y) \asymp \frac{e^{-\mathbf{m}_\alpha y}}{\Gamma(\alpha + 1)}$$

and these imply that

$$\int_0^\infty (1 \wedge y) \bar{u}_{\alpha, m}(y) U_{\alpha, m}(y) dy < \infty.$$

Secondly, (4.21) and (4.22) allow via integration by parts to obtain that

$$\begin{aligned} \left| \int_0^\infty (1 \wedge y) \left(\overline{\Pi} \star U_{\alpha, m}(y) \right)' dy \right| &= \int_0^1 \overline{\Pi} \star U_{\alpha, m}(y) dy \\ &= \int_0^1 U_{\alpha, m}(y) \int_y^1 \overline{\Pi}(r - y) dr dy \\ &\leq \int_0^1 U_{\alpha, m}(y) dy \int_0^1 \overline{\Pi}(y) dy < \infty. \end{aligned}$$

Indeed, the finiteness of $\int_0^1 U_{\alpha, m}(y) dy$ has been discussed above, whereas $\int_0^1 \overline{\Pi}(y) dy < \infty$ is (4.18). Thus $\int_0^\infty (1 \wedge y) f(y) dy < \infty$. Therefore, $\Phi_{\alpha, m}$ is a Bernstein function with Lévy measure $f(y)dy$. When $\sigma^2 = 0$ we have that $\bar{u}_{\alpha, m}(y) = -m, \forall y > 0$, see (4.19). Using this from the first inequality of (4.25) we then get that, for any $0 < A < 1$,

$$\begin{aligned} &\left(\overline{\Pi} \star U_{\alpha, m}(y) \right)' - \bar{u}_{\alpha, m}(y) U_{\alpha, m}(y) \\ &\leq \left(\overline{\Pi}(\bar{y}_A) + m \right) U_{\alpha, m}(\bar{y}_A) \left(\frac{\overline{\Pi}(y) + m}{\overline{\Pi}(\bar{y}_A) + m} - 1 + \frac{U_{\alpha, m}(y)}{U_{\alpha, m}(\bar{y}_A)} \right). \end{aligned}$$

Since $f(y) = - \left(\overline{\Pi} \star U_{\alpha, m}(y) \right)' + \bar{u}_{\alpha, m}(y) U_{\alpha, m}(y), \forall y > 0$, then $\int_0^\infty (1 \wedge y) |f(y)| dy < \infty$ follows as in the case $\sigma^2 > 0$. It defines a Lévy measure if the third factor above is non-positive for all $y > 0$. However, by assumption for every $y > 0$ we can choose $A \in (0, 1)$ such that

$$\left(\frac{\overline{\Pi}(y) + m}{\overline{\Pi}(\bar{y}_A) + m} - 1 + \frac{U_{\alpha, m}(y)}{U_{\alpha, m}(\bar{y}_A)} \right) \leq 0$$

and therefore $f(y) \leq 0$. □

4.3. Useful estimates of Bernstein functions on \mathbb{C}_+

In this part we derive estimates for some functionals of Bernstein functions. We introduce the notation

$$\Delta_b f(a) = f(a + ib) - f(a),$$

$$\Re(f(a + ib) - f(a)) + i \Im(f(a + ib) - f(a)) = \Delta_b^{\Re} f(a) + i \Delta_b^{\Im} f(a),$$

and, we recall that, for any $k \geq 1$, $f^{(k)}(x) = \frac{d^k}{dx^k} f(x)$.

LEMMA 4.3.1. *Let $\phi \in \mathcal{B}$.*

(1) *Let $b \in \mathbb{R}$ and $a > 0$. Then,*

$$(4.26) \quad 0 \leq \Delta_b^{\Re} \phi(a) \leq \frac{b^2}{2} |\phi''(a)| \quad \text{and} \quad |\Delta_b^{\Im} \phi(a)| \leq |b| |\phi'(a)|,$$

and, for $k \geq 1$,

$$(4.27) \quad \left| \Delta_b^{\Re} \phi^{(k)}(a) \right| \leq 2 \left| \phi^{(k)}(a) \right| \quad \text{and} \quad \left| \Delta_b^{\Im} \phi^{(k)}(a) \right| \leq \left| \phi^{(k)}(a) \right|.$$

(2) *Finally, we have, for any $u > 0$,*

$$(4.28) \quad \begin{aligned} \int_u^\infty \frac{|\phi''(y+ib)|}{|\phi(y+ib)|} dy &\leq \sqrt{10} \int_u^\infty \frac{|\phi''(y)|}{\phi(y)} dy \leq \frac{2\sqrt{10}}{u}, \\ \int_u^\infty \frac{|\phi'(y+ib)|^2}{|\phi(y+ib)|^2} dy &\leq 10 \int_u^\infty \left(\frac{\phi'(y)}{\phi(y)} \right)^2 dy \leq \frac{10}{u}. \end{aligned}$$

Proof. First, using the inequality $1 - \cos(y) \leq \frac{y^2}{2}$, we get that

$$\begin{aligned} \Delta_b^{\Re} \phi(a) &= \int_0^\infty (1 - \cos(by)) e^{-ay} \mu(dy) \\ &\leq \frac{b^2}{2} \int_0^\infty y^2 e^{-ay} \mu(dy) = \frac{b^2}{2} |\phi''(a)| \end{aligned}$$

and the first inequality in (4.26) follows. Similarly, for any $k \geq 1$, we have

$$\begin{aligned} \left| \Delta_b^{\Re} \phi^{(k)}(a) \right| &= \int_0^\infty (1 - \cos(by)) y^k e^{-ay} \mu(dy) \leq 2 \int_0^\infty y^k e^{-ay} \mu(dy) \\ &\leq 2 \left| \phi^{(k)}(a) \right|, \end{aligned}$$

which provides the first claim of (4.27) since, for $k \geq 1$, $y^k e^{-ay} \mu(dy)$ is integrable on \mathbb{R}_+ . The imaginary part estimates, that is the second claims of (4.26) and (4.27), follow by similar computations completing the proof of the first item. The second inequality in the first and second lines of (4.28) follows from $\phi'(a)/\phi(a) \leq 1/a$ and $\phi''(a)/\phi(a) \leq 2/a^2$ according to (4.4) and (4.5). Let now $b \neq 0$ and write

$$\frac{|\phi''(a+ib)|}{|\phi(a+ib)|} = \frac{|\phi''(a)| \left| 1 + \frac{\Delta_b \phi''(a)}{\phi''(a)} \right|}{\phi(a) \left| 1 + \frac{\Delta_b \phi(a)}{\phi(a)} \right|} \stackrel{(4.27)}{\leq} \frac{|\phi''(a)|}{\phi(a)} \frac{\sqrt{10}}{1 + \frac{\Delta_b^{\Re} \phi(a)}{\phi(a)}} \stackrel{(4.26)}{\leq} \sqrt{10} \frac{|\phi''(a)|}{\phi(a)}.$$

This gives the first inequality in the first line of (4.28). To conclude the proof of the lemma for the first inequality of the second line of (4.28) we use (4.27) to get that

$$\frac{|\phi'(a+ib)|}{|\phi(a+ib)|} = \frac{\phi'(a) \left| 1 + \frac{\Delta_b \phi'(a)}{\phi'(a)} \right|}{\phi(a) \left| 1 + \frac{\Delta_b \phi(a)}{\phi(a)} \right|} \stackrel{(4.27)}{\leq} \frac{\phi'(a)}{\phi(a)} \frac{\sqrt{10}}{1 + \frac{\Delta_b^{\Re} \phi(a)}{\phi(a)}} \stackrel{(4.26)}{\leq} \sqrt{10} \frac{\phi'(a)}{\phi(a)}.$$

□

The next result provides additional estimates about some specific quantities.

LEMMA 4.3.2. *Let $\phi \in \mathcal{B}_{\mathcal{N}}$. Then, for $a > 0$, $b > 0$, and some constants $0 < C < D < \infty$,*

$$\begin{aligned} \sigma^2 b + (e^{-\pi a} - e^{-2\pi a}) \int_0^{\frac{\pi}{b}} \sin(by) \bar{\Pi}(y) dy &\leq \Delta_b^{\Im} \phi(ba) \leq \sigma^2 b + b \int_0^{\frac{\pi}{b}} y \bar{\Pi}(y) dy, \\ \frac{Cb^2}{e^a} \int_0^{\frac{1}{b}} y^2 \bar{\Pi}(y) dy &\leq \Delta_b^{\Re} \phi(ba) \leq Db \int_0^{\frac{\pi}{b}} y \bar{\Pi}(y) dy + \frac{2}{e^a} \bar{\Pi}\left(\frac{1}{b}\right), \\ (4.29) \quad \Delta_b^{\Re} \phi(ba) &\leq b^2 |\phi''(ba)| \quad \text{and} \quad \Delta_b^{\Im} \phi(ba) \leq b\phi'(ba). \end{aligned}$$

Proof. We set $\sigma^2 = 0$ as otherwise we simply add $\sigma^2 b$ in the first line of (4.29). Then, splitting in the periods of $\sin(by)$ and using the fact that $\bar{\Pi}$ is non-increasing we get that

$$\begin{aligned} \Delta_b^{\Im} \phi(ba) &= \int_0^{\infty} \sin(by) e^{-bya} \bar{\Pi}(y) dy \\ &= \sum_{k=0}^{\infty} e^{-2k\pi a} \int_0^{\frac{\pi}{b}} \frac{\sin(by)}{e^{bya}} \left(\bar{\Pi}\left(y + \frac{2k\pi}{b}\right) - \frac{\bar{\Pi}\left(y + \frac{(2k+1)\pi}{b}\right)}{e^{\pi a}} \right) dy \\ &\geq 0. \end{aligned}$$

Furthermore, picking the term when $k = 0$ proves the left-hand side of (4.29) since $\bar{\Pi}$ is a non-increasing function. For the upper bound, we use the following estimates obtained from the expression above by using the properties of $\bar{\Pi}$,

$$\begin{aligned} \Delta_b^{\Im} \phi(ba) &\leq \sum_{k=0}^{\infty} \int_0^{\frac{\pi}{b}} \frac{\sin(by)}{e^{bya}} \left(\frac{\bar{\Pi}\left(y + \frac{2k\pi}{b}\right)}{e^{2k\pi a}} - \frac{\bar{\Pi}\left(y + \frac{(2k+2)\pi}{b}\right)}{e^{2(k+1)\pi a}} \right) dy \\ &= \int_0^{\frac{\pi}{b}} \sin(by) e^{-bya} \bar{\Pi}(y) dy \leq \int_0^{\frac{\pi}{b}} \sin(by) \bar{\Pi}(y) dy \\ &\leq b \int_0^{\frac{\pi}{b}} y \bar{\Pi}(y) dy \end{aligned}$$

and we achieve the first part of (4.29). The second part is trivial. The last statement follows by just considering (4.26) with a replaced by ab . \square

Fine properties of the density of the invariant measure

The development of the spectral expansion of gL semigroups requires a variety of detailed information on the density ν of the invariant measure. For instance, the existence of co-eigenfunctions, that is when does, for some $n \in \mathbb{N}$, $\mathcal{V}_n(x) = \frac{(x^n \nu(x))^{(n)}}{n! \nu(x)} \in L^2(\nu)$?, hinges on smoothness properties and precise estimates for the large and small asymptotic behaviour of ν along with its successive derivatives. Unfortunately, the only information on the invariant measure that one can easily extract from the literature is the expression of its entire moments, see (5.14) below, from which it seems delicate to derive the sought fine distributional properties. To overcome this difficulty, we shall point out, see Proposition 5.1.1 below, that the set of invariant measures of gL semigroups is in fact closely connected to a subset of the class of distributions of the positive self-decomposable variables, a substantial family of random variables which has been thoroughly studied in the literature. We are going to take advantage of this relationship to derive for ν some of the properties mentioned above. However, for our purpose, we shall need to deepen in different directions the study of this subset of self-decomposable variables obtaining results of independent interest. More specifically, in this Chapter, we derive new fine distributional properties including the small and large asymptotic behaviour of the densities along with the successive derivatives of this subset of self-decomposable variables. We mention that the results presented here will be used at several places throughout the rest of the paper, justifying our choice to gather them in one chapter.

We start by stating the following series of substantial results on the density ν of the distribution of the positive variable $V_\psi, \psi \in \mathcal{N}$, whose law is, according to Proposition 2.3.1(1), the invariant measure of the associated generalized Laguerre (gL) semigroup. The Mellin transform of V_ψ is denoted by \mathcal{M}_{V_ψ} , i.e. $\mathcal{M}_{V_\psi}(z) = \int_0^\infty x^{z-1} \nu(x) dx, z \in 1 + i\mathbb{R}$. We also recall that $\tau = \phi(\infty)$ with $\phi(\infty) = \infty$ if $\sigma^2 > 0$ or $\overline{\Pi}(0^+) = \infty$, or $\phi(\infty) = \overline{\Pi}(0^+) + m$ otherwise, $N_\tau = \left\lceil \frac{\overline{\Pi}(0^+)}{\tau} \right\rceil - 1 \in [0, \infty]$ with the convention that $N_\tau = \infty$ when $\overline{\Pi}(0^+) = \infty$ and $d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\}$.

THEOREM 5.0.1. *Let $\psi \in \mathcal{N}$ and recall that $\psi(z) = z\phi(z)$ with $\phi \in \mathcal{B}_\mathcal{N}$.*

- (1) $\mathcal{M}_{V_\psi} \in \mathcal{A}_{(d_\phi, \infty)}$ and \mathcal{M}_{V_ψ} is a solution to the functional equation
- $$(5.1) \quad \mathcal{M}_{V_\psi}(z+1) = \phi(z) \mathcal{M}_{V_\psi}(z), \text{ with } \mathcal{M}_{V_\psi}(1) = 1,$$
- valid for $z \in \mathbb{C}_{(d_\phi, \infty)}$.

- (2) The mapping of moments of order greater than -1 , that is $u \mapsto \mathcal{M}_{V_\psi}(u)$, is the unique positive, log-convex (i.e. $\log \mathcal{M}_{V_\psi}$ is convex) solution to the functional equation (5.1) on \mathbb{R}_+ and

$$(5.2) \quad \mathcal{M}_{V_\psi}(u+1) \underset{\approx}{\sim} C_\psi \sqrt{\phi(u)} e^{G(u)},$$

where $C_\psi > 0$ and $G(u) = \int_1^u \ln \phi(r) dr$.

- (3) Finally, for any real number $u \leq \max(N_\tau - 1, 0)$ and any $a > d_\phi$, we have that

$$(5.3) \quad |\mathcal{M}_{V_\psi}(a+ib)| \underset{\pm\infty}{\sim} o(|b|^{-u}).$$

The proof of items (1) and (2) of Theorem 5.0.1 is given in Section 5.1.1 whereas the proof of Theorem 5.0.1 (3) is postponed to Section 5.3. We proceed with the following results regarding the smoothness properties of ν .

THEOREM 5.0.2. (1) *Supp $V_\psi = [0, \tau]$ and $\nu > 0$ on $(0, \tau)$.*

- (2) (a) *If $\psi \in \mathcal{N}_\infty$ then $\nu \in \mathcal{C}_0^\infty(\mathbb{R}_+)$.*

- (b) *If $\psi \in \mathcal{N}_\infty^c$ with $N_\tau = \left\lceil \frac{\overline{\Pi}(0^+)}{\tau} \right\rceil - 1 > 1$ then $\nu \in \mathcal{C}_0^{N_\tau-1}(\mathbb{R}_+)$, and, in any case, $\nu^{(N_\tau)} \in \mathcal{C}(\mathbb{R}_+ \setminus \{\tau\})$ and the mapping $x \mapsto (\tau-x)\nu^{(N_\tau)}(x) \in \mathcal{C}(\mathbb{R}_+)$ with $\lim_{x \rightarrow \tau} (\tau-x)\nu^{(N_\tau)}(x) = 0$. Consequently, for any $\psi \in \mathcal{N}$, $\nu \in \mathcal{C}^{N_\tau}(0, \tau)$.*

- (c) *Moreover, if $0 \leq N_\tau < \infty$, then for any $n = 0, 1, \dots, N_\tau$*

$$(5.4) \quad \nu^{(n)}(x) \underset{\approx}{\sim} C(\tau-x)^{\frac{\overline{\Pi}(0^+)}{\tau}-n-1} l(\tau-x) \mathbb{I}_{\{x < \tau\}},$$

where $C > 0$ and l is a slowly varying function at 0.

- (3) *If $\psi \in \mathcal{N}_\Theta$, then $\nu \in \mathcal{A}(\Theta_\phi)$, i.e. it is holomorphic in the sector $\mathbb{C}(\Theta_\phi) = \{z \in \mathbb{C}; |\arg z| < \Theta_\phi\}$. In particular, if $\psi \in \mathcal{N}_P$ then $\nu \in \mathcal{A}_{(0, \infty)} = \mathcal{A}\left(\frac{\pi}{2}\right)$.*

REMARK 5.0.3. The proof of item (3), which follows directly from a classical argument on Mellin transform described in (1.37), requires the estimate (6.47) along imaginary lines of the Mellin transform of ν which is given in Proposition 6.6.1. Note also that the last claim of item (3) is deduced from the previous one combined with Theorem 6.5.1(2(b)i). Although the proofs of these estimates are given in Remark 6.6.2 for sake of completeness and clarity, we state the analyticity property of ν here.

Items (1), (2b) and (2c) are deduced in Section 5.2.3 whereas item (2a) is settled in Section 5.3.2. The next result describes small time bounds and in some cases small asymptotic behaviour of ν .

THEOREM 5.0.4. *Let $\psi \in \mathcal{N}$. For any $\underline{a} < d_\phi$, $\underline{A} \in (0, \tau)$, there exists $C_{\underline{a}, \underline{A}} > 0$ such that*

$$(5.5) \quad \nu(x) \geq C_{\underline{a}, \underline{A}} x^{-\underline{a}}, \quad x \in (0, \underline{A}).$$

Moreover, if $m = \phi(0) = 0$ and $\phi'(0^+) < \infty$, then there exists $C > 0$ such that

$$(5.6) \quad \nu(x) \stackrel{0}{\sim} C = \nu(0^+).$$

The proof of Theorem 5.0.4 can be found in Section 5.2.5.

THEOREM 5.0.5. *Let $\psi \in \mathcal{N}_{\infty, \infty} = \{\psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \bar{\Pi}(0^+) = \infty\}$. Writing $\varphi : [m = \phi(0), \infty) \mapsto [0, \infty)$ for the continuous inverse of the continuous increasing function ϕ , i.e. $\varphi(\phi(u)) = u$, then there exists $C_\psi > 0$, such that for any $n \in \mathbb{N}$,*

$$(5.7) \quad \nu^{(n)}(x) \stackrel{\infty}{\sim} (-x)^{-n} \frac{C_\psi}{\sqrt{2\pi}} \sqrt{\varphi'(x)} \varphi^n(x) e^{-\int_m^x \varphi(y) \frac{dy}{y}}.$$

In particular,

(1) if $\psi \in \mathcal{N}_P$ then with $\bar{\sigma} := \sigma^{-2} > 0$ there exists $C_{\psi, \bar{\sigma}} > 0$ such that

$$(5.8) \quad \nu^{(n)}(x) \stackrel{\infty}{\sim} (-1)^n \frac{C_{\psi, \bar{\sigma}}}{\sqrt{2\pi}} \bar{\sigma}^{n+\frac{1}{2}} x^{m\bar{\sigma}} e^{-\bar{\sigma}x} e^{\bar{\sigma} \int_m^x \varphi_o(y) \frac{dy}{y}},$$

where $0 \leq \varphi_o(y) \stackrel{\infty}{\sim} o(y)$. If $\bar{\Pi} \in RV_{1+\alpha}(0)$, $\alpha \in (0, 1)$, see (1.35) for a definition, then

$$\varphi_o(y) \stackrel{\infty}{\sim} \bar{\sigma}^\alpha \alpha \Gamma(1 - \alpha) y^\alpha l(y)$$

and if $\bar{\Pi}(0^+) < \infty$ then $\varphi_o(y) \stackrel{\infty}{\sim} \bar{\Pi}(0^+)$;

(2) if $\psi \in \mathcal{N}_\alpha$, i.e. $\psi(u) \stackrel{\infty}{\sim} C_\alpha u^{\alpha+1}$, $C_\alpha > 0$, $\alpha \in (0, 1)$, then, there exists $C_{\psi, \alpha} > 0$ such that

$$(5.9) \quad \nu^{(n)}(x) \stackrel{\infty}{\sim} (-1)^n \frac{C_{\psi, \alpha}}{\sqrt{2\pi}} C_\alpha^{-\frac{1}{\alpha}(n+\frac{1}{2})} x^{n(\frac{1}{\alpha}-1)+\frac{1}{2\alpha}-\frac{1}{2}} e^{-\alpha C_\alpha^{-\frac{1}{\alpha}} x^{\frac{1}{\alpha}}(1+h(x))},$$

where $C_\alpha^{-\frac{1}{\alpha}} x^{\frac{1}{\alpha}} h(x) = \int_m^x \left(\varphi(y) - C_\alpha^{-\frac{1}{\alpha}} y^{\frac{1}{\alpha}} \right) \frac{dy}{y}$ and $h(x) = o(1)$.

The proof of these claims is given in Section 5.6 and it hinges on a generalization of a non-classical Tauberian theorem which was originally derived by Balkema et al. [**BKS93**] and we establish its new version in Proposition 5.4.3.

REMARK 5.0.6. Let $\hat{\nu}$ be the density of the positive self-decomposable law discussed in Section 5.1 below. Our result (5.7) is in fact a consequence of the small time asymptotic for $\hat{\nu}$ as presented in (5.48) of Theorem 5.4.1 via $\hat{\nu}_1(x) = x^{-2}\nu(x^{-1})$, see (5.12).

REMARK 5.0.7. Note that when $\phi(u) = u$ then $\nu(x) = e^{-x}$, $x > 0$, which is consistent with (5.7) with $C_\psi = \sqrt{2\pi}$. In general, it is not clear how to compute C_ψ precisely. This fact is due to the unknown constant appearing in [**Web97**, Theorem 6.3.].

As mentioned above the proofs of these results rely on a connection between the distribution of V_ψ and the one of a positive self-decomposable variable that we now describe.

5.1. A connection with remarkable self-decomposable variables

We recall that a (real-valued) variable X is self-decomposable, or of class \mathfrak{L} , if for any $a \in (0, 1)$, there exists an independent random variable X_a such that the following random affine equation

$$X \stackrel{(d)}{=} aX + X_a$$

holds. This class of variables plays a substantial role in probability theory as they arise in limit theorems for (properly normalized) sums of independent (not necessarily identically distributed) random variables. There is an important literature devoted to the study of their fine distributional properties and we refer to Sato's book [Sat99] and the paper of Sato and Yamazato [SY78], and the references therein, for a thorough account. In particular, in [SY78], a deep analysis of their probability distribution functions, such as smoothness properties, asymptotic behaviour at the lower end of their support, ultimate log-concavity property of the density, is carried out. This part aims to complement significantly this analysis for specific subclasses of \mathfrak{L} to the benefit of our spectral-theoretical study.

In [Sat99, Corollary 15.11], another interesting characterization of the class \mathcal{L} is presented as a subclass of the infinitely divisible random variables (recall that a variable Y is infinitely divisible if for every $n \in \mathbb{N} \setminus \{0\}$, there exists a sequence $(Y_{(k)})_{1 \leq k \leq n}$ of independent and identically distributed variables such that $Y \stackrel{(d)}{=} \sum_{k=1}^n Y_{(k)}$). For our purpose, we simply focus on the subset \mathfrak{L}_+ of positive self-decomposable variables whose Laplace transform takes the form, for any $u \geq 0$,

$$(5.10) \quad -\log \mathbb{E}[e^{-uX}] = \widehat{\phi}(u) = \widehat{\delta}u + \int_0^\infty (1 - e^{-uy}) \frac{\widehat{\kappa}(y)}{y} dy,$$

where $\widehat{\delta} \geq 0$ and $\widehat{\kappa}$ is a non-negative and non-increasing function such that it is true that $\int_0^\infty \widehat{\kappa}(y) dy < \infty$. Since $y \mapsto \frac{\widehat{\kappa}(y)}{y}$ is non-increasing therefore $\widehat{\phi} \in \mathcal{B}_{\mathcal{N}}$. Before stating our results, we introduce some further notation. We denote by

$$\mathcal{N}(m) = \{\psi \in \mathcal{N}; \psi'(0^+) = \phi(0) = m > 0\}.$$

Then, we set

$$(5.11) \quad I_\psi = \int_0^\infty e^{-\xi_t} dt$$

where $(\xi_t)_{t \geq 0}$ is a spectrally negative Lévy process with Laplace exponent $\psi \in \mathcal{N}(m)$ and $I_\psi < \infty$ a.s. since from the strong law of large numbers $\lim_{t \rightarrow \infty} \frac{\xi_t}{t} = m > 0$ a.s., see e.g. [BY05, Proposition 1]. This positive variable is called the exponential functional of the Lévy process ξ and has been the object of intense research over the last two decades. A review including motivation for its study is given in Section 6.1. Our interest in considering the variable I_ψ stems from the following result which explains its intimate connection to V_ψ .

PROPOSITION 5.1.1. *(1) For any $\psi \in \mathcal{N}(m)$, $I_\psi \in \mathfrak{L}_+$. Thus, the law of I_ψ is absolutely continuous with density denoted by $\widehat{\nu}$.*

(2) For any $\psi \in \mathcal{N}$, we have, for any $x > 0$,

$$(5.12) \quad \nu(x) = \frac{1}{x^2} \widehat{\nu}_1 \left(\frac{1}{x} \right),$$

with $\widehat{\nu}_1$ the density of $I_{\mathcal{T}_1\psi}$, where we recall, from Proposition 4.1.1(9), that $\mathcal{T}_1\psi(u) = u\phi(u+1) \in \mathcal{N}(\phi(1))$. Moreover,

$$\lim_{u \rightarrow \infty} \frac{\mathcal{T}_1\psi(u)}{u} = \lim_{u \rightarrow \infty} \phi(u+1) = \phi(\infty)$$

and with the notation Π_1 for the Lévy measure of $\mathcal{T}_1\psi$ then $\overline{\Pi}_1(0^+) = \overline{\Pi}(0^+)$.

(3) For any $\psi \in \mathcal{N}$, both variables $I_{\mathcal{T}_1\psi}$ and $X_{\mathcal{T}_1\psi} = \ln I_{\mathcal{T}_1\psi}$ are infinitely divisible.

REMARK 5.1.2. Note that item (3) reveals a remarkable property that is enjoyed by the class of positive self-decomposable variables considered in this Chapter.

Proof. Although the first claim is a well known fact, for sake of completeness, we provide its short proof. Writing, for any $a > 0$, $T_a = \inf\{t > 0; \xi_t \geq a\}$ and observing by absence of positive jumps for ξ that $\xi_{T_a} = a$ a.s. ($T_a < \infty$ a.s. as $\lim_{t \rightarrow \infty} \xi_t = \infty$), we get after performing a change of variables

$$\begin{aligned} I_\psi &= \int_0^{T_a} e^{-\xi_t} dt + \int_{T_a}^\infty e^{-\xi_t} dt \\ &= \int_0^{T_a} e^{-\xi_t} dt + e^{-a} \int_0^\infty e^{-(\xi_{T_a+t} - \xi_{T_a})} dt \\ &\stackrel{(d)}{=} \int_0^{T_a} e^{-\xi_t} dt + e^{-a} I_\psi, \end{aligned}$$

where I_ψ on the right-hand side is independent of $\int_0^{T_a} e^{-\xi_t} dt$ as from the strong Markov property for Lévy processes, see e.g. [Ber96, Proposition I.6], the process $(\xi_{T_a+t} - \xi_{T_a})_{t \geq 0}$ is a Lévy process distributed as ξ and independent of $(\xi_t)_{0 \leq t \leq T_a}$, see Section 5.2.1 for details. Hence $I_\psi \in \mathfrak{L}_+$. Thus, its law is absolutely continuous on \mathbb{R}_+ , see e.g. [Sat99, Theorem 27.13]. For the proof of item (2) we invoke [BY02b, Proposition 2] to get that

$$(5.13) \quad \mathbb{E} \left[I_\psi^{-n} \right] = \psi'(0^+) \frac{\prod_{k=1}^{n-1} \psi(k)}{(n-1)!} = \phi(0) W_\phi(n), \text{ for any } n = 1, 2, \dots,$$

where for the last identity we have used that

$$\psi'(0^+) = \lim_{u \downarrow 0} \frac{\psi(u)}{u} = \lim_{u \downarrow 0} \phi(u) = \phi(0)$$

and the definition of W_ϕ , see (1.18). Next, note that $\mathcal{T}_1\psi(u) = u\phi(u+1)$ and thus the claim $\lim_{u \rightarrow \infty} \frac{\mathcal{T}_1\psi(u)}{u} = \phi(\infty)$, in (2), is obvious. From (4.9), we have that $\overline{\Pi}_1(y) = \int_y^\infty e^{-r} (\overline{\Pi}(r) dr + \Pi(dr)) = e^{-y} \overline{\Pi}(y)$, whereby we deduce the very last claim. Finally, observe, from (2.20) for the first identity and from (5.13) for the

last one, where we use the notation $\phi_1(u) = \phi(u+1)$, that, for any $n = 1, 2, \dots$, we have that

$$(5.14) \quad \mathbb{E} \left[V_\psi^n \right] = W_\phi(n+1) = \phi(1) \prod_{k=1}^{n-1} \phi(k+1) = \phi_1(0) W_{\phi_1}(n) = \mathbb{E} \left[I_{\mathcal{T}_1\psi}^{-n} \right].$$

Since V_ψ is moment determinate, the proof of (5.12) and thus of (2) is completed. The last claim follows easily from the fact $I_{\mathcal{T}_1\psi}$ is self-decomposable and hence infinitely divisible and from [Urb95] where it is shown that the variable V_ψ is a multiplicative infinitely divisible variable, that is, in particular, $X_{\mathcal{T}_1\psi}$ is infinitely divisible. \square

5.1.1. Proof of Theorem 5.0.1(1) and (2). In the proof of [BY02b, Proposition 2], the authors show that for any $n \in \mathbb{N}$

$$(5.15) \quad \mathbb{E} \left[I_{\mathcal{T}_1\psi}^{-n-1} \right] = \frac{\mathcal{T}_1\psi(n)}{n} \mathbb{E} \left[I_{\mathcal{T}_1\psi}^{-n} \right] = \frac{n\phi(n+1)}{n} \mathbb{E} \left[I_{\mathcal{T}_1\psi}^{-n} \right] = \phi(n+1) \mathbb{E} \left[I_{\mathcal{T}_1\psi}^{-n} \right].$$

For completeness we replicate their proof valid also for $z \in \mathbb{C}_{(d_\phi-1, \infty)}$, i.e. $\Re(z) > d_\phi - 1$ or when $\phi(z+1)$ is well-defined. Set $I_t = \int_t^\infty e^{-\xi s} ds$, $t \geq 0$, where ξ is a Lévy process with exponent $\mathcal{T}_1\psi$, see Proposition 4.1.1(9) and (5.21). For any $z \in \mathbb{C}$,

$$I_t^{-z} - I_0^{-z} = z \int_0^t e^{-\xi s} I_s^{-z-1} ds.$$

Moreover, from Section 5.2.1, $I_t = e^{-\xi t} \int_0^\infty e^{-(\xi_t+s-\xi_t)} ds \stackrel{(d)}{=} e^{-\xi t} I_0$ with $I_0 \stackrel{(d)}{=} I_{\mathcal{T}_1\psi}$ independent of $e^{-\xi t}$. From (5.21) $\mathbb{E} [e^{z\xi t}] = e^{t\mathcal{T}_1\psi(z)}$ and it is finite whenever $\mathcal{T}_1\psi(z) = z\phi(z+1)$ exists, i.e. $z \in \mathbb{C}_{(d_\phi-1, \infty)}$. Then upon taking expectations we get

$$\begin{aligned} \mathbb{E} [I_t^{-z} - I_0^{-z}] &= (\mathbb{E} [e^{z\xi t}] - 1) \mathbb{E} [I_0^{-z}] = z \mathbb{E} [I_0^{-z-1}] \int_0^t \mathbb{E} [e^{z\xi s}] ds \\ &= \frac{z}{\mathcal{T}_1\psi(z)} \left(e^{t\mathcal{T}_1\psi(z)} - 1 \right) \mathbb{E} [I_0^{-z-1}] \\ &= \frac{z}{\mathcal{T}_1\psi(z)} (\mathbb{E} [e^{z\xi t}] - 1) \mathbb{E} [I_0^{-z-1}], \end{aligned}$$

that is the complex version of (5.15). Since $\mathbb{E} [I_{\mathcal{T}_1\psi}^0] = 1$ and (5.15) links the moments recurrently, $\mathbb{E} [I_{\mathcal{T}_1\psi}^{-n}] < \infty$, for all $n \in \mathbb{N}$. By moment determinacy the identity $\mathbb{E} [V_\psi^n] = \mathbb{E} [I_{\mathcal{T}_1\psi}^{-n}]$ in (5.14) extends to $z \in \mathbb{C}_{(d_\phi-1, \infty)}$ which together with (5.15) give (5.1), i.e. item (1) with $\mathcal{M}_{V_\psi}(1) = \mathbb{E} [V_\psi^0] = 1$. Next, recall that the mapping $u \mapsto \mathcal{M}_{V_\psi}(u) = \mathbb{E} [V_\psi^{u-1}]$, as the moments of order greater than -1 of a positive random variable, is a positive and log-convex function on \mathbb{R}_+ . However, from items (2) and (6) of Proposition 4.1.1 for any $\phi \in \mathcal{B}$, the mapping $u \mapsto \phi(u)$ is positive and strictly log-concave on $(0, \infty)$ and the limit in (4.6) holds. This means that the multiplier ϕ in (5.1) satisfies the conditions of [Web97, Theorem 7.1], which states that the functional equation (5.1) has a unique, positive, log-convex solution on \mathbb{R}_+ . However, the same conditions on ϕ trigger the validity of [Web97, Theorem 6.3] and we complete (5.2) and Theorem 5.0.1(2).

5.2. Fine distributional properties of I_ψ

In this part we state some properties on the density of the distribution of the variable I_ψ - one regarding its smoothness and the other its large asymptotic behaviour. Since the proof of both results hinges on techniques based on the fluctuation and excursion theory of spectrally negative Lévy processes, we proceed by recalling some essential facts on this topic which will make the proofs more legible. For any $\psi \in \mathcal{N}(m)$, write $X_\psi = -\ln I_\psi$ and denote by χ its density, that is

$$(5.16) \quad \chi(x) = e^{-x} \widehat{\nu}(e^{-x}), \quad x \in \mathbb{R}.$$

Next, recall that $\mathcal{N}_\infty = \{\psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \overline{\Pi}(0^+) = \infty\}$ and

$$\mathcal{N}_{\infty, \infty} = \left\{ \psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \overline{\overline{\Pi}}(0^+) = \infty \right\}.$$

We use in the sequel

$$\mathcal{N}_\infty(m) = \mathcal{N}_\infty \cap \mathcal{N}(m), \quad \mathcal{N}_\infty^c(m) = \mathcal{N}_\infty^c \cap \mathcal{N}(m) \text{ and } \mathcal{N}_{\infty, \infty}^c(m) = \mathcal{N}(m) \setminus \mathcal{N}_{\infty, \infty}(m).$$

PROPOSITION 5.2.1. *Let $\psi \in \mathcal{N}(m)$ and recall that $\psi(u) = u\phi(u)$, with $\phi \in \mathcal{B}_\mathcal{N}$. Then, since $I_\psi \in \mathfrak{L}_+$, with the notation of (5.10), we have, for any $u \geq 0$,*

$$(5.17) \quad -\log \mathbb{E}[e^{-uI_\psi}] = \widehat{\phi}(u) = \widehat{\delta}u + \int_0^\infty (1 - e^{-uy}) \frac{\widehat{\kappa}(y)}{y} dy,$$

where

(1) $\widehat{\delta} = \frac{1}{\mathfrak{r}} \geq 0$, where we recall that $0 < \mathfrak{r} = \phi(\infty) = \overline{\overline{\Pi}}(0^+) + m \leq \infty$. Thus, $\widehat{\delta} > 0$ if and only if $\psi \in \mathcal{N}_{\infty, \infty}^c(m)$. Moreover, $\widehat{\kappa}(0^+) = \infty$ (resp. $0 < \widehat{\kappa}(0^+) = \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} < \infty$) if and only if $\psi \in \mathcal{N}_\infty(m)$ (resp. $\psi \in \mathcal{N}_\infty^c(m)$).

(2) $\text{Supp } I_\psi = [\frac{1}{\mathfrak{r}}, \infty)$.

(3) Next, $\widehat{\nu} \in \mathcal{C}^\infty(\mathbb{R})$ if and only if $\psi \in \mathcal{N}_\infty(m)$. Otherwise, if $1 < N_\mathfrak{r} < \infty$, $\widehat{\nu} \in \mathcal{C}^{N_\mathfrak{r}-1}(\mathbb{R})$, where we recall that $N_\mathfrak{r} = \left\lceil \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} \right\rceil - 1$ (with $N_\mathfrak{r} = \infty$ whenever $\overline{\Pi}(0^+) = \infty$), and, in any case, $\widehat{\nu}^{(N_\mathfrak{r})} \in \mathcal{C}(\mathbb{R} \setminus \{\frac{1}{\mathfrak{r}}\})$ and the mapping $x \mapsto (x - \frac{1}{\mathfrak{r}}) \widehat{\nu}^{(N_\mathfrak{r})}(x) \in \mathcal{C}(\mathbb{R})$ with $\lim_{x \rightarrow \frac{1}{\mathfrak{r}}} (x - \frac{1}{\mathfrak{r}}) \widehat{\nu}^{(N_\mathfrak{r})}(x) = 0$.

(4) Moreover, if $N_\mathfrak{r} < \infty$, then for any $n = 0, 1, \dots, N_\mathfrak{r}$

$$(5.18) \quad \widehat{\nu}^{(n)}(x) \underset{\frac{1}{\mathfrak{r}}}{\sim} C x_\mathfrak{r}^{\frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - n - 1} l(x_\mathfrak{r}),$$

where $C > 0$, $x_\mathfrak{r} = x - \frac{1}{\mathfrak{r}}$ and l is a slowly varying function at 0.

(5) The statements concerning the support and smoothness properties on $\widehat{\nu}$ hold in a similar way for χ as defined in (5.16).

We proceed by providing very useful results concerning the asymptotic behaviour of $\widehat{\nu}$ at infinity. Recall, from (1.19), that $d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\} \leq 0$.

PROPOSITION 5.2.2. *Let $\psi \in \mathcal{N}(m)$.*

(1) *For any $a < d_\phi$, $A > \frac{1}{\mathfrak{r}}$, there exists a constant $C_{a,A} > 0$ such that*

$$(5.19) \quad \widehat{\nu}(x) \geq C_{a,A} x^{a-1}, \quad x \in (A, \infty).$$

(2) Assume that there is $u_0 > 0$ such that $\int_1^\infty ye^{u_0 y} \Pi(dy) < \infty$ and $\psi(-u_0) = \phi(-u_0) = 0$, then there exists a constant $C_{u_0} > 0$ such that

$$(5.20) \quad \widehat{\nu}(x) \approx C_{u_0} x^{-u_0-1}.$$

REMARK 5.2.3. If C_{u_0} in (5.20) can be made explicit then it is an evaluation of the so-called Kesten's constant, see [Kes73], for an affine random equation solved by I_ψ . We have achieved this [PS18a]. However, for our purposes, the weaker estimate (5.19) suffices.

We prove the Proposition 5.2.1 and Proposition 5.2.2 in section 5.2.2 and 5.2.4 respectively. In what follows, we review several aspects of Lévy processes which play a central role in their proofs and are also useful throughout the rest of the paper. We refer to the monograph [Ber96] for a nice account on Lévy processes.

5.2.1. Spectrally negative Lévy processes, fluctuation and excursions theory. A spectrally negative Lévy process is a real-valued stochastic process, $\xi = (\xi_t)_{t \geq 0}$, defined on the probability space $(\Omega, \mathcal{F}, \mathbb{P})$, which can jump *downwards* only and possesses stationary and independent increments, i.e. $\xi_t - \xi_s \stackrel{(d)}{=} \xi_{t-s}$, for $0 \leq s < t$, and $\xi_t - \xi_s$ is independent of $(\xi_u)_{u \leq s}$. Every general Lévy process and in particular every spectrally negative one, has the Lévy-Itô decomposition $\xi_t = mt + \sigma B_t + Z_t$, where $m \in \mathbb{R}$ in accordance with (1.3), $B = (B_t)_{t \geq 0}$ is a Brownian motion independent of the pure jump process $Z = (Z_t)_{t \geq 0}$. There is a natural bijection between the subclass of negative definite functions \mathcal{N} , see (1.4), and a large subclass of spectrally negative Lévy processes via (1.3), where σ^2 is the variance of B and Π describes the intensity and the size of the jumps of Z . For the class \mathcal{N} , see (1.4), which is the focus of our paper, we have the bijection between $\psi \in \mathcal{N}$ of the form

$$(5.21) \quad \ln \mathbb{E} [e^{z\xi_1}] = \psi(z) = mz + \frac{\sigma^2}{2} z^2 + \int_0^\infty (e^{-zy} - 1 + zy) \Pi(dy), \quad z \in i\mathbb{R},$$

and ξ spectrally negative Lévy process with $m = \psi'(0^+) = \mathbb{E} [\xi_1] \in [0, \infty)$. In the setting of Lévy processes $m > 0$ gives $\lim_{t \rightarrow \infty} \xi_t = \infty$ \mathbb{P} -almost surely (a.s.) and $m = 0$ leads to $\overline{\lim}_{t \rightarrow \infty} \xi_t = -\underline{\lim}_{t \rightarrow \infty} \xi_t = \infty$ a.s.. It is clear from (5.21) that $\psi \in \mathcal{A}_{[0, \infty)}$.

Moreover, for $a > 0$, $\psi \in \mathcal{A}_{(-a, \infty)}$ if and only if $\forall u \in (-a, 0)$, $|\mathbb{E} [e^{u\xi_1}]| < \infty$ which is equivalent to

$$(5.22) \quad \left| \int_{y>1} e^{-uy} \Pi(dy) \right| < \infty,$$

see [Ber96, Chap. I]. The restriction of ψ on the real interval $(-a, \infty)$ is clearly a convex function and ψ is zero free on $(0, \infty)$. The analytical form of the Wiener-Hopf factorization for $\psi \in \mathcal{N}$ reads off as follows

$$(5.23) \quad \psi(z) = z\phi(z), \quad z \in \mathbb{C}_{[0, \infty)},$$

that is (1.9). It has the following probabilistic interpretation through the identities

$$(5.24) \quad \ln \mathbb{E} [e^{z\eta_1^+}] = z \quad \text{and} \quad \ln \mathbb{E} [e^{-z\eta_1^-}] = -\phi(z),$$

where $\eta^- = (\eta_t^-)_{t \geq 0}$ is a possibly killed subordinator, that is a non-decreasing Lévy process possibly killed at an independent exponential time, known as the descending ladder height process, and, the ascending ladder height process $\eta^+ = (\eta_t^+ = t)_{t \geq 0}$ is a pure drift process. We note that for $\psi \in \mathcal{N}$, η^+ is never killed. We record that

$$(5.25) \quad \phi(z) = m + \sigma^2 z + \int_0^\infty (1 - e^{-zy}) \bar{\Pi}(y) dy,$$

that is the classical Bernstein function already defined in (1.10). Hence $\sigma^2 \geq 0$, $\int_0^\infty (y \wedge 1) \bar{\Pi}(y) dy < \infty$ and $m \geq 0$ is the killing term of η^- since $\overline{\lim}_{t \rightarrow \infty} \xi_t = \infty$ which in turn is due to $\psi \in \mathcal{N}$. Finally, $\phi(\infty) < \infty$ if and only if $\bar{\Pi}(\mathbb{R}_+) < \infty$ if and only if η^- is a compound Poisson process, see [Ber96, Chap I]. Next, define $\bar{\xi} = (\bar{\xi}_t)_{t \geq 0} = (\sup_{s \leq t} \xi_s - \xi_t)_{t \geq 0}$, which is known as the reflected at the supremum Lévy process. It possesses a local time at 0, i.e. a non-decreasing (in $t \geq 0$) family of continuous functionals $\bar{\ell} = (\bar{\ell}_t)_{t \geq 0}$ such that $\bar{\ell}$ increases only on the closure of the set $\{s \geq 0; \bar{\xi}_s = 0\}$, that is on the closure of the set of times when a new running supremum is attained for ξ . The inverse local time at 0 of $\bar{\xi}$ is defined, for any $t \geq 0$, as $\bar{\eta}_t = \inf\{s \geq 0; \bar{\ell}_s > t\}$. We know that $(\xi \circ \bar{\eta}, \bar{\eta}) = (\eta^+, \bar{\eta})$ a.s. and thus the process $(\xi \circ \bar{\eta}, \bar{\eta})$ determines a bivariate subordinator. Furthermore, with the reflected process $\bar{\xi}$ one associates excursions away from the supremum (loosely speaking the piece of path of ξ between successive suprema) in the following manner. Write, for any $t > 0$, $\mathbf{e}_t = (\bar{\xi}_{\bar{\eta}_{t-} + s})_{0 \leq s < \zeta}$, where $\zeta = \zeta(t) = \inf\{s > 0; \bar{\xi}_{\bar{\eta}_{t-} + s} = 0\}$. $\mathbf{e} = (\mathbf{e}_t)_{t \geq 0}$ is the excursion process which forms a Poisson point process in the space of right-continuous functions with left limits (for short rcll). Its Poisson measure \mathbf{n} lives on the sets of rcll functions and is referred to as the excursion measure. ζ is called a lifetime of an excursion. We use the formula $\mathbf{n}(F(\mathbf{e})) = \int_{\mathfrak{W}} F(\omega) \mathbf{n}(d\omega)$ for computing various functionals on the space of excursions \mathfrak{W} (for more information on excursion theory of Lévy processes, we refer to [Ber96, Chap. IV and Chap. VI]). Finally, when $\psi \in \mathcal{N}$ then ξ is a spectrally negative Lévy process drifting with $\overline{\lim}_{t \rightarrow \infty} \xi_t = \infty$. Thus, recalling that $T_t = \inf\{s > 0; \xi_s > t\}$, we have, see [Ber96, Theorem 1, Chapter VII], for any $t, u > 0$,

$$\mathbb{E} [e^{-u\bar{\eta}_t}] = \mathbb{E} [e^{-uT_t}] = e^{-t\bar{\phi}(u)},$$

where $\psi(\bar{\phi})(u) = u$, with $\bar{\phi} \in \mathcal{B}$ and from [Ber96, Theorem 8, Chapter IV],

$$(5.26) \quad \bar{\phi}(u) = \bar{\delta}u + u \int_0^\infty e^{-uy} \mathbf{n}(\zeta > y) dy,$$

with $\bar{\delta} \geq 0$ and \mathbf{n} the excursion measure of $\bar{\xi}$.

5.2.2. Proof of Proposition 5.2.1. The proof of Proposition 5.2.1 requires the following two intermediate results which characterize respectively the quantities $\widehat{\delta}$ and $\widehat{\kappa}$ in (5.17).

LEMMA 5.2.4. $\widehat{\delta} = \frac{1}{\nu} \geq 0$.

Proof. From [Riv12, Lemma 2], $\widehat{\delta}$ in (5.17) is the drift of the subordinator $\bar{\eta} = (\bar{\eta}_t)_{t \geq 0}$, the inverse local time at 0 of $\bar{\xi}$, i.e. $\bar{\delta}$ in (5.26). Hence, from Proposition

4.1.1(3), we get that

$$(5.27) \quad \widehat{\delta} = \bar{\delta} = \lim_{u \rightarrow \infty} \frac{\bar{\phi}(u)}{u} = \lim_{u \rightarrow \infty} \frac{u}{\psi(u)} = \frac{1}{\phi(\infty)} = \frac{1}{\mathfrak{r}} \geq 0,$$

which completes the proof of the Lemma. \square

We proceed the proof with a statement which, in particular, characterizes $\widehat{\kappa}$.

LEMMA 5.2.5. (1) We have, for any $y > 0$,

$$(5.28) \quad \widehat{\kappa}(y) = \mathbf{n} \left(\int_0^\zeta e^{\varepsilon_s} ds > y \right),$$

where we recall that ζ is the lifetime of a given excursion $(\varepsilon_s)_{s \geq 0}$.

(2) Moreover, $\widehat{\kappa}(0^+) = \infty$ (resp. $\widehat{\kappa}(0^+) = \frac{\bar{\Pi}(0^+)}{\mathfrak{r}}$) if and only if $\psi \in \mathcal{N}_\infty(m)$ (resp. $\psi \in \mathcal{N}_\infty^c(m)$).

REMARK 5.2.6. The subset of the class \mathfrak{L}_+ we consider can be characterized as the one associated to a subordinator η^* , as defined in (5.29) below, whose tail of the Lévy measure is associated via (5.28) to the excursion measure \mathbf{n} of $\bar{\xi}$. Clearly, our class does not include the positive self-decomposable laws associated to subordinators with an atomic Lévy measure, since these can not be related to such an excursion measure.

Proof. We start with the following well-known integral representation of $I_\psi \in \mathfrak{L}_+$

$$(5.29) \quad I_\psi = \int_0^\infty e^{-t} d\eta_t^*,$$

where $\eta^* = (\eta_t^*)_{t \geq 0}$ is a subordinator whose tail of the Lévy measure is $\widehat{\kappa}$, see e.g. [Sat99, Theorem 17.5, Example 17.10]. The first identity (5.28) then follows from [Riv12, Lemma 2] (note that the exponential functional is defined in [Riv12] as $\int_0^\infty e^{\varepsilon_s} ds$ whereas we have $I_\psi = \int_0^\infty e^{-\varepsilon_s} ds$). Since $\zeta > 0$ and $\varepsilon_s > 0$, for all $0 \leq s < \zeta$, we get that

$$\mathbf{n} \left(\int_0^\zeta e^{\varepsilon_s} ds > 0 \right) = \mathbf{n}(\zeta > 0).$$

This identity combined with (5.28) yield that $\widehat{\kappa}(0^+) = \infty$ (resp. $\widehat{\kappa}(0^+) < \infty$) if and only if $\mathbf{n}(\zeta > 0) = \infty$ (resp. $\mathbf{n}(\zeta > 0) < \infty$) which is equivalent to $\psi \in \mathcal{N}_\infty(m)$ (resp. $\mathcal{N}_\infty^c(m)$), see e.g. [Don07, Proposition 15 (iv)]. It thus remains to compute $\widehat{\kappa}(0^+)$ when $\psi \in \mathcal{N}_\infty^c(m)$. Recall from (5.26) that $\mathbf{n}(\zeta > 0)$ is the total mass of the Lévy measure associated to $\bar{\phi}$. When $\psi \in \mathcal{N}_\infty^c(m)$ from (5.27), $\widehat{\delta} = \bar{\delta} = \frac{1}{\mathfrak{r}} = \frac{1}{\phi(\infty)} >$

0, a standard computation leads to

$$\begin{aligned} \mathbf{n}(\zeta > 0) &= \lim_{u \rightarrow \infty} (\bar{\phi}(u) - \bar{\delta}u) = \lim_{u \rightarrow \infty} \left(u - \widehat{\delta}\psi(u) \right) = \lim_{u \rightarrow \infty} \frac{1}{\phi(\infty)} u (\phi(\infty) - \phi(u)) \\ &= \frac{1}{\phi(\infty)} \lim_{u \rightarrow \infty} u \left(\int_0^\infty \bar{\Pi}(y) dy - \int_0^\infty (1 - e^{-uy}) \bar{\Pi}(y) dy \right) \\ &= \frac{\bar{\Pi}(0^+)}{\mathfrak{r}}, \end{aligned}$$

where we have used the identity (5.25) to express ϕ . \square

The two previous lemmas provide the proof of Proposition 5.2.1(1). Next, from [HR12, Lemma 2.1], $\text{Supp } I_\psi = [0, \infty)$ if and only if $\widehat{\delta} = 0$, which is equivalent from (5.27), to $\psi \in \mathcal{N}_{\infty, \infty}(m)$. Otherwise, according to [Riv12, Remark 1, p. 9] it holds that $\text{Supp } I_\psi = [\widehat{\delta}, \infty)$ since η^* in (5.29) is a subordinator with drift $\widehat{\delta}$ (note that the constant c in [Riv12] is 1 which is clear here from the identity $\psi(u) = u\phi(u)$). Thus, we have proved the statement for the support of $\widehat{\nu}$, i.e. item (2). Next, if $\psi \in \mathcal{N}_\infty(m)$ (resp. $\psi \in \mathcal{N}_\infty^c(m)$), then, since from Lemma 5.2.5(2), we have $\widehat{\kappa}(0^+) = \infty$ (resp. $\widehat{\kappa}(0^+) = \frac{\bar{\Pi}(0^+)}{\mathfrak{r}}$) and $\widehat{\delta} = 0$ (resp. $\widehat{\delta} > 0$), we derive from [Sat99, Theorem 28.4(ii)] (resp. [Sat99, Theorem 28.4(i)]) that $\widehat{\nu} \in \mathcal{C}^\infty(\mathbb{R})$ (resp. $\widehat{\nu} \in \mathcal{C}^{\mathbf{N}_\mathfrak{r}-1}(\mathbb{R})$, when $\frac{\bar{\Pi}(0^+)}{\mathfrak{r}} > 1$, i.e. $\mathbf{N}_\mathfrak{r} \geq 1$). The claim (4) is given in [SY78, Theorem 1.6]. The proof of Proposition 5.2.1(5) follows immediately from the definition of χ in (5.16), which completes the proof of Proposition 5.2.1.

5.2.3. Proofs of Theorem 5.0.2(1), (2b) and (2c). First, we recall from Proposition 5.1.1(2) that, for any $\psi \in \mathcal{N}$, $\nu(x) = x^{-2}\widehat{\nu}_1(x^{-1})$ where $\widehat{\nu}_1$ stands for the density of $I_{\mathcal{T}_1\psi}$ with $\mathcal{T}_1\psi(u) = u\phi(u+1) \in \mathcal{N}(\phi(1))$ and $\bar{\Pi}_1(0^+) = \bar{\Pi}(0^+)$. Thus, the positivity follows from the fact that $\widehat{\nu}_1 > 0$ on $(\frac{1}{\mathfrak{r}}, \infty)$, see [PS12a, SY78], and the expression of the support of V_ψ is deduced from Proposition 5.2.1(2). Next, the smoothness property of ν stated in Theorem 5.0.2(2b) and the ensuing Lemma 5.2.7 follow readily from Proposition 5.2.1(3). Finally, the asymptotic behaviour at \mathfrak{r} in the item (2c) follows also readily from Proposition 5.2.1(4) and $\nu(x) = x^{-2}\widehat{\nu}_1(x^{-1})$.

Proposition 5.2.1(3) and Proposition 5.1.1(2) give immediately the following claim.

LEMMA 5.2.7. *If $\psi \in \mathcal{N}_\infty$ then $\nu \in \mathcal{C}^\infty(\mathbb{R}_+)$.*

To prove Proposition 5.2.2 and for sake of completeness we restate, adapted to our setting, a shortened version of [SY78, Theorem 2.1].

THEOREM 5.2.8. *Let $\psi \in \mathcal{N}(m)$. Then*

$$(5.30) \quad x\widehat{\nu}(x) = \int_0^x \widehat{\nu}(x-y)\widehat{\kappa}(y)dy + \widehat{\delta}\widehat{\nu}(x), \quad x > \frac{1}{\mathfrak{r}}.$$

For any $x > \frac{1}{\mathfrak{r}}$ and with $x \rightarrow \infty$ in the last relation

$$(5.31) \quad x\widehat{\nu}(x) \geq \widehat{\kappa}(x) \int_{\frac{1}{\mathfrak{r}}}^x \widehat{\nu}(y)dy = \mathbf{n} \left(\int_0^\zeta e^{\varepsilon s} ds > x \right) \int_{\frac{1}{\mathfrak{r}}}^x \widehat{\nu}(y)dy \approx \mathbf{n} \left(\int_0^\zeta e^{\varepsilon s} ds > x \right).$$

PROOF OF THEOREM 5.2.8. Since thanks to the link in (5.28), I_ψ is a positive self-decomposable law with $\int_0^\infty \widehat{\kappa}(y)dy < \infty$ then in the notation of [SY78, (1.5)] we have

$$\log \mathbb{E} [e^{iuI_\psi}] = iu\gamma_0 - \frac{\sigma^2 u^2}{2} + \int_0^\infty (e^{iuy} - 1) \frac{k(y)}{y} dy.$$

Comparing this to (5.10) yields in our notation that $\gamma_0 = \widehat{\delta} = \mathfrak{r}^{-1}$, $\sigma^2 = 0$ and $\widehat{\kappa} \equiv k$. Then, the second identity of [SY78, Theorem 2.1, (2.1)] in their notation takes the form

$$(x - \gamma) f(x) = \int_0^x f(x-y)k(y)dy - f(x) \int_0^\infty \frac{k(y)}{1+y^2} dy$$

Since [SY78, (1.6)] gives $\gamma_0 = \gamma - \int_0^\infty \frac{k(y)}{1+y^2} dy = \widehat{\delta}$ and $f = \widehat{\nu}$ we deduct (5.30) and (5.31) since $\widehat{\kappa}$ is non-increasing. \square

5.2.4. Proof of Proposition 5.2.2. For item 1 from the last relation of (5.31) we need to discuss $\mathfrak{n} \left(\int_0^\zeta e^{\epsilon s} ds > x \right)$ only. We recall that to each $\psi \in \mathcal{N}(m)$ there is an associated spectrally negative Lévy process ξ such that $\lim_{t \rightarrow \infty} \xi_t = \infty$, see the discussion succeeding (5.21). We set, for any $a > 0$,

$$I_\psi(\bar{T}_a) = \int_0^{\bar{T}_a} e^{-\xi_s} ds,$$

where $\bar{T}_a = \inf\{t \geq 0; \xi_t = a\}$ is the first hitting time of a for ξ . Also, it is well-known from [Ber96, Chapter IV, p.117] that, for any $a, x > 0$,

$$\mathfrak{n} \left(\int_1^\zeta e^{\epsilon s} ds > x \mid \mathfrak{e}_1 = a, \zeta > 1 \right) = \mathbb{P} (e^a I_\psi(\bar{T}_a) > x),$$

which reflects the fact that under $\mathfrak{n} \left(\cdot \mid \mathfrak{e}_1 = a, \zeta > 1 \right)$, $(\mathfrak{e}_s)_{1 \leq s < \zeta} = (a + \mathfrak{e}_s - \mathfrak{e}_1)_{1 \leq s < \zeta}$ has the same law as $(a - \xi_s)_{0 \leq s < \bar{T}_a}$, since $\mathfrak{e}_1 = a$ means that being a units beneath the running supremum the process ξ needs to stride upwards those a units to attain a new maximum. Using this and trivial estimates yield, for $x > 0$, that

$$\begin{aligned} \mathfrak{n} \left(\int_0^\zeta e^{\epsilon s} ds > x \right) &\geq \mathfrak{n} \left(\int_1^\zeta e^{\epsilon s} ds > x, \zeta > 1, \mathfrak{e}_1 > e \right) \\ &= \int_{a>e} \mathbb{P} (e^a I_\psi(\bar{T}_a) > x) \mathfrak{n}(\mathfrak{e}_1 \in da, \zeta > 1) \\ (5.32) \quad &\geq \mathbb{P} (I_\psi(\bar{T}_1) > x) \mathfrak{n}(\mathfrak{e}_1 > e, \zeta > 1). \end{aligned}$$

Next, let us denote by $\xi^\epsilon = (\xi_t^\epsilon)_{t \geq 0} = \left(\xi_t - \sum_{s \leq t} (\xi_s - \xi_{s-}) \mathbb{I}_{\{(\xi_s - \xi_{s-}) < -\epsilon^{-1}\}} \right)_{t \geq 0}$ the Lévy process constructed from ξ by removing all its jumps smaller than $-\epsilon^{-1} < 0$. Thus, pathwise, $\xi_t \leq \xi_t^\epsilon$, for any $t > 0$, $\bar{T}_1 \geq \bar{T}_1^\epsilon$ and therefore, $I_\psi(\bar{T}_1) \geq I_\psi(\bar{T}_1^\epsilon)$. Choose furthermore ϵ small enough such that $\Pi(0, \epsilon^{-1}) > 0$, that is ξ^ϵ is not a pure drift process or phrased otherwise it still possesses negative jumps. For any $a > e, x > e$, from the strong Markov property of Lévy processes, conditionally on the event $\{\bar{T}_{-\ln xa}^\epsilon < \infty\}$, the process $\check{\xi}^\epsilon = \left(\check{\xi}_t^\epsilon = \xi_{t+\bar{T}_{-\ln xa}^\epsilon}^\epsilon + \ln xa \right)_{t \geq 0}$ is a Lévy

process issued forth from 0 and independent of $(\xi_t^\epsilon)_{0 \leq t \leq \bar{T}_{-\ln(xa)}^\epsilon}$ and

$$\begin{aligned}
I_\psi(\bar{T}_1) &\geq I_{\psi^\epsilon}(\bar{T}_1^\epsilon) = \int_0^{\bar{T}_1^\epsilon} e^{-\xi_s^\epsilon} ds \geq \int_{\bar{T}_{-\ln(xa)}^\epsilon}^{\bar{T}_1^\epsilon} e^{-\xi_s^\epsilon} ds \mathbb{I}_{\{\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\}} \\
&= xa \int_0^{\bar{T}_1^\epsilon + \ln xa} e^{-\tilde{\xi}_s^\epsilon} ds \mathbb{I}_{\{\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\}} \\
(5.33) \quad &\geq xa \int_0^{\bar{T}_{\ln a}^\epsilon} e^{-\tilde{\xi}_s^\epsilon} ds \mathbb{I}_{\{\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\}},
\end{aligned}$$

where \bar{T}_1^ϵ stands for the hitting times of $\tilde{\xi}^\epsilon$. Hence, using the independence of $\tilde{\xi}^\epsilon$ and ξ^ϵ , $\xi^\epsilon = \tilde{\xi}^\epsilon$ in law and the fact that $\tilde{\xi}_t^\epsilon \leq \ln a$ on $\{t < \bar{T}_{\ln a}^\epsilon\}$ we obtain from (5.33) that

$$\begin{aligned}
\mathbb{P}(I_\psi(\bar{T}_1) > x) &\geq \mathbb{P}(I_{\psi^\epsilon}(\bar{T}_1^\epsilon) > x) \geq \mathbb{P}\left(a \int_0^{\bar{T}_{\ln a}^\epsilon} e^{-\tilde{\xi}_s^\epsilon} ds > 1\right) \mathbb{P}\left(\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\right) \\
&\geq \mathbb{P}\left(\bar{T}_{\ln(a)}^\epsilon > 1\right) \mathbb{P}\left(\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\right) \\
(5.34) \quad &= \mathbb{P}\left(\bar{T}_{\ln(a)}^\epsilon > 1\right) \mathbb{P}\left(\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\right).
\end{aligned}$$

From $\lim_{t \rightarrow \infty} \xi_t^\epsilon \geq \lim_{t \rightarrow \infty} \xi_t = \infty$, ξ^ϵ only jumps down and is not a pure drift process we deduce that $\mathbb{P}\left(\bar{T}_{\ln(a)}^\epsilon > 1\right) > 0$. Next,

$$\begin{aligned}
\mathbb{P}\left(\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\right) &= 1 - \mathbb{P}\left(\bar{T}_1^\epsilon \leq \bar{T}_{-\ln(xa)}^\epsilon\right) \\
(5.35) \quad &= 1 - \frac{S_\epsilon(\ln(xa))}{S_\epsilon(\ln(xa) + 1)} = \frac{S_\epsilon(\ln(xae)) - S_\epsilon(\ln(xa))}{S_\epsilon(\ln(xae))},
\end{aligned}$$

where S_ϵ is the scale function related to the spectrally negative Lévy process ξ^ϵ , see [Ber96, Chap. VII, Sec. 2]. Let $m^\epsilon = \mathbb{E}[\xi_1^\epsilon]$, see the discussion after (5.21). From $\xi_1 \leq \xi_1^\epsilon$ then $0 < m \leq m^\epsilon$. Since in the construction of ξ^ϵ we truncate jumps smaller than $-\epsilon^{-1}$ we have that, for all $\epsilon < w$, $0 < m \leq m^\epsilon \leq m^w$ and plainly ψ^ϵ, ψ^w are analytic in \mathbb{C} , see the discussion around (5.22). Moreover, by differentiating twice ψ^ϵ is easily seen to be strictly convex in \mathbb{R} and for any $u \in \mathbb{R}$, ψ^ϵ has the form, see (5.21),

$$\psi^\epsilon(u) = m^\epsilon u + \frac{\sigma^2}{2} u^2 + \int_0^{\frac{1}{\epsilon}} (e^{-uy} - 1 + uy) \Pi(dy),$$

where the truncation of jumps is reflected in the measure Π and hence in the integration bounds. Hence, for $\epsilon < w$, using $e^{-x} + x - 1 \geq 0$, $x \leq 0$, we get that, for $u < 0$,

$$\psi^\epsilon(u) - \psi^w(u) = (m^\epsilon - m^w)u + \int_{\frac{1}{w}}^{\frac{1}{\epsilon}} (e^{-uy} - 1 + uy) \Pi(dy) \geq 0.$$

Also from (5.21),

$$e^{\psi^\epsilon(u)} = \mathbb{E}\left[e^{u\xi_1^\epsilon}\right] \geq \mathbb{E}\left[e^{u\xi_1^\epsilon} \mathbb{I}_{\{\xi_1^\epsilon < 0\}}\right]$$

and we get that $\lim_{u \rightarrow -\infty} \psi^\epsilon(u) = \infty$. Hence using the latter, the fact that $(\psi^\epsilon)'(0^+) = \mathbb{E}[\xi_1^\epsilon] = m^\epsilon > 0$ and ψ^ϵ is strictly convex with $\psi^\epsilon(0) = 0$ we deduce that there exists unique $u_0^\epsilon < 0$ such that $\psi^\epsilon(u_0^\epsilon) = 0$. From $\psi^\epsilon(u) - \psi^w(u) \geq 0$, for

$0 < \epsilon < w$ and $u < 0$, and by convexity, we get $u_0^w \leq u_0^\epsilon < 0$, where $\psi^w(u_0^w) = 0$. We proceed to show that $u_0^* = \lim_{\epsilon \rightarrow 0} u_0^\epsilon = d_\phi$. Recall that $\psi(u) = u\phi(u)$ and if $u < d_\phi \leq 0$ then $\psi(u) \in (0, \infty]$, where $\psi(u) = \infty$ if $\int_0^\infty (e^{-uy} - 1 + uy) \Pi(dy) = \infty$, see (5.21). It is clear that $u_0^* \geq d_\phi$, as otherwise, for any $u \in (u_0^*, d_\phi)$, we must have $\lim_{\epsilon \rightarrow 0} \psi^\epsilon(u) = \psi(u) \in (0, \infty]$, whilst the convexity of ψ^ϵ and $\psi^\epsilon(u_0^\epsilon) = \psi^\epsilon(0) = 0$ lead to $\psi^\epsilon < 0$ on $(u_0^\epsilon, 0)$ and thus $\lim_{\epsilon \rightarrow 0} \psi^\epsilon(u) \leq 0$ if $u \in (u_0^*, d_\phi)$. The case $d_\phi = 0$ follows as $u_0^\epsilon < 0$. Assume that $d_\phi < u_0^* \leq 0$. Then for any $u \in (d_\phi, u_0^*)$, we have that $\lim_{\epsilon \rightarrow 0} \psi^\epsilon(u) = \psi(u) < 0$ which contradicts $u_0^\epsilon \uparrow u_0^*$, $\psi^\epsilon > 0$ on $(-\infty, u_0^\epsilon)$. Thus, $d_\phi = u_0^* = \lim_{\epsilon \rightarrow 0} u_0^\epsilon$. Finally, since according to [Ber96, Ch. VII, Sec. 2], we have $S_\epsilon(x) = \frac{1}{m^\epsilon} \mathbb{P}\left(-\underline{\xi}_\infty^\epsilon \leq x\right)$, where we let $\underline{\xi}_\infty^\epsilon = \inf_{t \geq 0} \xi_t^\epsilon$ be the global infimum of ξ^ϵ , we deduce that

$$S_\epsilon(\ln(xae)) - S_\epsilon(\ln(xa)) = \frac{1}{m^\epsilon} \left(\mathbb{P}\left(-\underline{\xi}_\infty^\epsilon > \ln(xa)\right) - \mathbb{P}\left(-\underline{\xi}_\infty^\epsilon > \ln(xae)\right) \right),$$

which together with $\mathbb{P}(-\underline{\xi}_\infty^\epsilon > x) \approx C^\epsilon e^{u_0^\epsilon x}$, see e.g. [BD94], yield, recalling that $u_0^\epsilon < 0$,

$$x^{-u_0^\epsilon} (S_\epsilon(\ln(xae)) - S_\epsilon(\ln(xa))) \approx \frac{C^\epsilon}{m^\epsilon} a^{u_0^\epsilon} (1 - e^{u_0^\epsilon}).$$

Then, since plainly

$$S_\epsilon(\infty) = \lim_{x \rightarrow \infty} S_\epsilon(\ln(xae)) = \frac{1}{m^\epsilon} \lim_{x \rightarrow \infty} \mathbb{P}\left(-\underline{\xi}_\infty^\epsilon \leq \ln(xae)\right) = \frac{1}{m^\epsilon},$$

we deduce from (5.35) that

$$x^{-u_0^\epsilon} \mathbb{P}\left(\bar{T}_{-\ln(xa)}^\epsilon < \bar{T}_1^\epsilon\right) = x^{-u_0^\epsilon} \left(\frac{S_\epsilon(\ln(xae)) - S_\epsilon(\ln(xa))}{S_\epsilon(\ln(xae))} \right) \approx C^\epsilon a^{u_0^\epsilon} (1 - e^{u_0^\epsilon}).$$

Therefore, since $a > e$ is fixed, the latter asymptotic relation allows the successive usage of (5.34), (5.32), (5.31) together with $\widehat{\nu} > 0$ on $(\mathfrak{r}^{-1}, \infty)$ and $\widehat{\nu} \in \mathfrak{C}((\mathfrak{r}^{-1}, \infty))$, see Proposition 5.2.1(3) to derive (5.19) as $u_0^\epsilon \uparrow d_\phi$ when $\epsilon \rightarrow 0$. The second statement follows readily from [Riv05, Lemma 4], which states that $\mathbb{P}(I_\psi > x) = \int_x^\infty \widehat{\nu}(y) dy \approx Cx^{-u_0}$ and an application of the monotone decreasing density theorem which yields $\widehat{\nu}(x) \approx Cu_0 x^{-u_0-1}$ and is valid since $\widehat{\nu}$ is ultimately monotone as a unimodal distribution.

5.2.5. Proof of Theorem 5.0.4. The first statement follows by combining (5.19) with (5.12), i.e. $\nu(x) = x^{-2} \widehat{\nu}_1(x^{-1})$ and $\widehat{\nu}_1$ the density of $I_{\mathcal{T}_1\psi}$, $\mathcal{T}_1\psi(u) = u\phi_1(u) = u\phi(u+1)$, and observing that (5.19) holds for $a < d_{\phi_1} = d_\phi - 1$. The second statement is deduced from (5.20) with $u_0 = 1$ since $\mathcal{T}_1\psi(-1) = -\phi(0) = 0$ and $(\mathcal{T}_1\psi)'(-1^+) = -\phi'(0^+) < \infty$.

5.3. Proof of Theorem 5.0.1(3)

The proof of this claim requires a combination of several types of results. Since some of them will be useful later in this work, we state them separately. We start by discussing various further functional properties of $\widehat{\nu}$.

5.3.1. Some useful facts on positive self-decomposable laws. The following lemma is essentially due to Sato and Yamazato [SY78] where we denote by \mathcal{F}_f the bilateral moment generating function (resp. Fourier transform) of a function $f : \mathbb{R} \mapsto \mathbb{R}$, i.e. for some real u (resp. some imaginary number $u = ib, b \in \mathbb{R}$.)

$$\mathcal{F}_f(u) = \int_{\mathbb{R}} e^{uy} f(y) dy.$$

LEMMA 5.3.1. (1) Let $\psi \in \mathcal{N}_{\infty}(m)$. Then there exists a decreasing positive sequence $\underline{a}_{\widehat{\nu}} = (a_n)_{n \geq 0}$ with $a_0 = \infty > a_1 > \dots > a_n \dots > \frac{1}{\tau}$, such that, for every $n \geq 1$,

$$\widehat{\nu}^{(n)} > 0 \text{ on } \left(\frac{1}{\tau}, a_n\right) \text{ and } \widehat{\nu}^{(n)} < 0 \text{ on } (a_n, a_{n-1}).$$

Moreover, for any $n \geq 0$, we have $a_n > \beta_n + \frac{1}{\tau}$ where $\beta_n = \sup\{y > 0; \widehat{\kappa}(y) \geq n\}$.

(2) Let $\psi \in \mathcal{N}_{\infty}^c(m)$. Then there exists a sequence $\underline{a}_{\widehat{\nu}} = (a_n)_{0 \leq n \leq N_{\tau}}$, such that $a_0 = \infty > a_1 > \dots > a_{N_{\tau}} > \frac{1}{\tau}$ enjoying the same properties as the sequence in (1) above if and only if $N_{\tau} \geq 1$, where recall that $N_{\tau} = \lceil \frac{\overline{\Pi}(0^+)}{\tau} \rceil - 1$.

(3) Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. Then, we have, for any $n \geq 0$ and $x > 0$,

$$(5.36) \quad x\widehat{\nu}^{(n+1)}(x) = -(n+1)\widehat{\nu}^{(n)}(x) + \int_0^{\infty} \left(\widehat{\nu}^{(n)}(x-y) - \widehat{\nu}^{(n)}(x)\right) d\widehat{\kappa}(y).$$

(4) For any $0 \leq u < \frac{\overline{\Pi}(0^+)}{\tau}$, one has

$$(5.37) \quad |\mathcal{F}_{\widehat{\nu}}(ib)| \stackrel{\pm\infty}{\asymp} o(|b|^{-u}),$$

with the convention that $\frac{\overline{\Pi}(0^+)}{\tau} = \infty$ when $\overline{\Pi}(0^+) = \infty$.

Proof. Let us first assume that $\psi \in \mathcal{N}_{\infty}(m)$. According to Lemma 5.2.5(2) we have that $\widehat{\kappa}(0^+) = \infty$ and hence $I_{\psi} \in \mathbf{I}_6$ in the sense of [SY78, p.275]. We derive from [SY78, Theorem 5.1(ii)] the existence of the sequence $\underline{a}_{\widehat{\nu}}$ with the inequality $a_n > \beta_n + \frac{1}{\tau}$. Otherwise, if $\psi \in \mathcal{N}_{\infty}^c(m)$ and $\frac{\overline{\Pi}(0^+)}{\tau} > 1$ the existence of the sequence $\underline{a}_{\widehat{\nu}}$ and β_n , for $1 \leq n \leq N_{\tau}$, follows from [SY78, Theorem 5.1(i)]. Next from [SY78, p.297, (5.4)] equation (5.36) holds for every $n \geq 0$ when $\widehat{\kappa}(0^+) = \infty$ and $\frac{1}{\tau} = 0$ (i.e. $\gamma_0 = 0$ in [SY78, p.297, (5.4)]), which is the case once $\psi \in \mathcal{N}_{\infty, \infty}(m)$. Item (4) follows from [SY78, Lemma 2.4, (2.17)]. \square

We continue with the following substantial but easy estimates which will be used at several places in this work.

LEMMA 5.3.2. Let $\psi \in \mathcal{N}(m)$. If $N_{\tau} \geq 1$ then, for all $n < N_{\tau} \leq \infty$ and $k < N_{\tau} - n \leq \infty$,

$$(5.38) \quad \lim_{x \downarrow \frac{1}{\tau}} \left(x - \frac{1}{\tau}\right)^{-k} \widehat{\nu}^{(n)}(x) = \lim_{x \downarrow 0} x^{-k} \widehat{\nu}^{(n)}\left(x + \frac{1}{\tau}\right) = 0.$$

Moreover, for any $D \in (0, 1)$, $n < N_{\mathfrak{r}}$ and $x \in (0, D(a_{n+1} - \frac{1}{\mathfrak{r}}))$, with $\underline{a}_{\widehat{\nu}} = (a_n)_{n \leq N_{\mathfrak{r}}}$ as in Lemma 5.3.1, we have that

$$(5.39) \quad \left| \widehat{\nu}^{(n)} \left(\frac{1}{\mathfrak{r}} + x \right) \right| \leq \left(\frac{D}{1-D} \right)^n n! x^{-n} \widehat{\nu} \left(\frac{1}{\mathfrak{r}} + \frac{x}{D} \right).$$

Proof. Let first $\psi \in \mathcal{N}_{\infty}(m)$. We first prove (5.38) for $k = 0$ and $n \in \mathbb{N}$ which follows easily from the fact that, since $\overline{\Pi}(0^+) = \infty$ and hence $N_{\mathfrak{r}} = \infty$ in this case, $\widehat{\nu} \in \mathcal{C}^{\infty}(\mathbb{R})$ and $\text{Supp } I_{\psi} = [0, \infty]$, see Proposition 5.2.1. Let now set $n = 0$ and $k \geq 1$. The existence of the sequence $\underline{a}_{\widehat{\nu}}$ defined in Lemma 5.3.1, implies that $\widehat{\nu}^{(l)}$, $l \geq 1$, increases on $(\frac{1}{\mathfrak{r}}, a_{l+1})$. Therefore,

$$\widehat{\nu}^{(l-1)}(x) = \int_{\frac{1}{\mathfrak{r}}}^x \widehat{\nu}^{(l)}(y) dy \leq \left(x - \frac{1}{\mathfrak{r}} \right) \widehat{\nu}^{(l)}(x), \text{ for } x \in \left(\frac{1}{\mathfrak{r}}, a_{l+1} \right),$$

since from (5.38) with $k = 0$ and $l \in \mathbb{N}$ we have that

$$(5.40) \quad \widehat{\nu}^{(l)} \left(\frac{1}{\mathfrak{r}} \right) = 0.$$

Repeating this argument k times, we get, for any $x \in (\frac{1}{\mathfrak{r}}, a_{k+1})$, that

$$\left(x - \frac{1}{\mathfrak{r}} \right)^{-k} \widehat{\nu}(x) \leq \widehat{\nu}^{(k)}(x),$$

which proves (5.38) for $k \geq 1$ and $n = 0$ using (5.40). Before proving the remaining case, we establish (5.39). From (5.40), for any $n \geq 0$, $x \in (\frac{1}{\mathfrak{r}}, a_{n+1})$ and $D \in (0, 1)$,

$$\begin{aligned} \widehat{\nu}(x) &= \int_{\frac{1}{\mathfrak{r}}}^x \int_{\frac{1}{\mathfrak{r}}}^{y_1} \dots \int_{\frac{1}{\mathfrak{r}}}^{y_{n-1}} \widehat{\nu}^{(n)}(y_n) dy_n \dots dy_1 \\ &\geq \int_{\frac{1}{\mathfrak{r}}+D(x-\frac{1}{\mathfrak{r}})}^x \int_{\frac{1}{\mathfrak{r}}+D(x-\frac{1}{\mathfrak{r}})}^{y_1} \dots \int_{\frac{1}{\mathfrak{r}}+D(x-\frac{1}{\mathfrak{r}})}^{y_{n-1}} \widehat{\nu}^{(n)}(y_n) dy_n \dots dy_1 \\ &\geq \frac{(1-D)^n}{n!} \left(x - \frac{1}{\mathfrak{r}} \right)^n \widehat{\nu}^{(n)} \left(\frac{1}{\mathfrak{r}} + D \left(x - \frac{1}{\mathfrak{r}} \right) \right), \end{aligned}$$

where we have employed that $\widehat{\nu}^{(n)}$ is increasing on $(\frac{1}{\mathfrak{r}}, a_{n+1})$, see Lemma 5.3.1(1). This is (5.39). Finally, the proof of (5.38) for $n, k > 0$, follows from the case $n = 0$ combined with (5.39). Next, when $\psi \in \mathcal{N}_{\infty}^c(m)$ and $N_{\mathfrak{r}} \geq 1$, (5.38) is immediate thanks to (5.18) of Proposition 5.2.1(4) which furnishes an explicit asymptotic of $\widehat{\nu}$ at $1/\mathfrak{r}$ and $\text{Supp } I_{\psi} = [\frac{1}{\mathfrak{r}}, \infty)$, which was proved in Proposition 5.2.5, and, Proposition 5.2.1(3) which asserts that in this case $\widehat{\nu} \in \mathcal{C}^{N_{\mathfrak{r}}-1}(\mathbb{R})$. \square

Recall that $W_{\phi}(p) = \prod_{j=1}^n \phi(j)$. We proceed with the following lemmas.

LEMMA 5.3.3. (1) Let $\psi \in \mathcal{N}(m)$. Then, for any $p = 1, 2, \dots$, we have that

$$\widehat{\nu}_p(x) = \frac{\widehat{\nu}_{p-1}(x)}{\phi(p-1)x} = \dots = \frac{\widehat{\nu}(x)}{mW_{\phi}(p)x^p}, \quad x > 0,$$

is the density of the variable $I_{\mathcal{T}_p\psi}$ where $\mathcal{T}_p\psi(u) = u\phi(u+p) \in \mathcal{N}(\phi(p))$ and we set $\widehat{\nu}_0 = \widehat{\nu}$. Moreover, $\lim_{u \rightarrow \infty} \frac{\mathcal{T}_p\psi(u)}{u} = \lim_{u \rightarrow \infty} \phi(u+p) = \phi(\infty)$, with the obvious notation $\overline{\Pi}_p(0^+) = \overline{\Pi}(0^+)$ and $\text{Supp } I_{\mathcal{T}_p\psi} = [\frac{1}{\mathfrak{r}}, \infty)$.

(2) For any $\psi \in \mathcal{N}$, we have, writing, for any $p = 0, 1, \dots$, ν_p the density of the variable $V_{\mathcal{T}_p\psi}$, $\nu_p(x) = x^{-2}\widehat{\nu}_{p+1}(x^{-1})$, for any $p = 1, 2, \dots$,

$$\nu_p(x) = \frac{x\nu_{p-1}(x)}{\phi(p)} = \dots = \frac{x^p\nu(x)}{W_\phi(p+1)},$$

where we set $\nu = \nu_0$, and, $\text{Supp } V_{\mathcal{T}_p\psi} = [0, \tau]$.

Proof. First, we recall from Proposition 4.1.1(9), that, for any $p \in \mathbb{N}$, $\mathcal{T}_p\psi(u) = u\phi(u+p) \in \mathcal{N}(\phi(p))$, where we recall that $\phi(0) = m > 0$ as $\psi = \mathcal{T}_0\psi \in \mathcal{N}(m)$. Next, from the expression of the entire moment of $I_{\mathcal{T}_p\psi}$ given in (5.13), we have that, for any $n = 1, 2, \dots$,

$$\mathbb{E} \left[I_{\mathcal{T}_p\psi}^{-n} \right] = \phi(p) \prod_{k=1}^{n-1} \phi(k+p) = \prod_{k=1}^n \phi(k+p-1) = \frac{1}{\phi(p-1)} \mathbb{E} \left[I_{\mathcal{T}_{p-1}\psi}^{-n-1} \right]$$

from where we get the first statement by remarking that the mappings $(\mathcal{T}_p)_{p \geq 0}$ form a semigroup on \mathcal{N} , i.e. for any $\psi \in \mathcal{N}$, $p, r \geq 0$,

$$\mathcal{T}_p \circ \mathcal{T}_r\psi(u) = \mathcal{T}_p \left(\frac{\cdot}{\cdot+r}\psi(\cdot+r) \right) (u) = \frac{u}{u+r+p}\psi(u+r+p) = \mathcal{T}_{p+r}\psi(u) \in \mathcal{N}$$

and $I_{\mathcal{T}_1\psi}^{-1} \stackrel{(d)}{=} V_\psi$, see (5.14), is moment determinate and hence $I_{\mathcal{T}_p\psi}^{-1}$, $p = 1, 2, \dots$, are moment determinate. Also a simple integration of (4.9) yields that $\bar{\Pi}_p(y) = e^{-py}\bar{\Pi}(y)$, $y > 0$. Finally, since $\phi_p(\infty) = \phi(\infty) = \tau$ we conclude $\text{Supp } I_{\mathcal{T}_p\psi} = [\frac{1}{\tau}, \infty)$ from Proposition 5.2.1(2) and item (1) is thus settled. Item (2) follows readily from the previous one combined with the identity $\nu(x) = \frac{1}{x^2}\widehat{\nu}_1(\frac{1}{x})$ that is (5.12) of Proposition 5.1.1. \square

LEMMA 5.3.4. Let $\psi \in \mathcal{N}(m)$ and $N_\tau \geq 1$. Then, with the notation of Lemma 5.3.3, we have, for all $n \leq N_\tau - 1$, $p = 0, 1, \dots$, and $a > \frac{1}{2}$,

$$(5.41) \quad \int_0^\infty \left| x^{-a}\widehat{\nu}_p^{(n)}(x) \right| dx < \infty,$$

where we recall that $\widehat{\nu}_0 = \widehat{\nu}$.

Proof. Let us prove that (5.41) holds for $p = 0$. The proof for $p = 1, 2, \dots$ follows since $\mathcal{T}_p\psi \in \mathcal{N}(\phi(p)) \subset \mathcal{N}(m)$. From Parseval's identity and the fact that, for all $0 \leq n \leq N_\tau - 1$, $\widehat{\nu}^{(n)} \in \mathcal{C}(\mathbb{R})$, see items (1) and (3) of Proposition 5.2.1, one gets, for $n \leq N_\tau - 1$, that

$$(5.42) \quad \int_{\frac{1}{\tau}}^\infty \left| \widehat{\nu}^{(n)}(x) \right|^2 dx = \int_0^\infty \left| \widehat{\nu}^{(n)}(x) \right|^2 dx = \int_{-\infty}^\infty b^{2n} |\mathcal{F}_{\widehat{\nu}}(ib)|^2 db < \infty,$$

where the finiteness follows from (5.37). To prove (5.41) in the case $N_\tau = \infty$ and $p = 0$, we use the notation of Lemma 5.3.2 with $D_n = \frac{1}{\tau} + D(a_{n+1} - \frac{1}{\tau})$, and, for the first inequality below, the estimate (5.39) and the continuity of $\widehat{\nu}^{(n)}$, and, for the second one the classical Cauchy-Schwarz inequality, to get, for any $a > \frac{1}{2}$, $n \geq 0$,

that

$$\begin{aligned}
\int_{\frac{1}{\mathfrak{r}}}^{\infty} |x^{-a} \widehat{\nu}^{(n)}(x)| dx &= \int_{\frac{1}{\mathfrak{r}}}^{D_n} |x^{-a} \widehat{\nu}^{(n)}(x)| dx + \int_{D_n}^{\infty} |x^{-a} \widehat{\nu}^{(n)}(x)| dx \\
&\leq \int_0^{D_n - \frac{1}{\mathfrak{r}}} \left| \left(x + \frac{1}{\mathfrak{r}} \right)^{-a} \widehat{\nu}^{(n)} \left(x + \frac{1}{\mathfrak{r}} \right) \right| dx + \int_{D_n}^{\infty} |x^{-a} \widehat{\nu}^{(n)}(x)| dx \\
&\leq C_n \int_0^{D_n - \frac{1}{\mathfrak{r}}} x^{-n-a} \widehat{\nu} \left(\frac{x}{D} + \frac{1}{\mathfrak{r}} \right) dx + C_{n,a} \left(\int_0^{\infty} |\widehat{\nu}^{(n)}(x)|^2 dx \right)^{\frac{1}{2}} \\
&< \infty,
\end{aligned}$$

where $C_n > 0$, $C_{n,a} > 0$ and the finiteness follows from (5.38) and the estimate (5.42). To prove (5.41) in the case $N_{\mathfrak{r}} < \infty$, we follow the same line of reasoning by recalling, from Proposition 5.2.1(3), that, for any $n \leq N_{\mathfrak{r}} - 1 < \infty$, $\widehat{\nu}^{(n)} \in \mathcal{C}(\mathbb{R})$, and (5.42) holds for $n \leq N_{\mathfrak{r}} - 1$. The only modification, thanks to $\frac{1}{\mathfrak{r}} > 0$, lies in the estimate

$$\int_{\frac{1}{\mathfrak{r}}}^{\infty} |x^{-a} \widehat{\nu}^{(n)}(x)| dx \leq \mathfrak{r}^a \int_{\frac{1}{\mathfrak{r}}}^{D_n} |\widehat{\nu}^{(n)}(x)| dx + \int_{D_n}^{\infty} |x^{-a} \widehat{\nu}^{(n)}(x)| dx < \infty.$$

□

We are ready to prove the polynomial decay of the modulus of the Mellin transform of ν along imaginary lines, that is (5.3) in Theorem 5.0.1(3). With the notation and the claims of Lemma 5.3.3, for any $p = 0, 1, \dots$, we write

$$\chi_{p+1}(y) = e^{-y} \widehat{\nu}_{p+1}(e^{-y}) = \frac{\widehat{\nu}_p(e^{-y})}{\phi(p)} = e^y \nu_p(e^y)$$

and we have $N_{\mathfrak{r}}$ and $\frac{\overline{\Pi}(0^+)}{\mathfrak{r}}$ invariant in p for any ν_p , $\widehat{\nu}_p$ and $\text{Supp } \chi_{p+1} = (-\infty, \ln \mathfrak{r}]$, if $\mathfrak{r} < \infty$, and $\text{Supp } \chi_{p+1} = (-\infty, \infty)$, otherwise. Then, we get, from Proposition 5.2.1(3), that $\chi_{p+1} \in \mathcal{C}^{N_{\mathfrak{r}}}(\mathbb{R} \setminus \{\ln \mathfrak{r}\})$ if $\mathfrak{r} < \infty$ and $\chi_{p+1} \in \mathcal{C}^{\infty}(\mathbb{R})$, otherwise. Thus, for any $a \geq n$, applying the Faa di Bruno formula to $\widehat{\nu}_p(e^{-y})$ we get with some $C_n > 0$,

$$\begin{aligned}
\int_{-\infty}^{\ln \mathfrak{r}} e^{ay} |\chi_{p+1}^{(n)}(y)| dy &\leq \frac{C_n}{\phi(p)} \int_{-\infty}^{\ln \mathfrak{r}} e^{ay} \sum_{k=0}^n |e^{-ky} \widehat{\nu}_p^{(k)}(e^{-y})| dy \\
&= \frac{C_n}{\phi(p)} \int_{\frac{1}{\mathfrak{r}}}^{\infty} \sum_{k=0}^n |x^{k-1-a} \widehat{\nu}_p^{(k)}(x)| dx < \infty,
\end{aligned}$$

where the finiteness follows from (5.41) since from $a \geq n$ then $a + 1 - k \geq 1$, $0 \leq k \leq n$, and $n \leq N_{\mathfrak{r}} - 1$. Thus, we have for all $p = 0, 1, \dots$ and $0 \leq n \leq N_{\mathfrak{r}} - 1$, that, for all $a \geq n$, the mapping $y \mapsto \chi_{p+1,a}^{(n)}(y) = (e^{ay} \chi_{p+1}(y))^{(n)} \in L^1(\mathbb{R})$ and by the Riemann-Lebesgue lemma applied to the Fourier transform of $(e^{ay} \chi_{p+1}(y))^{(n)}$ we get that

$$(5.43) \quad |\mathcal{F}_{\chi_{p+1,a}}(ib)| \stackrel{\pm\infty}{=} o(|b|^{-n}).$$

On the other hand, writing $z = a + ib$, $a \geq n$, $b \in \mathbb{R}$, we have because $\chi_{p+1}(y) = e^y \nu_p(e^y)$

$$\mathcal{F}_{\chi_{p+1,a}}(ib) = \mathcal{F}_{\chi_{p+1}}(z) = \int_{\mathbb{R}} e^{zy} \chi_{p+1}(y) dy = \int_0^{\infty} x^z \nu_p(x) dx = \mathcal{M}_{V_{\mathfrak{r}_p \psi}}(z + 1).$$

From Lemma 5.3.3(2) and then from the functional equation (5.1) we deduce that

$$\mathcal{M}_{V_{\tau_p\psi}}(z+1) = \frac{\mathcal{M}_{V_\psi}(z+p+1)}{W_\phi(p+1)} = \frac{\phi(z+p)\dots\phi(z)}{W_\phi(p+1)}\mathcal{M}_{V_\psi}(z).$$

Henceforth, we get that, for any $p = 0, 1, \dots, n \leq N_\tau - 1$ and $a \geq n$,

$$(5.44) \quad \lim_{|b| \rightarrow \infty} |b|^n \frac{\prod_{j=0}^p |\phi(a+j+ib)|}{W_\phi(p+1)} |\mathcal{M}_{V_\psi}(a+ib)| = \lim_{|b| \rightarrow \infty} |b|^n |\mathcal{F}_{\chi_{p+1,a}}(ib)| = 0,$$

where the last identity follows from in (5.43). Fix $a' > d_\phi$ and $p' \in \mathbb{N}$, $p' \geq 1$, such that $a = a' + p' > n$. We have, from (5.1), that

$$\mathcal{M}_{V_\psi}(p' + a' + ib) = \phi(p' + a' - 1 + ib) \cdots \phi(a' + ib) \mathcal{M}_{V_\psi}(a' + ib).$$

Then applying (5.44) with $a = a' + p' > n$ and $p = p'$ we arrive at

$$(5.45) \quad \lim_{|b| \rightarrow \infty} |b|^n \frac{\prod_{j=0}^{2p'} |\phi(a+j+ib)|}{W_\phi(p+1)} |\mathcal{M}_{V_\psi}(a' + ib)| = 0.$$

Since $\phi \in \mathcal{B}_\mathcal{N}$ from Proposition 4.1.1(8) the Lévy measure associated to ϕ is $\bar{\Pi}(y)dy$, $y > 0$, that is, absolutely continuous with respect to the Lebesgue measure. Hence, if in addition $\bar{\Pi}(0^+) = \int_0^\infty \bar{\Pi}(y)dy = \infty$, then [Sat99, Theorem 27.7] implies that the underlying descending ladder height process η^- , see (5.24), is absolutely continuous with respect to the Lebesgue measure. Therefore, from (5.24) and the Riemann-Lebesgue Theorem, for any fixed $a' > d_\phi$,

$$\lim_{|b| \rightarrow \infty} e^{-\phi(a'+ib)} = \lim_{|b| \rightarrow \infty} \mathbb{E} \left[e^{-(a'+ib)\eta_1^-} \right] = \lim_{|b| \rightarrow \infty} \int_0^\infty e^{-a'x-ibx} \mathbb{P}(\eta_1^- \in dx) = 0$$

and we conclude that $\lim_{|b| \rightarrow \infty} |\phi(a' + ib)| = \infty$. Otherwise, if $\bar{\Pi}(0^+) < \infty$, i.e. $\phi(\infty) = \tau < \infty$, then elementary from the representation of ϕ in Proposition 4.1.1(8) and the Riemann-Lebesgue Theorem $\lim_{|b| \rightarrow \infty} \phi(a' + ib) = \phi(\infty) = \tau$. In both cases, the limit is non-zero and we deduct from (5.45) that

$$\lim_{|b| \rightarrow \infty} |b|^n |\mathcal{M}_{V_\psi}(a' + ib)| = 0,$$

which furnishes the proof of our claim for $n \leq N_\tau - 1$.

5.3.2. Proof of Theorem 5.0.2(2a). We are ready to complete Theorem 5.0.2(2a). From Lemma 5.2.7 it remains to show that if $\psi \in \mathcal{N}_{\infty, \infty}$ then, for any $n \in \mathbb{N}$, $\lim_{x \rightarrow \infty} \nu^{(n)}(x) = 0$. Note that the case when $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$ is settled since ν has support $(0, \tau)$, $\tau < \infty$, in this case. The following estimate which holds in a more general context provides this property.

LEMMA 5.3.5. *Let $\psi \in \mathcal{N}$ with $N_\tau > 2$. Then, for any $x > 0$, $n, k \in \mathbb{N}$, with $k < N_\tau - 2$, and any $\bar{a} > d_\phi$, there exists a constant $C = C_{n,k,\bar{a}} > 0$ such that for any $x > 0$*

$$(5.46) \quad \left| (x^n \nu(x))^{(k)} \right| \leq C x^{n-k-\bar{a}}.$$

REMARK 5.3.6. Note that the bounds at ∞ presented in Lemma 9.2.1 below are far more precise as they depend on ν . However, (5.46) are uniform on \mathbb{R}_+ and allow us to have some grip on the behaviour at zero of ν and its derivatives.

Proof. Recall the Pochhammer symbol $(n - z)_k = \prod_{j=0}^{k-1} (n - j - z)$. Fix $\bar{a} > d_\phi$ and $k < N_\tau - 2$. Then from Theorem 5.0.1(3), $|z|^k |\mathcal{M}_{V_\psi}(z)| = O(|z|^{-2})$ along $\mathbb{C}_{\bar{a}}$. Therefore, $|(n - z)_k \mathcal{M}_{V_\psi}(z)|$ is integrable along $\mathbb{C}_{\bar{a}}$ and multiplying by x^n the Mellin inversion formula for ν , see (1.36), and differentiating it k times to get

$$\begin{aligned} |(x^n \nu(x))^{(k)}| &= \left| \frac{1}{2\pi i} \int_{\bar{a}-i\infty}^{\bar{a}+i\infty} x^{n-k-z} (n-z)_k \mathcal{M}_{V_\psi}(z) dz \right| \\ &\leq x^{n-k-\bar{a}} \frac{1}{2\pi} \int_{-\infty}^{\infty} |(n-\bar{a}-ib)_k \mathcal{M}_{V_\psi}(\bar{a}+ib)| db \end{aligned}$$

and the finiteness of the last integral gives (5.46). \square

5.4. Small asymptotic behaviour of $\hat{\nu}$ and of its successive derivatives

Finally, we provide an extremely precise asymptotic results for $\hat{\nu}(x)$ (resp. $\chi(x)$) and its successive derivatives when x tends to 0 (resp. ∞). We stress that, for only some isolated cases, one can find in the literature information about the behaviour of $\ln \int_0^x \hat{\nu}(y) dy$, see [SY78, Theorem 5.2]. On the other hand, we are not aware of any instances of a class of probability density functions for which such a precise asymptotic estimate has been provided. The novelty of our approach seems to come from the fact that we are able to describe the asymptotic behaviour of the Mellin transform of $\hat{\nu}$, i.e. \mathcal{M}_{I_ψ} , along the negative real line and imaginary lines, together with some fine distributional properties such as log-concavity of $\hat{\nu}$ and related to it quantities at 0.

THEOREM 5.4.1. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m) = \mathcal{N}_{\infty, \infty} \cap \mathcal{N}(m)$. Recalling that $\varphi : [m, \infty) \mapsto [0, \infty)$ stands for the inverse function of ϕ , we have, with $\tilde{C}_\psi > 0$, that*

$$(5.47) \quad \hat{\nu}(x) \underset{0}{\sim} \frac{\tilde{C}_\psi m}{\sqrt{2\pi}} \frac{1}{x} \sqrt{\varphi' \left(\frac{1}{x} \right)} e^{-\int_m^{\frac{1}{x}} \varphi(r) \frac{dr}{r}}.$$

Moreover, for any $n \geq 0$, we have that

$$(5.48) \quad \hat{\nu}^{(n)}(x) \underset{0}{\sim} \frac{\tilde{C}_\psi m}{\sqrt{2\pi}} \frac{1}{x^{n+1}} \varphi^n \left(\frac{1}{x} \right) \sqrt{\varphi' \left(\frac{1}{x} \right)} e^{-\int_m^{\frac{1}{x}} \varphi(r) \frac{dr}{r}}.$$

Finally, the following relation holds true

$$(5.49) \quad \chi(x) \underset{\infty}{\sim} \frac{\tilde{C}_\psi m}{\sqrt{2\pi}} \sqrt{\varphi'(e^x)} e^{-\int_m^{e^x} \varphi(r) \frac{dr}{r}}.$$

REMARK 5.4.2. We note that when $\psi \in \mathcal{N}_{\infty, \infty}(m)$ for the positive self-decomposable variable I_ψ we have that $\hat{\kappa}(0^+) = \infty$, recall (5.10) for the definition of $\hat{\kappa}$. This is strictly beyond [Sat99, Theorem 53.6] which discusses only the case when $0 < \hat{\kappa}(0^+) < \infty$. Our case, i.e. $\hat{\kappa}(0^+) = \infty$ and $\hat{\kappa}(0^-) = 0$, seems to have been studied only in [Wat96, Lemma 2.5] but merely on the *log*-scale and when $\hat{\kappa} \in RV_\alpha(0)$, that is when $\hat{\kappa}$ is regularly varying at zero.

5.4.1. A non-classical Tauberian theorem. The proof of Theorem 5.4.1 is based on an improved version of a Tauberian theorem that was originally proved by Balkema et al. in [BKS93, Theorem 4.4]. It is a non-classical Tauberian in the sense that it relates the upper tail behaviour of the bilateral Laplace transform to the upper tail behaviour of the associated probability density function. For the sake of clarity, we state and prove below a slight generalization and an adapted version of this Tauberian theorem which is more suitable to our context and allows its application for the successive derivatives of the density function. In particular, the original result [BKS93, Theorem 4.4] is stated for the density of a probability distribution and a minor device allows us to extend it to real-valued functions which are ultimately positive. We proceed by introducing some notation and terminologies. Let $G : \mathbb{R} \mapsto \mathbb{R}$ be a convex function. Then

$$(5.50) \quad G^*(y) = \sup_{u \in \mathbb{R}} \{yu - G(u)\}$$

is called the Legendre transform or the complex conjugate of G . If $G^{(2)} > 0$ on \mathbb{R} then the supremum is attained at u such that $y = G'(u)$. We say that G is *asymptotically parabolic* if $G^{(2)} > 0$ on \mathbb{R} and its scale function $s_G = 1/\sqrt{G^{(2)}}$ is *self-neglecting*, i.e.

$$(5.51) \quad \lim_{u \rightarrow \infty} \frac{s_G(u + \mathfrak{a}s_G(u))}{s_G(u)} = 1$$

uniformly on bounded intervals of the real variable \mathfrak{a} .

Next, we say that a function F has a *very thin tail* if there exists $y_0 \in \mathbb{R}$ such that

$$(5.52) \quad F(y) > 0 \quad \text{for all } y > y_0,$$

$$(5.53) \quad \lim_{y \rightarrow \infty} F(y)e^{ky} = 0 \quad \text{for all } k \in \mathbb{N}.$$

Finally, we recall the notation for a function f and some real u ,

$$\mathcal{F}_f(u) = \int_{\mathbb{R}} e^{uy} f(y) dy.$$

We are ready to state and prove the following adapted version of [BKS93, Theorem 4.4].

PROPOSITION 5.4.3. *Let us assume that the following conditions hold.*

(a) *Let $f : \mathbb{R} \rightarrow \mathbb{R}$ such that $f(y) > 0$, for all $y > a \in \mathbb{R}$. Set $f_a(y) = f(y)\mathbb{I}_{\{y>a\}}$ and assume that $F_a(x) = \int_x^\infty f_a(y) dy$ has a very thin tail and*

$$\mathcal{F}_f(u) \approx \mathcal{F}_{f_a}(u).$$

(b) *f is log-concave in a neighbourhood of ∞ .*

(c) *We have*

$$\mathcal{F}_f(u) = \int_{\mathbb{R}} e^{uy} f(y) dy \approx \beta(u)e^{G(u)},$$

with G being asymptotically parabolic and

$$(5.54) \quad \lim_{u \rightarrow \infty} \frac{\beta(u + \mathfrak{a}s_G(u))}{\beta(u)} = 1$$

uniformly on bounded intervals of the real variable \mathbf{a} .

Then

$$(5.55) \quad f(y) \approx \frac{1}{\sqrt{2\pi}} \frac{\beta^*(y)}{s_{G^*}(y)} e^{-G^*(y)},$$

where G^* is the Legendre transform of G , s_{G^*} is its own scale function and $\beta^*(y) = \beta(u)$ is defined via the following relation between y and u , that is $y = G'(u)$.

Proof. Let us first assume that f is a probability density function. We show that all the conditions of the statement imply (5.55) via an application of [BKS93, Theorem 4.4]. First, we show that f is of Gaussian tail in the sense of [BKS93, Definition on p.389], i.e. f is of very thin tail itself and $f(y) \approx e^{-\tilde{G}(y)}$ with some asymptotically parabolic function \tilde{G} . This would trigger the first important condition, that is [BKS93, Theorem 4.4(1)]. To prove that f is of Gaussian tail we will invoke [BKS93, Theorem 2.2]. To do so we see that the log-concavity of f implies the condition (2.1) of [BKS93, Theorem 2.2]. The other condition of [BKS93, Theorem 2.2], i.e. U_τ , see [BKS93, (1.6)] for definition, to be asymptotically normal, i.e. [BKS93, (1.10)] to be fulfilled, holds true thanks to the assumption for the self-neglecting property of s_G and the behaviour of \mathcal{F}_f at infinity, that is the condition (c) above, which is enough for [BKS93, Theorem 1.2, (1.10)] to be true, as G is asymptotically parabolic and β satisfies (5.54), i.e. β is flat with respect to G in the sense of [BKS93]. All this verifies that f is of Gaussian tail and the proof follows in this case since the other conditions of [BKS93, Theorem 4.4] follow from the assumptions of our theorem. It is clear that in the end [BKS93, Theorem 4.4.] serves its purpose only to elucidate the form of \tilde{G} . Now, with

$$0 < C_a^{-1} = \int_a^\infty f_a(y) dy = \int_a^\infty f(y) dy < \infty,$$

we have that $\bar{f}_a(y) = C_a f_a(y)$ is a probability density function. Moreover, under the condition (a) about the asymptotic behaviour of \mathcal{F}_{f_a} , we can apply the previous reasoning to \bar{f}_a , after checking easily that all conditions are satisfied, to get,

$$\bar{f}_a(y) \approx \frac{C_a}{\sqrt{2\pi}} \frac{\beta^*(y)}{s_{G^*}(y)} e^{-G^*(y)}$$

which, from the definition of \bar{f}_a and f_a , completes the proof. \square

5.5. Proof of Theorem 5.4.1

Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$ and thus $\hat{\nu} \in \mathcal{C}^\infty(\mathbb{R})$ from Proposition 5.2.1(3). We aim at applying the Tauberian result of Proposition 5.4.3 to the continuous functions f_n , defined, for any $n \geq 0$, by

$$(5.56) \quad f_n(y) = e^y v_n(e^y) = e^{-ny} \hat{\nu}^{(n)}(e^{-y}), \quad y \in \mathbb{R}.$$

Throughout this proof we use the Pochhammer notation $(u)_n = \frac{\Gamma(u+1)}{\Gamma(u+1-n)} = u(u-1)\dots(u-n+1)$, $u \geq 0$, $n \in \mathbb{N}$. We start with the log-concavity property of f_n in a neighborhood of ∞ , that is condition (b), and postpone to the next sections the proof of the two remaining Tauberian conditions, namely (a) and (c).

5.5.1. Condition (b): the log-concavity property of $\widehat{\nu}^{(n)}$ and f_n .

PROPOSITION 5.5.1. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. There exists a decreasing positive sequence $\underline{a}_{\widehat{\nu}} = (a_n)_{n \geq 0}$ with $a_0 = \infty > a_1 > \dots > a_n \dots > 0$, such that, for every $n \geq 0$,*

$$x \mapsto \widehat{\nu}^{(n)}(x) \text{ is log-concave on } (0, a_{n+1}).$$

Even more, for any $n \geq 0$, the mapping

$$y \mapsto f_n(y) = e^{-ny} \widehat{\nu}^{(n)}(e^{-y}) \text{ is also log-concave on } (-\log(a_{n+1}), \infty).$$

REMARK 5.5.2. The statement that $\widehat{\nu}$ is log-concave on $(0, a_1)$ is proved in [SY78, Theorem 1.3 (vii)] and improved to strict log-concavity in [SY78, Theorem 4.2]. Clearly, the log-concavity of $y \mapsto \widehat{\nu}^{(n)}(e^{-y})$ is a stronger assertion than the log-concavity of the derivatives $\widehat{\nu}^{(n)}(x)$ and the proof using Lemma 5.5.7 below suggests that it is the more natural result.

For the reader's convenience we split the (lengthy) proof of Proposition 5.5.1 into several intermediate results.

LEMMA 5.5.3. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$ and thus $\widehat{\kappa}(0^+) = \infty$. For any $j \geq 1$, set*

$$\widehat{\kappa}_j(y) = \widehat{\kappa}(\max(y, j^{-1})), \quad y > 0.$$

Then, for any $j \geq 1$, there exists $I_j \in \mathfrak{L}_+$ such that its law is absolutely continuous with a density v_j which is characterized, for any $b \in \mathbb{R}$, by

$$(5.57) \quad \mathcal{F}_{v_j}(ib) = \int_0^\infty e^{ibx} v_j(x) dx = e^{-\phi_j(-ib)} = e^{\int_0^\infty (e^{iby} - 1) \frac{\widehat{\kappa}_j(y)}{y} dy}.$$

Moreover, for any $j \geq 1$, the following holds.

- (1) $\text{Supp } I_j = [0, \infty)$.
- (2) $v_j \in \mathcal{C}^{N_j}(\mathbb{R}_+) \cap \mathcal{C}^{N_j-1}(\mathbb{R})$, where $N_j = \lceil \widehat{\kappa}_j(0^+) \rceil - 1 = \lceil \widehat{\kappa}(j^{-1}) \rceil - 1$.
- (3) *Then there exists a decreasing positive sequence $\underline{a}_{v_j} = (a_n(j))_{1 \leq n \leq N_j}$ with $a_0(j) = \infty > a_1(j) > \dots > a_{N_j}(j) > 0$, such that, for every $0 \leq n \leq N_j$,*

$$v_j^{(n)} > 0 \text{ on } (0, a_n(j)) \text{ and } v_j^{(n)} < 0 \text{ on } (a_n(j), a_{n-1}(j)).$$

- (4) *Let, for $n \leq N_j$, $\beta_n(j) = \sup \{y > 0; \widehat{\kappa}_j(y) \geq n\}$. Also, for any $n \geq 1$,*

$$(5.58) \quad \underline{a}_n := \varliminf_{j \rightarrow \infty} a_n(j) \geq \beta_n = \sup \{y > 0; \widehat{\kappa}(y) \geq n\} > 0,$$

- (5) *We have, for any $x > 0$,*

$$(5.59) \quad xv_j^{(1)}(x) = (\widehat{\kappa}_j(0^+) - 1)v_j(x) + \int_0^x v_j(x-y) d\widehat{\kappa}_j(y).$$

REMARK 5.5.4. We stress that I_j does not belong to the subclass of positive self-decomposable laws arising from $\psi \in \mathcal{N}_{\infty}^c(m)$, that have support $[\frac{1}{\tau}, \infty)$, $\tau < \infty$, see Proposition 5.2.1(2).

Proof. First, since for every $j \geq 1$, $\widehat{\kappa}_j$ defines a non-increasing function on \mathbb{R}_+ with $\widehat{\kappa}_j(0^+) = \widehat{\kappa}(j^{-1}) < \infty$, the right-hand side of (5.57) is the Fourier transform of a random variable $I_j \in \mathfrak{L}_+$, whose distributions, as mentioned earlier, are absolutely continuous. Furthermore, since $\widehat{\kappa}_j(0^+) < \infty$, therefore $I_j \in \mathbf{I}_5$ in the sense of [SY78, p.275]. Since $\widehat{\delta}_j = 0$ then $\text{Supp } I_j = [0, \infty)$ follows from [SY78, Theorem 1.3, (vii)]. Also from [SY78, Theorem 1.2] then $v_j \in \mathcal{C}^{\mathbf{N}_j}(\mathbb{R}_+)$, from [SY78, Theorem 5.1, (i)] then $v_j \in \mathcal{C}^{\mathbf{N}_j-1}(\mathbb{R})$ and from [SY78, Theorem 5.1, (i)] there exists $\underline{a}_{v_j} = (a_n(j))_{1 \leq n \leq \mathbf{N}_j}$ with $a_0(j) = \infty > a_1(j) > \dots > a_{\mathbf{N}_j}(j) > 0$ such that $v_j^{(n)}$ is positive on $(0, a_n(j))$ and negative on $(a_n(j), a_{n-1}(j))$, for $1 \leq n \leq \mathbf{N}_j$. Thanks to [SY78, Theorem 5.1, (5.2) and (5.3)],

$$a_n(j) > \beta_n(j) = \sup \{y > 0; \widehat{\kappa}_j(y) \geq n\}$$

and the fact that $\beta_n(j) = \beta_n = \sup \{y > 0; \widehat{\kappa}(y) \geq n\}$ for all j big enough, for $n \leq \mathbf{N}_j$ gives (5.58). Item (5) is [SY78, Corollary 2.1 (2.34)], wherein $\gamma_0 = 0$ since we have no drift in the exponent of (5.57) and the integration is on $(0, \infty)$. \square

LEMMA 5.5.5. Fix $n \in \mathbb{N}$. If $\widehat{\kappa}_j(0^+) > 2n + 5$, then $v_j^{(n)}$ is log-concave on $(0, a_{n+1}(j))$.

Proof. Let $\widehat{\kappa}_j(0^+) > 2n+5$. Since $v_j \in \mathcal{C}^{\mathbf{N}_j-1}(\mathbb{R})$ and $d\widehat{\kappa}_j(y) \equiv 0$, $y \in (0, j^{-1})$, see Lemma 5.5.3 by differentiating (5.59), we get for (at least) $k \leq \mathbf{N}_j - 2 = \lceil \widehat{\kappa}_j(0^+) \rceil - 3$ and any $x > 0$,

$$(5.60) \quad xv_j^{(2)}(x) = (\widehat{\kappa}_j(0^+) - 2)v_j^{(1)}(x) + \int_0^x v_j^{(1)}(x-y)d\widehat{\kappa}_j(y),$$

$$(5.61) \quad xv_j^{(k+2)}(x) = (\widehat{\kappa}_j(0^+) - 2 - k)v_j^{(k+1)}(x) + \int_0^x v_j^{(k+1)}(x-y)d\widehat{\kappa}_j(y).$$

Also (5.60) and (5.61) are elementary consequence to [SY78, p. 297, l.-1 and (5.4)] and a simple integration. Next, note that $d\widehat{\kappa}_j(y) = 0 dy$, for $y \in (0, j^{-1})$ and therefore on $x < j^{-1}$ (5.60) reduces to the simple differential equation

$$xv_j^{(2)}(x) = (\widehat{\kappa}_j(0^+) - 2)v_j^{(1)}(x).$$

Hence $v_j(x) = Cx^{\widehat{\kappa}_j(0^+)-1}$ for some $C > 0$ and $x \in (0, j^{-1})$. Therefore setting

$$V_j^{(n)}(x) := v_j^{(n+2)}(x)v_j^{(n)}(x) - (v_j^{(n+1)}(x))^2 = \left(v_j^{(n)}(x)\right)^2 \frac{d}{dx} \left(\frac{v_j^{(n+1)}(x)}{v_j^{(n)}(x)}\right).$$

we get then that $V_j^{(n)} < 0$ on $(0, j^{-1})$. Moreover, since $\widehat{\kappa}_j(0^+) > 2n+5$, $x \mapsto v_j^{(n)}(x)$ is strictly log-concave and increasing for $x \in (0, j^{-1})$ and thus from Lemma 5.5.3(3) and the properties of the sequence \underline{a}_{v_j} we conclude $a_{n+1}(j) \geq j^{-1}$. Assume now $a_{n+1}(j) > j^{-1}$ and that there exists $x_0 \in (j^{-1}, a_{n+1}(j))$ such that $V_j^{(n)}(x_0) = 0$ and $V_j^{(n)}(x) < 0$ for $x \in (0, x_0)$. Multiplying first (5.61) for $k = n$ at the point x_0 by $v_j^{(n)}(x_0)$ and then for $k = n - 1$ by $-v_j^{(n+1)}(x_0)$ at the point x_0 and adding the

two expressions, we get the identity

$$(5.62) \quad \begin{aligned} x_0 V_j^{(n)}(x_0) &= -v_j^{(n)}(x_0)v_j^{(n+1)}(x_0) \\ &+ \int_{j^{-1}}^{x_0} v_j^{(n)}(x_0 - y)v_j^{(n)}(x_0) \left(\frac{v_j^{(n+1)}(x_0 - y)}{v_j^{(n)}(x_0 - y)} - \frac{v_j^{(n+1)}(x_0)}{v_j^{(n)}(x_0)} \right) d\widehat{\kappa}_j(y). \end{aligned}$$

From $x_0 < a_{n+1}(j)$ and $x \mapsto v_j^{(n)}(x)$ is increasing on $(0, a_{n+1}(j))$ as $x \mapsto v_j^{(n+1)}(x)$ is positive on $(0, a_{n+1}(j))$, see Lemma 5.5.3(3), then $-v_j^{(n)}(x_0)v_j^{(n+1)}(x_0) < 0$ in (5.62). Also $V_j^{(n)} < 0$ on $(0, x_0)$ means that $x \mapsto \frac{v_j^{(n+1)}(x)}{v_j^{(n)}(x)}$ decreases on $(0, x_0)$. Thus the integrand in (5.62) is positive. However, as $\widehat{\kappa}_j$ is non-increasing on \mathbb{R}_+ , $d\widehat{\kappa}_j(y)$ defines a negative measure for $y \in [j^{-1}, \infty)$ and $d\widehat{\kappa}_j(y) = 0 dy$, for $y \in (0, j^{-1})$, we deduce that the integral in (5.62) is negative too. Therefore $V_j^{(n)}(x_0) < 0$ which is in contradiction with the definition of x_0 , i.e. $V_j^{(n)}(x_0) = 0$. Thus, if $\widehat{\kappa}_j(0^+) > 2n + 5$, $v_j^{(n)}$ is log-concave on $(0, a_{n+1}(j))$, completing the proof. \square

LEMMA 5.5.6. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. Then, for all $n \geq 0$, $\lim_{j \rightarrow \infty} v_j^{(n)} = \widehat{\nu}^{(n)}$ uniformly on \mathbb{R}_+ . Consequently, $\widehat{\nu}^{(n)}$ is log-concave on $(0, a_{n+1})$.*

Proof. Since $\psi \in \mathcal{N}_{\infty, \infty}(m) \subset \mathcal{N}_{\infty}(m)$ then $\widehat{\kappa}(0^+) = \infty$, see Proposition 5.2.1(1). Pick j such that $\widehat{\kappa}_j(0^+) > 2n + 5$ and v_j corresponding to it as in Lemma 5.5.3. From [SY78, Lemma 2.4 (2.17)] for the first relation and (5.37) for the second we get that

$$(5.63) \quad |\mathcal{F}_{v_j}(ib)| \stackrel{\pm\infty}{\asymp} o(|b|^{-2n-5}) \quad \text{and} \quad |\mathcal{F}_{\widehat{\nu}}(ib)| \stackrel{\pm\infty}{\asymp} o(|b|^{-u}), \quad \text{for any } u > 0.$$

Therefore, for any $b \in \mathbb{R}$ and $k \leq n + 2$,

$$\int_0^{\infty} e^{ibx} v_j^{(k)}(y) dy = i^k b^k \mathcal{F}_{v_j}(ib) \quad \text{and} \quad \int_0^{\infty} e^{ibx} \widehat{\nu}^{(k)}(y) dy = i^k b^k \mathcal{F}_{\widehat{\nu}}(ib).$$

Therefore, by Fourier inversion, for any $k \leq n + 2$,

$$\sup_{x>0} \left| v_j^{(k)}(x) - \widehat{\nu}^{(k)}(x) \right| \leq \frac{1}{2\pi} \int_{-\infty}^{\infty} |b|^k |\mathcal{F}_{v_j}(ib) - \mathcal{F}_{\widehat{\nu}}(ib)| db.$$

However, (5.63) implies, for $|b| > 1$ and some $C > 0$, that $|\mathcal{F}_{v_j}(ib) - \mathcal{F}_{\widehat{\nu}}(ib)| \leq C|b|^{-2n-5}$. Also, from the definition $\widehat{\kappa}_j(y) = \widehat{\kappa}(\max(y, j^{-1}))$, (5.10) and (5.57), we have $\lim_{j \rightarrow \infty} \mathcal{F}_{v_j}(ib) = \mathcal{F}_{\widehat{\nu}}(ib)$ pointwise. Therefore the dominated convergence theorem applies and yields $v_j^{(k)} \rightarrow \widehat{\nu}^{(k)}$ uniformly on \mathbb{R}_+ , for any $k \leq n + 2$.

This allows to show that $(\log v_j^{(n)})'' \rightarrow (\log \widehat{\nu}^{(n)})''$ for $x \in (0, \underline{a}_{n+1})$ with $\underline{a}_{n+1} = \varliminf_{j \rightarrow \infty} a_{n+1}(j)$ defined in (5.58). Since from Lemma 5.5.5 we have that $v_j^{(n)}$ is log-concave on $(0, a_{n+1}(j))$ then $\widehat{\nu}^{(n)}$ is log-concave on $(0, \underline{a}_{n+1})$. Next, we show that $\underline{a}_{n+1} \geq a_{n+1}$, $n > 0$. From [SY78, Lemma 4.4], we conclude that

$$\lim_{j \rightarrow \infty} a_1(j) = a_1 = \underline{a}_1.$$

Assume that $\underline{a}_{n+1} < a_{n+1}$, $n > 0$. Then Lemma 5.5.3(3) and j such that $\widehat{\kappa}_j(0^+) > 2n + 5$ imply that $v_j^{(n+1)}(a_{n+1}(j)) = 0$. Choose $(a_{n+1}(j))_{l \geq 1}$ such that

$$\lim_{l \rightarrow \infty} a_{n+1}(j_l) = \underline{a}_{n+1}.$$

Then the uniform convergence of $v_j^{(n+1)}$ to $\widehat{\nu}^{(n+1)}$, $\widehat{\nu}^{(n+1)} \in \mathcal{C}^\infty(\mathbb{R})$, see Proposition 5.2.1(3) and $\widehat{\nu}^{(n+1)} > 0$ on $(0, a_{n+1})$, see Lemma 5.3.1(1) yield with the help of $v_j^{(n+1)}(a_{n+1}(j)) = 0$ that

$$0 = \lim_{l \rightarrow \infty} v_{j_l}^{(n+1)}(a_{n+1}(j_l)) = \widehat{\nu}^{(n+1)}(\underline{a}_{n+1}) > 0,$$

that is a contradiction to $\underline{a}_{n+1} < a_{n+1}$. Therefore $\widehat{\nu}^{(n)}$ is log-concave on $(0, a_{n+1})$. \square

Finally, to prove the condition (b), it remains to show that, for all $n \geq 0$,

$$(5.64) \quad y \mapsto f_n(y) = e^{-ny} \widehat{\nu}^{(n)}(e^{-y}) \text{ is log-concave on } (-\log a_{n+1}, \infty).$$

Since multiplication by e^{-ny} does not alter the log-concavity property we check that

$$(5.65) \quad y \mapsto \widehat{\nu}^{(n)}(e^{-y}) \text{ is log-concave on } (-\log a_{n+1}, \infty).$$

Relation (5.65) comes from Lemma 5.5.7 which can be seen to be applicable to $\widehat{\nu}^{(n)}$ as follows. First set $v = \widehat{\nu}^{(n)}$ and choose $A = a_{n+1} > 0$ and note, from Lemma 5.3.1(1), that $v > 0$ on $(0, a_n) \supseteq (0, a_{n+1})$. Next, clearly from Proposition 5.2.1(3), $v \in \mathcal{C}^\infty(\mathbb{R})$ and $v = v' = v'' = 0$ on $(-\infty, 0]$. Also Lemma 5.5.6 gives that v is log-concave on $(0, a_{n+1})$. Finally, from (5.36), since (5.66) and (5.67) hold with $u = n + 1$ and with $\kappa = \widehat{\kappa}$ and $\int_0^\infty \widehat{\kappa}(y) dy < \infty$, we conclude that $\left| \int_0^1 y d\widehat{\kappa}(y) \right| < \infty$ and therefore (5.65) indeed holds true.

LEMMA 5.5.7. *Let $A > 0$, $v : (0, A) \mapsto \mathbb{R}^+$ and $v, v', v'' : [-\infty, A] \mapsto \mathbb{R}$. Assume that $v(x) = v'(x) = v''(x) = 0$, for $x \leq 0$, $v \in \mathcal{C}^2((-\infty, A])$ and v is log-concave on $(0, A)$. Moreover, let for any $x \in (0, A)$ and some $u \in \mathbb{R}_+$, the following equations hold*

$$(5.66) \quad xv'(x) = -uv(x) + \int_0^\infty (v(x-y) - v(x)) d\kappa(y),$$

$$(5.67) \quad xv''(x) = -(u+1)v'(x) + \int_0^\infty (v'(x-y) - v'(x)) d\kappa(y),$$

where $\kappa : \mathbb{R}_+ \mapsto \mathbb{R}_+$ is non-increasing and $\left| \int_0^1 y d\kappa(y) \right| < \infty$. Then, even $y \mapsto v(e^{-y})$ is log-concave on $(-\ln A, \infty)$.

Proof. Differentiating $y \mapsto \log v(e^{-y})$ twice and performing a change of variables it suffices to show that for any $x \in (0, A)$

$$V(x) = xv(x)v'(x) + x^2v''(x)v(x) - x^2(v'(x))^2 \leq 0.$$

Multiplying (5.66) (resp. (5.67)) by $v(x)$ (resp. $xv(x)$) and then (5.66) separately by $-xv'(x)$, we evaluate V by substitution to get that for any $x \in (0, A)$,

$$\begin{aligned}
V(x) &= -uv^2(x) + v(x) \int_0^\infty v(x-y) - v(x) d\kappa(y) \\
&\quad - (u+1)v'(x)xv(x) + xv(x) \int_0^\infty v'(x-y) - v'(x) d\kappa(y) \\
&\quad + uxv'(x)v(x) - xv'(x) \int_0^\infty (v(x-y) - v(x)) d\kappa(y) \\
&= -xv'(x)v(x) - uv^2(x) + v(x) \int_0^\infty (v(x-y) - v(x)) d\kappa(y) \\
&\quad + xv(x) \int_0^\infty (v'(x-y) - v'(x)) d\kappa(y) - xv'(x) \int_0^\infty (v(x-y) - v(x)) d\kappa(y).
\end{aligned}$$

Note that from the equation (5.66) we also get that

$$-v'(x)xv(x) - uv^2(x) = -v(x) \int_0^\infty (v(x-y) - v(x)) d\kappa(y)$$

and thus

$$V(x) = xv(x) \int_0^\infty (v'(x-y) - v'(x)) d\kappa(y) - xv'(x) \int_0^\infty (v(x-y) - v(x)) d\kappa(y).$$

Note, that since $v > 0$ on $(0, A)$ and $v = v' = 0$ on $(-\infty, 0]$, we have that, for $x \in (0, A)$,

$$\begin{aligned}
V(x) &= x \left(\int_0^x v(x-y)v'(x-y) - v'(x)v(x-y) d\kappa(y) \right) \\
&= x \left(\int_0^x v(x-y)v'(x-y) \left(\frac{v'(x-y)}{v(x-y)} - \frac{v'(x)}{v(x)} \right) d\kappa(y) \right) \leq 0,
\end{aligned}$$

since by assumption v is log-concave and thus $x \mapsto (\log v)'(x)$ is non-increasing on $(0, A)$, $v \geq 0$ and $d\kappa(y)$ defines a negative measure because κ is non-increasing. This settles the claim that $y \mapsto v(e^{-y})$ is log-concave on $(-\ln A, \infty)$. \square

5.5.2. Condition (a). For any $n \geq 0$, set $\bar{a}_n = \min(a_n, \frac{1}{2})$ (where we recall that the a'_n 's are defined in Lemma 5.3.1(1)) and

$$(5.68) \quad f_{a,n}(y) = f_n(y) \mathbb{I}_{\{y > \ln \bar{a}_n^{-1}\}}, \quad y \in \mathbb{R},$$

where from (5.64), $f_n(y) = e^{-ny} \hat{\nu}^{(n)}(e^{-y})$. Then, according to Lemma 5.3.1, since $\hat{\nu}^{(n)} > 0$ on $(0, a_n)$ we have clearly $f_{a,n} > 0$ on $(\ln \bar{a}_n^{-1}, \infty)$. We shall check that condition (a) holds for $f_{a,n}$. We start with the following result.

LEMMA 5.5.8. *For any $n \geq 0$ both $f_{a,n}$ and $F_{a,n}(x) = \int_x^\infty f_{a,n}(y) dy$ have a very thin tail.*

Proof. Fix $n \geq 0$. First, plainly $f_{a,n}$ satisfies (5.52). Moreover, from (5.38) of Lemma 5.3.2, we have for any $k \geq 1$,

$$(5.69) \quad \lim_{y \rightarrow \infty} e^{ky} f_{a,n}(y) = \lim_{y \rightarrow \infty} e^{(k-n)y} \hat{\nu}^{(n)}(e^{-y}) = \lim_{x \rightarrow 0} x^{-k+n} \hat{\nu}^{(n)}(x) = 0,$$

i.e. (5.53) holds and therefore $f_{a,n}$ has a very thin tail. We now turn our attention to $F_{a,n}$. The fact that $F_{a,n}(x) = \int_x^\infty f_{a,n}(y)dy$ is ultimately positive is clear. Finally, using again the definition of $f_{a,n}$, we observe that, for any $k \geq 1$ and for $x > -\ln a_n$,

$$\begin{aligned} e^{kx} F_{a,n}(x) &= e^{kx} \int_x^\infty e^{-ny} \widehat{\nu}^{(n)}(e^{-y}) dy = e^{kx} \int_0^{e^{-x}} u^{n-1} \widehat{\nu}^{(n)}(u) du \\ &\leq e^{kx-nx} \widehat{\nu}^{(n-1)}(e^{-x}). \end{aligned}$$

We stress that for the last inequality we have used that $\widehat{\nu}^{(n-1)}$ is non-decreasing on $(0, a_n)$, see Lemma 5.3.1. This together with (5.69) completes the proof. \square

To prove that condition (a) holds, it remains to check the asymptotic equivalence of the bilateral moment generating functions, i.e. $\mathcal{F}_{f_n}(u) \approx \mathcal{F}_{f_{a,n}}(u)$. We stress that this equivalence requires a specific behaviour of the Mellin transform of $\widehat{\nu}$ along imaginary lines. We split this verification into several intermediate lemmas as their results are required for other purposes. We shall need the following technical lemma which allows to study f_n via v_n , as defined in (5.56).

LEMMA 5.5.9. *Let $\psi \in \mathcal{N}(m)$ and $N_\tau > 2$. Then, for any $x > 0$, $n \in \mathbb{N}$ with $n < N_\tau - 2$ and for any $\underline{a} < 1 - d_\phi$, there exists a constant $C = C_{n,\underline{a}} > 0$*

$$(5.70) \quad \left| \widehat{\nu}^{(n)}(x) \right| \leq C x^{-\underline{a}-n}.$$

Proof. First, from Theorem 5.0.1(1), we have $\mathcal{M}_{V_\psi}(z) \in \mathcal{A}_{(d_\phi, \infty)}$. Then, from (5.12) of Proposition 5.1.1 and

$$\widehat{\nu}_1(x) = \frac{1}{\phi(0)} \frac{\widehat{\nu}(x)}{x}$$

which is contained in Lemma 5.3.3(1), one can deduce the following relationship between Mellin transforms

$$(5.71) \quad \mathcal{M}_{V_\psi}(z) = \mathcal{M}_{I_{\tau_1\psi}}(2-z) = \frac{1}{m} \mathcal{M}_{I_\psi}(1-z), \quad z \in \mathbb{C}_{(d_\phi, \infty)},$$

where we recall that, for any $\psi \in \mathcal{N}(m)$ we have $\phi(0) = m > 0$, $\mathcal{T}_1\psi(u) = u\phi(u+1) \in \mathcal{N}(\phi(1))$ and $\mathcal{M}_{I_\psi}(1-z) = \mathbb{E} \left[I_\psi^{-z} \right]$. The requirement $N_\tau > 2$ is to ensure that the Mellin transform \mathcal{M}_{V_ψ} (resp. \mathcal{M}_{I_ψ}) is absolutely integrable against any polynomial of order $k < N_\tau - 2$ on \mathbb{C}_a with $a > d_\phi$ (resp. smaller than $1 - d_\phi$), see (5.3). The estimate is obtained by following a similar line of computation as for the proof of Lemma 5.3.5. \square

Our next Lemma specifies some preliminary properties of the bilateral moment transforms to interval of \mathbb{R} . Recall that $(u)_n = \frac{\Gamma(u+1)}{\Gamma(u+1-n)}$.

LEMMA 5.5.10. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. For all $n \geq 0$, we have, for $u > 0$, that*

$$(5.72) \quad \mathcal{F}_{f_n}(u) = \mathcal{M}_{v_n}(u+1) = m (u)_n \mathcal{M}_{V_\psi}(u+1) = (u)_n \mathcal{M}_{I_\psi}(-u),$$

where \mathcal{M}_{v_n} is the moment transform of v_n , recalling from (5.56) that

$$f_n(y) = e^y v_n(e^y) = e^{-ny} \widehat{\nu}^{(n)}(e^{-y}), \quad y \in \mathbb{R}.$$

Proof. From (5.71) we have that $\mathcal{M}_{I_\psi} \in \mathcal{A}_{(-\infty, 1)}$ and since

$$v_n(x) = x^{-n-1} \widehat{\nu}^{(n)}(x^{-1})$$

we get that, for $u > 0$,

$$(5.73) \quad \mathcal{M}_{v_n}(u+1) = \int_0^\infty y^{-u-1} y^n \widehat{\nu}^{(n)}(y) dy.$$

To use integrations by parts for the last integral, we need to verify in the process that

$$\lim_{x \rightarrow 0} x^{n-u-l+1} \widehat{\nu}^{(n-l)}(x) = \lim_{x \rightarrow \infty} x^{n-u-l+1} \widehat{\nu}^{(n-l)}(x) = 0.$$

The first limit follows from (5.38). The second one is deduced from (5.70) since

$$\left| x^{n-u-l+1} \widehat{\nu}^{(n-l)}(x) \right| \leq C x^{-u+1-\underline{a}}$$

holds for any $x > 0$, $\underline{a} < 1 - d_\phi$ and we choose $\underline{a} \in (1 - u, 1 - d_\phi)$. Therefore, we get that

$$\mathcal{M}_{v_n}(u+1) = (u)_n \mathcal{M}_{I_\psi}(-u) = m(u)_n \mathcal{M}_{V_\psi}(u+1),$$

where for the latter we have invoked (5.71). Finally, from the definition (5.56) we conclude (5.72) via the computation

$$\begin{aligned} \mathcal{F}_{f_n}(u) &= \int_{\mathbb{R}} e^{uy} f_n(y) dy = \int_0^\infty x^{u-n-1} \widehat{\nu}^{(n)}(x^{-1}) dx = \int_0^\infty x^u v_n(x) dx \\ &= \mathcal{M}_{v_n}(u+1). \end{aligned}$$

□

We now check the asymptotic equivalence of the upper tail of the moment generating functions of f_n and $f_{a,n}$, see (5.68), which completes the verification of condition (a).

LEMMA 5.5.11. *Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. Then, for any $n \geq 0$, we have*

$$(5.74) \quad \mathcal{F}_{f_{a,n}}(u) \approx \mathcal{F}_{f_n}(u).$$

Proof. Set

$$\bar{f}_{a,n}(y) = f_n(y) - f_{a,n}(y) = f_n(y) - f_n(y) \mathbb{I}_{\{y > \ln \bar{a}_n^{-1}\}}, \quad y \in \mathbb{R}, \quad n \geq 0,$$

see (5.68), and note that, for $u > 0$,

$$\begin{aligned} \mathcal{F}_{f_n}(u) &= \mathcal{F}_{f_{a,n}}(u) + \mathcal{F}_{\bar{f}_{a,n}}(u) = \mathcal{M}_{v_n}(u+1) \\ &= \int_0^{\bar{a}_{n+1}} y^{n-1-u} \widehat{\nu}^{(n)}(y) dy + \int_{\bar{a}_{n+1}}^\infty y^{n-1-u} \widehat{\nu}^{(n)}(y) dy, \end{aligned}$$

where we recall that $\bar{a}_n = \min(a_n, \frac{1}{2})$ and we have used (5.72) and (5.73) for the second and third identities. Moreover, we get, with $u > 2n$, that

$$(5.75) \quad \left| \int_{\bar{a}_{n+1}}^\infty y^{n-1-u} \widehat{\nu}^{(n)}(y) dy \right| \leq C \bar{a}_{n+1}^{n+1-u} \int_{\bar{a}_{n+1}}^\infty y^{-2} |\widehat{\nu}^{(n)}(y)| dy$$

with $C > 0$ throughout being a generic constant depending at most on n . However, it is immediate from (5.41) of Lemma 5.3.4 with $p = 0$, $a = 2$ that

$$\int_{\bar{a}_{n+1}}^{\infty} y^{-2} |\widehat{\nu}^{(n)}(y)| dy \leq C,$$

and thus from (5.75) we get that

$$(5.76) \quad \left| \mathcal{F}_{\bar{f}_{a,n}}(u) \right| = \left| \int_{\bar{a}_{n+1}}^{\infty} x^{n-1-u} \widehat{\nu}^{(n)}(x) dx \right| \leq C \bar{a}_{n+1}^{n+1-u}.$$

Furthermore, using that $\widehat{\nu}^{(n)}$ is non-decreasing and positive on $(0, a_{n+1})$, see Lemma 5.3.1(1), $\widehat{\nu}^{(n-1)}(0) = 0$, see (5.40) with $\frac{1}{\tau} = 0$ and $n-1-u < -n < 0$ when $u > 2n$, we obtain that

$$\begin{aligned} \mathcal{F}_{f_{a,n}}(u) &= \int_0^{\bar{a}_{n+1}} x^{n-1-u} \widehat{\nu}^{(n)}(x) dx \geq \int_0^{\bar{a}_{n+1}^2} x^{n-1-u} \widehat{\nu}^{(n)}(x) dx \\ &\geq \bar{a}_{n+1}^{2n-2-2u} \int_0^{\bar{a}_{n+1}^2} \widehat{\nu}^{(n)}(x) dx = \bar{a}_{n+1}^{2n-2-2u} \widehat{\nu}^{(n-1)}(\bar{a}_{n+1}^2). \end{aligned}$$

Clearly, then, from (5.76), we have that

$$\left| \frac{\mathcal{F}_{\bar{f}_{a,n}}(u)}{\mathcal{F}_{f_{a,n}}(u)} \right| \leq \frac{C}{\widehat{\nu}^{(n-1)}(\bar{a}_{n+1}^2)} \frac{\bar{a}_{n+1}^{n+1-u}}{\bar{a}_{n+1}^{2n-2-2u}} = \frac{C}{\widehat{\nu}^{(n-1)}(\bar{a}_{n+1}^2)} \bar{a}_{n+1}^{u-n+3} \stackrel{u \rightarrow \infty}{\asymp} o(1),$$

since $\bar{a}_{n+1} \leq \frac{1}{2}$. Therefore, $\mathcal{F}_{f_n}(u) \approx \mathcal{F}_{f_{a,n}}(u)$. \square

5.5.3. Condition (c). Since we have shown that the conditions (a) and (b) hold, we proceed by discussing the asymptotic properties of $\mathcal{F}_{f_n}(u)$. The next result is in fact a restatement of (5.2) of Theorem 5.0.1 for f_n .

LEMMA 5.5.12. *For any $n \geq 0$ and $\psi \in \mathcal{N}(m)$*

$$(5.77) \quad \mathcal{F}_{f_n}(u) \approx C_{\psi} m(u)_n^+ \sqrt{\phi(u)} e^{G(u)} = \beta(u) e^{G(u)},$$

where $(u)_n^+ = (u)_n \vee 1$, $C_{\psi} > 0$ and $G(u) = \int_1^u \ln \phi(r) dr$.

Proof. Relation (5.77) follows immediately from (5.72) and (5.2). \square

The last lemma shows that the functional form of the asymptotic of $\mathcal{F}_{f_n}(u)$, $u \rightarrow \infty$, is as required in condition (c). In order to check that the terms appearing in the large asymptotic behaviour (5.77) of \mathcal{F}_{f_n} satisfy the conditions in (c), we need to modify the function G which is simply defined on \mathbb{R}_+ . To this end, we note that, for $u > 0$,

$$(5.78) \quad G^{(2)}(u) = \frac{\phi'(u)}{\phi(u)}.$$

Since $G^{(2)}(1) > 0$, we may for instance consider the function

$$\begin{aligned} G_1(u) &= G(u) \mathbb{I}_{\{u \geq 1\}} \\ &+ \left(\frac{G^{(2)}(1)}{2} u^2 + \left(G'(1) - G^{(2)}(1) \right) u + G(1) - G'(1) + \frac{G^{(2)}(1)}{2} \right) \mathbb{I}_{\{u < 1\}}, \end{aligned}$$

which is strictly convex on the real line. However, since the remaining conditions involve only the large asymptotic behaviour of G_1 or related functionals at ∞ , there is no loss of generality of keeping the notation G throughout. In the same vein we simply write

$$(5.79) \quad s_G(u) = \frac{1}{\sqrt{G^{(2)}(u)}} = \sqrt{\frac{\phi(u)}{\phi'(u)}}.$$

We first have the following result.

LEMMA 5.5.13. *For any $\psi \in \mathcal{N}(m)$ it is true that $\lim_{u \rightarrow \infty} s_G(u) = \infty$.*

Proof. Since $\phi \in \mathcal{B}_{\mathcal{N}}$ then Proposition 4.1.1(8) gives $\mu(dy) = \bar{\Pi}(y)dy$. Then, from item (4) and (4.3) of Proposition 4.1.1 we have that $s_G^{-2}(u) = \frac{\phi'(u)}{\phi(u)}$ is the Laplace transform on \mathbb{R}_+ of the measure $(\sigma^2\delta_0 + \tilde{\mu}) \star \Upsilon$ with $\tilde{\mu}(dy) = y\mu(dy)$, $y > 0$. Therefore $\lim_{u \rightarrow \infty} s_G^{-2}(u) = 0$ since $(\sigma^2\delta_0 + \tilde{\mu}) \star \Upsilon$ is integrable in a neighbourhood of zero. \square

Our first result provides a necessary and sufficient condition in terms of ϕ for s_G to be self-neglecting, see (5.51) for a definition, and $\beta(u) = (u)_n^+ \sqrt{\phi(u)}$ to be flat with respect to G , see (5.54). This together with (5.77) will ensure that \mathcal{F}_{f_n} satisfies (c).

PROPOSITION 5.5.14. *Let $\phi \in \mathcal{B}_{\mathcal{N}}$. Then,*

$$(5.80) \quad \lim_{u \rightarrow \infty} \left(\frac{u^2 \phi'(u)}{\phi(u)} \right)^{\frac{1}{2}} = \lim_{u \rightarrow \infty} \frac{u}{s_G(u)} = \left(\frac{\bar{\Pi}(0^+)}{\mathfrak{r}} \right)^{\frac{1}{2}} \in (0, \infty]$$

and the latter limit is infinite if and only if $\psi \in \mathcal{N}_{\infty}(m)$. As a consequence

$$(5.81) \quad \lim_{u \rightarrow \infty} \frac{\phi(u + \mathfrak{a}s_G(u))}{\phi(u)} = 1$$

uniformly for \mathfrak{a} -compact intervals of $\left(-\left(\frac{\mathfrak{r}}{\bar{\Pi}(0^+)} \right)^{\frac{1}{2}}, \infty \right)$ and

$$(5.82) \quad \lim_{u \rightarrow \infty} \frac{\phi'(u + \mathfrak{a}s_G(u))}{\phi'(u)} = 1$$

uniformly for \mathfrak{a} -compact intervals of $\mathbb{R} \iff \psi \in \mathcal{N}_{\infty}(m)$. Thus, $\psi \in \mathcal{N}_{\infty}(m)$ is a necessary and sufficient condition for \mathcal{F}_{f_n} , $n \geq 0$, to satisfy the condition (c). Consequently, for any $\psi \in \mathcal{N}_{\infty}(m)$, (c) is satisfied.

Proof. From (5.79) note that (5.80) is a restatement of Proposition 4.1.2 which also asserts that the limit of (5.80) is infinite if and only if $\psi \in \mathcal{N}_{\infty}(m)$. Let $\mathfrak{a}_0 > 0$. Then, for any $\mathfrak{a} \in (0, \mathfrak{a}_0)$,

$$\phi(u + \mathfrak{a}s_G(u)) = \phi(u) + \int_u^{u + \mathfrak{a}s_G(u)} \phi'(s) ds \leq \phi(u) + \mathfrak{a}_0 s_G(u) \phi'(u),$$

where the last inequality follows from the fact that ϕ' is non-increasing due to Proposition 4.1.1(2). Therefore, as ϕ is also non-decreasing, we have that

$$1 \leq \frac{\phi(u + \mathfrak{a}s_G(u))}{\phi(u)} \leq 1 + \mathfrak{a}_0 s_G(u) \frac{\phi'(u)}{\phi(u)} = 1 + \frac{\mathfrak{a}_0}{s_G(u)}.$$

Since from Lemma 5.5.13, $\lim_{u \rightarrow \infty} s_G(u) = \infty$, we get (5.81) for compact intervals of \mathbb{R}_+ . Choose arbitrary but fixed $\mathfrak{a}_0 \in \left(0, \left(\frac{\tau}{\bar{\Pi}(0^+)}\right)^{\frac{1}{2}}\right)$. From (5.80) there exists $u_0 = u(\mathfrak{a}_0) > 0$ such that for all $u > u_0$, $0 < 1 - \mathfrak{a}_0 \frac{s_G(u)}{u} < 1$. This, together with $\phi'' \leq 0, \phi''' \geq 0$ on \mathbb{R}^+ , see (4.3), which implies that ϕ is a concave function on \mathbb{R}_+ , gives for all $u > u_0$ and $\mathfrak{a} \in (0, \mathfrak{a}_0)$ that

$$1 \geq \frac{\phi(u - \mathfrak{a}s_G(u))}{\phi(u)} \geq 1 - \mathfrak{a}_0 \frac{\phi'(u)}{\phi(u)} s_G(u) + \frac{\mathfrak{a}_0^2}{2} s_G^2(u) \frac{\phi''(u)}{\phi(u)} = 1 - \frac{\mathfrak{a}_0}{s_G(u)} + \frac{\mathfrak{a}_0^2}{2} \frac{\phi''(u)}{\phi'(u)}.$$

Since $\lim_{u \rightarrow \infty} \mathfrak{a}/s_G(u) = 0$ to show that (5.81) holds it suffices to show that

$$\lim_{u \rightarrow \infty} \frac{\phi''(u)}{\phi'(u)} = \lim_{u \rightarrow \infty} \frac{-\int_0^\infty y^2 e^{-uy} \bar{\Pi}(y) dy}{\sigma^2 + \int_0^\infty y e^{-uy} \bar{\Pi}(y) dy} = 0,$$

where we have used Proposition 4.1.1 for the form of ϕ', ϕ'' . When $\sigma^2 > 0$ it is immediate whereas otherwise the limit follows a standard routing based on the fact that $y^2 \bar{\Pi}(y) \stackrel{0}{=} o(y \bar{\Pi}(y))$. Next, $s_G > 0$ on \mathbb{R}_+ and we observe, for any $n \geq 0$, that

$$\frac{(u + \mathfrak{a}s_G(u))_{(n)}^+}{u_{(n)}^+} \approx \left(1 + \mathfrak{a} \frac{s_G(u)}{u}\right)^n = \left(1 + \mathfrak{a} \sqrt{\frac{\phi(u)}{u^2 \phi'(u)}}\right)^n.$$

Hence, $\psi \in \mathcal{N}_\infty(m)$ is necessary and sufficient for $\beta(u) = (u)_{(n)}^+ \sqrt{\phi(u)}$ to be flat, i.e. that

$$(5.83) \quad \lim_{u \rightarrow \infty} \frac{(u + \mathfrak{a}s_G(u))_{(n)}^+}{u_{(n)}^+} \sqrt{\frac{\phi(u + \mathfrak{a}s_G(u))}{\phi(u)}} = 1$$

uniformly for \mathfrak{a} -compact intervals of \mathbb{R} . We proceed by showing that (5.80) is sufficient for s_G to be self-neglecting which by the definition of a self-neglecting function, see (5.51), the form of s_G , see (5.79), and (5.81) is equivalent to showing (5.82). Let us assume that $\sigma^2 > 0$. Then (5.82) follows from $\phi(u) \approx \sigma^2 u$ and $\phi'(u) \approx \sigma^2$, i.e. Proposition 4.1.1(3). Let $\sigma^2 = 0$. Pick $\mathfrak{a} \in (0, \mathfrak{a}_0)$. Then we again use that ϕ' is non-increasing and the second representation of ϕ' in (4.3) to get

$$\begin{aligned} 1 &\geq \frac{\phi'(u + \mathfrak{a}s_G(u))}{\phi'(u)} = \frac{\int_0^\infty e^{-y(u + \mathfrak{a}s_G(u))} y \bar{\Pi}(y) dy}{\int_0^\infty e^{-yu} y \bar{\Pi}(y) dy} \\ &= \left(1 + \mathfrak{a} \frac{s_G(u)}{u}\right)^{-2} \frac{\int_0^\infty e^{-uy} y \bar{\Pi}\left(\frac{y}{(1 + \mathfrak{a}u^{-1}s_G(u))}\right) dy}{\int_0^\infty e^{-uy} y \bar{\Pi}(y) dy} \\ &\geq \left(1 + \mathfrak{a} \frac{s_G(u)}{u}\right)^{-2} \frac{\int_0^\infty e^{-uy} y \bar{\Pi}(y) dy}{\int_0^\infty e^{-uy} y \bar{\Pi}(y) dy} \geq \left(1 + \mathfrak{a}_0 \frac{s_G(u)}{u}\right)^{-2}, \end{aligned}$$

where we have also used that $\bar{\Pi}$ is non-increasing and $1 + \mathfrak{a}u^{-1}s_G(u) > 1$. Clearly, then the right-hand side converges to 1 if and only if $\psi \in \mathcal{N}_\infty(m)$. Similar computations give

$$\begin{aligned} 1 &\leq \frac{\phi'(u - \mathfrak{a}s_G(u))}{\phi'(u)} = \frac{\int_0^\infty e^{-y(u - \mathfrak{a}s_G(u))} y \bar{\Pi}(y) dy}{\phi'(u)} \\ &\leq \left(1 - \mathfrak{a} \frac{s_G(u)}{u}\right)^{-2} \frac{\int_0^\infty e^{-uy} y \bar{\Pi}(y) dy}{\phi'(u)} \leq \left(1 - \mathfrak{a}_0 \frac{s_G(u)}{u}\right)^{-2} \end{aligned}$$

and we conclude the proof of (5.82). Therefore, from (5.77), \mathcal{F}_{f_n} satisfies the condition (c) holds if and only if $\psi \in \mathcal{N}_\infty(m)$. The proof is thus completed. \square

5.5.4. End of the proof of Theorem 5.4.1. Let $\psi \in \mathcal{N}_{\infty, \infty}(m)$. We apply Proposition 5.4.3 to the functions f_n via their truncation $f_{a,n} > 0$ on $(\ln \bar{a}_n)^{-1}, \infty)$, see (5.68). Lemma 5.5.8 gives the first part of Proposition 5.4.3(a) that is $F_{a,n}(x) = \int_x^\infty f_{a,n}(y) dy$ and even $f_{a,n}$ are of very thin tail. The fact that $\mathcal{F}_{f_n}(u) \approx \mathcal{F}_{f_{a,n}}(u)$ follows from Lemma 5.5.11 which proves (a). Condition (b) for f_n follows from the second claim of Lemma 5.5.1. The first condition of item (c) is Lemma 5.5.12. When $\psi \in \mathcal{N}_{\infty, \infty}(m)$ the properties flatness and asymptotic parabolicity of β, G follow from Proposition 5.5.14 after observing from Lemma 5.5.12 that

$$G(u) = \int_1^u \ln \phi(r) dr = u \ln \phi(u) - \int_1^u \frac{r \phi'(r)}{\phi(r)} dr - \ln \phi(1)$$

and $\beta(u) = C_\psi m(u)_{(n)} \sqrt{\phi(u)}$. Therefore, we conclude that (5.55) holds, that is

$$f_n(y) \approx \frac{1}{\sqrt{2\pi} s_{G^*}(y)} \beta^*(y) e^{-G^*(y)}.$$

To get G^* and β^* we use the definition of the Legendre transform (5.50). Henceforth,

$$G^*(y) = \sup_{u \in \mathbb{R}} \{yu - G(u)\} = y\varphi(e^y) - G(\varphi(e^y)) = \int_m^{e^y} \frac{\varphi(r)}{r} dr - \int_m^{\phi(1)} \frac{\varphi(r)}{r} dr + \ln \phi(1),$$

since the relevant supremum is attained for $u = \varphi(e^y)$ satisfying $y = G'(u) = \ln \phi(u)$ and $\varphi : [m, \infty) \mapsto [0, \infty)$ is the inverse function of ϕ . To compute s_{G^*} we just differentiate twice G^* to get

$$s_{G^*}(y) = \frac{1}{\sqrt{(G^*)''(y)}} = \frac{1}{\sqrt{e^y \varphi'(e^y)}} = \sqrt{e^{-y} \phi'(\varphi(e^y))}.$$

Finally, since $\beta^*(y) = \beta(u)$, $y = G'(u)$, see Proposition 5.4.3, we get that

$$\beta^*(y) = \beta(\varphi(e^y)) = C_\psi m(\varphi(e^y))_n \sqrt{\phi(\varphi(e^y))} = C_\psi m(\varphi(e^y))_n \sqrt{e^y}$$

and clearly

$$\beta^*(y) \approx C_\psi m \sqrt{e^y} \varphi^n(e^y).$$

Therefore, collecting all the results so far and applying (5.55) we get that

$$f_n(y) \approx \frac{1}{\sqrt{2\pi} s_{G^*}(y)} \beta^*(y) e^{-G^*(y)} \approx \frac{\tilde{C}_\psi m}{\sqrt{2\pi}} e^y \varphi^n(e^y) \sqrt{\phi'(\varphi(e^y))} e^{-\int_m^{e^y} \frac{\varphi(r)}{r} dr}.$$

From (5.56) we know that $f_n(y) = e^y v_n(e^y) = e^{-ny} \widehat{\nu}^{(n)}(e^{-y})$ and hence

$$\widehat{\nu}^{(n)}(e^{-y}) \approx \frac{\tilde{C}_\psi m}{\sqrt{2\pi}} e^{(n+1)y} \varphi^n(e^y) \sqrt{\varphi'(e^y)} e^{-\int_m^{e^y} \frac{\varphi(r)}{r} dr}.$$

Thus (5.47) and (5.48) hold. Relation (5.49) then follows from (5.16). This concludes the lengthy proof of Theorem 5.4.1.

5.6. Proof of Theorem 5.0.5

Let us recall, from Proposition 5.1.1, that

$$(5.84) \quad \nu(x) = \frac{1}{x^2} \widehat{\nu}_1\left(\frac{1}{x}\right),$$

with $\widehat{\nu}_1$ the density of $I_{\mathcal{T}_1\psi}$, where $\mathcal{T}_1\psi(u) = u\phi_1(u) = u\phi(u+1) \in \mathcal{N}(\phi(1))$. Since, from (5.48) of Theorem 5.4.1, we know the small asymptotic behaviour of any derivative of $\widehat{\nu}_1$ an application of the Faa di Bruno formula with $\tilde{k}_n = \sum_{j=1}^n j k_j$ and $\bar{k}_n = \sum_{j=1}^n k_j$ gives that

$$\begin{aligned} & \left(\widehat{\nu}_1\left(\frac{1}{x}\right)\right)^{(n)} \\ &= \sum_{\tilde{k}_n=n; \bar{k}_n=k} \frac{n! \widehat{\nu}_1^{(k)}(x^{-1})}{\prod_{j=1}^n k_j! \prod_{j=1}^n (j!)^{k_j}} \prod_{j=1}^n \left(p_1(x^{-1})^{(j)}\right)^{k_j} \\ (5.85) \quad &= (-1)^n \sum_{\tilde{k}_n=n; \bar{k}_n=k} \frac{n! \widehat{\nu}_1^{(k)}(x^{-1}) x^{-n-k}}{\prod_{j=1}^n k_j!} \\ &\approx (-1)^n \frac{\tilde{C}_{\mathcal{T}_1\psi} m_1}{\sqrt{2\pi}} x^{-n+1} \sqrt{\varphi'_1(x)} e^{-\int_{m_1}^x \frac{\varphi_1(y)}{y} dy} \sum_{\tilde{k}_n=n; \bar{k}_n=k} \frac{n! \varphi_1^k(x)}{k_1! k_2! \dots k_n!}, \end{aligned}$$

where as usual $p_1(x) = x$, $\varphi_1(\phi(u+1)) = u$ is the inverse function of $\phi_1(u) = \phi(u+1)$, $m_1 = \phi_1(0) = \phi(1)$ and we have used that under the sum with $\tilde{k}_n = n$, $\bar{k}_n = k$, we have $\prod_{j=1}^n (p_1(x^{-1})^{(j)})^{k_j} = (-1)^n x^{-n-k} \prod_{j=1}^n (j!)^{k_j}$. Since $\lim_{u \rightarrow \infty} \varphi_1(u) = \lim_{u \rightarrow \infty} \phi^{-1}(u+1) = \infty$, the leading term in the last sum of (5.85) is for $k_1 = n$ and $k_2 = \dots = k_n = 0$ or equivalently $\tilde{k}_n = \bar{k}_n = n$. Recall that $\varphi(\phi(u)) = u$ and $m = \phi(0)$. Thus, we get that

$$\begin{aligned} \varphi_1(r) &= \varphi(r) - 1 \approx \varphi(r), \quad \varphi'_1(r) = \varphi'(r), \\ \int_{m_1}^x \frac{\varphi_1(r)}{r} dr &= \int_m^x \frac{\varphi(r)}{r} dr - \int_m^{m_1} \frac{\varphi(r)}{r} dr - \ln x + \ln(m_1). \end{aligned}$$

Plugging these relations in (5.85) we get with $C_\psi = \tilde{C}_{\mathcal{T}_1\psi} e^{\int_m^{m_1} \varphi(y) \frac{dy}{y}}$ that

$$(5.86) \quad \widehat{\nu}_1^{(n)}\left(\frac{1}{x}\right) \approx (-1)^n \frac{C_\psi}{\sqrt{2\pi}} x^{-n+2} \sqrt{\varphi'(x)} \varphi^n(x) e^{-\int_m^x \frac{\varphi(r)}{r} dr}.$$

Then differentiating n times the relation (5.84), we get employing (5.86) that

$$\begin{aligned}
\nu^{(n)}(x) &\approx \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{(k+1)!}{x^{k+2}} \widehat{\mathcal{V}}_1^{(n-k)} \left(\frac{1}{x} \right) \\
(5.87) \quad &\approx (-1)^n \frac{C_\psi}{\sqrt{2\pi}} \frac{1}{x^n} \sqrt{\varphi'(x)} e^{-\int_m^x \frac{\varphi(r)}{r} dr} \sum_{k=0}^n \binom{n}{k} (k+1)! \varphi^{n-k}(x) \\
&\approx (-1)^n \frac{C_\psi}{\sqrt{2\pi}} \frac{1}{x^n} \sqrt{\varphi'(x)} \varphi^n(x) e^{-\int_m^x \frac{\varphi(r)}{r} dr},
\end{aligned}$$

where for the last sum we have used again the fact that $\varphi^n(x)$ dominates all other terms. This proves (5.7) of Theorem 5.0.5. Next, assume that $\sigma^2 > 0$, and note that, since $\phi'(u) \approx \sigma^2$, see Proposition 4.1.1(3), we have from $\phi'(\varphi(y)) \varphi'(y) = 1$, $y > m$, that $\varphi'(y) \approx \sigma^{-2}$ and $\varphi(y) \approx \sigma^{-2}y$. Moreover, from the first expression of (4.2) with $\bar{\mu} = \bar{\Pi}$, since $\phi \in \mathcal{B}_N$, φ solves on (m, ∞)

$$y = \phi(\varphi(y)) = m + \sigma^2 \varphi(y) + \varphi(y) \int_0^\infty e^{-\varphi(y)u} \bar{\Pi}(u) du = m + \sigma^2 \varphi(y) + \varphi(y) H(y),$$

where the last identity sets a notation. Re-expressing

$$y \approx \sigma^2 \varphi(y) = y - \varphi(y) H(y) - m,$$

we obtain that $H(y) = o(1)$ and $\varphi_o(y) := \varphi(y) H(y) = o(y)$. Thus, (5.8) follows from substituting $\varphi(y) = \frac{1}{\sigma^2} (y - \varphi_o(y) - m)$ in the exponent of (5.7) and applying $\varphi'(y) \approx \sigma^{-2}$ and $\varphi(y) \approx \sigma^{-2}y$. Assume in addition that $\bar{\Pi} \in RV_{1+\alpha}(0)$. We refer to [BGT87] for a detailed account on this set of functions and to (1.35) for their explicit definition. Then a classical Tauberian theorem gives that $\bar{\Pi}(y) \stackrel{0}{\sim} \alpha y^{-\alpha} \ell(y)$ with ℓ a slowly varying function and subsequently that

$$H(\phi(y)) = \int_0^\infty e^{-y u} \bar{\Pi}(u) du \approx \alpha \Gamma(1-\alpha) y^{\alpha-1} \ell(y).$$

Thus since $\varphi(y) \approx \sigma^{-2}y$

$$\varphi_o(y) = \varphi(y) H(y) \approx \alpha \Gamma(1-\alpha) \varphi^\alpha(y) \ell(\varphi(y)) \approx \alpha \Gamma(1-\alpha) \sigma^{-2\alpha} y^\alpha \ell(y),$$

which concludes the proof of this case. If now $\bar{\Pi}(0^+) < \infty$, then from the identity

$$\varphi_o(y) = \varphi(y) H(y) = \varphi(y) \int_0^\infty e^{-\varphi(y)u} \bar{\Pi}(u) du = \int_0^\infty (1 - e^{-\varphi(y)u}) \bar{\Pi}(u) du$$

combined with $\varphi(y) \approx \sigma^{-2}y$ we finally get $\lim_{y \rightarrow \infty} \varphi(y) H(y) = \bar{\Pi}(0^+)$. Next, we assume that $\sigma^2 = 0$ and $\psi \in \mathcal{N}_\alpha$ with $\alpha \in (0, 1)$, that is $\psi(y) \approx C_\alpha y^{1+\alpha}$ and thus $\phi(y) \approx C_\alpha y^\alpha$. A standard result from [BGT87] tells us that $\varphi(y) \approx C_\alpha^{-\frac{1}{\alpha}} y^{\frac{1}{\alpha}}$. Finally, from the monotone density theorem thanks to the monotonicity of ϕ' , see Proposition 4.1.1(2), $\phi'(y) \approx \alpha C_\alpha y^{\alpha-1}$ and from the identity $1 = \phi'(\varphi(y)) \varphi'(y)$ we conclude the statement as we have $\varphi'(y) \approx \alpha^{-1} C_\alpha^{-\frac{1}{\alpha}} y^{\frac{1}{\alpha}-1}$. This completes the proof of Theorem 5.0.5.

5.7. End of proof of Theorem 1.1.1

We proceed by describing the dual semigroup. We start by recalling that for any $\psi \in \mathcal{N}$, it is shown in [BY02a, Lemma 2], that, for any $f, g \in \mathbf{B}_b(\mathbb{R}_+)$, writing here $(f, g) = \int_0^\infty f(x)g(x)dx$, one has the following weak duality relationship

$$(K_t f, g) = (f, \widehat{K}_t g)$$

where $(\widehat{K}_t)_{t \geq 0}$ is the (minimal) Feller semigroup of the self-similar process (with 0 an accessible absorbing boundary when $m > 0$) associated via the Lamperti mapping to the dual Lévy process $\hat{\xi} = -\xi$ with Laplace exponent $\widehat{\psi}(u) = -\psi(-u)$, $u < 0$. Note that, actually, in the aforementioned paper, this duality is stated for $m > 0$ but the case $m = 0$ follows from their proof without any modification. In all cases the Lebesgue measure $\xi(dx) = dx$ is an excessive reference measure, i.e. $\xi \widehat{K}_t g \leq \xi g$, $g \geq 0$. Then, from (2.18), we deduce the identities

$$(5.88) \quad (P_t f, g) = e^t (K_{1-e^{-t}} f, d_{e^t} g) = e^t (f, \widehat{K}_{1-e^{-t}} d_{e^t} g) = e^t (f, \widehat{P}_t g)$$

where the last identity is used to set a notation. By following the same computation than above, one easily shows that \widehat{P} is also a Feller semigroup on $(0, \infty)$ and $\xi \widehat{P}_t g = \xi \widehat{K}_{1-e^{-t}} d_{e^t} g \leq e^{-t} \xi g$. The identity (5.88) combined with the invariance property of the measure $\nu(x)dx$ yields

$$e^t (f, \widehat{P}_t \nu) = (P_t f, \nu) = (f, \nu),$$

that is the positive and continuous function ν on $(0, \mathfrak{r})$ is a -1 -invariant function for \widehat{P} . Thus, we may define for any $x \in (0, \mathfrak{r})$, $P_t^* g(x) = \frac{e^t}{\nu(x)} \widehat{P}_t \nu g(x)$, i.e. P^* is a Doob-h transform of \widehat{P} , which defines at least the semigroup of a standard process, see [KW65]. Then, with the obvious notation, we deduce the weak duality relationship, with $\nu(x)dx$ as reference measure,

$$(5.89) \quad (P_t f, g)_\nu = e^t (f, \widehat{P}_t g)_\nu = (f, P_t^* g)_\nu,$$

and $\nu P_t^* g = \int_0^\infty \frac{e^t}{\nu(x)} \widehat{P}_t g \nu(x) \nu(x) dx \leq \nu g$. Finally, from Theorem 5.0.2(2a), we get that $\nu \in \mathbf{C}_0(\mathbb{R}_+)$ for any $N_\mathfrak{r} \geq 1$. Since \widehat{P} is a Feller semigroup on $(0, \infty)$ and, for any $g \in \mathbf{C}_b(\mathbb{R}_+)$, $\nu g \in \mathbf{C}_0(\mathbb{R}_+)$, we deduce that $\widehat{P}_t \nu g \in \mathbf{C}_0(\mathbb{R}_+)$ and thus $P_t^* g \in \mathbf{C}(0, \mathfrak{r})$. Moreover, the inequality, valid for any $x \in (0, \mathfrak{r})$,

$$|P_t^* g(x)| \leq \|g\|_\infty \frac{e^t}{\nu(x)} \widehat{P}_t \nu(x) = \|g\|_\infty$$

implies that P^* is a Feller-Dynkin semigroup on $(0, \mathfrak{r})$ for any $N_\mathfrak{r} > 1$. In all cases, by considering in (5.89) the extension $P \in \mathbf{B}(L^2(\nu))$ gives that P^* admits a contraction semigroup extension in $L^2(\nu)$ which completes the proof of Theorem 1.1.1(4). It simply remains to provide the expression of the generator associated to the Feller-Dynkin dual semigroup when $N_\mathfrak{r} > 2$ to complete the proof of Theorem 1.1.1. Proceeding as in (2.19), we get, with the obvious notation, that for any $f : \mathbb{R}_+ \mapsto \mathbb{R}$ such that $y \mapsto f_e(y) = f(e^y) \in \mathbf{C}^2([-\infty, \infty])$, $\widehat{\mathbf{G}}f(x) = \widehat{\mathbf{G}}_0 f(x) + x f'(x)$, $x > 0$, where $\widehat{\mathbf{G}}_0$ is the generator of the self-similar semigroup \widehat{K} associated via the Lamperti mapping to $\widehat{\psi}(u) = \psi(-u)$ and its expression can be found in [Lam72]. Now, since for $N_\mathfrak{r} > 2$, we have from Theorem 5.0.2 that $\nu \in \mathbf{C}_0^2(\mathbb{R}_+)$, and thus for any $f : \mathbb{R}_+ \mapsto \mathbb{R}$ such that $(f\nu)_e \in \mathbf{C}^2([-\infty, \infty])$, we have $\mathbf{G}^* f(x) = \frac{1}{\nu(x)} \widehat{\mathbf{G}}\nu f(x)$. Some easy algebra completes the proof of item (5). Finally the last item follows first

from the fact that when $\Pi \equiv 0$, then P is the semigroup of a diffusion which is well known to be self-adjoint in $L^2(\nu)$, where $\nu(x)dx$ is the so-called speed measure, see Example 3.0.1. Otherwise, the conclusion is obvious as the semigroup P is associated to a process having only negative jumps whereas its adjoint have only positive jumps, see (5.88) and the fact that the direction of jumps is left invariant by h -transform.

Bernstein-Weierstrass products and Mellin transforms

In this Chapter, we study in depth the Mellin transform of the variables V_ψ and I_ϕ defining, when $\psi \in \mathcal{N}$ and $\phi \in \mathcal{B}_\mathcal{N}$, the invariant measure and the intertwining operator, respectively. We point out that we actually focus on a generalization of these sets of random variables, i.e. for $\phi \in \mathcal{B}$, see (4.1) for definition. Indeed, our approach is comprehensive enough to extend the results without any specific efforts to the general framework. We point out that these variables have been the focus of interest of a number of intense studies in the probabilistic and harmonic analysis literature over the last decade, see Section 6.1 below for a review on these classes of variables. This part also complements to the right-half plane, that is $\mathbb{C}_{(0,\infty)}$, a very interesting research developed by Webster [Web97] on the positive real line for functional equations of the form (6.6) below. This study is completed in [PS18a]. To state the main results of this Chapter, we introduce the following notation. For a function $\phi : \mathbb{C} \rightarrow \mathbb{C}$, we write formally the generalized Weierstrass product

$$(6.1) \quad W_\phi(z) = \frac{e^{-\gamma_\phi z}}{\phi(z)} \prod_{k=1}^{\infty} \frac{\phi(k)}{\phi(k+z)} e^{\frac{\phi'(k)}{\phi(k)} z},$$

where also formally

$$(6.2) \quad \gamma_\phi = \lim_{n \rightarrow \infty} \left(\sum_{k=1}^n \frac{\phi'(k)}{\phi(k)} - \ln \phi(n) \right).$$

Note that if $\phi(z) = z$, then W_ϕ corresponds to the Weierstrass product representation of the Gamma function Γ , valid on $\mathbb{C} \setminus \{0, -1, -2, \dots\}$, and $\gamma_\phi = \gamma$ is the Euler-Mascheroni constant, see e.g. [Leb72], justifying both the terminology and notation. We recall (1.18), i.e. for any $\phi \in \mathcal{B}$, $W_\phi(1) = 1$ and for any $n \geq 1$,

$$W_\phi(n+1) = \prod_{k=1}^n \phi(k).$$

We point out that the definition (1.19) extends to any $\phi \in \mathcal{B}$, i.e.

$$d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\} \in (-\infty, 0].$$

Here are the two main results of this Chapter.

THEOREM 6.0.1. *Let $\phi \in \mathcal{B}$.*

(1) *There exists a positive random variable V_ϕ , whose law is determined by its entire moments*

$$(6.3) \quad \mathbb{E} [V_\phi^n] = W_\phi(n+1), \quad n \in \mathbb{N}.$$

If in addition $\phi \in \mathcal{B}_\mathcal{N}$, then the following identity in law holds

$$(6.4) \quad V_\phi \stackrel{(d)}{=} V_\psi,$$

where $\psi(u) = u\phi(u) \in \mathcal{N}$ and recall that V_ψ was defined in Proposition 2.3.1(1) (note that since $\mathcal{N} \cap \mathcal{B} = \emptyset$ the notation is consistent).

(2) *The product W_ϕ in (6.1) is absolutely convergent on $\mathbb{C}_{(d_\phi, \infty)}$ with $0 \leq \gamma_\phi + \ln \phi(1) \leq \frac{\phi'(1)}{\phi(1)} < \infty$, and, on $\mathbb{C}_{[0, \infty)}$ if $\phi(0) > 0$ and $d_\phi = 0$. Thus, always $W_\phi \in \mathcal{A}_{(d_\phi, \infty)}$ and if addition $\phi(0) > 0$ and $d_\phi = 0$ then $W_\phi \in \mathcal{A}_{[0, \infty)}$.*

(3) *Finally, writing $\mathcal{M}_{V_\phi}(z) = \mathbb{E} [V_\phi^{z-1}]$ for the Mellin transform of V_ϕ , we have on $\mathbb{C}_{(d_\phi, \infty)}$*

$$(6.5) \quad \mathcal{M}_{V_\phi}(z) = W_\phi(z),$$

and, W_ϕ is the unique solution, in the space of Mellin transforms of probability measures, to the following functional equation with initial condition

$$(6.6) \quad \mathcal{M}_\phi(z+1) = \phi(z)\mathcal{M}_\phi(z) \text{ with } \mathcal{M}_\phi(1) = 1,$$

valid on $\mathbb{C}_{(d_\phi, \infty)}$ and on $\mathbb{C}_{[0, \infty)}$ provided that $\phi(0) > 0$ and $d_\phi = 0$.

Since the proof of Theorem 6.0.1(1) is straightforward, we supply it here. The relation (6.3) can be found in [BY01, Proposition 1] whereas the identity in law (6.4) can be deduced, by moments identification, from the expression (2.20), i.e. $\mathcal{M}_{V_\psi}(n+1) = W_\phi(n+1)$, $n \geq 0$, when $\psi \in \mathcal{N}$ and hence $\phi \in \mathcal{B}_\mathcal{N}$. The remaining claims are proved in Section 6.3.

We proceed by stating bounds on W_ϕ for $\phi \in \mathcal{B}$ along with detailed information regarding the asymptotic behaviour of W_ϕ for $\phi \in \mathcal{B}_\mathcal{N}$, which, according to (6.4) and (6.5), boils down to the Mellin transform \mathcal{M}_{V_ψ} with $\psi(u) = u\phi(u) \in \mathcal{N}$. We recall that the asymptotic behaviour of \mathcal{M}_{V_ψ} , on the positive real line has been presented in (5.2) of Theorem 5.0.1. Another purpose of the next proposition is to elucidate the role played by the classes $\mathcal{N}_\Theta = \{\psi \in \mathcal{N}; \Theta_\psi > 0\}$, where $\Theta_\psi = \lim_{|b| \rightarrow \infty} \int_0^\infty \ln \left(\frac{|\phi(by+ib)|}{\phi(by)} \right) dy$, and \mathcal{N}_∞ introduced in the Table 1 and 2 in Chapter 1.

More detailed statements about these asymptotic estimates together with further examples are provided in the sections 6.5–6.8 below, which also contain the proof of the following assertion.

THEOREM 6.0.2. (1) *Let $\phi \in \mathcal{B}$. Then, for any $a > 0$ and $b \in \mathbb{R}$,*

$$(6.7) \quad |W_\phi(a+ib)| \asymp \frac{\sqrt{\phi(a)}W_\phi(a)}{\sqrt{|\phi(a+ib)|}} e^{-|b|\Theta_\phi(a,|b|)}$$

where $\Theta_\phi(a, |b|) = \int_{\frac{a}{|b|}}^\infty \ln \left(\frac{|\phi(|b|y+i|b|)|}{\phi(|b|y)} \right) dy$ and for any $b \in \mathbb{R}$,

$$0 \leq \Theta_\phi(a, |b|) \leq \frac{\pi}{2}.$$

(2) Let $\psi \in \mathcal{N}$ and hence $\phi(u) = \frac{\psi(u)}{u} \in \mathcal{B}_{\mathcal{N}}$.

(a) For any $p = 0, 1, \dots$, any real number $u \leq \max(N_{\tau} - 1, 0)$, recalling from (1.8) that $N_{\tau} = \left\lceil \frac{\overline{\Pi}(0^+)}{\tau} \right\rceil - 1$ with $N_{\tau} = \infty$ if $\overline{\Pi}(0^+) = \infty$, and $a > d_{\phi}$, we have that

$$(6.8) \quad \lim_{|b| \rightarrow \infty} |b|^u |W_{\phi}(a + ib)| = 0.$$

In particular, $\psi \in \mathcal{N}_{\infty}$ if and only if (6.8) holds for all $u \geq 0$ and $a > d_{\phi}$.

(b) We have $\underline{\Theta}_{\phi} \in [0, \frac{\pi}{2}]$ with the following specific values.

(i) if $\psi \in \mathcal{N}_P$ then $\underline{\Theta}_{\phi} = \frac{\pi}{2}$ and $\psi \in \mathcal{N}_{\underline{\Theta}}$,

(ii) if $\psi \in \mathcal{N}_{\alpha}$, $\alpha \in (0, 1)$, then $\underline{\Theta}_{\phi} = \frac{\pi}{2}\alpha$ and $\psi \in \mathcal{N}_{\underline{\Theta}}$,

(iii) if $\overline{\Pi}(y) = y^{-1} |\ln y|^2 \mathbb{I}_{[0, 1/2]}(y)$ then $\Theta_{\phi}(|b|) \approx \ln^3 b$ and $\psi \in \mathcal{N}_{\infty} \setminus \mathcal{N}_{\underline{\Theta}}$.

Theorem 6.0.2(1) is proved in section 6.5.1, the proof of Theorem 6.0.2(2(b)i), 6.0.2(2(b)ii) and Theorem 6.0.2(2(b)iii)) is given in section 6.8.1, 6.8.2 and 6.8.3 respectively. Finally, the proof of Theorem 6.0.2(2a) is discussed in the following remark.

REMARK 6.0.3. Identities (6.4) and (6.5) yield that $W_{\phi} = \mathcal{M}_{V_{\phi}}$ and the necessary and sufficient conditions for the polynomial decay of W_{ϕ} along imaginary lines stated in (6.8) follows readily from the estimates (5.3) on $\mathcal{M}_{V_{\phi}}$. Recall that (5.3) was obtained by means of the Riemann-Lebesgue Lemma combined with the general theory of the class of self-decomposable variables.

REMARK 6.0.4. We emphasize that our developments on Bernstein-Weierstrass products offer a unified and comprehensive treatment of some well known and substantial special functions. In order to illustrate this fact, without aiming at being exhaustive, we simply mention two examples. First, note that, with $\sigma^2 = m = 0$ and $\mu(dy) = \frac{\delta - \ln q(y)}{1-q}$, $0 < q < 1$ in (4.2), then $\phi_q(u) = \frac{1-q^u}{1-q} \in \mathcal{B} \setminus \mathcal{B}_{\mathcal{N}}$, and $W_{\phi_q} = \Gamma_q$ is the q -gamma function, which has been studied intermittently for over a century and originally introduced by Thomae in 1869. Moreover for $\phi(u) = \phi_{\alpha, 1 - \frac{1}{\alpha}}^R\left(\frac{u}{\alpha}\right) = \frac{\Gamma(u+\alpha)}{\Gamma(u)} \in \mathcal{B}_{\mathcal{N}}$, see (1.15) for the definition of $\phi_{\alpha, 1 - \frac{1}{\alpha}}^R$, then W_{ϕ} boils down to the ratio of the Barnes gamma function. Note also that our work reveals some interesting connections between these special functions and the spectral theory of some non-self-adjoint contraction semigroups.

REMARK 6.0.5. The upper bound for $\Theta_{\phi}(a, |b|)$ in (6.7), i.e. $\frac{\pi}{2}$, is attained precisely when $\phi(u) = u$ and $b = \infty$. This is the case when $\psi(u) = u^2$ and V_{ϕ} is an exponential random variable, whose Mellin transform is simply the gamma function. In this sense the gamma function envelops from below the rate of decay of all Mellin transforms $\mathcal{M}_{V_{\phi}}(z)$ along complex lines.

REMARK 6.0.6. We mention that to obtain the bounds (6.7) valid along imaginary lines, we resort to the Weierstrass representation of W_{ϕ} and elaborate upon an

approach which has been used to derive estimates of the gamma function and can be found for instance in [Olv74, Chap. 8].

The proofs of Theorem 6.0.1 and Theorem 6.0.2 are given after a short review on exponential functionals of subordinators which appear in definition (1.13) of the intertwining operator.

6.1. Exponential functional of subordinators

We present here some basic facts on the so-called exponential functional of Lévy processes, a random variable which has been intensively studied during the last two decades, something which seems to be attributed to its close connection to numerous mathematical fields, such as probability theory, harmonic and real analysis and mathematical physics, to name but a few. We refer the interested reader to the survey paper of Bertoin and Yor [BY05] and the paper by Berg and Durán [BD04]. Its original study, in the case of subordinators, traces back to the work of Urbanik [Urb95] and the characterization of its distribution through its Mellin transform for the entire class of Lévy processes has been achieved by the authors. This has been announced without proofs in [PS13b] and we refer also to [PPS12] and [PS12a] for previous developments in this direction. We refer to [PS18a] for the detailed proofs of this comprehensive result and for the most recent developments regarding the exponential functionals. Let us now proceed with the definition of this random variable when the Lévy process is a subordinator, that is, a non-decreasing Lévy process, and we mention that the case when it is a spectrally negative was already introduced in Chapter 5(5.1). For any $\phi \in \mathcal{B}$, that is from (4.2),

$$(6.9) \quad \phi(u) = m + \sigma^2 u + \int_0^\infty (1 - e^{-uy}) \mu(dy), \quad u \geq 0,$$

where $m, z, \sigma^2 \geq 0$ and μ is a Lévy measure such that $\int_0^\infty (1 \wedge y) \mu(dy) < \infty$, we write

$$I_\phi = \int_0^\infty e^{-\eta_t} dt$$

with $(\eta_t)_{t \geq 0}$ a (possibly) killed subordinator with Laplace exponent ϕ (see Chapter 4 for exhaustive information on Bernstein functions). Since from the strong law of large numbers, $\lim_{t \rightarrow \infty} \frac{\eta_t}{t} = \mathbb{E}[\eta_1] \in (0, \infty]$ a.s., we have $I_\phi < \infty$ a.s., see e.g. [BY05, Proposition 1]. We now collect some basic known properties that are useful in our context.

PROPOSITION 6.1.1. *Let $\phi \in \mathcal{B}$.*

- (1) *The law of I_ϕ is absolutely continuous with density denoted by ι and with support $[0, \sigma^{-2}]$ if $\sigma^2 > 0$ and $[0, \infty)$ otherwise.*
- (2) *Its Mellin transform \mathcal{M}_{I_ϕ} satisfies the following functional equation, with initial condition $\mathcal{M}_{I_\phi}(1) = 1$,*

$$(6.10) \quad \overline{\mathcal{M}}_\phi(z+1) = \frac{z}{\phi(z)} \overline{\mathcal{M}}_\phi(z), \quad \overline{\mathcal{M}}_\phi(1) = 1,$$

valid on $\mathbb{C}_{(0, \infty)}$ and $\mathcal{M}_{I_\phi} \in \mathcal{A}_{(0, \infty)}$.

(3) Its law is moment determinate, and, for any $n \geq 0$, we have that

$$(6.11) \quad \mathcal{M}_{I_\phi}(n+1) = \frac{n!}{W_\phi(n+1)}.$$

(4) Finally, $\mathbb{E}[e^{uI_\phi}] < \infty$, for any $u < \phi(\infty)$, where $\phi(\infty)$ is finite if and only if $\sigma^2 = 0$ and $\bar{\mu}(0^+) < \infty$, see (6.9) for the definition of the measure μ .

Proof. The absolute continuity of the law of I_ϕ is shown in the proof of [PS12a, Theorem 2.4] whereas its support is derived in [HR12, Lemma 2.1]. The characterization of \mathcal{M}_{I_ϕ} as a solution to the functional equation (6.10) is given in [MZ06, Theorem 1]. $\mathbb{E}[I_\phi^u] = \mathcal{M}_{I_\phi}(u+1) < \infty$ for $u \in (-1, \infty)$ follows from (6.10). The expression of the integer moments of I_ϕ (6.11) can be found in [BY05]. Finally, the claim $\mathbb{E}[e^{uI_\phi}] < \infty$ for any $u < \phi(\infty)$ is immediate from (6.11), $W_\phi(n+1) = \prod_{k=1}^n \phi(k)$, the monotonicity of ϕ , see Proposition 4.1.1(2), and the Taylor expansion of the exponential function. The very last assertion follows from (6.9) with $u = \infty$. \square

We end this part with the following statement which provides additional insights on the variable I_ϕ and its proof is postponed to section 6.3.

PROPOSITION 6.1.2. *Let $\phi \in \mathcal{B}$. Then $z \mapsto \mathcal{M}_{I_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)} \in \mathcal{A}_{(0, \infty)}$ is the unique solution in the space of Mellin transforms of probability measures to the functional equation (6.10) valid on the strip $\mathbb{C}_{(0, \infty)}$. If, in addition, $\phi(0) = m = 0$, ϕ does not vanish on $i\mathbb{R} \setminus \{0\}$ and $\phi'(0^+) < \infty$, then $z \mapsto \mathcal{M}_{I_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)} \in \mathcal{A}_{[0, \infty)}$ is the unique solution on $\mathbb{C}_{[0, \infty)}$.*

6.2. The functional equations (6.6) and (6.10) on \mathbb{R}_+

We start the proof of the main results of this Chapter with the following result which can be seen as a generalization of the Bohr-Mollerup-Artin classical characterization of the Gamma function. It is essentially due to Webster [Web97].

LEMMA 6.2.1. *Let $\phi \in \mathcal{B}$. Then, we have, for all $u > 0$,*

$$(6.12) \quad \mathcal{M}_{V_\phi}(u) = W_\phi(u),$$

with $0 \leq \gamma_\phi + \ln \phi(1) \leq \frac{\phi'(1)}{\phi(1)} < \infty$. Moreover, \mathcal{M}_{V_ϕ} is the unique positive log-convex solution to (6.6) on \mathbb{R}_+ . Similarly, we have, for all $u > 0$,

$$(6.13) \quad \mathcal{M}_{I_\phi}(u) = \frac{\Gamma(u)}{W_\phi(u)},$$

and, \mathcal{M}_{I_ϕ} is the unique positive log-convex solution to (6.10) on \mathbb{R}_+ .

Proof. From the items (2) and (6) of Proposition 4.1.1, for any $\phi \in \mathcal{B}$, the mapping $u \mapsto \phi(u)$ is positive and strictly log-concave on $(0, \infty)$ and the limit in (4.6) holds. This means that the multiplier ϕ in (6.6) satisfies the conditions of [Web97, Theorem 7.1], which states that the limit defining γ_ϕ in (6.2) exists and

W_ϕ in (6.1) is the unique positive log-convex solution of (6.6) on \mathbb{R}_+ . We complete the proof of (6.12) by observing that $\mathcal{M}_{V_\phi}(1) = 1$, for any $n \in \mathbb{N} \setminus \{0\}$

$$(6.14) \quad \mathcal{M}_{V_\phi}(n+1) = W_\phi(n+1) = \prod_{k=1}^n \phi(k) = \phi(n)\mathcal{M}_{V_\phi}(n),$$

and, by recalling that the mapping $u \mapsto \mathcal{M}_{V_\phi}(u)$, as the moments of order greater than -1 of a positive random variable, is a positive and log-convex function on \mathbb{R}_+ . Thus, necessarily $\mathcal{M}_{V_\phi}(u) = W_\phi(u)$ on \mathbb{R}_+ . The bounds for γ_ϕ are provided in the comments preceding [**Web97**, Theorem 7.1]. Next, note, from (4.6), that $\lim_{x \rightarrow \infty} \frac{(x+u)\phi(x)}{x\phi(x+u)} = 1$, for all $u > 0$, and from Proposition 4.1.1(5) that $u \mapsto \frac{u}{\phi(u)}$ is log-concave and positive on \mathbb{R}_+ . Thus, we can apply again [**Web97**, Theorem 7.1] to the functional equation with initial condition (6.10) to conclude that it has, on \mathbb{R}_+ , a unique positive log-convex solution say $\overline{\mathcal{M}}_\phi$ given by the following infinite product representation

$$(6.15) \quad \overline{\mathcal{M}}_\phi(u) = e^{-(\gamma-\gamma_\phi)u} \frac{\phi(u)}{u} \prod_{k=1}^{\infty} \frac{\phi(u+k)k}{\phi(k)(u+k)} e^{\left(\frac{1}{k} - \frac{\phi'(k)}{\phi(k)}\right)u}, \text{ for any } u > 0,$$

where with $g(u) = u/\phi(u)$, $\gamma_g = \gamma - \gamma_\phi$ in [**Web97**, Theorem 7.1] comes from the computation

$$\begin{aligned} \gamma_g &= \lim_{n \rightarrow \infty} \left(\sum_{k=1}^n \frac{g'(k)}{g(k)} - \ln g(n) \right) \\ &= \lim_{n \rightarrow \infty} \left(\sum_{k=1}^n \frac{1}{k} - \ln n - \left(\sum_{k=1}^n \frac{\phi'(k)}{\phi(k)} - \ln \phi(n) \right) \right) \\ &= \gamma - \gamma_\phi. \end{aligned}$$

Therefore, one may rewrite (6.15) using (6.1), for any $u > 0$, as follows

$$(6.16) \quad \overline{\mathcal{M}}_\phi(u) = \frac{\Gamma(u)}{W_\phi(u)}.$$

By proceeding as in the previous case, we conclude the proof of (6.13) after recalling that $\mathcal{M}_{I_\phi}(n+1) = \frac{n}{\phi(n)}\mathcal{M}_{I_\phi}(n)$ for $n \in \mathbb{N} \setminus \{0\}$. \square

As the Mellin transform of some positive variables which have all moments of order greater than -1 finite, see (6.12) and (6.13), it is well-known that both \mathcal{M}_{V_ϕ} and \mathcal{M}_{I_ϕ} admit an analytical extension to (at least) $\mathbb{C}_{(0,\infty)}$. Thus, the uniqueness property stated in Theorem 6.0.1 (resp. Proposition 6.1.2), that is $W_\phi = \mathcal{M}_{V_\phi}$ (resp. $\frac{\Gamma}{W_\phi} = \mathcal{M}_{I_\phi}$) is the unique solution in the space of Mellin transforms of probability measures of the functional equation (6.6) (resp. (6.10)), is an immediate consequence of an analytical extension argument combined with the uniqueness property established on the positive real line in (6.12) and (6.13) of Lemma (6.2.1). The next two sections contain the justification that the Bernstein-Weierstrass representation of \mathcal{M}_{V_ϕ} and \mathcal{M}_{I_ϕ} is also valid on $\mathbb{C}_{(0,\infty)}$.

6.3. Proof of Theorem 6.0.1

We start by recalling that the proof of Theorem 6.0.1(1) was given right after its statement and the proof of Theorem 6.0.1(3) follows, directly from the discussion above, once the analyticity property stated in Theorem 6.0.1(2), that is $W_\phi \in \mathcal{A}_{(0,\infty)}$, is justified. To complete the proof of Theorem 6.0.1, it remains therefore to extend analytically the function W_ϕ to the positive half-plane. Let now $\phi \in \mathcal{B}$, and, for the reader's convenience, we recall, from Section 4.3, for any $a > 0$ and $b \in \mathbb{R}$, the following notation

$$\Delta_b^{\Re} \phi(a) = \Re(\Delta_b \phi(a)) = \Re(\phi(a+ib) - \phi(a)) = \int_0^\infty (1 - \cos(by)) e^{-ay} \mu(dy),$$

$$\Delta_b^{\Im} \phi(a) = \Im(\Delta_b \phi(a)) = \Im(\phi(a+ib) - \phi(a)) = \sigma^2 b + \int_0^\infty \sin(by) e^{-ay} \mu(dy).$$

For $z \in \mathbb{C}_{(0,\infty)}$, i.e. $z = a + ib$, with $a > 0$ and $b \in \mathbb{R}$ fixed, we write formally

$$(6.17) \quad Z_\phi(z) = \prod_{k=1}^{\infty} \frac{\phi(k+a)}{\phi(k+z)} e^{\frac{\phi'(k)}{\phi(k)} ib},$$

and, for any $k \geq 1$,

$$A_k = \frac{\phi(z_k)}{\phi(a_k)} e^{-\frac{\phi'(k)}{\phi(k)} ib},$$

where, for any number x and integer l , we set $x_l = x + l$. Then, we get that

$$\frac{\phi(k)}{\phi(z_k)} e^{\frac{\phi'(k)}{\phi(k)} z} = \frac{\phi(k)}{\phi(a_k)} e^{\frac{\phi'(k)}{\phi(k)} a} A_k^{-1}.$$

Thus, from (6.1) and W_ϕ solution on \mathbb{R}_+ to (6.6), see Lemma 6.2.1, then on $\mathbb{C}_{(0,\infty)}$

$$(6.18) \quad \begin{aligned} W_\phi(z+1) &= \phi(z)W_\phi(z) = \phi(a)W_\phi(a) e^{-i\gamma_\phi b} \prod_{k=1}^{\infty} A_k^{-1} \\ &= \phi(a)W_\phi(a) e^{-i\gamma_\phi b} Z_\phi(z), \end{aligned}$$

provided the last infinite product is absolutely convergent on $\mathbb{C}_{(0,\infty)}$ which by a standard result in complex analysis will follow if $\sum_{k=1}^{\infty} |A_k^{-1} - 1| < \infty$. To prove the latter it suffices to show that $\sum_{k=1}^{\infty} |A_k - 1| < \infty$. To achieve this, first, observe that

$$(6.19) \quad A_k - 1 = \left(1 + \frac{\Delta_b^{\Re} \phi(a_k) + i \Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} \right) e^{-i \frac{\phi'(k)}{\phi(k)} b} - 1.$$

Since from (4.26), $0 \leq \Delta_b^{\Re} \phi(a_k) \leq \frac{b^2}{2} |\phi''(a_k)|$ and $a \mapsto \left| \frac{\phi''(a)}{\phi(a)} \right|$ is plainly non-increasing on \mathbb{R}_+ because ϕ' is completely monotone and ϕ is non-decreasing, see Proposition 4.1.1(2), we get with the help of (4.5) and (4.28) for the last inequality

below that

$$\begin{aligned}
\sum_{k=1}^{\infty} \left| \frac{\Delta_b^{\Re} \phi(a_k)}{\phi(a_k)} \right| &\leq \frac{b^2}{2} \sum_{k=1}^{\infty} \left| \frac{\phi''(a_k)}{\phi(a_k)} \right| \\
&\leq \frac{b^2}{2} \left(\frac{|\phi''(1+a)|}{\phi(1+a)} + \int_{a+1}^{\infty} \left| \frac{\phi''(u)}{\phi(u)} \right| du \right) \\
(6.20) \quad &\leq \frac{2b^2}{(1+a)^2} + \frac{2\sqrt{10}b^2}{1+a} < \infty.
\end{aligned}$$

Thus, it remains to show that $\sum_{k=1}^{\infty} |\bar{A}_k| < \infty$, where

$$\begin{aligned}
\bar{A}_k &= \left(1 + \frac{i\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} \right) e^{-\frac{\phi'(k)}{\phi(k)} ib} - 1 \\
(6.21) \quad &= (C_k(b) - 1) \left(1 + i\frac{\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} \right) + i \left(\frac{\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} - S_k(b) \right) + \frac{\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} S_k(b),
\end{aligned}$$

with $C_k(b) = \cos\left(\frac{\phi'(k)}{\phi(k)}b\right)$ and $S_k(b) = \sin\left(\frac{\phi'(k)}{\phi(k)}b\right)$. We have, using the estimates (4.26) in the second inequality and (4.4) in the second and third inequality, that

$$\begin{aligned}
\sum_{k=1}^{\infty} \left| (1 - C_k(b)) \left(1 + \frac{i\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} \right) \right| &\leq \sum_{k=1}^{\infty} b^2 \left(\frac{\phi'(k)}{\phi(k)} \right)^2 \left(1 + \left| \frac{\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} \right| \right) \\
&\leq \sum_{k=1}^{\infty} \frac{b^2}{k^2} \left(1 + \left| b \frac{\phi'(a_k)}{\phi(a_k)} \right| \right) \\
(6.22) \quad &\leq \sum_{k=1}^{\infty} \frac{b^2}{k^2} \left(1 + \left| \frac{b}{a_k} \right| \right) < \infty.
\end{aligned}$$

Let us now consider the second term on the right-hand side of (6.21). Since from (4.4) we have that $0 \leq \lim_{a \rightarrow \infty} \frac{\phi'(a)}{\phi(a)} \leq \lim_{a \rightarrow \infty} \frac{1}{a} = 0$, then for any $k \geq k_0(b)$ large enough,

$$S_k(b) = b \frac{\phi'(k)}{\phi(k)} + O\left(\left(\frac{\phi'(k)}{\phi(k)}\right)^2\right).$$

Therefore, as from (4.4) $\sum_{k=1}^{\infty} \left(\frac{\phi'(k)}{\phi(k)}\right)^2 \leq \sum_{k=1}^{\infty} \frac{1}{k^2} < \infty$, to prove the summability of the second term in (6.21) it remains to prove the finiteness of the sum

$$\begin{aligned}
\sum_{k=1}^{\infty} \left| \frac{\Delta_b^{\Im} \phi(a_k)}{\phi(a_k)} - b \frac{\phi'(k)}{\phi(k)} \right| &\leq \sum_{k=1}^{\infty} |b| \left(\frac{\phi'(k)}{\phi(k)} - \frac{\phi'(a_k)}{\phi(a_k)} \right) \\
(6.23) \quad &+ \sum_{k=1}^{\infty} \frac{|\Delta_b^{\Im} \phi(a_k) - b\phi'(a_k)|}{\phi(a_k)}.
\end{aligned}$$

Since from Proposition 4.1.1(2), $u \mapsto \frac{\phi'(u)}{\phi(u)}$ is non-increasing on \mathbb{R}_+ , we get that

$$(6.24) \quad \begin{aligned} \sum_{k=1}^{\infty} \left(\frac{\phi'(k)}{\phi(k)} - \frac{\phi'(a_k)}{\phi(a_k)} \right) &\leq \sum_{k=1}^{\infty} \left(\frac{\phi'(k)}{\phi(k)} - \frac{\phi'(k + \lceil a \rceil + 1)}{\phi(k + \lceil a \rceil + 1)} \right) \\ &= \sum_{k=1}^{\lceil a \rceil + 1} \frac{\phi'(k)}{\phi(k)} \leq 1 + \lceil a \rceil < \infty. \end{aligned}$$

On the other hand, the form of $\Delta_b^{\mathfrak{S}}$ and the second expression for ϕ' in (4.3) yield that

$$\begin{aligned} |\Delta_b^{\mathfrak{S}} \phi(a_k) - b\phi'(a_k)| &\leq \int_0^{\infty} |\sin(by) - by| e^{-a_k y} \mu(dy) \\ &\leq \left(|b|^3 \int_0^{\frac{1}{|b|}} y^3 e^{-a_k y} \mu(dy) + \int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-a_k y} \mu(dy) \right) \\ &\leq \left(b^2 |\phi''(a_k)| + \int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-a_k y} \mu(dy) \right), \end{aligned}$$

where for the last inequality we have used the upper bound

$$|b| \int_0^{\frac{1}{|b|}} y^3 e^{-a_k y} \mu(dy) \leq \int_0^{\frac{1}{|b|}} y^2 e^{-a_k y} \mu(dy) \leq \int_0^{\infty} y^2 e^{-a_k y} \mu(dy) = |\phi''(a_k)|.$$

Thus, as $a \mapsto \frac{|\phi''(a)|}{\phi(a)}$ is non-increasing on \mathbb{R}_+ , with the help of (6.20) and (4.28) for the third inequality, we get, writing $\bar{C}_{ab} = b^2 \left(\frac{2}{(a+1)^2} + \frac{2\sqrt{10}}{a+1} \right)$ and recalling that $a_k = a + k$,

$$(6.25) \quad \begin{aligned} \sum_{k=1}^{\infty} \frac{|\Delta_b^{\mathfrak{S}} \phi(a_k) - b\phi'(a_k)|}{\phi(a_k)} &\leq b^2 \sum_{k=1}^{\infty} \frac{|\phi''(a_k)|}{\phi(a_k)} + \sum_{k=1}^{\infty} \frac{\int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-a_k y} \mu(dy)}{\phi(a_k)} \\ &\leq \bar{C}_{ab} + \sum_{k=1}^{\infty} \frac{\int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-a_k y} \mu(dy)}{\phi(1+a)} \\ &\leq \bar{C}_{ab} + \frac{\int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-ay} \frac{e^{-y}}{1-e^{-y}} \mu(dy)}{\phi(1+a)} \\ &\leq \bar{C}_{ab} + \frac{\int_{\frac{1}{|b|}}^{\infty} (1 + |b|y) e^{-ay} e^{-y} \mu(dy)}{\phi(1+a) \left(1 - e^{-\frac{1}{|b|}}\right)} \\ &\leq \bar{C}_{ab} + \frac{|b| e^{-\frac{|b|-1}{|b|}} \int_{\frac{1}{|b|}}^{\infty} e^{-ay} \mu(dy)}{\phi(1+a) \left(1 - e^{-\frac{1}{|b|}}\right)} < \infty. \end{aligned}$$

For the third term in (6.21), we have, using the second relation in (4.26) and $|\sin(y)| \leq |y|$, $y \in \mathbb{R}$, in the first inequality and (4.4) in the second, that

$$(6.26) \quad \sum_{k=1}^{\infty} \left| \frac{\Delta_b^{\mathfrak{S}} \phi(a_k)}{\phi(a_k)} S_k(b) \right| \leq b^2 \sum_{k=1}^{\infty} \frac{|\phi'(a_k)|}{\phi(a_k)} \frac{|\phi'(k)|}{\phi(k)} \leq b^2 \sum_{k=1}^{\infty} \frac{1}{ka_k} < \infty.$$

Then, we put (6.24) and (6.25) in (6.23) which together with (6.22) and (6.26) is used in (6.21) to confirm that $\sum_{k=1}^{\infty} |\bar{A}_k| < \infty$. The latter triggers with the help of (6.20) and (6.19) $\sum_{k=1}^{\infty} |A_k - 1| < \infty$ and hence $\sum_{k=1}^{\infty} |A_k^{-1} - 1| < \infty$. An application of Montel's Theorem in (6.18) yields that the right-hand side of (6.1) defines a holomorphic function on $\mathbb{C}_{(0,\infty)}$ and since $W_\phi(u+1) = \mathcal{M}_{V_\phi}(u+1)$ for $u > -1$, see (6.12), we have proved Theorem 6.0.1 for $z \in \mathbb{C}_{(0,\infty)}$. When $\phi(0) = m > 0$, we check that all computations above extend to the imaginary line $i\mathbb{R}$ as (6.20), (6.22), (6.24), (6.25) and (6.26) do not explode for $a = 0$, and $W_\phi(0) = m^{-1} \in (0, \infty)$ thanks to (6.6). Next, if $d_\phi < 0$ then plainly $\phi > 0$ on $(d_\phi, 0)$. Clearly, then the infinite product in (6.1) extends holomorphically to $\mathbb{C}_{(d_\phi, \infty)}$. Indeed, only at most $\lceil -d_\phi \rceil + 1$ of its terms are with argument whose real part is in $(d_\phi, 0)$ and their product defines an analytic function. The rest of the product is absolutely convergent and defines a holomorphic function as already proved above. Since W_ϕ extends analytically to $\mathbb{C}_{(d_\phi, \infty)}$ the proof of Theorem 6.0.1(2) and hence of Theorem 6.0.1 is therefore completed.

6.4. Proof of Proposition 6.1.2

Proceeding as in the proof of Theorem 6.0.1, to show that $\bar{\mathcal{M}}_\phi$ in (6.16) extends to a holomorphic function in $\mathbb{C}_{(0,\infty)}$, we set for any $z = a + ib$, $a > 0$,

$$\begin{aligned} \bar{\mathcal{M}}_\phi(z) &= \frac{\Gamma(z) \phi(z)}{\Gamma(a) \phi(a)} \mathcal{M}_{I_\phi}(a) e^{-i\gamma_\phi b} \prod_{k=1}^{\infty} \frac{\phi(z+k)}{\phi(a+k)} e^{-i \frac{\phi'(k)}{\phi(k)} b} \\ &= \frac{\Gamma(z) \phi(z)}{\Gamma(a) \phi(a)} \mathcal{M}_{I_\phi}(a) e^{-i\gamma_\phi b} \prod_{k=1}^{\infty} A_k. \end{aligned}$$

This completes the proof of Proposition 6.1.2 for $z \in \mathbb{C}_{(0,\infty)}$ since $\bar{\mathcal{M}}_\phi = \mathcal{M}_{I_\phi}$ on \mathbb{R}_+ from (6.13) and $\prod_{k=1}^{\infty} A_k < \infty$ from the proof of Theorem 6.0.1, where the uniqueness argument was also discussed. If $\phi(0) = 0$, ϕ does not vanish on $i\mathbb{R} \setminus \{0\}$ and $\phi'(0^+) = \sigma^2 + \int_0^\infty \bar{\mu}(y) dy < \infty$ then since $\phi \not\equiv 0$ we have that $\phi'(0^+) > 0$ and $\frac{z}{\phi(z)}$ extends to $i\mathbb{R}$. Hence, through (6.10) and (6.15) the left- and right-hand side of $\mathcal{M}_{I_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)}$ extend to $\mathbb{C}_{[0,\infty)}$.

We mention that, as from Theorem 6.0.1, $\mathcal{M}_{V_\phi} = W_\phi$, the results regarding estimates of the Mellin transform \mathcal{M}_{V_ϕ} , throughout the rest of this chapter will be stated in terms of W_ϕ .

6.5. Proof of Theorem 6.0.2(1): Bounds for W_ϕ

In this part, we derive general bounds for the absolute value of W_ϕ along imaginary lines which provide some bounds for the Mellin transforms \mathcal{M}_{V_ϕ} and \mathcal{M}_{I_ϕ} . This will serve us to obtain exact asymptotic estimates along imaginary lines for $|W_\phi|$ and $|\mathcal{M}_{I_\phi}|$, that we shall exploit to establish smoothness properties of the invariant density as well as existence and smoothness properties of the sequence of co-eigenfunctions. To obtain our results, we resort to the Bernstein-Weierstrass product representation of the Mellin transform and extend an approach which has been used to derive estimates of the gamma function, which can be found for instance in [Olv74, Chap. 8]. Before stating the next result, we revisit and introduce

some notation. Let us recall that, for any function $\phi \in \mathcal{B}$ and numbers a and $b > 0$, we write

$$(6.27) \quad \Theta_\phi(a, b) = \int_{\frac{a}{b}}^{\infty} \ln \left(\frac{|\phi(bu + ib)|}{\phi(bu)} \right) du.$$

Furthermore, we set formally, for $a, b \in \mathbb{R}$,

$$E_\phi(a, b) = e^{-\frac{1}{8} \int_a^\infty \frac{|\phi''(u+ib)|}{|\phi(u+ib)|} + \frac{|\phi'(u+ib)|^2}{|\phi(u+ib)|^2} du - \frac{1}{8} \frac{\phi'(a)}{\phi(a)}}$$

and

$$(6.28) \quad Z_\phi^l(a, b) = \sqrt{\frac{|\phi(a + ib + l)|}{\phi(a + l)}} \prod_{k=0}^{l-1} \frac{|\phi(k + a + ib)|}{\phi(k + a)},$$

where we use the convention $\prod_{k=0}^{-1} = 1$. We are ready to state the following result.

PROPOSITION 6.5.1. *Let $\phi \in \mathcal{B}$. For any $b \in \mathbb{R} \setminus \{0\}$ and $a > 0$, we have*

$$(1) \quad 0 \leq \Theta_\phi(a, |b|) \leq \frac{\pi}{2},$$

$$(2) \quad e^{-\frac{19}{8a}} \leq E_\phi(a, |b|) \leq e^{\frac{19}{8a}}.$$

(3) *Moreover, if $a > d_\phi$, $l \in \mathbb{N}$ such that $a + l > 0$,*

$$(6.29) \quad E_\phi(a + l, |b|) \frac{W_\phi(a)}{Z_\phi^l(a, |b|)} \leq \frac{|W_\phi(a + i|b|)|}{e^{-|b|\Theta_\phi(a+l, |b|)}} \leq \frac{W_\phi(a)}{Z_\phi^l(a, |b|)E_\phi(a + l, |b|)}.$$

If $\phi(0) = m > 0$ and $d_\phi = 0$ then (6.29) extends to $a = 0$.

(4) *With $a > 0$ and $l \in \mathbb{N}$, we have that*

$$(6.30) \quad \frac{1}{E_\phi(a + l, |b|)} \leq \frac{e^{-|b|\Theta_\phi(a+l, |b|)} |\mathcal{M}_{I_\phi}(a + i|b|)|}{W_\phi(a)Z_\phi^l(a, |b|)|\Gamma(a + i|b|)} \leq E_\phi(a + l, |b|).$$

REMARK 6.5.2. The upper bound for $\Theta_\phi(a, |b|)$, that is, $\frac{\pi}{2}$, is attained when $\phi(u) = u$ and $b = \infty$. In this case $W_\phi(z) = \Gamma(z)$.

Proof. Note that for $\phi \in \mathcal{B}$, $z = a + ib$ with $a, b > 0$, and with the usual notation for Δ_b , Δ_b^{\Re} and Δ_b^{\Im} used in Section 6.3, we have that

$$(6.31) \quad \frac{|\phi(a + ib)|}{\phi(a)} = \left| 1 + \frac{\Delta_b \phi(a)}{\phi(a)} \right| \geq \left(1 + \frac{\Delta_b^{\Re} \phi(a)}{\phi(a)} \right) \geq 1,$$

since, from (4.26), $\Delta_b^{\Re} \phi(a) \geq 0$. Thus, we get that $\Theta_\phi(a, b) \geq 0$. Next, using the first identity in (4.2), we deduce that

$$(6.32) \quad \begin{aligned} |\phi(bu + ib)| &\leq m + \sigma^2 |bu + ib| + |bu + ib| \int_0^\infty e^{-buy} \bar{\mu}(y) dy \\ &\leq \sqrt{1 + \frac{1}{u^2}} \left(m + \sigma^2 bu + bu \int_0^\infty e^{-buy} \bar{\mu}(y) dy \right) \\ &= \sqrt{1 + \frac{1}{u^2}} \phi(bu). \end{aligned}$$

Hence, after performing an obvious change of variables, we get that

$$\Theta_\phi(a, b) \leq \int_0^\infty \ln \left(\sqrt{1 + \frac{1}{u^2}} \right) du = \frac{\pi}{2},$$

which concludes item (1). From (6.18), we have that

$$(6.33) \quad |W_\phi(z)| = W_\phi(a) \frac{\phi(a)}{|\phi(z)|} \prod_{k=1}^{\infty} \left| \frac{\phi(a+k)}{\phi(z+k)} \right| = W_\phi(a) \frac{\phi(a)}{|\phi(z)|} |Z_\phi(z)|.$$

Let in the sequel $z = a + ib$, $b > 0$, $a > d_\phi$ or $a \geq 0$, if $\phi(0) = m > 0$ and $d_\phi = 0$, and $l \in \mathbb{N}$ such that $a_l = a + l > 0$. Then, we denote by

$$(6.34) \quad \frac{\phi(a)}{|\phi(z)|} |Z_\phi(z)| = \prod_{k=0}^l \left| \frac{\phi(a_k)}{\phi(a+k+ib)} \right| \lim_{n \rightarrow \infty} e^{S_n(a_l, b)},$$

where we recall that $a_k = a + k$ and we have set

$$S_n(a_l, b) = \ln \left(\prod_{k=1}^n \frac{\phi(a_{k+l})}{|\phi(a_{k+l} + ib)|} \right).$$

Therefore, from (6.33) and (6.34) it remains to estimate $S_n(a_l, b)$, as $n \rightarrow \infty$. For this purpose we rewrite

$$\begin{aligned} S_n(a_l, b) &= \ln \left(\prod_{k=1}^n \frac{\phi(a_{k+l})}{|\phi(a_{k+l} + ib)|} \right) \\ &= - \sum_{k=0}^n (\ln |\phi(a_{k+l} + ib)| - \ln \phi(a_{k+l})) + \ln \frac{|\phi(a_l + ib)|}{\phi(a_l)}. \end{aligned}$$

Then the approximation techniques developed in [Olv74, Section 8.2(2.01)] for the gamma function are valid for any twice differentiable function f such that $\int_0^n |f''(u)| du < \infty$ and take the form

$$\sum_{k=0}^n f(k) = \int_0^n f(u) du + \frac{1}{2} f(0) + \frac{1}{2} f(n) + \int_0^n \frac{B_2 - B_2(u - [u])}{2} f''(u) du,$$

where B_2 is the Bernoulli number and $B_2(u)$ the Bernoulli polynomial as defined in [Olv74, Section 8.1]. Applying this to the sum in $S_n(a_l, b)$ above with $f(u) = \ln \left(\frac{|\phi(a_l + ib + u)|}{\phi(a_l + u)} \right)$, $u \geq 0$, we get

$$(6.35) \quad \begin{aligned} S_n(a_l, b) &= - \int_0^n \ln \left(\frac{|\phi(u + a_l + ib)|}{\phi(u + a_l)} \right) du \\ &\quad - \frac{1}{2} \ln \left(\frac{|\phi(a_l + n + ib)| \phi(a_l)}{|\phi(a_l + ib)| \phi(a_l + n)} \right) + E_\phi^B(n, a_l) - E_\phi^B(n, a_l + ib), \end{aligned}$$

where formally $E_\phi^B(n, z) = \int_0^n \frac{B_2 - B_2(u-[u])}{2} (\ln |\phi(u+z)|)'' du$. Next, note that

$$\begin{aligned}
& \ln \left(\frac{|\phi(a_{l+n} + ib)|}{\phi(a_{l+n})} \right) \\
&= \ln \left(\left| 1 + \frac{\Delta_b^{\Re} \phi(a_{l+n}) + i \Delta_b^{\Im} \phi(a_{l+n})}{\phi(a_{l+n})} \right| \right) \\
&= \frac{1}{2} \ln \left(1 + 2 \frac{\Delta_b^{\Re} \phi(a_{l+n})}{\phi(a_{l+n})} + \left(\frac{\Delta_b^{\Re} \phi(a_{l+n})}{\phi(a_{l+n})} \right)^2 + \left(\frac{\Delta_b^{\Im} \phi(a_{l+n})}{\phi(a_{l+n})} \right)^2 \right) \\
(6.36) \quad &\leq \frac{1}{2} \ln \left(1 + b^2 \frac{|\phi''(a_{l+n})|}{\phi(a_{l+n})} + \frac{b^4}{4} \frac{|\phi''(a_{l+n})|^2}{\phi^2(a_{l+n})} + b^2 \frac{(\phi'(a_{l+n}))^2}{\phi^2(a_{l+n})} \right),
\end{aligned}$$

where we have used (4.26) for the inequality. Thus, from (6.31) for the first inequality and (4.4) and (4.5) for the second, we deduce that

$$(6.37) \quad 0 \leq \lim_{n \rightarrow \infty} \ln \left(\frac{|\phi(a_{l+n} + ib)|}{\phi(a_{l+n})} \right) \leq \frac{1}{2} \lim_{n \rightarrow \infty} \ln \left(1 + \frac{2b^2}{(l+n)} + \frac{b^4 + b^2}{(l+n)^2} \right) = 0.$$

Using the same inequality as in (6.36) to the integral term appearing on the right-hand side of (6.35), and the estimates in (4.28), we conclude, thanks to the dominated convergence theorem, that, for any $a_l, b > 0$,

$$(6.38) \quad \lim_{n \rightarrow \infty} \int_0^n \ln \left(\frac{|\phi(u + a_l + ib)|}{\phi(u + a_l)} \right) du = \int_0^\infty \ln \left(\frac{|\phi(u + a_l + ib)|}{\phi(u + a_l)} \right) du.$$

Since ϕ is log-concave on \mathbb{R}_+ , see Proposition 4.1.1(2), which means that its second derivative does not change sign, then by the estimate [Olv74, Section 8.2(2.04)] (with $m = 1$ and $f(u) = \ln \phi(a_l + u)$ in the notation of [Olv74, Section 8.2(2.04)]), we get that

$$(6.39) \quad \lim_{n \rightarrow \infty} |E_\phi^B(n, a_l)| \leq -\frac{1}{8} \lim_{n \rightarrow \infty} \int_0^n (\ln \phi(a_l + u))'' du = \frac{1}{8} \frac{\phi'(a_l)}{\phi(a_l)} \leq \frac{1}{8a_l},$$

where, we have used that (4.4) implies $\lim_{n \rightarrow \infty} \frac{\phi'(a_{l+n})}{\phi(a_{l+n})} = 0$ and the last inequality itself. Next, we focus on $E_\phi^B(n, a_l + ib)$. Since $(\ln |\phi(u + ib)|)'' \leq |(\ln \phi(u + ib))''|$, as $(\ln \phi(u + ib))'' = (\ln |\phi(u + ib)|)'' + i(\arg(\phi(u + ib)))''$, we have according to [Olv74, Section 8.2(2.04)] again that

$$\begin{aligned}
& \lim_{n \rightarrow \infty} |E_\phi^B(n, a_l + ib)| \\
(6.40) \quad &\leq \frac{1}{8} \int_0^\infty |(\ln \phi(u + a_l + ib))''| du \\
&= \frac{1}{8} \int_0^\infty \left| \frac{\phi(u + a_l + ib) \phi''(u + a_l + ib) - (\phi'(u + a_l + ib))^2}{(\phi(u + a_l + ib))^2} \right| du \\
&\leq \frac{1}{8} \int_{a_l}^\infty \left(\frac{|\phi''(u + ib)|}{|\phi(u + ib)|} + \left(\frac{|\phi'(u + ib)|}{|\phi(u + ib)|} \right)^2 \right) du \\
(6.41) \quad &\leq \frac{1}{8} \int_{a_l}^\infty \left(\sqrt{10} \frac{|\phi''(u)|}{\phi(u)} + 10 \left(\frac{\phi'(u)}{\phi(u)} \right)^2 \right) du \leq \frac{2\sqrt{10} + 10}{8a_l},
\end{aligned}$$

where for the final two inequalities we have employed all the estimates of (4.28). Collecting (6.37), (6.38), (6.39) and (6.40) we immediately deduce the upper bound

$$\begin{aligned} \lim_{n \rightarrow \infty} S_n(a_l, b) &\leq \frac{1}{8} \int_{a_l}^{\infty} \left(\frac{|\phi''(u+ib)|}{|\phi(u+ib)|} + \left(\frac{|\phi'(u+ib)|}{|\phi(u+ib)|} \right)^2 \right) du + \frac{1}{8} \frac{\phi'(a_l)}{\phi(a_l)} \\ &\quad - \int_0^{\infty} \ln \left(\frac{|\phi(u+a_l+ib)|}{\phi(u+a_l)} \right) du + \frac{1}{2} \ln \left(\frac{|\phi(a_l+ib)|}{\phi(a_l)} \right), \\ &= \ln \left(\frac{1}{E_\phi(a_l, |b|)} \right) - b\Theta_\phi(a_l, b) + \frac{1}{2} \ln \left(\frac{|\phi(a_l+ib)|}{\phi(a_l)} \right). \end{aligned}$$

The lower bound is the same but with error term $\ln(E_\phi(a_l, |b|))$. Therefore, from (6.34), the definition of Z_ϕ^l , see (6.28) and the definition of Z_ϕ , see (6.17), we get that

$$\begin{aligned} E_\phi(a_l, b) \frac{e^{-|b|\Theta_\phi(a_l, |b|)}}{Z_\phi^l(a, z)} &\leq \frac{\phi(a)}{|\phi(z)|} |Z_\phi(z)| = \frac{\phi(a)}{|\phi(z)|} \prod_{k=1}^{\infty} \left| \frac{\phi(k+a)}{\phi(k+z)} \right| \\ (6.42) \qquad \qquad \qquad &\leq \frac{1}{E_\phi(a_l, b)} \frac{e^{-|b|\Theta_\phi(a_l, |b|)}}{Z_\phi^l(a, z)}. \end{aligned}$$

Then (6.42) via (6.33) proves (6.29), which establishes the whole item (3) since we have allowed for all possible values of a claimed in item (3). The case when $b < 0$ is dealt with in the same manner since, for all $a > 0$ and $b \in \mathbb{R}$,

$$(6.43) \qquad \Im\phi(a-ib) = -\Im\phi(a+ib) \quad \text{and} \quad \Re\phi(a+ib) = \Re\phi(a-ib).$$

Item (2) follows by gathering the very last upper bounds in (6.39) and (6.40). The bounds (6.30) follow from Proposition 6.1.2, since $\mathcal{M}_{I_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)}$ and the previous result. \square

6.5.1. Proof of Theorem 6.0.2(1). Choosing $l = 0$ in (6.29) and using the definition of $Z_\phi^0(a, |b|)$ given in (6.28), one gets, for any $b \in \mathbb{R}$ and $a > 0$,

$$\begin{aligned} E_\phi(a, |b|) e^{-|b|\Theta_\phi(a, |b|)} \sqrt{\frac{\phi(a)}{|\phi(a+ib)|}} W_\phi(a) &\leq |W_\phi(a+ib)| \\ &\leq \sqrt{\frac{\phi(a)}{|\phi(a+ib)|} \frac{W_\phi(a) e^{-|b|\Theta_\phi(a, |b|)}}{E_\phi(a, |b|)}}. \end{aligned}$$

We complete the proof of item (1) by using the bounds for $E_\phi(a, |b|)$ stated in Proposition 6.5.1(2).

6.6. Large asymptotic behaviours of W_ϕ along imaginary lines

From the bounds (6.29), which hold for any $\phi \in \mathcal{B}$, we are able to derive precise information regarding the decay of $|W_\phi|$ and $|\mathcal{M}_{I_\phi}|$ along imaginary lines. We emphasize that these asymptotic estimates are interesting in their own right as we offer comprehensive statements for the entire class of Bernstein-Weierstrass products which encompasses many substantial special functions. However, our motivation to investigate in depth and accurately this large asymptotic behaviour comes from several substantial issues that arise later in this work, such as the existence and

uniform asymptotic bounds of the sequence of co-eigenfunctions, and, that can be solved by means of Mellin transform techniques.

6.6.1. Necessary and sufficient conditions for exponential decay of W_ϕ . We start by providing necessary and sufficient conditions for the exponential decay of $|W_\phi|$ along imaginary lines in terms of $\phi \in \mathcal{B}$ and its characteristic triplet (m, σ^2, μ) . Below the notation

$$f(a+ib) \overset{\infty}{\approx} g(a+ib) \quad (\text{resp. } f(a+ib) \overset{\infty}{\lesssim} g(a+ib)),$$

for $z = a+ib \in \mathbb{C}$ means that there exist positive finite constants $c_-(a)$ and $c_+(a)$ (resp. a positive finite constant $c_+(a)$) such that

$$(6.44) \quad c_-(a) \leq \liminf_{|b| \rightarrow \infty} \left| \frac{f(a+ib)}{g(a+ib)} \right| \leq \overline{\lim}_{|b| \rightarrow \infty} \left| \frac{f(a+ib)}{g(a+ib)} \right| \leq c_+(a)$$

(resp. $\overline{\lim}_{|b| \rightarrow \infty} \left| \frac{f(a+ib)}{g(a+ib)} \right| \leq c_+(a)$).

PROPOSITION 6.6.1. *Let $\phi \in \mathcal{B}$.*

- (1) *For any $a \geq 0$, writing $\underline{\Theta}_\phi(a) = \liminf_{|b| \rightarrow \infty} \Theta_\phi(a, |b|)$ and simply $\underline{\Theta}_\phi = \underline{\Theta}_\phi(0)$, we have that $\underline{\Theta}_\phi = \underline{\Theta}_\phi(a)$ and*

$$(6.45) \quad \underline{\Theta}_\phi \in \left[0, \frac{\pi}{2}\right].$$

- (2) *For any $\phi \in \mathcal{B}$ and $a > 0$, we have that*

$$(6.46) \quad \frac{e^{-\frac{\pi}{2}|b|}}{|b|^{\frac{1}{2}}} \overset{\infty}{\lesssim} \frac{e^{-\frac{\pi}{2}|b|}}{\sqrt{|\phi(a+ib)|}} \overset{\infty}{\lesssim} |W_\phi(a+ib)|.$$

- (3) *$\psi \in \mathcal{N}_{\underline{\Theta}}$, i.e. $\underline{\Theta}_\phi > 0$, if and only if, for any $a > 0$ and $\epsilon > 0$, we have that*

$$(6.47) \quad \frac{e^{-\frac{\pi}{2}|b|}}{|b|^{\frac{1}{2}}} \overset{\infty}{\lesssim} \frac{e^{-\frac{\pi}{2}|b|}}{\sqrt{|\phi(a+ib)|}} \overset{\infty}{\lesssim} |W_\phi(a+ib)| \overset{\infty}{\lesssim} \frac{e^{-(\underline{\Theta}_\phi - \epsilon)|b|}}{\sqrt{|\phi(a+ib)|}}.$$

If $\phi(0) = m > 0$ and $d_\phi = 0$ then the estimate extends to $a = 0$.

In all cases, we can retrieve, from Proposition 6.5.1, the positive constants $c_-(a)$ and $c_+(a)$ defined in (6.44).

REMARK 6.6.2. Estimate (6.47) together with $\mathcal{M}_{V_\psi} = W_\phi$ being the Mellin transform of ν yield the validity of (1.37) which in turn shows that $\nu \in \mathcal{A}(\underline{\Theta}_\phi)$, i.e. it is holomorphic in the sector $\mathbb{C}(\underline{\Theta}_\phi) = \{z \in \mathbb{C}; |\arg z| < \underline{\Theta}_\phi\}$ and verifies Theorem 5.0.2(3), i.e. $\psi \in \mathcal{N}_{\underline{\Theta}}$ implies that $\nu \in \mathcal{A}(\underline{\Theta}_\phi)$.

REMARK 6.6.3. We mention that in some statements here and also below we focus on the case, $\psi \in \mathcal{N}$, i.e. $\phi \in \mathcal{B}_{\mathcal{N}}$. However, up to some minor and obvious modifications, they can easily be extended to the general case $\phi \in \mathcal{B}$. We also point out that the estimate (6.48) below can serve for more precise study of the asymptotic.

Proof. Recall the definition of Θ_ϕ , (6.27), and put $\Theta_\phi(0, |b|) = \Theta_\phi(|b|)$. First, from (6.31) in the first inequality and (6.32) in the second, we have that

$$\begin{aligned}
\Theta_\phi(|b|) &\geq \Theta_\phi(a, |b|) = \Theta_\phi(|b|) - \int_0^{\frac{a}{|b|}} \ln \left(\frac{|\phi(|b|(u+i))|}{\phi(|b|u)} \right) du \\
&\geq \Theta_\phi(|b|) - \int_0^{\frac{a}{|b|}} \ln \left(\sqrt{1 + \frac{1}{u^2}} \right) du \\
(6.48) \quad &= \Theta_\phi(|b|) - \frac{a}{|b|} \ln \left(1 + \frac{b^2}{a^2} \right) - \arctan \left(\frac{a}{|b|} \right) \cong \Theta_\phi(|b|) - o(1).
\end{aligned}$$

Hence, from Proposition 6.5.1(1), we deduce that

$$\underline{\Theta}_\phi = \liminf_{|b| \rightarrow \infty} \Theta_\phi(a, |b|) = \liminf_{|b| \rightarrow \infty} \Theta_\phi(|b|) \in \left[0, \frac{\pi}{2} \right],$$

and, the proof of (6.45) is completed. We prove (6.46) by using the lower bound in (6.29) with $l = 0$ since from the second claim of Proposition 4.1.1(3) we have that $|\phi(z)| \cong \sigma^2 |z| + o(|z|)$, as $z = a + ib$, $a > 0$ fixed and as $|b| \rightarrow \infty$. Next, let $\underline{\Theta}_\phi > 0$ then (6.47) follows from the upper bound of (6.29) with $l = 0$ and (6.48). Since all constants in Proposition 6.5.1 are explicit we can recover the constants $c_\pm(a)$. This ends the proof. \square

We proceed by providing sufficient conditions for $\underline{\Theta}_\phi > 0$, that is for the exponential decay of the Mellin transforms along imaginary lines which are stated in Theorem 1.2.1(1).

6.7. Proof of Theorem 1.2.1(1)

Let us first prove that $\psi \in \mathcal{N}_\Theta$ implies $\mathfrak{r} = \phi(\infty) = \infty$. Thanks to (6.29) with $l = 0$ it will suffice to show that for some $a > 0$, $b\Theta(a, b) = o(b)$, as $b \rightarrow \infty$, since then we get $\liminf_{b \rightarrow \infty} \Theta(a, b) = \underline{\Theta}_\phi = 0$, that is a contradiction with the definition of \mathcal{N}_Θ ,

see Table 2. However, $\psi(u) = u\phi(u)$ with $\phi \in \mathcal{B}_\mathcal{N}$, and if $\phi(\infty) = \overline{\Pi}(0^+) + m = \int_0^\infty \overline{\Pi}(y) dy + m < \infty$, then $\sigma^2 = 0$ and for any fixed $a > 0$

$$\lim_{b \rightarrow \infty} \frac{\phi(ba + ib)}{\phi(ba)} = 1 + \lim_{b \rightarrow \infty} \frac{\int_0^\infty (1 - e^{iby}) e^{-bay} \overline{\Pi}(y) dy}{\phi(ba)} = 1,$$

since $\lim_{b \rightarrow \infty} \phi(ba) = \phi(\infty) < \infty$ and

$$\lim_{b \rightarrow \infty} \left| \int_0^\infty (1 - e^{iby}) e^{-bay} \overline{\Pi}(y) dy \right| \leq 2 \lim_{b \rightarrow \infty} \int_0^\infty e^{-bay} \overline{\Pi}(y) dy = 0.$$

Therefore, from the definition of $\Theta(a, b)$, see Theorem 6.0.2(1) and (6.32) which ensures the application of the dominated convergence theorem we deduce that $b\Theta(a, b) = o(b)$. Next, we show that if $\psi \in \mathcal{N}_G$, i.e. $\liminf_{u \rightarrow \infty} \frac{\overline{\Pi}(\frac{1}{u})}{\psi(u)} > 0$ or $\sigma^2 > 0$ then $\underline{\Theta}_\phi > 0$, i.e. $\psi \in \mathcal{N}_\Theta$. Let first assume that $\sigma^2 > 0$ and recall that $\phi(ba + ib) - \phi(ba) = \Delta_b^{\Re} \phi(ba) + i \Delta_b^{\Im} \phi(ba)$, see Section 4.3 for more information on these functions. The positivity of the integrand that defines $\Theta_\phi(a, b)$, see (6.31) and (6.27), the fact that

$\Delta_b^{\Re} \phi(ba) \geq 0$, the definition of $\underline{\Theta}_\phi$, see Proposition 6.6.1(1) and Fatou's lemma yield that

$$2\underline{\Theta}_\phi = 2 \liminf_{b \rightarrow \infty} \int_0^\infty \ln \left(\frac{|\phi(b(a+i))|}{\phi(ba)} \right) da \geq \int_1^\infty \liminf_{b \rightarrow \infty} \ln \left(1 + \left(\frac{\Delta_b^{\Im} \phi(ba)}{\phi(ba)} \right)^2 \right) da.$$

By means of Proposition 4.1.1(3) we have that $\phi(ba) = \sigma^2 ba + o(ba)$. This coupled with the first set of inequalities in (4.29) of Lemma 4.3.2, which is applicable since $\phi \in \mathcal{B}_{\mathcal{N}}$, gives that for all b big enough and $a \geq 1$

$$\begin{aligned} \frac{\Delta_b^{\Im} \phi(ba)}{\phi(ba)} &\geq \frac{\sigma^2 b + (e^{-\pi a} - e^{-2\pi a}) \int_0^{\frac{\pi}{b}} \sin(by) \overline{\Pi}(y) dy}{2\sigma^2 ba} \\ &\geq \frac{\sigma^2 + C_1 e^{-\pi a} (1 - e^{-\pi a}) \int_0^{\frac{1}{b}} y \overline{\Pi}(y) dy}{2\sigma^2 a}, \end{aligned}$$

where for the second inequality we have used $\sin(by) \geq C_1 by$, for some $C_1 > 0$ valid on $y \in (0, \frac{1}{b})$. Since $\lim_{b \rightarrow \infty} \int_0^{\frac{1}{b}} y \overline{\Pi}(y) dy = 0$, as $\int_0^1 y^2 \Pi(dy) < \infty$, see (4.18), we get the bound

$$(6.49) \quad \underline{\Theta}_\phi \geq \frac{1}{2} \int_1^\infty \ln \left(1 + \frac{1}{4a^2} \right) da > 0.$$

This completes the proof for $b \rightarrow \infty$ and $\sigma^2 > 0$. If $b \rightarrow -\infty$, $\sigma^2 > 0$, the arguments around (6.43) provide the same lower bound. Assume that $\sigma^2 = 0$, $b > 0$. Observe that

$$(6.50) \quad \underline{\Theta}_\phi = \liminf_{b \rightarrow \infty} \int_0^\infty \ln \left(\frac{|\phi(ba + ib)|}{\phi(ba)} \right) da \geq \int_0^\infty \liminf_{b \rightarrow \infty} \ln \left(1 + \frac{\Delta_b^{\Re} \phi(ba)}{\phi(ba)} \right) da,$$

thanks to (6.31) and the Fatou's lemma. Note that, for any $a > 0$,

$$b \Delta_b^{\Re} \phi(ba) = \int_0^\infty (1 - \cos(y)) e^{-ay} \overline{\Pi} \left(\frac{y}{b} \right) dy.$$

From the inequality $1 - \cos(y) = 2(\sin(y/2))^2 \geq 2(\frac{y}{\pi})^2$ valid on $y \in (0, 1)$, we get that

$$(6.51) \quad b \Delta_b^{\Re} \phi(ba) \geq 2 \frac{e^{-a}}{\pi^2} \int_0^1 y^2 \overline{\Pi} \left(\frac{y}{b} \right) dy \geq \frac{2}{3\pi^2} e^{-a} \overline{\Pi} \left(\frac{1}{b} \right).$$

Recall that $\Delta_b^{\Re} \phi(ba) \geq 0$. Thus, from (6.50), the positivity of the integrand therein and the fact that ϕ is non-decreasing, we confirm that

$$(6.52) \quad \underline{\Theta}_\phi = \liminf_{b \rightarrow \infty} \underline{\Theta}_\phi(|b|) \geq \int_0^1 \liminf_{b \rightarrow \infty} \ln \left(1 + \frac{2}{3\pi^2} e^{-a} \frac{\overline{\Pi}(\frac{1}{b})}{b\phi(b)} \right) da.$$

Since $b\phi(b) = \psi(b)$ and $\psi \in \mathcal{N}_G$ with $\sigma^2 = 0$, we deduce, with $\liminf_{b \rightarrow \infty} \frac{\overline{\Pi}(\frac{1}{b})}{\psi(b)} > C_2 > 0$, that the following inequality holds

$$\underline{\Theta}_\phi \geq \int_0^1 \ln \left(1 + \frac{2C_2}{3\pi^2} e^{-a} \right) da > 0,$$

which gives the proof for $b \rightarrow \infty$, $\psi \in \mathcal{N}_G$ and $\sigma^2 = 0$. The case $b \rightarrow -\infty$ comes from (6.43). Relation (1.16) follows from the global asymptotic (4.7) for ϕ given in Proposition 4.1.1(7).

6.8. Proof of Theorem 6.0.2 (2b): Examples of large asymptotic estimates of $|W_\phi|$

For the eigenvalues expansions of the gL semigroups, it is important to provide precise bounds for the norm of the sequence of co-eigenfunctions. Among the different strategies we implement to get such estimates is the Mellin transform technique. For this reason, in this part, we deepen our analysis on the asymptotic estimate of $|W_\phi(a+ib)|$ by either computing Θ_ϕ for substantial subclasses of $\mathcal{B}_\mathcal{N}$, or, detailing the exact subexponential decay for some subclasses. The results presented below extend with minor modifications to $|W_\phi(a+ib)|$ for $\phi \in \mathcal{B}$.

6.8.1. Proof of Theorem 6.0.2(2b)i): The case $\psi \in \mathcal{N}_P$. In the case $\sigma^2 > 0$, i.e. $\psi \in \mathcal{N}_P$, we obtain precise bounds particularly when in addition $\bar{\Pi}(0^+) < \infty$, or, if we have a good control on the tail $\bar{\Pi}$. For this purpose we introduce the following notation. For any $a > 0$, $b \in \mathbb{R}$, we set

$$(6.53) \quad L(a, |b|) = \frac{1}{\sigma^2} \int_0^1 e^{-ay} \left| \sin \left(\frac{|b|y}{2} \right) \right| \bar{\Pi}(y) \frac{dy}{y}$$

and

$$L(a, |b|) \leq \bar{L}(|b|) = \sup_{r \leq |b|} L(0, r).$$

PROPOSITION 6.8.1. *Let $\psi \in \mathcal{N}_P$. Fix any $a > d_\phi = \sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\}$. Then*

$$(6.54) \quad \frac{1}{|b|^{\frac{1}{2}}} e^{-\frac{\pi}{2}|b|} \lesssim |W_\phi(a+ib)| \lesssim |b|^{\frac{4m}{\sigma^2} + a - \frac{1}{2}} e^{-\frac{\pi}{2}|b| + \bar{L}(|b|)},$$

with $\bar{L}(|b|) \cong o(|b|)$. Moreover, when $\bar{\Pi}(0^+) < \infty$, we have $\bar{L}(|b|) \cong O(\ln(|b|))$, and, if for all $y \in (0, 1)$, $\bar{\Pi}(y) \leq Cy^{-\alpha}$ with $\alpha \in (0, 1)$, then

$$(6.55) \quad \bar{L}(|b|) \cong O(|b|^\alpha).$$

The proof of Proposition 6.8.1 is postponed after the following lemma whose statement requires some further notation. Set $\phi(z) = zg(z)$ and write formally

$$(6.56) \quad \bar{W}_g(z) = \frac{e^{-\gamma_g z}}{g(z)} \prod_{k=1}^{\infty} \frac{g(k)}{g(k+z)} e^{\frac{g'(k)}{g(k)} z},$$

where $\gamma_g = \gamma_\phi - \gamma$ and γ is the Euler-Mascheroni constant. We first show that $\bar{W}_g \in \mathcal{A}_{(d_\phi, \infty)}$ and give bounds on $|\bar{W}_g(z)|$.

LEMMA 6.8.2. *Let $\psi \in \mathcal{N}_P$.*

(1) *Then $\bar{W}_g \in \mathcal{A}_{(d_\phi, \infty)}$ and we have, on $\mathbb{C}_{(d_\phi, \infty)}$, that*

$$(6.57) \quad W_\phi(z) = \Gamma(z) \bar{W}_g(z).$$

(2) *There exists $k_0 = k_0(m, \sigma^2) \in \mathbb{N} \setminus \{0\}$ such that, with $z = a + ib$, $a > 0$,*

$$(6.58) \quad \frac{e^{-\mathfrak{W}_{k_0}(a+ib) - L(a, |b|)}}{C_{k_0}(a, |b|)} \leq \frac{|\bar{W}_g(a+ib)|}{\bar{W}_g(a)} \leq C_{k_0}(a, |b|) e^{\mathfrak{W}_{k_0}(a+ib) + L(a, |b|)},$$

where $L(a, |b|)$ is defined in (6.53),

$$(6.59) \quad \mathfrak{W}_{k_0}(z) = \ln \left| \prod_{k=0}^{k_0} \frac{g(k+a)}{g(k+z)} \right|,$$

$$\lim_{|z| \rightarrow \infty} \mathfrak{W}_{k_0}(z) = \lim_{|z| \rightarrow \infty} \ln \left| \prod_{k=0}^{k_0} \frac{g(k+a)}{g(k+z)} \right| = \frac{\prod_{k=0}^{k_0} g(k+a)}{\sigma^{2k_0+2}} \text{ and}$$

$$C_{k_0}(a, |b|) = \left(1 + \frac{|b|}{1+a}\right)^{\frac{4m}{\sigma^2}} e^{\frac{4e}{\sigma^2} \int_1^\infty \frac{e^{-(k_0+a)y}}{1-e^{-y}} \bar{\Pi}(y) dy}.$$

REMARK 6.8.3. The mapping $z \mapsto \frac{1}{g(z)} = \frac{z}{\phi(z)}$ has a root at zero if $m > 0$, which compensates the pole coming from $\Gamma(z)$ in (6.57) to ensure that $W_\phi(0) = \mathbb{E} \left[V_\phi^{-1} \right] = \phi(0) = m$. If W_ϕ extends further to the left then a zero of $\frac{1}{g}$ at zero coming from a term in the product in (6.56) cancels the impact of the poles of $\Gamma(z)$ at $z = -1, -2, \dots$.

Proof. Since $g(z) = \frac{\phi(z)}{z}$, we get, from (6.1) and the Weirstrass product representation of the gamma function, that (6.57) formally holds provided \bar{W}_g is holomorphic. Due to its definition in (6.56), this will follow if the involved infinite product is absolutely convergent. However, since the products defining W_ϕ and Γ are both absolutely convergent on $\mathbb{C}_{(0, \infty)}$ from (6.56), \bar{W}_g is absolutely convergent on $\mathbb{C}_{(0, \infty)}$ and the analyticity follows. When $d_\phi < 0$, (6.57) follows from an argument involving the recurrent equation (6.6). In the sequel we provide the bounds for $|\bar{W}_g(z)|$. To this end put $z = a + ib$, $a > d_\phi$, $b > 0$, and set

$$\begin{aligned} \bar{W}_g(z) &= \frac{e^{-\gamma_g a}}{g(a)} \prod_{k=1}^{\infty} \frac{g(k)}{g(k+a)} e^{\frac{g'(k)}{g(k)} a} \times \frac{g(a) e^{-i\gamma_g b}}{g(z)} \prod_{k=1}^{\infty} \frac{g(k+a)}{g(k+z)} e^{i \frac{g'(k)}{g(k)} b} \\ &= \bar{W}_g(a) \bar{W}_g^{(a)}(z), \end{aligned}$$

and, we proceed to estimate $|\bar{W}_g^{(a)}(z)|$. Since $\phi \in \mathcal{B}_N$, then Proposition 4.1.1(8) gives that $\bar{\mu}(y) = \bar{\Pi}(y)$, $y > 0$, and then the first identity of (4.2) gives on $\mathbb{C}_{(d_\phi, \infty)} \setminus \{0\}$ that

$$g(z) = \frac{\phi(z)}{z} = \frac{m}{z} + \sigma^2 + \int_0^\infty e^{-zy} \bar{\Pi}(y) dy.$$

For any $u = a + k > 0$ and $b > 0$,

$$\begin{aligned} \frac{g(u)}{g(u+ib)} &= 1 + \sigma^{-2} \frac{\frac{mib}{u(u+ib)} + \int_0^\infty e^{-uy} (1 - e^{-iby}) \bar{\Pi}(y) dy}{1 + \sigma^{-2} \left(\frac{m}{u+ib} + \int_0^\infty e^{-uy - iby} \bar{\Pi}(y) dy \right)} \\ &= 1 + \frac{\rho_m(u, b) + \rho_\Pi(u, b)}{1 + \tilde{\rho}_m(u, b) + \tilde{\rho}_\Pi(u, b)}, \end{aligned}$$

where each term ρ or $\tilde{\rho}$ in the last expression matches the corresponding term in the middle one. Set $\bar{C} = m\sigma^{-2}$. We have the following bounds $|\tilde{\rho}_\Pi(u, b)| \leq \sigma^{-2} \int_0^\infty e^{-uy} \bar{\Pi}(y) dy$,

$$(6.60) \quad |\tilde{\rho}_m(u, b)| \leq \frac{\bar{C}}{u+b} \quad \text{and} \quad |\rho_m(u, b)| \leq \frac{\bar{C}b}{u(u+b)}$$

which all tend to zero as $u \rightarrow \infty$. Thus, for $u \geq u_0 > 0$, where $u_0 = u_0(m, \sigma^2)$, such that

$$1 - |\tilde{\rho}_m(u, b)| - |\tilde{\rho}_\Pi(u, b)| \geq \frac{1}{2},$$

we have using $|1 - e^{-iby}| = 2 \left| \sin\left(\frac{by}{2}\right) \right|$ that

$$\begin{aligned} & \left| \frac{g(u)}{g(u+ib)} - 1 \right| \\ & \leq \frac{2\bar{C}b}{u(u+b)} + \frac{2}{\sigma^2} \int_0^\infty e^{-uy} |1 - e^{-iby}| \bar{\Pi}(y) dy \\ (6.61) \quad & \leq \frac{4\bar{C}b}{u(u+b)} + \frac{4}{\sigma^2} \left(\int_1^\infty e^{-uy} \bar{\Pi}(y) dy + \int_0^1 e^{-uy} \left| \sin\left(\frac{by}{2}\right) \right| \bar{\Pi}(y) dy \right). \end{aligned}$$

Splitting the product that defines $\bar{W}_g^{(a)}$ at $k_0 = k_0(m, \sigma^2) = \lceil u_0 \rceil + 2 + |a|$, using (6.61), (6.59) for the definition of \mathfrak{W}_{k_0} and $\ln(1+x) \leq x$, $x \geq 0$, we get that

$$\begin{aligned} & \ln \left| \bar{W}_g^{(a)}(a+ib) \right| - \ln \mathfrak{W}_{k_0}(a+ib) \\ & = \sum_{k=k_0}^\infty \ln \left| \frac{g(k+a)}{g(k+a+ib)} \right| \\ & \leq \sum_{k=k_0}^\infty \left| \frac{g(k+a)}{g(k+a+ib)} - 1 \right| \\ & \leq \sum_{k=k_0}^\infty \frac{4\bar{C}b}{(k+a)(k+a+b)} \\ & \quad + \sum_{k=k_0}^\infty \frac{4}{\sigma^2} \left(\int_1^\infty \frac{\bar{\Pi}(y)}{e^{(k+a)y}} dy + \int_0^1 e^{-(k+a)y} \left| \sin\left(\frac{by}{2}\right) \right| \bar{\Pi}(y) dy \right) \\ & \leq 4\bar{C} \left(\sum_{k \geq 2} \frac{b}{(k+a)(k+a+b)} \right) + \frac{4e}{\sigma^2} \left(\int_1^\infty \frac{e^{-(k_0+a)y}}{1-e^{-y}} \bar{\Pi}(y) dy + L(a, b) \right), \end{aligned}$$

where implicitly in the last term of the inequality to get $L(a, b)$ we have used $1 - e^{-y} \geq y/e$ on $(0, 1)$. Furthermore, since

$$\sum_{k \geq 2} \frac{b}{(k+a)(k+a+b)} \leq \lim_{A \rightarrow \infty} \int_1^A \left(\frac{1}{r+a} - \frac{1}{r+a+b} \right) dr = \ln \left(1 + \frac{b}{1+a} \right),$$

we obtain, with $C_{k_0}(a, b)$ defined in the statement of the lemma, that

$$\frac{e^{-\mathfrak{W}_{k_0}(a+ib) - L(a, b)}}{C_{k_0}(a, b)} \leq \left| \bar{W}_g^{(a)}(a+ib) \right| \leq C_{k_0}(a, b) e^{\mathfrak{W}_{k_0}(a+ib) + L(a, b)}.$$

The bounds (6.58) follow for $b > 0$. The case $b < 0$ is dealt with similarly as in (6.43). \square

Proof of Proposition 6.8.1. The lower bound in (6.54) is the lower bound in (6.47). We get from (6.57) that $|W_\phi(z)| = |\Gamma(z)\bar{W}_g(z)|$. For $z = a+ib$, $a > 0$ and

$|b|$ large enough the upper bound in (6.54) follows from the classical asymptotic for the gamma function, i.e. for $a > 0$ fixed,

$$(6.62) \quad |\Gamma(a + i|b|)| \approx C_a |b|^{a-\frac{1}{2}} e^{-\frac{\pi}{2}|b|},$$

where $C_a > 0$, combined with (6.58) and taking into account all polynomial dependence on $|b|$. If $d_\phi < 0$ and $a \in (d_\phi, 0]$ then the functional equation (6.6) relates the asymptotic of $|W_\phi(a + ib)|$ to that of $|W_\phi(a + [-d_\phi - 1] + ib)|$. The polynomial decay in (6.54) is again $|b|^{\frac{4m}{\sigma^2} + a - \frac{1}{2}}$ as on each iteration of (6.6) we collect a term of the type, for a fixed a , $|\phi(a + ib)| \approx \sigma^2 |b|$, see Proposition 4.1.1(3). Next, recall the definition of $L(a, |b|)$, see (6.53), and therefore, observe, using $|\sin(y)| \leq |y| \wedge 1$, that, for any $0 < \epsilon < 1$,

$$\begin{aligned} 0 \leq \lim_{|b| \rightarrow \infty} \frac{L(a, |b|)}{|b|} &\leq \lim_{|b| \rightarrow \infty} \frac{\sup_{r \leq |b|} L(0, r)}{|b|} = \lim_{|b| \rightarrow \infty} \frac{\bar{L}(|b|)}{|b|} \\ &\leq \int_0^\epsilon \bar{\Pi}(y) dy + \lim_{|b| \rightarrow \infty} \frac{\int_\epsilon^1 \bar{\Pi}(y) \frac{dy}{y}}{b} = \int_0^\epsilon \bar{\Pi}(y) dy, \end{aligned}$$

which shows that $\bar{L}(|b|) \approx o(|b|)$ since $\int_0^1 \bar{\Pi}(y) dy < \infty$, see (4.18). Finally, it remains to study $\bar{L}(|b|)$ for specific instances. First, when $\bar{\Pi}(0^+) < \infty$, we have that

$$\sup_{r \leq |b|} L(0, r) = \bar{L}(|b|) \leq \bar{\Pi}(0^+) \int_0^{|b|} \frac{|\sin(r)|}{r} dr \leq \bar{\Pi}(0^+) \ln(|b|).$$

Then, if for some $\alpha \in (0, 1)$, $\bar{\Pi}(y) \stackrel{0}{=} O(y^{-\alpha})$, then trivially

$$(6.63) \quad L(|b|) \leq C_1 |b|^\alpha \int_0^{|b|} \left| \sin\left(\frac{r}{2}\right) \right| \frac{dr}{r^{1+\alpha}} \leq C_2 |b|^\alpha,$$

which completes the proof of our Proposition 6.8.1. \square

6.8.2. Proof of Theorem 6.0.2(2(b)ii): The case $\psi \in \mathcal{N}_\alpha$. We continue to apply the theory developed above to functions ψ which are regularly varying at infinity without aiming at the most general case of regular variation. We simply attempt to illustrate the tractability of our approach with the aim to compute explicitly and easily Θ_ϕ . We now show that for $\psi \in \mathcal{N}_\alpha$, with $\alpha \in (0, 1)$, we have $\Theta_\phi = \frac{\pi}{2}\alpha > 0$. To this end set $\bar{C} = C_\alpha \Gamma(\alpha + 1)$. Since $\psi \in \mathcal{N}_\alpha$, $\phi \in \mathcal{B}_\mathcal{N}$ from Proposition 4.1.1(8) we get that $\sigma^2 = 0$, $\bar{\mu}(y) = \bar{\Pi}(y)$, $y > 0$, and $\bar{\Pi}(y) \stackrel{0}{\sim} \bar{C} y^{-\alpha}$, $\alpha \in (0, 1)$. The latter and a standard Tauberian theorem imply that Proposition 4.1.1(7) can be augmented to $\phi(a|b|) \approx \bar{C} \Gamma(1 - \alpha) a^\alpha |b|^\alpha$. Next, for fixed $a > 0$ and $|b| \rightarrow \infty$,

we get from the first expression for ϕ in (4.2) and $\bar{\mu}(y) = \bar{\Pi}(y) \stackrel{0}{\sim} \bar{C}y^{-\alpha}$ that

$$\begin{aligned}
& \phi(|b|a + i|b|) \\
&= m + (|b|a + i|b|) \int_0^\infty e^{-i|b|y - |b|ay} \bar{\Pi}(y) dy \\
&= m + (a + i) \left(\int_0^\infty \cos(y) e^{-ya} \bar{\Pi} \left(\frac{y}{|b|} \right) dy - i \int_0^\infty \sin(y) e^{-ya} \bar{\Pi} \left(\frac{y}{|b|} \right) dy \right) \\
&\approx \bar{C}|b|^\alpha (a + i) \left(\int_0^\infty \cos(y) e^{-ya} \frac{dy}{y^\alpha} - i \int_0^\infty \sin(y) e^{-ya} \frac{dy}{y^\alpha} \right) \\
&= \bar{C}|b|^\alpha (a + i) \int_0^\infty e^{-y(a+i)} \frac{dy}{y^\alpha} = \bar{C}\Gamma(1 - \alpha)(a + i)^\alpha |b|^\alpha.
\end{aligned}$$

In the last integral we have used Cauchy's theorem on the closed contour $[0, \beta a + i\beta], [\beta a + i\beta, \beta], [\beta, 0]$ with $\beta \rightarrow \infty$ to the function $e^{-z}z^{-\alpha}$ which is holomorphic on $\mathbb{C}_{(0, \infty)}$. Since Proposition 6.6.1(1) holds we compute from (6.27) with $a = 1$ and the asymptotic relations above

$$\Theta_\phi = \lim_{|b| \rightarrow \infty} \int_{\frac{1}{b}}^\infty \ln \left(\frac{|\phi(bua + ib)|}{\phi(ba)} \right) da = \int_0^\infty \lim_{|b| \rightarrow \infty} \ln \left(\frac{|\phi(ba + ib)|}{\phi(ba)} \right) da = \frac{\pi}{2} \alpha.$$

The dominated convergence theorem holds thanks to inequality (6.32)

6.8.3. Proof of Theorem 6.0.2(2)(b)iii): Examples of subexponential decay of $|W_\phi|$. We further illustrate our approach by detailing more examples that reveal again that for several important families of Lévy measures, we can derive explicit bounds for the rate of decay of $|W_\phi(z)|$.

PROPOSITION 6.8.4. *Let $\psi \in \mathcal{N} \setminus \mathcal{N}_P$, i.e. $\sigma^2 = 0$, and fix $a > d_\phi$.*

(1) *Assume that for some $\alpha \in (0, 1)$, $\lim_{y \rightarrow 0} y^\alpha \bar{\Pi}(y) > 0$ and $\phi(\infty) = \mathfrak{r} < \infty$.*

Then, there exists $C_{a, \alpha} > 0$ such that

$$|W_\phi(a + ib)| \lesssim^\infty e^{-C_{a, \alpha} |b|^\alpha}.$$

(2) *Let us assume that $\lim_{y \rightarrow 0} y \bar{\Pi}(y) > 0$, then there exists $C_a > 0$ such that*

$$|W_\phi(a + ib)| \lesssim^\infty e^{-C_a \frac{|b|}{\phi(b)}}.$$

(3) *Let us assume that $\bar{\Pi}(y) = y^{-1} |\ln y|^2 1_{[0, 1/2]}(y)$ then*

$$|W_\phi(a + ib)| \lesssim^\infty e^{-(\ln b)^3}.$$

Proof. First, we observe that by combining (6.51) with (6.52), we get, for $b \in \mathbb{R}$,

$$(6.64) \quad \begin{aligned} |b|\Theta_\phi(|b|) &\geq |b| \int_0^1 \ln \left(1 + Ce^{-y} \frac{\bar{\Pi}\left(\frac{1}{|b|}\right)}{|b|\phi(|b|)} \right) dy \\ &\geq |b| \ln \left(1 + C_1 \frac{\bar{\Pi}\left(\frac{1}{|b|}\right)}{|b|\phi(|b|)} \right), \end{aligned}$$

where $C_1 = Ce^{-1} > 0$. The two first statements are direct applications of (6.64). Indeed, if $\phi(\infty) < \infty$ then from (4.8) $\infty > \bar{\Pi}(0+) \geq \int_0^a \bar{\Pi}(y)dy \geq a\bar{\Pi}(a)$, for any $a \in (0, 1)$. Thus $\lim_{a \rightarrow 0} a\bar{\Pi}(a) = 0$ and so with $\ln(1+x) \stackrel{0}{\sim} x$ in (6.64) we conclude the claim using $\liminf_{y \rightarrow 0} y^\alpha \bar{\Pi}(y) > 0$. Item (2) follows from (6.64) and $\liminf_{y \rightarrow 0} y\bar{\Pi}(y) > 0$. Item (3) follows similarly by observing that for $u > 1/2$, $\phi(u) = \frac{u}{\ln(u)}$ and plugging this expression in (6.64). \square

6.8.4. Asymptotic of the Mellin transform in the case $\psi \in \mathcal{N}_\infty^c$. Finally, we study the case when $\psi \in \mathcal{N}_\infty^c$, i.e. $\bar{\Pi}(0+) < \infty$ and $\sigma^2 = 0$, which corresponds to the Laplace exponent of a spectrally negative compound Poisson process with a positive drift.

LEMMA 6.8.5. *Let $\psi \in \mathcal{N}_\infty^c$. Then, $\phi(u) = -\int_0^\infty e^{-uy}\bar{\Pi}(y)dy + \mathfrak{r}$, where recall that $\mathfrak{r} = \phi(\infty) = \bar{\Pi}(0+) + m$. Then, for $z \in \mathbb{C}_{(d_\phi, \infty)}$,*

$$(6.65) \quad W_\phi(z) = \frac{\mathfrak{r}^z}{\phi(z)} \prod_{k=1}^{\infty} \frac{\phi(k)}{\phi(k+z)}.$$

On the real line $W_\phi(u)\mathfrak{r}^{-u}$ is bounded and decreasing, as $u \rightarrow \infty$, and for fixed $a > 0$,

$$(6.66) \quad \liminf_{|b| \rightarrow \infty} |W_\phi(a+ib)| = 0.$$

Proof. Since $\lim_{n \rightarrow \infty} \ln \phi(n) = \ln \mathfrak{r}$ then from (6.2) $\lim_{n \rightarrow \infty} \sum_{k=1}^n \frac{\phi'(k)}{\phi(k)} = \ln \mathfrak{r} + \gamma_\phi$ and (6.65) follows immediately from (6.1). To conclude the other statements we study the product in (6.65). As

$$u \mapsto \bar{W}_\phi(u) = \prod_{k=1}^{\infty} \frac{\phi(k)}{\phi(k+u)} \text{ is bounded and decreasing on } \mathbb{R}_+,$$

we deduce that $\bar{W}_\phi(u) \in [0, 1]$. It remains to show (6.66). For any $a > 0$, we employ (6.29) with $l = 0$. Substituting the expression $\phi(z) = -\int_0^\infty e^{-zy}\bar{\Pi}(y)dy + \mathfrak{r}$ in (6.27) we get that

$$\begin{aligned} b\Theta_\phi(a, |b|) &= b \int_{\frac{a}{b}}^{\infty} \ln \left| \frac{\mathfrak{r} - \int_0^\infty e^{-buy - iby}\bar{\Pi}(y)dy}{\mathfrak{r} - \int_0^\infty e^{-buy}\bar{\Pi}(y)dy} \right| du \\ &= b \int_{\frac{a}{b}}^{\infty} \ln \left| 1 + \frac{\int_0^\infty e^{-byu} (1 - e^{-iby}) \bar{\Pi}(y)dy}{\phi(bu)} \right| du. \end{aligned}$$

Since ϕ is non-decreasing choose $a > 0$ big enough so that

$$\sup_{u \geq \frac{a}{b}} \frac{\left| \int_0^\infty e^{-byu} (1 - e^{-iby}) \bar{\Pi}(y) dy \right|}{\phi(bu)} \leq 2 \frac{\bar{\Pi}(0^+)}{a\phi(a)} < \frac{1}{2}.$$

Using, $\ln|1+z| \geq C|z|$, with some $C > 0$, for all $|z| < \frac{1}{2}$, we get that

$$b\Theta_\phi(a, |b|) \geq \frac{C}{\phi(a)} b \int_{\frac{a}{b}}^\infty \left| \int_0^\infty e^{-byu} (1 - e^{-iby}) \bar{\Pi}(y) dy \right| du.$$

From Lemma 4.3.2 $\int_0^\infty e^{-byu} \sin(by) \bar{\Pi}(y) dy = \Delta_b^{\Im} \phi(bu) \geq 0$. Therefore, we complete the proof by estimating from below with the real part of the expression above that is

$$\begin{aligned} & \lim_{b \rightarrow \infty} b \int_{\frac{a}{b}}^\infty \int_0^\infty e^{-byu} (1 - \cos(by)) \bar{\Pi}(y) dy du \\ &= \lim_{b \rightarrow \infty} \int_0^\infty e^{-ay} (1 - \cos(by)) \bar{\Pi}(y) \frac{dy}{y} \\ &= \lim_{b \rightarrow \infty} \int_0^\infty e^{-\frac{a}{b}y} (1 - \cos(y)) \bar{\Pi}\left(\frac{y}{b}\right) \frac{dy}{y} \\ &= \infty. \end{aligned}$$

Thus, $\lim_{b \rightarrow \infty} b\Theta_\phi(a, |b|) = \infty$ and from (6.29) with $l = 0$ (6.66) holds for all a big enough since all other quantities in (6.29) with $l = 0$ are bounded when $\bar{\Pi}(0^+) < \infty$. However, from (6.6) we also deduce that (6.66) holds for all $a > d_\phi$. \square

Intertwining relations and a set of eigenfunctions

In this chapter, we elaborate and exploit intertwining relations that relate the entire class of generalized Laguerre semigroups to the classical Laguerre semigroup of order 0. This commutation relation between NSA semigroups and a self-adjoint semigroup is the central concept in our development of the spectral decomposition of these operators. In this perspective, it is proved to be very useful to characterize a set of eigenfunctions of the gL semigroups and to provide properties of this set regarded as a sequence in a Hilbert space.

Let us recall that $Q = (Q_t)_{t \geq 0}$ denotes the classical Laguerre semigroup associated via the Lamperti bijection to $\psi(u) = u^2$, i.e. $\sigma^2 = 1$, $m = 0$ and $\Pi \equiv 0$ in (1.3), and its associate Feller process is a diffusion. Basic facts about this semigroup are reviewed in Example 3.0.1 where, for instance, one finds that it is self-adjoint in $L^2(\varepsilon)$ where

$$\varepsilon(x) = e^{-x}, \quad x > 0,$$

is the density of an exponential distribution of parameter 1. We also recall that, for any $\phi \in \mathcal{B}_{\mathcal{N}}$, the Markov multiplicative operator \mathcal{I}_ϕ is defined, for at least any $f \in \mathbf{C}_b(\mathbb{R}_+)$, by

$$(7.1) \quad \mathcal{I}_\phi f(x) = \mathbb{E}[f(xI_\phi)], \quad x > 0,$$

where, with $\eta = (\eta_t)_{t \geq 0}$ a subordinator with Laplace exponent ϕ , I_ϕ is the positive random variable

$$(7.2) \quad I_\phi = \int_0^\infty e^{-\eta_t} dt.$$

We have all ingredients to state our first main result of this chapter.

THEOREM 7.0.1. *Let $\psi \in \mathcal{N}$ and recall that $\phi(u) = \frac{\psi(u)}{u} \in \mathcal{B}_{\mathcal{N}}$. Then, for any $t \geq 0$, we have the following intertwining identity*

$$(7.3) \quad P_t \mathcal{I}_\phi f = \mathcal{I}_\phi Q_t f$$

valid for all $f \in L^2(\varepsilon)$. Moreover, the following properties of the intertwining operator hold.

- (1) $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu)) \cap \mathbf{B}(\mathbf{C}_0(\overline{\mathbb{R}_+}))$.
- (2) $\overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$. However, there exists $C > 0$ such that for all $f \in L^2(\nu)$,

$$\|\mathcal{I}_\phi f\|_\nu \geq C \|f\|_\varepsilon$$

if and only if $\psi(u) = \sigma^2 u^2$, $\sigma^2 > 0$.

(3) \mathcal{I}_ϕ is a Markov operator determined by its integer moments, and we have, for any $z \in \mathbb{C}_{(-1, \infty)}$ and $x > 0$, recalling that $p_z(x) = x^z$,

$$(7.4) \quad \mathcal{I}_\phi p_z(x) = \frac{\Gamma(z+1)}{W_\phi(z+1)} p_z(x).$$

REMARK 7.0.2. We point out that the intertwining identity (7.3) generalizes to all $\psi \in \mathcal{N}$ the relation obtained by Carmona et al. [CPY98] between the so-called saw-teeth process and the Bessel process, that is when $\psi(u) = u^{\frac{u+1-a}{u+b}}$, $0 < a < 1 < a+b$.

To state the next result concerning the characterization and properties of a set of eigenfunctions, we introduce notions discussed in detail in Chapter 11. We say that a sequence $(P_n)_{n \geq 0}$ in the Hilbert space $L^2(\nu)$ is a Bessel sequence if there exists $A > 0$ such that the inequality

$$(7.5) \quad \sum_{n=0}^{\infty} |\langle f, P_n \rangle_\nu|^2 \leq A \|f\|_\nu^2$$

holds for all $f \in L^2(\nu)$. If in addition $\overline{\text{Span}}(P_n) = L^2(\nu)$ and there exists $B > 0$ such that for all finite scalar sequences $(c_n)_{n \geq 0}$

$$(7.6) \quad B \sum_{n=0}^{\infty} c_n^2 \leq \left\| \sum_{n=0}^{\infty} c_n P_n \right\|_\nu^2$$

then $(P_n)_{n \geq 0}$ is a Riesz basis in $L^2(\nu)$. Finally, we recall the notation $\mathcal{P}_0(x) = 1$ and for $n \in \mathbb{N}$, \mathcal{P}_n is the polynomial defined by

$$(7.7) \quad \mathcal{P}_n(x) = \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{W_\phi(k+1)} x^k,$$

where from (1.18) $W_\phi(n+1) = \prod_{k=1}^n \phi(k)$. The polynomials $(\mathcal{P}_n)_{n \geq 0}$ can be seen as the Laguerre polynomials perturbed by the Stieltjes moment sequence of the intertwining operator \mathcal{I}_ϕ , see (3). To state the next result we need the following terminology which is due to Blumenthal and Gettoor [BG61]. For a Bernstein function ϕ , we define its lower index as follows

$$(7.8) \quad \underline{\phi} = \sup\{a > 0; \lim_{u \rightarrow \infty} u^{-a} \phi(u) = \infty\} = \liminf_{u \rightarrow \infty} \frac{\ln \phi(u)}{\ln u} \in [0, 1],$$

with the usual convention $\sup\{\emptyset\} = 0$. This quantity appears in substantial path properties of the associated subordinators, and, for instance, we point out that $\underline{\phi}$ corresponds to the Hausdorff dimension of their range and is related to their packing dimension, see e.g. [Ber96, Chap. III] and [Sav14].

THEOREM 7.0.3. *Let $\psi \in \mathcal{N}$.*

(1) *For any $n \geq 0, t \geq 0$, \mathcal{P}_n is an eigenfunction for P_t associated to the eigenvalue e^{-nt} , i.e. $\mathcal{P}_n \in L^2(\nu)$ and*

$$(7.9) \quad P_t \mathcal{P}_n(x) = e^{-nt} \mathcal{P}_n(x).$$

We also have, for any $n \in \mathbb{N}$,

$$(7.10) \quad \|\mathcal{P}_n\|_\nu \leq 1.$$

(2) Moreover $\overline{\text{Span}(\mathcal{P}_n)} = L^2(\nu)$ and $(\mathcal{P}_n)_{n \geq 0}$ is a Bessel sequence but it is not a Riesz basis in $L^2(\nu)$.

(3) The sequence of polynomials $(\mathcal{P}_n)_{n \geq 0}$ satisfies, for any $n \geq 2$, the following three term perturbed recurrence relation

$$(7.11) \quad \mathcal{P}_n(x) = \left(2 - \frac{1}{n}\right) \mathcal{P}_{n-1}(x) - \frac{x}{n\phi(1)} \mathcal{P}_{n-1}^{(\mathcal{T}_1)}(x) - \left(1 - \frac{1}{n}\right) \mathcal{P}_{n-2}(x),$$

where

$$\mathcal{P}_n^{(\mathcal{T}_1)}(x) = \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{W_{\mathcal{T}_1\phi}(k+1)} x^k$$

and the transform $\mathcal{T}_1\phi(u) = \frac{u}{u+1}\phi(u+1) \in \mathcal{B}_{\mathcal{N}}$ is discussed in Proposition 4.1.1(9).

(4) The sequence $(\mathcal{P}_n)_{n \geq 0}$ is formed of orthogonal polynomials in some weighted L^2 space if and only if $\psi(u) = \sigma^2 u^2 + mu$, $\sigma^2 > 0$, $m \geq 0$, i.e. $(\mathcal{P}_n)_{n \geq 0}$ is the sequence of Laguerre polynomials of order $m \geq 0$, see (3.3) and (3.8) for definition.

(5) The polynomials $(\mathcal{P}_n(-x))_{n \geq 0}$ are the Jensen polynomials associated to the entire function

$$\mathfrak{J}_{W_\phi}(x) = \sum_{n=0}^{\infty} \frac{1}{W_\phi(n+1)} \frac{x^n}{n!},$$

i.e. for any $x, t \in \mathbb{R}$, we have that

$$(7.12) \quad e^t \mathfrak{J}_{W_\phi}(xt) = \sum_{n=0}^{\infty} \mathcal{P}_n(-x) \frac{t^n}{n!}.$$

Moreover, denoting \mathfrak{o}_ϕ (resp. \mathfrak{t}_ϕ) the order (resp. the type) of the entire function \mathfrak{J}_{W_ϕ} , we have

$$\mathfrak{o}_\phi = \frac{1}{1 + \underline{\phi}} \in \left[\frac{1}{2}, 1\right]$$

and

$$\mathfrak{t}_\phi \geq (1 + \underline{\phi}) e^{-\frac{\phi}{1+\underline{\phi}}} \frac{1}{\varliminf_{n \rightarrow \infty} \phi(n) n^{-\underline{\phi}}}.$$

Note that $\mathfrak{o}_\phi = 1$ if and only if $\underline{\phi} = 0$, and, in this case, $\mathfrak{t}_\phi = \frac{1}{\tau}$, with the usual convention $\frac{1}{\infty} = 0$. Finally, we have, for large n , any $x > 0$ and any integer p ,

$$(7.13) \quad \mathcal{P}_n^{(p)}(-x) = O\left(\mathcal{E}_\phi(nx) e^{-n} (\mathfrak{J}_\Gamma(n) + \mathfrak{J}_\Gamma(-n))\right) = O\left(n^{p+\frac{1}{2}} \mathcal{E}_\phi(nx)\right),$$

where, for any $\epsilon > 0$, we set

$$\mathcal{E}_\phi(x) = e^{\mathfrak{t}_\phi x^{\mathfrak{o}_\phi}} \mathbb{I}_{\{0 < \mathfrak{t}_\phi < \infty\}} + e^{\epsilon x^{\mathfrak{o}_\phi}} \mathbb{I}_{\{\mathfrak{t}_\phi = 0\}} + e^{x^{\mathfrak{o}_\phi + \epsilon}} \mathbb{I}_{\{\mathfrak{t}_\phi = \infty\}}$$

and \mathfrak{J}_Γ is the modified Bessel function of order 0.

REMARK 7.0.4. The entire function \mathfrak{J}_{W_ϕ} was introduced by the first author in [Pat09] as the increasing 1-invariant function of the self-similar semigroup K and boils down, when $\psi(u) = u^2$, i.e. $W_\phi = \Gamma$, to \mathfrak{J}_Γ , which explains the notation.

REMARK 7.0.5. Although we shall provide another proof, we mention that property (4) regarding the necessary and sufficient conditions for the orthogonality of $(\mathcal{P}_n)_{n \geq 0}$ can be deduced from an elegant result of Chihara [Chi68] stating that the Laguerre polynomials are the only sequence of orthogonal polynomials generating the so-called Brenke type function of the form (7.12).

The remaining part of this chapter contains the proof of these two statements as well as the uniqueness of the invariant measure.

7.1. Proof of Theorem 7.0.1

7.1.1. Factorization of invariant measures. To prove the identity

$$(7.14) \quad P_t \mathcal{I}_\phi f = \mathcal{I}_\phi Q_t f, \text{ for any } f \in L^2(\varepsilon) \text{ and } t \geq 0,$$

i.e. the intertwining relation of Theorem 7.0.1(7.3) we proceed in two steps. First, we establish this identity in the space $\mathcal{C}_0(\overline{\mathbb{R}}_+)$. Then, we show that it extends to $L^2(\varepsilon)$. For the first step, by recalling from Definition 1.0.1 the notation

$$K_t f = P_{\ln(t+1)} d_{t+1} f, \quad t \geq 0, \quad f \in \mathcal{C}_0(\overline{\mathbb{R}}_+),$$

i.e. K is the Feller semigroup of a conservative self-similar Markov process on $[0, \infty)$, we observe that relation (7.14) considered on $\mathcal{C}_0(\overline{\mathbb{R}}_+)$ is equivalent, because of the deterministic space-time transformation, to the intertwining relation, for any $f \in \mathcal{C}_0(\overline{\mathbb{R}}_+)$,

$$(7.15) \quad K_t \mathcal{I}_\phi f = \mathcal{I}_\phi K_t^{(0)} f, \quad t \geq 0,$$

where we recall that $K^{(0)}$ stands for the semigroup of a Bessel process of dimension 2, see (3.7). Now, to prove (7.15) we resort to a criterion which has been provided by Carmona et al. [CPY97, Proposition 3.2]. More precisely, they showed that the intertwining relation (7.15) is valid on $\mathcal{C}_0(\overline{\mathbb{R}}_+)$ whenever the following two conditions are satisfied.

- (1) The ensuing factorization of entrance laws, which characterizes the invariant measures in our setting,

$$K_1 \mathcal{I}_\phi f(0) = K_1^{(0)} f(0),$$

holds. Since according to (2.22) with $t = 1$ we have that

$$K_1 f(0) = \int_0^\infty f(x) \nu(x) dx = \mathcal{V}_\psi f(1),$$

we emphasize that, in our notation, this factorization translates to

$$(7.16) \quad \mathcal{V}_\psi \mathcal{I}_\phi f(1) = \mathcal{E} f(1),$$

where the Markov operators are defined in Lemma 7.1.1 below.

- (2) The operator associated to the entrance law of the semigroup K is injective in $\mathcal{C}_0(\overline{\mathbb{R}}_+)$. More specifically, by means of the self-similar property of index 1 of the semigroup K , this means that

$$(7.17) \quad \text{for any } f, g \in \mathcal{C}_0(\overline{\mathbb{R}}_+), \text{ if for all } t > 0, \mathcal{V}_\psi f(t) = \mathcal{V}_\psi g(t) \text{ then } f = g.$$

Factorization (7.16) is the purpose of Lemma 7.1.1. Lemma 7.1.4 provides the condition (7.17).

LEMMA 7.1.1. *Let $\psi \in \mathcal{N}$ and recall that $\psi(u) = u\phi(u)$. Then, we have the following factorization of the multiplicative Markov operator \mathcal{E} associated to an exponential random variable \mathbf{e} of parameter 1, i.e. with*

$$\mathcal{E}f(x) = \mathbb{E}[f(x\mathbf{e})] = \int_0^\infty f(xy)e^{-y}dy, \quad f \in \mathbf{C}_0(\overline{\mathbb{R}}_+),$$

it holds that

$$(7.18) \quad \mathcal{E}f(x) = \mathcal{V}_\psi \mathcal{I}_\phi f(x) = \mathcal{I}_\phi \mathcal{V}_\psi f(x), \quad x > 0,$$

where we have set $\mathcal{V}_\psi f(x) = \mathbb{E}[f(xV_\psi)]$ with V_ψ defined in Proposition 2.3.1(1). Let $g \in L^2(\varepsilon)$ then (7.18) holds with $f = g^2$ at $x = 1$.

REMARK 7.1.2. We mention that, in the special case $\phi(0) = 0$, identity (7.18) can be reformulated as the factorization of the exponential law identified by Bertoin and Yor [BY02b]. We also point out that in Lemma 10.1.2 below, under some additional conditions, we shall provide another factorization of this type which is useful for proving the completeness property of the sequence of co-eigenfunctions $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$.

Proof. Since $\psi \in \mathcal{N}$, $\phi \in \mathcal{B}_\mathcal{N}$, and we know from (6.4) of Theorem 6.0.1(1) and Proposition 6.1.1(3) that both operators \mathcal{I}_ϕ and \mathcal{V}_ψ are moment determinate and more precisely we have, for any $n \geq 0$, $x \geq 0$, $p_n(x) = x^n$, using (6.11) and (2.20), that

$$(7.19) \quad \mathcal{I}_\phi p_n(x) = \frac{n!}{W_\phi(n+1)} p_n(x) \quad \text{and} \quad \mathcal{V}_\psi p_n(x) = W_\phi(n+1) p_n(x).$$

Hence,

$$\mathcal{I}_\phi \mathcal{V}_\psi p_n(x) = \mathcal{V}_\psi \mathcal{I}_\phi p_n(x) = n! p_n(x) = \mathcal{E}p_n(x)$$

and identity (7.18) follows from the fact that the law of the exponential random variable \mathbf{e} associated to \mathcal{E} is also moment determinate. If $g \in L^2(\varepsilon)$ then

$$\|g\|_\varepsilon^2 = \mathcal{E}g^2(1) = \mathcal{V}_\psi \mathcal{I}_\phi g^2(1) = \mathcal{I}_\phi \mathcal{V}_\psi g^2(1)$$

follows from elementary application of Fubini's theorem. \square

We proceed with the following corollary which is a consequence of the Bernstein-Weierstrass representation of the Mellin transform developed in Chapter 6.

COROLLARY 7.1.3. *For any $\phi \in \mathcal{B}$, we have $\mathcal{M}_{V_\phi}(z) \neq 0$, for any $z \in \mathbb{C}_{(d_\phi, \infty)}$, and, if in addition $\phi(0) = m > 0$, $d_\phi = 0$, then $\mathcal{M}_{V_\phi}(z) \neq 0$, for any $z \in \mathbb{C}_{[0, \infty)}$. Similarly, $\mathcal{M}_{I_\phi}(z) \neq 0$ on the strip $\mathbb{C}_{(0, \infty)}$. And, if in addition $\phi(0) = m = 0$, ϕ does not vanish on $\mathbb{C}_0 \setminus \{0\}$ and $\phi'(0^+) < \infty$ then $\mathcal{M}_{I_\phi}(z) \neq 0$ on the strip $\mathbb{C}_{[0, \infty)}$.*

Proof. The proof follows from the one of Theorem 6.0.1 (resp. Proposition 6.1.2) of the absolute convergence of the complex logarithm of \mathcal{M}_{V_ϕ} (resp. \mathcal{M}_{I_ϕ}), see from (6.18) onwards. \square

Recall that $L^\infty(A)$ stands for the set of measurable a.e. bounded functions on $A \subseteq \mathbb{R}$. The next result proves the second condition, that is (7.17).

LEMMA 7.1.4. *For any $\psi \in \mathcal{N}$, $\mathcal{V}_\psi \in \mathbf{B}(\mathbf{C}_0(\overline{\mathbb{R}}_+))$ and is one-to-one in $L^\infty(\mathbb{R}_+)$.*

Proof. From the definition of a Markov operator, for any $f \in L^\infty(\mathbb{R}_+)$ and $y \in \mathbb{R}$,

$$(7.20) \quad \mathcal{V}_\psi f(e^y) = \mathbb{E}[f(V_\psi e^y)] = \mathbb{E}[f_e(y + \ln V_\psi)] = f_e \star \tilde{\nu}_e(-y),$$

where we have set

$$\tilde{\nu}_e(y) = e^{-y}\nu(e^{-y}), f_e(y) = f(e^y) \in L^\infty(\mathbb{R})$$

and \star stands for the standard additive convolution. Assume that there exists $g \in L^\infty(\mathbb{R}_+)$ such that

$$\mathcal{V}_\psi g(e^y) = g_e \star \tilde{\nu}_e(-y) = 0, \text{ for all } y \in \mathbb{R}.$$

Then, according to the Wiener's Theorem, see e.g. [BGT87, Theorem 4.8.4(ii)], we have, for some $b \in \mathbb{R}$, that

$$(7.21) \quad \mathcal{M}_{V_\psi}(1 + ib) = \int_0^\infty x^{ib}\nu(x)dx = \int_{\mathbb{R}} e^{-iby}\tilde{\nu}_e(y)dy = 0,$$

which contradicts Corollary 7.1.3, where it is stated that \mathcal{M}_{V_ψ} is zero-free at least on $z \in \mathbb{C}_{(0,\infty)}$ and we know from (6.4) that for $\phi \in \mathcal{B}_\mathcal{N}$, $V_\phi \stackrel{(d)}{=} V_\psi$ and thus $\mathcal{M}_{V_\phi} = \mathcal{M}_{V_\psi}$. \square

Now it remains to show that the intertwining relation (7.14) extends from $\mathbf{C}_0(\overline{\mathbb{R}_+})$ to $L^2(\varepsilon)$. To this end, we first prove the following.

7.1.2. Proof of Theorem 7.0.1(1). First, observe that \mathcal{I}_ϕ is plainly linear. Next, let $f \in L^2(\varepsilon)$ and using Hölder's inequality and the very last claim of Lemma 7.1.1, we get that

$$(7.22) \quad \begin{aligned} \|\mathcal{I}_\phi f\|_\nu^2 &\leq \int_0^\infty \mathcal{I}_\phi f^2(x)\nu(x)dx = \mathcal{V}_\psi \mathcal{I}_\phi f^2(1) \\ &= \mathcal{E}f^2(1) = \|f\|_\varepsilon^2 < \infty, \end{aligned}$$

which shows that $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$. The fact that $\mathcal{I}_\phi \in \mathbf{B}(\mathbf{C}_0(\overline{\mathbb{R}_+}))$ follows immediately by the dominated convergence which completes the proof of Theorem 7.0.1(1).

7.2. End of the proof of the intertwining relation (7.3)

We recall that $\mathbf{C}_c(\mathbb{R}_+)$, the space of continuous functions with compact support in \mathbb{R}_+ , is dense in $L^2(\varepsilon)$ and since $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, and, from Theorem 1.1.1(3), for all $t \geq 0$, $Q_t \in \mathbf{B}(L^2(\varepsilon))$ and $P_t \in \mathbf{B}(L^2(\nu))$, we conclude the extension of (7.14) from $\mathbf{C}_0(\overline{\mathbb{R}_+})$ to $L^2(\varepsilon)$ by a density argument.

7.3. Proofs of Theorem 7.0.1(2) and (3)

Let $\psi \in \mathcal{N}$ and thus $\phi \in \mathcal{B}_\mathcal{N}$. Next, we recall that, for any $n \geq 0$, $p_n(x) = x^n$, and from (7.19) we have immediately that $p_n \in L^2(\nu)$, $n \geq 0$. Also from $\psi(u) = u\phi(u)$ and (7.19) the monomials are eigenfunctions for \mathcal{I}_ϕ in $L^2(\nu)$ and

$$\mathcal{I}_\phi p_n(x) = \frac{n!}{W_\phi(n+1)} = \frac{(n!)^2}{\prod_{k=1}^n \psi(k)} p_n(x) = \lambda_n p_n(x),$$

for all $n \geq 0$. Moreover, as the probability measure $\nu(x)dx$ of the positive random variable V_ψ is moment determinate, see Theorem 6.0.1(1), the polynomials are dense

in $L^2(\nu)$, see [Akh65, Chap. 2, Cor. 2.3.3, p. 45], which proves $\overline{\text{Ran}}(\mathcal{I}_\phi) = L^2(\nu)$ in $L^2(\nu)$. To prove the last claim of item (2), we check with the notation above and $\psi(u) = u\phi(u)$ that, for all $n \geq 0$,

$$\begin{aligned} \|\mathcal{I}_\phi p_n\|_\nu^2 &= \lambda_n^2 \|p_n\|_\nu^2 = \lambda_n^2 W_\phi(2n+1) \\ &= \Gamma(2n+1) \frac{\prod_{k=n+1}^{2n} \left(\frac{\psi(k)}{k^2}\right)}{\prod_{k=1}^n \left(\frac{\psi(k)}{k^2}\right)}. \end{aligned}$$

Since $\|p_n\|_\varepsilon^2 = \Gamma(2n+1)$, we get that \mathcal{I}_ϕ is not bounded from below if and only if

$$\begin{aligned} \lim_{n \rightarrow \infty} \frac{\|\mathcal{I}_\phi p_n\|_\nu^2}{\|p_n\|_\varepsilon^2} &= \lim_{n \rightarrow \infty} \frac{\prod_{k=n+1}^{2n} \left(\frac{\psi(k)}{k^2}\right)}{\prod_{k=1}^n \left(\frac{\psi(k)}{k^2}\right)} \\ (7.23) \qquad &= \lim_{n \rightarrow \infty} \frac{\prod_{k=1}^n \left(\frac{\psi(k+n)}{(k+n)^2}\right)}{\prod_{k=1}^n \left(\frac{\psi(k)}{k^2}\right)} = 0. \end{aligned}$$

Next, note from (1.3) that an integration by parts yields that

$$\frac{\psi(k)}{k^2} = \sigma^2 + \frac{m}{k} + \int_0^\infty e^{-ky} \overline{\overline{\Pi}}(y) dy,$$

where we recall that $\overline{\overline{\Pi}}(y) = \int_y^\infty \overline{\overline{\Pi}}(r) dr$, see (1.5). If $\sigma^2 > 0$, and, assume without a loss of generality that $\sigma^2 = 1$, then from (7.23) it suffices to show that

$$\begin{aligned} \lim_{n \rightarrow \infty} \sum_{k=1}^n \ln \left(1 - \frac{\frac{mn}{k(n+k)} + \overline{\varphi}_k(n)}{1 + \frac{m}{k} + \int_0^\infty e^{-ky} \overline{\overline{\Pi}}(y) dy} \right) &\asymp - \lim_{n \rightarrow \infty} \sum_{k=1}^n \left(\frac{mn}{k(n+k)} + \overline{\varphi}_k(n) \right) \\ &= -\infty, \end{aligned}$$

where we have set

$$\overline{\varphi}_k(n) = \int_0^\infty e^{-ky} (1 - e^{-ny}) \overline{\overline{\Pi}}(y) dy.$$

If $m > 0$ the claim is clear. If $m = 0$, unless $\Pi(dy) \equiv 0 dy$, $y > 0$, which holds if and only if

$$\overline{\overline{\Pi}}(0^+) = \int_0^\infty y \Pi(dy) = 0,$$

we have since $\overline{\overline{\Pi}}$ is non-increasing on \mathbb{R}_+ and $\overline{\overline{\Pi}}(0^+) > 0$

$$- \lim_{n \rightarrow \infty} \sum_{k=1}^n \int_0^\infty e^{-ky} (1 - e^{-ny}) \overline{\overline{\Pi}}(y) dy = - \int_0^\infty \frac{e^{-y}}{1 - e^{-y}} \overline{\overline{\Pi}}(y) dy = -\infty.$$

If $\sigma^2 = 0$ the claim is obvious by noting that from the first equality in (4.2) of Proposition 4.1.1 we get that $k^{-2}\psi(k) = k^{-1}\phi(k)$ is non-increasing to zero, as $k \rightarrow \infty$, and therefore the last product in (7.23) is bounded from above by

$$\frac{\psi(n+1)}{(n+1)^2 \psi^2(1)} = \frac{\phi(n+1)}{(n+1) \psi^2(1)} \stackrel{\infty}{=} o(1).$$

Thus \mathcal{I}_ϕ is not bounded from below whenever Π is not the zero measure on \mathbb{R}_+ or $m > 0$. Otherwise, $\phi(u) = \sigma^2 u$ and by moment identification from (7.19) we have that $\nu(x) = \sigma^{-2} e^{-\sigma^{-2}x}$, $x > 0$, and $I_\phi \stackrel{(d)}{=} \sigma^{-2}$. That is for all $f \in L^2(\nu)$, we have that

$$\|\mathcal{I}_\phi f\|_\nu^2 = \int_0^\infty f^2(\sigma^{-2}x) \sigma^{-2} e^{-\sigma^{-2}x} dx = \|f\|_\varepsilon^2$$

and the claim follows. This completes the proof of Theorem 7.0.1 since item (3) is easily deduced from the first identity in (7.19).

7.4. Proof of the uniqueness of the invariant measure

The intertwining relation allows also to prove the uniqueness of the invariant measure, that is $\nu(x)dx$ is the unique invariant measure of the Feller semigroup $(P_t)_{t \geq 0}$. Indeed, assume that there exists a measure $\tilde{\nu}(dx) \neq \nu(x)dx$ such that for all $f \in \mathbf{C}_0(\overline{\mathbb{R}_+})$,

$$\tilde{\nu} P_t f = \tilde{\nu} f = \int_0^\infty f(x) \tilde{\nu}(dx).$$

From Theorem 7.0.1(1) we have that $\mathcal{I}_\phi \in \mathbf{B}(\mathbf{C}_0(\overline{\mathbb{R}_+}))$. Since $\mathbf{C}_0(\overline{\mathbb{R}_+}) \subseteq L^2(\varepsilon)$ then the intertwining relation (7.3) holds in $\mathbf{C}_0(\overline{\mathbb{R}_+})$ and we obtain that

$$\tilde{\nu} \mathcal{I}_\phi Q_t f = \tilde{\nu} P_t \mathcal{I}_\phi f = \tilde{\nu} \mathcal{I}_\phi f.$$

Therefore,

$$\tilde{\nu}(x)dx = \int_0^\infty \iota\left(\frac{x}{y}\right) \frac{\tilde{\nu}(dy)}{y} dx,$$

where ι is the density of I_ϕ , see Proposition 6.1.1(1), is an invariant measure for the classical Laguerre semigroup Q , and, thus by uniqueness of its invariant measure $\bar{\nu}(x) = \varepsilon(x) = e^{-x}$, $x > 0$. Hence, the factorization (7.18) holds with \mathcal{V}_ψ replaced by $\check{\mathcal{V}}$ that is the Markov operator associated to the variable \check{V} with law $\tilde{\nu}$. By taking Mellin transform on both sides of this new factorization and noting from Corollary 7.1.3 that $\mathcal{M}_{\mathcal{I}_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)}$ is zero-free on $\mathbb{C}_{(0,\infty)}$ we get that $\mathcal{M}_{V_\psi}(z) = \mathcal{M}_{\tilde{\nu}}(z)$, i.e. $\tilde{\nu}(dx) = \nu(x)dx$.

7.5. Proof of Theorem 7.0.3

Let $\psi \in \mathcal{N}$. Using the form of the Laguerre polynomials, see (3.4), the first identity of (7.19) and the linearity of \mathcal{I}_ϕ , we first note, for any $n \geq 0$, that

$$(7.24) \quad \mathcal{I}_\phi \mathcal{L}_n(x) = \sum_{k=0}^n (-1)^k \binom{n}{k} \mathbb{E}[I_\phi^k] \frac{x^k}{k!}$$

$$(7.25) \quad = \sum_{k=0}^n (-1)^k \frac{\binom{n}{k}}{W_\phi(k+1)} x^k = \mathcal{P}_n(x).$$

As, for all $n \geq 0$, $\mathcal{L}_n \in L^2(\varepsilon)$ and $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, we get that $\mathcal{P}_n \in L^2(\nu)$, and,

$$P_t \mathcal{I}_\phi \mathcal{L}_n(x) = \mathcal{I}_\phi Q_t \mathcal{L}_n(x) = e^{-nt} \mathcal{I}_\phi \mathcal{L}_n(x),$$

where we have used successively the intertwining relation (7.14), the eigenfunction property of the Laguerre polynomials, see (3.6), and, again the linearity property

of \mathcal{I}_ϕ . This proves the first claim of Theorem 7.0.3(1). The second one, i.e. the bound (7.10), is obtained by observing that, for any $n \in \mathbb{N}$,

$$\|\mathcal{P}_n\|_\nu = \|\mathcal{I}_\phi \mathcal{L}_n\|_\nu \leq \|\mathcal{L}_n\|_\varepsilon = 1,$$

where we used that $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ is a contraction, see (7.22), and the Laguerre polynomials form an orthonormal basis of $L^2(\varepsilon)$. Next, using the fact that V_ψ is moment determinate, see Theorem 6.0.1(1), the polynomials are dense in $L^2(\nu)$, see [Akh65, Chap. 2, Cor. 2.3.3, p. 45], which gives that $\overline{\text{Span}(\mathcal{P}_n)} = L^2(\nu)$ in $L^2(\nu)$. Moreover, for any $f \in L^2(\nu)$, the Bessel property of the sequence $(\mathcal{P}_n)_{n \geq 0}$ is obtained from the following relations

$$(7.26) \quad \begin{aligned} \sum_{n=0}^{\infty} |\langle f, \mathcal{P}_n \rangle_\nu|^2 &= \sum_{n=0}^{\infty} |\langle f, \mathcal{I}_\phi \mathcal{L}_n \rangle_\nu|^2 = \sum_{n=0}^{\infty} |\langle \mathcal{I}_\phi^* f, \mathcal{L}_n \rangle_\varepsilon|^2 = \|\mathcal{I}_\phi^* f\|_\varepsilon^2 \\ &\leq \|f\|_\nu^2, \end{aligned}$$

where we have used the Parseval identity for the Laguerre polynomials $(\mathcal{L}_n)_{n \geq 0}$ in $L^2(\varepsilon)$ and that $\mathcal{I}_\phi^* \in \mathbf{B}(L^2(\nu), L^2(\varepsilon))$ is a contraction as the adjoint of the contraction $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, see Theorem 7.0.1(1). Next, assume that $(\mathcal{P}_n)_{n \geq 0}$ is a Riesz basis, then being already a complete Bessel sequence it means that for every sequence $(s_n)_{n \geq 0}$ in $\ell^2(\mathbb{N})$ there exists a unique $f \in L^2(\nu)$ such that, for any $n \geq 0$,

$$s_n = \langle \mathcal{P}_n, f \rangle_\nu.$$

Since $(\mathcal{L}_n)_{n \geq 0}$ is an orthonormal basis in $L^2(\varepsilon)$, we have that $(s_n)_{n \geq 0} \in \ell^2(\mathbb{N})$ if and only if there exists a unique $g \in L^2(\varepsilon)$ such that for each $n \in \mathbb{N}$, $s_n = \langle \mathcal{L}_n, g \rangle_\varepsilon$. Therefore,

$$s_n = \langle \mathcal{P}_n, f \rangle_\nu = \langle \mathcal{I}_\phi \mathcal{L}_n, f \rangle_\nu = \langle \mathcal{L}_n, \mathcal{I}_\phi^* f \rangle_\varepsilon$$

implies that the equation $\mathcal{I}_\phi^* f = g$ must have a unique solution for any $g \in L^2(\varepsilon)$. In turn this means from the open mapping theorem that, with the obvious notation, $\mathcal{I}_\phi = \mathcal{I}_\phi^{**}$ is bounded from below which from Theorem 7.0.1(2) provides a contradiction and completes the proof of item (2). Next, from the three term recurrence relation satisfied by the Laguerre polynomials, see (3.5), combined with the identity (7.24), we get easily, by linearity, that, for any $n \geq 2$,

$$\begin{aligned} \mathcal{P}_n(x) &= \mathcal{I}_\phi \mathcal{L}_n(x) = \mathcal{I}_\phi \left(\left(2 + \frac{-1-x}{n} \right) \mathcal{L}_{n-1}(x) - \left(1 - \frac{1}{n} \right) \mathcal{L}_{n-2}(x) \right) \\ &= \left(2 - \frac{1}{n} \right) \mathcal{P}_{n-1}(x) - \frac{x}{n} \sum_{k=0}^{n-1} (-1)^k \binom{n-1}{k} \mathbb{E} \left[I_\phi^{k+1} \right] \frac{x^k}{k!} \\ &\quad - \left(1 - \frac{1}{n} \right) \mathcal{P}_{n-2}(x). \end{aligned}$$

The recurrence relationship (7.11) follows after observing that

$$\mathbb{E} \left[I_\phi^{k+1} \right] = \frac{1}{\phi(1)} \mathbb{E} \left[I_{\mathcal{T}_1 \phi}^k \right],$$

where the transformation $\mathcal{T}_1 \phi(u) = \frac{u}{u+1} \phi(u+1) \in \mathcal{B}$ was introduced in Proposition 4.1.1(9). Finally, from a classical result in the theory of orthogonal polynomials, see e.g. [Sze75, Theorem 3.2.1], the sequence of polynomials $(\mathcal{P}_n)_{n \geq 0}$ is orthogonal in

some weighted L^2 space if and only if there exist real constants A_n, B_n, C_n , $n \geq 2$, such that,

$$(7.27) \quad \mathcal{P}_n(x) = (A_n x + B_n) \mathcal{P}_{n-1}(x) + C_n \mathcal{P}_{n-2}(x).$$

First, equating the coefficients of the monomials $x^n, 1$ of (7.27) we get, since $W_\phi(n+1) = \prod_{k=1}^n \phi(k)$, that

$$A_n = -\frac{1}{\phi(n)}, \quad B_n + C_n = 1.$$

Then, equating the monomials x, x^{n-1} in (7.27) and from the expressions for A_n, C_n we get that

$$B_n = 2 - \frac{\phi(1)}{\phi(n)} = n - \frac{\phi(n-1)}{\phi(n)}(n-1).$$

This allows us to check that relation (7.27) holds if and only if the following equation is satisfied

$$(n-2)\phi(n) - (n-1)\phi(n-1) + \phi(1) = 0.$$

It is easy to verify that $\phi(n) = \sigma^2 n + m$, with $\sigma^2, m \geq 0$, is a solution to this equation. To show that there do not exist other solutions, we set

$$g(u) = \frac{\phi(3) - \phi(1)}{\phi(u+2) - \phi(1)}, \quad u \geq 0,$$

and, observe that g is a solution to the functional equation

$$g(n+1) = \frac{n+1}{n+2} g(n), \quad n \geq 1,$$

with $g(1) = 1$. Hence,

$$g(n) = \frac{\phi(3) - \phi(1)}{\phi(n+2) - \phi(1)} = \frac{2}{n+1},$$

which completes the proof of (4). Next, from [CC89, Proposition 2.1(ii)], easy algebra yields the identity (7.12), i.e.

$$e^t \mathfrak{J}_{W_\phi}(xt) = \sum_{n=0}^{\infty} \mathcal{P}_n(-x) \frac{t^n}{n!},$$

for any $x, t \in \mathbb{R}$. Furthermore, we observe, for any $p = 1, \dots, n-1$, and $x \in \mathbb{R}$, and, modifying slightly here the notation to emphasize the dependency on ϕ , that

$$\begin{aligned} (\mathcal{P}_n^\phi(-x))^{(p)} &= \sum_{k=p}^n \frac{\Gamma(k+1)}{\Gamma(k-p+1)} \frac{\binom{n}{k}}{W_\phi(k+1)} x^{k-p} \\ &= \frac{\Gamma(n+1)}{\Gamma(n-p+1) W_\phi(p+1)} \sum_{k=0}^{n-p} \frac{\binom{n-p}{k}}{W_{\phi_p}(k+1)} x^k \\ (7.28) \quad &= \frac{\Gamma(n+1)}{\Gamma(n-p+1) W_\phi(p+1)} \mathcal{P}_{n-p}^{\phi_p}(-x), \end{aligned}$$

where we recall that since $\phi_p(u) = \phi(u+p)$, we have, from (6.14),

$$W_{\phi_p}(k+1) = \phi(k+p) W_{\phi_p}(k) = \frac{W_\phi(k+p+1)}{W_\phi(p+1)}.$$

The expression

$$\mathfrak{o}_\phi = \left(1 + \underline{\lim}_{n \rightarrow \infty} \frac{\log \phi(n)}{\log n}\right)^{-1} = \frac{1}{1 + \underline{\phi}} \in \left[\frac{1}{2}, 1\right]$$

of the order of the entire function \mathfrak{J}_{W_ϕ} as well as the lower bound for its type

$$\mathfrak{t}_\phi \geq (1 + \underline{\phi}) e^{-\frac{\phi}{1+\underline{\phi}}} \frac{1}{\underline{\lim}_{n \rightarrow \infty} \phi(n) n^{-\underline{\phi}}},$$

have been found in [BP14, Proposition 5.2]. We simply reproduce here the proof of the expression of the type when $\mathfrak{o}_\phi = 1$ or equivalently when $\underline{\phi} = 0$. We recall that

$$\mathfrak{J}_{W_\phi}(x) = \sum_{n=0}^{\infty} \frac{1}{W_\phi(n+1)} \frac{x^n}{n!}.$$

From the classical formula of the type of an entire function, see e.g. [Lev96, Chap. 1, Theorem 3], with $\mathfrak{o}_\phi = 1$, using the asymptotic equivalent of

$$W_\phi(n+1) \asymp C_\phi \sqrt{\phi(n)} e^{G(n)}, \quad G(n) = \int_1^n \ln \phi(u) du,$$

in (5.2) and the classical Stirling estimate

$$\Gamma(n+1) \asymp \frac{1}{\sqrt{2\pi}} n^{n-\frac{1}{2}} e^{-n},$$

we get that

$$\mathfrak{t}_\phi = \frac{1}{e} \overline{\lim}_{n \rightarrow \infty} \left(\frac{n^n}{\Gamma(n+1) W_\phi(n+1)} \right)^{\frac{1}{n}} = \overline{\lim}_{n \rightarrow \infty} \exp \left(- \left(\frac{\ln \phi(n)}{2n} + \frac{G(n)}{n} \right) \right) = \frac{1}{\mathfrak{r}},$$

where we have used the fact that $\ln \phi(n) = O(\ln n)$, which follows from Proposition 4.1.1(3), and, by l'Hôpital's rule,

$$\lim_{n \rightarrow \infty} \frac{G(n)}{n} = \lim_{n \rightarrow \infty} \ln \phi(n) = \ln \phi(\infty) = \ln \mathfrak{r}.$$

Next, observe that, for any $p \geq 0$, with $\phi_p(u) = \phi(u+p)$ as above,

$$\mathfrak{o}_{\phi_p} = \left(1 + \underline{\lim}_{n \rightarrow \infty} \frac{\log \phi(n+p)}{\log n}\right)^{-1} = \left(1 + \underline{\lim}_{n \rightarrow \infty} \frac{\log \phi(n) \log \phi(n+p)}{\log \phi(n) \log n}\right)^{-1} = \mathfrak{o}_\phi,$$

where the last identity holds when $\phi(\infty) = \infty$ using

$$\phi(n) \leq \phi(n+p) \leq \left(\frac{p}{n} + 1\right) \phi(n),$$

which follows from (4.4) and Proposition 4.1.1(2), and is obvious when $\phi(\infty) < \infty$. Also using [Lev96, Chap. 1, Theorem 3] it can be shown that $\mathfrak{t}_\phi = \mathfrak{t}_{\phi_p}$. Next, from (7.12), we get, after performing a change of variables, that, for all $n, x > 0$,

$$\begin{aligned} \mathcal{P}_n(-x) &= \frac{n!}{2\pi i} \oint_n e^z \mathfrak{J}_{W_\phi}(zx) \frac{dz}{z^{n+1}} \\ &= \frac{n!}{2\pi i} x^n \oint_{nx} e^{\frac{z}{x}} \mathfrak{J}_{W_\phi}(z) \frac{dz}{z^{n+1}}, \end{aligned}$$

where the last contour is a circle centered at 0 with radius $nx > 0$. Then, the definition of the order, see [Lev96, Chapter 1], yields that for any $x > 0$ and for large n ,

$$\max_{|z|=nx} |\mathfrak{J}_{W_\phi}(z)| \leq \mathcal{E}_\phi(nx),$$

where we recall that, for any $\epsilon > 0$,

$$\mathcal{E}_\phi(x) = e^{\mathfrak{t}_\phi x^{\circ_\phi}} \mathbb{I}_{\{0 < \mathfrak{t}_\phi < \infty\}} + e^{\epsilon x^{\circ_\phi}} \mathbb{I}_{\{\mathfrak{t}_\phi = 0\}} + e^{x^{\circ_\phi + \epsilon}} \mathbb{I}_{\{\mathfrak{t}_\phi = \infty\}}.$$

Hence, we get, that for large n and for all $x > 0$,

$$\begin{aligned} |\mathcal{P}_n(-x)| &\leq \mathcal{E}_\phi(nx) \frac{n!}{2\pi} e^{-n \ln n} \int_0^{2\pi} e^{n \cos \theta} d\theta \\ &= \mathcal{E}_\phi(nx) \frac{n! e^{-n \ln n}}{2} (\mathfrak{J}_\Gamma(n) + \mathfrak{J}_\Gamma(-n)) \\ &= O\left(\mathcal{E}_\phi(nx) e^{\frac{1}{2} \ln n}\right), \end{aligned}$$

where we have used the integral representation of the modified Bessel function

$$\mathfrak{J}_\Gamma(n) = \frac{1}{\pi} \int_0^\pi e^{n \cos \theta} d\theta,$$

see e.g. [Leb72, (5.7.4) and (5.10.8)], and, for the last inequality the bound $n! \leq e^{1-n} n^{n+\frac{1}{2}}$. Finally, from (7.28),

$$\frac{\Gamma(n+1)}{\Gamma(n-p+1)} \approx n^p,$$

$\mathfrak{t}_\phi = \mathfrak{t}_{\phi_p}$ and $\mathfrak{o}_{\phi_p} = \mathfrak{o}_\phi$ we complete (7.13) since then $\mathcal{E}_\phi(x) = \mathcal{E}_{\phi_p}(x)$. The proof of Theorem 7.0.3 is thus concluded.

Co-eigenfunctions: existence and characterization

In this chapter, we initiate our lengthy study, that will also cover the two following chapters, on the eigenfunctions of the adjoint semigroup P^* , which we recall is discussed in Theorem 1.1.1. We also name them the co-eigenfunctions of P . More specifically, we say that, for some $t > 0$ and $n \geq 0$, \mathcal{V}_n is a co-eigenfunction for P_t , or equivalently, an eigenfunction for its adjoint P_t^* in $L^2(\nu)$, associated to the eigenvalue e^{-nt} if $\mathcal{V}_n \in L^2(\nu)$ and

$$P_t^* \mathcal{V}_n = e^{-nt} \mathcal{V}_n,$$

which can be rephrased as, for any $f \in L^2(\nu)$,

$$(8.1) \quad \langle f, P_t^* \mathcal{V}_n \rangle_\nu = \langle P_t f, \mathcal{V}_n \rangle_\nu = e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu.$$

The investigation of co-eigenfunctions turns out to be more delicate than the one of eigenfunctions. Indeed, from the adjoint intertwining relation, see (8.3), they are defined, whenever they exist, as the image of the Laguerre polynomials by an unbounded from below operator, see (8.4). This forces us to develop several strategies to provide relevant information regarding the co-eigenfunctions such as existence, characterization, completeness and norm estimates. The next theorem, which is the main result of this chapter, deals with the two first issues. Before stating it, we recall, from Chapter 1, that

$$\mathcal{N}_\infty = \{ \psi \in \mathcal{N}; \sigma^2 > 0 \text{ or } \bar{\Pi}(0^+) = \infty \}, \quad \mathfrak{r} = \infty \mathbb{I}_{\{\sigma^2 > 0\}} + m + \bar{\bar{\Pi}}(0^+)$$

and $N_\mathfrak{r} = \left\lceil \frac{\bar{\Pi}(0^+)}{\mathfrak{r}} \right\rceil - 1$ with $N_\mathfrak{r} = \infty$ whenever $\bar{\Pi}(0^+) = \infty$.

THEOREM 8.0.1. *We have that \mathcal{V}_n defined for any $x > 0$ by*

$$(8.2) \quad \mathcal{V}_n(x) = \frac{\mathcal{R}^{(n)} \nu(x)}{\nu(x)} = \frac{(x^n \nu(x))^{(n)}}{n! \nu(x)} = \frac{w_n(x)}{\nu(x)}$$

is a co-eigenfunction associated to the eigenvalue e^{-nt} in each of the following cases:

(1) $\psi \in \mathcal{N}_\infty$ and $n \geq 0$, and, in this case,

$$\mathcal{V}_n \in \mathbf{C}^\infty((0, \mathfrak{r})),$$

(2) $\psi \in \mathcal{N}_\infty^c$ and $0 \leq n < \frac{\bar{\Pi}(0^+)}{2\mathfrak{r}}$.

If $\psi \in \mathcal{N}_\infty^c$, we have, for $n \leq N_\mathfrak{r}$,

$$\mathcal{V}_n \in \mathbf{C}^{N_\mathfrak{r} - n}((0, \mathfrak{r})).$$

Finally, if $N_\mathfrak{r} \geq 1$ then for any $n > \frac{\bar{\Pi}(0^+)}{2\mathfrak{r}}$,

$$\mathcal{V}_n \notin L^2(\nu).$$

REMARK 8.0.2. Note that the final claim elucidates that $\mathcal{V}_n \notin L^2(\nu)$ for all $n > \frac{\bar{\Pi}(0^+)}{2\tau}$ with $N_\tau \geq 1$. This phenomenon tells us that then the corresponding eigenvalues e^{-nt} of the adjoint semigroup P are part of the residual spectrum. It seems rare in the literature to quantify such a cut-off between the point and the residual spectrum for such a general class of operators.

REMARK 8.0.3. Note that the condition in item (2), for any $n < \frac{\bar{\Pi}(0^+)}{2\tau}$, is sharp in the sense that for the existence of co-eigenfunctions of a specific gL semigroups, see Example 3.0.4, this is a necessary and sufficient condition.

The rest of this chapter is devoted to the proof of this theorem. We start with the following claims on the adjoint intertwining relation.

PROPOSITION 8.0.4. *Let $\psi \in \mathcal{N}$ and recall that $\phi(u) = \frac{\psi(u)}{u} \in \mathcal{B}_\mathcal{N}$.*

(1) *For any $t \geq 0$ and $g \in L^2(\nu)$, we have the following intertwining relation*

$$(8.3) \quad Q_t \mathcal{I}_\phi^* g = \mathcal{I}_\phi^* P_t^* g,$$

where $\mathcal{I}_\phi^ \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ is the adjoint of \mathcal{I}_ϕ .*

(2) *$\text{Ker}(\mathcal{I}_\phi^*) = \{0\}$. However, $\text{Ran}(\mathcal{I}_\phi^*) = L^2(\varepsilon)$ if and only if $\psi(u) = Cu^2$.*

(3) *For any $n \in \mathbb{N}$, the equation*

$$(8.4) \quad \mathcal{I}_\phi^* g = \mathcal{L}_n$$

has at most one solution in $L^2(\nu)$. Moreover, if, for some $n \in \mathbb{N}$, $\mathcal{V}_n \in L^2(\nu)$ is a solution to the equation (8.4) then \mathcal{V}_n is a co-eigenfunction for P_t associated to the eigenvalue e^{-nt} .

Proof. First note that the claim $\mathcal{I}_\phi^* \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ follows readily from Theorem 7.0.1(1). Then, from the intertwining relation stated in Theorem 7.0.1(7.3), and the self-adjoint property of Q_t , we get, for any $t \geq 0$, $f \in L^2(\varepsilon)$ and $g \in L^2(\nu)$, that

$$(8.5) \quad \langle f, \mathcal{I}_\phi^* P_t^* g \rangle_\varepsilon = \langle P_t \mathcal{I}_\phi f, g \rangle_\nu = \langle \mathcal{I}_\phi Q_t f, g \rangle_\varepsilon = \langle f, Q_t \mathcal{I}_\phi g \rangle_\varepsilon,$$

which proves the identity (8.3). Next, the claim $\text{Ker}(\mathcal{I}_\phi^*) = \{0\}$ of the second item is obtained from the fact that $\overline{\text{Span}(\mathcal{P}_n)} = \overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$, see Theorem 7.0.3(2). The fact that $\text{Ran}(\mathcal{I}_\phi^*) = L^2(\varepsilon)$ if and only if $\psi(u) = Cu^2$ follows readily from Proposition 7.0.1(2). Since $\text{Ker}(\mathcal{I}_\phi^*) = \{0\}$, the equation (8.4) has at most one solution in $L^2(\nu)$. Hence, if \mathcal{V}_n is a solution of (8.4) then $\mathcal{V}_n = (\mathcal{I}_\phi^*)^{-1} \mathcal{L}_n$, where $(\mathcal{I}_\phi^*)^{-1}$ is the unbounded inverse of \mathcal{I}_ϕ^* . Therefore, since from the intertwining (8.3), we observe that $Q_t \mathcal{I}_\phi^* \mathcal{V}_n \in \text{Ran}(\mathcal{I}_\phi^*)$, we get that

$$(8.6) \quad \begin{aligned} P_t^* \mathcal{V}_n &= (\mathcal{I}_\phi^*)^{-1} Q_t \mathcal{I}_\phi^* \mathcal{V}_n = (\mathcal{I}_\phi^*)^{-1} Q_t \mathcal{L}_n \\ &= e^{-nt} (\mathcal{I}_\phi^*)^{-1} \mathcal{L}_n = e^{-nt} \mathcal{V}_n, \end{aligned}$$

which concludes the proof. \square

The description of the range of \mathcal{I}_ϕ^* and of its unbounded inverse seem to be very difficult problems. Thus, to identify and to provide conditions for the existence of co-eigenfunctions, we implement the following two-steps program. First, by considering the formal adjoint of \mathcal{I}_ϕ in $L^2(\mathbb{R}_+)$, we transfer equation (8.4) defined in $L^2(\nu)$ into a Mellin convolution equation that can be studied in $L^2(\mathbb{R}_+)$ or even in the sense of Mellin distribution, see (8.11) below for a precise statement. Then, by means of Mellin transform techniques we study (8.11) to obtain, in distributional sense, necessary and sufficient conditions for existence, uniqueness and description of its solution. In particular, we get a characterization in terms of the Rodrigues operator acting on the density of the invariant measures ν . Then, applying the variety of results on ν (smoothness, positivity, small and large asymptotic equivalents or bounds) developed in Chapter 5, we obtain (almost) necessary and sufficient conditions for the existence of a unique solution to the original equation (8.4) considered in the Hilbert space $L^2(\nu)$.

8.1. Mellin convolution equations: distributional and classical solutions

Let us start by introducing the necessary notation to formulate the distributional setting of Mellin transforms and we refer to [ML86, Chap. 11] for a concise description. We denote by $E_{p,q}$ (resp. $E'_{p,q}$), with $p < q$ reals, the linear space of infinitely differentiable functions f defined on \mathbb{R}_+ such that there exist two strictly positive numbers a and a' for which, for all $k \in \mathbb{N}$,

$$f^{(k)}(x) \stackrel{0}{=} o(x^{p+a-k-1}) \quad \text{and} \quad f^{(k)}(x) \stackrel{\infty}{=} o(x^{q-a'-k-1}),$$

(resp. the linear space of continuous linear functionals on $E_{p,q}$ endowed with a structure of a countably multinormed space as described in [ML86, p. 231]). We simply write

$$(8.7) \quad E = \cup_{q>0} E_{0,q}$$

with E' standing for the corresponding linear space of continuous linear functionals on E . We also recall, from Proposition 6.1.1(1), that $\mathbb{P}(I_\phi \in dy) = \iota(y)dy$, $y > 0$, and, $\mathcal{I}_\phi f(x) = \int_0^\infty f(xy)\iota(y)dy$ for a smooth function f .

LEMMA 8.1.1. *Let $\psi(u) = u\phi(u) \in \mathcal{N}$.*

- (1) $\mathcal{I}_\phi \in \mathbf{B}(L^2(\mathbf{p}_\alpha))$ where we recall that $\mathbf{p}_\alpha(x) = x^{-\alpha}$, $x > 0$, with $\alpha \in (0, 1)$. If $m = 0$ and $\phi'(0^+) < \infty$ then one may choose $\alpha = 0$. Moreover, denoting by $\widehat{\mathcal{I}}_\phi^\alpha$ the adjoint of \mathcal{I}_ϕ in $L^2(\mathbf{p}_\alpha)$, we have for any $f \in L^2(\mathbf{p}_\alpha)$ that, for almost every (a.e.) $x > 0$,

$$(8.8) \quad \widehat{\mathcal{I}}_\phi^\alpha f(x) = \int_0^\infty f(xy)\widehat{\iota}(y)\mathbf{p}_\alpha(y)dy,$$

where we recall that $\widehat{\iota}(y) = \iota(1/y)1/y$ with ι the density of I_ϕ . The adjoint when $\alpha = 0$ can be formally defined through the right-hand side of (8.8), in the case $m > 0$, for any $f \in L^2(\mathbb{R}_+)$ such that $\int_0^\infty |f(xy)|\widehat{\iota}(y)dy < \infty$ for a.e. $x > 0$.

(2) We have that $\mathcal{I}_\phi^* \in \mathbf{B}(\mathbf{L}^2(\nu), \mathbf{L}^2(\varepsilon))$. It is the linear operator characterized, for any $f \in \mathbf{L}^2(\nu)$ and a.e. $x > 0$, by

$$(8.9) \quad \mathcal{I}_\phi^* f(x) = \frac{1}{\varepsilon(x)} \int_0^\infty f(xy) \nu(xy) \widehat{\iota}(y) dy = \frac{1}{\varepsilon(x)} \widehat{\mathcal{I}}_\phi f \nu(x)$$

where

$$\widehat{\mathcal{I}}_\phi f(x) = \widehat{\mathcal{I}}_\phi^0 f(x) = \int_0^\infty f(xy) \widehat{\iota}(y) dy.$$

(3) Moreover, we have $\iota \in \mathbf{E}'_{0,q}$ for every $q > 0$ and $\widehat{\iota} \in \mathbf{E}'_{p,1}$, for all $p < 1$. Consequently, for any $w \in \mathbf{E}'_{0,q}$, with $q > 0$, we have that

$$(8.10) \quad \langle \mathcal{I}_\phi^\vee w, f \rangle_{\mathbf{E}'_{0,q}, \mathbf{E}_{0,q}} = \langle w, \mathcal{I}_\phi f \rangle_{\mathbf{E}'_{0,q}, \mathbf{E}_{0,q}}, \quad \forall f \in \mathbf{E}_{0,q},$$

where we have set $\mathcal{I}_\phi^\vee w = w \sqrt{\iota}$ and $w \sqrt{\iota}$ is the Mellin convolution operator in the space of distributions, see [ML86, Chapter 11.11] for definition and notation.

REMARK 8.1.2. It is worth pointing out that the left-hand side of (8.9) implies

$$\int_0^\infty |f(xy)| \nu(xy) \widehat{\iota}(y) dy < \infty$$

simply by the virtue of the fact that $\mathcal{I}_\phi \in \mathbf{B}(\mathbf{L}^2(\varepsilon), \mathbf{L}^2(\nu))$ whereas $\widehat{\mathcal{I}}_\phi$ may not belong to $\mathbf{B}(\mathbf{L}^2(\mathbb{R}_+), \mathbf{L}^2(\mathbb{R}_+))$, see item (1).

REMARK 8.1.3. Note that for $w \in \mathbf{L}^1(\widehat{\iota})$, we have the identities

$$\mathcal{I}_\phi^\vee w(x) = w \sqrt{\iota}(x) = \int_0^\infty w\left(\frac{x}{y}\right) \iota(y) \frac{dy}{y} = \int_0^\infty w(xy) \widehat{\iota}(y) dy = \widehat{\mathcal{I}}_\phi w(x),$$

which justifies the notation above.

Proof. Plainly, \mathcal{I}_ϕ is a linear operator. Next, since for any $\alpha \in (0, 1)$, $\mathcal{M}_{I_\phi}(\alpha) = \mathbb{E} \left[I_\phi^{\alpha-1} \right] < \infty$, see Proposition 6.1.2, we have, from the Cauchy-Schwarz inequality and a change of variables, that, for any $f \in \mathbf{L}^2(\mathbf{p}_\alpha)$,

$$\begin{aligned} \|\mathcal{I}_\phi f\|_{\mathbf{p}_\alpha}^2 &\leq \mathbb{E} \left[\int_0^\infty f^2(x I_\phi) \mathbf{p}_\alpha(x) dx \right] = \mathcal{M}_{I_\phi}(\alpha) \int_0^\infty f^2(x) \mathbf{p}_\alpha(x) dx \\ &= \mathcal{M}_{I_\phi}(\alpha) \|f\|_{\mathbf{p}_\alpha}^2, \end{aligned}$$

which proves $\mathcal{I}_\phi \in \mathbf{B}(\mathbf{L}^2(\mathbf{p}_\alpha))$. The latter is valid for $\alpha = 0$ whenever $\mathcal{M}_{I_\phi}(0) < \infty$ which, according to Proposition 6.1.2, is true if $m = 0$ and $\phi'(0^+) < \infty$. For any $f, g \in \mathbf{L}^2(\mathbf{p}_\alpha)$, we have, after performing a change of variables and using Fubini's Theorem, that

$$\begin{aligned} \langle \mathcal{I}_\phi f, g \rangle_{\mathbf{p}_\alpha} &= \int_0^\infty \int_0^\infty f(xy) \iota(y) dy g(x) \mathbf{p}_\alpha(x) dx \\ &= \int_0^\infty \int_0^\infty g(ry) \frac{\widehat{\iota}(y) dy}{y^{\alpha-1}} f(r) \mathbf{p}_\alpha(r) dr \\ &= \langle f, \widehat{\mathcal{I}}_\phi^\alpha g \rangle_{\mathbf{p}_\alpha}, \end{aligned}$$

which yields (8.8) and item (1) as the last claim is obvious. Since, from Proposition 7.0.1(1), $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ then $\mathcal{I}_\phi^* \in \mathbf{B}(L^2(\nu), L^2(\varepsilon))$. For any $f \in L^2(\varepsilon)$, $f \geq 0$ and $g \in L^2(\nu)$, $g \geq 0$,

$$\begin{aligned} \langle \mathcal{I}_\phi f, g \rangle_\nu &= \int_0^\infty \int_0^\infty f(xy) \iota(y) dy g(x) \nu(x) dx \\ &= \int_0^\infty f(r) \varepsilon^{-1}(r) \int_0^\infty \iota(r/x) g(x) \nu(x) / x dx \varepsilon(r) dr \\ &= \int_0^\infty f(r) \varepsilon^{-1}(r) \int_0^\infty g(rv) \nu(rv) \widehat{\iota}(v) dv \varepsilon(r) dr \\ &= \langle f, \varepsilon^{-1} \widehat{\mathcal{I}}_\phi g \nu \rangle_\varepsilon. \end{aligned}$$

Thus, we conclude (8.9). Let us consider item (3). Since the mapping

$$z \mapsto \mathcal{M}_\iota(z) = \mathcal{M}_{I_\phi}(z) = \mathcal{M}_\iota(1-z) \in \mathcal{A}_{(0,\infty)}$$

and $|\mathcal{M}_{I_\phi}(z)| \leq \mathcal{M}_{I_\phi}(\Re(z)) < \infty$, for $z \in \mathbb{C}_{(0,\infty)}$, see Proposition 6.1.2, we deduce from [ML86, Theorem 11.10.1] that $\iota \in E'_{0,q}$, for every $q > 0$ and $\widehat{\iota} \in E'_{p,1}$ for every $p < 1$. The proof of (8.10) is immediate from [ML86, 11.11.1] checking that, for any $f \in E_{0,q}$, $q > 0$,

$$\mathcal{I}_\phi f(x) = \int_0^\infty f(xy) \iota(y) dy = \langle \iota, f(x \cdot) \rangle_{E'_{0,q}, E_{0,q}},$$

and thus

$$\langle \mathcal{I}_\phi^\vee w, f \rangle_{E'_{0,q}, E_{0,q}} = \langle w \sqrt{\iota}, f \rangle_{E'_{0,q}, E_{0,q}} = \langle w, \mathcal{I}_\phi f \rangle_{E'_{0,q}, E_{0,q}},$$

where the last relation is simply [ML86, Definition 11.11.1]. \square

Recall that the Rodrigues operator is defined as $\mathcal{R}^{(n)} f(x) = \frac{1}{n!} (x^n f(x))^{(n)}$.

PROPOSITION 8.1.4. *Let $\psi \in \mathcal{N}$, and, for any $n \in \mathbb{N}$, we write*

$$\varepsilon_n(x) = \mathcal{L}_n(x) \varepsilon(x) = \mathcal{R}^{(n)} \varepsilon(x).$$

(1) *Then, for any $n \in \mathbb{N}$ the Mellin convolution equation*

$$(8.11) \quad \mathcal{I}_\phi^\vee \widehat{f}(x) = \varepsilon_n(x)$$

has a unique solution, in the sense of distributions, given by

$$w_n(x) = \mathcal{R}^{(n)} \nu(x) = \frac{1}{n!} (x^n \nu(x))^{(n)} \in E'.$$

Its Mellin transform is given, for any $z \in \mathbb{C}_{(0,\infty)}$, by

$$(8.12) \quad \mathcal{M}_{w_n}(z) = \frac{(-1)^n}{\Gamma(n+1)} \frac{\Gamma(z)}{\Gamma(z-n)} \mathcal{M}_{V_\psi}(z).$$

(2) *In fact, writing $\mathbf{p}_\alpha(x) = x^{-\alpha}$, $x > 0$, with $-\infty < \alpha < 1 - 2d_\phi$,*

(a) *if $\psi \in \mathcal{N}_\infty$ then, for any $n \geq 0$, $w_n \in L^2(\mathbf{p}_\alpha) \cap \mathcal{C}_0^\infty(\mathbb{R}_+)$, and,*

(b) if $\psi \in \mathcal{N}_\infty^c$ we have the following cases.

If $n \in N_2 = \left\{ p \in \mathbb{N}; \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - \frac{1}{2} > 0 \text{ and } 0 \leq p < \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - \frac{1}{2} \right\}$ then

$$(8.13) \quad w_n \in L^2(\mathbf{p}_\alpha),$$

whereas,

if $n \in \overline{N}_2 = \left\{ p \in \mathbb{N}; \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - \frac{1}{2} < 0 \text{ and } p = 0, \text{ or, } p > \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - \frac{1}{2} \right\}$ then

$$w_n \notin L^2(\mathbf{p}_\alpha).$$

Moreover, for any $n \leq N_\mathfrak{r}$, $w_n \in \mathbf{C}^{N_\mathfrak{r}-n}((0, \mathfrak{r}))$ and if in addition $N_\mathfrak{r} \geq 1$ then $w_n \in \mathbf{C}_0^{N_\mathfrak{r}-n-1}(\mathbb{R}_+)$.

(3) Finally, if $N_\mathfrak{r} > 2$ we have, for any $n < N_\mathfrak{r}$ and $x > 0$, that

$$(8.14) \quad \widehat{\mathcal{I}}_\phi w_n(x) = \int_0^\infty w_n(xy) \widehat{l}(y) dy = \varepsilon_n(x).$$

REMARK 8.1.5. It is possible that $n = \frac{\overline{\Pi}(0^+)}{\mathfrak{r}} - \frac{1}{2}$ in item (2). Then whether $w_n \in L^2(\mathbb{R}_+)$ or not can be extracted from the behaviour of the slowly varying function l that appears in Theorem 5.0.2(5.4), and, from [SY78, (1.9)], which can in turn be expressed in terms of the excursion measure \mathbf{n} defined in (5.26). Clearly, though, in any case, $w_n \notin L^2(\mathbb{R}_+)$, if $n \geq N_\mathfrak{r} + 1$.

Proof. Let $\psi \in \mathcal{N}$. Recall from Lemma 8.1.1(3) that $\mathcal{I}_\phi^\vee w = w\sqrt{l}$. Take $w \in E'_{0,q}$, $q > 0$, and, with $z \in \mathbb{C}_{(0,q)}$, note that $p_z(x) = x^z \in E_{0,q}$. Then, we have that

$$\mathcal{M}_{\mathcal{I}_\phi^\vee w}(z) = \langle w\sqrt{l}, p_{z-1} \rangle_{E'_{0,q}, E_{0,q}} = \langle w, \mathcal{I}_\phi p_{z-1} \rangle_{E'_{0,q}, E_{0,q}} = \mathcal{M}_{\mathcal{I}_\phi}(z) \mathcal{M}_w(z),$$

where we have used that

$$\mathcal{I}_\phi p_{z-1}(x) = p_{z-1}(x) \mathcal{M}_{\mathcal{I}_\phi}(z),$$

see (7.19). However, since, for any $n \in \mathbb{N}$, $\mathcal{L}_n(x) = \frac{\mathcal{R}^{(n)} \varepsilon(x)}{\varepsilon(x)}$, see (3.3), that is $\varepsilon_n(x) = \mathcal{R}^{(n)} \varepsilon(x)$ we get, from [ML86, 11.7.7], that

$$\mathcal{M}_{\varepsilon_n}(z) = \frac{(-1)^n}{n!} \frac{\Gamma(z)}{\Gamma(z-n)} \Gamma(z).$$

Putting pieces together, we get that the Mellin transform of a solution to (8.11) takes the form

$$(8.15) \quad \mathcal{M}_{\widehat{f}}(z) = \frac{(-1)^n}{n!} \frac{\Gamma(z)}{\Gamma(z-n)} \frac{\Gamma(z)}{\mathcal{M}_{\mathcal{I}_\phi}(z)} = \frac{(-1)^n}{\Gamma(n+1)} \frac{\Gamma(z)}{\Gamma(z-n)} \mathcal{M}_{V_\psi}(z),$$

where for the last identity we have used from Proposition 6.1.2 that $\mathcal{M}_{\mathcal{I}_\phi}(z) = \frac{\Gamma(z)}{W_\phi(z)}$, and, from (6.4) and (6.5), that $\mathcal{M}_{V_\psi}(z) = W_\phi(z)$. Next, since, from Theorem 6.0.1(3), we have that, for $z \in \mathbb{C}_{(0,\infty)}$, $z \mapsto \mathcal{M}_{V_\psi}(z)$ is analytical with

$$|\mathcal{M}_{V_\psi}(z)| \leq \mathcal{M}_{V_\psi}(\Re(z)) < \infty,$$

we deduce, from [ML86, Theorem 11.10.1] that $\nu \in E'_{0,q}$, for any $q > 0$. Hence, by means of [ML86, 11.7.7], we have that with $\hat{f} = w_n = \mathcal{R}^{(n)}\nu$, $\hat{f} \in E'_{0,q}$ and \hat{f} is a solution to (8.11). The uniqueness of the solution and thus (8.12) follow from the uniqueness of Mellin transforms in the distributional sense. For item (2), from the expression of \mathcal{M}_{w_n} in (8.12), we first observe, from Theorem 6.0.1(3) again, that for all $n \geq 0$, $\mathcal{M}_{w_n} \in \mathcal{A}_{(0,\infty)}$. Then, the classical estimates of the gamma function (9.7) combined with the asymptotic behaviour of \mathcal{M}_{V_ψ} in (5.3), yield for any $\underline{\alpha} < 1 - 2d_\phi$ and any $u \leq \max(N_\tau - 1, 0)$ that, for large $|b|$,

$$\left| \mathcal{M}_{w_n} \left(\frac{1-\underline{\alpha}}{2} + ib \right) \right| = O \left(|b|^n \left| \mathcal{M}_{V_\psi} \left(\frac{1-\underline{\alpha}}{2} + ib \right) \right| \right) = o(|b|^{n-u}).$$

Since $N_\tau = \infty$ when $\psi \in \mathcal{N}_\infty$, the proof, in this case, follows easily from this estimate by appealing, for the $\mathbf{C}_0^\infty(\mathbb{R}_+)$ property to a standard Mellin inversion argument, explained in (1.36), and, for the $L^2(\mathbf{p}_\alpha)$ property, to the Parseval identity (1.39), since for all $n \in \mathbb{N}$, recalling that $p_a(x) = x^a$,

$$b \mapsto \left| \mathcal{M}_{p_{-\frac{\alpha}{2}} w_n} \left(\frac{1}{2} + ib \right) \right| = \left| \mathcal{M}_{w_n} \left(\frac{1-\alpha}{2} + ib \right) \right| \in L^2(\mathbb{R})$$

and hence $x \mapsto x^{-\frac{\alpha}{2}} w_n(x) \in L^2(\mathbb{R}_+)$, that is $w_n \in L^2(\mathbf{p}_\alpha)$. For $\psi \in \mathcal{N}_\infty^c$ it can be checked immediately from [PS18a, Theorem 5.2] that, if $n < \frac{\overline{\Pi}(0^+)}{\tau} - \frac{1}{2}$,

$$b \mapsto \left| b^n \mathcal{M}_{V_\psi} \left(\frac{1-\alpha}{2} + ib \right) \right| \in L^2(\mathbb{R})$$

and fails to be square integrable provided $n > \frac{\overline{\Pi}(0^+)}{\tau} - \frac{1}{2}$. Therefore, the Parseval identity (1.39) gives the same condition to distinguish whether $p_{-\frac{\alpha}{2}} w_n$ belongs to $L^2(\mathbb{R}_+)$ or not. Next recalling, from Theorem 5.0.2(2), that $\nu \in \mathbf{C}^{N_\tau}((0, \tau))$, and, for $N_\tau \geq 1$, $\nu \in \mathbf{C}_0^{N_\tau-1}(\mathbb{R}_+)$, we get the stated equivalent properties for w_n and complete the proof of the item (2). We proceed with item (3) and thus assume that $N_\tau > 2$. To prove (8.14) we first note that

$$\int_0^\infty |w_n(xy)| \widehat{\iota}(y) dy = \int_0^\infty |w_n(xy)| \frac{1}{y} \iota \left(\frac{1}{y} \right) dy.$$

Then (5.46) yields that for any $n < N_\tau - 2$ and any $\bar{a} > d_\phi$ there exists a constant $C = C_{n,\bar{a}} > 0$ such that for any $x > 0$

$$\left| (x^n \nu(x))^{(n)} \right| \leq C x^{-\bar{a}}.$$

Then we get for an such \bar{a} that

$$\begin{aligned} \widehat{\mathcal{I}}_\phi |w_n|(x) &\leq \int_0^\infty |w_n(xy)| \frac{1}{y} \iota \left(\frac{1}{y} \right) dy \\ &\leq \frac{C_{n,\bar{a}}}{x^{\bar{a}}} \int_0^\infty y^{\bar{a}-1} \frac{1}{y^2} \iota \left(\frac{1}{y} \right) dy \\ &= \frac{C_{n,\bar{a}}}{x^{\bar{a}}} \mathcal{M}_{I_\phi}(\bar{a}) < \infty, \end{aligned}$$

since $\mathcal{M}_{I_\phi} \in \mathcal{A}_{(0,\infty)}$ in any case, see Proposition 6.1.2. Hence, $\widehat{\mathcal{I}}_\phi w_n$ is well defined in the sense of (8.8) with $\alpha = 0$ and it is a proper product convolution. Clearly,

then, for any $n < N_{\mathfrak{r}} - 2$,

$$\mathcal{M}_{\widehat{\mathcal{I}}_{\phi} w_n}(z) = \mathcal{M}_{w_n}(z) \mathcal{M}_{I_{\phi}}(z) = \mathcal{M}_{\varepsilon_n}(z),$$

where the latter follows precisely as in (8.15) above. \square

8.2. Existence of co-eigenfunctions: Proof of Theorem 8.0.1

From Proposition 8.1.4(2) and $\nu > 0$ on $(0, \mathfrak{r})$, Theorem 5.0.2(1), we conclude from the definition of $\mathcal{V}_n = w_n/\nu$, see (8.2), that $\mathcal{V}_n \in \mathbb{C}^{N_{\mathfrak{r}}-n}((0, \mathfrak{r}))$. Assume for a moment that for some $n \in \mathbb{N}$ we have that $\mathcal{V}_n \in L^2(\nu)$. Then from (8.9) $\mathcal{I}_{\phi}^* \mathcal{V}_n = \frac{1}{\varepsilon} \widehat{\mathcal{I}}_{\phi} w_n$ and $\widehat{\mathcal{I}}_{\phi} |w_n| = \int_0^{\infty} |w_n(xy)| \hat{i}(y) dy < \infty$. Hence, (8.14) holds and yields that

$$\mathcal{I}_{\phi}^* \mathcal{V}_n = \frac{1}{\varepsilon} \widehat{\mathcal{I}}_{\phi} w_n = \frac{\varepsilon_n}{\varepsilon} = \mathcal{L}_n.$$

Finally, (8.4) implies that \mathcal{V}_n is a co-eigenfunction of P_t associated to the eigenvalue e^{-nt} . Therefore, it remains to show that $\mathcal{V}_n \in L^2(\nu)$. This is a difficult question which requires, in particular, several of the very precise estimates developed in Chapter 5. We now observe that, for some $n \in \mathbb{N}$, $\mathcal{V}_n \in L^2(\nu)$ if and only if the function

$$(8.16) \quad F_n(x) = (n!)^2 \mathcal{V}_n^2(x) \nu(x) = \frac{\left((x^n \nu(x))^{(n)} \right)^2}{\nu(x)} = \frac{w_n^2(x)}{\nu(x)} \in L^1((0, \mathfrak{r})).$$

Since $\nu > 0$ on $(0, \mathfrak{r})$ with $\nu \in \mathbb{C}^{N_{\mathfrak{r}}}(0, \mathfrak{r})$, see Theorem 5.0.2(1), for $n \leq N_{\mathfrak{r}}$, it suffices to check the integrability of F_n in neighbourhoods of 0 and \mathfrak{r} . In the case $1 \leq N_{\mathfrak{r}} < \infty$, for $n > N_{\mathfrak{r}}$, we shall show by other means that $\mathcal{V}_n \notin L^2(\nu)$.

We state the following result after recalling, from the bound (5.5) stated in Theorem 5.0.4, that for any $\psi \in \mathcal{N}$ and any $\underline{a} < d_{\phi}$, $\underline{A} \in (0, \mathfrak{r})$, there exists $C_{\underline{a}, \underline{A}} > 0$ such that for all $x \in (0, \underline{A})$

$$(8.17) \quad \nu(x) \geq C_{\underline{a}, \underline{A}} x^{-\underline{a}}.$$

LEMMA 8.2.1. *For $\psi \in \mathcal{N}_{\infty}$ and $n \in \mathbb{N}$ or $\psi \in \mathcal{N}_{\infty}^c$ and $n \in N_2$, see (8.13), we have that $x \mapsto F_n(x) \mathbb{I}_{\{x < A\}}$ is integrable on \mathbb{R}_+ for any $A \in (0, \mathfrak{r})$. However, if $\psi \in \mathcal{N}_{\infty}^c$ and $x \mapsto x w_n(x) \mathbb{I}_{\{x < A\}} \notin L^2(\mathbb{R}_+)$ for some $A \in (0, \mathfrak{r})$ then $F_n(x) \mathbb{I}_{\{x < A\}}$ is not integrable on \mathbb{R}_+ .*

Proof. From (5.12) we have that $\nu(x) = x^{-2} \widehat{\nu}_1(x^{-1})$ with $\widehat{\nu}_1$ the density of a self-decomposable random variable which is known to be unimodal and hence $\widehat{\nu}_1$ is ultimately non-increasing, see [SY78, Theorem 1.1]. Therefore, for any $x < A < \mathfrak{r}$ we have that $\nu(x) \leq \underline{C}^{-1} x^{-2}$, for some $\underline{C} = \underline{C}(A) > 0$. This bound and the one recalled in (8.17), yield that, for any $n \in \mathbb{N}$, for any $\underline{a} < d_{\phi}$ and $x \in (0, A)$,

$$(8.18) \quad \underline{C} x^2 w_n^2(x) \mathbb{I}_{\{x < A\}} \leq F_n(x) \mathbb{I}_{\{x < A\}} = \frac{w_n^2(x)}{\nu(x)} \mathbb{I}_{\{x < A\}} \leq C_{\underline{a}, \underline{A}} x^{\underline{a}} w_n^2(x).$$

The upper bound with $\underline{a} \in (2d_{\phi} - 1, d_{\phi})$ together with Proposition 8.1.4(2) gives the first claim for the integrability on \mathbb{R}_+ of $x \mapsto F_n(x) \mathbb{I}_{\{x < A\}}$, for any $A \in (0, \mathfrak{r})$. The lower bound proves the second claim for the non-integrability. \square

Upon the classes considered, the study of the integrability property of F_n at \mathfrak{r} requires different analysis that we split into several parts.

8.3. The case $\psi \in \mathcal{N}_{\infty, \infty}$.

We start by investigating the case when $\psi \in \mathcal{N}_{\infty, \infty}$, that is when $N_{\mathfrak{r}} = \mathfrak{r} = \infty$, and, from Theorem 5.0.5, we obtain the following.

LEMMA 8.3.1. *Let $\psi \in \mathcal{N}_{\infty, \infty}$. Then, for any $n \in \mathbb{N}$, there exists $C_\psi > 0$ such that*

$$(8.19) \quad (x^n \nu(x))^{(n)} \approx (-1)^n \frac{C_\psi}{\sqrt{2\pi}} \sqrt{\varphi'(x)} \varphi^n(x) e^{-\int_m^x \varphi(y) \frac{dy}{y}},$$

where $\varphi(\phi(x)) = x$. Consequently, for any $n \in \mathbb{N}$, and $A > 0$, $F_n(x) \mathbb{I}_{\{x > A\}}$ is integrable on \mathbb{R}_+ .

Proof. Recall that (5.7) claims that, for any $n \in \mathbb{N}$,

$$(8.20) \quad \nu^{(n)}(x) \approx (-1)^{-n} \frac{C_\psi}{\sqrt{2\pi}} \sqrt{\varphi'(x)} \frac{\varphi^n(x)}{x^n} e^{-\int_m^x \varphi(y) \frac{dy}{y}},$$

from which we deduce that

$$(x^n \nu(x))^{(n)} \approx \frac{C_\psi}{\sqrt{2\pi}} e^{-\int_m^x \varphi(y) \frac{dy}{y}} \sqrt{\varphi'(x)} \sum_{k=0}^n \binom{n}{k} \frac{\Gamma(n+1)}{\Gamma(n-k+1)} (-1)^{n-k} \varphi^{n-k}(x).$$

Since $\lim_{x \rightarrow \infty} \varphi(x) = \infty$, hence $\varphi^n(x)$ is the leading term and (8.19) holds. From the third equality in (8.16), using the asymptotic relation (8.20) for $n = 0$ and twice (8.19), we get that

$$\begin{aligned} F_n(x) &\approx \frac{\left((x^n \nu(x))^{(n)} \right)^2}{\nu(x)} \\ &\approx \frac{C_\psi}{\sqrt{2\pi}} \sqrt{\varphi'(x)} \varphi^{2n}(x) e^{-\int_m^x \varphi(y) \frac{dy}{y}} \approx (x^{2n} \nu(x))^{(2n)}. \end{aligned}$$

Finally, from the upper bound (5.46), we get, with $k = 2n$ and $\bar{a} > 1$ therein, that for any $x > 0$, there exists a constant $C = C_{2n, \bar{a}} > 0$ such that

$$\left| (x^{2n} \nu(x))^{(2n)} \right| \leq C x^{-\bar{a}},$$

which shows that $F_n(x) \mathbb{I}_{\{x > A\}}$ is integrable on \mathbb{R}_+ and completes the proof of the lemma. \square

Lemma 8.2.1 and Lemma 8.3.1 complete the proof of Theorem 8.0.1(1) when $\psi \in \mathcal{N}_{\infty, \infty}$. We proceed with the remaining case of 8.0.1(1), that is $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$.

8.4. The case $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$

We now discuss the case $\mathfrak{r} < \infty$ and $N_\mathfrak{r} = \infty$. Expanding the third expression in (8.16), to obtain the integrability of F_n in a neighbourhood of $\mathfrak{r} < \infty$ it suffices to show, for any $0 \leq i + j < \infty$, that

$$(8.21) \quad \lim_{x \uparrow \mathfrak{r}} \frac{\nu^{(j)}(x)\nu^{(i)}(x)}{\nu(x)} = 0.$$

To this end we introduce the linear space

$$\mathbf{C}_\mathfrak{r}^\infty(\mathbb{R}_+) = \{f \in \mathbf{C}^\infty(\mathbb{R}_+); \forall l \in \mathbb{N}, r_{f^{(l)}} < \mathfrak{r} \text{ and } \text{Supp} f = [0, \mathfrak{r}]\},$$

where

$$(8.22) \quad r_f = \sup \{0 \leq x < \mathfrak{r}; f(x) = 0\} \in [0, \mathfrak{r}],$$

and prove the following.

LEMMA 8.4.1. *If $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$ then $\nu \in \mathbf{C}_\mathfrak{r}^\infty(\mathbb{R}_+)$.*

Proof. Since from Theorem 5.0.2(2a), we have $\nu \in \mathbf{C}^\infty(\mathbb{R}^+)$, we simply need to check that, for all $n \in \mathbb{N}$, $r_{\nu^{(n)}} < \mathfrak{r}$. From (5.84) and the middle identity of (5.85) on $(0, \mathfrak{r})$, we get, recalling that $\tilde{k}_n = \sum_{j=1}^n j k_j$ and $\bar{k}_n = \sum_{j=1}^n k_j$, that, for all $n \in \mathbb{N}$,

$$(8.23) \quad \begin{aligned} \nu^{(n)}(x) &= \left(\frac{1}{x^2} \widehat{\nu}_1 \left(\frac{1}{x} \right) \right)^{(n)} \\ &= (-1)^n \sum_{m=0}^n \binom{n}{m} \frac{(m+1)!}{x^{2+m}} \sum_{\substack{\bar{k}_n = n-m \\ \tilde{k}_n = k}} \frac{(n-m)! \widehat{\nu}_1^{(k)}(x^{-1}) x^{-n+m-k}}{\prod_{j=1}^{n-m} k_j!}, \end{aligned}$$

with $\widehat{\nu}_1$ the density of $I_{\mathcal{T}_1} \psi$ and $\mathcal{T}_1 \psi(u) = u\phi(u+1) \in \mathcal{N}(\phi(1))$ has associated Lévy measure Π_1 . From Proposition 5.1.1(2), $\bar{\Pi}_1(0^+) = \bar{\Pi}(0^+) = \infty$, and clearly

$$\mathfrak{r}_1 = \lim_{u \rightarrow \infty} \phi(u+1) = \phi(\infty) = \mathfrak{r}.$$

Hence

$$\left\lceil \frac{\bar{\Pi}_1(0^+)}{\mathfrak{r}_1} \right\rceil - 1 = \infty.$$

From Lemma 5.3.1(1), for each $n \in \mathbb{N}$, there exists $a_n^{(l)} > \frac{1}{\mathfrak{r}}$ such that, for all $0 \leq l \leq n$, $\widehat{\nu}_1^{(l)} > 0$ on $(\frac{1}{\mathfrak{r}}, a_n^{(l)})$. Thus, from (8.23), $r_{\nu^{(n)}} < \mathfrak{r}$ for all $n \in \mathbb{N}$, which completes the proof. \square

LEMMA 8.4.2. *If $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$ then for any $i, j \in \mathbb{N}$, we have*

$$(8.24) \quad \lim_{x \uparrow \mathfrak{r}} \frac{\nu^{(j)}(x)\nu^{(i)}(x)}{\nu(x)} = 0.$$

Proof. We work by induction in $i + j$ for the whole class $\mathbf{C}_\mathfrak{r}^\infty(\mathbb{R}_+)$. When $i + j = 0$, for any $f \in \mathbf{C}_\mathfrak{r}^\infty(\mathbb{R}_+)$ (8.24) reduces to $\lim_{x \uparrow \mathfrak{r}} f(x) = 0$ which holds since by definition $\mathbf{C}_\mathfrak{r}^\infty(\mathbb{R}_+) \subseteq \mathbf{C}^\infty(\mathbb{R}_+)$ and $\text{Supp} f = [0, \mathfrak{r}]$. Assume (8.24) is valid

for some $n = i + j$ and all $f \in \mathcal{C}_\tau^\infty(\mathbb{R}_+)$. Consider $i + j = n + 1$. Then, by the L'Hôpital's rule

$$(8.25) \quad \lim_{x \uparrow \tau} \frac{f^{(i)}(x)f^{(j)}(x)}{f(x)} = \lim_{x \uparrow \tau} \left(\frac{g^{(i)}(x)g^{(j-1)}(x)}{g(x)} + \frac{g^{(i-1)}(x)g^{(j)}(x)}{g(x)} \right),$$

where $g = f'$. Since $g \in \mathcal{C}_\tau^\infty(\mathbb{R}_+)$ we see by the induction hypothesis that the limit to the right vanishes. This proves (8.24) as $\nu \in \mathcal{C}_\tau^\infty(\mathbb{R}_+)$ according to Lemma 8.4.1. \square

Lemmas 8.2.1, 8.4.1 and 8.4.2 conclude Theorem 8.0.1(1) in the remaining case, i.e. $\psi \in \mathcal{N}_\infty \setminus \mathcal{N}_{\infty, \infty}$. We proceed with Theorem 8.0.1(2) and the final claim of Theorem 8.0.1.

8.5. The case $\psi \in \mathcal{N}_\infty^c$.

In this case we have that $\tau < \infty$ and $N_\tau < \infty$.

LEMMA 8.5.1. *Let $\psi \in \mathcal{N}_\infty^c$. Then, for any $A \in (0, \tau)$, $F_n(x)\mathbb{I}_{\{x>A\}}$ is integrable (resp. not integrable) on \mathbb{R}_+ if $0 \leq n < \frac{\bar{\Pi}(0^+)}{2\tau}$ (resp. $\frac{\bar{\Pi}(0^+)}{2\tau} < n \leq N_\tau$, $N_\tau \geq 1$).*

REMARK 8.5.2. The case $\frac{\bar{\Pi}(0^+)}{2\tau} = n$ depends on the slowly varying function appearing in (8.26).

Proof. From Theorem 5.0.2(2c), we deduce, after expanding the right-hand side of (8.16) that, for any $n = 1, \dots, N_\tau$,

$$(8.26) \quad F_n(x) \underset{\tau}{\sim} C_\tau \frac{(\nu^{(n)}(x))^2}{\nu(x)} \mathbb{I}_{\{x < \tau\}} \underset{\tau}{\sim} C(\tau - x)^{\frac{\bar{\Pi}(0^+)}{\tau} - 2n - 1} l(\tau - x) \mathbb{I}_{\{x < \tau\}},$$

where $C_\tau, C > 0$ and l a slowly varying function at 0. This completes the proof. \square

Note that always $\mathcal{V}_0 \equiv 1 \in L^2(\nu)$. Next, we assume first that

$$N_\tau = \lceil \frac{\bar{\Pi}(0^+)}{\tau} \rceil - 1 \geq 1$$

which implies that if $n \in \mathbb{N}$ and $n < \frac{\bar{\Pi}(0^+)}{2\tau}$ then $n \in N_2$, see (8.13). Then, Lemmas 8.2.1 and 8.5.1 give that $\mathcal{V}_n \in L^2(\nu)$ for any $0 \leq n < \frac{\bar{\Pi}(0^+)}{2\tau}$. Also, Lemma 8.5.1 shows that $\mathcal{V}_n \notin L^2(\nu)$ if $\frac{\bar{\Pi}(0^+)}{2\tau} < n \leq N_\tau$, $N_\tau \geq 1$. Assume from now on that

$$n \geq N_\tau + 1 > \frac{\bar{\Pi}(0^+)}{\tau} - \frac{1}{2}$$

and $N_\tau \geq 0$. Then necessarily, from Proposition 8.1.4(2) we have $n \in \bar{N}_2$ and hence $w_n \notin L^2(\mathbb{R}_+)$. The latter may be due to either w_n^2 is not measurable and in this case the proof is completed since from (8.2), $\mathcal{V}_n^2 = \frac{w_n^2}{\nu}$, or the measurable function w_n^2 is not integrable on an open set $E \subseteq (0, \tau)$ with the possible alternatives

- i) $0, \tau \notin E$ or/and
- ii) w_n^2 is not integrable at τ or/and
- iii) w_n^2 is not integrable at 0.

If i) holds then, since $\nu > 0$ and continuous on $(0, \tau)$, see Theorem 5.0.2(1), then plainly $w_n \notin L^2(\mathbb{R}_+)$ implies that $\mathcal{V}_n \notin L^2(\nu)$. If ii) holds, then since from Theorem 5.0.2(2c)

$$\nu(x) \stackrel{\tau}{\sim} C(\tau - x)^{\frac{\bar{\pi}(0^+)}{\tau} - 1} l(\tau - x) \mathbb{I}_{\{x < \tau\}},$$

where $C > 0$ is throughout generic and l a slowly varying function at 0. Then, assuming first that $N_\tau = \lceil \frac{\bar{\pi}(0^+)}{\tau} \rceil - 1 \geq 1$ leads to $\frac{\bar{\pi}(0^+)}{\tau} > 1$ and hence there exists $C, A > 0$ such that

$$\|\mathcal{V}_n\|_\nu^2 \geq \int_0^A \frac{w_n^2(x)}{\nu(x)} dx + C \int_A^\tau w_n^2(x) dx \geq C \int_A^\tau w_n^2(x) dx = \infty,$$

which gives the statement in the case $\mathcal{N}_\tau \geq 1$.

The case iii) is dealt with as follows. Since $w_n \notin L^2(\mathbb{R}_+)$ then from Proposition 8.1.4(2) we have that even $x \mapsto xw_n(x) \notin L^2(\mathbb{R}_+)$ and if $x \mapsto x^2w_n^2(x)$ fails to be integrable at 0 then Lemma 8.2.1 furnishes that $\mathcal{V}_n \notin L^2(\nu)$. Otherwise, clearly, at least case i) and/or case ii) must be valid for w_n and hence $\mathcal{V}_n \notin L^2(\nu)$. This, completes the proof of Theorem 8.0.1(2). Next, we discuss the final claim of Theorem 8.0.1. The claim $\mathcal{V}_n \in \mathcal{C}^{N_\tau - n}((0, \tau))$, $n \leq N_\tau$, follows from Proposition 8.1.4(2), (8.2) and $\nu > 0$ on $(0, \tau)$, see Theorem 5.0.2(1).

We close this chapter with the following additional properties of the adjoint intertwining operator.

LEMMA 8.5.3. *Let $\psi \in \mathcal{N}_\infty$.*

- (1) *Then $\overline{\text{Ran}}(\mathcal{I}_\phi^*) = L^2(\varepsilon)$ and $\text{Ker}(\mathcal{I}_\phi) = \{0\}$.*
- (2) *For any $\alpha \in (0, 1)$ we have that $\overline{\text{Ran}}(\mathcal{I}_\phi)_{L^2(\mathbf{p}_\alpha)} = L^2(\mathbf{p}_\alpha)$ whereas when $\psi \in \mathcal{N}_\infty$ with $m = 0$ and $\phi'(0^+) < \infty$, this extends to $\alpha = 0$, i.e. to $L^2(\mathbb{R}_+)$.*

Proof. From Theorem 8.0.1(1) and (8.4) we have, for any $\psi \in \mathcal{N}_\infty$ and $n \in \mathbb{N}$, $\mathcal{V}_n \in L^2(\nu)$ and $\mathcal{I}_\phi^* \mathcal{V}_n = \mathcal{L}_n$. Hence, $\overline{\text{Span}}(\mathcal{L}_n) = L^2(\varepsilon)$ entails that $\overline{\text{Ran}}(\mathcal{I}_\phi^*) = L^2(\varepsilon)$ and

$$\text{Ker}(\mathcal{I}_\phi) = \overline{\text{Ran}}(\mathcal{I}_\phi^*)^\perp = \{0\}.$$

This proves item (1). Next, let us consider, for any $\beta \geq 0$, the power series

$$\begin{aligned} f_\beta(x) &= x^\beta \sum_{n=0}^{\infty} \frac{W_\phi(n+1+\beta)}{\Gamma(n+1+\beta)n!} (-1)^n x^n \\ &= x^\beta \bar{f}_\beta(x). \end{aligned}$$

From (2.20) and Proposition 4.1.1(3), we get that

$$\lim_{n \rightarrow \infty} \frac{W_\phi(n+1+\beta)}{W_\phi(n+\beta)n(n+\beta)} = \lim_{n \rightarrow \infty} \frac{\phi(n+\beta)}{n(n+\beta)} = 0.$$

Thus, \bar{f}_β is an entire function. Clearly, from Theorem 6.0.1(2) $W_\phi \in \mathcal{A}_{(0, \infty)}$ and hence

$$(8.27) \quad \mathcal{M}_\beta(z) = \Gamma(z) \frac{W_\phi(-z+1+\beta)}{\Gamma(-z+1+\beta)} \in \mathcal{A}_{(0, 1+\beta)}.$$

To show that \mathcal{M}_β is the Mellin transform of \bar{f}_β we proceed as follows. First, for any $n \in \mathbb{N}$, $b \in \mathbb{R}$, the recurrence relation of the gamma function yields

$$\left| \frac{\Gamma(-n - \frac{1}{2} + ib)}{\Gamma(n + \frac{3}{2} + \beta - ib)} \right| = \frac{1}{\prod_{j=0}^{2n} |-n - \frac{1}{2} + ib + j|} \left| \frac{\Gamma(n + \frac{1}{2} + ib)}{\Gamma(n + \frac{3}{2} + \beta + ib)} \right|.$$

Then, from the latter and [PK01, (2.1.16)] applied with $z = n + \frac{1}{2} + ib$ we get uniformly for $b \in \mathbb{R}$

$$\begin{aligned} \left| \frac{\Gamma(-n - \frac{1}{2} + ib)}{\Gamma(n + \frac{3}{2} + \beta - ib)} \right| &\leq \frac{1}{\prod_{j=0}^{2n} |-n - \frac{1}{2} + ib + j|} \frac{1}{|n + \frac{1}{2} + ib|^{1+\beta}} \\ &\leq \frac{\Gamma^2(\frac{1}{2})}{\Gamma^2(n - \frac{1}{2})} \frac{1}{|n + \frac{1}{2} + ib|^{1+\beta} |n - \frac{1}{2} + ib|}. \end{aligned}$$

Moreover, from (6.5), W_ϕ is a Mellin transform of a nonnegative random variable and therefore

$$\left| W_\phi\left(n + \frac{3}{2} + \beta + ib\right) \right| \leq W_\phi\left(n + \frac{3}{2} + \beta\right).$$

From the estimates above and the fact that \mathcal{M}_β in (8.27) extends to a meromorphic function on $\mathbb{C}_{(-\infty, 0]}$, we obtain

$$(8.28) \quad \left| \mathcal{M}_\beta\left(-n - \frac{1}{2} + ib\right) \right| \leq C \frac{W_\phi\left(n + \frac{3}{2} + \beta\right)}{\Gamma^2\left(n - \frac{1}{2}\right)} \frac{1}{|n + \frac{1}{2} + ib|^{1+\beta} |n - \frac{1}{2} + ib|}.$$

Since $\beta \geq 0$ the right-hand side of (8.28) is integrable along any line $-n - \frac{1}{2} + ib$, $n \in \mathbb{N}$. Inspection of the steps above shows that similar integrable bound for $|\mathcal{M}_\beta|$, whose proof is based directly on the decay of W_ϕ since $N_\tau = \infty > 2$, see (5.3), holds even with $n = -1$. Therefore, by Mellin inversion, a shift of contour and an application of the Cauchy Theorem, we get, for any $M \in \mathbb{N}$, that for some function $h : \mathbb{R}_+ \mapsto \mathbb{R}$

$$\begin{aligned} h(x) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} x^{-\frac{1}{2}-ib} \mathcal{M}_\beta\left(\frac{1}{2} + ib\right) db \\ (8.29) \quad &= \sum_{n=0}^M \frac{W_\phi(n+1+\beta)}{\Gamma(n+1+\beta)n!} (-1)^n x^n \\ &\quad + \frac{1}{2\pi} \int_{-\infty}^{\infty} x^{M+\frac{1}{2}-ib} \mathcal{M}_\beta\left(-M - \frac{1}{2} + ib\right) db. \end{aligned}$$

To conclude that $h = \bar{f}_\beta$, see (8.27), it remains to show that the very last term vanishes as $M \rightarrow \infty$. First, we invoke (5.2) and Proposition 4.1.1(3) to check that for large M ,

$$W_\phi\left(M + \frac{3}{2} + \beta\right) \lesssim M^{1/2} e^{M(\ln M + \ln \sigma^2 + o(1))}.$$

Second, the Stirling asymptotic (9.8) yields

$$\Gamma^2\left(M + \frac{1}{2}\right) \gtrsim e^{\frac{3}{2}M \ln(M)}.$$

Hence, for any fixed $x > 0$

$$\lim_{M \rightarrow \infty} x^M \frac{W_\phi\left(M + \frac{3}{2} + \beta\right)}{\Gamma^2\left(M + \frac{1}{2}\right)} \leq \lim_{M \rightarrow \infty} M^{\frac{1}{2}} x^M e^{M(\ln M + \ln \sigma^2 + o(1)) - \frac{3}{2}M \ln(M)} = 0$$

and using (8.28) we get that

$$\left| \mathcal{M}_\beta \left(-n - \frac{1}{2} + ib \right) \right| \lesssim \frac{1}{|b|^{2+\beta}} e^{-\frac{1}{3}n \ln(n)}$$

and hence the integral in (8.30) goes to zero as $M \rightarrow \infty$. Thus, \mathcal{M}_β is the Mellin transform of \bar{f}_β and

$$\mathcal{M}_\beta \left(\frac{1}{2} + ib \right) = \frac{\Gamma(\frac{1}{2} + ib) W_\phi(1 + \beta - ib)}{\Gamma(1 + \beta - ib)}.$$

From the asymptotic estimates (9.7) for the gamma function combined with (5.3) with $u = 2$, since $N_\tau = \infty$, we get that, for $|b|$ large enough,

$$|\mathcal{M}_\beta(1/2 + ib)| \leq C|b|^{-\beta-2}.$$

Thus, for any $\beta \geq 0$,

$$|\mathcal{M}_\beta(1/2 + ib)| \in L^2(\mathbb{R})$$

and from the Parseval identity for Mellin transform, see (1.39), we deduce that $\bar{f}_\beta \in L^2(\mathbb{R}_+)$. Choosing $\beta = \alpha/2$, $\alpha \in [0, 1)$, then $f_{\alpha/2}(x) = x^{\alpha/2} \bar{f}_{\alpha/2}(x)$ and therefore $f_{\alpha/2} \in L^2(\mathbf{p}_\alpha)$. For $\lambda, x > 0$, set

$$f_{\alpha/2}^\lambda(x) = f_{\alpha/2}(\lambda x).$$

Then $f_{\alpha/2}^\lambda \in L^2(\mathbf{p}_\alpha)$ since $f_{\alpha/2} \in L^2(\mathbf{p}_\alpha)$. By dominated convergence, one shows using (8.27) and Proposition 6.1.2 that

$$\begin{aligned} \mathcal{I}_\phi f_{\alpha/2}^\lambda(x) &= \mathcal{I}_\phi \left(\sum_{n=0}^{\infty} \frac{W_\phi(n+1+\alpha/2)}{\Gamma(n+1+\alpha/2)n!} (-1)^n (\lambda x)^{n+\alpha/2} \right) = (\lambda x)^{\alpha/2} e^{-\lambda x} \\ &= p_{\alpha/2}(\lambda x) e^{-\lambda x}. \end{aligned}$$

Hence $\mathcal{I}_\phi f_{\alpha/2}^\lambda \in L^2(\mathbf{p}_\alpha)$ and $p_{-\alpha/2} \mathcal{I}_\phi f_{\alpha/2}^\lambda \in L^2(\mathbb{R}_+)$. Since, with $e_{-\lambda}(x) = e^{-\lambda x}$, $x > 0$, the linear hull of $\{e_{-\lambda}, \lambda > 0\}$ is dense in $L^2(\mathbb{R}_+)$ and $\mathcal{I}_\phi \in \mathbf{B}(L^2(\mathbf{p}_\alpha))$, $\alpha \in (0, 1)$, see Lemma 8.1.1(1), we get that

$$\overline{\text{Ran}(\mathcal{I}_\phi)}_{L^2(\mathbf{p}_\alpha)} = L^2(\mathbf{p}_\alpha).$$

Finally, when $\psi \in \mathcal{N}$ with $m = 0$ then $\mathcal{I}_\phi \in \mathbf{B}(L^2(\mathbb{R}_+))$ and this shows that in this case $\overline{\text{Ran}(\mathcal{I}_\phi)}_{L^2(\mathbb{R}_+)} = L^2(\mathbb{R}_+)$. \square

Uniform and norms estimates of the co-eigenfunctions

With the aim of characterizing the domain of the spectral operator, we continue our study on the co-eigenfunctions by focusing on their asymptotic estimates for large n considered in different topologies. We emphasize that the problem of estimating the norm of \mathcal{V}_n in $L^2(\nu)$ for large n seems extremely delicate as one is dealing with a weighted Hilbert which, in general, should require uniform asymptotic estimates for the co-eigenfunctions $\mathcal{V}_n(x)$. We point out that there is a rich and fascinating literature on uniform asymptotic expansions of the Laguerre polynomials which reveals already that this issue even in a simple case with explicit expression and several representation at hand is far from being trivial, see e.g. [Sze75] and [Tem90] for a thorough description. In this direction, we also mention our recent work where such uniform asymptotic analysis is conducted for some generalized Laguerre polynomials leading to the concept of reference semigroup that we elaborate in the next chapter. We carry out three different approaches that give the bounds for w_n or \mathcal{V}_n stated below, two are of complex analytical nature whereas the last one relies on probabilistic arguments. Recall that, with

$$\underline{\Theta}_\phi = \lim_{|b| \rightarrow \infty} \int_0^\infty \ln \left(\frac{|\phi(by + ib)|}{\phi(by)} \right) dy,$$

we have the class

$$\mathcal{N}_{\underline{\Theta}} = \{\psi \in \mathcal{N}; \underline{\Theta}_\psi > 0\} \text{ and } d_\phi = \{u \leq 0; \phi(u) = 0 \text{ or } \phi(u) = -\infty\}.$$

Also recall that

$$w_n(x) = \frac{(x^n \nu(x))^{(n)}}{n!}.$$

THEOREM 9.0.1. *(1) Let us assume that $\psi \in \mathcal{N}_{\underline{\Theta}}$. Then, we have for any $a > d_\phi$ (or $a \geq 0$ if $\phi(0) > 0$ and $d_\phi = 0$), for any $\epsilon > 0$ and n large, uniformly on $x > 0$,*

$$(9.1) \quad x^a |w_n(x)| = O\left(n^{\frac{1}{2}-a} e^{(T_{\underline{\Theta}_\phi} + \epsilon)n}\right),$$

where we recall that $T_{\underline{\Theta}_\phi} = -\ln \sin \underline{\Theta}_\phi$. Therefore, recalling that $\mathbf{p}_\alpha(x) = x^{-\alpha}$, $x > 0$, $\alpha \in (0, 1)$, we have, for large n , that

$$(9.2) \quad \|w_n\| = O\left(n^{\frac{1}{2}-a} e^{(T_{\underline{\Theta}_\phi} + \epsilon)n}\right),$$

$$(9.3) \quad \left\| \frac{w_n}{\mathbf{p}_\alpha} \right\|_{\mathbf{p}_\alpha} = O\left(n^{\frac{1}{2}-a} e^{(T_{\underline{\Theta}_\phi} + \epsilon)n}\right),$$

for any $d_\phi < a < \frac{1}{2}$ in (9.2) and $d_\phi < a < \frac{\alpha+1}{2}$ in (9.3).

(2) Let $\psi \in \mathcal{N}_\alpha$, i.e. $\psi(u) \approx C_\alpha u^{\alpha+1}$, $\alpha \in (0, 1)$, $C_\alpha > 0$. Then the following estimates hold.

(a) For any $a > d_\phi$ and $\epsilon > 0$,

$$(9.4) \quad \left\| \frac{w_n}{\bar{\nu}_\gamma} \right\|_{\bar{\nu}_\gamma} = O\left(n^{\frac{1}{2}-a} e^{(T_{\pi_\alpha} + \epsilon)n}\right),$$

where we recall that $T_{\pi_\alpha} = -\ln \sin\left(\frac{\pi}{2}\alpha\right)$ and

$$\bar{\nu}_\gamma(x) = e^{-x^{\frac{1}{\gamma}}}, \quad x > 0, \gamma > 1 + \alpha.$$

(b) Recalling that $T_{\pi_\alpha, \rho_\alpha} = \max\left(T_{\pi_\alpha}, 1 + \frac{1}{\rho_\alpha}\right)$, where ρ_α is the largest solution to the equation $(1 - \rho)^{\frac{1}{\alpha}} \cos\left(\frac{\arcsin(\rho)}{\alpha}\right) = \frac{1}{2}$, we have for large n and for any $\epsilon > 0$,

$$(9.5) \quad \|\mathcal{V}_n\|_\nu = O\left(e^{(T_{\pi_\alpha, \rho_\alpha} + \epsilon)n}\right).$$

REMARK 9.0.2. We point out that a better pointwise handle of $w_n(x)$ in the scenario of item (2) is attained in (9.19) but this does not allow for the improvement of the norm estimate.

REMARK 9.0.3. To derive the estimates (9.5), we resort to the Phragmenn-Lindelöf principle. This approach requires that the asymptotic behaviour of ν is characterized in terms of a conformal mapping forcing us to specialize to the class \mathcal{N}_α .

9.1. Proof of Theorem 9.0.1(1) via a classical saddle point method

Fix $n \in \mathbb{N}$, $a > d_\phi$, $-a \notin \mathbb{N}$ and $z = a + ib$. As $\psi \in \mathcal{N}_\Theta \subset \mathcal{N}_{\infty, \infty}$, then $N_\tau = \infty$, see (1.8), and there exists a constant $C_a > 0$ such that for any $b \in \mathbb{R}$

$$(9.6) \quad |\mathcal{M}_{V_\psi}(a + ib)| \leq C_a e^{-\Theta_\phi |b| + o(|b|)},$$

with $\Theta_\phi \in (0, \frac{\pi}{2}]$, see Theorem 6.0.2(2b). Note that (8.12) gives the Mellin transform of w_n

$$\mathcal{M}_{w_n}(z) = \frac{(-1)^n \Gamma(z)}{\Gamma(n+1)\Gamma(z-n)} \mathcal{M}_{V_\psi}(z).$$

Since $\mathcal{M}_{V_\psi} \in \mathcal{A}_{(d_\phi, \infty)}$, see Theorem 6.0.1(3), then plainly $\mathcal{M}_{w_n} \in \mathcal{A}_{(d_\phi, \infty)}$. Thanks to (9.6) and Lemma 5.3.5 with $N_\tau = \infty$, Mellin inversion, see (1.36), yields, for $x > 0$, that

$$w_n(x) = \frac{(-1)^n}{2\pi i} \int_{a-i\infty}^{a+i\infty} x^{-z} \frac{\Gamma(z)}{\Gamma(n+1)\Gamma(z-n)} \mathcal{M}_{V_\psi}(z) dz,$$

and, since for the integrand $|f(a+ib)| = |f(a-ib)|$, we get that

$$|w_n(x)| \leq C x^{-a} \int_0^\infty \frac{|\Gamma(a+ib)|}{\Gamma(n+1)|\Gamma(a+ib-n)|} |\mathcal{M}_{V_\psi}(a+ib)| db,$$

where throughout C stands for a generic positive constant. Recalling the formula

$$|\Gamma(a+ib-n)| |\Gamma(n+1-a+ib)| = \frac{\pi}{|\sin(\pi(a-n-ib))|},$$

and the uniform bound

$$|\sin(\pi(a + ib))| \leq C e^{\pi|b|}$$

together with the exponential decay in (9.6) we get, for any $a > d_\phi$, that

$$|w_n(x)| \leq C x^{-a} \int_0^\infty \frac{|\Gamma(a + ib)| |\Gamma(n + 1 - a + ib)|}{\Gamma(n + 1)} e^{(\pi - \Theta_\phi)b + o(b)} db.$$

Next, using the asymptotic relation (9.7), we get, for any $0 < \epsilon < \Theta_\phi$, that

$$|w_n(x)| \leq C x^{-a} \int_{-\infty}^\infty \frac{|\Gamma(n + 1 - a + ib)|}{\Gamma(n + 1)} e^{(\frac{\pi}{2} - \Theta_\phi + \epsilon)b} db.$$

Hence, using [PK01, Lemma 2.6], we obtain, for large n , that

$$|w_n(x)| \leq C x^{-a} n^{1-a} \frac{e^{n \ln n - n}}{\Gamma(n + 1)} \sec\left(\frac{\pi}{2} - \Theta_\phi + \epsilon\right)^{n + \frac{3}{2} - a}.$$

The Stirling approximation for $\Gamma(n + 1)$ in (9.8) shows that (9.1) holds. To prove (9.2) it suffices to integrate (9.1) for $a > \frac{1}{2}$ in a neighbourhood of infinity and to integrate (9.1) for $d_\phi < a < \frac{1}{2}$ in a neighbourhood of zero. Finally, it is trivial to extend our estimate (9.1) for $a = 0$. Similar arguments give (9.3) completing the proof of Theorem 9.0.1(1).

For the convenience of the reader, we recall the following auxiliary well-known result on the asymptotic of the gamma function that was used in the previous proof, see e.g. [PK01, (2.1.8)].

LEMMA 9.1.1. *For any fixed $a \in \mathbb{R}$*

$$(9.7) \quad |\Gamma(a + ib)| \approx C_a |b|^{a - \frac{1}{2}} e^{-\frac{\pi}{2}|b|},$$

where $C_a > 0$. We also have the Stirling approximation

$$(9.8) \quad \Gamma(n + 1) \approx \sqrt{2\pi n} n^{n + \frac{1}{2}} e^{-n}.$$

9.2. Proof of Theorem 9.0.1(2) via the asymptotic behaviour of zeros of the derivatives of ν

We start with the following useful Lemma.

LEMMA 9.2.1. *Let $\psi \in \mathcal{N}_{\infty, \infty}$. Then, with the notation of Lemma 5.3.1, there exists a sequence $\bar{a}_\nu = (\bar{a}_n = a_n^{-1} > 0)_{n \geq 0}$, such that for any $n \in \mathbb{N}$, for any $D \in (0, 1)$ and $x > D^{-1} \bar{a}_{n+1} > 0$*

$$(9.9) \quad \left| \nu^{(n)}(x) \right| \leq \left(\frac{D^n}{(1 - D)^n} \vee 1 \right) (n!)^2 e^{o(n)} \frac{\nu(xD)}{x^n},$$

$$(9.10) \quad \left| (x^n \nu(x))^{(n)} \right| \leq \left(\frac{D^n}{(1 - D)^n} \vee 1 \right) (n!)^2 e^{o(n)} \nu(xD).$$

Proof. Let $\psi \in \mathcal{N}_{\infty, \infty}$, i.e. $\mathfrak{r} = \phi(\infty) = \infty$. Recall that (5.84) states $\nu(x) = x^{-2} \widehat{\nu}_1(x^{-1})$. Then, from (5.85), we get, recalling the notation $\bar{k}_n = \sum_{j=1}^n k_j$ and $\tilde{k}_n = \sum_{j=1}^n j k_j$, that

$$\left(\widehat{\nu}_1 \left(\frac{1}{x} \right) \right)^{(n)} = \frac{(-1)^n}{x^n} \sum_{\bar{k}_n = n; \bar{k}_n = k} \frac{n!}{k_1! k_2! \dots k_n!} \frac{1}{x^k} \widehat{\nu}_1^{(k)} \left(\frac{1}{x} \right).$$

Next, we use the estimates (5.39) applied to $\widehat{\nu}_1$ with $1/\tau = 0$. Therefore, for $k \leq n$ we get

$$x > (Da_{n+1})^{-1} = D^{-1}\bar{a}_{n+1} > D^{-1}\bar{a}_{k+1} = \frac{1}{Da_{k+1}},$$

where $D \in (0, 1)$ and hence with $\bar{D} = \frac{D}{1-D}$ we arrive at

$$\widehat{\nu}_1^{(k)}(x^{-1}) \leq k! \bar{D}^k x^k \widehat{\nu}_1((Dx)^{-1}).$$

Thus from the expression above we obtain with $\mathbf{p}(n)$ the partition function of n , that, for all $x > D^{-1}\bar{a}_{n+1}$,

$$(9.11) \quad \left| (\widehat{\nu}_1(x^{-1}))^{(n)} \right| \leq (\bar{D}^n \vee 1) (n!)^2 \mathbf{p}(n) x^{-n} \widehat{\nu}_1((Dx)^{-1}).$$

Thus, for $x > D^{-1}\bar{a}_{n+1}$, from (9.11) and $\nu(x) = x^{-2} \widehat{\nu}_1(x^{-1})$ we get (9.9) via

$$\begin{aligned} & \left| \nu^{(n)}(x) \right| \\ & \leq \widehat{\nu}_1((Dx)^{-1}) \sum_{m=0}^n \binom{n}{m} \frac{(m+1)!}{x^{m+2}} x^{-n+m} (\bar{D}^{n-m} \vee 1) ((n-m)!)^2 \mathbf{p}(n-m) \\ & \leq x^{-n} \nu(xD) D^2 (\bar{D}^n \vee 1) \sum_{m=0}^n \binom{n}{m} (m+1)! ((n-m)!)^2 \mathbf{p}(n-m) \\ & \leq x^{-n} \nu(xD) (\bar{D}^n \vee 1) (n!)^2 n^2 \mathbf{p}(n) \\ & = x^{-n} \nu(xD) (\bar{D}^n \vee 1) (n!)^2 e^{o(n)}, \end{aligned}$$

where we have used the celebrated asymptotic formula

$$\mathbf{p}(n) \approx \frac{e^{\pi\sqrt{2n/3}}}{4n\sqrt{3}},$$

which can be found in [Usp20]. Similarly, for $x > D^{-1}\bar{a}_{n+1}$, using (9.11) again, we get, for $n \geq 2$, that

$$\begin{aligned} & \left| (x^n \nu(x))^{(n)} \right| \\ & = \left| (x^{n-2} \widehat{\nu}_1(x^{-1}))^{(n)} \right| \\ & \leq \frac{\widehat{\nu}_1((Dx)^{-1})}{x^2} \sum_{m=0}^{n-2} \binom{n}{m} \frac{(n-2)!}{(n-2-m)!} (\bar{D}^{n-m} \vee 1) ((n-m)!)^2 \mathbf{p}(n-m) \\ & \leq D^2 \nu(xD) (\bar{D}^n \vee 1) n! \mathbf{p}(n) \sum_{m=0}^{n-2} \frac{(n-m)!}{m!} \frac{(n-2)!}{(n-2-m)!} \\ & \leq \nu(xD) \mathbf{p}(n) (n!)^2 (\bar{D}^n \vee 1) \sum_{m=0}^{n-2} \frac{1}{m!} \\ & \leq e \nu(xD) (\bar{D}^n \vee 1) (n!)^2 e^{o(n)}, \end{aligned}$$

which completes the proof of (9.10). When $n = 1$ a similar inequality holds. \square

We are ready to prove Theorem 9.0.1(2). Recall from Lemma 5.2.5 that

$$\widehat{\kappa}(y) = \mathbf{n} \left(\int_0^\zeta e^{\epsilon s} ds > y \right), \text{ for } y > 0,$$

with $\widehat{\kappa}$ defined in Proposition 5.2.1 and $\mathbf{n}, \epsilon, \zeta$ in (5.2.1). Note that, for $0 < \epsilon < 1$,

$$\mathbf{n}(\zeta > y) \leq \mathbf{n} \left(\int_0^\zeta e^{\epsilon s} ds > y \right) \stackrel{0}{=} O(\mathbf{n}(\zeta > \epsilon y)),$$

where the first inequality is absolute since $\epsilon_s \geq 0$. For the second observe that

$$(9.12) \quad \mathbf{n} \left(\int_0^\zeta e^{\epsilon s} ds > y \right) \leq \mathbf{n} \left(\sup_{0 \leq s < \zeta} \epsilon_s > -\ln \epsilon \right) + \mathbf{n}(\zeta > \epsilon y),$$

and, note that since $\psi \in \mathcal{N}_\alpha \subset \mathcal{N}_{\infty, \infty}$ then

$$\lim_{y \rightarrow 0} \mathbf{n}(\zeta > \epsilon y) = \infty,$$

see Lemma 5.2.5(2), whereas

$$\mathbf{n} \left(\sup_{0 \leq s < \zeta} \epsilon_s > -\ln \epsilon \right) < \infty.$$

Next, recall, from (5.26), that $\mathbf{n}(\zeta > y)$ is the tail of the Lévy measure associated to $\bar{\phi} \in \mathcal{B}$ and

$$\psi(\bar{\phi}(u)) = \bar{\phi}(u)\phi(\bar{\phi}(u)) = u, \quad u > 0.$$

Put $\alpha_1 = \alpha + 1$. Since $\psi \in \mathcal{N}_\alpha$, $\psi(u) \stackrel{\infty}{\sim} C_\alpha u^{\alpha_1}$ from a classical result, see e.g. [BGT87, Proposition 1.5.15], we have $\bar{\phi}(u) \stackrel{\infty}{\sim} C_\alpha^{-\alpha_1} u^{\frac{1}{\alpha_1}}$ and thus

$$\widehat{\kappa}(y) = \mathbf{n}(\zeta > y) \stackrel{0}{\sim} C_\alpha^{-\alpha_1} \frac{\alpha}{\alpha_1} y^{-\frac{1}{\alpha_1}}.$$

Since $\psi \in \mathcal{N}_{\infty, \infty}$ then $\tau = \infty$ in Lemma 5.3.1(1) and with $a_n > \beta_n = \sup\{y > 0; \widehat{\kappa}(y) \geq n\}$ therein, we deduce that

$$a_n > \beta_n \asymp n^{-\alpha_1}.$$

We also have, from (9.10), that, for any fixed $D < 1$ and $x > D^{-1}a_{n+1}^{-1}$,

$$(9.13) \quad |w_n(x)| = \frac{|(x^n \nu(x))^{(n)}|}{n!} \leq \left(\frac{D^n}{(1-D)^n} \vee 1 \right) n! e^{o(n)} \nu(Dx).$$

Thus, recalling that $\bar{\nu}_\gamma(x) = e^{-x^{\frac{1}{\gamma}}}$, $x > 0$, with $\gamma > \alpha_1 = \alpha + 1$, and, since $a_n > \beta_n \asymp n^{-\alpha_1}$ we can assume that $a_n^{-1} \leq n^{\alpha_1}$ for all n big enough to get that

$$\int_0^{2a_n^{-1}} \frac{w_n^2(x)}{\bar{\nu}_\gamma(x)} dx \leq \int_0^{2n^{\alpha_1}} \frac{w_n^2(x)}{\bar{\nu}_\gamma(x)} dx \leq \frac{\|w_n\|^2}{\bar{\nu}_\gamma((2n)^{\alpha_1})} = e^{o(n)} \|w_n\|^2.$$

Next, the bound (9.13) with $D = \frac{1}{2}$, for $x > 2a_n^{-1}$ and the asymptotic equivalent (5.9) for $\psi \in \mathcal{N}_\alpha$, yield for large n and with some generic constant $C > 0$ that

$$\begin{aligned} \int_{2n^{\alpha_1}}^{\infty} \frac{w_n^2(x)}{\bar{\nu}_\gamma(x)} dx &\leq \frac{1}{(n!)^2} \int_{2n^{\alpha_1}}^{\infty} \frac{((x^n \nu(x))^{(n)})^2}{\bar{\nu}_\gamma(x)} dx \\ &\leq (n!)^2 e^{o(n)} \int_{2n^{\alpha_1}}^{\infty} \frac{\nu^2\left(\frac{x}{2}\right)}{\bar{\nu}_\gamma(x)} dx \\ &\leq (n!)^2 e^{o(n)} \int_{2n^{\alpha_1}}^{\infty} e^{-Cx^{\frac{1}{\alpha}} + x^{\frac{1}{\gamma}}} dx \leq (n!)^2 e^{o(n)} e^{-\frac{C}{2} n^{\frac{\alpha_1}{\alpha}}} \\ &\leq e^{-\frac{C}{3} n^{1+\frac{1}{\alpha}}}. \end{aligned}$$

Therefore,

$$\left\| \frac{w_n}{\bar{\nu}_\gamma} \right\|_{\bar{\nu}_\gamma}^2 \leq e^{o(n)} \|w_n\|^2 + o(1)$$

and (9.2) provides (9.4).

9.3. Proof of Theorem 9.0.1(2b) through Phragmén-Lindelöf principle

This part aims to provide norm estimates of the co-eigenfunctions \mathcal{V}_n in the weighted Hilbert space $L^2(\nu)$ when $\psi \in \mathcal{N}_\alpha$. The approach we develop here is based on the Phragmén-Lindelöf principle which allows us to obtain precise bounds on a sector of the complex plane for the invariant density ν and its derivatives. In this vein we denote by θ_ν the angle of analyticity of ν , i.e. $\nu \in \mathcal{A}(\theta_\nu)$. We first state the following bound, which is an extension of (5.46).

LEMMA 9.3.1. *Let $\psi \in \mathcal{N}_{\underline{\Theta}}$. Then, for all $n, k \in \mathbb{N}$ and $a > d_\phi$ (or $a \geq 0$ if $d_\phi = 0$ and $\phi(0) > 0$), we have, for all $z \in \mathbb{C}(\underline{\Theta}_\phi)$, that*

$$(9.14) \quad \left| (z^n \nu(z))^{(k)} \right| \leq C(|\arg z|) |z|^{n-k-a},$$

where $C(|\arg z|) = C(n, k, a, |\arg z|) > 0$ is increasing in $|\arg z|$ for fixed n, k, a .

Proof. Recall that since $\psi \in \mathcal{N}_{\underline{\Theta}}$ then $\nu \in \mathcal{A}(\underline{\Theta}_\phi)$ from Theorem 5.0.2(3), and, from Proposition 6.0.1(2) combined with the upper bound in (6.47), that, for all $\epsilon > 0$ and a as in the statement, its Mellin transform is for large b ,

$$|\mathcal{M}_{V_\psi}(a + ib)| = |W_\phi(a + ib)| = O\left(e^{-(\underline{\Theta}_\phi - \epsilon)|b|}\right).$$

This enables to use the inversion formula (1.38), to get that for all $n, k \in \mathbb{N}$, a as above and $x > 0$,

$$\begin{aligned} (x^n \nu(x))^{(k)} &= \frac{(-1)^k}{2\pi i} \int_{a+n-k-i\infty}^{a+n-k+i\infty} x^{-s} (s-k)_k \mathcal{M}_{V_\psi}(s+n-k) ds \\ (9.15) \quad &= \frac{(-1)^k}{2\pi i} x^{n-k} \int_{a-i\infty}^{a+i\infty} x^{-s} (s-n)_k \mathcal{M}_{V_\psi}(s) ds. \end{aligned}$$

Let now $z \in \mathbb{C}(\underline{\Theta}_\phi)$ and choosing $\epsilon > 0$ above such that $\theta_\epsilon = \underline{\Theta}_\phi - \epsilon - \arg z > 0$, we get that

$$\left| \int_{a-i\infty}^{a+i\infty} z^{-s} (s-n)_k \mathcal{M}_{V_\psi}(s) ds \right| \leq \frac{|z|^{-a}}{2\pi} \int_{-\infty}^{\infty} |(a+ib-n)_k| e^{-\theta_\epsilon |b|} db < \infty.$$

The upper bound (9.14) is then easily derived by means of this latter bound combined with an obvious extension of the identity (9.15). Finally, the monotonicity in $|\arg z|$ of $|z^{-a-ib}| = |z|^{-a}e^{|\arg z||b|}$ yields the same property for $C(n, k, a, |\arg z|)$ for fixed n, k, a . \square

We proceed with the following estimates.

LEMMA 9.3.2. *Let $\psi \in \mathcal{N}_\alpha$, i.e. $\psi(u) \lesssim C_\alpha u^{\alpha+1}$, $\alpha \in (0, 1)$. Then, for any $n \geq 0$, $z \in \mathbb{C}(\frac{\pi}{2}\alpha)$, we have that for any $0 < \epsilon < \frac{\pi}{2}$, there exists $C = C(\epsilon, n) > 0$, such that with $\bar{c}_\alpha = \alpha C_\alpha^{-\frac{1}{\alpha}}$*

$$(9.16) \quad \left| \nu^{(n)}(1+z) \right| \leq C e^{-\bar{c}_\alpha \cos(\frac{\arg z}{\alpha} + \epsilon) |z|^{\frac{1}{\alpha}}}.$$

Proof. We first consider the case $n = 0$ and we set, for $0 < \epsilon < \frac{\pi}{2}$ and $x > 0$,

$$V_\epsilon(x) = \nu(1+x^\alpha) e_\epsilon(x) = \nu(1+x^\alpha) e^{\bar{c}_\alpha x e^{i\epsilon}}.$$

Then, the asymptotic bound (5.9) for $\nu(1+x^\alpha)$ entails, that, for large $x > 0$,

$$|V_\epsilon(x)| = |\nu(1+x^\alpha) e_\epsilon(x)| \leq e^{-\bar{c}_\alpha(1-\cos \epsilon)x + o(x)}.$$

Therefore since $\nu \in \mathcal{C}_0(\mathbb{R}_+)$, see Theorem 5.0.2(2a), there exists a constant $C_\epsilon > 0$ such that

$$(9.17) \quad \sup_{x \geq 0} |V_\epsilon(x)| < C_\epsilon.$$

On the other hand, Theorem 6.0.2(2(b)ii) implies that $\psi \in \mathcal{N}_\alpha$ with $\Theta_\psi = \frac{\pi}{2}\alpha$ and, thus thanks to Theorem 5.0.2(3), $\nu \in \mathcal{A}(\frac{\pi}{2}\alpha)$. Hence, $V_\epsilon \in \mathcal{A}(\frac{\pi}{2})$. If $\bar{z}_\epsilon \in \mathbb{C}$ with $\arg \bar{z}_\epsilon = \frac{\pi}{2} - \epsilon$, then $1 + \bar{z}_\epsilon^\alpha \in \mathbb{C}(\frac{\pi}{2}\alpha)$. Thus, estimate (9.14) applied to $\nu(1 + \bar{z}_\epsilon^\alpha)$ with $a > 0, k = n = 0$, yields

$$(9.18) \quad \begin{aligned} |V_\epsilon(\bar{z}_\epsilon)| &= |\nu(1 + \bar{z}_\epsilon^\alpha) e_\epsilon(\bar{z}_\epsilon)| \\ &= |\nu(1 + \bar{z}_\epsilon^\alpha)| \leq \underline{C}_\epsilon |1 + \bar{z}_\epsilon^\alpha|^{-a} \leq \tilde{C}_\epsilon, \end{aligned}$$

where

$$\underline{C}_\epsilon = \underline{C}_\epsilon \left(\frac{\pi}{2} - \epsilon \right), \tilde{C}_\epsilon = \tilde{C}_\epsilon \left(\frac{\pi}{2} - \epsilon \right) > 0$$

and we have used that since $\arg \bar{z}_\epsilon = \frac{\pi}{2} - \epsilon$ then

$$|e_\epsilon(\bar{z}_\epsilon)| = \left| e^{\bar{c}_\alpha |\bar{z}_\epsilon| e^{i(\frac{\pi}{2}-\epsilon)}} e^{i\epsilon} \right| = \left| e^{i\bar{c}_\alpha |\bar{z}_\epsilon|} \right| = 1.$$

Therefore, we deduce from the bounds (9.17) and (9.18) that the function V_ϵ is bounded on the boundary of the sector

$$\mathbb{C}^+ \left(\frac{\pi}{2} - \epsilon \right) = \left\{ z \in \mathbb{C}; 0 < \arg(z) < \frac{\pi}{2} - \epsilon \right\}.$$

It is obvious, again from estimate (9.14) with $n = k = 0$ and some $a > 0$, that, for any $z \in \mathbb{C}^+ \left(\frac{\pi}{2} - \epsilon \right)$, we have that

$$|V_\epsilon(z)| = \left| \nu(1+z^\alpha) e^{\bar{c}_\alpha |z| e^{i(\arg z + \epsilon)}} \right| \leq C e^{\bar{c}_\alpha \cos(\arg z + \epsilon) |z|}$$

where $C = C(a, \alpha \left(\frac{\pi}{2} - \epsilon \right)) > 0$ is the constant in Lemma 9.3.1 and $C(a, \cdot)$ is increasing. Thus

$$|V_\epsilon(z)| \leq C e^{\bar{c}_\alpha \cos(\arg z + \epsilon) |z|}$$

for all $z \in \mathbb{C}^+ \left(\frac{\pi}{2} - \epsilon \right)$. Since, clearly for any $0 < \epsilon < \frac{\pi}{2}$, then $1 < \frac{\pi}{\frac{\pi}{2} - \epsilon}$, an application of the Phragmen-Lindelöf principle, see e.g. [Tit39], yields that, for any $z \in \mathbb{C}^+ \left(\frac{\pi}{2} - \epsilon \right)$,

$$|V_\epsilon(z)| \leq \bar{C}_\epsilon$$

where

$$\bar{C}_\epsilon = \max(C_\epsilon, \tilde{C}_\epsilon) > 0.$$

Thus, for any $z \in \mathbb{C}^+ \left(\frac{\pi}{2} - \epsilon \right)$,

$$|\nu(1 + z^\alpha)| \leq \bar{C}_\epsilon e^{-\bar{c}_\alpha \cos(\arg z + \epsilon)|z|},$$

and we conclude that

$$|\nu(1 + z)| \leq \bar{C}_\epsilon e^{-\bar{c}_\alpha \cos\left(\frac{\arg z}{\alpha} + \epsilon\right)|z|^{\frac{1}{\alpha}}}.$$

The case $\arg z \in \left(-\left(\frac{\pi}{2} - \epsilon\right), 0\right)$ is identical. This proves the statement for $n = 0$. The case $n \geq 1$ follows similarly using the asymptotic equivalent (5.9) valid for $\nu^{(n)}$ for all $n \geq 1$. \square

Our next lemma gives the bounds to complete the proof of the estimates in Theorem 9.0.1(2b).

LEMMA 9.3.3. *Let $\psi \in \mathcal{N}_\alpha$. Then, for every $x > 2$, we have that for any $\rho \in (0, \sin(\frac{\pi}{2}\alpha))$, $\epsilon > 0$ and $n \geq 0$, there exists $C = C(\rho, \epsilon) > 0$ such that*

$$(9.19) \quad |w_n(x)| \leq C R_\rho^n (x-1) e^{-\bar{C}_{\alpha,\rho}(x-1)^{\frac{1}{\alpha}}}$$

where

$$R_\rho(x) = 1 + \frac{1}{\rho} + \frac{1}{\rho(x-1)}$$

and

$$\bar{C}_{\alpha,\rho} = \bar{c}_\alpha (1-\rho)^{\frac{1}{\alpha}} \cos\left(\frac{\arcsin(\rho)}{\alpha} + \epsilon\right).$$

Proof. Recall that since $\psi \in \mathcal{N}_\alpha$ then $\nu \in \mathcal{A}\left(\frac{\pi}{2}\alpha\right)$. Then, an application of the Cauchy integral representation yields, for $\rho \in (0, \sin(\frac{\pi}{2}\alpha))$ and any $x > 1$ and $x_1 = x - 1$, that

$$\begin{aligned} |w_n(x)| &= \left| \frac{(x^n \nu(x))^{(n)}}{n!} \right| = \left| \frac{1}{2\pi i} \oint \frac{\zeta^n \nu(\zeta)}{(\zeta - x)^{n+1}} d\zeta \right| \\ &\leq \frac{|x + \rho x_1|^n}{\rho^n x_1^n} \sup_{\theta \in [0, 2\pi]} |\nu(1 + x_1(1 + \rho e^{i\theta}))|, \end{aligned}$$

where the integral is over the circle

$$\mathcal{C} = \{z \in \mathbb{C}; |z - x| < \rho x_1\} \subset \mathbb{C}\left(\frac{\pi}{2}\alpha\right).$$

Next, observe on the one hand, that

$$\left| \frac{x + \rho x_1}{\rho x_1} \right|^n \leq R_\rho^n(x_1) = R_\rho^n(x-1)$$

and, on the other hand that

$$\begin{aligned} \sup_{\theta \in [0, 2\pi]} |\nu(1 + x_1(1 + \rho e^{i\theta}))| &\leq \sup_{z \in \mathcal{C}} |\nu(1 + z)| \\ &\leq \sup_{z \in \mathbb{C}_{\rho, x_1}} |\nu(1 + z)| \end{aligned}$$

where

$$\mathcal{C} \subseteq \mathbb{C}_{\rho, x_1} = \{z \in \mathbb{C}; \Re(z) > (1 - \rho)x_1 \text{ and } |\arg z| < \arcsin \rho\}.$$

To complete the proof, one bounds the last term by means of estimate (9.16) applied to $|\nu(1 + z)|$ with $|z| = (1 - \rho)x_1$ and $\arg z = \arcsin(\rho)$. \square

We are now ready to derive the bound Theorem 9.0.1(2b). To this end, for any $A > 2$, $n \in \mathbb{N}$, we write

$$\begin{aligned} \mathcal{V}_n(x) &= \mathcal{V}_n(x)\mathbb{I}_{\{0 < x \leq A\}} + \mathcal{V}_n(x)\mathbb{I}_{\{x > A\}} \\ &= \underline{v}_n(x) + \bar{v}_n(x). \end{aligned}$$

Next, choose $\rho \in (0, \sin(\frac{\pi}{2}\alpha))$ such that

$$(1 - \rho)^{\frac{1}{\alpha}} \cos\left(\frac{\arcsin(\rho)}{\alpha}\right) = \frac{1}{2} + h,$$

for some $h \in (0, \frac{1}{2})$, and choose ϵ, ϵ_1 so small enough that

$$T(\rho, h, \epsilon, \epsilon_1) = (1 + 2h) \cos(\epsilon) - (1 + \epsilon_1) - 2(1 - \rho)^{\frac{1}{\alpha}} \sin\left(\left(\frac{\arcsin(\rho)}{\alpha}\right)\right) \sin(\epsilon) > 0.$$

Since from (5.9) there exist $C_A > 0$ such that for $x > A$,

$$\nu(x) \geq C_A e^{-\bar{c}_\alpha(x-1)^{\frac{1}{\alpha}}(1+\epsilon_1)},$$

the bound (9.19) with ϵ combined with an expansion of

$$\cos(x + y) = \cos(x) \cos(y) - \sin(x) \sin(y)$$

gives, for some constant $C = C_{h, \epsilon, A, \rho} > 0$, that

$$\begin{aligned} \|\bar{v}_n\|_{\nu}^2 &= \int_A^\infty \frac{w_n^2(x)}{\nu(x)} dx \\ &\leq CR_\rho^{2n}(A-1) \int_A^\infty e^{-\bar{c}_\alpha(x-1)^{\frac{1}{\alpha}}} \left(2(1-\rho)^{\frac{1}{\alpha}} \cos\left(\frac{\arcsin(\rho)}{\alpha} + \epsilon\right) - (1+\epsilon_1)\right) dx \\ &= CR_\rho^{2n}(A-1) \int_A^\infty e^{-\bar{c}_\alpha(x-1)^{\frac{1}{\alpha}}} T(\rho, h, \epsilon, \epsilon_1) dx \\ &\leq C'R_\rho^{2n}(A-1), \end{aligned}$$

since $R_\rho^n(x)$ is decreasing in x . Thanks to (9.1), for any $x > 0$, $a > d_\phi$ and $\epsilon \in (0, \frac{\pi}{2}\alpha)$,

$$w_n^2(x) = O\left(n^{1-2a} \sin\left(\frac{\pi}{2}\alpha - \epsilon\right)^{-2n} x^{-2a}\right),$$

where recall that $d_\phi \leq 0$, see (1.19). Theorem 5.0.4(5.5) yields that for $x < A$ and any $a = -d_\phi + \epsilon$,

$$\nu(x) \geq C_a x^a$$

and henceforth

$$\int_0^A \frac{w_n^2(x)}{\nu(x)} dx = O\left(n^{1-2a} \sin\left(\frac{\pi}{2}\alpha - \epsilon\right)^{-2n} \int_0^A x^{-2d_\phi - 2\epsilon + d_\phi - \epsilon} dx\right)$$

as long as $\epsilon < 1/3$. Therefore, we conclude in this case as well that

$$\|\underline{v}_n\|_\nu^2 = O\left(n^{1-2a} \sin\left(\frac{\pi}{2}\alpha - \epsilon\right)^{-2n}\right).$$

Combining the estimates for $\|\underline{v}_n\|_\nu^2$ and $\|\bar{v}_n\|_\nu^2$ we complete the proof of (9.5).

The concept of reference semigroups: $L^2(\nu)$ -norm estimates and completeness of the set of co-eigenfunctions

We develop here the concept of reference semigroups whose main underlying idea can be explained as follows. It consists in identifying gL semigroups \overline{P} which satisfy the following two criteria. First, their special structure permits to study their spectral decomposition in detail. In particular, one has accurate information on their sequence of co-eigenfunctions such as precise norm estimates and completeness. Furthermore, there should exist a subclass of gL semigroups such that for each element P in this class we have the adjoint intertwining relation $P_t^* \Lambda^* = \Lambda^* \overline{P}_t^*$, where Λ^* is the adjoint of a bounded operator. Indeed, under these conditions one has, with the obvious notation, $\mathcal{V}_n = \Lambda^* \overline{\mathcal{V}}_n$, providing readily the existence as well as bounds of the norms of \mathcal{V}_n in the $L^2(\nu)$ topology. Another delicate issue that is of great interest concerns the completeness properties of the sequence $(\mathcal{V}_n)_{n \geq 0}$. Assuming that this property holds for the sequence of co-eigenfunctions of the reference semigroup \overline{P} , one may deduce this property for $(\mathcal{V}_n)_{n \geq 0}$ as soon as the bounded operator Λ^* has a dense range. Although this approach may be extended to more general classes, we present below two different reference semigroups, namely the one-parametric class of self-adjoint Laguerre semigroups and the two-parameter family of Gauss-Laguerre semigroups whose detailed spectral analysis has been conducted by the authors in [PS17] and are reviewed in Example 3.0.3 above. In particular, this approach enables us to deal with the spectral expansion in the full Hilbert space of the perturbation class \mathcal{N}_P , that is when $\sigma^2 > 0$. It is worth pointing out that one of the main technical difficulty in implementing this approach is to identify intertwining operators which are bounded in the appropriate Hilbert spaces. One of the key steps in achieving this aim is the new development on Bernstein functions presented in Proposition 4.2.2, namely that a subset of the cone of Bernstein functions is invariant under multiplication. Recalling the notation $\mathfrak{c}_n(a) = \frac{\Gamma(n+1)\Gamma(a+1)}{\Gamma(n+a+1)}$, for $n \in \mathbb{N}$ and $a > 0$, we are now ready to state the following.

THEOREM 10.0.1. (1) *Let $\psi \in \mathcal{N}_P$. Then, for any $\epsilon > 0$ and large n ,*

$$(10.1) \quad \|\mathcal{V}_n\|_\nu = O(e^{\epsilon n}).$$

If in addition $\overline{\Pi}(0^+) < \infty$ then, recalling that $\underline{\mathfrak{m}} = \frac{m + \overline{\Pi}(0^+)}{\sigma^2}$, we have for large n ,

$$(10.2) \quad \|\mathcal{V}_n\|_\nu = O(n^{\underline{\mathfrak{m}}}),$$

and the sequence $(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})}\mathcal{V}_n)_{n \geq 0}$ is a Bessel sequence in $L^2(\nu)$. Finally, if $-d_\phi > 0$, then for any $\epsilon > 0$ such that $\mathbf{d}_\epsilon = -d_\phi - \epsilon > 0$, the sequence $\left(\frac{\mathcal{P}_n}{\sqrt{\mathbf{c}_n(\mathbf{d}_\epsilon)}}\right)_{n \geq 0}$ is a Bessel sequence in $L^2(\nu)$.

(2) Let $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$ with $(\alpha, \mathbf{m}) \in \mathfrak{R} = \{(\alpha, \mathbf{m}); \alpha \in (0, 1] \text{ and } \mathbf{m} \geq 1 - \frac{1}{\alpha}\}$. Then, for large n , with $T_\alpha = -\ln(2^\alpha - 1)$, we have that

$$(10.3) \quad \|\mathcal{V}_n\|_\nu = O(e^{T_\alpha n}).$$

(3) Recall from (1.34) that $\mathcal{N}_{\alpha, \mathbf{m}}^{d_\phi} = \{\psi \in \mathcal{N}_{\alpha, \mathbf{m}}; d_\phi < 1 - \frac{\mathbf{m}}{2} - \frac{1}{2\alpha}\}$ and assume that $\psi \in \mathcal{N}_P \cup \mathcal{N}_{\alpha, \mathbf{m}}^{d_\phi}$. Then,

$$\overline{\text{Span}}(\mathcal{V}_n) = L^2(\nu).$$

The rest of the chapter is devoted to the proof of this Theorem and is structured as follows. Estimates (10.1) and (10.3) are settled in 10.1.2; the completeness of \mathcal{V}_n , i.e. item (2) is proved in 10.2; estimate (10.2) and the succeeding claims of (1) are proved in 10.1.1.

10.1. Estimates for the $L^2(\nu)$ norm of \mathcal{V}_n

10.1.1. The small perturbation case. Recall that, for any $\alpha \in (0, 1]$ and $\mathbf{m} \geq 1 - \frac{1}{\alpha}$

$$u \mapsto \phi_{\alpha, \mathbf{m}}^R(u) = \frac{\Gamma(\alpha u + \alpha \mathbf{m} + 1)}{\Gamma(\alpha u + \alpha \mathbf{m} + 1 - \alpha)} \in \mathcal{B}_\mathcal{N},$$

see (1.15), and, put $\phi_{\mathbf{m}}^R(u) = \phi_{1, \mathbf{m}}^R(u) = u + \mathbf{m}$. Note also that, for any $z \in \mathbb{C}_{(-\mathbf{m} - \frac{1}{\alpha}, \infty)}$,

$$(10.4) \quad W_{\phi_{\alpha, \mathbf{m}}^R}(z + 1) = \frac{\Gamma(\alpha z + \alpha \mathbf{m} + 1)}{\Gamma(\alpha \mathbf{m} + 1)}.$$

LEMMA 10.1.1. *Let $\psi \in \mathcal{N}_P$ with $\overline{\Pi}(0^+) < \infty$. Then, for any $\mathbf{m} > \underline{\mathbf{m}} = \frac{\mathbf{m} + \overline{\Pi}(0^+)}{\sigma^2}$, $\phi_{\mathbf{m}}(u) = \frac{\phi(u)}{u + \mathbf{m}} \in \mathcal{B}$. With $\mathcal{V}_{\phi_{\mathbf{m}}}$ (resp. $\mathcal{E}_{\mathbf{m}}$) denoting the Markov operator associated to the variable $V_{\phi_{\mathbf{m}}}$ (resp. $V_{\phi_{\mathbf{m}}^R}$), we have the following factorization of multiplicative Markov operators*

$$(10.5) \quad \mathcal{E}_{\mathbf{m}} \mathcal{V}_{\phi_{\mathbf{m}}} = \mathcal{V}_{\phi_{\mathbf{m}}} \mathcal{E}_{\mathbf{m}} = \mathcal{V}_\psi.$$

Moreover, $\mathcal{V}_{\phi_{\mathbf{m}}} \in \mathbf{B}(L^2(\nu), L^2(\varepsilon_{\mathbf{m}}))$, $\varepsilon_{\mathbf{m}}(x) = \frac{x^{\mathbf{m}}}{\Gamma(\mathbf{m} + 1)} e^{-x} dx$, $x > 0$, and the following intertwining relation

$$(10.6) \quad Q_t^{(\mathbf{m})} \mathcal{V}_{\phi_{\mathbf{m}}} = \mathcal{V}_{\phi_{\mathbf{m}}} P_t$$

holds on $L^2(\nu)$ for all $t \geq 0$, where $Q^{(\mathbf{m})} = (Q_t^{(\mathbf{m})})_{t \geq 0}$ is the classical Laguerre semigroup of order \mathbf{m} , see Chapter 3. Consequently, the sequence $(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})}\mathcal{V}_n)_{n \geq 0}$, where $\mathbf{c}_n(\underline{\mathbf{m}}) = \frac{\Gamma(n+1)\Gamma(\mathbf{m}+1)}{\Gamma(n+\mathbf{m}+1)}$, is a Bessel sequence in $L^2(\nu)$ with bound 1 and for any $n \geq 0$,

$$(10.7) \quad 1 \leq \|\mathcal{V}_n\|_\nu \leq \mathbf{c}_n^{-1}(\underline{\mathbf{m}}) \cong O(n^{\underline{\mathbf{m}}}).$$

Proof. Note that $\phi_m \in \mathcal{B}$ is given in Proposition 4.2.2. Next, observing, from (7.19), (6.3), (10.4) with $\alpha = 1$ and

$$\mathbb{E} [V_{\phi_m}^n] = W_{\phi_m}(n+1) = \prod_{k=1}^n \frac{\phi(k)}{m+k}$$

that for any $n \geq 0$ with $p_n(x) = x^n$,

$$\begin{aligned} \mathcal{E}_m \mathcal{V}_{\phi_m} p_n(1) &= \mathbb{E} [V_{\phi_m}^n] \mathbb{E} [V_{\phi_m^R}^n] = \frac{\Gamma(n+1+m)}{\Gamma(m+1)} \prod_{k=1}^n \frac{\phi(k)}{m+k} = W_{\phi}(n+1) \\ &= \mathcal{V}_{\psi} p_n(1), \end{aligned}$$

we get the factorization (10.5) by moment determinacy of the involved operators. Replicating the bound (7.22) in the proof of Proposition 7.0.1, we get that $\mathcal{V}_{\phi_m} \in \mathbf{B}(L^2(\nu), L^2(\varepsilon_m))$ and \mathcal{V}_{ϕ_m} is a contraction. To prove (10.6) note, with the help of (3.8) and (7.7), that

$$\begin{aligned} \mathcal{V}_{\phi_m} \mathcal{P}_n(x) &= \sum_{k=0}^n (-1)^k \frac{\mathbb{E} [V_{\phi_m}^k] \binom{n}{k}}{W_{\phi}(k+1)} x^k = \sum_{k=0}^n (-1)^k \frac{\Gamma(m+1) \binom{n}{k}}{\Gamma(m+k+1)} x^k \\ (10.8) \quad &= \mathbf{c}_n(\mathbf{m}) \mathcal{L}_n^{(\mathbf{m})}(x). \end{aligned}$$

We refer to Example 3.0.1 in Chapter 3 for a detailed description of the Laguerre semigroup of order $m \geq 0$. Then, from the eigenfunction property of \mathcal{P}_n , we have, for any $t, x > 0$,

$$\begin{aligned} \mathcal{V}_{\phi_m} P_t \mathcal{P}_n(x) &= e^{-nt} \mathcal{V}_{\phi_m} \mathcal{P}_n(x) = e^{-nt} \mathbf{c}_n(\mathbf{m}) \mathcal{L}_n^{(\mathbf{m})}(x) = Q_t^{(\mathbf{m})} \mathbf{c}_n(\mathbf{m}) \mathcal{L}_n^{(\mathbf{m})}(x) \\ &= Q_t^{(\mathbf{m})} \mathcal{V}_{\phi_m} \mathcal{P}_n(x). \end{aligned}$$

Since V_{ψ} is moment determinate, see Theorem 6.0.1, we have $\overline{\text{Span}(\mathcal{P}_n)} = L^2(\nu)$, see [Akh65]. Thus, the intertwining (10.6) holds on a dense subset, and, by continuity of the involved operators, on $L^2(\nu)$. Finally, repeating the computation (8.5) with the intertwining relation (10.6) and using that $Q^{(\mathbf{m})}$ is self-adjoint with coeigenfunctions $(\mathcal{L}_n^{(\mathbf{m})})_{n \geq 0}$, see Example 3.0.1, we deduce, with the obvious notation, that

$$\mathcal{V}_{\phi_m}^* \mathcal{L}_n^{(\mathbf{m})}(x) = \mathcal{V}_n(x).$$

The last relation concludes that the sequence $(\sqrt{\mathbf{c}_n(\mathbf{m})} \mathcal{V}_n)_{n \geq 0}$ is a Bessel sequence and that (10.7) holds by replicating the computation (7.26) and utilizing that $(\sqrt{\mathbf{c}_n(\mathbf{m})} \mathcal{L}_n^{(\mathbf{m})})_{n \geq 0}$ is an orthonormal sequence in $L^2(\varepsilon_m)$, $\mathcal{V}_{\phi_m}^* \in \mathbf{B}(L^2(\varepsilon_m), L^2(\nu))$ is a contraction and $\mathbf{m} > \underline{\mathbf{m}}$. \square

We proceed by an additional factorization of the entrance law. It will be useful for getting a better upper bound for the spectral operator norm and also later in this chapter to prove completeness of the sequence of co-eigenfunctions. Recall that $\varepsilon_m(x) = \frac{x^m}{\Gamma(m+1)} e^{-x}$, $x > 0, m \geq 0$, and $\mathcal{E}_m f(x) = \mathbb{E}[f(x\mathbf{e}_m)]$, with simply $\mathcal{E} = \mathcal{E}_0$, is the Markov operator associated to \mathbf{e}_m , that is the gamma random variable with density function ε_m . Let

$$(10.9) \quad \mathfrak{L}_a f(x) = \int_0^{\infty} f(y) y^a e^{-yx} dy$$

stand for the weighted Laplace transform with weight $a \in \mathbb{R}$ and write simply $\mathfrak{L} = \mathfrak{L}_0$.

LEMMA 10.1.2. *Let $\psi \in \mathcal{N}$ with $-d_\phi > 0$. Then, for any $\epsilon \in (0, -d_\phi)$, writing $\mathbf{d}_\epsilon = -d_\phi - \epsilon > 0$, we have that*

$$\phi_{\mathbf{d}_\epsilon}(u) = \frac{u}{u + \mathbf{d}_\epsilon} \phi(u) \in \mathcal{B}.$$

For every function $f : \mathbb{R}_+ \mapsto \mathbb{R}$ and $x > 0$ such that $\mathcal{E}_{\mathbf{d}_\epsilon}|f|(x) < \infty$, we have that

$$(10.10) \quad \mathcal{E}_{\mathbf{d}_\epsilon} f(x) = \mathcal{V}_\psi \mathcal{I}_{\phi_{\mathbf{d}_\epsilon}} f(x) = \mathcal{I}_{\phi_{\mathbf{d}_\epsilon}} \mathcal{V}_\psi f(x),$$

$g_x(v) = \mathcal{V}_\psi f(xv)$ is $\iota_{\mathbf{d}_\epsilon}(v)dv$ -a.e. finite, where $\iota_{\mathbf{d}_\epsilon}$ is the density of $\mathcal{I}_{\phi_{\mathbf{d}_\epsilon}}$ and

$$(10.11) \quad \begin{aligned} \mathcal{E}_{\mathbf{d}_\epsilon} f(x) &= \int_0^\infty f(xy) \frac{y^{\mathbf{d}_\epsilon} e^{-y}}{\Gamma(\mathbf{d}_\epsilon + 1)} dy = x^{-1-\mathbf{d}_\epsilon} \int_0^\infty f(y) \frac{y^{\mathbf{d}_\epsilon} e^{-\frac{y}{x}}}{\Gamma(\mathbf{d}_\epsilon + 1)} dy \\ &= \frac{x^{-1-\mathbf{d}_\epsilon} \mathfrak{L}_{\mathbf{d}_\epsilon} f\left(\frac{1}{x}\right)}{\Gamma(\mathbf{d}_\epsilon + 1)}. \end{aligned}$$

All claims above hold with $\mathbf{d}_\epsilon = 0$ if $d_\phi = 0$.

Proof. Since $\psi \in \mathcal{N}$ with

$$-d_\phi = -\sup\{u \leq 0; \phi(u) = -\infty \text{ or } \phi(u) = 0\} > 0$$

then $m = \phi(0) > 0$, $\phi > 0$ on $(d_\phi, 0)$ and thus, for any $\epsilon \in (0, -d_\phi)$, the mapping $u \mapsto \tilde{\phi}(u) = \phi(u - \mathbf{d}_\epsilon) \in \mathcal{B}$ since $\tilde{\phi}(0) = \phi(-\mathbf{d}_\epsilon) \geq 0$ and clearly $\tilde{\phi}'$ is completely monotone. Hence,

$$\phi_{\mathbf{d}_\epsilon}(u) = \frac{u}{u + \mathbf{d}_\epsilon} \phi(u) = \frac{u}{u + \mathbf{d}_\epsilon} \tilde{\phi}(u + \mathbf{d}_\epsilon) \in \mathcal{B},$$

see Proposition 4.1.1(9). On the other hand, as, for all $n \geq 0$,

$$W_{\phi_{\mathbf{d}_\epsilon}}(n+1) = \prod_{k=1}^n \frac{k}{k + \mathbf{d}_\epsilon} \phi(k) = \frac{\Gamma(\mathbf{d}_\epsilon + 1)}{\Gamma(n + \mathbf{d}_\epsilon + 1)} n! W_\phi(n+1),$$

we deduce that

$$\mathcal{V}_\psi \mathcal{I}_{\phi_{\mathbf{d}_\epsilon}} p_n(x) = \frac{n!}{W_{\phi_{\mathbf{d}_\epsilon}}(n+1)} W_\phi(n+1) p_n(x) = \frac{\Gamma(n + \mathbf{d}_\epsilon + 1)}{\Gamma(\mathbf{d}_\epsilon + 1)} p_n(x) = \mathcal{E}_{\mathbf{d}_\epsilon} p_n(x).$$

The latter by moment determinacy completes the proof of (10.10) since we conclude that

$$\varepsilon_{\mathbf{d}_\epsilon}(x) = \frac{x^{\mathbf{d}_\epsilon}}{\Gamma(\mathbf{d}_\epsilon + 1)} e^{-x} = \int_0^\infty \nu\left(\frac{x}{y}\right) \iota_{\mathbf{d}_\epsilon}(y) \frac{dy}{y}.$$

All other claims, by a Fubini argument, follow via the classical decomposition $f = f_+ - f_-$ provided that $\mathcal{E}_{\mathbf{d}_\epsilon}|f|(x) < \infty$, for some $x > 0$. The case $d_\phi = 0$ is the same. \square

LEMMA 10.1.3. *Let $\psi \in \mathcal{N}_P$ with $\overline{\overline{\Pi}}(0^+) < \infty$. Then $-d_\phi \in [0, \frac{m}{\sigma^2}]$ with $-d_\phi = \frac{m}{\sigma^2}$ only when $\overline{\overline{\Pi}}(0^+) = 0$. Moreover, if $-d_\phi > 0$, then for any $0 < \epsilon$ such that $\mathbf{d}_\epsilon = -d_\phi - \epsilon > 0$, with the notation of Lemma 10.1.2, $\mathcal{I}_{\phi_{\mathbf{d}_\epsilon}} \in \mathbf{B}(\mathbf{L}^2(\nu), \mathbf{L}^2(\varepsilon_{\mathbf{d}_\epsilon}))$ and, the following intertwining*

$$(10.12) \quad \mathcal{I}_{\phi_{\mathbf{d}_\epsilon}} Q_t^{(\mathbf{d}_\epsilon)} = P_t \mathcal{I}_{\phi_{\mathbf{d}_\epsilon}}$$

holds on $L^2(\varepsilon_{\mathbf{d}_\epsilon})$ for all $t \geq 0$, where $Q^{(\mathbf{d}_\epsilon)} = (Q_t^{(\mathbf{d}_\epsilon)})_{t \geq 0}$ is the classical Laguerre semigroup of order \mathbf{d}_ϵ , see Chapter 3. Consequently, $(\mathfrak{c}_n^{-\frac{1}{2}}(\mathbf{d}_\epsilon)\mathcal{P}_n)_{n \geq 0}$, is a Bessel sequence in $L^2(\nu)$.

Proof. The fact that $-d_\phi \in [0, \frac{m}{\sigma^2}]$ follows directly from Proposition 4.2.2(1) and $\phi > 0$ on $(d_\phi, 0)$, see (1.19). By using the factorization (10.10) and following a line of reasoning similar to the proof of Lemma 10.1.1 and Theorem 7.0.3 we settle easily the claims. \square

10.1.2. The large perturbation case and beyond.

LEMMA 10.1.4. *Let $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$ with $(\alpha, \mathbf{m}) \in \mathfrak{R}$ and $\Phi_{\alpha, \mathbf{m}}(u) = \frac{\phi(u)}{\phi_{\alpha, \mathbf{m}}^R(u)}$.*

(1) *Then, there exists a Markov operator $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}$ associated to the variable $V_{\Phi_{\alpha, \mathbf{m}}}$, such that the factorization of operators*

$$\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} \mathcal{V}_{\phi_{\alpha, \mathbf{m}}^R} = \mathcal{V}_\psi$$

holds. Moreover, $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} \in \mathbf{B}(L^2(\nu), L^2(\mathbf{e}_{\alpha, \mathbf{m}}))$ and the following intertwining relation

$$P_t^{(\alpha, \mathbf{m})} \mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} = \mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} P_t$$

holds on $L^2(\nu)$ for all $t \geq 0$, where $P^{(\alpha, \mathbf{m})} = (P_t^{(\alpha, \mathbf{m})})_{t \geq 0}$ is the Gauss-Laguerre semigroup described in Example 3.0.3.

(2) *Consequently, we have the following estimate, for large n ,*

$$(10.13) \quad \|\mathcal{V}_n\|_\nu = \mathcal{O}(e^{nT_\alpha}).$$

In particular, if $\psi \in \mathcal{N}_P$, then for any $\epsilon > 0$ and large n ,

$$(10.14) \quad \|\mathcal{V}_n\|_\nu = \mathcal{O}(e^{\epsilon n}).$$

(3) *For any $(\alpha, \mathbf{m}) \in \mathfrak{R}$, $\mathcal{N}_{\alpha, \mathbf{m}} \subset \mathcal{N}_{\underline{\Theta}}$ with $\frac{\pi}{2}\alpha \leq \underline{\Theta}_\phi \leq \frac{\pi}{2}$.*

Proof. Since $\psi \in \mathcal{N}_{\alpha, \mathbf{m}}$, by definition, see Table 2 in Chapter 1, the mapping $u \mapsto \frac{1}{\Phi_{\alpha, \mathbf{m}}(u)}$ is completely monotone. Thus, according to [BD04, Theorem 1.3], the sequence $(a_n)_{n \geq 0}$, where $a_0 = 1$ and $a_n = \prod_{k=1}^n \Phi_{\alpha, \mathbf{m}}(k)$, $n \geq 1$, is a Stieltjes moment sequence, which entails that there exists a unique Markov operator $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}$ such that, for $n, x \geq 0$,

$$(10.15) \quad \begin{aligned} \mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} p_n(x) &= \mathbb{E} \left[V_{\Phi_{\alpha, \mathbf{m}}}^n \right] p_n(x) = \frac{W_\phi(n+1)\Gamma(\alpha\mathbf{m}+1)}{\Gamma(\alpha n + \alpha\mathbf{m}+1)} p_n(x) \\ &= \frac{W_\phi(n+1)}{W_{\phi_{\alpha, \mathbf{m}}^R}(n+1)} p_n(x) \end{aligned}$$

where we used for the last equality (10.4). From this characterization, by following a line of reasoning similar to the proof of Lemma 10.1.1, we easily derive the factorization of Markov operators, the continuity of $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}$, the intertwining relation as well as its dual version. Estimate (10.13) is deduced from the dual intertwining relation which yields in this case that

$$\|\mathcal{V}_n\|_\nu = \|\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}^* \mathcal{V}_n^{(\alpha, \mathbf{m})}\|_\nu \leq \|\mathcal{V}_n^{(\alpha, \mathbf{m})}\|_{\mathbf{e}_{\alpha, \mathbf{m}}} = \mathcal{O}(e^{nT_\alpha}),$$

where the last bound can be found in Example 3.0.3,(3.11). For the second item, since according to Proposition 4.2.2(2), we have for any $\psi \in \mathcal{N}_P$ with $\overline{\Pi}(0^+) = \infty$, that for any $0 < \alpha < 1$, there exists $\mathbf{m} > 0$ such that $\Phi_{\alpha, \mathbf{m}} = \frac{\phi}{\phi_{\alpha, \mathbf{m}}} \in \mathcal{B}$ and hence $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}$ is a Markov operator. We deduce (10.14) from (10.13) which holds for all $\alpha \in (0, 1)$ in this case and

$$\lim_{\alpha \uparrow 1} T_\alpha = \lim_{\alpha \uparrow 1} \ln(2^\alpha - 1) = 0.$$

The case $\psi \in \mathcal{N}_P$ with $\overline{\Pi}(0^+) < \infty$ was treated in Lemma 10.1.1. Finally, from (10.15), we deduce that, for any $z \in \mathbb{C}_{(d_{\Phi_{\alpha, \mathbf{m}}, \infty})}$ with $d_{\Phi_{\alpha, \mathbf{m}}} = d_\phi \vee (-\mathbf{m} - \frac{1}{\alpha})$,

$$\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}} p_z(x) = \frac{W_\phi(z+1)}{W_{\phi_{\alpha, \mathbf{m}}^R}(z+1)} p_z(x).$$

Thus, since $\mathcal{V}_{\Phi_{\alpha, \mathbf{m}}}$ is a Markov operator, we have for any $z = a + ib$ with $a > d_{\Phi_{\alpha, \mathbf{m}}}$,

$$\left| \frac{W_\phi(z+1)}{W_{\phi_{\alpha, \mathbf{m}}^R}(z+1)} \right| \leq \frac{W_\phi(a+1)}{W_{\phi_{\alpha, \mathbf{m}}^R}(a+1)}$$

or equivalently $|W_\phi(z+1)| \leq \frac{W_\phi(a+1)}{W_{\phi_{\alpha, \mathbf{m}}^R}(a+1)} |W_{\phi_{\alpha, \mathbf{m}}^R}(z+1)|$. Next, note that the expression (10.4) combined with the Stirling approximation (9.7) yield that, for any $a > -\mathbf{m} - \frac{1}{\alpha}$,

$$\left| W_{\phi_{\alpha, \mathbf{m}}^R}(a+ib+1) \right| \approx C_a |b|^{a+\alpha\mathbf{m}+\frac{1}{2}} e^{-\frac{\pi}{2}\alpha|b|}, C_a > 0,$$

and hence, from the condition (6.47), we get that $W_{\phi_{\alpha, \mathbf{m}}^R} \in \mathcal{N}_{\frac{\pi}{2}\alpha}$ which completes the proof as from Theorem 6.0.2(2b), we have $\Theta_\phi \leq \frac{\pi}{2}$. \square

10.2. Completeness of $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$

We shall take two different paths to prove the completeness of $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$. The first one, when $\psi \in \mathcal{N}_P$, i.e. $\sigma^2 > 0$, is based on the factorization identity (7.18) which allows to derive a non-trivial injectivity property of the operator \mathcal{V}_ψ in the weighted Hilbert space $L^2(\nu)$. As we can not show this property beyond this case, we resort to another approach for the case $\psi \in \mathcal{N}_R \setminus \mathcal{N}_P$, for which an analytical extension property for the Mellin transform of the invariant density is needed. However, we stress that both approaches stem from the concept of reference semigroup as they require the precise estimate of $\|\mathcal{V}_n\|_\nu$ derived in the previous chapter.

10.2.1. The case $\psi \in \mathcal{N}_P$. We split the proof into several intermediate results. We start with a general statement, that will also be used later in Chapter 12, which provides, in particular, pointwise bounds for $|w_n(y)|$ which depend solely on the analyticity of ν .

PROPOSITION 10.2.1. (1) *Let $\nu \in \mathcal{A}(\Theta)$ with $\Theta \in (0, \frac{\pi}{2}]$, i.e. ν is analytic in $\mathbb{C}(\Theta) = \{z \in \mathbb{C}; |\arg z| < \Theta\}$. Then, writing $\mathfrak{z} = \frac{1}{1-z}$ and $T_\Theta = -\ln \sin \Theta$, we have, for any $y > 0$, that*

$$(10.16) \quad \mathfrak{z} d_3 \nu(y) = \frac{1}{1-z} \nu\left(\frac{y}{1-z}\right) = \sum_{n=0}^{\infty} w_n(y) z^n, \quad |z| < e^{-T_\Theta},$$

where we recall that $d_c f(x) = f(cx)$. Moreover, for any $t > T_\Theta$, $y > 0$, there exists $F(y, t) > 0$ such that, for any $n \geq 1$,

$$(10.17) \quad |w_n(y)| \leq F(y, t)e^{tn}.$$

Finally, as a function $(y, t) \mapsto F(y, t)$ is locally uniformly bounded on $\mathbb{R}_+ \times (T_\Theta, \infty)$.

(2) Let $\psi \in \mathcal{N}_P$. Then, for any $f \in L^2(\nu)$ and any $|z| < 1$,

$$\mathcal{V}_\psi f(1-z) = \sum_{n=0}^{\infty} \langle f, \mathcal{V}_n \rangle_\nu z^n.$$

Proof. Let $0 < \theta < \Theta$ and let $\mathcal{C} = \mathcal{C}_y(y \sin \theta)$ be the circle with centre $y > 0$ and radius $y \sin \theta < y \sin \Theta$ enclosing the ball $\mathring{\mathcal{C}}$. Then $\mathcal{C} \subset \mathbb{C}(\Theta)$. Thus, for any $y > 0$ and z , we have $\mathfrak{z}y \in \mathring{\mathcal{C}}$ if and only if

$$|\mathfrak{z}z| = \frac{z}{1-z} < \sin \theta.$$

Choosing such z , for any $y > 0$, using twice Cauchy's theorem over \mathcal{C} , we get, for any $M \in \mathbb{N}$, that

$$\begin{aligned} \sum_{n=0}^M w_n(y) z^n &= \sum_{n=0}^M (-1)^n \frac{(y^n \nu(y))^{(n)}}{n!} z^n = \frac{1}{2\pi i} \sum_{n=0}^M \left(\oint \frac{\zeta^n \nu(\zeta)}{(\zeta - y)^{n+1}} d\zeta \right) z^n \\ &= \frac{\mathfrak{z}}{2\pi i} \oint \left(1 - \left(\frac{z\zeta}{\zeta - y} \right)^{M+1} \right) \frac{\nu(\zeta)}{\zeta - \mathfrak{z}y} d\zeta \\ &= \mathfrak{z} d_{\mathfrak{z}} \nu(y) - \frac{\mathfrak{z}}{2\pi i} \oint \left(\frac{z\zeta}{\zeta - y} \right)^{M+1} \frac{\nu(\zeta)}{\zeta - \mathfrak{z}y} d\zeta, \end{aligned}$$

where

$$\frac{\mathfrak{z}}{2\pi i} \oint \frac{\nu(\zeta)}{\zeta - \mathfrak{z}y} d\zeta = \mathfrak{z} d_{\mathfrak{z}} \nu(y)$$

merely because $\mathfrak{z}y \in \mathring{\mathcal{C}}$. However, if furthermore, $|z| < \frac{\sin \theta}{1 + \sin \theta}$ then $\left| \frac{z\zeta}{\zeta - y} \right| < 1$, $\zeta \in \mathcal{C}$, and letting M tend to ∞ we conclude (10.16) for $|z| < \frac{\sin \theta}{1 + \sin \theta}$. Next, since $\nu \in \mathcal{A}(\Theta)$ then $\mathfrak{z} d_{\mathfrak{z}} \nu(y)$ is analytic for $|z| < e^{-T_\Theta}$, that is when $|\arg(y\mathfrak{z})| = |\arg(\mathfrak{z})| < \Theta$, and then the convergence of the series $\sum_{n=0}^{\infty} w_n(y) z^n$ extends on this region via (10.16). Then, (10.17) follows by trivially estimating the representation

$$|w_n(y)| = \frac{1}{2\pi} \left| \oint \frac{\mathfrak{z} d_{\mathfrak{z}} \nu(y)}{z^{n+1}} dz \right|,$$

over the circle $\mathcal{C}_0(e^{-t})$, $t > T_\Theta$. The final claim of the first item follows since $\nu \in \mathcal{A}(\Theta)$. When $\psi \in \mathcal{N}_P$, i.e. $\sigma^2 > 0$, (10.16) holds with $T_\Theta = 0$ because thanks to Theorem 6.0.2(2(b)i) and Lemma 9.3.1 we have that $\nu \in \mathcal{A}(\frac{\pi}{2})$, that is ν is analytic in $\mathbb{C}_{(0, \infty)}$. Moreover, the bound (10.14) yields that, for any $f \in L^2(\nu)$ and any $\epsilon > 0$,

$$\int_0^\infty |f(y)| |w_n(y)| dy = \langle |f|, |\mathcal{V}_n| \rangle_\nu \leq \|f\|_\nu \|\mathcal{V}_n\|_\nu \leq C e^{\epsilon n}.$$

Thus, we get by an application of Fubini's Theorem that, for any $|z| < 1$,

$$\begin{aligned} \mathcal{V}_\psi f(1-z) &= \int_0^\infty f(y) \nu(zy) dy \\ &= \int_0^\infty f(y) \sum_{n=0}^\infty w_n(y) z^n dy \\ &= \sum_{n=0}^\infty \langle f, \mathcal{V}_n \rangle_\nu z^n. \end{aligned}$$

□

Before proving the main ingredient for the completeness of the sequence $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$ we collect additional information.

PROPOSITION 10.2.2. *Let $\psi \in \mathcal{N}_P$ and fix $x \in (0, 2\sigma^2)$. If $d_\phi = 0$ then for any $f \in L^2(\nu)$, $\mathcal{E}_{\mathbf{d}^*}|f|(x) < \infty$ and if $d_\phi < 0$ then, with $\mathbf{d}_\epsilon = -d_\phi - \epsilon$, $\mathcal{E}_{\mathbf{d}_\epsilon}|f|(x) < \infty$ if $\epsilon \in (0, \frac{1-d_\phi}{3})$.*

Proof. Let $\psi \in \mathcal{N}_P$, $f \in L^2(\nu)$ and choose $\mathbf{d}^* \geq 0$. From (10.11)

$$(10.18) \quad \mathcal{E}_{\mathbf{d}^*}|f|\left(\frac{1}{x}\right) = \int_0^\infty |f(y)| y^{\mathbf{d}^*} e^{-\frac{y}{x}} dy < \infty \iff \mathcal{E}_{\mathbf{d}^*}|f|(x) < \infty.$$

Let $0 < \frac{1}{x} = \frac{1}{2\sigma^2} + h$, for some $h > 0$, that is $x \in (0, 2\sigma^2)$. Then

$$\left(\int_1^\infty |f(y)| y^{\mathbf{d}^*} e^{-\frac{y}{x}} dy \right)^2 \leq \int_1^\infty f^2(y) e^{-\frac{y}{\sigma^2} - hy} dy \int_1^\infty y^{2\mathbf{d}^*} e^{-hy} dy < \infty$$

since from (5.8) $\nu(y) \asymp e^{-\frac{y}{\sigma^2} + o(y)}$ and $f \in L^2(\nu)$. From $\nu > 0$ on \mathbb{R}_+ , see Theorem 5.0.2(1),

$$\left(\int_0^1 y^{\mathbf{d}^*} |f(y)| dy \right)^2 \leq \int_0^1 f^2(y) \nu(y) dy \int_0^1 \frac{y^{2\mathbf{d}^*}}{\nu(y)} dy \leq \|f\|_\nu \int_0^1 \frac{y^{2\mathbf{d}^*}}{\nu(y)} dy.$$

Since $\nu(y) \geq C_{\underline{a},1} y^{-\underline{a}}$ where $\underline{a} < d_\phi$, see (5.5) of Theorem 5.0.4, then

$$\frac{y^{2\mathbf{d}^*}}{\nu(y)} \leq C_{\underline{a},1}^{-1} y^{2\mathbf{d}^* + \underline{a}}.$$

Putting $\mathbf{d}^* = \mathbf{d}_\epsilon = -d_\phi - \epsilon$ the last integral is finite if $\epsilon \in (0, \frac{1-d_\phi}{3})$ and \underline{a} is close to $-d_\phi$. □

LEMMA 10.2.3. *Let $\psi \in \mathcal{N}_P$ and $f \in L^2(\nu)$. If $\mathcal{V}_\psi f(x) = 0$ for any $x \in (0, 2)$ then, for a.e. $y > 0$, $f(y) = 0$.*

Proof. Let $\psi \in \mathcal{N}_P$ with $\mathbf{d}^* = 0$ when $d_\phi = 0$ and $\mathbf{d}^* = \mathbf{d}_\epsilon = -d_\phi - \epsilon$, for some $\epsilon \in (0, \frac{1-d_\phi}{3})$, when $d_\phi < 0$. Next, fix $x < 2\sigma^2$. Therefore, from Proposition 10.2.2 we get that $\mathcal{E}_{\mathbf{d}^*}|f|(x) < \infty$. Recall that $\phi_{\mathbf{d}^*}(u) = \frac{u}{u+\mathbf{d}^*} \phi(u) \in \mathcal{B}$, see Lemma 10.1.2. Since $\Psi \in \mathcal{N}_P$ then $\phi(u) \asymp \sigma^2 u$, see Proposition 4.1.1(3). Hence, $\phi_{\mathbf{d}^*}(u) \asymp \sigma^2 u$ and

we conclude that $\sigma_{\mathfrak{d}^*}^2 = \sigma^2$. Thus, from $\mathcal{E}_{\mathfrak{d}^*}|f|(x) < \infty$ then (10.10) is applicable and yields

$$\mathcal{E}_{\mathfrak{d}^*} f(x) = \int_0^{\frac{1}{\sigma^2}} \mathcal{V}_\psi f(xy) \iota_{\phi_{\mathfrak{d}^*}}(y) dy = 0,$$

since $\iota_{\mathfrak{d}^*}$ has support on $[0, \sigma^{-2}]$, see Proposition 6.1.1(1), and by assumption $\mathcal{V}_\psi f(xy) = 0$ when $xy < 2\sigma^2 \frac{1}{\sigma^2} = 2$. Therefore, from (10.11), for all $x > \frac{1}{2\sigma^2}$, $\mathcal{E}_{\mathfrak{d}^*} f(x) = 0$. Choose $x_0 > \frac{1}{2\sigma^2}$. The same arguments as in the proof of Proposition 10.2.2 for $f \in L^2(\nu)$ lead to

$$\tilde{f}(y) = y^{\mathfrak{d}^*} e^{-x_0 y} f(y) \in L^2(\mathbb{R}_+).$$

Since for all $x > x_0$, $\mathcal{L}\tilde{f}(x - x_0) = \mathcal{E}_{\mathfrak{d}^*} f(x) = 0$, the injectivity property of the Laplace transform \mathcal{L} entails that for a.e. $y > 0$, $f(y) = \tilde{f}(y) = 0$, which completes the proof. \square

Next we finish the proof of the completeness of the sequence $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$. Assume that $f \in L^2(\nu)$ and $\langle f, \mathcal{V}_n \rangle_\nu = 0, \forall n \geq 0$. From Proposition 10.2.1 we deduce that $\mathcal{V}_\psi f = 0$ on $(0, 2)$. The statement follows by invoking Lemma 10.2.3 which gives $f = 0$ a.e. and the fact that in a separable Hilbert space, the concept of total sequence and complete sequence coincide.

10.2.2. The case $\psi \in \mathcal{N}_{\alpha, \mathfrak{m}}^{d_\phi} \setminus \mathcal{N}_P$. Assume that $\psi \in \mathcal{N}_{\alpha, \mathfrak{m}}^{d_\phi} \setminus \mathcal{N}_P$, i.e. $\sigma^2 = 0$ and from (1.34), $\psi \in \mathcal{N}_{\alpha, \mathfrak{m}}$ with $d_\phi < 1 - \frac{\mathfrak{m}}{2} - \frac{1}{2\alpha}$. Recall that $\mathbf{e}_{\alpha, \mathfrak{m}}(x) = \frac{x^{\mathfrak{m} + \frac{1}{\alpha} - 1} e^{-x \frac{1}{\alpha}}}{\Gamma(\alpha \mathfrak{m} + 1)}$, $x > 0$, see Example 3.0.3(3.10). To prove the completeness of the sequence $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$, we argue as follows using the notation and results of Lemma 10.1.4. First, note that from the identity $\mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^* \mathcal{V}_n^{(\alpha, \mathfrak{m})}(x) = \mathcal{V}_n(x)$, $n \geq 0$, then $(\mathcal{V}_n)_{n \geq 0} \subset \mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^*(L^2(\mathbf{e}_{\alpha, \mathfrak{m}}))$. Second, $\overline{\text{Span}(\mathcal{V}_n^{(\alpha, \mathfrak{m})})} = L^2(\mathbf{e}_{\alpha, \mathfrak{m}})$, see [PS17, Proposition 2.1(2)], and $\mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^* \in \mathbf{B}(L^2(\mathbf{e}_{\alpha, \mathfrak{m}}), L^2(\nu))$. Third, if $\overline{\text{Ran}(\mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^*)} = L^2(\nu)$ then classical approximation yields the completeness of $(\mathcal{V}_n)_{n \geq 0}$ in $L^2(\nu)$. It remains to prove the latter. By linearity and density of the set of polynomials in $L^2(\nu)$, it is furnished if there exist $F_n \in L^2(\mathbf{e}_{\alpha, \mathfrak{m}})$, for all $n \geq 0$, such that $\mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^* F_n(x) = p_n(x) = x^n$. We consider the following convolution equation

$$(10.19) \quad \widehat{\mathcal{V}}_{\Phi_{\alpha, \mathfrak{m}}} \widehat{F}_n(x) = \nu(x) p_n(x).$$

Following a reasoning similar to the proof of (8.9) of Lemma 8.1.1, we have, for a.e. $x > 0$ and integrable $f, \hat{f} = f e_{\alpha, \mathfrak{m}}$, that

$$(10.20) \quad \begin{aligned} \mathcal{V}_{\Phi_{\alpha, \mathfrak{m}}}^* f(x) &= \frac{1}{\nu(x)} \widehat{\mathcal{V}}_{\Phi_{\alpha, \mathfrak{m}}} f \mathbf{e}_{\alpha, \mathfrak{m}}(x) = \frac{1}{\nu(x)} \widehat{\mathcal{V}}_{\Phi_{\alpha, \mathfrak{m}}} \hat{f}(x) \\ &= \frac{1}{\nu(x)} \int_0^\infty f(xy) \mathbf{e}_{\alpha, \mathfrak{m}}(xy) \nu_{\alpha, \mathfrak{m}}(y) \frac{dy}{y}, \end{aligned}$$

with $\nu_{\alpha, \mathfrak{m}}$ the density of the variable $V_{\Phi_{\alpha, \mathfrak{m}}}$. We solve (10.19) from the perspective of Mellin distributions as described in Proposition 8.1.4. Henceforth, by means of the Mellin transform of (10.19) and the functional equation (6.6) satisfied by W_ϕ , that

is $W_\phi(z+1) = \phi(z)W_\phi(z)$, one gets, for $\Re(z) > d_\phi^{\alpha, \mathbf{m}} = \max(d_\phi, 1 - \frac{1}{\alpha} - \mathbf{m})$,

$$\mathcal{M}_{\hat{F}_n}(z) = \frac{W_\phi(z+n)}{W_\phi(z)} W_{\Phi_{\alpha, \mathbf{m}}}(z) = \prod_{j=0}^{n-1} \phi(z+j) \frac{\Gamma(\alpha z - \alpha + \alpha \mathbf{m} + 1)}{\Gamma(\alpha \mathbf{m} + 1)},$$

where $W_\phi(z+n) = \int_0^\infty x^{z-1} \nu(x) p_n(x) dx$ and

$$\int_0^\infty x^{z-1} \nu_{\alpha, \mathbf{m}}(x) dx = W_\phi(z) \frac{\Gamma(\alpha \mathbf{m} + 1)}{\Gamma(\alpha z - \alpha + \alpha \mathbf{m} + 1)},$$

see (10.15) for the latter. Thus, for any $\epsilon > 0$, $a > d_\phi^{\alpha, \mathbf{m}}$ and large $|b|$, since $|\phi(z)| = O(|z|)$ for $|z| = |a + ib|$, $|b|$ large enough and a fixed, see Proposition 4.1.1(3), we have that

$$\begin{aligned} |\mathcal{M}_{\hat{F}_n}(a+ib)| &\leq |\phi(a+n+ib)|^n |\Gamma(\alpha(a+ib) - \alpha + \alpha \mathbf{m} + 1)| \\ &\leq C_{a,n} |b|^n e^{-\frac{\pi}{2}(\alpha-\epsilon)|b|} \end{aligned}$$

for some $C_{a,n} > 0$. Hence, by the Parseval identity for the Mellin transform, see (5.42), we have that there exists $\hat{F}_n \in L^2(\mathbb{R}_+)$ a solution to the convolution equation (10.19) above. From (10.20) it remains to show that, for any $n \geq 0$, $F_n = \frac{\hat{F}_n}{\mathbf{e}_{\alpha, \mathbf{m}}} \in L^2(\mathbf{e}_{\alpha, \mathbf{m}})$. However, by Mellin inversion, we get, for all $x > 0$ and $a > d_\phi^{\alpha, \mathbf{m}}$, that

$$\hat{F}_n(x) = \frac{1}{2\pi i} \int_{a-i\infty}^{a+i\infty} x^{-z} \frac{W_\phi(z+n)}{W_\phi(z)} \frac{\Gamma(\alpha z - \alpha + \alpha \mathbf{m} + 1)}{\Gamma(\alpha \mathbf{m} + 1)} dz.$$

Thus, for all $x > 0$, $a > d_\phi^{\alpha, \mathbf{m}}$ and $\epsilon < \alpha$,

$$(10.21) \quad |\hat{F}_n(x)| \leq C x^{-a} \int_{-\infty}^{\infty} e^{\epsilon|b|} |\Gamma(\alpha ib + (a-1)\alpha + \alpha \mathbf{m} + 1)| db \leq C_a x^{-a}.$$

Moreover, after an obvious change of variables in [PK01, Lemma 2.6] we get that

$$|\hat{F}_n(x)| \leq C x^M e^{-\cos(\frac{\epsilon}{\alpha})x^{\frac{1}{\alpha}}}$$

for some $M > 0$, ϵ small and all $x \geq 1$. From this estimate and the notation

$$\mathbf{e}_{\alpha, \mathbf{m}}(x) = \frac{x^{\mathbf{m} + \frac{1}{\alpha} - 1} e^{-x^{\frac{1}{\alpha}}}}{\Gamma(\alpha \mathbf{m} + 1)}, \quad x > 0,$$

$$\frac{\hat{F}_n(x)}{\mathbf{e}_{\alpha, \mathbf{m}}(x)} \mathbb{1}_{\{x \geq 1\}} \in L^2(\mathbf{e}_{\alpha, \mathbf{m}})$$

if $\epsilon < \frac{\alpha\pi}{3}$. For $\frac{\hat{F}_n(x)}{\mathbf{e}_{\alpha, \mathbf{m}}(x)} \mathbb{1}_{\{x \leq 1\}} \in L^2(\mathbf{e}_{\alpha, \mathbf{m}})$ consider (10.21) with ϵ small enough that $a = 1 - \frac{\mathbf{m}}{2} - \frac{1}{2\alpha} - \epsilon > d_\phi^{\alpha, \mathbf{m}}$. Then $\hat{F}_n^2(x) \mathbf{e}_{\alpha, \mathbf{m}}^{-1}(x) \leq x^{-1+\epsilon}$, $x \leq 1$, and thus, we deduce that $\frac{\hat{F}_n}{\mathbf{e}_{\alpha, \mathbf{m}}} \in L^2(\mathbf{e}_{\alpha, \mathbf{m}})$, which completes the proof.

Hilbert sequences, intertwining and spectrum

This part, which ends by the proof of Theorem 1.5.1, aims at establishing some interesting and new connections between three different concepts: intertwining relation, Hilbert sequences arising in non-harmonic analysis and spectrum of non self-adjoint operators. We present and prove several results, sometimes in a slightly more general context than the one of the current work. For sake of clarity, we proceed by recalling a few definitions that were introduced in Section 11, concerning the spectrum of linear operators. First, a complex number $\lambda \in S(P)$ belongs to the spectrum of the linear operator $P \in \mathbf{B}(L^2(\nu))$, if $P - \lambda\mathbf{I}$ does not have an inverse in $L^2(\nu)$ with the following three distinctions. $\lambda \in S_p(P)$, the point spectrum, if $\text{Ker}(P - \lambda\mathbf{I}) \neq \{0\}$, $\lambda \in S_c(P)$, the continuous spectrum, if $\text{Ker}(P - \lambda\mathbf{I}) = \{0\}$ and $\overline{\text{Ran}}(P - \lambda\mathbf{I}) = L^2(\nu)$ but $\text{Ran}(P - \lambda\mathbf{I}) \subsetneq L^2(\nu)$, and, finally $\lambda \in S_r(P)$, the residual spectrum, if $\text{Ker}(P - \lambda\mathbf{I}) = \{0\}$ and $\overline{\text{Ran}}(P - \lambda\mathbf{I}) \subsetneq L^2(\nu)$. Clearly,

$$S(P) = S_p(P) \cup S_c(P) \cup S_r(P).$$

Let $\lambda \in S_p(P)$ be an isolated eigenvalue. Then its geometric multiplicity, denoted by $\mathfrak{M}_g(\lambda, P)$ is computed as follows

$$(11.1) \quad \mathfrak{M}_g(\lambda, P) = \dim \text{Ker}(P - \lambda I),$$

that is the dimension of the corresponding eigenspace. Its algebraic multiplicity, denoted by $\mathfrak{M}_a(\lambda, P)$ is defined by

$$(11.2) \quad \mathfrak{M}_a(\lambda, P) = \dim \bigcup_{k=1}^{\infty} \text{Ker}(P - \lambda I)^k.$$

Note that always $\mathfrak{M}_g(\lambda, P) \leq \mathfrak{M}_a(\lambda, P)$. Next, keeping notation similar to the rest of the paper, let us assume that there exists $\mathcal{I} \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, such that for any $f \in L^2(\nu)$,

$$(11.3) \quad P\mathcal{I}f = \mathcal{I}Qf,$$

where $P \in \mathbf{B}(L^2(\nu))$ and $Q \in \mathbf{B}(L^2(\varepsilon))$ is self-adjoint. Moreover, we suppose that $S(Q) = S_p(Q) = (\lambda_n \in \mathbb{C})_{n \geq 0}$ and write $(\mathcal{L}_n)_{n \geq 0}$ for the associated sequence of eigenfunctions, with $\overline{\text{Span}}(\mathcal{L}_n) = L^2(\varepsilon)$. Although most of the results presented below may also hold in a more general settings, for sake of clarity, we focus on the conditions of the present work. We say that \mathcal{I} is an intertwining operator. It is immediate that the adjoint intertwining relation holds, that is with $\mathcal{I}^* \in \mathbf{B}(L^2(\nu), L^2(\varepsilon))$ the adjoint of \mathcal{I} , we have that, for any $g \in L^2(\nu)$

$$(11.4) \quad \mathcal{I}^* P^* g = Q \mathcal{I}^* g$$

where $P^* \in \mathbf{B}(L^2(\nu))$ stands for the adjoint of P . There is a substantial literature devoted to the study of intertwining relations. A natural problem is to understand how the spectral properties of an operator are preserved under such type of transformation. In our context, this issue is still unclear since without any additional assumptions one can find examples where the spectrum of one operator may not intersect the spectrum of the other one. A natural requirement is that the intertwining operator \mathcal{I} is an affinity, that is a bounded operator admitting a bounded inverse. In such case, Q and P are called similar and the two spectra coincide. In this direction, we mention the recent paper of Inoue and Trapani [IT14] where it is proved that a closed operator admits a non-self-adjoint resolutions of the identity if and only if it is similar to a self-adjoint operator. There exists an intermediate notion, called quasi-similarity, that is when $\mathcal{I} \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ is one-to-one and has a dense range, which was first introduced by Sz.-Nagy and Foias, see [SNF67], in their theory considering an infinite dimensional analogue of the Jordan form for certain classes of operators; it replaces the familiar notion of similarity which is the appropriate equivalence relation to use with finite dimensional Jordan forms. However, two operators can be quasi-similar and yet have unequal spectra, see e.g. [Hoo72]. For normal operators this cannot happen: it follows from the Putnam-Fuglede commutativity theorem that if two normal operators are quasisimilar, they are actually unitary equivalent, see [Dou69, Lemma 4.1], and therefore have equal spectra. Finally, we refer to [AT13] for a recent account of similar and quasi-similar operators.

We are now ready to state the following relations between point spectra and multiplicity of eigenvalues, recalling that $S(Q) = S_p(Q) = (\lambda_n)_{n \geq 0}$ and we write $(\mathcal{L}_n)_{n \geq 0}$ the associated set of eigenfunctions.

PROPOSITION 11.0.1. *Assume that the intertwining relation (11.3) holds with $\mathcal{I} \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$.*

- (1) *If for some $n \geq 0$, $\mathcal{L}_n \notin \text{Ker}(\mathcal{I})$, then $\lambda_n \in S_p(P)$.*
- (2) *Finally, if \mathcal{I} is one-to-one with $\mathcal{I}\mathbf{1} \equiv \mathbf{1}$ then $S_p(Q) \subseteq S_p(P)$ and for any $n \geq 0$, $\mathfrak{M}_g(Q, \lambda_n) \leq \mathfrak{M}_g(P, \lambda_n)$ and $\mathfrak{M}_a(Q, \lambda_n) \leq \mathfrak{M}_a(P, \lambda_n)$.*

REMARK 11.0.2. Note that it is possible that $S_p(Q) \subsetneq S_p(P)$. Indeed there may exist $P_n \in L^2(\nu) \setminus \text{Ran}(\mathcal{I})$ an eigenfunction of P associated to the eigenvalue $\bar{\lambda}_n \in \mathbb{C} \setminus S_p(Q)$.

Proof. Since by assumption $\mathcal{P}_n = \mathcal{I}\mathcal{L}_n \neq 0$, we deduce from (11.3) together with the linearity of \mathcal{I} that

$$P\mathcal{P}_n = \mathcal{I}Q\mathcal{L}_n = \lambda_n\mathcal{I}\mathcal{L}_n = \lambda_n\mathcal{P}_n,$$

which completes the proof of the first statement. Both facts $S_p(Q) \subseteq S_p(P)$ and for any $n \geq 0$, $\mathfrak{M}_g(Q, \lambda_n) \leq \mathfrak{M}_g(P, \lambda_n)$ follow by a similar line of reasoning and employing the fact that \mathcal{I} is one-to-one. Finally, if for some $n \in \mathbb{N}$, there exists an integer p and $L_n \neq 0 \in L^2(\varepsilon)$ such that

$$(Q - \lambda_n\mathbf{1})^{p+1}L_n = 0 \text{ and } (Q - \lambda_n\mathbf{1})^pL_n \neq 0$$

then, since \mathcal{I} is one-to-one with $\mathcal{I}\mathbf{1} \equiv \mathbf{1}$,

$$(P - \lambda_n \mathbf{1})^{p+1} \mathcal{I}L_n = 0 \text{ and } (P - \lambda_n \mathbf{1})^p \mathcal{I}L_n \neq 0,$$

which completes the proof of the fact that $\mathfrak{M}_a(Q, \lambda_n) \leq \mathfrak{M}_a(P, \lambda_n)$. \square

The next result discusses the consequence of the completeness property of a set of eigenfunctions on the residual spectrum and on the (geometric and algebraic) multiplicity of the eigenvalues of the adjoint operator. Note that the proof of the latter result is based on a substantial property of biorthogonal sequences in Hilbert space that we had recalled above, namely that a complete sequence in a separable Hilbert space admits at most one biorthogonal sequence. We also recall that a sequence that admits a biorthogonal sequence is said to be minimal and a sequence that is both minimal and complete, in the sense that its linear span is dense in $L^2(\nu)$, will be called exact.

PROPOSITION 11.0.3. *Let us assume that there exists a sequence of eigenfunctions $(\mathcal{P}_n)_{n \geq 0}$ of P associated to the sequence of distinct eigenvalues $(\lambda_n)_{n \geq 0}$ such that $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$.*

(1) $S_r(P) = \emptyset$. If in addition $(\mathcal{P}_n)_{n \geq 0}$ is minimal, then for all $n \geq 0$, $\lambda_n \in S_p(P^*)$ and

$$\mathfrak{M}_a(\lambda_n, P^*) = \mathfrak{M}_g(\lambda_n, P^*) = 1.$$

(2) If for all n , $\mathcal{I}\mathcal{L}_n = \mathcal{P}_n$ and for some \bar{n} , $\mathcal{L}_{\bar{n}} \notin \text{Ran}(\mathcal{I}^*)$ then $\lambda_{\bar{n}} \in S_r(P^*)$.

Proof. Let us assume that there exists a complex number $q \in S_r(P)$, that is $P - q\mathbf{1}$ is one-to-one but does not have a dense range. Since $q \notin S_p(P)$, we have for all $n \geq 0$,

$$(P - q\mathbf{1})\mathcal{P}_n = (\lambda_n - q)\mathcal{P}_n \neq 0,$$

which yields to a contradiction as $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$ and proves the first claim. Next, for any $n, m \geq 0$, using that $(\mathcal{V}_n)_{n \geq 0}$ is biorthogonal to $(\mathcal{P}_n)_{n \geq 0}$ and that $\mathcal{V}_n \neq 0$, we have that

$$(11.5) \quad \langle P^*\mathcal{V}_n - \lambda_n \mathcal{V}_n, \mathcal{P}_m \rangle_\nu = \langle \mathcal{V}_n, P\mathcal{P}_m \rangle_\nu - \lambda_n \delta_{nm} = (\lambda_m - \lambda_n) \delta_{nm} = 0.$$

That is

$$P^*\mathcal{V}_n - \lambda_n \mathcal{V}_n \in \text{Span}(\mathcal{P}_n)^\perp = \{0\}$$

since by assumption $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$. Thus, we deduce that $\lambda_n \in S_p(P^*)$. Assume now that there exists $n_0 \in \mathbb{N}$ such that $\mathfrak{M}_g(\lambda_{n_0}, P^*) = 2$. That is there exists $\bar{v}_{n_0} \in L^2(\nu)$ such that $\bar{v}_{n_0} \neq 0$, $\bar{v}_{n_0} \neq \mathcal{V}_{n_0}$ and $P^*\bar{v}_{n_0} = \lambda_{n_0} \bar{v}_{n_0}$. Thus, for all $m \geq 0$, we have

$$(11.6) \quad \lambda_{n_0} \langle \mathcal{P}_m, \bar{v}_{n_0} \rangle_\nu = \langle \mathcal{P}_m, P^*\bar{v}_{n_0} \rangle_\nu = \langle P\mathcal{P}_m, \bar{v}_{n_0} \rangle_\nu = \lambda_m \langle \mathcal{P}_m, \bar{v}_{n_0} \rangle_\nu$$

that is for some $C \in \mathbb{R}$, $\langle \mathcal{P}_m, \bar{v}_{n_0} \rangle_\nu = C \delta_{mn_0}$. Note that we may choose \bar{v}_{n_0} such that $C = 1$. Indeed as P^* is linear, $C \neq 0$ because otherwise $\bar{v}_{n_0} = 0$ as $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$. Thus the sequence $(\tilde{v}_n)_{n \geq 0}$ defined by $\tilde{v}_{n_0} = \frac{1}{2}v_{n_0} + \frac{1}{2}\bar{v}_{n_0}$ and otherwise $\tilde{v}_n = v_n$ is another biorthogonal sequence to $(\mathcal{P}_n)_{n \geq 0}$. However, as mentioned before the statement, this is impossible as $\overline{\text{Span}}(\mathcal{P}_n) = L^2(\nu)$ which implies the uniqueness of the biorthogonal sequence. Hence for all $n \geq 0$,

$$\mathfrak{M}_g(\lambda_n, P^*) = 1.$$

Next, let $n \geq 0$ and note that if $\dim \text{Ker}(P^* - \lambda_n \mathbf{1})^2 = 1$ then

$$\mathfrak{M}_a(\lambda_n, P^*) = \lim_{k \rightarrow \infty} \dim \text{Ker}(P^* - \lambda_n \mathbf{1})^k = \dim \text{Ker}(P^* - \lambda_n \mathbf{1}) = 1.$$

Indeed, assume that there exists $V_n \in L^2(\nu)$ such that $(P^* - \lambda_n \mathbf{1})^3 V_n = 0$, then necessarily $(P^* - \lambda_n \mathbf{1})V_n = 0$, that is $V_n \in \text{Ker}(P^* - \lambda_n \mathbf{1})$, since the converse inclusion always holds, this gives the statement for $k = 3$. A recurrence argument yields the claim for all $k \geq 2$. Hence it remains to show that

$$\text{Ker}(P^* - \lambda_n \mathbf{1})^2 = \text{Ker}(P^* - \lambda_n \mathbf{1}).$$

To this end, assume that there exists

$$V_n \neq \mathcal{V}_n \in \text{Ker}(P^* - \lambda_n \mathbf{1})^2,$$

that is since we have

$$\begin{aligned} \dim \text{Ker}(P^* - \lambda_n \mathbf{1}) &= \mathfrak{M}_g(\lambda_n, P^*) = 1, \\ (P^* - \lambda_n \mathbf{1})V_n &= CV_n \end{aligned}$$

for some real constant $C \neq 0$. Thus

$$\langle P^* V_n, \mathcal{P}_n \rangle_\nu - \lambda_n \langle V_n, \mathcal{P}_n \rangle_\nu = C \langle \mathcal{V}_n, \mathcal{P}_n \rangle_\nu.$$

Using the eigenfunction property of \mathcal{P}_n and the biorthogonality property, we get

$$0 = \lambda_n \langle V_n, \mathcal{P}_n \rangle_\nu - \lambda_n \langle V_n, \mathcal{P}_n \rangle_\nu = C \neq 0,$$

which completes the proof of the first item through a contradiction argument. To prove the second item we assume that $\lambda_{\bar{n}} \in S_p(P^*)$, that is there exists $\mathcal{V}_{\bar{n}} \in L^2(\nu)$ such that $P^* \mathcal{V}_{\bar{n}} = \lambda_{\bar{n}} \mathcal{V}_{\bar{n}}$. Next, since $\mathcal{I}\mathcal{L}_n = \mathcal{P}_n$, for all $n \geq 0$, we easily deduce that \mathcal{I} has dense range in $L^2(\nu)$ and thus \mathcal{I}^* is one-to-one in $L^2(\nu)$. Thus, we deduce from the intertwining relation (11.4) and Proposition 11.0.1 that $\mathcal{I}^* \mathcal{V}_{\bar{n}} = \mathcal{L}_{\bar{n}}$ which yields to a contradiction as it is assumed that $\mathcal{L}_{\bar{n}} \notin \text{Ran}(\mathcal{I}^*)$. Hence $\lambda_{\bar{n}} \in S_p(P) \setminus S_p(P^*)$ which implies that $\text{Ker}(P^* - \lambda_{\bar{n}} \mathbf{1}) = \{\emptyset\}$ with

$$\text{Ran}(P^* - \lambda_{\bar{n}} \mathbf{1})^\perp = \text{Ker}(P - \lambda_{\bar{n}} \mathbf{1}) \neq \{\emptyset\},$$

that is, $\lambda_{\bar{n}} \in S_r(P^*)$, which completes the proof. \square

We complete this part with the following result which provides, in the context of intertwining relation, a set of sufficient conditions for the characterization of the point and residual spectra. Note that this type of condition seems to be weaker than the one appearing in the intertwining literature. Indeed, we are able to extract this information without the assumption that either the intertwining operator has a bounded inverse or the linear operators P belong to some special classes of linear operators.

PROPOSITION 11.0.4. *Assume that the intertwining relation (11.3) holds with $\mathcal{I} \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, $\overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$ and assume that there exists a sequence $(\mathcal{P}_n)_{n \geq 0}$ of eigenfunctions for P associated to the set of distinct eigenvalues $(\lambda_n)_{n \geq 0}$. If in addition the sequence $(\mathcal{P}_n)_{n \geq 0}$ is exact and its biorthogonal sequence $(\mathcal{V}_n)_{n \geq 0}$ is complete, then*

$$S_p(P^*) = S_p(Q) = S_p(P), \quad S_r(P^*) = S_r(Q) = S_r(P) = \emptyset$$

and

$$\mathfrak{M}_a(\lambda_n, P) = \mathfrak{M}_g(\lambda_n, P) = \mathfrak{M}_a(\lambda_n, P^*) = \mathfrak{M}_g(\lambda_n, P^*) = 1.$$

Proof. Since $\overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$ and thus $\text{Ker}(\mathcal{I}_\phi^*) = \{\emptyset\}$, we get, by means of the intertwining relation (11.3), its adjoint version (11.4), the minimality of $(\mathcal{P}_n)_{n \geq 0}$ and Proposition 11.0.1 that

$$S_p(P^*) \subseteq S_p(Q) \subseteq S_p(P).$$

Moreover, as both sequences $(\mathcal{P}_n)_{n \geq 0}$ and $(\mathcal{V}_n)_{n \geq 0}$ are complete we have from Proposition 11.0.3 that

$$S_r(P) = S_r(P^*) = \emptyset.$$

Next, from the identity $\text{Ker}(P - \lambda \mathbf{1}) = \text{Ran}(P^* - \lambda \mathbf{1})^\perp$, we easily deduce the inclusions

$$S_r(P^*) \subseteq S_p(P) \subseteq S_r(P^*) \cup S_p(P^*),$$

which provides that $S_p(P^*) = S_p(P)$ and completes the proof of the first two claims. Finally the last one follows from Proposition 11.0.3. \square

11.1. Proof of Theorem 1.5.1

We recall that the item (1) was stated in Theorem 7.0.3(2) and proved in that Chapter. Let now $\psi \in \mathcal{N}_\infty$. Then, from Theorem 8.0.1 we have that, for all $n \geq 0$, $\mathcal{V}_n \in L^2(\nu)$ and the equation (8.4) in Proposition 8.0.4 entails that $\mathcal{L}_n = \mathcal{I}_\phi^* \mathcal{V}_n$. Hence, since from (7.24), $\mathcal{P}_m = \mathcal{I}_\phi \mathcal{L}_m$, $m \geq 0$, we get that

$$\langle \mathcal{P}_m, \mathcal{V}_n \rangle_\nu = \langle \mathcal{I}_\phi \mathcal{L}_m, \mathcal{V}_n \rangle_\nu = \langle \mathcal{L}_m, \mathcal{I}_\phi^* \mathcal{V}_n \rangle_\varepsilon = \langle \mathcal{L}_m, \mathcal{L}_n \rangle_\varepsilon = \delta_{nm},$$

which proves that $(\mathcal{V}_n)_{n \geq 0}$ is minimal. Its uniqueness property follows since the sequence $(\mathcal{P}_n)_{n \geq 0}$ is exact which itself is a direct consequence of item (1). We complete the proof of this item (2) by means of item (1) combined with the fact that a sequence which is biorthogonal to a Bessel sequence is a Riesz-Fischer sequence, see e.g. [CCLL02, Proposition 2.3(ii)]. Assume now that $\psi \in \mathcal{N}_P \cup \mathcal{N}_{\alpha, m}^{d_\phi}$. Then, the claims of item (3) follow readily from the fact that $\overline{\text{Span}(\mathcal{V}_n)} = L^2(\nu)$ which is found in Theorem 10.0.1(3). For the last item (4), assuming first that $\psi \in \mathcal{N}$, the inclusion $\mathcal{E}_t \subseteq S_p(P_t)$ is given in Theorem 7.0.3(1), whereas the last claim of (4a), i.e. $S_r(P_t) = \emptyset$, is deduced from Proposition 11.0.3(1). Theorem 8.0.1 (resp. Proposition 11.0.3(1)) gives the first (resp. second) claim of the item (4b). Item (4c) follows by combining the previous item (3) with Proposition 11.0.4. Finally, the last item (4d) is immediate from Theorem 8.1.4.

Proof of Theorems 1.3.1, 1.3.7 and 1.4.1

Let first $\psi \in \mathcal{N}$. Then the expressions of \mathcal{P}_n as stated in (1.21) have been derived in Theorem 7.0.3(1). Next, from the intertwining relationship (7.3) and the expansion (3.2) of the Laguerre semigroup of order $m = 0$, we get, in the $L^2(\nu)$ topology, that for any $g \in L^2(\varepsilon)$, $t > 0$,

$$P_t \mathcal{I}_\phi g = \mathcal{I}_\phi Q_t g = \mathcal{I}_\phi \sum_{n=0}^{\infty} e^{-nt} \langle g, \mathcal{L}_n \rangle_\varepsilon \mathcal{L}_n = \sum_{n=0}^{\infty} e^{-nt} \langle g, \mathcal{L}_n \rangle_\varepsilon \mathcal{P}_n,$$

where the last identity is justified by the Bessel property of the sequence $(\mathcal{P}_n)_{n \geq 0}$ combined with $(\langle g, \mathcal{L}_n \rangle_\varepsilon)_{n \geq 0} \in \ell^2(\mathbb{N})$, for any $f \in L^2(\varepsilon)$. Next, since from Theorem 7.0.1(1) $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ and from Theorem 7.0.1(2), $\overline{\text{Ran}(\mathcal{I}_\phi)} = L^2(\nu)$, we have that its pseudo-inverse \mathcal{I}_ϕ^\dagger , see [BIG02, p.234] for definition, is densely defined from $L^2(\nu)$ into $L^2(\varepsilon)$. Thus, for any $f \in \text{Ran}(\mathcal{I}_\phi)$, $t \geq 0$,

$$(12.1) \quad P_t f = \sum_{n=0}^{\infty} e^{-nt} \langle \mathcal{I}_\phi^\dagger f, \mathcal{L}_n \rangle_\varepsilon \mathcal{P}_n \quad \text{in } L^2(\nu),$$

and, the two linear operators in (12.1) coincide on a dense subset of $L^2(\nu)$.

12.1. Proof of Theorem 1.3.1(2)

We have here $\psi \in \mathcal{N}_\infty^c$ and in particular $\mathfrak{r} < \infty$. First, note that identity (12.1) proves (1.24) of Theorem 1.3.1(2). Next, from (7.13) and the notation there, we get that, for all $p \in \mathbb{N}$ and large n

$$(12.2) \quad \left| \mathcal{P}_n^{(p)}(x) \right| \leq \mathcal{P}_n^{(p)}(-x) \leq n^{p+\frac{1}{2}} \mathcal{E}_\phi(nx) = O\left(n^{p+\frac{1}{2}} e^{-\frac{x}{\mathfrak{r}}n}\right).$$

Then, writing $c_n^\dagger(f) = \langle \mathcal{I}_\phi^\dagger f, \mathcal{L}_n \rangle_\varepsilon$, we get, with $x \in (0, K)$, for any $K > 0$, any $f \in \text{Ran}(\mathcal{I}_\phi)$ and $t \geq 0$

$$\begin{aligned} \left| \sum_{n=p}^{\infty} e^{-nt} c_n^\dagger(f) \mathcal{P}_n^{(p)}(x) \right| &\leq C \sum_{n=p}^{\infty} n^{p+\frac{1}{2}} e^{-n(t-\frac{x}{\mathfrak{r}})} c_n^\dagger(f) \\ &\leq C \sum_{n=p}^{\infty} n^{p+\frac{1}{2}} e^{-n(t-\frac{K}{\mathfrak{r}})} c_n^\dagger(f) \end{aligned}$$

where $C > 0$. Since, from the discussion above, $(c_n^\dagger(f))_{n \geq 0} \in \ell^2(\mathbb{N})$ for all $f \in \text{Ran}(\mathcal{I}_\phi)$, we get that for all $k, p \in \mathbb{N}$

$$\frac{d^k}{dt^k} (P_t f)^{(p)}(x) = \sum_{n=p}^{\infty} (-n)^k e^{-nt} c_n^\dagger(f) \mathcal{P}_n^{(p)}(x),$$

where the series is locally uniformly convergent in $\mathfrak{C}_t = \{(t, x) \in \mathbb{R}_+^2; x \leq \mathfrak{r}t\}$, which completes the proof of Theorem 1.3.1(2).

12.2. Proof of Theorem 1.3.1(1)

We prove Theorem 1.3.1(1a) for the domain
(12.3)

$$\mathcal{D}_T = \mathcal{D}_0^{\mathcal{N}_P}(\mathbb{L}^2(\nu)) \cup \mathcal{D}_{T_{\mathfrak{E}\phi}}^{\mathcal{N}_{\mathfrak{E}\phi}}(\mathbb{L}^2(\mathfrak{p}_\alpha)) \cup \mathcal{D}_{T_{\pi\alpha}}^{\mathcal{N}_\alpha}(\mathbb{L}^2(\bar{\nu}_\gamma)) \cup \mathcal{D}_{T_{\pi\alpha, \rho\alpha}}^{\mathcal{N}_\alpha}(\mathbb{L}^2(\nu)) \cup \mathcal{D}_{T_{\bar{\alpha}}}^{\mathcal{N}_R}(\mathbb{L}^2(\nu)).$$

We make the following observations first. For any $t > 0$ consider the linear space $\mathbf{D}(S_t) = \{f \in \mathbb{L}^2(\nu); \langle P_t f, \mathcal{V}_n \rangle_\nu \in \ell^2(\mathbb{N})\}$. Let $S_t : \mathbf{D}(S_t) \mapsto \mathbb{L}^2(\nu)$ be the spectral operator defined by

$$S_t f(x) := \sum_{n=0}^{\infty} \langle P_t f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n(x).$$

Note that $P_t = S_t$ on $\text{Ran}(\mathcal{I}_\phi) \subset \mathbf{D}(S_t)$ and thus $\overline{\mathbf{D}(S_t)} = \mathbb{L}^2(\nu)$. Indeed, $\psi \in \mathcal{N}_\infty$ implies $\mathcal{V}_n \in \mathbb{L}^2(\nu)$, $n \geq 0$, see Theorem 8.0.1(1) with $\mathcal{I}_\phi^* \mathcal{V}_n = \mathcal{L}_n$. Then, for any $f \in \text{Ran}(\mathcal{I}_\phi)$, $\langle \mathcal{I}_\phi^\dagger f, \mathcal{L}_n \rangle_\varepsilon = \langle f, \mathcal{V}_n \rangle_\nu$ and from (12.1), for any $f \in \text{Ran}(\mathcal{I}_\phi)$, $t > 0$,

$$(12.4) \quad P_t f(x) = S_t f(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n(x) \quad \text{in } \mathbb{L}^2(\nu).$$

This incidentally proves (1.20) of Theorem 1.3.1(1a) for $\mathcal{D}_0^{\mathcal{N}_\infty}(\text{Ran}(\mathcal{I}_\phi))$. Since $\overline{\text{Ran}(\mathcal{I}_\phi)} = \mathbb{L}^2(\nu)$, by the bounded linear transformation theorem, we get that P_t is the unique linear continuous extension of S_t in $\mathbb{L}^2(\nu)$. In order to characterize the domain of the spectral operator and show that each member of \mathcal{D}_T in (12.3) is a subset of this domain, we present several methodologies to show that for some L a linear subspace of $\mathbb{L}^2(\nu)$ we have, for all $f \in L$, $(\langle P_t f, \mathcal{V}_n \rangle_\nu)_{n \geq 0} \in \ell^2(\mathbb{N})$ and as a consequence $P_t f = S_t f$ in $\mathbb{L}^2(\nu)$. Although, we know that $P_t = S_t$ on $\text{Ran}(\mathcal{I}_\phi)$ and also $\overline{\text{Ran}(\mathcal{I}_\phi)} = \mathbb{L}^2(\nu)$, it may not be trivial to show, even when $S_t f \in \mathbb{L}^2(\nu)$ for $f \in L$, that S_t is also continuous in $\mathbb{L}^2(\nu)$ on this subspace.

12.2.1. Hilbert sequence argument. Recall that $\psi \in \mathcal{N}_\infty$. The following statement holds.

PROPOSITION 12.2.1. *Write $V^\perp = \text{Span}(\mathcal{V}_n)^\perp$. Then, for any $t \geq 0$, $P_t V^\perp \subseteq V^\perp$ and $\text{Ker}(P_t) \subseteq V^\perp = \text{Ker}(S_t)$. Moreover, if $\overline{\text{Span}(\mathcal{V}_n)} = \mathbb{L}^2(\nu)$ then $\text{Ker}(P_t) = \{0\}$ and for any $f \in \mathbb{L}^2(\nu)$ and $t > 0$ such that $(\langle P_t f, \mathcal{V}_n \rangle_\nu)_{n \geq 0} \in \ell^2(\mathbb{N})$ we have that $P_t f = S_t f$ in $\mathbb{L}^2(\nu)$.*

Proof. If $f \in V^\perp$ then from (8.1), we have for any $n \in \mathbb{N}$, $\langle P_t f, \mathcal{V}_n \rangle_\nu = e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu = 0$ that is $P_t f \in V^\perp$. A similar argument yields $\text{Ker}(P_t) \subseteq V^\perp$. With still $f \in V^\perp$, we have that

$$S_t f = \sum_{n=0}^{\infty} \langle P_t f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n = 0 \in \mathbb{L}^2(\nu),$$

that is $V^\perp \subseteq \text{Ker}(S_t)$. Next, let $f \in \text{Ker}(S_t)$. Then, as $(\mathcal{V}_n, \mathcal{P}_n)_{n \geq 0}$ is a biorthogonal sequence, see Theorem 1.5.1(2), we have, for any $m \geq 0$, that

$$0 = \langle S_t f, \mathcal{V}_m \rangle_\nu = e^{-mt} \langle f, \mathcal{V}_m \rangle_\nu,$$

that is $\text{Ker}(S_t) \subseteq V^\perp$. Assume now that $\overline{\text{Span}}(\mathcal{V}_n) = L^2(\nu)$ and f, t are such that $(\langle P_t f, \mathcal{V}_n \rangle_\nu)_{n \geq 0} \in \ell^2(\mathbb{N})$. Since $(\mathcal{P}_n)_{n \geq 0}$ is a Bessel sequence, see Theorem 7.0.3(2), then

$$S_t f = \sum_{n=0}^{\infty} \langle P_t f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n \in L^2(\nu).$$

Thus, as $(\mathcal{V}_n, \mathcal{P}_n)_{n \geq 0}$ is a biorthogonal sequence, we have, for any $m \geq 0$, that

$$\langle S_t f, \mathcal{V}_m \rangle_\nu = \langle P_t f, \mathcal{V}_m \rangle_\nu,$$

which provides the statement as in a Hilbert space the notions of complete and total sequences coincide, i.e. $V^\perp = \{0\}$ if and only if $\overline{\text{Span}}(\mathcal{V}_n) = L^2(\nu)$. \square

12.2.2. A density argument via the intertwining operator. As we can prove the completeness of the sequence of co-eigenfunctions for some classes and thus invoke Proposition 12.2.1, in this part we develop another technique, which relies on a density argument, to show that the two operators coincide on a subspace. More specifically, consider the Hilbert space L with norm $\|\cdot\|_L$ generated by a non-negative measure on \mathbb{R}^+ with density L and $L \subseteq L^2(\nu)$ such that $\mathcal{I}_\phi \in \mathbf{B}(L)$ with $\overline{\text{Ran}}(\mathcal{I}_\phi)_{\|\cdot\|_L} = L$. Note that in this case $L \subseteq L^2(\nu)$ implies that ν is absolutely continuous with respect to L and the $C\|\cdot\|_L \geq \|\cdot\|_\nu$, for some $C \in (0, \infty)$. Assume further that for any $f \in L$, we have, for some functions $V_n \in L$, that

$$(12.5) \quad \langle f, \mathcal{V}_n \rangle_\nu \leq \|f\|_L \|V_n\|_L \text{ with } \|V_n\|_L = O(n^a e^{nT}),$$

for some $T, a > 0$ and n large. This implies that, for all $t > T$, $(e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu)_{n \geq 0} \in \ell^2(\mathbb{N})$. Since $(\mathcal{P}_n)_{n \geq 0}$ forms a Bessel sequence with Bessel bound 1, see Theorem 7.0.3(2) and (7.26), then for all $f \in L$ and $t > T$,

$$S_t f = \sum_{n=0}^{\infty} \langle P_t f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n = \mathcal{S}((e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu)_{n \geq 0}) \in L^2(\nu),$$

where

$$\mathcal{S} : (c_n)_{n \geq 0} \in \ell^2(\mathbb{N}) \rightarrow \sum_{n=0}^{\infty} c_n \mathcal{P}_n \in L^2(\nu),$$

see (1.31), is norm bounded by 1. We aim to show that $S_t f = P_t f$. From the assumption $\overline{\text{Ran}}(\mathcal{I}_\phi)_{\|\cdot\|_L} = L$ there is $(f_m)_{m \geq 0} \in L$ such that $\lim_{m \rightarrow \infty} \mathcal{I}_\phi f_m = f$ in L . Putting for any $n, m \in \mathbb{N}$, $g \in L^2(\nu)$,

$$c_{n,t}(g) = e^{-nt} \langle g, \mathcal{V}_n \rangle_\nu,$$

we have, from (12.4), for all $t > 0$, and in particular for all $t > T$ and $m \geq 0$, that

$$P_t \mathcal{I}_\phi f_m(x) = \sum_{n=0}^{\infty} c_{n,t}(\mathcal{I}_\phi f_m) \mathcal{P}_n(x) \quad \text{in } L^2(\nu).$$

Therefore, we get that, for $t > T$,

$$\begin{aligned} \|P_t \mathcal{I}_\phi f_m - S_t f\|_\nu^2 &= \|\mathcal{S}(c_{n,t}(\mathcal{I}_\phi f_m - f))\|_\nu^2 \\ &\leq \sum_{n=0}^{\infty} c_{n,t}^2 (\mathcal{I}_\phi f_m - f) \leq \|\mathcal{I}_\phi f_m - f\|_{\mathbb{L}}^2 \sum_{n=0}^{\infty} e^{-2nt} \|V_n\|_{\mathbb{L}}^2 \\ &\leq C_t \|\mathcal{I}_\phi f_m - f\|_{\mathbb{L}}^2, \end{aligned}$$

where $\infty > C_t > 0$ and we have used (12.5) for the two last inequalities. Thus $\lim_{m \rightarrow \infty} P_t \mathcal{I}_\phi f_m = S_t f$ in $L^2(\nu)$. However, as P_t is a contraction in $L^2(\nu)$ and $\lim_{m \rightarrow \infty} \mathcal{I}_\phi f_m = f$ in $L^2(\nu)$, we conclude that $P_t f = S_t f$

We now apply the previous results to all domains in (12.3).

First, if $\psi \in \mathcal{N}_P$ (resp. $\psi \in \mathcal{N}_R$), since from (10.1) (resp. (10.3)), we have for large n and $\epsilon > 0$, $\|\mathcal{V}_n\|_\nu = O(e^{\epsilon n})$ (resp. $\|\mathcal{V}_n\|_\nu = O(e^{T_{\bar{\alpha}} n})$), we deduce, from the density argument presented in 12.2.2 with $L = L^2(\nu)$, the convergence in $L^2(\nu)$ of the spectral operator for the domain

$$\mathbf{D}_0^{\mathcal{N}_P}(L^2(\nu)) \cup \mathbf{D}_{T_{\bar{\alpha}}}^{\mathcal{N}_R}(L^2(\nu))$$

and thus (1.20) of Theorem 1.3.1(1a) for those. If, furthermore, $\psi \in \mathcal{N}_{\alpha, m}^{d_\phi}$, see (1.34) for the definition, the Hilbert sequence approach developed in the subsection 12.2.1 can be used as in this case $\overline{\text{Span}}(\mathcal{V}_n) = L^2(\nu)$, see Theorem 1.5.1(3). We continue with $\psi \in \mathcal{N}_\alpha$ and hence the domains

$$\mathcal{D}_{T_{\pi_\alpha}}^{\mathcal{N}_\alpha}(L^2(\bar{\nu}_\gamma)) \cup \mathcal{D}_{T_{\pi_\alpha, \rho_\alpha}}^{\mathcal{N}_\alpha}(L^2(\nu))$$

from (12.3). Note that for any $f \in L^2(\bar{\nu}_\gamma)$, $\bar{\nu}_\gamma(x) = \max\left(\nu(x), e^{-x^{\frac{1}{\gamma}}}\right)$, $x > 0$, $\gamma > \alpha + 1$, the Cauchy-Schwarz inequality yields that with $\mathcal{V}_n \nu = w_n$, $n \geq 0$,

$$\begin{aligned} |\langle f, \mathcal{V}_n \rangle_\nu| &= \left| \left\langle f, \frac{w_n}{\bar{\nu}_\gamma} \right\rangle_{\bar{\nu}_\gamma} \right| \\ &\leq \|f\|_{\bar{\nu}_\gamma} \left\| \frac{w_n}{\bar{\nu}_\gamma} \right\|_{\bar{\nu}_\gamma} = O\left(e^{(T_{\pi_\alpha} + \epsilon)n}\right), \end{aligned}$$

where, for large n , we have used the estimate (9.4) with $a > \frac{3}{2}$ and any $\epsilon > 0$. Moreover as plainly $\frac{\bar{\nu}_\gamma(x)}{\Gamma(\gamma+1)} dx$ is a probability measure on \mathbb{R}_+ with $\int_0^\infty e^{\epsilon x} \bar{\nu}_\gamma(x) dx < \infty$ for some $\epsilon > 0$, we have that it is moment determinate and thus the set of polynomials is dense in $L^2(\bar{\nu}_\gamma)$, see [Akh65]. Hence

$$\overline{\text{Ran}}(\mathcal{I}_\phi)_{\|\cdot\|_{\bar{\nu}_\gamma}} = L^2(\bar{\nu}_\gamma).$$

By resorting again to the density argument outlined in the subsection 12.2.2 with $L = L^2(\bar{\nu}_\gamma) \subseteq L^2(\nu)$ or $L = L^2(\nu)$ and using in this latter case the estimate (9.5), we get that $P_t = S_t$ thus (1.20) of Theorem 1.3.1(1a) for the domains $\mathbf{D}_{T_{\pi_\alpha}}^{\mathcal{N}_\alpha}(L^2(\bar{\nu}_\gamma))$ and $\mathbf{D}_{T_{\pi_\alpha, \rho_\alpha}}^{\mathcal{N}_\alpha}(L^2(\nu))$ with

$$T_{\pi_\alpha} = -\ln \sin\left(\frac{\pi}{2}\alpha\right) \leq T_{\pi_\alpha, \rho_\alpha} = \max(T_{\pi_\alpha}, 1 + \rho_\alpha^{-1}),$$

see Theorem 9.0.1(2b). Note that often $T_{\pi_\alpha, \rho_\alpha} > T_{\pi_\alpha}$ which makes the two different domains valuable.

Finally, by a similar token, for any $\psi \in \mathcal{N}_{\underline{\Theta}_\phi}$ and $f \in L^2(\mathbf{p}_\alpha)$, $\mathbf{p}_\alpha(x) = x^{-\alpha}$, $x > 0$, $\alpha \in (0, 1)$, using the estimate (9.2), we get, for any $\varepsilon > 0$, that

$$\begin{aligned} |\langle f, \mathcal{V}_n \rangle_\nu| &= \left| \left\langle f, \frac{w_n}{\mathbf{p}_\alpha} \right\rangle_{\mathbf{p}_\alpha} \right| \\ &\leq \|f\|_{\mathbf{p}_\alpha} \left\| \frac{w_n}{\mathbf{p}_\alpha} \right\|_{\mathbf{p}_\alpha} = O\left(e^{(T_{\underline{\Theta}_\phi} + \varepsilon)n}\right). \end{aligned}$$

This combined with the fact that $\overline{\text{Ran}(\mathcal{I}_\phi)}_{\|\cdot\|_{\mathbf{p}_\alpha}} = L^2(\mathbf{p}_\alpha)$, see Lemma 8.5.3(2), and $L^2(\mathbf{p}_\alpha) \subseteq L^2(\nu)$, thanks to Lemma 5.3.5 being applicable with $n = k = 0$, $\bar{a} = \alpha$, yield the validity of (1.20) of Theorem 1.3.1(1a) for the domain $\mathbf{D}_{T_{\underline{\Theta}_\phi}}^{\mathcal{N}_{\underline{\Theta}_\phi}}(L^2(\mathbf{p}_\alpha))$ and hence we have established the whole of Theorem 1.3.1(1a). Theorem 1.3.1(1b) is a direct consequence of (12.4) and (12.2). Theorem 1.3.1(1c) also follows the same way noting that when $\mathfrak{r} = \infty$, (12.2) gives that

$$\left| \mathcal{P}_n^{(p)}(x) \right| = O(e^{\varepsilon n}),$$

for all $\varepsilon > 0$. Theorem 1.3.1(1d) is then trivial.

12.3. Heat kernel expansion

This part is devoted to the proof of Theorem 1.3.1(3). To this end, we assume that $\psi \in \mathcal{N}_{\underline{\Theta}}$ and recall, from Theorem 1.2.1(1) that, in this case, $\mathfrak{r} = \infty$, and, from Theorem 1.2.1(3c), that $\mathcal{N}_R \subset \mathcal{N}_{\underline{\Theta}}$. Moreover, we set

$$T = T_{\underline{\Theta}_\phi} = -\ln \sin \underline{\Theta}_\phi \text{ if } \psi \in \mathcal{N}_{\underline{\Theta}_\phi} \text{ or } T = T_{\bar{\alpha}} \wedge T_{\underline{\Theta}_\phi} \text{ if } \psi \in \mathcal{N}_R,$$

where we recall that $T_{\bar{\alpha}} = -\ln(2^{\bar{\alpha}} - 1)$ with where for any $\psi \in \mathcal{N}_R$ we have defined

$$\bar{\alpha} = \sup \left\{ 0 < \alpha \leq 1; \exists \mathfrak{m} > \frac{1}{\alpha} - 1 \text{ and } \psi \in \mathcal{N}_{\alpha, \mathfrak{m}} \right\}.$$

With this notation, we introduce the linear operator \bar{S}_t defined formally, for any $t > T$, by

$$\bar{S}_t f(x) = \int_0^\infty f(y) \left(\sum_{n=0}^\infty e^{-nt} \mathcal{P}_n(x) w_n(y) \right) dy = \int_0^\infty f(y) S_t \delta_y(x).$$

Our ultimate aim is to show that $\bar{S}_t f = P_t f$ for functions f with compact support in order to obtain the expression of the heat kernel. To do so we first recall the pointwise bounds for $|w_n(y)|$, which depend solely on the analyticity of ν , that were stated in Proposition 10.2.1(10.17). It says that if $\nu \in \mathcal{A}(\Theta)$ then for any $t > T_\Theta$, $y > 0$, and $n \geq 1$,

$$(12.6) \quad |w_n(y)| \leq F(y, t) e^{tn},$$

where the function $(y, t) \mapsto F(y, t)$ is locally uniformly bounded on $\mathbb{R}_+ \times (T_{\underline{\Theta}_\phi}, \infty)$. The next result shows that for $f \in \mathbf{C}_c(\mathbb{R}_+)$, we have that $S_t f = \bar{S}_t f$ for all t big enough.

PROPOSITION 12.3.1. *Let $\nu \in \mathcal{A}(\Theta)$ then for $|z| < e^{-T_\Theta}$, with $T_\Theta = -\ln \sin \Theta$, we have that, for any $f \in \mathcal{C}_c(\mathbb{R}_+)$,*

$$S_t f(x) = \bar{S}_t f(x) = \sum_{n=0}^{\infty} e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n(x),$$

where the series converges locally uniformly in (t, x) on $(T_\Theta, \infty) \times \mathbb{R}_+$. Finally, $\bar{S}_t f \in L^2(\nu)$ for $t > T_\Theta$.

Proof. We show that for any $f \in \mathcal{C}_c(\mathbb{R}_+)$, we have, for $t > T_\Theta$, $S_t f(x) = \bar{S}_t f(x)$. Let $K = \text{Supp } f \subsetneq (0, \infty)$. Then, since $\mathfrak{r} = \infty$, from (12.2), which gives that

$$\left| \mathcal{P}_n^{(p)}(x) \right| = O(e^{\epsilon n}),$$

for all $\epsilon > 0$, and (12.6) we deduce that for any $t > t_0 > T_\Theta$, with t_0 fixed, and any $\epsilon \in (0, t - t_0)$

$$\begin{aligned} & \sum_{n=0}^{\infty} e^{-nt} \int_0^\infty |f(y)| |w_n(y)| dy |\mathcal{P}_n(x)| \\ & \leq \sup_{y \in K} C(y, t_0) \sum_{n=0}^{\infty} e^{-n(t-t_0)} \int_K |f(y)| dy |\mathcal{P}_n(x)| \\ & \leq C \sum_{n=0}^{\infty} e^{-n(t-t_0)} |\mathcal{P}_n(x)| \\ (12.7) \quad & \leq C \sum_{n=0}^{\infty} e^{-n(t-t_0-\epsilon)} < \infty. \end{aligned}$$

Then an application of Fubini's theorem yields, that for any $f \in \mathcal{C}_c(\mathbb{R}_+)$ and $t > T_\Theta$,

$$\begin{aligned} \bar{S}_t f(x) &= \int_0^\infty f(y) \sum_{n=0}^{\infty} e^{-nt} \mathcal{P}_n(x) w_n(y) dy \\ &= \sum_{n=0}^{\infty} e^{-nt} \mathcal{P}_n(x) \int_0^\infty f(y) w_n(y) dy \\ &= S_t f(x). \end{aligned}$$

Finally, $\bar{S}_t f \in L^2(\nu)$, for $t > T_\Theta$, follows from $(\mathcal{P}_n)_{n \geq 0}$ being a Bessel sequence and $(e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu) \in \ell^2(\mathbb{N})$ which in turn is a consequence of the first computation in (12.7). \square

From the proof of Theorem 1.3.1(1) we have that $P_t f = S_t f$ for the domain $\mathcal{D}_{T_{\Theta_\phi}}^{\mathcal{N}_{\Theta_\phi}}(L^2(\mathbf{p}_\alpha))$ that is for all $f \in L^2(\mathbf{p}_\alpha)$, $t > T_{\Theta_\phi}$, when $\psi \in \mathcal{N}_{\Theta_\phi}$. However, as in the proof of Lemma 9.3.1, $\nu \in \mathcal{A}(\Theta_\phi)$ and therefore from Proposition 12.3.1 we have that $\bar{S}_t f = S_t f$, for $t > T_{\Theta_\phi}$ and $f \in \mathcal{C}_c(\mathbb{R}_+)$. As $\mathcal{C}_c(\mathbb{R}_+) \subseteq L^2(\mathbf{p}_\alpha)$ we conclude that the semi-group is absolutely continuous with density

$$P_t(x, y) = \sum_{n=0}^{\infty} e^{-nt} \mathcal{P}_n(x) w_n(y),$$

for $t > T_{\Theta_\phi}$, $(x, y) \in \mathbb{R}^+ \times \mathbb{R}^+$. Then the expression (1.25), that is

$$\frac{d^k}{dt^k} P_t^{(p,q)}(x, y) = \sum_{n=p}^{\infty} (-n)^k e^{-nt} \mathcal{P}_n^{(p)}(x) w_n^{(q)}(y)$$

follows from (12.2) with $\mathfrak{r} = \infty$, which gives that $|\mathcal{P}_n^{(p)}(x)| = O(e^{\epsilon n})$, for all $\epsilon > 0$, and an easy extension of (10.17) which yields $|w_n^{(q)}(y)| \leq F_q(y, t)e^{tn}$, $t > T_{\bar{\alpha}}$, where F_q has the same properties of F . For the final part of the proof of Theorem 1.3.1(3) assume that $\nu \in \mathcal{A}(\frac{\pi}{2})$. Therefore, Proposition 12.3.1 gives that $S_t f = \bar{S}_t f \in L^2(\nu)$, for all $t > 0$ and $f \in \mathcal{C}_c(\mathbb{R}_+)$. We prove the following result which relates \bar{S} to the semigroup P .

PROPOSITION 12.3.2. *Assume that $S_t f = \bar{S}_t f \in L^2(\nu)$, for all $t > 0$, $f \in \mathcal{C}_c(\mathbb{R}_+)$ and that, for $t > t_0(f)$, $P_t f = S_t f$. Then*

$$P_t f = S_t f = \bar{S}_t f,$$

for all $t > 0$.

Proof. Fix $a > 0$, write, for any $M \in \mathbb{N}$,

$$S_t^M f = \sum_{n=0}^M e^{-nt} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n$$

and $\tilde{S}_t^M f = S_t f - S_t^M f$. Let now $f \in \mathcal{C}_c(\mathbb{R}_+)$ and note that with any $t > 0$, $M \in \mathbb{N}$,

$$\begin{aligned} \|P_t S_a f - S_{a+t} f\|_\nu &\leq \|P_t S_a^M f - S_{a+t}^M f\|_\nu + \|\tilde{S}_{a+t}^M f\|_\nu + \|P_t \tilde{S}_a^M f\|_\nu \\ &= \|\tilde{S}_{a+t}^M f\|_\nu + \|P_t \tilde{S}_a^M f\|_\nu \end{aligned}$$

since

$$\begin{aligned} P_t S_a^M f &= P_t \sum_{n=0}^M e^{-na} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n \\ &= \sum_{n=0}^M e^{-n(a+t)} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n = S_{a+t}^M f. \end{aligned}$$

Now, as $P_t \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$, we get that

$$\|P_t S_a f - S_{a+t} f\|_\nu \leq \|\tilde{S}_{a+t}^M f\|_\nu + \|\tilde{S}_a^M f\|_\nu.$$

Letting $M \rightarrow \infty$ and using the fact that $S_a f, S_{a+t} f \in L^2(\nu)$ we conclude that

$$P_t S_a f = S_{a+t} f,$$

for all $t > 0$. Therefore, from the assumptions, for $t + a > t_0(f)$ we have that $S_{a+t} f = P_{a+t} f$ and thus

$$P_t P_a f = P_t S_a f$$

and hence $P_a f - S_a f \in \text{Ker}(P_t)$. However, if we conclude that $P_a f - S_a f \in \mathcal{D}(\mathbf{G}_\nu)$, the domain of the infinitesimal generator (in $L^2(\nu)$) \mathbf{G}_ν of P , then a classical result of strongly continuous semigroups yields $P_a f = S_a f$. Choose furthermore $f \in \mathcal{C}_c^\infty(\mathbb{R}_+)$. Then, since the gL processes are Feller processes, see Definition 1.0.1, we know that $f \in \mathcal{D}(\mathbf{G}) \subset \mathcal{C}_0(\mathbb{R}_+)$, where \mathbf{G} is the generator of the Feller semigroup of the gL process, as a semigroup in $\mathcal{C}_0(\mathbb{R}_+)$. However, from (1.6) and $f \in \mathcal{C}_c^\infty(\mathbb{R}_+)$

it can be easily seen that $\|\mathbf{G}f\|_\infty < \infty$ and hence $f \in \mathcal{D}(\mathbf{G}_\nu) \subset L^2(\nu)$, where $\mathcal{D}(\mathbf{G}_\nu)$ is the domain of the extension of \mathbf{G} to \mathbf{G}_ν . Therefore,

$$P_a f \in \mathcal{D}(\mathbf{G}_\nu) \subset L^2(\nu).$$

It remains to show that $S_a f \in \mathcal{D}(\mathbf{G}_\nu)$. Clearly,

$$\begin{aligned} \lim_{t \rightarrow 0} \frac{P_t S_a f - S_a f}{t} &= \lim_{t \rightarrow 0} \frac{S_{a+tf} - S_a f}{t} \\ &= \lim_{t \rightarrow 0} \sum_{n=0}^{\infty} e^{-na} \frac{(1 - e^{-nt})}{t} \langle f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n \end{aligned}$$

and we can take the limit under the sum to verify that $S_a f \in \mathcal{D}(\mathbf{G}_\nu) \subset L^2(\nu)$ thanks to $(\mathcal{P}_n)_{n \geq 0}$ is a Bessel sequence and

$$\left(e^{-na} \frac{(1 - e^{-nt})}{t} \langle f, \mathcal{V}_n \rangle_\nu \right)_{n \geq 0} \in \ell^2(\mathbb{N})$$

since $S_a f \in L^2(\nu)$. Thus, $S_a f = P_a f$ for $f \in \mathbf{C}_c^\infty(\mathbb{R}_+)$. For, $f \in \mathbf{C}_c(\mathbb{R}_+)$ we can approximate. \square

We can now conclude the claim of Theorem 1.3.1(3). Indeed, from Proposition 12.3.1 $\bar{S}_t f = S_t f$, $t > 0$, for any $f \in \mathbf{C}_c(\mathbb{R}_+)$, since $\nu \in \mathcal{A}(\frac{\pi}{2})$. However, since when $\psi \in \mathcal{N}_{\underline{\Theta}}$ and $(\psi, t, f) \in \mathcal{D}_T$, for $t > T_{\underline{\Theta}_\psi}$, we have that $P_t f = S_t f$. Thus, Proposition 12.3.2 implies that

$$P_t f = S_t f = \bar{S}_t f,$$

for all $t > 0$, and the kernel expansion follows immediately. Besides, when $\psi \in \mathcal{N}_P$ then $\nu \in \mathcal{A}(\frac{\pi}{2})$, see Lemma 9.3.1, Lemma 12.3.3 gives the second such condition in item (3).

12.3.1. Laguerre type expansions of the invariant density.

LEMMA 12.3.3. *Let $\psi \in \mathcal{N} \setminus \mathcal{N}_P$. Then, writing*

$$\bar{R}_\phi = - \lim_{n \rightarrow \infty} \frac{\ln \left| \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{W_\phi(k+1)}{\Gamma(k+1)} \left(\frac{W_\phi}{\Gamma} \right) \right|}{2\sqrt{n}},$$

we have that ν is holomorphic in the parabolic domain

$$\mathbb{C}_{\mathcal{P}_{\bar{R}_\phi}} = \{z = a + ib \in \mathbb{C}; b^2 < 4\bar{R}_\phi^2(a + \bar{R}_\phi^2)\}.$$

In particular, if $\bar{R}_\phi = \infty$, then $\nu \in \mathcal{A}(\frac{\pi}{2})$.

Proof. We first prove that when $\sigma^2 = 0$ then $\frac{\nu^2}{\varepsilon} \in L^2(\mathbb{R}_+)$ and hence $\frac{\nu}{\varepsilon} \in L^2(\varepsilon)$. The fact that $\frac{\nu^2}{\varepsilon}$ is integrable on $(0, 1)$ is clear from Theorem 5.3.5 applied with $n = k = 0$ and $a \in (0, 1)$. For the interval $(1, \infty)$ we use Theorem 5.0.5(5.7), that is

$$\nu(x) \approx \frac{C_\psi}{\sqrt{2\pi}} \sqrt{\varphi'(x)} e^{-\int_m^x \varphi(y) \frac{dy}{y}}.$$

Since $\sigma^2 = 0$ then Proposition 4.1.1(3) gives that $\phi(x) \cong o(x)$, $\phi'(x) = o(1)$ and thus $\lim_{x \rightarrow \infty} x^{-1} \varphi(x) = \infty$ and

$$e^{-\int_m^x \varphi(y) \frac{dy}{y}} \lesssim e^{-x-m} \lesssim e^{-x}.$$

Then, for $x > 1$ with some constant $C > 0$

$$\begin{aligned} \nu^2(x) &\leq C\varphi'(x)e^{-2\int_m^x \varphi(y)\frac{dy}{y}} \\ &\leq C\varphi'(x)e^{-\int_m^x \varphi(y)\frac{dy}{y}}\varepsilon(x). \end{aligned}$$

Hence the integrability on $(1, \infty)$ of $\frac{\nu^2}{\varepsilon}$ follows easily. Thus, in $L^2(\varepsilon)$,

$$\frac{\nu(x)}{\varepsilon(x)} = \sum_{n=0}^{\infty} \left\langle \frac{\nu}{\varepsilon}, \mathcal{L}_n \right\rangle_{\varepsilon} \mathcal{L}_n(x).$$

Next, observe that, for any $n \geq 0$,

$$\begin{aligned} \left\langle \frac{\nu}{\varepsilon}, \mathcal{L}_n \right\rangle_{\varepsilon} &= \int_0^{\infty} \nu(x)\mathcal{L}_n(x)dx = \mathcal{V}_{\psi}\mathcal{L}_n(1) \\ &= \sum_{k=0}^n (-1)^k \binom{n}{k} \frac{W_{\phi}(k+1)}{\Gamma(k+1)}, \end{aligned}$$

where we have used the linearity of the Markov operator \mathcal{V}_{ψ} and the expression (2.20) of its moments. The statement regarding the analyticity of ν follows from [Pol47] and the last claim follows readily. \square

12.4. Expansion of the adjoint semigroup: Proof of Theorem 1.3.7

Let $\psi \in \mathcal{N}_{\Theta}$, from (1.25) and the duality property, we have, for any $t > T_{\Theta} = -\ln \sin(\Theta)$, that

$$P_t^*(y, x) = \sum_{n=0}^{\infty} e^{-nt} \mathcal{V}_n(y) \mathcal{P}_n(x) \nu(x)$$

where the convergence is locally uniform on $\mathbb{R}_+ \times \mathbb{R}_+$. Next, note that for any $g \in L^2(\nu)$, an application of the Cauchy-Schwarz inequality yields, for any $n \in \mathbb{N}$,

$$(12.8) \quad | \langle |g|, |\mathcal{P}_n| \rangle_{\nu} | \leq \|g\|_{\nu} \|\mathcal{P}_n\|_{\nu} \leq \|g\|_{\nu},$$

where $\|\mathcal{P}_n\|_{\nu} \leq 1$ from (7.10). This combined with bound (12.6) allows an application of Fubini's theorem to get, for any $t > T_{\Theta}$, that

$$\begin{aligned} P_t^*g(y) &= \int_0^{\infty} g(x) \sum_{n=0}^{\infty} e^{-nt} \mathcal{V}_n(y) \mathcal{P}_n(x) \nu(x) dx \\ (12.9) \quad &= \sum_{n=0}^{\infty} \langle g, \mathcal{P}_n \rangle_{\nu} e^{-nt} \mathcal{V}_n(y), \end{aligned}$$

where the series converges uniformly on \mathbb{R}_+ and completes the proof of item (1).

For item (2), as $\psi \in \mathcal{N}_P \subset \mathcal{N}_{\Theta}$ with $0 < \underline{\mathbf{m}} = \frac{m+\overline{\Pi}(0^+)}{\sigma^2} < \infty$, from Theorem 10.0.1(1), we have that $(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})}\mathcal{V}_n)_{n \geq 0}$ is a Bessel sequence, where

$$\mathbf{c}_n(\underline{\mathbf{m}}) = \frac{\Gamma(n+1)\Gamma(\underline{\mathbf{m}}+1)}{\Gamma(n+\underline{\mathbf{m}}+1)} \cong O(n^{-\underline{\mathbf{m}}}).$$

Hence, with (12.8), we have that

$$\left(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})}^{-1} e^{-nt} \langle g, \mathcal{P}_n \rangle_{\nu} \right)_{n \geq 0} \in \ell^2(\mathbb{N})$$

for all $g \in L^2(\nu)$. Therefore, from the synthesis operator

$$\mathcal{S} : (c_n)_{n \geq 0} \mapsto \mathcal{S}((c_n)) = \sum_{n=0}^{\infty} c_n \left(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})} \mathcal{V}_n \right)$$

associated to the Bessel sequence $(\sqrt{\mathbf{c}_n(\underline{\mathbf{m}})} \mathcal{V}_n)_{n \geq 0}$, see (1.31), we get that the expansion (12.9) converges in $L^2(\nu)$ and coincides with $P_t^* g$, for all $t > 0$.

12.5. Proof of of Theorem 1.4.1: Rate of convergence to equilibrium.

Let first $\psi \in \mathcal{N}$. Then, note that for any $f \in \text{Ran}(\mathcal{I}_\phi)$, writing $f = \mathcal{I}_\phi \mathbf{f}$, $\mathbf{f} \in L^2(\varepsilon)$, we have from the factorization (7.18), that for any $x > 0$,

$$\nu f = \nu \mathcal{I}_\phi \mathbf{f} = \mathcal{V}_\psi \mathcal{I}_\phi \mathbf{f}(1) = \mathcal{E} \mathbf{f}(1) = \varepsilon \mathbf{f} = \mathcal{I}_\phi \varepsilon \mathbf{f}(x),$$

where the last identity follows from the facts that $\varepsilon \mathbf{f}$ is a constant and \mathcal{I}_ϕ is a Markov operator. Hence, assuming in addition that f is not a constant, we have, for any $t > 0$,

$$\begin{aligned} \|P_t f - \nu f\|_\nu^2 &= \|\mathcal{I}_\phi(Q_t \mathbf{f} - \varepsilon \mathbf{f})\|_\nu^2 \leq \|\mathcal{I}_\phi\|^2 \|Q_t \mathbf{f} - \varepsilon \mathbf{f}\|_\varepsilon^2 \\ &\leq \|\mathcal{I}_\phi\|^2 \frac{\|\mathbf{f} - \varepsilon \mathbf{f}\|_\varepsilon^2}{\|f - \nu f\|_\nu^2} e^{-2t} \|f - \nu f\|_\nu^2, \end{aligned}$$

where we have used successively the intertwining relationship (11.3), the fact that $\mathcal{I}_\phi \in \mathbf{B}(L^2(\varepsilon), L^2(\nu))$ is a contraction, see (7.22), and the exponential decay of the Laguerre semigroup, see [BGL14, Chap. 4]. Moreover since

$$\|f - \nu f\|_\nu^2 \leq \|\mathcal{I}_\phi\|^2 \|\mathbf{f} - \varepsilon \mathbf{f}\|_\varepsilon^2,$$

we get the inequality

$$\|\mathcal{I}_\phi\|^2 \frac{\|\mathbf{f} - \varepsilon \mathbf{f}\|_\varepsilon^2}{\|f - \nu f\|_\nu^2} \geq 1.$$

This ends the proof of item (1). We proceed with item (2). Next, for any $(\psi, f, t) \in \mathcal{D}_T$ such that there exists, for any $n \geq 0$, $V_n \in L$,

$$\|\langle f, \mathcal{V}_n \rangle_\nu\| \leq \|f\|_L \|V_n\|_L \text{ with } \|V_n\|_L = O(n^a e^{nT}),$$

we have, for any $t > T > 0$, writing $\bar{c}_{n,t}(f) = \langle P_t f, \mathcal{P}_n \rangle_\nu$ for $n \geq 1$ and $\bar{c}_{0,t}(f) = 0$, and $k = [2a]$,

$$\begin{aligned} \|P_t f - \nu f\|_\nu^2 &= \left\| \sum_{n=1}^{\infty} \langle P_t f, \mathcal{V}_n \rangle_\nu \mathcal{P}_n(x) \right\|_\nu^2 = \|\mathcal{S}(\bar{c}_{n,t}(f))\|_\nu^2 \\ &\leq \sum_{n=1}^{\infty} |\langle P_t f, \mathcal{V}_n \rangle_\nu|^2 = \sum_{n=1}^{\infty} e^{-2nt} |\langle f, \mathcal{V}_n \rangle_\nu|^2 \\ (12.10) \quad &= \sum_{n=1}^{\infty} e^{-2nt} |\langle f - \nu f, \mathcal{V}_n \rangle_\nu|^2 \\ &\leq \|f - \nu f\|_L^2 \sum_{n=1}^{\infty} e^{-2nt} \|V_n\|_L^2 \leq C_L \|f - \nu f\|_L^2 \sum_{n=1}^{\infty} n^k e^{-2n(t-T)} \\ &\leq 2^{\frac{k}{2}} C_L \sqrt{\left(\frac{1}{e^{2(t-T)} - 1} \right)^{(k)}} \|f - \nu f\|_L^2, \end{aligned}$$

where $C_L > 0$ and we used the synthesis operator associated to the Bessel sequence $(\mathcal{P}_n)_{n \geq 0}$, see (1.31). When $\psi \in \mathcal{N}_R$ note from (10.13) that $\|V_n\|_L \leq C_\nu e^{nT_\alpha}$, $k = 0$ and we conclude the proof of Theorem 1.4.1(2). The first inequality of Theorem 1.4.1(3) follows similarly from (10.14) with $C_L = C_{\nu, \epsilon}$, for any $t > \epsilon > 0$ and any $\epsilon > 0$. Finally, let $\psi \in \mathcal{N}_P$ with $\overline{\Pi}(0^+) < \infty$ and from, Lemma 10.1.3, recall that

$$\mathbf{m} = \frac{m + \overline{\Pi}(0^+)}{\sigma^2} > \mathbf{d}_\epsilon^+ = \mathbf{d}_\epsilon \mathbb{I}_{\{-d_\phi > 0\}},$$

where $\mathbf{d}_\epsilon = -d_\phi - \epsilon > 0$, for some $\epsilon > 0$. Then, from Lemma 10.1.1 and when $-d_\phi > 0$, Lemma 10.1.3, we have that both sequences $\left(\frac{\mathcal{P}_n}{\sqrt{\mathbf{c}_n(\mathbf{d}_\epsilon^+)}}\right)_{n \geq 0}$ and $(\sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n)_{n \geq 0}$, where $\mathbf{c}_n(m) = \frac{\Gamma(n+1)\Gamma(m+1)}{\Gamma(n+m+1)}$, are Bessel sequences in $L^2(\nu)$ with bound 1. Next, observe that since

$$\frac{e^{-2(n-1)t}\mathbf{c}_n(\mathbf{d}_\epsilon^+)\mathbf{c}_1(\mathbf{m})}{\mathbf{c}_1(\mathbf{d}_\epsilon^+)\mathbf{c}_n(\mathbf{m})} = \prod_{j=0}^{n-2} e^{-2t} \left(\frac{n + \mathbf{m} - j}{n + \mathbf{d}_\epsilon - j} \right)$$

we get that

$$(12.11) \quad \sup_{n \geq 1} \frac{e^{-2(n-1)t}\mathbf{c}_n(\mathbf{d}_\epsilon^+)\mathbf{c}_1(\mathbf{m})}{\mathbf{c}_1(\mathbf{d}_\epsilon^+)\mathbf{c}_n(\mathbf{m})} \leq 1 \iff e^{-2t} \left(\frac{\mathbf{m} + 2}{\mathbf{d}_\epsilon + 2} \right) \leq 1,$$

which holds if and only if $t > T_{\mathbf{m}} = \frac{1}{2} \ln \left(\frac{\mathbf{m} + 2}{\mathbf{d}_\epsilon + 2} \right)$. Hence, for any $f \in L^2(\nu)$ and $t > T_{\mathbf{m}} = \frac{1}{2} \ln \left(\frac{\mathbf{m} + 2}{\mathbf{d}_\epsilon + 2} \right)$, we have

$$\begin{aligned} \|P_t f - \nu f\|_\nu^2 &= \|S_t f - \nu f\|_\nu^2 = \|\mathcal{S}(\bar{c}_{n,t}(f))\|_\nu^2 \\ &\leq \sum_{n=1}^{\infty} \frac{\mathbf{c}_n(\mathbf{d}_\epsilon^+)}{\mathbf{c}_n(\mathbf{m})} \left| \left\langle P_t f, \sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n \right\rangle_\nu \right|^2 \\ &= e^{-2t} \sum_{n=1}^{\infty} \frac{e^{-2(n-1)t}\mathbf{c}_n(\mathbf{d}_\epsilon^+)}{\mathbf{c}_n(\mathbf{m})} \left| \left\langle f, \sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n \right\rangle_\nu \right|^2 \\ &= \frac{e^{-2t}\mathbf{c}_1(\mathbf{d}_\epsilon^+)}{\mathbf{c}_1(\mathbf{m})} \sum_{n=1}^{\infty} \frac{e^{-2(n-1)t}\mathbf{c}_n(\mathbf{d}_\epsilon^+)\mathbf{c}_1(\mathbf{m})}{\mathbf{c}_1(\mathbf{d}_\epsilon^+)\mathbf{c}_n(\mathbf{m})} \left| \left\langle f - \nu f, \sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n \right\rangle_\nu \right|^2 \\ &\leq \frac{\mathbf{m} + 1}{\mathbf{d}_\epsilon^+ + 1} e^{-2t} \sum_{n=1}^{\infty} \left| \left\langle f - \nu f, \sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n \right\rangle_\nu \right|^2 \\ &\leq \frac{\mathbf{m} + 1}{\mathbf{d}_\epsilon^+ + 1} e^{-2t} \|f - \nu f\|_\nu^2 \end{aligned}$$

where we used, for the third equality, the biorthogonality of the sequences $(\mathcal{P}_n)_{n \geq 0}$ and $(\mathcal{V}_n)_{n \geq 0}$ and the fact that $\mathcal{P}_0 \equiv 1$, the equivalence (12.11) for the second inequality and, to obtain the last bound the synthesis operator associated to the Bessel sequence $(\sqrt{\mathbf{c}_n(\mathbf{m})}\mathcal{V}_n)_{n \geq 0}$. This proves (1.28) for $t > T_{\mathbf{m}}$. Finally, since for any $t \leq T_{\mathbf{m}}$,

$$\frac{\mathbf{m} + 1}{\mathbf{d}_\epsilon^+ + 1} e^{-2t} \geq \frac{\mathbf{d}_\epsilon^+ + 2}{\mathbf{d}_\epsilon^+ + 1} \frac{\mathbf{m} + 1}{\mathbf{m} + 2} \geq 1$$

as $\mathbf{m} > \mathbf{d}_\epsilon^+$, invoking that P_t is contractive, we complete the proof of this statement and hence of Theorem 1.4.1.

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