Co-operative electrostatics for understanding crown ether hydration patterns

Shridhar R. Gadre* and Subhash S. Pingale

Department of Chemistry, University of Pune, Pune 411 007, India

A prototype problem in supramolecular chemistry is: Would it be possible to understand qualitative trends in hydration process of crown ether from simple considerations at the molecular level? An answer is offered for the hydration of 18-crown-6 (18C6) in terms of co-operative electrostatic effects. These effects are monitored by mapping the molecular electrostatic potential topography of 18-crown-6 as well as various intermediate hydrated species, followed by electrostatic modelling. All model calculations have been done at ab-initio HF/6-31G** level. The trends of these hydration patterns are in good agreement with the corresponding fully optimized ab-initio ones. Final structure of 18C6.4H₂O is quite similar to the corresponding experimental crystal structure. Such an electrostatics-based method seems to be an excellent general tool for understanding weak interactions in supramolecular chemistry.

Supramolecular chemistry, defined by Lehn' as the chemistry of the intermolecular bond, deals with entities formed by complexation of two or more molecular units. This subject covers a wide variety of molecular recognition processes including enzyme-receptor binding', hydration of proteins and DNA, catalyst-substrate as well as cation- π interactions² etc. The importance of electrostatics in these processes has been widely recognized²⁻⁴. For instance, ab-initio⁵ studies as well as computer simulations, despite being computer-intensive methods, are being employed for studying the hydration processes. Although a great deal of energetic and dynamic information is thereby engendered, the understanding of primitive patterns at molecular level is not directly possible by the use of these methods. We propose to present such a general approach, for gaining insights at molecular level, with a test example of crown ether hydration.

The hydration of 18-crown-6 (18C6) has been extensively studied⁶⁻⁸ experimentally as well as theoretically as a prototype problem in supramolecular chemistry. By constructing the melting diagrams of 18C6·nH₂O, the existence of several binary hydrates has been verified⁶. The detailed crystal structures of these hydrates have been probed⁶ by X-ray diffraction. The D_{3d} conformation of hydrated crown ether has also been confirmed by Raman spectroscopy⁷. Earlier investigations have indi-

cated that in aqueous solution, the predominant conformation is the centrosymmetric D_{3d} , which is favourable for complexation with water molecules.

Computational studies⁸ on the hydration of 18C6 include the Monte Carlo simulations of conformationdependent hydration reported by Ranghino et al.8 as well as the recent work by Thompson⁸. It was noted by Ranghino et al.8 that for the D_{3d} conformation, there is a typical co-operativity between two water molecules on each face. They attempted, quite unsuccessfully, to correlate the solute-water interaction energy with the dipole moment of water. Some other correlations with the accessible molecular surface as well as the deepest energy potential around 18C6 were also tried out, although the success there was also rather limited. It was concluded that the results of Monte Carlo simulations are qualitatively explicable but hardly predictable and could not be extrapolated to other systems. Earlier, Hori et al. had concluded from the HF/STO-3G calculations that the charge transfer and electrostatics are significant for cation—crown ether interactions. However, considering that no ionic species or strong polarization effects are involved, it is intuitively felt that for the hydration process, electrostatics alone would dominate.

Recent molecular dynamic simulations by Kowall and Geiger⁸ have also indicated that a water molecule is strongly bound to ether oxygens, two on each face of the ring, whereas the third ether oxygen, one on each face is coordinated by a singly-bound and rapidly-fluctuating water molecule. Thus the question asked by Ranghino et al.⁸ is still unanswered viz. would it be possible to understand (or predict) at least trends in hydration processes from simple considerations?

It is the purpose of this communication to analyse, qualitatively as well as semiquantitatively, the pattern of hydration of 18C6, based on the cooperative electrostatic effects due to the lone pairs of oxygens of 18C6, and those of water molecules. We shall limit the present work to $18C6.4H_2O$, although further extensions to the higher hydrates are readily possible.

The lock and key mechanism for understanding the hydration process of D_{3d} conformers of 18C6 is offered by mapping of the topography of the Molecular Electrostatic Potential (MESP) of 18C6 as well as the various intermediate hydrated species. MESP has been widely used^{9,10} in the chemical literature as a tool for probing electron-rich sites, although MESP topography-mapping is a more recent development. The ab-initio calculations reported in this work are carried out with the help of GAMESS¹¹ and GAUSSIAN¹¹ packages at HF/6-31G** level. The individual MESP features of 18C6 and water have been mapped by employing indigenously developed packages UNIPROP and UNIVIS¹² (Figure 1). It may be expected that the 18C6···H₂O interaction is stronger than the H₂O···H₂O interaction, since 18C6 MESP

^{*}For correspondence.

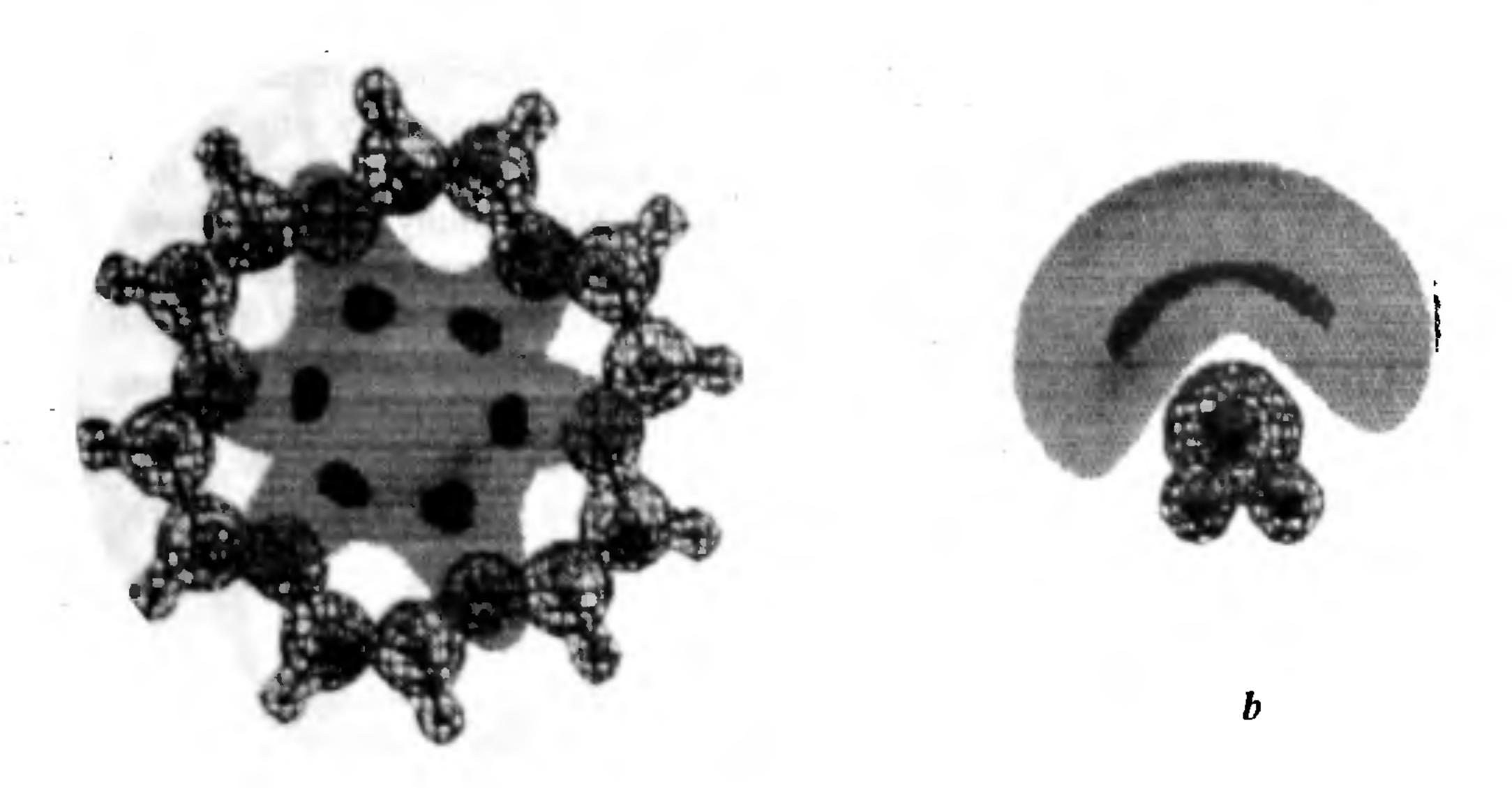


Figure 1. a. 1575.3 kJ mol⁻¹ MESP mesh wire plot and -157.5 kJ mol⁻¹ (light) and -288.8 kJ mol⁻¹ (dark) isosurfaces of 18C6. The MESP minima values are -318.9 kJ mol⁻¹; b. 2362.9 kJ mol⁻¹ MESP mesh wire plot and -131.3 kJ mol⁻¹ (light) and -252.1 kJ mol⁻¹ (dark) isosurfaces for the water molecule. The MESP minima correspond to -261.2 kJ mol⁻¹.

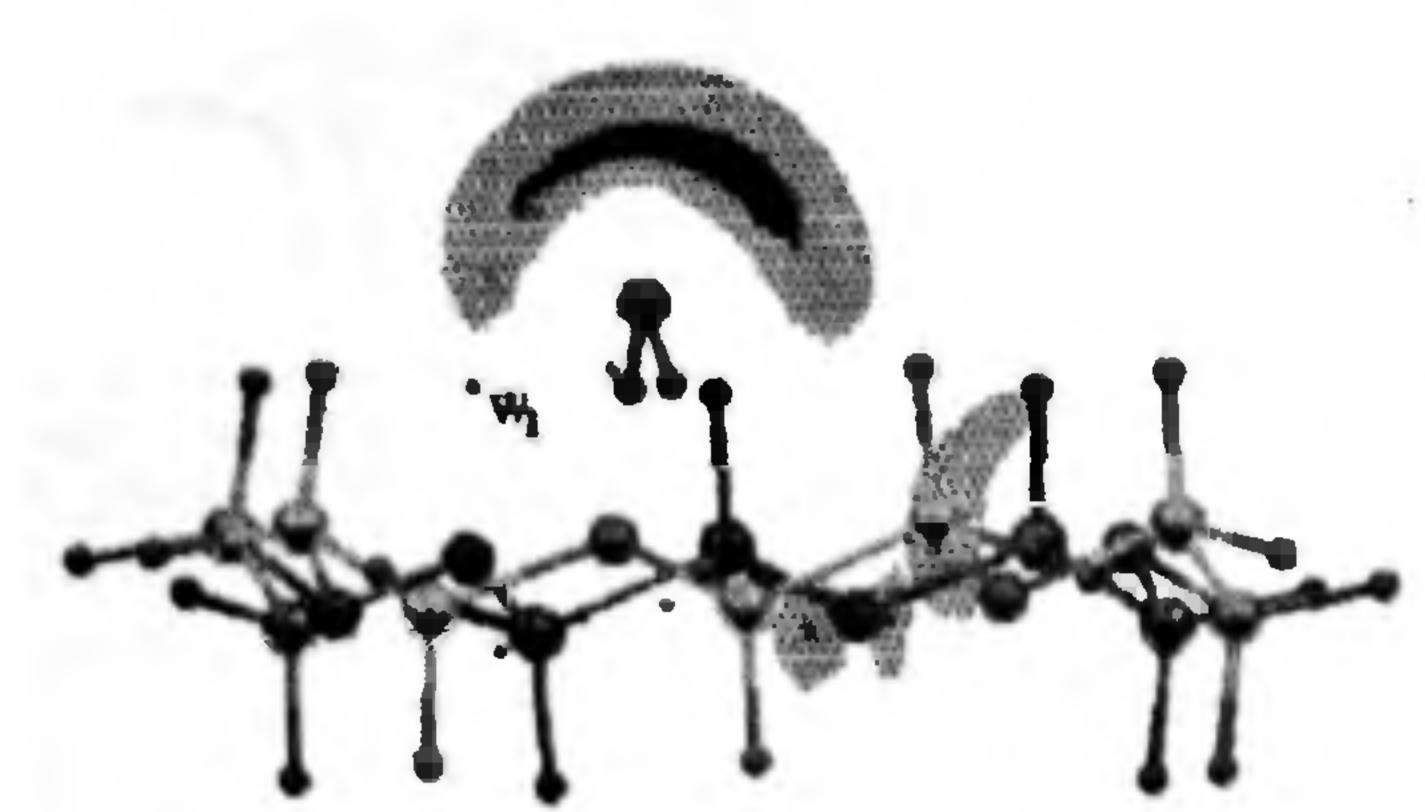


Figure 2. Ab-initio optimized structure of 18C6·H₂O (the water molecule shown as W₁), with crown ether MESP CPs shown as dots, MESP isosurfaces of values -270.4 (light) and -359.7 (dark) kJ mol⁻¹. The two deepest lone pair MESP minima for W₁ are -386.5 and -369.7 kJ mol⁻¹.

minimum is -318.9 kJ mol⁻¹ as against the corresponding value of -261.2 kJ mol⁻¹ for water.

Where the first water molecule binds to 18C6 is qualitatively decided by the MESP topography. An electrostatic model developed in the author's laboratory viz. Electrostatic Potential for Intermolecular Complexation (EPIC)¹³ has been employed for obtaining structures and energetics of 18C6 hydrates. This model utilizes the entire three-dimensional MESP distribution as well as atom-centered charges derived from the respective MESP's¹⁴. The binding energy of the first water molecule (W₁) to 18C6 within EPIC model is 41.0 kJ mol⁻¹, and an ab-initio optimization carried out starting from this geometry yields an interaction energy of 49.4 kJ mol-(all ab-initio values refer to HF/6-31G** level). On the basis set superposition energy (BSSE) correction it turns out to be 36.1 kJ mol⁻¹, 73% of that without the BSSE correction. The MESP topography of 18C6-H₂O has been

calculated using the block diagonal density matrix obtained by patching the corresponding individual abinitio density matrices. Figure 2 presents MESP isosurfaces for 18C6·H₂O complex optimized at ab-initio level. The most noteworthy observation is that the MESP minimum of the bonded water molecule becomes much deeper, viz. -362.3 kJ mol⁻¹ -386.5 kJ mol⁻¹ for EPICpatched and ab-initio optimized structures respectively. This feature has earlier been noticed¹⁵ for the lone pairs in the case of hydrogen-bonded DNA base pairs. The remaining 18C6 oxygen minima show somewhat higher value compared to the 18C6 minima values (typical value -288.8 kJ mol⁻¹). A significant implication of this observation is that the incoming second water molecule (W₂) perhaps binds to the deepest minimum of the first water molecule with its other hydrogen interacting with the 18C6 oxygen minimum on the same face. The EPIC and ab-initio interaction energies of 18C6-H₂O with the second water molecule (second molecule coming from the side of first water molecule of 18C6, see Figure 3) turn out to be 45.7 kJ mol⁻¹ and 52.2 kJ mol⁻¹ respectively. The *ab-initio* interaction energy for the same with BSSE is 40.0 kJ mol⁻¹. The interaction energy, when the second water molecule attacks 18C6·H₂O from the side opposite to that of the first water molecule, turns out to be 39.1 kJ mol⁻¹ and 42.8 kJ mol⁻¹ with EPIC and *ab-initio* level respectively. The *ab-initio* value on BSSE correction becomes 31.6 kJ mol⁻¹. These

interaction energies reveal that the second water molecule preferentially attaches from the side of the first water molecule, according to the predictions based on the MESP study. As seen from Figure 3, the MESP of W₂ becomes substantially deeper (CP values of -333.4 kJ mol⁻¹ and -379.0 kJ mol⁻¹ at EPIC and abinitio level respectively).

An interesting question now is: where will the next incoming water molecule (W₃) bind? Since all the sites on the top face (where the first and second water

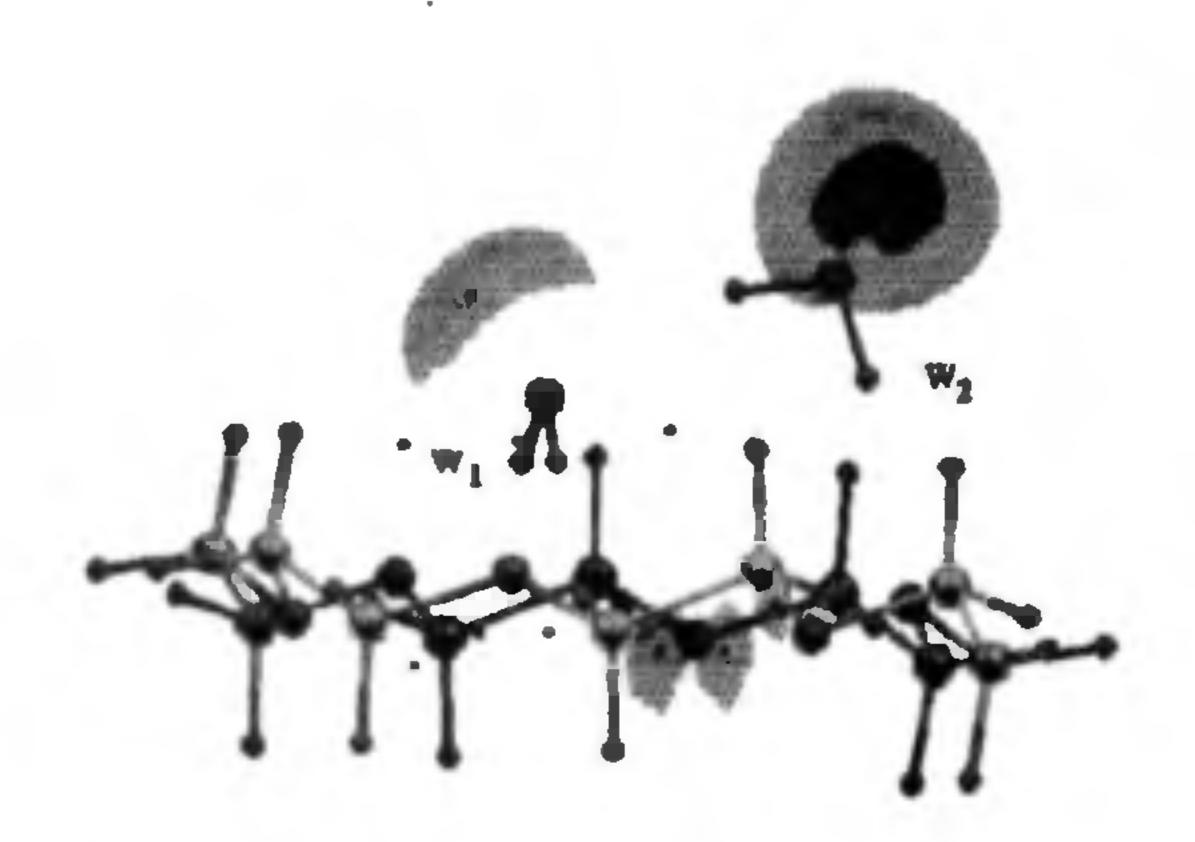


Figure 3. Ab-initio optimized structure of $18C6.2H_2O$ (water molecules as W_1 and W_2), with crown ether MESP CPs shown as dots, MESP isosurfaces of values -244.2 (light) and -341.3 (dark) kJ mol⁻¹. The two deepest lone pair MESP minima of W_2 are -379.0 and -377.8 kJ mol⁻¹.



Figure 4. Ab-initio optimized structure of 18C8-3H₂O (water molecule denoted as W₁, W₂ and W₃), with crown ether MESP CPs shown as dots, MESP isosurfaces of values -199.5 (light) and -325.6 (dark) kJ mol⁻¹. The two lone pair minima for W₂ are -370.6 and -369.3 kJ mol⁻¹ and the two lone pair MESP minima for W₃ are -351.5 and -336.6 kJ mol⁻¹.

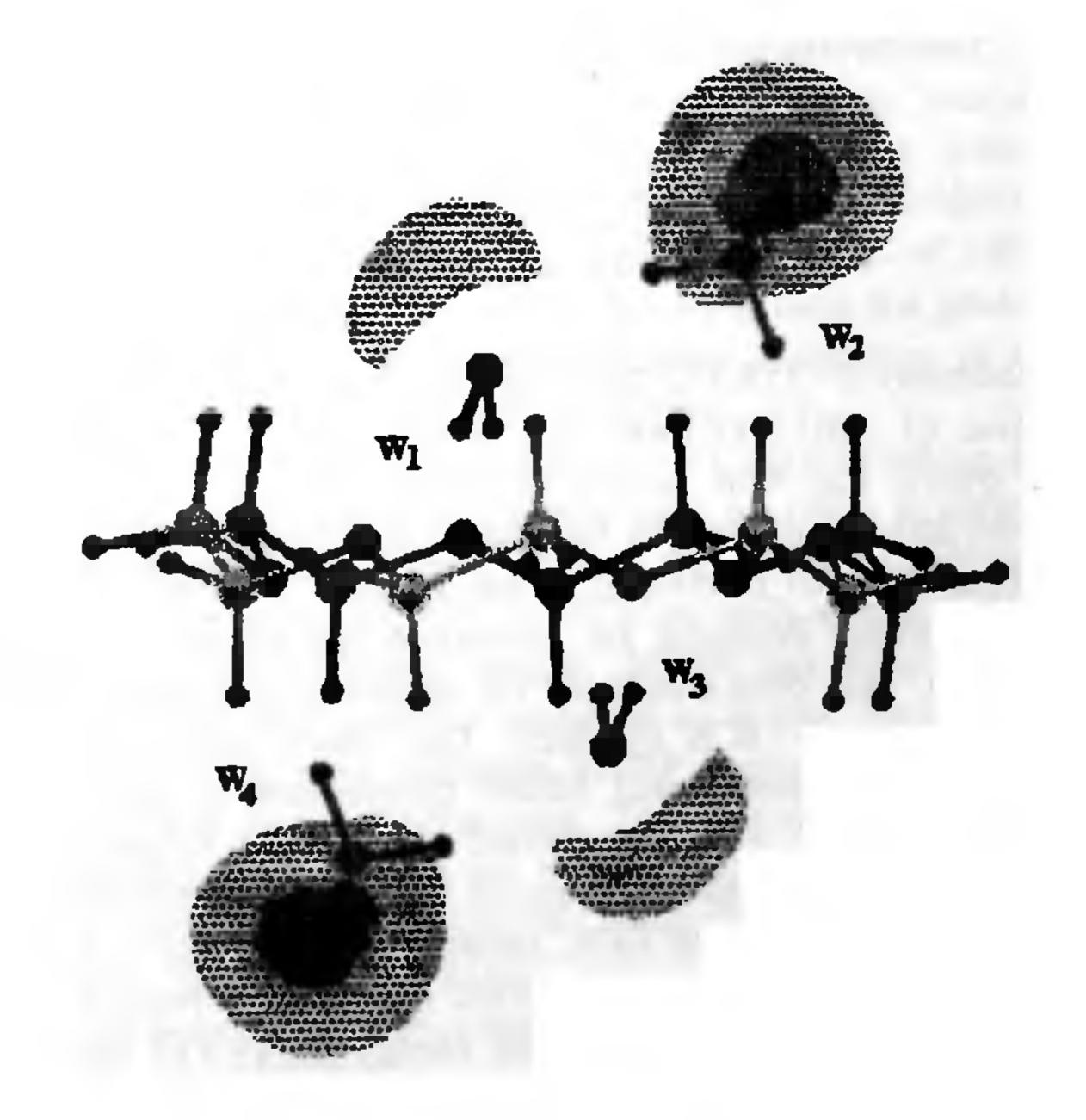


Figure 5. Ab-initio optimized structure of $18C6\cdot 4H_2O$ (water molecules represented as W_1 , W_2 , W_3 and W_4), with MESP isosurfaces of values -183.8 (light) and -315.1 (dark) kJ mol⁻¹. The four deepest lone pair MESP minima for W_2 and W_4 are -356.7 kJ mol⁻¹.

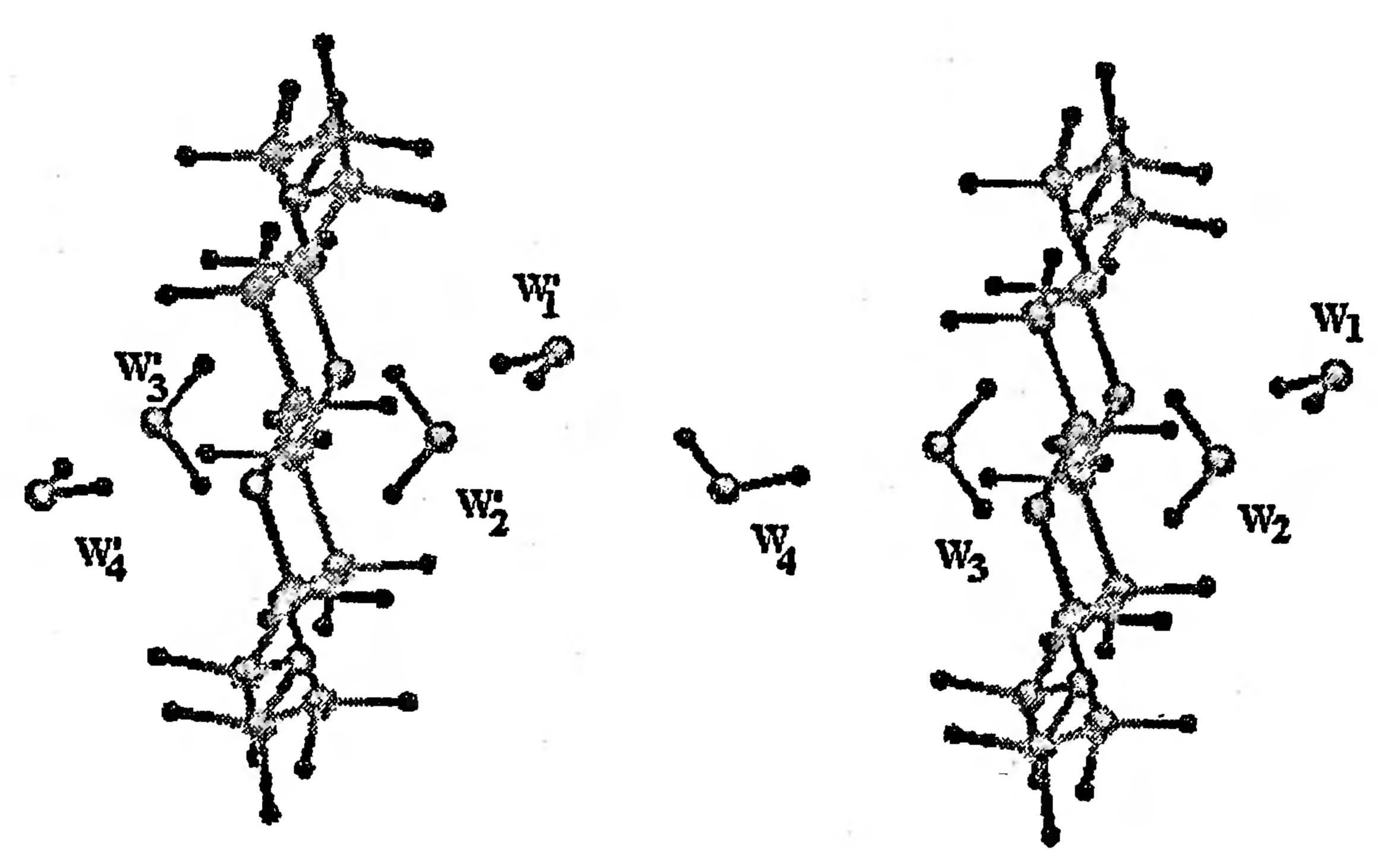


Figure 6. EPIC model optimized structure of dimer of tetrahydrate of 18C6 with flipped water molecule connecting two units viz. 18C6'4H₂O' and 18C6·4H₂O.

molecules bind) of 18C6 are blocked and W, can anchor to only one oxygen minimum of a water molecule, it is therefore expected that W₃ will bind on the other face of 18C6 with a slightly reduced interaction energy compared to that of the first water because the 18C6 oxygen minima of the other face is more positive (typical value -262.6 kJ mol⁻¹). The interaction energies of W₁ with 18C6·2H₂O indeed turn out to be 40.1 kJ mol⁻¹ and 43.5 kJ mol⁻¹ within EPIC and ab-initio frameworks respectively (see Figure 4), the latter becomes 30.2 kJ mol⁻¹ with BSSE correction. It is also found that three water molecules cannot bind on the top face of 18C6 in a cooperative manner as strongly as the above-mentioned arrangement. Continuing this story, the next water molecule (W₄) prefers to bind with one hydrogen, capturing an MESP minimum of W₃ molecule $(-320.3 \text{ kJ mol}^{-1} \text{ and } -351.8 \text{ kJ mol}^{-1} \text{ at EPIC and } ab$ initio levels respectively) and with another hydrogen getting attached to 18C6 minimum (cf. Figure 5). The binding energies of W₄ to 18C6·3H₂O are 43.7 and 49.3 kJ mol⁻¹ at EPIC and ab-initio levels respectively, being 37.43 kJ mol⁻¹ with BSSE correction. Although the ab-initio SCF interaction energies without BSSE correction are typically 70 to 80% for 18C6·nH₂O (n=1, 2, 3) and 4), the qualitative picture of hydration pattern of 18C6 does not change. The MESP minimum value of crown ether to which this fourth water molecule binds is about -242.6 kJ mol⁻¹. Further details of energetics and structures of $18C6 \cdot nH_2O$ for n = 1 to 6 are being published¹⁶ elsewhere.

The question now is: if two such crown ether moieties are brought together, will this fourth water molecule.

flip and get attached to the MESP minimum of a water molecule bonded to the other 18C6.4H,O? The answer to this question can be sought by bringing an 18C6.4H₂O and an 18C6-3H,O optimally together along with a water molecule in between them with a hydrogen 'sticky end', binding to the deepest MESP minimum of 18C6.4H₂O (viz. -356.7 kJ mol⁻¹). EPIC model has been employed for this purpose and the binding energy of this water molecule turns out to be 56.8 kJ mol⁻¹ in contrast to that for the fourth water molecule in the 18C6.4H,O viz. 43.7 kJ mol⁻¹. The corresponding EPIC structure is shown in Figure 6. It is gratifying to note that this structure is in excellent agreement with the crystal structure recently reported by Albert et al.6. In conclusion, the lock and key arrangement engendered by the MESP topography followed by electrostatic modelling seems to offer an excellent tool for understanding weak intermolecular interactions in supramolecular chemistry.

^{1.} See, for a lucid and comprehensive review, Lehn, J.-M. Nobel Lecture delivered on 8 December 1987 taken from *Nobel Lectures in Chemistry 1981-1990* (ed. Malström, Bo G.), World Scientific, Singapore, 1992.

^{2.} Dougherty, D. A., Science, 1996, 271, 163-168; Dougherty, D. A., in Comprehensive Supramolecular Chemistry (ed. Lehn, J.-M), Pergamon, New York, 1997; Mecozzi, S., West, A. P. and Dougherty, D. A., J. Am. Chem. Soc., 1996, 118, 2307-2308.

^{3.} For a collection of related articles, cf. Modeling the Hydrogen Bond (ed. Smith, D. A.), American Chemical Society, Washington DC, 1994. See, in particular, the articles by Gao, J., pp. 7-21; Keith, T. A. and Frisch, M. J., pp. 22-35 and Ryan, M. D., pp. 36-59.

^{4.} Jeffrey, G. A., An Introduction to Hydrogen Bond, Oxford University Press, New York, 1997.

- 5 Gorb, L. and Leszczynski, J., J. Am. Chem. Soc., 1998 (in press). Zhanpeisov, N. U. and Leszczynski, J., Int. J. Quantum Chem., 1998 (in press).
- Albert, A. and Mootz, D. Z., Naturforsch., 1997, 52, 615-619, Mootz, D., Albert, A., Schaefgen, S. and Staben, D., J. Am. Chem. Soc., 1994, 116, 12045–12046.
- 7. Fukuhara, K., Ikeda, K. and Matsuura, H., Spectrochim. Acta., 1994, 50, 1619–1628. Fukuhara, K., Tachikake, M., Matsumoto. S. and Matsuura, H., J. Phys. Chem., 1995, 99, 8617–8623 and references therein.
- 8 Ranghino, G., Romano, S., Lehn, J.-M. and Wipff, G., J. Am. Chem. Soc., 1985, 117, 7873-7877; Thompson, M. A., J. Phys. Chem., 1995, 99, 4794-4804; Hori, K., Yamada, H. and Yamabe, T., Tetrahedron, 1983, 39, 67-73; Kowall, T. A. and Geiger, A., J. Phys. Chem., 1994, 98, 6216-6224.
- Scrocco, E. and Tomasi, J., Topics in Curr. Chem., 1973, 42, 75;
 Scrocco, E. and Tomasi, J., in Advances in Quantum Chemistry (ed. Lowdin, P.-O.). Academic Press, vol. II, 1975; Politzer, P. and Trublar, D. G. (eds). Chemical Applications of Atomic and Molecular Electrostatic Potentials, Plenum, New York, 1980; Pullman, B., Int. J. Quantum Chem. Quant. Biol. Symp., 1990, 17, 81.
- 10. Tomasi, J., Bonaccorsi, R. and Cammii, R. in Theoretical Models of Chemical Bonding (ed. Maksic, Z. B.), Springer, Berlin, 1991; Naray-Szabo, G. and Ferenczy, G. G., Chem. Rev., 1995, 95, 829; Murray, J. and Sen, K. D. (eds), Molecular Electrostatic Potentials: Theory and Applications, Elsevier, Amsterdam, 1996.
- GAMESS Schmidt, M. W., Baldridge, K. K., Boatz, J. A., Elbert, S. T., Gordon, M. S., Jensen, S., Koseki, S., Matsunaga, N., Niguyen, K. A., Su, S. J., Windus, T. L., Dupuis, M. and Montgomery, J. A., J. Comput. Chem., 1993, 4, 1347-1363; GAUSSIAN 94: Frisch, M. J., Trucks, G. W., Schlegel, H. B., Gill, P. M. W., Johnson, B. G., Robb, M. A., Cheeseman, J. R., Keith, T., Peterson, G. A., Montgomery, J. A., Raghavachari, K., Al-Laham, M. A., Zakrzewski, V. G., Oritz, J. V., Foresman, J. B., Ciosłowski, J., Stefanov, B. B., Nanayakkara, A., Challacombe, M., Peng, C. Y., Ayala, P. Y., Chen, W., Wong, M. W., Andres, J. L., Replogle, E. S., Gomperts, R., Martin, R. L., Fox, D. J., Binkley, J. S., Defrees, D. J., Baker, J., Stewart, J. J. P., Head-Gordon, M., Gonzalez, C. and Pople, J. A., GAUSSIAN Inc., Pittsburgh, PA, 1995.
- 12. The package UNIPROP developed by Gadre et al., see Shirsat, R. N., Bapat, S. V. and Gadre, S. R., Chem. Phys. Lett., 1992, 200, 373-378. The MESP, V(r) at a point r is defined as

$$V(\mathbf{r}) = \sum_{A} \frac{Z_A}{|\mathbf{R}_A - \mathbf{r}|} \int \frac{\rho(\mathbf{r}')}{|\mathbf{r}' - \mathbf{r}|} d^3 \mathbf{r}',$$

where Z_A is charge of nucleus A located at R_A and $\rho(r)$ denotes the molecular charge density. Negative value of V(r) provides information regarding electron-rich sites. The topographical analysis of this scalar field involves locating and characterizing the critical points (CPs), where the first-order partial derivatives of this field vanish. The nature of the function at these CPs is known by inspecting the signs of the eigenvalues of the respective Hessian matrix. These CPs are identified by the notation (rank, signature), where rank is number of nonzero eigenvalues of Hessian matrix, and signature is algebraic sum of the signs of the eigenvalues. The nondegenerate CPs of rank 3 could be any one of (3, +3) (minimum), (3, -3) (maximum) or saddles of the types (3, +1) and (3, -1). The package UNIVIS developed by Gadre and coworkers; see Limaye, A., Dattawadkar, S., Inamdar, P. V. and Gadre, S. R., J. Mol. Graph., 1996, 14-19.

13. Gadre, S. R., Bhadane, P. K., Pundlik, S. S. and Pingale, S. S., in *Molecular Electrostatic Potential: Concepts and Applications* (eds Murray, J. and Sen, K. D.) Elsevier, Amsterdam, 1996, vol. 3,

pp. 219-255; In the EPIC model described in this work, rigid species A and B are moved relative to each other so as to minimize interaction energy

$$E_{\text{int}} = \frac{1}{2} + \sum_{i} V_{4,i} q_{R,i} + \sum_{j} V_{B,j} q_{A,j} \},$$

where, $V_{A,i}$ and $V_{B,j}$ are the MESPs due to species A and B respectively at the i^{th} and j^{th} nuclear site in species B and A respectively. Here, the MESP-derived charges are $q_{B,i}$ and $q_{A,j}$ respectively. The movement is constrained such that the van der Waals surfaces of A and B do not interpenetrate.

- 14. The FORTRAN program performing charge fitting to molecular electrostatic potentials was made available from C. Chipot, University De Nancy I, France. We thank Dr C. Chipot for supplying us this code. See also, Chipot, C., Ángyán, J. A., Ferenczy, G. G. and Scheraga, H. A., J. Phys. Chem., 1993, 97, 6628-6636.
- Gadre, S. R. and Pundlik, S. S., J. Phys. Chem., 1997, B101, 3298-3303.
- 16. Pingale, S. S. and Gadre, S. R. To be published, 1998.

ACKNOWLEDGEMENTS. We thank CSIR, New Delhi, for financial assistance.

Received 16 July 1998; revised accepted 17 September 1998