

A coupled physical-biological-chemical model for the Indian Ocean

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A coupled physical-biological-chemical model has been developed at C-MMACS. for studying the time-variation of primary productivity and air-sea carbon-dioxide exchange in the Indian Ocean. The physical model is based on the Modular Ocean Model, Version 2 (MOM2) and the biological model describes the nonlinear dynamics of a 7-component marine ecosystem. The chemical model includes dynamical equation for the evolution of dissolved inorganic carbon and total alkalinity. The interaction between the biological and chemical model is through the Redfield ratio. The partial pressure of carbon dioxide ($p\text{CO}_2$) of the surface layer is obtained from the chemical equilibrium equations of Peng *et al* 1987. Transfer coefficients for air-sea exchange of CO_2 are computed dynamically based on the wind speeds. The coupled model reproduces the high productivity observed in the Arabian Sea off the Somali and Omani coasts during the Southwest (SW) monsoon. The entire Arabian Sea is an outgassing region for CO_2 in spite of high productivity with transfer rates as high as $80 \text{ m-mol C/m}^2/\text{day}$ during SW monsoon near the Somali Coast on account of strong winds.

1. Introduction

One of the major aims of the Joint Global Ocean Flux Study (JGOFS) programme is the quantification of CO_2 fluxes across the ocean-air interface on a basin-wide scale. The reason for this is quite clear; the sources and sinks for the carbon cycle are required with reasonable accuracy as inputs for global change studies. While there has been an intense observational programme in the Arabian Sea in the past 10 years, reliable estimates of CO_2 fluxes and carbon sequestering by the deeper oceans on a basin-wide scale still require the use of ocean general circulation models coupled to biological and chemical models.

It is well known that the Arabian Sea is one of the highest biologically productive regions of the world, especially in the SW monsoon season. Upwelling of deeper ocean waters off the Somali Coast tends to bring up an abundance of nutrients essential to the growth of phytoplankton. In addition, inorganic carbon is also pumped into the upper layers. Photosynthesis fixes

inorganic carbon reducing the partial pressure of CO_2 . Thus there are two competing influences; physics which increases the available inorganic carbon and biology which consumes it. Only a model which incorporates all three; i.e., physics, biology and chemistry can answer the question posed by JGOFS.

Till date there have been very few studies which have coupled physical and biological models in the Indian Ocean (McCreary *et al* 1996, Ryabchenko *et al* 1998, Keen *et al* 1997) on a basin-wide scale. Often these studies employ a 2.5 or 3 layer model for physics similar to the ones developed by either McCreary (McCreary *et al* 1993) or Hasselman (Hasselmann 1982). The top layer is representative of the mixed-layer in which the coupled equations are solved. On the other hand, there have been several such studies for the Atlantic Ocean (Sarmiento *et al* 1993, Oschlies and Garcon 1998, 1999, Drange 1994). These models have been quite successful in reproducing several observed features in the Atlantic. Of all the models cited above, only Drange has coupled chemistry on a

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basin-wide scale. The authors are not aware of such studies in the Indian Ocean.

The computational complexity and large demand on computer resources necessitate simplification of the biological and chemical components. Except Drange's and Sarmiento's studies, all the others use a simplified 4 component (phytoplankton, zooplankton, nitrate and detritus) biology model. There is considerable debate even today on how many of such compartments are needed to model the ecosystem with sufficient accuracy. However, as our aim is to get the carbon flux on a large scale, we will not dwell too much on this aspect.

2. Coupled physical-chemical-biological model

2.1 Physics

Our physical model is based on Modular Ocean Model Ver 2 (Pacanowski 1995; Bryan 1969), which is perhaps the most versatile model in existence today. It is a fully 3-dimensional model with good representation of bottom topography and multiple subgrid scale mixing options. The Boussinesq approximation is invoked and the vertical momentum equation is approximated by the hydrostatic equation. The numerics of MOM splits the velocities into two modes; the barotropic or external mode and baroclinic or internal mode. The former is solved by integrating the momentum equations over depth and applying a rigid lid condition on top to yield a volume transport stream function. The curl of the vertically integrated momentum equations to remove unknown surface pressure yields a Poisson equation which is solved by the application of a conjugate gradient algorithm. Islands in the domain are handled by the application of Stokes theorem. The three-dimensional baroclinic equations are solved by explicit time-stepping with a leap-frog scheme with a time-mixing euler backward step at every 17th time step. In addition to the momentum equations, tracer equations for temperature and salinity are solved concurrently and these quantities together with the UNESCO equation of state determine the density used in the pressure gradient term in the momentum equations. The advection terms in the equations are normally handled with central differencing, although we use upwind differencing for tracers as will be described in the next section. The diffusion terms are much more complex; for horizontal diffusion we use Laplacian term with constant viscosity while the vertical diffusion is handled by the application of the Pacanowski-Philander scheme (Pacanowski and Philander 1981). In this scheme, the stability of the water column is first determined based on the gradient Richardson number and the vertical diffusion coefficient

for momentum and tracers are determined according to the following formulae:

$$\kappa_m = \frac{C}{(1 + 5Ri)^2} + \kappa_{bm}, \quad (1)$$

$$\kappa_T = \frac{C}{(1 + 5Ri)^3} + \kappa_{bT} \quad (2)$$

where C is a constant ($50 \text{ cm}^2/\text{s}$), the background coefficients κ_{bm} and κ_{bT} are $1 \text{ cm}^2/\text{s}$ and $0.1 \text{ cm}^2/\text{s}$, respectively. The Richardson number in the above equations are defined as

$$Ri = \frac{-\frac{g}{\rho} \frac{\partial \rho}{\partial z}}{\left(\frac{\partial u}{\partial z}\right)^2 + \left(\frac{\partial v}{\partial z}\right)^2} \quad (3)$$

where u and v are the zonal and meridional velocities, g is the acceleration due to gravity, and the denominator is indicative of the velocity shear. High values of the denominator accompanied by marginal static stability (numerator) results in very small values of Ri . However, the constant C in equations (1) and (2) ensures that the diffusion coefficients do not exceed $50 \text{ cm}^2/\text{s}$. Although this scheme is supposed to take care of large shears which occur in the upper ocean, the limit on the diffusion coefficient does not permit the deepening of mixed layer as much as a scheme which employs a fully turbulent closure model such as Mellor-Yamada (Mellor and Yamada 1982) or the Gaspar scheme (Gaspar *et al* 1990; Blanke and Delecluse 1993) may allow. In any case, mixing in the ocean is an area of intense study (Gent and McWilliams 1990, Large *et al* 1994) and newer parameterisations are constantly being evolved. For this study, Pacanowski-Philander mixing scheme does a reasonable job and is sufficient for the goal we have in mind i.e. to estimate basin-wide CO_2 efflux from or to the ocean. In addition, gravitational instability is handled after time-stepping by an explicit convective mixing scheme. The model scans through the water column at each horizontal grid point and checks the stability of the water column. Unstable adjacent vertical regions are mixed instantaneously. The scan is repeated several times until the whole column is stable. The alternative to this scheme is to do the vertical diffusion implicitly with very high values for the vertical diffusion coefficient in the unstable regions. Tests by Pacanowski have shown that the explicit convective scheme is quick and reliable (Pacanowski 1995).

The details of the MOM run are as follows: The domain is 15°S to 26°N and 38°E to 100°E . The grid resolution is 0.5° in the longitude. In the latitude, the region between 15°S and 5°S has a resolution of 1° , the region between 5°S and the Equator has a resolution of 0.43° and the region northwards of the Equator has a resolution of 0.33° . There are 20 vertical levels with 10 in the top 100 metres. The total number of grid points is $126 \times 88 \times 20 (= 221,760)$. On the southern boundary between 15°S and 5°S we apply a sponge

which restores the temperature and salinity values to interpolated Levitus values with a time scale of 15 days. The boundaries on the east and west are closed; on the east there is no provision for the Indonesian Throughflow. One of the main reasons for this restrictive domain is the computational cost when additional tracers for biology and chemistry are introduced in the model. It is expected that the model will perform quite well away from the southern and eastern boundaries on not too long time scales.

The horizontal mixing is parametrized with a Laplacian operator and a constant diffusion coefficient of $1.0E8 \text{ cm}^2/\text{s}$ for momentum and half this value for temperature and salinity. The time step of integration is 2700 seconds which is well below the CFL limit for explicit time-stepping. The bathymetry of the Indian Ocean is taken from Scripps 1° topography data after smoothing the topographic roughness to eliminate breakdown of the external mode. Annual mean temperatures and salinities from the Levitus data base (Levitus 1982) are interpolated to each grid point and the model is started from rest. It is forced on the surface by Hellerman and Rosenstein (Hellerman and Rosenstein 1983) monthly wind stresses interpolated to each surface grid at each time. The solar flux of the heat flux component at the surface is explicitly applied from the Oberhuber monthly climatology (Oberhuber 1988). For the penetration of solar radiation into the upper layers of the ocean, we apply the absorption coefficients described by Paulson and Simpson (Paulson and Simpson 1977, Pacanowski 1995). For the other components of the heat flux, we apply a restoring condition, i.e., Newtonian damping with a time scale of 50 days and a depth scale of 10 m, where an equivalent flux is computed based on the difference between the first level model temperatures and Levitus SST (Levitus 1982) values. A similar procedure is applied to salinity. The physical model is integrated for 35 years to reach an acceptable annual cycle before biology and chemistry are introduced. The model is integrated for a further 11 years with biology and with chemistry in the last two.

2.2 Biology

The biological model is very similar to the one of Sarmiento *et al.* (Sarmiento *et al.* 1993), and Sharada and Yajnik (Sharada and Yajnik 1997) but with some subtle differences. The model has seven compartments, phytoplankton, zooplankton, bacteria, nitrate, ammonia, particulate organic nitrogen and dissolved organic nitrogen. The common unit of currency is moles of nitrogen/ m^3 ; in other words, it is a nitrogen-based ecosystem model. There is still considerable debate about the use of nitrate exclusively as the nutrient. There are studies which use iron, phosphates and silicates as the limiting nutrients (Peng *et al.* 1987) but for our purposes, nitrogen is adequate. The

biological quantities are introduced in MOM2 as tracers with the following equation:

$$\frac{\partial T}{\partial t} + L(T) = \frac{\partial}{\partial z} \left(\kappa_T \frac{\partial T}{\partial z} \right) + A_h \nabla_h^2 T + SMS(T) \quad (4)$$

where T stands for any tracer (temperature, salinity, biological or chemical species), $L(T)$ is the advective term, the first term on the right is the vertical diffusion term and the second is the horizontal diffusion term and $SMS(T)$ is the source minus sink term. The advective term is handled by central differencing in the momentum equations to ensure balances in long term integration but this is not normally acceptable with tracer equations as it causes too many oscillations in the solution. We first invoked a flux-corrected algorithm (FCT) which is a blend of upwind and central differencing schemes (Pacanowski 1995) but this proved to be too demanding computationally as we have 11 tracers in all. So we modified MOM code to do simple upwinding. Although this is better than central differencing, the scheme is quite diffusive as shown by Oschlies and Garcon (Oschlies and Garcon 1999). There is one term which is not shown explicitly in the above equation and that is the convective adjustment term in the presence of gravitational instability. The SMS terms for each of the biological components are identical to those in Sarmiento (Sarmiento *et al.* 1993) but are given again below for future reference in the article.

$$SMS(P) = (1 - \gamma_1) \bar{J}(z, t) [Q_1(N_n, N_r) + Q_2(N_r)] P - G_1 - \mu_1 P, \quad (5)$$

$$SMS(Z) = \gamma_2 (G_1 + G_2 + G_3) - (\mu_2 + \mu_5) Z, \quad (6)$$

$$SMS(B) = U_1 + U_2 - G_2 - \mu_3 B, \quad (7)$$

$$SMS(N_n) = -\bar{J}(z, t) Q_1(N_n, N_r) P, \quad (8)$$

$$SMS(N_r) = [\gamma_3 \mu_2 + (1 - \gamma_4) \mu_5] Z + \mu_3 B - \bar{J}(z, t) Q_2(N_r) P - U_2, \quad (9)$$

$$SMS(N_d) = \gamma_1 \bar{J}(z, t) [Q_1(N_n, N_r) + Q_2(N_r)] P + (1 - \gamma_3) \mu_2 Z + \mu_4 N_p - U_1, \quad (10)$$

$$SMS(N_p) = \mu_1 P + (1 - \gamma_2) (G_1 + G_2 + G_3) - G_3 - \mu_4 N_p - w_s \frac{\partial N_p}{\partial z} \quad (11)$$

where P stands for phytoplankton, Z for zooplankton, B for bacteria, N_n is nitrate, N_r is ammonium, N_d is dissolved organic nitrogen and N_p is detritus. There are two nutrient sources for primary production, nitrate and ammonium, and the uptake functions are given by

$$Q_1(N_n, N_r) = \frac{N_n}{K_1 + N_n} e^{-\psi N_r}, \quad (12)$$

$$Q_2(N_r) = \frac{N_r}{K_2 + N_r}. \quad (13)$$

The exponential term in equation (12) limits the uptake of nitrate in the presence of ammonium. The average irradiance is given by

$$\bar{J}(z, t) = \frac{1}{z_k - z_{k+1}} \int_{z_{k+1}}^{z_k} J(z, t) dz \quad (14)$$

where z_k 's are the depths of the vertical boxes and $J(z, t)$ is given by

$$J(z, t) = \frac{V_p \alpha I(z, t)}{[V_p^2 + (\alpha I(z, t))^2]^{0.5}} \quad (15)$$

and the irradiance $I(z, t)$ is given by

$$I(z, t) = I(0, t) \exp\left(k_w z - \int_z^0 k_c P dz\right) PAR \quad (16)$$

and where $V_p = ab^{cT}$ and $I(0, t)$, the irradiance at the top of the ocean taken from Oberhuber's solar radiation atlas, PAR is the fraction of solar radiation in the photosynthetically active region. Note that we account for absorption by both water as well as phytoplankton in equation (16). All the parameters are defined with their numerical values in table 1. The main difference between Sarmiento *et al* (1993) and this study is the way we have taken solar radiation into account. We do not resolve the diurnal variation and we do not make a distinction between direct and

diffuse contributions. The average solar radiation is taken to account for both.

Zooplankton in the model grazes on phytoplankton, bacteria and detritus. We do not allow any self-grazing unlike McCreary (McCreary *et al* 1996). The preference for the three food sources are given by

$$G_j = gZ \frac{p_j C_j}{K_3 + \sum_{k=1}^3 p_k C_k} \quad (17)$$

where $j = 1, 2, 3$ and $C_1 = P$, $C_2 = B$ and $C_3 = N_p$. The uptake of the three forms are related to their relative abundance as follows:

$$p_j = \frac{\varrho_j C_j}{\sum_{k=1}^3 \varrho_k C_k} \quad (18)$$

by means of the preferences ϱ_j . These are again given in table 1. There is a slight modification to the bacterial uptake terms based on a suggestion by Drange (Drange 1994) as follows:

$$U_1 = \frac{V_b B \frac{S}{\eta}}{K_4 + S \left(1 + \frac{1}{\eta}\right)} \quad (19)$$

$$U_2 = \eta U_1 \quad (20)$$

and $S = \min(N_r, \eta N_d)$ where η is the ratio of ammonium to dissolved organic nitrogen that is

Table 1. *Parameters in the biology model.*

Parameter	Symbol	Value	Units
Exudation fraction of P	γ_1	0.05	
Specific mortality rate of P	μ_1	0.04	d^{-1}
Half-saturation constants for nutrient uptake of P	K_1, K_2	0.5	mmol/m^3
Ammonium inhibition parameter	Ψ	1.5	$(\text{mmol/m}^3)^{-1}$
Initial slope of PI curve	α	0.025	$d^{-1}/(\text{Wm}^{-2})$
Light attenuation due to water	k_w	0.04	m^{-1}
Light attenuation due to phytoplankton	k_c	0.03	$\text{m}^{-1}(\text{mmol/m}^3)^{-1}$
Photosynthetically active radiation fraction	PAR	0.4	
Maximum growth rate parameters of P	a	0.6	d^{-1}
	b	1.066	
	c	1.0	C^{-1}
Zooplankton assimilation efficiency	γ_2	0.4	
Ammonium fraction of Z excretion	γ_3	0.75	
Detrital fraction of Z mortality	γ_4	0.33	
Z specific excretion rate	μ_2	0.1	d^{-1}
Z mortality rate	μ_5	0.05	d^{-1}
Z maximum growth rate	g	1.0	d^{-1}
Half saturation for ingestion of Z	K_3	1.0	mmol/m^3
Relative preference of Z for phytoplankton	ϱ_1	0.4	
Relative preference of Z for bacteria	ϱ_2	0.3	
Relative preference of Z for PON	ϱ_3	0.3	
Specific bacterial excretion rate	μ_3	0.05	d^{-1}
Maximum bacterial growth rate	V_b	2.0	d^{-1}
Half-saturation rate for uptake by bacteria	K_4	0.5	mmol/m^3
Ammonium/DON uptake ratio for bacteria	η	0.6	
Detrital breakdown rate	μ_4	0.05	d^{-1}
Detrital sinking velocity	w_s	-4.0	m/d

Most of the above parameters are identical to Sarmiento *et al*, 1993. However there are some notable changes; The assimilation efficiency of Z , γ_2 is 0.4 to limit the run-away growth of Z . In addition, the relative preference of Z for P , ϱ_1 is reduced to 0.4 with slight increases for B and PON. Finally the sinking velocity for PON, w_s is 4 m/d instead of 10 in Sarmiento. The regeneration equation for $z > 123\text{m}$ are identical to Sarmiento.

required for bacteria to obtain enough nitrogen to consume carbon in dissolved organic matter, and V_b is the growth rate. The parameters are listed in table 1. Equations (5–20) are valid in the euphotic zone which is taken to be 123 m. Below this depth, we apply remineralisation equations which are identical to Sarmiento's equation (7). (Sarmiento *et al* 1993).

We prescribe initial nitrate values from Levitus (Conkright *et al* 1994) climatological atlas at every grid point. For phytoplankton we prescribe a small surface value of 0.14 mmol-N/m^3 , for the others 0.014 mmol-N/m^3 ; all of which decrease exponentially with a scale depth of 100m. We do not allow any transfer of the ecosystem components across the air-sea interface. It was found that all the system lost all its memory of initial values very rapidly and settled quickly into an annual cycle with a small long-term trend after two years of integration.

2.3 Chemistry

The chemical component is quite similar to Drange's (Drange 1994) with some minor differences. The two components are dissolved inorganic carbon (DIC) and total alkalinity (TALK). Phytoplankton is the only carbon fixing component in the whole system. The tracer equations are again given by equation (4) with the *SMS* terms being:

$$\begin{aligned} SMS(DIC) = & -(1 - \gamma_1)\bar{J}(z, t)[Q_1(N_n, N_r) \\ & + Q_2(N_r)]PR_p - \frac{\partial}{\partial t}CaCO_3 + \mu_3BR_b \\ & + [\gamma_3\mu_2 + (1 - \gamma_4)\mu_5]ZR_p, \end{aligned} \quad (21)$$

$$\begin{aligned} SMS(TALK) = & -RHS(Eq.8) + RHS(Eq.9) \\ & - 2\frac{\partial}{\partial t}CaCO_3. \end{aligned} \quad (22)$$

In the above equations the $\frac{\partial}{\partial t}CaCO_3$ term accounts for the formation of biogenic calcium carbonate. Although this is not modelled explicitly, what we have done is to compute the net downward flux of organic matter at a depth of 123m. It is assumed that the biogenic $CaCO_3$, which is proportional to the downward flux of organic matter at 123m dissolved in the water column above (Drange, 1998). The flux of organic matter at the base of the euphotic zone is computed by summing the vertically advective and diffusive fluxes of phytoplankton, zooplankton and detrital matter. Twenty per cent of this flux is redistributed over the euphotic zone in an exponentially decreasing manner with an e-folding depth of 4000m (Drange 1994). The treatment of biogenic $CaCO_3$ in this study is rudimentary but was included on the advice of Drange. It was seen later that the effect of $CaCO_3$ is almost insignificant and needed to be improved substantially. In equation (21), R_p and R_b denote the Redfield ratios of phytoplankton and bacteria, respectively which convert the ecosystem

variables from nitrogen to carbon units. The value of R_p is taken to be 7 and R_b is taken to be 5. It is assumed that phytoplankton and zooplankton have the same Redfield ratio.

The main difference between Drange's formulation and ours is the last term in equation (21). While Drange has a quadratic term for zooplankton excretion and mortality, we have used the form consistent with equations (6) and (9) in the biology equation. The initialisation of the DIC and TALK fields have been taken from the objectively interpolated Atlas of Myrmehl and Drange (Myrmehl and Drange 1998). However, unlike Levitus's nitrate atlas, this atlas is very coarse ($2^\circ \times 2^\circ$) with considerable errors near continental boundaries.

2.3.1 Thermodynamics of carbonic acid

We need to estimate the pCO_2 of the surface waters in order to determine the transfer of CO_2 across the air-sea interface. The DIC that we have computed needs to be partitioned into three species, the dissolved CO_2 component (pCO_2), HCO_3^- and CO_3^{2-} components. Boric acid which is proportional to the salinity dissociates into $H_2BO_3^-$ and undissociated boric acid. We have employed the chemical equilibrium model of Peng (Peng *et al* 1987, Drange 1994) in our calculations. Unlike Peng, we do not include equations for phosphoric acid and silicic acid; we retain only carbonic acid, boric acid and water. The equations are identical to those found in the Appendix of Peng's paper. Drange (Drange 1994) has compared several other formulations and found Peng's formulation one of the most accurate. In a nutshell, we have four equations which correspond to two disassociations of carbonic acid, and one each for water and boric acid. Using in addition three mass balance equations for carbon, boron and total alkalinity, we get a system of 7 nonlinear equations to solve. These are solved by the application of HYBRD1 subroutine from the MINPACK library. This is the most compute intensive part of the whole calculation. In order to save time, the equilibrium calculations are made only for the surface grid points.

2.3.2 Transfer of CO_2 across the air-sea interface

Once the pCO_2 of the surface water is found the flux of CO_2 across the air-sea interface is found by the application of the following equation:

$$F = k_l\alpha_s(pCO_2^{\text{air}} - pCO_2^{\text{sea}}) \quad (23)$$

where F is the flux of CO_2 from the atmosphere to the ocean, k_l is the piston or transfer velocity and α_s is the solubility of CO_2 in sea water (Peng *et al* 1987). It must be noted here that MOM uses CGS units in its representation and we have to make sure that the units are consistent in equation (23). The pCO_2^{air} has

to be adjusted for 100 relative humidity. The piston velocity is regulated by the turbulence in the air-sea interface and the chemical reactions in the liquid phase. We apply the formulation suggested by Wanninkhof (Wanninkhof 1992) as follows:

$$k_l = 0.35U_{10}^2 \left(\frac{Sc}{660}\right)^{-0.5} \tag{24}$$

where k_l is in cm/s, U_{10} is the air speed in m/s at 10m (taken from Oberhuber’s monthly mean Atlas) and Sc is the Schmidt number given by

$$Sc = 2073.1 - 125.62T + 3.6276T^2 - 0.043219T^3 \tag{25}$$

where T is the temperature in degrees Celsius. It is not clear what value should be used for pCO_2^{air} as it is

increasing quite rapidly. We have done the calculations for two values, 345 ppm and 360 ppm. In contrast to other biological and chemical species, the flux computed by equation. (23) is added only to the top grid box while solving the DIC equation.

3. Results and discussion

The parameters used in the biological model are given in table 1. There are more than 30 parameters to be fixed for the biological model alone. It is not presently clear how well some of these are known. It is not unlikely that some of the parameters are specific to the location. Although the parameters have been revised a few times in the past depending on measurements in the Sub-Arctic Pacific Ocean

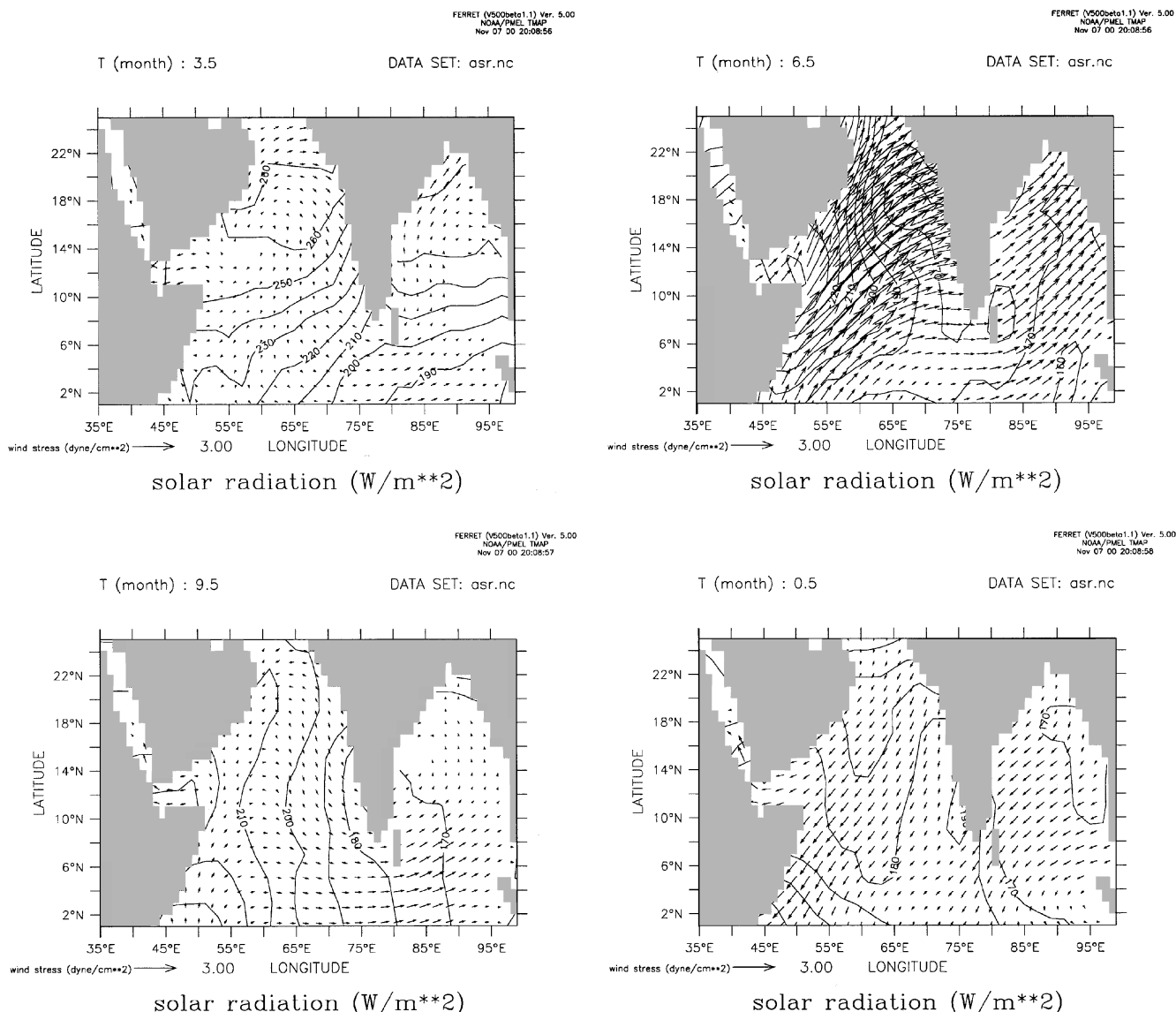


Figure 1. Solar radiation at the top of the ocean from Oberhuber’s Atlas and wind stresses from Hellerman and Rosenstein’s Atlas.

(Welschmeyer *et al* 1993) we have chosen to go with most of them identical to the ones in Sarmiento (Sarmiento *et al* 1993). The differences between the two (listed on the footnote of table 1) are concerned with the growth of *Z*. One of the major problems was the runaway growth of *Z*. In order to contain this the assimilation efficiency of *Z* was reduced to 0.4. In order to permit the growth of phytoplankton, the relative preference of *Z* for *P* was reduced to 0.4 from 0.5 in Sarmiento. Sharada and Yajnik (Sharada and Yajnik 1997) have done a careful sensitivity analysis of the basic biology model in comparison with some *in situ* data in the Central Arabian Sea and the numbers chosen here are in accordance with their study.

The solar radiation on the ocean surface (taken from monthly mean Atlas of Oberhuber 1988) and the windstresses (taken from Hellerman and Rosenstein

1983) at four typical months (April, July, October and January) are shown in figure 1. We will concentrate on the results on the northern Indian Ocean as the region south of the equator is modelled very coarsely. Notice the large reduction in solar flux during July on account of clouds but still the values are reasonably high (greater than 200 W/m^2) in the western Arabian Sea. Also notice the effect of SW monsoon in July and the NE monsoon in January. The wind speeds at 10m (taken from Oberhuber's Atlas) are shown in figure 2. Notice that the winds reach speeds of over 10m/s in most of the Arabian Sea during July and over 5m/s in January.

Fifteen day time averages of the horizontal velocities at 5m (the first vertical grid point in the model) and the vertical velocities at 20m are shown in figure 3 corresponding to the same periods shown in figures 1

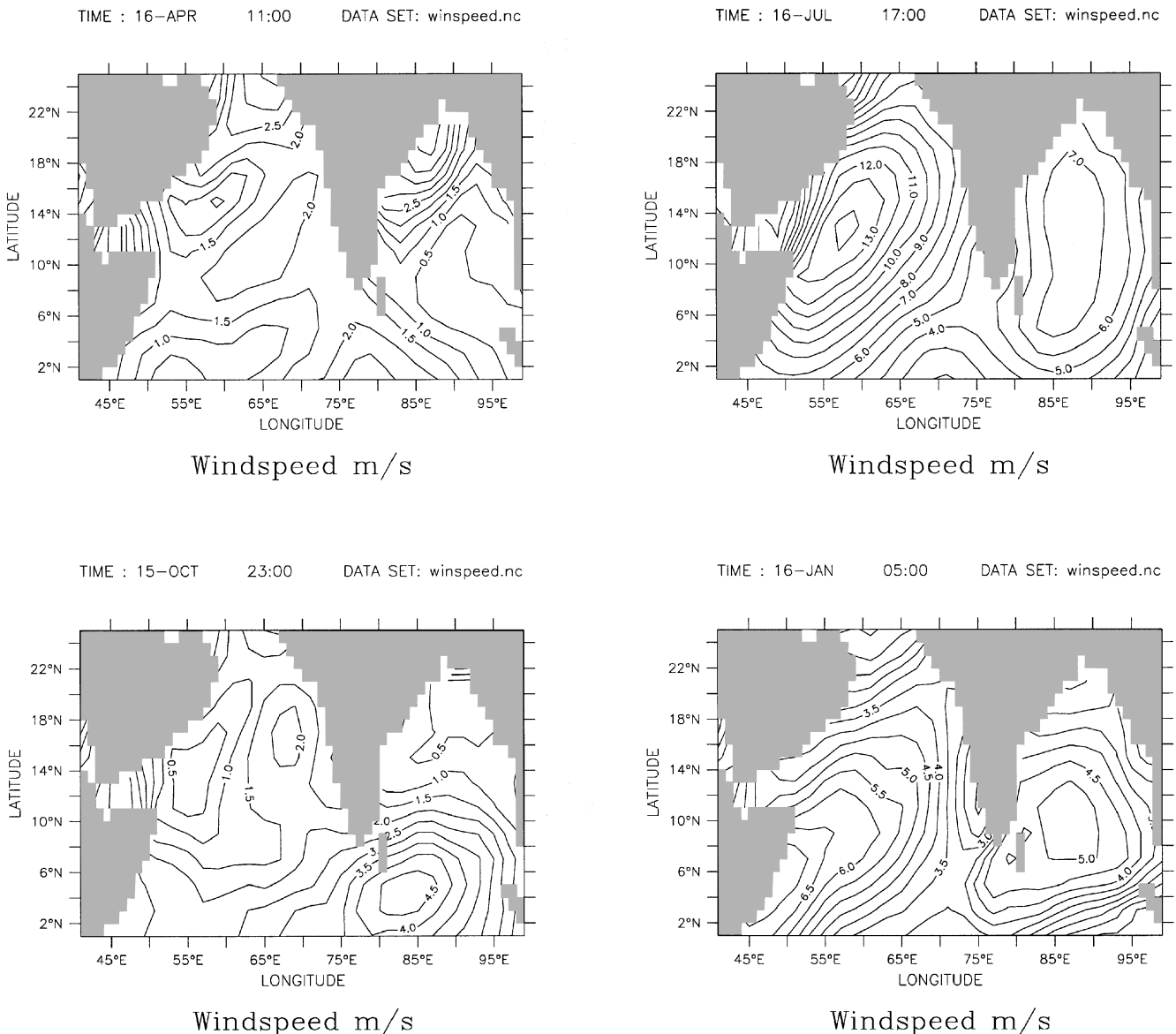


Figure 2. Wind speed at 10 metres from Oberhuber's Atlas.

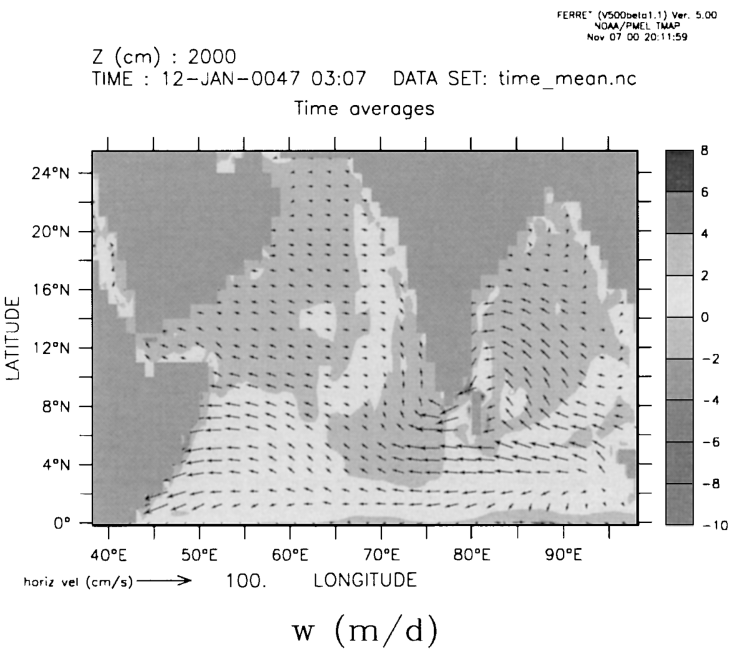
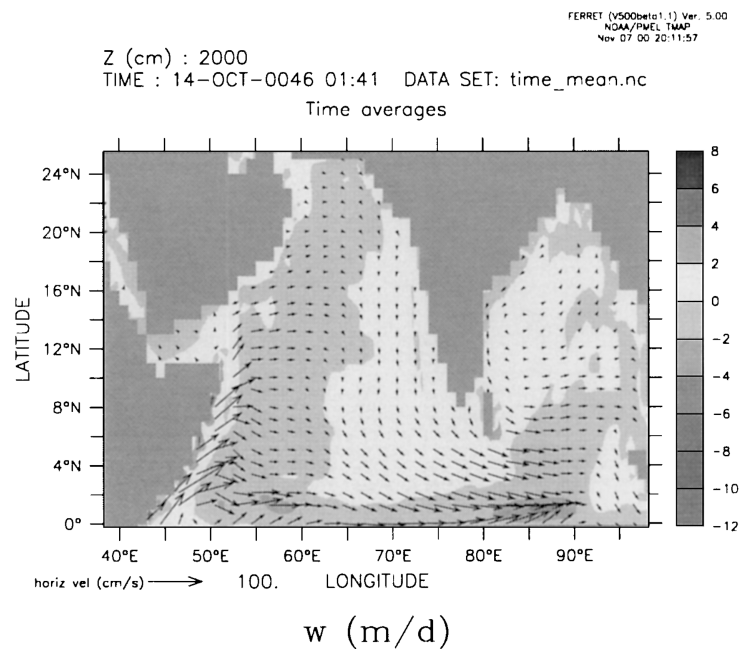
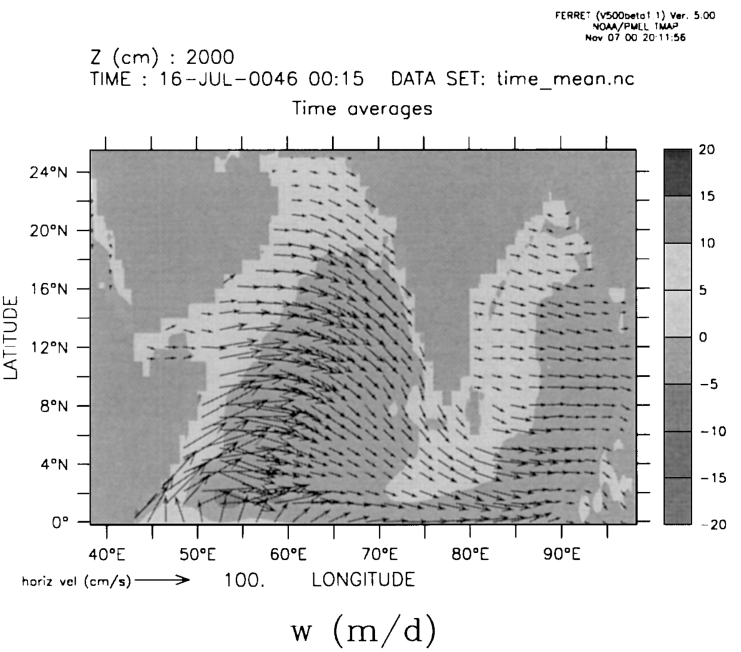
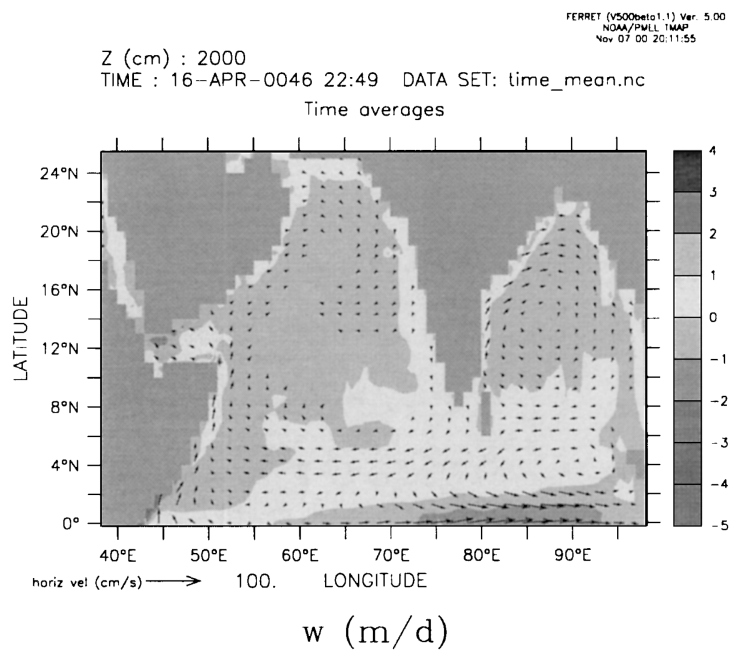


Figure 3. Fifteen day averages of horizontal velocities at 5 m depth (in cm/s) and vertical velocities at 20 m (in m/d).

and 2. The units for w are m/d while the units for horizontal currents is cm/s. Notice the strong upwelling (velocities greater than 10 m/d) that occurs off the Somali and Omani coasts and strong horizontal currents (speeds greater than 200 cm/s) which occur during SW monsoon. Most of central Arabian Sea shows weak upwelling or downwelling. Also notice the change in the direction of the North Equatorial Current between the SW and NE monsoon seasons as well as the change in direction of Somali currents. In addition, the region off the southern coast of India is mostly a weak upwelling region except in January.

The computed P distributions averaged over the top 25m of the ocean (first 3 vertical grids) are shown in figure 4. Keep in mind the differences in scale while viewing the picture. Notice the strong blooms off the Somali and Arabian Coasts (with values greater than 5 mmol-N/m^3) and blooms off the southern and eastern coasts of India during SW monsoon. This picture has a good qualitative resemblance to the CZCS climatological map of the same month. The mild blooms which appear between 50 and 60E and 0 and 4N during April and January seems to be artifacts. It is not clear at present what deficiencies in the model have produced these features. The time evolution of P along zonal sections are shown in figure 5. The last two years of integration are shown so that one may make inferences about the steady annual cycle that the model has reached. The onset of the bloom in the western Arabian Sea in June and its disappearance in September/October can be clearly seen. The vertical sections along the same latitudes in August can be seen in figure 6. Intense production can be seen in the first 30-50m in the bloom regions. Similar figures for Z are shown in figures 7-9. In contrast to P , Z reaches a maximum in September/October. This is a typical result of the predator-prey type of behaviour expected in such models. The reason for abundance of Z even in the dearth of P is the switching behaviour of Z to feed on bacteria and particulate organic nitrogen. The depth of zooplankton production is much larger than P precisely for the same reason; there is no light limitation as in the case of P growth.

The availability of nitrate in the upper ocean is shown in figure 10. Notice from the biological equations that nitrate has no source in the euphotic zone; it is only consumed. Remineralisation takes place only below 123m. Therefore nitrates which reach the surface are exclusively due to physics. This can be clearly seen in the April and July figures; from a nitrate concentration of less than 0.5 mmol-N/m^3 it reaches more than 10 mmol-N/m^3 in July near the Somali Coast indicating the strong upwelling which has brought up nitrate rich waters from below. The nitrate seen in October is completely consumed by January. Figure 11 shows the time series at four zonal sections for nitrate similar to figures 5 and 8. Figure

12 shows the nitrate distribution in the top 110m similar to figures. 7 and 9. The tongues of high nitrate concentrations clearly demarcate the regions where it is being upwelled from the deeper waters.

We will not discuss the dynamics of other biological compartments as they have no direct bearing on the CO_2 transfer. However, what is perhaps more edifying is a break up of the terms in the tracer equation (equation 4). We take the term balances in the region near the Somali Coast (52-56E,5-9N,0-50m). The term balances for nitrate are shown in figure 13. At the beginning of June there is a steep rise in the rate of change of nitrate and the reason is the combined effect of zonal and meridional advection in spite of the large consumption by phytoplankton. The diffusion terms do not contribute much. By August, consumption by P overcomes advective effects and we have negative rates of increase.

Similar terms for the P equation are shown in figure 14. The rate of change peaks in the beginning of June with advection and SMS terms being the main contributors. The largest negative contributor is the convective mixing component. The reason is that most of the phytoplankton is produced at the surface and the explicit convective adjustment scheme which mixes heavier surface waters with lighter waters at depth is very effective in removing excess P . By the beginning of June, the growth of Z further reduces the SMS term and growth becomes negative. It is not until November when Z is depleted that P shows a small positive growth. A more detailed break up of the growth term is shown in figure 15. In this figure, the source is broken up into a growth term, a grazing term and a mortality term. The growth term reaches a maximum of $0.75 \text{ mmol-N/m}^3/\text{d}$ in July but from August onwards it is smaller than the grazing term. The mortality term plays a relatively minor role. The terms which contribute to P growth are shown in figure 16. The steep drop in the availability of light (\bar{J}) after May is due to the self-shading effect of phytoplankton. Also shown in the same figure are the terms Q_1 and Q_2 in equation (5). The former indicates new production while the latter, regenerated production using ammonium as the nutrient. The ratio of new production to total production is termed as the f-ratio. It may be seen that once the production of Z picks up and consequently a source in the ammonium equation, regenerated production is the main contributor. In the months of May-July, it is new production which dominates with f-ratio reaching a maximum of 0.8 in June. The break up of the SMS terms for Z are shown in figure 17. Notice that growth exceeds mortality till September after which the reverse is true.

One point within this region (55.3 E, 7.4N) was chosen to perform a simulated XBT study (term balances at each grid point in the vertical column). Figures 18 and 19 show the terms in the equation for

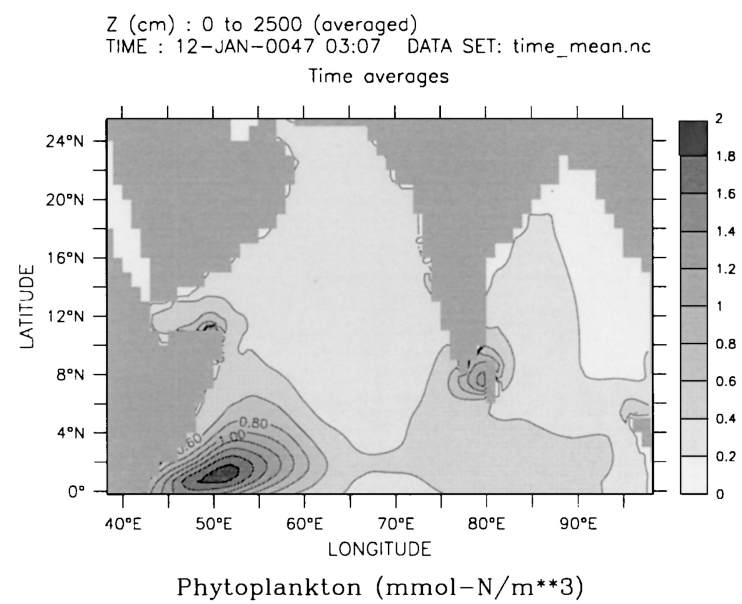
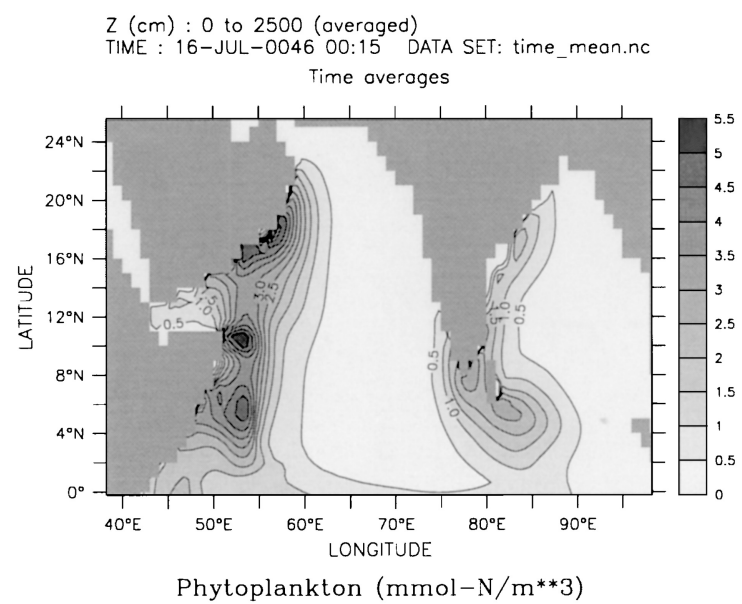
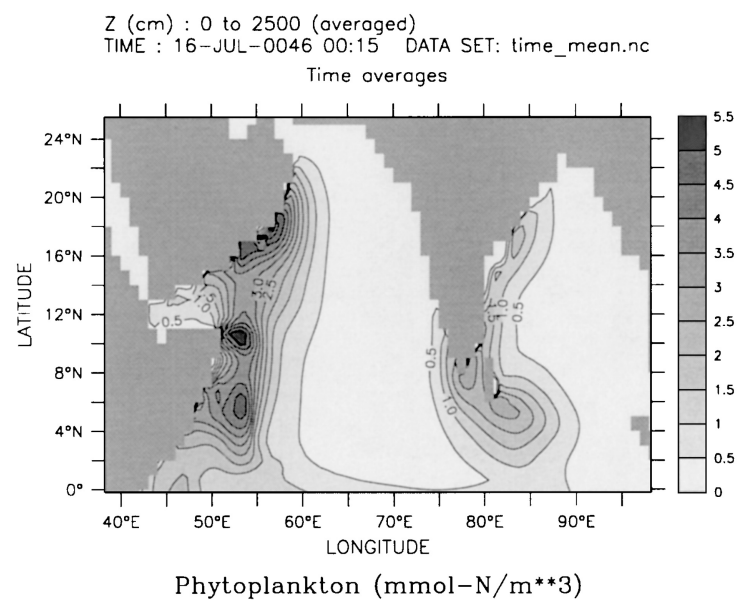
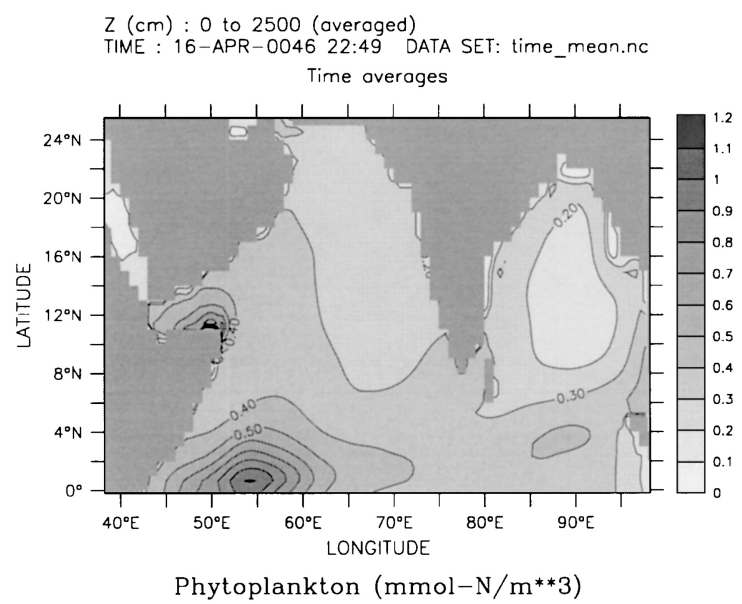


Figure 4. Time averages of phytoplankton distribution averaged in the top 25 m. Note the change of scale in different frames.

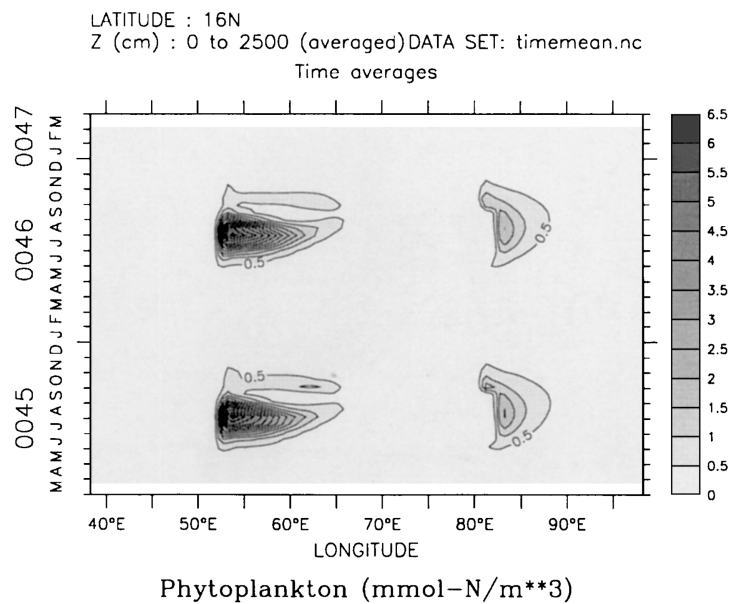
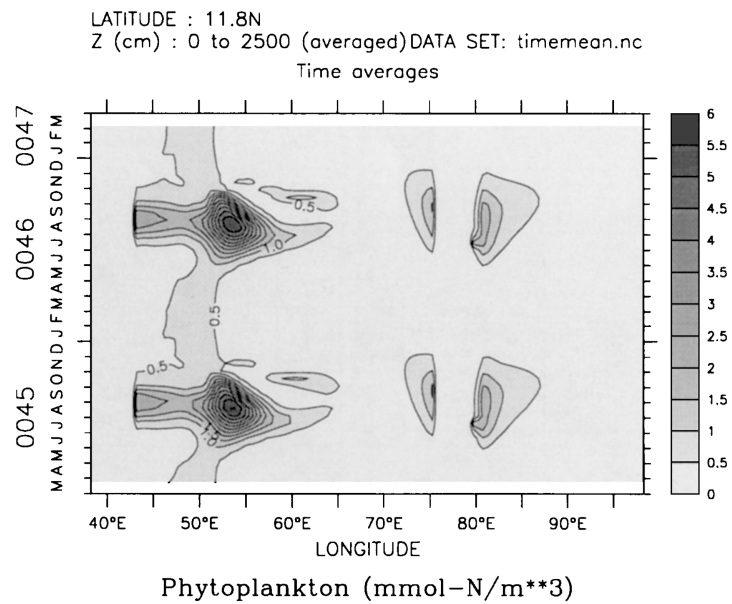
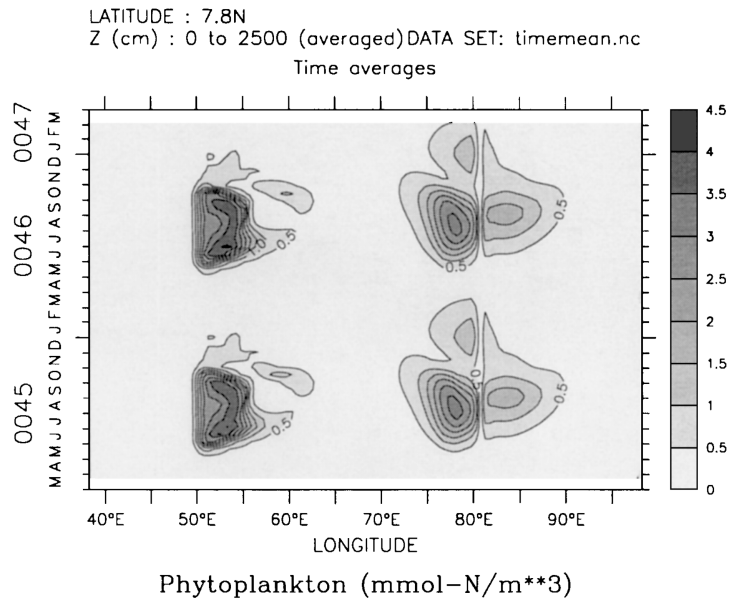
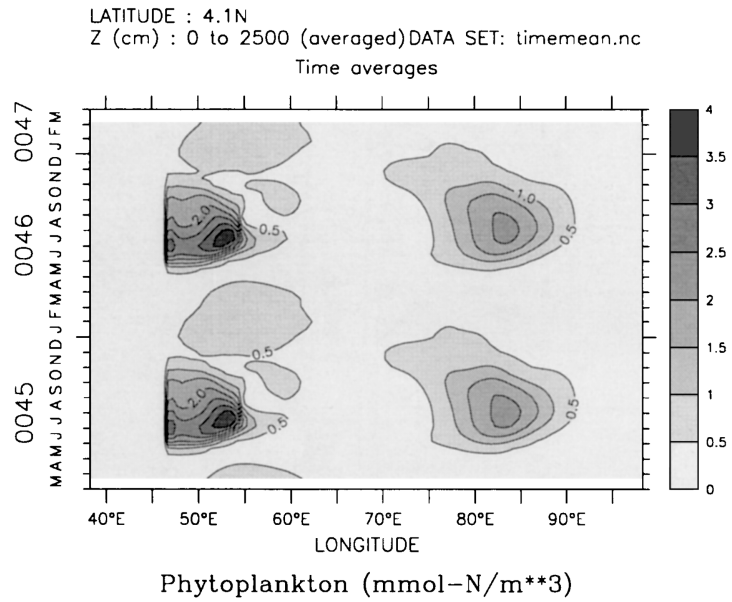


Figure 5. Time series of average phytoplankton distribution in the top 25 m at four latitudinal sections.

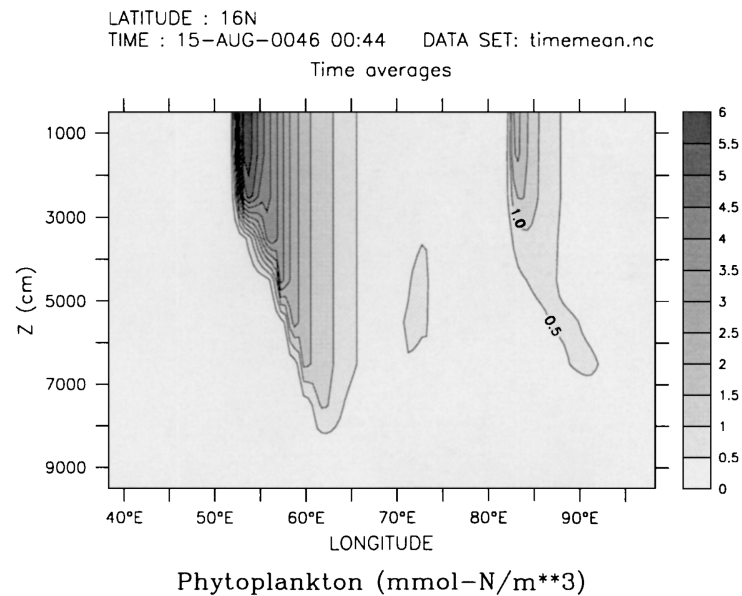
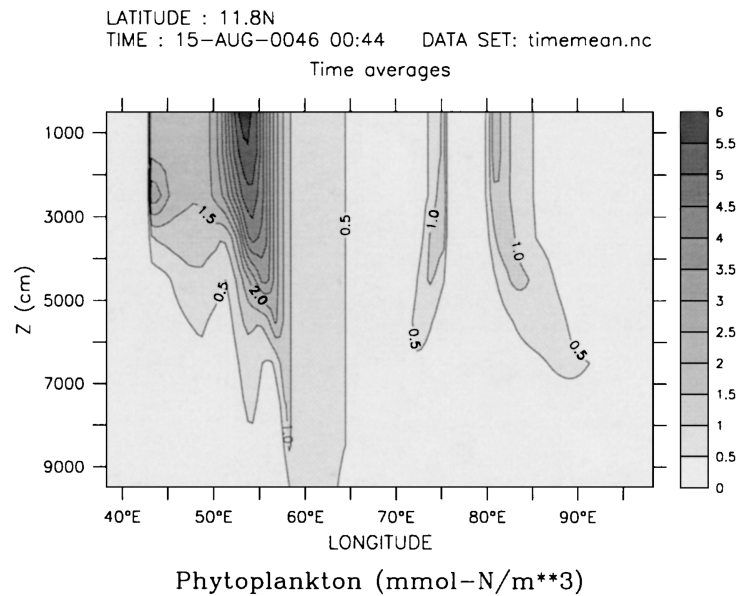
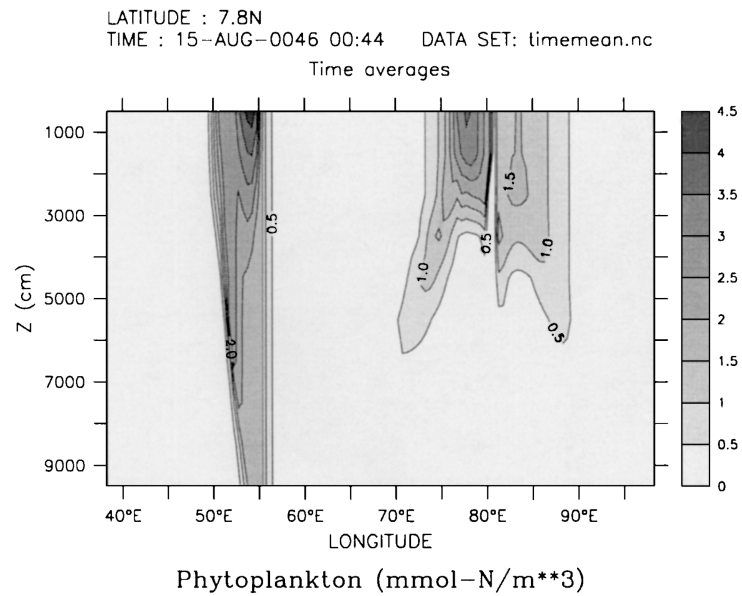
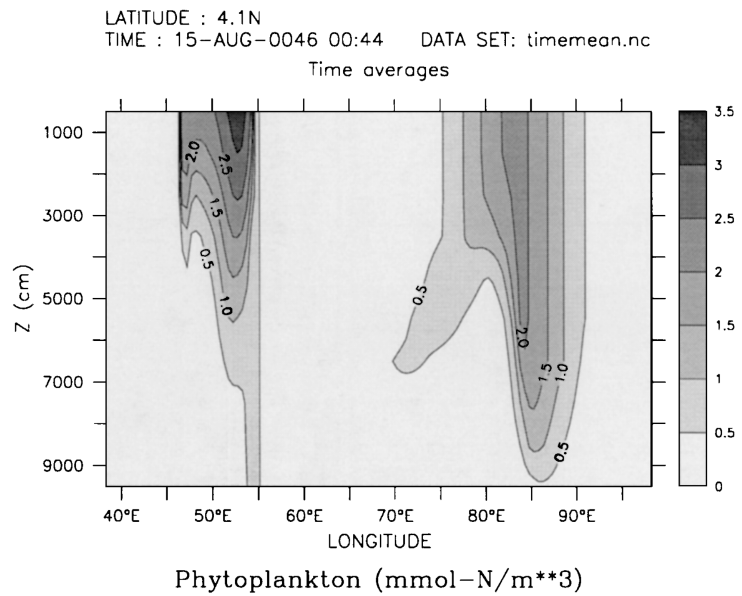


Figure 6. Depth-longitude sections of phytoplankton distribution in August at four latitudinal sections.

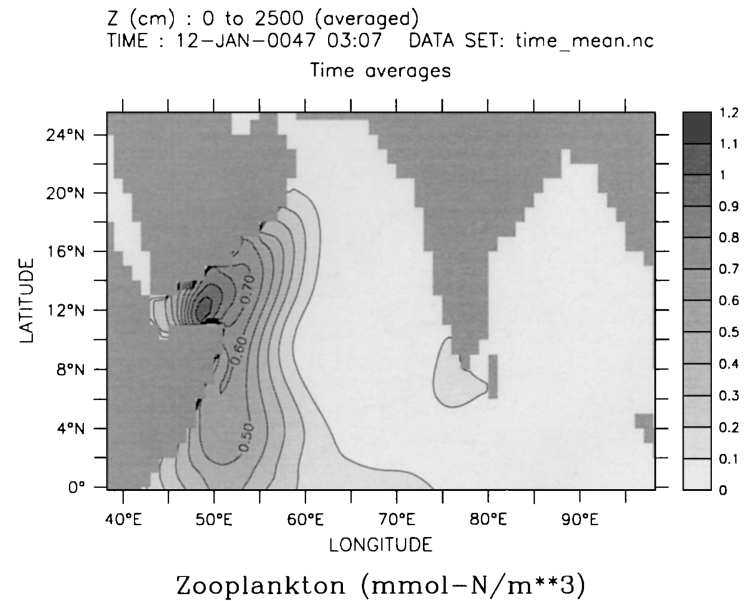
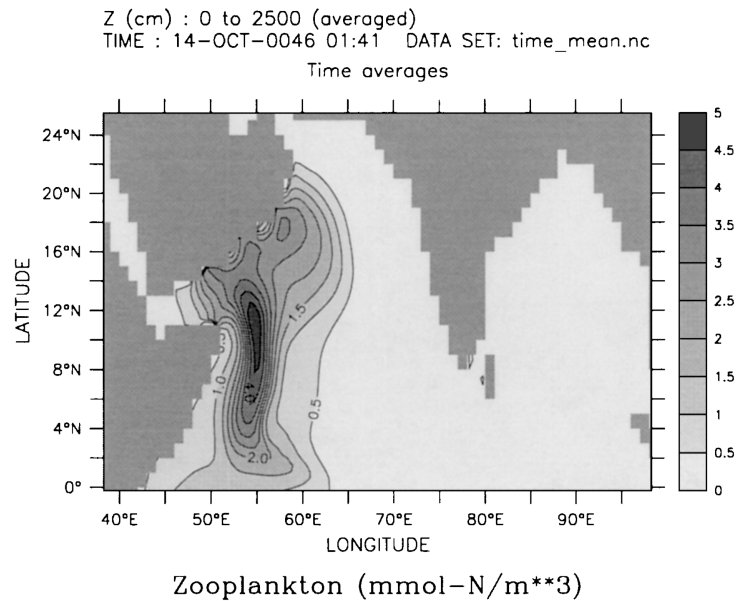
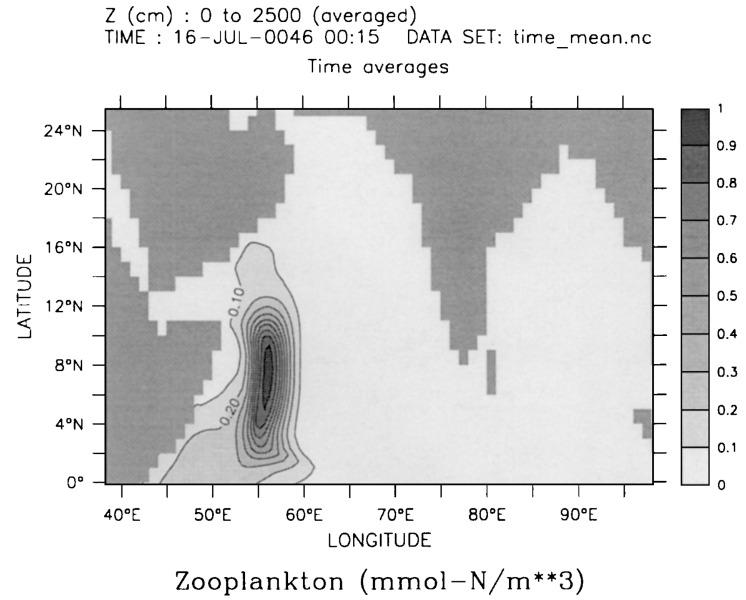
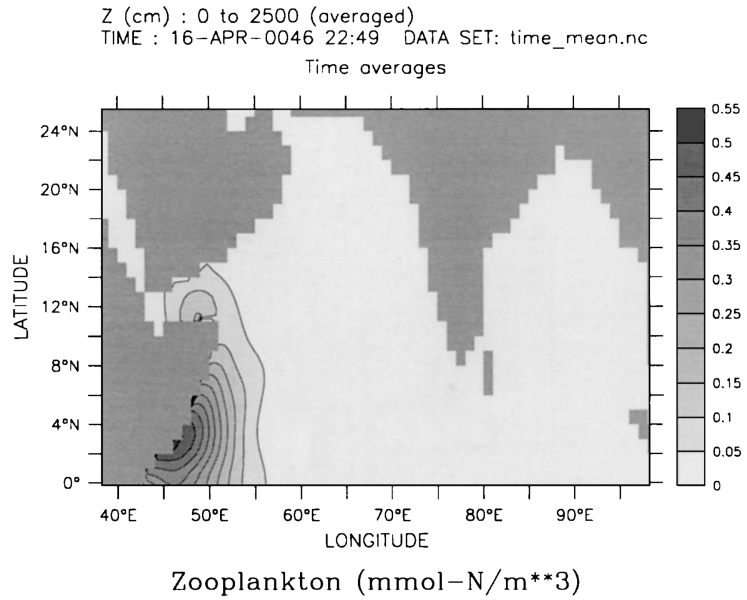


Figure 7. Time averages of zooplankton distribution averaged in the top 25 m. Note the change of scale in different frames.

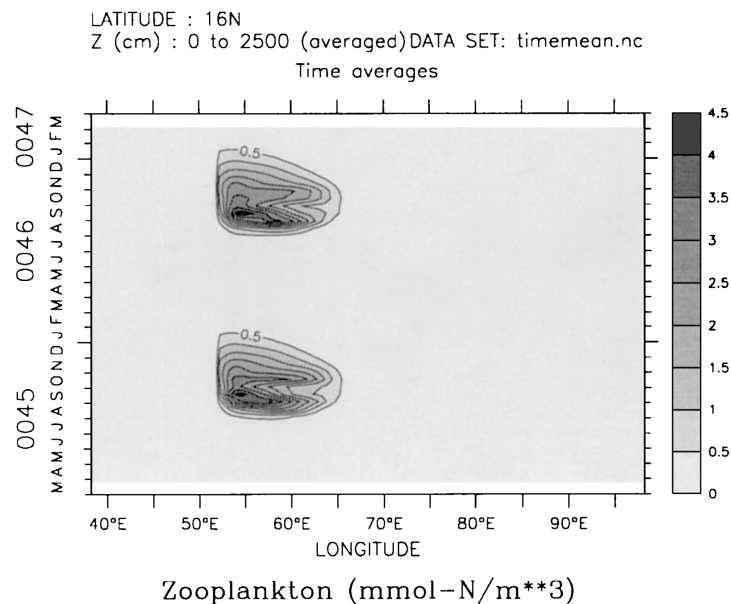
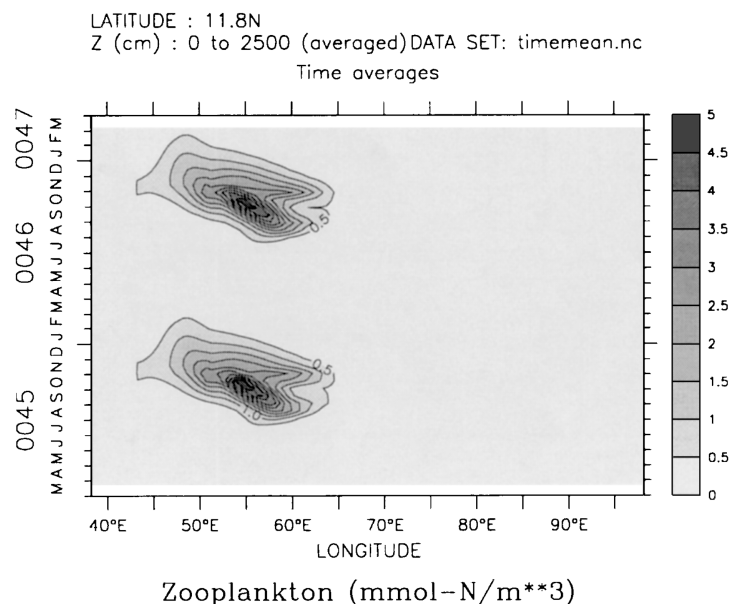
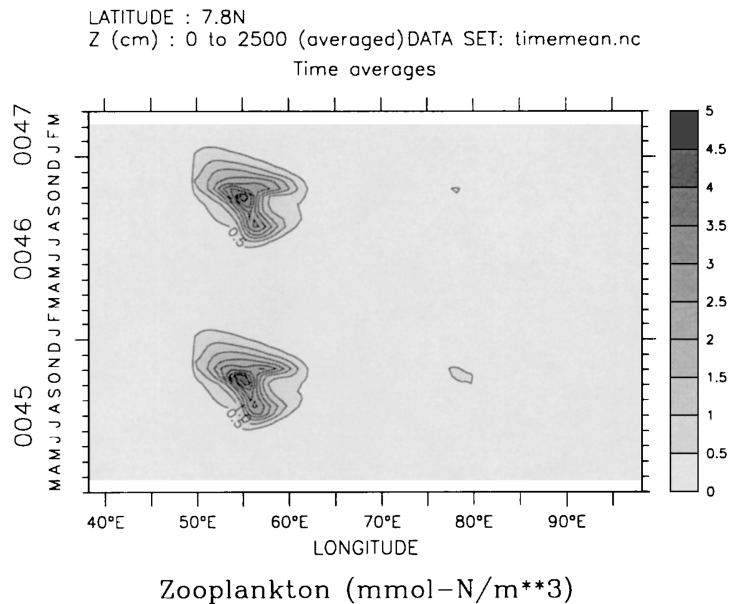
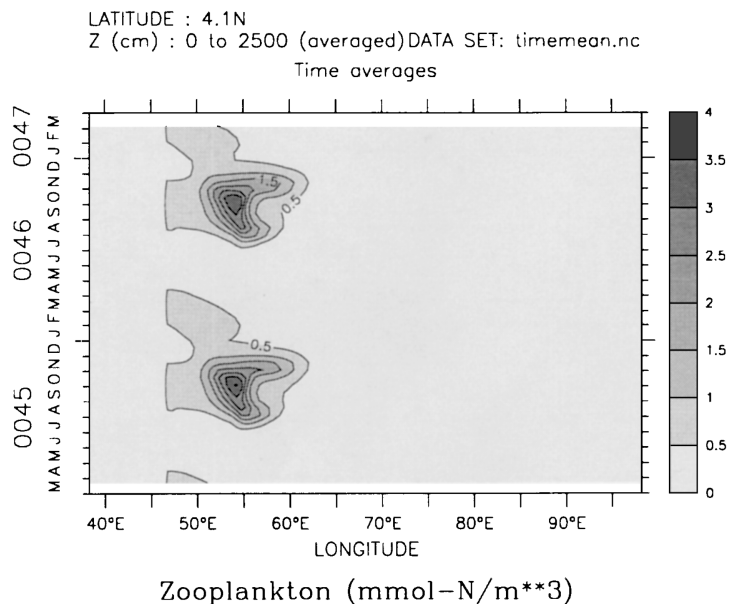


Figure 8. Time series of average zooplankton distribution in the top 25 m at four latitudinal sections.

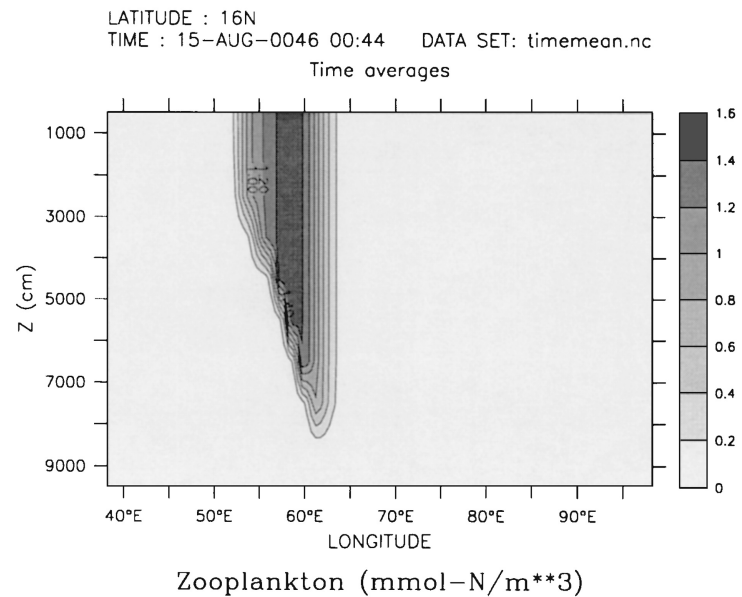
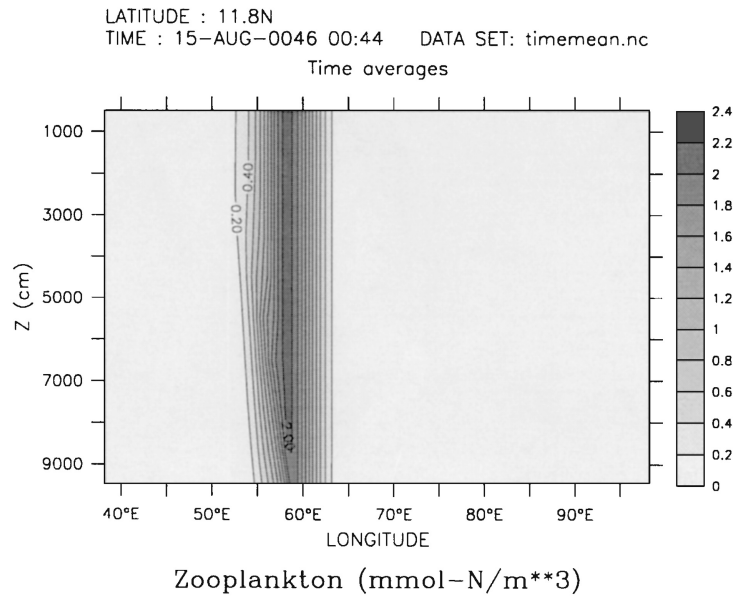
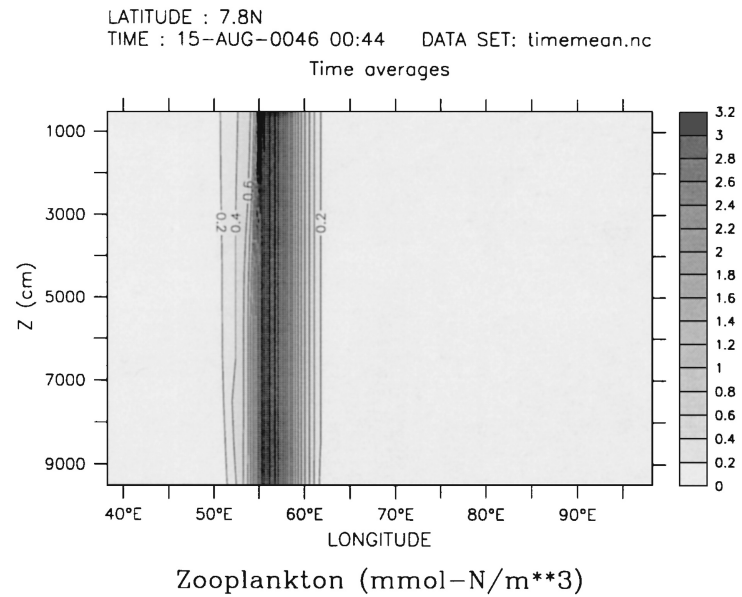
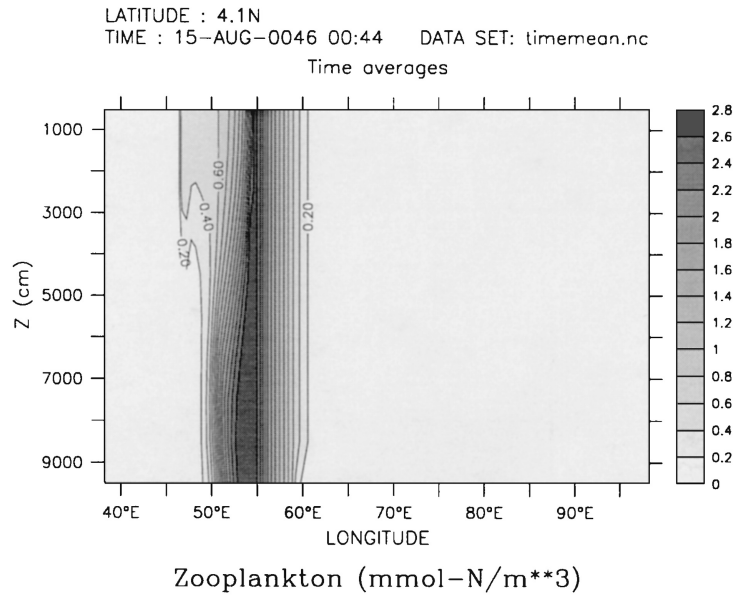
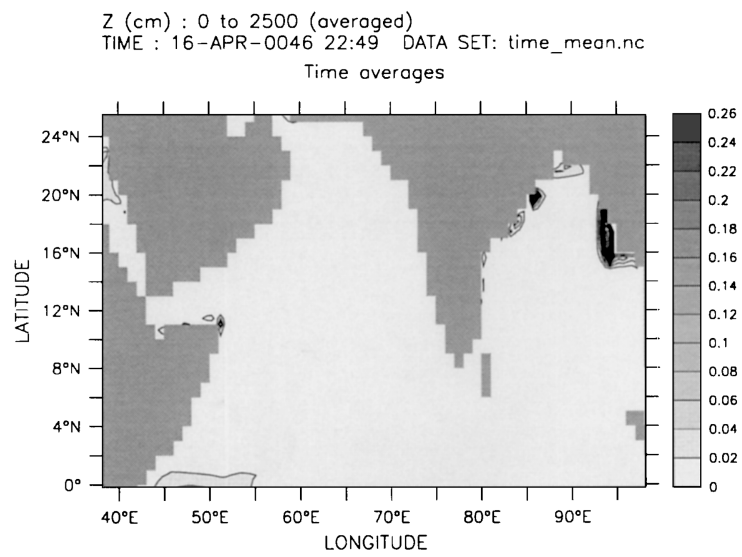
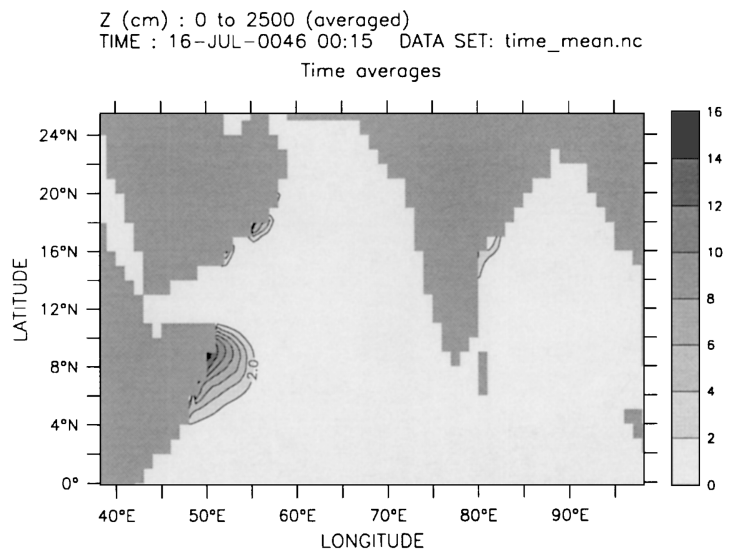


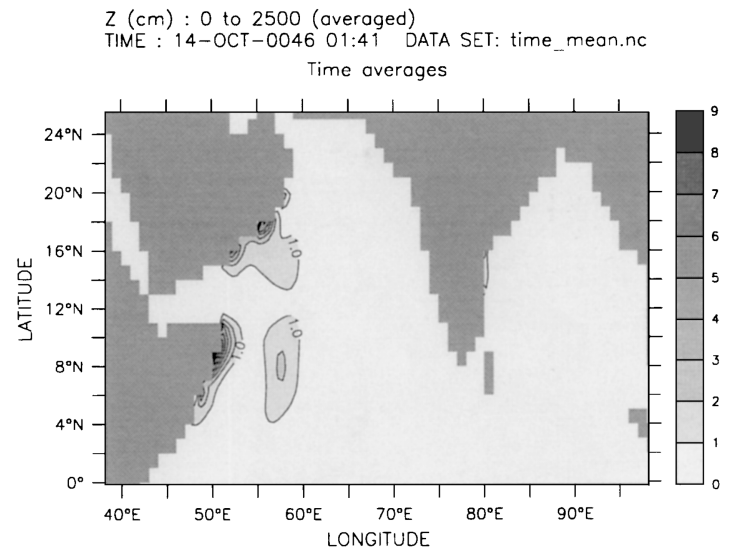
Figure 9. Depth-longitude sections of zooplankton distribution in August at four latitudinal sections.



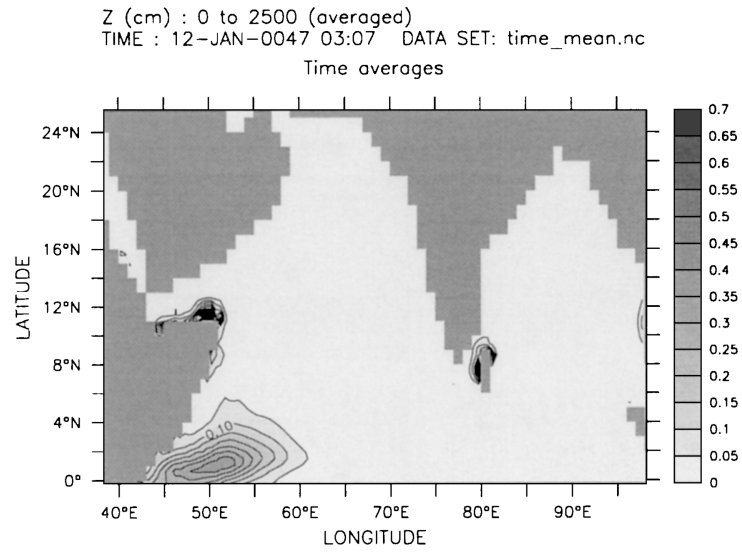
Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)

Figure 10. Time averages of nitrate distribution averaged in the top 25 m. Note the change of scale in different frames.

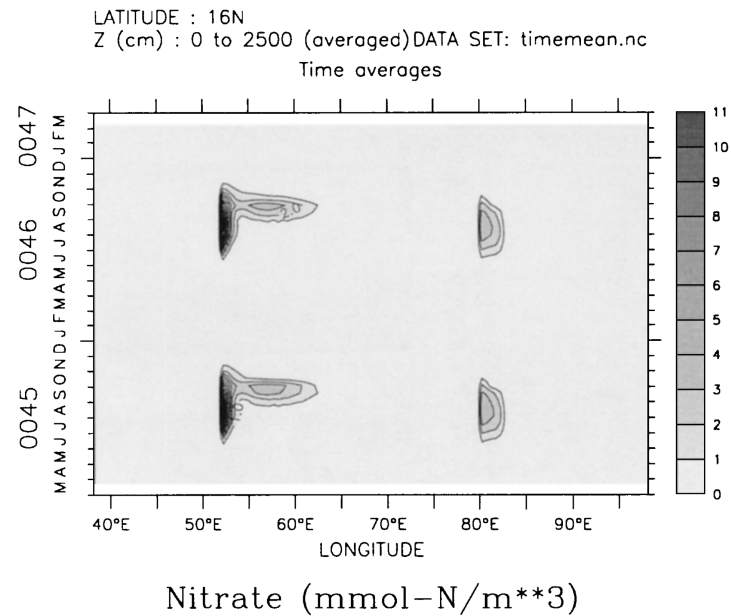
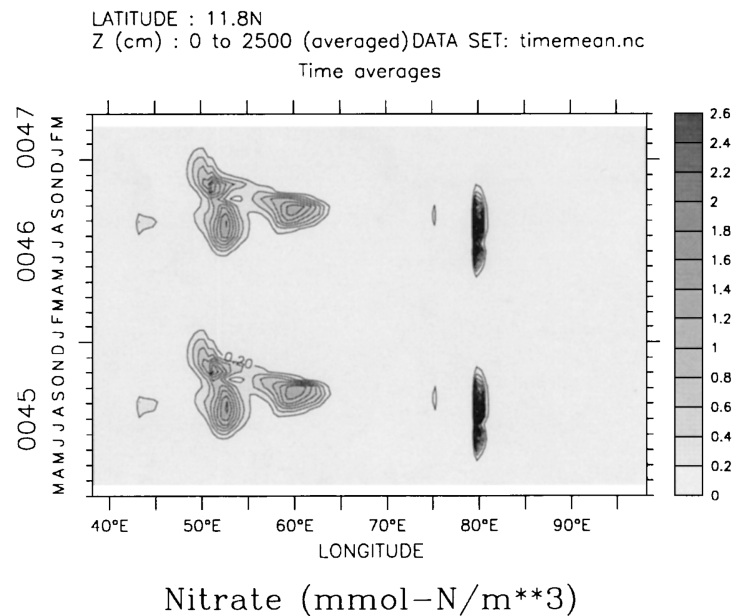
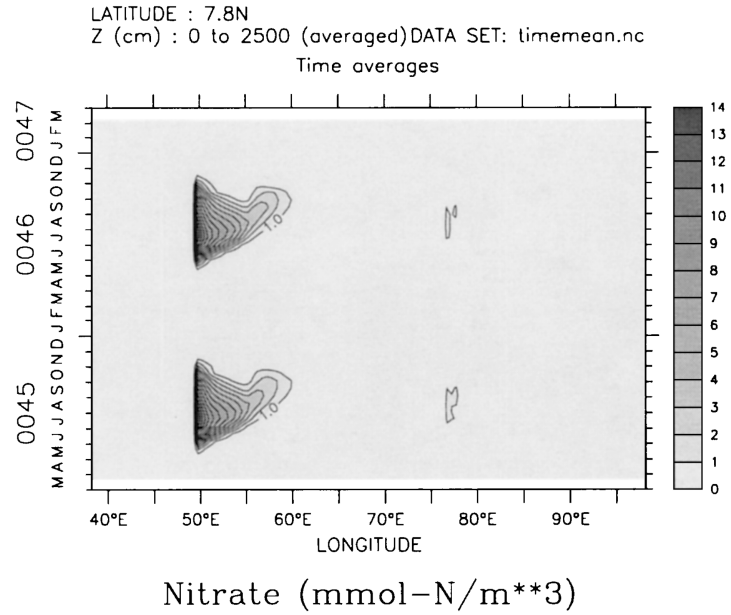
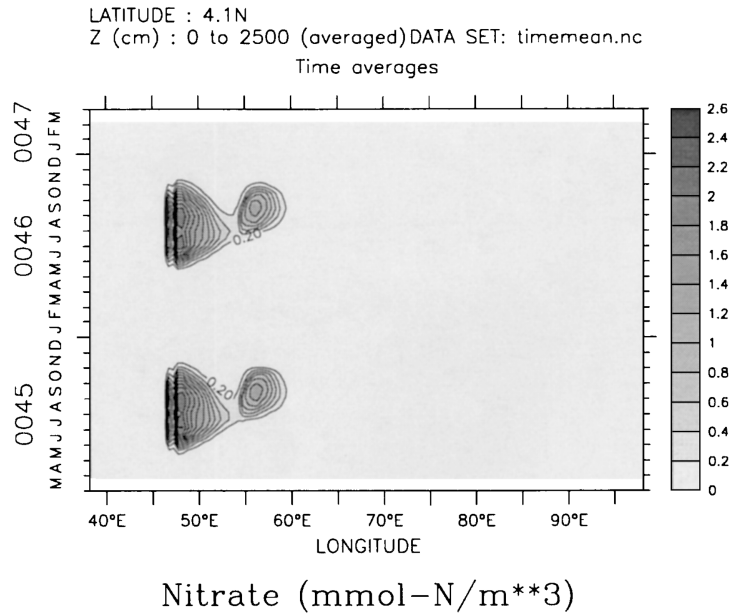
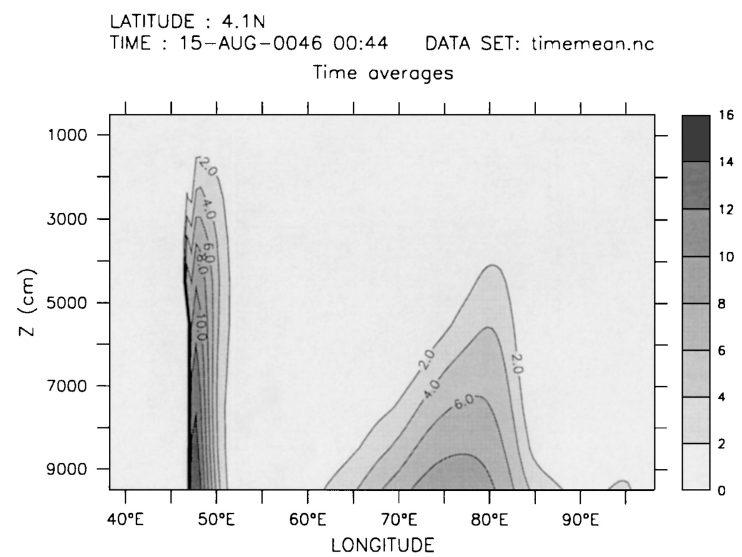
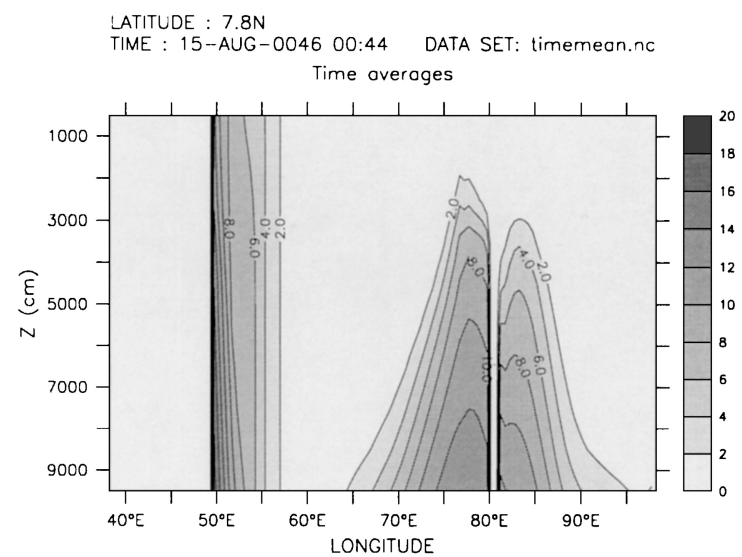


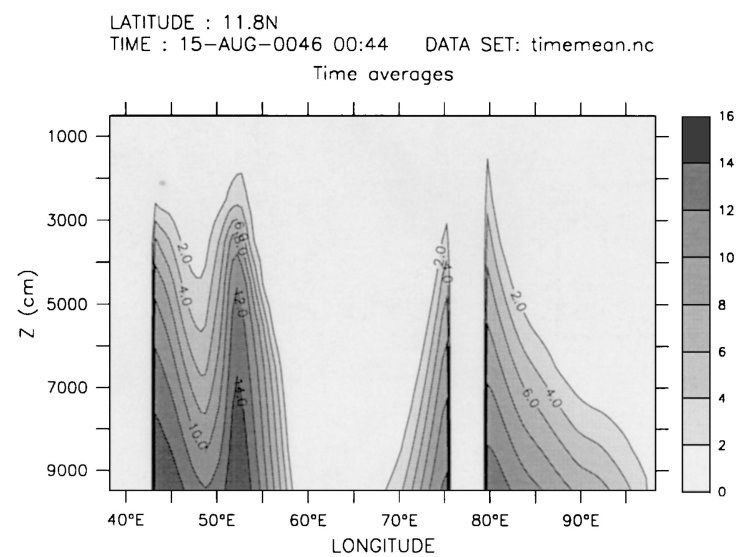
Figure 11. Time series of average nitrate distribution in the top 25 m at four latitudinal sections.



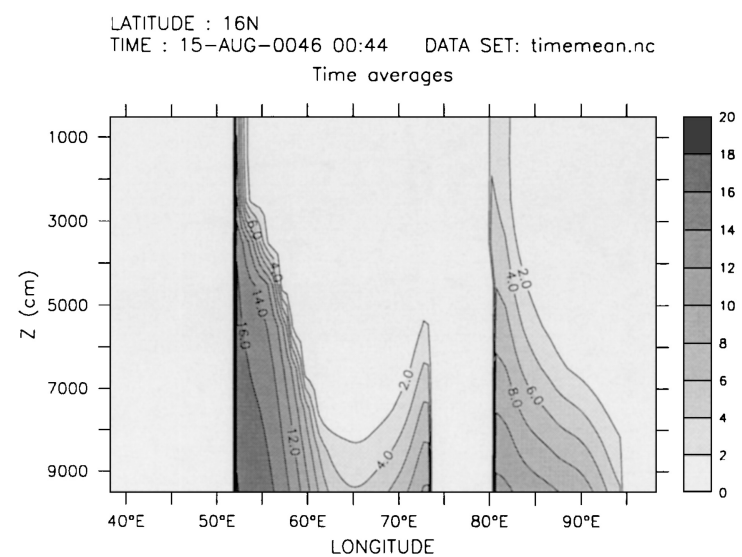
Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)



Nitrate (mmol-N/m**3)

Figure 12. Depth-longitude sections of nitrate distribution in August at four latitudinal sections.

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 NOAA/PMEL TMAP
 Nov 05 00 16:52:20

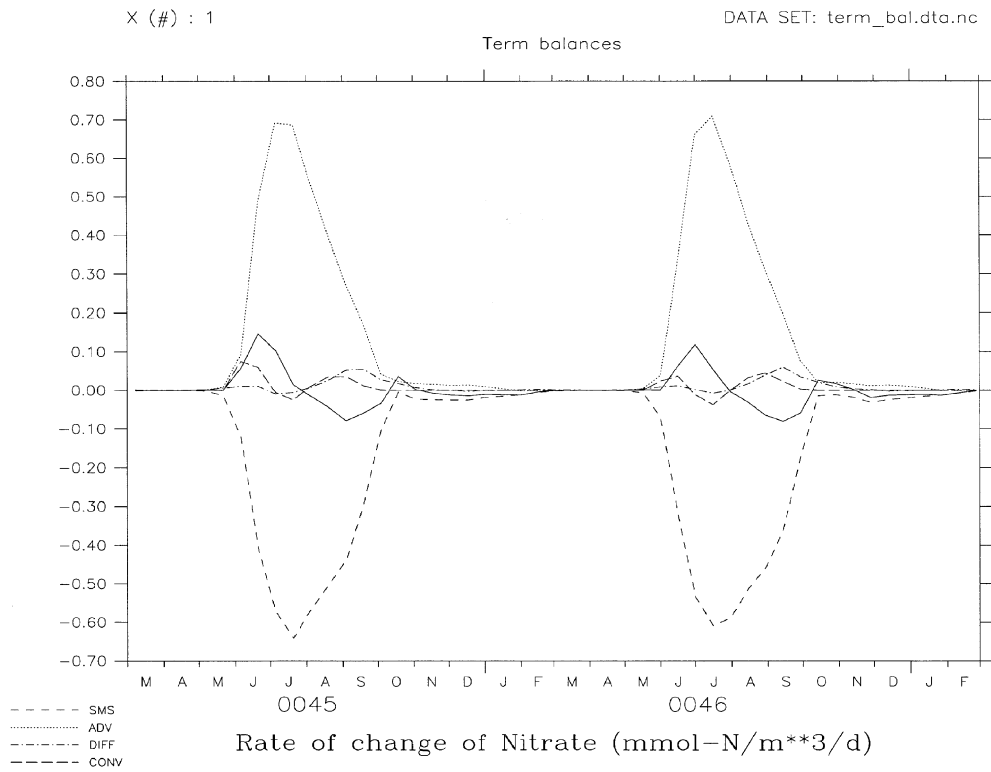


Figure 13. Term balances of equation (4) for nitrate averaged over the region (52E-56E, 5N-9N and 0-50m depth). The rate of change $\frac{\partial N_t}{\partial t}$ is shown by the solid black line while the other terms are indicated in the legend. Note that CONV stands for convective mixing due to gravitational instability which is not indicated explicitly in equation (4). The last two years of integration (years 45 and 46) are shown to get an idea of the annual cycle behaviour of the coupled model.

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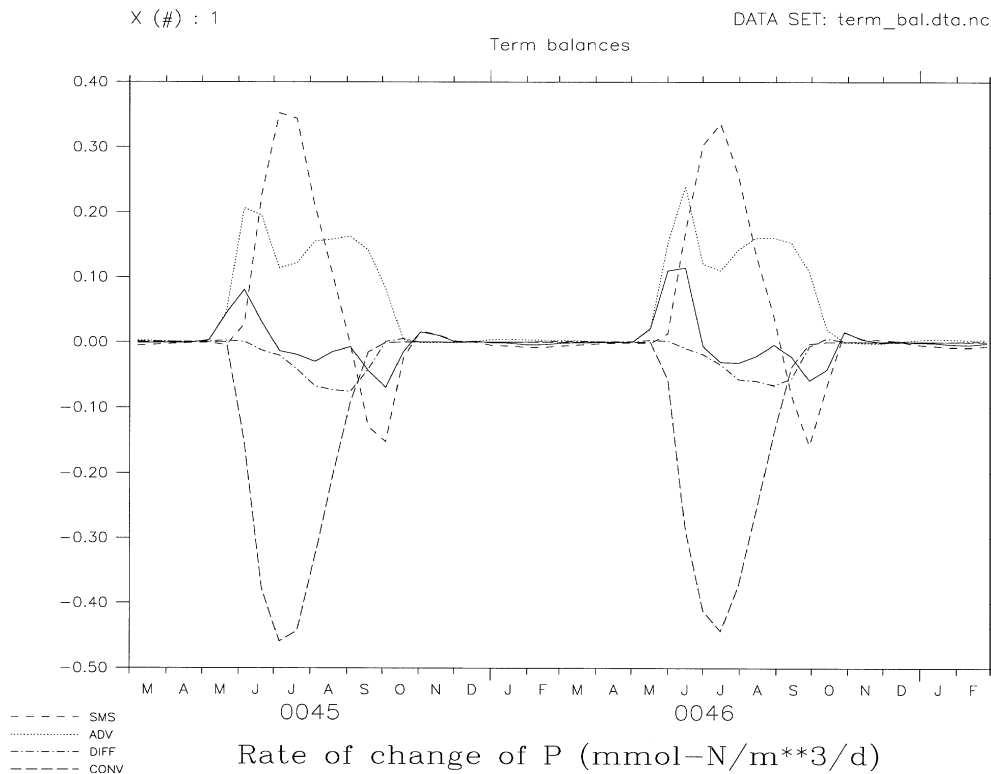


Figure 14. Term balances similar to figure 13 for phytoplankton. The rate of change $\frac{\partial P}{\partial t}$ is shown by the solid black line.

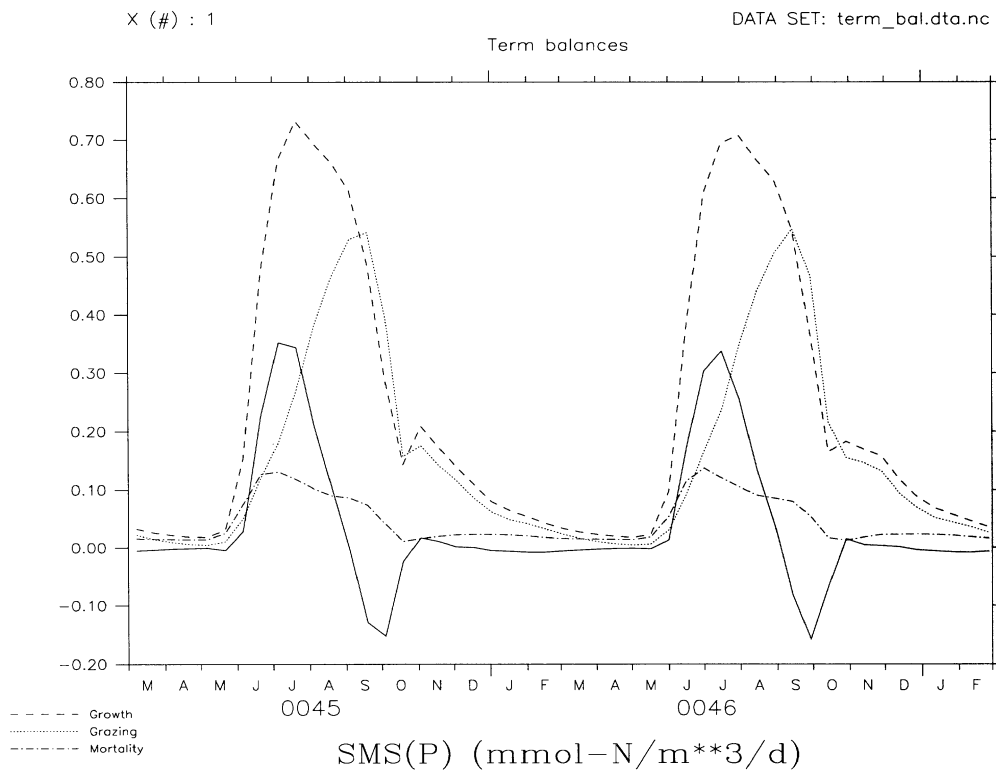


Figure 15. The averaged source term (SMS(P)) and its components in equation (5) for the region (52E–56E, 5N–9N, 0–50m depth). The growth term is the first term on the right hand side of equation (5), the grazing term is G_1 and the mortality term is $\mu_1 P$.

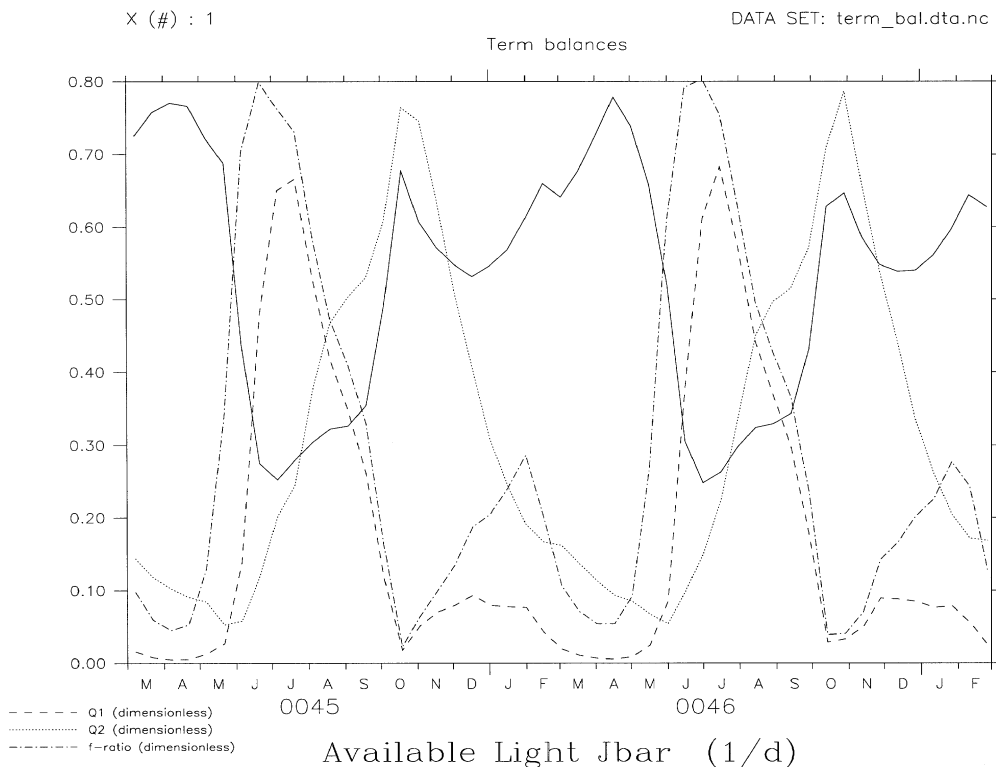


Figure 16. The components of phytoplankton growth term corresponding to figures 14 and 15. The available light term (shown in solid black) is the term in equation (14) averaged over the region indicated in figure 15. The Q_1 and Q_2 terms are given by equations (12) and (13), respectively. F-ratio is $Q_1/(Q_1 + Q_2)$.

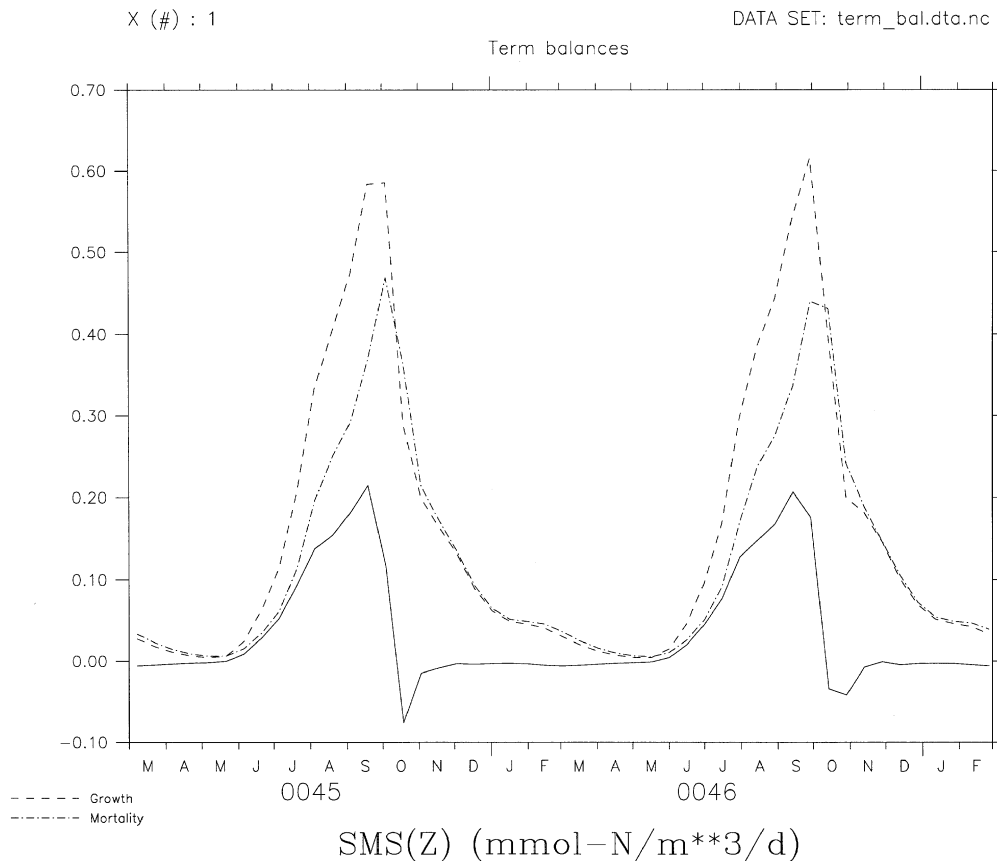


Figure 17. The averaged source term (SMS(Z)) and its components in equation (6) for the region (52E–56E, 5N–9N, 0–50m depth). The growth term is the first term on the right hand side of equation (6), and the mortality term is $(\mu_2 + \mu_5)Z$.

P. There are two distinct blooms during SW monsoon, in June and August, after which zooplankton grazing tends to dominate. Of the advective terms it is the zonal advection which is predominant. The vertical diffusion term dominates in mixing and the negative rates in monsoon are due to shear-induced mixing. The convection term is also large and negative. The light penetration, and the f-ratio are shown in figure 20. Notice the reduction in light penetration during the bloom due to self shading effects. As seen earlier, new production dominates in early monsoon and regenerated production after August. The XBT results for nitrate are in figures 21 and 22. The physical processes transporting nitrates are quite similar to what was observed with *P*, but the convective mixing term has a different sign as deeper waters contain more nitrates unlike *P*.

The DIC values at four typical months is shown in figure 23. The increase in the month of July near the Somali Coast is due mainly to upwelling. It appears that physics which brings up DIC from the deeper ocean wins over biology which fixes DIC. The higher values of DIC in the western Arabian Sea persist in

October but are removed by January. The time series of DIC at latitudinal sections are shown in figure 24. The increased concentration of DIC corresponding to the SW monsoon months can be clearly seen. The depth-longitude sections at the same latitudes (figure 25) in August reveals a structure similar to the one seen for nitrate in figure 11. The term balances in the region 52–56E, 5–9N and 0–50m are shown in figure 26. The effect of photosynthesis is clearly seen in the June and July months. However this negative influence is offset by large zonal advection. The convective contribution is also quite significant with large positive rates as deeper waters contain more DIC. The combined influence of all the terms shows a rapid growth of DIC in June followed by a rapid decline reaching negative growth rates in July and August. The depletion of *P* (by grazing by *Z*) allows DIC to grow again in September and October. The distribution of TALK in the upper waters is shown in figure 27. The higher values in the Arabian Sea are due to higher salinities compared with the waters of Bay of Bengal. The range of variation, however, is much smaller than DIC.

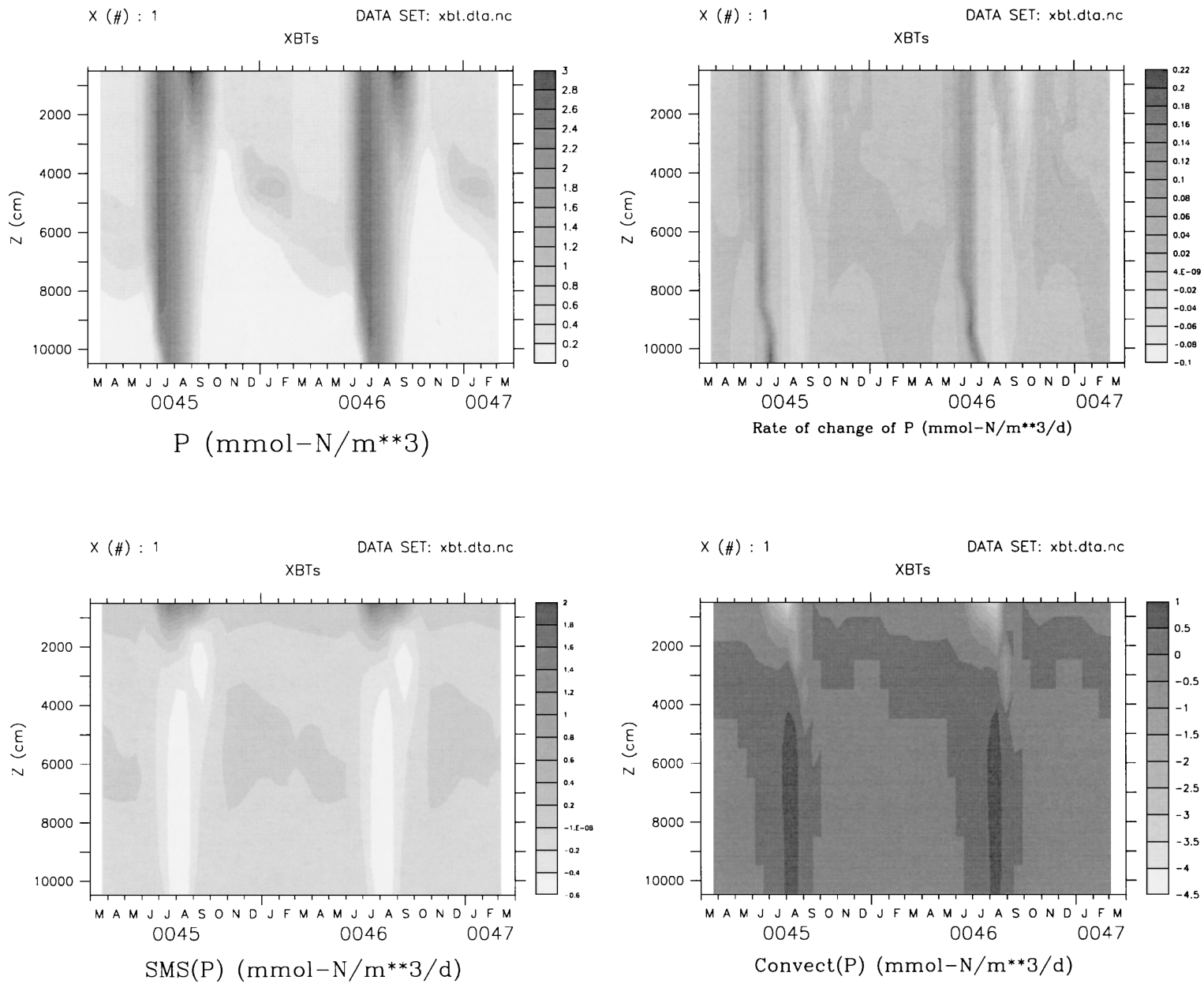


Figure 18. Simulated XBT depth profile at 55.3E, 7.4N of term balances in equation (4) for phytoplankton. Clockwise from top left: P distribution, $\partial P/\partial t$, SMS and convection terms.

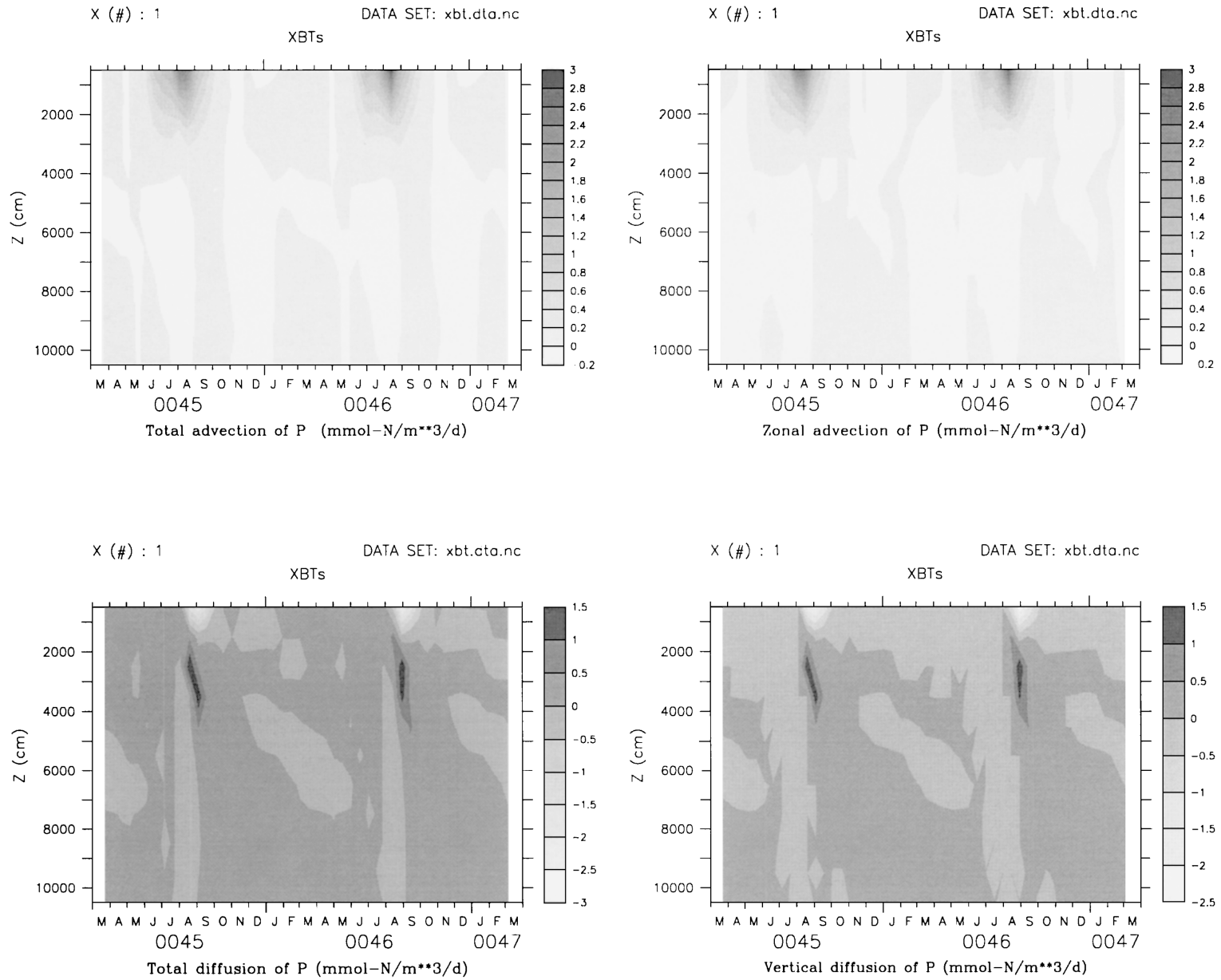


Figure 19. Advective and diffusive components of the XBT profile in figure 18. Notice that the advective component is dominated by zonal advection and the diffusive component is dominated by vertical diffusion.

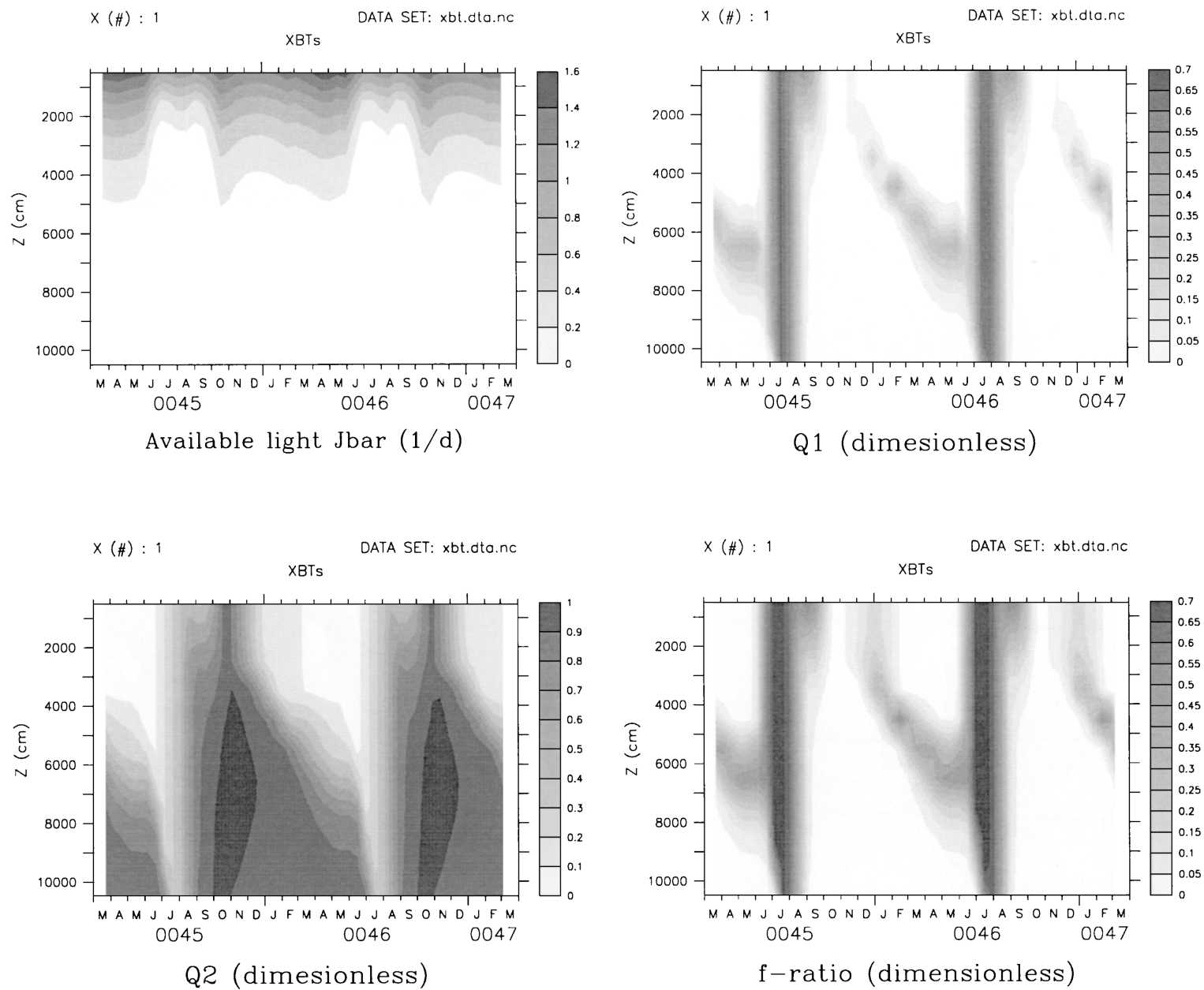


Figure 20. The components of phytoplankton growth similar to figure 16 but for the XBT profile at 55.3E, 7.4N.

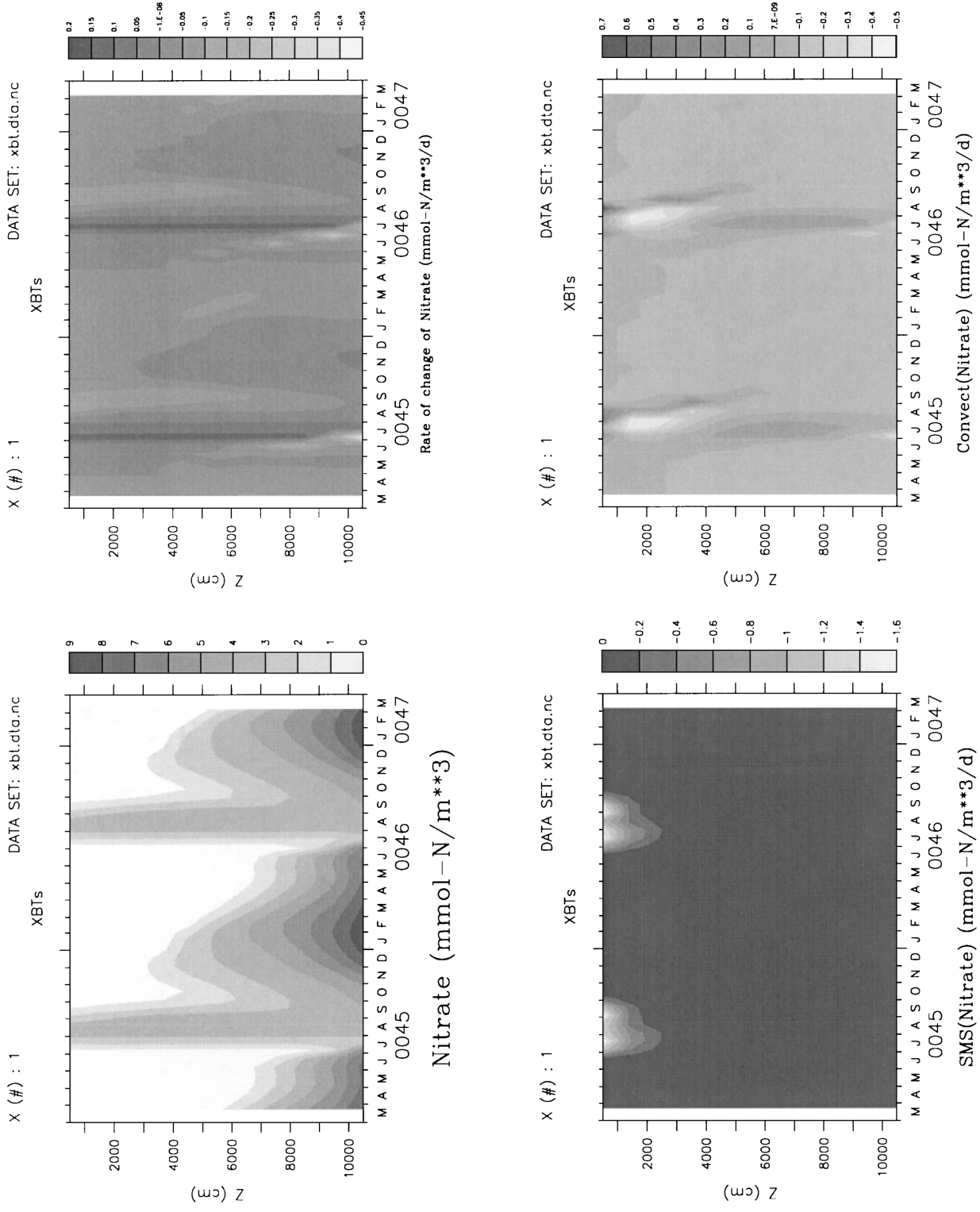


Figure 21. XBT profile for nitrate similar to figure 18.

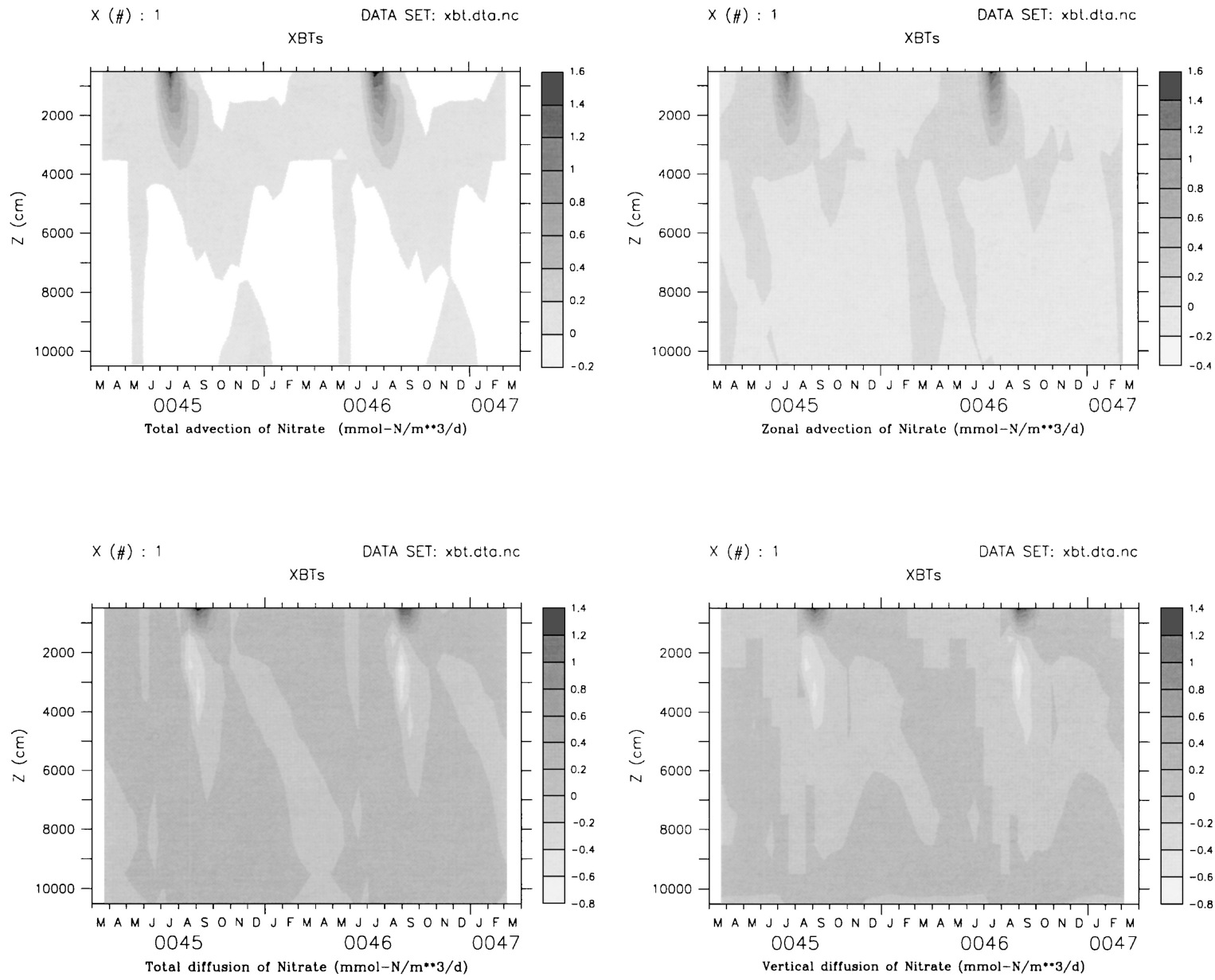


Figure 22. Advective and diffusive components for nitrate similar to figure 19.

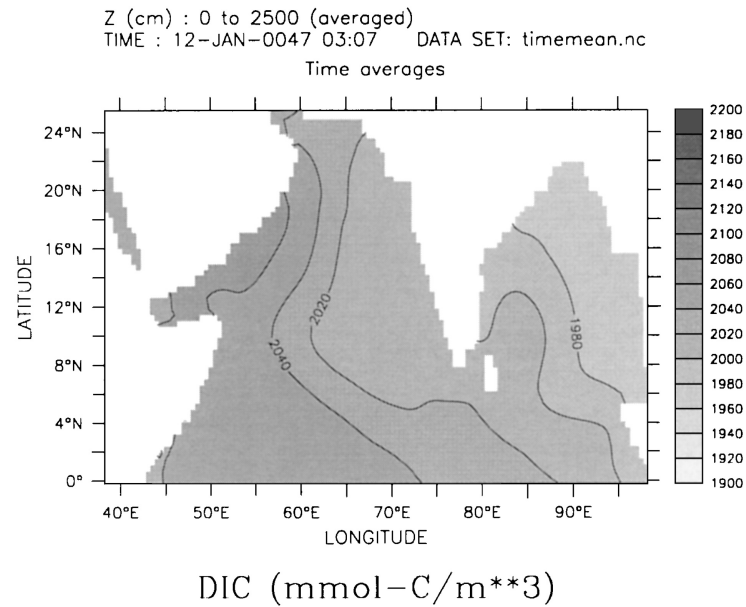
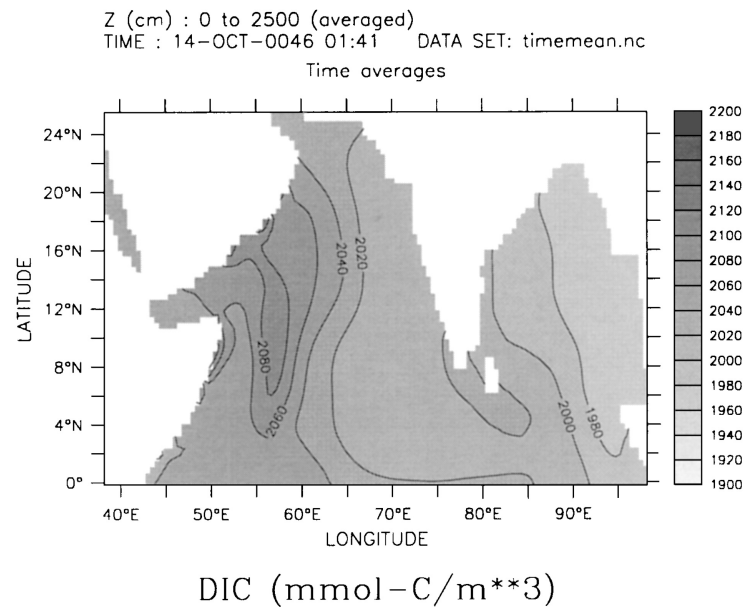
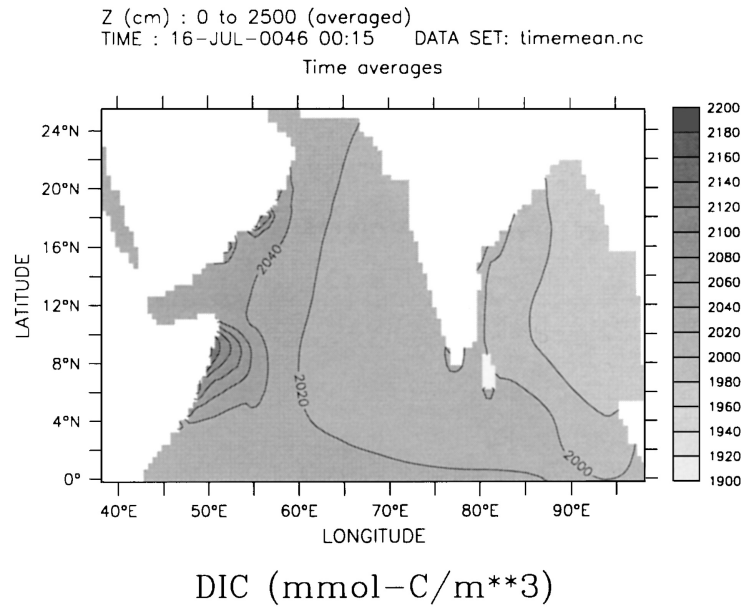
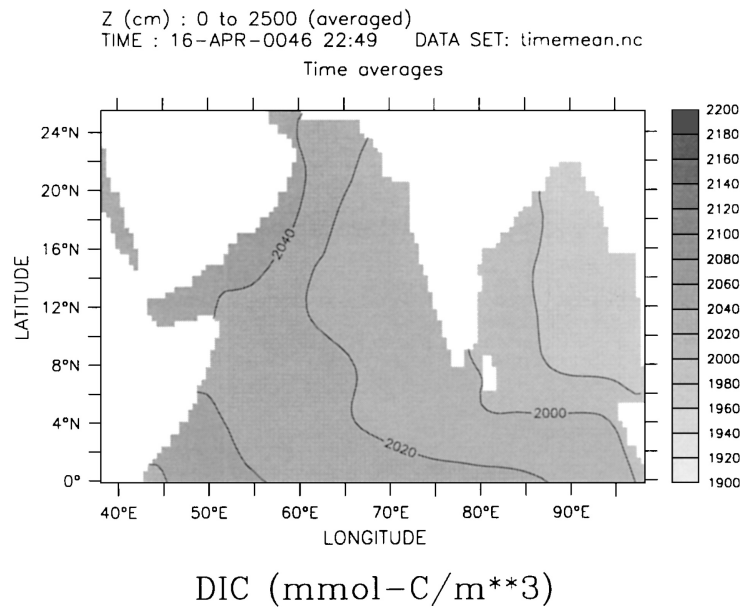
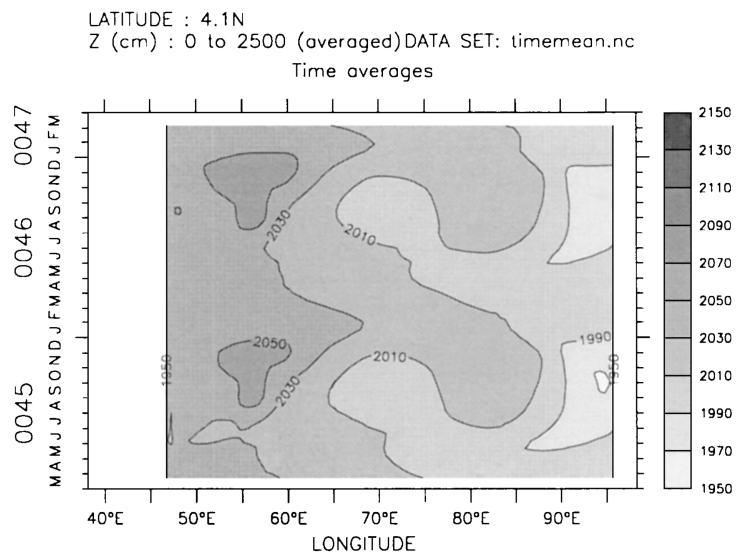
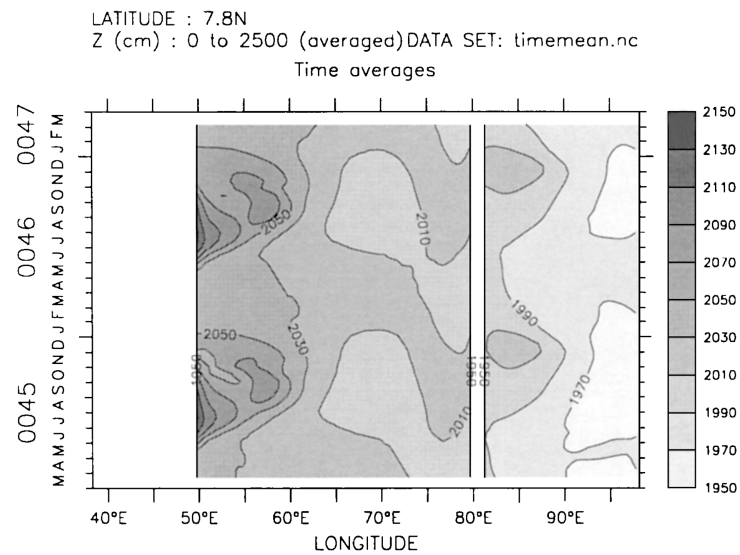


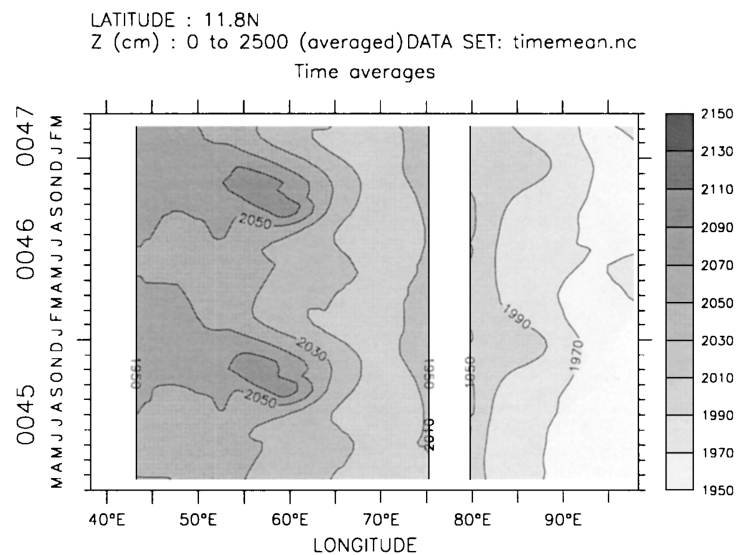
Figure 23. Time averages of DIC distribution averaged in the top 25 m.



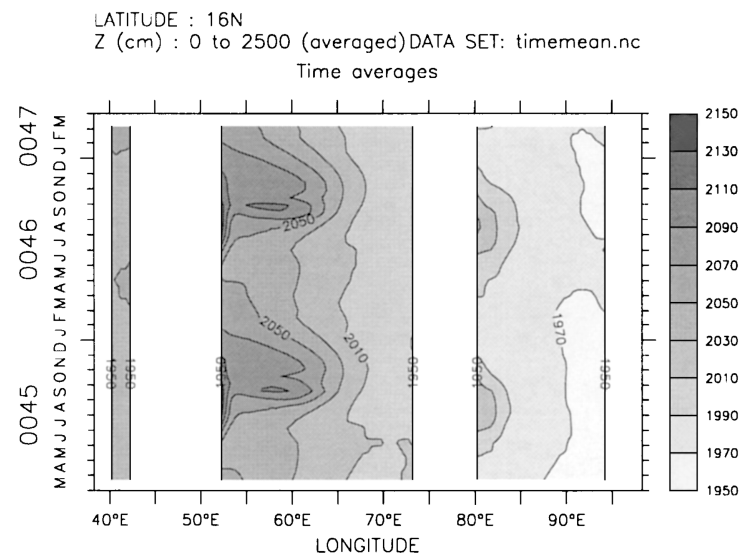
DIC (mmol-C/m**3)



DIC (mmol-C/m**3)



DIC (mmol-C/m**3)



DIC (mmol-C/m**3)

Figure 24. Time series of average DIC distribution in the top 25 m at four latitudinal sections.

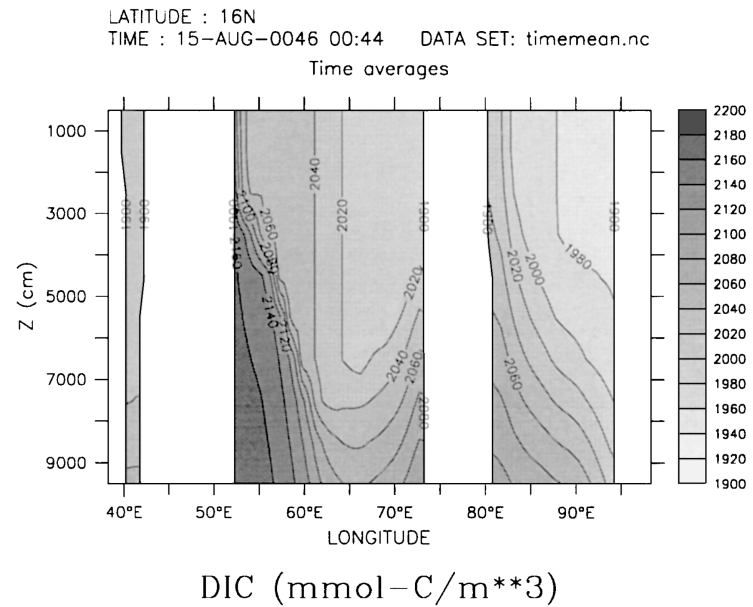
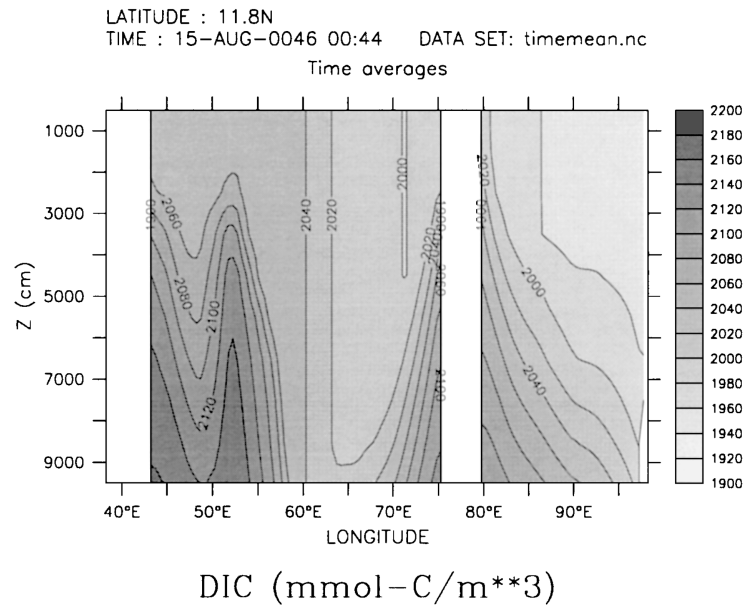
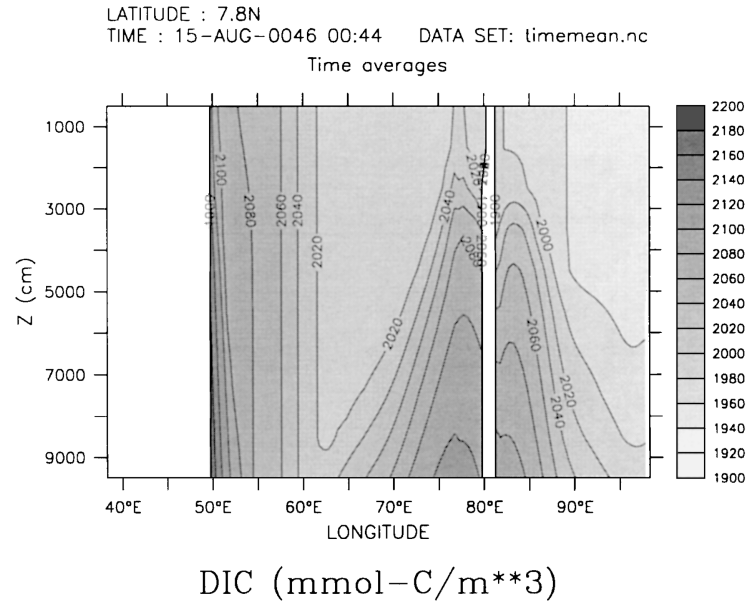
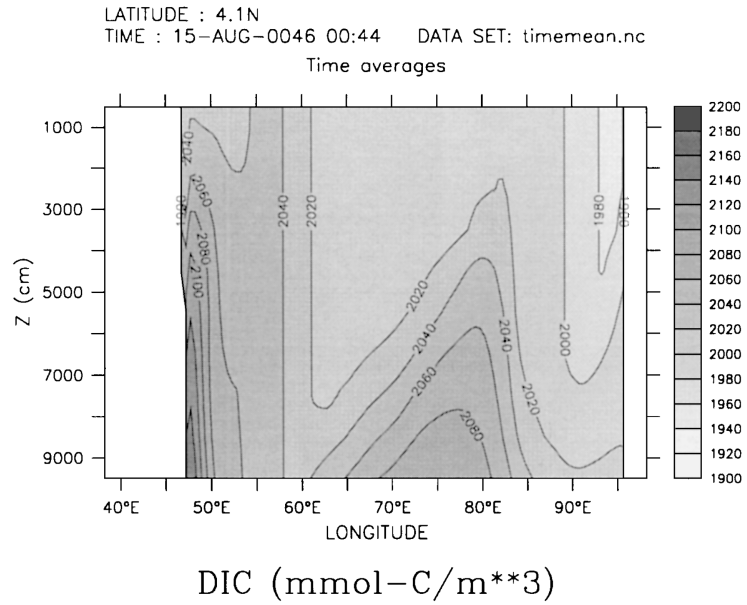


Figure 25. Depth-longitude sections of DIC distribution in August at four latitudinal sections.

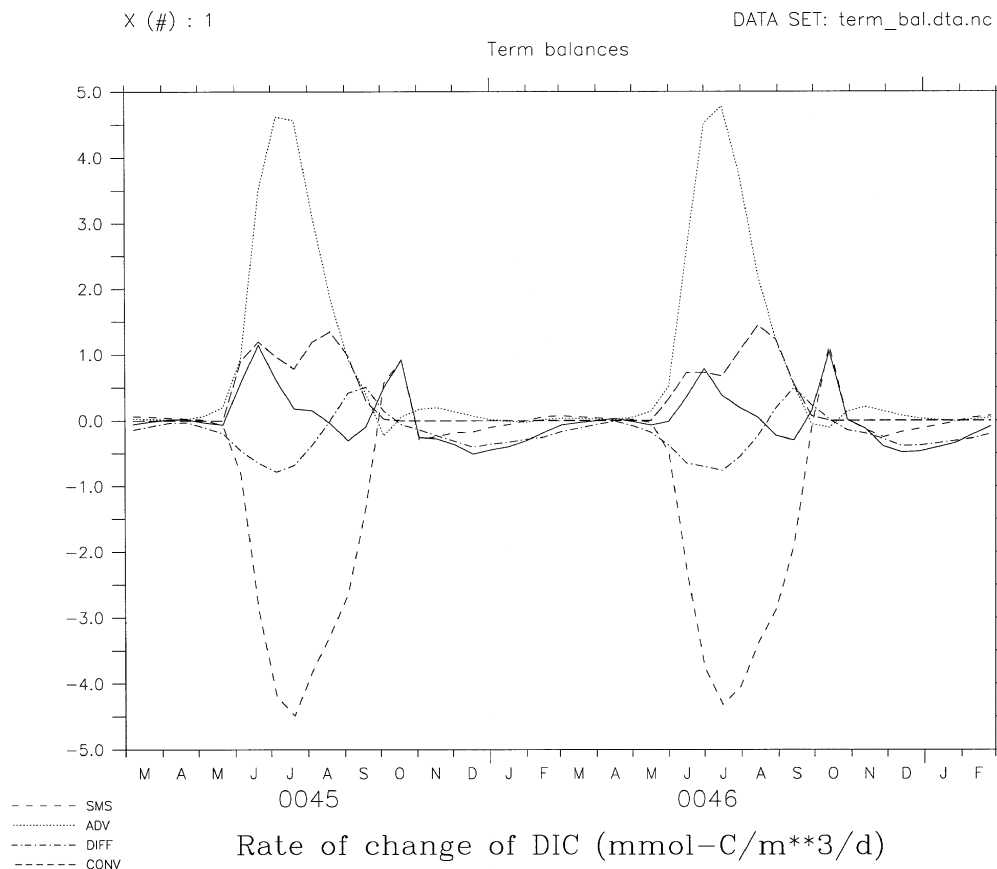


Figure 26. Term balances of equation (4) for DIC averaged over the region (52E-56E, 5N-9N and 0-50m depth). The rate of change $\partial/\partial t$ DIC is shown by the solid black line.

The partial pressure of CO_2 in the surface waters for assumed $\text{pCO}_2^{\text{air}}$ of 345 ppm and 360 ppm are shown in figures 28 and 29, respectively. The two are more or less alike with differences of 20 ppm (the former is lower) in certain regions. There are some regions, especially to the west, near the Arabian Coast where pCO_2 is much higher than observed. Part of the reason for this is the closing off of the Persian Gulf in the model. There are other areas where there are systematic deficiencies but the overall coverage seems to be reasonable. It must be noted that pCO_2 on the sea surface is greater than the atmospheric value everywhere in the Indian Ocean. The Arabian Sea in general has higher values of pCO_2 than the Bay of Bengal due to physics as well as high salinity. The air-sea flux of CO_2 for atmospheric pCO_2 of 345 ppm (left panels) and 360 ppm (right panels) during the two monsoons are shown in figure 30. The other months are not shown because the air-sea flux is less than $-5(\text{mmol-C/m}^2/\text{d})$ in contrast to the high rates seen in the monsoons. A negative value of the flux indicates that the flux is from the ocean and into the atmosphere. The differences between the fluxes for

345 and 360 ppm atmospheric pCO_2 does not seem significant. (Also see figures 28 and 29 where the pCO_2 in the surface waters is more than 420 ppm in the Arabian Sea even during SW monsoon). The maximum outgassing flux is nearly $80 \text{ mmol-C/m}^2/\text{day}$ which occurs near the Somali tip.

4. Conclusions

For the first time, a large scale basin model of the Indian Ocean incorporating physics, chemistry and biology has been developed to study the cycling of carbon in the Indian Ocean. The effect of different terms in the transport equations have been illustrated in detail. The model captures the SW monsoon blooms of phytoplankton near the Somali Coast reasonably well. The competition between physics and biology in maintaining the surface pCO_2 concentration has been modelled reasonably well. The results show that in spite of the large primary production in the Arabian Sea, physics still prevails by pumping in sufficient dissolved inorganic carbon

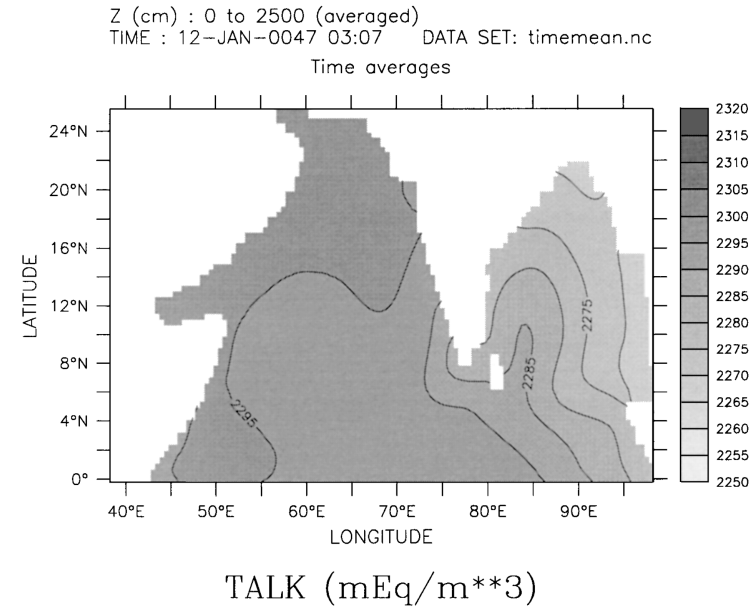
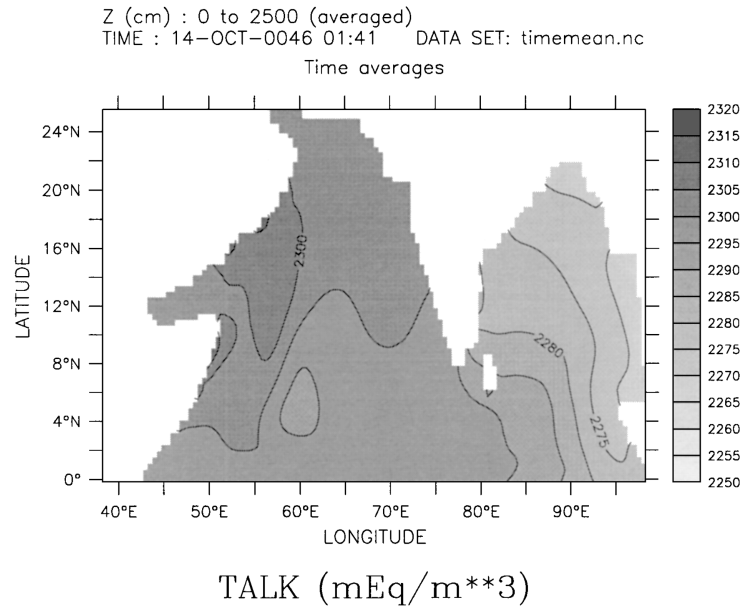
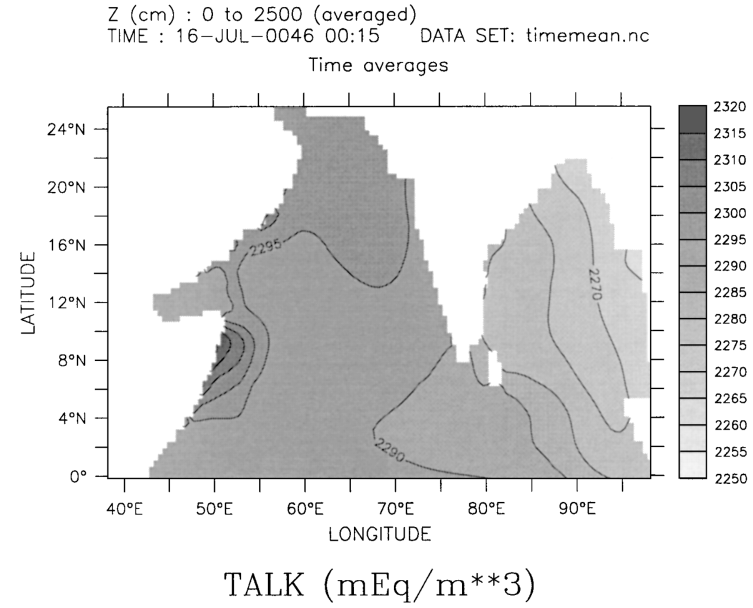
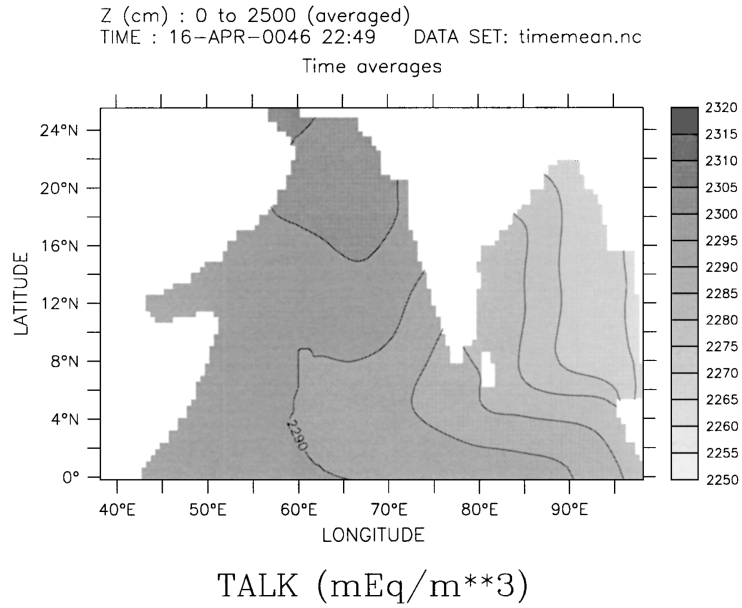


Figure 27. Time averages of TALK distribution averaged in top 25 m.

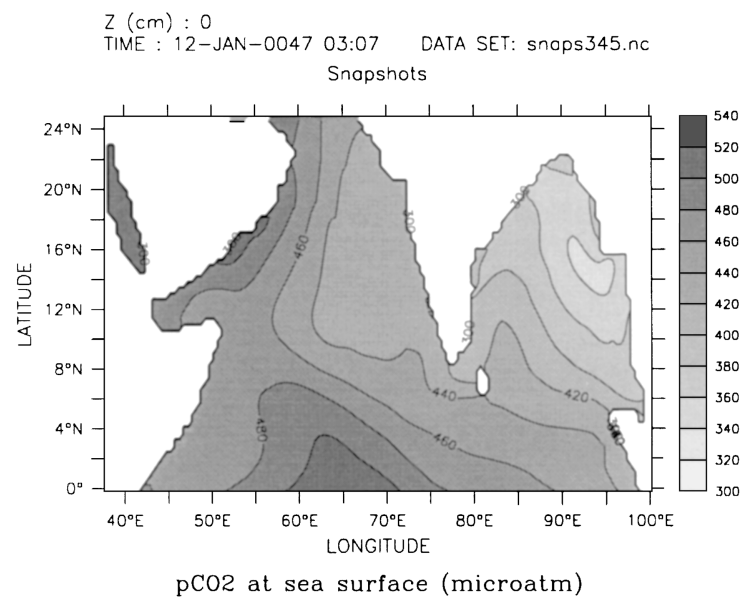
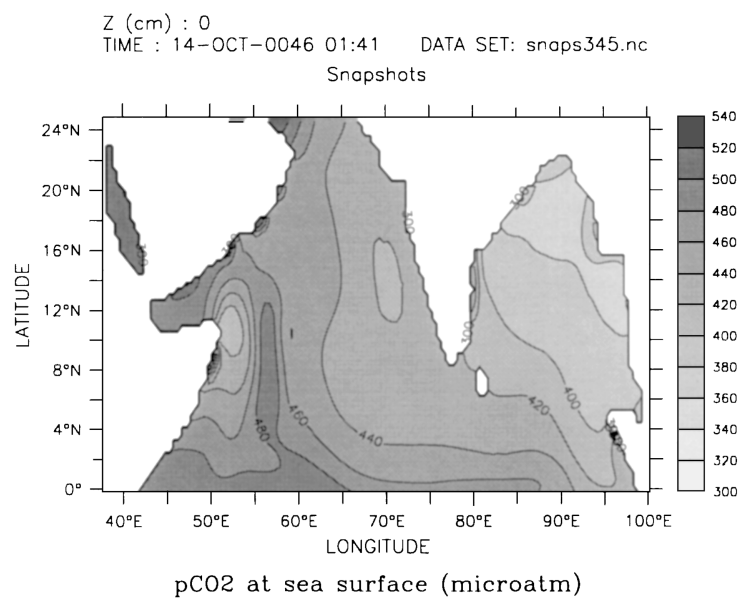
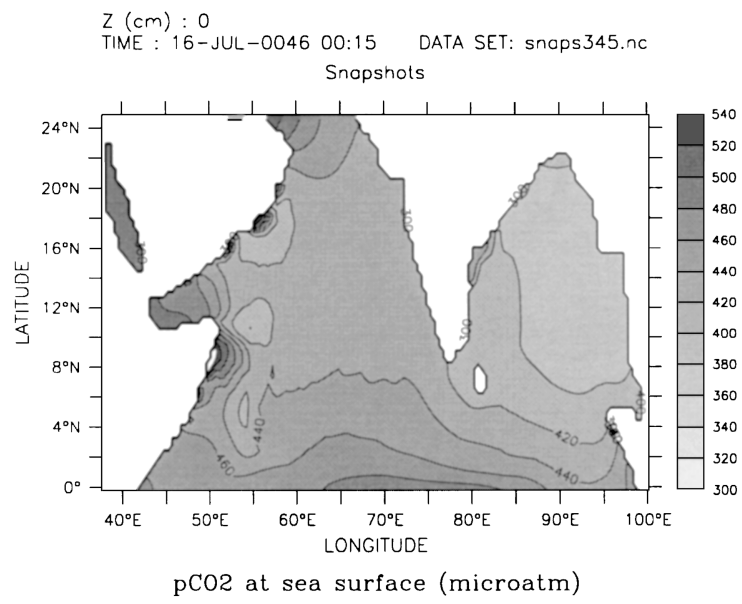
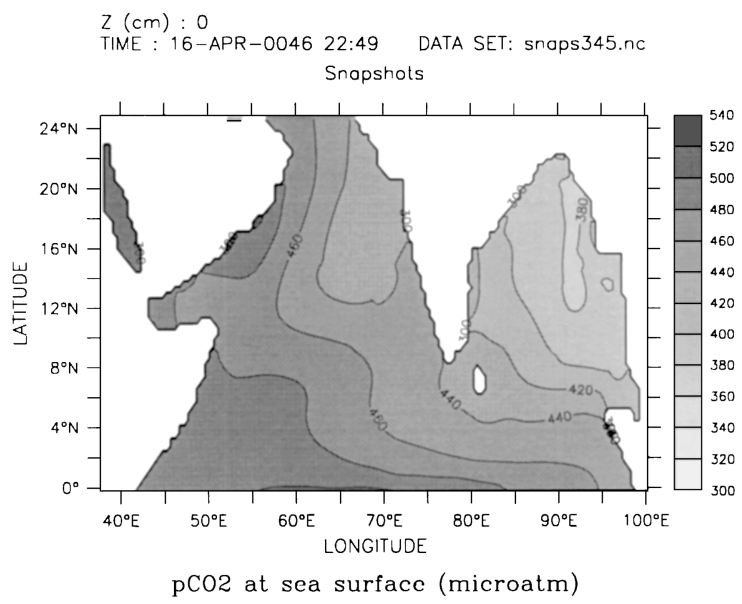


Figure 28. Partial pressure of CO₂ in the surface waters for assumed pCO₂^{air} of 345 ppm.

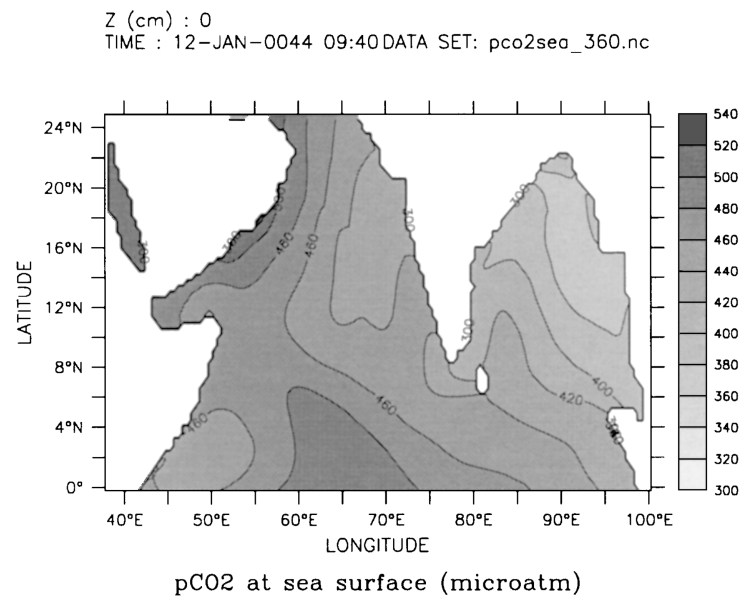
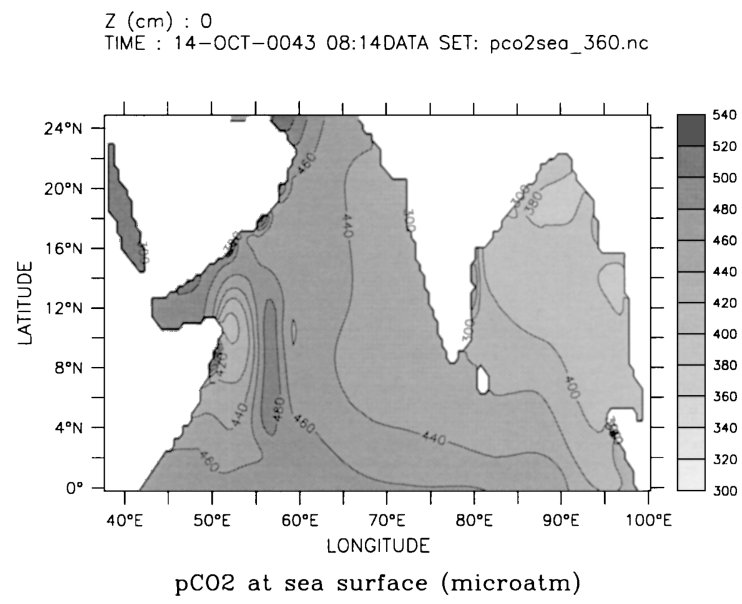
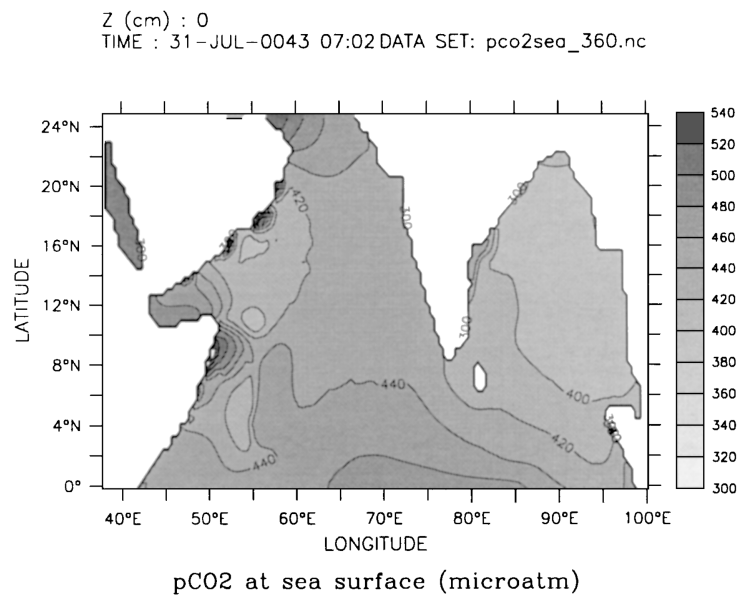
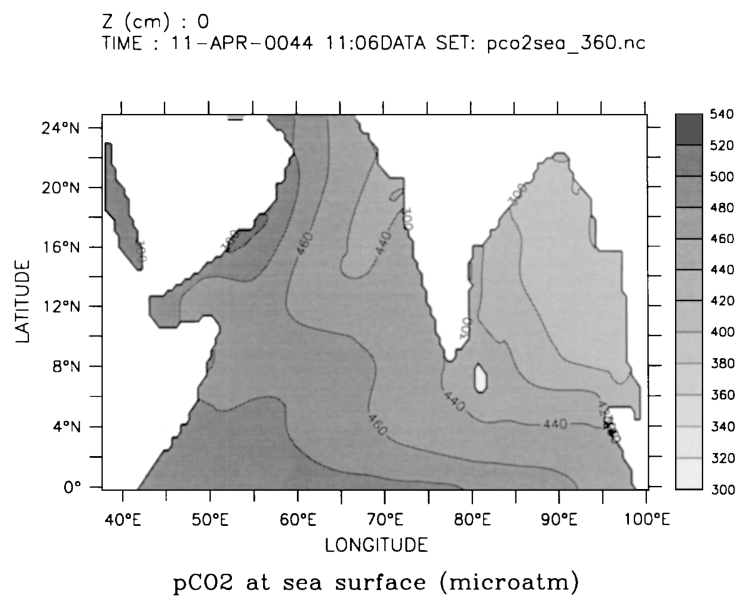
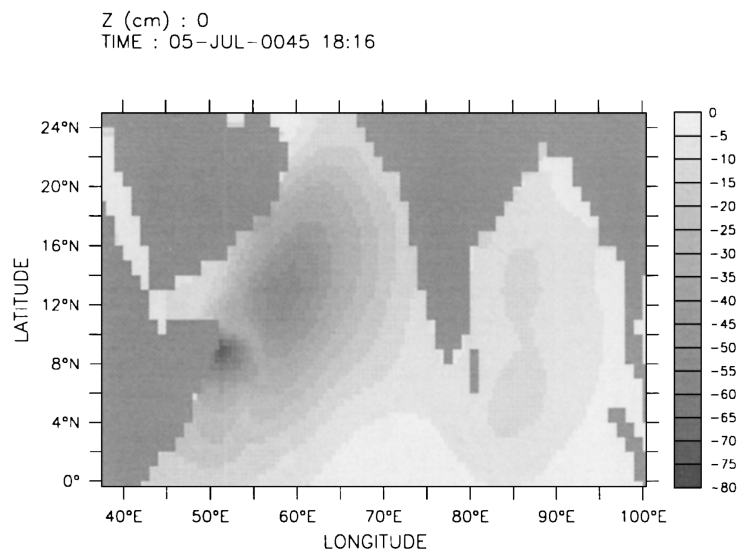
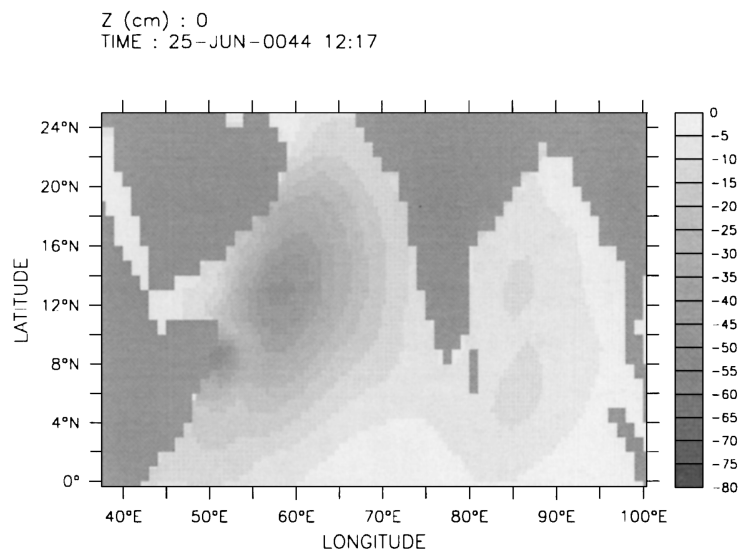


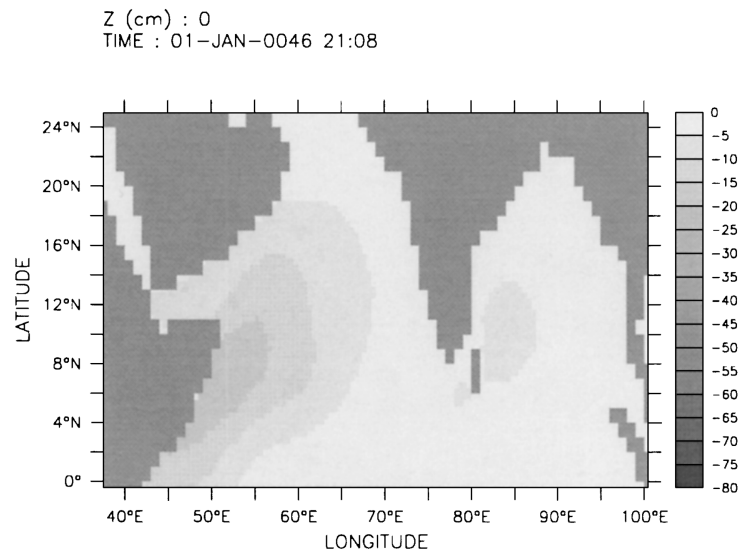
Figure 29. Partial pressure of CO₂ in the surface waters for assumed pCO₂^{air} of 360 ppm.



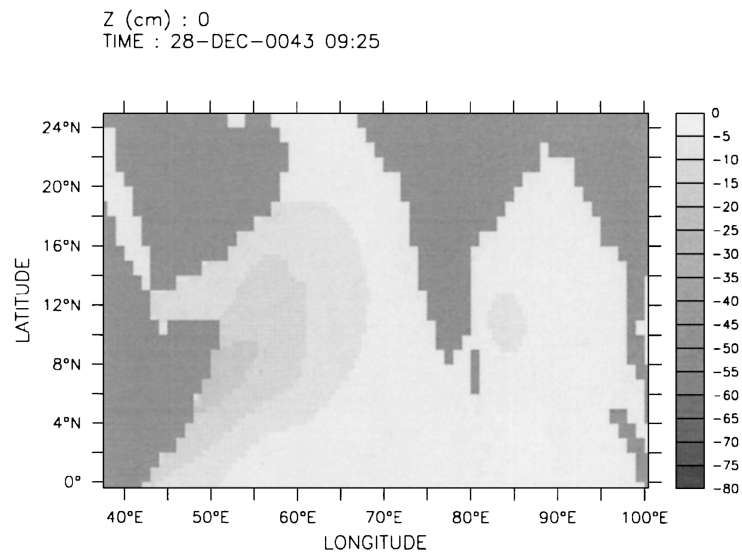
C Flux (mmol/m**2/d)



C Flux (mmol/m**2/d)



C Flux (mmol/m**2/d)



C Flux (mmol/m**2/d)

Figure 30. Air-sea flux of CO₂ (negative values indicate outgassing from the sea). The left panels are for pCO₂^{air} of 345 ppm while the right panels are for 360 ppm.

from the deeper waters to make the entire Indian Ocean an outgassing region.

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