The Thermal Width of the Glueball at Non-Zero Temperature

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We use SU(3) anisotropic lattice QCD at quenched level to study the $0^{++}$ glueball correlator at various temperature taking into account the possible existence of the thermal width in the ground-state peak. For this purpose, we adopt the Breit-Wigner ansatz for the appropriate fit-function for the lattice data obtained with 5,500–9,900 gauge configurations at each $T$. The results show the significant thermal width broadening as $\Gamma(T_c) \approx 300$ MeV with a reduction in the peak center as $\Delta\omega_0(T_c) \approx 100$ MeV near the critical temperature $T_c$.

1. INTRODUCTION

At finite temperature/density, the vacuum properties of QCD are expected to change such as the reduction of the string tension, the partial chiral restoration, etc. Such changes of the vacuum properties should be followed by the changes of the hadron properties, since hadrons are composite particles consisting of quarks and gluons. The hadronic pole-mass shifts are thus considered to serve as the important precritical phenomena of the QCD phase transition near the critical temperature $T_c$, and were extensively studied by using various QCD-motivated low-energy effective theories [1–4]. These studies suggested the pole-mass reductions of charmoniums, light $q\bar{q}$ mesons and the glueball near the critical temperature. Recently, the accurate pole-masses measurement at finite temperature with the lattice QCD became possible by means of the anisotropic lattice [5,6]. Quenched-level Monte Carlo calculations showed that the pole-masses of the $q\bar{q}$ mesons are almost unchanged from their zero-temperature values in the confinement phase [6,7], while the pole-mass of the $0^{++}$ glueball shows the 300 MeV reduction near the critical temperature [8]. These tendencies are consistent with the recent lattice studies on the screening mass at finite temperature [9,10]. In these analysis, the bound-state peaks are assumed to be sufficiently narrow. However, each bound-state acquires the thermal width through the interaction with the heat bath, and the thermal width is expected to grow up with temperature, which may leads to a possible collapse of the narrow-peak assumption in some cases. In this paper, we first discuss what is the expected consequence of the thermal width broadening. Then, we propose the Breit-Wigner ansatz for the fit-function for the temporal correlator at finite temperature. We finally show the results of the Breit-Wigner analysis of the $0^{++}$ glueball correlator at finite temperature [11].

2. THE BREIT-WIGNER ANSATZ

We consider the temporal correlator $G(\tau) \equiv Z(\beta)^{-1}\text{Tr}(e^{-\beta H}\phi(\tau)\phi(0))$. Its spectral representation is given as

$$G(\tau) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{\cosh(\omega(\beta/2 - \tau))}{2 \sinh(\beta\omega/2)} \rho(\omega),$$  \hspace{1cm} (1)

where $H$ denotes the QCD Hamiltonian, $Z(\beta) \equiv \text{Tr}(e^{-\beta H})$ the partition function, $\phi(\tau)$ is the zero-momentum projected glueball operator [8] repre-
sented in the imaginary-time Heisenberg picture as $\phi(\tau) \equiv e^{-\beta H}\phi(0)e^{-\tau H}$. Here, $\rho(\omega)$ denotes the spectral function

$$
\rho(\omega) = \sum_{m,n} \frac{|\langle n|\phi|m\rangle|^2}{Z(\beta)} e^{-\beta E_m} \times 2\pi \left( \delta(\omega - \Delta E_{nm}) - \delta(\omega - \Delta E_{mn}) \right),
$$

(2)

where $E_n$ denotes the energy of $n$th excited states, and $\Delta E_{mn} \equiv E_m - E_n$. Note that $\rho(\omega)$ is odd in $\omega$ reflecting the bosonic nature of the glueball. By adopting the appropriate ansatz for $\rho(\omega)$, we can extract various physical quantities such as the pole-mass and the width through the spectral representation Eq. (1).

We first consider the case where the bound-state peak is narrow [6–8]. In this case, by introducing the temperature-dependent pole-mass $m(T)$, $\rho(\omega)$ can be parameterized as

$$
\rho(\omega) \approx 2\pi A \left( \delta(\omega - m(T)) - \delta(\omega + m(T)) \right),
$$

(3)

where $A$ represents the strength. The second delta-function is introduced to respect the odd-function nature of $\rho(\omega)$. Since the corresponding $G(\tau)$ reduces to a single hyperbolic cosine, the pole-mass measurement at finite temperature can be performed in the same way as the standard mass measurement at zero temperature [12].

We next consider the case where the thermal width is wide. In this case, the peak center $\omega_0$ of $\rho(\omega)$ represents the observed “mass” of the thermal hadron. What follows the narrow-peak assumption now? To consider this, we notice that $G(\tau)$ can be thought of as a weighted average of hyperbolic cosines with the weight as

$$
W(\omega) \equiv \frac{\rho(\omega)}{2 \sinh(\beta \omega/2)}.
$$

(4)

Here, $2 \sinh(\beta \omega/2)$ in the denominator works as the biased factor, which enhances the smaller $\omega$ region while suppressing the larger $\omega$ region. Consequently, the pole-mass $m(T)$, which is approximated with the peak position of $W(\omega)$, is smaller than the peak center $\omega_0$ of the spectral function $\rho(\omega)$, i.e., the observed hadron “mass” [11].

What is the appropriate functional form of the fit-function? To find this, we consider the retarded Green function $G_R(\omega)$. At $T = 0$, bound-state poles of $G_R(\omega)$ are located on the real $\omega$-axis. At $T > 0$, bound-state poles are moving into the complex $\omega$-plane with increasing temperature. Suppose that a bound-state pole is located at $\omega = \omega_0 - i\Gamma$ as

$$
G_R(\omega) = \frac{A}{\omega - \omega_0 + i\Gamma} + \cdots,
$$

(5)

where $A$ represents the residue at the pole, and “…” the non-singular terms around the pole. Since the spectral function is the imaginary part of the retarded Green function, the contribution of this complex pole is expressed in the form of Lorentzian as

$$
\rho(\omega) = -2\text{Im} \left( G_R(\omega) \right).
$$

(6)

Here, $\delta(\omega - m(T)) - \delta(\omega + m(T))$ is a smeared delta-function with the width $\epsilon > 0$. The second term in Eq. (6) is introduced to respect the odd-function nature of $\rho(\omega)$. In the limit $\Gamma \to 0$, Eq. (6) reduces to Eq. (3). We thus see that the appropriate fit-function, which takes into account the effect of the non-zero thermal width, is the following Breit-Wigner type as

$$
g(\tau) \equiv \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{\cosh(\omega(\beta/2 - \tau))}{2 \sinh(\beta \omega/2)} \times 2\pi \tilde{A} \left( \delta(\omega - \omega_0) - \delta(\omega + \omega_0) \right),
$$

(7)

where $\tilde{A}$, $\Gamma$ and $\omega_0$ are understood as fit-parameters, corresponding to the residue, the thermal width and the peak center, respectively. Note that $g(\tau)$ is a generalization of the ordinary single hyperbolic cosine fit-function. In order to use Eqs. (3) and (6), it is essential to suppress the higher spectral contributions. This is usually achieved by appropriate choices of the fit-ranges and also by improving the glueball operator, for instance, with the smearing method [11].

3. NUMERICAL RESULT

We use the SU(3) anisotropic lattice plaquette action as $S_G = \frac{\beta_{lat}}{N_c \gamma_G} \sum_{s,i<j\leq3} \text{ReTr}(1 -
Figure 1. The center $\omega_0$ and the thermal width $\Gamma$ of the lowest $0^{++}$ glueball peak plotted against temperature $T$. The vertical dotted lines indicate the critical temperature $T_c \simeq 280$ MeV.

The peak center and the thermal width of thermal glueball are obtained by the best-fit analysis of the suitably smeared glueball correlator $G(\tau)/G(0)$ with Eq. (7) at various temperatures. In Fig.1, the peak center $\omega_0$ and the thermal width $\Gamma$ are plotted against temperature. While narrow-peak ansatz leads to the pole-mass reduction of 300 MeV near $T_c$ in Ref.[8], the Breit-Wigner analysis indicates a small reduction in the peak center as $\Delta \omega_0(T_c) \sim 100$ MeV. Instead, we observe a significant thermal width broadening as $\Gamma(T_c) \sim 300$ MeV.

4. SUMMARY

We have studied the temporal $0^{++}$ glueball correlator at finite temperature using SU(3) anisotropic lattice QCD at quenched level with 5,500 to 9,900 gauge configurations at each temperature. We have proposed the Breit-Wigner ansatz for the fit-function to take into account the effect of the non-zero thermal width at finite temperature. We have applied the Breit-Wigner analysis to the temporal glueball correlator at finite temperature, and have observed a slight reduction of the peak center as $\Delta \omega_0(T_c) \sim 100$ MeV and the significant broadening of the thermal width as $\Gamma(T_c) \sim 300$ MeV in the vicinity of the critical temperature $T_c$.

REFERENCES

7. T. Umeda et al., Int. J. Mod. Phys. A16 (’01) 2215.
10. S. Datta et al., hep-lat/0208001.
11. N. Ishii et al., hep-lat/0206020.