The general form of the $\star$-commutator on the Grassman algebra

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Abstract

We study the general form of the $\star$-commutator treated as a deformation of the Poisson bracket on the Grassman algebra. We show that, up to a similarity transformation, there are other deformations of the Poisson bracket in addition to the Moyal commutator (one at even and one at odd $n$, $n$ is the number of the generators of the Grassman algebra) which are not reduced to the Moyal commutator by a similarity transformation.
1 Introduction

As is well known, the major hopes for the construction of quantum mechanics on nontrivial manifolds are connected with the so-called geometric or deformation quantizations ([1] – [5]). The functions on phase space are put into correspondence with the operators and the product of the operators and their commutator are described by an associative \( * \)-product and \( * \)-commutator of functions, which represent deformations of a usual “pointwise” product and the Poisson bracket. On even manifolds, at least locally, the \( * \)-commutator coincides with usual commutator in algebra with the associative \( * \)-product. It is interesting to find out, what is the situation in the case, when the phase space is a supermanifold. In present paper we investigate a general form of the deformation of the nonsingular Poisson bracket on the Grassman algebra. We show that on the Grassman algebra in addition to the \( * \)-commutators which are equivalent to the Moyal commutator [6], there are also other deformations which are not reduced to the Moyal commutator by a similarity transformation.

The paper is organized as follows. In Sect. 2 we formulate the problem. In Sect. 3 the solution of the Jacobi identity considered as an equation to the \( * \)-commutator is found in the lowest approximation in deformation parameter. In Sect. 4 the higher deformations are considered and the main result of the paper is formulated. In two Appendices solutions of the equations considered in the main text are presented.

Notations and conventions.

\( \xi^\alpha, \alpha = 1, \ldots, n \) are odd anticommuting generators of the Grassman algebra:

\[
\varepsilon(\xi^\alpha) = 1, \quad \xi^\alpha \xi^\beta + \xi^\beta \xi^\alpha = 0, \quad \partial_\alpha \equiv \frac{\partial}{\partial \xi^\alpha},
\]

\( \varepsilon(A) \) denotes the Grassman parity of \( A \):

\[
[\xi^\alpha]^0 \equiv 1, \quad [\xi^\alpha]^k \equiv \xi^\alpha_1 \cdots \xi^\alpha_k, 1 \leq k \leq n, \quad [\xi^\alpha]^k = 0, k > n,
\]

\[
[\partial_\alpha]^0 \equiv 1, \quad [\partial_\alpha]^k \equiv \partial_\alpha_1 \cdots \partial_\alpha_k, 1 \leq k \leq n, \quad [\partial_\alpha]^k = 0, k > n,
\]

\[
[T_{...\alpha_k...}]_{...\alpha_i...} = T_{...\alpha_1...\alpha_k...}, \quad T_{...\alpha_i...\alpha_{i+1}...} = -T_{...\alpha_{i+1}...\alpha_i...}, \quad i = 1, \ldots, k - 1.
\]

2 Formulation of the problem

We consider the Grassman algebra \( G_K \) over the field \( K = \mathbb{C} \) or \( \mathbb{R} \), with the generators \( \xi^\alpha, \alpha = 1, 2, \ldots, n \). The generic element \( f \) of the algebra (a function of the generators) is

\[
f \equiv f(\xi) = \sum_{k=0}^{n} \frac{1}{k!} f_{[\alpha]_k} [\xi^\alpha]^k, \quad f_{[\alpha]_0} \equiv f_0 = \text{const}, \quad f_{[\alpha]_k} \in K.
\]

Let the nonsingular Poisson bracket

\[
\{f_1, f_2\}(\xi) \equiv [f_1, f_2]_\omega(\xi) \equiv c_0(\xi) f_1 f_2 = f_1(\xi) \overleftarrow{\partial_\alpha} \omega^\alpha(\xi) \partial_\beta f_2(\xi)
\]
be given on the Grassman algebra $G_K$ where $\omega^{\alpha\beta}(\xi)$ is the symplectic metric, the even symmetric tensor function satisfying the Jacobi identity:

$$\omega^{\beta\alpha} = \omega^{\alpha\beta}, \quad \omega^{\alpha\beta}\partial_\beta \omega^{\beta\gamma} + \omega^{\gamma\delta}\partial_\delta \omega^{\alpha\beta} + \omega^{\beta\delta}\partial_\delta \omega^{\gamma\alpha} = 0.$$

We are interested in the general form of the $*$–commutator, i.e. the bilinear mapping $G_K \times G_K \to G_K$:

$$f_3(\xi) \equiv [f_1, f_2]_*(\xi) = C(\xi|f_1, f_2),$$

with the following properties:

i) The mapping is even

$$\varepsilon(C(\xi|f_1, f_2)) = \varepsilon(f_1) + \varepsilon(f_2).$$

ii) Antisymmetry

$$C(\xi|f_2, f_1) = -(-1)^{\varepsilon(f_1)\varepsilon(f_2)}C(\xi|f_1, f_2).$$

iii) The Jacobi identity

$$( -1)^{\varepsilon(f_1)\varepsilon(f_3)} [f_1, [f_2, f_3]]_* + \text{cycle}(1, 2, 3) = 0, \quad \forall f_1, f_2, f_3,$$

or, equivalently,

$$( -1)^{\varepsilon(f_1)\varepsilon(f_3)} C(\xi|f_1, f_23) + \text{cycle}(1, 2, 3) = 0, \quad \forall f_1, f_2, f_3,$$

where we denote $f_{ij}(\xi) \equiv C(\xi|f_i, f_j)$.

iv) It is supposed, that the $*$–commutator is defined in terms of the series in the deformation parameter $\hbar^2$,

$$[f_1, f_2]_*(\xi) = \sum_{k=0} h^{2k} [f_1, f_2]_{*k}(\xi) = \sum_{k=0} h^{2k} c_k(\xi|f_1, f_2),$$

where $[f_1, f_2]_{*0} = c_0(\xi|f_1, f_2)$ coincides with the Poisson bracket (1) (the boundary condition or the correspondence principle). Differently, we want to establish the general form of a deformation of the Poisson bracket on the Grassman algebra. The deformation parameter is denoted by $\hbar^2$ and, without loss of generality, it is possible to consider $\hbar$ as real and positive number.

We treat condition (2) (or, equivalently, (3)) as an equation on possible structure of the $*$–commutator. Note that $CT$ (or, equivalently, the $*T$–commutator),

$$CT(\xi|f_1, f_2) = T^{-1}C(\xi|Tf_1, Tf_2),$$

$T$ is a nonsingular even linear mapping $G_K \to G_K$, satisfies the properties i) - iii), if $C(\xi|f_1, f_2)$ satisfies these properties. However, the boundary condition can change. In general, the operator $T$ can be represented as

$$T = T'(1 + \hbar^2 t_1 + \ldots), \quad T'|_{\hbar=0}.$$

Let $T' = T_\eta$ be the operator responsible for the change of generators:

$$T_\eta \xi^\alpha = \eta^\alpha(\xi), \quad T_\eta^{-1} \xi^\alpha = T_\zeta \xi^\alpha = \zeta^\alpha(\xi), \quad T_\eta f(\xi) = f(\eta(\xi)), \quad T_\eta^{-1} f(\xi) = T_\zeta f(\xi) = f(\zeta(\xi)),$$
where $\zeta^\alpha(\xi)$ is inverse change, $\eta^\alpha(\zeta(\xi)) = \xi^\alpha$, $\varepsilon(\eta^\alpha) = \varepsilon(\zeta^\alpha) = 1$. The $*_T$–commutator obey the property iv) but now with modified Poisson bracket:

$$\omega^{\alpha\beta}(\xi) \quad \longrightarrow \quad \omega^{\alpha\beta}_{*T}(\xi) = \xi^\alpha \frac{\partial}{\partial \xi^\alpha} \omega^{\epsilon\beta}(\zeta(\xi)) \frac{\partial}{\partial \xi^\beta}.$$ 

We say, that the $*_T$–commutator and $*$–commutator are related by a similarity transformation, or call them equivalent. The $*$–commutator obtained from the Poisson bracket by a similarity transformation is a trivial deformation of the Poisson bracket.

We assume, that after appropriate similarity transformation the Poisson bracket is of canonical form

$$\{f_1, f_2\}(\xi) \equiv [f_1, [f_2, f_3]_{\alpha_0}\}(\xi) \equiv c_0(\xi|f_1; f_2) = f_1(\xi) \frac{\partial}{\partial \xi^\alpha} \lambda_\alpha \partial_{\alpha} f_2(\xi), \quad \lambda_\alpha^2 = 1, \quad (7)$$

where $\lambda_\alpha = 1$ in the case of the Grassman algebra over complex numbers ($\mathcal{G}_C$) and $\lambda_\alpha = \pm 1$ in the case of the Grassman algebra over real numbers ($\mathcal{G}_R$).

Eq. (2) (or equivalent eq. (3)) will be solved in terms of expansion in deformation parameter $h^2$. It is obviously satisfied in the zeroth approximation. In the first order one gets

$$(-1)^{\varepsilon(f_1)\varepsilon(f_3)}[f_1, [f_2, f_3]_{\alpha_1}\]_{\alpha_0} + (-1)^{\varepsilon(f_1)\varepsilon(f_3)}[f_1, [f_2, f_3]_{\alpha_0}\]_{\alpha_1} + \text{cycle}(1, 2, 3) = 0, \quad (8)$$

or, equivalently,

$$(-1)^{\varepsilon(f_1)\varepsilon(f_3)}[f_1, c_1(f_2, f_3)]_{\alpha_0}(\xi) + (-1)^{\varepsilon(f_1)\varepsilon(f_3)}c_1(\xi|f_1, [f_2, f_3]_{\alpha_0}\) + \text{cycle}(1, 2, 3) = 0. \quad (9)$$

Besides, the conditions i) and ii) should be satisfied, or

$$c_1(\xi|f_2, f_1) = -(1)^{\varepsilon(f_1)\varepsilon(f_2)}c_1(\xi|f_1, f_2), \quad \varepsilon(c_1(\xi|f_1, f_2)) = \varepsilon(f_1) + \varepsilon(f_2).$$

The obvious solution of equation (8) or (9) is a bilinear functional $c_{1\text{triv}}$, obtained from the Poisson bracket by the similarity transformation $T f = T(\xi|f) = (1 + h t_1 + O(h^2))f(\xi) = f(\xi) + h t_1(\xi|f) + O(h^2)$, $\varepsilon(T f) = \varepsilon(t_1 f) = \varepsilon(f)$:

$$c_{1\text{triv}}(\xi|f_1, f_2) = [f_1, t_1 f_2]_{\alpha_0}(\xi) = t_1(\xi|[f_1, f_2]_{\alpha_0}\) + [t_1 f_1, f_2]_{\alpha_0}(\xi). \quad (10)$$

Consider the algebra $\mathcal{B}_k = \bigoplus \mathcal{B}_k$ of even $k$–linear functionals $\Phi_k(\xi|f_1, \ldots, f_k) \in \mathcal{B}_k$ ($k$–linear mappings $(\mathcal{G} \times)^k \to \mathcal{G}$), $\varepsilon(\Phi_k(\xi|f_1, \ldots, f_k) = \varepsilon(f_1) + \cdots + \varepsilon(f_k)$ (mod 2), $k \geq 1$, $\mathcal{B}_0 = K$, with the property

$$\Phi_k(\xi|q_1, \ldots, f_{i-1}, f_i, f_{i+1}, \ldots, f_n) = -(1)^{\varepsilon(f_1)\varepsilon(f_{i+1})}\Phi_k(\xi|q_1, \ldots, f_{i-1}, f_{i+1}, f_i, \ldots, f_n).$$

The product $\circ$ in this algebra (a mapping $\mathcal{B}_k \times \mathcal{B}_l \to \mathcal{B}_{k+l}$) is defined as follows:

$$\Phi_k \circ \Phi_l(\xi|f_1, \ldots, f_{k+l}) = \frac{1}{(k+l)!} \sum_P \text{sign}(P)\Phi_k(\xi|f_1, \ldots, f_{k+l})\Phi_l(\xi|f_{k+l+1}, \ldots, f_{k+l+1}),$$

where the sum is taken over all permutations $P$ of numbers $1, 2, \ldots, k+l$; $\text{sign}(P)$ is the parity of the permutation $P$ calculated by the following rule: the permutation of the neighboring
indexes \( i_p \) and \( i_q \) gives the sign factor \((-1)^{1+\varepsilon(f_p)\varepsilon(f_q)}\). In this algebra there is a natural grading \( g(\Phi_k) = k \), turning the algebra \( B \) into the graded algebra, \( g(B_k) = k \), and the linear operator \( d_S \), the differential coboundary Chevalley operator \( B_k \to B_{k+1} \), \( g(d_H) = 1 \), acting according to the rule:

\[
d_S \Phi_k(f_1, \ldots, f_{k+1}) = \sum_{i=1}^{k+1} (-1)^{i+1+\varepsilon(f_i)} \sum_{i=1}^{i+1} \varepsilon(f_i) f_i(\xi) *_0 \Phi_k(\xi|f_1, \ldots, f_{i-1}, f_{i+1}, \ldots, f_{k+1}) +
\]

\[
+ \sum_{1 \leq i < j \leq k+1} (-1)^{i+1+\varepsilon(f_i)}(\varepsilon(f_{i+1}) + \cdots + \varepsilon(f_{j-1})) \Phi_k(\xi|f_1, \ldots, f_{i-1}, f_i*0f_j, f_{i+1}, \ldots, f_{j-1}, f_{j+1}, \ldots, f_{k+1}),
\]

\[
d_S B_0 = 0.
\]

It is easy to prove that

\[
d_S^2 = 0, \quad d_S (\Phi_k \circ \Phi_l) = (d_S \Phi_k) \circ \Phi_l + (-1)^{g(\Phi_k)} \Phi_k \circ d_S \Phi_l.
\]

Eqs. (9) and (10) can be rewritten in terms of the Chevalley operator \( d_S \) as

\[
d_S c_1(\xi|f_1, f_2) = 0,
\]

\[
c_{1\text{triv}}(\xi|f_1, f_2) = d_S t_1(\xi|f_1, f_2),
\]

i.e. it means that \( c_1 \) belongs to the second Chevalley cogomology group, while the first order trivial deformations of the Poisson bracket are coboundaries, i.e. they belong to the zeroth Chevalley cogomology. Thus, the problem of description of all solutions of eq. (9) can be now formulated as a problem of calculation of the second Chevalley cogomology group \( H^2(d_S, B) \).

### 3 The first order deformation

It is convenient to use the momentum representation of the \(*\)-commutator (any bilinear functional can be represented in such a form)

\[
c_1(\xi|f_1, f_2) = f_1(\xi) \sum_{k,l=0}^{n} [\overline{\partial}_a]^k c_1^{[\alpha]k[\beta]l}(\xi)[\partial_\beta]^l f_2(\xi).
\]

The properties i) and ii) give:

\[
\varepsilon(c_1^{[\alpha]k[\beta]l}) = k + l \pmod{2}, \quad c_1^{[\beta]l[\alpha]k} = -(-1)^{kl} c_1^{[\alpha]k[\beta]l}.
\]  

(11)
The equation for the coefficient functions $c_1^{[\alpha_1][\beta_1]}(\xi)$ follows from eq. (9)
\begin{align*}
  f_1(\xi) &\quad \partial_\gamma \lambda_\gamma \partial_\gamma \left( f_2(\xi) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_3(\xi) \right) + \\
  +(-1)^{\varepsilon(f_2)\varepsilon(f_3)} &\quad \left( f_1(\xi) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_3(\xi) \right) \partial_\gamma \lambda_\gamma \partial_\gamma f_2(\xi) - \\
  -\left( f_1(\xi) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_2(\xi) \right) \partial_\gamma \lambda_\gamma \partial_\gamma f_3(\xi) + \\
  + f_1(\xi) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_3(\xi) \right) f_2(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_3(\xi) + \\
  +(-1)^{\varepsilon(f_2)\varepsilon(f_3)} &\quad \left( f_1(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_3(\xi) \right) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_2(\xi) - \\
  -\left( f_1(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_2(\xi) \right) \sum_{k,l=0}^n \left[\partial_\alpha\right]^k c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_3(\xi) = 0.
\end{align*}

We also write the similarity transformation in the momentum representation
\begin{align*}
  t_1(\xi|f) = \sum_{k=0}^n t_1^{[\alpha_1]}(\xi)[\partial_\alpha]^k f(\xi), &\quad \varepsilon(t_1^{[\alpha_1]}) = k \ (mod 2), \\
  c_{1triv}(\xi|f_1, f_2) = f_1(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma \left( \sum_{k=0}^n t_1^{[\alpha_1]}(\xi)[\partial_\alpha]^k f_2(\xi) \right) - \\
  -\sum_{k=0}^n t_1^{[\alpha_1]}(\xi)[\partial_\alpha]^k \left( f_1(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_2(\xi) \right) + \left( \sum_{k=0}^n t_1^{[\alpha_1]}(\xi)[\partial_\alpha]^k f_1(\xi) \right) \partial_\gamma \lambda_\gamma \partial_\gamma f_2(\xi).
\end{align*}

We solve eq. (12) studying the coefficients in front of different degrees of the derivatives of the functions $f_i$. From now on it is assumed that $n \geq 2$.

**Step 0** Consider the factor in eq. (12) in front of $f_1(\xi)$ (without derivatives):
\begin{align*}
  (-1)^{\varepsilon(f_2)\varepsilon(f_3)} &\quad \sum_{l=0}^n c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_3(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_2(\xi) - \\
  -\sum_{l=0}^n c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l f_2(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_3(\xi) + \sum_{l=0}^n c_1^{[\alpha_1][\beta_1]}(\xi)[\partial_\beta]^l \left( f_2(\xi) \partial_\gamma \lambda_\gamma \partial_\gamma f_3(\xi) \right) = 0.
\end{align*}

From antisymmetry condition (11) it follows that $c_1^{[0][0]}(\xi) = 0$.

**0a** The coefficient in front of $\partial_\alpha \partial_\beta f_2(\xi)$ gives:
\begin{align*}
  \sum_{l=1}^n (-1)^l c_1^{[\alpha_1][\beta_1]}(\xi) \partial_\alpha \lambda_\alpha \partial_\beta [\partial_\beta]^l f_3(\xi) - c_1^{[0][0]}(\xi) \partial_\alpha \lambda_\alpha f_3(\xi) = 0.
\end{align*}

From these equations it follows (see Appendix 1) that:
\begin{align*}
  \lambda_\alpha \partial_\alpha c_1^{[0][0]}(\xi) + \lambda_\beta \partial_\beta c_1^{[0][0]}(\xi) = 0 &\quad \implies \ c_1^{[0][0]}(\xi) = \lambda_\alpha \partial_\alpha c_{10}(\xi), \quad \varepsilon(c_{10}) = 0, \quad (14) \\
  \partial_\alpha c_1^{[0][0]}[\partial_\beta]^l f_3(\xi) = 0 &\quad \implies \ c_1^{[0][0]}(\xi) = const., \quad l \geq 2.
\end{align*}

**0b** The coefficient in front of $\partial_\alpha \partial_\beta f_2(\xi)$, $f_3(\xi) = \exp(\xi^i p_i)$, $p_\alpha = p_\beta + p_\gamma + p_\delta = 0$, provides:
\begin{align*}
  \lambda_\alpha c_1^{[0][0]}[\partial_\gamma]^l p_\beta - \lambda_\beta c_1^{[0][0]}[\partial_\gamma]^l p_\alpha = 0, \quad l \geq 2.
\end{align*}
The general solution of eq. (15) is (see Appendix 1):
\[
c_1^{[0][\alpha][l]} = 0, \quad 2 \leq l \leq n - 1, \quad c_1^{[0][\alpha][n]} = c_{1n} \varepsilon^{[\alpha][n]}, \quad c_{1n} = \text{const,} \quad \varepsilon(c_{1n}) = n \pmod{2}.
\]

The contribution from the similarity transformation to the factor \(c_1^{[0][\alpha][k]}\) arises only from the term \(t_1^0\) in the expression (13) and is equal to
\[
c_{1\text{triv}}^{[0][\alpha]} = -\lambda_\alpha \partial_\alpha t_1^0(\xi), \quad c_{1\text{triv}}^{[0][\alpha][k]} = 0, \quad k \geq 2.
\]

Performing the similarity transformation with \(t_1^0(\xi) = c_{10}(\xi)\), we obtain the expression for \(*\)-commutator with \(c_1^{[0][\alpha][k]} = c_1^{[\alpha][k][0]} = 0, \quad k \leq n - 1, \quad \text{and}
\[
c_1^{[0][\alpha][n]} = c_{1n} \varepsilon^{[\alpha][n]}, \quad c_1^{[\alpha][n][0]} = -c_{1n} \varepsilon^{[\alpha][n]}, \quad c_{1n} = \text{const,} \quad \varepsilon(c_{1n}) = n \pmod{2}.
\]

Note, that in the case of the Grassmann algebras over \(\mathbb{C}\) or \(\mathbb{R}\) the coefficient \(c_{1n}\) can be different from zero only at even \(n\). Furthermore, if one requires, that the functions \(f(\xi) = \text{const} \ast\)-commute with any function, then \(c_{1n} = 0\).

**Step 1** The coefficient in eq. (12) in front of \(\partial_\gamma f_1(\xi)\) gives:
\[
\lambda_\gamma f_2(\xi) \sum_{k,l=1}^{n} (-1)^k [\tilde{\partial}_\alpha c_1^{[\alpha][l][\beta]}(\xi) [\partial_\beta] f_3(\xi) - \sum_{l=1}^{n} \partial_\alpha c_1^{[\alpha][l][\beta]}(\xi) [\partial_\beta] f_2(\xi) \lambda_\gamma \partial_\alpha f_3(\xi) + +(-1)^{\varepsilon(f_2)} \lambda_\alpha \partial_\alpha f_2(\xi) \sum_{l=1}^{n} \partial_\alpha c_1^{[\alpha][l][\beta]}(\xi) [\partial_\beta] f_3(\xi) - 2c_{1n} \varepsilon^n [\partial_\beta] f_2(\xi) \lambda_\gamma \partial_\gamma f_3(\xi) + +(-1)^{\varepsilon(f_2)} 2\lambda_\alpha \partial_\gamma f_2(\xi) c_{1n} \varepsilon^n [\partial_\beta] f_3(\xi) - \sum_{l=0}^{n} c_1^{[\alpha][l][\beta][\gamma]} [\partial_\beta] f_2(\xi) \lambda_\alpha \partial_\alpha f_3(\xi) f_{2\xi}(\xi) = 0.
\]

1a The coefficient in front of \(\partial_\alpha f_2(\xi), \partial_\gamma)^k f_3(\xi)\) provides:
\[
\lambda_\alpha \partial_\alpha c_1^{[\alpha][\gamma][k]}(\xi) + \lambda_\beta \partial_\beta c_1^{[\alpha][\gamma][k]}(\xi) + \delta_{1k} \lambda_\gamma \partial_\gamma c_1^{[\alpha][\gamma]}(\xi) + \delta_{kn} 2c_{1n} \lambda_\alpha \delta_{\alpha \beta} \varepsilon^{[\gamma][n]} = 0, \quad k = 1, \ldots, n.
\]

\[1_{a1} \quad k = 1\]
\[
\lambda_\alpha \partial_\alpha c_1^{[\alpha][\gamma]}(\xi) + \lambda_\beta \partial_\beta c_1^{[\alpha][\gamma]}(\xi) + \lambda_\gamma \partial_\gamma c_1^{[\alpha][\gamma]}(\xi) = 0 \implies c_1^{[\alpha][\beta]}(\xi) = \lambda_\alpha \partial_\alpha c_1^{[\beta][\alpha]}(\xi) + \lambda_\beta \partial_\beta c_1^{[\alpha][\beta]}(\xi), \quad (17)
\]

with some function \(c_1^{[\alpha][\beta]}(\xi)\) (see Appendix 1); here it was taken into account that from antisymmetry properties (11) of the coefficients \(c_1^{[\alpha][k][\beta][l]}\) it follows that \(c_1^{[0][\alpha][\beta]} = c_1^{[\beta][\alpha]}\).

\[1_{a2} \quad 2 \leq k \leq n - 1\]
\[
\lambda_\alpha \partial_\alpha c_1^{[\alpha][\gamma][k]}(\xi) + \lambda_\beta \partial_\beta c_1^{[\alpha][\gamma][k]}(\xi) = 0, \implies c_1^{[\alpha][\beta][k]}(\xi) = \lambda_\alpha \partial_\alpha c_1^{[\beta][\alpha][k]}(\xi), \quad k = 2, \ldots, n - 1, \quad (18)
\]

with some functions \(c_1^{[\beta][k]}(\xi)\) (see Appendix 1).

\[1_{a2} \quad k = n\]
\[
\lambda_\alpha \partial_\alpha c_1^{[\alpha][\gamma][n]}(\xi) + \lambda_\beta \partial_\beta c_1^{[\alpha][\gamma][n]}(\xi) = \lambda_\alpha c_{1n} \delta_{\alpha \beta} \varepsilon^{[\gamma][n]} \implies (19)
\]
\[
c_1^{[\alpha][\beta][n]}(\xi) = (\lambda_\alpha \partial_\alpha c_1^{[n]}(\xi) - \frac{1}{2} \varepsilon^{[\alpha]} c_{1n}) \varepsilon^{[\beta][n]}.
\]
with some function $c'_{1n}(\xi)$. Perform the similarity transformation with $t_1^0(\xi) = 0, t_1^{[a]k}(\xi) = -c_1^{[a]k}(\xi), k = 1, \ldots, n - 1, t_1^{[a]n} = -c'_1(\xi)\varepsilon^{[a]n}$. It does not change the coefficients $c^{[a]k}$ and reduces $c_1^{[\beta][k]}$ to the form

$$c_1^{[\beta][k]}(\xi) = 0, \quad k \leq n - 1, \quad c_1^{[\beta][n]}(\xi) = -\frac{1}{2}\xi^{[\beta]}c_{1n}\varepsilon^{[\beta]n}.$$

1b The coefficients in front of $[\partial_\alpha]^k f_2(\xi)[\partial_\beta]^l f_3(\xi), k, l \geq 2$, give:

$$\lambda_\gamma f_2(\xi)[\overline{\partial}_\alpha]^k[\overline{\partial}_\beta]^l c_1^{[\beta][k]}(\xi)[\partial_\beta]^l f_3(\xi) -$$

$$-\delta_{k+l,n+2}(-1)^{(l-1)}C_n^{k-1}\xi^{[\alpha]k-1}\lambda_\alpha[\partial_\alpha][\partial_\beta][\partial_\gamma][\partial_\delta][\partial_\epsilon][\partial_\sigma]^l f_2(\xi)\partial_\sigma[\partial_\delta]^l f_3(\xi) = 0. \quad (20)$$

It is easy to see that second term in eq. (20) is identically equal to zero, so from eq. (20) it follows

$$c_1^{[\alpha][k][\beta][l]}(\xi) = c_1^{[\alpha][k][\beta][l]} = \text{const}, \quad \forall k, l \geq 2.$$

For $n = 2$ we have $c_1^{[\alpha][2][\beta][2]} = c_1^{(2)}\varepsilon^{[\alpha][2][\beta][2]}$. Antisymmetry condition (11) gives $c_1^{(2)} = -c_1^{(2)} = 0$. Thus one obtains for $n = 2$

$$[f_1, f_2]_{1n}(\xi) = d_{st}t_1(\xi)[f_1, f_2] +$$

$$+c_1\left(f_1(\xi)(1 - \frac{1}{2}\overline{\partial}_\alpha\xi^{[\beta]}[\overline{\partial}_\beta]^2) f_2(\xi) f_1(\xi)[\overline{\partial}_\alpha]^2\varepsilon^{[\alpha][2]}(1 - \frac{1}{2}\xi^{[\beta]}\overline{\partial}_\beta) f_2(\xi)\right). \quad (21)$$

In what follows $n \geq 3$.

2 Antisymmetry condition (11) gives $c_1^{(2)} = -c_1^{(2)} = 0$. Thus one obtains for $n = 2$

$$k\left(\lambda_\gamma f_2(\xi)[\overline{\partial}_\alpha]^k[\overline{\partial}_\beta]^l c_1^{[\beta][k-1][\gamma][l-1]} - \lambda_\beta f_2(\xi)[\overline{\partial}_\alpha]^k[\overline{\partial}_\gamma]^l c_1^{[\alpha][k-1][\beta][l]}\right)[\partial_\delta]^l f_3(\xi) +$$

$$+f_2(\xi)[\overline{\partial}_\gamma]^k\left(c_1^{[\gamma][k][\beta][l-1]} - \lambda_\alpha[\partial_\alpha][\partial_\beta][\partial_\gamma][\partial_\delta][\partial_\epsilon][\partial_\sigma]^l f_3(\xi) - c_1^{[\gamma][k][\alpha][l-1][\beta][l]}\lambda_\beta[\partial_\alpha][\partial_\gamma][\partial_\delta][\partial_\epsilon][\partial_\sigma]^l f_3(\xi)\right) -$$

$$-2(-1)^{k+l}C_n^{k-1}\xi^{[\gamma][k-1][\alpha][l-1][\beta][l]}[\partial_\delta]^l f_3(\xi) = 0, \quad k, l \geq 2. \quad (22)$$

$$f_2(\xi) = \xi\gamma[\xi^\sigma], \quad f_3(\xi) = \exp(\xi[\delta]p_\delta):$$

$$\left(\delta_{\alpha\beta}\lambda_\gamma c_1^{[\gamma][l][\beta][l]} - \delta_{\gamma\alpha}\lambda_\beta c_1^{[\alpha][l][\gamma][l]} - \delta_{\gamma\beta}\lambda_\alpha c_1^{[\alpha][l][\gamma][l]} + \delta_{\sigma\alpha}\lambda_\gamma c_1^{[\gamma][l][\sigma][l]}\right) p_\delta =$$

$$= l\left(\delta_{\sigma\delta}c_1^{[\sigma][l][\delta][l-1]} - \delta_{\beta\delta}c_1^{[\alpha][l][\sigma][l-1]} - \delta_{\gamma\delta}c_1^{[\alpha][l][\sigma][l-1]} + \delta_{\sigma\delta}c_1^{[\gamma][l][\sigma][l-1]}\right)\lambda_\delta p_\delta = 0. \quad (23)$$

The general solution of eq. (23) is (see Appendix 1):

$$c_1^{[\alpha][\beta][l]}[p_\delta]^l = \left(\delta_{\alpha\delta}c_1^{[\beta][l][\delta][l-1]} - \delta_{\beta\delta}c_1^{[\alpha][l][\beta][l-1]}\right)\lambda_\delta p_\delta [p_\delta]^l, \quad l \geq 2, \quad (24)$$

where tensor $a_1^{[\alpha][\beta][\gamma][l-1]} = a_1^{[\alpha][l]}$ is totally antisymmetric. Note that according to formula (24), $c_1^{[\alpha][l]} = 0$.

Perform the similarity transformation with $t_1^{[a]k}(\xi) = (2/k)a_1^{[a]k}, k \geq 2$. According to eq. (13) one gets

$$c_{1\text{triv}}(\xi|f_1, f_2) = -f_1(\xi)[\overline{\partial}_\alpha][\overline{\partial}_\beta] \sum_{l=2} (\delta_{\alpha\delta}a_1^{[\beta][l][\delta][l-1]} - \delta_{\beta\delta}a_1^{[\alpha][l][\beta][l-1]}\lambda_\delta p_\delta [p_\delta]^l f_2(\xi) + \ldots,$$
where dots mean the terms containing the third and higher derivatives of the function \( f_1(\xi) \).
This transformation cancels the contribution to the \( \ast \)-commutator containing the terms \( c_1^{[\beta][\delta]} \), \( 2 \leq l < n \), i.e., one may put
\[
c_1^{[\alpha][\delta]} = 0, \quad 2 \leq l \leq n.
\]

**Step 3** Now eq. (22) is simplified because the third term in the left hand side is equal to zero (including the cases \( k + l \geq n + 2 \)):
\[
k \left( \lambda_\alpha f_2(\xi) \overline{\partial}_\alpha \overline{\partial}_\gamma k^{-1} c_1^{[\gamma][\delta]} - \lambda_\beta f_2(\xi) \overline{\partial}_\beta \overline{\partial}_\gamma k^{-1} c_1^{[\gamma][\delta]} \right) |\partial_\delta|^l f_3(\xi) +
+ l f_2(\xi) \overline{\partial}_\gamma \left( c_1^{[\gamma][\delta]} k^{-1} \alpha \partial_\delta |\partial_\delta|^{l-1} f_3(\xi) - c_1^{[\gamma][\delta]} k^{-1} \lambda_\beta \partial_\beta |\partial_\delta|^{l-1} f_3(\xi) \right) = 0, \quad k, l \geq 2.
\]

3a The coefficients in front of \( |\partial_\gamma|^n f_2(\xi), f_3(\xi) = \exp (\xi^\delta p_\delta) \). Introducing the notation \( c_1^{[\alpha][\gamma][\delta]} = c_1^{[\alpha][\gamma]} c_1^{[\delta]} \), \( c_1^{[\gamma][\delta]} = -(1)^n c_1^{[\gamma][\delta]} \) and using the relation
\[
\overline{\partial}_\alpha \overline{\partial}_\gamma \|^{n-1} \varepsilon^{[\gamma][\delta]} = \frac{1}{n} \delta^{[\alpha][\beta]} \left( \overline{\partial}_\gamma \right)^n \varepsilon^{[\gamma][\delta]},
\]
we obtain (the first term in eq. (25) is identically equal to zero)
\[
c_1^{[\gamma][\delta]} = c_1^{[\gamma][\delta]} \left[ -\lambda_\beta p_\beta \right] = c_1^{[\gamma][\delta]} \lambda_\beta p_\beta |\partial_\delta|^{l-1}.
\]
From eq. (26) it follows that (see Appendix 1) \( c_1^{[\alpha]} = 0, \quad k < n \), i.e.
\[
c_1^{[\alpha][\gamma][\delta]} = c_1^{[\alpha][\gamma]} c_1^{[\delta]} = 0, \quad k, l < n.
\]
For \( k = l = n \) one has:
\[
c_1^{[\alpha][\gamma][\delta]} = \varepsilon^{[\alpha]} n \varepsilon^{[\gamma]} n \varepsilon^{[\delta]} n, \quad \varepsilon^{(n)} = 0, \quad c_1^{(n)} = -(1)^n c_1^{(n)},
\]
the last equality in (27) follows from antisymmetry condition (11) and that means that the coefficients \( c_1^{(n)} \) can be nonzero only for odd \( n \).

It is useful to point out that the operator \( (\overline{\partial}_\alpha)^n \varepsilon^{[\alpha]} (\overline{\partial}_\alpha)^n \varepsilon^{[\gamma]} (\overline{\partial}_\alpha)^n \varepsilon^{[\delta]} (\overline{\partial}_\alpha)^n \) can be also represented as
\[
\left( \overline{\partial}_\alpha \right)^n \varepsilon^{[\alpha]} c_1^{(n)} \varepsilon^{[\gamma]} c_1^{(n)} \varepsilon^{[\delta]} c_1^{(n)} = c_1^{(n)} \left( \overline{\partial}_\alpha \lambda_\alpha \partial_\alpha \right)^n, \quad c_1^{(n)} = -(1)^p n! \lambda_1 \cdots \lambda_n c_1^{(n)}.
\]

3b \( k = l = 3 \) (\( f_2(\xi) = \xi^{\alpha_1} \xi^{\alpha_2} \xi^{\alpha_3}, f_3(\xi) = \xi^{\beta_1} \xi^{\beta_2} \xi^{\beta_3} \)).
\[
\delta_{\alpha_1} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} + \delta_{\alpha_2} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} + \delta_{\alpha_3} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} -
- \delta_{\beta_1} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} + \delta_{\beta_2} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} + \delta_{\beta_3} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} +
+ \delta_{\alpha_1} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} + \delta_{\alpha_2} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} + \delta_{\alpha_3} \lambda_\alpha \lambda_\alpha c_1^{[\gamma][\delta]} -
- \delta_{\beta_1} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} + \delta_{\beta_2} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} + \delta_{\beta_3} \lambda_\beta \lambda_\beta c_1^{[\gamma][\delta]} = 0.
\]
The general solution of eq. (29) is (see Appendix 1)
\[
c_1^{[\alpha][\gamma][\delta]} = (1)^n \lambda_\alpha \lambda_\alpha \lambda_\alpha \sum \left( \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} \delta_{\alpha_4} \delta_{\alpha_5} \delta_{\alpha_6} \delta_{\alpha_7} \delta_{\alpha_8} \delta_{\alpha_9} \delta_{\alpha_10} \right)
\]
\[
\delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} \delta_{\alpha_4} \delta_{\alpha_5} \delta_{\alpha_6} \delta_{\alpha_7} \delta_{\alpha_8} \delta_{\alpha_9} \delta_{\alpha_10} -
- \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} - \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} - \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} -
- \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3} - \delta_{\alpha_1} \delta_{\alpha_2} \delta_{\alpha_3}.
\]
The following general structure:

It follows from eq. (32) that (see Appendix 1):

\[ f_2(q_1) + q_2 \partial + q_3 \partial^2 q_1 = p_1 \partial + p_2q_1 = q_0 \xi + \xi^2q_0 = 0, \]

(Compare with expression (28)).

The coefficients in eq. (12) in front of \( \partial \alpha_1 \partial \alpha_2 \partial \alpha_3 f_1(\xi) \) and derivatives of \( f_2(\xi) \) and \( f_3(\xi) \) of orders \( \geq 3 \).

Thus, the structure of all coefficients \( c_1^{(k)} \) are numbers they can be nonzero only for even \( n \).

The general solution of this equation is (see Appendix 1):

\[ c_1^{(k)} = -c_1^{(n-k)}. \]

Since \( c_1^{(k)} \) are numbers, they can be nonzero only for even \( n \).

The first three terms in eq. (12) are equal to zero. The other three give an equation of the following general structure:

\[ c_1^{[\beta]3}[\gamma]_{k+i-2}[\partial]^k f_2[\partial]^l f_3 + c_1^{[\beta]k}[\gamma]_{l+i-1}[\partial]^k f_2[\partial]^l f_3 + c_1^{[\beta]k+1}[\gamma]_{l}[\partial]^k f_2[\partial]^l f_3 = 0. \]

In this case, the first two terms in eq. (34) are equal to zero and we obtain:

\[ (\delta_{\alpha_1, \beta_1} c_1^{[\alpha_1][\beta_1]i-2} + \delta_{\alpha_2, \beta_1} c_1^{[\alpha_2][\beta_1]i-1} + \delta_{\alpha_3, \beta_1} c_1^{[\alpha_3][\beta_1]i} + \delta_{\alpha_4, \beta_2} c_1^{[\alpha_4][\beta_2]i} ) \gamma_{\xi} \beta_3 q_3 [\beta_3] \gamma_{l-2} = 0. \]

The general solution of this equation is (see Appendix 1)

\[ c_1^{[\beta]i}[\gamma] \gamma_i = 0. \]

Thus, the structure of all coefficients \( c_1^{[\alpha][\beta][\gamma]} \) is determined for odd \( n \).

\( n = 2m \) is even, \( k = m - 1, l = m, f_2(\xi) = \exp(\xi_3 \gamma), f_3(\xi) = \exp(\xi_4 \gamma) \).

Eq. (34) gives (with already established structure of the coefficients \( c_1^{[\alpha][\beta][\gamma]2m-k}, k \neq m, see formulas (33)):

\[ 3C_m^{-2} c_1^{i} (3 \epsilon_{\alpha_1 \alpha_2 \alpha_3 \beta_1} \gamma_{m-1} \lambda_{\beta} q_3 [\beta_3] \gamma_{m-2} p_3 [\beta_3] \gamma_{m-1} + \lambda_{\beta} q_3 [\beta_3] \gamma_{m-2} p_3 [\beta_3] \gamma_{m-1} + \epsilon_{\alpha_1 \alpha_2 \alpha_3 \beta_1} \gamma_{m-1} \delta_{\alpha_3, \gamma} + \epsilon_{\alpha_1 \alpha_2 \alpha_3 \beta_1} \gamma_{m-1} \lambda_{\beta} p_3 [\beta_3] \gamma_{m-1} - \epsilon_{\alpha_1 \alpha_2 \alpha_3 \beta_1} \gamma_{m-1} \lambda_{\beta} q_3 [\beta_3] \gamma_{m-2} [\beta_3] \gamma_{m-2} = 0. \]

It follows from eq. (36) (see Appendix 1)

\[ c_1^{[\alpha][\beta][\gamma]i} = c_1^{(m)} \epsilon^{[\alpha][\beta][\gamma]i}. \]
But then we have (compare with (33))

\[ c_1^{(m)} = -c_1^{(m)} = 0, \]

i.e.

\[ c_1^{[\alpha]_{m}[\beta]_{m}} = 0. \]

**4c**  The coefficients in eq. (34) in front of \([\partial]^{k}f_2, \ [\partial]^{l}f_3, \ k + l = n - 4.\)

In this case the number of indices for all coefficients \(c_1^{[\alpha]_{k}[\beta]_{k_2}} \) in eq. (12) is equal to \(n \) \((k_1 + k_2 = n). \) All these coefficients have the structure

\[ c_1^{[\alpha]_{k}[\beta]_{n-k}} = c_1^{(k)} \varepsilon[\alpha]_{k}[\beta]_{n-k}, \quad c_1^{(n-k)} = -c_1^{(k)}, \quad c_1^{(m)} = 0. \]

Eq. (34) acquires the form

\[
3C_{n-3}^{k-1}c_1^{(3)}\varepsilon^{\alpha\beta}c_1^{[\alpha]_{k-1}[\beta]_{n-k}}\lambda_\delta q_\delta[q_\beta]^{k-1}p_\gamma[p_\gamma]^{l-1} - (-1)^kC_{n-4}^{2}c_1^{(k)}\left(\varepsilon^{\alpha\beta}c_1^{[\alpha]_{k-1}[\beta]_{n-k}}\lambda_\delta q_\delta[q_\beta]^{k-1}p_\gamma[p_\gamma]^{l-1} - C_{n-4}^{2}c_1^{(k)}\left(\varepsilon^{\alpha\beta}c_1^{[\alpha]_{k-1}[\beta]_{n-k}}\lambda_\delta q_\delta[q_\beta]^{k-1}p_\gamma[p_\gamma]^{l-1} = 0. \right) \]

It follows from eq. (37) (see Appendix 1)

\[ c_1^{(k)} = 0, \forall k. \]

Hence, we have established, that the general solution of eq. (12) has the following form for \(n \geq 3: \)

\[
[f_1, f_2]_{1}(\xi) = \sigma f_1(\xi)(\overline{\partial}_\alpha\lambda_\alpha\partial_\alpha)^3f_2(\xi) + c_1^{(n)}f_1(\xi)[\overline{\partial}_\alpha]^n\varepsilon^{[\alpha]_{n}}\varepsilon^{[\beta]_{n}}[\partial_\beta]^nf_2(\xi) + c_1n \left( f_1(\xi)(1 - \frac{1}{2}\overline{\partial}_\alpha\xi_\alpha)\varepsilon^{[\beta]_{n}}[\partial_\beta]^nf_2(\xi) - f_1(\xi)[\overline{\partial}_\alpha]^n\varepsilon^{[\alpha]_{n}}(1 - \frac{1}{2}\xi_\beta\partial_\beta)f_2(\xi) \right) + d_st_1(\xi)[f_1, f_2], \]

Remember that one may put \(c_1^{(3)} = 0\) according to (28) and we shall accept this in what follows. By direct substitution of expression (38) into eq. (12) we see that it is satisfied without any further restrictions on the coefficients.

Now let us check that the first three terms in the r.h.s. of eq. (38) cannot be represented as \(c_{triv} = d_st. \)

To do this we should solve the equation

\[
\sigma f_1(\xi)(\overline{\partial}_\alpha\lambda_\alpha\partial_\alpha)^3f_2(\xi) + c_1^{(n)}f_1(\xi)[\overline{\partial}_\alpha]^n\varepsilon^{[\alpha]_{n}}\varepsilon^{[\beta]_{n}}[\partial_\beta]^nf_2(\xi) + c_1n \left( f_1(\xi)(1 - \frac{1}{2}\overline{\partial}_\alpha\xi_\alpha)\varepsilon^{[\beta]_{n}}[\partial_\beta]^nf_2(\xi) - f_1(\xi)[\overline{\partial}_\alpha]^n\varepsilon^{[\alpha]_{n}}(1 - \frac{1}{2}\xi_\beta\partial_\beta)f_2(\xi) \right) = d_st(\xi)[f_1, f_2]. \]

However, we prefer to solve a more general equation

\[
\text{l.h.s. of eq. (39)} = T_{n}^{-1}(d_st(\xi)[T_{n}f_1, T_{n}f_2]), \]

\[ (40) \]
where $T_n$ is an operator of changing the generators (see eq. (6)). Eq. (40) can be rewritten as:

$$
\sigma f_1(\xi) \left( \overline{\partial_\lambda} \lambda_\alpha \partial_\alpha \right)^3 f_2(\xi) + c^{(n)} f_1(\xi) \left[ \overline{\partial_\alpha} \right]^{[\alpha][\alpha][\alpha]} [\partial_\beta] f_2(\xi) + c_n \left( f_1(\xi) \left( 1 - \frac{2}{3} \varepsilon^{[\alpha][\beta]} [\partial_\alpha]^{[\alpha][\alpha][\alpha]} \right) f_2(\xi) - f_1(\xi) \left[ \overline{\partial_\alpha} \right]^{[\alpha][\alpha][\alpha]} (1 - \frac{1}{3} \varepsilon^{[\beta][\beta]} [\partial_\beta]) f_2(\xi) \right) = 
$$

$$
= \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} \left( f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) \right) - f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma \left( \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} f_1(\xi) \right) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) - \left( \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} f_1(\xi) \right) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) 
$$

(41)

Remember that we consider $c^{(3)}_1 = 0$.

First of all, note that the highest order of the derivatives acting on both functions in the l.h.s. of eq. (41) is $2n$, while the highest order of the derivatives in the r.h.s. is $n + 2$. The equality $2n = n + 2$ can not be valid for $n \geq 3$ so it follows that $c^{(n)} = 0, n \geq 2$. Further, the term in the l.h.s. of eq. (41) containing $f_1(\xi)$ without derivatives contains also $[\partial_\alpha]^{[\alpha][\alpha][\alpha]} f_2(\xi)$. At the same time the term in the r.h.s. of this equation containing $f_1(\xi)$ without derivatives contains only the first derivative of the function $f_2(\xi)$ from what it follows that $c_n = 0, n \geq 2$. As a result, eq. (41) for $n \geq 3$ accepts the form

$$
\sigma f_1(\xi) \left( \overline{\partial_\lambda} \lambda_\alpha \partial_\alpha \right)^3 f_2(\xi) = \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} \left( f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) \right) - f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma \left( \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} f_1(\xi) \right) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) 
$$

or, in the equivalent form,

$$
\sigma f_1(\xi) \left( \overline{\partial_\lambda} \lambda_\alpha \partial_\alpha \right)^3 f_2(\xi) = \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} \left( f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi) \right) - f_1(\xi) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma \left( \sum_{k=0}^n t^{[1]} [\xi(\xi)] [\partial_\xi]^{[1]} f_1(\xi) \right) \overline{\partial_\gamma} \lambda_\gamma \partial_\gamma f_2(\xi). 
$$

(42)

**Step 0'** The coefficient in eq. (42) in front of $f_1(\xi)$ without derivatives gives:

$$
\partial_\gamma t^0(\xi) \lambda_\gamma \partial_\gamma f_2(\xi) = 0, \quad \implies \quad t^0(\xi) = t_0 = \text{const.}
$$

**Step 1'** The coefficients in eq. (43) at $\partial_\alpha f_1(\xi)$, $[\partial_\beta]^{k} f_2(\xi)$, $k \geq 4$ give:

$$
\partial_\alpha t^{[\beta]}(\xi) = 0 \quad \implies \quad t^{[\beta]}(\xi) = t^{[\beta]} = \text{const}, \quad k \geq 4. 
$$

(44)

**Step 2'** The coefficient in eq. (43) in front of $\partial_\alpha \partial_\beta f_1(\xi)$, $f_2(\xi) = \exp(\xi^{[\alpha]} \partial_\alpha)$ provide:

$$
\left( t^{[\alpha]} \right)^{[\alpha]} \delta^{[\beta]} - t^{[\beta]} \partial^{[\alpha]} \delta^{[\alpha]} \right) \lambda_\gamma p_\gamma [p_\gamma]^{k} = 0, \quad k \geq 4. 
$$

(45)

The general solution of eq. (45) is (see Appendix 1)

$$
t^{[\alpha]} = 0, \quad 4 \leq k \leq n - 1, \quad t^{[\alpha]} = t_n \varepsilon^{[\alpha]}. 
$$

(46)
Eq. (43) now accepts the form
\[
\sigma f_1(\xi)(\frac{\partial}{\partial \eta^\alpha} \lambda_\alpha \partial_{\eta^\alpha})^3 f_2(\xi) = \sum_{k=0}^{3} t^{[\alpha]k}(\xi)[\partial_\alpha]^k \left( f_1(\xi) \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha f_2(\xi) \right) - f_1(\xi) \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha \left( \sum_{k=0}^{3} t^{[\alpha]k}(\xi)[\partial_\alpha]^k f_2(\xi) \right) - \left( \sum_{k=0}^{3} t^{[\alpha]k}(\xi)[\partial_\alpha]^k f_1(\xi) \right) \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha f_2(\xi).
\]

(47)

**Step 3'** The coefficient in eq. (47) in front of $[\partial_\alpha]^3 f_1(\xi)$, $[\partial_\beta]^3 f_2(\xi)$ provides:
\[
\sigma A_{\gamma_1 \gamma_2 \gamma_3}^{[\alpha]3} \lambda_{\gamma_1} \lambda_{\gamma_2} \lambda_{\gamma_3} A_{\gamma_1 \gamma_2 \gamma_3}^{[\beta]3} = 0, \quad \Rightarrow \quad \sigma = 0.
\]

So, representations (39) and (40) are exist only in the case $\sigma = \sigma^{(n)} = c_n = 0$, i.e., in particular, the second group of the Chevalley ego-mology is described by the parameters $\sigma$, $\sigma^{(n)}$, $c_n$. For completeness, we present the general solution of the equation
\[
d_{st}(\xi|f_1, f_2) = 0.
\]
It has the form
\[
t(\xi|f) = t_0(1 - \frac{1}{2} \xi^\alpha \partial_\alpha) f(\xi) + t(\xi) \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha f(\xi) + t_n \varepsilon^{[\alpha]}n [\partial_\alpha]^n f(\xi).
\]

We have shown thus, that the general solution of eq. (12) for $n \geq 3$ has form (38). Assuming, that the appropriate similarity transformation eliminating the terms of the type $d_{st1}(\xi|f_1, f_2)$ from the expression for the *-commutator is performed, we present the *-commutator for $n \geq 3$ in the form:
\[
C(\xi|f_1, f_2) = [f_1, f_2]_{M(h(\kappa_1 \lambda))}(\xi) + h^2 c_1^{(n)} f_1(\xi) [\partial_\alpha]^{[\alpha]}n \varepsilon^{[\alpha]}n [\partial_\beta]^{[\alpha]}n f_2(\xi) + h^2 c_1^n \left( f_1(\xi)(1 - \frac{1}{2} \partial_\alpha \xi^\alpha) \varepsilon^{[\beta]}n [\partial_\beta]^{[\alpha]}n f_2(\xi) - f_1(\xi) [\partial_\alpha]^{[\alpha]}n (1 - \frac{1}{2} \xi^\beta \partial_\beta) f_2(\xi) \right) + h^4 c_2(\xi|f_1, f_2) + O(h^6),
\]

(48)

\[
c_2(\xi|f_1, f_2) = f_1(\xi) \sum_{k,l=0}^{n} [\partial_\alpha]^{[\alpha]}k c_2^{[\beta]l}[\partial_\beta]^{[\alpha]}l (\xi)[\partial_\beta]^{[\alpha]}l f_2(\xi),
\]

\[
\varepsilon^{[\alpha]k}[\alpha][\beta]l = k + l \pmod{2}, \quad c_2^{[\beta]l}[\alpha][\beta]k = (-1)^{kl} c_2^{[\alpha]k}[\beta]l,
\]

where the Moyal bracket $[f_1, f_2]_{M(h(\kappa_1 \lambda))}(\xi)$, $\kappa_1 = (6\sigma_1)^{1/2}$, is defined as follows:
\[
[f_1, f_2]_{M(h(\kappa_1 \lambda))}(\xi) = \frac{1}{h \kappa_1} f_1(\xi) \sinh (h \kappa_1 \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha) f_2(\xi).
\]

(49)

The Moyal bracket (49) satisfies the Jacobi identity (2), (3). Indeed, consider the associative *-product
\[
f_1 *_{G(h(\kappa \lambda))} f_2(\xi) = f_1(\xi) \exp (h \kappa \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha) f_2(\xi)
\]
for any complex $\kappa$. It is easy to verify, that it has the property
\[
f_2 *_{G(h(\kappa \lambda))} f_1(\xi) = (-1)^{\varepsilon(f_1)\varepsilon(f_2)} f_2(\xi) \exp (h \kappa \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha) f_2(\xi).
\]

Therefore, the following representation is valid
\[
2f_1(\xi) \sinh (h \kappa \frac{\partial}{\partial \lambda_\alpha} \partial_\alpha) f_2(\xi) = f_1 *_{G(h(\kappa \lambda))} f_2(\xi) - (-1)^{\varepsilon(f_1)\varepsilon(f_2)} f_2 *_{G(h(\kappa \lambda))} f_1(\xi),
\]
from which it follows that the Moyal bracket $[f_1, f_2]_{M(h(\kappa \lambda))}(\xi)$ satisfies the Jacobi identity.
4 Higher order deformations

Now it is useful to point out, that in order to construct general solution of eq. (12), one needs equations only for the coefficients in front of

$$[\partial_\alpha]^k f_1(\xi), \quad k = 0, 1, 2,$$

and in front of

$$[\partial_\alpha]^3 f_1(\xi) [\partial_\alpha]^k f_2(\xi) [\partial_\alpha]^l f_3, \quad k, l < n \quad (\text{for } n \geq 4).$$

Substitute representation (48) into the Jacobi identity (2) which will be satisfied in the zeroth and first orders in $\hbar^2$. In the $\hbar^4$ order we have:

$$(-1)^{\varepsilon(f_1)\varepsilon(f_3)} d_{S\varepsilon}c_2(\xi|f_1, f_2, f_3) =$$

$$= 2c_1\sigma_1 \left( (-1)^{\varepsilon(f_1)\varepsilon(f_3)} f_1(\xi)[\partial_{\alpha}]^{[\alpha]}f_2(\xi)(\partial_\beta\lambda_\beta\partial_\beta)^3f_3(\xi) + \text{cycle}(1, 2, 3) \right). \quad (50)$$

In eq. (50) it was taken into account that $c_{1n}^{(n)} = 0$, $\forall n$, since the factors $c_1^{(n)}$ are equal to zero for even $n$, and the factors $c_{1n}$ are equal to zero for odd $n$.

Let $n \geq 3$.

Consider the equations arising from vanishing of the coefficients in front of $[\partial_\alpha]^k f_1(\xi)$, $k = 0, 1, 2$, and in front of $[\partial_\alpha]^3 f_1(\xi) [\partial_\alpha]^k f_2(\xi) [\partial_\alpha]^l f_3(\xi)$, $k, l < n$, in eq. (50). The r.h.s. of eq. (50) does not give the contribution to these equations, i.e., they coincide with the corresponding equations arising from the solution of homogeneous equation (12). As was already noticed, it is enough for complete definition of all coefficients of the functional $c_2(\xi|f_1, f_2)$. Furthermore, the following equality will be obviously satisfied

$$d_{S\varepsilon}c_2(\xi|f_1, f_2, f_3) = 0.$$

Then it follows from eq. (50) that

$$c_{1n}\sigma_1 = 0 \implies \text{either } c_{1n} = 0, \text{ or } \sigma_1 = 0.$$

By induction, we obtain that up to a similarity transformation, the satisfying the conditions formulated in Sect. 2 nonsingular deformations of the Poisson bracket on the Grassman algebra for $n \geq 3$ are given by the following expressions:

i) $n = 2k$, $k \geq 2$

$$[f_1, f_2]_s^{(1)}(\xi) = [f_1, f_2]_{M(h\kappa\lambda)}(\xi), \quad (51)$$

$$[f_1, f_2]_s^{(2)}(\xi) = [f_1, f_2]_{s0}(\xi) + \hbar^2 c_{2n} \left( f_1(\xi)(1 - \frac{1}{2}[\partial_\alpha]\xi^\alpha)[\partial_\beta]^n f_2(\xi) - f_1(\xi)[\partial_{\alpha}]^{[\alpha]}(1 - \frac{1}{2}\xi^\beta\partial_\beta)f_2(\xi) \right). \quad (52)$$

ii) $n = 2k + 1$, $k \geq 1$

$$[f_1, f_2]_s^{(1)}(\xi) = [f_1, f_2]_{M(h\kappa\lambda)}(\xi) + \hbar^2 c_{2n}^{(n)} f_1(\xi)[\partial_{\alpha}]^{[\alpha]}[\partial_\beta]^n f_2(\xi), \quad (53)$$

$$[f_1, f_2]_s^{(2)}(\xi) = [f_1, f_2]_{s0}(\xi) + \hbar^2 c_{2n}^{(n)} f_1(\xi)[\partial_{\alpha}]^{[\alpha]}[\partial_\beta]^n f_2(\xi). \quad (54)$$
Certainly, *(2)–commutator (54) is a limiting case (for \( \kappa \to 0 \)) of \(*\)–commutator (53). The parameters \( \kappa, c_n, c^{(n)} \) can depend on \( \hbar \). If one requires that the functions \( f(\xi) = \text{const} \) \(*\)–commute with any function, then \( c_n = 0 \).

In expressions (51) – (53) one can put \( \kappa = \lambda_\alpha = 1, c_n = 0, 1 \) in the case of the Grassman algebra \( \mathcal{G}_C \), and \( \kappa = 1, i, \lambda_\alpha = \pm 1, c_n = 0, \pm 1 \) in the case of the Grassman algebra \( \mathcal{G}_R \). Besides, in eq. (54) we may set \( c^{(n)} = 0, 1 \) in the case of the Grassman algebra \( \mathcal{G}_C \) and \( c^{(n)} = 0, \pm 1 \) in the case of the Grassman algebra \( \mathcal{G}_R \).

Indeed, after the similarity transformation of \(*\)–commutators (51) – (53) by the operator \( T = T^{(\mu)} \),

\[
T^{(\mu)} f(\xi) = \frac{1}{\mu^2} f(\mu \xi), \quad \mu \in K, \tag{55}
\]

the transformed \(*\)–commutators will preserve structure (51) – (53) (in particular, the zeroth approximation in \( \hbar^2 \) does not vary), but with new parameters:

\[
\kappa \rightarrow \kappa \mu^2, \quad c^{(n)} \rightarrow c^{(n)} \mu^{2(n-1)}, \quad c_n \rightarrow c_n \mu^{n-2}. \tag{56}
\]

It is possible to fix \( \mu \) imposing \( \kappa \mu^2 = 1 \) in eqs. (51) and (53) and \( c_n \mu^{n-2} = 1 \) for \( c_n \neq 0 \) in eq. (52) for the Grassman algebra \( \mathcal{G}_C \). In the case of the Grassman algebra \( \mathcal{G}_R \) it is possible to satisfy \( \kappa \mu^2 = 1 \) in formulas (51) and (53) for \( \sigma_1 > 0 \). If \( \sigma_1 < 0 \), then \( \kappa \) is pure imaginary and it is possible to fix \( \mu \) from the condition \( \kappa \mu^2 = i \). Note that in this case the Moyal bracket cannot be presented as a commutator in the algebra with associative product. The parameter \( c_n \) in formula (52) can be normalized to \( +1, 0, -1 \) for \( c_n > 0, = 0, < 0 \), respectively. If the parameter \( \kappa \) is fixed in eq. (53), then we already do not have a possibility to change the parameter \( c^{(n)} \) and it remains arbitrary. However, in formula (54) parameter \( c^{(n)} \) can be normalized to \( 0, 1 \) for \( c^{(n)} = 0, \neq 0 \) in the case of the Grassman algebra \( \mathcal{G}_C \) and to \( +1, 0, -1 \) for \( c^{(n)} > 0, = 0, < 0 \) in the case of the Grassman algebra \( \mathcal{G}_R \).

Consider now the remaining cases \( n = 1 \) and \( n = 2 \).

1) \( n = 1 \)

\[
f = a + b \xi, \quad [f_1, f_2]_s(\xi) = \lambda f_1(\xi) \overleftarrow{\partial} f_1(\xi) = \lambda b_1 b_2, \quad \lambda = \pm 1,
\]

\[
[f_1, f_2]_s(\xi) = \lambda f_1(\xi) \overleftarrow{\partial} f_1(\xi) + h^2 c_1(\xi) (-1)^{\varepsilon(f_1)} \left( f_1(\xi) \partial f_2(\xi) + f_1(\xi) \overleftarrow{\partial} f_2(\xi) \right) + h^2 c_2(\xi) f_1(\xi) \overleftarrow{\partial} f_1(\xi), \quad \varepsilon(c_1(\xi)) = 1 \Rightarrow c_1(\xi) = c_1, \quad \varepsilon(c_2(\xi)) = 0 \Rightarrow c_2(\xi) = c_2, \quad c_1, c_2 = \text{const}.
\]

After the similarity transformation generated by \( \xi \rightarrow \xi' = (1 + h^2 c_2/\lambda)^{-1/2} \xi \), we obtain the following expression for the \(*\)–commutator:

\[
[f_1, f_2]_s(\xi) = \lambda f_1(\xi) \overleftarrow{\partial} f_1(\xi) + h^2 c_1(\xi) (-1)^{\varepsilon(f_1)} \left( f_1(\xi) \partial f_2(\xi) + f_1(\xi) \overleftarrow{\partial} f_2(\xi) \right) = \lambda b_1 b_2 + h^2 c_1(\xi)(a_1 b_2 - b_1 a_2).
\]

Take \( f_1(\xi) = a_1, \varepsilon(f_1) = 0, f_{2,3}(\xi) = b_{2,3} \xi, \varepsilon(f_{2,3}) = 1 \) and consider the Jacobi identity (2) for these functions:

\[
0 = [f_1, [f_2, f_3]]_s + [f_2, [f_3, f_1]]_s - [f_3, [f_1, f_2]]_s = -2h^2 \lambda c_1 a_1 b_2 b_3, \quad \Rightarrow \quad c_1 = 0.
\]
Thus, the nontrivial deformation of the nonsingular Poisson bracket are absent in the case of the Grassman algebra with one generator.

ii) $n = 2$

Using the previous results (see, in particular, eqs. (50) and (21)), we obtain:

Up to the similarity transformations, the possible deformations of the nonsingular Poisson bracket on the Grassman algebra with two generators are

$$[f_1, f_2]_*(\xi) = [f_1, f_2]_{*0}(\xi) + h^2 c_2 \left( f_1(\xi) (1 - \frac{1}{2} \tilde{\partial}_\alpha \xi^\alpha) \varepsilon^{[\beta]2} [\tilde{\partial}_\beta] f_2(\xi) - f_1(\xi) \tilde{\partial}_\alpha [\tilde{\partial}_\beta]^{2} \varepsilon^{[\alpha]2} (1 - \frac{1}{2} \xi^\beta \partial_\beta) f_2(\xi) \right),$$

$c_2$ can depend on $\hbar$, i.e., are given in fact by expression (52). Note, that formally formulas (51) – (54) are valid for all $n$, since the Moyal bracket for $n = 1, 2$ reduces to the Poisson bracket, and the second terms in the r.h.s. of eqs. (53) and (54) are proportional to the Poisson bracket for $n = 1$.

Thus, we have shown that an arbitrary $*$-commutator (satisfying conditions formulated in Sect. 2) can be reduced to form (51) – (54) with the help of the similarity transformations. Is it also possible to reduce expressions (51) – (54) to the Poisson bracket by the similarity transformation? The answer to this question is in negative. It is enough to consider the case of the Grassman algebra $G_C$. Present the operator $T$ generating the similarity transformation as

$$T = T'(1 + h^2 t_1 + O(h^4)), \quad T' = T|_{\hbar = 0}$$

($T'$ is nonsingular operator). Due to boundary condition (4) the similarity transformation generated by the operator $T'$ should retain the zeroth order in $h^2$ approximation to the $*$-commutator to be the Poisson bracket (maybe with a different symplectic metric):

$$T'^{-1} \left( [T' f_1, T' f_2]_{*0} \right)(\xi) = [f_1, f_2]_{*0}(\xi) = T'^{-1} \left( [T\eta f_1, T\eta f_2]_{*0} \right)(\xi),$$

where $T\eta$ is an operator of the change of generator, reducing the Poisson bracket $*_{\eta'}$ into the canonical form $*_{0}$. The operator $T'' = T^{-1}_\eta T'$ satisfies the equation

$$T''[f_1, f_2]_{*0}(\xi) = [T'' f_1, T'' f_2]_{*0}(\xi). \quad (57)$$

The general form of such operators is found in Appendix 2:

$$T' = T^{(\mu)} T^{(c)} T^{(n)},$$

$$T^{(c)} f(\xi) = f(\eta(\xi)), \quad T^{(n)} = (1 + t_1 \varepsilon^{[\alpha][\beta]} [\partial_\alpha]_{\eta^\beta}), \quad t_n = \text{const}, \quad \varepsilon(t_n) = n(\text{mod} \ 2),$$

$\varepsilon(\eta^\alpha) = 1$, the change $\xi^\alpha \rightarrow \eta^\alpha(\xi)$ is canonical, the operator $T^{(\mu)}$ is determined by (55). Then the operator $T'$ is given by

$$T' = T^{(\mu)} T^{(c)} T^{(n)},$$

where $T^{(c)}$ is an operator of some change of the generators.

It is easy to verify that the similarity transformation generated by the operator $T^{(n)}$ leaves invariant not only the $*_{\eta}$-commutator (the Poisson bracket), but also the exact $*-$commutator given by formulas (51) – (53). The similarity transformation generated by the
operator $T^{(\mu)}$ results only in multiplication of parameters in eqs. (51) – (53) by nonzero factors (see eq. (56)). Thus, the possibility of reducing the $\ast$–commutator to the $\ast_0$–commutator in the first $h^2$ order depends on the validity of eq. (40), which, as has been shown, has no solutions for the nonzero parameters $\kappa$, $c^{(n)}$, $c_n$.

We have finally shown, that the following theorem is valid

**Theorem**

Up to similarity transformation, the satisfying the conditions formulated in Sect. 2 deformations of the nonsingular Poisson bracket on the Grassman algebra are given by expressions (51) – (54) and $\kappa = 1$, $c_n = 0, 1$ in the case of the Grassman algebra $\mathcal{G}_C$ and $\kappa = 1, i$, $c_n = 0, \pm 1$ in the case of the Grassman algebra $\mathcal{G}_R$; $c^{(n)}$ is an arbitrary complex number in the case of the algebra $\mathcal{G}_C$ and an arbitrary real number in the case of the algebra $\mathcal{G}_R$ in formula (53), and $c^{(n)} = 0, 1$ in the case of the algebra $\mathcal{G}_C$ and $c^{(n)} = 0, \pm 1$ in the case of the algebra $\mathcal{G}_R$ in formula (54). Furthermore note, that $\ast$–commutators (51) – (54) with different values of the parameter $\hbar$ can also be connected by the similarity transformation generated by the operator $T^{(\mu)}$ with suitable $\mu$.

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**Appendix 1**

In this appendix the solutions of some equations from the main text are presented.

**Equations (14), (18) and (17).**

Consider equation

$$d\Omega(x, \xi) = 0, \quad d = \xi^\alpha \lambda_\alpha \frac{\partial}{\partial x^\alpha}, \quad d^2 = 0, \quad (A.1)$$

where $x^\alpha$ are the set of even variables (generators), $x^\alpha x^\beta - x^\beta x^\alpha = x^\alpha \xi^\beta - \xi^\beta x^\alpha = 0$. Introduce the operators

$$\gamma = x^\alpha \lambda_\alpha \frac{\partial}{\partial \xi^\alpha}, \quad N = x^\alpha \frac{\partial}{\partial x^\alpha} + \xi^\alpha \frac{\partial}{\partial \xi^\alpha},$$

generating the algebra

$$\gamma^2 = 0, \quad d\gamma + \gamma d = 0, \quad [d, N] = [\gamma, N] = 0.$$

It is easy to see, that if the function $\Omega(x, \xi)$ satisfies the eq. (A.1) and the operator $N^{-1}$ is defined on $\Omega(x, \xi)$ (it is necessary to have $\Omega(0, 0) = 0$ for this), then the function $\Omega(x, \xi)$ can be presented as

$$\Omega(x, \xi) = d\Xi(x, \xi), \quad \Xi(x, \xi) = \gamma \frac{1}{N} \Omega(x, \xi),$$

and the general solution of eq. (A.1) is

$$\Omega(x, \xi) = d\Xi(x, \xi) + \text{const.} \quad (A.2)$$
Let $\Omega(x, \xi) = x^\alpha c^\alpha(\xi)$ also be solution of eq. (A.1). Then it follows from eqs. (A.1) and (A.2) that

$$\lambda_\alpha \partial_\alpha c^\beta(\xi) + \lambda_\beta \partial_\beta c^\alpha(\xi) = 0, \quad c^\alpha(\xi) = \lambda_\alpha \partial_\alpha c(\xi),$$

It gives the solutions of eqs. (14) and (18). Similarly, choosing $\Omega(x, \xi) = x^\alpha x^\beta c_{\alpha\beta}(\xi)$, we obtain the solution of eq. (17).

**Equations (15), (26) and (45).**

All these equations can be reduced to the form

$$\left( a^{\alpha[k]}(\xi)\delta_{\beta\gamma} - a^{\beta[k]}(\xi)\delta_{\alpha\gamma}\right)\lambda_\gamma \partial_\gamma[p_\gamma]^{k-1} = 0, \quad k \geq 2,$$

where $a^{\alpha[k]} = c^{[a][k]}_1 + c^{[a][k]}_2$. Acting by the operator $\lambda_\beta \partial_\beta p_\beta$ on eq. (A.3) one gets:

$$(n - k)a^{\alpha[k]} = 0 \implies a^{\alpha[k]} = 0, \quad 2 \leq k \leq n - 1.$$  

Note that $a^{\alpha[n]} = a_{\alpha}\varepsilon^{[\alpha][n]}$ satisfies eq. (A.3) identically.

**Equation (23)**

$$\left(\delta_{\gamma\lambda} \lambda_\gamma c^{[\sigma][\delta]}_1 - \delta_{\beta\lambda} \lambda_\beta c^{[\beta][\delta]}_1 + \delta_{\sigma\alpha} \lambda_\alpha c^{[\gamma][\delta]}_1\right)[p_\delta]^l =$$

$$= l\left(\delta_{\sigma\delta} c^{[\beta][\delta]}_1 - \delta_{\delta\beta} c^{[\beta][\delta]}_1 - \delta_{\gamma\delta} c^{[\sigma][\delta]}_1 + \delta_{\sigma\delta} c^{[\gamma][\delta]}_1\right)\lambda_\delta [p_\delta]^{l+1}.$$  

Acting by the operator $\lambda_\alpha \partial_\alpha$ on eq. (A.4) we obtain:

$$(n - l)c^{[\sigma][\delta]}_1[p_\delta]^{l+1} = \lambda_\delta\left(\delta_{\beta\sigma} \Delta^{[\gamma][\delta]}_i - \delta_{\beta\delta} \Delta^{[\sigma][\delta]}_i\right)[p_\delta]^{l+1} +$$

$$+ (l - 1)\left(\delta_{\sigma\delta} \Delta^{[\beta][\delta]}_i - \delta_{\beta\delta} \Delta^{[\beta][\delta]}_i\right)\lambda_\delta [p_\delta]^{l+2}, \quad \Delta^{[\sigma][\delta]}_i = 0 \forall \alpha, \beta, \gamma.$$  

Now, acting by the operator $p_\sigma \lambda_\beta \partial_\beta$ on eq. (A.5) one gets:

$$(l - 2)p_\gamma \Delta^{[p_\delta]}_i - [p_\delta]^{l+2} = 0, \quad \Delta^{[p_\delta]}_i = 0 \forall \alpha, \beta, \gamma.$$  

from where it follows that $\Delta^{[\sigma]}_i$ can be nonzero only for $l = 2$ or $l = 2 = n$. Since $l \leq n$, we get that $\Delta^{[\sigma]}_i = 0$ for $l \neq 2$.

Multiplying eq. (A.5) by $\lambda_\gamma \partial_\gamma$ one obtains:

$$\Delta^{[\sigma][\delta]}_i[p_\delta]^{l+1} + \Delta^{[\delta][p_\delta]}_i[p_\delta]^{l+2} = \Delta^{[p_\delta]}_i - \lambda_\gamma p_\gamma[p_\delta]^{l+2}.$$  

Applying the operator $\lambda_\gamma \partial_\gamma$ to eq. (A.6) we find that $\Delta^{[\sigma]}_i = 0$ also if $l = 2$. Eq. (A.6) acquires the form:

$$\Delta^{[\sigma][\delta]}_i[p_\delta]^{l+1} + \Delta^{[\delta][p_\delta]}_i[p_\delta]^{l+2} = 0.$$  

Differentiating eq. (A.7) over $p_\sigma$ we obtain

$$l\Delta^{[\sigma][\delta]}_i[p_\delta]^{l+2} = (l - 2)(\Delta^{[\delta][p_\delta]}_i[p_\delta]^{l+3} + \Delta^{[\sigma][\delta]}_i[p_\delta]^{l+2} = 0,$$

from where it follows that $\Delta^{[\sigma]}_i$ are totally antisymmetric in all indices:

$$\Delta^{[\sigma]}_i = \Delta^{[\sigma]}_i.$$
Multiplying eq. (A.5) for \( l < n \) by \( p_\beta \) we finally find

\[
c_1^{\alpha|\beta|\gamma}[p_\delta]^l = (\delta_{\alpha\delta}a_1^{\beta|\delta|-1} - \delta_{\beta\delta}a_1^{\alpha|\delta|-1})\lambda_\delta p_\delta[p_\delta]^{l-1}, \quad a_1^{\alpha|\delta|-1} \equiv -\frac{l}{n-l}\delta^{[\alpha]}, \quad 2 \leq l < n. \tag{A.8}
\]

The r.h.s. of eq. (A.4) vanishes for \( l = n \) (to see this, it is necessary to use the equality \( c_1^{\alpha|\beta|\gamma} = c_1^{\alpha|\beta|\gamma} \) and the identity \( \varepsilon^{\beta|\delta|-1}p_\alpha[p_\delta]^{n-1} = (1/n)\delta_{\alpha\beta}\varepsilon^{\beta|\delta|-1}p_\delta^n \)). Multiplying eq. (A.4) by \( \lambda_\alpha\delta_{\alpha\gamma} \):

\[
(2-n)c_1^{\beta|\gamma|\delta}[p_\delta]^n = 0,
\]

we obtain \( c_1^{\beta|\gamma|\delta} = 0 \). Note that representation (A.8) is valid also for \( l = n \).

Equation (29)

\[
d_{\alpha\alpha_1}\lambda_\alpha c_1^{\beta_1\beta_2\alpha_1\beta_3} - \delta_{\alpha\beta_1}\lambda_\beta c_1^{\beta_1\beta_2\alpha_1\beta_3} + \delta_{\alpha\beta_2}\lambda_\beta c_1^{\beta_1\beta_2\alpha_1\beta_3} - \delta_{\alpha\beta_1}\lambda_\beta c_1^{\beta_1\beta_2\alpha_1\beta_3} + \delta_{\alpha\beta_2}\lambda_\beta c_1^{\beta_1\beta_2\alpha_1\beta_3} = \lambda_\alpha c_1^{\alpha_1\alpha_2|\beta_1\beta_2}.
\tag{A.9}
\]

Multiply eq. (A.9) by \( \lambda_\alpha\delta_{\alpha\alpha_1} \) (and then replace \( \beta \) by \( \alpha_1 \)):

\[
(n - 2)c_1^{\alpha_1\alpha_2|\beta_1\beta_2} = \nu\alpha_1\left(\delta_{\alpha_1\beta_1}\Sigma_{\alpha_2|\beta_1\beta_2} + \delta_{\alpha_2\beta_1}\Sigma_{\alpha_1\beta_1\beta_2} + \Sigma_{\alpha_1\alpha_2|\beta_1\beta_2}\right) = \lambda_\gamma c_1^{\alpha_1\alpha_2|\beta_1\beta_2} = \Sigma_{\beta_1\beta_2|\alpha_1\alpha_2}.
\tag{A.10}
\]

The symmetry properties of the tensor \( \Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} \) follow from antisymmetry condition (11). Multiplying eq. (A.10) by \( \lambda_\alpha\delta_{\alpha_2\beta_1} \) one has:

\[
(n - 1)\Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} - \nu\alpha_1\left[\Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} + \Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} + \Sigma_{\alpha_1\alpha_2|\beta_1\beta_2}\right] = \lambda_\alpha\left(\delta_{\alpha_1\beta_1}\Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} - \delta_{\alpha_1\beta_1}\Sigma_{\alpha_1\alpha_2|\beta_1\beta_2}\right),
\tag{A.11}
\]

\[
\Sigma_{\alpha|\beta} = \lambda_\gamma\Sigma_{\gamma|\gamma} = \Sigma_{\alpha|\beta}.
\]

The multiplication of eq. (A.10) by \( \lambda_\alpha\delta_{\alpha_1\alpha_2} \) gives

\[
\Sigma_{\alpha_1\beta_1|\beta_2\beta_3} + \Sigma_{\alpha_2\beta_1|\beta_2\beta_3} + \Sigma_{\alpha_3\beta_1|\beta_2\beta_3} = 0 \tag{A.12}
\]

Finally, we multiply eq. (A.11) by \( \lambda_\alpha\delta_{\alpha_2\beta_1} \):

\[
\Sigma_{\alpha|\beta} = \frac{1}{n}\lambda_\alpha\delta_{\alpha_2\beta_1}\Sigma, \tag{A.13}
\]

\[
\Sigma = \lambda_\gamma\Sigma_{\gamma|\gamma}, \quad \varepsilon(\Sigma) = 0.
\]

Substituting expression (A.13) into eq. (A.11) and using condition (A.12), we obtain

\[
\Sigma_{\alpha_1\alpha_2|\beta_1\beta_2} = -\frac{1}{n(n - 1)}\lambda_\alpha_1\lambda_\alpha_2(\delta_{\alpha_1\beta_1}\delta_{\alpha_2\beta_2} - \delta_{\alpha_1\beta_2}\delta_{\alpha_2\beta_1})\Sigma. \tag{A.14}
\]
Represent $c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3}$ as

$$c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = -\frac{\lambda_{\alpha_1}\lambda_{\alpha_2}\lambda_{\alpha_3}\Sigma}{n(n-1)(n-2)} \left( \delta_{\alpha_1\beta_1} \delta_{\alpha_2\beta_2} \delta_{\alpha_3\beta_3} + \delta_{\alpha_1\beta_3} \delta_{\alpha_2\beta_1} \delta_{\alpha_3\beta_2} + \delta_{\alpha_1\beta_2} \delta_{\alpha_2\beta_3} \delta_{\alpha_3\beta_1} - \delta_{\alpha_1\beta_2} \delta_{\alpha_2\beta_1} \delta_{\alpha_3\beta_3} + \delta_{\alpha_1\beta_3} \delta_{\alpha_2\beta_2} \delta_{\alpha_3\beta_1} - \delta_{\alpha_1\beta_1} \delta_{\alpha_2\beta_3} \delta_{\alpha_3\beta_2} \right) + c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3}.$$ 

Substituting this representation into eq. (A.10) we have

$$(n - 2)c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = c_{1}^{\alpha_2\alpha_3\alpha_1|\beta_1\beta_2\beta_3} - c_{1}^{\alpha_2\alpha_3\beta_1|\beta_2\beta_3} + c_{1}^{\alpha_2\alpha_3\beta_2|\beta_2\alpha_1} - c_{1}^{\alpha_2\alpha_3\beta_3|\alpha_1\beta_2\beta_3}. \quad (A.15)$$

The r.h.s. of eq. (A.15) is antisymmetric under the permutation of the indices $\alpha_1$ and $\beta_1$, therefore, its l.h.s., the tensor $c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3}$, changes a sign under the permutation of $\alpha_1$ and $\beta_1$. Hence the tensor $c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3}$ is totally antisymmetric in all indices. In particular,

$$c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = -c_{1}^{\beta_1\beta_2\beta_3|\alpha_1\alpha_2\alpha_3}.$$ 

However, from the condition (11) it follows $c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = c_{1}^{\beta_1\beta_2\beta_3|\alpha_1\alpha_2\alpha_3}$, i.e. $c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = 0$. Thus, we finally obtain

$$c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3} = \frac{\lambda_{\alpha_1}\lambda_{\alpha_2}\lambda_{\alpha_3}\Sigma}{n(n-1)(n-2)} \left( \delta_{\alpha_1\beta_1} \delta_{\alpha_2\beta_2} \delta_{\alpha_3\beta_3} + \delta_{\alpha_1\beta_3} \delta_{\alpha_2\beta_1} \delta_{\alpha_3\beta_2} + \delta_{\alpha_1\beta_2} \delta_{\alpha_2\beta_3} \delta_{\alpha_3\beta_1} - \delta_{\alpha_1\beta_2} \delta_{\alpha_2\beta_1} \delta_{\alpha_3\beta_3} + \delta_{\alpha_1\beta_3} \delta_{\alpha_2\beta_2} \delta_{\alpha_3\beta_1} - \delta_{\alpha_1\beta_1} \delta_{\alpha_2\beta_3} \delta_{\alpha_3\beta_2} \right) + c_{1}^{\alpha_1\alpha_2\alpha_3|\beta_1\beta_2\beta_3}.$$ 

**Equation (32)**

$$k \left( \lambda_{\alpha} q_{\gamma} [q_{\gamma}]^{k-1} c_{1}^{\gamma|\kappa-1|\delta} - \lambda_{\beta} q_{\delta} [q_{\delta}]^{k-1} c_{1}^{\alpha|\kappa-1|\delta} \right) [p_{\delta}]^{l} + l [q_{\gamma}]^{k} \left( c_{1}^{\gamma|\beta|\delta|l-1} \lambda_{\alpha} p_{\alpha} [p_{\alpha}]^{l-1} - c_{1}^{\gamma|\alpha|\delta|l-1} \lambda_{\beta} p_{\beta} [p_{\beta}]^{l-1} \right) = 0, \quad 3 \leq k, l < n. \quad (A.16)$$

Acting on eq. (A.16) by the operator $\lambda_{\beta} \partial/\partial q_{\beta}$ we get:

$$(n - l)[q_{\gamma}]^{k} c_{1}^{\gamma|\kappa|\alpha|\delta|l-1} + k q_{\delta} [q_{\delta}]^{k-1} c_{1}^{\alpha|\kappa|\delta|l-1} + (-1)^{k} k \lambda_{\alpha} q_{\alpha} [q_{\gamma}]^{k-1} \Delta^{\gamma|\kappa-1|\delta|l-1}, \quad (A.17)$$

$$\Delta^{\gamma|\kappa-1|\delta|l-1} \equiv \lambda_{\gamma} c_{1}^{\gamma|\kappa-1|\delta|l-1}.$$ 

Act on eq. (A.17) by the operator $\lambda_{\alpha} \partial/\partial q_{\alpha}$ one has:

$$(k - l) \Delta^{\gamma|\kappa-1|\delta|l-1} = 0, \quad \Delta^{\gamma|\kappa-1|\delta|l-1} = 0, \quad k \neq l.$$ 

Take $k \neq l$ in eq. (A.17). Acting on this equation by the operator $\partial/\partial q_{\alpha}$ we obtain:

$$[q_{\gamma}]^{k-1} \left( (n - l) c_{1}^{\gamma|\kappa-1|\beta|\delta|l-1} + c_{1}^{\beta|\kappa-1|\alpha|\delta|l-1} \right) = (k - 1) q_{\delta} [q_{\delta}]^{k-2} c_{1}^{\beta|\kappa-2|\alpha|\delta|l-1}.$$ 

Combining this expression with the expression with permuted indices $\alpha$ and $\beta$ one gets:

$$c_{1}^{\alpha|\kappa-1|\beta|\delta|l-1} + c_{1}^{\beta|\kappa-1|\alpha|\delta|l-1} = 0,$$
Thus, we finally have
\[ (n - k - l) c_1^{[a] [\beta ]_l} = 0. \]

Thus, we finally have
\[ c_1^{[a] [\beta ]_l} = 0, \quad 3 \leq k, l < n, \quad k \neq l, \quad k + l \neq n, \]
\[ c_1^{[a] [\beta ]_{n-k}} = c_1^{(k)} c_1^{[a] [\beta ]_{n-k}}, \quad k, n - k \geq 4 \quad 2k \neq n. \]

**Equation (35)**
\[ (\delta_{\alpha_3 \beta} c_1^{\alpha_2 \alpha_3 [\beta]_l [\gamma]_l} + \delta_{\alpha_2 \beta} c_1^{\alpha_3 \alpha_1 [\beta]_l [\gamma]_l} + \delta_{\alpha_3 \beta} c_1^{\alpha_1 \alpha_2 [\beta]_l [\gamma]_l}) \lambda_{\beta} q_{\beta} [q_{\beta}]^{l-2} = 0. \]  
(A.18)

Acting on eq. (A.18) by the operator \( \lambda_{\alpha_3} \partial / \partial q_{\alpha_3} \) we obtain (after cancelation of \( (n - 2) \))
\[ c_1^{\alpha_1 \alpha_2 [\beta]_l [\gamma]_l} = 0. \]

**Equation (36)**
\[
\begin{aligned}
&3 C_{2m-3}^{m-2} c_1^{(3)} (m-1) c_1^{(m-1)} \lambda_{\beta} q_{\beta} [q_{\beta}]^{m-2} p_{\beta} [p_{\gamma}]^{m-1} + \\
&+ (-1)^m C_{m+1}^2 c_1^{(m-1)} \delta_{\alpha_3 \beta} + \text{cycle}(1, 2, 3) [q_{\beta}]^{m-1} \lambda_{\gamma} p_{\gamma} [p_{\gamma}]^{m-1} - \\
&- C_{m}^2 (c_1^{\alpha_1 \alpha_2 [\beta]_l [\gamma]_l} + \text{cycle}(1, 2, 3)) \lambda_{\beta} q_{\beta} [q_{\beta}]^{m-2} [p_{\gamma}]^{m} = 0.
\end{aligned}
\]  
(A.19)

Acting on eq. (A.19) by the operator \( \lambda_{\alpha_3} \partial / \partial q_{\alpha_3} \) one has:
\[
\begin{aligned}
&\left( (-1)^m 3 C_{2m-3}^{m-2} c_1^{(3)} + (m - 1) C_{m+1}^2 c_1^{(m-1)} \right) \varepsilon_{\alpha_1 \alpha_2 [\beta]_l [\gamma]_l} [q_{\beta}]^{m-2} [p_{\gamma}]^{m} - \\
&- m C_{m}^2 c_1^{\alpha_1 \alpha_2 [\beta]_l [\gamma]_l} [q_{\beta}]^{m-2} [p_{\gamma}]^{m} = 0,
\end{aligned}
\]
from where it follows
\[ c_1^{[a] [\beta]_m} = c_1^{(m)} c_1^{[a] [\beta]_m}, \quad c_1^{(m)} = \frac{1}{m C_{m}^2} \left((-1)^m 3 C_{2m-3}^{m-2} c_1^{(3)} + C_{m+1}^2 c_1^{(m-1)}\right). \]

**Equation (37), n = 2m.**
\[
\begin{aligned}
&3 C_{n-3}^{k-1} c_1^{(3)} \lambda_{\beta} q_{\beta} [q_{\beta}]^{k-1} p_{\beta} [p_{\gamma}]^{l-1} - \\
&- (-1)^k C_{n-k}^2 c_1^{(l)} \lambda_{\beta} q_{\beta} [q_{\beta}]^{k-1} [p_{\gamma}]^{l-1} - \\
&- C_{n-k}^2 c_1^{(k)} \lambda_{\beta} q_{\beta} [q_{\beta}]^{k-1} [p_{\gamma}]^{l-1} = 0.
\end{aligned}
\]  
(A.20)

Acting on eq. (A.20) by the operator \( \lambda_{\alpha_3} \partial / \partial q_{\alpha_3} \) we have:
\[
3 C_{n-3}^{k-1} c_1^{(3)} - (-1)^k k C_{n-k}^2 c_1^{(k)} + (-1)^k l C_{n-l}^2 c_1^{(l)} = 0.
\]  
(A.21)

Eq. (A.21) is simplified, if one rewrites it in terms of \( d_k \):
\[
d_k \equiv (-1)^k k! (n - k)! c_1^{(k)}, \quad d_{n-k} = -d_k,
\]
\[
d_3 + d_k - d_{k+1} = 0.
\]  
(A.22)

It follows from eq. (A.22)
\[ d_k = (k - 2) d_3. \]

Remembering that \( 0 = c_1^{(m)} = d_m \), we obtain
\[ d_3 = 0 \quad \implies \quad c_1^{(k)} = 0 \quad \forall k. \]
Appendix 2

In this Appendix we find the general form of the operators $T''$ generating similarity transformations, which leave invariant the Poisson bracket on the Grassman algebra $\mathcal{G}_C$. The operators $T''$ satisfy the condition (see (57))

$$T''[f_1, f_2]_{*0}(\xi) = [T'' f_1, T'' f_2]_{*0}(\xi).$$  \hspace{1cm} \text{(A.23)}

We use the momentum representation for the operator $T''$:

$$T'' = \sum_{k=0}^{n} t^{[\alpha]k}(\xi) [\partial_\alpha]^k, \quad \varepsilon(t^{[\alpha]k}) = k \, (\text{mod} \, 2),$$

In such terms eq. (A.23) reads:

$$\sum_{k=0}^{n} t^{[\alpha]k}(\xi) [\partial_\alpha]^k \left( \partial_\gamma f_1(\xi) \lambda_\gamma \partial_\gamma f_2(\xi) \right) =$$

$$= \left( \sum_{k_1=0}^{n} \partial_\gamma t^{[\alpha]k_1}(\xi) [\partial_\alpha]^{k_1} + \sum_{k_1=0}^{n} t^{[\alpha]k_1}(\xi) [\partial_\alpha]^{k_1} \partial_\gamma \right) f_1(\xi) \lambda_\gamma \times$$

$$\times \left( \sum_{k_2=0}^{n} \partial_\gamma t^{[\beta]k_2}(\xi) [\partial_\beta]^{k_2} + \sum_{k_2=0}^{n} t^{[\beta]k_2}(\xi) [\partial_\beta]^{k_2} \partial_\gamma \right) f_2(\xi).$$  \hspace{1cm} \text{(A.24)}

\textbf{Step 0} \quad The coefficient in eq. (A.24) in front of $f_1(\xi)$ gives:

$$\partial_\gamma t^0(\xi) \lambda_\gamma \partial_\gamma \left( \sum_{k=0}^{n} t^{[\beta]k}(\xi) [\partial_\beta]^{k} f_2(\xi) \right) = 0 \implies t^0(\xi) = t_0 = \text{const.}$$

From nonsingularity of the operator $T''$ it follows that $t_0 \neq 0$. Present the operator $T''$ as

$$T'' = T^{(\mu)} T^{(1)},$$

$$T^{(\mu)} f(\xi) = \frac{1}{\mu^2} f(\mu \xi),$$  \hspace{1cm} \text{(A.25)}

where $\mu = t_0^{1/2}$. It is easy to verify, that the similarity transformation generated by the operator $T^{(\mu)}$ (and by its inverse), leaves the $*_{0}$–commutator invariant. Hence, the operator $T^{(1)}$ and its coefficients (for which we retain previous notation $t^{[\alpha]k}$) satisfy eq. (A.23) and eq. (A.24), respectively. Furthermore, the operator $T^{(1)}$ has $t_0 = 1$.

\textbf{Step 1} \quad The coefficient in (A.24) in front of $\partial_\alpha f_1(\xi), \partial_\beta f_2(\xi)$ ($t_0 = 1$) gives:

$$\{ \eta^\alpha(\xi), \eta^\beta(\xi) \} = \lambda_\alpha \delta_\alpha \beta = \{ \xi^\alpha, \xi^\beta \},$$

where the notation $\eta^\alpha(\xi) \equiv t^\alpha(\xi) + \xi^\alpha$ is used. Thus, the transformation of generators $\xi^\alpha \rightarrow \eta^\alpha(\xi)$ is canonical. Let us present the operator $T''$ as

$$T'' = T^{(\mu)} T^{(c)} T^{(2)},$$

where the operator $T^{(\mu)}$ is defined by eq. (A.25) and the operator $T^{(c)}$ describes a canonical transformation of the generators: $T^{(c)} f(\xi) = f(\eta(\xi))$. It is easy to verify, that the similarity
transformation generated by the operator $T^{(c)}_η$, leaves the $*_0$–commutator invariant. Hence, the similarity transformation generated by the operator $T^{(2)}$ also leaves the $*_0$–commutator invariant. Besides the operator $T^{(2)}$ has $t^0(ξ) = 1$, $t^α(ξ) = 0$ (we again retain the previous notation $t^{(a)} k$ for the coefficients of the operator $T^{(2)}$). In what follows we consider eq. (A.24) as an equation for the coefficients of the operator $T^{(2)}$.

**Step 1'** The coefficients in (A.24) in front of $∂_α f_1(ξ), [∂_β]^k f_2(ξ), k ≥ 2$, give:

$$∂_α t^{[β]} k(ξ) = 0 \implies t^{[α]} k(ξ) = t^{[α]} k = \text{const.}$$

**Step 2** The coefficient in (A.24) in front of $∂_α ∂_β f_1(ξ), f_2(ξ) = \exp(ξ^α p_α)$, gives:

$$(δ_α γ t^{[γ]} k − 1 − δ_β γ t^{[γ]} k − 1)ʃ γ p_γ [p_γ]^{k−1} = 0.$$  

The general solution of this equation (see the solution of eq. (A.3)) is

$$t^{[α]} k = 0, \ 2 ≤ k ≤ n − 1, \ t^{[α]} n = t_n \varepsilon^{[α]} n,$$

$t_n \neq 0$ only at even $n$, from where it follows

$$T^{(2)} ≡ T^{(n)} = 1 + t_n \varepsilon^{[α]} n [∂_α]^n.$$  

Thus, the general form of the operators $T''$ is:

$$T'' = T^{(μ)} T^{(c)} T^{(n)}.$$  

**References**


