## F A S C I C U L I M A T H E M A T I C I <br> Nr 50

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ON THE DYNAMICS OF THE RECURSIVE SEQUENCE

$$
x_{n+1}=\frac{\alpha x_{n-1}}{\beta+\gamma \sum_{k=1}^{t} x_{n-2 k}^{p} \prod_{k=1}^{t} x_{n-2 k}^{q}} *
$$

Abstract. In this paper, we investigate the global behavior of the difference equation

$$
x_{n+1}=\frac{\alpha x_{n-1}}{\beta+\gamma \sum_{k=1}^{t} x_{n-2 k}^{p} \prod_{k=1}^{t} x_{n-2 k}^{q}}, \quad n=0,1, \ldots,
$$

where $\beta$ is a positive parameter and $\alpha, \gamma$ are non-negative parameters and non-negative initial conditions.
KEY words: difference equations, recursive sequences, oscillation, global asymptotic behavior, period two solutions, semicycles.
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## 1. Introduction

Consider the higher-order difference equation

$$
\begin{equation*}
x_{n+1}=\frac{\alpha x_{n-1}}{\beta+\gamma \sum_{k=1}^{t} x_{n-2 k}^{p} \prod_{k=1}^{t} x_{n-2 k}^{q}}, \quad n=0,1, \ldots \tag{1}
\end{equation*}
$$

where the parameters, $\beta$ is positive and $\alpha, \gamma$ are non-negative real numbers and the initial conditions $x_{-2 t}, \ldots, x_{-2}, x_{-1}$ and $x_{0}$ are non-negative real numbers such that

$$
0<\beta+\gamma \sum_{k=1}^{t} x_{n-2 k}^{p} \prod_{k=1}^{t} x_{n-2 k}^{q}, \quad n=0,1, \ldots
$$

and if $\alpha=0$ the equation $x_{n+1}=0$ is trivial, if $\gamma=0$ the equation $x_{n+1}=$ $\frac{\alpha}{\beta} x_{n-1}$ is linear. We assume that all parameters in equations are positive.

[^0]We investigate the global asymptotic behavior and the periodic character of the solutions of the difference (1), by generalizing the results due to El-Owaidy et al. [1] corresponding to the difference equation

$$
x_{n+1}=\frac{\alpha x_{n-1}}{\beta+\gamma x_{n-2}^{p}}, \quad n=0,1, \ldots
$$

where the parameters $\alpha, \beta$ and $\gamma$ are positive real numbers and the initial conditions $x_{-2}, x_{-1}$ and $x_{0}$ are arbitrary non-negative real numbers. Similar recursive sequences were studied previously; for example, see Refs. [1-22].

We need the following definitions and theorem [23]:
Definition 1. Let $I$ be an interval of the real numbers and let $f$ : $I^{2 t+1} \longrightarrow I$ be a continuously differantiable function. Consider the difference equation

$$
\begin{equation*}
x_{n+1}=f\left(x_{n}, x_{n-1}, x_{n-2}, \ldots, x_{n-2 t}\right), \quad n=0,1, \ldots . \tag{2}
\end{equation*}
$$

with $x_{-i}$ for $i=0,1, \ldots, 2 t \in I$. Let $\bar{x}$ be the equilibrium point of (2). The linearized equation of (2) about the equilibrium point $\bar{x}$ is

$$
\begin{equation*}
y_{n+1}=c_{1} y_{n}+c_{2} y_{n-1}+\ldots+c_{2 t+1} y_{n-2 t}, \quad n=0,1, \ldots \tag{3}
\end{equation*}
$$

where

$$
\begin{aligned}
c_{1} & =\frac{\partial f}{\partial x_{n}}(\bar{x}, \bar{x}, \ldots, \bar{x}), \\
c_{2} & =\frac{\partial f}{\partial x_{n-1}}(\bar{x}, \bar{x}, \ldots, \bar{x}), \\
\vdots & \\
c_{2 t+1} & =\frac{\partial f}{\partial x_{n-2 t}}(\bar{x}, \bar{x}, \ldots, \bar{x}) .
\end{aligned}
$$

The characteristic equation of (3) is

$$
\begin{equation*}
\lambda^{2 t+1}-c_{1} \lambda^{2 t}-\ldots-c_{2 t-1} \lambda^{2}-c_{2 t} \lambda-c_{2 t+1}=0 \tag{4}
\end{equation*}
$$

Definition 2. Let $\bar{x}$ be an equilibrium point of (2).
(a) The equilibrium $\bar{x}$ is called locally stable if for every $\varepsilon>0$, there exists $\delta>0$ such that if $x_{0}, \ldots, x_{-2 t} \in I$ and $\left|x_{0}-\bar{x}\right|+\cdots+\left|x_{-2 t}-\bar{x}\right|<\delta$, then $\left|x_{n}-\bar{x}\right|<\varepsilon$, for all $n \geq-2 t$.
(b) The equilibrium $\bar{x}$ is called locally asymptotically stable if it is locally stable and if there exists $\gamma>0$ such that if $x_{0}, \ldots, x_{-2 t} \in I$ and $\left|x_{0}-\bar{x}\right|+$ $\cdots+\left|x_{-2 t}-\bar{x}\right|<\gamma$, then $\lim _{n \rightarrow \infty} x_{n}=\bar{x}$.
(c) The equilibrium $\bar{x}$ is called global attractor if for every $x_{0}, \ldots, x_{-2 t} \in$ $I$ we have $\lim _{n \rightarrow \infty} x_{n}=\bar{x}$.
(d) The equilibrium $\bar{x}$ is called globally asymptotically stable if it is locally stable and is a global attractor.

Definition 3. A positive semicycle of $\left\{x_{n}\right\}_{n=-2 t}^{\infty}$ of (2) consists of a 'string' of terms $\left\{x_{l}, x_{l+1}, \ldots, x_{m}\right\}$, all greater than or equal to $\bar{x}$, with $l \geq$ $-2 t$ and $m<\infty$ and such that either $l=-2 t$ or $l>-2 t$ and $x_{l-1}<\bar{x}$ and either $m=\infty$ or $m<\infty$ and $x_{m+1}<\bar{x}$.

A negative semicycle of $\left\{x_{n}\right\}_{n=-2 t}^{\infty}$ of (2) consists of a 'string' of terms $\left\{x_{l}, x_{l+1}, \ldots, x_{m}\right\}$ all less than $\bar{x}$, with $l \geq-2 t$ and $m<\infty$ and such that either $l=-2 t$ or $l>-2 t$ and $x_{l-1} \geq \bar{x}$ and either $m=\infty$ or $m<\infty$ and $x_{m+1} \geq \bar{x}$.

Definition 4. A solution $\left\{x_{n}\right\}_{n=-2 t}^{\infty}$ of (2) is called nonoscillatory if there exists $N \geq-2 t$ such that either

$$
x_{n}>\bar{x} \text { or } x_{n}<\bar{x} \quad \text { for } \forall n \geq N
$$

and it is called oscillatory if it is not nonoscillatory.
Theorem 1. (i) If all roots of (4) have absolute values less than one, then the equilibrium point $\bar{x}$ of (2) is locally asymptotically stable.
(ii) If at least one of the roots of (4) has absolute value greater than one, then the equilibrium point $\bar{x}$ of (2) is unstable.
(iii) The equilibrium point $\bar{x}$ of (2) is called saddle point if (4) has roots both inside and outside the unit disk.

## 2. Dynamics of equation (1)

In this section, we investigate the dynamics of (1) under the assumptions that all parameters in the equation are positive and the initial conditions are non-negative.

The change of variables $x_{n}=(\beta / \gamma)^{1 / q t+p} y_{n}$ reduces (1) to the difference equation

$$
\begin{equation*}
y_{n+1}=\frac{r y_{n-1}}{1+y_{n-2}^{q+p} y_{n-4}^{q} \ldots y_{n-2 t}^{q}+y_{n-2}^{q} y_{n-4}^{q+p} y_{n-2 t}^{q}+\ldots+y_{n-2}^{q} y_{n-4}^{q} \ldots y_{n-2 t}^{q+p}} \tag{5}
\end{equation*}
$$

where $r=\alpha / \beta>0$ and $n=0,1, \ldots$.
Note that $\bar{y}_{1}=0$ is always an equilibrium point of (5). When $r>1$, (5) also possesses the unique positive equilibrium $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$.

Theorem 2. The following statements are true.
(i) If $r<1$, then the equilibrium point $\bar{y}_{1}=0$ of (5) is locally asymptotically stable.
(ii) If $r>1$, then the equilibrium point $\bar{y}_{1}=0$ of (5) is a saddle point.
(iii) When $r>1$, then the positive equilibrium point $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$ of (5) is unstable.

Proof. The linearized equation of (5) about the equilibrium point $\bar{y}_{1}=0$ is

$$
z_{n+1}=r z_{n-1}, \quad n=0,1, \ldots
$$

so, the characteristic equation of (5) about the equilibrium point $\bar{y}_{1}=0$ is

$$
\lambda^{2 t+1}-r \lambda^{2 t-1}=0
$$

hence the proof of $(i)$ and (ii) follows Theorem 1.
For (iii) we assume that $r>1$; then the linearized equation of (5) about the equilibrium point $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$ has the form

$$
\begin{aligned}
z_{n+1}= & z_{n-1}-\frac{(q t+p)}{t} \frac{(r-1)}{r} z_{n-2}-\frac{(q t+p)}{t} \frac{(r-1)}{r} z_{n-4} \\
& -\ldots-\frac{(q t+p)}{t} \frac{(r-1)}{r} z_{n-2 t}=0
\end{aligned}
$$

where $n=0,1, \ldots$. So the characteristic equation of (5) about the equilibrium point $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$ is

$$
\begin{equation*}
\lambda^{2 t+1}-\lambda^{2 t-1}+\frac{(q t+p)}{t} \frac{(r-1)}{r} \lambda^{2 t-2}+\ldots+\frac{(q t+p)}{t} \frac{(r-1)}{r}=0 . \tag{6}
\end{equation*}
$$

It is clear that (6) has a root in the interval $(-\infty,-1)$ and so $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$ is an unstable equilibrium point. This completes the proof.

Theorem 3. Assume that $r>1$. Let $\left\{y_{n}\right\}_{n=-2 t}^{\infty}$ be a solution of (5) such that

$$
\begin{equation*}
y_{-2 t}, \ldots y_{0} \geq \bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}, \quad y_{-2 t+1}, \ldots, y_{-1}<\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p} \tag{7}
\end{equation*}
$$

or
(8) $y_{-2 t}, \ldots y_{0}<\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}, \quad y_{-2 t+1}, \ldots, y_{-1} \geq \bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$

Then $\left\{y_{n}\right\}_{n=-2 t}^{\infty}$ oscillates about $\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p}$ with semicycle of length 1 .
Proof. Assume that (7) holds. (The case where (8) holds is similar and will be omitted.) Then,

$$
\begin{aligned}
y_{1} & =\frac{r y_{-1}}{1+y_{-2}^{q+p} y_{-4}^{q} \ldots y_{-2 t}^{q}+y_{-2}^{q} y_{-4}^{q+p} \ldots y_{-2 t}^{q}+y_{-2}^{q} y_{-4}^{q} \ldots y_{-2 t}^{q+p}} \\
& <\frac{r \bar{y}_{2}}{1+t \overline{y_{2}^{q t+p}}}=\frac{r \bar{y}_{2}}{1+r-1}=\bar{y}_{2} \\
& <\bar{y}_{2}
\end{aligned}
$$

and

$$
\begin{aligned}
y_{2} & =\frac{r y_{0}}{1+y_{-1}^{q+p} y_{-3}^{q} \ldots y_{-2 t+1}^{q}+y_{-1}^{q} y_{-3}^{q+p} \ldots y_{-2 t+1}^{q}+y_{-1}^{q} y_{-3}^{q} \ldots y_{-2 t+1}^{q+p}} \\
& \geq \frac{r \bar{y}_{2}}{1+t \bar{y}_{2}^{q t+p}}=\frac{r \bar{y}_{2}}{1+r-1}=\bar{y}_{2}, \\
& \geq \bar{y}_{2}
\end{aligned}
$$

then the proof follows by induction.
Theorem 4. Assume that $r<1$; then the equilibrium point $\bar{y}_{1}=0$ of (5) is globally asymptotically stable.

Proof. We know by Theorem 2 that the equilibrium point $\bar{y}_{1}=0$ of (5) is locally asymptotically stable. So let $\left\{y_{n}\right\}_{n=-2 t}^{\infty}$ be a solution of (5). It suffices to show that

$$
\lim _{n \rightarrow \infty} y_{n}=0
$$

Since

$$
\begin{aligned}
y_{n+1} & =\frac{r y_{n-1}}{1+y_{n-2}^{q+p} y_{n-4}^{q} \cdots y_{n-2 t}^{q}+y_{n-2}^{q} y_{n-4}^{q+\ldots} y_{n-2 t}^{q}+y_{n-2}^{q} y_{n-4}^{q} \ldots y_{n-2 t}^{q+p}} \\
y_{2 n-1} & <r^{n} y_{-1} \text { and } y_{2 n}<r^{n} y_{0}
\end{aligned}
$$

we have

$$
\lim _{n \rightarrow \infty} y_{n}=0
$$

This completes the proof.
Theorem 5. Assume that $r=1$; then (5) possesses the prime period 2 solutions

$$
\begin{equation*}
\ldots, \Phi, \Psi, \Phi, \Psi, \ldots \tag{9}
\end{equation*}
$$

with $\Phi>0$. Furthermore, every solution of (5) converges to a period 2 solution (9) with $\Phi \geq 0$.

Proof. Let

$$
\ldots, \Phi, \Psi, \Phi, \Psi, \ldots
$$

be a period two solution of (5). Then

$$
\Phi=\frac{r \Phi}{1+\Psi^{q+p} \Psi^{q} \ldots \Psi^{q}+\Psi^{q} \Psi^{q+p} \ldots \Psi^{q}+\Psi^{q} \Psi^{q} \ldots \Psi^{q+p}}
$$

and

$$
\Psi=\frac{r \Psi}{1+\Phi^{q+p} \Phi^{q} \ldots \Phi^{q}+\Phi^{q} \Phi^{q+p} \ldots \Phi^{q}+\Phi^{q} \Phi^{q} \ldots \Phi^{q+p}} .
$$

So

$$
t \Phi \Psi=\frac{(\Phi-\Psi)(r-1)}{\Psi^{q t+p-1}-\Phi^{q t+p-1}} \geq 0
$$

which implies that $r-1 \leq 0$.
If $r<1$, then this implies that $\Phi<0$ or $\Psi<0$, which is impossible, so $r=1$. If $r>1$, then this implies that $\Phi=\Psi=\left(\frac{r-1}{t}\right)^{1 / q t+p} \neq 0$, which contradicts that
$\Phi \neq \Psi$, so $r=1$. To complete the proof, assume that $r=1$ and let $\left\{y_{n}\right\}_{n=-2 t}^{\infty}$ be a solution of (5); then

$$
\begin{aligned}
& y_{n+1}-y_{n-1} \\
& =\frac{-y_{n-1} y_{n-2}^{q+p} y_{n-4}^{q} \ldots y_{n-2 t}^{q}-y_{n-1} y_{n-2}^{q} y_{n-4}^{q+p} \ldots, y_{n-2 t}^{q}-y_{n-1} y_{n-2}^{q} y_{n-4}^{q} \ldots y_{n-2 t}^{q+p}}{1+y_{n-2}^{q+p} y_{n-4}^{q} \ldots y_{n-2 t}^{q}+y_{n-2}^{q} y_{n-4}^{q+p} \ldots, y_{n-2 t}^{q}+y_{n-2}^{q} y_{n-4}^{q} \ldots y_{n-2 t}^{q+p}} \\
& y_{n+1}-y_{n-1} \leq 0 .
\end{aligned}
$$

So, the even terms of this solution decrease to a limit (say $\Phi \geq 0$ ) and the odd terms decrease to a limit ( say $\Psi \geq 0$ ). Thus

$$
\Phi=\frac{\Phi}{1+\Psi^{q+p} \Psi^{q} \ldots \Psi^{q}+\Psi^{q} \Psi^{q+p} \ldots \Psi^{q}+\Psi^{q} \Psi^{q} \ldots \Psi^{q+p}}
$$

and

$$
\Psi=\frac{\Psi}{1+\Phi^{q+p} \Phi^{q} \ldots \Phi^{q}+\Phi^{q} \Phi^{q+p} \ldots \Phi^{q}+\Phi^{q} \Phi^{q} \ldots \Phi^{q+p}},
$$

which implies that

$$
t \Phi \Psi^{q t+p}=0 \text { and } t \Psi \Phi^{q t+p}=0 .
$$

This completes the proof.
Theorem 6. Assume that $r>1$; then (5) possesses an unbounded solution.
Proof. From Theorem 3, we can assume without loss of generality that the solution $\left\{y_{n}\right\}_{n=-2 t}^{\infty}$ of (5) is such that

$$
y_{2 n-1}<\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p} \quad \text { and } \quad y_{2 n}>\bar{y}_{2}=\left(\frac{r-1}{t}\right)^{1 / q t+p},
$$

for $n \geq 0$. Then

$$
\begin{aligned}
& y_{2 n+2} \\
& =\frac{r y_{2 n}}{1+y_{2 n-1}^{q+p} y_{2 n-3}^{q} \ldots y_{2 n-2 t+1}^{q}+y_{2 n-1}^{q} y_{2 n-3}^{q+p} \ldots y_{2 n-2 t+1}^{q}+y_{2 n-1}^{q} y_{2 n-3}^{q} \ldots y_{2 n-2 t+1}^{q+p}} \\
& y_{2 n+2}>\frac{r y_{2 n}}{1+(r-1)}=y_{2 n}
\end{aligned}
$$

and
$y_{2 n+3}=\frac{r y_{2 n+1}}{1+y_{2 n}^{q+p} y_{2 n-2}^{q} \ldots y_{2 n-2 t+2}^{q}+y_{2 n}^{q} y_{2 n-2}^{q+p} \ldots y_{2 n-2 t+2}^{q}+y_{2 n}^{q} y_{2 n-2}^{q} \ldots y_{2 n-2 t+2}^{q+p}}$
$y_{2 n+3}<\frac{r y_{2 n+1}}{1+(r-1)}=y_{2 n+1}$
from which it follows that

$$
\lim n \rightarrow \infty y_{2 n}=\infty \quad \text { and } \quad \lim _{n \rightarrow \infty} y_{2 n+1}=0
$$

Then, the proof is complete.

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