

# Involute Curves Of Timelike Biharmonic Reeb Curves $(LCS)_3$ - Manifolds

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**Abstract:** In this paper, we study involute timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold. We characterize curvatures of timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold. We obtain parametric equation involute curves of the timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold.

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## 1. Introduction

A smooth map  $\phi : N \rightarrow M$  is said to be biharmonic if it is a critical point of the bienergy functional:

$$E_2(\phi) = \int_N \frac{1}{2} |\mathcal{T}(\phi)|^2 dv_h,$$

where  $\mathcal{T}(\phi) := \text{tr} \nabla^\phi d\phi$  is the tension field of  $\phi$

The Euler–Lagrange equation of the bienergy is given by  $\mathcal{T}_2(\phi) = 0$ . Here the section  $\mathcal{T}_2(\phi)$  is defined by

$$\mathcal{T}_2(\phi) = -\Delta_\phi \mathcal{T}(\phi) + \text{tr} R(\mathcal{T}(\phi), d\phi) d\phi, \quad (1)$$

and called the bitension field of  $\phi$ . Non-harmonic biharmonic maps are called proper biharmonic maps.

In this paper, we study timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold. We characterize curvatures of timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold. Several

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interesting results on a  $(LCS)_3$ -manifold are obtained in terms of timelike biharmonic Reeb curves. Finally, we obtain parametric equation of the timelike biharmonic Reeb curves in  $(LCS)_3$ -manifold.

## 2. $(LCS)_3$ -Manifolds

**Definition 2.1.** In a Lorentzian manifold  $(M, g)$  a vector field  $P$  defined by

$$g(X, P) = A(X)$$

for any  $X \in \chi(M)$  is said to be a concircular vector field if

$$(\nabla_X A)(Y) = \alpha\{g(X, Y) + \omega(X)A(Y)\},$$

where  $\alpha$  is a non-zero scalar and  $\omega$  is a closed 1-form [15].

Let  $M^3$  be a Lorentzian manifold admitting a unit timelike concircular vector field  $\xi$ , called the characteristic vector field of the manifold. Then we have

$$g(\xi, \xi) = -1. \quad (2)$$

Since  $\xi$  is a unit concircular vector field, it follows that there exists a non-zero 1-form  $\eta$  such that for

$$g(X, \xi) = \eta(X), \quad (3)$$

the equation of the following form holds

$$(\nabla_X \eta)(Y) = \alpha\{g(X, Y) + \eta(X)\eta(Y)\} \quad (\alpha \neq 0)$$

for all vector fields  $X, Y$ , where  $\nabla$  denotes the operator of covariant differentiation with respect to the Lorentzian metric  $g$  and  $\alpha$  is a non-zero scalar function satisfies

$$\nabla_X \alpha = (X\alpha) = d\alpha(X) = \rho\eta(X),$$

$\rho$  being a certain scalar function given by  $\rho = -(\xi\alpha)$ . If we put

$$\phi X = \frac{1}{\alpha} \nabla_X \xi,$$

then from (2.3) and (2.5) we have

$$\phi X = X + \eta(X)\xi,$$

from which it follows that  $\phi$  is a symmetric (1) tensor and called the structure tensor of the manifold. Thus the Lorentzian manifold  $M^3$  together with the unit timelike concircular vector field  $\xi$ , its associated 1-form  $\eta$  and (1) tensor field  $\phi$  is said to be a Lorentzian

conircular structure manifold (briefly  $(LCS)_3$ - manifold) [16]. In a  $(LCS)_3$ -manifold, the following relations hold [15]:

$$\begin{aligned}\eta(\xi) &= -1, \\ \phi\xi &= 0, \\ \eta(\phi X) &= 0,\end{aligned}\tag{4}$$

$$g(\phi X, \phi Y) = g(X, Y) + \eta(X)\eta(Y),\tag{5}$$

$$\eta(R(X, Y)Z) = (\alpha^2 - \rho)[g(Y, Z)\eta(X) - g(X, Z)\eta(Y)],\tag{6}$$

$$S(X, \xi) = (n - 1)(\alpha^2 - \rho)\eta(X),\tag{7}$$

$$R(X, Y)\xi = (\alpha^2 - \rho)[\eta(Y)X - \eta(X)Y],\tag{8}$$

$$(\nabla_X\phi)(Y) = \alpha[g(X, Y)\xi + 2\eta(X)\eta(Y)\xi + \eta(Y)X],\tag{9}$$

for all vector fields  $X, Y, Z$ , where  $R, S$  denote respectively the curvature tensor and the Ricci tensor of the manifold.

### 3. Biharmonic Reeb Curves in the $(LCS)_3$ -Manifold

Let  $\gamma$  be a timelike curve on the  $(LCS)_3$ - manifold parametrized by arc length. Let  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  be the Frenet frame fields tangent to the 3-dimensional Kenmotsu manifold along  $\gamma$  defined as follows:

$\mathbf{T}$  is the unit vector field  $\gamma'$  tangent to  $\gamma$ ,  $\mathbf{N}$  is the unit vector field in the direction of  $\nabla_{\mathbf{T}}\mathbf{T}$  (normal to  $\gamma$ ), and  $\mathbf{B}$  is chosen so that  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

$$\begin{aligned}\nabla_{\mathbf{T}}\mathbf{T} &= \kappa\mathbf{N}, \\ \nabla_{\mathbf{T}}\mathbf{N} &= \kappa\mathbf{T} + \tau\mathbf{B}, \\ \nabla_{\mathbf{T}}\mathbf{B} &= -\tau\mathbf{N},\end{aligned}\tag{10}$$

where  $\kappa$  is the curvature of  $\gamma$  and  $\tau$  its torsion and

$$\begin{aligned}g(\mathbf{T}, \mathbf{T}) &= -1, \quad g(\mathbf{N}, \mathbf{N}) = 1, \quad g(\mathbf{B}, \mathbf{B}) = 1, \\ g(\mathbf{T}, \mathbf{N}) &= g(\mathbf{T}, \mathbf{B}) = g(\mathbf{N}, \mathbf{B}) = 0.\end{aligned}\tag{11}$$

**Theorem 3.1.** [13] *Let  $(M, \phi, \xi, \eta, g)$  be an 3-dimensional  $(LCS)_3$ - manifold and unit vector field  $X$  orthogonal to the Reeb vector field  $\xi$ . Then,*

$$R(\xi, X)\xi = (\alpha^2 - \rho)X,\tag{12}$$

$$R(X, \xi)X = -(\alpha^2 - \rho)\xi.\tag{13}$$

**Theorem 3.2.** [13]  *$\gamma$  is a timelike biharmonic Reeb curve which are either tangent or normal to the Reeb vector field in  $(LCS)_3$ - manifold then*

$$\begin{aligned}\kappa &= \text{constant} \neq 0, \\ \kappa^2 - \tau^2 &= \alpha^2 - \rho, \\ \tau &= \text{constant}.\end{aligned}\tag{14}$$

**Proof.** Using (12) and Frenet formulas (10), we have (16).

**Corollary 3.3.** [13] *If  $\gamma$  is a timelike biharmonic Reeb curve which orthogonal to the Reeb vector field  $\xi$  in  $(LCS)_3$ -manifold, then  $\gamma$  is a helix.*

We consider the three-dimensional  $(LCS)_3$ -manifold

$$\mathbb{M} = \left\{ (x, y, z) \in \mathbb{R}^3 : x \neq \pm\sqrt{2}z^2, x \neq 0, z \neq 0 \right\},$$

where  $(x, y, z)$  are the standard coordinates in  $\mathbb{R}^3$ . Let  $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$  be linearly independent global frame on  $\mathbb{M}$  given by [17]

$$\mathbf{e}_1 = z \frac{\partial}{\partial x}, \quad \mathbf{e}_2 = zx \frac{\partial}{\partial y}, \quad \mathbf{e}_3 = \frac{\partial}{\partial z}. \quad (15)$$

Let  $g$  be the Riemannian metric defined by

$$\begin{aligned} g(\mathbf{e}_1, \mathbf{e}_1) &= g(\mathbf{e}_2, \mathbf{e}_2) = 1, & g(\mathbf{e}_3, \mathbf{e}_3) &= -1, \\ g(\mathbf{e}_1, \mathbf{e}_2) &= g(\mathbf{e}_2, \mathbf{e}_3) = g(\mathbf{e}_1, \mathbf{e}_3) = 0. \end{aligned}$$

Let  $\eta$  be the 1-form defined by

$$\eta(U) = g(U, \mathbf{e}_3) \text{ for any } U \in \chi(\mathbb{M}).$$

Let  $\phi$  be the (1) tensor field defined by

$$\phi(\mathbf{e}_1) = \mathbf{e}_1, \quad \phi(\mathbf{e}_2) = \mathbf{e}_2, \quad \phi(\mathbf{e}_3) = 0.$$

Then using the linearity of  $\phi$  and  $g$  we have

$$\begin{aligned} \eta(\mathbf{e}_3) &= -1, \\ \phi^2(U) &= U - \eta(U)\mathbf{e}_3, \\ g(\phi U, \phi W) &= g(U, W) - \eta(U)\eta(W), \end{aligned}$$

for any  $U, W \in \chi(\mathbb{M})$ .

Let  $\nabla$  be the Levi-Civita connection with respect to  $g$ . Then, we have

$$[\mathbf{e}_1, \mathbf{e}_2] = \frac{z}{x}\mathbf{e}_2, \quad [\mathbf{e}_1, \mathbf{e}_3] = -\frac{1}{z}\mathbf{e}_1, \quad [\mathbf{e}_2, \mathbf{e}_3] = -\frac{1}{z}\mathbf{e}_2.$$

Taking  $\mathbf{e}_3 = \xi$  and using the Koszul's formula, we obtain

$$\begin{aligned} \nabla_{\mathbf{e}_1}\mathbf{e}_1 &= -\frac{1}{z}\mathbf{e}_3, & \nabla_{\mathbf{e}_1}\mathbf{e}_2 &= 0, & \nabla_{\mathbf{e}_1}\mathbf{e}_3 &= -\frac{1}{z}\mathbf{e}_1, \\ \nabla_{\mathbf{e}_2}\mathbf{e}_1 &= -\frac{z}{x}\mathbf{e}_2, & \nabla_{\mathbf{e}_2}\mathbf{e}_2 &= \frac{z}{x}\mathbf{e}_1 - \frac{1}{z}\mathbf{e}_3, & \nabla_{\mathbf{e}_2}\mathbf{e}_3 &= -\frac{1}{z}\mathbf{e}_2, \\ \nabla_{\mathbf{e}_3}\mathbf{e}_1 &= 0, & \nabla_{\mathbf{e}_3}\mathbf{e}_2 &= 0, & \nabla_{\mathbf{e}_3}\mathbf{e}_3 &= 0. \end{aligned} \quad (16)$$

Moreover we put

$$R_{ijk} = R(\mathbf{e}_i, \mathbf{e}_j)\mathbf{e}_k,$$

where the indices  $i, j, k$  take the values 1, 2 and 3.

$$R_{122} = -[2(\frac{z}{x})^2 - \frac{1}{z^2}]\mathbf{e}_1, \quad R_{133} = -\frac{2}{z^2}\mathbf{e}_1, \quad R_{233} = -\frac{2}{z^2}\mathbf{e}_2.$$

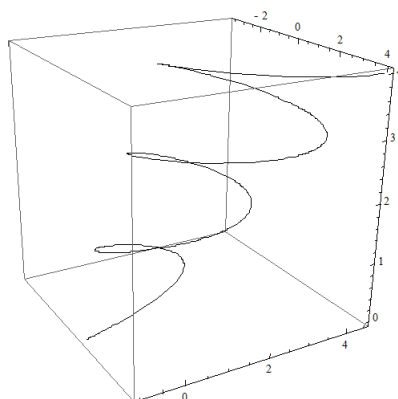
**Theorem 3.4.** [13] Let  $\gamma : I \rightarrow \mathbb{M}$  be a unit speed timelike biharmonic Reeb curve which orthogonal to the Reeb vector field  $\xi$  in  $(LCS)_3$ - manifold  $\mathbb{M}$ . Then, the parametric equations of  $\gamma$  are

$$\begin{aligned} x(s) &= \frac{\sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \\ &\quad + \frac{1}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) + c_2, \\ y(s) &= \frac{1}{12\wp^4} (\sinh^2 \varphi - 1) (2\wp^3 s(3c_1^2 + 3\sinh \varphi c_1 s + \sinh^2 \varphi s^2) \\ &\quad - 6\wp \sinh \varphi (\sinh \varphi s + c_1) \cos[2(\wp s + \sigma)] \\ &\quad - 3(\wp^2 c_1^2 + 2\wp^2 \sinh \varphi c_1 s + \sinh^2 \varphi (-1 + \wp^2 s^2)) \sin[2(\wp s + \sigma)]) \\ &\quad + \frac{c_2 \sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \\ &\quad - \frac{c_2}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) + c_3, \end{aligned} \quad (17)$$

$$z(s) = \sinh \varphi s + c_1,$$

where  $\sigma, c_1, c_2, c_3$  are constants of integration.

Using Mathematica in above theorem, we have



#### 4. Involute Curves of Biharmonic Reeb Curves in the $(LCS)_3$ -Manifold

**Definition 4.1.** Let unit speed timelike curve  $\gamma : I \rightarrow \mathbb{M}$  and the curve  $\Theta : I \rightarrow \mathbb{M}$  be given. For  $\forall s \in I$ , then the curve  $\Theta$  is called the involute of the curve  $\gamma$ , if the tangent

at the point  $\gamma(s)$  to the curve  $\gamma$  passes through the tangent at the point  $\Theta(s)$  to the curve  $\Theta$  and

$$g(\mathbf{T}^*(s), \mathbf{T}(s)) = 0. \quad (18)$$

Let the Frenet-Serret frames of the curves  $\gamma$  and  $\zeta$  be  $\{\mathbf{T}, \mathbf{N}, \mathbf{B}\}$  and  $\{\mathbf{T}^*, \mathbf{N}^*, \mathbf{B}^*\}$ , respectively.

**Theorem 4.2.** *Let  $\gamma : I \rightarrow \mathbb{M}$  be a unit speed timelike biharmonic Reeb curve which are either tangent or normal to the Reeb vector field in  $(LCS)_3$ - manifold  $\mathbb{M}$ ,  $\Theta$  its involute curve. Then, the parametric equations of  $\Theta$  are*

$$\begin{aligned} x(s) &= \frac{\sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \\ &\quad + \frac{1}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \\ &\quad + (C - s) (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) + c_2, \\ y(s) &= \frac{1}{12\wp^4} (\sinh^2 \varphi - 1) (2\wp^3 s(3c_1^2 + 3 \sinh \varphi c_1 s + \sinh^2 \varphi s^2) \\ &\quad - 6\wp \sinh \varphi (\sinh \varphi s + c_1) \cos[2(\wp s + \sigma)] \\ &\quad - 3(\wp^2 c_1^2 + 2\wp^2 \sinh \varphi c_1 s + \sinh^2 \varphi (-1 + \wp^2 s^2)) \sin[2(\wp s + \sigma)]) \\ &\quad + \frac{c_2 \sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \\ &\quad - \frac{c_2}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \\ &\quad + (C - s) (\sinh \varphi s + c_1) \left( \frac{\sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \right. \\ &\quad \left. + \frac{1}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \right. \\ &\quad \left. + c_2 \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) + c_3, \right. \\ z(s) &= C \sinh \varphi + c_1, \end{aligned} \quad (19)$$

where  $\sigma, c_1, c_2, c_3$  are constants of integration.

**Proof.** We assume that  $\gamma : I \rightarrow \text{Heis}^3$  be a unit speed timelike biharmonic curve and  $\Theta$  its involute curve on  $\text{Heis}^3$ . We find that the parametric equations of  $\Theta$ .

The involute curve of timelike biharmonic curve may be given as

$$\Theta(s) = \gamma(s) + u(s)\mathbf{T}(s). \quad (20)$$

From (4.3), then we have

$$\Theta'(s) = (1 + u'(s)) \mathbf{T}(s) + u(s)\kappa(s) \mathbf{N}(s). \quad (21)$$

Since the curve  $\Theta$  is involute of the curve  $\gamma$ ,  $g(\mathbf{T}^*(s), \mathbf{T}(s)) = 0$ . Then, we get

$$1 + u'(s) = 0 \text{ or } u(s) = C - s, \quad (22)$$

where  $C$  is constant of integration.

Substituting (22) into (20), we get

$$\Theta(s) = \gamma(s) + (C - s) \mathbf{T}(s). \quad (23)$$

On the other hand, (19) and (23), imply

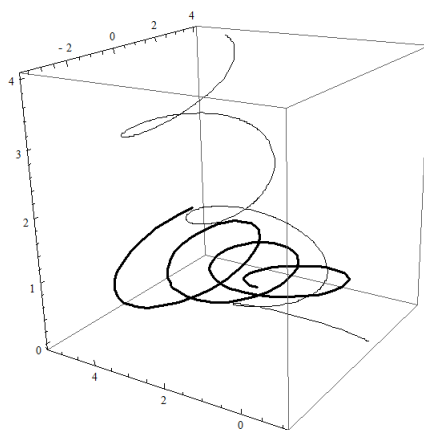
$$\mathbf{T} = \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \mathbf{e}_1 + \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \mathbf{e}_2 + \sinh \varphi \mathbf{e}_3. \quad (24)$$

Therefore, from (23) and (24) we have

$$\begin{aligned} \mathbf{T} = & ((\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma), \\ & (\sinh \varphi s + c_1) \left( \frac{\sinh \varphi}{\wp^2} \sqrt{\sinh^2 \varphi - 1} \cos(\wp s + \sigma) \right. \\ & \left. + \frac{1}{\wp} (\sinh \varphi s + c_1) \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma) \right. \\ & \left. + c_2) \sqrt{\sinh^2 \varphi - 1} \sin(\wp s + \sigma), \sinh \varphi). \end{aligned} \quad (25)$$

If we substitute (17) and (25) into (23), we have (19). This concludes the proof of Theorem.

We show that  $\gamma$  and  $\Theta$  in terms of Mathematica as follows:



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