

1965

The effect of molecular weight on the compressibility of bisphenol a polycarbonate

Allan L. Holmstrom
Lehigh University

Follow this and additional works at: <https://preserve.lehigh.edu/etd>

 Part of the [Chemical Engineering Commons](#)

Recommended Citation

Holmstrom, Allan L., "The effect of molecular weight on the compressibility of bisphenol a polycarbonate" (1965). *Theses and Dissertations*. 5043.
<https://preserve.lehigh.edu/etd/5043>

This Thesis is brought to you for free and open access by Lehigh Preserve. It has been accepted for inclusion in Theses and Dissertations by an authorized administrator of Lehigh Preserve. For more information, please contact preserve@lehigh.edu.

June 4, 1965

Dear Dr. Moyle,

Indeed our letters did cross in the mail. I presume you understood my anxiety as the semester drew to a close. I did, however, finish on time and as of about ten minutes ago was cleared financially for Graduation.

I do want to personally thank you for all you have done for me, you've been most kind and I appreciate it.

Enclosed is the copy of the report you requested. In a seminar I gave for the Materials Research Group, the only question that came up of consequence was whether the strain on the sample was isotropic. It would be isotropic I believe if the Poisson ratio were 0.5, however, I believe it has a value more like 0.25. It may be more appropriate to therefore call the bulk modulus the "apparent" bulk modulus to keep the terminology clean. - The solution to the above I guess is to suspend a spherical sample in an incompressible fluid and make compressibility measurements; but, we'd never do it on our budget. - If you have any comment I'd be interested in hearing it.

I'm leaving for North Carolina Sunday morning and hope to be able to rent a house in short order so we can have our things moved down. I can see a big month of indcision coming up.

Please say Hello to your family from us.

Sincerely,

Al Kolman

THE EFFECT OF MOLECULAR WEIGHT ON THE
COMPRESSIBILITY OF BISPHENOL A POLYCARBONATE

by Allan L. Holmstrom

A Research Report
Presented to the Graduate Faculty
of Lehigh University
in Candidacy for the Degree of
Master of Science

Lehigh University
Bethlehem, Pennsylvania

1965

CERTIFICATE OF APPROVAL

This research report is accepted and approved in partial fulfillment of the requirements for the degree of Master of Science in Chemical Engineering

May 20 1965
(date)

Morton P. Doyle
(Professor in Charge)

(Head of the Department)

Acknowledgements

I wish to express my appreciation to Dr. Morton P. Moyle, my advisor, who offered invaluable advice and encouragement to me; to Joseph Hojsak for his part in constructing the apparatus; and to fellow students who were so willing to give assistance when needed.

Thanks are also in order to General Electric Company's Polycarbonate Research and development Department, the Polymer Corporation of Pennsylvania, Westlake Plastic Company, Cadillac Plastic Company, and Dow Chemical Company for donating samples of the material used in this project, as well as data on physical properties.

Finally, to my wife, Marilyn, who shared in the successes and failures of the work, my thanks for typing this report.

TABLE OF CONTENTS

		<u>Page</u>
I	Certificate Of Approval	1
II	Acknowledgements	11
III	List Of Figures	1
IV	List Of Tables	3
V	Abstract	4
VI	Background	6
VII	Introduction	7
VIII	Equipment	10
IX	Procedure	25
X	Equipment Test	27
XI	Results	33
XII	Discussion Of Results	70
XIII	Conclusions	73
XIV	Nomenclature	74
XV	Bibliography	75
XVI	Vita	77

List of Figures

<u>Figure</u>	<u>Title</u>	<u>Page</u>
1	Detailed Compressibility Cylinder Drawing	11
2	Detailed Plunger Drawing	12
3	Detailed Bottom Plate Drawing	14
4	Schematic Diagram of Bottom Plate Thermocouple Mounting	15
5	Detailed Dynamometer Drawing	16
6	Dynamometer Strain Versus Press Strain	18
7	Load Versus Dynamometer Strain	19
8	Temperature Control Wiring Diagram	21
9	Schematic Diagram of Assembled Compressibility Apparatus	23
10	Assembled Apparatus	24
11	Specific Volume Versus Pressure for Polystyrene	31
12	Specific Volume Versus Pressure for Polymethyl methacrylate	32
13	Compressibility of Lexan Versus Pressure at 50°F	36
14	Compressibility of Lexan Versus Pressure at 100°F	38
15	Compressibility of Lexan Versus Pressure at 150°F	40
16	Compressibility of Lexan Versus Pressure at 200°F	42
17	Compressibility of Lexan Versus Pressure at 250°F	44
18	Compressibility of Lexan Versus Pressure at 300°F	47

List of Figures Continued

19	Compressibility of Lexan Versus Pressure at 350°F	50
20	Possible Interaction Units	53
21	$\sqrt{-T \Delta P / \Delta V}$ Versus Pressure for Lexan at 300°F	56
22	$\sqrt{-T \Delta P / \Delta V}$ Versus Pressure for Lexan at 350°F	57
23	Specific Volume Versus Pressure for Lexan at 300°F	63
24	Specific Volume Versus Pressure for Lexan at 350°F	64
25	Bulk Modulus Versus Pressure for Lexan at 300°F	66
26	Bulk Modulus Versus Pressure for Lexan at 350°F	67

List of Tables

<u>Table</u>	<u>Subject</u>	<u>Page</u>
I	Pressure-Volume Data for Polystyrene	29
II	Pressure-Volume Data for Polymethyl methacrylate	30
III	Compressibility Data for Lexan at 50°F	34
IV	Compressibility Data for Lexan at 100°F	37
V	Compressibility Data for Lexan at 150°F	39
VI	Compressibility Data for Lexan at 200°F	41
VII	Compressibility Data for Lexan at 250°F	43
VIII	Compressibility Data for Lexan at 300°F	45
IX	Compressibility Data for Lexan at 350°F	48
X	Data for $\sqrt{-T\Delta P/\Delta V}$ Versus Pressure	55
XI	Internal Pressure Data for Lexan at 300°F	58
XII	Internal Pressure Data for Lexan at 350°F	59
XIII	Experimental Pressure-Volume Equilibrium Values	61
XIV	Pressure Versus Bulk Modulus Data	65

Abstract

The compressibilities of Lexan ~~130-112~~ ¹⁰¹⁻¹¹² and Lexan ~~130-112~~ ¹³⁰⁻¹¹², having average molecular weights of ~~40,000~~ ^{35,000} and ~~35,000~~ ^{40,000}, respectively, have been investigated over the pressure range from ~~14.7~~ ^{atmospheric @ 30,000} to ~~psia.~~ and over the temperature range from 50°F to 350°F.

The compressibility of Lexan above the glass transition temperature (~~295°F~~ ^{295°F}) was found to be independent of molecular weight; ~~but~~ ^{however, the polymer with} below 200°F the higher average molecular weight ~~40,000~~ ^{40,000} material exhibited a compressibility 25 percent higher than the ~~lower molecular weight material.~~ ^{polymer with lower molecular weight.}

A modified van der Waal's equation of the form

$$(P + 40,000)(V - 0.6487) = 20.8T$$

was ~~shown~~ ^{found} to represent the experimental P-V-T ~~values~~ ^{data} obtained from the compressibility measurements within ± 7%.
Differentiation and
Rearrangement of the Tait equation ~~into terms of bulk~~ ^{yield a relationship between} modulus and pressure ^{a plot of K versus pressure} shows a discontinuity at about 15,000 psi. when ~~plotted~~ ^{a plot of K versus pressure}, indicating the transition of ~~Lexan~~ ^{Lexan} from a rubbery material to a glassy material in the 300°F to 350°F range.

MISSING

PAGES

Background

There is very little data in the literature on the compressibility and P-V-T relationships for polymers. The majority of data reported has consisted primarily of temperature-volume relationships; very little attention has been given to the equally important pressure-volume relationships.

The fabrication of polymers utilizing extruders and injection molding equipment, requires that the polymer be formed under conditions which are often 10-20,000 psi. and 100-300°C above their normal conditions of use. An urgent practical need for P-V-T work on polymers is thus created in relating proper mold size at fabricating conditions to the proper size of the finished product at its normal conditions.

The need for this type of work does not end on a practical basis, however, for it is hoped, that when enough P-V-T data has been accumulated, a general equation of state for polymeric materials will be formulated.

It is the object of this work to explore the P-V-T relationships for the bis-phenol A polycarbonate with particular attention to the effect of molecular weight on P-V-T behavior and to illustrate some of the information that can be obtained directly from a P-V-T investigation.

compressibility data of high polymers
investigations ().

Introduction

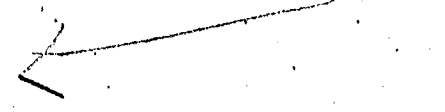
An exhaustive search of the literature failed to reveal any data concerning the effect of molecular weight on the compressibility of high polymers. Spencer and Gilmore have reported equilibrium P-V-T data for Polystyrenes (9,11) having molecular weights of 321,000 and 360,000.

Calculations show that the compressibility of the higher molecular weight sample is greater than the low molecular weight sample for external pressures greater than 14,000 psi. Below 14,000 psi. external pressure there is no statistically significant effect of molecular weight.

No other published data offers information from which a further relationship may be derived for compressibility as a function of molecular weight. Spencer and Gilmore do offer equilibrium P-V-T data at one temperature and one molecular weight for polymethyl methacrylate, ethyl cellulose, cellulose acetate butyrate, and polyethylene at a single temperature and molecular weight level.

The main purpose of Spencer and Gilmore's work was to develop an equation of state for the polymers they worked with. The type of equation they proposed was a modified van der Waal's equation of the form

$$(P + \pi)(V - w) = \frac{RT}{M}$$



where P is the pressure, V the volume and π the "internal" pressure, ω the volume at 0°K , R the gas constant, T the absolute temperature and M the molecular weight of the interaction unit which in many cases is the monomeric repeat unit.

Dynamic compressibility data on several polymers in the frequency range of 0.1 cps. to 60 kc/sec. has been published by Heydemann.⁽⁴⁾ Polyesters having different side groups which alters the molecular weight appear to show a decrease in compressibility with increasing molecular weight. Of course, this is different than increasing molecular weight by increasing the degree of polymerization. Simha has pointed out that isothermal or isobaric volume changes are brought about by two mechanisms, one instantaneous and the other time dependent.⁽⁸⁾ Dynamic data does not take into account the time dependency that can be substantial in viscoelastic materials.

Weir has proposed an equation of the virial type⁽¹³⁾ having 12 constants to describe the P-V-T relationships of ~~fluorocarbon) materials~~ Teflon over a range of 100°C . Although this is an excellent equation for accuracy, its usefulness is very limited since derivatives of the equation cannot be used in the prediction of other thermodynamic coefficients, its size prohibits quick calculation on anything less than a computer, and its temperature range is quite restricted.

An equation of state has been proposed for rubber through thermodynamic reasoning and the statistical model of the rubber molecule. The form of the equation is:

$$\tau = \frac{g \nu R T}{V} \left(1 - \frac{2 M_c}{\text{Mol. Wt.}}\right) \left[\frac{l}{l_0} - \left(\frac{l_0}{l}\right)^2\right]$$

where τ tension, g is a correction factor, ν is number of moles of segments present and corresponds to the n in the ideal gas equation, M_c is the average molecular weight of a segment, l extension length and l_0 is the original unstretched equilibrium length. (3) As indicated, this equation is remarkable similar to the ideal gas equation and has been found to be well suited to rubber.

Recently, Wood of the National Bureau of Standards has shown that the Tait equation, proposed to relate the variables of pressure and volume for sea water, is useful for high polymers (14). The Tait equation

$$\frac{V_0 - V}{V_0} = c \log_{10} \left(1 + \frac{P}{b}\right)$$

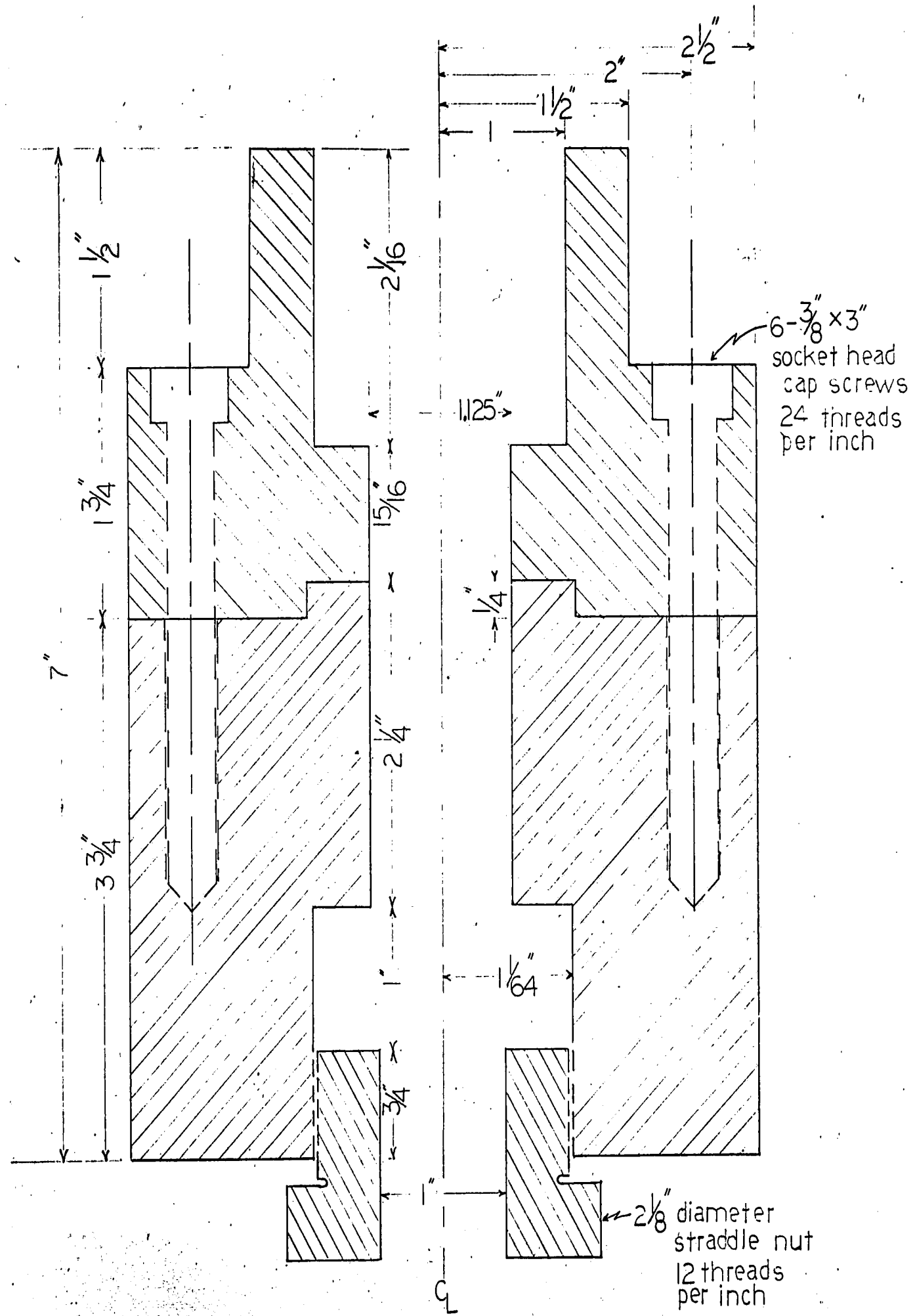
where V is the volume, P is the pressure, V_0 the volume at zero pressure, b and c being constants, is not very useful in this general form, however, rearranging it into the bulk modulus, K , and plotting it against pressure, P , it has been found to predict the second order transition pressure when the polymer is above the glass transition temperature. Second order transition is defined as the point where a polymer changes from a glassy state to a rubbery state, or vice versa. The glass transition temperature is the second order transition temperature under a pressure of one atmosphere.

Equipment

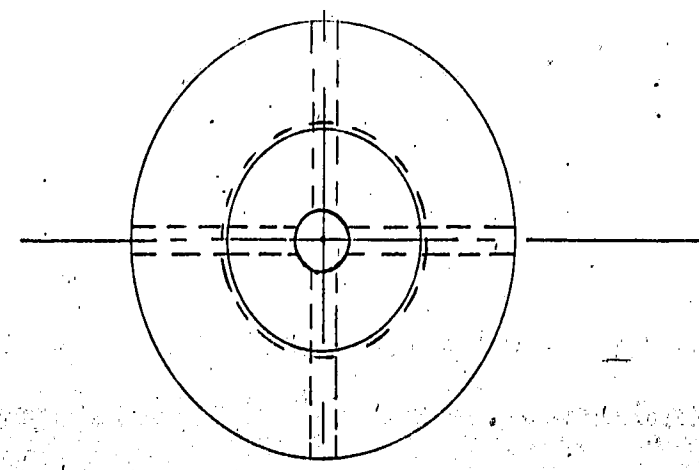
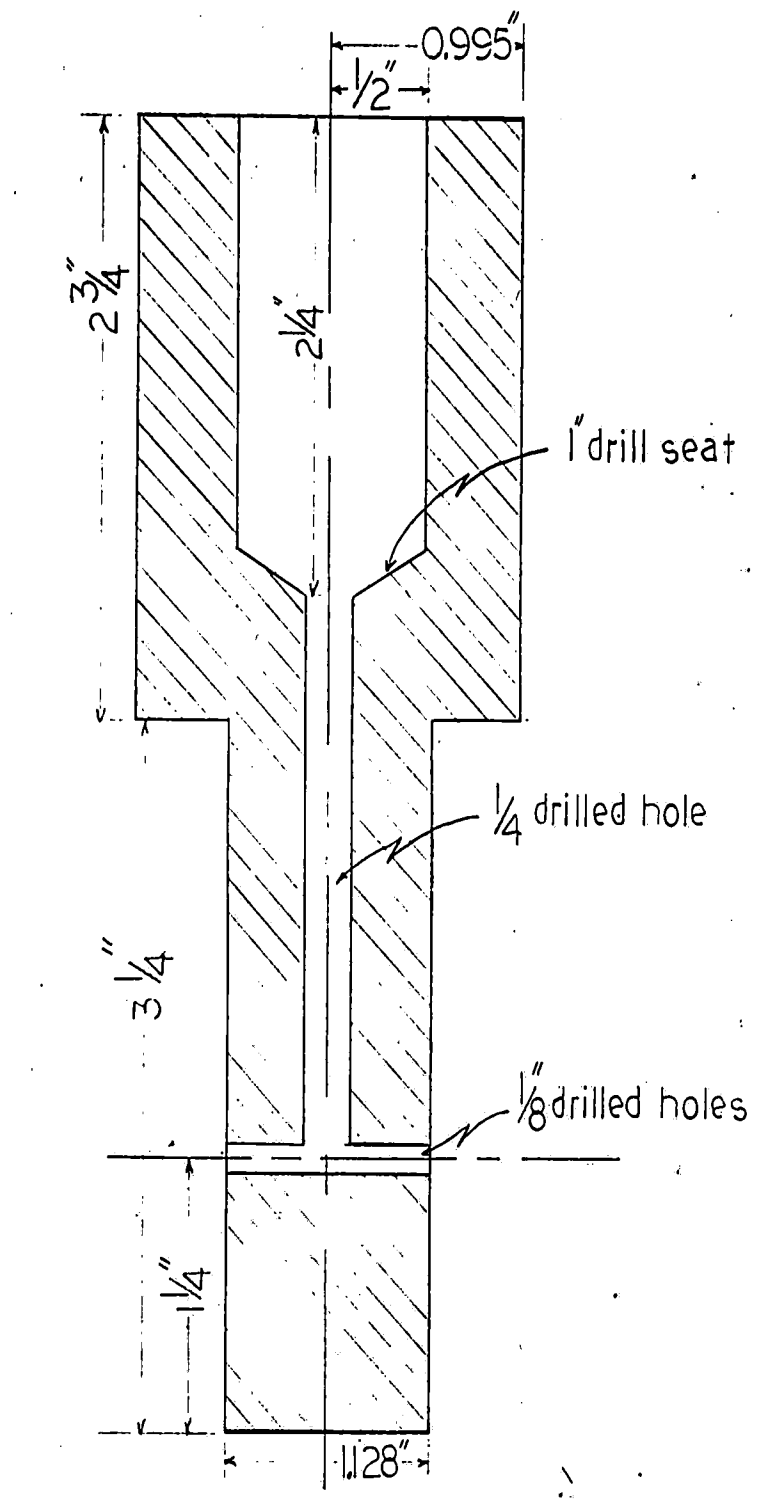
To obtain P-V-T and compressibility data requires constant temperature, constant pressure, volume measurements. For this purpose ^A a compression chamber was constructed from 4140 steel consisting of three basic parts: cylinder, plunger, and bottom plate (figure 1).

The cylinder was turned down in two pieces to the dimensions shown in Figure 1. The reason for two pieces was to allow the equipment to be modified to handle a liquid polymer seal if the need arose. After the cylinder was heat treated to harden it, the inner surface ^{of the cylinder} was honed out to a diameter of 1.135 inches.

In Figure 2 the ^{drawings} specifications for the plunger ~~are~~ given. The top of the plunger is bored to ~~a one-inch diameter~~ ^{so that} to allow for a piston to force molten polymer down the central ~~drilled~~ channel ~~and~~ to the outer surface of the plunger ^{to form} the liquid polymer seal, ~~reference to~~. The lower part of the plunger, having the smallest diameter, was then plated with hard chrome ^{and polished} to ~~about~~ 1.135 inches in diameter. Cylinder and plunger were then worked together with a small quantity of valve grinding compound until a close fit was obtained, but not so tight as to restrict vertical movement of the plunger in the cylinder. ~~The plunger was not hardened so that it can easily be worked to modify it for a polymer seal piston.~~



CYLINDER
 Figure 1

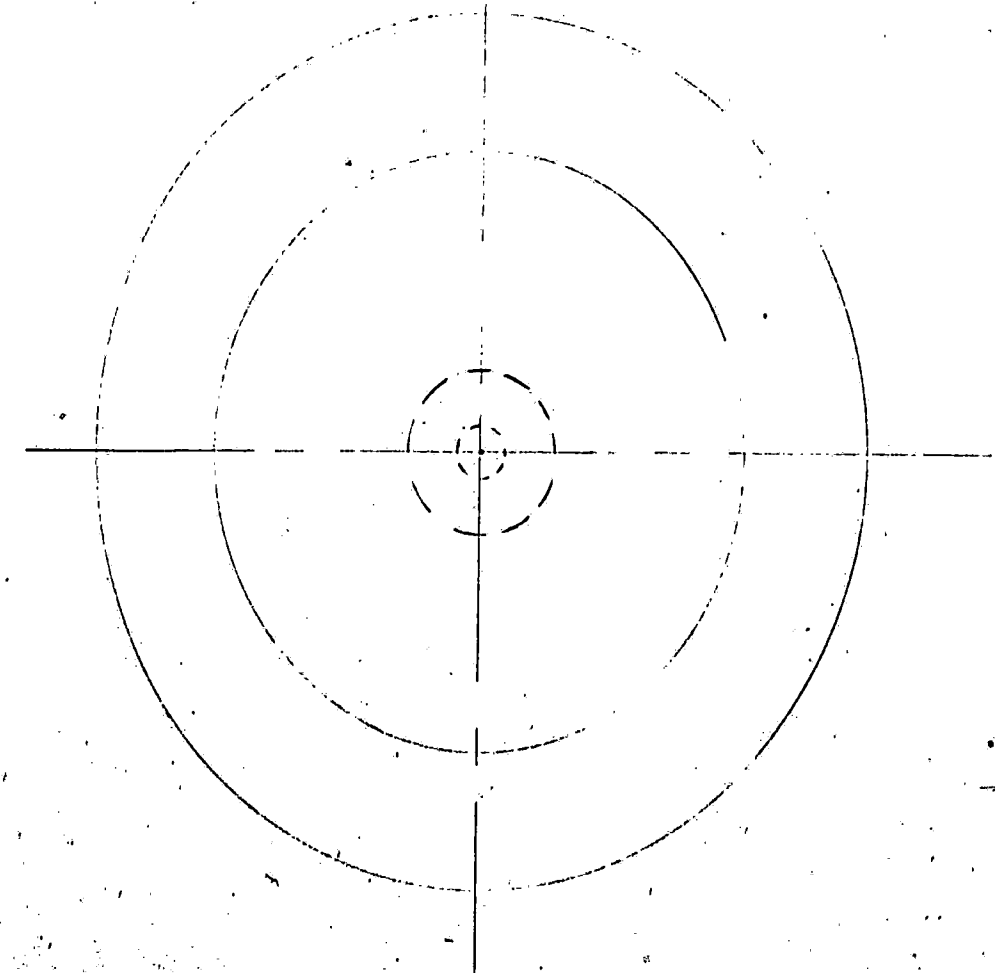
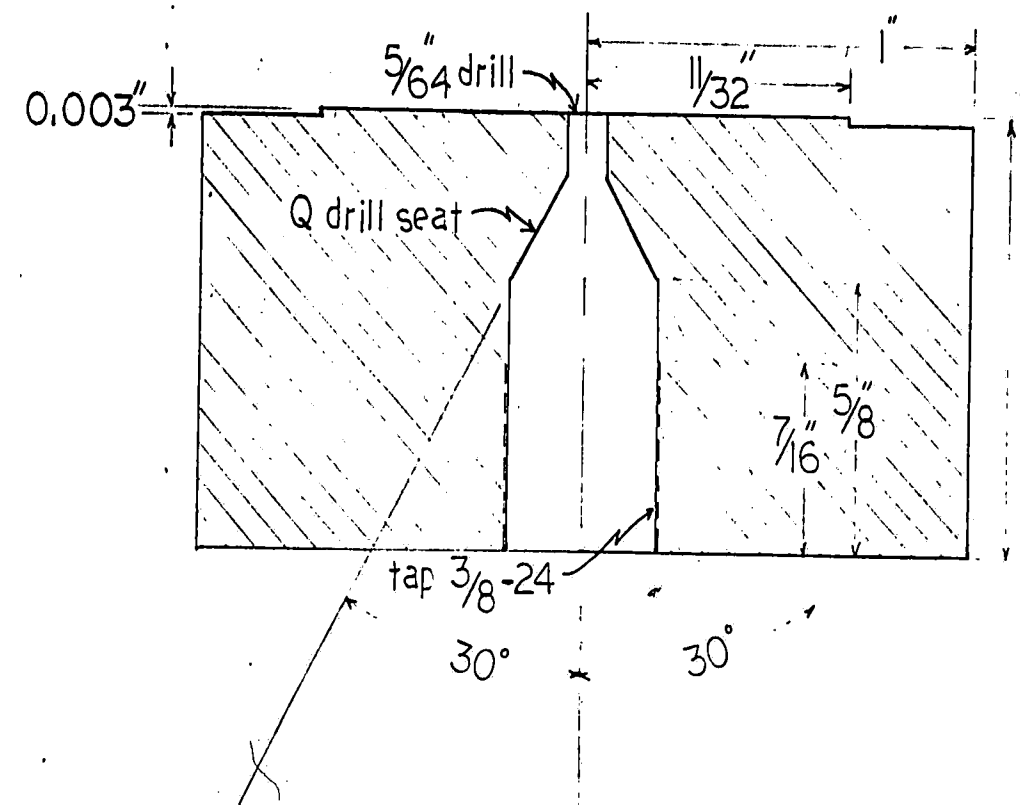


PLUNGER
 Figure 2

^{applies}
~~At the bottom of the cylinder is a plate, shown in~~
~~Figure 3, which retains~~ the thermocouple mounting arrangement
 shown in Figure 4. A 0.003 inch layer was removed from the
 edges of the face of the bottom plate to reduce the surface
 area acting on the lead washer seal between this plate and
 the bottom face of the cylinder.

The ~~thermocouple~~, ^{an} Aero Research Instrument Company, ^{was constant}
~~model T-91-M-4AJ8A16, has a 0.010 inch diameter 304 stainless~~
~~steel sheath with iron-constantan thermocouple wire and MgO~~
~~insulation.~~ The hot junction ^{was} is integrally welded to the
 end of the sheath. The soapstone and steel follower are
~~standard high-pressure fittings made by the American Instru-~~
~~ment Company.~~ Epoxy resin was used at the tip since it is
 an adequate filler as well as an insulator.

In Figure 5 the details of the dynamometer used to
 measure the pressure are shown. A Tinius-Olson Testing
 Machine was used to calibrate the dynamometer so that the
 strain as read by an SR-4 strain indicator, could be directly
 related to the pressure exerted on the polymer sample. Since
 the dynamometer ^{was} would be subjected to a temperature gradient
 it was desirable to examine the effect of the temperature
 gradient on the dynamometer. ^{was investigated.} As a result, an independent
 system was constructed by placing two strain gages on the
 supporting beam of the press. This part of the press was
 always at room temperature. The dynamometer was then
 calibrated against this second system at room temperature,



BOTTOM PLATE
 Figure 3

SCHMATIC DIAGRAM OF BOTTOM PLATE
THERMOCOUPLE MOUNTING

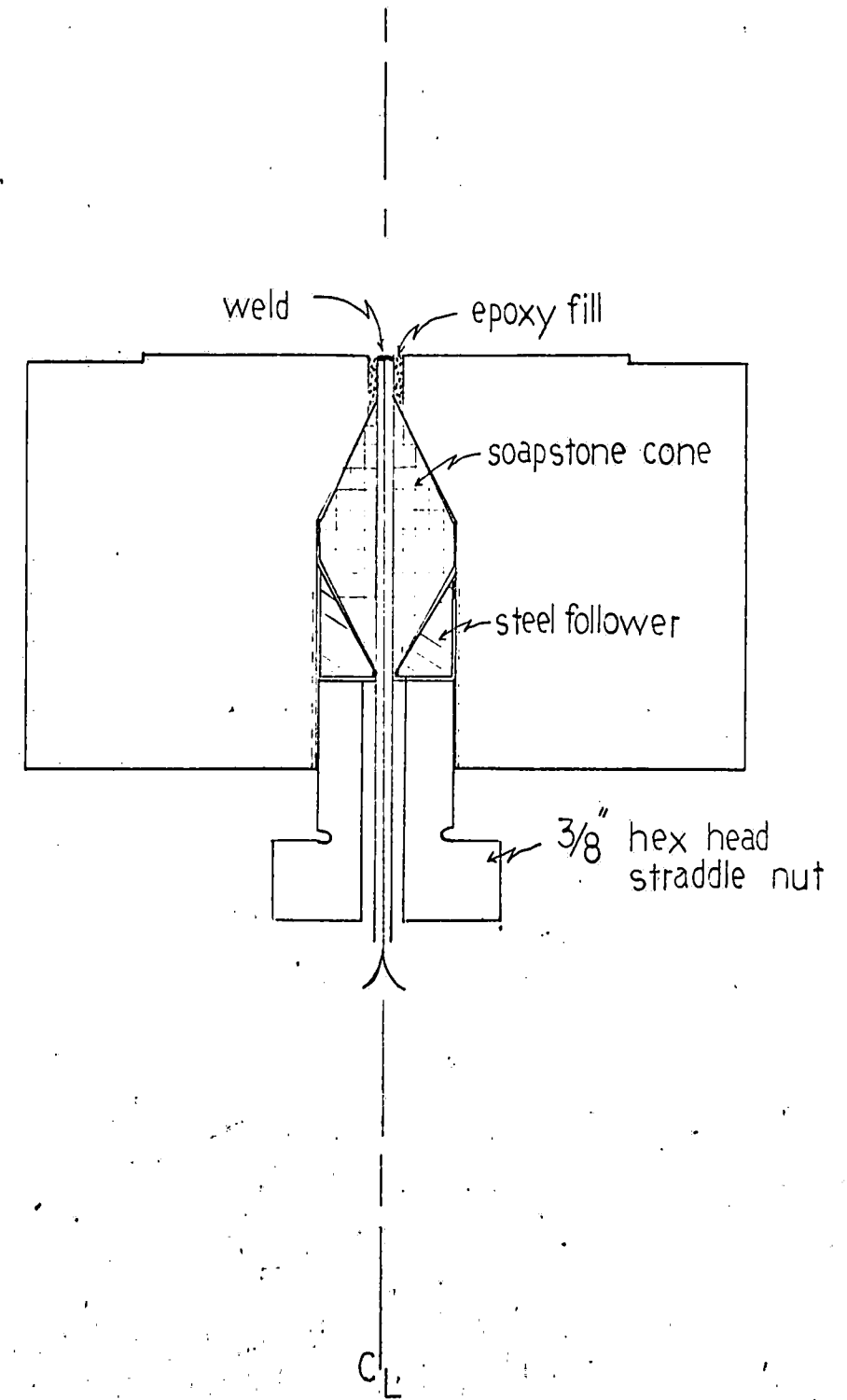


Figure 4

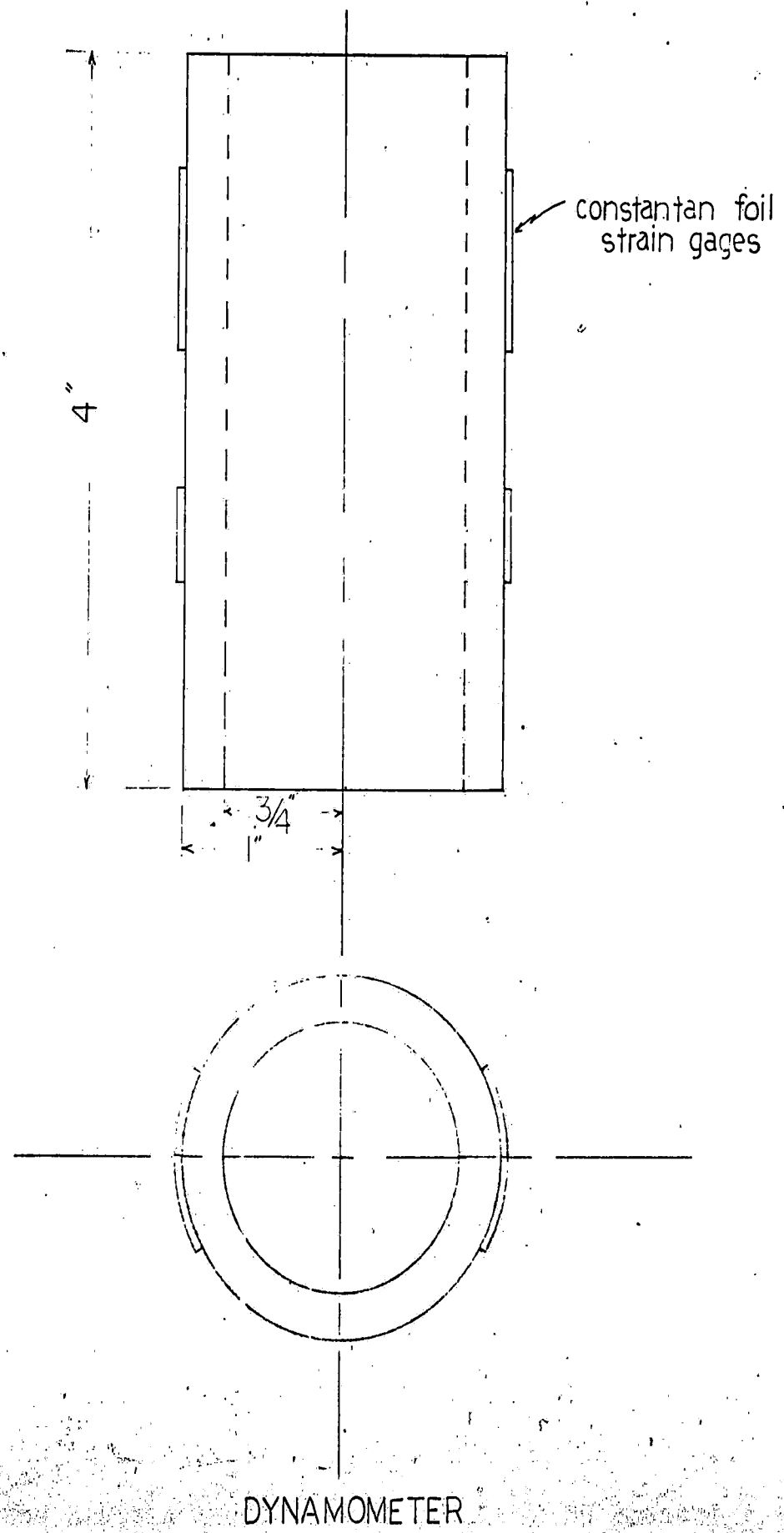


Figure 5

and with a 250°F gradient. In Figure 6 the results are plotted; it can be seen that the zero point is translated, however, the slope is constant which indicates that the true pressure was recorded.

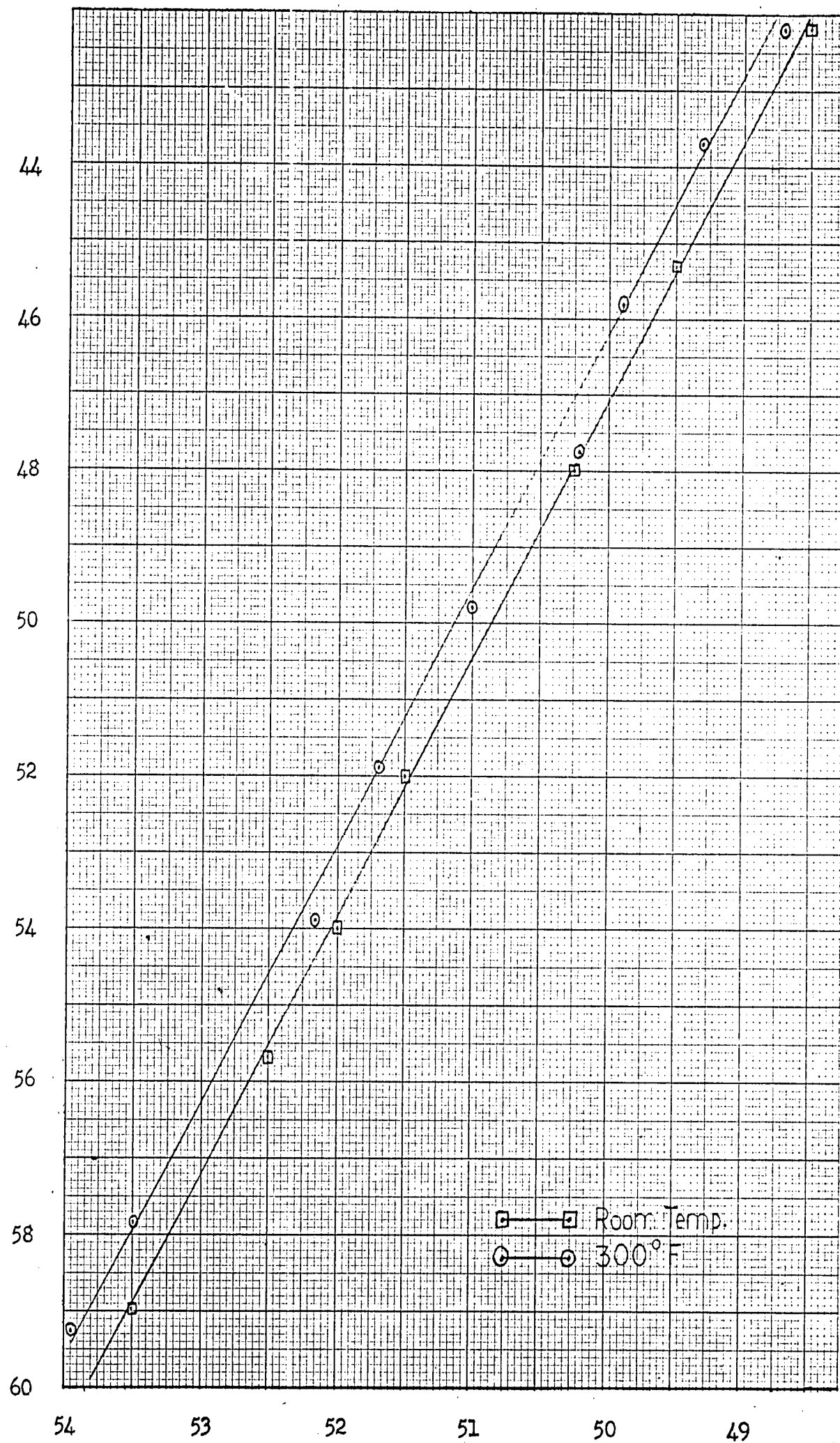
Figure 7 shows the strain versus pressure calibration curve. The actual calibration curve used was plotted on a 2 feet by 4 feet graph.

Heat was applied to the cylinder by two five inch Chromalox band heaters, each capable of delivering 825 watts. In addition, Briskeat flexible heating tape has been placed around the upper part of the plunger and around the two inch straddle nut to decrease the temperature differential between, for example, the plunger in contact with the sample, and the plunger in contact with the surroundings. This reduces the temperature gradient in the plunger and bottom plate and prevents "heat leak" in the axial direction.

One inch thick asbestos pipe insulation encased the bulk of the cylinder. A sheet of aluminum foil was placed on the outside of the insulation and served to reduce radiation losses at higher temperatures and restrict convection currents around the cylinder.

A Wheelco on-off controller was used to control the temperature. When the controller ^{was} ~~is~~ off, a two position relay switch ^{was} ~~is~~ unactivated and allows a Variac to feed power to the heaters in the amount of about 85 percent

DYNAMOMETER STRAIN (μ -in. $\times 10^{-2}$)



Press Strain (μ -in. $\times 10^{-2}$)

Figure 6

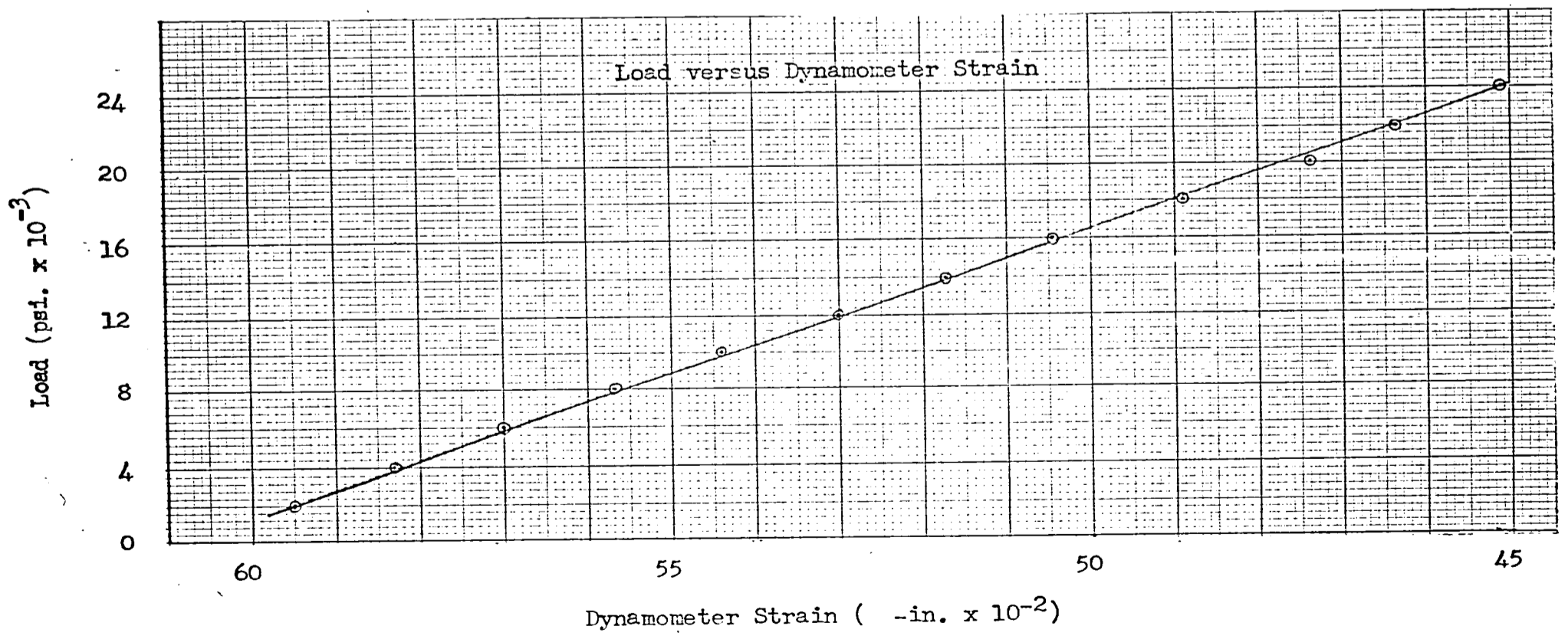


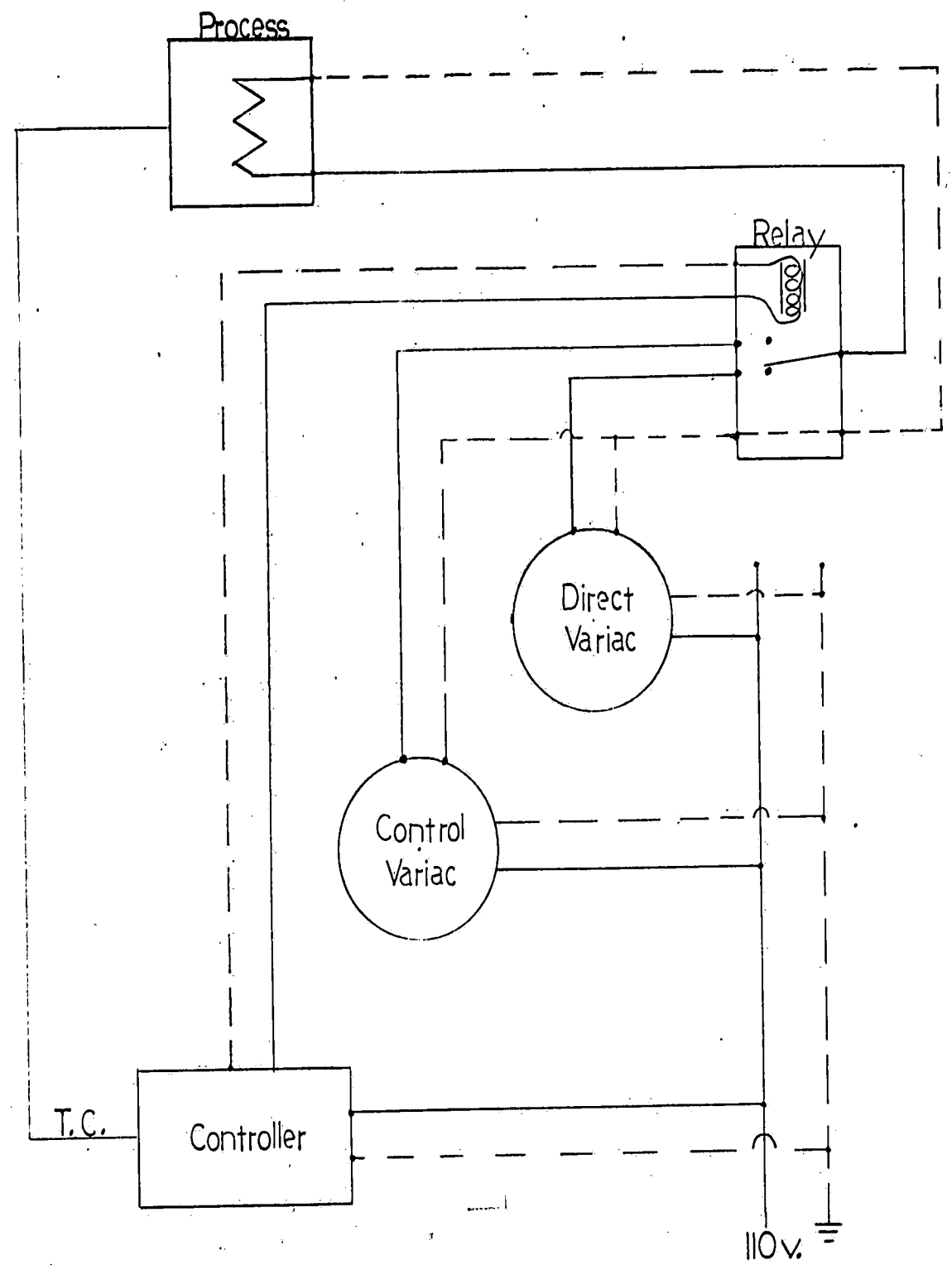
Figure 7

needed to maintain the set point. When the controller ~~was~~ ^{was} on it activates the relay and permits another Variac to supply power of about 110 percent of that needed to maintain the set point. The surface temperature of the cylinder between the two band heaters was used as the control temperature and this arrangement ~~will~~ controlled the sample temperature to within $\pm 2^{\circ}\text{F}$. A ~~wiring diagram~~ of the control and heating elements is contained in ~~Figure 8~~.

Besides the thermocouple in the bottom plate there ~~are~~ ^{were} thermocouples at the bottom and top of the cylinder, between the two heating bands, ambient, and one down the center channel of the plunger. These thermocouples yielded a profile of the temperature of the entire cylinder and ~~the~~ millivolt readings were printed out on a six point Speedomax recorder with a 10 millivolt range. A ^{portable} potentiometer was used to ~~make~~ ^{take} accurate measurements of the temperature of the sample through the bottom plate thermocouple. The potentiometer ~~will~~ read the thermocouple in contact with the sample to $\pm 1^{\circ}\text{F}$ accuracy.

Pressure was applied to the sample ~~with~~ ^{via} a hydraulic press having a 60,000 psi. capacity. Hydraulic leakage was ~~very small~~ ^{imperceptible} and did not pose a problem in maintaining ~~steady~~ pressure.

Volume changes were indicated by a dial micrometer which measured the distance the plunger traveled as the



TEMPERATURE CONTROL WIRING DIAGRAM

Figure 8

sample was compressed.

~~See~~ Figure 9 ^{shown} ~~for~~ a schematic diagram of the compression cylinder, and ~~Figure 10 for the view of the assembled apparatus and control equipment.~~

Calculation of the compressibility of the plunger and the temperature effects on the plunger and cylinder ^{will} yield about ^{a maximum of ±1%} 1 percent error under the most extreme conditions.

~~Hence~~, This factor ^{was} ~~is~~ ignored since the accuracy of the data will only be ±1 percent due to errors in pressure and compression measurements.

SCHMATIC DIAGRAM OF ASSEMBLED
COMPRESSIBILITY APPARATUS

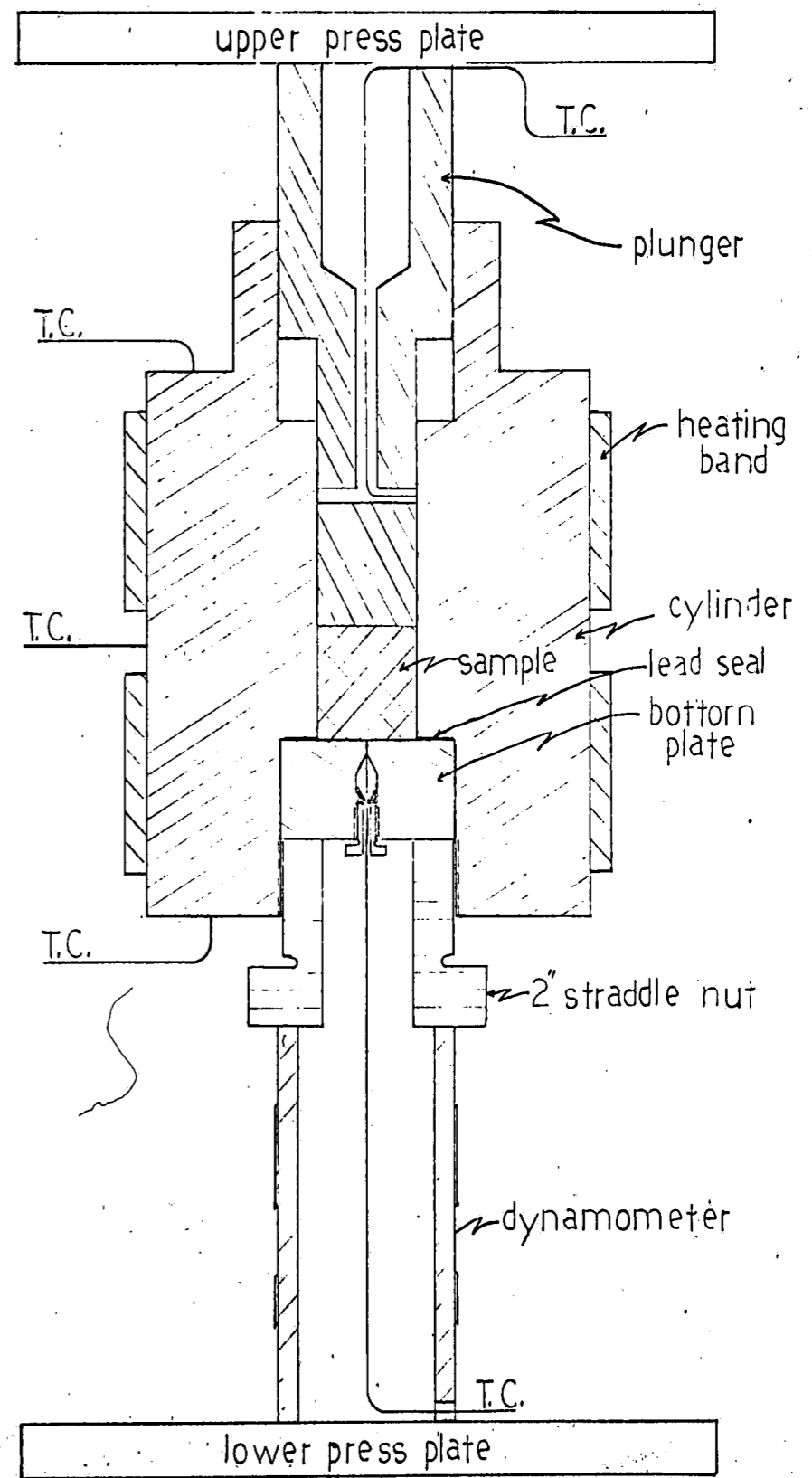


Figure 9

ASSEMBLED APPARATUS

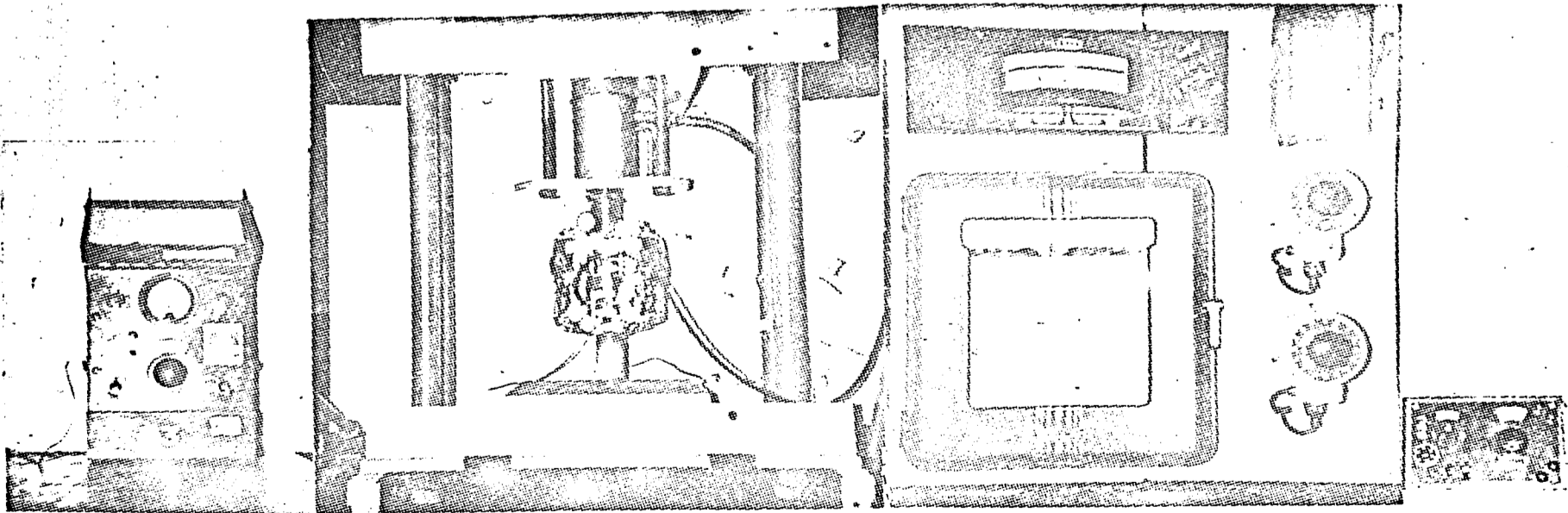


Figure 10

Procedure

Test specimens were prepared by carefully machining them from one and a half inch extruded rods of the sample material. The diameter of the specimen was 1.135 inches. This diameter produced a snug fit in the cylinder requiring the help of the hydraulic press to fully seat it, yet, without the problem of having the leading edge of the cylinder scraping material off the sides of the specimen.

Drying of the specimen was necessary to remove absorbed moisture which would yield a weight slightly higher than that of the dry specimen. Twelve hours at 175°F proved to be sufficient drying condition. After removing from the drying oven the specimen was kept in a desiccator until it was placed in the compressibility apparatus.

Careful measurements with a micrometer determined the initial volume of the specimen. The weight of the specimen was obtained from an analytical balance.

When the apparatus had been assembled containing the specimen, the setpoint of the controller was set at a desired temperature and the system allowed to achieve thermal equilibrium. At the equilibrium temperature a desired pressure was applied to the plunger by the hydraulic press. When the volume became constant, as indicated by the dial micrometer showing that the plunger had stopped its forward progress, temperature, pressure, and volume

readings were recorded. A new pressure was then set and when equilibrium readings were obtained they were recorded. When the entire range of pressure had been examined a new temperature was assigned and the procedure repeated.

The increments used were as follows:

1. Pressure: 3,000 psi. intervals over the range of 0 to 30,000 psi.
2. Temperature: 50°F intervals from room temperature to 350°F.

Identical procedure was used for both samples Lexan 130-112 and Lexan 101-112, a polycarbonate prepared from bisphenol-A having weight average molecular weights of 40,000 and 35,000, respectively.

At the end of a run the specimen was removed and examined for thermal damage and weighed to determine if any weight was lost.

Water acts as a catalyst for the decomposition of polycarbonates. Because of the long times the sample was required to be at a temperature and the fact that it is difficult to keep small amounts of water from the sample, the maximum workable temperature was 350°F. Above this temperature decomposition occurred too rapidly to reliably evaluate the properties of the material.

Equipment Test

Spencer and Gilmore's P-V-T work with five polymers made it possible to test the apparatus to see if the results obtained would be consistent with the work already performed. Since polymethyl methacrylate and polystyrene samples were available, and had been examined by Spencer and Gilmore, they were both tested in this apparatus.

The equation Spencer and Gilmore suggested for polystyrene was:

$$(P + 27,000) (V - 0.882) = 11.6T$$

Since the equation was obtained from data having temperatures of 100 to 300°F above the second order transition temperature, the equation can only be applied above this transition temperature. Thus, 200°F was chosen for the comparison and the equation is plotted in Figure 11 at this temperature.

To obtain the volume as a function of pressure, one must first know the starting volume of the sample. If the room temperature volume is determined from micrometer measurements, the initial volume at an elevated temperature can be calculated by adding to the room temperature volume a volume contribution from thermal expansion. For an amorphous material where there are no preferential thermal expansion directions due to crystal considerations this volume can be determined utilizing the fact that the coefficient of thermal expansion will be the same in three dimensions.

Once the initial volume has been determined, the volume at any pressure is the initial volume less the decrease in volume by compression. In Table I are listed the volumes predicted by Spencer and Gilmore's equation for selected pressures, and the results this equipment produces. The experimental points are also plotted in Figure 11.

Data for polymethyl methacrylate derived in an analogous manner to those for polystyrene are given in Table II. Spencer and Gilmore's equation for polymethyl methacrylate was:

$$(P+31,300)(V-0.735) = 12.05T$$

These results are plotted in Figure 12. Equilibrium time was considered to be approximately three and a half hours for both materials.

In both cases the experimental results are seen to be in agreement with the predictions of the equations. What deviation does exist arises from two sources: one, the work of Spencer and Gilmore was carried out at a temperature about 200°F higher; and two, this work and Spencer and Gilmore's work only claim one to two percent accuracy due to normal experimental errors.

TABLE I.

Polystyrene

At 200°F

$$V = \frac{(11.6)(366)}{(P+27,000)} = 0.822$$

<u>Pressure (psi.)</u>	<u>Volume (cc/g)</u>
0	0.9795
5,000	0.9549
10,000	0.9369
15,000	0.9232
20,000	0.9125
25,000	0.9038
30,000	0.8966

Experimental Results

Sample weight = 23.833g.

<u>Pressure (psi.)</u>	<u>Height (in)</u>	<u>Volume (in.³)</u>	<u>Volume (cc/g)</u>
0	1.3840	1.4028	0.9646
3,100	1.3680	1.3866	0.9534
6,510	1.3532	1.3716	0.9431
10,060	1.3416	1.3598	0.9351
12,300	1.3340	1.3521	0.9297
21,820	1.3175	1.3354	0.9182
29,820	1.3036	1.3216	0.9087

TABLE II.

Polymethyl methacrylate

At 200°F

$$V = \frac{(12.05)(366)}{(P + 31,300)} = 0.734$$

<u>Pressure (psi.)</u>	<u>Volume (cc/g)</u>
0	0.8750
5,000	0.8550
10,000	0.8408
15,000	0.8293
20,000	0.8200
25,000	0.8123
30,000	0.8060

Experimental Results

Sample weight = 26.650 g.

<u>Pressure (psi.)</u>	<u>Height (in)</u>	<u>Volume (in.³)</u>	<u>Volume (cc/g.)</u>
0	1.4027	1.4218	0.8743
2,950	1.3862	1.4050	0.8639
5,820	1.3720	1.3906	0.8551
10,550	1.3598	1.3783	0.8476
14,760	1.3498	1.3682	0.8414
18,260	1.3396	1.3572	0.8349
21,800	1.3309	1.3490	0.8296
25,850	1.3207	1.3387	0.8232
30,470	1.3128	1.3306	0.8183

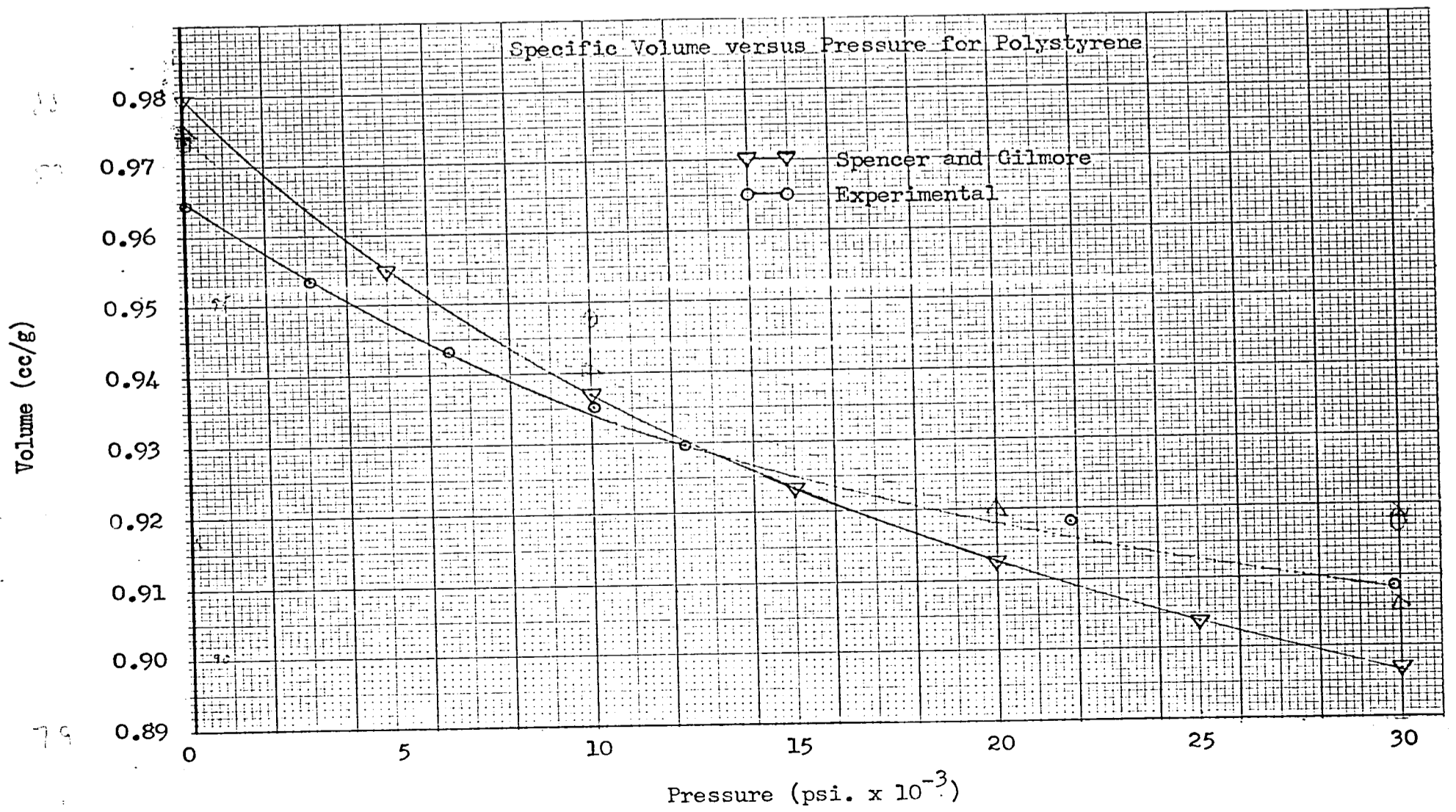


Figure 11

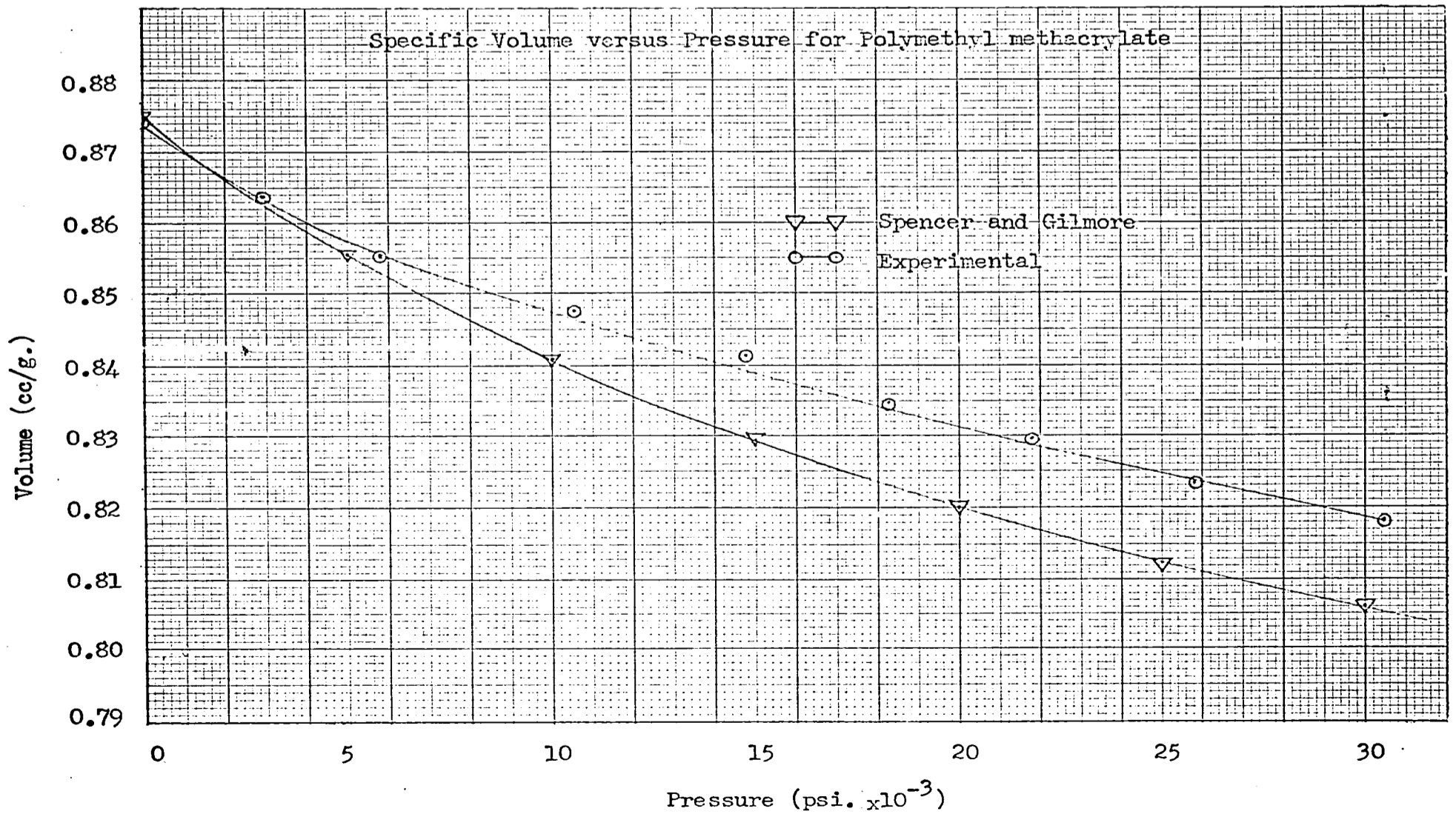


Figure 12

Results

The compressibility (β) is defined as $-\frac{1}{V_0} \frac{\partial V}{\partial P}$ where V_0 is the volume at zero pressure and ∂V and ∂T are differential changes in volume and temperature, respectively. V_0 can be obtained by extrapolating a plot of volume versus pressure, however, the difference between the volume at zero pressure and the volume at one atmosphere is so small that we can use the volume at one atmosphere without introducing appreciable error. Furthermore, $\frac{\partial V}{\partial P}$ can be approximated quite well by $\frac{\Delta V}{\Delta P}$ as long as the straight line relationship implied by these finite difference quantities closely approximates the actual curve. Then, at constant temperature, compressibility may be calculated by evaluating the volume of a sample at several pressures.

$$\Delta V = V_b - V_a, \quad \Delta P = P_b - P_a, \quad V_1 \text{ volume @ one atm.}, \quad \beta = -\frac{1}{V_1} \frac{\Delta V}{\Delta P}$$

The data collected from this experimental work is tabulated in Tables III through IX. They include readings every 50°F from 50°F to 350°F, each through a pressure range of 30,000 psi. The values of compressibility are plotted against the average pressure $((P_b + P_a)/2)$ in Figures 13 through 19.

From these plots it can be seen that at low temperatures, well below the glass transition temperature, the Lexan 130-112, the higher average molecular weight material is more compressible than the Lexan 101-112 of lower average molecular weight. This trend continues until

TABLE III.

Lexan at 50°F

Pressure a (psi.)	Pressure b (psi.)	Average Pressure F (psi.)	Pressure Increment ΔP (psi.)	Volume Change ΔV (in. ³ x 10 ³)	Compressibility, $\beta = -1/V(\Delta V/\Delta P)$ (psi. ⁻¹ x 10 ⁶)
Lexan 101-112 $1/V = 0.7220$					
0	5,120	2,560	5,120	7.222	1.020
5,120	9,120	7,165	4,090	4.999	0.882
9,120	12,420	10,815	3,210	4.444	1.001
12,420	16,610	14,515	4,190	4.747	0.821
16,610	20,140	18,375	3,530	5.252	1.078
20,140	24,160	22,150	4,020	6.514	1.172
24,160	27,230	25,695	3,070	4.494	1.049
27,230	31,390	29,310	4,160	7.020	1.216
31,390					
Lexan 130-112 $1/V = 0.7220$					
4,720	7,710	6,215	2,990	5.151	1.244
9,780	15,330	12,555	5,550	9.393	1.223
15,330	19,080	17,205	3,750	6.666	1.288
19,080	22,130	20,605	3,050	6.464	1.526
21,800	26,970	24,385	5,170	9.999	1.396
26,970	29,660	28,315	2,690	5.959	1.606

TABLE III. CONT.

<u>P(a)</u>	<u>P(b)</u>	<u>\bar{P}</u>	<u>ΔP</u>	<u>ΔV</u>	<u>β</u>
Lexan 130-112 $1/V = 0.7220$					
3,380 / ^{1.3850} 1.3767	7,990	5,685	4,610	8.282	1.288 .776
7,990 / 1.3709	11,390	9,690	3,400	5.858	1.244 .804
11,390 / 1.3657	14,610	13,000	3,220	5.151	1.157 .824
14,610 / 1.3571	19,225	16,917	4,615	8.585	1.346 .793
19,225 / 1.3425	26,700	22,962	7,475	14.645	1.418 .705
26,700 / 1.3348	30,280	28,490	3,580	7.676	1.548 .690

$\frac{1}{\beta}$

MISSING

PAGES

TABLE IV.

Lexan at 100°F

Pressure a (psi.)	Pressure b (psi.)	Average Pressure P (psi.)	Pressure Increment ΔP (psi.)	Volume Change ΔV (in.)X10 ³	Compressibility ₁ - -1/V(ΔV/ΔP) (psi.)X10 ⁶
Lexan 101-112 1/V = 0.7180					
1.3928					
2,150/1.3921	7,150	4,640	4,980	6.943	1.001 10.0
7,130/1.3918	9,680	8,405	2,550	3.345	0.942 10.615
9,680/1.3913	13,490	11,585	3,810	4.663	0.879 11.376
13,490/1.3910	15,720	14,605	2,230	3.345	1.077 9.285
15,720/1.3905	19,530	17,625	3,810	4.460	0.840 11.905
19,530/1.3900	23,600	21,565	4,070	5.068	0.894 11.186
23,600/1.3895	26,970	25,285	3,370	5.271	1.123 8.905
26,970/1.3888	30,925	29,948	3,955	6.538	1.187 8.425
Lexan 130-112 1/V = 0.7162					
1.3963					
2,280/1.3939	6,500	4,390	4,220	2.438	0.422 13.69
6,500/1.3901	9,000	7,750	2,500	3.750	1.077 9.29
9,000/1.3853	12,300	10,515	3,030	4.815	1.141 8.76
12,030/1.3807	15,420	13,735	3,390	4.612	0.976 10.25
15,420/1.3744	18,000	17,210	3,580	6.335	1.270 7.77
18,000/1.3683	21,200	19,600	3,200	6.087	1.366 7.32
21,200/1.3619	24,900	23,050	3,700	6.740	1.307 7.65
24,900/1.3553	27,750	26,325	2,850	6.487	1.634 6.12
27,750/1.3492	31,150	29,450	3,400	6.087	1.285 7.78

MISSING

PAGES

TABLE V.

Lexan at 150°F

Pressure a (psi.)	Pressure b (psi.)	Average Pressure P (psi.)	Pressure Increment ΔP (psi.)	Volume Change, ΔV (in. ³ x 10 ³)	Compressibility $\beta = -1/V(\Delta V/\Delta P)$ (psi. ⁻¹ x 10 ⁶)
Lexan 101-112 $1/V = 0.7138$					
0	5,410	2,705	5,410	8.058	1.063 9.467
5,410	7,730	6,750	2,320	6.234	1.991 5.022
7,730	11,230	9,480	3,500	5.017	1.023 9.775
11,230	15,075	13,152	3,845	4.581	0.850 11.765
15,075	19,380	17,227	4,305	5.757	0.955 10.471
19,380	21,970	20,675	2,590	3.953	1.099 9.099
21,970	24,420	23,195	2,450	3.801	0.898 11.136
24,420	27,360	25,990	2,940	5.119	1.242 8.052
24,420	30,900	27,660	6,480	10.389	1.144 8.741
Lexan 130-112 $1/V = 0.7138$					
2,310	8,010	5,160	5,700	8.565	1.073 .932
8,010	10,450	9,230	2,440	8.159	2.384 .419
10,450	13,330	11,890	2,880	5.980	1.482 .675
13,330	18,160	15,745	4,830	12.467	1.842 .543
18,160	21,360	19,760	3,200	6.082	1.356 .735
21,360	27,170	24,265	5,810	10.846	1.333 .750
27,170	29,410	28,290	2,240	4.967	1.582 .632

MISSING

PAGES

TABLE VI.

Lexan at 200°F

Pressure a (psi.)	Pressure b (psi.)	Average Pressure P (psi.)	Pressure Increment ΔP (psi.)	Volume Change ΔV (in ³ x10 ³)	Compressibility $\beta = -1/V(\Delta V/\Delta P)$ (psi. ⁻¹ x10 ⁶)
		Lexan 101-112		$1/V = 0.7098$	
0	3,050	1,525	3,050	7.800	1.816
3,050	6,350	4,700	3,300	4.511	0.970
6,350	9,650	8,000	3,300	4.815	1.036
9,650	12,570	11,100	2,920	5.17	1.257
12,570	16,130	14,350	3,560	5.372	1.072
16,130	19,030	17,580	2,900	6.791	1.662
19,030	22,150	21,590	3,120	6.943	1.580
22,150	25,570	23,860	3,420	5.423	1.126
25,570	28,550	27,060	2,980	7.197	1.714
28,550	31,530	28,550	5,960	13.481	1.497
28,550	31,530	30,040	2,980	6.284	1.606
		Lexan 130-112		$1/V = 0.7099$	
3,460	7,100	5,330	3,740	5.271	1.000
7,100	10,680	8,890	3,580	7.044	1.397
10,680	14,950	12,815	4,270	6.740	1.120
14,950	18,550	16,750	3,600	7.450	1.468
18,550	24,460	21,505	5,910	8.109	0.973
24,460	26,180	24,320	3,720	8.818	1.682
26,180	30,560	28,370	4,380	9.883	1.601

MISSING

PAGES

TABLE VII.
Lexan at 250°F

Pressure a (psi.)	Pressure b (psi.)	Average Pressure P (psi.)	Pressure Increment ΔP (psi.)	Volume Change ΔV (in. $\times 10^3$)	Compressibility $\beta = -1/V(\Delta V/\Delta P)$ (psi. $^{-1} \times 10^6$)
Lexan 101-112 $1/V = 0.7058$					
0	2,130	1,065	2,130	5.473	1.813
2,130	5,820	3,975	3,690	7.450	1.425
5,820	10,300	8,060	4,480	8.109	1.271
10,300	14,100	12,200	3,800	7.146	1.328
14,100	17,650	15,875	3,550	7.298	1.452
17,650	21,520	19,585	3,870	8.970	1.636
21,520	25,900	23,710	4,380	9.274	1.495
25,900	29,550	27,725	3,650	7.602	1.471
17,650	25,520	21,585	7,870	18.240	1.637
Lexan 130-112 $1/V = 0.7059$					
2,450	6,740	4,595	4,290	8.413	1.384
6,740	10,750	8,745	4,010	7.450	1.312
10,750	13,740	12,245	2,990	6.639	1.567
13,740	19,110	16,425	5,370	9.477	1.246
19,110	22,440	20,775	3,330	7.602	1.612
22,440	25,820	24,130	3,380	7.805	1.629
25,820	30,275	28,048	4,455	9.376	1.485

MISSING

PAGES

TABLE VIII.

Lexan at 300°F

<u>Pressure a</u> (psi.)	<u>Pressure b</u> (psi.)	<u>Average Pressure</u> <u>F (psi.)</u>	<u>Pressure Increment</u> <u>ΔP (psi.)</u>	<u>Volume Change</u> <u>ΔV (in.³ x10³)</u>	<u>Compressibility</u> <u>β = -1/V(ΔV/ΔP) (psi.⁻¹ x10⁶)</u>
Lexan 101-112 1/V = 0.7018					
3,050	5,640	4,345	2,590	16.015	4.341
5,640	10,000	7,820	4,360	19.157	3.085
10,000	13,380	11,690	3,380	8.818	1.831
13,380	17,630	15,505	4,250	10.947	1.808
17,630	19,380	18,505	1,750	4.663	1.871
19,380	24,670	22,025	5,290	10.541	1.399
24,670	28,550	26,610	3,880	7.298	1.321
19,380	23,480	21,430	4,100	9.122	1.563
Lexan 130-112 1/V = 0.7019					
350	2,880	1,615	2,530	20.576	5.708
2,880	6,150	4,515	3,270	16.826	3.611
6,150	9,630	7,890	3,480	13.532	2.729
9,630	13,990	11,810	4,360	20.525	3.304
13,990	19,110	16,550	5,120	8.616	1.181
19,110	21,600	20,355	2,490	10.136	2.857
21,600	25,300	23,450	3,700	8.869	1.683
25,300	29,360	27,330	4,060	7.247	1.253

TABLE VIII. CONT.

<u>P(a)</u>	<u>P (b)</u>	<u>\bar{P}</u>	<u>ΔP</u>	<u>ΔV</u>	<u>β</u>
Lexan 130-112 $1/V = 0.7009$					
1,660	4,150	2,905	2,490	14.342	4.037 <i>.298</i>
4,150	7,950	6,050	3,800	16.015	2.954 <i>.359</i>
7,950	10,550	9,250	2,600	10.440	2.815 <i>.458</i>
10,550	13,800	12,175	3,250	7.602	1.640 <i>.616</i>
13,800	16,550	15,175	2,750	7.196	1.834 <i>.545</i>
16,550	19,420	17,985	2,870	6.791	1.658 <i>.603</i>
19,420	22,560	20,990	3,140	8.565	1.912 <i>.528</i>
22,560	25,610	24,085	3,050	6.943	1.595 <i>.627</i>
25,610	29,730	27,670	4,120	7.45	1.267 <i>.789</i>

MISSING

PAGES

TABLE IX.
Lexan at 350°F

<u>Pressure a</u> (psi.)	<u>Pressure b</u> (psi.)	<u>Average Pressure</u> <u>P</u> (psi.)	<u>Pressure Increment</u> <u>ΔP</u> (psi.)	<u>Volume Change</u> <u>ΔV</u> (in. ³ × 10 ³)	<u>Compressibility</u> <u>β</u> = -1/V(ΔV/ΔP) (psi. ⁻¹ × 10 ⁶)
Lexan 101-112 1/V = 0.6980					
1,900 / 1.4327	5,270	3,585	3,370	21.691	4.493
5,270 / 1.3937	8,440	6,855	3,170	17.282	3.806
8,440 / 1.3784	11,650	10,045	3,210	15.356	3.131
11,650 / 1.3668	13,950	12,900	3,210	11.555	3.507
13,950 / 1.3576	17,800	15,875	3,850	15.255	2.766
17,800 / 1.3392	21,000	19,400	3,200	12.366	2.698
21,000 / 1.3295	26,600	23,800	5,600	9.680	1.207
26,600 / 1.3238	28,300	27,450	1,700	5.676	2.331
21,000 / 1.3084	28,300	24,650	7,300	15.356	1.468
Lexan 130-112 1/V = 0.6979					
490 / 1.4329	3,130	1,810	2,640	20.982	5.548
3,130 / 1.3959	6,150	4,640	3,020	15.964	3.689
6,150 / 1.3767	10,480	8,315	4,330	19.208	3.097
10,480 / 1.3625	14,100	12,290	3,620	14.190	2.736
14,100 / 1.3487	17,360	15,730	3,260	13.785	2.952
17,360 / 1.3389	21,770	19,565	4,410	9.832	1.556
21,770 / 1.3312	25,000	23,385	3,230	7.754	1.676

TABLE IX. CONT.

<u>P(a)</u>	<u>P(b)</u>	<u>\bar{P}</u>	<u>P</u>	<u>V</u>	<u>β</u>
Lexan 130-112 $1/V = 0.6973$					
9,000 ^{1.4341} / 1.4226	11,660	10,330	2,660	11.500	3.014
11,600 / 1.4088	14,620	13,640	2,960	13.836	3.265
14,620 / 1.3954	18,000	16,310	3,380	13.380	2.760
18,000 / 1.3840	21,420	19,710	3,420	11.352	2.314
21,420 / 1.3761	24,330	22,875	2,910	7.906	1.896
24,330 / 1.3696	27,200	25,765	2,870	6.487	1.576
27,200 / 1.3618	30,860	29,030	3,660	7.805	1.487

Compressibility ($\text{psi.}^{-1} \times 10^{-6}$)

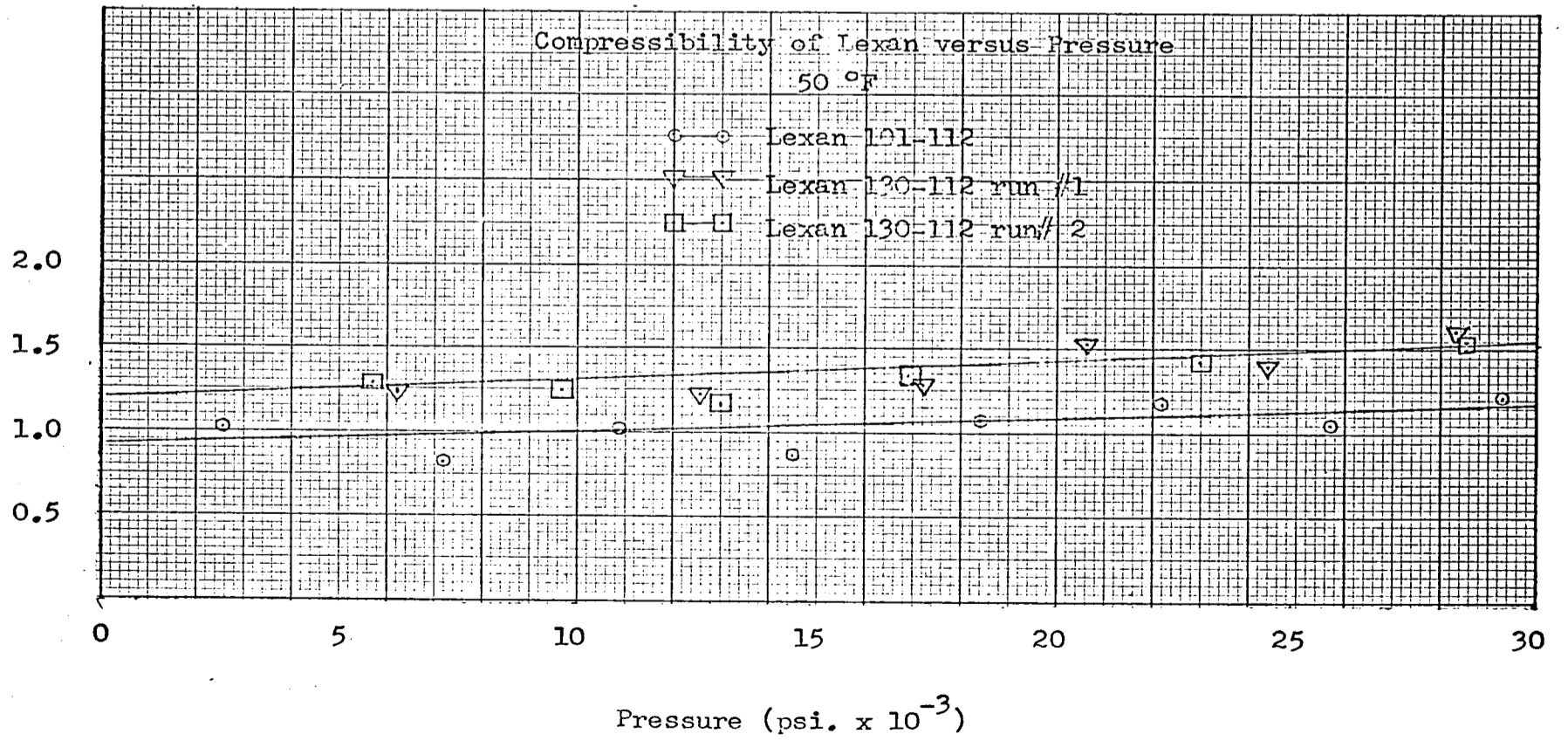


Figure 13

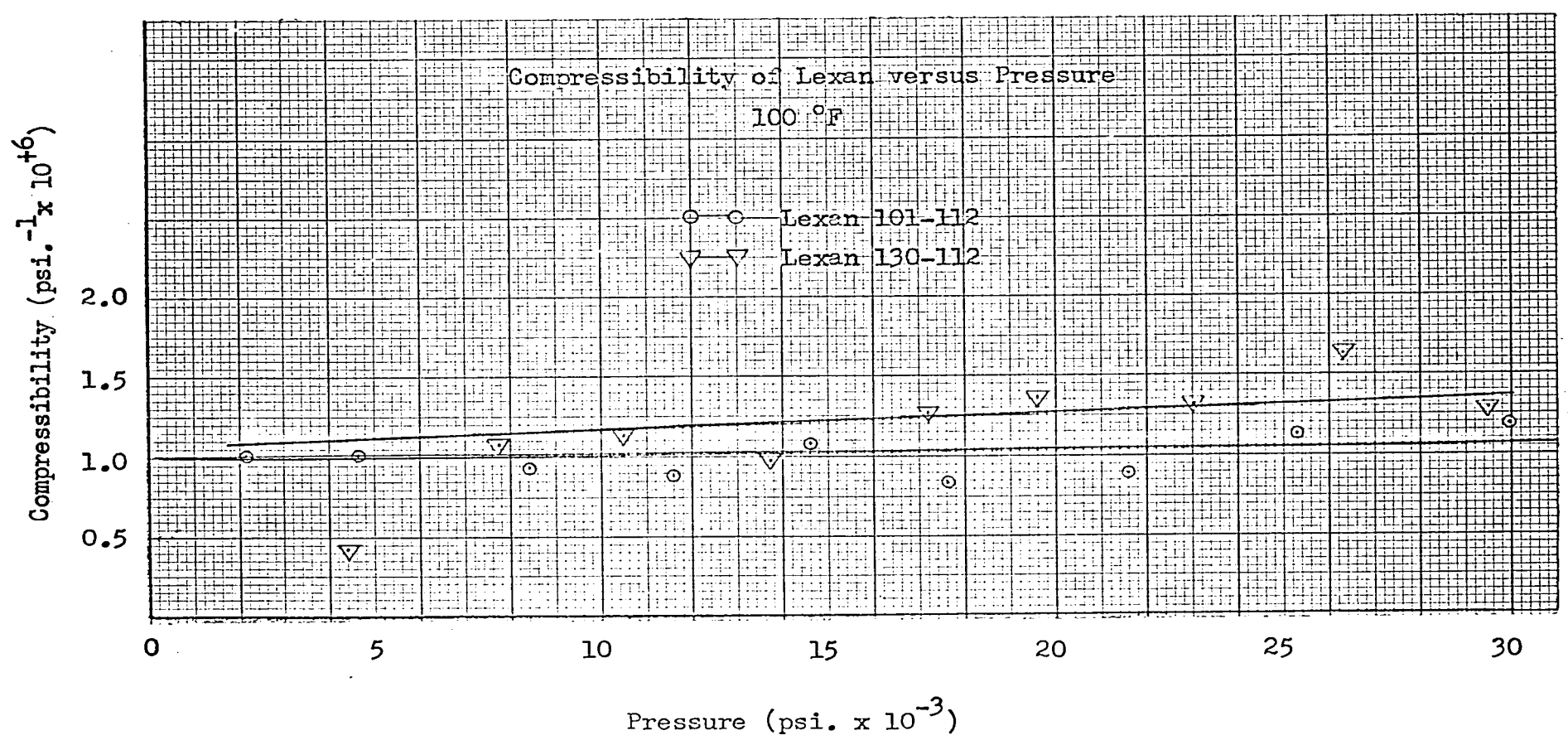


Figure 14

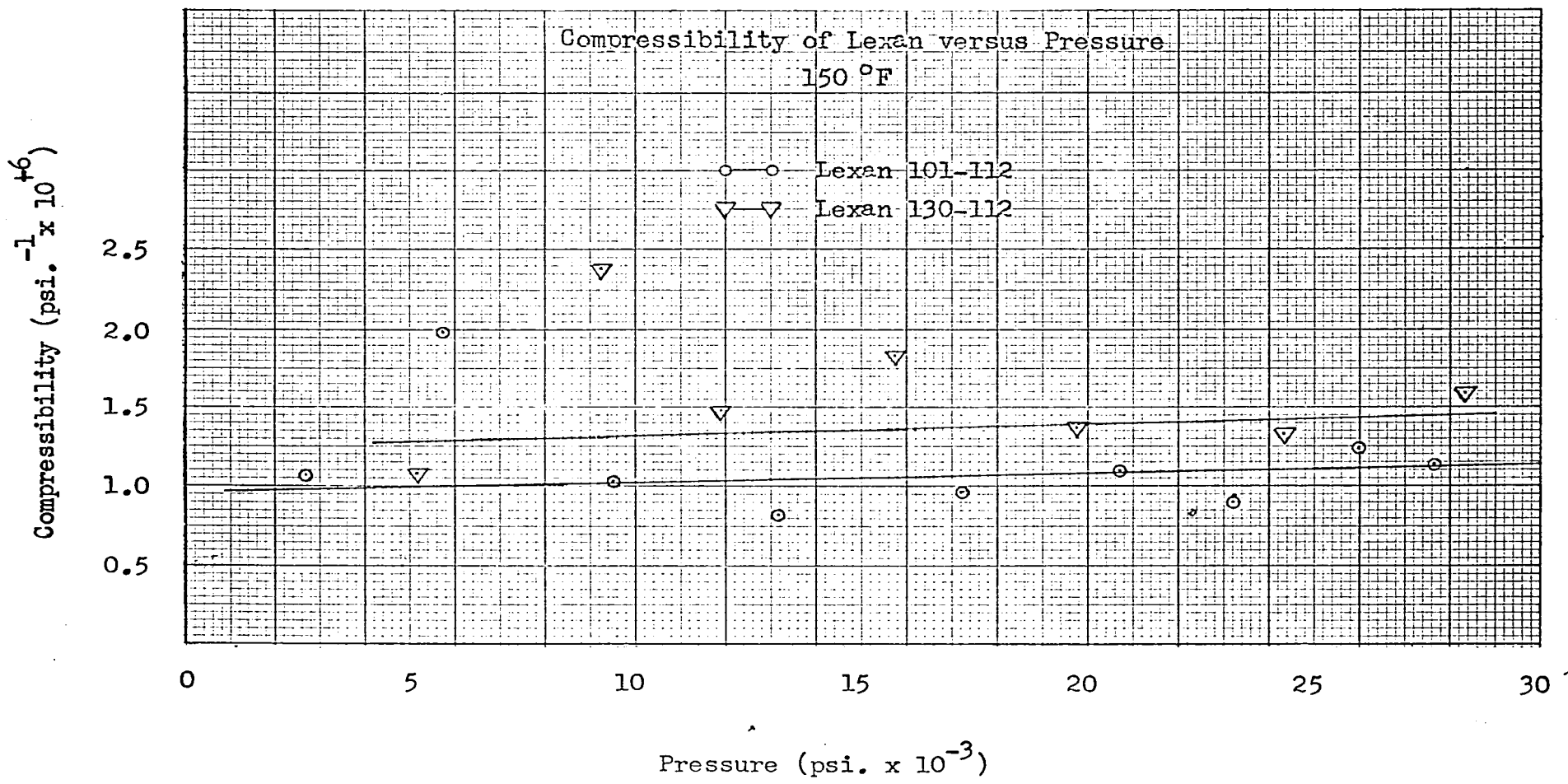


Figure 15

Compressibility ($\text{psi.}^{-1} \times 10^{+6}$)

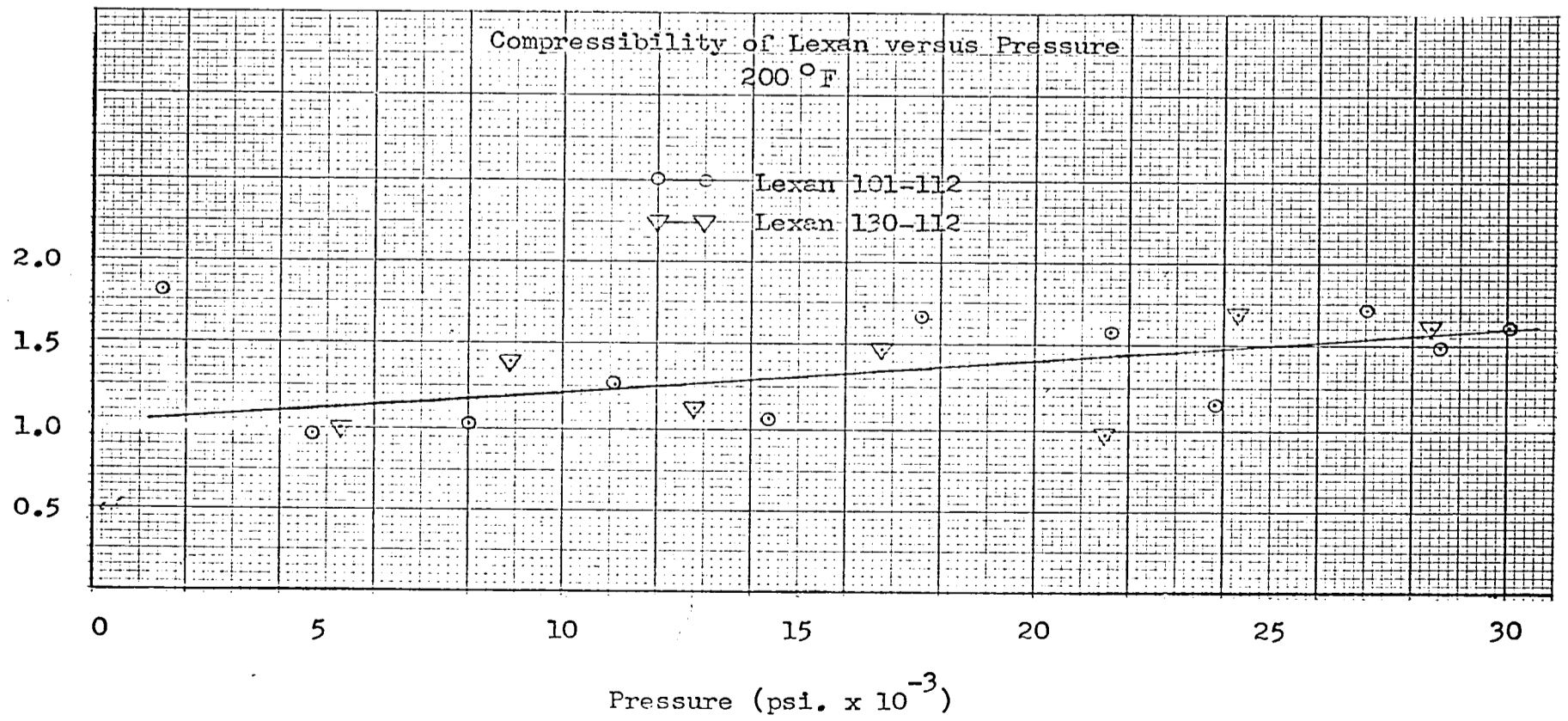


Figure 16

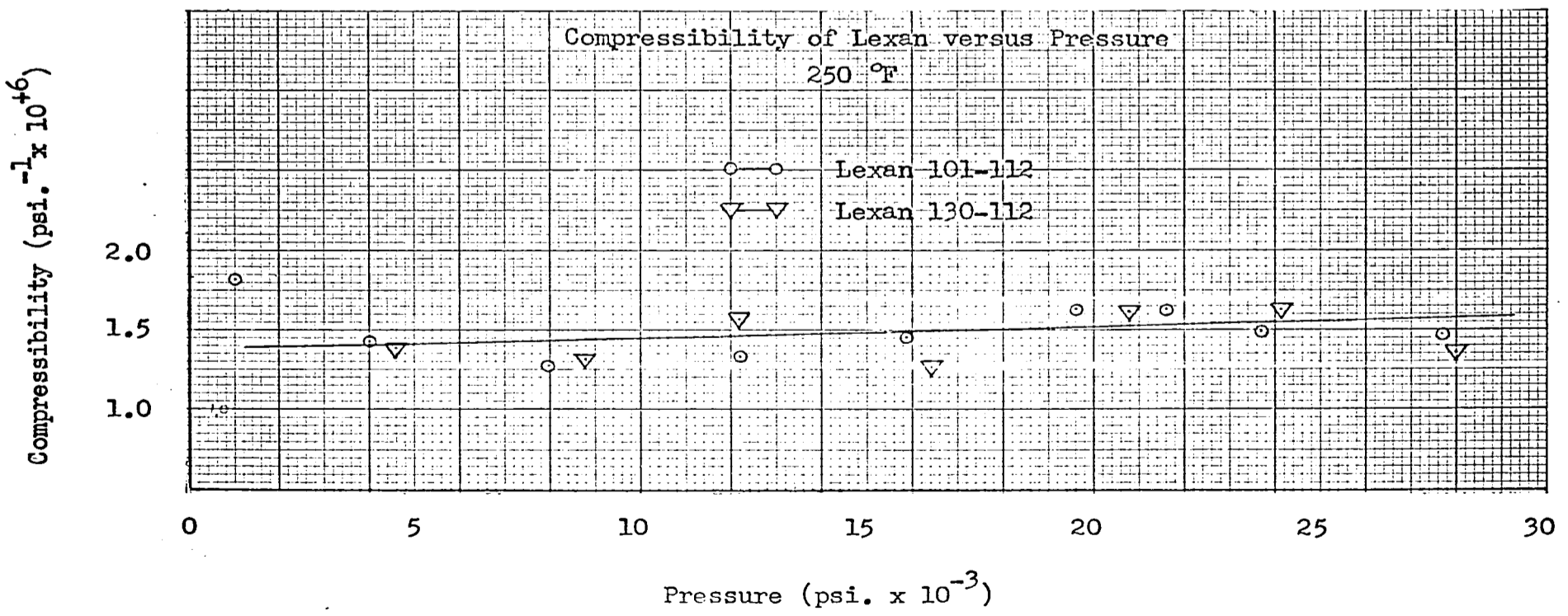


Figure 17

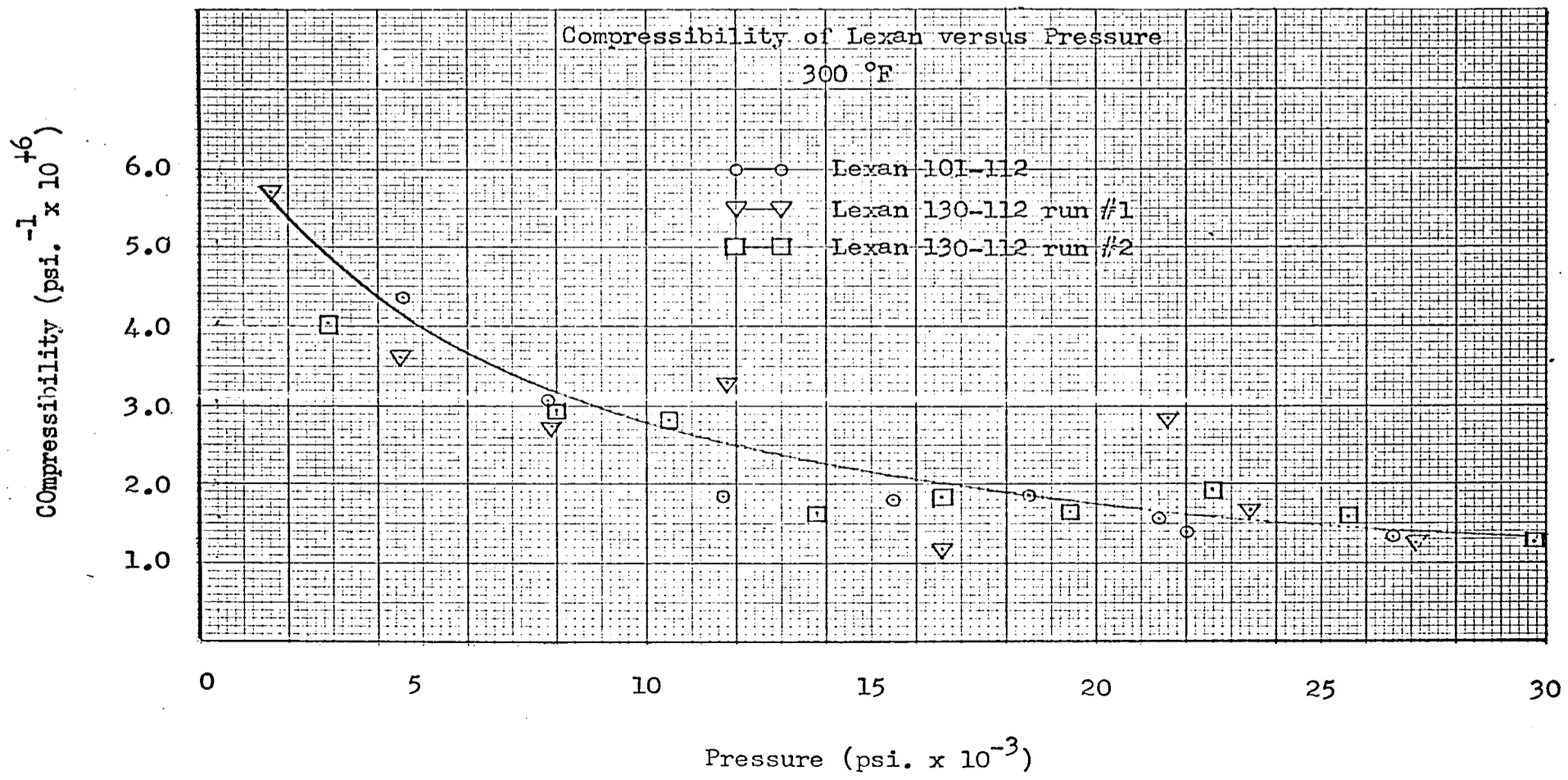


Figure 18

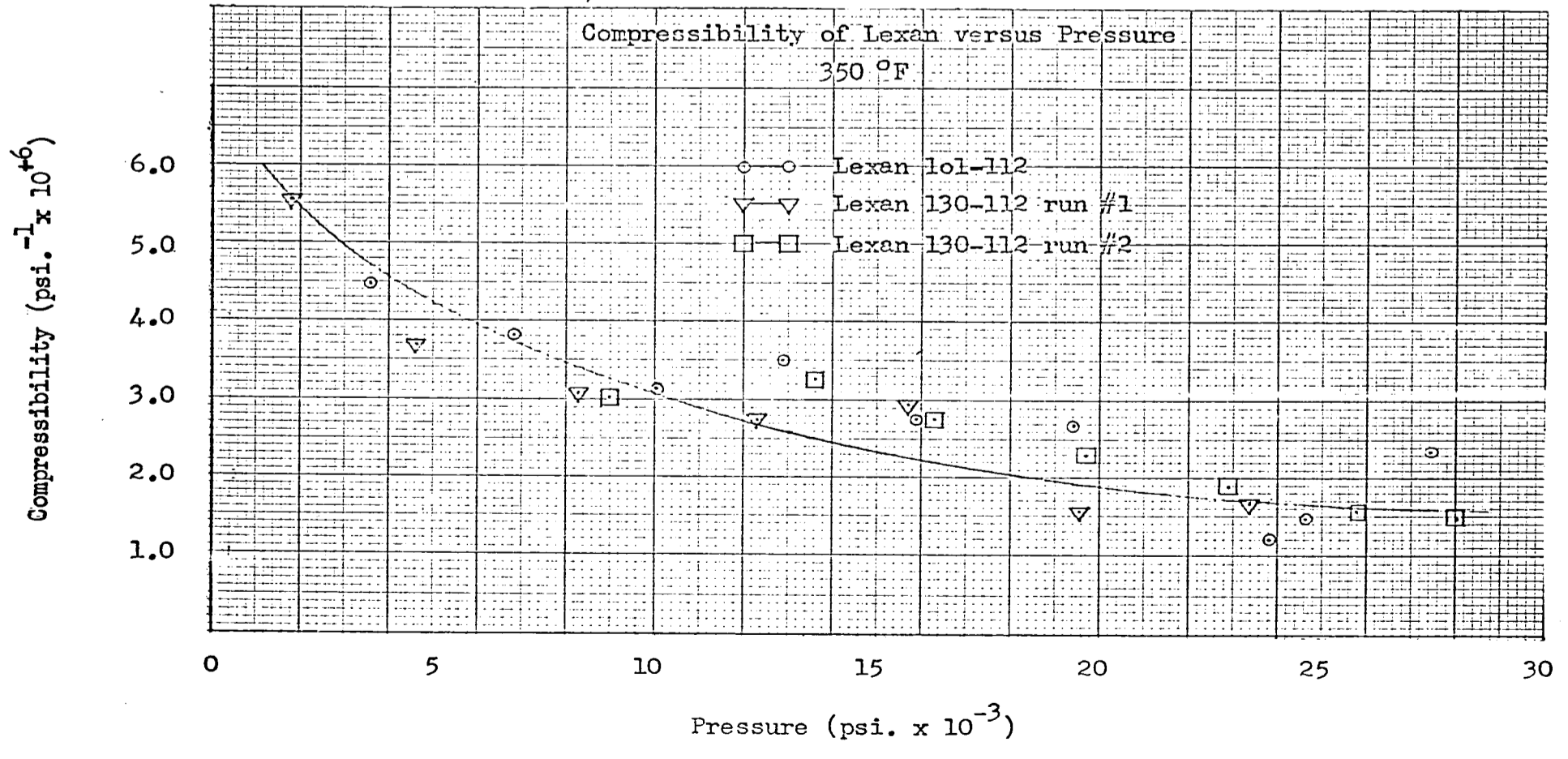


Figure 19

MISSING

PAGES

about 200°F where they become indistinguishable from each other. At 295°F the second order transition temperature is exceeded and the result is a departure from the straight line relations of the lower temperatures to a condition where the material is highly compressible up to about 10,000 Psi. (in the case of the 300°F run) trailing off and approaching the compressibility values obtained below the glass transition temperature. As the temperature is increased the higher compressibilities extend out to higher pressure. Note, however, that even at 350°F the material approaches a compressibility value close to that at lower temperature.

Spencer and Gilmore proposed the modified van der Waal's equation of the form $(P+\pi)(V-\omega)=RT$. It is necessary to calculate two constants; π , the internal pressure in psi., and ω , the specific volume at absolute zero temperature. If the density at one temperature, and the volume coefficient of expansion, $\frac{dV}{dT}$, are known under equilibrium conditions, then ω can be calculated. From Schnell's thermal data on the bisphenol A polycarbonate we find the specific volume to be 0.8852 cc/g. at 200°C and the volume coefficient of expansion to be 5.0×10^{-4} cc/g.°K. The specific volume is then:

$$\omega = 0.8852 \text{ cc/g} - [5.0 \times 10^{-4} \text{ cc/g}^\circ\text{K} (473^\circ\text{K})]$$

or $\omega = 0.6487 \text{ cc/g.}$

Since there are two mechanisms responsible for the volume coefficient, i. e., change due to thermal energy of the molecule and the change due to the configuration of the molecule or the change in position of molecular segments,

there is some question as to which value of the coefficient to use, the instantaneous value or the equilibrium value. Because the values dealt with here are all equilibrium values, it seems consistent to use the equilibrium coefficient of volume expansion.

Rearranging the proposed equation,

A $P = \frac{RT}{(V-\omega)} - \pi$ at constant temperature. Differentiating with respect to volume,

$$\left(\frac{dP}{dV}\right)_T = -\frac{RT}{(V-\omega)^2}$$

B $\sqrt{(V-\omega)^2} = (V-\omega) = \sqrt{-RT\left(\frac{dV}{dP}\right)_T}$

Rearranging A and substituting equation B

C $\pi = \sqrt{-RT\left(\frac{dP}{dV}\right)_T} - P$

Thus, for constant temperature measurements, π may be evaluated directly without extrapolating values.

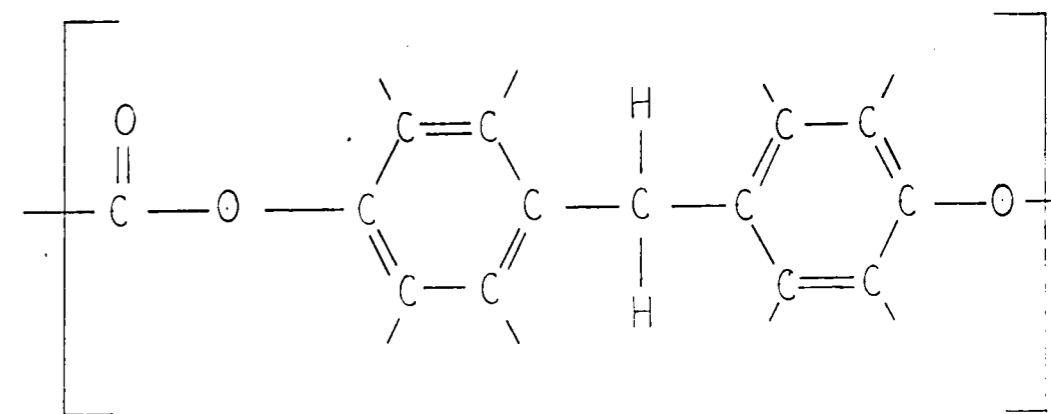
In the unique set of units used for this equation $R=1206/M$, where M is defined as the monomeric "Interaction unit." Since the bisphenol A polycarbonate unit has no side groups it seems reasonable to assume that M is the weight of the repeat unit, in which case:

$$R=1206/226=5.34$$

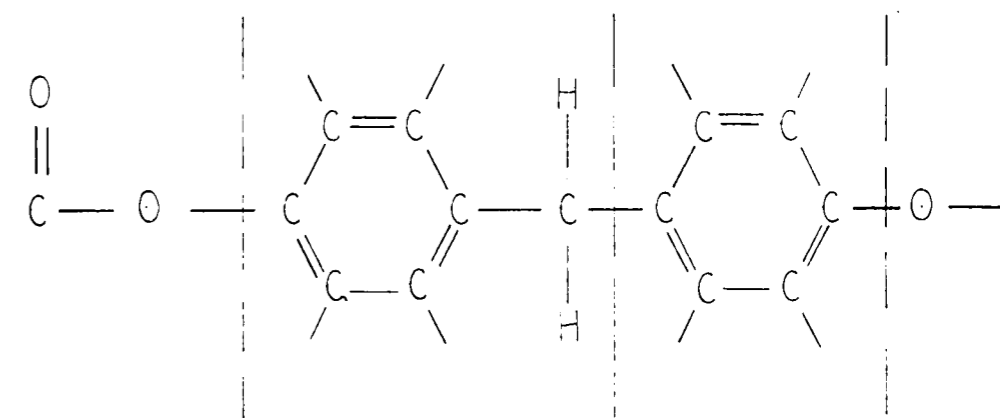
The other possibility, see Figure 20, is that a group about the size of a benzene ring acts as the interaction unit in which case we can predict:

$$R=1206/78=15.5$$

POSSIBLE INTERACTION UNITS



Bisphenol A Polycarbonate Repeat Unit



Benzene Ring Interaction Unit Approximation

Figure 20

TABLE X.

\bar{P} (psi.)	$\Delta P/\Delta V$ (psi./in. ³) $\times 10^{-5}$	$-T\Delta P/\Delta V$ (°K psi./in. ³) $\times 10^{-8}$	$\sqrt{-T\Delta P/\Delta V}$ $\times 10^{-4}$
Lexan at 300°F (423°K)			
27,600	5.531	2.339 3,826	1.530 1.95
24,085	4.394	1.858 3,040	1.364 1.74
22,025	5.018	2.122 3,472	1.458 1.86
18,000	4.226	1.787 2,924	1.337 1.71
15,175	3.821	1.616 2,644	1.272 1.62
11,690	3.833	1.621 2,652	1.621 1.63
9,200	2.491	1.053 1,723	1.027 1.31
4,500	1.944	0.822 1,345	0.907 1.16
Lexan at 350°F (450°K)			
29,030	4.688	2.109 3,450	1.454 1.86
25,765	4.425	1.991 3,257	1.412 1.81
23,050	3.788	1.704 2,788	1.306 1.67
19,710	3.013	1.355 2,217	1.165 1.49
15,875	2.524	1.136 1,858	1.066 1.36
12,290	2.551	1.148 1,878	1.072 1.37
10,045	2.330	1.003 1,641	1.002 1.28
6,855	1.834	0.825 1,350	0.910 1.16
3,585	1.554	0.699 1,144	0.837 1.07

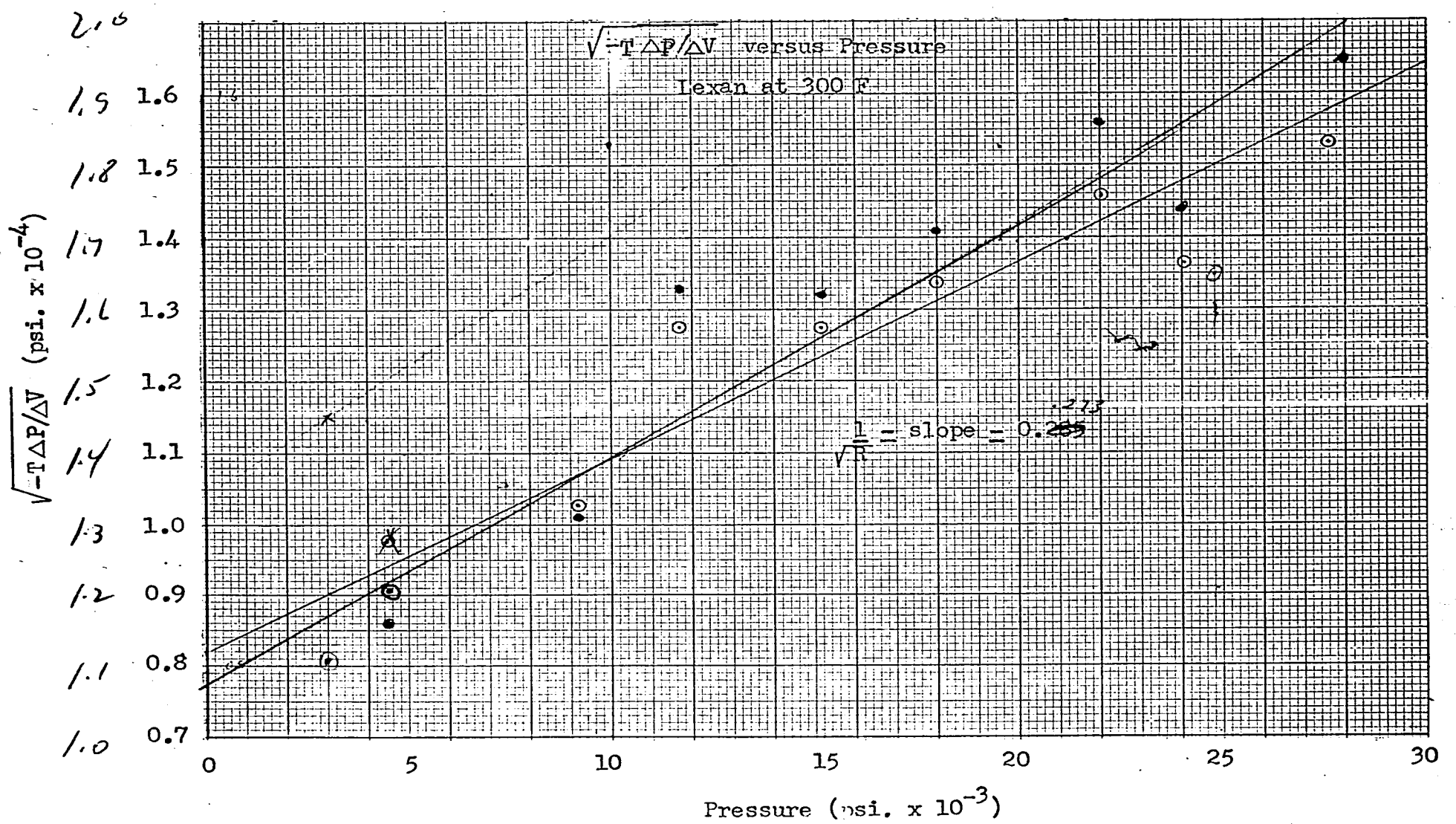


Figure 21

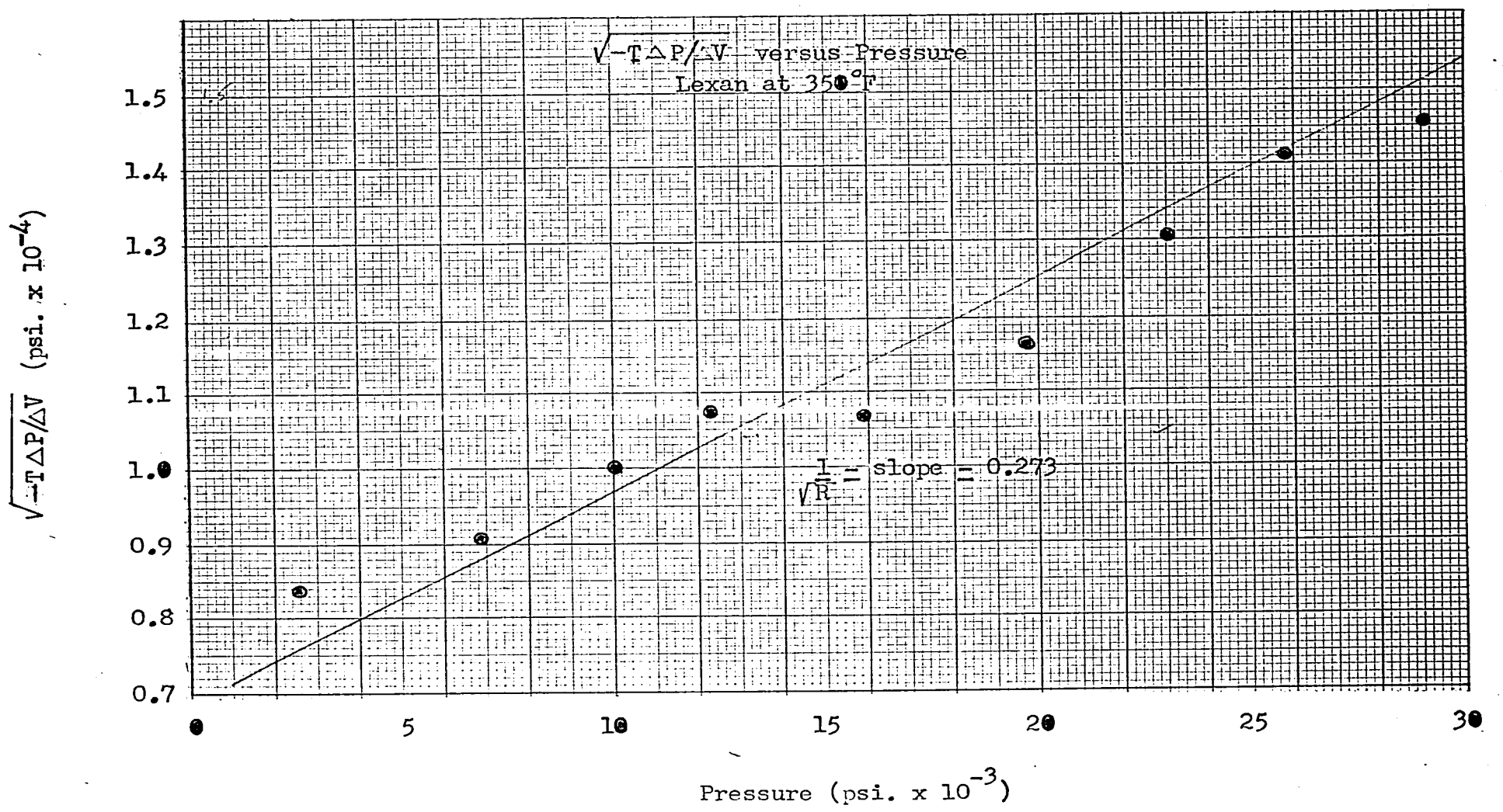


Figure 22

MISSING

PAGES

Taking equation C and dividing by \sqrt{R}

$$\frac{\pi}{\sqrt{R}} = \sqrt{-T \left(\frac{dP}{dV} \right)_T} - \frac{P}{\sqrt{R}}$$

For small intervals $\left(\frac{\Delta P}{\Delta V} \right)_T$ may replace $\left(\frac{dP}{dV} \right)_T$ and P then becomes \bar{P} , the average pressure over the interval.

$$D \quad \frac{\pi}{\sqrt{R}} = \sqrt{-T \left(\frac{\Delta P}{\Delta V} \right)_T} - \frac{\bar{P}}{\sqrt{R}}$$

Plotting $\sqrt{-T \frac{\Delta P}{\Delta V}}$ versus \bar{P} enables R to be calculated empirically from the slope of this plot.

In Table X a number of points are listed for the 300°F and 350°F runs. The points are plotted in Figures 21 and 22 and the least squares line is shown. R, evaluated from the slope, shows that approximately 3 groups the size of the benzene group constitute the "interaction unit" since $R=12.9$.

It should be noted that one equation of this type will not describe the P-V-T relationships both above and below the second order transition temperature. Because our main interest is in determining quantities such as mold size, the equation of interest should then describe the material above this transition temperature as has been started.

In Table XI the resulting values of π are listed. π was found to have an average value of 26,602. The final form of the equation is:

MISSING

PAGES

TABLE XI.
Calculation of π for Lexan at 300°F

$\sqrt{RT\Delta P/\Delta V}$ (psi.)	\bar{P} (psi.)	π (psi.)
For Lexan 101-112		
29,750 ³⁷⁴⁰⁰	4,345	25,415 ~ 33,000
35,300	7,820	27,480
45,810	11,690	34,120
46,120	15,505	30,615
45,320	18,505	26,815
52,390	22,025	30,365
53,980 ⁶²⁸⁰⁰	26,610	27,370 ³⁶⁰⁰⁰
For Lexan 130-112		
25,920	1,615	24,305
32,680	4,515	28,165
37,570	7,890	29,680
47,800	23,450	24,350
55,380	27,330	27,050
30,810	2,905	27,905
36,010	6,050	29,960
36,900	9,250	27,650
48,400	12,175	36,250
45,800	15,175	30,625
48,090	17,985	30,105
45,180	20,990	24,190
49,050	24,085	24,965
55,000	27,670	27,330
Σ all $\pi = 593,610$		
average $\bar{\pi} = 28,267$		

TABLE XI. CONT.

Calculation of $\bar{\Pi}$ for Lexan at 350° F

$\sqrt{RT\Delta P/\Delta V}$ (psi.)	\bar{P} (psi.)	$\bar{\Pi}$ (psi.)
For Lexan 101-112		
30,400	3,585	26,815
32,620	6,855	25,765
37,100	10,045	27,055
34,000	12,900	21,100
38,300	18,875	19,425
38,690	19,400	19,290
57,990	23,800	34,190
For Lexan 130-112		
27,030	1,810	25,220
33,140	4,640	28,500
36,190	8,315	27,875
38,490	12,290	26,200
37,040	15,730	21,310
51,010	19,565	31,545
49,180	23,385	25,795
36,650	10,330	26,320
35,410	13,640	21,950
38,310	16,310	22,000
41,830	19,710	22,120
46,260	22,875	23,385
50,720	25,765	24,655
52,200	29,030	23,170

$$\sum \text{all } \bar{\Pi} = 523,675$$

$$\text{average } \bar{\Pi} = 24,937$$

$$(P + 26,602)(V - 0.6487) = 12.9T$$

Plotting the above equation with the data obtained experimentally reveals the fact that the equation does not satisfactorily represent the data, in fact, it deviates from the experimental values by as much as seven percent of the total volume. As a result the variable parameters M , the interaction unit, and $\hat{\pi}$, the internal pressure, were adjusted to yield a good fit. The final value of M was chosen to yield a value of $R = 20.8$, and the final value of $\hat{\pi}$, the internal pressure, was 40,000. The resulting adjusted equation is:

$$(P + 40,000)(V - 0.6487) = 20.8T$$

In Table XII the experimentally determined volumes at various pressures are tabulated. In Figures 23 and 24 the adjusted equation is plotted along with the experimental points to compare the fit obtained.

The Tait equation:

I. $\frac{V_0 - V}{V_0} = c \text{ Log } (1 + P/b)$ is not particularly useful in this form. Differentiating for the $\frac{\partial V}{\partial P}$ we obtain the following expression:

$$\text{II. } \frac{\partial V}{\partial P} = -V_0 (0.4343c) \frac{1}{b+P}$$

upon rearranging,

TABLE XII.

Experimental Pressure - Volume Equilibrium Values

Cross-sectional area of cylinder = 1.0136"

Pressure (psi.)	Sample Height (in.)	Sample Volume (in ³)	Specific Volume (cc/g)
Lexan 101-112 at 300°F		Sample weight = 27.197 g.	
0	1.4056	1.4248	0.8584
3,050	1.3860	1.4048	0.8464
5,640	1.3702	1.3888	0.8369
10,000	1.3513	1.3697	0.8253
13,500	1.3426	1.3609	0.8199
17,630	1.3318	1.3499	0.8133
19,380	1.3272	1.3452	0.8105
24,670	1.3168	1.3347	0.8042
28,550	1.3096	1.3274	0.8001
Lexan 130-112 at 300°F		Sample weight = 26.821 g.	
0	1.4054	1.4245	0.8704
2,880	1.3820	1.4008	0.8559
6,150	1.3655	1.3841	0.8458
9,630	1.3521	1.3705	0.8373
13,990	1.3419	1.3601	0.8310
19,110	1.3333	1.3514	0.8256
21,600	1.3233	1.3413	0.8196
25,300	1.3146	1.3325	0.8142
29,360	1.3075	1.3253	0.8097
Lexan 130-112 at 300°F		Sample weight = 27.034 g.	
0	1.4074	1.4265	0.8648
1,660	1.3902	1.4091	0.8541
4,150	1.3761	1.3948	0.8454
7,950	1.3603	1.3788	0.8358
10,550	1.3500	1.3684	0.8296
13,800	1.3420	1.3603	0.8246
16,550	1.3349	1.3531	0.8202
19,420	1.3282	1.3463	0.8161
22,560	1.3198	1.3377	0.8109
25,610	1.3129	1.3308	0.8068
29,730	1.3056	1.3234	0.8022

TABLE XII. CONT.

<u>Pressure</u>	<u>Sample Height</u>	<u>Sample Volume</u>	<u>Specific Volume</u>
Lexan 101-112 at 350°F		Sample weight = 27.197 g.	
0	1.4234	1.4328	0.8633
1,900	1.4086	1.4277	0.8601
5,270	1.3872	1.4061	0.8472
8,440	1.3701	1.3887	0.8367
11,650	1.3550	1.3734	0.8276
13,950	1.3436	1.3619	0.8205
17,800	1.3285	1.3461	0.8111
21,000	1.3163	1.3342	0.8039
26,600	1.3068	1.3246	0.7981
28,300	1.3012	1.3189	0.7946
Lexan 130-112 at 350°F		Sample weight = 26.821 g.	
0	1.4234	1.4428	0.8815
3,130	1.3971	1.4161	0.8653
6,150	1.3812	1.3999	0.8553
10,480	1.3623	1.3808	0.8436
14,100	1.3483	1.3666	0.8350
17,360	1.3347	1.3528	0.8266
21,770	1.3250	1.3430	0.8205
25,000	1.3173	1.3352	0.8158
Lexan 130-112 at 350°F		Sample weight = 27.034 g.	
0	1.4149	1.4342	0.8694
9,000	1.3695	1.3881	0.8415
11,600	1.3586	1.3771	0.8348
14,620	1.3459	1.3642	0.8269
18,000	1.3327	1.3508	0.8189
21,420	1.3215	1.3395	0.8120
24,330	1.3137	1.3316	0.8073
27,200	1.3073	1.3251	0.8033
30,860	1.2986	1.3163	0.7979

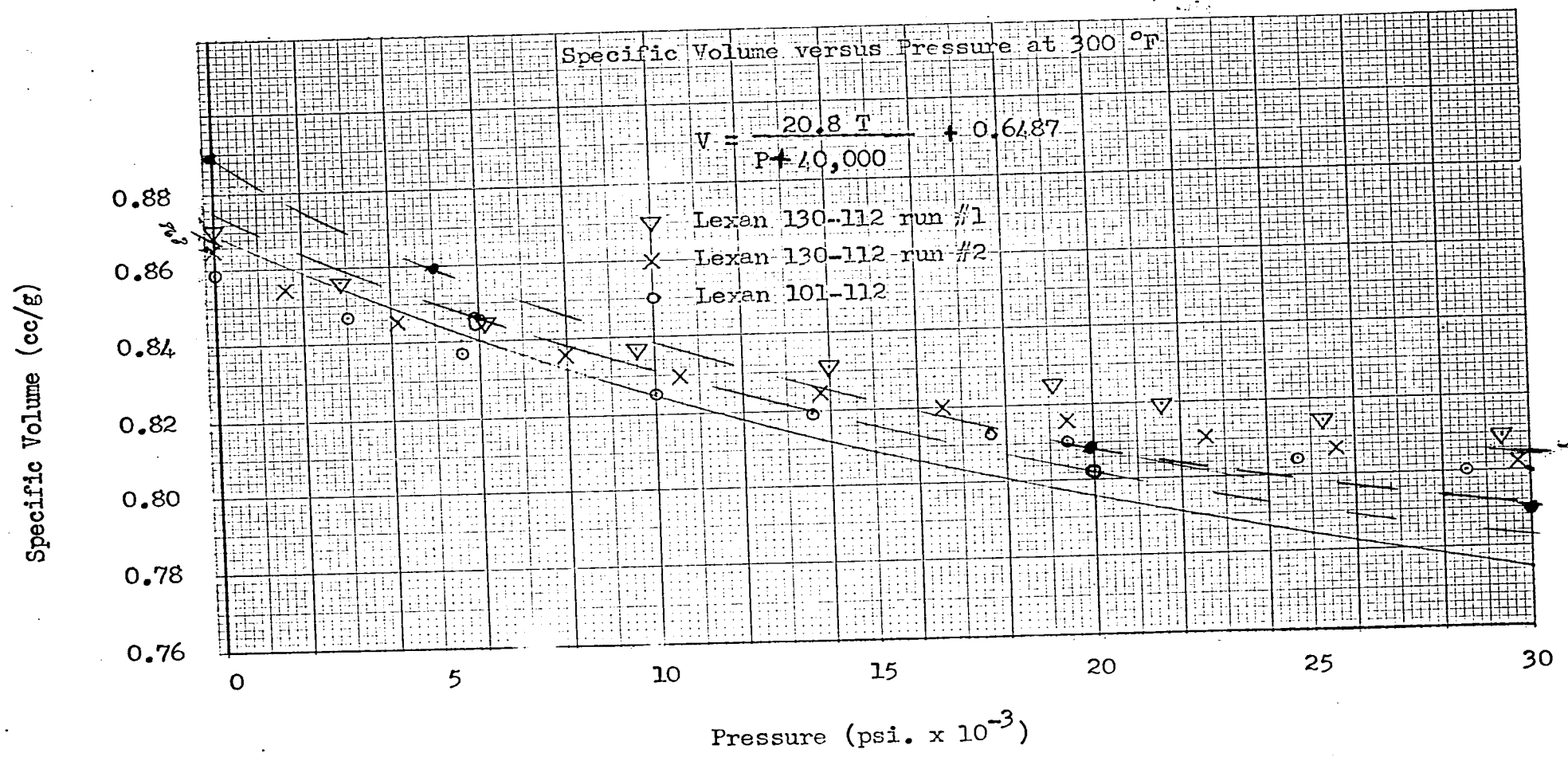


Figure 23

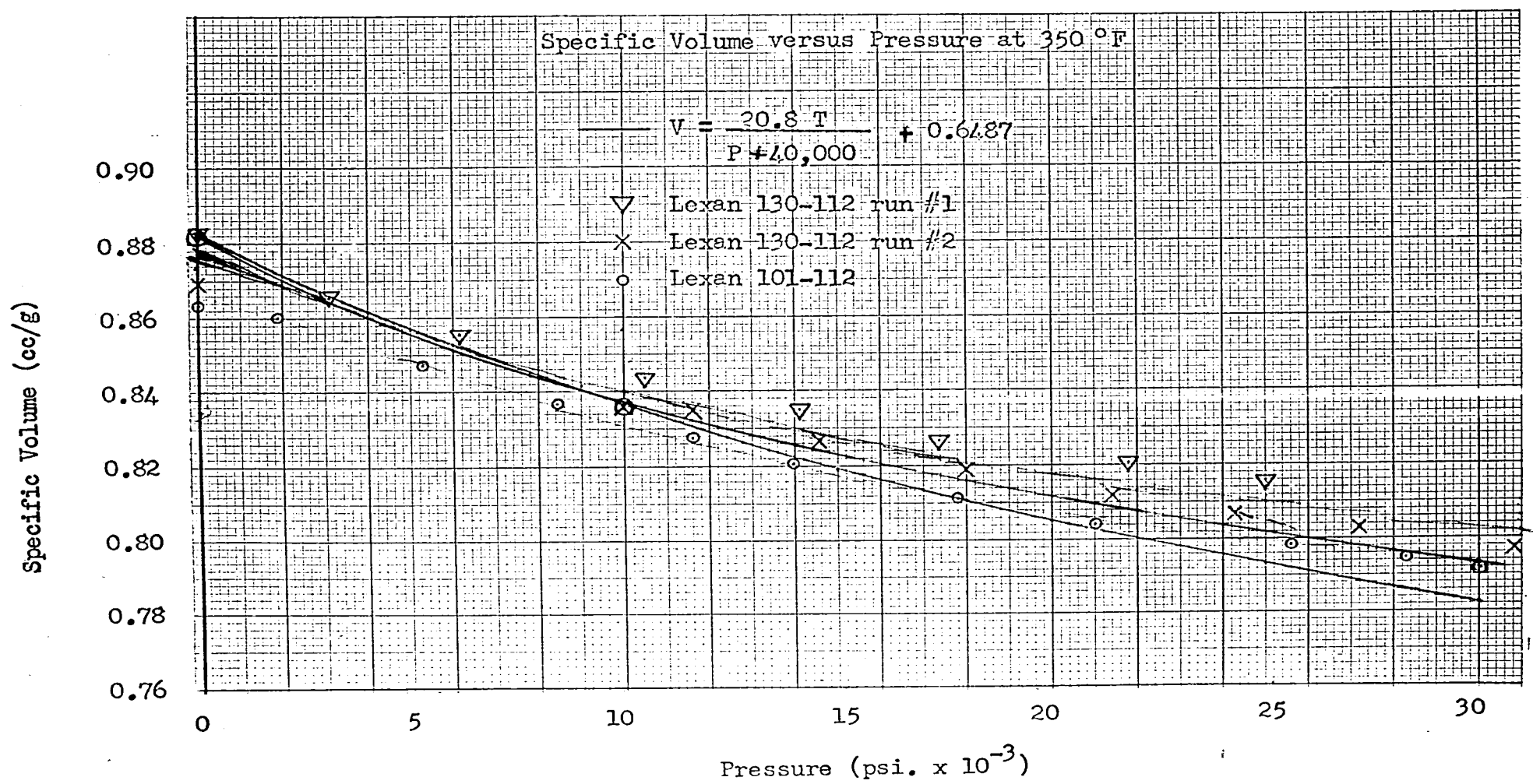


Figure 24

62

MISSING

PAGES

TABLE XIII.

Pressure versus Bulk Modulus

300°F		350°F	
Pressure \bar{P} (psi.)	Bulk Modulus K $1/\beta = \text{psi.} \times 10^{-5}$	Pressure \bar{P} (psi.)	Bulk Modulus K $1/\beta = \text{psi.} \times 10^{-5}$
Lexan 101-112			
4,345	2.304	3,585	2.226
7,820	3.241	6,855	2.627
11,690	5.461	10,045	3.194
15,505	5.531	12,900	2.851
18,505	5.345	15,875	3.615
22,025	7.148	19,400	3.706
26,610	7.570	23,800	8.285
21,430	6.398	27,450	4.290
		24,650	6.812
		23,050	5.426
Lexan 130-112			
1,615	1.752	1,810	1.802
4,515	2.769	4,640	2.711
7,890	3.664	8,315	3.229
11,810	3.027	12,290	3.655
16,550	8.467	15,730	3.388
20,355	3.500	19,565	6.427
23,450	5.942	23,385	5.966
27,330	7.981		
2,905	2.477	10,330	3.318
6,050	3.385	13,640	3.063
9,250	3.552	16,310	3.623
12,175	6.097	19,710	4.321
15,175	5.452	22,875	5.274
17,985	6.031	25,765	6.345
20,990	5.230	29,030	6.725
24,085	6.270		
27,670	7.893		

MISSING

PAGES

$$\text{III } \frac{-1}{V_0} \left(\frac{\partial V}{\partial P} \right) = \frac{0.4343c}{b+P} = \beta$$

where β is the compressibility. Taking the inverse of equation III,

$$\text{IV } K = 1/\beta = \frac{b}{0.4343c} + \frac{P}{0.4343c}, \quad \text{we obtain an equation of}$$

a straight line for the bulk modulus, K . Obviously, plotting the bulk modulus, K , against the pressure P , should yield a straight line by which the constant C could be evaluated from the slope of the plot and b could be calculated from the intercept. Note, however, that the equation predicts only one straight line. What happens if there is a discontinuity, i. e., a change in slope and a change in intercept from a certain group of points? The obvious conclusion is that the K versus P relationship or mechanism has changed. It is this result that is of importance to us here since it allows one to approximate the second order transition pressure for a material above the glass transition temperature.

In Table XIII the bulk modulus and corresponding pressures are tabulated for all points obtained on the Lexan above the second order transition temperature. Figures 25 and 26 contain plots of bulk modulus versus pressure for temperatures of 300 and 350°F, respectively. In both figures it can be seen that there is a discontinuity at about 15,000 psi. It is about this pressure where the material changes from its

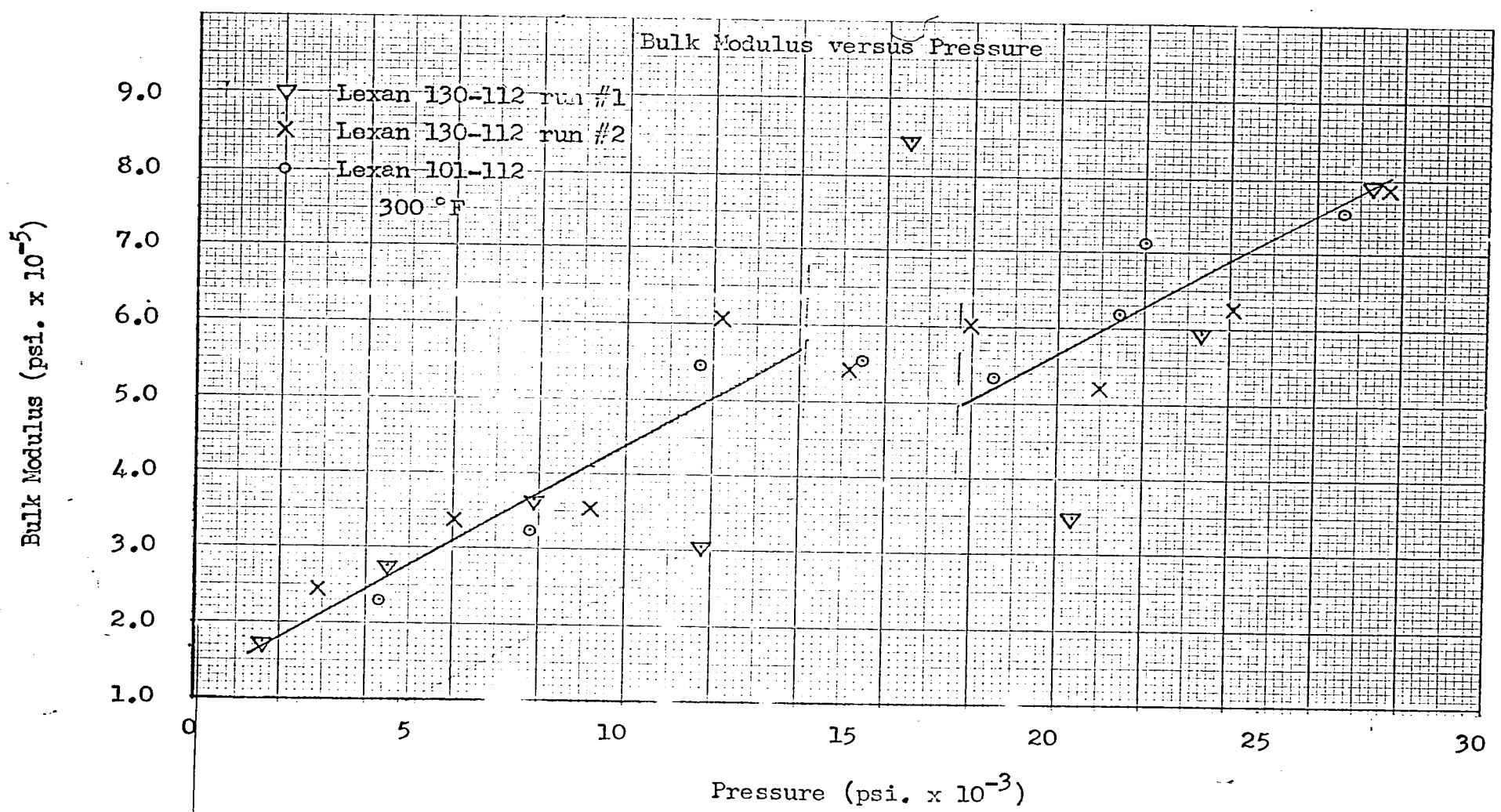


Figure 25

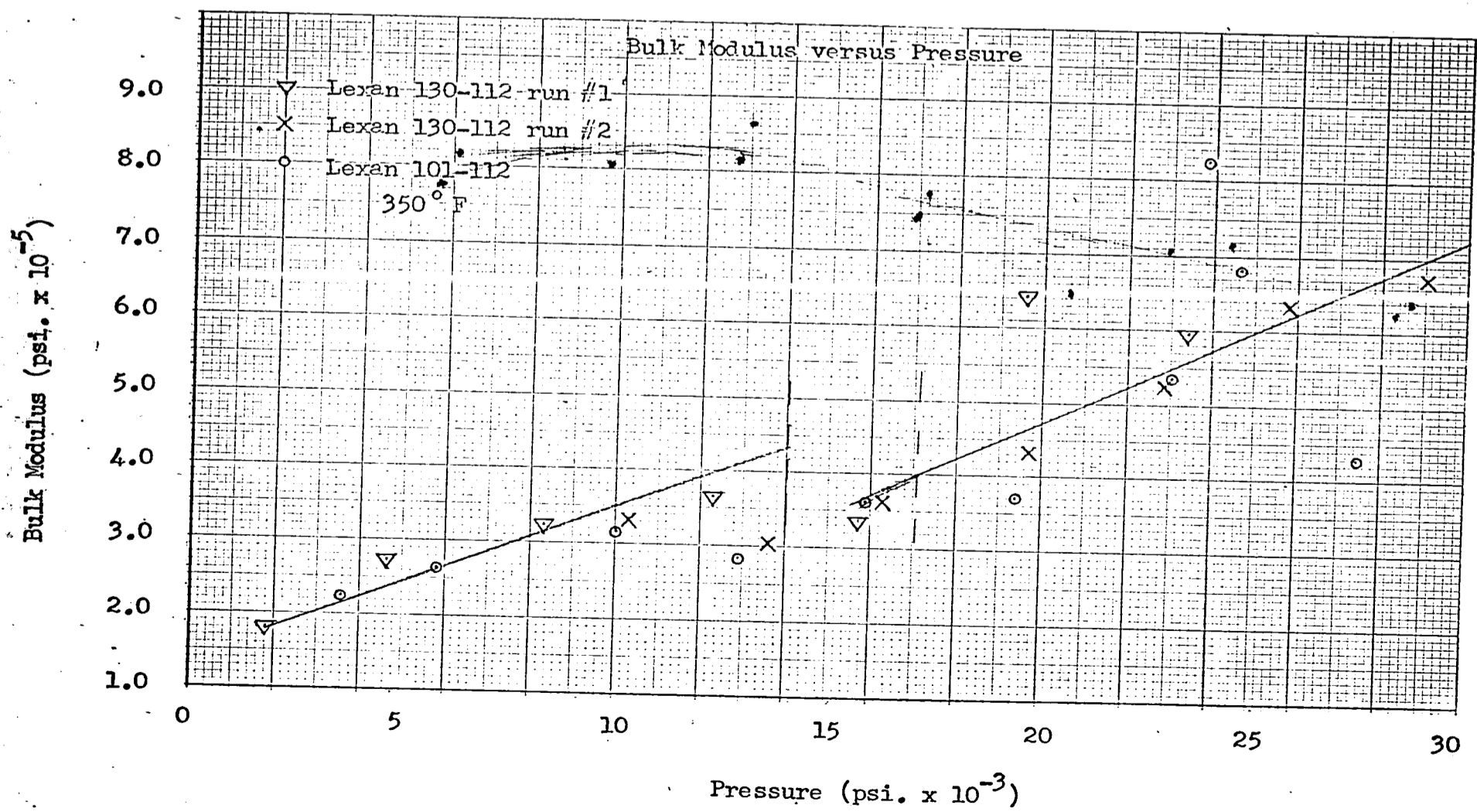
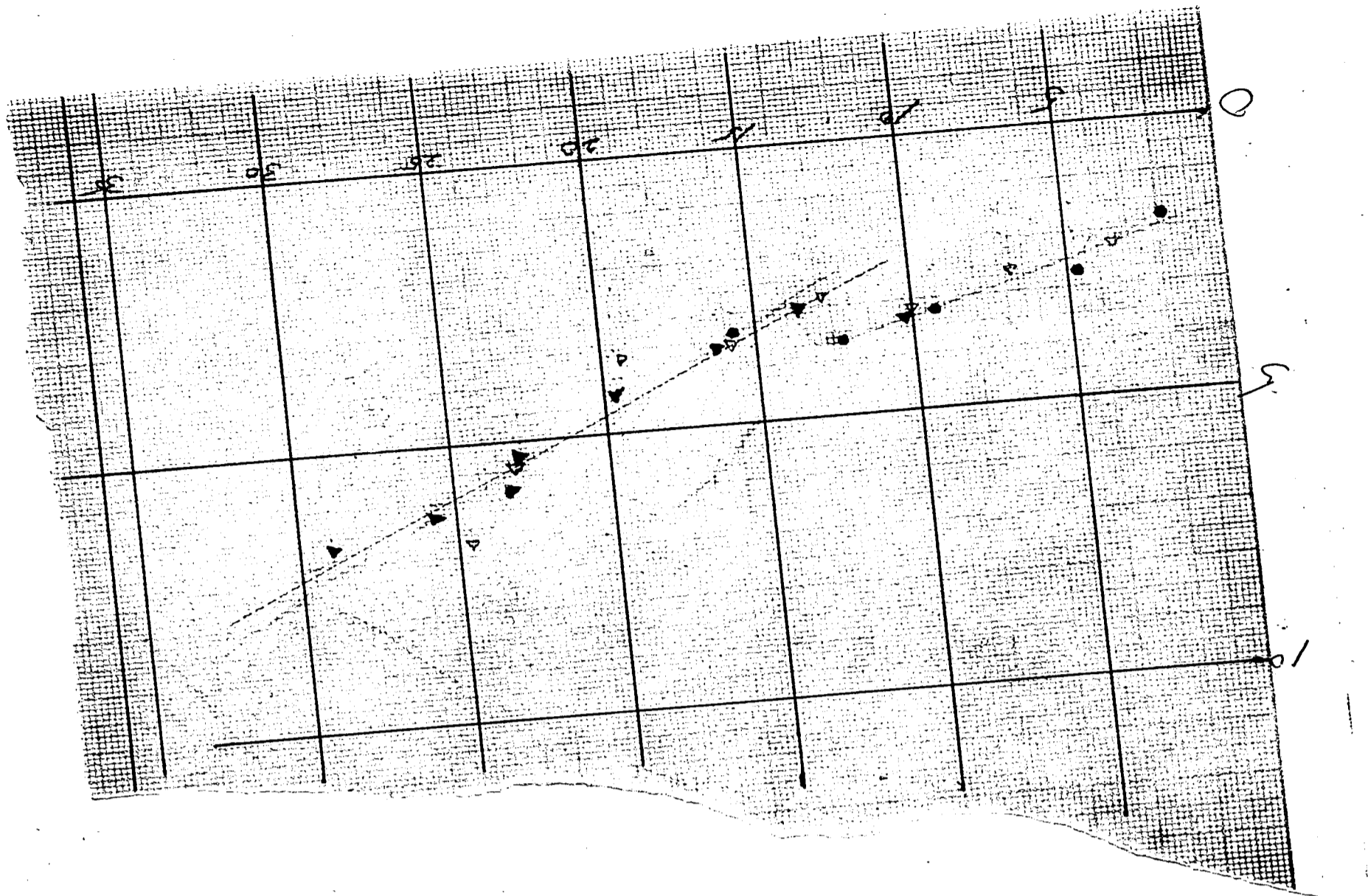


Figure 26

rubbery state to its glassy state. To exactly pinpoint the transition pressure several points obtained near this pressure are needed to eliminate the range of uncertainty.

Discussion of Results

The experimental work performed on Lexan tends to substantiate the prior work performed on high molecular weight amorphous polymers. The compressibility is readily seen to be a function of more than one mechanism as proposed by Simha. One mechanism is that of molecular size due to thermal energy acting similar to a molecule in a lattice sight. This mechanism is predominate at lower temperatures and accounts for the difference in the compressibility with average molecular weight. As the temperature increases each molecule becomes more mobile and begins to acquire enough agility to actually rearrange segments of itself to make itself small under increasing pressure. With increasing temperature the molecule tends to become larger; with increasing pressure to become smaller. For Lexan, a temperature of about 200°F gives the molecule enough mobility so that the difference in compressibility with different molecular weight samples approaches zero.

Of course, above the glass transition temperature the type of mechanism that predominates is that of molecular rearrangement. In the lower pressure ranges above this temperature the large effects of this mechanism are clearly demonstrated with values of compressibility four to five times that at the lower temperatures.

The possibility of a change in state with increasing pressure is substantiated when the data is applied to the Tait relationship. There is no question that there are two lines formed each having a different slope and intercept which can only indicate that some internal change of mechanism has occurred. In this case, a change from a rubbery state at low pressures to a glassy state at high pressures.

The inability of the derived van der waal's equation to satisfactorily fit the data demonstrates the superficial nature of the model. The model is far too simple to accurately represent the complex interaction of a polymeric material. The adjusted equation of this type was able to approximate the volume-pressure relationship sufficiently for engineering applications. In the critical pressure range, i. e., 10,000 to 20,000 psi. where the majority of Lexan is processed, the accuracy of the equation is quite good. At higher temperatures the fit could be improved.

Molecular weight has been shown to have a negligible effect on the P-V-T properties at elevated temperature. This does not imply, however, that a Lexan polymer of extremely high molecular weight (or low molecular weight) will follow this equation. This equation was developed for the commercially obtainable grades of bisphenol A polycarbonate. These grades generally include average molecular weights of 32,000 to 42,000. Outside this weight range the properties of the material make

it difficult to fabricate and hence these extremes are used only for very special applications. Generally speaking, the equation derived here, and those developed earlier by Spencer and Gilmore, should be valid if the material is within ± 15 percent of the molecular weights tested, which includes most of the commercially available materials.

Observed, but not investigated during the course of the experiment, was a "hysteresis effect" on the volume of the sample depending on whether or not the volume at a point was taken while increasing the pressure or decreasing the pressure. The volume of the polymer sample during a decreasing pressure run was found to be from 30 percent to 40 percent of the total volume change due to 30,000 psi. greater than at a corresponding point during an increasing pressure measurement. A similar run was made on a cylindrical sample of brass and the difference had a maximum of 4 percent difference based on the total volume change over the entire pressure range. The polymer molecule which tends to rearrange itself upon compression and become smaller apparently does not contain sufficient energy to extend itself when part of the pressure is released to attain the state achieved with increasing pressure measurement. It does exhibit a "memory" since it returns to its original volume when all pressure is removed.

Conclusions

From the P-V-T data gathered in this project the following conclusions can be drawn.

1.) For temperatures above the glass transition temperature the compressibility of commercially available Lexan is essentially independent of molecular weight.

2.) Combined with the work of Spencer and Gilmore on polystyrene, one may deduce a trend that indicates that most commercial high molecular weight amorphous polymers will have compressibilities independent of molecular weight at temperatures above the second order transition temperature.

3.) At temperatures well below the glass transition temperature the Lexan having the highest average molecular weight had a compressibility higher than the compressibility for the lower molecular weight material. The difference averages to approximately 25 percent.

4.) The modified van der Waal's equation $(P+40,00)(V-0.6487) = 20.8T$ will describe the behavior of the bisphenol A polycarbonate above the glass transition temperature.

5.) It has been verified that at temperatures above the glass transition temperature, the Tait Equation, when rearranged in a form relating the bulk modulus to the pressure, will indicate a pressure at which Lexan changes from a rubbery material to a glassy material. This second order transition pressure was found to be about 15,000 psi. for temperatures between 300 and 350°F.

Nomenclature

b	- constant
c	- constant
g	- stress-strain correction factor constant
K	- bulk modulus
l	- stretched length
l_0	- unstretched length
M	- molecular interaction unit
M_c	- molecular weight of a repeat unit
P	- pressure (psi.)
R	- gas constant
T	- temperature (absolute)
v	- number of moles of molecular interaction segments
V	- volume (cc/g)
V_0	- volume at zero pressure
V_1	- volume at one atmosphere pressure
β	- compressibility (psi. ⁻¹)
π	- "internal" pressure constant (psi.)
τ	- tension
ω	- volume at zero degrees Kelvin

BIBLIOGRAPHY

- 1.) Billmeyer, Fred W. Textbook of Polymer Science.
New York: Interscience Publishers, 1962.
- 2.) Christopher, W. and Fox, D. Polycarbonates.
New York: Reinhold Publishing Company, 1962.
- 3.) Gordon, M. High Polymers. London: ILIFFE Books
Limited, 1963.
- 4.) Heydemann, P. "The dynamic compressibility of high
polymers in the frequency range of 0.1 cps.
to 60 Kcs." Acoustica, IX (1959), 446-452.
- 5.) Nason, H.K. "High temperature, high pressure
rheometer for plastics!" Journal of Applied
Physics, XVI (June, 1945), 338-343.
- 6.) Schnell, Herman. Chemistry and Physics of Poly-
carbonates. New York: Interscience Publishers,
1964.
- 7.) Scott, A.H. "Effect of pressure on dielectric constant,
power factor, and conductivity of rubber-sulphur
compounds." Journal of Research, National
Bureau of Standards, XIV (1935), 99.
- 8.) Simha, R. "On anomalies of elasticity and flow and
their interpretation." Journal of Physical
Chemistry, LIII (April, 1943) 297.
- 9.) Spencer, R.S. and Gilmore, G.D. "Equation of
state for polystyrene." Journal of Applied
Physics, XX (June, 1949) 502-506.
- 10.) Spencer, R., Gilmore, G., and Wiley, R. "Behavior
of granulated polymers under pressure."
Journal of Applied Physics, XXI (1950), 527-521.
- 11.) Spencer, R. and Gilmore, G. "Equation of state for
high polymers!" Journal of Applied Physics,
XXI (1950), 523-527.
- 12.) Spencer, R. and Boyer, R. "Thermal expansion and
second order thermal effects in high polymers."
Journal of Applied Physics, XVII (1946), 398-403.

- 13.) Weir, G.E. "Transitions and Phases of Polytetrafluoroethylene". Journal of Research, National Bureau of Standards, L (February, 1953) 95-97.
- 14.) Wood, L.A. "Utility of the Tait Equation Relating Volume and Pressure In the Study of Transitions In Polymers." Journal of Polymer Science, II (1964) 703-707

Vita

The author was born in Jamestown, New York, on May 6, 1940. He graduated from Jamestown High School in June, 1958.

He enrolled in Bucknell Universities' five year engineering program for the B.A. and B.S. degrees in Chemical Engineering and was graduated in June, 1963.

In the fall of 1963 he entered Lehigh University to pursue a course of study leading to the Master of Science Degree in Chemical Engineering.